



Saltwater exposure accelerates ice grain growth and may increase fracture vulnerability

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Abstract:

Natural ices often fail at stresses much lower than those measured in laboratory settings, complicating our understanding of glacial failure and icy moon crustal fracture. This may be because strength models, which depend on the size of individual ice grains, do not account for the saltwater commonly found in terrestrial and planetary ices. We conducted grain growth experiments, finding that saltwater always modifies grain growth compared to pure ice, and that increasing volume of saltwater introduces a pinning effect limiting this growth. Ice grain size therefore depends directly on liquid fraction, controlled by salinity and temperature. As such, we introduce an equation for grain growth parameters that is dependent on brine fraction and temperature relative to the liquidus. We then use our grain growth models to predict how saltwater infiltration affects ice strength. Modeled effects of grain growth on tensile strength following saltwater infiltration find that low-salinity water can reduce ice strength by up to 46% within 24 hours, narrowing gaps between observations and experiments.



1. Introduction

The strength evolution of icy materials is controlled by microstructural processes. Ice grain size, for example, is a crucial parameter for predicting future glacial mass balance (Alley, 1992; Alley et al., 1986; Behn et al., 2021; Faria et al., 2014a; Kingslake et al., 2022; Ranganathan & Minchew, 2024; Rignot et al., 2011) and the extent of heat released to maintain subsurface oceans in icy moons (Kihoulou et al., 2025; Nimmo, 2025; Nimmo & Manga, 2017; Poirier et al., 1983). As cryospheric environments like glaciers and icy moon interiors experience stress, they can deform both viscously and brittly. It has been observed in laboratory settings that coarse-grained ices are more resistant to flow (Goldsby & Kohlstedt, 1997, 2001), yet also more susceptible to fracture (Currier & Schulson, 1982; Lee & Schulson, 1988; Schulson et al., 1984). Fracture yields macroscopic damage, such as crevassing and calving, that relieves buttressing stresses which otherwise resist fast glacial flow and mass loss at Earth's poles (Chudley et al., 2025), and allows faults or vents to propagate through the crusts of icy moons. Grain size evolution is therefore key to understanding the future of cryospheric environments. However, we do not yet have a good understanding of how grain growth proceeds if liquid water is present.

Grain size evolution in a single-phase material is commonly treated as a dynamic competition between processes that facilitate grain size reduction, such as dynamic recrystallization, and processes that drive normal or “static” grain growth driven by reduction of grain-boundary energy (Austin & Evans, 2007; Cross et al., 2015; De Bresser et al., 1998; Roessiger et al., 2011). Static grain growth is often modeled using an Arrhenius (or Hillert) relationship to predict grain size, d , as a function of time, t , such that

$$d(t)^p = d(t_0)^p + kt$$

(Eq. 1)

$$k = k_0 \exp(-Q/RT)$$

(Eq. 2)

where p is the grain-growth exponent, $d(t_0)$ is the initial grain size, k_0 is a constant, Q is activation enthalpy, R is the ideal gas constant, and T is temperature in Kelvin. The theoretical value of p (also called m or n in some literature) is 2 for pure, single-phase materials. For ice, p has been found to vary between 2 – 20 depending on the presence of bubbles, local conditions, and impurity content (Alley et al., 1986; Azuma et al., 2012; Faria et al., 2014a, 2014b; Jellinek & Gouda, 1969).



Laboratory studies show that below the eutectic temperature of a water-salt mixture, 251.95 K (−21.2°C), where only solid phases are present, grain growth is inhibited relative to pure ice (De Achaval et al., 1987; Jellinek & Gouda, 1969; Wang et al., 2024). This is because impurities preferentially segregate to grain boundaries, where they exert drag and reduce boundary mobility by orders of magnitude (Gleiter, 1969). Conversely, at temperatures above the minimum eutectic temperature of a water-salt mixture where saline water is present between grains (the concentration of which is determined by relative distance from a liquidus curve; Figure S1), grain growth is *enhanced* by the presence of salts (De Achaval et al., 1987; Jellinek & Gouda, 1969). This observation is often attributed to the fact that diffusion is energetically easier through an interstitial liquid phase than along a solid-solid interface (e.g., Cooper & Kohlstedt, 1984). However, the influence of salinity and liquid fraction have not been determined in isolation; it is not yet clear whether grain growth enhancement in the presence of a brine is due to the presence of the liquid alone, or if the salinity and volume of liquid further affect grain growth.

To address uncertainty relating temperature, salinity, and grain size, we conducted static grain-growth experiments on ice of varying salinities at temperatures above the H₂O-NaCl eutectic minimum. We characterized grain sizes and dihedral angles through microstructural imaging and hand tracing. In doing so, we found that grain growth was always higher in salty ice than in pure ice. However, we also found that increasing salinity, and therefore an increasing amount of saltwater, inhibited grain growth. We provide a physical explanation for this observation, establish a constitutive relationship for grain growth parameters as a function of both salinity and temperature, and discuss the implications for ice dynamics in cases of tensile failure.

2. Methods

We surveyed a range of salinities and temperatures in order to explore different melt fractions in the NaCl-H₂O eutectic system (Figure 1A). By holding salinity or temperature constant, we were able to explore variations in melt fraction and the “vertical lever” (Figure 1B), a quantity which has been shown to simplify variations between systems in properties like acoustic velocities (McCarthy et al., 2019).

