

# Constraining a Radiative Transfer Model with Satellite Retrievals: Contrasts between Cirrus Formed via Homogeneous and Heterogeneous Freezing and their Implications for Cirrus Cloud Thinning

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Ehsan Erfani<sup>1</sup> and David L. Mitchell<sup>1</sup>

<sup>1</sup>Division of Atmospheric Sciences, Desert Research Institute, Reno, Nevada, USA

Correspondence to: Ehsan Erfani, ([Ehsan.Erfani@dri.edu](mailto:Ehsan.Erfani@dri.edu))

## 10 Abstract

The efficacy of the climate intervention method known as cirrus cloud thinning (CCT) is difficult to evaluate in climate models, largely due to uncertainties governing the relative contributions of homogeneous and heterogeneous ice nucleation. Here we take a different approach by employing recent satellite retrievals from the Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observation (CALIPSO) which provide estimates of the fraction of cirrus clouds dominated by homogeneous and heterogeneous ice nucleation and their associated physical properties. We employ a radiative transfer model (RTM) to quantify the cloud radiative effect for homogeneous and heterogeneous cirrus clouds at the top of atmosphere (TOA), Earth's surface, and within the atmosphere. The RTM experiments are initialized using cirrus microphysical profiles derived from CALIPSO retrievals for cirrus clouds dominated by homogeneous and heterogeneous ice nucleation across different regions (Arctic, Antarctic, and midlatitude) and surface types (ocean and land). We define two bounds: the lower bound assumes a full microphysical transition from the observed composition of homogeneous- and heterogeneous-dominated cirrus to only heterogeneous cirrus and production of new cirrus. The upper bound assumes production of new cirrus and that the atmospheric dynamics enables homogeneous freezing nucleation to occur regardless of the concentration of ice nucleating particles. Based on these bounds, we estimate an instantaneous surface effect ranging from  $-0.5$  to  $+0.6$  W m<sup>-2</sup> and a TOA effect from  $-0.9$  to  $+1.1$  W m<sup>-2</sup>, respectively, showing the possibility of both cooling and warming. Recommendations are provided to improve the treatment of cirrus clouds in climate models.

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# 1 Introduction

Cirrus clouds are a critical component of the Earth's radiation budget; the global annual mean coverage of these clouds ranges from 17-20% (Matus and L'Ecuyer, 2017; Sassen et al., 2009) to 35% (Hong et al., 2016) with high spatial variability. Cirrus cloud coverage is about 30% in mid-latitudes and about 60-80% in the tropics (Guignard et al., 2012; Stubenrauch et al., 2006). In addition, cirrus clouds are more frequent during the winter seasons in the mid and high latitudes (Mitchell et al., 2018; Zhu et al., 2024). They significantly absorb and scatter incoming solar radiation and absorb outgoing thermal radiation from the Earth's surface and low-level clouds. Although these two effect counteract each other, it is estimated that on global annual averages, these clouds warm the planet by approximately  $5 \text{ W m}^{-2}$  (Gasparini and Lohmann, 2016). Despite their significant impacts on radiation and climate, uncertainty exists in measuring, retrieving, and modeling cirrus clouds partly because the processes involved in their formation are poorly understood (Heymsfield et al., 2017) or are not represented in climate models (Lyu and Liu, 2023). This complexity has left many important questions unanswered (Kärcher, 2017; Kay et al., 2012). In particular, our understanding of the mechanisms of cirrus cloud development and their microphysical properties, such as ice crystal shape and size distribution remain insufficient (Krämer et al., 2016; Lawson et al., 2019). Cirrus clouds exhibit diverse geometric features (Fig. 1), which reflect their varied microphysical and macrophysical properties.

One of the main uncertainties in modeling cirrus clouds is related to insufficient knowledge of the relative contribution of homogeneous and heterogeneous ice nucleations in cirrus clouds (Heymsfield et al., 2017). Homogeneous ice nucleation happens when liquid solution droplets (haze or cloud droplets) freeze spontaneously, with no ice nucleating particles (INPs) to initiate freezing. This is when the temperature ( $T$ ) is colder than  $-38 \text{ }^{\circ}\text{C}$  and supersaturation (quantified by relative humidity with respect to ice or  $\text{RH}_i$ ) is greater than 140-150%. In contrast, heterogeneous ice nucleation requires INPs to initiate freezing at  $T < 0 \text{ }^{\circ}\text{C}$  and lower  $\text{RH}_i$  values (Heymsfield et al., 2017; Kanji et al., 2017). Since INP concentrations are generally much lower than solution droplet concentrations, heterogeneous cirrus usually have fewer and larger ice particles, and therefore are optically thinner, whereas homogeneous cirrus generally contain higher ice particle concentrations of smaller size, and are optically thicker (Krämer et al., 2016; Mitchell and Garnier,

2024). With such distinct microphysical properties, these two types of cirrus clouds demonstrate significantly different radiative effects, and this makes it crucial to investigate their contributions.

There are different methods to retrieve cirrus cloud properties using satellite instruments such as infrared radiometers (Magurno et al., 2020; Mitchell et al., 2018; Nazaryan et al., 2008; Stubenrauch et al., 2008; Yue et al., 2020), visible radiometers (Gao et al., 2002; Wang et al., 2019), microwave radiometers (Evans et al., 2012; Jiang et al., 2019; Wu et al., 2014), and a combination of instruments (Yorks et al., 2023). Satellite microwave radiometers have been used widely to retrieve cirrus clouds, however, their coarse spatial (Wang et al., 2001) and temporal (Jiang et al., 2019) resolutions, the sensitivity of the retrievals to surface reflectivity (Wang et al., 2001), and the need for ancillary information from the surface to properly estimate the surface albedo (Jiang et al., 2019) limit their ability for studying the cirrus clouds. Visible retrievals also have limitations such as low sensitivity to detecting cirrus clouds (especially, thin ones since they have low reflectivity and absorption in the visible range) and contamination of land surface reflectance (Schläpfer et al., 2020). On the other hand, infrared retrievals have a much lower sensitivity to surface reflectivity and can detect thin cirrus clouds using water vapor absorption bands (Roskovensky and Liou, 2003).

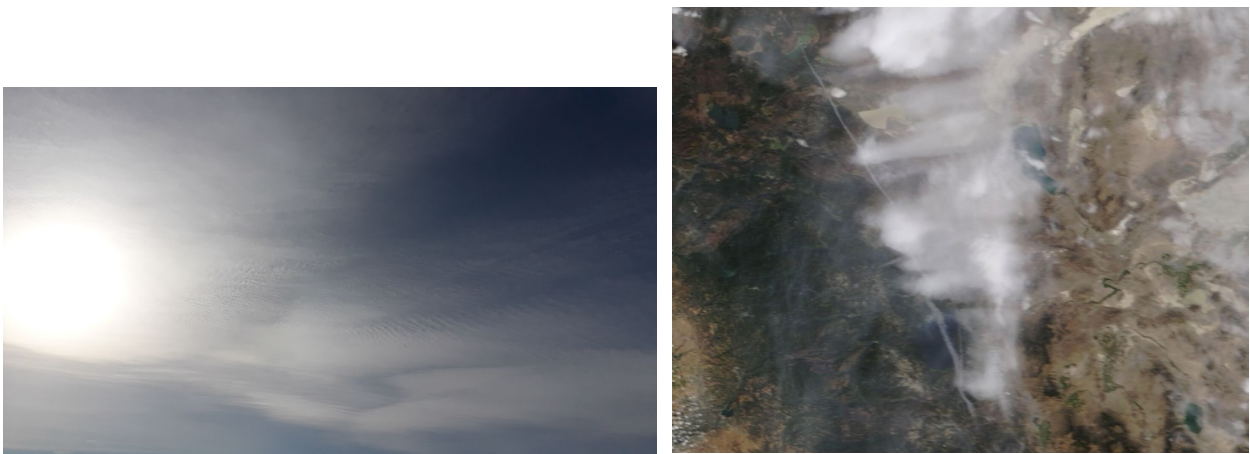


Figure 1. Left: Photography of sky over Reno, Nevada, USA on 25 Sep. 2023, showing cirrus clouds with various geometric features (e.g., thin and thick) (Photo taken by Ehsan Erfani). Right: Satellite imagery showing the same types of cirrus on the same day. Reno is located between Lake Tahoe and Pyramid Lake and is covered by clouds. Note that the two photos do not correspond to the same time, but provide general cloud patterns on the same day (the satellite image provided by MODIS instrument onboard NASA Terra satellite and taken from NASA Worldview website: <https://worldview.earthdata.nasa.gov/>).

The Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observations (CALIPSO) dataset has been used to study cirrus cloud properties (Li and Groß, 2021; Sassen et al., 2009). It also has some limitations; for instance, lidar-radar (DARDAR) retrievals of the ice particle number concentration ( $N_i$ ) are based on assumptions about the shape of the ice particle size distribution, which can lead to uncertainties in the retrieved values (Sourdeval et al., 2018). Despite this, the CALIPSO dataset remains a valuable tool for studying cirrus clouds and their radiative impacts on climate. Recently, Mitchell and Garnier (2024) expanded on Mitchell et al. (2024) work and developed a CALIPSO retrieval to quantify homogeneous and heterogeneous cirrus on a global scale (note that the accurate terms would be “dominated by homogeneous” and “dominated by heterogeneous” ice nucleation regimes, but for simplicity, we use the terms homogeneous and heterogeneous in this study). The data from two Infrared Imaging Radiometer (IIR) channels, 10.6  $\mu\text{m}$  and 12  $\mu\text{m}$ , along with CALIPSO lidar measurements pertaining to cloud top and cloud base, were used to calculate ice optical and microphysical properties, such as  $N_i$ , ice water content (IWC), effective diameter ( $D_e$ ), and shortwave extinction coefficient ( $\alpha_{ext}$ ) using ice particle mass-dimension relationships from Erfani and Mitchell (2016). To establish a threshold transition between homogeneous and heterogeneous cirrus regimes (henceforth, referred to as cirrus regimes), they considered the  $D_e$  maximum in the  $\alpha_{ext} - D_e$  plane as this threshold (note that high  $N_i$  should limit ice particle growth and  $D_e$  due to increased competition for water vapor). In particular, they showed that although heterogeneous cirrus is dominant in most regions and seasons, the homogeneous fraction weighted by cloud optical depth contributes more than 50% during the winter in the extratropics.

The findings by Mitchell and Garnier (2024) have important implications for a *climate intervention* technique called *cirrus cloud thinning* (CCT). Climate change has disastrous effects on humans, the environment, and society, and such effects exacerbate as global  $\text{CO}_2$  level and sea surface temperature (SST) increase (IPCC report, 2021). The last time with  $\text{CO}_2$  concentrations near 400 ppm was during the mid-Pliocene (3.25 million years ago) when global SST was 4.1°C warmer than the preindustrial period (Tierney et al., 2025). Global climate models (GCMs) project that global warming will continue in the next decades (IPCC report, 2021), and even in the unlikely scenario where global greenhouse gas (GHG) emissions are eliminated by 2050 (Forster et al., 2021; Hansen et al., 2023; 2025), the global mean temperature would remain around its 2050 value

for centuries unless atmospheric GHG concentrations were decreased somehow. This has prompted some to advocate for a threefold solution: (1) GHG emission reductions, (2) GHG concentration reduction, and (3) climate interventions to cool the planet (Baiman et al., 2024). Solution (3) would take only several years to act, whereas solutions (1) and (2) would take several decades and thus risk triggering tipping points in the climate system (e.g., Steffen et al., 2018). Therefore, various climate intervention methods, including CCT (Mitchell and Finnegan, 2009; Storelvmo et al., 2014), have been proposed to cool the planet (NASEM report, 2021). It is important to conduct comprehensive research on climate intervention methods in order to quantify their efficacy, cost, risks, and limitations. Climate intervention methods, if proven effective, are not replacements for but rather complement GHG emission reduction and removal.

CCT is a proposed climate intervention method often considered under the Solar Radiation Modification (SRM) category and is suggested to deliberately slow down the warming of the planet by injecting proper aerosols that act as ice nucleating particles (INPs) in the upper troposphere to reduce the thickness and coverage of cirrus clouds (Mitchell and Finnegan, 2009).

CCT can be efficient and cool the planet if the homogeneous cirrus is abundant, leading to a “transition from homogeneous to heterogeneous cirrus” (Note: throughout this study, this phrase refers to the concept that the presence of INPs, either through deliberate injection for CCT purposes or through natural and anthropogenic aerosols, can shift the ice nucleation pathway from homogeneous toward heterogeneous, potentially modifying cirrus radiative effects). Heterogeneous cirrus is considered to be dominant outside of tropics (Cziczo et al., 2013; Froyd et al., 2022), but recent satellite retrievals (Gryspeerd et al., 2018; Mitchell et al., 2018; Mitchell and Garnier, 2024) have shown that homogeneous cirrus might have been underestimated. The effectiveness of CCT might surpass previous estimates, considering that the cooling efficacy of CCT depends on the fraction of homogeneous cirrus. CCT should be most impactful in the high latitudes during the period having relatively less daylight because the cirrus longwave (LW) cloud radiative effect (CRE) is significantly stronger than shortwave (SW) CRE, and therefore significant surface cooling could happen. Efficient CCT has the potential to reduce the thawing of Arctic permafrost and to enhance the sea ice cover (Storelvmo et al., 2014), and thus enhance the Atlantic Meridional Overturning Current (AMOC) by cooling sea surface temperatures to promote downwelling just south of Greenland. Note that the AMOC is a climate tipping point (Steffen et

al., 2018). Moreover, CCT could slow down Arctic amplification (AA), a phenomenon characterized by warming of the Arctic at a rate two to four times faster than the rest of the globe mainly because of sea ice loss (Rantanen et al., 2022; Screen and Simmonds, 2010).

