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Revisited heat budget and probability distributions of turbulent heat fluxes in the Mediterranean Sea

Mahmud Hasan Ghani¹, Nadia Pinardi^{1,2}, Antonio Navarra^{2,3}, Lorenzo Mentaschi^{1,2},
Silvia Bianconcini⁴, Francesco Maicu², and Francesco Trotta²

¹ Dept. of Physics and Astronomy, University of Bologna, Bologna, Italy
² CMCC Foundation - Euro-Mediterranean Center on Climate Change, Italy
³ Dept. of Biological, Geological and Environmental Sciences, University of Bologna, Bologna, Italy
⁴ Dept. of Statistics, University of Bologna, Bologna, Italy

Corresponding author: Mahmud Hasan Ghani, mahmudhasan.ghani2@unibo.it

Abstract

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Understanding the surface heat budget of the Mediterranean Sea is essential for assessing its role in regional climate and ocean circulation. Under the steady-state heat budget closure hypothesis, the Mediterranean should exhibit a net surface heat loss to balance the heat gained through the inflow of warm Atlantic water at the Gibraltar Strait. However, literature estimates of the net heat flux vary widely, raising questions about the accuracy of existing reanalysis products. In this study, we compute the net surface heat flux over the Mediterranean using two atmospheric datasets: high-resolution (0.125°) ECMWF analysis and lower-resolution (0.25°) ERA5 reanalysis. By applying the same sea surface temperature fields and bulk formulas in both cases, we isolate the impact of atmospheric resolution and data quality. We find that the ECMWF analysis yields a basin-averaged net heat flux of $-3.6 \pm 1.3 \text{ W m}^{-2}$, consistent with the closure hypothesis, while ERA5 gives a spurious positive flux of $+5 \pm 1.2 \text{ W m}^{-2}$. Furthermore, beyond simply assessing the net heat budget, this study delves into the probability distributions of air-sea heat fluxes, aiming to gain a deeper understanding of associated uncertainties and extreme values in turbulent heat fluxes. The probability distributions for turbulent heat flux components exhibit characteristics such as skewness and kurtosis, respectively, varying across the basin. To assess the influence of extremes, we apply the Interquartile Range (IQR) method within statistical models that account for the skewed nature of turbulent heat flux distributions, enabling a consistent treatment of outliers. Our results reveal that extreme negative heat flux events play a critical role in determining the net heat flux direction; excluding these extremes leads to a spurious positive heat budget. Only ECMWF fields are consistent with the heat budget closure hypothesis. Furthermore, we demonstrate that the Mediterranean heat budget closure hypothesis is connected to extreme heat loss events occurring in key regions of the basin, such as the Gulf of Lion, the Adriatic Sea, the Aegean Sea, and the southern Turkish coasts.

KEYWORDS: Heat fluxes, Mediterranean net heat budget, Fluxes probability distributions, Heat flux extremes

41 **1. Introduction**
42

43 The exchange of momentum, water, and heat between the atmosphere and ocean plays a pivotal role in connecting
44 their dynamics (Kara et al., 2000). These fluxes, influenced by atmospheric surface variables and Sea Surface
45 Temperature (SST), drive ocean circulation (Large and Yeager, 2009; Small et al., 2019).

46 Our study focuses on the Mediterranean Sea, a unique semi-enclosed anti-estuarine basin where heat, water, and
47 momentum fluxes intertwine to fuel a robust vertical circulation (Pinaridi et al., 2019). We aim to reassess the long
48 term mean net heat flux of the basin since this flux is a source of energy for the basin wide circulation (Cessi et al.,
49 2014). Moreover, the Mediterranean net heat budget comprises of several terms that show a considerable range of
50 uncertainties (Jorda' et al., 2017).

51 Understanding the heat balance in the Mediterranean Sea has long been a formidable task (Bignami et al., 1995;
52 Castellari et al., 1998; Matsoukas et al., 2005; Pettenuzzo et al., 2010; Sanchez-Gomez et al., 2011; Criado-
53 Aldeanueva et al., 2012; Jordà et al., 2017), whether through numerical models or observational data analysis. The
54 fundamental problem of in-situ observations is the limited space-time extension of the data sets, while for numerical
55 modelling, the limitations lie in the semi-empirical approach of the air-sea bulk formulas. Numerous endeavours
56 have been undertaken (Large and Yeager, 2009) to calculate air-sea heat fluxes using atmospheric state variables
57 obtained from in-situ observations, remote sensing data, or numerical model outputs. In our study, we utilize
58 atmospheric analysis and reanalysis data, which provide an optimal reconstruction of past atmospheric surface state
59 variables using models and observations. Furthermore, the estimate of the Mediterranean Sea heat budget from
60 ECMWF meteorological analysis data sets has not been done before.

61 Numerous past studies have employed well-established bulk transfer formulas to estimate air-sea fluxes (e.g., Fairall
62 et al., 2003; Pettenuzzo et al., 2010; Cronin et al., 2019). The turbulent heat flux components-latent and sensible
63 heat flux are commonly derived from surface wind speed, sea surface temperature, near-surface air temperature,
64 and humidity (Large and Yeager, 2009). However, Gulev and Belyaev (2012) noted that global heat flux products
65 often vary significantly, mainly due to differences in the bulk formulations and input variables adopted across
66 studies.

67 At the Gibraltar Strait, the Mediterranean Sea exchanges water with the Atlantic through a characteristic two-layer
68 flow: warm, relatively fresh Atlantic water enters at the surface, while colder, saltier Mediterranean water exits at
69 depth. This arrangement leads to a net gain of heat for the Mediterranean basin, since the incoming surface water
70 carries more thermal energy than the colder outflow. To maintain a long-term heat balance, this lateral heat gain
71 must be compensated by a net loss of heat at the sea surface. In other words, the basin-average surface heat flux
72 should be negative-a constraint known as the heat budget closure hypothesis. Accurately estimating this surface
73 heat flux remains a challenge due to limited data and uncertainties in flux parameterizations. A benchmark estimates
74 of the net heat budget, -7 W m^{-2} , was proposed by Béthoux et al. (1998), though it is based on data from the 1970s
75 and 1980s and may not reflect present-day conditions under a changing climate (Criado-Aldeanueva et al., 2012;
76 Marullo et al., 2021). We realise that assuming perfect balance between lateral and vertical heat fluxes, even in the
77 Mediterranean Sea, is an approximation. Being heat clearly entering the Mediterranean Sea through Gibraltar, we
78 search for a negative net heat flux, which we call the closure hypothesis. How negative such net heat flux is, we do

79 not know but searching for a negative value is a conservative assumption aligned with current scientific
80 understanding.

81 Recent studies highlight significant uncertainty in the estimated long-term net heat budget of the Mediterranean Sea,
82 with some even reporting positive values. Song and Yu (2017), presented an ensemble climatology of surface heat
83 fluxes, reporting a net heat budget of $2 \pm 12 \text{ W m}^{-2}$ and noting a warm bias in this ensemble estimate. Utilizing an
84 ensemble of high-resolution regional climate models (RCMs), Sanchez-Gomez et al. (2011) found that individual
85 RCMs did not achieve a heat budget closure, but the ensemble mean heat flux was $-7 \pm 21 \text{ W m}^{-2}$. Using downscaled
86 NCEP/NCAR global reanalysis of $1/2^\circ \times 1/2^\circ$ resolution, Ruiz et al. (2008) computed a heat budget of -1 W m^{-2} .
87 However, their heat flux components values are not close to most of the literature values (for instance, the major
88 difference was in the value for net short wave with 84 W m^{-2}). Merullo et al. (2021) recently analysed several
89 atmospheric data sets, revealing a significant net heat flux variability ranging between 1.6 and 40 W m^{-2} . They
90 attributed this variability primarily to longwave radiation fluxes uncertainties. In addition to these challenges, past
91 studies of air-sea fluxes have primarily focused on establishing mean and variance, leaving limited knowledge about
92 their statistical distributions (Korolev et al., 2015; Tian et al., 2017). Understanding the probability distributions of
93 air-sea fluxes and their higher moments could provide insights into the uncertainties associated with air-sea physics.
94 Also, the analysis of probability distributions can help to assess skills of different reanalyses to replicate extreme
95 fluxes (Gulev and Belayaev, 2012).

96 In this study we investigate two very different aspects of the net surface heat budget closure problem of the
97 Mediterranean Sea. First, we employ two different by high quality surface atmospheric variable data sets at different
98 horizontal resolution and we calculate the heat fluxes with the same bulk formula and the same SST. This isolates
99 the impact of atmospheric model resolution and quality as the sole source of variation in the heat flux estimates.
100 Therefore, we answer the question: is the Mediterranean Sea in the past 15 years still losing heat at the surface?

101 Secondly, we study the statistical distributions of the heat flux components, utilizing the atmospheric analysis dataset
102 which are used to produce weather forecast by ECMWF (European Centre for Medium-Range Weather Forecasts).
103 Knowing the skewness and kurtosis distributions across the basin, we analyse the extremes of the net heat budget,
104 and we determine the specific importance of the extreme heat losses to the long-term mean. The second question
105 we answer is: what are the underlying causes of the net heat budget closure problem?

