

Convection-generated gravity waves in the tropical lower stratosphere from Aeolus wind profiling, GNSS-RO and ERA5 reanalysis

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Abstract. The European Space Agency's Aeolus satellite, equipped with the Atmospheric LAser Doppler INstrument (ALADIN), provides near-global wind profiles from the surface to about 30 km altitude. These wind measurements enable the investigation of atmospheric dynamics, including gravity waves (GWs) in the upper troposphere and lower stratosphere (UTLS). This study analyzes ALADIN wind observations and ERA5 reanalysis, by deriving GWs kinetic energy (E_k) distributions, examining their temporal and spatial variability throughout the tropical UTLS. A prominent hotspot of enhanced GW activity is identified by Aeolus, migrating from the Indian Ocean in Boreal Summer to the Maritime Continent in Boreal Winter, closely matching outgoing longwave radiation minima and thus highlighting convective origins. Results show that ERA5 consistently underestimates E_k in convective regions, especially over the Indian Ocean, where conventional wind measurements are sparse. Additional comparisons with Global Navigation Satellite System Radio Occultation (GNSS-RO) measurements of GW potential energy (E_p) corroborate these findings and suggest significant underrepresentation of convection-driven wave activity in reanalyses. A multi-instrumental exploratory analysis also allows to verify the empirical grounding of the established E_k to E_p ratio. By providing direct wind measurements in otherwise data-sparse regions, Aeolus offers a valuable dataset for evaluating and potentially improving the representation of GWs in reanalyses, particularly in remote tropical areas. The combination of Aeolus and GNSS-RO data allows for an observationally-based examination of the partitioning between kinetic and potential energy, highlighting discrepancies with reanalysis products that could inform future model parameterization development.

1. Introduction

Atmospheric reanalyses like ERA5, a global atmospheric dataset produced by the European Centre for Medium-range Weather Forecasts (ECMWF), are essential for climate assessments and atmospheric research (Hersbach et al., 2020). By integrating observational data with state-of-the-art general circulation models and data assimilation methods, reanalyses provide comprehensive atmospheric snapshots for a variety of meteorological research (Muñoz-Sabater et al., 2021).

However, one significant limitation of these datasets, including ERA5, is their reliance predominantly on temperature measurements for data assimilation, with wind measurements being notably sparse (Campos et al., 2022; Podglajen et al., 2014). Because of this, ERA5 tends to underestimate low-level wind speeds in certain regions, compared to radiosonde measurements (Munday et al., 2022). Having said that, only a relatively few radiosonde and cloud-tracked wind measurements directly constrain wind variability: Radiosonde measurements are notably sparse over oceans, as they are typically launched from land-based stations, leaving vast oceanic regions under-sampled (Baker et al., 2014; Ladstädter et al., 2011). While some ship-based radiosonde launches occur, they are infrequent and cover limited areas. Satellite cloud-tracking methods, such as Atmospheric Motion Vectors (AMVs), provide wind data by tracking cloud movements (Bedka et al., 2009). However, these methods have limitations: they cannot retrieve wind profiles in clear-sky conditions and often lack detailed vertical resolution. This results in significant observational gaps in wind measurements over oceans and clear-sky regions. This limitation is particularly critical when considering atmospheric waves, such as gravity waves, which manifest themselves in temperature and wind vertical profiles.

Gravity waves (GW) play a crucial role in the dynamics of the Earth's atmosphere. Generated by mechanisms such as flow over orography, convection, and flow deformation, these waves are instrumental in transporting momentum and energy, influencing atmospheric regions far from their origin points (Fritts and Alexander, 2003). While Rossby waves are well represented due to their quasi-geostrophic nature, divergent wave modes like gravity waves, Kelvin waves, Rossby-gravity waves, and inertia-gravity waves are not sufficiently characterized and must often be parametrized internally by the models (Plougonven and Zhang, 2014). The underrepresentation of gravity waves with long horizontal and short vertical scales in ERA5 has been highlighted previously (Bramberger et al., 2022).

For the study period, ERA5 utilizes the non-orographic gravity wave drag (GWD) scheme described by Orr et al., (2010), which is based on a spectral approach (Scinocca, 2003 ; Referred to as S03 in Orr et al., 2010). This scheme does not explicitly resolve convectively generated waves based on model-diagnosed convection; instead, it launches a globally uniform and constant spectrum of waves from the troposphere. The momentum deposition occurs as these waves propagate vertically and interact with the resolved flow via critical-level filtering and nonlinear dissipation. While this parameterization improves the

middle atmosphere climate compared to simpler schemes, evaluations have shown it has limitations in fully capturing the required wave forcing, particularly for the Quasi-Biennial Oscillation (QBO) in the tropics (Pahlavan et al., 2021).

Furthermore, even with improvements in reanalysis products, challenges in accurately representing tropical winds persist. Studies of previous-generation reanalyses identified significant errors in tropical regions (Podglajen et al., 2014), and recent work shows that even ERA5's accuracy is highly site-dependent, with notable errors in locations influenced by warm currents (Campos et al., 2022). This is compounded by difficulties in data assimilation systems, such as 4-D var and perfect model scenarios, which struggle to extract circulation information from high-resolution temperature data (Žagar et al., 2004). Despite advancements in the quality of tropical forecasts and analyses, the evidence suggests that radio occultation (RO) data could potentially enable effective long-term monitoring of wind fields globally (Danzer et al., 2023). However, the overall lack of direct wind observations continues to pose significant challenges (Baker et al., 2014).

Historically, most GW studies have relied on ground-based or single-use instruments like radiosondes (Zhang and Yi, 2005), rockets (Wüst and Bittner, 2008), or global coverage measurements from the Global Navigation Satellite System Radio Occultation (GNSS-RO). While GNSS-RO provides high-resolution temperature profiling, effectively characterizing GW potential energy (E_p) (Fröhlich et al., 2007; Khaykin et al., 2015; Schmidt et al., 2016), it does not capture kinetic energy (E_k), which requires precise wind profiling.

In an effort to bridge many gaps within the observational world, the 2018 launch of the European Space Agency's Aeolus satellite changed our ability to capture atmospheric dynamics, particularly in the upper troposphere and lower stratosphere (UTLS). The UTLS is a region marked by a dramatic increase in static stability at the tropopause, where gravity waves are refracted to shorter vertical wavelengths (Dhaka et al., 2006; Geldenhuys et al., 2023). These waves with short vertical wavelengths (typically 2-10 km) are primarily lower-frequency gravity waves, as dictated by the dispersion relation, and exhibit relatively large amplitude wind variability. The Aeolus satellite, equipped with its Atmospheric LAsER Doppler INstrument (ALADIN), is able to measure global wind profiles up to an altitude of 30 km, providing insights into the behavior of gravity waves with vertical wavelengths down to ~ 1.5 -2 km in these critical atmospheric layers (Banyard et al., 2021; Rennie et al., 2021; Ratynski et al., 2023).

In this context, this study aims at utilizing Aeolus's global wind profiling capabilities to derive a tropics-wide distribution and variability of the kinetic energy of gravity waves, addressing a gap not typically captured in ERA5 reanalysis. By comparing direct measurements with ERA5 data, we reveal certain limitations in the reanalysis's ability to represent tropical gravity wave dynamics. We will look at the most recent reprocessed Aeolus baseline 2B16, providing data from June 2019 to August 2022. Additionally, our study aims at exploring a broader set of analyses, aiming to contextualize the Aeolus wind observations within a multi-instrument framework. By comparing Aeolus-derived kinetic energy of GWs with the potential energy estimates

67 from GNSS-RO, we assess the consistency of independent data sources and examine the ratio of kinetic to potential energy in
68 real-world atmospheric conditions. With this study, we provide the first observationally-based, tropics-wide estimate of gravity
69 wave kinetic energy from June 2019 to August 2022, directly linking its variability to deep convective sources.

70

71 The paper will be organized as follows: In Sect. 2, we will discuss the data as well as the methods. It includes a description of
72 the Aeolus, ERA5 and the GNSS-RO datasets, but also explains the horizontal detrending method with its potential and
73 limitations. In Sect. 3, we will analyze the wave activity in terms of kinetic energy using Aeolus Rayleigh wind profiling and
74 directly comparing it with ERA5. Additionally, in Sect.4, we broaden our analyses to contextualize Aeolus observations against
75 GNSS-RO data and criticize the ratio between both elements. Finally, the results are discussed in Sect. 5, followed by the
76 conclusions in Sect. 6

77 **2. Data and Methods**

78 **2.1 Instruments and Datasets**

79 The Aeolus satellite, with its ALADIN Doppler wind lidar, orbited Earth at a 97-degree inclination and 320 km altitude. Its
80 data consists of 24 vertical range bins that divide the atmosphere, allowing wind profiling between 0 and 30 km. Laser pulses
81 and two receivers—Rayleigh and Mie channels—detect the atmosphere's Doppler shifts through molecular and particle
82 backscatter, respectively. The data, organized into atmospheric scenes, cloudy or clear, has an 87 km along-track integration
83 and a vertical resolution varying between 0.25 to 2 km. The distribution of these range bins is determined by a dedicated range
84 bin setting (RBS), which can be adjusted to cater to specific needs, such as enhanced sampling at certain heights. This study
85 uses the Level 2B Rayleigh clear product, with the latest Baseline 2B16 at the time of submission, offering the horizontal line
86 of sight (HLOS) wind components. The HLOS wind speed is derived using Aeolus NWP Impact Experiments guidance, with
87 the vertical wind speed assumed to be negligible. A complete description of the instrument, its measurement principles, range
88 bin settings, and data products can be found in Rennie and Isaksen. (2024). The angle θ denotes the azimuth of the target-to-
89 satellite pointing vector, being around 100.5° over the tropics. When injecting the azimuth value into Eq. 1, it becomes apparent
90 that the HLOS wind over the tropics is quasi-zonal.

91

$$92 \quad v_{HLOS} = -u \sin(\theta) - v \cos(\theta) \quad (1)$$

93

94 The ERA5 reanalysis dataset, a ECMWF product, offers comprehensive atmospheric, land-surface, and ocean-wave
95 parameters at hourly resolution and global coverage (Hersbach et al., 2020). Its exceptional horizontal resolution of
96 approximately 33 km at the equator (corresponding to 0.3° latitude/longitude), the best among widely used reanalysis products,
97 enables it to resolve gravity waves with horizontal wavelengths as small as ~ 100 km (Wright and Hindley, 2018, their table

1). The data products used in this study were retrieved from the ECMWF archive on a regular $0.25^\circ \times 0.25^\circ$ latitude-longitude grid. Additionally, its higher vertical resolution in the troposphere, with 137 vertical levels reaching up to 0.01 hPa, makes it particularly adept at capturing gravity waves with vertical wavelengths down to $\sim 1\text{--}2$ km. ERA5 also incorporates advanced modelling features such as sponge layers and hyperdiffusion to attenuate artificial wave reflections and stabilize the model numerically, allowing for efficient modelling of large-scale phenomena, notably simulating gravity waves with wavelengths greater than 400 km (Stephan and Mariaccia, 2021). It is therefore the best candidate to use as a benchmark for Aeolus' performances. For this study, wind components are retrieved on the native 137 model levels and then interpolated to our high-resolution analysis grid using geopotential height, temperature, and humidity.

The GNSS-RO method offers many advantages for studying atmospheric dynamics, particularly GW activity and parameters. The first RO-derived GW estimates date back from the early 2000s and several missions have since provided data for further studies, focusing on potential energy as a proxy for retrieving GW activity (Tsuda et al., 2000; Fröhlich et al., 2007; Wang and Alexander, 2010; Luna et al., 2013; Schmidt et al., 2016). The Radio Occultation Meteorology Satellite Application Facility (ROMSAF) provides global GNSS-RO datasets. For the study period of June 2019 to August 2022 these datasets are dominated by the Metop constellation: Metop-B and Metop-C throughout, with Metop-A contributing until its retirement in November 2021 (von Engel et al., 2011). These datasets are derived from the bending angles of GNSS signals as they pass through the Earth's atmosphere and are observed by low Earth-orbiting satellites. It provides global coverage with a high vertical resolution, sub-Kelvin accuracy, full diurnal coverage, and all-weather capability. The vertical resolution of GNSS-RO temperature profiles is fundamentally limited by diffraction and varies with altitude, typically ranging from ~ 0.5 km in the lower troposphere to ~ 1.4 km in the middle atmosphere (Kursinski et al., 1997). While sharp vertical gradients in refractivity (e.g., due to temperature inversions or strong humidity gradients) can be detected, the effective resolution for resolving distinct atmospheric layers is constrained by these diffraction limits. The along-track horizontal resolution is typically around 200-300 km. Marquart and Healy (2005) showed that small-scale fluctuations in dry temperature RO profiles could be attributed to GWs with vertical wavelengths equal to or greater than 2 kilometers. Alexander et al. (2008b) suggested analyzing data below 30 kilometers in altitude to maintain the signal-to-noise ratio for temperature fluctuations above the detection threshold, which also happens to be Aeolus' maximal capability. Most GW parameters can be derived from single RO temperature profiles. However, estimating momentum flux requires knowledge of the horizontal wave number or wavelength, which cannot be deduced from a single temperature profile. To determine the horizontal structure of GWs, it is necessary to analyze clusters of three or more profiles adjacent in space and time.

This study specifically utilizes Aeolus Level 2B Rayleigh clear HLOS winds, ERA5 wind components, and GNSS-RO temperature profiles, all brought to a standard interpolated grid to facilitate the accurate comparison and integration of data from the different sources. The chosen grid has a vertical resolution of 100 meters and spans a range from 0 to 30 km altitude. This approach preserves the maximum vertical detail from each dataset before analysis.

131

132 The choice to compare Aeolus measurements with the ERA5 reanalysis, which does not assimilate Aeolus winds, is intentional.
133 This approach allows for a direct assessment of Aeolus's potential contribution by treating it as an independent dataset, thereby
134 highlighting regions where its direct wind measurements might fill observational gaps present in the conventional observing
135 system assimilated by ERA5.

136

137 2.2 Methods and Limitations

138 The following section discusses the retrieval of GW kinetic energy, E_k . A primary challenge in this retrieval, particularly in
139 the tropical UTLS, is the robust separation of GWs from other dominant, synoptic-to-planetary scale equatorial waves, such
140 as Kelvin waves. Observational studies using GNSS-RO data have consistently shown that Kelvin waves, with typical vertical
141 wavelengths in the range of $\sim 4\text{--}8$ km (Randel et al., 2021; Randel and Wu, 2005), are a prominent feature of the tropical
142 temperature and wind fields. This presents a potential for spectral overlap with the longer vertical wavelength portion of the
143 GW spectrum that this study aims to capture.

144 Several methods exist for background state determination and large-scale process separation. These broadly fall into two
145 categories: Vertical Detrending (VD), often applied to single profiles from instruments like lidars and radiosondes (Gubenko
146 et al., 2012; Khaykin et al., 2015), and Horizontal Detrending (HD). HD requires spatially resolved datasets like satellite
147 observations or model reanalyses and typically involves spatio-temporal averaging to define the background (Alexander et al.,
148 2008b; Khaykin et al., 2015). A comparative study by John and Kumar., (2013) highlighted significant discrepancies in derived
149 E_p magnitudes depending on the chosen method.

150 The choice of detrending method is particularly critical in the tropics due to the presence of waves like Kelvin waves. VD
151 methods, if not carefully designed, may inadvertently remove GWs with long vertical wavelengths or, conversely, retain short
152 vertical wavelength components of planetary-scale waves (e.g., Kelvin waves observed with vertical wavelengths as short as
153 3 km, as noted by Alexander and Ortland., (2010) and Cao et al., (2022)). Consequently, Schmidt et al. (2016) strongly
154 recommend using HD for satellite data, as VD may overestimate GW activity by including remnant signals from synoptic and
155 planetary waves that possess significant vertical structure in the tropics. Given these considerations, our study employs an HD
156 approach, calculating the background profile within a fixed spatio-temporal grid (20° longitude $\times$ 5° latitude over 7 days),
157 which we deem best suited for retrieving GW energy information from the Aeolus, GNSS-RO, and ERA5 datasets.

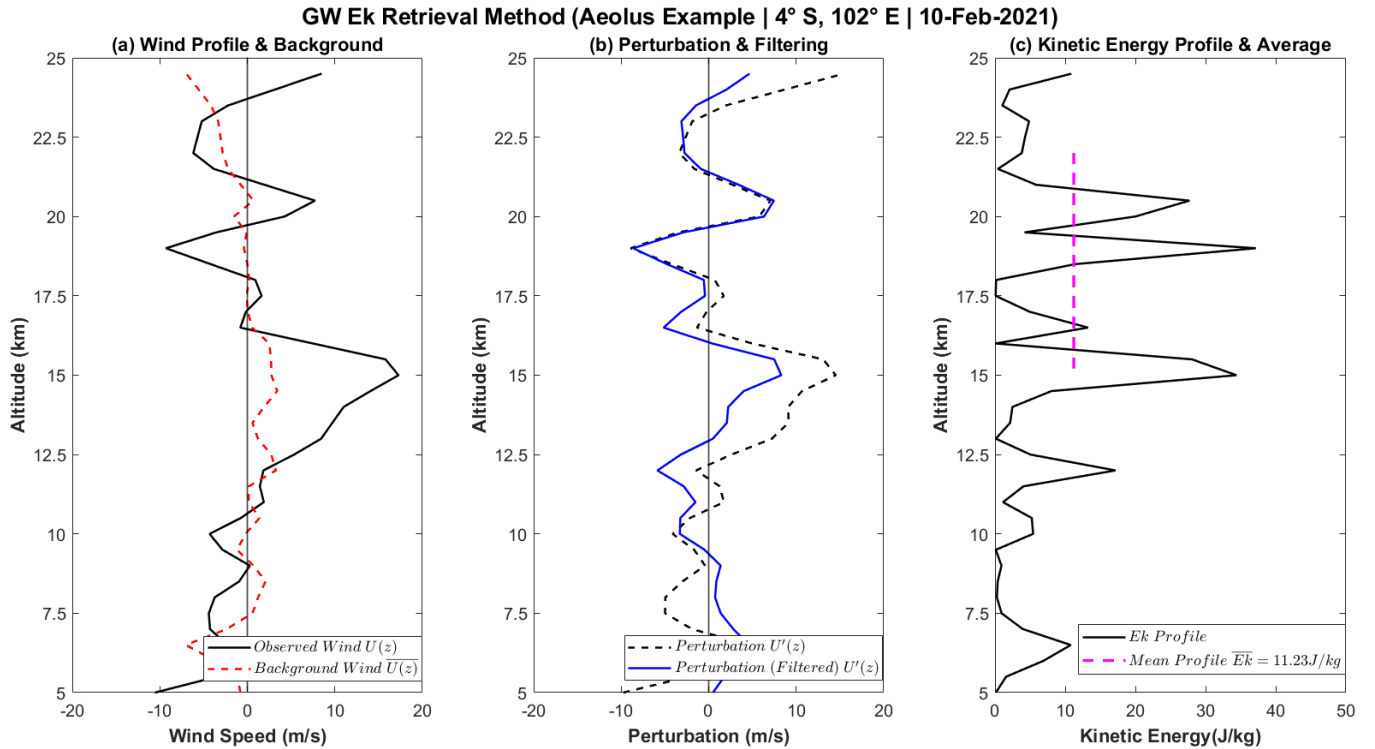
158 The separation of the wind or temperature profile into a background state and perturbations using HD is intended to isolate
159 fluctuations characteristic of gravity waves by filtering out larger-scale and slower-evolving processes like the mean
160 components of Rossby and Kelvin waves. This selection relies on the distinct scale and structural characteristics of GW
161 perturbations. However, the work by Randel et al., (2021) using dense COSMIC-2 RO data reveals further complexities. They
162 found that "residual" small-scale temperature variances (analogous to our perturbation fields) exhibit coherent maxima in the
163 longitudinal and vertical shear zones of large-scale Kelvin waves. This suggests that the local atmospheric environment shaped

164 by Kelvin waves, particularly variations in static stability (N^2), can modulate the amplitude of smaller-scale variability,
 165 potentially including GWs. Furthermore, the assimilation of Aeolus wind data itself has been shown to directly impact the
 166 representation of vertically propagating Kelvin waves in numerical weather prediction models, especially in regions of strong
 167 vertical wind shear (Žagar et al., 2021). This implies that Kelvin waves are indeed present in the Aeolus observations and that
 168 their characteristics might differ from those in reanalyses not assimilating Aeolus data.

