Snow Particle Fragmentation Enhances Snow Sublimation

Ning Huang^{1,3}, Jiacheng Bao^{1,3}, Hongxiang Yu², and Guang Li²

Correspondence: Hongxiang Yu (yuhx2023@lzu.edu.cn); Guang Li(lg@lzu.edu.cn)

Abstract. Drifting and blowing snow (DBS) affects the mass and energy balance of snow cover through the transportation and sublimation of snow particles, playing a crucial role in atmospheric and hydrological processes. During DBS events, Fragmentation of snow particles, where dendritic snowflakes transform into roundedspherical shapes upon impact with surface and other particles during drifting and blowing snow events, plays a critical role in shaping snow dynamics. This phenomenon is important because it influences the size distribution and concentration of snow particles, affecting mass flux and sublimation rates. However, existing Currently, prevailing drifting and blowing snow models of drifting and blowing snow often neglect the effects of ignore the snow particles fragmentation, introducing leading to heightened uncertainties in the prediction of predicting flow dynamics structures and sublimation rates. In this study, we incorporate a snow particle fragmentation model into a well-developed wind blowing snow model to quantitatively investigate the influence of fragmentation under varying wind conditions Here, we aim to quantitatively investigate the impact of snow particle fragmentation on sublimation. We establish a drifting and blowing snow model considering the snow particle fragmentation process and investigated the effects of fragmentation on drifting and blowing snow. Our results reveal that fragmentation within the saltation layer generates smaller particles, leading to an increase in mass flux and subsequently enhancing sublimation rates in drifting and blowing snowWe found that fragmentation enhances the sublimation of blowing snow and changes the airborne particle size distribution. Notably, the effects of fragmentation on sublimation are more pronounced for suspension particles than saltation particles, particularly under low wind conditions The sublimation rate of saltating snow particles increases 11 % on average, and that of suspension snow particles increases 76 % on average, when the friction wind speed is between 0.3 m/s to 0.5 m/s. This work highlights the critical role of collision-induced fragmentation, wherein dendritic snowflakes are shattered into smaller particles during transportation. This quantitative assessment of fragmentation impact on snow sublimation underscores its importance for improving the physical representation of drifting and blowing snow. These findings have important implications for improving the snow transportphysical dynamic models, with potential applications of drifting and blowing snow, which may contribute to predictions in snow hydrology and climate modeling models.

¹College of Civil Engineering and Mechanics, Lanzhou University, Lanzhou, China

²College of Atmospheric Sciences, Lanzhou University, Lanzhou, China

³The Ministry of Educational Department, Key Laboratory of Mechanics on Disaster and Environment in Western China, Lanzhou, China

25 1 Introduction

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Snow plays an important role in Earth's climate system because of its wide coverage and seasonal variation, leading to variable surface conditions. Sublimation is a significant processway for snow surface to exchangetransfer the heat, mass, and energy with the atmosphere. Snow sublimation includes static surfacesnow cover sublimation and dynamical airbrone particlesnow particles sublimation—in the air. The latter process usually happens incontains drifting and blowing snow (DBS)sublimation (DSS)or blowing snow sublimation (BSS), in which snow particles follow the air flow, driven by the wind. Water vapor transport created by snow sublimationDSS has a significant influence on the local hydrological cycle and distribution, especially in the polar region and high alpine regions. For example, in the coastal area of Antarcticaeoastline area, ice sheet masswater loss caused by DBSdrifting and blowing snow reaches 18.3 % of the whole DBSdrifting and blowing snow amount each year (Pomeroy and Jones, 1996). In Antarctica, snow sublimation depleted approximately 17 - 20 % of its annual precipitation (Déry and Yau, 2001). In Mongolia, snow sublimation depleted 20.3 - 21.6 % of annual snowfall (Zhang et al., 2008). On In the Tibetan Tibet Plateau, due to its extremely dry, cold, and windyspecial environmental condition, the sublimation amount is very severely high, up to and it takes about 50 % of the amount of snow cover every year (Ueno et al., 2007).

DBSWindblown snow sublimation fluxes (40 - 60 W/m²) are more than twice as high as surface sublimation fluxes (20 - 30 W/m²) (Pomeroy and Essery, 1999). DBS sublimation is obviously stronger than surface snow sublimation for: 1) The turbulence is stronger during DBS events. 2) Aerodynamically entrained particles from the surface enlarge the contact surface with air. 3) The relative humidity decreases during DBS, which promotes a faster sublimation process. Therefore, investigating the role of sublimation in DBS is the forehead to accurately assess the water equivalent and understand the interaction between land surface and atmosphere in cold areas, especially for polar regions.

In DBS, mass, momentum, and energy are transferred between the surface and atmosphere accompanied by snow particle movement. Snowfall is the initial source of snow particles on the ground. Once the snow particles deposit on the surface, they either start rolling on the ground, saltating near the surface entrained by the wind force, or suspending in high air when wind speed exceeds the threshold value. During particle saltating, they jump on the ground surface and may bring up more particles to start to move, which is called splash. The above processes have been well described in the current numerical models of DBS (Pomeroy and Male, 1992; Taylor, 1998; Lehning et al., 2008; Vionnet et al., 2014; Sigmund et al., 2021; Yu et al., 2022; Melo et al., 2022). Early saltation models are usually empirical mass transport equations, which are functions transportation functions related to surface shear stress (Pomeroy et al., 1993; Déry and Yau, 2001). These models are susceptible to the empirical parameters. JDoorschot et al. (2004) developed a numerical model for steady-state saltation by considering the aerodynamic entrainment and rebound processes, which shows a more physical picture. Nemoto and Nishimura (2004) developed a new numerical model for saltation and suspension that considers aerodynamic entrainment, grain-bed collision, and wind modification processes with a distribution of grain sizes. Based on their model, a few improved works by

Zhang and Huang (2008), Wang and Huang (2017), Yu et al. (2022), Hames et al. (2022), and Melo et al. (2022) are carried out. However, the role of sublimation in drifting snow has not been demonstrated in these models.