Despite the cooling potential of CCT from theory (e.g., Lohmann and Gasparini, 2017; Mitchell and Finnegan, 2009), the results of modeling studies on CCT are not conclusive as some CCT simulations indicated that CCT cooling is negligible (Gasparini & Lohmann, 2016; Penner et al., 2015; Tully et al., 2022) while others (Gruber et al., 2019; Storelvmo et al., 2013, 2014) showed that such cooling is significant. GCMs and regional climate models (RCMs) have significant uncertainties in predicting the microphysical properties of cirrus clouds largely because of limitations in capturing the complicated set of under-resolved physical mechanisms associated with cirrus clouds and their interactions with aerosols (Eliasson et al., 2011; Kay et al., 2012; Maciel et al., 2023; Patnaude et al., 2021). Some possible ways for improving the treatment of CCT in GCMs are described in Mitchell and Garnier (2024) and in Sect. 5 of this study. For this reason, it is important to constrain models with observations to achieve a better understanding of cirrus clouds in general and CCT in particular.

An additional concern in the context of CCT is the risk of “overseeding,” where excessive injections of INPs could lead to too many small ice crystals, increasing the optical thickness and the lifetime of cirrus clouds, and thus causing a net warming effect instead of cooling (Gasparini and Lohmann, 2016; Penner et al., 2015). Another potential aspect of overseeding is the formation of “new cirrus” due to INPs injected into clear-sky ice-supersaturated regions (Tan et al., 2016). Observational evidence indicates that stratospheric plumes of enriched INP concentration from volcanic eruptions, upon entering the troposphere, can increase cirrus cloud cover by about 20% (Lin et al., 2025; Sporre et al., 2022), suggesting that CCT seeding may have a similar impact. The extent to which this “new cirrus” effect might offset or even dominate the intended homogeneous-to-heterogeneous transition remains unknown. However, in this study, we address this potential counteracting mechanism.

To evaluate CCT’s cooling potential without the use of climate models, a radiative transfer model (RTM) is employed in this study. Over the past decades, RTMs have been used extensively to study the radiative properties of cirrus, contrail, and mixed-phase clouds, since RTMs are the most

175 accurate tools for calculating radiative fluxes when ice cloud microphysical fields are measured  
 (which is difficult to reproduce in a complex GCM). RTMs have been used to determine heating  
 rates and/or the radiative effect of ice clouds, with their microphysical characteristics sometimes  
 measured during aircraft field campaigns (Marsing et al., 2023), retrieved from satellite  
 measurements (Hong et al., 2016; Sun et al., 2011), or simulated by models such as box-models  
 180 (Cirisan et al., 2013) or models used as stochastic cloud generators (Fauchez et al., 2017; Zhou et  
 al., 2017) or a mesoscale cloud model complex (Khvorostyanov and Sassen, 1998). RTM  
 simulations of cirrus clouds show that their radiative effects are highly sensitive to cloud  
 microphysical characteristics such as ice water path (Córdoba-Jabonero et al., 2020; Fu and Liou,  
 1993), and ice particle shape and size (Macke et al., 1998; Takano et al., 1992; Zhang et al., 1999).  
 185 A few studies (e.g., Schumann et al., 2012; Wolf et al., 2023) considered multiple microphysical  
 and environmental parameters (e.g., temperature, surface albedo, solar zenith angle) when  
 computing the radiative effect of cirrus and contrails. Despite significant progress in calculating  
 cirrus cloud radiative properties by using an RTM, the contribution of homogeneous and  
 heterogeneous cirrus to the total cirrus CRE and the efficacy of CCT has not been studied yet.

190 This study aims to combine new advances in satellite remote sensing and radiative transfer  
 modeling to develop a conceptual platform for studying different types of cirrus clouds and their  
 impact on Earth's energy budget. We use the novel CALIPSO satellite retrievals from Mitchell et  
 al. (2024) to infer the microphysical properties of cirrus clouds (e.g., IWC and  $D_e$ ) and then employ  
 those as inputs to an RTM to calculate cirrus CREs. This is done by calculating the vertical profiles  
 195 of IWC and  $D_e$  for two types of cirrus clouds (homogeneous and heterogeneous) and different  
 environmental conditions (latitude bands, surface types, seasons) based on CALIPSO retrievals.  
 These are then used in an RTM to calculate cirrus cloud CRE at the surface (Sfc), at top of the  
 atmosphere (TOA), and in the column of atmosphere (Atm). By investigating the difference in  
 CRE between homogeneous and heterogeneous cirrus, this study provides estimated bounds of the  
 200 efficacy of CCT as a first estimate, with implications for improving GCMs. This study is  
 specifically focused on the Arctic and Antarctic during the cold season because these are  
 conditions which (i) homogeneous cirrus occurrence is highest, and (ii) the CCT intervention is  
 expected to have the largest radiative impact due to zero or very weak solar radiation. This targeted  
 design within an RTM framework was intended to support a process-level understanding of cirrus

radiative effects and the implications for CCT. The rest of this paper is organized as follows: in Section 2, a description of the observational data and RTM experimental design is presented; the main RTM results are explained in Section 3 for relevant geographical conditions; the sensitivity to thermodynamic profiles, low clouds, and aerosols are explored in Section 4; suggestions for improving cirrus cloud modeling of CCT is provided in Section 5; and finally, conclusions are presented in Section 6.

## 2 Methodology

### 2.1 Data

The RTM requires the vertical profiles of atmospheric variables and trace gases as inputs and by default, uses available standard profiles for the tropics, mid-latitude, sub-arctic, and U.S. regions for winter and summer seasons and from surface to 120 km provided by Air Force Geophysical Laboratory (AFGL) atmospheric constituent dataset (Anderson et al., 1986). The radiative impacts of trace gases are small, so we use the standard vertical profiles of trace gases. However, the cirrus cloud properties are closely related to thermodynamic profiles, in particular temperature ( $T$ ). Therefore, to force the RTM with realistic thermodynamic profiles, we replace the standard vertical profiles of  $T$  and water vapor mixing ratio ( $q_v$ ) with those extracted from Modern-Era Retrospective Analysis for Research and Applications, version two (MERRA2; Gelaro et al., 2017) reanalysis dataset with a spatial resolution of  $0.5 \times 0.625^\circ$ , 72 vertical levels, and a temporal resolution of 1 month. Using this dataset is preferred because it was also used in the CALIPSO satellite retrievals of homogeneous and heterogeneous cirrus clouds. The RTM requires air density ( $\rho_a$ ) to be consistent with thermodynamic profiles, therefore, we calculate  $\rho_a$  based on MERRA2  $T$  and pressure ( $P$ ) following the ideal gas law:  $\rho_a = P/kT$ , where  $k$  is Boltzmann constant. This new  $\rho_a$  then replaces the default  $\rho_a$ . The area-weighted averages of  $T$ ,  $q_v$ , and  $\rho_a$  profiles are calculated for grid points in the Arctic ( $60$ - $90^\circ\text{N}$ ), Antarctic ( $90$ - $60^\circ\text{S}$ ), and the Northern Hemisphere (NH) mid-latitude ( $30$ - $60^\circ\text{N}$ ), and for winter seasons of the same years as the CALIPSO retrievals (2008, 2010, 2012, and 2013). In addition, maximum and minimum profiles in each region are calculated

as a range of change in thermodynamic variables (Fig. 2). Using RTM standard sub-arctic profiles are not justified, because they over-estimate the cold and dry profiles over the Arctic.

The CALIPSO satellite retrievals based on the methodology of Mitchell et al. (2024) and Mitchell and Garnier (2024) are used to create cirrus cloud property statistics (e.g., median and 25<sup>th</sup> and 75<sup>th</sup> percentiles) for each season, latitude band, and surface type (land or ocean). In addition, the data is grouped into homogeneous and heterogeneous cirrus categories, based on temperature-dependent  $\alpha_{ext}$  thresholds derived from  $D_e$  maxima (related to the  $\alpha_{ext}$ ) as established by those studies. The reader is advised to check Mitchell and Garnier (2024) for a detailed explanation of the method for discriminating between heterogeneous and homogeneous cirrus clouds, but we can say that the microphysical properties of the latter are strongly affected by homogeneous nucleation.

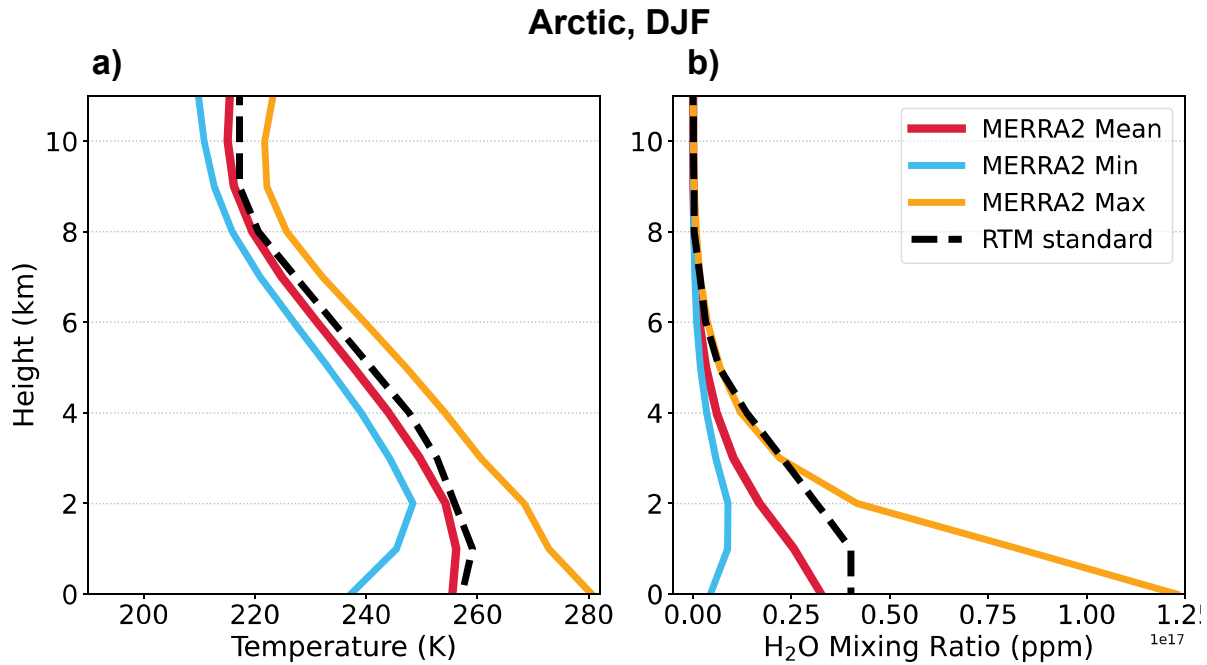


Figure 2. Vertical profiles of a) temperature and b) water mixing ratio for wintertime. The libRadtran RTM standard profiles are for subarctic (no Arctic/Antarctic profile provided), whereas MERRA2 profiles are for the Arctic region (60-90°N) during the boreal winter of 2008, 2010, 2012, and 2013. Mean refers to area-weighted average over all grid points in this region.

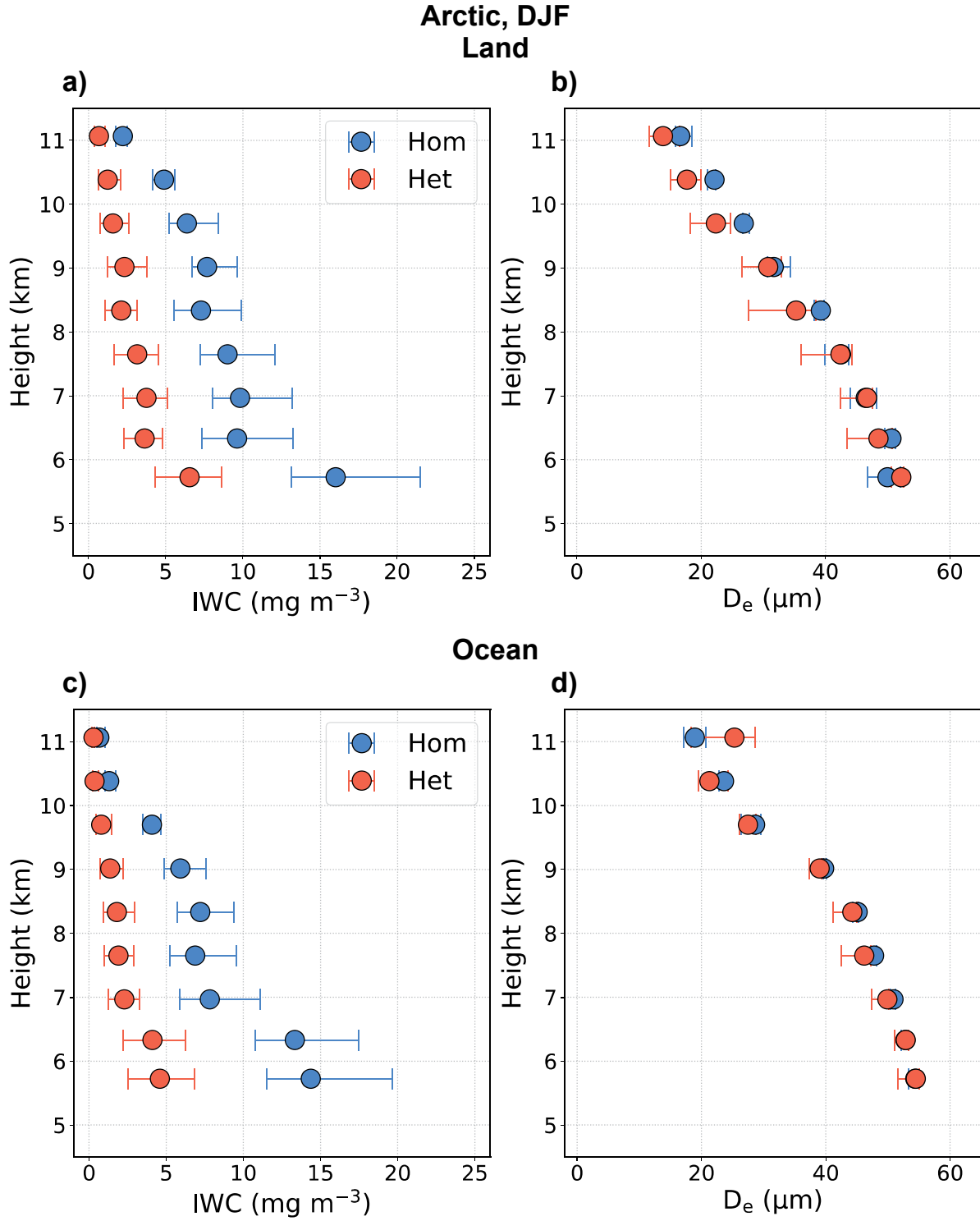


Figure 3. Microphysical properties of cirrus clouds from CALIPSO retrievals: a&c) IWC vs. height and b&d)  $D_e$  vs. height for two cirrus regimes (homogeneous and heterogeneous). The results are for Arctic (60-90°N) during boreal winter (DJF) of 2008, 2010, 2012, and 2013 and for two different surface types: a-b) land and c-d) ocean. Markers show median values, whereas error bars show 25<sup>th</sup> and 75<sup>th</sup> percentiles.

