

 Abstract: The South China Sea (SCS) is the largest semi-enclosed marginal sea in western Pacific. It exhibits a unique vertically rotating cyclonic, anticyclonic, and cyclonic circulation in its upper, middle, and deep layers. Over slope topography, these layered currents interact and significantly shape the structure and intensity of the basin circulation. In this study, we employ process-oriented numerical simulations to investigate how upper-layer processes, characterized by greater magnitude and variability, influence the layered circulation over the irregular topographic slope. The simulations reveal that stronger upper intrusion from open ocean directly enhances upper layer circulation, which subsequently strengthens the middle and the deep slope currents. Vorticity dynamics illustrate that changes in the middle and deep slope 26 current are largely related to the vertical stretching $(\zeta_D IV)$ induced by bottom geostrophic 27 cross-isobath transport (CGT_b) . As the upper-layer cyclonic slope current intensifies, it modulates the bottom pressure distribution, resulting in stronger negative ζ_DIV predominantly over the northwestern slope to intensify the middle anticyclone slope current. Similarly, for the 30 deep cyclonic slope current, the CGT_b maintains downward cascading in the northern part and upwelling over the southern slope. Over the southern slope, the strengthening of the positive CGT_b is induced by the increment of the advection of relative vorticity and planetary vorticity in water column, in which the middle layer provides approximately 40% of the total strengthening trend, but the upper layer has a minimal impact. Conversely, on the northern 35 slope, the strengthening of the negative CGT_b is primarily influenced by the upper layer. **Keywords**: Circulation dynamics; Layered circulation; Process-oriented simulation;

- Topographic modulation; Vertical coupling
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1. **Introduction**

 Marginal seas circulation plays crucial role in the mass transport and regional climate dynamics, with circulation patterns being strongly influenced by their unique topographic features (Millot 1999, Omstedt, Elken et al. 2004, Oey, Ezer et al. 2005, Gan, Li et al. 2006, Johns and Sofianos 2012). This topographic regulation results in intricate layered slope currents that flow in different directions at various depths (Yuan 2002, Wang, Xie et al. 2011, Lan, Zhang et al. 2013, Gan, Kung et al. 2022). These layered currents interact over slope topography, significantly shaping the structure and intensity of the overall circulation (Gan, Liu et al. 2016, Quan and Xue 2018). For instance, in the Gulf of Mexico, loop current eddies influence not just the upper-layer circulation but also the deep flow, illustrating the importance of vertical coupling in transferring energy and variability to deeper currents (Tenreiro, Candela et al. 2018). Understanding these dynamics helps elucidate the behavior of marginal sea's circulation (Liang, Spall et al. 2017, Zhu and Liang 2020, Olvera-Prado, Moreles et al. 2023).

 The South China Sea (SCS) is the largest semi-enclosed marginal sea located in the tropical region. It exhibits a unique mean cyclonic, anticyclonic, and cyclonic (CAC) circulation in its upper, middle, and deep layers (Wang, Xie et al. 2011, Shu, Xue et al. 2014, Lan, Wang et al. 2015, Gan, Liu et al. 2016, Zhu, Sun et al. 2017). The strong upper layer circulation is driven by the Asia monsoon and Kuroshio intrusion from Luzon Strait (LS), while the middle and deep circulations are maintained by the outflux and deep intrusion through LS, respectively. Over the slope topography, the CAC circulation affects each other through the vertical coupling among them (Shu, Wang et al. 2018). For example, the study of Quan and Xue (2018) demonstrated that perturbations in the upper layer can transit through the water column and affect deeper circulation by altering layer thickness. This coupling is particularly evident over the strong boundary current that in the northern basin and along the western boundary. Numerical experiments by Wang et al. (2018) suggest that, although the basin-scale circulation is primarily driven by the LS overflow, upwelling patterns are essential in shaping the detailed structure of the deep-water circulation. Similar processes have been observed in other marginal seas as well (e.g., Testor et al., 2018; Wang et al., 2024; Xu et al., 2009; Olvera-Prado, Moreles et al. 2023).

 Despite advances in understanding the major features of mean layered circulation, the response of this circulation to changes in external forcings—particularly the strong upper ocean processes—and the specific mechanisms involved remain poorly understood. A comprehensive understanding of how topographically modulated vertical coupling over basin slopes influences the structure and intensity of mean circulation is crucial for exploring the long-term evolution of marginal sea circulation. For instance, Kuroshio intrusion has been observed to weaken with

 decadal variations, while upper ocean currents have shown acceleration in response to a warming climate (Peng et al., 2022; Chen et al., 2019; Nan et al., 2013). How the layered circulation in marginal seas responds to these changes remains unclear. This study employs process-oriented simulations to elucidate how upper-layer processes, which exhibit greater intensity and variability, influence layered circulation over the meandering bottom slope in the 81 SCS. The paper is organized as follows: Section 2 explains the configuration of the simulation. Section 3 offers the response of layered slope current intensity to the changes of the upper circulation. Section 4 explores the underlying physical mechanisms. Finally, Section 5 provides a summary of our findings.

