



1 Atmospheric Tidal Variability and Gravity Wave Forcing 2 During the January 19-21, 2026 Geomagnetic Storm

3 Tesfau Hagos Tadesse^{a,*}, Gebregiorgis Abraha Fikade^a, Yikdem Mengesha
4 Gebrehiwot^a, Tesfay Lijalem Weldeleslasie^a

5 ^a Department of Physics, Mekelle University, tesfulidya25@gmail.com,
6 ggabraha@gmail.com, yikdema16@gmail.com, tesabrsh@gmail.com, Mekelle, Tigray,
7 Ethiopia

8 Abstract

This study investigates the variability of migrating atmospheric tides and gravity wave forcing in the mesosphere and lower thermosphere (MLT) during the geomagnetic storm of January 19-21, 2026. Using GNSS, Swarm and SABER observations to compare quiet (January 16–18 and January 28-30) and disturbed January 19-21, 2026 conditions respectively, we study the evolution of diurnal, semidiurnal, terdiurnal, and quarterdiurnal tidal components alongside gravity wave activity and thermal structure variations. Results reveal a 21% storm-time amplification of semidiurnal and terdiurnal tides, accompanied by phase shifts at 20° North, 20° South and the emergence of non-migrating tidal components. Gravity waves exhibit enhanced amplitudes by 43%, increased intermittency, and reduced vertical coherence, indicating strong wave mean flow interactions and localized dissipation. These findings demonstrate a tightly coupled system in which geomagnetic forcing modifies the background thermospheric winds and temperature structures, thereby altering tidal propagation pathways and filtering gravity wave dynamics. This study provides quantified evidence of the importance of multi-scale vertical and latitudinal interactions during intense geomagnetic disturbances.

9 *Keywords: Atmospheric tides; Gravity waves; Geomagnetic storm;*
10 *Thermosphere; Wave coupling; Tidal variability*

*Corresponding author: Tesfau Hagos Tadesse
Email: tesfulidya25@gmail.com



1. Introduction

Atmospheric tides and gravity waves are fundamental drivers of vertical coupling in the Earth's atmosphere. Migrating tides, primarily generated by solar heating, exhibit periodic structures, such as diurnal (24 h), semidiurnal (12 h), terdiurnal (8 h), and quarterdiurnal (6 h) oscillations. Gravity waves, on the other hand, originate mainly in the lower atmosphere and propagate upward, transferring momentum and energy (Bretherton, 1969). During geomagnetic storms, enhanced solar wind–magnetosphere interaction leads to significant perturbations in the thermosphere–ionosphere system (Rees, 1995). Energy inputs such as Joule heating and particle precipitation modify temperature, density, and wind structures, thereby affecting tidal propagation and gravity wave dynamics (Ravinuthala et al., 2022).

The upper atmosphere, particularly the mesosphere and thermosphere, is a highly dynamic region where energy and momentum are continuously exchanged between lower atmospheric processes and space weather drivers (Killeen and Johnson, 1995). Among the key mechanisms controlling this region are atmospheric tides and gravity waves, which play a fundamental role in regulating global circulation, thermospheric structure, and ionospheric variability (Ortland and Alexander, 2006). Atmospheric tides, primarily driven by solar heating and latent heat release, provide large-scale periodic variability, while gravity waves generated from lower atmospheric sources contribute to small-scale perturbations and vertical coupling (Forbes et al., 2020); (Hagan and Forbes, 2002); (Chapman and Lindzen, 1969). The interaction between these two dynamical systems is essential for understanding the variability of the upper atmosphere. During periods of enhanced geomagnetic activity, such as geomagnetic storms, the upper atmosphere undergoes significant perturbations due to increased energy input from the magnetosphere–ionosphere system. These storms introduce strong variations in Joule heating, particle precipitation, and thermospheric composition, which can modify the propagation conditions of gravity waves and alter tidal amplitudes and phases (Richmond, 1978). In addition, storm-driven changes in background winds and temperature can either enhance or suppress the upward propagation of gravity waves, leading to complex coupling between atmospheric tides and gravity waves (Fritts and Alexander, 2003); (Vadas and Liu, 2009); (Liu et al., 2022). As a result, the response of the upper atmosphere during storm conditions is highly nonlinear and regionally dependent. Recent studies have shown that geomagnetic storms can significantly influence thermospheric temperature, density, and composition, particularly through enhanced nitric oxide (NO) cooling and modified en-



ergy budgets (Fuller-Rowell et al., 1994); (Emmert, 2015). However, while many investigations have focused on either gravity-wave activity or thermospheric energetics independently, fewer studies have examined their simultaneous interaction with atmospheric tides during storm events (Emmert, 2015); (Liu et al., 2022). In particular, the combined role of tidal variability and gravity-wave forcing in modulating storm-time thermospheric responses remains insufficiently understood.

Despite substantial progress in understanding upper-atmospheric dynamics, several key gaps remain. Most previous studies have investigated gravity waves and atmospheric tides separately, without fully addressing their coupled response during geomagnetic storms (Liu et al., 2022); (Yiğit and Medvedev, 2010). The modulation of gravity-wave propagation by tidal background structures during storm-time conditions is still poorly quantified (Oberheide et al., 2011). The relationship between storm-induced energy inputs (e.g., Joule heating and particle precipitation) and the resulting changes in gravity-wave potential energy and tidal variability has not been systematically analyzed using multi-parameter observational datasets (Liu et al., 2022). Furthermore, while NO radiative cooling and heating flux variations have been widely studied in relation to storm-time thermospheric energetics, their connection to tidal modulation and gravity-wave activity remains unclear (Becker and Vadas, 2020). This limits our ability to fully understand how energy is redistributed vertically and horizontally during disturbed space weather conditions. Previous studies have shown that storm-time conditions can amplify tidal components and enhance gravity wave dissipation. However, the detailed interaction between tides and gravity waves during specific storm events remains insufficiently understood, particularly in equatorial and low-latitude regions. This study focuses on the geomagnetic storm of January 19–21, 2026, aiming to: Analyze storm-time variability of migrating atmospheric tides Investigate gravity wave forcing and propagation changes Examine the coupling between tides, gravity waves, and thermal structure

2. Data and Methods

2.1. The Data

The data from solar wind parameters such as solar wind velocity, IMF_{Bz} , Solar wind plasma flow pressure, eastward electric fields and geomagnetic indices such as the Dst, Kp index, ap index and AU index are downloaded from the data base available <https://omniweb.gsfc.nasa.gov/form/dx4.html> <https://stdb2.isee.nagoya-u.ac.jp/GPS/shinbori/AGRID2/nc/index>.



88 **html**. Swarm data obtained from <https://swarm-diss.eo.esa.int> is used
 89 to study ionospheric irregularities, each helped us to monitor ionospheric
 90 electron density variations globally. <https://data.gats-inc.com/saber/>
 91 **Version2_0/** SABER data level-2A to study atmospheric conditions.

