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Elucidating the boundary layer turbulence dissipation rate using high-resolution measurements from a radar wind profiler network over the Tibetan Plateau

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25

Abstract

26 The planetary boundary layer (PBL) over the Tibetan Plateau (TP) exerts a significant
27 influence on regional and global climate, while its vertical structures of turbulence and
28 evolution features remain poorly understood, largely due to the scarcity of observation.
29 This study examines the vertical profile and daytime variation of turbulence dissipation
30 rate (ε) in the PBL over the TP using the high-resolution (6 min and 120 m) measurements
31 from the radar wind profiler (RWP) network, combined with the hourly data from the
32 ERA5 reanalysis. Observational analyses show that the magnitude of ε below 3km under
33 all-sky conditions exhibits large spatial discrepancy over the six RWP sites over the TP.
34 Particularly, the values of ε at Minfeng and Jiuquan over the northern TP and Dingri over
35 the southern TP are roughly an order of magnitude greater than those at Lijiang, Ganzi and
36 Hongyuan over the eastern TP. This could be partially attributed to the difference of land
37 cover across the six RWP sites. In terms of the diurnal variation, ε rapidly intensifies from
38 0900 local standard time (LST) to 1400 LST, and then gradually levels off in the late
39 afternoon. Under clear-sky conditions, both ε and planetary boundary layer height (z_i) are
40 greater, compared with cloudy-sky conditions. This reveals that clouds would suppress the
41 turbulence development and deduce z_i . In the lower PBL ($0.2 < z/z_i < 0.5$, where z is the
42 height above ground level), the dominant influential factor for the development of
43 turbulence is the surface-air temperature difference ($T_s - T_a$). By comparison, in the
44 upper PBL ($0.6 < z/z_i < 1.0$), both the $T_s - T_a$ and vertical wind shear (VWS) affect the
45 development of turbulence. Above the PBL ($1.0 < z/z_i < 2.0$), the shear production resulting
46 from VWS dominates the variation of turbulence. Under cloudy-sky conditions, clouds are
47 found to decrease the surface total solar radiation, thereby reducing $T_s - T_a$ and surface
48 sensible heat flux. This weakened sensible heat flux tends to inhibit the turbulent motion
49 within PBL especially in the lower PBL and decrease the growth rate of z_i . On the other
50 hand, the strong VWS induced by clouds enhances the turbulence above the PBL. The
51 findings obtained here underscore the importance of RWP network in revealing the fine-
52 scale structures of the PBL over the TP and gaining new insight into the PBL evolution.

53



54 1. Introduction

55 Turbulence ranks among the most intricate phenomena within the atmosphere, ensuring
56 that the planetary boundary layer (PBL) remains thoroughly mixed during daylight hours
57 (Li et al., 2023). As a result, the structure of the PBL is, to a considerable extent, governed
58 by the evolution of turbulence (Teixeira et al., 2021). Turbulence dissipation rate (ε)
59 reflects the amount of turbulent kinetic energy (TKE) that is converted into heat at the
60 Kolmogorov scale and is a measure of the turbulence intensity (McCaffrey et al., 2017;
61 Muñoz-Esparza et al., 2018). Proper parameterizations of the turbulence dissipation term
62 with the aid of observations have great impact on the model forecast skill for the weather
63 and climate, as ε strongly affects vertical turbulent mixing through its influence on TKE
64 (Yang et al., 2017). Accurate estimation of ε is crucial for understanding the structure of
65 turbulence in the PBL. To date, a variety of instruments have been used to observe or
66 retrieve the vertical profiles of ε , including sodar, radar wind profiler (RWP), radiosonde,
67 Doppler wind lidar (DWL) and ultrasonic anemometer (Dodson and Griswold, 2021;
68 Jacoby-Koaly et al., 2002; Kotthaus et al., 2023; Lv et al., 2021). Compared with the DWL,
69 the RWP exhibit better capability in capturing the turbulence structures in the cloudy sky.
70 Furthermore, it is hard for the radiosonde and ultrasonic anemometer to get the temporal
71 continuous measurements of the atmospheric turbulence, due to the high costs.

72 The Tibetan Plateau (TP), with an averaged elevation above 4000 m and an area of
73 approximately 2.5 million km², is towering into the lower and middle troposphere (Huang
74 et al., 2023). By receiving a greater amount of solar shortwave radiation, the surface layer
75 of the TP could transfer more heat through the PBL to the free atmosphere (Ma et al., 2023;
76 Wang et al., 2015). The PBL over the TP exhibits strong convective bubbling and upward
77 motions due to the lower air density and buoyancy effect, which results in significant
78 turbulence motions and turbulence-convection interactions with "popcorn" cloud structures
79 (Xu et al., 2002; Xu et al., 2023). Understand the statistical behavior of ε is key to
80 revealing the vertical structure and evolution of PBL turbulence, which could improve the
81 parameterization of PBL processes over the TP (Ma et al., 2023; Wang et al., 2015; Xu et
82 al., 2019; Zhao et al., 2019). However, due to the limited observations of turbulence



83 profiles, the daytime variation characteristics of ε over the TP and its main influential
84 mechanisms remains poorly understood.

85 A vast range of previous studies have attempted to figure out the mechanisms behind the
86 observed turbulence, but most of them are based on radiosonde measurements or model
87 simulation or reanalysis data (e.g., Banerjee et al., 2018; Che and Zhao, 2021; Wang et al.,
88 2023a). A myriad of driving mechanisms is proposed to account for the PBL development
89 over the TP, such as surface thermal and dynamic forcing, atmosphere stability, among
90 others (Chechin et al., 2023; Chen et al., 2016; Lai et al., 2021; Wang et al., 2023a; Wang
91 and Zhang, 2022). It has been demonstrated that the buoyancy term contribution on the
92 southern slope of the TP is significantly larger than that on the southeastern edge of the TP
93 (Wang et al., 2015). A larger surface-air temperature difference ($T_s - T_a$) and sensible heat
94 flux could promote the rapid growth of deep PBL in the western and southern TP (Chen et
95 al., 2013, 2016; Li et al., 2017a; Wang et al., 2016; Zhang et al., 2022). Chen et al (2016)
96 found that the weak atmosphere stability at the top of the mixed layer is a key factor
97 contributing to the rapid growth of the deep turbulence in winter over the TP.

98 Also, cloud radiative effects are found to be another significant factor to modulate the
99 evolution of daytime PBL turbulence (Bodenschatz et., 2010). The TP is characterized by
100 a high frequency of cumulus clouds which is about five times the regional mean over the
101 other areas of China (Wang et al., 2015), and the occurrence frequency of clouds over the
102 TP shows large diurnal and spatial variability, with the maxima in the afternoon in the
103 eastern TP (Wan et al., 2022). Guo et al. (2019) has revealed that the cloud tends to suppress
104 the development of summer PBL in the early afternoon across China using the fine-
105 resolution radiosonde observations. Under continuous cloudy-sky conditions, the
106 convective PBL develops slowly due to the smaller surface sensible heat compared to the
107 clear-sky conditions (Wang and Zhang, 2022). The turbulence motion and dynamic
108 structure in the PBL is contributed to the formation and development of the popcorn-like
109 convective clouds (Wang et al., 2020; Xu et al., 2002). Based on surface observation,
110 radiosonde, satellite and reanalysis data, Wang et al. (2020) pointed out that higher PBL
111 height and lower lifting condensation level due to lower temperature and lower atmospheric
112 density may enhance low cloud occurrence in the afternoon, and in turn influencing PBL
113 development over the TP. However, the turbulence structure differences between clear-sky



114 and cloudy-sky conditions are rarely explored, and the possible mechanism influencing the
115 cloud topped PBL turbulence evolution remains elusive. To the best of our knowledge,
116 most of the above-mentioned studies over the TP lack high-temporal resolution turbulence
117 profile observations. Coincidentally, there exists a RWP network of China, which provides
118 us an invaluable opportunity to characterize the PBL turbulence structure over the TP (Guo
119 et al., 2021a). Therefore, the main objective of this study is to resolve the above issues over
120 the TP, by using observations from the RWP network together with other ground-based
121 meteorological measurements and the ERA5 data. We also analyze the joint effect of
122 thermodynamic and dynamic on ε structure in the daytime (0900–1700 LST) PBL
123 through $T_s - T_a$ and vertical wind shear (VWS).

124 The remainder of this manuscript proceeds as follows, Section 2 describes the data and
125 methods used in this study. In Section 3, we analyze the spatio-temporal characteristics and
126 daytime pattern of ε over the TP and investigate the possible thermodynamic and dynamic
127 effect on PBL turbulence under clear-sky and cloudy-sky conditions. The summary and
128 conclusions are given in section 4.

129

130 **2. Data and methods**

131 ***2.1 The RWP network over the TP***

132 In this study, we use the vertical measurements of RWP data with a vertical resolution
133 of 120 m and a temporal resolution of 6-min from the RWP network over the TP, which
134 contains six operational sites (Minfeng, Jiuquan, Hongyuan, Ganzi, Lijiang and Dingri)
135 operated by the China Meteorological Administration (CMA) during the period from
136 September 1, 2022 to November 31, 2023. The spatial distribution of RWP network over
137 the TP is shown in Fig. 1, and the detailed information for each RWP site, including
138 longitude, latitude, elevation, and land cover type is given in Table 1. Among these six
139 RWP sites, the Dingri site is located in the foothills of the Himalayas with an elevation
140 more than 4,300m above sea level (AGL), dominated by the land cover of bare and alpine
141 grassland. The Lijiang site is located in the southeastern TP characterized by complex



142 terrain with an elevation of about 2,400m AGL. The Ganzi and Hongyuan sites are situated
143 in the eastern TP, with elevations ranging from 3,300 to 3,500m AGL, and whose underlay
144 is mainly alpine grassland. The Minfen and Dunhuang sites are situated in arid and semi-
145 arid zones to the north of the TP, with elevations ranging from 1,400 to 1,500m, and their
146 dominant underlying land cover is mainly bare land. Therefore, these two sites are well
147 representative of the northern TP.

148 The RWP has the capability to obtain the high-temporal resolution atmospheric
149 turbulence and wind profiles over the TP compared to the radiosonde and reanalysis, which
150 makes it possible to analyze the fine PBL structures. The low and medium detection modes
151 of RWPs can acquire the wind field and turbulence information bellow 5 km AGL
152 (McCaffrey et al., 2017; Ruan et al., 2014). The RWP provides the radial observations
153 (marked as RAD subset), including profiles of radial velocity, doppler spectral width, and
154 signal noise ratio. Also provided by the RWP is the real-time sampling data (marked as
155 ROBS subset), including the profiles of horizontal wind (direction and speed), vertical
156 velocity, and refractive index structure constant (Liu et al., 2020). There exist large
157 uncertainties in the profiling measurements from RWP, thus the quality control for both
158 RAD and ROBS subsets are indispensable before retrieving related dynamic variables over
159 the TP (Liu et al., 2020; Wang et al., 2023). For instance, the profiling measurements highly
160 deviate from the truth below 0.5 km AGL and above 5 km AGL, which are attributed to
161 the near-surface clutter and significant beam attenuation, respectively (Guo et al., 2023).
162 Thus, here only the RWP measurements at heights from 0.5 km to 5 km are utilized for
163 analysis.

164

165 ***2.2 Other miscellaneous meteorological data***

166 In this study, the hourly ground-based meteorological variables, including 2m air
167 temperature (T_a), ground surface temperature (T_s), pressure and cloud cover, are derived
168 from the six automatic weather sites over the TP. Also, 1-min rainfall observations from
169 rain gauges are used to minimize the potential influence of rainfall on the profiling
170 measurements from RWP. All these meteorological datasets are subjected to strict data-
171 quality control by the National Meteorological Information Center (NMIC) of the China



172 Meteorological Administration (Wang et al., 2023b). In addition, the hourly temperature
173 data at pressure levels from the ERA5 reanalysis data is also used in this study (Hersbach
174 et al., 2020).

