

Reactions of Carbonyl Oxide with Aldehydes: Accurate Electronic Structure Methods, Kinetic Insights, and Atmospheric Implications

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Abstract: Carbonyl oxide (CH_2OO) is paramount in atmospheric oxidation chemistry, yet quantitative kinetics data for its bimolecular reactions are very limited and even unknown. Here we establish a computational framework to obtain quantitative kinetics from small to large reaction systems. For $\text{CH}_2\text{OO} + \text{HCHO}$, we develop electronic structure methods to reach CCSDTQ/CBS accuracy for its activation enthalpies at 0 K. For $\text{CH}_2\text{OO} + \text{aldehydes (RCHO; R = CH}_3\text{-C}_5\text{H}_{11}, \text{CH}_2\text{F, CHF}_2, \text{CF}_3)$, we introduce two strategies that recover CCSDTQ/CBS-quality activation enthalpies at 0 K. A dual-level strategy has been used to calculate their kinetics. The calculated rate constants show excellent agreement with available experimental data for $\text{CH}_2\text{OO} + \text{RCHO (R = CH}_3\text{-C}_3\text{H}_7)$, which validates the designed computational framework. We find that fluorination leads to exceptional rate enhancement, with reactions of CHF_2CHO and CF_3CHO exceeding $10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ over 200–320 K, approaching the collision limit. We also find that fluorination-driven reactivity enhancement originates predominantly from lower-level electronic effects than that of post-CCSD(T). Incorporation of the kinetics into a global chemical transport model uncovers previously unrecognized atmospheric impacts, with $\text{CH}_2\text{OO} + \text{HCHO}$ reducing nighttime CH_2OO and gas-phase sulfate concentrations by 25.3% in Antarctica and 12.2% over Canada, respectively. The present findings address a long-term challenge in how to obtain quantitative kinetics for large molecular systems, where post-CCSD(T) calculations are prohibitive and provide new insights into the chemical transformation of CH_2OO and fluorinated aldehydes in the atmosphere.

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22 **1 Introduction**

23 Aldehydes are a major class of oxygenated volatile organic compounds (OVOCs) that substantially influence atmospheric
24 oxidative capacity, secondary organic aerosol (SOA) formation, and air quality (Lary and Shallcross, 2000; Liu et al., 2022;
25 Zhao et al., 2024; Li et al., 2024; Mellouki et al., 2015; Bao et al., 2025; Zhang et al., 2012; Bari and Kindzierski, 2018;
26 Edwards et al., 2014; Yang et al., 2018). They originate from both direct emissions—including biomass and fossil-fuel
27 combustion, biogenic sources, and vehicle exhaust—and secondary production via VOC oxidation (Zhao et al., 2024; Knote
28 et al., 2014; Parrish et al., 2012; Chen et al., 2014; Luecken et al., 2012; Grosjean et al., 1983). Their atmospheric removal is
29 governed primarily by photolysis and OH reactions during daytime, whereas fluorinated aldehydes exhibit notably reduced
30 OH reactivity (Wenger, 2006; Jiménez et al., 2007; Atkinson and Pitts, 1978; Lily et al., 2021; Sellevåg et al., 2005; Scollard
31 et al., 1993; Thévenet et al., 2000; D'anna et al., 2001). NO₃ reactions constitute a nighttime sink but proceed extremely slow,
32 highlighting the need to identify alternative nocturnal loss pathways (Cabañas et al., 2001; Bossmeyer et al., 2006; Papagni et
33 al., 2000).

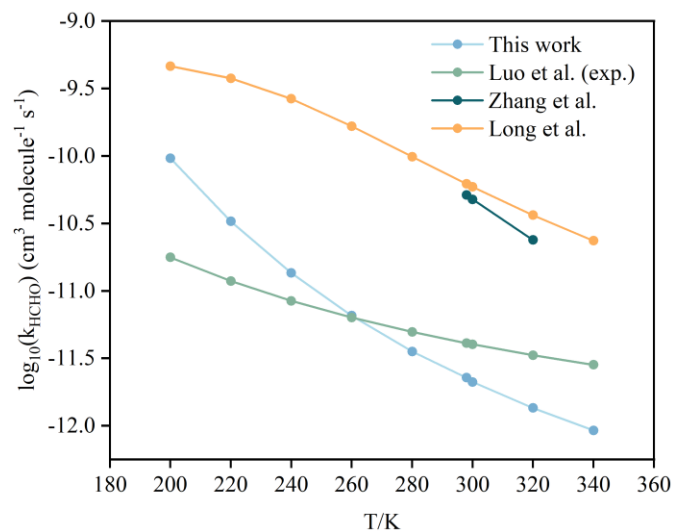
34 Stabilized Criegee intermediates (sCIs), key intermediate species of O₃-initiated alkene ozonolysis (Criegee, 1975;
35 Criegee and Wenner, 1949), play critical roles in atmospheric oxidation and SOA formation (Khan et al., 2018; Novelli et al.,
36 2014; Percival et al., 2013; Chhantyal-Pun et al., 2020) and react rapidly with acids (Cabezas and Endo, 2019; Chung et al.,
37 2019; Peltola et al., 2020; Foreman et al., 2016; Raghunath et al., 2017), amides (Wei et al., 2022; Long et al., 2025), and SO₂
38 (Berndt et al., 2014; Boy et al., 2013; Manonmani et al., 2023; Kukui et al., 2021). Accurate kinetics for their bimolecular
39 reactions are therefore essential for constraining their atmospheric fate.

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45 **Table 1.** Rate constants of CH₂OO + HCHO by previous investigation at different temperatures and pressures.

Reaction	P (Torr)	T (K)	$k(T)$ (cm ³ molecule ⁻¹ s ⁻¹)	Ref.
Exp.	56	296	$(4.11 \pm 0.25) \times 10^{-12}$	(Luo et al., 2023)
	78	275	$(4.84 \pm 0.41) \times 10^{-12}$	(Enders et al., 2024)
		295	$(3.50 \pm 0.35) \times 10^{-12}$	
Theory	10	213	3.28×10^{-9}	(Zhang et al., 2023)
	202	230	1.29×10^{-9}	
	406	259	3.52×10^{-10}	
	760	296	5.51×10^{-10}	
		295	5.71×10^{-10}	(Long et al., 2021)
		296	6.52×10^{-11}	
	275	1.11×10^{-10}		
	295	6.68×10^{-11}	This work	
	296	3.01×10^{-11}		
	275	5.62×10^{-11}		
	295	3.10×10^{-11}		

46 Despite numerous studies on CH₂OO + aldehydes, important gaps remain (Table 1 and Figure 1) for CH₂OO + HCHO,
 47 theoretical and experimental rate constants differ by an order of magnitude (Luo et al., 2023; Enders et al., 2024; Long et al.,
 48 2021; Zhang et al., 2023); prior work on CH₂OO + CH₃CHO/C₂H₅CHO/C₃H₇CHO (Tables 2-3) relied primarily on CCSD(T)
 49 despite evidence that higher-level excitations are required (Taatjes et al., 2012; Elsamra et al., 2016; Stone et al., 2014; Berndt
 50 et al., 2015; Jiang et al., 2024; Kaipara and Rajakumar, 2018; Liu et al., 2020; Liu et al., 2023; Cornwell et al., 2023; Debnath
 51 and Rajakumar, 2024); and key effects such as anharmonicity, torsional anharmonicity, and recrossing were generally neglected
 52 (Luo et al., 2023; Enders et al., 2024; Kaipara and Rajakumar, 2018; Debnath and Rajakumar, 2024; Jalan et al., 2013).



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Figure 1. A comparison of reported rate constants for the $\text{CH}_2\text{OO} + \text{HCHO}$ reaction from previous studies at different temperatures and high-pressure limit.

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Table 2. Rate constants of $\text{CH}_2\text{OO} + \text{CH}_3\text{CHO}$ by previous investigation at different temperatures and pressures.

Reaction	P (Torr)	T (K)	$k(T)$ ($\text{cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$)	Ref.
Exp.	4	293	$(9.50 \pm 0.70) \times 10^{-13}$	(Taatjes et al., 2012)
	25	298	$(1.20 \pm 0.20) \times 10^{-12}$	(Elsamra et al., 2016)
		340	$(8.00 \pm 1.10) \times 10^{-13}$	
	4	298	$(1.10 \pm 0.10) \times 10^{-12}$	
	50		$(1.30 \pm 0.20) \times 10^{-12}$	
	25	295	$(1.48 \pm 0.04) \times 10^{-12}$	(Stone et al., 2014)
	760	297	$(1.70 \pm 0.50) \times 10^{-12}$	(Berndt et al., 2015)
	78	275	$(2.37 \pm 0.21) \times 10^{-12}$	(Enders et al., 2024)
		295	$(1.61 \pm 0.14) \times 10^{-12}$	
	50	280	$(2.57 \pm 0.46) \times 10^{-12}$	(Jiang et al., 2024)
		298	$(2.13 \pm 0.38) \times 10^{-12}$	
	5.5	298	$(1.73 \pm 0.32) \times 10^{-12}$	
	10		$(2.08 \pm 0.38) \times 10^{-12}$	
	30		$(2.10 \pm 0.38) \times 10^{-12}$	
	50		$(2.13 \pm 0.38) \times 10^{-12}$	
100	$(2.16 \pm 0.38) \times 10^{-12}$			
80	275	$(10.20 \pm 0.80) \times 10^{-13}$	(Cornwell et al., 2023)	

		295	$(8.00 \pm 0.70) \times 10^{-13}$	
Theory	760	275	4.63×10^{-12}	This work
		280	4.02×10^{-12}	
		293	2.83×10^{-12}	
		295	2.69×10^{-12}	
		297	2.56×10^{-12}	
		298	2.50×10^{-12}	

58 Moreover, no kinetic data exist for reactions with larger or fluorinated aldehydes, including pentanal, hexanal, CH₂FCHO,
59 CHF₂CHO, and CF₃CHO. To address these gaps, atmospheric models have effectively utilized rate constants derived from
60 empirical structure–reactivity relationships (SRRs)—such as those proposed by Jenkin et al. (Jenkin et al., 2018)—which
61 provide a practical and robust framework for large-scale modeling. Given the inherent complexity of computing atmospheric
62 kinetics, these empirical methods remain a primary tool for estimation.

