



1           **Process-based evaluation of ENSO simulation sensitivity to**  
2           **horizontal resolution in the Chinese Academy of Sciences FGOALS-**  
3           **f3 Climate System Model**

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19 **Abstract.** El Niño-Southern Oscillation (ENSO) is the most prominent interannual climate variability,  
20 hence its simulation performance represents a critical benchmark for evaluating the fidelity of coupled  
21 climate models. Increasing model resolution is an effective approach to improve the model simulation;  
22 however, the impact of refining horizontal resolution from the hundred-kilometer scale to the tens-of-  
23 kilometer scale on ENSO simulation, as well as the underlying mechanisms, remains unclear. This study  
24 provides a process-based evaluation of ENSO behaviour in two versions of the Chinese Academy of  
25 Sciences Flexible Global Ocean–Atmosphere–Land System Finite-Volume version 3 (FGOALS-f3)  
26 climate system model: a low-resolution configuration (~100 km; FGOALS-f3-L, hereafter f3-L) and a  
27 high-resolution configuration (~25 km; FGOALS-f3-H, hereafter f3-H). Using a reproducible diagnostic  
28 framework, we assess how horizontal resolution influences ENSO amplitude, oscillation characteristics,  
29 key air–sea coupling processes, and high-frequency (HF) atmospheric variability. The low-resolution  
30 severely overestimates ENSO amplitude, whereas f3-H produces amplitude closer to the observation.  
31 Process-based diagnostics show that this improvement arises from the more realistic representation of  
32 thermocline and zonal advection feedback processes in f3-H, which arises from the more realistic  
33 representation of the meridional structure of ENSO-related zonal wind stress anomalies over equatorial  
34 Pacific in f3-H and can be traced back to its improved horizontal resolution. The ENSO cycle in f3-L  
35 exhibits excessive regularity, featuring periodic warm-cold transitions; while f3-H reproduces an  
36 irregular oscillation resembling the observation. The excessive regularity in f3-L is attributed to its  
37 coarser resolution, which limits the simulation performance of tropical cyclones and consequently  
38 weakens high-frequency westerly wind activity over the tropical Pacific. The feeble stochastic forcing in  
39 f3-L is insufficient to disrupt its overly intense ENSO cycle, yielding an overly regular oscillation. By  
40 identifying the structural sources of ENSO biases across resolutions, this study provides a reproducible  
41 and model-agnostic framework for diagnosing resolution effects on ENSO performance in climate  
42 models and informs future development of FGOALS-f3 model family.

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## 44 **1 Introduction**

45 El Niño-Southern Oscillation (ENSO), as one of the most prominent interannual variabilities in the  
46 Earth's climate system, exerts a profound influence on regional and global climate (McPhaden et al.,  
47 2006; Cai et al., 2021). Therefore, the ability to accurately simulate ENSO phenomenon serves as a  
48 fundamental benchmark for evaluating the fidelity of coupled climate models (Timmermann et al., 2018).  
49 Despite significant progress in the development of climate models, large inter-model spread remains in  
50 ENSO characteristics across Coupled Model Intercomparison Project (CMIP) generations (Zhang et al.,  
51 2020; Planton et al., 2021). Common biases include inaccuracies in simulating ENSO amplitude (Planton  
52 et al., 2021), period (Lu et al., 2018), seasonal phase locking (Liao et al., 2023; Yan and Sun, 2024),  
53 spatial distribution (Jiang et al., 2021), and intensity asymmetry (Zhao and Sun, 2022), as well as overly  
54 regular ENSO oscillation (Chen et al., 2016; Guilyardi et al., 2020). These deficiencies reflect systematic  
55 structural biases within models, which directly limit the predictive skill of dynamical models (Barnston  
56 et al., 2012) and hinder the reliability of climate projection regarding how ENSO may evolve under future  
57 climate change (Jiang et al., 2020a). Therefore, it is imperative to improve the fidelity of ENSO  
58 simulation in climate models.

59 Horizontal resolution has long been recognized as an important factor influencing model behavior  
60 (Yu et al., 2024). Motivated by this, the latest CMIP6 launched a dedicated High-Resolution Model Inter-  
61 comparison Project (HighResMIP) (Eyring et al., 2016) to systematically assess the benefits of increased  
62 resolution on model simulations. Previous studies have shown that as model resolution increases,  
63 simulation capabilities for both the climate mean state and variability exhibit improvements (Dawson et  
64 al., 2013; Chang et al., 2020). For example, high-resolution models demonstrate superior performance  
65 over low-resolution counterparts in simulating the Asian summer monsoon (He et al., 2025a), the impact  
66 of Tibetan Plateau thermal forcing on Asian summer monsoon (He et al., 2025b), the precipitation in  
67 southern China (Zi et al., 2024), the spatial distribution and frequency of tropical cyclones (Kreussler et  
68 al., 2021; Li et al., 2021), air-sea turbulent flux (Small et al., 2019), heat transport by boundary currents  
69 (Docquier et al., 2019), ocean mesoscale eddies (Hallberg, 2013), and Antarctic sea ice (Docquier et al.,  
70 2019). This is partly because higher resolution allows models to explicitly resolve finer-scale physical  
71 processes, reducing their dependence on parameterization schemes. Furthermore, high-resolution models  
72 facilitate a more accurate simulation of topographically sensitive regions, therefore improving the



73 simulation of atmosphere-ocean processes related to complex terrain (Bacmeister et al., 2014; Hewitt et  
74 al., 2016). Then a pertinent inquiry arises: does the ENSO simulation benefit from the increased  
75 resolution?

76       Regarding the impact of increased horizontal resolution on ENSO characteristics, previous studies  
77 found that when the atmospheric horizontal resolution in climate models was increased from 3.8° (T30)  
78 or 2.8° (T42) to 1° (T106), significant improvements can be found in the simulated ENSO period  
79 (Guilyardi et al., 2004; Navarra et al., 2008) and ENSO amplitude (Hua et al., 2018). However, the  
80 aforementioned studies were primarily based on the comparisons from ~400 km (or 300 km) to 100 km  
81 scales. In the latest climate models participated in the HighResMIP, model resolutions have advanced  
82 substantially, with some achieving 25 km resolution in atmospheric component and 10 km in oceanic  
83 component. This higher resolution enables a more realistic reproduction of climate variability, such as  
84 tropical cyclones (TCs) (Li et al., 2021) and tropical instability waves (Li et al., 2023). Yet, a critical  
85 question remains unresolved: when model resolution reaches ~25 km, a scale that can reasonably  
86 simulate weather-scale systems like TCs, does ENSO simulation also improve further?

87       ENSO is fundamentally driven by a number of coupled ocean-atmosphere feedbacks (Li, 1997; Jin  
88 et al., 2006; Chen et al., 2015b; Chen et al., 2016); however, some analyses regarding ENSO simulations  
89 are often result-oriented, focusing only on changes in statistical indices while lacking sufficient  
90 diagnostic analysis of the key air-sea feedback processes that shape ENSO's properties. Therefore, it is  
91 necessary to conduct a process-oriented evaluation when comparing the ENSO simulation between  
92 different resolution versions.

93       It is worth noting that the observational evidences have suggested the significant influence of  
94 atmospheric "noise" on the development and evolution of ENSO (Chen et al., 2015a; Fedorov et al.,  
95 2015). Here the atmospheric "noise" primarily refers to high-frequency (HF) wind activities such as  
96 westerly wind bursts (Harrison and Vecchi, 1997; Fedorov, 2002), including synoptic and intra-seasonal  
97 scales. TC and Madden-Julian Oscillation (MJO; Madden and Julian, 1971; Madden and Julian, 1972)  
98 have been recognized as key sources of HF wind activities in recent studies (Ying et al., 2019; Liang and  
99 Fedorov, 2021). Considering the atmospheric components in HighResMIP models can reach a horizontal  
100 resolution of 50 km or finer (which is sufficient to reasonably reproduce TC features) and that the  
101 simulation performance for TC and MJO activities improves with increasing resolution (Davis, 2018;



102 Tang et al., 2022; Roberts et al., 2025), a further question arises: does the improved simulation of HF  
103 activities like TCs and MJO in high-resolution models contribute to improved ENSO simulation?

104 Motivated by these open questions, this study will conduct a process-oriented evaluation of ENSO  
105 simulation for a Chinese climate model, the Flexible Global Ocean–Atmosphere–Land System Finite-  
106 Volume version 3 (FGOALS-f3) climate system model that was developed by the Institute of  
107 Atmospheric Physics, Chinese Academy of Sciences (IAP-CAS). FGOALS-f3 participated in  
108 HighResMIP with both a low-resolution (~100 km; FGOALS-f3-L, hereafter f3-L) and a high-resolution  
109 version (~25 km; FGOALS-f3-H, hereafter f3-H) (An et al., 2022; Bao et al., 2020). The high-resolution  
110 version (f3-H) has been shown to have superior simulation performances over the low-resolution  
111 counterpart (f3-L) in representing mesoscale vortices (An et al., 2022), tropical instability waves (Li et  
112 al., 2023), TCs (Li et al., 2021), East Asian summer monsoon precipitation (Zi et al., 2024), and the  
113 climatological mean states (Yu et al., 2024). Therefore, this study employs the CAS-developed f3-H and  
114 f3-L to reveal the impact of model resolution on ENSO simulation performance and the underlying air-  
115 sea coupling processes.

116 This model evaluation study will provide insights into the resolution sensitivity of ENSO-related  
117 processes in FGOALS-f3 and establishes a diagnostic framework that can be applied to other coupled  
118 models participating in CMIP6 and future CMIP phases. The remainder of this paper is organized as  
119 follows. Section 2 describes the model configurations, observational datasets, and diagnostic framework.  
120 Section 3 presents an overview of the ENSO characteristics in two versions of FGOALS-f3 model.  
121 Section 4 and Section 5 demonstrate how the model horizontal resolution impacts the ENSO simulation.  
122 Section 6 summarizes the findings and discusses implications for model development and evaluation.

## 123 **2. Model configurations, datasets, and diagnostic framework**

### 124 **2.1 Model configurations**

125 FGOALS-f3 is a fully coupled climate system model developed by State Key Laboratory of  
126 Numerical Modeling for Atmospheric Sciences and Geophysical Fluid Dynamics (LASG), IAP-CAS,  
127 which couples four component models using the CPL7 coupler (Craig et al., 2012). The four component  
128 models are the atmospheric model FAMIL2.2 (He et al., 2019; Li et al., 2021), the ocean model  
129 LICOM3.0 (Li et al., 2020), the land model CLM4.0 (Lawrence et al., 2011), and the sea ice model



130 CICE4.0 (Hunke and Lipscomb, 2010).

131 The atmospheric component FAMIL2.2 is the last version of the Finite-volume Atmospheric Model  
132 developed by the LASG-IAP (FAMIL) (Li et al., 2021). FAMIL2.2 utilizes a finite-volume dynamical  
133 core constructed on a cubed sphere grid that is globally partitioned into six tiles (Zhou et al., 2015). In  
134 the vertical direction, the model use hybrid coordinates over 32 layers, and the model top is 1 hPa. While  
135 the horizontal resolution ranges from C96 (about 100 km) to C384 (about 25 km) across the different  
136 resolution version. The oceanic component LICOM 3.0 is the third version of LASG-IAP Climate System  
137 Ocean Model (LICOM) (Yu et al., 2018). LICOM 3.0 updated a new advection scheme and employed a  
138 tripolar grid based on orthogonal curvilinear coordinates. The horizontal resolution of LICOM 3.0 can  
139 vary flexibly between 1° and 1/20°. Sub-grid parametrization schemes employed in LICOM 3.0 include  
140 the tidal mixing scheme, a buoyancy frequency related thickness diffusivity scheme, a vertical viscosity  
141 and diffusion scheme, and a chlorophyll-a dependent solar penetration scheme, etc. A comprehensive  
142 description of the physical package in LICOM3.0 can be found in Li et al. (2020). The land component  
143 in FGOALS-f3 model is Community Land Model (CLM) version 4.0. This advanced model simulates  
144 the water and momentum balances at the land surface and incorporates interactive carbon and nitrogen  
145 cycles, allowing for a more realistic representation of vegetation dynamics and ecosystem processes  
146 (Lawrence et al., 2011). In FGOALS-f3 model, sea ice is simulated using the Los Alamos Sea Ice Model  
147 version 4.0 (CICE 4.0). This is a dynamic-thermodynamic model that simulates the evolution of sea ice  
148 thickness, concentration, and velocity. It features multiple ice thickness categories and an elastic-viscous-  
149 plastic (EVP) rheology to model ice deformation and dynamics (Hunke and Lipscomb, 2010).

150 The FGOALS-f3 model includes two versions: f3-L and f3-H (He et al., 2019; An et al., 2022). Both  
151 models are participating in HighResMIP of CMIP6 and have successfully completed the Tier-1 and Tier-  
152 3 experiments. These two models have the same components and physical processes. The sole distinction  
153 between f3-H and f3-L lies in their horizontal resolution and the corresponding time steps within their  
154 finite-volume cubed-sphere dynamical core (FV3). To maintain the stability of the integration for the  
155 dynamical core, the two parameters ( $k\_split$  and  $n\_split$ ) in FV3 are set to 6 and 15 in f3-H, respectively  
156 (they are 2 and 6, respectively, in f3-L). The specific resolutions of each component of f3-L and f3-H  
157 models are shown in Table 1.

158



159 **Table 1.** Resolution comparison of each component model between the FGOALS-f3-L model and the FGOALS-f3-  
160 H model

	FGOALS-f3-L	FGOALS-f3-H
Atmosphere	FAMIL2.2 (1°×1°, 32 levels)	FAMIL2.2 (0.25°×0.25°, 32 levels)
Ocean	LICOM3.0 (1°×1°, 30 levels)	LICOM3.0 (0.1°×0.1°, 55 levels)
Land	CLIM4.0 (1.25°×0.9°)	CLIM4.0 (0.31°×0.23°)
Sea Ice	CICE4.0 (1°×1°)	CICE4.0 (0.1°×0.1°)

161 The model data used in this study are obtained from the historical experiment outputs of f3-L and  
162 f3-H model. Considering the experiment outputs of f3-H (highresSST-present) start from 1950, the period  
163 during 1950-2014 for the two models are analyzed. The main variables used in this study include monthly  
164 sea surface temperature (SST), three-dimensional ocean currents (uo, vo, wo), oceanic potential  
165 temperature (thetao), surface wind stress (tauu, tauv), and net radiation flux (rsds, rsus, rlds, rlus, hfss,  
166 hfls); daily surface wind (uas, vas), precipitation (pr), six-hourly sea level pressure (SLP), 850 hPa wind  
167 (ua, va), 300 hPa and 500 hPa temperature (ta). All data is detrended before analyzing.