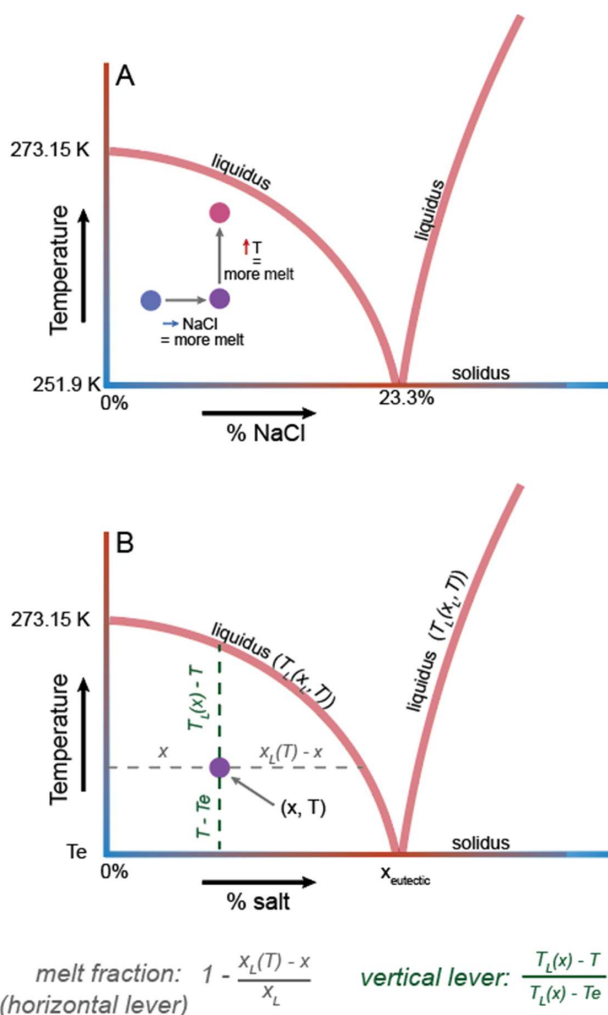


Fig. 1. Eutectic diagrams in a saline-H₂O system. A) Diagram specific to the NaCl-H₂O system, showing how increasing temperature and increasing salinity both add more melt by drawing a salty ice aggregate closer to the liquidus curve. B) Generic eutectic diagram indicating the calculation of melt fraction and the “vertical lever.”

2.1 Sample Fabrication

Individual ice samples with “high” salinity (HS), “low” salinity (LS), and no salinity (NS) were prepared using the freeze-flood “standard ice” method of (Cole, 1979), adapted for eutectic mixtures as described by Ravikumar et al. (2025). First, bubble-free ice made from pure, deionized water was ground and sieved to obtain ice particles of 106–250 μm diameter. This ice powder was packed into a die with internal dimensions of 50.8 \times 50.8 \times 19.8 mm, which was sealed with vacuum grease, and placed in a saltwater



bath at 268 K (HS/LS) or 272 K (NS) or below. This temperature exposure does not cause ice grains to grow, as the die preserves the temperature of 259 K or below used in the freezer during preparation. The die was then placed under a vacuum relative pressure of -0.98 bar for at least 30 minutes to remove any pore space between ice particles. Next, the die was flooded with either saline water of 5 wt.% (HS) or 2.9 wt.% (LS) NaCl, or pure deionized water (NS). Vacuum was maintained until this “flood” water percolated through the entire sample, at which point the sample was promptly disconnected from vacuum, and placed on a copper plate in a chest freezer at 248 K for at least 12 hours to initiate unidirectional, bottom-up freezing expelling any trapped air.

To characterize the initial microstructures, a flat surface of each sample was prepared via cryomicrotome, then imaged using a reflected-light microscope in a walk-in freezer at temperatures below the minimum H_2O -NaCl eutectic temperature (251.8 K). Grains within the NS samples grew tightly packed during freezing, and so had slightly larger initial grain sizes than the HS/LS samples. Pieces of each sample were also collected for analysis of their salinity using a handheld refractive-index meter and the calibration curves (Ravikumar et al., 2025). The effective bulk salinities of the HS samples fell between 1.31–2.17 wt% NaCl, the salinities of LS samples fell between 0.34–0.74 wt% NaCl, and the NS samples had 0 wt% NaCl (or other impurities). For context, natural sea ices cover a similar range of salinities, up to 3.5% (Weeks, 2010).

2.2 Grain growth procedure

Each sample was cut into rectangular pieces of $16.9 \times 50.8 \times 19.8$ mm, and vacuum sealed in plastic bags. Experiments were named by the convention xSyT, where x was H/L/N to indicate salinity, and y was H/L to indicate high/low temperature (268 K and 260 K, respectively), followed by a number (1/2) to indicate the sample prep number (from the first or second batch of samples) and a letter (A/B/C/D/E) to indicate the duration of annealing at the given temp (3.16hr/7hr/31hr/100hr/316hr). For example, sample LSHT2.E was a low-salinity sample annealed at relatively high temperature, came from the second batch of LSHT samples, and was annealed for 316 hours. Samples were placed in an insulated box inside each freezer that stabilized temperatures by reducing the duty cycle effect (Figure S1). After annealing, samples were quenched using liquid nitrogen to capture the microstructure at the exact time of removal, and then stored in a freezer at 218 K for at least two hours.

