



1 **Intense concentric gravity waves from the stratosphere to the**
2 **thermosphere and a massive red sprite outbreak triggered by severe**
3 **thunderstorms South of the Himalayas**

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16 **Abstract**

17 Severe thunderstorms drive vertical coupling through electrical processes
18 associated with transient luminous events (TLEs) and dynamical processes associated
19 with atmospheric gravity waves (AGWs). On 19 May 2022, an exceptionally intense
20 thunderstorm outbreak occurred over the northern Bay of Bengal–Himalayan foothill
21 region. The storm system produced more than 100 red sprites, 16 secondary gigantic
22 jets, and at least four ghosts, accompanied by visible concentric airglow ripples. This
23 event represented the first large-scale detection of a massive red sprite outbreak in
24 China. From a dynamical perspective, concentric gravity waves (CGWs) were observed
25 from the stratosphere to the thermosphere using coordinated satellite and ground-based
26 airglow observations. AIRS measurements revealed concentric wave structures in the
27 stratosphere above the convective source region, while ground-based airglow imagers
28 detected CGWs in the OH, OI 557.7 nm, and OI 630.0 nm emission layers. CGWs in
29 the mesosphere and lower thermosphere (MLT) propagated northeastward for more
30 than 3000 km, suggesting propagation within favorable MLT ducting structures. In
31 contrast, thermospheric CGWs observed in the OI 630.0 nm emission exhibited
32 horizontal phase speeds of 260–280 m s⁻¹. Reverse ray-tracing analysis indicates a
33 preferred source altitude near 150 km, suggesting that these waves were likely
34 secondary gravity waves generated by the dissipation and breaking of upward-
35 propagating primary gravity waves. Their asymmetric occurrence is likely attributable
36 to Doppler-induced dissipation associated with thermospheric background winds.
37 These observations provide a rare example of simultaneous electrical and dynamical



38 coupling from the troposphere to the thermosphere during an extreme convective event.



39 1. Introduction

40 Severe thunderstorms influence the upper atmosphere through both dynamical and
41 electrical pathways. Dynamically, latent heat release and overshooting convective
42 motions generate atmospheric gravity waves (AGWs) that can propagate from the
43 troposphere into the mesosphere and thermosphere (Fovell et al., 1992; Alexander and
44 Holton, 2004; Heale et al., 2019, 2022; Smith et al., 2020; Nyassor et al., 2022; Li et
45 al., 2025). As AGWs propagate upward through the decreasing-density atmosphere,
46 their amplitudes increase and may eventually become convectively or dynamically
47 unstable. Subsequent wave breaking deposits momentum and energy into the
48 background atmosphere and can generate secondary gravity waves that continue
49 propagating into higher atmospheric regions (Holton, 1983; Vadas et al., 2003; Smith
50 et al., 2013; Li et al., 2022). Through these processes, AGWs transfer energy and
51 momentum across atmospheric layers and constitute one of the primary mechanisms
52 responsible for vertical atmospheric coupling (Fritts and Alexander, 2003; Plane et al.,
53 2023).

54 Electrically, intense thunderstorms can produce transient luminous events (TLEs),
55 including red sprites, gigantic jets, and other high-altitude optical emissions above
56 active storm systems (Pasko et al., 1997; Sentman et al., 2003; Chou et al., 2016; Wang
57 et al., 2019; Lai et al., 2022; Maurya et al., 2022). Red sprites are among the most
58 common TLEs and typically occur in the mesosphere above powerful positive cloud-
59 to-ground lightning discharges. Particularly intense sprite outbreaks may also be
60 accompanied by rare phenomena such as secondary gigantic jets and ghosts (van der



61 Velde et al., 2019; Passas-Varo et al., 2023). Although both TLEs and AGWs originate
62 from the same convective source, observations simultaneously revealing the electrical
63 and dynamical impacts of extreme thunderstorms across multiple atmospheric layers
64 remain relatively rare.

65 The northern Bay of Bengal–Himalayan foothill region is a major hotspot of deep
66 convection and atmospheric disturbances in the Earth system (Maurya et al., 2022; Peng
67 et al., 2024). On the night of 19 May 2022, an exceptionally intense thunderstorm
68 outbreak occurred over the northern Bay of Bengal–Himalayan foothill region. The
69 storm system produced more than 100 red sprites, 16 secondary gigantic jets, and four
70 ghosts, representing the largest documented sprite outbreak observed from China.
71 Simultaneously, distinct concentric airglow ripples were recorded, suggesting strong
72 gravity-wave activity associated with the convective system. Subsequent observations
73 revealed concentric gravity waves (CGWs) spanning from the stratosphere to the
74 thermosphere, providing a unique opportunity to investigate both the electrical and
75 dynamical impacts of an extreme convective event on the middle and upper atmosphere.

76 In this study, we use coordinated observations from ground-based airglow imagers,
77 AIRS stratospheric measurements, TIMED/SABER observations, and Meteosat/
78 SEVIRI imagery to investigate the CGWs associated with this exceptional
79 thunderstorm outbreak. Particular attention is given to (1) the long-distance propagation
80 of CGWs in the mesosphere and lower thermosphere (MLT), (2) the generation
81 mechanism of thermospheric CGWs, and (3) the role of thermospheric background
82 winds in modulating CGW observability. The observations provide a rare example of

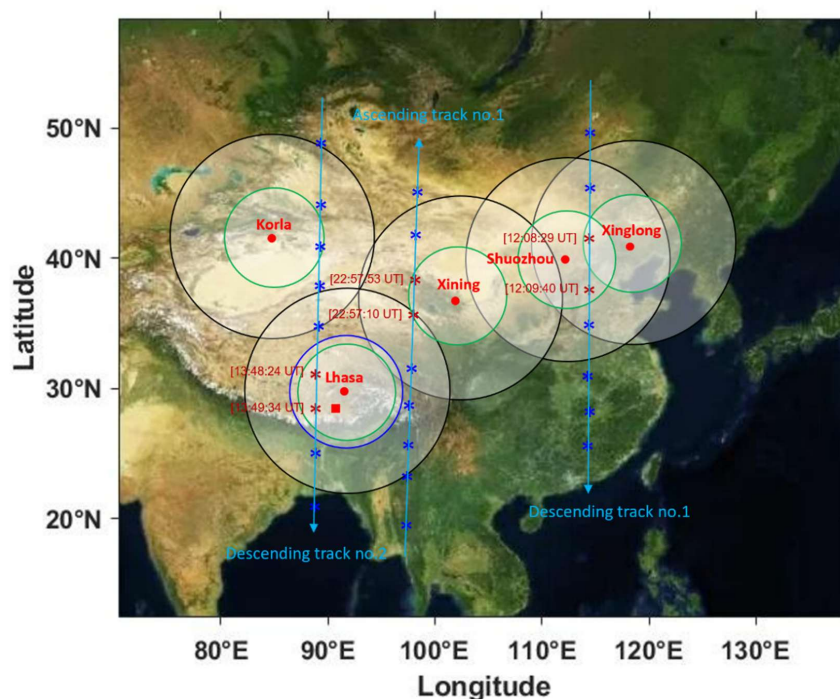


83 deep-convection-driven vertical coupling extending from the troposphere to the
84 thermosphere and reveal distinct propagation characteristics of mesospheric and
85 thermospheric CGWs.

86 **2. Observations**

87 **2.1 Ground-based airglow observations and amateur photography imaging**

88 Figure 1 presents the location map of airglow observation stations at Lhasa
89 (30.1°N, 90.6°E), Xining (36.5°N, 102.1°E), Shuozhou (39.8°N, 112.1°E), Xinglong
90 (40.2°N, 117.4°E), and Korla (41.7°N, 86.2°E). The airglow observation station in
91 Lhasa is equipped with three dedicated airglow imagers for multi-layer monitoring. The
92 three imagers observe distinct emission layers: one observes the hydroxyl (OH) airglow
93 near 87 km altitude with a 63-second temporal resolution; a second is configured for
94 atomic oxygen emission at 557.7 nm (OI 557.7 nm) originating from approximately 96
95 km; and a third detects the OI 630.0 nm emission near 250 km altitude. The two oxygen
96 emission imagers (OI 557.7 nm and OI 630.0 nm) both operate with a 92-second
97 temporal resolution. This coordinated instrument array enables simultaneous
98 observations of mesospheric (OH and OI 557.7 nm) and thermospheric (OI 630.0 nm)
99 airglow emissions. The Xining, Shuozhou, Xinglong, and Korla stations are equipped
100 with OH and OI 630.0 nm imagers identical to those deployed at Lhasa, but do not
101 include the OI 557.7 nm channel. In addition to the scientific airglow imaging
102 instruments, supplementary time-lapse photographs of the sprite event were obtained
103 using three Sony Alpha 7RIII cameras equipped with 14 mm, 35 mm, and 135 mm
104 lenses, respectively, with 5 s exposures at f/1.4.