Figure 3 shows an example of this analysis for IWC and  $D_e$  vs. height over the Arctic during the December-January-February (DJF) period. Note that each panel presents a compilation of numerous cirrus cloud samples for various heights, grid points, and days, and therefore, it is not correct to assume that it represents a single cirrus from the lowest to highest height shown. For practical purposes, the IWC and  $D_e$  apparent “profiles” from the lowest to highest height for each cirrus regime are divided into 4 clouds each having a thickness of  $\sim 1.3$  km (typical thickness of cirrus clouds; Dowling and Radke, 1990; Gouveia et al., 2017), but with different cloud base and top heights (CBHs and CTHs). Each of these clouds with their respective IWC and  $D_e$  profiles (with an approximate vertical resolution of 50 m) are then used as input to an RTM to simulate the radiative properties for that cloud.

## 2.2 Radiative Transfer Model (RTM)

In this study, the calculations of various thermal or LW fluxes and solar or SW fluxes are conducted using an RTM termed library for Radiative transfer (libRadtran), which employs "uvspec" as its main core (Emde et al., 2016). For simplicity, we refer to libRadtran uvspec as RTM in the rest of this paper. The RTM solver is selected to be the one-dimensional Discrete Ordinate Radiative Transfer model (DISORT; Stamnes et al., 2000; Buras et al., 2011) with six streams. The spectral wavelength range is from 0.25  $\mu\text{m}$  to 5  $\mu\text{m}$  for SW and from 3.1  $\mu\text{m}$  to 100  $\mu\text{m}$  for LW radiation. In addition, the REPTRAN parameterization with fine resolution is selected to account for molecular absorption (Gasteiger et al., 2014).

The RTM has the option to calculate the radiative impact of clouds based on the vertical profiles of cloud water content and effective radius ( $r_e$ ) which are provided as inputs. Ice and liquid cloud properties need to be specified separately in the RTM input files. To calculate the cloud optical properties from IWC and  $r_e$  in the RTM, we specify the Baum parameterization (Baum et al., 2005) with the assumption of a general habit mixture (GHM). The GHM consists of a mixture of different ice particle shapes or habits (e.g. columns, plates, bullet rosettes, aggregates) that vary with particle size. This allows for a more realistic representation of the ice particles since cirrus clouds consist of a wide range of ice habits and sizes (Erfani and Mitchell, 2016, 2017; Lawson et al., 2019). The liquid cloud parameterization of RTM follows the method of Hu and Stamnes (1993). The preparation of variables required for the atmospheric profile file is explained in Sect. 2.1.

By turning on the aerosols option in the RTM, we select the fall-winter season and the maritime haze for the atmosphere below 2 km (as boundary layer or BL) and the background for the atmosphere above 2 km (as free troposphere or FT), following the aerosol model of Shettle (1989) for the main RTM simulations. The broadband thermal emissivity ( $\varepsilon$ ) varies based on the surface type. Although the  $\varepsilon$  value of snow and ice surfaces is very close to that of a blackbody (equal to unity), it is approximately 0.99 for ocean and forest, and lower for surface types such as cropland, shrubland, and deserts (Wilber et al., 1999). Nonetheless, the sensitivity of LW fluxes to  $\varepsilon$  is much smaller than that to temperature based on Stefan–Boltzmann law. Therefore, we use an  $\varepsilon$  value of unity throughout this study but conduct simulations to investigate the sensitivity to temperature.

Table 1. A summary of RTM runs conducted in this study.

Experiment	Region	Season	Surface type	Radiation	Cirrus cloud regimes	Number of simulations
Main runs using CALIPSO IWC and $D_e$ (median, upper quartile, and lower quartile profiles)	Arctic	DJF	Land	LW	Hom, Het, Clr	25
	Arctic	DJF	Ocean	LW	Hom, Het, Clr	24
	Antarctic	JJA	Land	LW	Hom, Het, Clr	25
	Antarctic	JJA	Ocean	LW	Hom, Het, Clr	24
	NH midlatitude	DJF	Land	LW	Hom, Het, Clr	25
	NH midlatitude	DJF	Land	SW	Hom, Het, Clr	25
Sensitivity to meteorology (min and max $T$ and $q_v$ profiles)	Arctic	DJF	Land	LW	Hom, Het, Clr	16
Sensitivity to low clouds (with three LWC values)	Arctic	DJF	Land	LW	Hom, Het, Clr	24
Sensitivity to aerosols (two BL and two FT options)	Arctic	DJF	Land	LW	Hom, Het, Clr	32
						<b>Total:</b> 220

A summary of RTM experiments in this study is provided in Table 1. A total of 220 simulations are conducted for various regions (Arctic, Antarctic, NH midlatitude), surface type (land and ocean), and different upper-level cloud conditions (homogeneous, heterogeneous, and clear sky).

Furthermore, we explore sensitivity to low liquid clouds, thermodynamic profiles, and atmospheric aerosols. In order to test the impact of low liquid cloud, we add a layer from 500 m to 1100 m (thickness of 600 m) with cloud droplet  $r_e$  of 7  $\mu\text{m}$ . These values are consistent with field measurements of low clouds over the Arctic Ocean and Greenland (Järvinen et al., 2023). Three low liquid clouds are tested by varying liquid water content (LWC): 0.01, 0.03, and 0.05  $\text{g m}^{-3}$ . To investigate the effect of thermodynamic profiles, we use the maximum and minimum  $T$  and  $q_v$  profiles in the Arctic during the winter (Fig. 2) and conduct RTM sensitivity tests. Also, four different aerosol options are explored for RTM sensitivity to aerosols: “marine haze, low volcanic”, “urban haze, low volcanic”, “marine haze, high volcanic”, and “urban haze, high volcanic”.

### 2.3 Cloud Radiative Effect

The change in radiative fluxes caused by cirrus clouds is quantified by the CRE following Loeb et al. (2009):

$$\text{CRE}_{\text{LW}_z} = (\text{LW } \downarrow_{z_{\text{cld}}} - \text{LW } \uparrow_{z_{\text{cld}}}) - (\text{LW } \downarrow_{z_{\text{clr}}} - \text{LW } \uparrow_{z_{\text{clr}}}), \quad (1)$$

where  $z$  refers to a specific height (which is either TOA or Sfc in this study), arrows indicate upward or downward fluxes, “cld” refers to the cloudy condition, and “clr” refers to the clear-sky condition. Each term is in units of  $\text{W m}^{-2}$  and all the radiative fluxes in the right-hand side of the above equation are the outputs of the RTM. As shown in Eq. (1), we consider downward fluxes as positive and vice versa throughout this study. The CRE in the Atm is calculated as:

$$\text{CRE}_{\text{LW}_{\text{Atm}}} = \text{CRE}_{\text{LW}_{\text{TOA}}} - \text{CRE}_{\text{LW}_{\text{Sfc}}}. \quad (2)$$

A similar set of equations is used to derive the SW CRE. In our RTM study, we use  $\text{CRE}_{\text{LW}_{\text{Sfc}}}$  to estimate the instantaneous effect of cirrus clouds, while  $\text{CRE}_{\text{LW}_{\text{Atm}}}$  represents the cirrus effect that could potentially influence the surface over longer timescales through adjustment and feedback processes. The net CRE is defined as:

$$\text{CRE}_{\text{net}_z} = \text{CRE}_{\text{LW}_z} + \text{CRE}_{\text{SW}_z}, \quad (3)$$

which can be calculated for the TOA, Sfc, or Atm.

### 2.3.1 CCT under ideal microphysical change

In this study, we define the lower bound of CCT efficacy (cooling effect) under the assumption of a complete microphysical transition from the observed mixture of homogeneous and heterogeneous cirrus clouds to heterogeneous cirrus. This bound represents an idealized condition where an increase in available INPs due to seeding enables heterogeneous freezing to completely suppress homogeneous nucleation. We assume that the cirrus clouds then form under the microphysical conditions typically associated with natural heterogeneous cirrus, e.g., conditions that generally result in lower IWC than in homogeneous cirrus. The derived IWC and  $D_e$  profiles for heterogeneous and homogeneous regimes are based on CALIPSO retrievals (Fig. 3). This idealized bound enables us to quantify the maximum cooling impact of CCT, using the net CRE difference between these two regimes. We calculate this as:

$$\Delta\text{CRE} = \langle \text{CRE}_{\text{net},\text{het}} - \text{CRE}_{\text{net},\text{hom}} \rangle, \quad (4)$$

where angle brackets show the average for the four cirrus clouds at 4 different altitudes, as explained in Sect. 2.1. Note that  $\Delta\text{CRE}$  is based on the ideal assumptions that cirrus cloud overcast condition exists. Therefore, correction factors are required to estimate a more realistic impact:

$$\Delta\text{CRE}_{\text{max}} = \Delta\text{CRE} \times \text{CF}_{\text{cirrus}} \times F_{\text{hom}}, \quad (5)$$

where  $\Delta\text{CRE}_{\text{max}}$  indicates that new cirrus cloud formation is not accounted for, and  $\text{CF}_{\text{cirrus}}$  is cirrus cloud fraction and  $F_{\text{hom}}$  is fraction of homogeneous cirrus clouds. The CALIPSO cirrus cloud analysis of Mitchell and Garnier (2024) does not explicitly provide values of  $\text{CF}_{\text{cirrus}}$ .

Therefore, we use a typical value of 35% for extratropical regions (Gasparini et al., 2023). This estimate may be conservative for the polar regions during winter when ice cloud coverage is greater than in other seasons (Hong et al., 2016; Mitchell et al., 2018; Sassen et al., 2009). The retrievals provide vertical profiles of the homogeneous fraction (defined as the number of homogeneous cirrus pixels divided by the total number of cirrus pixels) for different regions and seasons as shown in Fig. 4. Strong variability is seen in homogeneous fraction with height, region (Arctic, Antarctic, and midlatitude), and surface type (land and ocean) and this makes it important to conduct a different RTM simulation for each of those geographical conditions. We use the IWC-weighted average of the homogeneous fraction to calculate  $F_{\text{hom}}$ . In this study, Sfc  $\Delta\text{CRE}_{\text{max}}$  is

used to estimate the instantaneous efficacy of CCT, while  $\text{Atm } \Delta\text{CRE}_{\text{max}}$  represents the potential CCT effect, that is, the extent to which changes in atmospheric cooling due to CCT could ultimately influence the surface through climatic feedback processes.

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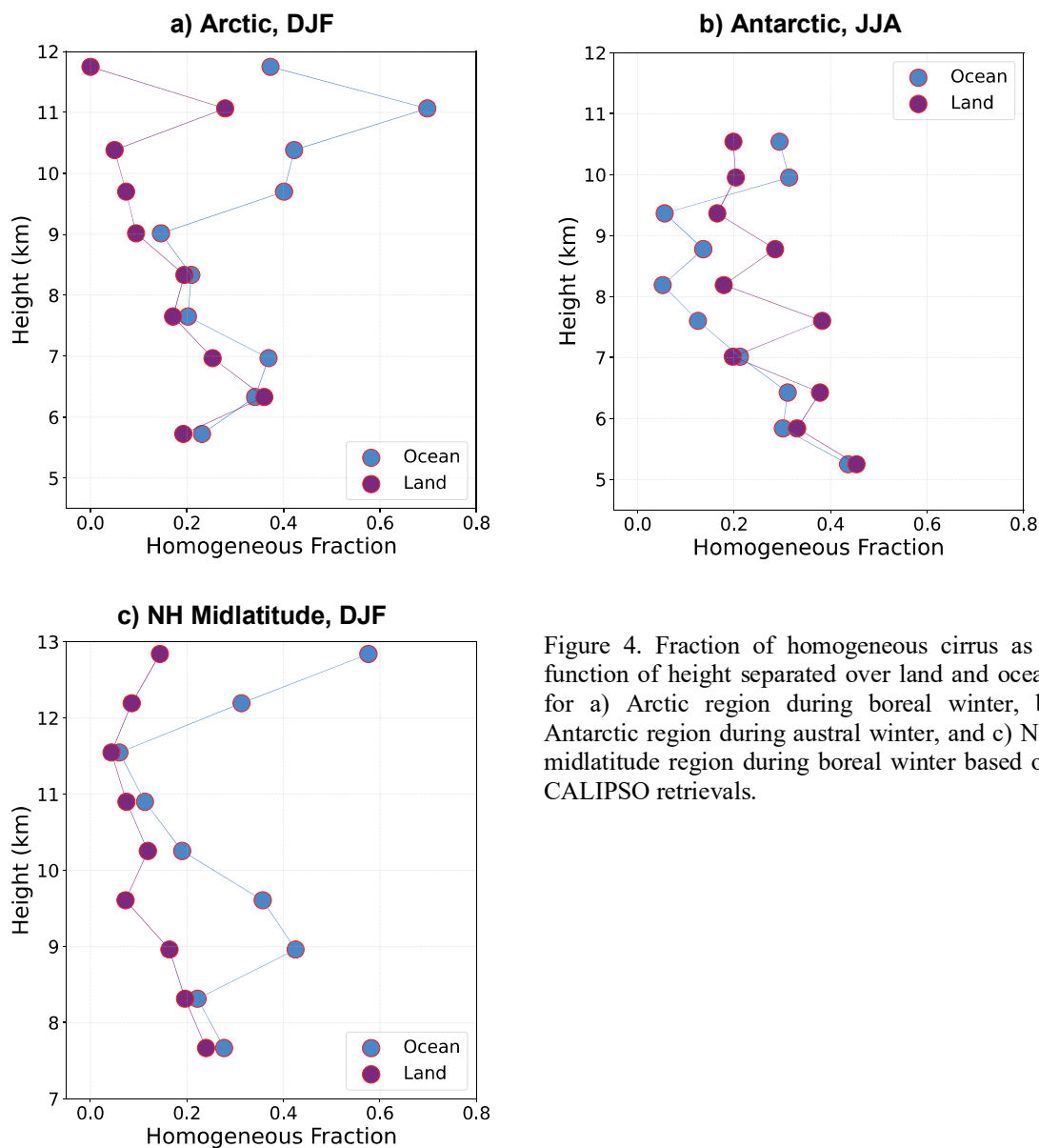


Figure 4. Fraction of homogeneous cirrus as a function of height separated over land and ocean for a) Arctic region during boreal winter, b) Antarctic region during austral winter, and c) NH midlatitude region during boreal winter based on CALIPSO retrievals.

