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Revisited heat budget and probability distributions of turbulent heat fluxes in the Mediterranean Sea

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Abstract

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19 Understanding the surface heat budget of the Mediterranean Sea is essential for assessing its role in regional climate
20 and ocean circulation. Under the steady-state heat budget closure hypothesis, the Mediterranean should exhibit a
21 net surface heat loss to balance the heat gained through the inflow of warm Atlantic water at the Gibraltar Strait.
22 However, literature estimates of the net heat flux vary widely, raising questions about the accuracy of existing
23 reanalysis products. In this study, we compute the net surface heat flux over the Mediterranean using two
24 atmospheric datasets: high-resolution (0.125°) ECMWF analysis and lower-resolution (0.25°) ERA5 reanalysis. By
25 applying the same sea surface temperature fields and bulk formulas in both cases, we isolate the impact of
26 atmospheric resolution and data quality. We find that the ECMWF analysis yields a basin-averaged net heat flux
27 of $-3.6 \pm 1.3 \text{ W m}^{-2}$, consistent with the closure hypothesis, while ERA5 gives a spurious positive flux of
28 $+5 \pm 1.2 \text{ W m}^{-2}$. Furthermore, beyond simply assessing the net heat budget, this study delves into the probability
29 distributions of air-sea heat fluxes, aiming to gain a deeper understanding of associated uncertainties and extreme
30 values in turbulent heat fluxes. The probability distributions for turbulent heat flux components exhibit
31 characteristics such as skewness and kurtosis, respectively, varying across the basin. To assess the influence of
32 extremes, we apply the Interquartile Range (IQR) method within statistical models that account for the skewed
33 nature of turbulent heat flux distributions, enabling a consistent treatment of outliers. Our results reveal that extreme
34 negative heat flux events play a critical role in determining the net heat flux direction; excluding these extremes
35 leads to a spurious positive heat budget. Only ECMWF fields are consistent with the heat budget closure hypothesis.
36 Furthermore, we demonstrate that the Mediterranean heat budget closure hypothesis is connected to extreme heat
37 loss events occurring in key regions of the basin, such as the Gulf of Lion, the Adriatic Sea, the Aegean Sea, and
38 the southern Turkish coasts.

39 KEYWORDS: Air-Sea Heat Fluxes, Heat Budget, Fluxes Probability Distributions, Extreme Heat Fluxes,
40 Mediterranean Sea

41

42 **1. Introduction**
43

44 The exchange of momentum, water, and heat between the atmosphere and ocean plays a pivotal role in connecting
45 their dynamics (Kara et al., 2000). These fluxes, influenced by atmospheric surface variables and Sea Surface
46 Temperature (SST), drive ocean circulation (Large and Yeager, 2009; Small et al., 2019).

47 Our study focuses on the Mediterranean Sea, a unique semi-enclosed anti-estuarine basin where heat, water, and
48 momentum fluxes intertwine to fuel a robust vertical circulation (Pinardi et al., 2019). The Mediterranean heat
49 budget comprises of two main terms- the basin averaged surface and lateral boundary heat fluxes. Large
50 uncertainties are associated with surface heat fluxes at different temporal scales (Jordà et al., 2017). We aim to
51 reassess the long term mean net heat flux of the basin since this flux is a source of energy for the basin wide
52 circulation (Cessi et al., 2014).

53 Understanding the heat budget in the Mediterranean Sea has long been a formidable task (Bignami et al., 1995;
54 Castellari et al., 1998; Matsoukas et al., 2005; Pettenuzzo et al., 2010; Sanchez-Gomez et al., 2011; Criado-
55 Aldeanueva et al., 2012; Jordà et al., 2017), whether through numerical models or observational data analysis. The
56 fundamental challenge of in-situ observations is their limited spatial and temporal coverage, while numerical
57 modelling is constrained by the semi-empirical nature of air-sea bulk formulas. Numerous endeavours have been
58 undertaken (Large and Yeager, 2009) to calculate air-sea heat fluxes using atmospheric state variables obtained from
59 in-situ observations, remote sensing data, or numerical model outputs. In this study, we utilize atmospheric analysis
60 and reanalysis data, which provide an optimal reconstruction of past atmospheric surface state variables using
61 models and observations. Furthermore, the estimate of the Mediterranean Sea heat budget from ECMWF
62 meteorological analysis data sets has not been done before.

63 Numerous past studies have employed well-established bulk transfer formulas to estimate air-sea fluxes (e.g., Fairall
64 et al., 2003; Pettenuzzo et al., 2010; Cronin et al., 2019). The turbulent heat flux components (latent and sensible
65 heat flux) are commonly derived from surface wind speed, sea surface temperature, near-surface air temperature,
66 and humidity (Large and Yeager, 2009). Gulev and Belyaev (2012) noted that global heat flux products often vary
67 significantly, mainly due to differences in the bulk formulations and input variables adopted across studies.

68 At the Gibraltar Strait, the Mediterranean Sea exchanges water with the Atlantic through a characteristic two-layer
69 flow: warm, relatively fresh Atlantic water enters at the surface, while colder, saltier Mediterranean water exits at
70 depth leading to a net heat gain for the Mediterranean basin. To maintain a steady state balance in the basin averaged
71 heat balance, this lateral heat gain must be compensated by a net loss of heat at the sea surface. In other words, the
72 basin-average surface heat flux should be negative-a constraint known as the heat budget closure hypothesis.
73 Accurately estimating this surface heat flux remains a challenge due to limited data and uncertainties in flux
74 parameterizations. A benchmark estimates of the net heat budget, -7 W m^{-2} , was proposed by Béthoux et al. (1998),
75 though it is based on data from the 1970s and 1980s and may not reflect present-day conditions under a changing
76 climate (Criado-Aldeanueva et al., 2012; Marullo et al., 2021). We realise that assuming perfect balance between
77 lateral and vertical heat fluxes, even in the Mediterranean Sea, is an approximation. Being heat clearly entering the
78 Mediterranean Sea through Gibraltar, we search for a negative net heat flux, which we call the closure hypothesis.

79 How negative such net heat flux is, we do not know but searching for a negative value is a conservative assumption
80 aligned with current scientific understanding.

81 Recent studies highlight significant uncertainty in the estimated long-term net heat budget of the Mediterranean Sea,
82 with some even reporting positive values. Song and Yu (2017), presented an ensemble climatology of surface heat
83 fluxes, reporting a net heat budget of $2 \pm 12 \text{ W m}^{-2}$ and noting a warm bias in this ensemble estimate. Utilizing an
84 ensemble of high-resolution regional climate models (RCMs), Sanchez-Gomez et al. (2011) found that individual
85 RCMs did not achieve a heat budget closure, but the ensemble mean heat flux was $-7 \pm 21 \text{ W m}^{-2}$. Using downscaled
86 NCEP/NCAR global reanalysis of $\frac{1}{2}^\circ \times \frac{1}{2}^\circ$ resolution, Ruiz et al. (2008) computed a net heat budget of -1 W m^{-2} .
87 However, their heat flux components values are not close to most of the literature values (for instance, the major
88 difference was in the value for net short wave with 84 W m^{-2}). Marullo et al. (2021) recently analysed several
89 atmospheric data sets, revealing a significant net heat flux variability ranging between 1.6 and 40 W m^{-2} . They
90 attributed this variability primarily to longwave radiation fluxes uncertainties. In addition to these challenges, past
91 studies of air-sea fluxes have primarily focused on establishing mean and variance, leaving limited knowledge about
92 their statistical distributions (Korolev et al., 2015; Tian et al., 2017). Understanding the probability distributions of
93 air-sea fluxes and their higher moments could provide insights into the uncertainties associated with air-sea physics.
94 Also, the analysis of probability distributions can help to assess skills of different reanalyses to replicate extreme
95 fluxes (Gulev and Belayaev, 2012).

96 In this study we investigate two different aspects of the surface heat budget closure hypothesis. First, we employ
97 two different high quality surface atmospheric variable data sets at different horizontal resolution and we calculate
98 the heat fluxes with the same bulk formula and the same SST. This isolates the impact of atmospheric model
99 resolution and quality as the source of variation in the heat flux estimates. Therefore, we answer the question: is
100 the Mediterranean Sea in the past 15 years still losing heat at the surface?

101 Secondly, we study the statistical distributions of the heat flux components, utilizing the atmospheric analysis dataset
102 from ECMWF (European Centre for Medium-Range Weather Forecasts). Knowing the skewness and kurtosis
103 distributions across the basin, we analyse the extremes of the net heat budget, and we determine the specific
104 importance of the extreme heat losses to the long-term mean. The second question we address is: what is the cause
105 of the Mediterranean Sea negative long-term mean heat budget?

106 The paper is structured into the following sections. Section 2 presents the atmospheric analysis and reanalysis
107 datasets from ECMWF, along with satellite SST data and the bulk formula used in the estimation of the fluxes. In
108 Section 3, we present the new values of the heat budget closure problem, compared to the literature. In section 4,
109 we analyse the probability distributions of turbulent heat fluxes. In Section 5, we determine the causes of the long
110 term mean net heat budget values. Finally, Section 6 summarizes the findings and highlights key insights gleaned
111 from this research.