## 168 2.2 Observation and reanalysis datasets

169 The observational and reanalysis data used in this study include: 1) the monthly SST data obtained  
170 from the Hadley Centre Sea Ice and Sea Surface Temperature version 1.1 dataset (HadISST v1,1), with  
171 a horizontal resolution of 1°×1° (Rayner et al., 2003); 2) the monthly, daily, and hourly 10m wind fields  
172 are from the European Centre for Medium-Range Weather Forecasts (ECMWF) reanalysis data Fifth  
173 Generation (ERA5), with a horizontal resolution of 0.25°× 0.25° (Hersbach et al., 2020); 3) the daily  
174 precipitation are from the Global Precipitation Climate Program version 1.3 (GPCP v1.3), with a  
175 horizontal resolution of 2.5°× 2.5° (Adler et al., 2003); 4) the monthly sea surface wind stress data are  
176 provided by the ECMWF Ocean Reanalysis Data (ORAS5), with a horizontal resolution of 1°× 1° (Zuo  
177 et al., 2019) and America Ocean Data Assimilation Data Set version 2.2.4 (SODA v2.2.4), with a  
178 horizontal resolution of 0.5°× 0.5° (Carton and Giese, 2008); 5) the historical TC data are from the China  
179 Meteorological Administration (CMA) tropical cyclone best track dataset (Ying et al., 2014; Lu et al.,  
180 2021). All data is detrended before analyzing.



## 181 2.3 Diagnostic framework

182 This study employs a reproducible, model-agnostic diagnostic framework for evaluating resolution-  
183 dependent ENSO behavior. Diagnostics include ENSO amplitude and spectrum, Bjerknes feedback  
184 decomposition, thermocline and current responses, oceanic zonal current decomposition, and high-  
185 frequency (HF) atmospheric variability.

### 186 2.3.1 ENSO indices and spectral characteristics

187 Monthly anomalies of Niño3.4 (5°S–5°N, 170°–120°W) is computed after removing the  
188 climatological seasonal cycle. Power spectra are derived using the multi-taper method following  
189 Thomson (1982). ENSO irregularity is diagnosed via spectral width and cycle-to-cycle variability.

### 190 2.3.2 Bjerknes feedback decomposition

191 In order to quantitatively analyze the air-sea coupling processes responsible for ENSO amplitude in  
192 FGOALS-f3 models, the diagnostic framework of BJ index (Kim and Jin, 2011a, 2011b) is performed in  
193 this study. The specific formulation is listed as follows:

$$194 \quad BJ = \frac{R-\varepsilon}{2} \quad (1),$$

$$195 \quad R = -\left(\alpha_1 \frac{\langle \Delta \bar{w} \rangle_E}{L_x} + \alpha_2 \frac{\langle \Delta \bar{v} \rangle_E}{L_y}\right) - \alpha_s + \mu_a \beta_u \left\langle -\frac{\partial T}{\partial x} \right\rangle_E + \mu_a \beta_w \left\langle -\frac{\partial T}{\partial z} \right\rangle_E + \mu_a \beta_h a_h \left\langle \frac{\bar{w}}{H_1} \right\rangle_E \quad (2),$$

$$196 \quad \frac{\partial \langle h \rangle_w}{\partial t} = -\varepsilon \langle h \rangle_w - F[\tau_x] \quad (3),$$

$$197 \quad \langle Q \rangle_E = -\alpha_s \langle T \rangle_E \quad (4),$$

$$198 \quad \langle H(\bar{w})T_{sub} \rangle_E = a_h \langle T \rangle_E \quad (5),$$

$$199 \quad \langle h \rangle_E - \langle h \rangle_w = \beta_h [\tau_x] \quad (6),$$

$$200 \quad \langle H(\bar{w})w \rangle_E = \beta_w [\tau_x] \quad (7),$$

$$201 \quad \langle u \rangle_E = \beta_u [\tau_x] + \beta_{uh} \langle h \rangle_w \quad (8),$$

$$202 \quad [\tau_x] = \mu_a \langle T \rangle_E \quad (9),$$

203 where variables with overbar represent the climatological fields and others denote the anomalous fields  
204 obtained by removing the seasonal cycle;  $u$ ,  $v$ , and  $w$  are the three dimensional oceanic current velocity;  
205  $T$  denotes sea surface temperature;  $\langle \rangle_E$  and  $\langle \rangle_w$  represent averaged volume of the eastern box region  
206 and western box region, respectively, from the ocean surface to the mixed layer depth;  $L_x$  and  $L_y$  denote  
207 the longitudinal and latitudinal lengths of the eastern box, respectively;  $H_1$  indicates the constant mixed



208 layer depth (65m). The coefficients  $a_1$  and  $a_2$  are obtained through linear regression using SST anomalies  
209 (SSTA) zonally or meridionally averaged at the boundaries and area-averaged SSTA over the box. the  
210 " $\Delta$ " in equation (2) represents the differences of ocean current between the eastern and western  
211 boundaries. In equation (3), the first term on the right-hand side are the damping process at the rate of  $\varepsilon$ ;  
212 the second term represents the Sverdrup transport across the equatorial Pacific basin. In Equations (4) to  
213 (9),  $\alpha_s$  indicates the response of the thermodynamic damping to SSTA;  $\mu_a$  represents the response of  
214 surface wind stress anomalies ( $\tau_x'$ ) to SSTA;  $\beta_u$  represents the response of upper ocean current  
215 anomalies ( $u_o'$ ) to  $\tau_x'$ ;  $\beta_h$  indicates the response of the anomalous zonal slope of the equatorial Pacific  
216 thermocline to  $\tau_x'$ ;  $\beta_w$  denotes the response of ocean upwelling anomalies ( $w_o'$ ) to  $\tau_x'$ ;  $a_h$  shows the  
217 effect of thermocline depth change on subsurface temperature anomalies. In this study, the computational  
218 region is (5°S–5°N, 130°E–90°W).

219 Based on the above equations, the main contributing terms of the BJ index include two negative  
220 feedback processes: the mean advection (MA;  $-\left(a_1 \frac{\langle \Delta \bar{u} \rangle_E}{L_x} + a_2 \frac{\langle \Delta \bar{v} \rangle_E}{L_y}\right)$ ), the thermodynamic feedback  
221 (TD;  $-\alpha_s$ ), and three positive feedback processes: the zonal advection feedback (ZA;  $\mu_a \beta_u \langle -\frac{\partial \bar{T}}{\partial x} \rangle_E$ ), the  
222 thermocline feedback (TH;  $\mu_a \beta_h a_h \langle \frac{\bar{w}}{H_1} \rangle_E$ ), and the Ekman feedback (EK;  $\mu_a \beta_w \langle -\frac{\partial \bar{T}}{\partial z} \rangle_E$ ).

### 223 2.3.3 HF westerly (easterly) wind index

224 An index regarding the HF zonal wind anomalies is used to quantitatively measure the intensity of  
225 HF westerly (easterly) wind activity in the f3-L and f3-H models. This index (named WWI index and  
226 EWI index) is defined as the integration of the averaged zonal HF westerly (easterly) wind over a fixed  
227 region (5°S–5°N, 120°E–180°). The formula is as follows:

$$228 \quad WWI \text{ index} = \int_{time1}^{time2} u_{HF}' dt, \quad u_{HF}' > 1m/s \quad (10),$$

$$229 \quad EWI \text{ index} = \int_{time1}^{time2} u_{HF}' dt, \quad u_{HF}' < -1m/s \quad (11),$$

230 where  $u_{HF}'$  represents the daily zonal wind anomaly after applying a 90-day high-pass filter, *time1* and  
231 *time2* represent the start and end time of the integrated period, respectively. In this study, the development  
232 period of ENSO event (January to August) is selected as the integrated period.



### 233 2.3.4 Oceanic zonal current decomposition

234 ENSO-related zonal current anomaly ( $u'_o$ ) in the equatorial region is composed of anomalous  
235 zonal geostrophic currents ( $u'_g$ ) and anomalous Ekman currents ( $u'_e$ ) (Su et al., 2010; Su et al., 2014).

236 The specific formulas of  $u'_g$  and  $u'_e$  are listed below:

$$237 \quad u'_g = -\frac{g\partial^2 D'}{\beta\partial y^2} \quad (12),$$

$$238 \quad u'_e = \frac{1}{\rho H} \frac{r_s \tau'_x + \beta y \tau'_y}{r_s^2 + (\beta y)^2} \quad (13),$$

239 where  $\beta$  and  $r_s$  are the planetary vorticity gradient and Rayleigh dissipation rate ( $0.5 \text{ day}^{-1}$ ),  
240 respectively;  $\tau'_y$  denotes the sea meridional wind stress anomaly;  $H$  denotes the mean mixed layer depth,  
241 and  $\rho$  is the seawater density.

### 242 2.3.5 TC detection and metrics

243 The TC detection algorithm developed by the Geophysical Fluid Dynamics Laboratory  
244 (<https://www.gfdl.noaa.gov/tstorms/>; Zhao et al., 2023) was used to detect TC activities in FGOALS-f3  
245 model. The basic identification criteria and steps are: 1) a local minimum SLP is found within a maximum  
246 distance of 3,000 km; 2) the wind speed at 850 hPa for this vortex exceed 17 m/s; 3) the absolute value  
247 of the vorticity is greater than  $1.5 \times 10^{-4} \text{ s}^{-1}$ ; 4) The temperature within 1200 km of the vortex center is  
248 higher than that in the 1200-2400 km radius at 300-500 hPa, indicating a warm core structure; 5) The  
249 lifespan of this TC should be at least three days (72h).

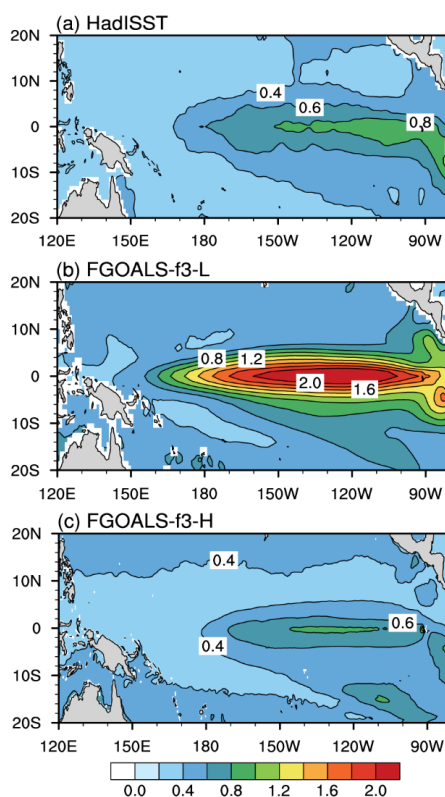
250 We further employ two metrics, i.e., accumulated cyclone energy (ACE) and TC track density  
251 (TCTD), to assess the TC activity in FGOALS-f3 models. The ACE index (Bell et al., 2000) in each grid  
252 cell is defined as the sum of the squares of the maximum surface wind speeds for all TCs occurring within  
253 a  $4^\circ \times 4^\circ$  grid cell over all 6-hourly periods (i.e.,  $ACE = \sum_i V_{max}^2$ , where  $i$  denotes the  $i$ th TC in a grid cell  
254 and  $V_{max}$  denotes its maximum surface wind speed). The TCTD in each grid is defined as the sum of the  
255 number of TCs that have passed through the region within a distance of  $4^\circ \times 4^\circ$  from the grid center.

## 256 3. Evaluation of ENSO Characteristics in f3-L and f3-H

257 Figure 1 shows the spatial distribution of the standard deviation (STD) of SSTA over the tropical  
258 Pacific. In both the high- and low-resolution versions of FGOALS-f3, the interannual variability of SSTA  
259 is concentrated in the central-eastern equatorial Pacific (Fig. 1b-c), which is spatially consistent with the



260 observation (Fig. 1a). However, significant differences exist in the ENSO amplitude. Specifically, the  
261 STD of the Niño3.4 index is 0.69°C in the observation, 1.53°C in f3-L, and 0.62°C in f3-H, respectively.  
262 This indicates that f3-L severely overestimates the ENSO amplitude, whereas f3-H's simulation is much  
263 closer to the observation. Considering that the ENSO amplitude simulated in f3-L is approximately 2.5  
264 times that of f3-H, this study will firstly address the causes of the stronger ENSO amplitude in f3-L  
265 compared to its counterpart in f3-H.

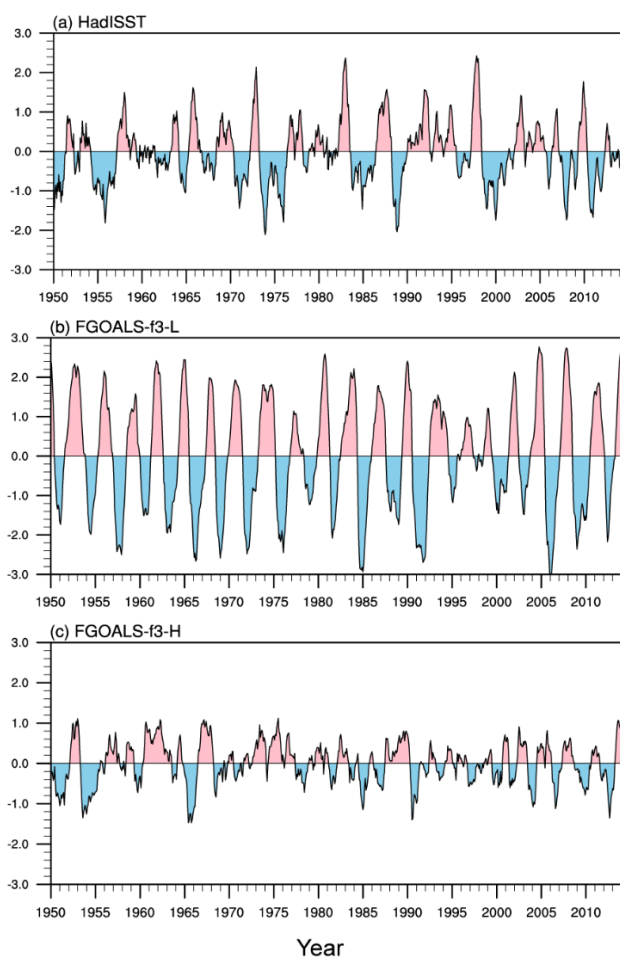


267 **Figure 1.** The standard deviation of SSTA (shading, units: °C) for (a) Observation (HadISST), (b) f3-L and (c) f3-  
268 H.

