

# ClimKern v1.1.2: a new Python package and kernel repository for calculating radiative feedbacks

Tyler P. Janoski<sup>1-5</sup>, Ivan Mitevski<sup>6</sup>, Ryan J. Kramer<sup>7</sup>, Michael Previdi<sup>2</sup>, and Lorenzo M. Polvani<sup>1-2,8</sup>

<sup>1</sup>Dept. of Earth and Environmental Sciences, Columbia University, New York, NY, USA

<sup>2</sup>Lamont-Doherty Earth Observatory, Columbia University, Palisades, NY, USA

<sup>3</sup>NOAA Center for Earth System Science and Remote Sensing Technologies (CESSRST-II), New York, NY, USA

<sup>4</sup>City College of New York, New York, NY, USA

<sup>5</sup>NOAA National Severe Storms Laboratory, Norman, OK, USA

<sup>6</sup>Dept. of Geosciences, Princeton University, Princeton, NJ, USA

<sup>7</sup>NOAA Geophysical Fluid Dynamics Laboratory, Princeton, NJ, USA

<sup>8</sup>Dept. of Applied Physics and Mathematics, Columbia University, New York, NY, USA

**Correspondence:** Tyler P. Janoski (tjanoski@ccny.cuny.edu)

## Abstract.

Climate feedbacks are a significant source of uncertainty in future climate projections and need to be quantified accurately and robustly. The radiative kernel method is commonly used to efficiently compute individual climate feedbacks from climate model or reanalysis output. Despite its popularity, it suffers from complications, including difficult-to-locate radiative kernels, inconsistent kernel properties, and a lack of standardized assumptions in radiative feedback calculations, limiting the robustness and reproducibility of climate feedback computations. We designed the ClimKern project to address these issues with a kernel repository and a separate but complementary Python package of the same name. We selected eleven sets of radiative kernels and gave them a common nomenclature and data structure. The ClimKern Python package provides easy access to the kernel repository and functions to compute feedbacks, sometimes with a single line of code. ~~The ClimKern~~ functions contain helpful optional parameters while maintaining standard practices between calculations.

After documenting the kernels and ClimKern package, we test it with sample climate model output ~~from an abrupt 2×CO<sub>2</sub> experiment~~ to explore the sensitivity of feedback calculations to kernel choice. Interkernel spread ~~shows exhibits~~ considerable spatial heterogeneity, with the greatest spread in the ~~surface albedo and cloud feedbacks occurring in the~~ Arctic and ~~over the~~ Southern Ocean. ~~Considerable sensitivity to kernel choice is found even in the global means, with the surface albedo and cloud feedbacks showing the greatest spread across different kernels~~ ~~In the global mean, the Planck and surface albedo feedbacks show the greatest interkernel variability.~~ Our results highlight the importance of using ~~more than one radiative kernel~~ ~~multiple radiative kernels~~ and standardizing feedback calculations ~~, like those offered by ClimKern,~~ in climate feedback, climate-sensitivity, and polar amplification studies. As ClimKern continues to evolve, we hope others will contribute to its development to make it ~~even more useful to the~~ ~~an even greater tool for the radiative~~ feedback community.

## 20 1 Introduction

One of the fundamental questions in climate science is how much the surface will warm in response to the radiative forcing imposed by increasing CO<sub>2</sub> concentrations. A typical framework for answering this question is expressing the top-of-the-atmosphere (TOA) radiative imbalance,  $\Delta R$ , as

$$\Delta R = \Delta F + \lambda \Delta T, \quad (1)$$

25 where  $\Delta F$  is the radiative forcing,  $\lambda$  is the net climate feedback parameter, and  $\Delta T$  is the global mean surface temperature response. The feedback parameter  $\lambda$  is the increase in outgoing radiation per degree warming with units of W m<sup>-2</sup> K<sup>-1</sup> and represents the effects of all global average radiative feedbacks combined. Using this forcing-feedback framework, we can compute the equilibrium climate sensitivity (ECS), which is the global mean surface temperature response needed to restore the TOA imbalance to zero after doubling CO<sub>2</sub>, (Sherwood et al., 2020), as

$$30 \text{ ECS} = \frac{\Delta F_{2 \times \text{CO}_2}}{-\lambda}. \quad (2)$$

The complexity of the climate system and observational uncertainty lead to large uncertainties in estimates of ECS, with the climate feedback parameter  $\lambda$  considered a greater source of uncertainty in ECS than the forcing  $\Delta F$  (Sherwood et al., 2020). The uncertainty in  $\lambda$  stems from the significant uncertainty in its components, notably the cloud and water vapor feedbacks (Roe and Baker, 2007; Andrews et al., 2012; Vial et al., 2013; Sherwood et al., 2020). Feedback uncertainty is also important  
35 on regional scales. For instance, the Arctic, which is warming faster than the global average in a phenomenon known as Arctic amplification (AA), ~~is characterized by~~ faces considerable feedback uncertainty, making it ~~diffieult~~ challenging to attribute warming to individual feedbacks (Pithan and Mauritsen, 2014; Hahn et al., 2021; Shi and Lohmann, 2024).

The net feedback parameter can be linearly decomposed into a sum of individual feedbacks:  $\lambda = \sum_i \lambda_i$ , where  $\lambda_i$  represents the contributions of individual feedbacks: lapse rate, Planck, water vapor, surface albedo, and cloud feedbacks. There are two  
40 caveats to this decomposition worth noting. First, representing  $\lambda$  as a linear combination of individual feedbacks ignores the interaction *between* feedbacks, which can be ~~important~~ crucial, especially on local scales (~~Feldl and Roe, 2013; Knutti and Rugenstein, 2015;~~ Feldl and Roe, 2013; Knutti and Rugenstein, 2015; Feldl et al., 2017; Huang et al., 2021; Bonan et al., 2025). Second,  $\lambda$  and its individual components are likely not constant, varying with the climate state and with the pattern of surface temperature change (Knutti and Rugenstein, 2015; Gregory and Andrews, 2016; Dong et al., 2019; Meyssignac et al., 2023). Even with  
45 these caveats, the linear decomposition of feedbacks remains a commonly used framework.

The most common way to calculate individual radiative feedbacks is ~~by~~ using radiative kernels (Soden et al., 2008). Radiative kernels are the pre-calculated radiative sensitivities at some vertical level, often the TOA, to incremental changes in climate variables, such as temperature, water vapor, and surface albedo. The TOA radiative imbalance due to feedbacks,  $\Delta R_\lambda$  (equivalent to  $\lambda \Delta T$ , see Eq. 1), is decomposed as

$$50 \quad \Delta R_\lambda = \sum_i \frac{\partial R_i}{\partial x_i} \Delta x_i, \quad (3)$$

where  $\frac{\partial R}{\partial x}$  is the radiative kernel, and  $\Delta x$  is the change in a climate variable (e.g., following  $2\times\text{CO}_2$ ). The radiative kernel method offers several advantages over other methods of calculating radiative feedbacks. For example, radiative kernels can be applied to virtually any gridded data (e.g., climate model output, reanalysis products, etc.) as long as standard variables ~~such as temperature and~~ (temperature, specific humidity, etc.) are available. Using existing radiative kernels also alleviates the need to perform computationally expensive partial radiative perturbation calculations or run offline radiative transfer models (Wetherald and Manabe, 1988; Colman and McAvaney, 2011; Smith et al., 2020). Another use of radiative kernels is the decomposition of the effective radiative forcing into individual components, allowing for the separation and quantification of specific adjustments, such as changes in cloud properties or aerosol concentrations (Larson and Portmann, 2016).

The underlying assumption of the radiative kernel method is that ~~differences between kernels produced using variations in~~ kernels produced with different models are minor compared to ~~differences-discrepancies~~ in the climate responses ~~between across~~ models. This is because interkernel variation stems only from differences in radiative transfer models and model base states (Pincus et al., 2020), ~~both of~~ which are, ideally, physically reasonable representations of the real world (Soden et al., 2008). ~~This-The~~ assumption of minor differences ~~between-across~~ kernels enables intermodel feedback comparisons and allows for ~~the use of using~~ virtually any radiative kernel to calculate feedbacks.

A question naturally follows: are the differences between kernels actually minor? Only a few studies have addressed this question. Zelinka et al. (2020) assessed the variability of global radiative feedbacks across six kernels (their Fig. S2). Soden et al. (2008) found that among ~~three-four~~ kernels calculated using ~~three~~ different models, vertically integrated zonal mean kernels varied by  $\sim 10\%$  ~~except for in~~, except for the Southern Ocean, where they varied by  $\sim 30\%$ . ~~In the context of the global mean~~, Global mean temperature and water vapor kernels ~~only varied by~~ varied by less than  $\sim 5\%$ , although the surface albedo kernel varied considerably more ( $\sim 18\%$ ). In a more recent study, Hahn et al. (2021) found ~~considerable spread in global and regional surface albedo and cloud feedbacks calculated from different kernels~~ that the relative importance of feedbacks as polar amplification mechanisms shows kernel dependence. Huang and Huang (2023) documented a new set of kernels and found agreement in the global mean TOA feedbacks among ~~the seven kernels they considered~~ seven sets of kernels but notable differences in ~~the surface feedbacks~~ feedbacks at the surface. Although we do not seek to ~~completely~~ answer the question of the importance of interkernel differences ~~here we note that given the popularity of the radiative kernel method, completely, we feel that~~ it deserves more attention ~~however given the method's popularity. However~~, the current research environment makes intercomparing ~~different~~ radiative kernels difficult.

Using different sets of kernels introduces uncertainty that can limit the reproducibility and robustness of climate feedback studies. First, although many kernels have been produced since the early studies of Soden et al. (2008) and Shell et al. (2008), they are scattered among different research groups and institutions, making them difficult to locate; even after accessing a kernel, there is often little to no guidance on their proper usage. Second, kernels vary considerably in their properties, such as horizontal and vertical grids, model tops, sign conventions, and nomenclature, which may introduce calculation discrepancies across studies. Lastly, using kernels to calculate radiative feedbacks requires several choices and assumptions; examples include what base temperature to use when calculating the specific humidity increase from a ~~1K~~ 1 K increase in atmo-

85 spheric temperature and how to handle vertical integration to the surface while accounting for surface pressure and terrain  
(Pendergrass, 2019; Huang and Huang, 2023). These three factors make comparing results between feedback studies difficult,  
even when studies may use the same radiative kernels.

To standardize radiative feedback calculations and establish a central kernel repository, we created the ClimKern project.  
This project consists of two distinct parts: the ClimKern Python package, an open-source library for computing radiative  
90 feedbacks, and the ClimKern repository, which provides easy access to 11 sets of radiative kernels ~~, as of this writing,~~ computed  
from various climate models, reanalyses, and satellite observations. The package provides functions for calculating radiative  
feedbacks using any of the radiative kernels in just one or two lines of code per feedback. The package greatly enhances the  
reproducibility of feedback studies by standardizing the assumptions and choices. It also enables straightforward interkernel  
comparisons to better understand the role of kernel choice in ~~these~~ feedback studies.

95 The remaining sections are organized as follows: section 2 provides detailed information about ClimKern radiative kernels  
and the sample data we included for demonstration purposes. Section 3 covers the methodological choices made in crafting  
the feedback calculation functions. Section 4 shows the results of using the package with the sample climate model output to  
calculate feedbacks. In section 5, we put our package and the sample results in the context of the greater climate feedback and  
sensitivity community.