106 The paper is structured into the following sections. Section 2 presents the atmospheric analysis and reanalysis
107 datasets from ECMWF, along with satellite SST data and the bulk formula used in the estimation of the fluxes. In
108 Section 3, we present the new values of the heat budget closure problem, compared to the literature. In section 4,
109 we analyse the probability distributions of turbulent heat fluxes. In Section 5, we determine the causes of the long
110 term mean net heat budget values. Finally, Section 6 summarizes the findings and highlights key insights gleaned
111 from this research.

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115 2. Methodology and datasets

116 2.1 Air-sea physics in the Mediterranean Sea

117 For the Mediterranean Sea, several formulations have been established over the past decades through extensive
118 studies. In this section, we present these adopted formulations, beginning with the net heat flux formula, followed
119 by the specific heat flux components utilized in this study.

120 The net surface heat flux, Q_{net} comprises the net shortwave radiation Q_{SW} , net longwave radiation Q_{LW} , and surface
121 turbulent flux components, which encompass the latent heat flux of evaporation Q_{LH} and sensible heat flux Q_{SH}
122 (Cronin et al., 2019; Pettenuzzo et al., 2010).

$$123 \quad \quad \quad 124 \quad \quad \quad Q_{net} = Q_{SW} + Q_{LW} + Q_{lat} + Q_{sen} \quad (1)$$

125
126 Here, we use the convention that positive heat fluxes denote heat gain by the ocean. We did not use directly the
127 atmospheric model heat flux values since we wanted to intercompare two different atmospheric data sets in terms
128 of their quality and resolution not on the basis of the specific bulk parametrizations and SST used. Thus, we used
129 the same bulk formula and SST for both ECMWF and ERA5 surface variables that are described in section 2.2.

130

131 2.1.1 Shortwave radiation flux

132 The shortwave radiation flux (SW) is derived from the formulation proposed by Rosati and Miyakoda (1988). The
133 largest heat flux component is the solar radiation which is reduced by the cloud coverage and partially reflected by
134 the sea surface (albedo). The shortwave heat flux formula is therefore expressed as:

$$135 \quad \quad \quad 136 \quad \quad \quad Q_{SW} = Q_{TOT} (1 - 0.62 C + 0.0019 \beta)(1 - \alpha) \quad \text{if } C \geq 0.3 \quad (2)$$

$$137 \quad \quad \quad Q_{SW} = Q_{TOT} (1 - \alpha) \quad \quad \quad \text{if } C < 0.3$$

138

139 where Q_{TOT} indicates the clear sky solar radiation calculated with astronomical formulae, C (%) is the fractional
140 cloud coverage, β is the noon solar altitude in degrees and α is the ocean surface albedo varying month wise values
141 taken from Payne (1972). The incoming solar radiation varies on locations with sun zenith angel and Q_{TOT} reaches
142 at the ocean surface after diffusion can be represented by the components: the sum of the direct solar radiation Q_{DIR}
143 for direct solar radiation and Q_{DIF} for downward diffused radiation. Then net solar radiation Q_{TOT} can be
144 represented by the summation of components Q_{DIR} and Q_{DIF} :

$$145 \quad \quad \quad Q_{TOT} = Q_{DIR} + Q_{DIF}$$
$$146 \quad \quad \quad = Q_0 \tau^{secz} + [(1 - A_a) Q_0 - Q_0 \tau^{secz}] * 0.5$$

147 Here Q_0 is the solar radiation at the top of atmosphere, τ is equal to 0.7 and is the atmospheric transmission
 148 coefficient, A_a is a constant value (0.09) and z is the sun zenith angle.

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152 **2.1.2 Longwave radiation flux**

153 The longwave surface radiation flux (LW) is the difference between the upward infrared radiation (IR) emitted from
 154 the ocean surface (LU) and the atmospheric downwelling infrared radiation (LD). The LD component is adapted
 155 from Bignami et al. (1995), and the longwave radiation flux is written as:

156

$$157 \quad Q_{LW} = Q_{LU} + Q_{LD} \quad (3)$$

$$158 \quad Q_{LU} = -\epsilon \sigma_{SB} T_S^4 \quad (4)$$

$$159 \quad Q_{LD} = [\sigma_{SB} T_A^4 (0.653 + 0.00535 e_A)](1 + 0.1762 C^2) \quad (5)$$

160

161 where: T_S and T_A indicate the sea surface temperature and air temperature in degrees Kelvin, σ_{SB} is the Stefan-
 162 Boltzmann constant, ϵ is the ocean emissivity set to 1 according to Large and Yager (2009) and e_A is the atmospheric
 163 vapor pressure computed from the mixing ratio of the air W_{air} (Wallace and Hobbs, 2006).

$$164 \quad W_{air} = \frac{q_A}{1 - q_A} \quad (6)$$

$$165 \quad e_A = \frac{W_{air}}{(W_{air} + \gamma)} p \quad (7)$$

166 and q_A is the specific humidity of air, p is the surface air pressure, and γ is a constant (0.622).

167 The specific humidity (q_A) saturated at the T_A is computed using the following equation (Large, 2006), where $\rho =$
 168 1.22 kg m^{-3} is the air density and

$$169 \quad q_A = \rho^{-1} 640,380 \exp(-5107.4/T_D) \quad (8)$$

170 where, T_D is the dew point temperature retrieved from the atmospheric model outputs.

171

172 **2.1.3 Turbulent heat fluxes**

173 The turbulent heat flux is composed of sensible heat Q_{SH} and latent heat Q_{LH} given by the following formula:

$$174 \quad Q_{SH} = -\rho_A C_P C_H |\vec{V}| (T_S - T_A) \quad (9)$$

$$175 \quad Q_{LH} = -\rho_A L_E C_E |\vec{V}| (q_S - q_A) \quad (10)$$

177 where $|\vec{V}|$ is the wind speed, ρ_A is the density of moist air, C_p is the specific heat capacity ($1005 \text{ J g}^{-1}\cdot\text{K}$), C_H and
 178 C_E are turbulent exchange coefficients for temperature and humidity, L_E is the latent heat of vaporization, q_A is
 179 defined in (8) and q_S , which is the specific humidity of air saturated at the sea surface temperature T_S , is calculated
 180 with (8) using T_S instead of T_D , and applying a 0.98 factor (Sverdrup, 1942). Since the average wind speed in the
 181 Mediterranean is 5 m/s, Pettenuzzo et al. (2010) suggested using constant turbulent exchange coefficients such as
 182 $C_H = 1.3 \cdot 10^{-3}$ and $C_E = 1.5 \cdot 10^{-3}$.

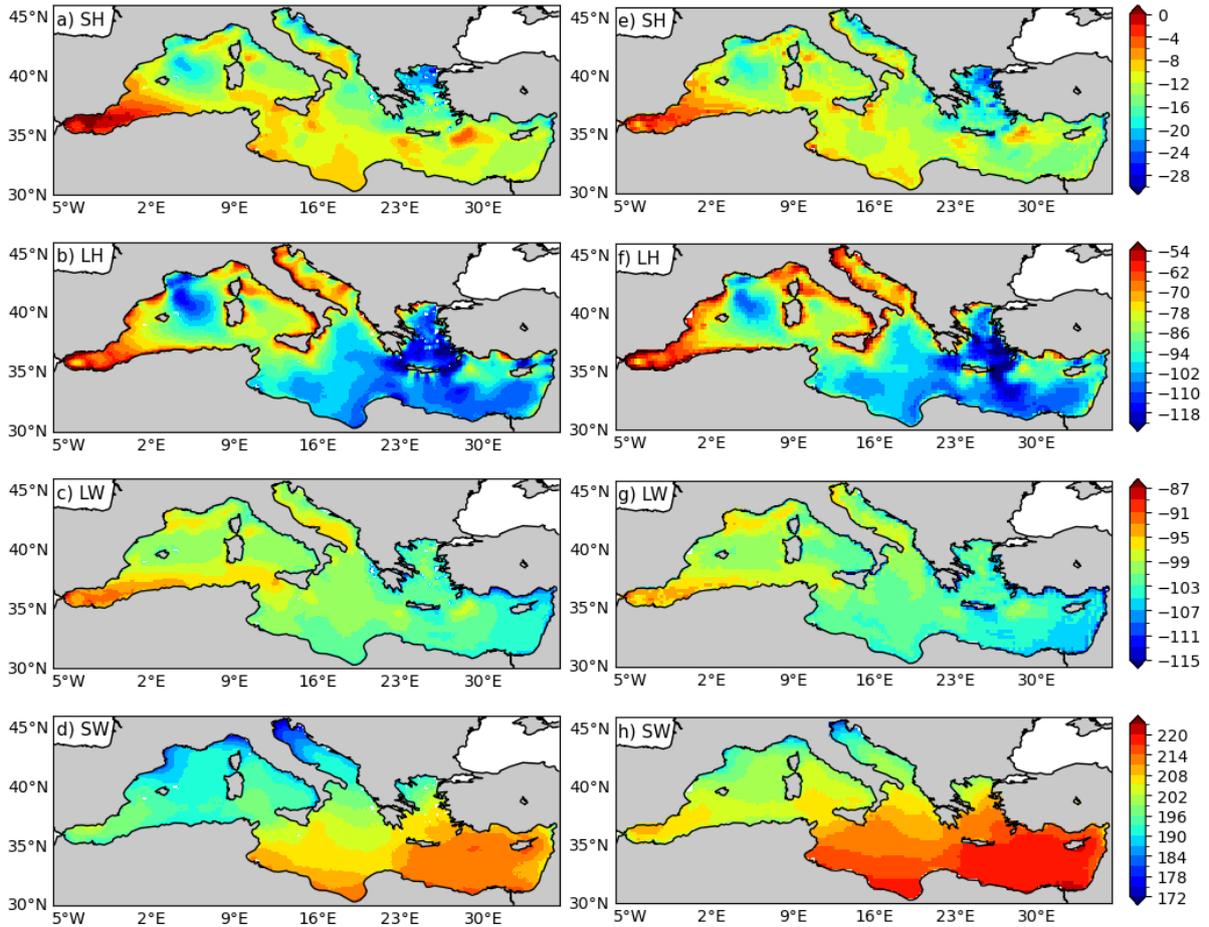
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184 2.2 Datasets