112

113

114

115 2. Methodology and datasets

116 2.1 Air-sea physics in the Mediterranean Sea

117 For the Mediterranean Sea, several formulations have been established over the past decades through extensive
118 studies. In this section, we present these adopted formulations, beginning with the net heat flux formula, followed
119 by the specific heat flux components utilized in this study.

120 The net surface heat flux, Q_{net} comprises the net shortwave radiation Q_{SW} , net longwave radiation Q_{LW} , and surface
121 turbulent flux components, which encompass the latent heat flux of evaporation Q_{LH} and sensible heat flux Q_{SH}
122 (Cronin et al., 2019; Pettenuzzo et al., 2010).

$$123$$
$$124 \quad Q_{net} = Q_{SW} + Q_{LW} + Q_{lat} + Q_{sen} \quad (1)$$
$$125$$

126 Here, we use the convention that positive heat fluxes denote heat gain by the ocean. We did not use directly the
127 atmospheric model heat flux values since we wanted to intercompare two different atmospheric data sets in terms
128 of their quality and resolution not on the basis of the specific bulk parametrizations and SST used. Thus, we used
129 the same bulk formula and SST for both ECMWF and ERA5 surface variables that are described in section 2.2.

130

131 2.1.1 Shortwave radiation flux

132 The shortwave radiation flux (SW) is derived from the formulation proposed by Rosati and Miyakoda (1988). The
133 largest heat flux component is the solar radiation which is reduced by the cloud coverage and partially reflected by
134 the sea surface (albedo). The shortwave heat flux formula is therefore expressed as:

$$135$$
$$136 \quad Q_{SW} = Q_{TOT} (1 - 0.62 C + 0.0019 \beta)(1 - \alpha) \quad \text{if } C \geq 0.3 \quad (2)$$

$$137 \quad Q_{SW} = Q_{TOT} (1 - \alpha) \quad \text{if } C < 0.3$$

138 where Q_{TOT} indicates the clear sky solar radiation calculated with astronomical formulae, β is the noon solar altitude
139 in degrees and α is the ocean surface albedo varying month wise values taken from Payne (1972). For the cloud
140 cover, we follow the Reed (1977) formula, where the threshold cloud fraction 0.3 is in tenths, indicating 30% cloud
141 coverage. The incoming solar radiation varies on locations with sun zenith angel and Q_{TOT} reaches at the ocean
142 surface after diffusion can be represented by the components: the sum of the direct solar radiation Q_{DIR} for direct
143 solar radiation and Q_{DIF} for downward diffused radiation. Then net solar radiation Q_{TOT} can be represented by the
144 summation of components Q_{DIR} and Q_{DIF} :

145

$$146 \quad Q_{TOT} = Q_{DIR} + Q_{DIF}$$
$$147 \quad = Q_0 \tau^{\sec z} + [(1 - A_a) Q_0 - Q_0 \tau^{\sec z}] * 0.5$$

148 Here Q_0 is the solar radiation at the top of atmosphere, τ is equal to 0.7 and is the atmospheric transmission
 149 coefficient, A_a is a constant value (0.09) and z is the sun zenith angle.

150

151 2.1.2 Longwave radiation flux

152 The longwave surface radiation flux is the difference between the upward infrared radiation (IR) emitted from the
 153 ocean surface (LU) and the atmospheric downwelling infrared radiation (LD). The LD component is adapted from
 154 Bignami et al. (1995), and the longwave radiation flux is written as:

155

$$156 \quad Q_{LW} = Q_{LU} + Q_{LD} \quad (3)$$

$$157 \quad Q_{LU} = -\epsilon \sigma_{SB} T_S^4 \quad (4)$$

$$158 \quad Q_{LD} = [\sigma_{SB} T_A^4 (0.653 + 0.00535 e_A)](1 + 0.1762 C^2) \quad (5)$$

159

160 where: T_S and T_A indicate the sea surface temperature and air temperature in degrees Kelvin, σ_{SB} is the Stefan-
 161 Boltzmann constant, ϵ is the ocean emissivity set to 1 according to Large and Yager (2009) and e_A is the atmospheric
 162 vapor pressure computed from the mixing ratio of the air W_{air} (Wallace and Hobbs, 2006).

$$163 \quad W_{air} = \frac{q_A}{1 - q_A} \quad (6)$$

$$164 \quad e_A = \frac{W_{air}}{(W_{air} + \gamma)} p \quad (7)$$

165 and q_A is the specific humidity of air, p is the surface air pressure, and γ is a constant (0.622).

166 The specific humidity (q_A) saturated at the T_A is computed using the following equation (Large, 2006), where $\rho =$
 167 1.22 kg m^{-3} is the air density:

$$168 \quad q_A = \rho^{-1} 640,380 \exp(-5107.4/T_D) \quad (8)$$

169 where, T_D is the dew point temperature retrieved from the atmospheric model outputs.

170

171 2.1.3 Turbulent heat fluxes

172 The turbulent heat flux is composed of sensible heat Q_{SH} and latent heat Q_{LH} given by the following formula:

$$173 \quad Q_{SH} = -\rho_A C_P C_H |\vec{V}| (T_S - T_A) \quad (09)$$

$$174 \quad Q_{LH} = -\rho_A L_E C_E |\vec{V}| (q_S - q_A) \quad (10)$$

175

176 where $|\vec{V}|$ is the wind speed, ρ_A is the density of moist air, C_p is the specific heat capacity ($1005 \text{ J g}^{-1}\cdot\text{K}$), C_H and
177 C_E are turbulent exchange coefficients for temperature and humidity, L_E is the latent heat of vaporization, q_A is
178 defined in (8) and q_S , which is the specific humidity of air saturated at the sea surface temperature T_S , is calculated
179 with (8) using T_S instead of T_D , and applying a 0.98 factor (Sverdrup, 1942). Since the average wind speed in the
180 Mediterranean is 5 m/s, Pettenuzzo et al. (2010) suggested using constant turbulent exchange coefficients such as
181 $C_H = 1.3 \cdot 10^{-3}$ and $C_E = 1.5 \cdot 10^{-3}$.

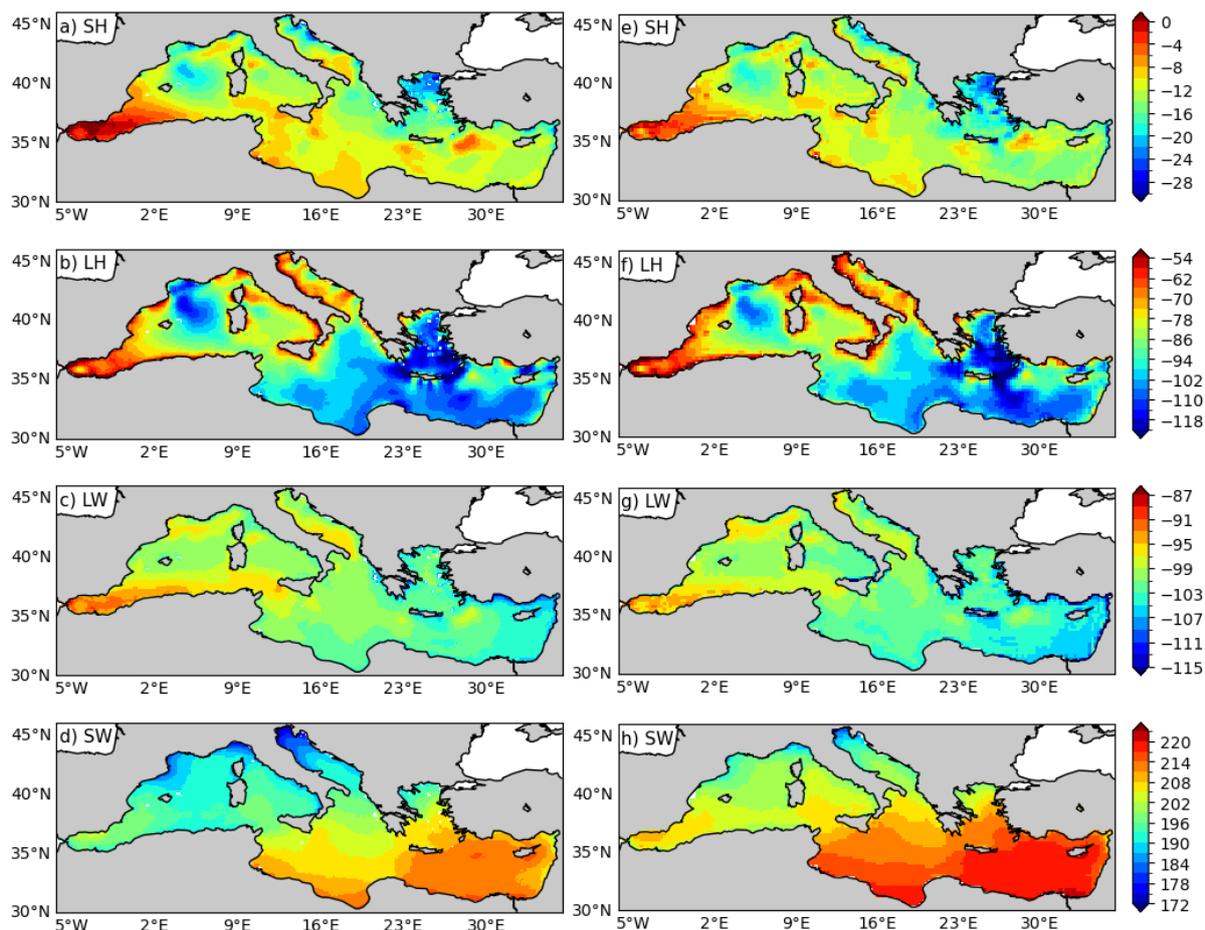
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183 **2.2 Datasets**