 Figure1. Location and bathymetry (km) of the South China Sea, showing Luzon Strait, Taiwan Strait, Mindoro Strait, and Karimata Strait.

2. Methodology

 Based on the Regional Ocean Modeling System (ROMS) (Shchepetkin & McWilliams, 2005), we performed a three-dimensional process-oriented simulation in this study (Cai et al., 2023). The model setup included a circular basin on the west representing the SCS and a rectangular basin on the east representing the Western Pacific Ocean, connected by a narrow strait (representing the Luzon Strait, LS) with a depth of 2500 m. The slope in the SCS increases from 400 to 4000 m and the central basin has the depth of 4000 m.

The model employed a uniform horizontal grid with a 5 km resolution, and the vertical

 structure utilized stretched, terrain-following coordinates with 30 layers (Song and Haidvogel 1994) were adopted with 30 layers. In the upper layer, the influx/outflux was defined at the southeastern (25 Sv) and northeastern (20 Sv) boundaries of the open ocean, and the SCS had an opening to the south to allow for outflux (Figure 2a), consequently, the intrusion from open 102 ocean into SCS will be generated. The model incorporated variable mixing coefficients (K_n) to simulate differences in exchange currents between the SCS and the Western Pacific Ocean, particularly in the middle and deep parts of the LS (Tian, Yang et al. 2009, Yang, Zhao et al. 105 2016). In the SCS and the western half of the LS, K_v values were set as $5 \cdot 10^{-4} m^2 s^{-1}$ between 106 500 and 1,500 m, $2 \cdot 10^{-3} m^2 s^{-1}$ below 1,500 m, and $5 \cdot 10^{-3} m^2 s^{-1}$ in the layer 500 m from the bottom. In the upper 500 m of SCS and LS, and throughout the entire water column of the open 108 ocean, a background K_v of $2 \cdot 10^{-5} m^2 s^{-1}$ was applied. A zonal wind stress with meridional variations was applied over the open ocean, while the atmospheric buoyancy flux over the SCS was simplified by setting it to zero.

- The model was initialized with horizontally uniform temperature and salinity profiles derived from the mean data of the World Ocean Atlas, averaged over the region west of the LS. The simulation ran for 25 years, and the analysis was conducted on the results from the final 5 years.
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 Figure 2. (a) The topography and geometry used in the idealized process-oriented simulation. The black lines show the realistic surrounding coastline. S1 and S2 are two vertical sections for showing slope current. (b) Distribution of the wind stress.

3. Results

 In the LS, the influx occurs in the upper and deep layers, while the outflux occurs in the middle layer between them (Figure 3b). The upper intrusion is intrinsically linked to the western

 boundary current in the open ocean, i.e. Kuroshio Current, while the exchange currents in the middle and bottom layers are primarily related to the density differences between the South China Sea (SCS) and the Pacific (Zhu et al., 2017, 2019; Cai, Chen et al. 2023; Zhou et al., 2023). Under the density difference, the westward pressure gradient was formed that drives the deep intrusion from the open ocean towards the SCS (Wang, Xie et al. 2011). Associated with the stimulated layered exchanging current, the horizontally averaged vorticity features with the positive-negative-positive values in the respective depth of the basin, which represent the cyclonic, anticyclonic and cyclonic slope current (Figure 3b).

134 Figure 3. (a) Vertical profile of domain averaged density $ρ$ (kg·m⁻³) in SCS and Pacific between
135 1000m to 3000m. (b) Black line shows the vertical profile of the SCS basin-averaged vorticity 1000m to 3000m. (b) Black line shows the vertical profile of the SCS basin-averaged vorticity 136 (s⁻¹) and the red line shows the flux through LS $(m^2 \cdot s^{-1})$. The positive/negative vorticity represents the cyclonic/anticyclonic circulation, and the positive/negative value flux represents the outflux/influx through LS.