Note that  $\Delta\text{CRE}_{\text{max}}$  emphasizes the first CCT scenario, namely the transition from homogeneous cirrus to heterogeneous cirrus. To account for new cirrus clouds formed by INPs injected into clear-sky ice-supersaturated regions, we only consider the heterogeneous cirrus CRE, because homogeneous cirrus clouds depend primarily on  $\text{RH}_i$  and would form regardless. As mentioned,

370 Lin et al. (2025) and Sporre et al. (2022) report that elevated INP concentrations within volcanic plumes entering the troposphere increased the cirrus coverage or fraction by approximately 20% (denoted here as  $\Delta CF_{\text{new cirrus}}$ ). This provides a means of estimating the CRE of new cirrus:

$$CRE_{\text{new cirrus}} = \Delta CF_{\text{new cirrus}} \times CF_{\text{cirrus}} \times \langle CRE_{\text{net},\text{het}} \rangle. \quad (6)$$

Finally, the total  $\Delta CRE$  can be calculated as:

$$375 \quad \Delta CRE_{\text{tot,lb}} = \Delta CRE_{\text{max}} + CRE_{\text{new cirrus}}, \quad (7)$$

where lb refers to lower bound. This calculation provides a lower-bound estimate for CCT-induced radiative impact by assuming full microphysical change under ideal meteorological conditions for heterogeneous cirrus formation.

### 2.3.2 CCT under minimal microphysical change

380 To complement the lower-bound condition, we also define a conceptual upper bound for CCT efficacy by assuming that the change in microphysical conditions is minimal after seeding, such that the seeded cirrus cloud IWC and  $D_e$  remain identical to those of homogeneous cirrus. This would correspond to conditions where cloud updrafts were sufficiently strong to render seeding effects within homogeneous cirrus clouds as impotent, and where INP seeding produces new cirrus  
385 clouds. An example might be cirrus formed over steep mountains by orographic gravity waves (OGWs). Since these IWC and  $D_e$  are the same RTM inputs as for homogeneous cirrus, this bounding condition means that  $\Delta CRE$  from Eq. (4) and  $\Delta CRE_{\text{max}}$  from Eq. (5) are zero. This framing provides a physically plausible upper limit for the efficacy of CCT and acknowledges that not all seeding events will produce sufficient microphysical changes to yield meaningful cooling.

390 The total  $\Delta CRE$  can be calculated as:

$$\Delta CRE_{\text{tot,ub}} = CRE_{\text{new cirrus}}, \quad (8)$$

where ub refers to upper bound.

Together, the upper- and lower-bounds define a range of possible radiative outcomes from CCT interventions, constrained by satellite observations and calculated within an RTM that assumes  
395 fixed cloud profiles and instantaneous radiative changes, without time-dependent feedbacks.

### 3 Main RTM simulations

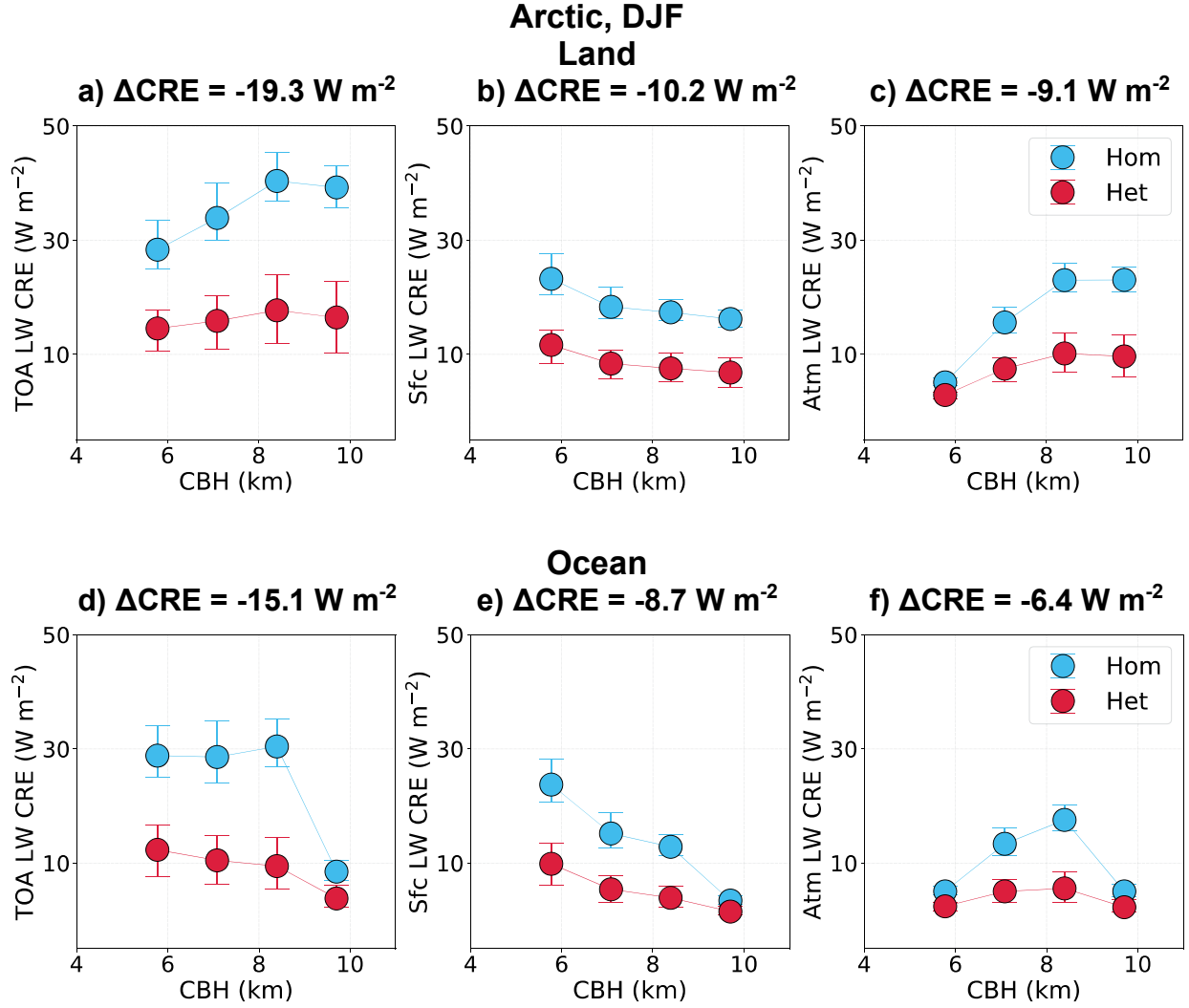
#### 3.1 Arctic region

The RTM simulations are conducted using mean thermodynamic profiles from MERRA2 for the Arctic during the boreal winter (Fig. 2) and ice cloud properties using the median, 25<sup>th</sup> and 75<sup>th</sup> percentile IWC and  $D_e$  from CALIPSO satellite retrievals, as shown in Fig. 3. A general pattern of cirrus cloud properties is seen in Fig. 3 (e.g., a decrease in both IWC and  $D_e$  with height, which is characteristic of cirrus clouds). The difference in IWC between homogeneous and heterogeneous cirrus is distinct, as homogeneous cirrus in our CALIPSO retrievals have much larger median IWC than heterogeneous cirrus at the same temperature, in agreement with previous observational studies conducted over Europe and Africa (Krämer et al., 2016, 2020) and over the Americas and Pacific Ocean (Ngo et al., 2024; Patnaude et al., 2021; Patnaude and Diao, 2020). Mitchell et al. (2024) showed that the CALIPSO retrievals generally agree well with aircraft measurements from Krämer et al. (2020). See the former for a more detailed discussion on the similarities and differences between satellite and aircraft-based observation techniques.

Despite the distinct pattern in median IWC among homogeneous and heterogeneous cirrus,  $D_e$  values are similar in both cirrus regimes, which results from the criteria applied to define heterogeneous and homogeneous cirrus clouds in Mitchell and Garnier (2024). That is, when  $D_e$  is plotted against either the SW  $\alpha_{ext}$  or IWC as shown in Fig. S1, there is generally a  $D_e$  maximum that divides the two cirrus regimes for a given  $T$ . The maximum in the number of CALIPSO cirrus cloud samples when related to  $\alpha_{ext}$  or IWC tends to coincide with this  $D_e$  maximum, resulting in similar mean  $D_e$  values for each cirrus regime. But as  $\alpha_{ext}$  or IWC increases beyond this  $D_e$  maximum,  $D_e$  decreases, which is consistent with conventional knowledge that an increase in homogeneous ice nucleation activity will act to increase  $N_i$  and decrease particle sizes due to water vapor competition effects.

Due to different cloud properties over land and ocean, different RTM simulations are conducted for land and ocean. Figure 5 shows TOA, Sfc, and Atm LW CRE calculated from RTM simulations using Eqs. (1) and (2). Note that no RTM simulation is conducted for SW range because of the absence of solar radiation in this region during the winter. As such, these results serve as net CRE

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Figure 5. Results of RTM simulations showing LW CRE as a function of CBH over the Arctic during the boreal winter for 4 cirrus clouds separated based on surface types (land and ocean) and cirrus regimes (homogeneous and heterogeneous). The CRE is calculated at the TOA, at the surface, and within the column of atmosphere. The  $\Delta\text{CRE}$  at the top of each panel represent the transition from homogeneous to heterogeneous cirrus based on Eq. (5). A total of 48 RTM simulations are shown in this figure with markers and error bars referring to simulations based on CALIPSO profiles in Fig. 4.

(Eq. 3). LW CRE in Fig. 5 varies with CBH, highlighting the effects of cirrus cloud altitude as well as microphysical properties. The LW CRE at the surface generally decreases with CBH because colder clouds at higher altitudes emit less LW radiation compared to warmer clouds at lower altitudes, based on the Stefan–Boltzmann law. Note that cirrus cloud altitude is closely related to cirrus cloud temperature, since both are connected via the vertical temperature profile. In addition, cirrus clouds at higher altitudes often have lower IWC (Fig. 3), and this makes them optically thinner. In contrast, smaller  $D_e$  in cirrus at higher altitudes could lead to stronger LW

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CRE (Fu and Liou, 1993). At the TOA, LW CRE depends on the difference between the cloud's LW emission and the emission from the Earth's surface (Corti and Peter, 2009), and such difference is larger for cirrus at higher altitudes.

Also seen in Fig. 5 is significantly larger LW CRE at the TOA, at the Sfc, and within the Atm for homogeneous cirrus than that for heterogeneous cirrus of the same altitude. This is mainly due to higher IWC values for homogeneous cirrus (Fig. 3), which leads to optically thicker cirrus (Krämer et al., 2016, 2020). When both cirrus regimes have comparable IWC, as seen for the highest altitude over the ocean, their LW CRE is comparable. This highlights the critical role of IWC in determining the radiative impact of cirrus clouds.

For cirrus overcast conditions over land, the lower bound of CCT efficacy (defined as the cooling resulting from the transition of natural cirrus clouds to pure heterogeneous cirrus, and quantified by  $\Delta\text{CRE}$  in Eq. (4)), has a TOA cooling effect of  $-19.3 \text{ W m}^{-2}$  (the mean value of the four clouds considered), with a corresponding Sfc cooling of  $-10.2 \text{ W m}^{-2}$  and atmospheric column cooling of  $-9.1 \text{ W m}^{-2}$  (Figs. 5a-c). Considering that the typical cirrus cloud cover over the Arctic is 35% and that the IWC-weighted average of the homogeneous fraction is 0.21 (Fig. 4a), Eq. 5 gives the maximum cooling effect  $\Delta\text{CRE}_{\text{max}}$  at the TOA, Sfc, and Atm as  $\sim -1.4$ ,  $-0.7$ , and  $-0.7 \text{ W m}^{-2}$ , respectively (Table 2). After accounting for the impact of new cirrus formation (Eq. 6), the lower bound of total cloud effect  $\Delta\text{CRE}_{\text{tot,lb}}$  (Eq. 7) at the TOA, Sfc, and Atm is  $\sim -0.3$ ,  $-0.2$ , and  $-0.1 \text{ W m}^{-2}$ , respectively (negative values indicate a cooling effect). The upper bound ( $\Delta\text{CRE}_{\text{tot,ub}}$ ; Eq. 8), however, results in a warming of 1.1, 0.5, and  $0.6 \text{ W m}^{-2}$  at TOA, Sfc, and Atm, respectively. Of particular importance for CCT is the cooling at the surface but it should be noted that the RTM provides instantaneous values only. For the atmospheric column, the CRE is similar to the surface CRE. This might have implications for long-term feedback processes and possibly impact of AA, as the lower-bound atmospheric column cooling could lead to lower geopotential thickness over the Arctic, which in turn might affect meridional  $T$  gradients, thermal winds, and the extratropical jet stream (Cohen et al., 2020). The upper bound implies the opposite, e.g., warming in both Sfc and Atm CRE, which might lead to enhanced AA. A careful GCM study is required to evaluate the sign and magnitude of CCT and the corresponding feedbacks.