185 Two atmospheric datasets have been selected for this study. The first dataset is the ECMWF (European Centre for
 186 Medium-Range Weather Forecasts) high-resolution analysis dataset (Rabier et al., 2000) at six-hourly temporal
 187 resolution and 0.125 degrees of spatial resolution. It's worth noting that the original operational dataset, from which
 188 the atmospheric fields have been extracted, underwent changes between 1991 and 2006 in terms of model resolution
 189 and the assimilated number of observations. For consistency, we opted to utilize the dataset with approximately
 190 uniform model resolution and physics spanning a 15-year period from 2006 to 2020. The second dataset employed
 191 in this study is ERA5 reanalysis (Hersbach et al., 2020). This dataset is available at hourly intervals. However, it
 192 features a horizontal resolution of 0.25 degrees.

193 To mitigate unresolved atmospheric temperature daily cycles in ECMWF and make the two data sets consistent for
 194 the time variability, the ECMWF and ERA5 fields are further processed into daily mean values for the entire period.
 195 Comparisons conducted with daily and six-hourly input fields indicated minimal differences in the probability
 196 distributions of the heat fluxes, leading us to prioritize filtering out daily variability to the greatest extent possible.

197 To compute the heat fluxes the following atmospheric surface variables are extracted from the two datasets: the 10-
 198 meter wind components (U for the zonal direction and V for the meridional direction), mean sea level pressure, dew
 199 point temperature, total cloud coverage, and 2-meter air temperature.



200
 201 **Figure 1: Mean annual heat flux components for the period of 2006-2020 computed from ECMWF (left**
 202 **panel) and ERA5 (right panel) daily time series. The corresponding ECMWF time series is shown in**
 203 **supplementary material, Fig. S1 and S2.**

204
 205 For the oceanic SST data, we utilized the satellite dataset distributed by the Copernicus Marine Environment Service
 206 (CMEMS). This SST dataset is a blended product from multiple satellite sensors, categorized as L4, with a
 207 horizontal resolution of $0.05^\circ \times 0.05^\circ$. To align the SST data with the atmospheric analysis and reanalysis dataset
 208 grids, we applied an interpolation and extrapolation method known as the 'sea-over land' (De Dominicis et al., 2013).
 209 This method involves an iterative process to extrapolate sea values over land before interpolating, thus not allowing
 210 the contamination of land values on the interpolation.

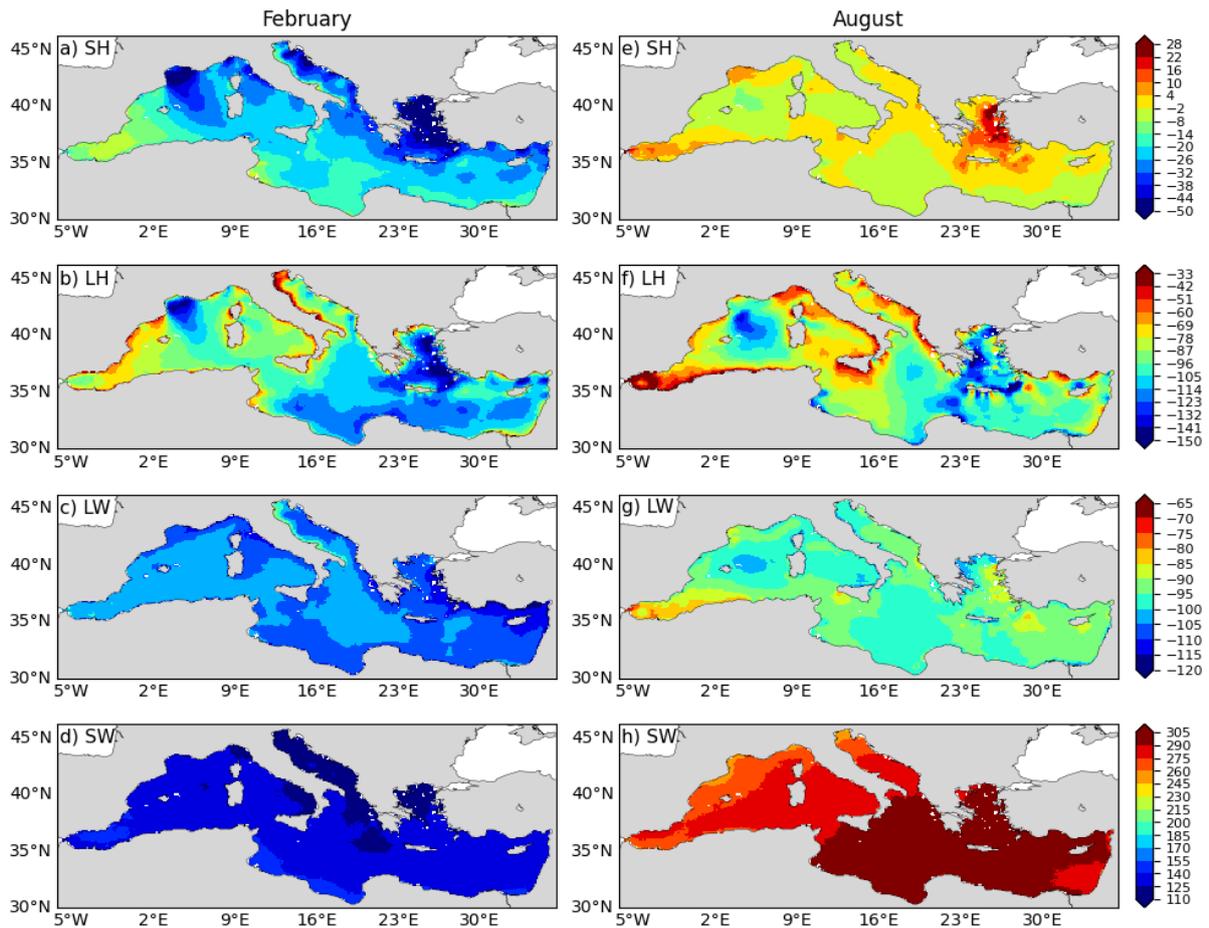
211
 212 **3. Heat budget closure problem revisited**

213 **3.1 Analysis of the heat budget components**
 214

215 We compute the heat fluxes for the 15-year period, 2006–2020, using the ERA5 dataset and compare them with the
 216 fluxes computed using the ECMWF dataset (Fig. 1). In Fig. 1 we show the results for the 15-year mean of each
 217 heat budget components. We start describing the ECMWF patterns and then we detail the differences.

218 Turbulent heat fluxes exhibit distinct sub-basin-scale patterns, varying between the eastern and western
219 Mediterranean Seas as well as the Central Mediterranean region. The smallest mean sensible heat loss is observed
220 in the whole Alboran sea area with absolute value range of 0-6 W m^{-2} , while the Aegean Sea and the centre of Gulf
221 of Lion loses more heat in the maximum value of 25 W m^{-2} . Similarly, the highest absolute values of LH are
222 recorded in the Gulf of Lion and the Aegean and Levantine Seas, attributed to the influence of strong and cold winds
223 like the Mistral and Etesian in the north-western and eastern Mediterranean regions, respectively. The eastern
224 Mediterranean emerges as the region with the highest evaporation, reaching approximately 122 W m^{-2} in absolute
225 value. Notably, along the south-eastern coastline, a wide range of maximum absolute values (102-122 W m^{-2}) in
226 the evaporation is observed. The turbulent heat fluxes show limited differences between the ECMWF and ERA5
227 datasets.

