Reconstructing landscapes : an adjoint model of the Stream Power and diffusion erosion equation

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Abstract. We simulate landscape evolution using a diffusion–advection equation, where the advection velocity is derived from the erodibility parameters of the Stream Power Law. This formulation allows for forward modeling within a finite-element framework, and enables the use of adjoint methods for sensitivity analysis and parameter inversion, specifically for spatially variable erodibility and diffusion coefficients. When considered individually, model parameters such as the diffusion coefficient, erodibility, initial topography, and time-dependent uplift can be inverted using constraints from final topography, sediment flux, or cumulative denudation at specific locations. Sensitivity analysis on a real landscape reveals that sensitivity to erosion parameters is higher in steep, high-relief areas and that hillslope diffusion and fluvial incision affect the model differently. We apply the adjoint model to two natural cases: (1) reconstructing the pre-incision topography of the southeastern French Massif Central, which appears as a smooth, flat footwall bounded by a linear escarpment along a major lithological boundary; and (2) estimating the Quaternary uplift rate along the Wasatch Range, USA, where our model suggests a significant increase in uplift from 0.2 to 1 $\mathrm{mm.yr^{-1}}$ over the last $\sim$2 million years.

1 Introduction

Tectonic events coupled to climatic variations shape the Earth's surface. These processes result from thousands to millions of years of evolution. Although modern climate, topography, and fluvial activity can be recorded and monitored, reconstructing past landscape evolution remains challenging. To establish relationships between geological constraints like lithological variations and uplift, and erosion-sedimentation mechanisms that quantify sensitivity to fluvial incision or hillslope processes, Landscape Evolution Models (LEMs) have been developed (e.g., Tucker and Slingerland, 1994; Tucker et al., 2001; Davy and Lague, 2009; Tucker and Hancock, 2010; Braun and Willett, 2013; Carretier et al., 2016; Salles and Hardiman, 2016).

The majority of Landscape Evolution Models (LEMs) are forward models: they rely on assumptions about initial and boundary conditions and are validated retrospectively by comparing their results to available data on surface morphometry, sedimentation, exposure age dating, or low-temperature thermochronology. These models often incorporate a wide range of parameters to describe how interactions between climate, lithology, sediment granulometry, isostasy, and/or tectonics influence landscape evolution (Carretier et al., 2016; Salles, 2016; Braun and Willett, 2013; Tucker et al., 2001). As the number of parameters increases to account for various forcings and physical processes, it can become difficult to isolate the influence of each one on the

25   final result. Moreover, the significance of geomorphological observations in constraining model outputs, i.e., identifying which observations are most informative with respect to the model parameters, is generally unknown. As different sets of parameters can produce similar "final" topographies, another difficulty lies in the non-uniqueness of the model. Additionally, the model implementation itself—including the choice of more or less complex physical laws for erosion, transport, sedimentation, and tectonics—can affect model results and performance. These differences in implementation also complicate sensitivity analysis

30   and model validation (Barnhart et al., 2020b).

    The questions that arise from being able to "read" the landscape and understand its past and present behavior can be addressed by inverse models. Although perhaps less prevalent in geomorphology than in other geoscience disciplines, inverse approaches have been widely developed in recent decades. Neighbourhood algorithms have proven useful to constrain erosion parameters, uplift history, or to reconstruct ancient (pre-glacial) landscapes by fitting the present-day topography and/or sed-

35   iment fluxes (e.g., Croissant and Braun, 2014; Pedersen et al., 2018). Some studies apply them in more complex modelling frameworks that couple erosion, cosmogenic isotope production, and low-temperature thermochronology to provide additional constraints on landscape evolution scenarios (e.g., Glotzbach, 2015). Bayesian approaches have been employed to infer erosional histories from sedimentary records (Yuan et al., 2019), and to constrain temporal variations in model parameters as well as associated uncertainties (Chandra et al., 2019). All these approaches rely on extensive sampling of the parameter space

40   and, consequently, on running large numbers of forward models. These forward models can incorporate various processes, including fluvial and hillslope erosion, sediment deposition, and lithospheric isostatic response, and are sometimes coupled with thermokinematic models to integrate low-temperature thermochronology data as constraints (Braun, 2003; Braun et al., 2012), which may increase their computational cost.

    Among inverse models, many studies have been devoted to inferring uplift rates or climatic histories from river longitudinal

45   profiles (Goren, 2016; Roberts et al., 2012; Roberts and White, 2010; Petit et al., 2017). These specific landscape features are particularly interesting because it has been demonstrated that river channels progressively adjust to external forcings, until eventually reaching a steady state (Willett and Brandon, 2002). The response time, defined as the period required for the channel profile to adjust to new conditions (Armitage et al., 2013; Godard et al., 2013), controls how long river channels can preserve a record of past conditions. This understanding has led to numerous applications, notably in explaining large-scale landscape fea-

50   tures and major channel knickzones, such as those observed in Colorado, Australia, and Africa, as the result of long-term uplift rate variations (Roberts et al., 2012; Roberts and White, 2010; Rudge et al., 2015). The advantage of river profile analyses is that, provided the uplift rate is the only unknown and its spatial distribution is uniform, there is a direct relationship between the channel slope and the uplift rate (Royden and Perron, 2013), making the inversion relatively straightforward. All these studies highlight the potential of inverse modeling to advance our understanding of landscape evolution, improve parameter estimation,

55   and quantify uncertainties. However, aside from river longitudinal profile inversion, most inverse modeling approaches rely on extensive sampling of the parameter space. The same limitation applies to sensitivity analysis, which aims to quantify the relative influence of model parameters or different model formulations on the simulation outcome. For instance, a sensitivity analysis of the CAESAR-Lisflood landscape evolution model (LEM) was performed using the Morris method, which consists of computing elementary effects on the solution by changing one parameter at a time, over very short timescales (approxi-

mately 10 years) (Skinner et al., 2018). This approach involved simplifying the model behavior into model response functions, enabling the estimation of how individual parameters affect the output. Regarding model formulation, Barnhart et al. (2020b, c) also applied this method in a multi-model analysis framework over longer time scales (13 ka), assessing how the inclusion of various processes influenced the model results and how increasing model complexity affected performance. In another study, Armitage et al. (2018) evaluate the long-term (several Myr) sensitivity of classical landscape evolution models to changes in precipitation rates and show that the choice of the transport law (diffusion or incision) influences the model response time to perturbations and the outgoing sediment flux. While these methods offer insights into global model sensitivity at different time scales, they often come at a high computational cost.

In contrast to the previously cited gradient-free methods, the adjoint approach enables efficient sensitivity analysis and the inversion of unknown parameters by a cost-effective, gradient-based approach (e.g., Hu and Kozlowski, 2020). Despite their potential, the use of adjoint models in geomorphology remains limited, with few exceptions in related fields such as turbidity currents modeling (Parkinson et al., 2017), terrain correction for weather forecasting (Tao et al., 2019) or hydrodynamics (Clare et al., 2022). The adjoint method involves solving an additional partial differential equation (PDE), called the adjoint equation, which provides information about how changes in the model's output variables affect the objective function (or cost function) that is being optimized (e.g., Givoli, 2021). By solving both the forward and adjoint problems, one can efficiently compute the gradient of the cost function with respect to all input parameters, regardless of their number. The resulting sensitivity is local, meaning it represents the gradient of the cost functional in the vicinity of the current parameter state. In this paper, we present a new modelling approach that simulates fluvial incision and hillslope processes as a diffusion-advection equation using a finite-element (FE) formulation. The adjoint equation is used to estimate model sensitivity to erosion parameters and to reconstruct the initial conditions, source term or erosion coefficients in different settings. The method combines components of the Landlab toolbox (Hobley et al., 2017; Barnhart et al., 2020a; Hutton et al., 2020) to remove local minima from the initial topography and compute the drainage area, with the FE discretization, numerical solution and adjoint provided by the Firedrake software (Ham et al., 2023).