105

106 **Figure 1** Map showing the geographic location of the airglow imagers station in Lhasa, Xining,
107 Shuozhou, Xinglong, and Korla (red dots) and red-sprite observation site documented by amateur
108 photographers (red square). The circle on the map gives the effective observation ranges of OH
109 airglow imagers and red-line (630 nm) imagers with a 164° field of view. The red asterisks denote
110 the TIMED/SABER descending track footprints passing over Lhasa on 19 May 2022.

111 2.2 Aqua/AIRS satellite observations

112 Mounted on NASA's Aqua satellite, the AIRS (Aumann et al., 2003) employs
113 infrared spectroscopy for atmospheric sounding. The instrument conducts continuous
114 scans perpendicular to the satellite track, sequentially sampling atmospheric radiance
115 within footprints measuring approximately 13–14 km across at nadir. This scanning
116 pattern yields a swath width of ~1765 km (Hoffmann et al., 2014), with data typically
117 aggregated into 6-minute granules. AIRS exhibits sensitivity to atmospheric
118 temperature fluctuations induced by GWs, specifically detecting perturbations



119 associated with vertical wavelengths exceeding 10–15 km and horizontal wavelengths
120 between ~50 and 500 km (Hoffmann and Alexander, 2010). In this study, radiance
121 measurements within the 4.3 μm CO₂ fundamental band are employed. This spectral
122 region is highly sensitive to temperature perturbations in the stratosphere, particularly
123 between altitudes of 30–40 km, making it well suited for quantifying stratospheric
124 temperature perturbations.

125 **2.3 TIMED/SABER satellite observations**

126 Aboard NASA's TIMED satellite (launched Dec. 7, 2001), SABER is a 10-channel
127 broadband limb-scanning infrared radiometer (1.27–17 μm) (Russell et al., 1999) that
128 provides key observations for investigating energy transfer processes in the
129 Mesosphere–Lower Thermosphere–Ionosphere (MLTI) region. SABER measures
130 kinetic temperature (± 2 –5 K) using CO₂ emissions, including the 15 μm LTE channel
131 below ~90 km and the 4.3 μm non-LTE channel at higher altitudes. By simultaneously
132 observing emissions from O₃ (9.6 μm), OH (1.6–2.0 μm), and O₂ (1.27 μm), SABER
133 quantifies key thermal processes, including radiative cooling (reaching 150 K day⁻¹)
134 and chemical heating (~8 K day⁻¹). These measurements, with a 2–4 km vertical
135 resolution, characterize the thermal structure and background state of the MLT region.

136 **2.4 Meteosat/SEVIRI satellite observations**

137 The Meteosat series, operated by EUMETSAT, is Europe's primary geostationary
138 meteorological satellite system, providing continuous observations of atmospheric and
139 convective processes over Europe, Africa, and adjacent regions. The primary
140 instrument aboard the Meteosat Second Generation (MSG) satellites (Schmetz et al.,



141 2002) is SEVIRI, which acquires full-disk images every 15 minutes across 12 spectral
142 bands.

143 **2.5 SD-WACCM-X and TIE-GCM 3.0 simulations**

144 SD-WACCM-X (the Specified Dynamics Whole Atmosphere Community
145 Climate Model with thermosphere-ionosphere eXtension) is a whole atmosphere model
146 developed within the NCAR Community Earth System Model (CESM) framework
147 (Hurrell et al., 2013; Liu et al., 2018). Its domain extends from the Earth's surface to
148 the upper thermosphere. The specified dynamics configuration enables improved
149 representation of upward-propagating atmospheric waves. TIE-GCM 3.0 (the
150 Thermosphere–Ionosphere–Electrodynamics General Circulation Model, Version 3.0)
151 is a first-principles numerical model that self-consistently simulates the coupled
152 thermosphere-ionosphere system in the upper atmosphere (approximately 97–600 km
153 altitude) (Richmond et al., 1992; Günzkofer et al., 2025). The model includes self-
154 consistent electrodynamics and ion-neutral coupling processes that are important for
155 thermosphere-ionosphere dynamics.

156 **3. Results and Discussion**

157 The extreme convective event produced both electrical and dynamical responses
158 in the middle and upper atmosphere. The electrical response was manifested by an
159 exceptional TLE outbreak, including red sprites, secondary gigantic jets, and ghosts,
160 reflecting intense lightning activity extending into the middle atmosphere. The
161 dynamical response involved the generation, propagation, and secondary excitation of
162 gravity waves from the stratosphere to the thermosphere.



163 **3.1 Convective evolution and transient luminous event outbreak**

164 Figure 2 shows continuous SEVIRI 10.8 μm brightness temperature imagery from
165 the Meteosat Second Generation (MSG) satellite during 10:00–21:00 UT on 19 May
166 2022. The imagery reveals the evolution of a deep convective system located along the
167 southern foothills of the Himalayas. At 10:00 UT, the system initially appeared as
168 multiple isolated convective cells with cold cloud tops distributed over approximately
169 23–27°N and 85–93°E. Between 11:00 and 12:00 UT, these individual cells rapidly
170 intensified and merged, forming a more organized mesoscale convective system (MCS)
171 with substantial eastward expansion. This period was characterized by vigorous deep
172 convection and enhanced latent heat release, conditions favorable for the excitation of
173 atmospheric GWs. The convective system also served as the parent storm of the
174 subsequent transient luminous event outbreak.

175 Between 13:00 and 15:00 UT, the MCS reached its mature stage, characterized by
176 a large and consolidated cold cloud shield extending across approximately 22–28°N
177 and 85–90°E. Persistently low brightness temperatures indicate strong vertical
178 convective development penetrating into the upper troposphere and lower stratosphere.
179 From 16:00 to 18:00 UT, the convective system gradually weakened and fragmented,
180 accompanied by contraction of the cold cloud shield. By 19:00–21:00 UT, only
181 scattered remnants of stratiform cloud structures remained over 22–28°N and 88–94°E,
182 indicating the dissipation of the previously organized convective system. This
183 exceptionally intense convective system subsequently produced one of the largest
184 transient luminous event outbreaks observed from China.

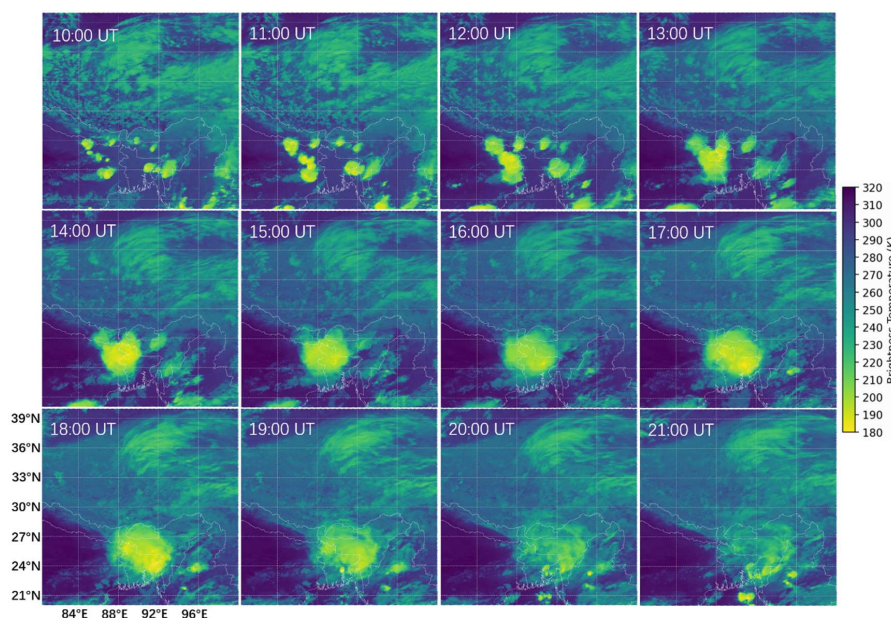


Figure 2 SEVIRI 10.8 μm channel brightness temperature composite from Meteosat Second Generation (MSG) during 10:00–21:00 UT on 19 May 2022.