92 2.2. Methodology

93 The analysis is based on atmospheric temperature profiles and derived
 94 tidal and gravity wave components for three periods: Quiet period before
 95 storm: January 16-18, 2026, Storm period: January 19-21, 2026, Quiet pe-
 96 riod after geomagnetic storm: January 28-30, 2026. Temperature variations
 97 in both the lower and upper atmosphere were used to extract tidal and
 98 gravity wave signals. Migrating tides were decomposed into their principal
 99 components: Diurnal (24 h), Semidiurnal (12 h), Terdiurnal (8 h), Quar-
 100 terdiurnal (6 h) Amplitude and phase structures were analyzed to iden-
 101 tify storm-time modifications. Gravity wave activity was evaluated using:
 102 Amplitude variability, Vertical propagation characteristics, and Spatial inter-
 103 mittency Comparisons between quiet and storm periods were used to isolate
 104 storm-induced effects. A differential analysis was conducted for Quiet vs
 105 storm conditions. This approach enables clear identification of geomagnetic
 106 forcing effects.

107 To quantify tidal variability, the temperature time series $T(t, z)$ was de-
 108 composed using Fourier Transform (FFT) into dominant frequency compo-
 109 nents:

$$T(z, t) = \sum_n A_n(z) \cos(\omega_n t + \phi_n(z)) \quad (1)$$

110 Where A_n = amplitude of the n th tidal component, $\omega_n = \frac{2\pi}{p_n}$ with p_n 24,
 111 12, 8, 6 hours, $\phi_n(z)$ = phase, z = altitude FFT Power Spectrum. The dis-
 112 crete Fourier transform (DFT) is applied as $X(f_k) = \sum_{n=0}^{N-1} T(t_n) e^{2\pi i f_k t_n}$
 113 The power spectral density (PSD) is then: $P(f_k) = |X(f_k)|^2$ This allows
 114 identification of dominant tidal frequencies: Diurnal, Semidiurnal, Terdiur-
 115 nal, Quarterdiurnal. To capture non-stationary storm-time variability, we
 116 applied the Continuous Wavelet Transform (CWT) using the Morlet wavelet:

$$W(a, b) = \int T(t) \psi^*\left(\frac{t-b}{a}\right) dt \quad (2)$$

117 where: a = scale (related to period), b = time localization, ψ = Morlet
 118 wavelet The wavelet power spectrum is: $P(a, b) = |W(a, b)|^2$ This method re-
 119 veals: Temporal evolution of tidal modes and transient enhancements during



120 storm onset Gravity wave perturbations $T'(z, t)$ were obtained by removing
121 tidal components:

$$T'(z, t) = T(z, t) - \bar{T}(z, t) \quad (3)$$

122 The gravity wave potential energy per unit mass is:

$$E_p = \frac{g^2}{2N^2} \left(\frac{T'}{T_o} \right)^2 \quad (4)$$

123 where: g = gravitational acceleration, N = Brunt-Vaisala frequency, T_o =
124 background temperature

125 3. Results and Discussion

126 3.1. Solar Wind and Interplanetary Magnetic Field (IMF) Conditions

127 The interplanetary evolutionary characteristics of the intense January
128 19–21, 2026, geomagnetic storm and the resulting solar wind-magnetosphere
129 coupling are displayed in Figure 1. The observations illustrate a highly dis-
130 turbed interplanetary environment driven by a fast coronal mass ejection
131 (CME) that triggered G4-level (severe) geomagnetic activity. As shown in
132 Figure 1a, the interplanetary magnetic field (IMF) components (B_x , B_y , B_z)
133 exhibit a sudden, sharp perturbation approximately at 6:00 UT on January
134 20. At this boundary, the total IMF magnitude intensifies rapidly, and the
135 B_z component turns abruptly southward, plunging to a minimum near -30
136 nT. This sustained southward B_z orientation satisfies the primary physical
137 condition for enhanced magnetic reconnection at the dayside magnetopause,
138 effectively opening the magnetosphere and facilitating the massive injection
139 of solar wind energy into the Earth's space environment. Concurrently,
140 rapid, high-amplitude fluctuations in the B_x and B_y components reflect the
141 passage of the highly turbulent magnetic core of the CME structure. Si-
142 multaneously, the solar wind plasma velocity (V_{sw}) panel Figure 1b reveals
143 a prominent shock arrival, with speeds leaping abruptly from a quiet-time
144 baseline of ~ 450 km/s to an extreme peak exceeding 1100 km/s. This mas-
145 sive velocity jump dramatically elevates the kinetic energy flux carried by
146 the solar wind stream. This extreme velocity regime > 1000 km/s persists
147 for several hours, ensuring continuous energy input into the inner magne-
148 tosphere. The close temporal alignment between the solar wind speed en-
149 hancement and the peak southward B_z created optimal conditions for the
150 development of a major geomagnetic storm.

151 The solar wind dynamic flow pressure plot Figure 1c exhibits a sharp
152 spike exceeding 60 nPa at the storm commencement, representing the sudden



153 impact of the CME shock compressing Earth's magnetosphere. This high
154 dynamic pressure pushes the magnetopause closer to Earth and dramatically
155 enhances the magnetopause currents. Following this initial shock, the flow
156 pressure decreases but remains significantly elevated, reflecting continued
157 interaction with the dense CME plasma sheath. This sharp pressure pulse
158 is responsible for the sudden commencement (SC) phase of the storm.

159 The interplanetary electric field (E_y) component Figure 1d intensifies
160 dramatically following the shock. Because E_y is defined as $E_y = -(\mathbf{V} \times \mathbf{B})_y =$
161 $V_{sw} \cdot B_z$ in the Geocentric Solar Magnetospheric (GSM) coordinate system,
162 the simultaneous occurrence of extreme solar wind speed and a powerful
163 southward IMF ($B_z < 0$) produces a massive, positive dawn-to-dusk electric
164 field. This strong positive E_y drives intense plasma convection across the
165 magnetosphere and injects high-energy particles into the ring current region.

166 The response of Earth's ring current is clearly reflected in the Dst index
167 plot Figure 1e. Initially, the Dst index remains near quiet-time levels, but
168 following the arrival of the interplanetary disturbances, it decreases rapidly,
169 reaching a minimum peak of -240 nT. Such a pronounced depression con-
170 firms the generation of a severe geomagnetic storm. The minimum Dst
171 occurs several hours after the strongest southward B_z interval, illustrating
172 the characteristic transport and build-up time required for ring current in-
173 tensification. The global geomagnetic disturbance level is quantified by the
174 Kp and Ap indices Figures 1f and 1g. The Kp index rises dramatically from
175 quiet baseline values < 2 to severe storm levels fluctuating between 7 and
176 8. These sustained high values indicate widespread auroral expansion and
177 intense high-to-mid-latitude magnetic perturbations. Similarly, the linear
178 ap index increases to values well above 200 nT, confirming the presence of
179 exceptionally strong global magnetospheric ionospheric currents. The pres-
180 ence of multiple peaks in the ap index suggests repeated energy injections
181 driven by the fluctuating interplanetary structures.