## 100 2 Data

### 2.1 Radiative kernels

We acquired 11 sets of all-sky (with clouds) and clear-sky (cloudless) TOA radiative kernels that were ~~either~~ publicly available  
or ~~made available~~ provided to us by the creators. To be included in ClimKern, a kernel product must have 4-dimensional  
water vapor and air temperature kernels, as well as 3-dimensional surface temperature and surface albedo kernels. The kernels  
105 must be monthly averages to capture the seasonal variations in TOA radiative fluxes and must be on horizontal latitude-  
longitude grids. The above requirements were chosen to ensure ease of use and that feedback calculations using different  
kernels are directly comparable. In this ~~first~~ version of ClimKern, we excluded radiative kernels that require nontraditional  
(i.e., considerably different from Soden et al. (2008)) variables to compute feedbacks; examples include the cloud kernels from  
Zelinka et al. (2012), which require ~~satellite-simulator-produced~~ satellite-simulator-produced output, and new kernels from the  
110 NASA Goddard Institute for Space Studies that use column precipitable water and sea ice fraction variables (Zhang, 2023). We  
also excluded band-by-band or “spectral” kernels, such as those in Bani Shahabadi and Huang (2014) and Huang et al. (2024).

Seven of the 11 TOA kernel sets had corresponding surface kernels for calculating radiative feedbacks from a surface perspec-  
tive, as in Pithan and Mauritsen (2014) and Laîné et al. (2016). Although they are included in the repository for ease of access,  
surface feedback calculations have not been implemented in ClimKern, and our discussion exclusively focuses on TOA kernels  
115 and feedbacks. Future versions of ClimKern may expand compatibility to surface kernels and other kernel types.

**Table 1.** The 11 radiative kernel sets included in the ClimKern repository. ~~From left to right, the~~ The table contains the kernel names, ~~the~~ horizontal resolution, ~~the~~ number of vertical levels and highest pressure level in ~~the~~ their ClimKern version ~~of the kernel, and the highest~~ pressure level of the kernel main data source used. Additionally, ~~and the paper that first documents~~ reference documenting the kernel ~~is~~ provided.

<del>kernel name</del> <u>Kernel</u>	<del>horizontal resolution (lat×lon)</del> <u>Res. (lat×lon)</u>	<del>number of vertical levels</del> <u>Vert. Levels</u>	<del>highest pressure level (hPa)</del> <u>Min P (hPa)</u>
BMRC	3.2° <del>x</del> <u>°</u> ×5.6° <del>°</del> <u>°</u>	19	1
CAM3	2.8° <del>x</del> <u>°</u> ×2.8° <del>°</del> <u>°</u>	17	10
CAM5	0.94° <del>x</del> <u>°</u> ×1.25° <del>°</del> <u>°</u>	22	3.64
CERES	0.5° <del>x</del> <u>°</u> ×1° <del>°</del> <u>°</u>	30	0.1
CloudSat	2° <del>x</del> <u>°</u> ×2.5° <del>°</del> <u>°</u>	17	10
ECHAM6	1.88° <del>x</del> <u>°</u> ×1.88° <del>°</del> <u>°</u>	19	1
ECMWF-RRTM	2.5° <del>x</del> <u>°</u> ×2.5° <del>°</del> <u>°</u>	24	1
ERA5	2.5° <del>x</del> <u>°</u> ×2.5° <del>°</del> <u>°</u>	37	1
GFDL	2° <del>x</del> <u>°</u> ×2.5° <del>°</del> <u>°</u>	17	10
HadGEM2	1.25° <del>x</del> <u>°</u> ×1.88° <del>°</del> <u>°</u>	19	1
HadGEM3-GA7.1	1.25° <del>x</del> <u>°</u> ×1.9° <del>°</del> <u>°</u>	39	3

<sup>a</sup> For kernels with a satellite data source, reanalysis data were used to supplement calculations.

Details about each kernel set can be found in Table 1. These 11 kernel sets were developed independently using ~~a variety of~~ various data sources for their base states, ~~including:~~ climate model output, reanalysis data, and satellite observations (Soden et al., 2008; Huang et al., 2017; Kramer et al., 2019). Horizontal resolutions range from several degrees to under one degree in latitude and longitude. Nearly all the kernels were already available on standard pressure levels, the desired vertical coordinate to ensure ~~kernel compatibility with various~~ compatibility with climate model output to calculate feedbacks. Kernels available on their native model grids (i.e., CAM5 & HadGEM3-GA7.1) were linearly interpolated to pressure levels. The native CAM5 kernels were available on hybrid sigma-pressure vertical coordinates; isobaric levels in the upper troposphere were unchanged, while hybrid levels in the lower and mid-troposphere were converted to the standard pressure levels used in the Coupled Model Intercomparison Project Phase 6 (CMIP6) (Eyring et al., 2016). The native HadGEM3-GA7.1 kernels were on a pure sigma vertical coordinate that lacks isobaric surfaces. Because they were specifically developed with a high model top and enhanced vertical resolution to capture stratospheric adjustments (Smith et al., 2020), they were interpolated to 39 pressure levels, the highest standard CMIP6 vertical resolution. ~~For further details about individual kernels, see the~~ We also included in Table 1 information regarding the data used to generate the kernel sets. See the corresponding citing papers in Table 1 ~~for additional information about each kernel set.~~

After ~~we collected the kernels and performed the necessary regridding~~ collecting and regridding the kernels, we combined each kernel set into one netCDF file per ~~kernel~~ source. The native kernel variables were renamed ~~so as~~ to have a standardized

set of variables, and their units or other metadata were altered for consistency and accuracy. For example, all surface albedo kernels had their units changed to  $\text{Wm}^{-2}\%$  if needed. We then inspected the kernels to find inconsistencies with their sign conventions, which were corrected. The resulting dataset was uploaded to Zenodo (Janoski et al., 2024a), ~~from where it~~ which can be downloaded manually or via a built-in script in the ClimKern Python package.

## 2.2 Sample climate model output

We also provide a tutorial dataset within the package to calculate, for verification purposes, the feedbacks given in Table 2. The package includes a function that accesses the sample data derived from pre-industrial and abrupt- $2\times\text{CO}_2$  fully coupled runs using the Large Ensemble version of the Community Earth System Model 1 (CESM1-LE). The CESM1-LE model incorporates the Community Atmosphere Model version 5 (CAM5) with 30 vertical levels and the Parallel Ocean Program version 2 (POP2) with 60 vertical levels. The model operates at a horizontal resolution of  $1^\circ$  across all components (Kay et al., 2015). These experiments have been extensively documented in prior studies (Mitevski et al., 2021, 2022, 2023). We also provide the effective radiative forcing (ERF) ~~from these experiments, calculated from simulations with prescribed,~~ calculated as the difference between the global mean net TOA flux between an abrupt  $2\times\text{CO}_2$  run and pre-industrial ~~sea surface temperatures and sea ice concentrations as in Forster et al. (2016).~~ Lastly, we provide the instantaneous radiative forcing (IRF) for the same dataset ~~calculated offline with the radiative transfer model SOCRATES (Edwards and Slingo, 1996; Manners, 2015). Note that because SOCRATES is not the radiative transfer scheme used in CESM, it may not yield perfect energy budget closure (e.g., following Eq. 1) even with correct kernel decomposition.~~ control run where the sea-surface temperatures and sea-ice concentrations are fixed to pre-industrial conditions in both runs (Forster et al., 2016).

## 3 Feedback calculations

The ClimKern Python package, hereafter referred to as simply “ClimKern,” contains many built-in functions for calculating radiative feedbacks and other valuable quantities of interest. Note that all ~~output outputs~~ from feedback functions ~~is in the form of the~~ are TOA radiative perturbations from the feedback in units of  $\text{Wm}^{-2}$ ; if the user wishes to express feedback values ~~as~~ per unit temperature ( $\text{Wm}^{-2}\text{K}^{-1}$ ), that ~~is usually can be~~ achieved by dividing or regressing by the surface temperature response. We avoided incorporating this step into the functions as there are ~~other several~~ ways of expressing feedbacks, such as in the form of warming contributions (Pithan and Mauritsen, 2014; Goosse et al., 2018; Previdi et al., 2020; Janoski et al., 2023), and the radiative perturbations in  $\text{Wm}^{-2}$  are ~~useful to~~ helpful in calculating rapid adjustments to radiative forcing ~~(Smith et al., 2018).~~ (Vial et al., 2013; Block and Mauritsen, 2013; Smith et al., 2018). All functions can compute all-sky and clear-sky feedbacks depending on the sky argument, either "all-sky" or "clear-sky". Additionally, because radiative kernels are typically only available as monthly means and monthly means are standard climate model output, ClimKern currently only accepts monthly mean input fields. Below, we document the required user input for the feedback calculation functions and provide details on their methodologies. This is not an exhaustive list of functions available in ClimKern, and specifics are

subject to change in future versions. Still, we hope it will prove helpful to discuss the philosophy behind the design of each function.

### 165 3.1 Temperature feedbacks

Temperature feedbacks refer to the radiative perturbations at the TOA from changes in the surface and atmospheric temperatures. Traditionally, the total temperature feedback is decomposed into the Planck feedback (or Planck response, depending on the specific definition of “feedback” used) and the lapse rate feedback (Soden and Held, 2006; Bony et al., 2006; Soden et al., 2008). The Planck feedback is the radiative response to a vertically uniform temperature change of equal magnitude  
170 to that of the surface; it is the most fundamental response of the radiative budget to a change in temperature, following the Stefan-Boltzmann law (Previdi et al., 2021). The lapse rate feedback ~~differs in that it reflects~~ is instead the *deviation* from vertically uniform warming to quantify the radiative effects of an altered tropospheric lapse rate.

ClimKern provides the `calc_T_feedbacks` function that computes the tropospheric Planck and lapse rate feedbacks using user-provided 4-D air temperature and 3-D surface temperature and pressure fields from two climate model simulations: a  
175 control simulation (representing a baseline state of the climate system) and a perturbed simulation ~~;~~ -(representing the climate system under changed conditions, such as increased greenhouse gas or aerosol concentrations), the difference of which is used to calculate the temperature response. In the tutorial data provided with ClimKern, these are a  $1\times\text{CO}_2$  and a  $2\times\text{CO}_2$  (relative to preindustrial levels) simulation, respectively. Reanalysis data can be used ~~in a similar fashion~~ similarly by separating data into two time periods for comparison.

180 First, ClimKern checks the input to ensure its format is compatible, including checking the time dimensions and units; then, the function will either proceed, issuing a warning to the user if any assumptions are made for missing metadata ~~;~~ or return an error for major incompatibilities (e.g., not providing input in the form of an Xarray DataArray (Hoyer and Hamman, 2017)). If the user did not provide an optional model- or user-defined tropopause, ClimKern will create a tropopause defined as 100 hPa at the Equator and linearly increasing with the cosine of latitude to 300 hPa at the poles. It will also read in the user-selected  
185 temperature and surface temperature kernels from locally stored package data. Using the xESMF module (Zhuang et al., 2023), the kernels are horizontally regridded using bilinear interpolation with periodic boundary conditions to match the resolution of the input model data. We elected to horizontally regrid to the input data’s resolution so that the user always receives output on the same horizontal grid as the input.

Following this setup, ClimKern creates a monthly climatology from the control simulation surface and atmospheric tempera-  
190 tures and subtracts it from the perturbed simulation fields, yielding a surface and air temperature response. ClimKern uses the control simulation’s climatological surface pressure to mask values below the surface for the air temperature response. The air temperature response is linearly interpolated to match the vertical kernel resolution; subsequent testing for tropospheric feedbacks at the TOA demonstrates little difference if the input vertical resolution is used (not shown). Layer thicknesses are then calculated ~~to be used in the~~ for the subsequent vertical integration of the temperature feedbacks ~~in the subsequent step.~~



195 ~~Note that the~~ The user-supplied perturbed simulation pressure and optional tropopause height are used when calculating the layer thicknesses to ensure that the vertical integration only extends from the surface to the tropopause.

