

1
2 **Elucidating the boundary layer turbulence dissipation**
3 **rate using high-resolution measurements from a radar**
4 **wind profiler network over the Tibetan Plateau**
5
6

7 Deli Meng^{1,2}, Jianping Guo^{1,3*} Xiaoran Guo^{1*}, Yinjun Wang¹, Ning Li¹, Yuping Sun¹,
8 Zhen Zhang¹, Na Tang¹, Haoran Li¹, Fan Zhang¹, Bing Tong³, Hui Xu¹, Tianmeng Chen¹

9
10
11 *¹State Key Laboratory of Severe Weather, Chinese Academy of Meteorological Sciences,*
12 *Beijing 100081, China*

13 *²High Impact Weather Key Laboratory of CMA, Changsha 410073, China*

14 *³Fujian Key Laboratory of Severe Weather, Fujian Institute of Meteorological Sciences,*
15 *Fuzhou 350028, China*

16 *⁴State Key Laboratory of Urban and Regional Ecology, Research Center for Eco-*
17 *Environmental Sciences, Chinese Academy of Sciences, Beijing 100085, China*

18
19
20
21
22
23
24 *Correspondence to:* Dr/ Prof. Jianping Guo (Email: jpguocams@gmail.com)

25 Dr. Xiaoran Guo (Email: guoxiaoran2018@hotmail.com)

Abstract

The planetary boundary layer (PBL) over the Tibetan Plateau (TP) exerts a significant influence on regional and global climate, while its vertical structures of turbulence and evolution features remain poorly understood, largely due to the scarcity of observation. This study examines the vertical profile and daytime variation of turbulence dissipation rate (ε) in the PBL and free troposphere over the TP using the high-resolution (6 min and 120 m) measurements from the radar wind profiler (RWP) network, combined with the hourly data from the ERA5 reanalysis during the period from September 1, 2022 to October 31, 2023. Observational analyses show that the magnitude of ε below 3 km under all-sky conditions exhibits a large spatial discrepancy over the six RWP stations over the TP. Particularly, the values of ε at Minfeng and Jiuquan over the northern TP and Dingri over the southern TP are roughly an order of magnitude greater than those at Lijiang, Ganzi and Hongyuan over the eastern TP. This could be partially attributed to the difference in land cover across the six RWP stations. In terms of the diurnal variation, ε rapidly intensifies from 0900 local standard time (LST) to 1400 LST, and then gradually levels off in the late afternoon. Under clear-sky conditions, both ε and planetary boundary layer height (z_i) are greater, compared with cloudy-sky conditions, which could be due to the cooling effect by cloud that reduces the solar irradiation reaching the surface. In the lower PBL ($0.3 \leq z/z_i \leq 0.5$), where z is the height above ground level, the dominant influential factor for the development of turbulence is the surface-air temperature difference ($T_s - T_a$). By comparison, in the upper PBL ($0.6 \leq z/z_i \leq 1.0$), both the $T_s - T_a$ and vertical wind shear (VWS) affect the development of turbulence. Above the PBL ($1.0 < z/z_i \leq 2.0$), the shear production resulting from VWS dominates the variation of turbulence. Under cloudy-sky conditions, the reduced $T_s - T_a$ and weakened surface sensible heat flux tend to inhibit the turbulent motion in the PBL. On the other hand, the strong VWS induced by clouds enhances the turbulence above the PBL. The findings obtained here underscore the importance of RWP network in revealing the fine-scale structures of the PBL over the TP and gaining new insight into the PBL evolution.

56 1. Introduction

57 Turbulence ranks among the most intricate phenomena within the atmosphere, ensuring
58 that the planetary boundary layer (PBL) remains thoroughly mixed during daylight hours
59 (Li et al., 2023). As a result, the structure of the PBL is, to a considerable extent, governed
60 by the evolution of turbulence (Teixeira et al., 2021). Turbulence dissipation rate (ε)
61 reflects the amount of turbulent kinetic energy (TKE) that is converted into heat at the
62 Kolmogorov scale and is a measure of the turbulence intensity (McCaffrey et al., 2017;
63 Muñoz-Esparza et al., 2018). Proper parameterizations of the turbulence dissipation term
64 with the aid of observations have great impacts on the forecast skill of weather and climate
65 models, as ε strongly affects vertical turbulent mixing through its influence on TKE (Yang
66 et al., 2017). Accurate estimation of ε is crucial for understanding the structure of
67 turbulence in the PBL. To date, a variety of instruments have been used to observe or
68 retrieve the vertical profiles of ε , including sodar, radar wind profiler (RWP), radiosonde,
69 Doppler wind lidar (DWL) and ultrasonic anemometer (Jacoby-Koaly et al., 2002; Dodson
70 and Griswold, 2021; Lv et al., 2021; Kotthaus et al., 2023). Compared with the DWL, the
71 RWP exhibits better capability in capturing the turbulence structures in the cloudy sky.
72 Furthermore, it is hard for radiosondes and ultrasonic anemometers to get the temporal
73 continuous measurements of atmospheric turbulence, due to the high costs.

74 The Tibetan Plateau (TP), with an averaged elevation greater than 4,000m above sea
75 level (ASL) and an area of approximately 2.5 million km², is towering into the lower and
76 middle troposphere (Huang et al., 2023). By receiving a greater amount of solar shortwave
77 radiation, the surface layer of the TP can transfer more heat through the PBL to the free
78 atmosphere (Wang et al., 2015; Ma et al., 2023). The PBL over the TP exhibits strong
79 convective thermals of warm air and upward motions due to the lower air density and
80 buoyancy effect, which results in significant turbulence motions and turbulence-convection
81 interactions with "popcorn" cloud structures (Xu et al., 2002; Xu et al., 2023).
82 Understanding the statistical behavior of ε is key to revealing the vertical structure and
83 evolution of PBL turbulence, which could improve the parameterization of PBL processes
84 over the TP (Wang et al., 2015; Xu et al., 2019; Zhao et al., 2019; Ma et al., 2023). However,

85 due to the limited observations of turbulence profiles, the daytime variation characteristics
86 of ε over the TP and its main influencing mechanisms remain poorly understood.

87 A vast range of previous studies have attempted to figure out the mechanisms behind the
88 turbulence, but most of them are based on radiosonde measurements or model simulation
89 or reanalysis data (e.g., Banerjee et al., 2018; Che and Zhao, 2021; Wang et al., 2023a). A
90 myriad of driving mechanisms is proposed to account for the PBL development over the
91 TP, such as surface thermal and dynamic forcing, atmosphere stability, among others (Chen
92 et al., 2016; Lai et al., 2021; Wang and Zhang, 2022; Chechin et al., 2023; Wang et al.,
93 2023a). It has been demonstrated that the buoyancy term contribution on the southern slope
94 of the TP is significantly larger than that on the southeastern edge of the TP (Wang et al.,
95 2015). A larger surface-air temperature difference ($T_s - T_a$) and sensible heat flux
96 promotes the rapid growth of deep PBL in the western and southern TP (Chen et al., 2013,
97 2016; Wang et al., 2016; Li et al., 2017a; Zhang et al., 2022).

98 Except for the above-mentioned thermal and dynamic effects, cloud radiative effect is
99 found to be another significant factor that can dramatically modulate the evolution of
100 daytime PBL turbulence (Bodenschatz et al., 2010; Davis et al., 2020). For instance, cloud
101 radiative forcing accounts for the rapid morning transition from stable to unstable PBL,
102 thereby notably affecting the diurnal variation of the PBL (Su et al., 2023). Notably,
103 longwave radiative cooling at the top of stratocumulus clouds can enhance turbulent
104 diffusion within the stratocumulus topped PBL (Sun et al., 2016). A recent observational
105 study suggests that cloud radiative cooling contributed about 32% to turbulent mixing even
106 near the surface (Huang et al., 2020). In other words, cloud radiative processes, including
107 entrainment and radiative cooling, can affect the TKE in the atmosphere (Nicholls et al.,
108 1986; Sedlar et al., 2022; Chechin et al., 2023).

109 The TP is characterized by a high frequency of cumulus clouds which is about five times
110 the regional mean over the other areas of China (Wang et al., 2015), and the occurrence
111 frequency of clouds over the TP shows large diurnal and spatial variability, with the
112 maxima in the afternoon in the eastern TP (Wan et al., 2022). The clouds have been found
113 to significantly suppress the development of summer PBL in the early afternoon across
114 China using fine-resolution radiosonde observations (Guo et al., 2019). Under continuous

115 cloudy-sky conditions, the convective PBL develops slowly due to the smaller surface
116 sensible heat compared to clear-sky conditions (Wang and Zhang, 2022). The turbulence
117 motion in the PBL and its dynamic structure contribute to the formation and development
118 of the popcorn-like convective clouds (Xu et al., 2002; Wang et al., 2020). Compared with
119 eastern China, the more occurrence of low cloud in the afternoon over the TP is found to
120 facilitate the PBL development, mainly owing to the lower atmospheric density (Wang et
121 al., 2020).

122 However, the differences of turbulence vertical structures between clear-sky and cloudy-
123 sky conditions are rarely explored, and the possible mechanism influencing the cloud
124 topped PBL turbulence evolution remains unclear. To the best of our knowledge, most of
125 the above-mentioned studies over the TP lack high-temporal resolution turbulence profile
126 observations. Coincidentally, the RWP network in China provides us a valuable opportunity
127 to characterize the PBL turbulence structure over the TP (Guo et al., 2021a). Therefore, the
128 main objective of this study is to resolve the above issues over the TP, by using
129 observations from the RWP network together with other ground-based meteorological
130 measurements and the ERA5 data. We also analyze the joint effect of thermodynamic and
131 dynamic on ε structure in the daytime (0900–1700 local standard time, LST) PBL through
132 $T_s - T_a$ and VWS.

133 The remainder of this manuscript proceeds as follows, Section 2 describes the data and
134 methods used in this study. In Section 3, we analyze the spatio-temporal characteristics and
135 daytime pattern of ε over the TP and investigate the possible thermodynamic and dynamic
136 effect on PBL turbulence under clear-sky and cloudy-sky conditions. The summary and
137 conclusions are given in section 4.

138

139 **2. Data and methods**

140 ***2.1 The RWP network over the TP***

141 In this study, we use the vertical measurements of RWP data with a vertical resolution
142 of 120 m and a temporal resolution of 6 min from the RWP network over the TP, which

143 contains six operational stations (Minfeng, Jiuquan, Hongyuan, Ganzi, Lijiang and Dingri)
144 operated by the China Meteorological Administration (CMA) during the period from
145 September 1, 2022 to October 31, 2023. The spatial distribution of the RWP network over
146 the TP is shown in Fig. 1, and detailed information for each RWP station, including
147 longitude, latitude, elevation, and land cover type is given in Table 1. Among these six
148 RWP stations, the Dingri station is located in the foothills of the Himalayas with an
149 elevation more than 4,300m ASL, dominated by the land cover of bare and alpine grassland.
150 The Lijiang station is located in the southeastern TP characterized by complex terrain with
151 an elevation of about 2,400m ASL. The Ganzi and Hongyuan stations are situated in the
152 eastern TP, with elevations ranging from 3,300 to 3,500m ASL, and whose underlay is
153 mainly alpine grassland. The Minfen and Jiuquan stations are situated in arid and semi-arid
154 zones to the north of the TP, with elevations ranging from 1,400 to 1,500m, and their
155 dominant underlying land cover is mainly bare land. Therefore, these two stations are well
156 representative of the northern TP.

157 The RWP has the capability to obtain the high-temporal resolution atmospheric
158 turbulence and wind profiles over the TP compared to the radiosonde and reanalysis, which
159 makes it possible to analyze the fine PBL structures. The low and medium detection modes
160 of RWPs can acquire the wind field and turbulence information bellow 5.0 km above
161 ground level (AGL) (McCaffrey et al., 2017; Ruan et al., 2014). The RWP provides the
162 radial observations (marked as RAD subset), including profiles of radial velocity, doppler
163 spectral width, and signal-to-noise ratio (SNR). Also provided by the RWP is the real-time
164 sampling data (marked as ROBS subset), including the profiles of horizontal wind
165 (direction and speed), vertical velocity, and refractive index structure constant (Liu et al.,
166 2020). There exist large uncertainties in the profiling measurements from RWP, thus the
167 quality control for both RAD and ROBS subsets are indispensable before retrieving related
168 dynamic variables over the TP (Liu et al., 2020; Wang et al., 2023). For instance, the
169 profiling measurements highly deviate from the truth below 0.5 km AGL and above 5.0
170 km AGL, which are attributed to the near-surface clutter and significant beam attenuation,
171 respectively (Guo et al., 2023). Thus, here only the RWP measurements at heights from
172 0.5 to 5.0 km are utilized for analysis.

173

174 **2.2 Miscellaneous meteorological data**

175 In this study, the hourly ground-based meteorological variables, including 2m air
176 temperature (T_a), ground surface temperature (T_s), pressure and cloud cover, are derived
177 from the six automatic weather stations over the TP. Also, 1 min rainfall observations from
178 rain gauges are used to minimize the potential influence of rainfall on the profiling
179 measurements from RWP. All these meteorological datasets are subjected to strict data-
180 quality control by the National Meteorological Information Center (NMIC) of the CMA
181 (Wang et al., 2023b). In addition, the hourly temperature data at pressure levels from the
182 ERA5 reanalysis data is also used in this study (Hersbach et al., 2020).

183

184 **2.3 Methods**

185 *2.3.1 Retrieval of turbulence dissipation rate*

186 As a widely used ground-based equipment for detecting atmospheric wind profile (Liu
187 et al., 2020), RWP has the advantage to estimate ε since it could measure Doppler velocity
188 spectrum in the radar volume where the turbulence parcel motion accounts for the spectral
189 width broadening (Jacoby-Koaly et al., 2002; White, 1999). In this study, the spectral width
190 method is applied to retrieve ε from the RAD subset based on the underlying assumption
191 that turbulence is isotropic, and the contributions to the spectral width from turbulent and
192 non-turbulent process are independent of each other (Solanki et al., 2022; White, 1999).

