

1 **Polarity and direction dependence of energetic cross-frontal eddy transport in the Southern**  
2 **Ocean's Pacific sector**

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29 **Abstract:**

30 Mesoscale eddies play a critical role in mediating meridional transport across the Antarctic Circumpolar Current  
31 (ACC), yet the dynamics of cross-frontal eddies (CFEs) and their energy exchanges with frontal jets remain  
32 inadequately quantified. This study presents a systematic analysis of CFE characteristics, kinetic energy evolution,  
33 and thermohaline transport effects in the Pacific sector of the Southern Ocean, utilizing 23 years (2000–2022) of  
34 satellite altimetry and Argo float data. Our results reveal a fundamental polarity- and direction-dependent  
35 asymmetry in CFE dynamics. Equatorward-propagating cyclonic eddies (CEs) dominate CFE activity, followed by  
36 poleward-moving anticyclonic eddies (AEs). These dominant CFE types exhibit superior energetic characteristics,  
37 including significantly higher eddy kinetic energy (*EKE*) and stronger nonlinearity compared to their counterparts.  
38 Complete CFEs experience polarity- and direction-selective energy gains during frontal crossing, with equatorward  
39 CEs and poleward AEs extracting energy from eastward frontal jets, while their counterparts lose energy. This  
40 energization mechanism has intensified over the past two decades, with both polarity CFEs showing substantial  
41 *EKE* increases that substantially exceed previous basin-scale estimates. Hydrographic analysis demonstrates that  
42 CEs and AEs transport distinct water masses across frontal boundaries, creating sharp thermohaline contrasts within  
43 interfrontal zones. Our findings establish CFEs as crucial regulators that buffer wind- and warming- induced  
44 baroclinicity increases through compensatory heat transport, thereby maintaining the Southern Ocean's thermal  
45 equilibrium and modulating the ACC's response to external forcing in a changing climate.

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48 **Keywords:** Cross-frontal eddies; Eddy kinetic energy; Meridional transport; Antarctic Circumpolar Current; the  
49 Pacific sector of the Southern Ocean

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## 59 **1 Introduction**

60 Mesoscale eddies are ubiquitous in the Southern Ocean (SO), play a vital role in the zonal and meridional  
61 transport of quantities including heat and momentum across the Antarctic Circumpolar Current (ACC), and also  
62 influence the uptake of heat and carbon dioxide from the atmosphere (Moreau et al., 2017; Patel et al., 2019; Sallé  
63 et al., 2008; Sokolov and Rintoul, 2007) and the transport and connectivity of marine species (e.g., Duan et al.,  
64 2020; Zhu et al., 2025). The ACC comprises multiple zonal fronts, where oceanic jets exit. Here, a front refers to a  
65 boundary between distinct water masses, characterized by strong horizontal density gradients, while a jet denotes  
66 the narrow, swift current that flows along the axis of such a front. Together, they form the dynamic core of the  
67 ACC. At these fronts, mesoscale activity is enhanced, with higher eddy kinetic energy (*EKE*) and more frequent  
68 eddy generation and dissipation (Barthel et al., 2017; Gille, 1994; Hughes, 1995; Hughes and Ash, 2001; Morrow  
69 et al., 1994; Sokolov and Rintoul, 2002). In turn, eddies impact the fronts' structure, intensity, and location. For  
70 instance, eddies may accelerate the jets, and cyclonic (anticyclonic) eddies cause the equatorward (poleward)  
71 deviation of frontal meanders in some cases (Chapman et al., 2020; Duan et al., 2016; Frenger et al., 2015; Sprintall,  
72 2003). These interactions between mesoscale eddies and oceanic fronts can shape local thermohaline structures,  
73 exert profound influences on large-scale circulation and vertical flux processes. They also have tremendous  
74 implications for the redistribution and survival of marine species and the stability of the climate system.

75 In the SO, the transition from the warm subtropical waters to the cold Antarctic waters is not smooth but  
76 concentrated along a series of fronts (Orsi et al., 1995; Belkin and Gordon, 1996), often corresponding to the  
77 locations of narrow, high-speed currents known as “jets” (Sokolov and Rintoul, 2002, 2007). These fronts delineate  
78 the boundaries of distinct water masses, each with unique environmental characteristics (Orsi et al., 1995). The  
79 existence of fronts hinders meridional exchanges of heat and tracers (Chapman and Sallé, 2017; Naveira Garabato  
80 et al., 2011; Thompson and Sallé, 2012). At the same time, eddies enable cross-frontal transport and serve as  
81 primary carriers for meridional water mass properties, including heat (De Szoek and Levine, 1981; Foppert et al.,  
82 2017). Cross-frontal eddies (CFEs) must overcome intense geostrophic shear to achieve material transport and  
83 render their dynamical contributions to meridional transport, which is significantly more pronounced compared to  
84 other eddy types (Thompson and Sallé, 2012).

85 Eddies' capability of trapping materials and achieving long-distance cross-frontal transport helps in mitigating  
86 sharp meridional hydrographic gradients, facilitating new water formation and carbon transport, and also enhancing  
87 subsurface temperature extremes in the SO. Holte et al. (2013) presented that cross-frontal exchanges by eddies can  
88 penetrate strong potential vorticity gradients associated with the Subantarctic Front (SAF) and facilitate the

89 downstream evolution of Subantarctic Mode Water by transporting cold, low-salinity water across the ACC from  
90 the Polar Front Zone (Holte et al., 2013). In a study of a cold eddy in the southwest Indian Ocean, Swart et al.  
91 (2008) found that the eddy displaced temperature and salinity anomalies by  $1.5^\circ$  towards lower latitudes. This  
92 single eddy contributed 2.5% of the annual northward flux of Antarctic Surface Water in the southwest Indian  
93 sector (Swart et al., 2008). In addition, eddies induce carbon transport across the ACC, which alters the carbon  
94 properties and budget of the Subantarctic Zone waters (Moreau et al., 2017). Patel et al. (2019) proposed that about  
95 21% of the heat transported across the SAF to the Subantarctic Zone south of Tasmania is achieved by cyclonic  
96 eddies. He et al. (2023) demonstrated that nearly half of the subsurface temperature extremes in the Southern Ocean  
97 occur within eddies, with cross-frontal eddies (CFEs) generating extremely high-temperature events on the cold  
98 side of the ACC and extremely low-temperature events on the warm side. These extremes eventually impact marine  
99 organisms and ecosystems. For instance, *Electrona carlsbergi* in the high-latitude Antarctic region may be  
100 transported across the fronts from the Argentine Basin by the poleward eddy activity (Saunders et al., 2017; Zhu et  
101 al., 2025).

102 Eddies in the SO can moderate the ACC's response to surface wind forcing changes, namely the "eddy  
103 saturation" hypothesis (Hallberg and Gnanadesikan, 2001, 2006; Straub, 1993). Reanalysis of data since 1972 show  
104 an increasing trend in wind stress (associated with a positive trend of Southern Annular Mode) over the Pacific  
105 sector that dominates the basin-wide wind stress variability, driving enhanced eddy activity responses in this sector,  
106 with *EKE* intensifying at a rate of  $14.9 \pm 4.1 \text{ m}^4 \text{ s}^{-2}$  per decade (Duan et al., 2016; Hogg et al., 2015; Menna et al.,  
107 2020; Morrow et al., 2010). Recent work by Zhang et al. (2021) demonstrates that *EKE* intensification is not  
108 spatially homogeneous in the SO but concentrated south of New Zealand and downstream of the Campbell Plateau  
109 in the Pacific sector. This localized enhancement likely stems from the release of available potential energy stored  
110 in tilted isopycnals, thus acting to moderate the eastward flow in the ACC, which has significantly intensified  
111 between  $48^\circ\text{S}$  and  $58^\circ\text{S}$  mainly due to buoyancy forcing (Shi et al., 2021). Mesoscale energy gain from mean  
112 flows is achieved through baroclinic (primary) and barotropic (secondary) pathways (Fu et al., 2023). Regarding  
113 topographic effects, previous studies have established that interactions between ACC and seafloor topography  
114 intensify oceanic eddy mixing by enhancing downstream baroclinic shear. This process enhances eddy generation  
115 and increases *EKE* downstream of major topographic features (Frenger et al., 2015; Morrow et al., 1992; Park et al.,  
116 1993; Thompson and Sallée, 2012). Consequently, ACC frontal jets with strong geostrophic characteristics  
117 experience mesoscale eddy modulation near prominent topographies (Kim and Orsi, 2014; Thompson et al., 2010).

118 Despite extensive research on basin-scale *EKE* modulations and case studies of CFE transport, and the

119 well-established asymmetric eddy distribution on both sides of the ocean fronts, fundamental questions remain  
120 regarding how eddy-jet interactions vary based on eddy characteristics and directional approach in the SO.  
121 Specifically, it is essential to understand: (1) the polarity and direction preferences of CFEs during frontal crossing;  
122 (2) the magnitude and pattern of kinetic energy change within eddies following frontal crossing; (3) the resultant  
123 hydrographic property redistribution achieved by CFEs in the interfrontal zones. Motivated by these research gaps,  
124 we conducted a systematic assessment of cross-frontal mesoscale eddies in the Pacific sector to elucidate their role  
125 in regional ocean dynamics and hydrographic redistribution. Utilizing 23 years (2000–2022) of satellite altimetry  
126 data, we characterize the spatiotemporal variability, *EKE* patterns, and eddy-jet interactions of CFEs in the Pacific  
127 sector. We complement these analyses with Argo (Array for Real-time Geostrophic Oceanography) float profiles to  
128 quantify normalized hydrographic differences between cyclonic and anticyclonic eddies within the interfrontal  
129 zones. These approaches aim to improve our understanding of the dynamic characteristics of CFEs and their role in  
130 mediating meridional transport across the ACC in this sector.

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## 149 2 Data and methods

### 150 2.1 Data

151 This study focuses on the SO's Pacific sector in the range of 150°E–110°W, 35°S–80°S. Prominent  
152 topographic features within this region include the Campbell Plateau, Pacific-Antarctic Ridge, Eltanin Fracture  
153 Zone, and Udintsev Fracture Zone (Figure 1). We utilized the gridded satellite altimeter data for eddy detection and  
154 tracking. This dataset is merged from multiple satellites and provided by the Copernicus Marine Service as the  
155 CMEMS all-satellite L4 SLA product, SEA LEVEL\_GLO\_PHY\_L4\_MY\_008\_047 (Copernicus Marine Service,  
156 2024). The data were accessed/downloaded on 18 December 2025 from the Copernicus Marine Service Information  
157 (DOI: <https://doi.org/10.48670/moi-00148>). It includes daily Sea Level Anomaly (*SLA*) and sea surface geostrophic  
158 velocity anomalies ( $u'$ ,  $v'$ ) data during 2000–2022 with a spatial resolution of  $0.125^\circ \times 0.125^\circ$ . The *SLA*  
159 represents the sea surface height anomaly relative to the mean sea surface from 1993 to 2012. The  
160 corresponding geostrophic velocity anomalies ( $u'$ ,  $v'$ ) are derived from *SLA* based on the geostrophic  
161 balance (Pujol and Grassi, 2025).

162 The geographical positions of the ACC's fronts and boundaries used in this study are from the synthesis of  
163 Park et al. (2019). This dataset provides the most updated mapping of the ACC frontal system and its associated  
164 boundaries, derived from satellite altimetry and independently validated against extensive subsurface observations,  
165 including Argo float profiles (2001–2017) and dedicated CTD surveys (2016–2017). As shown in Figure 1, the  
166 dataset defines five major streamlines from north to south: the Northern Boundary (NB), the Subantarctic Front  
167 (SAF), the Polar Front (PF), the Southern ACC Front (SACCF), and the Southern Boundary (SB). Specifically, the  
168 NB represents the northern dynamical limit of the ACC and coincides with the northern expression of the  
169 Subantarctic Front system (SAF-N) in this region. The SAF, PF, and SACCF correspond to the core frontal jets.

