



1 **Wave effect mechanisms enhancing sea–air CO<sub>2</sub> exchange**  
2 **and modulating seawater carbonate–pH adaptation in the**  
3 **POP2–waves coupled model**

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8 **Abstract**

9 Wave and bubble mechanisms have demonstrated their impact on sea–air CO<sub>2</sub> flux  
10 by enhancing gas transfer velocity ( $K_w$ ) through significant wave height ( $H_s$ ).  
11 Neglecting wave and bubble processes may lead to an underestimation of CO<sub>2</sub> flux  
12 under high 10-m wind speeds ( $U_{10}$ ) in most state-of-the-art climate models. In this study,  
13 a waves module from the Princeton Ocean Model (POM) has been incorporated into  
14 the Parallel Ocean Program version 2 (POP2), referred to as POP2–waves, in the  
15 Community Earth System Model version 1.2.2 (CESM1.2.2) framework. The POP2–  
16 waves and a control run of CESM1.2.2 (B–CTL) CO<sub>2</sub> flux simulations are compared  
17 with the National Oceanic and Atmospheric Administration’s (NOAA) CarbonTracker,  
18 version 2022 (CT2022) data. Overall, bubbles contribute up to 41.3% to the total sea–  
19 air CO<sub>2</sub> flux, consistent with recent studies, and POP2–waves exhibits a stronger CO<sub>2</sub>  
20 flux than B–CTL under high  $U_{10}$ . Likewise, the spatial distribution of POP2–waves CO<sub>2</sub>  
21 flux is broadly agrees with that of NOAA CT2022, although some discrepancies remain.  
22 Under the sea–air partial pressure differences ( $dp\text{CO}_2$ ) negative feedback associated  
23 with the interaction between CO<sub>2</sub> fluxes and the carbonate–pH system, POP2–waves  
24 show increases of 11.8%, 41.6%, and 1.8% in the CO<sub>2</sub> sink, source, and global average,  
25 respectively, compared to the B–CTL. The  $dp\text{CO}_2$  (pH) exhibits the strongest positive  
26 (negative) regression coefficient with CO<sub>2</sub> flux across the global ocean. Additionally,  
27  $K_w$  shows a positive (negative) regression coefficient with CO<sub>2</sub> flux in source (sink)  
28 regions, while SST displays the opposite pattern relative to  $K_w$ .



## 29 1. Introduction

30 The exchange of CO<sub>2</sub> between the air and sea is a crucial component of the global  
 31 carbon cycle, carrying significant implications for Earth's climate (Bange et al., 2024;  
 32 Friedlingstein et al., 2022; McKinley et al., 2020; Müller et al., 2023; Shutler et al.,  
 33 2019). Traditionally, the exchange of CO<sub>2</sub> across the sea–air interface can be  
 34 characterized using the following bulk formula (e.g., Wanninkhof, 1992, 2014;  
 35 McKinley et al., 2020; Dong et al., 2022; Fay et al., 2021; Zhou et al., 2023; Heimdal  
 36 et al., 2024):

$$37 FCO_2 = K_w \cdot k_0 \cdot (pCO_2^w - pCO_2^a) \cdot (1 - ice) \quad (1)$$

38 where the sea–air CO<sub>2</sub> flux ( $FCO_2$ ) is determined by the gas transfer velocity ( $K_w$ ), the  
 39 difference in partial pressure of CO<sub>2</sub> ( $pCO_2$ ) between the seawater and atmosphere  
 40 (indicated by superscripts "w" and "a," respectively), and the solubility constant ( $k_0$ ),  
 41 which varies with salinity and temperature (Weiss, 1974). Moreover, the fluxes are  
 42 weighted by 1 minus the ice fraction ( $ice$ ). The estimation of bulk CO<sub>2</sub> fluxes in  $K_w$  is  
 43 typically similar to the Wanninkhof (1992) equation:

$$44 K_w = 0.251 \langle U_{10}^2 \rangle (Sc/660)^{-0.5} \quad (2)$$

45 where  $\langle U_{10}^2 \rangle$  represents the mean squared wind speed at 10 meters, and  $Sc$  denotes the  
 46 Schmidt number. The flux estimate error can reach as high as 34% due to model  
 47 limitations and uncertainties surrounding bulk parameters. (Signorini and McClain,  
 48 2009). However, significant uncertainties persist in estimates of sea–air CO<sub>2</sub> flux,  
 49 primarily stemming from an incomplete understanding of the spatiotemporal variability  
 50 in the governing mechanisms (Shutler et al., 2019).

51 Some studies have highlighted that estimating CO<sub>2</sub> fluxes involves considerations  
 52 beyond just wind speed and solubility factors. Monahan and Spillane (1984) previously  
 53 regarded whitecaps as "low impedance vents", effectively "shortening" the water-side  
 54 transfer resistance. Without wave breaking, transport between the ocean and the



55 atmosphere occurs through slow conduction and molecular diffusion, while breaking  
 56 acts as a transitional process from laminar to turbulent flow (Deike, 2022). Bubbles  
 57 offer an extra surface for gas transfer, rising through the water-side mass boundary layer  
 58 of the water surface and bursting at the water surface, enhancing near-surface  
 59 turbulence (Soloviev and Lukas, 2010; Bell et al., 2017; Deike and Melville, 2018;  
 60 Krall et al., 2019; Czerski et al., 2022). Gutiérrez-Loza et al. (2022) pointed out that sea  
 61 spray transported into the atmosphere and ultimately evaporating serves an efficient  
 62 mechanism for sea–air CO<sub>2</sub> exchange during high and intermediate wind speeds (above  
 63 6–8 m s<sup>−1</sup>), whereas during low wind speeds (<6 m s<sup>−1</sup>), water-side convection was  
 64 identified as a significant control mechanism. Several studies corroborate the  
 65 significance of breaking waves and bubble injection mechanisms in facilitating CO<sub>2</sub>  
 66 transfer (e.g., Andreas et al., 2016; Brumer et al., 2017; Blomquist et al., 2017; Reichl  
 67 and Deike, 2020; Li et al., 2023; Zhou et al., 2023). Deike and Melville (2018)  
 68 demonstrated that bubble contribution to CO<sub>2</sub> transfer accounts for more than 40%  
 69 when the 10-meter wind speed ( $U_{10}$ ) is  $\geq 10$  m s<sup>−1</sup>, becomes dominant at  $U_{10} \geq 15$  m s<sup>−1</sup>,  
 70 and reaches 60% at 20 m s<sup>−1</sup>.

71 Monitoring the sea surface CO<sub>2</sub> content is crucial for comprehending the Earth  
 72 system, as climate change has begun to impact the ocean's carbon uptake capacity  
 73 (Behncke et al., 2024). However, indirect estimates of  $p\text{CO}_2$  derived from pH and  
 74 salinity measurements still entail uncertainties (e.g., Williams et al., 2017; Gray et al.,  
 75 2018; Coggins et al., 2023; Behncke et al., 2024; Fay et al., 2024; Yang et al., 2024).  
 76 Direct observation is the most desirable tool for understanding the real world, but the  
 77 scarcity of  $p\text{CO}_2$  and sea–air CO<sub>2</sub> flux measurements results in relatively short time  
 78 series. The direct measurements involved discrete samples with an equilibrator system  
 79 and the use of closed-path CO<sub>2</sub> gas analyzers to measure  $p\text{CO}_2$  (Sutton et al., 2014,  
 80 2021; Bakker et al., 2016; Sabine et al., 2020; Akhand et al., 2021; Wu and Qi, 2023).



81 In addition to closed-path CO<sub>2</sub> measurements, direct measurements of the open-path  
 82 CO<sub>2</sub> flux using the eddy covariance technique were conducted over the open ocean,  
 83 revealing CO<sub>2</sub> fluxes were within a factor of 40 of those estimated by Eq. (1) (Edson et  
 84 al., 2011; Tokoro et al., 2014; Bell et al., 2017; Dong et al., 2021; Van Dam et al., 2021).  
 85 Local CO<sub>2</sub> flux observations present challenges for validating global Earth system  
 86 models (ESMs) due to discrepancies in time and space scales, implementing a data  
 87 assimilation system is crucial for effective model evaluation.

88 Several recent studies have estimated the sea–air CO<sub>2</sub> flux using Eq. (1), while  
 89 current parameterizations in oceanic and atmospheric models still rely exclusively on  
 90 wind speed (Long et al., 2013; Moore et al., 2013; Couldrey et al., 2016; Jin et al., 2017;  
 91 Lovenduski et al., 2019; Ziehn et al., 2020; Chikamoto and DiNezio, 2021) and through  
 92 data analysis of non-model simulations (Wanninkhof et al., 1992, 2014; Fay et al., 2021;  
 93 Zhou et al., 2023). The uptake of anthropogenic CO<sub>2</sub> by the oceans is lowering pH and  
 94 altering carbonate chemistry in nonlinear ways, further reducing the oceans' capacity to  
 95 absorb additional CO<sub>2</sub> (Moore et al., 2013). The mechanisms of breaking waves and  
 96 bubble injection in facilitating  $K_w$  have been examined primarily through data analysis  
 97 of non-model simulations (e.g., Andreas et al., 2016; Blomquist et al., 2017; Reichl and  
 98 Deike, 2020; Zhou et al., 2023). However, there is relatively limited discussion on the  
 99 impact of wave effect mechanisms in estimating the sea–air CO<sub>2</sub> flux and seawater  
 100 acidification within coupling the ocean biogeochemistry with wave models, whereas  
 101 local processes are more tightly controlled by the sea state.