Figure 3 presents representative images of the TLE outbreak associated with the convective system shown in Fig. 2. The images were captured using time-lapse photography by photography enthusiasts from the shores of Lake Pumoyongcuo in Shannan, Tibet, on the night of 19 May 2022. The outbreak persisted for approximately three hours, from 13:00 to 16:00 UT, coinciding with the mature stage of the parent mesoscale convective system. More than 100 red sprites, 16 secondary gigantic jets, and four ghosts were recorded during this period, making it the largest documented sprite outbreak observed from China (Huang et al., 2025). The images shown in Fig. 3 represent only a subset of the observed activity. Various sprite morphologies, including carrot-shaped and columniform structures, were observed, often occurring in large clusters over a broad horizontal region above the parent storm. Particularly notable are



199 the secondary gigantic jets visible at 13:16 UT and the extensive sprite clusters
200 observed around 13:19–13:42 UT. The simultaneous occurrence of numerous sprites
201 together with multiple secondary gigantic jets and ghosts indicates exceptionally
202 intense electrical activity within the parent convective system.



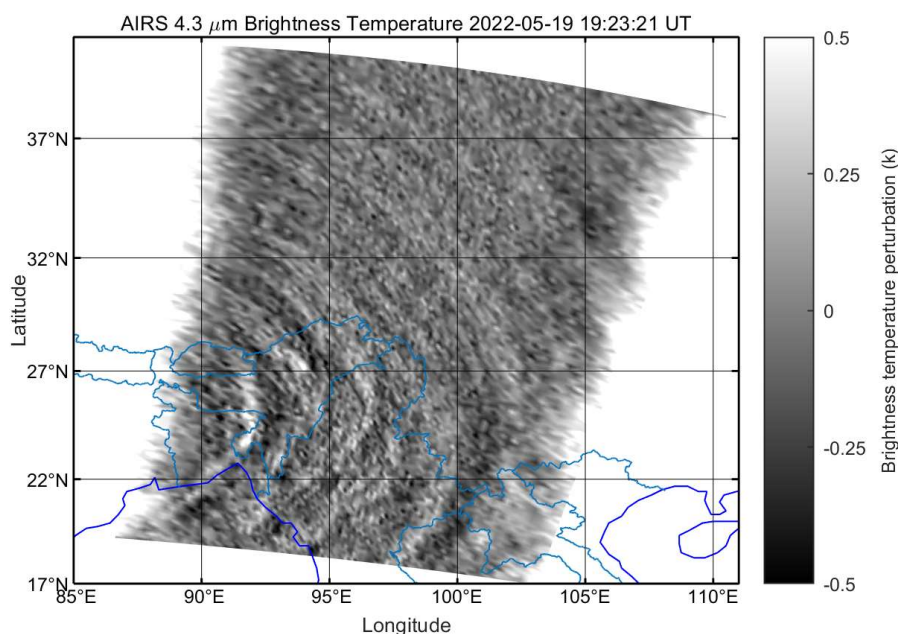
203
204 **Figure 3** Transient luminous events (TLEs) captured by amateur photographer Angel An over the
205 southern Tibetan Plateau on the night of 19 May 2022, showcasing a variety of red sprite
206 morphologies. Photographs © Angel An.

207 **3.2 Stratospheric concentric gravity waves from the AIRS/Aqua**

208 Figure 4 shows annular GWs in AIRS 4.3 μm brightness temperature perturbations
209 (19:23:21 UT, 19 May 2022), exhibiting concentric wave fronts centered over
210 northeastern Bangladesh ($\sim 25^\circ\text{N}$, 88°E). These waves display a distinct radial
211 wavelength gradient, with horizontal wavelengths increasing from ~ 50 km near the
212 central region to >150 km at the periphery. The observed concentric structures are
213 spatially associated with the previously identified deep convective system and suggest
214 active wave excitation in the stratosphere during the late stage of convective evolution.



215 Although the primary convective cloud shield had weakened by this time, localized
216 convective activity may still have contributed to the generation of the observed CGWs.



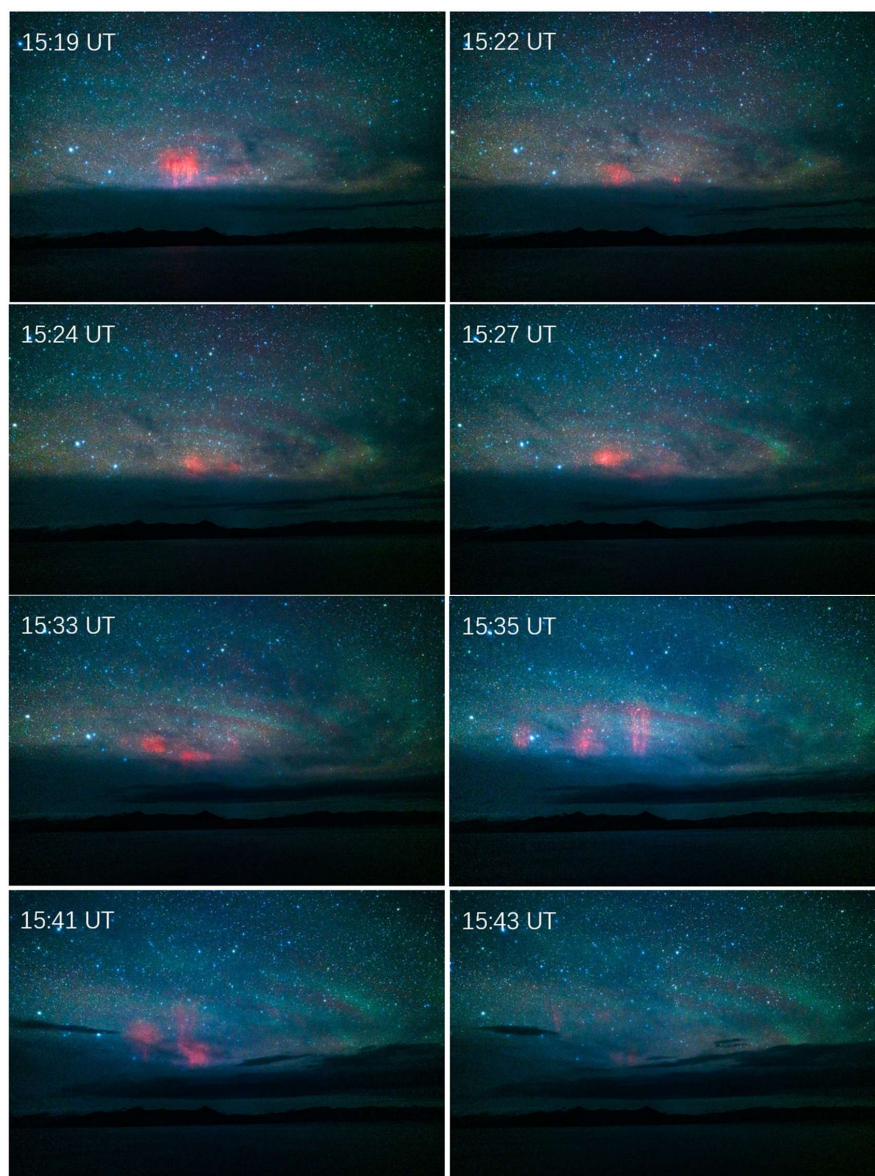
217
218 **Figure 4** CGWs observed by Aqua/ARIS at 4.3 μm brightness temperature on 19 May 2022 at
219 19:23:21 UT.

