



1	Topographic modulation on the layered circulation in South China Sea
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17 Abstract: The South China Sea (SCS) is the largest semi-enclosed marginal sea in western 18 Pacific. It exhibits a unique vertically rotating cyclonic, anticyclonic, and cyclonic circulation 19 in its upper, middle, and deep layers. Over slope topography, these layered currents interact and 20 significantly shape the structure and intensity of the basin circulation. In this study, we employ process-oriented numerical simulations to investigate how upper-layer processes, characterized 21 22 by greater magnitude and variability, influence the layered circulation over the irregular topographic slope. The simulations reveal that stronger upper intrusion from open ocean 23 24 directly enhances upper layer circulation, which subsequently strengthens the middle and the 25 deep slope currents. Vorticity dynamics illustrate that changes in the middle and deep slope 26 current are largely related to the vertical stretching (ζ _DIV) induced by bottom geostrophic 27 cross-isobath transport (CGT_b). As the upper-layer cyclonic slope current intensifies, it 28 modulates the bottom pressure distribution, resulting in stronger negative ζ DIV predominantly 29 over the northwestern slope to intensify the middle anticyclone slope current. Similarly, for the deep cyclonic slope current, the CGT_b maintains downward cascading in the northern part and 30 upwelling over the southern slope. Over the southern slope, the strengthening of the positive 31 CGT_b is induced by the increment of the advection of relative vorticity and planetary vorticity 32 33 in water column, in which the middle layer provides approximately 40% of the total 34 strengthening trend, but the upper layer has a minimal impact. Conversely, on the northern 35 slope, the strengthening of the negative CGT_b is primarily influenced by the upper layer. 36 Keywords: Circulation dynamics; Layered circulation; Process-oriented simulation; 37

- 38 Topographic modulation; Vertical coupling
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40 1. Introduction

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42 Marginal seas circulation plays crucial role in the mass transport and regional climate dynamics, with circulation patterns being strongly influenced by their unique topographic 43 44 features (Millot 1999, Omstedt, Elken et al. 2004, Oey, Ezer et al. 2005, Gan, Li et al. 2006, 45 Johns and Sofianos 2012). This topographic regulation results in intricate layered slope currents that flow in different directions at various depths (Yuan 2002, Wang, Xie et al. 2011, Lan, Zhang 46 47 et al. 2013, Gan, Kung et al. 2022). These layered currents interact over slope topography, 48 significantly shaping the structure and intensity of the overall circulation (Gan, Liu et al. 2016, 49 Quan and Xue 2018). For instance, in the Gulf of Mexico, loop current eddies influence not 50 just the upper-layer circulation but also the deep flow, illustrating the importance of vertical 51 coupling in transferring energy and variability to deeper currents (Tenreiro, Candela et al. 2018). 52 Understanding these dynamics helps elucidate the behavior of marginal sea's circulation (Liang, Spall et al. 2017, Zhu and Liang 2020, Olvera-Prado, Moreles et al. 2023). 53

The South China Sea (SCS) is the largest semi-enclosed marginal sea located in the 54 55 tropical region. It exhibits a unique mean cyclonic, anticyclonic, and cyclonic (CAC) circulation in its upper, middle, and deep layers (Wang, Xie et al. 2011, Shu, Xue et al. 2014, 56 57 Lan, Wang et al. 2015, Gan, Liu et al. 2016, Zhu, Sun et al. 2017). The strong upper layer 58 circulation is driven by the Asia monsoon and Kuroshio intrusion from Luzon Strait (LS), while 59 the middle and deep circulations are maintained by the outflux and deep intrusion through LS, 60 respectively. Over the slope topography, the CAC circulation affects each other through the vertical coupling among them (Shu, Wang et al. 2018). For example, the study of Quan and 61 Xue (2018) demonstrated that perturbations in the upper layer can transit through the water 62 column and affect deeper circulation by altering layer thickness. This coupling is particularly 63 evident over the strong boundary current that in the northern basin and along the western 64 boundary. Numerical experiments by Wang et al. (2018) suggest that, although the basin-scale 65 circulation is primarily driven by the LS overflow, upwelling patterns are essential in shaping 66 the detailed structure of the deep-water circulation. Similar processes have been observed in 67 other marginal seas as well (e.g., Testor et al., 2018; Wang et al., 2024; Xu et al., 2009; Olvera-68 Prado, Moreles et al. 2023). 69