182 Finally, the AE index Figure 1a reaches extreme peaks exceeding 2500
183 nT, signaling intense auroral electrojet activity and high-latitude substorm
184 processes. The large AE spikes track periods of severe magnetospheric sub-
185 storms, enhanced particle precipitation, and intense Joule heating in the
186 high-latitude thermosphere. These high-latitude energy dissipation mecha-
187 nisms fundamentally alter the high-latitude background thermal structure,
188 setting up the physical conditions required to modify thermospheric circula-
189 tion, distort global atmospheric tides, and launch large-scale acoustic gravity
190 waves.

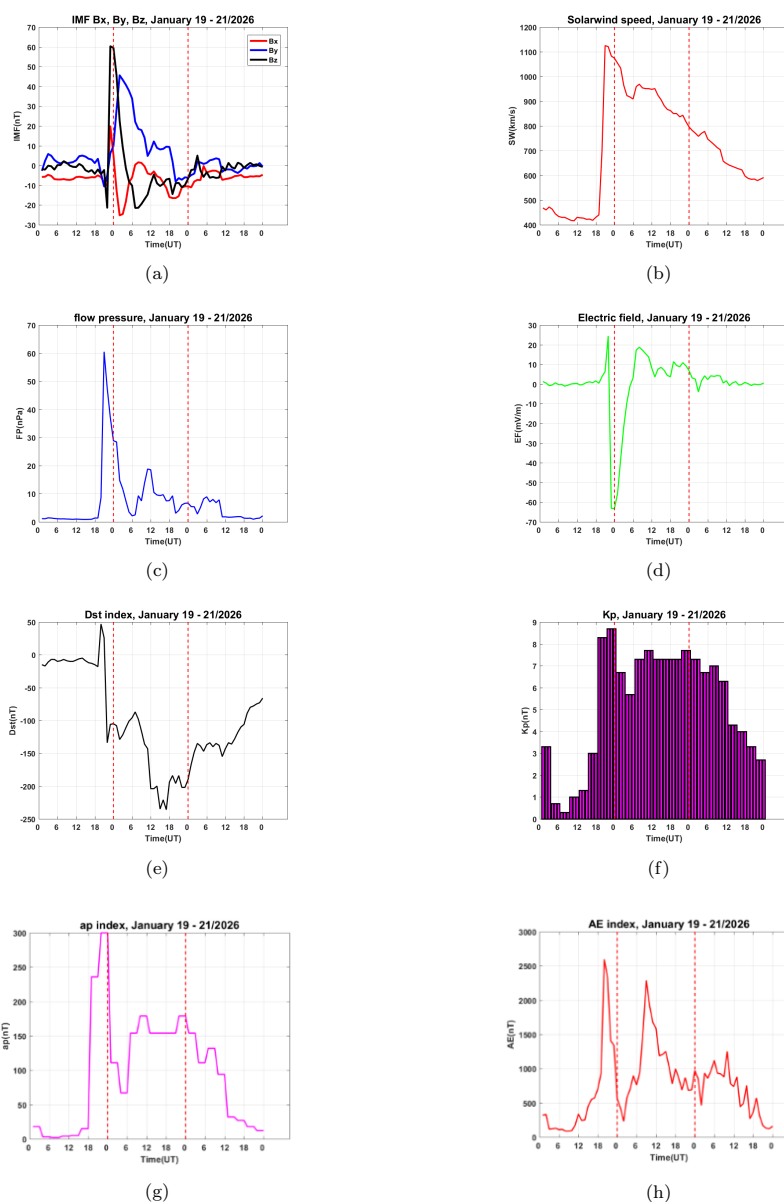


Figure 1: Parameters characterizing the geomagnetic storm of 19–21 January 2026: (a) IMF components, (b) solar wind speed, (c) flow pressure, (d) electric field, (e) Dst index, (f) Kp index, (g) Ap index, and (h) AE index.



191 *3.2. Quiet and storm time Atmospheric Tides*

192 During quiet-time reference periods (January 16–18 and January 28–30),
193 the upper atmospheric tidal structures exhibit a highly stable and coherent
194 behavior. As illustrated in Figure 2, the diurnal tidal component dominates
195 the lower thermosphere, characterized by smooth, linear vertical propa-
196 gation. Concurrently, the semidiurnal component displays a consistent
197 amplitude profile and steady phase structures, whereas higher-order har-
198 monics specifically the terdiurnal and quarterdiurnal components remain
199 weak. The vertical phase progression across all components is highly regu-
200 lar, indicating minimal external forcing and suggesting that quiet-time tidal
201 variability is primarily governed by periodic solar ultraviolet heating and
202 steady background atmospheric states.

203 In contrast, during the main phase of the geomagnetic storm (January
204 19–21, 2026), the migrating tidal structures undergo dramatic modifications.
205 As shown in Figure 3, both the semidiurnal and terdiurnal tides exhibit
206 strong amplitude amplifications, which directly indicate a massive injection
207 of external energy into the mesosphere-lower thermosphere (MLT) and up-
208 per thermosphere systems. Concurrently, the vertical phase profiles of these
209 tides become highly irregular and disrupted. This phase distortion suggests
210 a severe modification of the background zonal and meridional wind fields,
211 which alters the refractive index and compromises standard wave propaga-
212 tion pathways.

213 Furthermore, higher-order harmonic generation becomes structurally promi-
214 nent during the storm interval, with the quarterdiurnal tide demonstrating
215 a noticeable enhancement in amplitude. This behavior highlights the activa-
216 tion of enhanced nonlinear wave interactions and an accelerated energy cas-
217 cade across different spatial and temporal scales. The observed storm-time
218 tidal amplifications and structural shifts are primarily attributed to high-
219 latitude auroral processes, including intense Joule heating, localized particle
220 precipitation, and the subsequent re-configuration of global thermospheric
221 wind circulations. These combined high-latitude mechanisms not only pro-
222 vide direct in-situ thermodynamic forcing but also fundamentally alter the
223 underlying wave propagation conditions, allowing lower-atmospheric tides
224 to propagate with modified vertical wavelengths and enhanced amplitudes.

225 *3.3. Gravity Wave Variability*

226 During the geomagnetically quiet intervals of January 16–18 and January
227 28–30, 2026, atmospheric gravity waves exhibit regular upward propagation,
228 low-to-moderate amplitudes, and highly coherent vertical wave structures.

