Overview & general comments

This paper presents the results of an initial, and idealized, modelling study of the impact of the 2-dimensional spatial structure of broken sea ice and water surfaces on surface fluxes of momentum and heat for a fixed ice fraction. This is an important, but to date, little studied topic.

During much of the year, except for the summer melt, the skin temperatures of open water and sea ice can be very different, resulting in sudden changes in surface stability as air flows the boundaries between ice and water surfaces. Notably the air temperature – as in this study – may be at a temperature between that of the ice and water, so that stability changes from unstable over water to stable over ice. Depending on the length scale of continuous stretches of ice or water, the near-surface air may never have chance to adjust completely and for turbulence to reach a steady state. This has important implications for surface flux parameterizations in models, since the basis of these is Monin-Obukhov similarity theory, which relies on an assumption of statistical stationarity of turbulence and horizontal homogeneity of the surface.

Most surface flux parameterizations over sea ice are functions of ice fraction only (e.g. Elvidge et al. 2016, 2021), with some (e.g. Lüpkes et al. 2012, in its most complete form) including information about ice properties such as thickness, ridging, etc., which affect its roughness. Lüpkes and Gryaniak (2015) have also attempted to address, from a theoretical perspective, the issue of changing stability with flow between ice and water surfaces.

This study does not directly address the issue of changing stability, although it is an inherent feature of the model set up, but examines a consequence – the fact that surface exchange, vertical profiles of turbulent fluxes, and the evolution of the vertical structure of the boundary layer, become functions of the spatial distribution of ice and water, even when the total fractions remain constant.

The results demonstrate that this dependence can be strong, and depends also on the orientation of the mean wind with respect to any anisotropy of the spatial patterns of ice and water. A key result is the importance of the diffusive flux contributions to the total flux – a result of spatially coherent correlations between vertical air motions and other quantities after time averaging.

One caveat that I think worth your discussing. While the results make clear the potential importance of the spatial distribution of ice/water surfaces, and the consequent potential importance of diffusive fluxes resulting from organised circulations induced by that spatial distribution, I think the cases presented here are likely to be extremes. The LES configuration has been chosen to ensure a surface flow parallel to the x-axis/east-west and to one axis of symmetry of the simple ice/water surface patterns. This ensures a nice, simple, well defined – and with periodic boundaries – repeating pattern of surface forcing. There is also a strong difference in surface temperatures for ice/water and hence the stability over the two surfaces. In all cases, part of the flow will encounter only a water surface (at north/south edges of domain), never flowing over ice, and thus always convective surface forcing, while the remaining flow will encounter varying
fractions of alternating ice and water surfaces. This ought to encourage some sort of organised circulation to develop, forcing spatially coherent features in the time-averaged fields, and hence promoting the diffusive fluxes. This is unlikely to occur in this way in the real world. It doesn’t invalidate the results, but it may make them rather extreme cases, and this should be noted.

Thank you for your overall positive assessment and valuable comments. Regarding the last paragraph, the periodicity in our domain is chosen such that it represents the MIZ properly; in other words, the LES surface is an infinitely long domain with patches. Regarding the comment that these may be extreme cases, they indeed are by design. Our aim, which may not have been clear in version 1, is to show the maximum extent to which the spatial configuration of ice patches is important to the flow. We clarify in the revised version that in many real MIZ setups the effect would be smaller.

We also note that simulations in a previous paper (Fogarty & Bou-Zeid 2023) looked at circulations in two cases, one with consistently alternating strips of ice and water, and one with two parallel infinitely long patches of ice and water. In these cases, that convective surface forcing was not always present. Also in a forthcoming paper we will submit soon, we will be looking at real ice maps and examining how the configuration metrics developed here can help us estimate the effect of surface heterogeneity.

**Recommendation: publish with minor revisions**


Thank you for your responses and comments. The above citations have been added, and the comments below will be addressed.

**Detailed comments**

**Line 16:** “...shows persistent biases in coarse-resolution...” – need some additional detail, biases in...what? Ice fraction, thickness,...?