Despite advances in understanding the major features of mean layered circulation, the response of this circulation to changes in external forcings—particularly the strong upper ocean processes—and the specific mechanisms involved remain poorly understood. A comprehensive understanding of how topographically modulated vertical coupling over basin slopes influences the structure and intensity of mean circulation is crucial for exploring the long-term evolution of marginal sea circulation. For instance, Kuroshio intrusion has been observed to weaken with





76 decadal variations, while upper ocean currents have shown acceleration in response to a 77 warming climate (Peng et al., 2022; Chen et al., 2019; Nan et al., 2013). How the layered 78 circulation in marginal seas responds to these changes remains unclear. This study employs 79 process-oriented simulations to elucidate how upper-layer processes, which exhibit greater 80 intensity and variability, influence layered circulation over the meandering bottom slope in the 81 SCS. The paper is organized as follows: Section 2 explains the configuration of the simulation. 82 Section 3 offers the response of layered slope current intensity to the changes of the upper 83 circulation. Section 4 explores the underlying physical mechanisms. Finally, Section 5 provides 84 a summary of our findings.



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Figure1. Location and bathymetry (km) of the South China Sea, showing Luzon Strait, Taiwan
Strait, Mindoro Strait, and Karimata Strait.

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90 2. Methodology

Based on the Regional Ocean Modeling System (ROMS) (Shchepetkin & McWilliams,
2005), we performed a three-dimensional process-oriented simulation in this study (Cai et al.,
2023). The model setup included a circular basin on the west representing the SCS and a
rectangular basin on the east representing the Western Pacific Ocean, connected by a narrow
strait (representing the Luzon Strait, LS) with a depth of 2500 m. The slope in the SCS increases
from 400 to 4000 m and the central basin has the depth of 4000 m.

The model employed a uniform horizontal grid with a 5 km resolution, and the vertical





98 structure utilized stretched, terrain-following coordinates with 30 layers (Song and Haidvogel 99 1994) were adopted with 30 layers. In the upper layer, the influx/outflux was defined at the 100 southeastern (25 Sv) and northeastern (20 Sv) boundaries of the open ocean, and the SCS had an opening to the south to allow for outflux (Figure 2a), consequently, the intrusion from open 101 102 ocean into SCS will be generated. The model incorporated variable mixing coefficients (K_v) to simulate differences in exchange currents between the SCS and the Western Pacific Ocean, 103 particularly in the middle and deep parts of the LS (Tian, Yang et al. 2009, Yang, Zhao et al. 104 2016). In the SCS and the western half of the LS, K_v values were set as $5 \cdot 10^{-4} m^2 s^{-1}$ between 105 500 and 1,500 m, $2 \cdot 10^{-3} m^2 s^{-1}$ below 1,500 m, and $5 \cdot 10^{-3} m^2 s^{-1}$ in the layer 500 m from the 106 bottom. In the upper 500 m of SCS and LS, and throughout the entire water column of the open 107 ocean, a background K_v of $2 \cdot 10^{-5} m^2 s^{-1}$ was applied. A zonal wind stress with meridional 108 109 variations was applied over the open ocean, while the atmospheric buoyancy flux over the SCS was simplified by setting it to zero. 110

111 The model was initialized with horizontally uniform temperature and salinity profiles 112 derived from the mean data of the World Ocean Atlas, averaged over the region west of the LS. 113 The simulation ran for 25 years, and the analysis was conducted on the results from the final 5 114 years.

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Figure 2. (a) The topography and geometry used in the idealized process-oriented simulation.
The black lines show the realistic surrounding coastline. S1 and S2 are two vertical sections for
showing slope current. (b) Distribution of the wind stress.

