

## Response to reviewer 1, author comments given in bold text:

### General comment

The paper presents investigations of the dipole-magnetotail transition region by means of global hybrid-Vlasov simulations of Earth's magnetosphere. The present run of the employed Vlasiator code is merged with an ionospheric solver, and the ionospheric field-aligned currents are related to the magnetospheric vorticity, as a proxy to auroral dynamics. The focus of the paper is on a wave-like density structure that appeared in the transition region after magnetotail reconnection. The wave-like structure is formed by earthward flows with ion vortices on their sides. The authors attribute the wave-like structure to ballooning/interchange activity.

**-The authors thank the referee for the comments and suggestions. We will revise the paper according to the comments and have added several new references as suggested by the reviewer. These additional references will give a more thorough view of previous work into similar topics. The manuscript will be improved by the implementation of these suggestions.**

The simulations clearly reveal a development of a Bz/entropy ridge at about  $-10 R_E$  (Figure 4b,f), which is apparently the source of further earthward low-entropy (bubble) intrusions (Figure 4c,g,d,h) due to an interchange process. This is indeed similar to the results of recent global high-resolution (down to  $\sim 300\text{km}$ ) MHD simulations by Sorathia et al., 2020 (10.1029/2020GL088227). At the same, due to multiple differences (e.g. significantly larger scales and velocities of the present low-entropy intrusions), the present simulation better matches the Rice Convection Model simulations of sawtooth events by Sazykin et al., 2002 (10.1029/2001GL014416), Yang et al., 2008 (10.1029/2008JA013635) and Sun et al., 2021 (10.1029/2021GL094097), where interchange instability operates during storms or substorm, unlike quiet growth phase in simulations of Sorathia et al., 2020 (10.1029/2020GL088227). The RCM simulations show that a wide injection boundary around the geosynchronous orbit may break up into multiple injection channels with the local time separation of about 1–2 h, similarly to present simulations.

**Thank you for these new references, we had not come across the RCM studies previously, but they are very relevant to this study, and we will add them to the manuscript. There are indeed several differences between our study and the Sorathia et al. 2020 study. While the instability appears in some ways to be similar in our work and Sorathia's, the onset mechanism is different. In the Sorathia 2020 paper the**

instability is triggered in the growth phase of a synthetic substorm (prior to any reconnection in the tail), while in our case the instability is closely related to the large-scale onset of reconnection in the magnetotail. In both cases the instability is governed by changes in entropy and  $B_z$ , but otherwise the mechanism is different.

Our results are more similar to previous MHD studies(e.g. Lapenta et al., 2011 <https://doi.org/10.1029/2011GL047742>), and RCM studies (e.g. Yu et al., 2017 <https://doi.org/10.1002/2017JA024168> and the new papers suggested by the referee). We will expand on this in the revised manuscript, and clarify the differences between our work and the Sorathia study.

Even more so, the authors attribute the appearance of the interchange-unstable magnetotail configuration ( $B_z$ /entropy ridge at  $-10 R_E$ ) to reconnection and loss of density via plasmoid release, which would be similar to the results of Birn et al., 2011 (10.1029/2010JA016083). This is also a different mechanism, as opposed to the mechanism that is based on flux return to the dayside (Hsieh and Otto, 2015, 10.1002/2014JA020925), which was identified to operate in the run of Sorathia et al., 2020 (10.1029/2020GL088227).

We will clarify the differences and similarities between our results and previous work in the revised manuscript. Our interpretation is that the onset of reconnection, similarly to Birn et al., 2011 results in the lowering of entropy, and eventually a plasmoid release. In our simulation this low entropy region spreads over the magnetotail (as the reconnection spreads over the tail), resulting in the region becoming unstable to the ballooning/interchange instability. Dayside flux depletion also occurs in our simulation (Tao et al., 2025), contributing to the tail current sheet thinning. The entropy depletion in the tail coincides with the low density, fast Earthward flow regions caused by reconnection in the tail. Thus we attribute the entropy depletion and the instability to the nightside dynamics.

Thus we have a combination of both BBF-like dynamics, and the ballooning/interchange instability. As we model the ion-kinetic physics of the magnetotail, we can self-consistently capture these phenomena, and go beyond the MHD and RCM descriptions.

Tao, S., Alho, M., Zaitsev, I., Turc, L., Battarbee, M., Ganse, U., Pfau-Kempf, Y., and Palmroth, M.: Magnetospheric convection in a hybrid-Vlasov simulation, EGU sphere [preprint], <https://doi.org/10.5194/egusphere-2025-1340>, 2025.

The above major points need to be carefully addressed before publication of the paper. In addition to them I also list below a number of minor suggestions, which may help improve the paper.

#### Specific comments

A clarifying comment on what leads to reconnection triggering in the Vlasator would be useful.

**We will comment on this in the manuscript. The reconnection is driven both by numerical diffusion and a tearing-type instability.**

Line 31: reference to Sitnov may not be the best one here, and some auroral paper could be cited instead.