102 In this study, we aim to investigate the significance of breaking waves and bubble  
 103 injection mechanisms through ocean geochemistry response on the sea–air CO<sub>2</sub> flux  
 104 and ocean's buffering capacity through the carbonate–pH system, using the Parallel  
 105 Ocean Program version 2 (POP2; Smith et al., 2010) of the Community Earth System  
 106 Model version 1.2.2 (CESM1.2.2; Hurrell et al., 2013) coupled with waves module



107 (Mellor et al., 2008) of the Princeton Ocean Model (POM; Blumberg and Mellor, 1987).

108 The coupled model is referred to as POP2–waves.

109 The structure of this paper is organized as follows. Section 2 introduces the model,  
110 data, methodology, and experiments employed in this study. The performance of the  
111 POP2–waves coupled model in simulating sea–air CO<sub>2</sub> flux and interaction between  
112 CO<sub>2</sub> flux and carbonate–pH system in Section 3, while Section 4 focuses on the impact  
113 of different configurations of waves and bubble effect on ocean geochemistry response  
114 and uncertainty arising from the absence of interactions between ocean CO<sub>2</sub> flux and  
115 the carbonate–pH system. Finally, Section 5 presents the conclusions.

116

## 117 **2. Data, model experiments, and methodology**

### 118 **2.1 Observational data**

119 The spatial patterns of sea–air CO<sub>2</sub> flux simulated by the coupled CESM1 and  
120 POP2–waves coupled model are compared with estimates from the National Oceanic  
121 and Atmospheric Administration’s (NOAA) CarbonTracker modeling and data  
122 assimilation system version 2022 (CT2022; Jacobson et al., 2023), which provides  
123 global surface-atmosphere CO<sub>2</sub> flux estimates spanning January 2000 through  
124 December 2020. We analyzed the seasonal variations and mean states of CO<sub>2</sub> flux  
125 globally, especially in primary sources and sinks of CO<sub>2</sub> flux. Some uncertainties in the  
126 assimilation dataset arise from factors such as anthropogenic influences and ice melt;  
127 therefore, following Fay et al. (2024), this study excludes the coastal ocean and high-  
128 latitude sea regions from the model comparison.

129

### 130 **2.2 Description of the model framework and experiments**

131 In this study, we investigated the role of waves and bubble mechanisms through  
132 ocean geochemistry response between control POP2 and POP2–waves coupled model



133 in CESM1.2.2 framework. The CESM simulation is conducted using a fully coupled  
 134 component set over a 30-year period, with well-mixed greenhouse gases (CO<sub>2</sub>, CH<sub>4</sub>,  
 135 N<sub>2</sub>O, etc.), ozone, and aerosols fixed to year-2000 values (B\_2000\_CAM5 compsets).  
 136 It has a horizontal resolution of 1.9° × 2.5° in Community Atmosphere Model 5.3  
 137 (CAM5.3) and a nominal 1° horizontal resolution with 60 vertical layers in POP2. A  
 138 control CO<sub>2</sub> flux and seawater acidification simulation (B-CTL) uses Eq. (1), where  
 139  $K_w$  is the Wanninkhof (1992) equation (Eq. 2). We use the present-day scenario (starting  
 140 from the year 2000) to estimate sea-air CO<sub>2</sub> flux,  $pCO_2^w - pCO_2^a$  and pH across both  
 141 B-CTL and POP2-waves experiments, allowing us to focus on the impact of wave  
 142 effects on carbonate chemistry and to compare the results with the existing observations.  
 143 In the wave effect mechanism simulations, we coupled the waves module (Mellor et al.,  
 144 2008) of the POM model into POP2 (POP2-waves experiment) and incorporated  
 145 bubble-mediated gas transfer, computed from a mechanistic model for air bubble  
 146 entrainment at the breaking wave scale (Deike and Melville, 2018), combined with the  
 147 POP2 geochemistry response.

148 POP2-waves experiment adopts the parameterization proposed by Deike and  
 149 Melville (2018), Reichl and Deike (2020), and Deike (2022) for the bubble-mediated  
 150 gas transfer velocity as a function of friction velocity,  $u_*$  (m s<sup>-1</sup>) and significant wave  
 151 height,  $H_s$  (m). This parameterization captures the main wave effect, as well as  
 152 solubility and diffusivity, effectively consolidating for CO<sub>2</sub> by expressing the gas  
 153 transfer velocity as the sum of the non-bubble ( $k_{wNB}$ ) and bubble ( $k_{wB}$ ) components,  
 154 following Deike and Melville (2018):

$$155 \quad K_w = K_{wNB} + K_{wB} = A_{NB} u_* \left( \frac{Sc}{660} \right)^{-1/2} + \frac{A_B}{K_0 R T_0} u_*^{5/3} (g H_s)^{2/3} \quad (3)$$

156 where  $A_{NB}$  is an empirical, nondimensional coefficient given as  $1.55 \times 10^{-4}$ ,  $A_B$  is a  
 157 dimensional fitting coefficient ( $1.2 \times 10^{-5} \text{ s}^2 \text{ m}^{-2}$ ),  $R$  is the ideal gas constant (0.082 L



158  $\text{atm mol}^{-1} \text{K}^{-1}$ ),  $T_0$  is the sea surface temperature, SST (K), and  $g$  is gravitational  
 159 acceleration ( $9.806 \text{ m s}^{-1}$ ). In addition, Section 3.3 will examine the contribution ratios  
 160 of the non-bubble and bubble components in simulating the sea–air  $\text{CO}_2$  flux.

161

## 162 **2.3 Coupling POP2–waves Methodology**

163 The architecture for the POP2–waves experiment in the CESM1.2.2 framework is  
 164 shown in Fig. 1. In addition to the original component relationships in CESM1.2.2, we  
 165 obtain zonal and meridional surface winds, as well as friction velocity over the ice  
 166 fraction from the coupler (CPL) to calculate the wave’s property and  $K_w$ . To obtain  
 167 significant wave height from the wave’s module, we (1) interpolate the depth-variable  
 168 currents of the POP2 z-coordinates (60 levels of zonal and meridional horizontal  
 169 velocity) into the POM (waves) sigma-coordinates (21 levels) for calculating the depth-  
 170 dependent wave radiation equation and the specified spectrum, and (2) horizontal  
 171 progression from one grid to the adjacent grid to prevent discontinuities in wave energy  
 172 at block boundaries.

173

## 174 **3. Results**

### 175 **3.1 The climatological mean and seasonal variations of ocean $\text{CO}_2$ flux**

176 We use the traditional sea–air  $\text{CO}_2$  flux bulk formula (Eq. 1) in CESM1.2.2 while  
 177 also incorporating the effects of waves and bubble-mediated processes with  $K_w$ , which  
 178 is parameterized based on friction velocity, SST, and  $H_s$  (Eq. 3). Figure 2a–c compares  
 179 the climatological mean ocean  $\text{CO}_2$  flux ( $\text{mol m}^{-2} \text{yr}^{-1}$ ; shaded) and the monthly  
 180 standard deviation (SD; contours) among B–CTL, POP2–waves, and NOAA CT2022.  
 181 Overall, the spatial distribution of the model-simulated  $\text{CO}_2$  flux is broadly consistent  
 182 with that of CT2022, though some discrepancies persist. The primary sources and sinks  
 183 of average  $\text{CO}_2$  flux in the Pacific simulated by the POP2–waves experiment more



184 closely resemble those in NOAA CT2022 than those in B-CTL, although the POP2-  
 185 waves experiment exhibits a larger monthly SD than B-CTL over both the Pacific and  
 186 Indian Oceans. Figure 2c highlights the primary sources and sinks of CO<sub>2</sub> flux in the  
 187 Pacific using red dashed lines, while black dashed lines denote other significant regions.  
 188 Based on these highlighted areas, Fig. 2d presents 12 prominent source and sink regions  
 189 identified from NOAA CT2022, selected from the 30 regions defined in the NOAA  
 190 atmospheric CO<sub>2</sub> inversion model. These regions include: (1) Pacific Ocean –  
 191 Northwestern Pacific (NWP), Eastern Equatorial Pacific (EEP), Northeastern Pacific  
 192 (NEP), South Pacific Ocean 1 (SPO1), and South Pacific Ocean 2 (SPO2); (2) Indian  
 193 Ocean – North Indian Ocean (NIO), South Indian Ocean 1 (SIO1), and South Indian  
 194 Ocean 2 (SIO2); and (3) Atlantic Ocean – North Atlantic Ocean 1 (NAO1), North  
 195 Atlantic Ocean 2 (NAO2), Equatorial Atlantic Ocean (EAO), and South Atlantic Ocean  
 196 (SAO).