122 **3. Results**

123 In the LS, the influx occurs in the upper and deep layers, while the outflux occurs in the 124 middle layer between them (Figure 3b). The upper intrusion is intrinsically linked to the western





125 boundary current in the open ocean, i.e. Kuroshio Current, while the exchange currents in the middle and bottom layers are primarily related to the density differences between the South 126 127 China Sea (SCS) and the Pacific (Zhu et al., 2017, 2019; Cai, Chen et al. 2023; Zhou et al., 2023). Under the density difference, the westward pressure gradient was formed that drives the 128 129 deep intrusion from the open ocean towards the SCS (Wang, Xie et al. 2011). Associated with the stimulated layered exchanging current, the horizontally averaged vorticity features with the 130 131 positive-negative-positive values in the respective depth of the basin, which represent the cyclonic, anticyclonic and cyclonic slope current (Figure 3b). 132



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Figure 3. (a) Vertical profile of domain averaged density ρ (kg·m⁻³) in SCS and Pacific between 1000m to 3000m. (b) Black line shows the vertical profile of the SCS basin-averaged vorticity (s⁻¹) and the red line shows the flux through LS (m²·s⁻¹). The positive/negative vorticity represents the cyclonic/anticyclonic circulation, and the positive/negative value flux represents the outflux/influx through LS.

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141 To investigate how the changes in upper layer current modulates the layered circulation, additional numerical experiments were conducted by adjusting the northeastern outflux to 22Sv, 142 143 23Sv and 24Sv, thus the upper LS intrusion was adjusted to 3Sv (Case U3), 2Sv (Case U2) 144 and 1Sv (Case U1). In these scenarios, the middle and deep exchanging currents, which were 145 primarily sustained by the contrasting densities between the SCS and the open ocean, remained 146 relatively unchanged (Figure 4a). Inside the basin, the enhancement of the upper layer inflow 147 directly intensified the upper layer cyclonic circulation (Figure 4 c-f). And subsequently 148 strengthened the middle anticyclonic slope current, particularly in the northwestern part of the basin (transect S1, Figure 4c, d). Additionally, although not directly connected to the upper 149 layer, the deep cyclonic slope current also exhibited increased strength (Figure 4c-f). Using the 150 151 domain-averaged vorticity as an indicator, it was observed that the intensification of the upper layer circulation resulted in a 10% increase in the intensity of the middle anticyclonic 152





- 153 circulation and a 27% increase in the deep cyclonic circulation (Figure 4b). Given that only the
- 154 upper layer circulation is directly amplified by the upper intrusion, the remote influence from
- 155 upper layer processes over the basin slope significantly impacts the intensity of the layered
- 156 circulation, which was explored below.





Figure 4. (a) Vertical profile flux across the LS under varying upper layer influx. (b) Changes in the basin-averaged vorticity between 0-500 m (Upper), 1000-2000 m (Middle) and 3000-4000 m (Deep) with different upper influx. (c-f) The along-slope current $(m \cdot s^{-1})$ over the section S1 and S1 in Standard case (upper intrusion is 5 Sv) and Case_U1 (upper LS intrusion is 1Sv), respectively. The locations of the two transects are shown in Figure 1a. The transect locations are indicated in Figure 1a. Positive values represent anticyclonic flow, while negative values indicate cyclonic flow.

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169 4. Discussion

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171 To understand how the upper layer processes modulate the layered current over the slope,

172 we first investigated the layered-integrated vorticity budget for each layer (Gan, Liu et al. 2016,

173 Cai, Chen et al. 2023):

$$174 \qquad \overbrace{\int_{A} [\nabla \times \int_{layer}^{\Box} (\overrightarrow{PGF}) dz] dA}^{\zeta _PGF} + \overbrace{\int_{A} [\nabla \times \int_{layer}^{\Box} (\overrightarrow{ADV}) dz] dA}^{\zeta _ADV} + \underbrace{\underbrace{\int_{A} \left[-f \nabla \cdot \int_{layer}^{\Box} (\overrightarrow{V}_{h}) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _BETA} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _BETA} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _BETA} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _BETA} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}_{\zeta _DIV} + \underbrace{\int_{A} \left[-\beta \cdot \int_{layer}^{\Box} (v) dz \right] dA}$$

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$$\overline{\int_{A} [\nabla \times \int_{laver}^{\Box} (\overline{VIS}) dz] dA} = 0$$
 (1)

where PGF is pressure gradient force, ADV is nonlinear advection, VIS is the turbulent

