

# 1 JUICE-MAJIS Earth observations during the 2024 gravity assist: first analysis 2 and comparison with PRISMA data

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20

## 21 **Abstract**

22

23 The *Jupiter ICy moons Explorer* spacecraft (JUICE) performed a Lunar-Earth gravity assist  
24 maneuver on 20th August 2024, during which the scientific instruments were turned on to  
25 test their functionality. In the time of the Earth flyby, the *Moon and Jupiter Imaging*  
26 *Spectrometer* (MAJIS) on board JUICE acquired a sequence of multispectral images over  
27 the Western Pacific Ocean at tropical latitudes. In parallel, an observing campaign was also  
28 conducted by the Earth-orbiting PRISMA imaging spectrometer, with the purpose of  
29 validating MAJIS spectral observations with independent measurements of the same kind.

30 These two datasets are here exploited to investigate and compare several atmospheric and  
31 cloud properties, including composition, temperatures, and atmospheric gravity waves. In the  
32 MAJIS spectral range, covering the 500-5560 nm wavelengths, we identified major and  
33 minor atmospheric gases, including O<sub>2</sub>, H<sub>2</sub>O, CO<sub>2</sub>, O<sub>3</sub>, CH<sub>4</sub>, N<sub>2</sub>O. Since MAJIS observations  
34 mostly covered diffuse cloudiness over the ocean, our analysis mainly focused on the  
35 discrimination of clouds' features and altitudes. We verified that ice particles are widespread  
36 in the data, allowing for an investigation of their properties (e.g. crystallinity) through different  
37 spectral signatures. The only land features identified in MAJIS data are not observed in  
38 daylight, hence only a thermal emission analysis is presented. Finally, the coverage of the  
39 4300 nm CO<sub>2</sub> band enables the identification of high altitude structures, revealing the  
40 presence of several atmospheric wave packets, likely induced by convective events, or  
41 lightning strikes known to have occurred at the time of the flyby. The present analysis  
42 demonstrates how MAJIS data can contribute to the scientific investigation of an  
43 atmospheric environment, and provide the first benchmark in the analysis of water ice,  
44 whose characterization in the Jovian system will be of primary importance for the JUICE  
45 mission.

46

## 47 **1. Introduction**

48

49 The JUICE mission is conceived for the investigation of Jupiter's icy satellites' surfaces and  
50 interiors, but also for the characterization of the giant planet's atmosphere and  
51 magnetosphere. These scientific objectives will be achieved thanks to a payload consisting  
52 of several remote sensing and in-situ instruments, including an altimeter, a magnetometer, a  
53 gravity experiment, a radio instrument, neutral/energetic particles and plasma detectors and  
54 an ultraviolet spectrograph. Moreover, the visible-thermal infrared spectral range will be  
55 investigated by a visible camera (JANUS) and by the *Moons and Jupiter Imaging*  
56 *Spectrometer* (MAJIS, Poulet et al., 2024a), which in particular will allow the spectroscopic  
57 investigation of Jupiter's atmosphere, moons and rings system.

58 On the 20th of August 2024 the JUICE spacecraft performed a *Lunar-Earth Gravity Assist*  
59 (LEGA) and is now headed for a second *Earth Gravity Assist* (EGA) happening in September  
60 2026. In this study we will focus on 2024 EGA data acquired by MAJIS which, along with  
61 JANUS (see Hueso et al., this issue), was turned on providing its very first observations of a  
62 planetary target (for a general overview of the flyby refer to Poulet et al., this issue, while  
63 valuable information about MAJIS operations, functioning and performances is given in  
64 Langevin et al. and Seignovert et al., this issue). During the flyby, different Earth observing  
65 spectrometers were coordinated to provide spatially and temporally comparable  
66 observations (Poulet et al., this issue). A companion paper by Guerlet et al. (this issue) is  
67 focused on MAJIS IR channel's data comparison with co-located acquisitions by the IASI  
68 thermal Fourier spectrometer onboard the EUMETSAT/Metop satellite. Instead, we exploit  
69 PRISMA spectrometer data as a proxy to compare with MAJIS VISNIR channel observations  
70 (Sect. 2.1), even if the different times and regions of acquisition prevent a direct comparison  
71 of the scans (see Sect. 2). PRISMA (Sect. 2.2) is a technology demonstrator mission  
72 completely funded by the Italian Space Agency (ASI) and devoted to the qualification of a  
73 panchromatic/hyperspectral technology for monitoring the Earth at visual-near infrared  
74 wavelengths at moderate spectral resolution and high spatial resolution (Pignatti et al.,  
75 2013).

76 During the EGA, JUICE flew over Western Pacific Ocean at tropical latitudes, moving  
77 approximately from Sumatra to Hawaii islands and spanning local times from about 03:00 to  
78 10:30 (see Table 1). The majority of these measurements took place over the ocean,  
79 allowing a broad characterization of atmospheric gaseous composition and structure (Sect.  
80 2.3.2). Land features are only marginally detected in a couple of observations mainly in the  
81 thermal range (Sect. 4.4).

82 Given the widespread presence of clouds and the early local times of acquisition (Sect. 2),  
83 ice is observed in almost all MAJIS scans (Sect. 4.1), allowing benchmarking of the  
84 spectrometer's response to this observable in view of Jupiter's icy satellites investigation.  
85 Also atmospheric waves (whose role is fundamental in regulating the middle-atmosphere  
86 circulation, e.g. Hamilton, 1996; Fritts and Alexander, 2003) are detected in many MAJIS  
87 observations. Such phenomena are either linked with orography (Queney 1948; Kim et al.  
88 2003) or with the occurrence of thunderstorms (Taylor and Hapgood, 1988; Dewan et al.  
89 1998). Given that MAJIS EGA observations mainly targeted the ocean, we investigate the  
90 waves' connection with strong convective events or lightning strikes (Sect. 4.4).

91 The manuscript is arranged in sections describing the data (Sect. 2), the methods for their  
92 investigation (Sect. 3) and the obtained results (Sect. 4). Such a wide ensemble of

93 atmospheric observable features is finally discussed in the context of Jupiter science in Sect.  
94 5.

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## 96 2. Observations

97

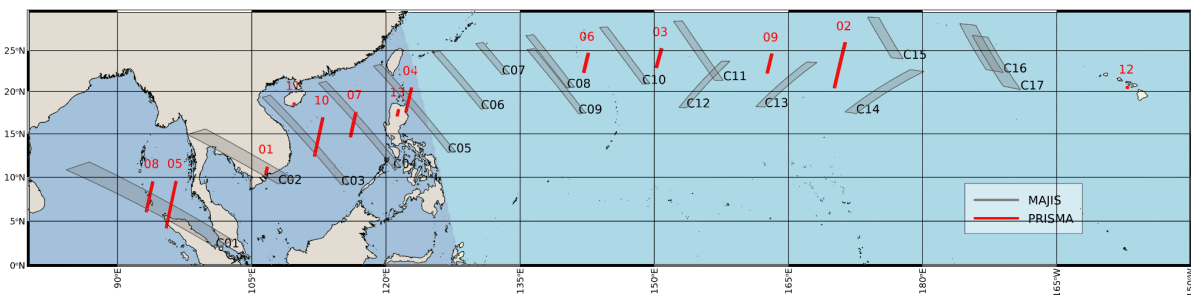
### 98 2.1. MAJIS EGA Data

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100 MAJIS is a dispersion grating imaging spectrometer operating between 500 and 5560 nm by  
101 means of two spectral channels (Poulet et al., 2024a). The first channel (VISNIR, 500–2350  
102 nm) is characterized by nominal spectral resolution and sampling of 2.9-4.6 nm and 3.5-3.8  
103 nm/band respectively, while the second (IR, 2270–5560 nm) works with a spectral resolution  
104 of 5.5-7.0 nm and a sampling of 5.9-6.9 nm/band. The nominal instrument’s instantaneous  
105 field of view (IFOV) is 150  $\mu$ rad/pixel. MAJIS concept has been optimized for the  
106 characterization of the surface and near-surface environment of Jupiter’s icy moons (Poulet  
107 et al., 2024a), as well as for the investigation of Jupiter’s atmosphere (Fletcher et al., 2023).  
108 Detailed descriptions of the instrument functioning, operations and calibration are given in  
109 Haffoud et al. (2024), Langevin et al. (2024), Poulet et al. (2024b), Filacchione et al. (2024),  
110 Rodriguez et al. (2024), Vincendon et al. (2024), and Stefani et al.(2025). Scene geometry is  
111 reconstructed via the SPICE-NAIF toolkit (Acton, 1996; Acton et al., 2018) and kernels  
112 provided by ESA (“JUICE SPICE Kernel Dataset,” 2019).

113 Figure 1 and Table 1 summarize footprint locations and main basic properties of the 17  
114 MAJIS EGA data investigated in this work (see Poulet et al., this issue, for further  
115 instrumental parameters). Two additional cubes, targeted off-limb for calibration purposes  
116 (Poulet et al., this issue), are not considered here. Each MAJIS acquisition consists of  
117 hyperspectral *cubes* (i.e. 2D spatial frames with a third spectral dimension) collected as  
118 pushbroom spectral scans via internal mirror rotation, with different widths and lengths.

119



120

121 **Figure 1:** Geographical coverage of the investigated observations, MAJIS in grey color,  
122 PRISMA in red color. The darker area westward of the Philippines indicates the nightside at  
123 the time of the terminator crossing of MAJIS observations (2024-Aug-20 21:30 UTC). Instead  
124 all PRISMA footprints are in daylight, at local time  $\sim$ 10:30. Coastlines are from  
125 OpenStreetMap, available under the Open Database License.

126

127 The first 4 cubes (C1 to C4) pointed to the Earth surface at nighttime and contain a  
128 significant signal only in the thermal part of the spectrum ( $\lambda > 3000$  nm). The only exception  
129 is C1, where a lightning emission is identifiable at visible wavelengths (D’Aversa et al., this

130 issue). C5 is straddling the terminator and is the first cube containing information on the  
 131 dayside ocean and clouds. Some coastlines are identifiable in C4 and C5 at thermal  
 132 wavelengths, as it will be discussed in Sect. 4.4. All the subsequent cubes (C6 to C17) are  
 133 acquired in daylight and hence the full spectrum can be investigated, even if they only cover  
 134 the ocean mostly under cloudy/stormy conditions.

135 Cubes C11, C12 and C13 have been acquired with longer integration times, with the purpose  
 136 of testing the instrument response. This leads to signal saturation in many regions  
 137 (especially at visual wavelengths over clouds, see Sect. 2.3.1), that have been removed from  
 138 our analysis. The spatial resolution in this dataset is quite stable (about 1.4 km per pixel,  
 139 slightly affected by motion smearing) and is suited for the investigation of both homogeneous  
 140 and localized cloud structures. On the other hand, the IFOV is affected by unresolved  
 141 cloudiness (likely widespread) which dilutes the low reflectivity of deep water hence  
 142 preventing the acquisition of clear-sky ocean (Sect. 2.3.1, Fig. 2).

143

144 **Table 1:** MAJIS observing parameters during EGA. Phase angle is always close to 90°.

	ID	target	incidence angle (°)	emission angle (°)	local time (h)	instantaneous resolution (km/px)
C1	20240820212509	surface night	115-130	28-42	03:00 – 04:18	1.80
C2	20240820212818	surface night	106-116	17-27	03:54 – 04:48	1.55
C3	20240820213029	surface night	100-106	12-19	04:30 – 05:12	1.50
C4	20240820213208	surface night	93-100	6-13	05:00 – 05:36	1.45
C5	20240820213347	surface terminator	87-93	0-11	05:24 – 06:00	1.40
C6	20240820213530	surface day	82-87	6-11	05:54 – 06:18	1.35
C7	20240820213644	surface day	79-82	11-14	06:12 – 06:30	1.30
C8	20240820213731	surface day	72-77	17-20	06:36 – 07:00	1.30
C9	20240820213840	surface day	71-76	14-20	06:36 – 07:06	1.30
C10	20240820214003	surface day	64-69	24-27	07:12 – 07:36	1.30
C11	20240820214117	surface day	56-61	32-37	07:48 – 08:12	1.30
C12	20240820214231	surface day	55-60	29-34	07:48 – 08:12	1.25
C13	20240820214350	surface day	46-52	39-45	08:24 – 08:54	1.30
C14	20240820214509	surface day	34-42	50-58	09:06 – 09:42	1.35
C15	20240820214628	surface day	36-41	49-53	09:18 – 09:36	1.30
C16	20240820214720	surface day	26-32	60-65	10:00 – 10:18	1.40
C17	20240820214813	surface day	23-31	60-66	10:06 – 10:30	1.40

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146

## 2.2. PRISMA data

147

148 An observing campaign coordinated to the EGA was conducted by the mission PRISMA  
 149 (*PR*ecursore *I*perSpettrale della *M*issione *A*pplicativa), managed by the Italian Space  
 150 Agency. The mission hosts a visible and near-infrared imaging spectrometer, covering a  
 151 range (400-2500 nm) compatible with the MAJIS-VISNIR channel but having a coarser  
 152 spectral resolution (~12 nm) in turn compensated by a higher spatial resolution (~30  
 153 m/pixel). Details about the instrument and the mission can be found in Pignatti et al. (2013),  
 154 while mission characteristics, access, products, calibration, geometry navigation and data  
 155 policy are fully described in Lopinto et al. (2021).

156 PRISMA sequences (13 in total, red rectangles in Fig. 1, main parameters summarized in  
 157 Table 2) consist of a variable number of 30 x 30 km hyperspectral cubes, each composed of  
 158 1000 x 1000 spatial pixels. Due to the PRISMA orbit (Sun-Synchronous-Low-Earth-Orbit),  
 159 observations are acquired at a fixed solar local time (~10:30), making it impossible to  
 160 achieve spatial/temporal coincidence with MAJIS ones (see next section).

161

162 **Table 2:** PRISMA observations acquired in coordination with JUICE.

PRISMA sequence	Num cubes	Start UTC	Solar zenith angle (°)	Emission angle (°)	Cloud coverage (%)	$\Delta t$ (h) (PRISMA-MAJIS)
01	3	2024-08-17 03:34	20.3	14.6	14	-90.9
02	21	2024-08-18 23:13	23.4	17.1	8	-46.45
03	9	2024-08-19 00:50	20.6	20.7	9	-44.83
04	11	2024-08-19 02:22	21.4	4.2	100	-43.19
05	21	2024-08-19 04:08	22.7	1.2	73	-41.52
06	9	2024-08-20 01:07	23.6	16.6	2	-20.55
07	11	2024-08-20 02:46	22.5	18.4	18	-18.90
08	13	2024-08-20 04:25	21.0	16.0	98	-17.24
09	9	2024-08-20 23:46	23.2	12.3	1	2.11
10	17	2024-08-21 03:03	22.0	12.0	5	5.39
11	3	2024-08-22 03:19	21.7	5.6	20	29.66
12	1	2024-08-22 21:07	22.2	7.6	16	47.45
13	3	2024-08-25 02:32	21.6	4.3	7	100.88

163

### 164 2.3. General comparison overview

165

166 Both MAJIS and PRISMA acquired multispectral data covering the same kinds of structures,  
 167 offering a useful benchmark for checking MAJIS capabilities in detecting and analyzing  
 168 specific features of scientific interest. In the following section we investigate how the spectral  
 169 signatures of the main atmospheric gases and of clouds are affected by the different  
 170 spatial/spectral resolutions and observing conditions. When reflectances are discussed,  
 171 they are obtained for both instruments as  $I/F$ , where  $I$  is the observed spectral radiance and

172 *F* is Kurucz solar spectral radiance (Kurucz et al., 1984; Kurucz, 1995) available at the  
 173 website <https://earth.gsfc.nasa.gov/climate/projects/solar-irradiance/data> (accessed  
 174 February, 10, 2026). For readability, in the following we will refer to spectral radiance simply  
 175 as *radiance*.

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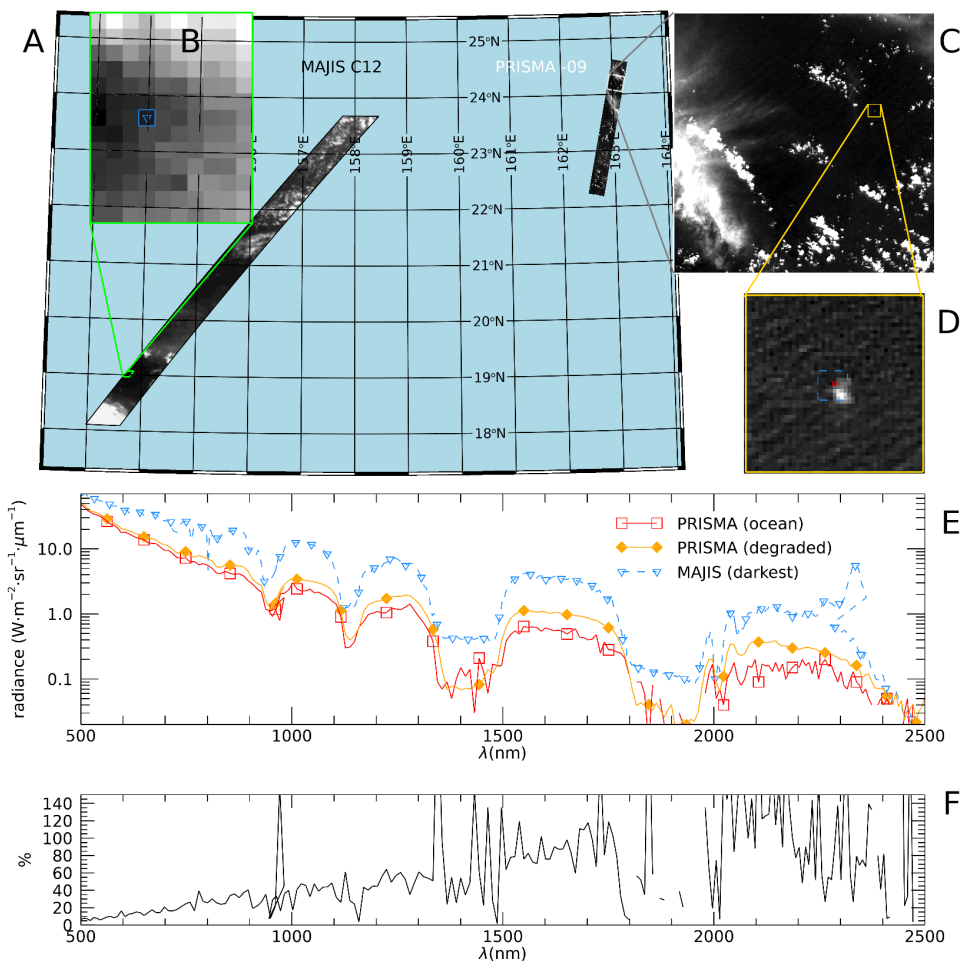
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### 2.3.1. Ocean/clouds spectra first comparison

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179 Figure 2 shows the two PRISMA and MAJIS cubes that are closest from both a spatial and  
 180 temporal point of view (~550 km and ~2 h apart), covering open ocean areas overlaid by a  
 181 different amount of clouds. In this framework, the most robust radiance comparison should  
 182 consider ocean cloud-free spectra, expected to be quite stable in space and time and very  
 183 dark at visual wavelengths (given the very low ocean albedo, ~4%). However, the  
 184 comparison between the two instruments (Fig. 2E) highlights that the darkest MAJIS signals  
 185 are still brighter than those from PRISMA, possibly suggesting enhanced cloud/aerosol  
 186 content. Indeed, the higher spatial resolution of PRISMA data reveals a number of  
 187 small-scale structures, likely unresolved by MAJIS, yet affecting its signal. For instance, the  
 188 small bright feature imaged by PRISMA in Fig. 2D, covering only a portion of a MAJIS pixel  
 189 footprint, may induce spectral variations of the ocean spectrum up to 50% (Fig. 2F) once  
 190 observed at the MAJIS resolution scale.

