Technical Note:

Spectral slopes in deep, weakly-stratified ocean and coupling between sub-mesoscale motions and small-scale mechanisms

by Hans van Haren

Royal Netherlands Institute for Sea Research (NIOZ), P.O. Box 59, 1790 AB Den Burg, the Netherlands.
e-mail: hans.van.haren@nioz.nl

Short summary. Large ocean circulations include small-scale physical processes like transport by sub-mesoscale eddies and turbulence by internal wave breaking. Knowledge is lacking on precise interaction between different processes. In deep weakly stratified waters, continuous spectral slopes are observed that extend from sub-mesoscales across the internal wave band to turbulence range. Such cross-spectral correspondence is suggested a potential feedback mechanism stabilizing large-scale ocean circulations.

Abstract. Large, basin-wide ocean circulations are complex nonlinear dynamical systems. They include small-scale physical processes such as, for example, transport by sub-mesoscale eddies and turbulence-generating breaking of internal waves. To date however, knowledge is lacking on precise interaction between different processes. In this note, a potential contributor to interaction is investigated using spectra from deep-sea moored observations. In weakly stratified waters, continuous spectral slopes are observed that extend from sub-mesoscales across the internal wave band to turbulence range. In the latter, the governing slope can be distinctly different from the inertial subrange of shear turbulence and is described as the buoyancy subrange of convection turbulence. At sub-inertial frequencies, the slope's extension either describes quasi-gyroscopic waves or sub-mesoscale eddies. Such cross-spectral correspondence is suggested a potential feedback mechanism stabilizing large-scale ocean circulations.

1 Introduction

The extent of anthropogenic influence on the Earth's climate warrants studies of the ocean as a major player. Large, basin-wide ocean circulations are important for transporting properties like heat, carbon and nutrients. Schematically, the Atlantic(-Ocean) Meridional Overturning Circulation (AMOC) is depicted to transport heat from the equator to the poles near the surface and carbon in the abyssal return (e.g., Aldama-Campino et al., 2023). It includes physical processes like 'deep dense-water formation' in the polar region. Recent mathematical and numerical modelling such as based on varying single parameters like sea-

surface temperature (e.g., Ditlevsen and Ditlevsen, 2023) and freshwater influx (e.g., van Westen et al., 2024) suggest a potential future collapse of the AMOC. It is argued that this may have consequences for Northwest-European climate.

Whilst the modelling might be robust mathematically, it lacks physical processes of the drivers of the AMOC and observational evidence thereof. This will have consequences for the feedback mechanisms at work in the nonlinear dynamical system of ocean circulation. As has been reviewed for AMOC numerical models (Gent, 2018), important feedback mechanisms are vertical turbulent mixing, sub-mesoscale gyre 'eddy' transport, and the coupling with the atmosphere. Here we elaborate on the importance of turbulence induced by internal wave breaking, possibly coupling with sub-mesoscale eddies (e.g., Chunchuzov et al., 2021), and stability variations in vertical density stratification for such feedback, by reviewing insights from recent modeling and deep-sea observations. In particular as an example for complexity of dynamical system interactions, the core of ocean motions is spectrally investigated focusing on most energetic mesoscale, internal wave, and turbulence scales, for deep weakly stratified waters.

In contrast with the atmosphere, the ocean is not an effective heat engine (Wunsch and Ferrari, 2004) despite its heat transportation. As a result, the AMOC is not predominantly buoyancy-driven via push by deep dense-water formation near the poles (Marshall and Schott, 1999; Marotzke and Scott, 1999), which notably occurs in sporadic pulses rather than continuously. Instead, the AMOC is mainly wind-steered (e.g., Liu et al., 2024) and tide-driven, with turbulent mixing by internal wave breaking, and possibly associated upwelling close to boundaries (Ferrari et al., 2016; McDougall and Ferrari, 2017), being considered an important physics process of pull that dominates over push by a heat engine. Winds, near the ocean surface, and tides, via interaction with seafloor topography deeper down, contribute about equally to generate internal waves that are found everywhere in the ocean interior. Such waves break predominantly at ubiquitous underwater seamounts and continental slopes.

Without turbulent mixing, the AMOC would be confined to a 100-m thick near-surface layer and the deep-ocean would be a stagnant pool of cold water (Munk and Wunsch,

1998). This is not the case however, and the solar heat is mixed from the surface downward so that the ocean is stably stratified in density all the way into its deepest trenches, as has been shown in hydrographic deep-ocean observations (Taira et al., 2005; van Haren et al., 2021a). Although turbulent mixing by internal wave breaking in the ocean-interior is insufficient by at least a factor of two to maintain the vertical density stratification (e.g., Gregg, 1989, Polzin et al., 1997), such breaking along ocean boundaries has been suggested to be more than sufficient (Munk, 1966; Polzin et al., 1997). Especially large internal wave breaking is expected to occur above steeply sloping topography (Eriksen, 1982; Thorpe, 1987; Sarkar and Scotti, 2017). Because there are more and larger seamounts than mountains on land, equally abundant sloping seafloors lead to abundant turbulent mixing, as has been charted from recent observations and modelling results summarized below.

As recent observations (van Haren and Dijkstra, 2021; van Haren et al., 2024) demonstrate that breaking waves can lead to considerable buoyancy driven convection turbulence, this note attempts further understanding of a little studied deep-sea complex process and its potential interaction with sub-mesoscale motions. The sub-inertial range of sub-mesoscale motions has rarely been a subject of oceanographic spectral observations. Knowledge about such small-scale processes and their interactions may be vital for understanding potential feedback mechanisms affecting the stability of large-scale ocean circulations.