The total air temperature response is decomposed into a vertically uniform component and ~~a deviation therefrom~~ deviation to calculate the Planck and lapse rate feedbacks separately. Both feedbacks are calculated by multiplying the respective temperature response component, temperature kernel, and layer thickness array and taking a sum along the pressure axis. In the case of the Planck feedback, the surface temperature response is multiplied by the surface temperature kernel and added to this sum. The function then returns both feedbacks. To the user, all of this culminates in two lines of code:

---

```
1: import climkern as ck
2: LR, Planck = ck.calc_T_feedbacks(ctrl.T, ctrl.TS, ctrl.PS,
205 3: ----- pert.T, pert.TS, pert.PS, pert.TROP_P, kern="GFDL")
4: ~~~~~ pert.T, pert.TS, pert.PS,
5: ~~~~~ pert.TROP_P, kern="GFDL",
6: ~~~~~ sky="all-sky", fixRH=False)
```

---

210 where LR and Planck are the vertically integrated, monthly- and spatially varying lapse rate and Planck feedbacks, respectively, ctrl and pert are Xarray Datasets (Hoyer and Hamman, 2017) ~~holding~~ containing the control and perturbed ~~simulations~~ simulation output, T is the 4-dimensional air temperature, TS is the 3-dimensional surface temperature, PS is the 3-dimensional surface pressure, TROP\_P is the optional 3-dimensional tropopause height, and kern is the optional kernel choice argument. All feedback calculations share this kern argument, which defaults to "GFDL" ~~;~~ to specify which of the 11 kernels ClimKern should use.

215 This function contains several ~~options~~ optional arguments, including the kernel name, tropopause heights, whether to calculate the all-sky or clear-sky feedbacks, and whether to ~~calculate the feedbacks using~~ use relative humidity as a state variable, as in Held and Shell (2012). Further details about the computations and optional parameters can be found in the source code ~~located~~ in Janoski et al. (2024b).

## 220 3.2 Water vapor feedback

ClimKern also offers a `calc_q_feedbacks` function to compute water vapor feedbacks:

```
q_lw, q_sw = ck.calc_q_feedbacks(ctrl.Q, ctrl.T, ctrl.PS,
                                pert.Q, pert.PS, pert.TROP_P,
                                kern="GFDL", method=1)
```

225 where `q_lw` and `q_sw` are the TOA radiative perturbations from the longwave and shortwave water vapor feedbacks, respectively, `Q` is the 4-dimensional specific humidity, and all other variables are as they are in Section 3.1. Note that a “control” air temperature variable is required because water vapor kernels are traditionally calculated not using a unit increase in specific



humidity but rather the specific humidity change corresponding to a 1K increase in temperature with constant relative humidity (Shell et al., 2008); consequently, the units of the water vapor kernels are  $\text{Wm}^{-2}\text{K}^{-1}$ .

230 The ~~basic~~ flow of the function is similar to that of the temperature feedbacks: first, ClimKern checks all input data and tries to identify proper units. If the user did not provide a DataArray with tropopause pressure, ClimKern constructs a default one. Next, ClimKern produces a monthly climatology of the control simulation surface pressure, specific humidity, and air temperature, ~~masking values below the surface. ClimKern then computes the specific humidity response using a linear or logarithmic approach according to the method argument discussed below. As for the temperature feedbacks. Like the temperature feedback~~  
 235 ~~function,~~ the kernels are regridded to the horizontal grid of the input data ~~while,~~ and the climatologies of the specific humidity ~~and air temperature,~~ air temperature, and the specific humidity response are put on kernel pressure levels. Values on pressure levels below the control simulations's climatological surface pressure are masked and not included in further calculations. The product of the kernel, specific humidity response, and layer thickness is vertically integrated over the troposphere; a normalization factor must also be included, as discussed below.

240 ~~Using a temperature perturbation to produce the~~ Water vapor feedbacks are commonly computed using the change in the natural log of specific humidity because water vapor's absorption of longwave radiation is roughly proportional to the logarithm of its concentration (Shell et al., 2008; Lacis et al., 2013; Colman and Soden, 2021). The change in the natural log of specific humidity can be written as:

$$\Delta \ln(q) = \ln(q_{\text{pert}}) - \ln(q_{\text{ctrl}}), \quad (4)$$

245 ~~where  $q_{\text{pert}}$  and  $q_{\text{ctrl}}$  are the perturbed and control specific humidities, respectively. Using logarithm properties, this can be expressed as:~~

$$\Delta \ln(q) = \ln\left(1 + \frac{\Delta q}{q_{\text{ctrl}}}\right), \quad (5)$$

~~where  $\Delta q = q_{\text{pert}} - q_{\text{ctrl}}$ . For small changes in  $q$  such that  $(\Delta q/q_{\text{ctrl}}) \ll 1$ , the natural logarithm can be approximated using a first-order Taylor expansion:~~

250 
$$\ln\left(1 + \frac{\Delta q}{q_{\text{ctrl}}}\right) \approx \frac{\Delta q}{q_{\text{ctrl}}}, \quad (6)$$

~~leading to the fractional approximation used in Pendergrass (2019):~~

$$\Delta \ln(q) \approx \frac{\Delta q}{q_{\text{ctrl}}}. \quad (7)$$

~~ClimKern allows the user to choose whether to use this legacy fractional approximation or use the more precise, actual difference in natural logs of specific humidity for water vapor feedback calculations via the method parameter in the `calc_q_feedback` function, outlined below. Alternatively, users may instead use the linear change in specific humidity, i.e.,~~  
 255

$$\Delta q = q_{\text{pert}} - q_{\text{ctrl}}. \quad (8)$$

~~The water vapor kernels requires the water vapor kernels to~~ must be normalized by the change in specific humidity per unit temperature ~~before calculating the feedback itself~~ increase. Ideally, one would ~~use the change in specific humidity per unit temperature from the kernel calculation to normalize the kernel~~ have this field from the kernel-producing simulation, as in  
 260 Shell et al. (2008) and Pendergrass (2019), but, in practice, this quantity is rarely included with the distributed kernels. Given ~~that we have little information~~ the little information available about the base states used in the individual kernel calculations, ClimKern utilizes the climatological air temperature from the user-provided control simulation to produce a water vapor kernel normalization factor using the Buck (1981) empirical formula for saturation vapor pressure.

~~The calc\_q\_feedbacks function also contains a unique “method” parameter that accepts one of four numeric arguments (1-4). Past studies vary in the way they compute the water vapor feedback—namely, whether to use the natural logarithm of water vapor concentration and, if so, whether the log is approximated as the fractional change in water vapor. It is most common to use the natural log of specific humidity. Note that the change in water vapor feedback calculations because the absorption of longwave radiation by water vapor is roughly proportional to the logarithm of its concentration (Shell et al., 2008; Lacis et al., 2013; Colman et al., 2013); however, we also included an option to use the linear change in specific humidity, as in Pendergrass (2019). The difference in natural logarithms of specific humidity has sometimes been approximated as the fractional change in specific humidity. ClimKern includes that option for the specific humidity response and the normalization factors~~ specific humidity per unit temperature increase can similarly either be the logarithmic or linear change and, in the case of the former, use the fractional approximation.  
 270

The calc\_q\_feedbacks function contains four “method” options to accommodate the variations in the literature described above. The options and corresponding numeric arguments are:  
 275

1. Uses the actual logarithm for both the specific humidity response and ~~the~~ normalization factor.
2. Uses the ~~fractional change approximation of logarithms only in the normalization factor, with the actual logarithm used in actual logarithm for~~ the specific humidity response and fractional approximation for the normalization factor.
3. Uses the fractional ~~change approximation of logarithms approximation~~ in the specific humidity response ~~& and~~ normal-  
 280 ization factor.
4. Uses the linear change ~~in specific humidity for both~~ for both the specific humidity response and normalization factor.

The function defaults to option 1. Further details can be found in the function’s docstring (Janoski et al., 2024b).

### 3.3 Surface albedo feedback

The `calc_alb_feedback` function, which computes surface albedo feedback, is relatively straightforward; it requires the  
 285 user to provide the upwelling and downwelling shortwave radiation at the surface from the control and perturbed simulations. The first step is to compute the surface albedo as the ratio of surface upwelling to downwelling radiation while masking areas with a downwelling radiation value of  $0 \text{ Wm}^{-2}$ . ClimKern then takes the difference between the perturbed simulation’s albedo

and the control simulation’s monthly climatological albedo. The desired albedo kernel is loaded from memory, regridded to the input horizontal resolution, and multiplied by the albedo response to produce the surface albedo feedback. ~~As in the other~~  
 290 ~~ClimKern feedback functions, users may specify the kernel to use and whether to compute the all-sky or clear-sky feedback.~~

### 3.4 Cloud feedbacks

Cloud feedbacks are comparatively more complicated than the other feedbacks, owing to nonlinearities in kernel computations and the vertical overlapping of clouds (Soden and Held, 2006; Soden et al., 2008; Shell et al., 2008). Consequently, ~~most~~  
 295 cloud feedbacks — most commonly, the residual and adjustment methods. ClimKern contains a function for each method, which we will detail below. Note that for both methods, ClimKern ~~requires one or more~~ optionally accepts radiative forcing terms that will vary with the ~~given~~ experimental setup (i.e., control and perturbation simulations). In other words, there is not a precise type of radiative forcing quantity ~~, including ERF and IRF,~~ that will suit every scenario. ClimKern avoids making assumptions regarding the forcing, ~~and it is up to the user to ensure that all terms in the radiative budget are being properly~~  
 300 ~~accounted for. In~~; if the user does not provide it, cloud feedback functions assume the radiative forcing terms are zero by default. We use the ERF to compute cloud feedback in our sample results (Section 4), ~~we use the ERF to compute cloud~~  
~~feedbacks.~~

#### 3.4.1 Residual method

In the residual method, the cloud feedbacks are computed as a residual of the TOA energy budget, that is:

$$305 \quad \Delta R_{cloud} = \Delta R_{all-sky} - \Delta F - \sum_i \Delta R_i \quad (9)$$

where  $\Delta R_{cloud}$  is the TOA radiative perturbation from the cloud feedback,  ~~$\Delta R_{all-sky}$  is the~~  $\Delta R$  is the all-sky net TOA radiative imbalance,  $\Delta F$  is the radiative forcing (e.g., from  $\text{CO}_2$ ), and  $\sum_i \Delta R_i$  is the sum of the TOA radiative perturbations from other non-cloud feedbacks (Soden and Held, 2006; Zhang et al., 2018; Zhu et al., 2019). Put another way, the cloud feedback is assumed to be the missing piece in the TOA radiative budget after accounting for other terms. Although this method provides a  
 310 “clean” approach that fully closes the radiative budget in a kernel feedback decomposition, it carries two main drawbacks. First, it is highly sensitive to uncertainties in the other terms and ~~, especially, in~~ the often unavailable radiative forcing,  $\Delta F$  (Soden et al., 2008). Second, since the cloud feedbacks are assumed to close the radiative budget, feedback decompositions using this method yield no separate error estimate, which is sometimes ~~useful in~~ valuable for evaluating radiative kernels. Despite these disadvantages, the residual method is still ~~widely~~ used.

