

1 **Contrasting Aerosol Mixing States at Inland and Coastal Sites: An**  
2 **Entropy-Based Metric for CCN Activity**

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25 **Abstract**

26 Simplified assumptions of aerosol mixing states in modeling studies often  
27 introduce substantial uncertainties in estimating cloud condensation nuclei (CCN)  
28 concentrations ( $N_{\text{CCN}}$ ) and their climatic impacts. This study systematically investigates  
29 the contrasting relationships between mixing states and CCN activity by combining  
30 field measurements of probability distribution function of the hygroscopicity with the  
31 algorithm of entropy at two inland and coastal sites. We show distinct seasonal  
32 variations of aerosol mixing state. In winter, externally-mixed particles dominated both  
33 sites, with comparable mixing state indices ( $\chi$ ) of  $0.38 \pm 0.12$  and  $0.39 \pm 0.09$  respectively  
34 for coastal air and inland air. However, summer measurements showed pronounced  
35 differences: photochemical processes promoted significantly higher internal mixing in  
36 coastal aerosols ( $\chi = 0.69 \pm 0.19$ ), whereas inland  $\chi$  values only increased moderately to  
37  $0.47 \pm 0.12$ . Aerosol mixing state is largely influenced by primary emissions and  
38 secondary formation process. Externally-mixed particles originate chiefly from  
39 anthropogenic emissions in inland or sea salt in coastal. During the aging process,  
40 particles become more internally-mixed as the enhanced fraction of more-hygroscopic  
41 mode. A universal logarithmic correlation was identified between the critical diameter  
42 ( $D_{\text{crit}}$ ) characterizing CCN activity and  $\chi$  ( $D_{\text{crit}} = -32.15 \ln(\chi) + 84.71$ , Pearson  $r = -0.74$ ),  
43 but with distinct decrement rates for coastal vs. inland aerosols. Further analysis reveals  
44 a 0.1 increase in  $\chi$  enhanced winter  $N_{\text{CCN}}$  by 39–65% at the supersaturation of 0.2%,  
45 whereas this effect diminished to ~9% in summer. These results underscore that mixing

46 states exert different control over  $N_{CCN}$  in diverse environments. Our work provides  
47 critical constraints for parameterizing fine aerosols CCN activity in climate models.

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## 51 **1. Introduction**

52 Atmospheric cloud condensation nuclei (CCNs) are complex mixtures of organic  
53 and inorganic components. Their chemical and physical properties make quantifying  
54 aerosol-cloud interactions challenging (Liu et al., 2018; Rosenfeld et al., 2019; Xu et  
55 al., 2022, 2024; Virtanen et al., 2025), introducing uncertainties into climate effect  
56 assessments (Charlson et al., 1992; Shrivastava et al., 2017; IPCC, 2021; Chen et al.,  
57 2022a; Manavi et al., 2025). Accurate climate model predictions of aerosol impacts  
58 require understanding aerosol mixing states under different atmospheric conditions and  
59 their effects on CCN activity (Ching et al., 2016; Zheng et al., 2021a). Current models  
60 often oversimplify mixing states by assuming pure internal or external mixing (Winkler,  
61 1973; Stevens et al., 2019; Riemer et al., 2019; Zheng et al., 2021b). This is problematic  
62 because mixing states directly determine particle hygroscopicity distribution and CCN  
63 estimates (Wang et al., 2010; Tao et al., 2024). For example, internal-mixed aerosol  
64 particles have unimodal hygroscopicity distribution, while the external-mixed particles  
65 are characterized by the bimodal/trimodal or partly overlapping structures (Spitieri et  
66 al., 2023; Liu et al., 2025). Such simplifications can lead to significant errors, e.g.,  
67 Sotiropoulou et al. (2007) found that mixing state assumptions caused two-fold  $N_{CCN}$

68 estimation errors in global models.

69        Systematic observations across diverse environments are critical because aerosol  
70 mixing states exhibit pronounced spatial-temporal variations (Ye et al., 2018; Hughes  
71 et al., 2018; Liu et al., 2025). For example, continental and coastal regions present  
72 contrasting scenarios (Ramachandran et al., 2016). The continental areas are dominated  
73 by anthropogenic emissions, where aerosol aging is driven by industrial and traffic-  
74 related pollutants (Huang et al., 2014; Ren et al., 2023). Particles here undergo  
75 progressive internal mixing via photochemical process and heterogenous reactions,  
76 altering their hygroscopic properties (Ervens et al., 2010; Tao et al., 2021). While the  
77 coastal regions feature dynamic interactions between marine aerosols (e.g., sea salt) and  
78 continental pollutants (Schill et al., 2015; Collins et al., 2013; Cheung et al., 2020).  
79 Seasonal shifts in air mass sources (e.g., marine vs. continental dominance) might  
80 create unique mixing state patterns (Xu et al., 2020, 2021a). For instance, summer  
81 photochemical aging and heterogenous processes in coastal areas can enhance the  
82 degree of internal mixing, while winter often retains more external mixing due to the  
83 presence of the sea-salt particles with less-hygroscopic organic matter.

84        The continental aerosols influence regional cloud formation, while coastal  
85 aerosols may provide insights into the characteristics of marine aerosols in region. The  
86 properties of marine aerosols are significantly different from those of continental  
87 aerosols, and therefore have distinct climate feedback mechanisms (Bellouin et al.,  
88 2020; Xu et al., 2024; Liu et al., 2024). However, the current models lack regional-  
89 specific mixing state parameters and usually assume uniform mixing in both

90 environments. This could lead to large uncertainties in predicting CCN concentrations,  
91 highlighting the need for site-specific observations. For example, Ren et al. (2018)  
92 found that the impact of aerosol mixing state on CCN activation characteristics ranged  
93 from -34% to +16 % in urban atmosphere. Comparison between a fully internal mixture  
94 assumption and using the mixing state index from the particle-resolved model, Ching  
95 et al. (2017) found the obvious overestimation in CCN concentration estimation.  
96 Especially in the regions eg., Amazon Basin, Central Africa and Indonesia, the particles  
97 appeared to be more external, errors in CCN concentration would increase up to 100%  
98 (Hughes et al., 2018). A detailed exploration of mixing state on CCN concentration in  
99 global scale was conducted by Zheng et al. (2021a), and the results showed that the  
100 mixing state varied spatially with more externally mixed over the North Atlantic Ocean,  
101 off the coasts of Southern Africa, and Australia. Thus, assuming particles with  
102 internally-mixed would introduce errors in CCN concentration of 50-100%.

103       Therefore, for quantifying the aerosol mixing state in the ambient atmosphere,  
104 we apply the algorithm of entropy proposed by Riemer and West (2013) to investigate  
105 the aerosol heterogeneity. This index has been applied to quantify the mixing state more  
106 reasonably both in field campaigns (Zhao et al., 2021; Yuan et al., 2023) and model  
107 simulations (Ching et al., 2016; Zheng et al., 2021a). However, most studies focused  
108 on quantifying the particle heterogeneity in composition (Ching et al., 2019; Fierce et  
109 al., 2020; Zhao et al., 2021). Here we concentrated on evaluating the heterogeneity in  
110 aerosol hygroscopicity for sub-micron particles, which directly related to CCN budget.  
111 By refereeing to Yuan et al. 2023, the heterogeneity in hygroscopicity was investigated

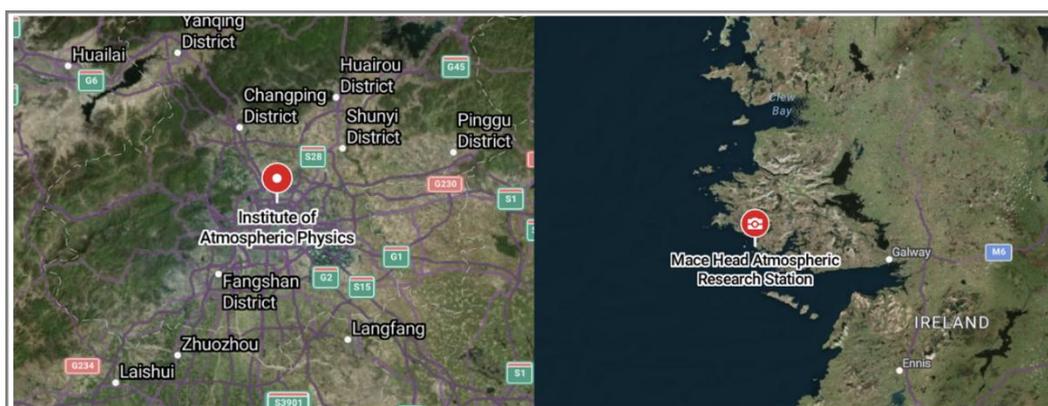
112 by combining in-situ measurements of probability distribution function of the  
113 hygroscopicity with the algorithm of entropy. Briefly, the mixing state index  $\chi$ , is  
114 devised based on the concept of information entropy concerning the distribution of  
115 hygroscopicity across the aerosol population. It varies between 0 (external mixing  
116 completely) and 1 (internal mixing completely). By integrating inland and coastal  
117 measurements, this study will focus on addressing two key gaps, (1) How continental  
118 vs. marine-dominated environments shape aerosol mixing states and CCN activity; (2)  
119 Whether  $\chi$ -based CCN parameterizations show regional dependencies, providing  
120 critical constraints for climate models.