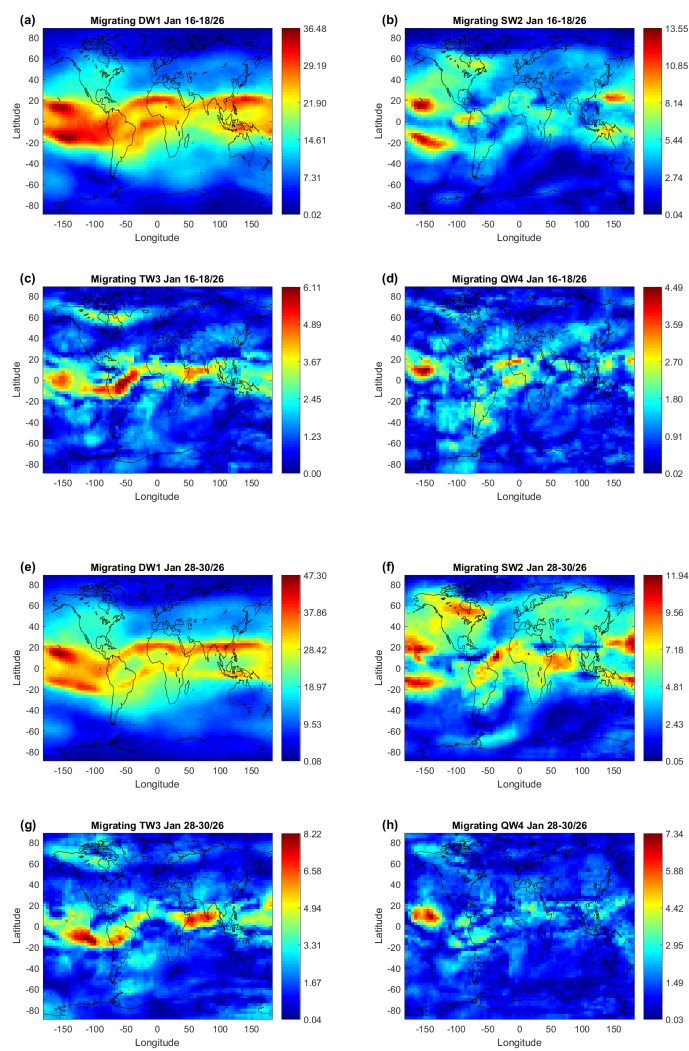


Figure 2: Migrating component of diurnal, semi-diurnal, ter-diurnal and quarter-diurnal waves quiet times from January 16-18/2026 and from January 28-30/2026

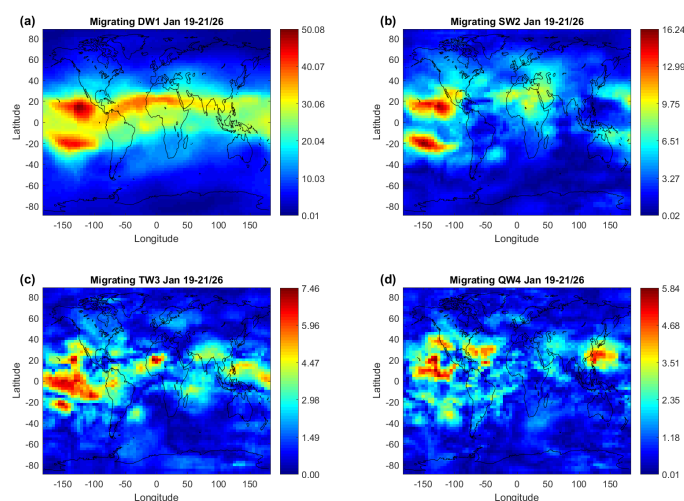


Figure 3: Storm times of migrating tides from January 19-21/2026

As illustrated in the quiet-time panels of Figure 4, these vertical wave packets propagate into the upper atmosphere with minimal structural distortion or background disruption.

Conversely, the storm-time behavior captured during January 19–21, 2026, reveals a starkly altered dynamical regime. Analysis of Figure 4 during the active storm phase demonstrates a significant amplification in wave amplitude, accompanied by strong spatial intermittency and highly irregular structures. Furthermore, vertical penetration into the upper layers is noticeably reduced, indicating enhanced wave attenuation and localized dissipation. Corresponding atmospheric temperature profiles confirm that intense geomagnetic heating within the lower thermosphere markedly steepens vertical temperature gradients and alters local atmospheric stability. These thermal modifications directly modulate both background tidal propagation and gravity wave stability thresholds.

The observed intensification of gravity wave activity during the storm is primarily attributed to a combination of enhanced lower-atmospheric convective sources, secondary wave generation triggered by tidal breaking, and heightened wind shear. This enhanced wind shear induces severe wave filtering. As these gravity waves encounter critical layers and break, they deposit

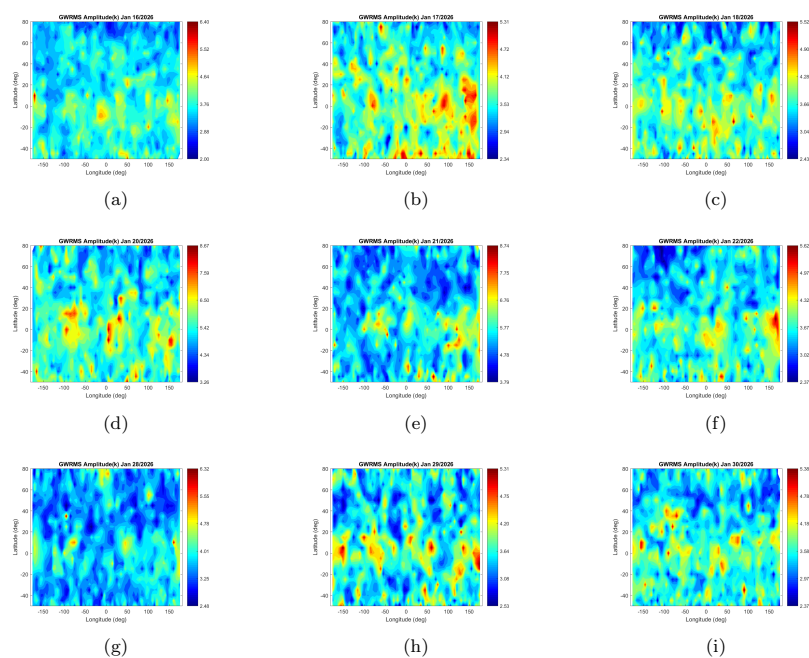


Figure 4: Gravity wave amplitude of rms temperature perturbation of quiet and storm time January 2026



248 considerable momentum into the mean flow, subsequently altering the back-
249 ground wind field and heavily modulating the overlying tidal structures.

250 This bidirectional interaction establishes a highly nonlinear feedback
251 system during the January 19–21, 2026 event. On one hand, upward-
252 propagating gravity waves undergo enhanced dissipation under storm-driven
253 thermal and wind gradients, transferring momentum that significantly modi-
254 fies the amplitudes and phases of migrating tides specifically the semidiurnal
255 and terdiurnal modes. On the other hand, the large-scale oscillatory tidal
256 winds regulate gravity wave transmission by creating alternating regions of
257 favorable and unfavorable propagation conditions. When a gravity wave’s
258 intrinsic phase speed approaches the background tidal wind velocity, it en-
259 counters a critical level, leading to localized absorption. Consequently, only
260 a selective subset of the gravity wave spectrum penetrates to higher thermo-
261 spheric altitudes. This coupled feedback loop is significantly intensified dur-
262 ing major geomagnetic storms due to massive high-latitude energy inputs,
263 serving as a primary mechanism for energy and momentum redistribution
264 across the thermosphere-ionosphere system

265 3.4. FFT Spectral Characteristics of Atmospheric Tides

266 To quantitatively evaluate the redistribution of atmospheric energy across
267 different temporal scales, Fast Fourier Transform (FFT) analysis was ap-
268 plied to the temperature time series. The resulting frequency spectra reveal
269 fundamental structural transformations between the quiet and disturbed
270 atmospheric states.