228 SW fields show the well-known meridional gradients with larger gradient values arising from the ECMWF dataset.
229 The mean SW exhibits a gradual decline from the eastern to western Mediterranean, influenced by the variation of
230 the solar zenith angle with longitudes. The first reason for SW differences between Western Mediterranean and
231 Eastern Mediterranean is the latitudinal position of each sub-basin. Furthermore, SW differences using ECMWF
232 and ERA5 datasets are connected to different cloud cover schemes (not shown), leading to a larger heat gain in the
233 Eastern Mediterranean. Notably, the northern Adriatic region stands out with a distinct distribution, suggesting it
234 receives relatively less annual solar radiation compared to other areas. In contrast, the mean longwave (LW)
235 radiation distribution maintains a relatively consistent range of absolute values between 87–113 W m^{-2} across the
236 entire domain with absolute minimum values in the Alboran Sea, presumably due to the warm Atlantic surface water
237 inflow. Overall, while the turbulent heat fluxes show limited differences between the ECMWF and ERA5 datasets,
238 significant discrepancies are observed in radiative heat fluxes. Additionally, Fig. 1 shows the noisiness of the
239 fluxes due to the ERA5 low resolution with respect to ECMWF while retaining an overall consistency.



240
 241 **Figure 2: Seasonal variations of heat flux components: Left column is the monthly average values for**
 242 **February and right column the average for August for the period 2006-2020 (ECMWF data).**

243
 244 Figure 2 shows seasonal variations in heat flux components for February and August using ECMWF data. Both SH
 245 and latent heat (LH) fluxes exhibit a greater spatial gradient in February compared to August. In winter, the SH loss
 246 is larger, especially in the Gulf of Lion, Aegean, and parts of the Adriatic, with stronger spatial gradients compared
 247 to summer. In August, SH flux becomes positive in the Aegean and the Alboran Sea. LH loss is highest in February
 248 in the whole eastern Mediterranean and the Gulf of Lion. In August, LH losses decrease in the western
 249 Mediterranean, with absolute value minima in the Alboran and Adriatic Sea, remaining largely negative in the lower
 250 part of in the Eastern Mediterranean. SW fields show the strongest seasonal cycle as expected, with the absolute
 251 maximum of 260-305 $W m^{-2}$ in summer and in the Eastern Mediterranean. LW is largest in absolute value in winter
 252 showing a small seasonal cycle. Significant seasonal variations are observed in the distribution range for radiative
 253 heat fluxes, low in February and high in August across the entire domain. These patterns are quite similar to the
 254 ones reported in the literature.

255
 256
 257

258 3.2 Net heat budget estimation

259
260 The net surface heat flux Q_{net} is depicted in Figure 3 for ERA5 and ECMWF and basin-average 15 year mean
261 values are listed together with the literature in Table 1.

262 Fig. 3 shows that the Gulf of Lion and the Aegean Sea are the areas of maximum heat losses while the basin gains
263 heat in the Alboran Sea, in some areas of the Levantine basin, and in the shelf areas around the Italian peninsula.

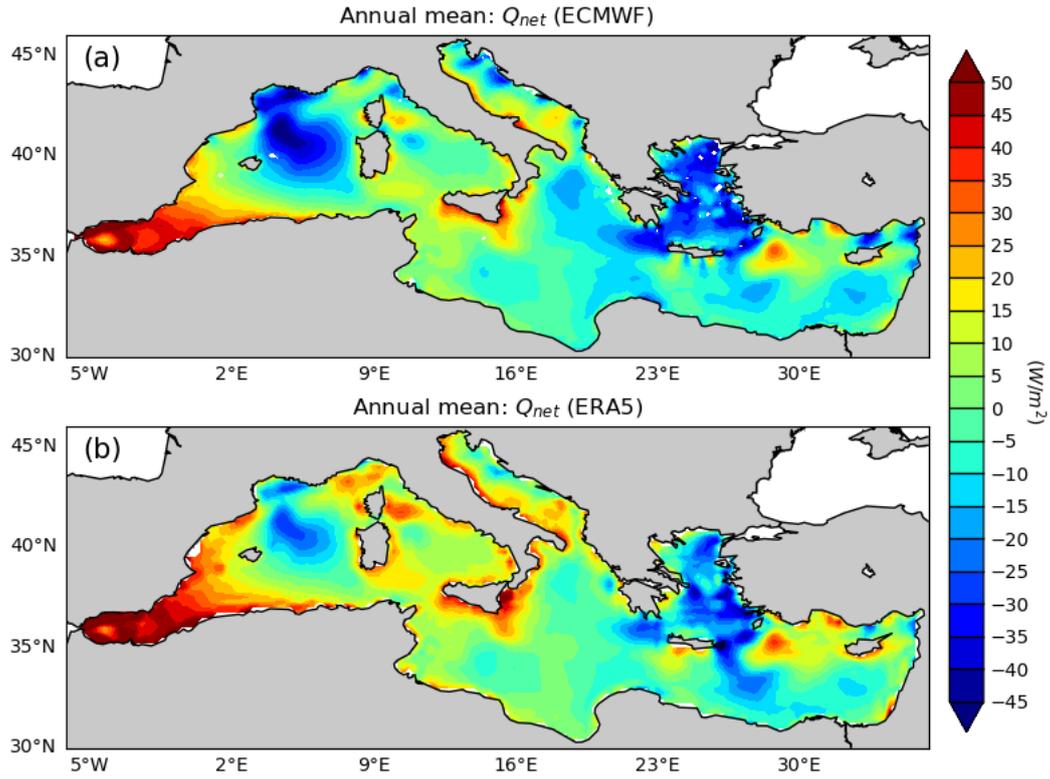
264 The Mediterranean Sea gains comparatively more heat with the ERA5 inputs. Besides the difference in surface
265 domain for Q_{net} , for both cases, air-sea flux dynamics is strongly visible in the Alboran Sea for net heat gain, in
266 the Gulf of Lion for heat loss due to the continental cold wind (Mistral wind), and in the Aegean Sea due to the
267 strong wind (Etesian wind) that blows during the summer period. Using ECMWF inputs, Q_{net} is -3.6 W m^{-2} , a
268 value consistent with previous estimates for the Mediterranean Sea domain and for ERA5, it is 5 W m^{-2} (Table 1).

269 Errors in Q_{net} mean value are determined by a bootstrapping method where Q_{net} time series is resampled 5000
270 times to compute a standard deviation around the mean of the resampled time series (Tibshirani & Efron, 1993).

271 We argue that our results show that the negative heat budget is achieved by using only ECMWF fields at high
272 resolution, i.e. 0.125 degrees. Higher resolution implies differences in all atmospheric fields used to compute the
273 fluxes. Furthermore, ERA5 and ECMWF model physics and dynamics is different contributing to the differences in
274 the mean heat budget. However, both datasets use observations, and we argue that the most relevant difference
275 between the analysis and the reanalysis data set is the resolution due to the peculiar geometry of the Mediterranean
276 Sea.

277 Since all the literature datasets are coarser, this is most likely the reason of the failure to determine the correct heat
278 budget closure value. In Pettenuzzo et al. (2010) several ad-hoc corrections were made to the surface atmospheric
279 fields to obtain the negative heat flux budget while in Sanchez-Gomes et al., (2011) they used an ensemble of
280 deterministically downscaled ERA40 fluxes giving rise to a very large uncertainty. Considering a recent literature,
281 our resulted Q_{net} is closely matched with the computed net heat budget of $-3 \pm 8 \text{ W m}^{-2}$ from Jordà et al., (2017), but
282 their result was associated to large temporal uncertainties from the surface fluxes through Gibraltar Strait.

283



284
 285 **Figure 3: Comparison of the annual Q_{net} ($W m^{-2}$) computed from, a) ECMWF and b) ERA-5 input datasets.**
 286

287 **Table 1: Computed heat flux components and net heat fluxes (Q_{net}), and values from the references**

Authors	SH	LH	LW	SW	Net Flux (Q_{net})
Bethoux (1979)	-13	-120	-68	195	-6
Bunker (1982-1)	-13	-101	-68	202	20
Bunker et al (1982-2)	-11	-130	-68	202	20
May (1986)	-11	-130	-68	193	2
Garret et al. (1993)	-7	-99	-67	202	29
Matsoukas et al. (2005)	-11	-122	186	-63	22
Ruiz et al. (2008)	-8	-88	-73	168	-1
Pettenuzzo et al. (2010)	-14	-90	-79	178	-7
Sanchez-Gomez et al. (2011)	-13±5	-100±13	-75±6	181±18	-7±21
Criado-Aldeanueva et al. (2012)	-15.1	-93.5	-76.9	186.3	0.73
Song & Yoy (2017)	-13±4	-98±10	-78±13	192±19	2±12
Jordá, et al., 2017	-	-	-	-	-3±8
ECMWF analyses	-12.1±4	-92±16	-100.5±3	201±8	-3.6±1.3
ERA5 reanalysis	-13±3	-89±14	-101±3	208±8	5±1.2

288 Spatially, the mean Q_{net} distribution generally shows a heat loss across much of the Eastern Mediterranean. Overall,
 289 distributions of more positive net heat budget values for the western Mediterranean and negative for the eastern
 290 Mediterranean have matched with the similar result from Criado-Aldeanueva et al. (2012). Strong spatial gradients
 291 are evident, particularly in the Aegean Sea, although a few patches displaying net heat loss (negative Q_{net}) are also
 292 noticeable in this vicinity. Conversely, the western Mediterranean exhibits a stronger heat gain area, which appears
 293 particularly concentrated zone in the Gulf of Lion region and this feature is apparent in results from both atmospheric
 294 datasets. Such a spatial related uncertainty in Q_{net} represents a significant challenge for accurately closing regional
 295 heat budgets as well as validating existing ocean circulation models within the complex Mediterranean basin.