## 121 **2. Data and Methods**

### 122 **2.1 Field Campaigns**

123 The inland atmospheric measurements were conducted for two campaigns from  
124 16 November to 6 December 2016 and 29 May to 13 June 2017 as a part of the Air  
125 Pollution and Human Health (APHH) project (Shi et al., 2019), at the Institute of  
126 Atmospheric Physics, Chinese Academy of Sciences (IAP, 39.97° N, 116.37° E) in  
127 urban Beijing. The campaigns were complemented by the hygroscopicity and CCN  
128 observations and were conducive to provide information on the aerosol hygroscopicity  
129 affecting urban pollutions. This urban site exhibited highly variable aerosol populations  
130 dominated by local anthropogenic sources including vehicular, cooking emissions, and  
131 residential heating. Coastal measurements were performed at the Mace Head  
132 atmospheric research station (MHD, 53.33° N, 9.90° W) from 1 November 2009 to 30

133 January 2010, and summer periods from 11 to 31 August 2009 and July 2010, which  
134 located on the west coast of Ireland. Aerosol particles here experience alternating  
135 influences from polluted continental and clean marine atmospheres. The map of the  
136 sites was shown in Figure 1. More details about the campaigns were given in Fan et al.  
137 (2020) and Xu et al. (2021a).



138  
139 **Fig 1.** Map of the sites in the Inland of the Institute of Atmospheric Physics (IAP) and  
140 Coastal of Mace Head (MHD). (© Google Maps, <https://maps.google.com/>, last access:  
141 2 April 2025).

## 142 **2.2 Instrumentation**

### 143 **Hygroscopicity measurements**

144 The particle hygroscopicity at both sites was characterized using the humidified  
145 tandem differential mobility analyzer (HTDMA). The hygroscopic growth factor (Gf),  
146 defined as the ratio of the particle diameter at the fixed RH (90%) and dry diameter set  
147 in this study for 40, 80, 110, 150, 200 nm at IAP and 35, 50, 75, 110 and 165 nm at  
148 MHD, respectively. The Gf probability density function (Gf-PDF) was derived using  
149 the TDMAinv algorithm (Gysel et al., 2009). The number fraction (NF) of near-

150 hydrophobic mode (NH:  $Gf \leq 1.21$ ), and more hygroscopic mode (MH:  $Gf > 1.21$ ) in  
151 IAP site was referred from Chen et al. (2022b). It was integrated into three modes for  
152 the MHD site with the near-hydrophobic mode (NH:  $1 < Gf < 1.3$ ), more hygroscopic  
153 mode (MH:  $1.3 \leq Gf < 1.85$ ) and sea salt mode (SS:  $Gf \geq 1.85$ ) for further examination  
154 (Xu et al., 2021a).

155 Here for each particle size, the hygroscopicity parameter  $\kappa$  can be subsequently  
156 calculated using  $\kappa$ -Köhler theory (Petters and Kreidenweis, 2007):

$$157 \quad \kappa = (Gf^3 - 1) \cdot \left[ \frac{1}{RH} \exp\left(\frac{4\sigma_s M_w}{RT\rho_w D_d Gf}\right) - 1 \right] \quad (1)$$

158 where RH is the HTDMA relative humidity (90% set in the instrument),  $\sigma_{s/a}$  is the  
159 surface tension of pure water ( $0.072 \text{ mN m}^{-1}$ ),  $M_w$  and  $\rho_w$  are the molecular weight and  
160 the density of pure water,  $R$  is the gas constant, and  $T$  is the absolute temperature,  $D_d$  is  
161 the droplet diameter.

162 Then, the  $\kappa$ -PDF is obtained and normalized as  $\int_0^\infty c(\kappa)d\kappa = 1$ , where  $c(\kappa)$  is  
163 normalized as  $\kappa$ -PDF. Further it was used to calculate the particle population  
164 heterogeneity (Calculation seen in Section 2.3).

## 165 **Chemical components**

166 For the inland atmospheric measurements, the non-refractory submicron aerosol  
167 (smaller than  $1\mu\text{m}$ , NR-PM<sub>1</sub>) chemical composition was quantitatively characterized  
168 using the Aerodyne High-Resolution Time-of-Flight Aerosol Mass Spectrometer (HR-  
169 ToF-AMS) (DeCarlo et al., 2006), including sulfate ( $\text{SO}_4^{2-}$ ), nitrate ( $\text{NO}_3^-$ ), ammonium  
170 ( $\text{NH}_4^+$ ), chloride (Chl) and organics (Org). The black carbon (BC) mass concentration  
171 was determined from the light absorption with a seven-wavelength aethalometer (AE33,

172 Magee Scientific Corp.).

173 Measurements of  $PM_{10}$  in the coastal atmosphere were also performed by the HR-  
174 ToF-AMS, including major inorganic salts (non-sea-salt sulfate,  $nss-SO_4^{2-}$ ;  
175 methanesulfonic acid, MSA;  $NO_3^-$ ;  $NH_4^+$ ) and organic matter. The instrument operation  
176 and calibration have been described in previous studies (Ovadnevaite et al., 2014; Xu  
177 et al., 2019).

### 178 **Aerosol number size distribution and CCN number concentration**

179 Particle number size distributions (PNSD) were measured using an integrated  
180 system consisting of a Differential Mobility Analyzer (DMA; model 3081, TSI Inc.)  
181 coupled with a Condensation Particle Counter (CPC; model 3772, TSI Inc.). During the  
182 measurements at IAP, the PNSD covered the size range of 10-550 nm with a 5-minute  
183 time resolution. It scanned size range of 20-500 nm at MHD with a 10-minute temporal  
184 resolution. The CCN number concentrations were quantified at both sites using a  
185 Droplet Measurement Technologies CCN counter (DMT-CCNc) (Lance et al., 2006).  
186 The instrument's supersaturation (SS) settings were carefully calibrated before and after  
187 each campaign using ammonium sulfate aerosol following Rose et al. (2008). Four  
188 effective supersaturations (SS) were 0.14%, 0.23%, 0.40% and 0.76% at IAP site. Four  
189 SS levels were 0.25%, 0.5%, 0.75% and 1% at MHD site with an uncertainty of  $\pm 0.03\%$ .  
190 Using measurements at set supersaturation of 0.2% as an example explores the CCN  
191 activity in the following discussions.

## 192 2.3 Calculation the heterogeneity for aerosol particles

193 To characterize the heterogeneous distribution of the hygroscopic and non-  
194 hygroscopic components in populations (Chen et al., 2022b), we calculated the mixing  
195 state index ( $\chi$ ) using the  $\kappa$ -PDF, following the methodology of Yuan et al. (2023). Two  
196 surrogate groups in a population of  $N$  aerosol particles were assumed (Zheng et al.,  
197 2021a). One surrogate group consists the non- and/or slightly hygroscopic species with  
198  $\kappa_N$  of  $<0.05$  and another group contains the more hygroscopic species with  $\kappa_H$  of 0.5-  
199 0.6 (Yuan et al., 2023, referred inorganics). Ambient particles typically contain one or  
200 two of the components and the  $\kappa$  lies between 0 and 0.6 at IAP or 0.8 at MHD as shown  
201 in Figure S1. Taking into account the enhanced hydrophilicity of marine aerosols at  
202 MHD site, calculation assuming  $\kappa_H$  values of 0.7 and 0.8 were shown in Fig.  
203 S2. While these variations in  $\kappa_H$  introduced a mean uncertainty of 8% in  $\chi$  values, it  
204 did not significantly affect the seasonal or site comparisons. The volume fraction of two  
205 surrogate groups can be calculated based on the total  $\kappa$  according to the Zdanovskii-  
206 Stokes-Robinson (ZSR) mixing rule (Zdanovskii, 1948; Stokes et al., 1966).

207 The mixing state index  $\chi$  is defined as the affine ratio of the average particle species  
208 diversity ( $D\alpha$ ) and population species diversity ( $D\gamma$ ) as:

$$209 \chi = \frac{D\alpha-1}{D\gamma-1} \quad (2)$$

210 The average per-particle species diversity  $D\alpha$  can be calculated as follows. First,  
211 the mixing entropies at bin  $i$  ( $H_i$ ) are determined according to equation (3),

$$212 H_i = -P_{i,N} \times \ln P_{i,N} - P_{i,H} \times \ln P_{i,H} \quad (3)$$

213 where  $P_{i,N}$  and  $P_{i,H}$  are the volume fraction of each group for the  $\kappa$ -PDF with X bins at

214 bin  $i$  ( $i=1,2,\dots,X$ ), and can be determined from the  $P_{i,N} + P_{i,H} = 1$  and  $P_{i,N} \times \kappa_N +$   
 215  $P_{i,H} \times \kappa_H = \kappa_i$ . Here  $\kappa_N = 0.01, \kappa_H = 0.6$ ;  $\kappa_i$  represents the hygroscopicity parameter  
 216 at bin  $i$ .

