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30 — ~~runup model~~. The prediction models performed well both in the training and testing stages, with higher  $R^2$

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31	values. The developed model was, and were validated by comparing the results with via established	Formatted
32	statistical indices and by performing sensitivity and parametric analysis, highlighting	Formatted
33	that wave amplitude is greatly influenced by. The prediction models highlighted a cumulative 90% contribution of impact velocity, cliff height, and the	Formatted
34	number of fragments, contributing approximately 90% to on the wave amplitude. In	Formatted
35	comparison, runup is was greatly influenced by bank slope, angle, impact velocity, cliff mass,	Formatted
36	and height. The experimental results and developed prediction models can provide the	Formatted
37	basis for understanding the rotational cliff collapse-induced waves and can help with	Formatted
38	disaster mitigation and risk assessment by effectively predicting the wave amplitude	Formatted
39	and runup.	Formatted
40	<b>Keywords:</b> Rotational Cliff fragmentation; landslide tsunami; prediction models; rotational cliff collapse; wave amplitude, and runup; cliff fragmentation.	Formatted: Normal, Left, Indent: First line: 0.74 cm, Line spacing: 1.5 lines, No bullets or numbering, Tab stops: Not at 5.87 ch
41	prediction model.	Formatted
42	<b>1. Introduction</b>	Formatted: Font: (Default) Times New Roman, 小四, Bold, Not Expanded by / Condensed by
43	The phenomenon of cliff overturning is common along rivers and reservoirs (glacial	Formatted: List Paragraph, Left, Line spacing: 1.5 lines, Outline numbered + Level: 1 + Numbering Style: 1, 2, 3, ... + Start at: 1 + Alignment: Left + Aligned at: 0 cm + Indent at: 0.63 cm, Tab stops: Not at 5.87 ch
44	lakes, recreational lakes), and has been captured by various people around the globe.	Formatted
45	). The cliffs around these lakes are weathered due to climate change and wave action (Ró	Formatted
46	and Cerkowniak, 2024; Young et al., 2021) and can no longer be supported by the parent	Formatted
47	rock. When these initially intact, weathered cliffs fall into water, they usually fragment	Formatted: Font: (Default) Times New Roman, Font color: Black, Not Expanded by / Condensed by
48	upon impact with the water surface, and as a result, induce an impulse water wave.	Formatted
49	Upon impact, the energy of gravitational mass is transferred to the water body, resulting	Formatted
50	in a huge splash and a wave train, propagating away from the point of impact. In the	Formatted
51	reservoirs and water channels located in mountainous regions, such as glacial lakes,	Formatted
52	dams, and a river flowing through valleys, these waves do not travel a long distance	Formatted
53	before reaching obstacles, opposite shores, or other infrastructure. As the waves retain	Formatted
54	most of their energy, size, and strength, the impact can cause significant damage to the	Formatted
55	population and infrastructure located along the banks of the reservoir. Historically,	Formatted
56	extreme impulse wave heights have been observed induced by landslides in events of	Formatted
57	1958 Lituya Bay, USA, which caused a wave height of 524 m (Boulton et al., 2006;	Formatted
58	Franco et al., 2020; Miller, 1960a, 1960), 2007 Chehalis Lake, Canada, induced a wave of 38	Formatted
59	m (Wang et al., 2015), 2015 Taan Fjord, USA, caused a wave of 193 m (Higman et al.,	Formatted
60	2018), and 2014 Lake Askaja (Gylfadóttir et al., 2017). More recently, a volcanic activity in 2018 at Anak Krakatoa, Indonesia, triggered a tsunami reaching up to a height of 13 m, and in 2023, the Dickson Fjord ice-rock avalanche caused a wave.	Formatted

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Moreover, the highly energetic gravity waves are capable of overtopping the dam wall, especially where the freeboard is just a few meters. The overtopping can result in dam failure and can lead to catastrophic events, such as caused by 1963 Vajont rock slide, in North Italy, where a  $250 \text{ Mm}^3$  of rock mass slid into the dam reservoir and induced a huge wave that ran to a height of 200 m at the opposite bank (Franci (Svennevig et al., 2024).

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2020; Heller and Ruffini, 2023; Ward, N. Steve and Day, Simon, 2011; Zhao et al., 2016), the resultant wave overtopped the dam and destroyed an entire village downstream. Similarly, in 2003 Qianjiangping landslide with a volume of  $24 \text{ Mm}^3$ , and in 2008 Gongjiafang landslide with a volume of  $0.38 \text{ Mm}^3$  in Three Gorges dam reservoir area induced a water wave that had an amplitude of 30 m and 32 m respectively (Wang et al., 2021). Gongjiafang landslide induced wave ran up to a height of 12.4 m on opposite bank (Huang et al., 2012).

These incidents highlight the need for predicting the subsequent energy transfer of such cliff collapses for disaster mitigation. The wave amplitude and runup height are of great importance. In contrast, the cases mentioned above are extreme, whereas the small phenomena of sliding,

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toppling, and falling events of cliffs in small lakes and reservoirs can induce a wave of comparably small amplitude but capable of causing substantial damage to densely populated areas along the shoreline. Particularly, in the case of glacial lakes, recreational lakes, and lakes formed by previous landslides are prone to cause major disasters as they are considerably smaller compared to dam reservoirs (Gardezi et al., 2021). The phenomena of cliff overturning and falling are common around these lakes quite frequent and have been captured by various people around the globe. Fig. 1 (a, b, and c) indicates

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a rotational (topple) cliff collapse in Furnas Lake, Brazil, on 8 January 2022, killing 10 people (Maciel et al., 2023; Sun et al., 2024). As a result of the collapse, a huge splash and induced waves can be seen in Fig.

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1 (c). Though scientists have analyzed the amplitude and runup of the waves induced by sliding masses, the literature lacks in providing detailed information. Despite recurring events of cliff collapse along the water banks, the rotational failure of cliffs accompanied by fragmentation upon impact with the water surface remains poorly understood. Recent studies on the

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formation and propagation of the wave induced by rotational cliff collapse. Moreover, the literature also lacks in elaborating water waves have focused on the shape of the induced splash. Effect of cliff

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fragmentation on the induced wave, as can be seen in Fig. 1 (a, b, and c), the falling cliff was still intact and broke under its own weight upon impact with the water surface

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91 and induced a huge splash. Furthermore, though there are numerous prediction models  
92 available for the amplitude and runup of landslide induced water waves, the prediction  
93 models for water waves induced by rotational fall of cliffs considering fragmentation  
94 are nonexistent.

95 Field data related to historical events is critical for disaster mitigation, but due to  
96 their occurrence in remote areas, the unavailability of measuring devices makes it  
97 difficult, leaving the physical modeling as the only source for understanding the wave  
98 generation, and propagation phenomena (Bellotti and Romano, 2017; Grilli et al., 2017;  
99 Takabatake et al., 2022; Wang et al., 2017a; Watts, 1998a). Previously, scientists have  
100 performed both two, and three-dimensional physical modeling for landslide induced  
101 water waves using either block slide (Heinrich, 1992; Heller and Spinneken, 2013;  
102 Najafi-Jilani and Ataie-Ashtiani, 2008; Sælevik et al., 2009), (M. Di Risio et al., 2009;  
103 Marcello Di Risio et al., 2009; Lindström et al., 2014; Montagna et al., 2011; Panizzo  
104 et al., 2005; Wang et al., 2016) block slides, translational slides, or granular slide (Fritz et al.,  
2003a, 2003b; Lindström,

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105 2016; Miller et al., 2016; Zweifel et al., 2006), (flows, where the mass moves along a predefined  
basal plane. Moreover, physical modeling has also been carried out to analyse the amplitude and  
runup of the landslides-induced water waves either by using block slide or granular slide (Heller  
and Spinneken, 2015; Huang, 2013; Lindström, 2016; Lindström et  
106 al., 2014; McFall and Fritz, 2016; Miller, 1960; Mohammed and Fritz, 2012; Montagna et al., 2011;

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Romano et al., 2023, 2020; Sælevik et al., 2009; Yin et al., 2015). However, the rotational collapse is different  
from the type of failure mentioned above. In rotational collapse, the cliff rotates along the base and falls into  
water, thus inducing a huge splash and fragmenting upon impact. Romano et al., 2023),

107 But none have developed a physical model to quantify the amplitude and runup of the  
108 waves induced by rotational cliff collapse, incorporating cliff fragmentation.

109 Along with the physical modeling, the wave generation and propagation  
phenomena have also been analyzed using While a few relevant studies provide partial understandings on  
the phenomena, such as as Liu et al. (2025) numerically analyzed the waves induced by different types of mass  
movements, considering different shapes using smoothed particles, similarly, Heller et al. (2021) experimentally  
analyzed the waves induced by iceburg calving and Yin et al. (2015) studied the potential cliff collapse of  
Jianchuandong rock mass in Three Gorges dam. While these studies contribute to the broader understanding of  
the impact induced by water waves, they do not consider the combined effect of rotational collapse and  
fragmentation. Moreover, the shape of the induced splash, as observed in Lake Furnas, has also not been properly  
explored either experimentally or numerically, as can be seen in Fig. 1 (a, b, and c), the falling cliff was still intact  
and broke under its own weight upon impact with the water surface and induced a huge splash.

110 Parallel advancements in numerical modeling by using have enhanced our understanding of  
landslide induced water waves thorough, computational fluid dynamics (CFD), Eulerian and  
111 Lagrangian methods, employing depth-averaged model, nonlinear shallow water,

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112 — Navier-Stokes model, or Boussinesq equation, for both two- and three-dimensional

113 — modeling (Ceeioni modelling (Franci et al., 2014;2020; Grilli et al., 2019; Guan and Shi, 2023;

114 — Heidarzadeh et al., 2020; Lovholt et

115 — al., 2005; Ruffini et al., 2019; Watts et al., 2003; Whittaker et al., 2017; Yavari and

116 — Ataie-Ashtiani, 2017). Moreover, numerous scientists have also used computational

117 — fluid dynamics (CFD) methods to analyze wave phenomena, just like experimental

118 — modeling, considering the sliding phase as solid material and water as the fluid phase

119 — (Abadie et al., 2010; Chen et al., 2020; Clous and Abadie, 2019; Franci et al., 2020b;

120 — Guan and Shi, 2023; Heller et al., 2016; Kim et al., 2020; Ma et al., 2015; Montagna et

al., 2011; Mulligan et al., 2020; Paris et al., 2021; Rauter et al., 2022; Ruffini et al., 2019). Though these

methods have successfully analysed the wave generation, propagation, and wave dynamics either induced by

granular slide or block slide, their direct application to rotational cliff collapse remains limited. 2022; Romano et

al.,

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2020; Shi et al., 2016).

Furthermore, scientists have also developed empirical and regression-based hybrid prediction models have been developed for

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landslide-induced water waves by considering a combination of several parameters, i.e.,

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geometric, geological, and kinematic characteristics of slides that contribute to wave

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generation parameters, as shown in Table 1. Scientists have M. M. Das and Wiegel (1972) proposed

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that the sliding velocity of the sliding material and water depth are the main components

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factors affecting the wave amplitude of the waves. Watts (1998) emphasised Stated that the role of slope angle, length,

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and mass of the slide are major factors influencing the amplitude of the wave, while Fritz et

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al. (2003) highlighted stated that the role of landslide mass thickness mainly drives governing the amplitude of the

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induced wave. The empirical relations mentioned in Table 1 are mainly for the

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landslide induced water waves, not for cliff collapse induced water waves. Since this

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study is related to the wave induced by rotational cliff collapse, not the granular slide,

the contributing parameters should also be different. Here in this study, we have

considered seven parameters for developing a prediction model, i.e., water depth, fall

height of the cliff, number of fragments, runup slope angle, height of the cliff, and

impact velocity. In this study, we have incorporated a new parameter, i.e., the number

of fragments, as the induced waves from fragments better replicate actual geohazard

events.

Since the experimental and numerical models are expensive, laborious, time-

consuming, and require a lot of expertise, to overcome While these problems, there is a need

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for models that are quick and require less effort and cost. Consequently, the use of AI provide important

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parametric correlations, they were developed for translational or granular slides and show limited application for complex rotational cliff collapses involving fragmentation (Dai et al., 2023; Dignan et al., 2023; Esposti Ongaro et al., 2021).

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and ML based models is gaining fame in the field of engineering. Previous prediction

models for wave amplitude and runup employ simple regression analysis, which is

insufficient for complex problems involving multiple parameters, but recently Recently, scientists

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have shifted towards more advanced ML models to machine learning (ML) approaches for predicting wave dynamics (Bujak et al., 2023; Cesario et al.,

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2024; Li et al., 2024, 2023a2023; Romano et al., 2009; Tarwidi et al., 2023; Tian et al., 2025;

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Wang et al., 2017b2017; Wiguna, 2022). While these prediction models have shown improved

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performance over traditional regression and hybrid models, there is a need to develop an ML. Though scientists have used machine learning for

wave amplitude and runup prediction modeling induced by various types of gravity

flows, the prediction model for the waves induced by the rotational collapse of the cliff

involving fragmentation is nonexistent to the authors' knowledge. Here in this study,

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151 we have developed prediction models for wave amplitude and runup using genetic  
152 programming (GP).  
153 GP-based models have recently gained traction for prediction, and multi-framework for predicting  
amplitude and runup of the waves induced by rotational cliff collapse, considering fragmentation.