The overall pattern of CRE change over the ocean is consistent with that over land, but the cooling effect over the ocean is slightly weaker, with a TOA  $\Delta\text{CRE}$  of -15.1, a Sfc  $\Delta\text{CRE}$  of -8.7 W m<sup>-2</sup> and an Atm  $\Delta\text{CRE}$  of -6.4 W m<sup>-2</sup> (Figs. 5d-f). With a typical cirrus cloud cover value of 35% and IWC-weighted mean homogeneous fraction of 0.29 (Fig. 4a) over the ocean, TOA, Sfc, and Atm  $\Delta\text{CRE}_{\text{max}}$  are approximately -1.5, -0.9, and -0.6 W m<sup>-2</sup>, respectively (Table 2). These values are higher than  $\Delta\text{CRE}_{\text{max}}$  over land because of the higher homogeneous fraction over the ocean. In addition, the lower and upper bounds of  $\Delta\text{CRE}_{\text{tot}}$  at TOA, Sfc, and Atm are approximately [-0.9, 0.6], [-0.5, 0.4], and [-0.4, 0.2] W m<sup>-2</sup>, respectively.

Note that in Mitchell and Garnier (2024), regions consisting of sea ice are considered as land. As shown in Fig. S2a, the higher sea ice fraction in winter along with the pure land fraction constitutes a much larger area than water surfaces. As such,  $\Delta\text{CRE}_{\text{tot}}$  over the ocean makes a smaller impact. Nevertheless, we conduct analysis for both land and ocean for a more comprehensive analysis. As the climate continues to warm, the ocean fraction of the winter Arctic will likely increase.

### 3.2 Antarctic region

The RTM simulations for the Antarctic are conducted similarly to those for the Arctic, using mean thermodynamic profiles from MERRA2 (not shown) and median, 25<sup>th</sup> and 75<sup>th</sup> percentile IWC and  $D_e$  profiles from CALIPSO satellite retrievals (Fig. 6) during the austral winter for this region. While the general patterns of IWC and  $D_e$  profiles for homogeneous and heterogeneous cirrus are similar to those in the Arctic, the specific values and details differ between the two regions. Simulations are performed for both land and ocean, and the LW CRE (equivalent to net CRE due to the absence of SW radiation during austral winter) is calculated at the TOA, Sfc, and Atm, as shown in Fig. 7.

The TOA CRE over Antarctic land for cirrus overcast conditions is weaker than that over the Arctic for cirrus clouds at the same altitude, particularly for homogeneous cirrus at the two lowest altitudes. This is likely due to lower IWC in the lowest altitudes over the Antarctic compared to the Arctic (Figs. 3 and 6). As a result, the transition from homogeneous to heterogeneous cirrus,

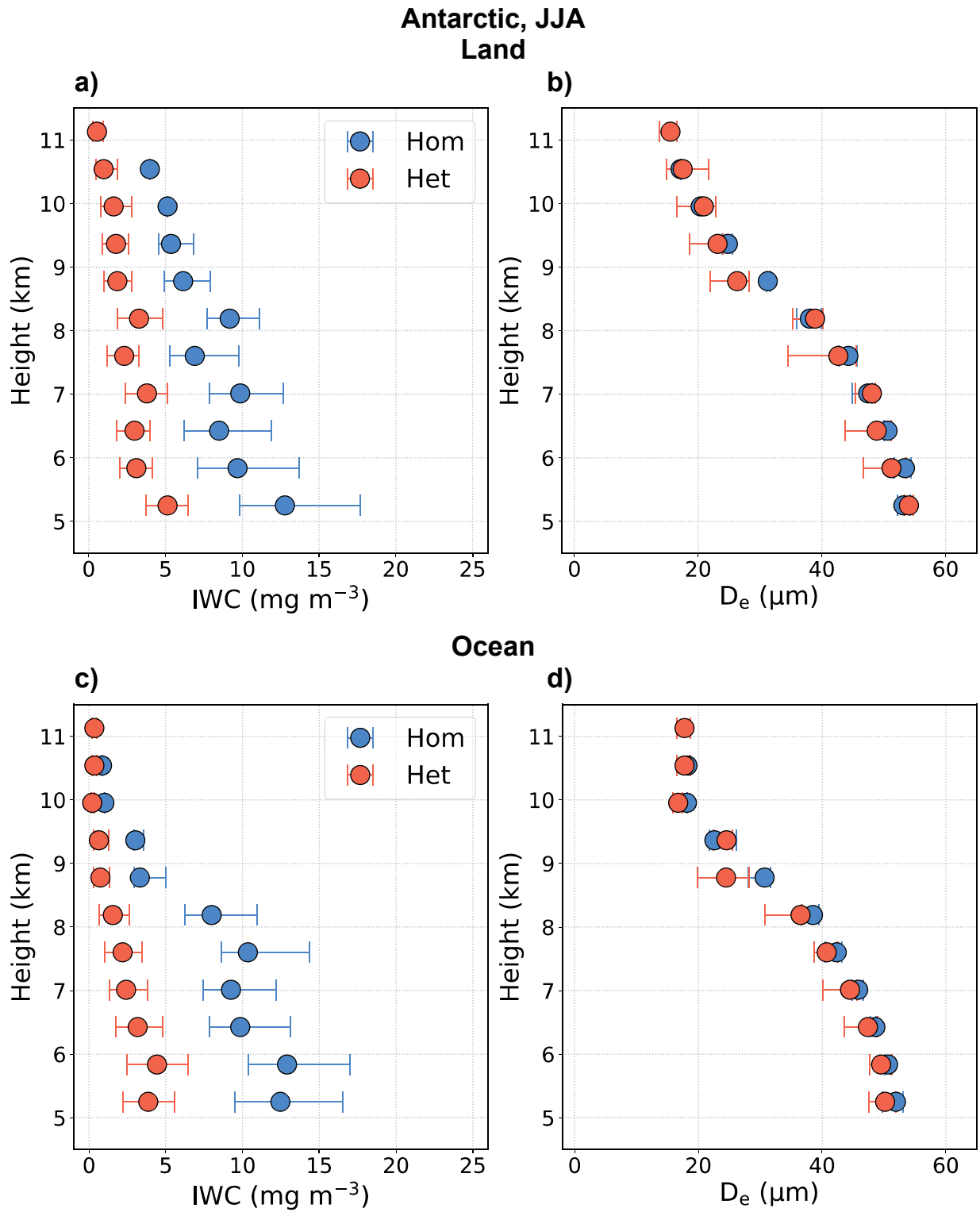


Figure 6. As in Fig. 3, but the results are for Antarctic (90-60°S) during austral winter (JJA).

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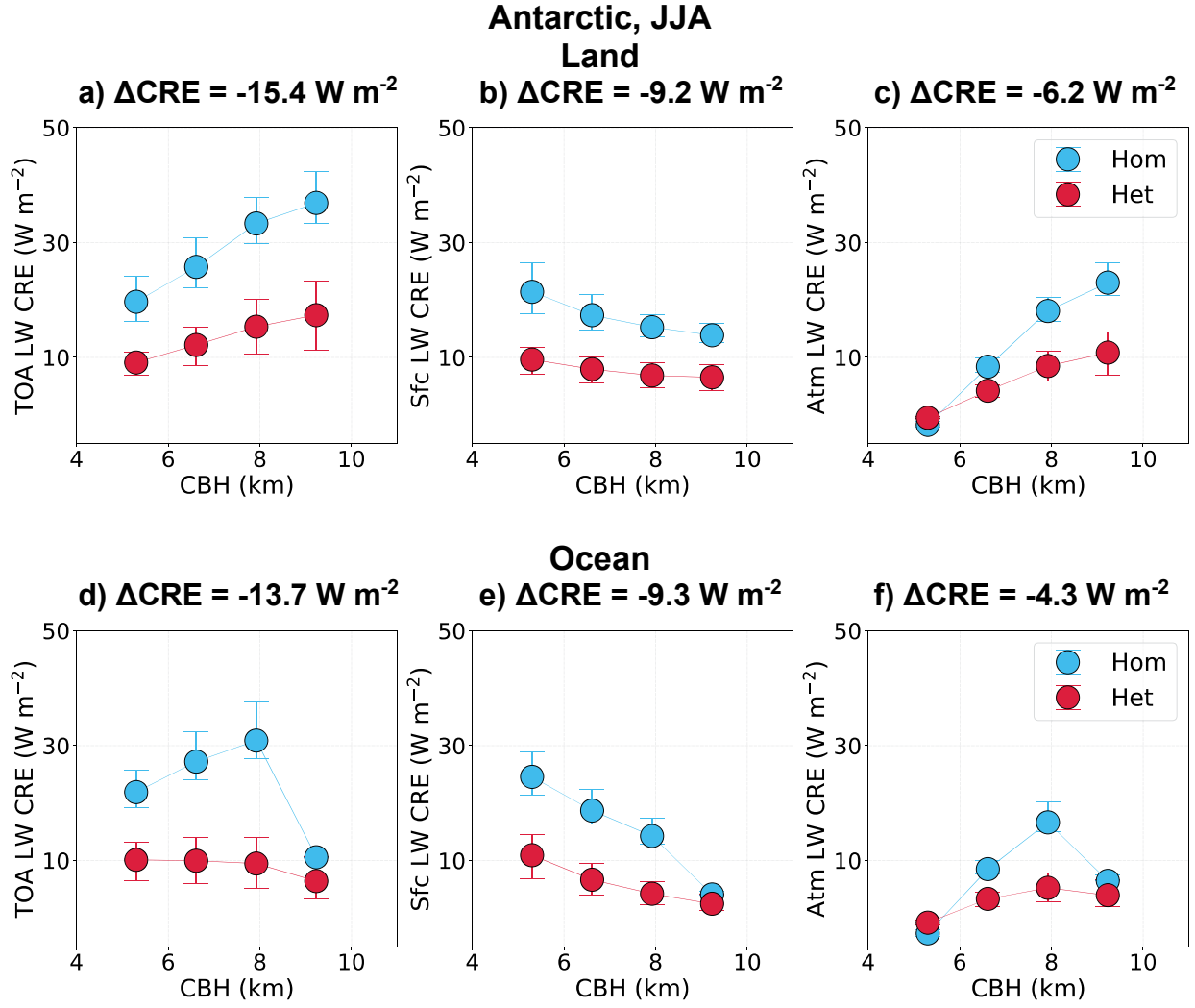


Figure 7. As in Fig. 5, but the results are RTM simulations for Antarctic during austral winter.

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Table 2. Quantifying the transition from homogeneous to heterogeneous cirrus (for overcast skies) using the change in their cloud radiative effect ( $\Delta\text{CRE}$ ) and its maximum value that assumes 35% cloud coverage ( $\Delta\text{CRE}_{\text{max}}$ ) at various levels based on Eq. (5) for different regions, seasons, and surface types. In addition, total values for lower bound ( $\Delta\text{CRE}_{\text{tot,lb}}$ ) and upper bound ( $\Delta\text{CRE}_{\text{tot,ub}}$ ) are provided based on Eqs. (7) and (8) to account for the new cirrus formation.

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Region	Season	Surface type	$F_{\text{hom}}$	$\Delta\text{CRE} \text{ (W m}^{-2}\text{)}$			$\Delta\text{CRE}_{\text{max}} \text{ (W m}^{-2}\text{)}$			$\Delta\text{CRE}_{\text{tot,lb}} \text{ (W m}^{-2}\text{)}$			$\Delta\text{CRE}_{\text{tot,ub}} \text{ (W m}^{-2}\text{)}$		
				TOA	Sfc	Atm	TOA	Sfc	Atm	TOA	Sfc	Atm	TOA	Sfc	Atm
Arctic	DJF	Land	0.21	-19.3	-10.2	-9.1	-1.4	-0.7	-0.7	-0.3	-0.2	-0.1	1.1	0.5	0.6
		Ocean	0.29	-15.1	-8.7	-6.4	-1.5	-0.9	-0.6	-0.9	-0.5	-0.4	0.6	0.4	0.2
Antarctic	JJA	Land	0.3	-15.4	-9.2	-6.2	-1.6	-1.0	-0.6	-0.7	-0.4	-0.3	0.9	0.6	0.3
		Ocean	0.24	-13.7	-9.3	-4.3	-1.2	-0.8	-0.4	-0.5	-0.3	-0.2	0.7	0.5	0.2
NH midlat	DJF	Land	0.15	-22.9	+0.2	-23.1	-1.2	0.0	-1.2	+0.3	+0.0	+0.3	1.5	0	1.5

quantified by  $\Delta\text{CRE}$ , leads to a TOA cooling of  $-15.4 \text{ W m}^{-2}$ , which is roughly 20% weaker than the  $\Delta\text{CRE}$  over Arctic land. The Sfc and Atm  $\Delta\text{CRE}$  values are  $-9.2 \text{ W m}^{-2}$  ( $\sim 10\%$  weaker than that over the Arctic land), and  $-6.2 \text{ W m}^{-2}$  ( $\sim 40\%$  weaker than that over the Arctic land), respectively. Despite the lower IWC for homogeneous cirrus over the Antarctic, the homogeneous fraction is significantly higher (IWC-weighted average is 0.30), resulting in stronger maximum cooling over the Antarctic land than over the Arctic land; the maximum cooling effects ( $\Delta\text{CRE}_{\text{max}}$ ) at the TOA, Sfc, and Atm are approximately  $-1.6$ ,  $-1.0$ , and  $-0.6 \text{ W m}^{-2}$ , respectively, and the lower and upper bounds of  $\Delta\text{CRE}_{\text{tot}}$  at the TOA, Sfc, and Atm are approximately  $[-0.7, 0.9]$ ,  $[-0.4, 0.6]$ , and  $[-0.3, 0.3] \text{ W m}^{-2}$ , respectively (Table 2).