For a single snow particle sublimation process, the sublimation rate is well described in the model proposed by Thorpe and Mason (1966), in which the sublimation rate is related to the particle size and environmental conditions. This T-M model has been implemented in the current snow model to estimate the sublimation in DBS, to assess the amount of sublimation amount caused by DBS on a small scale (meters) (Groot Zwaaftink et al., 2011; Huang and Shi, 2017; Dai and Huang, 2014; Vionnet et al., 2014; Sharma et al., 2018) and middle scale (kilometers) (Sharma et al., 2023; Gadde and Berg, 2024). The numerical simulation results from the above models all show that snow sublimation is an important snow physical process that cannot be ignored in DBS. In these small-scaled snow models, they track each single snow particle's trajectory, by considering particle diameter is unchanged. Also, the parameterization of sublimation in those middle-scaled snow models is from these small-scaled snow models, based on the same assumption.

In fact, snow particles, one of the most easily deformable and fragile granular systems, inevitably undergo fragmentation - a process that occurs during saltating, where particle-particle and particle-surface collision cause snowflakes to break apart and deform into smaller particles. This fragmentation not only alters their dynamic behavior by changing particle size (Sato et al., 2008; Walter et al., 2024) but also significantly impacts the sublimation rate, as for the sublimation of snow particles is closely linkedrelated to their particle size, shape, and specific surface area of particle (Domine et al., 2009). For a moving particle in the air, the reduction in size of a moving particle in the air is a joint effect of breakage and sublimation. In turn, the dynamically varying size of snow particles will affect their moving rules such as changing their trajectories, which further influences the mass flux and sublimation rate. However, this mutual physical feedback caused by fragmentation in DBS has never been reported, and the relevant model is still lacking, when snow particles move in the air, their size will not only be diminished due to sublimation but also

Sublimation changes with particle size, meanwhile, the sublimation rate is affected by the particle size. When snow particles move in the air, their size will not only be diminished due to sublimation but also from breaking due to collision with the surface and other particles. The phenomenon wherein particles collide with a surface during saltation, resulting in the separation and deformation of the snowflakes into smaller-sized snow particles, is referred to as fragmentation. Fragmentation not only changes the size distribution of the snow granular system but also changes the snowpack's physical properties such as the surface albedo of snow cover (Domine et al., 2006).

To date, tThere is only one model (Comola et al., 2017) considering the fragmentation of the snow particles during the drifting snow. That study, This work using a statistical mechanics model, calculatesealeulated the fragmented number of particles from the perspective of energy and mass balance and analyzes the effect of fragmentation on the particle size distribution. However, it did not explore the impact of fragmentation on drifting snow flux or the subsequent sublimation of snow particles. and simply analyzed the effects of fragmentation on the particle size

distribution. However, in that work, the influence of fragmentation on the drifting snow flux and subsequent snow sublimation has not been investigated.

In this work, we introduce the snow fragmentation model into the drifting snow model, enabling a more realistic representation of the movement and dynamic size changes of individual particles in the air. This advancement allows for a more accurate prediction of snow particle sublimation rates, offering critical insights into the micro-scale processes that govern snow-atmosphere interactions.

It is therefore essential to incorporate the fragmentation process into a physical numerical model to quantitatively assess its effect on snow particle sublimation, especially from the micro-scale perspective of individual particles experiencing dynamic size changes during their movement.

In previous numerical drifting snow models, snowflakes have typically been represented as solid spheres with a constant diameter. However, given that snowflakes are an easily deformable and fragile granular material, their diameters continuously change during saltation in the air (Sato et al., 2008; Walter et al., 2024). Consequently, the impact of the dynamically evolving size distribution of snow particles cannot be ignored in the numerical simulation. This is particularly crucial for accurately estimating sublimation rates and quantities, as a precise characterization of the dynamically changing size of particles is fundamentally important.

2 Model description

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The model by Huang and Shi (2017) is a simulation framework for DBS, which is able to simultaneously describe the behavior of both saltation and suspension particles. In our model, the saltation particles are described using an Euler-Lagrangian tracking method, which captures the saltating motion of particles. For suspension particles, which are typically smaller in size, we employ a dispersion function to characterize their movement dynamics. A threshold grain size was used to separate the saltating and suspended particles. The Thorpe and Mason model (Thorpe and Mason, 1966) is used to calculate the sublimation of DBS. The feedback of particle motion and particle sublimation to the wind field, air temperature, and air humidity are also considered. A threshold grain size iswas used to separatejudge the saltating and suspended particles. T-M model iswas used to calculate the sublimation of DBS. Additionally, the feedback of particle motion and particle sublimation to the wind field, air temperature, and air humidity arewere also considered. However, the particle fragmentation and its interaction with the wind field remain unexplored.

Comola's model, as the only one drifting model that considers snow particle fragmentation, is a one-dimensional, non-CFD (computational fluid dynamics) statistical approach. While it incorporates particle fragmentation, it does not couple the particles with the wind field. As a result, the effects of fragmentation on the wind field cannot be evaluated in their framework.

Building on the previous model developed by Huang and Shi (2017), we incorporated the fragmentation model proposed by Comola et al. (2017) and set up a comprehensive DBS model. This new model addresses the limitations

of Comola's model and provides a more comprehensive description of the interactions between particle fragmentation and the wind field. The particle fragmentation is newly taken into consideration in the saltation splash process. By explicitly considering particle fragmentation in the splash process, our model offers a more complete representation of the dynamics of DBS.

Here, we reintroducereintroduced them briefly.

