



# Surface buoyancy control of millennial-scale variations of the Atlantic meridional ocean circulation

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## Abstract.

Dansgaard-Oeschger (DO) events are a pervasive feature of glacial climates. It is widely accepted that the associated changes in climate, which are most pronounced in the North Atlantic region, are caused by abrupt changes in the strength and/or latitude reach of the Atlantic meridional overturning circulation (AMOC), possibly originating from spontaneous transitions in the ocean-sea-ice-atmosphere system. Here we use an Earth System Model that produces DO-like events to show that the climate conditions under which millennial-scale AMOC variations occur are controlled by the surface ocean buoyancy flux. In particular, we find that the present day-like convection pattern with deep water formation in the Labrador and Nordic Seas becomes unstable when the buoyancy flux integrated over the northern North Atlantic turns from negative to positive. It is in the proximity of this point that the model produces transitions between different convection patterns associated with strong and weak AMOC states. The buoyancy flux depends on the surface freshwater and heat fluxes and on sea surface temperature through the temperature dependence of the thermal expansion coefficient of seawater. We find that larger ice sheets tend to stabilize convection by decreasing the net freshwater flux while CO<sub>2</sub>-induced cooling decreases buoyancy loss and destabilizes convection. These results help to explain the conditions under which DO events appear, and are a step towards an improved understanding of the mechanisms of abrupt climate changes.

## 15 1 Introduction

Most of the Quaternary, except for interglacial states and full glacial states, is characterized by numerous abrupt climate change events (Dansgaard-Oeschger or DO events) (Dansgaard et al., 1993; Rahmstorf, 2002) that are most pronounced in the North Atlantic realm (Hodell et al., 2023; Hoff et al., 2016; Bond et al., 1993) and clearly seen in Greenland ice core data (Andersen et al., 2004), which suggest annual air temperature variations of 6–16°C (Kindler et al., 2014). Atmospheric CO<sub>2</sub> concentration (Zhang et al., 2021) as well as global ice volume and orbital parameters (Mitsui and Crucifix, 2017; Lohmann and Ditlevsen, 2018) all affect the occurrence and characteristics of DO events. All DO events have a similar time evolution, with an abrupt warming followed by a slow cooling and then a rather abrupt return to stadial (cold) conditions, but their duration varies from several hundred to several thousand years (Lohmann and Ditlevsen, 2019).



It is now generally accepted that the temporal and spatial dynamics of DO events can be explained by abrupt transitions between two modes of the Atlantic Meridional Overturning circulation (or AMOC) (Ganopolski and Rahmstorf, 2001), consistent with paleoclimate records (Henry et al., 2016; Bohm et al., 2015; Keigwin and Boyle, 1999; Skinner and Elderfield, 2007). Noise has been argued to be important in triggering the transition between the different AMOC states, either by amplifying the response to a small periodic forcing (Braun et al., 2005; Rahmstorf, 2003) through stochastic resonance (Ganopolski and Rahmstorf, 2002; Alley et al., 2001; Vélez-Belchí et al., 2001), or by exciting and reorganizing internal modes of variability through the coherence resonance mechanism (Timmermann et al., 2003; Pikovsky and Kurths, 1997; Ditlevsen et al., 2005). DO events could therefore originate from internal oscillations within the climate system (Li and Born, 2019; Menviel et al., 2020; Boers et al., 2018), and an increasing number of simplified and general circulation models (GCMs) have revealed spontaneous oscillations resembling DO events with a typical periodicity of 1000 years under a wide range of glacial boundary conditions (Malmierca-Vallet and Sime, 2023). Some GCMs produce internal oscillations for present-day ice sheets but low CO<sub>2</sub> (Brown and Galbraith, 2016; Klockmann et al., 2018), others under full glacial conditions (Peltier and Vettoretti, 2014; Romé et al., 2022; Prange et al., 2023) or some combination of mid-glacial (MIS3) conditions (Armstrong et al., 2022; Kuniyoshi et al., 2022; Zhang et al., 2021; Vettoretti et al., 2022) in terms of ice sheets, atmospheric CO<sub>2</sub> and orbital parameters.

However, while the concept of a ‘sweet spot’ for the occurrence of DO-like variability has recently gained considerable attention, what physical conditions control where it is located in the ice sheet–CO<sub>2</sub>–orbit space in the different models has remained largely unexplained. Clarifying this is important in order to understand why in reality the DO events occurred under a broad range of glacial climate and boundary conditions, but not during interglacials and peak glacial conditions. Here we use a large number of simulations with an Earth system model to address this question.

## 2 Methods

### 2.1 Earth system model

We use the CLIMBER-X (Willeit et al., 2022) Earth system model in a climate-only setup, including a frictional-geostrophic 3D ocean model, a semi-empirical statistical-dynamical atmosphere model, a dynamic-thermodynamic sea ice model and a land surface model with interactive vegetation. All components of the climate model have a horizontal resolution of 5°x5°. The model is described in detail in Willeit et al. (2022) and in general shows performances that are comparable with state-of-the-art CMIP6 models under different forcings and boundary conditions. In particular, the simulated present-day AMOC overturning profile at 26°N in the Atlantic agrees well with observations and is within the range produced by CMIP6 models (Willeit et al., 2022).

### 2.2 Model experiments

To explore the impact of different aspects of climate on AMOC strength and variability we run an ensemble of climate model simulations using CLIMBER-X with different prescribed ice sheet configurations and atmospheric CO<sub>2</sub> concentrations. We



55 performed equilibrium simulations for three different ice sheet configurations, representative of interglacial conditions (present-day, Fig. 1a), full glacial conditions (Last Glacial Maximum GLAC-1D reconstruction (Tarasov et al., 2012), Fig. 1e) and mid-glacial conditions (GLAC-1D reconstruction (Tarasov et al., 2012) for 12 ka, Fig. 1c). The 12 ka GLAC-1D ice sheet reconstructions is similar to the 35 ka ice sheets from PaleoMIST (Gowan et al., 2021), which has been suggested as boundary condition for a DO intercomparison project (Malmierca-Vallet and Sime, 2023), but with a slightly larger Fennoscandian ice sheet. For all three ice sheet configurations we run a set of experiments for 8000 years with prescribed constant atmospheric CO<sub>2</sub> concentrations ranging from a pre-industrial value of 280 ppm to 150 ppm, at steps of 10 ppm. The first 3000 years of the simulations are treated as spinup and the following 5000 years are used in the analysis. To isolate the effects of ice sheets and CO<sub>2</sub>, in all simulations we use present-day orbital parameters.

65 Since we are not changing the concentration of other greenhouse gases (GHGs) such as CH<sub>4</sub> and N<sub>2</sub>O, these CO<sub>2</sub> concentrations should be considered in terms of a CO<sub>2</sub> equivalent that implicitly includes the radiative effect of other GHG changes relative to pre-industrial. The CO<sub>2</sub> equivalent is roughly ~20 ppm lower than the actual CO<sub>2</sub> concentration for most of the last glacial cycle (Appendix A), with minimum values of ~160 ppm during glacial maxima (Fig. A2).

70 Ice sheets and bedrock topography are prescribed in the simulations, resulting in different surface elevation, bathymetry and land-sea mask for each different ice sheet configuration. The runoff routing directions are automatically derived using a steepest surface gradient approach. In order to conserve water in the climate system, the melting of land ice is inhibited and snow accumulating over the ice sheets is cut off at a maximum thickness of 4 m with the excess being routed as ‘frozen water’ to the ocean, where the latent heat of fusion is also accounted for.

75 Since the model does not explicitly resolve synoptic-scale and interannual variability in the atmosphere and ocean, to mimic the effect of weather on annual mean AMOC strength we applied perturbations to the surface ocean freshwater flux in the North Atlantic in the latitudinal belt between 50–80°N. We apply a Gaussian white noise with a standard deviation of 0.5 kg m<sup>-2</sup> day<sup>-1</sup> uniformly over the area and constant over each year. See Appendix B for further details. For mid-glacial ice sheets we also run the simulations with different amplitudes of the noise in the freshwater flux in the North Atlantic, with standard deviations of 0, 0.0625, 0.125, 0.25 and 1 kg m<sup>-2</sup> day<sup>-1</sup>, in addition to the reference simulations with 0.5 kg m<sup>-2</sup> day<sup>-1</sup>. In all experiments the model is initialized from a 10,000 years long pre-industrial equilibrium spinup with 80 280 ppm of atmospheric CO<sub>2</sub> and present-day ice sheets.

85 For each ice sheet configuration we additionally performed transient simulations with slowly varying CO<sub>2</sub> concentrations: (i) starting at 280 ppm and gradually decreasing CO<sub>2</sub> down to 150 ppm and (ii) starting from 150 ppm and gradually increasing CO<sub>2</sub> up to 280 ppm. In both cases the rate of change of CO<sub>2</sub> is 3 ppm kyr<sup>-1</sup> implying a total simulation length of ~43,000 years. The initial state for these simulations is an equilibrium with either 280 or 150 ppm of atmospheric CO<sub>2</sub>. Simulations (i) and (ii) were performed both with no noise in the freshwater flux in the North Atlantic and with the reference noise of 0.5 kg m<sup>-2</sup> day<sup>-1</sup> and for different diapycnal diffusivities in the upper ocean (Appendix C and Fig. C1).

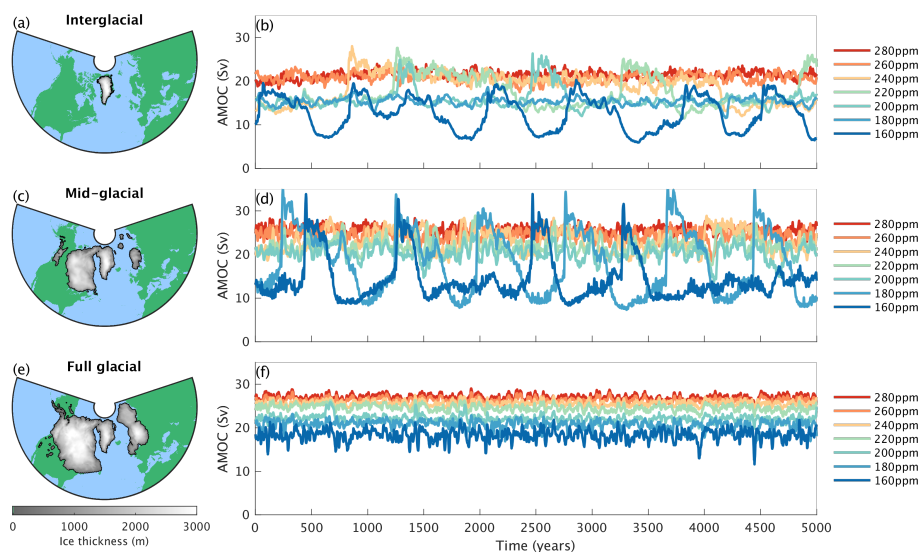
An additional experiment with freshwater hosing with slowly increasing the freshwater flux from -0.2 and 0.2 Sv in the North Atlantic (50-70°N) was performed to explore when convective instability is triggered by the addition of freshwater. This experiment was run only without noise and for present-day ice sheets.



