

# Response to reviewer R1 comments

August 22, 2025

Thank you for your detailed and helpful review. In this document, reviewer comments are in **black** and our comments are in **red**. New text added to the manuscript is in **blue**.

## 1 Reviewer 1

This is a clearly written paper with nice figures describing nice analysis of an extraordinarily rare and hard to obtain dataset. The manuscript should be published.

Thank you for your positive assessment of our paper, we will address your individual points below

I do have a number of comments, questions and morsels for thought that I list below in the order in which I read. The majority are (very) minor, amounting to text and grammar nits, but some are more substantive. In particular I would like to see

- more supporting evidence behind the claim that mixing is weak (for the reasons given in the final comment below),

We show that the median TKE dissipation rate is  $10^{-11}$  to  $10^{-10}$ , which are very low values, comparable to the background TKE dissipation rate in the ocean (see Figures 6 and 7 in Waterhouse et al. (2014))

terhouse et al. (2014) does not give average values for epsilon, instead we can compare values for kappa. In our study median values of kappa range between  $0.2 \times 10^{-4} m^2 s^{-2} - 1.1 \times 10^{-4} m^2 s^{-2}$ . Waterhouse et al. (2014) gives average deep ocean (depth between 1000 m and the bottom) values of kappa as  $4.3(0.4 - 11.5) \times 10^{-4} m^2 s^{-1}$ , with the values in parenthesis the 95-th percentile bootstrap confidence range. This indicates that our values of kappa lie within the lower range or just below the global distribution for the deep ocean. We will add a reference to global average values of mixing in Waterhouse et al. (2014). Additionally, we will include a new figure showing the distribution of measured epsilon in different ice shelf cavities. Our values lie within the range of previous observations as we have thus rephrased our abstract to remove the reference to “low mixing”, the sentence in question now reads “Rates of background mixing are  $\varepsilon \approx 10^{-10} W kg^{-1}$  with patches of higher mixing of  $\varepsilon \approx 10^{-8} W kg^{-1}$ .”

- better figure 3 and 4, which currently mixes aspect ratios, has the reader going back and forth and does not allow direct comparisons of the most relevant quantities - specifically epsilon and the different instability indicators

We have combined Figures 3 and 4 to the new Figure 3. All panels now have the same aspect ratio.

- quantification of the ADCP vertical wavenumber response and hence justification of the numerical values of Ri presented (or alternately toning down the reference to specific values such as  $Ri = 1/4$  given the estimates are noisy and not fully resolved),

We are unsure what the reviewer is asking about here. Do you refer to the vertical wave number response that Polzin et al. (2002) refer to when estimating turbulent mixing processes from vertical shear in the ADCP? We do not use the ADCP to calculate mixing, we only use it to get information on the horizontal velocity in the vicinity of our microstructure shear measurements. The Richardson number is calcu-

53       lated from the vertical shear between successive 8 m tall (in the vertical)  
54       ADCP bins, but this is not used for the turbulent shear calculations.

- 55       • justification for use of median versus mean

56       We use the median as it is less impacted by outliers or non-normal  
57       distribution of values. If the data is normally distributed the median  
58       and mean are identical, so there is no negative effect of using the median  
59       as the default method for averaging values.

- 60       • and finally and perhaps most substantively, an explanation for why the  
61       turbulent heat fluxes just above the bottom are important to measure.  
62       Ie, is that the water that will eventually meet the grounding line, or  
63       should the study have been done nearer the top of the mCDW water-  
64       mass where the gradients and heat losses are much stronger?

65       Have added information to the text to clarify that we measure heat  
66       fluxes close to the bottom to capture the effect of topography roughness  
67       on the flow, to capture the mixing where the bottom intensified warm  
68       inflow interacts with the seabed and due to practical constraints (the  
69       ALR needs to stay within 100 m of the seabed to allow for accurate  
70       dead-reckoning and bottom tracking). The schematic we added to the  
71       Introduction (see your comment below) should also make the reason for  
72       our interest in the lower mCDW clearer. We have added the following  
73       sentence to the introduction: “our study targets the current of warm  
74       mCDW flowing into the ice shelf cavity and maintains a dive track  
75       close to the seabed. We investigate the circulation and mixing in the  
76       mCDW inflow close to the bed of the cavity to understand the effect  
77       of bathymetry on mixing and circulation. We quantify the upward  
78       heat transport that cools the mCDW in the deepest part of the cavity  
79       whilst warming the overlying mCDW (which can access the grounding  
80       line and the ice shelf base; Figure 1), and investigate drivers for the  
81       observed mixing. ”

82       Good luck. I enjoyed reading the paper and hope that these comments are

83 useful.

84 11: topography, turbulent or both not resolved?

85 We have clarified this sentence to confer that turbulent mixing is not re-  
86 solved in models and topography is not resolved in bathymetry products or  
87 models. The sentence now reads: “We show a highly complex spatial pat-  
88 tern of turbulent mixing and of bottom topography. The bottom topography  
89 is currently not resolved in bathymetry products and both the topography  
90 and turbulent mixing are currently not resolved in models of ice-shelf–ocean  
91 interactions.”

92 26: awkward

93 The sentence now reads “ The mCDW can cause melting at the grounding  
94 line, leading to basal mass loss and grounding line retreat.”

95 35, 53: “this” is a weak reference. Please reword; see Strunk and White if  
96 needed.

97 the sentences now read “ The depth at which meltwater enters the ocean is  
98 influenced by where melt predominantly occurs. ”and “Due to the remote  
99 location and difficult access, measuring turbulent kinetic energy dissipation  
100 rate in ice shelf cavities is only now starting to become feasible.”, respectively.

101 48: Melt rates two words?

102 We have corrected this

103 52: This statement is actually not true: epsilon is the dissipation rate and  
104 further assumptions must be invoked to infer the mixing. This needs to be  
105 corrected and expanded upon.

106 We have clarified that  $\varepsilon$  is only a measure of turbulence if the turbulence is  
107 isotropic. The sentence now reads “The turbulent kinetic energy dissipation  
108 rate,  $\varepsilon$ , is the rate at which molecular viscosity dampens isotropic turbulence  
109 generated at large scales by e.g. vertical or lateral shear, and is used to

110 quantify turbulent mixing.”

111 55: This would be a good place to distinguish what is different about this  
112 study from the other two.

113 Thank you for your comment, the paragraph in question now reads: “To  
114 our knowledge, there exist two published studies of mixing in an ice-shelf  
115 cavity measured by an underwater vehicle, one under Pine Island Glacier  
116 (Kimura et al., 2016), and one under the Filchner Ronne Ice Shelf (Davis  
117 et al., 2022). We present a third such study, targeting DIS. DIS and Pine  
118 Island Ice Shelf experience low tidal flows, whereas Filchner Ronne Ice Shelf  
119 experiences strong tidal flows. Unlike Davis et al. (2022) and Kimura et al.  
120 (2016), our study targets the current of warm mCDW flowing into the ice  
121 shelf cavity and maintains a dive track close to the seabed. We investigate the  
122 circulation and mixing in the mCDW inflow close to the bed of the cavity to  
123 understand the effect of bathymetry on mixing and circulation. We quantify  
124 the upward heat transport that cools the mCDW in the deepest part of the  
125 cavity whilst warming the overlying mCDW (which can access the grounding  
126 line and the ice shelf base; Figure 1), and investigate drivers for the observed  
127 mixing. ”

128 56: which  $\rightarrow$  that. Also, is this the only reason mixing is important to  
129 know for these situations?

130 Thank you, we have corrected that. Mixing at the seabed – ocean interface  
131 is also important for nutrient transport, such as the transport of iron from  
132 sedimentary sources to the euphotic zone. We refer to such processes in  
133 the paragraph above: “The input of meltwater to the Amundsen Sea is also  
134 important for biological activity in the region. The flow of mCDW along the  
135 seafloor on its way into the DIS cavity enriches the mCDW in dissolved iron  
136 and manganese while the meltwater from the ice shelf itself is a source of  
137 particulate iron and manganese (van Manen et al., 2022). The addition of  
138 glacial meltwater makes the outflowing mCDW more buoyant than the dense  
139 mCDW inflow, transporting iron and manganese to the surface ocean (van  
140 Manen et al., 2022) where they are important micronutrients for primary

141 producers (Twining & Baines, 2013).”