193 The major steps for ε retrieval can be summarized as follows: (1) the spectral width
194 variance consisting of the turbulence and non-turbulence variance is obtained from the
195 spectral width measurements. (2) The non-turbulence broadening variances are
196 decomposed into beam broadening variance due to the finite width of the beam, shear
197 broadening variance generated by the presence of a wind gradient, and broadening variance
198 arising from data processing, among others (Nastrom, 1997). (3) The turbulent broadening
199 variance (σ_t^2) is extracted from the spectral width variance by excluding the above-
200 mentioned non-turbulence broadening variances. (4) ε is estimated from σ_t^2 (White et al.,
201 1999). For more details about the spectral width method, refer to the references (Jacoby-
202 Koaly et al., 2002; McCaffrey et al., 2017; Nastrom, 1997; Solanki et al., 2021).

203 One caveat of the above-mentioned methods used to estimate ε lies in its sensitivity to
204 the uncertainty in measuring horizontal wind speed, and the occurrence of negative value
205 of σ_t^2 , resulting in negative ε (i.e., invalid retrieval and should be discarded), which is
206 previously documented (e.g., Chen et al., 2021; McCaffrey et al., 2017). It's noteworthy
207 that ε estimates derived from the RWP lacks validation against in situ ε measurements
208 from sonic anemometer in the aircraft or tower. This is another factor causing uncertainties
209 that needs to be addressed in the future.

210 *2.3.2 Estimation of planetary boundary layer height*

211 The PBL height (hereafter referred to as z_i) is an important parameter for characterizing
212 fine vertical structure of the PBL, which has important implications for the air mass
213 exchange between the Earth's surface and the atmosphere aloft, thus affecting cloud
214 development and air pollutant dispersion (Dai et al., 2014; Dodson and Griswold, 2021;
215 Guo et al., 2021a; Li et al., 2017b; Wang et al., 2022).

216 Here daytime z_i at each RWP station is retrieved from the original SNR profiles from
217 the RAD subset based on the improved threshold method (ITM), which is originally
218 proposed by Liu et al. (2019). The steps are briefly outlined as follows. First of all, the
219 original SNR profiles are normalized, leading to the profile of normalized SNR (NSNR),
220 which is expected to avoid instrumental inconsistencies. Secondly, the NSNR threshold is
221 set to 0.75 based on the z_i estimated by the radiosonde measurements at the same station.
222 Thirdly, the profile of NSNR is scanned downward from the top to the ground surface.
223 Finally, z_i is determined as the height where the NSNR profile is greater than 0.75 for the
224 first time. For more details for the ITM, refer to Liu et al. (2019).

225 It is not optimal to retrieve z_i directly from the RWP measurements during nighttime,
226 when the turbulence is weaker and SNR is stronger, leading to an overestimation of z_i
227 (Duncan et al., 2022). The accuracy of the SNR data from RWP directly affects the
228 accuracy of z_i . The z_i estimation for the ITM is particularly applicable in the daytime
229 PBL (Bianco et al., 2008; Collaud Coen et al., 2014). The presence of clouds is proved to
230 bring about uncertainty in z_i retrievals from the ITM, due to the challenge in identifying
231 the peak from the NSNR profile (Angel et al., 2024). Notably, a convective cloud is
232 accompanied by strong turbulence, which results in its boundary being misjudged as z_i .

233 *2.3.3 Vertical wind shear*

234 The ROBS subset is used to calculate VWS, which is an important parameter that
235 presents the dynamical effect on the development of PBL (Zhang et al., 2020). VWS is
236 given by

$$\text{VWS} = \left[\left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2 \right]^{1/2} \quad (1)$$

237 where u and v denote zonal and meridional wind component, respectively, z denotes the
238 sample height AGL.

239 *2.3.4 Classification of cloudy- and clear-sky conditions*

240 Using RWP combined with the ground-based cloud cover observations at each station,
241 the effect of clouds on daytime variations of PBL turbulence and z_i over the TP are
242 investigated. Firstly, the 1 min precipitation and 6 min RWP data are time-matched to
243 remove the profile data half an hour before and after the precipitation to obtain non-
244 precipitation data (Wu et al, 2023). Then, all-sky conditions are defined as non-
245 precipitation hours. Finally, the clear-sky (cloudy-sky) conditions are identified as hours
246 with the cloud fraction less (greater) than 30% (80%), respectively (Guo et al., 2016;
247 Solanki et al., 2021).

248 *2.3.5 Calculation of the gradient Richardson number*

249 The evolution of turbulence in the PBL has been previously recognized to be closely
250 associated with atmospheric stability (Chechin et al., 2023; Chen et al, 2013; Lai et al.,
251 2021; Muhsin et al., 2016). Therefore, we take the gradient Richardson number (Ri) as a
252 variable to characterize atmospheric stability and the formation of turbulence over the TP.
253 Following Stull (1988), Ri is formulated as follows:

$$Ri = \frac{g}{\theta_v} \frac{\partial \theta_v / \partial z}{(\partial u / \partial z)^2 + (\partial v / \partial z)^2} \quad (2)$$

254 where θ_v is the virtual potential temperature from ERA5, u and v are the hourly zonal
255 and meridional wind components derived from RWP, respectively, g is the gravitational
256 acceleration, and z represents the sample height AGL.

257

258 **3. Results and discussion**

259 *3.1 Spatio-temporal distributions of daytime PBL turbulence dissipation rate*

260 Both the PBL turbulence dissipation rate and z_i have significant diurnal variations over
261 mountain and urban areas (Adler et al., 2014; Liu et al., 2019; Solanki et al., 2021; Yang
262 et al., 2023). Since the longitude of the six stations over the TP ranges from 82.7°E to
263 102.6°E, it is necessary to use the LST to accurately capture the daytime variations of the
264 PBL and make a comparison between different stations.

265 Figure 2 presents a comprehensive overview of the ε profile at 6 min intervals and
266 hourly averaged z_i in lower troposphere at heights from 0.5 to 3.0 km for six RWP stations
267 over the TP during the period from September 1, 2022 to October 31, 2023. As shown in
268 the right panels of Fig. 2, ε generally decrease with increased height at all six RWP
269 stations. The magnitude of ε and its vertical structures during the daytime at both Minfeng
270 and Jiuquan stations over the northern TP and at Dingri station over the southern TP stand
271 in stark contrast to those RWP stations (i.e., Lijiang, Ganzi and Hongyuan) in the eastern
272 TP. It is apparent that ε exhibits a large spatial discrepancy. In terms of the latitudinal
273 variation, the one-year averaged ε at the RWP stations in the east part of TP is smaller
274 than in the western part of TP. In terms of the meridional variation, ε at the two RWP
275 stations in the northern TP have a significantly larger magnitude than the other four stations.
276 In particularly, the maximum mean value of daytime ε is found at Minfeng and Jiuquan
277 in the northern TP, which reaches values up to $10^{-3.59} \text{ m}^2 \text{ s}^{-3}$ and $10^{-3.73} \text{ m}^2 \text{ s}^{-3}$, respectively.
278 By comparison, the lowest value of ε is found in the eastern TP, with the mean values of
279 $10^{-4.06} \text{ m}^2 \text{ s}^{-3}$, $10^{-4.30} \text{ m}^2 \text{ s}^{-3}$ and $10^{-4.22} \text{ m}^2 \text{ s}^{-3}$ at Lijiang, Hongyuan and Ganzi, respectively.
280 Meanwhile, the mean magnitude of ε at Dingri in the southern TP lies between the
281 magnitude of ε in the northern and eastern TP, which is $10^{-3.88} \text{ m}^2 \text{ s}^{-3}$.

282 Overall, the spatial distribution of the z_i at all six RWP stations is clearly dependent on
283 geographical location (Fig. 2), which resembles that of the ε . The geographic pattern of
284 z_i from RWP agrees well with those from radiosonde measurements (Che and Zhao, 2021)
285 and reanalysis (Slättberg, 2022). Of the six RWP stations, Dingri is located in the northern
286 foothills of the Himalayas with an altitude of over 4300 m, where the bare land type results

287 in a large surface sensible heat flux. This, together with the lowest atmospheric density,
288 leads to the highest daytime mean value of z_i up to 2.10 km (Wang et al., 2015). The land
289 surfaces at the Minfeng and Jiuquan stations in the northern TP are dominated by barren
290 and relatively homogenous terrain, in sharp contrast to the highly vegetated terrain at the
291 Ganzi and Hongyuan stations in the eastern TP (Fig. 1). The sparse vegetation in the
292 northern TP generally comes with large Bowen ratio during the daytime, which tends to
293 produce larger sensible heat flux compared to that in the eastern TP. The increased
294 turbulence intensity in the PBL is generally associated with larger sensible heat flux, which
295 has been reported by previous studies (Wang et al., 2016; Zhang et al., 2022). Therefore,
296 the spatial and temporal variation of daytime ε over the TP are affected by the underlying
297 surface type and air density.

298 Regarding the daytime pattern of turbulence (all six panels with color shading in Fig. 2),
299 the turbulence over the TP shows a pronounced signature of single-peak variability. During
300 the period 0900–1100 LST, the magnitude of ε at all six RWP stations is relatively weak.
301 From 1100 LST onward, with the increase of downward solar shortwave radiation, surface
302 sensible heat flux gradually rises, which leads to acceleration of turbulence mixing
303 processes. Then, ε reaches peak in the early afternoon (1300–1500 LST). Afterwards,
304 during the later afternoon (1500–1700 LST), ε diminishes gradually. Likewise, z_i almost
305 follows the same daytime variation pattern of ε .

306 On the seasonal scale, the turbulence at the six RWP stations is characterized by
307 significant variability, which is shown in Fig. S1. To be more specific, ε reaches the
308 maximum in summer with the highest z_i , while touches the minimum in winter at Minfeng
309 and Jiuquan. At the remaining four stations, the strongest ε is found in spring, as opposed
310 to the weakest ε in autumn.

311 The above-mentioned findings imply that the turbulence intensity at the RWP stations
312 over the northern and western TP is about one order of magnitude greater than that in the
313 eastern TP. To further investigate the possible reasons for this significant difference in ε ,
314 the relationships between $T_s - T_a$ and ε for different regions are presented in Fig. 3. The
315 mean value of $T_s - T_a$ in the northern and southern TP is 14.29°C, which is greater than
316 that of eastern TP with the value of 11.26°C (Fig. 3a). The mean daytime ε for the two

317 regions reaches up to $10^{-3.74} \text{ m}^2 \text{ s}^{-3}$ and $10^{-4.20} \text{ m}^2 \text{ s}^{-3}$, respectively (Fig. 3b). Additionally, ε
318 is significantly and positively correlated with $T_s - T_a$ ($R > 0.35$, $p < 0.005$), which
319 illustrates that the thermal forcing makes an important contribution to turbulence
320 development in the TP (Fig. 3c-d). As shown in Fig. S2, there is a positive correlation
321 between the $T_s - T_a$ and ε , indicating that the thermal effect of the $T_s - T_a$ can promote
322 the development of turbulence in at heights from 0.5 to 3.0 km under all-sky conditions.
323 However, the relationship varies significantly between each RWP station. The slope values
324 of the regression coefficients for the other five RWP stations, except for Hongyuan are all
325 greater than 0.015. The maximum slope values are observed at Lijiang (0.029) and Dingri
326 (0.027) in the southern TP, as compared with the minimum slope of 0.007 at Hongyuan.
327 This suggests that near-surface thermal properties have nothing to do with ε at Hongyuan
328 in the eastern TP.

329 Besides, the potential impact of VWS on ε is also examined, which is shown Fig. S3.
330 Overall, VWS is found to positively correlate with ε at heights from 0.5 to 3.0 km under
331 all-sky conditions, differing by RWP stations. The maximum slope values are observed at
332 Lijiang (79.34) and Hongyuan (68.56), as compared with the minimum slope of 1.15 at
333 Minfeng. Therefore, it can be inferred that atmospheric dynamic effect induced by VWS
334 dominates the variability of ε at regions with the same underlying terrain and land over
335 as Hongyuan.

336 ***3.2 Characteristics of daytime PBL turbulence dissipation rate under clear- and*** 337 ***cloudy-sky conditions***

338 The influence of clouds on the PBL properties has been discussed and analyzed in
339 previous studies (e.g., Guo et al., 2016; Huang et al., 2023; Ma et al., 2023; Schumann et
340 al., 1991; Yu et al., 2004). To reveal the potential impact of clouds on the PBL ε over the
341 TP, the comparison analyses between clear- and cloudy-sky conditions are presented in
342 this section. Figure 4 shows the daytime cycle of mean ε profile and z_i averaged over the
343 six RWP stations under all-, clear- and cloudy-sky conditions. Overall, both the profile of
344 ε and z_i under all-sky conditions over the TP present distinct single-peak variations, and
345 their peaks approximately occur at 1400 LST (Fig. 4a). The daytime averaged ε below 3.0

346 km AGL is $10^{-3.95} \text{ m}^2 \text{ s}^{-3}$, and mean z_i is 1.47 km, respectively. There is a significant
347 positive correlation between ε and z_i during the daytime ($R=0.63$, $p<0.01$).

348 Under clear-sky condition, the daytime mean ε is $10^{-3.88} \text{ m}^2 \text{ s}^{-3}$ (Fig. 4b). During the
349 period 1300–1500 LST, ε ranges from $10^{-3.43}$ to $10^{-2.82} \text{ m}^2 \text{ s}^{-3}$ ($10^{-4.17}$ to $10^{-3.40} \text{ m}^2 \text{ s}^{-3}$) at
350 heights from 0.5 km (1.0 km) to 1.0 km (2.0 km) in lower (upper) PBL. Thus, the well-
351 mixed turbulence maintains the development of PBL in the early afternoon. By comparison,
352 under cloudy-sky condition (Fig. 4c), the daytime mean value of z_i can reach up to 1.4 km,
353 which is 0.12 km lower than that of clear-sky conditions. This means that the clouds would
354 suppress the development of the PBL turbulence in the early afternoon which has been
355 observed by the radiosonde observations described in Guo et al. (2016).

356 It is well known that there exists diurnal variation in PBL. To better reveal the
357 mechanism how a myriad of geophysical parameters affect turbulence, the height-revolved
358 ε retrievals are further normalized by the average PBL height. As noted above, the valid
359 minimum altitude of the RWP is 0.5 km at 120 m vertical resolution, and the maximum z_i
360 is approximately 2.0 km (Figs. 2&4). z is normalized by z_i to provide a nondimensional
361 vertical coordinate for ε . It follows that z/z_i is great than 0.25, and the range of z/z_i is
362 set from 0.3 to 2.0 for the following analyses.