**We will change this reference to Partamies et al., 2015**  
(<https://doi.org/10.1002/2015JA021217>)

Line 32: Additional reference could be added here:

Baumjohann, W., G. Paschmann, and H. Lühr (1990), Characteristics of High-Speed Ion Flows in the Plasma Sheet, *J. Geophys. Res.*, 95, 3801–3809

Line 34: Additional references could be added here:

Baumjohann, W., Hesse, M., Kokubun, S., Mukai, T., Nagai, T., & Petrukovich, A. A. (1999). Substorm dipolarization and recovery. *Journal of Geophysical Research*, 104, 24995–25000.

Baumjohann, W. (2002), Modes of convection in the magnetotail, *Phys. Plasmas*, 9, 3665–3667, doi:10.1063/1.1499116

Ohtani, S., Singer, H. J., & Mukai, T. (2006). Effects of the fast plasma sheet flow on the geosynchronous magnetic configuration: Geotail and GOES coordinated study. *Journal of Geophysical Research*, 111, A01204. <https://doi.org/10.1029/2005JA011383>

Merkin, V. G., Panov, E. V., Sorathia, K., & Ukhorskiy, A. Y. (2019). Contribution of bursty bulk flows to the global dipolarization of the magnetotail during an isolated substorm. *Journal of Geophysical Research: Space Physics*, 124, 8647–8668. <https://doi.org/10.1029/2019JA026872>

Line 35: Additional references could be added here:

Angelopoulos, V., et al. (1996), Multipoint analysis of a bursty bulk flow event on April 11, 1985, *J. Geophys. Res.*, 101, 4967–4989.

Sergeev, V. A., V. Angelopoulos, J. T. Gosling, C. A. Cattell, and C. T. Russell (1996), Detection of localized, plasma-depleted flux tubes or bubbles in the midtail plasma sheet, *J. Geophys. Res.*, 101, 10,817– 10,826, doi:10.1029/96JA00460

Line 37: Additional reference could be added here:

Nakamura, R., Baumjohann, W., Klecker, B., Bogdanova, Y., Balogh, A., Rème, H., Bosqued, J. M., Dandouras, I., Sauvaud, J. A., Glassmeier, K.-H., Kistler, L., Mouikis, C., Zhang, T. L., Eichelberger, H., and Runov, A. (2002). Motion of the dipolarization front during a flow burst event observed by Cluster. *Geophys. Res. Lett.*, 29:1942

Line 39: Additional reference could be added here:

Shiokawa, K., W. Baumjohann, and G. Haerendel (1997), Braking of highspeed flows in the near-Earth tail, *Geophys. Res. Lett.*, 24, 1179–1182, doi:10.1029/97GL01062.

Line 40: Additional reference could be added here:

Ohtani, S., Y. Miyashita, H. Singer, and T. Mukai (2009), Tailward flows with positive B<sub>Z</sub> in the near-Earth plasma sheet, *J. Geophys. Res.*, 114, A06218, doi:10.1029/2009JA014159.

Panov, E. V., et al. (2010), Plasma sheet thickness during a bursty bulk flow reversal, *J. Geophys. Res.*, 115, A05213,

doi:10.1029/2009JA014743.

The reference to Panov, E. V., et al. (2010) on Multiple overshoot and rebound of a bursty bulk flow (10.1029/2009GL041971) belongs together with Birn et al., 2011.

Also, the following two references could be placed next to Birn et al., 2011 in this line.

Keika, K., et al. (2009), Observations of plasma vortices in the vicinity of flow-braking: A case study, *Ann. Geophys.*, 27, 3009–3017.

Keiling, A., et al. (2009), Substorm current wedge driven by plasma flow vortices: THEMIS observations, *J. Geophys. Res.*, 114, A00C22,

doi:10.1029/2009JA014114.

Line 47: Additional reference could be added here:

Baumjohann, W., Pellinen, R. J., Opgenoorth, H. J., & Nielsen, E. (1981). Joint two-dimensional observations of ground magnetic and ionospheric electric fields associated with auroral zone currents—Current systems associated with local auroral break-ups. *Planetary and Space Science*, 29, 431–435.

Birn, J., & Hesse, M. (2014). The substorm current wedge: Further insights from MHD simulations. *Journal of Geophysical Research: Space Physics*, 119, 3503–3513.

<https://doi.org/10.1002/2014JA019863>

McPherron, R. L., Nakamura, R., Kokubun, S., Kamide, Y., Shiokawa, K., Yumoto, K., Mukai, T., Saito, Y., Hayashi, K., Nagai, T., Ables, S., Baker, D. N., Friis-Christensen, E., Fraser, B., Hughes, T., Reeves, G., & Singer, H. (1997). Fields and flows at GEOTAIL during a moderate substorm. *Advances in Space Research*, 20, 923–931.

