









 **Abstract.** Atmospheric planetary waves has long been studied, focusing on the equator and the middle and low latitudes. However, variation of polar planetary wave activity, especially temperature, temporal, and interannual variations of polar planetary waves, remain largely unexplored. In this study, we use MERRA-2 dataset to investigate the impact of atmospheric variations on polar planetary waves there during austral winter from 2002 to 2022. The temperature amplitude and wave periods of each polar planetary wave event were determined using 2-D least-squares fitting. Our results show that, as the zonal wavenumber increases, E1, E2, E3, and E4 occur earlier with weaker peak amplitudes and shorter wave periods. 24 The phase velocities of E1, E2, E3 and E4 are similar to  $\sim$  40 m/s in the polar atmosphere. In order to elucidate the variations in the observed polar planetary waves and its possible propagation and amplification mechanism, we carry out the diagnostic analyses with MERRA-2 reanalysis data from the surface up to ∼ 80 km. Results indicate that the polar planetary waves 29 can be amplified in instability of the mesosphere  $\sim$  50–60 $\degree$ S and the 30 stratosphere  $\sim$ 70–80°S. The mean flow instability of the strong polar planetary wave during the austral winter mesosphere is greater than that of the weak wave. The selective generation, propagation, and amplification of planetary waves with varying zonal wavenumbers due to variations in background zonal winds in the polar atmosphere. The temperature amplitude of polar planetary wave correlates with solar activity F10.7. The correlation between zonal wavenumber, wave period, and phase velocity implies that polar planetary waves propagate as fixed-phase wave packets.





- These results can suggest that, the background atmospheric conditions in the polar regions play a crucial role in modulating the generation, propagation, and amplification of planetary waves. Overall, we analyze their dynamics variation of eastward planetary waves in the polar atmosphere during the 2002-2022 austral winter periods and statistically analyze the interannual variation.
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## **1 Introduction**

 The winter stratosphere is primarily influenced by dominant large- scale planetary waves, with their interaction with the zonal mean flow serving as the main driving force for winter stratospheric dynamics. Planetary wave activity significantly influences the polar stratosphere thermal and dynamic structure, causing significant changes in wind, temperature, and composition. They also influence the horizontal transport and distribution of substances such as O3, H2O, and CH4 (Coy et al., 2003; Manney et al., 1998; Allen et al., 1997). Recent studies have specifically examined the significant eastward-propagating planetary waves with periods of approximately 2 and 4 days in the polar stratosphere and mesosphere (Tang et al., 2021; Rhodes et al., 2021; Lu et al., 2017; Lu et al., 2013; Alexander and Shepherd, 2010; Watanabe et al., 2009; Tunbridge and Mitchell, 2009; Sandford et al., 2008; Pancheva et al., 2008; Baumgaertner et al., 2008; Merzlyakov and Pancheva, 2007). Two types of zonal mean flow instabilities have been proposed as the sources of these planetary waves: the barotropic instability of the stratospheric polar night jet (polar night jet) and the barotropic/baroclinic instability of the double- jet structure (subtropical mesospheric jet). The winter polar night jet 64 consists of strong eastward zonal winds centered  $\sim$  50°S–60°S in the upper stratosphere, while the subtropical mesosphere jet is typically located at  $\sim$  30°S between ~50 and ~70 km.





 The polar planetary waves with zonal wavenumbers -1 (E1), -2 (E2), 68 -3 (E3), and -4 (E4) correspond to  $\sim$ 4,  $\sim$ 2,  $\sim$ 1.3, and  $\sim$ 1-day waves respectively. These waves can induce significant changes in wind and temperature in the polar stratosphere and mesosphere (Merzlyakov and Pancheva, 2007; Manney et al., 1998; Lawrence et al., 1995; Manney and Randel, 1993; Fraser et al., 1993; Venne and Stanford, 1979). They analyzed data from the Arctic Esrange (68°N, 21°E) meteor radar and the Microwave Limb Sounder (MLS) instrument on the EOS Aura satellite to investigate polar waves in the winter stratosphere, mesosphere, and lower thermosphere (Sandford et al., 2008). The E2 reaches maximum at the stratopause and may be generated by instability in the polar night jet. They believed that planetary waves are crucial for dynamically coupling different atmospheric layers and influencing ionospheric variability in polar atmosphere.

 They employed a middle atmosphere General Circulation Model 82 (GCM) to simulate the characteristics of the E1  $(-4-day)$  wave in the Antarctic winter mesosphere (Watanabe et al., 2009). They found that the eastward forcing from the E1 wave occurred within the double-jet structure, counteracting a portion of the westward forcing from gravity waves. This phenomenon contributes to stabilizing the mean airflow structure of the Antarctic winter mesosphere. They investigated atmospheric diurnal variability with periods ranging from 1.5 to 5 days using neutral meridional





 wind data from Esrange meteor radar, as well as atmospheric temperature and pressure data from SABER (Merzlyakov and Pancheva, 2007). Their findings showed that the Eliassen-Palm (EP) fluxes of waves from jet instability are predominantly directed downward. This suggests that dynamic influences from the mesosphere may impact the lower atmosphere. They observed that the eastward propagating planetary waves were predominantly confined to high latitudes during winter (Lu et al., 2013). This limitation could be attributed to the evanescent wave 97 characteristics resulting from a negative refractive index of  $\sim45^{\circ}$ S in the equatoria, thereby impeding their propagation towards lower latitudes. They investigated the seasonal variation of eastward E2 in 2007 using temperature and wind data from the Whole Atmosphere Community Climate Model + Data Assimilation Research Testbed (WACCM + DART) (Gu et al., 2017). Their findings showed that E2 mainly occurs in winter, with peak amplitudes in the stratosphere. Temperature amplitude, zonal 104 wind, and meridional wind in southern winters can reach amplitudes of  $\sim$ 10 105 K,  $\sim$  20 m/s, and  $\sim$  30 m/s respectively. Conversely, in northern winters, E2 exhibit only one-third of this strength in temperature amplitude compared to their southern counterparts. They employed WACCM simulations to investigate eastward planetary waves during a significant sudden stratospheric warming (SSW) event occurring in January 2009 within boreal (Rhodes et al., 2021). They believed that planetary and gravity





 waves caused eastward wind maxima in the stratosphere and mesosphere before the SSW in the stratosphere. This dual-maxima wind configuration promotes the planetary wave growth through over-reflection due to wind shear instability. They investigated the global variation of eastward propagating wavenumbers E1, E2, E3, and E4 planetary waves in the polar atmosphere using temperature and wind data from the 2019 Modern-Era Retrospective Research Analysis for Research and Applications (MERRA- 2) (Tang et al., 2021). Their findings revealed a slightly larger maximum amplitude in the southern hemisphere compared to the northern hemisphere. They observed that as the wavenumber increases, polar planetary wave peak at lower latitudes with smaller amplitudes. Additionally, diagnostic analysis suggested that mean flow instability in the upper stratosphere and upper mesosphere may contribute to enhanced polar planetary waves.

 The interannual variation of polar planetary waves in the polar atmosphere lacks comprehensive research. We analyzed polar planetary waves E1, E2, E3, and E4 during austral winter using MERRA-2 data from 2002 to 2022. Section 2 summarizes the MERRA-2 data and analysis methods used in this work. Focusing specifically on occurrence dates, peak amplitudes, and wave periods for each respective event (Sections 3.1 through 3.4). Comparison of E1, E2, E3, and E4 (Section 3.5). In Section 4, we discuss the variations in temperature amplitude, temporal, and





interannual of polar planetary waves and their dynamics. Section 5

provides the summary.

