

<sup>1</sup> **Setting up the physical principles of resilience in a model of the  
2 Earth System**

<sup>3</sup> Orfeu Bertolami<sup>1,\*</sup> and Magnus Nyström<sup>2,†</sup>

<sup>4</sup> *<sup>1</sup>Departamento de Física e Astronomia,*

<sup>5</sup> *Faculdade de Ciências da Universidade do Porto,*

<sup>6</sup> *Rua do Campo Alegre s/n, 4169-007, Porto,*

<sup>7</sup> *Portugal; Centro de Física das Universidades do Minho e do Porto,*

<sup>8</sup> *Rua do Campo Alegre s/n, 4169-007, Porto, Portugal.*

<sup>9</sup> *<sup>2</sup>Stockholm Resilience Centre, University of Stockholm,*

<sup>10</sup> *Albanovägen 28, 106 91, Stockholm, Sweden.*

<sup>11</sup> (Dated: August 14, 2025)

## Abstract

Resilience is a property of social, ecological, social-ecological and biophysical systems. It describes the capacity of a system to cope with, adapt to and innovate in response to a changing surrounding. Given the current climate change crisis, ensuring conditions for a sustainable future for the habitability on the planet is fundamentally dependent on Earth System (ES) resilience. It is thus particularly relevant to establish a model that captures and frames resilience of the ES, **most particularly in physical terms that can be altered by the restoration of ecosystems, adaptation, mitigation and other strategies.** In this work we propose that resilience can serve as a theoretical foundation when unpacking and describing metastable states of equilibrium and energy dissipation in any dynamic description of the variables that characterise the ES. Since the impact of the human activities can be suitably gauged by the planetary boundaries (PBs) and the planet's temperature is the net result of multiple PBs interactions, such as  $CO_2$  concentration and radiative forcing, atmospheric aerosol loading, atmospheric ozone depletion, etc, then resilience features arise once conditions to avoid **the ES runaway to a state where the average temperature is much higher than the current one.** In this work it is shown that this runaway can be provided by the presence of metastable states and dynamic friction built out of the interaction among the PB variables **once suitable conditions are satisfied. In this work these conditions are specified.** As humanity moves away from Holocene conditions, we argue that resilience features arising from metastable states might be crucial for the ES to follow sustainable trajectories in the Anthropocene **that prevent it run into a much hotter equilibrium state.**

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<sup>\*</sup>Electronic address: [orfeu.bertolami@fc.up.pt](mailto:orfeu.bertolami@fc.up.pt)

<sup>†</sup>Electronic address: [Magnus.Nystrom@su.se](mailto:Magnus.Nystrom@su.se)

12 **I. INTRODUCTION**

13 Over the past decades the human imprint on the Earth System (ES) has been exceptional  
14 (Jouffray et al. 2020, Steffen et al. 2015a). While the mass of humans is only about 0.01% of  
15 the total biomass, we have become a dominant force in shaping the face of Earth, including  
16 its atmosphere, biosphere, hydrosphere and lithosphere (Crutzen 2002, Ellis 2011, Foley  
17 2011, Nyström et al. 2019, Vitousek et al. 1997), and as of 2020 the global human-made  
18 mass surpasses the dry-weight of all living biomass (Elhacham et al. 2020). Thus, humans  
19 have become a hyper-keystone species (Worm & Paine 2016), which rivals geological forces  
20 in influencing the trajectory of the ES (Steffen et al. 2018)

21 A major concern of these changes is the risk of crossing of so-called tipping-points, which  
22 refer to the critical threshold at which a small change or event triggers a significant and  
23 potentially irreversible (regime) shift in a system (Lenton et al. 2008). Tipping-points have  
24 been observed in various systems, such as ecosystems (e.g. food webs, benthic communi-  
25 ties), social systems (e.g. norms, policy), economic systems (e.g. market-based economy)  
26 and technological systems (e.g. steam engine, smartphone, artificial intelligence) (Nyborg  
27 et al. 2016, Scheffer 2009, Scheffer et al. 2001). Over the past couple of decades there have  
28 been raising concerns around the existence of tipping-elements, which are large-scale com-  
29 ponents (subsystems) of the ES that may transgress a tipping-point (Barnosky et al. 2012,  
30 Lenton et al. 2008). Example of such tipping-elements include, the Greenland Ice Sheet, the  
31 Atlantic Meridional Overturning Circulation (AMOC), permafrost, monsoon systems, and  
32 the Amazon rainforest. Importantly, these tipping-elements interact, which **may** lead to a  
33 cascading behaviour of the entire ES (Wunderling et al. 2024). The consequences of these  
34 dynamics for humanity **could** be colossal (Steffen et al. 2018).

35 Clearly, knowledge about tipping-points, where they are located, when they are ap-  
36 proached and identifying ways to navigate away from them, are key challenges for humanity  
37 (Barnosky et al. 2012, Scheffer et al. 2012). Two broad frameworks that could help assist  
38 in this regard are planetary boundaries and resilience theory. The two are complementary  
39 in the sense that the planetary boundaries provide a quantitative assessment whereas the  
40 resilience framework adds a strong theoretical underpinning.