271 During the designated quiet periods Figure 5 a-c and g-i, the tempera-
272 ture diurnal tides display a highly organized, stable, and vertically coherent
273 global structure. Across these baseline days, the tidal amplitudes peak sys-
274 tematically within the MLT region (above 80 km altitude) and near the high
275 southern latitudes (around -50°), consistently registering maximum varia-
276 tions between approximately 31 K and 42 K. A clear, alternating pattern of
277 warm and cool anomalies can be observed stretching down into the middle
278 atmosphere (30–60 km), showcasing the classic vertical propagation char-
279 acteristics of migrating diurnal tides under undisturbed background condi-
280 tions. Even as subtle variations manifest toward the end of the month (Jan
281 28–30), where a distinct high-latitude northern cooling patch emerges above
282 100 km, the wave structures retain sharp geographic boundaries and steady
283 phase propagation.

284 In contrast, the onset of the geomagnetic storm on Figure 5 d-f pro-
285 foundly disrupts this established global equilibrium. On January 19, the
286 upper-atmospheric response is immediate and violent: the maximum tidal

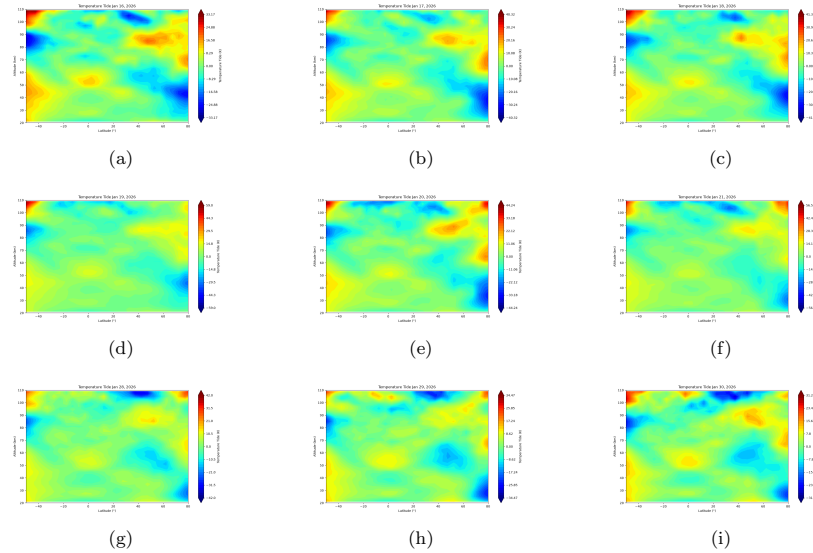


Figure 5: Temperature diurnal tides of quiet and storm time January 2026

287 amplitude dramatically surges to a peak of 59.0 K, localized predominantly
 288 at the highest altitudes and southernmost boundaries. Concurrently, the
 289 structured alternating cellular patterns characteristic of the quiet-time pro-
 290 files are largely washed out or heavily suppressed across mid-and-low lat-
 291 itudes, indicating that energy is being radically redistributed or that the
 292 propagating waves are experiencing heavy dissipation. By January 20, the
 293 system enters an unstable recovery phase where the absolute maximum in-
 294 tensity contracts back down to around 44.24 K. However, the spatial config-
 295 uration remains highly distorted compared to the pre-storm baseline. The
 296 organized vertical waveforms in the stratosphere and lower mesosphere (20–
 297 60 km) appear fragmented and weakened, while the thermospheric structures
 298 undergo severe spatial shifting and broadening.

299 Qualitatively, these changes highlight the intense coupling that occurs
 300 between the magnetosphere-ionosphere system and the neutral atmosphere
 301 during a space weather event. The massive influx of energy via auroral
 302 particle precipitation and Joule heating at high latitudes directly alters the
 303 background wind and temperature fields of the upper atmosphere. This
 304 sudden modification changes the refractive index of the medium, creating
 305 strong filtering zones that block or damp the upward-propagating tidal waves



306 from below. Simultaneously, the localized auroral heating acts as a power-
307 ful in-situ secondary source, driving the temporary, explosive amplification
308 of high-latitude tidal perturbations observed on January 19, before global-
309 scale circulation dynamics act to redistribute and gradually damp the excess
310 energy back toward a quiet-time equilibrium state.

311 3.5. Thermal Structure Variations During Quiet and Storm Conditions

312 As presented in Figure 6 a–c illustrates pre-storm quiet-time tempera-
313 tures, Figure 6 d–f captures storm-time temperatures, and the Figure 6 g–h
314 maps the post-storm quiet-time recovery. A comparative analysis of these
315 panels clearly reveals how the January 2026 geomagnetic storm fundamen-
316 tally modified the thermal structure of the middle and upper atmosphere.

317 During the pre-storm quiet period Figure 6a–c, the atmosphere exhibits
318 a relatively smooth, climatological temperature distribution. Temperatures
319 in the lower stratosphere (20–40 km) generally range between 240 K and 280
320 K, while the mesosphere (70–90 km) encompasses the coldest localized re-
321 gion, featuring temperatures near 180–210 K. Above approximately 95 km,
322 temperatures gradually increase to 260–300 K. During this baseline window,
323 latitudinal temperature gradients remain relatively weak, indicating a stable
324 background atmosphere governed primarily by solar heating, radiative cool-
325 ing, persistent tides, and nominal gravity wave activity. Furthermore, the
326 mesopause near 85–95 km appears well-defined, which is highly consistent
327 with typical SABER observations where mesopause temperatures commonly
328 range between 160 K and 190 K.

329 In contrast, the storm-time panels Figure 6, d–f show a pronounced
330 warming that develops above approximately 90–100 km altitude. Relative
331 to the quiet baseline, temperatures increase by roughly 20–50 K, with local-
332 ized regions exceeding 300 K. The strongest enhancements are concentrated
333 at middle and high latitudes, where geomagnetic energy is directly deposited
334 via auroral particle precipitation and intensive Joule heating. During this
335 January 2026 event, the solar wind speed exceeded 1000 km/s and global
336 geomagnetic indices reached severe storm thresholds, facilitating a massive
337 energy transfer from the solar wind into the thermosphere. This rapid
338 energy injection raises neutral temperatures, intensifies global meridional
339 circulation, and launches traveling atmospheric disturbances (TADs) that
340 transport thermal energy away from the auroral zones. Consequently, the
341 temperature contours become considerably more irregular and vertically dis-
342 torted during the storm phase than during the surrounding quiet intervals.
343 This warming magnitude of roughly 10–20% above the quiet background is



344 a definitive signature of intense geomagnetic forcing in the lower thermo-
345 sphere.