315 ClimKern contains separate `calc_cloud_LW_res` and `calc_cloud_SW_res` functions to calculate the longwave and shortwave cloud feedbacks, respectively. For the longwave, ClimKern requires net longwave radiative flux at the TOA from the control and perturbed simulations, the longwave all-sky radiative forcing, and the radiative perturbations from the total temperature and longwave water vapor feedbacks. The shortwave function instead requires the net shortwave radiative flux at

the TOA in the control and perturbed simulations, the shortwave all-sky radiative forcing, and the radiative perturbations from the surface albedo and shortwave water vapor feedbacks. From there, both functions compute the cloud feedback ~~is computed~~ using Eq. ~~4 in both functions~~9.

### 3.4.2 Adjustment method

The adjustment method for calculating cloud feedbacks is named as such because the change in cloud radiative effect (CRE) is “adjusted” for masking by other feedbacks and the radiative forcing to produce a cloud feedback:

$$\Delta R_{cloud} = \Delta CRE + \sum_i (\Delta R_{i,clear-sky}^o - \Delta R_{i,all-sky}^i) + (\Delta F_{clear-sky}^o - \Delta F_{all-sky}). \quad (10)$$

where  $\Delta CRE$  is the CRE response,  ~~$\Delta R_{i,clear-sky}$  and  $\Delta R_{i,all-sky}$~~   $\Delta R_i^o$  and  $\Delta R_i^i$  are the clear-sky and all-sky radiative feedbacks, and  ~~$\Delta F_{clear-sky}$  and  $\Delta F_{all-sky}$~~   $\Delta F^o$  and  $\Delta F$  are the clear-sky and all-sky radiative forcings (Soden et al., 2008; Zhang et al., 2018).  $\Delta CRE$  is computed as  ~~$\Delta R_{all-sky} - \Delta R_{clear-sky}$ , i.e.~~  $\Delta R - \Delta R^o$ , i.e. the difference in the all-sky and clear-sky TOA radiative flux. The adjustment method is considered less sensitive to uncertainties in the other terms, especially the forcing term (Soden et al., 2008). Additionally, since the resulting cloud feedback is not computed as a residual, it allows one to separately quantify the error in closing the TOA radiative budget.

The longwave and shortwave adjustment-method cloud feedbacks can be computed via the `calc_cloud_LW` and `calc_cloud_SW` functions. The longwave function ~~uses as input~~ accepts the change in the longwave CRE, ~~along with~~ and the all-sky and clear-sky radiative perturbations at the TOA from the total temperature feedback, longwave water vapor feedback, and long-wave radiative forcing. The shortwave function uses the shortwave versions of the longwave function input, except that it uses the surface albedo feedback instead of the temperature feedback. ClimKern includes separate `calc_dCRE_LW` and `calc_dCRE_SW` functions that evaluate the change in longwave and shortwave CRE and that require several radiative fields from the user, including the TOA all-sky and clear-sky LW or SW radiative fluxes in the control and perturbation simulations. After reading in all the necessary input, the adjustment method cloud feedback functions calculate the differences between the all-sky and clear-sky perturbations from non-cloud terms and combine them with the change in CRE to return the desired cloud feedback.

### 3.5 Other functions

We included several other utility functions in ClimKern. First, there are stratosphere versions of the temperature and water vapor feedback functions, named `calc_strato_T` and `calc_strato_q`, respectively. They are mostly analogous to their tropospheric counterparts, but the vertical integration is performed from the tropopause to the TOA. Next, ClimKern provides a `calc_RH_feedback` function to calculate the relative humidity feedback following Shell et al. (2008), Held and Shell (2012), and Zelinka et al. (2020). Typically, the relative humidity feedback would be a component of a radiative feedback decomposition if the user calculated the temperature feedbacks with the `fixRH` option. Finally, the `spat_avg` function

computes the spatial average of a DataArray while weighting for the cosine of latitude. We refer the reader to Janoski et al.  
350 (2024b) for additional documentation.

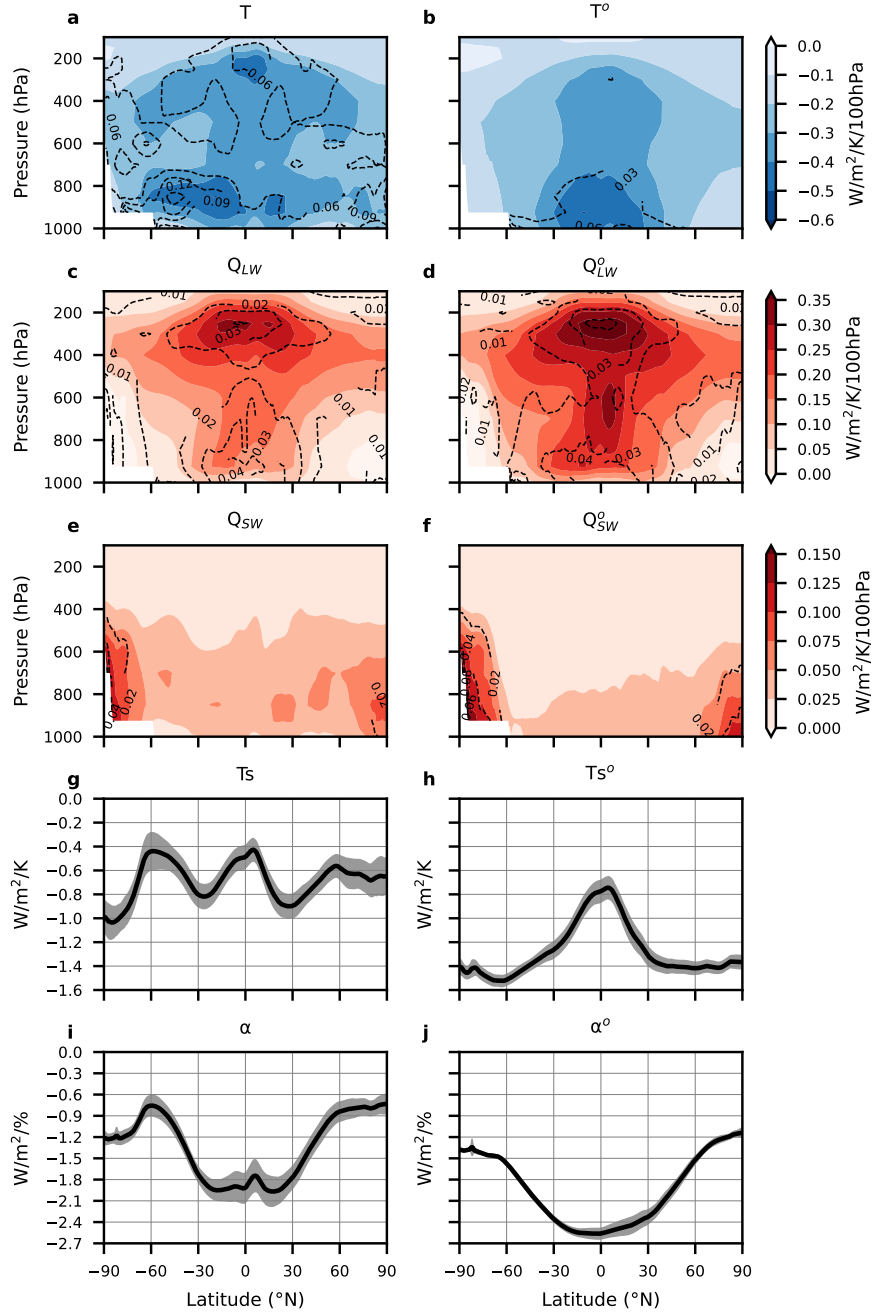
## 4 Results with sample data

### 4.1 Radiative kernels

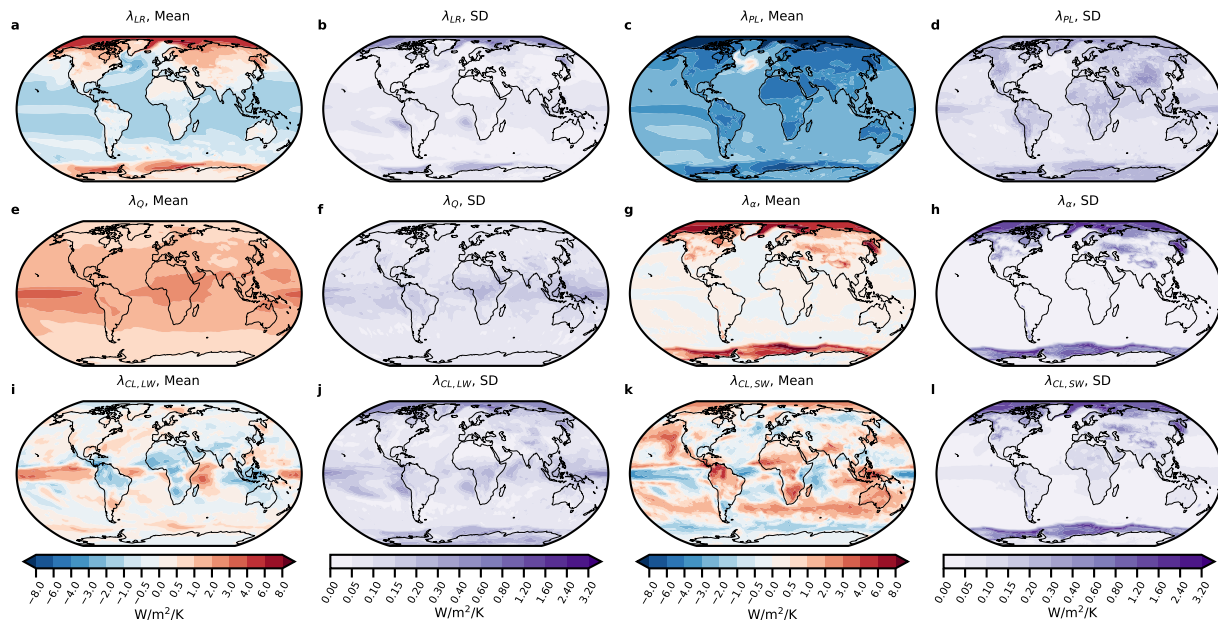
Having outlined the data and functions packaged with ClimKern, we now focus on the characteristics of the TOA radiative kernels. Fig. 1 shows the annual- and zonal-average mean kernel values and two across-kernel standard deviation ranges after  
355 linearly interpolating to a common 17 standard pressure levels. ~~Each kernel exhibits remarkably different spatial patterns in the kernel mean and standard deviations and even, in some cases, between the all-sky and clear-sky versions of the same kernel.~~ The mean and standard deviation of the all-sky air temperature kernels (Fig. 1a) have two local maxima in magnitude, one in the equatorial upper troposphere and the other in the mid-to-high-latitudes lower troposphere in the Southern Hemisphere (SH). In the clear-sky kernels, the lower tropospheric maximum is located over the Equator rather than the extratropical SH  
360 (Fig. 1b). Because the all-sky and clear-sky kernels differ only by the existence of cloud effects in their calculations, the different maxima locations are likely a result of clouds, which exert considerable influence on temperature kernels via cloud-top height temperatures (Kramer et al., 2019). ~~Overall, the~~ for example, high, cold clouds reduce outgoing longwave radiation efficiency by altering the effective emission height, affecting temperature kernel values (Kramer et al., 2019). The clear-sky air temperature kernels ~~have considerably less spread~~ exhibit less spread ( $\max \sim 0.06 \text{ Wm}^{-2}\text{K}^{-1}100\text{hPa}^{-1}$ ) than the all-  
365 sky kernels ( $\max \sim 0.13 \text{ Wm}^{-2}\text{K}^{-1}100\text{hPa}^{-1}$ ) (Fig. 1a-b), ~~highlighting implicating~~ the uncertainty introduced by clouds in radiative schemes.