41 The planetary boundaries (PB) framework (Richardson et al. 2023, Rockström et al. 2009,  
42 Steffen et al. 2015b) has been used to define global and regional limits in biophysical processes

43 – ‘safe operating space’ – that must not be crossed if humanity is to stay away from systemic  
44 and potentially irreversible shifts in the ES. As such, the planetary boundaries framework  
45 serves as a ”global dashboard,” tracking humanity’s collective impact on key environmental  
46 factors that threaten the Earth’s ability to sustain human life. More recently, focus has  
47 been directed towards exploration of how different boundaries can interact and potentially  
48 cascade, thereby shrinking the safe operating space for future human impacts on the ES  
49 ([Lade et al. 2020](#)).

50 The resilience concept describes the extent to which a system can resist and develop (e.g.  
51 ecosystems or the the entire ES) with change by absorbing recurrent perturbations, deal with  
52 uncertainty and risk, and still sustain its key properties ([Folke 2006](#), [Holling 2001](#)). It links  
53 to the planetary boundaries framework **as the latter can signal** the existence of tipping-  
54 points (or thresholds), multiple states (or regimes) and self-reinforcing feedback mechanisms  
55 (i.e. hysteresis). Resilience has also been suggested as a conceptual framework that could  
56 assist in developing paths towards sustainability ([Folke et al. 2016](#)). Hence, it can serve as a  
57 theoretical and practical foundation for the planetary boundaries framework. An important  
58 point to bear in mind, however is that resilience is a property of a system and is neither  
59 ”good” nor ”bad” per se. It can help maintain the current state of a system no matter  
60 whether it is deemed desirable or undesirable. **However, the Holocene epoch has been**  
61 **marked by an unusually stable climate compared to previous geological periods.**  
62 **This has allowed for the development of agriculture, permanent settlements,**  
63 **and the emergence of complex human societies. Hence, from this perspective, a**  
64 **Holocene(-like) state can be deemed desirable, and the safeguarding of resilience**  
65 **of this state of critical importance for humanity.**

66 Bearing in mind the resilience concept and its importance we aim in this work to specify,  
67 in the context of a thermodynamical model of the ES, what are the physical properties  
68 that manifest themselves collectively as resilience features of the ES. Our starting point is  
69 a thermodynamical model of the ES from Holocene state conditions to other potentially  
70 stable states, which can be regarded as phase transitions and admit a description through  
71 the Landau-Ginzburg Theory (LGT) ([Barbosa et al. 2020](#), [Bertolami & Francisco 2018](#),  
72 [2019](#)). The LGT is a theoretical framework used in physics to describe phase transitions,  
73 such as when a material changes from a solid to a liquid state or a magnetic material loses  
74 its magnetism. Here we use the LGT to describe the transitions the ES has gone throughout

75 the history of Earth.

76 As we shall review in the next section, this framework allows for determining the equilib-  
77 rium states of the ES in terms of the planet's biophysical subsystems or processes that are,  
78 due to the impact of the human activities, the driving forces that dominate its evolution. In  
79 the Anthropocene, here collectively denoted by  $H$ . In the phase-transition model discussed  
80 in Refs. (Barbosa et al. 2020, Bertolami & Francisco 2018, 2019),  $H$  was considered an exter-  
81 nal field, however, in the present work, we admit that **through large scale restoration of**  
82 **ecosystems, adaptation, mitigation and geo-engineering, the dynamic features**  
83 **of the ES can be altered so to modify the topographic landscape of possible**  
84 **Anthropocene trajectories.**

85 As previously discussed, the proposed Landau-Ginzburg model allows for getting the  
86 evolution equation of the ES, the so-called Anthropocene equation, and to associate the sharp  
87 rise of the physical parameters that characterise the ES to the great acceleration of the human  
88 activities (Bertolami & Francisco 2018), which became conspicuous from the second half of  
89 the 20th century and onwards (Steffen et al. 2015a). **However, the original model did**  
90 **not exhibit explicit features that resemble resilience.** This is the main purpose  
91 **of the present work.** As the model is based on thermodynamical arguments, one  
92 must seek for physical properties that would lead to a more resilient behaviour  
93 **of the ES.** In the context of the model, resilience is regarded as the resistance  
94 the ES shows in changing from one equilibrium state to another. At the present  
95 transient period, the Anthropocene, one infers from a multitude of observations  
96 that the ES is moving away from the Holocene equilibrium state to a new state,  
97 most likely a Hothouse Earth state (Steffen et al. 2018). As we shall see, our  
98 results show that resilience is associated to the existence of metastable states  
99 and explicit dissipation of energy that prevent the ES to runaway towards the  
100 **Hothouse Earth state.**

101 A pleasing feature of the proposed description is that it allows for drawing trajectories of  
102 the ES in the phase space of model's variables. By considering that the PBs and the ensued  
103 temperature display dynamics **that are affected by PBs self-interactions which are**  
104 **shown to be non-vanishing** (Barbosa et al. 2020), two well defined and distinct sets of  
105 trajectories were identified upon assumptions about the evolution of the PB: a linear growth  
106 of the human activities,  $H(T) = H_0 t$ , where  $H_0$  is an arbitrary constant, from which follows

107 that all ES trajectories starting at the Holocene are led to Hothouse Earth state (Steffen  
108 et al. 2018) with a necessarily higher temperature than the Holocene average temperature  
109 (Bertolami & Francisco 2019); if instead, the increase of the human activities impact on the  
110 ES obey a discrete logistic map (Jakobson 1981, Kingsland 1995, May 1976), trajectories  
111 can display bifurcations or chaotic behaviour (Bernardini et al. 2025). Of course, as human  
112 activities are bounded by the finiteness of resources, the logistic map might be a more  
113 accurate description of its behaviour, although it is not quite clear what is the time span  
114 elapsed between successive steps of the logistic map. In any case, it is relevant to keep in  
115 mind that a too fast increase might give origin to trajectory bifurcations or even chaotic  
116 behaviour, which, of course, precludes predictions and control measures on the evolution of  
117 the ES.