191



192

193 **Figure 2: A)** MAJIS observation C12 and PRISMA sequence P09 (~2h apart) shown at 875  
194 nm in an equal-area projection. **B)** Blow-up of the darkest area in the MAJIS image,  
195 highlighting individual pixels' size (the pixel with the lowest signal is highlighted in blue). **C)**  
196 The second cube of the PRISMA sequence is shown in its full extension of 1000x1000  
197 pixels. **D)** Blow-up of an area of PRISMA data encompassing a small bright cloud. The blue  
198 dashed box shows the approximate size of a MAJIS pixel. **E)** Single-pixel spectra from the  
199 darkest pixels of MAJIS (blue color, triangle symbol in B) and PRISMA (red curve, red  
200 square in D). The orange curve represents a PRISMA spectrum degraded to MAJIS spatial  
201 resolution (average inside the blue box of panel D). The MAJIS spectrum is multiplied by the  
202 ratio of solar incidence cosines ( $=1.82$ ) to achieve a radiance level comparable with  
203 PRISMA. **F)** Spectral effect of the spatial degradation in PRISMA data, shown as the relative  
204 difference between the red and orange curves of panel E.

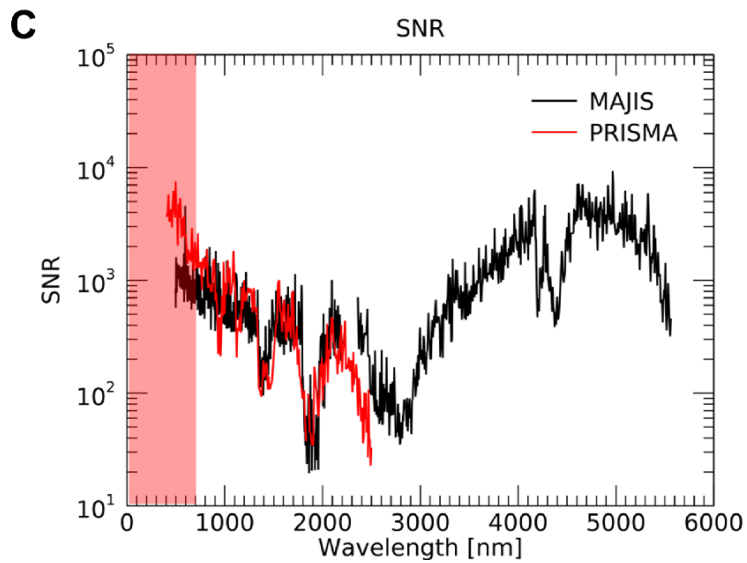
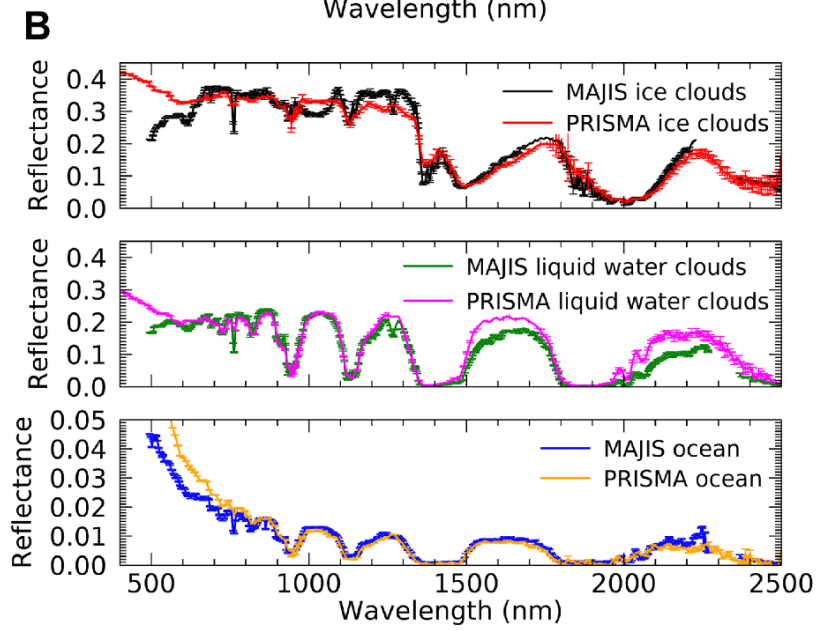
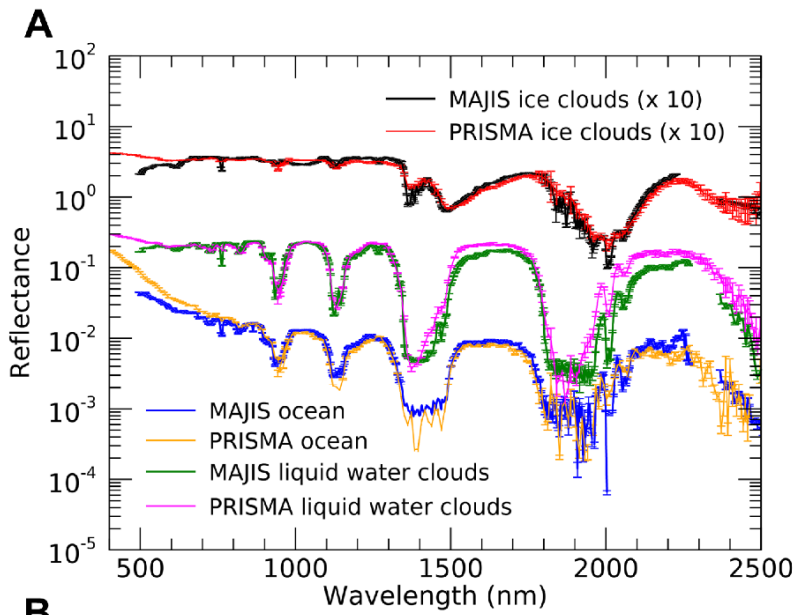
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206 Most of the spectral variability in both datasets is driven by changes in the H<sub>2</sub>O absorption  
207 bands. Besides the general low reflectivity, ocean spectra are characterized by the presence  
208 of large and often saturated water absorption bands. On the other hand, H<sub>2</sub>O clouds (either  
209 composed of liquid droplets, ice crystals or a mixture) can easily be identified through RGB  
210 imaging from both datasets due to their bright appearance (Sect. 3.1). H<sub>2</sub>O bands are less  
211 saturated over clouds, where light scattering prevents photons from reaching the  
212 underneath, more absorbing, atmospheric layers. Ice clouds' discrimination is basically  
213 driven by the spectral shift of absorption bands between solid and liquid H<sub>2</sub>O phase (Sect.  
214 3.1). The comparison of spectral signatures related to the ocean and clouds (main spectral  
215 endmembers for both instruments) is shown in Fig. 3A-B in log and linear scale respectively  
216 (refer to Fig. 4 and 5 for the gaseous features identification). This comparison should be  
217 considered as qualitative, since spectra acquired at different locations, geometries and local  
218 times are being considered (see Tables 1 and 2). Therefore, clouds are likely characterized  
219 by different vertical distributions and microphysical properties, driven by a radiative forcing  
220 that is changing between early and mid-morning. Also, differential sun-glint effects  
221 (dependent on geometry and wind strength) could produce differences in the overall  
222 reflectivity of the ocean (Cox and Munk, 1954). All these effects (straylight could also have  
223 an impact here, see Langevin et al, this issue) are likely to contribute to non-linear offsets in  
224 the continuum below about 700 nm (e.g. Zinner et al., 2016), and slightly different depth and  
225 shape of water absorption bands, not ascribable solely to differences in spectral resolution.

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228



230 **Figure 3:** comparison between MAJIS and PRISMA reflectances in log (A) and linear (B)  
231 scales related to ocean, liquid water clouds and ice water clouds (the latter multiplied by 10  
232 for clarity in panel A). PRISMA spectra are selected from two orbits in session 7, MAJIS  
233 ones from orbits C7 (ice clouds) and C10 (ocean and liquid water clouds). Panel C shows  
234 the SNR estimated for the two instruments (cube C15 for MAJIS, one cube of session 07 for  
235 PRISMA) as described in Sect. 2.3.1. The red shaded area indicates the spectral region  
236 possibly affected by straylight contamination, not yet fully assessed in both datasets.

237

238 The three endmembers in Fig. 3 show similar trends in reflectivity, with the main absorption  
239 bands' shape correctly reproduced, even if the probed atmospheric structure is probably not  
240 the same. For example, MAJIS liquid water clouds spectrum shows wider wings and a flatter  
241 bottom for the bands at 1400 and 1900 nm, suggesting different scattering properties in the  
242 atmospheric column for the two cases (see Sect. 4.2.4). A slightly flatter bands' bottom is  
243 also observed in the ocean spectrum (blue compared to the orange PRISMA spectrum). On  
244 one side, this could indicate that early-morning thin clouds in the mid-high troposphere are  
245 mixed in MAJIS footprint, preventing the formation of the narrower water lines inside the  
246 bands (MAJIS spectrum refers to 7:30 local time, when the presence of unresolved hazes is  
247 likely). On the other hand, such low signals could reach the instrument noise equivalent  
248 spectral radiance (NESR), hence explaining the featureless bands' bottom. We derived an  
249 upper limit for the NESR by investigating the darkest ocean region in the selected MAJIS  
250 cube (C10), resulting in about  $10^{-3}$  W/(m<sup>2</sup> μm sr) at 1900 nm. This value corresponds to  
251 reflectances of  $10^{-4}$ , about one order of magnitude below the ocean signal at that wavelength  
252 (Fig. 3A), hence making the mixed-footprint hypothesis more likely. The occurrence of  
253 saturation in some parts of MAJIS spectrum is highlighted in the ice clouds comparison,  
254 evident as a broad absorption between 900 and 1100 nm in Fig. 3B. MAJIS uncertainties are  
255 extensively discussed in the paper by Poulet et al. (this issue), but here we attempt an a  
256 *posteriori* estimation of the spectral signal to noise ratios (SNR) for both instruments by  
257 performing a statistical analysis of spatial fluctuations computed in 5x5 pixels boxes (Fig.  
258 3C). For each wavelength (excluding saturated regions) we select those regions producing  
259 the minimum relative error, hence representing both noise statistics and true variations in the  
260 observed scene. As a result, the spectral SNRs in Fig. 3C refer to wavelength-dependent  
261 locations in the respective cubes, rather than to a single region. This means that the high  
262 frequency oscillations in the red and black lines are mostly driven by spatial differences  
263 between the selected boxes (at the scale covered by the respective cubes). Values below  
264 ~700 nm (red shaded area in Fig. 3C) could be driven by differences in clouds/aerosols'  
265 properties (in turn impacting the intensity of Rayleigh scattering), and are possibly  
266 contaminated by the presence of straylight affecting the actual trend of the SNR for both  
267 instruments.

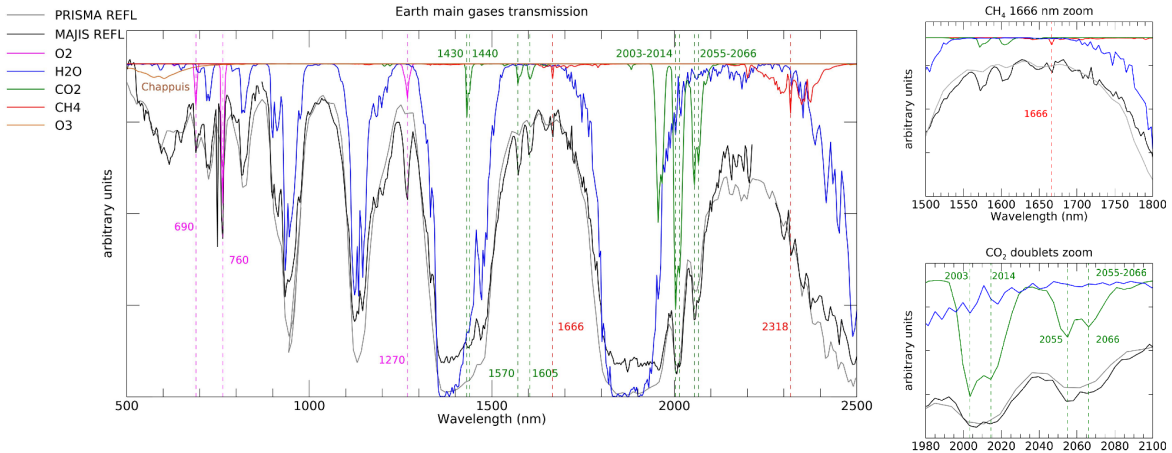
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### 2.3.2. Gaseous compounds

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271 Figure 4 compares sample MAJIS/PRISMA liquid water clouds reflectance spectra with  
272 two-way vertical transmission due to O<sub>2</sub>, H<sub>2</sub>O, CO<sub>2</sub>, CH<sub>4</sub>, N<sub>2</sub>O and O<sub>3</sub>, based on an average  
273 vertical structure from Efremenko and Kokahnovsky (2021) and calculated through the  
274 line-by-line method with line parameters from the HITRAN database (Gordon et al., 2022)  
275 and from O<sub>3</sub> cross sections by Gorshchev et al. (2014) and Serdyuchenko et al. (2014).  
276 Finally, transmissions are convolved at the MAJIS spectral resolution.



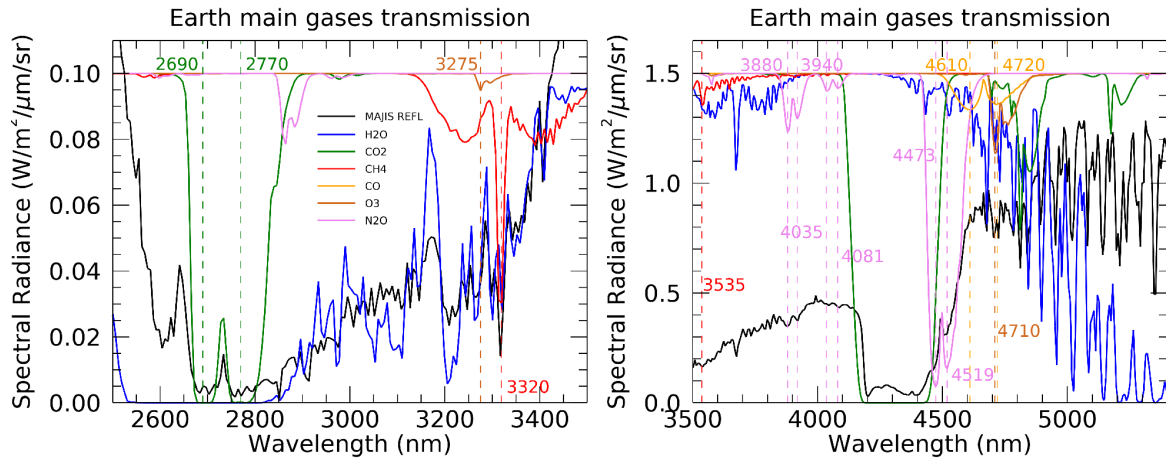
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278 **Figure 4:** MAJIS (black, taken from C7) and PRISMA (grey, taken from session 7)  
 279 normalized reflectances (both pertaining to liquid water cloud scenarios) compared to main  
 280 Earth's atmospheric gases two-way transmissions convolved on MAJIS spectral grid. Vertical  
 281 dashed lines indicate the main non-H<sub>2</sub>O molecular lines identifiable in the observations.  
 282 Zooms related to the CH<sub>4</sub> 1666 nm absorption line, and the CO<sub>2</sub> doublets at 2003-2014 nm /  
 283 2055-2066 nm are shown in the upper and lower right panels respectively.

284

285 In their common spectral range, both instruments allow to identify the main absorption  
 286 features of H<sub>2</sub>O, O<sub>2</sub> and CO<sub>2</sub> (Fig. 4, see also Poulet et al., this issue). The reduced spectral  
 287 resolution of PRISMA makes it difficult to resolve narrow features like the methane  
 288 absorption at 1666 nm (Fig. 4, upper right panel), the close doublets of CO<sub>2</sub> at 2003-2014  
 289 nm and 2055-2066 nm (Fig. 4, lower right panel), or shallower lines of water. On the other  
 290 hand, the PRISMA spatial resolution is expected to reduce the spatial mixing of different  
 291 types of surfaces or aerosols, allowing a more robust tracking of localized and transient  
 292 phenomena (e.g. smog layers, ice patches, oil spills, CO<sub>2</sub> emissions, etc.). At wavelengths  
 293 around 600 nm a broad absorption possibly matching the O<sub>3</sub> Chappuis band appears in both  
 294 datasets. In MAJIS, this is enhanced over thick clouds and in particular in grazing  
 295 illumination conditions (Sect. 4.4) in which the atmospheric column above ~20 km is directly  
 296 illuminated resulting in a very long photon path length that increases the absorption from O<sub>3</sub>  
 297 in the scattered light (most of terrestrial ozone resides between altitudes of 20 and 40 km).  
 298 Nevertheless, a better quantification of this feature requires a more rigorous assessment of  
 299 the straylight contamination (Langevin et al., this issue).

300 Besides the better spectral resolution, MAJIS also has the advantage of an extended  
 301 spectral range covering wavelengths from 2500 nm up to 5560 nm. In this range, thermal  
 302 emission dominates and provides information on the temperature of the sampled  
 303 atmospheric layers, or of the ocean and clouds. This interval is characterized by several H<sub>2</sub>O  
 304 absorption bands (the stronger one centered at about 2700 nm), strong and saturated CO<sub>2</sub>  
 305 ones at 2690, 2770 and 4300 nm, and weaker CH<sub>4</sub>, O<sub>3</sub>, CO and N<sub>2</sub>O signatures (Fig. 5, see  
 306 also Guerlet et al., this issue). In particular, the strong CO<sub>2</sub> absorption (and emission) at  
 307 4300 nm, can be used for the estimation of the vertical structure of atmospheric  
 308 temperatures (see Poulet et al., this issue).



309

310 **Figure 5:** MAJIS (black) spectral radiance compared to main Earth's atmospheric gases  
 311 two-way spectral transmissions (offset for clarity) in the 2500 - 3500 nm range (left) and  
 312 3500-5400 nm range (right). Thermal emission is not considered in the transmission  
 313 computation and all spectra are convolved to the MAJIS spectral grid.

314

### 315 3. Methods

316

317 In this section we describe different methods for investigating the information content in the  
 318 data, including surface/cloud features identification (Sect. 3.1), ice characterization (Sect.  
 319 3.2), clouds' altitude estimation (Sect. 3.3) and high-altitude features investigation (Sect.  
 320 3.4).