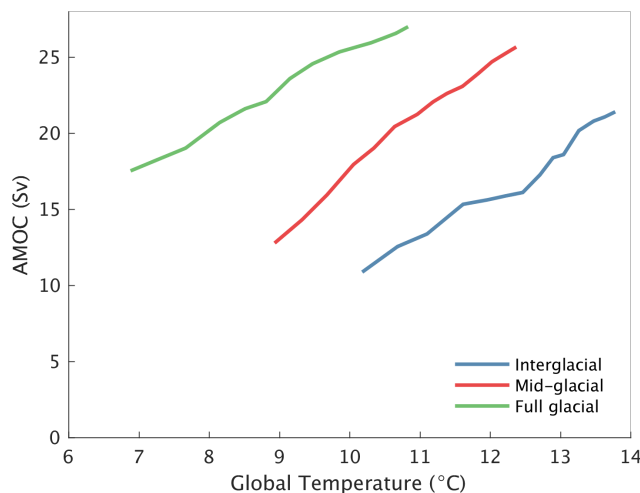
## 90 3 Results

### 3.1 The effect of ice sheets and CO<sub>2</sub> on AMOC

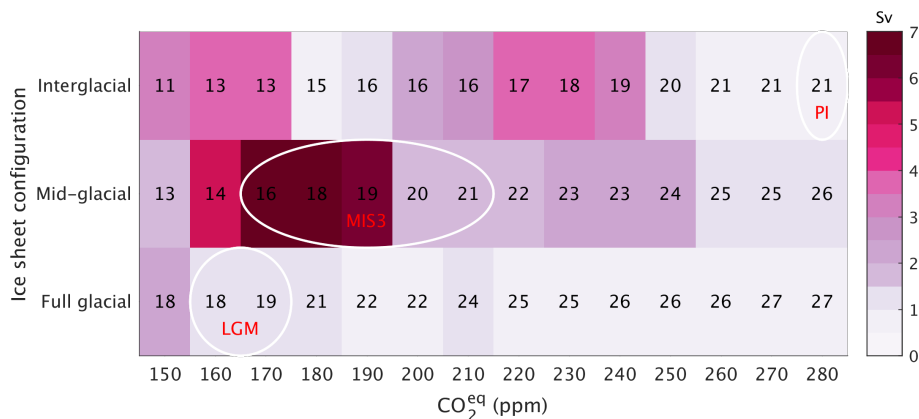
For all three ice sheet configurations, when averaged over 5000 years, the AMOC tends to be stronger for higher CO<sub>2</sub> concentrations and weakens as CO<sub>2</sub> decreases (Fig. 1 and Fig. 2), in line with other models (Brown and Galbraith, 2016; Galbraith and de Lavergne, 2019; Oka et al., 2012; Klockmann et al., 2018; Stouffer and Manabe, 2003; Prange et al., 2023). Moreover, for a given CO<sub>2</sub> concentration, the AMOC is generally stronger in experiments with more extensive land ice cover (Fig. 1, Fig. 2 and Fig. 3), again in agreement with other models (Klockmann et al., 2018; Brown and Galbraith, 2016). Therefore, larger ice sheets act to strengthen the AMOC in the model, while lower CO<sub>2</sub> concentrations weaken it. This does not contradict to the expected AMOC weakening due to anthropogenic CO<sub>2</sub> increase, because the latter is an inherently transient phenomenon (Bonan et al., 2022; Stouffer and Manabe, 2003). The net effect of these two counteracting factors leads to a shallower and somewhat weaker AMOC being simulated under full glacial conditions compared to the pre-industrial (Fig. 4i,j). This is in agreement with paleoclimate data (McManus et al., 2004; Bohm et al., 2015; Pöppelmeier et al., 2023), in contrast to most PMIP models which show a tendency towards a deeper and stronger AMOC at the last glacial maximum compared to the present (Kageyama et al., 2021).



**Figure 1.** AMOC response to ice sheets and CO<sub>2</sub>. Maximum strength of the Atlantic meridional overturning circulation in model simulations (b,d,f) for the three different ice sheet configurations in (a,c,e) and different prescribed constant equivalent atmospheric CO<sub>2</sub> concentrations as shown in the legends on the right.



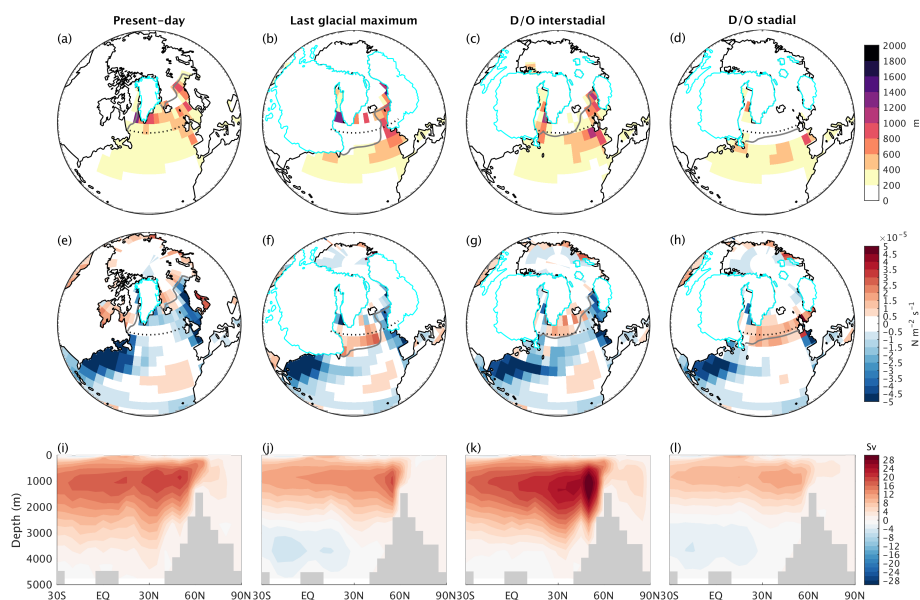
**Figure 2.** AMOC strength as a function of global temperature for the three different ice sheet configurations derived from the ensemble of model simulations with CO<sub>2</sub> equivalent ranging between 150 and 280 ppm.



**Figure 3.** AMOC mean and variability. Interannual standard deviation of AMOC time series (shading) and mean AMOC strength (numbers) for different combinations of ice sheet configurations and CO<sub>2</sub> concentrations. Conditions representative for the pre-industrial (PI), MIS3 and LGM conditions are indicated.

### 3.2 Internal AMOC variability

105 Under some combinations of ice sheet configuration and CO<sub>2</sub> concentration, the model produces spontaneous oscillations of the AMOC (Fig. 1 and Fig. 3). This is in line with a growing number of models of different complexity producing self-sustained oscillations for specific combinations of ice sheets, CO<sub>2</sub> and orbital parameters (Brown and Galbraith, 2016; Klockmann et al., 2018; Romé et al., 2022; Peltier and Vettoretti, 2014; Kuniyoshi et al., 2022; Zhang et al., 2021; Armstrong et al., 2022;

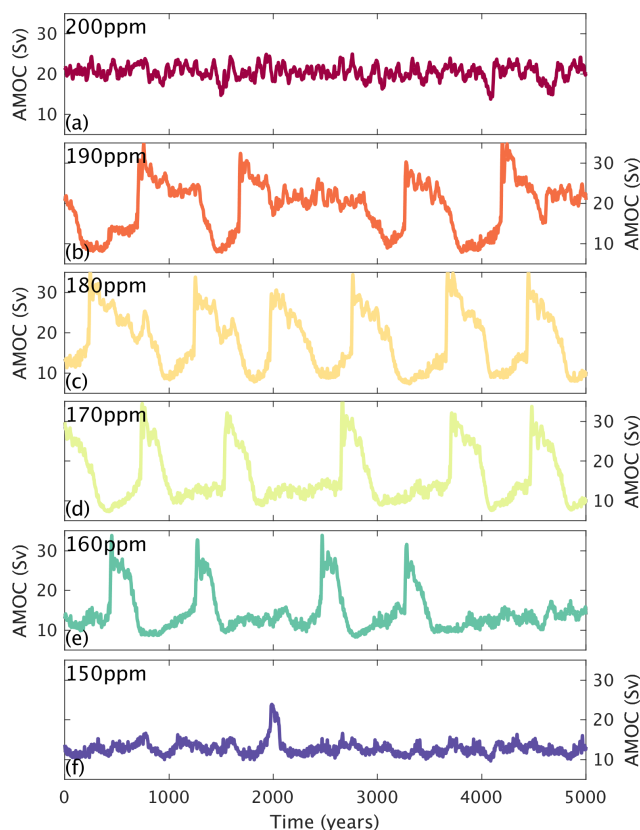


**Figure 4.** Maximum monthly mixed layer depth (a-d), annual mean surface ocean buoyancy flux (e-h) and Atlantic meridional overturning streamfunction (i-l) for different conditions: (left to right) present-day (year 2000 CE) from a transient historical simulation, last glacial maximum (5000 years simulation with LGM boundary conditions following the PMIP protocol (Kageyama et al., 2017) with GLAC-1D ice sheets), interstadial and stadial conditions from the simulation with mid-glacial ice sheets and a  $\text{CO}_2$  concentration of 180 ppm. The grey line in a-h shows the maximum sea ice extent. The cyan contours indicate the ice sheet extent and the dotted line in panels a-h marks the  $55^\circ\text{N}$  latitude.

Friedrich et al., 2010; Schulz et al., 2007; Sakai and Peltier, 1997). The appearance of internal AMOC oscillations within a window of  $\text{CO}_2$  concentrations is fully consistent with recent results from a general circulation model (Vettoretti et al., 2022).

The amplitude and period of the oscillations depends on the boundary conditions. Notably, for the mid-glacial ice sheet configuration and  $\text{CO}_2$  between 150 and 200 ppm, the model produces large amplitude AMOC oscillations that qualitatively resemble DO events recorded in Greenland ice core data both in terms of shape and millennial-scale periodicity (Fig. 1d and Fig. 5). The sea-saw pattern between Greenland and Antarctic temperatures is also well reproduced, with Antarctic temperature maxima lagging Greenland temperature maxima by  $\sim 100$  years (Fig. 6b), in agreement with observations (Svensson et al., 2020). The amplitude of the simulated temperature variations over Greenland is underestimated compared to ice core reconstructions, but sea surface temperatures at the Iberian margin vary by  $\sim 1.5^\circ\text{C}$  between stadial and interstadial (Fig. 6c), in very good agreement with proxy records (Martrat et al., 2007).

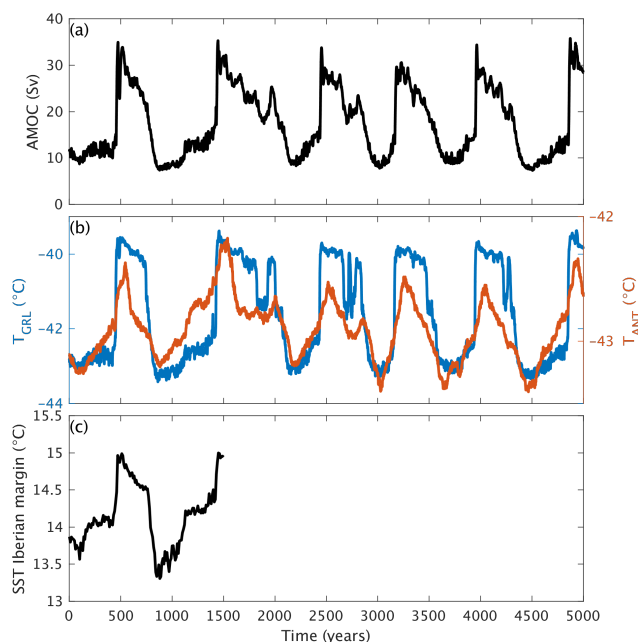
During the stadial phase deep water is formed south of  $55^\circ\text{N}$  (Fig. 4d). This is consistent with the ice core data showing no significant temperature difference over Greenland between stadials with or without Heinrich events, which indicates that the stadial AMOC does not reach far enough north to warm Greenland. The stadial AMOC is also weaker and shallower than at present (Fig. 4i,l), while during the interstadial phase deep convection occurs in the Labrador Sea and the Nordic Seas (Fig. 4c)



**Figure 5.** Time series of maximum strength of the AMOC streamfunction in simulations with mid-glacial ice sheets and different constant atmospheric CO<sub>2</sub> concentrations, decreasing from top to bottom.

resulting in an AMOC that is stronger than at present (Fig. 4k). The heat transport reaching further north during interstadials compared to stadials is the main reason for the warming over Greenland. The northward shift of the deep water formation sites during the interstadial and the associated retreat of sea ice results in annual temperatures up to 15 °C warmer in the Nordic Seas than during stadial conditions (Fig. 7). The simulated change in sea ice cover in the Nordic Seas (Fig. 4c,d and Fig. 7) is in good agreement with proxy-based estimates showing extensive winter sea ice cover over the area during stadials and ice-free conditions during interstadials (Sadatzki et al., 2019, 2020; Hoff et al., 2016; Dokken et al., 2013). A shift from perennial sea ice during stadials to seasonal sea ice during interstadials is simulated in the Labrador Sea and the Baffin Bay (Fig. 7) and is consistent with reconstructions (Scoto et al., 2022).