Over the ocean, the TOA cooling effect ( $\Delta\text{CRE}$ ) is weaker compared to all previous results in this study. The TOA, Sfc, and Atm  $\Delta\text{CRE}$  values are estimated to be  $-13.7$ ,  $-9.3$ , and  $-4.3 \text{ W m}^{-2}$ , respectively. With an IWC-weighted average homogeneous fraction of 0.24,  $\Delta\text{CRE}_{\text{max}}$  at the TOA, Sfc, and Atm are approximately  $-1.2$ ,  $-0.8$ , and  $-0.4 \text{ W m}^{-2}$ , respectively, and the lower and upper bounds of  $\Delta\text{CRE}_{\text{tot}}$  at the TOA, Sfc, and Atm are approximately  $[-0.5, 0.7]$ ,  $[-0.3, 0.5]$ , and  $[-0.2, 0.2] \text{ W m}^{-2}$ , respectively (Table 2). These values are slightly weaker than those for Antarctic land. However, for the Antarctic, the CCT cooling effect over the ocean is much smaller than that over land, given that the surface water fraction is much smaller than the fraction of sea ice and the Antarctic land mass during austral winter (Fig. S2b).

To the best of our knowledge, no previous study has used an RTM to estimate the efficacy of CCT. Although the instantaneous surface cooling in our study for both polar regions and over land and ocean (Sfc  $\Delta\text{CRE}_{\text{max}}$ :  $-0.7$  to  $-1.0 \text{ W m}^{-2}$  and Sfc  $\Delta\text{CRE}_{\text{tot,lb}}$ :  $-0.2$  to  $-0.5 \text{ W m}^{-2}$ ) and the TOA cooling (TOA  $\Delta\text{CRE}_{\text{max}}$ :  $-1.2$  to  $-1.6 \text{ W m}^{-2}$  and TOA  $\Delta\text{CRE}_{\text{tot,lb}}$ :  $-0.3$  to  $-0.9 \text{ W m}^{-2}$ ) are much weaker than the potential cooling of  $-2.8 \text{ W m}^{-2}$  suggested by Mitchell and Finnegan (2009), they fall within the range of maximum CCT cooling from previous GCM studies, from  $-0.25 \text{ W m}^{-2}$  (Gasparini and Lohmann, 2016) to  $-2 \text{ W m}^{-2}$  (Storelvmo et al., 2013; Storelvmo and Herger, 2014). We acknowledge that this is not a direct comparison, as GCMs calculate global CREs while accounting for feedback processes. However, we note that CCT in the polar regions during winter could be as effective as CCT applied globally throughout the year, largely because LW trapping by cirrus clouds outside the polar regions is counteracted by SW scattering (Storelvmo et al., 2014).

### 3.3 North hemispheric mid-latitude region

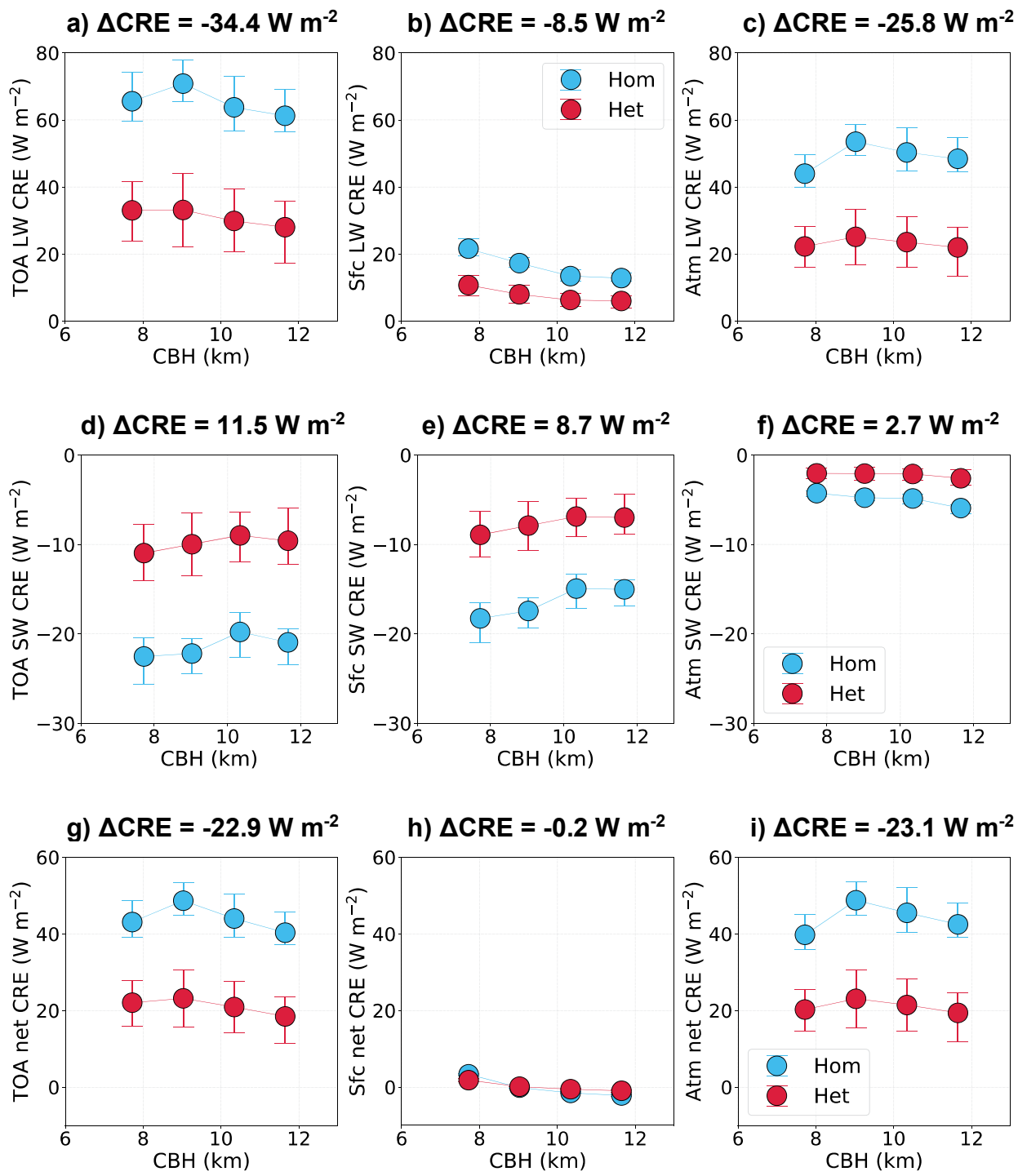
Mid-latitude regions ( $30^{\circ}\text{N}$  to  $60^{\circ}\text{N}$  and  $-60^{\circ}\text{S}$  to  $-30^{\circ}\text{S}$  latitude bands) comprise approximately 37% of the Earth's surface, which is about three times the area of the high latitudes. This makes it important to evaluate the potential efficacy of CCT in these regions. During winter, the SW impact of cirrus clouds is minimized due to shorter days and higher solar zenith angles (SZA). The SZA, which is the angle between the Sun's rays and a line perpendicular to the Earth's surface at a specific location (ranging from  $0^{\circ}$  at the equator at midday during an equinox to  $90^{\circ}$  at sunrise and sunset) (Aktaş and Kirçiçek, 2021), has a daytime average of  $73^{\circ}$  at  $45^{\circ}\text{N}$  latitude during the winter solstice (Hartmann, 2016). In addition to LW RTM simulations, we conduct SW simulations for a daytime average winter solstice mid-latitude scenario:  $45^{\circ}\text{N}$  latitude, a surface albedo of 0.3, and a SZA of  $73^{\circ}$ . The RTM is forced with mean thermodynamic profiles from MERRA2 (not shown) and median, 25th, and 75th percentile IWC and  $D_e$  profiles from CALIPSO satellite retrievals (Fig. S3) during the boreal winter for NH mid-latitude land.

The results of the RTM simulations for various CREs are shown in Fig. 8. The LW CRE at the TOA over mid-latitudes is significantly larger than that over polar regions for cirrus clouds of the same regime (homogeneous or heterogeneous) and at the same altitude. This is likely due to higher IWC within cirrus clouds (Fig. S3) and a warmer temperature profile for midlatitudes compared to polar regions. Cirrus clouds with higher IWC retain more LW radiation, resulting in stronger LW CRE (Fu and Liou, 1993). Furthermore, the warmer atmospheric column and in particular warmer surface in mid-latitudes emit more LW radiation toward the upper troposphere, which is absorbed and re-emitted at colder temperatures by cirrus clouds. This causes a stronger difference between LW radiation emitted by cirrus cloud and Earth's surface and enhances the TOA LW CRE (Corti and Peter, 2009).

The SW CRE (Figs. 8d–f) is calculated to provide daily-mean values. To account for the diurnal cycle of SW radiation, the SW CRE from Eqs. (1) and (2) is multiplied by a factor of 0.37, representing the ratio of daytime hours (8.8 hours) to 24 hours at  $45^{\circ}\text{N}$  latitude during the winter solstice. This post-simulation factor, combined with the daytime-average SZA used in the RTM simulations, averages the SW CRE at  $45^{\circ}\text{N}$  over a full 24-hour period, consistent with the LW CRE calculations. All SW CRE values are negative, indicating the cooling effect of cirrus clouds

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NH Midlatitude Winter over Land



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Figure 8. As in Fig. 5, but the results are RTM simulations for LW, SW, and net CRE over NH midlatitude land with a total of 50 RTM simulations.

at different altitudes and with various microphysical properties due to the absorption and scattering of solar radiation. Homogeneous cirrus clouds exhibit significantly stronger SW cooling effects than heterogeneous cirrus clouds at the TOA and Sfc, as they contain higher IWC, which corresponds to greater scattering and absorption by ice particles (Fu and Liou, 1993). The change in SW CRE with cloud altitude depends on changes in  $\alpha_{ext}$ , where  $\alpha_{ext} = 3 \text{ IWC}/(\rho_i D_e)$ , and  $\rho_i$  is bulk density of ice. As cloud altitude increases, both IWC and  $D_e$  decrease, resulting in a relatively slow decrease in  $\alpha_{ext}$  with increasing altitude (Fu and Liou, 1993; Stephens et al., 1990).

For cirrus overcast conditions at the TOA, the strong difference in LW CRE between the two cirrus regimes results in significant LW cooling ( $\Delta\text{CRE} = -34.4 \text{ W m}^{-2}$ ), which is partially offset by SW warming ( $\Delta\text{CRE} = 11.5 \text{ W m}^{-2}$ ), yielding a net TOA cooling of  $-22.9 \text{ W m}^{-2}$  (Fig. 8g). The transition from homogeneous to heterogeneous cirrus results in a surface LW cooling ( $\Delta\text{CRE}$ ) of  $-8.5 \text{ W m}^{-2}$ , which is largely offset by SW warming ( $\Delta\text{CRE} = 8.7 \text{ W m}^{-2}$ ), leading to a relatively small net surface  $\Delta\text{CRE}$  of  $-0.2 \text{ W m}^{-2}$  (Fig. 8h). Within the atmospheric column, a significant net cooling of  $-23.1 \text{ W m}^{-2}$  occurs (Fig. 8i). Considering an IWC-weighted average homogeneous fraction of 0.15 (Fig. 4c) and a cirrus cloud cover of 35%, the maximum net cooling effects ( $\Delta\text{CRE}_{\text{max}}$ ) at the TOA, Sfc, and Atm are approximately  $-1.2$ ,  $0.0$ , and  $-1.2 \text{ W m}^{-2}$ , respectively (Table 2). These results demonstrate that in the absence of new cirrus formation, while the instantaneous cooling efficacy of CCT (Sfc net  $\Delta\text{CRE}_{\text{max}}$ ) in mid-latitudes during winter is negligible, CCT could still be effective if its impact on the atmospheric column (Atm net  $\Delta\text{CRE}_{\text{max}}$ ) can reach the surface through feedback processes. However, after accounting for new cirrus formation and the bounds of change in microphysical conditions (from full change to no change), the lower and upper bounds of  $\Delta\text{CRE}_{\text{tot}}$  at the TOA, Sfc, and Atm are  $\sim [+0.3, 1.5]$ ,  $[0.0, 0.0]$ , and  $[+0.3, 1.5] \text{ W m}^{-2}$ , respectively (Table 2), indicating a warming effect in the TOA and Atm, and suggesting that CCT could even result in net warming in this season and latitude band.

## 4 Sensitivity tests

### 4.1 Sensitivity to thermodynamic profiles

The impact of temperature and humidity on cirrus LW CRE is evaluated using minimum and maximum air  $T$  and  $q_v$  profiles (referred to as  $T_{\text{min}}$  and  $T_{\text{max}}$  for brevity) from MERRA2 data for Arctic land during the winter (Fig. 9). TOA LW CRE significantly increases with an increase in  $T$

and  $q_v$ . In particular, Earth's surface plays an important role because it typically acts as a blackbody (its  $\varepsilon$  is very close to unity), and even a rather small surface warming can significantly enhance LW radiation emitted from the surface, as described by Stefan–Boltzmann law. With unchanged cirrus temperature and LW emission, the enhanced upward LW radiation from the Earth's surface creates a stronger LW contrast, resulting in a stronger TOA LW CRE (Corti and Peter, 2009).

