



1 **Conditional Modulation of Extreme Internal Solitary Waves in the**
2 **Lombok Strait**

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22 **Abstract.**

23 Using 13 months of continuous mooring observations at the Lombok Strait
24 (November 2023–December 2024), we identify 302 internal solitary wave (ISW)
25 packets and 618 individual solitons with amplitudes up to 104 m. ISW occurrence
26 peaks in boreal winter–spring (March: 45 packets/month) and drops to near-zero in
27 summer (July: 3 packets/month). The observations reveal an amplitude-dependent
28 dynamical structure. ISW occurrence for both amplitude classes is primarily
29 modulated by the ITF: normal-amplitude ISWs (25–65 m) are seasonally suppressed
30 during enhanced ITF periods—when thermocline stratification is concurrently
31 elevated—indicating that the apparent negative N^2 –ISW correlation is a seasonal
32 covariance with ITF strength rather than a direct stratification effect. Extreme ISWs
33 (>65 m) show negligible stratification dependence and instead tend to coincide with
34 transient ITF weakening as a precondition for nonlinear amplification. Low-pass
35 filtered acoustic Doppler current profiler records show that extreme events
36 consistently coincide with northward velocity pulses in the 50–200 m layer that
37 substantially reduce Doppler retardation, consistent with the existence of episodic
38 dynamic windows favorable for nonlinear wave amplification.



39 **1 Introduction**

40 Internal solitary waves (ISWs) are large-amplitude, nonlinear disturbances that
41 propagate along oceanic pycnoclines, transporting momentum and energy across vast
42 distances and driving intense diapycnal mixing that influences nutrient cycling,
43 biological productivity, and heat redistribution (Lamb, 2014; Alford et al., 2015).
44 Although ISWs are ubiquitous in stratified marginal seas and ocean straits, the
45 Indonesian Archipelago stands out as a global hotspot: it constitutes the only low-
46 latitude conduit connecting the Pacific and Indian Oceans via the Indonesian
47 Throughflow (ITF), and its complex system of narrow passages and steep sills
48 generates some of the most energetic ISW fields on Earth (Gordon, 2005; Sprintall et
49 al., 2009). The Lombok Strait—separating the islands of Bali and Lombok at
50 approximately 8.5°S—is one of the three principal ITF exit passages and among the
51 most active ISW-generation sites globally. Semidiurnal tidal currents interact with the
52 strait's shallow sill to produce large-amplitude ISW packets that propagate
53 predominantly northward into the Bali Sea, with isopycnal displacements of 40–
54 100 m and phase speeds of 1.5–2.5 m s⁻¹ (Susanto et al., 2005; Purwandana et al.,
55 2021; Gong et al., 2022). These waves exert important impacts on mixing, biological
56 productivity, acoustic propagation, and submarine navigation safety—a concern
57 underscored by the loss of the Indonesian naval submarine KRI Nanggala-402 in
58 these waters in 2021.

59 The seasonal variability of Lombok Strait ISWs has been studied almost exclusively
60 through satellite remote sensing and numerical modeling. Mitnik et al. (2000) and



61 Susanto et al. (2005) first characterized ISW distributions using SAR imagery.

62 Matthews et al. (2011) combined PALSAR/ERS SAR observations with numerical

63 modeling to demonstrate that northward-propagating ISWs are far more frequent

64 during the Northwest Monsoon (boreal winter) than the Southeast Monsoon (boreal

65 summer), and proposed that the strong southward ITF during the Southeast Monsoon

66 physically inhibits embryonic wave growth at the sill by constraining the northward

67 tidal intrusion—a hydraulic-control mechanism confirmed by 2D non-hydrostatic

68 simulations. Chonnaniyah et al. (2024) subsequently showed from multi-satellite

69 observations that soliton numbers and phase speeds are higher during the Northwest

70 Monsoon, consistent with shallower thermocline conditions. More recently, Zhuang et

71 al. (2024) compiled 858 satellite images (2018–2022) and confirmed that ISW

72 occurrence peaks in October (29% detection frequency) and is lowest in June (6%),

73 with enhanced southward ITF corresponding to suppressed ISW occurrence and

74 reduced northward propagation speed; three-dimensional MITgcm experiments

75 further reproduced the suppressive effect of stronger background flow via hydraulic-

76 control delay of ISW generation. Limited in situ data from coastal acoustic

77 tomography (Syamsudin et al., 2019) and shipboard ADCP surveys during research

78 cruises (Purwandana et al., 2021) have corroborated the presence of these waves, but

79 no continuous subsurface mooring record spanning a full seasonal cycle has yet been

80 reported.

81 Despite these advances, three critical gaps persist in our understanding of Lombok

82 Strait ISW dynamics. First, all existing seasonal analyses rely on satellite remote



83 sensing, which captures only surface manifestations and cannot resolve subsurface
84 isopycnal displacement amplitude, individual event timing, or the full water-column
85 velocity structure. Second, the relative contributions of thermocline depth, local
86 stratification strength (buoyancy frequency N^2), and background ITF velocity to ISW
87 seasonal variability have never been disentangled from direct in situ observations.
88 Third, and most critically, no study has explicitly examined whether ISWs of different
89 amplitudes may be influenced by distinct dynamical mechanisms. Prior work
90 implicitly treats ISW activity as controlled by a single set of environmental conditions,
91 yet nonlinear wave theory predicts that large-amplitude waves may undergo
92 substantial steepening and amplification through weakly nonlinear wave processes
93 that are dynamically distinct from the linear wave generation that governs smaller
94 events. Weakly nonlinear theory is adopted here because ISW amplitudes observed in
95 the Lombok Strait (25–104 m, i.e., roughly 10–60% of the thermocline depth) place
96 most events within the weakly-to-moderately nonlinear regime where amplitude-to-
97 thermocline-depth ratios $a/h_1 \leq 0.5$, and the qualitative predictions of weakly
98 nonlinear theory regarding the sign and direction of nonlinear steepening remain valid
99 (Lamb, 2014). Fully nonlinear approaches (e.g., Dubreil-Jacotin–Long equation)
100 would provide more accurate amplitude estimates for the largest events ($a/h_1 \approx 0.5$ –
101 1.0) but require complete density profiles unavailable from the present thermistor
102 chain; this limitation is noted in the discussion (Sect. 4.4). Here we present the first
103 analysis of ISW amplitude statistics and their amplitude-dependent seasonal dynamics
104 from a 13-month continuous mooring observation at the Lombok Strait.