169 To further refine the isolation of GWs and address the potential aliasing from such equatorial waves, our HD approach is
 170 combined with a vertical band-pass filter applied to the perturbation profiles. This filter targets vertical wavelengths between
 171 1.5 km and 9 km. The lower limit is chosen based on the effective vertical resolution of the instruments (particularly Aeolus),
 172 while the upper limit of 9 km is selected to be slightly above the typical dominant vertical wavelengths reported for Kelvin
 173 waves in the UTLS, thereby further reducing their potential contribution. This combined HD and vertical filtering methodology
 174 has been widely used for retrieving GW Ep from temperature data (Alexander et al., 2008b; Schmidt et al., 2008; Šácha et al.,
 175 2014; Khaykin et al., 2015), and the availability of Aeolus wind profiles allows us to apply a consistent approach for GW Ek.

176 Nevertheless, it is acknowledged that a perfect separation is difficult. The sharpest vertical gradients associated with Kelvin
 177 waves, or localized enhancements of GW activity within Kelvin wave-modified environments as suggested by Randel et al.,
 178 (2021) might still contribute to the derived GW energy. The interpretation of our GW Ek and Ep must therefore consider this
 179 context, particularly when analyzing variability in regions known for strong equatorial wave activity.

180



181 **Figure 1. Derivation of GW energy profiles from wind measurements (a) Observed wind profile and the corresponding background**
 182 **state profile. (b) Wind perturbation profile alongside its filtered counterpart. (c) Resulting in Ek, smoothed and then averaged within**
 183 **the given altitude range.**

184 Based on the linear theory of GW, the measured wind profile $U(z)$ shown in Fig.1a is divided into a background wind $\bar{U}(z)$
 185 also present in Fig.1a and a perturbation $U'(z)$ depicted in Fig.1b. The background is obtained by averaging all individual wind
 186 profiles for kinetic energy retrieval, within a spatiotemporal grid box of 20° longitude $\times$ 5° latitude over 7 days. While this
 187 horizontal detrending method was originally demonstrated using temperature profiles in Alexander et al., (2008b), its
 188 application to wind profiles is theoretically sound. Linear gravity wave theory dictates that wind and temperature perturbations
 189 are coupled manifestations of the same wave phenomena, and thus the principle of separating smaller-scale waves from the
 190 large-scale background flow via spatiotemporal averaging is equally valid for both fields. Following the arguments presented
 191 in Alexander et al., (2008b), this choice is justified by the need to ensure a sufficient number of profiles per grid cell, which
 192 minimizes random noise while preserving meaningful variability in the data. Shorter temporal windows would lead to
 193 insufficient sampling, while longer windows would smooth out critical small-scale wave features. The grid size is also designed
 194 to preserve the spatiotemporal variability of mesoscale gravity waves and equatorially trapped structures, in an attempt to
 195 separate the background and perturbation components without introducing significant biases.

196
 197 Finally, this configuration mitigates errors in the definition of the $\bar{U}(z)$ profile, ensuring reliable kinetic energy calculations
 198 and robust separation of gravity wave perturbations. We performed sensitivity tests with varying grid sizes and temporal
 199 windows to confirm that this configuration provides the best possible background state when prioritizing Aeolus retrieval (see
 200 Fig. A1 in Appendix A). The average number of profiles used for the background state determination is 55 for Aeolus, 20 for
 201 GNSS-RO and 1400 for ERA5.

202
 203 The next step involves subtracting the background profile from its corresponding individual profile, eliminating most large-
 204 scale waves (Planetary Waves, Kelvin Waves, Rossby Waves). This yields the perturbation profile $U'(z)$, which is then
 205 subjected to Welch-windowing, which is done in order to mitigate spectral leakage (Alexander et al., 2008a; 2008b; Khaykin
 206 et al., 2015). A prior study also applied a similar windowing function (half cosine), aiming to counteract the "effects of the
 207 edge of the height range" (Hei et al., 2008). After said windowing, a band-pass filter designed to retain vertical wavelengths
 208 between 1.5 km and 9 km. is applied to the perturbation profile, as seen in Fig.1b and 1c. The upper limit of 9 km isolates
 209 GWs from larger-scale planetary waves, consistent with our background removal strategy. The lower limit of 1.5 km is chosen
 210 to reflect the effective vertical resolution of the Aeolus instrument (Ratynski et al., 2023) and ensures that our comparison is
 211 restricted to wave scales reliably resolved by all datasets (Banyard et al., 2021). This procedure provides a methodologically
 212 consistent basis for comparing GW energy across the different instruments.

213
 214 The GW Ek can be derived from the variance of wind components as follows:

215

$$216 \quad E_k = \frac{1}{2} (\overline{u'^2} + \overline{v'^2} + \overline{w'^2}), \quad (2)$$

217

218 where u , v , and w represent the zonal, meridional, and vertical wind components, respectively. Considering that Aeolus's
 219 viewing geometry in the tropics makes its HLOS wind primarily sensitive to the zonal component (as shown in Eq. 1), we will
 220 note all mentions of retrieved speed as u for clarity. In our case, since the vertical wind speed is neglected and the satellite is
 221 not able to distinguish between zonal and meridional wind, it is necessary to provide a new formalism for the retrieved metric:

$$222 \quad E_{k \text{ HLOS}} = \frac{1}{2} (\overline{u'^2_{\text{HLOS}}}) \quad (3)$$

223

224 The resulting profile, which is essentially the perturbation squared, is cut to keep the data between one kilometer below the
 225 tropopause and 22 km. The altitude range is chosen considering Aeolus' limitations, such as increasing error at higher altitudes
 226 due to lack of backscatter signal (Ratynski et al., 2023, their Fig.3). The lower bound is set one kilometer below the tropopause
 227 to focus on events extending beyond it, balancing Aeolus' resolution with our interest in upper-end dynamics. For consistency,
 228 the tropopause height is derived directly from the ERA5 dataset for all analyses. The profile is then averaged over the selected
 229 range, representing the E_k , as seen in Fig.1c.

230 We acknowledge that including the layer just below the tropopause presents a potential challenge, as strong, non-wave
 231 divergent outflow from deep convection could be partially aliased into our derived kinetic energy (Stephan et al., 2021). To
 232 rigorously test the robustness of our results against this potential contamination, we have performed a comprehensive
 233 sensitivity analysis by recalculating the kinetic energy fields using two more conservative averaging layers, starting from 1
 234 km and 2 km above the tropopause, respectively (see Fig.B1 and B2 in Appendix B). The detailed results of this analysis
 235 confirm that our main conclusions are not sensitive to this choice.

236

237 Although the above steps focus on retrieving GW E_k from Aeolus wind measurements, the same procedure can be applied to
 238 temperature-based observations such as GNSS-RO for E_p . The main difference lies in substituting temperature $T(z)$ for wind
 239 $U(z)$ throughout the background-perturbation decomposition, which means using $T'(z)$ rather than $U'(z)$. The Welch window
 240 was applied to all perturbation profiles (wind and temperature) before filtering to mitigate spectral leakage. The same band-
 241 pass filtering strategy and vertical averaging then provide the E_p profile from the temperature perturbations. In this case, the
 242 GW E_p is calculated using this formula:

$$243 \quad E_p = \frac{1}{2} \left(\frac{g}{N} \right)^2 \left(\frac{T'}{\bar{T}} \right)^2 \quad (4)$$

244

245 The $E_{k \text{ HLOS}}$ metric derived from Aeolus (Eq. 3) represents the kinetic energy projected onto the instrument's line of sight.
 246 Since our study focuses on the tropical UTLS region, the meridional wind component will have a minor contribution compared

247 to the zonal component. Therefore, the $E_{k\ HLOS}$ energy represents primarily the zonal activity, meaning that we are missing a
248 non-negligible proportion of wave activity. To evaluate the contribution of v' to the total kinetic energy we use ERA5 data
249 and compute the ratio between total E_k (derived from u' and v') and $E_{k\ HLOS}$ (as it is observed by ALADIN).

250

251 This ratio (Fig. C1 in Appendix C1) exhibits significant geographic variability, which can be linked to physical mechanisms
252 that create wave anisotropy. For instance, over regions like the Indian Ocean, the ratio is relatively low (~ 1.5), suggesting a
253 predominantly zonal orientation of wave energy. This is physically plausible, as persistent surface winds like the trade winds
254 can influence the tropopause-level wave field through two main processes. Firstly, flow over orography can preferentially
255 generate zonally-oriented waves (Kruse et al., 2023). Secondly, the background wind profile itself acts as a directional filter,
256 selectively allowing waves propagating in certain directions to reach the UTLS while attenuating others through critical-level
257 interactions (Plougonven et al., 2017; Achatz et al., 2024).

258

259 When averaged over the mission period and focused on the equatorial band (10°S – 10°N), the ratio settles at approximately
260 1.6. This implies that $E_{k\ HLOS}$ accounts for around 62.5% of total E_k , the remainder being undetectable due to HLOS
261 projection. The meridional component, less significant in this specific geographical area for Aeolus, contributes the remaining
262 37.5% of E_k not considered by $E_{k\ HLOS}$. Although not dominant, $E_{k\ HLOS}$ represent a substantial contribution to E_k . This
263 scaling factor is used when discussing the implications of our Aeolus findings for the total GW kinetic energy budget. The
264 details of the spatio-temporal variability of this ratio are provided in the Appendix C.

265

266 The methods employed constrain the analysis to a specific range of horizontal and vertical wavelengths. Aeolus' RBS
267 determines the spacing between sampling points, impacting the vertical and horizontal resolution and maximal detectable
268 wavelength. The vertical wavelength analysis is constrained by the 9 km upper band-pass, representing roughly half the average
269 profile length in the tropics, after limiting the profile to the optimal range and especially considering the dynamic lower bound.
270 Profiles generally extend to heights between 23km and 26km. In the horizontal dimension, since a $20^\circ \times 5^\circ$ degrees grid is
271 used for the background removal and the wind is supposed quasi-zonal, the zonal wavelengths, therefore, reside below 2220
272 km.

273

274 Additionally, Aeolus can be prone to errors alternating the quality of wind profiles. Amongst the most notable ones are dark
275 currents in the charge-coupled devices ("hot pixels"), potentially leading to errors of up to several meters per second (Weiler
276 et al., 2021). Another identified issue is the oscillating perturbations, parasitic deformations of the signal, yet to be attributed
277 to a cause, which can be mistaken for GW-induced signals (Ratynski et al., 2023). While corrections were implemented for
278 the first issue (Weiler et al., 2021), the overall signal random error varies with time, with a general tendency to increase due
279 to instrument degradation. Aeolus' HLOS wind variance is inherently linked to the measurement noise (i.e., random error). In

other words, the observed wind variance is a sum of the variance due to waves (detected using the given data and method) and the variance due to ALADIN noise, i.e., its random error squared.

$$\overline{u'^2_{HLOS}} = \overline{u'^2_{GW}} + \overline{u'^2_{I.N}} \quad (5)$$

283

with $\overline{u'^2_{GW}}$ representing the variance contribution from gravity waves and $\overline{u'^2_{I.N}}$ the contribution from instrument noise. Since kinetic energy is proportional to variance, this relationship holds for kinetic energy as well. The observed Aeolus kinetic energy $\overline{Ek_{Aeolus HLOS}}$ is therefore the sum of the true geophysical signal and a noise component which increases over the mission lifetime. To isolate the true gravity wave energy, this time-varying noise component must be estimated and removed. This correction is performed at the kinetic energy level. While radiosondes provide a valuable independent reference, their sparse coverage in the tropics makes them unsuitable for creating a globally consistent correction field. We therefore use the ERA5 reanalysis as a temporally stable global reference to estimate the Aeolus instrument noise. The core principle is to produce a corrected dataset, denoted as $\overline{Ek_{Aeolus HLOS*}}$ by subtracting our best estimate of the noise energy $\overline{Ek_{I.N}}$ from the observed energy:

$$\overline{Ek_{Aeolus HLOS*}} = \overline{Ek_{Aeolus HLOS}} - \overline{Ek_{I.N}} \quad (6)$$

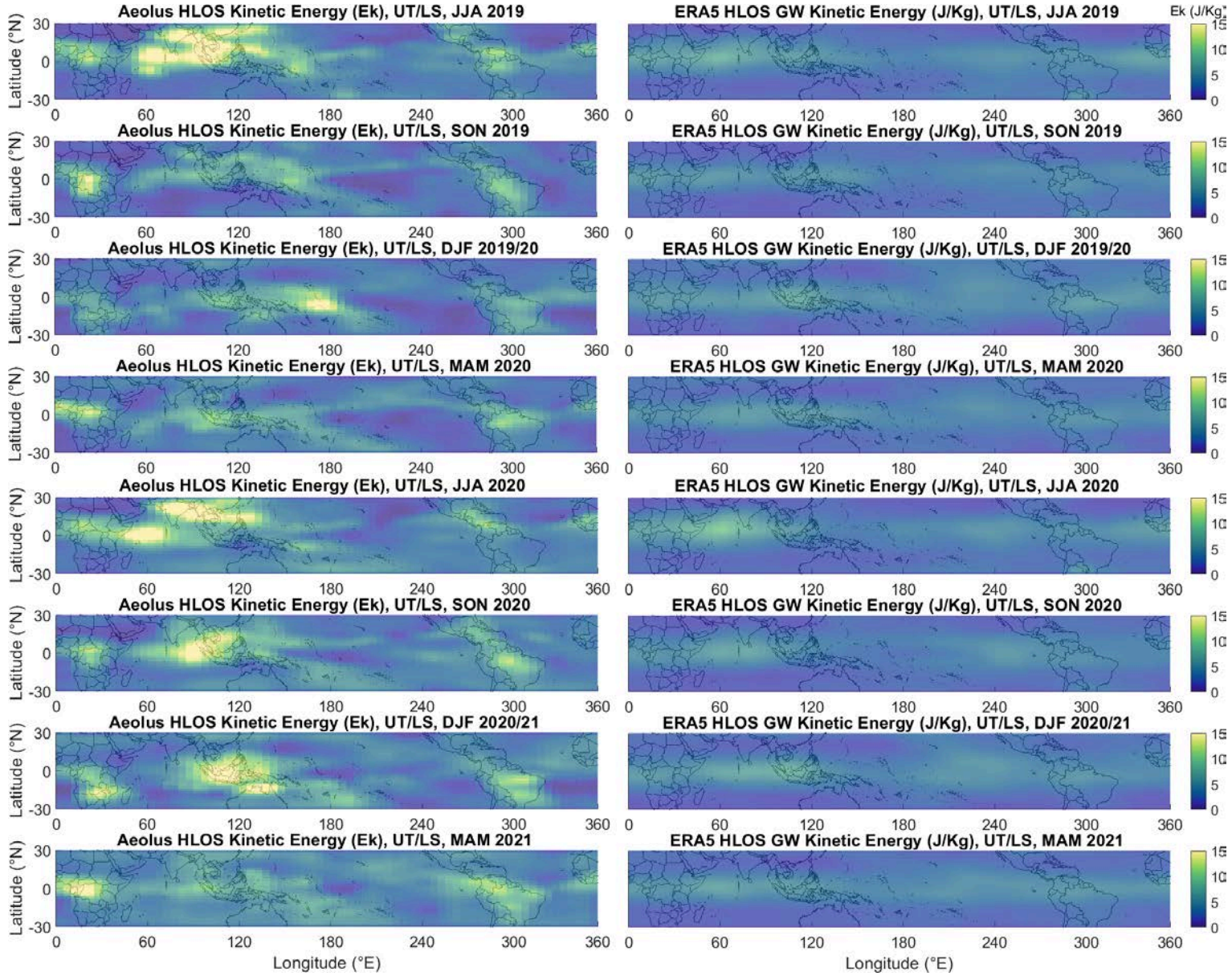
The estimation of the noise term, $\overline{Ek_{I.N}}$, is not a simple subtraction. It is derived using a spatio-temporally adaptive algorithm that blends an additive offset (representing baseline instrument noise, dominant in quiescent atmospheric regions) with a multiplicative scaling factor (more influential in active convective regions where noise effects might scale with the signal). This adaptive approach ensures that instrumental artefacts are removed without suppressing the high-energy gravity wave hotspots uniquely captured by Aeolus, or over-correcting areas of low variance.

An additional refinement step is required for seasonally averaged geographical maps. To produce these maps, individual profile energy values were first binned into a 5° longitude by 2° latitude grid. A second stage adapts the noise correction derived from the tropical 10°S–10°N band for application to the broader latitudinal extent of the maps (e.g., 30°S–30°N), accounting for latitudinal variations in energy and ensuring physically consistent, non-negative results. To reduce noise and highlight large-scale patterns, a 3-point median filter followed by a 3-point moving average filter was applied sequentially in both the zonal and meridional directions.

The full mathematical derivation of the adaptive estimation of $\overline{Ek_{I.N}}$, the details of the map-specific refinement, and a series of diagnostic plots, including comparisons of Aeolus data before and after correction to validate the assumptions made, are provided in the Appendix D.

308 3. Results

309 3.1 Seasonal variation of GW Ek



310 **Figure 2. Comparison between $E_{k_{Aeolus\ HLOS}}$ (left column) and $E_{k_{ERA5\ HLOS}}$ (right column). Each line corresponds to a season, from**
 311 **June-July-August 2019 to March-April-May 2021. The UTLS altitudes are defined between one kilometer below the tropopause and**
 312 **22 km. The tropopause is determined from the ERA5 reanalysis. The maps are smoothed using a combination of median and moving**
 313 **average filters as described in the Methods section.**

314 Figure 2 displays the $E_{k_{HLOS}}$ distribution from the Boreal Spring of 2019 to the Austral summer of 2020, derived from the
315 corrected Aeolus observations and the ERA5 reanalysis. This comparison reveals both key similarities in large-scale patterns
316 and significant differences in the representation of regional wave activity.

317 Both datasets consistently show that the majority of GW kinetic energy is confined to the equatorial belt (approximately 15°S–
318 15°N). This observation aligns with the expectation that deep tropical convection, concentrated within the Intertropical
319 Convergence Zone (ITCZ), is a primary source of the observed waves. A clear seasonal cycle is evident in both Aeolus and
320 ERA5. During Boreal summer (JJA), enhanced Ek is prominent over Central Africa and the Indian Ocean. This corresponds
321 to the active phases of the African and Indian monsoon systems, which provide a persistent, large-scale environment favorable
322 for the development of organized, deep convective systems known to be efficient gravity wave generators (Forbes et al., 2022).
323 During Boreal winter (DJF), the focus of activity shifts eastward towards the Maritime Continent and the Western Pacific,
324 coinciding with that region's primary convective season. These general patterns are also consistent with previous climatologies
325 of GW potential energy derived from temperature measurements Alexander et al. (2008b, their Fig.3 and Fig.4). The GNSS-
326 RO derived Ep values, which range from 0 to 6.6 J/kg at 15 km and 0 to 4.4 J/kg at 22 km (Alexander et al., 2008c), after
327 applying the usual Ek/Ep ratio of 1.6, are generally aligned with our observations.

328

329 It is also necessary to clarify the interpretation of the wave activity observed at the subtropical edges of our analysis domain
330 (near 30°N/S). While our study focuses on convectively generated waves originating from the deep tropics, the kinetic energy
331 measured in the subtropics is likely dominated by different, local sources. The strong subtropical jets and associated frontal
332 systems are potent generators of inertia-gravity waves through mechanisms of geostrophic adjustment and shear instability
333 (Kruse et al., 2023; Plougonven and Zhang, 2014). A recent case study has confirmed that such jet-merging events can produce
334 significant, large-scale GW fields (Woiwode et al., 2023). These jet- and front-generated waves typically have sub-weekly
335 periods and significant wind perturbations, meaning they fall within the detection window of our filtering methodology (Achatz
336 et al., 2024). Therefore, the enhanced energy often visible near 30°N and 30°S in our seasonal maps should be interpreted as
337 stemming primarily from these midlatitude dynamical processes, rather than from the poleward propagation of the equatorial
338 convective waves.