2 Landscape evolution model

2.1 Simulating landscape evolution with a diffusion-advection equation

Landscape evolution models in their simplest configuration aim at solving (e.g., Perron et al., 2008, and references therein):

$$\frac{\partial h}{\partial t} = -[\nabla \cdot (\mathbf{Q_h}) + \nabla \cdot (\mathbf{Q_f})] + U, \tag{1}$$

where h is the elevation [L], $\frac{\partial h}{\partial t}$ is the rate of change of elevation over time [L.T^{-1}], $\mathbf{Q_h}$ and $\mathbf{Q_f}$ are sediment fluxes per unit surface [L^2.T^{-1}] coming from hillslope processes and fluvial incision, respectively, and U is the uplift rate [L.T^{-1}] accounting for tectonic and/or isostatic vertical movements. The flux $\mathbf{Q_h}$ represents hillslope processes, which are responsible for the progressive smoothing of the landscape with time. It writes:

$$\mathbf{Q_h} = -\kappa \nabla h, \tag{2}$$

3

with κ the diffusion coefficient [$L^2.T^{-1}$]. Physical and empirical relationships have permitted to relate the fluvial incision rate arising from the divergence of the fluvial sediment flux $\mathbf{Q_f}$ to the local channel slope S and the upstream drainage area A [L^2]:

$$\epsilon = \nabla \cdot \mathbf{Q_f} = K_f A^m S^n. \tag{3}$$

Here, K_f represents the rock erodibility [$L^{1-2m}.T^{-1}$], m is an exponent related to the area, and the exponent n controls the non-linearity of fluvial erosion (e.g., Howard et al., 1994). This equation simulates the incision of bedrock river channels and is known as the Stream Power Law (SPL) (e.g., Seidl et al., 1994; Whipple and Tucker, 1999). It assumes that river incision is detachment-limited, meaning the erosive power of the fluvial system is controlled solely by the flow's capacity to detach material from the bedrock. Ideally, the local channel slope S is equal to the norm of the topographic gradient vector ∇h such that:

$$\epsilon = K_f A^m \|\nabla h\|^n. \tag{4}$$

Substituting Eqs. 2 and 3 into Eq. 1 we get:

$$\frac{\partial h}{\partial t} = \nabla \cdot (\kappa \nabla h) - K_f A^m \|\nabla h\|^n + U. \tag{5}$$

Eq. (5) can be cast into a diffusion-advection equation (e.g., Perron et al., 2008; Braun and Willett, 2013) considering a few modifications:

1. The exponent n is set to 1, as usually done (e.g., Seidl et al., 1994), in order to ensure a linear dependence of the advection term to the topographic gradient.

2. The scalar quantity $K_f A^m$ is interpreted as the norm of a velocity vector $K_f A^m \mathbf{u}$, where $\mathbf{u}$ is a unit vector oriented along the topographic gradient. This velocity vector can be seen, for example, as the rate of knickpoint propagation within the drainage network (Whipple and Tucker, 1999).

3. The advection velocity $\mathbf{c}$ is defined as:

$$\mathbf{c} = K_f A^m \mathbf{u}. \tag{6}$$

These modifications of Eq. (5) lead to the following equation:

$$\frac{\partial h}{\partial t} - \nabla \cdot (\kappa \nabla h) + \mathbf{c} \cdot \nabla h = U, \tag{7}$$

The dot product $\mathbf{c} \cdot \nabla h$ can be expressed as:

$$\mathbf{c} \cdot \nabla h = \|\mathbf{c}\| \|\nabla h\| \cos \theta, \tag{8}$$

Where θ is the angle between both vectors; since $\mathbf{u}$ is defined to be parallel to the topographic gradient ∇h, Eq. 8 simplifies to $\mathbf{c} \cdot \nabla h = \|\mathbf{c}\| \|\nabla h\|$. Replacing the norm of the velocity vector $\mathbf{c}$ from Eq.(6) and keeping $n = 1$, we recover the expression of the Stream Power Law given in Eq.(3). Compared to the general form of a diffusion-advection equation:

$$\frac{\partial h}{\partial t} - \nabla \cdot (\kappa \nabla h) + \nabla \cdot (\mathbf{c} h) = U, \tag{9}$$

this formulation intrinsically assumes an incompressible velocity field, i.e., $\nabla \cdot \mathbf{c} = 0$. This assumption is justified because $\mathbf{c}$ does not represent a physical transport of mass or energy (as in heat or pollutant advection), but rather the propagation of topographic information (e.g., incision fronts or knickpoints), which is not subject to creation or destruction. It also assumes that the erodibility is the same for bedrock and for sediment produced by hillslope processes. This is a reasonable approximation only when the removal of hillslope-derived sediment represents a minor contribution to the overall fluvial erosion compared to direct bedrock incision. Otherwise, it could lead to spurious channel morphologies and steepnesses (Litwin et al., 2025).

Assuming that the local channel slope is equal to the norm of the topographic gradient requires that the direction of the gradient is accurately respected when computing the river network. Significant approximations in the gradient direction would limit the validity of Eq. (5) when interpreted as a diffusion-advection equation. In many LEMs however, due to the discretization of the topographic grid and the use of single-flow direction water routing algorithms, the computed drainage direction is not always co-linear with the topographic gradient as it should be. Certain routing algorithms, such as the D_∞ method (Tarboton, 1997), distribute water flux between two nodes with links oriented closely to the gradient direction, which permits to minimize the orientation discrepancy between the drainage direction and the true topographic gradient. The D_∞ method also permits divergent drainage to occur, which is impossible with single-flow direction methods.

2.2 Solving the forward problem

Solving equation Eq. (7) as a classical diffusion-advection PDE offers several advantages. First, in a diffusion-advection equation, the quantity which is involved in diffusive and advective processes is expected to be the same i.e., the gradient ∇h. In contrast, in Eq. (5) the diffusive term uses the actual topographic gradient, while the fluvial incision term uses the local channel slope, which, depending on the local problem and mesh resolution, may or may not correspond to the topographic gradient. Moreover, this new approach naturally propagates information upstream by advecting topographic characteristics with a velocity proportional to the drainage area in the direction of the gradient resulting in two important features: (1) along channel sides, although slow, the advection velocity can be nonzero, allowing channel widening with time and (2) along channel paths, despite small slopes, the advection is more pronounced thanks to the large drainage area.

Furthermore, the classical coupling of 2D hillslope diffusion with 1D (vertical) fluvial incision in Eq.(5) may introduce scaling issues and strong resolution dependency, arising from the implicitly fixed river width (Hergarten, 2020). Reformulating the model as Eq.(7) helps reduce these issues: while resolution still affects drainage area computation and topographic gradient estimation, the assumption of a fixed channel width is no longer required, since vertical incision is reinterpreted as horizontal advection. Finally, this formulation provides a solid basis for deriving and solving the adjoint equation associated with Eq. (7).

In this study, we present a new method to both solve the forward (Eq. 7) and adjoint PDEs describing landscape evolution as a classical diffusion-advection equation, with a time- and space-variable advection velocity. At each time step k, the topography from the previous step $k-1$ is used to compute the drainage area using Landlab's D_∞ flow routing algorithm, then the drainage area value is copied to a corresponding Firedrake space function. The topographic gradient ∇h is calculated using Firedrake's derivative operators, normalized to obtain a unit vector $\mathbf{u}$, and the advection velocity vector function $\mathbf{c}$ is finally computed as $\mathbf{c} = K_f A^m \mathbf{u}$. This updated advection velocity is then incorporated into the weak formulation of Eq.(7), which is then solved using Firedrake's finite element framework with Q_1 elements.

To solve Eq. (7), we assume the system is closed with the initial and boundary conditions:

$$h = h_0 \qquad\qquad\qquad \forall \mathbf{x} \in \Omega \text{ and } t = 0,$$

$$h = \bar{h} \qquad\qquad\qquad \forall \mathbf{x} \in \partial\Omega \text{ and } \forall \mathbf{t} \in [0, T],$$

where Ω is the model domain, $\partial\Omega$ its boundary, h_0 the initial topography, and $\bar{h}$ the boundary conditions. The latter are of Dirichlet type (fixed elevation), and can be updated over time to account for progressive uplift at the domain boundaries.

During the forward model run, the time-dependent sediment flux is recorded as the cumulative volume of eroded material, which can subsequently be used in the adjoint model as an additional constraint for parameter inversion. Local denudation, defined as the cumulative height of eroded material at each node over time, can also be recorded.

3 Adjoint method

3.1 Governing equations

We consider the PDE describing the evolution of surface elevation h as in Eq. (7). We rewrite this equation in the residual form $Lh - f = 0$, where: $L := L(p)$ is a differential operator acting on h, such as:

$$Lh = \frac{\partial h}{\partial t} + \mathbf{c} \cdot \nabla h - \nabla \cdot (\kappa \nabla h), \tag{10}$$

and f is the source term, which in this case corresponds to the uplift rate U. The solution $h := h(p)$ of this PDE is controlled by a set of parameters p (initial conditions, h_0, velocity, $\mathbf{c}$, diffusivity, κ, and source term, U).

We define a cost functional $J := J(h)$ that has to be minimized:

$$J = \int_\Omega \phi(h) \mathrm{d}V, \tag{11}$$

where $\phi(h)$ can be defined as the residual between output topography at time T and observations h_{obs} such that:

$$\phi(h_T) = \frac{1}{2} (h_T - h_{obs})^2, \tag{12}$$

and its derivative with respect to h is:

$$\phi'(h) = \frac{\mathrm{d}\phi}{\mathrm{d}h} = h - h_{obs}. \tag{13}$$

6

The objective is to evaluate the gradient $\frac{dJ}{dp}$, which will then allows us either to minimize J (for parameter optimization) or to analyze the influence of certain parameters on the model result (for sensitivity analysis). To this end, we first solve the forward problem using trial parameter values p, and then compute the adjoint variable λ by solving the adjoint equation:

$$L^\dagger \lambda = \phi', \qquad\qquad\qquad \forall \lambda \in \Omega, \tag{14}$$

$$\lambda = 0 \qquad\qquad\qquad \forall \lambda \in \partial\Omega. \tag{15}$$

Here $L^\dagger$ is the formal adjoint of the operator L. For diffusion-advection equations, the adjoint operator has the same form as the forward operator, except that the advection velocity $\mathbf{c}$ is reversed (e.g., Celia et al., 1990):

$$L^\dagger \lambda = \frac{\partial \lambda}{\partial t} - \mathbf{c} \cdot \nabla \lambda - \nabla \cdot (\kappa \nabla \lambda). \tag{16}$$

First, we have to solve for λ in Equation 14, where the forcing term is ϕ'. Once the adjoint variable λ is computed, we use it to evaluate the gradient of the cost functional with respect to the parameters p:

$$\frac{\mathrm{d}J}{\mathrm{d}p} = -\int_\Omega h \frac{\partial L^\dagger}{\partial p} \lambda \mathrm{d}V. \tag{17}$$

The partial derivative $\frac{\partial L}{\partial p}$ can be computed analytically when possible, or numerically. The functional J may also include time-dependent controls, in which case it will integrate the misfit between models and observations over the model duration T.