2.1 Air flow

Considering the steady state of saltation, the horizontal wind field satisfies the following equations (Nemoto and Nishimura, 2004):

$$\frac{\partial}{\partial z}(\rho_f \kappa^2 z^2 | \frac{du}{dz} | \frac{du}{dz}) + f_p = 0 \tag{1}$$

where z is height aboveover the surface, ρ_f is the air density, κ is the von Karman constant, u is the wind speed, and f_p is the feedback force of the airborne snow particles.

The air temperature and humidity equations are Eq. (2) and (3), respectively, satisfying the horizontal uniformity condition are formulated according to follow (Bintanja, 2000):

$$\frac{\partial \theta}{\partial t} = \frac{\partial}{\partial z} (K_{\theta} \frac{\partial \theta}{\partial z}) + S_{\theta} \tag{2}$$

where θ is the air potential temperature, $K_{\theta} = \kappa u_* z + K_T$ is the turbulent heat exchangeheat turbulent coefficient, K_T is the molecular diffusion coefficients of heat, and S_{θ} is the sublimation heat feedback of the airborne snow particles u_* is the friction velocity.

$$\frac{\partial q_v}{\partial t} = \frac{\partial}{\partial z} (K_q \frac{\partial q_v}{\partial z}) + S_q \tag{3}$$

where q_v is the water vapor mixing ratio, $K_q = \kappa u_* z + K_v$ is the water vapor turbulent exchange coefficient, K_v is the molecular diffusion coefficient of water vapor, S_q is the sublimation humidity feedback of the airborne snow particles.

2.2 Snow saltation

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The motion of saltating snow particles can be described as five sub-processes, which are aerodynamic entrainment, particle trajectory, splash function, sublimation, and feedback to air.÷

(1) Aerodynamic entrainment

The snow particles start to move when the wind speed reaches a critical errain value (namely fluid threshold, usually presented by friction velocity) for a given errain snow surface; this, which is called aerodynamic entrainment.

The rateratio of aerodynamic entrainment is known as a linear function of the surface shear stress τ (Anderson and Haff, 1991):

$$155 \quad N_a = A(\tau - \tau_t) \tag{4}$$

where N_a is the number of aerodynamic entrainment particles per unit area per unit timesecond, $A[N^{-1}s^{-1}]$ is an empirical coefficient, and τ_t is threshold surface shear stress. The particle size distribution f(d) follows athe Gamma Distribution:

$$f(d) = \frac{d^{\alpha - 1}}{\beta^{\alpha} \Gamma(\alpha)} e^{-d/\beta} \tag{5}$$

160 where α and β are shape and inverse scale parameters, d is particle diameter.

There are two ways to describejudge the motion pattern of snow particles. The most used way is to define $\frac{1}{2}$ at threshold height to divide the saltation $\frac{1}{2}$ and suspension layerslayer, which is easy to apply. However, the threshold value is empirical and varies significantly in the existing $\frac{1}{2}$ big differences among different literature $\frac{1}{2}$ literatures. The other way is determined by the particle size, which is based on the traceability followability. The lifting velocity of aerodynamic entrainment particles is set to $\frac{1}{2}$ (Dai and Huang, 2014), where g is gravitational acceleration, which is not sensitive to the steady state of saltation.

(2) Particle's trajectory

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After the snow particles are lifted into the air, their ballistic trajectories can be described by Newton's second law:

$$170 \quad m\frac{d\mathbf{r}}{dt} = \mathbf{f_d} - m\mathbf{g} \tag{6}$$

where m is the mass of the snow particle, r(x,z) is the location of the particle, f_d is the drag force by fluid, g is the gravitational acceleration.

(3) Splash function

In this model, we use the probability functions to describe the movement after particlesparticle impacting the surface (Sugiura and Maeno, 2000). The restitution coefficient in the vertical direction $S_v(e_v)$, the restitution coefficient in the horizontal direction $S_h(e_h)$ and the number of particles ejected from the surface n_e are defined as followsfollows these functions:

$$S_v(e_v) = \frac{e_v^{a-1}}{b^a \Gamma(a)} e^{-\frac{e_v}{b}} \tag{7}$$

$$S_h(e_h) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{(e_h - \mu)^2}{2\sigma^2}}$$
(8)

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$$n_e = C_m^{n_e} p^{n_e} (1-p)^{m-n_e}$$
 (9)

whereinwhich v_{ix} is the horizontal component velocity of the impacting grain, v_{iy} is the vertical component velocity of the impacting grain, v_{ex} is the horizontal component velocity of the ejected grain, and v_{ey} is the vertical component velocity of the ejected grain, v_{ex} is the horizontal recovery coefficient, v_{ex} is the number of ejected grain grains. v_{ex} is the horizontal recovery coefficient, v_{ex} is the number of ejected grain grains. v_{ex} is the probability distribution function of v_{ex} is the grain grain, v_{ex} is the probability distribution function of v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the probability distribution function of v_{ex} is the probability distribution function of v_{ex} is the grain v_{ex} is the vertical grain, v_{ex} is the ve

Particle fragmentation was not considered in Sugiura and Maeno's splash function. We added it when a particle-bed collision happens. Thethe probability for particles breaking break after impact (Comola et al., 2017) is calculated according to Comola et al. (2017):as

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$$p(f) = 1 - \frac{1}{\sqrt{1 + \frac{\sigma^2}{w_s}}}$$
 (10)

wherein which $\sigma^2 = 2.5u_*^2$ is the turbulence velocity variance which refers to the fluctuations in velocity caused by turbulent eddies (Stull, 1988). w_s is the settling velocity of snow particles. When a snow particle falls back to the ground (initial velocity $v_i > 0.5 \text{ m/s} m/s$), use Equation 10 is used to determine whether it breaksis broken, and then calculate the number of snow particles N is calculated by Eq. (11), and λ is the ratio of particle size before and after fragmentation, again following Comola et al. (2017), by Eq. (12):as (Comola et al., 2017):

$$\begin{cases} N = 15v_i - 2.5, & 0.5 < v_i < 1.5 \\ N = \frac{5}{7}(6v_i + 19), & v_i > 1.5 \end{cases}$$
(11)

$$\begin{cases} \lambda = -0.4v_i + 0.7, & 0.5 < v_i < 1.0 \\ \lambda = -0.1v_i + 0.4, & 1.0 < v_i < 1.5 \\ \lambda = 0.25, & v_i > 1.5 \end{cases}$$
(12)

The velocity and the direction angle of the newly produced snow particles after crushing is keptthe same as that the size of the original snow particles. that are broken, and the direction angle of the exit velocity is the same as that of the original snow particles.