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154 expression programming To address this gap, the current study uses Genetic Programming (GP),  
and specifically Multi-Expression Programming (MEP) and genetic expression programming  
(GEP) are the  
155 most advanced, sophisticated, and widely used models. Both models are gene type  
156 programming models that form tree-like models. These models are ), to develop a data-driven  
prediction model for the wave amplitude and runup. This model is similar to living  
157 organisms, which can learn, adapt, and modify their composition, size, and shape  
158 (Gardezi et al., 2024); Usama et al., 2023), MEP is a cutting-edge, advanced form of  
159 GEP GP that adopts a demonstrative model for programming and uses linear chromosomes  
160 to determine optimum population size, mutation probability, and evolutionary model.  
161 Compared to other ML models, it can produce more precise results even when the  
162 problem complexity is unknown (Usama et al., 2023).

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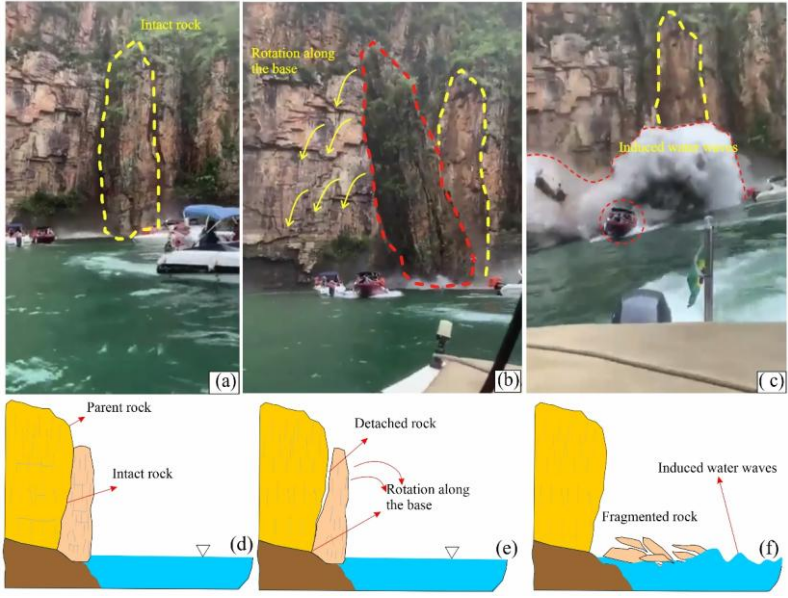


Fig. 1: (a, b, and c), waves induced by a cliff collapse in Lake Furnas, Brazil. (d, e, and f) sketch diagram indicating the detachment and rotational fall process.

Table 1: Historical overview of the prediction models for wave amplitude

Authors	Predictive model
Kamphuis and Bowering	$v_s = 0.7 \cdot l_s$

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(Fritz, 1970)	$A_m = \left( \frac{V_s}{\sqrt{gh}} \right)^{0.31} (0.31 + 0.2 \log \left( \frac{2.2 + 0.35e^{-0.08(x/h)}}{0.31 + 0.2 \log \left( \frac{V_s}{\sqrt{gh}} \right)} \right) + 0.35e^{-0.08(x/h)})$
(Noda et al., 1970)	$A_m = 1.32 \frac{V_s}{\sqrt{gh}}$
(Huber and Hager, 1997)	$\frac{H}{h} = 2 \times 0.88 \sin \theta \cos^2 \left( \frac{2\alpha}{3} \right) \left( \frac{V_s}{\sqrt{gh}} \right)^{0.25} \left( \frac{V}{wh^2} \right)^{0.5} \left( \frac{l}{h} \right)^2$
(Fritz et al., 2004)	$A_m = 0.25 \left( \frac{V_s}{\sqrt{gh}} \right)^{1.4} \left( \frac{l}{h} \right)^{0.8}$
(Panizzo et al., 2005)	$\frac{Hm}{h} = 0.07 \left( \frac{T_s h^2}{ws} \right)^{-0.45} (\sin \alpha)^{-0.88} e^{0.6 \cos \theta} \left( \frac{l}{h} \right)^{-0.44}$
(Heller, 2007)	$A_m = \frac{4}{9} \left[ F \left( \frac{s}{h} \right)^{1/2} \rho^{1/4} \left( \cos \frac{6\alpha}{7} \right)^{2/5} \right]$ $A_m = \max(A_{c1}, A_{c2})$
(Mohammed and Fritz, 2012)	$A_{c1} = 0.3 F^{2.1} \left( \frac{s}{h} \right)^{0.6} \left( \frac{l}{h} \right)^{-0.45} \left( \frac{V}{wh^2} \right)^{-0.02} \left( \frac{V}{wh^2} \right)^{-0.33/h}$ $A_{c2} = 1.0 F S^{0.8} \left( \frac{w}{h} \right)^{-0.4} \left( \frac{l}{h} \right)^{0.25} \left( \frac{V}{wh^2} \right)^{0.45} \left( \frac{V}{wh^2} \right)^{-0.07} \left( \frac{V}{wh^2} \right)^{-0.3} \cos^2 \alpha$
(Wang et al., 2016)	$A_m = 1.17 F \left( \frac{s}{bh} \right)^{0.25} \left( \frac{w}{h} \right)^{0.45} (\sin^2 \alpha + 0.6 \cos^2 \alpha)$
(Li et al., 2023)	$A_m = 0.59 \sqrt{\frac{2H(1-f \cot \alpha)}{h}} \left( \frac{swl}{h^3} \right)^{N-0.11} \left( \frac{l}{h} \right)^{-0.43} \cos^2 \left( \frac{2}{3} \alpha \right)$

Note:  $l$  is the landslide length;  $s$  is the landslide thickness;  $w$  is the landslide width;  $m$  is the landslide mass weight;  $V$  is the landslide volume;  $H$  is the landslide height;  $T_s$  time for motion of slide;  $b$  is the river width;  $h$  is the still water depth;  $x(r)$  is the offshore distance from the bank slope;  $\alpha$  is the slope angle;  $\theta$  is the angular direction;  $V_s$  is the impact velocity.

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Fritz et al. (2004)	$A_m = 0.25 \left( \frac{V_s}{\sqrt{gh}} \right)^{1.4} \left( \frac{l}{h} \right)^{0.8}$
Panizzo et al. (2005)	$\frac{Hm}{h} = 0.07 \left( \frac{T_s h^2}{ws} \right)^{-0.45} (\sin \alpha)^{-0.88} e^{0.6 \cos \theta} \left( \frac{l}{h} \right)^{-0.44}$
Heller (2007)	$A_m = \frac{4}{9} \left[ F \left( \frac{s}{h} \right)^{1/2} \rho^{1/4} \left( \cos \frac{6\alpha}{7} \right)^{2/5} \right]$ $A_m = \max(A_{c1}, A_{c2})$
Mohammed and Fritz (2012)	$A_{c1} = 0.3 F^{2.1} \left( \frac{s}{h} \right)^{0.6} \left( \frac{l}{h} \right)^{-0.45} \left( \frac{V}{wh^2} \right)^{-0.02} \left( \frac{V}{wh^2} \right)^{-0.33/h}$ $A_{c2} = 1.0 F S^{0.8} \left( \frac{w}{h} \right)^{-0.4} \left( \frac{l}{h} \right)^{0.25} \left( \frac{V}{wh^2} \right)^{0.45} \left( \frac{V}{wh^2} \right)^{-0.07} \left( \frac{V}{wh^2} \right)^{-0.3} \cos^2 \alpha$
Wang et al. (2016)	$A_m = 1.17 F \left( \frac{s}{bh} \right)^{0.25} \left( \frac{w}{h} \right)^{0.45} (\sin^2 \alpha + 0.6 \cos^2 \alpha)$
Li et al. (2023)	$A_m = 0.59 \sqrt{\frac{2H(1-f \cot \alpha)}{h}} \left( \frac{swl}{h^3} \right)^{N-0.11} \left( \frac{l}{h} \right)^{-0.43} \cos^2 \left( \frac{2}{3} \alpha \right)$

Note: Note:  $l$  is the landslide length;  $s$  is the landslide thickness;  $w$  is the landslide width;  $m$  is the landslide mass weight;  $V$  is the landslide volume;  $H$  is the landslide height;  $b$  is the river width;  $h$  is the still water depth;  $x(r)$  is the offshore distance from the bank slope;  $\alpha$  is the slope angle;  $\theta$  is the angular direction;  $V_s$  is the impact velocity.

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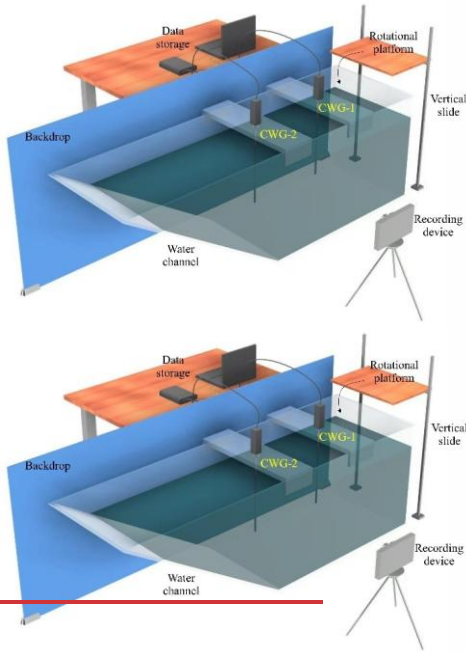
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collapse was carried out in a three-dimensional water flume made up of plexiglass, as shown in Figs. 2 and 3. One end of the flume is vertical at 90°, whereas the other end is inclined and fixed at 30° (Fig. 3a and b). The flume is 0.55 m high, 0.5 m wide, and 1.4 m long along the base and 2.35 m long at the top. Furthermore, to measure the runup of induced water waves at various slope angles, two sliding rails were installed towards the inclined end at 45° and 60°. So, upon insertion of the gate at 45° and 60°, the top length of the flume was further reduced according to the Pythagoras theorem. To induce the rotational cliff collapse, a 0.55 m wide and 0.6 m high movable platform was prepared designed, which can move in the vertical direction and can also rotate about its axis. The rotational motion was induced by pulling the hinge; the release ensured a pure rotational motion, which was visually verified by video analysis. The flume was marked with a vertical scale to measure the water depth. The wave amplitude was measured using capacitance-type wave gauges with an accuracy of ± 0.5 mm, placed along the centerline at specified intervals. The runup height was measured using a graduated paper attached to the inclined surface. The entire process was recorded using a digital camera (240 fps, 720p resolution) placed perpendicular to the experimental flume; the velocity of the falling cliff was verified by frame-by-frame video analysis using Particle Image Velocimeter (PIV).

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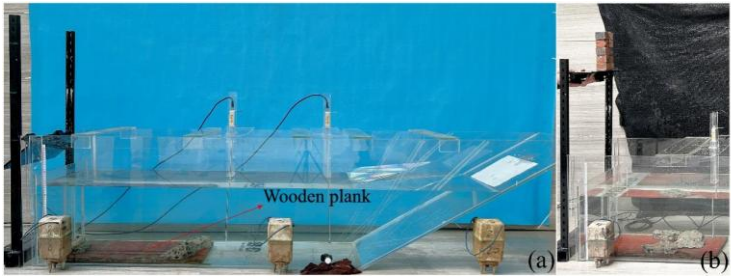
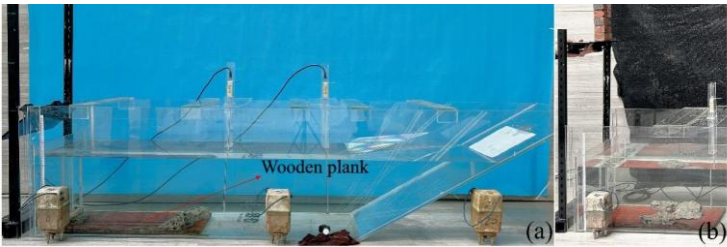
208 Fig. 2: Illustration of experimental setup including wave gauges, rotational platform,  
209 recording, and data storage devices.

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211 Fig. 3: Photographs of the setup, (a) Experimental flume, (b) platform for inducing  
212 rotational cliff collapse.

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213 2.2 Test preparation and materials

214 Physical experiments were carried out by varying the water depth, fall height,

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215 number of fragments, bank slope angle, mass of falling rock, cliff height, and impact

216 velocity. The tests were carried out for three water depths, i.e., 0.34 m, 0.27 m, and 0.20

217 m, and three fall heights, i.e., 0.64 m, 0.44 m, and 0.245 m from the surface of the water

218 level. Furthermore, the number of blocks was also varied, i.e., 6, 10, and 12 blocks

219 having combined weights of 1.445 kg, 2.29 kg, and 2.82 kg, respectively. At the same

220 time, the impact velocity changed by changing the fall height. The wave runup was

221 measured by varying the bank slope angle, i.e., 30°, 45° and 60°.