220 **3.3 Multi-layers ground-based optical observations of concentric gravity waves** 221 **from the mesosphere to the thermosphere**

222 Figure 5 presents a sequence of time-lapse photographs acquired from the shores
223 of Lake Pumoyongcuo between 15:19 and 15:43 UT on 19 May 2022. In addition to
224 the transient luminous events, a series of faint concentric airglow ripples can be
225 identified across the night sky. These ripples appeared as multiple arc-shaped bands
226 extending over a broad horizontal region and exhibited a clear concentric geometry,
227 suggesting the presence of concentric gravity waves (CGWs) propagating through the
228 mesopause airglow layer. The wave fronts remained visible for more than 20 min and



229 gradually expanded across the field of view. Simultaneously, sprite activity continued
230 above the distant convective system, indicating that both the electrical phenomena and
231 the airglow disturbances were associated with the same parent thunderstorm. Although
232 the photographic observations do not permit quantitative determination of wave
233 parameters, they provide direct visual evidence of large-scale concentric wave
234 structures accompanying the exceptional convective outbreak.



235

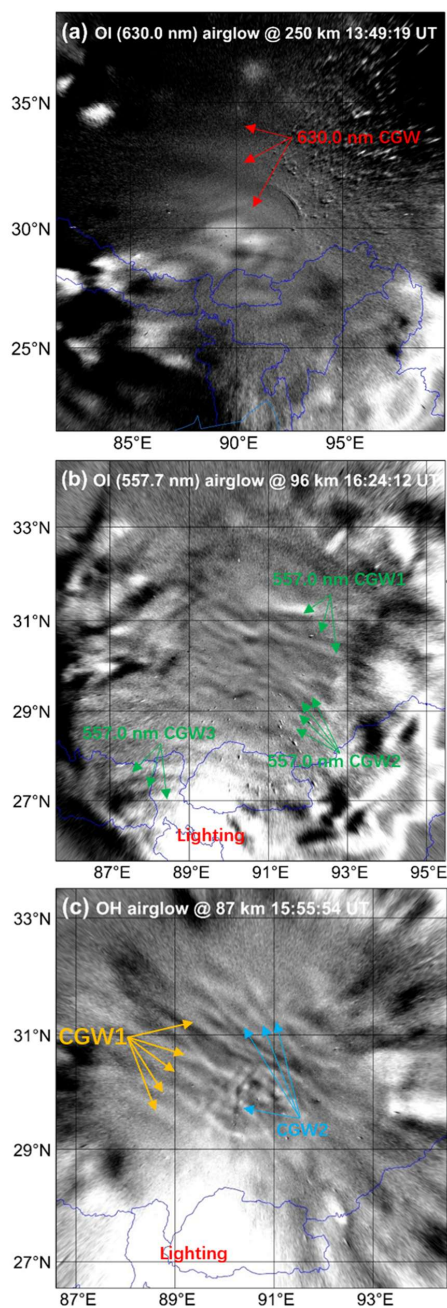
236 **Figure 5.** Temporal sequence of concentric gravity waves (CGWs) observed during the red sprite
237 outbreak over the southern Tibetan Plateau from 15:19 to 15:43 UT on 19 May 2022. Photographs
238 © Angel An.

239 Figure 6 presents multi-layer airglow observations at the Lhasa station, revealing
240 concentric gravity wave structures observed in the mesosphere, mesopause region, and



241 thermosphere. Figure 6a shows OI (630.0 nm) airglow observations at ~250 km altitude,
242 where large-scale concentric wave fronts are visible in the thermospheric F-region.
243 Figure 6b displays OI (557.7 nm) airglow observations near ~96 km, showing multiple
244 concentric gravity wave packets in the mesopause region. Figure 6c depicts OH airglow
245 observations at ~87 km, revealing pronounced concentric wave structures in the MLT
246 region.

247 The wave morphology and propagation characteristics exhibit notable differences
248 among the three emission layers. The OH and OI (557.7 nm) emissions display clearer
249 and more coherent concentric wave fronts in the MLT region, whereas the OI (630.0
250 nm) emission exhibits broader and more diffuse thermospheric wave structures. Despite
251 these differences, the observed wave patterns are spatially associated with the same
252 deep convective system to the south of the Himalayas. The simultaneous detection of
253 concentric gravity wave signatures over a broad altitude range suggests strong vertical
254 coupling between the lower atmosphere and the upper atmosphere during this event.



255

256 **Figure 6** Multi-layer airglow observations of CGWs at Lhasa station, showing emissions from three
 257 different altitudes: (a) OI 630.0 nm (thermospheric airglow, ~250 km); (b) OI 557.7 nm
 258 (mesospheric airglow, ~97 km); (c) OH band (mesospheric airglow, ~87 km).



259 To estimate the wave parameters of the CGWs shown in Fig. 5 captured by
260 astrophotography enthusiasts, we used a star chart to precisely determine the elevation
261 (el) and azimuth angle (az) of the stars in the image. The horizontal distances between
262 the observation point and the airglow layer can be calculated using the following
263 equation (Yue and Lyons, 2015):

$$s = R * \left[\frac{\pi}{2} - el - \sin^{-1} \left(\frac{R * \sin(el + \frac{\pi}{2})}{R + H} \right) \right], \quad (1)$$

264 where R is the Earth's radius, and H is the altitude of the airglow layer.

265 The horizontal coordinates (x, y) within the airglow layer can then be expressed as
266 follows,

$$x = s * \sin(az), \quad (2)$$

$$y = s * \cos(az), \quad (3)$$

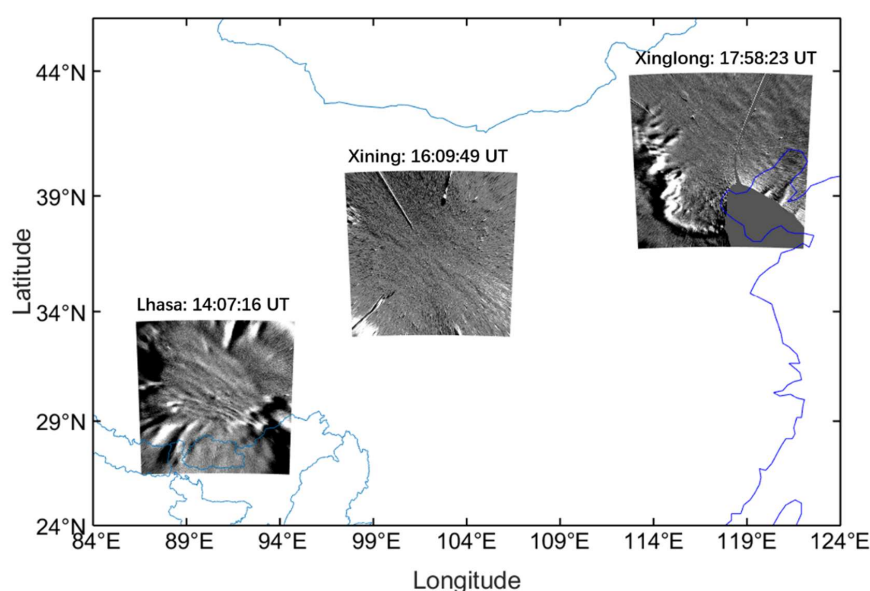
270 The horizontal wavelength of the concentric wave structures was estimated to be
271 approximately 38 ± 5 km. This value is consistent with the horizontal wavelength of
272 CGW1 observed at 15:55:54 UT in Fig. 6c, suggesting that the ripple structures
273 captured by the amateur photographer and the CGW observed by the ground-based OH
274 airglow imager likely correspond to the same event. This agreement provides
275 independent confirmation that the concentric airglow ripples visible in the photographs
276 correspond to the same mesospheric CGW system detected by the OH imager.