#### **2 Data and Analysis**

 The air temperature, potential temperature, zonal wind, meridional wind, vertical wind, air density, and model mid-layer height output were obtained from the MERRA-2 dataset with a 3-hourly temporal resolution 139 and a horizontal resolution of  $0.5^{\circ} \times 0.625^{\circ}$ . Recently upgraded using the Goddard Earth Observing System Model Version 5 (GEOS-5) data assimilation system (Gelaro et al., 2017; Molod et al., 2014; Molod et al., 2012), this dataset provides information from 72 different model levels above the surface, with particular focus on levels ranging from 9 to 15. The pressure at the topmost level of the model is set at 0.01 hPa. MERRA-2 has been extensively utilized in various studies investigating phenomena such as planetary waves in polar atmospheres, global thermal tides, climate variability, aerosol and ozone trends and processes (Zamora et al., 2022; Bahramvash Shams et al., 2022; Tang et al., 2021; Ukhov et al., 2020; Bali et al., 2019; Lu et al., 2013). Numerous recent studies have demonstrated that utilizing MERRA-2 data is feasible for our present study.

 Time windows of 10, 6, 4, and 4 days were selected to analyze planetary waves E1, E2, E3, and E4 using the least-squares method to obtain fluctuation information for each window. Subsequently, characteristic properties of the planetary waves, such as peak amplitude,







$$
y = A \cos[2\pi(\sigma \cdot t + s \cdot \lambda)] + B \sin[2\pi(\sigma \cdot t + s \cdot \lambda)] + C \qquad (1)
$$
  
160 The values of A, B, and C in Equation (1) were obtained by least-  
squares fitting. The values of A, B, and C in Formula (1) were obtained by  
162 least-square fitting. The frequency and zonal wavenumber are represented  
163 by  $\sigma$  and s. The longitude of the satellite samples and UT time are  
164 defined by  $\lambda$  and t. Planetary wave amplitude R can be expressed by  
165  $R = \sqrt{A^2 + B^2}$ .

166 The baroclinic/barotropic instability in the atmospheric structure 167 arises from the simultaneous equalization of the negative latitude gradient 168 and quasi-geostrophic potential vorticity. The angular speed of the Earth's 169 rotation (Ω), latitude (φ), zonal mean zonal wind  $(ū)$ , Earth radius (*a*), air 170 density ( $\rho$ ), Coriolis parameter ( $f$ ), buoyancy frequency ( $N$ ), subscripts 171 vertical (z) and latitudinal ( $\varphi$ ) gradients are denoted as in Equation (2).

172 
$$
\overline{q_{\varphi}} = 2\Omega \cos \varphi - \left(\frac{(\overline{u} \cos \varphi)_{\varphi}}{a \cos \varphi}\right)_{\varphi} - \frac{a}{\rho} \left(\frac{f^2}{N^2} \rho \overline{u}_z\right)_z
$$
(2)

 The Eliassen-Palm (EP) flux vectors (F) can be used to calculate the properties of planetary wave propagation (Equation 3). The planetary wave 175 perturbations in the zonal and meridional wind are represented by  $u'$  and  $v'$ , respectively;  $\theta'$  represents the potential temperature, while w'





177 represents the vertical wind.

178 
$$
F = \rho a \cos \varphi \left[ \left[ f - \frac{\frac{\overline{u}_z \overline{v'} \theta'}{\overline{\theta}_z} - \overline{v'} u'}{a \cos \varphi} \right] \frac{\overline{v'} \theta'}{\overline{\theta}_z} - \overline{w'} u' \right]
$$
(3)

179 The propagation of planetary waves is favorable only when the square 180 of the refractive index  $m^2$  is positive. The refractive index squared serves 181 as the waveguide for planetary waves, where *s* represents the zonal 182 wavenumber,  $c$  denotes the phase speed, and  $H$  stands for the scale 183 height. The equatorial linear velocity, denoted as  $v_0$ , T is related to the 184 wave period.

185 
$$
m^2 = \frac{\bar{q}_{\varphi}}{a(\bar{u}-c)} - \frac{s^2}{(a\cos\varphi)^2} - \frac{f^2}{4N^2H^2}
$$
 (4)

$$
c = -v_0 \cos\left(\frac{\varphi \pi}{180}\right) / sT \tag{5}
$$

#### 187 **3 Results**

 Figure 1 illustrates the temporal variations of polar planetary waves E1, E2, E3, and E4 observed during the 2002 austral winter. Figures 1a-1d illustrate the variations of polar planetary waves at different latitudes. The 191 observations show E1 at  $\sim$ 70-80°S and  $\sim$ 50 km, E2 at  $\sim$ 60-70°S and  $\sim$ 50 192 km, E3 at ~60-70°S km and ~50 km, and E4 at ~50-60°S km and ~50 km. The largest peak of E1 occurs on days 202–212 (∼3.125 day/~75 hr, ∼7 K), while the other three peaks occur on days 154–164 (∼3.75 day/~90 hr, ∼4 K), 184–194 (∼3.46 day/~83 hr, ∼4 K), and 260–270 (∼6.5 day/156 hr, ∼6 K). The maximum peak of E2 occurs on days 186–192 (∼32 hr, ∼4 K), with smaller peaks on days 198–204 (∼48 hr, ∼3 K) and days 208–214







 We presented the analysis results only for 2002 due to the representative of the wave activities for the entire range of 2002–2022. We thoroughly examined the wavenumber-period spectra for each event to ensure accuracy and validate our findings. Subsequently, we statistically constrained their occurrence date, wave period, and peak amplitude during 2002-2022. The temperature amplitude and wave periods of each event were determined using 2-D least-squares fitting.

## **3.1 E1**

214 Table 1 lists the 107 E1 events, occurring at  $\sim$ 70–80°S and  $\sim$ 50 km during the austral winter from 2002 to 2022. The E1 appeared earliest on days 142-152 in 2005, 2006, 2009, 2010, and 2021, while latest on days 258-268 in 2011, 2015, and 2018. The shortest wave period observed for E1 was ~53 hr between days 190-200 in 2012, while the longest was ~171 hr between days 242-252 in 2015. E1 shows the smallest peak amplitude





 (~4-4.5 K) on days 160-170 in 2003, 158-168 in 2010, and 142-152 in 2021. The largest peak amplitude of E1 reached ~13.5-14 K between days 214- 224 in 2003 and 194-204 in 2017. The scatter diagram in Figure 2 shows E1 variations in date-amplitude, period-amplitude, and date-period for the austral winter from 2002-2022. Larger E1 events occur in the middle of the austral winter, while smaller ones occur in the early and late. E1 events 226 with temperature amplitudes exceeding  $\sim$  8 K mainly occur within  $\sim$  72-120 227 hr, while those lasting longer than  $\sim$  144 hr have stronger temperature amplitudes but lower frequency. The E1 events with longer periods mainly occur in early austral winter, while those with shorter periods tend to happen later. Most temperature amplitudes associated with E1 events 231 surpass  $\sim$  5 K (Figures 2a, 2d, and 2g).