217 Based on the assumption that particles in the same diameter have the same mixing  
 218 entropy  $H_\alpha = \sum_{j=1}^N P_j \times H_j$ ,  $P_j = \frac{V_j}{V_{total}} = \frac{1}{N}$ ; the per-particle mixing entropies  $H_\alpha$  is  
 219 determined according to equation (4),

$$220 \quad H_\alpha = \sum_{i=1}^X H_i \times c(\kappa)_i \times \Delta\kappa \quad (4)$$

221 where  $c(\kappa)_i$  is the probability density of the normalized  $\kappa$ -PDF at bin  $i$ , and  $\Delta\kappa$   
 222 represents the bin width. Then, the average per-particle species diversity  $D_\alpha$  can be  
 223 determined as  $D_\alpha = e^{H_\alpha}$ ;

224 The bulk population species diversity  $D_\gamma$  can be calculated as follows. First, the  
 225 aerosol population of the mixing entropy can be calculated as equation (5):

$$226 \quad H_\gamma = -P_N \times \ln P_N - P_H \times \ln P_H \quad (5)$$

227 where  $P_N$  and  $P_H$  are the volume fraction of the non-hygroscopic and hygroscopic  
 228 components in the population, and can be calculated by equation (6) and (7):

$$229 \quad P_N = \sum_{i=1}^X P_{i,N} \times c(\kappa)_i \times \Delta\kappa \quad (6)$$

$$230 \quad P_H = \sum_{i=1}^X P_{i,H} \times c(\kappa)_i \times \Delta\kappa \quad (7)$$

231 Then, the bulk population species diversity  $D_\gamma$  can be determined as  $D_\gamma = e^{H_\gamma}$ .

232 Here, the definition of surrogate species as supersets encompassing hygroscopicity  
 233 heterogeneity implies that the heterogeneity parameter  $\chi$  ranges from 0 to 1. When the  
 234 mixing index  $\chi$  approaches 0, it indicates a completely segregated state where  
 235 hygroscopic and non-hygroscopic species reside in distinct particles. While for the case

236 the mixing index  $\chi$  to be 1 represents that the non-hygroscopic and hygroscopic species  
237 distributing homogeneously throughout the aerosol population.

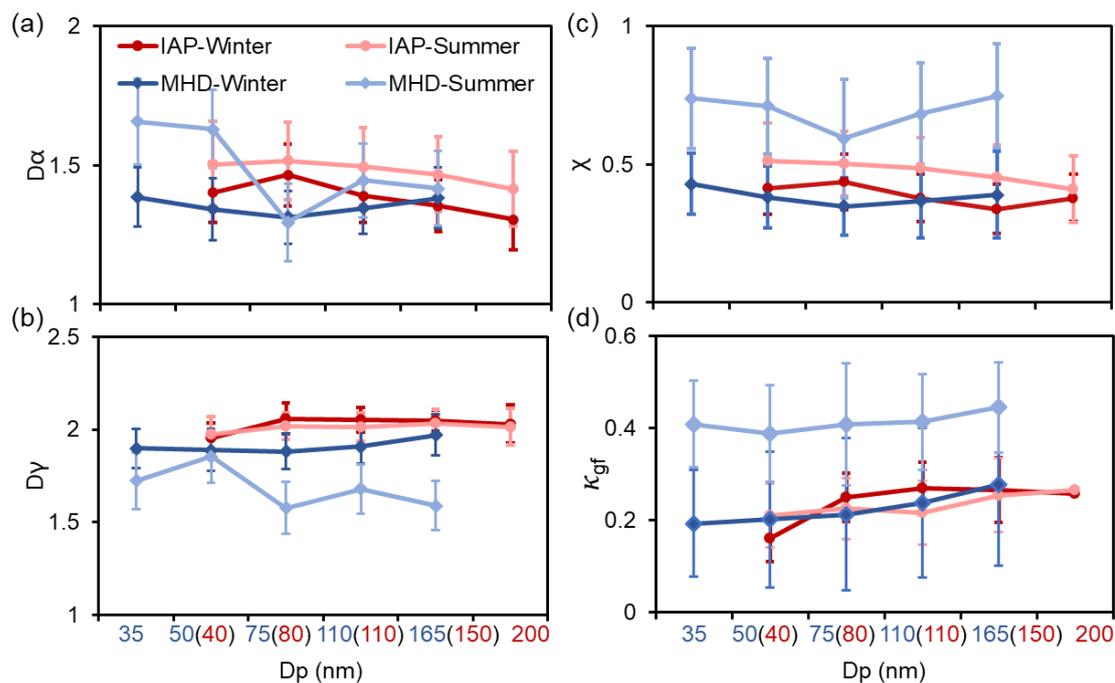
### 238 **3. Result and Discussion**

#### 239 **3.1 Comparison of the heterogeneity in the inland and coastal atmosphere**

240 To characterize the hygroscopic heterogeneity of atmospheric aerosols, Figure 2  
241 depicts variations in mixing state metrics ( $D\alpha$ ,  $D\gamma$ ,  $\chi$ ) and the hygroscopic parameter  
242 ( $\kappa_{gf}$ ) across particle size distributions. The  $D\alpha$  and  $\chi$  decrease with increasing particle  
243 diameter, accompanied by higher  $\kappa_{gf}$  values at IAP site. This trend indicates that inland  
244 particle populations tend to homogenize into hygroscopic compositions through  
245 primary particle aging or secondary formation processes (Liu et al., 2025; Chen et al.,  
246 2022b; Zhong et al., 2022). In contrast, particles exhibit a non-monotonic pattern at  
247 MHD site:  $D\alpha$  and  $\chi$  decrease for Aitken-mode particles (<100 nm) but increase for  
248 accumulation-mode particles. The  $\kappa_{gf}$  shows consistent size-dependent increases in  
249 both winter and summer campaigns.

250 Notably, the mixing state metrics exhibit a pronounced minimum at 75 nm  
251 particles, influenced by distinct mechanisms: winter minima reflect the high sea salt  
252 fraction, while summer minima are driven by anthropogenic organic matter (Cheung et  
253 al., 2020; Xu et al., 2021a). Lower winter  $\chi$  values—coupled with broader  $\kappa$ -PDF  
254 distributions—indicate stronger external mixing and compositional diversity compared  
255 to summer (Fig. S1). Seasonal  $\chi$  and  $\kappa_{gf}$  disparities are more pronounced at MHD site,  
256 primarily driven by the seasonal alternation of marine and anthropogenic emission

257 sources.

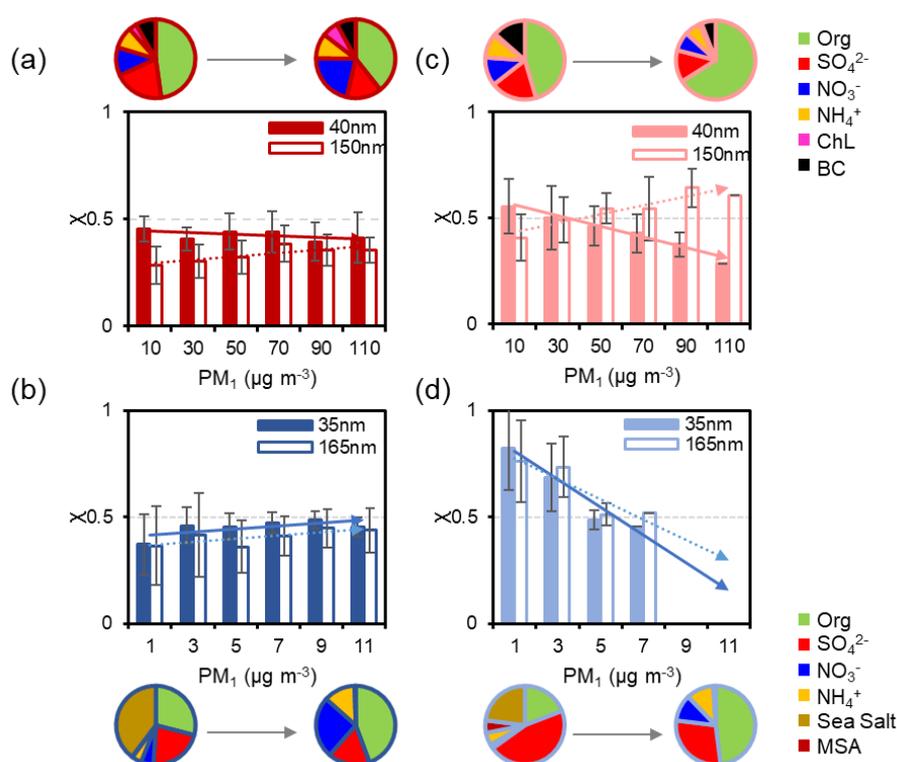


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259 **Fig 2.** Mean values of the  $D\alpha$  (a),  $D\gamma$  (b),  $\chi$  (c) and  $\kappa_{gf}$  (d) for aerosols of five diameters  
260 during winter and summer periods at IAP and Mace Head sites.