197 Figure 3 presents the nine characteristic regions of model-simulated CO<sub>2</sub> flux. The  
 198 seasonal variations in these regions are generally consistent with those in CT2022 over  
 199 the NWP, EEP, NEP, SPO1, NIO, SIO1, NAO1, EAO, and SAO. The black dashed line  
 200 represents B-CTL, the red solid line represents POP2-waves, and the blue solid line  
 201 represents NOAA CT2022. The bar graph shows the monthly standard deviation as a  
 202 percentage of the overall range (95% minus 5%). Black hollow bars correspond to B-  
 203 CTL, red bars to POP2-waves, and light blue hollow bars to NOAA CT2022. Fay et al.  
 204 (2024) attributed seasonal variations in sea-air fluxes to a combination of factors,  
 205 including fluctuations in SST, biological uptake of carbon dioxide, and variations in  
 206 mixing and wind speeds. The percentage of monthly SD is generally within 30%, with  
 207 higher values observed during the boreal winter, corresponding to strong oceanic CO<sub>2</sub>  
 208 sinks in high-latitude regions of the Northern Hemisphere, such as NWP, NEP, and  
 209 NAO1. In contrast, the Southern Hemisphere exhibits large monthly standard





210 deviations alongside relatively strong CO<sub>2</sub> sinks during the boreal summer in regions  
 211 such as SPO1, SIO1, and SAO. Both B–CTL and POP2–waves simulations in the EEP  
 212 display minimal seasonal variability in CO<sub>2</sub> flux, maintaining a persistent year-round  
 213 efflux, consistent with the findings of Fay et al. (2024) in the equatorial Pacific. In  
 214 contrast, NOAA CT2022 shows a greater degree of seasonal variability in the EEP.

215 Three regions—SPO2, SIO2, and NAO2—exhibit model discrepancies when  
 216 compared with NOAA CT2022 (Fig. 4), which are not attributable to wave effects or  
 217 the CO<sub>2</sub> flux bulk formula. In NAO2, both B–CTL and POP2–waves show seasonal  
 218 variations similar to those in NAO1, with strong oceanic CO<sub>2</sub> sinks during the boreal  
 219 winter, whereas NOAA CT2022 displays the opposite seasonal pattern in this region. A  
 220 similar pattern is observed in SPO2, where the seasonal variations in both B–CTL and  
 221 POP2–waves resemble those in SPO1 but contrast with the trend shown by NOAA  
 222 CT2022.

223

### 224 **3.2 The relationship between $U_{10}$ and $K_w$ over the areas of the Pacific's primary** 225 **sources and sinks of CO<sub>2</sub> flux**

226 The monthly gas transfer velocities for both the POP2–waves and B–CTL  
 227 experiments over the smaller area of the Pacific's primary sources and sinks of CO<sub>2</sub>  
 228 flux (Fig. 2c) is shown in Fig. 5 for (a) WP (160–180°E, 35–40°N, Fig. 5a) and (b) EP  
 229 (230–250°E, 0–5°S, Fig. 5b). Gray open circles represent  $K_w$  calculated using the  
 230 Wanninkhof (1992) 10-m wind speed equation (B–CTL) with their simple linear  
 231 regression (SLR) and squared correlation coefficient ( $R^2$ ) shown in black line and text.  
 232 Green open squares represent POP2–waves data for  $H_s < 1.5$ m with the orange solid  
 233 line and text indicating its SLR and  $R^2$ , while blue dots denote cases where  $H_s > 1.5$  m  
 234 with the red solid line and text depicting their SLR and  $R^2$ . Generally, the B–CTL data  
 235 shows higher  $R^2$  but lower  $K_w$  values compared to POP2–waves across all WP and EP.



POP2-waves exhibits a slightly lower  $R^2$  than B-CTL due to greater deviations between  $K_w$  and  $U_{10}$ , particularly at higher surface wind speeds, consistent with the findings of Gutiérrez-Loza et al. (2022) under conditions where  $H_s > 1.5$  m. Zhou et al. (2023) indicated that the wind-only formula for gas transfer velocity tends to underestimate values when  $U_{10}$  exceeds  $10 \text{ m s}^{-1}$ , and that enhanced gas transfer velocity frequently occurs under conditions of strong wave breaking and high significant wave height. These deviations occur when  $U_{10} > 12 \text{ m s}^{-1}$  over the WP (Fig. 5a) and around  $U_{10} = 9 \text{ m s}^{-1}$  (Fig. 5b) over the EP.

Figs. 5c and 5d display the monthly mean values of  $K_w$  for WP and EP, respectively. Red, orange solid, and black dotted lines/texts indicate POP2-waves ( $H_s > 1.5$  m), average of all POP2-waves data, and B-CTL, respectively. The bar graph illustrates the monthly SD as a percentage of the overall ranking range (95% minus 5%). The annual mean value of POP2-waves  $K_w$  in WP is greater than that in EP. Figure 5c shows that WP cases with  $H_s > 1.5$  m have monthly data available except in August and October, and their  $K_w$  values are higher than those from POP2-waves and B-CTL. In contrast, EP shows  $H_s > 1.5$  m data only during the summer months (July to October), and the corresponding  $K_w$  values are also higher than those from POP2-waves and B-CTL. The percentage of monthly SD in the bar graph indicates no significant differences between  $H_s > 1.5$  m data and POP2-waves experiments in the WP, while the  $K_w$  difference between the two datasets from January to May remains small because the average  $H_s$  during this period exceeds 1.5 m. However, in EP, a significant difference is observed between the  $H_s > 1.5$  m condition and the average in POP2-waves from July to September.

### 3.3 Regression of $\text{CO}_2$ flux against $U_{10}$ under high-wind conditions



261 The wind-only parameterization of gas transfer velocity tends to underestimate  
 262 values when  $U_{10} > 10 \text{ m s}^{-1}$  or  $H_s > 1.5 \text{ m}$  (e.g., Gutiérrez-Loza et al., 2022; Zhou et  
 263 al., 2023). Accordingly, Fig. 5 presents scatter plots of  $\text{CO}_2$  flux from B-CTL and  
 264 POP2-waves experiment against  $U_{10}$  (for  $U_{10} > 10 \text{ m s}^{-1}$ ) over the WP region (160–  
 265 180°E, 35–40°N) to compare fluxes estimated with wind-only  $K_w$  and wave-  
 266 modified  $K_w$ . High-wind conditions over the WP occur primarily in winter (42%  
 267 valid data), and the associated low ocean surface temperatures further enhance  $\text{CO}_2$   
 268 solubility. The slopes of the regression equations (–0.77 in B-CTL and –1.40 in  
 269 POP2-waves) and the mean  $\text{CO}_2$  fluxes ( $Y_{\text{avg}}$ : –6.35 in B-CTL and –6.97 in POP2–  
 270 waves) both indicate that POP2-waves represents a stronger  $\text{CO}_2$  sink in the WP  
 271 region under higher wind speeds. The correlation coefficient ( $R$ ) of POP2-waves is  
 272 higher than that of B-CTL, indicating that the sea state-dependent  $K_w$  provides a  
 273 more suitable estimate of  $\text{CO}_2$  flux than the wind-only  $K_w$  under high-wind  
 274 conditions, consistent with the findings of Gutiérrez-Loza et al. (2022) and Zhou et  
 275 al. (2023).

276

### 277 **3.4 Surface climate fields and $K_w$ component differences between POP2-waves** 278 **and B-CTL**

279 Figures 7a–b present the surface circulation along with sea level pressure (PSL),  
 280  $U_{10}$ , and  $H_s$ , all of which are factors influencing  $K_w$  in the model. The region with  $H_s >$   
 281 1.5 m is generally situated where  $U_{10} > 8 \text{ m s}^{-1}$  and lies south of the high PSL region.  
 282 Johansson et al. (2022) reported that strong winds can either fill or empty small sub-  
 283 basins while simultaneously generating high waves. Similarly, the ratio of  $K_{wB}/K_w$  in  
 284 POP2-waves greater than 0.5 (yellow shaded areas in Fig. 7b) mostly occurs in regions  
 285 where  $H_s > 1.5 \text{ m}$ . However, high  $H_s$  does not always correspond to a higher  $K_{wB}/K_w$  in  
 286 POP2-waves, for example, in the North Pacific. Deike and Melville (2018)



287 demonstrated that the bubble contribution to CO<sub>2</sub> transfer exceeds 40% when the  $U_{10}$  is  
 288  $\geq 10 \text{ m s}^{-1}$ . The bubble contribution to the total gas transfer velocity ( $K_{wB}/K_w$ ) in POP2–  
 289 waves is approximately 38%, which is slightly higher than ~30% reported by Reichl  
 290 and Deike (2020). This discrepancy may be attributed to differences in the  $A_B$   
 291 coefficients used in this study compared to previous work. The coefficient in the sea-  
 292 state formulation ( $A_B$  in Eq. (3)) is derived from field data, where gas transfer velocity  
 293 is estimated from flux measurements using eddy covariance methods (Reichl and Deike,  
 294 2020; Brumer et al., 2017).

295 B–CTL  $K_w$  and POP2–waves  $K_{wNB}$  exhibit similar spatial distributions, with values  
 296 exceeding  $0.006 \text{ cm s}^{-1}$  and  $0.004 \text{ cm s}^{-1}$ , respectively, despite using different input  
 297 parameters ( $U_{10}$  vs.  $u^*$ ). Aside from the polar regions,  $K_w$  in the POP2–waves coupled  
 298 model is greater than that in B–CTL with the largest differences occurring in the tropics  
 299 (Fig. 7c). However, the values of  $K_{wNB}$  are approximately 30–50% lower than  $K_w$  in B–  
 300 CTL (Fig. 7d).

