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2 **Parameterization and Evaluation of Nonhydrostatic Effect in the**
3 **Orographic Gravity Wave Drag in China Meteorological**
4 **Administration Global Forecast System (CMA-GFS) v4.0 Model**

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Abstract

21 The China Meteorological Administration Global Forecast System (CMA-GFS)
22 v4.0 model was upgraded to a higher resolution of 0.125° in May 2023. To be compatible
23 with its fine resolution, the parameterization scheme of orographic gravity wave drag
24 (OGWD) in CAM-GFS is revised herein by accounting for the nonhydrostatic effect (NHE)
25 on the wave momentum flux of subgrid-scale orographic gravity waves. The performance
26 of the revised OGWD scheme is then evaluated for the 10-day medium-range forecast in
27 December 2023. Results show that the revised OGWD scheme can better capture the large-
28 scale circulation in the Northern Hemisphere (NH), particularly in the high latitudes. The
29 easterly (westerly) wind biases in the NH polar stratosphere (troposphere) are decreased.
30 The underestimation of East Asia subtropical jet is also alleviated. Quantitative evaluation
31 shows that the revised OGWD scheme reduces both the mean bias and root mean square
32 error of 500-hPa geopotential height in the NH after the 6th forecast day, reaching 11.59%
33 and 5.06%, respectively, by day 10. The decrease of easterly biases in the polar stratosphere
34 is owing to the weakening of stratospheric zonal OGWD by the NHE. It is also contributed
35 by the secondary circulation induced by the weakened OGWD, which increases the thermal
36 contrast between middle and high latitudes by adiabatic warming/cooling. For the decrease
37 of westerly biases in the NH polar troposphere, it is due to the fact that the enhanced
38 stratospheric winds suppress the upward propagation of Rossby waves into the stratosphere,
39 resulting in greater convergence of Eliassen-Palm flux in the mid-upper troposphere.

40 **Keywords:** orographic gravity wave drag, parameterization, global NWP medium-range
41 forecast

42



43 **1. Introduction**

44 Orographic gravity wave drag (OGWD) is an important process in atmospheric
45 dynamics, arising from the interaction of airflow with complex terrain (Kim et al., 2003;
46 Teixeira, 2014). When airstream flows over mountains, it generates orographic gravity
47 waves (OGWs) which propagate vertically, when breaking transferring momentum
48 from the surface to higher levels (Fritts and Alexander, 2003). This momentum transfer
49 plays a crucial role in driving the atmospheric circulation, influencing both tropospheric
50 and stratospheric dynamics (Alexander et al., 2010).

51 In numerical weather prediction (NWP) and climate models, the OGWD is
52 typically a subgrid-scale process which needs to be parameterized. Various OGWD
53 parameterization schemes have been developed over the past few decades (e.g., Palmer
54 et al., 1986; McFarlane, 1987; Kim and Arakawa, 1995; Lott and Miller, 1997; Scinocca
55 and McFarlane, 2000; Kim and Doyle, 2005) based on both linear and nonlinear OGW
56 dynamics. Their implementation has been shown to help alleviate the systematic biases
57 in both NWP and climate models, ranging from general circulation to regional climate
58 and weather (e.g., Kim, 2007; McLandress et al., 2012; Choi and Hong, 2015; Zhong
59 and Chen, 2015; Chen et al., 2016; Lu et al., 2020; Zhang et al., 2020; Li et al., 2023;
60 Xu et al., 2023; Wei et al., 2025).

61 While accurate representation of OGWD is essential for weather forecast and
62 climate simulation/projection, traditional OGWD parameterization schemes rely on the
63 assumption of hydrostatic balance which can significantly simplify the formulae of
64 OGW dynamics. This assumption is appropriate for coarse-resolution numerical models
65 where subgrid-scale OGWs are dominated by hydrostatic GWs as the dominant subgrid



66 “mountains” are large. However, as the model resolution increases, the hydrostatic
67 assumption becomes less valid because the subgrid-scale orography (SSO) becomes
68 smaller, so that the unresolved GWs have shorter wavelengths. In this situation, the
69 nonhydrostatic effects (NHEs) will exert remarkable influences on the subgrid-scale
70 OGWs. [Note that NHE should also affect the OGWD in coarser-resolution models but
71 only has a weaker impact. This is because subgrid OGWs are generally in hydrostatic
72 balance in coarser-resolution models, with very few momentum fluxes being carried by
73 nonhydrostatic OGWs.] For example, nonhydrostatic OGWs propagate more laterally
74 than hydrostatic ones. The horizontal dispersion of wave energy can reduce the wave
75 amplitude (Smith, 1979; Klemp and Durran, 1983; Zängl, 2003), which substantially
76 suppresses gravity wave breaking and the deposition of wave momentum into the mean
77 flow. These NHEs are thus critical for accurately representing the dynamical impacts
78 of subgrid-scale OGWs in high-resolution models, e.g., the *state-of-the-art* global NWP
79 models.

80 Recently, Xu et al. (2021) theoretically derived the analytical expressions for
81 the surface wave momentum flux (WMF) of nonhydrostatic OGWs generated by
82 idealized three-dimensional orography. They found that the degree of nonhydrostaticity
83 can be measured by a nondimensional parameter of Froude number which is equal to
84 the wind speed over the mountain half width and buoyancy frequency. Physically, this
85 parameter represents the ratio between the period of buoyancy oscillation and the time
86 for airflow travelling through the mountain. The larger the horizontal Froude number,
87 the more important the NHE is. Based upon the theoretical study, Xu et al. (2023, 2024)
88 revised the OGWD parameterization scheme developed by Kim and Doyle (2005,



89 hereafter KD05) by accounting for the NHE on the surface WMF of upward-
90 propagating OGWs. Then the new OGWD scheme was implemented in the Model for
91 Prediction Across Scales (MPAS), which was shown to improve the seasonal
92 simulation of the stratospheric polar night jet and reduce the wet biases over the western
93 Tibetan Plateau in winter. Li et al. (2024) further evaluated the revised OGWD scheme
94 in the Weather Research and Forecasting (WRF) model for short-range forecast of
95 Northeast China cold vortices (NECVs). The underestimation of the NECV intensity
96 (in terms of minimum 500-hPa geopotential height) is alleviated as the NHE decreases
97 the lower-tropospheric OGWD.

98 The China Meteorological Administration Global Forecast System (CMA-GFS)
99 is a rename of the Global/Regional Assimilation and Prediction System (GRAPES)
100 developed in early 2000s (Shen et al., 2017). Operationally running at the CMA Earth
101 System Modeling and Prediction Center, the CMA-GFS was upgraded to a new version
102 of v4.0 in May 2023, with its horizontal resolution increasing from about 0.25° to 0.125°
103 (about 13 km), along with many other improvements in the model dynamics and physics.
104 For instance, the convective triggering function and quasi-equilibrium closure
105 conditions are improved to reduce biases in tropics and enhance forecast skill of
106 precipitation in East Asia. In CMA-GFS, three components of subgrid orographic
107 effects have been implemented, i.e., the blocking-flow drag (BFD), the OGWD and
108 turbulent orographic form drag (TOFD). The TOFD scheme was implemented CMA-
109 GFS based on Beljaars et al. (2004) and Xue et al. (2011) and will not be further
110 discussed as it is not relevant to the NHE effect studied in this study. The BFD
111 component based on Lott and Millor (1997, hereafter LM97) and OGWD component



112 from Kim and Arakawa (1995, hereafter KA95) were implemented in CMA-GFS
113 similar to that in Alpert (2004) and was described in Chen et al. (2016). It is noticed
114 that the NHE effects were partially considered in KA95 by partitioning the momentum
115 stress with the Scorer parameter when model grid mainly locates at the downstream of
116 the subgrid orography. However, the surface OGWD remains with the original KA95
117 scheme. In order to be compatible with its high resolution, this hydrostatic OGWD
118 parameterization scheme needs to be upgraded to account for the NHE.