175

176 **2.3 Methods**

177 *2.3.1 Retrieval of turbulence dissipation rate*

178 As a widely used ground-based equipment for detecting atmospheric wind profile (Liu
179 et al., 2020), RWP has the advantage to estimate ε since it could measure Doppler velocity
180 spectrum in the radar volume where the turbulence parcel motion accounts for the spectral
181 width broadening (Jacoby-Koaly et al., 2002; White, 1999). In this study, the spectral width
182 method is applied to retrieve ε from the RAD subset based on the hypothesis that
183 turbulence is isotropic, and the contributions to the spectral width from turbulent and non-
184 turbulent process are independent of each other (Solanki et al., 2021; White, 1999).

185 The major steps for ε retrieval can be summarized as follows: (1) the spectral width
186 variance consisting of the turbulence and non-turbulence variance is obtained from the
187 spectral width measurements. (2) The non-turbulence broadening variances are
188 decomposed into beam broadening variance due to the finite width of the beam, shear
189 broadening variance generated by the presence of a wind gradient, and broadening variance
190 arising from data processing, among others (Nastrom, 1997). (3) The turbulent broadening
191 variance is extracted from the spectral width variance by excluding above mentioned non-
192 turbulence broadening variances. (4) ε is estimated from turbulent broadening variance
193 with the main assumption of isotropic and homogeneous turbulence, as well as Gaussian
194 antenna symmetric illumination function and Gaussian radial response of the receiver
195 (White et al., 1999). For more details for the spectral width method, refer to the references
196 (Jacoby-Koaly et al., 2002; McCaffrey et al., 2017; Nastrom, 1997; Solanki et al., 2021).

197 *2.3.2 Estimation of planetary boundary layer height*

198 The planetary boundary layer height (hereafter referred to as z_i) is an important
199 parameter for characterizing fine vertical structure of the PBL, which has important
200 implications for the air mass exchange between Earth's surface and the atmosphere aloft,



201 thus affecting the cloud development and air pollutant dispersion (Dai et al., 2014; Dodson
202 and Griswold, 2021; Guo et al., 2021a; Li et al., 2017b; Wang et al., 2022).

203 Here daytime z_i at each RWP site is derived from original signal-to-noise ratio (SNR)
204 profiles from the RAD subset based on the improved threshold method (ITM). The steps
205 are briefly outlined as follows. First of all, here we use the profile of normalized SNR
206 (NSNR) to avoid instrumental inconsistencies. Secondly, the NSNR threshold is set to 0.75
207 based on the z_i estimated by the radiosonde measurements at the same site. Thirdly, the
208 profile of NSNR is scanned downwardly from the top to the ground surface. Finally, z_i
209 ultimately is determined as the height where the NSNR profile greater than 0.75 for the
210 first time. For more details for the ITM, refer to the references (Liu et al., 2019).

211 2.3.3 Vertical wind shear

212 The ROBS subset is used to calculate vertical wind shear (VWS), which is an important
213 parameter that present the dynamical effect on the development of PBL (Zhang et al., 2020).
214 VWS can be calculated as follows.

$$\text{VWS} = \left[\left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2 \right]^{1/2} \quad (1)$$

215 where u and v denote zonal and meridional wind component, respectively, z denotes the
216 sample height AGL.

217 2.3.4 Classification of cloudy- and clear-sky conditions

218 Using RWP combined with the ground-based cloud cover observations at each site, the
219 effect of clouds on daytime variations of PBL turbulence and z_i over the TP are
220 investigated. Firstly, the 1-min precipitation and 6-min RWP data are time-matched to
221 remove the profile data half an hour before and after the precipitation to obtain non-
222 precipitation data (Wu et al, 2023). Then, all-sky conditions are defined as non-
223 precipitation hours. Finally, the clear-sky (cloudy-sky) conditions are identified as hours
224 with the cloud fraction less (greater) than 30% (80%), respectively (Guo et al., 2016;
225 Solanki et al., 2021).



226 *2.3.5 Calculation of the gradient Richardson number*

227 The evolution of turbulence in the PBL has been previously recognized to be closely
228 associated with atmospheric stability (Chechin et al., 2023; Chen et al, 2013; Lai et al.,
229 2021; Muhsin et al., 2016). Therefore, we take the gradient Richardson number (Ri) as a
230 variable to characterize atmospheric stability and the formation of turbulence over the TP.
231 Following Stull (1988), Ri is formulated as follows:

$$Ri = \frac{g}{\theta_v} \frac{\partial \theta_v / \partial z}{(\partial u / \partial z)^2 + (\partial v / \partial z)^2} \quad (2)$$

232 where θ_v is the virtual potential temperature from ERA5, u and v are the hourly zonal
233 and meridional wind components derived from RWP, respectively, g is the gravitational
234 acceleration, and z represents the AGL.

235

236 **3. Results and discussion**

237 *3.1 Spatio-temporal distributions of daytime PBL turbulence dissipation rate*

238 Both the PBL turbulence dissipation rate and z_i have significant diurnal variations over
239 mountain and urban areas (Adler et al., 2014; Liu et al., 2019; Solanki et al., 2021; Yang
240 et al., 2023). Since the longitude of the six sites over the TP is ranged from 82.7°E to
241 102.6°E, it is necessary to use the Local Standard Time (LST) to accurately capture the
242 daytime variations of the PBL and make a comparison between different sites.

243 Figure 2 presents a comprehensive overview of the ε profile at 6 min intervals and
244 hourly averaged z_i in lower troposphere at heights from 0.5 km to 3.0 km for six RWP
245 sites over the TP during the period from September 1, 2022 to November 31, 2023. The
246 magnitude of ε and its vertical structures during the daytime at both Minfeng and Jiuquan
247 sites over the northern TP and at Dingri site over the southern TP stand in stark contrast to
248 those RWP sites (i.e., Lijiang, Ganzi and Hongyuan) in the eastern TP. As shown in the
249 right panels of Fig. 2, ε generally decrease with increased height at all six RWP sites. It is
250 apparent that ε exhibits a larger west-east and north-southern spatial difference under all-
251 sky conditions. In terms of the latitudinal variation, ε exhibits a decreasing trend from west



252 to east at both Minfeng and Jiuquan sites along the altitude belt of 38°N, so does the RWP
253 sites of Dingri, Lijiang, Ganzi and Hongyuan along the altitude belt of 30°N. In terms of
254 the meridional variation, ε at the two RWP sites in the northern TP have a significantly
255 larger magnitude than the other four sites. In particular, the maximum mean value of
256 daytime ε in the height range of 0.5 to 3.0 km is found at Minfeng and Jiuquan in the
257 northern TP, which reaches up to $10^{-3.59} \text{ m}^2 \text{ s}^{-3}$ and $10^{-3.73} \text{ m}^2 \text{ s}^{-3}$, respectively. By
258 comparison, the least magnitude of ε is found in the eastern TP, with the mean values of
259 $10^{-4.06} \text{ m}^2 \text{ s}^{-3}$, $10^{-4.30} \text{ m}^2 \text{ s}^{-3}$ and $10^{-4.22} \text{ m}^2 \text{ s}^{-3}$ at Lijiang, Hongyuan and Ganzi, respectively.
260 Meanwhile, the mean magnitude of ε at Dingri in the southern TP lies between the
261 magnitude of ε in the northern and eastern TP, which is $10^{-3.88} \text{ m}^2 \text{ s}^{-3}$.

262 The above results imply that the turbulence intensity at the RWP sites over the northern
263 and western TP is about one order of magnitude greater than that in the eastern TP. To
264 further investigate the possible reasons for this significant difference in ε , the relationships
265 between $T_s - T_a$ and ε for different regions are presented in Fig. 3. The mean value of
266 $T_s - T_a$ in the northern and southern TP is 14.29°C, which is greater than that of eastern
267 TP with the value of 11.26°C (Fig. 3a). The mean daytime ε for the two regions reaches
268 up to $10^{-3.74} \text{ m}^2 \text{ s}^{-3}$ and $10^{-4.20} \text{ m}^2 \text{ s}^{-3}$, respectively (Fig. 3b). Additionally, ε is significantly
269 and positively correlated with $T_s - T_a$ ($R > 0.35$, $p < 0.005$), which illustrates that the
270 thermal forcing makes an important contribution to turbulence development in the TP (Figs.
271 3c and 3d).

272 Overall, the spatial distribution of the z_i at all six RWP sites is clearly dependent on
273 geographical location (Fig. 2), which resembles that of the ε . The geographic pattern of
274 z_i from RWP agrees well with those from radiosonde measurements (Che and Zhao, 2021)
275 and reanalysis (Slättberg, 2022). The land surfaces at both Minfeng and Jiuquan sites in
276 the northern TP are dominated by barren and relatively homogenous terrain, in sharp
277 contrast to the highly vegetated underlying terrain at both Ganzi and Hongyuan sites in the
278 eastern TP (Fig. 1). The sparse vegetation in the northern TP generally comes with large
279 Bowen ratio during the daytime, which tends to produce larger sensible heat flux compared
280 to that in the eastern TP. The increased turbulence intensity in PBL is generally associated
281 with larger sensible heat flux, which has been reported by previous studies (Wang et al.,



282 2016; Zhang et al., 2022). Thus, the spatial and magnitude of ε over the TP are most likely
283 relevant to the underlying surface type.

284 Regarding the daytime pattern of turbulence (all six panels with color shading in Fig. 2),
285 the turbulence over the TP shows a pronounced signature of single-peak variability. During
286 the period 0900–1100 LST, the magnitude of ε at all six RWP sites is relatively weak.
287 From 1100 LST onward, with the increase of downward solar shortwave radiation, surface
288 sensible heat flux gradually rises, which leads to acceleration of turbulence mixing
289 processes. Then, ε reaches peak in the early afternoon (1300–1500 LST). Afterwards,
290 during the later afternoon (1500–1700 LST), ε diminishes gradually. Likewise, z_i almost
291 follows the same daytime variation pattern of ε .

292

293 ***3.2 Characteristics of daytime PBL turbulence dissipation rate under clear- and*** 294 ***cloudy-sky conditions***

295 The influence of clouds on PBL has been discussed and analyzed in previous studies
296 (e.g., Guo et al., 2016; Huang et al., 2023; Ma et al., 2023; Schumann et al., 1991; Yu et
297 al., 2004). To reveal the potential impact of clouds on PBL ε over the TP, the comparison
298 analyses between clear- and cloudy-sky conditions are presented in this section. Figure 4
299 shows the daytime cycle of mean ε profile and z_i averaged over the six RWP sites under
300 all-, clear- and cloudy-sky conditions. Overall, both the profile of ε and z_i over the TP
301 present distinct single-peak variations, and their peaks approximately occur at 1400 LST
302 (Fig. 4a). The daytime averaged ε below 3.0 km AGL is $10^{-3.95} \text{ m}^2 \text{ s}^{-3}$, and mean z_i is
303 1472 m, respectively. There is a significant positive correlation between ε and z_i during
304 the daytime ($R=0.63$, $p<0.01$).

305 Under clear-sky condition, the daytime mean ε is $10^{-3.88} \text{ m}^2 \text{ s}^{-3}$ (Fig. 4b). During the
306 period 1300–1500 LST, ε is ranged from $10^{-3.43}$ to $10^{-2.82} \text{ m}^2 \text{ s}^{-3}$ ($10^{-4.17}$ to $10^{-3.40} \text{ m}^2 \text{ s}^{-3}$) at
307 heights from 0.5 km (1.0 km) to 1.0 km (2.0 km) in lower (upper) PBL. Thus, the well-
308 mixed turbulence maintains the development of PBL in the early afternoon. Under cloudy-
309 sky condition (Fig. 4c), the daytime mean value of z_i can reach up to 1415 m, which is
310 117 m lower than that of clear-sky conditions. This means that the clouds would suppress



311 the development of the PBL turbulence in the early afternoon which has been observed by
312 the radiosonde observations described in Guo et al (2016).