170 Furthermore, a total of 1094 quality-controlled Argo profiles (0–2000 m; <https://argo.ucsd.edu/data/>) located  
171 within detected eddies were utilized to analyze the internal thermohaline structure of cyclonic (CEs) and  
172 anticyclonic eddies (AEs). The potential temperature ( $\theta$ ) and salinity ( $S$ ) within each eddy were normalized radially  
173 by binning profiles according to their relative distance from the eddy center (normalized by the eddy radius,  $R$ ), at  
174 an average interval of  $0.03R$ . Owing to limited spatial coverage, profiles from the SB-SACCF zone and areas south  
175 of the SB were excluded. Consequently, the analysis focused on the northern inter-frontal zones of SAF-NB,  
176 PF-SAF, and SACCF-PF, which contained 400, 150, and 252 profiles, respectively. Furthermore, temporal  
177 variability (e.g., interannual and seasonal) was not considered in this composite analysis due to the uneven  
178 distribution of profiles over time.

## 179 2.2 Eddy detection, tracking and CFE categorization

180 We combined the Okubo-Weiss (OW) parameter method with the outermost closed contour of *SLA* to detect  
 181 eddies. As a widely used eddy detection method, the OW parameter method was developed based on flow field  
 182 deformation by high vorticity or high strain (Okubo, 1970; Weiss, 1991). The OW parameter is defined as:

$$183 \quad W = s_n^2 + s_s^2 - \omega^2, \quad (1)$$

184 where  $s_n = \frac{\partial u'}{\partial x} - \frac{\partial v'}{\partial y}$  and  $s_s = \frac{\partial v'}{\partial x} + \frac{\partial u'}{\partial y}$  are the normal and shear components of strain, respectively, and

185  $\omega = \frac{\partial v'}{\partial x} - \frac{\partial u'}{\partial y}$  is the relative vorticity of the flow. The sign of  $W$  determines a region to be strain-dominated ( $W > 0$ )

186 or vorticity-dominated ( $W < 0$ ). Eddies are highly vorticity-dominated circulations, thus corresponding to coherent  
 187 negative- $W$  areas (for both CEs and AEs), and the negative  $W$  must be larger than that for the background field  
 188 (Henson and Thomas, 2008).

189 To identify physically consistent eddy boundaries, we adopt the hybrid geometric–physical approach validated  
 190 by Saraceno and Provost (2012), which avoids common biases associated with fixed  $W$  thresholds. A threshold of  
 191  $W < -0.2 \sigma_w$  is often used to delineate eddy boundaries, with  $\sigma_w$  being the standard deviation of  $W$  over the entire  
 192 region (e.g., Henson and Thomas, 2008; Isern-Fontanet et al., 2006; Frenger et al., 2015). However, this method  
 193 can underestimate eddy area in certain regions (Matsuoka et al., 2016) and misidentify meanders as eddies in  
 194 energetic frontal zones (Saraceno and Provost, 2012). Our approach proceeds as follows: after detecting the eddy  
 195 center using the OW-based method, we identify the outermost closed *SLA* contour that encloses this point. This  
 196 contour defines the eddy boundary, and the center was then recalculated as its geometric centroid. Eddy radius was  
 197 computed as the radius of a circle of equivalent area, and eddy amplitude is the absolute *SLA* difference between  
 198 the center and along the contour.

199 For eddy tracking, the algorithm identifies eddies at time  $t+1$  that meet the following criteria relative to time  $t$ :  
 200 (1) minimal centroid distance, (2) identical polarity (i.e., rotation direction), and (3) the minimum radius variation.  
 201 If no eddy at  $t+1$  satisfies these proximity thresholds for a given eddy at  $t$ , the eddy is considered dissipated.  
 202 Conversely, if an eddy detected at  $t+1$  does not match any eddy at  $t$ , it is classified as a newly generated eddy.

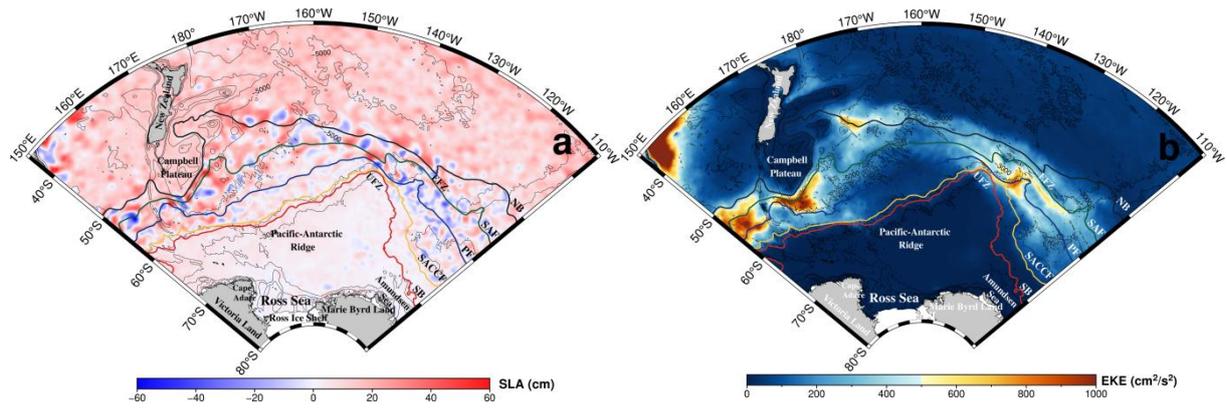
203 To ensure statistical robustness, our analysis focused exclusively on significant mesoscale eddies, which were  
 204 defined as well-resolved, energetic eddies with sufficient temporal coherence. Eddies meeting all of the following  
 205 criteria were retained as the core dataset for subsequent analyses: (1) radius  $> 30$  km; (2) amplitude  $> 5$  cm; and (3)  
 206 lifespan  $> 14$  days.

207 The *EKE* was computed from geostrophic velocity anomalies using the equation  $EKE=(u^2+v^2)/2$ . In this  
 208 study, except for Figure 1, all analyses of *EKE* variation during cross-frontal processes were based on the total *EKE*  
 209 within each eddy interior ( $EKE_T$ ) for better tracking *EKE* changes in specific eddies, calculated as  $EKE_T =$   
 210  $\sum_{i=1}^N EKE_i \cdot ds$ , where  $EKE_i$  is the *EKE* for grid  $i$ ,  $N$  is the grid amount within an eddy, and  $ds$  is the grid area. The  
 211 eddy nonlinearity parameter ( $\beta$ ) was computed based on  $\beta = U/C$ , where  $U$  is the maximum circum-average  
 212 geostrophic velocity within the eddy, and  $C$  represents the eddy's transporting speed (Chelton et al., 2011). The  
 213 eddy is nonlinear when  $\beta > 1$ , indicating the presence of trapped fluid parcels advected with the eddy movement.

214 While climatological fronts define the ACC's mean structure (Park et al., 2019), their positions exhibit  
 215 meridional variability influenced by both bathymetry and eddy activity (Kim and Orsi, 2014; Thompson et al.,  
 216 2010). Fronts stabilize over major bathymetric features (e.g., the Pacific-Antarctic Ridge) but show maximum  
 217 variability in flat basins. Due to eddy-mean flow interaction processes, frontal zones become greater downstream of  
 218 topographic obstacles like the Campbell Plateau. Based on the observed frontal variability in the Pacific sector, the  
 219 maximum total meridional frontal drift during 1993–2010 is approximately 80 km southward (at 150 °E), and  
 220 annual cycle amplitude is  $< 40$  km (Kim and Orsi, 2014). To account for these frontal displacements, we first  
 221 defined a baseline frontal zone as a  $\pm 15$  km strap in the normal direction from each climatological front. Then, to  
 222 objectively identify eddy-front interactions, we applied a geometric criterion: an eddy was considered interacting  
 223 when its boundary contacted the frontal zone. Since all analyzed eddies have a radius ( $R$ )  $> 30$  km, this criterion  
 224 effectively creates a dynamic interaction zone with a minimum half-width of 45 km (15 km + 30 km), which  
 225 comfortably exceeds the observed ranges of frontal variability. To further ensure robustness, we conducted a  
 226 sensitivity analysis by expanding the frontal zones to a  $\pm 25$  km strap (see Supplementary Materials), which  
 227 confirmed that all key findings are insensitive to the exact zone definition.

228 The eddy-front interaction was then divided into three sequential phases: pre-cross-frontal, cross-frontal, and  
 229 post-cross-frontal. CFEs were further categorized into four types: (1) Front-generated eddies, generated within the  
 230 dynamic interaction zone and subsequently propagating away, (2) Front-dissipated eddies, propagating into the  
 231 dynamic interaction zone and dissipated there, (3) Transient frontal eddies, both generated and dissipated within the  
 232 same interaction zone, and (4) Complete CFEs, undergoing all three phases (pre-, cross-, and post-frontal) relative  
 233 to the dynamic interaction zone. Both types (1) and (1) eddies were collectively classified as partial CFEs.  
 234 Hereafter, all frontal zones refer to dynamic interaction zones. Notably, according to the definition, the partial  
 235 frontal crossing eddies of Type 1 include rings pinched off from the meandering structures of a front.

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Figure 1. Study region. (a) Sea Level Anomaly (SLA) distribution on January 20<sup>th</sup>, 2022. (b) Spatial distribution of mean eddy kinetic energy (EKE) during 2000–2022. Thick colored lines from north to south represent the northern boundary of ACC (NB), the Subantarctic Front (SAF), the Polar Front (PF), the Southern Antarctic Circumpolar Current Front (SACCF) and the southern boundary of ACC (SB; Park et al., 2019). Significant seafloor topographies have been labeled, with UFZ denoting the Udintsev Fracture Zone. The abbreviation EFZ denotes the Eltanin Fracture Zone.

## 272 **3 Results**

### 273 **3.1 Analysis of CFE characteristics**

274 CFE transport is active across all major ACC frontal zones in the Pacific sector (Figure 2). At each front,  
 275 equatorward-moving CFEs consistently outnumber poleward-moving eddies (Figure 2b). CEs dominate  
 276 equatorward CFEs, with CEs outnumbering AEs by a factor of  $\geq 1.5$ , while AEs prevail in poleward CFE motions.  
 277 The resulting hierarchy of CFE prevalence is as follows: equatorward CEs are the most frequent (36% of total  
 278 CFEs), followed by poleward AEs (30%) and equatorward AEs (18%), with poleward CEs (16%) being the least  
 279 frequent. Among the different frontal zones, the SAF hosts the most eddy occurrences (30% of total CFEs),  
 280 followed by the PF (27%), reflecting intense eddy-mean flow interactions around these two fronts. The  
 281 northernmost NB (23% of total CFEs) and the SACCF (16%) exhibit comparable and moderate CFE levels, while  
 282 the southernmost SB (7%) displays the lowest CFE exchanges.

283 The frontal system exhibits strong meandering patterns due to topographic steering, accompanied by spatially  
 284 heterogeneous CFE distributions. CFE occurrence peaks downstream of prominent topographic features,  
 285 particularly near the Campbell Plateau (150 E–180 E; 41% of total CFEs) and downstream of the Udintsev  
 286 Fracture Zone (125 W–160 W; 38%), where multiple fronts converge (Figure 2a). Eddies may cross multiple  
 287 fronts sequentially at these frontal convergent regions. The majority of eddies cross a single front (Figure 2c).  
 288 Double-frontal crossings (total 411) occur preferentially at southern fronts (SACCF/SB; > 50% of cases; Figure 2d).  
 289 Triple-frontal crossings are rare and primarily limited to the PF/SACCF/ SB system (Figure 2e), and no instances  
 290 of quadruple-frontal crossings were observed.