142 66-68: Please give order of magnitude of the clock offsets before correction  
143 and the precision of the alignment afterwards.

144 We have added this information in the revised manuscript. The paragraph  
145 now reads: “A clock offset of approximately 2 minutes between the ALR  
146 CTD and the MicroRider was resolved by calculating lagged correlations  
147 between the MicroRider pressure sensor and the CTD pressure sensor to find  
148 the offset, then correcting for the identified clock offset and drift. ”.

149 74: Please explain why you used median instead of mean?

150 See our explanation above.

151 95: Could indicate this is likely because of  $F = ma$ ; ie the same force on the  
152 huge autos produces much smaller accelerations.

153 Thank you for this prompt, the revised sentence now reads: “ Unlike mi-  
154 crostructure measurements performed with a small, light-weight AUV (e.g.  
155 Kolås et al., 2022), the shear microstructure recorded on AutoSub Long  
156 Range was not critically impacted by vehicle vibrations, possibly due to its  
157 greater mass.”

158 105: on which this study focuses.

159 Thank you, we have made the correction.

160 105 general: is this the first paper that presents the details of shear mi-  
161 crostructure from Autosub? Surprising if so but if true, you might consider  
162 showing a few spectra and additional details, possibly in an appendix, so  
163 that future work can cite this paper.

164 This is not the first such paper, we refer the reader to Davis et al. (2022) for  
165 information of the spectral response of the shear probes on ALR. We have  
166 added the sentence “The shear power spectra from a MicroRider mounted  
167 on an ALR have been described in detail in Davis et al. (2022).” to the

168 manuscript.

169 111: Shih et al is a very bad reference for this! They find a  $Re_b$ -dependent  
170 Gamma. Suggest just citing Osborn (1980). There are also now a handful of  
171 observational references supporting the assertion that  $\Gamma = 0.2$ .

172 Thank you for pointing this out, we have removed the reference to Shih.  
173 The sentence now reads “ $\Gamma = 0.2$  is the mixing efficiency, a measure of the  
174 amount of available turbulent kinetic energy that is permanently converted  
175 to potential energy by turbulent mixing, which is generally set to 0.2 (Osborn,  
176 1980)”

177 113: How close to the bottom of the ice is the shallowest CTD measurement  
178 shown? The very strong gradients at the very top of the cavity CTD casts  
179 (Fig 2 black) are interesting.

180 The CTD cast goes right to the ice – ocean interface. We refer the reader  
181 to A. Wåhlin et al. (2024) for a discussion of the CTD measurements at the  
182 interface.

183 123 and throughout: I believe units should be in roman, not italicized, font.

184 We have corrected this where we found such instances, all remaining format-  
185 ting will be finalized in the copy editing process.

186 136: Suggest reformatting the equation.

187 We have reformatted the equation.

188 140: Please make it very clear that  $Ri$  (under the ice at least) is based on a  
189 single  $N^2$  profile whereas the shear is a function of location and time. This  
190 is OK, but appropriate caveats as to its governing local instabilities without  
191 in-situ  $N^2$  should be given.

192 Thank you for this comment, we have added the following words to the  
193 paragraph describing  $Ri$ : “Thus,  $Ri$  is calculated from a constant value of  $N^2$ ,  
194 based on a single profile in the cavity, and shear is a function of space and

195 time along the track of the ALR. Variations of Ri due to variations in  $N^2$   
196 are not captured. For constant  $N^2$ , Ri is low in areas of high shear.”

197 173: Generally, avoid “there is” in favor of more active language such as  
198 “flow is to the ...”

199 We will change some of our wording where we deem appropriate in the revised  
200 manuscript. The sentence in question here has been reworded to “A bottom  
201 intensified southward current flows into the cavity in the east, between the  
202 400 m and 900 m isobaths, and a shallower, bottom intensified northward  
203 current flows out of the cavity in the west (Figure 3c).”

204 177: High compared to what?

205 This sentence has been rephrased to read “Below 500 m depth, turbulent ki-  
206 netic energy dissipation is elevated (compared to other areas below 500 m  
207 along the ice front) in the inflow. Turbulent kinetic energy dissipation is  
208  $\approx 10^{-8} \text{ W kg}^{-1}$  in the inflow over an area approximately 7 km wide and 200 m  
209 high (Figure 3d; turbulent kinetic energy dissipation rate is elevated between  
210 38 km and 45 km of the ice front and  $\sim 200$  m above the seabed).”

211 177: runon sentence.

212 In addition to adding context (see above), this sentence has been split into  
213 shorter sentences.

214 Figure 3, lines 2 and 4 of caption: runon sentences. Also, the dots are said  
215 to indicate the starting locations - but they are a continuous line. I’d have  
216 thought there would just be two starting locations, one for center and one  
217 for east? Please clarify.

218 The new Figure 3 has shorter sentences in the caption and the dots are  
219 described as: “10-minute medians of the values measured by the ALR are  
220 shown as coloured dots in panels a-d. The two dots with bold outlines show  
221 the starting locations of the ALR east and centre short dive tracks into the  
222 cavity.”



223 Figure 4: Personally I think it would be better to keep the aspect ratio  
 224 constant between Fig 3 and 4. Also, sine you already plotted velocity in  
 225 Figure 3, suggest including a panel of  $N_2$ . The aspect ratio is all the more  
 226 a problem later when the authors are comparing epsilon to the different  
 227 instability indicators - but the reader must go back and forth between figure  
 228 3 and 4. Suggest standardizing the aspect ratio and including an epsilon  
 229 panel in Figure 4. Possibly even adding Ri contours to the epsilon panel or  
 230 epsilon contours to the Ri panel since the authors are trying to demonstrate  
 231 correspondence between the two quantities.

232 Thank you for this feedback, we have combined Figure 3 and 4 into the  
 233 new Figure 3, in which all panels have the same aspect ratio. We have also  
 234 included a panel of  $N_2$  at the ice front. We have not plotted Ri contours on  
 235 the epsilon panel, as that proved to be confusing (switching between density  
 236 contours and Ri contours).

237 Also, the Ri panel is just a big sea of red. Consider plotting something else  
 238 to highlight the unstable regions such as  $Ri^{-1}$  or  $Fr = Uz/N$ .

239 The Ri panel is mainly red due to the choice of colourbar. We chose to plot  
 240  $Ri < 1/4$ ,  $1 > Ri > 1/4$ , and  $Ri > 1$  as three different colours in keeping  
 241 with established practice (e.g. Dotto et al., 2025) to distinguish along criteria  
 242 for instability. Plotting  $1/Ri$  would make it less obvious where  $Ri < 1/4$ . We  
 243 would like to avoid plotting additional instability metrics such as the Froude  
 244 number to avoid confusion.

245 182: Doesn't negative PV mean unstable? The whole water column is unsta-  
 246 ble? Is it backwards in the southern hemisphere? Some statements to clarify  
 247 would be useful.

248 We have clarified this in the text by adding the sentence: "Instabilities may  
 249 develop when potential vorticity and  $f$  have opposite signs, as  $f$  is nega-  
 250 tive in the southern hemisphere, potential vorticity  $> 0$  indicates conditions  
 251 favourable to instability. "

252 188: I don't agree with this statement - the high dissipation does not appear  
253 to me to line up at all with for  $Ri$ . Furthermore, given the ADCP's finite  
254 vertical resolution and noise, some additional detail needs to be given on  
255 how seriously we are to take the numerical value of  $Ri$ . I think that either  
256 some wavenumber spectra and transfer functions a la Polzin 2002 need to be  
257 included, or  $Ri$  used as a qualitative indicator.