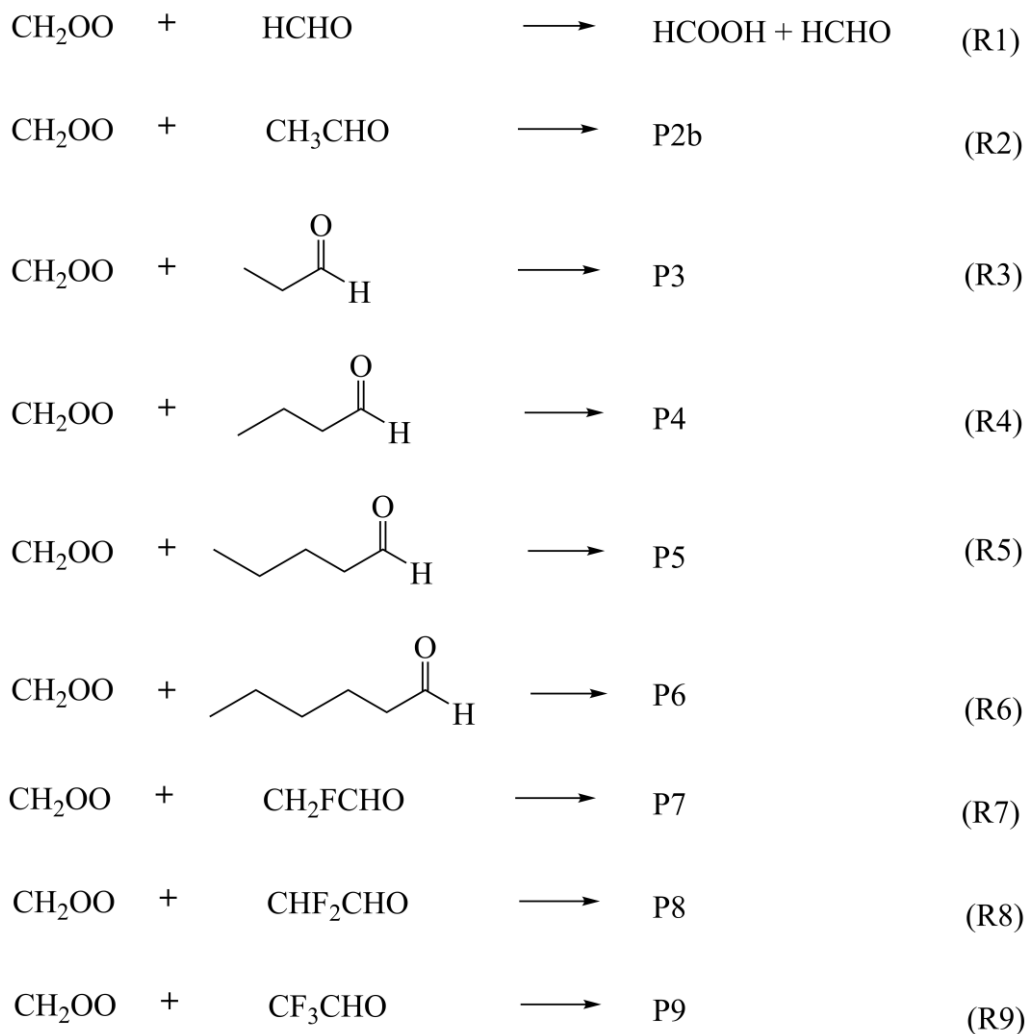
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64 **Table 3.** Rate constants of CH₂OO + RCHO (R = C₂H₅/C₃H₇) by previous investigation at different temperatures and pressures.

Reaction		P	T	$k(T)$ (cm ³ molecule ⁻¹ s ⁻¹)	Ref.	
C ₂ H ₅ CHO	Exp.	50 Torr	283 K	$(3.55 \pm 0.50) \times 10^{-12}$	(Liu et al., 2020)	
			298 K	$(3.12 \pm 0.44) \times 10^{-12}$		
		5 Torr	298 K	$(2.39 \pm 0.22) \times 10^{-12}$		
		5.2 Torr		$(2.52 \pm 0.24) \times 10^{-12}$		
		10 Torr		$(3.07 \pm 0.20) \times 10^{-12}$		
		25 Torr		$(2.12 \pm 0.19) \times 10^{-12}$		
		75 Torr		$(3.30 \pm 0.20) \times 10^{-12}$		
		100 Torr		$(3.08 \pm 0.19) \times 10^{-12}$		
		150 Torr		$(3.18 \pm 0.19) \times 10^{-12}$		
		200 Torr		$(3.19 \pm 0.21) \times 10^{-12}$		
	78 Torr	275 K	$(4.35 \pm 0.38) \times 10^{-12}$	(Enders et al., 2024)		
		295 K	$(3.29 \pm 0.29) \times 10^{-12}$			
	Theory	HPL		283 K	2.29×10^{-12}	(Kaipara and Rajakumar, 2018)
				298 K	1.51×10^{-12}	
				275 K	2.92×10^{-12}	
				295 K	1.63×10^{-12}	
283 K				4.49×10^{-12}	This work	
298 K				3.11×10^{-12}		

			275 K	5.57×10^{-12}	
			295 K	3.33×10^{-12}	
C ₃ H ₇ CHO	Exp.	50 Torr	253 K	$(4.20 \pm 0.10) \times 10^{-12}$	(Debnath and Rajakumar, 2024)
			268 K	$(3.61 \pm 0.10) \times 10^{-12}$	
			283 K	$(2.99 \pm 0.22) \times 10^{-12}$	
			298 K	$(2.63 \pm 0.14) \times 10^{-12}$	
	Theory	HPL	253 K	8.83×10^{-12}	This work
			268 K	5.30×10^{-12}	
			283 K	3.38×10^{-12}	
			298 K	2.27×10^{-12}	

65 Here, we investigate CH₂OO reactions with nine aldehydes (RCHO; R = H, CH₃, C₂H₅, C₃H₇, C₄H₉, C₅H₁₁, CH₂F, CHF₂,
66 CF₃) to obtain quantitative rate constants and to establish a general high-accuracy computational protocol applicable from
67 small benchmark systems to large atmospheric molecules. For the prototypical CH₂OO + HCHO reaction, we develop the
68 GMM(Q).L4 composite scheme that approaches full-CI accuracy, and for the broader reaction suite we devise a scalable
69 strategy capable of delivering near-full-CI activation energies. Dual-level strategy calculations accounting for all major
70 anharmonic and dynamical effects yield benchmark-quality rate constants, which are subsequently implemented in GEOS-
71 Chem to quantify their atmospheric impacts. This work provides a broadly extensible computational framework and
72 significantly advances the understanding of CH₂OO-aldehyde chemistry.



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74 **Scheme 1.** Reactions of CH₂OO with aldehydes75 **2 Computational methods and strategies**76 **2.1 Electronic structure best estimates for the CH₂OO + HCHO reaction**

77 Accurate electronic-structure data are essential for quantitative kinetics. All geometries and harmonic frequencies were
 78 optimized at the CCSD(T)-F12a/cc-pVTZ-F12 level (Adler et al., 2007; Knizia et al., 2009; Bischoff et al., 2009). To approach
 79 the full-CI limit for single-point energies, we developed a composite protocol, GMMQ.L4, which effectively reproduces

80 CCSDTQ/CBS quality:

$$81 E_{\text{GMMQ.L4}} = E_{\text{MW2-F12.L}} + \Delta E_{\text{T-(T)}} + \Delta E_{\text{(Q)-T}} + \Delta E_{\text{Q-(Q)}} \quad (1)$$

82 Here, $E_{\text{MW2-F12.L}}$ is obtained from the previously validated MW2-F12.L scheme which detailed in Table S7 (Long et al.,
83 2021). $\Delta E_{\text{T-(T)}}$ (CCSDT-CCSD(T)) and $\Delta E_{\text{(Q)-T}}$ (CCSDT(Q)-CCSDT) are extrapolated to the CBS limit (cc-pVDZ \rightarrow cc-
84 pVTZ and cc-pVDZ \rightarrow VTZ(d)) using

$$85 \Delta E_{\text{L}} = \Delta E_{\text{CBS}} + \frac{A}{L^3} \quad (2)$$

86 with $L=2$ for cc-pVDZ and 3 for cc-pVTZ and VTZ(d).

87 The final correction, $\Delta E_{\text{Q-(Q)}}$, is evaluated at the CCSDTQ-CCSDT(Q) level using the VDZ(NP) basis set. VTZ(d) employs
88 H(s) and heavy-atom(sp), while VDZ(NP) uses H(s) and heavy-atom(sp) functions (Chan and Radom, 2015).

89 Coupled-cluster theory converges systematically toward Full configuration interaction (Full-CI), but the steep scaling
90 necessitates truncation. Previous studies have established rapid basis-set convergence for both CCSDT(Q)-CCSDT and
91 CCSDT-CCSD(T) (Long et al., 2021; Long et al., 2019; Xia et al., 2025). Consistently, the CCSDTQ-CCSDT(Q) contribution
92 in our system is only 0.096 kcal mol⁻¹, indicating that excitations beyond quadruples contribute <0.10 kcal mol⁻¹ in Table S1.
93 Thus, GMMQ.L4//CCSD(T)-F12a/cc-pVTZ-F12 serves as the benchmark level in our dual-level kinetics framework.

94 We further compared GMMQ.L4 with the W3X-L composite method (Chan and Radom, 2015) for reaction R1. Although
95 both protocols include identical post-CCSD(T) contributions, GMMQ.L4 employs the MW2-F12.L component, whereas
96 W3X-L is based on W2X. Detailed comparisons are provided in Tables S1, S7, and S8. The observed deviation of 0.24 kcal
97 mol⁻¹ indicates that W3X-L does not achieve quantitatively reliable barrier heights for this system. Our analysis shows that
98 this discrepancy primarily originates from the difference between MW2-F12.L and W2X. Specifically, MW2-F12.L includes
99 HF energies, ΔCCSD and $\Delta(\text{T})$ correlation contributions, core-valence ($\Delta(\text{C+V})$) corrections, and scalar relativistic ($\Delta(\text{C+R})$)
100 effects, all evaluated with larger basis sets. In contrast, W2X comprises analogous HF, ΔCCSD , $\Delta(\text{T})$, and $\Delta(\text{C+R})$ terms, but
101 these are computed using smaller basis sets. The calculated results showed the difference of 0.24 kcal mol⁻¹ comes from the
102 $\Delta(\text{C+V})$ and $\Delta(\text{C+R})$ terms, which differ by 0.19 kcal/mol and 0.12 kcal/mol, respectively. Additionally, CCSD(T)-F12
103 convergence was verified by comparing W2X energies computed with cc-pVTZ-F12 and cc-pVDZ-F12 geometries; the

104 difference of only 0.04 kcal mol⁻¹ confirms near-CBS performance of CCSD(T)-F12 for structural and vibrational data (See
105 Table 4).