301

## 302 4. Discussion

### 303 4.1 Mean State and model differences in $p\text{CO}_2^w$ and surface pH

304 Previous studies have not clearly identified the discrepancy in CO<sub>2</sub> flux estimates  
 305 between bulk formulas (e.g., Wanninkhof, 1992, 2014) and approaches that incorporate  
 306 breaking waves and bubble injection to parameterize  $K_w$  in an Earth System Model.  
 307 Some studies incorporate wave effects following Deike and Melville (2018), using the  
 308 the Surface Ocean CO<sub>2</sub> Atlas (SOCAT) database (Bakker et al., 2016) to estimate CO<sub>2</sub>  
 309 fluxes. In these approaches, CO<sub>2</sub> flux is treated independently of  $p\text{CO}_2^w$  values and  
 310 ocean pH, with no feedback between them. Our study highlights that the ocean CO<sub>2</sub>  
 311 system is influenced by the wave-induced  $K_w$ , affecting parameters such as alkalinity,  
 312 pH, and carbonate species, which in turn feedback into the dynamics of CO<sub>2</sub> flux. In



313 other words, perturbations in  $K_w$  influence both  $dpCO_2$  and  $CO_2$  flux in the Earth  
 314 System Model, differing from approaches that treat wave-induced  $K_w$  and  $dpCO_2$   
 315 independently when estimating  $CO_2$  flux (e.g., Reichl and Deike, 2020).

316 Figure 8 illustrates the means and differences in the spatial distribution of  $pCO_2^w$   
 317 and pH when using the  $U_{10}$  bulk formula for  $K_w$  versus the wave-induced  $K_w$ . In B–  
 318 CTL, high  $pCO_2^w$  is concentrated in the equatorial eastern Pacific (EEP), contributing  
 319 to a larger oceanic  $CO_2$  source (Fig. 2a). In contrast, low  $pCO_2^w$  is found in the high-  
 320 latitude regions of the Pacific and Atlantic (e.g., NWP, NEP, NAO1, and NAO2),  
 321 resulting in a strong oceanic  $CO_2$  sink (Fig. 2a). The standard deviation of B–CTL  
 322  $pCO_2^w$  varies between 10 and 30 ppm over time (Fig. 8a).

323 Figure 2 shows that the  $CO_2$  flux in POP2–waves is stronger than that in B–CTL  
 324 across both source and sink regions, consistent with the higher  $K_w$  indicated by contours  
 325 in Fig. 7c. Notably, POP2–waves reduces the high  $pCO_2^w$  by more than 10 ppm over  
 326 the equatorial western Pacific and increases the low  $pCO_2^w$  in other regions (Fig. 8b).  
 327 In the CESM1 framework, atmospheric  $pCO_2^a$  is held constant. Consequently, in  
 328 source (sink) regions, the POP2–waves coupled model decreases (increases)  $pCO_2^w$ ,  
 329 thereby reducing  $dpCO_2$ . This mitigates the large increase in  $CO_2$  flux simulated by the  
 330 POP2–waves model—a process referred to as 'negative feedback from  $dpCO_2$ —and is  
 331 discussed further in Section 4.3. The SD of the  $pCO_2^w$  difference between POP2–waves  
 332 and B–CTL varies significantly in EEP, while SD changes in other regions are less  
 333 pronounced.

334 Carbonic acid is a weak acid that can dissociate into bicarbonate ( $HCO_3^-$ ) and  
 335 hydrogen ions ( $H^+$ ), thereby affecting the pH of the solution. The average pH of the  
 336 ocean is currently around 8.1 (Fig. 8c), which is slightly alkaline; however, as the ocean  
 337 continues to absorb more  $CO_2$ , its pH decreases, making it more acidic. Ocean pH and  
 338  $pCO_2^w$  are strongly negatively correlated (Macovei et al. 2021), with low pH values



339 observed in the equatorial eastern Pacific (Fig. 8c), resulting from high  
 340  $pCO_2^w$  concentrations in that region (Fig. 8a). Slightly higher pH values are observed in  
 341 regions with lower  $pCO_2^w$ , such as the NWP, NEP, NAO1, and NAO2. Accounting for  
 342 wave effects results in simulated pH differences of approximately + (-) 0.01 in ocean  
 343 source (sink) regions. This suggests that including this mechanism mitigates  
 344 acidification in the equatorial Pacific but enhances ocean acidification in the NWP, NEP,  
 345 NAO1, and NAO2 regions (Fig. 8d).

346

#### 347 **4.2 Impact of sea–air $dpCO_2$ , $K_w$ , SST, and pH on $CO_2$ flux**

348 To determine which factor is directly associated with oceanic  $CO_2$  flux, Fig. 9  
 349 presents the linear regression coefficients (LRC) between  $CO_2$  flux and various  
 350 normalized variables:  $dpCO_2$  (Fig. 9a),  $K_w$  (Fig. 9b), SST (Fig. 9c), and pH (Fig. 9d).  
 351 Here, we do not delve into the complex chemical and biochemical mechanisms in  
 352 carbonate species; although, the influence of carbonate species composition, including  
 353 alkalinity (ALK), and dissolved inorganic carbon (DIC), on  $pCO_2$  is greater than that  
 354 of solubility changes driven by sea surface salinity and SST (Koseki et al., 2023).

355 Among these factors,  $dpCO_2$  and  $CO_2$  flux exhibit the strongest positive LRC  
 356 (0.6–0.8) across the global ocean, with the ocean acting as a  $CO_2$  source or sink  
 357 depending on whether the sea–air  $dpCO_2$  is positive or negative. Additionally, wave  
 358 effects increase the LRC by approximately 0.1 to 0.2, particularly in regions with strong  
 359 oceanic  $CO_2$  sources and sinks. There is a clear positive (negative) LRC between  $K_w$   
 360 and  $CO_2$  flux in oceanic  $CO_2$  source (sink) regions. Furthermore, wave effects reduce  
 361 the absolute LRC between  $K_w$  and  $CO_2$  flux in both oceanic  $CO_2$  source and sink regions  
 362 (Fig. 9b), indicating that wave-influenced  $CO_2$  flux is less correlated with  $K_w$  and more  
 363 closely linked to  $dpCO_2$ , after considering interactions between modeled  $CO_2$  flux and  
 364 the ocean carbonate-pH system.



365 In contrast to the former, there is a clear positive (negative) LRC between SST and  
 366 CO<sub>2</sub> flux in oceanic CO<sub>2</sub> sink (source) regions (Fig. 9c). This pattern arises because  
 367 CO<sub>2</sub> solubility increases at lower temperatures (e.g., high-latitude regions), enhancing  
 368 oceanic uptake of atmospheric CO<sub>2</sub>, whereas higher oceanic temperatures (e.g., tropical  
 369 regions) reduce solubility, leading to decreased absorption and greater CO<sub>2</sub> release from  
 370 the ocean to the atmosphere. Wave effects only slightly reduce the LRC between SST  
 371 and CO<sub>2</sub> flux—by approximately 0.1—in the tropics, with negligible differences  
 372 observed in other regions. The strong negative correlation between ocean pH and  
 373  $pCO_2^w$  is attributed to the acid dissociation constants of H<sub>2</sub>CO<sub>3</sub> (Fig. 8a and 8c), which  
 374 govern the relationship between dissolved CO<sub>2</sub> and hydrogen ion concentration  
 375 (Williams et al., 2017). pH and CO<sub>2</sub> flux exhibit strong negative LRC ( $< -0.9$ ) across  
 376 the global ocean, with only a few exceptions in the EEP and along the Arctic margins  
 377 (Fig. 9d). Wave effects again only slightly reduce the LRC between pH and CO<sub>2</sub> flux  
 378 in the EEP—by approximately 0.1.

379

#### 380 **4.3 Uncertainty arising from the absence of interactions between CO<sub>2</sub> flux and the** 381 **ocean carbonate–pH system**

382 To assess the uncertainty in CO<sub>2</sub> flux resulting from the absence of interactions  
 383 between CO<sub>2</sub> flux and the ocean carbonate–pH system (i.e., the lack of dpCO<sub>2</sub>-driven  
 384 negative feedback), Fig. 10a presents the CO<sub>2</sub> flux calculated using  $dpCO_2$  values (from  
 385 B–CTL) and  $K_w$  incorporating wave effects (from POP2–waves), with shading  
 386 representing the mean (labeled as "no dpCO<sub>2</sub> feedback") and contours indicating their  
 387 SD. Only data significant at a 95% confidence level of CO<sub>2</sub> flux difference between "no  
 388 dpCO<sub>2</sub> feedback" and POP2–waves are shown in Fig. 10b. Compared to Fig. 2b, the  
 389 CO<sub>2</sub> flux in the "no dpCO<sub>2</sub> feedback" case is stronger than that in the POP2–waves  
 390 simulation across both source and sink regions, and is accompanied by a larger SD. The



391 most pronounced differences occur in the Pacific source and sink regions; however,  
 392 some areas (e.g., NEP, SIO<sub>2</sub>, and SPO<sub>2</sub>) do not reach the 95% confidence threshold.