(4) Sublimation

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The drifting snow sublimation is calculated using by the Thorpe and Mason model (Thorpe and Mason, 1966):

$$\frac{dm}{dt} = \frac{\pi d(RH - 1) - \frac{Q_r}{KNuT} (\frac{L_s}{R_vT} - 1)}{\frac{L_s}{KNuT} (\frac{L_s}{R_vT} - 1) + \frac{R_vT}{ShDe_s}}$$
(13)

where m is the particle mass, d is the particle diameter, T is the air temperature, RH is the air relative humidity, Q_r is the solar radiation which snow particles absorbabsorbed, K is the heat conductivity, R_v is the gas constant of water vapor (461.5 J/kg//K), D is the molecular diffusivity of water vapor, e_s is the saturated vapor pressure relative to the ice surface, Nu is the Nussel number and Sh is the Sherwoods number.

(5) Feedback to air

The airborne particles have a significant effect on airflow. To consider this effect, we used an equivalent body force one grid is calculated as:

$$f_{p} = -\frac{1}{V} \sum_{i=1}^{N} f_{di} \tag{14}$$

where f_p is the feedback body force of the airborne particles in each control volume mesha vertical grid, V is the volume of the this vertical grid, N is the total number of the airborne particles in the this vertical grid, f_{di} is the drag force of the i_{th} particle in the grid.

The volume sublimation rate in each control volume meshgrid can be calculated as:

$$S = -\frac{1}{V} \sum_{i=1}^{N} \frac{dm_i}{dt} \tag{15}$$

where m_i is the mass of the i_{th} particle in the grid.

Then, the sublimation feedback to air temperature and humidity are:

$$220 \quad S_{\theta} = -\frac{L_s S}{\rho_f C} \tag{16}$$

$$S_q = \frac{S}{\rho_f} \tag{17}$$

where $L_s = 2.84 \times 10^6$ J/kg is the latent heat of sublimation, S is the volume DBS sublimation rate, and C = 1006 J/kg/K is the specific heat of air.

225 2.3 Snow suspension

In the simulation, we define a diameterradius threshold for distinguishingjudging suspension and saltation particles (Huang and Shi, 2017), the threshold diameterradius is calcualted based onby the Rouse number:

$$R_N = \frac{w_s}{\kappa u_*} \tag{18}$$

in which, u_* is the friction velocity and w_s is the particle deposit terminal velocity. Therefore, the conditions for determining the saltation and suspension snow particles are (Scott, 1995):

$$\begin{cases} R_N > 1, saltation \\ R_N \leq 1, suspension \end{cases}$$

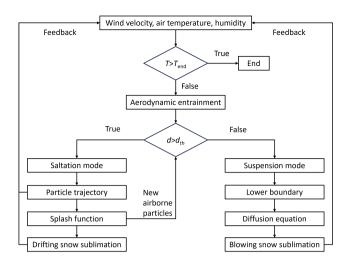


Figure 1. Flowchart of the new drifting and blowing snow model.

The suspended snowsnows follow the vertical diffusion equation (Déry and Yau, 2002):

$$\frac{\partial q_s}{\partial t} = \frac{\partial}{\partial z} (K_s \frac{\partial q_s}{\partial z} + w_s q_s) + S \tag{19}$$

where q_s is the suspended snow particle mixing ratio, $K_s = \delta \kappa u_* z$ is the diffusion coefficient of suspended particles, w_s is the terminal velocity. The δ follows (Csanady, 1963):

$$\delta = \frac{1}{\sqrt{1 + \frac{\beta^2 f^2}{w'^2}}} \tag{20}$$

where $\beta = 1$ is the proportionality constant, w' is the vertical turbulent fluid velocity, and we set $\overline{w'^2} = u_*^2$. A flowchart of the model is shown as Fig. 1 to illustract the present how this new splash function works.

2.4 Model verification

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To verify the model, we compared simulated particle size distribution, sublimation rate, and mass concentration to observational data. We first compare the particle size distribution at all heights with field observation data of Gordon and Taylor (2009) and Nishimura and Nemoto (2005), shown in Fig. 2. The simulated results are in consistent with size distribution range and variation trend, compared to field observation. Overall, the size distribution variation with height deviates more significantly in the near-surface (0.02 m to 0.1 m), compared to that in higher space (0.12 m to 1.13 m). The proportion of smaller-sized particles increases when considering the fragmentation (blue columns in Fig. 2), which is closer to the observation results (white columns in Fig. 2). For particles in height between 0.12 m to 1.13 m, particle size is in a narrow range of 0-90 μ m. In contrast, within the saltation layer (up to 0.1 m in height), particle sizes display a broader distribution, ranging from 50-450 μ m, and the simulated average particle size decreases with increasing height. However, this trend is not evident in the field, which might be due

to the complexities of the field environment compared to ideal simulation, as well as limitations in the accuracy of measurement sensors.

