

# Sensitivity of bi-spectral retrievals to the fixed effective variance assumption in ship tracks

Iarla Boyce<sup>1</sup>, Alice Cicirello<sup>1</sup>, and Edward Gryspeerdt<sup>2</sup>

<sup>1</sup>Centre for Climate Repair, University of Cambridge, Cambridge, UK

<sup>2</sup>Department of Physics, Imperial College London, London, UK

**Correspondence:** Iarla Boyce (ib541@cam.ac.uk)

**Abstract.** Ship tracks, bright lines in clouds formed by ship exhaust, serve as “natural laboratories” for investigating aerosol-cloud interactions, one of the largest sources of uncertainty in the human forcing of the climate. Observing ship tracks has been used to help constrain the effect of anthropogenic aerosols on cloud brightness, amount and water content. The validity of these constraints relies, in part, on the accuracy of satellite retrieval algorithms used to measure cloud properties. A known source of uncertainty in these algorithms is the representation of the droplet size distribution. Standard bi-spectral retrievals (e.g. MODIS) rely on a fixed effective variance ( $v_{\text{eff}}$ ) for the modified gamma distribution used to model cloud droplet dispersion. The introduction of aerosols into clean, marine clouds produces not only smaller droplets but also a narrower size distribution, contradicting this fixed assumption. This study presents a controlled, synthetic retrieval experiment that isolates the sensitivity of cloud property retrievals, and the derived aerosol-cloud interaction metrics, to this single assumption. This study uses idealised ship tracks as the test case because they provide the strongest realistic contrast between the assumed and true distribution widths. The results produced indicate that neglecting the narrowing of the droplet size distribution causes retrieved effective radius ( $r_e$ ) to differ systematically between clean and polluted regimes. The polluted branch is overestimated relative to the near-unbiased clean branch by approximately 2.3–3.4%, depending on the assumed retrieval baseline. Optical depth ( $\tau$ ) is virtually unaffected in either regime. LWP is retrieved with little bias in either regime, so the clean-to-polluted LWP contrast is largely preserved.  $N_d$  shows the opposite problem. The apparent clean-to-polluted increase in  $N_d$  is overstated by 22–23% under both operational assumptions. This discrepancy is driven by the inverse dependence of  $N_d$  on the spectral width parameter  $k$ , inflating the droplet count in narrow polluted distributions while underestimating it in broader clean ones. If the aerosol-driven narrowing of the droplet size distribution assumed here is representative of real ship tracks, this contrast inflation could exaggerate the apparent susceptibility of clouds to aerosols, contributing to an overstatement of the Twomey effect in observation-based estimates reliant on data from ship tracks. Under the same conditions, satellite-based monitoring of climate intervention efforts, such as marine cloud brightening, could similarly overestimate their efficacy.

## 1 Introduction

Aerosols influence cloud radiative properties through an instantaneous microphysical change followed by a series of time-dependent adjustments. The primary instantaneous effect occurs as aerosols act as cloud condensation nuclei (CCN), increas-

25 ing the droplet number concentration ( $N_d$ ). For a constant liquid water path (LWP), this shift toward more numerous, smaller droplets increases the cloud's total surface area, thereby enhancing its optical depth ( $\tau$ ) and albedo (Twomey, 1974). Beyond this initial response, several rapid adjustments further modify the clouds' microphysical state. A key mechanism is precipitation suppression, where smaller droplets inhibit collision-coalescence, potentially increasing both LWP and cloud lifetime (Albrecht, 1989). However, the initial Twomey cooling is often offset by rapid cloud adjustments. For example, altered micro-

30 physics can accelerate the entrainment of dry air into the cloud top, which may enhance evaporation and lead to a subsequent decrease in LWP (Wood, 2007; Ackerman et al., 2004; Bretherton et al., 2007). Furthermore, light-absorbing aerosols like black carbon can introduce a semi-direct effect, heating the surrounding air and further inducing cloud dissipation (Koch and Del Genio, 2010). While these specific responses counteract the initial brightening, the overall contribution of aerosol-cloud interactions induces a net cooling effect, the magnitude of which is poorly constrained (Glassmeier et al., 2021; Chun et al.,

35 2023; Christensen et al., 2021).

Ship tracks serve as an observational laboratory for disentangling these competing cloud adjustments. By providing a localised aerosol perturbation against a clear control case of unpolluted cloud, ship tracks allow the isolation of these microphysical processes from the background meteorological conditions. Constraining the magnitude of this aerosol-induced cooling is essential not only for greater accuracy in radiative forcing estimates but also for ascertaining the feasibility of climate engineering techniques such as MCB (Feingold et al., 2024). While the instantaneous Twomey effect is well documented in

40 observational studies (Noone et al., 2000a, b), the subsequent LWP response is significantly less constrained (Chen et al., 2012; Gryspeerdt et al., 2021). This uncertainty arises because the net LWP signal is often weak (Toll et al., 2019; Tippet et al., 2024) and highly dependent on both the local meteorology (Goren et al., 2025; Zhang and Feingold, 2023) and the evolving timescales of the perturbation (Gryspeerdt et al., 2021).

45 To date, observational studies of ship tracks have primarily utilised satellite remote sensing data from instruments such as the Moderate Resolution Imaging Spectroradiometer (MODIS). MODIS uses a multi-spectral imager that collects reflectance data at multiple relevant wavelengths in the visible and infra-red. For this study, bi-spectral retrievals using the MODIS wavelength bands in the visible-near infrared and shortwave infrared are examined (King et al., 1997). These are sensitive to cloud optical thickness ( $\tau$ ) and droplet effective radius ( $r_e$ ), respectively, but not entirely independent (Nakajima and King, 1990). As the

50 cloud property retrieval algorithms are typically under-constrained, they use a variety of simplifying assumptions about the cloud microphysical properties. One of these is the assumption of a globally fixed width of the cloud droplet size distribution, or effective variance ( $v_{\text{eff}}$ ) (King et al., 1997). The effective variance is a measure of the diversity of droplet sizes, and its value is fundamentally altered by the aerosol perturbation in a ship track. In the clean background cloud, the process of droplets colliding and merging to form drizzle creates a broad distribution of sizes, resulting in a high  $v_{\text{eff}}$  (typically 0.10–0.15). The

55 introduction of ship-emitted aerosols narrows this distribution as the water in the cloud spreads to form smaller, uniform droplets around the aerosols (Martin et al., 1994), reducing the  $v_{\text{eff}}$  to as little as 0.05; this is further supported by the more recent aircraft-derived  $k$ - $N_d$  relationship of Lebsock and Witte (2023), discussed further in Sect. 2.3. This discrepancy is a recognised retrieval bias (Liu et al., 2008), as the relationship between reflectance and  $r_e$  is a function of  $v_{\text{eff}}$  (Nakajima and King, 1990). It has also been hypothesised by Gryspeerdt et al. (2021) that this is the cause of the apparent instantaneous LWP

60 response to ship aerosol, but this has yet to be properly quantified. While the sensitivity of retrieved cloud properties to the assumed droplet size distribution is established in the retrieval literature, this study quantifies the resulting bias specifically for the clean-to-polluted contrast characteristic of ship tracks, where the fixed-variance assumption is particularly likely to break down, and propagates it through to its consequences for  $N_d$  and derived susceptibility metrics, characterises its dependence on viewing geometry, and tests its robustness to the assumed strength of the underlying  $k$ - $N_d$  coupling. Susceptibility metrics of  
 65 this kind underpin observational constraints on the radiative forcing from aerosol-cloud interactions (RFaci). Biases in them therefore propagate directly into RFaci uncertainty.

Here, a synthetic retrieval framework designed to quantify this effect is presented, bounding the limits of the fixed effective variance assumption across core optical and microphysical cloud products. By forward-modelling top-of-atmosphere radiances for synthetic ship track scenes and inverting them using standard operational logic, the retrieval bias caused solely by the fixed  
 70  $v_{\text{eff}}$  assumption is isolated from all other components of the operational retrieval chain.. The goal is the clean-to-polluted contrast that ship-track studies rely on, not absolute retrieval accuracy in either regime alone. The errors introduced into direct measurements of effective radius and optical depth are first quantified, before examining how these errors propagate into derived liquid water path and droplet number concentrations. Finally, the importance of mitigating these biases is discussed in the context of accurately interpreting cloud susceptibility and assessing the feasibility of marine cloud brightening.