106 **Table 4.** Calculated enthalpies of activation at 0 K (ΔH_0^\ddagger in kcal/mol, relative to the bimolecular reactants) and unsigned
107 deviation (MUD) (in kcal/mol).

Methods	ΔH_0^\ddagger	
	TS1	UD
GMMQ.L4//CCSD(T)-F12a/cc-pVTZ-F12	-4.97	0.00
BE1//CCSD(T)-F12a/cc-pVTZ-F12	-4.97	0.00
BE2//CCSD(T)-F12a/cc-pVTZ-F12	-4.97	0.00
M11-L/MG3S	-5.16	0.19
W3X-L//CCSD(T)-F12a/cc-pVTZ-F12	-5.22	0.24
MW2-F12.L//CCSD(T)-F12a/cc-pVTZ-F12	-5.41	0.44
W2X//DF-CCSD(T)-F12a/jun-cc-pVDZ	-5.60	0.63
W2X//CCSD(T)-F12a/cc-pVTZ-F12	-5.62	0.64
W2X//CCSD(T)-F12a/cc-pVDZ-F12	-5.66	0.68
W2X//DF-CCSD(T)-F12b/VDZ(d)	-5.66	0.68
W2X//DF-CCSD(T)-F12a/cc-pVDZ	-5.72	0.74
W2X//DF-CCSD(T)-F12b/VDZ(NP)	-6.19	1.22

108 2.2 Electronic structure best estimates for R2-R9

109 **Geometrical optimization and frequency calculations.** Reliable optimized geometries and harmonic frequencies are
110 essential for obtaining quantitative 0 K activation enthalpies. For reaction R1, we verified that CCSD(T)-F12a/cc-pVDZ-F12
111 delivers results essentially identical to CCSD(T)-F12a/cc-pVTZ-F12, allowing us to employ the lower-cost cc-pVDZ-F12
112 basis for reaction R2. However, for larger CH₂OO + aldehyde systems, CCSD(T)-F12a/cc-pVDZ-F12 remains computationally
113 prohibitive. To overcome this limitation, we systematically benchmarked density-fitted F12 coupled-cluster methods (DF-
114 CCSD(T)-F12b) (Györfy and Werner, 2018) across a range of compact basis sets (Table 4). Remarkably, DF-CCSD(T)-
115 F12b/jun-cc-pVDZ (Parker et al., 2014) and DF-CCSD(T)-F12b/VDZ(d) exhibit exceptionally small mean unsigned
116 deviations of only 0.03 and 0.04 kcal mol⁻¹, respectively, relative to the best estimate for W2X reference (Table S2). This
117 identifies a new, computationally efficient F12 protocol capable of retaining sub-0.05 kcal mol⁻¹ accuracy for CH₂OO-

118 aldehyde reactions, representing a key methodological advance enabling routine treatment of larger Criegee intermediate–
119 carbonyl systems. Accordingly, we employed DF-CCSD(T)-F12b/jun-cc-pVDZ for R3–R5 and R7–R8, and DF-CCSD(T)-
120 F12b/VDZ(d) for R3–R6 to obtain geometries and vibrational frequencies with near-CBS accuracy at greatly reduced cost.

121 ***Single point energy calculations.*** To further reduce the cost of CCSDTQ/CBS-quality calculations, we developed a new
122 composite scheme, denoted BE1, which achieves near-GMMQ.L4 accuracy. The BE1 single-point energy is defined as

$$123 E_{\text{BE1}} = E_{\text{W2X}} + \Delta E_{(Q)-(T)} + \Delta E_{\text{SC1}} \quad (3)$$

124 where $\Delta E_{(Q)-(T)}$ is the CCSDT(Q) – CCSD(T) correction evaluated with the VDZ(NP) basis set for reactions R1–R8.

125 The term ΔE_{SC1} introduces a structure-specific correction and is given by

$$126 \Delta E_{\text{SC1}} = E_{\text{GMMQ.L4}}^{\text{TS1}} - E_{\text{W2X}}^{\text{TS1}} - [E_{\text{CCSDT(Q)/VDZ(NP)}}^{\text{TS1}} - E_{\text{CCSD(T)/VDZ(NP)}}^{\text{TS1}}] \quad (4)$$

127 This formulation anchors the composite energy to a single high-level reference transition state (TS1), ensuring the
128 transferability of the correction across the reaction series. The value of ΔE_{SC1} is 0.04 kcal/mol.

129 For comparison, we also employed our previously reported strategy, BE2 (Sun et al., 2024), which augments the W2X
130 energy with a constant post-CCSD(T) correction:

$$131 E_{\text{BE2}} = E_{\text{W2X}} + \Delta E_{\text{SC2}} \quad (5)$$

132 where ΔE_{SC2} is the GMMQ.L4 – W2X difference for TS1 (0.64 kcal mol⁻¹ in Table 4).

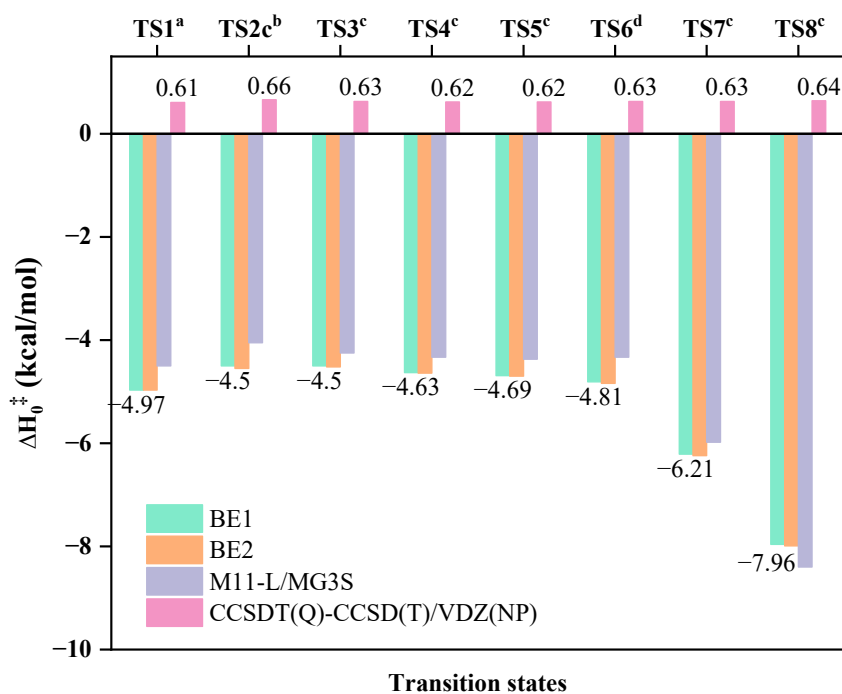
133 Both BE1 and BE2 offer computationally inexpensive routes to emulate CCSDTQ/CBS performance by incorporating
134 systematic, physically motivated corrections. In the present work, the BE1 protocol served as the high-level (HL) energy in
135 our dual-level kinetics strategy, with the underlying structures obtained from

- 136 • BE1//CCSD(T)-F12a/cc-pVDZ-F12 for R2,
- 137 • BE1//DF-CCSD(T)-F12b/jun-cc-pVDZ for R3–R5 and R7–R8, and
- 138 • BE1//DF-CCSD(T)-F12b/VDZ(d) for R6.

139 This composite strategy enables sub-kcal mol⁻¹ accuracy at a fraction of the cost of full GMMQ.L4 or CCSDTQ/CBS
140 calculations.

141 **2.3. Electronic structure density functional methods**

142 To enable efficient direct kinetics calculations for the full aldehyde series, we systematically evaluated a range of density
 143 functional methods against the BE1 benchmark. Among all tested functionals, M11-L (Peverati and Truhlar, 2012)/MG3S
 144 (Lynch et al., 2003) exhibits the best performance, yielding a remarkably small mean unsigned deviation (MUD) of 0.32 kcal
 145 mol⁻¹ cross reactions R1–R8 (Figure 2). This accuracy—well within sub-kcal mol⁻¹ agreement with the BE1 high-level
 146 reference—identifies M11-L/MG3S as a reliable and computationally economical low-level (LL) method for the dual-level
 147 kinetics framework. Accordingly, M11-L/MG3S was used for all direct kinetics calculations involving CH₂OO + aldehyde
 148 reactions. Standard vibrational scaling factors were applied as listed in Table S3.



149 **Figure 2.** Best estimate for reaction R1-R8 at different level.

150 ^aThe best estimate results by BE1//CCSD(T)-F12a/cc-pVTZ-F12 in the CH₂OO + HCHO reaction.

151 ^bThe best estimate results by BE1//CCSD(T)-F12a/cc-pVDZ-F12 in the CH₂OO + CH₃CHO reaction.

152 ^cThe best estimate results by BE1//DF-CCSD(T)-F12b/jun-cc-pVDZ in the CH₂OO + XCHO
 153 (X=C₂H₅/C₃H₇/C₄H₉/CH₂F/CHF₂) reaction.

154 ^dThe best estimate results by BE1//DF-CCSD(T)-F12b/VDZ(d) in the CH₂OO + C₅H₁₁CHO reaction.

155 Previous studies have suggested that standard scaling factors may be unsuitable for certain transition states, we explicitly

157 investigated the impact of anharmonicity. Using the method described by Long et al. (Long et al., 2023), we calculated specific
158 scaling factors (See Tables S4 and S5). However, we found that anharmonicity corrections to the zero-point energy (ZPE) were
159 negligible. Consequently, standard scaling factors are employed throughout this work. Full methodological details are provided
160 in the Supporting Information.

