



1 **Theoretical Framework for Measuring Cloud Effective**  
2 **Supersaturation Fluctuations with an Advanced Optical System**

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31 **Abstract**

32           Supersaturation is crucial in cloud physics, determining aerosol activation and influencing cloud  
33 droplet size distributions, yet its measurement remains challenging and poorly constrained. This study  
34 proposes a theoretical framework to simultaneously observe critical activation diameter and  
35 hygroscopicity of activated aerosols through direct measurements of scattering and water induced  
36 scattering enhancement of interstitial and activated aerosols, enabling effective supersaturation  
37 measurements. Advanced optical systems based on this framework allows minute- to second-level  
38 effective supersaturation measurements, capturing fluctuations vital to cloud microphysics. Although  
39 currently limited to clouds with supersaturations below  $\sim 0.2\%$  due to small scattering signals from  
40 sub-100 nm aerosols, advancements in optical sensors could extend its applicability. Its suitability for  
41 long-term measurements allows for climatological studies of fogs and mountain clouds. When  
42 equipped with aerial vehicles, the system could also measure aloft clouds. Therefore, the proposed  
43 theory serving a valuable way for both short-term and long-term cloud microphysics and aerosol-cloud  
44 interaction studies.

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## 59 1. Introduction

60 Clouds and fogs play critical roles in weather patterns and climate change, influencing both  
61 precipitation and the radiative balance of the Earth's atmosphere. As such, they are central to accurate  
62 weather and climate predictions. Despite their importance, representing clouds accurately in  
63 atmospheric models remains a significant challenge (Seinfeld et al., 2016). Supersaturation, defined  
64 as the difference between the actual water vapor pressure ( $e$ ) and the saturation vapor pressure ( $e_s$ )  
65 which is typically expressed as a dimensionless quantity  $(e - e_s)/e_s$ , is a key parameter that links  
66 aerosols to clouds through the process of aerosol activation, making it fundamental to cloud physics  
67 (Seinfeld and Pandis, 2016). Despite its importance, supersaturation is difficult to measure and remains  
68 poorly understood and constrained (Yang et al., 2019). Previous studies have highlighted that other  
69 than the mean supersaturation, supersaturation fluctuations also play critical roles in aerosol activation  
70 and cloud droplet growth, ultimately influencing the evolution of cloud droplet size distributions  
71 (Kaufman and Tanré, 1994;Sardina et al., 2015;Chandrakar et al., 2018;Chandrakar et al., 2020;Shaw  
72 et al., 2020). For instance, cloud chamber experiments have shown that supersaturation fluctuations  
73 promote aerosol activation and enhance aerosol activity (Shawon et al., 2021;Anderson et al., 2023),  
74 particularly when the magnitude of these fluctuations is comparable to the mean supersaturation  
75 (Prabhakaran et al., 2020). Both experimental and theoretical analyses suggest that supersaturation  
76 fluctuations can broaden cloud droplet size distributions (Chandrakar et al., 2016;Abade et al.,  
77 2018;Saito et al., 2019).

78 Supersaturation fluctuations arise not only from turbulent variations in the temperature and vapor  
79 pressure fields but also from the growth and evaporation of droplets, which drive mass and heat  
80 exchange between droplets and the surrounding air. Field measurements have demonstrated that the  
81 supersaturation is indeed a fluctuating quantity (Ditas et al., 2012;Siebert and Shaw, 2017). However,  
82 as noted by Shaw et al. (2020), measuring supersaturation remains a formidable challenge due to its  
83 extreme sensitivity to variations in  $e$  and temperature. Currently, cloud and fog supersaturation are  
84 typically retrieved from aerosol activation measurements (Ditas et al., 2012) or estimated from vertical  
85 velocity measurements (Siebert and Shaw, 2017). Direct measurements of water vapor pressure and  
86 temperature are generally used to estimate supersaturation fluctuations but do not provide precise  
87 direct supersaturation measurements (Ditas et al., 2012;Siebert and Shaw, 2017). Supersaturation  
88 parameterizations based on vertical velocity are common in models (Abdul-Razzak et al., 1998), while  
89 field measurements often rely on aerosol activation data to investigate supersaturation fluctuations and  
90 evolutions in clouds and fogs (Ditas et al., 2012;Hammer et al., 2014;Shen et al., 2018;Mazoyer et al.,



91 2019;Ziková et al., 2020;Wainwright et al., 2021;Kuang et al., 2024). In addition, supersaturations  
92 were also estimated using the closure between cloud droplet number and cloud condensation nuclei  
93 (CCN) measurements at various supersaturations (Yum et al., 1998;Sanchez et al., 2016;Sanchez et  
94 al., 2021;Saliba et al., 2023).