177 viscosity. COR is the Coriolis force and the domain-integrated ζ_{-COR} represents the lateral





178 planetary vorticity flux. Inside the basin, the ζ_{COR} can be further decomposed into the vertical stretching induced by the divergence of the horizontal flow (ζ _DIV) and β effect of the 179 meridional current (ζ_BETA). 180 181 The circulation in SCS is primarily governed by geostrophic balance, manifested as ζ _COR balanced by the ζ _PGF in both layers. It suggests that the mean cyclonic/anticyclonic 182 circulation is related to the lateral planetary vorticity influx/outflux (Figure 5) (Cai and Gan 183 184 2019, Cai, Chen et al. 2023; Zhu et al., 2017, 2019). Consequently, the strengthening of the 185 upper LS intrusion directly intensifies the upper layer cyclonic slope current by providing more positive planetary vorticity flux. In the semi-enclosed middle and deep layers, the vorticity 186 input is mainly provided by the ζ_DIV , as also highlighted in the previous studies (Zhu, Sun et 187 al. 2017, Wang, Du et al. 2018, Cai, Chen et al. 2023). The positive ζ_DIV in the upper and 188 deep layers reflects the downward and upward flux, which squeeze the middle layer and provide 189 the negative vertical flux (negative $\zeta_D IV$) for the anticyclonic circulation. Therefore, the 190 191 intensification of the middle and deep slope currents is largely attributed to vertical stretching 192 over the basin slope.

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Figure 5. The layered-integrated vorticity budget (Equation 1) for the standard case in the layers of (a) 0–500, (b) 1,000–2,000, and (c) 2,500–4,000 m.

200 a. Intensification of middle anticyclonic slope current

202 For the middle anticyclonic slope current, the negative middle-layered $\zeta_{\rm DIV}$ ($\zeta_{\rm DIV_M}$) 203 mainly occurs over the northwestern slope (Figure 6a), where a relatively strong anticyclonic slope current is observed. In this area, the upper layer generally exhibits positive ζ _DIV 204 (ζ_{DIV_U}) , indicating downward squeezing from the upper layer that helps maintain the middle 205 anticyclonic slope current (Figure 6 b). The increased squeezing from the upper slope current 206 207 is expected to intensify the middle anticyclonic slope current (Figure 4 c, d). Since the 208 stretching/squeezing imposed on the slope current is closely related to the vertical motion 209 (Liang, Spall et al. 2017), the mean vertical motions over the slope can be understood as follows:





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$$\overline{w} = -\overline{v}_h \cdot \int_{-H}^{z} \overline{\tilde{v}}_h dz \approx -\overline{v}_h \cdot \int_{-H}^{z} \overline{\tilde{v}}_{h_geo} dz = \underbrace{-(\overline{\tilde{v}}_{b_geo}H_y + \overline{\tilde{u}}_{b_geo}H_x)}_{Geo} + \underbrace{\frac{\beta}{f} \int_{-H}^{z} \overline{\tilde{v}}_{geo} dz}_{\beta \ effect}$$
(2)

211 where H is the topography, _geo represents the geostrophic component of the horizontal current, which is maintained by the distribution of bottom pressure. H_x and H_y are the 212 213 horizontal gradients of the bottom slope, $\overline{u}_{b_{geo}}$ and $\overline{v}_{b_{geo}}$ are the bottom geostrophic currents. 214 When the current flows over the slope, the bottom pressure distribution interacts with the 215 meandering topography and the resulting geostrophic cross-isobath transport (CGT_b) to influence the vertical motions over the slope. Generally, the CGT_h largely governs the mean 216 217 vertical motion over northwestern slope, which in turn squeezes and stretches the middle and 218 upper slope currents. This process is crucial in maintaining the vorticity input through ζ DIV (Figure 6 d), thereby sustaining the dynamics of the slope currents. 219

Associated with the intensification of the upper layer cyclonic slope current, the stronger slope current modulates the bottom pressure distribution over the slope, leading to stronger vertical squeezing (ζ _DIV) within the water column. The increase in squeezing gradually deepens the isopycnal surface and strengthens the middle anticyclonic slope current (Figure 6e).