232 The histogram shows the distribution of E1 information (Figures 2b, 233 2e, and 2h). E1 was more frequent at 200-220 days (24 times) and less so 234 at 140-160 days (17 times) and 240-260 days (11 times). E1 events were 235 recorded 60 times between days 180–240, accounting for ~56% of the total. 236 The wave period of E1 is distributed between  $~48$  and  $~192$  hr, with 84 237 times occurring between  $\sim$ 72 and  $\sim$ 120 hr, accounting for  $\sim$ 78.5% of the 238 total, among which 49 times occur between  $\sim$ 72 and  $\sim$ 96 hr. Longer wave 239 periods for E1 are scattered between  $\sim$ 120-192 hr (11 times). The wave 240 period of E1 decreases from  $\sim$ 96 hr on days  $\sim$ 160 to  $\sim$ 70 hr on days  $\sim$ 200. 241 The mean amplitude of E1 is smallest at  $140-160$  days ( $\sim$  6.5 K), reaching







**3.2 E2**

 Table 2 lists the dates, periods, and amplitudes of E2 polar planetary 250 wave events at ~60-70°S and ~50 km during austral winters from 2002 to 2022. The table records 99 E2 events. The E2 appeared earliest on days 142-148 in 2003, 2017, and 2022, while latest on days 256-262 in 2005 253 and 2011. The shortest wave period of E2 is  $\sim$ 30 hr during days 214-220 and 238-244 in 2017. The longest period of ~58 hr occur on days 256-262 in 2011 and 176-182 in 2012 for E2. The amplitude of E2 was smallest at  $\sim$  2.5 K on days 246-252 in 2017, and reached the largest amplitude of  $\sim$ 11 K during 174-180 days in 2007 and 220-226 in 2020. Larger E2 events occur in the middle of the austral winter, while smaller ones happen in the 259 early and late. E2 events with temperature amplitudes greater than  $~6 K$ 260 mainly occur in  $\sim$ 38-48 hr, while those other periods have weak temperature amplitudes. The E2 period is shorter in the middle of the austral winter but becomes longer in the early and late. Most E2 events 263 exhibit temperature amplitudes exceeding  $\sim$  4 K (Figures 3a, 3d, and 3g).







 Table 3 lists the dates, periods, and amplitudes of the E3 polar 282 planetary waves event at ~60-70°S and ~50 km during austral winters from 2002 to 2022. There were 74 E3 events, with the earliest occurring on days 142-146 in 2022 and the latest on days 260-264 in 2009 and 2018. The shortest wave period for E3 was ~21 hr between days 204-208 in 2006,







 The histogram illustrates the distribution of E3 information (Figures 4b, 4e, and 4h). The E3 events are most frequent (15 times) between days 160-180, less so (10 times) between days 220-240, and least frequent (9 times) on days 240-260. A total of 43 E3 events occurred between days 302 140-200, accounting for ~58%. The wave period of E3 from ~20 to ~44 hr, 303 with the majority (60 times) occurring between  $\sim$  24 and  $\sim$  32 hr, accounting 304 for  $\sim$ 81.1% of all E3 events. The peak frequency of E3 events is at  $\sim$ 24-28 hr (31 times), followed by ~28-32 hr (29 times). The long-period E3 event 306 occurs four times in  $\sim$ 36-44 hr. The wave period of E3 remains relatively stable at ~29 hr from days 140 to 160. The E3 wave period decreases to







#### **3.4 E4**

 The dates, periods, and amplitudes of E4 polar planetary wave events 317 at  $\sim$ 50–60°S and  $\sim$ 50 km during the austral winter period from 2002 to 2022 are listed in Table 4. The table contains a comprehensive list of sixty- nine E4 events, with the earliest occurring on days 142-146 in 2003, 2008, 2009, 2015, and 2021 and the latest on days 254-260 in 2006. The E4 event 321 had a shortest wave period of  $\sim$ 17 hr on days 208-212 in 2021, and the 322 longest wave period was also  $\sim$ 27 hr, occurring on days 224-228 in 2002, 323 142-144 in 2015, and 186-190 in 2018. The smallest E4 amplitude was  $\sim$ 1- 1.5 K on days 154-158 and 192-196 in 2002, while the largest amplitude reached ~6-6.5 K on days 230-234 in 2008 and 164-168 in 2017. Larger E4 events tend to occur earlier in the austral winter, while smaller ones 327 appear later. E4 events with temperature amplitudes above  $\sim$ 3 K mainly 328 occur between  $\sim$ 20 and  $\sim$ 26 hr. The E4 events show small temperature 329 variations and have low frequencies for periods longer than  $\sim$ 27 hr and





330 shorter than  $\sim$  20 hr. The wave period of E4 is shorter in the middle of the 331 austral winter but longer in the early and late. After  $\sim$ 200 days, the wave

332 period of the E4 event increased (Figures 5a, 5d, and 5g).

333 The E4 histogram shows a high frequency of occurrence (17 times) at 334 160-180 days, a low frequency of occurrence (8 times) at 220-240 days, 335 and another low frequency (4 times) at 240-260 days. There were 48 E4 336 events recorded between days 140-200, accounting for ~69.6% of all E4 337 events. The wave period was from  $~16-28$  hr, with the higher frequency 338 between  $\sim$ 20-26 hr (54 times), accounting for  $\sim$ 78.3% of all E4 events. 339 Long-period events in E4 are observed between  $\sim$ 26-28 hr (7 times), while 340 short-period ones occur between  $\sim$ 16-18 hr (1 time). The wave period of 341 E4 decreased from days 160 to 180 ( $\sim$ 22 to  $\sim$ 20 hr), and then remained 342 relatively stable at  $\sim$ 21 hr from day  $\sim$ 200 (Figures 5b, 5e, and 5h). The 343 mean amplitude of E4 peaks at  $\sim$ 4.5 K between 160-180 days, drops to a 344 minimum  $(\sim 3 K)$  at 140-160 and 240-260 days, and reaches a maximum of 345  $\sim$  6 K at  $\sim$ 160 days ( $\sim$ 26 hr) and  $\sim$ 4 K at  $\sim$ 200 days ( $\sim$ 22 hr). The mean 346 amplitude of the wave period for E4 from  $\sim$  16-28 hr is greater than or equal 347 to  $\sim$ 3 K (Figures 5c, 5f, and 5i).

# 348 **3.5 Comparison of E1, E2, E3, and E4**

349 To further investigate the amplitude and period variabilities of polar 350 planetary waves (E1, E2, E3, and E4) during the austral winter from 2002 351 to 2022, we statistically analyze their maximum amplitude, corresponding





wave period, and occurrence date.

 The occurrence date of planetary waves is considered to be an important indicator of atmospheric waves (Lu et al., 2019; Liu et al., 2019; Forbes and Zhang, 2015; Liu et al., 2013). Figures 6A1, 6B1, 6C1, and 6D1 illustrate the distribution of event dates for each wave. E1 events are mainly observed during days 200–220 (8 times), with a slightly lower frequency during days 140–160 (1 time), and only 5 times during days 220–240. The E2 event is most frequent between days 160 and 180 (6 times), followed by occurrences on days 220–240 (5 times) and days 180– 220 (4 times). There is only one E2 event between days 140 and 160. The majority of E3 events occur between days 140 and 200 (20 times), with only one occurring between days 220 and 240. The largest frequency of E4 events is recorded during days 160-180 (11 times), followed by slightly lower frequencies during both periods 180-200 and 200-220 (each occurring three times). Only 4 times were recorded during days 140-160 and 220-240. These results contrast with those of Tang (2021), which used 2019 MERRA2 data to analyze changes in the behavior of polar planetary waves.

 The mean amplitudes of dates for events E1, E2, E3, and E4 are illustrated in Figures 6A2, 6B2, 6C2, and 6D2. The diagrams show that the mean peak temperature amplitude of E1 exceeds ~10 K between 140-240 373 days, reaching  $\sim$  12 K from 180-200 days. For E2, the mean peak amplitude











396 ranges:  $\sim$ 24-28 hr and  $\sim$ 28-32 hr, each occurring ten times. Among the 397 eighteen strongest E4 events,  $\sim$ 42.9% have a period of  $\sim$ 22-24 hr, while 398 only one event has a period of  $\sim$ 18-20 hr. These results are consistent with 399 previous studies on planetary waves conducted in the polar atmosphere 400 (Tang et al., 2021; Lu et al., 2013).