296

297 **4. Probability distributions of the turbulent heat fluxes**

298 In this section, we analyse the probability distribution of turbulent heat fluxes computed using ECMWF data set
 299 only and for the anomaly heat fluxes. Recent studies by Gulev and Belyaev (2012) and Korolev et al. (2015) have
 300 analysed the statistical distributions of turbulent heat fluxes, and their findings are used here for comparison.
 301 Radiative flux components are excluded from this analysis, as they do not exhibit extremes of comparable magnitude
 302 to those of turbulent fluxes (Supplementary material, Fig S3). This suggests low skewness and kurtosis in their
 303 distributions, reducing the relevance of a detailed probability density function analysis for these components.

304

305 If we indicate the time series of each component of the heat budget with X_n we can define the heat flux climatology
 306 as:

307

$$308 \quad Q_t = \frac{1}{n} \sum_{j=1}^n X_{tj} \quad (11)$$

309

310 where ‘t’ indicates the day of the year, and ‘j’ is the number of years. The anomaly time series is computed by
 311 subtracting the long-term seasonal climatology Q_t from the observed heat flux time series X_{tj} and it will be indicated
 312 by:

313

$$314 \quad \tilde{X}_{tj} = X_{tj} - Q_t \quad (12)$$

315

316 **4.1 SH flux distribution**

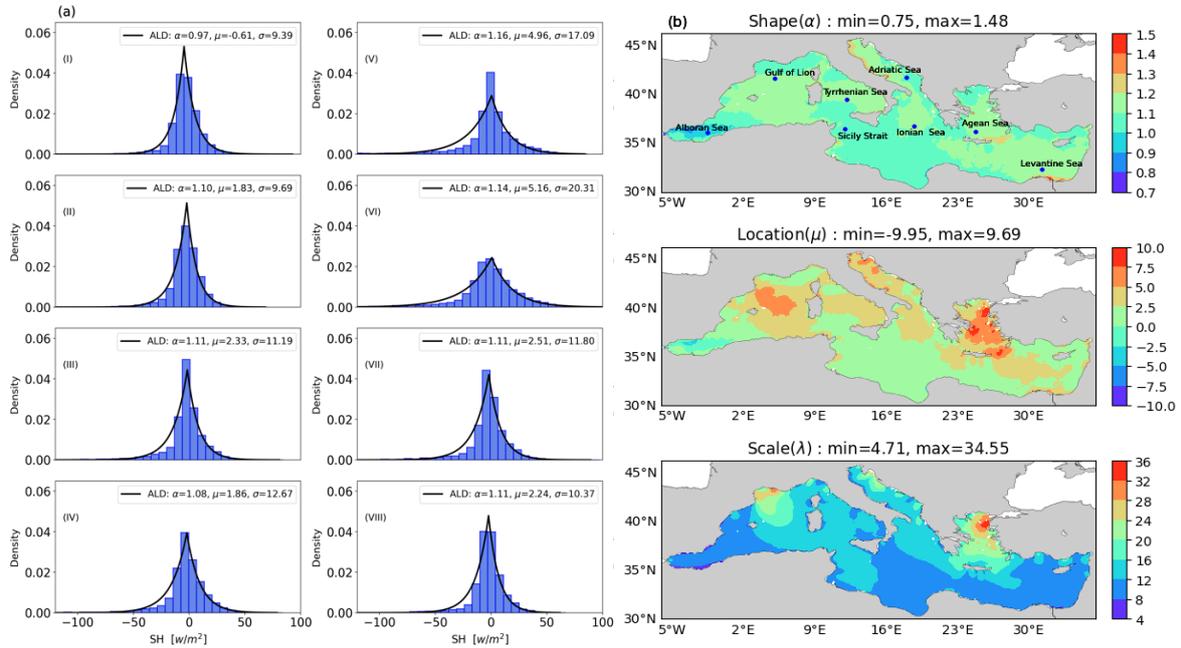
317 We found that gaussian or skew-normal distributions are not well fit for SH flux, as evident from the histograms at
 318 single grid points shown in Figure 4a. The histograms reveal a singularity around zero, indicating that the skew-
 319 normal distribution may not adequately capture the distribution of these values. This observation is consistent with
 320 findings by Gulev and Belyaev (2012), and we provide further explanation in the Appendix A.

321 The most common distribution with such near-discontinuous behaviour at the origin is the three-parameter
 322 Asymmetric Laplace Distribution (ALD) (Yu & Zhang, 2005) that we can defined as

323
$$F(x, \alpha, \mu, \lambda) = \frac{\lambda}{\alpha + \frac{1}{\alpha}} \begin{cases} \exp\left(\frac{\lambda}{\alpha}(x - \mu)\right) & \text{if } x < \mu \\ \exp(-\lambda \alpha(x - \mu)) & \text{if } x \geq \mu \end{cases} \quad (13)$$

324 where x is the random variable time series, α is the shape parameter, μ is the location and λ the scale.

325



326 **Figure: 4 a) The single grid point histograms for SH flux anomalies from the eight sampling locations for the**
 327 **period of 2006–2020, 4 b) The Asymmetric Laplace PDF parameter (α, μ, λ) distributions from computed SH**
 328 **flux anomaly for the observation period. [Sampling points: (I) Alboran Sea, (II) Gulf of Lion, (III) Tyrrhenian**
 329 **Sea, (IV) Sicily Strait, (V) Adriatic Sea, (VI) Ionian Sea, (VII) Aegean Sea, (VIII) Levantine Sea]**
 330

331

332 From the single grid point histogram, we have observed a one or two sharp peaks in the distribution that matches
 333 well with the Asymmetric Laplace Distribution (ALD) PDF. In accordance with findings by Yu and Zhang (2005),
 334 the distribution of the sensible heat (SH) flux anomaly time series exhibits characteristics of a double exponential
 335 distribution. This is evident from the histograms displaying both positive and negative skewness with long tails, as
 336 depicted in Figure 4a. The ALD parameters for the SH flux anomaly time series are illustrated in Figure 4b. The
 337 shape parameter (α) falls within the positive range of 0.73 to 1.48, indicating a moderate to high degree of peakiness
 338 in the distribution. Additionally, the location parameter (μ) exhibits mostly positive values while a small area in the
 339 Alboran Sea shows negative values, suggesting a shift in the central tendency of the distribution. Notably, the scale
 340 parameter (λ) displays a similar structure to the SH flux climatology depicted in Figure 1.

341 To check the quality of the fit, moments of both applied and theoretical PDF are compared (presented in
 342 supplementary materials, Fig. S4). The comparison shows the estimations of the three moments in the left panel for
 343 the observed SH flux and right panel for ALD PDF parameters. It can be seen that variances and skewness are
 344 similar in distribution while kurtosis differ at noticeable range. This observation is likely attributed to the fact that
 345 the kurtosis for the asymmetric Laplace distribution remains constant regardless of changes in the scale parameter.

346

347

348 4.2 LH flux distribution

349 In the case of the LH flux, no sharp exponential peaks were observed; instead, large skewness and long tails were
350 identified. Therefore, we applied the skew-normal PDF which is defined by $\alpha (\in \mathbb{R})$ as the shape parameter, $\mu (\in$
351 $\mathbb{R})$ the location parameter, and $\lambda > 0$ the scale parameter (Azzalini, 1985) and defined as:

$$352 f(x, \alpha, \mu, \lambda) = \frac{2}{\lambda} \phi\left(\frac{x-\mu}{\lambda}\right) \Phi\left(\alpha \frac{x-\mu}{\lambda}\right) \quad (14)$$