363 The probability density distribution (PDF) of ε in the PBL ($0.3 \leq z/z_i \leq 1.0$) and
364 above the PBL ($1.0 < z/z_i \leq 2.0$) under all-, clear- and cloudy-conditions are given in Fig. 5.
365 Overall, the mean ε are $10^{-3.82}$, $10^{-3.79}$ and $10^{-3.85} \text{ m}^2 \text{ s}^{-3}$ at the height range of $0.3 \leq z/z_i \leq 2.0$
366 under all-, clear- and cloudy-sky conditions, respectively (Fig. 5a). Within the PBL (Fig.
367 5b), the mean ε under clear-sky conditions ($10^{-3.27} \text{ m}^2 \text{ s}^{-3}$) is greater than that of under
368 cloudy-sky conditions ($10^{-3.36} \text{ m}^2 \text{ s}^{-3}$), and the standard deviation of ε under clear-sky
369 conditions is slightly greater than that under cloudy-sky conditions. This illustrates that
370 clouds can significantly inhibit the turbulence intensity in the PBL, with the value of $\Delta\varepsilon$
371 between clear- and cloudy-sky conditions is $-10^{-4.0} \text{ m}^2 \text{ s}^{-3}$. However, above the PBL (Fig.
372 5c), ε presents normal distribution characteristics, and there is no significant difference
373 between the mean ε under clear- and cloudy-sky conditions.

374 To examine the overall impact of clouds on the vertical structure of turbulence within
375 and above the PBL, Figure 5d shows the normalized contoured frequency by altitude
376 diagram (NCFAD) of the $\Delta\varepsilon$ for normalized (z/z_i) profiles of ε between cloudy-sky and
377 clear-sky conditions. Within the PBL, $\Delta\varepsilon$ is negative, and $|\Delta\varepsilon|$ generally decrease with
378 increased z/z_i , where $\Delta\varepsilon$ is $-10^{-4.3} \text{ m}^2 \text{ s}^{-3}$ at $z/z_i=0.5$, and $-10^{-5.0} \text{ m}^2 \text{ s}^{-3}$ at $z/z_i=1.0$,
379 respectively. This suggests that clouds may weaken turbulence within the PBL (Fig. 4b-c),
380 especially in the lower PBL ($z=820\text{m}$, $z/z_i<0.5$). Figure S4 further shows the distinct
381 spatial variability of cloud effect on ε across the six RWP stations. Particularly, the
382 turbulence is weakened by clouds within the PBL at Minfeng and Jiuquan in the northern
383 TP, as opposed to the enhanced ε within the PBL at Ganzi and Lijiang. This suggests that
384 the cloud impact on ε is much complicated than expected. One of the reasons could be
385 concerned with the cloud life stage, which is not dealt with in this present study. On top of
386 the life stage, the cloud impact on ε , in combination with $T_s - T_a$ and VWS , exhibits a
387 distinct altitude dependence, differing by RWP stations (Fig. S5).

388

389 ***3.3. Potential factors Influencing daytime PBL turbulence dissipation rate***

390 *3.3.1 Surface-air temperature difference*

391 The vertical structure of PBL ε and z_i over the TP show obvious spatial differences in
392 the context of a complex subsurface. The diverse land cover types lead to differences in
393 surface albedo and soil moisture, which in turn lead to distinctions in thermodynamic
394 characteristics such as sensible heat flux (Ma et al., 2023). Buoyant production driven by
395 solar heating from the surface is one of the dominant sources generating turbulence in the
396 convective PBL. The surface sensible heat flux is an important thermodynamic factor that
397 affects the buoyant convective processes (Stull, 1988). Meanwhile, previous studies (e.g.,
398 Wang et al., 2022; Yang et al., 2023) have suggested that $T_s - T_a$ can serve as a good
399 proxy for the sensible heat flux. There are not sensible heat flux measurements at six RWP
400 stations in this study, and thus we directly take $T_s - T_a$ as a proxy thermodynamic variable
401 to analyze its potential connection to variation of PBL turbulence.

402 Figure 6 shows the magnitude of ε varies as a function of $T_s - T_a$ for all six stations,
 403 within ($0.3 \leq z/z_i \leq 1.0$) and above ($1.0 < z/z_i \leq 2.0$) the PBL, under all, clear- and cloudy-
 404 sky conditions, respectively. $T_s - T_a$ are first classified into five bins, which are then
 405 statistically analyzed against the corresponding ε averaged for z/z_i values between 0.3
 406 and 2.0 to obtain regression equations incorporating slopes. Further, Table 2 shows the
 407 scatter plots between $\text{Log}_{10}\varepsilon$ (Figure 6) and $T_s - T_a$ (and VWS, Figure 7) at different
 408 altitude ranges under all-, clear- and cloudy-sky conditions. $\text{Log}_{10}\varepsilon$ is found to be linearly
 409 correlated with $T_s - T_a$ (and VWS) ($p < 0.05$). The surface sensible heat flux generally
 410 increases with increased $T_s - T_a$, thus the increased $T_s - T_a$ intensifies the turbulence in
 411 PBL ($0.3 \leq z/z_i \leq 1.0$), which is shown in Fig. 6b, e, h. Within the PBL, ε is also positively
 412 correlated with $T_s - T_a$ whose slope values are larger than those at $0.3 \leq z/z_i \leq 2.0$. As
 413 $T_s - T_a$ rises, the larger surface sensible heat flux would lead to enhanced buoyancy
 414 process and turbulent motion within the PBL. On the other hand, ε above the PBL is
 415 negatively correlated with $T_s - T_a$ (Figs. 6c, f, i). This suggests that $T_s - T_a$ dramatically
 416 affects the development of turbulence within the PBL, whereas it has little effect on the
 417 turbulence above the PBL.

418 Within the PBL, the magnitude of slope (slope=0.019) under clear-sky conditions is
 419 larger than that of under-cloudy conditions (slope = 0.015) as shown in Figs. 6e and 6h.
 420 This implies that $T_s - T_a$ is the governing parameter rather than cloud cover affecting the
 421 PBL turbulence, particularly under the clear-sky conditions. Given that turbulence in the
 422 mixed PBL over the TP is usually driven by convection (Xu et al., 2023), as $T_s - T_a$
 423 decreases when clouds are present, less heat is transferred from the surface to the
 424 atmosphere, reducing the buoyancy flux and leading to weaker turbulence in the PBL,
 425 especially for the lower PBL ($0.3 \leq z/z_i \leq 0.5$), as shown in Figures 4b and 4c. Consequently,
 426 the clouds tend to suppress the development of PBL (Fig. 5a) and reduce z_i .

427 3.2.2 Vertical wind shear

428 Besides $T_s - T_a$, VWS is another crucial dynamic parameter that is related to the
 429 mechanical turbulence within the PBL. Similar to Fig. 6, Figure 7 presents the relationship
 430 between ε and VWS (both normalized by z_i) within and above the PBL under all-, clear-

431 and cloudy-sky conditions, respectively. The near-surface clutter significantly increases
432 the uncertainty of RWP data, which leads to incapability of analyzing the effect of wind
433 shear on ε below 0.5 km AGL ($z/z_i \geq 0.3$) in the following sections.

434 Regardless of within or above the PBL, ε is positively correlated with VWS as shown
435 in Fig. 7a, d, g and Tabel 2, which indicates that larger VWS leads to stronger turbulence.
436 This suggests that the dynamic effect of VWS promotes the development of turbulence.
437 Within the PBL (Figs. 7b, e, h), the slope of ε against VWS are smaller than at $0.3 \leq$
438 $z/z_i \leq 2.0$ with values ranging from 9.5 to 10.3. Above the PBL (Figs. 7c, f, i), the values
439 of the slope are larger with values ranging from 10.7 to 18.1, which demonstrating that the
440 dynamical effects of VWS influence the development of turbulence both within the upper
441 PBL and above the PBL.

442 Under cloudy-sky conditions (Figs. 7h, i), the effect of VWS on turbulence within the
443 upper PBL (Slope=10.3) is weaker than above the PBL (Slope=18.1), significantly.
444 Compared to the clear-sky conditions (Figs. 7e, f), the values of the slopes are larger for
445 that of under cloudy-sky conditions (Figs. 7h, i) both within and above the PBL.
446 Remarkably, above the PBL, the effect of clouds on turbulence is more dramatic, as the
447 slope value under cloudy-sky conditions is nearly twice as large as that of under clear-sky
448 conditions. These results indicate the significant mechanical processes driven by VWS is
449 important in the development of turbulence. A larger VWS in the PBL corresponds to
450 stronger turbulence. Besides, above the PBL, the mechanical process of VWS is enhanced
451 under cloudy-sky conditions.

452

453 *3.2.3. Joint influence of $T_s - T_a$, VWS and atmosphere stability on ε*

454 It was stated that turbulence can be produced by buoyant convective processes (i.e.,
455 thermals of warm air rising) and by mechanical processes (i.e., wind shear). From the
456 previous section, it is known that $T_s - T_a$ and VWS both affect the development of PBL
457 turbulence. Figure 8 gives the slope profiles of ε against $T_s - T_a$ and VWS at normalized
458 heights (z/z_i) under all-, clear- and cloudy-sky conditions, respectively.

459 As inferred from the previous findings, $T_s - T_a$ primarily influences turbulence
460 development within the PBL, irrespective of clear-sky and cloudy-sky conditions (Fig. 6).
461 Figure 8a shows that the slope values within the PBL are predominantly positive, and the
462 slope value decreases rapidly with height, which indicates that the influence of $T_s - T_a$ on
463 PBL turbulence decreases with height. Interestingly, there is a nearly linear variation of the
464 slope from the lower PBL to the top of the PBL. Within the PBL, the slope is positive,
465 above the PBL, the slope becomes negative. This may be due to the linear decrease of heat
466 flux transport and buoyancy term in the convective PBL (Stull, 1988). Therefore, these
467 findings highlight the predominant thermal forcing of $T_s - T_a$ on turbulence development
468 within the lower PBL. Fig 8 clearly shows the influence of cloud cover on the $T_s - T_a$ and
469 the effect of the surface heating on the turbulence in the lower half of the PBL ($0.3 \leq z/z_i \leq 0.5$).
470 While there is little difference for the clear-sky and cloudy-sky conditions when
471 $z/z_i > 0.5$. Hence, under clear-sky conditions, the thermodynamic effect of $T_s - T_a$ is more
472 pronounced within the lower PBL.

473 As shown in Fig. 7, it is evident that VWS influences turbulence development within
474 and above the PBL. Figure 8b shows that when $0.3 \leq z/z_i \leq 2.0$, the slope values are
475 consistently positive, indicating that VWS predominantly affects turbulence development
476 within the mid-, upper- PBL and above the PBL. Moreover, when $0.3 \leq z/z_i \leq 1.2$, the
477 slope values increase with height. However, when $1.4 < z/z_i \leq 2.0$, the slope decreases with
478 height, which suggesting a diminishing influence of VWS. Additionally, within the PBL
479 ($0.3 \leq z/z_i \leq 0.7$), the slope values under clear-sky conditions are close to those under
480 cloudy-sky conditions, while the slope values under cloudy-sky conditions are even greater
481 when $0.7 < z/z_i \leq 2.0$. For instance, when $z/z_i = 1.4$, $\text{Slope}_{\text{Clear-sky}} = 14.6$, while $\text{Slope}_{\text{Cloudy-}}$
482 $\text{sky} = 27.0$, indicating that the latter is 1.8 times larger than the former. These results suggest
483 that clouds are primarily responsible for enhancing mechanical processes from VWS on
484 turbulence within the upper PBL and above the PBL.

485 Furthermore, it can be concluded that, $T_s - T_a$ is the thermodynamic factor influencing
486 turbulence development within the lower PBL ($0.3 \leq z/z_i \leq 0.5$), both $T_s - T_a$ and VWS
487 jointly strengthen turbulence development in the upper PBL ($0.6 \leq z/z_i \leq 1.0$), and VWS

488 emerges as the predominant factor affecting turbulence development above the PBL
489 ($1.0 < z/z_i \leq 2.0$) (Figs. 8a, b).

490 The previous sections have revealed that hours of both high $T_s - T_a$ and strong wind
491 shear would strengthen the turbulence within the PBL. Therefore, it's necessary to analyze
492 the combined influence of thermodynamics and dynamics factors on the development of
493 turbulence. Figure 9 presents the joint distribution of ε with $T_s - T_a$ and VWS within and
494 above the PBL under all-, clear- and cloudy-sky conditions. Within the PBL (Figs. 9b, e,
495 h), higher $T_s - T_a$ and VWS correspond to stronger turbulence (Fig. 8). In contrast, the
496 thermodynamic effect of $T_s - T_a$ on turbulence has diminished and is no longer a
497 dominant factor above the PBL, while the dynamical effect of VWS becomes the dominant
498 factor (Figs. 9c, f, i). Compared to clear-sky conditions, both $T_s - T_a$ and VWS decrease
499 under cloudy-sky conditions (Fig. 9h). This means that the weakening of both
500 thermodynamic and dynamic effects leads to a decrease in turbulence, thereby inhibiting
501 the development of turbulence within the PBL. Therefore, under cloudy-sky conditions,
502 although the VWS is reduced, the dynamical effect of VWS on turbulence is strengthened
503 (Figs. 7i and 8b), which in turn strengthens turbulence.

504 Since buoyant and mechanic forcing jointly influence the turbulence within the PBL,
505 and VWS only represents the dynamic driving effect, it cannot accurately portray the effect
506 of thermodynamic and dynamic effects on the PBL turbulence. The gradient Richardson
507 number (Ri), on the other hand, is one of the important parameters characterizing
508 atmospheric stability and can compare the buoyant turbulence production term and the
509 shear production term in the form of a dimensionless ratio.

510 Similar to Fig. 9, the joint distribution of ε with $T_s - T_a$ and Ri within and above the
511 PBL under all-sky, clear-sky and cloudy-sky conditions is given in Fig. 10. As shown in
512 Fig. 10b, e, h, it is evident that the turbulence in the PBL tend to be enhanced for unstable
513 conditions. Furthermore, under clear-sky conditions (Fig. 10e), the maximum number of
514 samples is found when $Ri < 1.0$ and $T_s - T_a > 21.1$ in strongly unstable conditions, which
515 may be caused by the buoyancy forcing driven by the larger $T_s - T_a$. By comparison, the
516 effect of Ri on turbulence is relatively weakened above the PBL (Figs. 10c, f, i).

