



1 Attenuation and Salinity Dominance of Submesoscale Surface Density 2 Gradients in the German Bight

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8 Abstract

9 Submesoscale horizontal density gradients at the ocean surface play a vital role in driving
10 upper-ocean mixing and air-sea exchange. However, horizontal density structures are often
11 assumed to directly reflect the underlying water dynamics. This study presents high-resolution
12 along-track observations collected by the autonomous surface vehicle HALOBATES in the
13 German Bight. The measurements capture co-located horizontal density gradients in both the
14 ocean skin layer, sampled at approximately 80 μm , and the near-surface layer (NSL) at 1 m
15 depth. The results reveal partial decoupling and significant attenuation of horizontal density
16 gradients between the skin and the NSL (slope = 0.41 ± 0.008 , $R^2 = 0.18$), alongside a persistent
17 baseline of skin layer spatial variability that produces a higher front-detection rate than in the
18 NSL (8.3% against 7.2%). Density fronts at the skin are disproportionately salinity-driven
19 rather than temperature-driven (44.1% salinity-only and 33.7% temperature-only), more so than
20 at the NSL (39.9% against 31.8%). Tidal phase significantly modulates gradient intensity at
21 both layers. Random Forest machine learning coupled with SHapley Additive exPlanations
22 identified skin temperature and salinity as the dominant predictors of gradient intensity, while
23 atmospheric pressure, wind speed, and air-sea temperature difference contributed substantially.
24 Nevertheless, predictive skill across the full range of gradient intensities remained limited (test
25 $R^2 = 0.19$). These findings demonstrate that submesoscale density gradients at the skin layer
26 can decouple substantially from underlying upper-ocean structures. This vertical disconnect
27 highlights potential biases in satellite validation and proxy-based approaches that rely on sea-
28 surface skin measurements to infer submesoscale dynamics in complex coastal waters.



1. Introduction

The oceanic skin layer, or sea surface microlayer, forms the dynamic boundary interface between the ocean and atmosphere (Ribas-Ribas et al., 2017; Wurl et al., 2024). It directly modulates air–sea exchanges of heat and gases, surfactant accumulation, and micro-scale wave dissipation. Despite its pivotal role in global biogeochemical cycles and upper-ocean thermodynamics, observational oceanography has historically relied on bulk or near-surface layer (NSL, 1 m depth) measurements collected at depths of 1 to 10 m via moorings, research vessels, or FerryBox automated monitoring networks (Grayek et al., 2011). While satellite remote sensing provides broad spatial coverage, it lacks the temporal continuity, cloud-free reliability, and vertical resolving power needed to continuously sample fine-scale gradients between the skin and underlying water column (Minnett et al., 2019).

Submesoscale processes operating on horizontal scales of 100 m to 50 km are characterized by intense spatial gradients, buoyant plume dynamics, lateral straining, and strong localized vertical motion (Capet et al., 2008; D’Asaro et al., 2011; McWilliams, 2016; Su et al., 2018). In shallow, tidally energetic coastal environments like the German Bight in the southeastern North Sea (Hagen et al., 2022), submesoscale variability is further modulated by complex interactions between tidal straining, seasonal thermal stratification, and major estuarine freshwater discharge from the Elbe, Weser, and Ems rivers (Simpson and Hunter, 1974). Existing submesoscale frameworks are typically defined under the implicit assumption that horizontal gradients observed within the NSL pass unimpeded to the absolute surface skin layer. This assumption is rarely tested directly because of the lack of instrumentation required to do so and having co-located, high-frequency skin and NSL sensors sampled along a moving platform that deliberately re-crosses the same feature. Satellite climatologies have been used to detect submesoscale fronts (Belkin et al., 2009; Belkin and O’Reilly, 2009), but they inherently treat the retrieved skin temperature as a proxy for the underlying frontal structure. This assumption is convenient for basin-scale mapping but has seldom been validated at the along-track, sub-kilometre resolution where submesoscale processes generate fronts. Ground-truthing this assumption requires a platform that can both sample the skin layer directly, at the few-hundreds of micrometres scale the skin layer occupies. The platform must also repeatedly transect an observed front rather than sampling it once in passing, a capability that most moored, drifting, or single-pass shipboard sensors cannot provide.

However, the air–sea interface is continuously subjected to latent evaporation, solar radiation heating, wind-driven shear stress, and organic surfactant activity (Donlon et al., 2002; Fairall et al., 1996; Wurl et al., 2011). These boundary-layer mechanisms establish a distinct thermal



70 and haline skin effect. Whether submesoscale horizontal density fronts in the NSL undergo
71 attenuation, structural deformation, or enhancement at the surface skin layer remains a
72 fundamental open question in physical oceanography. Furthermore, assessing whether these
73 shallow coastal density fronts are primarily governed by thermal or haline drivers, and
74 identifying the meteorological thresholds that control front formation is vital for refining
75 submesoscale ocean models and improving satellite calibration algorithms. A front that is
76 haline-driven carries no direct thermal signature and is therefore invisible to infrared satellite
77 retrieval regardless of how well that retrieval is calibrated against skin temperature. A front that
78 is thermally driven but strongly attenuated between the NSL and the skin will be under-
79 represented in satellite imagery relative to its true sub-surface intensity. Both possibilities have
80 direct consequences for how confidently a satellite-derived frontal product can be used to infer
81 sub-surface density structure, biological aggregation, or air–sea flux anomalies in a shallow
82 shelf sea. Similarly, if front formation and persistence are gated by identifiable, non-linear
83 thresholds in wind speed, tidal phase, or air–sea temperature difference, those thresholds
84 become directly actionable. They specify the meteorological conditions under which a satellite-
85 based or model-based frontal product should be trusted, and the conditions under which it
86 should not.

87 Our study utilizes high-resolution along-track data collected by the Autonomous Surface
88 Vehicle (ASV) HALOBATES (Wurl et al., 2024) across five research cruises over 54 cruise-
89 days in the German Bight to understand the density gradients of the sea surface, its dynamics
90 and to identify forcing factors. By analysing co-located skin and NSL measurements at 1-
91 minute temporal intervals and approximately 150 m along-track distance, we address three
92 primary research questions. First, how do horizontal gradients of density at the skin layer scale
93 with and diverge from co-located gradients within the NSL? Additionally, what are the
94 dominant thermohaline drivers governing horizontal density, salinity and temperature gradients
95 in the German Bight, and how are these regimes modulated across distinct tidal phases? Finally,
96 which meteorological and hydrographic variables exert the strongest predictive control on skin
97 layer density gradient, and what non-linear environmental thresholds govern front persistence
98 and dissipation?

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2. Methods

2.1 Study Area

Data were collected over 54 cruise-days during five research cruises in the German Bight (Figure 1b), south-eastern North Sea (Figure 1a). The study area is a shallow, tidally energetic shelf sea with an established history of tidal-front and FerryBox-based surface monitoring (Grayek et al., 2011). The expeditions were conducted on the research vessel (RV) Heincke with data from the cruises HE598 in May 2022 (Gassen et al., 2024a), HE609 in October 2022 (Gassen et al., 2025), HE614 in March 2023 (Gassen et al., 2024b), HE626 in July/August 2023 (Ayim et al., 2025a), HE644 in July 2024 (Ayim et al., 2026b). All data are available on the PANGAEA data repository with the references above for the individual cruises.