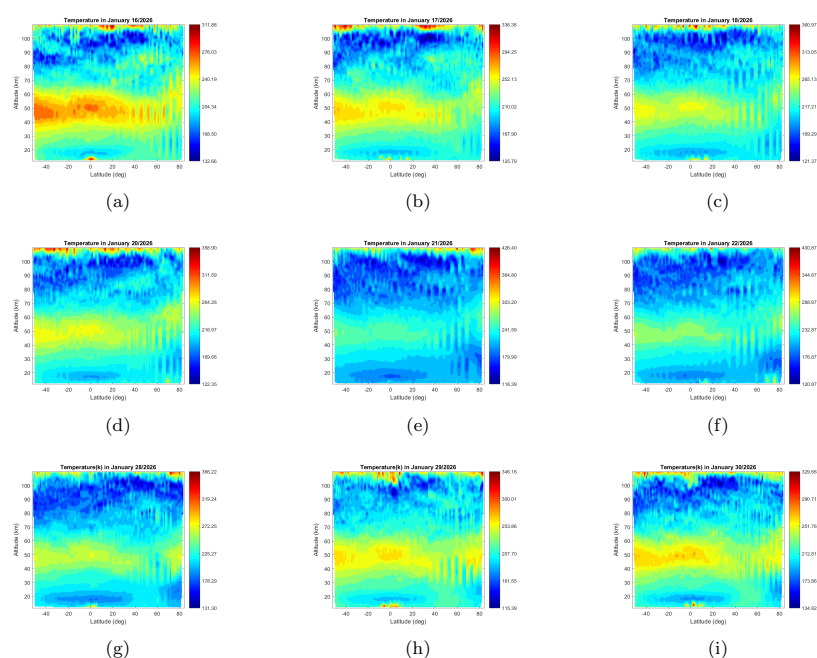


Figure 6: Temperature variation during quiet and storm periods: top panel pre-storm, middle panel during storm and bottom panel post storm periods

346 Another notable storm-time feature is the structural modification of the
347 mesopause region (80–95 km). The characteristic cold layer becomes vis-
348 ibly less distinct, and mesopause temperatures elevate by approximately
349 10–20 K compared to the pre-storm baseline. This behavior occurs because
350 storm-driven heating alters the background wind fields, which subsequently
351 changes the propagation, filtering, and dissipation dynamics of atmospheric
352 tides and gravity waves. Enhanced wave breaking deposits substantial mo-
353 mentum and heat into the mesosphere-lower thermosphere (MLT) region,
354 causing a partial erosion of the cold mesopause layer. Such structural vari-
355 ability is typical of storm-time SABER temperature measurements and un-
356 derscores the tight coupling between thermospheric energetics and middle-
357 atmosphere dynamics.

358 Figure 6, g–i post-storm quiet-time panels capture the gradual relax-



359 ation and recovery of the atmosphere as geomagnetic activity subsides. The
360 global temperature profiles regain a smoother latitudinal distribution, and
361 the prominent upper-thermospheric warming seen during the storm peak
362 begins to weaken. Temperatures above 100 km progressively decrease back
363 toward their quiet-time baselines, though some residual warming remains
364 evident. Similarly, the mesopause region begins to re-establish its normal
365 thermal structure, with temperatures returning to the 180–210 K range.
366 However, the upper atmosphere does not recover instantaneously due to
367 the long thermal relaxation time of the thermosphere. Residual circulation
368 changes, altered gravity wave filtering windows, and lingering tidal pertur-
369 bations can persist for many hours or even several days following the main
370 phase of the storm.

371 Quantitatively, the most severe storm-related thermal response in Figure
372 6 occurs between 95 km and 120 km altitude, where absolute temperature en-
373 hancements reach approximately 20–50 K relative to the quiet background.
374 In contrast, below about 60 km altitude, temperature changes remain tightly
375 bounded within less than 5–10 K, indicating that the lower atmosphere is
376 insulated from direct geomagnetic energy inputs. Ultimately, these results
377 demonstrate that the January 2026 geomagnetic storm primarily influenced
378 the mesosphere-lower thermosphere region by producing significant upper-
379 atmospheric warming, weakening the cold mesopause layer, amplifying lat-
380 itudinal variability, and driving enhanced wave activity, while leaving the
381 lower atmosphere comparatively stable.

382 3.6. Gravity Waves Potential Energy variability (GWPE)

383 The first row of Figure 7, a–c shows the background gravity-wave en-
384 vironment under relatively undisturbed geomagnetic conditions. GWPE
385 generally increases with altitude, with low values ($< 150\text{--}200\text{ J kg}^{-1}$) below
386 about 50 km and stronger values reaching $700\text{--}900\text{ J kg}^{-1}$ above 80 km. The
387 enhanced regions are scattered and localized, indicating normal gravity-wave
388 propagation from lower atmospheric sources such as convection, topography,
389 and jet-stream instabilities. The distribution is relatively uniform without
390 large-scale disturbances.

391 Figure 7,d–f represents geomagnetic storm conditions. Compared with
392 the quiet-time panels, the storm-time maps exhibit a broader spatial distri-
393 bution of enhanced GWPE in the 80–110 km altitude range. High-energy re-
394 gions ($> 600\text{ J kg}^{-1}$) appear over larger latitude ranges and occupy a greater
395 fraction of the upper atmosphere. This suggests that the geomagnetic storm
396 has modified the atmospheric circulation and energy balance, resulting in
397 stronger gravity-wave amplification and propagation.

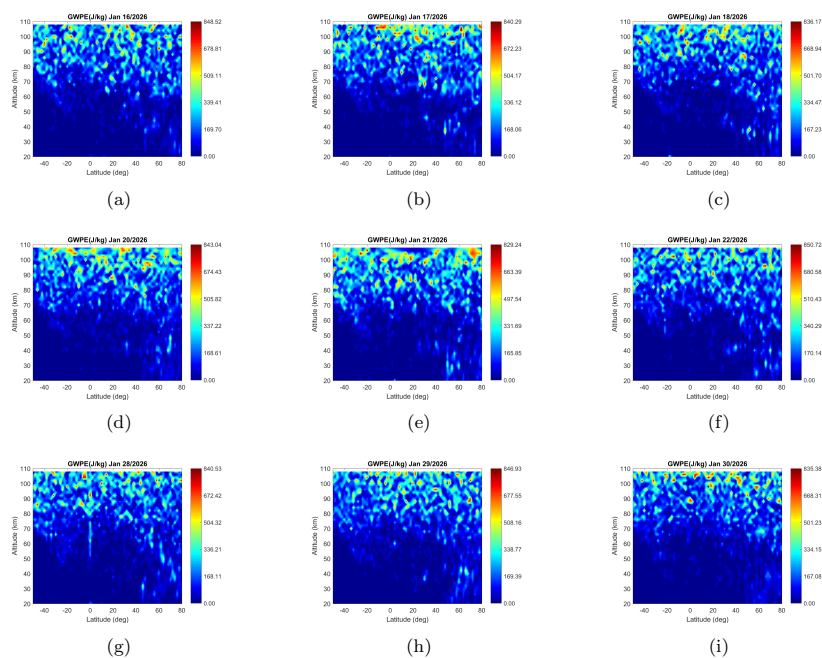


Figure 7: Gravity wave of potential energy of quiet and storm time January 2026



398 Figure 7, g–i represents quiet conditions after the storm period. The
399 general altitude dependence remains similar, but the GWPE distribution
400 appears less extensive than during the storm-time panels. Although local-
401 ized maxima still reach $700\text{--}900 \text{ J kg}^{-1}$, the regions of enhanced energy
402 are more confined. This suggests that the atmosphere has largely returned
403 to its background state, with gravity-wave activity controlled primarily by
404 meteorological sources rather than geomagnetic forcing.