161 2.4. Kinetics Methods

162 **High-pressure limited rate constants for R2-R6.** Dual-level strategy (Long et al., 2019; Sun et al., 2024; Long et al.,
163 2016) was employed, in which high-level (HL) conventional transition state theory (TST) provides the baseline rate constants,
164 whereas canonical variational transition state theory with small-curvature tunneling (CVT/SCT) at the low-level (LL) supplies
165 kinetic corrections. The high-pressure-limit rate constants were obtained according to eq 5:

$$166 \quad k = k_{\text{HL}}^{\text{TST}}(T) \kappa_{\text{LL}}(T) \Gamma_{\text{LL}}(T) F_{\text{fwd}}^{\text{MS-T,LL}}(T) \quad (5)$$

167 where $k_{\text{HL}}^{\text{TST}}$ is the rate constants calculated at HL. $\kappa_{\text{LL}}(T)$ and $\Gamma_{\text{LL}}(T)$ is tunneling and recrossing transmission coefficients
168 calculated at the LL level. $F_{\text{fwd}}^{\text{MS-T,LL}}(T)$ is referred to multi-structural anharmonic factor calculated by eqn (6) at the M11-
169 L/MG3S level

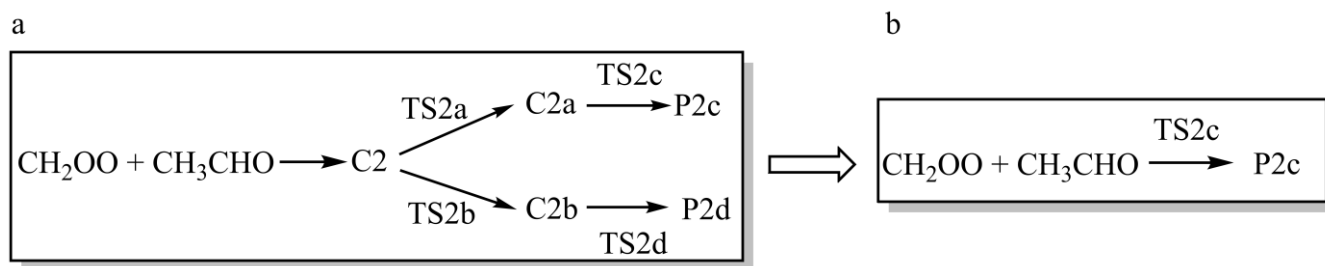
$$170 \quad F_{\text{fwd}}^{\text{MS-T,LL}} = \frac{F_{\text{TS}}^{\text{MS-T}}}{F_{\text{R}}^{\text{MS-T}}} \quad (6)$$

171 **High-pressure limited rate constants for R1 and R7-R9.** The rate constants of R1 and R7-R9 were calculated by
172 simultaneously considering both the loose transition state between reactants and the van der Waals complex, and the tight TS
173 between reactants and products. The rate constant for the loose TS (k_{loose}) was calculated using variable-reaction-coordinate
174 variational transition-state theory (VRC-VTST) (Georgievskii and Klippenstein, 2003; Zheng et al., 2008; Bao et al., 2016b)
175 with 500 configurations for Monte Carlo sampling. A single-faceted dividing surface was constructed with two pivot points,
176 following procedures validated in previous work (Long et al., 2021). One pivot point was placed along a vector at a distance
177 d from the center of mass (COM) of CH₂OO, oriented perpendicular to the CH₂OO plane, while the other was placed similarly
178 with respect to CH₂F/CHF₂/CF₃CHO. The pivot distance was set to $d=0.05$ Å. The reaction coordinate s was defined as the
179 separation between the two pivot points, ranging from 3.5 to 10 Å for R7, 3.9 to 10 Å for R8, 4.4 to 10 Å for R9 with increments

180 of 0.1 Å. The rate constant for the tight TS (k_{tight}) was calculated by using dual-level strategy presented above. The overall
181 rate constant was then obtained using the steady-state approximation (Garrett and Truhlar, 1982; Zhang et al., 2020; Long et
182 al., 2024) in equation (7).

$$183 \quad k = \frac{k_{\text{loose}}k_{\text{tight}}}{k_{\text{loose}}+k_{\text{tight}}} \quad (7)$$

184 **Pressure-dependent rate constant.** Master equation method with Rice–Ramsperger–Kassel–Marcus theory (ME/RRKM)
185 (Kenneth A. Holbrook, 1996; Fernández-Ramos et al., 2006; Georgievskii et al., 2013; Klippenstein, 2003) was used to
186 calculate pressure dependence of rate constants for the reactions of CH₂OO with HCHO and CH₃CHO. The calculation utilized
187 parameters from W3X-L//CCSD(T)-F12a/cc-pVTZ-F12 for reaction R1 and W2X//DF-CCSD(T)-F12b/jun-cc-pVDZ for
188 reaction R2. Both reactions were modeled with N₂ as the bath gas, employing Lennard-Jones parameters from Table S6 and
189 an average energy transfer parameter of $\langle\Delta E\rangle_{\text{down}} = 200 \text{ cm}^{-1}$. Within this framework, the pressure effect was approximated
190 as the quotient of the high-pressure limit and a pressure ratio. This ratio is defined as the value at 7.5×10^3 Torr relative to its
191 value at different pressures. We further inspect the simplification of reaction R2 in Scheme 2. The kinetic results for Schemes
192 2a and 2b demonstrate remarkable robustness, with the simplification introducing no statistically significant perturbations to
193 the calculated rate constants.



196 **Scheme 2.** The reaction mechanism for the CH₂OO + CH₃CHO reaction.

197 2.5. Atmospheric modeling

198 We performed two atmospheric simulations included reaction R1 and R2 to investigate the significance of these reactions
199 by observing the change of concentration globally in GEOS-Chem. This included: (1) a “base” model using default setting (2)
a “update1” model adding a new sink of HCHO in the base model, (3) a “update2” model adding a new sink of CH₃CHO in

200 the base model. These models include the meteorological data observations assimilated from the NASA Modern-Era
201 Retrospective Analysis for Research and Applications (MERRA-2) (Gelaro et al., 2017) and Emissions data from the default
202 Harmonized Emission Component (HEMCO) (Lin et al., 2021). For anthropogenic emissions, we used the Community
203 Emissions Data System (CEDS) (Hoesly et al., 2018). For biogenic emissions, we used offline VOC emissions computed from
204 the Model of Emissions of Gases and Aerosols from Nature (MEGAN) (Guenther et al., 2012). The simulation was carried out
205 with $2^\circ \times 2.5^\circ$ horizontal resolution at 47 vertical layers. The annual changes displayed are obtained from simulations that
206 employed meteorological data from February 1, 2018, to January 31, 2019, following a six-month model spin-up.

207 **2.6. Software.**

208 Density functional calculations were performed by using the Gaussian 16 (Frisch et al., 2016). The coupled cluster
209 calculations were performed by using the Molpro 2019 (Werner, et al., 2019) and MRCC codes (Kállay et al., 2020). Multi-
210 structural anharmonic calculations were performed in MSTor codes (Zheng et al., 2012). Rate constants were calculated using
211 the Polyrate 2017-C (Zheng et al., 2017b), Gaussrate 2017-B codes (Zheng et al., 2017a), and KiStHLP 2021 (Canneaux et al.,
212 2014). The master equation calculations were performed by utilizing the TUMME program (Zhang et al., 2022). Atmospheric
213 modeling was performed by using GEOS-Chem 14.4.2 (Bey et al., 2001, <http://www.geos-chem.org>, last access: 4 November
214 2025).

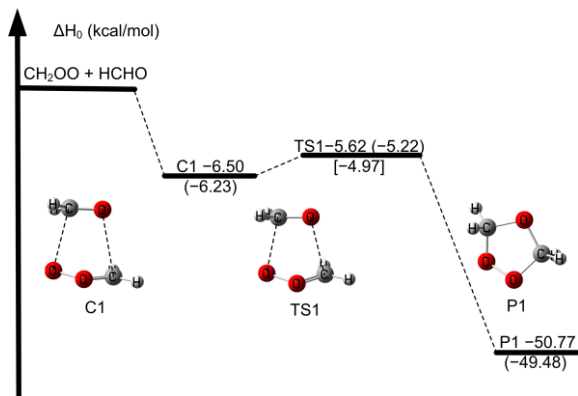
215 **3. RESULTS AND DISCUSSION**

216 The enthalpy of activation at 0 K (ΔH_0^\ddagger) is referred to the relative energies with zero-point energy between transition states
217 and reactants.

218 **3.1 The electronic structure of CH₂OO + HCHO**

219 The reaction mechanism examined here is consistent with that established in earlier studies (Luo et al., 2023; Long et al.,
220 2021; Jalan et al., 2013; Wang et al., 2022). The relative enthalpy profile for the CH₂OO + HCHO reaction is depicted in Figure
221 3, and the key data are summarized in Table 4. Notably, the activation enthalpy at 0 K obtained at the GMMQ.L4//CCSD(T)-

222 F12a/cc-pVTZ-F12 level (-4.97 kcal mol $^{-1}$) differs from that predicted by W3X-L//CCSD(T)-F12a/cc-pVTZ-F12 (-5.21 kcal
 223 mol $^{-1}$ in Table 4) and deviates even more substantially from the RCCSD(T)-F12a/VTZ-F12//B3LYP/MG3S value (-6.30 kcal
 224 mol $^{-1}$) (Jalan et al., 2013). These differences demonstrate the strong sensitivity of ΔH_0^\ddagger to the underlying electronic-structure
 225 treatment, thereby directly influencing predicted rate constants.



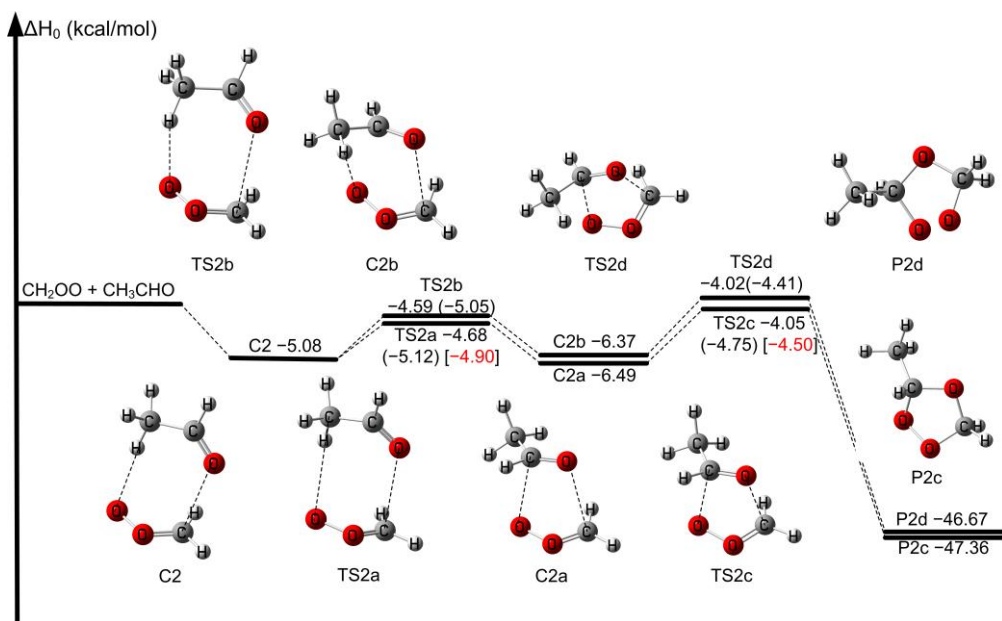
226
 227 **Figure 3.** The relative enthalpies at 0 K for the reaction of CH₂OO + HCHO. Values are given for all species as calculated by
 228 W2X//CCSD(T)-F12a/cc-pVTZ-F12, and in parentheses and bracket, values are given for the transition state TS1 as calculated
 229 by W3X-L//CCSD(T)-F12a/cc-pVTZ-F12 and GMMQ.L4//CCSD(T)-F12a/cc-pVTZ-F12.