119 Note that the present study extends the evaluation of the nonhydrostatic OGWD
120 scheme in the MPAS model presented in Xu et al. (2024) by transitioning it into a
121 practical, operational NWP. Although Xu et al. (2024) demonstrated the beneficial
122 impact of the revised OGWD scheme for long-term, large-scale circulation simulations,
123 its performance and value within a state-of-the-art, high-resolution global forecasting
124 system performing routine medium-range forecasts had not been tested. This work
125 provides this critical assessment using the CMA-GFS v4.0 model, with a specific focus
126 on forecast skill improvements for several quantitative metrics. The primary advance
127 lies in evaluating the revised OGWD scheme in reducing systematic forecast biases on
128 timescales directly relevant to weather prediction (1-10 days). Therefore, this research
129 bridges the gap between parameterization development and operational implementation,
130 demonstrating the tangible benefits of incorporating nonhydrostatic effects to improve
131 medium-range weather forecasting accuracy.

132 The remainder of the paper is organized as follows. Section 2 firstly introduces
133 the revision of the OGWD parameterization scheme and then describes the setup of the
134 numerical experiments. Section 3 gives an overall evaluation for the medium-range



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135 forecast of large-scale atmospheric circulation by the CMA-GFS model. The NHE
136 effects on parameterized OGWD and large-scale circulation are examined in section 4.
137 Finally, the paper is summarized in section 5 along with discussions.

138

139 **2. OGWD parameterization scheme and numerical experiments**

140 a. Revision of the KA95 OGWD parameterization scheme

141 The KA95 OGWD considers various aspects of the SSO including its standard
142 deviation (σ_h), orographic asymmetry (OA) and orographic convexity (OC). The
143 surface WMF of OGWs is given by

$$144 \quad \tau_0 = \rho_0 E \frac{m}{\lambda_{eff}} G \frac{|V_0|^3}{N_0}, \quad (1)$$

145 where

$$146 \quad E = (OA + 2)^{C_E \frac{Fr_0}{Fr_c}}, m = (1 + L_x)^{OA+1}, G = \frac{Fr_0^2}{Fr_0^2 + C_G OC^{-1}}, Fr_0 = \frac{\sigma_h N_L}{|V_L|}. \quad (2)$$

147 The variables ρ_0 , $|V_0|$ and N_0 are the low-level mean air density, horizontal wind speed
148 and buoyancy frequency. The parameter m denotes the “number of mountains” within
149 the grid cell, characterizing the total volume of SSO associated with the orography
150 length (L_x). The coefficient λ_{eff} acts as a tunable parameter denoting the effective grid
151 length. The transition function G is an enhancement factor for sharp mountains as OC
152 is large. This scheme also considers the effects of low-level wave breaking and/or lee
153 wave trapping through the factor E , which is a function of the shape and location of the
154 SSO within the model grid cell (i.e., OA), and the flow nonlinearity (i.e., Fr_0). The two
155 constants are set to $C_E = 0.8$ and $C_G = 0.5$, which are obtained empirically through a
156 series of numerical simulations conducted in Kim and Arakawa (1995).



157 The surface WMF is transported upward level by level until reaching the model
158 top or critical level. At each model level, airflow instability is checked according to the
159 hypothesis of wave saturation (Lindzen, 1981) and wave-breaking (Miller and Palmer,
160 1986).

161 Once the subgrid-scale OGWs saturate and break, they will exert a body force
162 on the mean flow through the deposition of wave momentum, namely,

$$163 \quad \frac{d\mathbf{V}(z)}{dt} = \frac{1}{\rho(z)} \frac{\partial \boldsymbol{\tau}(z)}{\partial z}. \quad (3)$$

164 where $\mathbf{V}(z)$ and $\boldsymbol{\tau}(z)$ are the mean flow velocity and the WMF at height z , respectively.
165 Note that both the directions of $\mathbf{V}(z)$ and $\boldsymbol{\tau}(z)$ are parallel to the mean flow at the low
166 level (i.e., \mathbf{V}_L). Readers are referred to KA95 for more details about the scheme.

167 The hydrostatic KA95 scheme above is revised by taking into account the NHE
168 on the surface WMF of OGWs, i.e.,

$$169 \quad \tilde{\boldsymbol{\tau}} = \tau_0 [1 + NHE(Fr)]. \quad (4)$$

170 where the NHE correction depends only on the horizontal Froude number $Fr = \frac{|\mathbf{V}_0|}{N_0 L_x}$

$$171 \quad NHE(Fr) = -\frac{9}{8} Fr^2 + e^{-2Fr^{-1}} \left(-\frac{5}{4} Fr^{-2} - \frac{1}{2} Fr^{-1} + \frac{5}{4} + \frac{9}{4} Fr + \frac{9}{8} Fr^2 \right). \quad (5)$$

172 Note that the horizontal Froude number is different from the traditional Froude number
173 Fr_0 above which measures the flow nonlinearity. Moreover, Eq. (5) is the NHE derived
174 for three-dimensional isotropic terrain. While Xu et al. (2021) have derived the fully
175 NHE for anisotropic terrain, the expressions involve complicated integrals that are not
176 suitable for practical use in parameterization [cf. their Eqs. (14) to (17)]. Xu et al. (2021)
177 demonstrated that terrain anisotropy only has a very weak influence on the NHE, so the
178 expression of NHE for isotropic terrain, i.e., Eq. (5), is used because of simplicity, since



179 only algebraic manipulations are involved. This is in agreement with Xu et al. (2023,
180 2024) and Li et al. (2024).

181

182 b. Setup of numerical experiments

183 The CMA-GFSv4.0 dynamical core is based upon the nonhydrostatic and
184 shallow atmospheric governing equations in spherical polar coordinates with full
185 physics packages (Shen et al., 2020). It adopts a regular latitude-longitude grid of
186 $0.125^\circ \times 0.125^\circ$ with *C*-grid staggering. In the vertical, a terrain-following height-based
187 grid is utilized along with Charney-Phillips staggering. There are 87 vertical levels,
188 with the model top located at 73 km. For model physics, the Liu-Ma microphysics
189 scheme (Ma et al., 2018) is employed, along with the RRTMG longwave and shortwave
190 radiation schemes (Morcrette et al., 2008), the MRF planetary boundary layer scheme
191 (Chen et al. 2020; Hong and Pan 1996), the CoLM land surface model (Dai et al., 2003),
192 NSAS convection scheme (Han and Pan, 2011; Liu et al., 2015), the combined KA95
193 OGWD and LM97 FBD parameterization scheme (Chen et al., 2016), and Beljaars's
194 TOFD scheme (Xue et al., 2011).

195 In order to examine the impact of the revised OGWD scheme on the medium-
196 range forecast, two sets of numerical simulations (i.e., EXP_CTL and EXP_NHE) are
197 conducted with different OGWD parameterization schemes, using the original KA95
198 scheme and the revised nonhydrostatic one, respectively. Both consist of 31 individual
199 runs with their start dates from 00UTC 1 December to 00UTC 31 December in 2023.
200 Herein, a winter month is chosen because the OGWD, which is mainly located in the
201 Northern Hemisphere (NH) given the pronounced orography, is the strongest in this



202 season (e.g., Xu et al. 2020; Lu et al. 2024). In each experiment, the CMA-GFS model
203 is integrated for 10 days, i.e., medium-range forecast, with 6-hr output interval. The
204 model initial conditions are derived from the $0.25^\circ \times 0.25^\circ$ ECMWF Reanalysis v5
205 (ERA5) dataset (Hersbach et al., 2020), which are also used as reference for the
206 evaluation of the CMA-GFS forecasts.