291 Consistent with the ACC dynamics, over 70% of cross-frontal CEs and AEs propagate eastward (Figure 3), a  
 292 notable contrast to the typical westward propagation of mesoscale eddies driven by Rossby waves in other ocean  
 293 basins (Frenger et al., 2015). Short-distance movements ( $< 2^\circ$ ) are more frequent for AEs at each front. In contrast,  
 294 CEs dominate long-distance propagation, particularly at the SAF and PF. The majority of CEs propagate northward  
 295 (over 64% of all CEs), while over 54% of AEs are oriented southward. The greater energetic content of CEs, as  
 296 indicated by their dominance in long-distance transport, aligns with their role as primary agents of meridional  
 297 exchange. These patterns highlight how ACC-steered eddy motions facilitate distinct transport pathways, with CEs  
 298 disproportionately driving long-distance exchanges, particularly at major fronts such as the SAF and PF.

299 Most CFEs (99%) exhibit nonlinear characteristics (Figure 3k), confirming their capability to trap and advect  
 300 water mass during their lifespan. In the nonlinearity regime ( $\beta > 1$ ), equatorward-moving CEs constitute 72.6% of  
 301 the total cross-frontal CEs, and 64.5% of the total cross-frontal AEs are poleward-moving ones (Figure 3k). In the

302 high nonlinearity regime ( $\beta > 5$ ), the proportion of CEs is notably higher than AEs, consistent with the greater  
 303 dynamic vigor of CEs observed in the above analyses. Therefore, the cross-frontal transport achieved by eddies,  
 304 primarily equatorward-moving CEs and poleward-moving AEs, can facilitate the redistribution of distinct source  
 305 water masses and reduce thermohaline gradients across frontal zones.

306 Cross-frontal CEs and AEs show similar distributions in lifespan, propagation distance, and size (Figure 4).  
 307 Both types show a steep decline in abundance with increasing lifespan. Eddies with lifespans  $\leq 50$  days dominate,  
 308 constituting 55% of the total eddy population, while only 3% exceed 200 days (Figure 4a). Propagation distances  
 309 are confined predominantly to  $\leq 400$  km (56% of total CFEs). CEs slightly outnumber AEs at longer distances  
 310 (400–1000 km; Figure 4b). Size distributions reveal that  $\sim 75\%$  of the total sample have mean radii of 30–50 km  
 311 (Figure 4c). Notably, CEs dominate at smaller radii ( $< 50$  km), while AEs prevail among larger eddies. This  
 312 distribution pattern is consistent with maximum radius statistics (Figure 4d). These CFE characteristics align with  
 313 previously reported eddies in the Pacific sector (Duan et al., 2016).

314 Subsequently, we found distinct characteristics among types regarding their behaviors, when dividing the  
 315 CFEs into partial CFEs, generated within and subsequently transported away (Type 1) and transported into and  
 316 dissipated within the frontal zones (Type 2); transient CFEs, both generated and dissipated within the same frontal  
 317 zone (Type 3); and complete CFEs, experiencing pre-crossing, crossing, and post-crossing phases (Type 4).  
 318 Transient CFEs dominate numerically, accounting for 48% of all CFEs, and partially generated and dissipated  
 319 CFEs constitute 23% and 20%, respectively (Figure 4a–d). These proportions collectively indicate that the frontal  
 320 zones primarily act as terminal/starting areas for eddy life cycles, rather than a simple transit pathway. The  
 321 proportion of transient CFEs falling within low-value parameter ranges is substantially higher than that of the other  
 322 types: 59% of these eddies have lifespans  $\leq 40$  days and propagation distances  $\leq 300$  km, as well as 58% have  
 323 mean radius  $\leq 43$  km and 63% have maximum radius  $\leq 60$  km. These values confirm the intrinsic nature of  
 324 transient eddies as “generated and dissipated locally”, and reveal the constraining role of the frontal system on eddy  
 325 evolution. In stark contrast, completely transported CFEs exhibit markedly different dynamical characteristics: 81%  
 326 have lifespans  $> 40$  days, 90% propagate  $> 300$  km, and 68% have maximum radii  $> 60$  km, indicating that these  
 327 eddies have completely escaped the constraints of the local frontal environment and possess the capability for  
 328 long-distance cross-frontal transport. Notably, among completely transported CFEs, small-scale CEs dominate  
 329 significantly, with CEs accounting for 63%, while AEs account for only 37% for eddies with mean radii of 30–50  
 330 km and maximum radii  $< 70$  km. This polarity bias suggests that small-scale CEs may possess higher transport  
 331 efficiency in cross-frontal material and energy exchange due to their unique dynamical structure.

332 Quantitative analysis of CFE types reveals distinct frontal-zone behaviors (Table 1). Transient eddies (Type 3)  
 333 account for the largest proportion overall ( $> 40\%$  by summing the transient AEs and CEs at each front), particularly  
 334 at the two weaker southern fronts (SACCF and SB). Their proportions are lower at the SAF and PF, indicating that  
 335 these two major fronts host more eddies that interact with areas outside the frontal zone during their lifecycle. For  
 336 partially front-generated eddies (Type 1), both AEs and CEs exhibit relatively high proportions at the SAF and PF.  
 337 At the southern SACCF and SB, however, the proportion of AEs drops markedly, whereas CEs show no such  
 338 reduction. Among partially front-dissipated eddies (Type 2), the three northern fronts consistently show a higher  
 339 proportion of CEs than AEs. At the two southern fronts, the pattern reverses, CE proportions decline sharply,  
 340 causing a higher proportion of AEs. This suggests that in the southern fronts, local cross-frontal CEs are more  
 341 readily generated and propagated outward, while being relatively resistant to dissipation.

342 Regarding the partial CFEs, the proportion of front-generated AEs consistently exceeds that of front-dissipated  
 343 AEs (except the SB), indicating that AEs are more likely to be generated within the fronts than to dissipate locally,  
 344 particularly at the three northern fronts. For complete CFEs (Type 4), the proportion of AEs decreases with frontal  
 345 latitude, while CEs reach their maximum proportions at the SAF and PF. These results demonstrate that different  
 346 CFE types exhibit distinct behaviors when interacting with each front, shaped by frontal dynamics and latitudinal  
 347 position.

348 Over the 23 years, the counts of poleward- and equatorward-moving eddies show pronounced interannual  
 349 variability (Figure 4e). The annual abundance hierarchy, equatorward CEs  $>$  poleward AEs  $>$  equatorward AEs  $>$   
 350 poleward CEs, mirrors the total distribution in Figure 2b. In terms of  $EKE_T$ , CEs consistently exhibit approximately  
 351 1.5-fold greater  $EKE_T$  than AEs (Figure 4f), consistent with their longer propagation distances and higher  
 352 nonlinearity (Figure 3). While the increasing trend in CEs'  $EKE_T$  is not statistically significant overall ( $1.77 \pm 1.81$ )  
 353  $\times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p = 0.051$ ), this result is influenced by an anomalously low value in 2017, coinciding with an  $EKE$   
 354 minimum reported by Fu et al. (2023) in the central Pacific sector. Excluding this outlier yields a significant trend  
 355 of  $(2.27 \pm 1.45) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p = 0.003$ ). In contrast, AEs'  $EKE_T$  displays a robust increase by  $(2.27 \pm 0.94) \times$   
 356  $10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p < 0.001$ ). These results indicate that both eddy polarities contribute to the long-term  $EKE_T$  rise,  
 357 with CEs exhibiting greater interannual variability. As established in Figure 2, equatorward-propagating CEs and  
 358 poleward-propagating AEs dominate cross-frontal eddy abundance. Their  $EKE_T$  signals are substantially stronger  
 359 than those of the overall CFE population, with significant increasing trends of  $(2.43 \pm 2.45) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p =$   
 360  $0.045$ ) for equatorward CEs and  $(2.64 \pm 1.39) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p < 0.001$ ) for poleward AEs. Thus, beyond their

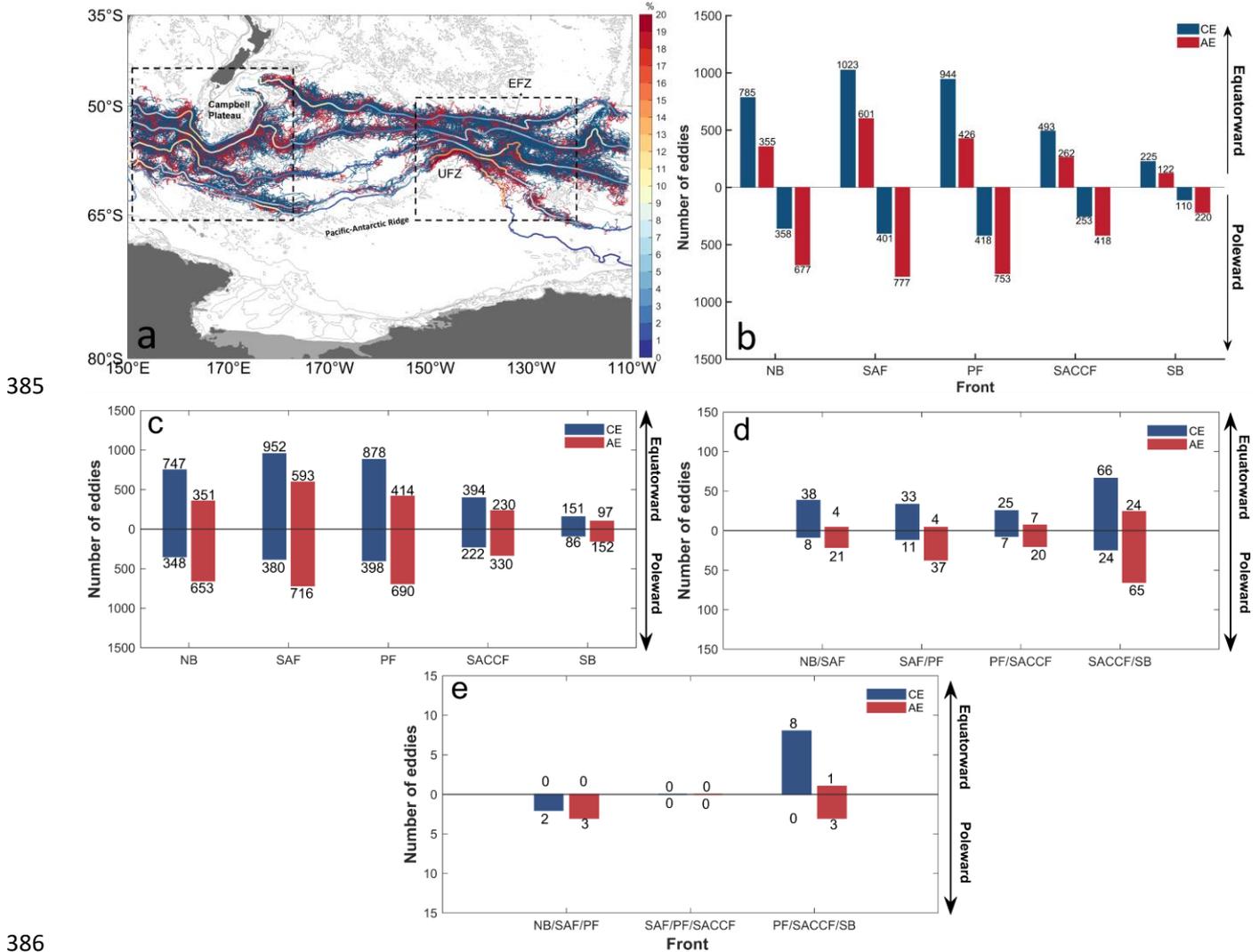
361 numerical dominance, these two subsets also govern the  $EKE$  level of CFEs and its intensification over the study  
 362 period.

363 More specifically, all CFE types exhibit pronounced interannual variability and follow the same abundance  
 364 hierarchy observed in Figure 4e in annual counts, with CEs dominating equatorward-moving eddies and AEs  
 365 prevailing among poleward-moving ones (Figure 5a–d). Among the four types, complete CFEs of both polarities  
 366 show the largest annual mean  $EKE_T$ , and display statistically significant increases over the study period with  $(4.16$   
 367  $\pm 3.54) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p = 0.024$ ) for CEs and  $(3.17 \pm 2.49) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  ( $p = 0.015$ ) for AEs, respectively  
 368 (Figure 5h). The dominant contributors to this enhancement are the same subsets that dominate abundance,  
 369 equatorward-moving CEs and poleward-moving AEs, underscoring their role as the primary  $EKE_T$  source for  
 370 complete CFEs. In contrast, partial and transient CFEs exhibit substantially lower  $EKE_T$  levels, with transient  
 371 eddies showing the weakest energy content. Most of their increasing trends are not statistically significant (Figure  
 372 5a–c), with the notable exception of transient AEs, which display a significant  $EKE_T$  increase (Figure 5c).