184 Two atmospheric datasets have been selected for this study. The first dataset is the ECMWF (European Centre for
185 Medium-Range Weather Forecasts) high-resolution analysis dataset (Rabier et al., 2000) at six-hourly temporal
186 resolution and 0.125 degrees of spatial resolution. It's worth noting that the original operational dataset, from which
187 the atmospheric fields have been extracted, underwent changes between 1991 and 2006 in terms of model resolution
188 and the assimilated number of observations. For consistency, we opted to utilize the dataset with approximately
189 uniform model resolution and physics spanning a 15-year period from 2006 to 2020. The second dataset employed
190 is ERA5 reanalysis (Hersbach et al., 2020). This dataset is available at hourly intervals. However, it features a
191 horizontal resolution of 0.25 degrees.

192 To mitigate unresolved atmospheric temperature daily cycles in ECMWF and make the two data sets consistent for
193 the time variability, the ECMWF and ERA5 fields are further processed into daily mean values for the entire period.
194 Comparisons conducted with daily and six-hourly input fields indicated minimal differences in the probability
195 distributions of the heat fluxes, leading us to prioritize filtering out daily variability to the greatest extent possible.

196 To compute the heat fluxes the following atmospheric surface variables are extracted from the two datasets: the 10-
197 meter wind components (U for the zonal direction and V for the meridional direction), mean sea level pressure, dew
198 point temperature, total cloud coverage, and 2-meter air temperature.



199
 200 **Figure 1: Annual means of heat flux components for the period of 2006-2020, computed from ECMWF (left**
 201 **panel) and ERA5 (right panel) daily time series. The corresponding ECMWF time series is shown in**
 202 **supplementary material, Fig. S1 and S2.**

203
 204 For the oceanic SST data, we utilized the satellite dataset distributed by the Copernicus Marine Environment Service
 205 (CMEMS). This SST dataset is a blended product from multiple satellite sensors, categorized as L4, with a
 206 horizontal resolution of $0.05^\circ \times 0.05^\circ$. To align the SST data with the atmospheric analysis and reanalysis dataset
 207 grids, we applied an interpolation and extrapolation method known as the 'sea-over land' (De Dominicis et al., 2013).
 208 This method involves an iterative process to extrapolate sea values over land before interpolating, thus not allowing
 209 the contamination of land values on the interpolation.

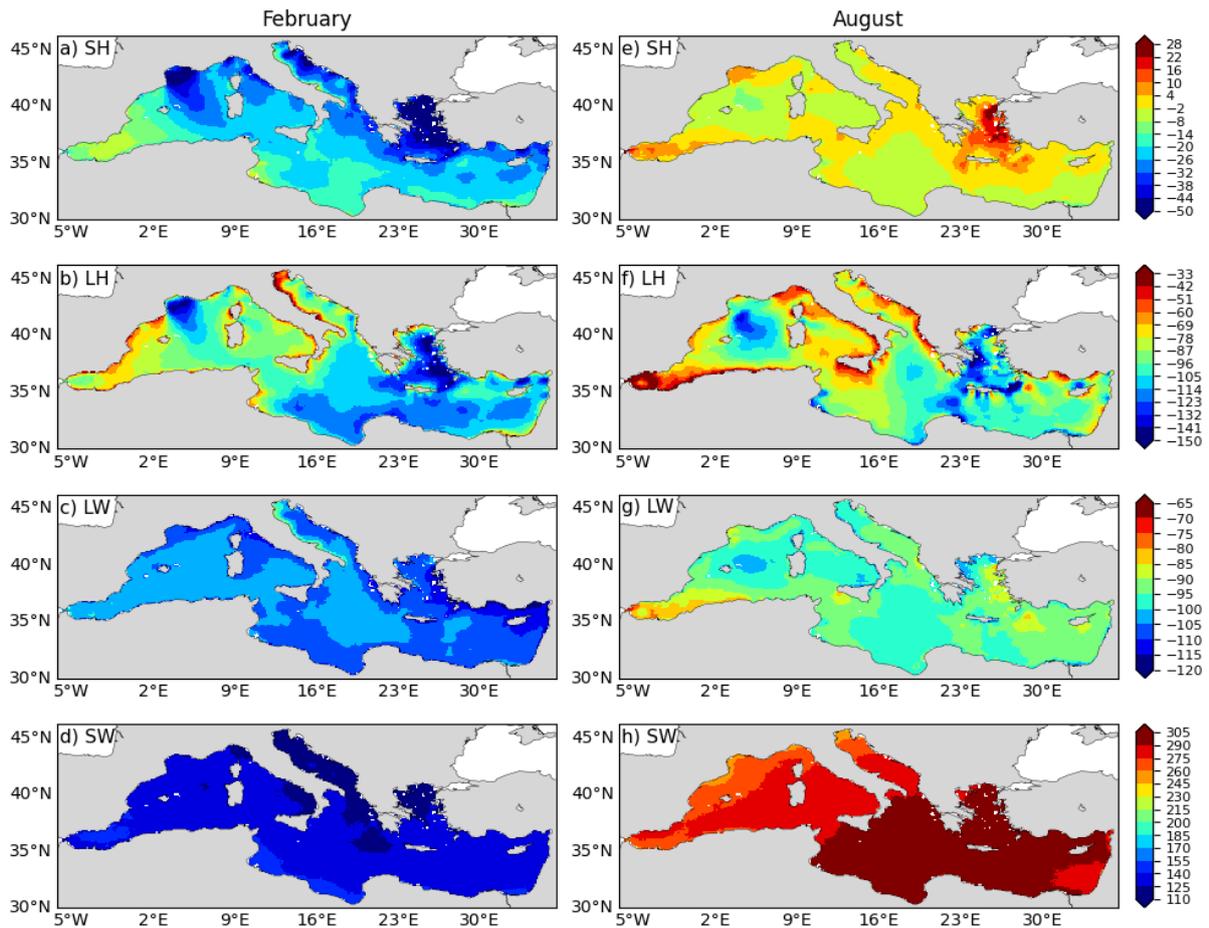
210
 211 **3. Heat budget closure problem revisited**

212 **3.1 Analysis of the heat budget components**
 213

214 We compute the heat fluxes for the 15-year period, 2006–2020, using the ERA5 dataset and compare them with the
 215 fluxes computed using the ECMWF dataset (Fig. 1). In Fig. 1 we show the results for the 15-year mean of each
 216 heat budget components. We start describing the ECMWF patterns and then we detail the differences.

217 Turbulent heat fluxes exhibit distinct sub-basin-scale patterns, varying between the eastern and western
218 Mediterranean Seas as well as the Central Mediterranean region. The smallest mean sensible heat loss is observed
219 in the whole Alboran sea area with absolute value range of 0-6 W m^{-2} , while the Aegean Sea and the centre of Gulf
220 of Lion loses more heat in the maximum value of 25 W m^{-2} . Similarly, the highest absolute values of LH are
221 recorded in the Gulf of Lion and the Aegean and Levantine Seas, attributed to the influence of strong and cold winds
222 like the Mistral and Etesian in the north-western and eastern Mediterranean regions, respectively. The eastern
223 Mediterranean emerges as the region with the highest evaporation, reaching approximately 122 W m^{-2} in absolute
224 value. Notably, along the south-eastern coastline, a wide range of maximum absolute values (102-122 W m^{-2}) in
225 the evaporation is observed. The turbulent heat fluxes show limited differences between the ECMWF and ERA5
226 datasets.

227 SW fields show the well-known meridional gradients with larger gradient values arising from the ECMWF dataset.
228 The mean SW exhibits a gradual decline from the eastern to western Mediterranean, influenced by the variation of
229 the solar zenith angle with longitudes. The first reason for SW differences between Western Mediterranean and
230 Eastern Mediterranean is the latitudinal position of each sub-basin. Furthermore, SW differences using ECMWF
231 and ERA5 datasets are connected to different cloud cover schemes (not shown), leading to a larger heat gain in the
232 Eastern Mediterranean. Notably, the northern Adriatic region stands out with a distinct distribution, suggesting it
233 receives relatively less annual solar radiation compared to other areas. In contrast, the mean longwave (LW)
234 radiation distribution maintains a relatively consistent range of absolute values between 87–113 W m^{-2} across the
235 entire domain with absolute minimum values in the Alboran Sea, presumably due to the warm Atlantic surface water
236 inflow. Overall, while the turbulent heat fluxes show limited differences between the ECMWF and ERA5 datasets,
237 significant discrepancies are observed in radiative heat fluxes. Additionally, Fig. 1 shows the noisiness of the fluxes
238 due to the ERA5 low resolution with respect to ECMWF while retaining an overall consistency.