230 Previous studies have shown that post-CCSD(T) correlation is essential for quantitative barriers in Criegee chemistry (Long
 231 et al., 2021; Long et al., 2016; Xia et al., 2022). For TS1, the unsigned deviation between GMMQ.L4 and MW2-F12.L is 0.40
 232 kcal mol $^{-1}$ —slightly different with the ~ 0.50 kcal mol $^{-1}$ benchmark established for post-CCSD(T) effects (Long et al., 2021)—
 233 reaffirming the need for high-level correlation to achieve quantitative accuracy. We further find that the post-CCSD(T)
 234 contribution through CCSDT(Q), quantified by the W3X-L – W2X difference, is 0.44 kcal mol $^{-1}$, in excellent agreement with
 235 the 0.40 kcal mol $^{-1}$ value. This concordance highlights the robustness of W3X-L in capturing post-CCSD(T) contributions
 236 (Table 4). The remaining 0.24 kcal mol $^{-1}$ discrepancy between GMMQ.L4 and W3X-L primarily reflects differences between
 237 the MW2-F12.L and W2X components of TS1 (Tables S7 and S8). The 0.21 kcal mol $^{-1}$ deviation between MW2-F12.L and
 238 W2X further illustrates that larger basis sets are required for fully quantitative predictions.

239 This present work provides a rigorously benchmarked assessment of ΔH_0^\ddagger for the CH₂OO + HCHO reaction, explicitly
 240 quantifying post-CCSD(T) contributions and revealing their decisive role in achieving sub-kcal mol $^{-1}$ accuracy. The systematic
 241 comparison among GMMQ.L4, MW2-F12.L, and W3X-L underscores the reliability of our calculated results.

242 3.2. The electronic structure of CH₂OO + CH₃CHO

243 We aim to demonstrate the feasibility of simplifying the reaction mechanism of larger aldehydes with CH₂OO in Scheme
 244 2. A partial reaction mechanism CH₂OO + CH₃CHO has been reported in our previous work (Wang et al., 2022). We first
 245 consider the seven-membered ring pre-reaction complex C2 formation in Figure 4, which is consistent with our previous results
 246 (Wang et al., 2022). However, due to two distinct orientations of the methyl group in CH₃CHO toward CH₂OO, there are two
 247 rotation transition states TS2a and TS2b connecting C2 to the five-membered ring complexes C2a and C2b, respectively.
 248 Therefore, the process is only the transformation of complex in the reaction processes. Then, C2a and C2b undergo the
 249 corresponding transition state TS2c and TS2d responsible for the formation of P2a and P2b. The mechanism was depicted in
 250 Scheme 2a. However, the enthalpies of activation at 0 K for TS2a and TS2b are lower than those of TS2c and TS2d by 0.64
 251 kcal/mol and 0.37 kcal/mol at W3X-L//CCSD(T)-F12a/cc-pVDZ-F12 in Figure 4, respectively. Therefore, TS2a and TS2b
 252 could be neglected from energetic point of view. We will also discuss it from the kinetics point of view.



253
 254 **Figure 4.** The relative enthalpies at 0 K for the reaction of CH₂OO + CH₃CHO. Values are given for all species as calculated
 255 by M11-L/MG3S, and in parentheses and bracket, values are given for the transition states as calculated by W3X-L//CCSD(T)-
 256 F12a/cc-pVDZ-F12 and BE1//CCSD(T)-F12a/cc-pVDZ-F12.

257 The five-membered ring complexes C2a and C2b can interconvert via TS2_{ISO} with C=O bond rotation, which lies 2.51

258 kcal mol⁻¹ above C2a at the M11-L/MG3S level (Figure S4), similar to the reaction between CH₂OO and FCHO (Xia et al.,
259 2024). For aldehydes with longer chains, the corresponding isomerization transition states of the five-membered ring
260 complexes (Figures S5–S6) exhibit similarly low barriers, indicating facile interconversion, which also verified from kinetics
261 perspective. Consequently, the complex mechanism can be effectively reduced to the straightforward reaction pathway b
262 depicted in Scheme 2. Accordingly, the mechanism for CH₂OO with larger aldehydes was simplified to consider only the
263 lowest-energy pathway corrected by torsional anharmonicity in kinetics calculations.

264 The ΔH_0^\ddagger for TS2c is -4.50 kcal mol⁻¹ at the BE1//CCSD(T)-F12a/cc-pVDZ-F12 level (See Table S9), which is 0.8 kcal
265 mol⁻¹ higher than the result reported by Jalan et al. at the RCCSD(T)-F12a/VTZ-F12//B3LYP/MG3S level and 0.19 kcal mol⁻¹
266 higher than that of Wang et al at the WMS//M11-L/MG3S level (Wang et al., 2022; Jalan et al., 2013). BE1 and BE2 for TS2c
267 agree well with each other in Figure 2 and Table S9, not only demonstrating the reliability of the computational protocol, but
268 also capturing the essential physical origin underlying the quantitative description of ΔH_0^\ddagger . The M11-L/MG3S has been chosen
269 for direct dynamics calculations due to the MUD of 0.81 kcal mol⁻¹ in Table S9.

270 The validity of the DF-CCSD(T)-F12/jun-cc-pVDZ and DF-CCSD(T)-F12b/VDZ(d) methods was also confirmed for
271 reaction R2. As shown in Table S9, these methods yielded mean unsigned deviations (MUD) of 0.05 and 0.02 kcal mol⁻¹,
272 respectively, relative to the CCSD(T)-F12a/cc-pVDZ-F12 benchmark.

273 3.3. Electronic structure of CH₂OO + RCHO (R = C₂H₅/C₃H₇/C₄H₉/C₅H₁₁)

274 The complexity of reactions R3–R6 increases with reactant system size owing to the presence of multiple conformers of
275 both reactants and transition states (Table S10). Conformers for each reactant and transition state were obtained by rotating the
276 dihedral angles listed in Table S10. Specifically, two conformers were identified for C₂H₅CHO, four for C₃H₇CHO, twelve for
277 C₄H₉CHO, and thirty-five for C₅H₁₁CHO, arising from C–C bond rotations. In contrast, conformational diversity is even more
278 pronounced for the transition states, with three conformers for TS3, eighteen for TS4, twenty-four for TS5, and seventy-nine
279 for TS6, primarily due to internal C=O and C–C bond rotations.

280 As the carbon chain prolongs, the change in ΔH_0^\ddagger for R1-R6 is not obvious, but it presents a trend. We find a slight
281 decrease in ΔH_0^\ddagger with the elongation of carbon chain for R2-R6 with the exception of R1. The ΔH_0^\ddagger calculated by best

282 estimate are -4.50 , -4.50 , -4.63 , -4.70 , and -4.80 kcal mol⁻¹ for R2-R6 (See Figure 2 and Table S11), which are about 3 kcal
283 mol⁻¹ below the reaction of the corresponding reactants with HO₂ (Gao et al., 2024; Long et al., 2022; Ding and Long, 2022).
284 Moreover, the influence of carbon chain length on enthalpy of activation for R2-R6 is analogue to the reaction of HO₂ and
285 aldehydes (Gao et al., 2024). Also, BE1 and BE2 for TS2c–TS6 (Figure 2 and Table S11) exhibit excellent mutual consistency.
286 This behavior can be attributed to the nearly invariant $(\text{CCSDT(Q)} - \text{CCSD(T)})/\text{VDZ(NP)}$ term (~ 0.6 kcal mol⁻¹) among
287 these transition states, demonstrating that the post-CCSD(T) contributions are almost uniform across this reaction series. These
288 observations provide compelling evidence that both alkyl substitution and carbon-chain elongation negligibly modulate the
289 magnitude of post-CCSD(T) corrections, implying that such higher-order correlation effects are intrinsically insensitive to
290 substituent-induced electronic and conformational changes.

291 3.4. Electronic structure of CH₂OO + RCHO (R = CH₂F/CHF₂/CF₃)

292 The electronic structure information was depicted in Figure 2 and Table S12. The activation enthalpies at 0 K decrease
293 significantly with the increasing number of fluorine substitutions in the methyl group of the aldehyde.

294 The ΔH_0^\ddagger for CH₂OO + CH₂FCHO (TS7) is -6.21 kcal mol⁻¹ by our best estimate, which is 1.24 kcal mol⁻¹ and 1.71
295 kcal mol⁻¹ lower than the reaction R1 and R2, respectively. Consequently, reaction R7 is expected to exhibit a significantly
296 larger rate constant compared to the CH₂OO + HCHO/CH₃CHO reactions. This reduction in ΔH_0^\ddagger indicates that fluorine
297 substitution enhances the reactivity of the aldehyde toward CH₂OO, which is similar to HO₂ + CF₃CHO (Long et al., 2022).
298 For the reaction of CH₂OO + CHF₂CHO (R8), the ΔH_0^\ddagger is -7.96 kcal mol⁻¹, which is 1.75 kcal mol⁻¹ lower than that of the
299 corresponding transition state, TS7. This value is close to that of CH₂OO + HCl (Foreman et al., 2016), which approaches the
300 bimolecular collision limit, suggesting that the reaction R8 through the tight transition state is not the rate-determining step.
301 Although fluorine substitution on the methyl group of the aldehyde leads to substantially enhanced reactivity toward CH₂OO,
302 the post-CCSD(T) contributions from the $(\text{CCSDT(Q)} - \text{CCSD(T)})/\text{VDZ(NP)}$ term (~ 0.6 kcal mol⁻¹) remain nearly identical
303 across the transition states as shown in Figure 2, revealing that the higher-order correlation effects are largely insensitive to
304 fluorination and establishing that the fluorination-driven reactivity enhancement originates primarily from lower-level
305 electronic effects than that of post-CCSD(T).