Fig. 3 shows the comparison of the total sublimation rate between the numerical simulation and the field observation and the field observation, and the gray curve is the simulation result of the sublimation rate under the same conditions. Comparing thesesthe two curves, we seecan find that the sublimation rate is the same on the order of magnitude, which shows that the model is suitable for reliable in calculating the sublimation rate of blowing snow. Fig. 4 shows Tthe mass concentration from temperature is shown in Fig. 4. The red point is the snow-blowing mass concentration data on (measured by Pomeroy and Male (1992) near Saskatoon), and the black curve is the result calculated in this paper.

The suspension sublimation rate from between this paper and from other sublimation models under the same conditions are were also compared (Fig. 5). The black line is the sublimation rate of suspension in the case of fragmentation of snow particles, and the other five curves are the simulation results of suspension sublimation of jump without breaking of snow particles (Xiao et al., 2000; Huang and Shi, 2017). The results demonstrate that drifting snow sublimation is important, particularly in the near-surface saltation layer. However, most previous models underestimate the sublimation rate near the surface, which significantly impacts the assessment of the drifting snow sublimation. Accounting for fragmentation increases the sublimation rate by approximately 1.3 times, which suggests that it is necessary to conclude snow particle fragmentation in DBS models. The results show that the fragmentation of snow particles enhances the sublimation of suspension snow particles.

3 Results

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3.1 Fragmentation effects Effects on particle Particle size Size distribution Distribution

The fragmentation of snow particles first leads to changes in their size distribution, releasing numerous smaller particles. Simulations were conducted to analyze the effect of fragmentation on particle size distribution.

Simulations are conducted with a same The friction velocity $u_* = 0.45$ m/s, and a same The initial mean particle size $\bar{d} = 200 \ \mu \text{m}$. As shown in Fig. 6, the size distribution pattern for particles, without considering fragmentation, follows the log-normal distribution function (blue bar). When considering fragmentation, the proportion of smaller-sized particles (< 100 μm) increases, while the overall proportion of larger-sized particles decreases. This results in a decrease in the average particle size.

3.2 Fragmentation effects on snow particle number Effects on Snow Particle Number

Fig. 7 presents the temporal evolution of the concentration of snow particles suspended in the air. It is observed that the number of saltation particles increases over time until reaching a steady state, regardless of the presence

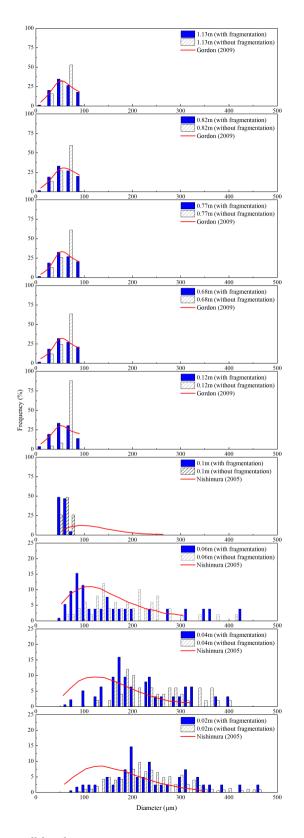


Figure 2. Particle size distribution at all heights.

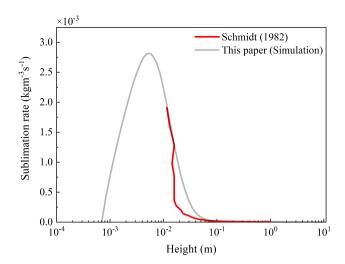


Figure 3. Comparison of total sublimation rate for this paper and field observations. (Schmidt, 1982) ($u_* = 0.63 \text{ m/s}, z_0 = 7 \times 10^{-4} \text{ m}, T = 267.45 \text{ K}$)

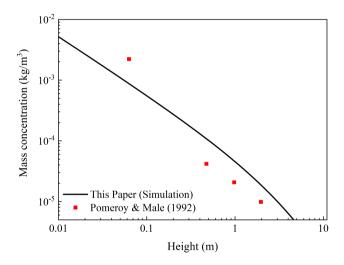


Figure 4. Comparison of mass concentration for this paper and field observations. Observation (Pomeroy and Male, 1992) ($u_* = 0.31 \text{ m/s}, T = 265 \text{ K}$)

of fragmentation. It is noted that when fragmentation processes are taken into consideration, the steady-state concentration of snow particles is consistently higher at all wind speeds. Under low wind speed ($u_* = 0.3 \text{ m/s}$), the amount of particle number increases by 42 %. Under high wind speed ($u_* = 0.5 \text{ m/s}$), the amount of particle number increases by 26 %. This suggests that the fragmentation contributes to the total number of snow particles in the air. Furthermore, Fig. 7 also reveals that the increase of particle number resulting from fragmentation is notably more pronounced at lower wind speeds. However, the increase rate of the particle number is lower than that of the

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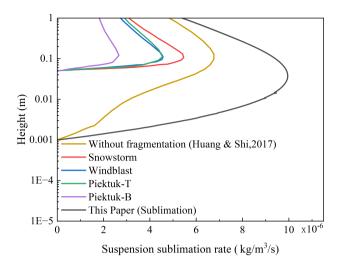


Figure 5. Comparison of suspension sublimation rates with other blowing snow models. (Xiao et al., 2000; Huang and Shi, 2017) ($u_* = 0.87 \text{ m/s}$, $z_0 = 0.001 \text{ m}$, $T_0 = 253.16 \text{ K}$)

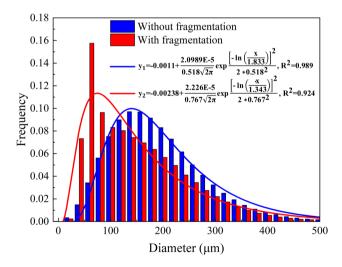


Figure 6. Size distribution with and without considering the particle fragmentation.