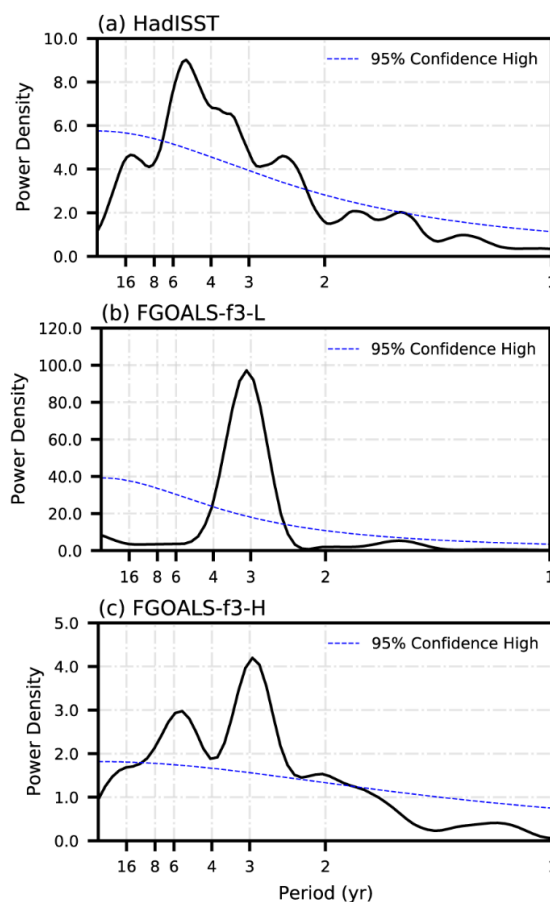
269 Figure 2 presents the time series of the Niño3.4 index for the observation and the two model versions.  
270 Obviously, the oscillation intensity of Niño3.4 index in f3-L is much stronger than that in the observation  
271 and f3-H, consistent with the results of the STD analysis (Fig. 1). Another notable difference lies in the  
272 regularity of the oscillation: ENSO events in f3-L exhibit highly periodic behavior, with a regular  
273 transition between warm and cold phases. In contrast, the ENSO evolution in both the observation and  
274 f3-H exhibit obvious irregularity. Furthermore, we conduct the power spectrum analysis on the Niño3.4



275 indices (Fig. 3). In the observation, the power spectrum is characterized by a broad band of 2–8 years.  
276 However, the power spectrum in f3-L shows a sharp and significant unimodal peak at approximately 3-  
277 year period, with excessive concentration of ENSO energy at this dominant period. In contrast, f3-H  
278 reproduces a broadband spectral distribution similar to the observation. Therefore, the second question  
279 this study will address is: what causes the overly regular oscillation in f3-L, while f3-H captures more  
280 realistic irregularity?



282 **Figure 2.** Temporal evolution of Niño3.4 index (the averaged SSTA in the Pacific Niño3.4 region (5°S–5°N, 170°W–  
283 120°W), units: °C) for (a) Observation (HadISST), (b) f3-L and (c) f3-H.



285 **Figure 3.** The power spectra of Niño3.4 index for (a) Observation (HadISST), (b) f3-L and (c) f3-H. The blue line  
286 in the plots is the 95% confidence level.

#### 287 **4. Process-based Diagnosis of ENSO Amplitude Differences**

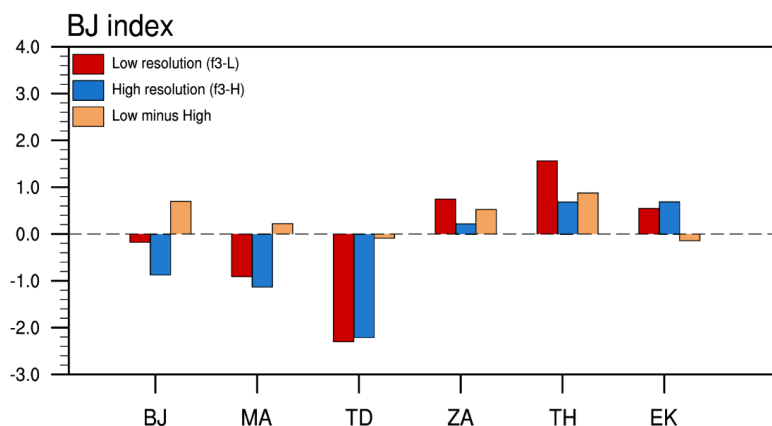
##### 288 **4.1 Diagnostic analysis based on the BJ index**

289 To investigate the physical mechanisms responsible for the ENSO amplitude difference between the  
290 high- and low-resolution versions, we apply the BJ index framework to quantitatively diagnose the  
291 stability of the coupled ocean-atmosphere system in both models. Figure 4 presents the BJ index and its  
292 five contributing terms for f3-L and f3-H, as well as their difference (f3-L minus f3-H). The results  
293 demonstrate that although both models show negative BJ indices, the index value for f3-L is significantly  
294 larger (i.e., a smaller negative value) than that for f3-H model. According to the physical interpretation  
295 of the BJ index (Jin et al., 2006; Kim and Jin, 2011a; Kim and Jin, 2011b), a smaller absolute value of



296 BJ index in f3-L indicates that the coupled air-sea system in f3-L is more unstable than that in f3-H. This  
 297 more unstable coupled system in f3-L is more favorable for ENSO's growth rate, thereby leading to a  
 298 stronger ENSO amplitude in f3-L than that in f3-H.

299 A further question arises: which physical processes contribute to the more unstable coupled system  
 300 in f3-L? By examining the differences in the five contributing terms of the BJ index (orange bars in Fig.  
 301 4), we find that the differences in the thermocline feedback (TH) term and the zonal advection feedback  
 302 (ZA) term are the decisive factors driving the difference in the BJ index (i.e., the system stability)  
 303 between the high- and low-resolution FGOALS-f3 versions. Therefore, the following analysis will focus  
 304 on the physical mechanisms responsible for the stronger TH and ZA terms in f3-L relative to f3-H.

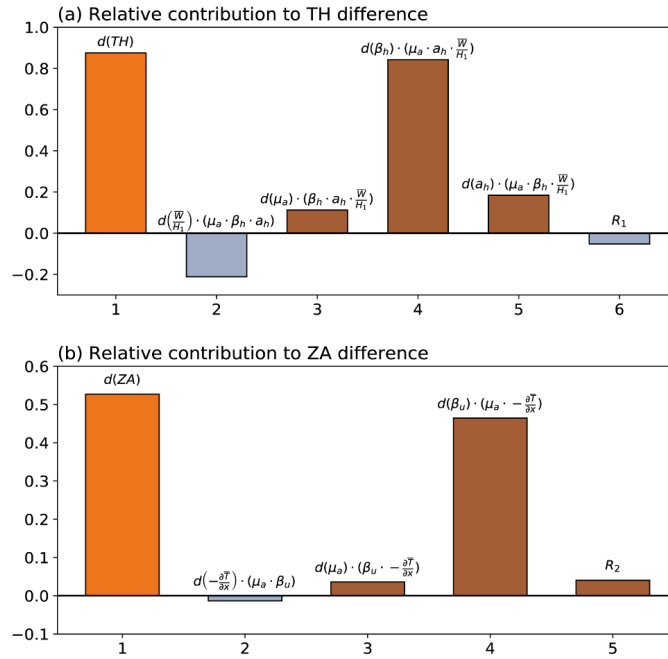


306 **Figure 4.** BJ index and the corresponding main contributing terms for the f3-L (red bars), f3-H (blue bars) and their  
 307 difference (f3-L minus f3-H, orange bars). The five contributing terms including mean advection feedback (MA),  
 308 thermodynamic damping feedback (TD), zonal advection feedback (ZA), thermocline feedback (TH) and Ekman  
 309 feedback (EK).

310 According to the definition of BJ index (see Eq. (2)), each ocean dynamic term consists of two  
 311 components: one related to the mean state and the other to air-sea feedback processes. Considering that  
 312 both components may contribute to the differences in ocean dynamic terms between f3-L and f3-H, we  
 313 first calculate the relative contributions of each component to the total difference in the TH and ZA terms  
 314 using a total derivative decomposition. As shown in Figure 5a, the stronger TH term in f3-L compared  
 315 to f3-H (bar 1) primarily comes from the difference in  $\beta_k$  (bar 4). The differences in  $\mu_a$  and  $a_k$   
 316 feedback processes make minor contributions, while the mean state differences and the covariance term  
 317 make negative contributions. Thus, the dominant factor causing the stronger TH term in f3-L is the  
 318 difference in  $\beta_k$ . Similarly, for the stronger ZA term in f3-L, the results (Fig. 5b) show that the difference



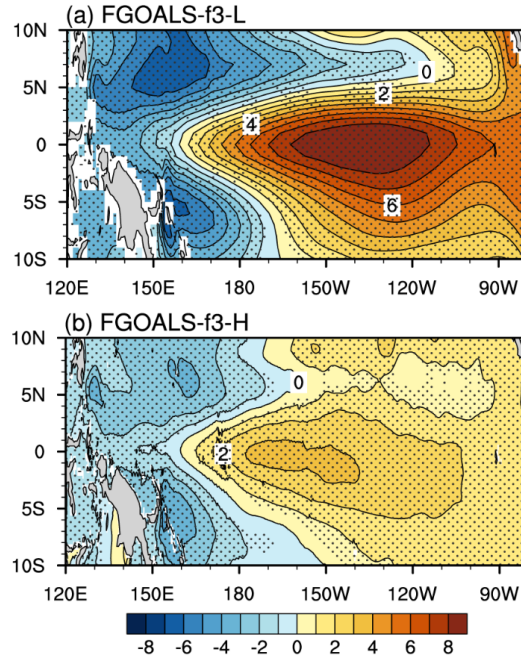
319 in the ZA term (bar 1) is primarily determined by the difference in  $\beta_u$  (bar 4), while the difference in  
 320  $a_h$  and the covariance change have marginal contribution.



322 **Figure 5.** The contribution of each terms in the thermocline feedback (TH). Bar 1 indicate the change of TH term  
 323 [f3-Lminus f3-H,  $d(TH)$ ]. Bar 2 denotes the contribution from the mean state part change [ $d(\bar{w}/H_1)$ ] to the change  
 324 of TH term. Bar 3, Bar 4 and Bar 5 indicate the contribution from the air-sea feedbacks [ $d(\mu_a)$ ,  $d(\beta_h)$  and  $d(a_h)$ ]  
 325 to the change of TH feedback, respectively. Bar 6 is the residual, which denotes the contribution from the covariant  
 326 changes of both the mean state part and the air-sea feedbacks. (b) The contribution of each terms in the zonal  
 327 advection feedback (ZA). Bar 1 indicate the change of ZA term [f3-L minus f3-H,  $d(ZA)$ ]. Bar 2 denotes the  
 328 contribution from the mean state part change [ $d(\partial T/\partial x)$ ] to the change of TH term. Bar 3 and Bar 4 indicate the  
 329 contribution from the air-sea feedbacks [ $d(\mu_a)$  and  $d(\beta_u)$ ] to the change of ZA feedback, respectively. Bar 5 is the  
 330 residual, which denotes the contribution from the covariant changes of both the mean state part and the air-sea  
 331 feedbacks.

#### 332 4.2 Resolution dependence of the $\beta_h$ feedback process difference

333 The air-sea feedback process  $\beta_h$  represents the response of the equatorial Pacific thermocline tilt  
 334 to  $\tau_x'$ . Figure 6 shows the regression of thermocline depth anomalies onto zonal wind stress anomalies  
 335 for the two models. Both models reproduce the expected pattern: in response to positive  $\tau_x'$  over the  
 336 equatorial Pacific, the anomalous thermocline depth ( $D$ ) deepens in the eastern equatorial Pacific and  
 337 shoals in the western equatorial Pacific. However, the response is much stronger in f3-L than in f3-H.



339 **Figure 6.** Distribution of the response of thermocline depth anomaly to zonal wind stress anomaly [ $\beta_h$ ; units:  $\text{m} (\text{N}$   
 340  $\text{m}^{-2})^{-1}$ ] for (a) f3-L and (b) f3-H. Based on the linear relationship between sea surface height anomaly (SSHA) and  
 341 thermocline depth anomaly ( $D'$ ), the SSHA is used as a proxy of  $D'$ . The  $D'$  response to zonal wind stress anomaly  
 342 is derived through regressing the SSHA field onto the averaged zonal wind stress anomaly over equatorial Pacific  
 343 ( $5^\circ\text{S}$ – $5^\circ\text{N}$ ,  $120^\circ\text{E}$ – $100^\circ\text{W}$ ). The stippling in this plot denotes the regression coefficient exceeding a confidence level  
 344 of 99% by using Student's  $t$  test.