2 Recent internal wave breaking results

Detailed observations and numerical modeling have revealed the extent of internal tide breaking processes above ocean topography (van Haren and Gostiaux, 2012; Winters, 2015; Wynne-Cattanach et al., 2024). Using high-resolution observations (e.g., van Haren and Gostiaux, 2012), internal tide breaking above steep deep-sea slopes is observed to generate spring-neap-average turbulent vertical diffusivity value of about 3×10^{-3} m^2 s^{-1}. This value is twice the value theoretically required to yield upwelling in a thin layer above sloping topography (McDougall and Ferrari, 2017). Such quantification of turbulent mixing shows

that it occurs with typical tidal-period-average values that are more than 100 times larger over super-critical slopes than open-ocean values. A super-critical seafloor slope is steeper than the slope of internal wave characteristics. While ocean-wide tides energetically dominate internal waves, not all seafloor slopes are super-critical for these waves. In contrast, nearly all seafloor slopes are super-critical for at least one component of secondary energetic near-inertial waves, which are generated via geostrophic adjustment following the passage or collapse of a disturbance such as fronts or atmospheric storms on the rotating Earth. Under common stratification, near-inertial waves are at the lowest frequency of freely propagating internal waves. The highest frequency propagating internal waves, near the buoyancy frequency, experience nearly vertical walls as super-critical seafloor slopes.

Within a tidal, or near-inertial, period, turbulence peaks in bursts of shorter duration than half an hour when highly nonlinear internal waves propagate as internal bores up a super-critical slope, once or twice a tidal cycle. The breaking of bores leads primarily to convection, buoyancy-driven turbulence, rather than frictional shear-turbulence over the sloping seafloor and occur at a wide variety of deep-sea and deep-ocean locations (e.g., van Haren et al., 2013; van Haren et al., 2024). Between bores, the turbulent mixing varies by an order of magnitude in intensity, with effects extending about 100 m vertically and several kilometers horizontally from the seafloor. Although intermittently occurring at a given position of the sloping seafloor and about 10% varying in arrival time, the turbulence is generated internally by the tide, for about 60% (Wunsch and Ferrari, 2004), and by winds, for about 40%, in a stratified ocean-environment. The turbulent bores also resuspend sediment and thereby replenish nutrients away from the seafloor (Hosegood et al., 2004), important for deep-sea life. Enhanced turbulent mixing above sloping boundaries has a demonstrated effect on the outcome of general ocean circulation models (e.g., Scott and Marotzke, 2002), with predicted subtle effects on upwelling near the seafloor (Ferrari et al., 2016).

The complexity of turbulence generation, mixing and restratification, are still subjects of deep-ocean research. While shear-induced turbulence has been relatively well studied in the stratified ocean, deviations such as convection-turbulence are little observed, with recent

exceptions (van Haren and Dijkstra, 2021; van Haren et al., 2024). Convection turbulence is dominant in the atmosphere especially during daytime, and has also been observed in the near-surface ocean during nighttime (e.g., Brainerd and Gregg, 1995), but it has never been uantitatively directly observed in deep dense-water formation zones (Thorpe, 2005) and above geothermal vents.

Deep dense-water formation does not only occur in polar seas, but occasionally also in the at least 10°C warmer Mediterranean (Gascard, 1978), with an important contribution of atmospheric exchange due to orographic generated winds affecting the preconditioning by cooling and drying of near-surface waters. Similarly, internal waves occur in oceans and in the Mediterranean under stratification conditions that vary over at least one order of magnitude in time and space, but tides are relatively weak in the Mediterranean, and yet 'sufficient' turbulent diapycnal mixing, sufficient for maintenance of deep-sea stratification and thereby driving overturning circulation, is generated via the breaking above topography of near-inertial motions mainly (van Haren et al., 2013). Further complications are expected from interactions of internal waves with sub-mesoscale eddies and potential consequences of varying intensity thereof, e.g., on seasonal scales.

3 Mediterranean observations as an example proxy for ocean conditions

In many physical oceanographic aspects of heat and salt budgets, large-scale water-flow circulation, strong boundary flow, eddies at sub-mesoscales, near-inertial motions including gyroscopic waves and internal wave turbulence, the Mediterranean Sea can be considered a sample for the state of the much larger oceans (e.g., Gascard, 1973; Crepon et al., 1982; Garrett, 1994; Milllot, 1999; van Haren and Millot, 2004; Testor and Gascard, 2006). Like in oceans, the Mediterranean seafloor reaches great depths and can be rugged with steep slopes in places, including continental slopes incised by deep canyons.

In the Northwest Mediterranean, vertical density stratification varies markedly with seasons and years, having relatively large near-surface values in summer and relatively low values in winter. The proximity of extensive mountain ranges on land generates highly

variable winds that can cool and dry surface waters. In winter in weaker stratified waters, this may lead to unstable conditions of buoyancy driven convection in an exchange of dense-water sinking down, and less dense-waters up. Like in the polar regions, such exchange can be observed daily in the upper 10 m from the sea-surface, regularly down to a few 100 m from the surface, and seldom, once every 5-8 years (e.g., Rhein, 1995; Mertens and Schott, 1998), down to the abyssal seafloor at about 2500 m. In contrast, horizontal density gradients associate with forcing of a dynamically unstable boundary current and eddies at multiple 1-100 km sub-mesoscales (e.g., Crepon et al., 1982; Testor and Gascard, 2006). These eddy motions may push relatively warm waters down, thereby increasing the weak stratification in the deep-sea.