261 Ultrafine particles (40 nm in IAP vs. 35 nm in MHD, Aitken mode) and larger  
262 particles (150 nm in IAP vs. 165 nm in MHD, accumulation mode) are selected to  
263 investigate distinct evolutionary processes of aerosol heterogeneity (Fig. 3 and Fig. S3).  
264 With the increasing of  $PM_{10}$  concentration during winter, the variation in  $\chi$  values exhibit  
265 only minor both at the IAP and MHD sites, generally fluctuating between  
266 approximately  $-0.04$  and  $0.08$  (Fig. 3a and b). Inland accumulation-mode particles  
267 show a modest increase in  $\chi$ , corresponding with a higher proportion of inorganic salts.  
268 Conversely, at MHD site, the composition fraction shifts from a sea-salt dominance  
269 toward organic matter, accompanied by a  $\sim 20\%$  increase in nitrate content (Fig. 3b). In  
270 summer, the variation of  $\chi$  with PM concentration becomes markedly pronounced at

271 both IAP and MHD stations. For example,  $\chi$  for 40 nm particles decreases as PM  
 272 increases at IAP site (Fig. 3c). The elevated particle heterogeneity mainly arises from  
 273 the locally primary emissions, corresponding to the enhanced primary organic  
 274 emissions as shown in Fig. S4. It appeared more pronounced during evening rush hours.  
 275 In contrast,  $\chi$  for 150 nm particles increases from  $\sim 0.40$  to  $\sim 0.60$  with rising PM,  
 276 reflecting enhanced secondary formation and internal mixing during pollution process  
 277 that render the particle population more homogeneous. At coastal sites,  $\chi$  declines with  
 278 rising PM by approximately 0.37 for 35 nm particles and 0.24 for 165 nm particles,  
 279 mirroring the shift in chemical composition makeup from inorganic dominance to  
 280 greater organic content (Fig. 3d).



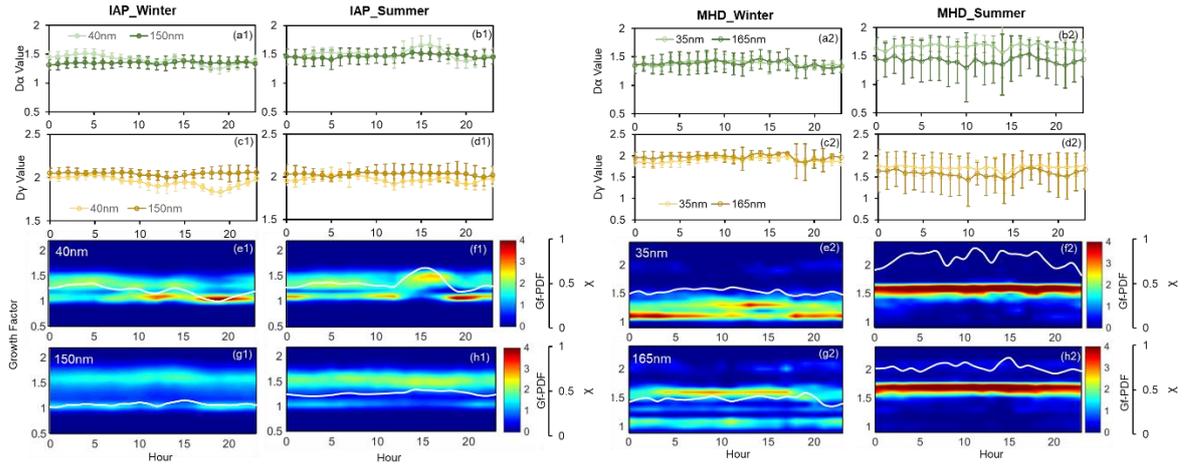
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282 **Fig 3.** Variation of the average  $\chi$  for 40 nm and 150 nm particles at IAP and 35 nm and

283 165 nm at MHD site with the particle mass concentration in IAP-winter (a), IAP-

284 summer (b), MHD-winter (c) and MHD-summer (d). The pie charts represent the  
285 average mass fraction during four field measurements.

286 Diurnal variations of mixing state metrics ( $D\alpha$ ,  $D\gamma$ , Gf-PDF and  $\chi$ ) at IAP and  
287 MHD sites are shown in Figure 4. In IAP-winter, particles exhibited steeper declines in  
288  $D\alpha$  and  $\chi$  during evening rush hours than summer, indicating a higher fraction of non-  
289 hygroscopic particles (40 nm) from fresh traffic emissions (Fig. 4a1 and S4).  
290 Concurrently, reduced  $D\gamma$  values suggest that the bulk population consists of uniformly  
291 distributed less-hygroscopic (LH) components (Fig. 4c1). Aitken mode particles  
292 showed bimodal and broader Gf-PDF distributions, corresponding to cooking activities  
293 (11:00–13:00 LT) and traffic peaks (17:00–20:00 LT) (Cai et al., 2020). Midday  
294 photochemical aging promoted more internally mixed aerosols (Yang et al., 2012; Liu  
295 et al., 2025), as evidenced by increasing  $D\alpha$  at the urban site (Fig. 4b1). Conversely, the  
296  $\chi$  for accumulation-mode particles showed minimal diurnal variations both in IAP-  
297 winter and IAP-summer. This is mainly due to the dominant hygroscopic mode for 150  
298 nm particles (Fig. 4g), especially during summer, which is mainly from secondary  
299 formation or aging of the primary particles (such as the transformation from primary  
300 organic aerosol (POA) to secondary organic aerosol (SOA) in Fig. S4) (Wang et al.,  
301 2019; Fan et al., 2020).



302

303 **Fig 4.** Diurnal variation of  $D\alpha$ ,  $D\gamma$ , Gf-PDF, and  $\chi$  during winter and summer periods  
 304 for 40 nm and 150 nm aerosols at IAP (a1-h1) and for 35 nm and 165 nm aerosols MHD  
 305 site (a2-h2).

306 For the coastal atmosphere, the mixing state metrics ( $D\alpha$ ,  $D\gamma$ , and  $\chi$ ) of Aitken and  
 307 accumulation mode particles in winter exhibited analogous diurnal patterns,  
 308 characterized by a descending trend at nightfall (Fig. 4a2-h2). This corresponds to an  
 309 enhanced modal distribution of near-hydrophobic (NH) particles at 35 nm and more-  
 310 hygroscopic (MH) particles at 165 nm. In summer,  $D\alpha$  and  $D\gamma$  both trended downward  
 311 during daytime, with the decline of  $D\gamma$  being more pronounced. A conspicuous seasonal  
 312 discrepancy between Aitken and accumulation mode particles was observed in this  
 313 region (Fig. 4a2–h2), where the mixing state index  $\chi$  increased incrementally from  
 314 winter to summer. Specifically, the mean  $\chi$  for 35 nm particles escalated from 0.42 to  
 315 0.80, and for 165 nm particles, it rose from 0.39 to 0.76. This trend demonstrates a  
 316 strong alignment with the spread factor (used as a measure of particle mixing state)  
 317 documented by Xu et al. (2021a).