For present-day ice sheets, millennial-scale oscillations are simulated for two different CO<sub>2</sub> ranges around 230 ppm and 160 ppm (Fig. 1b and Fig. 3). In both cases the AMOC variations are of lower amplitude than for mid-glacial ice sheets. No oscillations are produced by the model for full glacial ice sheets for any of the CO<sub>2</sub> concentrations considered (Fig. 1f).



**Figure 6.** (a) Maximum value of the AMOC streamfunction in model simulation with mid-glacial ice sheets and 180 ppm of atmospheric CO<sub>2</sub>. (b) Corresponding Greenland and Antarctic temperature evolution. Note the different y-axes range. (c) Simulated annual mean sea surface temperature at the Iberian margin. The Iberian margin temperature is available only for a part of the simulation that has frequent 2D SST output.

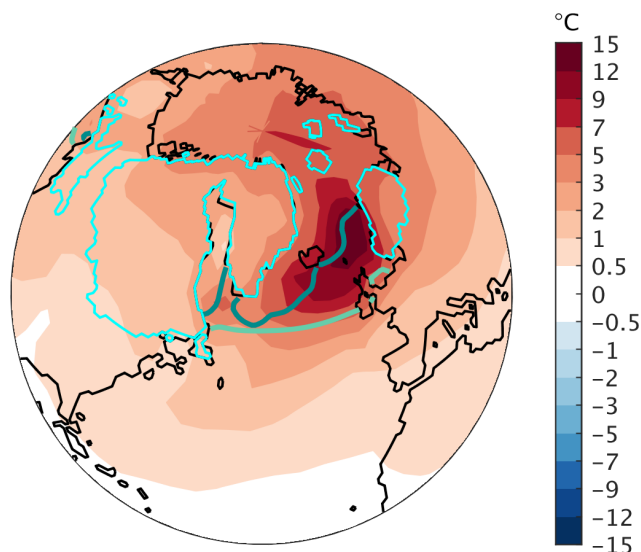
Overall, these results are in qualitative agreement with ice core data showing pronounced millennial-scale climate variability in the North Atlantic during intermediate glacial conditions (e.g. MIS3), but not during peak glacial conditions, such as the LGM, or during interglacials, such as the Holocene.

The range of boundary conditions under which the oscillations occur depends on the amplitude of the noise that is applied to the surface freshwater flux in the North Atlantic (Fig. 8), with larger amplitude of noise acting to broaden the range. For mid-glacial ice sheets, self-sustained oscillations are simulated even without noise (Fig. 9). Larger noise levels generally reduce the duration of the stadials, similarly to results obtained using a conceptual coupled climate model (Timmermann et al., 2003). While we find that in our model oscillations can be produced even without stochastic forcing, noise in the form of interannual variability intrinsic in the climate system could have played an important role in establishing the robust millennial-scale climate variability observed during the Quaternary. Our model has the advantage that it enables a separate investigation of the role of noise on DO dynamics, which cannot be done with GCMs resolving synoptic processes.

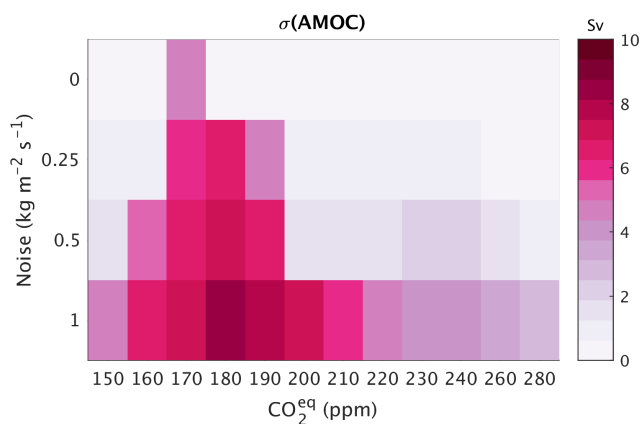
### 3.3 Convective instability and surface buoyancy flux

Transient simulations with slowly decreasing atmospheric CO<sub>2</sub> from 280 to 150 ppm help to elucidate the conditions under which AMOC oscillations occur in the model (Fig. 10). When these experiments are performed with noise in the surface





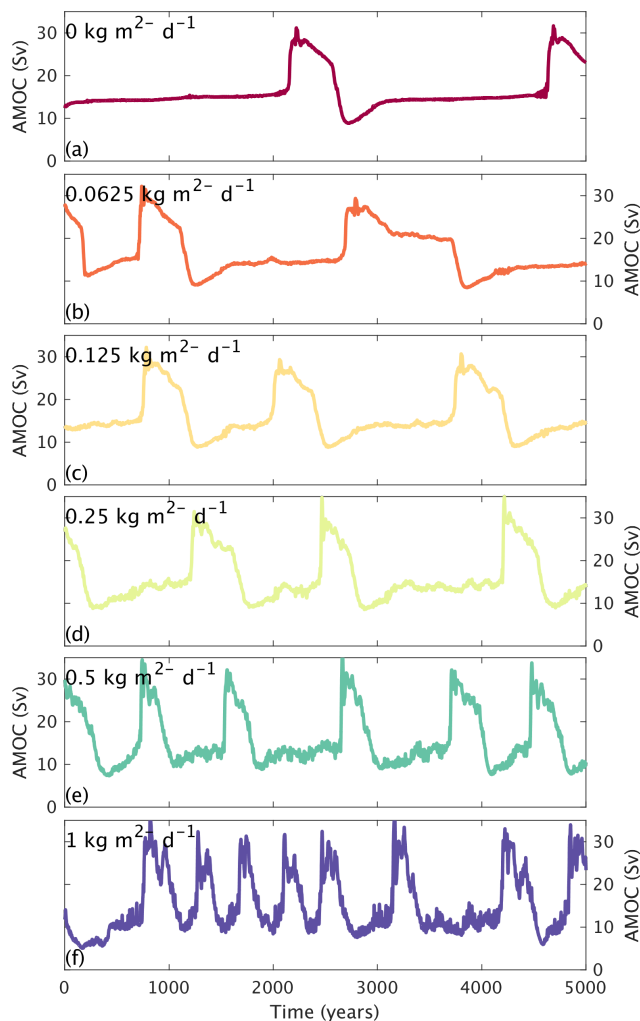
**Figure 7.** Difference in simulated annual mean near-surface air temperature between interstadial and stadial conditions. The annual mean sea ice extent for the stadial (green) and interstadial (grey) is also shown.



**Figure 8.** Interannual standard deviation of AMOC time series for mid-glacial ice sheets and different combinations of equivalent CO<sub>2</sub> concentration and amplitude of the noise applied to the freshwater flux into the North Atlantic.

freshwater flux in the North Atlantic, similar behavior is obtained as in the equilibrium simulations shown in Fig. 1. However, performing the same simulations without the imposed noise highlights the presence of discrete transitions in the AMOC at several critical CO<sub>2</sub> concentrations, which are specific to each ice sheet configuration. For a range of CO<sub>2</sub> values around these points the model shows internal oscillations when noise is applied.

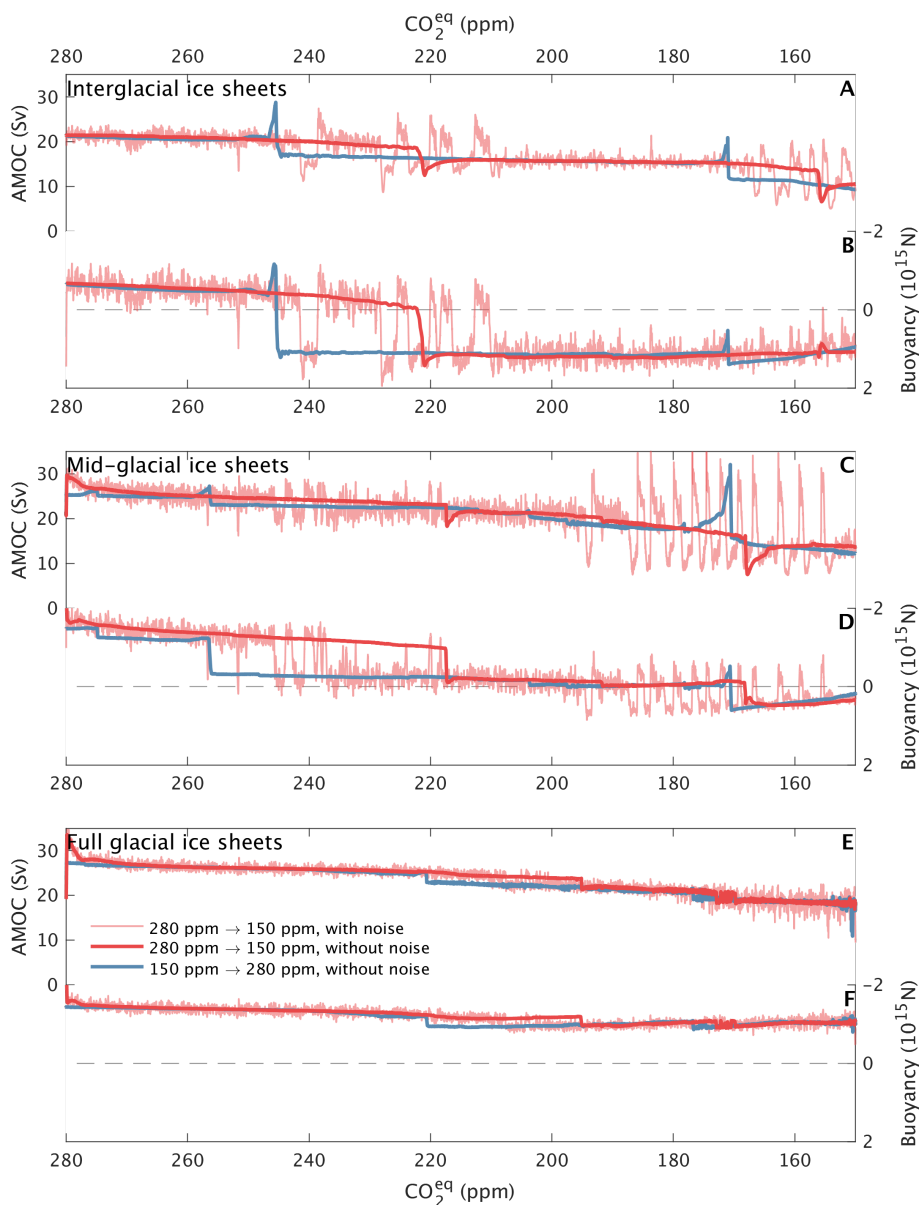
The different AMOC states are connected to the presence of qualitatively different stable convection patterns in the North Atlantic for different CO<sub>2</sub> values (Fig. 11). Under some conditions, e.g. for present day ice sheets and CO<sub>2</sub> between 220 and



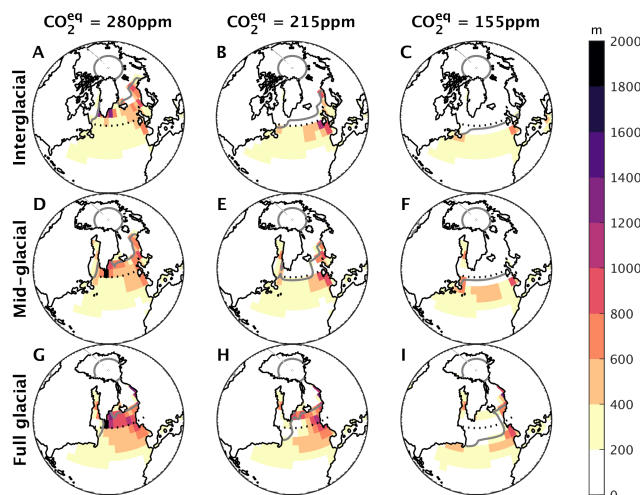
**Figure 9.** AMOC time series for mid-glacial ice sheets and different amplitudes of the noise applied to the freshwater flux into the North Atlantic, increasing from top to bottom as indicated in the panels.