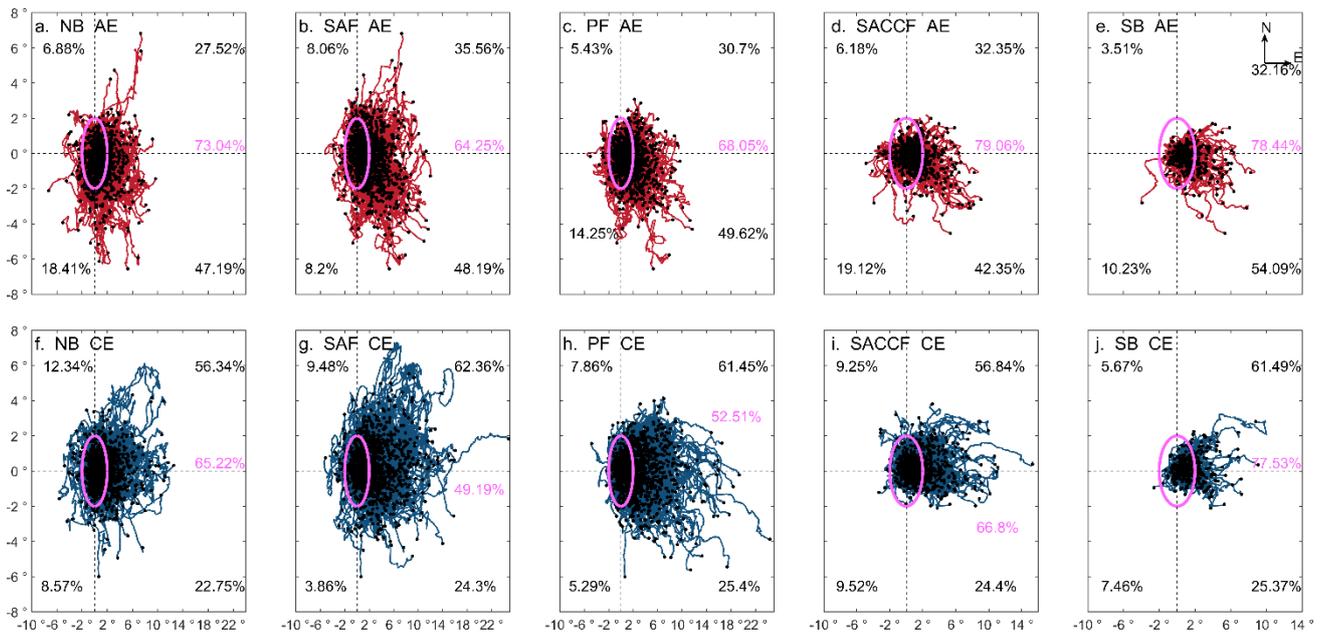
373 The relative contribution of each type to the total annual  $EKE_T$  reveals distinct energy compensation patterns  
 374 between AEs and CEs (Figure 5i, j). For AEs, the primary energy compensation occurs between Type 1  
 375 (front-generated) and Type 4 (complete frontal-crossing) eddies (correlation coefficient  $R = -0.59$ ;  $p = 0.003$ ),  
 376 indicating that enhanced activity of partially frontal-generated AEs tends to suppress complete frontal-crossing AEs,  
 377 and vice versa. A secondary compensation is between Type 2 and Type 3 ( $R = -0.46$ ;  $p = 0.027$ ). For CEs, the  
 378 dominant compensation is between Type 3 (transient) and Type 4 (complete) eddies ( $R = -0.66$ ;  $p < 0.001$ ), with a  
 379 modest compensation between Type 1 and Type 4 ( $R = -0.44$ ;  $p = 0.032$ ). These results suggest that more energetic  
 380 complete CFEs tend to coexist with reduced activity of either partially frontal-generated or transient eddies. This  
 381 compensatory relationship is critical for understanding mesoscale eddy-front interactions, particularly during the  
 382 period of elevated  $EKE_T$  in complete CFEs.

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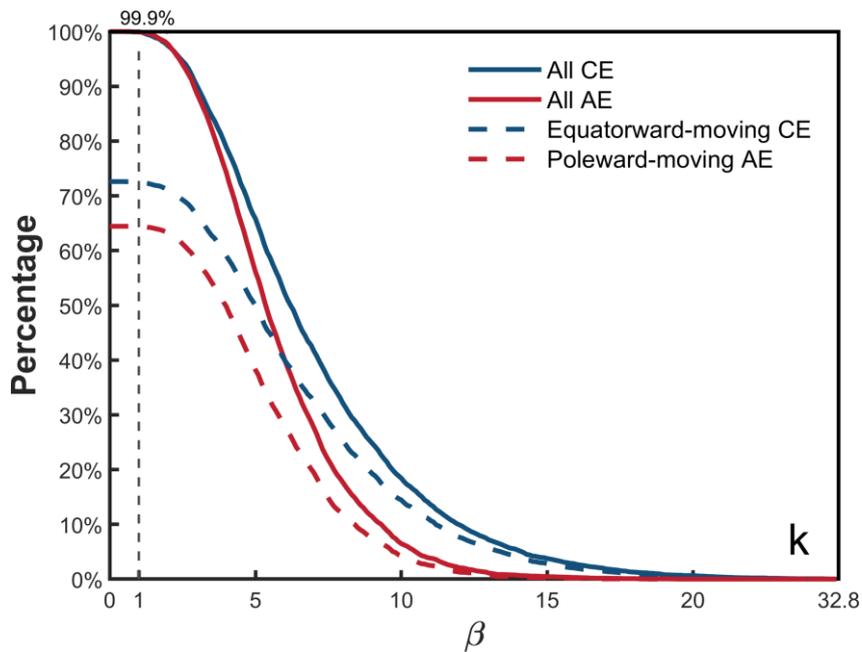


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403 **Figure 3. Relative movement trajectories of CFEs (a–j) and the percentage distribution of eddy nonlinearity  $\beta$  (k). In (a–j),**  
 404 **black percentages represent the proportion of eddies moving in different quadrant directions calculated based on the end point**  
 405 **of the trajectory, and purple percentages indicate the proportions of eddies with movement distances within a 2° range, with**  
 406 **the coordinate origin (0°, 0°) denoting the eddies' generation locations. Note that eddies crossing multiple fronts may appear**  
 407 **repeatedly at different frontal positions in this analysis.**

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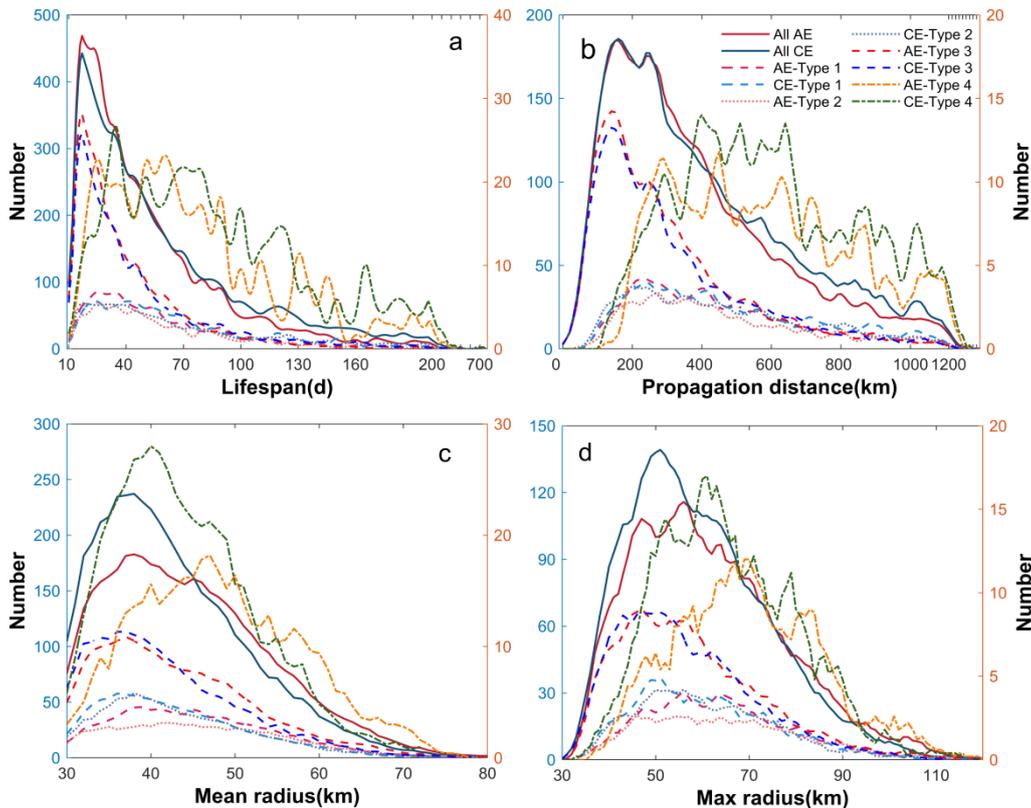
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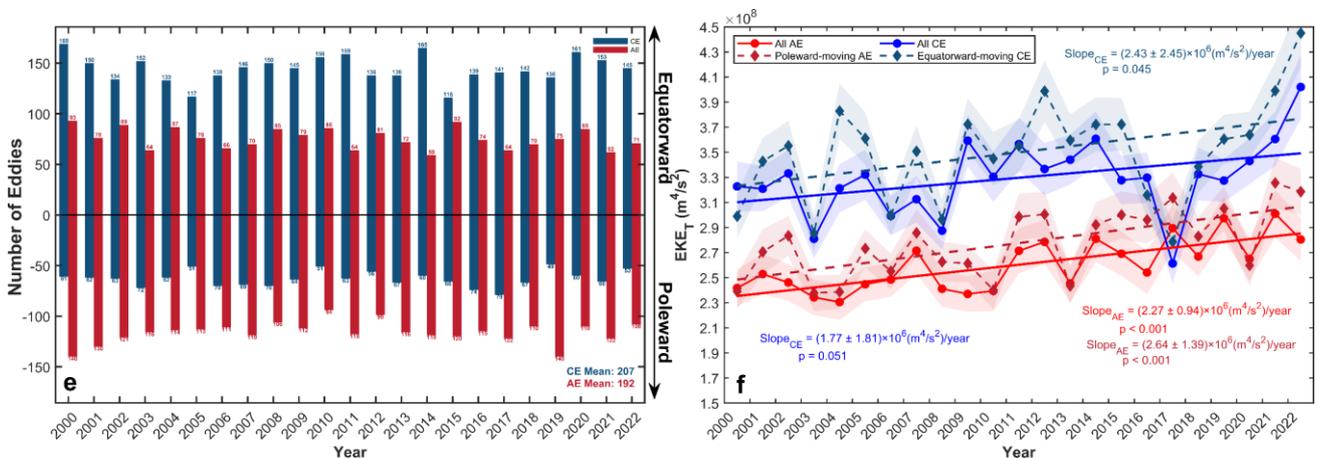
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417 **Figure 4. Statistical characteristics of different types of CFEs (a–d), time series of annual CFE counts (e), and annual mean**  
 418  **$EKE_T$  (f). Eddy counts according to (a) eddy lifespan, (b) propagation distance, (c) mean radius over lifespan, and (d)**  
 419 **maximum radius in the lifespan. In (a–d), “All” represents all CFEs, “Type 1” denotes eddies generated and**  
 420 **subsequently transported away, “Type 2” indicates eddies transported into the frontal zone and dissipated there, “Type 3”**  
 421 **represents eddies generated and dissipated within the same frontal zone, and “Type 4” shows complete CFEs experiencing**  
 422 **pre-crossing, crossing and post-crossing phases. The left Y-axis is for the first three subsets, and the right axis is for the Type 4**  
 423 **eddies to clarify their distribution. The x-axes in (a) and (b) are compressed at higher values. In (e), all-year mean counts of**  
 424 **CEs and AEs are indicated at the bottom right. In (f), the annual mean  $EKE_T$  for all AEs and CEs and the linear trends are**  
 425 **depicted by dots, red and blue solid lines, respectively. The extracted subsets of poleward-moving AEs and**  
 426 **equatorward-moving CEs are depicted by diamonds, red and blue dashed lines, respectively. Error shadings represent  $\pm 1$**   
 427 **standard deviation, and slope values are given with  $\pm 95\%$  confidence intervals.**