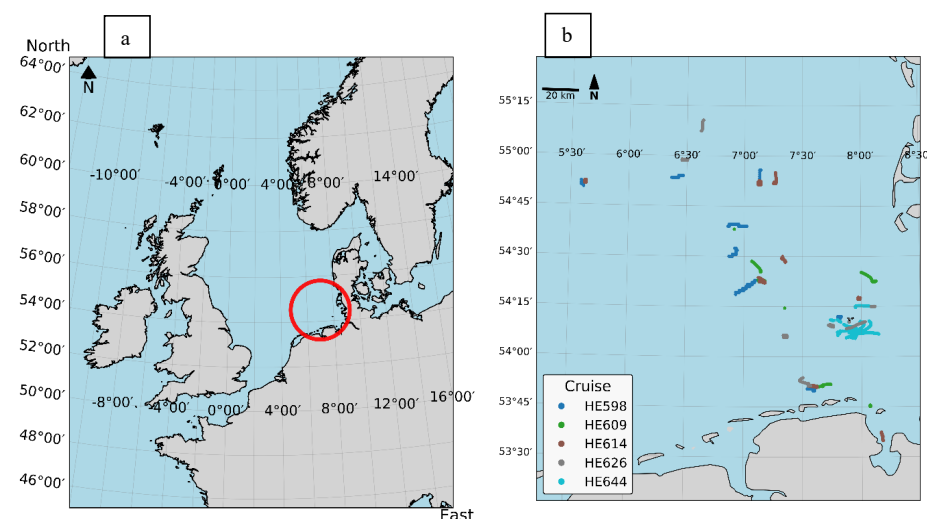


Figure 1: Study Area with stations across the 5 cruises in the German Bight. (a) is the map of the North Sea with the red circle highlighting the study area. (b) is the zoomed in study area showing the locations of the five research cruises.

2.2 Device and Data

HALOBATES is equipped with a rotating glass disc assembly that allows the sampling of the skin layer ($78 \pm 8.8 \mu\text{m}$) of the sea surface due to surface tension and adhesion (Ribas-Ribas et al., 2017; Shinki et al., 2012; Wurl et al., 2024). The glass disc assembly is covered with a tinted acrylic glass hood to avoid evaporative cooling at stronger winds as some biases without hood were observed during validation with an infrared radiometer (see figure 2 in (Wurl et al., 2019)). HALOBATES has a ladder at the bow with temperature and pressure sensors and sampling tubes for the flow-through conductivity-temperature-depth (CTD) sensors mounted on the

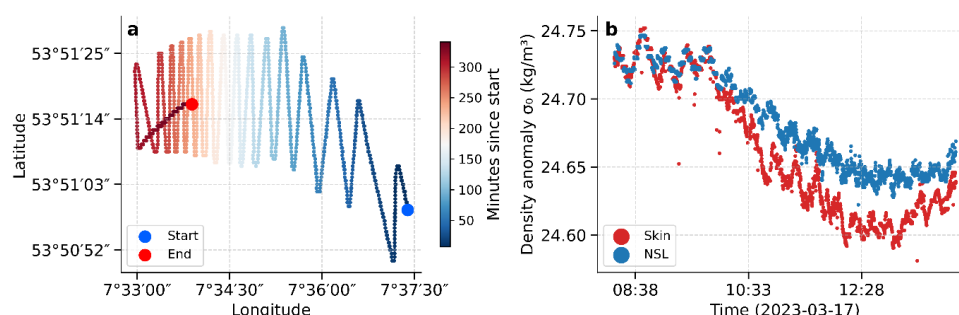


129 ASV. The maximum depth sampled was 1 meter. The meteorological sensors are mounted on
130 the bridge of the ASV at a height of 2 m measuring wind speed and direction, atmospheric
131 pressure, relative humidity and air temperature. Other meteorological sensors were mounted on
132 the ship which was averagely at a 500 m distance from the ASV, this includes net shortwave
133 and longwave radiation, precipitation intensity and amount.

134 The longitude, latitude and time were logged continuously. The raw dataset comprised 101,645
135 timestamped data points at 10 second resolution. Data spans various events such as front, rain,
136 slicks and days without any apparent surface phenomena. The CTDs (OS310, Idrounat, Italy)
137 had temperature and conductivity accuracy of ± 0.0015 °C and ± 0.003 mS/cm, respectively. A
138 detailed description of the design and capabilities of HALOBATES are reported in the technical
139 paper of the ASV (Wurl et al., 2024). On a typical front day, we observed the phenomenon by
140 going back and forth across the ocean feature (Figure 2). The TPXO-OSU model (Egbert et al.,
141 1994; Egbert and Erofeeva, 2002) was then used to obtain tidal record based on the date, time
142 and geolocation details of HALOBATES.

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146 *Figure 2: A typical track (a) and Density anomaly (b) as we moved back and forth across an observed front.*

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2.3 Data Quality

150 Raw samples were passed through a quality control (QC) scheme, with valid data row counts
151 logged at every stage. First, we took out October 12, 2022 as the skin layer salinity had long
152 periods of outliers and a few of the other sensors having no data. We also removed all rows and
153 time points where any of the four-core variables (skin and NSL temperature and salinity) had a
154 missing value. Due to the knowledge of the region's oceanography, a salinity minimum of 28
155 g kg⁻¹ was set to exclude false data by bubbles in the conductivity cell. Due to the focus on
156 gradients, we filtered out sensor spikes between consecutive measurements. To do so, we have
157 used a threshold of 0.03 °C and 0.03 g/kg for the temperature and salinity respectively. That is,



158 we have excluded values that exceed the previous value by this threshold within the 10 second
159 data interval. This threshold was chosen after testing different thresholds and evaluating the
160 effect on the whole dataset. Data were then resampled to a strict 1-minute grid, computed
161 independently within each day. The resampling used a first principle by taking the first value
162 of the minute and not the mean which will average the gradient intensities. Along-track distance
163 was computed via the haversine formula between consecutive 1-minute positions; a data row
164 was retained for gradient analysis if the time gap was ≤ 70 s and the along-track distance fell
165 between 0 and 360 m. The distance limits used is based on the maximum distance we expect
166 the ASV to cover in one minute.

167 After QC and resampling, 73,617 data points remained, and 14,446 one-minute data points were
168 produced, of which 13,029 (90.2%) were valid because the time and distance difference fulfilled
169 the criteria. Based on examination of the raw data and the pre-processed data, a targeted salinity-
170 spike guard was applied after gradient-intensity computation. Raw salinity gradient intensities
171 exceeding 0.2 g/kg, that is approximately 7 times the salinity front-detection threshold, were
172 considered implausible as a true salinity gradient between consecutive 1-minute measurements.
173 In addition, in cases where the skin and NSL salinity gradient intensity disagreed by more than
174 0.11 g/kg were treated as sensor artefacts. This threshold was chosen through visual inspection
175 of the time-series. Such events occurred in isolation without a smooth transition of the gradients
176 during such time periods and thus were treated as outliers. The salinity-spike guard and the
177 threshold for the difference between skin and NSL salinity flagged 101 data points (0.70%).

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180 **2.4 Density and Front Computations**

181 Density was computed from temperature and conductivity via the TEOS-10 gsw toolbox
182 (McDougall and Barker, 2011; UNESCO/IOC et al., 2010). Gradient intensity was defined as
183 the absolute difference between consecutive 1-minute measurements. We intentionally retained
184 this raw temporal difference rather than calculating a distance-normalized gradient, as linear
185 extrapolation across a full kilometre is physically invalid for fine submesoscale features
186 (McWilliams, 2016). Each step in our data corresponds to 1 minute and approximately 150 m.
187 We compute the density ratio R_p (Eq. 1) to assess the thermal and haline contribution of the
188 observed density gradients. The density ratio is based on the thermal expansion α and the β
189 saline contraction coefficients (Schmitt, 1981). To distinguish strong dominance from near-
190 compensated conditions, stratification was operationally categorized into haline-dominated (R_p
191 < 0.67), thermally-dominated ($R_p > 1.5$), and mixed ($0.67 \leq R_p \leq 1.5$) regimes. These represents



192 a minimum 1.5-fold (or 60%) relative contribution of one variable over the other (Ruddick,
193 1983; Rudnick and Martin, 2002).