258 As far as we understand Polzin et al., 2002 the vertical wavenumber re-  
259 sponse of the ADCP is relevant when calculating turbulent dissipation from  
260 the ADCP. We are not using the ADCP for turbulence. We use the VMP or  
261 microrider for shear microstructure and the LADCP and ADCP on the ALR  
262 to get an idea of the vertical and horizontal structure of the water column  
263 at much larger scales, a background value if you will.  $Ri$  and other stability  
264 criteria are frequently calculated from LADCP output with bin sizes of 8 m  
265 and used in comparisons with microstructure data (e.g. Dotto et al., 2025;  
266 Naveira Garabato et al., 2017; Naveira Garabato et al., 2019). We have clar-  
267 ified our reference to  $Ri$  and mixing, the paragraph now reads: "The region of  
268 high turbulent kinetic energy dissipation rate  $\varepsilon$  in the inflow (Figure 3d) co-  
269 incides with instances of  $R_i < 1/4$  captured at 40 km (Figure 3h), indicating  
270 conditions favourable to turbulent mixing. Turbulent kinetic energy dissi-  
271 pation rate is larger than  $10^{-8}$  here, one to two orders of magnitude higher  
272 than the background value (Figure 3d). Dotto et al. (2025) found similar  
273 results for the outflow of DIS. Although areas of high  $\varepsilon$  extend beyond areas  
274 of  $Ri < 1/4$ ,  $\varepsilon$  is higher and  $Ri$  is lower in the upper watercolumn and close to  
275 the seabed. We observe areas of low  $Ri$  and  $Ri < 1/4$  that are not associated  
276 with high values of  $\varepsilon$ , e.g. at 25 km along the transect. "

277 191: I disagree; elevated mixing is much broader than the regions of  $Ri < 1/4$   
278 - augmenting my previous point.

279 We will clarify that the high epsilon includes, but extends beyond, the region  
280 of low  $Ri$ . The relevant sentence now reads: "Although areas of high  $\varepsilon$  extend  
281 beyond areas of  $Ri < 1/4$ ,  $\varepsilon$  is higher and  $Ri$  is lower in the upper watercolumn  
282 and close to the seabed."

283 193: This statement is not justified. Epsilon appears surface intensified as  
284 well. And while it is bottom intensified, I do not think the statement that  
285 it is heightened over rough topography, shear or high currents (of which  
286 you generally must choose either high current or high shear, not both...) is  
287 supported. And as before, I don't think that high epsilon lines up with low  
288 Ri either. Either way, if this statement is retained, more analysis needs to  
289 be shown - scatter plots, binned averages, etc.

290 We have clarified that we are only considering epsilon below the Winter Water  
291 layer, thus we do not discuss high epsilon at the surface. We have included  
292 the following: "Below 500 m depth, turbulent kinetic energy dissipation is  
293 elevated in the inflow (compared to other areas below 500 m along the ice  
294 front). "With regards to the ice shelf front, we have changed our statement  
295 to read: "Our observations show turbulent mixing to be patchy, bottom  
296 intensified and to coincide with high velocities (Figure 3)." We maintain that  
297 in the cavity high epsilon is associated with high shear and low Ri and have  
298 added correlations plots that show this to the manuscript. See more below.

299 197: runon sentence. And seemingly unrelated sentences. Ri governs shear  
300 instability, not symmetric instability... (I understand they are highly corre-  
301 lated here, but they are different, so clarification is needed).

302 We will insert a paragraph break before "at the nearby Pine Island Ice  
303 Shelf...". The start of the new paragraph now reads: " At the nearby Pine  
304 Island Ice Shelf (PIIS) Naveira Garabato et al. (2017) conducted ADCP and  
305 VMP transects along the calving front. Naveira Garabato et al. (2017) do  
306 not detect a fast, narrow, turbulent inflow current, unlike what we observed  
307 at DIS (Figure 3). High rates of turbulent kinetic energy dissipation below  
308 the WW were mostly confined to the PIIS outflow. The PIIS is connected to  
309 another ice shelf cavity to the north and may receive some of its inflow from  
310 under this neighbouring ice shelf, which may decrease the inflow across the  
311 PIIS front and possibly the turbulent mixing there. Additionally, the ice shelf  
312 draft of the PIIS is deeper ( $\approx 400$  m) than the DIS ( $\approx 350$  m). The ice shelf  
313 draft induces a barotropic jump (an abrupt change in water column thick-

ness, blocking flow along constant lines of water column thickness) and limits barotropic inflow to the cavity (A. K. Wåhlin et al., 2020), thus decreasing inflow current velocities and possibly turbulent mixing. ”. The sentence regarding  $R_i$  has been removed, the relevant paragraph now reads “Symmetric instability is driven by high vertical current shear (Figure 3j). The region of high turbulent kinetic energy dissipation rate  $\varepsilon$  in the inflow (Figure 3d) coincides with instances of  $R_i < 1/4$  captured at 40 km (Figure 3h), indicating conditions favourable to turbulent mixing.”

202: What is a barotropical jump?

It is an oceanographic term for an abrupt change in water column thickness. This occurs at the ice shelf front, since ocean currents want to flow along lines of uniform water column thickness, the ice shelf draft poses a barrier to flow, even at depths deeper than its draft. We have added a parenthetical “The ice shelf draft induces a barotropic jump (an abrupt change in water column thickness, blocking flow along constant lines of water column thickness) and limits barotropic inflow to the cavity (A. K. Wåhlin et al., 2020), thus decreasing inflow current velocities and possibly turbulent mixing. ” to the sentence in question.

207: Please rewrite this passive and vague sentence.

We have rephrased this sentence, it now reads “Because the ALR measurements were not coincident in time with the LADCP section, the ALR may have failed to capture transient patches of high turbulent kinetic energy dissipation rate present in the LADCP section.”

204-210: Suggest moving this speculative bit to the discussion.

We originally had results and discussion split, but chose to integrate them to avoid duplicating information and to limit jumping back and forth between topics. We will retain this structure.

216: I think it would be nice to compare this to open ocean values at a similar depth and/or abyssal values, for context. Otherwise “weakly stable”

343 doesn't have meaning.

344 We have added typical open ocean values for  $N^2$  in the Southern Ocean.  
345 The sentence now reads: "We estimate  $N^2$  below a depth of 900 m to be  
346  $6 \times 10^{-7} \text{s}^{-1}$ . This is about three orders of magnitude lower than typical open  
347 ocean values for the southern ocean (King et al., 2012), indicating weakly  
348 stable stratification in the cavity. "

349 218: Style guides such as Strunk and White suggest avoiding "Figure x  
350 shows..." in favor of "statement x is true (Figure y)."

351 We have changed this sentence to read: "In the cavity, the ALR detected  
352 currents that flow predominantly southeastward with low vertical shear in  
353 the east dive track, and a more mixed pattern in the two centre dive tracks  
354 (Figures 5 and 6)."

355 223: Figure 6 and 5 -> Figures 5 and 6

356 Thank you for pointing out this typo, it has been corrected.

357 236: redundant. Suggest "Maximum values were" or "Values reached."

358 We have changed the sentence to read: "Maximum values were  $10^{-7} \text{W kg}^{-1}$   
359 ( $\epsilon$ ) and  $10^{-2} \text{m}^2 \text{s}^{-1}$  ( $\kappa$ ), respectively (Table 2)."

360 238: Again, I'm afraid I don't see this. There are counter examples where  
361 epsilon is high over flat bottoms. Please include plots that allow direct com-  
362 parison such as plotting epsilon with Ri, current speed or bathymetric slope  
363 over plotted, or scatter plots or binned averages (e.g. epsilon(Ri) etc) if you  
364 want to make this claim.

365 We have included scatter plots in our revised manuscript. We stress that  
366 our data are extremely noisy and thus correlation coefficients are low even if  
367 relationships are statistically significant to the 0.1% level. Additionally, the  
368 bathymetry gradient deeper in the cavity is affected by the low resolution  
369 of BedMachine, preventing us from fully resolving the relationship between  
370 bathymetry and epsilon.