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**Table 1. Proportions of numbers of different eddy types relative to the total number of CFEs at each frontal zone.**

<b>Type</b>	<b>Eddy polarity</b>	<b>NB</b>	<b>SAF</b>	<b>PF</b>	<b>SACCF</b>	<b>SB</b>
<b>(1) Front-generated eddies</b>	AE	11.03%	12.21%	12.00%	8.84%	7.68%
	CE	10.11%	12.81%	12.59%	12.69%	10.49%
<b>(2) Front-dissipated eddies</b>	AE	8.41%	9.17%	9.21%	7.22%	9.45%
	CE	12.78%	11.74%	13.58%	6.52%	5.91%
<b>(3) Transient frontal eddies</b>	AE	22.71%	22.66%	20.43%	29.18%	31.17%
	CE	24.47%	18.49%	20.54%	30.37%	31.31%
<b>(4) Complete CFEs</b>	AE	5.29%	5.14%	4.76%	2.45%	2.22%
	CE	5.20%	7.78%	6.89%	2.73%	1.77%

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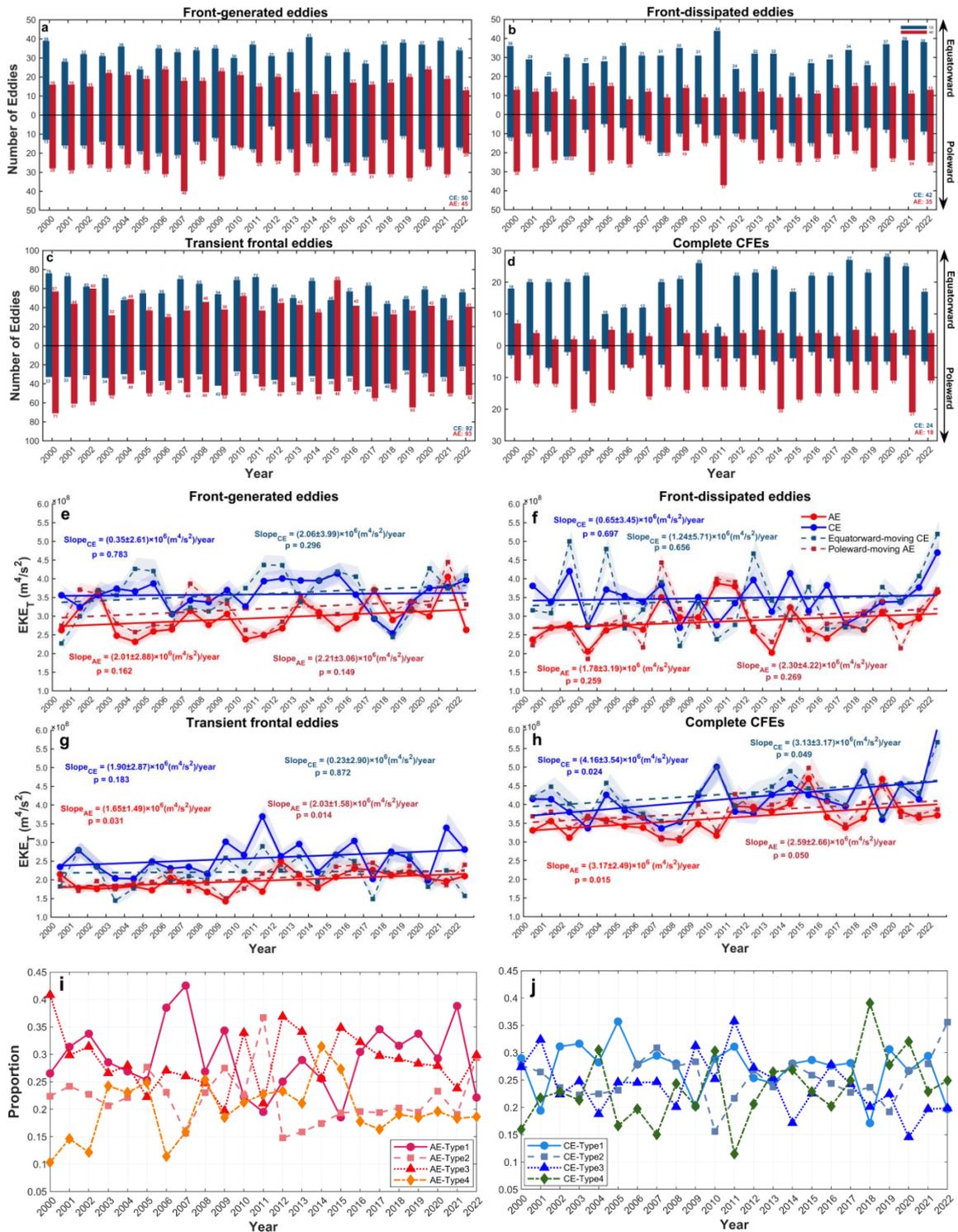
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446 **Figure 5.** Annual statistical characteristics of four types of CFES, including their counts, mean  $EKE_T$ , and proportions of the  
 447 summed  $EKE_T$  relative to all  $EKE_T$ . Time series of annual counts for (a) Front-generated eddies (Type 1), (b) Front-dissipated  
 448 eddies (Type 2), (c) Transient frontal eddies (Type 3), and (d) Complete CFES (Type 4). (e–h) Annual mean  $EKE_T$  for these  
 449 four types of eddies. (i, j) Annual summed  $EKE_T$  for each type relative to all  $EKE_T$  for total CFES. In (a–d), all-year mean  
 450 counts of CEs and AEs are indicated at the bottom right. In (e–h), the annual mean  $EKE_T$  for all AEs and CEs and the linear  
 451 trends are depicted by dots, red and blue solid lines, respectively. The extracted subsets of poleward-moving AEs and  
 452 equatorward-moving CEs are depicted by diamonds, red and blue dashed lines, respectively. Error shadings represent  $\pm 1$   
 453 standard deviation, and slope values are given with  $\pm 95\%$  confidence intervals.

### 454 3.2 $EKE_T$ evolution of complete CFEs during frontal crossing

455 Although only a small fraction of eddies complete full cross-frontal transport (Table 1; Figure 5), these  
 456 energetic features, originating in non-frontal zones and crossing entire frontal boundaries, likely dominate  
 457 long-distance heat and material exchange between inter-frontal zones. Their rising  $EKE_T$  underscores their  
 458 increasingly important dynamic role, motivating a closer examination of  $EKE_T$  evolution during frontal crossing.  
 459 The Southern Hemisphere's intrinsic vorticity asymmetry (clockwise CEs vs. counterclockwise AEs) creates  
 460 fundamental polarity differences in energy exchange when interacting with eastward frontal jets. Consequently,  
 461 eddies of opposing polarities and directions are expected to exhibit distinct patterns of  $EKE_T$  variability during the  
 462 cross-frontal transport.

463 Complete CFEs at the northern ACC fronts, NB, SAF, and PF, exhibit substantially higher  $EKE_T$  than those  
 464 at the southern fronts (SACCF, SB; Figures 6, 7), consistent with their more frequent occurrence at the northern  
 465 fronts (Table 1). For instance, mean  $EKE_T$  at the SAF ( $4.21 \times 10^8 \text{ m}^4 \text{ s}^{-2}$ ) exceeds that at the SB ( $9.51 \times 10^7 \text{ m}^4 \text{ s}^{-2}$ )  
 466 by  $3.26 \times 10^8 \text{ m}^4 \text{ s}^{-2}$ . During frontal crossing,  $EKE_T$  evolution exhibits clear polarity- and direction-dependence.  
 467 Poleward-moving AEs consistently gain kinetic energy during crossing (3.23–88.02% increase), and experience  
 468 further post-crossing amplification at the three northern fronts (43.52–71.76%), indicating sustained energy  
 469 extraction from the mean flow (Figure 6; Table 2). In contrast, poleward CEs subsequently lose energy both during  
 470 and after crossing, with reductions of 31.22–72.74% in the post-crossing phase. Equatorward-moving CFEs exhibit  
 471 opposing behaviors (Figure 7; Table 2). AEs consistently lose energy, showing reductions of 2.59–27.47% during  
 472 frontal crossing and further decline post-crossing (by 18.41–69.02%). In contrast, CEs generally gain energy, with  
 473 post-crossing increases of 48.13–76.70% at the four northern fronts. This pattern reverses at the SB, where CEs  
 474 show subsequent  $EKE_T$  loss (e.g., 43.02% decrease after crossing).

475 These results highlight fundamental asymmetries in eddy-front energy exchange governed by eddy polarity,  
 476 movement direction, and frontal latitude. They also elucidate the energy compensation observed between  
 477 frontal-generated partial/transient eddies and complete CFEs (as shown in Figure 5i, j): greater vorticity release to  
 478 complete CFEs intrinsically reduces local eddy generation. Thus, the energy extraction from frontal jets fuels the  
 479 complete frontal-crossing equatorward CEs and poleward AEs, contributing to their enhanced post-crossing  
 480 energetics, consistent with the mesoscale principle of potential vorticity conservation. The anomalous energization  
 481 of equatorward CEs at the SB, as well as lower post-crossing  $EKE_T$  relative to in-crossing values for poleward AEs  
 482 at SACCF and SB, is likely related to the weaker dynamical but stronger hydrographic characteristics of these  
 483 southernmost fronts (Park et al., 2019; Thorpe et al., 2002; Vereshchaka et al., 2021).