The longwave water vapor kernels (Fig. 1c-d) do not appear to show as large sensitivity to clouds as the air temperature kernels, ~~with the exception of the~~ except for the deep tropics between 800 and 400 hPa. The longwave water vapor kernel mean is largest in the Equatorial upper troposphere and decreases with latitude, consistent with the findings of Huang et al.  
370 (2007). The pattern of the standard deviation mostly follows that of the mean with Equatorial maxima in both the high and low ~~troposphere-tropospheres~~ (Fig. 1c-d), indicating that this region is particularly sensitive to the base state and physics used in kernel production. The shortwave water vapor kernels (Fig. 1e-f) exhibit an increase in mean and standard deviation with latitude, opposite to that of the longwave kernels. As suggested by Huang and Huang (2023), the higher shortwave reflectivity of land and ice surfaces vs. ocean surfaces likely causes this behavior. Interkernel spread in the shortwave water vapor kernel  
375 is ~~significantly~~-larger near the poles, which may be due to differences in the radiative characteristics of the surface (e.g., sea ice extent, snow cover, etc.) in the kernel base states.

The surface temperature and albedo kernels are 3-dimensional, so the annual- and zonal averages only vary with latitude (Fig. 1h-k). The surface temperature kernel is highly sensitive to clouds, especially in the extratropics, as evidenced by ~~the presence of~~ local maxima near  $60^\circ\text{N/S}$  in the clear-sky kernels only (Fig. 1h-i). Interkernel spread is relatively constant with lat-  
380 itude. ~~The~~ In the case of the all-sky surface albedo kernel ~~is also sensitive to clouds, especially near  $60^\circ\text{S}$  (Fig. 1j-k).~~ Interkernel



**Figure 1.** (a-b) The mean (shaded) and ~~twice the~~ standard deviation (contoured, dashed) of the all-sky (left) and clear-sky (right) temperature kernels, representing the response to a ~~1K-1 K~~ increase in temperature. (c-d) as in (a-b), but for the longwave water vapor kernels, which reflect specific humidity changes associated with a ~~1K-1 K~~ warming and fixed relative humidity. (e-f) as in (a-b), but for the shortwave water vapor kernels. (h-i-g-h) The mean (solid line) ~~and two represents  $\pm$  one standard deviations deviation~~ (shading) of the all-sky (left) and clear-sky (right) surface temperature kernels. (j-k-i-j) as in (h-i-g-h), but for the surface albedo kernels, corresponding to a 1% increase in surface albedo.



**Figure 2.** Left column First and third columns from left: The annual mean, kernel-mean kernel mean (a) lapse rate, (c) Planck, (e) water vapor, (g) surface albedo and, (i) longwave cloud feedbacks calculated using the adjustment method, and (Wm<sup>-2</sup>K<sup>-1</sup>) shortwave cloud feedbacks. Right column Second and fourth columns: as for the left column first and third columns, but showing twice the standard deviation (SD) among kernels. Feedbacks Cloud feedbacks were calculated using the adjustment method. Values are expressed in units of Wm<sup>-2</sup>K<sup>-1</sup> by normalizing by the global mean surface temperature response. Note the different color bar nonlinear colorbar scales between feedbacks for both the means and standard deviations used to make maps comparable.

spread in the all-sky surface albedo kernels is largest in the tropics and largely constant with latitude in the, interkernel spread is larger in the tropics ( $\sim 0.4$  W/m<sup>2</sup>/%) than the extratropics ( $\sim 0.2$  W/m<sup>2</sup>/%). The clear-sky kernels version of the kernel exhibits the greatest variability in the Northern Hemisphere tropics ( $\sim 0.3$  W/m<sup>2</sup>/%) and decreases with latitude. The clear-sky surface albedo interkernel spread is less than the all-sky spread at all latitudes, indicating that clouds are a main contributor to the latter. In the next section, we explore how the spatial structure of the kernel means and interkernel spread influence the resulting radiative feedbacks.

#### 4.1.1 Feedback results

#### 4.2 Feedback results

Having quantified the kernels themselves, we now focus on how the interkernel differences translate into differences in individual feedbacks. Recall that all feedbacks are calculated using identical methodologies between kernels and with the same CESM1-LE sample model output data described in Section 2.2, so all such that feedback variability is solely the result of



the kernel choice. ~~All feedbacks~~ Feedbacks were computed using the difference in the monthly climatology of the last 30 years of the preindustrial control and a standard 150-year-long  $2\times\text{CO}_2$  simulation. Temperature and water vapor feedbacks are vertically integrated from the surface to the model-derived tropopause in the  $2\times\text{CO}_2$  simulation via the TROP\_P function argument. Water vapor feedbacks were calculated using the method 1 option, which uses the actual change in the natural log of specific humidity in calculations, although we later show global average results with the other three methods as well. Please note that all results use a single model run from CESM1-LE to compute the climate response.

The choice of defining the response as the difference between abrupt  $2\times\text{CO}_2$  and pre-industrial control simulations was made to minimize the sample data size distributed with ClimKern while aligning with this study's goal of highlighting interkernel spread. Additionally, this approach is commonly used, as in Pithan and Mauritsen (2014); Goosse et al. (2018); Previdi et al. (2020); Hahn et al. (2021). A notable consequence of this choice is that feedback values presented here include rapid adjustments that occur after  $\text{CO}_2$  increase (Zelinka et al., 2020; Hahn et al., 2021).

Fig. 2 shows the kernel mean (~~left column~~ first and third columns) and standard deviation (~~right column~~ second and fourth columns) of the lapse rate, Planck, water vapor, surface albedo, ~~and cloud feedbacks, vertically integrated when required~~ longwave cloud, and shortwave cloud feedbacks. Generally, the lapse rate and Planck feedbacks' mean and standard deviation ~~are of greater magnitudes~~ magnitudes are greater at the poles (Fig. 2a-d). The strong latitudinal gradient and sign change in the lapse rate feedback (Fig. 2a) are well-recognized features in climate model simulations subjected to increasing  $\text{CO}_2$ . They are products of latitudinal differences in lower- and upper-tropospheric coupling, sea ice loss, and heat transport (Manabe and Wetherald, 1975; Graversen et al., 2014; Feldl et al., 2020; Colman and Soden, 2021; Previdi et al., 2021). Similarly, the kernel mean Planck feedback is most negative over the Arctic with values less than  $-12\text{ Wm}^{-2}\text{K}^{-1}$  (Fig. 2c). ~~This is~~ where the surface temperature increase is greatest via Arctic amplification, ~~leading to the producing~~ large increases in outgoing longwave radiation via the Stefan-Boltzmann law. ~~Another area of strong~~ Strongly negative Planck feedbacks ~~occurs of less than  $-8\text{ Wm}^{-2}\text{K}^{-1}$  occur~~ over the Southern Ocean (Fig. 2c). The ~~overlapping maxima in interkernel spread in the spatial distribution of~~ lapse rate and ~~Planck feedbacks over~~ Planck feedback standard deviations imply enhanced kernel sensitivity in the Arctic and Southern Ocean ~~indicate that these regions have the greatest sensitivity to kernel choice~~ (Fig. 2b,d); ~~however, standard deviation values are small compared to some other feedbacks, including the surface albedo and cloud feedbacks.~~

The water vapor feedback is most positive in the tropical Pacific, where with values ranging from 2 to  $4\text{ Wm}^{-2}\text{K}^{-1}$  (Fig. 2e). Here, the increase in water vapor concentration per degree of warming is greatest ~~, following via~~ the Clausius-Clayperon relationship (~~Fig. 2e~~). Interkernel spread is also maximized in the tropics but exhibits a ~~markedly different spatial~~ different longitudinal distribution, with the greatest variability located over the Western Pacific (Fig. 2f). The maximum in standard deviation in the Western Pacific extends along the Equator and to the southeast, suggesting indicating that this feature may be related to the double Intertropical Convergence Zone (ITCZ) bias present in many climate models (Lin, 2007; Tian and Dong, 2020). ~~Additionally, there is a local standard deviation maximum in the South Atlantic Convergence Zone off the southeastern coast of South America. These maxima suggest large differences in the base states within the tropics across the eleven water~~

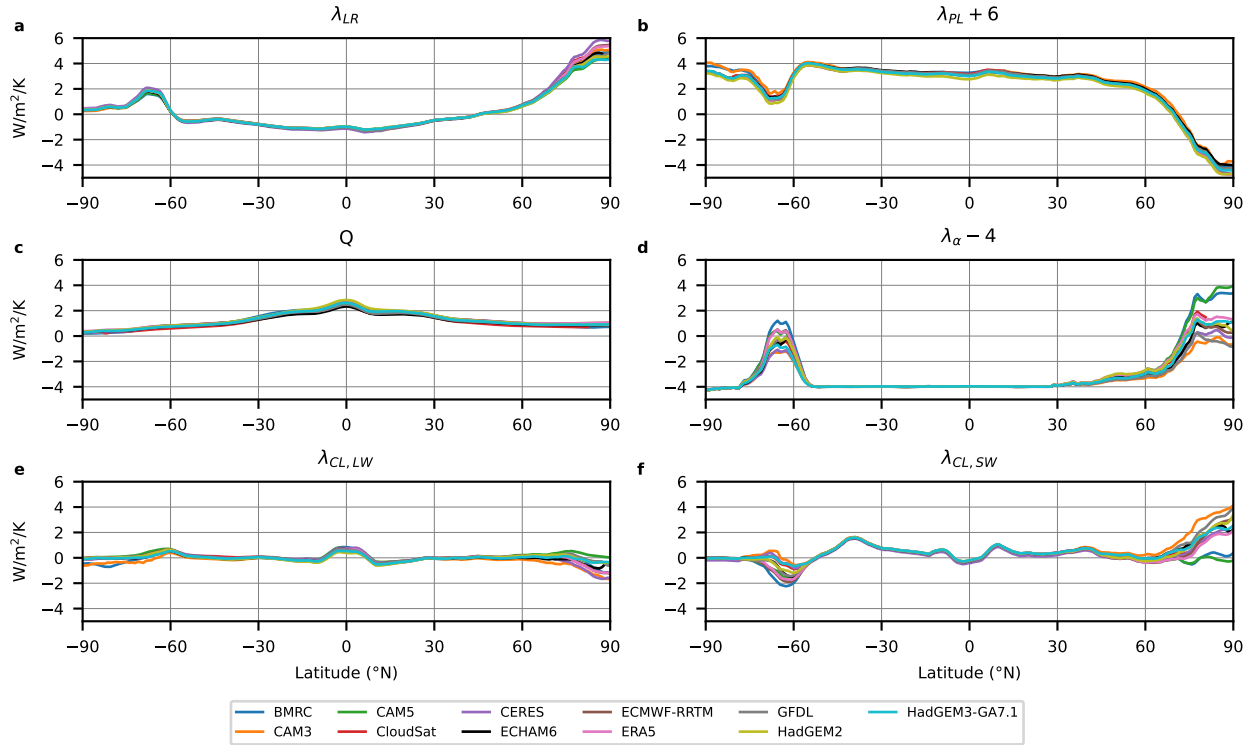
425 ~~vapor kernels.~~ As with the temperature feedbacks, interkernel spread is small ( $SD < 0.3 \text{ Wm}^{-2}\text{K}^{-1}$ ) relative to the feedbacks discussed next.