95 In summary, direct measurements of water vapor pressure and temperature are essential for  
96 quantifying supersaturations; however, they are nearly impossible with current technologies.  
97 Supersaturation measurements from aerosol and cloud microphysics monitoring often reflect an  
98 effective supersaturation that drives aerosol activation, which is indeed critical in cloud physics. The  
99 complexity of cloud formation and evolution and the central role of supersaturation in these processes  
100 underscore the need for precise measurement and representation of supersaturation. Advancements in  
101 measuring and understanding supersaturation are essential for improving the accuracy of models and  
102 reducing uncertainties in weather and climate predictions. In this study, we propose a theoretical  
103 framework for using optical methods to observe effective supersaturations based on aerosol activation  
104 in clouds and preliminarily validated utilizing data obtained from field campaigns. The feasibility of  
105 employing an advanced optical system to measure supersaturation fluctuations were also explored and  
106 discussed.

## 107 2. Methods and Materials

### 108 2.1 Observing effective supersaturations on the basis of $\kappa$ -Köhler theory

109 Fluctuations in supersaturation mean that the effective supersaturation, which directly affects  
110 aerosol activation, differs from the mean supersaturation. The  $\kappa$ -Köhler theory provides a framework  
111 for deriving effective supersaturation from aerosol activation measurements in clouds (Petters and  
112 Kreidenweis, 2007):

$$113 \quad S = \frac{D^3 - D_a^3}{D^3 - D_a^3(1 - \kappa)} \cdot \exp\left(\frac{4\sigma_{s/a} M_{water}}{R \cdot T \cdot D_p \cdot g \cdot \rho_w}\right) \quad (1)$$

114 where  $S$  is the saturation ratio,  $S - 1$  represents supersaturation,  $D$  is the droplet diameter,  $D_a$  is the dry  
115 diameter,  $\sigma_{s/a}$  is the surface tension of solution/air interface,  $T$  is the temperature,  $M_{water}$  is the  
116 molecular weight of water,  $R$  is the universal gas constant,  $\rho_w$  is the density of water, and  $\kappa$  is the  
117 hygroscopicity parameter. The  $\kappa$ -Köhler theory tells that if the critical diameter of aerosol activation  
118 ( $D_a$ ) and corresponding aerosol hygroscopicity parameter  $\kappa$  are known, the surrounding



119 supersaturation can be retrieved based on air temperature measurements and by assuming  $\sigma_{s/a}$  the  
120 surface tension of water (as shown in Fig.S1a).

121 However, the simultaneous measurements of  $D_a$  and  $\kappa$  of activated aerosols with diameters  
122 around  $D_a$  are indeed challenging. The direct measurements of size-resolved activation ratio (AR) in  
123 clouds are essential for  $D_a$  retrievals through the following equation:

$$124 \quad AR(D_p) = \frac{MAF}{2} \left( 1 + \operatorname{erf} \left( \frac{D_p - D_a}{\sqrt{2\pi}\sigma} \right) \right) \quad (2)$$

125 Where  $D_p$  is the particle diameter, MAF is the maximum activation fraction and  $D_a$  is critical  
126 activation diameter,  $\sigma$  is associated with the slope of the curve near  $D_a$ . This formula was previously  
127 proposed by Rose et al. (2008) to fit the AR measurements and widely used in AR parameterizations  
128 (Tao et al., 2018b). Therefore, it typically requires a unique inlet system and a suite of instruments that  
129 measure the aerosol size distribution of both interstitial and total aerosol populations (Hammer et al.,  
130 2014; Ziková et al., 2020). Consequently, this is rarely done, even in ground fog measurements. Instead,  
131  $D_a$  was usually estimated from aerosol measurements and fog droplet size distributions measurements  
132 which indirectly provides the number concentrations of activated aerosols therefore could be used in  
133 retrieving  $D_a$  through assuming that all aerosols larger than  $D_a$  are activated (Mazoyer et al.,  
134 2019; Wainwright et al., 2021; Shen et al., 2018) which brings uncertainty in  $D_a$  derivations due to that  
135 the maximum activation fraction of aerosols larger than  $D_a$  does not equal to unit although usually  
136 very close to (Tao et al., 2018b). For the effective supersaturation measured in aloft clouds, the aerosol  
137 number size distributions inside and outside the cloud as well as cloud droplet number concentrations  
138 were used by Ditas et al. (2012) to derive  $D_a$ , and other approaches were also used (Gong et al., 2023).  
139 The  $\kappa$  values were usually retrieved from size-resolved cloud condensation nuclei measurements under  
140 certain supersaturations (Hammer et al., 2014; Mazoyer et al., 2019) or from growth factor  
141 measurements (Wainwright et al., 2021) or sometime assumed due to the lack of measurements. The  
142  $\kappa$  of activated aerosols were not directly measured in these studies due to the difficulty of the direct  
143 sampling of activated aerosols as well as subsequent hygroscopicity measurements.