We mean biases towards predicted sea ice fraction and sea ice extent – this has been clarified in the text.
Table 1, and description of model setup (line 105-117) – most of this is fine, but I think a brief description of the initial air temperature is required. The way this is actually defined is technical, and sticking it in an appendix entirely appropriate, but I got lost when I reached the initial results. The formation of wind speed jets here requires a loss of frictional coupling to the surface to allow the low level wind to accelerate, and hence a stable stratification of the near-surface air...at which point I realized nothing had been said about the air temperature, and I had to go hunting through the text to see if I’d missed it. Eventually spotted the reference to appendix in Table 1, and could go and figure things out. But that really disrupts reading. Don’t need a detailed description in the main text, but a brief statement along the lines of ‘the initial surface air temperature is defined such that the area-averaged sensible heat flux is zero, and thus lies between that of the ice and water skin temperatures’ (table 1 implies it is constant with height...is that temperature or (I assume) potential temperature?)

Thank you for noting this disruption in reading, we have added: “The initial air temperature, a constant profile of potential temperature, $\theta$, is defined such that the area-averaged sensible heat flux is zero, and thus lies between that of the ice and water skin temperatures (see Appendix \ref{app:temp_init} for details),”

**Figure 4:** I assume the subscript ‘g’ on Mg is for ‘geostrophic’ – this isn’t noted in the text and should be stated for clarity.

Added clarifications on the normalization on all the subplots in Figure 4

Line 199: “All patterns except for Pattern5 developed a low-level jet” – it is pushing the definition of a low-level jet, but even Pattern5 does show a weak local maximum in the wind speed at $z/zi \sim 0.3$

We edited to include Pattern 5 as well: “All patterns developed a low-level jet (LLJ), which can be seen in Figure \ref{fig:lesresults_4panel}a, though the LLJs in Pattern1 and Pattern5 are weak \citep{tetzlaff_aircraft_2015, michealis_convective_2021}. But we are a bit uncertain whether the reviewer is also referring to the fact that these are not the canonical LLJ people observe, and that may emanate from the Holton or Blackadar mechanisms (Du and Rotunno, 2014). Indeed this may be true; however as the reviewer postulates in the next comment, a Blackadar-like mechanism may be at play where parcels of air advecting from warm to cold patches experience decoupling from the surface and accelerate (this plays out in time in the original Blackadar version). Unfortunately, delving further into these physics and possibilities would not be feasible for this paper.


Lines 200-202: discussion of low-level jets. The text states – “The LLJs seem to increase in Pattern2 and Pattern3, likely due to large swaths of ice in the direction of the geostrophic wind (and therefore little interruption by the unstable ocean surface).” – which is true, but I think it
would be beneficial to expand on the discussion a little. A reader from the sea ice community rather than boundary-layer meteorology might not immediately appreciate that here a low-level jet forms only because of stable stratification below it, decoupling the air from surface friction, allowing it to accelerate and form a jet; and hence that over water, where the stratification is unstable, convective mixing would prevent a jet forming (or inhibit/erode a jet formed over the ice). I think a brief explanation of why the jet forms is needed, and perhaps more emphasis given to the very different stability over the ice/water surfaces.

Added a brief explanation on the formation of jets: “The LLJs seem to increase in Pattern2 and Pattern3, likely due to large swaths of ice in the direction of the geostrophic wind; the stable stratification in these ice regions decouple the air from the surface friction, allowing low-level acceleration of the wind (conversely, over an unstable ocean surface, the convection produced by the relatively warmer water inhibits this phenomenon). This mechanism is similar to the one advanced by Blackadar (1957) for creation of a low level jet via an inertial oscillation in time as the ABL transition to stable at sunset. However, in this case the oscillation is in space as columns of air advect from a hot to a cold surface and decouple from the surface.”


Line 215 & figure 4: “…indicating significant differences in the wind and stress Ekman rotation with height” – might be useful to add a plot showing the wind direction profiles, to clearly show this.