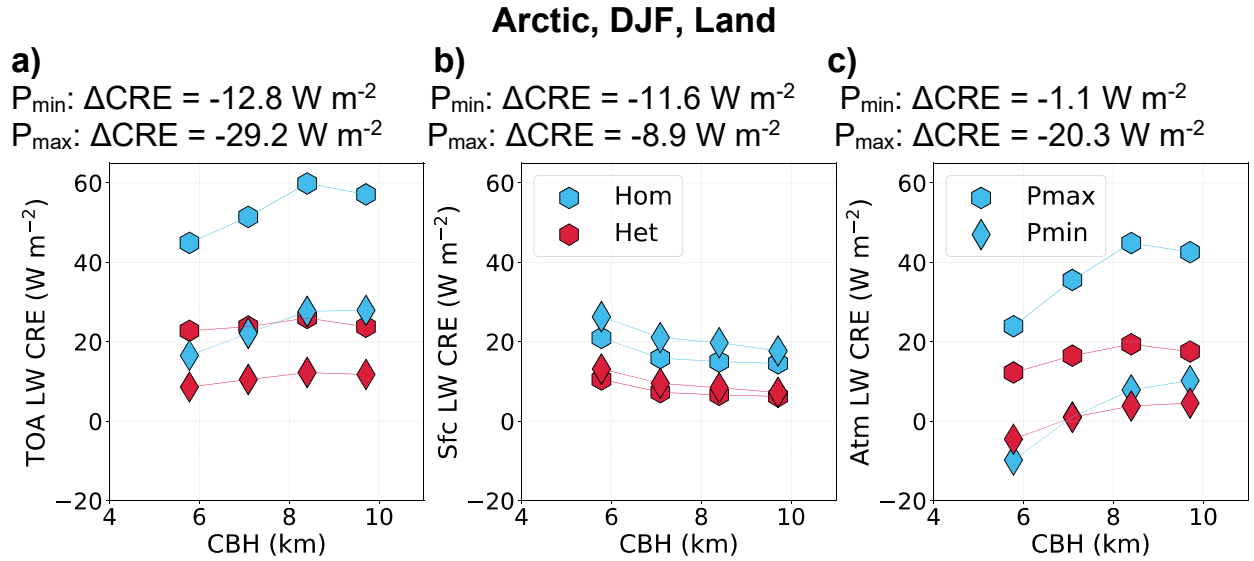


Figure 9. Sensitivity of RTM-simulated cirrus CRE to different thermodynamic profiles ( $P$ ) from MERRA2 minimum and maximum temperature and water mixing ratio (abbreviated as  $P_{\min}$  and  $P_{\max}$ ), as shown in Fig. 2.

At the surface, however, LW CRE is weakly sensitive to thermodynamic profiles (Fig. 9b). Profiles with lower  $T$  and  $q_v$  lead to slightly higher cirrus LW CRE at the surface, particularly for homogeneous cirrus. The surface LW CRE depends primarily on the downward LW radiation from cirrus clouds, rather than surface temperature (Eq. 1). Therefore, the lower surface LW CRE in maximum profiles compared to minimum profiles is due to higher water vapor in the atmosphere, which absorbs part of the downward LW radiation from cirrus clouds before it reaches the surface. This is consistent with the findings of Dupont and Haeffelin (2008).

Figure 9a shows that the transition from homogeneous to heterogeneous cirrus ( $\Delta\text{CRE}$ ) intensifies significantly with warmer and more humid thermodynamic profiles, particularly with higher surface temperatures. The  $\Delta\text{CRE}$  for minimum and maximum profiles is  $-12.8 \text{ W m}^{-2}$  and  $-29.2 \text{ W m}^{-2}$ , respectively. At the surface (Fig. 9b), the  $\Delta\text{CRE}$  for minimum and maximum profiles is  $-11.6 \text{ W m}^{-2}$  and  $-8.9 \text{ W m}^{-2}$ , respectively, indicating minimal sensitivity to thermodynamic profiles.

660 This consistency suggests that the instantaneous CCT efficacy is robust across different thermodynamic conditions. However, the atmospheric  $\Delta\text{CRE}$  (Fig. 9c) shows greater variability, ranging from  $-1.1 \text{ W m}^{-2}$  for the minimum thermodynamic profile to  $-20.3 \text{ W m}^{-2}$  for the maximum profile, highlighting the sensitivity of potential CCT efficacy to thermodynamic profiles.

## 4.2 Sensitivity to Arctic low clouds

665 Low clouds are frequent over the Arctic region and they have a significant impact on the radiation balance (Philipp et al., 2020). These clouds are controlled by many factors including atmospheric circulation and sea ice extent and in return, they impact the sea ice via an ice-albedo feedback (Huang et al., 2021). During the winter, low clouds retain outgoing longwave radiation and warm the surface, but during the summer, this effect is canceled by cooling from reflecting solar radiation  
670 (Maillard et al., 2021). Arctic low cloud cover varies by season and this variability is more distinct for higher latitudes of the Arctic (north of latitude 70) where low cloud cover changes from over 50% in summer to lower than 20% in winter (Eastman and Warren, 2010). Arctic low clouds tend to have higher cloud water path (CWP) over the open ocean and lower CWP over ice-covered areas (Yu et al., 2019) due to higher moisture availability over the ocean than ice (Monroe et al.,  
675 2021). The spatial distribution of arctic low clouds shows that over land their cover is typically around 35% in summer and around 15% in winter. Over the ocean, their cover is around 55% in summer, but drops below 30% on the Pacific side of the Arctic Ocean, meanwhile remains as high as 50% on the Atlantic side of the Arctic Ocean in winter (Huang et al., 2021).

Our RTM simulations explore the impact of low liquid clouds on cirrus CRE by introducing a low  
680 liquid cloud layer, as described in Sect. 2. Three low liquid clouds are tested by varying LWC (e.g., 0.01, 0.03, and  $0.05 \text{ g m}^{-3}$ ). To calculate cirrus CRE using Eq. (1), we consider the difference between an RTM run with both cirrus and low liquid cloud versus an RTM run with only low liquid cloud.

The results (Fig. 10) show that TOA LW CRE for cirrus clouds is not sensitive to the low liquid  
685 clouds. Over the Arctic, such clouds are close to the surface, and their temperature is very similar to that of the Earth's surface (due to inversion, mean profile of  $T$  in Fig. 2a varies slowly below 2 km). As a result, the LW radiation emitted by low liquid clouds is close to that emitted by Earth's

surface. Moreover, we only vary the LWC of low clouds, not their elevation, so their temperature remains constant. Consequently, the difference between cirrus LW radiation and the upward LW radiation from the underlying clouds and Earth's surface does not change significantly across the three sensitivity tests in this section when considering CRE at TOA.

### Arctic, DJF, Land

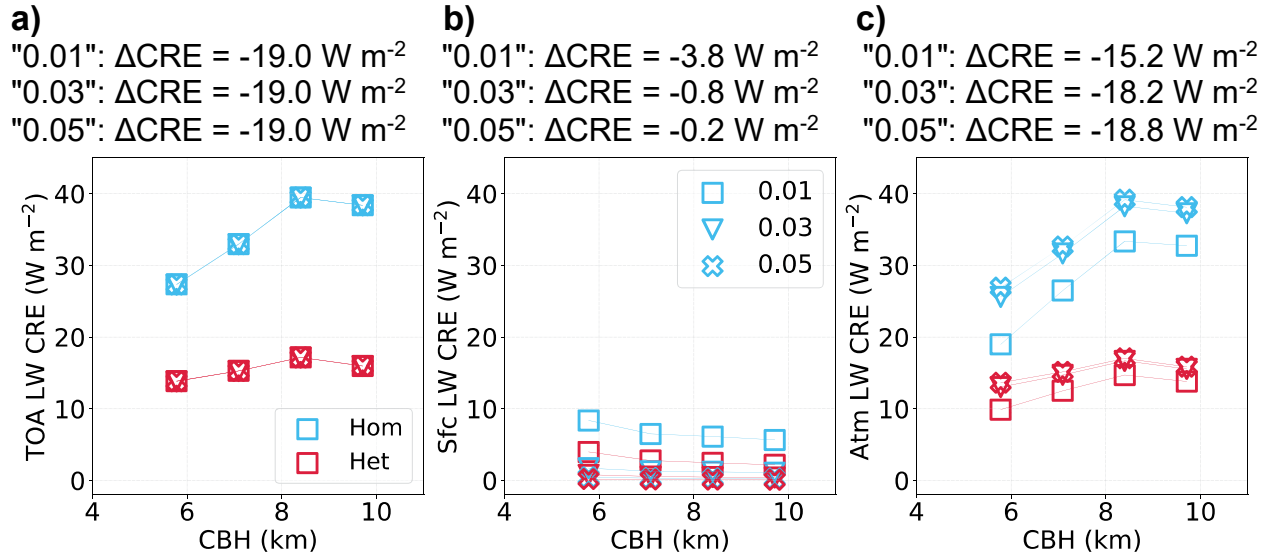


Figure 10. Sensitivity of RTM-simulated cirrus CRE to three different low liquid clouds with varying liquid water content (LWC) values of 0.01, 0.03, and 0.05  $\text{g m}^{-3}$ .

At the surface, however, cirrus LW CRE decreases rapidly as low cloud LWC increases. Note that the largest LWC selected here (0.05  $\text{g m}^{-3}$ ) is at the lower end of typical LWC values observed in the Arctic (Achtert et al., 2020). Our results demonstrate that low liquid clouds  $\sim 600 \text{ m}$  thick with a LWC greater than 0.05  $\text{g m}^{-3}$  act more like a “black body”, absorbing/emitting almost all the downward LW radiation emitted by cirrus clouds.

The presence of low clouds has little effect on the transition from homogeneous to heterogeneous cirrus at the TOA, with  $\Delta\text{CRE}$  remaining at  $-19.0 \text{ W m}^{-2}$ . However, it considerably reduces  $\Delta\text{CRE}$  at the surface, from  $-3.8 \text{ W m}^{-2}$  (for  $\text{LWC} = 0.01 \text{ g m}^{-3}$ ) to  $-0.2 \text{ W m}^{-2}$  (for  $\text{LWC} = 0.05 \text{ g m}^{-3}$ ). As a result, the atmospheric  $\Delta\text{CRE}$  remains between  $-15.2 \text{ W m}^{-2}$  and  $-18.8 \text{ W m}^{-2}$ . These results imply that while the instantaneous efficacy of CCT is negligible in the presence of low liquid clouds, its potential efficacy could still influence the surface through feedback processes over longer timescales.

### 4.3 Sensitivity to Arctic aerosols

In the past, the Arctic atmosphere was considered pristine, but over the past decades, it has been revealed that Arctic aerosols play an important role through aerosol-radiation interactions (Thorsen and Fu, 2015) and aerosol-cloud interactions (Creamean et al., 2021; Zamora et al., 2016). Both observations (Dagsson-Waldhauserova et al., 2019) and numerical simulations (Breider et al., 2014) showed that Arctic aerosol concentrations vary with season with the main peak in late winter and spring, and another peak in fall. The major peak is known as the Arctic haze, a phenomenon mainly caused by the transport of industrial anthropogenic aerosols from Europe and Asia that remain in the Arctic atmosphere due to a stable atmosphere and a lack of precipitation (Schmale et al., 2022). With the reduction of anthropogenic aerosols in summer, natural aerosols, including sea spray and organic compounds, dominate (Moschos et al., 2022). Another important aerosol type in the Arctic is dust with its maximum in late winter and early spring due to the long-range transport from Asia and Africa and its minimum in summer and fall predominantly because of local sources (Groot Zwaafink et al., 2016; Xie et al., 2022).

Our RTM simulations evaluate the sensitivity of cirrus CRE to different aerosol scenarios, as explained in Sect. 2. The results (Fig. S4) show that aerosol type and concentration have a relatively small impact on cirrus LW CRE. This finding is consistent with previous studies, which have demonstrated that while aerosols absorb SW radiation, they are weak absorbers of LW radiation (Bergstrom et al., 2007; Samset et al., 2018). As a result, the cooling effect of transitioning from homogeneous to heterogeneous cirrus is not sensitive to the choice of aerosol scenarios, with TOA  $\Delta$ CRE ranging from -19.3 to -19.8 W m<sup>-2</sup>, Sfc  $\Delta$ CRE from -10.2 to -10.4 W m<sup>-2</sup>, and Atm  $\Delta$ CRE from -9.1 to -9.4 W m<sup>-2</sup>. It is important to note that the modeling design here only accounts for the aerosol direct effect, as the RTM cannot simulate aerosol indirect effects. However, it would be possible to study such effect if cloud profiles are carefully explored and grouped based on aerosol loading.

## 5 Suggestions for improving cirrus cloud modeling

In previous sections, we implemented retrieved cloud microphysical products from satellite in an RTM to estimate the instantaneous cirrus CRE. RTMs have fewer degrees of freedom than GCMs, and this makes them more convenient for interpreting changes in cirrus radiative impacts.

However, GCMs are the ultimate tool for determining the global cirrus CRE since they account for climate feedback processes which can potentially increase or decrease the CRE predicted by an RTM. For example, the direct CCT polar cooling predicted by an RTM may promote coverage by snow and sea ice (Storelvmo et al., 2014), enhancing planetary albedo and thus cooling. Despite their advantages, GCMs face several challenges in accurately representing cirrus clouds. Below, we briefly discuss these issues and propose improvements based on recent research.