207

208 **3. Evaluation of the medium-range forecast**

209 *a. Atmospheric circulation*

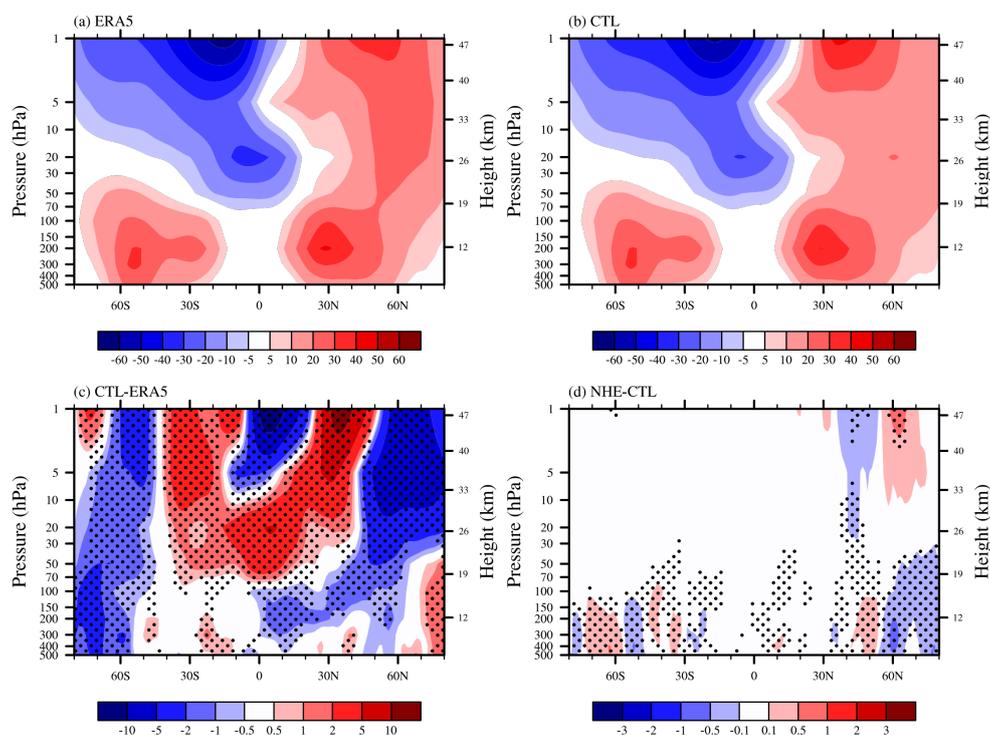
210 Figure 1 shows the zonal-mean zonal wind averaged during the period of 10
211 December 2023 to 10 January 2024 (i.e., the 10th day of the medium-range forecast).
212 In the winter of NH, as revealed by the ERA5 reanalysis (Fig. 1a), there is a subtropical
213 jet in the midlatitudes troposphere, with the jet core located at about 200 hPa and 30°N .
214 The maximum zonal wind speed exceeds 35 m s^{-1} . In the stratosphere of high latitudes,
215 there is another wind maxima near 1 hPa and 60°N , which is the well-known polar-
216 night jet (e.g., Kim, 2007). These two jets are separated around 70 hPa and 50°N , with
217 the stratospheric jet being stronger than its tropospheric counterpart. In the Southern
218 Hemisphere (SH), which is in summer, there is also a tropospheric jet but relatively
219 weaker and lower than in the NH. The SH jet core is located at about 50°S , with a
220 secondary one near 30°S . Given in the summer, the SH stratosphere is dominated by
221 easterlies which is distinctly different from that in the NH. With a center located near
222 1 hPa and 20°S , the stratospheric easterlies extend downward and equatorward to the
223 tropics. An easterly jet of over -30 m s^{-1} can be found near 20 hPa and 5°S .



224 Figure 1b is similar to Figure 1a but for the zonal-mean zonal winds obtained in
225 the EXP_CTL experiment. In general, the CMA-GFS model can capture the overall
226 pattern of the zonal-mean zonal winds, such as the tropospheric westerly jets in both
227 hemispheres and the easterly winds in the SH stratosphere. The polar night jet in the
228 NH, however, is underestimated by the model, with the jet core shifted southward by
229 about 10 latitudes. As shown in Fig. 1c, there are notable easterly biases of over -5 m
230 s⁻¹ in the stratosphere of the NH high latitudes (north of 50°N). By contrast, westerly
231 biases are present in the upper stratosphere of the NH midlatitudes (25°N-50°N) which
232 can exceed 10 m s⁻¹, extending downward and equatorward to the lower stratosphere of
233 the tropics. As in the NH troposphere, the zonal-mean zonal winds are overestimated
234 in the Arctic region, but with easterly biases in the mid-lower latitudes. In the SH, the
235 stratospheric easterlies are shifted northward, leading to westerly biases in the mid-
236 lower latitudes (40°S-5°S) and easterly biases in the tropics and lower latitudes of the
237 NH (5°S-20°N). In the high latitudes, there are predominantly easterly biases in both
238 troposphere and stratosphere except in the upper stratosphere of the Antarctic region.

239 When taking into account the NHE in the OGWD parameterization scheme, the
240 CMA-GFS model can better capture the large-scale circulation. Figure 1d presents the
241 differences between the zonal-mean zonal winds in the two numerical experiments (i.e.,
242 EXP_NHE minus EXP_CTL). In the stratosphere above ~10 hPa, positive and negative
243 wind differences are found to the north and south of 60°N, respectively, which are just
244 opposed to the wind biases in Figure 1c. It suggests that both the magnitude and location
245 of the polar night jet are improved. For example, the easterly biases in the mid-upper
246 stratosphere (~10 hPa to 1 hPa) are reduced by about 3%. The westerly biases in the

247 troposphere of Arctica region are also reduced, which reaches up to about 42% of the
248 total bias.



249

250 Figure 1. Vertical distributions of zonal-mean zonal wind (units: m s^{-1}) averaged in the
251 period of 10 December 2023 to 10 January 2024 obtained from (a) ERA5 and (b)
252 EXP_CTL, with their difference (i.e., EXP_CTL minus ERA5) given in (c). (d)
253 is similar to (c) but for the difference between the two experiments of EXP_NHE
254 EXP_CTL (i.e., EXP_NHE minus EXP_CTL). Stippling in (c) and (d) denote
255 differences statistically significant at the 95% confidence level.

256

257 For the zonal-mean zonal wind biases in the SH stratosphere, they are hardly
258 alleviated, however. This is because the stratospheric OGWD is very weak in summer
259 (see Fig. 5a of Xu et al. 2024; see also Fig. 3 below), given the absorption of OGWs at
260 the critical level (Booker and Bretherton, 1967), that is, the zero-wind level between