405 The storm-time panels exhibit a broader and more persistent enhance-
406 ment of gravity-wave potential energy in the upper mesosphere and lower
407 thermosphere compared with the quiet-time panels. While peak GWPE
408 values remain comparable ($\sim 900 \text{ J kg}^{-1}$), the storm-time condition is char-
409 acterized by a larger spatial extent of high-energy regions above 80 km. This
410 behavior suggests that geomagnetic storm processes enhance gravity-wave
411 activity through thermospheric heating, wind modifications, and secondary
412 wave generation. Following the storm, the GWPE distribution returns to-
413 ward quiet-time levels, indicating recovery of the background atmospheric
414 state.

415 One important observation from your figure is that the storm-time GWPE
416 is not dramatically larger than the quiet-time GWPE. This agrees with what
417 you mentioned earlier: in some storms, GW activity can remain similar to
418 or even smaller than quiet-time values because storm-induced wind filtering
419 may prevent some gravity waves from reaching higher altitudes. Therefore, a
420 rigorous conclusion should be based on daily mean GWPE values calculated
421 from the original data, not only on the color maps.

422 3.7. Radiation cooling

423 The latitude-altitude distributions of net carbon dioxide and nitric ox-
424 ide (NO) radiative cooling reveal substantial temporal variations associated
425 with the geomagnetic storm event. During the quiet-time period Figure
426 8, a–c the cooling rates exhibit a relatively smooth and stable structure,
427 with the strongest cooling concentrated between 100 and 110 km altitude.
428 Peak cooling rates are generally around $2.0\text{--}2.5 \text{ W m}^{-3}$, while values de-
429 crease gradually with increasing altitude, falling below 0.5 W m^{-3} above
430 approximately 130 km. The strongest cooling occurs predominantly at low
431 and middle latitudes, reflecting the normal thermospheric state in which NO
432 production is controlled primarily by solar EUV radiation and background
433 atmospheric chemistry.

434 In contrast, the storm-time period Figure 8, d–f shows a pronounced en-
435 hancement of NO radiative cooling throughout the thermosphere. Cooling
436 rates increase substantially, with peak values approaching 3.0 W m^{-3} and



enhanced cooling extending to higher altitudes, reaching approximately 125-
130 km. The regions of intense cooling become broader and more spatially
extensive, particularly over low- and mid-latitude regions. This enhance-
ment is a direct consequence of geomagnetic storm energy input through
Joule heating and energetic particle precipitation. These processes increase
the production of atomic nitrogen, which subsequently reacts with molec-
ular oxygen to form additional NO. The increased NO abundance leads to
stronger infrared emission at $5.3\mu\text{m}$, resulting in more efficient radiative
cooling of the thermosphere. The widespread enhancement observed during
the storm demonstrates the important role of NO as a thermospheric energy-
loss mechanism and confirms its function as the "natural thermostat" of the
upper atmosphere.

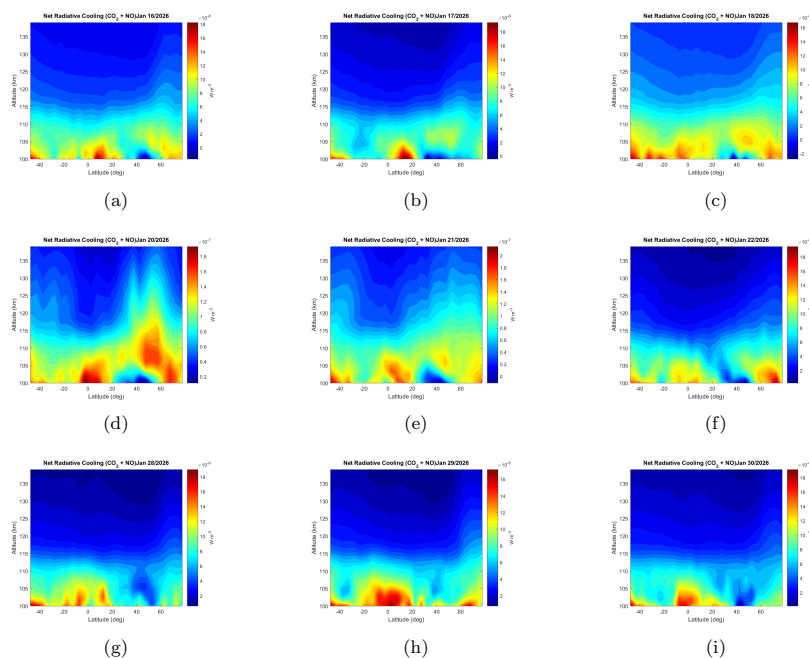


Figure 8: Net radiation cooling(Energy budgeting of atmosphere during quiet , storm time of January 2026

Following the storm, the recovery-phase panels Figure 8, g-i indicate a
gradual return toward quiet-time conditions. Although enhanced cooling
remains evident, particularly below 115 km altitude, the magnitude and



452 spatial extent of the cooling decrease compared with the storm peak. Max-
453 imum cooling rates generally return to values around $2.0\text{--}2.5\text{ W m}^{-3}$, and
454 the cooling layer becomes more confined to lower altitudes. This behavior
455 reflects the gradual reduction of storm-enhanced NO concentrations through
456 chemical loss processes and atmospheric transport. Nevertheless, residual
457 cooling enhancements persist for some time after the storm, indicating that
458 the thermosphere retains a memory of the storm-induced energy deposition
459 even after geomagnetic activity subsides.

460 Overall, the results demonstrate that geomagnetic storm conditions sig-
461 nificantly increase both the magnitude and vertical extent of NO radiative
462 cooling. Peak cooling rates increase by approximately 20-50% during the
463 storm compared with quiet conditions, while the cooling layer expands up-
464 ward by roughly 10-20 km. These observations provide clear evidence that
465 storm-time energy deposition strongly modifies thermospheric composition
466 and energetics through enhanced NO production, leading to more efficient
467 radiative cooling and contributing to the restoration of thermal balance in
468 the upper atmosphere.

469 3.8. Heat flux

470 The latitude–altitude distributions of atmospheric heating flux during
471 January 2026 show relatively modest variations compared with the NO cool-
472 ing results, with most heating flux values remaining within approximately
473 0.4×10^{-3} to $1.4 \times 10^{-3}\text{ W m}^2$. During the quiet-time period (top row), the
474 heating flux is generally distributed around $0.6\text{--}1.0 \times 10^{-3}\text{ W m}^2$, with only
475 isolated regions exceeding $1.2 \times 10^{-3}\text{ W m}^2$. The heating is fairly uniform
476 across latitude and altitude, indicating a thermosphere that is primarily
477 controlled by background solar heating and normal atmospheric dynamics.
478 No extensive regions of intense heating are evident, suggesting that external
479 energy inputs are relatively weak during these undisturbed conditions.