401 In Figures 6A4, 6B4, 6C4, and 6D4, the mean peak amplitude of E1 402 exceeds ~8 K within ~48-168 hr, reaching ~12 K in ~96-120 hr, with the 403 smallest peak amplitude occurring between  $\sim$ 48-72 hr ( $\sim$ 8 K). The 404 strongest peak amplitude for E2 is  $\sim$ 48 to  $\sim$ 53 hr, reaching  $\sim$ 9 K, 405 significantly stronger than neighboring intervals. The smallest peak 406 amplitude occurs between  $\sim$ 28 and  $\sim$ 33 hr, reaching  $\sim$ 6 K. Meanwhile, the 407 E2 amplitude reaches or exceeds  $\sim$  6 K from  $\sim$  33 to  $\sim$  58 hr. The mean peak 408 amplitude for E3 remains constant at  $\sim$  6 K in both time intervals of  $\sim$ 24-28 409 hr and  $\sim$ 28-32 hr. The peak amplitude for E4 exceeds  $\sim$ 3 K from  $\sim$ 18-28 hr, 410 peaking at  $\sim$  5 K between  $\sim$  22-24 hr. Our results also consistent with earlier 411 reports that the wave periods of planetary waves E1, E2, E3, and E4 during 412 the austral winter are primarily concentrated in  $\sim$  4 days,  $\sim$  2 days,  $\sim$  1.3 days, 413 and ~1 day (Lu et al., 2013; Merzlyakov and Pancheva, 2007).

414 The date-period distributions of E1, E2, E3, and E4 are shown in 415 Figures 6A5, 6B5, 6C5, and 6D5. For E1, there is a gradual shift from 416 predominantly occurring  $\sim$ 96 to  $\sim$ 72 hr before the date-period of 160 to 417 200 days, followed by a return to  $\sim$ 96 hr after the period exceeds 200 days.







# **4. Discussion**

 Based on the above statistical results, we find significant variations in temperature amplitude, temporal, and interannual variability of polar





 planetary waves in the polar stratosphere. We will discuss the dynamics of polar planetary waves through diagnostic analysis, studying the effects of mean flow instabilities, background zonal wind, critical layer, EP flux, and positive refractive index to reveal their propagation and amplification mechanism. Simultaneously, we investigate the interannual relationship between temperature amplitude of polar planetary waves and solar activity, and study the variation in phase velocity. We analyze the amplitude variations of the five events with the largest and smallest planetary waves from 2002 to 2022. Based on Figure 6A1, 6B1, 6C1, and 6D1, we examine the temporal variations of planetary waves and select events at different moments (early, middle, and late) for average analysis.

# **4.1 Amplitude variation of E1, E2, E3 and E4**

#### **4.1.1 Large and Small of E1**

 The background zonal wind and the instability of the polar atmosphere are key factors for the propagation and amplification of planetary waves in the polar atmosphere (Tang et al., 2021; Rhodes et al., 2021; Lu et al., 2019). Tang et al. (2021) investigated the 2019 MEARR2 data for the behavior of polar planetary waves and showed that the temperature amplitude of polar planetary waves is correlated with background zonal winds and instability. Figure 7 illustrates the mean spatial structure of E1 largest and smallest amplitudes (5 events), including polar planetary wave temperature, background zonal wind, positive





462 refractive index, mean flow instabilities, critical layer, and EP flux. E1 463 shows a large temperature amplitude of  $\sim$ 12 K at  $\sim$ 70-80°S and  $\sim$ 50 km, 464 with an additional peak of  $\sim 6.5$  K at  $\sim 65$  km. Small temperature amplitudes 465 of  $\sim$ 7.5 K and  $\sim$ 6.5 K were also observed at  $\sim$ 60 km and  $\sim$ 65 km, 466 respectively. The mean zonal wind for large E1 speeds is  $\sim 90$  m/s, while 467 for small E1 speeds is  $\sim 80$  m/s, both occurring between 40-60°S. The zonal 468 winds of small E1 at  $\sim$  40-60 km and  $\sim$  70-80°S were faster than those of 469 large E1. The results indicate that the peak speed of background zonal wind 470 in the mid and high latitudes of the Southern Hemisphere is close to the 471 lower polar atmosphere, which hinders E1 propagation and amplification. 472 Conversely, the weaker lower polar atmosphere zonal wind may favor the 473 propagation and amplification of E1.

 Large E1 positive refractive region indices have smoother characteristics than those with small E1 amplitudes, enhancing conditions for planetary wave propagation and amplification. EP fluxes are more likely to propagate during Southern Hemisphere winter and undergo significant amplification due to mean flow instabilities and appropriate 479 background winds in the polar and between  $~60$  km and  $~80$  km. The mean 480 flow instabilities and background winds at  $\sim$ 50-60°S and  $\sim$ 50-70 km contribute energy to propagate and amplify the EP flux into the lower atmosphere. The large E1 shows slightly weaker mean flow instabilities 483 and background wind at  $\sim$ 40-60 $\degree$ S and  $\sim$ 60-80 km compared to the small





484 E1. The wave-mean flow interaction near the critical layer (large:  $\sim$ 95 hr 485 and small:  $\sim$ 112 hr) along the green curve enhances E1. The critical layer absorbs or reflects planetary waves from the lower atmosphere as they propagate downward (Tang and Gu, 2023; Rhodes et al., 2021). The positive index region enclosed by the critical layer further facilitates E1 propagation. The positive refractive index region, stronger instability, and 490 weaker background wind at  $\sim 50-60^{\circ}$ S and  $\sim 60-70$  km provide sufficient energy for E1 amplification and propagation (Tang et al., 2021). In addition, 492 the strong instability and weak background wind at  $\sim$ 70–80°S and  $\sim$ 40–60 km provide sufficient energy for upward propagation and amplification of the EP flux (Figure 7d). The weak background winds and strong mean flow instabilities in the low polar atmosphere promote the upward propagation and amplification of EP flux. Similarly, favorable background winds and 497 mean flow instabilities at  $~60-80$  km in mid-latitudes also facilitate the propagation and amplification of EP fluxes into the lower atmosphere. Both instability and suitable background wind conditions are crucial for determining the propagation and amplification of E1. Additionally, EP fluxes can propagate and amplify to the upper polar, potentially explaining the significant E1 amplitude. Compared to the study of Tang et al. (2021) is that these events are representative based on multi-year average results, where there is more relative influence from background condition such as background zonal winds, instability, critical layers, and positive refractive





#### index.

#### **4.1.2 Large and Small of E2**

 Figure 8 illustrates the results of large and small E2. Large E2 shows 509 a large temperature amplitude of  $\sim 8.5$  K at  $\sim 60\text{-}70^{\circ}\text{S}$  and  $\sim 50$  km, with a 510 secondary peak appearing at  $~60$  km with a value of  $~7.5$  K. Small E2 511 shows temperature amplitude of  $\sim$  5 K and  $\sim$  4 K at  $\sim$  60-70°S,  $\sim$  50 km, and  $512 \sim 60$  km, respectively. The mean zonal wind reaches a maximum of  $\sim 90$ 513 m/s at  $\sim$ 50-60°S and  $\sim$ 40-60 km for large E2, and measures  $\sim$ 90 m/s for small E2. The mean zonal wind speed is greater in large E2 compared to 515 small E2 at  $\sim 60-80^\circ$ S and  $\sim 40-60$  km. The weaker background zonal wind confinement in the upper polar atmosphere hinders E2 propagation and amplification. Strong background zonal winds may promote the propagation and amplification of E2. These general results for E2 planetary wave are consistent with past study examining the 2011 Antarctic winter McMurdo (77.8° S, 166.7° E), Antarctica (Lu et al., 2013), where background zonal winds dominate the propagation and amplification of the planetary waves.