The computational workflow is as follows:

1. Solve the forward problem $Lh - f = 0$ for a given set of parameters p,

2. Solve the adjoint problem for λ using Equation (14),

3. Compute the gradient $\frac{\mathrm{d}J}{\mathrm{d}p}$ using Equation (17),

4. Update p using an optimization algorithm, and iterate until convergence.

In practice, the Firedrake package automates this process using automatic differentiation (AD) techniques. The user only needs to define the weak form of the PDE and specify the functional J.

Beyond parameter inversion, the adjoint directly provides the model relative (dimensionless) sensitivity to a parameter p through:

$$S_p^{\mathrm{rel}} = \frac{p}{J} \frac{\partial J}{\partial p}, \tag{18}$$

where all physical quantities are non-dimensionalized to ensure consistent scaling of both residuals and sensitivities. This relative sensitivity metric is dimensionless by construction, enabling direct comparison across parameters of different nature and scale, and supporting robust sensitivity analysis.

3.2 Definition of the cost functional

In our problem, the objective function J combines data misfit terms (comparing modeled and observed topography, sediment flux, or denudation rates) with regularization components. To recover solutions with sharp discontinuities while preserving smooth regions, we employ total variation (TV) regularization (Strong and Chan, 2003). Specifically, for the inversion of initial conditions, we replace the standard L_2-norm used in Eqs. (11) and (12) with an edge-preserving TV regularization scheme, adapted from image processing methods (Barcelos, 2002).

$$J = \int_{\Omega} \left[\alpha g \|\nabla h\| + \frac{1}{2}(1-g)(h - h_{\mathrm{obs}})^2 \right] dV, \tag{19}$$

$$g = \left(1 + k|\nabla G * h|^2 \right)^{-1}, \tag{20}$$

where $\alpha > 0$ controls the regularization strength, $g \in [0,1]$ is an edge-detection function using Gaussian-filtered (G) topography gradients, and k modulates edge sensitivity. We calibrate the parameters α and k to suppress high-frequency noise in reconstructed solutions and preserve geomorphologically significant sharp features (e.g., fault scarps, knickpoints). Finally, we define the functional J_{tot} which incorporates regularization terms and misfits to sediment flux and denudation values as follows:

$$J_{tot} = J + C \sum_{i=0}^{t_{max}} (q_i - q_{i_{obs}})^2 \Delta t + D \sum_{j=0}^{N} \sum_{i=0}^{t_{max}} (d_{ij} - d_{ij_{obs}})^2 \Delta t + \text{regularization terms} \tag{21}$$

The first term of the functional represents the misfit between modelled and observed topography integrated over the entire domain. The second term incorporates a penalization of the differences between the estimated outgoing sediment flux, q, and the true sediment flux, q_{obs}, summed over discrete time intervals corresponding to the model timestep Δt. The third term represents the difference between observed and modelled cumulated denudation, d, with time on some prescribed nodes n that could correspond to sampling points. Finally, the coefficients C and D control the relative importance of the sediment flux and denudation residual terms in the inversion. Depending on the parameter we want to analyse, coefficients C and/or D may be zero.

4 Results

Model n°	IC	Dimensions, Resolution m^2, elem	Type	Duration	κ $m^2.yr^{-1}$	K_f yr^{-1}	$U(t)$ $mm.yr^{-1}$	Figure n°
1	Square	$(19 \times 10^4)^2$, 128×128	Analysis	1 Ma	$0.5(\mathbf{1}) \times 10^{-2}$	5×10^{-7}	0	1b
2	Square	$(19 \times 10^4)^2$, 128×128	Analysis	1 Ma	10^{-2}	$1(\mathbf{0.5}) \times 10^{-6}$	0	1c
3	Square	$(19 \times 10^4)^2$, 128×128	Inversion	1 Ma	X	5×10^{-7}	0	2
4	Square	$(19 \times 10^4)^2$, 128×128	Inversion	1 Ma	10^{-3}	X	0	3
5	X (**Scarp**)	$(4 \times 10^4)^2$, 128×128	Inversion	1 Ma	5×10^{-3}	3×10^{-6}	0	4
6	Square	$(4 \times 10^4)^2$, 128×128	Inversion	4 Ma	5×10^{-3}	4×10^{-6}	X (**0.5**)	5
7	X	$(1.9 \times 10^5)^2$, 128×128	Inversion	4 Ma	10^{-3}	10^{-6}	0	6
8	Divide	$2.2 \times 10^4 \times 1.4 \times 10^4$, 128×80	Inversion	4 Ma	10^{-1}	3×10^{-6}	X	7
9	Square	$(4 \times 10^4)^2$, 128×128	Taylor test	0.5 Ma	10^{-2}	3×10^{-6}	0.5	8

Table 1. Summary of models and parameters. IC denotes initial conditions; resolution is the number of elements in the x and y dimensions. For inverse models, the letter X indicates which parameters has been inverted. For models 1 and 2, the parameter value indicate the guess value for which the sensitivity has been computed. Numbers in bold within parentheses indicate which values were used in the corresponding forward model (if relevant).

4.1 Sensitivity analysis and inversion

Before we consider the natural cases of the French Massif Central and Wasatch Range presented in section 4.1.4, we first conduct synthetic experiments by running first the forward FE model to generate synthetic data (e.g., topography, sediment flux, and/or cumulative denudation over time). We then use these data as controls for sensitivity analysis and/or inversion of the main model parameters: initial conditions, erodibility, diffusion coefficient, and source term. In the adjoint method, the forward model is executed once, starting with an initial guess that is slightly different from the value used for generating the synthetic data. In certain inversion procedures, we re-run the forward model using the inverted parameters to compare model outputs (topography, sediment flux, etc.) with synthetic data. Initial topographic conditions for the forward model can consist of simple geometric shapes (linear drainage divide, linear escarpment) or can be derived from an existing digital elevation model (DEM). For instance, to evaluate the model's ability to reconstruct spatial variations of κ and K_f, we use as initial conditions the same DEM as in section 4.1.4, i.e., the south-east border of the French Massif Central. This region provides a suitable test case for highlighting spatial model sensitivity to erosion processes, as it displays a contrasted topography with both flat, smooth surfaces and deeply incised slopes.

4.1.1 Erosion parameters

We evaluate the model sensitivity and its capability to reconstruct spatially variable erosion coefficients K_f and κ on a $190\,\mathrm{km} \times 190\,\mathrm{km}$ grid with 128×128 elements. The initial topography is similar to the target topography used in Section 4.1.4 for initial condition inversion — specifically, a dissected NE-SW linear escarpment marking the southeastern boundary of the French Massif Central (Fig. 1a, 2, and 3).

In the first test, we assume spatially uniform values of $K_f = 5 \times 10^{-7}$ yr^{-1} and $\kappa = 10^{-2}$ m^2.yr^{-1}, and independently compute the relative model sensitivity S_p^{rel} to each parameter, starting with trial values of 1×10^{-6} yr^{-1} and 5×10^{-3} m^2.yr^{-1}, respectively. The sensitivity maps in Figs. 1b and 1c represent the spatial variation of the gradient of the cost functional with respect to these 2 parameters at the point defined by the trial value in the parameter space. The sensitivity values for κ are generally low, except in two distinct regions: the NW corner of the domain, where the Cenozoic Cantal stratovolcano is located, and a NE-SW elongated zone corresponding to the southeastern escarpment of the Massif Central (Fig. 1b). The model sensitivity to the K_f coefficient is larger, especially in the Massif Central region, and is even most pronounced around the headwaters of the major catchments, where topographic slopes are steep (Fig. 1c). Overall, both tests suggest that the model's sensitivity to erosion coefficients is particularly large along the NE-SW trending southern Massif Central escarpment, and very small in the low altitude areas of the Rhône plain.