105 **2 Data and Methods**

106 **2.1 Mooring Configuration**

107 A subsurface mooring (LMK) was deployed at 8°13'31"S, 115°51'03"E in the
108 northern Lombok Strait (water depth ~1000 m) from 20 November 2023 to 11
109 December 2024, providing 13 months of continuous high-frequency observations
110 spanning a complete monsoon cycle. The mooring comprised three main components.
111 (1) A conductivity-temperature-depth (CTD) sensor nominally at 50 m depth recorded
112 temperature and conductivity at 1-minute intervals; the instrument was displaced to a
113 mean depth of ~110 m by the strong downward drag of the ITF. (2) Seven thermistors
114 at nominal depths of 70, 90, 110, 130, 150, 170, and 190 m recorded temperature at 1-
115 minute intervals; actual depths were estimated from mooring knockdown corrections
116 applied to pressure records. (3) An upward-looking Acoustic Doppler Current Profiler
117 (ADCP; 75 kHz) at a nominal depth of 500 m measured three-dimensional currents in
118 27 bins (~16 m vertical resolution, 40–500 m range) at 2-minute intervals. All times
119 are in Indonesian Western Time (UTC+7). Temperature and salinity observations
120 were interpolated onto a uniform 1-minute × 1-m grid; ADCP velocity profiles were
121 similarly gridded to 1-minute × 1-m over 40–500 m.

122 **2.2 ISW Identification**

123 ISWs were detected from the 18°C isotherm depth time series using an automated
124 algorithm adapted from Gong et al. (2023). A 6-hour low-pass Butterworth filter
125 (fourth-order) was applied to the isotherm depth to define the slowly varying



126 background thermocline state. When the 18°C isotherm was absent from the
127 observation range for more than 6 consecutive hours, detection was switched to the
128 16°C isotherm (invoked for <4% of events; sensitivity tests confirm this does not alter
129 amplitude statistics). Isotherm depth was obtained by linear interpolation between the
130 two bracketing thermistors; time steps where the target isotherm lay outside the 70–
131 190 m instrument range were set to missing, and gaps shorter than 6 h were bridged
132 by linear interpolation. Mooring knockdown was corrected using pressure records
133 from two independent sensors: the CTD pressure sensor (nominal depth 50 m, actual
134 mean depth ~110 m due to ITF-induced displacement) and the ADCP integrated
135 pressure sensor (nominal depth 500 m). Using two reference points is important
136 because the depth-dependent shear of the ITF causes upper instruments to experience
137 larger horizontal drag and hence larger knockdown than deeper instruments, so the
138 displacement profile is non-uniform with depth. The correction procedure is:
139 (i) Pressure at each sensor is converted to depth via $z = p/(\rho g)$, $\rho = 1025 \text{ kg m}^{-3}$,
140 $g = 9.81 \text{ m s}^{-2}$. (ii) Instantaneous knockdown displacement at each reference sensor:

$$\delta z_{\text{CTD}}(t) = z_{\text{CTD}}(t) - z_{\text{CTD}}^{\circ}$$

$$\delta z_{\text{ADCP}}(t) = z_{\text{ADCP}}(t) - z_{\text{ADCP}}^{\circ}$$

141 where $z_1(t)$ and $z_2(t)$ are pressure-derived depths at each time step and

142 $z_{\text{CTD}}^{\circ} \approx 50 \text{ m}$, $z_{\text{ADCP}}^{\circ} = 500 \text{ m}$ are the nominal depths.

143 (iii) Knockdown at each thermistor (nominal depths 70–190 m, all between the two

144 reference sensors) is linearly interpolated:

$$\delta z_i(t) = \delta z_{\text{CTD}}(t) + \frac{z_i^{\circ} - z_{\text{CTD}}^{\circ}}{z_{\text{ADCP}}^{\circ} - z_{\text{CTD}}^{\circ}} \times [\delta z_{\text{ADCP}}(t) - \delta z_{\text{CTD}}(t)]$$



145 giving corrected depth:

$$z_i(t) = z_i^{\circ} + \delta z_i(t)$$

146 where $z_{\text{CTD}}^{\circ} \approx 50$ m and $z_{\text{ADCP}}^{\circ} = 500$ m are the nominal deployment depths. Linear

147 interpolation is standard for taut wire moorings where wire stretch is negligible

148 (Dewey, 1999).

149 The CTD, nominally at 50 m, experienced a sustained mean displacement of ~ 60 m to

150 an actual mean depth of ~ 110 m due to the persistent horizontal drag of the strong ITF

151 current; additional dynamic fluctuations of ± 5 – 10 m around this mean occurred

152 during tidal and ISW forcing. The ADCP at 500 m, where currents are substantially

153 weaker, experienced dynamic displacements of 5–15 m. To separate ISW packets

154 from large-amplitude internal tides, events were required to show an isotherm

155 displacement rate exceeding 2.0 m min^{-1} and a peak-to-trough duration shorter than

156 30 min; these criteria effectively exclude the semidiurnal tidal signal (period ≈ 12.4 h)

157 while retaining all identified solitary-wave events. Only depression-type events

158 (downward isotherm displacement) were retained, consistent with the dominance of

159 mode-1 depression waves in the Lombok Strait (Susanto et al., 2005; Matthews et al.,

160 2011).

161 Two additional quality criteria were applied: (1) background isotherm stability within

162 ± 25 m of the 6-hour running mean; and (2) isotherm recovery to within 10 m of

163 background on both flanks of the detected anomaly. Individual solitons were

164 separated by a minimum gap of 15 min; consecutive packets required gaps of at least

165 90 min. All events were visually inspected via daily temperature Hovöller diagrams.

166 The final event catalog comprises 302 wave packets and 618 individual solitons. ISW



167 amplitude was defined as the maximum isotherm depression per soliton. For packet-
168 level statistical analyses (Sects. 3.4 and 4), each packet is represented by its peak-
169 amplitude soliton (the largest among all solitons in the packet); the 618 individual
170 soliton amplitudes are used for the amplitude distribution statistics only. Extreme
171 ISWs (>65 m) correspond to the sparse high-amplitude tail of the distribution (~9% of
172 packets), identified by a natural break in the amplitude histogram near 65 m
173 (Supplementary Fig. S1). A sensitivity analysis shows that varying this threshold
174 between 60 and 70 m does not qualitatively alter the conditional probability contrast
175 or Mann–Whitney results; the 65 m value is therefore retained throughout.