Figure 6. (a-c) ζ_{DIV} integrated between 1000-2000 m (ζ_{DIV_M}), 0-500 m (ζ_{DIV_U}), and *CGT_b* over the region of middle anticyclonic slope current between 1000-2000 m. The black line indicates the 2000 m isobath line. (d) Change of the ζ_{DIV} integrated between 100-2000 m (ζ_{DIV_M}), between 0-500 m (ζ_{DIV_U}) with the depth over the northwestern slope. (e) Changes of the ζ_{DIV_M} , ζ_{DIV_U} , CGT_b and the depth of isopycnal surface of 1027.4 kg/m³ with different





upper LS influx. (d) and (e) was plotted using the data averaged over the region of black box
 in (a)

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235 b. Intensification of the deep cyclonic slope current

For the deep cyclonic slope current, the deep $\zeta_{\rm DIV}$ ($\zeta_{\rm DIV}$) features a downward flux 236 237 (negative value) over the northern slope and upward flux (positive value) over the southern part (Figure 7a, d). This pattern aligns with the classical source-sink driven theory (Stommel and 238 Arons 1959; Yuan et al., 1997), which suggests the sinking from the northern basin and 239 upwelling from the southern basin would sustain the abyssal cyclonic circulation. Same as the 240 241 middle anticyclonic slope current, the ζ_{DIV_D} is largely controlled by the CGT_b that the 242 pressure distribution maintained the mean cascading in the northern side and the upwelling over the southern slope (Figure 7b, d). To explain the maintenance of the bottom pressure 243 244 distribution, the vertical integrated vorticity dynamic was employed (Gan, Liang et al. 2013, Cai and Gan 2021): 245

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$$\overbrace{\nabla \times \int_{-H}^{0} (\overrightarrow{PGF}) dz}^{\Omega_{_}PGF} + \overbrace{\nabla \times \int_{-H}^{0} (\overrightarrow{ADV}) dz}^{\Omega_{_}ADV} + \underbrace{\beta \int_{-H}^{0} v dz}_{\Omega_{_}BETA} - \overbrace{\nabla \times \int_{-H}^{0} (\overrightarrow{VIS}) dz}^{\Omega_{_}VIS} = 0$$
(3)

247 Where $\Omega_PGF = \nabla \times \int_{-H}^{0} (\overrightarrow{PGF}) dz = -(\overline{v}_{b_geo}H_y + \overline{u}_{b_geo}H_x)$, thus represents the effect of CGT_b . 248 It illustrates that the bottom pressure distribution is maintained by the nonlinear advection of 249 relative vorticity (Ω_ADV), advection of planetary vorticity (Ω_BETA), and net viscosity of the 250 current flowing over the slope topography. Among them, the Ω_BETA and Ω_ADV play the 251 major role in sustaining the bottom pressure distribution of the deep cyclonic slope current 252 (Figure 8 c, d).

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Figure 7 (a-c) ζ DIV integrated below 3000 m of the deep cyclonic slope current (ζ DIV_D), 256 CGT_{h} and $\Omega_{BETA} + \Omega_{ADV}$ over the region of middle anticyclonic slope current below 3000 257 258 m. The black line indicates the 4000 m isobath line. (d) Meridional changes of the $\zeta_{\rm DIV_D}$, 259 CGT_b and $\Omega_BETA + \Omega_ADV$ over the slope. 260

261 As the upper intrusion intensifies, the bottom pressure distribution adjusts in response to changes in the layered circulation, resulting in a gradual strengthening of the negative 262 263 CGT_b over the northern slope and positive CGT_b over the southern part (Figure 8 a-b). Over the 264 southern slope (Figure 8a), the strengthening of the positive CGT_b is induced by the increase of the $\Omega_{BETA} + \Omega_{ADV}$ in the water column, in which the middle layer $[\Omega_{BETA} + \Omega_{ADV} (M)]$ 265 266 provides approximately 40% of the total strengthening trend, while the upper layer [$\Omega_{-}BETA$ 267 $+ \Omega_A DV$ (U)] has a negligible impact. Conversely, over the northern slope, the strengthening of the $\Omega_BETA + \Omega_ADV$ is primarily influenced by the upper layer, with the middle layer has 268 the negative effect (Figure 8b). Thus, the intensification of downward CGT_b over the northern 269 270 slope is directly modulated by the strengthening of the upper cyclonic slope current, while the intensification of upward CGT_b over the southern slope is mainly driven by the strengthening 271 of the middle anticyclonic slope current. These changes in CGT_b enhance deep stretching, 272 273 thereby strengthening the deep cyclonic slope current accordingly (Figure 8c). 274