393 The ocean's buffering capacity — including the interactions between CO<sub>2</sub> flux  
 394 (estimated from  $K_w$ ) and the ocean carbonate–pH system — incorporating wave effects  
 395 mitigates pH acidification and reduces  $dp\text{CO}_2$  differences over CO<sub>2</sub> source regions,  
 396 while higher pH levels enhance the ocean's ability to absorb more CO<sub>2</sub> from the  
 397 atmosphere over CO<sub>2</sub> sink regions in the POP2–waves coupled model (Fig. 8b and 8d).  
 398 In contrast, the "lack of  $dp\text{CO}_2$ -driven negative feedback" case, which lacks real ocean  
 399 buffering effects, allows CO<sub>2</sub> source regions to release excess CO<sub>2</sub> into the atmosphere  
 400 and leads to increased excess CO<sub>2</sub> uptake in sink regions (Fig. 10b) compared to the  
 401 POP2–waves coupled model.

402

## 403 **5 Conclusions**

404 This study quantified oceanic sea–air CO<sub>2</sub> fluxes using both a control run of the  
 405 CESM1 model and a newly developed POP2–waves coupled model, which builds on  
 406 the CESM1 framework by incorporating the effects of wave dynamics, sea surface  
 407 temperature, pH, and gas transfer velocity, including interactions between CO<sub>2</sub> flux and  
 408 the ocean carbonate–pH system. The architecture of the POP2–waves coupled model  
 409 primarily addresses two challenges: (1) interpolating currents from the POP2 z-  
 410 coordinate system to the sigma-coordinate system for wave radiation equation  
 411 calculations, and (2) horizontal progression between adjacent grids is implemented to  
 412 prevent discontinuities in wave energy at block boundaries (Fig. 1). Several recent  
 413 studies have estimated the sea–air CO<sub>2</sub> flux using Eq. (1), although current  $K_w$   
 414 parameterizations in global models (Table 1) still rely exclusively on  $U_{10}$  (e.g.,  
 415 Chikamoto and DiNezio, 2021; Danabasoglu et al., 2020; Seland et al., 2020; Mauritsen  
 416 et al., 2019; Ziehn et al., 2020; Lovato et al., 2022; Sigmond et al., 2023).





417 In the POP2–waves coupled model,  $K_w$  follows the parameterization proposed by  
 418 Deike and Melville (2018), which relates the bubble-mediated  $K_w$  to friction velocity  
 419 and significant wave height. This parameterization captures the primary wave effects,  
 420 as well as solubility and diffusivity, and expresses  $K_w$  as the sum of non-bubble and  
 421 bubble components in Eq. (3), which is used to compare oceanic CO<sub>2</sub> fluxes with  
 422 estimates from NOAA CarbonTracker CT2022 data. Overall, the spatial distribution of  
 423 the model-simulated CO<sub>2</sub> flux is broadly consistent with that of CT2022, although some  
 424 discrepancies remain. The prominent source and sink regions of the average CO<sub>2</sub> flux  
 425 simulated by the POP2–waves experiment more closely resemble those in NOAA  
 426 CT2022 than those in B–CTL (Figs. 2 and 3), although notable discrepancies remain in  
 427 three regions—SPO2, SIO2, and NAO2—when compared with NOAA CT2022 (Fig.  
 428 4).

429 We find that larger deviations between POP2–waves  $K_w$  and  $U_{10}$  occur when  $H_s >$   
 430 1.5 m (Fig. 5a–b), consistent with Gutiérrez-Loza et al. (2022), indicating that POP2–  
 431 waves coupled model exhibits a stronger CO<sub>2</sub> flux than B–CTL under high  $U_{10}$  (Fig. 6).  
 432 In addition, the bubble contribution to the total gas transfer velocity ( $K_{wB}/K_w$ ) in POP2–  
 433 waves is approximately 38% over the entire space–time domain (Fig. 7(b)), slightly  
 434 higher than the ~30% reported by Reichl and Deike (2020). These high  $K_{wB}/K_w$  regions  
 435 coincide with areas of elevated  $H_s$  and  $U_{10}$ . Likewise, our results show that bubbles  
 436 account for up to 41.3% of the total sea–air CO<sub>2</sub> flux, consistent with Zhou et al. (2023)  
 437 and Reichl and Deike (2020), who reported contributions of about 40%.

438 POP2–waves reduce high  $p\text{CO}_2$  by more than 10 ppm over the equatorial western  
 439 Pacific and increase low  $p\text{CO}_2$  in other regions (Fig. 8b). Wave effects also result in  
 440 simulated pH differences of approximately positive (negative) 0.01 in ocean source  
 441 (sink) regions (Fig. 8d). To identify the factor most directly associated with oceanic  
 442 CO<sub>2</sub> flux, Figs. 9a and 9d shows that  $dp\text{CO}_2$  (pH) exhibits the strongest positive



443 (negative) regression coefficient across the global ocean. Additionally,  $K_w$  shows a  
 444 positive (negative) regression coefficient with  $\text{CO}_2$  flux in source (sink) regions,  
 445 whereas SST displays the opposite pattern relative to  $K_w$  (Figs. 9b and 9c).

446 Under the  $dp\text{CO}_2$ -driven negative feedback associated with the interaction  
 447 between  $\text{CO}_2$  fluxes and the carbonate–pH system, POP2–waves show increases of  
 448 11.8%, 41.6%, and 1.8% in the  $\text{CO}_2$  sink, source, and global average, respectively,  
 449 compared to B–CTL. However, in the absence of  $dp\text{CO}_2$ -driven negative feedback, both  
 450  $\text{CO}_2$  sink and source regions exhibit unreasonably stronger fluxes than those in POP2–  
 451 waves (Fig. 10b), indicating that the increase in  $K_w$  due to wave effects—when coupled  
 452 with  $dp\text{CO}_2$ -driven feedback—helps moderate the abnormal  $\text{CO}_2$  flux enhancement,  
 453 bringing the estimates closer to NOAA CT2022 (Fig. 3).

454 Finally, the effects of the  $K_w$  bulk formula (Wanninkhof et al., 1992) and wave  
 455 dynamics (Deike and Melville, 2018) on sea–air  $\text{CO}_2$  flux (arrows), surface pH (tube  
 456 colors indicator), and  $dp\text{CO}_2$  (circle markers) in the strongest sink (WP) and source (EP)  
 457 regions during the DJF and JJA seasons are schematically summarized in Fig. 11. B–  
 458 CTL is illustrated in gray with black text, while differences between POP2–waves and  
 459 B–CTL are highlighted in red. Each panel features the ocean–atmosphere interface, with  
 460 ocean color shading representing relative SST levels across seasons, and the SST value  
 461 displayed in the lower-left corner of each panel. The B–CTL sea–air  $\text{CO}_2$  flux shows a  
 462 strong sink in DJF but only a weak sink in JJA over the WP, indicating that reduced  
 463 solubility at higher SSTs weakens the oceanic  $\text{CO}_2$  sink. POP2–waves further diminish  
 464 this sink effect. The seasonal differences in  $dp\text{CO}_2$ , which correspond to changes in  $\text{CO}_2$   
 465 flux and are indicated by circle markers in Fig. 11a and 11c, reveal the interaction  
 466 between  $\text{CO}_2$  flux and the ocean carbonate–pH system. In contrast, the EP region  
 467 exhibits lower pH than the WP and acts as a strong  $\text{CO}_2$  source in both DJF and JJA



468 (Fig. 11b and 11d), with no obvious seasonal variation (consistent with Fay et al., 2024).

469 POP2–waves enhance the CO<sub>2</sub> flux but reduce the  $dp\text{CO}_2$  difference.

470

471 *Code and data availability.* The model code of POP2–waves coupled model is available  
472 at <https://doi.org/10.5281/zenodo.15795234>. Input data of POP2–waves using the  
473 climatological Hadley Centre Sea Ice and Sea Surface Temperature dataset and  
474 GODAS data forcing, including 30-year numerical experiments, are available at  
475 <https://doi.org/10.5281/zenodo.5510795>.