The surface albedo feedback is largest over high latitude oceans (Fig. 2h-i-g-h), driven by sea ice loss (Curry et al., 1995; Riihelä et al., 2021). This sea ice loss leads to large bottom-heavy warming in these regions, resulting in a strong positive lapse rate feedback, negative Planck feedback, and similarity in the spatial patterns of the lapse rate, Planck, and surface albedo  
430 feedbacks (Croll, 1875; Ingram et al., 1989; Previdi et al., 2021). It is important to note that in the Arctic and Southern Ocean, the ~~interkernel spread in standard deviation of~~ the surface albedo feedback is larger than ~~the spread that~~ of the lapse rate, Planck, and water vapor feedbacks, with ~~the albedo feedback standard deviation values~~ nearly as much as 50% of the ~~interkernel kernel~~ mean. Considering that there is little difference in the spread between the all-sky vs. the clear-sky albedo kernels (Fig. 1j-ki-j), it is likely not the clouds but the base-state sea ice conditions that produce the polar-amplified spread in albedo feedback.

435 The last ~~feedback we consider is the total cloud feedback~~ two feedbacks we consider are the longwave and shortwave cloud feedbacks in Fig. 2j-ki-l. The kernel mean longwave cloud feedback shows ~~considerable spatial inhomogeneity, tending to be more negative over the oceans and positive over land, and has maxima in Equatorial South America and Africa~~ spatial inhomogeneity with maxima in the tropical Pacific and western Indian Oceans and minima over northern South America, Africa, and Indonesia (Fig. 2i). The standard deviation of the longwave cloud feedback tends to increase poleward aside  
440 from relatively large values ( $\sim 0.4 \text{ Wm}^{-2}\text{K}^{-1}$ ) in the Equatorial Pacific (Fig. 2j). However, the interkernel spread in the cloud feedback is maximized Considerable spatial inhomogeneity is also found in the kernel mean shortwave cloud feedback, ranging from  $-4 \text{ Wm}^{-2}\text{K}^{-1}$  in the western Equatorial Pacific to  $7 \text{ Wm}^{-2}\text{K}^{-1}$  in northern South America (Fig. 2k). The shortwave cloud feedback standard deviation is largest over the Arctic and Southern Ocean (Fig. 2k). ~~We show below that this is mainly due to the spread of l).~~ The highest standard deviation values among all feedbacks are those of the surface albedo and shortwave  
445 cloud feedbacks in these regions, highlighting the importance of kernel choice near the poles.

Having analyzed the spatial distribution of interkernel spread, we focus on the differences between individual kernels in the zonal mean feedbacks in Fig. 3. The lapse rate and Planck feedbacks show minimal spread throughout the tropics and midlatitudes, with the greatest spread in the Arctic (Fig. 3a-b). The lapse rate feedback varies between  $4$  and  $6 \text{ Wm}^{-2}\text{K}^{-1}$  poleward of  $80^\circ\text{N}$  but varies by less than  $1 \text{ Wm}^{-2}\text{K}^{-1}$  elsewhere (Fig. 3a). For the Planck feedback, interkernel spread is greatest poleward  
450 of  $80^\circ$ , but generally less than  $1 \text{ Wm}^{-2}\text{K}^{-1}$  everywhere ~~and uniform in the zonal mean~~ (Fig. 3b). The zonal mean water vapor feedback ~~is most sensitive~~ shows little sensitivity to kernel choice but is most sensitive in the tropics with a spread of  $-0.5 \sim 0.5 \text{ Wm}^{-2}\text{K}^{-1}$  (Fig. 3c), similar to the spatial map (Fig. 2f).

~~We find very large interkernel~~ Interkernel spread in the surface albedo feedback (Fig. 3d). ~~The zonal mean Arctic surface albedo feedback spread is comparable to~~ is large relative to the kernel mean, especially in the Arctic. The area-average interkernel  
455 spread north of  $70^\circ\text{N}$  is  $2.8 \text{ Wm}^{-2}\text{K}^{-1}$ , roughly two-thirds of the kernel mean ~~feedback value. The Southern Ocean shows a similar but weaker interkernel~~ surface albedo feedback for the same region ( $4.2 \text{ Wm}^{-2}\text{K}^{-1}$ ). There is a similar, albeit weaker,



**Figure 3.** Annual zonal mean (a) lapse rate, (b) Planck, (c) total water vapor, (d) surface albedo, (e) longwave cloud, and (f) shortwave cloud feedbacks calculated using the adjustment method for each of the eleven kernels included in ClimKern in  $\text{Wm}^{-2}\text{K}^{-1}$ .  $6 \text{ Wm}^{-2}\text{K}^{-1}$  is added to the Planck feedback, and  $4 \text{ Wm}^{-2}\text{K}^{-1}$  is subtracted from the surface albedo feedback for visualization and comparison purposes.

interkernel spread in the surface albedo feedback spread in the Southern Ocean. These features are of particular importance particularly important in polar amplification studies, which we discuss in Section 5.

The cloud feedbacks are similarly-particularly sensitive to kernel choice because they are prone to uncertainties in the other feedback and radiative forcing terms, even when using the adjustment method (Soden et al., 2008). The interkernel variation-in the spread in the zonal mean longwave cloud feedback is largest at the poles, so-even-such that its sign in the Arctic depends on kernel choice (Fig. 3e). The shortwave cloud feedback (Fig. 3f) is likewise-similarly most sensitive to kernel choice in the high latitudes, with a zonal distribution of variability similar to that of the surface albedo feedback. This is likely-not a coincidence: the surface albedo feedback is used to calculate the shortwave cloud feedback via the adjustment method, so a large spread in the former translates to a large spread in the latter.

How do these zonal-average variations manifest in the global mean? We include the global-average-, annual mean feedback values for all 11 kernels in Table 2, along with the multikernel-multi-kernel mean and standard deviation: note-the-significant interkernel-spread-in-the-albedo-and-cloud-feedbacks-relative-to-their-mean-values. Our results, therefore, demonstrate the

**Table 2.** The global annual mean feedback values (in  $\text{Wm}^{-2}\text{K}^{-1}$ ), calculated using each kernel and the same sample CESM1-LE data; from ~~From~~ left to right, they are the lapse rate~~feedback~~, Planck~~feedback~~, total (longwave + shortwave) water vapor~~feedback~~, surface albedo~~feedback~~, ~~and the shortwave~~, longwave ~~cloud~~, ~~shortwave cloud~~, and total cloud feedbacks. ~~The cloud feedbacks are~~ calculated using the adjustment method. The last two rows contain the kernel mean and standard deviation of the feedbacks.

kernel	Lapse Rate $\lambda_{LR}$	Planck $\lambda_{PL}$	$Q_{total}$ $\lambda_Q$	Albedo $\lambda_a$	Cloud <sub>SW,adj</sub> $\lambda_{CL,LW}$	Cloud <sub>LW,adj</sub> $\lambda_{CL,SW}$	Cloud <sub>total</sub>
BMRC	-0.41	-3.07	1.52	0.57	<del>0.31</del> 0.00		0.31
CAM3	-0.40	-2.99	1.48	0.32	<del>0.51</del> -0.07		<u>0.51</u>
CAM5	-0.43	-3.16	1.48	0.54	<del>0.28</del> 0.10		<u>0.28</u>
CERES	-0.42	-3.14	1.54	0.36	<del>0.27</del> 0.06		<u>0.27</u>
CloudSat	-0.42	-3.02	1.34	0.43	<del>0.39</del> 0.02		<u>0.39</u>
ECHAM6	-0.39	-3.07	1.37	0.41	<del>0.38</del> -0.01		<u>0.38</u>
ECMWF-RRTM	-0.38	-3.21	1.53	0.51	<del>0.30</del> -0.00		0.30
ERA5	-0.37	-3.18	1.51	0.52	<del>0.33</del> 0.01		<u>0.33</u>
GFDL	-0.41	-3.12	1.44	0.38	<del>0.38</del> 0.03		<u>0.38</u>
HadGEM2	-0.39	-3.35	1.59	0.49	<del>0.06</del> -0.03		<del>0.03</del> <u>0.39</u>
HadGEM3-GA7.1	-0.39	-3.17	1.50	0.41	<del>0.44</del> -0.01		<u>0.44</u>
mean	<b>-0.40</b>	<b>-3.13</b>	<b>1.48</b>	<b>0.45</b>	<b><del>0.33</del> 0.01</b>		<b><u>0.36</u></b>
stdSD	<b>0.02</b>	<b>0.09</b>	<b>0.07</b>	<b>0.08</b>	<b><del>0.11</del> 0.04</b>		<b><del>0.11</del><u>0.07</u></b>

~~large sensitivity of feedback~~. While the lapse rate and Planck feedbacks are consistently negative and the water vapor, surface albedo, and shortwave cloud feedbacks are consistently positive, the sign of the longwave cloud feedback is kernel-dependent, with values ranging from -0.07 to 0.10  $\text{Wm}^{-2}\text{K}^{-1}$ . Interkernel variability, as gathered from the standard deviation, is greatest in the Planck feedback, followed by the surface albedo, water vapor, and shortwave cloud feedbacks (Table 2).

It is worth comparing how global, annual mean water vapor feedback values depend on the calculation method, as outlined in Section 3.2. Table S1 shows the global annual mean all-sky water vapor feedbacks for each kernel and method. Method 3 (fractional approximation for the specific humidity response and normalization factor) yields the greatest water vapor feedback values for all kernels, followed by methods 4, 1, and 2. This occurs because the fractional approximation systematically overestimates the change in natural logarithms for large perturbations such as  $2\times\text{CO}_2$ , while the linear change ignores the damping by the logarithm entirely. The range in kernel-mean water vapor feedback values introduced by method choice is 0.26  $\text{Wm}^{-2}\text{K}^{-1}$ , with kernel mean values ranging from 1.41 to kernel choice, a factor often overlooked by many climate feedback studies 1.67  $\text{Wm}^{-2}\text{K}^{-1}$ .

Following this finding, a natural question is whether a particular kernel set does a better job at closing the radiative budget than others. This study only uses two simulations from one model for sample data and is, thus, ill-equipped to answer this question

more generally; however, we may use clear-sky linearity tests to identify which kernel best closes the budget in this particular CESM1 experiment, as in Shell et al. (2008). Consider the clear-sky radiative budget at some point in time after  $2\times\text{CO}_2$ :

485 
$$\Delta R^o = \Delta F^o + \lambda^o \Delta T + Re, \tag{11}$$

where  $\Delta R^o$  is the clear-sky TOA radiative imbalance,  $\Delta F^o$  is the clear-sky  $2\times\text{CO}_2$  ERF,  $\lambda^o$  is the clear-sky total feedback parameter, and  $Re$  is the residual term accounting for feedback nonlinearities, kernel errors, etc. All terms are global annual means. Kernel sets that produce the smallest magnitude  $Re$  values do the best job closing the clear-sky radiative budget. Table S2 lists the residual values for all 11 kernel sets and four water vapor feedback calculation methods. The best-performing  
490 kernel set depends heavily on the water vapor feedback calculation method; for example, the BMRC kernel set does the best job of closing the radiative budget using water vapor feedback method 1 (no fractional approximation) with a residual of  $-0.01 \text{ Wm}^{-2}\text{K}^{-1}$ , but is the second-to-last kernel set in terms of budget closure when using method 3 (with fractional approximation) with a residual of  $-0.25 \text{ Wm}^{-2}\text{K}^{-1}$  (Table S2). Focusing only on the first method, which is the most physically  
495 at closing the clear-sky radiative budget, respectively. This result is, perhaps, unsurprising given that CAM5 is the atmosphere model included in CESM1, and CAM3 is an earlier version of CAM5. We avoid making claims regarding these kernel sets' performance when applied to other models and simulations and simply assert that kernel choice is potentially important for accurately decomposing a model simulation's radiative budget.

## 5 Conclusions

500 The radiative kernel method is a popular and efficient way of diagnosing radiative feedbacks in climate model simulations. We were motivated to develop ClimKern to streamline these sometimes complicated calculations, shed light on kernel choice's importance in feedback studies, and provide access to a growing collection of existing kernels. We used ClimKern to compute basic radiative feedbacks from a sample climate model output to quantify kernel differences, leading us to the following conclusions.