2.3 Imaging and analysis

Each sample was microtomed in a freezer below 251.8K, then imaged in reflected light microscopy. Refractive indices were measured at this point as well to compare with the starting material so as to confirm that no leaking or contamination had occurred during annealing or processing. Grains in each



image were then hand-traced and converted to equivalent mean diameter. In several samples, two perpendicular sides were polished in order to verify that no columnar growth occurred and that the assumption of roughly spherical particles was appropriate (Fig S2). We also took images of selected samples after one month in the 218K freezer, to ensure that no further grain growth happened while samples were in storage (Fig S2F). Grain size was then measured by hand-tracing reflected-light images, converting shapes to pixel area, and calculating the mean grain diameter for area-equivalent circles. We report arithmetic mean grain sizes (diameters) for consistency with previous studies of grain growth in saline ice (De Achaval et al., 1987; Wang et al., 2024).

Dihedral angles were also hand-traced and measured using reflected-light images following methodology described by (McCarthy et al., 2019) (Fig S3). The dihedral angle provides an indicator of melt network connectivity. For instance, at dihedral angles greater $> 60^\circ$, interconnected melt networks will not form and melt will, instead, be confined to isolated pockets at grain junctions (Bulau et al., 1979).

3. Results

3.1 Grain growth

Grain boundaries traced from high-salt (HS) and low-salt (LS) samples are shown in Figure 2. In both the high-temperature (HT) and low-temperatures (LT) experiments, grain boundaries became slightly straighter, while small grains were consumed by larger grains over time (Figures 2A–F, S4).

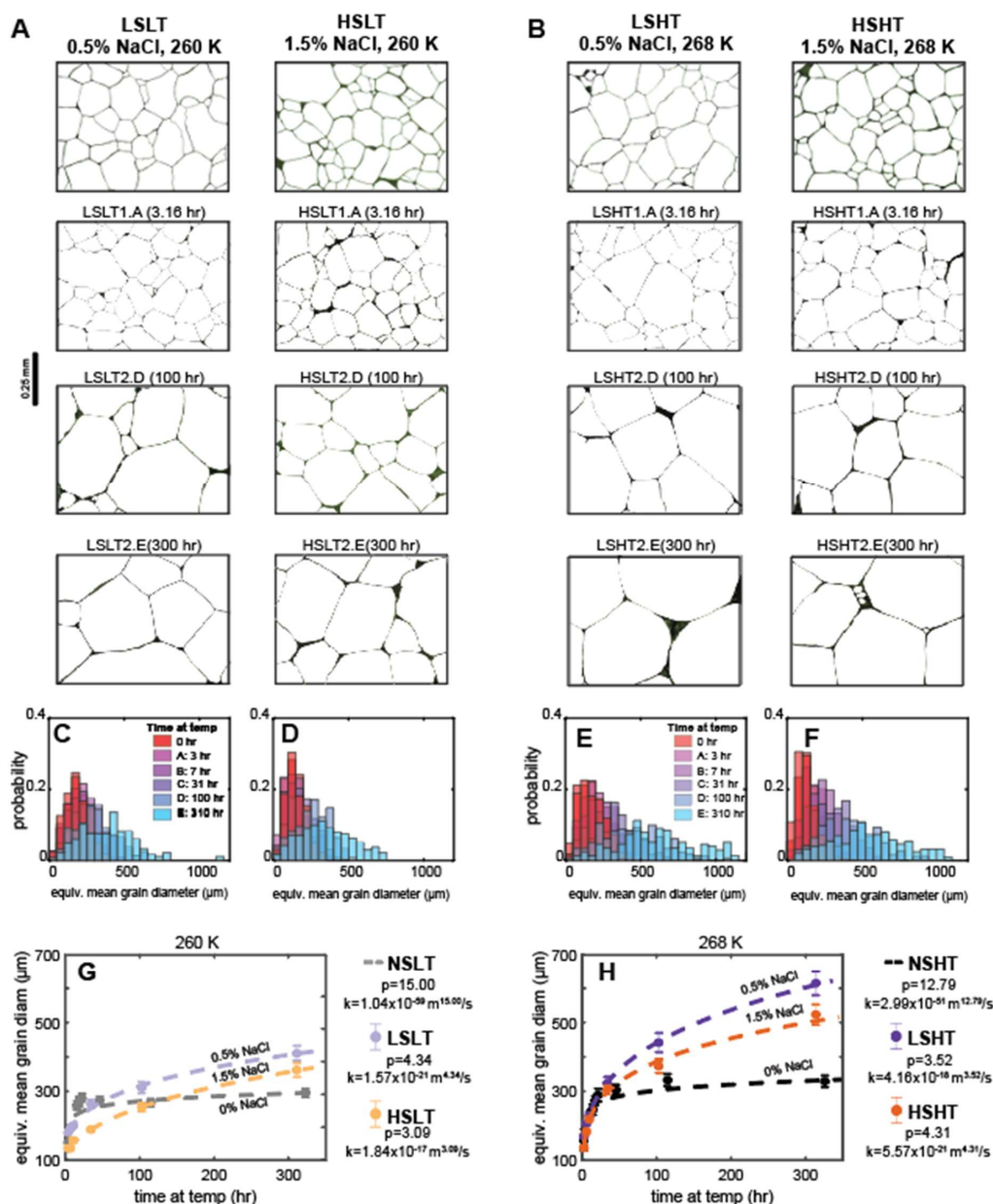


Fig. 2. Grain sizes from growth experiments on ice after saltwater exposure. Traced grain outlines from low-temperature (A) and high-temperature (B) grain growth tests, labeled by the duration of tests and the salinity. Top row of all images shows grain boundaries from starting material. Additional tracings can be found in Figure S1. Histograms (C – F) and grain law fits (G – H) show that the population of larger grains increases with time, and that the shape of the distribution across samples remains somewhat constant as a function of temperature. Lower-salinity



165 grains also consistently grow larger than higher-salinity grains at the same given temperature, and both exceed the
growth rate of salt-free ice.