We then define arbitrarily two simple polygons where these coefficients are different from the background value and then re-run the forward models for 1 Ma. The source term is null, meaning that no uplift is imposed to the surface topography. In the inversion tests, we fix the model duration, initial conditions, source term and either the erodibility or the diffusion coefficient as identical to that of the forward models. With these known values as a starting point, we perform separate inversions to estimate
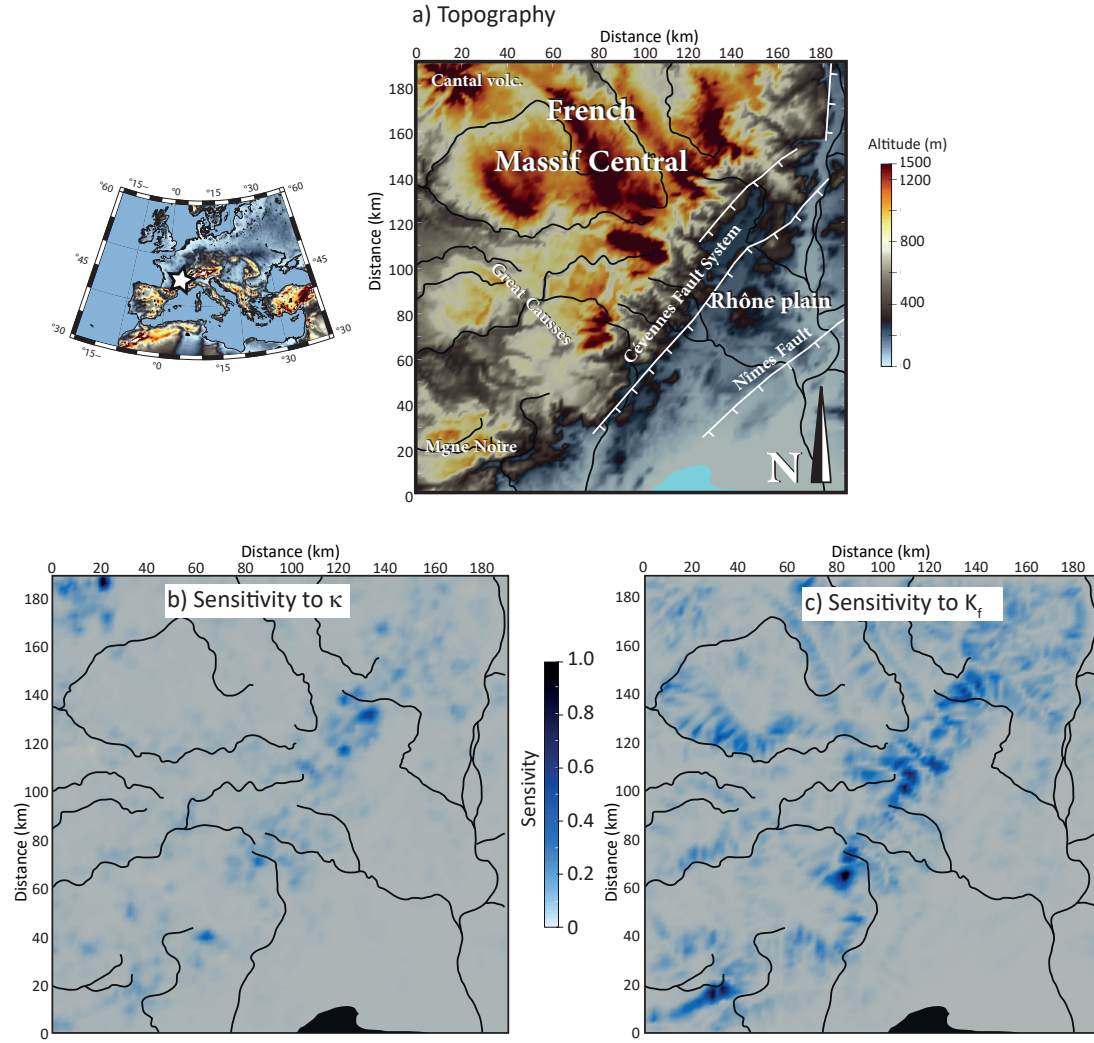
Figure 1. Sensitivity tests on a real topography for the diffusion coefficient κ and the erodibility coefficient K_f (models 1 and 2, Table 1). (a): Topography of the SE part of the French Massif Central with the main drainage systems (solid lines). White star on the left map indicates the location of the study area. (b): Relative model sensitivity (absolute value) to κ. (c): Relative model sensitivity (absolute value) to K_f.

the unknown parameter, either K_f or κ. The primary objective is to evaluate the model capability to accurately retrieve the expected values for these unknown parameters (see Table 1).

260   We carry out the inversion of the diffusion coefficient κ using only the final topography as the reference "observable" (h_{obs}) to control the inversion. Since hillslope diffusion is a local process, it would not be interesting to use other data representing long-distance sediment transport, such as the outgoing sediment flux, as a control. The "true" κ used in the forward model
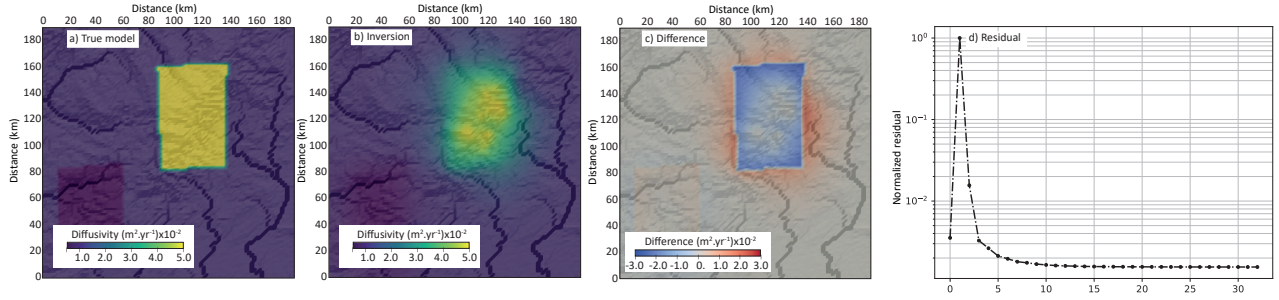
Figure 2. Inversion of the diffusion coefficient κ with the adjoint method (model 3, Table 1). (a): True spatial variations of κ used to build the forward model. (b): inverted κ values map. (c): Residual. (d): Inverse model convergence. Same initial topography as on Figure 1.

consists of a backgroud value of 10^{-2} m^2.yr^{-1} with two polygons of lower (5×10^{-3} m^2.yr^{-1}) and higher (2×10^{-2} m^2.yr^{-1}) values (Fig. 2a). At the beginning of the inversion procedure, we fix the initial value everywhere as identical to the background

265 value. The inversion of the diffusion coefficient κ converges after 33 iterations and the resulting κ map (Fig. 2b) shows that the inversion correctly retrieves the spatial variations of the κ coefficient, with some lateral smearing (Fig. 2c).

The inversion of the erodibility coefficient K_f is performed using both the final topography and the time-distributed outgoing sediment flux as observables. The forward model involves a background K_f value of 1×10^{-6} yr^{-1} with two polygons of larger values (2×10^{-6} and 3×10^{-6} yr^{-1}, Fig. 3a). Similarly to the previous case, the initial "guess" value for K_f is identical to

270 the background value. The inversion converges after 31 iterations and the K_f map shows that the areas of different coefficients are very well delineated, although the absolute value reconstructed by the inversion is slightly different and more variable than the actual one (Figs. 3b and 3c). In some areas, the model does not depart from the trial value, meaning that the inversion result is not well constrained. In the area of intermediate K_f value initially fixed at 2×10^{-6} yr^{-1}, the mean average value reconstructed by the inversion is too large, while it is correct in the area of initially large K_f value. Nevertheless, the areas with

275 initially different K_f coefficient are precisely reconstructed.

4.1.2   Initial conditions

In this section, we demonstrate the model ability to recover initial conditions when erosion parameters (K_f, κ, m), uplift rate and erosion duration are known. The process begins with a test simulation, lasting for 1 million years, on an initial topography featuring a linear escarpment associated with a random topographic noise with an amplitude of maximum 40 meters, and no

280 uplift (Fig. 4a). The final topography resulting from this simulation shows a strongly incised escarpment, though the surface trace of the fault is still relatively easy to follow. It is then used as the reference observable for the adjoint model, together with the evolution of the outgoing sediment flux with time.

The inversion process starts with an initial guess for the initial conditions, which consists of a linear topographic slope in the same direction as in the forward model, but devoid of any escarpment (Fig. 4b). Starting from this initial guess, the

285 inversion finally retrieves an initial topography close to the actual one (although slightly smoother, Fig. 4c). The inverted initial
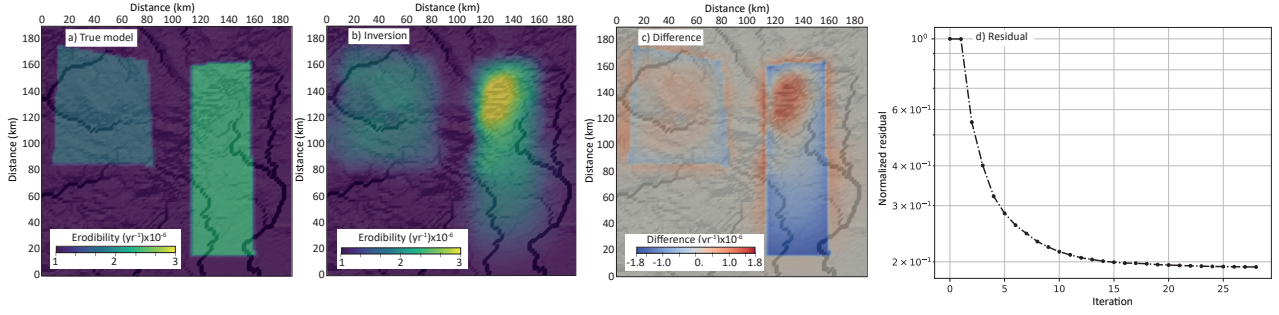
Figure 3. Inversion of the erodibility coefficient K_f with the adjoint method (model 4, Table 1). (a): True spatial variations of K_f used to build the forward model. (b): inverted K_f values map. (c): Residual. (d): Inverse model convergence. Same initial topography as on Figure 1.

topography deviates from the initial guess and consists of two flat-lying surfaces separated by a steep escarpment, which is located close to the actual one. The result is obtained after 57 iterations with a coefficient $\alpha = 2 \times 10^{-4}$ and an edge preservation coefficient $k = 100$ in Eq. (19). The normalized residual stabilizes around 3×10^{-7} and the modelled sediment flux is very close to the actual one, except in the first steps of the model evolution (Fig. 4d and e).