 The large E2 positive index region in the Southern Hemisphere has smoother features compared to the small E2, which facilitates the propagation and amplification of E2 in the polar atmosphere. It is shows that E2 is more likely to propagate during the Southern Hemisphere winter, 527 and mean flow instabilities in the mid-high latitudes between  $\sim$ 40 and  $\sim$ 80





 km significantly amplifies E2. As the EP flux moves towards the lower 529 atmosphere, it eventually propagates towards the equator at  $\sim$  50 km. E2 is amplified and propagated by wave-mean flow interaction near the critical layer (~48 hr) of the green curve, and promoted by a positive refractive index region. In addition, atmospheric instability and favorable background winds contribute to the energy for the propagation and amplification of the EP flux into the lower atmosphere. Based on the diagnostic analysis of small E2, it is found that E2 gains sufficient energy 536 from strong instability and background wind at  $\sim$ 50–60°S and  $\sim$ 60–70 km. This energy is then amplified and propagated to the lower atmosphere through the critical layer (~42 hr) and positive refractive index (Figure 8h). The atmosphere refractive index affects wave propagation characteristics widely (Yue et al., 2012; Chang et al., 2011; Liu et al., 2004). The 541 background wind at  $\sim$ 40-60°S and  $\sim$ 60-80 km is weaker for large E2 than 542 for small E2, while the instability of large E2 and small E2 at  $\sim$  50-60°S and ~60-70 km is similar. Our results show that u the planetary waves absorbed sufficient energy for amplification, under the background 545 conditions of large E2. TA similar region of instability  $({\sim}60-70^{\circ}S)$  and ~40 km) is also observed in (Watanabe et al., 2009; Manney et al., 1998; Lawrence et al., 1995; Allen et al., 1997), while no associated with significant EP flux.

#### **4.1.3 Large and Small of E3**







 The positive index region of large E3 exhibits smoother characteristics in the propagation area compared to that the small E3, which facilitates the propagation and amplification of E3. These regions are only in the Southern Hemisphere, indicating that E3 is more likely to propagate and amplify in the Southern Hemisphere (Tang et al., 2021; Lu et al., 2013; Fraser et al., 1993). The mean flow instabilities of E3 at ~50-70 km in mid- high latitudes, along with suitable background winds, significantly enhances the propagation of planetary waves. Additionally, the interaction





 $572$  near the critical layer at  $\sim$ 30 hr and within the positive index region further 573 enhances the E3 propagation. Notably, the weak atmospheric instability 574 and strong background winds of large E3 at  $\sim$ 40-60°S and  $\sim$ 60-80 km, 575 provide sufficient energy for the propagation and amplification of the EP 576 flux into the lower atmosphere, eventually pointing toward the equator at  $577 \sim 50$  km (Figure 9d). The EP flux of small E3 is directed towards the lower 578 atmosphere and amplified through interaction at the critical layer  $(\sim]26$  h). 579 In addition, strong instability and weak background winds at  $\sim$ 40–60°S and 580 ~60–80 km provides the energy to enhance E3 propagation (Figure 9h). 581 The EP flux of E3 eventually points to the equator at  $\sim$ 50 km. The stronger 582 the background wind at  $\sim 50-60^{\circ}$ S and  $\sim 60-70$  km, the greater the 583 temperature amplitude of E3. We propose that the background winds and 584 instability at the  $\sim$ 50–60°S and  $\sim$ 60–70 km are the primary factors driving 585 the propagation and amplification of EP flux into the lower atmosphere. 586 (Tang et al., 2021; Manney and Randel, 1993) also supports the conclusion 587 that equatorward momentum flux events appear to be consistent with the 588 instability of the double-jet structure.

#### 589 **4.1.4 Large and Small of E4**

590 Figure 10 shows the mean spatial structure of large and small E4. The 591 mean temperature amplitudes of large E4 are  $\sim$  5 K and  $\sim$  3.5 K at  $\sim$  50-60°S, 592 as well as  $\sim$  50 km and  $\sim$  70 km. Meanwhile, Small E4 exhibits temperature 593 amplitudes of  $\sim$ 2 K ( $\sim$ 50 km) and  $\sim$ 1.5 K ( $\sim$ 65 km) at  $\sim$ 50-60°S.





 Temperature amplitude measurements during planetary wave activity in the southern polar atmosphere indicate that amplitude decreases with increasing zonal wavenumber (Tang et al., 2021; Lu et al., 2013; Alexander and Shepherd, 2010; Watanabe et al., 2009; Merzlyakov and Pancheva, 2007; Fraser et al., 1993). The mean zonal wind speed of the large E4 599 reaches ~115 m/s at ~40-50°S, and ~60 km, while the small E4 has a mean 600 zonal wind speed of  $\sim$ 95 m/s at the same latitude and altitude. Additionally, 601 the zonal wind speed of large E4 exceeds that of small E4 at  $\sim$ 40-50°S and ~60 -80 km. A strong background zonal wind promotes the propagation and amplification of E4. The background zonal wind of large E4 is stronger than that of small E4 in the mid-high latitudes. This indicates that strong background zonal winds influence the propagation and amplification of E4 in the polar atmosphere. This is consistent with earlier studies in the polar atmosphere (Merzlyakov and Pancheva, 2007; Venne and Stanford, 1979), indicating that background zonal winds modulate the propagation and amplification of planetary waves. Previous studies of the polar planetary wave E4 have shown that the temperature amplitude and wave period are more similar to tidal waves s compared to others zonal wavenumber (Tang et al., 2021; Lu et al., 2013).

 Compared to the small E4, the positive refractive index region of the large E4 exhibits smoother features in the middle atmosphere at mid-latitudes, which facilitates the E4 propagation and amplification. The EP





 flux of E4 is similar to that of E3. For E4, in the middle and high latitudes, 617 between  $\sim 60$  and  $\sim 80$  km, mean flow instabilities significantly enhances E4. The EP flux of E4 propagating into the lower atmosphere and 619 eventually reaches the equator at  $\sim$  50 km. The interaction of waves near 620 the critical layer ( $\sim$ 23 hr and  $\sim$ 23 hr) amplifies and propagates E4, with the positive refractive index region providing support. The strong instability 622 and strong background winds of the large E4 at  $\sim$ 40-60°S and  $\sim$ 60-80 km provide sufficient energy for the propagation and amplification of the EP flux into the lower atmosphere. In addition, small E4 gains energy under 625 weak instability and weak background wind at  $\sim$ 40-60°S and  $\sim$ 60-80 km, then amplifies and propagates into the lower atmosphere. The background 627 wind at  $\sim$  40-60°S and  $\sim$  60-80 km is slightly stronger for large E4 than for 628 small E4, and the instability at  $\sim$ 40-60°S and  $\sim$ 60-80 km is also stronger for large E4 compared to small E4. Large E4 absorbs sufficient energy for amplification under background conditions, as shown in Figure 10a and 10e, resulting in a stronger temperature amplitude, which is consistent with the results of (Lu et al., 2013; Merzlyakov and Pancheva, 2007).