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Figure 8. (a) Changes of the CGT_b , $\Omega_{Beta} + \Omega_{Adv}$ in the total water column, $\Omega_{Beta} + \Omega_{Adv}$ in the layer between 0-500 m ($\Omega_{Beta_U} + \Omega_{Adv_U}$), $\Omega_{Beta} + \Omega_{Adv}$ in the between 1000-2000 m ($\Omega_{Beta_M} + \Omega_{Adv_M}$) over the southern slope (black box in Figure 8a). (b) same as (a) but over the northern slope (black box in Figure 8a). (c) Change of the CGT_b and ζ_DIV_D over the northern and southern slope.

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285 **5. Summary**

Marginal sea circulation plays a crucial role in mass transport and regional climate dynamics. Through process-oriented numerical simulations, we examined how upper-layer processes, which are characterized by greater intensity and variability, impact the layered circulation over the meandering bottom slope in SCS. The results underscore the intricate balance between topographical features and oceanic circulation, offering valuable insights for predicting the behavior of marginal sea circulation under varying forcings conditions

The numerical experiments show that the stimulated layered exchange currents through the LS introduce lateral planetary vorticity flux and facilitate the development of a layered basin circulation within the SCS. Inside the basin, the intensification of upper-layer inflow directly enhances the upper-layer cyclonic circulation, which in turn strengthens the middle anticyclonic slope current, particularly in the northwestern part of the basin. Furthermore, even though the deep cyclonic slope current is not directly connected to the upper layer, it also exhibits increased





298	strength. Using domain-averaged vorticity as an indicator, it was found that the intensification
299	of the upper-layer circulation resulted in a 10% increase in the intensity of the middle
300	anticyclonic circulation and a 27% increase in the deep cyclonic circulation.
301	The vorticity dynamics illustrate that the changes in the middle and deep slope current is
302	largely related to the vertical stretching over the slope ($\zeta_D IV$). When the current flows over
303	the slope, the bottom pressure distribution interacts with the meandering topography and the
304	resulting geostrophic cross-isobath transport (CGT_b) influences the vertical motions. For the
305	middle anticyclonic slope current, the negative middle $\zeta_{\rm DIV}$ predominantly occurs over the
306	northwestern slope, where a relatively strong anticyclonic slope current is present. As the upper-
307	layer cyclonic slope current intensifies, it modulates the bottom pressure distribution over the
308	slope, providing stronger vertical squeezing ζ _DIV within the water column, which gradually
309	strengthens the middle anticyclonic slope current. For the deep cyclonic slope current, the deep
310	$\zeta_{\rm DIV}$ is also largely controlled by the CGT _b over the slope, with the pressure distribution
311	maintaining cascading flows in the northern part and upwelling over the southern slope. Over
312	the southern slope, the strengthening of the positive CGT_b is induced by the increase of the
313	nonlinear advection of relative vorticity and planetary vorticity in the entire water column, in
314	which the middle layer provides approximately 40% of the total strengthening trend, while the
315	upper layer has a negligible impact. Conversely, on the northern slope, the strengthening of the
316	negative CGT_b is primarily influenced by the upper layer, with the middle layer has the negative
317	effect. The results deepen understanding of the intricate balance between topographical features
318	and layered circulation, which help to improve the predictions of the long-term behavior of
319	marginal sea circulation under varying climatic conditions.
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322	Data availability. All data used in this paper are available at
323	https://doi.org/10.5281/zenodo.13835538.
324	
325	Author contributions. QT conducted the investigation, developed the methodology, and carried
326	out the writing (original draft preparation). ZC was responsible for the conceptualization,
327	supervision, and writing (review and editing). ZL was responsible for the conceptualization and
328	writing (review and editing).
329	
330	Competing interests. The contact author has declared that none of the authors has any
331	competing interests.
332	
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