313 Furthermore, the probability density distribution (PDF) of ε in PBL ($0.2 < z/z_i < 1.0$) and
314 above the PBL ($1.0 < z/z_i < 2.0$) under all-, clear- and cloudy-conditions are given in Fig. 5.
315 Here AGL (z) is normalized by z_i to provide a nondimensional vertical coordinate of z/z_i .
316 Overall, the mean ε are $10^{-3.82}$, $10^{-3.79}$ and $10^{-3.85}$ $\text{m}^2 \text{s}^{-3}$ at height of $0.2 < z/z_i < 2.0$ under all-,
317 clear- and cloudy-sky conditions, respectively (Fig. 5a). Within the PBL (Fig. 5b), the mean
318 ε under clear-sky conditions ($10^{-3.27}$ $\text{m}^2 \text{s}^{-3}$) is greater than that of under cloudy-sky
319 conditions ($10^{-3.36}$ $\text{m}^2 \text{s}^{-3}$), and the standard deviation of ε under clear-sky conditions is
320 slightly greater than that under cloudy-sky conditions. This illustrates that clouds can
321 significantly inhibit the turbulence intensity in the PBL, with the value of $\Delta\varepsilon$ between
322 clear- and cloudy-sky conditions is $-10^{-4.0}$ $\text{m}^2 \text{s}^{-3}$. However, above the PBL (Fig. 5c), ε
323 presents normal distribution characteristics, and there is no significant difference between
324 the mean ε under clear- and cloudy-sky conditions.

325 To examine the impact of clouds on the vertical structure of turbulence within and above
326 the PBL, Figure 5d shows the normalized contoured frequency by altitude diagram
327 (NCFAD) of the $\Delta\varepsilon$ for normalized (z/z_i) profiles of ε between cloudy-sky and clear-sky
328 conditions. Within the PBL, $\Delta\varepsilon$ is negative, and $|\Delta\varepsilon|$ generally decrease with increased
329 z/z_i , where $\Delta\varepsilon$ is $-10^{-4.3}$ $\text{m}^2 \text{s}^{-3}$ at $z/z_i=0.5$, and $-10^{-5.0}$ $\text{m}^2 \text{s}^{-3}$ at $z/z_i=1.0$, respectively.
330 This suggests that clouds may weaken turbulence within the PBL (Figs. 4b and 4c),
331 especially in the lower PBL ($z=820\text{m}$, $z/z_i<0.5$).

332

333 **3.3. Potential factors Influencing daytime PBL turbulence dissipation rate**

334 **3.3.1 Surface-air temperature difference**

335 The vertical structure of PBL ε and z_i over the TP show obvious spatial differences in
336 the context of a complex subsurface. The diverse land cover types lead to differences in
337 surface albedo and soil moisture, which in turn lead to distinctions in thermodynamic
338 characteristics such as sensible heat flux (Ma et al., 2023). Buoyant production driven by
339 solar heating from the surface is one of the dominant sources generating turbulence in the



340 convective PBL. The surface sensible heat flux is an important thermodynamic factor that
341 affects the buoyant convective processes (Stull, 1988). Meanwhile, previous studies (e.g.,
342 Wang et al., 2022; Yang et al., 2023) have suggested that $T_s - T_a$ can serve as a good
343 proxy for the sensible heat flux. There are not sensible heat flux measurements at six RWP
344 sites in this study, and thus we directly take $T_s - T_a$ as a proxy thermodynamic variable to
345 analyze its potential connection to variation of PBL turbulence.

346 Figure 6 shows the magnitude of ε varies as a function of $T_s - T_a$ for all six sites,
347 within ($0.2 < z/z_i < 1.0$) and above ($1.0 < z/z_i < 2.0$) the PBL, under all, clear- and cloudy-sky
348 conditions, respectively. $T_s - T_a$ are first classified into five bins, which are then
349 statistically analyzed against the corresponding ε averaged for z/z_i values between 0.2
350 and 2.0 to obtain regression equations incorporating slopes. Further, Table 2 shows the
351 scatter plots between $\text{Log}_{10}\varepsilon$ and $T_s - T_a$ (and VWS) at different altitude ranges under
352 all-, clear- and cloudy-sky conditions. $\text{Log}_{10}\varepsilon$ is found to be linearly correlated with $T_s -$
353 T_a (and VWS) ($p < 0.05$). The surface sensible heat flux generally increases with increased
354 $T_s - T_a$, thus the increased $T_s - T_a$ intensifies the turbulence in PBL ($0 < z/z_i < 1.0$), which
355 is shown in Fig. 6b, e, h. Within the PBL, ε is also positively correlated with $T_s - T_a$
356 whose slope values are larger than those at $0.2 < z/z_i < 2.0$. As $T_s - T_a$ rises, the larger
357 surface sensible heat flux would lead to enhanced buoyancy process and turbulent motion
358 within the PBL. On the other hand, ε above the PBL is negatively correlated with $T_s -$
359 T_a (Figs. 6c, f, i). This suggests that $T_s - T_a$ dramatically affects the development of
360 turbulence within the PBL, whereas it has little effect on the turbulence above the PBL.

361 Within the PBL, the magnitude of slope (slope=0.019) under clear-sky conditions is
362 larger than that of under-cloudy conditions (slope = 0.015) as shown in Figs. 6e and 6h.
363 This implies that $T_s - T_a$ is one of the dominant factors affecting the PBL turbulence,
364 particularly under the clear-sky conditions. Given that turbulence in the mixed PBL over
365 the TP is usually driven by convection (Xu et al., 2023), as $T_s - T_a$ decreases when clouds
366 are present, less heat is transferred from the surface to the atmosphere, reducing the
367 buoyancy flux and leading to weaker turbulence in the PBL, especially for the lower PBL
368 ($0.2 < z/z_i < 0.5$), as shown in Figures 4b and 4c. Consequently, the clouds tend to suppress
369 the development of PBL (Fig. 5a) and reduce z_i .



370 *3.2.2 Vertical wind shear*

371 Besides $T_s - T_a$, VWS is another crucial dynamic parameter that is related to the
372 mechanical turbulence within the PBL. Similar to Fig. 6, Figure 7 presents the relationship
373 between ε and VWS (both normalized by z_i) within and above the PBL under all-, clear-
374 and cloudy-sky conditions, respectively. The near-surface clutter significantly increases
375 the uncertainty of RWP data, which leads to incapability of analyzing the effect of wind
376 shear on ε below 0.5 km AGL in the following sections.

377 Regardless of within or above the PBL, ε is positively correlated with VWS as shown
378 in Fig. 7a, d, g and Tabel 2, which indicates that larger VWS leads to stronger turbulence.
379 This suggests that the dynamic effect of VWS promotes the development of turbulence.
380 Within the PBL (Figs. 7b, e, h), the slope of ε against VWS are smaller than at
381 $0.2 < z/z_i < 2.0$ with values ranging from 9.5 to 10.3. Above the PBL (Figs. 7c, f, i), the
382 values of the slope are larger with values ranging from 10.7 to 18.1, which demonstrating
383 that the dynamical effects of VWS influence the development of turbulence both within
384 and above the PBL.

385 Under cloudy-sky conditions (Figs. 7h, i), the effect of VWS on turbulence within the
386 PBL (Slope = 10.3) is weaker than above the PBL (Slope = 18.1), significantly. Compared
387 to the clear-sky conditions (Figs. 7e, f), the values of the slopes are larger for that of under
388 cloudy-sky conditions (Figs. 7h, i) both within and above the PBL. Remarkably, above the
389 PBL, the effect of clouds on turbulence is more dramatic, as the slope value under cloudy-
390 sky conditions is nearly twice as large as that of under clear-sky conditions. These results
391 indicate the significant mechanical processes driven by VWS is important in the
392 development of turbulence. A larger VWS in the PBL corresponds to stronger turbulence.
393 Besides, above the PBL, the mechanical process of VWS is enhanced under cloudy-sky
394 conditions.

395

396 *3.2.3. Joint influence of $T_s - T_a$, VWS and atmosphere stability on ε*

397 It was stated that turbulence can be produced by buoyant convective processes (i.e.,
398 thermals of warm air rising) and by mechanical processes (i.e., wind shear). From the



399 previous section, it is known that $T_s - T_a$ and VWS both affect the development of PBL
400 turbulence. Figure 8 gives the slope profiles of ε against $T_s - T_a$ and VWS at normalized
401 heights (z/z_i) under all-, clear- and cloudy-sky conditions, respectively.

402 As inferred from the previous findings, $T_s - T_a$ primarily influences turbulence
403 development within the PBL, irrespective of clear-sky and cloudy-sky conditions (Fig. 6).
404 Figure 8a shows that the slope values within the PBL are predominantly positive, and the
405 slope value decreases rapidly with height, which indicates that the influence of $T_s - T_a$ on
406 PBL turbulence experiences decreasing trend with height. Interestingly, there is a nearly
407 linear of slope value from the lower PBL to a smaller positive value near the top of the
408 PBL. Above the PBL, the slope value becomes negative. This may be due to the linear
409 decrease of heat flux transport and buoyancy term in the convective PBL (Stull, 1988).
410 Therefore, these findings highlight the predominant thermal forcing of $T_s - T_a$ on
411 turbulence development within the lower PBL. In addition, when $0.2 < z/z_i < 0.5$, the slope
412 values are larger for clear-sky conditions than for cloudy-sky conditions, while there is
413 little difference for the clear-sky and cloudy-sky conditions when $0.5 < z/z_i$. Hence, under
414 clear-sky conditions, the thermodynamic effect of $T_s - T_a$ is more pronounced within the
415 lower PBL.

416 As shown in Fig. 7, it is evident that VWS influences turbulence development within
417 and above the PBL. Figure 8b shows that when $0.2 < z/z_i < 2.0$, the slope values are
418 consistently positive, indicating that VWS predominantly affects turbulence development
419 within the mid-, upper- PBL and above the PBL. Moreover, when $0.2 < z/z_i < 1.2$, the slope
420 values increase with height. However, when $1.4 < z/z_i < 2.0$, the slope values exhibit a
421 decreasing trend, which suggesting a diminishing influence of VWS. Additionally, within
422 the PBL ($0.2 < z/z_i < 0.7$), the slope values under clear-sky conditions are close to those
423 under cloudy-sky conditions, while the slope values under cloudy-sky conditions are even
424 greater when $0.7 < z/z_i < 2.0$. For instance, when $z/z_i = 1.4$, $\text{Slope}_{\text{Clear-sky}} = 14.6$, while
425 $\text{Slope}_{\text{Cloudy-sky}} = 27.0$, indicating that the latter is 1.8 times larger than the former. These
426 results suggest that clouds are primarily responsible for enhancing mechanical processes
427 from VWS on turbulence within the upper PBL and above the PBL.



428 Furthermore, it can be concluded that, $T_s - T_a$ is the thermodynamic factor influencing
429 turbulence development within the lower PBL ($0.2 < z/z_i < 0.5$), both $T_s - T_a$ and VWS
430 jointly strengthen turbulence development in the upper PBL ($0.6 < z/z_i < 1.0$), and VWS
431 emerges as the predominant factor affecting turbulence development above the PBL
432 ($1.0 < z/z_i < 2.0$) (Figs. 8a, b).

433 The previous sections have revealed that hours of both high $T_s - T_a$ and strong wind
434 shear would strengthen the turbulence within the PBL. Therefore, it's necessary to analyze
435 the combined influence of thermodynamics and dynamics factors on the development of
436 turbulence. Figure 9 presents the joint distribution of ε with $T_s - T_a$ and VWS within and
437 above the PBL under all-, clear- and cloudy-sky conditions. Within the PBL (Figs. 9b, e,
438 h), higher $T_s - T_a$ and VWS correspond to stronger turbulence (Fig. 8). In contrast, the
439 thermodynamic effect of $T_s - T_a$ on turbulence has diminished and is no longer a
440 dominant factor above the PBL, while the dynamical effect of VWS becomes the dominant
441 factor (Figs. 9c, f, i). Compared to clear-sky conditions, both $T_s - T_a$ and VWS decrease
442 under cloudy-sky conditions (Fig. 9h). This means that the weakening of both
443 thermodynamic and dynamic effects leads to a decrease in turbulence, thereby inhibiting
444 the development of turbulence within the PBL. Therefore, under cloudy-sky conditions,
445 although the VWS is reduced, the dynamical effect of VWS on turbulence is strengthened
446 (Figs. 7i and 8b), which in turn strengthens turbulence.