353

354 Where,

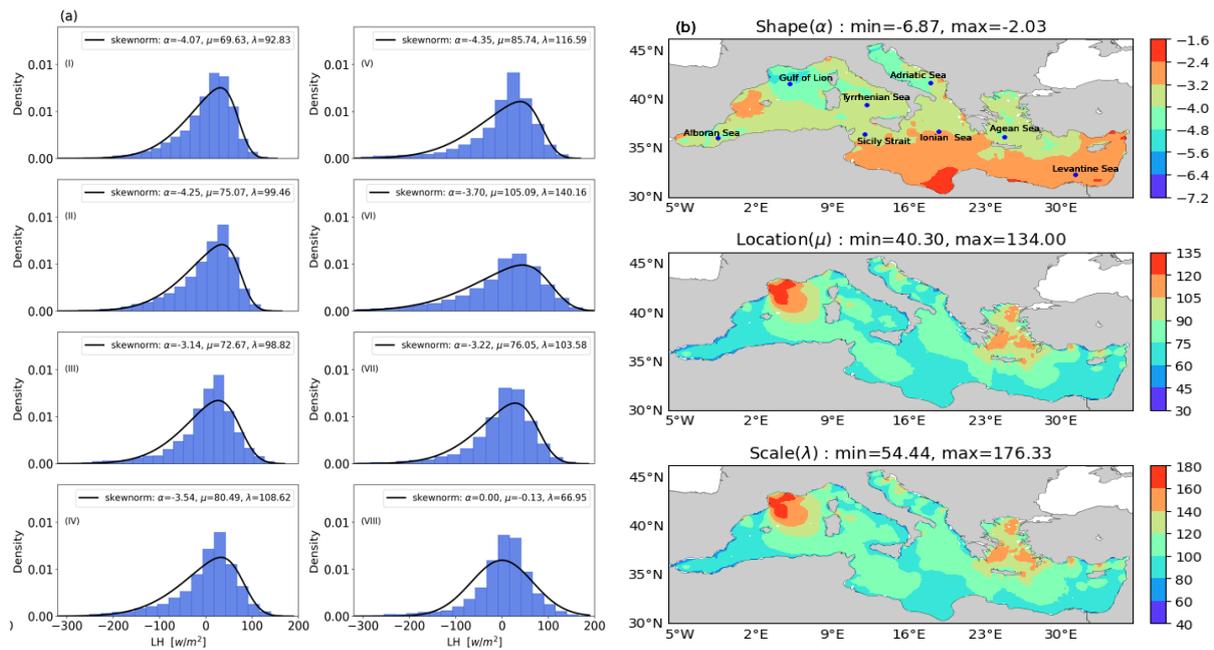
$$355 \phi\left(\frac{x-\mu}{\lambda}\right) = \frac{1}{\sqrt{2\pi}} e^{-\frac{1}{2} \frac{(x-\mu)^2}{\lambda^2}} \quad (15)$$

$$356 \Phi\left(\alpha \frac{x-\mu}{\lambda}\right) = \int_{-\infty}^{\alpha \frac{x-\mu}{\lambda}} \phi(t) dt \quad (16)$$

357 A skew-normal PDF is an extension of the normal distribution while covering the skewness and containing the
358 general characteristics of a Gaussian distribution (Flecher et al., 2010).

359 To examine visually the quality of PDF fit on LH flux anomaly values, histograms from eight sea locations were
360 fitted with the skew-normal PDF, as shown in Figure 5a. Figure 5b displays the parameter spatial variability. The
361 shape parameter distribution ranges from -6.83 to -2.5, with negative values observed across all points. This spatial
362 distribution of α , exhibiting a negative range, aligns with the negatively skewed pattern identified in the single grid
363 point PDF fitting test. Furthermore, the spatial distribution structure of the location and scale parameters
364 demonstrates a positive correlation across most locations.

365 In the Supplementary material, a comparison of statistical moments is conducted to qualitatively validate the fit
366 (supplementary materials, Fig S5). There is notable agreement in the variance distributions between the observed
367 LH flux anomaly and skew-normal PDF. While the skewness distributions mismatch at negligible level, with the
368 theoretical PDF skewness predominantly ranging from -0.9 to -0.3, whereas the observed skewness exhibits a
369 variation range spanning from -1.2 to over -0.3. Lastly, the kurtosis distribution of the skew-normal PDF differs in
370 the Aegean Sea, Alboran Sea and Gulf of Lion area.



372 **Figure 5 a) The single grid point histograms for LH flux anomalies at the eight sampling locations for the**
 373 **period of 2006-2020, 5 b) The skew-normal PDF parameter (α , μ , λ) distributions for computed LH Flux**
 374 **anomaly for the observation period [Sampling points: (I) Alboran Sea, (II) Gulf of Lion, (III) Tyrrhenian**
 375 **Sea, (IV) Sicily Strait, (V) Adriatic Sea, (VI) Ionian Sea, (VII) Aegean Sea, (VIII) Levantine Sea]**

376

377 4.3 Evaluation of the PDF fitting

378 In this section, we conducted a goodness of fit test to measure the distance between the empirical observed
 379 distribution and the fitted one. The objective of this evaluation test was to assess the degree of agreement between
 380 the applied theoretical distribution and the observed time series. The chi-squared method, a well-accepted test, was
 381 employed to measure the distance between two independent distributions.

382 We compared the results of the chi-squared test for the turbulent heat fluxes computed using the ECMWF and ERA5
 383 datasets. The decision rule for the χ^2 test was determined based on the level of significance, set at 0.05, and the
 384 degrees of freedom, defined as $DF = N - np$, where N represents the number of bins and np is the number of
 385 distribution parameters (i.e., 3 for both the ALD and skew-normal distribution). In the supplementary material we
 386 show the maps of Chi-square test statistics (Supplementary material, Fig. S6). The chi-squared results for the SH
 387 and LH fluxes computed using the ECMWF dataset indicate that almost all surface grid points are well-fitted with
 388 the applied theoretical PDFs, the ALD and skew-normal PDF. With the critical threshold of 33.92 (Elderton, 1902)
 389 for P values, we observed a very few mismatches location near the coasts.

390

391 5. How do heat loss extremes contribute to the heat budget closure?

392 The heat budget closure problem is associated with achieving a net negative heat flux, as discussed before. We test
 393 here the hypothesis that the negative long term mean negative heat budget of Table 1 for ECMWF data is correlated
 394 to the extremes in heat losses during autumn-winter.

395 Figure 6 illustrates the Q_{net} basin average daily time series, revealing a value range varying between 200 and -500
 396 $W m^{-2}$. Notably, the most pronounced extreme negative heat loss peaks, reaching up to 500 $W m^{-2}$ occur in the
 397 winters of 2011, 2015 and 2017. They approximately coincide with western Mediterranean Deep Water formation
 398 events, as documented in Escoudier et al. (2021). To identify and remove the potential extremes in our computed
 399 Q_{net} time series, we apply the Interquartile Range (IQR) method which measures the spread of a dataset and
 400 calculate the difference between the third quartile(Q3) and the first quartile (Q1). The IQR threshold is computed
 401 by the difference between the 1st quartile (Q1) and 3rd Quartile (Q3) of the observed dataset:

402

$$403 \quad IQR = Q3 - Q1 \quad (17)$$

$$404 \quad \text{Threshold} = Q1 - k * IQR \quad (18)$$

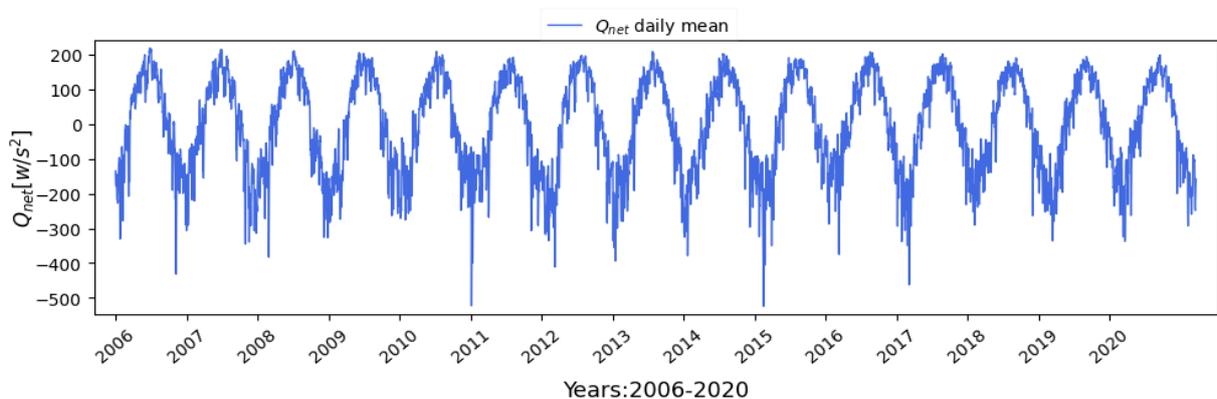
405

406 We used different values for k to exclude the negative extremes, which correspond to the maximum heat losses.
 407 These extreme values were replaced with long-term daily climatological values (using equation 11) to the respective
 408 days of extremes heat losses occurred and the long term mean neat heat budget Q_{net} is recomputed.

409 The resulting Q_{net} for different thresholds is displayed in Table 2 and the thresholds are shown in Fig. 7 together
 410 with the daily seasonal climatology. If compared with the long term mean heat budget in Table 1 ($-3.6 \pm 1.3 W m^{-2}$)
 411 we see that eliminating the winter extremes produces a smaller long term mean heat loss up to changing the sign to
 412 positive values. We argue that the ECMWF net negative heat extremes are the reason why ECMWF has a negative
 413 long term mean budget. Furthermore, if we calculate the yearly mean value of the seasonal climatology, we obtain
 414 the value of $+4 W m^{-2}$, which confirms again the importance of extremes in the heat budget closure of the
 415 Mediterranean Sea.