 To investigate how the changes in upper layer current modulates the layered circulation, additional numerical experiments were conducted by adjusting the northeastern outflux to 22Sv, 23Sv and 24Sv, thus the upper LS intrusion was adjusted to 3Sv (Case_U3), 2Sv (Case_U2) 144 and 1Sv (Case U1). In these scenarios, the middle and deep exchanging currents, which were primarily sustained by the contrasting densities between the SCS and the open ocean, remained relatively unchanged (Figure 4a). Inside the basin, the enhancement of the upper layer inflow directly intensified the upper layer cyclonic circulation (Figure 4 c-f). And subsequently strengthened the middle anticyclonic slope current, particularly in the northwestern part of the basin (transect S1, Figure 4c, d). Additionally, although not directly connected to the upper layer, the deep cyclonic slope current also exhibited increased strength (Figure 4c-f). Using the domain-averaged vorticity as an indicator, it was observed that the intensification of the upper layer circulation resulted in a 10% increase in the intensity of the middle anticyclonic

- 153 circulation and a 27% increase in the deep cyclonic circulation (Figure 4b). Given that only the
- 154 upper layer circulation is directly amplified by the upper intrusion, the remote influence from
- 155 upper layer processes over the basin slope significantly impacts the intensity of the layered
- 156 circulation, which was explored below.

160 **Figure 4**. (a) Vertical profile flux across the LS under varying upper layer influx. (b) Changes 161 in the basin-averaged vorticity between 0-500 m (Upper), 1000-2000 m (Middle) and 3000- 162 4000 m (Deep) with different upper influx. (c-f) The along-slope current $(m \cdot s^{-1})$ over the 163 section S1 and S1 in Standard case (upper intrusion is 5 Sv) and Case_U1 (upper LS intrusion 164 is 1Sv), respectively. The locations of the two transects are shown in Figure 1a. The transect 165 locations are indicated in Figure 1a. Positive values represent anticyclonic flow, while negative 166 values indicate cyclonic flow.

167 168

169 **4. Discussion**

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171 To understand how the upper layer processes modulate the layered current over the slope,

- 172 we first investigated the layered-integrated vorticity budget for each layer (Gan, Liu et al. 2016,
- 173 Cai, Chen et al. 2023):

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$$
\frac{\zeta_{\text{LOR}}}{\int_A [\nabla \times \int_{layer}^{\Box} (\overline{PGF}) \, dz] dA + \int_A [\nabla \times \int_{layer}^{\Box} (\overline{ADV}) \, dz] dA + \underbrace{\int_A \left[-f \nabla \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{LOIV}}} + \underbrace{\int_A \left[-f \nabla \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{LOIV}}} + \underbrace{\int_A \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta_{\text{DIIV}}} + \underbrace{\int_{\Box} \left[-f \cdot \int_{layer}^{\Box} (\overrightarrow{V}_h) \, dz \right] dA}_{\zeta
$$

175
$$
\int_{A} [\nabla \times \int_{layer} \overline{(VIS)} dz] dA = 0 \quad (1)
$$

176 where PGF is pressure gradient force, ADV is nonlinear advection, VIS is the turbulent 177 viscosity. COR is the Coriolis force and the domain-integrated ζ COR represents the lateral

178 planetary vorticity flux. Inside the basin, the ζ COR can be further decomposed into the 179 vertical stretching induced by the divergence of the horizontal flow $(\zeta_D IV)$ and β effect of the 180 meridional current $(\zeta$ _BETA). The circulation in SCS is primarily governed by geostrophic balance, manifested as 182 ζ COR balanced by the ζ PGF in both layers. It suggests that the mean cyclonic/anticyclonic circulation is related to the lateral planetary vorticity influx/outflux (Figure 5) (Cai and Gan 2019, Cai, Chen et al. 2023; Zhu et al., 2017, 2019). Consequently, the strengthening of the upper LS intrusion directly intensifies the upper layer cyclonic slope current by providing more positive planetary vorticity flux. In the semi-enclosed middle and deep layers, the vorticity 187 input is mainly provided by the ζ_D/V , as also highlighted in the previous studies (Zhu, Sun et 188 al. 2017, Wang, Du et al. 2018, Cai, Chen et al. 2023). The positive $\zeta_D IV$ in the upper and deep layers reflectsthe downward and upward flux, which squeeze the middle layer and provide 190 the negative vertical flux (negative $\zeta_D IV$) for the anticyclonic circulation. Therefore, the intensification of the middle and deep slope currents is largely attributed to vertical stretching over the basin slope.