We’ve added to Figure 4 the wind direction profiles, right below the wind magnitude, to better understand the differences in wind turning with height.

**Figure 5.** Should make clear that altitude is normalized, = z/zi just to avoid any confusion.

Updated Figure 5 and its description: “Normalized vertical profiles of normalized total heat flux...”

Line 280: “one would expect to find more ice-water edge instances, and thus more regions of stable-to-unstable stratification transition” – while it is implicit in the set up, it is perhaps worth noting here that increasing patch density doesn’t simply increase the number of stable-to-unstable stratification transitions, but also the unstable-to-stable, and importantly, reduces the time available over each consecutive surface type for the near-surface flow to adjust to the transition. At low PD, it is plausible that the flow approaches quasi-equilibrium, while for high PD that is never going to happen.

This is true. We’ve now noted this in the text: “As the PD of a sea ice surface increases, one would expect to find more ice-water edge instances, and thus more regions of stable-to-unstable and unstable-to-stable stratification transitions. We also note that PD increases the average time the
parcel spends over the stable (or unstable) surface, which affects how said parcel adjusts to the transition to this new stability regime."

However, at low PD one might have some large patches that generate strong secondary circulations, while at very high PD it is also possible that the patch scale becomes so small that the atmosphere essentially sees a homogeneous average surface. So the interpretations may be more complex.

Line 290: it’s not clear to me here exactly how PAFRAC is calculated...is the value just that of the gradient, k, in equation 8?

In the equation $A = k P^{2/D}$, it is $D$ that is PAFRAC. To keep it consistent with the other equations, we’ve replaced $D$ with PAFRAC in the equation and the text, and removed the variable ‘$D$’ from the manuscript

Line 300-307: This discussion presents the case for why you need to assess how the measures of ice/water distribution metrics behave when derived from images at different scales. I agree it is relevant and important to do this, but I think the rationale presented focuses on the wrong things.

The primary argument given here is that NWP models, and even most LES, have grid resolutions far lower than that of the high resolution imagery used here. True, but I think, irrelevant. The models don’t require the raw imagery, only the metrics derived from it over an area of one grid cell, and ultimately the resulting transfer coefficients. A more relevant issue, is that if this sort of information about ice/water spatial distribution is to be used operationally, then it will need to be updated on regular basis, and thus ideally draw upon all available imagery, each source of which may be at a different resolution. It is thus important to be able to achieve consistent results across different image sources.

Climate models are a different issue, since metrics must come from the sea ice model.

Thank you for your comment: overall, we think both ideas are true and important and we incorporated them since they bolster the ideas of the manuscript. We have thus added the following: “Furthermore, when considering the operational use of these metrics, the regular updating of these values would likely draw upon multiple satellite products with differing resolutions; thus metrics that are able to be extrapolated/interpolated between different grid cell sizes would allow for a consistent computation of metrics when standardized to a single weather model grid cell.”

Figure 6: not clear why a power law scaling function is fit to the ice fraction when this is stated in the text to be scale invariant.

Good point, since ice fraction is scale invariant, we do not need to fit a power law (as seen by the value of $D_p$ being so low). We’ve changed the figure to only fit a power law to PD, and we’ve changed the caption to reflect this.
Figure 7. This shows ice/water maps derived from high resolution satellite imagery. Why are there clear regions of mirror symmetry, both horizontal and vertical (though curiously not always extending the full length of the image). I’ve marked the symmetries as red dashed lines on the copy of the figure below. These may not affect the results (though ought to have a minor impact on the precise alignment of the eigenvectors), but they jump out as an oddity.

Some of the original maps obtained from the dataset described in Section 2.2 were not full squares of data, i.e. there were ice cells, water cells, and ‘no data’ cells. To run this in LES in a square domain, we “reflected” the pattern over an axis of symmetry to preserve the pattern as best as possible. This only affected small border areas of the images. All map metrics, calculations, and analyses have been done on the “reflected” patterns. You are right that this does not affect the results of the VIF analysis, but may slightly affect the results of eigenvectors. We added a discussion on this in Section 2.2 – “Sea Ice Data”, saying: “Some of the images did not fully cover this full extent, and thus in order to retain the real-world sea ice geometry, we “reflected” this onto the areas of no data. All metric calculations and analyses have been done on these modified surfaces.”