176 **2.3 Environmental Parameters and Phase Speed**

177 Four monthly-mean environmental parameters were computed. (1) Thermocline depth:
178 monthly-mean depths of the 18°C and 20°C isotherms by linear interpolation. (2)
179 Local stratification: buoyancy frequency N^2 averaged within ± 10 m of the 18°C
180 isotherm, computed as $N^2 = -(g/\rho_0)(\partial\rho/\partial z)$, where ρ is computed from temperature
181 and salinity using TEOS-10 (IOC et al., 2010) via the gsw library, with $g = 9.81 \text{ m s}^{-2}$
182 and $\rho_0 = 1025 \text{ kg m}^{-3}$. (3) Background ITF velocity: monthly-mean meridional
183 ADCP velocity averaged over 100–250 m (thermocline layer); ADCP data above
184 50 m were excluded due to sidelobe contamination. For time-series analysis, a 72-
185 hour low-pass Butterworth filter isolated background flow from tidal signals. ADCP
186 velocity records were available through 30 October 2024; for the November–
187 December 2024 period, background meridional velocity was supplemented using
188 daily GLORYS12 reanalysis fields (cmems_mod_glo_phy_my_0.083deg_PID-m;



189 Copernicus Marine Service), averaged over the two model grid points nearest the
190 mooring (115.833°E and 115.917°E, 8.333°S) and bias-corrected against the final 30-
191 day overlap with ADCP observations. (4) Effective phase speed: estimated using a
192 two-layer model, $c_0 = (g'h_1 - h_2 / (h_1 + h_2))^{1/2}$, where $g' = g(\rho_2 - \rho_1) / \rho_0$ is the
193 reduced gravity, h_1 is the monthly-mean 18°C isotherm depth (upper layer), $h_2 =$
194 $1000 - h_1$ (lower layer, total depth $H = 1000$ m), and densities ρ_1, ρ_2 were
195 computed with the TEOS-10 equation of state (IOC et al., 2010). The Doppler-
196 corrected effective phase speed is $c^{\text{eff}} = c_0 + V_{100-250}$, where negative (southward)
197 ITF velocity reduces c^{eff} and retards northward ISW propagation.

198 2.4 Statistical Analysis

199 Monthly ISW counts and environmental parameters were computed over 13 months
200 (December 2023–December 2024; November 2023 excluded as partial). Pearson
201 correlation coefficients (r) were computed between monthly ISW occurrence for two
202 amplitude classes—normal (25–65 m) and extreme (>65 m)—and each environmental
203 parameter, with significance assessed at $\alpha = 0.05$ (two-tailed). To quantify the
204 background flow anomaly during extreme ISW events, the 72-hour low-pass filtered
205 meridional velocity was averaged within ± 6 hours of each extreme ISW peak time
206 and compared against the record-mean. All analyses used Python (v3.9) with NumPy,
207 SciPy, and Pandas. To address the small monthly sample size ($n = 13$), robustness of
208 all monthly correlations was assessed using leave-one-out (LOO) analysis ($n = 12$ per
209 iteration) and non-parametric bootstrap resampling (5000 iterations). Bootstrap 95%

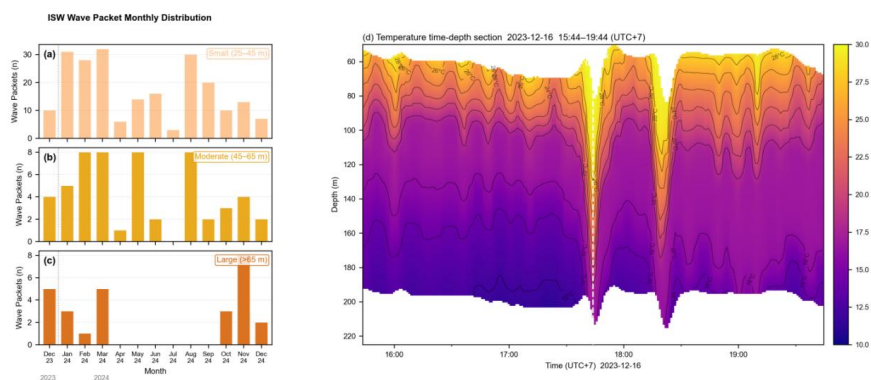


210 confidence intervals (CIs) and permutation p-values (5000 permutations) are reported
211 alongside parametric p-values.

212 **3 Results**

213 **3.1 ISW Occurrence and Amplitude Statistics**

214 Over the 13-month mooring record, 302 ISW wave packets comprising 618 individual
215 solitons were identified. Peak amplitudes ranged from 25 to 104 m, with a mean of
216 40.3 m. The amplitude distribution is strongly right-skewed: small ISWs (25–45 m)
217 account for 52.0% of all packets (n = 157), moderate ISWs (45–65 m) for 39.4%
218 (n = 119), and extreme ISWs (>65 m) for 8.6% (n = 26). Fig. 1a–c shows the monthly
219 distribution of each amplitude class. ISWs occurred in every month of the record,
220 with strong seasonal modulation. The maximum occurrence was in March 2024 (45
221 packets/month), followed by January 2024 (39 packets/month) and August 2024 (39
222 packets/month). The minimum occurred in July 2024 (3 packets/month). Extreme
223 ISWs were absent in April–July and September 2024, but reached 8 events in
224 November 2024, suggesting that the amplitude structure of the seasonal cycle differs
225 from that of total occurrence.
226



227

228 Figure 1. Monthly distribution of ISW wave packets and a representative ISW event.

229 (a)–(c) Bar charts showing the number of ISW wave packets per month in three

230 amplitude categories: (a) small ISWs (25–45 m; light orange), (b) moderate ISWs

231 (45–65 m; orange), and (c) extreme ISWs (>65 m; vermilion). The dotted vertical

232 line marks the boundary between the 2023 and 2024 portions of the record.

233 (d) Temperature time–depth section for the representative extreme ISW event of 16

234 December 2023 (15:44–19:44 UTC+7), plotted using the plasma colour scale (10–

235 30 °C). Isotherms are contoured every 2 °C. The white dashed vertical line marks the

236 ISW peak at 17:44.

237

238 A representative extreme ISW event observed on 16 December 2023 illustrates the

239 characteristic subsurface structure (Fig. 1d). Within a 4-hour window centered on

240 17:44 UTC+7, the 18°C isotherm depressed by ~80 m over the course of

241 approximately 20 minutes, with accompanying downward displacement of the 16°C

242 and 20°C isotherms. The temperature time–depth section reveals a classic depression

243 ISW signature: a sharp leading front followed by a gradual return to the background



244 state, with high-frequency oscillations trailing the primary wave packet. Such events
245 are clearly distinguishable from semidiurnal tidal oscillations in both their shorter
246 duration (<30 min for individual solitons) and steeper isotherm displacement rates
247 (>5 m min⁻¹).