239
 240 **Figure 2: Seasonal variations of heat flux components: Left column is the monthly average values for**
 241 **February and right column shows the average for August for the period 2006-2020 (ECMWF data).**

242
 243 Figure 2 shows seasonal variations in heat flux components for February and August using ECMWF data. Both SH
 244 and LH fluxes exhibit a greater spatial gradient in February compared to August. In winter, the SH loss is larger,
 245 especially in the Gulf of Lion, Aegean, and parts of the Adriatic, with stronger spatial gradients compared to summer.
 246 In August, SH flux becomes positive in the Aegean and the Alboran Sea. LH loss is highest in February in the whole
 247 eastern Mediterranean and the Gulf of Lion. In August, LH losses decrease in the western Mediterranean, with
 248 absolute value minima in the Alboran and Adriatic Sea, remaining largely negative in the lower part of in the Eastern
 249 Mediterranean. SW fields show the strongest seasonal cycle as expected, with the absolute maximum of 260-305
 250 $W m^{-2}$ in summer and in the Eastern Mediterranean. LW is largest in absolute value in winter showing a small
 251 seasonal cycle. Significant seasonal variations are observed in the distribution range for radiative heat fluxes, low
 252 in February and high in August across the entire domain. These patterns are quite similar to the ones reported in the
 253 literature.

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 255
 256

257 3.2 Net heat budget estimation

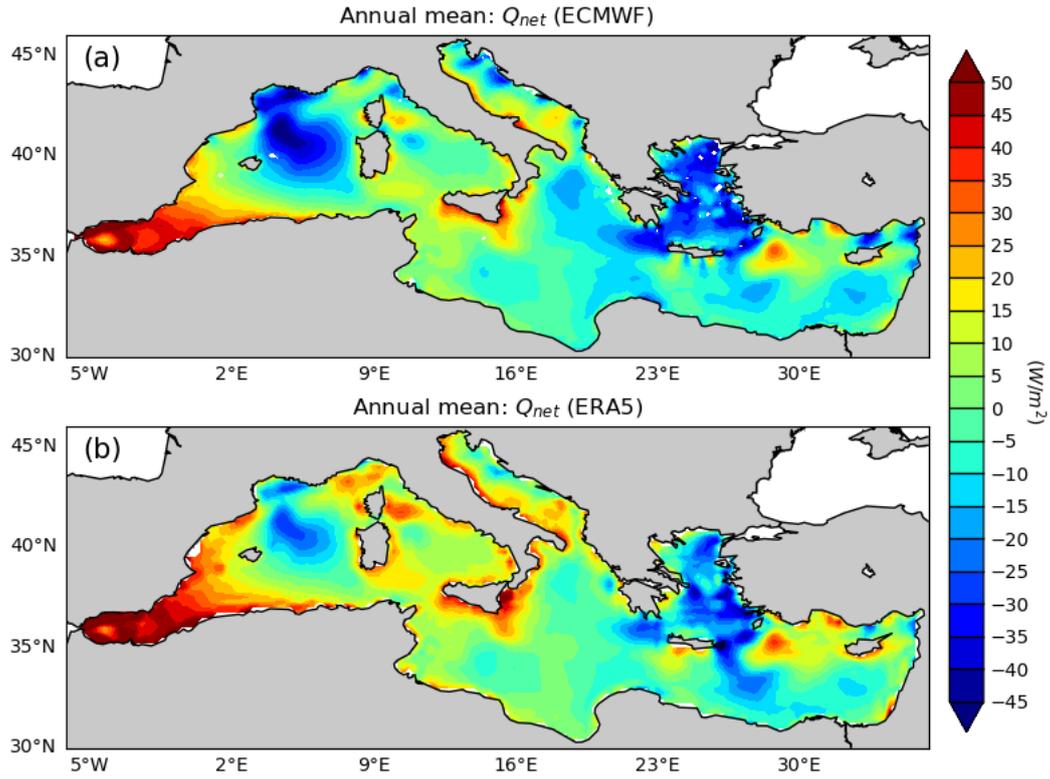
258
259 The net surface heat flux Q_{net} is depicted in Figure 3 for ERA5 and ECMWF and basin-average 15 year mean
260 values are listed together with the literature in Table 1.

261 Fig. 3 shows that the Gulf of Lion and the Aegean Sea are the areas of maximum heat losses while the basin gains
262 heat in the Alboran Sea, in some areas of the Levantine basin, and in the shelf areas around the Italian peninsula.

263 The Mediterranean Sea gains comparatively more heat with the ERA5 inputs. Besides the difference in surface
264 domain for Q_{net} , for both cases, air-sea flux dynamics is strongly visible in the Alboran Sea for net heat gain, in
265 the Gulf of Lion for heat loss due to the continental cold wind (Mistral wind), and in the Aegean Sea due to the
266 strong wind (Etesian wind) that blows during the summer period. Using ECMWF inputs, Q_{net} is -3.6 W m^{-2} , a value
267 consistent with previous estimates for the Mediterranean Sea domain and for ERA5, it is 5 W m^{-2} (Table 1). Errors
268 in Q_{net} mean value are determined by a bootstrapping method where Q_{net} time series is resampled 5000 times to
269 compute a standard deviation around the mean of the resampled time series (Tibshirani & Efron, 1993). We argue
270 that our results show that the negative heat budget is achieved by using only ECMWF fields at high resolution, i.e.
271 0.125 degrees. Higher resolution implies differences in all atmospheric fields used to compute the fluxes.
272 Furthermore, ERA5 and ECMWF model physics and dynamics is different contributing to the differences in the
273 mean heat budget. However, both datasets use observations, and we argue that the most relevant difference between
274 the analysis and the reanalysis data set is the resolution due to the peculiar geometry of the Mediterranean Sea.

275 Since all the datasets used in the literature are coarser, this is most likely the reason of the failure to determine the
276 correct heat budget closure value. In Pettenuzzo et al. (2010) several ad-hoc corrections were made to the surface
277 atmospheric fields to obtain the negative heat flux budget while in Sanchez-Gomes et al., (2011) they used an
278 ensemble of deterministically downscaled ERA40 fluxes giving rise to a very large uncertainty. Considering a
279 recent literature, our resulted Q_{net} is closely matched with the computed net heat budget of $-3 \pm 8 \text{ W m}^{-2}$ from Jordà
280 et al., (2017), but their result was associated to large temporal uncertainties from the surface fluxes through Gibraltar
281 Strait.

282



283

284 **Figure 3: Comparison of the annual Q_{net} ($W m^{-2}$) mean, computed from a) ECMWF and b) ERA-5 input**
 285 **datasets.**

286 **Table 1: Computed heat flux components and net heat fluxes (Q_{net}), and values from the references**

Authors	SH	LH	LW	SW	Net Flux (Q_{net})
Bethoux (1979)	-13	-120	-68	195	-6
Bunker (1982-1)	-13	-101	-68	202	20
Bunker et al (1982-2)	-11	-130	-68	202	20
May (1986)	-11	-130	-68	193	2
Garret et al. (1993)	-7	-99	-67	202	29
Matsoukas et al. (2005)	-11	-122	186	-63	22
Ruiz et al. (2008)	-8	-88	-73	168	-1
Pettenuzzo et al. (2010)	-14	-90	-79	178	-7
Sanchez-Gomez et al. (2011)	-13±5	-100±13	-75±6	181±18	-7±21
Criado-Aldeanueva et al. (2012)	-15.1	-93.5	-76.9	186.3	0.73
Song & Yoy (2017)	-13±4	-98±10	-78±13	192±19	2±12
Jordá, et al., 2017	-	-	-	-	-3±8
ECMWF analyses	-12.1±4	-92±16	-100.5±3	201±8	-3.6±1.3
ERA5 reanalysis	-13±3	-89±14	-101±3	208±8	5±1.2

287 Spatially, the mean Q_{net} distribution generally shows a heat loss across much of the Eastern Mediterranean. Overall,
 288 distributions of more positive net heat budget values for the western Mediterranean and negative for the eastern
 289 Mediterranean have matched with the similar result from Criado-Aldeanueva et al. (2012). Strong spatial gradients
 290 are evident, particularly in the Aegean Sea, although a few patches displaying net heat loss (negative Q_{net}) are also
 291 noticeable in this vicinity. Conversely, the western Mediterranean exhibits a stronger heat gain area, which appears
 292 particularly concentrated zone in the Gulf of Lion region and this feature is apparent in results from both atmospheric
 293 datasets. Such a spatial related uncertainty in Q_{net} represents a significant challenge for accurately closing regional
 294 heat budgets as well as validating existing ocean circulation models within the complex Mediterranean basin.