## 75 2 Methodology

### 2.1 The Modified Gamma Distribution

Satellite retrievals approximate the microphysics of clouds by relying on an assumed droplet size distribution. The droplet size distribution is typically described by a modified gamma distribution defined by Hansen and Travis (1974), and is given by:

$$n(r) = Ar^{\frac{1-3v_{\text{eff}}}{v_{\text{eff}}}} \exp\left(-\frac{r}{r_e v_{\text{eff}}}\right), \quad (1)$$

80 where  $n(r)$  is the number of droplets of radius  $r$  and  $A$  is a normalisation constant. This distribution is characterised by two parameters: effective radius ( $r_e$ ) and the effective variance ( $v_{\text{eff}}$ ). Physically,  $r_e$  is defined as the area-weighted mean radius of the droplet population:

$$r_e = \frac{\int_0^\infty r^3 n(r) dr}{\int_0^\infty r^2 n(r) dr}, \quad (2)$$

whereas  $v_{\text{eff}}$  is a dimensionless parameter representing the variance of the droplet size distribution normalised by  $r_e^2$ :

$$85 \quad v_{\text{eff}} = \frac{1}{Gr_e^2} \int_0^\infty (r - r_e)^2 \pi r^2 n(r) dr, \quad (3)$$

where  $G$  is the total geometric cross-sectional area per unit volume. While  $r_e$  is explicitly retrieved from short-wavelength infrared reflectance,  $v_{\text{eff}}$  must be assumed (Martin et al., 1994), with this study setting this at a constant value  $v_{\text{eff}} = 0.10$  to align with the operational assumption employed by the MODIS Collection 6 retrieval algorithm for liquid clouds. This fixed assumption introduces a systematic bias, as the  $v_{\text{eff}}$  of a cloud varies significantly with aerosol loading (Martin et al., 1994), entrainment (Lim and Hoffmann, 2026), and precipitation onset (Miles et al., 2000). This is especially problematic when retrieving polluted clouds, as the relationship between the radiative properties and  $N_d$  is dependent on the dispersion parameter,  $k$ , which relates the volume mean radius ( $r_v$ ) to  $r_e$  (Martin et al., 1994):

$$k = \left( \frac{r_v}{r_e} \right)^3 = (1 - v_{\text{eff}})(1 - 2v_{\text{eff}}) \quad (4)$$

As  $v_{\text{eff}}$  reduces,  $k$  increases, ultimately approaching unity. From Eq. (4) it is clear that by fixing  $v_{\text{eff}}$ ,  $k$  is also essentially fixed. This has implications for deriving  $N_d$ , typically calculated via the adiabatic assumption (Grosvenor et al., 2018):

$$N_d = \frac{3 \cdot \text{LWP}}{4\pi\rho_w k r_e^3 H}. \quad (5)$$

Here  $H = 500$  m is the fixed geometric cloud thickness, used consistently in both cloud generation and  $N_d$  derivation to convert LWP to mean liquid water content ( $\text{LWC} = \text{LWP}/H$ ). If the true  $k$  of polluted clouds exceeds the assumed  $k$  in Eq. (5), the retrieval will systematically overestimate  $N_d$ . In contrast, LWP is relatively insensitive to the  $v_{\text{eff}}$  assumption. LWP is formulated as (Stephens, 1978)

$$\text{LWP} = \frac{2}{3}\rho_w\tau r_e \quad (6)$$

and is thus linearly dependent on  $r_e$  and optical depth ( $\tau$ ) only. Therefore, LWP is only affected by secondary biases through  $r_e$ , and not directly biased by the  $v_{\text{eff}}$  assumption. As such, it is expected that LWP will be retrieved more robustly. Under the adiabatic assumption, Eq. (5) and Eq. (6) are mutually consistent: substituting Eq. (6) into Eq. (5) recovers  $N_d = \tau / (2\pi k r_e^2 H)$ . As a sensitivity comparison, the legacy pre-Collection 6 assumption of  $v_{\text{eff}} = 0.13$  is also evaluated throughout this study.

## 2.2 Retrieval Principle

The retrieval of  $\tau$  and  $r_e$  is commonly performed using the bi-spectral reflectance method set out by Nakajima and King (1990). This approach exploits the differing radiative properties of cloud droplets across the solar spectrum. A non-absorbing near-infrared band ( $0.86 \mu\text{m}$ ) is used to obtain  $\tau$ , as it is primarily sensitive to the cloud's total scattering cross section. For  $r_e$ , an absorbing shortwave infrared band ( $2.1 \mu\text{m}$ ) is employed, due to the sensitivity of liquid water absorption to droplet radius.

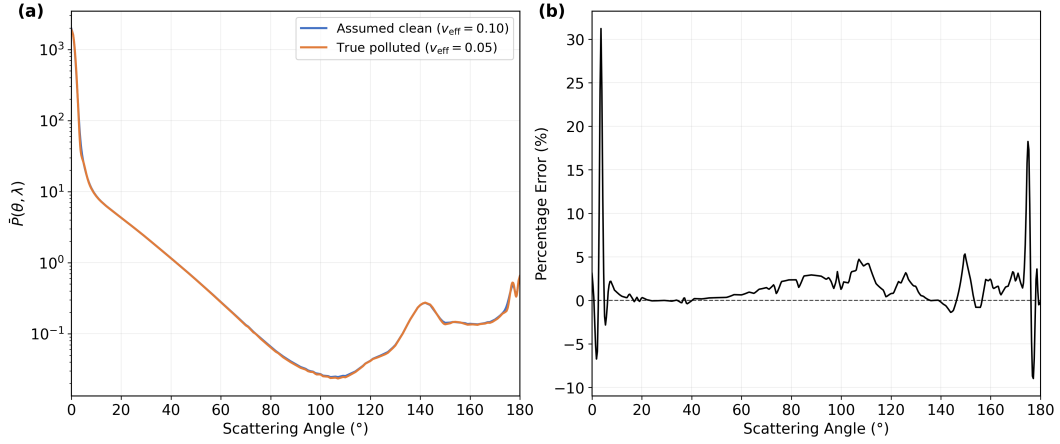
The fixed  $v_{\text{eff}}$  assumption leads to a discrepancy between the scattering phase function assumed in the look-up-table generation and the true scattering behaviour of the observed cloud. The physical origin of this discrepancy lies in the calculation of the cloud's bulk scattering phase function,  $\bar{P}(\theta, \lambda)$ , where  $\theta$  is the scattering angle and  $\lambda$  is the wavelength. This function is derived

by integrating the single-droplet Mie scattering phase function,  $P(\theta, r, \lambda)$ , over the entire cloud droplet size distribution,  $n(r)$ ,  
 115 described in Eq. (1). The bulk scattering phase function is calculated as (Hansen and Travis, 1974):

$$\bar{P}(\theta, \lambda) = \frac{\int_0^\infty \sigma_{scat}(r, \lambda) \cdot P(\theta, r, \lambda) \cdot n(r) dr}{\int_0^\infty \sigma_{scat}(r, \lambda) \cdot n(r) dr}, \quad (7)$$

where  $\sigma_{scat}$  is the scattering cross-section.

As shown in Eq. (1),  $n(r)$  is a function of both  $r_e$  and  $v_{eff}$ , the resulting bulk scattering phase function,  $\bar{P}(\theta, \lambda)$ , is also dependent on both parameters. This is illustrated in Figure 1, which compares the bulk scattering phase function calculated  
 120 across the plausible range of assumed retrieval  $v_{eff}$  (0.10, shown here, and 0.13) against the true polluted  $v_{eff} = 0.05$ ; the discrepancy scales with the assumed–true mismatch, growing larger as the assumed value moves further from the true polluted state. Across most scattering angles ( $20^\circ$ – $140^\circ$ ) the percentage error between the two phase functions remains below 5%, but it grows sharply toward the forward-scattering and backscattering limits, reaching approximately +31% near  $\Theta \approx 5^\circ$  and +18% near  $\Theta \approx 175^\circ$ , where diffraction and glory features are most sensitive to the underlying droplet size distribution. This  
 125 indicates that an algorithm with a fixed  $v_{eff}$  encountering a cloud with a different  $v_{eff}$  will assume a biased bulk phase function, thus misinterpreting the backscattered light intensity. This may lead to a systematic retrieval error.



**Figure 1.** (a) Comparison of the bulk scattering phase function,  $\bar{P}(\theta, \lambda)$ , evaluated at  $\lambda = 0.86 \mu m$  for a droplet effective radius of  $8 \mu m$ , for a true polluted cloud ( $v_{eff} = 0.05$ ) and the assumed clean cloud ( $v_{eff} = 0.10$ ). The y-axis is presented on a logarithmic scale. (b) The resulting percentage error in the phase function due to the fixed  $v_{eff}$  assumption.

### 2.3 Data generation and forward simulation

A population of 5000 clean clouds is generated with properties randomly sampled from ranges of typical non-drizzling marine stratocumulus clouds (Wood, 2012), with  $r_e$  in the range  $10$ – $18 \mu m$  and  $\tau$  in the range  $8$ – $30$ . This sample size ensures dense  
 130 coverage of the operational lookup table (LUT) space. The  $v_{eff}$  for these clouds is assigned one of  $v_{eff} = [0.1, 0.13, 0.15]$ , a

standard range for such clouds (Martin et al., 1994). For each clean case, a corresponding “polluted” case (representing a ship track) is generated by reducing  $r_e$  by a random factor between 15–35% (Grosvenor et al., 2018) while adjusting  $\tau$  in tandem to conserve LWP. This ensures that any change in radiative properties between the clean and polluted pair is due solely to the redistribution of water among a larger number of droplets. This is done to isolate the instantaneous Twomey effect from the adjustments. The polluted clouds are assigned a fixed  $v_{\text{eff}} = 0.05$  representing a strongly narrowed droplet size distribution as a lower bound to explore the maximum potential bias. While marine stratocumulus clouds generally exhibit broader distributions, the compilation of in-situ measurements by Miles et al. (2000) demonstrates that polluted air masses frequently produce highly narrowed distributions, including numerous recorded instances of the  $v_{\text{eff}}$  dropping below this bound. This choice is further supported by the empirical  $k$ - $N_d$  relationship derived from aircraft observations by Lebsock and Witte (2023): at a typical ship-track droplet concentration of  $N_d \approx 100 \text{ cm}^{-3}$ , their Combined Fit gives  $k \approx 0.84$ , which corresponds via Eq. (4) to  $v_{\text{eff}} \approx 0.05$ , with heavier pollution implying still narrower distributions.