222 To replicate the field density of the rocks, red gutka bricks having a density of

223 around 2000 kg/m<sup>3</sup> were used. The singleA singular block had a dimension of 0.55055x0.05x0.042

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224 m. m. The mass volume and dimension of all the blocks were unchanged to ensure consistency in the experiments. A combination of 6, 10, and 12 blocks of red gutka bricks were used to form a cliff	Formatted
225 and measure the wave amplitude and runup of induced waves. The blocks were joined	Formatted
226 together with the help of cement paste having a water-cement ratio W/C 0.8 and cured	Formatted
227 for 2 hours in front of an electric heater at 150 °C. To ensure the weak bond strength, several trials for bond strength were carried out after a curing period of 2 hours, and it was found to be in the range of 0.42-0.5 MPa. In contrast, the inertial stresses at the time of impact were several times higher, such that the bond is weak enough they caused the fragmentation of the cliff. This condition was purposely designed to imitate naturally fractured cliff materials, confirming that #	Formatted
228 fragments at the structure fragmented primarily along the joints upon impacting impact with the water surface, consistent with field observations of rotational cliff collapses. The bonded blocks were	Formatted
229 placed on the rotational platform at specific heights, i.e., 0.64 m, 0.44 m, and 0.245 m	Formatted
230 from the water level, and were allowed to rotate under their own weight by pulling the	Formatted
231 hinge, such that the placed block falls in the water following rotation motion along its	Formatted
232 base Fig. 3 (b). To avoid the slippage of blocks and to replicate field conditions, fine-	Formatted: Normal, Indent: First line: 0.74 cm, Don't adjust right indent when grid is defined, Space Before: 0 pt, Line spacing: 1.5 lines, No bullets or numbering, Don't adjust space between Latin and Asian text, Don't adjust space between Asian text and numbers, Tab stops: Not at 5.87 ch
233 grained ensure that it had sufficient frictional resistance needed for pure rotational motion of the	Formatted
simulated cliff, finely-grounded bricks of the same cliff material as the cliff were pasted on the rotational platform,	Formatted
thereby preventing translational motion or vertical free fall into the water.	Formatted
234 Furthermore, to reduce the impact of falling blocks on the base of the flume, a wooden	Formatted
235 plank weighing 2.69 kg and dimensions 0.65mx0.37mx0.01 m was placed at the	Formatted
236 point of impact inside the flume. Due to its large surface area and lighter density, it	Formatted: Normal, Indent: First line: 0.74 cm, Don't adjust right indent when grid is defined, Space Before: 0 pt, Line spacing: 1.5 lines, No bullets or numbering, Don't adjust space between Latin and Asian text, Don't adjust space between Asian text and numbers, Tab stops: Not at 5.87 ch
237 tends to float in the flume, so two blocks of concrete weighing 3.58 kg were placed on	Formatted
238 it, Fig. 3 (a). Since the fall height was small, no considerable local breakage was observed in the blocks,	Formatted
and the brief water contact minimised the water absorption effect.	Formatted: Font: (Default) Times New Roman, (Asian) +Body Asian (等线)
239 The induced wave amplitude was measured by placing the wave measuring gauges	Formatted
240 at a distance of 0.65 m and 0.135 m from the vertical face; the gauges were wired and	Formatted
241 connected to the laptop. At the same time, the runup was measured manually with the	Formatted
242 help of a scale by pasting a scaled paper on the slope. Furthermore, the experiments	Formatted
243 were also recorded with the help of a high-resolution camera for verification purposes.	Formatted: Normal, Indent: First line: 0.74 cm, Don't adjust right indent when grid is defined, Line spacing: 1.5 lines, No bullets or numbering, Don't adjust space between Latin and Asian text, Don't adjust space between Asian text and numbers, Tab stops: Not at 5.87 ch
244 2.3 Numerical Modeling	Formatted
245 Simulating multi-phase flows is challenging due to the constant deformation of the	Formatted
246 liquid-gas interface. Various numerical methods have been developed to model these	Formatted
247 flows, each offering unique advantages depending on the specific flow regime and	Formatted
248 characteristics of interest. In this study, the Volume of Fluid (VOF) method is utilized	Formatted
249 for its effectiveness in handling significant interface distortions and topological changes.	Formatted: Font: (Default) Times New Roman, Not Expanded by / Condensed by
250 The VOF method offers superior mass conservation, which is critical in high velocity	Formatted

252—(Brackbill (Backbill, et al., 1992; Hirt and Nichols, 1981). Alternative numerical schemes, such as  
the Front Tracking approach, are generally limited). Other approaches provide superior

accuracy in modeling interfaces and surface tension, but they struggle to manage handling complex scenarios (topological changes (Tryggvason et al., 2001; Liu and Liu, 2010; Monaghan, 1994; Yang and Kong, 2018)). Another approach is the Level Set method, but it suffers from mass conservation and convergence issues. The Lattice Boltzmann Method (LBM) is also common; however, its applicability to high velocity impact is rather limited (Aidun & Clausen, 2010). Given these trade-offs, the Volume of Fluid (VOF) method finds an optimal balance of computational efficiency, interface tracking capability, and proven reliability for modeling multiphase flow in the moderate-to-high velocity range relevant to this study. Therefore, a two-dimensional numerical model of a cliff, having the same properties as the experimental cliff mentioned previously, hitting the water surface is developed using the VOF method to accurately capture the liquid-gas interface.

In this approach, a volume fraction ( $\alpha$ ), ranging between 0 and 1, is applied across the entire computational domain. A value of  $\alpha=1$  indicates a control volume filled with liquid, while  $\alpha=0$  denotes a control volume filled with gas. The interface is represented by values where  $0 < \alpha < 1$ . In the Volume of Fluid (VOF) method, the momentum equation is solved across the entire computational domain, with the resulting velocity field shared by all phases. To account for surface tension effects, a continuum surface force (CSF) model is employed (Backbill et al., 1992). The normal vector  $n$  and interface mean curvature  $k$  are as follows, respectively:

interface mean curvature  $\kappa$  are as follows, respectively:

$$n = \frac{\nabla \alpha}{|\nabla \alpha|} \quad (1)$$

and

$$\kappa = \nabla \cdot \frac{\nabla \alpha}{|\nabla \alpha|} \quad (2)$$

The interface is maintained as sharp through the use of geometric reconstruction to ensure its clarity. The volume fraction ( $\alpha$ ) is discretized with the geo-reconstruct scheme, while the convective terms in the momentum equation are handled using a second-order upwind method. The PISO (Pressure-Implicit with Splitting of Operators) algorithm was employed for pressure-velocity coupling, which is well-suited for transient flows. Temporal discretization employs a second-order implicit scheme, and spatial gradients are calculated using the Least Squares Cell-Based method.

The boundary conditions were defined as follows: the bottom boundary was modeled as a no-slip wall, while the top boundary was set as a pressure outlet at atmospheric conditions, and the lateral sides were modeled as stationary walls to

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main. For accurate simulation of the rotational motion of the cliff through the air-water interface in a multi-phase flow environment, dynamic meshing was implemented within the ANSYS Fluent framework. This approach facilitated the adaptation of the

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282	computational mesh to accommodate the cliff's movement while maintaining the	Formatted
283	integrity of the liquid-gas interface captured by the Volume of Fluid (VOF) method.	Formatted
284	Dynamic meshing was critical for modeling the complex interactions between the	Formatted
285	falling cliff and the surrounding air and water phases, allowing the mesh to deform and	Formatted
286	adapt in response to the cliff's trajectory. In ANSYS Fluent, the dynamic meshing	Formatted
287	strategy employed a combination of mesh deformation and local remeshing techniques	Formatted
288	to handle the cliff's motion. Mesh deformation was applied to adjust the existing mesh	Formatted
289	nodes smoothly as the cliff moved, preserving mesh quality in regions experiencing	Formatted
290	moderate displacement. For areas near the cliff where significant deformation could	Formatted
291	lead to poor mesh quality, local remeshing was utilized to regenerate mesh elements for	Formatted
292	better numerical stability and accuracy. The smoothing and remeshing algorithms were	Formatted
293	configured to maintain high mesh quality, with a skewness threshold set to prevent	Formatted
294	excessive element distortion.	Formatted
295	The rotational cliff collapse was simulated using an in-house user-defined function	Formatted: Normal, Indent: First line: 0.74 cm, Line spacing: 1.5 lines, No bullets or numbering, Tab stops: Not at 5.87 ch
296	(UDF). This UDF interfaced with ANSYS Fluent to dynamically update the rock's	Formatted
297	position and velocity. To enhance computational efficiency, a dynamic mesh zone was	Formatted
298	defined around the cliff, with a finer mesh resolution near its surface to capture the	Formatted
299	sharp gradients in the flow field and interface dynamics. The mesh was gradually	Formatted
300	coarsened away from the rock to reduce computational cost while maintaining	Formatted
301	sufficient resolution in the far-field regions. The dynamic meshing process was	Formatted
302	synchronized synchronised with the transient flow solver, using a time step size determined through	Formatted
303	a time step independence study to balance accuracy and computational efficiency. It is also worth	Formatted
	mentioning that the numerical simulations were performed considering the rock as a unified mass. This approach	Formatted: Normal, Indent: First line: 0.74 cm, Space Before: 0 pt, Line spacing: 1.5 lines, No bullets or numbering, Tab stops: Not at 5.87 ch
	describes the slight differences between the experimental and numerical results, which are nonetheless within the	
	acceptable range.	
304	<b>2.4 Multi-expression programming</b>	Formatted
305	The MEP model was developed for predicting wave amplitude and runup using	Formatted
306	experimental data, as shown in Table 2. A dataset of 81 experiments was prepared by	Formatted
307	alternately varying seven different parameters, and the results for wave amplitude and	
308	runup were recorded. experimental results was used as an input to a machine learning model.	Formatted
	Furthermore, the data was divided into 70/30 ratios for training	Formatted
309	and validation purposes before developing the model. The model starts working by	Formatted
310	generating a random chromosome population, and it continues to generate the	Formatted
311	chromosomes until a terminal condition is achieved, generating an optimal expression	Formatted

from the data having input and output pairs over a certain number of generations, as shown in Fig. 4.

Based on a binary tournament process, parents are selected and then undergo a recombination process through a consistent crossover probability. This recombination produces two more offspring. These offspring go through mutation, and if these offspring perform better than the least fitting offspring in the current population, then the better offspring replace them. The illustrations used by MEP are similar to the ones used by C++ and Pascal compilers. The MEP chromosomes are comprised of numerous genes combined using various mathematical operators such as addition (+), subtraction (-), multiplication (x), and division (/), and these genes create expression trees (ETs) (Cheng et al., 2020). Moreover, there are several parameters/hyperparameters, such as code length, sub-population size and number, crossover probability, and other sets of various functions involved in the generation of MEP code, and they also govern the overall performance of the code. Among these parameters, the size of the population tells us about the number of programs being generated, whereas an increase or decrease in subpopulation size directly affects the complexity and computation time of the model. Moreover, the length of the developed model is controlled by varying the code length parameter. During model

development, prerequisite tuning procedures were applied to optimize these hyperparameters. This careful selection minimized the risk of premature convergence or underfitting while ensuring computational efficiency. Table 2: Input parameters Experimental dataset used for training and corresponding wave amplitude and runup height validation of the machine learning model.

S/No.	Water depth ( $d$ ) (m)	Drop/Fall height ( $H$ ) (m)	Fragments ( $N$ )	Angle ( $\alpha$ ) ( $^\circ$ )	Cliff Mass( $m$ ) Mass( $m$ ) (kg)	Cliff height( $h$ ) (m)	Velocity ( $v$ ) (m/s)	Amplitude ( $A$ ) (m)	Runup ( $R$ ) (m)
1	0.34	0.245	6	30	1.445	0.12	2.19	0.0225	0.051
2	0.34	0.445	6	30	1.445	0.12	2.95	0.0230	0.056
3	0.34	0.645	6	30	1.445	0.12	3.56	0.0365	0.068
4	0.34	0.245	6	45	1.445	0.12	2.19	0.0370	0.045
5	0.34	0.445	6	45	1.445	0.12	2.95	0.0425	0.051
6	0.34	0.645	6	45	1.445	0.12	3.56	0.0425	0.051
7	0.34	0.245	6	30	1.445	0.12	2.19	0.0225	0.051
8	0.34	0.445	6	30	1.445	0.12	2.95	0.0230	0.056
9	0.34	0.645	6	30	1.445	0.12	3.56	0.0365	0.068
10	0.34	0.245	6	45	1.445	0.12	2.19	0.0370	0.045
11	0.34	0.445	6	45	1.445	0.12	2.95	0.0425	0.051
12	0.34	0.645	6	45	1.445	0.12	3.56	0.0425	0.051
13	0.34	0.245	6	30	1.445	0.12	2.19	0.0225	0.051
14	0.34	0.445	6	30	1.445	0.12	2.95	0.0230	0.056
15	0.34	0.645	6	30	1.445	0.12	3.56	0.0365	0.068
16	0.34	0.245	6	45	1.445	0.12	2.19	0.0370	0.045
17	0.34	0.445	6	45	1.445	0.12	2.95	0.0425	0.051
18	0.34	0.645	6	45	1.445	0.12	3.56	0.0425	0.051
19	0.34	0.245	6	30	1.445	0.12	2.19	0.0225	0.051
20	0.34	0.445	6	30	1.445	0.12	2.95	0.0230	0.056
21	0.34	0.645	6	30	1.445	0.12	3.56	0.0365	0.068
22	0.34	0.245	6	45	1.445	0.12	2.19	0.0370	0.045
23	0.34	0.445	6	45	1.445	0.12	2.95	0.0425	0.051
24	0.34	0.645	6	45	1.445	0.12	3.56	0.0425	0.051
25	0.34	0.245	6	30	1.445	0.12	2.19	0.0225	0.051
26	0.34	0.445	6	30	1.445	0.12	2.95	0.0230	0.056
27	0.34	0.645	6	30	1.445	0.12	3.56	0.0365	0.068
28	0.34	0.245	6	45	1.445	0.12	2.19	0.0370	0.045
29	0.34	0.445	6	45	1.445	0.12	2.95	0.0425	0.051
30	0.34	0.645	6	45	1.445	0.12	3.56	0.0425	0.051
31	0.34	0.245	6	30	1.445	0.12	2.19	0.0225	0.051
32	0.34	0.445	6	30	1.445	0.12	2.95	0.0230	0.056
33	0.34	0.645	6	30	1.445	0.12	3.56	0.0365	0.068
34	0.34	0.245	6	45	1.445	0.12	2.19	0.0370	0.045
35	0.34	0.445	6	45	1.445	0.12	2.95	0.0425	0.051
36	0.34	0.645	6	45	1.445	0.12	3.56	0.0425	0.051
37	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
38	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
39	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
40	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
41	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
42	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
43	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
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45	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
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49	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
50	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
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52	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
53	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
54	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
55	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
56	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
57	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
58	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
59	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
60	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
61	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
62	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
63	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
64	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
65	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
66	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
67	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
68	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
69	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
70	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
71	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
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73	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
74	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
75	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
76	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
77	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
78	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
79	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
80	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
81	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
82	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
83	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
84	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
85	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
86	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
87	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
88	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
89	0.27	0.445	10	45	2.295	0.20	2.95	0.0523	0.102
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92	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
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94	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085
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96	0.27	0.645	10	45	2.295	0.20	3.56	0.0523	0.102
97	0.27	0.245	10	30	2.295	0.20	2.19	0.0431	0.116
98	0.27	0.445	10	30	2.295	0.20	2.95	0.0510	0.129
99	0.27	0.645	10	30	2.295	0.20	3.56	0.0685	0.141
100	0.27	0.245	10	45	2.295	0.20	2.19	0.0390	0.085

78	0.2	0.645	12	45	2.82	0.24	3.56	0.0733	0.146
79	0.2	0.245	12	60	2.82	0.24	2.19	0.0565	0.062
80	0.2	0.445	12	60	2.82	0.24	2.95	0.0636	0.083
81	0.2	0.645	12	60	2.82	0.24	3.56	0.0657	0.098

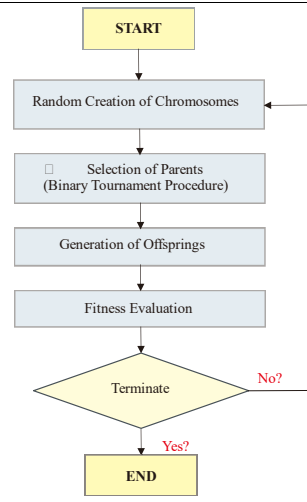


Fig.