All experiments shown here are well fit by the standard equation for grain growth (Eq 1; Figure 2E – F; Table S1). We used least-squares regression to derive the best-fit values of the grain-growth exponent, p (limited to values between 0 and 20), and grain-growth constant, k , in each example. In agreement with (Wang et al., 2024) we find that $p = 3\text{--}4.5$ provides the best fit, and that there is no clear relationship between salinity, temperature, and the value of p . This finding also agrees with observations, from silicates, that grain growth in the presence of extrinsic drag forces will have a value of $p > 2.5$ (Evans & Dresen, 1991); higher than the commonly cited $p = 2$ for idealized, single-material grain growth. For ease of comparison with earlier results, we normalized to $p = 4$. We observe that k -values, once normalized to growth exponent of $p = 4$, fall within the range of $1 - 10 \times 10^{-20}$, while those for pure ice at the same temperatures are an order of magnitude lower. Grain growth was thus usually much slower in pure ice samples than in any salt-and-melt-bearing sample at the same temperature (Figure 2G – H, Figure S4, Figure S5); however, p -values were much higher, indicating a quick plateau in grain growth. Furthermore, we note that grain size distribution moves from a lognormal to a more symmetric, normal distribution once grain growth begins to plateau (Figure 2C – F).

The presence of melt therefore modifies grain growth compared to pure ice at a given temperature, but higher salinities limit growth compared to lower salinities. This is sensible in the context of previous findings that salinity impedes crystal growth *below* the eutectic temperature (De Achaval et al., 1987; Jellinek & Gouda, 1969; Wang et al., 2024). Previous experimental work also noted that above the eutectic point in the NaCl-H₂O and KCl-H₂O systems, grain growth is faster than that of pure ice at the same temperature (Chatterjee & Jellinek, 1971; Jellinek & Gouda, 1969; Wang et al., 2024). A closer inspection of results by Wang et al., (Wang et al., 2024) who did not draw conclusions regarding the effects of salinity on grain growth above the eutectic, also shows that grain growth above the eutectic temperature is faster than in samples with a KCl concentration of 10^{-5} M than those with 10^{-2} M KCl (Figure 3). This agrees well with our results, as well as the increased grain boundary mobility inverse to KCl concentration seen by (Nasello et al., 2007).