240 ppm, two stable modes of the AMOC, which correspond to different convective patterns, are stable under the same CO<sub>2</sub> (Fig. 10a), but this is not a pre-requisite for the occurrence of internal oscillations.

The stability of the convection patterns, and therefore the AMOC transitions, can be directly linked to the annual mean surface ocean buoyancy flux integrated over the northern North Atlantic,  $M$  (Appendix D). The role of surface buoyancy for AMOC stability has previously been considered, but rather separately for the different deep water formation regions (Klockmann et al., 2018) and not in terms of an integral value over the whole northern North Atlantic. Our results show that a present day-like spatial organization of convection with deep water forming in the Labrador and Nordic Seas cannot be sustained if



**Figure 10.** Maximum AMOC and integrated surface buoyancy flux,  $M$ , for experiments with different ice sheet configurations (interglacial, mid-glacial and full glacial, from top to bottom) and slowly varying atmospheric  $\text{CO}_2$  concentration. The red lines are for simulations with a gradual  $\text{CO}_2$  decrease starting from pre-industrial conditions with 280 ppm, while the blue line is for a gradual increase of  $\text{CO}_2$  starting from 150 ppm. The thick lines are experiments without noise in the freshwater flux in the north Atlantic, while the thin line is from an experiment with noise, which is shown only for the  $\text{CO}_2$  decrease case. The imposed rate of change in  $\text{CO}_2$  is  $3\text{ppm kyr}^{-1}$  so that the total length of the simulations is  $\sim 43,000$  years.  $M$  is smoothed with a running-mean of 30 years.



**Figure 11.** Maximum of the monthly mean mixed layer depth in quasi-equilibrium conditions with different ice sheet configurations and  $\text{CO}_2$  concentrations from the simulations with slowly decreasing  $\text{CO}_2$  concentration and without noise in the freshwater flux (corresponding to the thick red lines in Fig. 10). The grey lines indicate the maximum sea ice extent.

$M$  integrated north of  $\sim 55^\circ\text{N}$  transitions from negative to positive (Fig. 10b,d), leading to a sudden shift in the deep-water formation to latitudes south of  $55^\circ\text{N}$  (Fig. 11a,b and Fig. 11e,f), with an associated weakening of the AMOC (Fig. 10a,c).

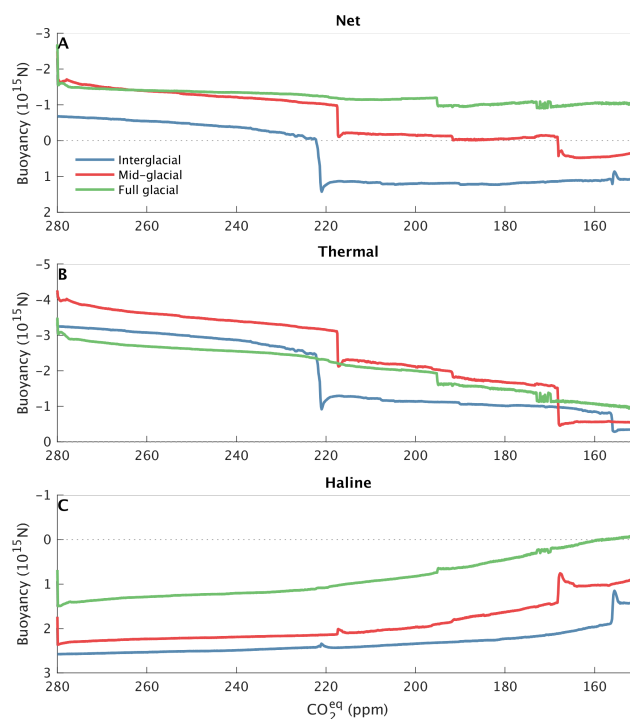
When noise is applied, the transitions between strong and weak AMOC states are associated with a change in the sign of the buoyancy measure  $M$  (Fig. 10b,d). During interstadials,  $M$  decreases in magnitude as the heat accumulated in the sub-surface during the stadial conditions is gradually released. The interstadial can only be sustained as long as  $M$  is negative, after which there is a rapid transition back to stadial conditions (Fig. 10a-d).

The surface buoyancy flux is tightly linked to the convection pattern, with pronounced buoyancy loss concentrated over areas of deep-water formation along the margins of perennial sea ice cover (Fig. 4e-h).

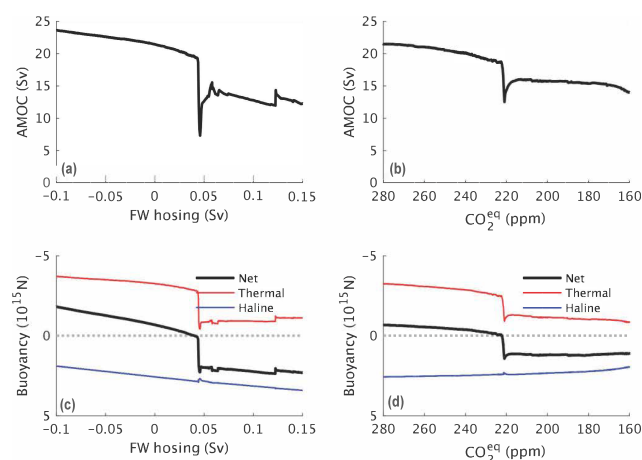
The response of  $M$  to decreasing  $\text{CO}_2$  shown in Fig. 10 results from a combination of factors: (i) the decrease of the thermal expansion coefficient of seawater with decreasing temperatures decreases the absolute value of the negative thermal component of the surface buoyancy flux and thus increases  $M$ , (ii) a weaker AMOC with an associated smaller meridional heat transport decreases the surface sensible heat loss to the atmosphere and therefore also increases  $M$ , while (iii) a decrease in net freshwater flux due to a weakening of the hydrological cycle acts to decrease  $M$  (Fig. 12 and Fig. 13). The net effect is that the buoyancy loss decreases as  $\text{CO}_2$  is gradually reduced (Fig. 13d and Fig. 12).

It is noteworthy that for mid-glacial ice sheets  $M$  is close to zero over a wide range of  $\text{CO}_2$  concentrations (Fig. 10d). This is the result of partially compensating effects of the thermal and haline components of the buoyancy flux (Fig. 12) and could help explain the observed ubiquitous appearance of DO events under mid-glacial conditions.

The presence of ice sheets does directly affect the surface buoyancy flux in the northern North Atlantic, with larger ice sheets resulting in increased buoyancy loss and consequently stronger AMOC for any given  $\text{CO}_2$  concentration (Fig. 10,



**Figure 12.** (a) Net buoyancy flux integrated north of  $55^{\circ}\text{N}$  in the Atlantic ( $M$ ) as a function of atmospheric  $\text{CO}_2$  concentration for different ice sheet configurations, with separation into (b) thermal and (c) haline components.

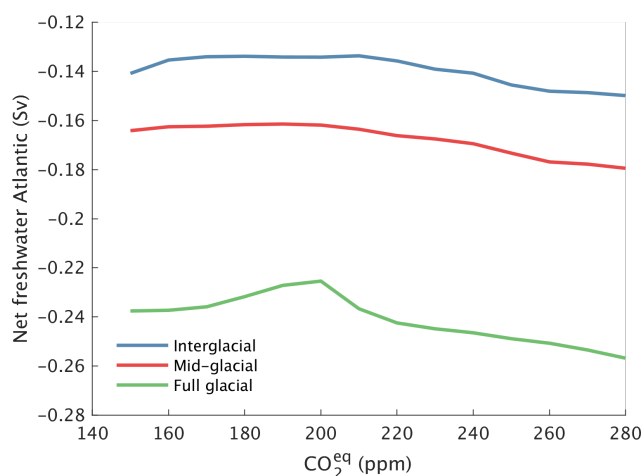


**Figure 13.** (a,b) AMOC strength and (c,d) integrated surface buoyancy flux,  $M$ , for model simulations with present-day ice sheets and (a,c) slowly increasing freshwater hosing applied to the northern North Atlantic ( $50\text{--}70^{\circ}\text{N}$ ) and (b,d) gradual atmospheric  $\text{CO}_2$  decrease. In the lower panels  $M$  is further separated into thermal and haline components.



180 Fig. 3). In our model, the presence of large Northern Hemisphere ice sheets reduces the net surface ocean freshwater flux into the Atlantic (Fig. 14). Prescribing LGM ice sheets leads to a decrease in the net Atlantic freshwater flux by  $\sim 0.1$  Sv compared to experiments with present-day ice sheets, almost independently of the  $\text{CO}_2$  concentration (Fig. 14), and is a result of the Laurentide ice sheet effectively blocking part of the Pacific-to-Atlantic atmospheric moisture transport. Most of the reduction in freshwater flux occurs in the northern North Atlantic, thereby decreasing the surface buoyancy flux in the

185 deep-water formation regions. CLIMBER-X and PMIP3/4 models show a generally similar  $M$  under pre-industrial and LGM conditions (Fig. 15). An increase in the thermal buoyancy component at the LGM relative to pre-industrial is by and large compensated by a decrease in the haline buoyancy component due to a reduced net surface freshwater flux at LGM compared to pre-industrial (Fig. 15).



**Figure 14.** Net surface freshwater flux into the Atlantic as a function of atmospheric  $\text{CO}_2$  concentration for different ice sheet configurations as indicated by the colored lines.

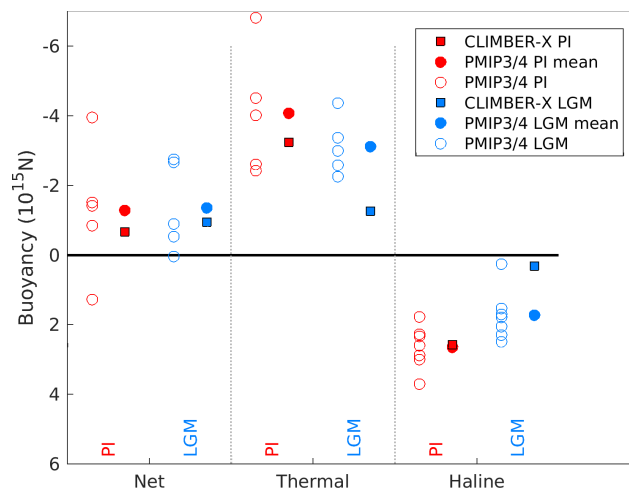
Convective instability can also be triggered by perturbing the surface freshwater balance, which directly affects  $M$  (Fig. 13a,c).