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486 **Table 2. Changes in mean  $EKE_T$  during different phases for complete cross-frontal eddies (CFEs) relative to pre-crossing**  
 487 **values (+: increase; -: decrease). Crossing phases represent when eddies are in the frontal zones, while post-crossing phases**  
 488 **indicate when eddies are moving away from the frontal zones. ‘-’ denotes no data.**

Direction	Eddy polarity	Phase	NB	SAF	PF	SACCF	SB
<b>poleward-moving</b>	AE	crossing	+51.80%	+61.92%	+3.23%	+88.02%	+67.03%
		post	+71.76%	+69.88%	+43.52%	+21.43%	+3.44%
	CE	crossing	-8.44%	-44.63%	-32.58%	-37.68%	-
		post	-31.22%	-72.74%	-57.63%	-69.77%	-
<b>equatorward-moving</b>	AE	crossing	-18.51%	-2.59%	-7.91%	-27.47%	-3.26%
		post	-58.00%	-46.10%	-18.41%	-69.02%	-27.91%
	CE	crossing	+64.28%	+17.91%	+47.86%	+54.70%	-10.89%
		post	+75.78%	+52.33%	+48.13%	+76.70%	-43.02%

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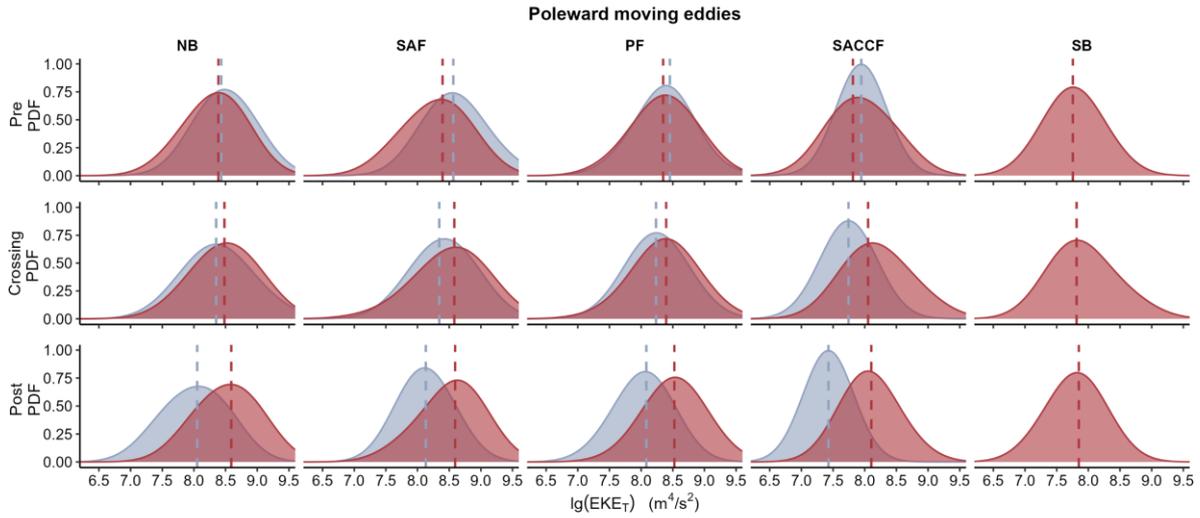
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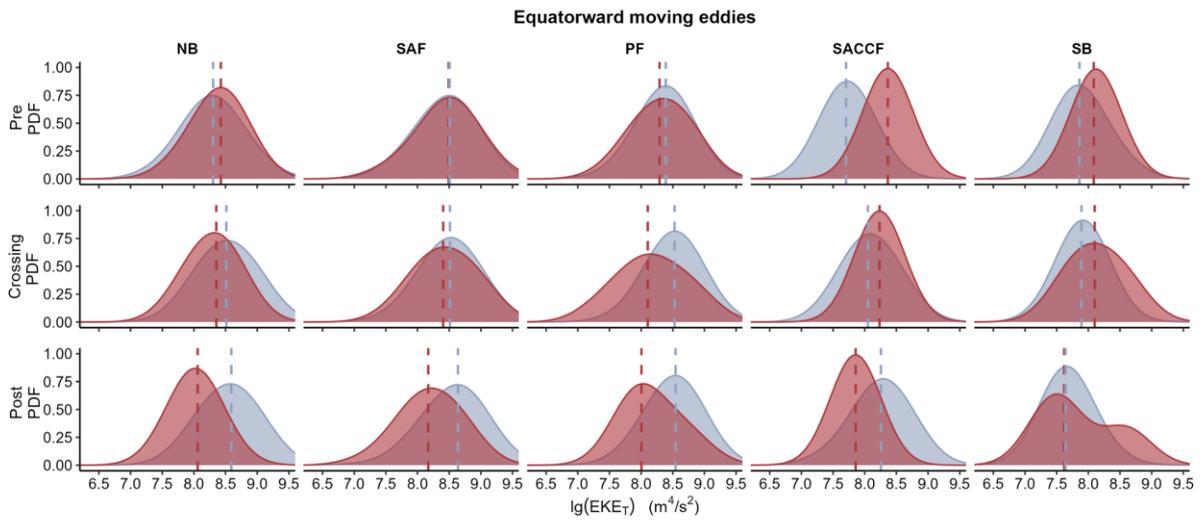
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500 **Figure 6. Probability density function (PDF) of  $EKE_T$  for poleward-moving CFEs in pre-crossing, crossing, and post-crossing**  
 501 **phases. Dashed lines indicate median  $EKE_T$  values. Blue and red colors represent CEs and AEs, respectively.**



502

503 **Figure 7. Probability density function (PDF) of  $EKE_T$  for equatorward-moving CFEs in pre-crossing, crossing, and**  
 504 **post-crossing phases. Dashed lines indicate median  $EKE_T$  values. Blue and red colors represent CEs and AEs, respectively.**

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### 513 3.3 Thermohaline transport effects of CFEs

514 Argo  $\theta$ - $S$  profiles (Figure 8) reveal that cyclonic eddies (CEs) and anticyclonic eddies (AEs) exhibit distinct,  
 515 polarity-dependent hydrographic signatures within the same interfrontal zones. CEs consistently contain colder,  
 516 fresher water with shallower isopycnals (upper 1000 dbar), whereas AEs are characterized by warmer, saltier water  
 517 and deeper isopycnals. These contrasts underscore the role of nonlinear eddies in mediating cross-frontal exchange.  
 518 However, a subset of AEs in the SACCF-PF zone trap anomalously cold, fresh polar waters ( $\theta_{min} = -1.76$  °C and  $S$   
 519  $< 34.0$  psu), similar to those observed in some AEs within the SB-SACCF zone, indicating that AEs can also  
 520 transport polar waters equatorward. Notably, eddy-induced vertical motions, upwelling in CEs and downwelling in  
 521 AEs, produce vertical displacements but do not alter  $\theta$ - $S$  properties of source water columns (Falkowski et al., 1991;  
 522 Li et al., 2022). Therefore, this mechanism can only account for overlapping  $\theta$ - $S$  signatures that arise from vertical  
 523 repositioning of the same water mass, rather than true cross-frontal modification.

524 Analysis of radius-normalized  $\theta$ - $S$  distributions reveals distinct water mass signatures in CEs and AEs across  
 525 northern interfrontal zones (SAF-NB, PF-SAF, and SACCF-PF). Core water masses, especially Subantarctic Mode  
 526 Water (SAMW) and Antarctic Intermediate Water (AAIW), are not circumpolarly uniform but exhibit substantial  
 527 regional variability (Bostock et al., 2013; Li et al., 2022). For instance, within the Pacific sector, the salinity  
 528 minimum of AAIW ranges from  $\sim 34.2$  in the southeast Pacific formation region to greater than 34.5 in the Tasman  
 529 Sea after mixing (Bostock et al., 2013). Similarly, SAMW exhibits distinct spatial patterns in its formation and  
 530 properties (Li et al., 2021). Accordingly, the ranges in Table 3 are intended as a practical guide for identifying  
 531 water masses within the specific Pacific sectoral context of this study.

532 Between the SAF and NB, well-defined layers of SAMW, AAIW, and Upper Circumpolar Deep Water  
 533 (UCDW) are observed from the upper to lower layer in the AE (Figure 9d, j), confirming their local origin within  
 534 the Antarctic Convergence Zone. Conversely, the CE in the same zone shows markedly different  $\theta$ - $S$  structure  
 535 (Figure 9a, g), with upper layers ( $< 1000$  dbar) lacking SAMW/AAIW signatures and instead containing colder,  
 536 fresher waters of southern origin. Neutral density ( $\gamma^n$ ) surfaces in the CE are approximately 200 dbar shallower than  
 537 in the AE, demonstrating that CEs effectively transport high-potential-energy southern waters into the SAF-NB  
 538 zone. This establishes strong mesoscale potential energy contrasts between the low-potential-energy waters in AEs  
 539 and the high-potential-energy waters in CEs, an energetic precondition for baroclinic instability via the release of  
 540 available potential energy (Fu et al., 2023).

541 In the PF-SAF region, both the CE and AE maintain thermohaline contrasts similar to those in the SAF-NB  
 542 zone but with reduced magnitude, preserving the characteristic warmer/saltier AE and colder/fresher CE signatures

543 (the middle panels of Figure 9). Notably, only the CE's upper layer exhibits distinct Winter Water (WW)  
544 characteristics, confirming their southern origins. Below this, the normalized CE sequentially displays UCDW and  
545 LCDW, while the AE shows only UCDW beneath the relatively warm and salty Antarctic Surface Water within the  
546 upper 2000 dbar. The vertical isopycnal structure reveals depth-dependent displacements: in the near-surface layer,  
547 the CE's isopycnal  $\gamma^{\sigma}=27.1 \text{ kg/m}^3$  is  $\sim 100$  dbar shallower than the AE, while at intermediate depths, the CE's  
548 isopycnal  $\gamma^{\sigma}=27.6 \text{ kg/m}^3$  (400–600 dbar) is  $\sim 500$  dbar shallower than the AE ( $\sim 1000$  dbar).

549 The thermohaline anomalies between CE and AE still exist in the SACCF-PF zone (the right panels of Figure  
550 9). In the CE, a subsurface WW layer overlies a warm UCDW core, with LCDW dominating below 1000 dbar,  
551 showing a characteristic of waters south of the SACCF (Aoki et al., 2013; Auger et al., 2021). While the AE also  
552 contains these water masses, they show weaker WW expression, a more pronounced  $\theta_{\max}$  core, and vertically  
553 extended UCDW, reflecting their relatively northern origins. Isopycnals in the CE remain consistently 300–400  
554 dbar shallower than in the AE throughout the water column.

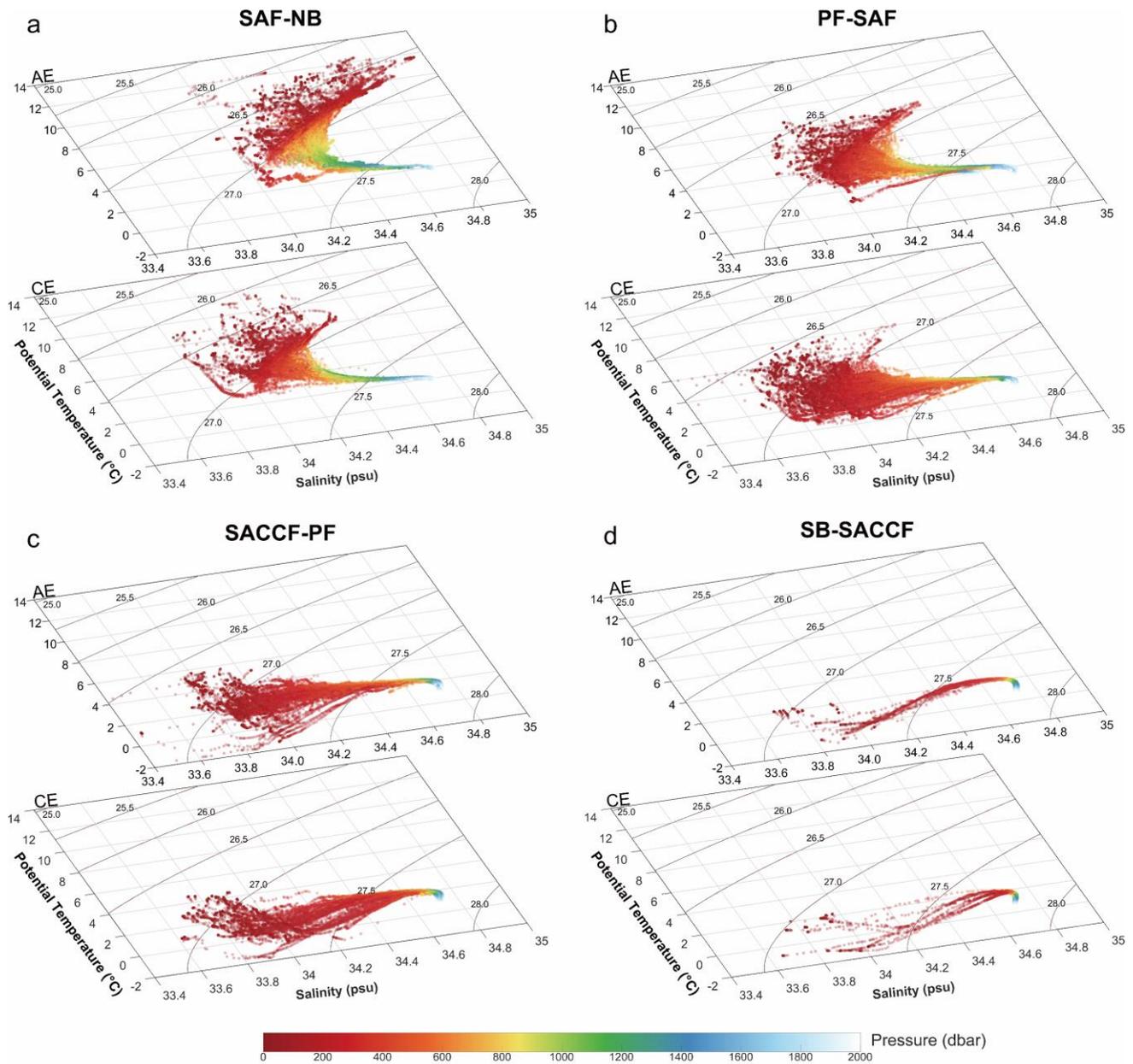
555 Therefore, the normalized AE and CE possess distinct water mass distributions within the same inter-frontal  
556 zones, marked by profound isopycnal thermohaline differences. AEs and CEs transport their respective source  
557 water masses into the same zones, amplifying mesoscale hydrographic variability. The above comparative analysis  
558 demonstrates that cross-frontal CEs play a dominant role in meridional water mass transport, particularly in the  
559 SAF-NB and SACCF-PF zones, consistent with their greater dynamical vigor. This cross-frontal exchange reduces  
560 baroclinicity between interfrontal zones while enhancing mesoscale available potential energy within individual  
561 zones.