345 Based on the Sverdrup balance relationship (Jin, 1997; Li, 1997), the response of the  $D'$  to  $\tau_x'$  is  
 346 primarily determined by the mean equatorial thermocline depth ( $\bar{H}$ ) and  $\tau_x'$ :

$$347 \quad \frac{\partial D'}{\partial x} = \frac{\tau_x'}{\rho g \bar{H}} \quad (14),$$

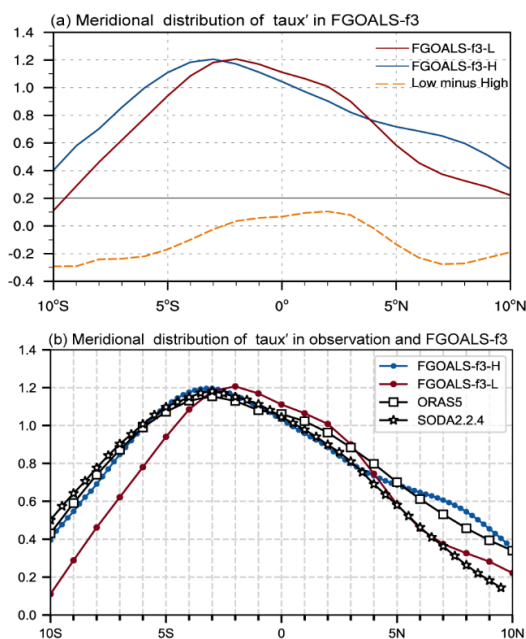
348 where  $\rho$  is the seawater density and  $g$  is the reduced gravity. We first examine the mean thermocline  
 349 depth in both models and find that the difference in  $\bar{H}$  between f3-L and f3-H is negligible (figure not  
 350 shown). Therefore, the difference in  $\bar{H}$  is not the primary cause.

351 Previous studies (Chen et al., 2015b; Chen et al., 2019) have pointed out that the meridional  
 352 structure of  $\tau_x'$  is another key factor influencing the strength of  $\beta_h$ . Figure 7a presents the meridional  
 353 structure of  $\tau_x'$  for the two models and their difference (orange, f3-L minus f3-H). A significant  
 354 difference exists: although the regression results show the same magnitude of  $\tau_x'$  within the equatorial  
 355 Pacific ( $5^\circ\text{S}$ – $5^\circ\text{N}$ ), the  $\tau_x'$  in f3-L is meridionally more concentrated near the equator ( $0^\circ$ ) than that in



356 f3-H. Since the  $\tau_x'$  closer to equator is more effective in influencing  $D'$  (Chen et al., 2015b; Chen et al.,  
 357 2019), this more equatorially-confined meridional structure of  $\tau_x'$  in f3-L inevitably results in a stronger  
 358 response of  $D'$  to  $\tau_x'$  (i.e., a larger  $\beta_h$ ).

359 Furthermore, we compare the meridional structure of ENSO-related  $\tau_x'$  in the two model versions  
 360 with two reanalysis datasets. As shown in Fig. 7b, the meridional structures of  $\tau_x'$  in the two reanalysis  
 361 products (black lines) are very similar between 10°S–5°N, while some discrepancies exist within 5°N–  
 362 10°N region. For the models, the f3-H (25 km atmospheric resolution) shows a  $\tau_x'$  meridional structure  
 363 in the equatorial region (5°S–5°N) that substantially resembles both reanalysis datasets. In contrast, the  
 364 meridional structure of  $\tau_x'$  within 5°S–5°N in f3-L (100 km atmospheric resolution) shows obvious  
 365 discrepancies from the two reanalysis datasets. This suggests that model horizontal resolution can  
 366 influence ENSO simulation by affecting the meridional distribution of  $\tau_x'$ . The higher horizontal  
 367 resolution is conducive to the more realistic representation of equatorial  $\tau_x'$  meridional structure,  
 368 thereby yielding a more reasonable thermocline feedback and ENSO amplitude.



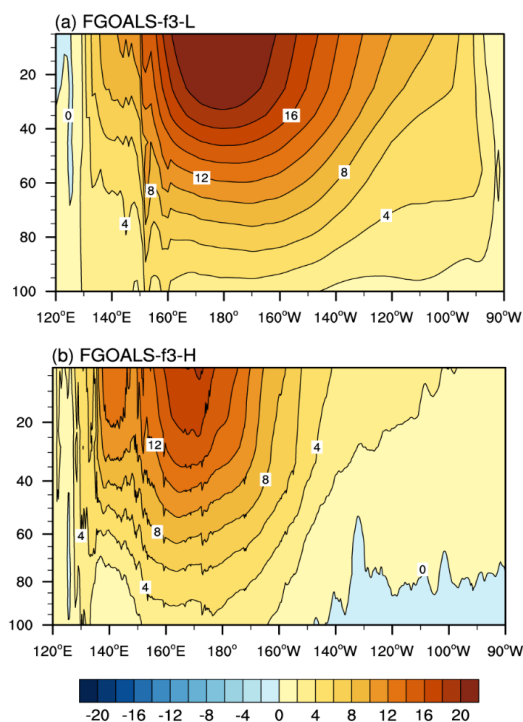
370 **Figure 7.** Meridional structure of normalized zonal wind stress anomalies [units:  $\text{N m}^{-2} (\text{N m}^{-2})^{-1}$ ] averaged over the  
 371 Niño4 longitude range (160°E–150°W) for (a) the result of f3-L (red solid line), f3-H (blue solid line) and their  
 372 difference (orange dash line, f3-L minus f3-H). In this plot, all the data are both interpolated onto a  $1^\circ \times 1^\circ$  grid to  
 373 facilitate the comparison; (b) the results of ORAS5 (black line marked by square, with a  $1^\circ \times 1^\circ$  grid), SODA2,2,4  
 374 (black line marked by asterisk, with a  $0.25^\circ \times 0.25^\circ$  grid), f3-L (red line marked by circle, with a  $1^\circ \times 1^\circ$  grid) and  
 375 f3-H (blue line marked by circle, with a  $0.25^\circ \times 0.25^\circ$  grid). The normalized zonal wind stress anomalies are obtained



376 by regressing the zonal wind stress anomalies field onto the Niño4 region (5°S–5°N, 160°E–150°W) averaged zonal  
 377 wind stress anomalies.

### 378 4.3 Resolution dependence of of the $\beta_u$ feedback process

379 The air-sea feedback process  $\beta_u$  represents the response of the equatorial Pacific upper ocean zonal  
 380 current anomaly ( $u'_o$ ) to  $\tau'_x$ . Figure 8 shows the equatorial profile for the response of  $u'_o$  to  $\tau'_x$  in both  
 381 models. Both models simulate eastward  $u'_o$  a unit eastward  $\tau'_x$  in equatorial Pacific, but the response  
 382 is significantly stronger in f3-L than in f3-H.

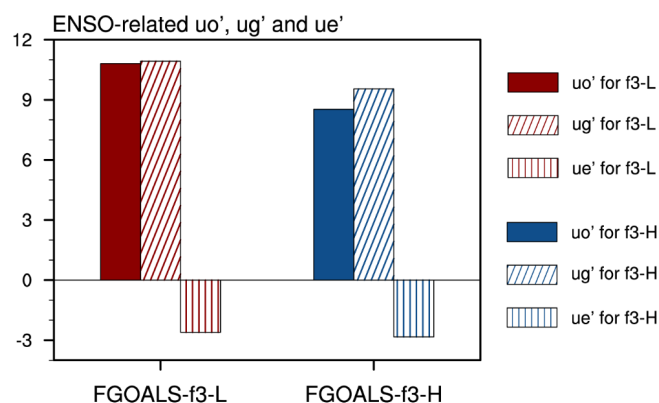


384 **Figure 8.** Vertical profile (averaged over 5°S–5°N) of the response of zonal ocean current anomaly ( $u'_o$ ) to the zonal  
 385 wind stress anomaly [ $\beta_u$ ; units:  $\text{m s}^{-1} (\text{N m}^{-2})^{-1}$ ] for (a) f3-L and (b) f3-H. The  $u'_o$  response to zonal wind stress  
 386 anomaly is derived through regressing the  $u'_o$  field onto the averaged zonal wind stress anomaly over equatorial  
 387 Pacific (5°S–5°N, 120°E–100°W).

388 Considering that ENSO-related  $u'_o$  in the equatorial region is composed of anomalous zonal  
 389 geostrophic currents ( $u'_g$ ) and anomalous Ekman currents ( $u'_e$ ) (see Section 2.3.4), we first diagnose the  
 390 response of  $u'_g$  and  $u'_e$  to wind stress anomalies (Fig. 9). In both versions, the response of  $u'_o$  to  $\tau'_x$   
 391 is primarily contributed by the response of  $u'_g$  to  $\tau'_x$ , while the contribution of  $u'_e$  is much smaller.  
 392 This indicates that the difference in  $u'_g$  response to  $\tau'_x$  is the main cause of the difference in  $\beta_u$



393 between the two model versions. Based on the geostrophic formula (Eq. (12)), the magnitude of  $u'_g$  is  
 394 related to the second-order meridional derivative of  $D'$ . Since the meridional structure of  $D'$  (Figure 6)  
 395 exhibits a parabolic shape (peaking at equator and decreasing poleward), a stronger  $D'$  in the equatorial  
 396 region corresponds to a larger value of second-order meridional derivative ( $-\frac{\partial^2 D'}{\partial y^2}$ ) and, consequently, a  
 397 stronger  $u'_g$ . As the  $D'$  response to  $\tau'_x$  is primarily modulated by the  $\beta_h$  process, the difference in the  $\beta_u$   
 398 between the f3-L and f3-H is also primarily attributed to difference in the  $\beta_h$  process. That is, the  
 399 stronger  $\beta_h$  in f3-L leads to a stronger  $D'$ , which in turn induces a stronger  $u'_g$  and  $u'_o$  and ultimately  
 400 resulting in a stronger  $\beta_u$ .



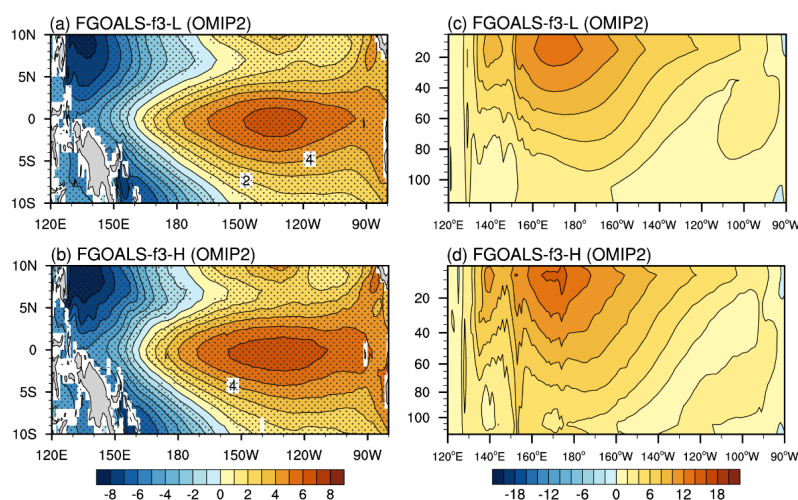
402 **Figure 9.** The ENSO-related zonal current anomalies ( $u'$ ) averaged over 0–65m, zonal geostrophic current anomalies  
 403 ( $u'_g$ ) and zonal Ekman current anomalies ( $u'_e$ ) along the equatorial eastern Pacific (2°S–2°N, 180°–80°W) for f3-L  
 404 (red bar) and f3-H (blue bar). The ENSO-related zonal current anomalies, zonal geostrophic current anomalies and  
 405 zonal Ekman current anomalies are obtained by regressing it anomalies field onto the zonal wind stress anomalies  
 406 averaged over Niño4 region (5°S–5°N, 160°E–150°W), respectively.

#### 407 4.4 Comparison with OMIP simulations

408 It should be noted that although our conclusions drawn from Coupled General Circulation Model  
 409 (CGCM) indicate that the improved ENSO simulation is closely linked to the increased atmospheric  
 410 resolution, the potential influence of oceanic horizontal resolution on ENSO simulation also merits brief  
 411 examination. For instance, from the perspective of the oceanic component of CGCM (OGCM), Li et al.  
 412 (2025) examined the oceanic zonal current and thermocline depth anomalies to zonal wind forcing by  
 413 analyzing the outputs from Ocean Model Intercomparison Project 2 (OMIP2, Griffies et al. 2016) and  
 414 found that the oceanic response to wind forcing is a key influencing ENSO simulation. According to the



415 CMIP6’s protocol, OMIP2 experiments are forced by the identical atmospheric reanalysis datasets, thus  
 416 providing an ideal framework to isolate biases originating from the oceanic component. Both f3-L and  
 417 f3-H have participated in OMIP2, with respective oceanic horizontal resolutions of approximately 1° and  
 418 0.1°. To further assess the potential impact of oceanic horizontal resolution on ENSO simulation, we  
 419 compare the key air-sea feedback terms ( $\beta_{\bar{h}}$  and  $\beta_u$ ) between f3-L and f3-H OMIP2 outputs. As shown  
 420 in figure 10, the response of zonal current and thermocline depth anomalies to the identical  $\tau'_x$  are  
 421 remarkably similar between the two ocean models. This indicates that the differences in OGCM  
 422 resolution itself may be not the primary driver of the differences in  $\beta_{\bar{h}}$  and  $\beta_u$ . Moreover, this finding  
 423 (Fig. 10) stands in stark contrast to the large differences in  $\beta_{\bar{h}}$  and  $\beta_u$  identified in the two coupled  
 424 simulations (Figures 6 and 8). Therefore, the evidence from the OMIP2 experiments corroborates our  
 425 main conclusion: the refined atmospheric horizontal resolution, which more realistically captures the  
 426 meridional structure of  $\tau'_x$ , plays a decisive role in improving the simulation of the key air-sea feedbacks  
 427 ( $\beta_{\bar{h}}$  and  $\beta_u$ ) and ENSO amplitude in FGOALS-f3 model. Nevertheless, a more comprehensive  
 428 investigation, potentially involving finer oceanic resolutions and their interactions with the atmosphere,  
 429 deserves further exploration in the future.



431 **Figure 10.** (a)–(b) and (c)–(d) are same as Figure 6 and 8, respectively, but for the OMIP2 simulations.