190 The noise that is applied to the surface freshwater flux in the model is thus also directly affecting the surface buoyancy flux and therefore facilitates the transition between different convection states. This also explains why larger noise amplitudes broaden the  $\text{CO}_2$  range over which oscillations are observed in the model (Fig. 8).

### 3.4 Sensitivity to model parameters

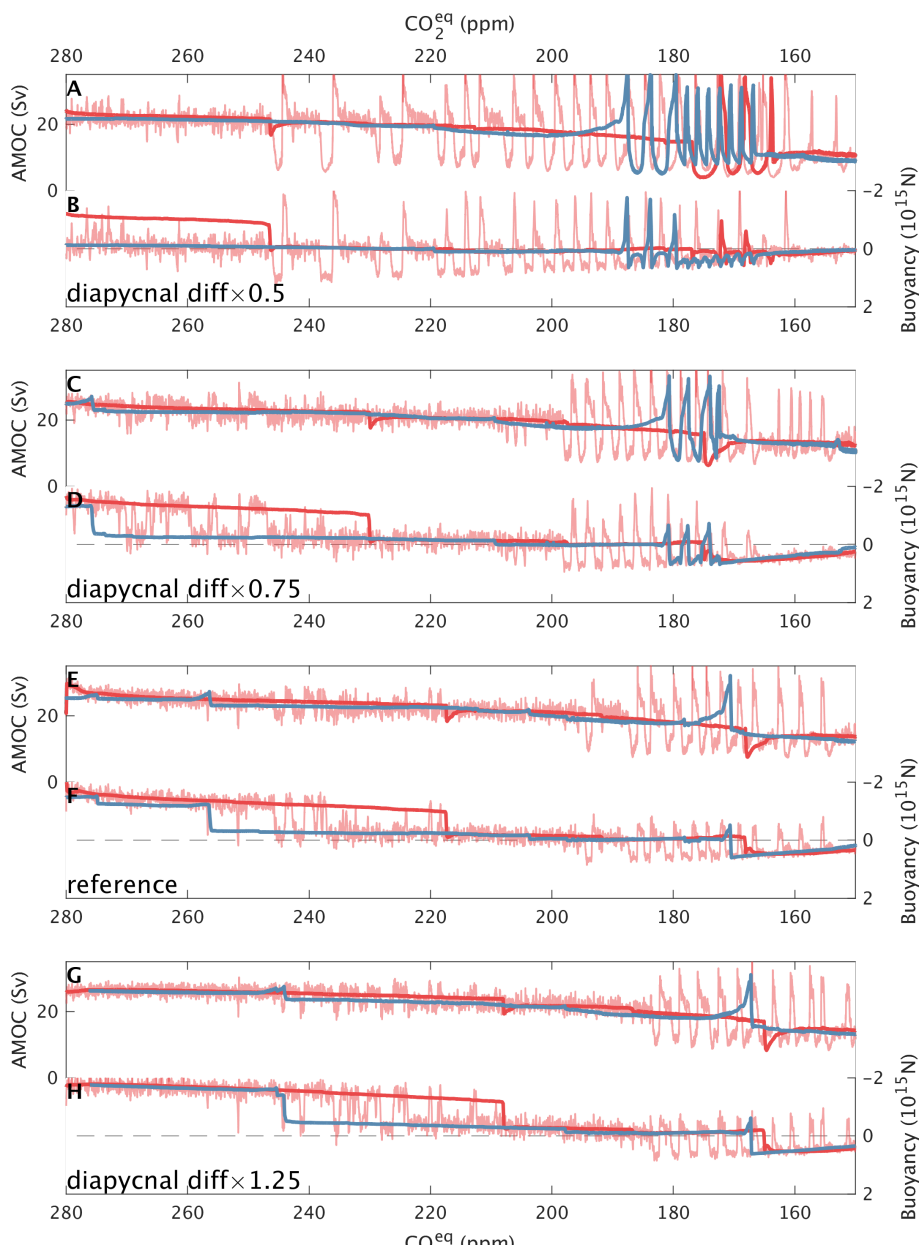
We have also tested the sensitivity of our results to model parameters, specifically the diapycnal diffusivity in the upper ocean

195 (Appendix C and Fig. C1). Larger diffusivities tend to make the AMOC stronger, therefore also increasing the northward heat transport with a consequent increase in surface sensible heat loss and a decrease in  $M$ . Consequently, decreasing diapycnal diffusivities leads to an increase in  $M$  and brings the system closer to convective instability. This is clearly seen in the simulated response to a slow  $\text{CO}_2$  decrease, where the critical thresholds for convective instability are systematically shifted to higher



**Figure 15.** Comparison of pre-industrial (PI) and LGM integrated surface buoyancy flux  $M$  into the Atlantic north of  $55^{\circ}\text{N}$  in CLIMBER-X and PMIP3/4 models. The separate contribution of the thermal and haline components of the buoyancy flux are also shown.

200  $\text{CO}_2$  values as diapycnal diffusivity decreases (Fig. 16). Moreover, smaller diapycnal diffusivities make the internal oscillations more robust, with pronounced oscillations simulated even in the absence of noise in the surface freshwater flux, and also extend the range of  $\text{CO}_2$  values over which millennial-scale variability is produced by the model (Fig. 16).



**Figure 16.** Maximum AMOC and integrated surface buoyancy flux,  $M$ , for experiments with mid-glacial ice sheets and slowly varying atmospheric  $\text{CO}_2$  concentration for different upper-ocean diapycnal diffusivities (Appendix C). The diffusivities are increasing from top to bottom and the values indicated in the panels specify the scaling factor applied to the value of diapycnal diffusivity at the surface. E and F are the same as panels C and D in Fig. 10. The red lines are for simulations with a gradual  $\text{CO}_2$  decrease starting from pre-industrial conditions with 280 ppm  $\text{CO}_2$ , while the blue line is for a gradual increase of  $\text{CO}_2$  starting from 150 ppm. The thick lines are experiments without noise in the freshwater flux in the north Atlantic, while the thin line is from an experiment with noise, which is shown only for the  $\text{CO}_2$  decrease case. The imposed rate of change in  $\text{CO}_2$  is  $3 \text{ ppm kyr}^{-1}$  so that the total length of the simulations is  $\sim 43,000$  years.





#### 4 Discussion and conclusions

The stability of the AMOC has historically often been considered in terms of advective instability (Stommel, 1961; Stocker and Wright, 1991; Wood et al., 2019; Rahmstorf, 1996; Hawkins et al., 2011; Hu et al., 2012; de Vries and Weber, 2005), which is controlled by the freshwater budget of the North Atlantic. However, the stability of deep water formation also plays a role in controlling the AMOC (Rahmstorf, 1994, 1995) and in Ganopolski and Rahmstorf (2001) it has been shown that DO events are actually best explained in terms of convective instability (Rahmstorf, 2001). Here we have shown that this convective instability is controlled by the annual mean integrated surface buoyancy flux in the northern North Atlantic.

With an *a-posteriori* knowledge of the geographical distribution of deep-water formation sites for different stable convection patterns, it is possible to define a latitude,  $\varphi_M$ , which spatially separates the convection sites corresponding to different AMOC modes. The surface buoyancy flux integrated north of  $\varphi_M$  ( $M$ ) is then a measure of whether the convection pattern characterized by deep water formation north of  $\varphi_M$  is stable or not. A negative value of  $M$  means that dense surface water is created, which can sustain convection and the formation of deep water north of  $\varphi_M$ . Since the net freshwater flux into the northern North Atlantic is positive, the necessary condition for having a negative  $M$  is that the surface cools sufficiently through heat loss to the atmosphere.

While temperature will only play a direct role for convective instability, a change in the freshwater flux into the North Atlantic impacts both advective and convective processes and has different effects on the AMOC depending on the region over which it occurs (Smith and Gregory, 2009; Ganopolski and Rahmstorf, 2001). If applied over deep water formation regions (e.g. 50–70°N) it directly affects  $M$  and can therefore trigger convective instability. This is not necessarily the case if the freshwater flux is applied to latitudes further south (e.g. 20–50°N), where it mainly affects the AMOC through its basin-wide effect on the advective salt feedback.

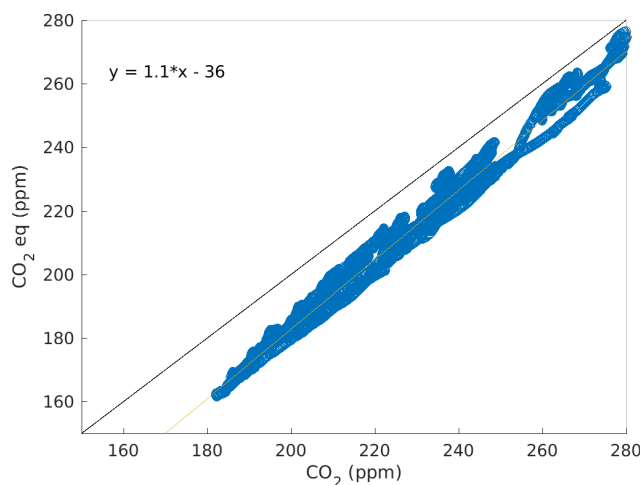
Here we have shown that the conditions under which convective instability occurs are controlled by the annual mean integrated surface buoyancy flux in the northern North Atlantic,  $M$ .  $M$  is also expected to be a useful measure for identifying conditions that could lead to the appearance of spontaneous oscillations in complex general circulation climate models. It should be considered as an approximate measure of the stability of convection, with deviations possible if freshwater export through the latitude  $\varphi_M$  by sea ice transport or ocean mixing are important. While the buoyancy criterion is expected to work best in quasi-equilibrium conditions, a further step will be to investigate whether  $M$  is also a suitable diagnostic to assess the stability of the AMOC under transient global warming scenarios.

*Code and data availability.* The CLIMBER-X model is freely available as open source code at <https://zenodo.org/record/7898797>. The time series of the model simulations shown in the paper are available from Zenodo: <https://doi.org/10.5281/zenodo.8372895>. CMIP6 model data are licensed under a Creative Commons Attribution-ShareAlike 4.0 International License (<https://creativecommons.org/licenses>) and can be accessed through the ESGF nodes (for instance [esgf-data.dkrz.de/search/cmip6-dkrz/](https://esgf-data.dkrz.de/search/cmip6-dkrz/)). Data from the RAPID AMOC monitoring project is funded by the Natural Environment Research Council and are freely available from [www.rapid.ac.uk/rapidmoc](http://www.rapid.ac.uk/rapidmoc).

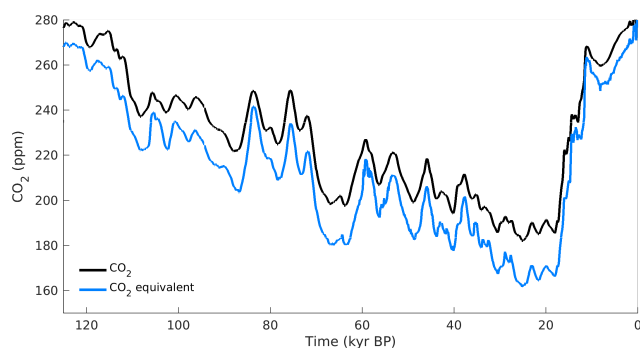


## Appendix A: Equivalent atmospheric CO<sub>2</sub> concentration

235 The equivalent CO<sub>2</sub> concentration is computed taking into account the radiative forcing from CH<sub>4</sub> and N<sub>2</sub>O gases relative to pre-industrial following Etminan et al. (2016) as described in Willeit et al. (2022) using ice core reconstructions of CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub>O from Köhler et al. (2017). The relation between actual and equivalent CO<sub>2</sub> concentration and their evolution over the last glacial cycle are shown in Fig.A1 and Fig. A2.