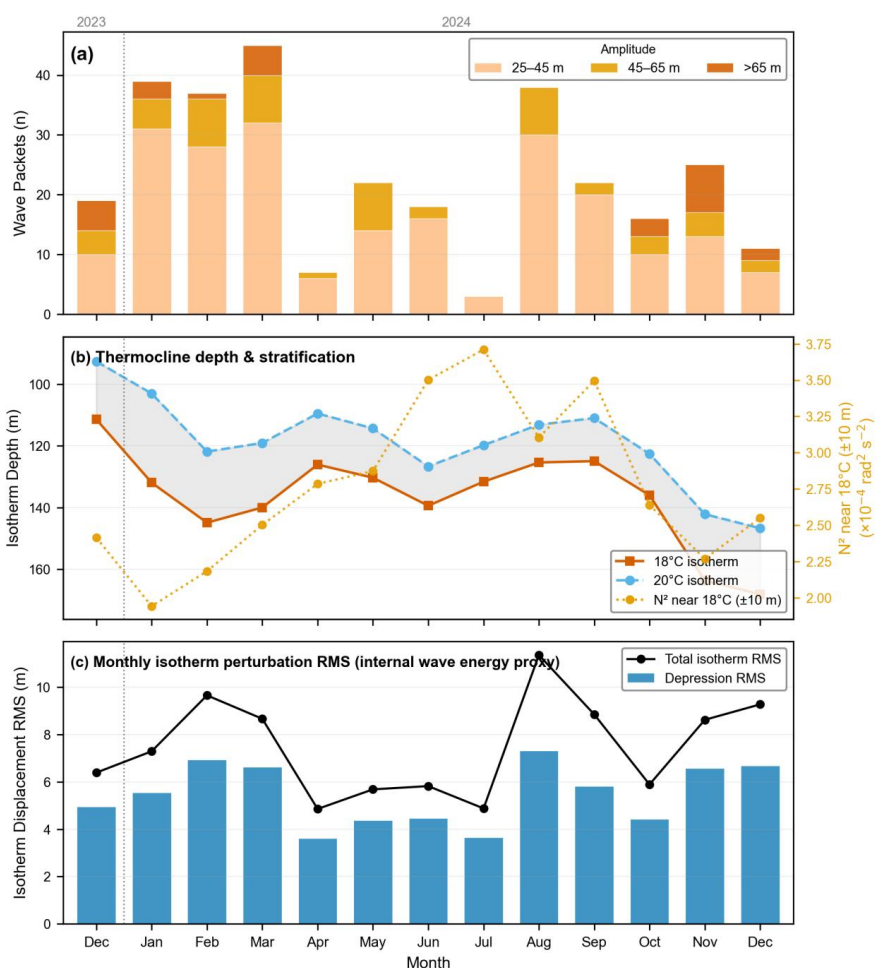
248 **3.2 Seasonal Variability and Environmental Parameters**

249 Figure 2 presents the monthly co-variability of ISW occurrence, thermocline state,
250 and internal wave energy. ISW wave packet counts (Fig. 2a) exhibit a bimodal pattern
251 with a primary maximum in boreal winter–spring (December 2023–March 2024) and
252 a secondary maximum in late summer (August 2024), separated by a pronounced
253 summer minimum (July 2024: 3 packets). The extreme ISW class (>65 m, dark bars)
254 is largely absent during April–September, even when total packet counts are moderate
255 (e.g., August 2024: 39 total packets but only 1 extreme event).

256 The 18°C and 20°C isotherm depths (Fig. 2b) show a progressive deepening from
257 December 2023 (~112 m and ~93 m, respectively) through December 2024 (~168 m
258 and ~147 m), a net descent of ~56 m over the 13-month record. Notably, the deepest
259 isotherm depths occur in November–December 2024 (Northwest Monsoon transition),
260 not during the Southeast Monsoon peak (June–August 2024): the July 2024 value
261 (~132 m) is shallower than November 2024 (~164 m), indicating that this record
262 reflects a combination of the seasonal ITF intensification cycle and an interannual
263 deepening trend that cannot be decomposed from a single 13-month observation. The
264 local N² within ±10 m of the 18°C isotherm (Fig. 2b right axis; TEOS-10) ranges from



265 1.9 to $3.7 \times 10^{-4} \text{ rad}^2 \text{ s}^{-2}$, with a maximum in July 2024 ($\sim 3.7 \times 10^{-4} \text{ rad}^2 \text{ s}^{-2}$) and
 266 minimum in January 2024 ($\sim 1.9 \times 10^{-4} \text{ rad}^2 \text{ s}^{-2}$). Notably, this window follows the
 267 moving isotherm; sensitivity analysis using fixed-depth and column-maximum
 268 stratification metrics yields non-significant correlations with ISW occurrence (see
 269 Sect. 4.1), indicating that this N^2 estimate should not be interpreted as evidence of
 270 direct stratification control on ISW variability.



271



272 Figure 2. Seasonal variability of ISW occurrence and environmental parameters.

273 (a) Monthly ISW wave packet counts, stacked by amplitude category: light orange =

274 small (25–45 m), orange = moderate (45–65 m), vermillion = extreme (>65 m).

275 (b) Monthly-mean depth of the 18 °C isotherm (vermillion squares, left axis) and

276 20 °C isotherm (sky-blue circles, left axis), with the grey shading indicating the

277 thermocline layer between them. The right axis (orange dashed line with circles)

278 shows local buoyancy frequency N^2 within ± 10 m of the 18 °C isotherm, computed

279 using TEOS-10. (c) Monthly root-mean-square (RMS) displacement of the 18 °C

280 isotherm (blue bars: depression-type perturbations; black line: total isotherm RMS).