505 **ClimKern makes radiative feedback calculations with kernels considerably easier while standardizing the underlying assumptions and methods.** The ClimKern ~~python~~ Python package contains straightforward, one-line commands for the most common calculations required for computing radiative feedbacks and can automatically load in data from the ClimKern data repository. The code is well-documented and easily accessible on GitHub and the Python Package Index for full transparency. Operations like vertical integration or horizontal regridding are consistent across the functions, even while offering the user  
510 different options. The repository similarly employs a standard and consistent nomenclature across all kernels, making it a practical resource for anyone wishing to compute radiative feedbacks.

**Kernel choice is a nonnegligible source of uncertainty in radiative feedback calculations, especially in the polar regions.**

In terms of global average feedbacks, the lapse rate ~~and Planck feedbacks appear~~ feedback appears to be the least sensitive to kernel choice. In contrast, the surface albedo and cloud feedbacks show considerably more sensitivity to the choice of kernel

515 (Table 2). Interkernel spread is horizontally and vertically inhomogeneous, with all but the water vapor feedback showing  
the greatest kernel sensitivity at the poles (Fig. 2); this spread may be a result of either differences in the base states or  
radiative schemes used to produce the radiative kernels. In the case of the The spread in the clear-sky temperature and surface  
albedo kernels , we note that the spread is not appreciably different between the is smaller than their all-sky and clear-sky  
versions counterparts, suggesting that clouds are ~~not the dominant cause of the spread across kernels~~ an important source of  
520 interkernel variability.

Polar amplification studies frequently use the radiative kernel method to compare surface warming contributions at the poles to  
the global or tropical average , and then rank the relative importance of the individual feedbacks (Pithan and Mauritsen, 2014;  
Stuecker et al., 2018; Previdi et al., 2020; Hahn et al., 2021; Janoski et al., 2023). ~~From the Arctic (>70°N) values in Fig. 3~~  
~~and the global average values in Table 2, the most important feedback contributing to Arctic amplification appears to be Two~~  
525 of the feedbacks most often identified as dominant polar amplification contributors, the lapse rate ~~or surface albedo feedback,~~  
~~depending on the kernel used. Therefore, kernel choice can affect the conclusions of polar amplification studies, and surface~~  
albedo feedbacks (Pithan and Mauritsen, 2014; Goosse et al., 2018; Previdi et al., 2020; Hahn et al., 2021; Previdi et al., 2021)  
, exhibit maximum interkernel variability in the Arctic (Fig. 2), further complicating comparisons between studies that use  
different radiative kernels.

530 Kernel choice can impact climate sensitivity studies that use global mean values. The mean Planck and surface albedo feedbacks  
show the greatest interkernel variability, while the sign of the mean longwave cloud feedback is kernel-dependent in this  
single CESM1 experiment. Although it is unclear how the variability introduced by kernel choice compares to that of other  
methodological choices, e.g., model or forcing scenario, it highlights its impact across spatial scales, leading to our final point.

**Future studies invoking the calculations of climate feedbacks can be more robust if they include a discussion of the sen-**  
535 **sitivity of the results to kernel choice.** One option would be to use multiple kernels from the ClimKern repository in a sensi-  
tivity analysis to explore this. Another option would be to ~~employ a kernel mean we provide here, instead of individual kernels,~~  
~~to limit the influence—and potential biases—of any one kernel. This option is particularly promising, given the consistent~~  
~~outperformance of multi-climate-model ensemble means over individual models in various metrics (Kharin and Zwiers, 2002; Tebaldi and~~  
take the kernel average of feedbacks instead of relying on individual kernels. Although a kernel average may dilute the benefits  
540 of using more advanced or better-performing kernels, it may reduce the sensitivity to individual kernel biases. This choice may  
be especially appropriate in studies using multimodel ensembles to avoid relying on a single kernel that may or may not align  
well with all models. Future work ~~may will~~ include comparing the sensitivity to kernel choice to other sources of uncertainty  
in climate studies and evaluating kernel mean performance compared to individual kernels in the computation of radiative  
feedbacks.

545 We intend for ClimKern to become a community-wide project and invite potential collaborators to contribute. The easiest way  
is to visit the ClimKern GitHub and fork the repository. New features and bug fixes can also be requested there.

*Code and data availability.* The ClimKern kernel and data repository is located at <https://zenodo.org/records/14743752>. The version of the ClimKern Python package documented in this work can be found at <https://zenodo.org/records/14743210>. Those interested in contributing to ClimKern or wishing to use the latest version should instead navigate to <https://github.com/tyfolino/climkern>. The Jupyter notebook containing the code to produce our figures is located at <https://zenodo.org/records/14757435> (Janoski, 2025).

*Author contributions.* TPJ and IM led the project development. TPJ performed most of the coding, conducted all data analysis, and generated the figures. TPJ also wrote most of the manuscript, with contributions from IM, who helped write Sections 1 and 2.2, co-conceptualized the project, and assisted in planning the manuscript and figures. IM also contributed to the package's development with some minor coding. RJK's earlier work provided foundational insights, and he, along with MP, offered technical support throughout the project. LMP contributed by offering broader guidance and helped refine the manuscript. All authors contributed to editing the manuscript.

*Competing interests.* We declare no competing interests.

*Acknowledgements.* TPJ is supported, in part, by a National Oceanic and Atmospheric Administration (NOAA) grant via the NOAA Center for Earth System Sciences and Remote Sensing Technologies. IM is supported by a Harry Hess post-doctoral fellowship from Princeton Geosciences. LMP and MP are supported, in part, by grants from the US National Science Foundation to Columbia University. We gratefully acknowledge the creators of the kernels used in this work for making their kernels available and the developers of Xarray and xESMF for creating software to make this work possible. We also thank Tyler Thorsen for providing the latest version of the CERES kernels. TPJ thanks the US Research Software Sustainability Institute and associated instructors for teaching him how to create a Python package. All authors thank Mark Zelinka and Max Coleman for their thoughtful reviews of this manuscript.



## References

- 565 Andrews, T., Gregory, J. M., Webb, M. J., and Taylor, K. E.: Forcing, feedbacks and climate sensitivity in CMIP5 coupled atmosphere-ocean climate models, *Geophysical research letters*, 39, 2012.
- Bani Shahabadi, M. and Huang, Y.: Logarithmic radiative effect of water vapor and spectral kernels, *Journal of Geophysical Research: Atmospheres*, 119, 6000–6008, 2014.
- Bellucci, A., Haarsma, R., Gualdi, S., Athanasiadis, P., Caian, M., Cassou, C., Fernandez, E., Germe, A., Jungclaus, J., Kröger, J., et al.: An  
570 assessment of a multi-model ensemble of decadal climate predictions, *Climate Dynamics*, 44, 2787–2806, 2015.
- Block, K. and Mauritsen, T.: Forcing and feedback in the MPI-ESM-LR coupled model under abruptly quadrupled CO<sub>2</sub>, *Journal of Advances in Modeling Earth Systems*, 5, 676–691, 2013.
- Bonan, D. B., Kay, J. E., Feldl, N., and Zelinka, M. D.: Mid-latitude clouds contribute to Arctic amplification via interactions with other climate feedbacks, *Environmental Research: Climate*, 2025.
- 575 Bony, S., Colman, R., Kattsov, V. M., Allan, R. P., Bretherton, C. S., Dufresne, J.-L., Hall, A., Hallegatte, S., Holland, M. M., Ingram, W., et al.: How well do we understand and evaluate climate change feedback processes?, *Journal of Climate*, 19, 3445–3482, 2006.
- Buck, A. L.: New equations for computing vapor pressure and enhancement factor, *Journal of Applied Meteorology (1962-1982)*, pp. 1527–1532, 1981.
- Colman, R. and McAvaney, B.: On tropospheric adjustment to forcing and climate feedbacks, *Climate dynamics*, 36, 1649–1658, 2011.
- 580 Colman, R. and Soden, B. J.: Water vapor and lapse rate feedbacks in the climate system, *Reviews of Modern Physics*, 93, 045 002, 2021.
- Croll, J.: Climate and time, *Nature*, 12, 329–329, 1875.
- Curry, J. A., Schramm, J. L., and Ebert, E. E.: Sea ice-albedo climate feedback mechanism, *Journal of Climate*, 8, 240–247, 1995.
- Dong, Y., Proistosescu, C., Armour, K. C., and Battisti, D. S.: Attributing Historical and Future Evolution of Radiative Feedbacks to Regional Warming Patterns using a Green’s Function Approach: The Preeminence of the Western Pacific, *Journal of Climate*, 32, 5471 – 5491,  
585 <https://doi.org/10.1175/JCLI-D-18-0843.1>, 2019.
- Edwards, J. and Slingo, A.: Studies with a flexible new radiation code. I: Choosing a configuration for a large-scale model, *Quarterly Journal of the Royal Meteorological Society*, 122, 689–719, 1996.
- Eyring, V., Bony, S., Meehl, G. A., Senior, C. A., Stevens, B., Stouffer, R. J., and Taylor, K. E.: Overview of the Coupled Model Intercomparison Project Phase 6 (CMIP6) experimental design and organization, *Geoscientific Model Development*, 9, 1937–1958, 2016.
- 590 Feldl, N. and Roe, G. H.: The nonlinear and nonlocal nature of climate feedbacks, *Journal of Climate*, 26, 8289–8304, 2013.
- Feldl, N., Bordoni, S., and Merlis, T. M.: Coupled high-latitude climate feedbacks and their impact on atmospheric heat transport, *Journal of Climate*, 30, 189–201, 2017.
- Feldl, N., Po-Chedley, S., Singh, H. K., Hay, S., and Kushner, P. J.: Sea ice and atmospheric circulation shape the high-latitude lapse rate feedback, *NPJ climate and atmospheric science*, 3, 41, 2020.
- 595 Forster, P., Richardson, T., Maycock, A. C., Smith, C. J., Samset, B. H., Myhre, G., Andrews, T., Pincus, R., and Schulz, M.: Recommendations for diagnosing effective radiative forcing from climate models for CMIP6, *Journal of Geophysical Research: Atmospheres*, 121, 12,460–12,475, <https://doi.org/https://doi.org/10.1002/2016JD025320>, 2016.
- Goosse, H., Kay, J. E., Armour, K. C., Bodas-Salcedo, A., Chepfer, H., Docquier, D., Jonko, A., Kushner, P. J., Lecomte, O., Massonnet, F., et al.: Quantifying climate feedbacks in polar regions, *Nature communications*, 9, 1919, 2018.
- 600 Graverson, R. G., Langen, P. L., and Mauritsen, T.: Polar amplification in CCSM4: Contributions from the lapse rate and surface albedo feedbacks, *Journal of Climate*, 27, 4433–4450, 2014.