### 3.1 Experimental results

4 Schematic representation of  
MEP workflow used in this  
study

## 3. Results and discussions

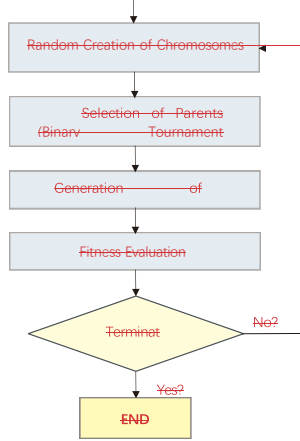


Fig. 4 MEP flowchart

334 The experimental results of the wave amplitude and runup, induced by rotational  
 335 cliff collapse, reveal complex hydrodynamic processes. As shown in Fig. 5, the failure  
 336 is initiated by the rotational fall of the cliff, leading to a significant amount of impact  
 337 energy upon hitting the water surface. The impact induced a huge splash, which is  
 338 evident from Fig. 5 (b, e & h). It was observed that the shape of the splash also varies  
 339 with water depth for all the cases; higher water depths resulted in a mushroom-shaped  
 340 splash, i.e., broader on the top, as the momentum dissipates before interacting with the bottom  
 341 surface, resulting in a vertical jet and the formation of a mushroom-shaped splash. as can be seen in  
 342 Fig. 5(h). The observed phenomena  
 343 perfectly align with the basic concepts of fluid dynamics, which state that greater depths  
 344 absorb more impact energy compared to shallow waters. Shallow waters produced a  
 345 vertically elongated splash as can be seen in Fig. 5 (b & e). It can be observed that as  
 the depth decreases, the splash becomes more elongated, and this is due to the fact that  
 as shallower depths intensify the upward momentum transfer, thus resulting in a more

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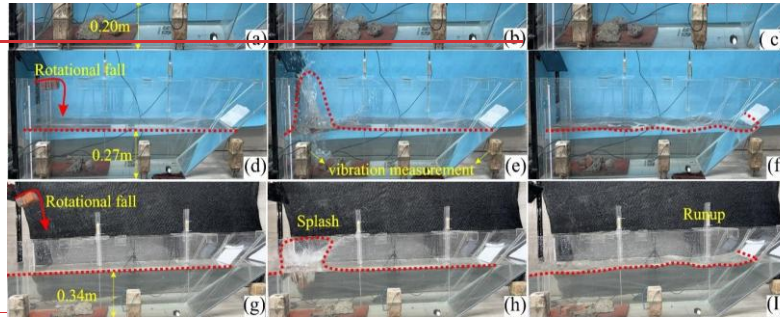
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346 elongated shape (Kubota and Mochizuki, 2009).

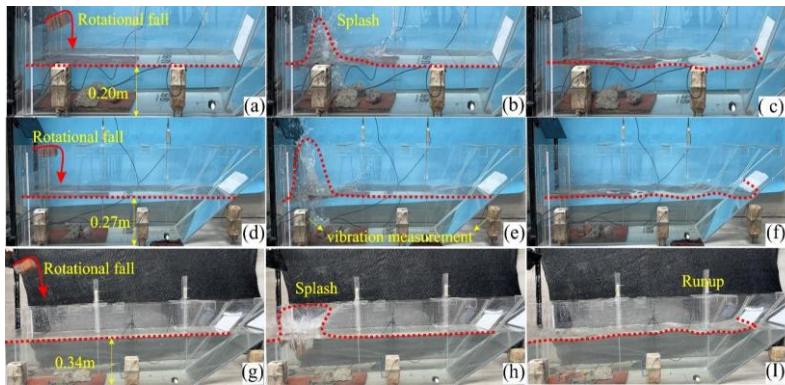
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347



348 Fig. 5: A pictorial display of the experimental setup for various water depths, i.e. 0.20  
349 m, 0.27m, and 0.34 m. (a, d & g) indicate rotational fall of the cliff, (b, e & h)  
350 showing splash as a result of cliff impact, (c, f & i) formation and propagation of  
351 induced wave and runup at various slope angles.

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### 352 3.1.1 Relation between energy and amplitude

353 Further, the relationship between impact energy and wave response was also  
354 investigated by establishing a dimensionless impact energy parameter ( $K.E/\rho gh^3$ ).  
355 Where  $K.E$  is the kinetic energy of the cliff,  $\rho$  is the density, and  $h$  is the water depth.  
356 The negative quadratic coefficient in Fig. 6(a) indicates a nonlinear response, such that  
357 at the start, the wave amplitude increases as the impact energy increases, but later it  
358 decreases due to reduced energy transfer at higher impact values. At higher impact values, the released energy  
was not fully used in the wave formation and propagation; instead, a part of the energy was dissipated in the  
formation of splash, and in the formation of air pockets and their subsequent collapse. Moreover, the  
359 coefficient of determination was found to be 77% indicating a good data fit.

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360 Moreover, the results for the relative wave amplitude and wave energy were  
361 analyzed for three water depths, i.e. 0.34 m, 0.27 m, and 0.20 m, as shown in Fig. 6(b).

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362	The results indicate a strong correlation for all three cases, with coefficients of	Formatted
363	determination around 0.96. The results indicate a direct relation between wave height	Formatted
364	and energy, whereas the decreasing slope values with the increasing water depth	Formatted
365	suggest that for deeper water the wave amplitude decreases at a slower rate with	Formatted
366	increasing wave energy, thus highlighting the impact of water depth on the wave	Formatted
367	dynamics, such that shallower water allows more amplification of waves for the same	Formatted

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energy level, and this is due to the increased non-linear interactions and enhanced energy concentrations in shallower depths (Myrhaug and Lader, 2019).

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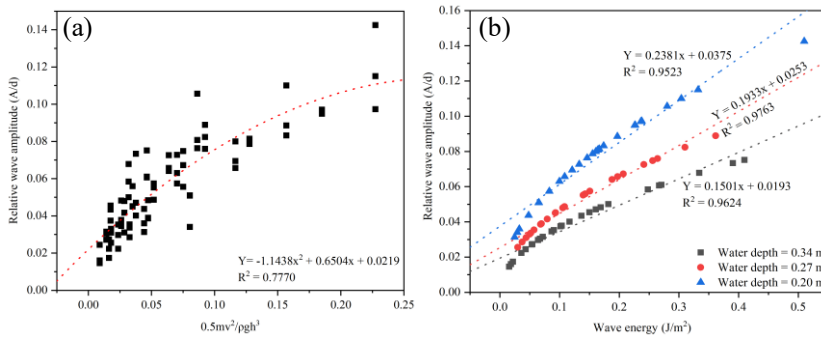
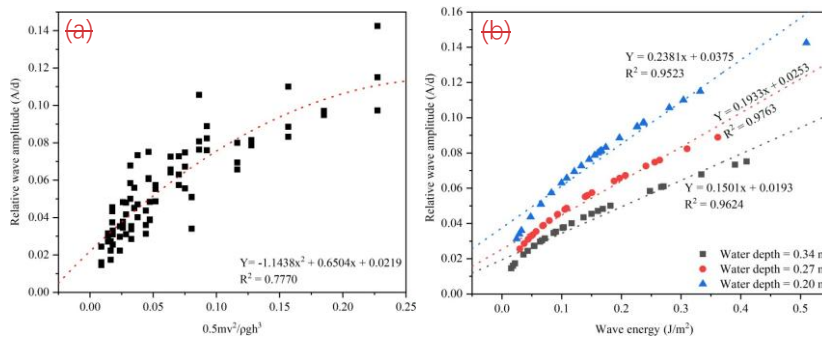
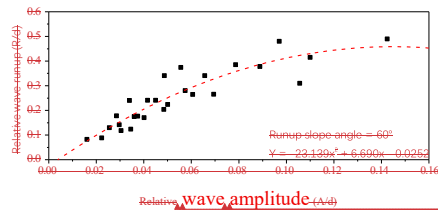


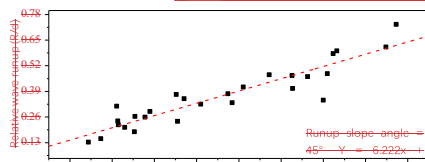
Fig. 6: (a) Impact Dimensionless impact energy ( $K.E/pgh^3$ ) vs relative wave amplitude, indicating a nonlinear trend, (b) Wave energy vs relative

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wave amplitude



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Relative, indicating higher wave amplitude (A/d) amplifications in shallow waters.

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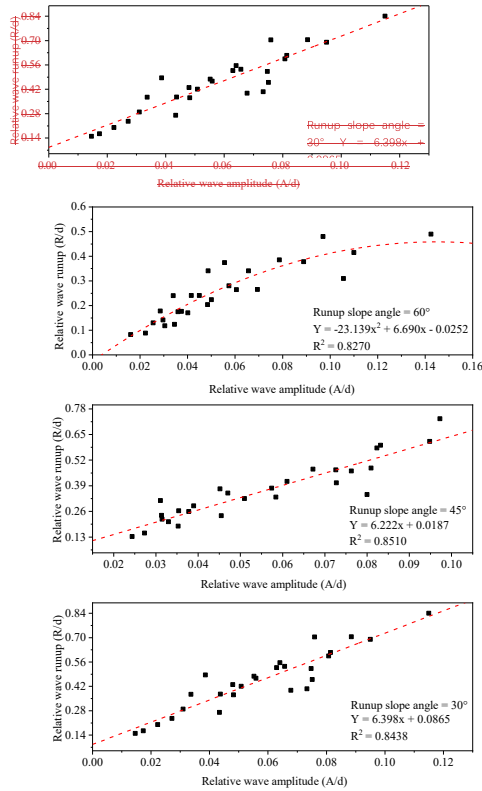


Fig. 7; Relative wave amplitude vs relative wave height runup at various slope angles and water depth.

The results for the relative wave height and runup for all three water depths and three runup slope angles are shown in Fig. 7. The relationship indicates a strong correlation between wave amplitude and runup for all three slope angles. The decreasing line-slope values with increasing runup slope angle indicate that wave runup increases at a slower rate for sharp slope angles compared to mild slopes. The trend highlights the effect of slope angle on the runup. The result also indicates that the mild slope angles help wave runup amplification, as they dissipate a very small amount of energy, whereas steeper angles result in lower runup heights because of higher energy losses (Wu et al., 2018).

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decreasing line slope values with increasing runup slope angle indicate that wave runup increases at a slower rate for sharp slope angles compared to mild slopes. The trend highlights the effect of slope angle on the runup. The result also indicates that the mild slope angles help wave runup amplification, as they dissipate a very small amount of energy, whereas steeper angles result in lower runup heights because of higher energy losses (Wu et al., 2018).

### 3.1.2 Impact Froude no. vs Relative wave amplitude

Fig. 8 indicates the relationship between the impact Froude number and relative wave amplitude ( $A/d$ ), under varying experimental conditions for the first gauge, i.e., near the impact zone. Since we are interested in the immediate response of the wave influenced by the impact Froude number. The results indicate that as the water depth decreases, the relative wave amplitude and impact Froude number increase, indicating a reduction in the dissipation of impact energy, causing pronounced surface turbulence and increased wave height. Additionally, the decreased water depth also increased the value of the impact Froude number by reducing its characteristic velocity, resulting in stronger wave generation upon impact. The calculations for Reynolds number for the experiments resulted in very high values, thus indicating a strong turbulent flow, which is also evident from Fig. 5, so viscous effects are very, very small and can be ignored, thus indicating the Froude dynamics similarity. The experimental results indicate the complex interaction between wave propagation, impact dynamics, and bathymetrical effects in waves induced by rotational cliff collapse. Moreover, upon impact, the cliff fragmentation distributes impact energy over a larger area of water, thus increasing wave height by enhanced turbulence and water splashing effects.