118 In this work we extend the previous studies of the ES model carried out in Refs. (Bar-  
119 bosa et al. 2020, Bernardini et al. 2025, Bertolami & Francisco 2018, 2019) on various  
120 aspects. More fundamentally, we consider **the dynamic features arising from the self-**  
121 **interactions of the 9 identified PBs, here generically denoted as  $h_i$ ,  $i = 1, \dots, 9$ ,**  
122 **and show the specific conditions to implement resilience in the the eleven di-**  
123 **mensional space  $(\psi, h_i, F(\psi, h_i))$ .** Resilience can be regarded as a set of measures that  
124 prevent or delay the evolution of the ES towards a Hothouse Earth state and ensuring that  
125 this state is as close as possible to the Holocene state<sup>1</sup>. This can be implemented by creating  
126 metastable states to avoid a runaway situation due to a barrier that arises as higher-order  
127 terms into the Helmholtz free energy are introduced (cf. discussion below). A further re-  
128 quirement is dynamic friction to restrict the change of state in the phase space. **This is**  
129 **a fairly natural condition as any realistic system dissipates energy.** The specific  
130 conditions for the ES to acquire effective resilience features will be discussed below. Trajec-  
131 tories of the ES without and with resilience are depicted in Figs. 1 and 2 respectively (cf. a  
132 detailed discussion below).

133 This paper is organised as follows: in section II we review the cardinal aspects of the  
134 LGT of the ES and discuss the most relevant features of the dynamical system emerging  
135 from the model; in section III, we discuss the implementation of the resilience features in  
136 the model and connect them to properties that any model of the ES should have. Finally, in  
137 section IV we present our conclusions and discuss how our work can be extended to address  
138 several issues concerning features and transformation of the global social-ecological system.

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<sup>1</sup> Notice that prior the Anthropocene, the equilibrium states of the ES correspond to cooler (glaciation) and hotter (Hothouse Earth) equilibrium states with respect to the Holocene. However, at the Anthropocene, human activities lead inevitably the ES towards a Hothouse Earth state due to the massive emission of greenhouse gases. This materialises in the minus sign of the linear term in Eq. (1) below.

<sup>139</sup> The social-ecological system (SES) framework (Berkes & Folke 1998) builds on  
<sup>140</sup> the notion that nature is no longer simply the backdrop for social interactions,  
<sup>141</sup> just as humans are not merely external forces acting upon ecosystems (Folke  
<sup>142</sup> et al. 2011). Instead, social-ecological systems represent fully integrated, inter-  
<sup>143</sup> dependent systems, where tightly linked feedbacks between social and ecological  
<sup>144</sup> components shape their overall behaviour and dynamics (Biggs et al. 2012).

## <sup>145</sup> II. A THERMODYNAMICAL MODEL FOR THE EARTH SYSTEM

<sup>146</sup> We first review the main features of the proposed model for the ES (Bertolami & Francisco  
<sup>147</sup> 2018) and discuss in the next section the conditions to extend it in order to explicitly exhibit  
<sup>148</sup> resilient properties.

The proposal of Ref. (Bertolami & Francisco 2018) is to regard transitions of the ES as phase transitions which can be described by the LGT through an order parameter,  $\psi$ , and natural parameters (astronomical, geophysical, internal). In the Anthropocene, the natural forces average out to zero and the system is driven by the strength of the human activities, collectively denoted by  $H$ . In this approach, the thermodynamic description of the system is obtained through the Helmholtz free energy,  $F$ , which can be written as an analytic function of an order parameter,  $\psi$ , which is chosen to be the reduced temperature relative to Holocene average temperature,  $\langle T_H \rangle$ ,  $\psi := (T - \langle T_H \rangle) / \langle T_H \rangle$ . Thus, in the Anthropocene, disregarding the spatial variation of  $\psi$ , one can write (Bertolami & Francisco 2018, 2019):

$$F(\psi, H) = F_0 + a\psi^2 + b\psi^4 - \gamma H\psi, \quad (1)$$

<sup>149</sup> where  $F_0$ ,  $a$ ,  $b$  and  $\gamma$  are constants. The linear term in  $\psi$  corresponds to the hu-  
<sup>150</sup> man activities, which at the Anthropocene can match the quadratic and quartic  
<sup>151</sup> contributions due to natural causes (astronomic, geological internal).