321

#### 322 3.1. Surface and clouds identification

323

324 In principle, Earth observations can encompass different types of surfaces, commonly  
 325 discriminated spectrally through indices expressed in the formalism of *Normalized Difference*  
 326 *spectral Indices* (NDIs, see Wolf, 2012 for a general review). Useful examples are given in  
 327 Hurley et al., 2014 (dealing with Rosetta/VIRTIS-M data, Coradini et al., 1999) and in Oliva  
 328 et al., 2017 (dealing with both Rosetta and Venus Express/VIRTIS-M data, Drossart et al.,  
 329 2004). Table 3 summarizes these indices (derived from spectral endmembers from  
 330 Rosetta/VIRTIS-M acquisitions, Fig. 6A-B, since MAJIS observations did not cover surface  
 331 features in daylight) that we test on PRISMA data (Fig. 6C-D) as a benchmark for the future  
 332 September 2026 EGA, in which Africa observations are planned. A new ocean index is also  
 333 defined specifically for MAJIS data, which do not cover all wavelengths of the nominal ocean  
 334 NDI. It is worth stressing that the ocean class should not be considered as representative of  
 335 clear-sky conditions as it may actually include some amount of aerosol opacity (Sect. 2.3.1).  
 336 No specific index has been adopted for generic clouds identification, but we rather assign to  
 337 this class all pixels that do not meet any of the surface classes' conditions. Indices  
 338 thresholds can be studied taking advantage of proxy images (e.g. the PRISMA one shown in  
 339 Fig. 6C, not pertaining to EGA sequence) in which the changing reflecting structures can be  
 340 clearly identified. The derived values depend on instrument features and require specific  
 341 tuning when switching between different datasets. Fig. 6C-D shows how the different types  
 342 of spectral classes can be reliably identified, even if, in this case, no ice clouds are present.  
 343 Other examples of application of the ocean, clouds and ice indices from Table 3 to MAJIS

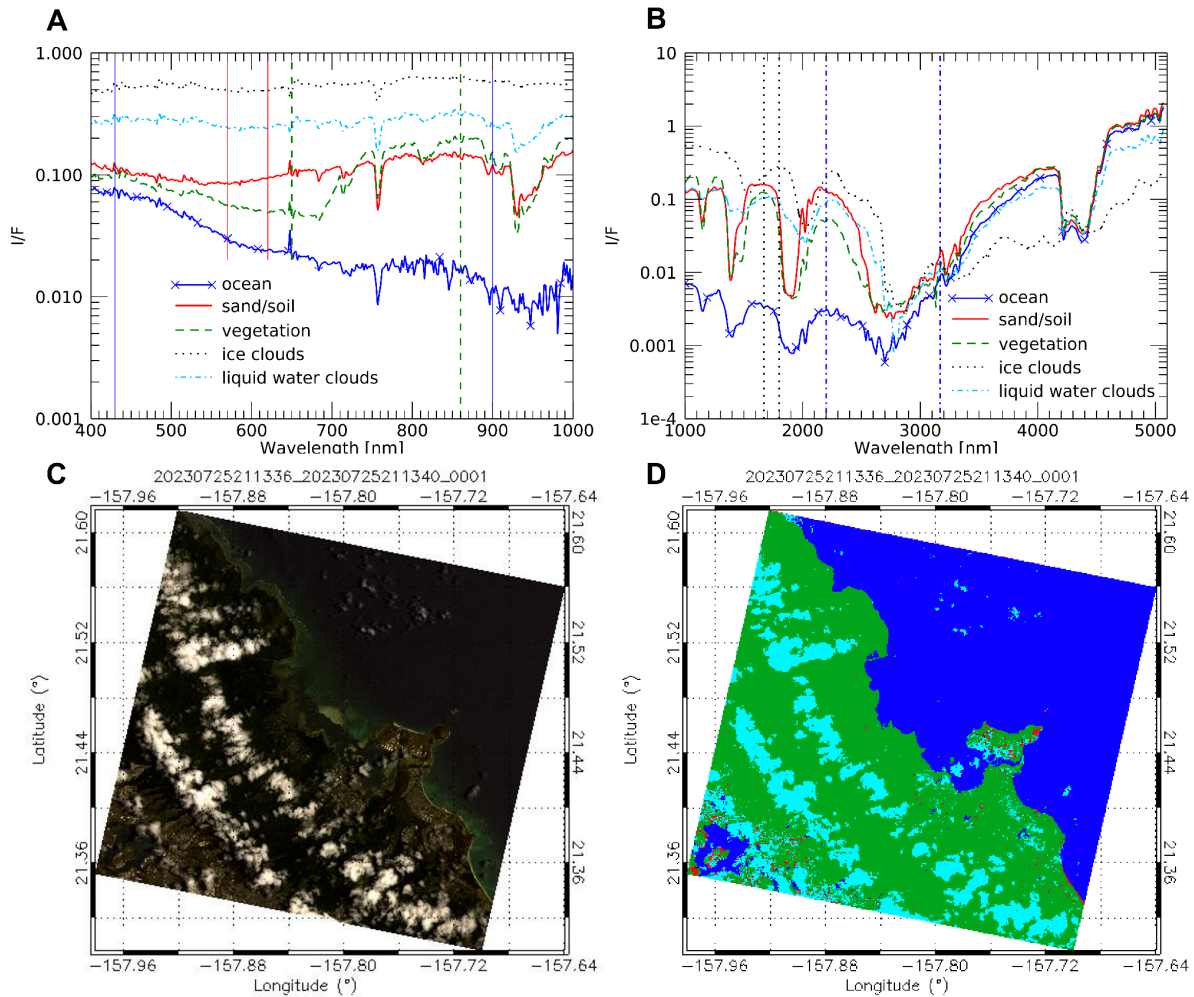
344 and PRISMA data are discussed in Sect. 4.1. Instead, the application of surface-related  
 345 indices to MAJIS data did not result in positive identification, since land features in MAJIS  
 346 data are not seen in daylight illumination, making NDIs not applicable.

347

SPECTRAL CLASS	SPECTRAL INDEX	SPECTRAL SIGNATURE	FIGURE
Vegetation	$NDVI: \frac{R_{860} - R_{650}}{R_{860} + R_{650}}$	Chlorophyll absorption in the red band	6D
Sand/Soil	$NDSI: \frac{R_{570} - R_{620}}{R_{570} + R_{620}}$	Enhanced contrast between the red and green bands	6D
Ocean	$NDWI: \frac{R_{430} - R_{900}}{R_{430} + R_{900}}$	Enhanced reflectivity in the blue with respect to NIR wavelengths	6D - 12D
<i>MAJIS Ocean</i>	$\frac{R_{2200}}{R_{3170}}$	Low solar I/F, large thermal emission	12C
Ice Clouds	$\frac{R_{1670}}{R_{1800}}$	Shift of the 1500 nm H <sub>2</sub> O ice absorption band to longer wavelengths with respect to the liquid phase (different arrangement of hydrogen bonds)	12C - 12D
Cloudy	pixels not assigned to surface types	/	6D - 12C - 12D

348

349 **Table 3.** Spectral indices for the identification of different spectral classes related to surfaces  
 350 and clouds. *R* indicates I/F and the subscript is the wavelength in nanometers.



351

352

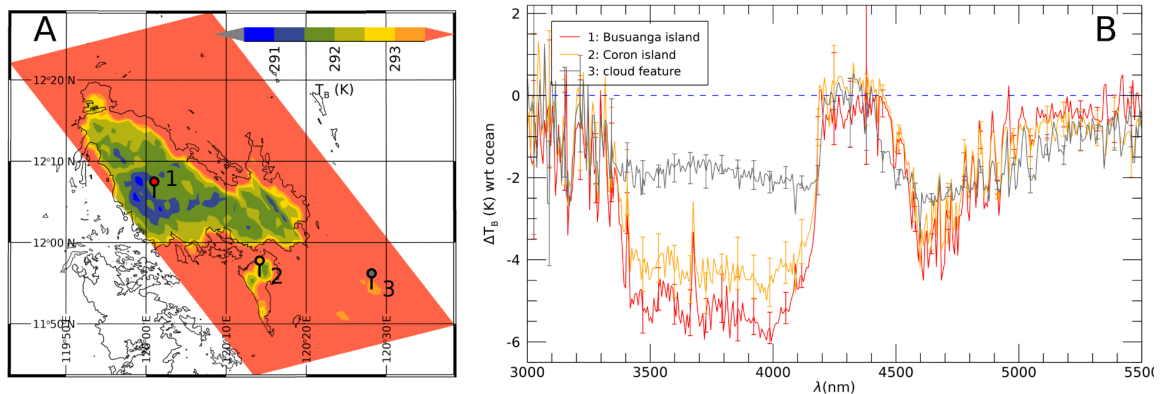
353 **Figure 6:** **A:** Reflectance endmembers of different classes of surface and clouds, derived  
 354 from Rosetta/VIRTIS-M VIS channel Earth observations (Oliva et al., 2017). Vertical lines  
 355 share the style and color of the corresponding spectral endmember and identify the  
 356 wavelengths adopted in the index definition (blue ones refer to the NDWI). **B:** same as in A  
 357 but I/F spectra from the NIR channel of VIRTIS-M are shown (vertical blue dashed-dotted  
 358 lines refer to the MAJIS ocean index). **C:** Example of a PRISMA RGB image covering  
 359 different surface types (data cube 20230725211336\_20230725211340\_0001) targeting the eastern  
 360 coastal line of Honolulu island ( $R=680$  nm;  $G=570$  nm,  $B=440$  nm). **D:** distribution of spectral  
 361 classes obtained from the spectral indices in Table 3. Green pixels indicate vegetation, red  
 362 ones are sand, cyan ones are clouds, blue ones indicate ocean/water (no ice clouds  
 363 present).

364

365 In the specific conditions of MAJIS EGA sequence, the most robust land identification must  
 366 rely on soil/ocean contrast in thermal emission (Sect. 4.4), triggered by the different thermal  
 367 inertia of the two classes. However, also the presence of clouds in the line of sight induces a  
 368 decrease of the observed brightness temperature ( $T_B$ ), hence land identification requires  
 369 matching the shapes of low  $T_B$  regions within known coastlines. The largest land region  
 370 emerging in this way is shown in Fig. 7 (cube C4), (Philippines's Busuanga and Coron  
 371 islands in cube C4), whose identification also allows a refinement of MAJIS pointing

372 reconstruction (Seignovert et al., this issue). The largest brightness temperature contrast for  
 373 both land/ocean and cloud/ocean cases occurs in the 3500-4000 nm and 4600-4800 nm  
 374 spectral ranges, which are less absorbed by atmospheric H<sub>2</sub>O and CO<sub>2</sub>. The application of  
 375 this method to other MAJIS data is illustrated in more detail in Sect. 4.4.

376



377

378 **Figure 7:** Land detection obtained by comparing the shapes of low brightness temperature  
 379 ( $T_B$ ) regions in MAJIS cube C4 with known coastlines. **A)** Identification of Busuanga and  
 380 Coron islands (markers 1 and 2 respectively), colder than the surrounding ocean, as well as  
 381 clouds (marker 3). **B)** Spectral contrast in brightness temperature ( $T_B$ ) with respect to the  
 382 ocean spectrum, measured over the islands (Busuanga in red, Coron in orange) and over a  
 383 thin cloud (grey curve). Coastlines data from OpenStreetMap, available under the Open  
 384 Database License.

385

### 386 3.2. Ice characterization

387

388 MAJIS and PRISMA data allow investigating the distribution of physical properties of ice and  
 389 how they relate, for example, to the altitude of the clouds where it is identified (see Sect. 4.1  
 390 and 4.2). The temperature, crystallinity, grain size, purity, and density affect the shape of ice  
 391 absorption bands (in particular the main ones at 1500 nm and 1900 nm) and of the  
 392 continuum. Since the long wavelength shoulder of the 1900 nm band encompasses the  
 393 noisy junction between the VISNIR and IR channels of MAJIS, we focus on the 1500 nm  
 394 band, spectrally well resolved in both MAJIS and PRISMA datasets. This band has a  
 395 characteristic asymmetry (due to its differential intensity with respect to the 1900 nm one,  
 396 e.g. Stephan et al., 2021) affecting the position and shape of the in-between transmission  
 397 window peak (~1700 nm) and has been used for the definition of the ice index in Table 3.  
 398 Within the 1500 nm band, the weaker 1650 nm absorption is present. Its strength is a proxy  
 399 for the degree of the ice crystallinity and temperature (Fink and Larson, 1975; Filacchione et  
 400 al., 2016). It is also observable in PRISMA, even if shallower and noisier due to the lower  
 401 spectral resolution (see zooms in Fig. 8A and B).

402 The 3000 - 4000 nm wavelength range, not accessible to PRISMA, hosts two ice reflection  
 403 peaks at around 3100 nm (the Fresnel peak) and 3700 nm (Fig. 8C). The former varies in  
 404 shape and intensity as a function of the ice crystallinity (Cartwright et al., 2025) while the  
 405 latter shifts to longer wavelengths as temperature increases (e.g. Filacchione et al., 2016,  
 406 see Sect. 3.3.3). Fresnel peak position variations are estimated in the data through  
 407 cross-correlating each ice spectrum with a constant shape (average peak shape in each  
 408 cube) which is rigidly shifted with a 0.1 nm sampling (hence allowing the estimation of the  
 409 peak with a sampling better than the nominal MAJIS one). On the other hand, the 3700 nm

410 peak position is obtained through fitting with a Gaussian function, reliably reproducing its  
411 shape.

412 Another proxy of the ice temperature is the intensity of its thermal emission, becoming  
413 significant at wavelengths larger than 4500 nm (Fig. 8D). However, in this range the emitted  
414 radiance is absorbed by a plethora of narrow bands of gaseous water, and therefore only a  
415 narrow transmission window around 4600 nm is suitable for this purpose. Table 4  
416 summarizes these ice spectral features, identifiable in MAJIS and PRISMA data. The  
417 average uncertainties  $\Delta$  are propagated taking into account the SNR estimates described in  
418 Sect. 2.3.1.

419

420

ICE PARAMETER	$\Delta$ MAJIS	$\Delta$ PRISMA	ICE PROPERTIES
1500 nm band depth	< 1 %	< 1 %	number density / grain size
1500 nm band asymmetry	< 2 %	< 3 %	grain size / crystallinity
1650 nm band depth	10 %	20 %	crystallinity
<i>Fresnel peak position</i> <i>Fresnel peak intensity</i>	<i>2 nm</i> <i>&lt; 1 %</i>	<i>/</i>	<i>temperature / crystallinity</i>
<i>3700 nm peak position</i> <i>3700 nm peak intensity</i>	<i>0.2 nm</i> <i>&lt; 1 %</i>	<i>/</i>	<i>temperature / crystallinity</i>
<i>4600 nm thermal intensity</i>	<i>&lt; 1 %</i>	<i>/</i>	<i>temperature</i>

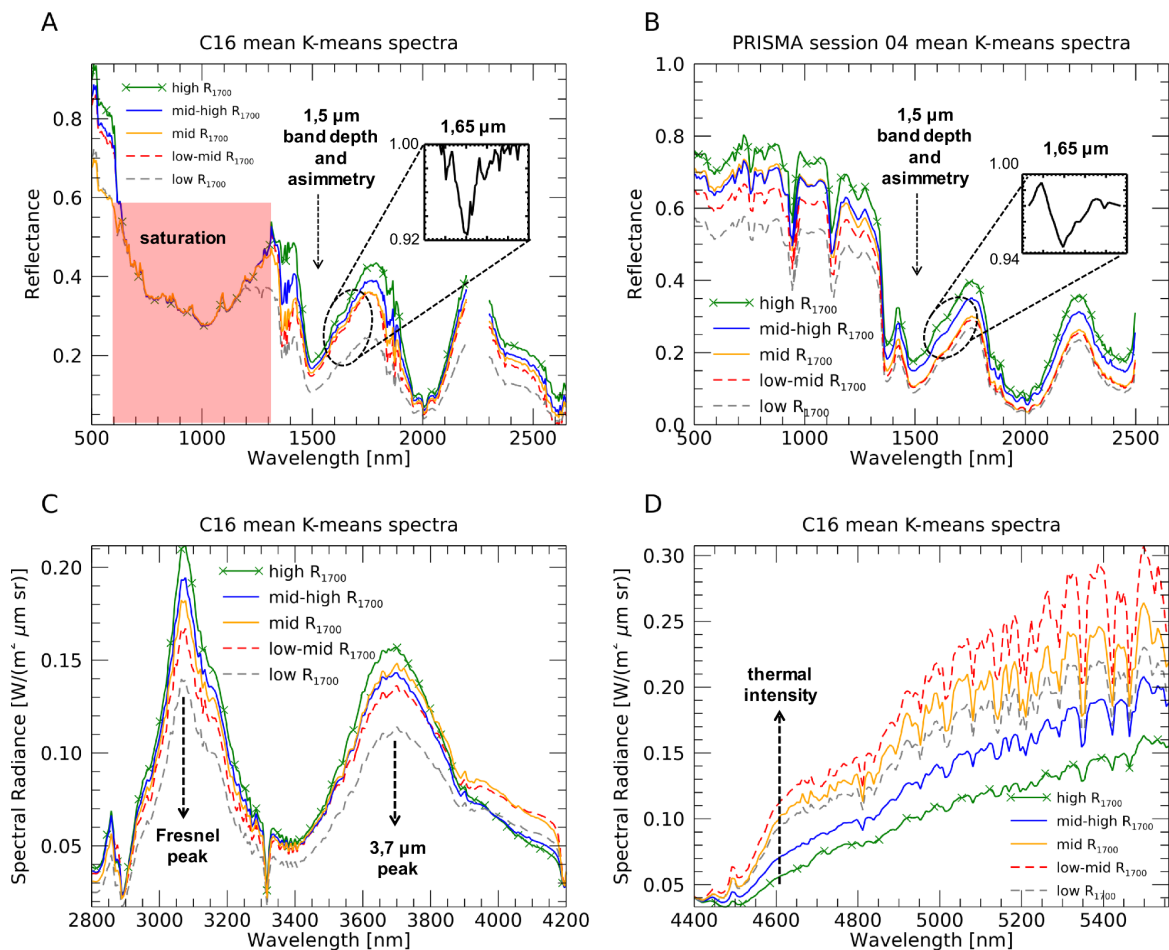
421 **Table 4:** investigated ice spectral parameters and related average uncertainties ( $\Delta$ ) and ice  
422 properties. Cells in light blue indicate parameters that only refer to MAJIS dataset.

423

424 As a first investigation of the ice spectral variability in MAJIS and PRISMA observations we  
425 take advantage of the unsupervised K-means classification algorithm included in the ENVI  
426 software package, version 6.0 (Exelis Visual Information Solutions, Boulder, CO, USA,  
427 <https://www.nv5geospatialsoftware.com/Products/ENVI>, accessed December, 15, 2025).  
428 This algorithm is capable of grouping the observations into an ensemble of “K”  
429 non-overlapping clusters, driven by spectral similarity, whose average spectra are  
430 representative of the main signatures in the dataset. This is done through an iterative  
431 minimum distance technique whose details are described in Tou and Gonzalez (1974). In  
432 this preliminary analysis, we arbitrarily set the algorithm to produce K=5 output average  
433 spectra, enough for visualizing the variability of the main ice diagnostic spectral features. It  
434 must be noted that, since we are also interested in features pertaining to infrared  
435 wavelengths, in MAJIS case the full VISNIR+IR spectral range is considered, and  
436 wavelengths longward of 2500 nm contribute to the clustering as well. As we will see in Sect.  
437 4.1, this also has an impact on the spatial distribution of the clusters. The resulting average  
438 spectra in the solar range are shown in Fig. 8A and Fig. 8B for MAJIS cube C16 and for one  
439 of PRISMA session 04 cubes respectively. These spectra result to be mainly driven by the  
440 changing intensity of the continuum, in turn providing information about the opacity of the ice  
441 clouds. The color scale is associated with increasing reflectance of the transmission window  
442 at 1700 nm (dashed grey, dashed red, orange, blue and green with crosses from low to high,  
443 indicating increasing opacity and variable crystal sizes). The same color scheme is retained

444 for the intensity of Fresnel peak at 3100 nm in MAJIS data (Fig. 8C), diagnostic of the ice  
 445 crystallinity. Instead, spectra with intermediate reflectances at 1700 nm switch order within  
 446 the 3700 nm ice reflectivity peak (dashed red to blue to orange from low to high, Fig. 8C)  
 447 indicating the increased weight of thermal emission on the overall signal in this range. At  
 448 MAJIS wavelengths larger than 4500 nm (Fig. 8D) the initial color scheme is totally  
 449 disrupted, due to the mixing of information about cloud emissivity, cloud temperature (i.e. the  
 450 altitude) and gaseous opacity. The combination of high NIR reflectances and low thermal  
 451 emission (green spectrum with crosses) suggests the presence of optically thick high-altitude  
 452 clouds, as confirmed by the shallower water absorption bands longward of 4900 nm. On the  
 453 other hand, large thermal radiance and deep water bands associated with intermediate NIR  
 454 reflectance (dashed red spectrum) indicate a population of moderate opacity clouds at quite  
 455 low altitudes. The other spectra present intermediate properties in the thermal range, not  
 456 strongly correlated with the NIR reflectance, calling for mixed-phase clouds of variable  
 457 microphysical properties and vertical structure.