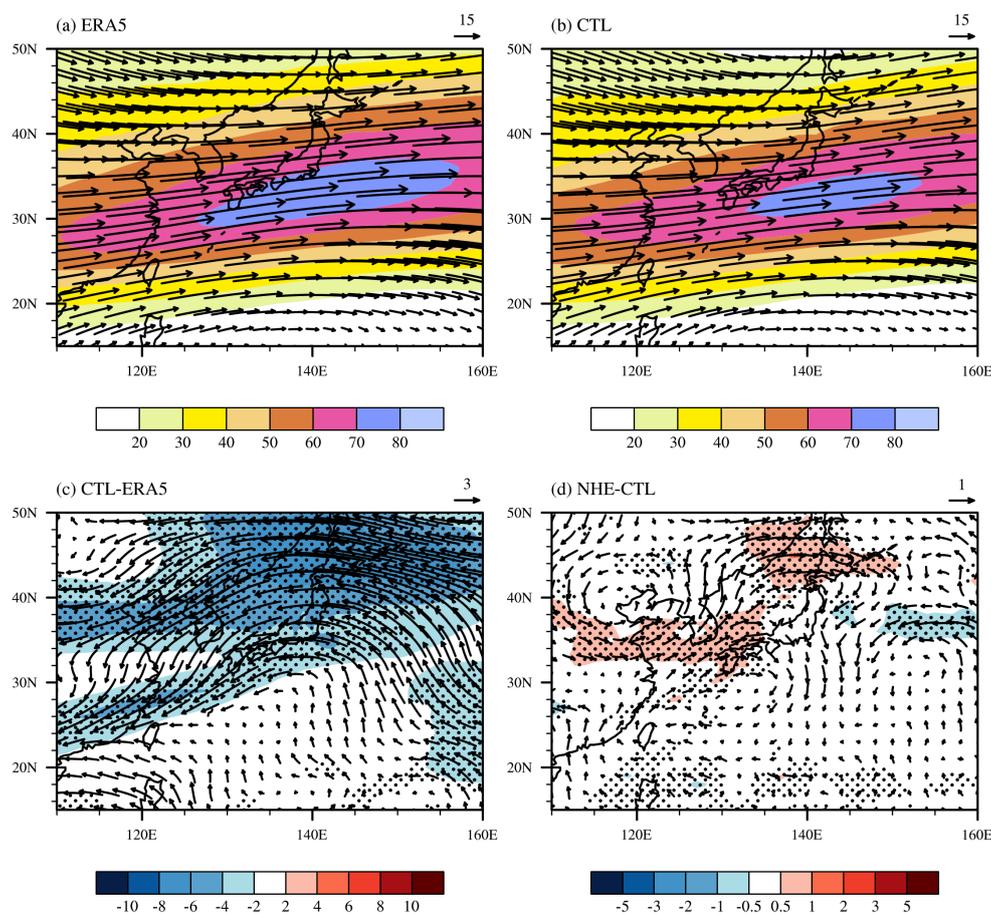


261 the tropospheric westerlies and the stratospheric easterlies (Fig. 1a). The easterly biases
262 in the troposphere of the high latitudes are reduced, especially around 60°S.

263 Although there is no significant improvement for the zonal-mean zonal wind in
264 the NH mid-lower latitudes (Fig. 1d), the wind circulation can be improved regionally,
265 such as the East Asia subtropical jet (EASJ). Figure 2a shows the horizontal wind field
266 and speed at 200 hPa averaged in the period of 10 December 2023 to 10 January 2024
267 from the ERA5 reanalysis. In boreal winter, the EASJ stretches from about 110°E to
268 160°E in the latitudes between about 25°N and 45°N. The jet core is located near Japan,
269 showing a high wind speed of over 70 m s⁻¹. The simulated EASJ in the EXP_CTL
270 experiment generally aligns well with the ERA5 reanalysis, in terms of its location and
271 orientation (Fig. 2b). But the jet intensity is underestimated, showing easterly biases of
272 over -4 m s⁻¹ (Fig. 2c). In the EXP_NHE experiment, the underestimation of the EASJ
273 is alleviated, with the mean bias decreased by about 6% (Fig. 2d).

274

275



276

277 Figure 2. Horizontal distributions of wind speed (shading; units: m s^{-1}) and wind field
278 (arrow) at 200 hPa averaged in the period of 10 December 2023 to 10 January 2024
279 obtained from (a) ERA5 and (b) EXP_CTL, with their difference (i.e., EXP_CTL minus
280 ERA5) given in (c). (d) is similar to (c) but for the difference between the two
281 experiments of EXP_NHE and EXP_CTL (i.e., EXP_NHE minus EXP_CTL).
282 Stippling in (c) and (d) denote differences statistically significant at the 95% confidence
283 level.

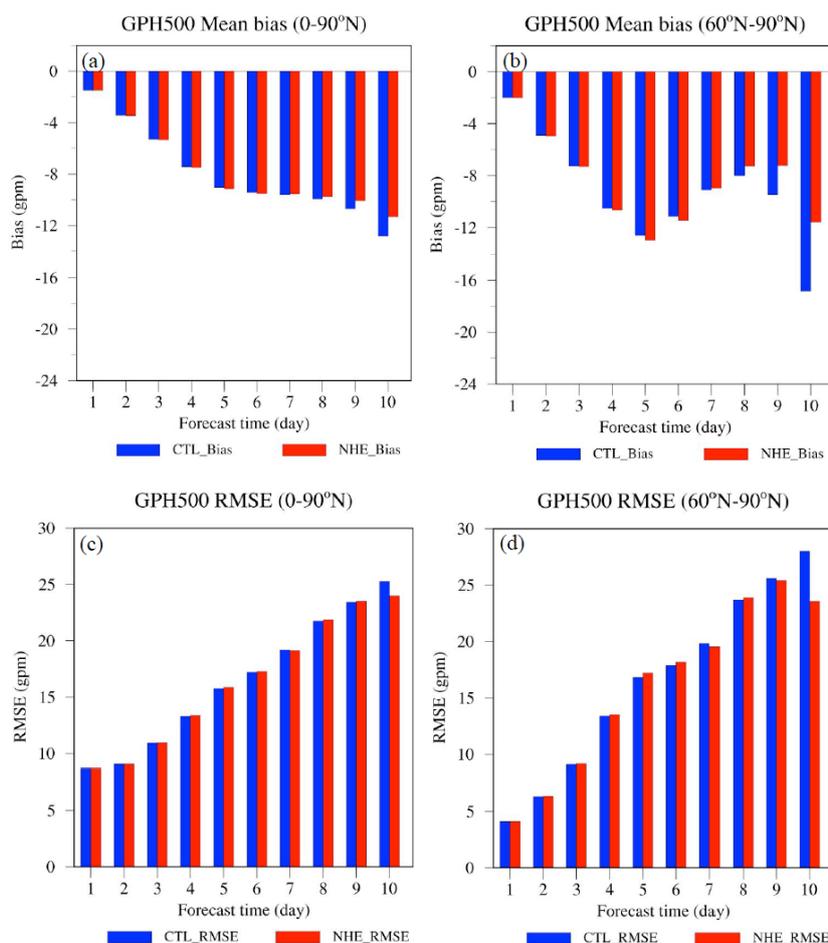
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285 *b. Quantitative evaluation of forecast skills*

286 For medium-range forecast, the wind circulation in the troposphere is often of
287 greater interest than in the stratosphere. To better evaluate the performance of the



288 revised OGWD scheme, Figures 3a and 3b depict the mean bias (MB) and root mean
289 square error (RMSE) for the 500 hPa geopotential height (GPH500) in the two
290 experiments. Herein, we are interested in the NH because of the relatively weak wind
291 differences in the SH (Fig. 1d) which is mainly covered by the ocean. Moreover, as will
292 be shown below, the parameterized OGWD changes little in the SH.



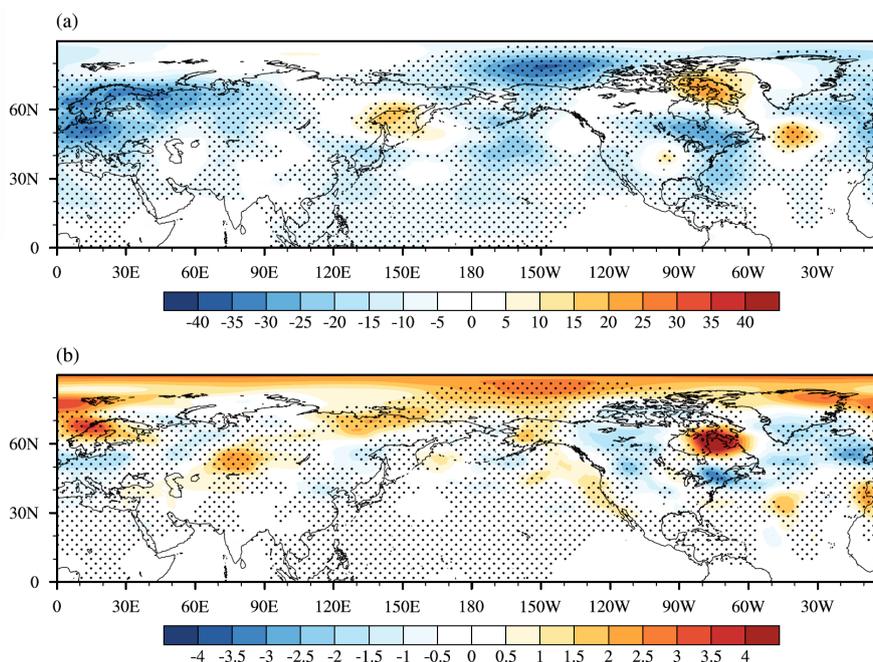
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294 Figure 3. (a) Mean bias and (c) RMSE of the GPH500 in the NH at different forecast
295 lead time. (b) and (d) are similar to (a) and (c) but in the region north of 60°N.