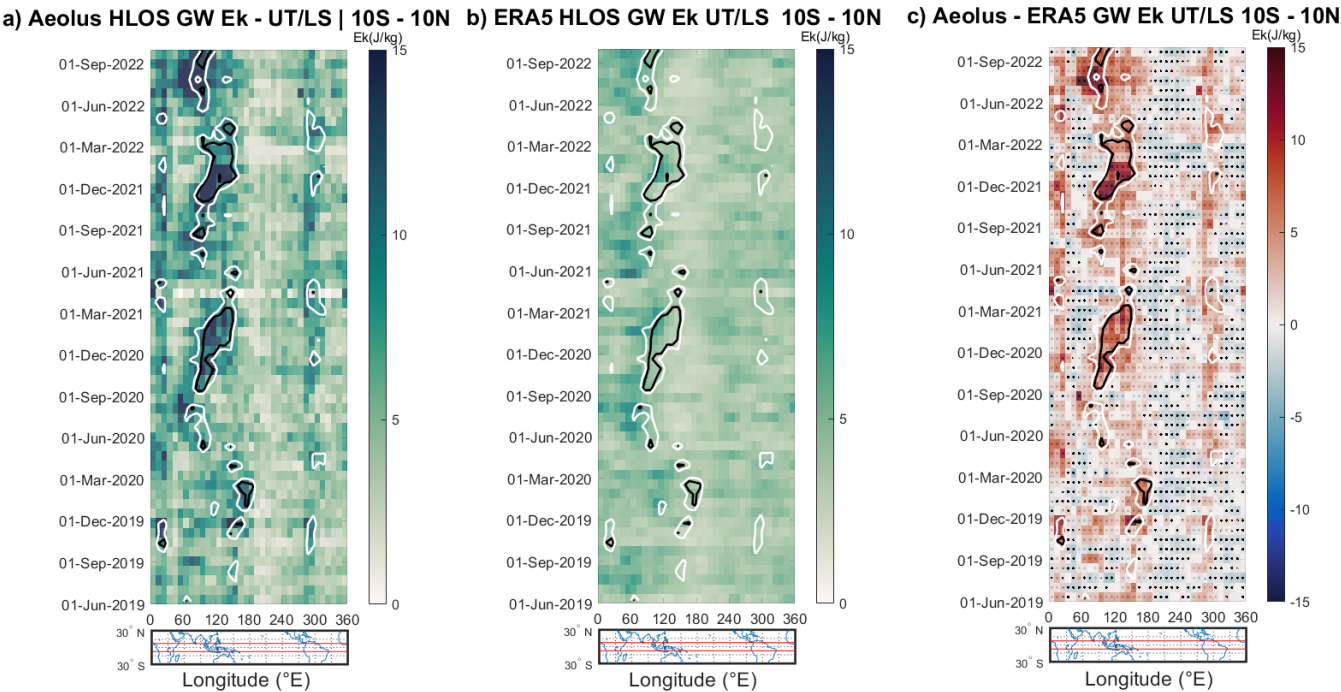
339 Despite these general agreements, a critical difference emerges in the structure and intensity of the energy hotspots. ERA5
340 tends to represent GW activity as a relatively smooth, zonally elongated band, with modest seasonal modulation. In stark
341 contrast, Aeolus reveals a picture of much more localized and intense Ek hotspots. For example, during JJA 2020 and SON
342 2020, Aeolus observes a well-defined hotspot over the Indian Ocean with Ek values exceeding 10-12 J/kg, whereas ERA5
343 shows only a diffuse enhancement in the same region with values rarely exceeding 5-7 J/kg. Similarly, the DJF 2020/21 hotspot
344 over the Maritime Continent is markedly stronger and more geographically confined in the Aeolus data.

345 This discrepancy suggests that while ERA5 captures the broad climatic envelope of convective GW activity, it significantly
346 underestimates the peak energy of waves generated by localized, intense convective systems. This is particularly evident in
347 regions where conventional wind observations are sparse, such as the Indian Ocean. The direct wind profiles from Aeolus
348 appear to capture magnitudes and structures of this convection-driven wave activity that are not present in the reanalysis.

349 The period from mid-2020 onward, which coincided with the development of La Niña conditions, exhibits the most
350 pronounced differences between the two datasets. While La Niña is known to enhance convection over the Maritime Continent,
351 the consistently higher energy levels observed by Aeolus across all regions during this period also correlate with a documented
352 increase in the satellite's instrumental random error (Ratynski et al., 2023, their Fig.6). Our adaptive noise correction (see
353 Appendix D) is designed to account for this degradation. However, it is challenging to perfectly disentangle the increased
354 geophysical signal (e.g., from La Niña) from the effects of increased instrument noise. Nevertheless, the geographical
355 consistency of the hotspots observed by Aeolus, which align with known convective centers, provides confidence that the
356 primary patterns represent true atmospheric phenomena that are underrepresented in the reanalysis.

357 Finally, regarding the strong latitudinal confinement of the signal, while this is primarily a physical feature, our noise-
358 correction methodology may also contribute to it. As detailed in the Appendix D (Part 2), the correction is weighted by the
359 latitudinal structure of the raw signal. This approach, designed to avoid over-correction in low-signal subtropical regions,
360 naturally sharpens the latitudinal gradient at the edges of the tropical belt.

361



365 **Figure 3. (a,b,c) Hovmöller diagram of $Ek_{\text{Aeolus HLOS}}$, $Ek_{\text{ERA5 HLOS}}$ and their difference. The contour plot represents the Outgoing**
366 **Longrange Radiation (OLR) for 210 and 220 W/m^2 (black and white, respectively). Each bin corresponds to an average of over 3**
367 **weeks and 10 degrees. The white bins represent the lack of satellite information in (a). The OLR measurements were obtained from**
368 **the Australian Bureau of Meteorology. The UTLS altitudes are defined between one kilometer below the tropopause and 22 km.**
369 **The tropopause is determined from the ERA5 reanalysis. Black stippling indicates regions where the difference between quantities**
370 **is statistically significant (two-sample t-test, $p < 0.05$).**

372 To assess the evolution and transition between the different seasons with greater precision, the Hovmöller diagrams in Fig.3
373 only show the observations within the deep tropics between 10° N and 10° S, as Fig.2 proves this region contains most of the
374 activity. Fig.3a shows $Ek_{\text{Aeolus HLOS}}$, where prominent hotspots of high Ek (often attaining 15 J/kg) are visible, with a broad
375 region of enhanced activity migrating eastward from the African continent (~0-60°E) towards the Indian Ocean and Maritime
376 Continent (~60-150°E) between June and March. This shift is recurring over multiple years and shows a relative consistency
377 between each year in terms of longitudinal and temporal range.

379 The presence of hotspots, represented by distinct shapes in the Ek patterns, is expected in regions with prevalent convective
380 activity. These can be attributed to multiple powerful wave generation mechanisms occurring at the scale of individual storms.
381 One primary mechanism is thermal forcing, where the pulsatile nature of latent heat release in a convective updraft acts like a
382 piston on the surrounding stable air, generating a broad spectrum of gravity waves (Beres et al., 2005). A second,

383 complementary mechanism is mechanical forcing, where the body of the strong updraft itself acts as a physical barrier to the
384 background wind. The flow forced over this "moving mountain" generates large-amplitude, low-phase-speed waves that are
385 stationary relative to the storm (Corcos et al., 2025; Wright et al., 2023). The intense kinetic energy observed by Aeolus is
386 likely the signature of both mechanisms operating within active convective systems. These convectively generated GWs can
387 propagate vertically and interact with the large-scale atmospheric circulation, transferring momentum and energy to the
388 background flow (Alexander et al., 2021).

389 In sharp contrast, Fig.3b presents a much more subdued and less dynamic picture for the corresponding ERA5 perspective.
390 While ERA5 shows some broad, low-amplitude enhancement of Ek that co-locates weakly with the seasonal convective cycle,
391 it completely fails to capture the intense, high-energy hotspots observed by Aeolus. The organized eastward propagation and
392 the high peak energy values are entirely absent in the reanalysis. For nearly all regions and time periods, the Ek in ERA5
393 remains below 7 J/kg.

394 The difference between the two datasets, shown in Fig.3c, quantifies this discrepancy. The plot is overwhelmingly positive,
395 indicating a systematic and significant underestimation of GW kinetic energy by ERA5 throughout the tropics. The regions of
396 greatest underestimation, where the difference exceeds 10 J/kg, align almost perfectly with areas of deep convection, as
397 identified by the low Outgoing Longwave Radiation (OLR) contours. The OLR represents the amount of terrestrial radiation
398 released into space and, by extension, the amount of cloud cover and water vapor that intercepts that radiation in the
399 atmosphere. It is a widely used and reliable proxy for deep convection due to its strong correlation with diabatic heating (Zhang
400 et al., 2017), reinforcing the conclusion that ERA5's primary weakness lies in representing convection-driven wave activity.

401 To confirm the robustness of this finding, a two-sample t-test was performed for each grid cell. The stippling in Fig.3c indicates
402 where the mean Ek from Aeolus is statistically significantly higher than that of ERA5 ($p < 0.05$). The pervasive stippling
403 across nearly all convective hotspots underscores that the observed differences are not random fluctuations but represent a
404 fundamental deficiency in the reanalysis. This finding suggests that without the assimilation of direct, high-resolution wind
405 profile data like that from Aeolus, reanalysis models may not fully resolve the kinetic energy associated with gravity waves
406 generated by localized, intense convective events. An alternative display of Fig.3c as a ratio, along with an F-test, can be found
407 in Appendix E.

408

409 **4. Comparison with Potential Energy: possibilities and limitations**

410 While a conservative analysis might prioritize directly comparable metrics, exploring less conventional approaches can reveal
411 patterns that remain hidden in traditional frameworks. A key goal of this study is to compare the kinetic energy from Aeolus

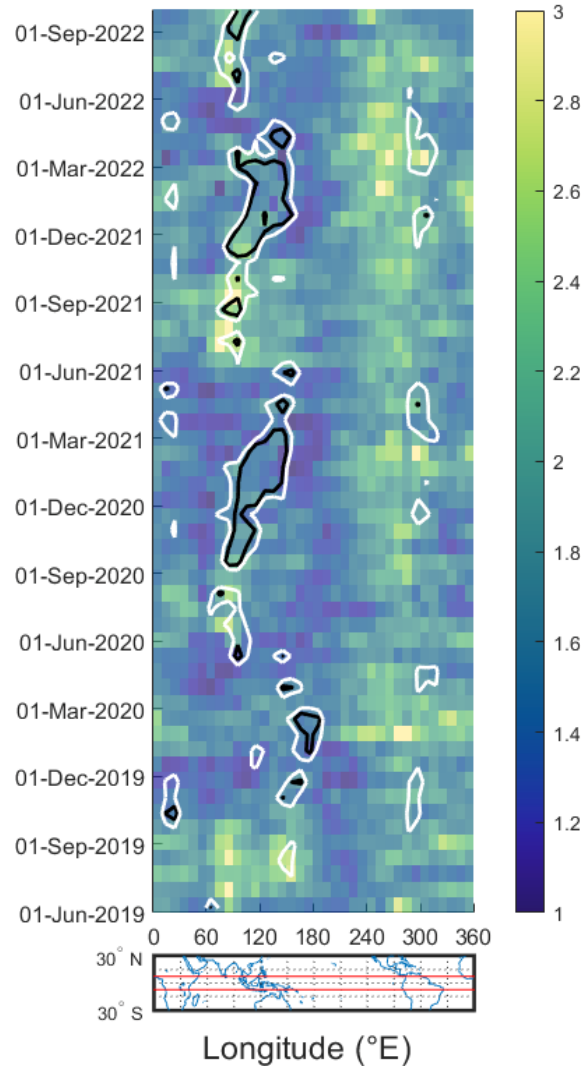
412 with potential energy, the most common metric for GW climatologies. However, this comparison is complicated by the
413 longstanding assumption of a constant E_k/E_p ratio.

414 Traditionally, linear gravity wave theory proposes a near-constant ratio of E_k to E_p , often quoted between 5/3 and 2.0 (Hei et
415 al., 2008; VanZandt, 1985). In stable, linear wave conditions, these theoretical predictions hold reasonably well (Nastrom et
416 al., 2000). However, a growing body of evidence from empirical studies reveals significant variability in this ratio, suggesting
417 that real-world atmospheric conditions often involve non-linear processes such as wave breaking or saturation, which are not
418 accounted for in the linear theory (Baumgarten et al., 2015; Guharay et al., 2010; Tsuda et al., 2004).

419

420

ERA5 EK/EP UT/LS 10S-10N



421

422 **Figure 4. Spatiotemporal distribution of the Ek/Ep ratio in the ERA5 reanalysis for the equatorial band (10°S–10°N). The UTLS**
 423 **altitudes are defined between one kilometer below the tropopause and 22 km. The contour plot represents the Outgoing Longrange**
 424 **Radiation (OLR) for 210 and 220 W/m² (black and white, respectively). Each bin corresponds to an average of over 3 weeks and 10**
 425 **degrees. The OLR measurements were obtained from the Australian Bureau of Meteorology. The tropopause is determined from**
 426 **the ERA5 reanalysis.**

427 To illustrate this complexity within a self-consistent framework, we first examine the Ek/Ep ratio derived entirely from the
 428 ERA5 reanalysis. Figure 4 presents the longitudinal and temporal variations of this ratio in the equatorial UTLS. The figure
 429 immediately reveals that the ratio is far from constant. It exhibits significant spatial and temporal variability, with values
 430 frequently exceeding the linear theory predictions (>2). Notably, distinct hotspots of high Ek/Ep ratios are present, particularly
 431 over the Indian Ocean at around 70° and the South American continent at 300°.

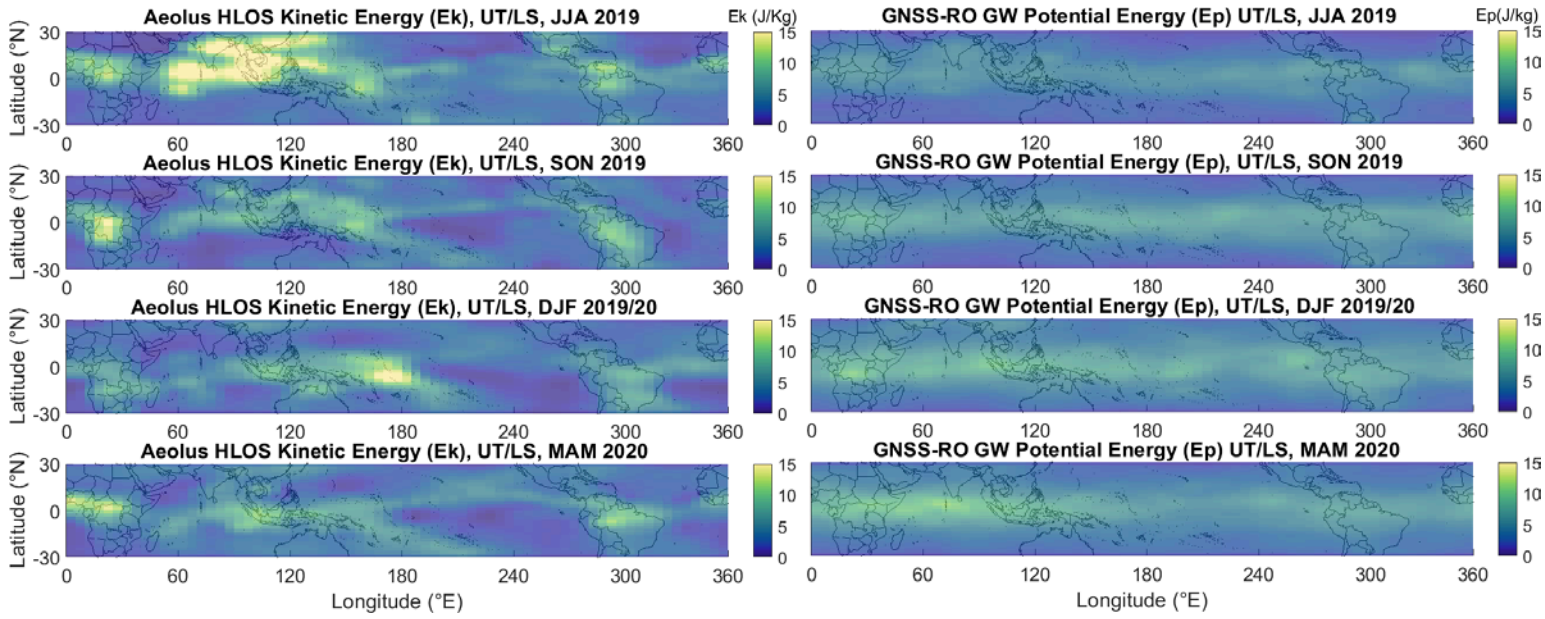
432 In regions outside of the most intense convective activity, where ERA5 does manage to represent some kinetic energy
433 enhancement, the agreement with Aeolus is often satisfactory. This is visible in Fig. 3c, where the same areas (120° and 300°)
434 show a correct correspondence. This suggests that when wave generation is not dominated by deep convection, ERA5 can
435 reproduce realistic GW structures. The strong agreement in these non-convective regions also reinforces the idea that the
436 dominant winds have a strong zonal component, as the quasi-zonal u_{HLOS} measurement from Aeolus is sufficient to capture
437 these features. The elongated white stripe during February–March 2020 comes from an intense intraseasonal disturbance, the
438 2020 Madden-Julian Oscillations (MJO), which can inject unusually strong gravity-wave energy into the upper troposphere
439 (Kumari et al., 2021).

440 The divergence between ERA5 and Aeolus becomes most pronounced precisely in the deep convective regions where Aeolus
441 observes its strongest Ek signals, inside the areas of low OLR. The discrepancy appears specifically linked to convection-
442 driven dynamics, which are either not properly represented or fail to trigger sufficient wave activity within the ERA5 model's
443 parameterizations. This suggests that the primary cause of ERA5's underestimation is not a simple mispartitioning between
444 the horizontal wind components (i.e., a directional bias in the line-of-sight projection) but rather a more fundamental, large-
445 scale underestimation of the total kinetic energy.

446 This model-internal result demonstrates that relying on a fixed ratio to infer one energy component from another is likely to
447 be inaccurate, especially in dynamically active regions. The partitioning of energy between kinetic and potential forms is itself
448 a key diagnostic of wave dynamics that requires further observational constraints.

449 One promising possibility of this study in providing deeper context lies in comparing the kinetic energy of gravity waves
450 observed by Aeolus with the potential energy derived from GNSS-RO data. GNSS-RO provides high-resolution temperature
451 profiles that are used to estimate the potential energy of gravity waves. Previous studies that looked into GW climatology all
452 relied on these estimate to base their observations on, as it was the only global instrumentation available (Schmidt et al., 2008;
453 Alexander et al., 2008b; Šácha et al., 2014; Khaykin et al., 2015). Hence, we will adopt this method of comparison as well.

454



455

456 **Figure 5. Comparison between $E_{k_{\text{Aeolus HLOS}^*}}$ (left column) and E_p GNSSRO (right column). Each line corresponds to a season,**
 457 **from June-July-August 2019 to March-April-May 2020. The white bins represent the lack of satellite information. The UTLS**
 458 **altitudes are defined between one kilometer below the tropopause and 22 km. The tropopause is determined from the ERA5**
 459 **reanalysis. The maps are smoothed using a combination of median and moving average filters as described in the Methods section.**

460 Fig.5 offers a side-by-side seasonal comparison of $E_{k_{\text{Aeolus HLOS}^*}}$ (left column) and E_p derived from GNSS-RO (right column),
 461 covering the period from June 2019 to May 2020. The figure highlights key spatial and temporal patterns of gravity wave
 462 activity detected by each instrument, with both datasets presenting clear seasonal variability.