The strength of the human activities are probed by their impact via the PBs (Rockström et al. 2009, Steffen et al. 2015b),  $h_i$ ,  $i = 1, 2, \dots, 9$  with respect to their Holocene values. Given that the PB can interact among themselves, the most general expression for  $H$  is given by (Bertolami & Francisco 2019):

$$H = \sum_{i=1}^9 h_i + \sum_{i,j=1}^9 g_{ij} h_i h_j + \sum_{i,j,k=1}^9 \alpha_{ijk} h_i h_j h_k + \dots, \quad (2)$$

152 where  $[g_{ij}]$  is a non-degenerate,  $\det[g_{ij}] \neq 0$   $9 \times 9$  matrix. Similar conditions should be  
 153 imposed on the coefficients  $\alpha_{ijk}$  and  $\beta_{ijkl}$  of the higher-order interaction terms. In principle,  
 154 these interactions terms are sub-dominating, however, their importance has to be established  
 155 empirically. As pointed out in Ref. (Bertolami & Francisco 2019), the interaction terms  
 156 may lead to new equilibrium states and suggest some mitigation strategies depending on  
 157 their sign and strength in the matrix entries (Bertolami & Francisco 2019). This will be  
 158 explicitly discussed in the next section. In Ref. (Barbosa et al. 2020), it was shown that the  
 159 interaction term between the climate change variable ( $CO_2$  concentration), say,  $h_1$ , and the  
 160 oceans acidity, say,  $h_2$ , was non-vanishing and contributed to about 10% of the value of the  
 161 individual contributions themselves.

162 **In order to introduce resilience features into the model, that is, resistance  
 163 to change from one equilibrium state into another**, we have to consider, contrary to  
 164 previous works (Barbosa et al. 2020, Bernardini et al. 2025, Bertolami & Francisco 2018,  
 165 2019), that the PBs are dynamical variables. **This allows us to project how the ES  
 166 would behave depending on its initial state and subsequent trajectory in the  
 167 phase space of the model, specified through the variables**  $(\psi, \dot{\psi}, h_i, \dot{h}_i)$ . Thus, for a  
 168 given set of initial conditions, corresponding to a state  $(\psi(0), \dot{\psi}(0), h_i(0), \dot{h}_i(0))$  in the phase  
 169 space, one can, in principle, obtain the trajectories, *orbits*, in the phase space after solving  
 170 the initial value problem through the evolution equations of the system. The equations  
 171 of motion are obtained through the Lagrangian or equivalently through the Hamiltonian  
 172 formalism. The latter, yielding to first order differential equations, is more suitable to  
 173 establish a dynamical system in its canonical form.

174 The Lagrangian function must include, besides the potential, which is given by the free  
 175 energy, a set of kinetic energy terms for the canonical coordinates. The simplest possible ki-  
 176 netic term is a quadratic term proportional to the squared first derivative of each coordinate.  
 177 Thus, we can write the following Lagrangian:

$$\mathcal{L}(q, \dot{q}) = \frac{\mu}{2} \dot{\psi}^2 + \frac{\nu}{2} \sum_{i=1}^9 \dot{h}_i^2 - F_0 - a\psi^2 - b\psi^4 + \gamma H\psi, \quad (3)$$

178 where  $\mu$  and  $\nu$  are arbitrary constants and the dots stand for time derivatives. The constant  
 179  $\nu$  is assumed to be the same for all PB variables.

Aiming to get the Hamiltonian function, we evince the relevant canonical conjugate mo-

menta associated to  $\psi$  and to a generic PB variable,  $h_i$ :

$$p_\psi = \frac{\partial \mathcal{L}}{\partial \dot{\psi}} = \mu \dot{\psi}, \quad (4)$$

$$p_{h_i} = \frac{\partial \mathcal{L}}{\partial \dot{h}_i} = \nu \dot{h}_i, \quad (5)$$

from which follows the Hamiltonian function

$$\mathcal{H}(\psi, p) = \frac{p_\psi^2}{2\mu} + \sum_{i=1}^9 \frac{p_{h_i}^2}{2\nu} + F_0 + a\psi^2 + b\psi^4 - \gamma H\psi, \quad (6)$$

and Hamilton's equations,

$$\dot{\psi} = \frac{\partial \mathcal{H}}{\partial p_\psi}, \quad p_\psi = -\frac{\partial \mathcal{H}}{\partial \psi}, \quad (7)$$

$$\dot{h}_i = \frac{\partial \mathcal{H}}{\partial p_{h_i}}, \quad p_{h_i} = -\frac{\partial \mathcal{H}}{\partial h_i}. \quad (8)$$

The equations of motion read, considering for while just the contribution from the lowest order terms in Eq. (2):

$$\mu \ddot{\psi} = -2a\psi - 4b\psi^3 + \gamma H \quad (9)$$

and

$$\nu \ddot{h}_i = \gamma \psi. \quad (10)$$

To exemplify the behaviour of variables  $\psi$  and  $h_i$ , let us obtain the resulting solutions for the simple case considered in Ref. (Bertolami & Francisco 2019). For  $b \simeq 0$ , we can neglect the cubic term in the equation of motion for  $\psi$  to get the equation of an harmonic oscillator under the action of an external force,  $H(t)$ . This yields for the simple case of an initial linear time evolution,

$$H(t) = H_0 t, \quad (11)$$

for an equilibrium initial state,  $\dot{\psi}(0) = 0$ , the analytical solution:

$$\psi(t) = \psi_0 \cos(\omega t) + \alpha t, \quad (12)$$

where  $\omega = \sqrt{2a/\mu}$  is an angular frequency,  $\alpha = \gamma H_0/2a$  and  $\psi_0$  is an arbitrary constant fixed by the initial conditions.