295

296 **4. Probability distributions of the turbulent heat fluxes**

297 In this section, we analyse the probability distribution of turbulent heat fluxes computed using ECMWF data set
 298 only and for the anomaly heat fluxes. Recent studies by Gulev and Belyaev (2012) and Korolev et al. (2015) have
 299 analysed the statistical distributions of turbulent heat fluxes, and their findings are used here for comparison.
 300 Radiative flux components are excluded from this analysis, as they do not exhibit extremes of comparable magnitude
 301 to those of turbulent fluxes (Supplementary material, Fig S3). This suggests low skewness and kurtosis in their
 302 distributions, reducing the relevance of a detailed probability density function analysis for these components.

303

304 If we indicate the time series of each component of the heat budget with X_n we can define the heat flux climatology
 305 as:

306

$$307 \quad Q_t = \frac{1}{n} \sum_{j=1}^n X_{tj} \quad (11)$$

308

309 where ‘t’ indicates the day of the year, and ‘j’ is the number of years. The anomaly time series is computed by
 310 subtracting the long-term seasonal climatology Q_t from the observed heat flux time series X_{tj} and it will be indicated
 311 by:

312

$$313 \quad \tilde{X}_{tj} = X_{tj} - Q_t \quad (12)$$

314

315 **4.1 SH flux distribution**

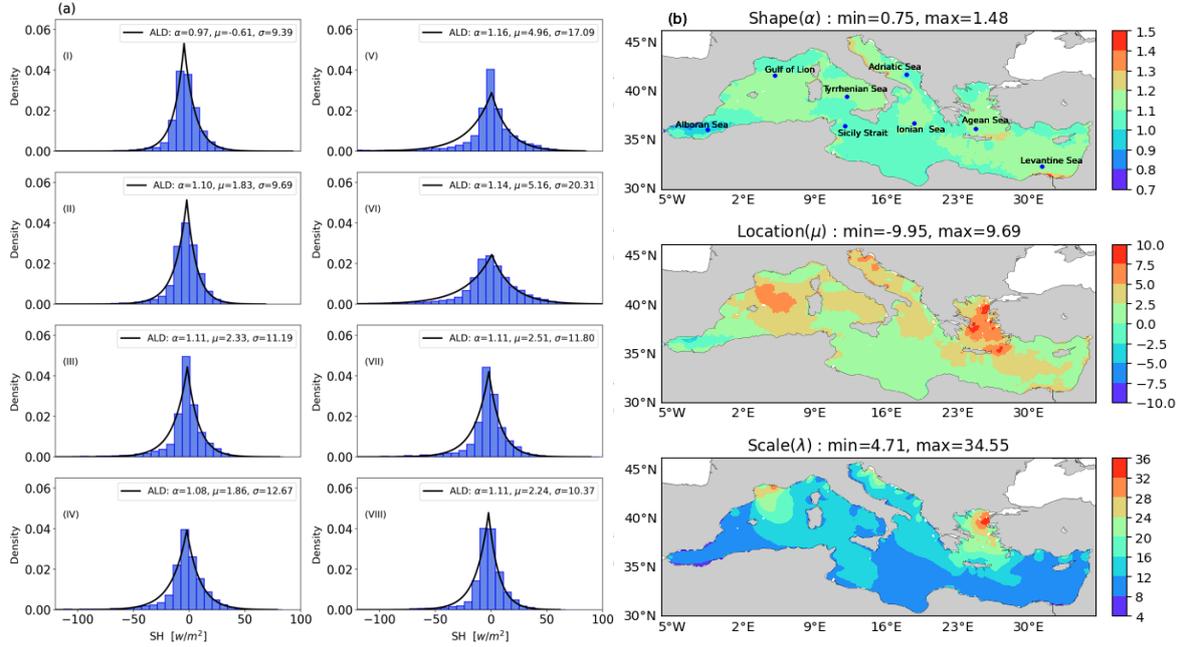
316 We found that gaussian or skew-normal distributions are not well fit for SH flux, as evident from the histograms at
 317 single grid points shown in Figure 4a. The histograms reveal a singularity around zero, indicating that the skew-
 318 normal distribution may not adequately capture the distribution of these values. This observation is consistent with
 319 findings by Gulev and Belyaev (2012), and we provide further explanation in the Appendix A.

320 The most common distribution with such near-discontinuous behaviour at the origin is the three-parameter
 321 Asymmetric Laplace Distribution (ALD) (Yu & Zhang, 2005) that we can defined as

322
$$F(x, \alpha, \mu, \lambda) = \frac{\lambda}{\alpha + \frac{1}{\alpha}} \begin{cases} \exp\left(\frac{\lambda}{\alpha}(x - \mu)\right) & \text{if } x < \mu \\ \exp(-\lambda \alpha(x - \mu)) & \text{if } x \geq \mu \end{cases} \quad (13)$$

323 where x is the random variable time series, α is the shape parameter, μ is the location and λ the scale.

324



325 **Figure 4: a) The single grid point histograms for SH flux anomalies from the eight sampling locations for the**
 326 **period of 2006–2020, b) The Asymmetric Laplace PDF parameter (α, μ, λ) distributions from computed SH**
 327 **flux anomaly for the observation period. [Sampling points: (I) Alboran Sea, (II) Gulf of Lion, (III) Tyrrhenian**
 328 **Sea, (IV) Sicily Strait, (V) Adriatic Sea, (VI) Ionian Sea, (VII) Aegean Sea, (VIII) Levantine Sea]**
 329

330
 331 From the single grid point histogram, we have observed a one or two sharp peaks in the distribution that matches
 332 well with the Asymmetric Laplace Distribution (ALD) PDF. In accordance with findings by Yu and Zhang (2005),
 333 the distribution of the sensible heat (SH) flux anomaly time series exhibits characteristics of a double exponential
 334 distribution. This is evident from the histograms displaying both positive and negative skewness with long tails, as
 335 depicted in Figure 4a. The ALD parameters for the SH flux anomaly time series are illustrated in Figure 4b. The
 336 shape parameter (α) falls within the positive range of 0.73 to 1.48, indicating a moderate to high degree of peakiness
 337 in the distribution. Additionally, the location parameter (μ) exhibits mostly positive values while a small area in the
 338 Alboran Sea shows negative values, suggesting a shift in the central tendency of the distribution. Notably, the scale
 339 parameter (λ) displays a similar structure to the SH flux climatology depicted in Figure 1.

340 To check the quality of the fit, moments of both applied and theoretical PDF are compared (presented in
 341 supplementary materials, Fig. S4). The comparison shows the estimations of the three moments in the left panel for
 342 the observed SH flux and right panel for ALD PDF parameters. It can be seen that variances and skewness are
 343 similar in distribution while kurtosis differ at noticeable range. This observation is likely attributed to the fact that
 344 the kurtosis for the asymmetric Laplace distribution remains constant regardless of changes in the scale parameter.

345

346 **4.2 LH flux distribution**

347 In the case of the LH flux, no sharp exponential peaks were observed; instead, large skewness and long tails were
348 identified. Therefore, we applied the skew-normal PDF which is defined by α ($\in \mathbb{R}$) as the shape parameter, μ (\in
349 \mathbb{R}) the location parameter, and $\lambda > 0$ the scale parameter (Azzalini, 1985) and defined as:

350
$$f(x, \alpha, \mu, \lambda) = \frac{2}{\lambda} \phi\left(\frac{x-\mu}{\lambda}\right) \Phi\left(\alpha \frac{x-\mu}{\lambda}\right) \quad (14)$$

351

352 Where,

353
$$\phi\left(\frac{x-\mu}{\lambda}\right) = \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2} \frac{(x-\mu)^2}{\lambda^2}} \quad (15)$$

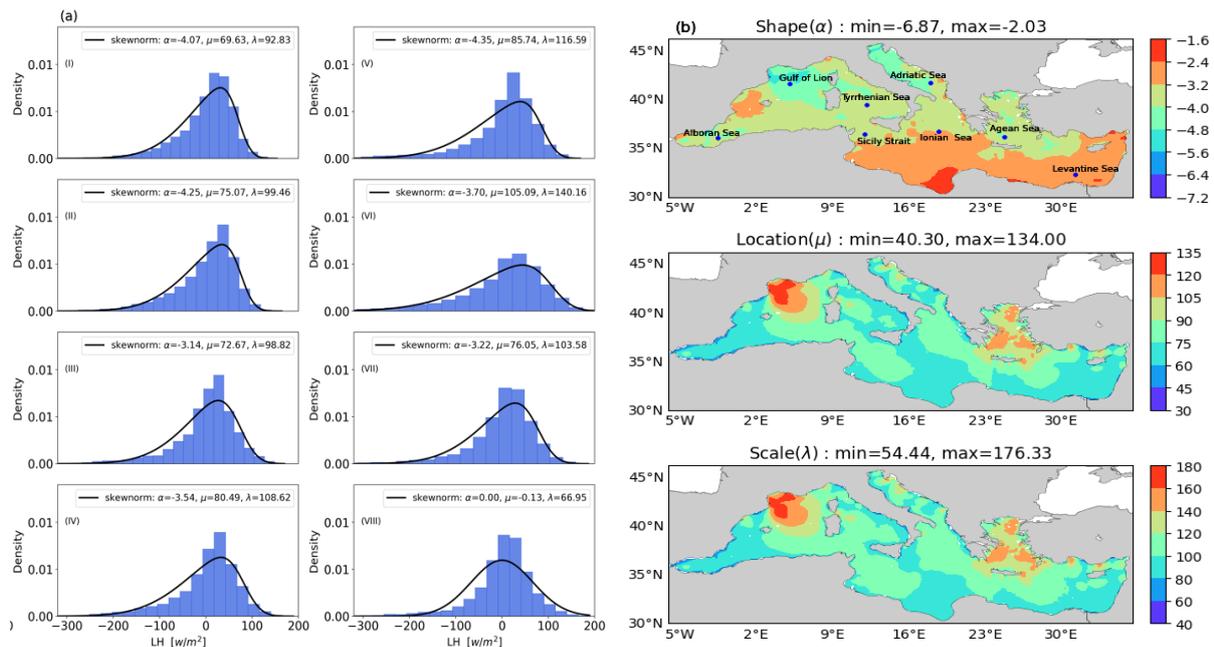
354
$$\Phi\left(\alpha \frac{x-\mu}{\lambda}\right) = \int_{-\infty}^{\alpha \frac{x-\mu}{\lambda}} \phi(t) dt \quad (16)$$