194

$$R\rho = \frac{\alpha|\Delta T|}{\beta|\Delta S|} \quad (1)$$

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196 In this study, a front is defined by the density gradient-intensity threshold ($\geq 0.02 \text{ kg m}^{-3}$) as a
197 fixed, reproducible cut-off. This threshold was based on a number of factors. First, we primarily
198 looked at the literature and found density gradient of $0.2 - 2.0 \text{ kg m}^{-3}$ per km in San Diego (Wu
199 et al., 2021). When this is scaled down to a 100 m distance, this gives density gradients of 0.02
200 $- 0.2 \text{ kg m}^{-3}$ per 100 m. We then visually inspected the days where we could clearly see a front
201 and found this threshold to be true for our dataset. Also, when we computed the percentiles,
202 this value sits around the 90th percentile. Density is the dynamically relevant quantity as it drives
203 the convergence or divergence associated with a front (Capet et al., 2008; McWilliams, 2016).
204 This makes it the appropriate primary detector, particularly given the well-documented
205 tendency of temperature and salinity gradients to compensate one another's effect on density
206 rather than simply add (Rudnick and Ferrari, 1999). Given a detected density front, temperature
207 ($\geq 0.03 \text{ }^{\circ}\text{C}$) and salinity ($\geq 0.02 \text{ g kg}^{-1}$) gradient conditions are then checked for co-
208 occurrence and classified as T and S (both conditions independently met), T only, S only, or
209 neither fulfilled the criteria. This detection approach is conceptually distinct from the widely
210 used gradient and histogram-based algorithms developed for satellite SST/chlorophyll imagery
211 (Belkin et al., 2009; Belkin and O'Reilly, 2009), but adapted here to a 1-D along-track
212 shipboard record rather than a 2-D image. Tidal phase was classified from the sign of the within-
213 day first difference of tidal height: positive indicates the rising tide, negative indicates falling
214 tide, and a small-magnitude of the tidal height difference $\Delta h < 1\text{e}^{-4} \text{ m}$ indicates a tidal slack.
215 The rising and falling tide are thus referred to as Flood and Ebb tide respectively.

216

217 **2.5 Machine Learning and SHapley Additive exPlanations attribution**

218 To explore meteorological and hydrographic predictor variables as best predictors for gradient
219 intensity, a random-forest regressor (Breiman, 2001) was trained to predict continuous skin
220 density gradient intensity. Two sets of variables were compared: the first set includes tidal
221 height, wind speed and direction, air-sea temperature difference, relative humidity, net
222 shortwave and longwave radiation, atmospheric pressure, skin temperature, skin salinity, and a



cyclic time-of-day encoding. The second set excludes shortwave and longwave radiation. Radiation fields carry substantially more missing values than the other variables, and an initial full-feature run found radiation had the lowest SHAP importance of any variable, so the second set is treated as the primary model, trading a small amount of potential signal for a much larger usable sample. Models were evaluated with a 75/25 train-test split, with performance assessed directly on the holdout test set using the coefficient of determination (R^2), root-mean-square error (RMSE) bias. SHAP (SHapley Additive exPlanations; (Lundberg and Lee, 2017)) values from a tree explainer were used to rank feature importance and to visualise partial dependence for the top predictors. We tested group differences across tidal phase with the Kruskal-Wallis test. The paired skin against NSL comparison were tested with the Wilcoxon signed-rank since skin and NSL gradient intensity are measured on the same along-track step and are therefore not independent data.

3. Results

For context, density gradient intensity is strongly heavy-tailed as shown in Figure 3. The median is 0.004 kg m^{-3} , but the 95th, 99th, and 99.9th percentiles are 0.028, 0.057, and 0.106 kg m^{-3} , respectively. The most intense data points (0.1% of all datapoints) are roughly 25 times the median. Most of the along-track record are below the fixed front-detection threshold of 0.02 kg m^{-3} (dashed line in Figure 3) highlighting that crossing into a true front is the exception rather than the rule, consistent with submesoscale straining events that are localised in space and time.

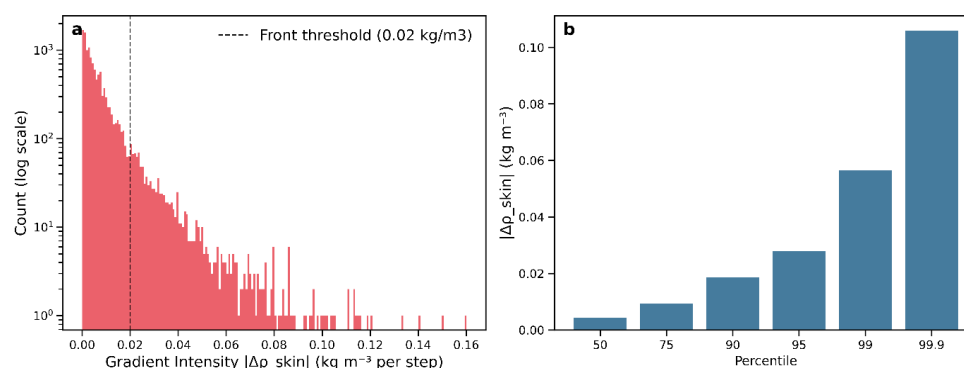


Figure 3:(a) Distribution of skin density gradient intensity (log-scale count), with the fixed front-detection threshold marked. (b) Percentiles of the same distribution, illustrating the heavy tail.



3.1 Difference between gradients at the skin and NSL

Using density as the primary front-detection variable, we identified 1192 skin fronts (8.3% of valid steps) and 1035 NSL fronts (7.2%), distributed across all five campaigns. To evaluate how the density gradients, differ from each other, we evaluated matched step-by-step observations by comparing the skin and NSL layers (Figure 4). We treat NSL gradient intensity as the independent (predictor) variable and skin gradient intensity as the dependent (response) variable, consistent with our physical framing of the skin layer as an attenuated surface expression of the underlying bulk signal, and quantify this relationship using Ordinary Least Squares (OLS) linear regression (Figure 4a). The regression slope ($m = 0.409 \pm 0.008$) falls well below the 1:1 line of equal intensity ($m = 1.0$), demonstrating substantial attenuation of the skin front compared to the sub-surface fronts at 1 meter depth. On average, a 1.00 kg m^{-3} increase in gradient intensity at the NSL corresponds to an increase of approximately 0.41 kg m^{-3} at the skin layer. The regression's positive intercept ($c = 0.005 \text{ kg m}^{-3}$) highlights a persistent background baseline level of spatial variability maintained at the skin layer even when NSL waters are completely homogeneous, while its low coefficient of determination ($R^2 = 0.179$) indicates that NSL dynamics explain less than 18% of the variance in skin gradient intensity. The Bland-Altman analysis reveals that this vertical coupling is regime dependent (Figure 4b). During times of low mean gradient $< 0.02 \text{ kg m}^{-3}$, paired differences are close to zero, indicating that skin and NSL layers track each other closely in weak gradient conditions. During a strong frontal regime (mean gradient $> 0.04 \text{ kg m}^{-3}$), the two layers diverge sharply with differences between layers ranging from -0.14 kg m^{-3} to 0.11 kg m^{-3} .