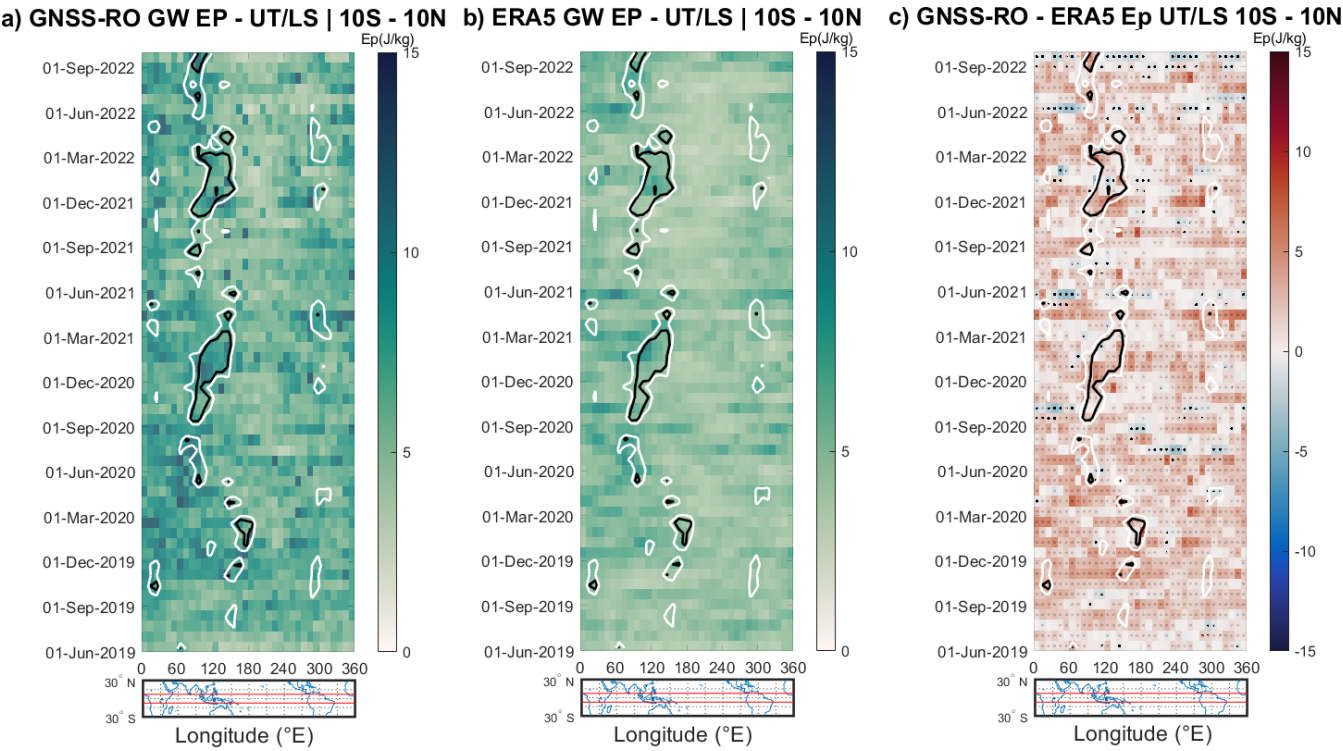
463

464 Although the ratios between E_k and E_p suggested by linear gravity wave theory generally range between 5/3 and 2.0, empirical
 465 observations show significant variability. This variability, which is influenced by geographical factors, nonlinear processes, or
 466 wave interactions, underscores the importance of examining these two forms of energy from different perspectives rather than
 467 seeking strict correspondences.

468

469 With that in mind, what stands out from this comparison is the overall consistency in detecting gravity wave hotspots,
 470 particularly within the tropical belt (The African land convection and Indian Ocean hotspot are consistent on both instruments).
 471 One notable aspect of the comparison is the seasonal shift in gravity wave activity between the two datasets, with both detecting
 472 enhanced wave activity during certain months (increased activity levels in DJF and MAM 2020). Because of inherent
 473 differences (different line of sight and signal projection, different physical quantities and their varying ratio that is empirically
 474 challenging the literature, different signal treatment and correction), direct one-to-one comparisons are not appropriate.
 475 Nonetheless, it allows us to draw parallels with Aeolus observations, where spatial and temporal correlation of hotspots should

476 follow the same disposition, allowing for an independent benchmark. Despite these methodological differences, both
 477 instruments align on the seasonal peaks and general distribution of wave activity, reinforcing the reliability of the data.
 478



479
 480
 481 **Figure 6. (a,b,c) Hovmöller diagram of $Ep_{\text{GNSS-RO}}$, Ep_{ERA5} and their difference. The contour plot represents the Outgoing**
 482 **Longrange Radiation (OLR) for 210 and 220 W/m^2 (black and white, respectively). Each bin corresponds to an average of over 3**
 483 **weeks and 10 degrees. The OLR measurements were obtained from the Australian Bureau of Meteorology. The UTLS altitudes are**
 484 **defined between one kilometer below the tropopause and 22 km. The tropopause is determined from the ERA5 reanalysis. Black**
 485 **stippling indicates regions where the difference between quantities is statistically significant (two-sample t-test, $p < 0.05$).**

486 The $Ep_{\text{GNSS-RO}}$ shown in Fig.6a does not closely follow the patterns of OLR activity. As the method employed removes most
 487 traces of kelvin waves in the signal, the remaining activity is mostly comprised of GWs. This suggests that Ep does not
 488 effectively capture GW activity in regions of deep convection, as indicated by the lowest OLR values. However, it is found
 489 that the lesser convective areas are seen both on instances of Ek and Ep , in Fig.3a and Fig.6a (with notable examples such as
 490 August 2020 around 100°E , as well as in May 2021 and 2022 near 50°E). This observation supports the notion that, in terms
 491 of GW activity, deep convective phenomena primarily generate Ek , while less intense convective events (indicated in Fig.6a
 492 as occurring in the neighbouring region outside the white contours) produce a more balanced distribution between both energy
 493 components. It would be incorrect to assume that no wave activity occurs in low OLR regions; previous studies have shown
 494 that Ep values peak at 15 km altitude around the maritime continent, where the Walker circulation rises under non-El Niño
 495 conditions (Ern et al., 2004; Yang et al., 2021).

496
497
498 Nonetheless, the $E_{p_{ERA5}}$ diagram shown in Fig.6b is generally consistent with the results shown in Fig.6a, if one admits that
499 the instrumental signal is prone to more noise and higher average values. Particularly in regions outside the primary convection
500 hotspots, in August 2020 around 100°E, we see coherent signals in both datasets. Similarly, in May 2021 near 50°E or in
501 February 2022 near 120°E, distinct patterns emerge in both datasets. These alignments indicate that when gravity waves have
502 a stronger potential energy component, both datasets capture these features, even outside the primary zones of low OLR. It
503 can also be noted that the patterns visible in Fig.6b strongly resemble the patterns presented by ERA5 in Fig.3b, a sign of
504 ERA5's tendency to rely on the existence of E_p to determine the presence of E_k .
505
506 The differences between ERA5 and GNSS-RO data, depicted in Fig.6c, are minimal, with a mean absolute difference of
507 approximately 1.5 J/kg. This close agreement is expected, given that ERA5 assimilates GNSS-RO temperature profiles. This
508 result demonstrates that the reanalysis system successfully reproduces the observed potential energy field. This finding stands
509 in sharp contrast to the significant underestimation of kinetic energy by ERA5 when compared to Aeolus observations (as
510 shown in Fig 4c). The combination of these results (good agreement in assimilated E_p and poor agreement in unassimilated
511 E_k) points to a specific limitation in the reanalysis's ability to model the kinetic component of convection-induced gravity
512 waves, particularly in regions where direct wind observations are sparse.
513
514 An alternative display of Fig.6c as a ratio, along with an F-test, can be found in Appendix E.

Aeolus Ek / GNSS-RO Ep | UT/LS | 10S - 10N

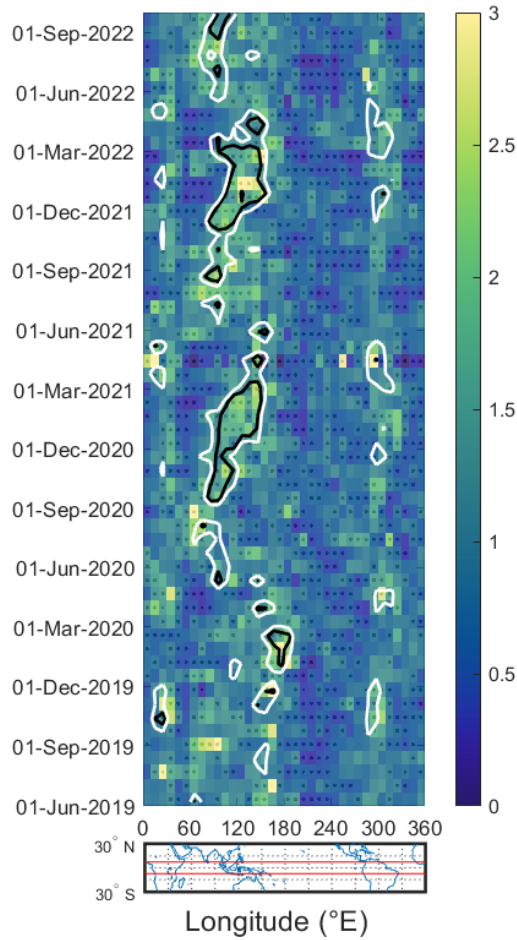


Figure 7. Relationship between $Ek_{\text{Aeolus HLOS}^*}$ and Ep from GNSSRO. The UTLS altitudes are defined between one kilometer below the tropopause and 22 km. The tropopause is determined from the ERA5 reanalysis. The tropopause is determined from the ERA5 reanalysis. Black stippling indicates regions where the difference between quantities is statistically significant (F-test, $p < 0.05$).

Fig. 7 presents the first observationally-derived long-term study of the Ek/Ep ratio, comparing Aeolus's HLOS Ek and GNSS-RO-derived Ep. It illustrates the longitudinal and temporal variations of the Ek/Ep ratio across the equatorial band (10°S to 10°N) from June 2019 to October 2022.

The regions with the highest ratio values are systematically co-located with areas of deep convection, as indicated by the low OLR contours. This is particularly evident over the Indian Ocean (e.g., September-June 2019/20, 2020/21, and 2021/20) and over the Western Pacific. This observation suggests that, in areas with similar seasonal characteristics, gravity waves tend to transport more kinetic energy during convective events, which amplifies their influence on the overall energy dynamics. The

528 periodic patterns observed in the data also hint at a seasonal component previously observed by Zhang et al. (2010), potentially
529 tied to atmospheric phenomena such as the shifting ITCZ or changes in jet stream dynamics (Hei et al., 2008). These seasonal
530 fluctuations in the Ek/Ep ratio further reinforce the notion that gravity wave behavior is not static but is influenced by broader
531 atmospheric cycles (Ern et al., 2018; Zhang et al., 2010), contrary to the traditional linear theory paradigm in the literature.
532 Statistical significance testing (represented by the black stippling) confirms that these hotspots of high, convection-linked
533 ratios are statistically significant features rather than artifacts.

534 A significant division between the Indian Ocean and the eastern Pacific, marked by a contrast around 200° longitude, can be
535 noted in both Fig.7 and Fig.4. This contrast reflects underlying geographic factors, including orographic influences and
536 convective activity. These two factors play a role in the generation and propagation of gravity waves, causing the distinct
537 variations in the ratio between the two energies.

538 This observational result stands in contrast to the picture presented by the ERA5 reanalysis in Fig.4. While ERA5 also shows
539 variability in its Ek/Ep ratio, its regions of highest ratio are often located outside the main convective centers. This suggests
540 that ERA5 misrepresents the physical link between deep convection and the partitioning of wave energy.

541 Given that ERA5 successfully assimilates GNSS-RO temperature data (and thus has a reasonable representation of Ep), this
542 discrepancy points to a fundamental difficulty in the reanalysis's ability to generate the corresponding kinetic energy
543 component in the right locations. Without direct wind profile assimilation in these data-sparse convective regions, the model's
544 parameterizations and background error covariances fail to create the intense, localized kinetic energy associated with
545 convective gravity waves.

546 However, it is noteworthy that in some specific regions and periods, such as over the Indian Ocean between June and
547 September of 2019, or in the longitude band around 300°E, a degree of correspondence between the model and observations
548 can be found. This suggests that for certain regimes, the reanalysis can approximate the energy partitioning, but it fails
549 systematically in the most intensely convective areas. These findings reinforce that direct kinetic energy measurements, as
550 provided by Aeolus, are essential for correcting model biases and improving our understanding of the gravity wave energy
551 budget.

552

553 **5. Discussion**

554 Overall, the results presented in this study allow us to discuss and address two main questions. The first consistent observation
555 made, was that ERA5 underestimates Ek distribution in such regions compared to the Aeolus-derived energy, particularly over

556 the Indian Ocean, where conventional radiosonde wind measurements are very sparse. That difference raised questions on the
557 potential reason for such discrepancies: Is this result an overestimation of Aeolus, due to its known increased noise and
558 inconsistent performance during its life-cycle, or an underestimation for ERA5, due to the lack of direct wind observations
559 assimilated?

560 The analysis of ALADIN wind profiling and ECMWF ERA5 reanalysis data, provided in Fig.2 and Fig.3, revealed enhanced
561 GW activity over the Indian Ocean during Boreal Summer, as well as over the western Pacific and maritime continent in Boreal
562 Winter. The enhanced GW activity migrating from eastern Africa to the Pacific maritime continent between June and
563 December is linked to convection, as suggested by the correlation between enhanced GW Ek and the regional minima in OLR.
564 The slow eastward propagation of these energy maxima suggests that the underlying wave sources are not random, but are
565 organized by planetary-scale atmospheric patterns. Indeed, the relation between OLR and the Madden-Julian Oscillation
566 (MJO) has been used before; It is a reliable index for analysis (Kiladis et al., 2014), and recent work has provided direct
567 observational evidence that the MJO modulates GW activity and momentum transport from the tropics to higher
568 latitudes (Zhou et al., 2024). The structures observed by Aeolus are therefore highly consistent with the kinetic energy signature
569 of gravity waves generated by the powerful thermal and mechanical forcing mechanisms (Beres et al., 2005; Corcos et al.,
570 2025) occurring within the large, organized convective superclusters of the MJO. This suggests that Aeolus is effectively
571 capturing convection-induced GWs that may be underrepresented in ERA5. One of the persistent features observed throughout
572 the study was the high-energy gravity wave hotspot over the African continent, which remained consistent across seasons and
573 years. This suggests a continuous mechanism of continental convection driving gravity wave activity in this region.

574 Another consideration in this study is whether the large Ek values observed by Aeolus, particularly over convective hotspots,
575 could be an artifact of misinterpreting non-wave tropospheric outflow rather than stratospheric gravity waves. Our sensitivity
576 analysis (see Fig. B1 and B2 in Appendix B) directly refutes this concern. By shifting the analysis layer upward to begin 1 km
577 and 2 km above the tropopause, we confirm that the geographical patterns of the energy hotspots are remarkably stable (spatial
578 correlation $r > 0.83$), and that the vast majority of the peak energy (~88-91%) persists well into the stratosphere. If the signal
579 were dominated by shallow tropospheric outflow, the energy peaks would have collapsed when the analysis layer was moved
580 above the tropopause. The fact that a strong, structured signal remains provides compelling evidence that we are observing
581 vertically propagating gravity waves that have penetrated the lower stratosphere. This validates our central conclusion: Aeolus
582 is capturing a significant field of convectively-generated stratospheric gravity wave kinetic energy that is largely absent in the
583 ERA5 reanalysis.

584 An additional tool at our disposal to solve the case is the global distribution of Ep, through the use of independent GNSS-RO
585 instruments. Our analysis confirms that the assimilation of GNSS-RO data in ERA5 is highly effective, with minimal
586 discrepancies observed between the reanalysis Ep and direct GNSS-RO observations (Fig.6c). This key finding allows us to

587 arbitrate between two potential causes for the Ek discrepancy: a lack of direct wind data assimilation versus inherent biases in
588 the model's physics (e.g., its GWD parameterization).

589 Several lines of evidence from our study point towards the lack of wind assimilation as the dominant cause. Firstly, the fact
590 that ERA5 accurately reproduces E_p fields demonstrates that the underlying model can represent the thermodynamic
591 signatures of wave activity when properly constrained. Conversely, the largest discrepancies are found in kinetic energy, a
592 purely wind-based quantity, and are concentrated over data-sparse regions like the Indian Ocean, precisely where Aeolus
593 provides direct wind profile measurements not available from other observing systems (Banyard et al., 2021).

594 Secondly, while ERA5's non-orographic GWD scheme has known limitations and is not directly forced by diagnosed
595 convection (Orr et al., 2010), it is unlikely to be the sole reason for the missing Ek. Such a parameterization bias would be
596 expected to manifest as a systematic error across different variables or regions, or as a persistent model drift requiring large,
597 ongoing corrections by the assimilation system (Dee, 2005). However, our findings show a targeted deficiency: the model
598 performs well on assimilated temperature (E_p) but poorly on unassimilated wind (Ek) in the very same locations. This sharp
599 contrast strongly suggests the problem is not a wholesale failure of the model's physics to generate wave energy, but rather its
600 inability to correctly partition that energy into kinetic and potential components without direct wind constraints.

601 In data-sparse areas, ERA5 must rely on its internal background error covariances to infer wind adjustments from the
602 assimilated mass field (Hersbach et al., 2020). These statistical relationships are primarily designed to represent large-scale,
603 quasi-balanced (rotational) flow and are known to be less effective at specifying the smaller-scale, divergent component of the
604 wind field to which convectively generated gravity waves belong, especially in the tropics (Žagar et al., 2004). Consequently,
605 while the assimilation of GNSS-RO constrains the thermodynamic (E_p) aspect of the wave, the system lacks the necessary
606 information and dynamic constraints to generate the corresponding divergent wind perturbations, leading to the
607 observed Ek deficit. This process evidently fails to capture the full spectrum of high-Ek wave modes generated by convection.

608 Overall, the findings presented here are in full agreement with the elements outlined in the introduction, suggesting that ERA5
609 is underestimating the Ek component. Indeed, ERA5 has several known shortcomings, such as its underrepresentation of
610 eastward-propagating inertio-gravity waves (Bramberger et al., 2022), its site-dependent errors in tropical regions (Campos et
611 al., 2022), and the broader limitations of data assimilation systems in capturing circulation dynamics, particularly in areas with
612 sparse wind observations (Podglajen et al., 2014; Žagar et al., 2004). These challenges are further emphasized by the QBO-
613 MJO modulation of the wave activity in the UTLS, possibly at play in the amplitude variability of the seasonal signals seen in
614 Fig.3a, where some years show an increased activity compared to others. An observational study revealed how gravity waves
615 generated during MJO phases interact with QBO-modulated wind patterns, influencing their dissipation and energy dispersion
616 (Kalisch et al., 2018). The QBO easterly phase (EQBO) has been shown to enhance MJO activity by strengthening convective
617 signals and reinforcing the propagation of Rossby and Kelvin waves in the UTLS, while the westerly phase (WQBO)
618 suppresses these dynamics (Song and Wu, 2020; Martin et al., 2021). However, the limitations of ERA5 in capturing QBO-

619 MJO interactions are evident, as reanalysis datasets often fail to fully reproduce the observed temperature and wind anomalies
620 associated with these processes, particularly in tropical regions (Lim and Son, 2022).

621 Another discussion enabled by Aeolus observations concerns the longstanding assumption of a constant E_k/E_p ratio in GW
622 studies. Specifically, the question arises: Is the conventional view of a constant ratio for inferring E_p from E_k (and vice versa)
623 still tenable? Or do the new data suggest that this ratio is no longer universally valid in real-world, often non-linear, atmospheric
624 conditions?

625 At first glance, using a fixed ratio appears straightforward for converting well-documented E_p (from temperature-based
626 instruments such as GNSS-RO) to E_k . Traditionally, linear GW theory proposes a near-constant ratio of E_k to E_p , often quoted
627 between 5/3 and 2.0 (VanZandt, 1985; Hei et al., 2008). In idealized models of linear wave behavior, the kinetic and potential
628 energies are expected to be comparable, leading to a ratio close to unity. This has been confirmed previously, in a study such
629 as Nastrom et al. (2000), which found that in stable, linear wave conditions, the energy ratios adhered closely to these
630 theoretical predictions.

631 However, a growing body of evidence challenges this simplification: Empirical work increasingly reveals significant
632 variability in this ratio, indicating non-linear effects in real-world atmospheric conditions (Baumgarten et al., 2015; Guharay
633 et al., 2010; Tsuda et al., 2004). When the observed energy ratios deviate significantly from this expected range, non-linear
634 processes may be at play. For instance, in situations where wave amplitudes are particularly large, wave-wave interactions,
635 such as those resulting from wave breaking or saturation, could lead to the observed discrepancies. This has been demonstrated
636 in earlier work by Mack and Jay. (1967), who found that under certain conditions, potential energy deviated markedly from
637 kinetic energy, suggesting non-linear effects. Similar findings have been reported by Fritts et al. (2009), who showed that
638 interactions between gravity waves and fine atmospheric structures can result in turbulence, thereby affecting the balance
639 between kinetic and potential energy.