The solution for the impact on the PB,  $h_i(t)$ , which initially behaves collectively as Eq. (11), that is  $\sum_{i=1}^9 h_i(t \simeq 0) = H_0$ , quickly evolves to a cubic growth in time:

$$h_i(t) = A \cos(\omega t) + B t^3 + \alpha_i t, \quad (13)$$

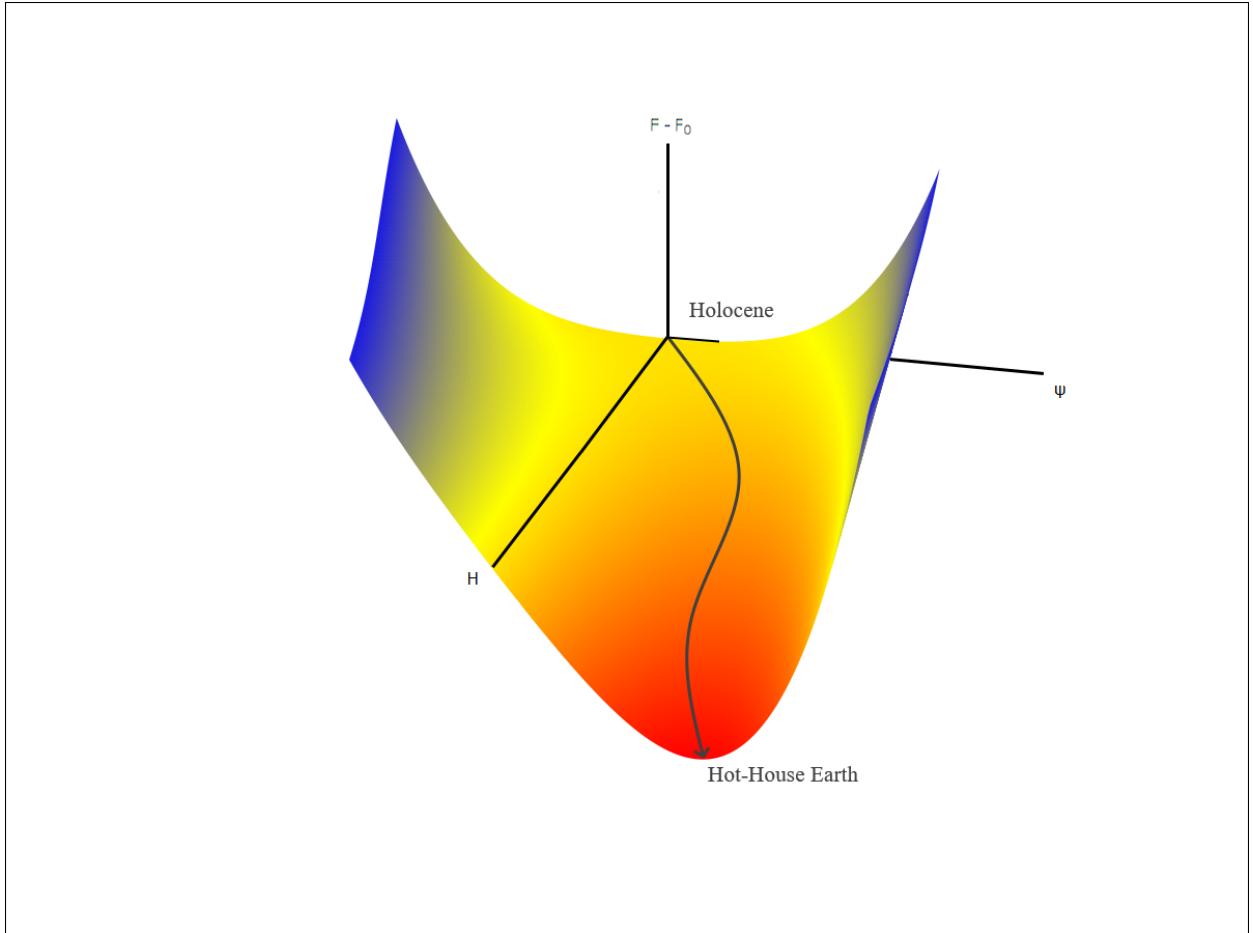


FIG. 1: Free energy in function of the temperature and of the intensity of the human impact on the PBs.

182 where  $A = -\gamma\psi_0/\nu\omega^2$ ,  $B = \alpha\gamma/6\nu$ , for an arbitrary  $\alpha_i$ .

183 These solutions clearly show that if the temperature  $\psi$  grows from an initial linear col-  
 184 lective behaviour of the PBs,  $H$ , then quickly turns the  $h_i$ s to have a cubic growth. Clearly,  
 185 this model shows no resilience features as depicted in Figure 1, where one clearly sees that  
 186 from the Holocene, Anthropocene trajectories inevitably evolve towards a Hothouse Earth  
 187 state.

188 In what follows we shall consider the introduction into the free energy function of a  
 189 cubic term for  $\psi$  and higher than linear order terms for the PBs as these will allow for  
 190 metastable states to arise, thus leading to bounded solutions for  $\psi$  and the PBs. **Metastable**  
 191 **states correspond to intermediate energy states between the Holocene state and**  
 192 **the Hothouse Earth least energy state. In the LGT, metastable states can be**  
 193 **considered and studied through cubic terms in the Helmholtz free energy.** The

194 conditions for the appearance of metastable states were already discussed in a completely  
 195 different context, namely in a proposal to classify rocky planets (Bertolami & Francisco  
 196 2022), using the ideas developed in Refs. (Barbosa et al. 2020, Bertolami & Francisco 2018,  
 197 2019) to describe the ES.