In summer, atmospheric disturbances are less intense, near-surface stratification is large due to solar heating, and eddy activity associated with some continental boundary flows is weaker (Albérola et al., 1995). This opens the possibility for detection of near-inertial wave dominance in kinetic energy. In relatively strong stratification, mainly gravity-driven parts of near-inertial waves generate largest vertical current differences ‘shear’ that destabilize stratification due to their relatively short vertical length-scale, not only in the Mediterranean but also as observed in the Atlantic Ocean (van Haren, 2007). This destabilization may lead to small-10-m vertical scale layering of near-homogeneous waters throughout seas and oceans. On larger-100-m vertical scales near-homogeneous waters occur in deep waters of the Mediterranean as well as of North-Atlantic basins like the Bay of Biscay and Canary Basin. In near-homogeneous water-layers with weak stratification, gyroscopic, Earth-rotation-driven, parts of near-inertial waves dominate and result in 0.1-1 km diameter smaller than sub-mesocale tubes of slantwise rather than vertical convection (Emanuel, 1994; Marshall and Schott, 1999; van Haren and Millot, 2004). Hence, one may expect frequency spectra of non-tidal dominated data from instruments moored in the Mediterranean reveal convection and thus deep transport under winter and summer conditions.

It is noted that ocean-spectra may show peaks such as at narrowband tidal and at, broader band, inertial frequencies, but they lack gaps. This lack of spectral gaps potentially

couples motions at sub-inertial with inertial-buoyancy internal wave with super-buoyancy turbulence frequency ranges. However, it is unclear how such a coupling may work as some motions represent two-dimensional '2D' eddies, some linear waves, some non-linear waves, some anisotropic stratified turbulence, and some isotropic 3D turbulence. This is investigated by renewed spectral analysis below, using, in analogy, slopes typical for investigating energy cascades in turbulence research.

4 Uncommon slopes in revisited spectra

Kinetic energy (KE) spectra from historic moored current meter observations down to mid-depth z = -1100 m in the Ligurian Sea under upper-sea strongly stratified 'summer' and weakly stratified 'winter' conditions surely lack gaps (Fig. 1). Year-round at z = -1100 m, the buoyancy frequency N, reflecting the square-root of vertical density stratification, is small N ~ O(f), f denoting the inertial frequency involving Earth rotation. This narrows the local internal wave band, while, especially in winter, sub-mesoscale activity is large in the area, and, occasionally, the few moored current meter temperature records showed inversions (van Haren and Millot, 2003). Although these hourly sampled data barely resolve the turbulence ranges at frequencies $\omega > N$, the internal wave continuum was suggested to scale like ω^p, with, on a log-log plot, 'spectral slope' $p = -2.2 \pm 0.4$, independent of location and season albeit with different KE (power) levels.

Within the uncertainty range, several possible explanations can be given for the observed spectral slope. Freely propagating internal gravity waves have been fitted to $p = -2 \pm 0.5$ but only for $f << \omega << N$ (Garrett and Munk, 1972). Considering that the data in Fig. 1 are from a site where locally $N = (3 \pm 2)f$, irrespective of season (van Haren and Millot, 2003), alternative explanations were sought for observed spectral slopes at sub-inertial frequencies 0.2 cpd $< \omega < f$. Cpd is short for 'cycles per day'. An obvious candidate is 'fine-structure contamination' of step functions passing sensors, which gives a theoretical vale of $p = -2$ (Phillips, 1971; Reid, 1971). For their winter data, van Haren and Millot (2003) attributed

231 such a slope to evidence intense mesoscale activity, because of the continuation of slope up to
232 ω = 5 cpd before rolling off near the Nyquist frequency. However, they did not elaborate.
233 Below, the data in Fig. 1 are re-analyzed from the perspective of convection-turbulence.

234 Theoretical considerations of non-zero-mean flow convection-turbulence suggest a
235 spectral scaling in the buoyancy subrange having p = -11/5 = -2.2 for KE, and p = -7/5 for an
236 active scalar quantity. This 'BO'-scaling follows atmospheric and theoretical works by
237 Bolgiano (1959) and Obukhov (1959). The scaling was set-up for a stably stratified
238 atmospheric environment for the anisotropic part in which turbulent kinetic energy is partially
239 transferred to potential energy leading to turbulent convection. Later works extended BO-
240 scaling to purely buoyancy-driven turbulence, e.g., for Rayleigh-Bénard convection (Lohse
241 and Xia, 2010) and Rayleigh-Taylor instabilities (Poujade, 2006; Celani et al., 2006).

242 Laboratory experiments on such gravitationally driven convection are inconclusive on
243 BO-scaling. On the one hand, this scaling is confirmed for both KE and temperature in
244 experiments by Ashkenazi and Steinberg (1999), while on the other hand it is only confirmed
245 for scalars by Pawar and Arakeri (2016) who found a slope of p = -5/3 for KE. The p = -5/3-
246 slope suggests dominance of shear-induced turbulence of the inertial subrange for equilibrium
247 isotropic turbulence cascade in the 'KO'-scaling (Kolmogorov, 1941; Obukhov, 1949) but
248 should also be found in spectra of scalars that are passive in this range. While Liot et al.
249 (2016) show KO-scaling in their model that may have to do with their Lagrangian data as
250 proper transfer brings the data closer to BO-scaling, Poujade (2006) and Cenari et al. (2006)
251 show clear BO-scaling in their models. This suggests particular conditions do affect the
252 dominance of shear- or convection-turbulence. It is noted that BO-scaling is also simply
253 considered as a significant deviation from KO-scaling, which is more commonly observed in
254 stratified shear flows.

255 Obviously, scalars cannot be passive and active at the same time and in the same
256 space. This discrepancy between types of scaling between scalars and KE may be because the
257 laboratory experiments of Pawar and Arakeri (2016) were in zero mean flow. Also, under
258 sufficiently stable conditions without shear, no inertial subrange is expected (Bolgiano, 1959).