517

518 4 Summary and concluding remarks

519 This study investigates the characteristics of spatio-temporal distribution of daytime
520 PBL turbulence dissipation rate (ε) based on more than one-year record (September 2022–
521 October 2023) of profiling measurements from a radar wind profilers (RWP) network on
522 the Tibet Plateau (TP). Also analyzed are the evolution of ε in the PBL and the possible
523 influential mechanisms.

524 First of all, ε is firstly retrieved from the vertical wind measurements from RWP using
525 the spectral width method. Afterwards, the spatial pattern of ε is examined. Results shows
526 that the values of ε at the Minfeng and Jiuquan stations in the northern TP, and at Dingri
527 over the southern TP are about one order of magnitude greater than those at the RWP
528 stations of Lijiang, Ganzi and Hongyuan over the eastern TP. Coincidentally, Minfeng and
529 Junquan are dominated by bare or semiarid land, as opposed to the highly vegetation-
530 covered land surface at Lijiang, Ganzi and Hongyuan. This suggests the spatial discrepancy
531 of ε over the TP is highly relevant to the types of underlying land cover.

532 Although ε exhibits a variety of magnitudes among the six RWPs, the daytime pattern
533 and vertical structure of ε are similar. Turbulence reaches the peak in the early afternoon
534 (1300–1500 LST), coinciding with the highest PBL top. Under cloudy-sky conditions, the
535 daytime mean value of ε is $10^{-4.02} \text{ m}^2 \text{ s}^{-3}$, and the daytime mean value of the PBL height
536 (z_i) can reach up to 1.40 km, which is 0.12 km lower than that of clear-sky conditions,
537 indicating that clouds would suppress the development of the PBL turbulence.

538 As far as both the thermodynamic and dynamic forcings are concerned, surface-air
539 temperature difference ($T_s - T_a$) and vertical wind shear (VWS) variables are examined by
540 performing correlation analysis with ε . The slope values of ε against $T_s - T_a$ under clear-
541 sky conditions is larger (slope=0.019) than under-cloudy conditions (slope = 0.013) within
542 the PBL, while those values are negative above the PBL. The slope values of ε against
543 VWS are positive regardless of within or above the PBL, where the largest value of 18.1 is
544 observed above the PBL under cloudy-sky conditions, and the smallest value of 9.5 is
545 observed in the PBL under clear-sky conditions.

546 Both the thermodynamic effect of $T_s - T_a$ and the dynamic effect of VWS enhance the
547 development of turbulence under clear-sky or cloudy-sky conditions in the PBL. In the
548 lower PBL ($0.3 \leq z/z_i \leq 0.5$), $T_s - T_a$ has a larger positive slope with ε , which suggests
549 that thermal forcing emerges as the dominant factor influencing development of the
550 turbulence and PBL. By comparison, in the upper PBL ($0.6 \leq z/z_i \leq 1.0$), $T_s - T_a$ and VWS
551 jointly influence the development of turbulence, with larger $T_s - T_a$ leading to unstable
552 atmospheric stability and stronger turbulence. Above the PBL ($1.0 < z/z_i \leq 2.0$), VWS
553 becomes the dominant factor influencing the development of turbulence. Compared to
554 clear-sky conditions, on one hand, clouds would diminish $T_s - T_a$, resulting in decreased
555 heat transfer from the surface to the PBL top, thereby weakening turbulence within the
556 lower PBL ($0.3 \leq z/z_i \leq 0.5$), inhibiting PBL development, and decreasing z_i . On the other
557 hand, the stronger wind shear process would enhance the turbulence above PBL under the
558 cloudy-sky conditions.

559 Although the above-mentioned findings of the PBL turbulence over the TP are the first
560 results from profiling network observations to the best of our knowledge, fine-resolution
561 spatial distribution remains unclear, largely due to the sparse distribution of RWP network
562 on the TP. On top of this, the role of roughness length, vertical velocity, and entrainment
563 remains unknown in the variation and evolution of atmospheric turbulence, which warrants
564 further in-depth studies based on intensive field campaigns, in combination with theoretical
565 analysis and numerical simulation experiments in the future.

566

567 **Data Availability**

568 The authors would like to acknowledge the National Meteorological Information Centre
569 (NMIC) of China Meteorological Administration (CMA) (<https://data.cma.cn>) for
570 providing the high-resolution radar wind profiler and ground-based meteorological data,
571 which can be only accessed via registration. We are grateful to ECMWF for providing
572 ERA5 hourly data (<https://www.ecmwf.int/en/forecasts/datasets/reanalysis-datasets/era5/>).

573 **Acknowledgments**

574 This work was jointly supported by the National Natural Science Foundation of China
575 under grant 42325501, the High Impact Weather Key Laboratory of CMA, and NSFC
576 under grants U2142209 and 42105090, the China Meteorological Administration Xiong'an
577 Atmospheric Boundary Layer Key Laboratory under grant 2023LABL-B06, and Chinese
578 Academy of Meteorological Sciences under grants 2021KJ008 and 2024Z003. Last but not
579 least, we appreciated tremendously the constructive comments and suggestions made by
580 the anonymous reviewers that significantly improved the quality of our manuscript.

581 **Author Contributions**

582 The study was completed with close cooperation between all authors. JG and XG conceived
583 of the idea for this work. DM performed the analysis, DM and JG drafted the original
584 manuscript with contributions from XG, NL, YS, ZZ and NT. YW, HL, FZ, BT, HX and
585 TC provided useful suggestions and comments for the study and helped revise the
586 manuscript.

587 **Completing interests**

588 The authors declare that they have no conflict of interest.

589 **References**

- 590 Adler, B. and Kalthoff, N.: Multi-scale transport processes observed in the boundary layer
591 over a mountainous island, *Bound.-Layer Meteor.*, 153, 515-537,
592 <https://doi.org/10.1007/s10546-014-9957-8>, 2014.
- 593 Angel A C, Manoj M G. A novel method of estimating atmospheric boundary layer height
594 using a 205 MHz VHF radar. *Sci. Total. Environ.*,
595 <https://doi.org/10.1016/j.scitotenv.2023.168109>, 907: 168109, 2024.
- 596 Banerjee, T., Brugger, P., De Roo, F., Kröniger, K., Yakir, D., Rotenberg, E., and Mauder,
597 M.: Turbulent transport of energy across a forest and a semiarid shrubland, *Atmos.*
598 *Chem. Phys.*, 18, 10025–10038, <https://doi.org/10.5194/acp-18-10025-2018>, 2018.

599 Bianco, L., Wilczak, J. M., and White, A. B.: Convective boundary layer depth estimation
600 from wind profilers: Statistical comparison between an automated algorithm and
601 expert estimations, *J. Atmos. Ocean. Technol.*, 25, 1397-1413,
602 <https://doi.org/10.1175/2008jtecha981.1>, 2008.

603 Bodenschatz, E., Malinowski, S. P., Shaw, R. A., and Stratmann, F.: Can we understand
604 clouds without turbulence? *Science*, 327, 970-971,
605 <https://doi.org/10.1126/science.1185138>, 2010.

606 Che, J. H. and Zhao, P.: Characteristics of the summer atmospheric boundary layer height
607 over the Tibetan Plateau and influential factors, *Atmos. Chem. Phys.*, 21, 5253-5268,
608 <https://doi.org/10.5194/acp-21-5253-2021>, 2021.

609 Chechin, D. G., Lüpkes, C., Hartmann, J., Ehrlich, A., and Wendisch, M.: Turbulent
610 structure of the Arctic boundary layer in early summer driven by stability, wind shear
611 and cloud-top radiative cooling: ACLOUD airborne observations, *Atmos. Chem.*
612 *Phys.*, 23, 4685-4707, <https://doi.org/10.5194/acp-23-4685-2023>, 2023.

613 Chen, X. L., Añel, J. A., Su, Z. B., de la Torre, L., Kelder, H., van Peet, J., and Ma, Y. M.:
614 The deep atmospheric boundary layer and its significance to the stratosphere and
615 troposphere exchange over the Tibetan Plateau, *PLoS One*, 8, 9,
616 <https://doi.org/10.1371/journal.pone.0056909>, 2013.

617 Chen, X. L., Skerlak, B., Rotach, M. W., Añel, J. A., Su, Z., Ma, Y. M., and Li, M. S.:
618 Reasons for the extremely high-ranging planetary boundary layer over the western
619 Tibetan Plateau in winter, *J. Atmos. Sci.*, 73, 2021-2038, [https://doi.org/10.1175/jas-](https://doi.org/10.1175/jas-d-15-0148.1)
620 [d-15-0148.1](https://doi.org/10.1175/jas-d-15-0148.1), 2016.

621 Chen, Z., Tian, Y., Wang, Y., Bi, Y., Wu, X., Huo, J., Pan, L., Wang, Y., and Lü, D.:
622 Turbulence parameters measured by the Beijing mesosphere–stratosphere–
623 troposphere radar in the troposphere and lower stratosphere with three models:
624 comparison and analyses, *Atmos. Meas. Tech.*, 15, 4785–4800,
625 <https://doi.org/10.5194/amt-15-4785-2022>, 2022.

626 Collaud Coen, M., Praz, C., Haeefe, A., Ruffieux, D., Kaufmann, P., and Calpini, B.:
627 Determination and climatology of the planetary boundary layer height above the
628 Swiss plateau by in situ and remote sensing measurements as well as by the COSMO-

629 2 model, *Atmos. Chem. Phys.*, 14, 13205–13221, [https://doi.org/10.5194/acp-14-](https://doi.org/10.5194/acp-14-13205-2014)
630 13205-2014, 2014.

631 Dai, C., Wang, Q., Kalogiros, J. A., Lenschow, D. H., Gao, Z., and Zhou, M.: Determining
632 boundary-layer height from aircraft measurements, *Bound.-Layer Meteor.*, 152, 277-
633 302, <https://doi.org/10.1007/s10546-014-9929-z>, 2014.

634 Davis E V, Rajeev K, and Mishra M K: Effect of clouds on the diurnal evolution of the
635 atmospheric boundary-layer height over a tropical coastal station. *Bound.-Layer*
636 *Meteor.*, 175: 135-152, <https://doi.org/10.1007/s10546-019-00497-6>, 2020.

637 Dodson, D. S. and Griswold, J. D. S.: Turbulent and boundary layer characteristics during
638 VOCALS-REx, *Atmos. Chem. Phys.*, 21, 1937-1961, [https://doi.org/10.5194/acp-21-](https://doi.org/10.5194/acp-21-1937-2021)
639 1937-2021, 2021.

640 Duncan Jr., J. B., Bianco, L., Adler, B., Bell, T., Djalalova, I. V., Riihimaki, L., Sedlar, J.,
641 Smith, E. N., Turner, D. D., Wagner, T. J., and Wilczak, J. M.: Evaluating convective
642 planetary boundary layer height estimations resolved by both active and passive
643 remote sensing instruments during the CHEESEHEAD19 field campaign, *Atmos.*
644 *Meas. Tech.*, 15, 2479–2502, <https://doi.org/10.5194/amt-15-2479-2022>, 2022.

645 Guo, J. P., Li, Y., Cohen, J. B., Li, J., Chen, D. D., Xu, H., Liu, L., Yin, J. F., Hu, K. X.,
646 and Zhai, P. M.: Shift in the temporal trend of boundary layer height in China using
647 long-term (1979-2016) radiosonde data, *Geophys. Res. Lett.*, 46, 6080-6089,
648 <https://doi.org/10.1029/2019gl082666>, 2019.

649 Guo, J. P., Miao, Y. C., Zhang, Y., Liu, H., Li, Z. Q., Zhang, W. C., He, J., Lou, M. Y.,
650 Yan, Y., Bian, L. G., and Zhai, P.: The climatology of planetary boundary layer height
651 in China derived from radiosonde and reanalysis data, *Atmos. Chem. Phys.*, 16,
652 13309-13319, <https://doi.org/10.5194/acp-16-13309-2016>, 2016.

653 Guo, J. P., Liu, B. M., Gong, W., Shi, L. J., Zhang, Y., Ma, Y. Y., Zhang, J., Chen, T. M.,
654 Bai, K. X., Stoffelen, A., de Leeuw, G., and Xu, X. F.: First comparison of wind
655 observations from ESA's satellite mission Aeolus and ground-based radar wind
656 profiler network of China, *Atmos. Chem. Phys.*, 21, 2945-2958,
657 <https://doi.org/10.5194/acp-21-2945-2021>, 2021a.

658 Guo, J. P., Zhang, J., Yang, K., Liao, H., Zhang, S. D., Huang, K. M., Lv, Y. M., Shao, J.,
659 Yu, T., Tong, B., Li, J., Su, T. N., Yim, S. H. L., Stoffelen, A., Zhai, P. M., and Xu,

660 X. F.: Investigation of near-global daytime boundary layer height using high-
661 resolution radiosondes: first results and comparison with ERA5, MERRA-2, JRA-55,
662 and NCEP-2 reanalyses, *Atmos. Chem. Phys.*, 21, 17079-17097,
663 <https://doi.org/10.5194/acp-21-17079-2021>, 2021b.

664 Guo, X. R., Guo, J. P., Zhang, D. L., and Yun, Y. X.: Vertical divergence profiles as
665 detected by two wind-profiler mesonets over East China: Implications for nowcasting
666 convective storms, *Q. J. R. Meteorol. Soc.*, 149, 1629-1649,
667 <https://doi.org/10.1002/qj.4474>, 2023.

668 Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., Nicolas,
669 J., Peubey, C., Radu, R., Schepers, D., Simmons, A., Soci, C., Abdalla, S., Abellan,
670 X., Balsamo, G., Bechtold, P., Biavati, G., Bidlot, J., Bonavita, M., De Chiara, G.,
671 Dahlgren, P., Dee, D., Diamantakis, M., Dragani, R., Flemming, J., Forbes, R.,
672 Fuentes, M., Geer, A., Haimberger, L., Healy, S., Hogan, R. J., Hólm, E., Janisková,
673 M., Keeley, S., Laloyaux, P., Lopez, P., Lupu, C., Radnoti, G., de Rosnay, P., Rozum,
674 I., Vamborg, F., Villaume, S., and Thépaut, J. N.: The ERA5 global reanalysis, *Q. J.
675 R. Meteorol. Soc.*, 146, 1999-2049, <https://doi.org/10.1002/qj.3803>, 2020.