4.1.3 Source term

We now aim to assess the adjoint method's ability to accurately estimate source term variations, assuming that the boundaries of the uplifting area are known. The forward model simulates a square mountain, similar to Model 1. In this scenario, the mountain experiences a constant uplift rate of 0.5 $\mathrm{mm.yr}^{-1}$ over 4 million years.

In the inversion process, we use the final topography, outgoing sediment flux and cumulative denudation on 3 specific model points as control points on the inversion (Fig. 5a). We run the forward model for a duration of 4 Myrs using a time step of 2,000 years. The topography modelled with the inverted uplift rate value after 4 Ma is very similar to the forward model 1 (Fig. 5a and b). The inversion results indicate that the uplift rate value is correctly estimated, with a value ranging from 0.48 $\mathrm{mm.yr}^{-1}$ to 0.51 $\mathrm{mm.yr}^{-1}$ at the end (i.e., with an error of 4 %), with a significant deviation from the initial guess value of 0.2 $\mathrm{mm.yr}^{-1}$ (Fig. 5d). The total outgoing sediment flux and point denudation with time are correctly reproduced (Figs. 5e and f).

4.1.4 Natural cases

The natural examples presented in this study were selected for their relatively simple geological settings and the availability of some data such as denudation, exhumation, or uplift data. These constraints allow us to test our modeling approach on real-world but controlled scenarios. Below, we present the scientific interest and the corresponding results for two such cases. In both cases, we extract the target topography from the Shuttle Radar Topography Mission database (NASA, 2013) and re-sample it at the desired model resolution.

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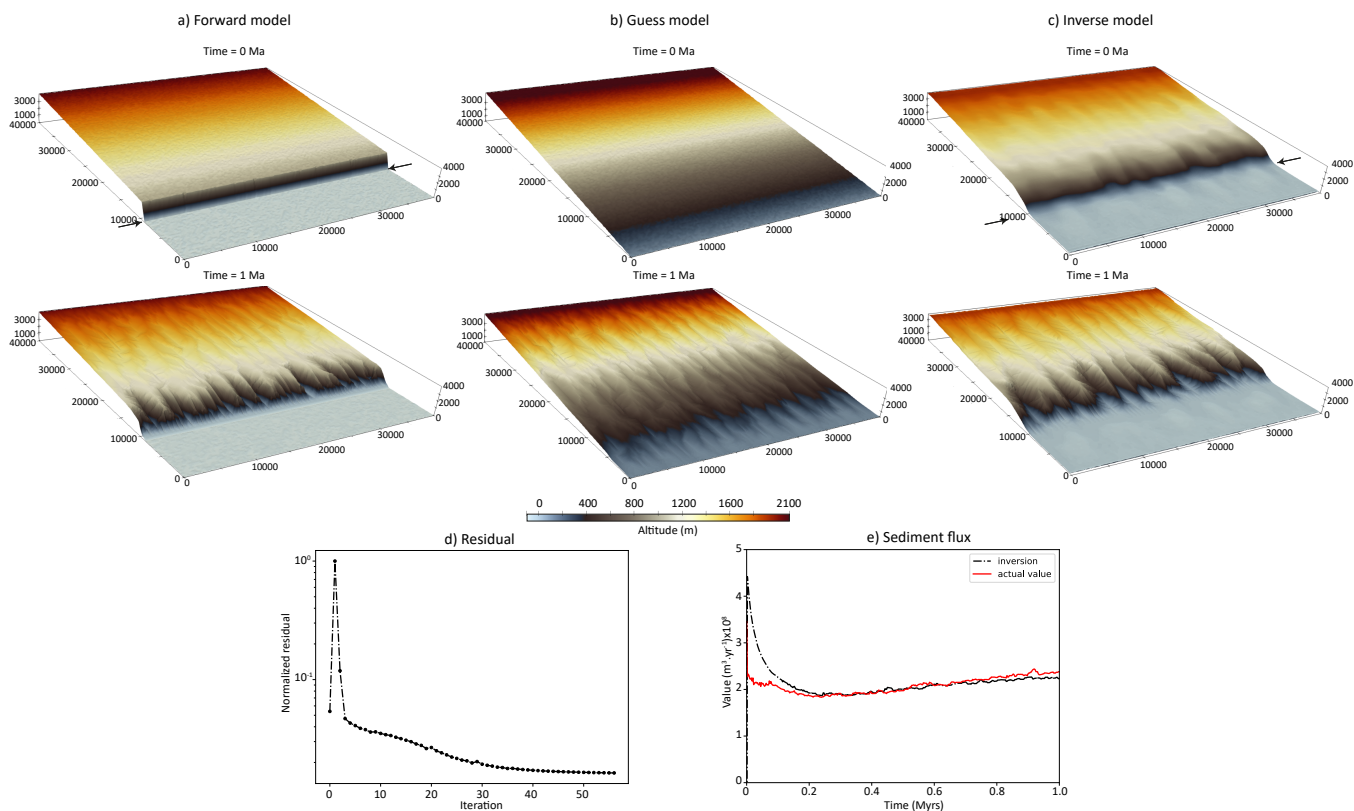
Figure 4. Inversion of initial conditions on a south-facing escaprment (model 5, Table 1). Top and middle panels: initial and final (1 Ma) topographies in the forward model, the initial guess and the inversion result (a, b and c, respectively). Bottom panels: Relative residual (d) and modelled sediment flux (e).

- *SE Border of the French Massif Central*

The SE border of the Massif Central is a former southeast-facing normal fault escarpment (the Cévennes Fault system) originated mainly from Mesozoic to Cenozoic extensional tectonics (Séranne et al., 2021). The uplift of the Massif Central occurred in several steps since the Paleogene, the last stage being a Late Miocene to Pliocene uplift which resulted into intense fluvial incision observed along its SE border (Olivetti et al., 2016; Fauquette et al., 2020). Denudation rates in the recent period range from 0.04 to 0.08 $\mathrm{mm.yr^{-1}}$ (Olivetti et al., 2016). The topography of the southeastern Massif Central prior to this latest phase of incision by tributaries of the Rhône River remains poorly constrained. Reconstructing the pre-incision landscape can provide insights into the original location and geometry of the main escarpment, and give a reliable initial condition to more sophisticated forward LEMs aiming at deciphering the mechanisms and extent of its retreat over the past few million years.
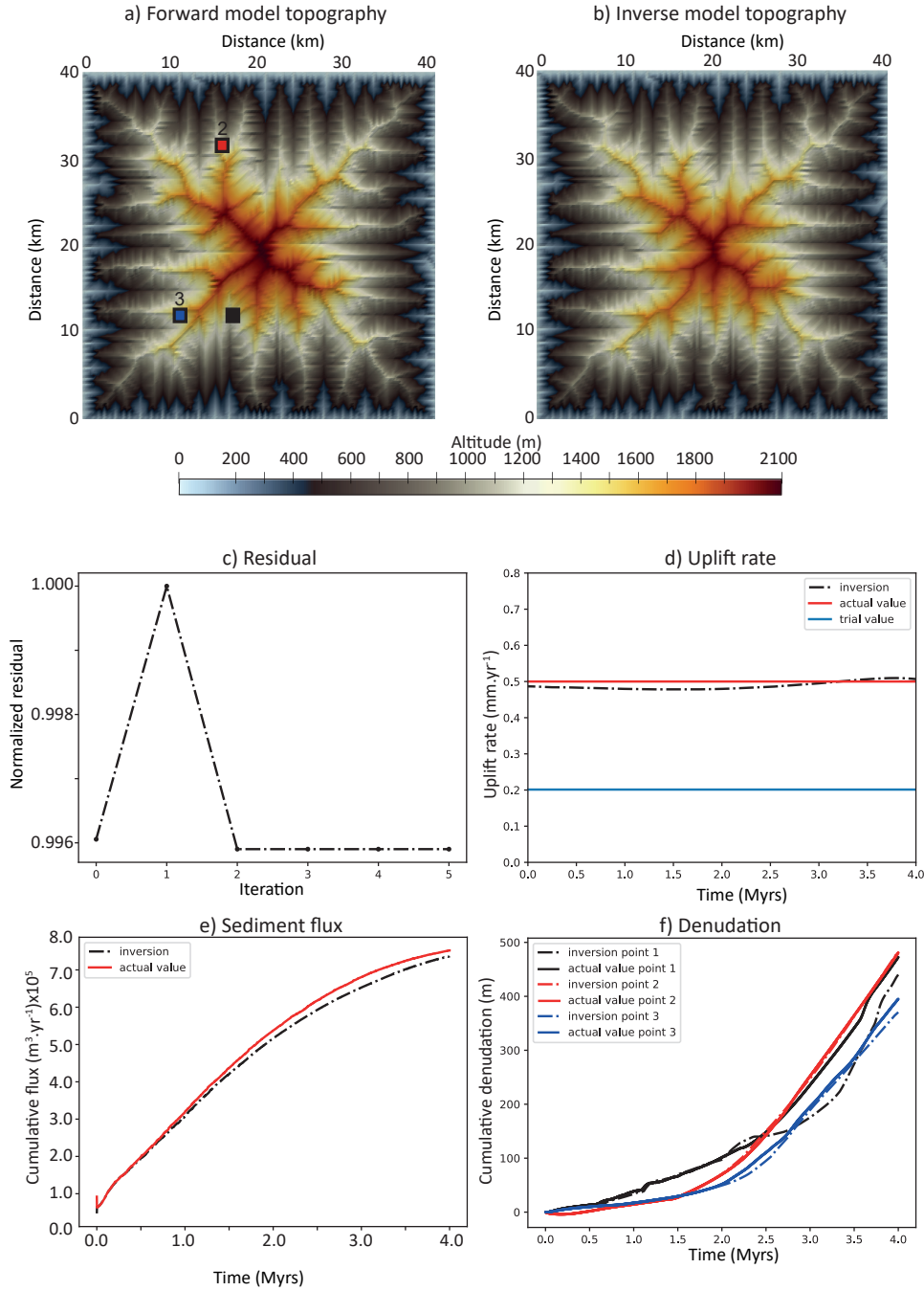
Figure 5. Results of source term (constant uplift rate) inversion (model 6, Table 1). Top: final topography with the actual (a) and inverted (b) uplift rates after 4 Myrs. Middle: normalized functional value (c) and uplift rate (d). Blue line on panel d indicates the trial uplift rate value used in the inversion procedure. (e): outgoing sediment flux from the forward ("actual value") and inverse model ("inversion"). (f): actual and modelled cumulative denudation on the 3 control points (see (a) for point location).