306 Given the demonstrated accuracy of the M11-L/MG3S method for reactions R7 and R8, this method was subsequently
307 applied to reaction R9, as depicted in Figure S3. Regarding $\text{CF}_3\text{CHO} + \text{CH}_2\text{OO}$ (R9), the ΔH_0^\ddagger further decreases to -9.74
308 kcal/mol at M11-L/MG3S level. However, this value is slightly higher than the activation enthalpies observed for the universal
309 mechanism of Criegee intermediates reacting with amides (Long et al., 2025), which are significantly submerged below the
310 reactants by approximately 9 to 11 kcal/mol. This shows that this tight transition state is not the rate-determining step for
311 reaction R9.

312 We further compare the calculated ΔH_0^\ddagger of the $\text{CH}_2\text{OO} + \text{RCHO}$ ($\text{R} = \text{CH}_2\text{F}, \text{CHF}_2, \text{CF}_3$) reactions with those of the
313 corresponding OH reactions. The ΔH_0^\ddagger for $\text{OH} + \text{CH}_2\text{FCHO}$ is $-1.15 \text{ kcal mol}^{-1}$ at the CCSD(T)/M06-2X/aug-cc-pVTZ level,
314 which is $5.06 \text{ kcal mol}^{-1}$ higher than that of R7. We also find that the ΔH_0^\ddagger for R8 by our best estimate is $8.19 \text{ kcal mol}^{-1}$ lower
315 than that of $\text{OH} + \text{CHF}_2\text{CHO}$, calculated at the CCSD(T)/aug-cc-pVDZ//MP2(FC)/aug-cc-pVDZ level. The ΔH_0^\ddagger for R9
316 calculated by M11-L/MG3S is 11.94 kcal/mol lower than that of $\text{OH} + \text{CF}_3\text{CHO}$ at QCISD(T)/6-311G(d,p) level (Chandra et
317 al., 2001). The present findings reveal that the much lower ΔH_0^\ddagger for R7-R9 leads to a much faster rate constant, indicating
318 that oxidation by CH_2OO contributes significantly to the atmospheric loss of fluorinated aldehydes relative to the OH-initiated
319 pathway from energetic point of view.

320 **3.5. Kinetics**

321 **3.5.1 Pressure-dependent rate constants.**

322 The pressure dependence of the rate constants for reactions R1 and R2 was evaluated using the ME/RRKM framework,
323 with the results summarized in Tables S13–S15. As shown in Table S13, reaction R1 exhibits no appreciable pressure
324 dependence over the conditions examined, indicating that pressure effects can be safely neglected for this channel. This
325 conclusion is fully consistent with the findings reported by Luo et al (Luo et al., 2023). For example, the falloff factor calculated
326 for the $\text{CH}_2\text{OO} + \text{HCHO}$ reaction at 298 K and 0.0316 bar is 1.34 (Table S13). This factor, defined as the ratio of the rate
327 constant at 1000 bar to that at 0.0316 bar, indicates only a weak pressure dependence for this system. We observed that at 295
328 K and 78 Torr, the pressure-dependent rate constant was $2.71 \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ in Table S13, which is 7.74 times

329 higher than the reported value $((3.50 \pm 0.35) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1})$ in Table 1 (Enders et al., 2024).

330 We assessed the validity of the simplified pathway by contrasting the full mechanism (Scheme 2a) with the model
331 (Scheme 2b) from a kinetic perspective as listed in Tables S13 and S14. The pressure-dependent rate constants obtained from
332 both models exhibit negligible deviations, thereby validating the simplified scheme as a computationally efficient strategy for
333 larger aldehydes. The calculated pressure-dependent rate constant for reaction R2 is $1.84 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ at 293 K
334 and 4 Torr in Table S14, in good agreement with the value of $(9.50 \pm 0.70) \times 10^{-13} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ reported by Taatjes et al
335 (Taatjes et al., 2012). Our pressure-dependent rate constant at 298 K and 25 Torr corroborates the experimental value of $(1.20$
336 $\pm 0.20) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ reported by Elsamra et al ($1.65 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ in Table S14) (Elsamra et al.,
337 2016). We found that the fall-off factor is only 1.36 (Table S14) for the reaction R2 at 298 K and 4 Torr, which also shown that
338 the rate constant of reaction R2 is negligibly pressure-dependent, which confirms the experimental results qualitatively (Enders
339 et al., 2024; Stone et al., 2014; Berndt et al., 2015; Jiang et al., 2024). In addition, there is experimental evidence that the
340 pressure effect is also insignificant for propionaldehyde and butyraldehyde (Liu et al., 2020; Debnath and Rajakumar, 2024).

341 3.5.2 High pressure limit rate constants

342 High-pressure limit rate constants for all reactions are summarized in Table 5, with additional details provided in Tables
343 S16–S24. The rate constants in the temperature range of 190–350 K were fitted using the four-parameter expression (Zheng
344 and Truhlar, 2012; Bao et al., 2016a):

$$345 k_{\infty} = A \left(\frac{T+T_0}{300} \right)^n \exp \left[-\frac{E(T+T_0)}{R(T^2+T_0^2)} \right] \quad (7)$$

346 Where R is the gas constant, T is temperature in K, the fitting parameters were listed in Table S25. The temperature dependence
347 of the Arrhenius activation energies was further calculated using the following expression:

$$348 E_a = -R \frac{d \ln k_{\infty}}{d(1/T)} \quad (8)$$

349
350 **Table 5.** The high-pressure limiting rate constants ($\times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$) of the $\text{CH}_2\text{OO} + \text{RCHO}$ (R =
351 H/CH₃/C₂H₅/C₃H₇/C₄H₉/C₅H₁₁/CH₂F/CHF₂/CF₃) reaction

T/K	k_1	k_2	k_3	k_4	k_5	k_6	k_7	k_8	k_9
-----	-------	-------	-------	-------	-------	-------	-------	-------	-------

200	426	94.9	109	119	110	214	430	451	740
220	297	32.4	37.3	42.1	37.1	71.9	248	416	688
240	171	13.4	15.5	18.1	15.3	29.5	115	381	652
260	90.4	64.9	7.48	8.98	6.95	13.1	51.1	328	626
280	48.0	3.52	4.05	5.00	3.74	7.09	24	252	607
298	28.3	2.20	2.52	3.19	2.34	4.46	13.1	182	594
300	26.8	2.10	2.40	3.05	2.28	4.41	12.3	174	593
320	15.9	1.35	1.53	2.00	1.41	2.67	6.90	110	583
340	10.0	0.19	1.04	1.39	0.95	1.80	4.17	67.4	576

352 **The reaction of CH₂OO + HCHO.** As summarized in Table 1 and Figure 1, a long-standing order-of-magnitude
353 discrepancy exists between previously reported experimental and theoretical rate constants for reaction R1. At 296 K, the rate
354 constant obtained in this work is $3.01 \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ in Table 1, which is 7.31 times larger than the experimental
355 value reported by Luo et al (Luo et al., 2023), but 1.83 and 2.17 times smaller than the theoretical predictions of Zhang et al.
356 (Zhang et al., 2023) and Long et al. (Long et al., 2021), respectively. We therefore consider two plausible explanations: The
357 experimental determination of CH₂OO kinetics may introduce systematic uncertainties. Alternatively, subtle dynamic effects
358 beyond conventional transition state theory (e.g., non-statistical dynamics or complex-forming behavior) may play a role and
359 require further investigation. Although the present value does not fully reconcile the experimental and theoretical results, it
360 substantially narrows the gap between the two, providing a quantitatively improved estimate for this key reaction.

361 Notably, the derived rate constant for R1 is approximately 8 times larger than that for the corresponding OH-initiated
362 reaction and more than two orders of magnitude larger than that for the HO₂-initiated pathway (Long et al., 2022; Sivakumaran
363 et al., 2003), highlighting the unexpectedly high reactivity of CH₂OO in this system. These findings underscore the need for
364 further high-precision experimental measurements and establish the present computational protocol as a robust framework for
365 resolving persistent discrepancies in atmospheric reaction kinetics.

366 **The reaction of CH₂OO + CH₃CHO.** To date, no theoretical kinetic studies have been reported for the CH₂OO + CH₃CHO
367 reaction in Table 2. The earliest experimental determination yielded a rate constant of $(9.50 \pm 0.25) \times 10^{-13} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$
368 at 293 K and 4 Torr, as measured by Taatjes et al. (Taatjes et al., 2012), which is a factor of 2.9 smaller than the present
369 theoretical prediction in Table 2. At 298 K, the calculated rate constant for reaction R2 is $2.20 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ in

370 Table 2, in excellent agreement with the experimental values reported by Elsamra et al. (Elsamra et al., 2016) and Jiang et al.
371 (Jiang et al., 2024) In addition, the value measured by Berndt et al. (Berndt et al., 2015) at 297 K, $(1.7 \pm 0.50) \times 10^{-12} \text{ cm}^3$
372 $\text{molecule}^{-1} \text{ s}^{-1}$, is fully consistent with our calculated result of $2.27 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ in Table 2. Overall, the rate
373 constants obtained in this work are in good agreement with the available experimental data (Elsamra et al., 2016; Stone et al.,
374 2014; Berndt et al., 2015; Jiang et al., 2024; Cornwell et al., 2023), providing the first reliable theoretical benchmark for the
375 kinetics of the $\text{CH}_2\text{OO} + \text{CH}_3\text{CHO}$ reaction. Notably, the rate constant for R2 is approximately 5.6 times smaller than that for
376 the corresponding OH-initiated reaction, yet nearly two orders of magnitude larger than that for the HO_2 -initiated pathway,
377 highlighting the distinct and non-negligible role of CH_2OO in aldehyde oxidation chemistry (Long et al., 2022; Zhu et al.,
378 2008). The five-membered ring species C2a and C2b readily interconvert, as the rate constant for the isomerization process is
379 approximately two orders of magnitude larger than that of the addition reaction (Table S26).