277 **3.4 Large-scale propagation of mesospheric and lower thermospheric concentric** 278 **gravity waves**

279 During a massive red sprite cluster over the Tibetan Plateau, coordinated OH
280 airglow imaging at Lhasa, Xining, and Xinglong shown in Fig. 7 revealed a large-scale



281 gravity wave propagating northeastward ($13.4^\circ \pm 3.7^\circ$ east of north) with a horizontal
282 wavelength of 33.4 ± 4.9 km and period of 13.4 ± 2.9 min. The wave exhibited a phase
283 speed of 42.8 ± 7.4 m s⁻¹ and spanned $>16^\circ$ in latitude ($\sim 30^\circ\text{N}$ – 46°N) and longitude,
284 though wave amplitude and wavelength at Xinglong were notably attenuated compared
285 to lower-latitude stations. The consistent propagation direction and wave characteristics
286 observed across multiple stations suggest that the CGWs originated from the same deep
287 convective system located to the south of the Himalayas. The coherent wave structures
288 observed across multiple stations indicate persistent northeastward propagation of the
289 concentric gravity waves over a large horizontal distance exceeding 3000 km. Such
290 extended horizontal propagation suggests that the waves were unlikely to propagate as
291 simple freely propagating GWs and may instead have experienced favorable
292 propagation conditions within the MLT region.



293
294 **Figure 7** CGWS observed by OH airglow imagers: Lhasa station at 14:07:16 UT, Xining station at



295 16:09:49 UT, and Xinglong station at 17:58:23 UT.

296 To investigate the propagation characteristics of the CGWs in the MLT region, we
297 examined the vertical wavenumber using the GW dispersion relation (Hines, 1960):

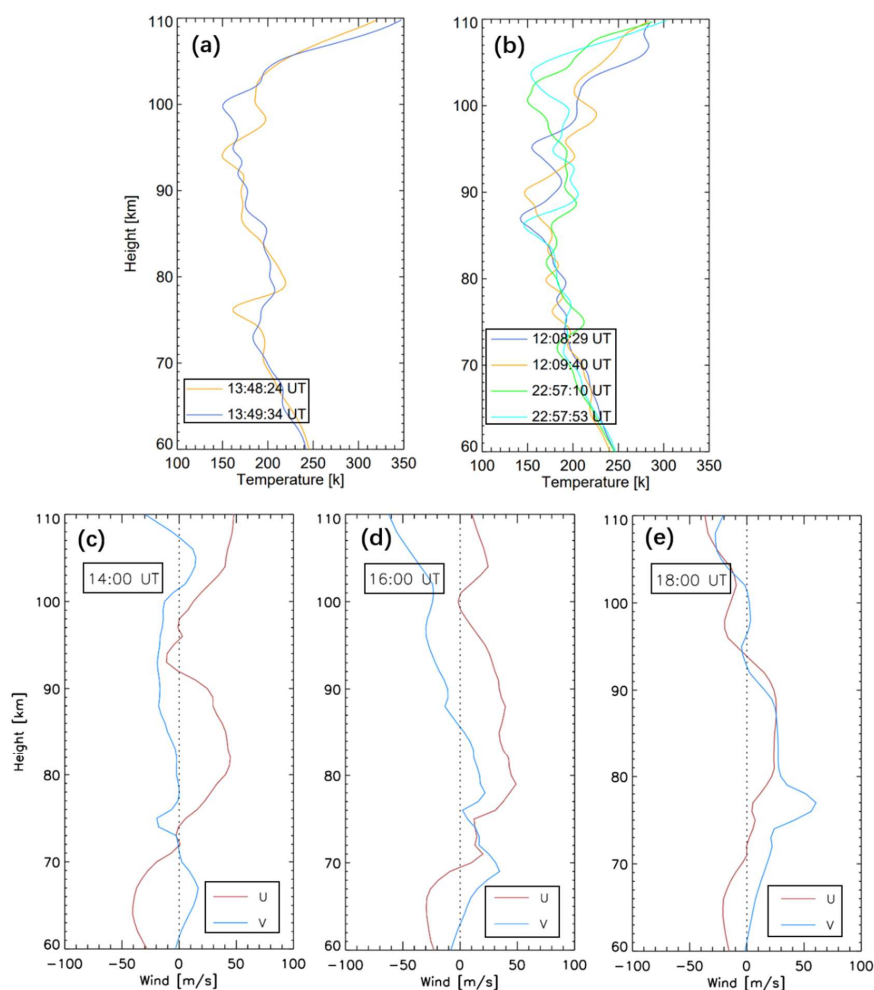
$$298 \quad m^2 = \frac{N^2}{(c-u)^2} - k_H^2 - \frac{1}{4H^2}, \quad (4)$$

299 Here, m and k_H represent the vertical and horizontal wavenumbers, respectively,
300 derived from the airglow imaging observations; c is the observed horizontal phase
301 speed of the gravity wave; u denotes the background wind speed in the direction of
302 wave propagation, obtained from the SD-WACCM-X wind fields; H is the scale height;
303 and N^2 is the Brunt–Väisälä frequency, defined as:

$$304 \quad N^2 = \frac{g}{T} \left(\frac{dT}{dz} + \frac{g}{c_p} \right), \quad (5)$$

305 where $g = 9.54 \text{ ms}^{-2}$ is gravitational acceleration, T is the temperature, $c_p = 1005$
306 $\text{J Kg}^{-1}\text{K}^{-1}$ is the specific heat at constant pressure, dT/dz represents the vertical
307 temperature gradient derived from TIMED/SABER observations.

308 The temperature profiles used in this analysis were obtained from SABER/TIMED
309 measurements, and the background wind conditions were characterized using SD-
310 WACCM-X wind fields, as shown in Fig. 8. Figure 8a and b show the temperature
311 profiles over the relevant locations on 19 May 2022, with the corresponding UTC times
312 and locations indicated in Fig. 1. Figure 8c, d and e illustrate the wind fields for Lhasa,
313 Xining, and Xinglong, respectively.



314

315 **Figure 8** Temperature profiles derived from SABER/TIMED measurements (panels a and b) and
 316 SD-WACCM-X wind fields (panels c–e). (a) and (b) show profiles measured on 19 May 2022 at the
 317 following UTC times: (a) 13:48:24 and 13:49:34; (b) 12:08:29, 12:09:40, 22:57:10, and 22:57:53
 318 (corresponding locations are marked in Fig. 1). (c), (d), and (e) present the wind fields over Lhasa,
 319 Xining, and Xinglong, respectively.

320 Figure 9 presents the vertical wavenumber profiles over (a) Lhasa, (b) Xining, and
 321 (c) Xinglong. The profiles reveal stable ducting layers at Lhasa (75–93 km), Xining
 322 (65–99 km), and Xinglong (70–94 km). These ducting layers suggest that the deep



323 convection-generated CGWs were able to propagate over horizontal distances
324 exceeding 3000 km within the MLT region. The presence of positive vertical
325 wavenumber regions bounded by evanescent layers indicates favorable waveguide
326 conditions that can effectively trap gravity-wave energy and suppress vertical leakage.
327 Consequently, wave energy can be transported horizontally over much greater distances
328 than would be expected for freely propagating gravity waves. The observed propagation
329 distance therefore highlights the critical role of MLT ducting structures in enabling
330 large-scale horizontal transport of gravity-wave energy and momentum.

331 Despite the shallower duct over Lhasa, the peak OH volume emission rate there
332 was approximately three times larger than those over Xining and Xinglong (Fig. 9).
333 Nevertheless, comparable CGW signatures were observed at all three stations. This
334 suggests that the visibility of the wave structures was controlled not only by the
335 background OH emission intensity but also by the ducting characteristics and the
336 vertical overlap between the trapped wave field and the OH emission layer. Therefore,
337 a deeper duct may favor long-distance horizontal propagation, whereas the observed
338 airglow perturbation amplitude is additionally influenced by local airglow conditions
339 and wave–airglow layer coupling.