Radiative transfer simulations are run using the DISORT algorithm (Stamnes et al., 2000) in the libRadtran package (Emde et al., 2016) driven by the pyLRT python wrapper (Gryspeerdt and Driver, 2024) for the 5000 clean/polluted cloud pairs. For each case, the top-of-atmosphere (TOA) reflectance function  $R$  is calculated following the definition of Nakajima and King (1990):

$$R(\tau_c; \mu, \mu_0, \phi) = \frac{\pi I(0, -\mu, \phi)}{\mu_0 F_0}, \quad (8)$$

where  $I(0, -\mu, \phi)$  is the reflected intensity at the TOA,  $\mu$  is the cosine of the viewing zenith angle,  $\phi$  is the relative azimuth angle,  $\mu_0$  is the cosine of the solar zenith angle, and  $F_0$  is the incident solar flux density. The simulations are performed for two standard MODIS bands centred at  $0.86 \mu\text{m}$  and  $2.1 \mu\text{m}$ , across a range of solar zenith angles ( $20^\circ$ ,  $40^\circ$ ,  $60^\circ$ ). Ranges of viewing zenith angle ( $0^\circ$ ,  $20^\circ$ ,  $40^\circ$ ,  $60^\circ$ ) and the relative azimuth angle ( $0^\circ$ ,  $45^\circ$ ,  $90^\circ$ ,  $135^\circ$ ,  $180^\circ$ ) are also chosen, to represent the full combination of observation geometries. The scattering angle,  $\Theta$ , is related to the solar zenith angle ( $\theta_0$ ), the viewing zenith angle ( $\theta_v$ ), and the relative azimuth angle ( $\phi$ ) by the spherical law of cosines:

$$\cos \Theta = -\cos \theta_0 \cos \theta_v + \sin \theta_0 \sin \theta_v \cos \phi. \quad (9)$$

The clouds are modelled as single, vertically homogeneous 1D plane-parallel slabs with a geometric thickness of 500 m, a reasonable value for marine stratocumulus (Wood, 2012). The microphysical properties are determined using pre-calculated lookup tables generated using the libRadtran Mie tool (Emde et al., 2016). For these calculations, the cloud droplet size distribution is modelled using the modified gamma distribution defined in Eq. (1). Within the libRadtran Mie tool, the shape of this distribution is dictated by the parameter  $\alpha$ , which corresponds to the exponent of the  $r$  term in Eq. (1) and is related to  $v_{\text{eff}}$  by (Hansen, 1971):

$$\alpha = \frac{1 - 3v_{\text{eff}}}{v_{\text{eff}}} \quad (10)$$

Each lookup table contains the bulk optical properties required to solve the radiative transfer equation: the extinction efficiency, single scattering albedo and the angular scattering phase function over a grid of  $r_e$  and wavelengths. Table 1 summarises the microphysical properties and viewing geometries used to generate the synthetic cloud population.

**Table 1.** Summary of microphysical properties and viewing geometries used to generate the synthetic cloud population. The polluted cases are derived directly from the clean control cases to conserve Liquid Water Path (LWP).

| Parameter                           | Clean (Background)          | Polluted (Ship Track)           |
|-------------------------------------|-----------------------------|---------------------------------|
| Sample Size ( $N$ )                 | 5000                        | 5000                            |
| Eff. Radius ( $r_e$ )               | Sampled 10–18 $\mu\text{m}$ | $r_e$ reduced by 15–35%         |
| Optical Depth ( $\tau$ )            | Sampled 8–30                | $\tau$ adjusted to conserve LWP |
| Eff. Variance ( $v_{\text{eff}}$ )  | 0.10, 0.13, 0.15            | Fixed 0.05                      |
| Solar Zenith Angle ( $\theta_0$ )   |                             | 20°, 40°, 60°                   |
| Viewing Zenith Angle ( $\theta_v$ ) |                             | 0°, 20°, 40°, 60°               |
| Relative Azimuth Angle ( $\phi$ )   |                             | 0°, 45°, 90°, 135°, 180°        |

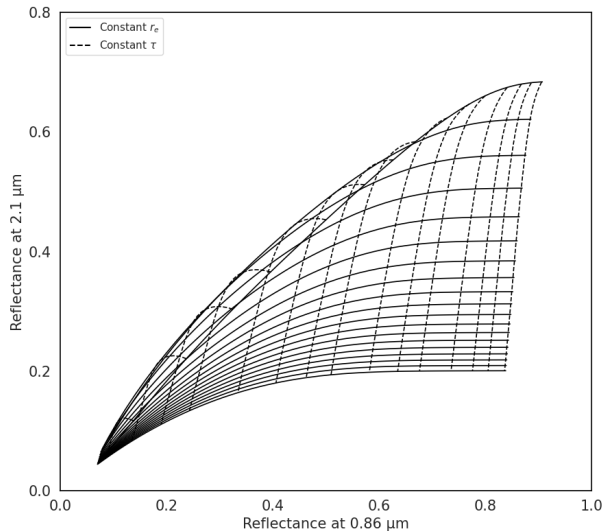
## 2.4 Retrieval Simulation

165 The simulated retrieval of cloud properties from the simulated reflectances mimics MODIS retrievals in two spectral bands (King et al., 1997). This is achieved by mapping TOA reflectances to  $r_e$  and  $\tau$  onto a Nakajima-King grid (Nakajima and King, 1990). For the simulated retrieval, the grid in Figure 2 is generated using the assumed  $v_{\text{eff}}$  of 0.10. The retrieval LUT spans effective radii from 2.0 to 20.0  $\mu\text{m}$  in steps of 0.1  $\mu\text{m}$ , and optical depths from 0.5 to 61, with a finer spacing of 0.25 between  $\tau = 0.5$  and  $\tau = 10$  to resolve the non-linear curvature of the Nakajima–King retrieval surface and coarser steps of 1.0

170 for  $\tau > 10$  where reflectance saturates, following the variable-resolution LUT design used operationally for other bi-spectral retrievals (Liu et al., 2023). The LUT covers the same angle grid used for the forward simulations (SZA: 20°, 40°, 60°; VZA: 0°–60°; RAZ: 0°–180°). Cloud properties are retrieved by trilinear interpolation across viewing geometry, followed by linear interpolation within the Delaunay triangulation of the two-dimensional reflectance grid at each geometry node. The fixed  $v_{\text{eff}}$  assumption creates a systematic physical inconsistency when the  $v_{\text{eff}}$  of clouds varies, especially in cases such as ship tracks

175 where  $v_{\text{eff}}$  can undergo a dramatic reduction (Martin et al., 1994; Miles et al., 2000). The percentage bias is calculated as the relative difference between the properties retrieved using the assumed  $v_{\text{eff}} = 0.10$  and the true properties used to generate each synthetic cloud. This study adopts  $v_{\text{eff}} = 0.10$  as its central retrieval baseline, in line with the current MODIS Collection 6 algorithm, which was revised from a previous value of 0.13 to better align with multi-angle observations (Platnick et al., 2016). The legacy  $v_{\text{eff}} = 0.13$  assumption is retained throughout as a secondary sensitivity comparison, since it maximises the

180 contrast between the assumed and true polluted state  $v_{\text{eff}}$  (0.05) and thus provides an upper-bound estimate for the retrieval bias relative to current operational algorithms. Importantly, this fixed-variance bias is bidirectional; while both the 0.10 and 0.13 baselines systematically bias the narrowed polluted clouds, tuning an algorithm to assume the narrowed 0.05 state would invert the problem, introducing a bias into the retrieval of the unpolluted background clouds.