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Figure 8. Potential temperature-salinity ( $\theta$ -S) diagrams in the eddy interiors observed in different inter-frontal zones. (a) SAF-NB; (b) PF-SAF; (c) SACCF-PF; (d) SB-SACCF.

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577 **Table 3. Criteria for the division of water masses according to potential temperature ( $\theta$ , °C), salinity ( $S$ , psu) and neutral**  
 578 **density ( $\gamma^n$ , kg/m<sup>3</sup>). Note: The ranges listed, particularly for SAMW and AAIW, are primarily representative of the Pacific**  
 579 **sector of the Southern Ocean.**

Water mass	$\theta$ (°C)	$S$ (psu)	$\gamma^n$ (kg/m <sup>3</sup> )	Reference
<b>SAMW</b>		34.35–34.60	27.00–27.20	Carter et al., 2022; Bostock et al., 2013; Herraiz-Borreguero and Rintoul, 2011
<b>AAIW</b>		$S_{\min}$ 34.28–34.40	27.10–27.60	Bostock et al., 2013; Xia et al., 2022; Valla et al., 2018
<b>UCDW</b>	$\theta_{\max}$		27.55–28.00	Naveira Garabato et al., 2002
<b>LCDW</b>		$S_{\max}$	28.00–28.27	
<b>WW</b>	$\theta_{\min}$ < 2.00			

580 \*SAMW, Subantarctic Mode Water; AAIW, Antarctic Intermediate Water; UCDW, Upper Circumpolar Deep Water; LCDW, Lower  
 581 Circumpolar Deep Water; WW, Winter Water.

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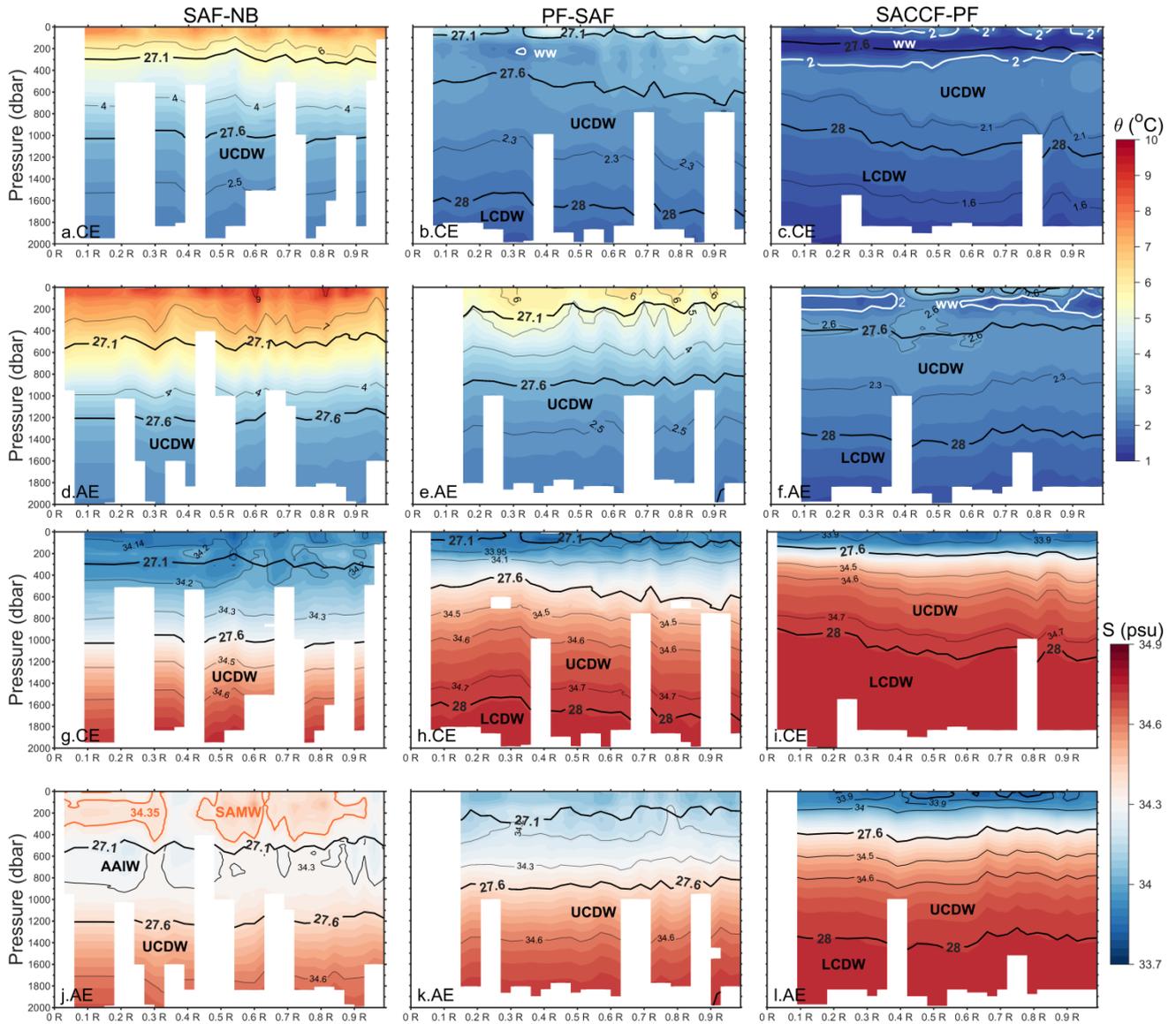
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595 **Figure 9. Sectional distributions of  $\theta$  (a–f) and  $S$  (g–l) along normalized eddy radius ( $R$ ) direction in the interfrontal zones of**  
596 **SAF-NB (the left panels), PF-SAF (the middle panels), and SACCF-PF (the right panels). Black thick contours indicate neutral**  
597 **density ( $\gamma^n$ ,  $\text{kg}/\text{m}^3$ ), thin contours represent  $\theta$  or  $S$  isolines, respectively.**  
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## 611 4 Discussion

612 This study reveals a fundamental polarity- and direction-dependent asymmetry in cross-frontal eddy (CFE)  
 613 dynamics within the Pacific sector of the SO. This asymmetry manifests in three key aspects: (1) a distinct  
 614 abundance hierarchy among CFE types, (2) contrasting *EKE* intensities and trends, and (3) polarity- and  
 615 direction-selective energy transfers during eddy-frontal jet interactions. The observed hierarchy, in which CEs  
 616 predominantly migrate equatorward and AEs poleward, aligns with established eddy dynamics in the SO (He et al.,  
 617 2023; Li et al., 2022; Patel et al., 2019). Beyond abundance, the dominant types (equatorward CEs and poleward  
 618 AEs) exhibit superior energetic characteristics, including higher *EKE* levels (Figure 4f), longer propagation  
 619 distances (Figure 3), and stronger nonlinearity compared to their counterparts. Our results demonstrate that  
 620 substantial energy gain during frontal crossing sustains the enhanced energetics of these two dominant complete  
 621 CFE types, as illustrated in Figure 11. Moreover, the energy compensation relationships (Figure 5i, j) suggest that  
 622 partial and transient CFEs of the same dominant polarity-direction combinations likely follow a similar  
 623 energization mechanism. Taken together, these findings indicate that polarity- and direction-dependent eddy-front  
 624 interactions fundamentally govern CFE energetics and, consequently, their capacity to drive meridional heat and  
 625 material transport across Southern Ocean frontal zones.

626 Although CFE abundance shows no significant trend over 2000–2022, both polarity groups experienced  
 627 substantial  $EKE_T$  intensification, with CEs gaining energy at  $(2.27 \pm 1.45) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$  (excluding the  
 628 anomalously low 2017 value) and AEs at  $(2.27 \pm 0.94) \times 10^6 \text{ m}^4 \text{ s}^{-2} \text{ yr}^{-1}$ . The increasing trends are also robust in  
 629 terms of annual area-weighted mean *EKE* ( $EKE = EKE_T/S$ , where  $S$  is the total area of an eddy; Figure S1 in the  
 630 Supplementary Materials), with trends of  $3.71 \pm 2.08 \text{ cm}^2 \text{ s}^{-2} \text{ yr}^{-1}$  for CEs and  $2.58 \pm 0.97 \text{ cm}^2 \text{ s}^{-2} \text{ yr}^{-1}$  for AEs. The  
 631 *EKE* enhancements are even more pronounced for the dominant subsets, equatorward CEs and poleward AEs  
 632 (Figures 4, S1). These rates substantially exceed previously reported *EKE* trends. Hogg et al. (2015) estimated a  
 633 regional mean *EKE* increase of  $14.9 \pm 4.1 \text{ cm}^2 \text{ s}^{-2} \text{ decade}^{-1}$  under intensifying westerlies in the Pacific sector  
 634 (1990–2015). Similarly, Zhang et al. (2021) documented an *EKE* increase of  $< 20 \text{ cm}^2 \text{ s}^{-2} \text{ decade}^{-1}$  south of New  
 635 Zealand and downstream of the Campbell Plateau, identified as the only region with significant *EKE* rise in the SO  
 636 during 1993–2018. The *EKE* trends for CFEs presented here are considerably larger than these basin-scale  
 637 estimates. In contrast, non-CFEs exhibit lower *EKE* levels and no comparable  $EKE_T$  increase (Figure 10). The  
 638 significant area-weighted mean *EKE* trend for non-frontal crossing CEs is also weaker (Figure S2). These contrasts  
 639 suggest that the overall *EKE* increase in the Pacific sector is primarily attributable to CFEs. Non-CFEs contribute  
 640 little to, or may even obscure, the observed *EKE* trends. This finding implies that eddy-front interactions, rather

641 than basin-scale wind forcing alone, may be the primary driver of recent *EKE* trends in the Pacific sector of the SO.

642 Our results suggest that the enhanced wind stress (Hogg et al., 2015; Menna et al., 2020) could preferentially  
643 energize cross-frontal activity, especially the CFEs achieving complete frontal crossing (Figure 5). The  
644 predominance of equatorward CEs aligns with intensified Ekman transport patterns reported by Shi et al. (2025),  
645 also suggesting wind-driven facilitation of meridional eddy migration. Building on Fu et al.'s (2023) framework of  
646 wind-driven energy pathways (baroclinic: mean kinetic energy→mean available potential energy→mesoscale  
647 available potential energy→EKE; barotropic: mean kinetic energy→EKE), we found that only equatorward CEs  
648 and poleward AEs gain kinetic energy from frontal jets (Table 2; Figures 6, 7). This energization likely arises from  
649 two synergistic mechanisms: (1) barotropic instability from enhanced horizontal shear when eddy rotation aligns  
650 with the eastward jet (Qiu et al., 2024), and (2) baroclinic instability triggered by potential energy release for  
651 enhanced hydrographic gradients with ambient waters (Fu et al., 2023). Conversely, significant energy losses of  
652 poleward CEs and equatorward AEs during frontal crossing, possibly due to counter-rotational turbulent dissipation  
653 (Dong et al., 2017; Jan et al., 2017) and upwelling (downwelling)-induced baroclinicity reduction with the ambient  
654 waters.