640 With everything in place to link these elements, the observed comparison in Fig.4 of the E_k/E_p ratios from ERA5, Aeolus, and
641 GNSS-RO confirms that the characteristics of gravity waves vary significantly across time and space. The observed ratios, 1.7
642 (+/- 0.38) for ERA5, 1.4 (+/- 0.54) for Aeolus/GNSS-RO, indicate that the waves encompass both linear and non-linear
643 processes. The frequent observation of ratios exceeding unity, aligning with trends identified in previous studies, suggests that
644 a substantial portion of the waves' energy is contained in kinetic form, often indicative of non-linear behavior. Because the
645 assumption of a constant ratio is increasingly challenged by empirical observations, it accentuates the need to shift the paradigm
646 from relying solely on temperature perturbations to directly deriving E_k . As such, directly measuring kinetic energy is a major
647 missing link for a comprehensive understanding of GW dynamics.

Beyond these considerations of gravity wave dynamics and energy ratios, we should also acknowledge the limitations of the Aeolus satellite. These include both its technical shortcomings and the constraints imposed by its HLOS projection, which directly impact the representativeness of its measurements. A 1.6 ratio was determined for $E_k/E_{k_{HLOS}}$ using ERA5, as seen in Fig.2. It reflects the efficiency with which HLOS winds from Aeolus can approximate the full kinetic energy field. The ratio indicates that HLOS winds account for approximately 62.5% of the total E_k , while the remaining 37.5% is undetectable due to the projection limitations of HLOS measurements. The discrepancy suggests that the HLOS winds alone cannot fully capture the energy contributions from multi-dimensional wave dynamics. However, this ratio can help estimate the full E_k indirectly with reasonable accuracy. While this approach introduces some assumptions, it can be further refined by cross-validating against comprehensive datasets from reanalyses like ERA5.

The spatial and temporal consistency observed between Aeolus and ERA5 datasets highlight the potential of Aeolus wind profiling for assimilation to improve our understanding of atmospheric dynamics in the tropical UTLS. However, there are limitations and uncertainties to consider; these include dark currents in its charge-coupled devices, known as "hot pixels", which can induce speed measurement errors of several meters per second (Weiler et al., 2021). Additionally, oscillating perturbations within the instrument can cause signal deformations, potentially misinterpreted as GW-induced signals (Ratynski et al., 2023). Even though corrections have been implemented, Fig.3 showed that the instrument's overall signal random error still fluctuates due to its degradation. While newer baselines such as the latest 2B16 used in this study, improve on the appearance rate of hot pixels, challenges remain in consistently identifying and correcting other issues such as oscillating perturbations, increased solar activity, or cloud contamination that can sporadically deform signals.

Understanding the vertical wavelength of convective GWs is an essential element for characterizing their dynamics. However, Aeolus is inherently limited in retrieving accurate vertical wavelengths due to its design. The placement of range bins was fixed at the time of observation, introducing inconsistencies in vertical resolution that affect the precise identification of wave peaks and troughs. Additionally, the N/P parameter, which controls the number of accumulated measurements (N) and pulses (P) per cycle, introduces variability in the horizontal resolution of Aeolus data. Changes to this setting, such as the transition from N=30 to N=5, improve horizontal resolution but exacerbate the misrepresentation of vertical wave structures. Furthermore, any spectral analysis of a finite vertical profile is inherently constrained. For geophysical spectra that are typically having more variance at longer wavelengths, a simple peak-finding method would likely identify a dominant wavelength that is an artifact of the analysis window or filtering choices. Given these limitations, we limit our analysis to the vertically-integrated energy within a defined passband (vertical wavelengths < 9 km), which is a more robust quantity.

Nevertheless, we can speculate that the high E_k values observed by Aeolus in convective regions are associated with shorter-wavelength waves. This interpretation is consistent with established physical mechanisms which state that waves with high E_k are typically generated in regions with strong convective updrafts and downdrafts, where the rapid vertical movement of

681 air masses creates intense small-scale disturbances. These localized and transient disturbances, arising from geostrophic
682 imbalance, generate GWs that carry energy away from the convective region, where strong forcing efficiently transfers energy
683 into the EK spectrum at shorter wavelengths (Waite and Snyder, 2009). The correlation between high EK and shorter
684 wavelengths is particularly pronounced in convective systems, as confirmed in both observational and numerical estimations
685 (Kalisch et al., 2016), especially in tropical regions and cyclones (Chane Ming et al., 2014). A definitive observational
686 confirmation of this from the satellite itself, however, remains a challenge due to the aforementioned limitations.

687 Looking forward, a critical application for such observations is the constraint of gravity wave momentum fluxes, which are
688 essential for global circulation models. However, deriving momentum flux estimates directly from single-
689 component wind measurements like those from Aeolus presents significant theoretical and observational challenges. The
690 vertical flux of horizontal momentum (e.g., $\langle u'w' \rangle$) fundamentally requires simultaneous knowledge of both horizontal (u') and
691 vertical (w') wind perturbations. Aeolus provides only a projection of the horizontal wind and, crucially, contains no direct
692 information on the vertical wind; in fact, w' is assumed to be negligible in the standard data processing (Krisch et al., 2022).
693 This represents the primary missing piece of information for a direct flux calculation.

694 A potential pathway to overcome this limitation involves creating synergistic datasets, for instance by combining Aeolus wind
695 data with simultaneous, collocated temperature measurements from instruments like GNSS-RO. In principle, gravity wave
696 polarization relations could then be used to infer the missing wind components. However, this approach is not a simple remedy
697 and relies on strong, often unverifiable, assumptions about unmeasured wave parameters, including the horizontal wavelength,
698 intrinsic frequency, and the stationarity of the wave field between measurements (Alexander et al., 2008a; Chen et al., 2022).

699 Therefore, while Aeolus does not directly measure momentum flux, its unprecedented global measurements of kinetic energy
700 provide an additional observational constraint. Such observations are a critical prerequisite for developing and testing the more
701 complex, multi-instrument techniques that will be required to eventually constrain the global gravity wave momentum budget.

702 **6. Conclusion**

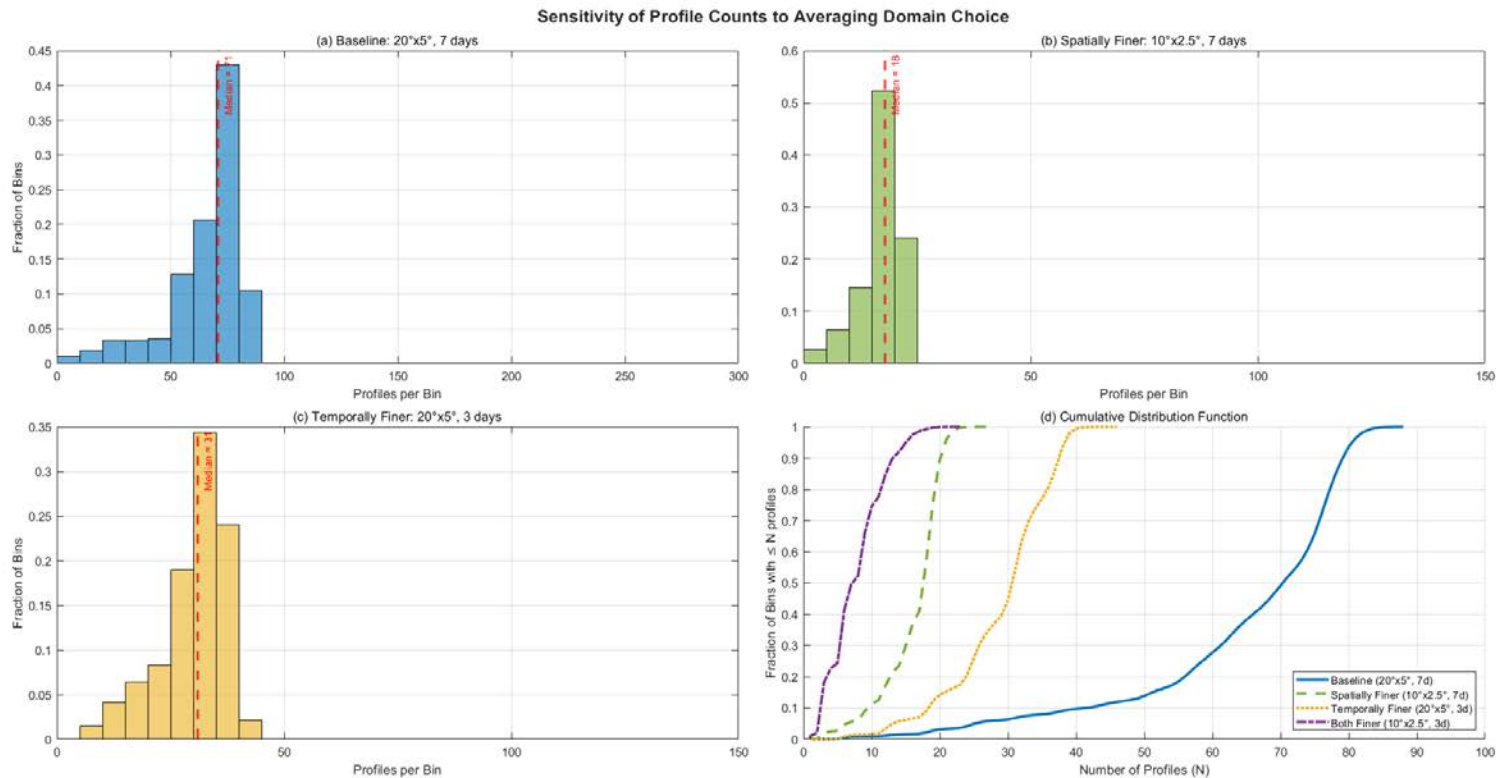
703 In this study, we examined the capacity of the Aeolus ALADIN instrument to capture and resolve GWs in tropical UTLS.
704 While this task might appear challenging at first, because of the data alteration issues Aeolus faced during its lifecycle, the
705 study proposed a noise correction process, which used ERA5 reanalysis as a reference to estimate and correct for Aeolus's
706 instrument-induced variance. This correction improved the retrieving of kinetic energy, and our comparison with collocated
707 radiosonde data further validated that approach. A key focus of our analysis was the ratio between kinetic and potential energies
708 (E_k/E_p), providing insights into the linear or non-linear nature of these waves.

709 The principal findings can be summarized as follows:

- Aeolus observations capture significant kinetic energy enhancements over tropical convection hotspots, particularly over the Indian Ocean, where ERA5 shows substantial underrepresentation due to sparse wind observations.
- Direct wind data from Aeolus could significantly enhance tropical UTLS reanalysis products, particularly in convection-driven GW regimes, reducing biases in Ek representation.
- In many regions with strong convective forcing, Aeolus data suggest a larger kinetic energy component, pointing to wave breaking, saturation, and other non-linear processes that depart from purely linear wave dynamics.
- While linear GW theory often prescribes an Ek/Ep ratio between ~ 1.6 and ~ 2.0 , our results show that this ratio can vary significantly, depending on location and season. This highlights the need for direct kinetic-energy measurements rather than relying solely on temperature-derived potential energy as a proxy.
- Aeolus also helps fill this gap. However, given its HLOS projection, Aeolus underestimates the total Ek if meridional components are significant, reinforcing that multi-instrument approaches are mandatory for accurately characterizing GW fields.

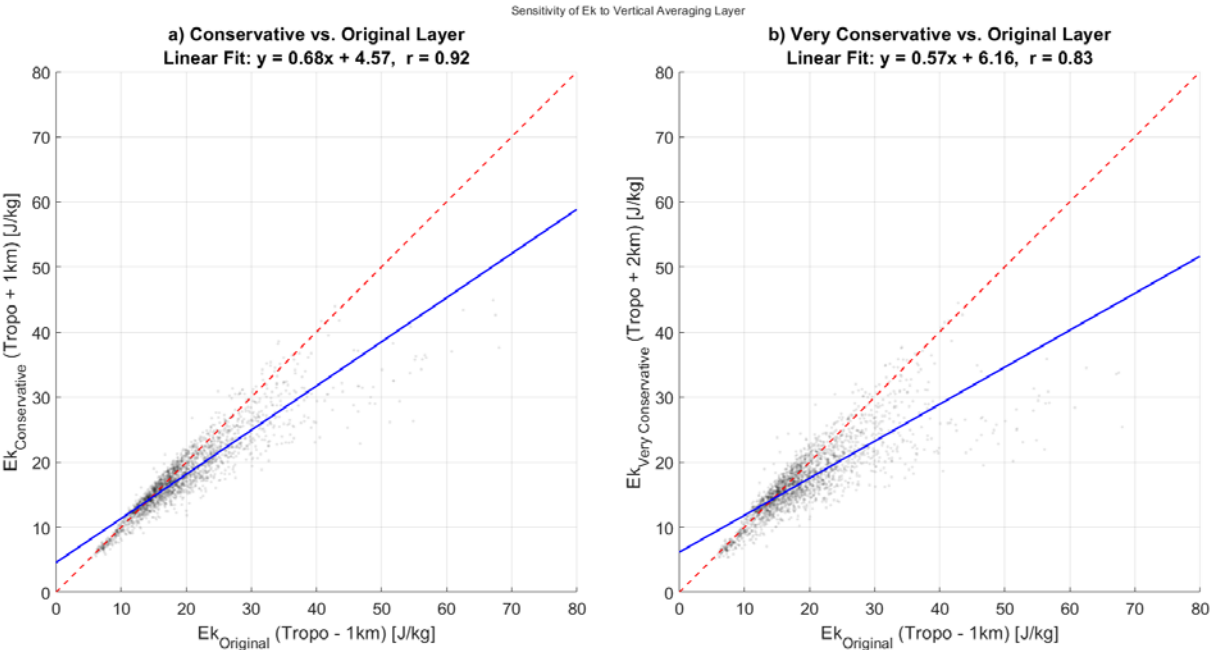
Thus, this study has demonstrated the value of Aeolus Rayleigh wind profiling for observing GWs in the tropical UTLS, despite the high and time-variable random error associated with its measurements. Our findings confirm that the annual and zonal variation of GW activity in the tropical tropopause layer and lower stratosphere is modulated by deep convection, as demonstrated by Dzambo et al. (2019) and Evan et al. (2020). Furthermore, Aeolus data expose a significant need for improving the reanalysis regarding the convective GW Ek. The lack of GW-derived Ek in ERA5 is most pronounced in the Indian Ocean region, where conventional radiosonde wind measurements are relatively sparse. It is highly likely that the missing Ek in ERA5 is due to the misrepresentation of convective processes. The results also indicate that standard assumptions about the Ek/Ep ratio do not always hold, particularly under convective or otherwise non-linear conditions. Aeolus' range-bin design and horizontal integration restrict its ability to determine wavelengths with accuracy, which poses a significant challenge for fully capturing the characteristics of GW. This limitation highlights the need for complementary datasets, which could be addressed in newer iterations of the instrument. While this study delivers some insights into UTLS GW activity and the benefits of global wind observation, future research should continue investigating the factors contributing to the discrepancies observed between Aeolus and ERA5 data. The kinetic energy constraints provided here represent a novel step, and future missions like Aeolus-2 will be essential for developing the synergistic techniques required to ultimately quantify the global momentum transport by these waves.

743



744 **Figure A1. Sensitivity of Aeolus profile counts per analysis bin to the choice of spatio-temporal averaging domain.**

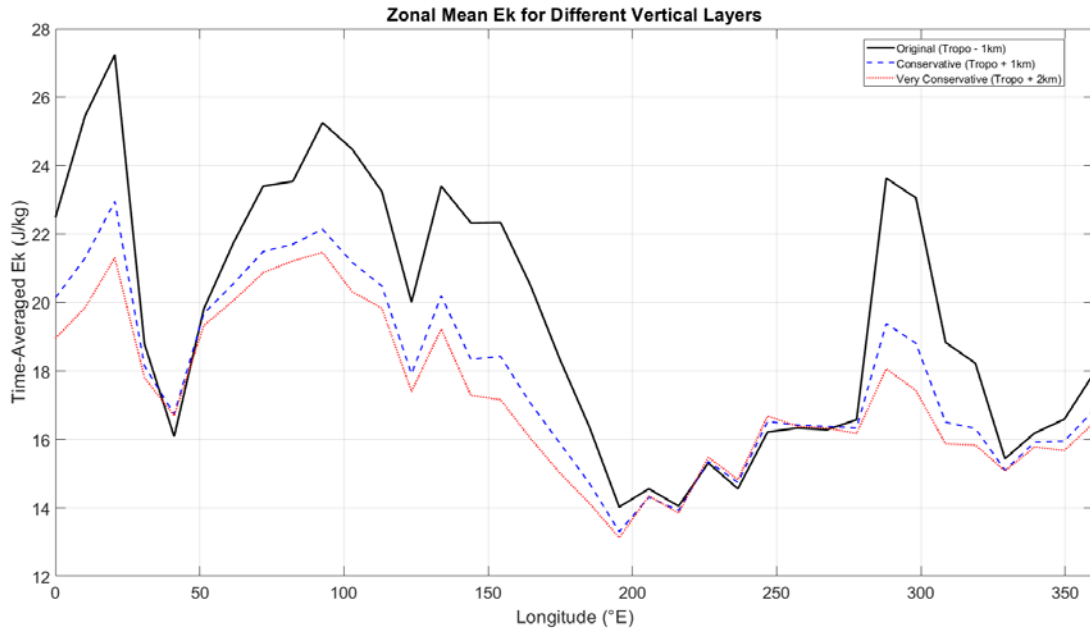
745 Figure A1 presents the histograms of the number of Aeolus profiles available per spatio-temporal bin for four different
746 averaging domain configurations: (a) the baseline 20°x5°, 7-day domain used in this study, (b) a spatially finer 10°x2.5°, 7-
747 day domain, and (c) a temporally finer 20°x5°, 3-day domain. Red dashed lines indicate the median profile count for each
748 configuration. Panel (d) shows the cumulative distribution function for all four configurations, including a domain that is finer
749 in both space and time. The results demonstrate the trade-off between domain size and sampling density. While finer domains
750 offer higher resolution, they significantly reduce the number of profiles available for robustly calculating the background state,
751 with a majority of bins containing fewer than 20 profiles. The baseline configuration (blue line) was chosen as it ensures a
752 sufficient number of profiles per bin (median of 71) to minimize noise and provide a stable background estimate, as discussed
753 in the main text.



755

756 **Figure B1. Scatter plot comparison of Aeolus-derived kinetic energy (Ek) calculated using different vertical averaging layers.**

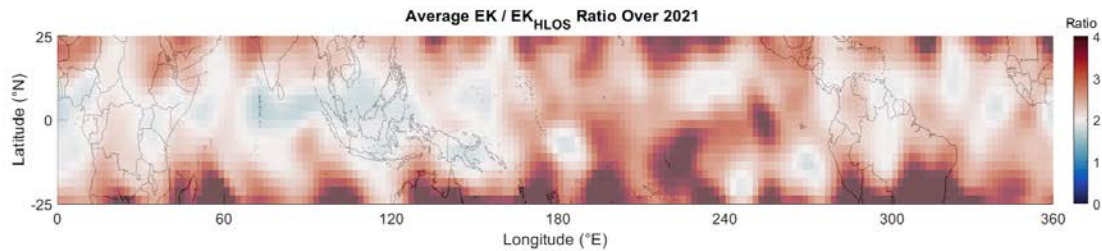
757 Figure B1 tests the sensitivity of the derived kinetic energy to the vertical averaging layer, addressing the potential
758 contamination from upper-tropospheric non-wave outflow. The x-axis on both panels represents the Ek calculated using our
759 original layer (starting 1 km below the tropopause). (a) The y-axis shows Ek calculated using a "Conservative" layer starting
760 1 km above the tropopause. (b) The y-axis shows Ek from a "Very Conservative" layer starting 2 km above the tropopause.
761 The strong linear correlation ($r = 0.92$ and $r = 0.83$, respectively) and the fact that the data points cluster near the 1:1 line (red
762 dashed line) demonstrate that the majority of the energy signal is retained when the upper troposphere is excluded. The linear
763 fit (blue solid line) shows a predictable reduction in magnitude but confirms that the underlying spatial patterns of high and
764 low energy are highly consistent across all layers.



765

766 **Figure B2. Zonal mean structure of time-averaged Aeolus kinetic energy (E_k) for different vertical averaging layers.**

767 Figure B2 shows the time-averaged, zonal mean E_k as a function of longitude for the three different vertical averaging layers
 768 defined in Figure B1. The black line represents our original layer, the blue dashed line is the "Conservative" layer (+1 km),
 769 and the red dotted line is the "Very Conservative" layer (+2 km). The plot confirms that while the absolute magnitude of E_k
 770 decreases as the layer is moved higher into the stratosphere (as expected due to wave dissipation), the geographical structure
 771 of the energy hotspots is remarkably stable. The primary peaks of high energy (e.g., over the African/Indian Ocean sector from
 772 ~60-120°E and the Americas/Atlantic sector from ~280-320°E) persist across all three calculations. This provides strong
 773 evidence that the observed energy hotspots are robust features originating from vertically propagating gravity waves that have
 774 reached the lower stratosphere, and not artifacts of shallow upper-tropospheric outflow.