355

356 A skew-normal PDF is an extension of the normal distribution while covering the skewness and containing the
357 general characteristics of a Gaussian distribution (Flecher et al., 2010).

358 To examine visually the quality of PDF fit on LH flux anomaly values, histograms from eight sea locations were
359 fitted with the skew-normal PDF, as shown in Figure 5a. Figure 5b displays the parameter spatial variability. The
360 shape parameter distribution ranges from -6.83 to -2.5, with negative values observed across all points. This spatial
361 distribution of α , exhibiting a negative range, aligns with the negatively skewed pattern identified in the single grid
362 point PDF fitting test. Furthermore, the spatial distribution structure of the location and scale parameters
363 demonstrates a positive correlation across most locations.

364 In the Supplementary material, a comparison of statistical moments is conducted to qualitatively validate the fit
365 (supplementary materials, Fig S5). There is notable agreement in the variance distributions between the observed
366 LH flux anomaly and skew-normal PDF. While the skewness distributions mismatch at negligible level, with the
367 theoretical PDF skewness predominantly ranging from -0.9 to -0.3, whereas the observed skewness exhibits a
368 variation range spanning from -1.2 to over -0.3. Lastly, the kurtosis distribution of the skew-normal PDF differs in
369 the Aegean Sea, Alboran Sea and Gulf of Lion area.



371 **Figure 5: a) The single grid point histograms for LH flux anomalies at the eight sampling locations for the**
 372 **period of 2006-2020, 5 b) The skew-normal PDF parameter (α, μ, λ) distributions for computed LH Flux**
 373 **anomaly for the observation period [Sampling points: (I) Alboran Sea, (II) Gulf of Lion, (III) Tyrrhenian**
 374 **Sea, (IV) Sicily Strait, (V) Adriatic Sea, (VI) Ionian Sea, (VII) Aegean Sea, (VIII) Levantine Sea]**

375

376 4.3 Evaluation of the PDF fitting

377 In this section, we conducted a goodness of fit test to measure the distance between the empirical distributions and
 378 the fitted ones. The objective of this evaluation test was to assess the degree of agreement between the applied
 379 theoretical distribution and the observed time series. The chi-squared method, a well-accepted test, was employed
 380 to measure the distance between two independent distributions.

381 We compared the results of the chi-squared test for the turbulent heat fluxes computed using the ECMWF and ERA5
 382 datasets. The decision rule for the χ^2 test was determined based on the level of significance, set at 0.05, and the
 383 degrees of freedom, defined as $DF = N - np$, where N represents the number of empirical histogram bins and np is
 384 the number of distribution parameters (i.e., 3 for both the ALD and skew-normal distributions). In the
 385 supplementary material we show the maps of Chi-square test statistics (Supplementary material, Fig. S6). The chi-
 386 squared results for the SH and LH fluxes computed using the ECMWF dataset indicate that almost all surface grid
 387 points are well-fitted with the applied theoretical PDFs. With the critical threshold of 33.92 (Elderton, 1902) for P
 388 values, we observed a very few mismatches, mainly located near the coasts.

389

390 5. How do heat loss extremes contribute to the heat budget closure hypothesis?

391 The heat budget closure problem is associated with achieving a net negative heat flux, as discussed before. We test
 392 here the hypothesis that the negative long term mean negative heat budget of Table 1 for ECMWF data is correlated
 393 to the extremes in heat losses during autumn-winter.

394 Figure 6 illustrates the Q_{net} basin average daily time series, revealing a value range varying between 200 and -500
 395 $W m^{-2}$. Notably, the most pronounced extreme negative heat losses, reaching up to $-500 W m^{-2}$ occur in the winters
 396 of 2011, 2015 and 2017. They approximately coincide with western Mediterranean Deep Water formation events,
 397 as documented in Escoudier et al. (2021). To identify and remove the potential extremes in our computed Q_{net} time
 398 series, we apply the Interquartile Range (IQR) method which measures the spread of a dataset and calculate the
 399 difference between the third quartile(Q3) and the first quartile (Q1). The IQR threshold is computed by the
 400 difference between the 1st quartile (Q1) and 3rd Quartile (Q3) of the observed dataset:

401

$$402 \quad IQR = Q3 - Q1 \quad (17)$$

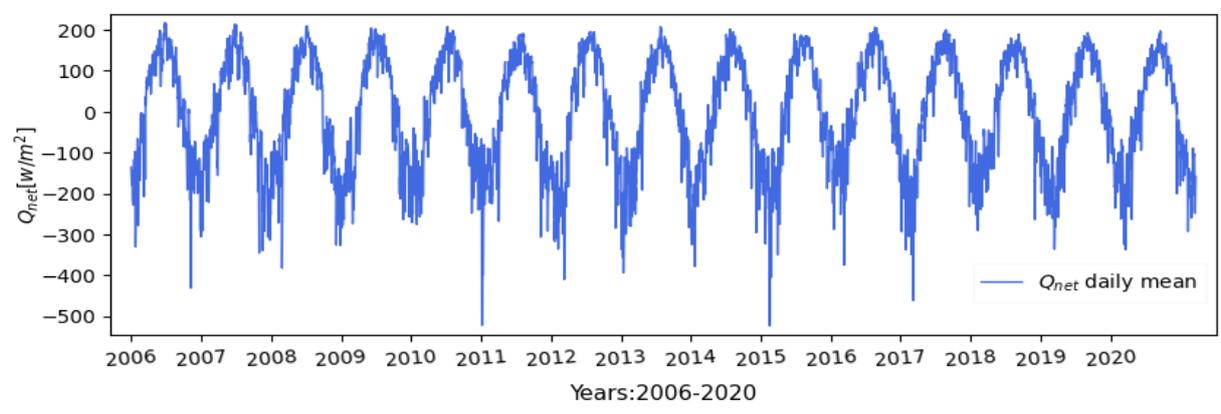
$$403 \quad \text{Threshold} = Q1 - k * IQR \quad (18)$$

404

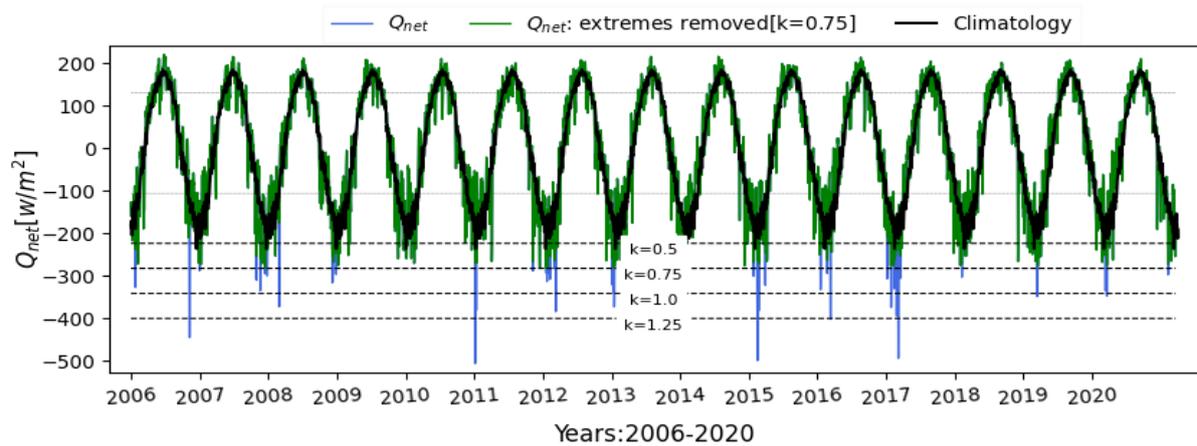
405 We used different values for k to exclude the negative extremes, which correspond to the largest heat losses. These
 406 extreme values were replaced with long-term daily climatological values (computed using eq. 11 and the long term
 407 mean neat heat budget Q_{net} is recomputed.