458  
 459



460

461 **Figure 8:** A: mean reflectance spectra from the K-means clustering algorithm for MAJIS  
 462 cube C16 ( $\lambda < 2500$  nm). The red shaded area indicates wavelengths that are saturated due  
 463 to the high reflectivity of clouds. Colors indicate different regimes of the continuum  
 464 reflectance, taken as reference at the 1700 nm transmission window ( $R_{1700}$ ). B: same as in  
 465 A but for PRISMA session 04 (full spectral range). The insets in A and B zoom between  
 466 1570 and 1780 nm to show the average 1650 nm band normalized to the continuum. C:

467 MAJIS radiances in the  $2800 < \lambda < 4200$  nm range, zooming on the Fresnel and 3700 nm ice  
468 reflectivity peaks. **D:** thermal part of the spectrum longward of 4400 nm. In all panels,  
469 dashed arrows highlight diagnostic spectral features of the ice.

470

### 471 **3.3. Estimation of clouds' altitude**

472

473 The most straightforward method for evaluating cloud altitudes involves the correlation of the  
474 brightness temperature at a given wavelength (e.g. 4610 nm, less affected by gaseous  
475 absorption in the MAJIS range) with a known vertical temperature profile. For ice clouds,  
476 temperatures can be derived from the 3700 nm peak position (Sect. 3.3.3). Other methods  
477 that we consider here are based on O<sub>2</sub> absorption bands' variability and on the analysis of  
478 clouds' shadows (Sect. 3.3.1 and 3.3.2). In this study, we rely on a fixed average  
479 temperature profile (Efremenko and Kokhanovsky, 2021), which may be not representative  
480 of the actual thermodynamic conditions of the atmosphere during the observations. As a  
481 consequence, all the methods that we adopt yield a range of results, each affected by their  
482 own intrinsic limitations. Although they appear quite consistent with each other, more  
483 quantitative investigations are postponed to future analyses.

484

#### 485 **3.3.1. O<sub>2</sub> band depth variability**

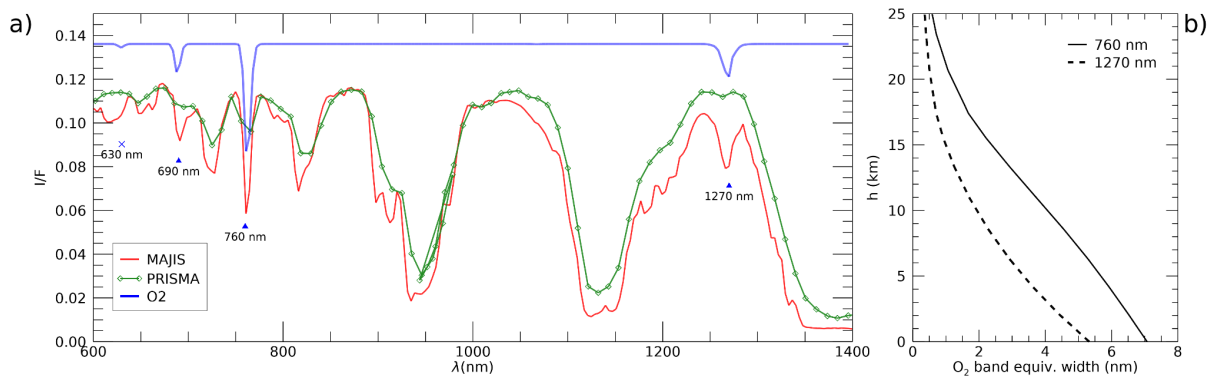
486

487 The O<sub>2</sub> spectral features covered by both MAJIS and PRISMA observations consist of the  
488 absorption bands at 630 nm, 690 nm, 760 nm and 1270 nm (Newnham & Ballard, 1998;  
489 Smith & Newnham, 1999). As we can see in Fig. 9A, MAJIS can resolve all bands except the  
490 630 nm one, while PRISMA data can only partially resolve the 760 nm one. The strongest  
491 760 nm band is the most used from satellite measurements in the near-infrared (e.g.  
492 GOSAT, Butz et al., 2011; SCIAMACHY, Bovensmann et al., 1999; TROPOMI, Veeffkind et  
493 al., 2012; OCO-2/3, Eldering et al., 2019) for inferring bulk atmospheric quantities like  
494 temperature profile, airmass (Stevens et al., 2017), aerosol and clouds properties (Geddes &  
495 Bösch, 2015). O<sub>2</sub> is a well-mixed component of the atmosphere, hence the curves of growth  
496 of its absorption bands with altitude in the presence of optically thick clouds can be  
497 translated into the altitude of the cloud top (e.g. Wei et al., 2024).

498 In our analysis we applied a simplified scheme for retrieving cloud top altitudes from the 760  
499 nm band in the PRISMA case and from both 760 and 1270 nm O<sub>2</sub> bands for MAJIS data.  
500 The different strength of the two bands implies a different curve of growth with altitude (Fig.  
501 9B), with the 1270 nm one less sensitive to higher clouds but more suitable for  
502 characterizing lower structures. The 630 and 690 nm bands, intrinsically weaker and more  
503 sensitive toward the surface, are not used in this analysis.

504 The comparison of a measured O<sub>2</sub> band depth with its theoretical curve of growth, evaluated  
505 for the actual airmass, allows us to directly retrieve the cloud top altitude (Sect. 4.2.1). It is  
506 worth stressing that although altitude, pressure and temperature of the cloud top are  
507 important atmospheric parameters (Nakajima et al., 2019), our simplified scheme neglects  
508 details of vertical distributions and scattering properties, introducing possible biases in the  
509 retrieved absolute values. Propagating the MAJIS uncertainties previously discussed (Sect.  
510 2.3.1) and assuming suitable model ones (~10% on the oxygen vertical profile induced by  
511 local changes in gaseous temperature, density, humidity), errors on cloud top altitude  
512 average to values of ~1 km, for both the 760 and 1270 nm bands. In addition, the 1270 nm  
513 band is known to contain a significant airglow emission feature that can alter the band depth  
514 and introduce further biases in the oxygen absorption evaluation (Kuang et al., 2002).

515



516

517 **Figure 9: A):** Typical appearance of O<sub>2</sub> features in the spectra of MAJIS (red) and PRISMA  
518 (green). Modeled spectral transmittance (in blue) highlights location and shape of the O<sub>2</sub>  
519 bands at 630, 690, 760, and 1270 nm. Only the last three can be detected in MAJIS spectra  
520 (red curve), while only the strongest 760 nm band is identifiable in PRISMA spectra (red  
521 curve). **B):** Examples of curves of growth of the O<sub>2</sub> absorption at 760 (solid curve) and 1270  
522 nm (dashed curve) in the standard clear-sky atmospheric column adopted in this work. The  
523 absorption is shown as band equivalent width for a 2-ways path with generic incidence and  
524 emission angles of 30°.

525

### 526 3.3.2. Cloud shadows analysis

527

528 The length of projected clouds shadows gives an estimate of their altitude, provided that the  
529 illumination geometry is well known. Significant lengths of projected shadows are more  
530 easily seen in the case of tall convective clouds in slant solar illumination. In the MAJIS case,  
531 clear shadows have been identified for strong convective events surrounded by widespread  
532 background clouds, hence their length can only give hints on relative altitudes (see Sect.  
533 4.2.2). Here we adopted the simplest assumption of homogeneous cylindrical shapes  
534 (without accounting for the actual three-dimensional density distribution) to estimate the  
535 height of clouds' top with purely geometrical considerations. Under these hypotheses the  
536 associated uncertainties are mainly driven by errors in edge detection (for both cloud and  
537 shadow edges) and, as a consequence, are limited by the spatial resolution. Errors on solar  
538 incidence angle may also play a role in very slant illumination, and the total relative  
539 uncertainties estimated in the conditions of MAJIS observations range between 6 and 10%.  
540 The application of the method is shown in Sect. 4.2.2.

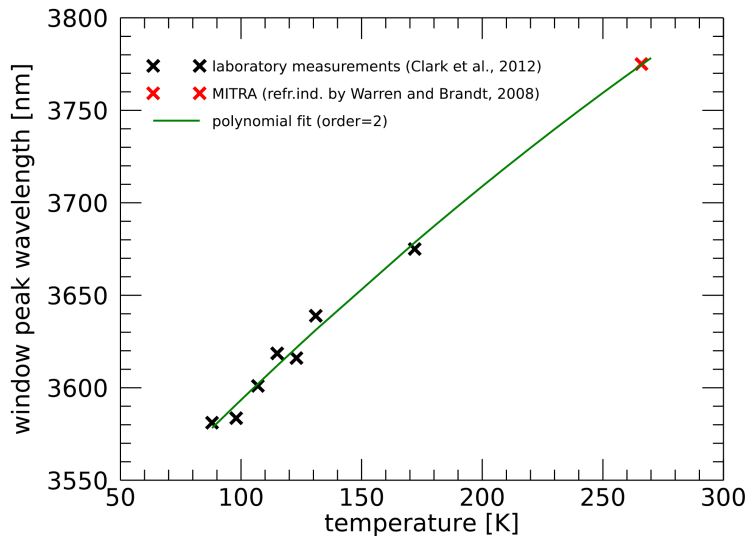
541

### 542 3.3.3. Derivation of clouds' altitude with the ice temperature

543

544 We apply to Earth's icy clouds the same method by Filacchione et al. (2016), who estimated  
545 the temperatures of Saturn's icy satellites surfaces from the displacement of the 3700 nm ice  
546 peak, deriving from a shift of the imaginary part of the ice refractive (Mastrapa et al., 2009).  
547 In that method, temperature-dependent peak reflectivities were derived from laboratory  
548 measurements by Clark et al. (2012), spanning between 88 and 172 K, a range too low to  
549 describe Earth troposphere where clouds are commonly observed. We extrapolate the  
550 peak-temperature dependence by also simulating the ice reflectivity at 266 K, i.e. the  
551 temperature of the optical constants by Warren and Brandt (2008). Since the ice grain size  
552 has little effect on the peak position (Filacchione et al., 2012) we assume an effective radius  
553 of 20 μm, representative of cirrus clouds (LeMone, 1988). The resulting trend covering from

554 88 to 266 K is shown in Fig. 10 (black and red crosses). It is reliably fit with a second-degree  
 555 polynomial (green line) and can be used for a qualitative estimation of the ice temperature in  
 556 MAJIS observations (Sect. 4.2.3).  
 557



558  
 559 **Figure 10:** correlation between ice temperature and 3700 nm reflectivity peak position. Black  
 560 crosses represent laboratory measurements by Clark et al. (2012), the red cross indicates an  
 561 RT simulation performed with ice grain size of 20  $\mu\text{m}$  and optical constants by Warren and  
 562 Brandt (2008), and the green line represents a second degree polynomial fit of all data (see  
 563 Sect. 3.3.3).  
 564

### 565 3.3.4. Forward RT modeling on liquid and ice H<sub>2</sub>O clouds;

566  
 567 The most accurate method for determining clouds' vertical distribution is through full RT  
 568 modeling. However, this would require a time-consuming retrieval of physical quantities that  
 569 is beyond the scope of this paper. Instead of spectral inversion, we here perform a  
 570 comparison of selected observations (i.e., those in Fig. 3) with forward RT models obtained  
 571 by manually tuning aerosols' physical parameters. The derived quantities are to be  
 572 considered as orders of magnitude of the altitude and microphysical properties of Earth's  
 573 clouds and aerosols. Forward models are produced with the MITRA RT tool (Oliva et al.,  
 574 2016; 2018; Sindoni et al., 2017; D'Aversa et al., 2022), adopting the optical constants from  
 575 Hale and Query (1973), Warren and Brandt (2008) and Kitamura et al. (2007) for computing  
 576 the scattering properties of liquid water, water ice and silicate minerals (assumed as  
 577 background aerosol), respectively. The spectral albedo of the ocean is taken from the  
 578 ASTER spectral library (Baldrige et al., 2009). In this simplified scheme, we neglect thermal  
 579 emission, discarding measurements longwards 3000 nm.

580 It is interesting to note that, even if beyond the scope of this paper, more accurate RT  
 581 modeling could also be considered for the evaluation of straylight contamination (studied for  
 582 MAJIS in Langevin et al., this issue), as it offers the possibility to extrapolate information  
 583 from the NIR part of the spectrum to visible wavelengths.  
 584

### 585 3.4. High altitude emissions and atmospheric waves identification

586

587 Among the many gaseous features observable in the 4000-5500 nm MAJIS range, two are  
588 particularly interesting, being observed as emission bands. These are the CO<sub>2</sub> double-peak  
589 at the bottom of the main 4300 nm band and an O<sub>3</sub> signature around 4700 nm. Both are  
590 evident above optically thick clouds at high altitudes, blocking the thermal contribution from  
591 the surface and lower (hotter) atmospheric layers. The CO<sub>2</sub> peak is radiometrically much  
592 more stable than other spectral features against variation of atmospheric structures (see  
593 Poulet et al., this issue). It is known to result from the combination of a LTE component  
594 induced by temperature increase in the stratosphere, and a non-LTE one due to the CO<sub>2</sub>  
595 excitation primarily induced by direct solar pumping occurring at even higher altitudes (where  
596 collisional quenching is no longer efficient, e.g. Cassini et al., 2025). The detailed analysis of  
597 this emission feature in MAJIS data, implying the evaluation of CO<sub>2</sub> vibrational temperature  
598 vertical profiles, is far beyond the purpose of this work. In any case, the spatial distribution of  
599 the CO<sub>2</sub> emission intensity can provide interesting insights about the probed layers, and we  
600 can indeed use it for detecting atmospheric waves and provide hints about their altitude and  
601 propagation (see Sect. 4.3.1). CO<sub>2</sub> emission can be identified already in MAJIS  
602 monochromatic frames at 4270 nm (i.e. the position of the main peak of the emission) but  
603 the integration of the band in a narrow spectral range is useful for reducing noise and  
604 enhancing the contrast in waves' investigation (Sect. 4.3). For the integration we consider  
605 wavelengths between 4254 and 4333 nm, which probe high altitudes in the atmosphere and  
606 are not affected by the thermal contribution from lower ones. Considering the SNR estimated  
607 at these wavelengths (Fig. 3C), we are able to detect waves whose relative intensity  
608 between crests and troughs is about 1%, assuming a 3-sigma uncertainty for the radiance at  
609 4270 nm.

610 On Earth, ozone has a maximum density in the lower stratosphere but its vertical distribution  
611 strongly depends on latitude (see for example Bekki and Lefevre, 2009). It is produced  
612 through a very fast and exothermic 3-body recombination reaction that includes O and O<sub>2</sub> in  
613 the presence of a catalytic species (either N<sub>2</sub> or O<sub>2</sub>). Aside from diagnostic bands at UV  
614 (outside MAJIS domain) and VIS wavelengths (the Chappuis band discussed in Sect. 2.3.2),  
615 the 4700 nm one is the strongest feature clearly detectable within the MAJIS range. This O<sub>3</sub>  
616 band is seen as either an absorption or emission feature in MAJIS nadir-looking  
617 observations, depending on the overall thermal emission of the atmospheric column. In clear  
618 sky conditions, when the emission from lower warmer layers is dominant, the O<sub>3</sub> 4700 nm  
619 band is hardly detectable being overcome by water absorption (as shown in Poulet et al., this  
620 issue), unless radiative transfer modeling is performed on the data (e.g. Guerlet et al., this  
621 issue). In the presence of mid-altitude clouds, a shallow O<sub>3</sub> band appears in absorption,  
622 while the obstruction of the densest part of the atmospheric column due to high-altitude  
623 clouds makes the O<sub>3</sub> band appear in emission. Given this phenomenology, in this preliminary  
624 study we investigate the O<sub>3</sub> emission amplitude through the difference between brightness  
625 temperatures estimated at 4717 nm (strongest O<sub>3</sub> line) and 4660 nm (outside O<sub>3</sub> band). Such  
626 a difference is positive when the O<sub>3</sub> is spectrally observed in emission, negative otherwise.

627

628

629

### 3.4.1. Atmospheric waves characterization

630

631 Atmospheric gravity waves are observed in almost all the MAJIS acquisitions (see examples  
632 in Sect. 4.3.1) at the wavelengths of the central peak of the 4300 nm CO<sub>2</sub> band. Due to the  
633 limited field of view, wave packets are usually not visible in their entirety and it is not possible

634 to identify the same wavy structures from one image to the other due to the large coverage  
635 gaps, preventing the study of the wave speed propagation. Nevertheless, we attempt to  
636 quantify wave properties and provide some hints on their altitude. The wave parameters -  
637 horizontal wavelength, total packet length, azimuthal extent, and packet width - are  
638 determined through visually processing each image. Automated methods were not possible  
639 because of the variability in images' contrast. After appropriate image stretching, the  
640 wavefronts are identified by tracing the crest lines. The horizontal wavelength is defined as  
641 the average distance between consecutive crests within each wave packet. The total packet  
642 length is measured as the distance between the first and last identified crests. The azimuthal  
643 extent is derived from the common orientation of the crests counted counterclockwise, while  
644 the packet width is defined as the maximum crest length among those identified. Taking into  
645 account spatial resolution and signal contrast, uncertainties in size estimation are of about 7  
646 km, while those on wavelengths are less than about 11 km.

647 Circular-wave patterns have been observed in some MAJIS images, likely resulting from the  
648 breaking of upward-propagating waves originating in sufficiently strong convective  
649 thunderstorms. Under this assumption, we attempt to infer the time delay between the  
650 wave-triggering event and its observation (Taylor and Hapgood, 1988; Dewan et al., 1998;  
651 see Sect. 4.3.1). This is done by neglecting wind transport and assuming a simplified  
652 isothermal dispersion relation (Hines, 1960) in which the wave speed is negligible with  
653 respect to the speed of sound. For circular waves we also measured maximum radius and  
654 expansion speed. Because only a portion of the circle is visible in the images, the wave  
655 radius is inferred from the following formula:

$$656 \quad r = f/2 + c^2/8f \quad (1)$$

657 where  $f$  is the sagitta and  $c$  is the chord, observed in the images in pixels and converted to  
658 km using the instantaneous resolution in km/px reported in Table 1. The waves are assumed  
659 to occur at 15 km height, based on estimations of the brightness temperature of CO<sub>2</sub>  
660 emission at 4300 nm, used as a proxy for these features. Following the formulation in Taylor  
661 and Hapgood (1988) and Dewan et al. (1998), the wave period  $\tau$  and expansion speed  $v_{gx}$   
662 can be obtained using the following formulas:

$$663 \quad r^2 = \tau_B^2 [(1 + 1/(\tan \phi)^2)] \quad (2)$$

$$664 \quad v_{gx} = (\lambda_x / \tau) [1 - (\tau_B / \tau)^2] \quad (3)$$

665 with  $\phi$  being the elevation angle identified by the wave propagation direction,  $\lambda_x =$   
666 wavelength of the propagating wave, and  $\tau_B =$  buoyancy period, the latter assumed to be  
667 equal to 5 min, which is a good approximation at stratospheric altitude (Dewan and Good,  
668 1986).