380 ***The reaction of $\text{CH}_2\text{OO} + \text{RCHO}$ ($\text{R}=\text{C}_2\text{H}_5/\text{C}_3\text{H}_7/\text{C}_4\text{H}_9/\text{C}_5\text{H}_{11}$).*** Rate constants for the reactions of CH_2OO with
381 $\text{C}_2\text{H}_5\text{CHO}$ have been reported previously from both experimental and theoretical studies (See Table 3) (Enders et al., 2024;
382 Kaipara and Rajakumar, 2018; Liu et al., 2020), whereas the reaction with $\text{C}_3\text{H}_7\text{CHO}$ has been examined only experimentally.
383 At 298 K, the calculated rate constant for $\text{CH}_2\text{OO} + \text{C}_2\text{H}_5\text{CHO}$ is $2.52 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ (Table 3), in excellent
384 agreement with the experimental value reported by Liu et al (Liu et al., 2020).

385 For $\text{CH}_2\text{OO} + \text{C}_3\text{H}_7\text{CHO}$, the calculated rate constant of $3.19 \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ (Table 3) closely reproduces the
386 experimental value of $(2.63 \pm 0.14) \times 10^{-12} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ (Debnath and Rajakumar, 2024), further validating the reliability
387 of the present computational protocol. To the best of our knowledge, no prior experimental or theoretical studies have reported
388 rate constants for the reactions of CH_2OO with pentanal or hexanal. Our calculations indicate that the rate constant for CH_2OO
389 + $\text{C}_4\text{H}_9\text{CHO}$ is comparable to that for CH_3CHO , whereas the rate constant for $\text{CH}_2\text{OO} + \text{C}_5\text{H}_{11}\text{CHO}$ is approximately twice as
390 large, yet remains within the same order of magnitude (Table 5). These results demonstrate that increasing alkyl chain length
391 exerts only a minor influence on the reaction kinetics of CH_2OO with aldehydes, revealing a weak and nonmonotonic size
392 dependence across the C_1 – C_5 series. This behavior is fully consistent with the computed activation enthalpies (See Figure 2)
393 and establishes a transferable structure–reactivity relationship for CH_2OO reactions with larger aldehydes. Overall, aside from
394 formaldehyde, the rate constants for CH_2OO reactions with alkyl-substituted aldehydes vary only modestly, underscoring the

395 limited role of substituent size in governing CH₂OO reactivity.

396 **The reaction of CH₂OO + RCHO (R=CH₂F/CHF₂/CHF₃).** A striking fluorination-induced reactivity enhancement
397 emerges upon substitution of hydrogen atoms on the methyl group. Introduction of fluorine leads to a pronounced increase in
398 the rate constants for CH₂OO + CH₃CHO reactions, revealing an unexpected structure–reactivity trend. At 298 K, the rate
399 constant for reaction R7 is $1.31 \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ (Table 5), which is about 6 times larger than that of R2 and about 5
400 times larger than the corresponding OH + CH₂FCHO reaction (Lily et al., 2021).

401 Even more dramatic behavior is observed for reactions R8 and R9. For R8, the calculated rate constants approach the
402 collision limit, decreasing slightly from $4.51 \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ at 200 K to $6.74 \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ at 340 K
403 in Table 5, indicating of a weak negative temperature dependence characteristic of barrierless processes. Notably, at 298 K the
404 reaction of CHF₂CHO with CH₂OO is more than two orders of magnitude faster than its reactions with OH [$(1.8 \pm 0.4) \times 10^{-12}$
405 $\text{cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$] (Sellevåg et al., 2005), underscoring the unusually high reactivity of CH₂OO toward fluorinated aldehydes.

406 The most pronounced effect is found for R9, for which the rate constant ranges from $7.40 \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ at
407 200 K to $5.76 \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ at 340 K in Table 5, fully approaching the collision limit and exceeding the
408 corresponding OH-initiated reaction rates by orders of magnitude. These results demonstrate that fluorination fundamentally
409 alters the reaction landscape of CH₂OO with aldehydes, transforming otherwise moderately fast bimolecular reactions into
410 near-collision-controlled processes.

411 3.6. Atmospheric Implications

412 The reaction of aldehydes with OH have been investigated extensively experimentally and theoretically. Here, we
413 considered the competition for aldehydes relative to CH₂OO and OH. The ratio of reaction rate was calculated by eqn (9):

$$414 \quad v_i = \frac{k_i[\text{CH}_2\text{OO}]}{k_{\text{OH},i}[\text{OH}]} \quad (9)$$

415 where the k_i is the rate constants for the reaction R2-R9, $k_{\text{OH},i}$ is the rate constant of OH + RCHO (R = CH₃, C₂H₅, C₃H₇, C₄H₉,
416 C₃H₁₁, CH₂F, CHF₂, CF₃), and i is referred to is equal to 2-9. The concentrations of CH₂OO and OH exhibit pronounced
417 geographical and spatial distributions. The concentration of OH is varied from 10^4 - $10^6 \text{ molecules cm}^{-3}$ (Khan et al., 2018; Ren
418 et al., 2003; Stone et al., 2012), and the estimated concentration for CH₂OO is range from 10^4 to $10^5 \text{ molecules cm}^{-3}$ (peaking

419 at 6×10^5 molecules cm^{-3}) (Lelieveld et al., 2016; Novelli et al., 2017) In contrast, the base-version model simulations yield
 420 CH_2OO concentrations approximately one order of magnitude lower than the estimated value. This discrepancy likely
 421 originates from (i) the adoption of relatively fast rate constants for CH_2OO loss via reactions with H_2O and $(\text{H}_2\text{O})_2$, and (ii) an
 422 incomplete representation of CH_2OO sources in the model framework. Consequently, the use of model-derived concentrations
 423 probably leads to an underestimation of the contribution of CH_2OO to aldehyde removal.

424 Our results demonstrate that for aliphatic aldehydes, reactions with CH_2OO constitute a negligible sink compared with OH
 425 oxidation, owing to both modest rate constants and low ambient CH_2OO concentration (See Tables S27–S29). Although
 426 fluorine substitution generally enhances reactivity, the increase in the rate constant for CH_2FCHO remains insufficient to
 427 meaningfully compete with the OH pathway. Effective competition is predicted only under highly specific conditions—namely,
 428 nighttime at ~ 10 km altitude over the Malaysian region (Table S29). In stark contrast, the reactions of highly fluorinated
 429 aldehydes with CH_2OO proceed at near-collision-limit rates. As a result, CH_2OO constitutes a major atmospheric sink for
 430 CHF_2CHO and CF_3CHO . As summarized in Table 6, CH_2OO competes effectively with OH for CHF_2CHO at night near the
 431 surface over Russia and the Arctic, influences its removal at 5 km over Russia and Indonesia, and contributes significantly at
 432 10 km over Indonesia. Notably, because the reaction of CF_3CHO with OH is intrinsically slow, CH_2OO dominates its
 433 atmospheric removal over Indonesia at all altitudes considered, while in the Russian region its influence is confined to 0 and
 434 5 km.

435 **Table 6.** rate concentration ratios CH_2OO to OH and the rate ratio at different heights from different region

Height	T/K	P/mBar	$[\text{CH}_2\text{OO}]/[\text{OH}]^a$	v_8^b	v_9^c
Gansu, China					
1	290.2	1013	2.48×10^{-4}	2.89×10^{-2}	2.29×10^{-1}
5	250.5	495.9	3.09×10^{-4}	6.25×10^{-2}	3.03×10^{-1}
10	215.6	242.8	3.51×10^{-5}	8.14×10^{-3}	3.77×10^{-2}
Russia					
1	290.2	1013	1.52×10^{-2}	1.77	14
5	250.5	495.9	6.39×10^{-3}	1.29	6.26
10	215.6	242.8	3.23×10^{-5}	7.48×10^{-3}	3.47×10^{-2}
Arctic					
1	290.2	1013	1.15×10^{-2}	1.33	10.6

5	250.5	495.9	5.16×10^{-4}	1.04×10^{-1}	5.05×10^{-1}
10	215.6	242.8	1.91×10^{-6}	4.43×10^{-4}	2.05×10^{-3}
Indonesia					
1	290.2	1013	3.16×10^{-3}	3.67×10^{-1}	2.91
5	250.5	495.9	5.85×10^{-3}	1.18	5.74
10	215.6	242.8	2.53×10^{-2}	5.87	27.2

436 ^aThe concentration ratio between CH₂OO and OH from GEOS-Chem.

437 ^bThe rate ratio between CH₂OO + CHF₂CHO and CHF₂CHO + OH.

438 ^cThe rate ratio between CH₂OO + CF₃CHO and CF₃CHO + OH.

439 Overall, these findings reveal a qualitative shift in aldehyde oxidation pathways upon heavy fluorination, identifying
 440 CH₂OO as a previously underappreciated but potentially dominant oxidant for highly fluorinated aldehydes under specific
 441 atmospheric regimes—an effect with important implications for the atmospheric lifetimes of emerging fluorinated oxygenated
 442 VOCs.

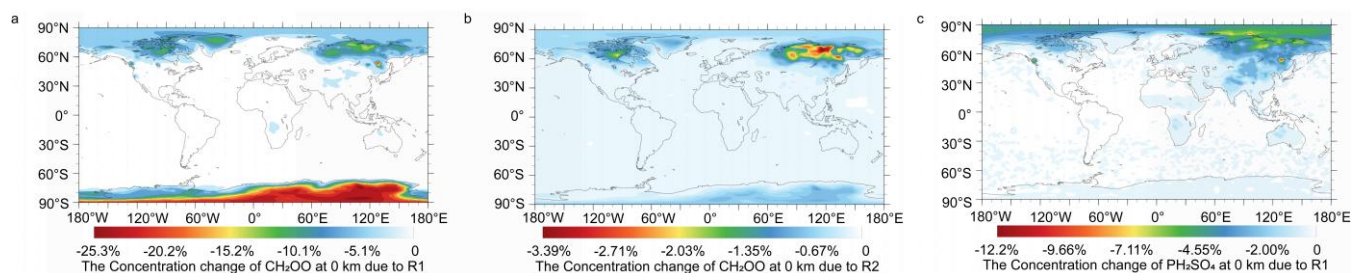
443 3.7. Atmospheric modelling

444 Model simulations were further performed to assess the atmospheric significance of nighttime reactions between CH₂OO
 445 and aldehydes. The Criegee intermediate (CI) chemistry implemented in the base model has been described in our previous
 446 work (Long et al., 2024). In this study, two targeted updates were introduced to isolate and quantify the impacts of newly
 447 identified CI–aldehyde reaction pathways. The first update incorporates the CH₂OO + HCHO reaction into the base mechanism,
 448 reflecting an improved understanding of CI removal under aldehyde-rich nighttime conditions. The second update further
 449 expands the CI sink by including the reaction between CH₂OO and CH₃CHO, thereby providing a more comprehensive
 450 representation of acetaldehyde-driven CI loss. The aldehyde chemistry employed in the model is summarized in Table S30.
 451 We do not consider the impact of CH₂OO on fluorinated aldehyde sinks by using GEOS-Chem, as fluorinated aldehydes are
 452 not involved in the current default GEOS-Chem version.