315   In the inversion procedure, we cannot accurately constrain time variations of the erosion parameters, which trade off with model duration (the largest the erodibility coefficient, the shortest the time needed to achieve a given amount of erosion). To address this ambiguity, we adjusted the erosion parameters to obtain denudation rates consistent with present-day estimates (Olivetti et al., 2016) as well as with long-term exhumation rates over several million years (Olivetti et al., 2020). We do not aim, as well, at reproducing specific morpho-tectonic events like the Messinian Salinity Crisis, which occurred at the end of

320 the Miocene, and we neglect the — presumably small — topographic uplift that could result from flexural isostatic response to ongoing erosion during the model duration. Consequently, the initial topography that we will obtain from the adjoint method likely resembles the pre-incision landscape, but its precise age remains unconstrained.

    The initial guess for the topography is either (1) a NE-SW trending divide, featuring a gently NW-dipping plane opposite to a steeper SE-dipping surface, or (2) the current topography. In the first case, the maximum altitude is 2000 m, located along

325 the drainage divide, while in the second case, it reaches only 1500 m in the NW corner of the model, which corresponds to the Miocene Cantal stratovolcano. The model runs for 3 Myrs with a diffusion coefficient of 10^{-2} m^2.yr^{-1} and an erodibility coefficient of 10^{-6} yr^{-1}. In this inversion, only the final topography serves as a reference observable to compute the functional J. For both initial guesses (divide or current topography), the inverted initial topography reveals a relatively smooth, poorly dissected region in the NW half of the model, representing the footwall of the normal fault system, and a lower-altitude plain

330 in the SE half. In both cases as well, these surfaces are delineated by a linear, well-marked escarpment.

    When the initial guess features a smooth topography with a linear drainage divide, the fault escarpment closely follows the observed surface fault traces, particularly in the NW corner of the map, and the average footwall elevation is higher than in the second case. In contrast, the model initialized with the current topography results in lower altitude contrasts and a slightly more sinuous fault escarpment, which is positioned further inland relative to the fault trace. This suggests that using the current

335 topography as an initial guess is less effective in annealing escarpment retreat caused by fluvial incision and predicts a slightly lower average topography compared to the model with a simple linear NE-SW divide.

    Geologically, the reconstructed initial conditions from the "divide" model reveal a contrasting landscape within the Cévennes fault system footwall. In the SW, the topography is characterized by a low-altitude (800 m) domain corresponding to the Mesozoic Great Causses and the Variscan Montagne Noire metamorphic massif. To the north lies a moderately elevated region

340 (1000-1400 m), encompassing the crystalline core of the Massif Central. This area features outcropping Variscan metamorphic rocks and granitoids locally overlain by Cenozoic volcanism. At the northeastern edge of the domain, a large circular massif corresponding to the cenozoic Cantal stratovolcano reaches altitudes exceeding 1400 m, however significantly lower than the 2500 m estimated from pollen assemblages prior to its Plio-Pleistocene erosion (Fauquette et al., 2020). In this model, the incision along the southeastern border of the Massif Central by steep, southeastward-flowing streams is largely erased, with

345 only a few notches remaining downstream. In fact, the morphological fault escarpment is not aligned with the main Cévennes fault, but runs along the boundary between the crystalline Massif Central and the sedimentary cover of the Rhône plain. This boundary is also marked by a major early Mesozoic normal fault. This suggests that prior to incision, the morphological fault escarpment did not coincide with the main trace of the Cévennes fault, but rather with the mixed tectonic and lithological contact lying between the Massif Central metamorphic and granitic rocks and the Rhône plain Meso-Cenozoic deposits. On

350   map, the main trace of the Cévennes fault system outcrops within marly Cretaceous deposits, and may never have constituted a major fault scarp.

    - Uplift rate of the Wasatch Fault, Utah

The Wasatch Fault, in the NE part of the Basin and Range is a well-known active normal fault, having been the focus of numerous tectonic and geomorphological studies since the 1990s. The rock uplift rate of the Wasatch Range in the footwall has

355   been estimated between 0.3 to 1.7 mm.yr^{-1} depending on the time scale, methodology, and specific fault segment considered (Ehlers et al., 2003; Friedrich et al., 2003; Mayo et al., 2009; Stock et al., 2009; Smith et al., 2024). These studies highlight the difficulty of extrapolating short-term uplift estimates to long-term trends. Over the time span from the Late Miocene to the Quaternary, fluctuations in uplift rates have also been observed, which may reflect variations in fault activity or changes in far-field tectonic velocities (Friedrich et al., 2003). Overall, despite the significant amount of data on this fault system, there

360   is no consensus on the long-term uplift rate of the footwall and its evolution through time. Although footwall uplift began in Miocene times (e.g., Armstrong et al., 2003; Ehlers et al., 2003), it is unlikely that the present landscape preserves traces of the earliest uplift events, as geomorphic processes, particularly regressive erosion, tend to erase older features over time. In this model, we treat the source term (uplift rate) as an unknown parameter and we only use the final topography to compute the functional J. The target topography corresponds to a ≈ 14 km-long section of the NS-trending Weber segment of the Wasatch

365   Fault, which separates the Wasatch Range to the east from the Great Salt Lake basin to the west (Fig. 7a). The initial topography reflects the general shape of the Wasatch Range in this area, featuring a NNW-SSE trending ridge with an elevation of 2500 m that overlooks a plain at an elevation of 1300 m (Fig. 7b). The model simulates 4 million years of landscape evolution (Fig. 7c), capturing the time frame during which river incision may still preserve a record of recent uplift history (Smith et al., 2024). We set the erosion parameters as follows: the diffusion coefficient κ is 0.1 m^2.yr^{-1}, and the erodibility coefficient K_f is 3.2×10^{-6}

370   yr^{-1}, close to the values used in previous inverse modeling studies by Smith et al. (2024). The initial guess for the uplift rate is 0.5 mm.yr^{-1}. The residual drops significantly within the first 10 iterations, and the model stops after 24 iterations (Fig. 7d). The estimated uplift rate shows an initial phase with relatively low values (less than 0.2 mm.yr^{-1}), gradually increasing to approximately 1 mm.yr^{-1} over 4 million years of evolution. This result aligns with the average Quaternary uplift rates obtained from river longitudinal profile inversion (Smith et al., 2024) (Fig. 7e) especially between 2 and 3.2 Ma of model

375   evolution where the mean uplift rate seems to increase regularly. In the last 1 Ma, river profiles tend to indicate a drop in uplift rate values from 0.8 to 0.5 mm.yr^{-1} while our inversion still predicts an increase, that is consistent with some short-term uplift rate estimates (Mayo et al., 2009) (Fig. 7e), but not with all. By dividing the modeled outgoing sediment flux by the area of the uplifted region, we estimate average denudation rates ranging from 0.2 to 0.35 mm.yr^{-1} (Fig. 7f), consistent with values derived from cosmogenic ^{10}Be measurements (Stock et al., 2009).