776

777 **Figure C1. Temporal and spatial variability of the ratio between E_{kHLOS} and E_k in the ERA5 model over the tropical region (30°S–**
778 **30°N) for the year 2021**

779 Figure C1 displays the ratio between E_{kHLOS} and E_k in the ERA5 model, over the tropical region for the UTLS. The ratio
780 values range from 1.5 to 3, between January 2021 and December 2021 included, depicting variations in how well the HLOS
781 measurements capture the total kinetic energy in this region.

782

783 This map illustrates the anisotropy of the gravity wave field as represented in the ERA5 model for the year 2021, chosen for
784 its significant wave activity which highlights the regional patterns of energy partitioning. The ratio quantifies the contribution
785 of the unobserved meridional wind component to the total kinetic energy. A ratio of 1 (blue colors) indicates purely zonal
786 wave energy, while higher values (red colors) signify an increasing contribution from meridional motions.

787 A distinct pattern emerges along the equatorial belt (10°S - 10°N). Over a vast longitudinal sector stretching from Africa,
788 across the Indian Ocean, and over the Maritime Continent (approximately 0°E to 160°E), the ratio remains low, generally
789 below 2.0. This indicates that wave energy is predominantly zonal in these regions, meaning a quasi-zonal HLOS measurement
790 like that from Aeolus is expected to capture a large fraction of the total kinetic energy. This aligns with the influence of
791 persistent large-scale zonal flows in this sector, such as easterly trade winds.

792 In contrast, the sector from approximately 160°E to 280°E, encompassing the central and eastern Pacific, shows significantly
793 higher ratios, often exceeding 2.5. This points to a wave field with a much stronger meridional component, where an HLOS
794 measurement would systematically underestimate the total kinetic energy. A third regime is observed over the Atlantic and
795 South America (from 280°E to 360°E), where the ratio is more variable but consistently elevated, indicating a mixed wave
796 field with a significant, though not always dominant, meridional energy component.

797

798 Appendix D : Aeolus Instrumental Noise Correction Methodology

799 This appendix provides a detailed, step-by-step description of the spatio-temporally adaptive algorithm used to correct for
800 time-varying instrumental noise in the Aeolus-derived gravity wave (GW) kinetic energy (Ek) data. The objective is to produce
801 a corrected dataset, $\overline{Ek}_{Aeolus\ HLOS^*}$, that removes instrumental bias while preserving the unique, high-fidelity geophysical
802 signals of GWs that Aeolus observes. The correction methodology involves two primary stages: first, a detailed correction
803 applied to time-longitude Hovmöller diagrams, and second, a refinement of this correction for application to broader
804 geographical maps.

805 Part 1: Correction of Time-Longitude Hovmöller Diagrams

806 As introduced in the main text, our best estimate of the noise energy is denoted as $\overline{Ek}_{I.N.}$. Within this appendix, where specific
807 spatio-temporal and component estimates of this noise are derived, we will use the "hat" notation to explicitly denote
808 these estimated noise values, such as $\widehat{Ek}_{I.N.}$ for the time-longitude noise estimate. The initial estimation of the noise term is
809 performed on the 2D Hovmöller data matrices of Aeolus $E_{k,Aeolus\ HLOS}(t,\phi)$ and ERA5 $E_{k,ERA5\ HLOS}(t,\phi)$ where t represents
810 time steps (3-week averages) and ϕ represents longitude bins. This correction is conducted within the 10°S–10°N latitude band
811 where deep tropical convection is most prominent. The algorithm adaptively estimates the instrumental noise component
812 $\widehat{Ek}_{I.N.}(t,\phi)$ which is then subtracted from the observed Aeolus energy as shown conceptually in Eq. (6) of the main text.

813 Step 1: Defining an Activity-Based Blending Weight

814 The core of the adaptive algorithm is a blending weight, $W(t,\phi)$, that determines the balance between an additive and a
815 multiplicative correction approach. This weight is derived from the ERA5 Ek, which serves as a proxy for the true level of
816 atmospheric GW activity. First, a background reference energy $E_{k,bg}$ (bg standing for background) defined as the median
817 ERA5 Ek within a pre-defined quiescent reference sector (200°E–250°E over the Pacific Ocean), chosen for its typically low
818 convective activity. An activity index, $x(t,L)$, is then computed for each point in the Hovmöller diagram:

$$819 \quad x(t,\phi) = \frac{E_{k,ERA5\ HLOS}(t,\phi) - E_{k,bg}}{E_{k,bg}} \quad (7)$$

820 From this index, a logistic function is used to create a preliminary activity map, $A(t,L)$, ranging from 0 (quiescent) to 1 (active):

$$821 \quad A(t,\phi) = \frac{1}{1 + e^{-x(t,\phi)}} \quad (8)$$

822

823 To ensure spatial coherence and prevent abrupt transitions, this activity map is smoothed using a 2D Gaussian filter with a
824 standard deviation of 2 pixels in the time dimension and 10 pixels in the longitude dimension. The final background-confidence
825 weight, $W(t,L)$, is then defined as:

$$W(t, \phi) = 1 - A_{smoothed}(t, \phi) \quad (9)$$

Thus, W is close to 1 in quiescent regions and approaches 0 in highly active regions.

Step 2: Deriving the Additive and Multiplicative Correction Components

Two candidate correction components are calculated. The additive component is designed to correct for the baseline instrumental noise offset. A time-dependent offset vector, $\delta_{noise}(t)$, is computed by taking the median difference between Aeolus and ERA5 Ek within the same quiet reference sector used in Step 1:

$$\delta_{noise}(t) = \text{median}_{\phi \in \text{ref. sector}} \left(E_{k, AeolusHLOS}(t, \phi) - E_{k, ERA5HLOS}(t, \phi) \right) \quad (10)$$

This offset vector, which captures the time-varying instrumental bias trend, is then smoothed temporally with a 7-point (21-week) moving median to reduce noise. The additively corrected energy field, $E_{k, bias-adj}$ is then defined as:

$$E_{k, bias-adj}(t, \phi) = E_{k, AeolusHLOS}(t, \phi) - \beta \cdot \delta_{noise}(t) \quad (11)$$

where β is an offset relaxation factor (set to 1.15) to empirically fine-tune the subtraction.

The multiplicative component is designed for active regions where noise effects might scale with the signal. A single, mission-period Activity Ratio Scalar, $\mathcal{R}_{\mathcal{A}/\mathcal{E}}$, is computed as the median ratio of Aeolus to ERA5 Ek over all grid points where the activity weight $A(t, \phi)$ from Eq. (8) is greater than 0.5:

$$\mathcal{R}_{\mathcal{A}/\mathcal{E}} = \text{median}_{(t, \phi) \in \text{active}} \left(\frac{E_{k, AeolusHLOS}(t, \phi)}{E_{k, ERA5HLOS}(t, \phi)} \right) \quad (12)$$

The Ratio-Scaled Energy, $E_{k, ratio-scaled}$, is then:

$$E_{k, ratio-scaled}(t, \phi) = \frac{E_{k, AeolusHLOS}(t, \phi)}{\mathcal{R}_{\mathcal{A}/\mathcal{E}}} \quad (13)$$

Step 3: Blending and Final Hovmöller Correction

The final corrected Hovmöller field, $Ek^*(t, \phi)$, is a weighted average of the two candidates, using the weight map W from Eq. (9):

$$E_k^*(t, \phi) = W(t, \phi) \cdot E_{k, \text{bias-adj}}(t, \phi) + [1 - W(t, \phi)] \cdot E_{k, \text{ratio-scaled}}(t, \phi) \quad (14)$$

Any resulting negative values are clipped to zero. The estimated noise component, $\overline{Ek}_{I,N}$, as referenced in Eq. (6) of the main text, is therefore equivalent to the total amount subtracted:

$$\overline{Ek}_{I,N} = \overline{Ek}_{\text{Aeolus HLOS}} - \overline{Ek}_{\text{Aeolus HLOS}^*} \quad (15)$$

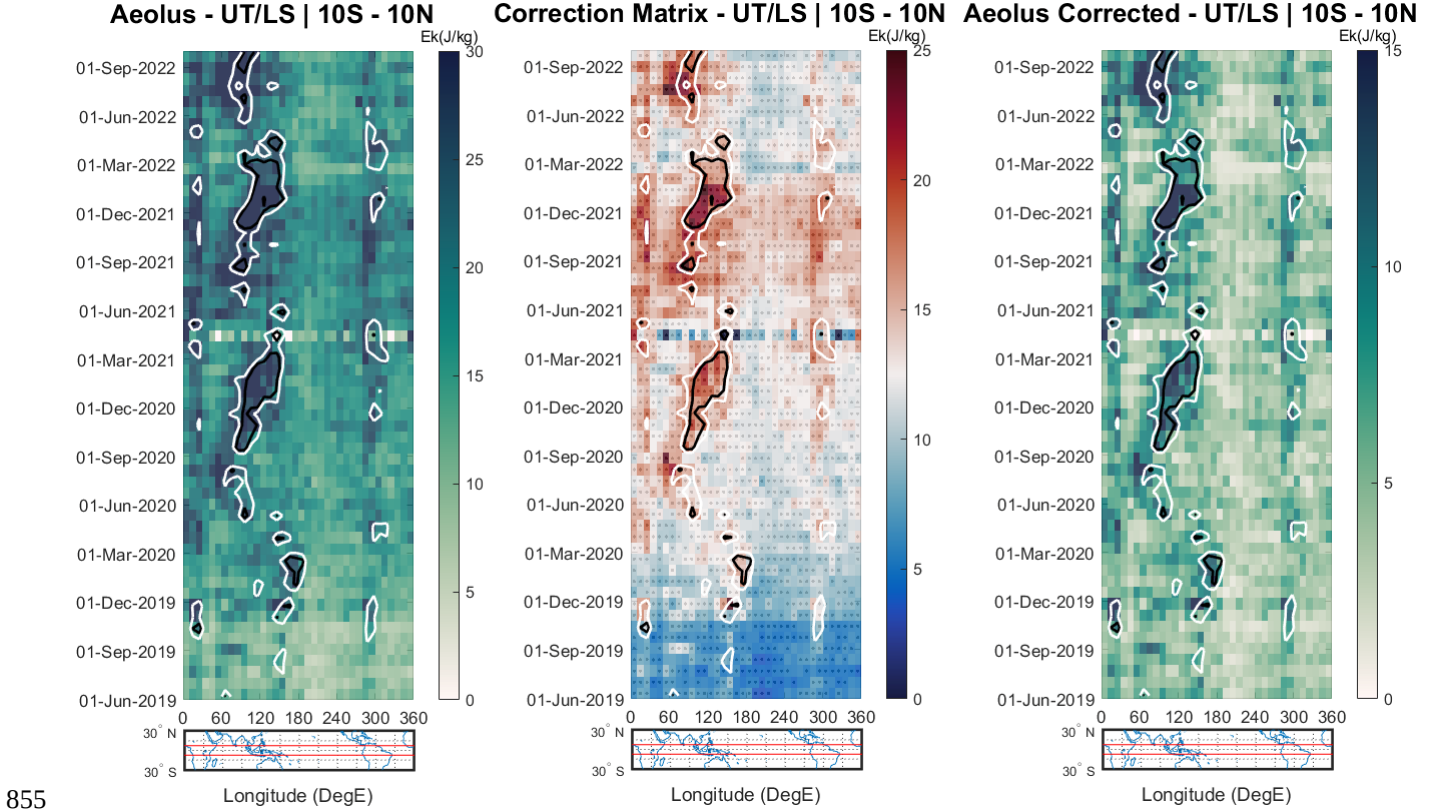


Figure D1. Breakdown of the Hovmöller diagram correction. (a) Raw Aeolus Ek before correction. (b) The estimated noise correction matrix, $Ek_{I,N}$, that is subtracted. Note the increasing trend over time, reflecting instrument degradation, and the spatial structure modulated by the blending of additive and multiplicative components. Stippling indicates regions where the corrected Aeolus Ek is statistically different from ERA5. Ek (F-test, $p < 0.05$). (c) Final corrected Aeolus Ek ($Ek_{\text{Hovmöller}^*}$). Background energy levels are significantly reduced while convective hotspots are preserved. White contours indicate low Outgoing Longwave radiation ($OLR < 220 \text{ W m}^{-2}$).

Figure D1 illustrates the core steps of the spatio-temporally adaptive noise correction. (a) The raw, uncorrected Aeolus Ek Hovmöller diagram, showing both geophysical signals and a clear increasing trend in background noise over the mission lifetime. (b) The estimated noise component derived using the adaptive algorithm. Note the temporal increase reflecting instrument degradation, as well as the spatial structure modulated by the blending of additive and multiplicative corrections. (c) The final corrected Aeolus Ek. The background energy levels have been significantly reduced, removing the artificial trend,

while the physically meaningful convective hotspots are preserved and sharpened. White contours indicate regions of deep convection ($OLR < 220 \text{ W m}^{-2}$).

Part 2: Refinement of Noise Correction for Seasonal Geographical Maps

For the seasonally averaged geographical maps, which cover a broader latitudinal range (0°S – 30°N , denoted by latitude θ) than the initial Hovmöller analysis, the noise estimation is refined. For each season (e.g., DJF, MAM, JJA, SON), the estimated noise field from the Hovmöller analysis $\overline{Ek_{I.N.}}$, is averaged over the relevant time steps to yield a Seasonal Mean Longitudinal Noise Profile $\overline{Ek_{I.N.}}$. This yields a seasonally representative longitudinal noise profile. This profile forms the basis for correcting the uncorrected, seasonally-averaged raw Aeolus Ek map, $E_{k,Aeolus,map}(\theta,\phi)$. The refinement proceeds in two steps: latitude-weighting and iterative adjustment.

Step 1: Latitude-Weighting of the Longitudinal Correction

The impact of instrumental noise may be proportionally larger in regions of higher true GW energy. To approximate this, the base longitudinal correction profile is weighted by the latitudinal structure of the raw Aeolus Ek map itself. First, the mean raw Aeolus Ek is calculated for each latitude band, $E_{k,Aeolus,map}$. This is normalized to create a latitude weight profile, W_θ :

$$W_\theta(\theta) = \frac{\overline{E_{k,Aeolus,map}}(\theta)}{\max(\overline{E_{k,Aeolus,map}}(\theta))} \quad (16)$$

This weight is clipped between 0 and 1. This latitudinal weight is then broadcast across the longitudinal correction profile to create an initial 2D correction field for the map, $\Delta E_{map,initial}$:

$$\Delta E_{map,initial}(\theta, \phi) = \overline{E_{k,I.N.}}(\phi) \cdot W_{lat}(\theta) \quad (17)$$

Step 2: Iterative Adjustment to Prevent Negative Kinetic Energy

A simple subtraction of $\Delta E_{map,initial}$ from $E_{k,Aeolus,map}$ could result in non-physical negative Ek values where the estimated correction is large relative to the observed Ek. To prevent this, a Columnar Clipping Factor, γ_{clip} , is calculated for each longitude column of the map. This factor is defined as the minimum ratio of the observed energy to the initial correction estimate across all latitudes in that column, ensuring the correction never exceeds the available energy:

$$\gamma_{clip}(\phi) = \max\left(0, \min\left(1, \min_{\theta} \left[\frac{E_{k,Aeolus,map}(\theta,\phi)}{\Delta E_{map,initial}(\theta,\phi)}\right]\right)\right) \quad (18)$$

The final 2D correction field applied to the map is then:

$$\Delta E_{map,final}(\theta, \phi) = \Delta E_{map,initial}(\theta, \phi) \cdot \gamma_{clip}(\phi) \quad (19)$$

Step 3: Final Corrected Map and Smoothing

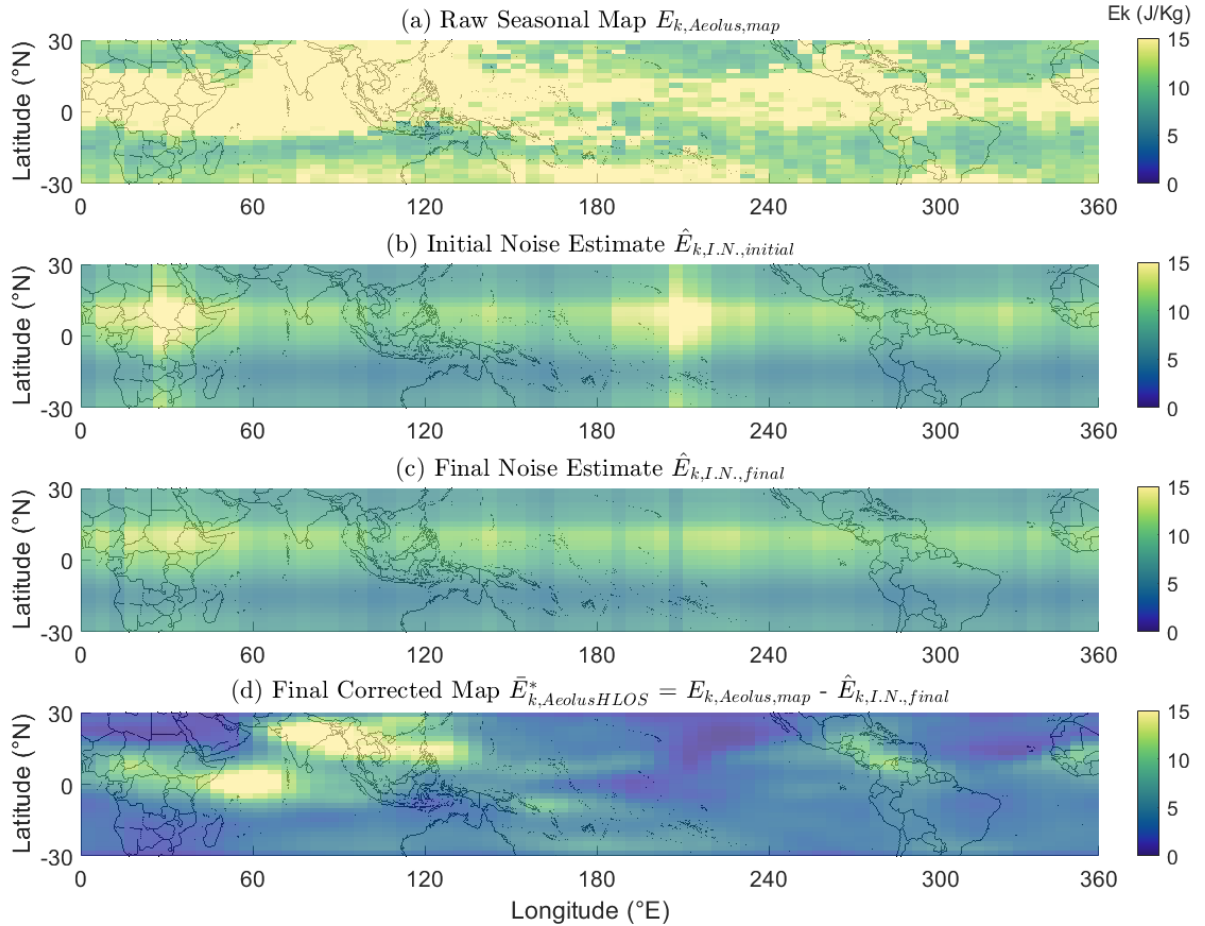
The final corrected seasonal map, denoted as $\overline{E_{k,AeolusHLOS}^*}$, is:

$$\overline{E_{k,AeolusHLOS}^*}(\theta, \phi) = E_{k,Aeolus,map}(\theta, \phi) - \Delta E_{map,final}(\theta, \phi) \quad (20)$$

895

896 Finally, a light 2D spatial smoothing (3x3 moving median followed by a 3x3 moving mean) is applied to the corrected map to

897 reduce pixel-scale noise introduced by the gridding and correction process.