296



297 In both experiments, the model GPH500 is lower than the ERA5 reanalysis, with
298 the MB and RMSE increasing with the forecast lead time (Figs. 3a, 3c). The GPH500
299 shows obvious underestimation in the EXP_CTL experiment which are significantly
300 reduced in EXP_NHE. This improvement is particularly evident in high latitudes north
301 of 60°N which is overall statistically significant at the 95% confidence level (Figs. 4a,
302 4b). Compared to EXP_CTL, EXP_NHE exhibits even greater MB and RMSE of
303 GPH500 till the 6th forecast day. Nevertheless, both the MB and RMSE of the GPH500
304 decrease quickly afterwards, which are reduced by 11.59% and 5.06%, respectively, at
305 the 10th day of the forecast (Table 1). Greater improvements are found for the GPH500
306 between 60°N and 90°N, the MB and RMSE of which are decreased by 31.18% and
307 15.93%, respectively, at the 10th forecast day (Table 1). All the values shown in Table
308 1 are statistically significant at the 95% confidence level. For the total 10 forecast days
309 as a whole, the MB and RMSE of GPH500 are separately reduced by 2.6% and 0.52 %,
310 indicating an overall improvement in the simulation of the large-scale circulation in the
311 NH when using the revised OGWD scheme (see Table 2).



312

313 Figure 4. (a) Mean biases of GPH500 (shading; units: gpm) in EXP_CTL experiment
 314 as compared to ERA5 averaged over the ten forecast days. (b) is similar to (a) but for
 315 the differences between the GPH500 (shading; units: gpm) in EXP_CTL and
 316 EXP_NHE experiments (i.e., EXP_NHE minus EXP_CTL). Stippling in (a) and (b)
 317 denote differences statistically significant at the 95% confidence level.

318

319 Table 1. Mean bias and RMSE of the 500-hPa geopotential height (GPH500) and sea
 320 level pressure (SLP) at the 10th forecast day

	Region	Mean Bias			RMSE		
		CTL	NHE	$\frac{NHE - CTL}{CTL} \times 100$	CTL	NHE	$\frac{NHE - CTL}{CTL} \times 100$
GPH500 (gpm)	0-90°N	-12.825	-11.338	-11.59	25.263	23.984	-5.06
	60°N-90°N	-16.841	-11.590	-31.18	28.017	23.554	-15.93
SLP (hPa)	0-90°N	-0.522	-0.343	-34.29	2.714	2.515	-8.33
	60°N-90°N	-1.712	-1.079	-36.97	3.049	2.517	-17.45

321

322 Table 2. Mean bias and RMSE of the 500-hPa geopotential height (GPH500) and sea
 323 level pressure (SLP) for the overall ten forecast days

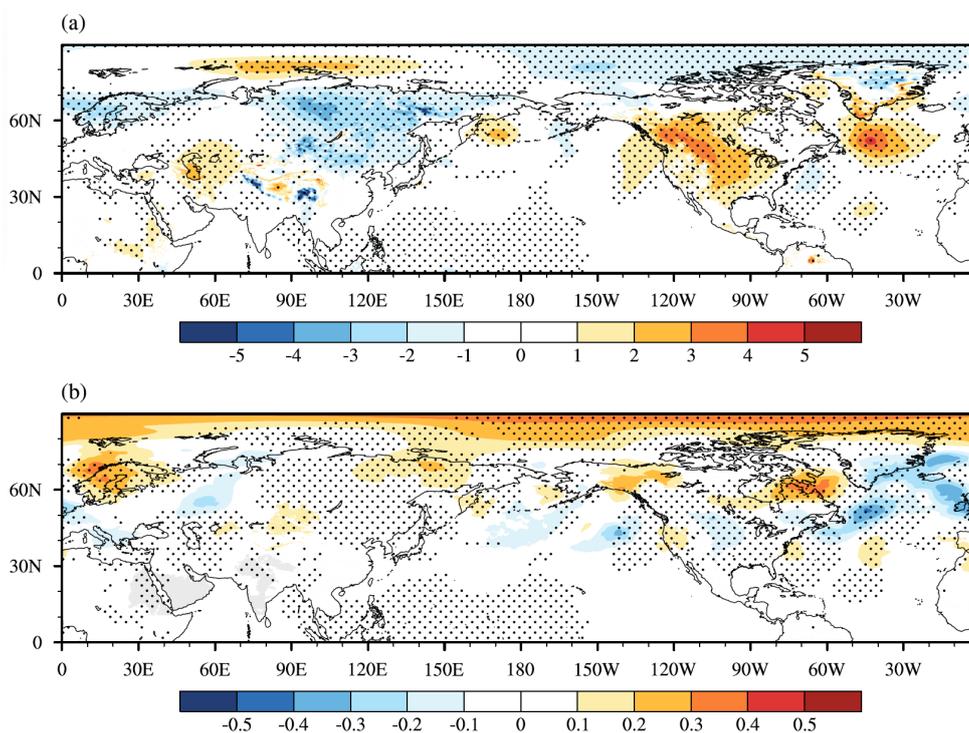
	Region	Mean Bias			RMSE		
		CTL	NHE	$\frac{NHE - CTL}{CTL} \times 100$	CTL	NHE	$\frac{NHE - CTL}{CTL} \times 100$
GPH500 (gpm)	0-90°N	-7.923	-7.717	-2.6	16.455	16.369	-0.52
	60°N-90°N	-9.187	-8.438	-8.15	16.470	16.089	-2.31
SLP (hPa)	0-90°N	-0.107	-0.082	-23.46	1.87	1.841	-1.53
	60°N-90°N	-0.498	-0.411	-17.47	1.598	1.532	-4.13

324

325 Besides the GPH500, the forecast skill is also examined for the SLP in the NH.
 326 As shown in Fig. 5a, the EXP_CTL experiment systematically underestimates the SLP
 327 over East Asia and the Arctic, while overestimating SLP across North America and the
 328 northern Atlantic. These biases are substantially corrected in the EXP_NHE experiment,
 329 particularly in high-latitude regions north of 60°N (Fig. 5b). Similarly, the EXP_NHE
 330 experiment firstly experiences a degradation in the early 6 days compared to EXP_CTL,
 331 but it shows significant improvement at the end of the 10th forecast day, with the MB
 332 reduced by 34.29% and the RMSE by 8.33%, respectively (Table 1). Greater decreases
 333 of 36.97% (for MB) and 17.45% (for RMSE) are found in the high latitudes north of
 334 60°N as well.

335

336



337

338 Figure 5. (a) Mean biases of SLP (shading; units: hPa) in the EXP_CTL experiment as
339 compared to ERA5 averaged over the ten forecast days. (b) is similar to (a) but for the
340 differences between the SLP (shading; units: hPa) in the EXP_CTL and EXP_NHE
341 experiments (i.e., EXP_NHE minus EXP_CTL). Stippling in (a) and (b) denote
342 differences statistically significant at the 95% confidence level.

343

344 From the above analyses, implementing the revised OGWD scheme in the high-
345 resolution CMA-GFS model can help improve the medium range forecast of the NH
346 large-scale circulation, especially in the high latitudes. In the next section, we will
347 examine the underlying mechanisms responsible for the improvement of the circulation.