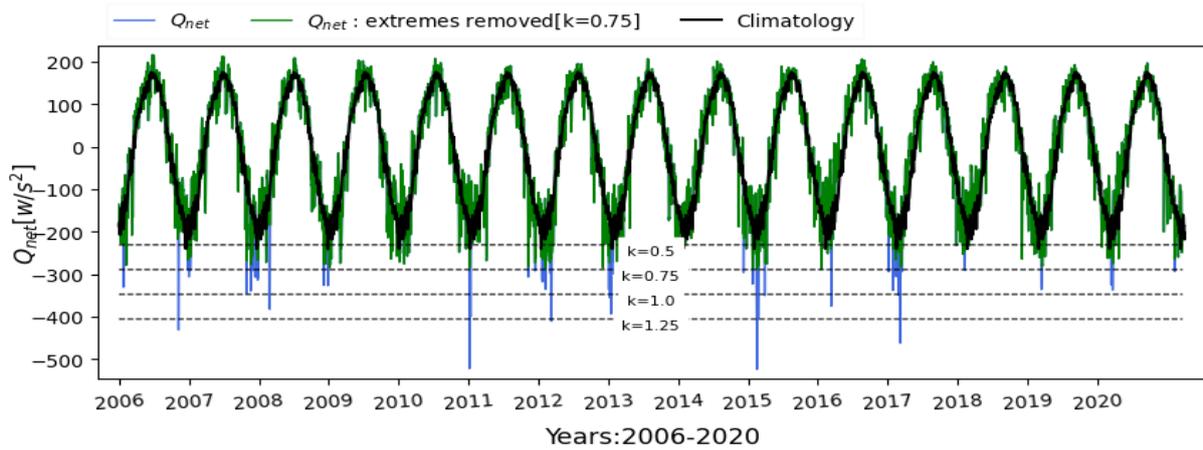
416 The Q_{net} could become an impact indicator of the Mediterranean for sea level trends in the basin. The net heat
 417 budget in fact relates to the sea level tendency (Pinardi et al., 2014) in the Mediterranean Sea and could be considered
 418 as a key indicator of climate impacts in the Mediterranean Sea.

419



420

421 **Figure 6: Basin averaged time series of the computed daily Q_{net} (units $W m^{-2}$) from the ECMWF computed**
 422 **heat fluxes, for the period 2006-2020.**



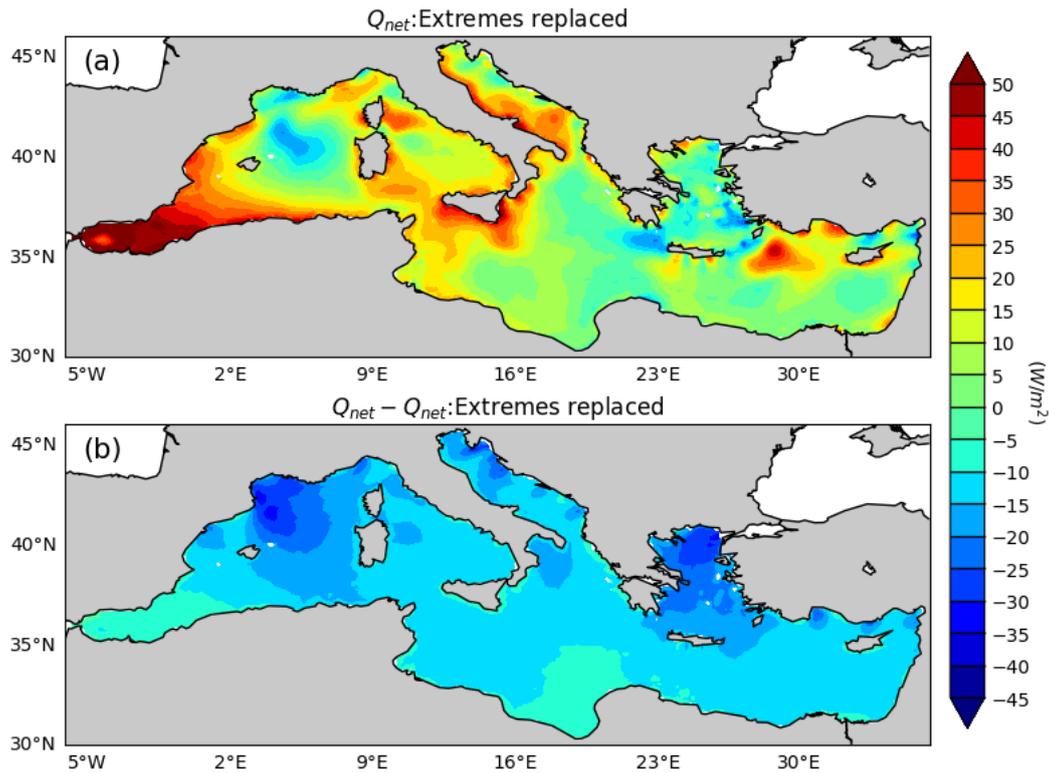
423
 424 **Figure 7: Time series of the basin averaged Q_{net} means, Q_{net} extremes removed, and long-term yearly**
 425 **climatology and four lower quantile boundary line marked with dashed lines using different k values [k=1.25,**
 426 **1,0.75, 0.5].**

427
 428 **Table 2: Different lower quantile boundary limits used to replace potential extremes and the resulting long-**
 429 **term mean basin averaged Q_{net} values**

IQR lower boundary limit	Threshold values ($W m^{-2}$)	$Q_{net}(W m^{-2})$
K=1.25	405	-3.2
K=1.0	347	-2.5
K=0.75	289	-1
K=0.5	231	2

430
 431 Figure 8 presents the new long-term mean spatial distribution of the surface heat budget after removing negative
 432 extreme values using a threshold of $-289 W m^{-2}$ ($K = 0.75$). The figure illustrates that these extreme events exert a
 433 substantial influence on the overall structure of the net heat budget in the Mediterranean Sea, with particularly
 434 pronounced effects in the Gulf of Lion, the Aegean Sea, the eastern Adriatic Sea, and along the southern Turkish
 435 shelves.

436



437
 438 **Figure 8: The annual mean after the removal of extremes with significant reduction of negative heat fluxes**
 439 **in the Gulf of Lion, Adriatic Sea and Aegean Sea regions.**

440

441

442 **6. Discussion and conclusions**

443

444 The primary motivation behind this investigation is to revisit the heat budget closure hypothesis from atmospheric
 445 consolidated data sets that are nowadays used frequently to drive ocean models. For this analysis, we covered a 15-
 446 year period from 2006 to 2020 with a daily time series frequency. The reason for the choice of this time range is
 447 that ECMWF analyses became quite stable starting from 2006 while before the model was at coarse resolution, like
 448 ERA5's model. Our strategy is to use the same bulk formula for the ERA5 and ECMWF data sets and compute the
 449 long term mean heat budget. Our strategy is to use the same SST and the same bulk formula but different atmospheric
 450 reanalysis and analysis surface variable data sets and compare the value of the long term mean heat budget in the
 451 Mediterranean Sea.

452 Firstly, the heat budget of the Mediterranean Sea was analysed to examine average annual mean and seasonal
 453 variations. The largest component of the heat budget is the net solar radiation (SW), followed by the latent heat
 454 (LH), longwave radiation (LW), and then sensible heat (SH), as shown in the literature. All heat flux components
 455 exhibit significant seasonality, as illustrated in Figure 3. Differences appear in the structure of the fluxes, especially
 456 the SW and LW, when different atmospheric data sets are used, a conclusion aligning with a suggestion from Marullo
 457 et al. (2021) on the sensitivity of LW estimates from the atmospheric dataset used to calculate fluxes. We compared

458 the ERA5-derived surface radiative fluxes with our own estimates and found that the ERA5 longwave (LW) fluxes
459 are substantially overestimated in absolute magnitude (Supplementary Material S7). The associated uncertainty is
460 comparable in order of magnitude to that reported by Marullo et al. (2021), who analysed an observational dataset
461 at a specific site. Nonetheless, compensating biases between the SW and LW components (Supplementary Material
462 S8) result in a net radiative balance that diverges primarily in the southern Mediterranean, where ERA5 exhibits
463 reduced LW flux values. These findings indicate that the ERA5 dataset may not adequately represent surface
464 radiative fluxes in the Mediterranean Sea

465
466 The basin-average net heat flux, Q_{net} , was calculated to be $-3.6 \pm 1.3 \text{ W m}^{-2}$ for ECMWF analysis data while it is
467 $5 \pm 1.2 \text{ W m}^{-2}$ for ERA5 (Table 1). This finding supports the conclusion that heat budget closure hypothesis cannot
468 be satisfied with a relatively coarse reanalysis atmospheric data set. Our initial question was: is the Mediterranean
469 Sea in the past 15 years still losing heat at the surface? The answer is yes if we use a high-resolution ECMWF
470 atmospheric analysis. Additionally, comparing the Q_{net} estimates derived from ERA5 and ECMWF with the same
471 bulk formulas demonstrates that the uncertainty peaks in the atmospheric forcing resolution and possibly cloud
472 cover. This uncertainty impacts the regional heat budget closure hypothesis.