380   ## 5   Discussion

These examples highlight the potential of solving a diffusion-advection equation for both forward and adjoint problems in landscape evolution. Notably, the ability to compute model sensitivity to key parameters such as erosion rates, source terms, or
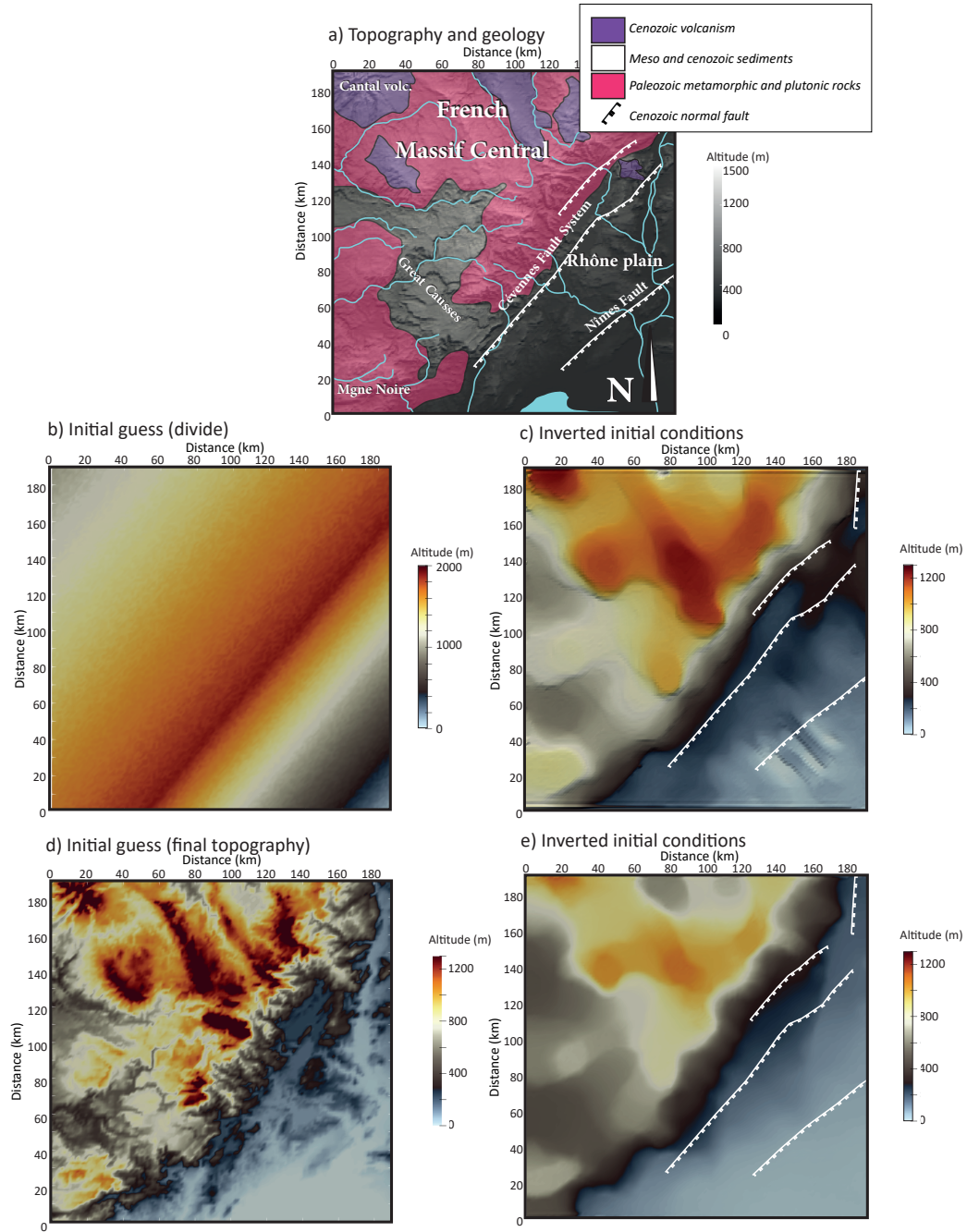
Figure 6. Inversion of initial conditions for the French Massif Central (model 7, Table 1). (a): Simplified geology (transparent colors) and present-day topography (grey shades) of the SE Massif Central with the main cenozoic faults (white lines). (b): guess initial conditions with a NE-SW asymmetric divide; (c): inverted initial conditions with (b) as guess; (d): guess initial conditions corresponding to the current topography; (e): inverted initial conditions with (d) as guess.
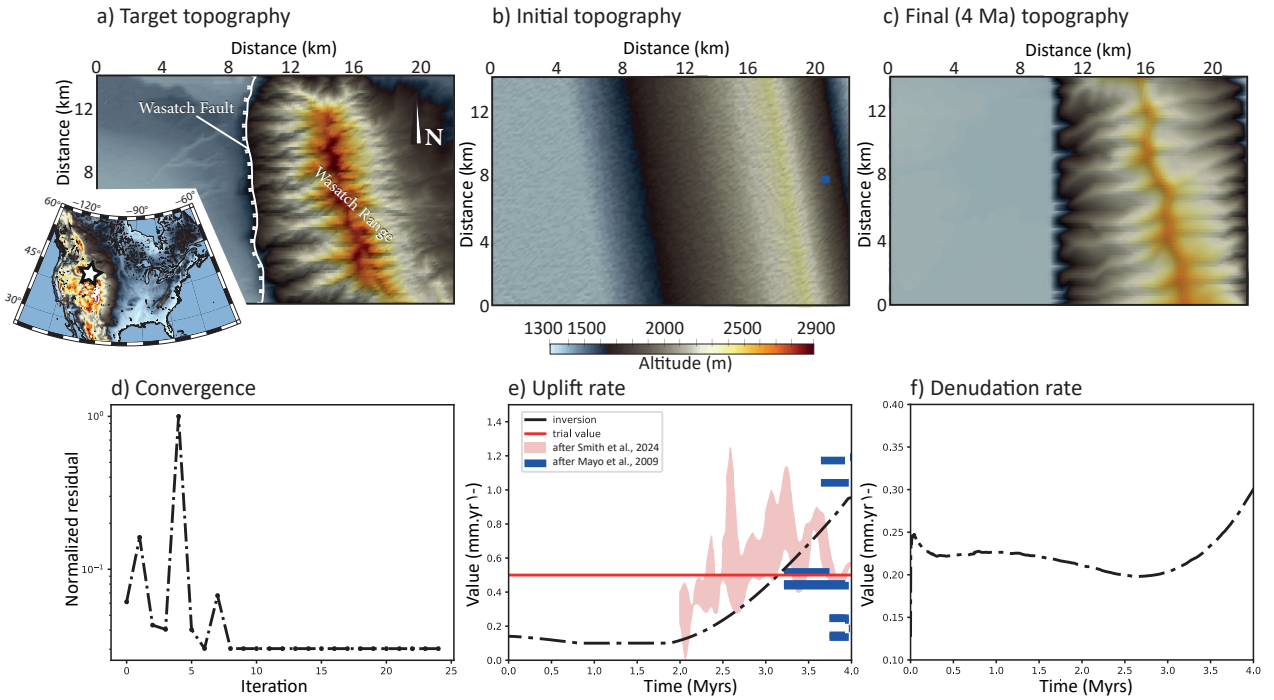
Figure 7. Inverse modelling of the Wasatch Range uplift rate (model 8, Table 1). (a): Topography of the Weber segment of the Wasatch fault used as the control parameter (white star on inset indicates its location in North America); (b): Initial topography conditions for the inverse model; (c): Modelled topography with uplift rates obtained by the inverse model; (d): model residuals; (e): Comparison between inverted uplift rates (this study) with those obtained from river longitudinal profile analysis (Smith et al., 2024), red enveloppe corresponding to 7 profiles, and from absolute dating of fluvial sediments in a cave (Mayo et al., 2009); (f): modelled denudation rate.

initial conditions provides important insights into the dominant drivers of landscape evolution in both generic and case-specific scenarios. Furthermore, this approach widens the scope of data-driven inverse modeling by enabling the integration of diverse

385 constraints, including time-dependent (e.g., denudation rates, sediment flux) and time-independent (e.g., final topography) observations into the inversion and to invert for a large number of parameters like the initial topography, or spatial variations κ or K_f.

For comparison, Yuan et al. (2019) conducted 100,000 forward model runs to constrain four spatially constant erosion parameters two of which define the minimum and maximum bounds of a temporally evolving fluvial incision parameter K_f

390 (which varies linearly over time). Pedersen et al. (2018) performed approximately 13,000 forward model runs to constrain between 2 and 5 free parameters related to uplift rate variations and erosion coefficient, the uplift parameters being defined within a prescribed spatial uplift zone and allowed to vary only across discrete time intervals. While such methods can identify the best-fit model at a reasonable computational cost, they lack the capacity to precisely quantify spatial sensitivity, that is, they cannot determine which specific topographic features most strongly influence the fit to observations.