198 Before concluding this discussion it is worth stressing once again that the behaviour of  
 199 the ES depends crucially on the assumptions about the evolution of the PB. Indeed, as  
 200 pointed out in the introduction, the supposition that human activities grow linearly as in  
 201 Eq. (11) implies, as exemplified above, that ES trajectories lead to the "Hothouse Earth"  
 202 state (Bertolami & Francisco 2019) as discussed by Ref. (Steffen et al. 2018). However, if the  
 203 human activities impact on the ES behaves as a discrete logistic map, as suggested in Ref.  
 204 (Bernardini et al. 2025), then trajectories will depend the rate of growth of human activities  
 205 as solutions admit regular trajectories as well as trajectories that present bifurcations and  
 206 even chaotic behaviour. **In the next section we shall consider the features that must**  
 207 **be introduced in the Helmholtz free energy and the conditions they must satisfy**  
 208 **in order to avoid the ES evolves towards the Hothouse Earth state.**

209 **III. SETTING UP THE PHYSICAL PRINCIPLES OF RESILIENCE**

As mentioned above, resilience features are associated to bounded trajectories in the Anthropocene and these ask for the existence of metastable states. In the LGT the metastable states arise by intruding cubic terms on the free energy. As pointed out in Ref. (Bertolami & Francisco 2022), the introduction of a cubic term allows for a richer variety of equilibrium states. Indeed, consider the free energy:

$$F(\psi, H) = F_0 + a\psi^2 - c|\psi|^3 + b\psi^4 - \gamma H\psi, \quad (14)$$

210 where we assume that constants  $b$ ,  $c$  and  $\gamma$  are positive, while constant  $a$  can be negative.

The existence of extrema is given by two conditions. The first one reads:

$$\frac{\partial F(\psi, H)}{\partial \psi} = 0 = 2a\psi - 3c\psi^2 + 4b\psi^3 - \gamma H. \quad (15)$$

211 The resulting cubic equation admits at least one real solution, say,  $\psi_M$ , meaning that there  
 212 are at least two metastable states,  $\psi_M$  and  $-\psi_M$ . Clearly,  $\psi_M \neq 0$  as far as  $H \neq 0$ .

213 However, the unboundedness of the evolution of the variables  $(\psi, h_i)$  is due to the un-  
 214 boundedness of the PBs. Recent assessment of the PBs has shown that 6 out of the 9 PBs

215 have gone beyond their Holocene values where they were at equilibrium, a state usually  
 216 referred to as Safe Operating Space (SOS).

The motion in the eleven-dimensional configuration space,  $(\psi, h_i, F(\psi, h_i))$ , is quite complex, so in order to simplify the analysis we consider one single generic PB,  $h_i$ , and assume that the remaining ones are unchanged<sup>2</sup>. The free energy can be written explicitly in terms of the high order contributions in  $H$  depicted in Eq. (2). Therefore, we get:

$$F(\psi, H) = \hat{F}_0 + a\psi^2 - c|\psi|^3 + b\psi^4 - \gamma(h_i + g_i h_i^2 + b_i h_i^3)\psi, \quad (16)$$

217 where we have aggregated all contributions to the quadratic and cubic terms in  $h_i$ , bf a  
 218 generic PB, within the constants  $g_i$  and  $b_i$ . To ensure boundedness it is necessary that  $g_i$  is  
 219 negative and that  $b_i$  is positive.

Thus, from Eq. (16), one gets the condition:

$$\frac{\partial F(\psi, h_i)}{\partial h_i} = 1 + 2g_i h_i + 3b_i h_i^2 = 0, \quad (17)$$

220 which admits real non-vanishing solutions,  $h_{iM}$ . as far as  $g_i^2 > 3b_i$  for  $b_i \neq 0$  or  $h_{iM} = \frac{1}{-2g_i}$   
 221 if  $b_i = 0$ .

The general conditions to ensure that the extremum  $(\psi_M, h_{iM})$  corresponds to a minimum and hence to a metastable state are given by:

$$\frac{\partial^2 F(\psi_M, h_{iM})}{\partial \psi^2} \frac{\partial^2 F(\psi_M, h_{iM})}{\partial h_i^2} - \left( \frac{\partial^2 F(\psi_M, h_{iM})}{\partial \psi \partial h_i} \right)^2 > 0. \quad (18)$$

and

$$\frac{\partial^2 F(\psi_M, h_{iM})}{\partial \psi^2} > 0, \quad (19)$$

which yield the following relationships:

$$g_i < -3b_i h_{iM} \quad (20)$$

and

$$2a - 6c|\psi_M| + 12b\psi_M^2 > 0. \quad (21)$$

222 Satisfying these conditions imply the ES can settle in a the metastable state,  $(\psi_M, h_{iM})$ ,  
 223 that is, the system shows resilience and does not runaway towards the "Hothouse Earth"  
 224 state as depicted in Figure 2 as far as  $3b_i < g_i^2 < 9b_i^2 h_{iM}$ . .