348

349



350 **4. Physical interpretation**

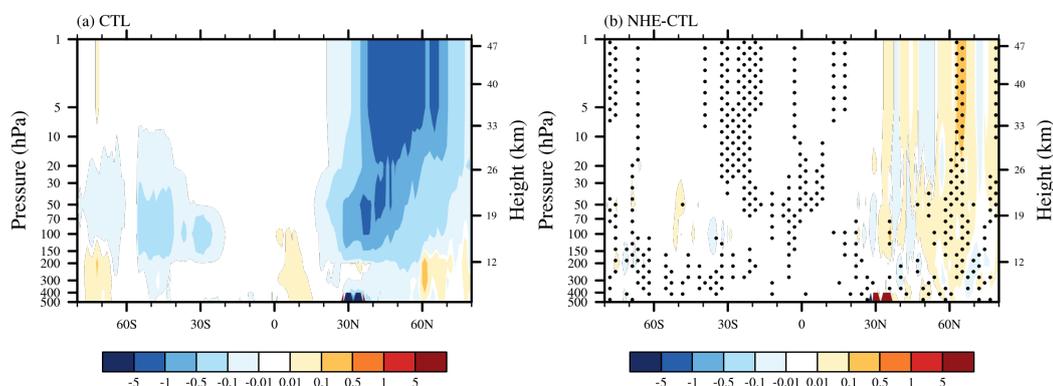
351 *a. Changes of parameterized OGWD*

352 The changes of wind circulation arise from the revision of the OGWD scheme.
353 It is thus straightforward to study the parameterized OGWD in these two experiments.
354 Figure 6a shows the vertical distribution of zonal-mean zonal OGWD averaged in the
355 period of 10 December 2023 to 10 January 2024 in the EXP_CTL experiment. The
356 differences between the parameterized OGWD in the two experiments (i.e., EXP_NHE
357 minus EXP_CTL) are given in Fig. 6b. In boreal winter, there exists prominent OGWD
358 in the stratosphere of the NH mid-to-high latitudes, which is favored by the decrease of
359 air density with height (e.g., Lindzen, 1981) and the relatively weak winds between the
360 tropospheric jet and polar night jet. The latter is called as the “valve layer” by Kruse et
361 al. (2016). The maximum OGWD occurs between about 20 hPa and 1 hPa. In contrast,
362 there is very weak OGWD in the SH stratosphere owing to the critical level absorption
363 as mentioned above.

364 Compared to that in EXP_CTL, the zonal-mean zonal OGWD in EXP_NHE is
365 generally weakened in the high latitudes of the NH (north of about 55°N), especially in
366 the stratosphere (Fig. 6b). Note that the positive difference indicates a decrease of
367 OGWD as the drag itself is negative. This is due to the fact that the NHE acts to decrease
368 the surface WMF [see Eqs. (4) and (5)], i.e., the source of parameterized OGWs which
369 determines the maximum wave momentum that can be deposited into the mean flow.
370 Therefore, the decrease of surface WMF is prone to reduce the OGWD. Note that there
371 is also enhancement of OGWD, e.g., in the stratosphere near 50°N. This is because the
372 decrease of surface WMF may suppress wave breaking in the troposphere due to

373 reduced gravity wave amplitude. In consequence, more WMF is transported to the
374 stratosphere where the OGWs break owing to the decay of air density and lead to
375 stronger OGWD. This is similar to the redistribution of WMF in the vertical owing to
376 selective critical level filtering of OGWs in directional shear flows (Shutts, 1995; Xu
377 et al., 2012, 2019; van Niekerk et al., 2023). Generally, the changes in the vertical
378 structure of parameterized OGWD in the CMA-GFS model are similar to that in the
379 MPAS seasonal simulations conducted in Xu et al. (2024) which also considers the
380 NHE correction to the OGWD (cf. their Fig. 5).

381



382

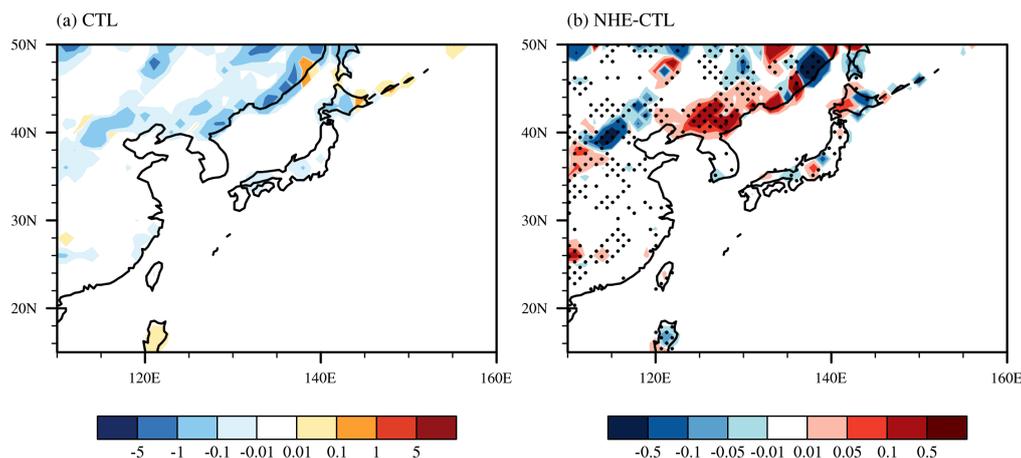
383 Figure 6. Vertical distributions of (a) zonal-mean zonal OGWD (units: $\text{m s}^{-1} \text{day}^{-1}$)
384 averaged in the period of 10 December 2023 to 10 January 2024 obtained from
385 EXP_CTL, and (b) zonal-mean OGWD difference between the two experiments of
386 EXP_NHE and EXP_CTL (i.e., EXP_NHE minus EXP_CTL). Stippling in (b) denotes
387 differences statistically significant at the 95% confidence level.

388

389 By the way, the parameterized OGWD in East Asia is also examined to explain
390 the changes of EASJ. As shown in Fig. 7a, there is notable westward OGWD at 200
391 hPa in the high latitudes of East Asia where the underlying terrain is very complex,

392 such as the Taihang, Yanshan, Changbai, Greater Khingan and Lesser Khingan
393 Mountains. Thus, the largest easterly biases of the EASJ occur in this region (Fig. 2c).
394 When accounting for the NHE in the OGWD parameterization scheme, the 200-hPa
395 OGWD is mainly reduced (Fig. 7b). Clearly, the enhanced EASJ is owing to the
396 weakening of the parameterized OGWD.

397



398

399 Figure 7. Horizontal distributions of (a) zonal OGWD (units: $\text{m s}^{-1} \text{ day}^{-1}$) at 200 hPa
400 averaged in the period of 10 December 2023 to 10 January 2024 obtained from
401 EXP_CTL, and (b) zonal OGWD difference between the two experiments of EXP_NHE
402 and EXP_CTL (i.e., EXP_NHE minus EXP_CTL) at 200 hPa. Stippling in (b) denotes
403 differences statistically significant at the 95% confidence level.

404

405 *b. Wave-induced secondary circulation: downward control mechanism*

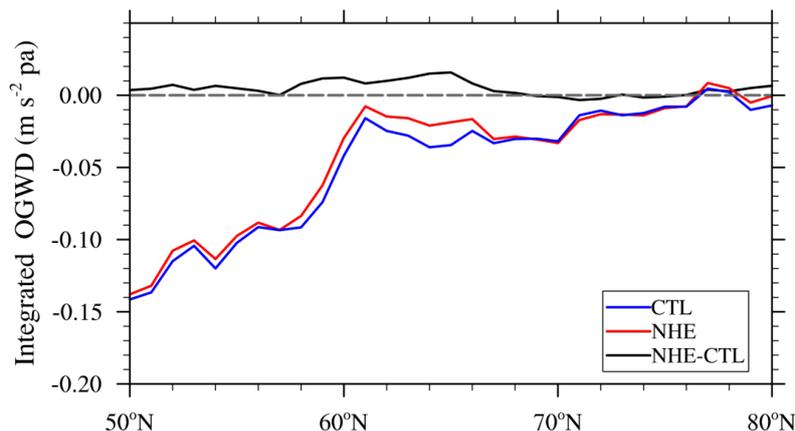
406 The weakening of stratospheric OGWD can directly increase the zonal wind in
407 the polar stratosphere of the NH, thus leading to a recover of the weakened westerly jet
408 (Figs. 1c, 1d). It can also affect the tropospheric circulation by inducing a secondary
409 circulation according to the “downward control” principle (Haynes et al., 1991), i.e.,

410
$$\omega(p) = \frac{1}{r \partial \varphi} \left[\frac{1}{f} \int_0^p \bar{F} dp \right]. \quad (6)$$

411 where r is the radius of the earth, φ is the latitude, f is the Coriolis parameter, and ω
412 is the vertical velocity in the pressure coordinate. The vertical velocity of the secondary
413 circulation at a given pressure level is proportional to the meridional gradient of the
414 vertically-integrated zonal wave forcing (\bar{F}) above this pressure level. Westward wave
415 forcing ($\bar{F} < 0$) acts to induce ascending ($\omega < 0$) and sinking ($\omega > 0$) motions on the
416 equatorward and poleward sides of the forcing, respectively. In contrast, the direction
417 of the meridional circulation is reversed in the case of eastward wave forcing ($\bar{F} > 0$).