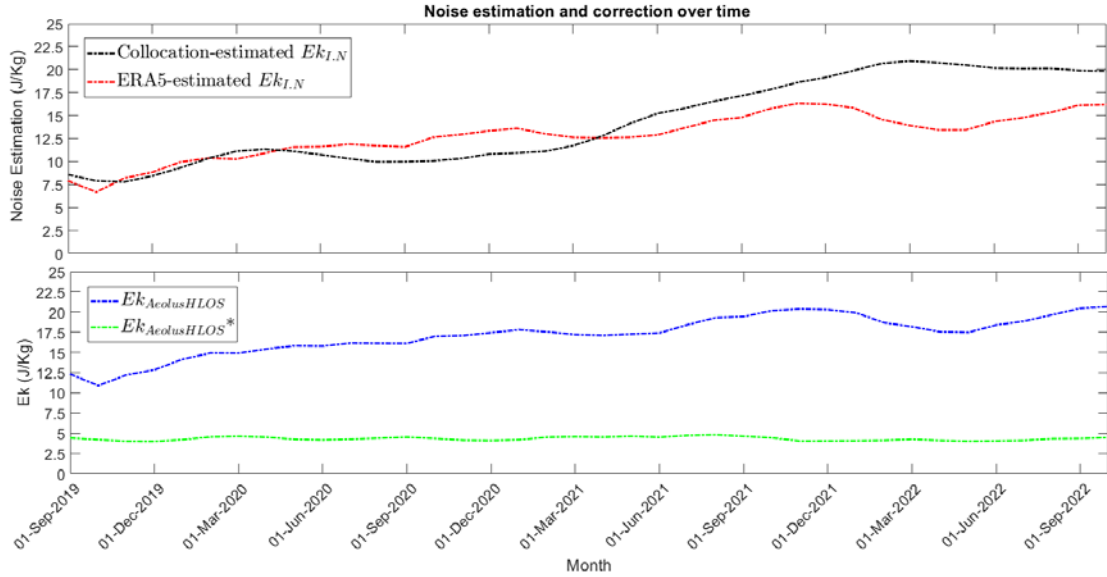


898

899 **Figure D2. Illustration of the noise correction refinement process for a seasonal geographical map (JJA 2021 shown as an**
900 **example).**

901 Figure D2 details the refinement steps used to adapt the noise correction for application to 2D seasonal maps. (a) The raw,
 902 uncorrected seasonally averaged Aeolus Ek map. (b) The initial noise estimate, created by applying a latitudinal weighting to
 903 the seasonal-mean longitudinal noise profile derived from the Hovmöller analysis. (c) The final noise estimate, after an iterative
 904 adjustment step that prevents non-physical negative energy values. (d) The final corrected seasonal map, representing the
 905 geophysical gravity wave kinetic energy field after the removal of instrumental artifacts. This process ensures that the
 906 correction is physically consistent across the entire map domain.

907 Diagnostic Validation of Noise Correction



908

909 **Figure D3. Difference between the Radiosonde-derived (black) and ERA5-derived (red) estimated noise correction, resulting in the**
 910 **difference between the uncorrected Aeolus HLOS GW Ek (blue) and the ERA5-corrected Aeolus HLOS GW Ek (green).**

911 Figure D3 provides an independent validation of the estimated noise trend by comparing the zonally-averaged $\delta_{noise}(t)$, term
 912 from Eq. (10) with noise estimates derived from collocated radiosonde measurements at La Réunion, as in Ratynski et al.
 913 (2023). It is intended to demonstrate that such a method of instrumental noise estimation is qualitatively consistent with the
 914 classical approach based on collocated reference measurements applied in (Ratynski et al., 2023).

915

916 The Météo-France upper-air soundings in La Réunion (Aéroport Gillot) was used for the conduct of this analysis. For each
 917 collocated radiosonde profile with an Aeolus overpass (within 200 km and +/- 6 hours), we downsampled the radiosonde
 918 profile resolution to be equivalent to ALADIN vertical bins. A point-wise difference is then calculated, and the standard

919 deviation of these differences is what we refer to as random error. In principle, if Aeolus would not experience any degradation
920 through its systems, this standard deviation would remain stable over the years and periods. However, since we observe an
921 increase, as reported by Ratynski et al. (2023, their Fig.6), a link can be made between the instrument degradation and this
922 increase, wrongly attributing signal-to-noise. Squaring this noise estimation provides a metric homogenous to the observed
923 Ek, representing the repercussions of noise on Ek estimation:

$$924 \quad Ek_{I.N} = \frac{1}{2} (\sigma_{Aeolus-Radiosondes})^2 \quad (21)$$

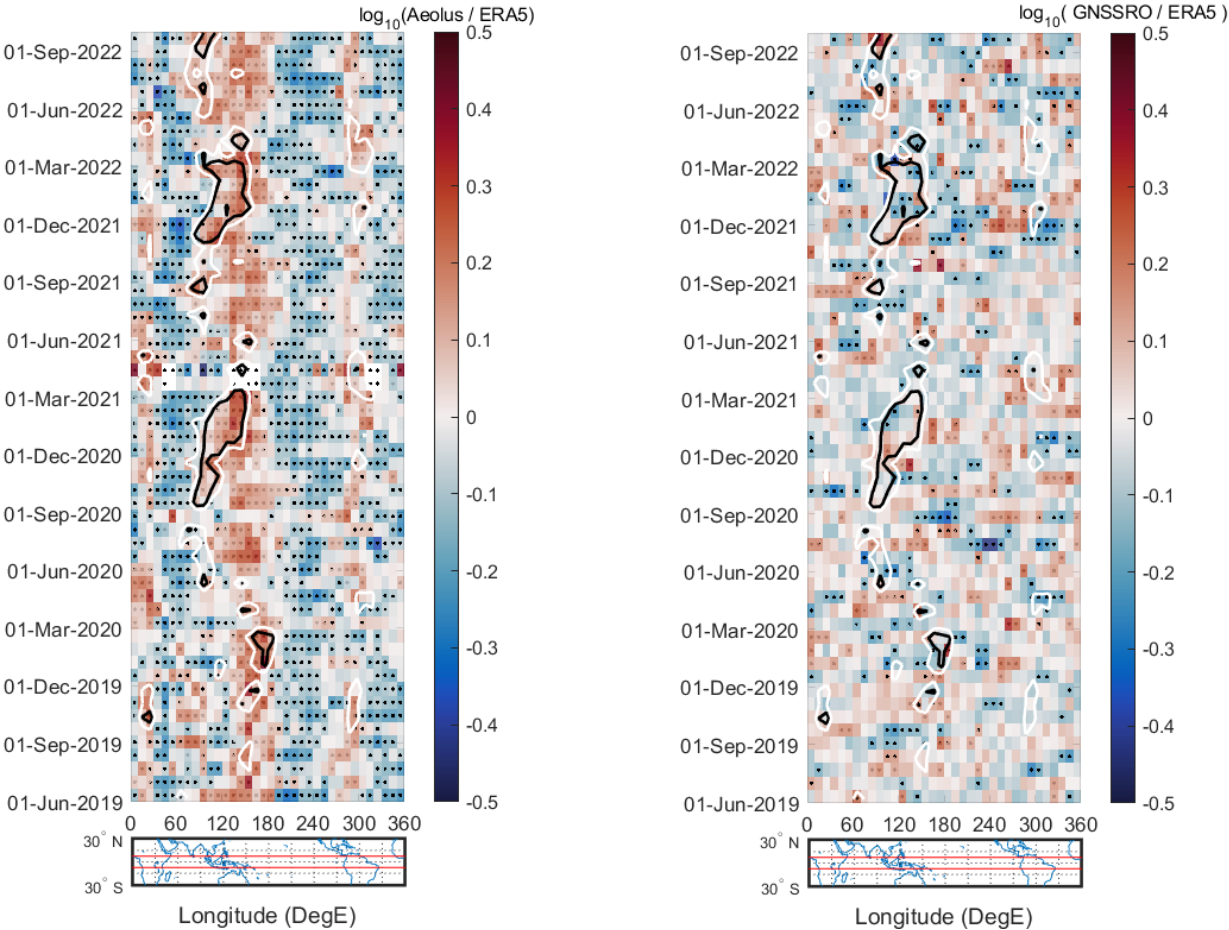
925

926 While both methods provide similar trends, the model approach remains the safest estimation when considering the potential
927 biases.

928

929

Aeolus / ERA5 GW Ek | UT/LS | 10S - 10N GNSS-RO / ERA5 Energy Ratio | UT/LS | 10S - 10N



931
932 **Figure E1. Ratio-based comparison of observed and reanalyzed gravity wave energy for the equatorial band (10°S–10°N), between**
933 **observational data and ERA5. The UTLS altitudes are defined between one kilometer below the tropopause and 22 km. The contour**
934 **plot represents the Outgoing Longrange Radiation (OLR) for 210 and 220 W/m² (black and white, respectively). Each bin**
935 **corresponds to an average of over 3 weeks and 10 degrees. The OLR measurements were obtained from the Australian Bureau of**
936 **Meteorology. The tropopause is determined from the ERA5 reanalysis. The panels show the base-10 logarithm of the ratio between**
937 **observational data and ERA5.**

938 Figure E1 provides an alternative, ratio-based view of the comparisons presented in the main text, complementing the
939 difference-based analysis. The color scale represents the base-10 logarithm of the ratio, where positive values (red) indicate
940 that the observational dataset has higher energy than ERA5, and negative values (blue) indicate the opposite.

941 The first panel reinforces the findings from the main text, showing a systematic and significant overestimation of kinetic energy
942 by Aeolus relative to ERA5 (predominantly red colors) within the convective regions identified by low OLR. The ratio
943 frequently exceeds 2 ($\log_{10}$ ratio > 0.3), particularly over the Indian Ocean and Maritime Continent, confirming that ERA5
944 substantially underestimates convection-driven kinetic energy.

945 In contrast to the kinetic energy, the potential energy ratio is much closer to unity ($\log_{10}$ ratio ≈ 0). The colors are predominantly
946 neutral or light shades of blue/red, indicating that ERA5 and GNSS-RO have very similar magnitudes of potential energy. This
947 is expected, as ERA5 assimilates GNSS-RO temperature data.

948 Taken together, these two panels provide evidence for the central conclusion of this study: the discrepancy between
949 observations and reanalysis is specific to the unassimilated kinetic energy component. While ERA5 successfully reproduces
950 the potential energy field it is constrained by, it fails to generate the corresponding kinetic energy associated with convection,
951 a gap that direct wind observations from Aeolus can fill. Black stippling indicates regions where the ratios are statistically
952 significant (F-test, $p < 0.05$).

953 **Data Availability:** Aeolus data are publicly available through the Aeolus online dissemination system (<https://aeolus->
954 ds.eo.esa.int/oads/access/). The dataset used for the realization of this study can be found at
955 <https://doi.org/10.5281/zenodo.8113261>
956

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961
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967 **References**

968 Achatz, U., Alexander, M. J., Becker, E., Chun, H.-Y., Dörnbrack, A., Holt, L., Plougonven, R., Polichtchouk, I., Sato, K.,
969 Sheshadri, A., Stephan, C. C., Niekerk, A. van, and Wright, C. J.: Atmospheric Gravity Waves: Processes and
970 Parameterization, <https://doi.org/10.1175/JAS-D-23-0210.1>, 2024.

- 971 Alexander, M. J. and Ortland, D. A.: Equatorial waves in High Resolution Dynamics Limb Sounder (HIRDLS) data, J.
972 Geophys. Res. Atmospheres, 115, <https://doi.org/10.1029/2010JD014782>, 2010.
- 973 Alexander, M. J., Gille, J., Cavanaugh, C., Coffey, M., Craig, C., Eden, T., Francis, G., Halvorson, C., Hannigan, J., Khosravi,
974 R., Kinnison, D., Lee, H., Massie, S., Nardi, B., Barnett, J., Hepplewhite, C., Lambert, A., and Dean, V.: Global estimates of
975 gravity wave momentum flux from High Resolution Dynamics Limb Sounder observations, J. Geophys. Res. Atmospheres,
976 113, <https://doi.org/10.1029/2007JD008807>, 2008a.
- 977 Alexander, M. J., Liu, C. C., Bacmeister, J., Bramberger, M., Hertzog, A., and Richter, J. H.: Observational Validation of
978 Parameterized Gravity Waves From Tropical Convection in the Whole Atmosphere Community Climate Model, J. Geophys.
979 Res. Atmospheres, 126, e2020JD033954, <https://doi.org/10.1029/2020JD033954>, 2021.
- 980 Alexander, S., Tsuda, T., and Kawatani, Y.: COSMIC GPS Observations of Northern Hemisphere winter stratospheric gravity
981 waves and comparisons with an atmospheric general circulation model, Geophys. Res. Lett. - GEOPHYS RES LETT, 35,
982 <https://doi.org/10.1029/2008GL033174>, 2008b.
- 983 Alexander, S. P., Tsuda, T., Kawatani, Y., and Takahashi, M.: Global distribution of atmospheric waves in the equatorial upper
984 troposphere and lower stratosphere: COSMIC observations of wave mean flow interactions, J. Geophys. Res. Atmospheres,
985 113, <https://doi.org/10.1029/2008JD010039>, 2008c.
- 986 Baker, W. E., Atlas, R., Cardinali, C., Clement, A., Emmitt, G. D., Gentry, B. M., Hardesty, R. M., Källén, E., Kavaya, M. J.,
987 Langland, R., Ma, Z., Masutani, M., McCarty, W., Pierce, R. B., Pu, Z., Riishojgaard, L. P., Ryan, J., Tucker, S., Weissmann,
988 M., and Yoe, J. G.: Lidar-Measured Wind Profiles: The Missing Link in the Global Observing System, Bull. Am. Meteorol.
989 Soc., 95, 543–564, <https://doi.org/10.1175/BAMS-D-12-00164.1>, 2014.
- 990 Banyard, T. P., Wright, C. J., Hindley, N. P., Halloran, G., Krisch, I., Kaifler, B., and Hoffmann, L.: Atmospheric Gravity
991 Waves in Aeolus Wind Lidar Observations, Geophys. Res. Lett., 48, e2021GL092756,
992 <https://doi.org/10.1029/2021GL092756>, 2021.
- 993 Baumgarten, G., Fiedler, J., Hildebrand, J., and Lübken, F.-J.: Inertia gravity wave in the stratosphere and mesosphere observed
994 by Doppler wind and temperature lidar, Geophys. Res. Lett., 42, 10,929–10,936, <https://doi.org/10.1002/2015GL066991>, 2015.
- 995 Bedka, K. M., Velden, C. S., Petersen, R. A., Feltz, W. F., and Mecikalski, J. R.: Comparisons of Satellite-Derived
996 Atmospheric Motion Vectors, Rawinsondes, and NOAA Wind Profiler Observations,
997 <https://doi.org/10.1175/2009JAMC1867.1>, 2009.
- 998 Beres, J. H., Garcia, R. R., Boville, B. A., and Sassi, F.: Implementation of a gravity wave source spectrum parameterization
999 dependent on the properties of convection in the Whole Atmosphere Community Climate Model (WACCM), J. Geophys. Res.
1000 Atmospheres, 110, <https://doi.org/10.1029/2004JD005504>, 2005.
- 1001 Bramberger, M., Alexander, M. J., Davis, S., Podglajen, A., Hertzog, A., Kalnajs, L., Deshler, T., Goetz, J. D., and Khaykin,
1002 S.: First Super-Pressure Balloon-Borne Fine-Vertical-Scale Profiles in the Upper TTL: Impacts of Atmospheric Waves on
1003 Cirrus Clouds and the QBO, Geophys. Res. Lett., 49, e2021GL097596, <https://doi.org/10.1029/2021GL097596>, 2022.
- 1004 Campos, R. M., Gramscianinov, C. B., de Camargo, R., and da Silva Dias, P. L.: Assessment and Calibration of ERA5 Severe
1005 Winds in the Atlantic Ocean Using Satellite Data, Remote Sens., 14, 4918, <https://doi.org/10.3390/rs14194918>, 2022.
- 1006 Cao, B., Haase, J. S., Murphy, M. J., Alexander, M. J., Bramberger, M., and Hertzog, A.: Equatorial waves resolved by balloon-
1007 borne Global Navigation Satellite System radio occultation in the Strateole-2 campaign, Atmospheric Chem. Phys., 22, 15379–
1008 15402, <https://doi.org/10.5194/acp-22-15379-2022>, 2022.

- 1009 Chane Ming, F., Ibrahim, C., Barthe, C., Jolivet, S., Keckhut, P., Liou, Y.-A., and Kuleshov, Y.: Observation and a numerical
1010 study of gravity waves during tropical cyclone Ivan (2008), *Atmospheric Chem. Phys.*, 14, 641–658,
1011 <https://doi.org/10.5194/acp-14-641-2014>, 2014.
- 1012 Chen, Q., Ntokas, K., Linder, B., Krasauskas, L., Ern, M., Preusse, P., Ungermann, J., Becker, E., Kaufmann, M., and Riese,
1013 M.: Satellite observations of gravity wave momentum flux in the mesosphere and lower thermosphere (MLT): feasibility and
1014 requirements, *Atmospheric Meas. Tech.*, 15, 7071–7103, <https://doi.org/10.5194/amt-15-7071-2022>, 2022.
- 1015 Corcos, M., Bramberger, M., Alexander, M. J., Hertzog, A., Liu, C., and Wright, C.: Observation of Gravity Waves Generated
1016 by Convection and the “Moving Mountain” Mechanism During Stratéole-2 Campaigns and Their Impact on the QBO, *J.*
1017 *Geophys. Res. Atmospheres*, 130, e2024JD041804, <https://doi.org/10.1029/2024JD041804>, 2025.
- 1018 Danzer, J., Pieler, M., and Kirchengast, G.: Closing the gap in the tropics: the added value of radio-occultation data for wind
1019 field monitoring across the equator, *Atmospheric Meas. Tech. Discuss.*, 1–19, <https://doi.org/10.5194/amt-2023-137>, 2023.
- 1020 Dee, D. P.: Bias and data assimilation, *Q. J. R. Meteorol. Soc.*, 131, 3323–3343, <https://doi.org/10.1256/qj.05.137>, 2005.
- 1021 Dhaka, S. K., Yamamoto, M. K., Shibagaki, Y., Hashiguchi, H., Fukao, S., and Chun, H.-Y.: Equatorial Atmosphere Radar
1022 observations of short vertical wavelength gravity waves in the upper troposphere and lower stratosphere region induced by
1023 localized convection, *Geophys. Res. Lett.*, 33, <https://doi.org/10.1029/2006GL027026>, 2006.
- 1024 Dzambo, A. M., Hitchman, M. H., and Chang, K.-W.: The Influence of Gravity Waves on Ice Saturation in the Tropical
1025 Tropopause Layer over Darwin, Australia, *Atmosphere*, 10, 778, <https://doi.org/10.3390/atmos10120778>, 2019.
- 1026 von Engeln, A., Andres, Y., Marquardt, C., and Sancho, F.: GRAS radio occultation on-board of Metop, *Adv. Space Res.*, 47,
1027 336–347, <https://doi.org/10.1016/j.asr.2010.07.028>, 2011.
- 1028 Ern, M., Preusse, P., Alexander, M. J., and Warner, C. D.: Absolute values of gravity wave momentum flux derived from
1029 satellite data, *J. Geophys. Res. Atmospheres*, 109, <https://doi.org/10.1029/2004JD004752>, 2004.
- 1030 Ern, M., Trinh, Q. T., Preusse, P., Gille, J. C., Mlynczak, M. G., Russell III, J. M., and Riese, M.: GRACILE: a comprehensive
1031 climatology of atmospheric gravity wave parameters based on satellite limb soundings, *Earth Syst. Sci. Data*, 10, 857–892,
1032 <https://doi.org/10.5194/essd-10-857-2018>, 2018.
- 1033 Evan, S., Brioude, J., Rosenlof, K., Davis, S. M., Vömel, H., Héron, D., Posny, F., Metzger, J.-M., Duflot, V., Payen, G.,
1034 Vérémes, H., Keckhut, P., and Cammas, J.-P.: Effect of deep convection on the tropical tropopause layer composition over the
1035 southwest Indian Ocean during austral summer, *Atmospheric Chem. Phys.*, 20, 10565–10586, <https://doi.org/10.5194/acp-20-10565-2020>, 2020.
- 1037 Forbes, J. M., Ern, M., and Zhang, X.: The Global Monsoon Convective System as Reflected in Upper Atmosphere Gravity
1038 Waves, *J. Geophys. Res. Space Phys.*, 127, e2022JA030572, <https://doi.org/10.1029/2022JA030572>, 2022.
- 1039 Fritts, D. C. and Alexander, M. J.: Gravity wave dynamics and effects in the middle atmosphere, *Rev. Geophys.*, 41,
1040 <https://doi.org/10.1029/2001RG000106>, 2003.
- 1041 Fritts, D. C., Wang, L., and Werne, J.: Gravity wave–fine structure interactions: A reservoir of small-scale and large-scale
1042 turbulence energy, *Geophys. Res. Lett.*, 36, <https://doi.org/10.1029/2009GL039501>, 2009.