281 The dotted vertical line separates 2023 and 2024.

282

283 The monthly isotherm displacement root-mean-square (RMS; Fig. 2c), used here as a

284 proxy for total internal wave energy in the thermocline, broadly follows the ISW

285 occurrence pattern but with important differences. The depression RMS peaks in

286 February–March 2024 (~10 m) and again in August 2024 (~7 m), while reaching a

287 minimum in July 2024 (~5 m). The fact that the total isotherm RMS (including both

288 depression and elevation displacements) is consistently elevated throughout the record

289 (5–9 m) indicates persistent internal wave activity even during the summer

290 suppression period, consistent with continuous tidal forcing. The relatively modest

291 suppression of total RMS in July (~5 m) compared to the near-complete absence of

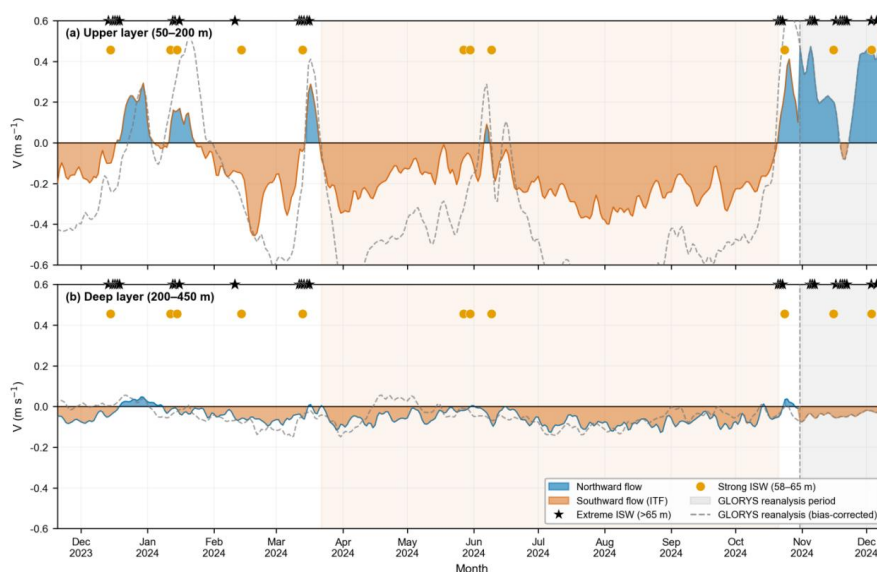
292 extreme ISWs suggests that the summer ITF does not suppress internal wave



293 generation entirely, but rather limits the nonlinear amplification required for extreme
294 wave development.

295 3.3 Background Flow Compensation and Extreme ISW Events

296 Fig. 3 shows the 72-hour low-pass filtered meridional velocity in the upper (50–200 m)
297 and deep (200–450 m) layers throughout the mooring record, with extreme ISW
298 events (>65 m, black stars) and strong ISW events (58–65 m, orange circles) indicated.
299 The upper layer (Fig. 3a) is dominated by southward (ITF) flow, with monthly-mean
300 velocities of -0.147 m s^{-1} across the full record. However, northward velocity pulses
301 of $0.1\text{--}0.4 \text{ m s}^{-1}$ occur episodically, most prominently in January 2024, March 2024,
302 and October–November 2024. These northward pulses are associated with weakening
303 or reversal of the ITF during the transition between monsoon seasons and under the
304 influence of intraseasonal Kelvin wave forcing from the Indian Ocean.



305



306 Figure 3. 72-hour low-pass filtered meridional velocity in the (a) upper (50–200 m)
307 and (b) deep (200–450 m) layers throughout the mooring record. Blue shading
308 indicates northward flow; vermillion shading indicates southward flow (Indonesian
309 Throughflow direction). The solid line shows the ADCP-observed velocity; the grey
310 dashed line shows bias-corrected GLORYS12 reanalysis velocity for comparison. The
311 grey shaded region after 30 October 2024 indicates the GLORYS-extended period.
312 Black stars mark extreme ISW events (amplitude >65 m); orange circles mark strong
313 ISW events (58–65 m). The vermillion background shading indicates the enhanced
314 ITF period (late March to late October 2024), determined from zero crossings of the
315 upper-layer low-pass velocity. The grey dashed vertical line indicates the ADCP data
316 cutoff (30 October 2024).

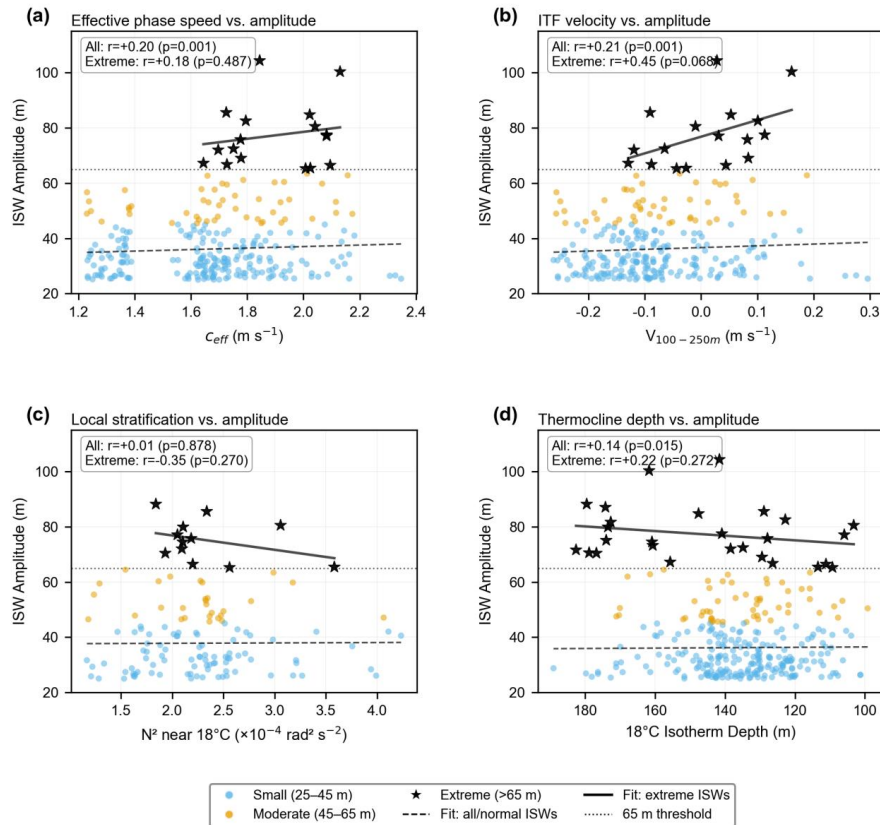
317

318 A clear correspondence is evident between northward velocity pulses and extreme
319 ISW occurrence: most extreme ISW events (black stars) coincide with periods when
320 the 50–200 m layer velocity is near-zero or weakly northward. During the enhanced
321 ITF period (approximately April–October 2024, as determined from the zero
322 crossings of the upper-layer low-pass velocity), the upper layer maintains persistent
323 southward velocities of -0.2 to -0.4 m s^{-1} , and extreme ISWs are essentially absent.
324 In contrast, the deep layer (200–450 m; Fig. 3b) shows much weaker variability
325 (record mean: -0.051 m s^{-1} , range: -0.14 to $+0.08$ m s^{-1}) and no clear association with
326 extreme ISW events, consistent with a compensation mechanism operating primarily
327 in the thermocline layer rather than at depth.