447 Since buoyant and mechanistic forcing jointly influence the turbulence within the PBL,
448 and VWS only represents the dynamic driving effect, it cannot accurately portray the effect
449 of thermodynamic and dynamic effects on the PBL turbulence. The gradient Richardson
450 number (Ri), on the other hand, is one of the important parameters characterizing
451 atmospheric stability and can compare the buoyant turbulence production term and the
452 shear production term in the form of a dimensionless ratio.

453 Similar to Fig. 9, the joint distribution of ε with $T_s - T_a$ and Ri within and above the
454 PBL under all-sky, clear-sky and cloudy-sky conditions is given in Fig. 10. Within the PBL
455 (Figs. 10b, e, h), it is evident that the turbulence in the PBL is enhancing for unstable
456 conditions. Furthermore, under clear-sky conditions (Fig. 10e), the maximum number of
457 samples is found when $Ri < 1.0$ and $T_s - T_a > 21.1$ in strongly unstable conditions caused



458 by the buoyancy forcing driven by the larger $T_s - T_a$. While the effect of Ri on turbulence
459 is relatively weakened above the PBL (Figs. 10c, f, i).

460 **4 Summary and concluding remarks**

461 This study investigates the characteristics of spatio-temporal distribution of daytime
462 PBL turbulence dissipation rate (ε) based on more than one-year record (September 2022–
463 October 2023) of profiling measurements from a radar wind profilers (RWP) network on
464 the Tibet Plateau (TP). Also analyzed are the evolution of ε in the PBL and the possible
465 influential mechanisms.

466 First of all, ε is firstly retrieved from the vertical wind measurements from RWP using
467 the spectral width method. Afterwards, the spatial pattern of ε is examined. Results shows
468 that the values of ε at both Minfeng and Jiuquan sites in the northern TP, and at Dingri
469 over the southern TP are about one order of magnitude greater than those at the RWP sites
470 of Lijiang, Ganzi and Hongyuan over the eastern TP. Coincidentally, Minfeng and Junquan
471 are dominated by bare or semiarid land, as opposed to the highly vegetation-covered land
472 surface at Lijiang, Ganzi and Hongyuan. This suggests the spatial discrepancy of ε over
473 the TP is highly relevant to the types of underlying land cover.

474 Although ε exhibits a variety of magnitudes among the six RWPs, the daytime pattern
475 and vertical structure of ε are similar. Turbulence reaches the peak in the early afternoon
476 (1300–1500 LST), coinciding with the highest PBL top. Under cloudy-sky conditions, the
477 daytime mean value of ε is $10^{-4.02} \text{ m}^2 \text{ s}^{-3}$, and the daytime mean value of the PBL height
478 (z_i) can reach up to 1415 m, which is 117 m lower than that of clear-sky conditions,
479 indicating that clouds would suppress the development of the PBL turbulence.

480 As far as both the thermodynamic and dynamic forcings are concerned, surface-air
481 temperature difference ($T_s - T_a$) and vertical wind shear (VWS) variables are examined by
482 performing correlation analysis with ε . The slope values of ε against $T_s - T_a$ under clear-
483 sky conditions is larger (slope=0.019) than under-cloudy conditions (slope = 0.013) within
484 the PBL, while those values are negative above the PBL. The slope values of ε against
485 VWS is positive regardless of within or above the PBL, where the largest value of 18.1 is



486 observed above the PBL under cloudy-sky conditions, and the smallest value of 9.5
487 observed in the PBL within clear-sky conditions.

488 Both the thermodynamic effect of $T_s - T_a$ and the dynamic effect of VWS enhance the
489 development of turbulence under clear-sky or cloudy-sky conditions in the PBL. In the
490 lower PBL ($0.2 < z/z_i < 0.5$), $T_s - T_a$ has a larger positive slope with ε , which suggests that
491 thermal forcing emerges as the dominant factor influencing development of the turbulence
492 and PBL. By comparison, in the upper PBL ($0.6 < z/z_i < 1.0$), $T_s - T_a$ and VWS jointly
493 influence the development of turbulence, with larger $T_s - T_a$ leading to unstable
494 atmospheric stability and stronger turbulence. Above the PBL ($1.0 < z/z_i < 2.0$), VWS
495 becomes the dominant factor influencing the development of turbulence. Compared to
496 clear-sky conditions, on one hand, clouds would diminish $T_s - T_a$, resulting in decreased
497 heat transfer from the surface to the PBL top, thereby weakening turbulence within the
498 lower PBL ($0.2 < z/z_i < 0.5$), inhibiting PBL development, and decreasing z_i . On the other
499 hand, the stronger wind shear process would enhance the turbulence above PBL under the
500 cloudy-sky conditions.

501 Although the above-mentioned findings of the PBL turbulence over the TP are the first
502 results from profiling network observations to the best of our knowledge, fine-resolution
503 spatial distribution remains unclear, largely due to the sparse distribution of RWP network
504 on the TP. On top of this, the role of roughness length remains known in the variation and
505 evolution of turbulence, especially in the lowest part of PBL, which warrants a field
506 campaign involved in the high-density turbulence observation network along with high-
507 resolution satellite images.

508

509 **Data Availability**

510 The authors would like to acknowledge the National Meteorological Information Centre
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520 **Author Contributions**

521 The study was completed with close cooperation between all authors. JG conceived of the
522 idea for this work. DM performed the analysis, DM and JG drafted the original manuscript
523 with contributions from XG, NL, YS, ZZ and NT. YW, HL, FZ, BT, HX and TC provided
524 useful suggestions and comments for the study and helped revise the manuscript.

525 **Completing interests**

526 The authors declare that they have no conflict of interest.

527 **References**

- 528 Adler, B. and Kalthoff, N.: Multi-scale transport processes observed in the boundary layer
529 over a mountainous island, *Bound.-Layer Meteor.*, 153, 515-537,
530 <https://doi.org/10.1007/s10546-014-9957-8>, 2014.
- 531 Banerjee, T., Brugger, P., De Roo, F., Kröniger, K., Yakir, D., Rotenberg, E., and Mauder,
532 M.: Turbulent transport of energy across a forest and a semiarid shrubland, *Atmos.*
533 *Chem. Phys.*, 18, 10025–10038, <https://doi.org/10.5194/acp-18-10025-2018>, 2018.
- 534 Bianco, L., Wilczak, J. M., and White, A. B.: Convective boundary layer depth estimation
535 from wind profilers: Statistical comparison between an automated algorithm and
536 expert estimations, *J. Atmos. Ocean. Technol.*, 25, 1397-1413,
537 <https://doi.org/10.1175/2008jtecha981.1>, 2008.



- 538 Bodenschatz, E., Malinowski, S. P., Shaw, R. A., and Stratmann, F.: Can we understand
539 clouds without turbulence? *Science*, 327, 970-971,
540 <https://doi.org/10.1126/science.1185138>, 2010.
- 541 Che, J. H. and Zhao, P.: Characteristics of the summer atmospheric boundary layer height
542 over the Tibetan Plateau and influential factors, *Atmos. Chem. Phys.*, 21, 5253-5268,
543 <https://doi.org/10.5194/acp-21-5253-2021>, 2021.
- 544 Chechin, D. G., Lüpkes, C., Hartmann, J., Ehrlich, A., and Wendisch, M.: Turbulent
545 structure of the Arctic boundary layer in early summer driven by stability, wind shear
546 and cloud-top radiative cooling: ALOUD airborne observations, *Atmos. Chem.*
547 *Phys.*, 23, 4685-4707, <https://doi.org/10.5194/acp-23-4685-2023>, 2023.
- 548 Chen, X. L., Añel, J. A., Su, Z. B., de la Torre, L., Kelder, H., van Peet, J., and Ma, Y. M.:
549 The deep atmospheric boundary layer and its significance to the stratosphere and
550 troposphere exchange over the Tibetan Plateau, *PLoS One*, 8, 9,
551 <https://doi.org/10.1371/journal.pone.0056909>, 2013.
- 552 Chen, X. L., Skerlak, B., Rotach, M. W., Añel, J. A., Su, Z., Ma, Y. M., and Li, M. S.:
553 Reasons for the extremely high-ranging planetary boundary layer over the western
554 Tibetan Plateau in winter, *J. Atmos. Sci.*, 73, 2021-2038, [https://doi.org/10.1175/jas-](https://doi.org/10.1175/jas-d-15-0148.1)
555 [d-15-0148.1](https://doi.org/10.1175/jas-d-15-0148.1), 2016.
- 556 Dai, C., Wang, Q., Kalogiros, J. A., Lenschow, D. H., Gao, Z., and Zhou, M.: Determining
557 boundary-layer height from aircraft measurements, *Bound.-Layer Meteor.*, 152, 277-
558 302, <https://doi.org/10.1007/s10546-014-9929-z>, 2014.
- 559 Dodson, D. S. and Griswold, J. D. S.: Turbulent and boundary layer characteristics during
560 VOCALS-REx, *Atmos. Chem. Phys.*, 21, 1937-1961, [https://doi.org/10.5194/acp-21-](https://doi.org/10.5194/acp-21-1937-2021)
561 [1937-2021](https://doi.org/10.5194/acp-21-1937-2021), 2021.
- 562 Guo, J. P., Li, Y., Cohen, J. B., Li, J., Chen, D. D., Xu, H., Liu, L., Yin, J. F., Hu, K. X.,
563 and Zhai, P. M.: Shift in the temporal trend of boundary layer height in China using
564 long-term (1979-2016) radiosonde data, *Geophys. Res. Lett.*, 46, 6080-6089,
565 <https://doi.org/10.1029/2019gl082666>, 2019.
- 566 Guo, J. P., Miao, Y. C., Zhang, Y., Liu, H., Li, Z. Q., Zhang, W. C., He, J., Lou, M. Y.,
567 Yan, Y., Bian, L. G., and Zhai, P.: The climatology of planetary boundary layer height



- 568 in China derived from radiosonde and reanalysis data, *Atmos. Chem. Phys.*, 16,
569 13309-13319, <https://doi.org/10.5194/acp-16-13309-2016>, 2016.
- 570 Guo, J. P., Liu, B. M., Gong, W., Shi, L. J., Zhang, Y., Ma, Y. Y., Zhang, J., Chen, T. M.,
571 Bai, K. X., Stoffelen, A., de Leeuw, G., and Xu, X. F.: First comparison of wind
572 observations from ESA's satellite mission Aeolus and ground-based radar wind
573 profiler network of China, *Atmos. Chem. Phys.*, 21, 2945-2958,
574 <https://doi.org/10.5194/acp-21-2945-2021>, 2021a.
- 575 Guo, J. P., Zhang, J., Yang, K., Liao, H., Zhang, S. D., Huang, K. M., Lv, Y. M., Shao, J.,
576 Yu, T., Tong, B., Li, J., Su, T. N., Yim, S. H. L., Stoffelen, A., Zhai, P. M., and Xu,
577 X. F.: Investigation of near-global daytime boundary layer height using high-
578 resolution radiosondes: first results and comparison with ERA5, MERRA-2, JRA-55,
579 and NCEP-2 reanalyses, *Atmos. Chem. Phys.*, 21, 17079-17097,
580 <https://doi.org/10.5194/acp-21-17079-2021>, 2021b.
- 581 Guo, X. R., Guo, J. P., Zhang, D. L., and Yun, Y. X.: Vertical divergence profiles as
582 detected by two wind-profiler mesonets over East China: Implications for nowcasting
583 convective storms, *Q. J. R. Meteorol. Soc.*, 149, 1629-1649,
584 <https://doi.org/10.1002/qj.4474>, 2023.
- 585 Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas,
586 J., Peubey, C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan,
587 X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara, G.,
588 Dahlgren, P., Dee, D., Diamantakis, M., Dragani, R., Flemming, J., Forbes, R.,
589 Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hogan, R. J., Hólm, E., Janisková,
590 M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay, P., Rozum,
591 I., Vamborg, F., Villaume, S., and Thépaut, J. N.: The ERA5 global reanalysis, *Q. J.*
592 *R. Meteorol. Soc.*, 146, 1999-2049, <https://doi.org/10.1002/qj.3803>, 2020.
- 593 Huang, J. P., Zhou, X. J., Wu, G. X., Xu, X. D., Zhao, Q. Y., Liu, Y. M., Duan, A. M., Xie,
594 Y. K., Ma, Y. M., Zhao, P., Yang, S., Yang, K., Yang, H. J., Bian, J. C., Fu, Y. F., Ge,
595 J. M., Liu, Y. Z., Wu, Q. G., Yu, H. P., Wang, B. B., Bao, Q., and Qie, K.: Global
596 climate impacts of land-surface and atmospheric processes over the Tibetan Plateau,
597 *Rev. Geophys.*, 61, 39, <https://doi.org/10.1029/2022rg000771>, 2023.