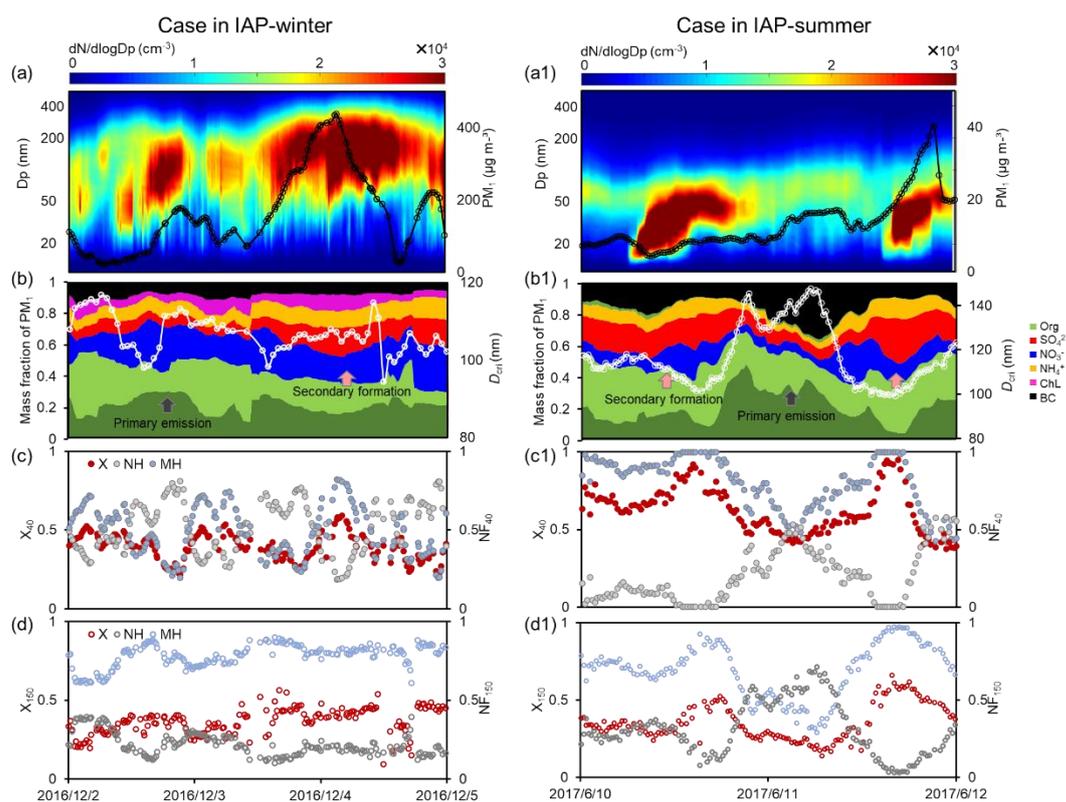
318 Similar to the mixing state  $\chi$ , a very clear seasonal pattern of the aerosol  
319 hygroscopic distribution was found (Fig. 4e-h). In winter, the Gf-PDF diurnal profiles  
320 of both Aitken and accumulation mode particles showed bimodal distribution (Fig. 4e2-  
321 g2) as evident by the number fraction of nearly-hydrophobic and more hygroscopic  
322 modes (Fig. S5). The NH mode was likely to be the anthropogenic organic matter and  
323 biogenic origin from marine mass (Xu et al., 2020), especially for the Aitken mode. The  
324 more hygroscopic and sea salt mode was mostly contributed from the nss-sulfate and  
325 sea salt in winter (Xu et al., 2021a). Analogously, accumulation mode particles with a  
326 higher proportion of MH and SS mode (Fig. S5) primarily attributed to the prevalence  
327 of non-sea-salt sulfate (nss-sulfate) and sea salt in the coastal atmosphere (Xu et al.,  
328 2020). The bimodal and broad of hygroscopic distribution suggested that particles were  
329 more diverse and external mixed, consistent with the lower  $\chi$  value in winter.

330 In contrast, summer observations revealed that Gf-PDFs of both Aitken and  
331 accumulation mode particles transitioned to unimodal distributions, signifying particles  
332 in summer had more homogeneous composition with a large extent of internal mixing  
333 particles (with higher  $\chi$ ). Such diurnal trend in Gf-PDFs was consistent along with the  
334 high number fraction of MH-mode and low NH-mode (Fig. S5). The higher  
335 hygroscopicity and MH mode in summer were largely driven by the enhancement of  
336 sulfate and decrease of organic matter (Fig. S5). And a clear shift from NH to MH mode  
337 at midday might further demonstrate the promotion of photochemical aging in summer  
338 (Xu et al., 2021a).

### 339 **3.2 Impacts of Primary Aerosol Emissions and Secondary Aerosol Formation on** 340 **Aerosol Mixing State**

341 As already noted above, changes in  $\chi$  were clearly associated with the chemical  
342 composition varying with site and season. The relationships between the mixing state  
343 index and the number fraction of hydrophobic and hygroscopic mode during four  
344 campaigns are presented in Figure S6. The  $\chi$  exhibited negative correlations with the  
345 fraction of hydrophobic mode but a positive relationship with the fraction of  
346 hygroscopic particles, highlighting the markedly different effects of the primary  
347 emissions and secondary formation on aerosol mixing state (Tao et al., 2024). To gain  
348 more insight on this effect between inland and coastal atmosphere, four case are  
349 analyzed (Fig. 5 and 6): case for IAP-winter, case for IAP-summer, case for MHD-  
350 winter and case for MHD-summer.

351 Case for IAP-winter is a heavy polluted event with the mean PM mass  
352 concentration increased from 22 to 437  $\mu\text{g m}^{-3}$  (Fig. 5a-d). The 40- and 150-nm  $\chi$   
353 patterns shifted quickly during the pollution periods. With the mass fraction of  
354 hydrophobic compounds (ie., POA) in  $\text{PM}_{10}$  increased, the  $\chi$  of 40-nm particles  
355 decreased from 0.5 to 0.2, that is, an enhanced NH mode and a weaken MH mode (Fig.  
356 5b-c). At this stage, large particles for 150 nm are mainly from aqueous formation with  
357 more proportion of nitrate. The corresponding  $\chi$  of 150 nm was higher. While with that  
358 the mass fractions of secondary organic and inorganic compositions increased, particles  
359 were more internal mixed with  $\chi$  increased to be 0.6 for 40-nm and 0.53 for 150-nm  
360 particles.

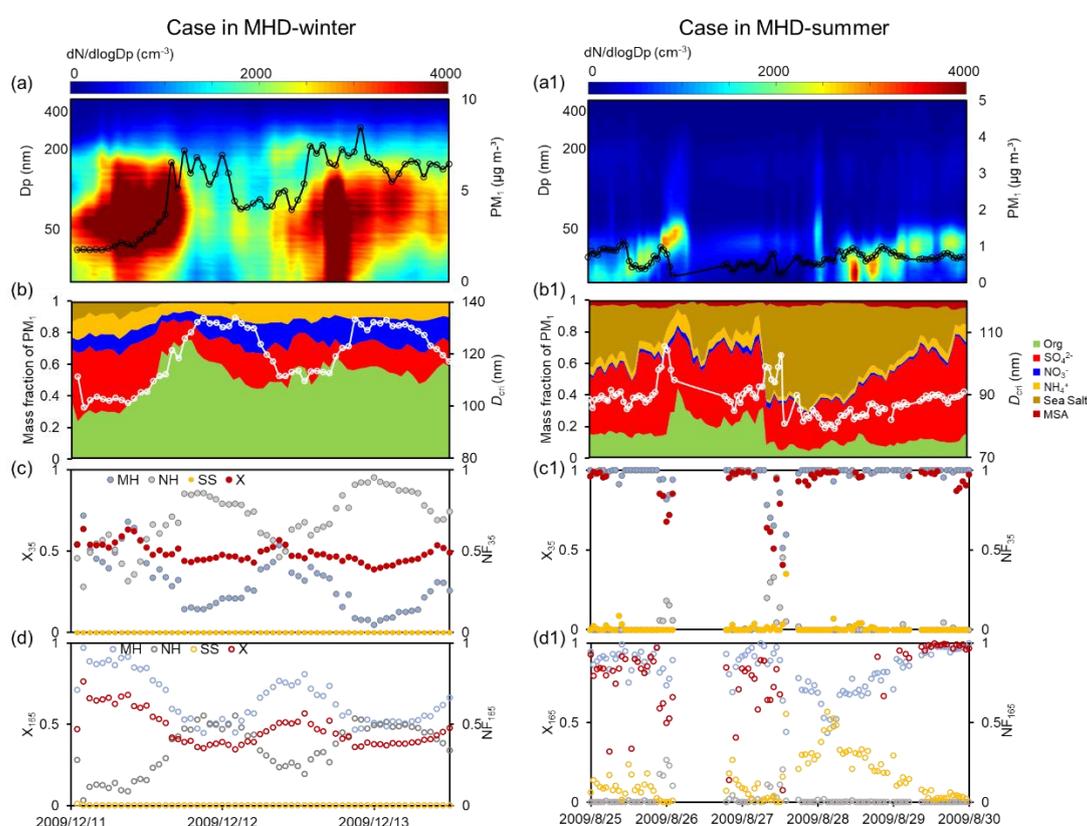


361

362 **Fig 5.** Case in IAP-winter and IAP-summer. Particle number size distribution and  $PM_{10}$   
 363 (a), mass fraction of the  $PM_{10}$  and the critical diameter ( $D_{crit}$ ) (b), mixing state index ( $\chi$ ),  
 364 number fraction of the nearly hydrophobic mode (NH) and more hygroscopic mode  
 365 (MH) for 40 nm particles (c),  $\chi$ , NH and MH for 150 nm particles (d).