328 Compositing the 72-hour low-pass velocity at times of extreme ISW events (± 6 h
329 around each peak) reveals a mean upper-layer velocity of $+0.005 \text{ m s}^{-1}$ during extreme
330 events, compared to a record mean of -0.147 m s^{-1} —a difference of $+0.153 \text{ m s}^{-1}$. The
331 deep layer shows negligible change ($+0.016 \text{ m s}^{-1}$). This layer-dependent pattern
332 suggests that the Doppler retardation of northward ISW propagation, quantified as the
333 reduction in effective phase speed $c^{\text{eff}} = c_0 + V$, may be substantially reduced during
334 extreme events. When $V \approx 0 \text{ m s}^{-1}$, the effective phase speed approaches the intrinsic
335 linear value c_0 ($1.47\text{--}2.22 \text{ m s}^{-1}$ across our record), partially alleviating the
336 propagation constraint and potentially opening a dynamic window for nonlinear wave
337 growth.

338 **3.4 Amplitude-Dependent Control Mechanisms**



339

340 Figure 4. Event-scale relationship of ISW peak amplitude versus four environmental

341 parameters ($n = 302$ wave packets, one per packet; $n = 265$ for panels (a)–(b) owing to

342 the subset with concurrent ADCP records). (a) Effective phase speed c^{eff} .

343 (b) Thermocline-layer ITF velocity $V_{100-250}$ (100–250 m; northward positive).

344 (c) Local N^2 within ± 10 m of the 18 °C isotherm. (d) Monthly-mean 18 °C

345 isotherm depth (x-axis inverted: right = deeper thermocline). Scatter point colours

346 indicate season: blue = December–February (Northwest Monsoon), orange = March–

347 May, vermilion = June–August (Southeast Monsoon), purple-pink = September–

348 November. Sky-blue circles denote small ISWs (25–45 m); orange circles denote

349 moderate ISWs (45–65 m); black five-pointed stars denote extreme ISWs (>65 m).



350 Solid lines show statistically significant linear fits ($p < 0.05$); dashed lines show fits
351 for the full dataset ($p \geq 0.05$). The horizontal dotted line marks the 65 m extreme ISW
352 threshold.

353

354 Figure 4 presents event-scale scatter plots of individual ISW peak amplitude ($n = 265$
355 for velocity parameters, $n = 302$ for thermocline depth) against four environmental
356 parameters measured at the time of each event: effective phase speed c^{eff} (Fig. 4a),
357 thermocline-layer ITF velocity $V_{100-250}$ (Fig. 4b), local stratification N^2 (Fig. 4c), and
358 18°C isotherm depth (Fig. 4d). Stars denote extreme ISW events (>65 m). This event-
359 scale approach avoids temporal aliasing and substantially increases statistical power.

360 Stratification shows no significant correlation with ISW amplitude ($r = +0.01$,
361 $p = 0.878$; Fig. 4c): extreme ISWs (stars) are distributed across the full range of
362 observed N^2 values, indicating that thermocline strength is not the primary control on
363 wave amplitude. In contrast, effective phase speed c^{eff} is positively correlated with
364 amplitude ($r = +0.20$, $p < 0.001$; Fig. 4a), and extreme ISWs cluster primarily at
365 $c^{\text{eff}} > 1.7 \text{ m s}^{-1}$. ITF thermocline velocity $V_{100-250}$ is similarly significantly correlated
366 with amplitude ($r = +0.21$, $p < 0.001$; Fig. 4b): no extreme ISW occurred when $V_{100-250}$
367 $< -0.2 \text{ m s}^{-1}$. Mann–Whitney U tests indicate that both c^{eff} and $V_{100-250}$ distributions
368 differ significantly between extreme and normal ISW events ($p < 0.001$). The 18°C
369 isotherm depth shows a weak positive correlation with amplitude ($r = +0.14$, $p = 0.015$;
370 Fig. 4d), likely reflecting seasonal covariance between deep thermocline and weak
371 ITF rather than a direct effect.



372 Although the Pearson correlations are modest ($r \approx 0.24$, explaining ~6% of amplitude
373 variance), reflecting the many factors controlling ISW amplitude (tidal phase,
374 topographic focusing, antecedent wave field), the Mann–Whitney results indicate that
375 the background flow difference between extreme and normal ISW events is robust
376 and statistically significant. This contrast—amplitude insensitive to stratification but
377 responsive to ITF-induced Doppler conditions—represents an important but
378 previously underexplored aspect of ISW dynamics in ITF-influenced passages.

379

380 **4 Discussion**

381 **4.1 Re-evaluation of the N^2 –ISW Relationship**

382 The apparent negative correlation between N^2 (computed within ± 10 m of the moving
383 18°C isotherm; Method A) and monthly ISW occurrence ($r = -0.60$, $p = 0.053$,
384 $n = 13$ months) requires careful re-examination. This computation window moves
385 with the isotherm as it descends from ~ 112 m in December 2023 to ~ 168 m in
386 December 2024, meaning that N^2 estimates from different months are computed at
387 different absolute depths. When stratification is instead measured at a fixed depth
388 range (100–200 m; Method B), the correlation is $r = -0.51$ ($p = 0.07$) with a bootstrap
389 95 % CI of $[-0.84, -0.09]$ that does not span zero (the apparent discrepancy between
390 the CI excluding zero and the parametric $p = 0.07$ reflects small-sample uncertainty
391 and differing assumptions: the bootstrap is non-parametric while the parametric test
392 assumes bivariate normality at $n = 13$). Regardless of the inference approach, this



393 negative correlation is consistent with seasonal ITF co-modulation: during the
394 Southeast Monsoon, the same ITF intensification that suppresses ISW occurrence via
395 hydraulic control and Doppler retardation also drives thermocline deepening,
396 elevating N^2 in the 100–200 m layer. The most direct evidence against an independent
397 stratification effect is the event-scale scatter ($r = +0.01$, $p = 0.878$, $n = 302$): when the
398 seasonal ITF background is removed and individual ISW events are compared, N^2 has
399 no predictive power for ISW amplitude whatsoever. Neither the column-maximum N^2
400 in 50–300 m (Method C: $r = -0.10$, $p = 0.76$) nor the two-layer effective phase speed
401 c_0 ($r = -0.39$, $p = 0.23$) shows a significant relationship with ISW occurrence, and
402 the 18°C isotherm depth itself is essentially uncorrelated with ISW counts ($r = +0.08$,
403 $p = 0.82$).