## 144 2.2 Field measurements

145 Kuang et al. [2024] developed an advanced aerosol-cloud sampling system designed to measure  
146 fog and cloud activation processes. This compact, integrated system can automatically switch between  
147 different inlets, including PM<sub>1</sub> (particles and droplets with an aerodynamic diameter < 1  $\mu\text{m}$ ), PM<sub>2.5</sub>



148 (particles and droplets with an aerodynamic diameter  $< 2.5 \mu\text{m}$ ) impactor, and Total Suspended  
149 Particles (TSP, encompassing all particles and droplets) (as shown in Fig. S2). When combined with  
150 instruments that measure aerosol physical, optical, and chemical properties, this system is well-suited  
151 for investigating cloud microphysics and chemistry. It was utilized in the AQ-SOFAR campaign,  
152 dedicated to studying AQueous Secondary aerOsol formation in Fogs and Aerosols and their Radiative  
153 effects in the North China Plain (Kuang et al., 2024).

154 During this campaign, several radiation fog events were observed, enabling the measurement of  
155 size-resolved AR curves, aerosol hygroscopicity as well as chemical compositions of interstitial and  
156 activated aerosols within fogs. These measurements provided insights into the evolution of  
157 supersaturations (Kuang et al., 2024). Notably, aerosol hygroscopicity was determined using a  
158 humidified nephelometer system, located downstream of the inlet system. This system measured  
159 multiwavelength scattering coefficients (450 nm, 525 nm, 635 nm) under both nearly dry ( $\text{RH}<20\%$ )  
160 and humid conditions ( $\text{RH}\sim 84\%$ ), offering aerosol hygroscopicity data based on the optical theory  
161 proposed by Kuang et al. (2017). The size-resolved AR curves and aerosol chemical compositions  
162 were obtained through the aerosol size distribution and the aerosol mass spectrometry measurements  
163 downstream of the inlet system. A schematic of the inlet system and associated instruments is provided  
164 in Fig. S1. Further details about the entire experimental setup, size-resolved AR calculations as well  
165 as data analysis about mass spectrometer measurements can be found in Kuang et al. (2024).

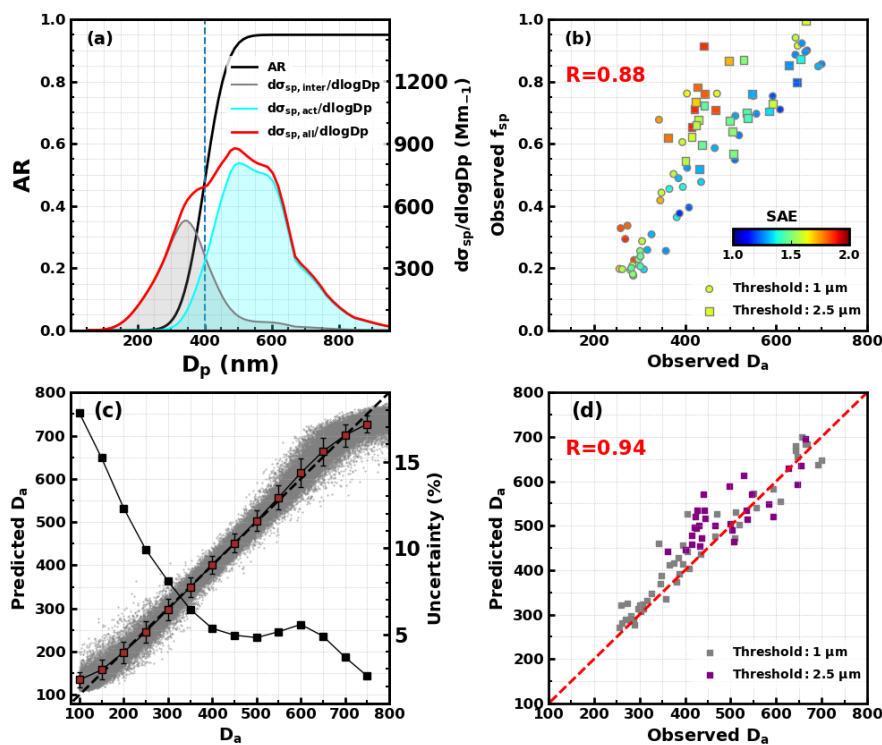
### 166 **3. Theoretical Framework and Concept Design of the Advanced Optical System**