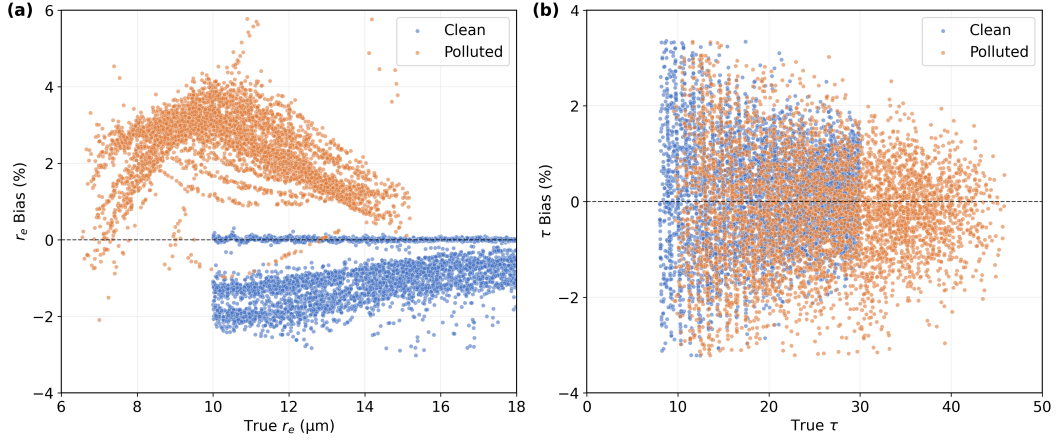
**Figure 2.** Nakajima-King retrieval grid representing the central operational assumption ( $v_{\text{eff}} = 0.10$ ), evaluated at  $\text{SZA} = 40^\circ$ ,  $\text{VZA} = 0^\circ$ ,  $\text{RAZ} = 0^\circ$ . The grid maps the top-of-atmosphere reflectances at  $0.86 \mu\text{m}$  and  $2.1 \mu\text{m}$  to a unique solution for  $\tau$  (dashed lines) and  $r_e$  (solid lines).

### 3 Results

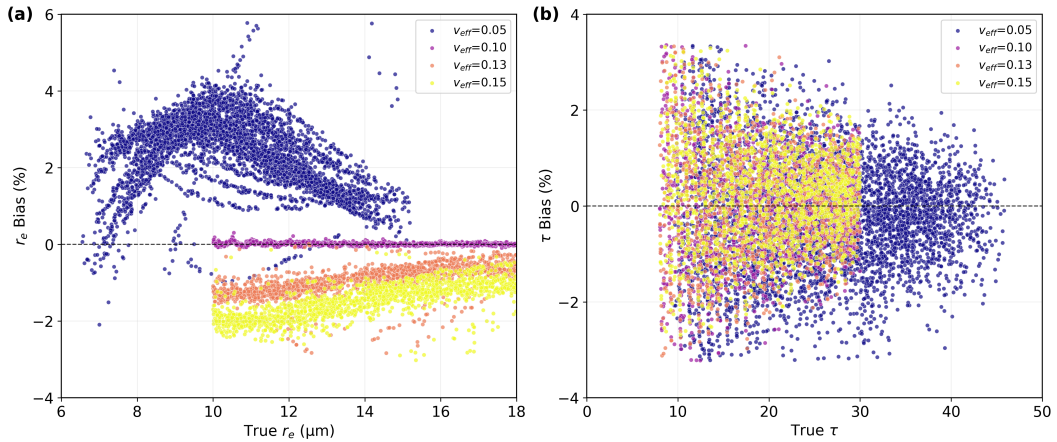
#### 185 3.1 Bias in Direct Retrieval Products ( $r_e$ , $\tau$ )

The retrieval bias is first quantified for the directly retrieved variables,  $r_e$  and  $\tau$ , using the central  $v_{\text{eff}} = 0.10$  retrieval baseline. Figure 3 displays the percentage bias retrieved  $r_e$  and  $\tau$  as a function of the true  $r_e$  and  $\tau$ , respectively. The two regimes, clean and polluted, are clearly separated by bias magnitude in  $r_e$ , but negligible bias is shown for  $\tau$  in both regimes. For  $\tau$ , the mean percentage bias in clean clouds is 0.07% and in polluted clouds -0.15%. It is shown in Figure 3 that optically thin clouds have a wider bias distribution in both regimes, with potentially compensating errors causing mean bias to be small. For  $r_e$  clean clouds, bias is negligible, with three distinct clusters.

Figure 4 demonstrates that the data points are clustered by  $v_{\text{eff}}$ , with the retrieval baseline 0.10 clustered around the 0% bias line and 0.13 and 0.15 biases calculated as -1.05% and -1.56%, respectively. Thus, the mean bias in the retrieved  $r_e$  of clean clouds is calculated to be -0.87% (averaged over the synthetic population as defined in Sect. 2.3). Polluted clouds, however, show a more substantial bias percentage with a mean bias of 2.30%. This occurs because a polluted cloud transition is assumed to involve a simultaneous increase in  $\tau$  to maintain a constant LWP as  $r_e$  decreases. The discrepancy in the true  $v_{\text{eff}}$  state explains why the bias profiles for the clean and polluted regimes remain distinct, even when compared to the same  $r_e$  value. As a sensitivity comparison, adopting the legacy  $v_{\text{eff}} = 0.13$  baseline instead nearly doubles the polluted  $r_e$  bias to 3.43%, since the contrast with the true polluted state ( $v_{\text{eff}} = 0.05$ ) is larger. It is shown in Figure 4b that the specific  $v_{\text{eff}}$  value has little effect on  $\tau$  bias, which is to be expected given  $\tau$  is known to be insensitive to the shape of the droplet size distribution.

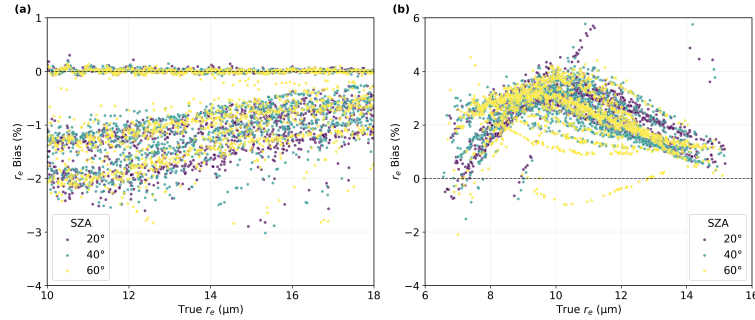


**Figure 3.** The percentage bias in retrieved (a)  $r_e$  and (b)  $\tau$  as a function of their true values, for clean (orange) and polluted (blue) clouds. All viewing geometries are included; for display clarity, points outside the 1st–99th percentile of each bias variable are omitted, while all statistics quoted in the text are computed on the unfiltered population.



**Figure 4.** The percentage bias in retrieved cloud properties as a function of their true values, coloured by effective variance ( $v_{\text{eff}}$ ). (a) effective radius ( $r_e$ ) and (b) optical thickness ( $\tau$ ). Note the distinct clustering of bias based on the underlying distribution width, with  $v_{\text{eff}} = 0.10$  (the retrieval assumption) centred on the zero-bias line. All viewing geometries are included; points outside the 1st–99th percentile of each bias variable are omitted for display clarity, while quoted statistics use the unfiltered population.

The sensitivity of the bias to SZA in the polluted  $r_e$  regime is further investigated, as shown in Figure 5. In Figure 5a the clean clouds are shown coloured by three chosen values of SZA, 20°, 40°, and 60°. The data shows that for clean clouds, the mean bias changes by 0.09% as the sun moves from 20° to 60°. Thus, it is concluded that there is no significant contribution to the bias from SZA. Figure 5b, showing the polluted regime, however, displays more separation due to SZA. Analysing the  
 205 contribution to the mean  $r_e$  of each SZA value shows this is a secondary contribution to the overall bias, with a mean bias of 2.47%, 2.41%, and 2.00% for SZA values 20°, 40°, and 60° respectively. This small SZA dependence arises from the change in scattering angle observed by the satellite, changing the discrepancy between the assumed and true phase functions. This demonstrates that lower SZA have an increased bias over higher SZA, but the mean effect is quite small and can be ignored in the context of the  $v_{\text{eff}}$  bias.

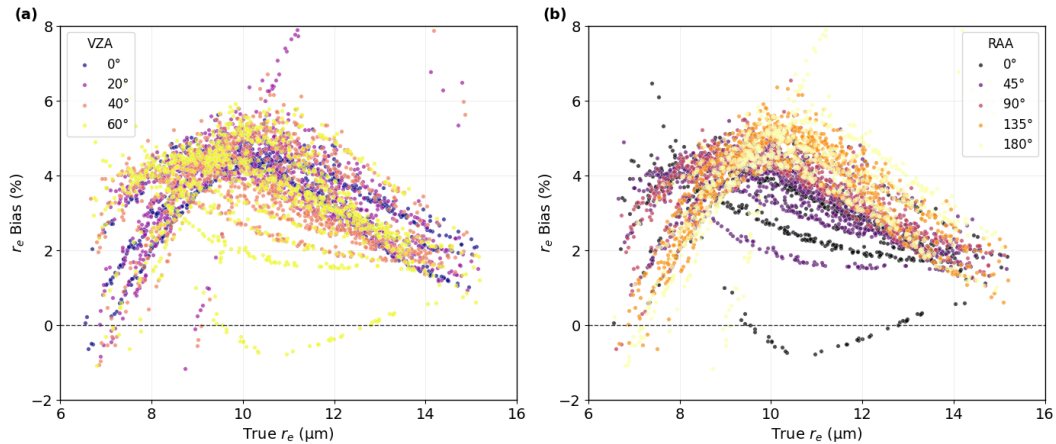


**Figure 5.** The percentage bias in retrieved effective radius ( $r_e$ ) coloured by Solar Zenith Angle (SZA). Clean clouds show negligible angular dependence. Polluted clouds show a secondary dependence on scattering geometry, with lower SZA exhibiting slightly higher bias than higher SZA. All viewing geometries are included; points outside the 1st–99th percentile of each bias variable are omitted for display clarity, while quoted statistics use the unfiltered population.