225 Notice that the conditions for the existence of a metastable state can be met if  $g_i < 0$   
 226 even if coefficients  $b_i$  vanish. This is quite welcome as these coefficients are associated to

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<sup>2</sup> Notice that the analysis of two-variables case is quite relevant as the Kolmogorov-Arnold representation theorem establishes that any continuous function of several variables can be constructed out of a finite sum of two-variable functions.

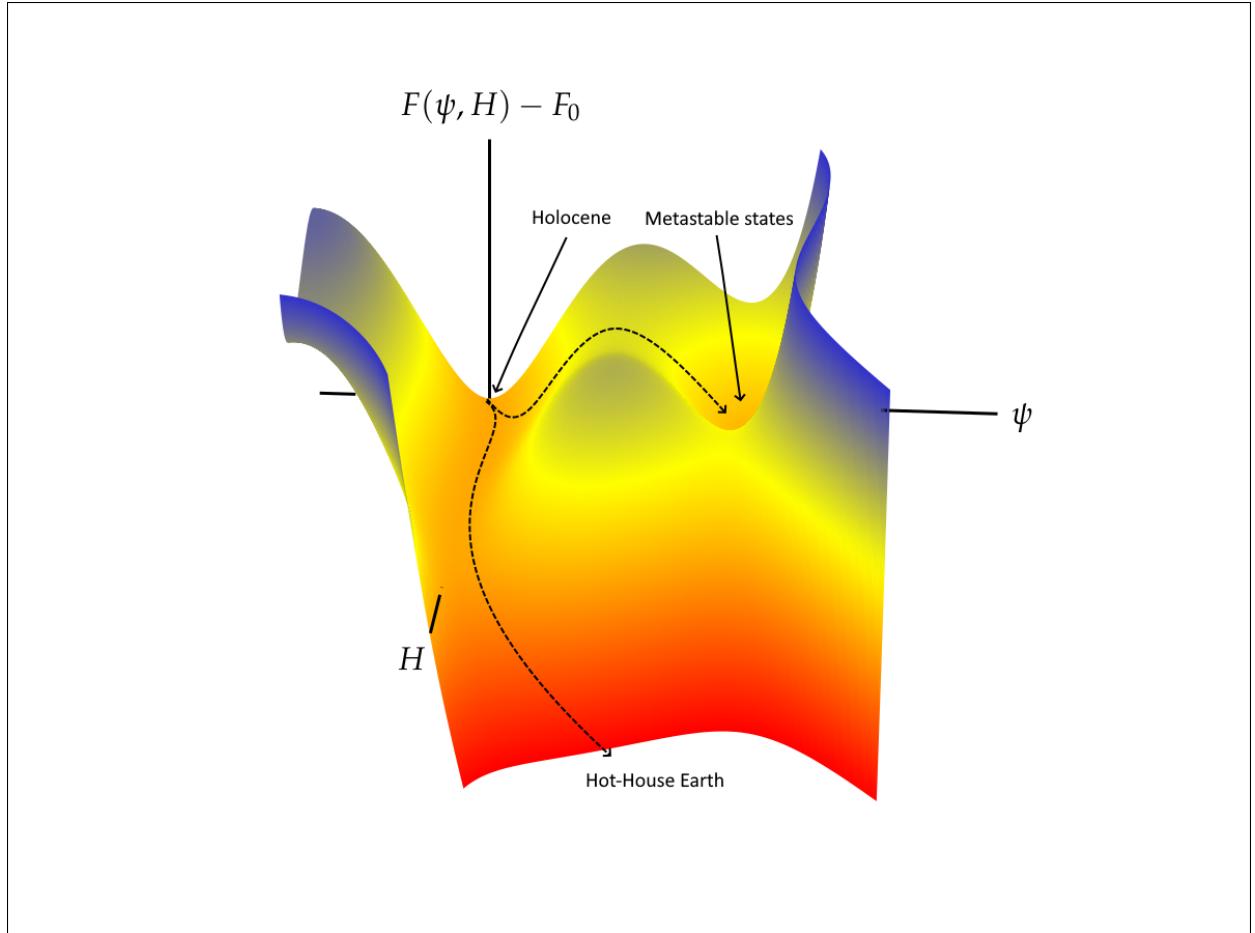


FIG. 2: Free energy in function of the temperature, planetary boundaries ( $H$ ) and resilience features (matastable state).

higher-order interaction terms, which from phenomenological considerations, are presumably small. On the other hand, a non-vanishing and negative contribution from the quadratic term  $h_i^2$  is absolutely necessary. Actually, the concrete case studied in Ref. (Barbosa et al. 2020) shows that this is indeed the case. Furthermore, condition Eq. (21) can be satisfied if  $a < 0$ .

Another feature associated to resilience is the “inertia” that the ES shows in changing from a given state to another. This feature can be identified with the ubiquitous dissipation of energy present in any physical system. Most often dynamical dissipation processes can be described through velocity-proportional frictional forces **which imply that just part of the free energy of a system is turned into kinetic energy, that is, motion of the system.** In the Lagrangian/Hamiltonian formalism for a particle, the effect of these forces can be accounted through the Rayleigh dissipation function,  $R = -\kappa p^2/m$ , where  $\kappa$

239 is a constant,  $p$  is the canonical conjugate momentum and  $m$  the mass of the particle.