371 241: Please remind reader that it's Ri computed from in-situ where and

372 The remainder of this sentence seems to be missing, please clarify

373 245: Again, please include transfer function and instrument response infor-  
374 mation if you wish to quantify the numerical value of Ri versus using it as a  
375 qualitative indication. Note as well that these transfer functions and hence  
376 the mapping of true to measured Ri will be different for the Autosub and the  
377 LADCP.

378 Can you clarify what you mean by the mapping of true to measured Ri and  
379 how that is influenced by the ADCP? It is common practice to calculate Ri  
380 from vertical shear from the 8 m binned ADCP and we do not use the ADCP  
381 to calculate fine scale microstructure.

382 257: Is it really necessary to use a package like this to compute a spatial  
383 gradient? More fundamentally I do not see a relationship between RMS  
384 bathymetric slope and dissipation rate.

385 We do not use a package to calculate the gradient. The bathymetry from ALR  
386 is only 1D, to get a 2D gradient we use BedMachine to get the bathymetry  
387 normal to and along the ALR dive track. Can you clarify what you mean by  
388 "RMS bathymetric slope"? We do not calculate RMS of the slope and to our  
389 eyes there is a clear relationship between the bathymetry and epsilon close to  
390 0 km on the east dive track. We have included scatter plots, linear fit lines,  
391 correlation coefficients and p-values for the relationship between bathymetric  
392 slope and epsilon which clearly show a strong connection.

393 264 onwards: consider moving all of this comparison to past work to the  
394 discussion, so that the results section just has your results?

395 A previous draft had results and discussion separated and the feedback from  
396 several readers was that this caused unnecessary confusion, duplication and  
397 jumping back and forth. We will keep the results and discussion merged.

398 270: I'm confused here, sorry. Weren't the ALR measurements entirely in

399 the warm inflow, since they were so deep?

400 We will clarify this sentence, you are correct that all our ALR measurements  
401 are in the warm layer of in the cavity, but we define the inflow as the narrow  
402 bottom intensified current along the 700 m isobath. The sentence now reads:  
403 “We observed our highest mixing values in the bottom intensified inflow to  
404 the cavity, whereas Kimura et al. (2016) observed the highest levels of mixing  
405 close to the grounding line. Our ALR dive tracks did not reach the grounding  
406 line, and the dive tracks of Kimura et al. (2016) did not cover the inflow of  
407 the PIIS, making comparison difficult. Naveira Garabato et al. (2017) did  
408 not find enhanced mixing in the PIIS inflow. ”

409 272: runon sentence.

410 The sentence in line 272 is not that long, are you sure that this is the line  
411 number you meant?

412 273: Due to what mechanism? This sentence has been modified to read:  
413 “Kimura et al. (2016) hypothesised that high (horizontal) density gradients  
414 driven by temperature differences and a bathymetric ridge can drive a baro-  
415 clinic current with strong vertical current shear. This high shear in turn  
416 drives high levels of turbulence at the ridge under PIIS. Our study shows  
417 that high density gradients are not a requirement for high levels of turbu-  
418 lence.”

419 281: Please change “this” to “their” to avoid confusing with your study.

420 This change has been made.

421 285: If you are going to state dissipation rates this low, I think you do need  
422 to demonstrate your minimum detectability threshold. Earlier you said it  
423 was  $1e-10$ . So how then do you get a median lower than this.

424 The detection limit is between  $10^{-11}$  and  $10^{-10}$  depending on the dive track.  
425 We never state that the detection limit is  $10^{-10}$  and have clarified the sen-  
426 tence you refer to. The paragraph now reads: “Smaller, narrower peaks at

427 frequencies below 10Hz in the accelerometer spectra are successfully removed  
 428 by the Goodman method for dissipation rates above  $1 \times 10^{-8} \text{ W kg}^{-1}$ . De-  
 429 viations from the fitted Nasmyth spectra remain for dissipation rates below  
 430  $1 \times 10^{-9}$ , arguing that quantitative estimates of dissipation rate in very qui-  
 431 escent regimes are not as reliable as estimates of high dissipation rates. Indi-  
 432 vidual dive tracks show good agreement between shear spectra and Nasmyth  
 433 spectra for dissipation rates lower than  $1 \times 10^{-10} \text{ W kg}^{-1}$ . Where dissipation  
 434 rates calculated from two orthogonal shear probes show good agreement, we  
 435 are confident in reporting dissipation rates down to  $1 \times 10^{-11} \text{ W kg}^{-1}$ . Ad-  
 436 ditionally, any signal in the shear spectra caused by the AUV motion, and  
 437 not removed by the Goodman filter, will have minimal effects on the spatio-  
 438 temporal pattern of high and low  $\varepsilon$  observed by the ALR or the qualitative  
 439 assessment of these patterns, on which this study focuses.”

440 Again, I think median should be avoided for all quantities unless there is a  
 441 good reason. Why not just use the mean?

442 We do not want to cause confusion by switching between mean and median for  
 443 data with and without outliers or non-normal distributions, since  $\text{median}(x)$   
 444  $= \text{mean}(x)$  when  $x$  is normally distributed we think median is a better choice.  
 445

446 333: The reason for these calculations is revealed here - suggest giving it ear-  
 447 lier to make the reader understand why they are being told all of this. More  
 448 fundamentally, is that the only reason turbulence is important to measure  
 449 under ice shelves? Ie, as a possible mitigator of the advective heat flux by  
 450 these warm flows?

451 We have added the following paragraph to the introduction together with a  
 452 schematic of the ice shelf cavity: “Basal melt under Dotson is highest close to  
 453 the grounding line of the Kohler East (often referred to as Smith West) and  
 454 Kohler West glaciers (Khazendar et al., 2016; Gourmelen et al., 2017). The  
 455 Kohler West grounding line lies at the southern end of the dashed path shown  
 456 in Figure 1a. A cross-section of the cavity along the path (Figure 1b) shows  
 457 an idealized view of the cavity circulation under the Dotson Ice Shelf. Warm



458 water entering the cavity in the east, and traveling along a path shallower  
459 than the 830 m deep sill (Jordan et al., 2020), can reach the grounding  
460 line. Warm water that reaches the grounding line causes high basal melt  
461 and grounding line retreat (Khazendar et al., 2016; Gourmelen et al., 2017).  
462 The sill may limit direct access of the deepest and warmest mCDW to the  
463 grounding line (Jordan et al., 2020; Khazendar et al., 2016). The addition of  
464 meltwater to the warm, salty mCDW forms a buoyant plume which travels  
465 along the underside of the ice before exiting the cavity in the west. Along its  
466 path, the water experiences turbulent mixing which can transport heat and  
467 salt upward, modifying the properties of the water which ultimately interacts  
468 with the grounding line, and the properties of the buoyant plume exiting the  
469 cavity.” As far as the ice melt rate and modelling efforts in the cavity are  
470 concerned the heat flux is the major concern. As we discuss above, mixing  
471 is also important for the trace metal and nutrient transport, however we do  
472 not have measurements of concentration gradients in the cavity and can not  
473 make a statement as to how the mixing influences them.

474 I, at least as a non ice sheet person, would like to see a cartoon (words  
475 or actual graphic) showing a cross section of the hypothesized warm water  
476 flow to the grounding line. The reason for this is that I don’t currently  
477 understand why the study focused so much on the near-bottom mixing. I’d  
478 think that the heat loss out of the mCDW would be better quantified near  
479 its upper edge. As the authors point out, the water near the bottom is  
480 very weakly stratified so the heat fluxes are expected to be small. Aloft  
481 nearer the interface, the gradients would be stronger, but also the distance  
482 from the topography which is presumably generating most of the turbulence  
483 (my comments above about that not having been adequately demonstrated  
484 notwithstanding). So, statements that mixing is weak such as on lines 356-  
485 258 should be tempered somewhat. And I think the cartoon or written  
486 description of the flow giving readers the sense of which depths are thought  
487 the most likely to eventually contact the ice would help inform this discussion,  
488 at least for me.