**Figure A1.** Equivalent atmospheric CO<sub>2</sub> concentration for radiation versus actual atmospheric CO<sub>2</sub> concentration (Köhler et al., 2017) for the last glacial cycle.

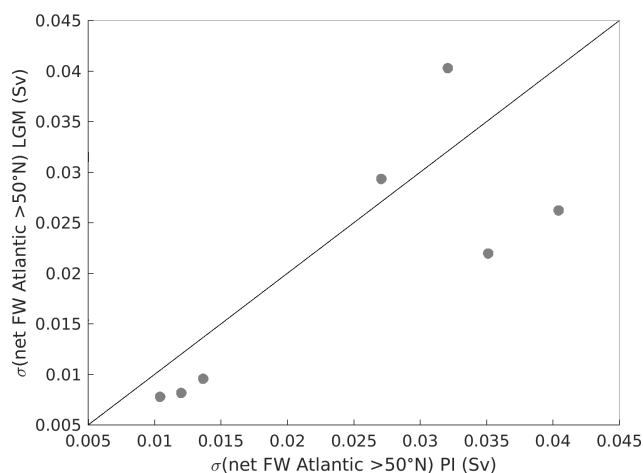


**Figure A2.** Comparison of atmospheric CO<sub>2</sub> from ice core data (Köhler et al., 2017) and equivalent CO<sub>2</sub> for radiation over the last glacial cycle.



## Appendix B: Noise in the surface freshwater flux

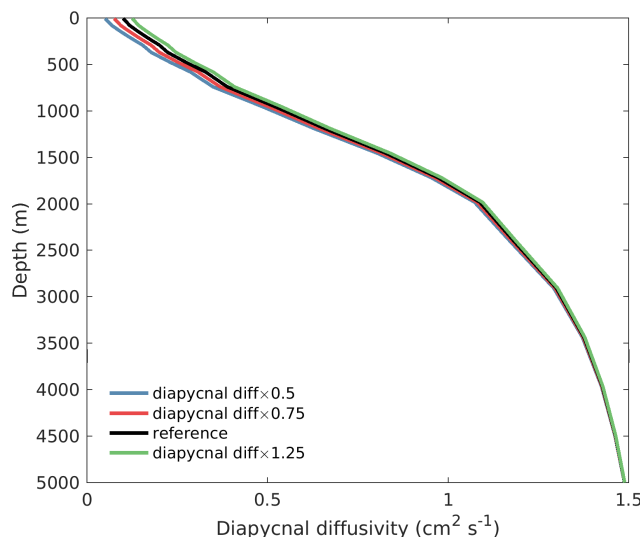
240 To mimic the effect of weather on interannual AMOC variability we applied perturbations to the surface ocean freshwater flux in the North Atlantic in the latitudinal belt between 50–80°N in the form of Gaussian white noise with a standard deviation of  $0.5 \text{ kg m}^{-2} \text{ day}^{-1}$  uniformly over the area and constant over each year. This corresponds to an integrated freshwater flux with a standard deviation of  $\sim 0.07 \text{ Sv}$ , which is roughly twice the variability simulated by PMIP3/4 models (Fig. B1). The larger variability used in our study is justified by the fact that the AMOC is affected also by variations in wind stress, sea ice cover and other factors whose interannual variability is not explicitly accounted for in our study. PMIP3/4 models indicate no clear differences between variability at present and at the LGM (Fig. B1), but freshwater variability in the glacial climate could have been larger due to the presence of surrounding ice sheets calving glaciers at irregular intervals. For pre-industrial conditions applying the noise to the freshwater flux results in an interannual AMOC variability with a standard deviation of  $\sim 1 \text{ Sv}$ , which is comparable to CMIP6 models (Kelson et al., 2022).



**Figure B1.** Interannual standard deviation of the net freshwater flux into the North Atlantic north of 50°N as simulated by different PMIP3/4 models at the LGM (y-axis) versus the pre-industrial (PI) (x-axis).

## 250 Appendix C: Diapycnal diffusivity profiles

Since diapycnal diffusivity is one of the most uncertain parameters in ocean models and often used for model calibration, we designed several different vertical diapycnal diffusivity profiles to test the sensitivity of our results to this important model parameter. In particular, we increased or decreased the diffusivity in the upper ocean by scaling the surface value up or down by factors of 0.5, 0.75 and 1.25. The resulting diapycnal diffusivity profiles are shown in Fig. C1.



**Figure C1.** Profiles of diapycnal diffusivity used in the parameter sensitivity tests shown in Fig. 16. The legend entries specify the scaling factor applied to the value of diapycnal diffusivity at the surface.

## 255 Appendix D: Integrated buoyancy flux

The surface buoyancy flux is computed from the UNESCO equation of state of (Millero and Poisson, 1981), which is used to compute seawater density in CLIMBER-X. However, for the purpose of diagnosing the surface buoyancy flux, using a simple quadratic equation of state including the temperature dependence of the thermal expansion coefficient is sufficiently accurate. Changes in seawater density can thus to a first approximation be written in terms of temperature and salinity changes, as:

$$260 \quad d\rho = -\alpha(T)dT + \beta dS \quad (D1)$$

where  $\alpha(T) = 0.052 + 0.012 \cdot T$  is the temperature-dependent thermal expansion coefficient ( $\text{kg m}^{-3} \text{ } ^\circ\text{C}^{-1}$ ),  $T$  is temperature in  $^\circ\text{C}$  and  $\beta = 0.8$  is the saline contraction coefficient ( $\text{kg m}^{-3} \text{ psu}^{-1}$ ), which is roughly constant. Surface seawater temperature and salinity are directly affected by the net surface heat and freshwater fluxes, and the surface buoyancy flux is a measure of the change in the buoyancy force resulting from the corresponding density changes. The surface buoyancy flux (in  $\text{N m}^{-2} \text{ s}^{-1}$ )

265 can be written as:

$$B = \frac{g}{c_p \rho_0} \alpha(T) \cdot Q + g \beta \frac{S_0}{\rho_0} \cdot F \quad (D2)$$

where  $Q$  is the net heat flux into the ocean ( $\text{W m}^{-2}$ ) and  $F$  is the net freshwater flux into the ocean ( $\text{kg m}^{-2} \text{ s}^{-1}$ ),  $g = 9.81 \text{ m s}^{-2}$  is the acceleration due to gravity,  $c_p = 4187 \text{ J kg}^{-1} \text{ } ^\circ\text{C}^{-1}$  is the specific heat of water at constant pressure,  $\rho_0 = 1000 \text{ kg m}^{-3}$  is a reference density and  $S_0 = 34.7 \text{ psu}$  is a reference salinity. The surface buoyancy flux is therefore sensitive to

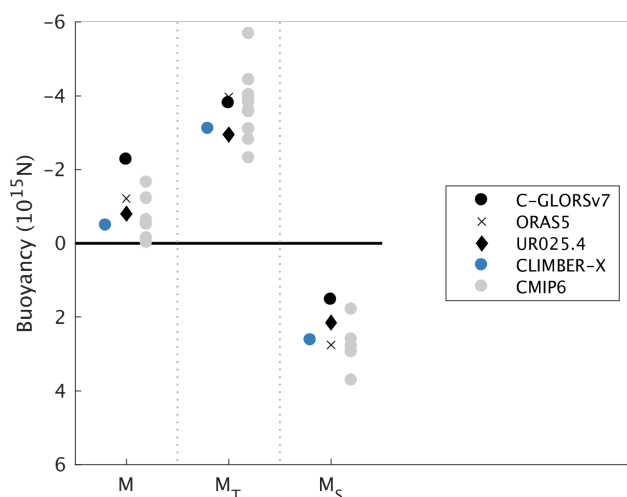
270 changes in the heat and freshwater fluxes, but also to changes in sea surface temperature, through the temperature-dependence of the thermal expansion coefficient  $\alpha$ .



An integrated measure of the buoyancy force (in N) resulting from the application of the buoyancy flux can then be obtained by integrating spatially over an area  $\Omega$  (in our case the whole North Atlantic north of  $55^\circ\text{N}$ ) and in time (over a full year  $\tau$ ):

$$M = \int_{\tau} \left[ \frac{g}{c_p \rho_0} \int_{\Omega} \alpha(T) Q d\Omega + g\beta \frac{S_0}{\rho_0} \int_{\Omega} F d\Omega \right] dt. \quad (\text{D3})$$

275 The integrated buoyancy flux  $M$  for present-day is compared to ocean reanalysis products (Storto and Masina, 2016; Zuo et al., 2019; Haines et al., 2013) and CMIP6 models in Fig. D1.



**Figure D1.** Integrated buoyancy flux  $M$  for the present day as simulated by CLIMBER-X compared to ocean reanalysis (Storto and Masina, 2016; Zuo et al., 2019; Haines et al., 2013) and CMIP6 models. The thermal ( $M_T$ ) and haline ( $M_S$ ) components are also shown separately.

*Author contributions.* MW and AG conceived the study. MW designed and performed the model simulations. AG developed the buoyancy framework, with contributions by MW. All authors contributed to the analysis and discussion of the results. MW produced the figures and wrote the manuscript with input from all co-authors .

280 *Competing interests.* Authors declare that they have no competing interests.