418 Figure 8 depicts the zonal-mean zonal OGWD in the mid-high latitudes of the
419 NH which are vertically integrated from 200 hPa to the model top. In both experiments,
420 the integrated zonal OGWD shows a decreasing trend from 50°N to 80°N. Compared
421 to the EXP_CTL experiment, the EXP_NHE experiment produces an overall weaker
422 westward OGWD between 50°N and 80°N, owing to the decrease of OGWD in the
423 stratosphere (Fig. 6b). Therefore, the meridional gradient of the zonal OGWD is
424 reduced, leading to greater upward motion poleward of about 70°N according to Eq.
425 (6). In response to this upward motion, the NH polar stratosphere cools adiabatically,
426 as can be seen clearly in Fig. 9. In contrast, the stratosphere warms adiabatically to the
427 south of 60°N where the descending branch of the wave-induced secondary circulation
428 exists. Given these temperature changes, the thermal contrast between the middle and
429 high latitudes of the NH amplifies, which enhances the zonal wind in the stratosphere
430 through the thermal wind relation for large-scale atmospheric motion.

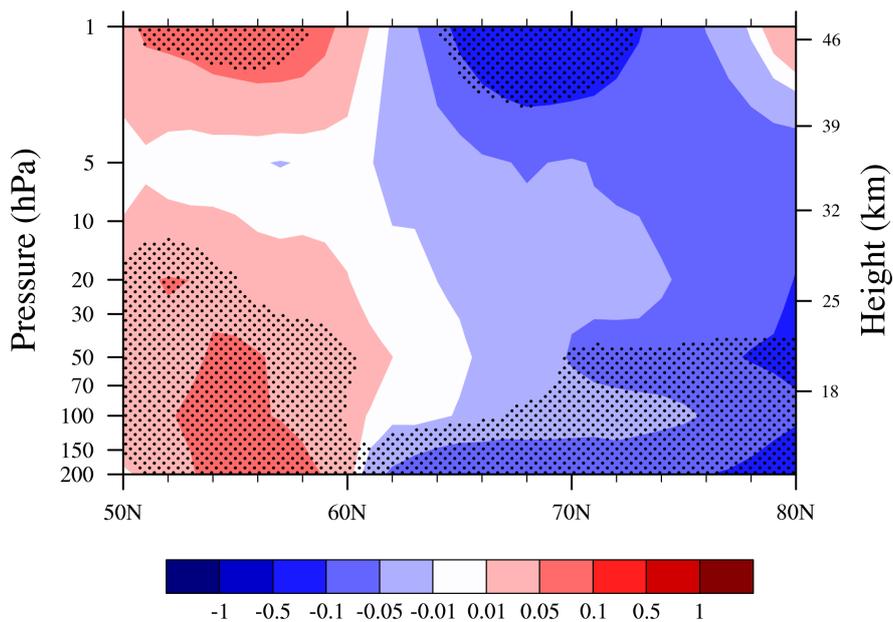
431



432

433 Figure 8. Zonal-mean zonal OGWD (units: m s⁻² Pa) integrated above 200 hPa averaged
434 in the period of 10 December 2023 to 10 January 2024 obtained from EXP_CTL (blue)
435 and EXP_NHE (red), along with the difference between the two experiments of
436 EXP_NHE and EXP_CTL (black).

437



438

439 Figure 9. Vertical distribution of zonal-mean temperature difference between the two
440 experiments of EXP_NHE and EXP_CTL (i.e., EXP_NHE minus EXP_CTL) averaged



441 in the period of 10 December 2023 to 10 January 2024 (units: K). Stippling denotes
442 differences statistically significant at the 95% confidence level.

443

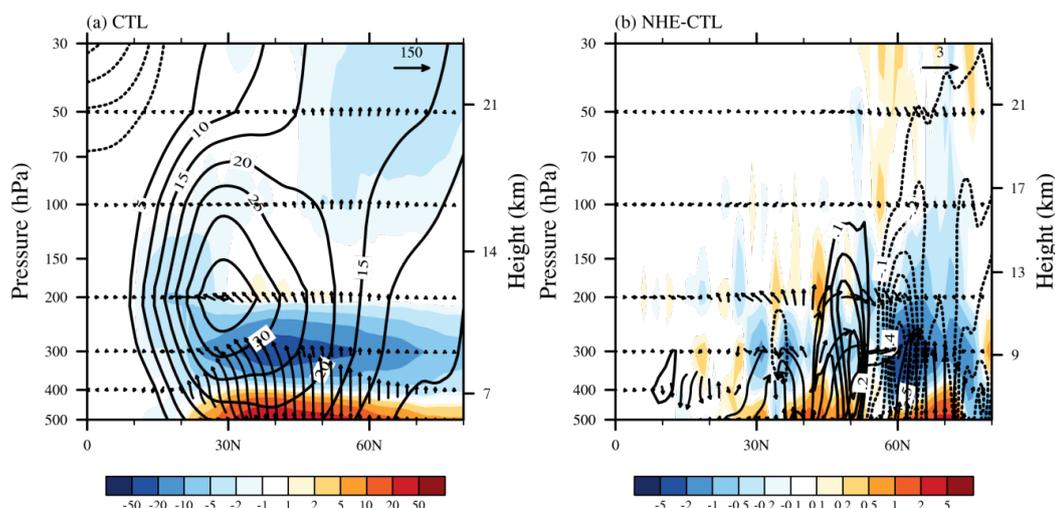
444 *c. Role of Rossby waves: wave-mean flow interaction*

445 While the polar stratospheric winds are increased, the polar tropospheric winds
446 are weakened (Fig. 1d). The zonal OGWD in the troposphere does not change as notable
447 as in the stratosphere (Fig. 6b). As shown below, the zonal-mean zonal winds in the
448 polar troposphere are mainly influenced by the changes in resolved large-scale Rossby
449 waves through the so-called catalytic wave-mean-flow positive feedback (White et al.,
450 2021).

451 It is well known that the impacts of large-scale Rossby waves on the mean flow
452 can be measured by the convergence of the zonal-mean Eliassen-Palm (EP) flux. Figure
453 10a illustrates the distribution of the zonal-mean EP flux in the EXP_CTL experiment
454 which is calculated following Edmon et al. (1980). In the NH middle-to-high latitudes,
455 Rossby waves originate from the lower troposphere which propagate upward and
456 converge in the upper troposphere. These waves split into two branches in the lower
457 stratosphere. The first branch turns to propagate equatorward across the tropospheric
458 jet. The other branch continues to propagate upward to the upper stratosphere where
459 the EP flux is converged and decelerates the mean flow. As in the SH, the Rossby waves
460 cannot propagate into the stratosphere (not shown) because of the presence of easterlies
461 there during austral summer (Fig. 1a).

462 Figure 10b is similar to Fig. 10a but gives the differences between the zonal-
463 mean EP fluxes in the two experiments (i.e., EXP_NHE minus EXP_CTL). Compared
464 to EXP_CTL, the upward propagation of Rossby waves into the stratosphere is

465 suppressed in the high latitudes of the NH, leading to an enhanced convergence of EP
466 flux in the troposphere. As a result, the zonal-mean zonal winds in the NH polar
467 troposphere are decelerated by the large-scale Rossby wave forcing, which contributes
468 to the alleviation of westerly biases there (Fig. 1c, 1d). This may explain why the
469 improvements in the GPH 500 forecast emerge after 6 days in EXP_NHE (Fig. 2). The
470 Rossby waves and their interaction with the mean flow require several days to develop
471 fully and influence the large-scale circulation.