Palin, L., Opgenoorth, H. J., Ågren, K., Zivkovic, T., Sergeev, V. A., Kubyshkina, M. V., Nikolaev, A., Kauristie, K., Kamp, M., Amm, O., Milan, S. E., Imber, S. M., Facskó, G., Palmroth, M., & Nakamura, R. (2016). Modulation of the substorm current wedge by bursty bulk flows: 8 September 2002—Revisited. *Journal of Geophysical Research: Space Physics*, 121, 4466–4482. <https://doi.org/10.1002/2015JA022262>

Panov, E. V., Baumjohann, W., Nakamura, R., Weygand, J. M., Giles, B. L., Russell, C. T., et al. (2019). Continent-wide R1/R2 current system and ohmic losses by broad dipolarization-

injection fronts. *Journal of Geophysical Research: Space Physics*, 124, 4064–4082.  
<https://doi.org/10.1029/2019JA026521>

Sergeev, V. A., Sauvaud, J.-A., Popescu, D., Kovrazhkin, R. A., Liou, K., Newell, P. T., Brittnacher, M., Parks, G., Nakamura, R., Mukai, T., & Reeves, G. D. (2000). Multiple-spacecraft observation of a narrow transient plasma jet in the Earth's plasma sheet. *Geophysical Research Letters*, 27, 851–854.

Line 81: Additional reference could be added here:

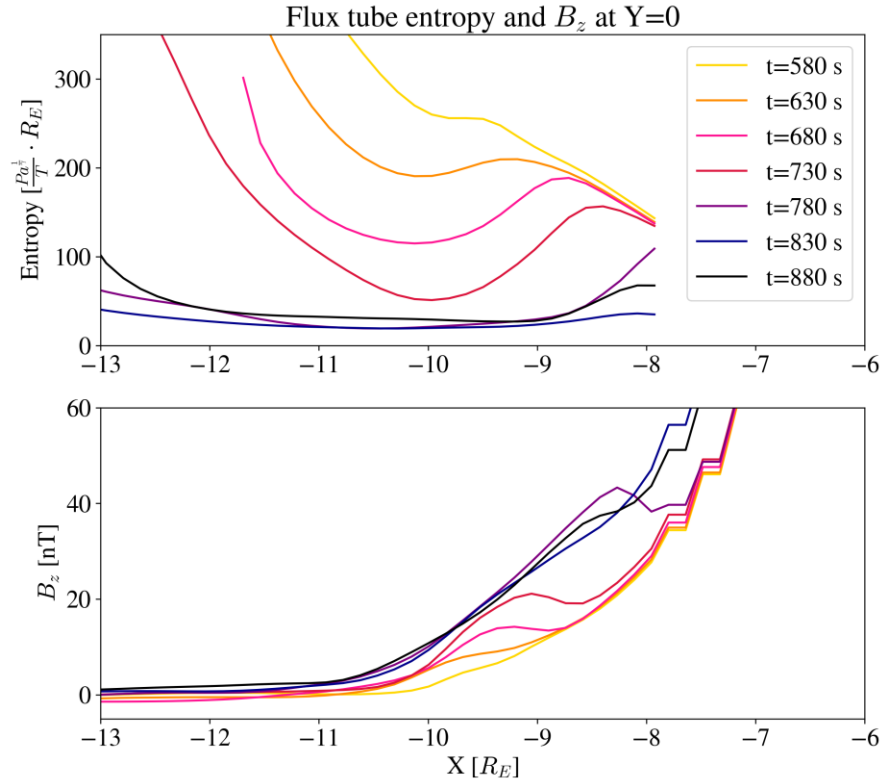
Pritchett, P. L., F. V. Coroniti, and Y. Nishimura (2014), The kinetic ballooning/interchange instability as a source of dipolarization fronts and auroral streamers, *J. Geophys. Res. Space Physics*, 119, 4723–4739, doi:10.1002/2014JA019890.

**We thank the referee for the new references, we will review them and add relevant references at the appropriate lines. The manuscript will benefit from this thorough background information.**

Line 222: Could specific time be indicated after “At the Earthward flows“?

**We will add a time here, it will be good to clarify that.**

Figure 4: A plot with the time evolution of the radial profiles of  $B_z/PV^\gamma$  could be shown here for the times around  $t=680$  s. This plot would show the growth/formation of the  $B_z$ /entropy ridge.



Indeed such a figure will be useful. We give the figure above, will add it to the manuscript and explain it further in the text. In the figure, we show the radial profiles of entropy and  $B_z$  every 50 seconds from 580 s to 880 s. There is a clear decrease in flux tube entropy, and a corresponding increase in  $B_z$  between  $t=580$  s and  $t=780$  s. After this, there is a slight increase in entropy, and a decrease in  $B_z$ , which is due to the wavy interchange structure increasing in azimuthal size, so that a higher entropy region moves to  $Y=0$  by 880 seconds (see Fig. 4g and Fig. 4h)

Figure 6 and associated text: Could the authors explain somewhere how the FAC was obtained?

**The FACs are determined from the curl of  $\mathbf{B}$  close to the inner boundary of the simulation. We will add this to the manuscript.**

Line 322: Midnight may be more appropriate as mid-tail sounds ambiguous when one considers radial distance instead of azimuthal.

**Midnight is indeed the better word to use here, we will make the change.**

Technical corrections

Line 117: It seems that in is missing between done and six dimensions.

**Yes, 'in' is missing from that sentence, that will be fixed.**