#### **4.2 Temporal variation of E1, E2, E3 and E4**

#### **4.2.1 Temporal variation of E1**

 Based on Figure 6A1, we selected E1 events occurring at 140-160 days (early E1), 200-220 days (mid E1), and 220-240 days (late E1) for average analysis (Figure 11). Consistent with previous studies, polar





 planetary waves are limited to high frequencies in certain temporal and low frequencies in others temporal (Tang et al., 2021; Lu et al., 2013; Merzlyakov and Pancheva, 2007). The mean zonal wind reaches a 641 maximum of ~100 m/s at ~30-40°S and ~50-70 km (early E1), while for 642 smaller amplitudes it measures ~90 m/s (mid E1) at ~50-60°S and ~30-40 643 km, and a maximum of ~95 m/s at ~30-40°S and ~50-70 km (late E1). The early E1 mean zonal winds exhibited strong velocities in the upper atmosphere at mid-latitudes. The strongest zonal winds gradually shifted to the lower polar atmosphere, as temporal into the middle and late. Strong background zonal wind in the middle latitudes inhibits E1 propagation and amplification, while weak wind favors it. The positive index region of middle E1 displays smoother features in the Southern Hemisphere compared to the early and late E1, promoting the propagation and amplification of E1 in polar atmosphere.

 The results indicate that E1 propagates more easily during the winter in the Southern Hemisphere, and that the mean flow instabilities at mid-654 high latitudes of  $~60$  to  $~70$  km significantly enhances E1. As the two EP fluxes move towards the lower and upper atmosphere. The early, middle, and late E1 waves are amplified and propagated near the mean critical layer 657 of the green curve ( $\sim$ 99 hr,  $\sim$ 93 hr, and  $\sim$ 95 hr) through wave-mean flow 658 interaction. The mean critical layer of E1 is close to  $\sim$  4 days ( $\sim$ 96 hr). They are also enhanced in the positive index region. In addition, the mean flow





 instabilities and favorable background winds in early, middle, and late E1 provide energy for the propagation and amplification of the EP flux into the lower atmosphere. The middle E1 shows weaker instability compared to early and late E1. The weak background zonal wind of early E1 in the lower polar atmosphere facilitated the propagation of EP flux to the upper atmosphere, while in the middle and late, the strong background zonal wind in the polar hindered this propagation to the upper polar atmosphere. The results indicate that the weak background wind and strong instability 668 at ~40-60°S and ~60-80 km favor the propagation and amplification of E1 to the polar lower atmosphere. The weak background zonal wind and strong instability of the lower atmosphere in the polar provide energy for E1 propagation to the upper atmosphere. The mid E1 shows greater instability in the polar lower atmosphere compared to early and late E1. Our findings suggest that the propagation and amplification of E1 are primarily influenced by the background zonal wind and instability in the polar atmosphere (Pancheva et al., 2008; Manney et al., 1998; Lawrence et al., 1995).

#### **4.2.2 Temporal variation of E2**

 Based on Figure 6B1, we selected E2 events occurring at 140-160 days (early E2), 160-180 days (mid E2), and 240-260 days (late E2) for 680 average analysis (Figure 12). The mean zonal wind reaches  $\sim$ 125 m/s (early 681 E2) and  $\sim$ 115 m/s (mid E2) in the  $\sim$ 30-40°S and  $\sim$ 60 km, and a maximum





682 of ~90 m/s at ~50-60°S and ~40 km (late E2). The mean zonal wind speed 683 of early E2 is greater than that of mid and late E2, at  $\sim$ 30-40°S and  $\sim$ 60 km. The results indicate that the strong background zonal wind in the upper atmosphere at mid-latitudes supports the E2 propagation and amplification. A weak background zonal wind does not create favorable conditions for E2 propagation and amplification. Early and mid E2 show stronger background zonal winds compared to late E2, suggesting that stronger zonal winds tend to influence E2 propagation in polar atmosphere. The positive refractive index region of the mid E2 hassmoother features at mid- latitudes, facilitating the E2 propagation and amplification (Tang et al., 2021; Lu et al., 2013).

 For E2, the mean flow instabilities in the mid-high latitudes at  $~60-80$  km significantly enhances E2. The E2 EP flux propagating into the lower 695 atmosphere and eventually reaches the equator at  $\sim$  50 km. E2 is amplified 696 and propagated through the interaction of waves near the critical layer  $(\sim42$ 697 hr,  $\sim$ 47 hr, and  $\sim$ 41 hr), with the positive refractive index region providing 698 support. The mean critical layer for mid E2 is closer to  $\sim$ 2 days ( $\sim$ 48 hr). 699 The stronger mean flow instabilities of early E2 at  $\sim$ 40-60°S and  $\sim$ 60-80 km compared to the mid and late. The instability in the upper atmosphere provides sufficient energy for E2 to propagation and amplification in the lower atmosphere. Early E2 strong instability in the mid-latitude upper atmosphere, but is hindered by a strong background zonal wind, which





 limits E2 propagation and amplification. Conversely, weak instability of E2 in the mid-latitude is offset by a favorable background zonal wind that promotes propagation and amplification. Late E2 encounters unfavorable conditions for propagation and amplification due to weak background 708 zonal winds. Mid E2 gains energy at  $~40-60^{\circ}$ S and  $~60-80$  km under suitable zonal wind conditions, then amplifies and propagates to the lower atmosphere. Our findings indicate that strong background zonal winds and polar instability primarily affect the E2 propagation and amplification, similar to the result by (Alexander and Shepherd, 2010; Tunbridge and Mitchell, 2009; Pancheva et al., 2008)

## **4.2.3 Temporal variation of E3**

 Based on Figure 6C1, we selected E3 events occurring at 140-160 days (early E3), 160-180 days (mid E3), and 220-240 days (late E3) for 717 average analysis (Figure 13). The mean zonal wind reaches  $\sim$ 105 m/s (early 718 E3) and  $\sim$ 120 m/s (mid E3) in the  $\sim$ 40-60°S and  $\sim$ 50-60 km. The mean zonal wind speed of late E3 is weaker than that of early and mid E3. The strong background zonal wind creates more favorable conditions for the propagation and amplification of E3 (Lu et al., 2013). This suggests that the propagation and amplification of E3 in polar atmosphere are influenced by strong background zonal wind. The positive index region of mid E3 shows smoother characteristics in the propagation compared to early and late E3, which promotes the propagation and amplification. The speed of







# **4.2.4 Temporal variation of E4**

 Based on Figure 6D1, we selected E4 events occurring at 140-160 days (early E4), 160-180 days (mid E4), and 220-240 days (late E4) for average analysis (Figure 14). The mean zonal wind speed reaches a 747 maximum of ~115 m/s at ~40-50°S and ~60 km (early E4), increases to





 $\sim$ 120 m/s in mid E4, and then drops to  $\sim$ 75 m/s in late E4. The mean zonal wind speed is greater in mid E4 compared to early and late at  $\sim$ 40-60 $\textdegree$ S and ~60-80 km. The strong background zonal wind facilitates the E4 propagation and amplification. A strong background zonal wind creates more favorable conditions for the propagation and amplification of E4. The positive index regions of mid E4 have smoother features in the Southern Hemisphere compared to the early and late E4 positive index regions, which facilitating the propagation and amplification of E4 in the polar atmosphere (Lu et al., 2013; Merzlyakov and Pancheva, 2007).

 The results indicate that E4 is more likely to propagate during the winter in the Southern Hemisphere. The mean flow instabilities at mid- high latitudes of ~60-80 km significantly enhances E4. As the EP flux moves towards the lower atmosphere, it eventually propagates towards the equator at  $\sim$  50 km. E4 is amplified and amplified near the critical layer 762 ( $\approx$ 23 hr,  $\approx$ 23 hr, and  $\approx$ 25 hr) of the green curve through wave-mean flow 763 interaction. The mean critical layer of E4 is close to  $\sim$ 1 day ( $\sim$ 24 hr). Atmospheric instability and favorable background winds provide energy for propagation and amplification the EP flux into the lower atmosphere. Through diagnostic analysis, it was discovered that E4 derived sufficient 767 energy from the strong instability and background wind at  $\sim$  50–60°S and  $768 \sim 60-70$  km. Our findings indicate that E4 can absorb sufficient energy to undergo amplification in the strong instability and favorable background





zonal winds, similar to that presented by (Tang et al., 2021).