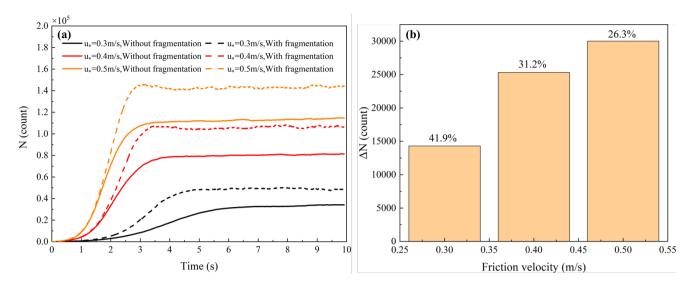


Figure 7. (a) Saltating particle number variation with time underin the wind conditions condition of $u_* = 0.3$, 0.4 and 0.5 m/s. (b) Increment number and ratios of saltation snow particles in the air.

high wind speed. This indicates that at higher wind speeds, the degree of fragmentation becomes more intense due to the transfer of kinetickinematic energy to the internal energy of the particles. Therefore, during the processes of particle-bed surface and particle-particle collisions, the release rate of particles from snow particle fragmentation is higher under stronger wind conditions.

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3.3 Fragmentation effects on mass concentration and mass flux Effects on Mass Concentration and Mass Flux

For near-surface saltation particles, it is illustrated in Fig. 8 illustrates that the variation in the mass concentration of saltating and suspended particles with height. The fragmentation of snow particles enhances the concentration of both saltating and suspended particles, at levels close to the surfaceespecially at lower altitudes. Fig. 8(a) depicts the changes in mass flux versus height above the surfacealong the vertical axis, showing that fragmentation enhances the transport of saltating particles near the ground-surface. This is because the fragmentation of snow particles increases the number of air-borne saltation particles, more saltation particles take part in the splash process, further increasing therefore the air saltation particle number increases further. When the friction velocity varies from 0.3 m/s to 0.5 m/s, the increment proportion of fragmentation mass concentration increases from 19 % to 3 %, which means the fragmentation only has strong effects on the mass concentration under weak wind conditions.

For particles suspended further alofthigh-air suspended particles, it is shown in Fig. 8(b) that the mass concentration of the suspendedsuspension snow particles at the same height in the air is higher and the overall suspension height is higher, when considering snow particle fragmentation. This is because of smaller and lighterit is small particles.

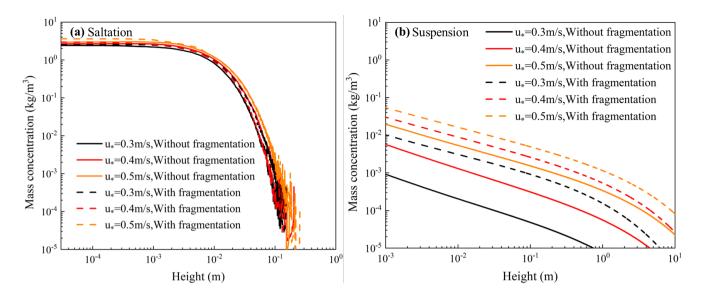


Figure 8. Mass concentration of (a) saltating and (b) suspension particles with/without considering fragmentation, under in the wind conditions conditions of $u_* = 0.3$, 0.4 and 0.5 m/s.

cles created by that can be released from the snow fragmentation, which have higher more possibility to be entrained enter and being suspended to in higher levels air.

Under the same friction velocity, the mass flux of near-surface (< 0.01 m) snow particles is larger with snow fragmentation higher than without that when considering the snow fragmentation, as is shown in Fig. 9. It is concluded that the fragmentation increases the particles transportation of snow particles transport.

3.4 Fragmentation effects on sublimation rateSublimation Rate

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The sublimation rates of saltating and suspended snow particles increase with fragmentation implemented in the model, as shown in Fig. 10. This enhancement is more significant at lower friction velocity wind speeds, indicating that snow particle fragmentation has a more profound effect on sublimation under such conditions.

When the wind speed is 0.3 m/s, the average sublimation rate of saltation particles increases by 20 % due to fragmentation, from $1.56\times10^{-2}~kg/m^2//s$ to $1.87\times10^{-2}~kg/m^2//s$. However, as the wind speed increases to 0.5 m/s, this increase drops to 3 %, from $4.37\times10^{-2}~kg/m^2//s$ to $4.49\times10^{-2}~kg/m^2//s$, indicating that the impact of fragmentation on the sublimation rate diminishes under stronger wind conditions. This trend can be attributed to the fact that higher wind speeds enhance particle transport and mixing, which reduces the relative contribution of fragmentation to the overall sublimation process.

A similar trend is observed for suspension particles at higher altitudes. At a wind speed of 0.3 m/s, the average sublimation rate of suspension particles increases by 8 times, from $1.09 \times 10^{-6} \ kg/m^2//s$ to $9.8 \times 10^{-6} \ kg/m^2//s$ when fragmentation is considered. As the wind speed increases to 0.5 m/s, this growth decreases to 50 %, from

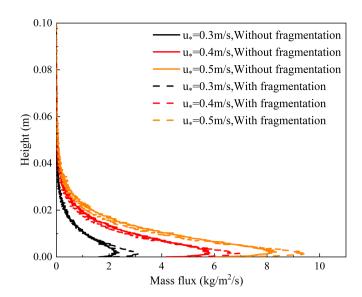


Figure 9. Vertical mass flux distribution of snow particles with/without considering fragmentation, under the wind conditions conditions of $u_* = 0.3, 0.4, \text{ and } 0.5 \text{ m/s}.$

320 $3.7 \times 10^{-6} \ kg/m^2//s$ to $5.7 \times 10^{-6} \ kg/m^2//s$. While the effect of fragmentation on sublimation remains significant at higher wind speeds, the reduction in growth indicates that other factors, such as increased turbulence and particle dispersion, may play a more prominent role in driving sublimation under these conditions.