473
474 Furthermore, we have demonstrated that the probability density of surface heat fluxes can be modelled and fitted
475 with a three-parameter PDF composed of a shape, a location, and a scale parameter. All the turbulent heat flux
476 components show asymmetric behaviour. There is encouraging agreement between the first two statistical moments
477 of the fitted PDF and the observed values. Kurtosis does not seem to be properly captured by the PDF used but our
478 time series is too short to arrive at a definitive conclusion. For the SH we demonstrate that the ALD PDF is generated
479 by the contributing distributions of wind speed (Weibull) and temperature difference, combined to form the heat
480 flux. We believe this is the first time that such kind of transformation is demonstrated.

481 Gulev and Belyaev (2012) applied the two-parameter Fisher–Tippett distribution (also known as the Gumbel
482 distribution) to monthly sensible and latent heat fluxes derived from NCEP–NCAR reanalysis fields. Their
483 approach focused on using the mean and standard deviation to estimate the distribution’s location and scale
484 parameters relevant to extreme events. However, the Gumbel distribution has a fixed skewness, limiting its ability
485 to capture the contribution of rare, asymmetric extremes. In contrast, our study analyses anomalies from the
486 seasonal cycle using full probability distributions that allow for variable skewness. This better reflects the nature
487 of atmospheric and oceanic variables, which are often inherently skewed (Sardeshmukh and Penland, 2015), and is
488 essential for understanding the influence of extremes on the surface heat budget. Our findings show that
489 incorporating a shape parameter is key to accurately capturing distribution structure and preserving asymmetric
490 tails. This analysis provides a useful framework for validating surface flux products and assessing their variability,
491 particularly important given that surface fluxes are the dominant source of uncertainty in the Mediterranean net heat
492 balance (Jordà et al., 2017). Correctly estimating skewness is crucial, as a small number of extreme outliers,
493 especially during intense winter events, can disproportionately affect the basin-wide mean and determine whether
494 heat budget closure is achieved.

495 For the first time, we have investigated the effects of extreme heat losses in the Mediterranean Sea in the long term
496 mean basin averaged heat budget. The northern basin areas are the site of the largest heat losses (Gulf of Lion and
497 the Aegean Sea, Adriatic Sea and the Turkish southern coasts). Exclusion of the negative extremes in these areas
498 resulted in a change in the sign of long term mean heat loss. The anomaly threshold value of -231 W m^{-2} (Table 2)
499 results in a long-term positive net heat flux, which is inconsistent with the basin's energy closure assumption,
500 thereby indicating the presence of long-term changes within the basin due to atmospheric forcing. Our second
501 initial question was: what is the cause of the Mediterranean Sea negative long-term mean heat budget? The answer
502 is that the long-term mean, basin averaged heat loss is due to winter extremes in the Northern regions of the
503 Mediterranean Sea.

504 In conclusion, understanding the characteristics and distributions of air-sea heat fluxes are crucial for gaining
505 insights into variations in the heat budget. Furthermore, the PDF analysis of turbulent heat fluxes will allow us to
506 have a better understanding of the extreme events and their contributions in the net negative heat budget. The next
507 steps could involve a machine learning study of air-sea flux bulk parametrizations for different atmospheric data
508 sets and coupled models, using as target the data set from this study.

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516 **Author contributions**

517
518 MHG: development of the concept, literature review, writing, methodology, coding, formal analysis, wiring,
519 visualization. NP: conceptualization, review, writing, methodology. AN: conceptualization, writing, review. LM:
520 review, writing. SB: methodology, review. FM: methodology, review. FT: methodology, coding.

521
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525 MHG and NP.

526

527 **Conflict of interest Statement**

528 The authors declare no conflicting interests.

529

APPENDIX A

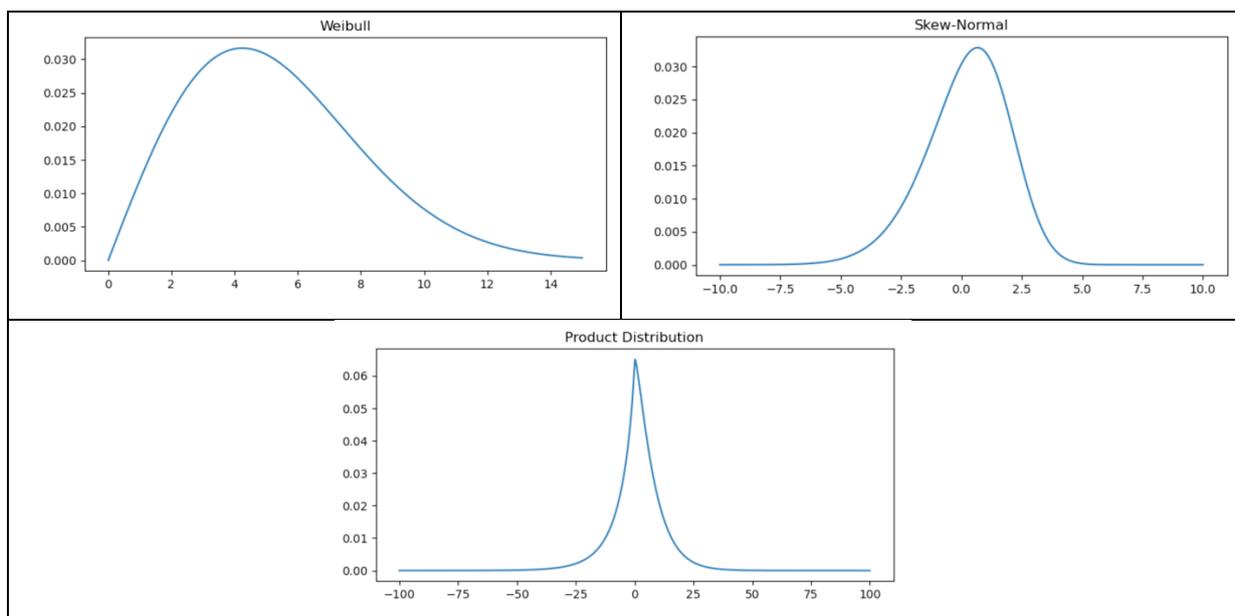
530

531 Here, we show that the characteristics of the SH flux distribution are due to the specific form of the heat flux as
532 given by (10), i.e. a multiplication of two distributions, wind speed and temperature differences, that are described
533 by a different distribution.

534

535 Let's indicate with $P(v*DT)$ the combined SH distribution of $Q(v)$ for wind speed and $R(DT)$ the temperature
536 difference as in equation (09). Assuming that the two distributions are independent, the combined distribution is
537 the product of Q and R . If we now define the variable $z=v*DT$, the new combined distribution on the heat flux
538 variable z is given by the Mellin transform and convolution, described in Papoulis, A., & Pillai, S. U. (2002).

539



540 Figure A1: Histograms presenting the two original distributions, $Q(v)$ (upper left quadrant, units wind speed) and
541 $R(DT)$ (upper right quadrant, units degrees C) and the combined distribution for SH flux in units of $W m^{-2}$. The
542 parameters used for the two original distributions are: $k = 2.0$ for the Weibull shape, $\lambda = 6.0$ for the Weibull
543 scale; $\alpha = - 2.0$ for the Skew Normal shape, $\mu = 2.0$ for the Skew Normal location and $\omega = 2.5$ for the
544 Skew Normal scale
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APPENDIX B

The statistical moments for the skew-normal PDF are given by:

$$E(x) = \mu + \lambda \delta \sqrt{\frac{2}{\pi}} \quad (\text{B1})$$

$$\sigma^2 = \lambda^2 \left(1 - \frac{2\delta^2}{\pi}\right) \quad (\text{B2})$$

$$\mu_3 = (4 - \pi) \frac{(\delta \sqrt{2/\pi})^3}{2(1 - 2\delta^2/\pi)^{3/2}} \quad (\text{B3})$$

$$\mu_4 = 2(\pi - 3) \frac{\left(\delta \sqrt{\frac{2}{\pi}}\right)^4}{\left(1 - \frac{2\delta^2}{\pi}\right)^2} \quad (\text{B4})$$

where $\delta = \frac{\alpha}{\sqrt{1 + \alpha^2}}$. Since the expected value of the time series is zero, we deduce that:

$$\mu = -\lambda \delta \sqrt{\frac{2}{\pi}} \quad (\text{B5})$$

In other words, location and shape parameters have opposite signs since the scale parameter, λ , is always positive.

SH flux anomaly distribution was analysed with ALD PDF and its' statistical moments are given by:

$$\text{mean} = \mu + \frac{1 - \alpha^2}{\lambda \alpha} \quad (\text{B6})$$

$$\text{variance} = \frac{1 + \alpha^2}{\lambda^2 \alpha^2} \quad (\text{B7})$$

$$\text{Skewness} = \frac{2(1 - \alpha^6)}{(\alpha^4 + 1)^{\frac{3}{2}}} \quad (\text{B8})$$

$$\text{Kurtosis} = \frac{6(1 + \alpha^3)}{(1 + \alpha^4)^2} \quad (\text{B9})$$

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