## **4.2.5 Variation of critical layer**

 Figure 15 illustrates the mean structure of the critical layers E1, E2, E3, and E4 during the early and late, and during the high frequency. The characteristics of E1, E2, E3, and E4 critical layers were analyzed within a 10-day window on different occurrence dates. Figure 15a, 15b, and 15c demonstrate the instability, background zonal wind, and critical layer structure of E1 during days 145-155, 205-215, and 245-255. E1, E2, E3, and E4 polar planetary waves have distinct wave periods at the same temporal, while their phase velocities are notably similar (Tang et al., 2021; Lu et al., 2013; Alexander and Shepherd, 2010; Merzlyakov and Pancheva, 2007). Figure 15a exhibits greater instability compared to Figure 15b and 15c. The 4-day critical layer of E1 in Figure 15a and 15b aligns with the edge of instability. In Figure 15c, the critical layer is distant from the instability in the upper atmosphere at mid-latitudes. Cannot supply sufficient energy for the E1 propagation and amplification. The atmospheric background of Figure 15b offers more favorable conditions for E1 propagation and amplification. Simultaneously, the wave-mean interaction near the critical layer of the green curve (~4-day) enhances E1. The absence of critical layer (~2 days/yellow curve, ~1.3 days/cyan curve, and ~1 day/magenta curve) indicates that these wave periods are unsuitable for occurrence, propagation, and amplification.





792 For E2 (Figure 15d-15f), The mean flow instabilities at  $\sim 60^\circ$ S and  $\sim 60$  km, which potentially facilitate E2. The amplification of E2 is enhanced 794 by wave-mean interaction near the critical layer of the yellow curve  $(\sim 2$ -795 day). Conversely, no critical layers ( $\sim$ 4-day,  $\sim$ 1.3-day, and  $\sim$ 1-day) are found in proximity to the instability region, hindering the occurrence, propagation, and amplification of these wave periods. Weaker background winds and instability also impede E2. For E3 and E4, we observe similar characteristics in their occurrence, propagation, and amplification as those observed for E2 within the atmospheric background on 165-175 days. The 801 wave-mean interaction near the critical layer  $(\sim 1.3$ -days/cyan curve and  $\sim$ 1-day/magenta curve) enhances the amplification of E3 and E4 (Figure 15g-15i and Figure 15j-15l). In general, polar planetary waves undergo selective amplification by the background atmosphere during their occurrence, propagation, and amplification processes at different zonal wavenumbers.

# **4.3 Interannual variation of E1, E2, E3 and E4**

 To put the interannual variation of polar planetary waves results into perspective, temperature amplitude, solar activity F10.7, and phase velocity data from 2002 to 2022, indicate that variations in polar planetary wave behavior are linked to F10.7, followed by a steady variation in phase 812 velocity over the years. The temperature amplitudes of the E1, E2, E3, and E4 polar planetary waves during austral winter periods from 2002 to 2022







 E1 temperature amplitude decreased significantly during solar minima in 2007, 2008, and 2020, while E1 amplitudes increased significantly during the solar maximum years of 2003, 2012, and 2022 compared to adjacent years like 2004, 2013, and 2021. Additionally, we found that E2 amplitude shows a negative correlation with solar activity.





 Notably, the exceptionally strong E2 activity in 2006-2009, 2019, and 2020 coincided with solar minimum conditions. Unlike E1, the amplitude of E3 remains relatively constant throughout the solar cycle. Additionally, there are similar variation characteristics between E4 and E1. Temperature amplitude of E1 is positively associated with solar activity. The temperature amplitude of E2 in 2012 was relatively weak due to strong solar activity in 2011 and 2013, indicating a negative correlation. The correlation between the temperature amplitude of E3 and solar activity is weak. The temperature amplitude of E4 is positively correlated with solar activity, but anomalies occurred in 2008 and 2017.

 Figure 16e shows the phase velocities of the strongest events (E1-E4) during the austral winter periods from 2002 to 2022. The phase velocity of 848 E1 was  $\sim$  24 m/s in 2006 and  $\sim$  58 m/s in 2020, resulting in a difference of 34 m/s between the maximum and minimum phase velocities. The 850 minimum phase velocity of E2 was  $\sim$ 33 m/s in 2012, while the maximum 851 phase velocity reached  $\sim$  59 m/s in 2002. As for E3, the phase velocity 852 reached  $\sim$  53 m/s in both 2016 and 2020 (maximum), and only  $\sim$  36 m/s in 2018 (minimum). The maximum and minimum phase velocities of E4 in 854 2002 and 2013 were  $\sim$ 35 m/s and  $\sim$ 50 m/s, respectively. The phase velocity change of E1 in 2006 was significant, while the phase velocity changes for E2, E3, and E4 years were relatively stable. The mean phase velocities for 857 E1, E2, E3, and E4 are  $\sim$ 41 m/s,  $\sim$ 46 m/s,  $\sim$ 45 m/s, and  $\sim$ 42 m/s respectively.





 We propose that polar planetary waves propagate as fixed-phase wave packets, selectively amplifying planetary waves with different zonal wavenumbers due to variations in the background atmospheres while maintaining similar phase velocities, similar to that presented by (Tang et al., 2021; Lu et al., 2013; Alexander and Shepherd, 2010).

**5. Summary**

 The objective of this study is to analyze the occurrence date, wave period, and peak amplitude distribution of planetary waves E1, E2, E3, and E4 in the stratosphere and mesosphere using the 2002-2022 MERRA2 data on wind and temperature, explore their amplitude and temporal variation, while also investigating their interannual variation. This is motivated by a large inventory of past studies that examined single year variation and related topics (SSWs) but did not have a comprehensive study of the propagation and amplification, which is a gap addressed by this study. An important result in this work is that the propagation and amplification characteristics of polar planetary waves during austral winter periods are influenced by background zonal winds and atmospheric instability, and that there is significant variability temperature amplitude, wave period, and temporal variations examined for a given zonal wavenumber. This point is crucial as it enables us to examine the correlation between polar planetary wave propagation, amplification, and zonal wavenumbers, which have been archived for this (Tang et al., 2021; Lu et al., 2013; Alexander and





 Shepherd, 2010; Merzlyakov and Pancheva, 2007) and other planetary waves. The key findings of this study are summarized below:

 Peaks of E1, E2, E3, and E4 decrease in amplitude and latitude with increasing wavenumber, while the timing of planetary wave events 884 advances. The wave period and amplitude of E1 is mainly from  $\sim$ 3 to  $\sim$ 5 885 days and ~12 K, while E2: is ~38 to ~48 hr and ~10 K, E3: ~24 to ~32 hr 886 and ~8 K, E4:  $\sim$  22 to ~24 hr and ~6 K. The wave periods of E1, E2, and E3 progressively shorten during early to middle austral winter and lengthen during late austral winter. Stratospheric dynamics still significantly influence the propagation and amplification of planetary waves in the polar. Planetary waves are confined to high latitudes due to the negative refractive 891 index at  $\sim$ 45°S in the equatorial direction, which results in their evanescent properties and prevents them from propagating to lower latitudes. EP flux and instability analysis indicate that the baroclinic instability between 894 ~50°S-60°S caused by the stratospheric polar night jet / "double jet" structure is likely responsible for generating these waves.