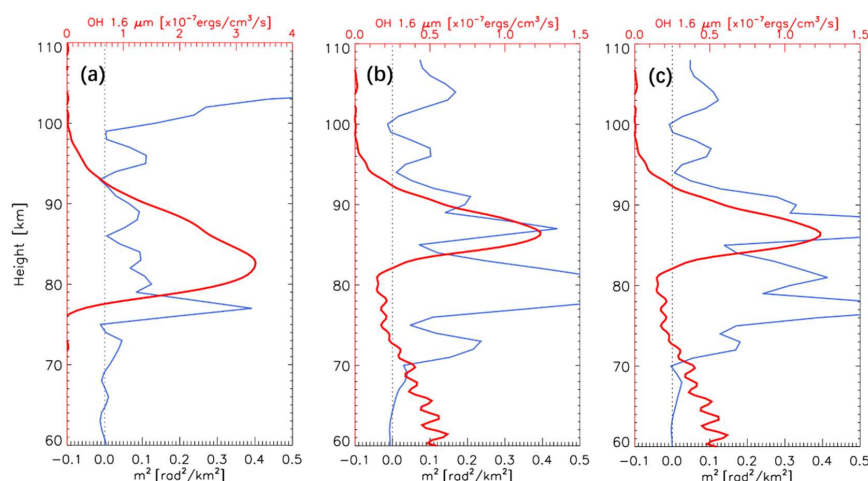


Figure 9 Vertical profiles of the vertical wave number squared (m^2) and OH $1.6 \mu m$ volume emission rates over the (a) Lhasa, (b) Xining, and (c) Xinglong stations. The profiles of m^2 are derived from temperature profiles (as shown in Fig. 8a–b) and SD-WACCM-X wind fields (as shown in Fig. 8c–e). The OH emission profiles are obtained from averaged SABER/TIMED measurements on 19 May 2022 at the following UTC times: (a) 13:48:24 and 13:49:34; (b) and (c) 12:08:29, 12:09:40, 22:57:10, and 22:57:53 (see Fig. 1 for the measurement locations).

3.5 Thermospheric concentric gravity waves by red-line (630 nm) airglow imagers

In addition to the CGWs observed in the stratosphere and MLT region, thermospheric CGWs were also identified in red-line (630.0 nm) airglow observations at the far-field Shuozhou station, as shown in Fig. 10. The observed wave fronts exhibit partial concentric arc structures propagating predominantly northeastward. The horizontal wavelength gradually increases with distance from the inferred source region, indicating characteristics consistent with wave excitation from a localized source. Fitting analysis suggests that the center of the concentric wave structures was approximately located at (31.6°N, 90.4°E).

In contrast, no comparable concentric gravity wave structures propagating toward



the north or northwest were detected in the red-line (630.0 nm) airglow observations at the Korla station. The distinct difference in wave occurrence between Shuozhou and Korla suggests pronounced directional asymmetry in the propagation of the thermospheric CGWs during this event.

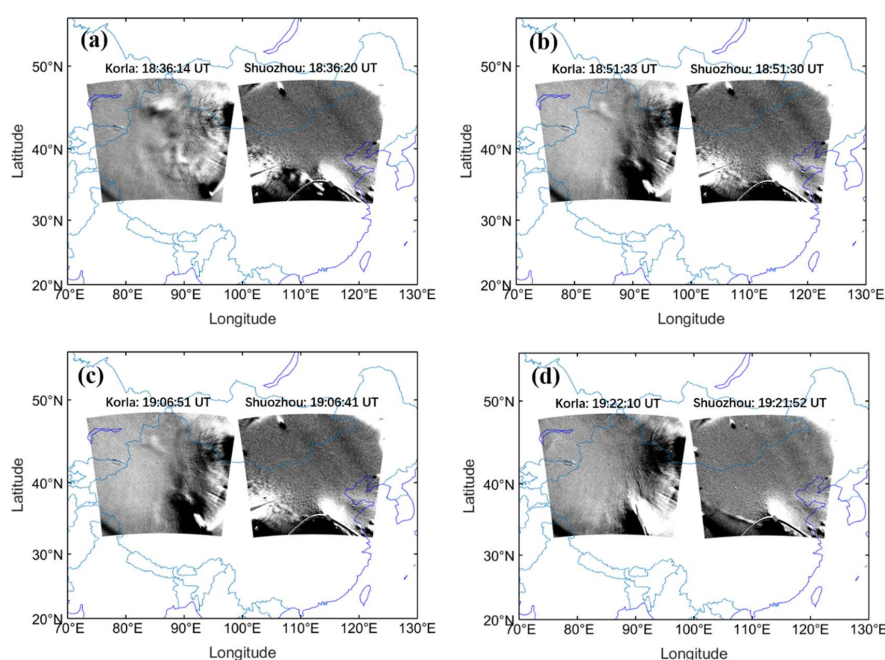


Figure 10 Red-line (630 nm) airglow observations from the Shuozhou and Korla stations. Thermospheric CGWs are clearly observed over Shuozhou, whereas no obvious CGW signatures are detected over Korla. Panels (a)–(d) correspond to 18:36, 18:51, 19:06, and 19:22 UT, respectively.

3.5.1 Potential mechanisms for the asymmetry of concentric gravity waves

Thermospheric background winds induce Doppler shifting that strongly influences the observability of AGWs through two primary mechanisms: cancellation effects within the finite airglow layer thickness and dissipative filtering. When AGWs propagate against the background wind, their intrinsic frequency increases, resulting in



371 longer vertical wavelengths. These longer vertical wavelengths reduce cancellation
372 effects within the airglow layer and are less susceptible to dissipative processes such as
373 molecular viscosity and thermal conduction, thereby enhancing wave detectability at
374 thermospheric altitudes. Conversely, AGWs propagating in the same direction as the
375 background wind experience reduced intrinsic frequencies and shorter vertical
376 wavelengths. Such shortened vertical wavelengths enhance cancellation effects within
377 the airglow layer and increase susceptibility to dissipative damping. As a result, these
378 waves are more readily attenuated before reaching observable amplitudes in airglow
379 measurements.

380 The westward zonal wind over the Xining region simulated by TIE-GCM (17:00-
381 19:00 UT, 19 May 2022) shown in Fig. 11 provides a plausible explanation for the
382 absence of CGW signatures at the Korla station. According to Doppler shift theory,
383 westward-propagating waves traveling with the background wind would experience
384 shortened vertical wavelengths, making them more vulnerable to both cancellation
385 effects and dissipative filtering. In contrast, the northeastward-propagating CGWs
386 observed at Shuozhou propagated partially against the prevailing wind, favoring longer
387 vertical wavelengths and improved observability. These results suggest that directional
388 filtering by thermospheric background winds likely contributed to the asymmetric
389 distribution of the observed thermospheric CGWs.

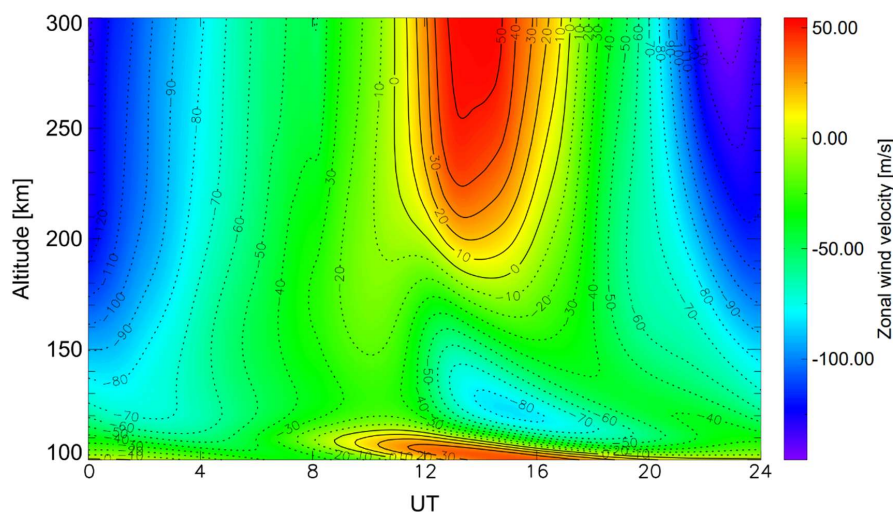


Figure 11 Zonal wind contour from the TIEGCM simulation over Xining station, from 00:00 to 24:00 UT on 19 May 2022.

3.5.2 Possible source of concentric gravity waves in the thermosphere

Figure 12 illustrates the background atmospheric conditions over Shuozhou at 18:00 UT. Specifically, panel (a) shows the temperature profile and the corresponding sound speed, while panel (b) presents the background wind fields. These composite profiles are constructed using SD-WACCM-X data for the altitude range of 0–106 km and TIE-GCM data for 106–300 km. The sound speed, $c_s = \sqrt{\gamma RT}$, near the turbopause is approximately 300 m/s (see Fig. 12a), where γ is ratio of specific heats, R is specific gas constant, and T is temperature.