480 During the storm-time period (middle row), the heating flux exhibits a
481 noticeable enhancement in both magnitude and spatial variability. Larger
482 areas attain values of approximately $1.0\text{--}1.4 \times 10^{-3}\text{ W m}^2$, with localized
483 maxima approaching $1.5 \times 10^{-3}\text{ W m}^2$. Compared with the quiet-time condi-
484 tions, this represents an increase of roughly 20–50% in the strongest heating
485 regions. The enhanced heating is distributed over a wider range of latitudes
486 and extends throughout much of the 100–140 km altitude range. Such in-
487 creases are consistent with storm-induced energy deposition through Joule
488 heating and particle precipitation, which inject additional energy into the
489 thermosphere and create localized regions of elevated heating flux. The in-



490 creased spatial variability during the storm also indicates a more disturbed
491 and dynamically active upper atmosphere.

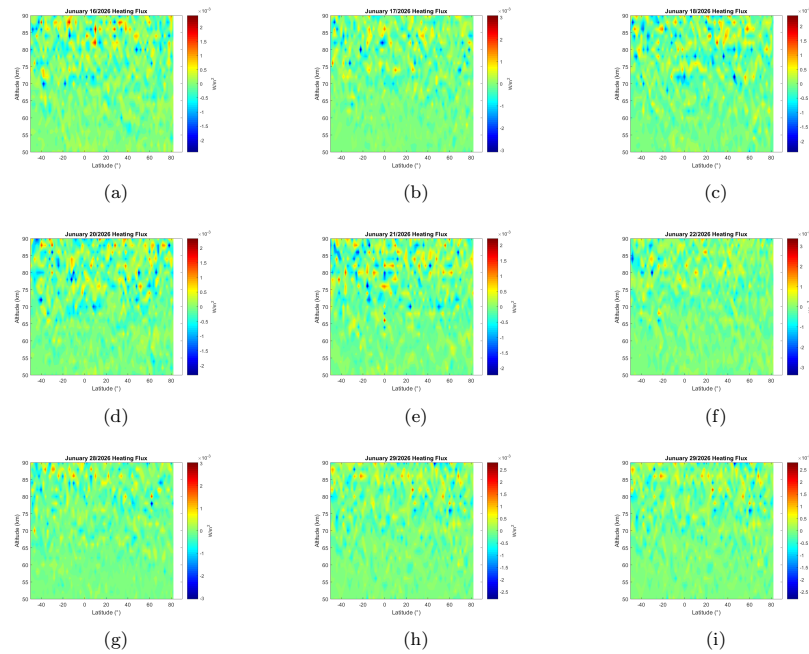


Figure 9: Heat flux of quiet , storm time of January 2026

492 Following the storm, the recovery-phase panels (bottom row) show a
493 gradual reduction in heating flux. Most values return to approximately $0.6 -$
494 $1.0 \times 10^{-3} Wm^2$, although some localized regions still exceed $1.2 \times 10^{-3} Wm^2$,
495 indicating that remnants of storm-time energy deposition persist. Compared
496 with the storm peak, both the intensity and spatial coverage of enhanced
497 heating decrease substantially. However, the atmosphere does not immedi-
498 ately return to its initial quiet-time state, reflecting the finite time required
499 for the thermosphere to dissipate excess energy and re-establish equilibrium.

500 Overall, the results indicate that geomagnetic storm activity enhances
501 atmospheric heating flux from typical quiet-time values of approximately
502 $0.6 - 1.0 \times 10^{-3} Wm^2$ to storm-time values reaching $1.4 - 1.5 \times 10^{-3} Wm^2$
503 . Although the absolute increase appears modest, the relative enhancement
504 of about 20–50% is significant for thermospheric energetics. The storm-time
505 increase in heating flux coincides with the enhanced NO cooling observed in



the previous figure, suggesting that the additional energy deposited during the storm is partially balanced by stronger infrared cooling. This simultaneous increase in heating and cooling demonstrates the thermosphere's response to geomagnetic forcing, whereby excess storm energy is redistributed and eventually removed through radiative processes, helping to restore the thermal balance of the upper atmosphere.

4. Conclusion

This study presents a comprehensive analysis of atmospheric tidal variability and gravity wave forcing during the January 19–21, 2026 geomagnetic storm. The key findings include: Significant amplification of semidiurnal and terdiurnal tides Emergence of higher-order tidal harmonics Enhanced gravity wave activity with strong intermittency Reduced vertical propagation due to increased dissipation Strong coupling between tides, gravity waves, and thermal structure The results demonstrate that geomagnetic storms drive complex multi-scale interactions that significantly modify atmospheric dynamics.

Strong vertical coupling from lower to upper atmosphere Sensitivity of atmospheric dynamics to geomagnetic forcing Importance of multi-scale interactions These findings are critical for improving models of space weather impacts. The present study demonstrates that geomagnetic storms exert a substantial influence on the upper atmosphere by altering its thermal structure, energy balance, and dynamical processes. During storm conditions, enhanced atmospheric heating flux indicates increased energy deposition through Joule heating and particle precipitation, while intensified NO radiative cooling reflects the thermosphere's response to dissipate excess energy and restore thermal equilibrium. Simultaneously, variations in gravity-wave potential energy reveal modifications in atmospheric dynamics and vertical energy transport. The observed increases in heating, cooling, and upper-atmospheric wave activity during storm periods compared with quiet conditions provide clear evidence of strong atmosphere–space weather coupling. Therefore, the results confirm that geomagnetic storms significantly affect the energetics and dynamics of the upper atmosphere, leading to measurable changes in thermospheric composition, temperature, and wave propagation characteristics.

The coordinated responses of GWPE, atmospheric heating flux, and NO radiative cooling demonstrate that geomagnetic storms are capable of significantly perturbing the upper atmosphere, modifying both its energy budget and dynamical behavior. These findings highlight the important role



544 of geomagnetic forcing in driving upper-atmospheric variability and
545 emphasize the need to consider storm-time effects when investigating thermo-
546 sphere–ionosphere coupling processes.

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556 stdb2.isee.nagoya-u.ac.jp/GPS/shinbori/AGRID2/nc/index.html. [https://](https://data.gats-inc.com/saber/Version2_0/)
557 data.gats-inc.com/saber/Version2_0/ SABER data level-2A to study
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563 Data Availability Statement

564 All data used in this study are publicly available at the following URLs:

- 565 • <https://omniweb.gsfc.nasa.gov/form/dx4.html>
- 566 • [https://stdb2.isee.nagoya-u.ac.jp/GPS/shinbori/AGRID2/nc/index](https://stdb2.isee.nagoya-u.ac.jp/GPS/shinbori/AGRID2/nc/index.html)
- 567 • <https://swarm-diss.eo.esa.int>
- 568 • https://data.gats-inc.com/saber/Version2_0/

569 **Declaration / Data Availability:** The authors declare that all data used
570 in this study are publicly available at the above URLs.



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