- 1043 Fröhlich, K., Schmidt, T., Ern, M., Preusse, P., de la Torre, A., Wickert, J., and Jacobi, Ch.: The global distribution of gravity
1044 wave energy in the lower stratosphere derived from GPS data and gravity wave modelling: Attempt and challenges, *J.*
1045 *Atmospheric Sol.-Terr. Phys.*, 69, 2238–2248, <https://doi.org/10.1016/j.jastp.2007.07.005>, 2007.
- 1046 Geldenhuys, M., Kaifler, B., Preusse, P., Ungermann, J., Alexander, P., Krasauskas, L., Rhode, S., Woiwode, W., Ern, M.,
1047 Rapp, M., and Riese, M.: Observations of Gravity Wave Refraction and Its Causes and Consequences, *J. Geophys. Res.*
1048 *Atmospheres*, 128, e2022JD036830, <https://doi.org/10.1029/2022JD036830>, 2023.
- 1049 Gubenko, V. N., Pavelyev, A. G., Salimzyanov, R. R., and Andreev, V. E.: A method for determination of internal gravity
1050 wave parameters from a vertical temperature or density profile measurement in the Earth’s atmosphere, *Cosm. Res.*, 50, 21–
1051 31, <https://doi.org/10.1134/S0010952512010029>, 2012.
- 1052 Guharay, A., Venkat Ratnam, M., Nath, D., and Dumka, U. C.: Investigation of saturated gravity waves in the tropical lower
1053 atmosphere using radiosonde observations, *Radio Sci.*, 45, <https://doi.org/10.1029/2010RS004372>, 2010.
- 1054 Hei, H., Tsuda, T., and Hirooka, T.: Characteristics of atmospheric gravity wave activity in the polar regions revealed by GPS
1055 radio occultation data with CHAMP, *J. Geophys. Res. Atmospheres*, 113, <https://doi.org/10.1029/2007JD008938>, 2008.
- 1056 Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas, J., Peubey, C., Radu, R., Schepers,
1057 D., Simmons, A., Soci, C., Abdalla, S., Abellan, X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara,
1058 G., Dahlgren, P., Dee, D., Diamantakis, M., Dragani, R., Flemming, J., Forbes, R., Fuentes, M., Geer, A., Haimberger, L.,
1059 Healy, S., Hogan, R. J., Hólm, E., Janisková, M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay, P.,
1060 Rozum, I., Vamborg, F., Villaume, S., and Thépaut, J.-N.: The ERA5 global reanalysis, *Q. J. R. Meteorol. Soc.*, 146, 1999–
1061 2049, <https://doi.org/10.1002/qj.3803>, 2020.
- 1062 John, S. R. and Kumar, K. K.: A discussion on the methods of extracting gravity wave perturbations from space-based
1063 measurements, *Geophys. Res. Lett.*, 40, 2406–2410, <https://doi.org/10.1002/grl.50451>, 2013.
- 1064 Kalisch, S., Chun, H.-Y., Ern, M., Preusse, P., Trinh, Q. T., Eckermann, S. D., and Riese, M.: Comparison of simulated and
1065 observed convective gravity waves, *J. Geophys. Res. Atmospheres*, 121, 13,474–13,492,
1066 <https://doi.org/10.1002/2016JD025235>, 2016.
- 1067 Kalisch, S., Kang, M.-J., and Chun, H.-Y.: Impact of Convective Gravity Waves on the Tropical Middle Atmosphere During
1068 the Madden-Julian Oscillation, *J. Geophys. Res. Atmospheres*, 123, 8975–8992, <https://doi.org/10.1029/2017JD028221>, 2018.
- 1069 Khaykin, S. M., Hauchecorne, A., Mzé, N., and Keckhut, P.: Seasonal variation of gravity wave activity at midlatitudes from
1070 7 years of COSMIC GPS and Rayleigh lidar temperature observations, *Geophys. Res. Lett.*, 42, 1251–1258,
1071 <https://doi.org/10.1002/2014GL062891>, 2015.
- 1072 Kiladis, G., Dias, J., Straub, K., Wheeler, M., Tulich, S., Kikuchi, K., Weickmann, K., and Ventrice, M.: A Comparison of
1073 OLR and Circulation-Based Indices for Tracking the MJO, *Mon. Weather Rev.*, 142, 1697–1715,
1074 <https://doi.org/10.1175/MWR-D-13-00301.1>, 2014.
- 1075 Krisch, I., Hindley, N. P., Reitebuch, O., and Wright, C. J.: On the derivation of zonal and meridional wind components from
1076 Aeolus horizontal line-of-sight wind, *Atmospheric Meas. Tech.*, 15, 3465–3479, <https://doi.org/10.5194/amt-15-3465-2022>,
1077 2022.
- 1078 Kruse, C. G., Richter, J. H., Alexander, M. J., Bacmeister, J. T., Heale, C., and Wei, J.: Gravity Wave Drag Parameterizations
1079 for Earth’s Atmosphere, in: *Fast Processes in Large-Scale Atmospheric Models*, American Geophysical Union (AGU), 229–
1080 256, <https://doi.org/10.1002/9781119529019.ch9>, 2023.

1081 Kumari, K., Wu, H., Long, A., Lu, X., and Oberheide, J.: Mechanism Studies of Madden-Julian Oscillation Coupling Into the
1082 Mesosphere/Lower Thermosphere Tides Using SABER, MERRA-2, and SD-WACCMX, *J. Geophys. Res. Atmospheres*, 126,
1083 e2021JD034595, <https://doi.org/10.1029/2021JD034595>, 2021.

1084 Kursinski, E. R., Hajj, G. A., Schofield, J. T., Linfield, R. P., and Hardy, K. R.: Observing Earth's atmosphere with radio
1085 occultation measurements using the Global Positioning System, *J. Geophys. Res. Atmospheres*, 102, 23429–23465,
1086 <https://doi.org/10.1029/97JD01569>, 1997.

1087 Ladstädter, F., Steiner, A. K., Foelsche, U., Haimberger, L., Tavolato, C., and Kirchengast, G.: An assessment of differences
1088 in lower stratospheric temperature records from (A)MSU, radiosondes, and GPS radio occultation, *Atmospheric Meas. Tech.*,
1089 4, 1965–1977, <https://doi.org/10.5194/amt-4-1965-2011>, 2011.

1090 Lim, Y. and Son, S.-W.: QBO Wind Influence on MJO-Induced Temperature Anomalies in the Upper Troposphere and Lower
1091 Stratosphere in an Idealized Model, <https://doi.org/10.1175/JAS-D-21-0296.1>, 2022.

1092 Luna, D., Alexander, P., and de la Torre, A.: Evaluation of uncertainty in gravity wave potential energy calculations through
1093 GPS radio occultation measurements, *Adv. Space Res.*, 52, 879–882, <https://doi.org/10.1016/j.asr.2013.05.015>, 2013.

1094 Mack, L. R. and Jay, B. E.: The partition of energy in standing gravity waves of finite amplitude, *J. Geophys. Res.* 1896-1977,
1095 72, 573–581, <https://doi.org/10.1029/JZ072i002p00573>, 1967.

1096 MARQUARDT, C. and Healy, S.: Measurement Noise and Stratospheric Gravity Wave Characteristics Obtained from GPS
1097 Occultation Data, *J. Meteorol. Soc. Jpn. - J METEOROL SOC JPN*, 83, 417–428, <https://doi.org/10.2151/jmsj.83.417>, 2005.

1098 Martin, Z., Son, S.-W., Butler, A., Hendon, H., Kim, H., Sobel, A., Yoden, S., and Zhang, C.: The influence of the quasi-
1099 biennial oscillation on the Madden–Julian oscillation, *Nat. Rev. Earth Environ.*, 2, 477–489, [https://doi.org/10.1038/s43017-](https://doi.org/10.1038/s43017-021-00173-9)
1100 021-00173-9, 2021.

1101 Munday, C., Engelstaedter, S., Ouma, G., Ogutu, G., Olago, D., Ong'ech, D., Lees, T., Wanguba, B., Nkatha, R., Ogolla, C.,
1102 Gàlgalo, R. A., Dokata, A. J., Kirui, E., Hope, R., and Washington, R.: Observations of the Turkana Jet and the East African
1103 Dry Tropics: The RIFTJet Field Campaign, *Bull. Am. Meteorol. Soc.*, 103, E1828–E1842, [https://doi.org/10.1175/BAMS-D-](https://doi.org/10.1175/BAMS-D-21-0214.1)
1104 21-0214.1, 2022.

1105 Muñoz-Sabater, J., Dutra, E., Agustí-Panareda, A., Albergel, C., Arduini, G., Balsamo, G., Boussetta, S., Choulga, M.,
1106 Harrigan, S., Hersbach, H., Martens, B., Miralles, D. G., Piles, M., Rodríguez-Fernández, N. J., Zsoter, E., Buontempo, C.,
1107 and Thépaut, J.-N.: ERA5-Land: a state-of-the-art global reanalysis dataset for land applications, *Earth Syst. Sci. Data*, 13,
1108 4349–4383, <https://doi.org/10.5194/essd-13-4349-2021>, 2021.

1109 Nastrom, G. D., Hansen, A. R., Tsuda, T., Nishida, M., and Ware, R.: A comparison of gravity wave energy observed by VHF
1110 radar and GPS/MET over central North America, *J. Geophys. Res. Atmospheres*, 105, 4685–4687,
1111 <https://doi.org/10.1029/1999JD901164>, 2000.

1112 Orr, A., Bechtold, P., Scinocca, J., Ern, M., and Janiskova, M.: Improved Middle Atmosphere Climate and Forecasts in the
1113 ECMWF Model through a Nonorographic Gravity Wave Drag Parameterization, <https://doi.org/10.1175/2010JCLI3490.1>,
1114 2010.

1115 Pahlavan, H. A., Fu, Q., Wallace, J. M., and Kiladis, G. N.: Revisiting the Quasi-Biennial Oscillation as Seen in ERA5. Part
1116 I: Description and Momentum Budget, <https://doi.org/10.1175/JAS-D-20-0248.1>, 2021.

- 1117 Plougonven, R. and Zhang, F.: Internal gravity waves from atmospheric jets and fronts, *Rev. Geophys.*, 52, 33–76,
1118 <https://doi.org/10.1002/2012RG000419>, 2014.
- 1119 Plougonven, R., Jewtoukoff, V., Cámara, A. de la, Lott, F., and Hertzog, A.: On the Relation between Gravity Waves and
1120 Wind Speed in the Lower Stratosphere over the Southern Ocean, <https://doi.org/10.1175/JAS-D-16-0096.1>, 2017.
- 1121 Podglajen, A., Hertzog, A., Plougonven, R., and Žagar, N.: Assessment of the accuracy of (re)analyses in the equatorial lower
1122 stratosphere, *J. Geophys. Res. Atmospheres*, 119, 11,166–11,188, <https://doi.org/10.1002/2014JD021849>, 2014.
- 1123 Randel, W. J. and Wu, F.: Kelvin wave variability near the equatorial tropopause observed in GPS radio occultation
1124 measurements, *J. Geophys. Res. Atmospheres*, 110, <https://doi.org/10.1029/2004JD005006>, 2005.
- 1125 Randel, W. J., Wu, F., and Podglajen, A.: Equatorial Waves, Diurnal Tides and Small-Scale Thermal Variability in the Tropical
1126 Lower Stratosphere From COSMIC-2 Radio Occultation, *J. Geophys. Res. Atmospheres*, 126, e2020JD033969,
1127 <https://doi.org/10.1029/2020JD033969>, 2021.
- 1128 Ratynski, M., Khaykin, S., Hauchecorne, A., Wing, R., Cammas, J.-P., Hello, Y., and Keckhut, P.: Validation of Aeolus wind
1129 profiles using ground-based lidar and radiosonde observations at Réunion island and the Observatoire de Haute-Provence,
1130 *Atmospheric Meas. Tech.*, 16, 997–1016, <https://doi.org/10.5194/amt-16-997-2023>, 2023.
- 1131 The NWP impact of Aeolus Level-2B winds at ECMWF: [https://www.ecmwf.int/en/elibrary/81546-nwp-impact-aeolus-level-](https://www.ecmwf.int/en/elibrary/81546-nwp-impact-aeolus-level-2b-winds-ecmwf)
1132 [2b-winds-ecmwf](https://www.ecmwf.int/en/elibrary/81546-nwp-impact-aeolus-level-2b-winds-ecmwf), last access: 28 June 2025.
- 1133 Rennie, M. P., Isaksen, L., Weiler, F., Jos de Kloe, Kanitz, T., and Reitebuch, O.: The impact of Aeolus wind retrievals on
1134 ECMWF global weather forecasts, *Q. J. R. Meteorol. Soc.*, 147, 3555–3586, <https://doi.org/10.1002/qj.4142>, 2021.
- 1135 Šácha, P., Foelsche, U., and Pišoft, P.: Analysis of internal gravity waves with GPS RO density profiles, *Atmospheric Meas.*
1136 *Tech.*, 7, 4123–4132, <https://doi.org/10.5194/amt-7-4123-2014>, 2014.
- 1137 Schmidt, T., de la Torre, A., and Wickert, J.: Global gravity wave activity in the tropopause region from CHAMP radio
1138 occultation data, *Geophys. Res. Lett.*, 35, <https://doi.org/10.1029/2008GL034986>, 2008.
- 1139 Schmidt, T., Alexander, P., and de la Torre, A.: Stratospheric gravity wave momentum flux from radio occultations, *J.*
1140 *Geophys. Res. Atmospheres*, 121, 4443–4467, <https://doi.org/10.1002/2015JD024135>, 2016.
- 1141 Scinocca, J. F.: An Accurate Spectral Nonorographic Gravity Wave Drag Parameterization for General Circulation Models,
1142 2003.
- 1143 Song, L. and Wu, R.: Modulation of the Westerly and Easterly Quasi-Biennial Oscillation Phases on the Connection between
1144 the Madden–Julian Oscillation and the Arctic Oscillation, *Atmosphere*, 11, 175, <https://doi.org/10.3390/atmos11020175>, 2020.
- 1145 Stephan, C. C. and Mariaccia, A.: The signature of the tropospheric gravity wave background in observed mesoscale motion,
1146 *Weather Clim. Dyn.*, 2, 359–372, <https://doi.org/10.5194/wcd-2-359-2021>, 2021.
- 1147 Stephan, C. C., Žagar, N., and Shepherd, T. G.: Waves and coherent flows in the tropical atmosphere: new opportunities, old
1148 challenges, *Q. J. R. Meteorol. Soc.*, 147, 2597–2624, <https://doi.org/10.1002/qj.4109>, 2021.
- 1149 Tsuda, T., Nishida, M., Rocken, C., and Ware, R. H.: A Global Morphology of Gravity Wave Activity in the Stratosphere
1150 Revealed by the GPS Occultation Data (GPS/MET), *J. Geophys. Res. Atmospheres*, 105, 7257–7273,
1151 <https://doi.org/10.1029/1999JD901005>, 2000.

- 1152 Tsuda, T., Ratnam, M. V., May, P. T., Alexander, M. J., Vincent, R. A., and MacKinnon, A.: Characteristics of gravity waves
1153 with short vertical wavelengths observed with radiosonde and GPS occultation during DAWEX (Darwin Area Wave
1154 Experiment), *J. Geophys. Res. Atmospheres*, 109, <https://doi.org/10.1029/2004JD004946>, 2004.
- 1155 VanZandt, T. E.: A model for gravity wave spectra observed by Doppler sounding systems, *Radio Sci.*, 20, 1323–1330,
1156 <https://doi.org/10.1029/RS020i006p01323>, 1985.
- 1157 Waite, M. L. and Snyder, C.: The Mesoscale Kinetic Energy Spectrum of a Baroclinic Life Cycle,
1158 <https://doi.org/10.1175/2008JAS2829.1>, 2009.
- 1159 Wang, L. and Alexander, M. J.: Global estimates of gravity wave parameters from GPS radio occultation temperature data, *J.*
1160 *Geophys. Res. Atmospheres*, 115, <https://doi.org/10.1029/2010JD013860>, 2010.
- 1161 Weiler, F., Kanitz, T., Wernham, D., Rennie, M., Huber, D., Schillinger, M., Saint-Pe, O., Bell, R., Parrinello, T., and
1162 Reitebuch, O.: Characterization of dark current signal measurements of the ACCDs used on board the Aeolus satellite,
1163 *Atmospheric Meas. Tech.*, 14, 5153–5177, <https://doi.org/10.5194/amt-14-5153-2021>, 2021.
- 1164 Woiwode, W., Dörnbrack, A., Geldenhuys, M., Friedl-Vallon, F., Giez, A., Gulde, T., Höpfner, M., Johansson, S., Kaifler, B.,
1165 Kleinert, A., Krasauskas, L., Kretschmer, E., Maucher, G., Neubert, T., Nordmeyer, H., Piesch, C., Preusse, P., Rapp, M.,
1166 Riese, M., Schumann, U., and Ungermann, J.: Non-Orographic Gravity Waves and Turbulence Caused by Merging Jet
1167 Streams, *J. Geophys. Res. Atmospheres*, 128, e2022JD038097, <https://doi.org/10.1029/2022JD038097>, 2023.
- 1168 Wright, C. J. and Hindley, N. P.: How well do stratospheric reanalyses reproduce high-resolution satellite temperature
1169 measurements?, *Atmospheric Chem. Phys.*, 18, 13703–13731, <https://doi.org/10.5194/acp-18-13703-2018>, 2018.
- 1170 Wright, C. J., Ungermann, J., Preusse, P., and Polichtchouk, I.: Using sub-limb observations to measure gravity waves excited
1171 by convection, *Npj Microgravity*, 9, 14, <https://doi.org/10.1038/s41526-023-00259-2>, 2023.
- 1172 Wüst, S. and Bittner, M.: Gravity wave reflection: Case study based on rocket data, *J. Atmospheric Sol.-Terr. Phys.*, 70, 742–
1173 755, <https://doi.org/10.1016/j.jastp.2007.10.010>, 2008.
- 1174 Yang, S.-S., Pan, C.-J., and Das, U.: Investigating the Spatio-Temporal Distribution of Gravity Wave Potential Energy over
1175 the Equatorial Region Using the ERA5 Reanalysis Data, *Atmosphere*, 12, 311, <https://doi.org/10.3390/atmos12030311>, 2021.
- 1176 Žagar, N., Gustafsson, N., and Källén, E.: Variational data assimilation in the tropics: The impact of a background-error
1177 constraint, *Q. J. R. Meteorol. Soc.*, 130, 103–125, <https://doi.org/10.1256/qj.03.13>, 2004.
- 1178 Žagar, N., Rennie, M., and Isaksen, L.: Uncertainties in Kelvin Waves in ECMWF Analyses and Forecasts: Insights From
1179 Aeolus Observing System Experiments, *Geophys. Res. Lett.*, 48, e2021GL094716, <https://doi.org/10.1029/2021GL094716>,
1180 2021.
- 1181 Zhang, K., Randel, W. J., and Fu, R.: Relationships between outgoing longwave radiation and diabatic heating in reanalyses,
1182 *Clim. Dyn.*, 49, 2911–2929, <https://doi.org/10.1007/s00382-016-3501-0>, 2017.
- 1183 Zhang, S. D. and Yi, F.: A statistical study of gravity waves from radiosonde observations at Wuhan (30° N, 114° E) China,
1184 *Ann. Geophys.*, 23, 665–673, <https://doi.org/10.5194/angeo-23-665-2005>, 2005.
- 1185 Zhang, S. D., Yi, F., Huang, C. M., and Zhou, Q.: Latitudinal and seasonal variations of lower atmospheric inertial gravity
1186 wave energy revealed by US radiosonde data, *Ann. Geophys.*, 28, 1065–1074, <https://doi.org/10.5194/angeo-28-1065-2010>,
1187 2010.

1188 Zhou, X., Liu, L., Yue, X., Chen, G., and Lu, X.: Observed Responses of Gravity Wave Momentum Fluxes to the Madden–
1189 Julian Oscillation Around the Extratropical Mesopause Using Mohe Meteor Radar Observations, *J. Geophys. Res.*
1190 *Atmospheres*, 129, e2024JD041447, <https://doi.org/10.1029/2024JD041447>, 2024.

1191