- 598 Jacoby-Koaly, S., Campistron, B., Bernard, S., Bénech, B., Girard-Ardhuin, F., Dessens,
599 J., Dupont, E., and Carissimo, B.: Turbulent dissipation rate in the boundary layer via
600 UHF wind profiler Doppler spectral width measurements, *Bound.-Layer Meteor.*, 103,
601 361-389, <https://doi.org/10.1023/a:1014985111855>, 2002.
- 602 Kotthaus, S., Bravo-Aranda, J. A., Coen, M. C., Guerrero-Rascado, J. L., Costa, M. J.,
603 Cimini, D., O'Connor, E. J., Hervo, M., Alados-Arboledas, L., Jiménez-Portaz, M.,
604 Mona, L., Ruffieux, D., Illingworth, A., and Haeffelin, M.: Atmospheric boundary
605 layer height from ground-based remote sensing: a review of capabilities and
606 limitations, *Atmos. Meas. Tech.*, 16, 433-479, [https://doi.org/10.5194/amt-16-433-](https://doi.org/10.5194/amt-16-433-2023)
607 2023, 2023.
- 608 Lai, Y., Chen, X. L., Ma, Y. M., Chen, D. L., and Zhaxi, S. L.: Impacts of the westerlies
609 on planetary boundary layer growth over a valley on the north side of the Central
610 Himalayas, *J. Geophys. Res.-Atmos.*, 126, 20, <https://doi.org/10.1029/2020jd033928>,
611 2021.
- 612 Li, Y., Y. Wu, J. Tang, P. Zhu, Z. Gao, and Y. Yang. Quantitative evaluation of wavelet
613 analysis method for turbulent flux calculation of non-stationary series, *Geophys. Res.*
614 *Let.*, 50(5), e2022GL101591, <http://dx.doi.org/10.1029/2022GL101591>, 2023.
- 615 Li, Z. G., Lyu, S. H., Wen, L. J., Zhao, L., Ao, Y. H., and Wang, S. Y.: Effect of a cold,
616 dry air incursion on atmospheric boundary layer processes over a high-altitude lake in
617 the Tibetan Plateau, *Atmos. Res.*, 185, 32-43,
618 <https://doi.org/10.1016/j.atmosres.2016.10.024>, 2017a.
- 619 Li, Z. Q., Guo, J. P., Ding, A. J., Liao, H., Liu, J. J., Sun, Y. L., Wang, T. J., Xue, H. W.,
620 Zhang, H. S., and Zhu, B.: Aerosol and boundary-layer interactions and impact on air
621 quality, *Natl. Sci. Rev.*, 4, 810-833, <https://doi.org/10.1016/10.1093/nsr/nwx117>,
622 2017b.
- 623 Liu, B. M., Guo, J. P., Gong, W., Shi, L. J., Zhang, Y., and Ma, Y. Y.: Characteristics and
624 performance of wind profiles as observed by the radar wind profiler network of China,
625 *Atmos. Meas. Tech.*, 13, 4589-4600, <https://doi.org/10.5194/amt-13-4589-2020>, 2020.
- 626 Liu, B. M., Ma, Y. Y., Guo, J. P., Gong, W., Zhang, Y., Mao, F. Y., Li, J., Guo, X. R., and
627 Shi, Y. F.: Boundary layer heights as derived from ground-based radar wind profiler



- 628 in Beijing, *IEEE Trans. Geosci. Remote Sensing*, 57, 8095-8104,
629 <https://doi.org/10.1109/tgrs.2019.2918301>, 2019.
- 630 Lv, Y. M., Guo, J. P., Li, J., Cao, L. J., Chen, T. M., Wang, D., Chen, D. D., Han, Y., Guo,
631 X. R., Xu, H., Liu, L., Solanki, R., and Huang, G.: Spatiotemporal characteristics of
632 atmospheric turbulence over China estimated using operational high-resolution
633 soundings, *Environ. Res. Lett.*, 16, 13, <https://doi.org/10.1088/1748-9326/abf461>,
634 2021.
- 635 Ma, Y. M., Yao, T. D., Zhong, L., Wang, B. B., Xu, X. D., Hu, Z. Y., Ma, W. Q., Sun, F.
636 L., Han, C. B., Li, M. S., Chen, X. L., Wang, J. M., Li, Y. Q., Gu, L. L., Xie, Z. P.,
637 Liu, L., Sun, G. H., Wang, S. J., Zhou, D. G., Zuo, H. C., Xu, C., Liu, X., Wang, Y.
638 J., and Wang, Z. Y.: Comprehensive study of energy and water exchange over the
639 Tibetan Plateau: A review and perspective: From GAME/Tibet and CAMP/Tibet to
640 TORP, TPEORP, and TPEITORP, *Earth-Sci. Rev.*, 104312, 2023.
- 641 McCaffrey, K., Bianco, L., and Wilczak, J. M.: Improved observations of turbulence
642 dissipation rates from wind profiling radars, *Atmos. Meas. Tech.*, 10, 2595-2611,
643 <https://doi.org/10.5194/amt-10-2595-2017>, 2017.
- 644 Muhsin, M., Sunilkumar, S. V., Ratnam, M. V., Parameswaran, K., Murthy, B. V. K.,
645 Ramkumar, G., and Rajeev, K.: Diurnal variation of atmospheric stability and
646 turbulence during different seasons in the troposphere and lower stratosphere derived
647 from simultaneous radiosonde observations at two tropical stations, in the Indian
648 Peninsula, *Atmos. Res.*, 180, 12-23, <https://doi.org/10.1016/j.atmosres.2016.04.021>,
649 2016.
- 650 Muñoz-Esparza, D., Sharman, R. D., and Lundquist, J. K.: Turbulence Dissipation Rate in
651 the Atmospheric Boundary Layer: Observations and WRF Mesoscale Modeling
652 during the XPIA Field Campaign, *Mon. Weather Rev.*, 146, 351-371,
653 <https://doi.org/10.1175/mwr-d-17-0186.1>, 2018.
- 654 Nastrom, G. D.: Doppler radar spectral width broadening due to beamwidth and wind shear,
655 *Ann. Geophys.-Atmos. Hydrospheres Space Sci.*, 15, 786-796,
656 <https://doi.org/10.1007/s00585-997-0786-7>, 1997.



- 657 Ruan, Z., Mu, R. Q., Wei, M., and Ge, R. S.: Spectrum analysis of wind profiling radar
658 measurements, *J. Meteorol. Res.*, 28, 656-667, [https://doi.org/10.1007/s13351-014-](https://doi.org/10.1007/s13351-014-3171-y)
659 3171-y, 2014.
- 660 Schumann, U. and Moeng, C. H.: Plume budgets in clear and cloudy convective boundary
661 layers, *J. Atmos. Sci.*, 48, 1758-1770, [https://doi.org/10.1175/1520-](https://doi.org/10.1175/1520-0469(1991)048<1758:Pbicac>2.0.Co;2)
662 0469(1991)048<1758:Pbicac>2.0.Co;2, 1991.
- 663 Slättberg, N., Lai, H. W., Chen, X. L., Ma, Y. M., and Chen, D. L.: Spatial and temporal
664 patterns of planetary boundary layer height during 1979-2018 over the Tibetan Plateau
665 using ERA5, *Int. J. Climatol.*, 42, 3360-3377, <https://doi.org/10.1002/joc.7420>, 2022.
- 666 Solanki, R., Guo, J. P., Lv, Y. M., Zhang, J., Wu, J. Y., Tong, B., and Li, J.: Elucidating
667 the atmospheric boundary layer turbulence by combining UHF radar wind profiler
668 and radiosonde measurements over urban area of Beijing, *Urban CLim.*, 43, 13,
669 <https://doi.org/10.1016/j.uclim.2022.101151>, 2022.
- 670 Solanki, R., Guo, J. P., Li, J., Singh, N., Guo, X. R., Han, Y., Lv, Y. M., Zhang, J., and Liu,
671 B. M.: Atmospheric-boundary-layer-height variation over mountainous and urban
672 sites in Beijing as derived from radar wind-profiler measurements, *Bound.-Layer*
673 *Meteor.*, 181, 125-144, <https://doi.org/10.1007/s10546-021-00639-9>, 2021.
- 674 Stull, R. B.: Mean Boundary Layer Characteristics, in: *An Introduction to boundary layer*
675 *meteorology*, edited by: Stull, R. B., Springer Netherlands, Dordrecht, 1–27,
676 https://doi.org/10.1007/978-94-009-3027-8_1, 1988.
- 677 Teixeira, J., Piepmeier, J. R., Nehrir, A. R., Ao, C. O., Chen, S. S., Clayson, C. A., Fridlind,
678 A. M., Lebsock, M., McCarty, W., Salmun, H., Santanello, J. A., Turner, D. D., Wang,
679 Z., and Zeng, X.: Toward a global planetary boundary layer observing system: the
680 NASA PBL incubation study team report, NASA PBL Incubation Study Team, 134
681 pp., available at: [https://science.nasa.gov/science-red/s3fs-](https://science.nasa.gov/science-red/s3fs-public/atoms/files/NASAPBLIncubationFinalReport.pdf)
682 [public/atoms/files/NASAPBLIncubationFinalReport.pdf](https://science.nasa.gov/science-red/s3fs-public/atoms/files/NASAPBLIncubationFinalReport.pdf), last access: 15 November
683 2021.
- 684 Wan, X., Zheng, J. F., Wan, R., Xu, G. R., Qin, J. F., and Yi, L.: Intercomparison of cloud
685 vertical structures over four different sites of the eastern slope of the Tibetan Plateau
686 in summer using Ka-band millimeter-wave radar measurements, *Remote Sens.*, 14,
687 19, <https://doi.org/10.3390/rs14153702>, 2022.