Overall, snow fragmentation has a more pronounced effect on sublimation rate of suspension particles than for saltation particles. This difference can be attributed to the longer residence time and greater exposure of suspension particles to the airflow, which amplifies the impact of fragmentation on their sublimation rates. In contrast, saltation particles, which are closer to the surface and are subject to more frequent impact and splash processes, experience a relatively weaker influence from fragmentation as wind speed increases.

The relative sublimation rate of saltation particles enhances increases 16 % to 6 % when the friction velocity increases from 0.3 m/s to 0.5 m/s, considering the fragmentation effects. For sublimation particles, the sublimation rate also decreases with friction velocitywind speeds. The sublimation rate of suspension particles increases from 81 % to 70 %, considering fragmentation, which is much higher compared to that snow fragmentation has a more dominant effectfundamental effects on the sublimation process for suspension than for salta

3.5 Effects of size distribution on fragmentation Size Distribution on Fragmentation

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335 TheHere, the averaged particle diameter is defined as $\overline{d} = \alpha \times \beta$, where α and β are the shape and scale parameter, respectively. α adjusts the peak position and β controls the width of the distribution (the higher the value, the wider the size distributes).. In which, α is the shape parameter. α adjusts the peak position of the curve and the steep

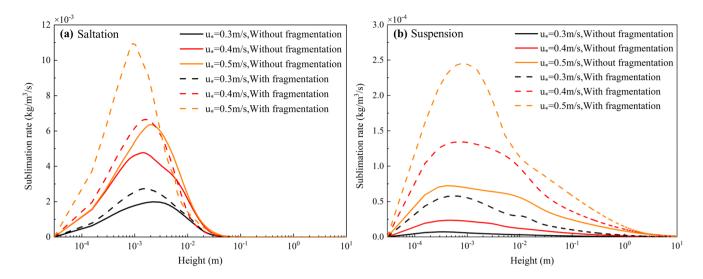


Figure 10. Sublimation rate versus varies with height. (a) Saltation particles. (b) Suspension particles.

extent of the curve. β is scale parameter, and they control width of size distribution (higher the value is, the wider the size distributes).

340 3.5.1 Average particle diameter Particle Diameter

As is shown in Fig. 11(a) and α (b), with the increasing value of α or β , the number of fragmented snow particles increases, and the fragmentation efficiency is higher. This indicates that the larger particles have a higher fragment extent, and this is because larger particles can produce more small snow particles.

3.5.2 Size proportion Proportion

We set up three cases with the same mean diameter $\bar{d}=200~\mu\mathrm{m}$ but different α and β , parameters. Here we chose three different values of β with the same mean diameter $\bar{d}=200~\mu\mathrm{m}$. These differences in α and β result in particle systems with varying size distributions. In these three cases, the proportion of particles with diameter larger than the threshold diameter is 73 % (blue), 85 % (green), and 96 % (purple), respectively. In these three cases, thethere proportion of the particles larger than the threshold diameter is 73 % (blue), 85 % (green), and 96 % (purple).

As shown in Fig. 11(c), the fragmentation is significant under the same mean averaged diameter when the large particles take a higher proportion in a granular system. The fragmentation number of snow particles with different particle size distribution increases almost linearly with the time. This is because, only snow particles with a larger size than threshold diameter, more snow particles will join the fragmentation.

4 Discussion

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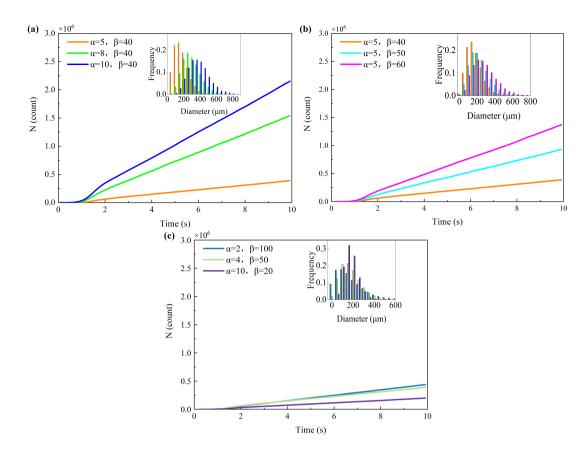


Figure 11. Number of snow particles as a function of varies with time. (a) variation of shape parameter α , scale parameter β = 40, const. Scale parameter β = 40. (b) variation of shape parameter β , scale parameter α = 5, const. Shape parameter α = 5. (c) Average diameter \bar{d} = 200 μ m.

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In this study, we investigated how particle fragmentation affects snow transport and sublimation processes. We found that when fragmentation is considered, the particle size distribution is modified, leading to an expected increase in snow sublimation during DBS. This occurs because fragmentation generates a greater number of smaller-sized particles, which increases the average specific surface area of transported snow. Since the sublimation rate is directly proportional to the specific surface area, the presence of smaller particles enhances sublimation. Additionally, the smaller particles produced by fragmentation reduce the averaged diameter of surface snow particles, leading to a lower threshold velocity of aerodynamic entrainment. As a result, more particles are lifted from the surface and transported by wind, further increasing the number of particles available for sublimation during transport. Moreover, the sublimation rate is directly proportional to the mass concentration of drifting snow particles, amplifying the overall sublimation effect.

However, fragmentation has a limited impact on the overall mass flux profile, with notable changes primarily observed in the near-surface layer. Specifically, when fragmentation is considered, the mass flux below 0.01 m

increases. This increase is mainly attributed to changes in mass concentration of suspension particles, for the mass concentration of saltation particles remains largely unchanged. This indicates that fragmentation predominantly occurs within the saltation layer but primarily contributes to the mass in the suspension layer. This is because smaller fragmented particles are more easily transported by wind forces, exhibiting greater flowability in air and reaching higher altitudes.