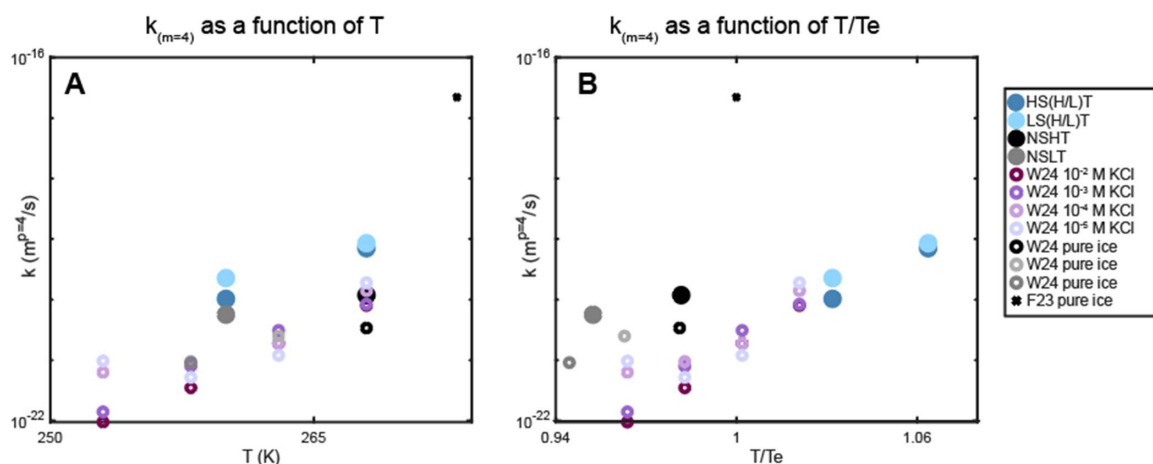


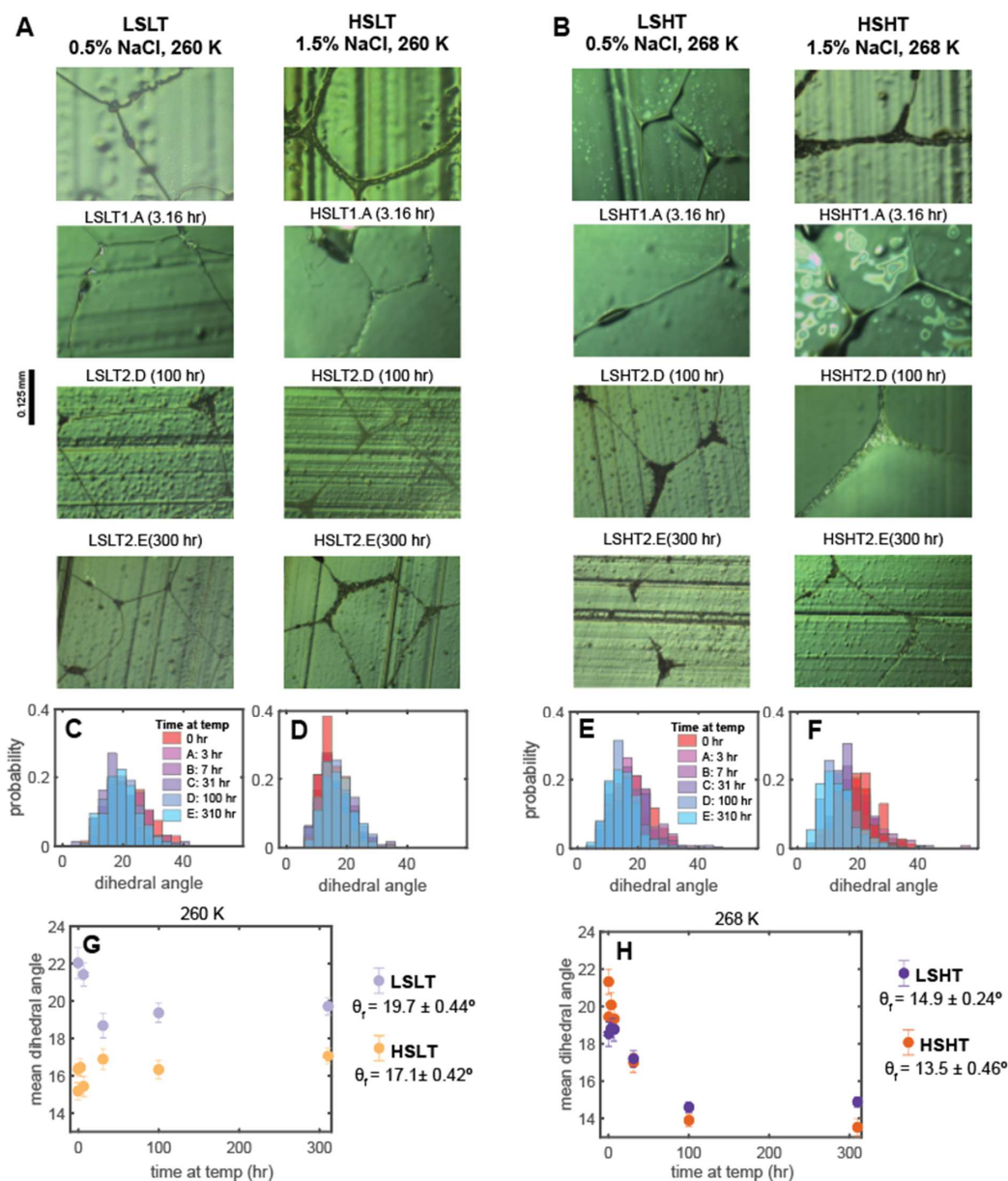
Fig. 3.

Values of grain growth parameter k (for a fixed $p = 4$) for chloride salts mixed with H_2O ice as a function of (A) absolute temperature, and (B) temperature relative to the eutectic temperature, using values from this study and Wang et al (2024), W24, and pure ice to the eutectic temperature from Fan et al (2023), F23. The relationship between salinity and k is not straightforward very close to the eutectic temperature, but below 0.96 T/T_e and above 1.03 T/T_e , higher-salinity samples have lower k values.

3.2 Dihedral angles

Dihedral angles varied modestly as a function of both salinity and temperature (Fig 4). The LS experiments uniformly showed higher dihedral angles than HS experiments, with melt tending to form between grains at triple junctions in LS experiments (Figures 4A – B), while HS experiments showed long melt veins between grains throughout.

All values of θ were below 60° , suggesting that saltwater within icy aggregates will nearly always form an interconnected melt network (Bulau et al., 1979). Dihedral angle also did not evolve appreciably with time in the LT (-13°C) experiments (Figures 4C–D, 4G) but did decrease significantly with time in the HT (-5°C) experiments (Figures 4E–F, 4H), where for both concentrations the θ values flatten to $\sim 14^\circ$. The decrease in θ with saltwater exposure time agrees with previous observations that grain growth may preferentially remove high-energy grain boundaries, and therefore higher dihedral angles (Xu et al., 2023).



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Fig. 4. Dihedral angles. Traced dihedral angles from low-temperature (A) and high-temperature (B) tests show that lower-salinity samples host liquid primarily at triple junctions, while higher-salinity samples host liquid along grain boundaries throughout. Top row of all images is the dihedral angle from starting material. Histograms (C – F) and plots of mean dihedral angle as a



function of time (**G – H**) show that lower-temperature dihedral angles remain roughly constant,
 while higher-temperature samples tend to decrease their dihedral angle with grain growth.

4. Discussion

We note that higher grain sizes are found in the lowest-salinity aggregates at both temperatures observed,
 such that an increased brine fraction limits grain growth. We suggest that increased saltwater therefore
 creates a pinning effect, which may lead to smaller grains and therefore change the effective strength of
 icy material.

4.1 Pinning effects of saltwater on grain growth

The extrinsic drag forces that limit grain growth at melt-grain interfaces can be modeled as a form
 of Zener pinning (see (Alley et al., 1986), wherein the surficial energy of a melt pocket exerts a limiting
 force on grain growth.

The interfacial energy between solid and liquid phases (γ_{SL}), and the interfacial energy between
 solid and solid phases (γ_{SS} , approximately 0.065 mJ/m^2 ; (Ketcham & Hobbs, 1969), are related by
 dihedral angle at a liquid-solid interface, θ , (Figure S5) such that

$$\gamma_{SL} = \frac{\gamma_{SS}}{2 \cos \left(\frac{\theta}{2} \right)}$$

(Eq. 3)

Pinning pressure is then expressed as

$$P_z = \frac{3f \gamma_{SL}}{2r_p} \psi$$

(Eq. 4)

where γ_{SL} is the grain-boundary energy at a solid-liquid interface, r_p is the radius of the melt pocket, and
 ψ is a shape correction term to modify the deviation from assumed spherical particles. This has units of
 J/m^3 , equivalent to Pascals. The dihedral angle θ is also closely associated with changes to geophysical
 properties like seismic velocities (Hier-Majumder & Abbott, 2010), crystallization history (Holness et al.,
 2005, 2012), and short-term anelastic energy dissipation (Yamauchi et al., 2024; Yamauchi & Takei,
 2024).