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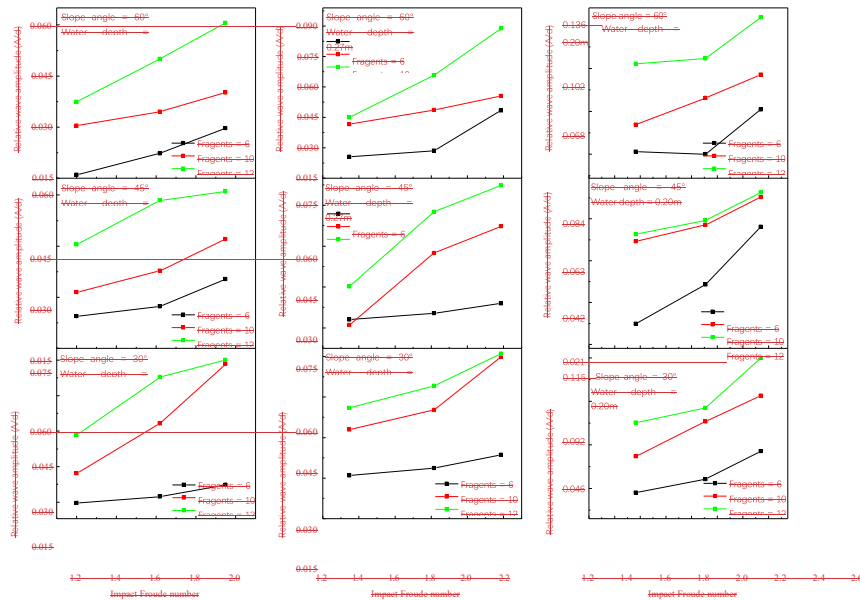
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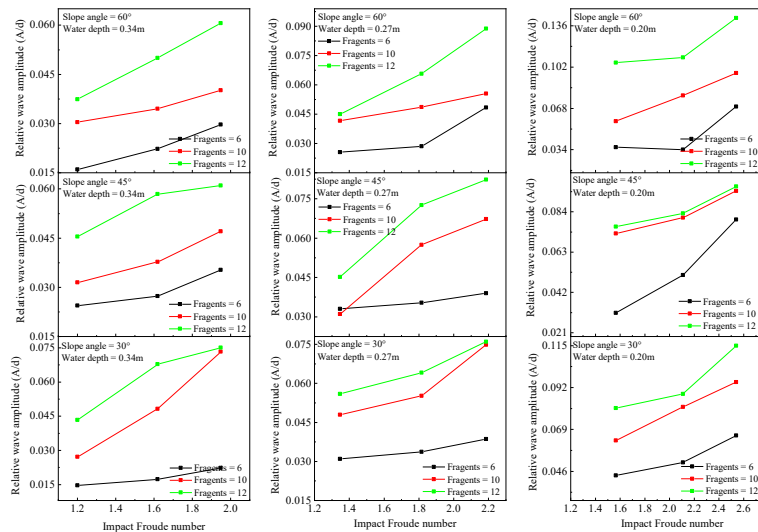


Fig. 8: Relationship between impact Froude number and relative wave height amplitude.

### 3.1.3 Wave amplitude results

The results for the wave amplitude for various parameters are shown in Figs. 9, 10,

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405 and 11. As mentioned earlier, two gauges were used to measure the induced wave

406 amplitude. Fig. 9 provides a detailed comparison of the wave amplitude recorded at

407 both gauges for a 60° runup slope angle and a 0.445m445 m fall height. It can be observed that

408 gauge-1, which is near to impact zone, has a higher relative amplitude compared to

409 gauge-2. Furthermore, the results for the relative wave amplitude against the

410 normalized time were also analyzedanalysed for all the water depths (0.20m20 m, 0.27m27 m, and

411 0.34m), fall height (0.245m245 m, 0.445m445 m, and 0.645m645 m), and cliff height (0.12m12 m,

412 0.20m20 m, and 0.24m).

413 24 m). The results indicate that the wave amplitude increases as the cliff height, impact

414 velocity, and number of fragments increase for all the water depths, as can be observed

415 in Fig. 10, thus demonstrating that the potential energy of the falling cliff plays a critical

role in the magnitude of the resulting wave.

416 Interestingly, comparing the wave amplitude induced by cliffs of various heights

417 falling from the same height revealed that the water depth and the wave have an inverse

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relationship. As shown in Fig. 10 (a, b, and c), the average wave amplitude for various cliff heights and the same fall height of 0.245 m at 0.20 m water depth is 26% more than the average wave amplitude induced by 0.27 m water depth and 50% more than the 0.34 m water depth wave amplitude. Similarly, Fig. 10 (d, e, and f) indicates that the average wave amplitude for 0.445 m fall height at 0.20 m water depth is 18% more than 0.027 m and 47% more than 0.34 m water depth, whereas, for 0.645 m fall height wave amplitude induced by 0.20 m water depth is 25% more than 0.27 m and 37% more than 0.34 m water depth (Fig. 10 g, h & i), thus suggesting that the deeper water dissipates the impact energy more effectively, as the deep water have more mass available to absorb and redistribute the impact energy, compared to shallower water thus reducing the overall amplitude of the induced wave. Moreover, a similar trend was observed for the wave amplitude involving 45° and 60° runup slope angle.

Later on, we performed another experiment by using granular material of equivalent mass as of cliff and slid it on a 30° slope, for all the water depths, and amplitude of the induced wave was measured as shown in Fig. 11. Fig. 11(a) indicates that the wave amplitude for 0.20 m water depth and 1.445 kg granular mass (equivalent to 0.12 m cliff height) was 15% more than 0.27 m water depth and 65% more than wave amplitude induced by 0.34 m water depth. Whereas for 2.29 kg and 2.82 kg granular mass equivalent to 0.20 m and 0.24 m cliff height similar trend was observed as shown in Fig. 11 (b and c), thus indicating that as the water depth increases, the wave amplitude decreases for all the equivalent granular masses, as happened in the case of cliff fall.

Furthermore, a comparison between the wave amplitude induced by a falling cliff and equivalent granular mass at various water depths indicates that the amplitude of the wave induced by an equivalent granular mass in 0.34 m, 0.27 m, and 0.20 m water depth was on average 28%, 35% and 42% less than the wave amplitude induced falling cliff. The substantial difference in wave amplitude highlights the importance of energy transfer in wave formation. The falling cliff following a rotational motion imparts a more sudden and concentrated impact that allows an efficient energy transfer to water, leading to higher wave amplitudes. On the other hand, granular flows, being more deformable and flowing along a slope, result in gradual energy transfer over a wide area,

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448 thus resulting in lower wave amplitudes. The results highlight that it's not only the total  
449 impact energy that affects the behavior of the induced wave, but the mode of energy  
450 transfer also plays a critical role, (Mohammed and Fritz, 2012; Wunnemann and Weiss,  
451 2015). Based on the experimental results for wave amplitude and runup induced by  
452 rotational cliff collapse that fragments upon impact with the water surface, a novel  
453 prediction model was prepared using multi-expression programming. The justifications  
454 for the use of MEP have been well explained in the previous sections.

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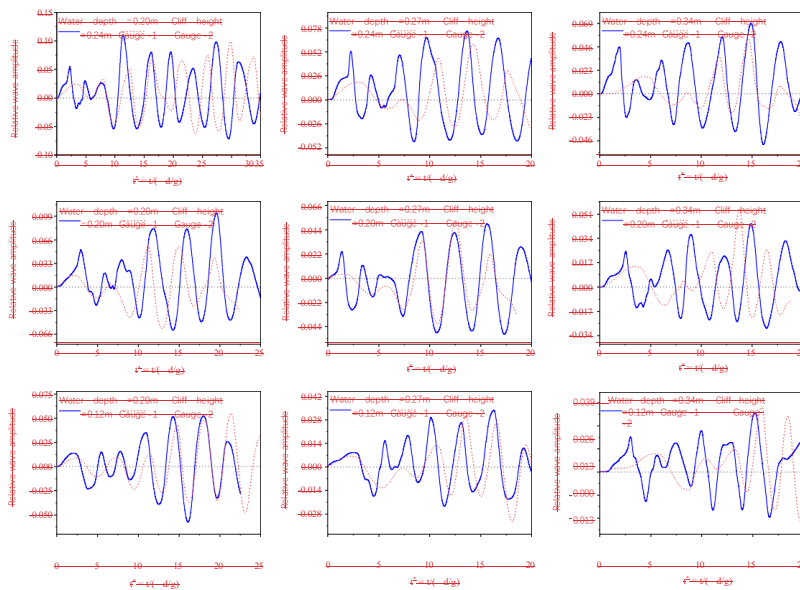
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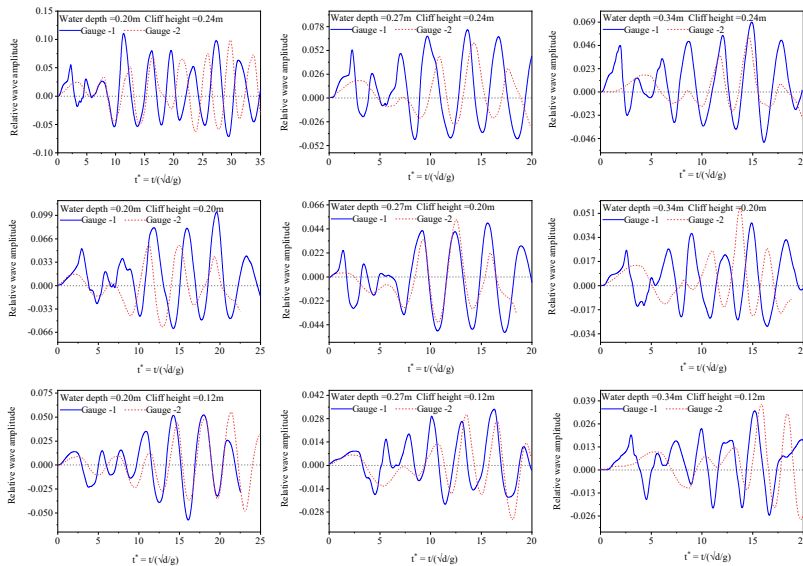


Fig. 9: A comparative display of the wave water waves recorded at gauge 1&2 for a 60° slope angle, and 0.445 m fall height.

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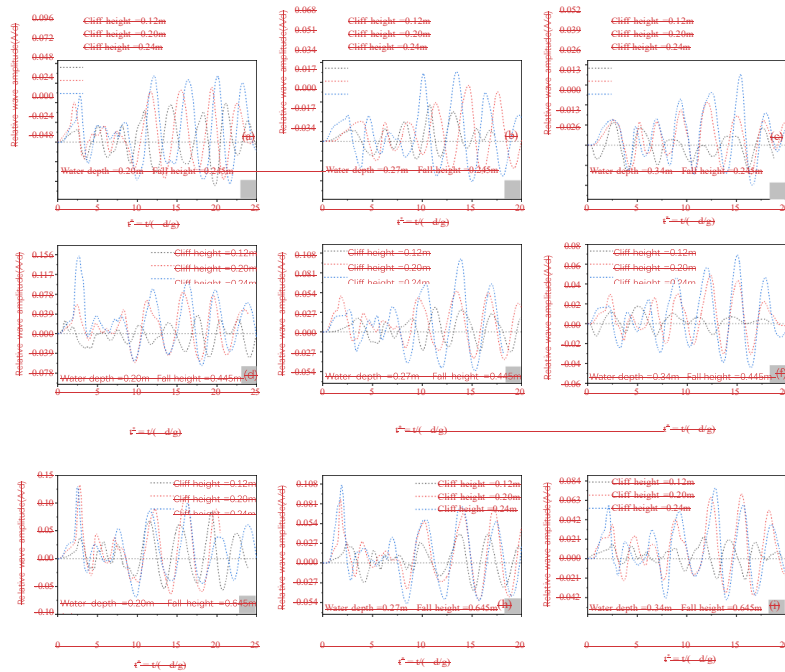
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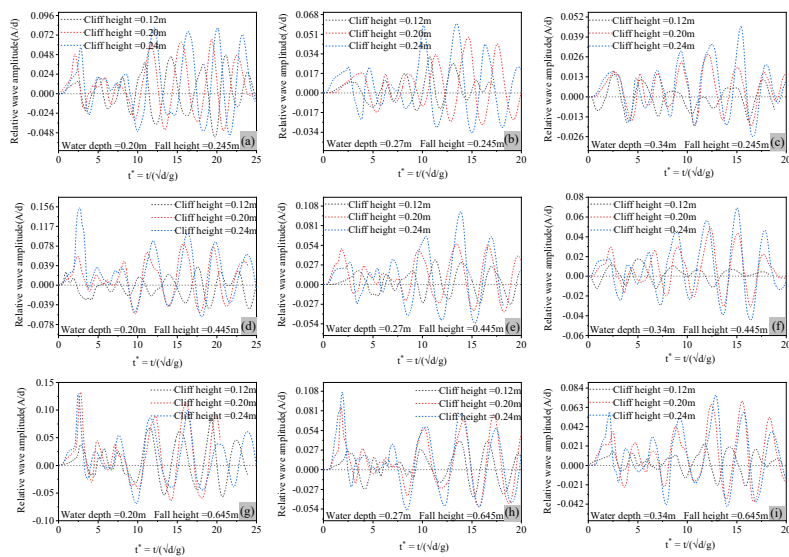
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Fig. 10: Relative wave amplitude for various water depths, cliff height, and fall height at 30°runup slope angle, (a, b&c) relative wave amplitude induced by 0.245m fall

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461 height, (d, e&f) relative wave amplitude induced by 0.445m fall height, (g, h&i)  
 462 relative wave amplitude induced by 0.645m fall height.