Lines 343, 358: At line 343 where the eigenvectors are introduced it is stated that “The principal eigendirection points in the direction of minimal variance”, but then at 358 “the principal eigendirection explains much more of the variance” contradicting the first statement.

Thank you for pointing this out, it should be the secondary eigendirection that contains more of the variance; we’ve edited this to say: “In theory, a sea ice map with a high \( POV(\lambda_0) \), and thus a low \( POV(\lambda_1) \), would be anisotropic, since the secondary eigendirection would contain much more of the variance than the principal eigendirection, and the surface thus has a preferential direction of variability”

Lines358-362: “a sea ice map with a high \( POV(\lambda_0) \) would be anisotropic”, “Conversely, a map with a low \( POV(\lambda_0) \) would be a fairly isotropic map”, a ‘low’ value tends to imply \( \rightarrow 0 \), doesn’t \( POV(\lambda_0) \rightarrow 0.5 \) for isotropic conditions as implied by the statement “By definition, \( POV(\lambda_0) > 0.5 \), since the \( POV(\lambda_0) \) is the POV for the principal eigendirection.”. What happens for a truly isotropic surface, where there is no preferential direction?

For a truly isotropic surface (one where, let’s say, the entire map is ice or the entire map is water), the program defaults to a left-to-right eigenvector, as it cannot pick a direction that stands out more than the others. And yes, the \( POV(\lambda_0) = 0.5 \) in this case – this has been clarified in the text.

Luckily, this doesn’t present an issue operationally, since in a truly isotropic surface, there would no need to find such a “principal direction.” We’ve added to the manuscript:

“By definition, \( POV(\lambda_0) \geq 0.5 \) (with \( POV(\lambda_0) = POV(\lambda_1) = 0.5 \) resulting for a truly isotropic surface), since \( POV(\lambda_0) \) is the POV for the principal eigendirection.”
Table 4 caption: “Note that these angles are not traditional meteorological wind angles, but are instead in Cartesian coordinates, as $0^\circ$ is a left-to-right westerly wind” – these are not wind angles (or rather, directions) at all. The important point is that the angles are stated in a Cartesian framework, increasing anticlockwise from the x-axis. They could be restated as compass headings...though it's not obvious here whether the ice maps are oriented north or with the field of view of the individual satellite orbits.

Yes, the wind angles here do not follow meteorological convention, as 0 is left-to-right westerly wind, and 90 would be southerly. This has been clarified a bit further in the caption: “Note that these angles are not traditional meteorological wind angles, but are instead in Cartesian coordinates; $0^\circ$ is a left-to-right westerly wind, and $90^\circ$ is a southerly wind.”

Line 457: “...such that the heat flux over the ice is equivalent to the heat flux coming from the water,” – it’s not clear from the text, but the implication of equation (B1) is that the initial temperature is chosen so that the mean heat flux is zero. ‘equivalent to’ is rather vague, be explicit.

Yes, the interpretation is correct and this has been made clearer in the text: “The initial temperature was chosen such that the mean heat flux over the entire domain is zero; in other words, the heat flux going into the ice is equivalent in magnitude to the heat flux coming from the water, based on the area fraction (ice fraction, in this case) of the domain.”

Minor typos & grammatical issues

All grammatical issues below have been resolved.

Line 6: “...such as those done in...” -> “...such as those used in...”

Line 28: “...it thus are...” -> “...it is thus...”

Line 42: “(as show for...)” -> “(as shown for...)”

Line 74: “...average pact compaction...” -> “...average patch compaction...” – I assume ‘patch’ but maybe something else was intended?

Line 130: “...here had already underwent...” -> “...here already underwent...”

Line 217: “...dispersive and turbulent counterparts” – ‘components’ or ‘contributions’ might be better words than ‘counterparts’ here.