- 688 Wang, C. X., Ma, Y. M., and Han, C. B.: Research on the atmospheric boundary layer
689 structure and its development mechanism in the Tibetan Plateau, *Adv. Atmos. Sci.*,
690 38, 414-428, 2023a.
- 691 Wang, M. Z. and Zhang, J. T.: The relationship among summer atmospheric boundary layer
692 height over the Taklimakan Desert, its land surface parameters and Eurasian
693 circulation, *Atmos. Sci. Lett.*, 23, 13, <https://doi.org/10.1002/asl.1122>, 2022.
- 694 Wang, M. Z., Lu, H., Ming, H., and Zhang, J. T.: Vertical structure of summer clear-sky
695 atmospheric boundary layer over the hinterland and southern margin of Taklamakan
696 Desert, *Meteorol. Appl.*, 23, 438-447, <https://doi.org/10.1002/met.1568>, 2016.
- 697 Wang, S. Q., Guo, J. P., Xian, T., Li, N., Meng, D. L., Li, H. J., and Cheng, W.:
698 Investigation of low-level supergeostrophic wind and Ekman spiral as observed by a
699 radar wind profiler in Beijing, *Front. Environ. Sci.*, 11, 9,
700 <https://doi.org/10.3389/fenvs.2023.1195750>, 2023b.
- 701 Wang, Y. J., Zeng, X. B., Xu, X. D., Xie, F. Q., and Zhao, Y.: Improving the estimate of
702 summer daytime planetary boundary layer height over land from GPS radio
703 occultation data, *Geophys. Res. Lett.*, 49, 9, <https://doi.org/10.1029/2021gl096304>,
704 2022.
- 705 Wang, Y. J., Xu, X. D., Zhao, T. L., Sun, J. H., Yao, W. Q., and Zhou, M. Y.: Structures
706 of convection and turbulent kinetic energy in boundary layer over the southeastern
707 edge of the Tibetan Plateau, *Sci. China-Earth Sci.*, 58, 1198-1209,
708 <https://doi.org/10.1007/s11430-015-5054-1>, 2015.
- 709 Wang, Y. J., Zeng, X. B., Xu, X. D., Welty, J., Lenschow, D. H., Zhou, M. Y., and Zhao,
710 Y.: Why are there more summer afternoon low clouds over the Tibetan Plateau
711 compared to eastern China? *Geophys. Res. Lett.*, 47, 10,
712 <https://doi.org/10.1029/2020gl089665>, 2020.
- 713 White, A. B., Lataitis, R. J., and Lawrence, R. S.: Space and time filtering of remotely
714 sensed velocity turbulence, *J. Atmos. Ocean. Technol.*, 16, 1967-1972,
715 [https://doi.org/10.1175/1520-0426\(1999\)016<1967:Satfor>2.0.Co;2](https://doi.org/10.1175/1520-0426(1999)016<1967:Satfor>2.0.Co;2), 1999.
- 716 Wu, J. Y., Guo, J. P., Yun, Y. X., Yang, R. F., Guo, X. R., Meng, D. L., Sun, Y. P., Zhang,
717 Z., Xu, H., and Chen, T. M.: Can ERA5 reanalysis data characterize the pre-storm



- 718 environment? *Atmos. Res.*, 297, 18, <https://doi.org/10.1016/j.atmosres.2023.107108>,
719 2024.
- 720 Xu, L. J., Liu, H. Z., Du, Q., and Xu, X. D.: The assessment of the planetary boundary
721 layer schemes in WRF over the central Tibetan Plateau, *Atmos. Res.*, 230, 12,
722 <https://doi.org/10.1016/j.atmosres.2019.104644>, 2019.
- 723 Xu, X. D., Tang, Y., Wang, Y. J., Zhang, H. S., Liu, R. X., and Zhou, M. Y.: Triggering
724 effects of large topography and boundary layer turbulence on convection over the
725 Tibetan Plateau, *Atmos. Chem. Phys.*, 23, 3299-3309, <https://doi.org/10.5194/acp-23-3299-2023>, 2023.
- 727 Xu, X. D., Zhou, M. Y., Chen, J. Y., Bian, L. G., Zhang, G. Z., Liu, H. Z., Li, S. M., Zhang,
728 H. S., Zhao, Y. J., Suolong, D.J., and Wang, J. Z.: A comprehensive physical pattern
729 of land-air dynamic and thermal structure on the Qinghai-Xizang Plateau, *Sci. China
730 Ser. D-Earth Sci.*, 45, 577-594, <https://doi.org/10.1360/02yd9060>, 2002.
- 731 Yang, B., Qian, Y., Berg, L. K., Ma, P. L., Wharton, S., Bulaevskaya, V., Yan, H. P., Hou,
732 Z. S., and Shaw, W. J.: Sensitivity of Turbine-Height Wind Speeds to Parameters in
733 Planetary Boundary-Layer and Surface-Layer Schemes in the Weather Research and
734 Forecasting Model, *Bound.-Layer Meteor.*, 162, 117-142,
735 <https://doi.org/10.1007/s10546-016-0185-2>, 2017.
- 736 Yang, R. F., Guo, J. P., Deng, W. L., Li, N., Fan, J. H., Meng, D. L., Liu, Z., Sun, Y. P.,
737 Zhang, G. L., and Liu, L. H.: Investigation of turbulent dissipation rate profiles from
738 two radar wind profilers at plateau and plain stations in the north China plain, *Remote
739 Sens.*, 15, 14, <https://doi.org/10.3390/rs15164103>, 2023.
- 740 Zhang, L., Zhang, H. S., Li, Q. H., Wei, W., Cai, X. H., Song, Y., Mamtimin, A., Wang,
741 M. Z., Yang, F., Wang, Y., and Zhou, C. L.: Turbulent mechanisms for the deep
742 convective boundary layer in the Taklimakan desert, *Geophys. Res. Lett.*, 49, 9,
743 <https://doi.org/10.1029/2022gl099447>, 2022.
- 744 Zhang, Y., Guo, J. P., Yang, Y. J., Wang, Y., and Yim, S. H. L.: Vertical wind shear
745 modulates particulate matter pollution: a perspective from radar wind profiler
746 observations in Beijing, China, *Remote Sens.*, 12, 17,
747 <https://doi.org/10.3390/rs12030546>, 2020.



748 Zhao, P., Li, Y. Q., Guo, X. L., Xu, X. D., Liu, Y. M., Tang, S. H., Xiao, W. M., Shi, C.
749 X., Ma, Y. M., Yu, X., Liu, H. Z., Jia, L., Chen, Y., Liu, Y. J., Li, J., Luo, D. B., Cao,
750 Y. C., Zheng, X. D., Chen, J. M., Xiao, A., Yuan, F., Chen, D. H., Pang, Y., Hu, Z.
751 Q., Zhang, S. J., Dong, L. X., Hu, J. Y., Han, S., and Zhou, X. J.: The Tibetan Plateau
752 surface-atmosphere coupling system and its weather and climate effects: the third
753 Tibetan Plateau atmospheric science experiment, *J. Meteorol. Res.*, 33, 375-399,
754 <https://doi.org/10.1007/s13351-019-8602-3>, 2019.
755



756 **Table list**

757 **Table 1.** Summary of the geographical conditions and land surface of the six radar wind
758 profiler (RWP) sites over the Tibet Plateau (TP).

RWP site	Latitude (°E)	Longitude (°N)	Elevation (m)	Land cover types
Minfeng	82.69	37.07	1408.9	Bare land
Jiuquan	98.49	39.77	1477.2	Bare land
Dingri	87.07	28.39	4326.0	Grassland
Ganzi	100.00	31.62	3353.0	Bare land, grassland
Hongyuan	102.55	32.79	3465.0	Bare land, grassland
Lijiang	100.22	26.85	2382.4	Bare land, grass land

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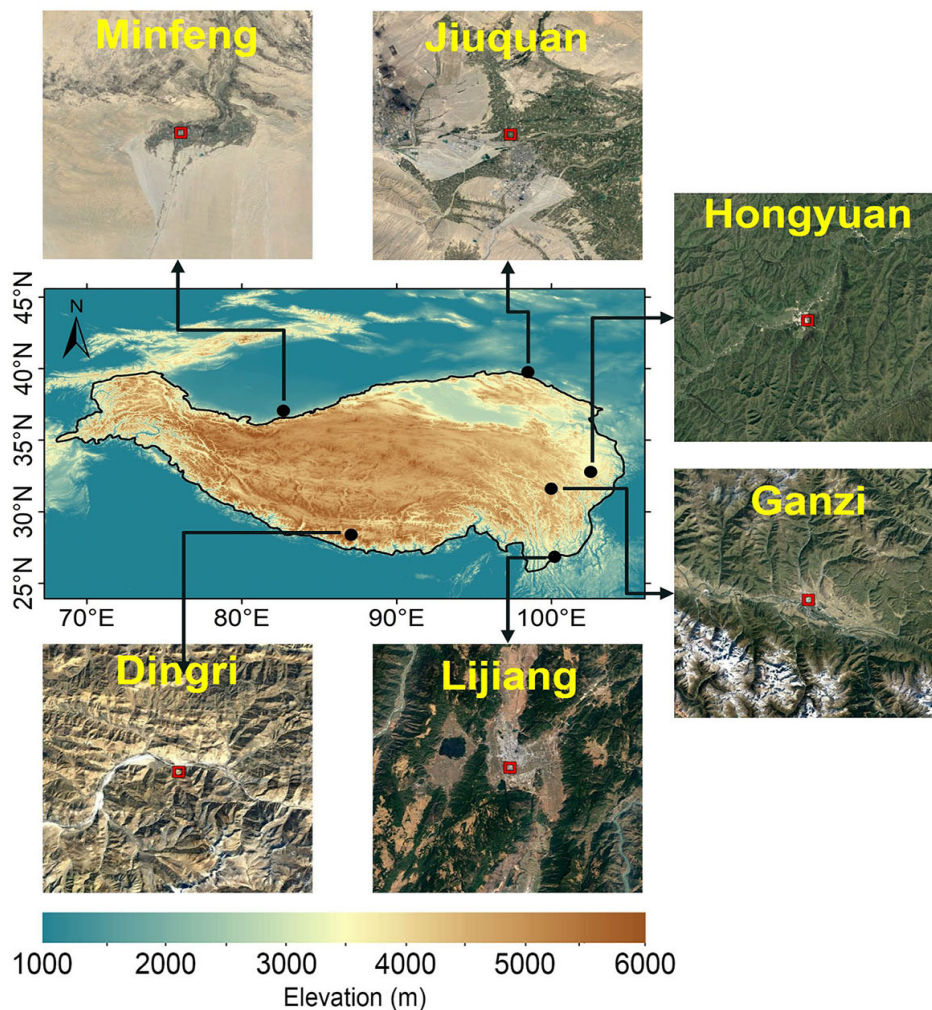
760 **Table 2.** Summary of the correlation of $\text{Log}_{10}\epsilon$ at different altitude ranges under all-,
 761 clear- and cloudy-sky conditions with $T_s - T_a$ and vertical wind shear (VWS) for all six
 762 RWP sites. The superscript * for R indicates that the regression slope is statistically
 763 significant at $p < 0.01$.

Conditions	$\text{Log}_{10}\epsilon$ VS $T_s - T_a$	$\text{Log}_{10}\epsilon$ VS VWS
all-sky, $0.2 < z/z_i < 2.0$	$y = 0.010x - 4.05$, $R = 0.21^*$	$y = 13.6x - 4.19$, $R = 0.29^*$
all-sky, $0.2 < z/z_i < 1.0$	$y = 0.018x - 3.70$, $R = 0.29^*$	$y = 13.2x - 3.77$, $R = 0.20^*$
all-sky, $1.0 < z/z_i < 2.0$	$y = -0.005x - 4.20$, $R = -0.09^*$	$y = 17.6x - 4.57$, $R = 0.36^*$
clear-sky, $0.2 < z/z_i < 2.0$	$y = 0.010x - 4.03$, $R = 0.23^*$	$y = 10.7x - 4.13$, $R = 0.26^*$
clear-sky, $0.2 < z/z_i < 1.0$	$y = 0.018x - 3.67$, $R = 0.30^*$	$y = 11.1x - 3.70$, $R = 0.17^*$
clear-sky, $1.0 < z/z_i < 2.0$	$y = -0.006x - 4.16$, $R = -0.12^*$	$y = 13.8x - 4.52$, $R = 0.34^*$
cloudy-sky, $0.2 < z/z_i < 2.0$	$y = 0.009x - 4.06$, $R = 0.17^*$	$y = 18.5x - 4.29$, $R = 0.33^*$
cloudy-sky, $0.2 < z/z_i < 1.0$	$y = 0.018x - 3.72$, $R = 0.26^*$	$y = 15.5x - 3.84$, $R = 0.23^*$
cloudy-sky, $1.0 < z/z_i < 2.0$	$y = -0.004x - 4.23$, $R = -0.08^*$	$y = 26.2x - 4.67$, $R = 0.42^*$