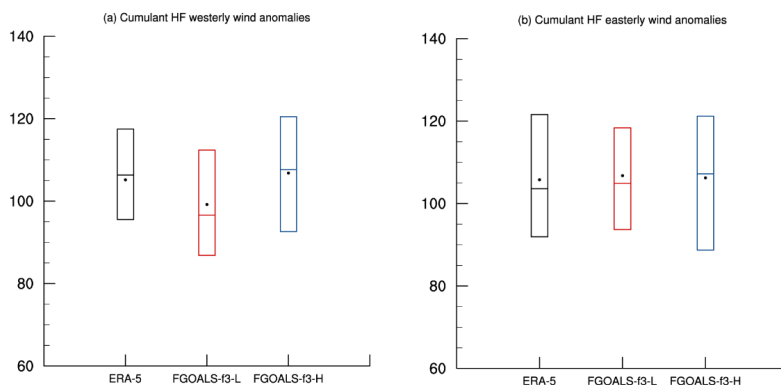


432 **5. Diagnostic of Oscillation Regularity**

433 **5.1 Resolution dependence of ENSO irregularity: Differences in HF wind activity**

434 A large body of studies have suggested that the HF zonal wind activity over the western and central  
435 equatorial Pacific (WCEP) plays a crucial role in the onset and development of ENSO events (Rong et  
436 al., 2011; Chen et al., 2017). In general, this HF zonal wind activity refers to westerly wind event (WWE)  
437 and easterly wind event (EWE) with time scales less than 90 days. Due to their transient, intense, and  
438 somewhat random nature, the HF zonal wind forcing is also considered as an important factor  
439 contributing to the diversity and irregularity of ENSO (Chen et al., 2015a; Fedorov et al., 2015).  
440 Motivated by this, we hypothesize that the difference in the ENSO regularity between f3-L and f3-H may  
441 arise from the differences in HF zonal wind activity.

442 To test this, we calculate the intensity of the HF zonal wind activity based on the Equations (10) and  
443 (11). Figure 11 shows the WWI and EWI index that measures the cumulative sum of HF westerly  
444 (easterly) wind anomalies over WCEP during the ENSO development period, for the observation and the  
445 two model versions. As shown in Figure 11a, the intensity of HF westerlies in f3-H is comparable to that  
446 in the observation, with similar mean and median values. However, f3-L simulates much weaker intensity  
447 of HF westerlies. For HF easterly activity (Fig. 11b), both models show comparable statistics to the  
448 observation, with no significant differences between f3-L and f3-H. Overall, these results indicate that  
449 f3-L strongly underestimates HF westerly wind activity relative to the f3-H and the observation.



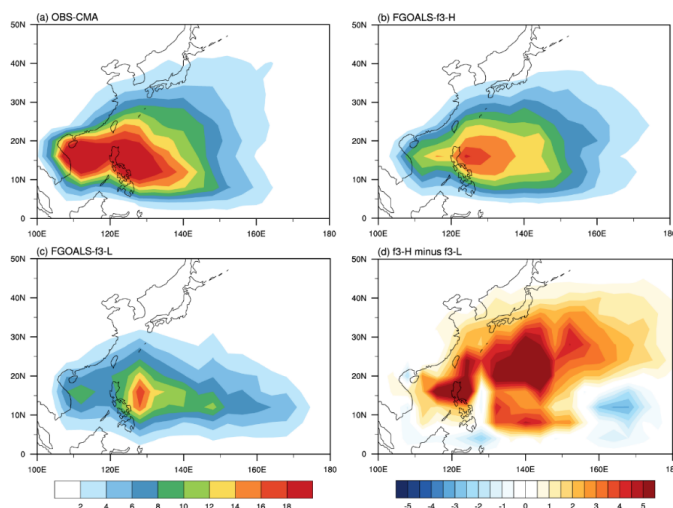
451 **Figure 11.** The boxplots of the (a) WWI index and (b) EWI index during the ENSO development phase (January to  
452 August). Red, blue and black boxes indicate the f3-L, f3-H and observations, respectively. The upper and lower  
453 boundaries of the box represent the 75th and 25th percentile values, respectively. The horizontal line in the box  
454 represents the median. The black dot represents the average value.



455 HF westerly wind activities are considered as semi-stochastic perturbations that modulate ENSO  
456 evolution (Gebbie et al., 2007; Gebbie and Tziperman, 2009). On the one hand, their occurrence is partly  
457 influenced by ENSO-related SSTA (Sun et al., 2020), but is also regarded as random noise independent  
458 of ENSO system (Moore and Kleeman, 1999). Most ENSO theoretical models treat HF westerly wind  
459 activity as external white noise forcing (Eisenman et al., 2005). This implies that the impact of HF activity  
460 on the coupled system depends, to some extent, on its relative magnitude relative to the ENSO amplitude.  
461 Notably, the aforementioned analysis based on the BJ index (a linear framework that excludes nonlinear  
462 processes like atmospheric "noise") has shown that the coupled system simulated by f3-L is more  
463 unstable than that of f3-H, hence the ENSO variability in f3-L is more prone to self-sustained oscillation.  
464 Considering that the HF wind activities in f3-L is also significantly weaker than that in f3-H, from the  
465 perspective of signal-to-noise ratio, this weaker "noise" is insufficient to "disrupt" the overly strong  
466 ENSO oscillation in f3-L, allowing its ENSO cycle to evolve in a regular and self-sustained manner.

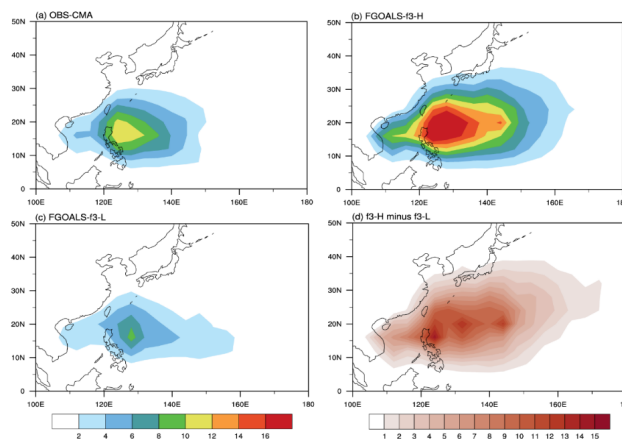
#### 467 **5.2 Sources of differences in HF westerly wind activity: evaluation of TC and MJO**

468 To further explore the origin of difference in HF westerly wind intensity between f3-L and f3-H, we  
469 compare the simulated performance of TC and MJO activities in these two models. The results show that  
470 the differences in the HF westerly wind intensity are primarily related to the models' ability to reproduce  
471 TC activity. Figure 12 shows the spatial distribution of TCTD over the western North Pacific (WNP).  
472 The spatial distributions of TCTD in both f3-H and f3-L are similar to the observation, with TC tracks  
473 primarily located in the southwestern quadrant of WNP. Although both versions can reasonably  
474 reproduce the spatial distribution characteristics of TC activities over the WNP, a significant difference  
475 exists in TC frequency: TC activity is much more frequent in f3-H than that in f3-L. The difference map  
476 (Fig. 12d, f3-H minus f3-L) shows positive values almost everywhere north of the equator. It is worth  
477 noting that TCTD in f3-H remains relatively weak compared to the observation, which is a common  
478 simulation bias in most current climate models (Nakamura et al., 2017; Tang et al., 2022). Although f3-  
479 H still underestimates TC activity compared to the observation, this improvement relative to f3-L is  
480 substantial.



482 **Figure 12.** The averaged TC track density (TCTD, units: counts year<sup>-1</sup>) over western North Pacific for (a)  
 483 Observation (CMA), (b) f3-H, (c) f3-L and (d) the difference between f3-H and f3-L (f3-H minus f3-L).

484 Furthermore, we compare the difference in the TC intensity between f3-H and f3-L. Figure 13 shows  
 485 the spatial distribution of ACE index over WNP in the observation and the models. Both models show  
 486 that strong TC activity is primarily concentrated east of the Philippines sea, consistent with the  
 487 observation (Ma et al., 2025). However, the ACE index in f3-H is significantly stronger than that in f3-  
 488 L, indicating stronger TC activity in f3-H. In summary, the TC activity over WNP is more frequent and  
 489 intense in f3-H than that in f3-L, which largely explains the stronger HF westerly wind anomalies in f3-  
 490 H.



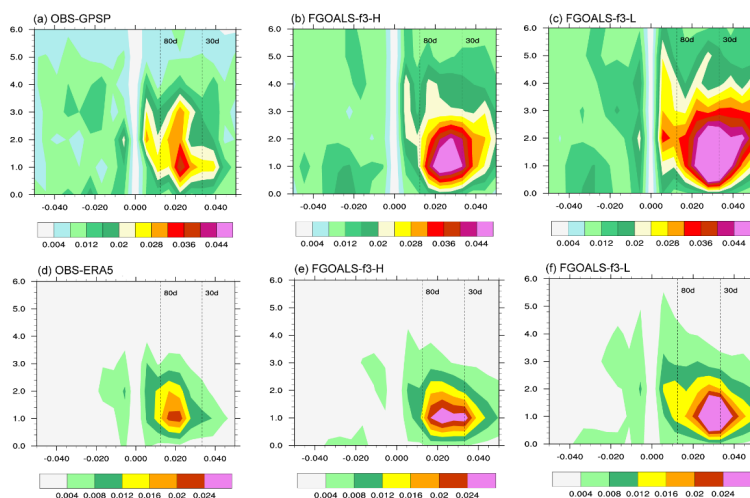
492 **Figure 13.** The averaged accumulated cyclone energy (ACE, units: 10<sup>4</sup> knots year<sup>-1</sup>) over western North Pacific for  
 493 (a) Observation (CMA), (b) f3-H, (c) f3-L and (d) the difference between f3-H and f3-L (f3-H minus f3-L).



494 Model horizontal resolution is a key factor in TC simulation (Tang et al., 2022). In general, climate  
495 models with coarse resolutions ( $\geq 100$  km) tend to only reproduce TC-like structures (Camargo et al.,  
496 2005; Camargo and Wing, 2016), with activity that is relatively weak and infrequent (Camargo, 2013;  
497 Nakamura et al., 2017). As resolution increases, models can simulate more frequent and more intense TC  
498 activity (Roberts et al., 2020a; Tang et al., 2022). In particular, when the horizontal resolution is increased  
499 to 25 km, the simulation of TC spatial structure and associated wind fields is significantly improved,  
500 yielding a more realistic characteristics of TC activity (Davis et al., 2018; Roberts et al., 2020b). The  
501 analysis results of the FGOALS-f3 models (Figs. 11 and 12) are consistent with these previous findings.  
502 Compared to the f3-L model (with a horizontal resolution of 100 km), the f3-H model (with a horizontal  
503 resolution of 25 km) provides a more realistic simulation of TC frequency and produces stronger TC  
504 intensities. This difference in TC simulation associated with model resolution modulates the intensity of  
505 HF westerly wind activity and hence influences the regularity of ENSO cycle in the two models.

506 On the other hand, we also examine the other primary source of HF westerly wind activity—the MJO.  
507 Figure 14 shows the wavenumber-frequency spectra of the space-time filtered precipitation and surface  
508 zonal wind. It can be seen that the spectral peak of precipitation and surface zonal wind in both models  
509 are concentrated at 30-80-day period, consistent with the observation. This indicates that both models  
510 reproduce the observed MJO timescales reasonably well. In terms of intensity, although the MJO-related  
511 precipitation and zonal wind fields are somewhat exaggerated in both models compared to the  
512 observation, the difference between f3-L and f3-H is small, especially for MJO-related zonal winds.  
513 Therefore, the difference in HF westerly wind activity is likely not directly linked to MJO activity.

514 Although finer resolution may improve model performance in certain aspects, these preliminary  
515 results show no significant improvement in MJO simulation from 100 km to 25 km in the case of these  
516 two FGOALS-f3 models. On the one hand, this may be because MJO simulation is heavily constrained  
517 by the accurate representation of physical processes like convection parameterization, boundary layer  
518 processes, and air-sea coupling. Thus, increased resolution must be combined with optimized physics  
519 schemes to effectively improve MJO simulation (Jiang et al., 2020b). On the other hand, recent studies  
520 suggest that significant MJO improvement can be seen when resolution increases to the kilometer-scale—  
521 the "convection-permitting resolution" (Savarin and Chen, 2022).



523 **Figure 14.** Wavenumber-frequency spectra of space-time filtered (a)–(c) precipitation (units: mm/day) and (d)–(f)  
 524 surface zonal wind fields (m/s) during the boreal winter (November to next August) for observation (left column),  
 525 f3-H (middle column) and f3-L (right column).

526 In summary, we find that the high-resolution version can better simulate TC activity, with more  
 527 frequent and stronger TCs over the WNP than the lower-resolution version. This difference results in  
 528 stronger HF westerly wind activity in f3-H than that in f3-L. Given that HF westerly wind activity acts  
 529 as a stochastic forcing on ENSO, the relatively weaker HF atmospheric noise in f3-L has a limited  
 530 randomizing effect on its stronger ENSO signal. In other words, the weaker HF zonal wind activity in  
 531 f3-L cannot overcome its inherently stronger ENSO signal, leading to overly regular oscillation in f3-L.  
 532 In contrast, f3-H has a weaker intrinsic ENSO signal but stronger "noise". Therefore, the ENSO cycle in  
 533 f3-L appears much more regular than that in f3-H.

## 534 6. Conclusions and implications

### 535 6.1 Conclusions

536 This study provides a process-based evaluation of how horizontal resolution influences ENSO  
 537 simulation in the CAS FGOALS-f3 climate system model. The comparison between its low-resolution  
 538 (f3-L) and high-resolution (f3-H) configurations reveals systematic and resolution-dependent differences  
 539 in ENSO amplitude, period, oscillation characteristics and underlying air-sea coupling processes.

540 A key structural source of bias in f3-L is the overly confined meridional structure of ENSO-related  
 541 zonal wind stress anomalies, which strengthens the thermocline and zonal-advection feedbacks and leads



542 to an exaggerated ENSO amplitude. The process-oriented diagnosis based on ENSO-related air-sea  
543 coupling processes demonstrates that these feedbacks can be directly attributed to resolution-sensitive  
544 wind stress structures, indicating that resolving the meridional structure of wind forcing is essential for  
545 realistic representation of ENSO amplitude.