489 Have added a schematic to the revised manuscript, thank you for the sug-

490 gestion. We have added the following text to the introduction: “Basal melt  
 491 under Dotson is highest close to the grounding line of the Kohler East (often  
 492 referred to as Smith West) and Kohler West glaciers (Khazendar et al., 2016;  
 493 Gourmelen et al., 2017). The Kohler West grounding line lies at the south-  
 494 ern end of the dashed path shown in Figure 1a. A cross-section of the cavity  
 495 along the path (Figure 1b) shows an idealized view of the cavity circulation  
 496 under the Dotson Ice Shelf. Warm water entering the cavity in the east,  
 497 and traveling along a path shallower than the 830 m deep sill (Jordan et al.,  
 498 2020), can reach the grounding line. Warm water that reaches the ground-  
 499 ing line causes high basal melt and grounding line retreat (Khazendar et al.,  
 500 2016; Gourmelen et al., 2017). The sill may limit direct access of the deepest  
 501 and warmest mCDW to the grounding line (Jordan et al., 2020; Khazendar  
 502 et al., 2016). The addition of meltwater to the warm, salty mCDW forms a  
 503 buoyant plume which travels along the underside of the ice before exiting the  
 504 cavity in the west. Along its path, the water experiences turbulent mixing  
 505 which can transport heat and salt upward, modifying the properties of the  
 506 water which ultimately interacts with the grounding line, and the properties  
 507 of the buoyant plume exiting the cavity.”

## 508 **2 Reviewer 2**

509 This paper discusses microstructure observations made beneath Dotson Ice  
 510 Shelf using an Autonomous Submersible Vehicle. The data appear to lack the  
 511 temporal and spatial coverage that would enable substantive conclusions to  
 512 be drawn about the role of turbulent mixing in the larger-scale processes that  
 513 operate beneath the ice shelf. They are, nevertheless, intrinsically interesting,  
 514 in that they represent some of the very few direct observations that we have  
 515 from within a sub-ice-shelf cavity. That remote part of the ocean plays a  
 516 pivotal role in setting the mass balance of the Antarctic Ice Sheet and hence  
 517 its impact on global sea level, so any observations are of value. I would  
 518 therefore recommend publication of the paper with only relatively minor  
 519 changes.

520 Thank you for your positive assessment of our manuscript. We agree that  
521 a greater spatial and temporal range of observations in ice shelf cavities is  
522 needed to gain a complete picture of the water mass transformations, heat  
523 and (fresh)water transport that influence Antarctic ice mass loss, grounding  
524 line retreat, sea level rise, deep water formation and nutrient transport. Until  
525 such measurements are routinely possible, we intend our manuscript to offer  
526 a glimpse at conditions and possible processes.

527 Suggested changes:

528 Title:

529 It's a minor point, but the current title does not reflect the content of the  
530 paper very well. It promises observations of ocean currents. While they are  
531 included there is very little discussion of them, and no more space is devoted  
532 to currents than to water properties.

533 We have changed the title of the paper to: "Observations of turbulent mixing  
534 in the Dotson Ice Shelf cavity"

535 Abstract:

536 Reflects the content of the paper and thus its main weakness, which is a  
537 lack of substantive conclusions. I accept that it is hard to put such detailed  
538 observations into a broader context, especially when they are made in such a  
539 data-poor region. However, I wonder if it might be possible to put an order  
540 of magnitude estimate on the cavity-wide mean vertical heat flux, given es-  
541 timates of inflow/outflow temperatures, residence time, melt rate, etc. That  
542 would put the numbers quoted in the abstract into a useful context.

543 Thank you for this suggestion. We have modified our abstract to read:

544 "Average vertical heat fluxes are on the order of  $0.1 \text{ W m}^{-2}$  and maxi-  
545 mum heat fluxes reach  $52 \text{ W m}^{-2}$ . This is compared to the  $59 \text{ W m}^{-2}$  to  
546  $176 \text{ W m}^{-2}$  needed to maintain observed average basal melt rates at DIS.  
547 Turbulent mixing is higher in the fast-flowing inflow region and over rough

topography. We show a highly complex spatial pattern of turbulent mixing and of bottom topography. The bottom topography is currently not resolved in bathymetry products and both the topography and turbulent mixing are currently not resolved in models of ice-shelf-ocean interactions. The levels of turbulent mixing experienced by the warm mCDW inflow to the DIS will lead to negligible loss of heat during its path to the grounding line, leaving plenty of heat available to melt the ice shelf base there. Higher average vertical heat fluxes than observed here must occur in areas of the cavity not resolved in this study. ”

We have expanded on your suggestion and added the following to our results section:

“We can estimate the DIS basal melt, assuming that the entire heat flux is used to melt ice at a depth of approximately 1000 m. With this assumption the melt rate  $m$  is

$$m = \frac{Q_t}{L_i \rho_i} \quad (1)$$

with  $L_i = 3.315 \times 10^5 \text{ J kg}^{-1}$  the latent heat of fusion at 1000 dbar, and  $\rho_i = 917 \text{ kg m}^{-3}$  the density of ice. This results in melt rate estimates of  $2 \times 10^{-3} \text{ m yr}^{-1}$  to  $11 \times 10^{-3} \text{ m yr}^{-1}$ . Published estimates for area averaged melt rates under DIS range from  $6 \text{ m yr}^{-1}$  to  $18 \text{ m yr}^{-1}$  (Gourmelen et al., 2017; Lilien et al., 2018; Robertson, 2013; Jenkins et al., 2018; Schodlok et al., 2012) with some estimates up to  $32 \text{ m yr}^{-1}$  (Jenkins et al., 2018). The low upward heat flux within the mCDW layer is thus not able to maintain the observed melt rates under DIS. To achieve the melt rate estimates from (Gourmelen et al., 2017; Lilien et al., 2018; Robertson, 2013; Schodlok et al., 2012; Jenkins et al., 2018) the vertical heat flux would need to be greater than  $59 \text{ W m}^{-2}$  to  $316 \text{ W m}^{-2}$ , values three to four orders of magnitude larger than our median estimates and up to six times our maximum estimate (Table 2).

Davis et al. (2025) showed elevated levels of  $\varepsilon$  in the ice-ocean boundary

layer under Thwaites, and Kimura et al. (2016) observed elevated values of  $\varepsilon$  close to the ice–ocean interface and over a bathymetric ridge in front of the PIIS grounding line. In these areas high turbulent kinetic energy dissipation rate and high vertical and horizontal temperature gradients lead to high temperature fluxes. Our study did not reach the ice–ocean boundary layer or the ridge limiting flow to the DIS grounding line (Figure 1) which may explain the underestimate of the area averaged ice shelf melt rate using the observed heat fluxes. The value for  $\kappa$  for stably stratified water, used in the ISOMIP+ protocol, matches our estimate of  $\kappa$ . Thus, modelled vertical heat transport, in regions for which this estimate is used, could also be too low to explain observed ice shelf melt rates. The low heat fluxes in the interior of ice shelf cavities would need to be offset by higher heat fluxes at the grounding line and in the ice–ocean boundary layer.

We can additionally estimate a melt rate for DIS from the temperature difference between the inflow and the outflow of the cavity and the average residence time of water within the cavity. We take the heat needed to warm the ice shelf to the freezing point temperature and the heat needed to warm the melt water to the temperature of the outflow into account. A back-of-the-envelope calculation for melt rate gives:

$$m = \frac{V_{in} C_p \bar{\rho} (\theta_{in} - \theta_{out})}{\rho_i A_{DIS} (C_i (\theta_f - \theta_{ice}) + L_i + C_p (\theta_{out} - \theta_f))} \quad (2)$$

with  $V_{in} = v_{cavity} A_{inflow}$ , the volume transport in the inflow;  $v_{cavity} = \frac{D}{t}$  the velocity of the inflow;  $D$  the distance water has to travel from the ice front to the grounding line and back;  $t$  the time the water takes to travel to the grounding line and back;  $A_{inflow}$  the area through which water flows into the cavity;  $A_{DIS}$  the area of the DIS;  $C_i$  the specific heat capacity of ice at  $-2^\circ\text{C}$  and 1000 db;  $\theta_f$  the freezing point temperature of seawater;  $\theta_{in}$  the average temperature of the inflow to DIS;  $\theta_{out}$  the average temperature of the outflow from DIS;  $\theta_{ice}$  the far-field internal temperature of the DIS.