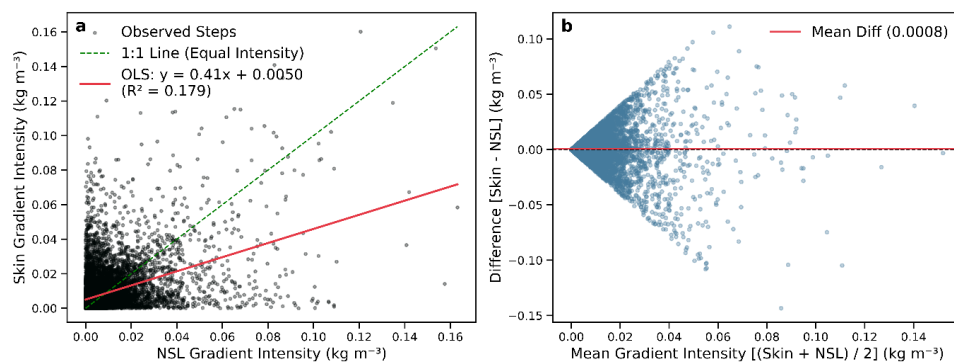


Figure 4: Scaling and alignment between ocean skin and NSL density gradient intensities. (a) Scatter plot and Ordinary Least Squares (OLS) linear regression. The solid red line demonstrates gradient attenuation relative to the 1:1 reference line (black dashed). (b) Bland-Altman plot of paired differences versus mean gradient intensity. Mean paired difference is shown by the red horizontal line.



3.2 Tidal Phase influence on gradients

Gradient intensities were compared across flood ($n = 6,836$), ebb ($n = 6,186$), and phases of slack water ($n = 183$), representing 13,205 one-minute observations (Table 1). Kruskal–Wallis tests revealed significant tidal effects on all skin layer gradients: temperature ($H(2) = 24.13$), salinity ($H(2) = 52.49$), and density ($H(2) = 27.25$) with all p -values $\ll 0.001$. Salinity exhibited the strongest tidal dependence, indicating that tidal modulation primarily influenced horizontal salinity structure (Figure 5). Flood tide produced the largest mean skin gradients for temperature ($0.020 \pm 0.032\text{ }^{\circ}\text{C}$), salinity ($0.008 \pm 0.012\text{ g kg}^{-1}$), and density ($0.008 \pm 0.012\text{ kg m}^{-3}$). Skin density gradients were approximately 22% larger than those at NSL (0.007 kg m^{-3}) during flood tide. Skin and NSL density gradients were identical with a value of 0.007 kg m^{-3} during ebb tide. Density fronts ($|\Delta\rho| \geq 0.02\text{ kg m}^{-3}$) occurred most frequently during flood tide (10.4%), compared with ebb (7.6%) and phases of slack water (6.6%). The vertical temperature difference ($\Delta T = \text{skin} - \text{NSL}$) remained positive throughout all tidal phases, with the strongest warming during flood tide ($0.079 \pm 0.273\text{ }^{\circ}\text{C}$). Dunn's post-hoc test showed that flood tide differed significantly from both ebb tide ($p < 0.001$) and phases of slack water ($p = 0.028$), but ebb and slack water were not significantly different ($p = 0.395$). It indicates that the principal tidal contrast was associated with enhanced density gradients during flood tides.

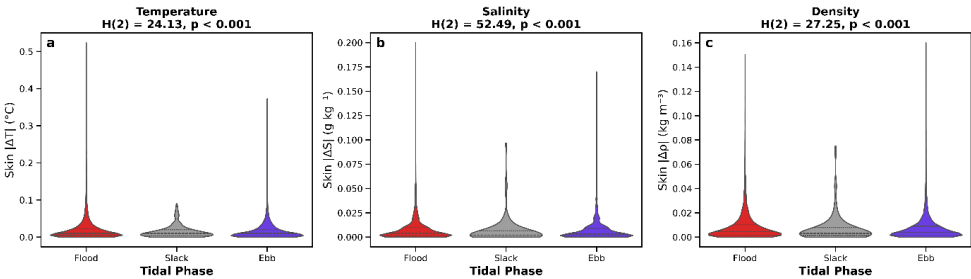


Figure 5: Violin plots of gradients in absolute values of skin (a) temperature, (b) salinity, and (c) density during flood, slack water, and ebb tidal phases ($N = 13,205$). Dashed and dotted lines denote the median and interquartile range. Kruskal–Wallis statistics demonstrate significant tidal dependence for all variables, with salinity exhibiting the strongest response. Summary statistics are provided in Table 1.



Table 1: Skin and NSL gradient properties across tidal phases. Values are reported as mean $\pm$ standard deviation.

Tidal phase	Skin $ \Delta T $ ($^{\circ}\text{C}$)	Skin $ \Delta S $ (g kg^{-1})	Skin $ \Delta \rho $ (kg m^{-3})	NSL $ \Delta \rho $ (kg m^{-3})	Front rate (Skin, %)	Front rate (NSL, %)	n
Flood	0.020 $\pm$ 0.032	0.008 $\pm$ 0.012	0.009 $\pm$ 0.012	0.007 $\pm$ 0.011	10.4	7.7	6,836
Slack	0.015 $\pm$ 0.017	0.006 $\pm$ 0.013	0.007 $\pm$ 0.010	0.005 $\pm$ 0.009	6.6	3.8	183
Ebb	0.016 $\pm$ 0.022	0.007 $\pm$ 0.012	0.007 $\pm$ 0.011	0.007 $\pm$ 0.010	7.6	8.1	6,186

3.3 Gradients and Environmental Factors

We assessed the thermal and haline contributions to the observed density gradients. We determine the percentage of density fronts occurring at the same time as salinity front, temperature front or neither, that is, cases where the density front was driven by sub-threshold contributions from both variables jointly. Detection thresholds for the co-occurrence of temperature and salinity fronts (Figure 6) reveals a clear asymmetry: at the skin layer, 44.1% of density fronts are salinity-driven, 33.7% are temperature-driven, 21.1% show both, and only 1.0% show neither individually crossing its threshold. At the NSL, the separation is closer: 39.9% salinity-only, 31.8% temperature-only, 27.8% both. Salinity gradients are therefore the dominant contributor to density fronts at the skin compared to the NSL in the North Sea.

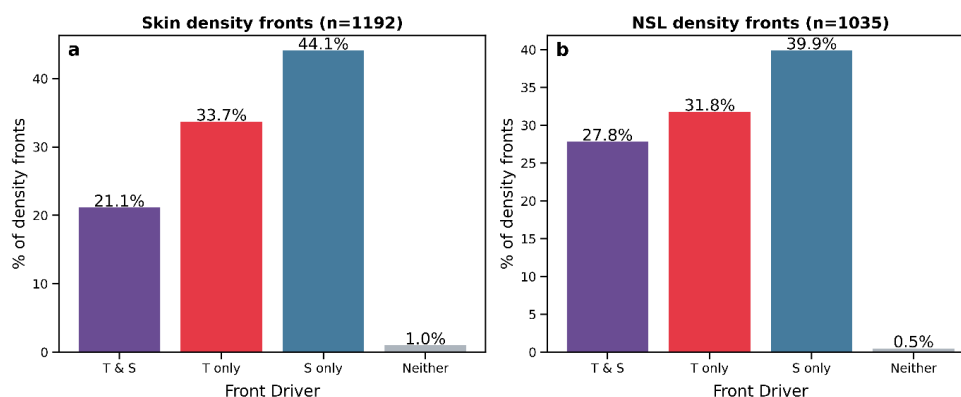


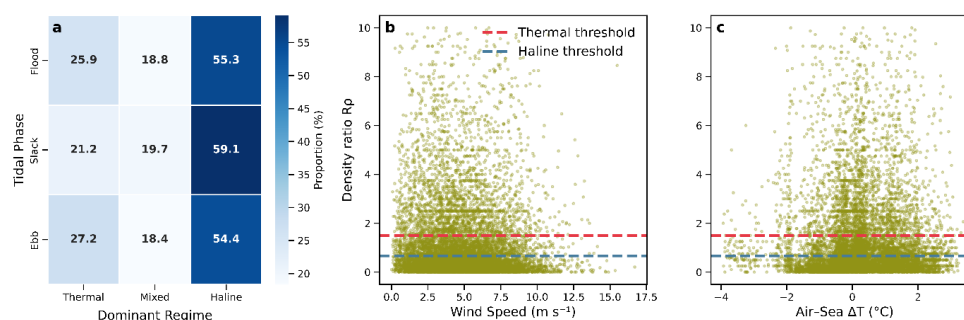
Figure 6: Temperature/salinity co-occurrence at detected density fronts, skin (a) vs. NSL (b).