240 For the ES, introducing dissipation through the Rayleigh function implies that the left  
241 hand side of the equations of motion (9) and (10) acquire extra terms  $-\kappa_\psi \dot{\psi}$  and  $-\kappa_{h_i} \dot{h}_i$ ,  
242 respectively. The effect of these terms is to reduce the amplitude of the motion of the ES  
243 once it goes from one state to another, thus acting as a resistance of the system to the  
244 change of its state. This can be clearly associated to resilience.

245 These considerations are sufficient for setting the physical conditions for the resilience  
246 of the ES. As we have seen, a metastable state corresponding to the solution  $(\psi_M, h_{iM})$  of  
247 equations (15), (18), and (19), whose free energy (16) coefficients satisfy the conditions (20)  
248 and (21) together with the unavoidable dynamic friction/**energy dissipation** that exists in  
249 any system are the physical properties that endow the ES for having a resilient behaviour.

250 Since the Holocene, the ES has been subjected to a great stress. From the Great Acceler-  
251 ation of the second half of the last century, which presumably sparked the Anthropocene, the  
252 hyper expansion of human activities resulted that the safe operating space has been crossed  
253 for 6 of the 9 PBs (Richardson et al. 2023) and created all sorts of tensions, whose ongoing  
254 climate change crisis is the most persistent consequence for the ES. The tipping of some  
255 of the major ecosystems that compose the ES, such the Amazon rainforest and the Pacific  
256 Coral reefs, are already visible. As to the question of knowing if we have already inflicted  
257 an irreversible damage on ES or are close to it, only the understanding the mechanisms of  
258 resilience **and how their boosting, through the PB interactions, can provide us**  
259 **with a knowledgeable answer. We hope that our work can provide a modest**  
260 **help in this respect.**

#### 261 IV. CONCLUSIONS

262 In this work we have considered the physical principles to ascertain the conditions of re-  
263 silience in a LGT model of the ES. In order to implement resilience features we have endowed  
264 and considered modifications of the free energy so to ensure the existence of metastable  
265 states. Furthermore, we have modelled the ES capability to remain in an equilibrium state  
266 by arguing that it can be suitably prevented to runaway towards the Hothouse Earth state by  
267 the presence of metastable states **whose existence conditions were explicitly shown**  
268 and the unavoidable dissipation of energy during the evolution of the relevant variables.

269     Indeed, we have shown that, thanks to the PBs interactions, a metastable state  $(\psi_M, h_{iM})$   
270     can exist if the conditions, Eqs. (20) and (21), for the coefficients of the free energy, Eq.  
271     (16), are satisfied. As pointed out in the above discussion, these conditions can be satisfied  
272     even if coefficients  $b_i$  vanish as far as  $g_i < 0$ .

273     Based on the observational data, it is possible to infer that the metastable state found  
274     above might correspond either to an actual state that the ES is close to reach or to a  
275     state that can be reached through **large scale restoration of ecosystems, adaptation,**  
276     **mitigation or** engineering measures designed to drive the ES away from the Anthropocene  
277     traps it seems to be currently entangled in (see. Ref. ([Søgaard Jørgensen et al. 2023](#)) for a  
278     description of the 14 major Anthropocene traps).

279     A recent assessment has shown that 6 out of the 9 PBs have been crossed ([Richardson](#)  
280     et al. 2023) meaning that the evolution of most of the PBs is uncontrolled. It is unclear  
281     if the ES has already reached a point of no return, but it is evident that urgent measures  
282     to reverse the current development are needed. In fact, no single set of measures seems  
283     to be sufficient to halt the evolution of the PBs beyond the safe operating state. Two of  
284     the PBs that deserve particular attention are climate change and biosphere integrity. Both  
285     are deemed “core” because their essential role in the ES. The climate system reflects the  
286     distribution and balance of energy at the Earth’s surface, while the controls material and  
287     energy flows, helping to strengthen the systems’s resilience against both rapid and long-term  
288     changes. This calls for a concerted action involving stewardship measures ([Bertolami 2022](#),  
289     [Steffen et al. 2015a, 2011](#)), **bringing into the economy (internalising) the workings of**  
290     **the ES** (see eg. Ref. ([Bertolami 2024](#))) and making them become part of revised economic  
291     paradigms ([Bertolami & Gonçalves 2024, 2025](#), [Sureth et al. 2023](#)), mitigation strategies  
292     that may include technological carbon sequestration (see e.g. ([Bertolami 2025](#), [Bertolami](#)  
293     & [de Matos 2024](#)) and refs. therein), and storage as means to curb climate overshoot, to  
294     avoid irreversible changes to the ES that will compromise the navigation space for the future  
295     generations.

296     **Acknowledgements**

297     O.B. would like to thank the kind hospitality extended to him during his stay at the  
298     Stockholm Resilience Centre in May 2024 where the initial discussions that led to this work

299 took place. He is also grateful for the fruitful discussions on various matters with Lan  
300 Wang-Erlandsson, Peter Søgaard Jørgensen, Dieter Gerten, Ricardo Elísio and Uno Svedin.

301 **V. AUTHOR CONTRIBUTIONS**

302 OB and MN conceptualised the study. OB performed the formula calculations. OB and  
303 MN wrote and edited the paper.

304 **VI. FINANCIAL SUPPORT**

305 The work of Magnus Nyström was supported by the Swedish Research Council grant  
306 (number 2020-04586).

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