210 The sensitivity of  $r_e$  to Viewing Zenith Angle (VZA) and Relative Azimuth Angle (RAA) in the polluted regime is demonstrated in Figure 6. It is found that the overestimation of  $r_e$  is consistent at approximately 2.4% at all VZA, aside from VZA= 60° at 1.84%. It is important to note that while the mean bias shows a small decrease at VZA= 60° the precision degrades, with the standard deviation increasing to 4.97% from 0.8–2.1% for the other geometries. The bias shows a small dependence on RAA, ranging from 2.01–2.64%. The precision of the retrievals is stable across most azimuths, although the  
 215 standard deviation increases to 5.97% at RAA= 180°, coinciding with the backscatter geometries ( $\Theta \rightarrow 180^\circ$ ) discussed further in Sect. 3.2; RAA remains a secondary driver of bias compared to the  $v_{\text{eff}}$  induced retrieval bias.

### 3.2 Bias in Derived Products (LWP, $N_d$ )

As discussed in Sect. 2.1, LWP and  $N_d$  products are derived from the directly retrieved products. The errors in  $r_e$  and  $\tau$  propagate to the derived variables LWP and  $N_d$  and are amplified or dampened depending on their specific functional form.  
 220 LWP is calculated using the linear relationship shown in Eq.(6). Due to the linear relationship between  $r_e$  and  $\tau$  in LWP, the bias is of a similar magnitude to  $r_e$  but is slightly damped due to the small negative error in  $\tau$ . Figure 7a illustrates the

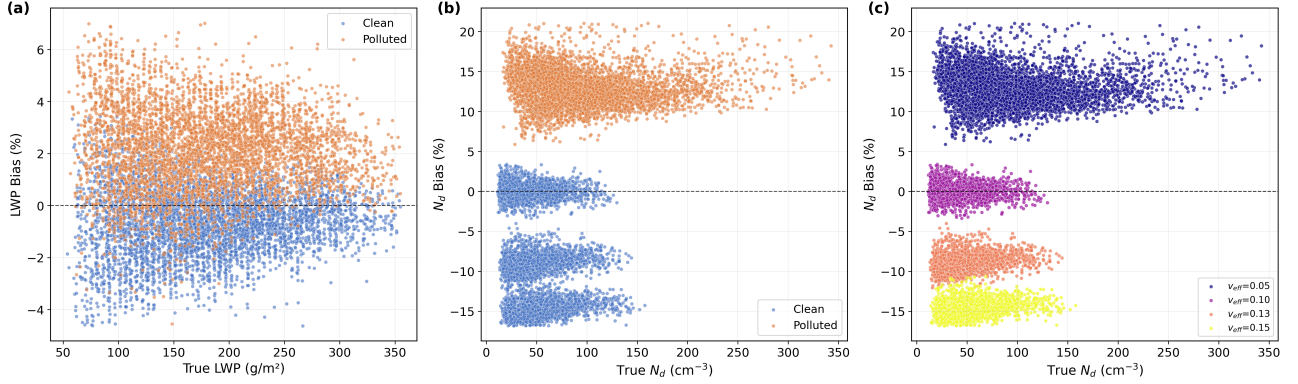


**Figure 6.** The percentage bias in retrieved effective radius ( $r_e$ ) coloured by Viewing Zenith Angle **(a)** and Relative Azimuth Angle **(b)**. **(a)** Clouds show little angular dependence with VZA= 60° showing higher variance than other angles. **(b)** Clouds show little mean dependence on RAA, though variance increases at RAA= 180°. Polluted clouds only; points outside the 1st–99th percentile of each bias variable are omitted for display clarity, while quoted statistics use the unfiltered population.

distribution of bias against the true LWP amount for clean and polluted clouds. The non-linearity in  $r_e$  seen in Figure 3a is lost through the propagation, with the distribution appearing unstructured. The mean bias in the polluted regime is calculated to be 2.15%, slightly less than the  $r_e$  bias due to the small compensating bias in  $\tau$ . This indicates that LWP is relatively robustly  
225 retrieved with respect to uncertainties in  $v_{\text{eff}}$ , with the wrongly assumed droplet size distribution having a small effect on bulk water content. The response of  $N_d$ , however, is much more severe due to its direct dependence on  $k$  defined in (4), which is, in turn, a function of  $v_{\text{eff}}$ .  $N_d$  is proportional to  $r_e^{-3}$ , meaning that the overestimation of  $r_e$  exerts a negative pressure on the derived  $N_d$ , causing an underestimation of the concentration.  $N_d$  is inversely proportional to the  $k$  parameter. The baseline retrieval assumes a distribution of 0.10 corresponding to a lower value of  $k$  ( $k_{\text{assumed}} \approx 0.72$ ). However, the true polluted cloud  
230 has a narrower distribution, corresponding to a  $k$  value closer to unity ( $k_{\text{true}} \approx 0.86$ ). The artificially small  $k_{\text{assumed}}$  causes the algorithm to overestimate the concentration of droplets due to the inverse dependence on  $k$ . This overestimation is illustrated in Figure 7b. A clear separation is seen between the clean and the polluted regime, with a substantial overestimation being seen in the polluted cases. This overestimation is calculated to be a mean bias of 13.91%, a large error. It is noted that the standard deviation of this bias is quite large at  $\pm 22.14$ , showing that the precision is relatively low, with the possibility of excessively  
235 high bias in some  $N_d$  retrievals. As a sensitivity comparison, adopting the legacy  $v_{\text{eff}} = 0.13$  baseline instead produces a substantially larger polluted  $N_d$  bias of 24.56% ( $\pm 23.75\%$ ), since the contrast with the true polluted state ( $v_{\text{eff}} = 0.05$ ) is greater; the central  $v_{\text{eff}} = 0.10$  baseline therefore roughly halves, but does not eliminate, the  $N_d$  overestimation relative to the legacy assumption.

Separation can also be seen in the clean regime of Figure 7b. Figure 7c displays  $N_d$  bias coloured by  $v_{\text{eff}}$ , and illustrates  
240 that there is not only substantive bias in the polluted regime, but also clean regime  $v_{\text{eff}}$  values that do not align with the

algorithm assumption. Clouds with a  $v_{\text{eff}}$  of 0.13 have their droplet number concentration underestimated by an average of 8.26%, whereas clouds with a  $v_{\text{eff}}$  of 0.15 are underestimated by 14.51%. This not only shows that fresh ship tracks with a very narrow distribution have their  $N_d$  significantly overestimated, but that clean clouds broader than the assumed baseline have their  $N_d$  correspondingly underestimated. Evaluating this range of  $v_{\text{eff}}$  values allows for a bounding of the bias associated with the fixed  $v_{\text{eff}}$  assumption. Although this bias range is substantial, this indicates that  $N_d$  is highly sensitive to discrepancies between the assumed and true  $v_{\text{eff}}$  values.



**Figure 7.** The percentage bias in derived cloud products as a function of their true values for clean (orange) and polluted (blue) regimes. (a) Liquid Water Path (LWP) shows a small, unstructured bias. (b) Droplet Number Concentration ( $N_d$ ) shows systematic overestimation in the polluted regime. (c) Percentage bias in retrieved  $N_d$  coloured by effective variance ( $v_{\text{eff}}$ ). Deviations from the assumed  $v_{\text{eff}} = 0.10$  drive the error, with narrow distributions ( $v_{\text{eff}} = 0.05$ , dark blue) causing significant overestimation, and broad distributions ( $v_{\text{eff}} = 0.15$ , yellow) causing underestimation. All viewing geometries are included; points outside the 1st–99th percentile of each bias variable (including the  $N_d$  outliers exceeding 1200% discussed in the text) are omitted for display clarity, while quoted statistics use the unfiltered population.

Further analysis of the  $N_d$  retrieval bias reveals the scattering angle ( $\Theta$ ) has a strong influence on the bias variance. For the vast majority of the simulated polluted clouds, the  $N_d$  overestimation varies between a minimum of +11.8% ( $\Theta \approx 152^\circ$ ) and a maximum of +36.6% at the backscatter peak ( $\Theta = 180^\circ$ ). At this geometry, the  $r_e$  bias drops to  $-0.30\%$ , removing the compensating bias effect of the  $r_e^{-3}$  term and maximising the error driven by  $k$ . Outlier analysis indicates that at extreme swath-edge geometries in the forward-scattering regime (e.g.,  $SZA = 60^\circ$ ,  $VZA = 60^\circ$ ,  $\Theta = 60^\circ$ ), the retrieval algorithm becomes highly unstable. At these specific angles the algorithm occasionally underestimates  $r_e$  by over 70%, which due to the inverse-cube dependence, yields overestimations of  $N_d$  exceeding 1200%. These extreme outliers are rare: across the full synthetic population of  $\sim 10\,000$  cases, approximately 0.01% (one case) exceeds the 1200% threshold, and excluding it shifts the mean polluted  $N_d$  bias by less than 0.3 percentage points (from 13.91% to 13.63%), confirming the mean estimate is robust to outlier removal. Nonetheless, the single outlier and the broader heavy tail drive the standard deviation to  $\pm 22.14\%$ . Standard quality-control filters on retrieved  $r_e$  and  $\tau$  do not remove this case, as it falls within physically plausible retrieval bounds. Dedicated geometric filtering of extreme forward-scattering angles is required. These geometries are not confined to the synthetic experiment. MODIS views up to viewing zenith angles of approximately  $60$ – $67^\circ$  at swath edge. Combined with

260 solar zenith angles around  $60^\circ$ , common outside near-noon overpasses, this range includes the forward-scattering geometry identified here. A practical filter would flag retrievals with both high viewing zenith angle and low scattering angle, matching the specific combination shown above to produce the largest bias.