546 The excessive regularity of ENSO in f3-L is another resolution-driven bias, arising from insufficient  
547 HF atmospheric variability. The high-resolution configuration produces more realistic TC activity and  
548 vigorous WWEs, which introduce stochastic forcing that disrupts the ENSO cycle and generates irregular  
549 variability. These results highlight that the representation of synoptic-scale atmospheric processes is  
550 integral to capturing realistic ENSO temporal evolution.

551 Overall, this study demonstrates that ENSO-related biases in FGOALS-f3 arise from identifiable,  
552 resolution-sensitive structural features in the coupled system. By quantifying feedback pathways and  
553 stochastic forcing, we provide a traceable explanation for how horizontal resolution modulates ENSO  
554 simulation. The evaluation framework developed here offers a practical tool for diagnosing ENSO  
555 performance in other climate models and can guide future development of the FGOALS-f3 model family.

## 556 **6.2 Implications**

557 The findings from this study yield several actionable insights for FGOALS-f3 development and for  
558 the broader CMIP-class modeling community. (1) Improving the atmospheric representation of  
559 equatorial wind stress gradients should be a development priority. (2) Resolution has complementary  
560 impacts on deterministic and stochastic ENSO dynamics: the deterministic aspect influencing ENSO  
561 simulation is via air-sea coupling processes, and the stochastic aspect influencing ENSO simulation is  
562 via high-frequency wind activity (e.g., WWE), whose dual role should be explicitly considered in future  
563 model evaluation framework. (3) The diagnostics applied here are model-agnostic and can serve as a  
564 reproducible framework for assessing resolution effects in other climate models participating in  
565 CMIP6/CMIP7. (4) The ~25 km atmosphere resolution of f3-H improves both the air-sea coupling  
566 processes and the stochastic forcing mechanisms. This supports the ongoing efforts toward next-  
567 generation high-resolution models.

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570 *Data and Code availability.* Source codes of the FGOALS-f3 model used in this study and the analysis  
571 scripts are archived at <https://doi.org/https://doi.org/10.5281/zenodo.17778266> (Song, 2025). The model  
572 output of FGOALS-f3 models described in this paper is distributed through the Earth System Grid  
573 Federation (ESGF) and is freely obtained via the ESGF data portals after registration  
574 (<https://aims2.llnl.gov/search>, last access: 27 November 2025). All the observation and reanalysis  
575 datasets are available online and publicly available as cited in the references. The ORAS5 and ERA-5  
576 datasets are obtained from <https://cds.climate.copernicus.eu/datasets>. The HadISST dataset is from  
577 <https://www.metoffice.gov.uk/hadobs/hadisst/data/download.html>. The GPCP dataset is available at  
578 <https://www.ncei.noaa.gov/data/global-precipitation-climatology-project-gpcp-daily/access>. The  
579 SODA2.2.4 dataset can be accessed at [http://apdrc.soest.hawaii.edu/datadoc/soda\\_2.2.4.php](http://apdrc.soest.hawaii.edu/datadoc/soda_2.2.4.php). The CMA  
580 TC best track dataset is from <https://tcddata.typhoon.org.cn/en/zjljsjj.html>. The description and data for  
581 HighResSST-present simulations of FGOALS-f3 can be found at  
582 <http://doi.org/10.22033/ESGF/CMIP6.3312> (Bao and He, 2019; Bao et al., 2020).

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586 the writing and revising of the manuscript. BA and HZ discussed the results and commented on the  
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602 **Reference**

603 Adler, R. F., Huffman, G. J., Chang, A., Ferraro, R., Xie, P.-P., Janowiak, J., Rudolf, B., Schneider, U.,  
604 Curtis, S., Bolvin, D., Gruber, A., Susskind, J., Arkin, P., and Nelkin, E.: The version-2 global  
605 precipitation climatology project (GPCP) monthly precipitation analysis (1979–present), *Journal of*  
606 *Hydrometeorology*, 4, 1147–1167, [https://doi.org/10.1175/1525-](https://doi.org/10.1175/1525-7541(2003)004%3C1147:TVGPCP%3E2.0.CO;2)  
607 [7541\(2003\)004%3C1147:TVGPCP%3E2.0.CO;2](https://doi.org/10.1175/1525-7541(2003)004%3C1147:TVGPCP%3E2.0.CO;2), 2003.

608 An, B., Yu, Y. Q., Bao, Q., He, B., Li, J. X., Luan, Y. H., Chen, K. J., and Zheng, W. P.: CAS FGOALS-  
609 f3-H dataset for the high-resolution model intercomparison project (HighResMIP) Tier 2, *Advances*  
610 *in Atmospheric Sciences*, 39, 1873–1884, <https://doi.org/10.1007/s00376-022-2030-5>, 2022.

611 Bacmeister, J. T., Wehner, M. F., Neale, R. B., Gettelman, A., Hannay, C., Lauritzen, P. H., Caron, J. M.,  
612 and Truesdale, J. E.: Exploratory high-resolution climate simulations using the community  
613 atmosphere model (CAM), *Journal of Climate*, 27, 3073–3099, [https://doi.org/10.1175/jcli-d-13-](https://doi.org/10.1175/jcli-d-13-00387.1)  
614 [00387.1](https://doi.org/10.1175/jcli-d-13-00387.1), 2014.

615 Bao, Q., Liu, Y. M., Wu, G. X., He, B., Li, J. X., Wang, L., Wu, X. F., Chen, K. J., Wang, X. C., Yang, J.,  
616 and Zhang, X. Q.: CAS FGOALS-f3-H and CAS FGOALS-f3-L outputs for the high-resolution  
617 model intercomparison project simulation of CMIP6, *Atmospheric and Oceanic Science Letters*, 13,  
618 576–581, <https://doi.org/10.1080/16742834.2020.1814675>, 2020.

619 Bao, Q. and He, B.: CAS FGOALS-f3-H model output prepared for CMIP6 HighResMIP highresSST-  
620 present, Earth System Grid Federation [data set], <https://doi.org/10.22033/ESGF/CMIP6.3312>,  
621 2019

622 Barnston, A. G., Tippett, M. K., L'Heureux, M. L., Li, S., and DeWitt, D. G.: Skill of real-time seasonal  
623 ENSO model predictions during 2002–11: Is our capability increasing?, *Bulletin of the American*  
624 *Meteorological Society*, 93, 631–651, <https://doi.org/10.1175/bams-d-11-00111.1>, 2012.

625 Bell, G. D., Halpert, M. S., Schnell, R. C., Higgins, R. W., Lawrimore, J., Kousky, V. E., Tinker, R.,



- 626 Thiaw, W., Chelliah, M., and Artusa, A.: Climate assessment for 1999, *Bulletin of the American*  
627 *Meteorological Society*, 81, s1–s50, [https://doi.org/10.1175/1520-0477\(2000\)81\[s1:CAF\]2.0.CO;2](https://doi.org/10.1175/1520-0477(2000)81[s1:CAF]2.0.CO;2),  
628 2000.
- 629 Cai, W. J., Santoso, A., Collins, M., Dewitte, B., Karamperidou, C., Kug, J.-S., Lengaigne, M., McPhaden,  
630 M. J., Stuecker, M. F., Taschetto, A. S., Timmermann, A., Wu, L. X., Yeh, S.-W., Wang, G., Ng, B.,  
631 Jia, F., Yang, Y., Ying, J., Zheng, X.-T., Bayr, T., Brown, J. R., Capotondi, A., Cobb, K. M., Gan, B.,  
632 Geng, T., Ham, Y.-G., Jin, F.-F., Jo, H.-S., Li, X., Lin, X., McGregor, S., Park, J.-H., Stein, K., Yang,  
633 K., Zhang, L., and Zhong, W. X.: Changing El Niño–Southern Oscillation in a warming climate,  
634 *Nature Reviews Earth and Environment*, 2, 628–644, <https://doi.org/10.1038/s43017-021-00199-z>,  
635 2021.
- 636 Camargo, S. J., Barnston, A. G., and Zebiak, S. E.: A statistical assessment of tropical cyclone activity in  
637 atmospheric general circulation models, *Tellus A: Dynamic Meteorology and Oceanography*, 57,  
638 589–604, <https://doi.org/10.3402/tellusa.v57i4.14705>, 2005.
- 639 Camargo, S. J.: Global and regional aspects of tropical cyclone activity in the CMIP5 models, *Journal of*  
640 *Climate*, 26, 9880–9902, <https://doi.org/10.1175/jcli-d-12-00549.1>, 2013.
- 641 Camargo, S. J. and Wing, A. A.: Tropical cyclones in climate models, *Wiley Interdisciplinary Reviews:*  
642 *Climate Change*, 7, 211–237, <https://doi.org/10.1002/wcc.373>, 2016.
- 643 Carton, J. A. and Giese, B. S.: A Reanalysis of Ocean Climate Using Simple Ocean Data Assimilation  
644 (SODA), *Monthly Weather Review*, 136, 2999–3017, <https://doi.org/10.1175/2007mwr1978.1>,  
645 2008.
- 646 Chang, P., Zhang, S. Q., Danabasoglu, G., Yeager, S. G., Fu, H. H., Wang, H., Castruccio, F. S., Chen,  
647 Y., Edwards, J., Fu, D., Jia, Y., Laurindo, L. C., Liu, X., Rosenbloom, N., Small, R. J., Xu, G. P.,  
648 Zeng, Y. H., Zhang, Q. Y., Bacmeister, J., Bailey, D. A., Duan, X., DuVivier, A. K., Li, D. P., Li, Y.  
649 X., Neale, R., Stössel, A., Wang, L., Zhuang, Y., Baker, A., Bates, S., Dennis, J., Diao, X., Gan, B.  
650 L., Gopal, A., Jia, D. N., Jing, Z., Ma, X., Saravanan, R., Strand, W. G., Tao, J., Yang, H. Y., Wang,  
651 X. Q., Wei, Z. Q., and Wu, L. X.: An unprecedented set of high-resolution earth system simulations  
652 for understanding multiscale interactions in climate variability and change, *Journal of Advances in*  
653 *Modeling Earth Systems*, 12, e2020MS002298, <https://doi.org/10.1029/2020ms002298>, 2020.
- 654 Chen, D. K., Lian, T., Fu, C., Cane, M. A., Tang, Y. M., Murtugudde, R., Song, X. S., Wu, Q. Y., and



- 655        Zhou, L.: Strong influence of westerly wind bursts on El Niño diversity, *Nature Geoscience*, 8, 339–  
656        345, <https://doi.org/10.1038/ngeo2399>, 2015a.
- 657        Chen, L., Li, T., and Yu, Y. Q.: Causes of strengthening and weakening of ENSO amplitude under global  
658        warming in four CMIP5 models, *Journal of Climate*, 28, 3250–3274, [https://doi.org/10.1175/jcli-d-](https://doi.org/10.1175/jcli-d-14-00439.1)  
659        14-00439.1, 2015b.
- 660        Chen, L., Yu, Y. Q., and Zheng, W. P.: Improved ENSO simulation from climate system model FGOALS-  
661        g1.0 to FGOALS-g2, *Climate Dynamics*, 47, 2617–2634, [https://doi.org/10.1007/s00382-016-](https://doi.org/10.1007/s00382-016-2988-8)  
662        2988-8, 2016.
- 663        Chen, L., Li, T., Wang, B., and Wang, L.: Formation mechanism for 2015/16 super El Niño, *Scientific*  
664        Reports, 7, 2975–2985, <https://doi.org/10.1038/s41598-017-02926-3>, 2017.
- 665        Chen, L., Wang, L., Li, T., and Liu, J.: Drivers of reduced ENSO variability in mid-Holocene in a coupled  
666        model, *Climate Dynamics*, 52, 5999–6014, <https://doi.org/10.1007/s00382-018-4496-5>, 2019.
- 667        Craig, A. P., Vertenstein, M., and Jacob, R.: A new flexible coupler for earth system modeling developed  
668        for CCSM4 and CESM1, *The International Journal of High Performance Computing Applications*,  
669        26, 31–42, <https://doi.org/10.1177/1094342011428141>, 2012.
- 670        Davis, C. A.: Resolving tropical cyclone intensity in models, *Geophysical Research Letters*, 45, 2082–  
671        2087, <https://doi.org/10.1002/2017GL076966>, 2018.
- 672        Dawson, A., Matthews, A. J., Stevens, D. P., Roberts, M. J., and Vidale, P. L.: Importance of oceanic  
673        resolution and mean state on the extra-tropical response to El Niño in a matrix of coupled models,  
674        *Climate Dynamics*, 41, 1439–1452, <https://doi.org/10.1007/s00382-012-1518-6>, 2013.
- 675        Docquier, D., Grist, J. P., Roberts, M. J., Roberts, C. D., Semmler, T., Ponsoni, L., Massonnet, F.,  
676        Sidorenko, D., Sein, D. V., Iovino, D., Bellucci, A., and Fichet, T.: Impact of model resolution on  
677        Arctic sea ice and North Atlantic Ocean heat transport, *Climate Dynamics*, 53, 4989–5017,  
678        <https://doi.org/10.1007/s00382-019-04840-y>, 2019.
- 679        Eisenman, I., Yu, L. S., and Tziperman, E.: Westerly wind bursts: ENSO’s tail rather than the dog?,  
680        *Journal of Climate*, 18, 5224–5237, <https://doi.org/10.1175/JCLI3588.1>, 2005.
- 681        Eyring, V., Bony, S., Meehl, G. A., Senior, C. A., Stevens, B., Stouffer, R. J., and Taylor, K. E.: Overview  
682        of the Coupled Model Intercomparison Project Phase 6 (CMIP6) experimental design and  
683        organization, *Geoscientific Model Development*, 9, 1937–1958, <https://doi.org/10.5194/gmd-9->