However, the spectral extent of BO-scaling is largely unknown albeit it is more generally found adjacent to higher-frequency inertial subrange. While KO-scaling is based on a forward cascade of energy, the direction of energy cascade is inconclusive for BO-scaling and may be partially forward and partially backward, at least as reasoned for pure buoyancy-driven convection-turbulence (Lohse and Xia, 2010). Probably, directions of cascade change with locality in the flow, and perhaps depend on scale, which would also imply that KO- and BO-scaling cannot be found at the same site.

Revisiting data from non-zero mean flow and weakly stratified deep-sea in Fig. 1 demonstrates the possibility of fit of $p = -11/5$ outside near-inertial harmonic peaks. In winter, such a fit is observed consistently through the entire range of $0.2 < \omega < 5$ cpd. In traditional terms, this frequency range covers the transition from mesoscale $\omega < f$, via internal wave $f < \omega < N$, to turbulence $\omega > N$ motions. In summer, the $p = -11/5$-slope is found at two different KE levels for bands $0.2 < \omega < \omega_{min}$ and $2\Omega < \omega < 5$ cpd at sub- and super-IGW frequencies, respectively. Here, $\omega_{min} \leq f$ denotes the minimum frequency bound for inertio-gravity waves IGW (LeBlond and Mysak, 1978), and Ω the Earth rotational frequency. Maximum IGW frequency is denoted by $\omega_{max} \geq 2\Omega, N$. The ω_{min} and ω_{max} are functions of N, latitude φ and direction of wave propagation (LeBlond and Mysak, 1978; Gerkema et al., 2008),

$$\omega_{max}, \omega_{min} = (A \pm (A^2 - B^2)^{1/2})^{1/2}/\sqrt{2}, \quad (1)$$

in which $A = N^2 + f^2 + f_s^2$, $B = 2fN$, and $f_s = f_h \sin\alpha$, α the angle to φ. For $f_s = 0$ or $N >> 2\Omega$, the traditional bounds [f, N] are retrieved from (1). The plotted IGW-bounds [ω_{min} ω_{max}] are for weakly stratified, near-homogeneous layers in which $N = f$. This weak stratification would lead to an impossible wave solution under the traditional approximation, but (1) allows wave propagation, albeit horizontally for one component (e.g., Gerkema et al., 2008).

The bridge between the KE-levels at sub- and super-IGW is formed by the finitely broad near-inertial peak. The base of this peak is proposed to slope like $p = -1$ reaching super-IGW BO-scaling at about $\omega \approx 4$ cpd $\approx N$. Such $p = -1$-slope has been observed for the KE-spectral continuum between [f N] from the deep Bay of Biscay, Northeast Atlantic Ocean

(van Haren et al., 2002). Theoretically, this slope represents spectral scaling of intermittency of a weakly chaotic nonlinear system (Schuster, 1984), i.e., 3D dynamical systems that evolve into self-organized critical structures of states which are minimally stable (Bak et al., 1987). Such a spectral bridge, or hump, is expected for turbulence in unstable stratification, as has been illustrated using atmospheric observations (Lin, 1969). It is attributed to the flow field absorbing energy from the scalar temperature field as potential energy is transferred to kinetic energy. It is not clear to what extent near-inertial internal waves contribute in a similar way to spectral redistribution of energy in our oceanographic data. As the observations from the central Ligurian Sea show similar results, the hump is unlikely associated with seafloor slopes matching the slope of near-inertial internal wave rays.

These spectral observations suggest a dominance of convection cascade from sub-meso- via IGW- to, probably because unresolved, turbulence-scales under high-energetic winter-conditions as they show a continuous slope across their frequency ranges. Such a cascade is also suggested under quieter summer conditions when, however, it is masked by IGW that lead a cascade at $\omega > \omega_{min}$. Especially the sub-inertial range of apparent BO-scaling seems out of the turbulence range, unless waters are near-homogeneous $N \rightarrow 0$ so that $\omega_{min} \rightarrow 0$, from (1). This would extend not only IGW, notably gyroscopic waves, but also turbulence, probably in the form of slantwise convection, to the sub-mesoscale range.

For the mesoscale range, the observations in Fig. 1 are supported by numerical modeling results that have suggested eddy-KE has a broad range of spectral slopes between $-3 < p < -5/3$ (Storer et al., 2022), and by satellite altimetry observations that indicated, after noise-correction and transfer to KE, a best-fit of $p = -2.28$ (Xu and Fu, 2012). No mention was made of BO-scaling, but the correspondence seems evident.

As the KE in Fig. 1 is at least one order of magnitude larger in winter than in summer, a near-inertial peak, if it exists, will be part of the spectral continuum during the former. The winter observations suggest a continuous spread of sub-mesoscale energy across the IGW band including inertial motions and into the turbulence range. In winter, near-surface

stratification is considerably weaker than in summer, so that local atmospheric-generated near-inertial motions will be smaller. It is noted that the signals near the Nyquist frequency not only contain instrumental white noise, but also unresolved turbulence motions, which are also larger in winter than in summer.

Inspired by Western Mediterranean observations, Saint-Guily (1972) proposed from theoretical work that winter-time inertial KE is spread over a broad featureless band, like quasi-gyroscopic waves that may be present between IGW-bounds (1) for $N \sim f$ (LeBlond and Mysak, 1978; Gerkema et al., 2008). However, observations from the year-round upper-layer-stratified central Western Mediterranean demonstrate that, also in deep homogeneous $N = 0$ waters, a near-inertial peak can be observed in KE-spectra (van Haren and Millot, 2004). This may be attributed to a year-round source of atmospheric-generated inertial waves that are the only internal waves that can propagate without attenuation from well-stratified to near-homogeneous layers and vice versa (van Haren, 2023b).