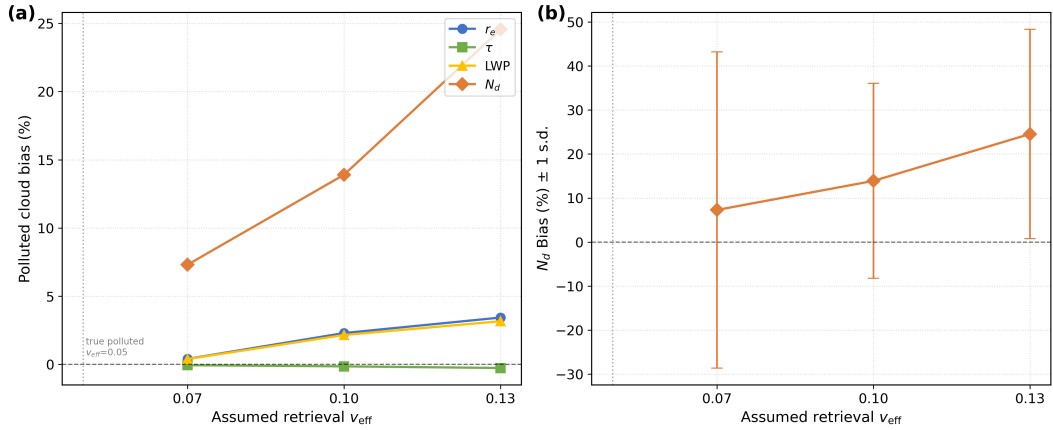
Table 2 summarises the bias statistics for each variable with their standard deviations, for both the central  $v_{\text{eff}} = 0.10$  baseline and the legacy  $v_{\text{eff}} = 0.13$  sensitivity comparison.

**Table 2.** Mean and standard deviation of percentage bias for polluted clouds, under the central  $v_{\text{eff}} = 0.10$  retrieval baseline and the legacy  $v_{\text{eff}} = 0.13$  sensitivity comparison. Statistics are computed on the full unfiltered polluted population across all viewing geometries, including outliers.

| Statistic                                               | $r_e$ Bias (%) | $\tau$ Bias (%) | LWP Bias (%) | $N_d$ Bias (%) |
|---------------------------------------------------------|----------------|-----------------|--------------|----------------|
| $v_{\text{eff}} = 0.10$ (central baseline)              |                |                 |              |                |
| Mean                                                    | 2.30           | -0.15           | 2.15         | 13.91          |
| Std. Dev.                                               | 2.81           | 1.43            | 3.24         | 22.14          |
| $v_{\text{eff}} = 0.13$ (legacy sensitivity comparison) |                |                 |              |                |
| Mean                                                    | 3.43           | -0.27           | 3.16         | 24.56          |
| Std. Dev.                                               | 3.13           | 1.46            | 3.75         | 23.75          |

### 265 3.3 Sensitivity to the Assumed Retrieval $v_{\text{eff}}$

The results above compare only two discrete points on the assumed- $v_{\text{eff}}$  axis: the current MODIS Collection 6 baseline (0.10) and the legacy pre-Collection 6 assumption (0.13). To characterise how the bias scales continuously as the assumed  $v_{\text{eff}}$  approaches the true polluted state ( $v_{\text{eff}} = 0.05$ ), the retrieval is repeated for an intermediate assumption of  $v_{\text{eff}} = 0.07$ , using the same synthetic cloud population and viewing geometry as the 0.10 and 0.13 cases. Figure 8a shows that all four bias metrics  
270 decrease monotonically as the assumed  $v_{\text{eff}}$  is reduced toward the true polluted value: the polluted  $N_d$  bias falls from 24.56% ( $v_{\text{eff}}=0.13$ ) to 13.91% ( $v_{\text{eff}}=0.10$ ) to 7.31% ( $v_{\text{eff}}=0.07$ ), while  $r_e$  and LWP bias fall from approximately 3.4% and 3.2% to 0.4% and 0.4%, respectively, over the same range. This confirms that the magnitude of the fixed- $v_{\text{eff}}$  bias scales with the discrepancy between the assumed and true polluted  $v_{\text{eff}}$ , rather than being an artefact specific to either baseline examined above. Notably, Figure 8b shows that the raw standard deviation of the  $N_d$  bias does not fall monotonically alongside the mean: it increases  
275 from  $\pm 22.14\%$  ( $v_{\text{eff}}=0.10$ ) to  $\pm 35.90\%$  ( $v_{\text{eff}}=0.07$ ) despite the smaller mean bias. This increase is driven almost entirely by the heavy tail of extreme-geometry cases discussed in Sect. 3.2, rather than by a broad loss of precision: when the 1st–99th percentile core of the distribution is considered, the spread in fact narrows monotonically, with standard deviations of 1.7%, 2.5%, and 3.4% for assumed  $v_{\text{eff}}$  of 0.07, 0.10, and 0.13, respectively. The amplification of the tail is consistent with  $k(v_{\text{eff}})$  becoming increasingly steep at low  $v_{\text{eff}}$  (Eq. (4)), such that the small number of severely biased  $r_e$  retrievals at extreme forward-scattering  
280 geometries translate into proportionally larger fluctuations in  $k$ , and hence  $N_d$ , as the assumed distribution narrows. Table 3 summarises the sweep.



**Figure 8.** Sensitivity of polluted-cloud retrieval bias to the assumed retrieval  $v_{\text{eff}}$  (0.07, 0.10, 0.13), relative to the fixed true polluted state  $v_{\text{eff}} = 0.05$  (dotted line). **(a)** Mean percentage bias in  $r_e$ ,  $\tau$ , LWP, and  $N_d$ . **(b)** Mean  $N_d$  bias  $\pm$  one standard deviation. The raw standard deviation increases toward low assumed  $v_{\text{eff}}$  due to amplification of the extreme-geometry tail, while the 1st–99th percentile core of the distribution narrows (see text). Statistics are computed on the full unfiltered polluted population across all viewing geometries, including outliers.

**Table 3.** Polluted-cloud bias statistics across the assumed retrieval  $v_{\text{eff}}$  sweep (true polluted  $v_{\text{eff}} = 0.05$  throughout). Statistics are computed on the full unfiltered polluted population across all viewing geometries, including outliers.

| Assumed $v_{\text{eff}}$ | $r_e$ Bias (%) | $\tau$ Bias (%) | LWP Bias (%) | $N_d$ Bias (%)        |
|--------------------------|----------------|-----------------|--------------|-----------------------|
| 0.07                     | 0.40           | -0.06           | 0.39         | 7.31 ( $\pm 35.90$ )  |
| 0.10                     | 2.30           | -0.15           | 2.15         | 13.91 ( $\pm 22.14$ ) |
| 0.13                     | 3.43           | -0.27           | 3.16         | 24.56 ( $\pm 23.75$ ) |

### 3.4 Sensitivity to Uncertainty in the True $v_{\text{eff}}-N_d$ Relationship

The analyses above, in common with the polluted-cloud population as a whole, assign every polluted case a fixed true  $v_{\text{eff}} = 0.05$ , regardless of its true  $N_d$ . This is a deliberate idealisation that isolates the retrieval bias from any assumption about how strongly droplet dispersion and concentration actually covary in nature, but it also means the reported bias magnitudes implicitly assume a perfect, noise-free correspondence between a cloud’s true  $N_d$  and its true  $v_{\text{eff}}$ . Empirical  $k-N_d$  relationships derived from aircraft data, including the fit underlying Lebsock and Witte (2023), are subject to considerable scatter about their central curve, and the strength of this coupling may be weaker than the fitted relationship alone would suggest. It is therefore necessary to investigate how robust the reported bias magnitudes are if the true  $v_{\text{eff}}-N_d$  coupling is much noisier than the deterministic case assumed above. The strength of this coupling may also depend on scale. This study’s paired ship-track construction implicitly assumes that factors other than the aerosol perturbation are held fixed between the clean and polluted cases. If this assumption does not hold in practice, for example if meteorology or background aerosol still varies somewhat

within a pair, the weak aggregate  $k$ - $N_d$  correlation seen in aircraft data may also be relevant at the individual ship-track scale, not only as a diluted, large-scale artifact. This motivates testing the weaker coupling case as a plausible scenario rather than only as a conservative bound.