676 Huang, J. P., Zhou, X. J., Wu, G. X., Xu, X. D., Zhao, Q. Y., Liu, Y. M., Duan, A. M., Xie,
677 Y. K., Ma, Y. M., Zhao, P., Yang, S., Yang, K., Yang, H. J., Bian, J. C., Fu, Y. F., Ge,
678 J. M., Liu, Y. Z., Wu, Q. G., Yu, H. P., Wang, B. B., Bao, Q., and Qie, K.: Global
679 climate impacts of land-surface and atmospheric processes over the Tibetan Plateau,
680 *Rev. Geophys.*, 61, 39, <https://doi.org/10.1029/2022rg000771>, 2023.

681 Huang, T., Yim, S. H. L., Yang, Y., Lee, O. S. M., Lam, D. H. Y., Cheng, J. C. H., and
682 Guo, J.: Observation of turbulent mixing characteristics in the typical daytime cloud-
683 topped boundary layer over Hong Kong in 2019, *Remote Sens.*, 12, 1533,
684 <https://doi.org/10.3390/RS12091533>, 2020.

685 Jacoby-Koaly, S., Campistron, B., Bernard, S., Bénech, B., Girard-Ardhuin, F., Dessens,
686 J., Dupont, E., and Carissimo, B.: Turbulent dissipation rate in the boundary layer via
687 UHF wind profiler Doppler spectral width measurements, *Bound.-Layer Meteor.*, 103,
688 361-389, <https://doi.org/10.1023/a:1014985111855>, 2002.

689 Kotthaus, S., Bravo-Aranda, J. A., Coen, M. C., Guerrero-Rascado, J. L., Costa, M. J.,
690 Cimini, D., O'Connor, E. J., Hervo, M., Alados-Arboledas, L., Jiménez-Portaz, M.,

691 Mona, L., Ruffieux, D., Illingworth, A., and Haeffelin, M.: Atmospheric boundary
692 layer height from ground-based remote sensing: a review of capabilities and
693 limitations, *Atmos. Meas. Tech.*, 16, 433-479, [https://doi.org/10.5194/amt-16-433-](https://doi.org/10.5194/amt-16-433-2023)
694 2023, 2023.

695 Lai, Y., Chen, X. L., Ma, Y. M., Chen, D. L., and Zhaxi, S. L.: Impacts of the westerlies
696 on planetary boundary layer growth over a valley on the north side of the Central
697 Himalayas, *J. Geophys. Res.-Atmos.*, 126, 20, <https://doi.org/10.1029/2020jd033928>,
698 2021.

699 Li, Y., Y. Wu, J. Tang, P. Zhu, Z. Gao, and Y. Yang. Quantitative evaluation of wavelet
700 analysis method for turbulent flux calculation of non-stationary series, *Geophys. Res.*
701 *Lett.*, 50(5), e2022GL101591, <http://dx.doi.org/10.1029/2022GL101591>, 2023.

702 Li, Z. G., Lyu, S. H., Wen, L. J., Zhao, L., Ao, Y. H., and Wang, S. Y.: Effect of a cold,
703 dry air incursion on atmospheric boundary layer processes over a high-altitude lake in
704 the Tibetan Plateau, *Atmos. Res.*, 185, 32-43,
705 <https://doi.org/10.1016/j.atmosres.2016.10.024>, 2017a.

706 Li, Z. Q., Guo, J. P., Ding, A. J., Liao, H., Liu, J. J., Sun, Y. L., Wang, T. J., Xue, H. W.,
707 Zhang, H. S., and Zhu, B.: Aerosol and boundary-layer interactions and impact on air
708 quality, *Natl. Sci. Rev.*, 4, 810-833, <https://doi.org/10.1016/10.1093/nsr/nwx117>,
709 2017b.

710 Liu, B. M., Guo, J. P., Gong, W., Shi, L. J., Zhang, Y., and Ma, Y. Y.: Characteristics and
711 performance of wind profiles as observed by the radar wind profiler network of China,
712 *Atmos. Meas. Tech.*, 13, 4589-4600, <https://doi.org/10.5194/amt-13-4589-2020>, 2020.

713 Liu, B. M., Ma, Y. Y., Guo, J. P., Gong, W., Zhang, Y., Mao, F. Y., Li, J., Guo, X. R., and
714 Shi, Y. F.: Boundary layer heights as derived from ground-based radar wind profiler
715 in Beijing, *IEEE Trans. Geosci. Remote Sensing*, 57, 8095-8104,
716 <https://doi.org/10.1109/tgrs.2019.2918301>, 2019.

717 Lv, Y. M., Guo, J. P., Li, J., Cao, L. J., Chen, T. M., Wang, D., Chen, D. D., Han, Y., Guo,
718 X. R., Xu, H., Liu, L., Solanki, R., and Huang, G.: Spatiotemporal characteristics of
719 atmospheric turbulence over China estimated using operational high-resolution
720 soundings, *Environ. Res. Lett.*, 16, 13, <https://doi.org/10.1088/1748-9326/abf461>,
721 2021.

722 Ma, Y. M., Yao, T. D., Zhong, L., Wang, B. B., Xu, X. D., Hu, Z. Y., Ma, W. Q., Sun, F.
723 L., Han, C. B., Li, M. S., Chen, X. L., Wang, J. M., Li, Y. Q., Gu, L. L., Xie, Z. P.,
724 Liu, L., Sun, G. H., Wang, S. J., Zhou, D. G., Zuo, H. C., Xu, C., Liu, X., Wang, Y.
725 J., and Wang, Z. Y.: Comprehensive study of energy and water exchange over the
726 Tibetan Plateau: A review and perspective: From GAME/Tibet and CAMP/Tibet to
727 TORP, TPEORP, and TPEITORP, *Earth-Sci. Rev.*, 104312, 2023.

728 McCaffrey, K., Bianco, L., and Wilczak, J. M.: Improved observations of turbulence
729 dissipation rates from wind profiling radars, *Atmos. Meas. Tech.*, 10, 2595-2611,
730 <https://doi.org/10.5194/amt-10-2595-2017>, 2017.

731 Muhsin, M., Sunilkumar, S. V., Ratnam, M. V., Parameswaran, K., Murthy, B. V. K.,
732 Ramkumar, G., and Rajeev, K.: Diurnal variation of atmospheric stability and
733 turbulence during different seasons in the troposphere and lower stratosphere derived
734 from simultaneous radiosonde observations at two tropical stations, in the Indian
735 Peninsula, *Atmos. Res.*, 180, 12-23, <https://doi.org/10.1016/j.atmosres.2016.04.021>,
736 2016.

737 Muñoz-Esparza, D., Sharman, R. D., and Lundquist, J. K.: Turbulence Dissipation Rate in
738 the Atmospheric Boundary Layer: Observations and WRF Mesoscale Modeling
739 during the XPIA Field Campaign, *Mon. Weather Rev.*, 146, 351-371,
740 <https://doi.org/10.1175/mwr-d-17-0186.1>, 2018.

741 Nastrom, G. D.: Doppler radar spectral width broadening due to beamwidth and wind shear,
742 *Ann. Geophys.-Atmos. Hydrospheres Space Sci.*, 15, 786-796,
743 <https://doi.org/10.1007/s00585-997-0786-7>, 1997.

744 Nicholls, S.: The dynamics of stratocumulus: Aircraft observations and comparisons with
745 a mixed layer model, *Q. J. Roy. Meteor. Soc.*, 110, 783-820,
746 <https://doi.org/10.1002/qj.49711046603>, 1984.

747 Ruan, Z., Mu, R. Q., Wei, M., and Ge, R. S.: Spectrum analysis of wind profiling radar
748 measurements, *J. Meteorol. Res.*, 28, 656-667, [https://doi.org/10.1007/s13351-014-](https://doi.org/10.1007/s13351-014-3171-y)
749 3171-y, 2014.

750 Sedlar, J., Riihimaki, L. D., Turner, D. D., Duncan, J., Adler, B., Bianco, L., Lantz, K., and
751 Wilczak, J.: Investigating the impacts of daytime boundary layer clouds on surface

752 energy fluxes and boundary layer structure during CHEESEHEAD19, *J. Geophys.*
753 *Res.-Atmos.*, 127, e2021JD036060, <https://doi.org/10.1029/2021JD036060>, 2022

754 Schumann, U. and Moeng, C. H.: Plume budgets in clear and cloudy convective boundary
755 layers, *J. Atmos. Sci.*, 48, 1758-1770, <https://doi.org/10.1175/1520->
756 0469(1991)048<1758:Pbicac>2.0.Co;2, 1991.

757 Slättberg, N., Lai, H. W., Chen, X. L., Ma, Y. M., and Chen, D. L.: Spatial and temporal
758 patterns of planetary boundary layer height during 1979-2018 over the Tibetan Plateau
759 using ERA5, *Int. J. Climatol.*, 42, 3360-3377, <https://doi.org/10.1002/joc.7420>, 2022.

760 Solanki, R., Guo, J. P., Lv, Y. M., Zhang, J., Wu, J. Y., Tong, B., and Li, J.: Elucidating
761 the atmospheric boundary layer turbulence by combining UHF radar wind profiler
762 and radiosonde measurements over urban area of Beijing, *Urban CLim.*, 43, 13,
763 <https://doi.org/10.1016/j.uclim.2022.101151>, 2022.

764 Solanki, R., Guo, J. P., Li, J., Singh, N., Guo, X. R., Han, Y., Lv, Y. M., Zhang, J., and Liu,
765 B. M.: Atmospheric-boundary-layer-height variation over mountainous and urban
766 stations in Beijing as derived from radar wind-profiler measurements, *Bound.-Layer*
767 *Meteor.*, 181, 125-144, <https://doi.org/10.1007/s10546-021-00639-9>, 2021.

768 Stull, R. B.: Mean Boundary Layer Characteristics, in: *An Introduction to boundary layer*
769 *meteorology*, edited by: Stull, R. B., Springer Netherlands, Dordrecht, 1–27,
770 https://doi.org/10.1007/978-94-009-3027-8_1, 1988.

771 Su, T. N., Li, Z. Q., and Zheng, Y. T.: Cloud-Surface Coupling Alters the Morning
772 Transition From Stable to Unstable Boundary Layer, *Geophys. Res. Lett.*, 50, 9,
773 <https://doi.org/10.1029/2022gl102256>, 2023.

774 Sun, W., Li, L., and Wang, B.: Reducing the biases in shortwave cloud radiative forcing in
775 tropical and subtropical regions from the perspective of boundary layer processes, *Sci.*
776 *China Earth Sci.*, 59, 1427–1439, <https://doi.org/10.1007/s11430-016-5290-z>, 2016.

777 Teixeira, J., Piepmeier, J. R., Nehrir, A. R., Ao, C. O., Chen, S. S., Clayson, C. A., Fridlind,
778 A. M., Lebsock, M., McCarty, W., Salmun, H., Santanello, J. A., Turner, D. D., Wang,
779 Z., and Zeng, X.: Toward a global planetary boundary layer observing system: the
780 NASA PBL incubation study team report, NASA PBL Incubation Study Team, 134
781 pp., available at: <https://science.nasa.gov/science-red/s3fs->

782 public/atoms/files/NASAPBLIncubationFinalReport.pdf, last access: 15 November
783 2021.

784 Wan, X., Zheng, J. F., Wan, R., Xu, G. R., Qin, J. F., and Yi, L.: Intercomparison of cloud
785 vertical structures over four different stations of the eastern slope of the Tibetan
786 Plateau in summer using Ka-band millimeter-wave radar measurements, *Remote*
787 *Sens.*, 14, 19, <https://doi.org/10.3390/rs14153702>, 2022.

788 Wang, C. X., Ma, Y. M., and Han, C. B.: Research on the atmospheric boundary layer
789 structure and its development mechanism in the Tibetan Plateau, *Adv. Atmos. Sci.*,
790 38, 414-428, 2023a.

791 Wang, M. Z. and Zhang, J. T.: The relationship among summer atmospheric boundary layer
792 height over the Taklimakan Desert, its land surface parameters and Eurasian
793 circulation, *Atmos. Sci. Lett.*, 23, 13, <https://doi.org/10.1002/asl.1122>, 2022.

794 Wang, M. Z., Lu, H., Ming, H., and Zhang, J. T.: Vertical structure of summer clear-sky
795 atmospheric boundary layer over the hinterland and southern margin of Taklamakan
796 Desert, *Meteorol. Appl.*, 23, 438-447, <https://doi.org/10.1002/met.1568>, 2016.

797 Wang, S. Q., Guo, J. P., Xian, T., Li, N., Meng, D. L., Li, H. J., and Cheng, W.:
798 Investigation of low-level supergeostrophic wind and Ekman spiral as observed by a
799 radar wind profiler in Beijing, *Front. Environ. Sci.*, 11, 9,
800 <https://doi.org/10.3389/fenvs.2023.1195750>, 2023b.

801 Wang, Y. J., Zeng, X. B., Xu, X. D., Xie, F. Q., and Zhao, Y.: Improving the estimate of
802 summer daytime planetary boundary layer height over land from GPS radio
803 occultation data, *Geophys. Res. Lett.*, 49, 9, <https://doi.org/10.1029/2021gl096304>,
804 2022.

805 Wang, Y. J., Xu, X. D., Zhao, T. L., Sun, J. H., Yao, W. Q., and Zhou, M. Y.: Structures
806 of convection and turbulent kinetic energy in boundary layer over the southeastern
807 edge of the Tibetan Plateau, *Sci. China-Earth Sci.*, 58, 1198-1209,
808 <https://doi.org/10.1007/s11430-015-5054-1>, 2015.

809 Wang, Y. J., Zeng, X. B., Xu, X. D., Welty, J., Lenschow, D. H., Zhou, M. Y., and Zhao,
810 Y.: Why are there more summer afternoon low clouds over the Tibetan Plateau
811 compared to eastern China? *Geophys. Res. Lett.*, 47, 10,
812 <https://doi.org/10.1029/2020gl089665>, 2020.

813 White, A. B., Lataitis, R. J., and Lawrence, R. S.: Space and time filtering of remotely
814 sensed velocity turbulence, *J. Atmos. Ocean. Technol.*, 16, 1967-1972,
815 [https://doi.org/10.1175/1520-0426\(1999\)016<1967:Satfor>2.0.Co;2](https://doi.org/10.1175/1520-0426(1999)016<1967:Satfor>2.0.Co;2), 1999.