#### 167 **3.1 Theory of Observing Critical Activation Diameter Using Scattering Measurements**

168 The typical shape of size-resolved AR curves observed in atmospheric fogs and clouds is  
169 illustrated in Fig. 1a (Ditas et al., 2012; Hammer et al., 2014; Zíková et al., 2020; Wainwright et al.,  
170 2021; Kuang et al., 2024). In clouds, aerosols can be classified as either activated aerosols, which form  
171 cloud droplets, or inactivated aerosols, which remain as interstitial aerosols. The critical diameter that  
172 distinguishes interstitial aerosols from cloud or fog droplets varies depending on the supersaturation  
173 (Kuang et al., 2024). A diameter of  $2.5 \mu\text{m}$  is typically suitable for surface fogs with relatively lower  
174 supersaturations ( $<0.1\%$ ), while  $1 \mu\text{m}$  is more appropriate for aloft clouds with higher supersaturations  
175 ( $>0.1\%$ ) (Mazoyer et al., 2019; Kuang et al., 2024; Lu et al., 2020). The typical AR curve shows that  
176 most aerosols larger than  $D_a$  are activated, while most smaller aerosols remain inactivated. As a result,  
177 the scattering properties, such as size-resolved scattering coefficients (Fig. 1a), the scattering Ångström



178 exponent (SAE) and its wavelength dependence, which are directly related to aerosol size distribution,  
 179 differ significantly between interstitial and activated aerosols.



**Figure 1.** (a) The typical shape of size-resolved aerosol activation ratio (AR) curve produced using the function of Eq.2, with the  $D_a$  of 400 nm, the MAF of 0.95 and the  $\sigma$  of 30, and the average PNSD observed in the North China Plain from six campaigns (Kuang et al., 2018) and the AR curve was used to simulate an example of the size-resolved aerosol scattering ( $\sigma_{sp}$ ) distributions of interstitial and activated aerosols at 525 nm; (b) Relations between observed  $D_a$  and  $f_{sp}$  during the AQ-SOFAR campaign using 1 and 2.5  $\mu\text{m}$  as the threshold of interstitial aerosols, with the scatter points are colored with corresponding SAE of total dry state  $\text{PM}_{10}$  aerosols. (c) Comparisons of all prescribed  $D_a$  and predicted  $D_a$  values represented by scatter points, they are further binned with interval of 50 nm, averages and standard deviations represented by purple squares and their error bars, black squares represent relative uncertainty of the right axis at each bin; (d) The comparisons of  $D_a$  retrieved using activation ratio observations and those predicted using scattering observations as inputs of the trained model, dashed lines represent 1:1 lines.

180 If we focus on  $\text{PM}_{10}$  of the total dry aerosol population (the reasoning for this is discussed in Sect.  
 181 S1 of the supplement), the scattering fraction of interstitial aerosols in the total dry  $\text{PM}_{10}$  population,  
 182 defined as  $f_{sp} = \sigma_{sp, \text{PM}_{10}, \text{inter}}(\text{dry}, 525 \text{ nm}) / \sigma_{sp, \text{PM}_{10}, \text{all}}(\text{dry}, 525 \text{ nm})$ , where  
 183  $\sigma_{sp, \text{PM}_{10}, \text{inter}}(\text{dry}, 525 \text{ nm})$  is the scattering coefficient of  $\text{PM}_{10}$  interstitial aerosols in a dry state at a  
 184 wavelength of 525 nm, and  $\sigma_{sp, \text{PM}_{10}, \text{all}}(\text{dry}, 525 \text{ nm})$  is that of all  $\text{PM}_{10}$  aerosols, is likely to be highly



185 correlated with  $D_a$ . Generally, the larger the  $D_a$ , the higher the  $f_{sp}$ . This relationship was directly  
186 confirmed using  $D_a$  and the scattering properties of dry  $PM_1$  interstitial and total aerosols during the  
187 AQ-SOFAR campaign, as shown in Fig. 1b, that observed  $D_a$  correlates highly with observed  $f_{sp}$   
188 ( $R=0.88$ ). However, at a given  $D_a$ ,  $f_{sp}$  can vary significantly, and these variations are closely related  
189 to the SAE of all dry  $PM_1$  aerosols, which are mainly determined by aerosol size distribution. In fact,  
190 aside from the size distribution of the total aerosol population that determines SAE, the shape of the  
191 AR curve also plays a significant role in the variations of  $f_{sp}$ .