Based on limited spectral observations, Gascard (1973) suggested the generation of 12-h stability waves, close to the buoyancy frequency of very weak stratification, which may briefly force dense-water formation, thereby implicitly suggesting a link between internal waves and sub-mesoscale eddies. As such eddies have estimated relative vorticity of $|\zeta| = f/2$ in the Western Mediterranean (Testor and Gascard, 2006), this addition to the planetary vorticity (f) automatically widens the 'effective' near-inertial band $0.5f < f_{eff} < 1.5f$, of which the bounds are close to IGW-bounds for $N = 0.8f$. One of the properties can be a modification of near-inertial frequency (Perkins, 1976), and trapping with downward propagation of near-inertial waves in anticyclonic eddies (Kunze, 1985; Voet et al., 2024). Such frequency modification may add to local physics of inertial wave caustics due to latitudinal variation (LeBlond and Mysak, 1978), which however can only lead up to 15% change in f in the Mediterranean. Although found to be limited to the rather flat KE-spectral dip in the immediate half-order-of-magnitude sub-inertial frequency band, standing vortical modes, i.e. low-frequency non-propagating motions, of vertical length-scale <10 m are suggested to be as energetic as internal waves (Polzin et al., 2003). Alternatively, it has been suggested for

North-Atlantic observations that vortical modes may interact with internal waves, affecting internal-wave shear that was peaking over O(10) m vertical scales at IGW-frequencies in a band with limits determined by weak stratification as in N = f (van Haren, 2007).

For hypothetical ω_{min} = 0.2 cpd, at which the observed spectral slope changes away from p = -11/5 (Fig. 1), one would require N = 0.21f, which is almost unmeasurable and non-existent for any prolonged period even in the deep Northwestern Mediterranean, to the knowledge of the author. However, it may reflect ω_{min} computed using f_{eff} = 0.5f and N = f_{eff}, noting that such conditions can only apply for part of the record. If so, it would reflect a direct coupling between sub-mesoscale and IGW-motions with slantwise convection (Marshall and Schott, 1999; van Haren and Millot, 2004; Gerkema et al., 2008). The p = -11/5 is significantly distinguishable from -2 over a frequency range of nearly two orders of magnitude, and from -5/3 over a range of just over half an order of magnitude (Fig. 1). The roll-off to noise (slope 0), for ω > 5 cpd, may partially be seen as following a slope of p = -5/3 before 0. The roll-off around 0.1 cpd suggests an unresolved broad mesoscale peak-value between 0.01 and 0.1 cpd. While these 1980's moored current meter data barely resolved the turbulence part of the KE-spectrum, and thus also not the p=-5/3 inertial subrange slope, their temperature sensors were too poor to simultaneously verify any spectral scaling for scalars.

About 40 years later, high-resolution and high-precision moored temperature sensor 'T-' data provided opportunity to verify scalar spectral scaling of turbulence energetic motions in the area. These T-data evidenced occasional warming of the deep Northwest Mediterranean seafloor (Fig. 2a), which, after comparison with data from higher-up appeared to be coming from above, or slanted sideways, under relatively stratified conditions, and from general non-vents geothermal heating from below (van Haren, 2023a). The data were collected during mid-fall, when near-surface waters were well stratified and no cold, dense-water production through convection was observed. Locally near the seafloor, the broad two-day warming around day 308 is most stratified, whilst during other periods waters are only weakly stratified, including the quasi-inertial variations between days 316 and 322. These

weakly stratified near-inertial, or near-buoyancy as $N \approx f$, temperature variations may evidence slantwise quasi-gyroscopic near-inertial waves, which can have a large vertical component (LeBlond and Mysak, 1978), as opposed to more common near-horizontal near-inertial waves in strongly stratified waters that are barely noticeable in temperature records.

The 18-day average spectrum of the 2-s sampled data poorly resolves sub-mesoscales but shows, near the seafloor, a well-resolved slope of $p = -1.4 \pm 0.025$ between a large range of $0.5 < \omega < 6000$ cpd, across the IGW band and well into the turbulence band (Fig. 2b). No transition to a -5/3-slope is observed before roll-off to noise, but this does not exclude an inertial subrange at higher frequencies hidden under white noise, although shear will be limited so close to the seafloor. The observed $p = -7/5$-slope is found significantly different from $p = -2$ and -5/3 over the indicated frequency range of four orders of magnitude and over the range between $100 < \omega < 10^4$ cpd thereby representing convection turbulence. Over a frequency range of half an order of magnitude the slope-error is about ±0.1. Albeit not greatly resolved, the range between $\omega_{max} < \omega < 10$ cpd falls-off more steeply roughly at $p = -2$ and the range between $10 < \omega < 100$ cpd shows a reduced variance that may partially be characterized by intermittency ($p = -1$; Schuster, 1984), but which is not yet explained. Here, it is observed to bridge between $p = -2$ and super-IGW BO-scaling $p = -7/5$. This would be further observation of a marginally ocean-state to the -1-scaling in KE-spectra (present Fig. 1 and van Haren et al., 2002) and in the continuum of the band [f N] in open-ocean T-spectra (van Haren and Gostiaux, 2009).

About 140 m above the seafloor, a less precise older-type T-sensor demonstrates $p = -7/5$ between a reduced range of about $10 < \omega < 1000$ cpd, with a suggestion for $p = -5/3$ around 10 cpd. This indicates convection can still dominate over shear extending O(100 m) above the seafloor, as has been shown in more detail for certain periods (van Haren, 2023c).