GCMs employ ice cloud parameterizations that are often based on limited observations and therefore, uncertainties could arise when generalizing those formulations (Eidhammer et al., 2017; Gettelman and Morrison, 2015). In particular, many field campaigns do not sample homogeneous cirrus clouds sufficiently (Mitchell et al., 2024). Also, in prognostic modeling frameworks, the competition between heterogeneous and homogeneous ice nucleation remains a complex process (Barahona and Nenes, 2009; Kärcher et al., 2022; Spichtinger and Cziczo, 2010; Spichtinger and Gierens, 2009). Current GCMs might underestimate the contribution of homogeneous nucleation, particularly outside the tropics during the winter season, when INP concentrations appear to be lower (Carlsen and David, 2022; Mitchell and Garnier, 2024). For example, the GCM simulations of Gasparini and Lohmann (2016) predict homogeneous nucleation dominating only below  $\sim 250$  hPa (above  $\sim 11$  km) when pre-existing ice was not considered, and the main CCT simulations in Tully et al. (2022, 2023) did not consider OGW induced cirrus clouds. This differs from the CALIPSO-derived results in Mitchell and Garnier (2024, Fig. 19) that show homogeneous cirrus clouds contributing significantly at all cirrus levels, with evidence that a substantial percentage of these homogeneous cirrus clouds are OGW cirrus clouds. This shortcoming in GCMs can lead to an underestimation of the radiative effects of cirrus clouds and the potential cooling efficacy of CCT. To address this, GCMs could use satellite retrievals of  $N_i$ ,  $D_e$ , and IWC when developing/constraining parameterizations that represent the two cirrus cloud regimes.

On the other hand, the GCM-CCT modeling study by Gasparini and Lohmann (2016) found that INP seeding affects mostly in situ cirrus clouds, with only minor impacts on cirrus clouds resulting from strong dynamical forcing, such as OGW cirrus clouds. While this has not been confirmed by observations (e.g., from a field experiment), it appears plausible that INP seeding may not sufficiently reduce the  $RH_i$  in the stronger OGW cirrus updrafts to prevent homogeneous freezing.

This factor may increase the value of the lower-bound  $\Delta\text{CRE}_{\text{tot}}$  estimates from this study (i.e., making them less negative).

A critical factor in modeling cirrus clouds is the treatment of pre-existing ice, which refers to ice particles already present before the formation of new ice particles. This treatment enhances the contribution of heterogeneous nucleation. Therefore, including pre-existing ice in GCMs significantly reduces  $N_i$ , as shown in simulations comparing models with and without pre-existing ice (Shi et al., 2015). As explained by Mitchell and Erfani (2025) and Mitchell and Garnier (2024), the current treatment of the pre-existing ice in GCMs leads to an overestimation of the pre-existing ice effect, which can bias the homogeneous and heterogeneous contributions and their radiative effects. Using models with higher vertical resolution, such as RCMs or large-eddy simulations (LES), can help mitigate the overestimation of pre-existing ice by better resolving vertical gradients of ice mass mixing ratio, temperature, and vertical velocity, which are critical for accurately capturing ice nucleation processes.

Another important factor in cirrus cloud modeling is the role of OGWs. OGWs are expected to promote homogeneous ice nucleation in cirrus clouds by increasing their updrafts and supersaturations. Recent studies have demonstrated that including OGWs in GCMs leads to stronger homogeneous ice nucleation, and thereby higher  $N_i$  and IWC and lower  $D_e$  (Lyu et al., 2023; Tully et al., 2022), highlighting the importance of OGWs in GCMs.

Furthermore, GCMs should account for complex processes for underlying mixed-phase clouds and their relationship with cirrus clouds. Through injecting INPs, CCT can modify cirrus cloud microphysics (e.g. reductions in  $N_i$  and increases in  $D_e$ ) which then affects the growth processes of ice particles in mixed-phase clouds that causes additional cooling (Gruber et al., 2019; Mitchell et al., 2020). This realization helped give birth to a new climate intervention method known as mixed-phase regime cloud thinning or MCT (Villanueva et al., 2022). In the CCT investigation described in Mitchell et al. (2020), most of the CCT CRE was due to mixed phase clouds that were affected by microphysical changes in the overlying cirrus clouds. This suggests that the glaciation of mixed phase clouds with subsequent CRE changes may be partly accomplished through CCT using INP concentrations on the order of  $10 \text{ L}^{-1}$  (Storelvmo et al., 2013; 2014) instead of the higher INP concentrations indicated in Villanueva et al. (2022), which were on the order of  $10^5 \text{ L}^{-1}$  in the

Arctic for producing a CRE change of  $-1 \text{ W m}^{-2}$ . This approach may also produce a CRE change or cooling effect greater than the CRE change produced by CCT or MCT alone.

Another significant gap in CCT research is the lack of process-based modeling using high vertical and/or horizontal resolutions such as LES and single column models. To the best of our knowledge, only one LES study has been conducted on CCT (Gruber et al., 2019). This limits our understanding of smaller-scale processes such as turbulence (Kärcher et al., 2025), convection, and cloud physics in cirrus clouds. In contrast, extensive LES research has been employed for another SRM method, called marine cloud brightening (MCB), in order to resolve those processes (Chun et al., 2023; Erfani et al., 2022, 2025). The knowledge gained from such studies can then be employed to improve the representation of MCB in GCMs. Similar efforts are needed for understanding processes related to CCT. In particular, two of the aforementioned issues, pre-existing ice treatment and OGW parameterization, should not be significant in high-resolution LES experiments.

## 6 Conclusions

This study investigates CCT as a climate intervention method by quantifying it as the transition from homogeneous to heterogeneous cirrus clouds. Considering the challenges of achieving rapid GHG emission reductions, it has been argued that climate intervention methods may be necessary to mitigate global warming (Baiman et al., 2024; Kriegler et al., 2018). However, modifying the environment involves many risks, including unintended consequences for air quality, weather, and climate (Blackstock et al., 2009; Pereira et al., 2021). For this reason, it is important to conduct comprehensive research in order to quantify the efficacy, risks, costs, and limitations of such methods. Even if these methods pass all necessary tests, they are not alternatives to GHG emission reduction; rather, they are intended to "buy time" for societies to avoid the worst consequences of climate change until GHG emissions (and concentrations perhaps) are reduced to safe levels.

GCMs are advantageous for identifying the global net forcing of cirrus clouds, while accounting for climate feedback processes. However, inaccurate cirrus cloud processes (e.g., INP concentrations and vertical motions at cirrus cloud levels) and resolution-dependent

parameterizations (e.g., pre-existing ice treatment) cause uncertainties in GCM simulations of  
830 CCT. For instance, GCMs that did not account for pre-existing ice predicted efficient CCT cooling  
(Storelvmo et al., 2013, 2014; Gasparini et al., 2020), while those that implemented pre-existing  
ice suggested minimal or adverse CCT effects (Gasparini and Lohmann, 2016; Tully et al., 2022,  
2023). In contrast, process-based models, such as the RTM used in this study, may more easily be  
835 constrained with satellite measurements of cirrus cloud properties and help isolate certain  
mechanisms. That knowledge can then be used to improve GCMs.

This study integrates the CALIPSO satellite retrievals described in Mitchell and Garnier (2024)  
with the libRadtran RTM to improve estimates of the radiative effects of homogeneous and  
heterogeneous cirrus clouds. Our results confirm that natural homogeneous cirrus clouds exert a  
significantly stronger CRE than natural heterogeneous cirrus, highlighting their distinct radiative  
840 properties in polar regions during winter. Building on this contrast, we estimate the instantaneous  
efficacy of CCT by defining two bounding cases: a lower bound assuming complete microphysical  
change from natural (observed) cirrus clouds to heterogeneous cirrus and formation of new cirrus,  
representing the idealized maximum cooling effect. The upper bound assumes that the atmospheric  
dynamics enable all naturally occurring homogeneous cirrus to form regardless of elevated INP  
845 concentrations from CCT, which produces warming (due to the INPs producing new cirrus clouds).  
 $\Delta\text{CRE}_{\text{max}}$  (i.e., CCT radiative effect without producing new cirrus clouds) yields surface cooling  
of  $-0.7$  to  $-1.0 \text{ W m}^{-2}$  and TOA cooling of  $-1.2$  to  $-1.6 \text{ W m}^{-2}$ , while inclusion of “new cirrus”  
formation from injected INPs in clear-sky ice-supersaturated regions partially offsets this effect,  
resulting in total surface cooling of  $-0.2$  to  $-0.5 \text{ W m}^{-2}$  and total TOA cooling of  $-0.3$  to  $-0.9 \text{ W}$   
850  $\text{m}^{-2}$  as the lower bound of CCT efficacy. These values fall within the cooling range of  $-0.25$  to  $-2$   
 $\text{W m}^{-2}$  estimated by previous GCM studies (Gasparini et al., 2020; Gasparini and Lohmann, 2016;  
Storelvmo et al., 2013; Storelvmo and Herger, 2014; Storelvmo et al., 2014). However, the upper  
bound (due to the exclusive formation of new cirrus clouds) yields a total surface warming of  $0.4$   
to  $0.6 \text{ W m}^{-2}$  and a total TOA warming of  $0.6$  to  $1.1 \text{ W m}^{-2}$ , consistent with studies reporting  
855 unexpected warming effects of CCT (Penner et al., 2015; Tully et al., 2022).

A major concern raised by previous CCT studies is *overseeding*, where injecting excessive INPs  
forms too many small ice particles through heterogeneous nucleation in cirrus clouds, leading to  
higher optical thickness, longer cloud lifetime, and ultimately a warming effect (Gasparini and

Lohmann, 2016; Penner et al., 2015; Storelvmo et al., 2013; Tully et al., 2022). A related seeding  
860 concern is the creation of new cirrus clouds in clear sky regions where the  $RH_i$  is above ice  
saturation and natural INP concentrations are relatively low. By nature, RTMs cannot directly test  
these side effects or any other adjustment or feedback process. However, regarding the latter,  
Gruber et al. (2019) investigated CCT for an Arctic case study using the ICON-ART modeling  
865 system with a horizontal resolution of 5 km and an integration time step of 25 s, and found that  
while seeding produced some new cirrus clouds, these new cirrus suppressed homogeneous  
nucleation downstream by lowering  $RH_i$  further downstream, with these two phenomena tending  
to cancel in terms of their radiative effect. And in regard to overseeding, this rarely occurred since  
homogeneous nucleation in natural cirrus was active throughout most of the model domain.  
Another concern is the potential impact of CCT on precipitation; however, this impact seems to  
870 be small as a change in global mean cirrus CRE caused by CCT was predicted to produce a global  
mean rainfall reduction of -1.3%, which is less than corresponding estimates for another climate  
engineering SRM method known as stratospheric aerosol injection (Storelvmo et al., 2014).

Over the mid-latitudes during winter, RTM simulations show CCT warming at the TOA and within  
the atmosphere and no significant impact at the surface due to competing LW and SW radiation  
875 effects: homogeneous cirrus absorbs/emits more LW radiation but also scatters more SW radiation  
than heterogeneous cirrus and these two effects cancel each other at the surface. This finding is  
consistent with Storelvmo et al. (2014), who suggested that conducting CCT globally is not more  
efficient than exclusively targeting high-latitude regions.

Sensitivity analyses reveal that the cooling efficacy of CCT is significantly affected by  
880 atmospheric thermodynamic profiles and the presence of low clouds. TOA cooling is sensitive to  
surface temperature, while surface cooling is less sensitive to changes in atmospheric water vapor.  
These findings align with previous studies (Corti and Peter, 2009; Dupont and Haeffelin, 2008),  
which demonstrated that cirrus CRE at the TOA depends on the temperature contrast between the  
Earth's surface and the cloud, whereas the cirrus CRE at the surface is reduced by a more humid  
885 atmosphere due to the absorption of downward LW radiation by water vapor. Furthermore, these  
results indicate that Arctic low clouds tend to strongly suppress the instantaneous efficacy of CCT  
by insulating the surface from the CCT atmospheric cooling. However, this strong atmospheric  
cooling suggests that CCT may still influence the surface through mixing and other feedback

mechanisms over longer timescales, even in the presence of low clouds. In addition, some studies  
890 indicated that winter-time Arctic low cloud cover has decreased in recent decades (Boccolari and  
Parmiggiani, 2018; Liu and Key, 2016; Schweiger, 2004; Wang and Key, 2003), which implies  
stronger potential for an instantaneous impact of CCT at the surface in the future.

Our study highlights the necessity of improving the representation of cirrus cloud processes in  
models, particularly the radiative contributions of cirrus clouds dominated by homogeneous and  
895 heterogeneous freezing nucleation. To more accurately quantify the efficacy of CCT, future work  
should focus on 1) using satellite retrievals of cirrus cloud properties to guide corresponding model  
parameterizations, 2) revisiting assumptions such as the treatment of pre-existing ice in GCMs, 3)  
including OGW cirrus clouds in GCMs, and 4) employing high-resolution LES experiments.  
While LES modeling has been widely used in studies of another climate intervention method (i.e.,  
900 MCB) to resolve smaller-scale processes (Chun et al., 2023; Erfani et al., 2022, 2025), its  
application to CCT remains limited to a single study (e.g., Gruber et al., 2019). Considering the  
persistent uncertainties in observing and modeling aerosol-cloud-precipitation interactions related  
to cirrus clouds, an integration of spatially and temporally high-resolution in-situ and/or remote  
sensing measurements may be essential for constraining parameterizations and for improving the  
905 representation of ice processes in LES and GCM modeling. In the future, we will incorporate  
CALIPSO retrievals of cirrus clouds into the NCAR GCM known as the Community Atmosphere  
Model, version 6 (CAM6) to quantify  $D_e$  as a function of IWC and temperature for heterogeneous  
freezing only and for observed cirrus cloud conditions (where both heterogeneous and  
homogeneous freezing are active), based on the same CALIPSO retrievals used here. This analysis  
910 will be region- and season-dependent.

**Data Availability Statement:** The MERRA2 reanalysis data is publicly available at  
<https://doi.org/10.5067/2E096JV59PK7> (Global Modeling and Assimilation Office (GMAO),  
2015). The libRadtran code is publicly accessible at <http://www.libradtran.org/doku.php> (Emde et  
915 al., 2016). The CALIPSO retrievals of IWC and  $D_e$  from Mitchell and Garnier (2024), and the  
RTM outputs in this study will be provided upon request.

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920 **Author contributions:** Both co-authors contributed to the conceptualization, methodology, and the interpretation of the results. EE developed the Python codes and conducted exploratory data analysis and RTM simulations. EE drafted the manuscript, and both co-authors provided edits and revisions.

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