192 The nephelometer measures the aerosol scattering coefficient at three wavelengths, enabling  
193 direct measurements of the SAE for both the total dry-state  $PM_1$  aerosols and the interstitial aerosols.  
194 Therefore, the relationship between  $f_{sp}$  and  $D_a$  can be further constrained by the SAE of interstitial  
195 and activated aerosols, as well as their wavelength dependence. This implies that a simple formulaic  
196 relationship between  $f_{sp}$  and  $D_a$  may not exist. However, the six scattering parameters  
197  $\sigma_{sp,PM_1,inter}(dry, \lambda)$ —at 450 nm, 525 nm, 635 nm, and  $\sigma_{sp,PM_1,all}(dry, \lambda)$  at 450 nm, 525 nm, 635  
198 nm—contain both the  $f_{sp}$  information and the SAE characteristics of both aerosol groups, thus  
199 potentially be used to accurately retrieve  $D_a$ . Machine learning techniques, which are well-suited for  
200 handling complex relationships, can be applied to this problem.

201 This assumption was tested using Mie theory, based on aerosol size distributions sampled during  
202 six campaigns conducted in the North China Plain region (Kuang et al., 2018). For each aerosol size  
203 distribution, we randomly assumed different activation curves using Eq.2. The details of this  
204 simulation are provided in Sect. S2 of the supplement. The simulation pairs of these six scattering  
205 parameters and  $D_a$  were used to train a random forest model (Kuang et al., 2018). To preliminarily  
206 validate this approach, we randomly selected 75% of the simulated data points for training the model,  
207 while the remaining 25% were used for validation. The results, shown in Fig. 1c, indicate that this  
208 approach could retrieve  $D_a$  with an uncertainty of less than 10% for  $D_a$  larger than 250 nm, and even  
209 as low as ~6% for  $D_a$  larger than 350 nm. However, the uncertainty increases as  $D_a$  decreases,  
210 particularly for  $D_a$  smaller than 250 nm. The larger uncertainty at smaller  $D_a$  is due to the fact that  
211 aerosols smaller than 250 nm typically contribute less than 10% to total scattering in the dry state,  
212 making  $f_{sp}$  less sensitive to variations in  $D_a$ . This issue becomes more pronounced when  $D_a$  is less  
213 than 100 nm, as aerosols smaller than 150 nm generally contribute negligibly to total aerosol scattering  
214 [Kuang et al., 2018]. This method was further validated using observations from the AQ-SOFAR  
215 campaign. In this validation,  $D_a$  values were first predicted using aerosol scattering observations with





216 the trained model and then compared with  $D_a$  values retrieved from size-resolved AR measurements,  
217 as shown in Fig. 1d. It should be noted that the impactor operates in a sequence of PM<sub>1</sub>, PM<sub>2.5</sub>, TSP,  
218 and then back to PM<sub>1</sub>, with the flow alternating between a thermodenuder and bypass every 10 minutes  
219 for each inlet. To calculate size-resolved AR curves, we assumed that aerosol populations remained  
220 unchanged during the 30-minute period (based on comparisons between PM<sub>1</sub>/PM<sub>2.5</sub> and TSP inlets),  
221 which can sometimes introduce significant uncertainties in the size-resolved AR calculations. When  
222 using PM<sub>2.5</sub> as the threshold, the much lower number concentrations of aerosols larger than 400 nm  
223 can introduce more uncertainty in  $D_a$  retrievals, partially explaining the lower performance in Fig. 1d  
224 when using the PM<sub>2.5</sub> threshold.