669

#### 670 4. Results and discussion

671

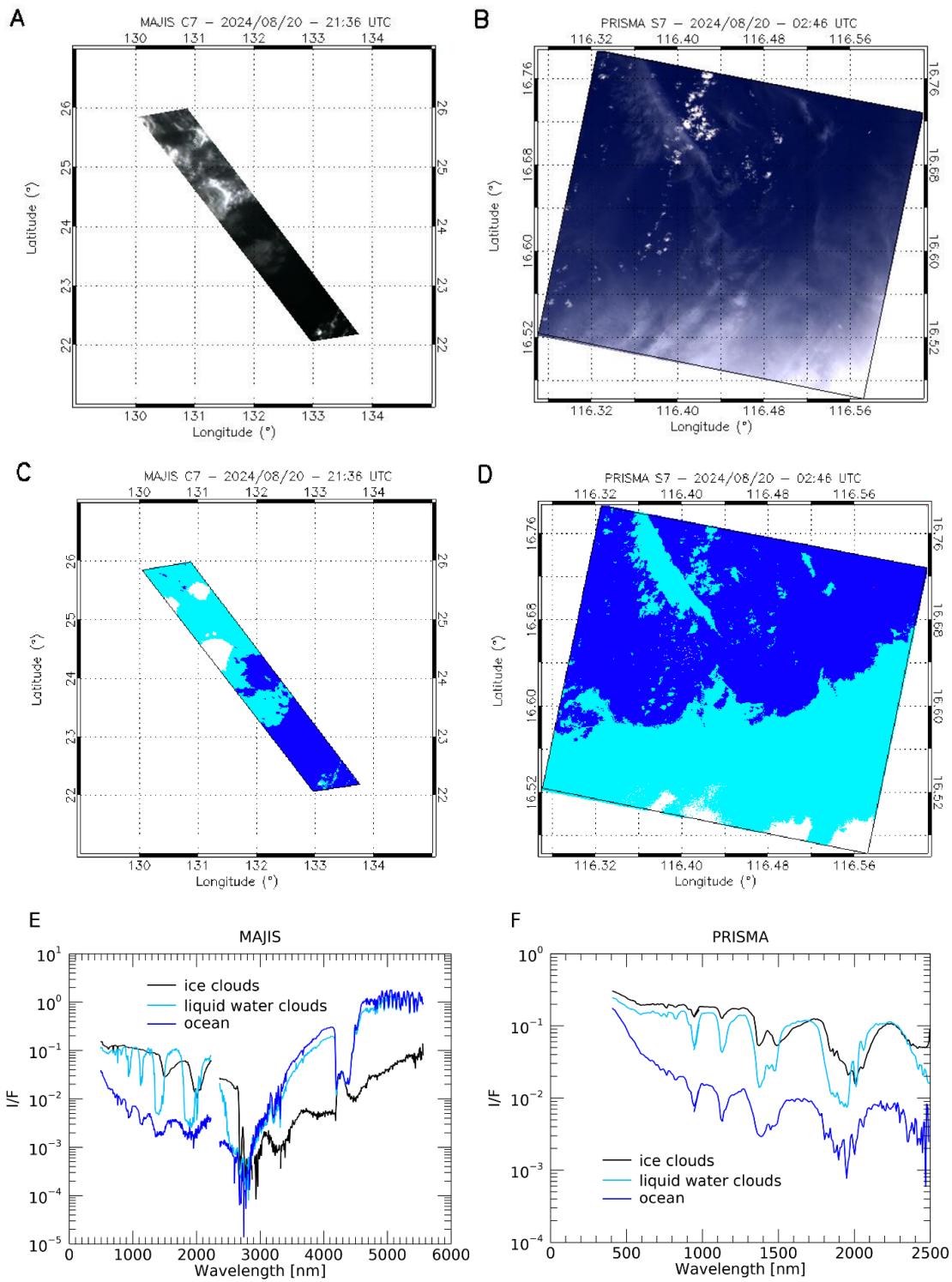
672 We now present the results we obtain through the application of the methods discussed in  
673 Sect. 3. Sect. 4.1 provides a discussion on ice properties, Sect. 4.2 focuses on the clouds'  
674 altitudes, Sect. 4.3 is devoted to high altitude features and Sect. 4.4 presents results on land  
675 features identification.

676

#### 677 4.1. Icy clouds properties

678

679 Examples of two MAJIS and PRISMA cubes containing ice clouds, identified through the ice  
680 spectral index in Table 3 (threshold  $< 1$ ), are given in Fig. 11. In the MAJIS case, ice is found  
681 in localized convective clouds (Fig. 11A and C), so high with respect to the background  
682 structures that they even cast well detectable shadows (see Sect. 4.2.2). Instead, in the  
683 PRISMA observation ice is detected both in diffuse bright clouds (e.g. at the southern east  
684 corner of Fig. 11B and D) and in thinner and less contrasted structures (probably identifiable  
685 as high altitude cirrus clouds, e.g. the white regions around longitude  $116.4^\circ$  - latitude  $16.5^\circ$ ,  
686 Fig. 11B and D) hence proving the effectiveness of the index with different regimes of ice  
687 optical depth. Sample spectra from the identified classes are shown in Fig. 11E and F for  
688 MAJIS and PRISMA respectively. It must be noted that the very low albedo of the ocean in  
689 MAJIS spectrum at visual wavelengths ( $< 1\%$ ) is due to the very slant illumination conditions  
690 for the selected observation (incidence angle of about  $80^\circ$  for cube C7, see Table 1). On the  
691 other hand, the spectra in the thermal range show consistency with the expected  
692 temperature regimes, with very cold ice clouds and the ocean hotter than liquid water clouds.

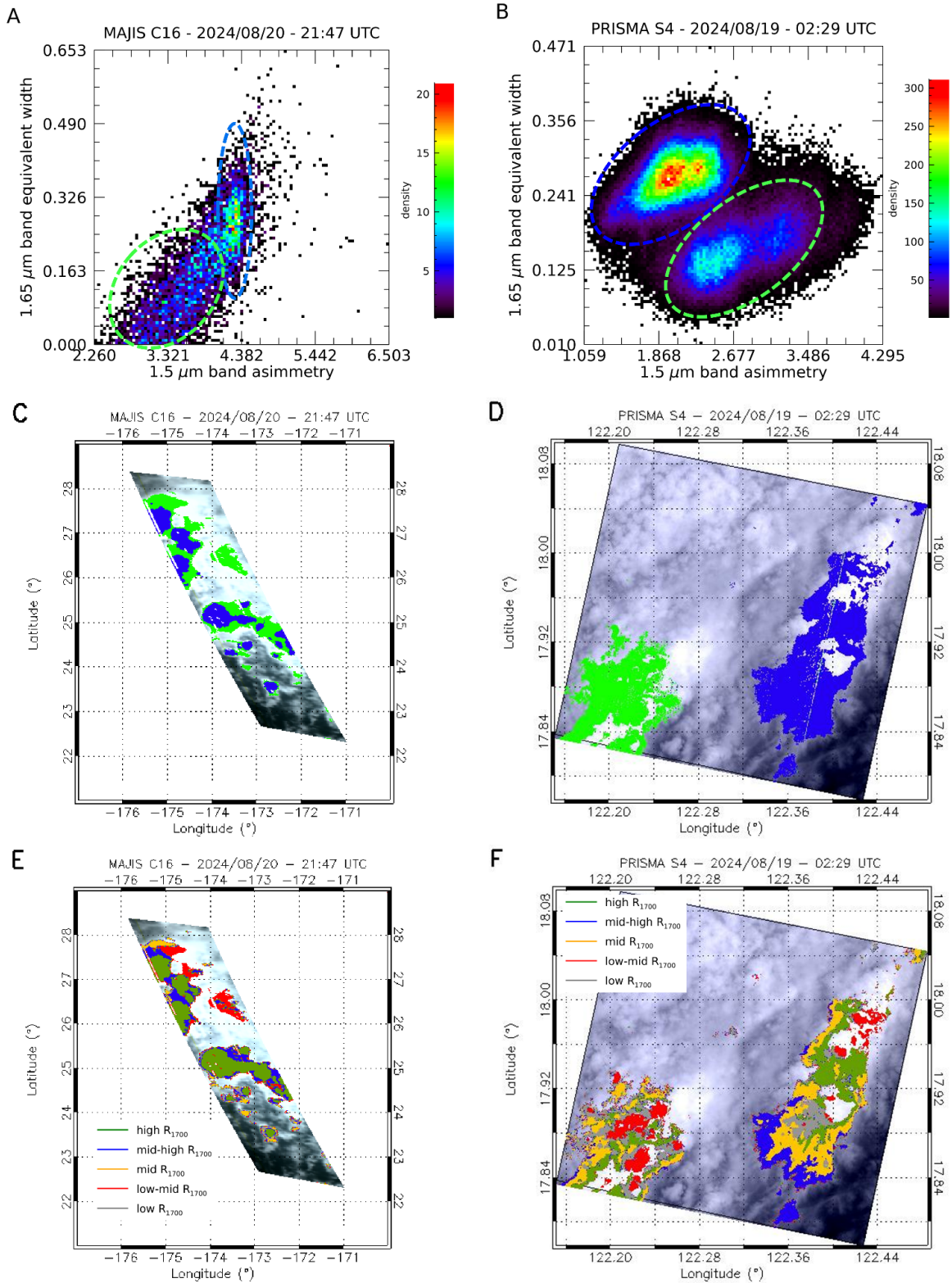


693

694 **Figure 11:** Panels **A** and **B** refer to MAJIS cube C7 and one of the PRISMA cubes from  
 695 session 07, respectively, displayed in RGB. Panels **C** and **D** show the masks for the  
 696 detection of ocean (blue), liquid water clouds (cyan, from the “cloudy” condition in Table 3)  
 697 and ice clouds (white) pixels related to the two cubes. Panels **E** and **F** display sample  
 698 spectra related to the different classes identified in MAJIS and PRISMA observations.

699

700 Ice is similarly widespread in other MAJIS and PRISMA data, so that some considerations  
701 on its distribution and correlations between its parameters can be made (Fig. 12). We  
702 compute the 1500 nm band asymmetry as a ratio of slopes, the first considered between  
703 1415 and 1500 nm (left wing) and the second between 1500 and 1790 nm (right wing). The  
704 asymmetry correlates with the strength of the 1650nm band (quantified as equivalent width,  
705 Fig. 12A and B), with higher values indicating increasingly crystalline ice (Mastrapa, 2008;  
706 Stephan et al. 2021; Grundy & Schmitt 1998). Different regimes of these two parameters  
707 map localized structures in MAJIS and PRISMA observations, as shown in Fig. 12C and D  
708 respectively where green and blue pixels refer to clusters contained within dashed ellipses  
709 sharing the same color in Fig. 12A and B. In MAJIS case, the blue cluster is characterized by  
710 an increasing 1650 nm equivalent width at constant 1500 nm band asymmetry. The green  
711 cluster, instead, shows a common trend of growth for the two parameters. On the other  
712 hand, the PRISMA ellipses identify well separated clusters of points within the two  
713 parameters' space.



714

715

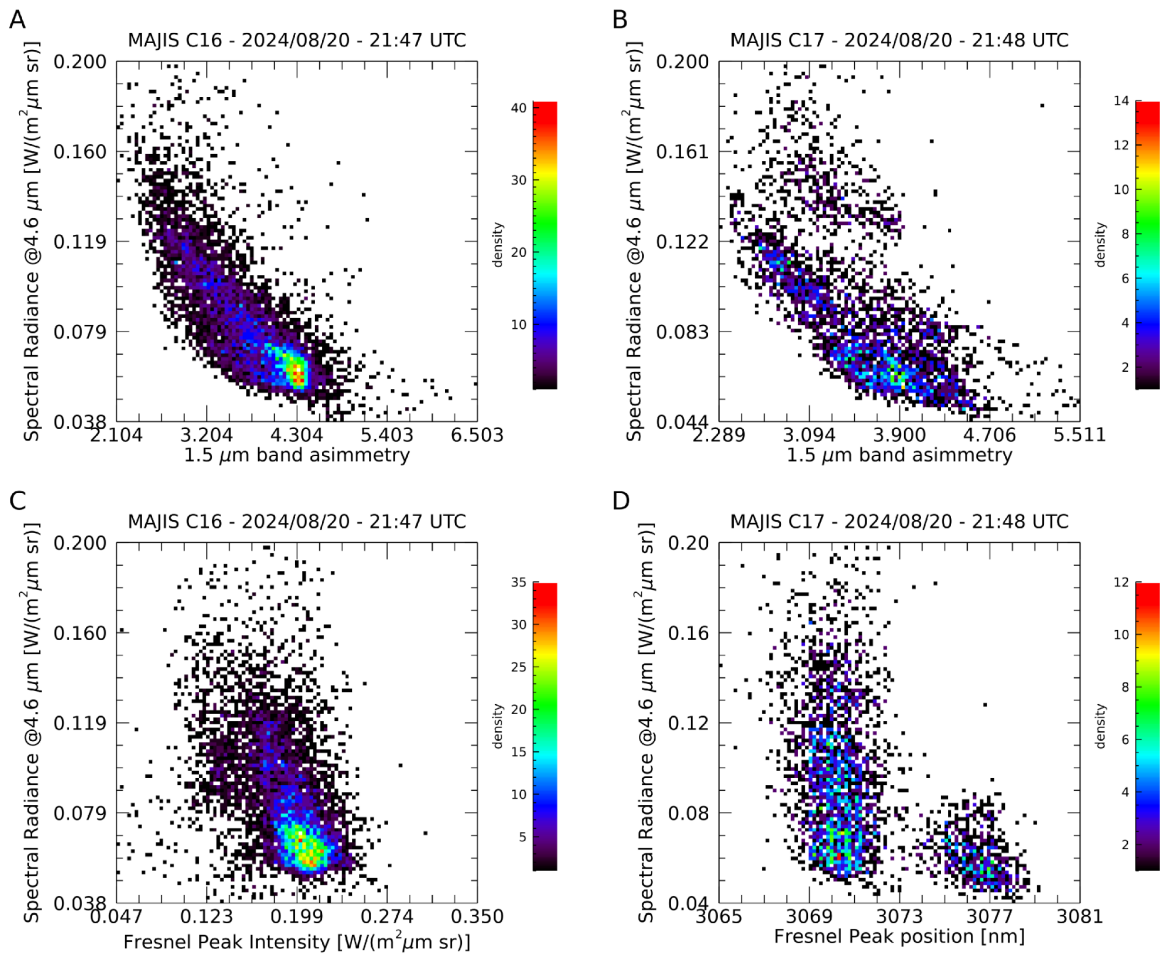
716 **Figure 12:** **A-B:** scatterplots of the 1500 nm band asymmetry and the 1650 nm band  
 717 equivalent width for MAJIS reflectance cube C16 and one of the PRISMA reflectance cubes  
 718 from session 04. The colored-dashed ellipses separate different regimes of the two  
 719 parameters (see Sect. 3.5). **C-D:** green and blue pixels map the clusters contained within the  
 720 respective ellipses in panels A and B. **E-F:** clustering of ice observations, obtained through

721 the K-means classification algorithm (see Sect. 3.2), grouped on the basis of the intensity of  
722 the reflectance at 1700 nm ( $R_{1700}$ ).

723

724 It is interesting to note that the correlation between these clusters and those obtained from  
725 the K-means classification discussed in Sect. 3.2 (Fig. 12E-F) is not straightforward. For  
726 MAJIS, ice spectra with high reflectivity in the solar part of the spectrum (green in Fig. 12E)  
727 are mostly correlated with the blue cluster in Fig. 12C. This trend is not observed in PRISMA,  
728 where all K-means clusters are equally distributed over both the blue and green clusters  
729 shown in Fig. 12D, suggesting variable ice densities and grain sizes within the same regimes  
730 of crystallinity. This difference derives from the fact that, as explained in Sect. 3.2, for MAJIS  
731 the thermal wavelengths contribute to the K-means classification of the spectra, hence  
732 providing information also on the temperature of the ice (see also Sect. 4.2.3). This is  
733 verified by the trend of the 1500 nm band asymmetry with the radiance in the thermal part of  
734 the spectrum, shown for MAJIS cubes C16 and C17 in Fig. 13A-B: more crystalline ice  
735 (larger asymmetry) is correlated with lower radiances (i.e. temperatures) at thermal  
736 wavelengths. In particular, orbit C17 also shows a detached cluster in the distribution of the  
737 thermal radiance suggesting different regimes of temperature (hence different clouds'  
738 altitude). Finally, we show the correlation between the ice crystallinity and its temperature in  
739 Fig. 13C-D, where the intensity and wavelength of the Fresnel peak are compared to MAJIS  
740 thermal radiances. Consistently with previous studies (e.g. Stephan et al., 2021), the  
741 intensity of Fresnel peak is higher when the temperature is low (Fig. 13C), indicating  
742 enhanced crystallinity (see also Poulet et al., this issue). The comparison in Fig. 13D shows  
743 two distinct regimes of the peak position, with the short wavelength cluster characterized by  
744 a larger spread of the thermal radiance (suggesting an enhanced temperature variability for  
745 a less crystalline ice, e.g. Stephan et al., 2021).

746



747

748 **Figure 13: A-B:** scatterplots of the 1500 nm band asymmetry and thermal radiances at 4600  
 749 nm for MAJIS orbit C16 and C17 respectively. **C:** scatterplot of the Fresnel peak intensity  
 750 with the thermal radiance at 4600 nm for MAJIS orbit C16. **D:** scatterplot of the Fresnel peak  
 751 wavelength with the thermal radiance at 4600 nm for MAJIS orbit C17.