- 684 1937-2016, 2016.
- 685 Fedorov, A. V.: The response of the coupled tropical ocean–atmosphere to westerly wind bursts,  
686 Quarterly Journal of the Royal Meteorological Society, 128, 1–23,  
687 <https://doi.org/10.1002/qj.200212857901>, 2002.
- 688 Fedorov, A. V., Hu, S. N., Lengaigne, M., and Guilyardi, E.: The impact of westerly wind bursts and  
689 ocean initial state on the development, and diversity of El Niño events, *Climate Dynamics*, 44,  
690 1381–1401, <https://doi.org/10.1007/s00382-014-2126-4>, 2015.
- 691 Gebbie, G., Eisenman, I., Wittenberg, A., and Tziperman, E.: Modulation of westerly wind bursts by sea  
692 surface temperature: A semistochastic feedback for ENSO, *Journal of the Atmospheric Sciences*, 64,  
693 3281–3295, <https://doi.org/10.1175/jas4029.1>, 2007.
- 694 Gebbie, G. and Tziperman, E.: Incorporating a semi-stochastic model of ocean-modulated westerly wind  
695 bursts into an ENSO prediction model, *Theoretical and Applied Climatology*, 97, 65–73,  
696 <https://doi.org/10.1007/s00704-008-0069-6>, 2009.
- 697 Griffies, S. M., Danabasoglu, G., Durack, P. J., Adcroft, A. J., Balaji, V., Böning, C. W., Chassignet, E.  
698 P., Curchitser, E., Deshayes, J., Drange, H., Fox-Kemper, B., Gleckler, P. J., Gregory, J. M., Haak,  
699 H., Hallberg, R. W., Heimbach, P., Hewitt, H. T., Holland, D. M., Ilyina, T., Jungclaus, J. H., Komuro,  
700 Y., Krasting, J. P., Large, W. G., Marsland, S. J., Masina, S., McDougall, T. J., Nurser, A. J. G., Orr,  
701 J. C., Pirani, A., Qiao, F. I., Stouffer, R. J., Taylor, K. E., Treguier, A. M., Tsujino, H., Uotila, P.,  
702 Valdivieso, M., Wang, Q., Winton, M., and Yeager, S. G.: OMIP contribution to CMIP6:  
703 experimental and diagnostic protocol for the physical component of the Ocean Model  
704 Intercomparison Project, *Geoscientific Model Development*, 9, 3231–3296,  
705 <https://doi.org/10.5194/gmd-9-3231-2016>, 2016.
- 706 Guilyardi, E., Gualdi, S., Slingo, J., Navarra, A., Delecluse, P., Cole, J., Madec, G., Roberts, M., Latif,  
707 M., and Terray, L.: Representing El Niño in coupled ocean–atmosphere GCMs: The dominant role  
708 of the atmospheric component, *Journal of Climate*, 17, 4623–4629, [https://doi.org/10.1175/JCLI-](https://doi.org/10.1175/JCLI-3260.1)  
709 3260.1, 2004.
- 710 Guilyardi, E., Capotondi, A., Lengaigne, M., Thual, S., and Wittenberg, A. T.: ENSO modeling: History,  
711 progress, and challenges, in: *El Niño Southern Oscillation in a Changing Climate*, edited by:  
712 McPhaden, M. J., Santoso, A., and Cai, W., American Geophysical Union, American, 199-226,



- 713 <https://doi.org/10.1002/9781119548164.ch9>, 2020.
- 714 Hallberg, R.: Using a resolution function to regulate parameterizations of oceanic mesoscale eddy effects,  
715 *Ocean Modelling*, 72, 92–103, <https://doi.org/10.1016/j.ocemod.2013.08.007>, 2013.
- 716 Harrison, D. E. and Vecchi, G. A.: Westerly wind events in the tropical Pacific, 1986–95\*, *Journal of*  
717 *Climate*, 10, 3131–3156, [https://doi.org/10.1175/1520-0442\(1997\)010<3131:Wweitt>2.0.Co;2](https://doi.org/10.1175/1520-0442(1997)010<3131:Wweitt>2.0.Co;2),  
718 1997.
- 719 He, B., Bao, Q., Wang, X. C., Zhou, L. J., Wu, X. F., Liu, Y. M., Wu, G. X., Chen, K. J., He, S. C., Hu,  
720 W. T., Li, J. D., Li, J. X., Nian, G. K., Wang, L., Yang, J., Zhang, M. H., and Zhang, X. Q.: CAS  
721 FGOALS-f3-L model datasets for CMIP6 historical atmospheric model intercomparison project  
722 simulation, *Advances in Atmospheric Sciences*, 36, 771–778, [https://doi.org/10.1007/s00376-019-](https://doi.org/10.1007/s00376-019-9027-8)  
723 9027-8, 2019.
- 724 He, B., He, X. Y., Liu, Y. M., Wu, G. X., Bao, Q., Hu, W. T., Sheng, C., and Feng, S. J.: Role of thermal  
725 and dynamical subdaily perturbations over the Tibetan Plateau in 30-day extended-range forecast  
726 of East Asian precipitation in early summer, *npj Climate and Atmospheric Science*, 8,  
727 <https://doi.org/10.1038/s41612-025-00931-2>, 2025a.
- 728 He, X.-Y., He, B., Bao, Q., Liu, Y.-M., Li, J.-D., Wang, X.-C., Chen, X.-C., and Wu, G.-X.: The sensitivity  
729 of the Asian summer monsoon simulation to horizontal resolution and air–sea coupling in the  
730 FGOALS-f climate system model, *Advances in Climate Change Research*, 16, 44 – 57,  
731 <https://doi.org/10.1016/j.accre.2025.01.008>, 2025b.
- 732 Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas, J., Peubey,  
733 C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan, X., Balsamo, G., Bechtold,  
734 P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara, G., Dahlgren, P., Dee, D., Diamantakis, M.,  
735 Dragani, R., Flemming, J., Forbes, R., Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hogan, R.  
736 J., Hólm, E., Janisková, M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay,  
737 P., Rozum, I., Vamborg, F., Villaume, S., and Thépaut, J. N.: The ERA5 global reanalysis, *Quarterly*  
738 *Journal of the Royal Meteorological Society*, 146, 1999–2049, <https://doi.org/10.1002/qj.3803>,  
739 2020.
- 740 Hewitt, H. T., Roberts, M. J., Hyder, P., Graham, T., Rae, J., Belcher, S. E., Bourdallé-Badie, R., Copsey,  
741 D., Coward, A., Guiavarch, C., Harris, C., Hill, R., Hirschi, J. J. M., Madec, G., Mizielinski, M. S.,



- 742 Neininger, E., New, A. L., Rioual, J.-C., Sinha, B., Storkey, D., Shelly, A., Thorpe, L., and Wood,  
743 R. A.: The impact of resolving the Rossby radius at mid-latitudes in the ocean: results from a high-  
744 resolution version of the Met Office GC2 coupled model, *Geoscientific Model Development*, 9,  
745 3655–3670, <https://doi.org/10.5194/gmd-9-3655-2016>, 2016.
- 746 Hua, L. J., Chen, L., Rong, X. Y., Su, J. Z., Wang, L., Li, T., and Yu, Y. Q.: Impact of atmospheric model  
747 resolution on simulation of ENSO feedback processes: a coupled model study, *Climate Dynamics*,  
748 51, 3077–3092, <https://doi.org/10.1007/s00382-017-4066-2>, 2018.
- 749 Hunke, E. C. and Lipscomb, W. H.: CICE: The Los Alamos Sea Ice Model, Documentation and Software  
750 User's Manual, Version 4.1, Tech. Rep. LA-CC-06-012, Los Alamos National Laboratory, Los  
751 Alamos, New Mexico, available at: <http://oceans11.lanl.gov/trac/CICE>, 2010.
- 752 Jiang, W. P., Huang, P., Li, G., and Huang, G.: Emergent constraint on the frequency of central Pacific  
753 El Niño under global warming by the equatorial Pacific cold tongue bias in CMIP5/6 models,  
754 *Geophysical Research Letters*, 47, <https://doi.org/10.1029/2020gl089519>, 2020a.
- 755 Jiang, W. P., Huang, P., Huang, G., and Ying, J.: Origins of the excessive westward extension of ENSO  
756 SST simulated in CMIP5 and CMIP6 models, *Journal of Climate*, 34, 2839–2851,  
757 <https://doi.org/10.1175/jcli-d-20-0551.1>, 2021.
- 758 Jiang, X. N., Adames, Á. F., Kim, D., Maloney, E. D., Lin, H., Kim, H., Zhang, C. D., DeMott, C. A.,  
759 and Klingaman, N. P.: Fifty years of research on the Madden-Julian Oscillation: Recent progress,  
760 challenges, and perspectives, *Journal of Geophysical Research: Atmospheres*, 125,  
761 <https://doi.org/10.1029/2019jd030911>, 2020b.
- 762 Jin, F.-F.: An equatorial ocean recharge paradigm for ENSO. Part I: Conceptual model, *Journal of the*  
763 *Atmospheric Sciences*, 54, 811–829, [https://doi.org/10.1175/1520-0469\(1997\)054<0811:Aeorpf>2.0.Co;2](https://doi.org/10.1175/1520-0469(1997)054<0811:Aeorpf>2.0.Co;2), 1997.
- 765 Jin, F. F., Kim, S. T., and Bejarano, L.: A coupled-stability index for ENSO, *Geophysical Research Letters*,  
766 33, <https://doi.org/10.1029/2006gl027221>, 2006.
- 767 Kim, S. T. and Jin, F.-F.: An ENSO stability analysis. Part II: Results from the twentieth and twenty-first  
768 century simulations of the CMIP3 models, *Climate Dynamics*, 36, 1609–1627,  
769 <https://doi.org/10.1007/s00382-010-0872-5>, 2011a.
- 770 Kim, S. T. and Jin, F.-F.: An ENSO stability analysis. Part I: results from a hybrid coupled model, *Climate*



- 771 Dynamics, 36, 1593–1607, <https://doi.org/10.1007/s00382-010-0796-0>, 2011b.
- 772 Kreussler, P., Caron, L. P., Wild, S., Loosveldt Tomas, S., Chauvin, F., Moine, M. P., Roberts, M. J.,  
773 Ruprich-Robert, Y., Seddon, J., Valcke, S., Vannière, B., and Vidale, P. L.: Tropical cyclone  
774 integrated kinetic energy in an ensemble of HighResMIP simulations, *Geophysical Research Letters*,  
775 48, <https://doi.org/10.1029/2020gl090963>, 2021.
- 776 Lawrence, D. M., Oleson, K. W., Flanner, M. G., Thornton, P. E., Swenson, S. C., Lawrence, P. J., Zeng,  
777 X., Yang, Z. L., Levis, S., and Sakaguchi, K.: Parameterization improvements and functional and  
778 structural advances in version 4 of the Community Land Model, *Journal of Advances in Modeling  
779 Earth Systems*, 3, <https://doi.org/10.1029/2011MS00045>, 2011.
- 780 Li, J., Yu, Y. Q., and Sun, D. Z.: Asymmetric Responses in the Equatorial Pacific to Wind Forcing in  
781 OMIP2 Experiments: Role of Zonal Currents, *Geophysical Research Letters*, 52, e2025GL114661,  
782 <https://doi.org/10.1029/2025gl114661>, 2025.
- 783 Li, J. X., Bao, Q., Liu, Y. M., Wang, L., Yang, J., Wu, G. X., Wu, X. F., He, B., Wang, X. C., Zhang, X.  
784 Q., Yang, Y. X., and Shen, Z. L.: Effect of horizontal resolution on the simulation of tropical  
785 cyclones in the Chinese Academy of Sciences FGOALS-f3 climate system model, *Geoscientific  
786 Model Development*, 14, 6113–6133, <https://doi.org/10.5194/gmd-14-6113-2021>, 2021.
- 787 Li, T.: Phase transition of the El Niño–Southern Oscillation: A stationary SST mode, *Journal of the  
788 atmospheric sciences*, 54, 2872–2887, [https://doi.org/10.1175/1520-  
0469\(1997\)054<2872:PTOTEN>2.0.CO;2](https://doi.org/10.1175/1520-<br/>789 0469(1997)054<2872:PTOTEN>2.0.CO;2), 1997.
- 790 Li, T. Y., Yu, Y. Q., An, B., Luan, Y. H., and Chen, K. J.: Tropical instability waves in a high-resolution  
791 oceanic and coupled GCM, *Ocean Modelling*, 182, <https://doi.org/10.1016/j.ocemod.2023.102169>,  
792 2023.
- 793 Li, Y., Liu, H., Ding, M., Lin, P., Yu, Z., Yu, Y., Meng, Y., Li, Y., Jian, X., and Jiang, J.: Eddy-resolving  
794 simulation of CAS-LICOM3 for phase 2 of the ocean model intercomparison project, *Advances in  
795 Atmospheric Sciences*, 37, 1067–1080, <https://doi.org/10.1007/s00376-020-0057-z>, 2020.
- 796 Liang, Y. and Fedorov, A. V.: Linking the Madden–Julian Oscillation, tropical cyclones and westerly  
797 wind bursts as part of El Niño development, *Climate Dynamics*, 57, 1039–1060,  
798 <https://doi.org/10.1007/s00382-021-05757-1>, 2021.
- 799 Liao, H. X., Cai, Z. C., Guo, J. S., and Song, Z. Y.: Effects of ITCZ poleward location bias on ENSO