### 225 3.2 Method of observing Hygroscopicity of Activated Aerosols

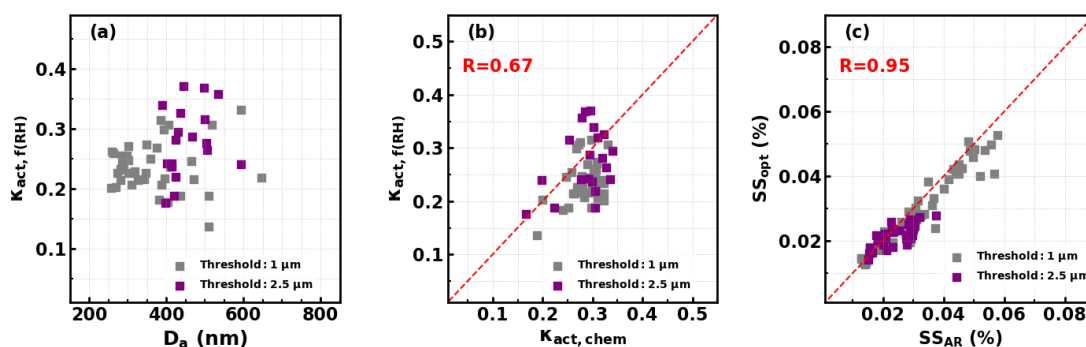
226 Measuring the hygroscopicity  $\kappa$  of activated aerosols at the critical activation diameter  $D_a$  under  
227 varying supersaturations is challenging, not only due to technical limitations but also because of the  
228 inherent variability in  $D_a$ . Kuang et al. (2017) introduced a novel optical method for observing aerosol  
229 hygroscopicity by using the aerosol light scattering enhancement factor  $f(RH)$  that associated with  
230 aerosol hygroscopic growth. This method is particularly suitable for the objectives outlined here. The  
231 method requires SAE and light scattering enhancement factors  $f(RH)$  of activated aerosols as inputs,  
232 and retrieved  $\kappa$  can be termed as  $\kappa_{act,f(RH)}$  which represents the overall hygroscopicity of activated  
233 aerosols and can be understood as the average  $\kappa$  of activated aerosols with the scattering contribution  
234 of each aerosol particle as the weight (Kuang et al., 2020). The scattering coefficients of activated  
235 aerosols at multiwavelength can be calculated as  $\sigma_{sp,PM_1,act}(dry, \lambda) = \sigma_{sp,PM_1,all}(dry, \lambda) -$   
236  $\sigma_{sp,PM_1,inter}(dry, \lambda)$ , therefore corresponding SAE can be obtained. The  $f(RH)$  of activated aerosols  
237 at 525 nm can be calculated as the following:

$$238 \quad f(RH)_{act} = \frac{\sigma_{sp,PM_1,all}(RH, 525 \text{ nm}) - \sigma_{sp,PM_1,inter}(RH, 525 \text{ nm})}{\sigma_{sp,PM_1,all}(dry, 525 \text{ nm}) - \sigma_{sp,PM_1,inter}(dry, 525 \text{ nm})}$$

239 During the AQ-SOFAR campaign, a humidified nephelometer system consisting of two  
240 nephelometers—one measuring aerosol scattering in the dry state and the other at a fixed RH of 84%—  
241 was placed downstream of the PM<sub>1</sub> impactor. This setup allows for the humidification of dry-state  
242 interstitial aerosols and total aerosol populations to a high RH (e.g., above 80%), facilitating the  
243 required measurements, therefore severs one choice. The Retrieved  $\kappa_{act,f(RH)}$  under different  $D_a$   
244 conditions are shown in Fig.2a, demonstrating significant variations in  $\kappa_{act,f(RH)}$  and its variations



245 need to constrained. Also, the derived  $\kappa_{act,f(RH)}$  are compared those estimated from aerosol chemical  
246 compositions measurements ( $\kappa_{act,chem}$ , details about calculation methods can refer to Kuang et al.  
247 (2020)), as shown in Fig.2b and in general agree. Note that the mass spectrometer could not identify  
248 all aerosol components, and assumptions about the mixing rule as well as densities of components  
249 would bring uncertainties (Kuang et al., 2021). The comparisons between effective supersaturations  
250 derived from size-resolved AR measurements as well as  $\kappa_{act,chem}$  and from the optical method are  
251 shown in Fig.2c. On average, 0.002% of SS bias are observed due to the bias of  $D_a$  which associated  
252 more with assumptions made in  $D_a$  retrievals as previously discussed. As demonstrated by Kuang et  
253 al. (2024), for the fog case in the campaign, the threshold of 2.5  $\mu\text{m}$  should be used, however, does not  
254 affect the comparisons here.