752

## 753 4.2. Clouds' altitude

754

755 We now discuss the altitudes of clouds derived with the different methods presented in Sect.  
 756 3.3.

757

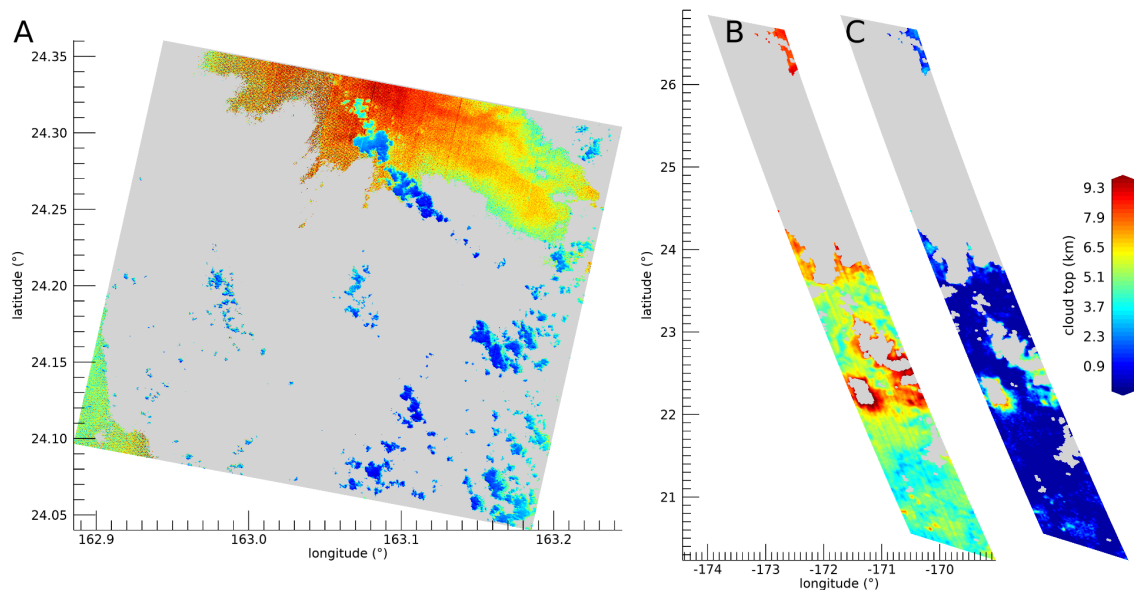
### 758 4.2.1. Altitudes from O<sub>2</sub> band depths

759

760 Figure 14 shows a comparison of cloud top altitude maps obtained by applying the O<sub>2</sub> bands'  
 761 investigation method (Sect. 3.3.1) to sample PRISMA and MAJIS cubes. In this case,  
 762 PRISMA retrievals (Fig.14A) show two main cloud layers, a higher one between 5 and 9 km  
 763 (yellow-red colors in the figure, modal value  $6.5 \pm 1$  km) and a lower one between less than 1  
 764 and 4 km (blue-cyan colors, modal value  $2.0 \pm 1$  km). MAJIS cloud tops (Fig.14B), whose  
 765 model value lies at  $4.8 \pm 1$  km, are in overall agreement with the PRISMA upper cloud deck,  
 766 even if the observing angles were very different in the two cases ( $>60^\circ$  for MAJIS,  $\sim 12^\circ$  for  
 767 PRISMA). On the other hand, the population of lower clouds detected by PRISMA appears  
 768 broken in a series of localized small structures that could remain unresolved if also present  
 769 in the MAJIS scene (see Sect. 2.3.1). A more systematic discrepancy is obtained when

770 comparing cloud top altitudes retrieved through the 760 nm and the 1270 nm bands  
 771 (Fig.14C, shown as reference for MAJIS data only), since the latter gives systematically  
 772 lower values (most clouds drop below 1.5 km altitude). This discrepancy reflects the smaller  
 773 sensitivity of this band to higher altitudes and the non optimal modeling assumptions  
 774 described in Sect. 3.3.1. However, such an issue can be resolved with a more complete  
 775 radiative transfer modeling as suggested by the benchmark presented in Sect. 4.2.4.  
 776

777

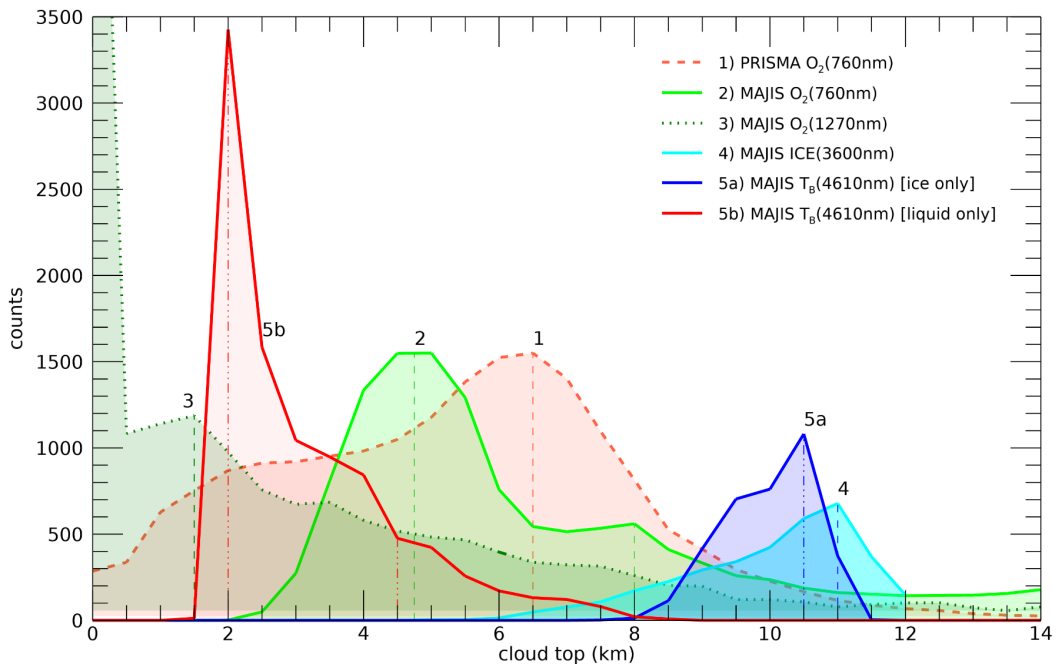


778

779 **Figure 14:** **A):** Map of cloud top altitude retrieved through the O<sub>2</sub> 760 nm band in a PRISMA  
 780 sequence 09 cube (20240820234657). Non-cloudy pixels or saturated ones, excluded from  
 781 the calculation, are shown in grey. **B):** the same as panel A but from a MAJIS data cube  
 782 C17. **C):** cloud top map for the same data in panel B (offset for clarity) but retrieved from the  
 783 1270 nm O<sub>2</sub> band. Uncertainties are of the order of 1 km (Sect. 3.3.1).  
 784

785 The counts distribution of cloud top altitudes derived from the maps in Fig. 14 is shown in  
 786 Fig. 15. The altitude ranges of the main cloud deck derived from the O<sub>2</sub> 760 nm band are in  
 787 good overall agreement between MAJIS and PRISMA (light green and dashed light red  
 788 curves), characterized by two broad peaks around 4.8 km and 6.5 km, respectively. The  
 789 displacement between these peaks is mainly driven by a true difference in the cloud  
 790 populations between the two observed scenes, and is further increased by the different  
 791 spatial/spectral resolutions. A much lower distribution, flattened towards the surface, is  
 792 indicated by the O<sub>2</sub> 1270 nm band, confirming its scarce usability to trace cloud altitudes.  
 793 The cloud heights derived from the 3700 nm ice spectral signature (cyan curve) only trace  
 794 icy pixels of MAJIS C17 cube (see lower panels of Fig. 17 in Sect. 4.2.3) and are distributed,  
 795 as expected, above the main cloud deck, with a peak around 11 km. The same behaviour is  
 796 also confirmed by the altitudes distribution evaluated through brightness temperatures'  
 797 estimation over the same icy spectra (blue curve, 10.5 km peak). The same estimation,

798 applied to non-icy spectra (red curve), gives altitudes significantly biased toward lower  
 799 levels. This can be mainly ascribed to the fact that the thermal part of the spectrum of thin  
 800 liquid water clouds observations is affected by an enhanced contribution from the ocean  
 801 thermal emission. Full radiative transfer calculations would be needed to quantitatively  
 802 assess this aspect and the assumptions on the specific emissivity of liquid/ice clouds (set to  
 803 1 for both in our calculations), but are not performed here as they are beyond the purpose of  
 804 this preliminary work.  
 805



806

807

808 **Figure 15:** Comparison of cloud top altitudes retrieved from PRISMA and MAJIS session 09  
 809 and C17 cubes respectively, through different methods. PRISMA counts are normalized to  
 810 the maximum value of MAJIS curve 2. Distributions derived from  $O_2$  band depths, related to  
 811 the maps in Fig. 14A, B, C, are shown in pink (dashed line), light green (solid line), and dark  
 812 green (dotted line) respectively. Cyan curve refers to ice clouds only (method described in  
 813 Sect. 3.3.3 and discussed in Sect. 4.2.3), while the distributions obtained from thermal  
 814 emission at 4610 nm are given separately for water ice (in blue) and liquid water (in red)  
 815 pixels. The heights of the main distribution peaks are highlighted by vertical dashed lines.  
 816 Even if the distributions are evaluated in 0.5 km altitude bins, an uncertainty of the order of 1  
 817 km must be considered.

818

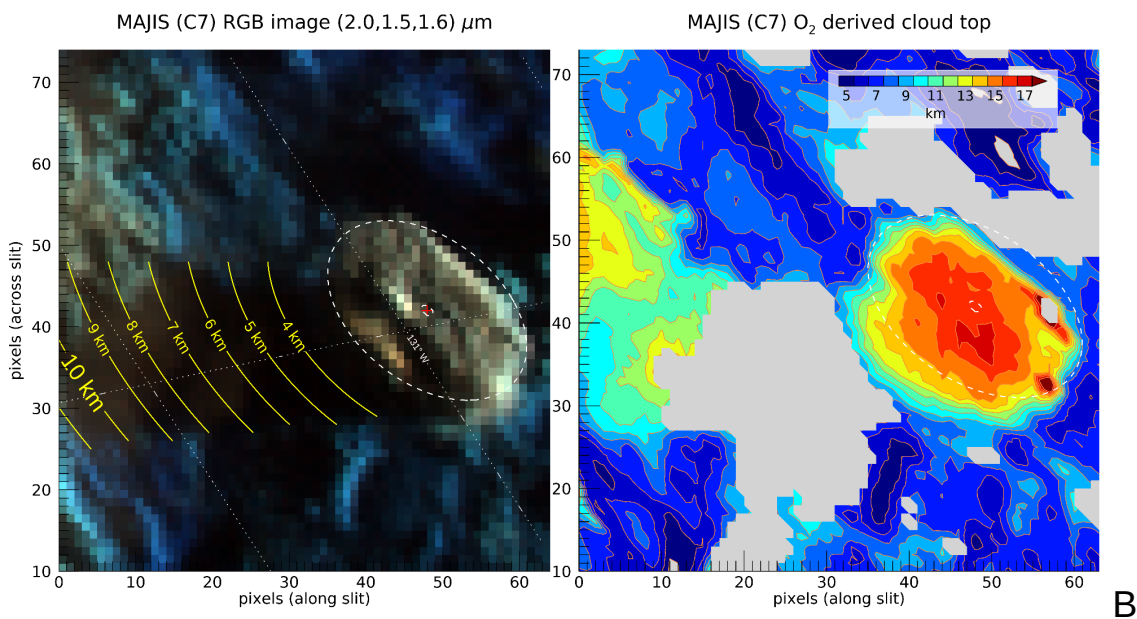
819

#### 4.2.2. Altitudes from clouds' shadows

820

821 An example of the results obtained from the method described in Sect. 3.3.2 is given in Fig.  
 822 16, where the shadows projected by high convective anvil clouds are clearly visible in MAJIS  
 823 data cube C7. The grazing illumination of the scene (incidence angle  $\sim 80^\circ$ ) enables a  
 824 vertical resolution of  $\sim 0.7$  km, inferred from uncertainties of  $\sim 0.5^\circ$  on incidence angles and  
 825 2.7 km on shadow length (about twice the horizontal spatial resolution). Within this

826 framework, the horizontal length of the shadow translates to a top altitude of about 10 km  
 827 (see yellow lines). Of course this value is not absolute but only an estimate relative to the  
 828 surrounding decks, whose altitudes can be qualitatively inferred through the estimation of the  
 829 O<sub>2</sub> 760 nm band depth (see previous section). The O<sub>2</sub>-derived elevations are shown in the  
 830 map of Fig. 16B, where the background structures appear to be located around 5 - 8 km,  
 831 while the anvil cloud top peaks at ~ 16 km. This implies a differential height of ~ 10 km  
 832 between the anvil and the surrounding clouds, in very good agreement with the estimated  
 833 shadow length. The absolute height of the cloud top can only be derived if multiple scattering  
 834 effects are accounted for in the reproduction of the 760 nm O<sub>2</sub> band (Sect. 4.2.4).  
 835 Nevertheless, the shadow analysis provides a quick and independent way for estimating the  
 836 relative height of isolated structures with respect to their background.



837 **A**  
 838 **Figure 16:** A): Example of cloud top altitude estimation based on projected shadow length in  
 839 the MAJIS data cube C7. The white dashed line indicates the approximate boundary of a  
 840 detached cloud (center indicated by the red dot). The yellow lines show how long the  
 841 expected shadow would be in the actual geometry by changing the cloud top altitude. The  
 842 shadow length observed in the background image matches a cloud about 10 km tall. B):  
 843 Cloud top altitudes retrieved in the same area from the O<sub>2</sub> 760 nm band (see Sect. 4.2.2),  
 844 shown for comparison. Gray-filled patches correspond to areas where no O<sub>2</sub> band is  
 845 measurable.

846

#### 847 4.2.3. Altitudes from ice temperature

848

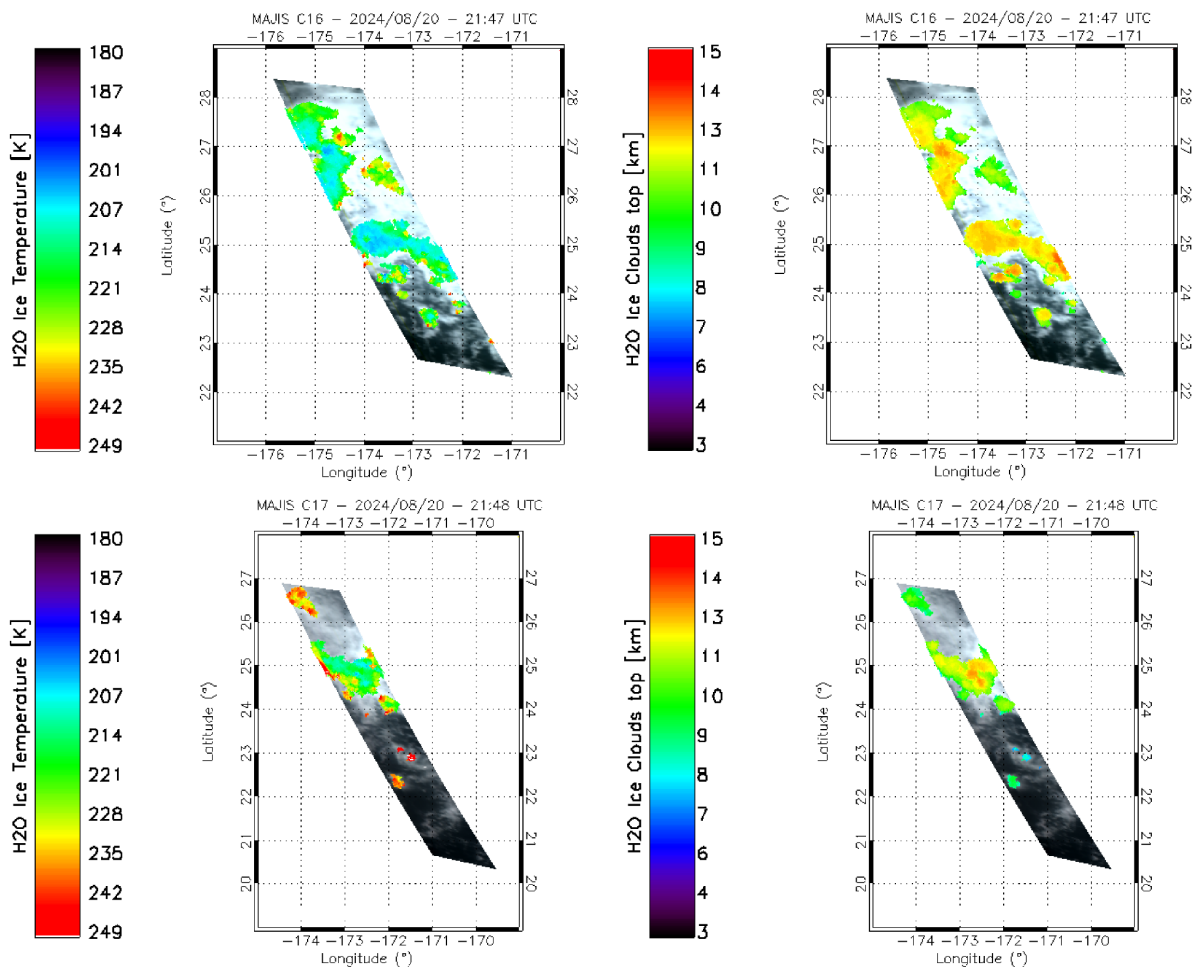
849 In Fig. 17 we show two examples of the temperature and altitude maps derived with the  
 850 method described in Sect. 3.3.3, for MAJIS cubes C16 (upper panels) and C17 (lower  
 851 panels).

852 Altitudes are derived by assuming that the clouds are in thermal equilibrium with the  
 853 surrounding air and reside within the troposphere, where the temperature vertical lapse rate  
 854 is positive. Altitudes' errors are of about 1 km (Sect. 3.3.1) while those related to  
 855 temperatures are propagated from the 3700 nm peak uncertainties (Table 4) and result of  
 856 about 1 K. Orbit C16 shows two main decks, placed respectively at  $z \sim 13$  km and  $z \sim 10$  km  
 857 which can be compared with the maps in Fig. 12C and E, where the 1650 nm band depth

858 and K-means clusters are shown. The higher deck at  $z \sim 13$  km correlates with the blue  
 859 cluster in Fig. 12C and the green one in Fig. 12E, suggesting increased opacity and  
 860 crystallinity at lower temperatures.

861 Similarly, two regimes of temperatures and altitudes are found in orbit C17, with higher  
 862 clouds at  $z \sim 13$  km ( $T \sim 205$  K) and lower ones at  $8 < z < 10$  km ( $215 < T < 250$  K). As  
 863 suggested by the scatterplot in Fig. 13D, these two decks are characterized by different ice  
 864 properties. Indeed, the short wavelength Fresnel peak cluster (i.e. reduced crystallinity,  
 865 Cartwright et al., 2025) shows a larger spread of temperatures, consistent with the lower  
 866 clouds discussed here. Instead, the long wavelength Fresnel peak cluster shows overall  
 867 lower thermal radiances, and hence temperatures, in agreement with the higher clouds  
 868 identified at  $\sim 13$  km (see also Poulet et al., this issue).

869



870

871 **Figure 17:** ice temperature (left) and inferred cloud altitude (right) mapped on MAJIS cube  
 872 C16 (upper panels) and C17 (lower panels). Ice is identified with a threshold  $< 1$  on the ice  
 873 clouds condition in Table 3.