408 The resulting Q_{net} for different thresholds is displayed in Table 2 and the thresholds are shown in Fig. 7 together
 409 with the daily climatology. If compared with the long term mean heat budget in Table 1 ($-3.6 \pm 1.3 W m^{-2}$), we see
 410 that eliminating the winter extremes produces a smaller long term mean heat loss up to changing the sign to positive
 411 values. We argue that the ECMWF net negative heat extremes determine the ECMWF negative long term mean
 412 heat budget. Furthermore, if we calculate the yearly mean value of the seasonal climatology, we obtain the value
 413 of $+4 W m^{-2}$, which confirms again the importance of extremes in the heat budget closure of the Mediterranean
 414 Sea.

415 The Q_{net} could become an impact indicator of the Mediterranean for sea level trends in the basin. The net heat
 416 budget in fact relates to the sea level tendency (Pinardi et al., 2014) in the Mediterranean Sea and could be considered
 417 as a key indicator of climate impacts in the Mediterranean Sea.



418
 419 **Figure 6: Basin averaged time series of the computed daily Q_{net} (units $W m^{-2}$) from the ECMWF computed**
 420 **heat fluxes, for the period 2006-2020.**



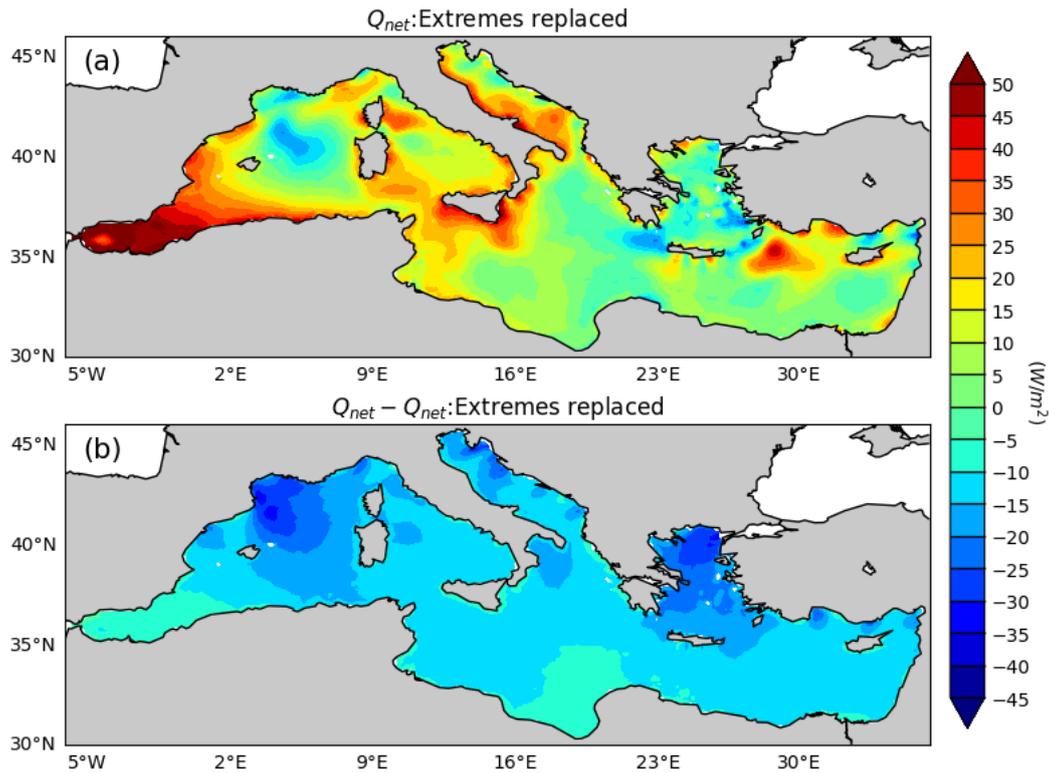
422
 423 **Figure 7: Time series of the basin averaged Q_{net} mean, Q_{net} extremes removed mean, and long-term yearly**
 424 **climatology and four lower quantile boundary line marked with dashed lines using different k values [k=1.25,**
 425 **1,0.75, 0.5].**

426
 427 **Table 2: Different lower quantile boundary limits used to replace potential extremes and the resulting long-**
 428 **term mean basin averaged Q_{net} values**

IQR lower boundary limit	Threshold values ($W m^{-2}$)	$Q_{net}(Wm^{-2})$
K=1.25	-405	-3.2
K=1.0	-347	-2.5
K=0.75	-289	-1
K=0.5	-231	2

429
 430 Figure 8 presents the new long-term mean spatial distribution of the surface heat budget after removing negative
 431 extreme values using a threshold of $-289 W m^{-2}$ ($K = 0.75$). The figure illustrates that these extreme events exert a
 432 substantial influence on the overall structure of the net heat budget in the Mediterranean Sea, with particularly
 433 pronounced effects in the Gulf of Lion, the Aegean Sea, the eastern Adriatic Sea, and along the southern Turkish
 434 shelves.

435



436
 437 **Figure 8: a) Long term mean net heat flux after the removal of extremes (Q_{net} :Extremes replaced) using**
 438 **threshold of -289 W m^{-2} ($K = 0.75$) b) The time mean distribution computed from the difference between**
 439 **Q_{net} (Fig. 3a) and “ Q_{net} : Extremes replaced” time series. It shows a significant reduction of heat losses in**
 440 **the Gulf of Lion, Adriatic Sea and Aegean Sea regions.**

441

442

443 6. Discussion and conclusions

444

445 The primary objective of this investigation is to revisit the heat budget closure hypothesis from atmospheric
 446 consolidated data sets that are nowadays used frequently to drive ocean models. Simultaneously, we intend to model
 447 the statistical distributions of turbulent heat fluxes and assess the contribution of extremes into the net heat budget
 448 closure of the Mediterranean Sea. For this analysis, we covered a 15-year period from 2006 to 2020 with a daily
 449 time series frequency. The reason for the choice of this time range is that ECMWF analyses became quite stable
 450 starting from 2006 while before the model was at coarse resolution, like ERA5’s model. Our strategy is to use the
 451 same SST and the same bulk formula but different atmospheric reanalysis and analysis surface variable data sets
 452 and compare the value of the long term mean heat budget in the Mediterranean Sea.

453 Firstly, the surface heat budget of the Mediterranean Sea was analysed to examine average annual mean and seasonal
 454 variations. The largest component of the heat budget is the net solar radiation (SW), followed by the latent heat
 455 (LH), longwave radiation (LW), and then sensible heat (SH), as shown in the literature (Table 1). All heat flux
 456 components exhibit significant seasonality, as illustrated in Figure 2. Differences appear in the structure of the
 457 fluxes, especially the SW and LW, when different atmospheric data sets are used, a conclusion aligning with a

458 suggestion from Marullo et al. (2021) on the sensitivity of LW estimates from the atmospheric dataset used to
459 calculate fluxes. We compared the ERA5-derived surface radiative fluxes with the computed ERA5 heat fluxes
460 presented in Figure 3b and found that the ERA5 longwave (LW) fluxes are substantially overestimated in absolute
461 magnitude (Supplementary Material S7). The associated uncertainty is comparable in order of magnitude to that
462 reported by Marullo et al. (2021), who analysed an observational dataset at a specific site. Nonetheless,
463 compensating biases between the SW and LW components (Supplementary Material S8) result in a net radiative
464 heat balance difference between ERA5-derived and computed heat fluxes that is large primarily in the southern
465 Mediterranean, where ERA5-derived exhibits reduced LW flux values.

466
467 The basin-average net heat flux, Q_{net} , was calculated to be $-3.6 \pm 1.3 \text{ W m}^{-2}$ for ECMWF analysis data while it is
468 $5 \pm 1.2 \text{ W m}^{-2}$ for ERA5 (Table 1). This finding supports the conclusion that heat budget closure hypothesis cannot
469 be satisfied with a relatively coarse reanalysis atmospheric data set. Our initial question was: is the Mediterranean
470 Sea in the past 15 years still losing heat at the surface? The answer is yes if we use a high-resolution ECMWF
471 atmospheric analysis. Additionally, comparing the Q_{net} estimates derived from ERA5 and ECMWF with the same
472 bulk formulas demonstrates that the uncertainty peaks in the atmospheric forcing resolution and possibly cloud
473 cover, the latter affecting the radiative components of the heat budget.

474
475 Furthermore, we have demonstrated that the probability density of surface heat fluxes can be modelled and fitted
476 with a three-parameter PDF composed of a shape, a location, and a scale parameter. All the turbulent heat flux
477 components show asymmetric behaviour. There is encouraging agreement between the first two statistical moments
478 of the fitted PDF and the observed values. Kurtosis does not seem to be properly captured by the PDF used but our
479 time series is too short to arrive at a definitive conclusion. For the SH we demonstrate that the Asymmetric Laplace
480 Distribution PDF is generated by the contributing distributions of wind speed (Weibull) and temperature difference
481 Skew Normal). We believe this is the first time that such kind of relationship is demonstrated.