This is corroborated by a continuous thermal/haline dominance-ratio classification using the same gradient-intensity terms (Figure 7a). Haline dominance represents the largest proportion of observations (approximately 55–60%) across flood and ebb tides and slack waters, followed



by thermal (20–25%) and mixed (15–20%) regimes. The distribution shows minor changes with tidal phases, indicating that the relative influence of temperature and salinity gradients is largely independent of tidal state. Thermally dominated fronts occur mainly when winds are $< 8 \text{ m s}^{-1}$, where the density ratio exhibits a broad range of values (Figure 7b). As wind speed increases, observations shift toward haline dominance, consistent with stronger turbulent mixing reducing horizontal thermal gradients while salinity gradients remain comparatively persistent. Most observations under a cooler atmosphere, that is air temperature is below sea surface temperature, fall within the haline and mixed regimes (Figure 7c). Thermal dominance occurs primarily when the atmosphere is slightly warmer than the sea surface (by approximately $0 - 2.5 \text{ }^{\circ}\text{C}$), consistent with conditions favouring stable oceanic near-surface thermal stratification. Overall, salinity is the dominant control on density gradients (Figure 7a), whereas thermal dominance occurs primarily at wind speed $< 8 \text{ m s}^{-1}$ (Figure 7b) and positive air-sea temperature differences (Figure 7c).

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339 *Figure 7: Thermal vs. haline dominance regime: (a) proportion by tidal phase, (b) regime vs. wind speed, (c) regime vs. air-*
340 *sea temperature difference, coloured by tidal height*

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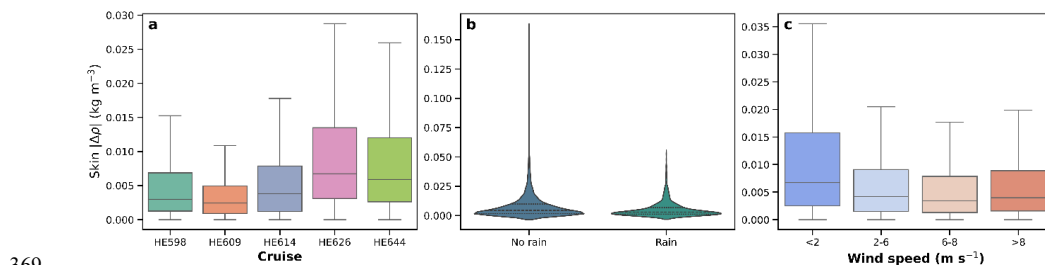
We binned wind speed into four ranges: $< 2 \text{ m s}^{-1}$ (calm, slick-favourable), $2-6 \text{ m s}^{-1}$ (transitional), $6-8 \text{ m s}^{-1}$ (approaching the wave-breaking threshold), and $> 8 \text{ m s}^{-1}$ (mixing-dominated). This partitioning combines field observations showing substantially reduced surface roughness in surfactant-rich coastal waters relative to the open ocean at wind speeds below 4 m s^{-1} (Laxague et al., 2024), consistent with visual observations during our own cruises of the sea surface losing its characteristic smoothness once wind speeds exceed $\sim 2 \text{ m s}^{-1}$, with threshold behaviour reported in laboratory wind-wave tank experiments (Engel et al., 2026), and with established thresholds for surfactant slick persistence (Romano, 1996; Wurl et al., 2011). The lowest wind speed bin ($< 2 \text{ m s}^{-1}$) shows a higher mean skin $|\Delta\rho|$ of 0.013 kg m^{-3} (Figure 8c), more variability and a front-detection rate of 19.4% which is roughly 2.3–3.4 times



352 higher than the 2–6 m s⁻¹ (8.5%), 6–8 m s⁻¹ (5.7%), and >8 m s⁻¹ (7.6%) bins. Rainfall was
353 flagged in 7.7% of valid one-minute steps (993 of 12,928). Skin density gradient intensity was
354 significantly lower during rain than during dry conditions (Mann-Whitney U = 4,886,105, p <
355 0.001). The median skin $|\Delta\rho|$ fell from 0.005 kg m⁻³ (no rain) to 0.003 kg m⁻³ (rain) (Figure 8b),
356 and the skin front-detection rate fell from 9.65% (no rain) to 4.03% (rain). This pattern held
357 within every tidal phase examined, with rain-conditional medians lower than dry-conditional
358 medians.

359 We also investigated whether the gradient intensities varied across the seasons, that is the
360 cruises (Figure 8a). Skin density gradient intensity differs sharply and significantly across
361 cruises (Kruskal–Wallis H = 969.4, p << 0.001). A Dunn's post-hoc test showed that four of the
362 five cruises are statistically distinguishable from nearly every other cruise with all pairwise p <
363 0.005 except between HE598 (May 2022) and HE614 (October 2022) with p = 0.13. The two
364 summer cruises, HE626 (July/August 2023) and HE644 (July 2024) showed the highest median
365 and mean skin gradient intensity (0.007 and 0.006 kg m⁻³ median, respectively) and the highest
366 front-detection rates (14.7% and 11.9%), while the autumn cruise HE609 (October 2022) shows
367 the lowest of both (0.002 kg m⁻³ median, 4.4% front-detection rate).

368



369

370 *Figure 8: Distribution of density gradients across the skin layer. (a) Boxplot of Skin $|\Delta\rho|$ across five research cruises (HE598*
371 *[May], HE609 [October], HE614 [March], and HE626, HE644 [July/August]). (b) Violin plot showing the probability density*
372 *distribution of Skin $|\Delta\rho|$ during dry (No rain) versus precipitating (Rain) conditions. (c) Boxplot of Skin $|\Delta\rho|$ grouped by*
373 *wind speed regimes. For the boxplots, horizontal lines denote the median, boxes span the interquartile range (IQR), and whiskers*
374 *represent the data spread. In the violin plots, dotted lines inside the distributions indicate the median and quartiles.*

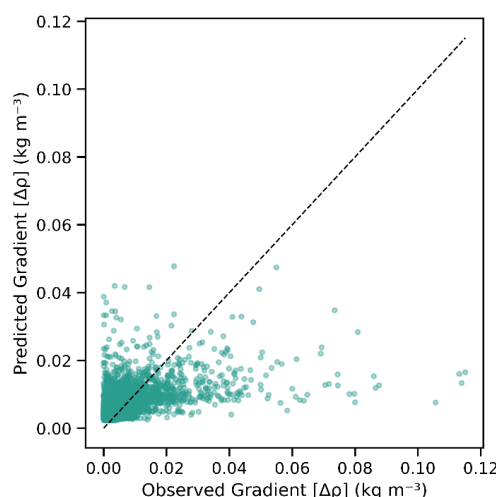
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376 The influence of radiation variables on gradient intensity prediction was initially evaluated to
377 capture potential diurnal warming effects. However, by including radiation parameters the data
378 set is substantially reduced from 12,367 to 4,003 observations due to incomplete data. Radiation
379 inclusion did not improve model performance. The full model achieved a low test $R^2 = 0.082$
380 compared with the model excluding radiation of $R^2 = 0.188$. Therefore, the no-radiation model



381 was selected as the primary model due to the larger number of observations ($N = 12,367$) and
382 improved predictive performance.