We assume the following values for these parameters:  $D = 240$  km (Figure 1);

606  $t = 2$  months (Milillo et al., 2022; Yang et al., 2022; Girton et al., 2019);  
 607  $A_{inflow} = 500 \text{ m} \times 15 \text{ km}$  (solid box in Figure 3);  $A_{DIS} = 5200 \text{ km}^2$  (Lilien  
 608 et al., 2018);  $\theta_f \approx -2^\circ\text{C}$ ;  $\theta_{in} = 0.2^\circ\text{C}$  (the average temperature in the solid  
 609 box in Figure 3);  $\theta_{out} = 0.17^\circ\text{C}$  (the average temperature in the dashed box  
 610 in Figure 3, the outflow extends to shallower depths in the water column  
 611 than the inflow due to the thinner ice shelf draft in the west of the DIS  
 612 (e.g. A. Wählin et al., 2024).);  $\theta_{ice} = -25^\circ\text{C}$ , an estimate of the far-field ice  
 613 temperature. Our estimate of melt rate and heat flux from in and outflow  
 614 temperatures is most sensitive to the area over which we average outflow  
 615 temperatures (Figure 3) and represents an order of magnitude estimate only.

616 Equation 2 results in an estimate of the melt rate of  $\sim 10 \pm 5 \text{ m yr}^{-1}$ , which  
 617 lies within the range of published values (e.g. Gourmelen et al., 2017; Lilien  
 618 et al., 2018; Robertson, 2013; Schodlok et al., 2012; Jenkins et al., 2018).  
 619 To maintain this melt rate the vertical heat flux in the cavity would need  
 620 to be  $100 \pm 50 \text{ W m}^{-2}$ , about three orders of magnitude higher than the  
 621 median values along the east dive track (Table 2). Rearranging Equation 2  
 622 allows us to estimate the percentage of the heat entering the ice shelf cavity  
 623 that is used to melt ice. We estimate that the inflow transports  $4 \pm 2 \text{ TW}$   
 624 into the cavity and the melt takes up  $0.6 \pm 0.4 \text{ TW}$ , thus, only  $\sim 15 \pm 9\%$   
 625 of the heat entering DIS is used to melt the ice shelf. Modelling studies  
 626 have estimated this value to be smaller, at 8 % (Jourdain et al., 2017), but  
 627 within our error range. Transport calculations by Jenkins et al. (2018) yield  
 628 the same range for heat flux into the cavity as our estimate does, however,  
 629 their calculated melt rate, derived from melt water fluxes, is significantly  
 630 higher ( $6 \text{ m yr}^{-1} - 33 \text{ m yr}^{-1}$ ). These melt rates would require heat fluxes of  
 631  $60 \text{ W m}^{-2} - 317 \text{ W m}^{-2}$ . We need significantly more measurements under ice  
 632 shelves to understand the role of mixing in different areas and regimes, and  
 633 its effect on ice shelf melt rate.”

634 Introduction:

635 The first paragraph talks about the DIS contribution to Amundsen Sea “mass  
 636 loss”, suggesting that the term refers to shrinkage of the ice sheet. However,

637 the latter part of the paragraph partitions “mass loss” for the ice shelves  
638 between calving and melting. In this instance the term does not refer to  
639 shrinkage of the ice shelves, but the contribution to the wastage side of the  
640 mass budget. Those are different concepts, and the distinction should be  
641 clarified.

642 Thank you for pointing out this ambiguity, we have rephrased this paragraph  
643 to now only refer to ice sheet mass loss and basal melt/thinning of the ice  
644 shelf. The revised paragraph now reads “Between 1979 and 2017 DIS con-  
645 tributed 0.6 mm to global eustatic sea level rise (Rignot et al., 2019). The rate  
646 of discharge across its grounding line has increased throughout the satellite  
647 record (Rignot et al., 2019; Mouginot et al., 2014) and the grounding line has  
648 retreated (Rignot et al., 2014; Scheuchl et al., 2016; Milillo et al., 2022). The  
649 increased ice flux across the Dotson grounding line, coupled with the stable  
650 ice flux across the calving front (Rignot et al., 2013; Mouginot et al., 2014)  
651 and the increased thinning of the ice shelf (Rignot et al., 2013; Mouginot et  
652 al., 2014; Gourmelen et al., 2017; Greene et al., 2022) leads to the conclu-  
653 sion that ocean thermal forcing has increased basal melt of the ice shelf (e.g.  
654 Mouginot et al., 2014). Dotson has thinned at a 37% higher rate than the  
655 average rate of thinning in the Amundsen Sea (Paolo et al., 2015). ”

656 The last paragraph states that there have been only two previous published  
657 studies of mixing beneath ice shelves, but that overlooks studies based on  
658 borehole data. The oversight is repeated in other parts of the manuscript.

659 This is a good point. We clarify the Introduction, by referring to autonomous  
660 vehicles: “To our knowledge, there exist two published studies of mixing in  
661 an ice-shelf cavity measured by an underwater vehicle, one under Pine Island  
662 Glacier (Kimura et al., 2016), and one under the Filchner Ronne Ice Shelf  
663 (Davis et al., 2022). We present a third such study, targeting DIS.” and  
664 discuss comparisons with borehole results in the Discussion section. To our  
665 knowledge there are three published studies of successful measurements of  
666 mixing through ice shelf boreholes in the Antarctic (Davis & Nicholls, 2019;  
667 Venables et al., 2014; Davis et al., 2025), but if you are aware of others we

668 have missed, we would happily include them. We have added the new Figure  
669 8 and the following paragraph to the revised manuscript:

670 “The highest levels of turbulent mixing occur in the inflow region at the ice  
671 front and in the east dive track, decreasing into the cavity (Figure 5 and  
672 Figure 6). The east dive track clearly shows the highest values for  $\varepsilon$  of the  
673 three ALR dive tracks at DIS (Figure 8). The range, maximum and median  
674 values of  $\varepsilon$  measured with the VMP at the ice front are higher than  
675 those observed in the cavity with the ALR, but ranges have a wide overlap.  
676 We compare our observations of turbulent kinetic energy dissipation rate  
677 with other observations under Ronne Ice Shelf (measured using a MicroR-  
678 ider mounted on an ALR; Davis et al., 2022), George VI Ice Shelf (measured  
679 with a VMP through a borehole; Venables et al., 2014), Thwaites Ice Shelf  
680 (measured with a VMP through a borehole; Davis et al., 2025) and Larsen C  
681 ice shelf (measured with a turbulence instrument cluster moored close to the  
682 ice–ocean interface; Davis & Nicholls, 2019). The distributions of  $\varepsilon$  under  
683 Ronne and George VI have similar shapes and ranges to our observations  
684 (Figure 8). The VMP observations do, however, show much higher maxi-  
685 mum values. This is likely caused by the greater vertical extent of the VMP  
686 measurements, which reach into the ice–ocean boundary layer where  $\varepsilon$  is el-  
687 evated (Davis et al., 2025). This is confirmed by the measurements 2.5 m  
688 and 13.5 m from ice–ocean interface under Larsen C, which show the highest  
689 average values of  $\varepsilon$  of the measurements included in Figure 8. Further studies  
690 are needed to establish whether observed differences between ice shelves are  
691 driven by different mixing regimes or different observation techniques. The  
692 current state of knowledge leads us to conclude that the measurements taken  
693 under Dotson agree remarkably well with available distributions of  $\varepsilon$  from  
694 other ice shelves, outside of the ice–ocean boundary layer. ”

695 Data and methods:

696 On line 127 there is a parenthetical note to authors that has not been ad-  
697 dressed.

698 We are very sorry to have missed that! Thank you for pointing it out, the



699 note has been removed.