482 Gulev and Belyaev (2012) applied the two-parameter Fisher–Tippett distribution (also known as the Gumbel
483 distribution) to monthly sensible and latent heat fluxes derived from NCEP–NCAR reanalysis fields. Their
484 approach focused on using the mean and standard deviation to estimate the distribution’s location and scale
485 parameters relevant to extreme events. However, the Gumbel distribution has a fixed skewness, limiting its ability
486 to capture the contribution of rare, asymmetric extremes. In contrast, our study analysis anomalies from the seasonal
487 cycle using full probability distributions that allow for variable skewness. This better reflects the nature of
488 atmospheric and oceanic variables, which are often inherently skewed (Sardeshmukh and Penland, 2015), and is
489 essential for understanding the influence of extremes on the surface heat budget. Our findings show that
490 incorporating a shape parameter is key to accurately capturing distribution structure and preserving asymmetric
491 tails. This analysis provides a useful framework for validating surface flux products and assessing their variability,
492 particularly important given that surface fluxes are the dominant source of uncertainty in the Mediterranean net heat
493 balance (Jordà et al., 2017). Correctly estimating skewness is crucial, as a small number of extreme outliers,
494 especially during intense winter events, can disproportionately affect the basin-wide mean and determine whether
495 heat budget closure is achieved.

496 For the first time, we have investigated the effects of extreme heat losses in the Mediterranean Sea in the long term
497 mean basin averaged heat budget. The northern basin areas are the site of the largest heat losses (Gulf of Lion and
498 the Aegean Sea, Adriatic Sea and the Turkish southern coasts). Exclusion of the negative extremes in these areas
499 resulted in a change in the sign of long term mean heat loss. The anomaly threshold value (-231 W m^{-2} , Table 2)
500 resulted a long-term positive net heat flux, which is inconsistent with the basin's heat flux closure hypothesis. Our
501 second initial question was: what is the cause of the Mediterranean Sea negative long-term mean heat budget? The
502 answer is that the long-term mean, basin averaged heat loss is due to winter extremes in the Northern regions of the
503 Mediterranean Sea.

504 In conclusion, understanding the characteristics and distributions of air-sea heat fluxes are crucial for gaining
505 insights into variations in the heat budget. Furthermore, the PDF analysis of turbulent heat fluxes will allow us to
506 have a better understanding of the extreme events and their contributions to the net negative heat budget. The next
507 steps could involve a machine learning study of air-sea flux bulk parametrizations for different atmospheric data
508 sets and coupled models, using as target the heat flux data set from this study.

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516 **Data availability**

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518 The ERA5 data available at Copernicus Climate Change Service, Climate Data Store, (2023): ERA5 hourly data
519 on single levels from 1940 to present. [https://cds.climate.copernicus.eu/datasets/reanalysis-era5-single-](https://cds.climate.copernicus.eu/datasets/reanalysis-era5-single-levels?tab=download)
520 [levels?tab=download](https://cds.climate.copernicus.eu/datasets/reanalysis-era5-single-levels?tab=download) . ECMWF analysis data can be arranged following the submission of a request.

521

522 **Author contributions**

523

524 MHG: development of the concept, literature review, writing, methodology, coding, formal analysis, wiring,
525 visualization. NP: conceptualization, review, writing, methodology. AN: conceptualization, writing, review. LM:
526 review, writing. SB: methodology, review. FM: methodology, review. FT: methodology, coding.

527

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531 MHG and NP.

532

533 **Conflict of interest Statement**

534 The authors declare no conflicting interests.

535

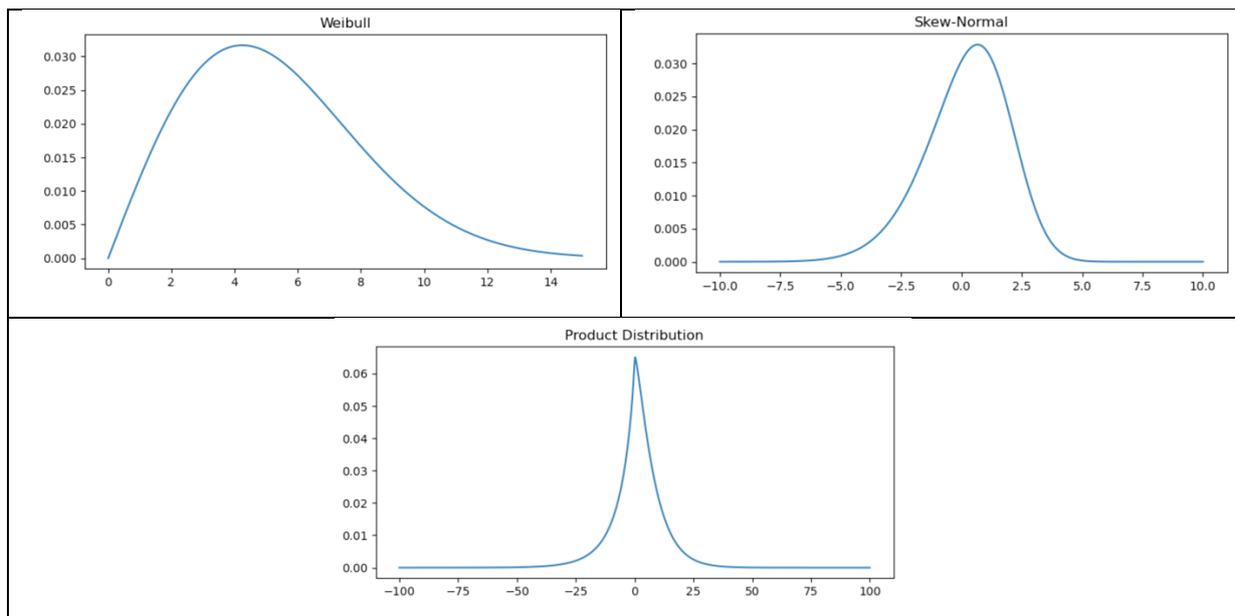
APPENDIX A

536

537 Here, we show that the characteristics of the SH flux distribution are due to the specific form of the heat flux as
538 given by (10), i.e. a multiplication of two distributions, wind speed, $Q(v)$, known to be a Weibull, and temperature
539 differences, $R(DT)$, taken to be a Skew Normal.

540 Let's indicate with $P(v*DT)$ the combined SH distribution of $Q(v)$ and $R(DT)$ the temperature difference as in
541 equation (09). Assuming that the two distributions are independent, the combined distribution is the product of Q
542 and R . If we now define the variable $z=v*DT$, the new combined distribution on the heat flux variable z is given
543 by the Mellin transform and convolution, described in Papoulis, A., & Pillai, S. U. (2002). The resulting P
544 distribution is an Asymmetric Laplace Distribution (ALD) like distribution which is similar to the one computed in
545 Fig.4.

546



547 Figure A1: Histograms presenting the two original distributions, $Q(v)$ (upper left quadrant, units wind speed) and
548 $R(DT)$ (upper right quadrant, units degrees C) and the combined distribution for SH flux in units of $W m^{-2}$. The
549 parameters used for the two original distributions are: $k = 2.0$ for the Weibull shape, $\lambda = 6.0$ for the Weibull
550 scale; $\alpha = - 2.0$ for the Skew Normal shape, $\mu = 2.0$ for the Skew Normal location and $\omega = 2.5$ for the
551 Skew Normal scale
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APPENDIX B

The statistical moments for the skew-normal PDF are given by:

$$E(x) = \mu + \lambda \delta \sqrt{\frac{2}{\pi}} \tag{B1}$$

$$\sigma^2 = \lambda^2 \left(1 - \frac{2\delta^2}{\pi} \right) \tag{B2}$$

$$\mu_3 = (4 - \pi) \frac{(\delta \sqrt{2/\pi})^3}{2(1 - 2\delta^2/\pi)^{3/2}} \tag{B3}$$

$$\mu_4 = 2(\pi - 3) \frac{\left(\delta \sqrt{\frac{2}{\pi}}\right)^4}{\left(1 - \frac{2\delta^2}{\pi}\right)^2} \tag{B4}$$

where $\delta = \frac{\alpha}{\sqrt{1 + \alpha^2}}$. Since the expected value, E, of the time series is zero, we deduce that:

$$\mu = -\lambda \delta \sqrt{\frac{2}{\pi}} \tag{B5}$$

In other words, location and shape parameters have opposite signs since the scale parameter, λ , is always positive.

SH flux anomaly distribution was analysed with the Asymmetric Laplace Distribution PDF, its statistical moments given by:

$$\text{mean} = \mu + \frac{1 - \alpha^2}{\lambda \alpha} \tag{B6}$$

$$\text{variance} = \frac{1 + \alpha^2}{\lambda^2 \alpha^2} \tag{B7}$$

$$\text{Skewness} = \frac{2(1 - \alpha^6)}{(\alpha^4 + 1)^{\frac{3}{2}}} \tag{B8}$$

$$\text{Kurtosis} = \frac{6(1 + \alpha^3)}{(1 + \alpha^4)^2} \tag{B9}$$

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