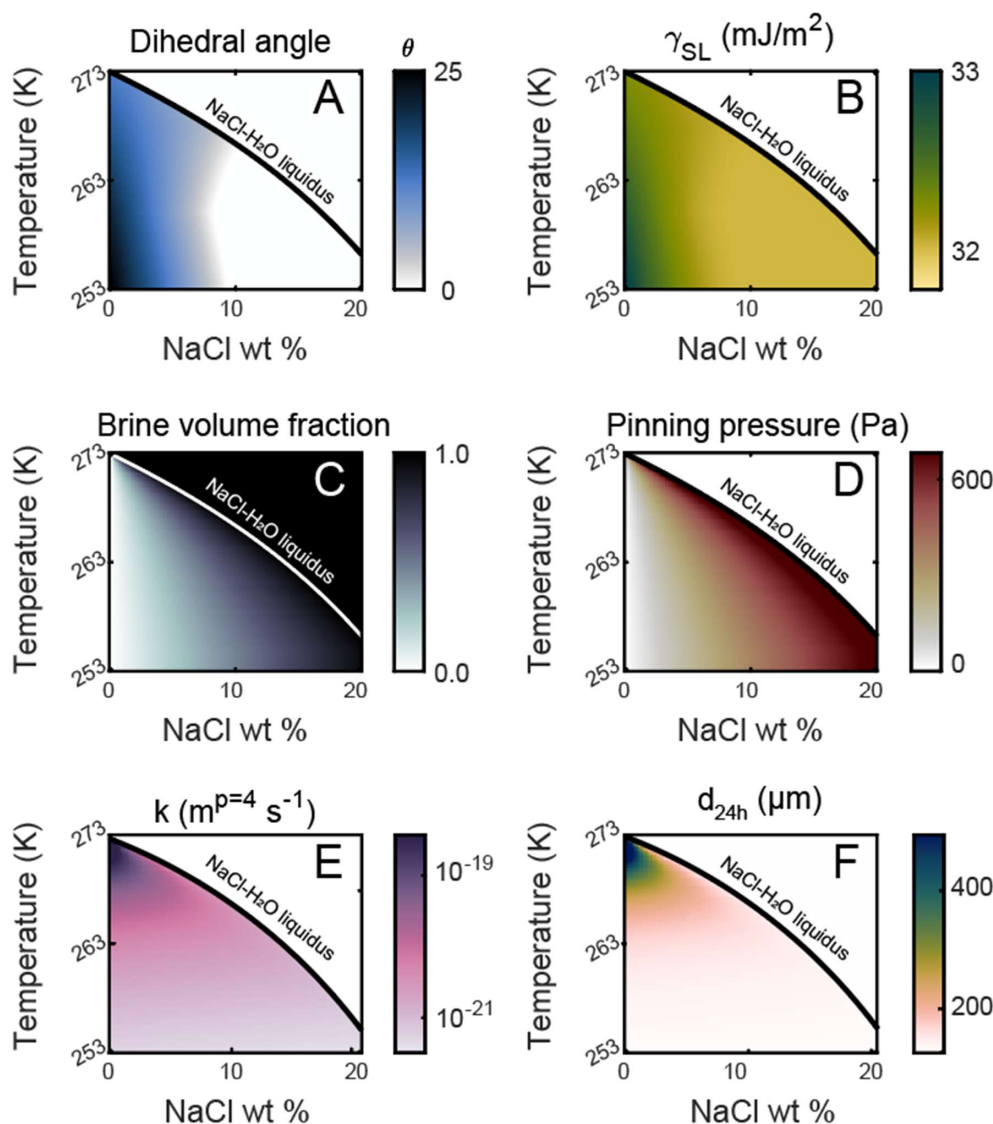


Fig. 5. Pinning calculations up to the NaCl-H₂O eutectic melting curve. **(A)** Dihedral angle extrapolated from experimental range of dihedral angles. Dihedral angles are always low enough to ensure connectivity, and above ~8 wt% NaCl, the dihedral angle falls close to 0, at which point all grain boundaries are fully lubricated. **(B)** γ_{SL} , surface energy at solid-liquid interfaces. Higher energy allows for enhanced grain growth. **(C)** Brine volume fraction as a function of temperature and NaCl wt %. **(D)** Zener pinning pressure of a saline water particle with radius of 50

260



265 μm . (E) Extrapolated k values for chloride salts as a function of weight percentage, temperature, and eutectic
temperature, using a fixed $p = 4$. (F) Projected grain size after 24 hours for a grain of initial diameter $d = 150\mu\text{m}$.



4.1.1 Grain growth parameters as a function of salinity and temperature

To explore all possible effects of saltwater on grain growth, we linearly interpolated dihedral angles in salinity-temperature space up to the NaCl-H₂O eutectic curve (Figure 5A) to calculate γ_{SL} (Eq. 3, Figure 5B), then used the lever rule to calculate effective melt fraction as a function of salinity and temperature (Figure 5C). We then calculated the effective local pinning pressure per unit micron of grain boundary (Eq. 4, Figure 5D). Pinning pressures generally increase with salinity such that, for a given temperature, lower salinities exhibit lower pinning pressure and therefore enhance grain growth compared to higher salinities.