472
473 Figure 10. Vertical distributions of zonal-mean EP flux (vectors) and its divergence
474 (shading; $\text{m s}^{-1} \text{ day}^{-1}$) due to resolved waves averaged in the period of 10 December
475 2023 to 10 January 2024 obtained from (a) EXP_CTL, and (b) the difference between
476 the two experiments of EXP_NHE and EXP_CTL (i.e., EXP_NHE minus EXP_CTL).
477 Contours are the corresponding zonal-mean zonal wind (units: m s^{-1}).
478

479 The suppressed upward propagation of Rossby waves can be understood from
480 the changes of the refractive index (*RFI*) that measures the ability of Rossby wave
481 propagation (e.g., Chen and Robinson, 1992; Hu et al., 2019), which is defined as

27

482
$$RFI = \left[\frac{\overline{q_\varphi}}{r\overline{U}} - \left(\frac{k}{r\cos\varphi} \right)^2 - \left(\frac{f}{2NH} \right)^2 \right] r^2, \quad (7)$$

483 where \overline{U} , k and H are the resolved zonal-mean zonal wind, zonal wave number and density
484 scale height, respectively. $\overline{q_\varphi}$ represents the meridional gradient of the potential vorticity,
485 which is expressed in the following form of

486
$$\overline{q_\varphi} = 2\Omega\cos\varphi - \left[\frac{(\overline{U}\cos\varphi)_\varphi}{a\cos\varphi} \right]_\varphi + \frac{af^2}{R_d} \left(\frac{p\theta}{T} \frac{\overline{U}_p}{\theta_p} \right)_p. \quad (8)$$

487 where Ω , θ and R_d being the Earth's angular frequency, potential temperature and dry gas
488 constant, respectively. The subscripts φ and p represent the partial derivatives with respect
489 to latitude and pressure, respectively. The overbars indicate temporal and zonal averages.
490 From Eq. (7), the increase of zonal wind will narrow the range of Rossby wave numbers
491 that can propagate into the stratosphere (Charney and Drazin, 1961; Xu et al., 2024).
492 As analyzed above, the weekend OGWD in the polar stratosphere enhances the zonal
493 winds there, thus inhabiting the upward propagation of Rossby waves.

494 **4. Summary and discussions**

495 The latest China Meteorological Administration Global Forecast System (CMA-
496 GFS) v4.0 model has been upgraded to a higher resolution of 0.125° in 2023. However,
497 this high-resolution global model still uses the parameterization scheme of orographic
498 gravity wave drag (OGWD) developed by Kim and Arakawa (1995; KA95) which is
499 based on hydrostatic orographic gravity wave (OGW) theory. In this study, the KA95
500 OGWD scheme is revised by taking into account the nonhydrostatic effect (NHE) on
501 the surface wave momentum flux (WMF) of OGWs, according to the nonhydrostatic



502 OGW theory derived in our earlier study of Xu et al. (2021). The performance of the
503 revised OGWD scheme is then evaluated for the medium-range forecast of the CMA-
504 GFSv4.0 model. Two sets of numerical experiments (i.e., EXP_CTL and EXP_NHE)
505 are conducted by using the original KA95 scheme and the revised NHE scheme,
506 respectively. In each numerical experiment, there are in total 31 forecasts of 10-day
507 forecasts which are initiated at 0000 UTC of each day in December 2023.

508 The results show that the revised OGWD scheme can improve the medium-
509 range forecast of large-scale circulation in the Northern Hemisphere (NH), especially
510 in the high latitudes. The easterly biases of zonal-mean zonal wind in the NH
511 stratosphere are reduced, with both the magnitude and location of the polar night jet
512 being better captured. The underestimation of the East Asia subtropical jet (EASJ) is
513 also alleviated. In contrast, the revised OGWD scheme shows little influence on the
514 stratospheric circulation in the Southern Hemisphere (SH). This is because, in boreal
515 winter (i.e., austral summer), there is hardly OGWD in the SH stratosphere owing to
516 the effect of critical-level absorption.

517 Quantitative evaluation is performed for the medium-range forecast skills in the
518 NH, taking the ERA5 reanalysis dataset as reference. Both experiments showed lower
519 geopotential height (GPH) at 500 hPa than ERA5 reanalysis, with the mean bias (MB)
520 and root mean square error (RMSE) increasing with forecast lead time. EXP_NHE
521 initially has greater MB and RMSE than EXP_CTL until the 6th forecast day, but these
522 metrics decrease more rapidly afterwards, resulting in an 11.59% reduction in MB and
523 a 5.06% reduction in RMSE by day 10. Over the total 10 forecast days, the MB and
524 RMSE of the GPH 500 in the NH are reduced by 2.6% and 0.52%, respectively,



525 indicating improved simulation of large-scale circulation in the NH. The improvement
526 is more noticeable in the high latitudes north of about 60°N, where the MB and RMSE
527 of GPH 500 are decreased by 31.18% and 15.93% at the 10th forecast day. Similar
528 results are found for the sea level pressure (SLP) in the NH, the MB and RMSE of
529 which are reduced by 34.29% and 8.33%, respectively, by day 10 in EXP_NHE.

530 The dynamics responsible for the improvement of large-scale wind circulation are
531 examined. The NHEs act to decrease the surface WMF of OGWs, which leads to a general
532 weakening of the parameterized OGWD in the high latitudes of the NH, especially in the
533 stratosphere. This weakening of stratospheric OGWD directly increases the zonal wind in
534 the polar stratosphere, reducing the easterly biases. It also induces a secondary circulation
535 according to the “downward control” principle, causing adiabatic cooling in the NH polar
536 stratosphere and warming to the south of 60°N. The resultant increase in thermal contrast
537 between the middle and high latitudes enhances the zonal wind in the stratosphere through
538 the thermal wind relation. The enhanced zonal wind in the polar stratosphere acts to narrow
539 the range of large-scale Rossby wave numbers that can propagate into the stratosphere
540 by reducing the refractive index of Rossby waves. The suppressed upward propagation of
541 Rossby waves leads to greater convergence of Eliassen-Palm (EP) flux (i.e., resolved wave
542 forcing) in the NH polar troposphere, which decelerates the zonal winds and thus alleviates
543 the westerly biases there.

544 To sum up, this study demonstrates that incorporating the NHE into the OGWD
545 parameterization scheme can help improve the simulation of large-scale atmospheric
546 circulation in high-resolution global NWP models, which is thus important for accurate
547 weather forecasts and/or climate simulations. However, there are still some limitations in



548 this work. For example, this study pays much attention to the medium-range forecast of
549 large-scale circulation in winter of NH. To further validate its effectiveness and robustness,
550 it is necessary to conduct more comprehensive evaluations of the nonhydrostatic OGWD
551 scheme in various numerical models focusing on different weather and climate systems.
552 Additionally, this work only takes into account the effects of NHE on the OGWD. Other
553 factors such as the vertical wind shear could also greatly affect the OGWD (e.g., Xu et al.,
554 2019; Xu et al., 2020; Zhang et al., 2025). While Xu et al. (2021) explored the
555 nonhydrostatic OGWs generated in a constant flow, it remains unclear how the vertical
556 wind shear affects the behavior of nonhydrostatic OGWs which needs further study.

557

558 *Code and data availability.* The China Meteorological Administration Global Forecast
559 System (CMA-GFS) model is archived on Zenodo under
560 <https://doi.org/10.5281/zenodo.18476721> (Zhang, 2026). The CMA-GFS outputs and the
561 ERA5 reanalysis data used for the comparison, along with the codes for orographic gravity
562 wave scheme are uploaded to <https://doi.org/10.5281/zenodo.18529537> (Zhang, 2026).

563

564 *Author contributions.* RRZ was responsible for formal analysis, writing the original draft
565 and visualization. ZZA carried out data analyses. XX developed the code and conducted
566 the analysis. HLX and QYC provided the resources and data curation.

567

568 *Competing interests.* The contact author has declared that none of the authors has any
569 competing interests.

570



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573

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578

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