816 Wu, J. Y., Guo, J. P., Yun, Y. X., Yang, R. F., Guo, X. R., Meng, D. L., Sun, Y. P., Zhang,
817 Z., Xu, H., and Chen, T. M.: Can ERA5 reanalysis data characterize the pre-storm
818 environment? *Atmos. Res.*, 297, 18, <https://doi.org/10.1016/j.atmosres.2023.107108>,
819 2024.

820 Xu, L. J., Liu, H. Z., Du, Q., and Xu, X. D.: The assessment of the planetary boundary
821 layer schemes in WRF over the central Tibetan Plateau, *Atmos. Res.*, 230, 12,
822 <https://doi.org/10.1016/j.atmosres.2019.104644>, 2019.

823 Xu, X. D., Tang, Y., Wang, Y. J., Zhang, H. S., Liu, R. X., and Zhou, M. Y.: Triggering
824 effects of large topography and boundary layer turbulence on convection over the
825 Tibetan Plateau, *Atmos. Chem. Phys.*, 23, 3299-3309, [https://doi.org/10.5194/acp-23-](https://doi.org/10.5194/acp-23-3299-2023)
826 [3299-2023](https://doi.org/10.5194/acp-23-3299-2023), 2023.

827 Xu, X. D., Zhou, M. Y., Chen, J. Y., Bian, L. G., Zhang, G. Z., Liu, H. Z., Li, S. M., Zhang,
828 H. S., Zhao, Y. J., Suolong, D.J., and Wang, J. Z.: A comprehensive physical pattern
829 of land-air dynamic and thermal structure on the Qinghai-Xizang Plateau, *Sci. China*
830 *Ser. D-Earth Sci.*, 45, 577-594, <https://doi.org/10.1360/02yd9060>, 2002.

831 Yang, B., Qian, Y., Berg, L. K., Ma, P. L., Wharton, S., Bulaevskaya, V., Yan, H. P., Hou,
832 Z. S., and Shaw, W. J.: Sensitivity of Turbine-Height Wind Speeds to Parameters in
833 Planetary Boundary-Layer and Surface-Layer Schemes in the Weather Research and
834 Forecasting Model, *Bound.-Layer Meteor.*, 162, 117-142,
835 <https://doi.org/10.1007/s10546-016-0185-2>, 2017.

836 Yang, R. F., Guo, J. P., Deng, W. L., Li, N., Fan, J. H., Meng, D. L., Liu, Z., Sun, Y. P.,
837 Zhang, G. L., and Liu, L. H.: Investigation of turbulent dissipation rate profiles from
838 two radar wind profilers at plateau and plain stations in the north China plain, *Remote*
839 *Sens.*, 15, 14, <https://doi.org/10.3390/rs15164103>, 2023.

840 Zhang, L., Zhang, H. S., Li, Q. H., Wei, W., Cai, X. H., Song, Y., Mamtimin, A., Wang,
841 M. Z., Yang, F., Wang, Y., and Zhou, C. L.: Turbulent mechanisms for the deep
842 convective boundary layer in the Taklimakan desert, *Geophys. Res. Lett.*, 49, 9,
843 <https://doi.org/10.1029/2022gl099447>, 2022.

844 Zhang, Y., Guo, J. P., Yang, Y. J., Wang, Y., and Yim, S. H. L.: Vertica wind shear
845 modulates particulate matter pollutions: a perspective from radar wind profiler
846 observations in Beijing, China, *Remote Sens.*, 12, 17,
847 <https://doi.org/10.3390/rs12030546>, 2020.

848 Zhao, P., Li, Y. Q., Guo, X. L., Xu, X. D., Liu, Y. M., Tang, S. H., Xiao, W. M., Shi, C.
849 X., Ma, Y. M., Yu, X., Liu, H. Z., Jia, L., Chen, Y., Liu, Y. J., Li, J., Luo, D. B., Cao,
850 Y. C., Zheng, X. D., Chen, J. M., Xiao, A., Yuan, F., Chen, D. H., Pang, Y., Hu, Z.
851 Q., Zhang, S. J., Dong, L. X., Hu, J. Y., Han, S., and Zhou, X. J.: The Tibetan Plateau
852 surface-atmosphere coupling system and its weather and climate effects: the third
853 Tibetan Plateau atmospheric science experiment, *J. Meteorol. Res.*, 33, 375-399,
854 <https://doi.org/10.1007/s13351-019-8602-3>, 2019.

855

856 **Table list**

857 **Table 1.** Summary of the geographical conditions and land surface of the six radar wind
 858 profiler (RWP) stations over the Tibet Plateau (TP).

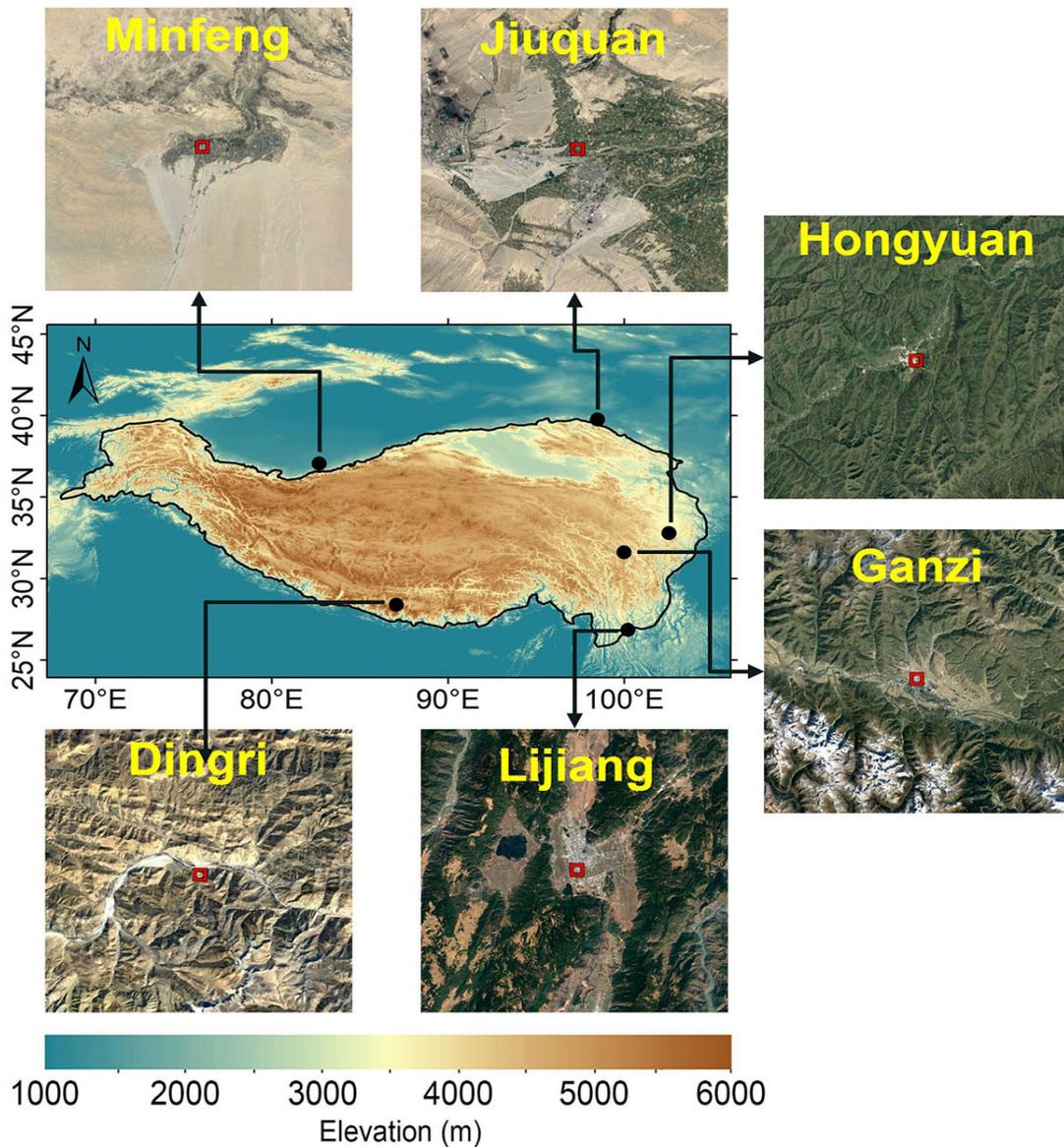
RWP station	Latitude (°E)	Longitude (°N)	Elevation (m)	Land cover types
Minfeng	82.69	37.07	1408.9	Bare land
Jiuquan	98.49	39.77	1477.2	Bare land
Dingri	87.07	28.39	4326.0	Grassland
Ganzi	100.00	31.62	3353.0	Bare land, grassland
Hongyuan	102.55	32.79	3465.0	Bare land, grassland
Lijiang	100.22	26.85	2382.4	Bare land, grass land

859

860

861 **Table 2.** Summary of the correlation of $\text{Log}_{10}\varepsilon$ at different altitude ranges under all-,
862 clear- and cloudy-sky conditions with $T_s - T_a$ and vertical wind shear (VWS) for all six
863 RWP stations. The superscript * for R indicates that the regression slope is statistically
864 significant at $p < 0.01$.

Conditions	$\text{Log}_{10}\varepsilon$ VS $T_s - T_a$	$\text{Log}_{10}\varepsilon$ VS VWS
all-sky, $0.3 \leq z/z_i \leq 2.0$	$y=0.010x-4.05$, R= 0.21*	$y=13.6x-4.19$, R=0.29*
all-sky, $0.3 \leq z/z_i \leq 1.0$	$y=0.018x-3.70$, R=0.29*	$y=13.2x-3.77$, R=0.20*
all-sky, $1.0 < z/z_i \leq 2.0$	$y=-0.004x-4.20$, R=-0.09*	$y=17.6x-4.57$, R=0.36*
clear-sky, $0.3 \leq z/z_i \leq 2.0$	$y=0.011x-4.04$, R=0.23*	$y=10.7x-4.13$, R=0.26*
clear -sky, $0.3 \leq z/z_i \leq 1.0$	$y=0.018x-3.67$, R=0.30*	$y=11.1x-3.70$, R=0.17*
clear -sky, $1.0 < z/z_i \leq 2.0$	$y=-0.005x-4.17$, R=-0.11*	$y=13.8x-4.52$, R=0.34*
cloudy-sky, $0.3 \leq z/z_i \leq 2.0$	$y=0.009x-4.06$, R=0.16*	$y=18.5x-4.29$, R=0.33*
cloudy-sky, $0.3 \leq z/z_i \leq 1.0$	$y=0.018x-3.71$, R=0.26*	$y=15.5x-3.84$, R=0.23*
cloudy-sky, $1.0 < z/z_i \leq 2.0$	$y=-0.004x-4.22$, R=-0.08*	$y=26.2x-4.67$, R=0.42*



866

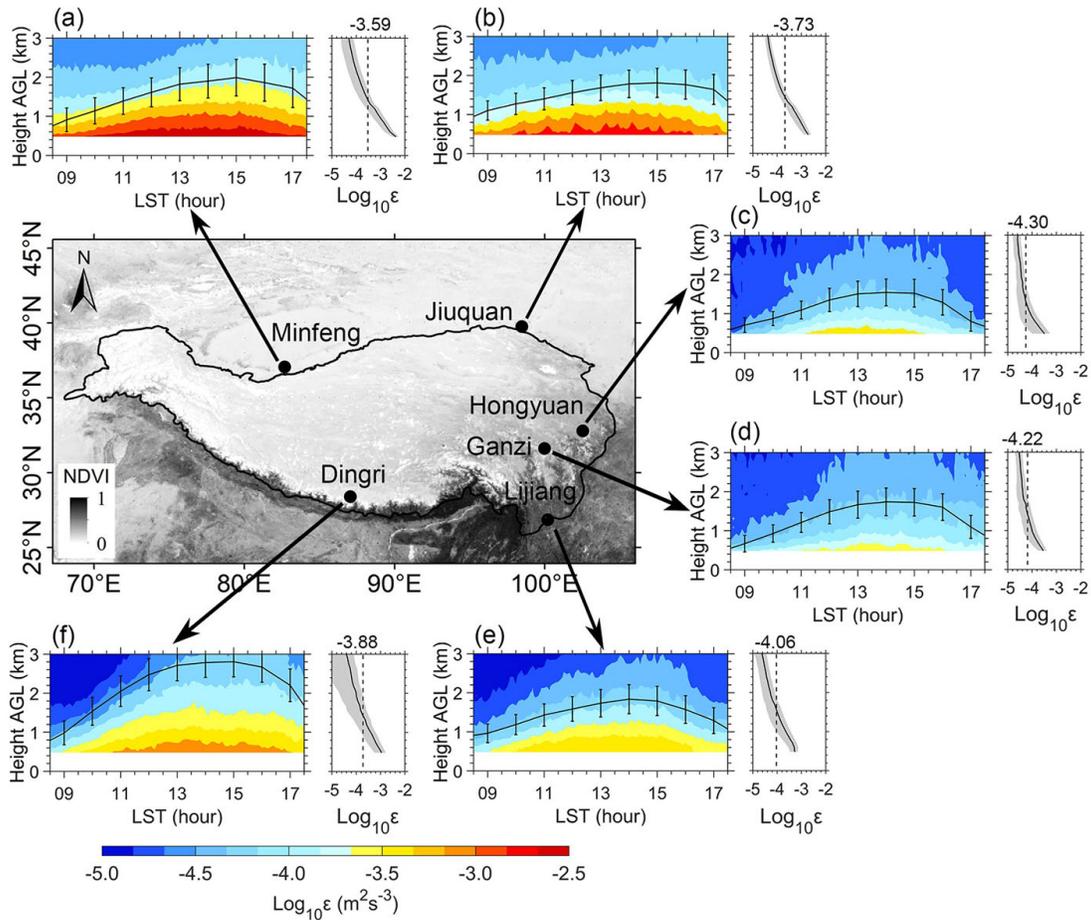
867 **Figure 1.** Spatial distribution of radar wind profiler (RWP) network comprised of six

868 stations (in black solid circles) on the Tibetan Plateau (TP). The inset map surrounding the

869 main frame denotes the RGB satellite image from Google Earth that is centered at each