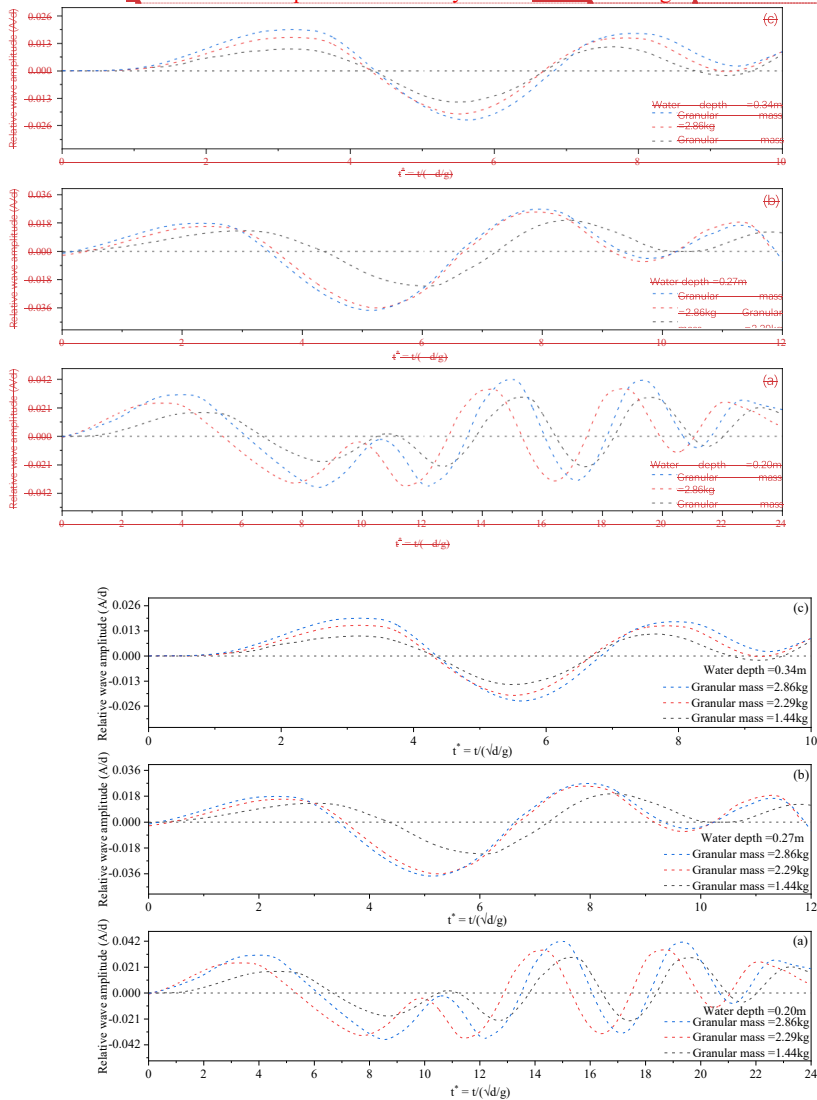


Fig. 11: WaveWater waves induced by equivalent granular mass at 30° slope angle

3.2 Numerical modeling results

The numerical simulations conducted in this study successfully captured key dynamic characteristics of the wave generated by the rotational cliff collapse, specifically the wave amplitude and wave runup, across a range of test cases. Moreover, the front velocity of the incident wave was also measured. The simulations were also focused on verifying the results obtained from the rotational cliff collapse in the experiments. To quantify the wave amplitude, runup, and velocity, a post-processing technique was employed. To establish the reliability of the wave front velocity measurements, the velocity was calculated at 5–7 distinct locations along the wave’s propagation path and at multiple time steps during the simulation. This multi-point sampling approach minimized errors due to spatial and temporal variations. Fig. 12 shows a representative case of wave formation and propagation in a water tank at a depth of  $d = 0.2$  m at various time frames.

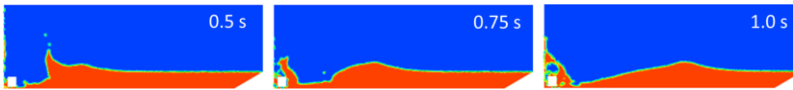


Fig. 12 Wave formation and propagation at water depth of  $d = 0.2$  m at various time frames.

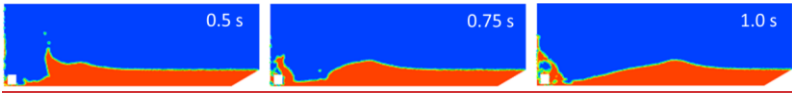


Fig. 12: Wave formation and propagation at water depth of  $d = 0.2$  m at various time frames.

The wave amplitude was defined as the peak vertical displacement of the liquid surface relative to the undisturbed free surface level. Fig. 13 illustrates a representative case, depicting the wave front propagation.

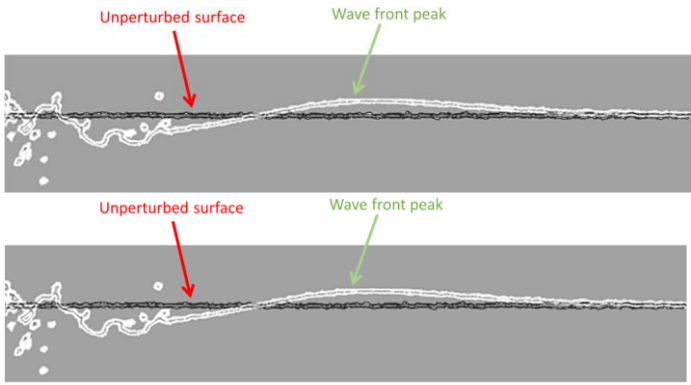


Fig. 13: Wave dynamics following a rotational cliff collapse in water depth  $d = 0.34$

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m. Stable liquid surface before impact (black line); wave front propagating away from the point of impact (white line).

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To validate the results of simulations, we compared the results of the runup height with the experimental values. Table 3 presents the runup values for various

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runup slope angles, i.e., 30°, 45°, and 60° for a water depth of  $d=0.27$  m. The comparison of simulated values was performed at this depth, as it lies in the middle of the experimental test range of water depths. Numerical modeling results indicate that for a fixed water depth, the runup values consistently decrease as the runup slope angle increases from 30° to 60°. At a water depth of  $d=0.27$  m, the runup decreases from 0.2 m at 30° to 0.17 m at 45°, and further to 0.11 m at 60°. This reduction is attributed to the changing momentum transfer dynamics with increasing slope angle. At less steep angles (closer to horizontal, e.g., 30°), the rock's momentum generates a stronger radial splash and greater upslope displacement of the liquid along the cliff. As the angle increases toward 60°, a larger component of the momentum is directed parallel to the cliff, reducing the vertical impulse. The experimental and numerical results agree well, and the difference lies within the acceptable range of 4-5%. The experimental results for the other two water depths also indicate similar behavior.

Table 3: Peak runup values along the various slope angles at a water depth of  $d=0.27$  m

Depth	Numerical- 30°	Exp- 30°	Numerical- 45°	Exp- 45°	Numerical- 60°	Exp- 60°
$d$ (m)						
0.27	0.20	0.19	0.17	0.16	0.11	0.102

Next, we measured the wave velocity through the numerical results, as it wasn't captured accurately through experimental images. Fig. 14 illustrates the simulated wave fronts at a time instant of  $t=1$  second following the impact of the solid rock on the liquid pool, for various water depths and a fixed slope angle of 30 degrees. These visualizations highlight the propagation of the waves from the impact zone. The slope angle was varied across simulations to assess its influence on wave characteristics. It was observed that changes in the slope angle induced only minor variations in both the wave front velocity and wave amplitude for a given pool depth. These perturbations were typically within 1–2% of the mean values. Consequently, to streamline the analysis and focus on dominant trends, the wave front velocity and height were averaged over the range of slope angles for each specific water depth.

However, variations in water depth exerted a pronounced effect on the wave

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dynamics, leading to significant alterations in both the propagation velocity and amplitude of the generated waves. This depth-dependent behavior is quantified in Table 4, which presents the averaged results from the numerical simulations. For a shallow water depth of  $d=0.2$  m, the average wave front velocity was computed as

1.48 m/s, with a corresponding average wave height of 0.11 m. As the pool depth increased to  $d=0.27$  m, the velocity rose to 1.58 m/s, while the wave height decreased to 0.07 m. Further deepening to  $d=0.34$  m yielded a velocity of 1.74 m/s and a reduced wave height/amplitude of 0.06 m. These trends indicate an approximately linear increase in velocity with depth, accompanied by an inverse relationship for wave amplitude.

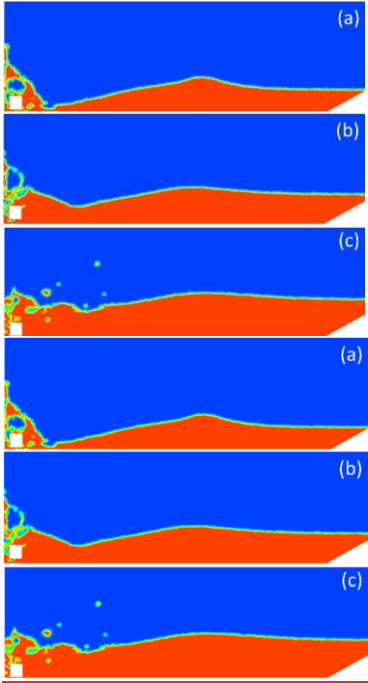


Fig. 14: Propagating wave fronts after the impact at time  $t = 1$  s for a slope angle of 30-degree. (a)  $d = 0.2$  m, (b)  $d = 0.27$  m, (c)  $d = 0.34$  m.

The observed depth dependence can be rationalized through fundamental principles of wave propagation in gravity-dominated, multi-phase flows. In the shallow water regime, given that the pool depths (0.2–0.34 m) are comparable to or smaller than the wavelengths of the generated waves, the phase velocity  $c$  of long gravity waves approximates  $\sqrt{gh}$ , where  $g$  is the gravitational acceleration ( $9.81 \text{ m/s}^2$ ) and  $h$  is the undisturbed water depth. This relation arises from the shallow water equations, where hydrostatic pressure balance and negligible vertical acceleration dominate, leading to a

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488	dispersionless incident wave speed that scales with the square root of depth.	Formatted
489	Substituting the water depths yields theoretical velocities of approximately 1.40 m/s for	Formatted
490	$d=0.2$ m, 1.63 m/s for $d=0.27$ m, and 1.83 m/s for $d=0.34$ m, which align closely with	Formatted
491	the simulated values (discrepancies of 7–10% may stem from viscous dissipation, non-	Formatted



### 3.3 MEP model results

The purpose was to develop a precise model for wave amplitude and runup induced by rotational cliff collapse. The predicted model is a function of seven variables, i.e., water depth, fall height, cliff mass, impact velocity, cliff height, runup slope angle, and number of fragments, and can be described as follows.

$$\text{Wave amplitude and runup} = f(d, H, m, v, h, \alpha, N_f) \quad (3)$$

The relation among the parameters was evaluated using Pearson's correlation to analyze the multicollinearity and interdependency between the parameters, as they can obscure the interpretation of the developed model. The model was developed by splitting the data into two subsets, i.e., training (70%) and testing (30%). The randomization was done by MEP itself. Following the criteria, 70% of the data, i.e., 57 data points, were taken as training data, whereas 30% of the data, i.e., 24 data points, were considered for validation of the model. The mathematical expression for MEP is obtained by solving the C++ code and representing it as per optimized hyperparameter settings, as shown in Table 5. The prediction model for wave amplitude and runup was developed by analyzing the MEP code in MATLAB, as shown in Equations 4 and 5.

Table 5: Parametric settings of the MEP algorithm for wave amplitude and runup

Sr.No.	Parameters	Wave amplitude	Wave runup
1	Number of sub-populations	125	85
2	Sub-population size	115	75
3	Crossover probability	0.85	0.60
4	Code length	35	25
5	Tournament size	30	10
6	Mutation probability	0.085	0.06
7	Number of generations	250	120
8	Crossover type	Uniform	Uniform
9	Error measure	Mean absolute error	Mean absolute error
10	Problem type	Regression	Regression
11	Function set	$+ \cdot \wedge$	$+ \cdot \wedge$
12	Terminal set	Problem Input	Problem Input
13	Operators	0.5	0.5
14	Simplified	Yes	Yes
15	Variables	0.5	0.5
16	Random seed	0	0
17	Constants	0	0

$$\text{Wave amplitude } A = d^{\frac{\alpha}{d(dGN_fGm)}} + \frac{2v h^k}{m + N_f + d(d + N_f + m)} + 2v h d^{\frac{\alpha}{d(dGN_fGm)}} \quad (4)$$

$$\text{Wave runup } R = \frac{A(h + (A(d \frac{B}{\alpha}))^{\frac{B}{\alpha}})^{\frac{B}{\alpha}}}{\alpha} \quad (5)$$

$$A = v + h^d$$

$$B = v + m + h^d$$

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Wave amplitude  $A = d^{\left(\frac{\alpha}{d(d+N_f+m)}\right)} + \frac{2vh^2}{m+N_f+d(d+N_f+m)} + 2vhd^{\left(\frac{\alpha}{d(d+N_f+m)}\right)}$  (4)

Wave runup  $R = \frac{A\left(h+\left(A\cdot\left(d-\frac{B}{\alpha}\right)\right)^{B/\alpha}\right)^A}{\alpha}$  (5)  $A = v + h^d$

$B = v + m + h^d$

529Whereas  $d$  is the water depth (m),  $m$  is the mass of the cliff (kg),  $v$  is the

530impact velocity (m/s),  $h$  is the cliff height (m),  $\alpha$  is the runup slope angle, and  $N_f$  is

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the number of fragments.