404 These results indicate that the Method A N^2 –ISW correlation is an artefact of the co-
405 movement between the tracking window and the ITF-driven thermocline. When the
406 southward ITF intensifies, it simultaneously depresses the 18°C isotherm (shifting the
407 ± 10 m window into the thermocline core and yielding higher local N^2) and suppresses
408 ISW occurrence via Doppler retardation and hydraulic control. The negative N^2 –ISW
409 relationship therefore reflects ITF co-modulation rather than a causal stratification
410 effect—a conclusion consistent with the satellite-based analyses of Matthews et al.
411 (2011) and Zhuang et al. (2024), who similarly attributed ISW suppression to ITF
412 dynamics rather than to stratification per se. This conclusion extends to both
413 amplitude classes: no stratification metric tested here shows a physically robust,
414 independent correlation with ISW occurrence or amplitude. The ITF thermocline-



415 layer velocity therefore emerges as the primary physical control for both amplitude
416 classes, operating through seasonal suppression for normal events and the episodic
417 background flow compensation mechanism for extreme events (Section 4.2)..

418 **4.2 Background Flow Compensation Mechanism for Extreme ISWs**

419 The decoupling of ISW amplitude from stratification ($r = +0.01$, $p = 0.878$), and its
420 conditional clustering relationship with ITF velocity ($r = +0.21$, $p < 0.001$, $n = 265$
421 events), suggests the presence of a distinct physical mechanism operating for the
422 largest wave events. We propose that this reflects a background flow compensation
423 process in which episodic reductions of the southward ITF—driven by intraseasonal
424 Kelvin waves propagating from the Indian Ocean (Matthews et al., 2011; Arief &
425 Murray, 1996) or by the seasonal ITF weakening during the Northwest Monsoon
426 transition—temporarily reduce the Doppler retardation that normally constrains
427 northward ISW propagation speed. When the thermocline-layer velocity approaches
428 zero ($c^{\text{eff}} \approx c_0$), the nonlinear steepening timescale, which scales as $(\alpha A)^{-1}$ in weakly
429 nonlinear theory where α is the nonlinearity parameter and A is wave amplitude,
430 becomes short relative to the propagation timescale over the strait width. Under such
431 conditions, initial tidal disturbances or baroclinic tidal beams may steepen into large-
432 amplitude solitons.

433 The quantitative observational evidence for this mechanism is broadly consistent with
434 the proposed mechanism. During extreme ISW events, the upper-layer background
435 velocity is $+0.005 \text{ m s}^{-1}$ (effectively zero), compared to the record mean of



436 -0.147 m s^{-1} . The corresponding increase in effective phase speed $\Delta c^{\text{eff}} \approx +0.15 \text{ m s}^{-1}$
437 is dynamically meaningful: given a typical intrinsic phase speed of $c_0 \approx 1.7$ –
438 2.2 m s^{-1} in our record, the mean ITF-induced retardation of 0.15 m s^{-1} represents a 7–
439 10% reduction in effective propagation speed. While modest in absolute terms, this
440 retardation is sufficient to alter the phase relationship between the generating tidal
441 forcing and the nascent wave disturbance, potentially modulating the efficiency of
442 energy extraction at the sill—consistent with the hydraulic-control timing mechanism
443 described by Zhuang et al. (2024). Furthermore, the absence of a significant deep-
444 layer (200–450 m) velocity anomaly during extreme events ($+0.016 \text{ m s}^{-1}$) indicates
445 that the compensation is confined to the thermocline layer, precisely where the ISW
446 orbital velocities and mode-1 baroclinic shear are strongest. Note that the monthly-
447 mean thermocline-layer velocity shows no significant correlation with monthly ISW
448 packet counts ($r = -0.06$, $p = 0.86$; LOO range $[-0.28, +0.09]$), confirming that the
449 ITF control operates at the event scale rather than through a simple monthly-mean
450 relationship. The amplitude-based event-scale analysis ($n = 265$; Sect. 3.4) therefore
451 provides the more direct and relevant test of the proposed mechanism.

452 **4.3 Implications for ISW Hazard Assessment**

453 The background flow compensation mechanism has practical implications for ISW
454 hazard prediction. Previous satellite-based approaches have relied on monsoon phase
455 or stratification indices as proxies (Zhuang et al., 2024; Chonnaniyah et al., 2024),
456 which are adequate for normal-amplitude ISWs but insufficient for anticipating
457 extreme events. Real-time monitoring of thermocline-layer ITF velocity—from



458 satellite altimetry or moored ADCPs—may potentially provide a more direct indicator
459 of elevated extreme ISW risk, though the limited record length (13 months) precludes
460 formal validation of predictive skill, particularly during the October–November
461 transition period when ITF weakening and extreme ISW occurrence most frequently
462 coincide (Gong et al., 2022).