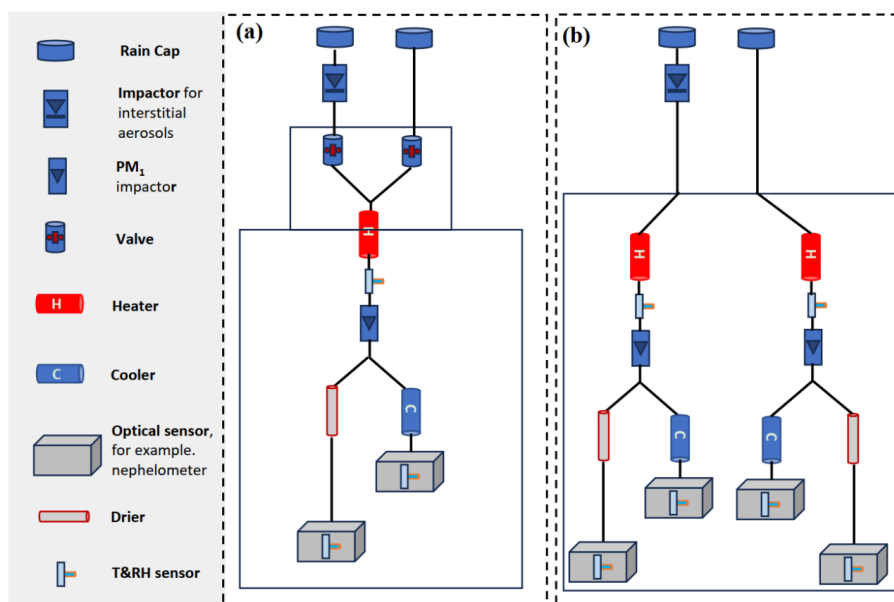


**Figure 2.** (a) Retrieved  $\kappa_{act,f(RH)}$  under different  $D_a$  conditions; (b) Comparison between  $\kappa$  of activated aerosols retrieved from the optical method ( $\kappa_{act,f(RH)}$ ) and estimated from aerosol chemical composition measurements ( $\kappa_{act,chem}$ ); (c) Comparisons between effective supersaturations (SSs) derived from size-resolved AR measurements as well as  $\kappa_{act,chem}$  ( $SS_{AR}$ ) and from the optical measurements ( $SS_{opt}$ ). Dashed red lines represent 1:1 lines.

255 Qiao et al. (2024) developed an advanced outdoor nephelometer system that measures aerosol  
256 dry scattering coefficients and scattering coefficients at nearly ambient RH without the need for  
257 humidifying the sample air by placing the entire nephelometer system in ambient air, with the  
258 instruments protected by a specially designed enclosure. This innovative design offers new insights  
259 into the hygroscopicity measurements of activated aerosols. Under cloud conditions, where the  
260 ambient RH is close to 100%, aerosol scattering under subsaturated conditions can be measured  
261 directly by applying heater.



262 **3.3 Concept Design of the Advanced Optical System for Measuring Effective**  
263 **Supersaturations**



**Figure 3.** Concept design of the advanced optical system with different number of optical sensors, (a) using two nephelometers or other optical sensors; (b) using four nephelometers or other optical sensors.

264 Based on the proposed optical methods for measuring  $D_a$  and  $\kappa_{act,f(RH)}$ , a conceptual design for  
265 outdoor instruments capable of measuring effective supersaturation with relatively high time  
266 resolution can be envisioned, as shown in Fig. 3a. The aerosol-cloud sampling system includes two  
267 inlets: one equipped with a PM<sub>1</sub> or PM<sub>2.5</sub> impactor (depending on cloud type) to sample interstitial  
268 aerosols, and another with a TSP inlet to sample both interstitial aerosols and cloud/fog droplets. A  
269 PM<sub>1</sub> impactor is placed downstream of the inlet system, where the RH of the sample air is reduced to  
270 70% (as discussed in Sect. S1 of the manuscript) through heater. Downstream of the PM<sub>1</sub> impactor,  
271 the sample flow is split into two streams: one is further dried to an RH below 10% before aerosol  
272 scattering coefficients are measured by the “dry” nephelometer, and the other is passed through an  
273 intelligent cooler to ensure the sample RH in the “wet” nephelometer remains close to 90%. The  
274 sample air is automatically switched between the interstitial inlet and the TSP inlet at set intervals,  
275 such as one minute for each inlet, enabling minute-level measurements of effective supersaturations.  
276 While the nephelometer can output scattering measurements every second, reliable data can only be  
277 achieved at intervals of around 30 seconds (exact values can be determined through future testing) due  
278 to the residence time of aerosols in the nephelometer and potential light source instability. If four



279 nephelometers are available, a more advanced optical system can be designed (Fig. 3b) that does not  
280 require switching between the interstitial inlet and the TSP inlet. Instead, two nephelometers would be  
281 placed downstream of the interstitial inlet and two downstream of the TSP inlet, enabling higher time  
282 resolution effective supersaturation measurements. Other types of optical instruments exist that can  
283 achieve stable second-level aerosol scattering or extinction measurements with a stable laser light  
284 source [Moise et al., 2015; Zhou et al., 2020]. Therefore, with the development of suitable optical  
285 instruments, it may be possible to achieve second-level effective supersaturation measurements.