700 On line 140 the dimensionless parameter could more precisely be referred to  
701 as a “gradient Richardson number”.

702 Yes, we will make that change.

703 On line 165 there is a mention of detiding LADCP data using CATS2008.  
704 Elsewhere it is stated that tides are unimportant, and CATS cannot be  
705 trusted because of the poor bathymetry in the model. One comment refers  
706 to (mainly) sub-ice data and the other to ice front data, but nevertheless  
707 the treatment seems inconsistent. If bathymetry is poor beneath the ice,  
708 won’t that influence currents at the ice front? If tides are weak enough to be  
709 ignored, why bother with detiding the LADCP data?

710 Thank you, we acknowledge that this is confusing for readers and have  
711 rephrased the text to make this clearer.

- 712 • As noted by the reviewer, the bathymetry is worse in the cavity than  
713 outside it due to a lack of observations in the cavity (as shown in Figure  
714 7), which is why we trust the CATS2008 model more at the ice front  
715 than in the cavity.
- 716 • At the ice front section we have ship ADCP for validation of the CATS  
717 model, making us more confident in the tide model solution for cor-  
718 recting the LADCP.
- 719 • The tidal currents at the ice front are indeed small  $O(1 \text{ cm s}^{-1})$  (Dotto  
720 et al., 2025).
- 721 • We expect the effect of tides in the cavity to be influenced by the  
722 barotropic jump at the ice front, something not well captured in the  
723 CATS model.
- 724 • We could not identify a tidal signal in our ALR ADCP time series in  
725 the cavity and thus we concluded that the error introduced by a faulty  
726 tidal model would likely be larger than the error caused by not detiding.

727 • On a practical point, the LADCP data were processed and detided by  
728 Dotto et al. (2025), and in order to stay consistent with Dotto et al.  
729 (2025) we use their detided dataset (Dotto, Tiago S et al., 2024) in our  
730 study.

731 We make these methodological approaches clearer in the revised paper. The  
732 relevant section of the methods now reads:

733 “Upward-looking and downward-looking LADCP measurements were pro-  
734 cessed with the LDEO\_IX toolbox, incorporating information from the vessel-  
735 mounted ADCP, CTD, GPS and bottom track from the LADCP (Thurnherr,  
736 2021). The processed data were averaged into 8-m vertical bins and detided  
737 using an updated version of the CATS2008 Antarctic tide model (Padman et  
738 al., 2002; Erofeeva et al., 2024). Modelled tidal current components are on  
739 the order of  $1 \text{ cm s}^{-1}$  at the ice front and the tide model agrees well with tides  
740 extracted from the shipboard ADCP data (Dotto et al., 2025). Conversely,  
741 the ALR ADCP data are not detided due to the ill-constrained bathymetry  
742 under DIS, the absence of a detectable tidal signal in a spectral analysis of  
743 the ALR ADCP currents in the cavity, and the risk of degrading the ADCP  
744 data quality with an ill-fitting tidal model. ”

745 Results and discussion:

746 In the title of section 3.1 and elsewhere in the manuscript the edge of the ice  
747 shelf is referred to as the “ice shelf front”. The correct term for that feature  
748 is the “ice front”.

749 We have made this change.

750 In figure 3, is there a “black line” showing the track of the ALR (third line  
751 of the caption)? I couldn’t see one.

752 Thank you for this feedback, we have removed the line (which is obscured in  
753 large parts by the coloured dots showing the ALR measurements along the  
754 track.) and the reference to it from the figure caption.

755 On lines 222-223 it is stated that water at the ice front is colder and lighter  
756 than that in the cavity. Does that refer only to measurements made with  
757 ALR? Was the warmer, saltier water apparent in the section observed with  
758 the ship? If not, I think it deserves some comment about where that warm,  
759 salty water may have come from? Waters in the cavity must be cooled and  
760 freshened, so the observation must say something about variability at the ice  
761 front. If, on the other hand, an equally warm, salty water mass is present in  
762 the ship CTD data, then the statement in the paper is a little misleading.

763 The statement on lines 222–223 refers to the ship CTD section along the ice  
764 front. We have made this clearer in the revised paper. The revised paragraph  
765 reads:

766 “Water at the ice front (measured with the ALR and the ship CTD) is  
767 colder but lighter than water found deeper in the cavity (Figure 6). The  
768 temperature (Figures 6 and 5) and salinity (not shown) in the cavity generally  
769 increase with depth. The presence of warmer, saltier, and denser water in the  
770 cavity than at the ice front may indicate seasonal or interannual variability  
771 in the properties of the water at the ice front (as described by Kim et al.  
772 (2021)) and thus of water flowing into the cavity.”

773 On lines 228-229, and elsewhere, it is stated that the observations reported  
774 in the paper are important for establishing mixing rates that can be “incor-  
775 porated into numerical models”. It is not clear to me how these data would  
776 be incorporated into a model. Perhaps the point could be clarified?

777 Thank you, we are happy to expand and clarify how our results may inform  
778 modelling efforts. We have removed the reference to parameterisations in  
779 lines 228–229 and adding the following to the manuscript at the end of sec-  
780 tion 3.2: “Maximum and median values of diapycnal diffusivity  $\kappa$ , vertical  
781 heat flux  $Q_T$ , and vertical salt flux  $Q_S$  from our observations under DIS are  
782 given in Table 2. Our median values of diapycnal diffusivity ( $O(10^{-4} \text{ m}^2 \text{ s}^{-1}$ –  
783  $O(10^{-5} \text{ m}^2 \text{ s}^{-1})$ ) are the same order of magnitude as globally-averaged ocean  
784 values (Waterhouse et al., 2014). The maximum values of diapycnal diffusiv-  
785 ity in our study ( $O(10^{-2} \text{ m}^2 \text{ s}^{-1}$ – $O(10^{-3} \text{ m}^2 \text{ s}^{-1})$ ) match values observed close

786 to the seabed over rough terrain or at ridges (Waterhouse et al., 2014).

787 Our observations under DIS provide valuable metrics against which turbulent  
788 mixing processes in numerical models could be assessed. Turbulent kinetic  
789 energy dissipation is not modelled or parameterised in regional or  
790 global models. Instead, diapycnal diffusivity  $\kappa$  is parameterised. A common  
791 parametrisation of diapycnal diffusivity in ice shelf cavities is the vertical  
792 profile method from Large et al. (1994) (e.g. in ROMS; Gwyther et al.  
793 (2015) or MITgcm; Nakayama et al. (2017)) which assumes higher values  
794 of  $\kappa$  in boundary layers than in the interior. The interior mixing is made  
795 up of contributions from internal waves (parameterised as a constant), from  
796 shear instability (parameterised from the gradient Richardson number), and  
797 from double diffusion (parameterised from the double diffusion density ratio)  
798 (Large et al., 1994). The ice base roughness has been shown to influence  
799 the ice–ocean boundary layer mixing and the heat and salt flux into the  
800 boundary layer, and thus the spatial and temporal distribution of ice shelf  
801 melt (Gwyther et al., 2015). We are not aware of studies investigating the  
802 effects of spatially variable bottom boundary layer roughness on mixing and  
803 basal melt in an ice shelf cavity. The range of values for  $\kappa$ , the spatial  
804 variability, and forcing mechanisms we discuss, can be compared to the values  
805 and variability of the  $\kappa$  profile parametrisation. This may allow a better  
806 understanding of the contribution of different drivers to mixing and of how  
807 realistic model mixing is.