To test this, the fixed true  $v_{\text{eff}} = 0.05$  assigned to each polluted case is replaced with a true  $k$  drawn from the Lebsock and Witte (2023) Combined Fit curve evaluated at that case’s true  $N_d$ , plus Gaussian scatter calibrated so that the resulting correlation between true  $k$  and true  $N_d$  falls to  $r \approx 0.10$ . This value is chosen as a bounding case rather than derived from a specific dataset; the qualitative conclusion below is not sensitive to the precise value chosen, only to the fact that the coupling is much weaker than deterministic. The scatter is clipped to a physically plausible range for cloud droplet spectra ( $v_{\text{eff}} \lesssim 0.3$ ) before recomputing the true  $N_d$  from the true  $r_e$  and  $\tau$ . This procedure holds the retrieved quantities ( $r_e$ ,  $\tau$ , and hence  $N_d$  retrieved) fixed at their values from the original forward-simulated retrieval, since these depend on the true top-of-atmosphere reflectances, which were generated at  $v_{\text{eff}} = 0.05$  and are not recomputed here. The result below should therefore be read as a bound on how sensitive the reported bias magnitudes are to the assumed strength of the  $v_{\text{eff}}$ - $N_d$  coupling, rather than as a fully independent re-simulation.

Table 4 compares the deterministic-truth bias statistics already reported in Table 3 against this weak-coupling case ( $r \approx 0.10$ , mean noisy true  $v_{\text{eff}} \approx 0.08$ , standard deviation  $\approx 0.08$ ). The mean polluted  $N_d$  bias falls substantially under the weaker coupling relative to the deterministic case at every retrieval baseline (13.91% to 3.80% at  $v_{\text{eff}} = 0.10$ ; 24.56% to 13.62% at  $v_{\text{eff}} = 0.13$ ), and the standard deviation increases by roughly 50–100% at every baseline. The Lebsock and Witte (2023) correction, which reduced the deterministic-truth bias to near zero (Sect. 4), instead systematically overcorrects by 5–11 percentage points under the weaker coupling. The apparent success in the deterministic case follows, in part, from testing the correction against a truth constructed from the same curve it corrects toward. This indicates that the polluted  $N_d$  bias magnitudes reported throughout this study, and the apparent effectiveness of the Lebsock and Witte (2023) correction, are best interpreted as an upper bound that assumes a tight coupling between droplet dispersion and concentration. If the true coupling is closer to the weaker correlation tested here, both the bias and the benefit of correcting for it are smaller and less certain than the deterministic case alone would suggest.

**Table 4.** Polluted-cloud  $N_d$  bias statistics under the deterministic true  $v_{\text{eff}} = 0.05$  assumption used throughout this study, compared against a weak-coupling case in which the true  $v_{\text{eff}}$ - $N_d$  relationship is calibrated to  $r \approx 0.10$  (see text). LW23 = after applying the Lebsock and Witte (2023) correction.

| Assumed $v_{\text{eff}}$ | Deterministic truth ( $v_{\text{eff}} = 0.05$ ) |                     | Weak coupling ( $r \approx 0.10$ ) |                        |
|--------------------------|-------------------------------------------------|---------------------|------------------------------------|------------------------|
|                          | $N_d$ Bias (%)                                  | LW23 $N_d$ Bias (%) | $N_d$ Bias (%)                     | LW23 $N_d$ Bias (%)    |
| 0.07                     | 7.31 ( $\pm 35.90$ )                            | –                   | -2.02 ( $\pm 47.44$ )              | -4.94 ( $\pm 41.29$ )  |
| 0.10                     | 13.91 ( $\pm 22.14$ )                           | 0.19                | 3.80 ( $\pm 35.31$ )               | -8.88 ( $\pm 28.96$ )  |
| 0.13                     | 24.56 ( $\pm 23.75$ )                           | -1.82               | 13.62 ( $\pm 38.52$ )              | -10.62 ( $\pm 28.35$ ) |

## 4 Discussion

The results of this study identify a systematic retrieval bias relevant to the ongoing challenge of constraining aerosol-cloud radiative forcing. While Forster et al. (2021) and Gryspeerdt et al. (2020) report improved convergence between observation-based and model-based ERF<sub>aci</sub> estimates, the uncertainty range remains large and retrieval assumptions such as fixed  $v_{\text{eff}}$  represent a tractable source of systematic error in observation-based assessments (Bellouin et al., 2020; Quaas et al., 2008). While the sensitivity of  $N_d$  retrievals to assumptions about the droplet size distribution has been identified as a potential source of uncertainty (Grosvenor et al., 2018; Brenguier et al., 2000), its impact on highly polluted regimes, such as ship tracks, has remained poorly constrained. The finding that the standard retrieval may overestimate  $N_d$  by approximately 14% in polluted regimes under the current MODIS Collection 6  $v_{\text{eff}} = 0.10$  assumption, rising to in excess of 24% under the legacy  $v_{\text{eff}} = 0.13$  assumption, suggests that, if the degree of spectral narrowing assumed here is representative, this bias could contribute to an overestimation of the microphysical sensitivity of clouds to aerosol perturbations in satellite derived datasets. Such a systematic bias would affect the interpretation of susceptibility metrics ( $S = d \ln N_d / d \ln \alpha$ ) used to constrain climate models (Quaas et al., 2009). This analysis suggests that in high-aerosol, narrow-variance regimes, this susceptibility could be artificially inflated by the fixed variance assumption used in standard operational algorithms (King et al., 1997). Correcting for this bias would be expected to flatten the derived susceptibility slope, potentially bringing observational constraints into closer alignment with more moderate estimates produced by models (Bellouin et al., 2020; Liu et al., 2008). As discussed previously, MODIS Collection 6 has transitioned from a  $v_{\text{eff}}$  value of 0.13 to 0.10, reducing the severity of this bias, but not eliminating it. As true droplet dispersion physically covaries with aerosol loading, applying any static  $v_{\text{eff}}$  assumption across a ship track would distort the derived cloud susceptibility.

If the coupling between droplet number and dispersion assumed in this study is representative of real ship tracks, an empirical correction of the kind proposed by Lebsock and Witte (2023) could in principle mitigate much of this bias. Lebsock and Witte derive a relationship between  $k$  and  $N_d$  from in-situ aircraft measurements,  $k(N_d) = (k_1 N^* + k_2 N_d) / (N_d + N^*)$ , which allows  $k$  to be estimated from the retrieved  $N_d$  itself rather than fixed a priori. Applying this correction (Combined Fit parameters,  $k_1 = 0.562$ ,  $k_2 = 0.974$ ,  $N^* = 48.5 \text{ cm}^{-3}$ ) to the synthetic retrievals here reduces the polluted  $N_d$  bias from 13.91% to 0.19% under the central  $v_{\text{eff}} = 0.10$  baseline, and from 24.56% to  $-1.82\%$  under the legacy  $v_{\text{eff}} = 0.13$  baseline, effectively removing the polluted bias identified above if the assumed  $k$ - $N_d$  coupling holds. However, the same correction applied uniformly to the clean background population does not improve, and in most cases worsens, the retrieval: the mean clean  $N_d$  bias shifts from  $-7.65\%$  to  $-11.26\%$  ( $v_{\text{eff}}=0.10$ ) and from  $+1.21\%$  to  $-12.83\%$  ( $v_{\text{eff}}=0.13$ ), with the overcorrection most pronounced for clean clouds with the broadest true distributions ( $v_{\text{eff}} = 0.15$ ). This indicates that, at least under the idealised conditions considered here, a blanket application of the Lebsock and Witte (2023) correction is not appropriate; its benefit is conditional on correctly identifying which pixels are genuinely narrow-distribution (i.e. polluted) cases, for instance via independent ship-track detection, rather than applying it indiscriminately across a scene. These results suggest a targeted fix: apply the aircraft-derived  $k$ - $N_d$  correction only within already-identified polluted regions, not scene-wide, and validate it against observational

350  $v_{\text{eff}}$  retrievals before wider adoption. Importantly, this estimate of the correction's benefit assumes a tight  $k$ - $N_d$  coupling, and Sect. 3.4 shows both the bias and the benefit shrink if the true coupling is weaker.

Furthermore, the precision of the retrieval of  $N_d$  in the polluted regime is severely degraded, with a standard deviation of  $\pm 22.14\%$  under the central  $v_{\text{eff}} = 0.10$  baseline ( $\pm 23.75\%$  under the legacy  $v_{\text{eff}} = 0.13$  comparison). This analysis indicates that this variance is largely driven by viewing geometry of  $r_e$  retrievals as illustrated by Figures 5 and 6. Due to  $N_d$  scaling  
355 with the inverse cube of  $r_e$  the geometry dependent variance in  $r_e$  is amplified. While the vast majority of the bias is bounded between 12–37%, at extreme geometries the retrieval algorithm can become unstable driving bias in  $N_d$  upwards of 1200%. These breakdowns are a consequence of the 1D plane-parallel cloud assumption breaking down. As demonstrated by Maddux et al. (2010), highly oblique VZA values significantly increase the satellite pixel footprint and optical pathlength, leading to a higher probability of partly cloudy fields and the observation of cloud sides. While standard observational methodologies  
360 typically filter out these extreme cases, discarding them does not solve the underlying variance discussed above.