 Figure 5. The layered-integrated vorticity budget (Equation 1) for the standard case in the layers of (a) 0–500, (b) 1,000–2,000, and (c) 2,500–4,000 m.

a. Intensification of middle anticyclonic slope current

202 For the middle anticyclonic slope current, the negative middle-layered $\zeta_D IV (\zeta_D IV_M)$ mainly occurs over the northwestern slope (Figure 6a), where a relatively strong anticyclonic 204 slope current is observed. In this area, the upper layer generally exhibits positive ζ_{D} DIV 205 (ζ _DIV_U), indicating downward squeezing from the upper layer that helps maintain the middle anticyclonic slope current (Figure 6 b). The increased squeezing from the upper slope current is expected to intensify the middle anticyclonic slope current (Figure 4 c, d). Since the stretching/squeezing imposed on the slope current is closely related to the vertical motion (Liang, Spall et al. 2017), the mean vertical motions over the slope can be understood as follows:

$$
210 \qquad \overline{w} = -\nabla_h \cdot \int_{-H}^{Z} \overline{V}_h dz \approx -\nabla_h \cdot \int_{-H}^{Z} \overline{V}_h \, \text{gen} \, dz = \overline{-(\tilde{v}_{b\text{gen}} H_y + \tilde{u}_{b\text{gen}} H_x)} + \frac{\beta}{f} \int_{-H}^{Z} \tilde{v}_{geo} \, dz \tag{2}
$$

212 current, which is maintained by the distribution of bottom pressure. H_x and H_y are the 211 where H is the topography, \emph{geo} represents the geostrophic component of the horizontal 213 horizontal gradients of the bottom slope, \overline{u}_{b_geo} and \overline{v}_{b_geo} are the bottom geostrophic currents. 214 When the current flows over the slope, the bottom pressure distribution interacts with the 215 meandering topography and the resulting geostrophic cross-isobath transport (CGT_b) to 216 influence the vertical motions over the slope. Generally, the CGT_h largely governs the mean 217 vertical motion over northwestern slope, which in turn squeezes and stretches the middle and 218 upper slope currents. This process is crucial in maintaining the vorticity input through ζ_DIV 219 (Figure 6 d), thereby sustaining the dynamics of the slope currents.

220 Associated with the intensification of the upper layer cyclonic slope current, the stronger 221 slope current modulates the bottom pressure distribution over the slope, leading to stronger 222 vertical squeezing $(\zeta_D IV)$ within the water column. The increase in squeezing gradually 223 deepens the isopycnal surface and strengthens the middle anticyclonic slope current (Figure 6e). 224

226 **Figure 6**. (a-c) ζ DIV integrated between 1000-2000 m (ζ DIV_M), 0-500 m (ζ DIV_U), and 227 CGT_b over the region of middle anticyclonic slope current between 1000-2000 m. The black 228 line indicates the 2000 m isobath line. (d) Change of the ζ_DIV integrated between 100-2000 m 229 ($\zeta_D IV_M$), between 0-500 m ($\zeta_D IV_U$) with the depth over the northwestern slope. (e) Changes
230 of the $\zeta_D IV_M$, $\zeta_D IV_U$, CGT_b and the depth of isopycnal surface of 1027.4 kg/m³ with different of the $\zeta_D IV_M$, $\zeta_D IV_U$, CGT_b and the depth of isopycnal surface of 1027.4 kg/m³ with different

231 upper LS influx. (d) and (e) was plotted using the data averaged over the region of black box in (a) $in (a)$

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235 **b. Intensification of the deep cyclonic slope current**

236 For the deep cyclonic slope current, the deep $\zeta_D IV(\zeta_D IV_p)$ features a downward flux (negative value) over the northern slope and upward flux (positive value) over the southern part (Figure 7a, d). This pattern aligns with the classical source-sink driven theory (Stommel and Arons 1959; Yuan et al., 1997), which suggests the sinking from the northern basin and upwelling from the southern basin would sustain the abyssal cyclonic circulation. Same as the 241 middle anticyclonic slope current, the $\zeta_D V_D$ is largely controlled by the CGT_b that the pressure distribution maintained the mean cascading in the northern side and the upwelling over the southern slope (Figure 7b, d). To explain the maintenance of the bottom pressure distribution, the vertical integrated vorticity dynamic was employed (Gan, Liang et al. 2013, Cai and Gan 2021):

$$
246 \qquad \overline{V} \times \int_{-H}^{0} \overline{(\overline{P}\overline{G}\overline{F})} \, dz + \overline{V} \times \int_{-H}^{0} \overline{(AD\overline{V})} \, dz + \underbrace{\beta \int_{-H}^{0} v dz}_{\Omega _ \overline{B} \overline{E} \overline{T}A} - \overline{V} \times \int_{-H}^{0} \overline{(V}\overline{I}\overline{S}) \, dz = 0 \tag{3}
$$