Whilst more extended work with longer data sets and more T-sensors is to be done, the extended continuous spectral slope from these high-resolution temperature observations suggests a direct coupling between sub-mesoscale motions, IGW motions, comprising

internal gravity and gyroscopic waves, and convection turbulence. The temperature spectra also show consistency with the limited KE-spectra of Fig. 1 from roughly the same area, and both indicate a dominance of non-isotropic, stratified-turbulence convection between sub-mesoscales and largest turbulent overturning scales in extended BO-scaling suggesting cross-spectral coupling. The discrepancy with KE-spectra in laboratory experiments of Pawar and Arakeri (2016) may be due to the difference of settings. In a non-zero-mean flow turbulence convection experiment near the gas-liquid critical point, BO-scaling was observed for both KE and temperature (Ashkenazi and Steinberg, 1999). We recall that our deep-sea conditions are non-zero-mean flow, weak tides, very high bulk Reynolds numbers $O(10^5)$ given the large scales, varying non-zero vertical density stratification, and our example spectra did not clearly resolve KO-scaling.

This 18-day T-sensor data set demonstrates dominant deviations from inertial subrange over several orders of magnitude of frequency range. The mesoscale-IGW-turbulence motions transport and locally mix warm waters with cooler surroundings outside a period of buoyancy-driven dense-water formation, which is thought to bring cooler waters downward during short periods of time.

5 How robust is the system of ocean circulation and stratification?

Any variation to the nonlinear system of ocean circulation may encounter several complex feedback mechanisms, of which the effects are not yet fully understood for the present-day ocean. Although stable density stratification hampers vertical exchange by turbulent mixing, it does not block it. While stratification supports internal waves and their destabilizing shear, turbulent mixing during particular phase of a wave may decrease or destroy it locally in time and space. However, a subsequent internal wave-phase will restratify the mixed patch, thereby maintaining its own support of stable stratification. Such a feedback system may be at work, for example when the ocean absorbs more heat.

Increased sea-surface temperature may lead to increased vertical density stratification, which may lead to less turbulent exchange as vertical overturning is suppressed.

However, it will also lead to more internal waves through the extension of their spectral band to higher frequencies, with the potential to increased interaction, non-linearity, and turbulence-generating wave breaking. As particular internal waves can propagate deep into the ocean interior away from their source, they can cause enhanced turbulent mixing elsewhere (e.g., Alford, 2003).

Limited observations have thus far not provided evidence for an inverse correspondence between changes in turbulent mixing and changes in temperature across the near-surface photic zone along a longitudinal section of the Northeast Atlantic Ocean (van Haren at al., 2021b). This lack of correspondence suggests a feedback mechanism at work mediating potential physical environment changes so that global warming may not affect vertical turbulent fluxes of heat, and thereby also of, e.g., carbon.

One such feedback mechanism may be convection-turbulence induced by internal waves and sub-mesoscale eddies. Renewed analysis of yearlong moored current meter data from the Irminger Sea, North-Atlantic Ocean, demonstrate a significant p = -11/5 spectral slope at sub- and at super-inertial frequencies (Fig. 3). As was outlined in van Haren (2007), the area showed an IGW-band (1), for N = f, with dominant sub-inertial shear at small 8-m vertical scales despite the dominant internal tidal KE. The correspondence with the Mediterranean data of Fig. 1 is striking, including the one order of magnitude change in KE between sub- and super-IGW p = -11/5-slopes with similar p = -1 bridge albeit uncertain crossing level, and similar heights of near-inertial peak despite the tidal peak in Fig. 3.

While few ocean observations have been presented of BO-scaling thus far in comparison with KO-scaling, perhaps also because of the lack of precision of standard oceanographic instrumentation, coupling has not been established between convection and stratified small-scale turbulence with mesoscale motions. Likewise, complicating factors are spectral interruption by internal waves. However, internal wave trapping by mesoscale eddies has been well described (e.g., Kunze, 1985; Voet et al., 2024), and thus provides an obvious coupling between these motions. It is expected that such coupling may lead to strong nonlinearity of the internal waves that leads to turbulent mixing produced by wave breaking.

Although such turbulent mixing is smaller than that induced by internal wave breaking above sloping topography, such coupling may be an important factor in downward transport of near-inertial energy that eventually breaks elsewhere, e.g., over topography.

As demonstrated using Mediterranean observations, not only convectively unstable cooler and/or saltier waters potentially lead to downward motions from the surface. Also, sub-mesoscale eddies and near-inertial waves can push stratified waters to the deep sea. Such a downward push can be fast to transport materials from surface to 2500-m deep seafloor in a day (van Haren et al., 2006), and which is of the same order of magnitude as attributed to dense-water convection (Schott et al., 1996). It can also be more turbulent compared to shear-induced motions in the stratified ocean-interior, whereby turbulence reaches the seafloor according to few observations from the abyssal Pacific (van Haren, 2020) and alpine freshwater Lake Garda (van Haren and Dijkstra, 2021). Further extended observational evidence is urgently needed, preferably resolving much larger scales.

Although the anthropogenic influence on the Earth's climate is without doubt, the impact on ocean circulation is not fully known because we lack sufficient, notably observational, information of the relevant processes that thus cannot be properly modeled yet. Therefore, we should be cautious in making predictions (e.g., Ditlevsen and Ditlevsen, 2023; van Westen et al., 2024) on future ocean circulation based on single parameters like ocean-surface temperature or fresh-water flux that are uncertain proxies. Because no observational (van Haren et al., 2021b), modeling (Little et al., 2020) or paleo-proxy validation (Cisneros et al., 2019) physics evidence exists that sea-surface temperature is a solid estimator of AMOC-strength variations, other properties like vertical density gradients (stratification), and turbulence intensity may be considered. Small-scale physical processes such as, for example, transport by sub-mesoscale eddies and turbulence-generating breaking of internal waves that are not incorporated in these models will alter such parameters, and thereby statistical analyses. This may lead to feedback mechanisms on property gradients such as density stratification so that large-scale ocean circulations like the AMOC may not collapse.