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### 3.3.1 Prediction performance of the developed model

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The robustness of the proposed model was evaluated by comparing it with well-established statistical indices, i.e., mean absolute error (MAE), root mean square error (RMSE), correlation coefficient (Cr), Nash–Sutcliffe efficiency (NSE), and performance index (PI). The indices can be represented by equation (6-10) (Khan et al., 2022).

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established statistical indices, i.e., mean absolute error (MAE), root mean square error (RMSE), correlation coefficient (Cr), Nash–Sutcliffe efficiency (NSE), and performance index (PI). The indices can be represented by equation (6-10) (Alavi et al., 2010; Khan et al., 2022).

$$MAE = \frac{\sum_{i=1}^n |e_i - p_i|}{n} \quad (6)$$

$$RMSE = \frac{\sum_{i=1}^n (e_i - p_i)^2}{n} \quad (7)$$

$$NSE = 1 - \frac{\sum_{i=1}^n (e_i - p_i)^2}{\sum_{i=1}^n (e_i - \bar{e})^2} \quad (8)$$

$$PI = \frac{RRMSE}{1+R} \quad (9)$$

$$R^2 = \frac{\sum_{i=1}^n (e_i - \bar{e})(p_i - \bar{p})}{\sum_{i=1}^n (e_i - \bar{e})^2 \sum_{i=1}^n (p_i - \bar{p})^2} \quad (10)$$

$$MAE = \frac{\sum_{i=1}^n |e_i - p_i|}{n} \quad (6)$$

$$RMSE = \frac{\sum_{i=1}^n (e_i - p_i)^2}{n} \quad (7)$$

$$NSE = 1 - \frac{\sum_{i=1}^n (e_i - p_i)^2}{\sum_{i=1}^n (e_i - \bar{e})^2} \quad (8)$$

$$PI = \frac{RRMSE}{1+R} \quad (9)$$

$$R^2 = \left( \frac{\sum_{i=1}^n (e_i - \bar{e})(p_i - \bar{p})}{\sum_{i=1}^n (e_i - \bar{e})^2 \sum_{i=1}^n (p_i - \bar{p})^2} \right)^2 \quad (10)$$

Whereas,  $\bar{e}$ ,  $\bar{p}$ , and  $\bar{p}$  are the average values of the experimental and predicted results, and  $e_i$  and  $p_i$  are  $i$ th values of the modeled and predicted results, for  $n$  total samples. It is good to consider the error indices while analyzing the predictive

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§46 capability of complex models. The wave runup model demonstrated a robust

§47 performance for both training and testing datasets. The lower values of RMSE and

§48 MAE indicate little deviation from experimental values, while RSE and RMSE values

§49 confirm lower normalized error, as shown in Table 6. The higher values of NSE and Cr

§50 further validated the model reliability for the training phase. Whereas for the validation

§51 dataset, i.e., the unseen data model displays even stronger performance with lower

§52 RMSE and MAE values compared to the training dataset. Moreover, higher Cr and

§53 lower performance index values highlight enhanced model efficiency. This suggests

§54 that the model works well for unseen data, making it suitable for predicting the wave

§55 runup induced by rotational cliff collapse (Gardezi et al., 2024).

§56 The predictive performance of the wave amplitude model in the case of training

§57 data demonstrated a strong correlation with high  $R^2$  values and low RMSE and MAE

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The accuracy of the proposed model can also be checked using residual error plots, which are obtained by subtracting experimental and predicted values (Alavi et al., 2013). The results indicate that the amplitude model has minimum and maximum values of 0.004 m and 0.0065 m, as shown in Fig. 16 (a), whereas for wave runup the minimum and maximum values are 0.01875 and 0.024 (Fig. 16b). Moreover, it can also be observed that error values are populated along the x axis, therefore, showing low error frequency, and accuracy of both the models.

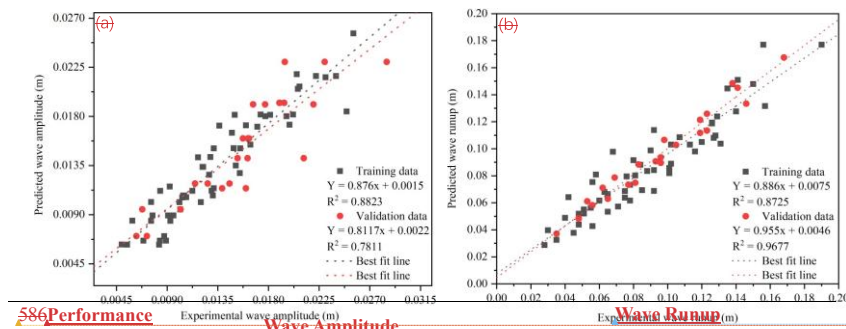


Fig. 15-  
Tracing the-  
experimental-  
results-by-  
predicted-  
values, (a)-  
wave-  
amplitude-  
and Training-  
data

Performance parameters	(a) wave-amplitude		(b) wave-runup	
	Training data	Validation data	Training data	Validation data
RSQ	0.8823	0.7811	0.8748	0.9691
RMSE	0.00178	0.0025	0.01327	0.00617
MAE	0.00135	0.00176	0.0108	0.00504
RSE	0.1180	0.2439	0.1306	0.0312
RRMSE	0.1314	0.1594	0.1472	0.0660
P. index	0.0698	0.0908	0.076	0.0333
NSE	0.8819	0.7560	0.8693	0.9687
C <sub>r</sub>	0.9393	0.8829	0.9353	0.9844

Previously, scientists have also used the slope of the regression line as a performance indicator for AI models, thus representing a correlation between experimental and predicted results. Fig. 15 (a & b) shows the regression line for our wave amplitude and runup model. For wave amplitude, the slope value for the training data set is 0.88, which is adequate, and 0.78 in validation, which is still greater than the minimum value of 0.7; it can happen

as the model involving numerous parameters and complex phenomena usually performs slower for the unseen data (Yarkoni and Westfall, 2019). Whereas, for wave runup, the model performed very well for both training and validation data sets with an  $R^2$  value of 0.87 and 0.96, respectively.

The accuracy of the proposed model can also be checked using residual error plots, which are obtained by subtracting experimental and predicted values. The results indicate that the amplitude model has minimum and maximum values of -0.004 m and 0.0065 m, as shown in Fig. 16 (a), whereas for wave runup the minimum and maximum values are -0.01875 and 0.024 (Fig. 16b).

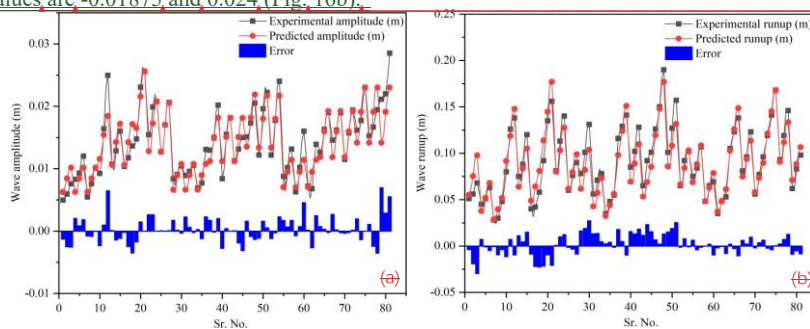


Fig. 16 Indicating error values between experimental and predicted model (a) wave amplitude, and (b) Wave runup

### 3.2 Validation of the developed model

The validation of the proposed model is an important feature in predictive modeling. It has been observed that sometimes the model performs very well for training data sets,

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Moreover, it can also be observed that error values are populated along the x-axis, therefore, showing low error frequency, and accuracy of both the models.

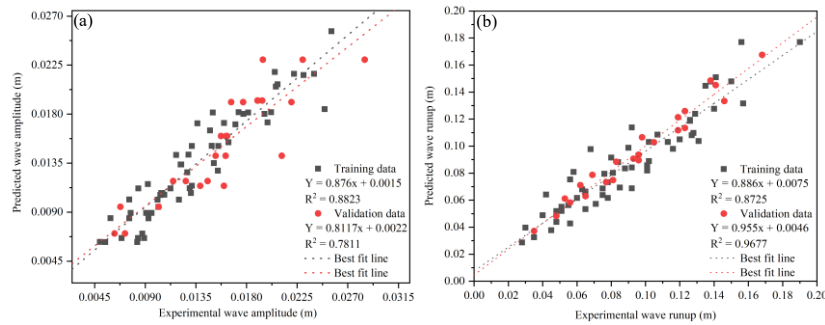


Fig. 15 Tracing the experimental results by predicted values, (a) wave amplitude and (b) wave runup

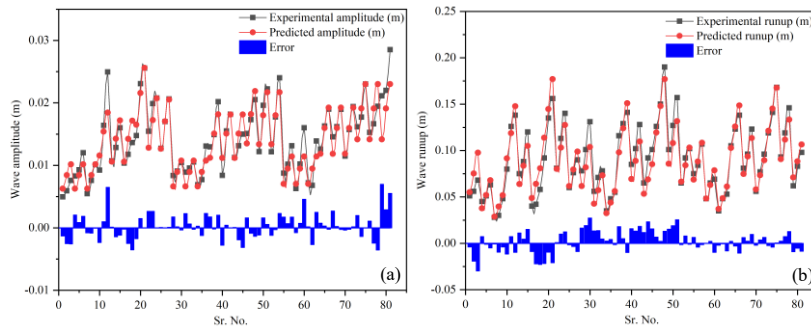


Fig. 16 Indicating error values between experimental and predicted model (a) wave amplitude, and (b) Wave runup

### 3.2 Validation of the developed model

The validation of the proposed model is an important feature in predictive modeling. It has been observed that sometimes the model performs very well for training data sets, but fails to perform during the validation stage for unseen data. So, the developed prediction model was further validated by conducting the sensitivity and parametric analysis for both the wave amplitude and runup.

#### 3.2.1 Sensitivity analysis

Sensitivity and parametric analysis play a vital role in determining the robustness of the proposed model. The sensitivity analysis (SA) of the developed prediction model for the entire dataset tells us how sensitive the model is to any changes in input parameters. So, for an independent parameter  $Y_i$ , the SA can be calculated using equations 11 and 12, which indicates that for any parameter, the values were varied between two extremes, and others were constant, at their average, and the outcome was found in the form of  $Y_i$ , and then the same process was repeated for all the remaining parameters.

$$R_k = f_{\max}(Y_k) - f_{\min}(Y_k) \quad (11)$$

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$f_{min}(Y_k)$  represent the minimum and maximum values of the model-based results grounded on the  $k^{th}$  domain of the input parameters in the above equation. Fig. 17 (a & b) shows the results of the sensitivity analysis of the developed prediction model for the wave amplitude and runup. Figure 17 (a) indicates that the wave amplitude is greatly influenced by the height of the cliff ( $h$ ) and has an effect of almost 51%. The water depth ( $d$ ) contributes 4.36% to wave amplitude, cliff mass ( $m$ ) contributes 4.69%, and impact velocity ( $v$ ) and number of fragments ( $N_f$ ) contribute 18% and 22% to the induced wave amplitude. Whereas the fall height ( $H$ ) and runup slope angle ( $\alpha$ ) do not affect the wave amplitude. Since the impact velocity parameters have already catered for the fall height that's why it is not visible in the proposed model. The model tells us that impact velocity, cliff height, and number of fragments contribute approximately 90% to the wave amplitude induced by the rotational fall of the cliff. It can be concluded that the effect of  $h \rightarrow N_f \rightarrow v \rightarrow m \rightarrow h \rightarrow N_f \rightarrow v \rightarrow m \rightarrow d$  on the induced wave amplitude.

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Similarly, the sensitivity analysis of wave runup (Fig. 17b) indicates that runup is

(9%) =

~~Rk~~\_\_\_\_\_

~~\_\_\_\_\_~~

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$$SA (\%) =$$

$$\frac{R_k}{\sum_{j=1}^n R_j} \times$$

10\_\_\_\_\_

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**Figure 1**

 $f(Y)$ 

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$$f = (Y_i)$$

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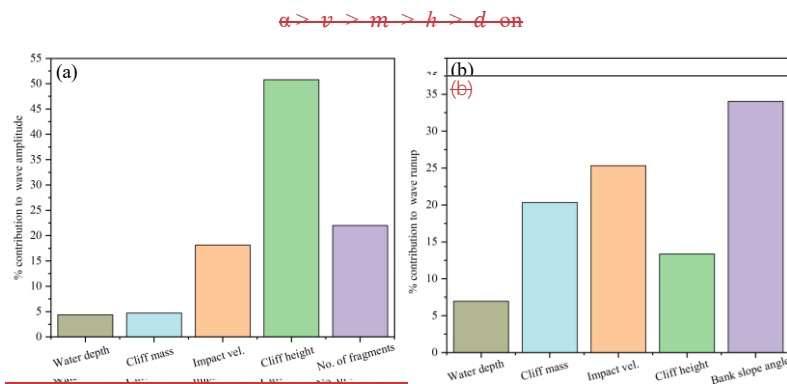
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625 greatly influenced by bank slope angle ( $\alpha$ ) and has an effect of 34%. Impact velocity  
626 ( $v$ ) contributes 25.3%, cliff mass ( $m$ ) 20.3%, cliff height ( $h$ ) 13.3%, and water depth  
627 ( $d$ ) contributes around 7% to wave runup. Whereas, the number of fragments and fall  
628 height that have already been catered in impact velocity don't contribute to wave runup.  
629 This suggests that wave runup is primarily governed by coastal geometry, i.e., bank  
630 slope angle and cliff height, and hydrodynamic forces, i.e., impact velocity, whereas  
631 water depth contributes a little to wave runup. It can also be concluded as the effect of  $\alpha > v > m > h > d$  on

$h > d$  on the induced wave amplitude.