870 RWP station.

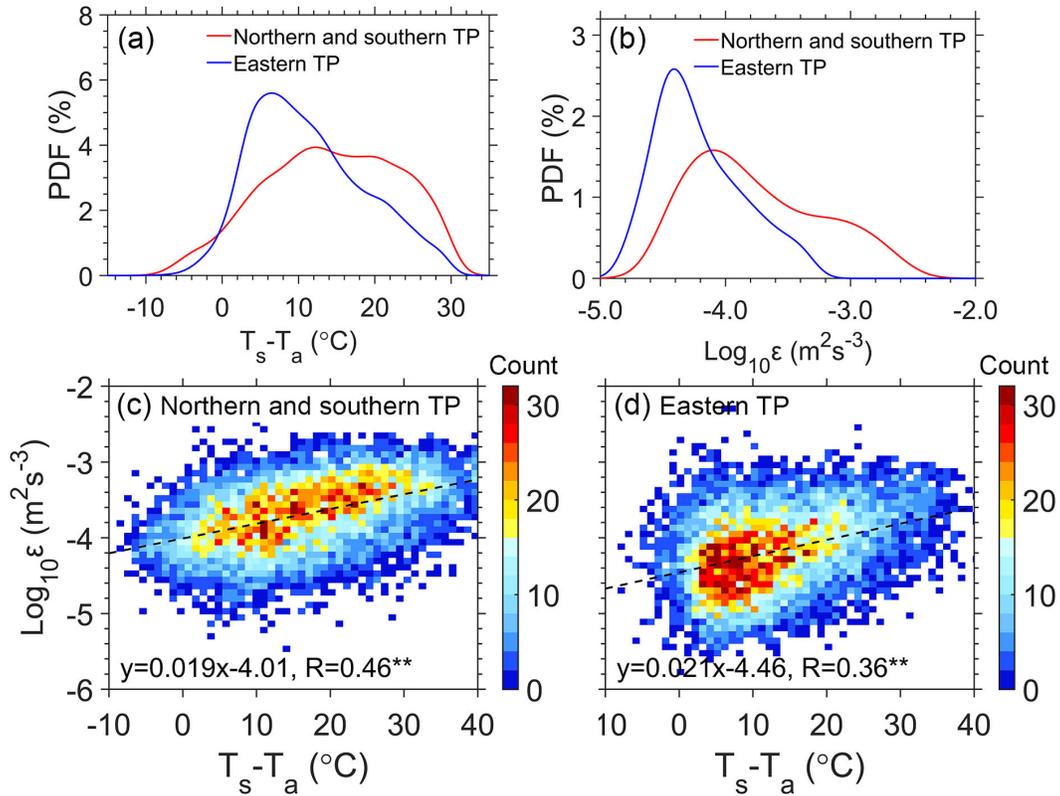
871



872

873 **Figure 2.** Spatial distribution of the diurnal evolution of the vertical profile of logarithmic
 874 turbulence dissipation rate ($Log_{10}\epsilon$ in color shading, unit: $m^2 s^{-3}$) at 120 m vertical
 875 resolution and 6 min intervals, and hourly mean planetary boundary layer height (z_i , black
 876 line, unit: km) during daytime under all-sky conditions from 0900 to 1700 LST for the
 877 period September 2022 to October 2023 as retrieved from the profiling measurements at
 878 six RWP stations over the TP. The vertical bars indicate the 0.5 standard deviations for z_i .
 879 Also shown on the right-hand side panel are temporally averaged vertical profile of ϵ (black
 880 line) and its corresponding one standard deviation (gray shading).

881



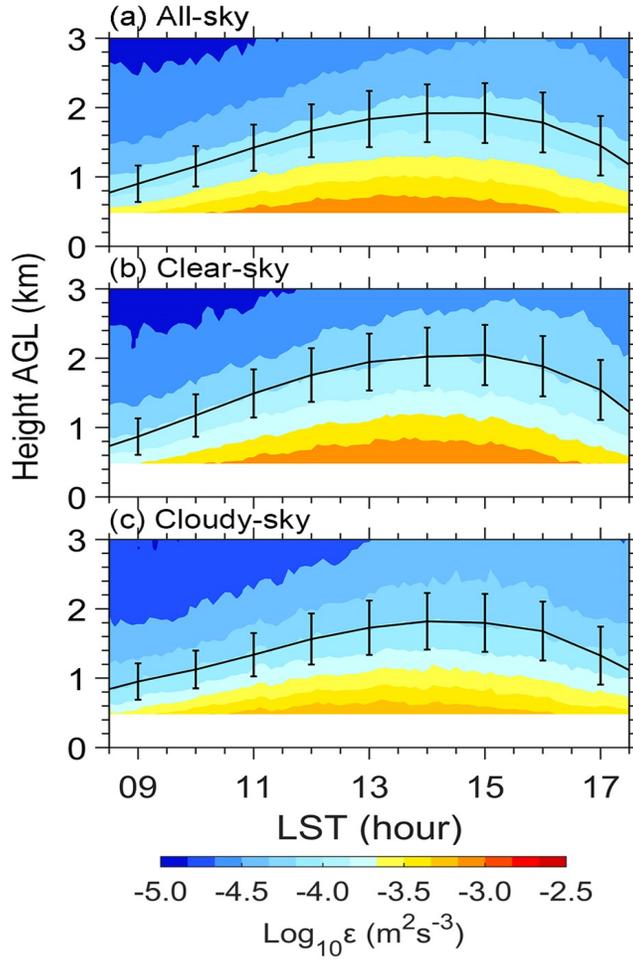
882

883 **Figure 3.** (a) PDF of surface-air temperature difference ($T_s - T_a$) for the northern and
 884 southern TP (red line) and eastern TP (blue line), (b) same as (a), but for PDF of $\text{Log}_{10}\epsilon$
 885 estimated from the measurements of radar wind profilers (RWPs) at the heights ranging
 886 from 0.5 to 3.0 km, (c) scatter plots of $\text{Log}_{10}\epsilon$ as a function of $T_s - T_a$ in the northern
 887 and southern TP, (d) same as (c), but for the eastern TP during daytime under all-sky
 888 conditions from 0900 to 1700 local standard time (LST) for the period September 2022 to
 889 October 2023. The superscript ** for R indicates that the regression slope is statistically
 890 significant at $p < 0.01$ level.

891

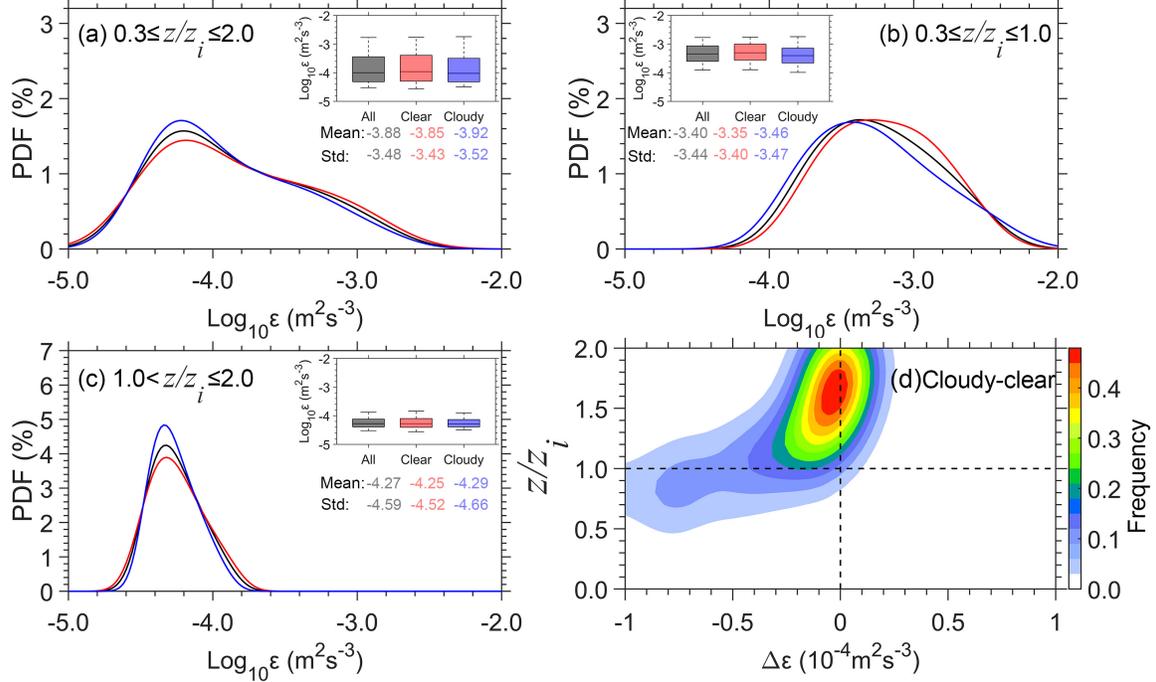
892

893



894

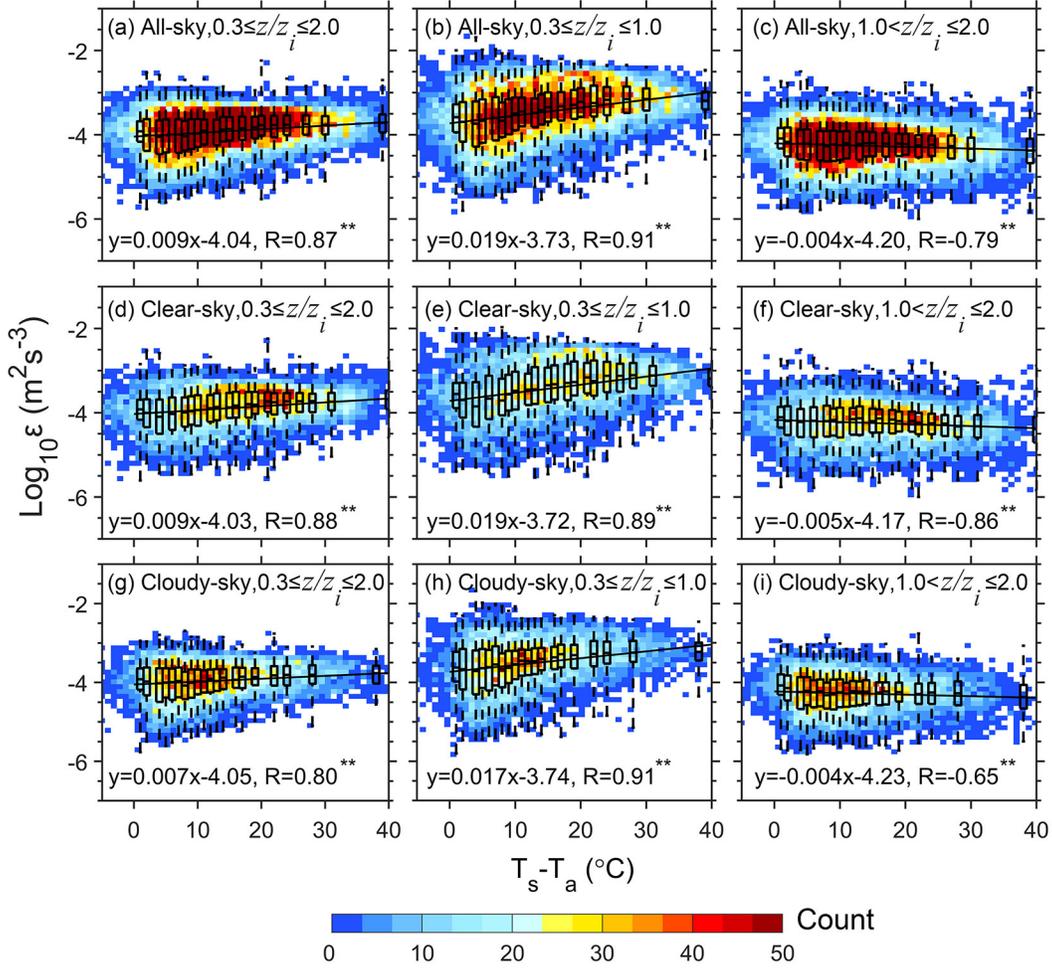
895 **Figure 4.** Diurnal evolution of the vertical profile of $\text{Log}_{10} \varepsilon$ (color shading, unit: $\text{m}^2 \text{s}^{-3}$)
 896 and z_i (solid line, unit: km) averaged over the six RWP stations over the TP during
 897 daytime from 0900 to 1700 LST for the period September 2022 to October 2023 for (a) all-
 898 sky conditions, (b) clear-sky conditions and (c) cloudy-sky conditions. The vertical bars
 899 indicate the 0.5 standard deviations.