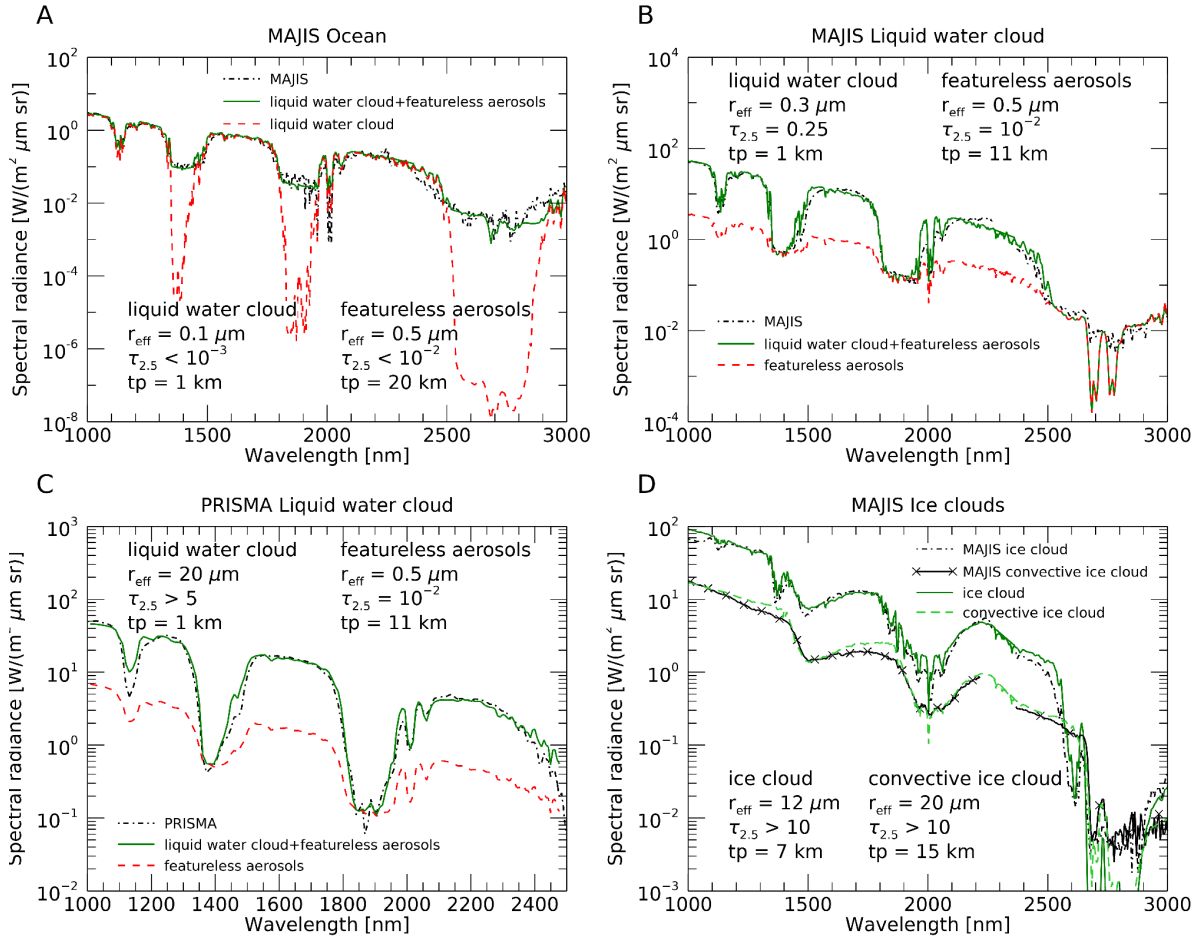
874

#### 875 4.2.4. Results from RT modeling

876

877 For our forward RT modeling (Sect. 3.3.4) we consider all MAJIS spectra and the PRISMA  
 878 liquid water cloud one from Fig. 3, as it is the one showing the most evident differences with

879 respect to its MAJIS counterpart. We also take into account a MAJIS ice cloud spectrum  
 880 related to one of the convective structures identified in Fig. 11A-C and studied in Sect. 4.2.2.  
 881



882

883 **Figure 18:** Panel A: MAJIS ocean spectrum from Fig. 3 is shown in dashed-dot black, its  
 884 forward RT fit is shown in green, while the contribution from the liquid water cloud in the  
 885 simulation is given in dashed red. Geometrical and microphysical parameters ( $r_{\text{eff}}$  is the  
 886 effective radius in  $\mu\text{m}$ ,  $\tau_{2.5}$  is the optical depth at  $2.5 \mu\text{m}$  and  $tp$  is the cloud top in km) of  
 887 aerosols involved in the fit are given in the figure. Panels B-C: same as in panel A, but liquid  
 888 water clouds observations by MAJIS and PRISMA from Fig. 3 are respectively fit. The  
 889 dashed red lines here refer to the contribution from the featureless aerosols in the model (i.e.  
 890 when no liquid water cloud is considered). D: MAJIS ice clouds forward RT fits (green and  
 891 dashed light green lines) related to MAJIS ice cloud spectrum from Fig. 3 (dashed-dot black)  
 892 and to a spectrum from the convective cloud identified in Fig. 11C (black line with crosses).  
 893

894 The best fits obtained with this approach are shown in Fig. 18. In general, grain sizes and  
 895 clouds' altitudes determine the shape and the signal of water absorption bands, while the  
 896 number density can be tweaked to match the intensity of the continuum. We assume that the  
 897 clouds are compact in vertical extent and only occupy a single layer of the atmospheric  
 898 profile. The ocean and liquid water clouds observations require two separate layers placed  
 899 at different altitudes in the atmosphere (Fig. 18A, B and C) suggesting that, as explained in  
 900 Sect. 2.3.1 and 3.1, also the ocean spectra we are investigating are partially obstructed by  
 901 non-resolved cloudy structures. The lower layer shapes the shoulders of water bands', in  
 902 which the atmospheric transmission is enough to probe down to the surface, while the upper

903 one is needed to correctly model the intensity of the bands' bottom. Indeed, if optically thick  
904 enough, high clouds prevent solar photons from reaching the underneath atmospheric  
905 layers, hence reducing the gaseous absorption. Such a differential effect in the models is  
906 shown as dashed red lines in Fig. 18A, B, C. In the ocean spectrum (Fig. 18A) the optically  
907 thin bottom layer ( $z = 1$  km,  $\tau < 10^{-3}$ ) with small grain sizes ( $r_{\text{eff}} = 0.1$   $\mu\text{m}$ ) is consistent with  
908 the average properties of maritime droplets ( $0 < z < 2$  km,  $5 \times 10^{-4} < \tau < 10^{-3}$ ,  $0.05 < r_{\text{eff}} < 1.5$   
909  $\mu\text{m}$ ) commonly observed above the surface of the ocean (Croft et al., 2021; Smirnov et al.,  
910 2002; Heintzenberg et al., 2000). On the other hand, the upper thin layer ( $\tau < 10^{-2}$ ) has  
911 slightly larger particles ( $r_{\text{eff}} = 0.5$   $\mu\text{m}$ ) and is placed at 20 km, in agreement with the presence  
912 of stratospheric background aerosols ( $15 < z < 25$  km,  $10^{-4} < \tau < 10^{-3}$ ,  $0.1 < r_{\text{eff}} < 1$   $\mu\text{m}$ ,  
913 Voudouri et al., 2023; Thomason et al., 2008). Such a configuration confirms the observation  
914 as a partially obstructed scenario.

915 The selected MAJIS and PRISMA liquid water clouds observations (Fig. 3 and Fig. 18B-C)  
916 show a good radiometric agreement but differences in water bands' shape that can be  
917 explained by changes in the aerosols' microphysical properties. Both observations are  
918 characterized by a high altitude, spectrally featureless, thin aerosol layer ( $z = 11$  km,  $\tau \sim 10^{-2}$ )  
919 that is required to reproduce the bottom of water bands. This indicates the presence of faint  
920 background stratospheric aerosols residing at the tropopause. Instead, the lower liquid water  
921 layer ( $z = 1$  km) is thin with small grains in the MAJIS case ( $\tau = 0.25$ ,  $r_{\text{eff}} = 0.3$   $\mu\text{m}$ )  
922 suggesting spray marine boundary layer aerosols (Sun et al., 2023; Zheng et al., 2018; Luo  
923 et al., 2014), and thicker with large grains in the PRISMA case ( $\tau > 5$ ,  $r_{\text{eff}} = 20$   $\mu\text{m}$ ),  
924 consistent with the presence of stratus clouds (Fu et al., 2022; Rossow and Shiffer, 1999;  
925 LeMone, 1988). Hence, different properties ensure the modeling of flatter (MAJIS) and  
926 sharper (PRISMA) bands in the two observations.

927 The two ice observations (Fig. 18D) are reproduced with a single cloud layer and do not  
928 require the lower one. This is because the ice clouds in the models have opacities so high ( $\tau$   
929  $> 10$ ) that they prevent observing the ocean and the atmospheric layers in between. In such  
930 conditions, the ice cloud in practice acts as a surface with high albedo, accounting for most  
931 of the spectral features in the observations. However, two different clouds' observations are  
932 considered here. The first one (black dashed-dot line in Fig. 18D) is related to a small  
933 structure identified around longitude  $133^\circ$  and latitude  $22^\circ$  in Fig. 11C. This cloud can be  
934 modelled with ice crystals of the order of  $10$   $\mu\text{m}$  in radius (green line). The altitude can be  
935 reliably tweaked by studying the depth of gaseous water absorption bands at  $1380$  nm and  
936  $2600$  nm, both identifiable in the observation. This means that the ice cloud is low enough to  
937 ensure some water absorption, before completely shielding the underneath atmosphere. As  
938 a result, our estimate is that it has its top at  $7$  km. These parameters suggest compatibility  
939 with the presence of a thick cirrus cloud ( $6 < z < 13$  km,  $\tau > 3$ ,  $10 < r_{\text{eff}} < 60$   $\mu\text{m}$ , Baran,  
940 2009; Zhou et al., 2017; LeMone, 1988).

941 The other ice cloud (black line with crosses in Fig. 18D) is selected on the larger convective  
942 structure identified in Fig. 11C. We already expect this to be higher in the atmosphere with  
943 respect to the other one (Sect. 4.2.2). Our model (dashed light green line) suggests that it is  
944 characterized by larger crystals ( $20$   $\mu\text{m}$ ) and reaches an altitude of at least  $15$  km, enough to  
945 prevent water absorption in the  $1380$  and  $2600$  nm bands (the model sensitivity to higher  
946 altitudes is reduced making this estimate a lower limit). These values indicate that in this  
947 observation MAJIS is probing the upper frozen top of a large convective cloud ( $8 < z < 16$   
948 km,  $\tau > 10$ ,  $10 < r_{\text{eff}} < 60$   $\mu\text{m}$ , Dolan et al., 2023; Krisna et al., 2018; van Diedenhoven et al.,  
949 2018).

950

951

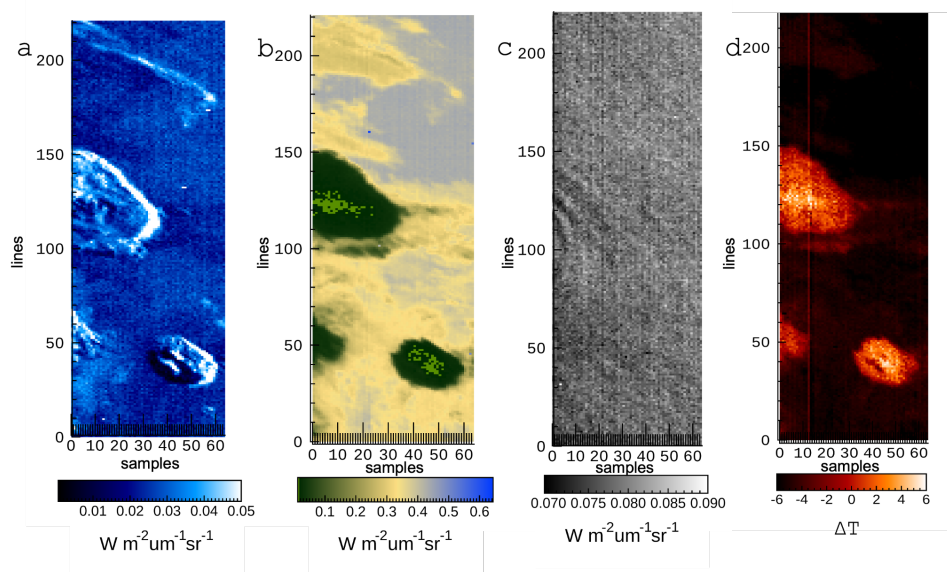
### 952 4.3. Upper atmosphere features

953

954 The CO<sub>2</sub> and O<sub>3</sub> emissions introduced in Sect. 3.4 have been studied in all MAJIS cubes,  
955 deriving maps like those shown in the examples of Fig. 19 and Fig. 20. In Fig. 19, panels A  
956 and B show MAJIS cube C7 displayed at 3100 and 4512 nm, whose anti-correlation  
957 highlights the presence of the convective clouds discussed in Sect. 4.1, 4.2.2 and 4.2.4.  
958 Panels C and D, instead, show the radiance of the peak of CO<sub>2</sub> emission at 4270 nm and the  
959 brightness temperature difference between the O<sub>3</sub> emission peak and its continuum (Sect.  
960 3.4). It is evident how wavy patterns can be seen in the CO<sub>2</sub> map and are uncorrelated with  
961 the clouds beneath. No wave patterns are spotted from the O<sub>3</sub> emission, whose positive  
962 values (and hence the emission) are only detectable above the convective structures. This  
963 suggests that, while both phenomena are likely happening above the clouds' top, waves are  
964 generated at different altitudes with respect to those pertaining to the O<sub>3</sub> emission. However,  
965 the actual heights are not investigated here, since a rigorous retrieval accounting for  
966 non-LTE effects (required for the assessment of these high-altitude emissions) is beyond the  
967 scope of the paper.

968

969



970

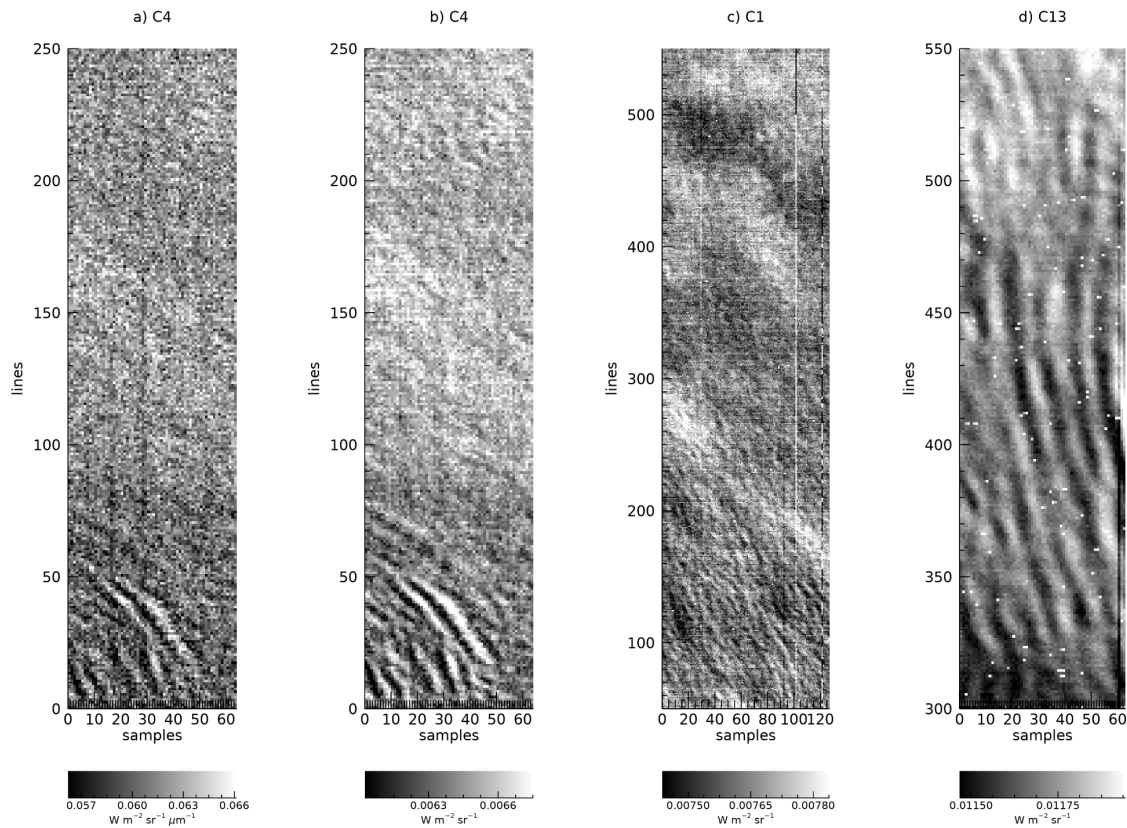
971

972 **Figure 19.** A-B: MAJIS cube C7 radiances at 3100 nm and 4512 nm respectively,  
973 highlighting the anti-correlation between enhanced ice content (A, i.e. larger reflectances of  
974 the Fresnel peak) and low thermal contribution (B). C: radiance of the CO<sub>2</sub> emission peak at  
975 4270 nm, in which the gravity wave pattern is identified. D: brightness temperature difference  
976 (in K) between the O<sub>3</sub> emission peak (4717 nm) and its continuum (4660 nm), showing  
977 positive values above the clouds. In all maps, 'samples' and "lines" indicate spatial pixel  
978 numbers in the direction along and across the instrument slit respectively.

979

980

981



982

983 **Figure 20.** A circular wave pattern is clearly observed in MAJIS C4 cube at 4270 nm (panel  
 984 a). Panel b shows the enhanced contrast achievable after spectral integration between 4254  
 985 and 4333 nm, which also improves detection of complex wave patterns in several MAJIS  
 986 observations, like in C1 (panel c) and C13 (panel d). Pixel scales are reported in Table 1.

987

988 Following the discussion in Sect. 3.4, in Fig. 20 we show the effect of the increased contrast  
 989 that can be achieved through the spectral integration of the CO<sub>2</sub> emission (right panel), with  
 990 respect to the single wavelength investigation (left panel). The integration reduces noise  
 991 hence allowing enhanced accuracy in detecting the wave patterns. Indeed, if the radiance  
 992 integrated in the band is considered, the detectable relative intensity drops from 1% to about  
 993 0.5%, which translates as an increased capability in characterizing the vertical structure of  
 994 the waves.

995

#### 4.3.1. Atmospheric waves properties

996

997  
 998 Examples of wavy structures identified in the MAJIS images at 4270 nm are provided in Fig.  
 999 20. The wave packets have characteristics different from one image to the other in terms of  
 1000 orientations and horizontal wavelengths. In some cases, a curved wavefront is observed  
 1001 (see Fig. 20 B, C, D) as well as a superposition between different packets (Fig. 20 D).

1002

1003

ID	Latitude (deg)	Packet length (km)	Packet width (km)	Horizontal wvl (km)	Azimuth (deg)
C1	9-10	157.6	36.1	27±7	163

C2	10-14	155	135.2	20±6	160
C4	20.85	107.9	94.7	21±6	162
C5	17.7-18.4	154.1	159.1	16±5	33.5
C6	22.9	74.5	94.8		133
C7	23.4-25.5	84.5	73.8	15±6	155
C13	19-22	134.6	88.1	24±8	123
C16	25-27	174.5	131	28±11	119

1004 **Table 5:** Summary of atmospheric waves parameters (packet length and width, horizontal  
1005 wavelength and azimuth) calculated from MAJIS data analysis. Columns indicate: image  
1006 cube, latitude (deg), packet length (km), packet width (km), horizontal wavelength (km),  
1007 azimuth (deg, see Sect. 3.4.1), respectively.

1008

1009 The values obtained from the method described in Sect. 3.4.1 are provided in Table 5. In the  
1010 observed waves, the measured wavelengths are in the range ~ 15-40 km, which can be  
1011 considered as short wavelengths. Similar waves can be generated by several sources and  
1012 are usually observed in the stratosphere. According to models, deep convection is the  
1013 principal source of forcing (Fovell et al. 1992; Piani et al. 2000; Lane et al. 2001) and is also  
1014 suggested to be responsible for circular wave fronts (alongside isolated thunderstorm  
1015 events, e.g. as observed from the Midcourse Space Experiment, Dewan et al. 1998).  
1016 Another source of gravity waves, related to wind flow over mountains, is orography (Fritts  
1017 and Alexander 2003; Kim et al. 2003). Depending on the topography, this can generate  
1018 waves with horizontal scales from a few to hundreds of kilometers (Nastrom and Fritts, 1992;  
1019 Dornbrack et al. 2002; Eckermann et al. 2007). However, as the majority of MAJIS EGA  
1020 observations occurred above open sea areas, a possible origin related to a thunderstorm  
1021 seems to be more realistic.

1022 For circular waves, we estimate the packets' properties and the time of occurrence of the  
1023 related thunderstorms (see Sect. 3.4.1). We assume storms occurring at an altitude of 15 km  
1024 (from thermal brightness estimations) and consider cubes C7 and C4 as examples. The  
1025 minimum and maximum radii, along with the expansion speed and wavelength derived from  
1026 the images, are as follows: for cube C7, these parameters are respectively 35 km, 50 km,  
1027 about 45 km/h and 15 km; for cube C4 they are 20 km, 110 km, about 100 km/h and 20 km.  
1028 In both cases, the thunderstorm-triggering events appear to occur approximately one hour  
1029 before the corresponding observations. This is compatible with the NASA Worldview archive,  
1030 where several thunderstorms have been registered over the areas observed by MAJIS at  
1031 around 05:00 local time. In particular, the wave detection in MAJIS C4 acquisition is located  
1032 about 80 km far from the coastline, and no significant orographic features are present along  
1033 the apparent direction of propagation. For this detection, the hypothesis of  
1034 thunderstorm-generated waves is also strengthened by the intense electrical activity  
1035 confirmed in D'Aversa et al. (this issue), where a lightning event has been detected in the  
1036 visible range of MAJIS cube C1 through the identification of neutral atomic oxygen and  
1037 nitrogen emission lines.