383 The model performance showed a low overall bias ($+0.0001 \text{ kg m}^{-3}$) and an RMSE of 0.0097
384 kg m^{-3} . Interpretation of these metrics must consider the strongly right-skewed distribution of
385 observed skin layer density gradients (Figure 3a), which exhibited a median of 0.0040 kg m^{-3}
386 and a mean of 0.0080 kg m^{-3} . The difference between mean and median reflects the influence
387 of relatively rare, high-intensity gradient events on the overall distribution. The model's mean
388 absolute error ($\text{MAE} = 0.0058 \text{ kg m}^{-3}$) remained below the observed mean gradient intensity,
389 indicating that typical prediction errors were comparable to the background variability of the
390 dataset. However, the RMSE exceeded both the observed mean (121%) and median (241%)
391 since squared-error metrics give disproportionate weight to large residuals associated with
392 extreme gradient events. Observed versus predicted comparisons (Figure 9) showed good
393 agreement along the 1:1 relationship under low-to-moderate gradient conditions ($\Delta\rho < 0.02 \text{ kg}$
394 m^{-3}), though scatter around the 1:1 line indicates that individual predictions in this regime can
395 deviate substantially even where aggregate error metrics are low, reflecting irreducible case-
396 by-case variability rather than systematic bias. However, the model systematically
397 underestimated the largest observed gradients exceeding 0.04 kg m^{-3} by more than 50%. This
398 limitation reflects the difficulty of predicting transient, localized but extreme skin layer density
399 gradient using atmospheric and hydrographic predictors.



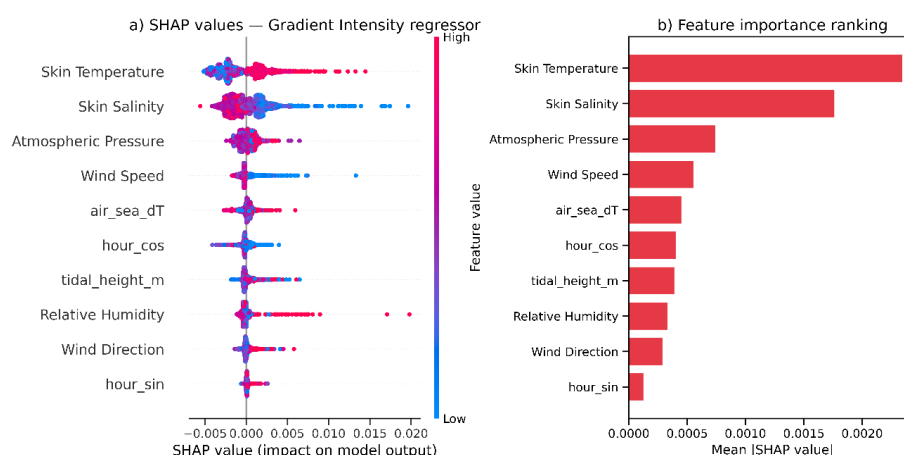
400

401 *Figure 9: Predicted vs. observed skin gradient intensity for the regressor model ($N = 3,092$ test samples). The black dashed*
402 *line represents the 1:1 line of perfect agreement.*



SHAP analysis identified surface hydrographic variables as the dominant control of density gradient intensity (Figure 10). Skin temperature was the most influential predictor (mean $|\text{SHAP}| \approx 0.0023 \text{ kg m}^{-3}$), followed by skin salinity ($\approx 0.0018 \text{ kg m}^{-3}$). Atmospheric pressure and wind speed were the most important atmospheric predictors. The analysis revealed that high skin temperatures generally increased predicted gradient intensity, while low salinity values produced strong positive SHAP contributions, indicating enhanced gradients under fresher conditions (Figure 10). In contrast, increasing wind speed generally reduced predicted gradient intensity, consistent with increased mixing and reduced persistence of sharp surface gradients. Partial dependence analysis (Figure 11) revealed distinct non-linear relationships between predictors and predicted gradient intensity. Skin temperature showed relatively weak effects between 5 and 17 °C, followed by a rapid increase in predicted gradient intensity above 17 °C and exceeding 0.022 kg m^{-3} at 22 °C. Skin salinity showed the strongest response under low-salinity conditions between 29 and 30 g kg^{-1} in which predicted gradients levelled up to approximately 0.033 kg m^{-3} . Density gradients declined toward a baseline of $\sim 0.010 \text{ kg m}^{-3}$ at higher salinities ($>33 \text{ g kg}^{-1}$). For atmospheric drivers, wind speed showed the clearest threshold behaviour. Gradient intensity was highest under near-zero winds ($>0.016 \text{ kg m}^{-3}$) and rapidly decreased between 2 and 4 m s^{-1} , reaching a stable minimum ($\sim 0.008 \text{ kg m}^{-3}$) above 5 m s^{-1} . Atmospheric pressure exhibited a weaker response, with elevated gradients under low-pressure conditions ($<1005 \text{ hPa}$) and minimum values between approximately 1012 and 1018 hPa.

422

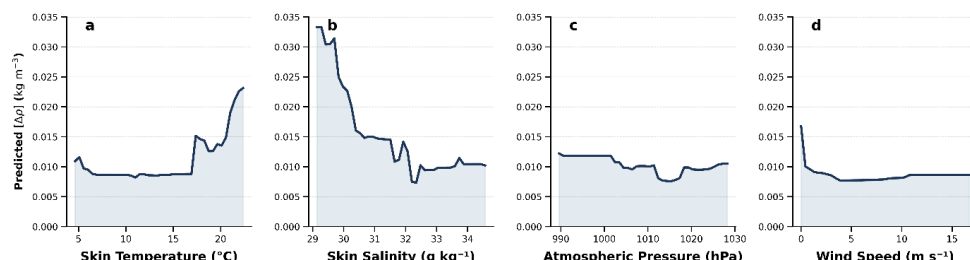


423

424 *Figure 10: Feature importance and directional attribution of environmental predictors governing density gradient intensity*
425 *using SHAP (SHapley Additive exPlanations) values. (a) SHAP summary plot showing feature value impact on model*
426 *predictions. Points are coloured by feature magnitude (red = high, blue = low). (b) Global feature importance ranking sorted*
427 *by mean absolute SHAP value.*



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Figure 11: Partial dependence plots displaying the marginal, non-linear effects of the top four environmental drivers on predicted density gradient intensity: Skin Temperature, Skin Salinity, Atmospheric Pressure, and Wind Speed.