Variability of the ocean in space and time is a key to its dynamics, but it is unclear how robust such variations can be, e.g., whether shifting sites for deep dense-water formation (Gou et al., 2024) may be part of the same system. Observational evidence verifying numerical simulations' outcome, not only predictions but also present-day, of ocean-state is needed. Observations are also required to demonstrate variability in relevant physics processes for model-implementation. Besides eddies and coupling with atmosphere (e.g., Gent, 2018), numerical models of complex nonlinear ocean circulation should contain internal-wave turbulence with appropriate space and time dependency. The importance of internal wave breaking leading to boundary mixing above sloping topography in general ocean circulation models has been acknowledged in various ways (Scott and Marotzke, 2002; Ferrari et al., 2016).

As for the ocean circulation in the horizontal plane near its surface with most impact on mankind, wind will remain the main driver. As long as the Earth rotation does not alter direction, wind will maintain its general course (Wunsch, 2004). The atmosphere remains the key player in the global heat transport across mid-latitudes rather than the ocean. Simultaneously, the importance of processes like stratification and turbulent mixing induced by, e.g., internal wave breaking with or without sub-mesoscale coupling cannot be underestimated for life near the ocean-surface as well as in the -deep, because it will come to a halt without such processes.

Data availability. No new data were created or analyzed in this study: replot and re-analysis of data presented in van Haren and Millot (2003), van Haren (2007) and van Haren (2023a).

Competing interests. The author declares that he has no conflict of interest.

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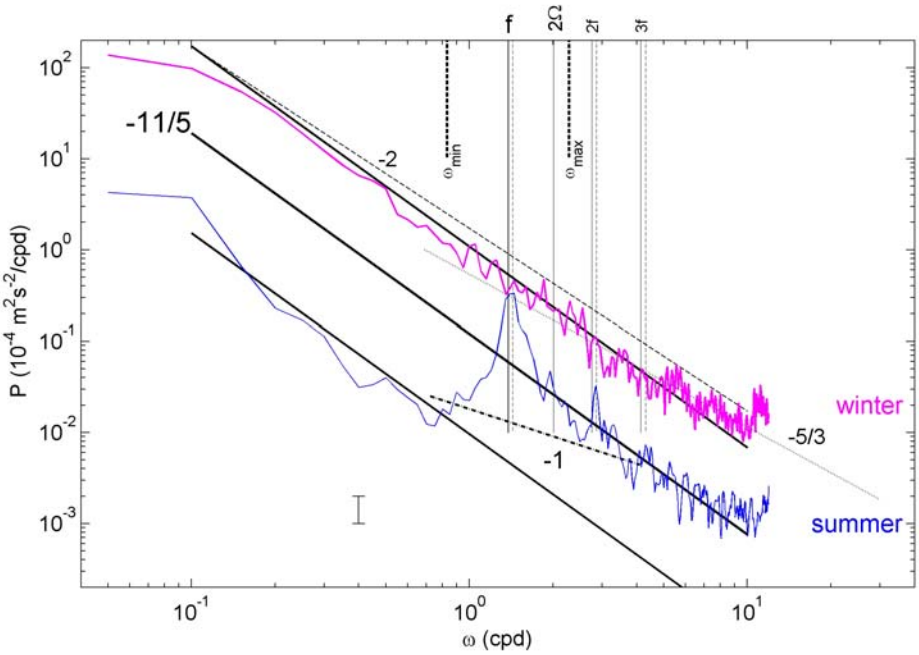


Fig. 1. Moderately smoothed (20 degrees of freedom, dof) kinetic energy (KE) spectra over 100 days of data from 3600-s sampled Aanderaa mechanical current meter moored in 1981/1982 at z = -1100 m well above the continental slope in the Ligurian Sea at 43° 28.32′ N, 7° 46.10′ E, 2250 m water depth. For details on these data, see van Haren and Millot (2003). The two spectra are not offset deliberately from each other; 'noise' also contains other signals near the Nyquist frequency. The 'summer' spectrum (blue) is an average from data between days 190 and 290 (in 1981), the 'winter' (magenta) between days 375 and 475 (adding +365 for days in 1982). Several frequencies are indicated including inertial frequency f, Earth rotational Ω and inertio-gravity wave bounds [$\omega_{min} \leq f$, $\omega_{max} \geq N, 2\Omega$] for buoyancy frequency N = f. The dashed lines indicate harmonics of 1.04f. Four spectral slopes ω^p are indicated by their exponent: p = -11/5 (solid slope in the log-log plot) for Bolgiano-Obukhov 'BO' scaling reflecting the buoyancy subrange of convection-turbulence (e.g., Pawar and Arakeri, 2016), p = -5/3 (dotted slope) for Kolmogorov-Obukhov 'KO' scaling reflecting the equilibrium inertial subrange for dominant shear-induced turbulence (Kolmogorov 1941; Obukhov, 1949), p = -1 (dash-dotted slope) for intermittency of self-organized criticality (Schuster, 1984; Bak et al., 1987) and p = -2 (dashed slope) for internal wave scaling (Garrett and Munk, 1972) or finestructure contamination (Phillips, 1971; Reid, 1971).