366 Case for IAP-summer is the typical new particle formation events (NPF) with the  
 367 mean  $PM_{10}$  of  $13 \mu\text{g m}^{-3}$  (Fig. 5a1-d1). With the evolution of NPF events, the  $\chi$  of 40-  
 368 and 150-nm particles increased to be 0.95 and 0.61 with the enhanced proportion of  
 369 more-hygroscopic components (ie., SOA,  $\text{NO}_3^-$ ,  $\text{SO}_4^{2-}$ ). The  $\chi$  pattern is opposite of that  
 370 the number fraction of NH mode and consistent with the variation of MH mode (Fig.  
 371 S6). Note that a sudden decrease in  $\chi$  on June 11th was disturbed by the strong primary  
 372 emission. The chemical mass fractions showed more POA and black carbon with an  
 373 enhanced NH mode and a weaker MH mode (Fig. 5b1-d1). The  $\chi$  of 40-nm particles

374 decreased to be 0.4 and that of the 150-nm particles decreased to be 0.2. The  $\chi$  patterns  
 375 appear to similar transitions for Aitken and accumulation-mode particles during haze  
 376 and NPF events. The increase in  $\chi$  is synchronous with the increase in MH mode from  
 377 secondary formation but opposite with that of LH mode from primary emissions. This  
 378 implies that the primary emissions would lead particles more external mixing while  
 379 secondary formation would promote aerosol more internal mixed in Inland atmosphere.



380  
 381 **Fig 6.** Case in MHD-winter and MHD-summer. Particle number size distribution and  
 382  $PM_{10}$  (a), mass fraction of the  $PM_{10}$  and the critical diameter ( $D_{cri}$ )(b), mixing state index  
 383 ( $\chi$ ), number fraction of the nearly hydrophobic mode (NH) and more hygroscopic mode  
 384 (MH) for 35 nm particles (c),  $\chi$ , NH and MH for 165 nm particles (d).

385 Case for MHD-winter is a high organic matter pollution event with the mean  $PM_{10}$

386 of  $5.2 \mu\text{g m}^{-3}$  and 52% mass fraction of organics (Fig. 6a-d). Larger presence of  
387 anthropogenic organic matter resulted the NH mode for 35-nm particles to be 95% and  
388 165-nm particles to be 53% (Fig. 6). The  $\chi$  of 35- and 165-nm particles decreased with  
389 the NH mode increased (Fig. S6), similar with the case for IAP site. There was a steady  
390 increase in  $\chi$  when the MH-mode particles started increasing with the increase in mass  
391 fraction inorganics, eg., 35 nm particles showed the mean  $\chi$  increasing from 0.43 to  
392 0.57 and 165 nm particles from 0.35 to 0.6. This indicated that the trend of aerosol  
393 mixing state closely followed the evolution emission and secondary formation.

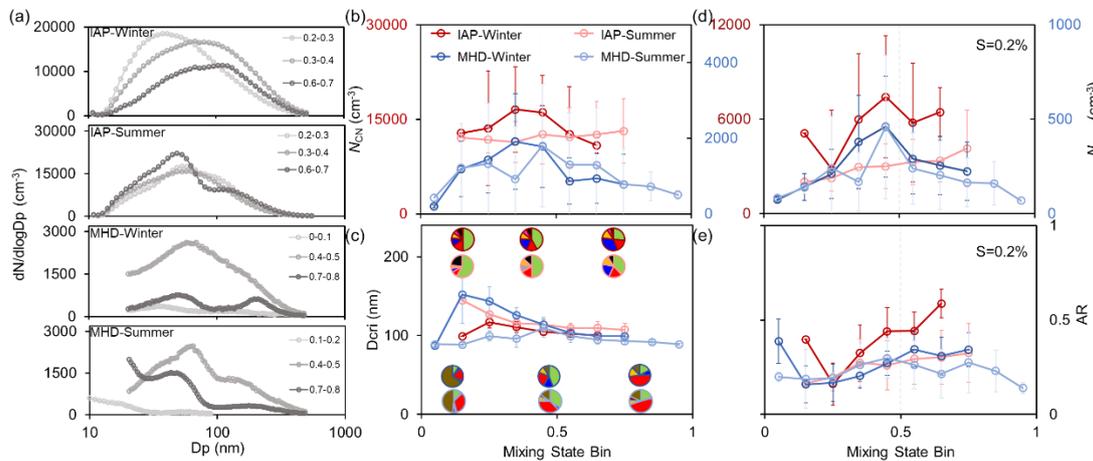
394 Case for MHD-summer is an extremely clean event with the mean  $\text{PM}_{10}$  of  $0.7 \mu\text{g}$   
395  $\text{m}^{-3}$  (Fig. 6a1-d1). The dominated MH mode was found throughout the case, which  
396 could be attributed from the high mass fraction of nss-sulfate (41% average). Compared  
397 with the case in MHD-winter, the mean proportion of organic has decreased to be 15%.  
398 Therefore, the  $\chi$  remains at a high value (mean  $\chi$  of 0.9 for 35-nm and 0.8 for 165-nm  
399 particles). Until August 28th, a stronger increase in the mass fraction of sea salt and  
400 accordingly SS mode in larger-size particles was observed. The  $\chi$  decreased rapidly  
401 with the decrease in MH mode and enhanced SS mode, especially for the accumulation  
402 mode particles, suggesting the sea spray production makes particles more externally  
403 mixed.

404 In summary, these results suggest that the primary emission and secondary  
405 formation drive the hygroscopic distribution and can result in significant variation of  
406 aerosol mixing state  $\chi$  both in Inland and coastal atmosphere. The pattern of  $\chi$  varied  
407 among site and season, highlighting the importance of considering the impact of mixing

408 state on CCN activity.

### 409 3.3 Impact of Mixing State on CCN Activity

410 The mixing state of particle populations undergoes dynamic transformations  
411 during atmospheric aging, profoundly influencing their CCN activity. Unlike prior  
412 studies that assumed mixing states based on chemical component fractions (Yang et al.,  
413 2012; Padró et al., 2012; Ren et al., 2018), this work employs the entropy-derived  
414 mixing state index  $\chi$ , which quantifies the distribution of hygroscopic and non-  
415 hygroscopic species (Zheng et al., 2021a; Ching et al., 2017). The variations of particle  
416 size and chemical composition with the increments of  $\chi$  (ranging from 0 to 1 with the  
417 step of 0.1) are illustrated in Figure 7, presenting key insights of two fundamental  
418 determinants of CCN activity (Dusek et al., 2006).



419

420 **Fig 7.** Comparison of the average particle number size distribution (PNSD) in different  
421 mixing state index ( $\chi$ ) (a), CN number concentration ( $N_{CN}$ ) as a function of  $\chi$  (b),  
422 Critical diameter ( $D_{crit}$ ) at  $S=0.2\%$  and mass fraction of chemical composition as a  
423 function of  $\chi$  (c), CCN number concentration ( $N_{CCN}$ ) (d) and activation ratio (AR) at

424  $S=0.2\%$  a function of  $\chi$  (e).

425 As  $\chi$  increases, the peak diameter ( $D_{\text{peak}}$ ) of the particle number size distribution  
426 (PNSD) shifts toward larger sizes (Fig. 7a and Fig. S7), while peak concentrations occur  
427 within the intermediate  $\chi$  range (0.3–0.6). This trend indicates that CN number  
428 concentration ( $N_{\text{CN}}$ ) first increases, driven by primary emissions and new particle  
429 formation, then decreases due to mixing and aging processes (Fig. 7b). Notably, new  
430 particle formation events frequently occurred in IAP-summer (Fig. S8), corresponding  
431 the gradually increase of  $\chi$ . And the  $\chi$  for Aitken-mode is significantly larger than the  
432 accumulation-mode particles during this period. Thus,  $N_{\text{CN}}$  exhibits a sustained slight  
433 increase as the degree of the internal mixing increases in IAP-summer.

434 The critical diameter ( $D_{\text{cri}}$ )—defined as the minimum size for activation at a given  
435 supersaturation—depends on aerosol hygroscopicity. This hygroscopicity is determined  
436 by both the hygroscopicity and the mass fraction of soluble components (Petters and  
437 Kreidenweis, 2007). Using measurements at supersaturation of 0.2% as an example,  
438 Fig. 7c shows that  $D_{\text{cri}}$  decreases with increasing highly hygroscopic inorganic  
439 components (e.g., sulfate, nitrate) in the inland atmosphere. In contrast, coastal  $D_{\text{cri}}$   
440 exhibits nonlinear variations with  $\chi$ : high external mixing (low  $\chi$ ) elevates  $D_{\text{cri}}$  due to  
441 dominant organic components, reducing sea salt particle fractions. As  $\chi$  increases, the  
442 mass fraction of non-sea-salt sulfate (nss-sulfate) rises, enhancing activation potential  
443 by decreasing  $D_{\text{cri}}$ .