655 Intensifying and poleward-shifting westerlies have emerged as a dominant dynamic forcing mechanism in the  
656 SO (Behrens and Bostock, 2023; Hogg et al., 2015). Meanwhile, buoyancy forcing from meridionally  
657 inhomogeneous ocean warming has been shown to accelerate the ACC at 48°S–58°S (Shi et al., 2021). Our  
658 results demonstrate that CFEs play a vital role in mediating the oceanic response to these forcings by compensating  
659 for heat transport (Figure 9). By facilitating equatorward cold-water transport via CEs and poleward warm-water  
660 transport via AEs, CFEs reduce cross-frontal water mass property gradients. This process effectively buffers wind-  
661 or warming-induced increases in baroclinicity, thereby maintaining the SO's thermal equilibrium and modulating  
662 the ACC's response to external forcing. These findings highlight the dual role of CFEs as both energy transporters  
663 and dynamical stabilizers in a changing climate.

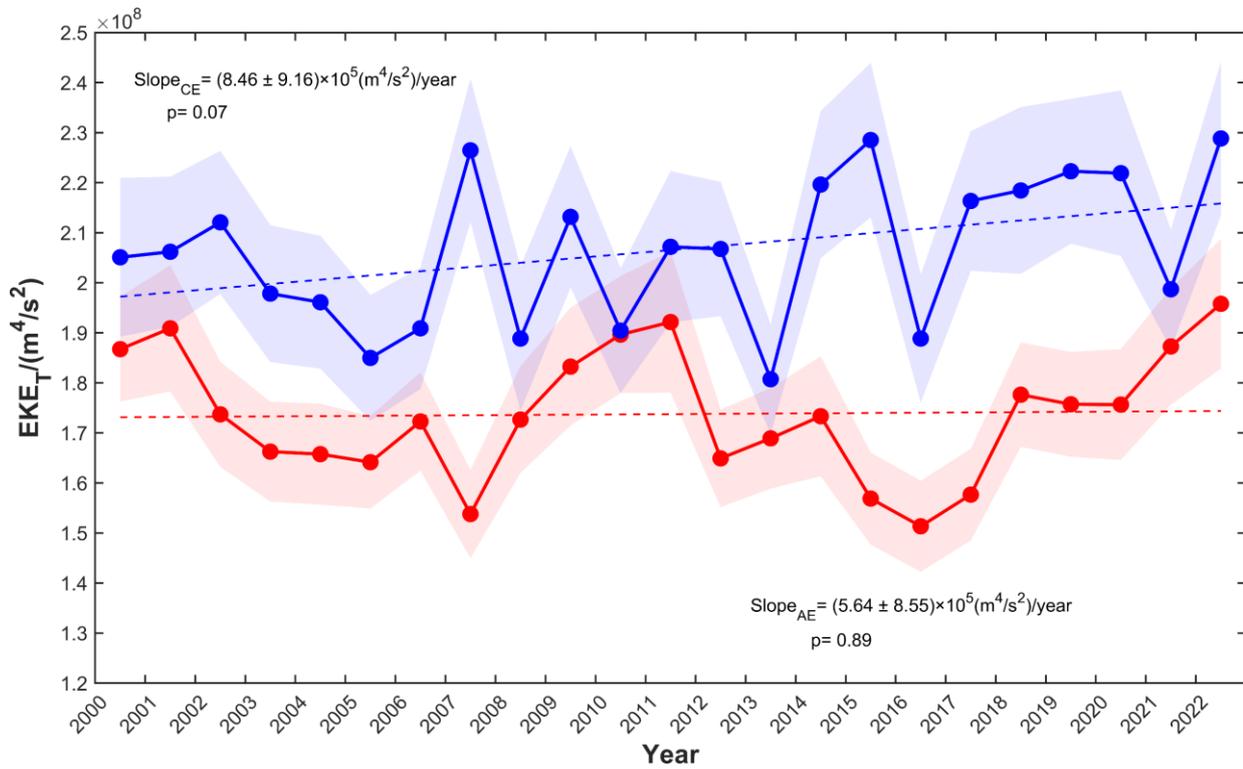
664 This study has several limitations that should be noted. First, the analysis does not account for potential  
665 interannual or seasonal variations in frontal positions. However, a comprehensive sensitivity analysis, in which the  
666 frontal zone half-width was expanded to  $\pm(25+R)$  km, confirms that all key conclusions remain qualitatively and  
667 statistically robust (see Supplementary Materials, Tables S1 and S2, and Figures S3–S6). Second, the hydrographic  
668 analysis utilized all qualified Argo profiles from 2000 to 2022 without segregating interannual or seasonal  
669 variability, which may introduce some uncertainty into the normalized water mass properties. Finally, while this  
670 study focuses on eddy characteristics, it does not evaluate the dynamic properties of the frontal jets themselves. A

671 detailed analysis of jet variability and energy transfer is essential for a more comprehensive mechanistic  
 672 understanding of eddy-jet interactions and represents an important direction for future research.

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677 **Figure 10.** Time series of annual mean  $EKE_T$  for eddies in the interfrontal zones.  $EKE_T$  is shown by blue solid line for CEs and  
 678 red solid line for AEs, with linear regression indicated by dashed lines, error shadings representing one standard deviation,  
 679 and slope values given with  $\pm 95\%$  confidence intervals.

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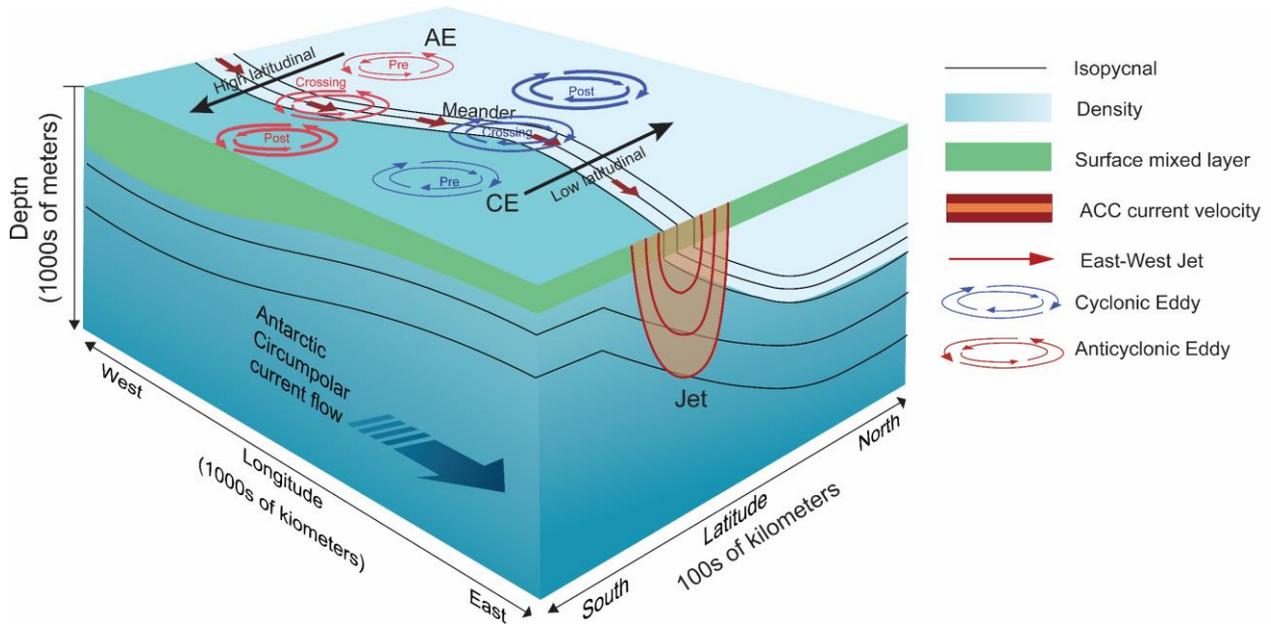
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693 **Figure 11. Illustrations of EKE variations during frontal crossing for poleward AEs and equatorward CEs (modified from**694 **Figure 1 in Chapman et al., 2020). The thickness of rotational velocity vectors represents relative flow intensity.**

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## 714 **5 Conclusions**

715 This study provides a comprehensive characterization of cross-frontal eddies (CFEs) in the Pacific sector of  
716 the Southern Ocean, revealing fundamental asymmetries that govern their role in meridional transport and energy  
717 exchange. Through the integration of multi-decadal satellite altimetry and in-situ hydrographic data, we have  
718 demonstrated that CFE activity is not random but follows a well-defined polarity- and  
719 direction-dependent hierarchy. The predominance of equatorward-moving cyclonic eddies  
720 (CEs) and poleward-moving anticyclonic eddies (AEs) reflects an intrinsic organization in cross-frontal exchange  
721 pathways that has profound implications for Southern Ocean dynamics.

722 Three key findings emerge from our analysis. First, the observed CFE abundance hierarchy is complemented  
723 by significant differences in energetic characteristics, with the dominant types (equatorward CEs and poleward AEs)  
724 exhibiting superior kinetic energy levels, longer propagation distances, and stronger nonlinearity. Second, these  
725 eddies experience sustained energization through polarity-selective energy transfers during frontal crossings,  
726 gaining kinetic energy from the eastward frontal jets while their counterparts experience energy dissipation. The  
727 intensification of CFE activity has occurred over the past two decades, with *EKE* trends substantially exceeding  
728 previous basin-scale estimates, suggesting that eddy-front interactions, rather than wind forcing alone, drive recent  
729 energetic changes in the region. Third, hydrographic analyses confirm that CFEs function as effective transporters  
730 of distinct water masses, with CEs and AEs maintaining sharp thermohaline contrasts within the same interfrontal  
731 zones. This cross-frontal exchange reduces large-scale baroclinicity while enhancing mesoscale available potential  
732 energy, creating a dynamic balance that regulates meridional heat transport.

733 As climate change intensifies westerly winds and modifies buoyancy forcing, CFEs are likely to play an  
734 increasingly important role in modulating the ACC's response to external forcing. Their ability to maintain thermal  
735 equilibrium across frontal zones highlights their significance for understanding future changes in Southern Ocean  
736 circulation, carbon uptake, and global climate feedbacks. Future research should focus on quantifying the precise  
737 contribution of CFEs to meridional heat and carbon fluxes and investigating how their stabilizing role might evolve  
738 under continued climate change.

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**744 Data availability**

745 The satellite altimeter data are available online at the Copernicus Marine Service (the CMEMS all-satellite L4  
746 SLA product at <https://doi.org/10.48670/moi-00148>). The data were accessed and downloaded on 18<sup>th</sup> December,  
747 2025. The frontal data used in this study were sourced from Park et al. (2019) at <https://doi.org/10.17882/59800>.  
748 Argo profiles are available at the website of <https://argo.ucsd.edu/data/>.

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**750 Author contributions**

751 Huimin Wang: Methodology, Software, Formal analysis, Investigation, Data curation, Writing-original draft,  
752 Visualization. Lingqiao Cheng: Conceptualization, Methodology, Resources, Writing-original draft,  
753 Writing-review & editing, Supervision, Project administration, Funding acquisition. Erik Behrens: Validation,  
754 Writing-review & editing. Zhuang Chen: Validation, Writing-review & editing. Jennifer Devine: Writing-review &  
755 editing. Guoping Zhu: Resources, Writing-review & editing, Funding acquisition.

756

**757 Competing Interests**

758 The authors declare that they have no conflict of interest.

759

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765

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