463 **4.4 Limitations and Future Work**

464 Several limitations of the present study warrant acknowledgement. First, our 13-
465 month record, while sufficient to capture a full seasonal cycle, precludes assessment
466 of interannual variability. ITF transport is known to vary significantly with ENSO and
467 the Indian Ocean Dipole (Matthews et al., 2011; Hu et al., 2022), and these
468 interannual signals are expected to modulate both the mean ITF strength and the
469 frequency of intraseasonal velocity pulses. Multi-year mooring records would be
470 required to assess whether the amplitude-dependent mechanisms identified here are
471 robust across ENSO phases. Second, with only a single mooring, we cannot directly
472 observe the ISW generation process at the sill or track individual wave packets from
473 generation to the observation point. The inferred hydraulic-control mechanism
474 remains physically plausible based on the velocity anomaly composites, but direct
475 verification would require an array of moorings spanning the sill and the northern
476 strait. Third, the two-layer model used to estimate effective phase speed is a
477 simplification; in reality, the vertical mode structure of ISWs in the Lombok Strait is
478 influenced by the continuously stratified water column and the depth-varying ITF
479 shear profile. More sophisticated estimates using continuously stratified weakly



480 nonlinear theory (e.g., Yang et al., 2021) could better quantify the nonlinear
481 amplification potential under varying background conditions. To assess robustness,
482 we extended the event-scale dataset to include 37 additional ISW events from
483 November–December 2024 using GLORYS12-supplemented background velocity
484 ($n = 337$ total; 10 additional extreme events). The correlations with $V_{100-250}$ ($r = +0.22$,
485 $p < 0.001$) and c^{eff} ($r = +0.26$, $p < 0.001$) remain highly significant, indicating that the
486 background flow compensation mechanism is not sensitive to the choice of data
487 source for the extended observation period.

488

489 **5 Conclusions**

490 We have presented the first continuous subsurface mooring analysis of internal
491 solitary wave (ISW) seasonal variability and amplitude-dependent dynamics in the
492 Lombok Strait, based on 13 months of concurrent thermistor and ADCP observations
493 (November 2023–December 2024). The key conclusions are as follows.

494 1. ISW occurrence is strongly seasonal, peaking in boreal winter–spring (March 2024:
495 45 wave packets/month) and reaching a minimum in boreal summer (July 2024: 3
496 packets/month). Over the full record, 302 wave packets and 618 individual solitons
497 were identified, with amplitudes ranging from 25 to 104 m.

498 2. A sensitivity analysis of stratification metrics reveals that the apparent N^2 –ISW
499 correlation from the moving-window estimate ($r = -0.60$, $p = 0.05$) is a
500 methodological artefact: when N^2 is computed at a fixed depth (100–200 m) or as the
501 column maximum (50–300 m), correlations with ISW occurrence are non-significant



502 ($r = -0.43$ and -0.10 respectively, both $p > 0.1$); the 18°C isotherm depth itself is also
503 uncorrelated with ISW counts ($r = +0.08$, $p = 0.82$). Stratification is therefore not an
504 independent control factor for ISW occurrence in either amplitude class. The moving-
505 window N^2 co-varies with ISW occurrence because the ITF simultaneously depresses
506 the isotherm (raising local N^2) and suppresses ISW generation via Doppler retardation,
507 producing a spurious correlation.

508 3. Extreme ISWs (>65 m, $\sim 9\%$ of all packets) appear dynamically decoupled from
509 stratification ($r = -0.25$, $p = 0.417$) and instead showed a tendency to coincide with
510 weakened ITF thermocline-layer velocity (monthly-scale $r = +0.57$, $p = 0.066$;
511 marginal). The event-scale scatter plot (Fig. 4b) yields $r = +0.45$ ($p = 0.07$) for extreme-
512 class ISWs only, measuring amplitude modulation by concurrent ITF velocity—a
513 distinct analysis from the monthly-scale correlation cited above. Seventy-two-hour
514 low-pass filtered ADCP records show that extreme events tend to coincide with near-
515 zero or northward velocity pulses in the 50–200 m layer (composite mean:
516 $+0.005$ m s^{-1} versus record mean of -0.147 m s^{-1}), while the deep layer (200–450 m)
517 shows negligible change.

518 4. We propose a background flow compensation mechanism in which episodic
519 weakening ITF is hypothesized to reduce Doppler retardation of northward ISW
520 propagation, potentially creating dynamic windows for nonlinear wave amplification.
521 The proposed mechanism appears to be amplitude-selective: it appears to govern the
522 generation of extreme events while leaving the production of normal-amplitude ISWs
523 — which is primarily suppressed by seasonal ITF intensification rather than by



524 episodic ITF weakening — largely unaffected by the compensation mechanism. This
525 duality—seasonal ITF suppression for normal ISWs, episodic ITF compensation for
526 extreme ISWs—represents an important but previously underexplored aspect of ISW
527 dynamics in ITF-influenced passages and has implications for ISW hazard assessment
528 and seasonal prediction in the Lombok Strait .

529

530 **Financial support.** This study is supported by the Strategic Priority Research
531 Program of the Chinese Academy of Sciences (XDB42000000), the MFA
532 Construction and Demonstration of Sino-Indonesian Ecological Marine Ranching
533 Program, the CAS-ANSO Sustainable Development Research Project (CAS-ANSO-
534 SDRP-2024-02), and the National Natural Science Foundation of China (41906012).

535

536 **Author contributions.** ZW conceived and designed the study, led the mooring
537 deployment and recovery, performed the data analysis, and wrote the manuscript. JJW
538 and GMF contributed to data processing and visualization. RL, WDM, DW, NPP, and
539 UA contributed to fieldwork and data collection. AP, DN, MFI, PDS, and MF led the
540 Indonesian research coordination and cruise operations. JNW and FW provided
541 scientific guidance and reviewed the manuscript. All authors reviewed and approved
542 the final manuscript.

543 **Acknowledgements**

544 We express our gratitude to the captain and crew on RV Geomarin III for their
545 outstanding performance and invaluable support throughout the cruises. We also want



546 to thank the IMPOLSE joint cruise between IOCAS and BRIN. This collaboration
547 work is facilitated by the platform of Sino-Indonesia Joint Laboratory for Marine
548 Sciences (SIMS). A.P. thanks the President’s International Fellowship Initiative (PIFI)
549 Program of the Chinese Academy of Sciences; D.N. thanks the ANSO visiting
550 fellowship program. Artificial intelligence tools (Claude) were used to assist with
551 language editing and text refinement. The authors are fully responsible for the
552 scientific content and interpretations presented.

553

554 **Code and data availability**

555 The mooring data (CTD, temperature and ADCP) and derived ISW detection results
556 used in this study can be requested at

557 <http://dx.doi.org/10.12157/IOCAS.20260428.001>. The GLORYS12 reanalysis data

558 (dataset: cmems_mod_glo_phy_my_0.083deg_P1D-m) used to extend the

559 background velocity time series for November – December 2024 are freely available

560 from the Copernicus Marine Service:

561 https://data.marine.copernicus.eu/product/GLOBAL_MULTIYEAR_PHY_001_030/d

562 escription, DOI: <https://doi.org/10.48670/moi-00021>

563

564 **Competing interests**

565 The authors declare that they have no conflict of interest.

566

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