476

477 *Author contributions.* YYL is the sole developer of the POP2–waves coupled model  
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479 analysis suggestions. WLL supports reorganization and offers analytical  
480 recommendations and SC offers a non-parallel wave module as part of the POM source  
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482

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498

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954 Table 1. Intercomparison of Earth System Model for estimating sea–air  
 955 CO<sub>2</sub> flux

Earth System Model	Ocean components	CO <sub>2</sub> flux consider waves	$K_w$ 's element	Reference
CESM1	POP2	no	$U_{10}$ (Wannikhof 1992)	Chikamoto and DiNezio, 2021; Chikamoto et al., 2023
CESM2	POP2	no	$U_{10}$ (Wannikhof 2014)	Danabasoglu et al., 2020
NorESM2	BLOM	no	$U_{10}$ (Wannikhof 2014)	Seland et al., 2020; Tjiputra et al., 2020
MPI-ESM1.2	MPIOM 1.6	no	$U_{10}$ (Wannikhof 2014)	Mauritsen et al., 2019; Liu et al., 2021
ACCESS-ESM1.5	MOM5	no	$U_{10}$ (Wannikhof 1992)	Ziehn et al., 2020
CMCC-ESM2	NEMO v3.6	no	$U_{10}$ (Wannikhof 1992)	Lovato et al., 2022
CanESM5	CanNEMO	no	$U_{10}$ (Wannikhof 1992)	Sigmond et al., 2023
GFDL-ESM4.1	MOM6	no	$U_{10}$ (Wannikhof 2014)	Stock et al., 2020; Roobaert et al., 2024
IPSL-CM6A-LR	NEMO-PISCES	no	$U_{10}$ (Wannikhof 1992)	Boucher et al., 2020
MIROC-ES2L	COCO 4.0	no	$U_{10}$ (Wannikhof 1992)	Hajima et al., 2020
UKESM1	NEMO v3.6	no	$U_{10}$ (Wannikhof 1992)	Yool et al., 2021
CNRM-ESM2-1	NEMO-PISCES	no	$U_{10}$ (Wannikhof 1992)	Séférian et al., 2019
CESM1-POP2–waves	POP2–wave coupled model	Yes	$H_s, u^*$ (Deike and Melville, 2018)	This study

956 Note: CESM1: the Community Earth System Model version 1 of the National Center  
 957 for Atmospheric Research (NCAR); CESM2: CESM version 2; POP2: the Parallel  
 958 Ocean Program version 2; NorESM2: The Norwegian Earth System Model version 2;  
 959 BLOM: Bergen Layered Ocean Model; MPI-ESM1.2: the Max Planck Institute for  
 960 Meteorology Earth System Model version 1.2; MPIOM 1.6: the Max-Planck Institute  
 961 Ocean Model version 1.6; ACCESS-ESM1.5: the Australian Community Climate and  
 962 Earth System Simulator form an Earth System Model version 1.5; MOM5: the GFDL  
 963 Modular Ocean Model version 5; CMCC-ESM2: the Euro-Mediterranean Centre on  
 964 Climate Change (CMCC) Earth System Model version 2; NEMO v3.6: Nucleus for  
 965 European Modelling of the Ocean version 3.6; CanESM5: the Canadian Earth System  
 966 Model version 5; CanNEMO: NEMO version 3.4 modified for CanESM; GFDL-  
 967 ESM4.1: the Geophysical Fluid Dynamics Laboratory's Earth System Model 4.1;  
 968 MOM6: the GFDL Modular Ocean Model version 6; IPSL-CM6A-LR :version 6 of



969 the Institut Pierre-Simon Laplace (IPSL) climate model; NEMO-PISCES: Nucleus for  
970 European Modelling of the Ocean, Pelagic Interaction Scheme for Carbon and  
971 Ecosystem Studies ocean general circulation and biogeochemistry model; MIROC-  
972 ES2L: the Model for Interdisciplinary Research on Climate, Earth System version 2;  
973 COCO 4.0: CCSR (Center for Climate System Research) Ocean Component Model  
974 version 4.0; UKESM1: the U.K. Earth System Model; CNRM-ESM2-1: the Earth  
975 system (ES) model of second generation developed by the Centre National de  
976 Recherches Météorologiques (CNRM); CESM1-POP2-wave: POP2-waves coupled  
977 model based on CESM1.  
978





979 **Figure List**

980 **Figure 1.** Architecture diagram for POP2–waves experiment in CESM1.2.2  
 981 framework, including components for the atmosphere [Community Atmosphere  
 982 Model version 5 (CAM5)], land [Community Land Model version 4 (CLM4)], ocean  
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 984 Passing Interface (MPI) utilities], sea ice [prognostic Los Alamos Sea Ice Model  
 985 (CICE)], and the coupler (CPL).

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 987 **Figure 2.** Climatological mean of ocean CO<sub>2</sub> flux (mol m<sup>-2</sup> yr<sup>-1</sup>; shaded) and monthly  
 988 standard deviation (SD: contour): (a) B–CTL; (b) POP2–waves; (c) NOAA CT2022:  
 989 red dashed lines highlight the primary sources and sinks of CO<sub>2</sub> flux in the Pacific,  
 990 while black dashed lines denote other significant areas; (d) Based on NOAA CT2022,  
 991 12 prominent source and sink regions have been identified among the 30 areas in the  
 992 NOAA atmospheric CO<sub>2</sub> inversion model. These include: (1) Pacific Ocean –  
 993 Northwestern Pacific (NWP), Eastern Equatorial Pacific (EEP), Northeastern Pacific  
 994 (NEP), South Pacific Ocean 1 (SPO1), and South Pacific Ocean 2 (SPO2); (2) Indian  
 995 Ocean – North Indian Ocean (NIO), South Indian Ocean 1 (SIO1), and South Indian  
 996 Ocean 2 (SIO2); (3) Atlantic Ocean – North Atlantic Ocean 1 (NAO1), North Atlantic  
 997 Ocean 2 (NAO2), Equatorial Atlantic Ocean (EAO), and South Atlantic Ocean  
 998 (SAO).

999  
 1000 **Figure 3.** The monthly mean values of the high-correlation climatological average of  
 1001 CO<sub>2</sub> flux with NOAA CT2022 are based on nine characteristic regions (NWP, EEP,  
 1002 NEP, SPO1, NIO, SIO1, NAO1, EAO, and SAO) shown in Fig. 2d. The black dashed  
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 1004 line represents NOAA CT2022. The bar graph illustrates the monthly standard  
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 1006 bars correspond to B–CTL, red solid bars to POP2–waves, and light blue hollow bars  
 1007 to NOAA CT2022.

1008  
 1009 **Figure 4.** Same as Fig. 3, but for the low-correlation climatological average of the  
 1010 simulated CO<sub>2</sub> flux compared with NOAA CT2022, based on the three characteristic  
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 1014 Pacific (WP; 160–180°E, 35–40°N) and (b) the Eastern Pacific (EP; 230–250°E, 0–  
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1017 Corresponding simple linear regressions (SLRs) and squared correlation coefficients  
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1025  
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 1028 for (a) B–CTL and (b) POP2–waves experiments. Black dots indicate monthly mean  
 1029 values from each experiment. Blue lines denote the 95% and 5% confidence limits of  
 1030 the mean response, while red lines denote the 95% and 5% prediction intervals. The  
 1031 regression equation ( $Y = mx + b$ ), correlation coefficient ( $R$ ), mean  $\text{CO}_2$  flux ( $Y_{\text{avg}}$ ),  
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 1035 and its corresponding ratio, as depicted in the figure: (a) Shaded areas represent 10-m  
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 1037 The ratio of bubble-mediated components ( $K_{wB}$ ) to POP2–waves  $K_w$  is shown in color,  
 1038 with contours representing significant wave height (m). (c) Shaded areas represent  $K_w$   
 1039 (B–CTL), with contours depicting the ratios of  $K_w$  ((POP2–waves–B–CTL)/B–CTL).  
 1040 (d) Shaded areas show the non-bubble-mediated component ( $K_{wNB}$ ), while contours  
 1041 represent the ratios of  $(K_{wNB} - K_w(\text{B–CTL}))/K_w(\text{B–CTL})$ .

1042  
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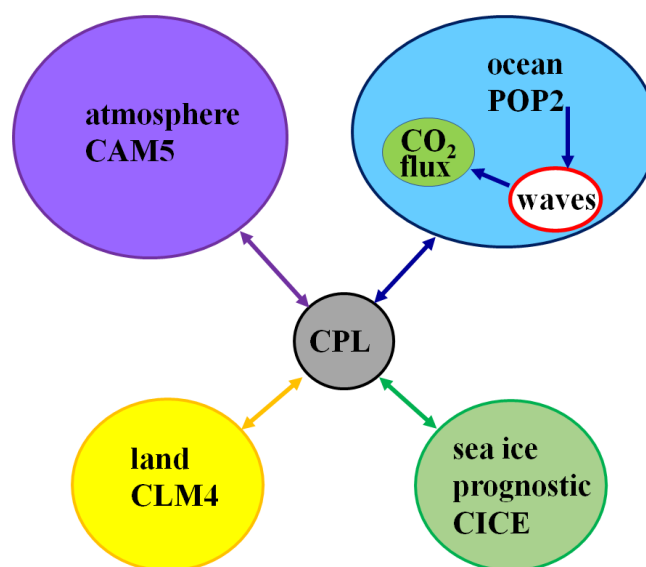
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1055 **Figure 10.** Estimated CO<sub>2</sub> flux assuming lack of  $dp\text{CO}_2$ -driven negative feedback. (a)  
 1056 CO<sub>2</sub> flux calculated using independent  $dp\text{CO}_2$  values (from B-CTL) and  $K_w$   
 1057 incorporating wave effects (labeled as "lack of  $dp\text{CO}_2$ -driven negative feedback"),  
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 1059 CO<sub>2</sub> flux between the "lack of  $dp\text{CO}_2$ -driven negative feedback" case and the POP2-  
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 1061 passing a 95% confidence level.