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4. Discussion

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4.1 Vertical Decoupling and Gradient Attenuation

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Our empirical findings of attenuated and weakly-coupled surface gradients that diverge under strong frontal forcing demonstrate a partial decoupling and significant attenuation of horizontal oceanographic gradients across the upper metre of the ocean column. Our analysis on paired linear regression between skin and NSL density gradient intensities (Figure 4a) indicates that, on average, a sub-surface horizontal density gradient is attenuated by roughly 60% at the sea surface skin layer. Furthermore, sub-surface gradient dynamics account for less than 18% of the total variance observed in surface skin gradients, indicating that the ocean skin layer operates as a dynamic, and only partially decoupled physical regime rather than a passive and separated top boundary. The Bland–Altman analysis confirms that this vertical alignment is strongly regime-dependent. In conditions of small horizontal gradients ($|\Delta\rho| < 0.02 \text{ kg m}^{-3}$), skin and NSL gradient intensities maintain tight equilibrium, tracking each other with near-zero mean bias. However, at stronger frontal regime ($|\Delta\rho| > 0.04 \text{ kg m}^{-3}$ per step), the two layers diverge sharply, with inter-layer differences spreading widely. This divergence during intense frontogenesis is explained by upper-ocean fluid dynamics and thermodynamic boundary processes (McWilliams, 2016; Thomas et al., 2008). For example, submesoscale fronts are characterized by convergent surface velocity fields that force opposing water masses together, driving downward subduction and sharp lateral shear in the bulk upper mixed layer (Capet et al., 2008; D’Asaro et al., 2011). At the skin layer, however, this lateral convergence competes directly with vertical air–sea heat fluxes, evaporation, and surfactant-induced Marangoni surface tension gradients (Donlon et al., 2002; Ribas-Ribas et al., 2017). These micro-surface processes damp localized density contrasts, resulting in a smoothed spatial profile relative to the underlying flow.



457 The positive regression intercept ($c = 0.005 \text{ kg m}^{-3}$, Figure 4a) underscores another key physical
458 behaviour. When bulk waters at 1 m depth are spatially homogeneous ($\Delta\rho \text{ NSL} = 0$), a baseline
459 level of horizontal spatial variability is maintained across the skin layer. This micro-spatial
460 baseline might be sustained by localized, patchy atmospheric exchanges, micro-scale wind
461 stress variations, and organic surfactant slicks that alter local evaporation rates independently
462 of sub-surface hydrodynamics (Ribas-Ribas et al., 2017). The attenuation is related to the higher
463 absolute front-detection rate at the skin (8.3%) than at the NSL (7.2%). The regression and
464 attenuation describe different aspects of the same regime-dependent coupling. The OLS slope
465 characterises the typical relationship between the layers, which is dominated by the very large
466 number of weak, near-zero steps where skin and NSL track closely as seen in Figure 4a. The
467 front-detection rate, by contrast, is a threshold-crossing statistic driven by the tail of the
468 distribution, where the positive intercept's baseline skin variability is large enough on its own
469 to push a skin gradient over the fixed 0.02 kg m^{-3} threshold even when the co-located NSL
470 gradient does not. That is, even though the skin layer understates the intensity of most gradients,
471 its constant micro-fluctuations under low sea states force it over low fixed thresholds more
472 frequently overall.

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474 **4.2 Tidal Hydrodynamics and Thermohaline Dominance Regimes**

475 Tidal forcing acts as a key hydrodynamic control on horizontal oceanographic structure across
476 the German Bight (Otto et al., 1990; Schrum, 1997). Non-parametric Kruskal–Wallis tests
477 confirmed significant variations across tidal phases for skin temperature, salinity, and density.
478 Salinity exhibited the strongest tidal sensitivity, underscoring that tidal modulation in this
479 shallow shelf sea primarily reorganizes freshwater boundaries rather than thermal fields. Flood
480 tides yielded the sharpest horizontal gradients and the highest density front-detection frequency.
481 During flood tide, denser offshore North Sea water flows landward, pushing directly against
482 buoyant estuarine plumes originating from the Elbe, Weser, and Ems rivers (Grayek et al., 2011;
483 Hagen et al., 2022). This opposing flow sharpens horizontal density boundaries via tidal
484 straining and advective frontogenesis. Conversely, during ebb tides, buoyant estuarine
485 freshwater flows offshore over denser open-sea water as an unconstrained surface plume,
486 leading to broader, more diffuse horizontal transitions and lower front detection rates (Horner-
487 Devine et al., 2015).

488 Evaluating thermohaline driver co-occurrence at detected density fronts revealed a strong
489 asymmetry (Figure 6). At the skin interface, majority of the density fronts were driven by



490 salinity only, followed by temperature and the combined thermohaline gradients. Haline
491 dominance consistently accounted for 55–60% of all valid observations across all tidal states,
492 whereas thermal dominance comprised 20–25% and mixed regimes 15–20%. These empirical
493 observations align with mixed-layer thermohaline compensation principles (Ferrari and
494 Rudnick, 2000; Rudnick and Ferrari, 1999). In shallow shelf seas, horizontal temperature
495 gradients are continually damped by air–sea heat exchange through latent, sensible, and
496 radiative heat fluxes (Donlon et al., 2002), whereas horizontal salinity gradients lack a rapid
497 restoring atmospheric flux (Ferrari and Rudnick, 2000), allowing fresh estuarine runoff to
498 maintain persistent spatial density contrasts.

499 **4.3 Environmental Thresholds and Predictive Machine Learning**

500 Skin density gradient intensity is governed by a set of identifiable, non-linear environmental
501 thresholds spanning wind speed, season, and rainfall, which together set the physical
502 background for the predictive machine-learning. Wind speed exerts the clearest threshold
503 behaviour. Binning wind speed showed that skin density gradient intensity and front-detection
504 rate are both far higher under calm conditions ($< 2 \text{ m s}^{-1}$) and decay rapidly as wind forcing
505 increases. This is consistent with field observations showing that surface roughness is
506 substantially suppressed in surfactant-rich waters below 4 m s^{-1} (Laxague et al., 2024), allowing
507 persistent, sharp-edged surfactant patches and their associated density signatures to survive
508 rather than being dispersed by wind-driven turbulence. The partitioning of the wind speeds
509 corresponds to independently documented physical thresholds for surfactant behaviour at the
510 sea surface. Field observations demonstrate that distinct bacterioneuston communities and
511 organic surfactant enrichments in the SML collapses above a sharp threshold of 5.6 m s^{-1} (Rahlff
512 et al., 2017). This agrees with wind-wave tank experiments showing a breakdown of
513 biopolymer SML enrichment at wind speeds above 6 m s^{-1} (Engel et al., 2026), corresponding
514 precisely to the 6–8 m s^{-1} wind bin in our dataset where gradient intensity and front detection
515 rates reached their lowest levels.

516 Below 5 m s^{-1} , natural surfactant films are known to damp capillary and short gravity waves
517 substantially, visually appearing as slicks (Romano, 1996; Wurl et al., 2011). A patch of
518 surfactant-damped surface water bounded by a sharp edge represent small-scale, non-advective
519 feature. This can inflate both the mean gradient intensity and the front-detection rate specifically
520 under low wind conditions, independent of the tidal and thermohaline forcing discussed above.
521 The regime classification supports this finding as the proportion of thermally-dominated
522 observations is highest under calm winds (33.2%) and declines under stronger winds. This is



523 consistent with thin, thermally stratified, surfactant-stabilised surface patches that persist only
524 under low to no turbulent conditions. We note that this dataset carries no independent optical
525 or visual slick flag, so the wind-bin evidence here is physically-motivated and consistent with,
526 rather than a direct confirmation of, slick presence.

527 The two summer deployments (HE626, July/August 2023; HE644, July 2024) showed the
528 highest median gradient intensity and front rates, and the autumn cruise (HE609, October 2022)
529 the lowest. This is consistent with the seasonality of stratification in the German Bight Region
530 of Freshwater Influence (ROFI) (Chegini et al., 2020). Strong stratification occurs in spring and
531 early summer, driven by peak Elbe and Weser discharge from snowmelt and seasonal rainfall
532 interacting with the spring–neap tidal cycle (Chegini et al., 2020). During these times, the along-
533 track horizontal density gradient across the Elbe–Weser plume boundary can locally reach
534 magnitudes of up to 1 kg m^{-3} over a few hundred metres under strong plume conditions (Kopte
535 et al., 2022). Conversely, the skin density gradient intensity was significantly lower during
536 rainfall than under dry conditions.