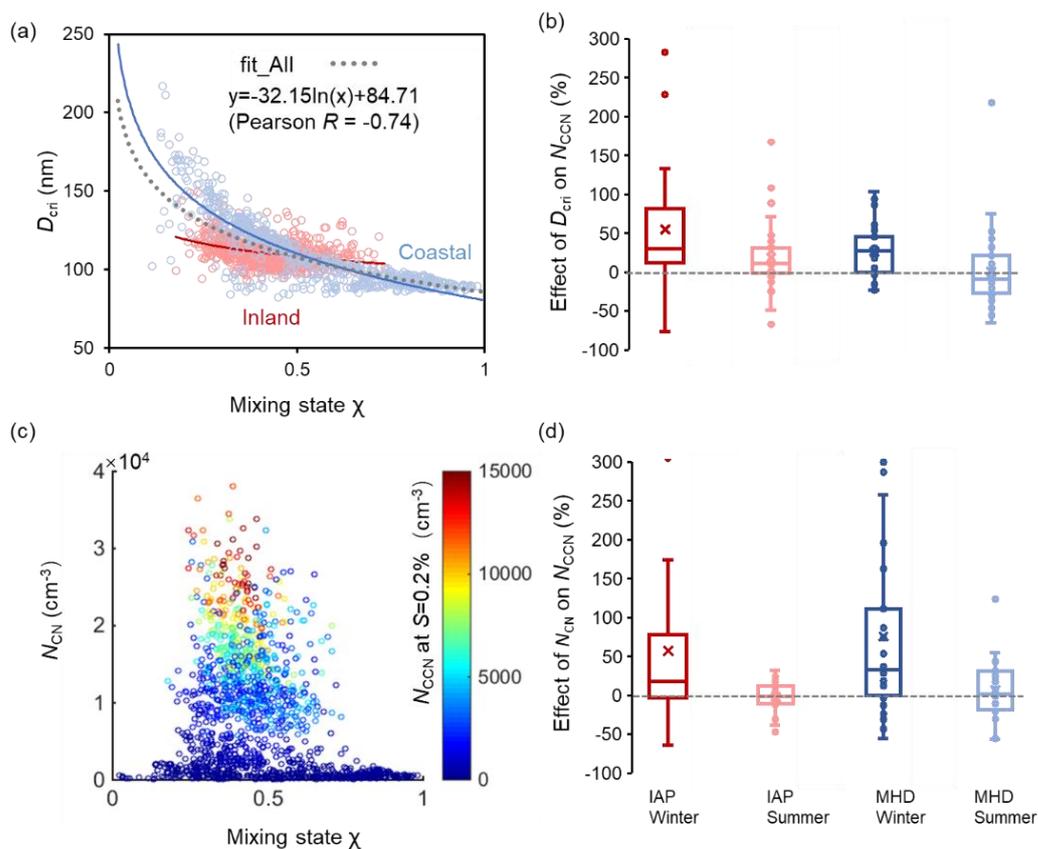
444 The dependence of CCN activity at 0.2% supersaturation on mixing state index  $\chi$   
445 reveals distinct inter-atmospheric differences, as shown in Fig. 7d-e. In the inland

446 atmosphere,  $N_{CCN}$  at  $S=0.2\%$  demonstrates a monotonic increasing trend with  $\chi$ ,  
447 attributed to the synergistic effects of rising  $N_{CN}$  and decreasing  $D_{cri}$  (Fig. S9). By  
448 contrast, coastal  $N_{CCN}$  follows a pattern analogous to  $N_{CN}$ , with peak concentrations  
449 shifting toward higher  $\chi$  values. This highlights the dominant role of particle size effects  
450 in enhancing CCN concentrations under marine-influenced conditions (Dusek et al.,  
451 2006).

452 Two distinct  $D_{cri}$ - $\chi$  trends underpin these disparities: one remains stable, driven by  
453 the inherent hygroscopicity of sea salt, while the other exhibits steep  $D_{cri}$  declines  
454 associated with anthropogenic pollution as internal mixing intensifies. These  
455 discrepancies are further manifested in the nonlinear  $D_{cri}$ - $\chi$  relationship. The activation  
456 ratio (AR)—quantifying aerosol cloud droplet formation potential at fixed  
457 supersaturation—also varies by site (Fig. 7e). Notably, AR shows a marked increase  
458 with  $\chi$  in IAP-winter, likely due to enhanced  $N_{CCN}$  from the elevated inorganic fraction  
459 under higher mixing states (Fig. 3). Conversely, the inorganic fraction decreases during  
460 other sampling periods, dampening AR growth.

461 To better interpret the impact of mixing state on CCN concentrations, Fig. 8  
462 quantifies the relative change in  $N_{CCN}$  at  $S=0.2\%$  as mixing state index  $\chi$  increases,  
463 contextualizing how CN concentration and chemical compositions (i.e.,  $D_{cri}$ ) evolve  
464 with mixing and aging across particle populations.  $D_{cri}$  demonstrates heightened  
465 sensitivity to minor  $\chi$  fluctuations at low mixing states ( $\chi < 0.5$ ; Fig. 8a), whereas further  
466 increases in internal mixing (higher  $\chi$ ) exert negligible influence on  $D_{cri}$  for already  
467 internally mixed particles. This behavior suggests that the  $D_{cri}$ - $\chi$  relationship may

468 enable a novel parameterization for  $D_{\text{cri}}$  estimation, a framework that is not yet reported  
 469 in prior literature.



470  
 471 **Fig 8.** Dependency of the critical diameter ( $D_{\text{cri}}$ ) on the  $\chi$  (a), relative change of CCN  
 472 number concentration ( $N_{\text{CCN}}$ ) at supersaturation  $S = 0.2\%$  with the reduction in  $D_{\text{cri}}$  (b);  
 473 Dependency of the CN number concentration ( $N_{\text{CN}}$ ) on the  $\chi$ , different colors represent  
 474 the  $N_{\text{CCN}}$  (c), relative change of  $N_{\text{CCN}}$  with the change in  $N_{\text{CN}}$  (d).

475 Aerosol data points at MHD site (blue dots) span a broad  $D_{\text{cri}}$  range (80–220 nm)  
 476 with  $\chi$  varying from 0.1 to 1, reflecting alternating influences of highly hygroscopic  
 477 inorganic salts (sea salt, sulfate) and less-hygroscopic organic matter. In contrast,  
 478 aerosols—dominated by anthropogenic pollutants—exhibit a narrower  $D_{\text{cri}}$  range (90–  
 479 150 nm) at IAP site. Both environments show negative  $D_{\text{cri}}-\chi$  correlations, but with

480 distinct functional forms: MHD aerosols feature an exceptional logarithmic fit ( $D_{\text{cri}} = -$   
481  $42.98\ln(\chi) + 80.36$ ,  $R^2 = 0.75$ ; Fig. 8a blue line), while IAP aerosols (red line) yield a  
482 shallower slope (-12.04). Pooling all data, we derive a generalized parameterization:  
483  $D_{\text{cri}} = -32.15\ln(\chi) + 84.71$  (Pearson  $r = -0.74$ ,  $R^2 = 0.54$ ). As already discussed above,  
484 strong impact of primary emission and secondary formation on aerosol mixing state  
485 was observed in both sites (Fig. 5 and 6). It also provides even more details on the  $D_{\text{cri}}$ -  
486  $\chi$  correlations. For example, the  $D_{\text{cri}}$  exhibited rapidly increased with the primary  
487 emissions (ie., mass fraction of POA enhanced) during polluted periods. The  $D_{\text{cri}}$  pattern  
488 appeared opposite with that of the mixing state index, especially for the accumulation-  
489 mode particles. More pronounced  $D_{\text{cri}}$ - $\chi$  correlations were observed during the new  
490 particle formation (Fig. 5a1-d1). The decreasing presence of  $D_{\text{cri}}$  matched the increasing  
491 proportion of  $\text{SO}_4^{2-}$  and SOA with the  $\chi$  increased during NPF events. Similar  
492 correlations between the critical diameter and mixing state index were also found in the  
493 coastal atmosphere, especially for the case of the enhanced anthropogenic organic  
494 matter and sea salt production (Fig.6). This implies that the relationship between the  
495  $D_{\text{cri}}$  and  $\chi$  might be disturbed by the variation of emission pollution and secondary  
496 formation processes, resulting in spatiotemporal differences.

497 Box plot analyses (Fig. S10) show that the mixing state reduces  $D_{\text{cri}}$  by 2.2–6.8%  
498 across campaigns, with the steepest winter decline. Changes in  $N_{\text{CN}}$  with differ starkly  
499 between environments: positive effects in polluted inland air (+9%) versus negative  
500 effects in coastal regions (-2%). Aerosols in IAP, frequently perturbed by primary  
501 emissions and new particle formation, exhibit elevated  $N_{\text{CN}}$  (peaking at  $\chi = 0.2$ –0.7),

502 while in MHD  $N_{\text{CN}}$  remains  $\sim 5000 \text{ cm}^{-3}$  across all  $\chi$ .