We also parameterized k for impure ice (k_{im}) as a function of both composition and temperature relative to the eutectic point. We compiled k values for a fixed grain growth exponent of $p = 4$ from this study, as well as previously reported values in the KCl-H₂O and MgSO₄-H₂O systems (Wang et al., 2024). We build from the ‘vertical lever rule’ (McCarthy et al., 2019), which relates the temperature, minimum eutectic temperature, and liquidus by

$$V_L(x, T, T_e) = \frac{T_L(x) - T}{T - T_e} \quad (\text{Eq. 5}),$$

where x is chloride salinity as a weight fraction, $T_L(x)$ is the liquidus temperature at a given composition (see Supporting Information; Fig S1B), T is temperature in Kelvin, and T_e is the eutectic temperature of a chloride system in Kelvin.

In doing so, we found that k_{im} can be calculated as a function of melt fraction (f_m), V_L , and T , as

$$k_{im}(f_m, V_L, T) = \exp(A * (1 - f_m) * (V_L + 1)^B) * k_{pure} * \exp\left(-\frac{Q}{RT}\right)$$

(Eq. 6),

where k_{pure} is $6.19 \times 10^{-8} \text{ m}^4/\text{s}$ and Q is $68 \pm 47 \text{ kJ}$, following (Wang et al., 2024)

Using a nonlinear least-squares regression for compiled k -values across systems, we find that $A = 8.209 \pm 0.001$ and $B = -4.75 \pm 0.41$ (Figure S6).



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4.1.2 Grain growth as a function of salinity and temperature

Using Eq. 6, we substitute k_{im} as k in Eq. 1 to calculate theoretical grain size after 24 hours (Figure 5F) in a new setting (i.e., 24 hours of melt exposure if above the eutectic temperature) using an initial $d = 150$ μm . This shows that grain size is indeed highly dependent on salinity and temperature relative to the eutectic temperature, and therefore the melt fraction. As seen in previous grain-growth studies, grain growth becomes considerably faster in the presence of liquid.

Using a modeled $p = 4$ matched results from a recent test on the grain-size-sensitive deformation of temperate ice with small degrees of meltwater present (Schohn et al., 2025), but only when ks are higher than those observed in saline systems (on the order of 10^{-16} – 10^{-14} ; Figure S7). This agrees well with our finding that increased salinity can create a pinning pressure above the eutectic point, and points to the need for increased understanding of grain growth close to 273.15 K.

4.2 Implications for tensile strength of natural ice

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Our finding that salinity and melt control the grain growth of temperate ice provides important insight into the physical properties of terrestrial and planetary cryospheres. For example, there is an orders-of-magnitude difference between measured laboratory strength (Currier & Schulson, 1982; Lee & Schulson, 1988; Litwin et al., 2012; Petrovic, 2003) and inferred natural strength from field observations (Grinsted et al., 2024; Ultee et al., 2020; Vaughan, 1993; Wells-Moran et al., 2025) of temperate ice. We follow the methodology of (Ranganathan & Robel, 2025) to model the dependence of tensile strength, σ_t on grain size, following the relationship

$$\sigma_t = K d^{-\frac{1}{2}} + \sigma_0$$

(Eq. 7)

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Where K and σ_0 are free parameters representing linear grain size dependence and theoretical initial tensile strength at a grain size approaching 0, and d is grain size in meters. We considered the relationship between a pure-ice mass of an ambient temperature, then model grain size as a function of time using modeled k_{im} (Figure 3E), using $K = 0.052 \text{ MPa}^{1/2}$ and $\sigma_0 = 0.5165$ (Ranganathan & Robel, 2025). The choice of K did not significantly change the modeled tensile strength of our experiments (Figure S8). This model simulates settings in which either salty water infiltrates through grain boundaries in ice masses, as has recently been shown to increase fracture susceptibility in glaciers (Chen et al., 2025) and is likely in icy moons containing interstitial brines, or where impure ice experiences warming.

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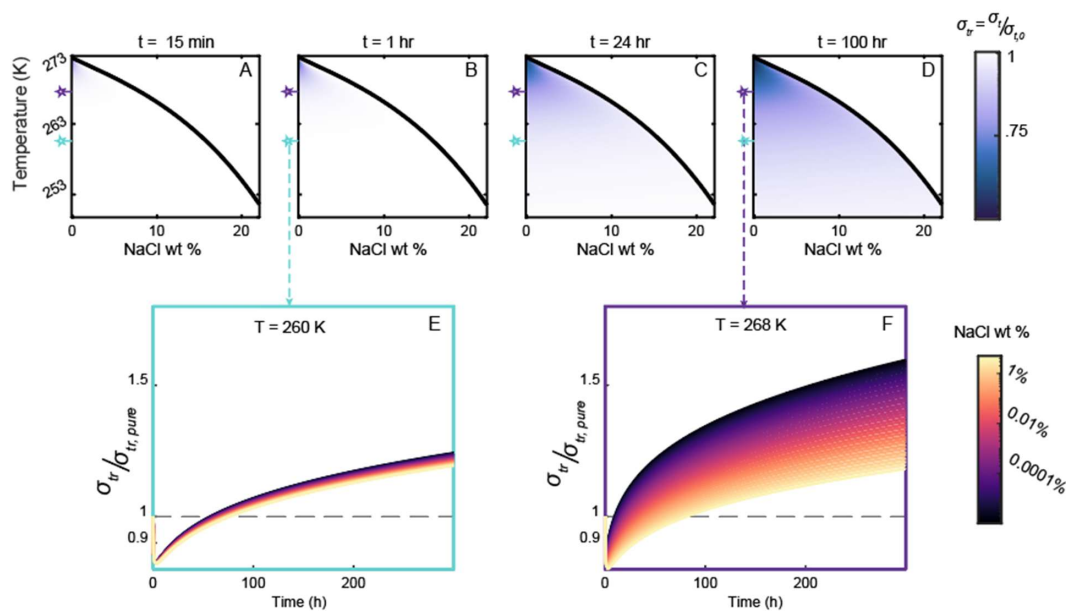