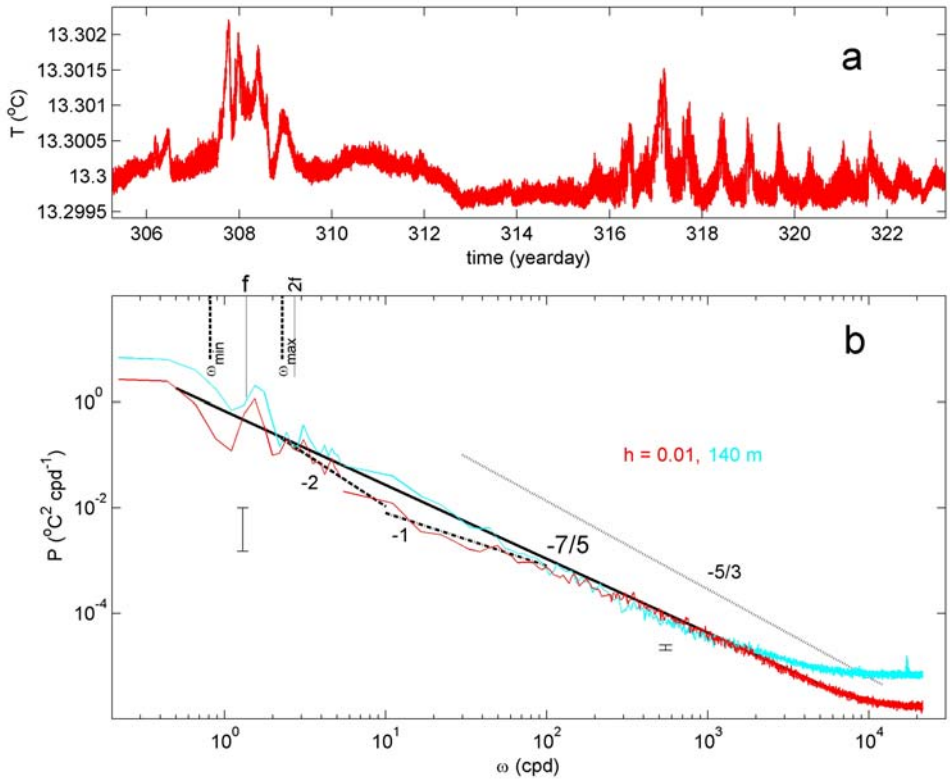
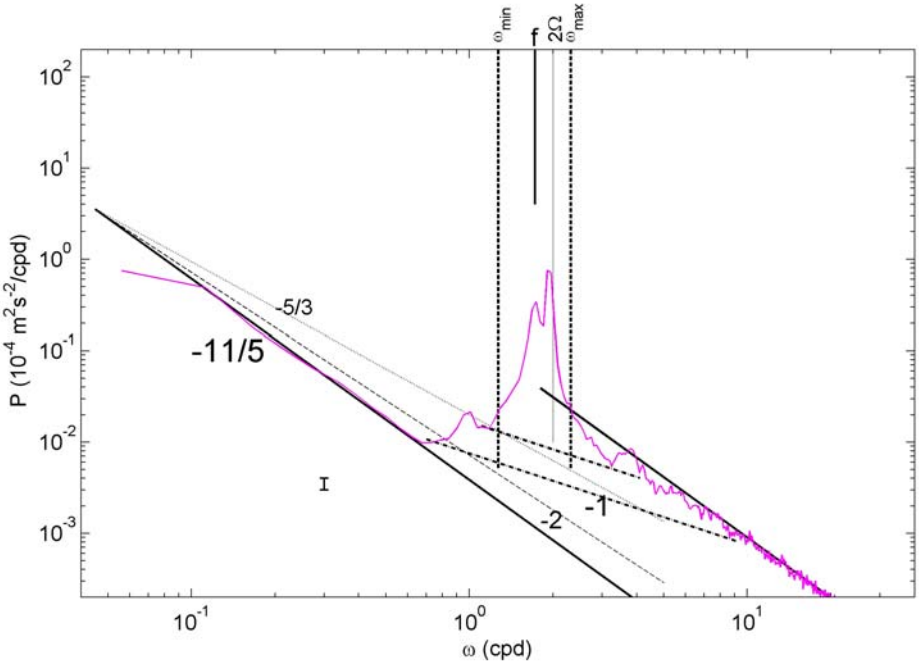


Fig. 2. Eighteen days of high-resolution 2-s sampled temperature 'T' data from a NIOZ T-sensor fallen off a mooring-line in 2020 and lying 0.01 m above a flat seafloor about 10 km south of the foot of the continental slope at 42° 49.50′ N, 6° 11.78′ E, 2458 m water depth, about 100 km WSW from the site in Fig. 1. For details on these data see van Haren (2023a). (a) Time series of 18 days of raw temperature data. (b) Temperature variance spectrum that is stitched together using two spectra with different smoothing. Weakly smoothed (10 dof; ω<5 cpd) and heavily smoothed (250 dof; ω>5 cpd) spectra of data in a., with bars showing the respective 95% confidence limits. For comparison, a spectrum is shown in cyan from data of a less precise T-sensor at a drag-parachute line stuck at 140 m above the seafloor. Frequency and spectral slope indications are as in Fig. 1, while -7/5 (solid slope) indicates BO-scaling of an active scalar (e.g., Pawar and Arakeri, 2016). Note the different axes-ranges compared with Fig. 1.



711
712 **Fig. 3**. Like Fig. 1 with the same axes ranges, but for strongly smoothed (50 dof) KE
713 spectra averaged over 400 days of data from 600-s sampled Valeport mechanical current
714 meter moored at z = -1000 m over the Mid-Atlantic Ridge at 58° 59.67′ N, 33° 56.12′ W,
715 2540 m water depth in 2003/2004, within the project discussed in van Haren (2007). The
716 small bar shows the 95% confidence interval.
717