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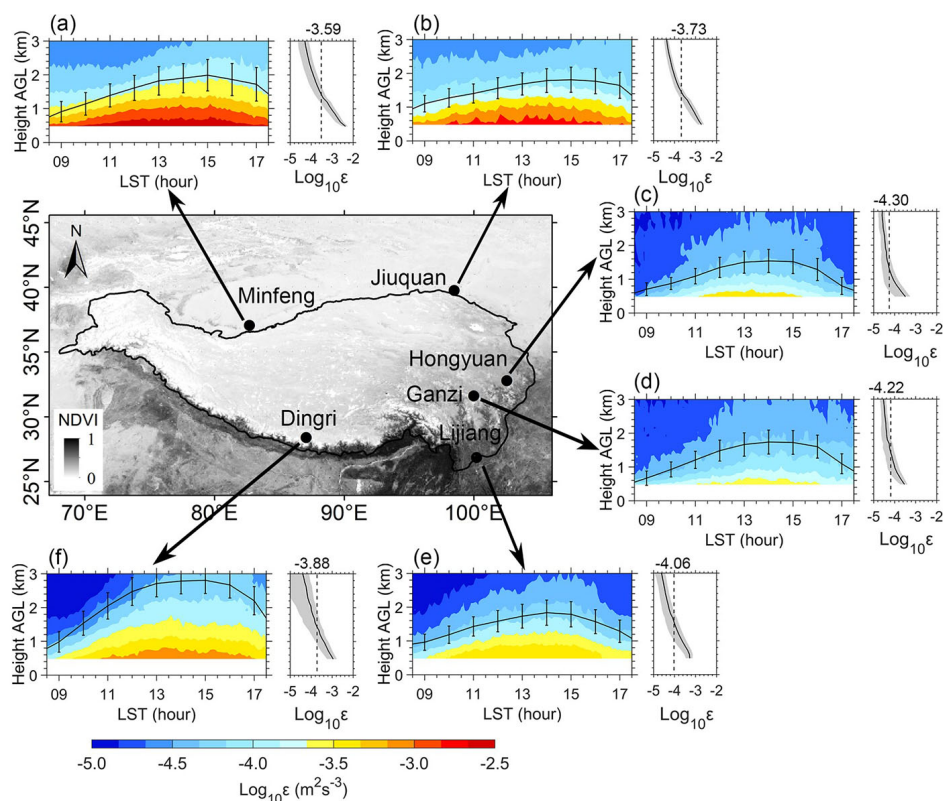
765 **Figures**



766

767 **Figure 1.** Spatial distribution of radar wind profiler (RWP) network comprised of six sites
768 (in black solid circles) on the Tibetan Plateau (TP). The inset map surrounding the main
769 frame denotes the RGB satellite image from © Google Earth that is centered at each RWP
770 site.

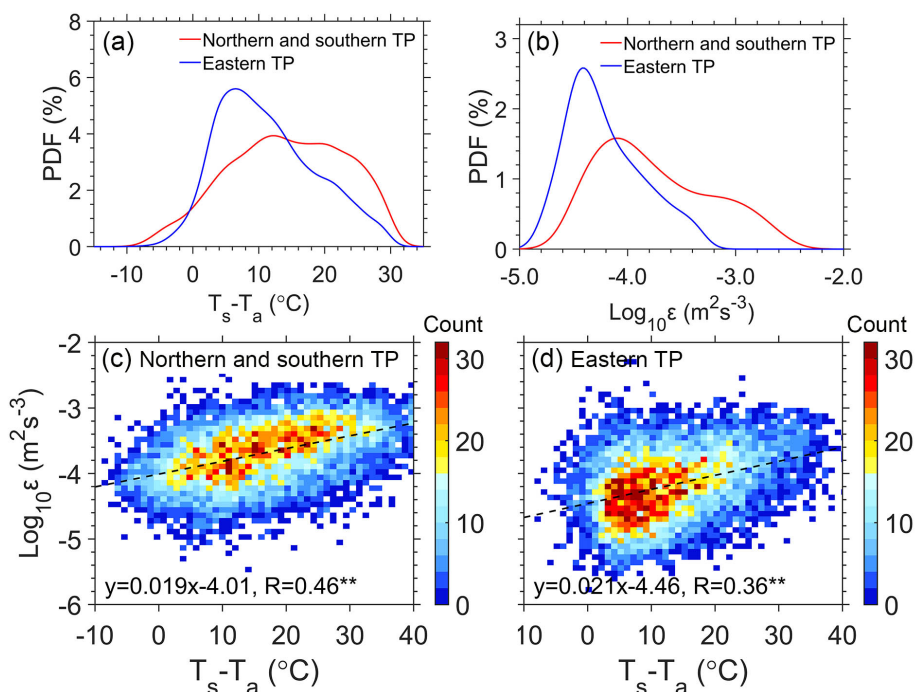
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773 **Figure 2.** Spatial distribution of the diurnal evolution of the vertical profile of logarithmic
 774 turbulence dissipation rate ($Log_{10}\epsilon$ in color shading, unit: $m^2 s^{-3}$) at 120 m vertical
 775 resolution and 6 min intervals, and hourly mean planetary boundary layer height (z_i , black
 776 line, unit: km) during daytime under all-sky conditions from 0900 to 1700 LST for the
 777 period September 2022 to October 2023 as retrieved from the profiling measurements at
 778 six RWP sites over the TP. The vertical bars indicate the 0.5 standard deviations for z_i .
 779 Also shown on the right-hand side panel are temporally averaged vertical profile of ϵ (black
 780 line) and its corresponding one standard deviation (gray shading).

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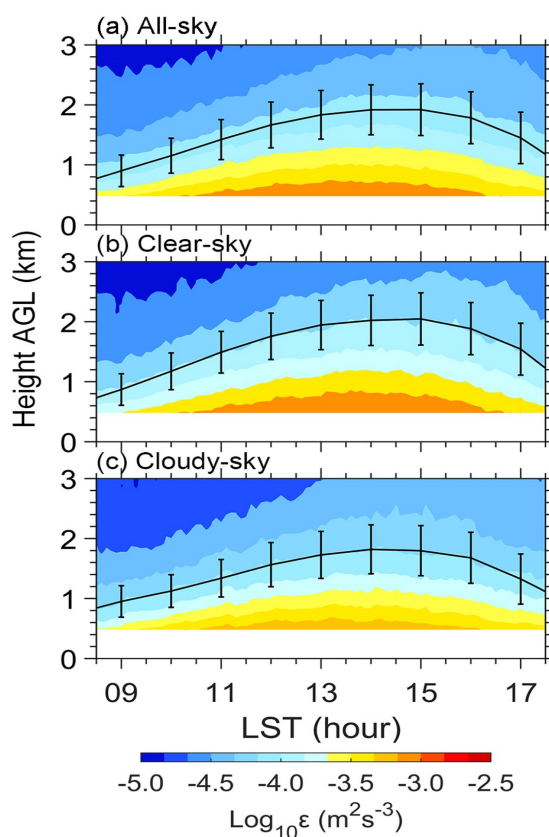
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783 **Figure 3.** (a) PDF of surface-air temperature difference ($T_s - T_a$) for the northern and
 784 southern TP (red line) and eastern TP (blue line), (b) same as (a), but for PDF of $\text{Log}_{10}\epsilon$
 785 estimated from the measurements of radar wind profilers (RWPs) at the height below 0.5
 786 km and above 5 km, (c) scatter plots of $\text{Log}_{10}\epsilon$ as a function of $T_s - T_a$ in the northern
 787 and southern TP, (d) same as (c), but for the eastern TP during daytime under all-sky
 788 conditions from 0900 to 1700 local standard time (LST) for the period September 2022 to
 789 October 2023. The superscript ** for R indicates that the regression slope is statistically
 790 significant at $p < 0.01$ level.

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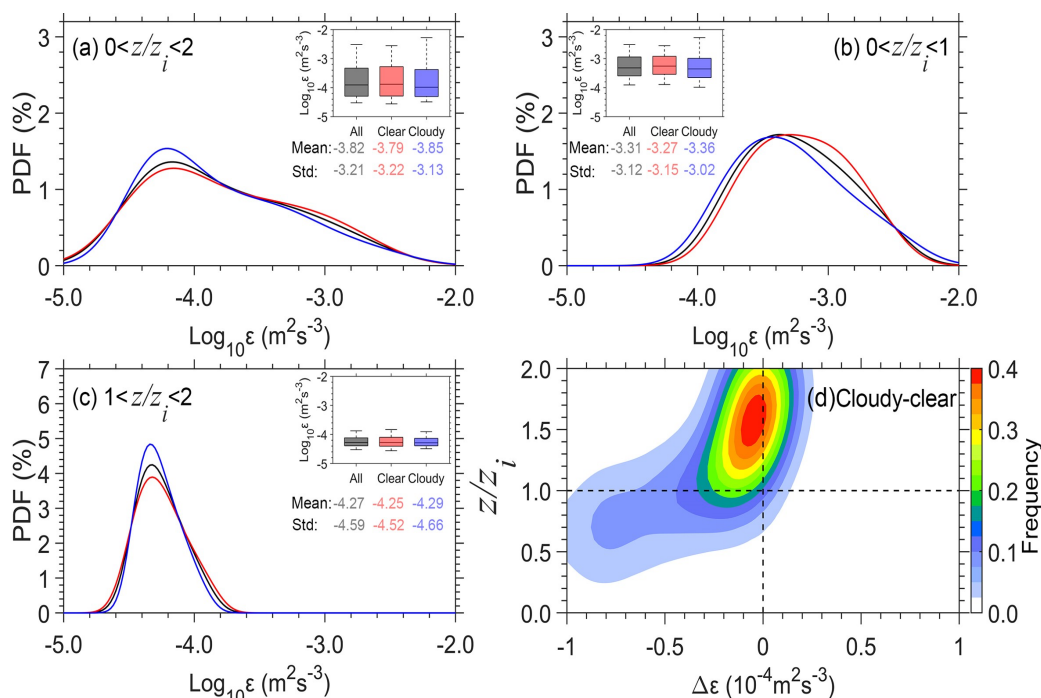


794

795 **Figure 4.** Diurnal evolution of the vertical profile of $\text{Log}_{10}\epsilon$ (color shading, unit: m^2s^{-3})
796 and z_i (solid line, unit: km) averaged over the six RWP sites over the TP during daytime
797 from 0900 to 1700 LST for the period September 2022 to October 2023 for (a) all-sky
798 conditions, (b) clear-sky conditions and (c) cloudy-sky conditions. The vertical bars
799 indicate the 0.5 standard deviations.



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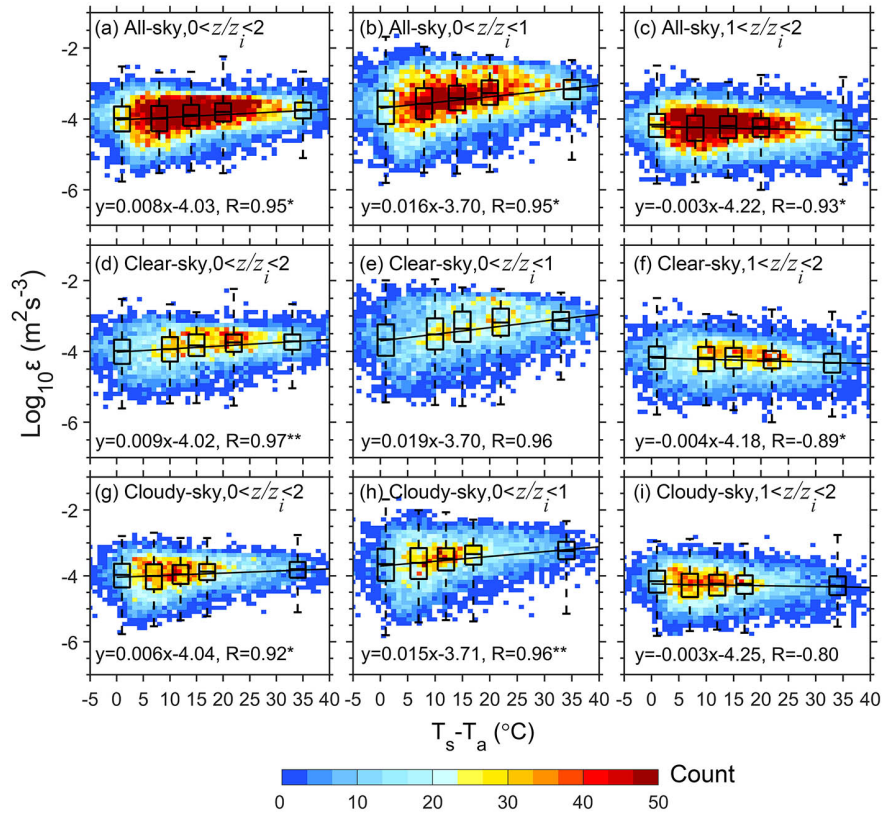
801

802 **Figure 5.** PDF of daytime $\text{Log}_{10} \varepsilon$ (a) in the whole lower troposphere ($0.2 < z/z_i < 2.0$), (b)
 803 in the PBL ($0.2 < z/z_i < 1.0$) and (c) above the PBL ($1.0 < z/z_i < 2.0$) over the TP under all-
 804 sky (black), clear-sky (red) and cloudy-sky (blue) conditions, respectively. (d) Normalized
 805 contoured frequency by altitude diagram (NCFAD) for the difference of ε between cloudy-
 806 sky and clear-sky conditions ($\Delta \varepsilon$) over the TP. Note that z_i denotes the depth of the PBL,
 807 the height (z) and turbulence dissipation rate (ε) is normalized by z_i in order to give a
 808 nondimensional vertical coordinate in the form of z/z_i .