- 800 seasonal phase-locking simulation in climate models, *Journal of Climate*, 36, 5233–5249,  
801 <https://doi.org/10.1175/jcli-d-22-0891.1>, 2023.
- 802 Lu, B., Jin, F.-F., and Ren, H.-L.: A coupled dynamic index for ENSO periodicity, *Journal of Climate*,  
803 31, 2361–2376, <https://doi.org/10.1175/jcli-d-17-0466.1>, 2018.
- 804 Lu, X. Q., Yu, H., Ying, M., Zhao, B. K., Zhang, S., Lin, L. M., Bai, L. N., and Wan, R. J.: Western north  
805 Pacific tropical cyclone database created by the China Meteorological Administration, *Advances in*  
806 *Atmospheric Sciences*, 38, 690–699, <https://doi.org/https://10.1007/s00376-020-0211-7>, 2021.
- 807 Ma, T., Yu, W.-D., Speich, S., Zhao, H.-K., Xin, R., Luo, H., and Wu, L.-G.: Philippine archipelago and  
808 South China Sea monsoon plus ocean cooling buffer Northwestern Pacific super typhoons, *Nature*  
809 *Communications*, 16, <https://doi.org/10.1038/s41467-025-62334-4>, 2025.
- 810 Madden, R. A. and Julian, P. R.: Detection of a 40–50 day oscillation in the zonal wind in the tropical  
811 Pacific, *Journal of Atmospheric Sciences*, 28, 702–708, [https://doi.org/10.1175/1520-0469\(1971\)028%3C0702:DOADOI%3E2.0.CO;2](https://doi.org/10.1175/1520-0469(1971)028%3C0702:DOADOI%3E2.0.CO;2), 1971.
- 813 Madden, R. A. and Julian, P. R.: Description of global-scale circulation cells in the tropics with a 40–50  
814 day period, *Journal of Atmospheric Sciences*, 29, 1109–1123, [https://doi.org/10.1175/1520-0469\(1972\)029%3C1109:DOGSCC%3E2.0.CO;2](https://doi.org/10.1175/1520-0469(1972)029%3C1109:DOGSCC%3E2.0.CO;2), 1972.
- 816 McPhaden, M. J., Zebiak, S. E., and Glantz, M. H.: ENSO as an integrating concept in earth science,  
817 *Science*, 314, 1740–1745, <https://doi.org/10.1126/science.1132588>, 2006.
- 818 Moore, A. M. and Kleeman, R.: Stochastic forcing of ENSO by the intraseasonal oscillation, *Journal of*  
819 *Climate*, 12, 1199–1220, [https://doi.org/10.1175/1520-0442\(1999\)012<1199:Sfoebt>2.0.Co;2](https://doi.org/10.1175/1520-0442(1999)012<1199:Sfoebt>2.0.Co;2),  
820 1999.
- 821 Nakamura, J., Camargo, S. J., Sobel, A. H., Henderson, N., Emanuel, K. A., Kumar, A., LaRow, T. E.,  
822 Murakami, H., Roberts, M. J., Scoccimarro, E., Vidale, P. L., Wang, H., Wehner, M. F., and Zhao,  
823 M.: Western North Pacific Tropical Cyclone Model Tracks in Present and Future Climates, *Journal*  
824 *of Geophysical Research: Atmospheres*, 122, 9721–9744, <https://doi.org/10.1002/2017jd027007>,  
825 2017.
- 826 Navarra, A., Gualdi, S., Masina, S., Behera, S., Luo, J. J., Masson, S., Guilyardi, E., Delecluse, P., and  
827 Yamagata, T.: Atmospheric horizontal resolution affects tropical climate variability in coupled  
828 models, *Journal of Climate*, 21, 730–750, <https://doi.org/10.1175/2007jcli1406.1>, 2008.



- 829 Planton, Y. Y., Guilyardi, E., Wittenberg, A. T., Lee, J., Gleckler, P. J., Bayr, T., McGregor, S., McPhaden,  
830 M. J., Power, S., Roehrig, R., Vialard, J., and Voltaire, A.: Evaluating climate models with the  
831 CLIVAR 2020 ENSO metrics package, *Bulletin of the American Meteorological Society*, 102,  
832 E193–E217, <https://doi.org/10.1175/bams-d-19-0337.1>, 2021.
- 833 Rayner, N. A., Parker, D. E., Horton, E. B., Folland, C. K., Alexander, L. V., Rowell, D. P., Kent, E. C.,  
834 and Kaplan, A.: Global analyses of sea surface temperature, sea ice, and night marine air  
835 temperature since the late nineteenth century, *Journal of Geophysical Research: Atmospheres*, 108,  
836 <https://doi.org/10.1029/2002jd002670>, 2003.
- 837 Roberts, M. J., Camp, J., Seddon, J., Vidale, P. L., Hodges, K., Vanniere, B., Mecking, J., Haarsma, R.,  
838 Bellucci, A., Scoccimarro, E., Caron, L.-P., Chauvin, F., Terray, L., Valcke, S., Moine, M.-P.,  
839 Putrasahan, D., Roberts, C., Senan, R., Zarzycki, C., and Ullrich, P.: Impact of model resolution on  
840 tropical cyclone simulation using the HighResMIP–PRIMAVERA multimodel ensemble, *Journal*  
841 *of Climate*, 33, 2557–2583, <https://doi.org/10.1175/jcli-d-19-0639.1>, 2020a.
- 842 Roberts, M. J., Camp, J., Seddon, J., Vidale, P. L., Hodges, K., Vannière, B., Mecking, J., Haarsma, R.,  
843 Bellucci, A., Scoccimarro, E., Caron, L. P., Chauvin, F., Terray, L., Valcke, S., Moine, M. P.,  
844 Putrasahan, D., Roberts, C. D., Senan, R., Zarzycki, C., Ullrich, P., Yamada, Y., Mizuta, R., Kodama,  
845 C., Fu, D., Zhang, Q., Danabasoglu, G., Rosenbloom, N., Wang, H., and Wu, L.: Projected future  
846 changes in tropical cyclones using the CMIP6 HighResMIP multimodel ensemble, *Geophysical*  
847 *Research Letters*, 47, e2020GL088662, <https://doi.org/10.1029/2020gl088662>, 2020b.
- 848 Roberts, M. J., Reed, K. A., Bao, Q., Barsugli, J. J., Camargo, S. J., Caron, L.-P., Chang, P., Chen, C.-T.,  
849 Christensen, H. M., Danabasoglu, G., Frenger, I., Fučkar, N. S., ul Hasson, S., Hewitt, H. T., Huang,  
850 H., Kim, D., Kodama, C., Lai, M., Leung, L.-Y. R., Mizuta, R., Nobre, P., Ortega, P., Paquin, D.,  
851 Roberts, C. D., Scoccimarro, E., Seddon, J., Treguier, A. M., Tu, C.-Y., Ullrich, P. A., Vidale, P. L.,  
852 Wehner, M. F., Zarzycki, C. M., Zhang, B., Zhang, W., and Zhao, M.: High-Resolution model  
853 intercomparison project phase 2 (HighResMIP2) towards CMIP7, *Geoscientific Model*  
854 *Development*, 18, 1307–1332, <https://doi.org/10.5194/gmd-18-1307-2025>, 2025.
- 855 Rong, X. Y., Zhang, R. H., Li, T., and Su, J. Z.: Upscale feedback of high-frequency winds to ENSO,  
856 *Quarterly Journal of the Royal Meteorological Society*, 137, 894–907,  
857 <https://doi.org/10.1002/qj.804>, 2011.



- 858 Savarin, A. and Chen, S. S.: Pathways to better prediction of the MJO:1. Effects of model resolution and  
859 moist physics on Atmospheric boundary layer and precipitation, *Journal of Advances in Modeling*  
860 *Earth Systems*, 14, e2021MS002928, <https://doi.org/10.1029/2021ms002929>, 2022.
- 861 Small, R. J., Bryan, F. O., Bishop, S. P., and Tomas, R. A.: Air–Sea turbulent heat fluxes in climate models  
862 and observational analyses: What drives their variability?, *Journal of Climate*, 32, 2397–2421,  
863 <https://doi.org/10.1175/jcli-d-18-0576.1>, 2019.
- 864 Song, M.: The code of FGOALS-f3 model study, Zenodo [code],  
865 <https://doi.org/10.5281/zenodo.17778266>, 2025.
- 866 Su, J. Z., Zhang, R. H., Li, T., Rong, X. Y., Kug, J. S., and Hong, C.-C.: Causes of the El Niño and La  
867 Niña amplitude asymmetry in the equatorial eastern Pacific, *Journal of Climate*, 23, 605–617,  
868 <https://doi.org/10.1175/2009jcli2894.1>, 2010.
- 869 Su, J. Z., Li, T., and Zhang, R. H.: The initiation and developing mechanisms of central Pacific El Niños,  
870 *Journal of Climate*, 27, 4473–4485, <https://doi.org/10.1175/jcli-d-13-00640.1>, 2014.
- 871 Sun, M., Li, T., and Chen, L.: El Niño phase-dependent high-frequency variability in western equatorial  
872 Pacific, *Climate Dynamics*, 55, 2165–2184, <https://doi.org/10.1007/s00382-020-05376-2>, 2020.
- 873 Tang, Y. L., HuangFu, J. I., Huang, R. H., and Chen, W.: Simulation and projection of tropical cyclone  
874 activities over the western North Pacific by CMIP6 HighResMIP, *Climate Dynamics*, 35, 7771–  
875 7794, <https://doi.org/10.1175/JCLI-D-21-0760.1>, 2022.
- 876 Thomson, D. J.: Spectrum estimation and harmonic analysis, *Proceedings of the IEEE*, 79, 1055–1096,  
877 <https://doi.org/10.1109/PROC.1982.12433>, 1982.
- 878 Timmermann, A., An, S. I., Kug, J. S., Jin, F. F., Cai, W., Capotondi, A., Cobb, K. M., Lengaigne, M.,  
879 McPhaden, M. J., Stuecker, M. F., Stein, K., Wittenberg, A. T., Yun, K. S., Bayr, T., Chen, H. C.,  
880 Chikamoto, Y., Dewitte, B., Dommenget, D., Grothe, P., Guilyardi, E., Ham, Y. G., Hayashi, M.,  
881 Ineson, S., Kang, D., Kim, S., Kim, W., Lee, J. Y., Li, T., Luo, J. J., McGregor, S., Planton, Y., Power,  
882 S., Rashid, H., Ren, H. L., Santoso, A., Takahashi, K., Todd, A., Wang, G., Wang, G., Xie, R., Yang,  
883 W. H., Yeh, S. W., Yoon, J., Zeller, E., and Zhang, X.: El Niño–Southern Oscillation complexity,  
884 *Nature*, 559, 535–545, <https://doi.org/10.1038/s41586-018-0252-6>, 2018.
- 885 Yan, Y. and Sun, D.-Z.: Phase-Locking of El Niño and La Niña events in CMIP6 models, *Atmosphere*,  
886 15, 882, <https://doi.org/10.3390/atmos15080882>, 2024.



- 887 Ying, J., Lian, T., Ren, H.-L., Zhang, C., Liu, T., and Tan, X.-X.: Effects of tropical cyclones on ENSO,  
888 *Journal of Climate*, 32, 6423–6443, <https://doi.org/10.1175/jcli-d-18-0821.1>, 2019.
- 889 Ying, M., Zhang, W., Yu, H., Lu, X. Q., Feng, J. X., Fan, Y. X., Zhu, Y. t., and Chen, D. Q.: An overview  
890 of the China Meteorological Administration tropical cyclone database, *Journal of Atmospheric and*  
891 *Oceanic Technology*, 31, 287–301, <https://doi.org/https://10.1175/JTECH-D-12-00119.1>, 2014.
- 892 Yu, Y. Q., An, B., Liu, H. L., Bao, Q., Lin, P. F., He, B., Zheng, W. P., Luan, Y. H., Bai, R. and Li, T. Y.:  
893 Review of the development and application of the high-resolution climate system model FGOALS,  
894 *Chinese Journal of Atmospheric Sciences (in Chinese)*, 48, 200-217,  
895 <https://doi.org/10.1007/BF02915571>, 2024.
- 896 Yu, Y. Q., Tang, S. L., Liu, H. L., Lin, P. F., and Li, X. L.: Development and evaluation of the dynamic  
897 framework of an ocean general circulation model with arbitrary orthogonal curvilinear coordinate,  
898 *Chinese Journal of Atmospheric Sciences (in Chinese)*, 42, 877–889,  
899 <https://doi.org/10.3878/j.issn.1006-9895.1805.17284>, 2018.
- 900 Zhang, R.-H., Yu, Y. Q., Song, Z. Y., Ren, H.-L., Tang, Y. M., Qiao, F. I., Wu, T. W., Gao, C., Hu, J. Y.,  
901 Tian, F., Zhu, Y. C., Chen, L., Liu, H. L., Lin, P. F., Wu, F. H., and Wang, L.: A review of progress  
902 in coupled ocean-atmosphere model developments for ENSO studies in China, *Journal of*  
903 *Oceanology and Limnology*, 38, 930–961, <https://doi.org/10.1007/s00343-020-0157-8>, 2020.
- 904 Zhao, J. W., Wang, F., Zhan, R. F., Guo, Y. P., Huang, X., and Liu, C.: How Does Tropical Cyclone  
905 Genesis Frequency Respond to a Changing Climate?, *Geophysical Research Letters*, 50,  
906 <https://doi.org/10.1029/2023gl102879>, 2023.
- 907 Zhao, Y. D. and Sun, D.-Z.: ENSO asymmetry in CMIP6 models, *Journal of Climate*, 35, 5555–5572,  
908 <https://doi.org/10.1175/jcli-d-21-0835.1>, 2022.
- 909 Zhou, L., Bao, Q., Liu, Y. M., Wu, G. X., Wang, W. C., Wang, X. C., He, B., Yu, H. Y., and Li, J. D.:  
910 Global energy and water balance: Characteristics from Finite-volume Atmospheric Model of the  
911 IAP/LASG (FAMIL 1), *Journal of Advances in Modeling Earth Systems*, 7, 1–20,  
912 <https://doi.org/10.1002/2014MS000349>, 2015.
- 913 Zi, P., Liu, Y. M., Li, J. D., Yang, R. W., He, B., and Bao, Q.: Reduced Spring Precipitation Bias and  
914 Associated Physical Causes over South China in FGOALS-f3 Climate Models: Experiments with  
915 the Horizontal Resolutions, *Journal of Meteorological Research*, 38, 784–804,



916            <https://doi.org/10.1007/s13351-024-3200-4>, 2024.

917    Zuo, H., Balmaseda, M. A., Tietsche, S., Mogensen, K., and Mayer, M.: The ECMWF operational

918            ensemble reanalysis–analysis system for ocean and sea ice: A description of the system and

919            assessment, *Ocean Science*, 15, 779–808, <https://doi.org/10.5194/os-15-779-2019>, 2019.