Moreover, the fragmentation of the snowflakes produces smaller-sized particles that remain in saltation and suspension, while deposited particles alter the surface size distribution. These changes in particle size distribution influence snow surface properties, such as albedo (Manninen et al., 2021), snow microstructure, static snow cover sublimation rate (Albert and Mcgilvary, 1992), and surface roughness. Larger snow particles reduce multiple scattering events because light travels longer paths within larger particles, leading to increased absorption, particularly in the near-infrared spectrum. Thus, larger snow particles have a low snow albedo. Smaller snow particles, with their higher specific surface area, allow light to undergo more scattering events within the snowpack and reduce the absorption of solar radiation, typically having a high albedo, especially in the visible spectrum. The variation of snow surface size distribution due to fragmentation in DBS influences the surface energy balance by changing the snow surface albedo. Additionally, smaller grains affect the snow thermal conductivity, mechanical stability, and the retention of impurities, which can further reduce albedo and accelerate snowmelt. Therefore, this variation of snow surface properties plays a critical role in determining the energy exchange between the snowpack and the atmosphere.

These findings highlight the importance of incorporating particle fragmentation processes into spatially distributed surface energy balance models. This can improve the accuracy of snowpack mass balance assessments, enhance predictions of seasonal snow dynamics, and better represent snow transport and sublimation processes in atmospheric and climate models.

5 Conclusion

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In this study, we carried out a numerical simulation to investigate the snow particle fragmentation and sublimation during DBS. The model is based on an Euler-Lagrangian method to track the trajectories of individual snow particles. To account for sublimation processes, we implemented the T-M model, which calculates the sublimation rate of snow particles based on their size, temperature, and the surrounding environmental conditions.

The simulation incorporates key physical processes in DBS, including particle-particle interactions, fragmentation due to collisions. The model was validated with experimental data from previous studies. This integrated method provides a detailed understanding of the dynamics of snow particles and their sublimation during snow particle transport.

In this study, we developed a drifting snow model that incorporates the snow particle fragmentation process. This model simultaneously accounts for both the dynamic processes, including the movement of saltation and suspension particles, and the thermodynamic processes, such as snow sublimation. The model was validated using experimental

data from previous studies, ensuring its reliability. This integrated approach, the model offers a comprehensive understanding of snow particle dynamics and sublimation during transport in DBS events.

Based on this model, this work investigates investigated the significant role of snow particle fragmentation in DBS drifting snow and blowing snow. We find found that fragmentation not only alters the particle size distribution but also increases the number, concentration, and mass flux of particles in DBS drifting and blowing snow. Subsequently, this phenomenon affects the sublimation rate of airborneair-moving snow particles. Specifically, Our results show that fragmentation reduces diminishes the average particle size on average, creating smaller particles that are more prone to sublimation which provides more opportunities for smaller particles to sublimate. The effects of fragmentation on sublimation are more pronounced for suspension particles than saltation particles, particularly under low wind conditions.

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CondiscringWith consideration of particle fragmentation, the amount of particle mass concentration and the sublimation rate increases by one order of magnitude indicatingwhen considering particle fragmentation, which indicates fragmentation of saltation snow particles is a key source forone of the fundamental sources of suspended snow particles. The sublimation rate of saltating snow particles increases by 11 % on average, and that of suspension snow particles by hinereases 76 % on average when the friction velocity is increased fromwind speed is between 0.3 m/s to 0.5 m/s. However,tThe differences in sublimation rate with/without considering fragmentation decrease with increasing wind speeds.

Investigating We also investigated the effects of particle size distribution of particles on the sublimation rates, we. We find found that the sublimation rate is enhanced for when the particles with a larger average diameter and a higher proportion of larger particles. These results underscore the importance of accurately representing fragmentation and size distributions in snow transport models.

Our simulation results are consistent with the previous observational data-observation results, suggesting-which indicates the validity of theour model. Furthermore, a comparison of simulation-compared with the calculating results considering or ignoring-from the model which does not consider fragmentation of snow particles, shows that our sublimation rates are 2-4 times higher than other-the previous model results. This is because fragmentation reduces diminishes the snow particle's size, and increases the number of airborne particles, which are more susceptible to sublimation and smaller snow particles are more prone to sublimate in the air. By integrating fragmentation into the numerical model, this study marks a significant step forward in understanding and quantifying the effects of particle dynamics on snow sublimation.

Our work provides insights into the complex dynamics of DBS. It provides a deeper understanding of the physical process of snow particle fragmentation during saltating/suspending in the air. This indicates the importance of fragmentation in the numerical modelsmodel of DBS. However, the usedthis model is a two-dimensional numerical model, which could not be applied to larger regions, especially for complex terrains. Therefore, the expansion implementation of this model into a three-dimensional drifting snow model in the future is necessary. Moreover,

crystal habits is another important factor in influencing the sublimation rate of snow particles, such as density, size, and specific surface area. Future numerical simulation should be carried out regarding crystal habits factors.

The simulation results provide detailed insights into the physical dynamic processes of particle-atmospheric momentum transfer, heat transfer, and mass transfer, from aeach single particle perspectiveaspect. This work provides supplies the theoretical foundation basement and prediction method methods for accurately assessing the accurate amount of snow sublimation during DBS. These findings have important implications for imporiving the representation of snow transport and sublimation processes in atmospheric and climate models, which can enhance predictions of snow mass balance and its broader environmental impacts., which should be further implemented in related atmospheric and climate models.

Author contributions. N H designed the conception, JC B contributed to the programming and numerical calculation, HX Y contributed to the conception, first draft and revision, G L designed the conception and revised the manuscript. HX Y contributed to the conception and first draft, JC B contributed to the programing and numerical calculation, N H and G L designed the conception and revised the manuscript.

Competing interests. All authors declare that no competing interests are present.

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