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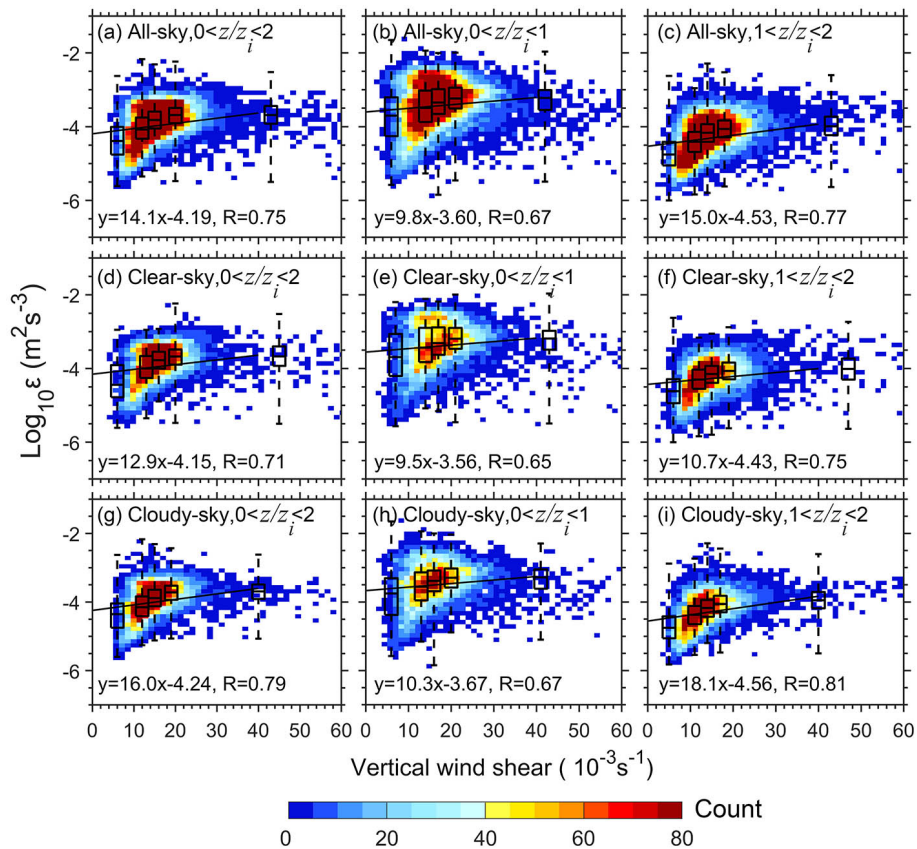
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813 **Figure 6.** Scatter plots (blue dots) of $\text{Log}_{10}\epsilon$ estimated from the measurements of RWPs
 814 in the whole lower troposphere ($0.2 < z/z_i < 2.0$, a, d, g), in the PBL ($0.2 < z/z_i < 1.0$, b, e, h)
 815 and above the PBL ($1.0 < z/z_i < 2.0$, c, f, i) over the TP as a function of $T_s - T_a$ under all-
 816 sky (a-c), clear-sky (d-f) and cloudy-sky conditions (g-i), respectively. Also overlaid are
 817 their corresponding box and whisker plots and regression linear equations and correlation
 818 coefficients in each panel, where all $T_s - T_a$ samples are divided into five bins, each of
 819 which has the same sample size. Note that the median is shown as a line whereas the outer
 820 boundaries of the boxes represent 25 and 75 quartiles and the dashed lines present
 821 interquartile range (IQR). The superscripts * and ** for R indicate that the regression slopes
 822 are statistically significant at $p < 0.05$ and $p < 0.01$ levels, respectively.

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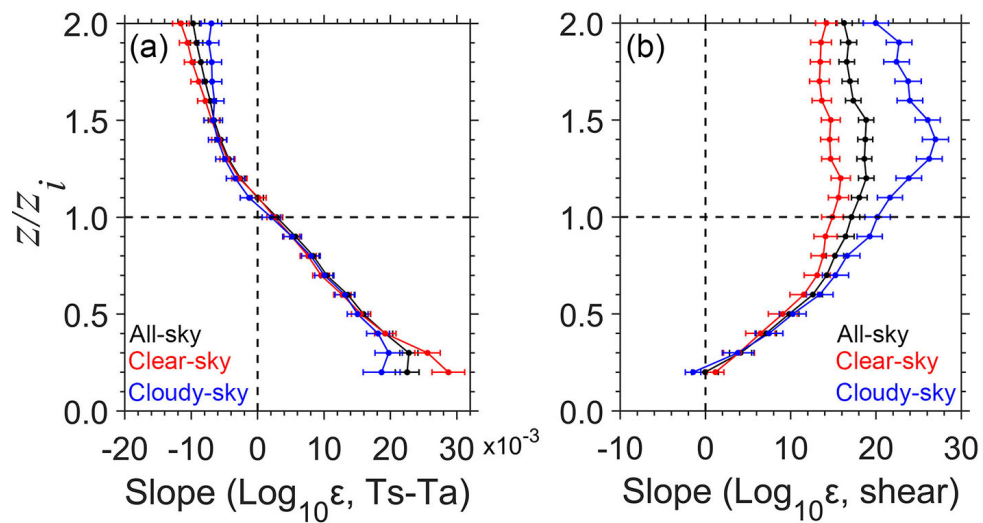


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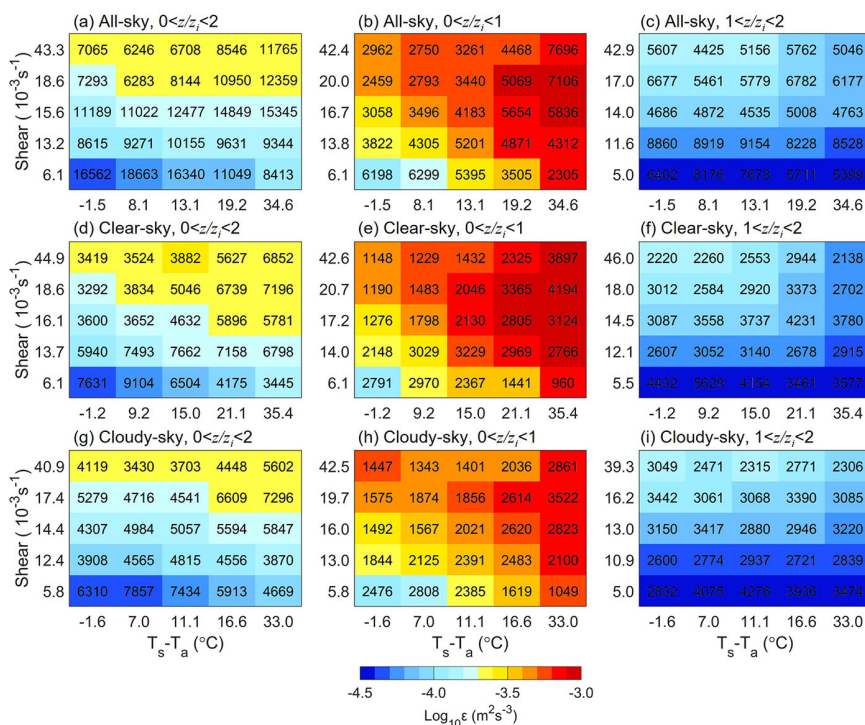
Figure 7. Same as Figure 5, but for $\text{Log}_{10}\epsilon$ as a function of vertical wind shear (VWS).

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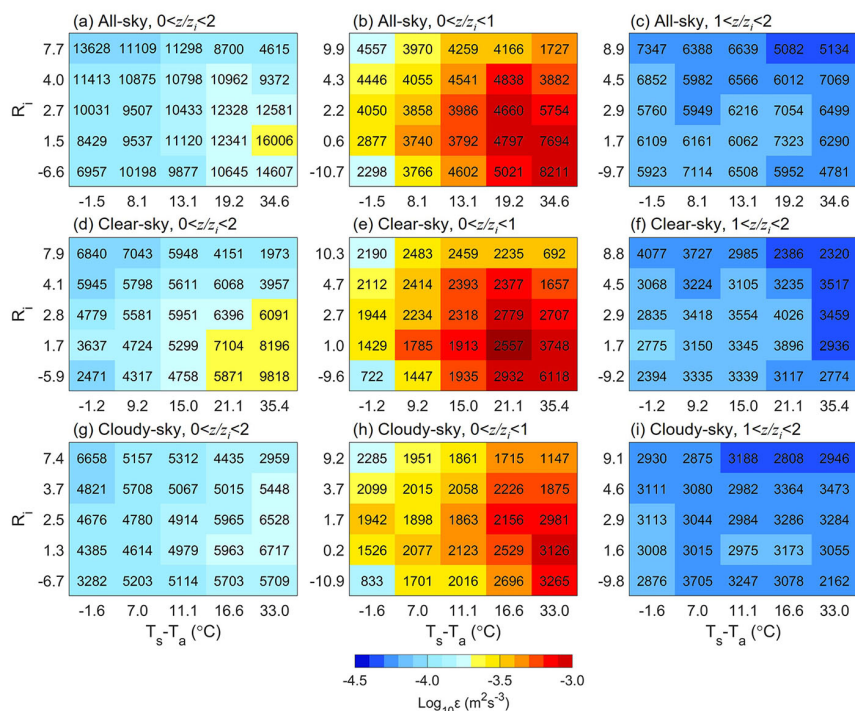
827

828 **Figure 8.** The vertical profiles of least squares regression slope between $\text{Log}_{10} \epsilon$ and
829 $T_s - T_a$ (a) and vertical wind shear (b) over the TP under all-sky (black), clear-sky (red)
830 and cloudy-sky (blue) conditions, respectively.



831

832 **Figure 9.** Joint dependence of $Log_{10}\epsilon$ (color shading) on the vertical wind shear and $T_s -$
 833 T_a within and above the PBL (a, d, g), in the PBL (b, e, h) and above the PBL (c, f, i) over
 834 the TP under all-sky (a-c), clear-sky (d-f) and cloudy-sky (g-i) conditions, respectively.
 835 The number given in each panel is the total number of samples used.



836

837 **Figure 10.** Joint dependence of $\text{Log}_{10}\epsilon$ (color shading) on the gradient Richardson
 838 number (Ri) and $T_s - T_a$ in and above the PBL (a, d, g), in the PBL (b, e, h) and above
 839 the PBL (c, f, i) over the TP under all-sky (a-c), clear-sky (d-f) and cloudy-sky (g-i)
 840 conditions, respectively. The number given in each panel is the total number of samples
 841 used.