808 Another common choice to parametrize mixing, used in the ISOMIP+ pro-  
809 tocol (Asay-Davis et al., 2016), is to prescribe constant values for  $\kappa$  in the  
810 vertical and horizontal, with higher values where the water column stratifi-  
811 cation is unstable. In stably stratified water, as under DIS, the ISOMIP+  
812 protocol sets as  $\kappa_{v,stable} = 5 \times 10^{-5} m^2 s^{-1}$  (Asay-Davis et al., 2016). The  
813 value of  $\kappa$  used in ISOMIP+ has the same order of magnitude as the me-  
814 dian value in the centre\_short dive track ( $2 \times 10^{-5} m^2 s^{-1}$ ), but is an order of  
815 magnitude lower than the median  $\kappa$  on the east dive track ( $1.1 \times 10^{-4} m^2 s^{-1}$ )  
816 and 2–3 orders of magnitude lower than the maximum values we find within  
817 the cavity (Table 2). Thus, the constant value of  $\kappa$  used in ISOMIP+ is a

818 good choice for slow flows with low shear over smooth topography, but may  
819 underestimate mixing in other areas which may in turn influence modelled  
820 ice-shelf melt. ”

821 On line 230, mention is made of a 100 m thick “melt layer” observed through a  
822 borehole. What feature are you referring to? The upper 100 m of the borehole  
823 data shown in Figure 2 appear to indicate the presence of less meltwater than  
824 deeper in the water column. Why is that? A shallow intrusion of WW along  
825 the ice shelf base?

826 You are correct that this water may show a shallow intrusion of WW, we  
827 have removed the reference to a melt layer.

828 Lines 197-203 draw comparisons with observations made at Pine Island Ice  
829 Front but point out differences in the physical setting. One difference that  
830 might be relevant, but which appears to have been overlooked, is that in  
831 the case of Pine Island there is a neighboring ice shelf to the north, so the  
832 northern sidewall of the channel confining the Pine Island Ice Shelf does not  
833 extend all the way to the ice front.

834 Thank you for pointing this out, we have changed the relevant paragraph  
835 in the revised manuscript to read: “At the nearby Pine Island Ice Shelf  
836 (PIIS) Naveira Garabato et al. (2017) conducted ADCP and VMP transects  
837 along the calving front. Naveira Garabato et al. (2017) do not detect a fast,  
838 narrow, turbulent inflow current, unlike what we observed at DIS (Figure 3).  
839 High rates of turbulent kinetic energy dissipation below the WW were mostly  
840 confined to the PIIS outflow. The PIIS is connected to another ice shelf cavity  
841 to the north and may receive some of its inflow from under this neighbouring  
842 ice shelf, which may decrease the inflow across the PIIS front and possibly  
843 the turbulent mixing there. Additionally, the ice shelf draft of the PIIS is  
844 deeper ( $\approx 400$  m) than the DIS ( $\approx 350$  m). The ice shelf draft induces a  
845 barotropic jump (an abrupt change in water column thickness, blocking flow  
846 along constant lines of water column thickness) and limits barotropic inflow  
847 to the cavity (A. K. Wåhlin et al., 2020), thus decreasing inflow current  
848 velocities and possibly turbulent mixing. ”

849 The last four paragraphs compare findings with other AUV based observa-  
850 tions of microstructure beneath ice shelves. However, elsewhere the manuscript  
851 highlights the differences between those regions. That makes the discussion  
852 feel like one that is motivated by common methodology rather than common  
853 physical setting. Why overlook borehole measurements of turbulence that  
854 have been made within cavities? Later in the section it is suggested that the  
855 AUV track beneath FRIS is 9 km long, but that does not seem to fit with  
856 the figures in the cited paper. At the end the of the section the text again  
857 talks about improving parameterisations of mixing in models, but again, I  
858 don't really see how you would use the data for that.

859 Yes, these are good points, thank you for the suggestions. We have strength-  
860 ened the discussion by adding comparisons with borehole data to the revised  
861 manuscript, we have given details on the revised text above. Apologies for  
862 the incorrect length of the FRIS dive track, this has been corrected. We  
863 have expanded our argument on how our observations can be used to inform  
864 modelling studies, the relevant text is included above.

865 Lines 334-335 suggest that the small vertical heat flux observed means that  
866 a lot of ocean heat can be used to melt ice at the grounding line. But how  
867 much is used there? The outflows at the ice front remain above the freezing  
868 point, so some ocean heat that enters the cavity exits it without being used  
869 for melting. Again, can you estimate some global budgets for the amount of  
870 heat used for melting and the overall average vertical heat flux that could  
871 put your spatially-limited observations in a DIS-cavity-relevant context?

872 These are good suggestions. We have added an estimation of the heat flux in  
873 DIS to the revised manuscript. The paragraph reads: "Equation 2 results in  
874 an estimate of the melt rate of  $\sim 10 \pm 5 \text{ m yr}^{-1}$ , which lies within the range of  
875 published values (e.g. Gourmelen et al., 2017; Lilien et al., 2018; Robertson,  
876 2013; Schodlok et al., 2012; Jenkins et al., 2018). To maintain this melt  
877 rate the vertical heat flux in the cavity would need to be  $100 \pm 50 \text{ W m}^{-2}$ ,  
878 about three orders of magnitude higher than the median values along the  
879 east dive track (Table 2). Rearranging Equation 2 allows us to estimate the

percentage of the heat entering the ice shelf cavity that is used to melt ice. We estimate that the inflow transports  $4 \pm 2$  TW into the cavity and the melt takes up  $0.6 \pm 0.4$  TW, thus, only  $\sim 15 \pm 9\%$  of the heat entering DIS is used to melt the ice shelf. Modelling studies have estimated this value to be smaller, at 8 % (Jourdain et al., 2017), but within our error range. Transport calculations by Jenkins et al. (2018) yield the same range for heat flux into the cavity as our estimate does, however, their calculated melt rate, derived from melt water fluxes, is significantly higher ( $6 \text{ m yr}^{-1} - 33 \text{ m yr}^{-1}$ ). These melt rates would require heat fluxes of  $60 \text{ W m}^{-2} - 317 \text{ W m}^{-2}$ . We need significantly more measurements under ice shelves to understand the role of mixing in different areas and regimes, and its effect on ice shelf melt rate. ”

### 3 Reviewer 3

This paper presents an interesting set of observations, in an environment difficult to access. The analysis is solid. It would be good to put these observations in the context of the previous work that has been done around Dotson - I understand that there are little observations in the cavity, but are the conditions along the face 'unusual'? It's hard to tell, and I acknowledge that this is not about long-term observations at Dotson, but it would be useful to put these observations in a broader context.

Thank you for your positive review of our manuscript. We have added a sentence that the ice front properties we observed in 2022 are within the usual range: “The temperature and salinity at the ice front are within the historic range of watermass distributions and properties at DIS (Kim et al., 2021).”

As pointed out by the other reviewers, some of the key results are a bit either overstated, or unclear.

We have addressed the concerns of the other two reviewers in our responses, which includes modifying some of our key points. The precise changes are

909 detailed in the responses to reviewer 1 and 2.

910 For example, on L193, one would be hard-pressed to directly identify the  
911 'enhanced mixing at the inflow' and it being over a larger area than that of  
912 the outflow from the section alone - it might useful to show a profile or two  
913 of dissipation rates. Sampling (station spacing) might be important when  
914 talking about "area", which is not discussed here.

915 We have changed the paragraph you refer to. It now reads: "Below 500 m  
916 depth, turbulent kinetic energy dissipation is elevated in the inflow (compared  
917 with other areas below 500 m along the ice front). Turbulent kinetic energy  
918 dissipation is  $\approx 10^{-8} \text{ W kg}^{-1}$  in the inflow over an area approximately 7 km  
919 wide and 200 m high (Figure 3d; turbulent kinetic energy dissipation rate is  
920 elevated between 38 km and 45 km of the ice front and  $\sim 200$  m above the  
921 seabed). "

922 Overall, I don't have many comments that were not captured by the other  
923 reviewers. This is an interesting paper and it should be published.

924 Thank you for your positive review.

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