900

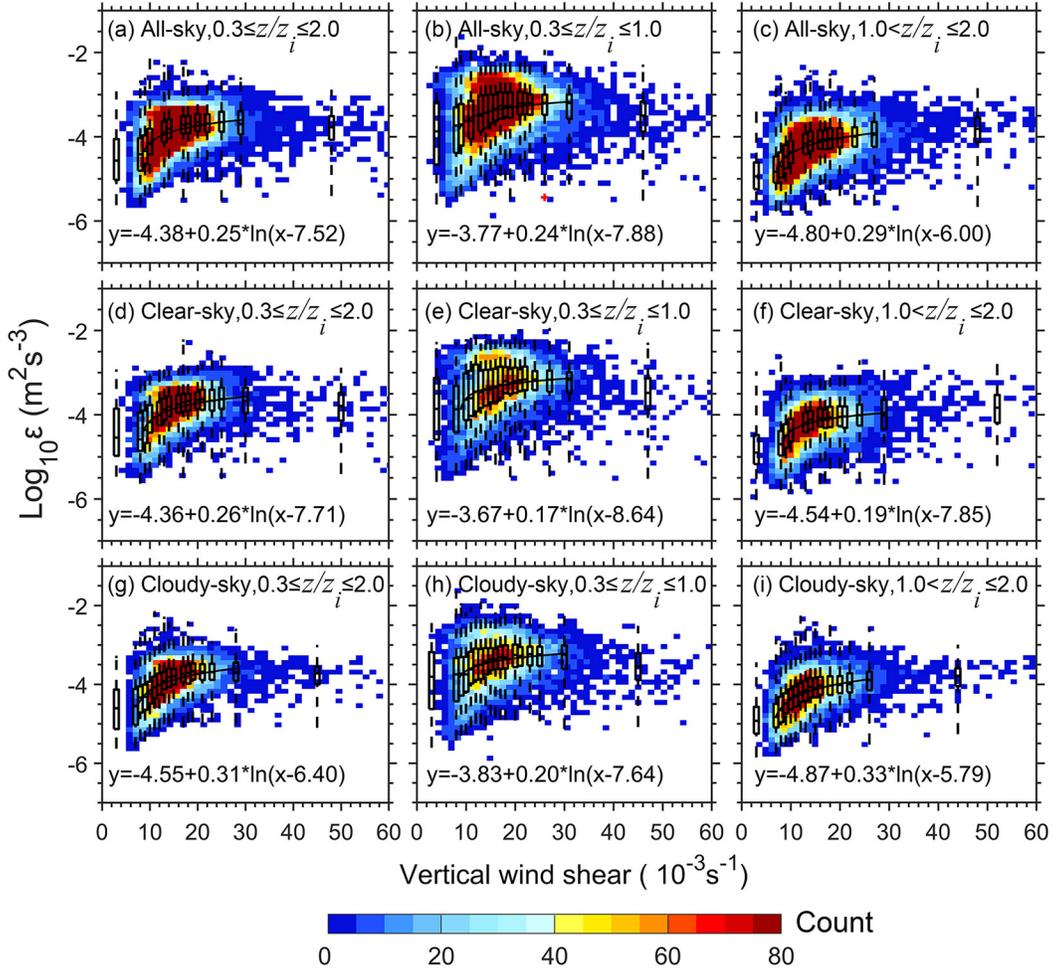
901 **Figure 5.** PDF of daytime $\text{Log}_{10}\epsilon$ (a) in the whole lower troposphere ($0.3 \leq z/z_i \leq 2.0$), (b)
 902 in the PBL ($0.3 \leq z/z_i \leq 1.0$) and (c) above the PBL ($1.0 < z/z_i \leq 2.0$) over the TP under all-
 903 sky (black), clear-sky (red) and cloudy-sky (blue) conditions, respectively. (d) Normalized
 904 contoured frequency by altitude diagram (NCFAD) for the difference of ϵ between cloudy-
 905 sky and clear-sky conditions ($\Delta\epsilon$) over the TP. Note that z_i denotes the depth of the PBL,
 906 the height (z) and turbulence dissipation rate (ϵ) is normalized by z_i in order to give a
 907 nondimensional vertical coordinate in the form of z/z_i .

908



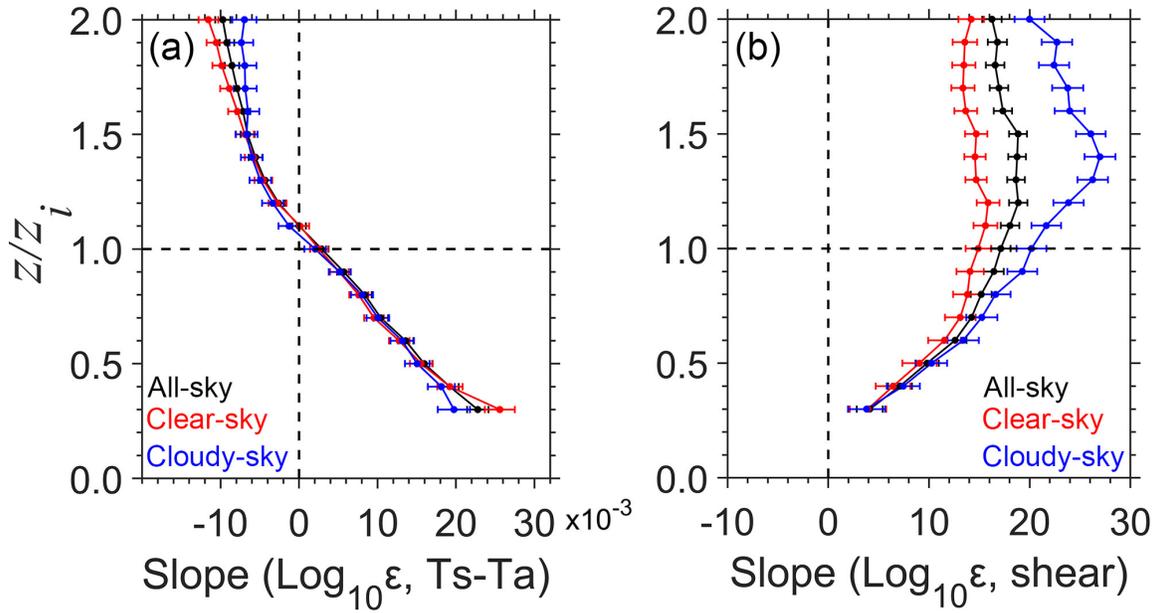
909
 910
 911
 912
 913
 914
 915
 916
 917
 918
 919
 920

Figure 6. Scatter plots (blue dots) of $\text{Log}_{10}\epsilon$ estimated from the measurements of RWPs in the whole lower troposphere ($0.3 \leq z/z_i \leq 2.0$, a, d, g), in the PBL ($0.3 \leq z/z_i \leq 1.0$, b, e, h) and above the PBL ($1.0 < z/z_i \leq 2.0$, c, f, i) over the TP as a function of $T_s - T_a$ under all-sky (a-c), clear-sky (d-f) and cloudy-sky conditions (g-i), respectively. Also overlaid are their corresponding box and whisker plots and regression linear equations and correlation coefficients in each panel, where all $T_s - T_a$ samples are divided into twenty bins, each of which has the same sample size. Note that the median is shown as a line whereas the outer boundaries of the boxes represent 25 and 75 quartiles and the dashed lines present interquartile range (IQR). The superscripts * and ** for R indicate that the regression slopes are statistically significant at $p < 0.05$ and $p < 0.01$ levels, respectively.



921
 922
 923
 924
 925
 926
 927
 928
 929
 930

Figure 7. Scatter plots of $\text{Log}_{10} \varepsilon$ estimated from the measurements of RWPs in the whole lower troposphere ($0.3 \leq z/z_i \leq 2.0$, a, d, g), in the PBL ($0.3 \leq z/z_i \leq 1.0$, b, e, h) and above the PBL ($1.0 < z/z_i \leq 2.0$, c, f, i) over the TP as a function of vertical wind shear (VWS) under all-sky (a-c), clear-sky (d-f) and cloudy-sky conditions (g-i), respectively. Also overlaid are their corresponding box and whisker plots and fitting equations in each panel, where all VWS samples are divided into twenty bins, each of which has the same sample size. Note that the median is shown as a line whereas the outer boundaries of the boxes represent 25 and 75 quartiles and the dashed lines present interquartile range (IQR).

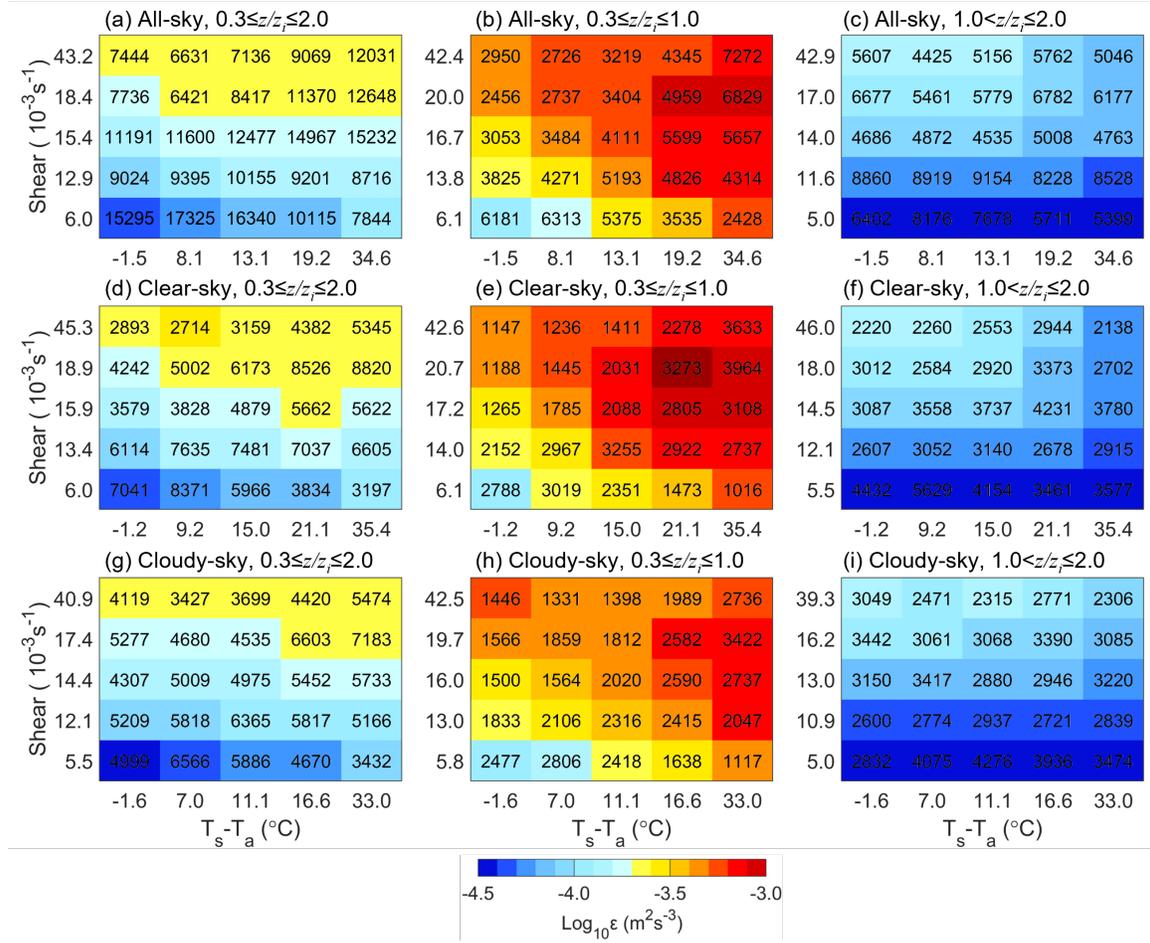


931

932 **Figure 8.** The vertical profiles of least squares regression slope between $\text{Log}_{10}\epsilon$ and

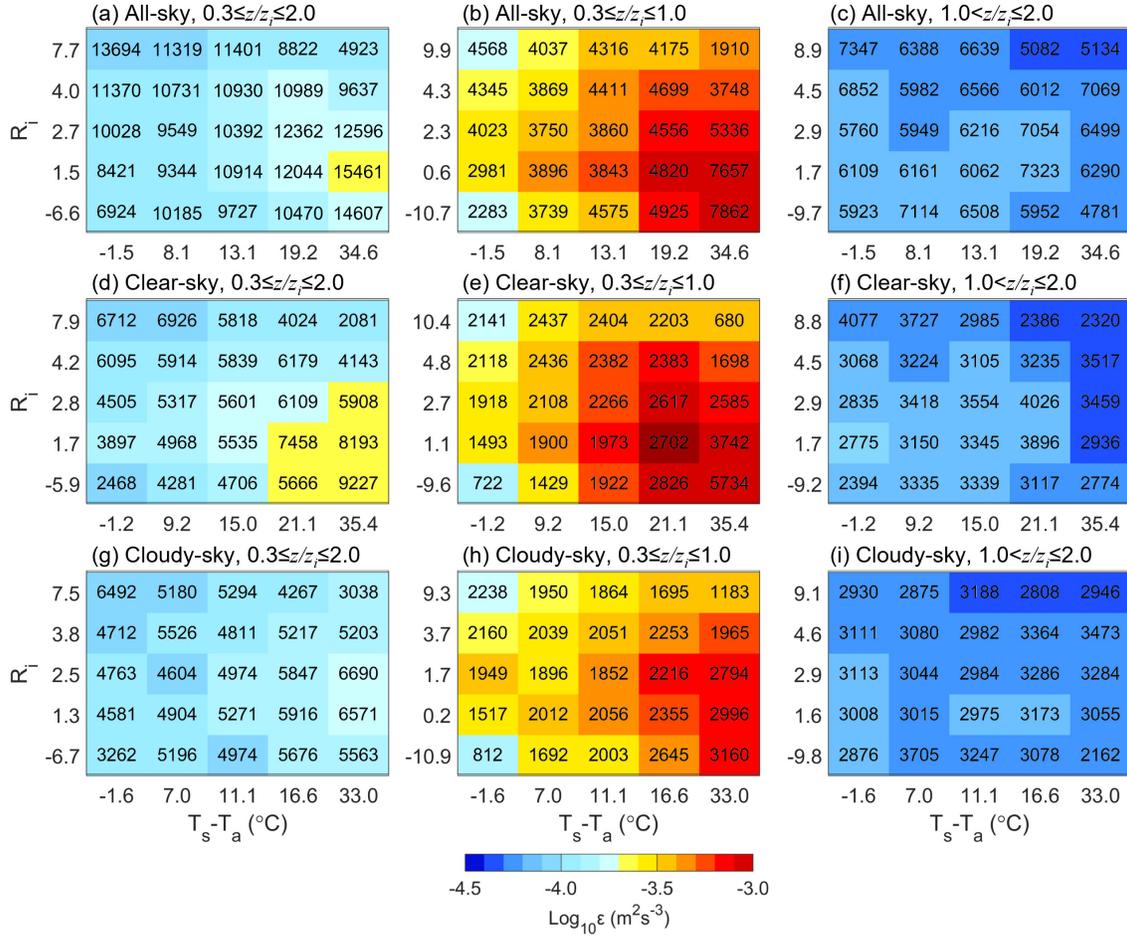
933 $T_s - T_a$ (a) and VWS (b) over the TP under all-sky (black), clear-sky (red) and cloudy-sky

934 (blue) conditions, respectively.



935

936 **Figure 9.** Joint dependence of $\text{Log}_{10}\epsilon$ (color shading) on the VWS and $T_s - T_a$ within
 937 and above the PBL (a, d, g), in the PBL (b, e, h) and above the PBL (c, f, i) over the TP
 938 under all-sky (a-c), clear-sky (d-f) and cloudy-sky (g-i) conditions, respectively. The
 939 number given in each panel is the total number of samples used.



940

941 **Figure 10.** Joint dependence of $\text{Log}_{10}\epsilon$ (color shading) on the gradient Richardson
 942 number (Ri) and $T_s - T_a$ in and above the PBL (a, d, g), in the PBL (b, e, h) and above
 943 the PBL (c, f, i) over the TP under all-sky (a-c), clear-sky (d-f) and cloudy-sky (g-i)
 944 conditions, respectively. The number given in each panel is the total number of samples
 945 used.