As demonstrated by Vadas and Crowley (2010), GWs cannot propagate vertically once their horizontal phase speed approaches or exceeds the local sound speed. This constraint arises because internal GWs depend on molecular interactions for their propagation and therefore cannot travel faster than the local sound speed. Under such conditions, the gravity wave transitions toward acoustic-wave behavior, preventing



vertical propagation as an internal gravity wave.

Based on this reasoning, Vadas et al. (2019) further derived the maximum

allowable intrinsic phase speed for gravity wave propagation, denoted as $\max(c_I)$:

$$\max(c_I) \approx \frac{2\sqrt{\gamma-1}}{\gamma} c_s, \quad (6)$$

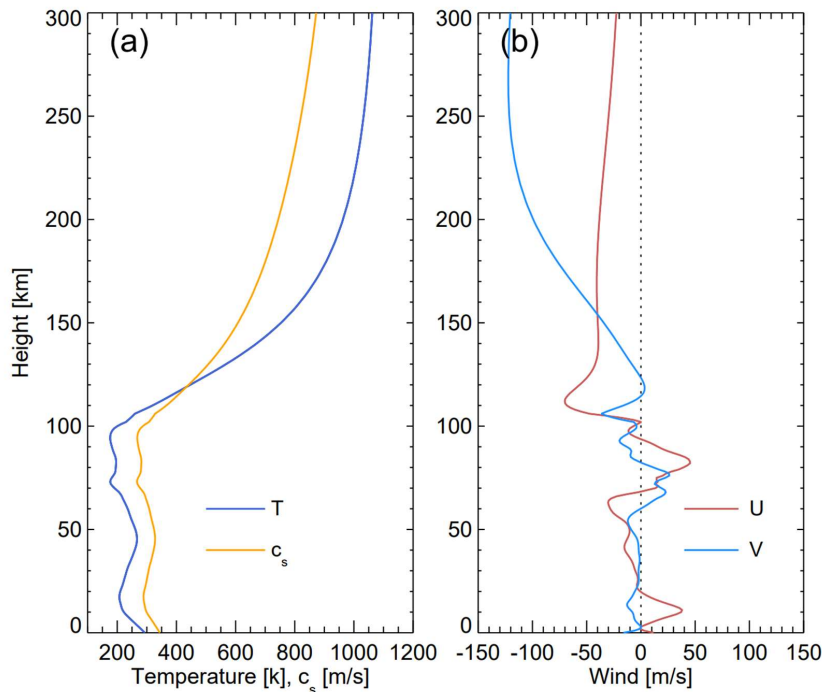


Figure 12 Background atmosphere over Shuozhou at 18:00 UT: (a) Temperature and corresponding sound speed, (b) Background atmospheric wind fields. The composite profiles are constructed using SD-WACCM-X data (0–106 km) and TIE-GCM data (106–300 km).

Based on the information from Fig. 13, which shows $c_I / \left[\frac{2\sqrt{\gamma-1}}{\gamma} c_s \right]$ as a

function of horizontal phase speed c and height for CGW, a gravity wave cannot

propagate when $c_I / \left[\frac{2\sqrt{\gamma-1}}{\gamma} c_s \right] \geq 1$. The thermosphere GWs observed at Shuozhou



station with $c > 250$ – 270 m/s must originate in the thermosphere and therefore cannot be primary GWs generated by deep convection.

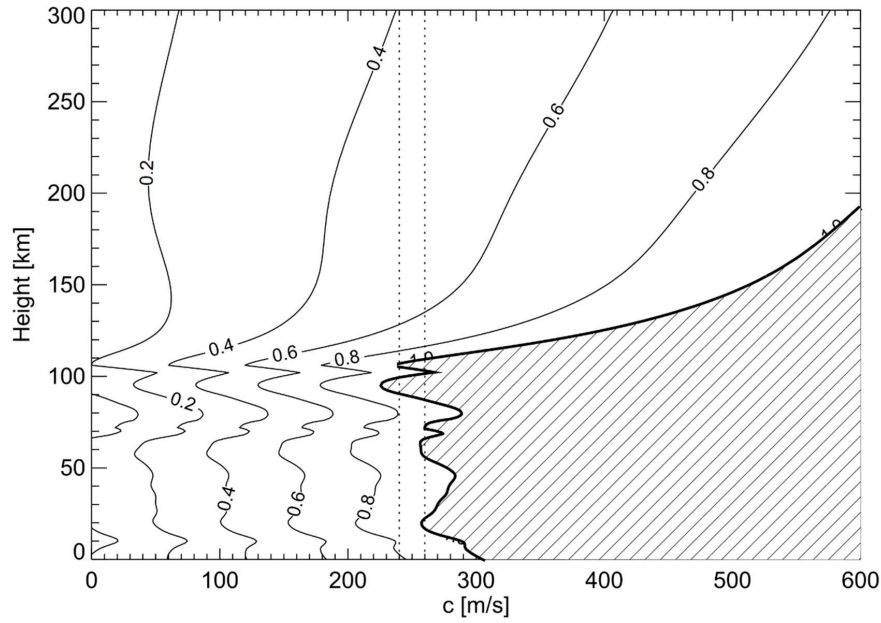


Figure 13 $c_l / \left[\frac{2\sqrt{\gamma-1}}{\gamma} c_s \right]$ as a function of c and H at 18:00 UT. The hatched regions show where $c_l / \left[\frac{2\sqrt{\gamma-1}}{\gamma} c_s \right] \geq 1$. The velocity range between the dashed lines represents the observed phase velocities of thermospheric CGWs.

We employed a ray-tracing method to estimate the source location of the thermospheric CGWs. The model is based on a dispersion relation that incorporates molecular viscosity and thermal diffusivity (Vadas, 2007), as given by Eq. (4). By solving this dispersion relation, we traced the propagation paths of gravity waves backward in order to identify their potential source regions in the upper atmosphere.

$$m^2 = \frac{k_H^2 N^2}{\omega_{lr}^2 (1 + \delta_+ + \delta^2 / \text{Pr})} \left[1 + \frac{\nu^2}{4\omega_{lr}^2} \left(k^2 - \frac{1}{4H^2} \right)^2 \frac{(1 - \text{Pr}^{-1})^2}{(1 + \delta_+ / 2)^2} \right]^{-1} - k_H^2 - \frac{1}{4H^2}, \quad (7)$$