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The adjoint method offers clear advantages for inverse modelling in landscape evolution. It provides an efficient and rigorous means of computing the sensitivity of a misfit function to all model parameters simultaneously, regardless of their number. This allows for a detailed understanding of how individual parameters influence model outcomes, without the need for extensive sampling of the parameter space through thousands of forward model runs, as required by gradient-free methods. Additionally, while incorporating more complex processes (e.g., sedimentation, isostatic rebound) significantly increases computational costs for gradient-free methods, the adjoint method remains computationally feasible as the forward model is only run once.

While solving erosion as a PDE with the adjoint method offers significant advantages, such as efficient gradient computation and sensitivity analysis, it has several limitations. In our case, when advection velocities become high, the drainage network can become numerically unstable, leading to hazardous propagation of local extrema. This issue stems from solving the stream power law as an advection equation rather than being specific to the adjoint method itself. A second constraint lies in the need for accurate, independent computation of drainage areas. The method requires that drainage directions align as closely as possible with the true topographic gradient, since upstream information propagation along channel networks depends on this alignment.

Another limitation of the adjoint approach is that it relies on gradient information, which means it is most effective when the initial parameter guess is reasonably close to the true solution. If the starting point is too far from the optimum, the method may converge slowly, get trapped in local minima, or fail to find a meaningful solution. In contrast, gradient-free methods, such as neighbourhood algorithms or Bayesian sampling, can explore the parameter space more globally without requiring a good initial guess, although at a higher computational cost.

Currently, our approach remains restricted to detachment-limited erosion scenarios where the SPL applies and where the slope exponent $n = 1$. Significant development work would be needed to extend this framework to more complex situations, such as transport-limited systems involving sediment deposition or flexural isostatic adjustment. Furthermore, incorporating additional constraints like cosmogenic exposure ages would enhance the method's applicability but still has to be implemented.

In this study, we have limited the model to simple inversions under two main assumptions: 1) only one type of parameter (even though spatially variable diffusivity or initial conditions involve as many parameters as there are grid nodes) is unknown, the other ones being perfectly determined; and 2) while the model can handle spatially variable erosion coefficients and initial topography, it cannot determine both temporal and spatial uplift rate variations. While such conditions are not always realistic in natural settings, this limitation is not unique to inverse problems; forward models of landscape evolution also face similar constraints. Like in forward models, another major tradeoff could arise from the under-determination of both the initial conditions and the uplift rate, or erosion coefficients and uplift rate in the case of an actively uplifting landscape, especially if the steady-state is not yet reached. For these reasons, joint inversion of multiple parameters such as the diffusion and erodibility coefficients, as well as source and initial conditions, could improve the procedure. The challenge, however, consists in incorporating sufficient data so that the joint inversion problem is not overly under-constrained.

Applying this method to natural cases demonstrates its potential to infer pre-incision topography or temporal variations in uplift rates, provided reasonable assumptions are made about the remaining parameters. In the case of the French Massif Central, while the age of the pre-incision topography remains uncertain, the recovered initial topography effectively highlights

430 a well-defined linear escarpment. This inversion proves relatively straightforward because the original fault escarpment has not been extensively degraded. Moreover, model sensitivity estimates highlight a domain with a large sensitivity both to diffusion and erodibility coefficients. This domain, located along the escarpment of the Mesozoic Cevennes fault system is presently characterized by intense seasonal rains called "Cevenol episodes" due to the conjonction of a topographic barrier to southerly winds carrying the humid and warm air of the mediterranean Sea (e.g., Delrieu et al., 2005). Assuming that the erodibility

435 coefficient K_f is directly related to the precipitation rate, we can deduce from the sensitivity maps that the Cévennes landscape evolution could be very sensitive to these extreme episodes. Finally, the inversion of temporal variations in uplift rates along a segment of the Wasatch Fault is feasible due to the strong dependency of the topography — particularly river longitudinal profiles — on the uplift rate. However, as regressive erosion progressively removes evidence of the earliest landscape history, it is virtually impossible to reconstruct the oldest phases of uplift. The observed discrepancies between uplift rates derived from

440 river profile inversion and our adjoint-based approach likely stem from several key factors. First, river profile methods rely on a steady-state assumption linking channel morphology to uplift history (Goren, 2016), an approximation that may not always be valid. Second, such approaches inherently focus on channel dynamics while neglecting hillslope processes and the rest of the landscape. Third, they typically assume temporally invariant drainage areas, ignoring catchment reorganization with time. In contrast, our adjoint method uses a reduced functional that minimizes the residual between modeled and observed topography,

445 ensuring consistency with the range-scale morphological evolution. While this approach better captures the average landscape response, it naturally reduces the importance of second-order features like knickpoints that record abrupt uplift changes.

6   Conclusion

A finite-element model that combines the Stream Power Law of fluvial incision with linear hillslope diffusion into a classic diffusion-advection equation offers several advantages. First, it enables a consistent treatment of the "slope" responsible for

450 both the downslope movement of the surficial altered layer (simulating hillslope diffusion) and water flow along channels (fluvial incision). In this approach, both processes rely on the same topographic gradient, providing a more physically realistic framework than the classical models that use a local channel slope derived from preferred drainage directions. Second, it enables the solution of the adjoint problem to invert unknown parameter values and assess the model sensitivity to these parameters, which is difficult to do with only forward models and gradient-free procedures.

455 Our results demonstrate that it is possible to invert spatial variations of the diffusion and erodibility coefficients, and to evaluate the landscape sensitivity to these parameters. Additionally, simple uplift scenarios involving constant or smoothly varying rock uplift rates can be effectively determined. When addressing the initial conditions of a degrading escarpment, some regularization parameters are required, yet this approach still yields consistent results with respect to the forward model.

Applications to real-world data, such as the pre-incision topography of the SE French Massif Central and the uplift rate of

460 the Wasatch Fault in Utah, suggest that this methodology holds potential for natural case studies. In particular, we show that the dissected, linear escarpment bounding the SE part of the French Massif Central is more sensitive to hillslope diffusion and river incision than the rest of the domain. In the same region, our results suggest that, prior to intense incision by SE flowing

rivers, the topography was delineated by a linear NE-SW escarpment that coincides with the trace of a major Mesozoic fault. Concerning the Wasatch fault, slip rate inversion depicts a first period of low slip rate, that has increased dramatically 2 Myrs ago.

7   Acknowledgements

8   Code availability

The code, parameter and data files are available on bitbucket (git@bitbucket.org:jourdon_anthony/laser.git).

9   Competing interests

The authors declare no competing interests.

10   Authors contribution

CP developed the code, analyzed the results and wrote the manuscript. AJ developed the code, discussed the results and participated in writing the manuscript. NC discussed the results and participated in writing the manuscript.

11   Appendix A

11.1   Taylor test of the adjoint model

We test the consistency of the adjoint model using a Taylor test on an initial forward model consisting of a $40 \text{ km} \times 40 \text{ km}$ domain, with a $20 \text{ km} \times 20 \text{ km}$ square mountain at the center uplifting at a constant rate of 0.5 mm.yr^{-1}. The Taylor test verifies that:

$$J(m + \epsilon \partial m) - J(m) - \epsilon \nabla J \cdot \partial m \to 0 \quad \text{at} \quad \mathcal{O}(\epsilon^2), \tag{22}$$

where m is the control parameter of interest, ∂m is a small perturbation of this parameter, and ϵ is a small value. The Taylor test consists of applying a perturbation to the initial erodibility coefficient K_f, computing the second-order Taylor remainder using Eq. (22), then successively decreasing ϵ and verifying that the residual decreases at the expected order of ϵ^2.

The results of the Taylor test applied to this model show that the residuals are consistent with theoretical predictions (Fig. 8), exhibiting a convergence rate close to 2 down to ϵ values of $\sim 1 \times 10^{-16}$. Beyond this threshold, the convergence deteriorates,
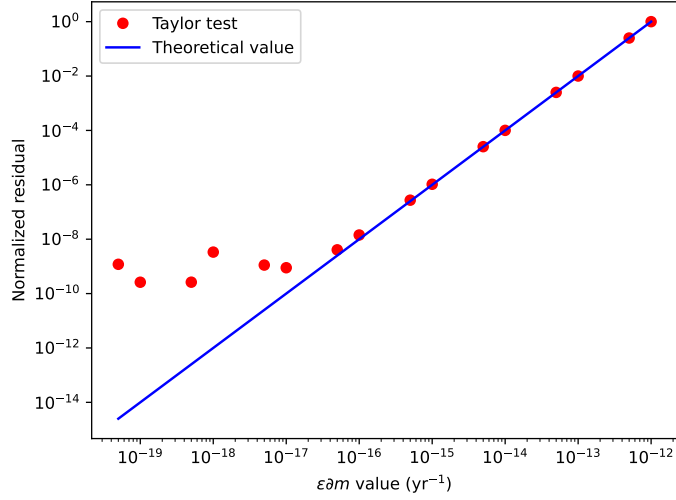
Figure 8. Taylor test on the erodibility coefficient K_f using an initial ϵ value of 1×10^{-6} (upper right point) with a constant ∂m value equal to $1 \times 10^{-6}\mathrm{yr}^{-1}$. Model 9, Table 1.

and the residual fluctuates between 10^{-10} and 10^{-8}. This result indicates that the adjoint model is robust down to machine precision (Fig. 8).

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