Where $\Omega_{\perp} PGF = \nabla \times \int_{-\pi}^{0} (\overrightarrow{PGF})$ 247 Where $\Omega_{\perp} PGF = \nabla \times \int_{-H}^{0} (PGF) dz = -(\bar{v}_{b_{}geo} H_y + \bar{u}_{b_{}geo} H_x)$, thus represents the effect of CGT_b . 248 It illustrates that the bottom pressure distribution is maintained by the nonlinear advection of 249 relative vorticity $(\Omega _ADV)$, advection of planetary vorticity $(\Omega _BETA)$, and net viscosity of the 250 current flowing over the slope topography. Among them, the Ω _{-BETA} and Ω _{-ADV} play the 251 major role in sustaining the bottom pressure distribution of the deep cyclonic slope current 252 (Figure 8 c, d).

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255

256 Figure 7 (a-c) ζ_DIV integrated below 3000 m of the deep cyclonic slope current (ζ_DIV_D),
257 *CGT_h* and Ω *BETA* + Ω *ADV* over the region of middle anticyclonic slope current below 3000 257 *CGT_b* and $Ω$ ₋*BETA* + $Ω$ ₋*ADV* over the region of middle anticyclonic slope current below 3000
258 m. The black line indicates the 4000 m isobath line. (d) Meridional changes of the $ζ$ -DIV_{*n*}, 258 m. The black line indicates the 4000 m isobath line. (d) Meridional changes of the $\zeta_D IV_D$,
259 CGT_h and $\Omega_BETA + \Omega_ADV$ over the slope. CGT_b and Ω _BETA + Ω _ADV over the slope. 260

261 As the upper intrusion intensifies, the bottom pressure distribution adjusts in response to 262 changes in the layered circulation, resulting in a gradual strengthening of the negative 263 CGT_b over the northern slope and positive CGT_b over the southern part (Figure 8 a-b). Over the 264 southern slope (Figure 8a), the strengthening of the positive CGT_b is induced by the increase of 265 the $Ω$ ₋BETA + $Ω$ ₋ADV in the water column, in which the middle layer $[Ω$ _{-BETA} + $Ω$ ₋ADV (M)] 266 provides approximately 40% of the total strengthening trend, while the upper layer $\left[\Omega_{\text{BETA}}\right]$ $267 + \Omega_A DV$ (U)] has a negligible impact. Conversely, over the northern slope, the strengthening 268 of the Ω _BETA + Ω _ADV is primarily influenced by the upper layer, with the middle layer has 269 the negative effect (Figure 8b). Thus, the intensification of downward CGT_b over the northern 270 slope is directly modulated by the strengthening of the upper cyclonic slope current, while the 271 intensification of upward CGT_b over the southern slope is mainly driven by the strengthening 272 of the middle anticyclonic slope current. These changes in CGT_b enhance deep stretching, 273 thereby strengthening the deep cyclonic slope current accordingly (Figure 8c). 274

Figure 8. (a) Changes of the CGT_b , $\Omega_{Beta} + \Omega_{Adv}$ in the total water column, $\Omega_{Beta} + \Omega_{Adv}$ in the layer between 0-500 m (Ω_{Beta} + Ω_{Adv} II), $\Omega_{Beta} + \Omega_{Adv}$ in the between 1000-2000 m layer between 0-500 m (Ω_{Beta_U} + Ω_{Adv_U}), Ω_{Beta} + Ω_{Adv} in the between 1000-2000 m 280 ($\Omega_{Beta_M} + \Omega_{Adv_M}$) over the southern slope (black box in Figure 8a). (b) same as (a) but over 281 the northern slope (black box in Figure 8a). (c) Change of the CGT_b and $\zeta_D IV_b$ over the 282 northern and southern slope. northern and southern slope.

5. Summary

 Marginal sea circulation plays a crucial role in mass transport and regional climate dynamics. Through process-oriented numerical simulations, we examined how upper-layer processes, which are characterized by greater intensity and variability, impact the layered circulation over the meandering bottom slope in SCS. The results underscore the intricate balance between topographical features and oceanic circulation, offering valuable insights for predicting the behavior of marginal sea circulation under varying forcings conditions

 The numerical experiments show that the stimulated layered exchange currents through the LS introduce lateral planetary vorticity flux and facilitate the development of a layered basin circulation within the SCS. Inside the basin, the intensification of upper-layer inflow directly enhances the upper-layer cyclonic circulation, which in turn strengthens the middle anticyclonic slope current, particularly in the northwestern part of the basin. Furthermore, even though the deep cyclonic slope current is not directly connected to the upper layer, it also exhibits increased

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