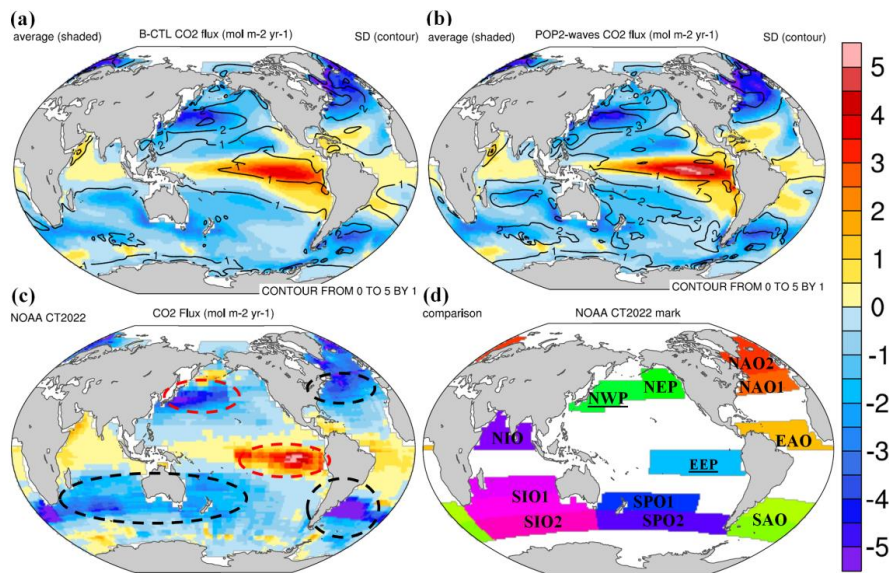
1062  
 1063 **Figure 11.** Schematic diagrams illustrate B-CTL (gray colors) and the differences  
 1064 between POP2-waves and B-CTL (red colors) in CO<sub>2</sub> flux (arrows), surface pH (tube  
 1065 colors indicator), and sea-air  $dp\text{CO}_2$  (circle markers) over the CO<sub>2</sub> sink (WP) and  
 1066 CO<sub>2</sub> source (EP). Each panel includes an ocean-atmosphere interface, with ocean  
 1067 color shading representing relative SST levels across different seasons. The lower-left  
 1068 corner displays the SST value, (a) WP in DJF seasonal mean, (b) EP in DJF seasonal  
 1069 mean, (c) WP in JJA seasonal mean, and (d) EP in JJA seasonal mean.



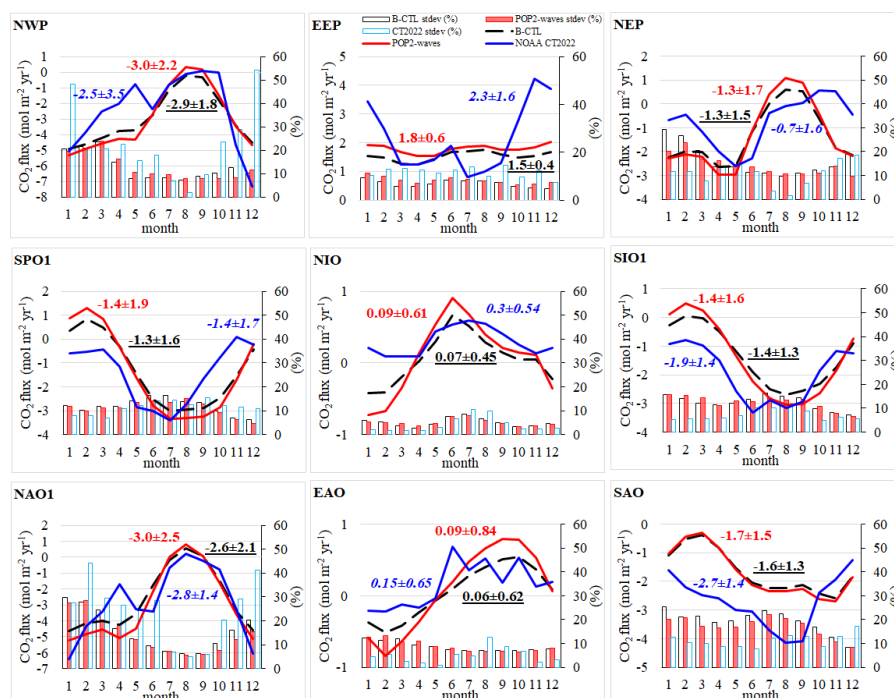
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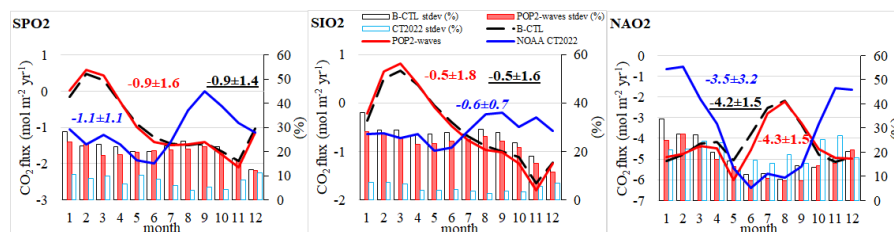


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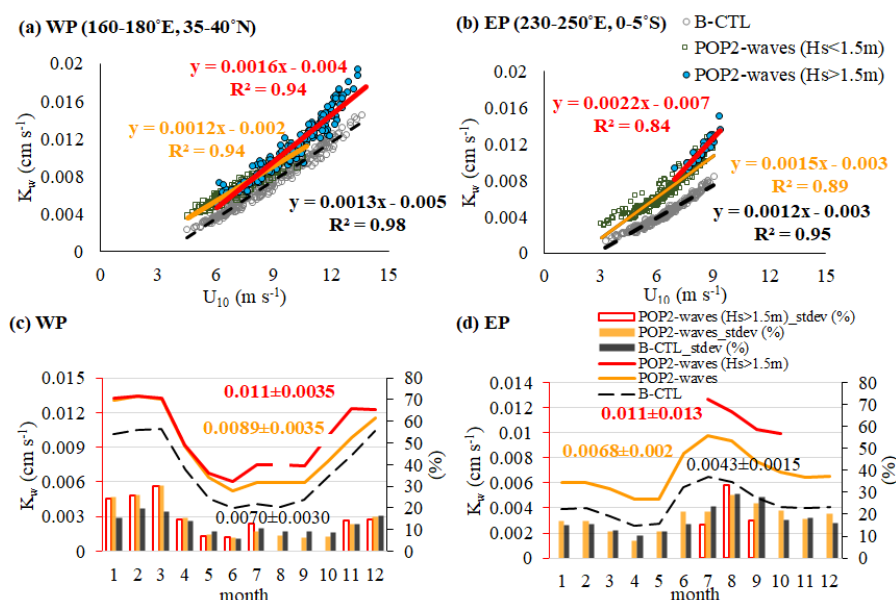
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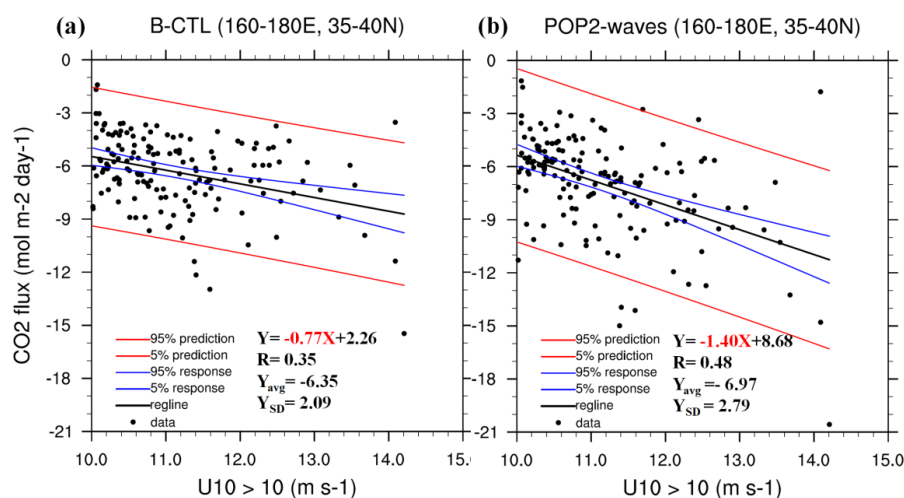
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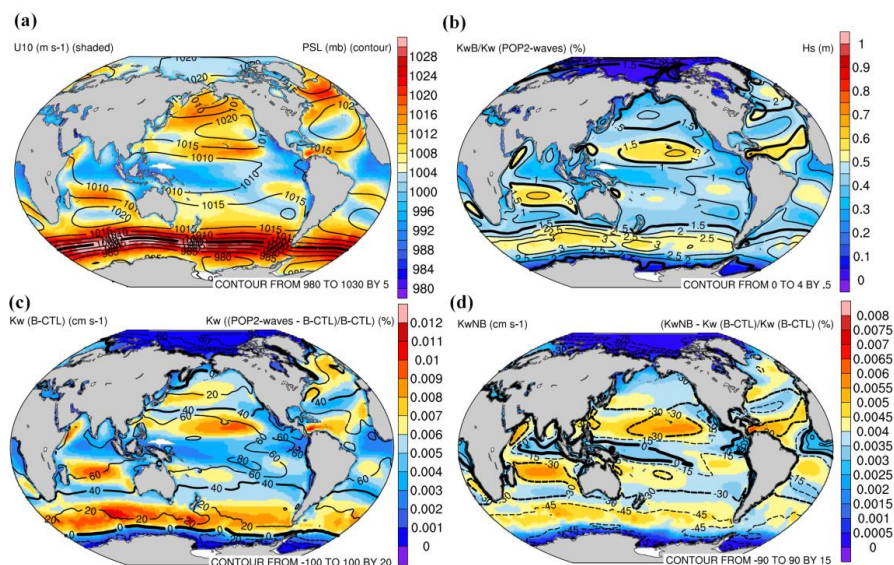
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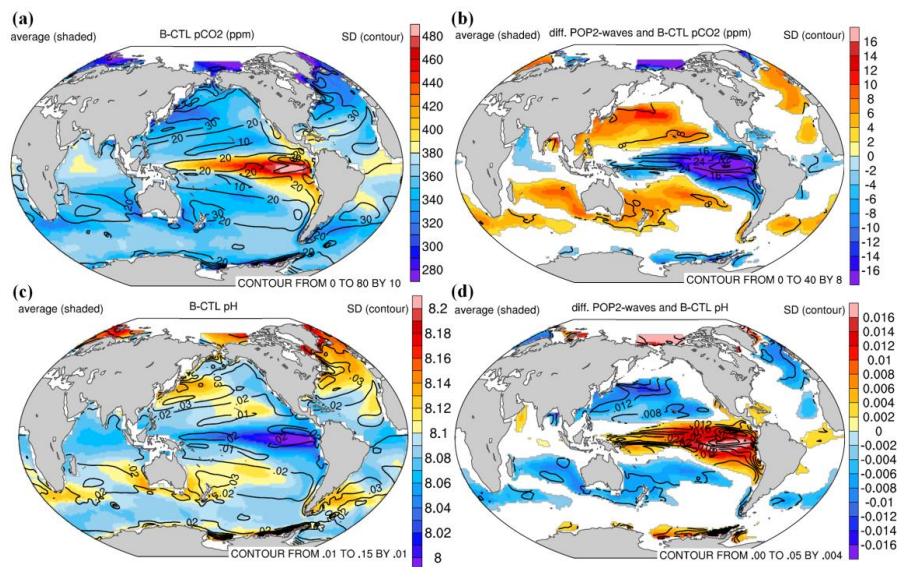