632 Fig. 17 Sensitivity analysis of the induced MEP-based wave amplitude and runup prediction model

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634 Fig. 17 Sensitivity analysis of the MEP-based wave amplitude and runup model

### 635 3.2.2 Parametric Analysis

636 Parametric analysis results for the input parameters for the wave amplitude used in  
637 this study are displayed in Fig. 18. The parametric analysis indicates that wave  
638 amplitude decreases as the water depth, number of fragments, and cliff mass increase,  
639 whereas it increases with the increase in cliff height and impact velocity. These trends

are in line with the fundamental physics principles (Bougouin et al., 2020; Lipiejko et al., 2023) deep waters dissipate more energy, and greater impact velocities and larger cliff heights impart more kinetic and potential energies to water bodies for wave generation. Whereas, the inverse relation between the number of fragments and wave amplitude proposes a potential threshold effect in which initial fragmentation contributes to wave formation, whereas excessive fragments contribute to energy dissipation owing to increased turbulence. The sensitivity analysis further quantified the effect of these parameters, classifying cliff height as a major contributing factor in

### 3.2.2 Parametric Analysis

Parametric analysis results for the input parameters for the wave amplitude used in this study are displayed in Fig. 18. The parametric analysis indicates that wave amplitude decreases as the water depth, number of fragments, and cliff mass increase, whereas it increases with the increase in cliff height and impact velocity. These trends are in line with the fundamental physics principles (Bougouin et al., 2020; Lipiejko et al., 2023). Deep waters dissipate more energy, and greater impact velocities and larger cliff heights impart more kinetic and potential energies to water bodies for wave generation. In contrast, the inverse relation between the number of fragments and wave amplitude proposes a potential threshold effect in which initial fragmentation contributes to wave formation, whereas excessive fragments contribute to energy dissipation owing to increased turbulence. The sensitivity analysis further quantified the effect of these parameters, classifying cliff height as a major contributing factor in wave amplitude variations, followed by impact velocity, number of fragments, water

depth, and mass of cliff. The strong influence of cliff height indicates its direct effect in determining the potential energy for wave generation. Moreover, the larger sensitivity value of fragments regardless of their inverse parametric relation shows a complex relation, where fragment count plays a considerable but context-dependent role in wave generation and propagation. The dominance of cliff height, impact velocity, and fragment count suggests that these parameters should be prioritized in future prediction models. These findings are important for developing predictive models for wave generations due to rotational cliff collapse.

The developed model for wave amplitude provides valuable insights into fundamental physics governing wave formation and propagation induced because of rotational cliff collapse. The strong height dependence of the model confirms the classical principle of conservation of potential energy, whereas the fragment count dependence reveals energy partitioning mechanisms. The results of performance indices and sensitivity, and parametric analysis increase our understanding of how geometric and dynamic characteristics govern the wave characteristics, with relevance to hazard assessment and disaster mitigation in coastal regions prone to cliff collapse following rotational motion.

The results of the parametric analysis for wave runup are presented in Fig. 19. It can be observed from Fig. 19 (a & e) that as the water depth and bank slope angle increase, the wave runup decreases due to energy dissipation and different wave breaking dynamics. Conversely, as the cliff mass, cliff height, and impact velocity increase, the wave runup increases, as greater kinetic energy and inertia impart greater uprush. Notably, all the parameters present a strong correlation with the runup (more than 97%), highlighting their statistical significance. The results agree with the general physics laws, where mild slopes and larger impact forces result in higher runups, whereas deep waters attenuate wave energy.

An important observation from parametric analysis of wave amplitude and runup, as shown in Fig. 18b, and 19c, indicates that cliff mass represents a nonlinear relation

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677— with wave amplitude and a linear relation with runup. This is due to the fact that the

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variations in wave amplitude are governed by a nonlinear energy dissipation, where hydrodynamic forces follow a quadratic dependence on the velocity. In the case of light cliff collapses, the dynamic responses result in complex absorption and distribution, whereas heavier cliff collapses promote wave reflection along with nonlinear effects of wave breaking and splash-induced turbulence, as can be observed in Fig. 5 (b, e&h). Conversely, the wave runup exhibits a linear relation with cliff mass, and this is due to the law of conservation of momentum, such that the resisting inertial force is directly proportional to cliff mass. The greater resistance to motion of heavier cliffs allows more energy to be conserved and utilized for higher wave runups before dissipation. The main difference between the two trends is that the wave amplitude is controlled by localized energy losses, whereas runup is governed by bulk momentum transfer rather than localized losses.

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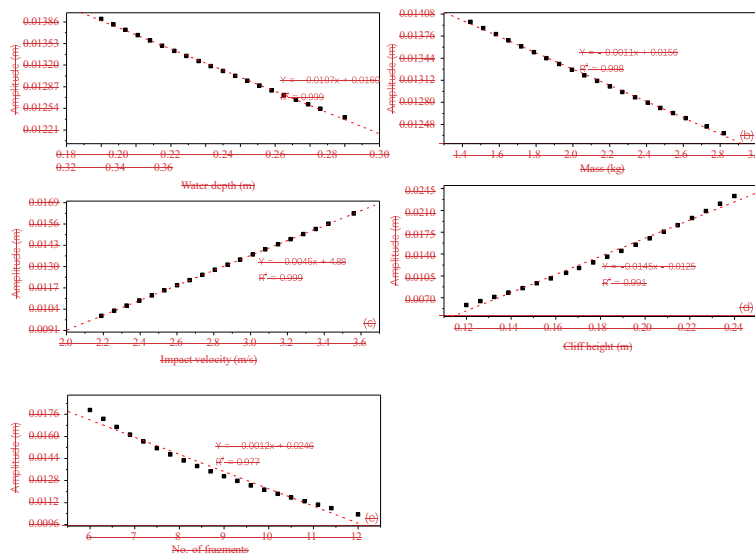
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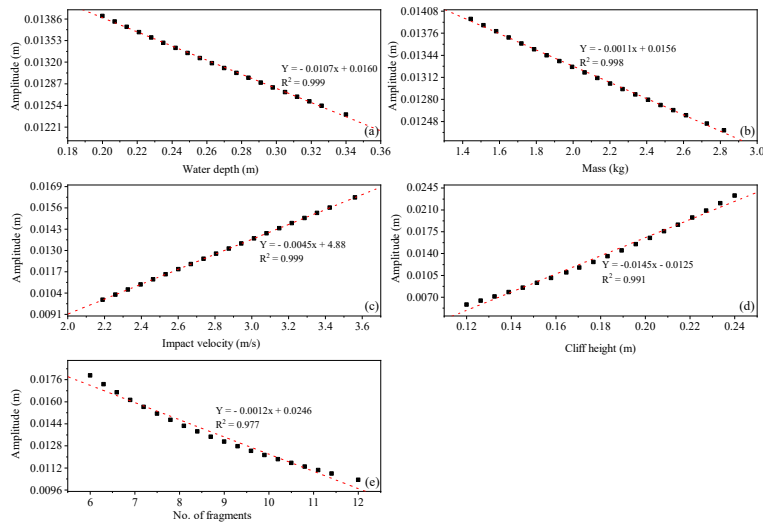
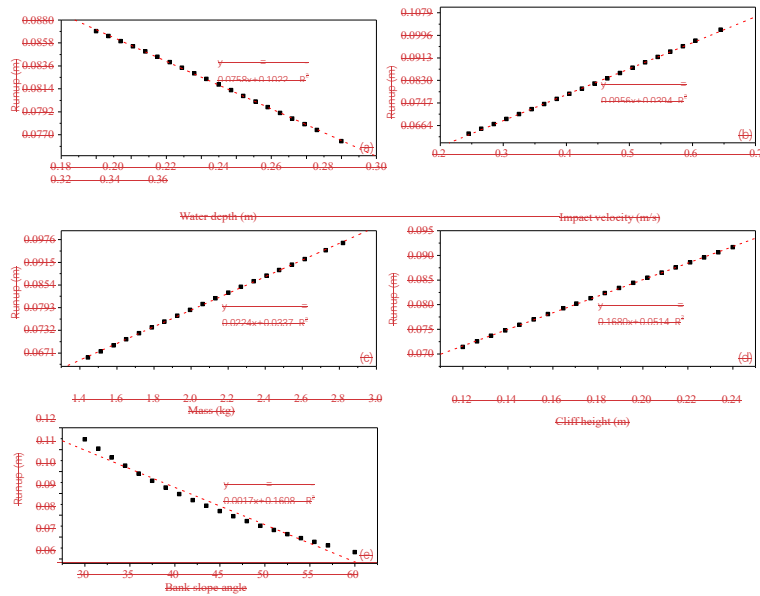
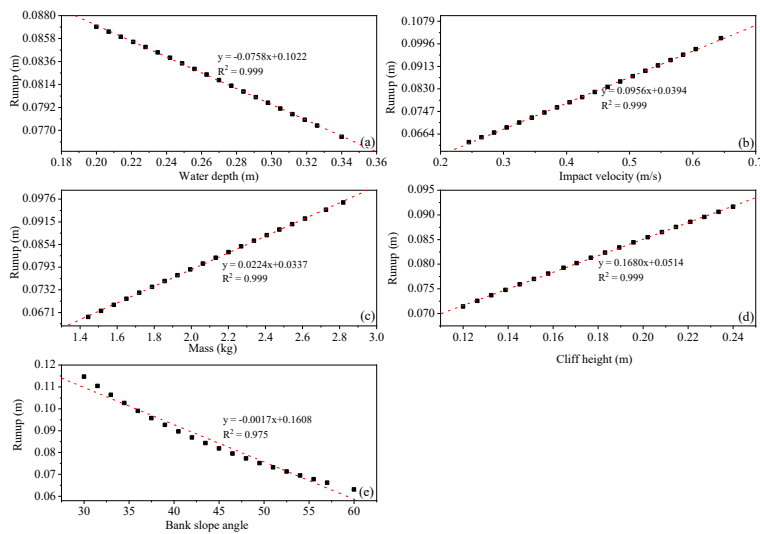


Fig. 18 Parametric analysis for wave amplitude (a) water depth, (b) cliff mass, (c) impact velocity, (d) cliff height, (e) number of fragments.



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694 Fig. 19 Parametric analysis for wave runup (a) water depth, (b) impact velocity, (c)  
695 mass of the cliff, (d) cliff height, (e) bank slope angle.

#### 6964 4. Conclusions

697 While designing wave protection structures along the banks of reservoirs, it is

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698 common to use the empirical relations developed for granular flows, i.e., landslides and

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699 avalanches, to predict the amplitude and runup of the waves. However, the waves

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700 induced by various types of slides behave differently and should be treated accordingly.

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701 The dynamics of the waves induced by falling cliffs are entirely different from the

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702 waves induced by continuous granular flows. Similarly, the dynamics of the waves

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703 induced by falling cliffs following different types of motion (translational, rotational)

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704 are also different. This study aimed to develop a novel wave amplitude and runup

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705 prediction model for waves induced by rotational fall of the cliff using a combination

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706 of seven governing parameters, and then compare it with the dynamics of the wave

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707 induced by continuous granular flows. Based on the results and discussions, the study

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708 concludes as follows,

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709 1. It was concluded that water depth strongly controls the shape of the induced splash depends on water depth;

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710 increased depth forms a and wave amplification. Shallow water induced elongated, tall splashes, and higher wave amplitudes; in contrast, deep water produced mushroom-shaped splash, whereas shallow water forms a

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711 vertically elongated splash. Moreover, shallow water allows more amplification of  
7121. waves for the same splashes with higher energy level compared to deep water dissipation and lower wave  
amplitudes.  
713 2. The effect of viscous forces is very, very small and can be ignored, since the  
714 Reynoldshigher values of Froude number ( $> 1.2$ ) for all the experiments is very high, thus  
leaving indicate that the viscous effects were negligible, so the Froude  
7152. number was selected as the best possiblemost suitable dynamic scaling factor. It is concluded that the  
Froude for describing the behaviour of the waves.  
716 number also increases as the water depth decreases.  
7173. 3. The study concludes that wave amplitude iswas greatly influenced by cliff height; (51 %), number  
of fragments (22 %), Impact velocity (18 %), cliff mass (4.69 %), and water depth (4.36 %). Whereas the  
wave runup was governed by the runup slope angle, impact velocity, and cliff mass.  
718 impact velocity, and the number of fragments. For all the cases, the deep water  
719 dissipated more energy compared to shallow waters, thus resulting in lower  
720 amplitudes.  
721 4. The amplitude of the wave induced by equivalent granular mass issliding on a 30° slope was  
28-42% lower than the  
722 waves induced by rotational cliff collapse, thus concluding that the mode of energy  
7234. transfer to the water body plays a critical role in wave dynamics.  
724 5. A second level validation of the developed model was performed by conducting  
725 sensitivity and parametric analysis. It is concluded that the amplitude is highly  
726 sensitive to any change in cliff height, impact velocity, and number of fragments.  
727 In contrast, runup greatly depends on runup slope angle, impact velocity, and mass  
728 of the cliff.  
5. A novel MEP-based prediction model was developed for wave amplitude and runup. The model showed  
great performance during the training and testing stage, and showed high sensitivity to the used parameters,  
thus confirming its reliability.  
6. Research findings highlight that accurate hazard assessment of the cliff collapse requires models that account  
for the rotational failure mode and the fragmentation upon impact with the water surface. Traditional granular  
slide models may result in an underestimation of the initial wave amplitude and energy transferred.  
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Author contributions

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Declarations

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Competing interests: The authors declare no competing interests.

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