453 The simulated aldehyde concentrations exhibit pronounced spatial and vertical heterogeneity. Surface-level HCHO
 454 concentrations reach up to 1.46×10^{11} molecules cm⁻³, while CH₃CHO attains maxima of 8.06×10^{10} molecules cm⁻³, with
 455 the highest abundances over Malaysia and Indonesia. These values are consistent with field observations, which report peak
 456 HCHO concentrations of up to 3.63×10^{11} molecules cm⁻³ (Hu et al., 2025), lending confidence to the model performance.

457 The simulated global mean surface concentration of CH_3CHO (5.89×10^9 molecules cm^{-3} , corresponding to ~ 200 ppt) is in
458 reasonable agreement with observational constraints and remains lower than values reported by Komazaki et al (Komazaki et
459 al., 1999; Tereszchuk and Bernath, 2011).

460 The contribution of HCHO to the reduction of CH_2OO has been assessed in our prior work and is once again validated
461 by model simulations (Long et al., 2021). Figure 5 shows the relative changes in annual mean surface-layer CH_2OO
462 concentrations resulting from the inclusion of the $\text{CH}_2\text{OO} + \text{HCHO}$ (R1) and $\text{CH}_2\text{OO} + \text{CH}_3\text{CHO}$ (R2) reactions. Incorporation
463 of the updated rate constant for R1 leads to a pronounced reduction in CH_2OO , with a maximum decrease of 25.3% over the
464 Antarctic region (Figure 5), highlighting the previously unrecognized importance of HCHO as a nighttime CI sink. In contrast,
465 R2 produces a more modest effect, with a maximum CH_2OO reduction of 3.39% over Russia in Figure 5.



466 **Figure 5.** Changes in global CH_2OO concentrations due to reaction R1 and R2 (a) reaction R1, (b) reaction R2, and (c) changes
467 in global sulfate concentrations due to reaction R1.
468

469 Despite the substantial impact on CI abundances, the direct effects on aldehyde concentrations remain small. As shown in
470 Figure S8, surface acetaldehyde decreases by only 0.12% in the Arctic. However, the influence on secondary oxygenated
471 products is more pronounced. As illustrated in Figure S9, inclusion of R1 enhances formic acid concentrations by up to 5.44%
472 over Canada and Russia, while acetic acid increases by as much as 0.69% in the Arctic. These results demonstrate that CI-
473 aldehyde reactions, while exerting limited feedback on aldehydes themselves, can make significant contribution to the sinks
474 of CH_2OO and the formation of atmospheric acids.

475 The potential implications of reaction R1 for regional air quality were also assessed, particularly regarding the mitigation of
476 gas-phase sulfate formation. We found that the concentration of gas-phase sulfate can reach 10^8 molecules cm^{-3} in Mexico
477 region in Figure S10. The inclusion of this reaction pathway effectively lowers the concentration of CH_2OO , thereby
478 diminishing its capacity to oxidize SO_2 into sulfuric acid precursors. This depletion of oxidative capacity leads to a marked

479 decrease in gas-phase sulfate concentration. The effect is geographical, with the reduction in gas-phase sulfate concentrations
480 estimated to be 12.2% in Canada and 6.01% in Russia during the nighttime in Figure 5c. While the relative changes might
481 initially imply a substantial regional sink for atmospheric sulfate aerosols, a detailed comparison of Figures 5c and S10 reveals
482 that the largest percentage changes in gas-phase sulfate predominantly occur in regions with low baseline concentrations.
483 Specifically, although peak concentrations over Canada and Russia reach $\sim 10^7$ molecules cm^{-3} , their regional averages remain
484 on the order of 10^5 molecules cm^{-3} . In contrast, regions with much higher absolute concentrations (e.g., $\sim 1 \times 10^8$ molecules
485 cm^{-3} over Mexico) exhibit only minimal relative changes. This indicates that modest absolute variations can produce large
486 percentage changes under low-background conditions, whereas comparable or even larger absolute changes appear
487 insignificant in high-concentration environments. Consequently, this reaction has a negligible impact on the global atmospheric
488 sulfate burden.

489 4. CONCLUSIONS

490 The present work establishes a transferable and systematically improvable theoretical framework for predicting quantitative
491 atmospheric reaction kinetics across molecular complexity, using the reactions of CH_2OO with a series of aldehydes as a
492 definitive test case. By explicitly approaching the full configuration interaction (CI) limit for the benchmark $\text{CH}_2\text{OO} + \text{HCHO}$
493 system, we delineate the accuracy requirements necessary for reliable kinetic predictions and provide a rigorous reference
494 against which lower-cost methods can be assessed. Energetic and kinetic analyses validate a simplified reaction mechanism,
495 attributed to the facile interconversion between complexes and the energetic preference for rotational transition states over
496 addition pathways.

497 Guided by the detailed electronic-structure insights obtained for $\text{CH}_2\text{OO} + \text{HCHO}$, we develop a computational protocol
498 that integrates optimized geometries, vibrational frequencies, and high-level single-point energies, enabling accurate kinetics
499 for larger systems at feasible computational cost. We find that DF-CCSD(T)-F12b/VDZ(d) and DF-CCSD(T)-F12b/jun-cc-
500 pVDZ can be used to reliably describe the optimized geometries and calculated frequencies. two generalizable strategies (BE1
501 and BE2) have been used to recover the CCSDTQ/CBS level single point energies, which provide new insight into how to

502 obtain the quantitative enthalpy of activation.

503 In kinetics calculations, for reactions with appreciable barriers (R2–R6), this dual-level strategy yields robust rate
504 constants, whereas for reactions characterized by exceptionally low or submerged barriers (R1 and R7–R9), the explicit
505 application of VRC-VTST proves essential for capturing the correct dynamical behavior. This demonstrates a practical pathway
506 for extending benchmark-level kinetics from small to chemically diverse, larger molecules.

507 The resulting kinetic trends reveal that alkyl-chain elongation exerts only a minor influence on reactivity, whereas fluorine
508 substitution dramatically enhances reaction rates, driving the $\text{CH}_2\text{OO} + \text{CHF}_2\text{CHO}$ and $\text{CH}_2\text{OO} + \text{CF}_3\text{CHO}$ reactions toward
509 the collision limit. All reactions exhibit negligible pressure dependence, underscoring their relevance under atmospheric
510 conditions. These high-precision rate constants provide a mechanistically grounded explanation for the increasingly important
511 role of Criegee intermediates in the oxidation of fluorinated aldehydes. We find that fluorine substitution on aldehydes
512 dramatically enhances their reactivity toward CH_2OO ; however, the post-CCSD(T) contributions remain almost equal across
513 the reaction series. This behavior indicates that fluorination-driven rate acceleration is governed primarily by lower-level
514 electronic effects rather than by higher-order electron correlation than CCSD(T). This observation also provides a fundamental
515 basis for the development of high-accuracy semiempirical correction schemes.

516 Beyond molecular-scale kinetics, global and regional modeling demonstrates that while reactions of CH_2OO with HCHO
517 and CH_3CHO contribute negligibly to aldehyde removal, HCHO constitutes a major global sink for Criegee intermediates,
518 accounting for a 25.3% reduction in the global CH_2OO burden during the night. In contrast, fluorination fundamentally alters
519 atmospheric fate: for CH_2FCHO , CH_2OO reactions become regionally significant (e.g., near 10 km altitude over Malaysia),
520 and for more heavily fluorinated aldehydes such as CHF_2CHO , CH_2OO overwhelmingly dominates over OH-initiated loss
521 pathways. The associated enhancement in acid formation, although modest, further highlights the chemical implications of
522 these processes. The inclusion of reaction R1 results in a reduction of gas-phase sulfate levels by 12.2% over Canada and 6.01%
523 over Russia. These present findings deliver a generalizable, benchmark-anchored strategy for quantitative kinetic prediction,
524 bridges electronic-structure theory with atmospheric modeling, and reveals how fluorination reshapes the atmospheric
525 relevance of Criegee intermediates—insights that are critical for atmospheric chemical mechanisms.

526

527 **Supplement.** The following information is provided in the Supplement: Details of reaction R9, enthalpies of binding and
528 activation and barrier height; vibrational frequency scale factors; Lennard-Jone parameters; Rate constants and rate constant
529 fits; Rate ratio; Absolute energies and the Cartesian coordinates and absolute energies; relative enthalpies for reaction of R3-
530 R9; Enthalpy profile for the conversion of pre-reaction complex; Changes in global CH₃CHO, HCOOH, and CH₃COOH
531 concentrations.

532

533 **Data and code availability.** Electronic structure calculations were performed using commercially available software
534 (Gaussian 16, Revision A.03 and Molpro 2019). Access to the software is subject to licensing terms. The MRCC and MStor
535 codes can be accessed at <https://www.mrcc.hu> and <https://comp.chem.umn.edu/mstor>, respectively. Polyrate 2017-C and
536 Gausstrate 2017-B are available at <https://comp.chem.umn.edu/polyrate> and <https://comp.chem.umn.edu/gausstrate>. KiSThIP
537 2021 is accessible at <http://kisthelp.univ-reims.fr>, and the TUMME program can be found at
538 <https://comp.chem.umn.edu/tumme>. The GEOS-Chem 14.4.2 is available at <http://www.geos-chem.org>. Optimized
539 geometries, and calculated energies are available in Supplement. Other data are available from the corresponding author
540 upon reasonable request.

541

542 **Author contributions.** Chaolu Xie carried out the calculations, analysed and interpretation of data, and wrote the
543 manuscript draft. Bo Long designed the project, analyzed and interpretation of data, and reviewed and edited the
544 manuscript.

545

546 **Competing interests.** The authors declare that they have no conflict of interest.

547

548 **Acknowledgements.** We also thank the Minnesota Supercomputing Institute for computational resources

549

550 **Financial support.** This work was supported in part by the National Natural Science Foundation of China (42120104007 and
551 41775125), by the Guizhou Provincial Science and Technology Projects, China (CXTD [2022]001 and GCC [2023]026), and
552 by the U.S. Department of Energy under Award DE-SC0015997, Guizhou Graduate Research Fund Project under Grant

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