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## References

- 290 Alley, R. B., Anandkrishnan, S., and Jung, P.: Stochastic resonance in the North Atlantic, *Paleoceanography*, 16, 190–198, <https://doi.org/10.1029/2000PA000518>, 2001.
- Andersen, K. K., Azuma, N., Barnola, J. M., Bigler, M., Biscaye, P., Caillon, N., Chappellaz, J., Clausen, H. B., Dahl-Jensen, D., Fischer, H., Flückiger, J., Fritzsche, D., Fujii, Y., Goto-Azuma, K., Grønvold, K., Gundestrup, N. S., Hansson, M., Huber, C., Hvidberg, C. S., Johnsen, S. J., Jonsell, U., Jouzel, J., Kipfstuhl, S., Landais, A., Leuenberger, M., Lorrain, R., Masson-Delmotte, V., Miller, H., Motoyama, H.,
- 295 Narita, H., Popp, T., Rasmussen, S. O., Raynaud, D., Rothlisberger, R., Ruth, U., Samyn, D., Schwander, J., Shoji, H., Siggard-Andersen, M. L., Steffensen, J. P., Stocker, T., Sveinbjörnsdóttir, A. E., Svensson, A., Takata, M., Tison, J. L., Thorsteinsson, T., Watanabe, O., Wilhelms, F., and White, J. W.: High-resolution record of Northern Hemisphere climate extending into the last interglacial period, *Nature*, 431, 147–151, <https://doi.org/10.1038/nature02805>, 2004.
- Armstrong, E., Izumi, K., and Valdes, P.: Identifying the mechanisms of DO-scale oscillations in a GCM: a salt oscillator triggered by the
- 300 Laurentide ice sheet, *Climate Dynamics*, <https://doi.org/10.1007/s00382-022-06564-y>, 2022.
- Boers, N., Ghil, M., and Rousseau, D.-D.: Ocean circulation, ice shelf, and sea ice interactions explain Dansgaard–Oeschger cycles, *Proceedings of the National Academy of Sciences*, 115, 201802573, <https://doi.org/10.1073/pnas.1802573115>, 2018.
- Bohm, E., Lippold, J., Gutjahr, M., Frank, M., Blaser, P., Antz, B., Fohlmeister, J., Frank, N., Andersen, M. B., and Deininger, M.: Strong and deep Atlantic meridional overturning circulation during the last glacial cycle, *Nature*, 517, 73–76, <https://doi.org/10.1038/nature14059>,
- 305 2015.
- Bonan, D. B., Thompson, A. F., Newsom, E. R., Sun, S., and Rugenstein, M.: Transient and Equilibrium Responses of the Atlantic Overturning Circulation to Warming in Coupled Climate Models: The Role of Temperature and Salinity, *Journal of Climate*, 35, 5173–5193, <https://doi.org/10.1175/JCLI-D-21-0912.1>, 2022.
- Bond, G., Broecker, W., Johnsen, S., McManus, J., Labeyrie, L., Jouzel, J., and Bonani, G.: Correlations between climate records from North
- 310 Atlantic sediments and Greenland ice, *Nature*, 365, 143–147, <https://doi.org/10.1038/365143a0>, 1993.
- Braun, H., Christl, M., Rahmstorf, S., Ganopolski, A., Mangini, A., Kubatzki, C., Roth, K., and Kromer, B.: Possible solar origin of the 1,470-year glacial climate cycle demonstrated in a coupled model., *Nature*, 438, 208–11, <https://doi.org/10.1038/nature04121>, 2005.
- Brown, N. and Galbraith, E. D.: Hosed vs. unhosed: Interruptions of the Atlantic Meridional Overturning Circulation in a global coupled model, with and without freshwater forcing, *Climate of the Past*, 12, 1663–1679, <https://doi.org/10.5194/cp-12-1663-2016>, 2016.
- 315 Dansgaard, W., Johnsen, S. J., Clausen, H. B., Dahl-Jensen, D., Gundestrup, N. S., Hammer, C. U., Hvidberg, C. S., Steffensen, J. P., Sveinbjörnsdóttir, A. E., Jouzel, J., and Bond, G.: Evidence for general instability of past climate from a 250-kyr ice-core record, *Nature*, 364, 218–220, <https://doi.org/10.1038/364218a0>, 1993.
- de Vries, P. and Weber, S. L.: The Atlantic freshwater budget as a diagnostic for the existence of a stable shut down of the meridional overturning circulation, *Geophysical Research Letters*, 32, 1–4, <https://doi.org/10.1029/2004GL021450>, 2005.
- 320 Ditlevsen, P. D., Kristensen, M. S., and Andersen, K. K.: The recurrence time of Dansgaard-Oeschger events and limits on the possible periodic component, *Journal of Climate*, 18, 2594–2603, <https://doi.org/10.1175/JCLI3437.1>, 2005.
- Dokken, T. M., Nisancioglu, K. H., Li, C., Battisti, D. S., and Kissel, C.: Dansgaard-Oeschger cycles: Interactions between ocean and sea ice intrinsic to the Nordic seas, *Paleoceanography*, 28, 491–502, <https://doi.org/10.1002/palo.20042>, 2013.



- 325 Etminan, M., Myhre, G., Highwood, E. J., and Shine, K. P.: Radiative forcing of carbon dioxide, methane, and nitrous oxide: A significant revision of the methane radiative forcing, *Geophysical Research Letters*, 43, 12,614–12,623, <https://doi.org/10.1002/2016GL071930>, 2016.
- Friedrich, T., Timmermann, A., Menviel, L., Elison Timm, O., Mouchet, A., and Roche, D. M.: The mechanism behind internally generated centennial-to-millennial scale climate variability in an earth system model of intermediate complexity, *Geoscientific Model Development*, 3, 377–389, <https://doi.org/10.5194/gmd-3-377-2010>, 2010.
- 330 Galbraith, E. and de Lavergne, C.: Response of a comprehensive climate model to a broad range of external forcings: relevance for deep ocean ventilation and the development of late Cenozoic ice ages, *Climate Dynamics*, 52, 653–679, <https://doi.org/10.1007/s00382-018-4157-8>, 2019.
- Ganopolski, A. and Rahmstorf, S.: Rapid changes of glacial climate simulated in a coupled climate model., *Nature*, 409, 153–8, <https://doi.org/10.1038/35051500>, 2001.
- 335 Ganopolski, A. and Rahmstorf, S.: Abrupt Glacial Climate Changes due to Stochastic Resonance, *Physical Review Letters*, 88, 038 501, <https://doi.org/10.1103/PhysRevLett.88.038501>, 2002.
- Gowan, E. J., Zhang, X., Khosravi, S., Rovere, A., Stocchi, P., Hughes, A. L., Gyllencreutz, R., Mangerud, J., Svendsen, J. I., and Lohmann, G.: A new global ice sheet reconstruction for the past 80 000 years, *Nature Communications*, 12, 1–9, <https://doi.org/10.1038/s41467-021-21469-w>, 2021.
- 340 Haines, K., Stepanov, V. N., Valdivieso, M., and Zuo, H.: Atlantic meridional heat transports in two ocean reanalyses evaluated against the RAPID array, *Geophysical Research Letters*, 40, 343–348, <https://doi.org/10.1029/2012GL054581>, 2013.
- Hawkins, E., Smith, R. S., Allison, L. C., Gregory, J. M., Woollings, T. J., Pohlmann, H., and De Cuevas, B.: Bistability of the Atlantic overturning circulation in a global climate model and links to ocean freshwater transport, *Geophysical Research Letters*, 38, 1–6, <https://doi.org/10.1029/2011GL047208>, 2011.
- 345 Henry, L. G., McManus, J. F., Curry, W. B., Roberts, N. L., Piotrowski, A. M., and Keigwin, L. D.: North Atlantic ocean circulation and abrupt climate change during the last glaciation, *Science*, 353, 470–474, <https://doi.org/10.1126/science.aaf5529>, 2016.
- Hodell, D. A., Crowhurst, S. J., Lourens, L., Margari, V., Nicolson, J., Rolfe, J. E., Skinner, L. C., Thomas, N. C., Tzedakis, P. C., Mleneck-Vautravers, M. J., and Wolff, E. W.: A 1.5-million-year record of orbital and millennial climate variability in the North Atlantic, *Climate of the Past*, 19, 607–636, <https://doi.org/10.5194/cp-19-607-2023>, 2023.
- 350 Hoff, U., Rasmussen, T. L., Stein, R., Ezat, M. M., and Fahl, K.: Sea ice and millennial-scale climate variability in the Nordic seas 90 kyr ago to present, *Nature Communications*, 7, 12 247, <https://doi.org/10.1038/ncomms12247>, 2016.
- Hu, A., Meehl, G. A., Han, W., Timmermann, A., Otto-Bliesner, B., Liu, Z., Washington, W. M., Large, W., Abe-Ouchi, A., Kimoto, M., Lambeck, K., and Wu, B.: Role of the Bering Strait on the hysteresis of the ocean conveyor belt circulation and glacial climate stability, *Proceedings of the National Academy of Sciences of the United States of America*, 109, 6417–6422, <https://doi.org/10.1073/pnas.1116014109>, 2012.
- 355 Kageyama, M., Albani, S., Braconnot, P., Harrison, S. P., Hopcroft, P. O., Ivanovic, R. F., Lambert, F., Marti, O., Richard Peltier, W., Peterschmitt, J. Y., Roche, D. M., Tarasov, L., Zhang, X., Brady, E. C., Haywood, A. M., Legrande, A. N., Lunt, D. J., Mahowald, N. M., Mikolajewicz, U., Nisancioglu, K. H., Otto-Bliesner, B. L., Renssen, H., Tomas, R. A., Zhang, Q., Abe-Ouchi, A., Bartlein, P. J., Cao, J., Li, Q., Lohmann, G., Ohgaito, R., Shi, X., Volodin, E., Yoshida, K., Zhang, X., and Zheng, W.: The PMIP4 contribution to CMIP6 - Part 4: Scientific objectives and experimental design of the PMIP4-CMIP6 Last Glacial Maximum experiments and PMIP4 sensitivity experiments, *Geoscientific Model Development*, 10, 4035–4055, <https://doi.org/10.5194/gmd-10-4035-2017>, 2017.





- 365 Kageyama, M., Harrison, S. P., Kapsch, M. L., Lofverstrom, M., Lora, J. M., Mikolajewicz, U., Sherriff-Tadano, S., Vadsaria, T., Abe-Ouchi, A., Bouttes, N., Chandan, D., Gregoire, L. J., Ivanovic, R. F., Izumi, K., Legrande, A. N., Lhardy, F., Lohmann, G., Morozova, P. A., Ohgaito, R., Paul, A., Richard Peltier, W., Poulsen, C. J., Quiquet, A., Roche, D. M., Shi, X., Tierney, J. E., Valdes, P. J., Volodin, E., and Zhu, J.: The PMIP4 Last Glacial Maximum experiments: Preliminary results and comparison with the PMIP3 simulations, *Climate of the Past*, 17, 1065–1089, <https://doi.org/10.5194/cp-17-1065-2021>, 2021.
- Keigwin, L. D. and Boyle, E. A.: Surface and deep ocean variability in the northern Sargasso Sea during marine isotope stage 3, *Paleoceanography*, 14, 164–170, <https://doi.org/10.1029/1998PA900026>, 1999.
- 370 Kelson, R. L., Straub, D. N., and Dufour, C. O.: Using CMIP6 Models to Assess the Significance of the Observed Trend in the Atlantic Meridional Overturning Circulation, *Geophysical Research Letters*, 49, <https://doi.org/10.1029/2022GL100202>, 2022.
- Kindler, P., Guillevic, M., Baumgartner, M., Schwander, J., Landais, A., and Leuenberger, M.: Temperature reconstruction from 10 to 120 kyr b2k from the NGRIP ice core, *Climate of the Past*, 10, 887–902, <https://doi.org/10.5194/cp-10-887-2014>, 2014.
- Klockmann, M., Mikolajewicz, U., and Marotzke, J.: Two AMOC states in response to decreasing greenhouse gas concentrations in the coupled climate model MPI-ESM, *Journal of Climate*, 31, 7969–7984, <https://doi.org/10.1175/JCLI-D-17-0859.1>, 2018.
- 375 Köhler, P., Nehrbass-Ahles, C., Schmitt, J., Stocker, T. F., and Fischer, H.: A 156 kyr smoothed history of the atmospheric greenhouse gases CO<sub>2</sub>, CH<sub>4</sub> and N<sub>2</sub>O and their radiative forcing, *Earth System Science Data*, 9, 363–387, <https://doi.org/10.5194/essd-9-363-2017>, 2017.
- Kuniyoshi, Y., Abe-Ouchi, A., Sherriff-Tadano, S., Chan, W. L., and Saito, F.: Effect of Climatic Precession on Dansgaard-Oeschger-Like Oscillations, *Geophysical Research Letters*, 49, 1–10, <https://doi.org/10.1029/2021GL095695>, 2022.
- Li, C. and Born, A.: Coupled atmosphere-ice-ocean dynamics in Dansgaard-Oeschger events, *Quaternary Science Reviews*, 203, 1–20, <https://doi.org/10.1016/j.quascirev.2018.10.031>, 2019.
- 380 Lohmann, J. and Ditlevsen, P. D.: Random and externally controlled occurrences of Dansgaard-Oeschger events, *Climate of the Past*, 14, 609–617, <https://doi.org/10.5194/cp-14-609-2018>, 2018.
- Lohmann, J. and Ditlevsen, P. D.: Objective extraction and analysis of statistical features of Dansgaard-Oeschger events, *Climate of the Past*, 15, 1771–1792, <https://doi.org/10.5194/cp-15-1771-2019>, 2019.
- 385 Malmierca-Vallet, I. and Sime, L. C.: Dansgaard-Oeschger events in climate models: review and baseline Marine Isotope Stage 3 (MIS3) protocol, *Climate of the Past*, 19, 915–942, <https://doi.org/10.5194/cp-19-915-2023>, 2023.
- Martrat, B., Grimalt, J. O., Shackleton, N. J., de Abreu, L., Hutterli, M. A., and Stocker, T. F.: Four Climate Cycles of Recurring Deep and Surface Water Destabilizations on the Iberian Margin, *Science*, 317, 502–507, <https://doi.org/10.1126/science.1139994>, 2007.
- McManus, J. F., Francois, R., Gherardi, J.-M., Keigwin, L. D., and Brown-Leger, S.: Collapse and rapid resumption of Atlantic meridional circulation linked to deglacial climate changes., *Nature*, 428, 834–7, <https://doi.org/10.1038/nature02494>, 2004.
- 390 Menviel, L. C., Skinner, L. C., Tarasov, L., and Tzedakis, P. C.: An ice–climate oscillatory framework for Dansgaard–Oeschger cycles, *Nature Reviews Earth & Environment*, 1, 677–693, <https://doi.org/10.1038/s43017-020-00106-y>, 2020.
- Millero, F. J. and Poisson, A.: International one-atmosphere equation of state of seawater, *Deep Sea Research Part A, Oceanographic Research Papers*, 28, 625–629, [https://doi.org/10.1016/0198-0149\(81\)90122-9](https://doi.org/10.1016/0198-0149(81)90122-9), 1981.
- 395 Mitsui, T. and Crucifix, M.: Influence of external forcings on abrupt millennial-scale climate changes: a statistical modelling study, *Climate Dynamics*, 48, 2729–2749, <https://doi.org/10.1007/s00382-016-3235-z>, 2017.
- Oka, A., Hasumi, H., and Abe-Ouchi, A.: The thermal threshold of the Atlantic meridional overturning circulation and its control by wind stress forcing during glacial climate, *Geophysical Research Letters*, 39, 1–6, <https://doi.org/10.1029/2012GL051421>, 2012.