#### 286 4. Discussions on Limitations and Advantages

287 The proposed theoretical framework enables simultaneous measurements of  $D_a$  and  $\kappa$  for  
288 activated aerosols, leveraging the high time resolution of optical instruments to potentially provide  
289 second-level measurements of supersaturation. However, several limitations should be discussed and  
290 might be improved upon: **(1) Shape of size-resolved AR curve:** Cloud chamber studies have shown  
291 that supersaturation fluctuations can lead to the coexistence of particles with the same critical  
292 supersaturation as both interstitial aerosols and cloud droplets (Shawon et al., 2021). This results in  
293 size-resolved AR curves deviate more from stepwise shape, a phenomenon also observed in some field  
294 measurements (Henning et al., 2004; Mertes et al., 2007). Despite this, a critical diameter  $D_a$  still exists,  
295 and such non-ideal curves can be treated a high standard deviation  $\sigma$  in the activation error function  
296 (Eq. 2), which does not fundamentally undermine the proposed framework, however, should be further  
297 checked for different cloud types. **(2) Measurement of  $\kappa$ :** Although the framework measures the  
298 overall  $\kappa$  of activated aerosols, the  $\kappa$  needed for supersaturation calculations is that of aerosols near  
299  $D_a$  ( $\kappa_{D_a}$ ). For  $D_a > \sim 200$  nm, the derived  $\kappa_{act,f(RH)}$  can provide a first-order estimate of  $\kappa_{D_a}$ , based on  
300 observed size-dependent characteristics of  $\kappa$  values (Liu et al., 2014; Shen et al., 2021; Wang et al.,  
301 2024), though more comprehensive evaluations are needed. Additionally,  $\kappa$  measured under  
302 subsaturated conditions differs from that under supersaturated conditions (Tao et al., 2023) might also  
303 bring some uncertainties. However, as shown in Fig. S1b, a bias of 0.1 in  $\kappa$  only results in a 0.01%  
304 bias in supersaturation retrievals, making the first-order estimates of  $\kappa_{D_a}$  from optical measurements  
305 suitable for supersaturation observations. **(3) Limitations in  $D_a$  Retrievals:** Current techniques using  
306 aerosol scattering measurements at visible wavelengths (e.g., nephelometers) are reliable only for  
307  $D_a > 100$  nm as shown in Fig. 1a, limiting effective supersaturation measurements to less than 0.21%  
308 (assuming a typical  $\kappa$  of 0.3). This restriction makes the technique most applicable to fog and stratus  
309 or stratocumulus cloud measurements. However, incorporating scattering measurements at ultraviolet



310 wavelengths could improve sensitivity to smaller  $D_a$  and lower  $\kappa$ , enabling measurements in  
311 conditions with higher effective supersaturation and a broader range of cloud types in the future.

312         Despite these limitations, the proposed theoretical framework represents the first system capable  
313 of directly providing high time resolution measurements of effective supersaturations using a single  
314 instrument. This system is particularly well-suited for surface fog and mountain cloud observations,  
315 and when coupled with aerial vehicles, it could also be employed for measurements in aloft clouds.  
316 The system offers several advantages for cloud and fog measurements: **(1) High-Resolution**  
317 **Supersaturation Measurements:** The system can provide measurements of effective supersaturations  
318 at even a second-level resolution, making it feasible for observing effective supersaturation  
319 fluctuations and supporting investigations into fog and cloud evolution mechanisms. **(2) Long-Term**  
320 **Measurement Capability:** The optical measurements, such as those from the nephelometer system,  
321 are well-suited for long-term observations, making it possible to acquire climatological data on the  
322 variability of fogs and mountain clouds. **(3) Comprehensive Aerosol and Cloud Data:** In addition to  
323 measuring effective supersaturations, the system directly captures the scattering and hygroscopic  
324 properties of both interstitial and activated aerosols. With further algorithm development, it could also  
325 retrieve the number concentrations of available cloud condensation nuclei (CCN) at certain  
326 supersaturations, as well as cloud droplet number concentrations, based on previous studies that have  
327 observed CCN using optical methods (Tao et al., 2018a). **(4) Monitoring Aerosol Hygroscopic**  
328 **Behavior:** The system continuously monitors aerosol hygroscopic behavior under subsaturated  
329 conditions along with the corresponding optical properties. This allows for clear documentation of the  
330 formation and dissipation of fog/cloud events, as well as the variation in aerosol optical and  
331 hygroscopic properties. Overall, the datasets generated by this system are well-suited for in-depth  
332 investigations of cloud physics and aerosol-cloud interactions. This system has the potential to  
333 significantly advance fundamental research on clouds and fogs. However, further theoretical studies  
334 are needed to refine and optimize this type of system.

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341

342 **Competing interests.** The authors declare that they have no conflict of interest.

343

344 **Data Availability.** All data presented in Figures of this manuscript are freely available at Kuang, Y.  
345 (2024), and more specific data will be made available on request.

346

347 **Author contribution.** YK conceived the theoretical framework and wrote the manuscript. JT, HX, L  
348 L, WX and WeX participated the field campaign and conducted measurements of aerosol chemical  
349 and physical properties. YS, PL, CZ reviewed and commented on the paper.

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