503 To isolate the impacts of critical diameter ( $D_{\text{cri}}$ ) and condensation nuclei number  
504 concentration ( $N_{\text{CN}}$ ) on CCN concentration, we categorized data into two groups: C1  
505 (particles within specific  $N_{\text{CN}}$  ranges) evaluates  $N_{\text{CCN}}$  variations driven by  $D_{\text{cri}}-\chi$   
506 relationships, while C2 (particles within fixed  $D_{\text{cri}}$  intervals) assesses  $N_{\text{CN}}-\chi$  effects (Fig.  
507 8b). Relative changes (RC) in  $D_{\text{cri}}$ ,  $N_{\text{CN}}$ , and  $N_{\text{CCN}}$  with  $\chi$  were calculated by comparing  
508 successive  $\chi$  increments ( $\chi_{i+1}$  vs.  $\chi_i$ ,  $i=0,0.1\dots 1$ ) within defined  $N_{\text{CN}}/D_{\text{cri}}$  windows.

509 Notably, change in  $N_{\text{CCN}}$  exerts more pronounced effects on externally mixed  
510 aerosols. For example, MHD-winter aerosols (high external mixing;  $\chi_{\text{mean}}=0.38\pm 0.12$ )  
511 showed  $N_{\text{CCN}}$  RCs of 23% (C1) and 72% (C2), whereas MHD-summer aerosols (high  
512 internal mixing;  $\chi_{\text{mean}}=0.69\pm 0.19$ ) exhibited negligible effects (-2.5% in C1, 0.9% in  
513 C2). Inland atmospheres, despite smaller seasonal  $\chi$  variations, showed analogous  
514 trends: winter  $N_{\text{CCN}}$  RCs (55% in C1, 57% in C2 for external mixing) exceeded summer  
515 values for more internally mixed populations (Fig. 8d). These results confirm that  
516 hygroscopic heterogeneity strongly influences  $N_{\text{CCN}}$  under external mixing, aligning  
517 with prior work (Ching et al., 2017).

518 With the variation in mixing state index  $\chi$ , changes in  $N_{\text{CCN}}$  are most pronounced  
519 during winter in both environments, attributed to heightened winter  $D_{\text{cri}}$  sensitivity to  $\chi$ :  
520 a 0.1  $\chi$  increase reduces  $D_{\text{cri}}$  by 5.2% (winter), boosting  $N_{\text{CCN}}$  by 39%, versus 2.4%  $D_{\text{cri}}$   
521 reduction (summer) yielding only 6%  $N_{\text{CCN}}$  enhancement. Concomitantly, winter  $N_{\text{CN}}-\chi$   
522 effects on  $N_{\text{CCN}}$  reach 65%, far exceeding summer responses.

523 Contrasting with prior evaluation methods that oversimplify mixing states (Ren et

524 al., 2018; Xu et al., 2021b), the entropy-based framework adopted herein enables  
525 explicit quantification of CCN activity evolution in response to mixing state transitions.  
526 Aerosols in IAP-winter are presumably shaped by intense urban pollution sources—  
527 including traffic emissions, residential heating, and cooking activities—thereby  
528 enriching the externally mixed particle fraction (Fan et al., 2020; Xie et al., 2020).  
529 Analogously, aerosols in MHD-winter exhibit dominant external mixing, consisting of  
530 near-hydrophobic and hydrophilic particle mixtures (Xu et al., 2021a). As illustrated in  
531 Fig. S1, winter aerosol populations display bimodal or multimodal  $\kappa$ -PDF distributions,  
532 evidencing high-degree external mixing with chemically diverse compositions. These  
533 results collectively highlight the pivotal role of mixing state heterogeneity in  
534 modulating CCN activity across environments.

#### 535 **4. Conclusions**

536 The mixing state of aerosol populations undergoes complex transformations  
537 during atmospheric aging, altering the distribution of hygroscopic and non-hygroscopic  
538 components and thus influencing CCN activity (Xu et al., 2021a; Ching et al., 2017).  
539 This study derived a mixing state index ( $\chi$ ) from field-measured hygroscopicity  
540 distributions, systematically investigating its impacts on CCN activity at two inland and  
541 coastal environments. Results provide field evidence that aerosol mixing states  
542 generally reside between purely internal and external extremes (Chen et al., 2022b).  
543 Aerosol mixing state is largely influenced by the primary emissions and secondary  
544 formation process. Externally-mixed particles with more hydrophobic-mode originate  
545 chiefly from primary emissions in IAP, while that of more sea-salt mode from sea spray

546 in MHD. While it becomes more internally-mixed as the enhanced fraction of more-  
547 hygroscopic mode and decreased of hydrophobic mode during the aging process. This  
548 highlights a dual regulatory mechanism of mixing state and its potential impact on  
549 hygroscopic distribution and CCN activity.

550 As  $\chi$  increases, CN number concentrations ( $N_{\text{CN}}$ ) first rise—driven by primary  
551 emissions and new particle formation—then decline due to condensation and  
552 coagulation during aging. Additionally, a logarithmic decreasing relationship between  
553 critical diameter ( $D_{\text{cri}}$ ) and  $\chi$  was identified for both inland and coastal particles,  
554 parameterized as  $D_{\text{cri}} = -32.15\ln(\chi) + 84.71$  (Pearson  $R = -0.74$ ,  $R^2 = 0.54$ ). This offers  
555 a practical approach to estimate  $D_{\text{cri}}$  from  $\chi$ , serving as a general framework for  
556 integrating mixing state effects on CCN activity in atmospheric models.

557 Entropy-based analyses confirm the pivotal role of mixing state in regulating  $N_{\text{CCN}}$ ,  
558 especially for externally mixed aerosols: a 0.1  $\chi$  increase can enhance  $N_{\text{CCN}}$  by 39–65%.  
559 Current models often oversimplify aerosol mixing states as purely internal or external  
560 (Stevens et al., 2019; Bauer et al., 2013), the latter being particularly sensitive to organic  
561 matter (Ren et al., 2018; Bhattu et al., 2015). Such simplifications introduce significant  
562 biases in  $N_{\text{CCN}}$  estimation (Riemer et al., 2019; Ching et al., 2019). The  $\chi$ - $D_{\text{cri}}$   
563 parameterization proposed here offers a novel approach to reduce model complexity in  
564 representing aerosol hygroscopicity and CCN activation, enabling more accurate  
565 simulations of aerosol CCN capacity. It is expected mitigate the underestimation in  
566 CCN compared with the complete external mixing assumption, while effectively  
567 alleviates the overestimation that arises from applying the complete internal mixing

568 assumption in regions characterized by high external mixing (Zheng et al., 2021a). This  
569 advancement improves our understanding of aerosol-cloud interactions (IPCC, 2021;  
570 Rosenfeld et al., 2019), critical for refining climate effect assessments.

#### 571 **Data availability**

572 All data used in the study are available at <https://doi.org/10.3974/geodb.2019.06.11.V1>  
573 (Fan et al., 2019) and <http://doi.org/10.17632/3dx6pnx869.1> (Xu et al., 2021a).

#### 574 **Author contributions**

575 RH and JR conceived the conceptual development of the paper. JR, FZ and WX directed  
576 and performed the experiments with YW and LC. FZ and YS provided the dataset in  
577 the inland site. JO, DC and CO provided the dataset in the coastal site. JR conducted  
578 the data analysis and wrote the draft. All authors edited and commented on the various  
579 sections of the paper.

#### 580 **Competing interests**

581 The contact author has declared that none of the authors has any competing interests.

#### 582 **Supporting Information**

583 Additional analysis results that were applied in this study. Mean values of the  $\kappa$ -  
584 PDF for aerosols of five diameters (Figure S1), sensitivity of the hygroscopic parameter  
585 for the group of the hygroscopic species on the mixing state index  $\chi$  (Figure S2), time  
586 series of the average per-particle species diversity  $Da$ , the bulk population species  
587 diversity  $D\gamma$ , and their affine ratio  $\chi$  (Figure S3), diurnal trend of particle size, chemical

588 mass fraction and number fraction (NF) of hydrophobic and hygroscopic mode in IAP  
589 (Figure S4) and in MHD (Figure S5), mixing state as a function of number fraction of  
590 hydrophobic and hygroscopic mode (Figure S6), variation of the peak diameter ( $D_{\text{peak}}$ )  
591 with the mixing state index (Figure S7), particle number size distribution and mixing  
592 state during new particle formation events (Figure S8), diurnal variation of  $\chi$  and CN  
593 concentration during winter and summer periods for 40 nm and 150 nm aerosols in  
594 inland and for 35 nm and 165 nm aerosols in coastal site (Figure S9), relative change  
595 of the critical diameter and CN concentration with the mixing state index  $\chi$  (Figure S10)  
596 (PDF).

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