537 These mechanistic interpretations are corroborated by our Random Forest regression and SHAP
538 attribution analysis (Breiman, 2001; Lundberg and Lee, 2017). Wind speed and atmospheric
539 pressure emerged as the two primary meteorological predictors of gradient intensity. Partial
540 dependence analysis predicted gradient intensity peaks under near-zero wind speed ($> 0.016 \text{ kg}$
541 m^{-3}), declines rapidly between 2 and 4 m s^{-1} , and stabilises at a low baseline above 5 m s^{-1} . This
542 5 m s^{-1} inflection matches the established threshold for rain-driven skin stratification disruption
543 (Gassen et al., 2024c) and the destabilization of the SML (Rahlff et al., 2017). Atmospheric
544 pressure showed a comparatively muted but non-monotonic relationship in the partial
545 dependence analysis, with predicted gradient intensity highest at lower pressures. In terms of
546 hydrographic predictors, skin temperature showed a steep non-linear increase in predicted
547 gradient intensity above 17°C , peaking near 22°C , consistent with the summer-cruise seasonal
548 signal identified above. Because surface tension decreases monotonically with temperature
549 (Nayar et al., 2014) localized solar warming in summer skin layers can generate thermal
550 Marangoni stresses (Soloviev and Lukas, 2014). Under low wind conditions, these surface-
551 tension gradients damp micro-turbulent dispersion and maintain sharp horizontal thermal
552 gradients across the skin interface, consistent with the observed increase in predicted gradient
553 intensity at higher skin temperatures. Skin salinity exerted its strongest effect in the least saline
554 water masses ($29\text{--}30 \text{ g kg}^{-1}$), consistent with proximity to the Elbe–Weser plume boundary
555 (Kopte et al., 2022). The complexity collaborated the convergence of the SHAP ranking, the



556 partial-dependence thresholds, and the independently-established physical thresholds for
557 surfactant collapse, rain-driven stratification, and plume-driven seasonality. Our findings
558 provide predictive results as a mechanistic foundation as it would not have as a standalone
559 statistical fit.

560 **4.4 Broader Scientific Implications for Remote Sensing and Air–Sea Fluxes**

561 These empirical findings have direct implications for satellite oceanography and climate
562 modelling. Infrared satellite sensors sample thermal radiation from the ocean skin layer, at a
563 depth of 10 – 20 μm , whereas microwave sensors integrate radiation over a much greater depth
564 of ~ 1 cm known as the subskin layer with signal decaying exponentially with depth (Donlon et
565 al., 2002). Front-detection algorithms applied to satellite imagery (Belkin and O’Reilly, 2009)
566 assume that detected surface gradients accurately reflect underlying water structures (Ayim et
567 al., 2025b, 2026a; Vazquez-Cuervo et al., 2019). Our results show that sub-surface fronts are
568 attenuated by roughly 60% at the skin interface and diverge significantly during strong frontal
569 processes. Consequently, satellite-derived gradient intensities likely underestimate the true
570 magnitude of underlying submesoscale fronts in shallow shelf seas. Furthermore, global climate
571 models parameterize air–sea heat and momentum exchange using bulk sea surface
572 measurements taken at 1 to 5 m depth (Fairall et al., 1996). The vertical partial decoupling
573 observed in our study confirms that the skin layer maintains its unique spatial variability.
574 Neglecting this micro-surface variability may introduce biases into calculated air–sea flux
575 budgets, particularly during low-wind regimes ($< 5 \text{ m s}^{-1}$) where skin gradients are most
576 pronounced.

577 The wind-speed threshold identifies the specific meteorological regime (calm, $< 2 \text{ m s}^{-1}$) in
578 which skin layer gradient structure is least representative of, and widely decoupled from, the
579 sub-surface field. Studies on satellite validations would typically use sub-surface characteristics
580 as ground truth due to the lack of “true” surface data. We have discussed the need of ocean skin
581 data for satellite validation in previous studies (Ayim et al., 2025b, 2026a). Validation
582 campaigns timed to calm, clear-sky conditions may therefore be systematically sampling the
583 regime of largest skin/NSL disagreement, and not the smallest. That is true in frontal zones as
584 shown in this study. The inter-cruise analysis shows that a single research cruise, even so using
585 state-of-the-art instrumentations, may not capture a representative tidal-phase relationship. This
586 is highlighted by the reversal observed during HE614 and the well-documented seasonal
587 modulation of German Bight ROFI stratification (Chegini et al., 2020; Kopte et al., 2022).



588 Multi-season validation campaigns, or explicit seasonal stratification of validation statistics, are
589 needed before generalising a tidal correction factor derived from any single deployment.

590

591 **5. Conclusion**

592 This study demonstrates that submesoscale horizontal density gradients undergo partial vertical
593 decoupling and attenuation between the ocean skin and the NSL. High-resolution along-track
594 measurements from the autonomous surface vehicle HALOBATES in the tidally energetic
595 German Bight reveal that surface skin gradients do not consistently mirror underlying water
596 dynamics, and that this decoupling is disproportionately carried by salinity rather than
597 temperature. Tidal phase is a significant modulator of gradient intensity, and machine learning
598 analysis via Random Forest models and SHAP attribution identified atmospheric pressure, wind
599 speed, and air-sea temperature difference as predictors of gradient intensity parallel to the
600 expected dominant hydrographic predictor of temperature and salinity.

601 Despite these key findings, several limitations should be considered. The observational dataset
602 is geographically restricted to the shallow, tidally energetic waters of the German Bight,
603 meaning the observed environmental thresholds may differ in open-ocean regimes or other
604 coastal environments. Sampling was limited to five specific deployment windows across two
605 years; the March 2023 cruise, notably, reversed the pooled tidal-phase pattern seen in the other
606 four cruises. This underscores that the tidal mechanism proposed here is likely conditional on
607 seasonal estuarine stratification state rather than universal and may not be fully captured by our
608 seasonal sampling design. The Random Forest model's low explanatory power ($R^2 = 0.19$) and
609 its systematic under-prediction of the largest observed gradients indicate that the SHAP-ranked
610 drivers should be considered as relative importance among measured variables rather than a
611 complete explanation of gradient intensity. Addressing these limitations through expanded
612 multi-platform campaigns, continuous vertical profiling, and multi-season deployment will be
613 essential for refining remote sensing validation strategies and improving submesoscale
614 frontogenesis models.

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621 **Data availability**

622 All data are freely available at the data repository PANGAEA as follows. Digital object
623 identifiers are available for processed data from the HALOBATES.

624 HE598; <https://doi.org/10.1594/PANGAEA.968799>.

625 HE609; <https://doi.org/10.1594/PANGAEA.968800>.

626 HE614; <https://doi.org/10.1594/PANGAEA.969378>.

627 HE626; <https://doi.org/10.1594/PANGAEA.972989>.

628 HE644; <https://doi.org/10.1594/PANGAEA.989737>.

629

630 **Author contributions**

631 OW and SMA conceptualized the ideas. SMA curated the data, conducted formal analysis and
632 investigation, developed the methodology and the prepared the original draft and visualisation.
633 OW acquired the funding, managed the project, provided the resources, supervised and
634 reviewed and edited the manuscript.

635

636 **Competing interests**

637 The authors declare that they have no conflict of interest.

638

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647

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