Fig. 6. Effective tensile strength of ices exposed to salty water, (A – D), as a percentage of initial tensile strength $\sigma_{t,0}$, over a period of (A) 15 minutes, (B) 1 hour, (C) 24 hours, and (D) 100 hours, as well as relative to the reduction in tensile strength due to temperature increases alone (normalized to 1 and indicated with a dashed line) in pure ice over the same periods at (E) 260 K and (F) 268 K, colored by NaCl weight percentage.

Saltwater significantly reduces the tensile strength of an icy aggregate in all cases, and lower-salinity infiltration yields the largest decreases in strength (Figure 6; Figure S8). Given a moderate initial tensile strength and 0.25% bulk salinity, these decreases can be up to $22.06 \pm 0.06\%$ after just one hour of saltwater exposure, and up to $40.76 \pm 0.04\%$ after 24 hours. If the salinity increases to 1%, the limiting effect of melt on grain growth caps this tensile strength reduction at $7.99 \pm 0.05\%$ in the first hour and $28.32 \pm 0.06\%$ after 24 hours. Modeled tensile strength reductions in ice exposed to heat, but not saltwater, are initially higher than those from ice exposed to saltwater (Figure 6E – F), but lower-salinity saltwater can match the effects of temperature alone within 12 hours, and cause reductions in tensile strength up to 1.67 times that of temperature alone. Models of tensile strength based on saltwater infiltration should thus predict a buttressing effect over hourslong timescales (relative to decreases from heat alone) that depends on the amount of saltwater present, but increasingly weaker ice masses with time.

Our results contribute to growing observational evidence that rheological and compositional differences in glacier ice impede or impel failure (Borstad et al., 2017; King et al., 2018; Kulesa et al.,



2019), and that liquid infiltration can cause severe ice shelf instability (Suganuma et al., 2025). Furthermore, the presence of subsurface brine sills like those hypothesized to exist on Europa (Barr & Showman, 2009; Chivers et al., 2023; Steinbrügge et al., 2020; Wolfenbarger et al., 2022) means that interiors of icy moon cryospheres are likely weaker than expected, explaining features like putative
 360 cryovolcanic vents and fracture surfaces. The co-evolution of grain size and dihedral angle together also suggest that most of the short-term energy dissipation seen as icy moons orbit their host planets will occur within 100 Earth hours of enhanced saltwater infiltration, as reduced dihedral angles and slowed grain growth indicate lower grain boundary energy. As saltwater infiltrates into icy crusts, melt networks can form and reduce local strengths and viscosities, leading to enhanced grain growth in regions of
 365 infiltration. Further study is required to understand the competition between grain growth and melt lubrication in determining the viscosity of ice after saltwater infiltration.

5 Conclusions

We find that the grain size and dihedral angle, two key parameters for interpreting the
 370 microstructural behavior of partially molten icy aggregates, are modified significantly in the presence of saltwater. While any amount of saltwater changes the rate of grain growth and therefore grain size compared to pure ice, we find that increased salinity of an aggregate can reduce grain size and dihedral angle compared to lower salinities at the same temperature. We also observed that elevated temperature may further reduce dihedral angles in both higher and lower salinity ice aggregates. Our results also agree
 375 with recent experimental studies showing that a grain growth exponent of $p \sim 4$ may be more appropriate for salty aggregates than the commonly used values of $p = 2$ for pure ice or $p > 6$ for bubbly ice. We also established a relationship of grain growth parameter k that depends on both salinity and temperature for chloride salts mixed with ice.

The limiting effect of increasing salinity and saltwater volume fraction is well explained by the
 380 physics of grain pinning, and can be extrapolated over a wide range of temperature-salinity conditions to better predict the dynamics of ice in contact with saltwater or containing saline inclusions.

We show that the presence of saltwater can induce a much stronger reduction in tensile strength compared to the effects of temperature alone, and that increasing salinity enhances this effect. Preliminary results show that grain growth in the presence of saltwater may reconcile orders-of-magnitude differences
 385 in tensile strength between lab and natural settings. Our finding that saltwater weakens icy aggregates by enhancing grain growth, but that increased salinity may suppress grain growth, suggests that future studies coupling local salinity and temperature gradients could be strong predictors of glacial mass loss.



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Author contribution:

Conceptualization: CS

Methodology: CS, CM

400 Investigation: CS, MB, NW

Validation: CS, MB

Writing – original draft: CS, CM, AJC, JM

Formal analysis: CS, MB

Supervision: CS, CM, AJC

405 **Competing interests:** Authors declare that they have no competing interests.

Code/data availability: A MATLAB structure storing all images and processed data used in this
work, along with a sample script for processing grain means and pinning pressures are available at
(Seltzer, 2025). Physical samples are stored at Lamont-Doherty Earth Observatory and may be
accessed on request.

410 **Supplementary Materials**

Supplementary Text

Figs. S1 to S8

Table S1



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