1038

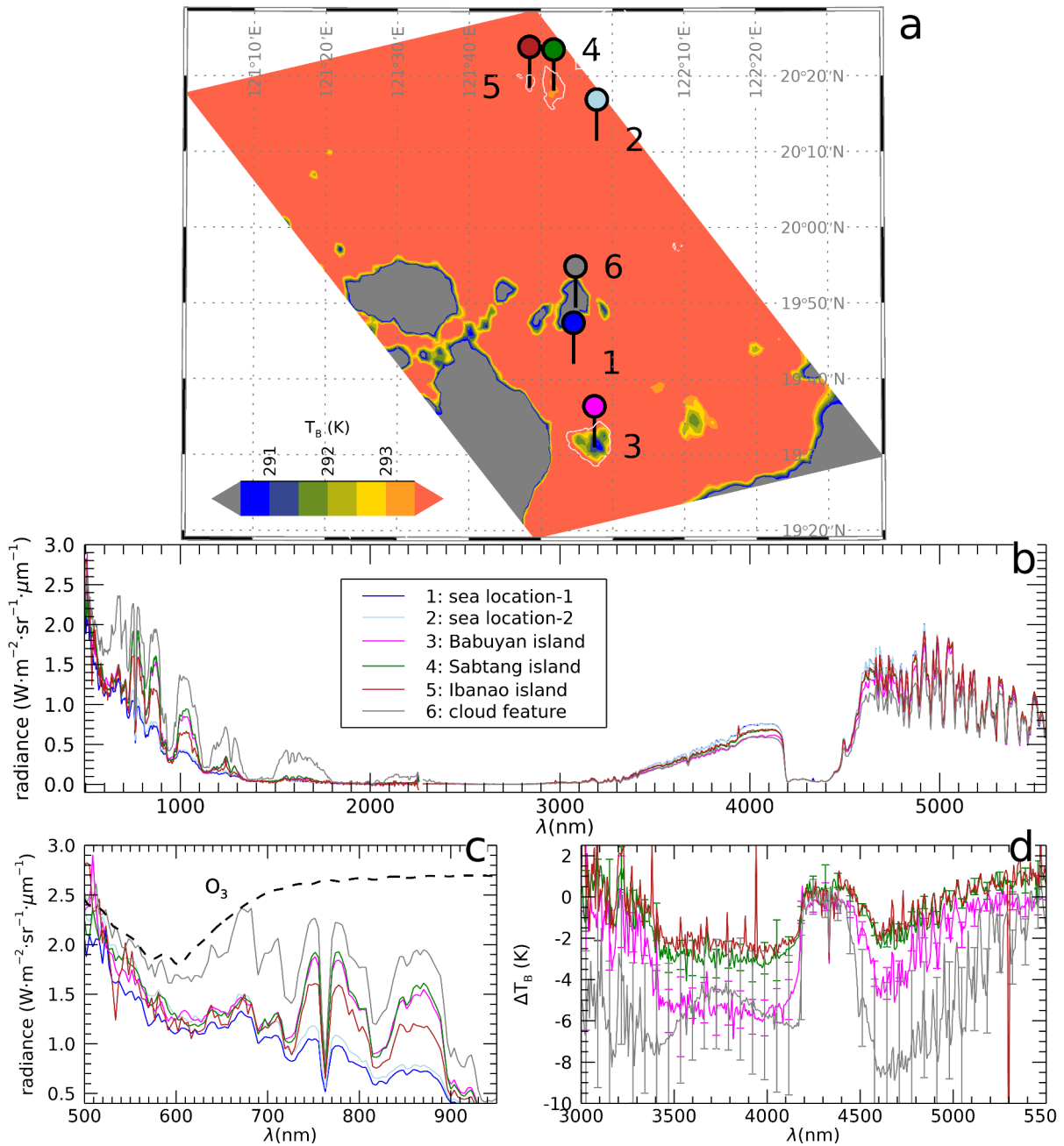
1039

#### 4.4. Land features

1040

1041 The land/ocean-contrast detection method described in Sect. 3.1 has been applied to all  
1042 MAJIS cubes, but only a few land features have been identified. The C1 and C2 cubes,  
1043 expected to cover large land areas at nighttime, encountered very thick and extended storm  
1044 systems that prevented any surface visibility. Hence, all observable land regions consist of  
1045 small islands seen in twilight illumination, colder than the surrounding sea surface but barely  
1046 observable at visible wavelengths. Besides the largest example (Fig. 7), other islands are  
1047 found in the cube C5 (Fig. 21): Babuyan (region 3), Sabtang (region 4), and the very small  
1048 Ibahos island (about 4 x 2.5 km wide, region 5), all part of the Batanes archipelago. The  
1049 nearby Dequey island, even smaller (~0.7x1 km), remains unresolved. With respect to the  
1050 ocean, the brightness temperatures measured over land and cloud areas (Fig. 21b) are  
1051 colder, with differences up to ~6 K and ~8 K respectively. Even if fully located beyond the  
1052 terminator (solar incidence angle ~90.8°), a significant signal is detectable also at visible  
1053 wavelengths, ascribable to light scattering in the upper illuminated portion of the atmospheric  
1054 column, and to multiple scattering effects in the lower part.

1055



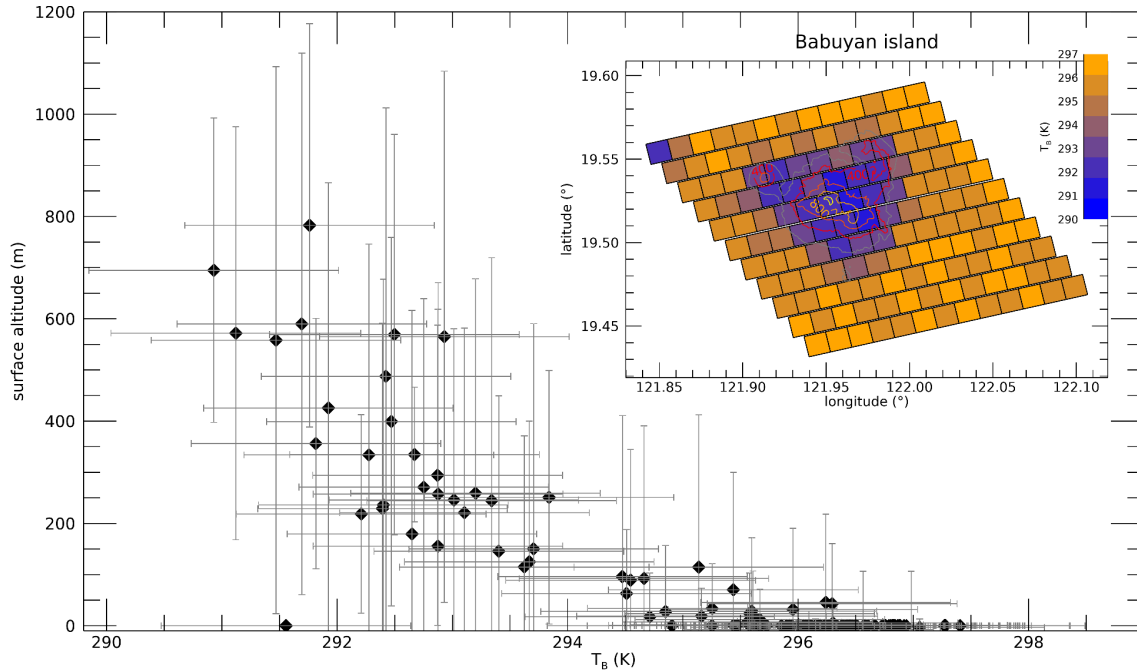
1056

1057 **Figure 21:** Land spectral features seen in twilight conditions in MAJIS cube C5. **a)** Regions  
 1058 of interest (ROIs, labeled 1 to 6) are selected over a 4610 nm brightness temperature map.  
 1059 The coldest areas (in gray) are identifiable as thick clouds, while land areas are slightly  
 1060 warmer (islands of Babuyan, Sabtang and Ibaños, regions 3, 4, 5 respectively), but still  
 1061 colder than the surrounding ocean (red-orange area). Coastlines, obtained from  
 1062 OpenStreetMap under the Open Database License, are shown as white lines. **b)** MAJIS  
 1063 full-range spectra over the ROIs. **c)** Blow-up of the visible spectral part, showing  $H_2O$  and  $O_2$   
 1064 absorption bands as well as a broad  $O_3$  absorption (see also Fig. 4). **d)** Blow-up of the  
 1065 infrared spectral part given as  $T_B$  difference with respect to the ocean spectrum.

1066

1067 The MAJIS sensitivity to temperature variations can be estimated from the signal fluctuations  
 1068 over cloud-free ocean regions. The resulting uncertainties in thermal brightness (at 4610 nm)  
 1069 vary between 0.5 and 1 K, which correspond to about 0.2% and 0.4% of the radiance at

1070 293K. This sensitivity appears sufficient to discriminate significant temperature variation not  
 1071 only between sea and land surfaces but also between different land regions. As an example,  
 1072 we show in Fig. 22 the variability of MAJIS brightness temperature inside the Babuyan  
 1073 island, which hosts a volcano of about 1 km in elevation (Babuyan Claro Volcano). Even if  
 1074 the spatial resolution is limited, a clear trend emerges with respect to the topographic  
 1075 altitude, suggesting that the MAJIS data are sensitive to the surface altimetric temperature  
 1076 change.



1077

1078 **Figure 22:** Thermal analysis of Babuyan island, as viewed in MAJIS data cube C5. The  
 1079 MAJIS-derived brightness temperature (at 4610 nm) is plotted against topographic altitude,  
 1080 stressing the detection of surface altimetric temperature change. Error bars on the x axis  
 1081 derived from signal fluctuation over sea surface around the island, while those on y axis  
 1082 represent the variability of surface altitude inside individual MAJIS pixels. Unit emissivity has  
 1083 been assumed everywhere. Topographic data are extracted from Google Earth Pro  
 1084 7.3.6.10441 (accessed September, 03, 2025).

1085

## 1086 5. Application to Jovian system science

1087

1088 This flyby represents the first acquisition of planetary data by MAJIS. Although the analysis  
 1089 presented here has been dedicated to Earth science, we can briefly identify and discuss  
 1090 different links to the MAJIS science that is foreseen at Jupiter and its icy satellites,  
 1091 highlighting the instrument capabilities in exploring different objects of the solar system.

1092

### 1093 5.1. From ice clouds to icy surfaces

1094

1095 The detection of terrestrial ice clouds described in Sect. 4.1 represents the first spectral  
 1096 observations of water ice performed by MAJIS, and is therefore the first approach to  
 1097 establish the potential outcomes from observations of Jovian icy satellites, in particular for  
 1098 Callisto and Ganymede.

1099 The investigation of ice properties possibly provides information on the differential evolution  
1100 these bodies underwent in the Jovian system environment. For example, Callisto's surface is  
1101 mainly covered by crystalline ice, while significant amorphous ice patches have been  
1102 observed on Ganymede (e.g. Tosi et al., 2024, Bockelée-Morvan et al., 2024; Cartwright et  
1103 al., 2024). These regions could indicate alteration through radiolysis induced by the  
1104 impinging of charged particles on the ice (Khurana et al., 2007), hence providing information  
1105 on the mechanisms connecting Jupiter's magnetic field lines and the moons' surfaces.  
1106 Moreover, while Callisto is characterized by an overall low ice content on the surface (~  
1107 50%) and presents a more ancient and stable scenario (Greeley et al., 2007), Ganymede's  
1108 fresh ice patches are indicative of more frequent ice resurfacing and cryo-volcanism events  
1109 (Ligier et al., 2019). Smaller ice crystals are observed at the poles, matching the distribution  
1110 of the fresher ice deposits and hence acting as a tracer of geologic activity. In this view, the  
1111 investigation of ice-related spectral parameters can be used to address many scientific goals  
1112 of the JUICE mission (Stephan et al., 2021a; Poulet et al., 2024a).

1113

## 1114 **5.2. Clouds**

1115

1116 Jupiter's atmosphere is thought to be dominated by the presence of three main cloud decks  
1117 residing at different altitudes and mixed by convective processes and atmospheric circulation  
1118 (Fletcher et al., 2023). From lower to higher heights these are respectively composed of a  
1119 H<sub>2</sub>O-NH<sub>3</sub> liquid solution, NH<sub>4</sub>SH solid aggregates, and NH<sub>3</sub> ice crystals (Atreya et al., 1999).  
1120 In particular, the NH<sub>4</sub>SH and NH<sub>3</sub> clouds can be responsible for the chromatic differences in  
1121 Jupiter's dark "belts" and bright "zones". Above these structures, hazes composed of  
1122 products of the photochemical disruption of CH<sub>4</sub> and NH<sub>3</sub> extend from the upper troposphere  
1123 to the stratosphere (e.g. Sindoni et al., 2017; Biagiotti et al., 2025). Such cloud complexity is  
1124 not present in Earth's atmosphere where water is the only condensible, aside from a variety  
1125 of aerosols of different origin (e.g. maritime, volcanic, smog, stratospheric). Nevertheless,  
1126 the study of EGA observations allows a first MAJIS data analysis devoted to disentangling  
1127 the spectral information related to different sources, like gases, clouds and, in this case, also  
1128 surfaces. In this manuscript we have investigated clouds under different points of view,  
1129 including their detection, water vapour phase identification, vertical structure assessment,  
1130 and microphysical properties estimation. All these techniques are applicable to Jupiter once  
1131 adapted to the different composition and structure of the giant planet. For example, the RT  
1132 modeling presented in Sect. 4.2.4 only dealt with the solar part of the spectrum, which would  
1133 only allow the investigation of Jupiter's hazes and the NH<sub>3</sub> deck (e.g. the recent work of  
1134 Biagiotti et al., 2025 on JUNO/Jiram data). The exploitation of the full MAJIS spectral range,  
1135 including thermal wavelengths, is instead mandatory for characterizing the deeper NH<sub>4</sub>SH  
1136 (Grassi et al., 2021) and H<sub>2</sub>O (Bjoraker et al., 2022) clouds, especially in "hot spot" regions.  
1137 The shadow technique for measuring cloud heights, commonly applied in planetary  
1138 high-resolution imaging analysis, is also applicable to Jupiter (e.g. Orton et al., 2017). For  
1139 instance, in observations acquired at the bottom of methane bands, Simon et al. (2015) were  
1140 able to measure shadows 45 km long, revealing wavy structures less than 1 km in amplitude.  
1141 In principle, MAJIS observations of Jupiter atmosphere will allow the application of this  
1142 technique to limited cases, mostly near the terminator and in polar regions when observed  
1143 from perijove. Maximum spatial resolutions of ~120 km/px achievable in these conditions  
1144 may enable detecting shadows related to vertical displacements of the order of 10 km.

1145

1146

### 1147 **5.3. High-altitude emissions**

1148

1149 The use of chemical atmospheric species as tracers for the atmospheric circulation,  
1150 including wind measurements and wave detections, is widely applied to the investigation of  
1151 both terrestrial (i.e. Hueso et al. 2008; Peralta et al. 2008) and giant planets (i.e.  
1152 Müller-Wodarg et al. 2019, Grassi et al., 2020). A similar approach is valid for the upcoming  
1153 MAJIS measurements at the Jovian system, whose upper atmospheric dynamical structure  
1154 can be investigated through the monitoring of the distribution (in latitude and local time) of  
1155 minor widespread species like  $\text{H}_3^+$  and hydrocarbons deriving from the photolysis of methane  
1156 (see Miller et al. 2020 for a thorough review) as demonstrated from both ground-based (see  
1157 for example O'Donoghue et al. 2016) and space-based data analyses (e.g. Moriconi et al.  
1158 2020). MAJIS IR channel will allow to spectrally discriminate the  $\text{CH}_4$  and  $\text{H}_3^+$  contributions in  
1159 the range 3000 - 4000 nm, where the two species present strong features (Castagnoli et al.,  
1160 2025) identifiable within the fundamental 3300 nm  $\text{CH}_4$  absorption band, similarly to the case  
1161 of the 4300 nm  $\text{CO}_2$  band in Earth's atmosphere (see Sect. 4.3). The study of  $\text{CH}_4$  and  $\text{H}_3^+$   
1162 (e.g. JWST data analysis, Melin et al., 2024) will give access to upper atmospheric layers  
1163 which are hardly probed otherwise. Altitudes from about 200 km above the 1-bar level are  
1164 typical of methane emission peak, while above 500 km the  $\text{H}_3^+$  emission seems to dominate,  
1165 as also shown by recent analyses of JIRAM-Juno data (Migliorini et al. 2023), where the two  
1166 species have been spatially separated.

1167

### 1168 **6. Summary and conclusions**

1169

1170 In this work we compare the observations of the MAJIS spectrometer on board the JUICE  
1171 spacecraft, acquired during the Earth gravity assist of 2024 (Sect. 1), with those registered  
1172 by the Italian Space Agency-led PRISMA spectrometer (Sect. 2). While no exact  
1173 temporal-spatial coincidence could be achieved, the comparison allowed testing MAJIS  
1174 spectral and radiometric response over ocean and clouds, the main targets observed during  
1175 this flyby. Clouds observations have been analyzed for the estimation of altitudes and  
1176 microphysical properties exploiting different methods (Sect. 4.2). Ice has been detected in  
1177 most of the observations, allowing a first benchmark of the study of its spectral properties  
1178 (Sect. 4.1) in view of Jupiter's icy satellites' exploration.

1179 High-altitude emissions from  $\text{CO}_2$  and  $\text{O}_3$  are also observed in MAJIS dataset, revealing the  
1180 presence of a significant number of atmospheric gravity waves, whose properties have been  
1181 derived (Sect. 4.3).

1182 While we discuss *ad hoc* spectral indices for the identification at VIS-NIR wavelengths of  
1183 different types of surfaces (in view of the next JUICE Earth flyby happening in September  
1184 2026) our investigation of land features is limited to the land/ocean temperature contrast or  
1185 to the changing surface altimetry (Sect. 4.4). Indeed, in the MAJIS 2024 EGA data no land  
1186 areas have been captured in daylight.

1187 This wide variety of scientific applications is finally put in the context of the Jupiter case,  
1188 taking into account the differences between our planet and the gaseous giant's atmosphere  
1189 and icy satellites (Sect. 5).

1190 In conclusion, EGA data provide the first scientific benchmark of MAJIS instrumental  
1191 response in a planetary environment, and give the first glimpse of the amount and quality of  
1192 spectral information we can expect in the Jovian system.

1193

## 1194 Author Contributions

1195 Conceptualization, F.O., E.D., A.Mi.; formal analysis, F.O., E.D., A.Mi.; Data Curation, F.O.,  
1196 E.D., A.Mi., F.P., Y.L., G.P., A.Z., M.G., E.L., G.S., C.P., S.R., B.S.; investigation, F.O., E.D.,  
1197 A.Mi.; methodology, F.O., E.D., A.Mi.; software, F.O., E.D., A.Mi.; supervision, F.O., E.D.,  
1198 A.Mi., G.P., F.P., Y.L., G.F., M.C., M.R., B.S., A.M., L.N.F., A.Z., M.G., E.L., G.S., C.P.;  
1199 validation, F.O., E.D., A.Mi.; writing—original draft, F.O., E.D., A.Mi.; writing—review &  
1200 editing, F.O., E.D., A.Mi., G.P., F.P., L.N.F., A.M.. All authors have read and agreed to this  
1201 version of the manuscript.

## 1202 Code availability

1203 The codes used in this manuscript have been developed by the authors and are available on  
1204 request.

## 1205 Competing interests

1206 The authors declare no competing interests in the production of this manuscript.

1207

## 1208 Acknowledgements & Data availability

1209

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1214 2023-6-HH.0.

1215 The MAJIS data acquired during the JUICE Moon–Earth flyby in August 2024 are currently  
1216 under the mission's cruise-phase proprietary period. These data will be made available  
1217 through the ESA Planetary Science Archive following the first Cruise Archive Delivery, which  
1218 is currently scheduled for six months after Earth Gravity Assist #3 in 2029.

1219 PRISMA products are generated by IAPS-INAF under a license from ASI Original PRISMA  
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