 Polar planetary waves exhibit two EP fluxes: one propagates from the lower to the upper atmosphere, while the other from the upper to the lower 898 atmosphere. EP fluxes ultimately deflect towards the equator at  $\sim$  50 km. The polar planetary wave amplifies and propagates through suitable background wind and instability in middle and high latitudes. The stronger atmospheric instability and favorable polar background winds promote the





 propagation and amplification of planetary waves, consequently leading to greater temperature amplitudes. The background atmosphere selectively amplifies planetary waves during their occurrence, propagation, and amplification at various zonal wavenumbers. The eastward planetary wave propagates in the polar atmosphere as a fixed-phase wave packet. E1, E2, E3, and E4 are selectively amplified by different background atmospheres while maintaining similar phase velocities. The polar planetary wave activity is linked to solar activity, with E1 showing a positive correlation to E4, E2 displaying a negative correlation, and E3 having no strong correlation. Overall, we analyze their dynamics variation of eastward planetary waves in the polar atmosphere during the 2002-2022 austral winter periods and statistically analyze the interannual variation.

 It is important to consider some limitations in this study to improve upon in future work, such as a need for more statistical data, a need for more detailed measured data, a lack of measurements relating to the effects of polar particle deposition and strong magnetic fields, and the influence of the unique atmospheric conditions and geographical features of the polar atmosphere on planetary wave structure. This work motivates continued attention to the polar planetary wave behavior and provides a special conclusion to enrich the polar atmospheric research for several decades.

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**Table 1.** Occurrence Date, Wave Period, and Peak Amplitude of the E1 polar planetary

waves During the 2002–2022 Austral Winter Period







 **Table 2.** Occurrence Date, Wave Period, and Peak Amplitude of the E2 polar planetary waves During the 2002–2022 Austral Winter Period

No.	Year	Day	Period(hr)	Amp(K)	No.	Year	Day	Period(hr)	Amp(K)
1	2002	186	32	6.2	51	2012	238	52	4.9
2	2002	198	35	3.2	52	2013	164	37	6.8
3	2002	208	39	5.1	53	2013	194	45	5.7
4	2003	142	42	4.5	54	2013	204	35	5.5
5	2003	162	44	5.5	55	2013	228	43	9.5
6	2003	182	40	8.1	56	2014	150	51	5.7
7	2003	208	42	6.9	57	2014	178	46	4.4
8	2003	228	38	7.5	58	2014	186	34	6.8
9	2003	250	47	4.3	59	2014	222	36	9.5
10	2004	164	41	7.5	60	2015	158	41	3.8
11	2004	192	45	6.6	61	2015	174	44	8.9
12	2004	206	31	6.0	62	2015	210	43	7.4
13	2004	254	56	5.2	63	2015	218	42	9.2
14	2005	160	46	6.1	64	2015	248	48	6.7
15	2005	178	39	5.8	65	2016	152	39	7.4
16	2005	192	37	6.2	66	2016	180	36	5.9
17	2005	226	39	5.1	67	2016	204	39	7.4
18	2005	256	47	3.7	68	2016	232	36	4.1
19	2006	148	52	5.1	69	2016	244	43	3.4
20	2006	176	42	4.3	70	2017	142	47	5.3
21	2006	188	44	6.4	71	2017	162	43	8.2
22	2006	220	46	4.9	72	2017	196	40	5.5







 **Table 3.** Occurrence Date, Wave Period, and Peak Amplitude of the E3 polar planetary waves During the 2002–2022 Austral Winter Period

No.	Year	Day	Period(hr)	Amp(K)	No.	Year	Day	Period(hr)	Amp(K)
1	2002	154	27	4.5	38	2013	190	22	4.3
2	2002	186	23	4.4	39	2013	204	36	3.3
3	2003	160	28	4.4	40	2013	226	24	3.0
4	2003	174	30	7.5	41	2014	148	29	6.8
5	2003	196	25	7.0	42	2014	176	26	5.5
6	2003	212	24	5.2	43	2014	190	26	6.5
7	2004	156	22	1.9	44	2014	218	23	3.9
8	2004	184	27	7.1	45	2015	178	27	6.4
9	2004	252	35	3.4	46	2015	202	27	5.5
10	2005	160	32	4.3	47	2015	230	30	6.2







 **Table 4.** Occurrence Date, Wave Period, and Peak Amplitude of the E4 polar planetary waves During the 2002–2022 Austral Winter Period

No.	Year	◡ Day	Period(hr)	Amp(K)	No.	Year	Day	Period(hr)	Amp(K)
1	2002	154	18	1.1	36	2012	210	18	2.0
2	2002	192	21	1.3	37	2013	156	21	2.7
3	2002	224	27	1.6	38	2013	172	19	3.2
4	2003	142	23	1.8	39	2013	190	18	2.9
5	2003	170	24	4.5	40	2014	166	23	3.9
6	2003	194	21	4.1	41	2015	142	27	5.7
$\overline{7}$	2004	166	22	4.4	42	2015	164	23	5.1
8	2004	186	21	2.0	43	2015	196	21	5.7
9	2005	150	26	3.6	44	2015	210	20	4.9
10	2005	164	24	3.7	45	2015	232	23	4.2
11	2005	194	20	2.8	46	2016	148	23	2.4























 **Figure 2.** The dispersion of E1 (a) date-amplitude, (d) period-amplitude, and (g) date-period variations during the austral winter period from 2002 to 2022. Statistical findings for the occurrence dates (b, e, and h) and mean amplitude (c, f, and i) of E1.



**Figure 3.** Same as Figure 2 but for E2.







**Figure 4.** Same as Figure 3 but for E3.



**Figure 5.** Same as Figure 4 but for E4.











- 1093 **Figure 6.** Statistical results of the strongest austral winter E1 (Column 1),
- 1094 E2 (Column 2), E3 (Column 3), and E4 (Column 4) events occurring
- 1095 between 2002 and 2022 are presented.







 **Figure 7.** The mean spatial structures of the E1 large and small amplitude events during the 2002 to 2022 are characterized by variations in temperature, zonal wind, positive refractive index and diagnostic analysis. The shaded regions in the diagnostic analysis indicate areas of instability, while the red arrows represent EP flux and the green line denotes the critical layer. The brown area signifies a positive refractive index.







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1104 **Figure 8.** Same as Figure 7 but for E2.







1106 **Figure 9.** Same as Figure 8 but for E3.







1108 **Figure 10.** Same as Figure 9 but for E4.







 **Figure 11.** The mean spatial structure of E1 early, middle, and late events from 2002 to 2022 is characterized by temperature, zonal wind, positive refractive index, and diagnostic analysis (early: 140-160 days, early: 200- 220 days, early: 220-240 days).







1115 **Figure 12.** Same as Figure 11 but for E2 (early: 140-160 days, early: 160-

1116 180 days, early: 240-260 days).







1118 **Figure 13.** Same as Figure 12 but for E3 (early: 140-160 days, early: 160-

1119 180 days, early: 220-240 days).







1121 **Figure 14.** Same as Figure 13 but for E4 (early: 140-160 days, early: 160-

1122 180 days, early: 220-240 days).







 **Figure 15.** The green, yellow, cyan, and magenta lines represent critical layers corresponding to the 4-day, 2-day, 1.3-day, and 1-day wave periods for planetary waves E1, E2, E3, and E4 during days 145–155 (a, d, g, and j), days 205–215 (b), days 165–175 (e, h, and k), and days 245–255 (c, f, i, and l) from 2002 to 2022.







 **Figure 16.** The strongest amplitudes of E1 (a), E2 (b), E3 (c), and E4 (d) during 2002 to 2022 are depicted. The blue line represents the mean 10.7 cm solar flux during the austral winter, while the orange line illustrates the strongest amplitude of E1, E2, E3, and E4. Description of phase velocities (e) for the strongest events (E1, E2, E3, and E4).