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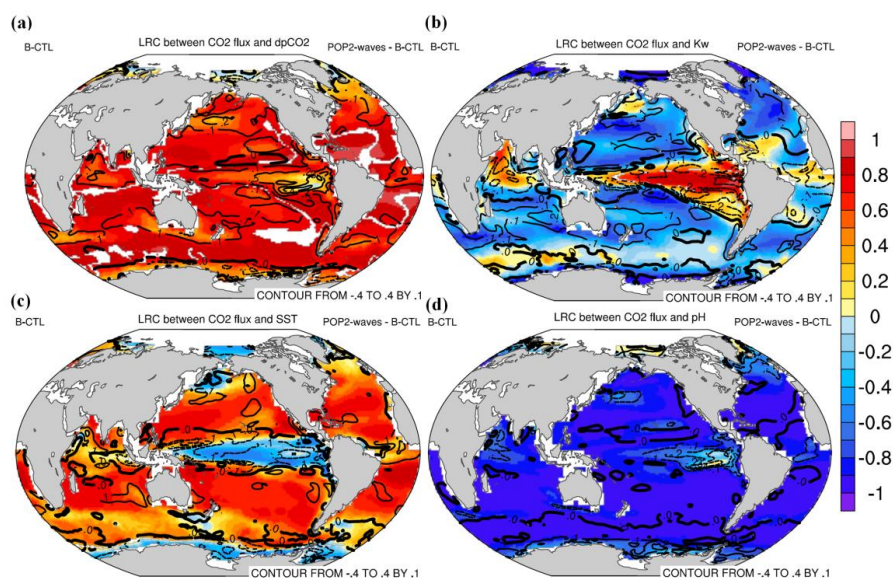


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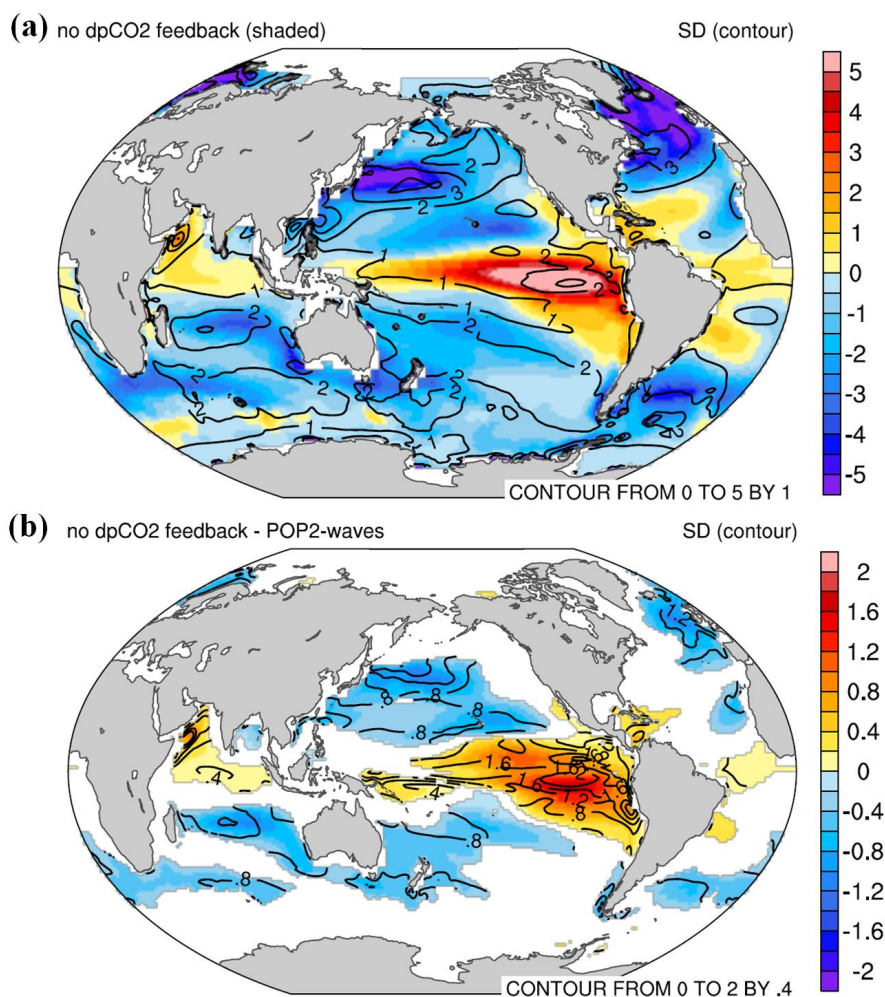


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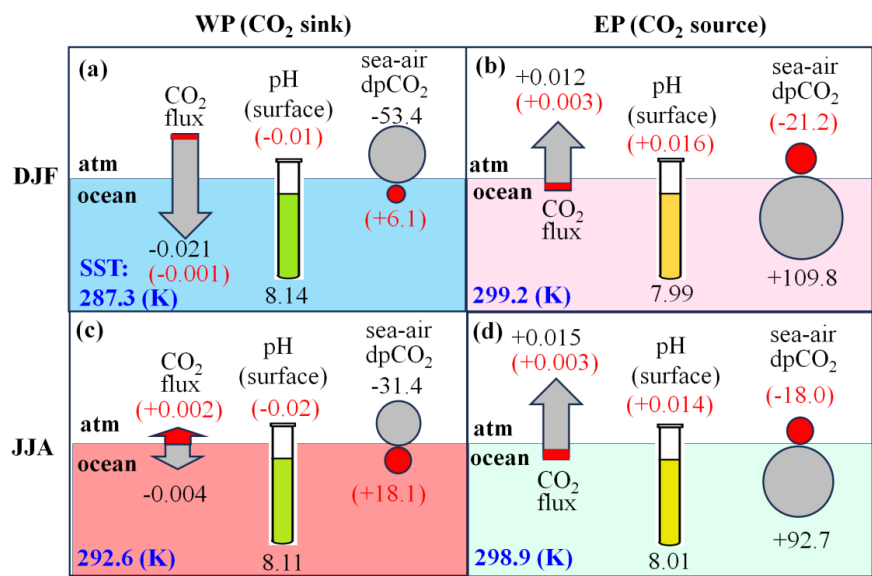
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1159 **Figure 11.** Schematic diagrams illustrate B-CTL (gray colors) and the differences  
 1160 between POP2-waves and B-CTL (red colors) in CO<sub>2</sub> flux (arrows), surface pH (tube  
 1161 colors indicator), and sea-air  $dpCO_2$  (circle markers) over the CO<sub>2</sub> sink (WP) and  
 1162 CO<sub>2</sub> source (EP). Each panel includes an ocean-atmosphere interface, with ocean  
 1163 color shading representing relative SST levels across different seasons. The lower-left  
 1164 corner displays the SST value, (a) WP in DJF seasonal mean, (b) EP in DJF seasonal  
 1165 mean, (c) WP in JJA seasonal mean, and (d) EP in JJA seasonal mean.