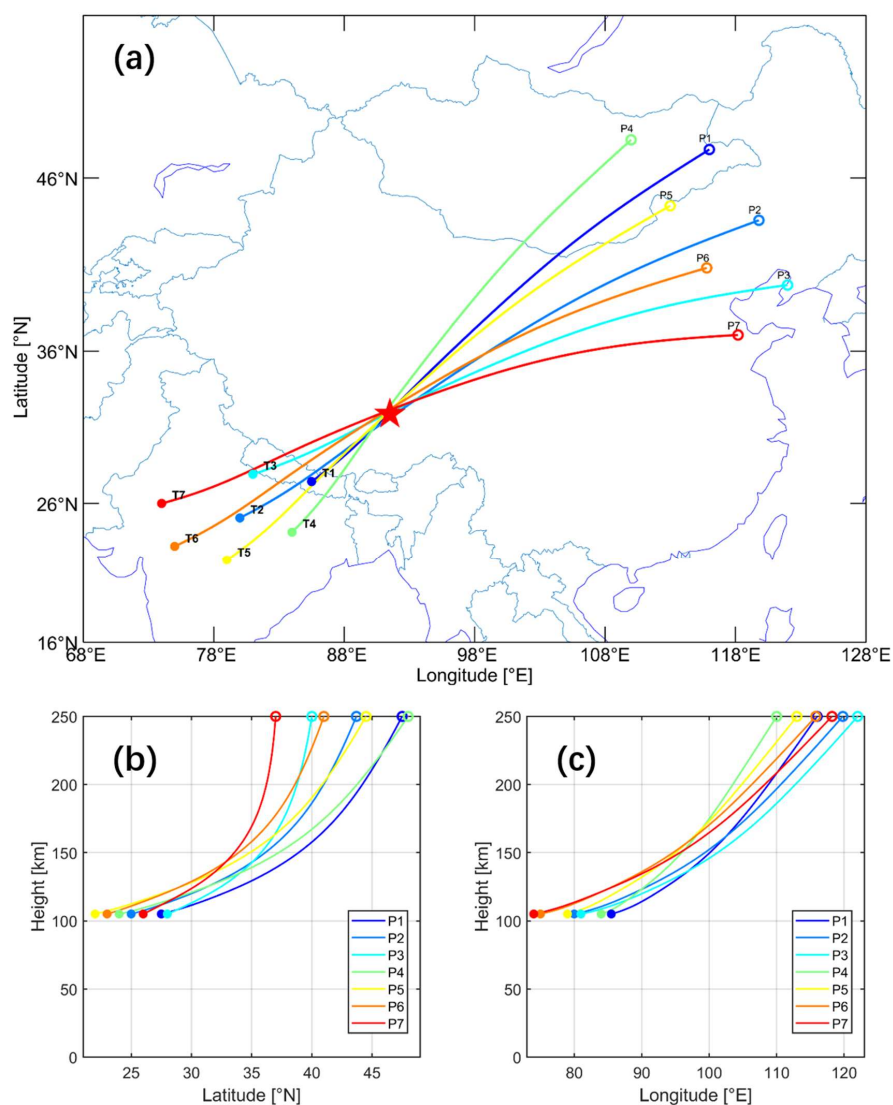
429 where ω_{Ir} is the intrinsic frequency; $k^2 = k_H^2 + m^2$, $k_H^2 = k^2 + l^2$; H is the scale height;
430 $\nu = \mu/\bar{\rho}$ is the kinematic viscosity where μ is the molecular viscosity and $\bar{\rho}$ is the
431 background density; $\delta = \nu m/H\omega_{Ir}$, $\delta_+ = \delta(1 + Pr^{-1})$, where Pr is the Prandtl number. k , l ,
432 and m are the zonal, meridional, and vertical wave number components of the GW,
433 respectively. The density $\bar{\rho}$ were obtained from the NRLMSISE-00 model (Picone et
434 al., 2002).

435 Figure 14 shows the results of the reverse ray tracing. Figure 14a presents the
436 latitude–longitude distribution of the rays, where p1–p7 (circles) represent the starting
437 points of the reverse ray tracing and T1–T7 (dots) denote the corresponding ray
438 endpoints. Figures 14b and 14c show the altitude variations of the ray trajectories as
439 functions of latitude and longitude, respectively. Although gravity wave generation
440 may occur anywhere along an individual ray path (Vadas and Azeem, 2021), the reverse
441 ray trajectories exhibit clear convergence near (32.2°N, 91.7°E) at an altitude of
442 approximately 150 km, suggesting a preferred source region for the observed
443 thermospheric CGWs in the lower thermosphere.

444 Primary GWs generated by deep convection typically undergo substantial
445 amplitude growth and dissipation as they propagate upward into the thermosphere
446 because of the rapidly decreasing atmospheric density. As these waves approach critical
447 levels or dissipative regimes, they can deposit momentum and energy into the
448 background atmosphere, thereby generating secondary GWs with distinct propagation
449 characteristics and phase speeds (Vadas et al., 2003; Vadas and Azeem, 2021). These
450 secondary waves are capable of propagating efficiently within the thermosphere and



451 can exhibit substantially larger horizontal phase speeds than the original primary waves.
452 The inferred source altitude near 150 km is therefore consistent with secondary wave
453 generation associated with the dissipation and breaking of upward-propagating primary
454 gravity waves.



455
456 **Figure 14** Reverse ray-tracing results for the thermospheric gravity wave observed at 250 km



457 altitude. (a) Horizontal propagation paths projected onto the geographic map. (b) Altitude–latitude
458 distributions of the ray trajectories. (c) Altitude–longitude distributions of the ray trajectories. The
459 ray paths converge toward a preferred source region near 150 km altitude.

460 **4 Conclusions**

461 Severe thunderstorms are important sources of energy and momentum for the
462 middle and upper atmosphere through electrical and dynamical pathways. Electrically,
463 intense thunderstorms generate TLEs, including red sprites, gigantic jets, and ghosts,
464 representing strong lightning activity extending into the middle atmosphere.
465 Dynamically, deep convection excites AGWs, which transport energy and momentum
466 upward and can generate secondary waves through dissipation and breaking. Together,
467 these processes provide important pathways for vertical atmospheric coupling.
468 However, simultaneous observations of electrical and dynamical responses of extreme
469 convective systems across multiple atmospheric layers remain rare.

470 To investigate the atmospheric response to this extreme convective event, we
471 combined ground-based optical imaging, amateur photography, satellite observations,
472 and numerical simulations. Multi-layer airglow observations characterized CGWs in
473 the mesosphere and thermosphere through OH, OI 557.7 nm, and OI 630.0 nm
474 emissions. Satellite observations from AIRS, TIMED/SABER, and Meteosat/SEVIRI,
475 together with SD-WACCM-X and TIE-GCM simulations, provided information on
476 wave activity, atmospheric conditions, and propagation environments.

477 The electrical response was characterized by an extraordinary TLE outbreak
478 observed on 19 May 2022 from Lake Pumoyongcuo, southern Tibet. Time-lapse
479 photography recorded more than 100 red sprites, 16 secondary gigantic jets, and at least



480 four ghosts between 13:00 and 16:00 UT, representing the largest documented red sprite
481 outbreak observed from China and one of the most remarkable TLE events in the
482 Himalayan–Bay of Bengal region. The simultaneous occurrence of multiple TLE types
483 indicates strong electrical coupling between the storm and the middle atmosphere.

484 The dynamical response from the stratosphere to the mesosphere–lower
485 thermosphere (MLT) was identified through satellite and ground-based CGW
486 observations. AIRS and airglow imagers revealed concentric gravity-wave structures,
487 with MLT CGWs propagating northeastward for more than 3000 km under favorable
488 ducting conditions. These observations highlight the importance of atmospheric
489 background conditions in sustaining long-distance gravity-wave propagation.

490 Thermospheric CGWs observed in the OI 630.0 nm layer exhibited horizontal
491 phase speeds of 260–280 m s⁻¹. Reverse ray tracing suggested a source altitude near
492 150 km, indicating that these waves were likely secondary gravity waves generated by
493 dissipation and breaking of primary waves. Their asymmetric distribution was likely
494 controlled by thermospheric winds through Doppler-induced vertical wavelength
495 changes and dissipative filtering. These results provide rare observational evidence of
496 multi-step gravity-wave coupling from the lower atmosphere to the thermosphere
497 during an extreme convective event.

498

499 **Data availability.** The multi-layer airglow network data are available from the Chinese
500 Meridian Project Data Center at <https://dcstatus.meridianproject.ac.cn> (MPDC, 2026).
501 TIMED/SABER temperature and emission intensity data are available from the



502 SABER data archive at <http://saber.gats-inc.com/data.php> (Russell et al., 1999).
503 MSG/SEVIRI data are available from the EUMETView service of the European
504 Organisation for the Exploitation of Meteorological Satellites (EUMETSAT) at
505 <https://user.eumetsat.int/data-access/eumetview> (EUMETSAT, 2026). AIRS Level 1B
506 radiance data are available from the NASA Goddard Earth Sciences Data and
507 Information Services Center (GES DISC) at <https://doi.org/10.5067/YZEXEVN4JGGJ>
508 (AIRS Project, 2007). The SD-WACCM-X and TIE-GCM 3.0 simulation outputs
509 generated in this study are available from the corresponding author upon reasonable
510 request.

511

512 **Author contributions.** QL conceived the study, conducted the data analysis,
513 interpreted the results, and wrote the manuscript. OAV contributed to the acquisition
514 and analysis of the SEVIRI data. AA contributed to the processing and interpretation of
515 the red sprite imagery. XW performed the TIE-GCM 3.0 simulations and contributed
516 to their interpretation. SG performed the SD-WACCM-X simulations and contributed
517 to the analysis of the model results. JX, WY, and YZ contributed to the interpretation
518 of the results and manuscript preparation. All authors discussed the results, revised the
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520

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523



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