Conversely, by demonstrating that LWP is retrieved robustly despite these variance assumptions, the presented results lend methodological confidence to the weak LWP responses documented by Toll et al. (2019) and Tippet et al. (2024). This indicates that the weak LWP adjustments observed in ship tracks are not a result of retrieval bias associated with the fixed  $v_{\text{eff}}$  assumption. Consequently, the competing mechanisms of darkening due to entrainment and brightening due to precipitation suppression,  
365 discussed by Wood (2007) and Chen et al. (2012), must be resolved physically or through the correction of other known biases. This finding supports the hypothesis that non-monotonic LWP responses are driven by environmental state variables rather than measurement error (Glassmeier et al., 2021). However, because the retrieval bias does not significantly perturb the derived LWP, it fails to resolve the instantaneous LWP decrease observed by Gryspeerdt et al. (2021), suggesting that this may be an unresolved microphysical adjustment. Notably, the finding that retrieval artefacts systematically overestimate  $N_d$  in polluted  
370 regimes aligns with the observational assessments of Gryspeerdt et al. (2023), which similarly identified that satellite-derived droplet concentrations are artificially inflated under high-aerosol conditions.

The feasibility of Marine Cloud Brightening (MCB) relies on identifying susceptible cloud decks where the injection of sea salt aerosols will yield a significant increase in albedo (Latham et al., 2012). These findings suggest that, to the extent that the assumed spectral narrowing occurs in practice, the background susceptibility of these clouds could be overstated. If the  
375 ship tracks used to calibrate susceptibility estimates (Christensen and Stephens, 2011) are subject to overestimation in their microphysical response due to spectral narrowing, the efficacy of interventions such as MCB may be correspondingly lower. The geometric sensitivity identified in this study also presents an operational obstacle. The apparent microphysical response of a perturbed cloud will artificially fluctuate simply due to the varying orbital position of the satellite across subsequent overpasses. An operational MCB program monitored by satellite would risk interpreting a retrieval artefact as a successful  
380 intervention.

## 4.1 Limitations

The results presented here are derived from a controlled, idealised sensitivity experiment isolating the effect of the fixed  $v_{\text{eff}}$  assumption. Several simplifications constrain the extent to which the specific bias magnitudes can be generalised to operational retrievals.

385 First, the clouds are modelled as single-layer, vertically homogeneous, plane-parallel slabs with a fixed geometric thickness of  $H = 500\text{m}$ . Real marine stratocumulus and ship-track clouds exhibit vertical structure in liquid water content and droplet size, as well as variable geometric thickness. Both are known to introduce additional retrieval biases beyond those considered here (Grosvenor et al., 2018). The sensitivity of the reported bias magnitudes to  $H$  has not been tested.

Second, this study does not reproduce the full operational MODIS retrieval chain, including cloud masking, multilayer and  
390 partly-cloudy pixel flags, and other quality-control logic. The experiment is deliberately restricted to the bias introduced by the  $v_{\text{eff}}$  assumption alone, isolated from these other sources of operational uncertainty.

Consequently, the reported bias magnitudes should be interpreted as a bound on the sensitivity to this one assumption under idealised conditions, rather than as a prediction of the net bias in any specific operational  $N_d$  product. In an operational setting this bias would act alongside other established uncertainty sources in bi-spectral retrievals, including three-dimensional  
395 radiative effects, sub-pixel heterogeneity, and instrument calibration (Grosvenor et al., 2018; Lai et al., 2019). A quantitative intercomparison of these contributions is beyond the scope of a single-assumption sensitivity test. The magnitude of the  $N_d$  biases reported here suggests that the fixed- $v_{\text{eff}}$  assumption warrants consideration alongside these established sources, particularly in ship-track studies where the clean-to-polluted contrast in  $v_{\text{eff}}$  is largest. While no published study isolates the fixed- $v_{\text{eff}}$  contribution to retrieval bias specifically for ship tracks, two recent aircraft-based evaluations provide broader context. Witte  
400 et al. (2018) find no significant bias in MODIS-retrieved  $r_e$  against in-situ measurements across nondrizzling-to-drizzling marine stratocumulus, consistent with the near-zero clean-regime  $r_e$  bias found here. Passer et al. (2025) report  $N_d$  discrepancies of  $\pm 50\%$  or more between MODIS and aircraft measurements in marine stratocumulus, indicating that the  $v_{\text{eff}}$ -driven  $N_d$  bias identified here (14–25%) is plausibly a substantial, but not dominant, contributor to the total observed retrieval uncertainty.

## 5 Conclusion

405 Satellite observational records provide the primary dataset for assessing aerosol-cloud interactions on a global scale. Yet, the reliability of these susceptibility estimates is inherently tethered to the mathematical assumptions embedded within satellite retrieval algorithms. A vulnerability of standard bi-spectral retrievals is their reliance on a fixed assumption for the droplet size distribution’s effective variance ( $v_{\text{eff}}$ ), set at  $v_{\text{eff}} = 0.1$  for MODIS Collection 6 and  $v_{\text{eff}} = 0.13$  for previous MODIS Collections. However, in polluted regimes such as ship tracks, aerosol injection not only reduces droplet size but also narrows  
410 the distribution. This narrowing of the distribution, set at  $v_{\text{eff}} = 0.05$  in this study, contradicts the fixed  $v_{\text{eff}}$  assumption. This study uses a synthetic retrieval experiment to quantify the systematic bias introduced by this assumption.

The results demonstrate that, within the constraints of the synthetic framework, assuming an inaccurately broad distribution for polluted clouds leads to a systematic overestimation of  $r_e$  by 2.30-3.43%, depending on the assumption (Figure 4). While the

retrieval LWP remains largely unaffected by the fixed assumption, due to the relative insensitivity of  $\tau$  to the distribution shape,  $N_d$  exhibits a strong positive bias, of approximately 23-24% when contrasted with the corresponding clean  $N_d$  (Figure 7). The substantial inflation of  $N_d$  is driven by its inverse dependence on the spectral width parameter ( $k$ ). As the droplet distribution narrows (increasing  $k$ ), the retrieval systematically overestimates droplet number due to  $k$  being held at an artificially low value. The retrieval of  $N_d$  is also found to be imprecise with a standard deviation of  $\pm 22.14 - 23.75\%$ , mainly driven by viewing geometry (Figures 5 and 6). These findings suggest that, if the aerosol-driven narrowing of the droplet size distribution assumed in this framework is representative of real ship tracks, the magnitude of the Twomey effect observed in ship track based observational studies may be overstated.

Under the same conditions, observational estimates of cloud susceptibility and the efficacy of marine cloud brightening could be exaggerated in current satellite products. Mitigation of these uncertainties will likely benefit from a combination of approaches. While global multi-angle and polarimetric datasets (e.g., POLDER, PACE) can constrain droplet size dispersion, their spatial resolutions are often insufficient to resolve fine-scale aerosol perturbations such as individual ship tracks. As high-resolution bi-spectral retrievals remain necessary for monitoring these localised interactions, future work should focus on computationally efficient correction strategies for these standard operational algorithms. This could include developing machine learning models capable of dynamically estimating  $v_{\text{eff}}$  from existing satellite radiances, or deploying adaptive algorithms that conditionally apply lower dispersion assumptions, or empirical  $k-N_d$  corrections of the kind proposed by Lebsock and Witte (2023), when specific features, such as ship tracks, are detected. The latter approach is supported by the finding here that the correction nearly eliminates the polluted-regime  $N_d$  bias when applied to identified ship-track pixels, but degrades the retrieval of the clean background when applied indiscriminately. This benefit itself depends on how tightly droplet dispersion and concentration are coupled in practice. Additional spectral channels may also help mitigate this bias. Liu and Teng (2026) show that shortwave-infrared channels beyond the standard bi-spectral pair can provide independent information about cloud microphysical structure not captured by current retrievals. Extending this principle to channels sensitive to the shape of the droplet size distribution could in principle allow  $v_{\text{eff}}$  to be constrained directly from the radiances. These findings indicate that the fixed effective-variance assumption is a relevant source of bias for ship-track-based estimates of aerosol-cloud interactions and climate intervention efficacy, and should be evaluated alongside other known retrieval uncertainties in future work.

*Code availability.* Radiative transfer code uses libRadtran (<https://doi.org/10.5194/gmd-9-1647-2016>, Emde et al. (2016)) and pyLRT (<https://doi.org/10.5281/zenodo.11626012>, Gryspeerd and Driver (2024)). Analysis code is available at <https://doi.org/10.5281/zenodo.20328197>.

*Data availability.* The generated data used in analysis can be found at <https://doi.org/10.5281/zenodo.19699485>

*Author contributions.* All authors contributed to designing the study. IB performed the analysis and wrote the paper. EG and AC assisted in the interpretation of the results and commented on the paper.

*Competing interests.* The contact author has declared that none of the authors has any competing interests.

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