- Peltier, W. R. and Vettoretti, G.: Dansgaard-Oeschger oscillations predicted in a comprehensive model of glacial climate: A "kicked" salt oscillator in the Atlantic, *Geophysical Research Letters*, 41, 7306–7313, <https://doi.org/10.1002/2014GL061413>, 2014.
- Pikovsky, A. S. and Kurths, J.: Coherence Resonance in a Noise-Driven Excitable System, *Physical Review Letters*, 78, 775–778, <https://doi.org/10.1103/PhysRevLett.78.775>, 1997.
- Pöppelmeier, F., Jeltsch-Thömmes, A., Lippold, J., Joos, F., and Stocker, T. F.: Multi-proxy constraints on Atlantic circulation dynamics since the last ice age, *Nature Geoscience*, 16, <https://doi.org/10.1038/s41561-023-01140-3>, 2023.
- 405 Prange, M., Jonkers, L., Merkel, U., Schulz, M., and Bakker, P.: A multicentennial mode of North Atlantic climate variability throughout the Last Glacial Maximum, *Science advances*, 9, eadh1106, <https://doi.org/10.1126/sciadv.adh1106>, 2023.
- Rahmstorf, S.: Rapid climate transitions in a coupled ocean-atmosphere model, <https://doi.org/10.1038/372082a0>, 1994.
- Rahmstorf, S.: Bifurcations of the Atlantic thermohaline circulation in response to changes in the hydrological cycle, *Nature*, 378, 145–149, <https://doi.org/10.1038/378145a0>, 1995.
- 410 Rahmstorf, S.: On the freshwater forcing and transport of the Atlantic thermohaline circulation, *Climate Dynamics*, 12, 799–811, <https://doi.org/10.1007/s003820050144>, 1996.
- Rahmstorf, S.: A simple model of seasonal open ocean convection - Part I: Theory, *Ocean Dynamics*, 52, 26–35, <https://doi.org/10.1007/s10236-001-8174-4>, 2001.
- Rahmstorf, S.: Ocean circulation and climate during the past 120,000 years., *Nature*, 419, 207–14, <https://doi.org/10.1038/nature01090>,  
415 2002.
- Rahmstorf, S.: Timing of abrupt climate change: A precise clock, *Geophysical Research Letters*, 30, n/a–n/a, <https://doi.org/10.1029/2003GL017115>, 2003.
- Romé, Y. M., Ivanovic, R. F., Gregoire, L. J., Sherriff-Tadano, S., and Valdes, P. J.: Millennial-Scale Climate Oscillations Triggered by Deglacial Meltwater Discharge in Last Glacial Maximum Simulations, *Paleoceanography and Paleoclimatology*, 37, <https://doi.org/10.1029/2022PA004451>, 2022.
- 420 Sadatzki, H., Dokken, T. M., Berben, S. M., Muschitiello, F., Stein, R., Fahl, K., Menviel, L., Timmermann, A., and Jansen, E.: Sea ice variability in the southern norwegian sea during glacial dansgaard-oeschger climate cycles, *Science Advances*, 5, 1–11, <https://doi.org/10.1126/sciadv.aau6174>, 2019.
- Sadatzki, H., Maffezzoli, N., Dokken, T. M., Simon, M. H., Berben, S. M., Fahl, K., Kjær, H. A., Spolaor, A., Stein, R., Vallelonga, P., Vinther, B. M., and Jansen, E.: Rapid reductions and millennial-scale variability in Nordic Seas sea ice cover during abrupt glacial climate changes, *Proceedings of the National Academy of Sciences of the United States of America*, 117, 29478–29486, <https://doi.org/10.1073/pnas.2005849117>, 2020.
- 425 P., Vinther, B. M., and Jansen, E.: Rapid reductions and millennial-scale variability in Nordic Seas sea ice cover during abrupt glacial climate changes, *Proceedings of the National Academy of Sciences of the United States of America*, 117, 29478–29486, <https://doi.org/10.1073/pnas.2005849117>, 2020.
- Sakai, K. and Peltier, W. R.: Dansgaard–Oeschger Oscillations in a Coupled Atmosphere–Ocean Climate Model, *Journal of Climate*, 10, 949–970, <http://www.jstor.org/stable/26243202>, 1997.
- 430 Schulz, M., Prange, M., and Klocker, A.: Low-frequency oscillations of the Atlantic Ocean meridional overturning circulation in a coupled climate model, *Climate of the Past*, 3, 97–107, <https://doi.org/10.5194/cp-3-97-2007>, 2007.
- Scoto, F., Sadatzki, H., Maffezzoli, N., Barbante, C., Gagliardi, A., Varin, C., Vallelonga, P., Gkinis, V., Dahl-Jensen, D., Kjær, H. A., Burgay, F., Saiz-Lopez, A., Stein, R., and Spolaor, A.: Sea ice fluctuations in the Baffin Bay and the Labrador Sea during glacial abrupt climate changes, *Proceedings of the National Academy of Sciences of the United States of America*, 119, 1–9, <https://doi.org/10.1073/pnas.2203468119>, 2022.
- 435



- Skinner, L. C. and Elderfield, H.: Rapid fluctuations in the deep North Atlantic heat budget during the last glacial period, *Paleoceanography*, 22, 1–9, <https://doi.org/10.1029/2006PA001338>, 2007.
- Smith, R. S. and Gregory, J. M.: A study of the sensitivity of ocean overturning circulation and climate to freshwater input in different regions of the North Atlantic, *Geophysical Research Letters*, 36, 1–5, <https://doi.org/10.1029/2009GL038607>, 2009.
- 440 Stocker, T. F. and Wright, D. G.: Rapid transitions of the ocean’s deep circulation induced by changes in surface water fluxes, *Nature*, 351, 729–732, <https://doi.org/10.1038/351729a0>, 1991.
- Stommel, H.: Thermohaline Convection with Two Stable Regimes of Flow, *Tellus*, 13, 224–230, <https://doi.org/10.1111/j.2153-3490.1961.tb00079.x>, 1961.
- Storto, A. and Masina, S.: C-GLORSv5: An improved multipurpose global ocean eddy-permitting physical reanalysis, *Earth System Science* 445 Data, 8, 679–696, <https://doi.org/10.5194/essd-8-679-2016>, 2016.
- Stouffer, R. J. and Manabe, S.: Equilibrium response of thermohaline circulation to large changes in atmospheric CO<sub>2</sub> concentration, *Climate Dynamics*, 20, 759–773, <https://doi.org/10.1007/s00382-002-0302-4>, 2003.
- Svensson, A., Dahl-Jensen, D., Steffensen, J. P., Blunier, T., Rasmussen, S. O., Vinther, B. M., Vallelonga, P., Capron, E., Gkinis, V., Cook, E., Astrid Kjær, H., Muscheler, R., Kipfstuhl, S., Wilhelms, F., Stocker, T. F., Fischer, H., Adolphi, F., Erhardt, T., Sigl, M., Landais, A., 450 Parrenin, F., Buizert, C., McConnell, J. R., Severi, M., Mulvaney, R., and Bigler, M.: Bipolar volcanic synchronization of abrupt climate change in Greenland and Antarctic ice cores during the last glacial period, *Climate of the Past*, 16, 1565–1580, <https://doi.org/10.5194/cp-16-1565-2020>, 2020.
- Tarasov, L., Dyke, A. S., Neal, R. M., and Peltier, W. R.: A data-calibrated distribution of deglacial chronologies for the North American ice complex from glaciological modeling, *Earth and Planetary Science Letters*, 315–316, 30–40, 455 <https://doi.org/https://doi.org/10.1016/j.epsl.2011.09.010>, 2012.
- Timmermann, A., Gildor, H., Schulz, M., and Tziperman, E.: Coherent resonant millennial-scale climate oscillations triggered by massive meltwater pulses, *Journal of Climate*, 16, 2569–2585, [https://doi.org/10.1175/1520-0442\(2003\)016<2569:CRMCOT>2.0.CO;2](https://doi.org/10.1175/1520-0442(2003)016<2569:CRMCOT>2.0.CO;2), 2003.
- Vélez-Belchí, P., Alvarez, A., Colet, P., Tintoré, J., and Haney, R. L.: Stochastic resonance in the thermohaline circulation, *Geophysical Research Letters*, 28, 2053–2056, <https://doi.org/10.1029/2000GL012091>, 2001.
- 460 Vettoretti, G., Ditlevsen, P., Jochum, M., and Rasmussen, S. O.: Atmospheric CO<sub>2</sub> control of spontaneous millennial-scale ice age climate oscillations, *Nature Geoscience*, 15, 300–306, <https://doi.org/10.1038/s41561-022-00920-7>, 2022.
- Willeit, M., Ganopolski, A., Robinson, A., and Edwards, N. R.: The Earth system model CLIMBER-X v1.0 – Part 1: Climate model description and validation, *Geoscientific Model Development*, 15, 5905–5948, <https://doi.org/10.5194/gmd-15-5905-2022>, 2022.
- Wood, R. A., Rodríguez, J. M., Smith, R. S., Jackson, L. C., and Hawkins, E.: Observable, low-order dynamical controls on thresholds of the 465 Atlantic meridional overturning circulation, *Climate Dynamics*, 53, 6815–6834, <https://doi.org/10.1007/s00382-019-04956-1>, 2019.
- Zhang, X., Barker, S., Knorr, G., Lohmann, G., Drysdale, R., Sun, Y., Hodell, D., and Chen, F.: Direct astronomical influence on abrupt climate variability, *Nature Geoscience*, 14, 819–826, <https://doi.org/10.1038/s41561-021-00846-6>, 2021.
- Zuo, H., Balmaseda, M. A., Tietsche, S., Mogensen, K., and Mayer, M.: The ECMWF operational ensemble reanalysis-analysis system for ocean and sea ice: A description of the system and assessment, *Ocean Science*, 15, 779–808, <https://doi.org/10.5194/os-15-779-2019>, 470 2019.