- Gregory, J. M. and Andrews, T.: Variation in climate sensitivity and feedback parameters during the historical period, *Geophysical Research Letters*, 43, 3911–3920, 2016.
- Hahn, L. C., Armour, K. C., Zelinka, M. D., Bitz, C. M., and Donohoe, A.: Contributions to polar amplification in CMIP5 and CMIP6 models, *Frontiers in Earth Science*, 9, 710036, 2021.
- 605 Held, I. M. and Shell, K. M.: Using relative humidity as a state variable in climate feedback analysis, *Journal of Climate*, 25, 2578–2582, 2012.
- Hoyer, S. and Hamman, J.: xarray: ND labeled arrays and datasets in Python, *Journal of Open Research Software*, 5, 10–10, 2017.
- Huang, H. and Huang, Y.: Radiative sensitivity quantified by a new set of radiation flux kernels based on the ECMWF Reanalysis v5 (ERA5), *Earth System Science Data*, 15, 3001–3021, 2023.
- 610 Huang, H., Huang, Y., Wei, Q., and Hu, Y.: Band-by-band spectral radiative kernels based on the ERA5 reanalysis, *Scientific Data*, 11, 237, 2024.
- Huang, Y., Ramaswamy, V., and Soden, B.: An investigation of the sensitivity of the clear-sky outgoing longwave radiation to atmospheric temperature and water vapor, *Journal of Geophysical Research: Atmospheres*, 112, 2007.
- 615 Huang, Y., Xia, Y., and Tan, X.: On the pattern of CO<sub>2</sub> radiative forcing and poleward energy transport, *Journal of Geophysical Research: Atmospheres*, 122, 10–578, 2017.
- Huang, Y., Huang, H., and Shakirova, A.: The nonlinear radiative feedback effects in the Arctic warming, *Frontiers in Earth Science*, 9, 693779, 2021.
- Ingram, W., Wilson, C., and Mitchell, J.: Modeling climate change: An assessment of sea ice and surface albedo feedbacks, *Journal of Geophysical Research: Atmospheres*, 94, 8609–8622, 1989.
- 620 Janoski, T.: climkern-analysis-and-plots: v1.1.0, <https://doi.org/10.5281/zenodo.14757435>, 2025.
- Janoski, T., Mitevski, I., and Kramer, R.: ClimKern Kernel & Data Repository, <https://doi.org/10.5281/zenodo.13287114>, 2024a.
- Janoski, T. P., Previdi, M., Chiodo, G., Smith, K. L., and Polvani, L. M.: Ultrafast Arctic amplification and its governing mechanisms, *Environmental Research: Climate*, 2, 035009, 2023.
- 625 Janoski, T. P., Mitevski, I., , and Wen, K.: ClimKern, <https://doi.org/10.5281/zenodo.10291284>, 2024b.
- Kay, J. E., Deser, C., Phillips, A., Mai, A., Hannay, C., Strand, G., Arblaster, J. M., Bates, S. C., Danabasoglu, G., Edwards, J., Holland, M., Kushner, P., Lamarque, J.-F., Lawrence, D., Lindsay, K., Middleton, A., Munoz, E., Neale, R., Oleson, K., Polvani, L., and Vertenstein, M.: The Community Earth System Model (CESM) Large Ensemble Project: A Community Resource for Studying Climate Change in the Presence of Internal Climate Variability, *Bulletin of the American Meteorological Society*, 96, 1333–1349, [https://doi.org/10.1175/BAMS-](https://doi.org/10.1175/BAMS-D-13-00255.1)
- 630 D-13-00255.1, 2015.
- Kharin, V. V. and Zwiers, F. W.: Climate predictions with multimodel ensembles, *Journal of Climate*, 15, 793–799, 2002.
- Knutti, R. and Rugenstein, M. A.: Feedbacks, climate sensitivity and the limits of linear models, *Philosophical Transactions of the Royal Society A: Mathematical, Physical and Engineering Sciences*, 373, 20150146, 2015.
- Kramer, R. J., Matus, A. V., Soden, B. J., and L’Ecuyer, T. S.: Observation-based radiative kernels from CloudSat/CALIPSO, *Journal of Geophysical Research: Atmospheres*, 124, 5431–5444, 2019.
- 635 Lacis, A. A., Hansen, J. E., Russell, G. L., Oinas, V., and Jonas, J.: The role of long-lived greenhouse gases as principal LW control knob that governs the global surface temperature for past and future climate change, *Tellus B: Chemical and Physical Meteorology*, 65, 19734, 2013.

- Laîné, A., Yoshimori, M., and Abe-Ouchi, A.: Surface Arctic amplification factors in CMIP5 models: land and oceanic surfaces and season-  
640 ality, *Journal of Climate*, 29, 3297–3316, 2016.
- Larson, E. J. and Portmann, R. W.: A temporal kernel method to compute effective radiative forcing in CMIP5 transient simulations, *Journal of Climate*, 29, 1497–1509, 2016.
- Lin, J.-L.: The double-ITCZ problem in IPCC AR4 coupled GCMs: Ocean–atmosphere feedback analysis, *Journal of Climate*, 20, 4497–4525, 2007.
- 645 Manabe, S. and Wetherald, R. T.: The effects of doubling the CO<sub>2</sub> concentration on the climate of a general circulation model, *Journal of Atmospheric Sciences*, 32, 3–15, 1975.
- Manners, J.: Socrates technical guide suite of community radiative transfer codes based on edwards and slingo, Exeter, UK, 2015.
- Meyssignac, B., Chenal, J., Loeb, N., Guillaume-Castel, R., and Ribes, A.: Time-variations of the climate feedback parameter  $\lambda$  are associated with the Pacific Decadal Oscillation, *Communications Earth & Environment*, 4, 241, 2023.
- 650 Mitevski, I., Orbe, C., Chemke, R., Nazarenko, L., and Polvani, L. M.: Non-Monotonic Response of the Climate System to Abrupt CO<sub>2</sub> Forcing, *Geophysical Research Letters*, 48, e2020GL090861, <https://doi.org/10.1029/2020GL090861>, 2021.
- Mitevski, I., Polvani, L. M., and Orbe, C.: Asymmetric Warming/Cooling Response to CO<sub>2</sub> Increase/Decrease Mainly Due To Non-Logarithmic Forcing, Not Feedbacks, *Geophysical Research Letters*, 49, e2021GL097133, <https://doi.org/10.1029/2021GL097133>, 2022.
- Mitevski, I., Dong, Y., Polvani, L. M., Rugenstein, M., and Orbe, C.: Non-Monotonic Feedback Dependence Under Abrupt CO<sub>2</sub> Forcing Due  
655 To a North Atlantic Pattern Effect, *Geophysical Research Letters*, 50, e2023GL103617, <https://doi.org/10.1029/2023GL103617>, 2023.
- Pendergrass, A. G.: apendergrass/cam5-kernels: Up to date codebase as of August 2019, <https://doi.org/10.5281/zenodo.3359041>, 2019.
- Pendergrass, A. G., Conley, A., and Vitt, F. M.: Surface and top-of-atmosphere radiative feedback kernels for CESM-CAM5, *Earth System Science Data*, 10, 317–324, 2018.
- Pincus, R., Buehler, S. A., Brath, M., Crevoisier, C., Jamil, O., Franklin Evans, K., Manners, J., Menzel, R. L., Mlawer, E. J., Paynter, D., et al.:  
660 Benchmark calculations of radiative forcing by greenhouse gases, *Journal of Geophysical Research: Atmospheres*, 125, e2020JD033483, 2020.
- Pithan, F. and Mauritsen, T.: Arctic amplification dominated by temperature feedbacks in contemporary climate models, *Nature geoscience*, 7, 181–184, 2014.
- Previdi, M., Janoski, T. P., Chiodo, G., Smith, K. L., and Polvani, L. M.: Arctic amplification: A rapid response to radiative forcing, *Geophysical Research Letters*, 47, e2020GL089933, 2020.
- 665 Previdi, M., Smith, K. L., and Polvani, L. M.: Arctic amplification of climate change: a review of underlying mechanisms, *Environmental Research Letters*, 16, 093003, 2021.
- Riihelä, A., Bright, R. M., and Anttila, K.: Recent strengthening of snow and ice albedo feedback driven by Antarctic sea-ice loss, *Nature Geoscience*, 14, 832–836, 2021.
- 670 Roe, G. H. and Baker, M. B.: Why is climate sensitivity so unpredictable?, *Science*, 318, 629–632, 2007.
- Shell, K. M., Kiehl, J. T., and Shields, C. A.: Using the radiative kernel technique to calculate climate feedbacks in NCAR’s Community Atmospheric Model, *Journal of Climate*, 21, 2269–2282, 2008.
- Sherwood, S. C., Webb, M. J., Annan, J. D., Armour, K. C., Forster, P. M., Hargreaves, J. C., Hegerl, G., Klein, S. A., Marvel, K. D., Rohling, E. J., Watanabe, M., Andrews, T., Braconnot, P., Bretherton, C. S., Foster, G. L., Hausfather, Z., Heydt, A. S. v. d., Knutti, R., Mauritsen, T., Norris, J. R., Proistosescu, C., Rugenstein, M., Schmidt, G. A., Tokarska, K. B., and Zelinka, M. D.: An assessment of Earth’s climate  
675 sensitivity using multiple lines of evidence, *Reviews of Geophysics*, <https://doi.org/https://doi.org/10.1029/2019RG000678>, 2020.

- Shi, J. and Lohmann, G.: Considerable uncertainty of simulated Arctic temperature change in the mid-Holocene due to initial ocean perturbation, *Geophysical Research Letters*, 51, e2023GL106 337, 2024.
- Smith, C., Kramer, R., Myhre, G., Forster, P., Soden, B., Andrews, T., Boucher, O., Faluvegi, G., Fläschner, D., Hodnebrog, Ø., et al.: Understanding rapid adjustments to diverse forcing agents, *Geophysical Research Letters*, 45, 12–023, 2018.
- Smith, C. J., Kramer, R. J., and Sima, A.: The HadGEM3-GA7. 1 radiative kernel: the importance of a well-resolved stratosphere, *Earth System Science Data*, 12, 2157–2168, 2020.
- Soden, B. J. and Held, I. M.: An assessment of climate feedbacks in coupled ocean–atmosphere models, *Journal of climate*, 19, 3354–3360, 2006.
- Soden, B. J., Held, I. M., Colman, R., Shell, K. M., Kiehl, J. T., and Shields, C. A.: Quantifying climate feedbacks using radiative kernels, *Journal of Climate*, 21, 3504–3520, 2008.
- Stuecker, M. F., Bitz, C. M., Armour, K. C., Proistosescu, C., Kang, S. M., Xie, S.-P., Kim, D., McGregor, S., Zhang, W., Zhao, S., et al.: Polar amplification dominated by local forcing and feedbacks, *Nature Climate Change*, 8, 1076–1081, 2018.
- Tebaldi, C. and Knutti, R.: The use of the multi-model ensemble in probabilistic climate projections, *Philosophical transactions of the royal society A: mathematical, physical and engineering sciences*, 365, 2053–2075, 2007.
- Thorsen, T. J., Kato, S., Loeb, N. G., and Rose, F. G.: Observation-based decomposition of radiative perturbations and radiative kernels, *Journal of climate*, 31, 10 039–10 058, 2018.
- Tian, B. and Dong, X.: The double-ITCZ bias in CMIP3, CMIP5, and CMIP6 models based on annual mean precipitation, *Geophysical Research Letters*, 47, e2020GL087 232, 2020.
- Vial, J., Dufresne, J.-L., and Bony, S.: On the interpretation of inter-model spread in CMIP5 climate sensitivity estimates, *Climate Dynamics*, 41, 3339–3362, 2013.
- Wetherald, R. and Manabe, S.: Cloud feedback processes in a general circulation model, *Journal of the Atmospheric Sciences*, 45, 1397–1416, 1988.
- Zelinka, M. D., Klein, S. A., and Hartmann, D. L.: Computing and partitioning cloud feedbacks using cloud property histograms. Part II: Attribution to changes in cloud amount, altitude, and optical depth, *Journal of Climate*, 25, 3736–3754, 2012.
- Zelinka, M. D., Myers, T. A., McCoy, D. T., Po-Chedley, S., Caldwell, P. M., Ceppi, P., Klein, S. A., and Taylor, K. E.: Causes of higher climate sensitivity in CMIP6 models, *Geophysical Research Letters*, 47, e2019GL085 782, 2020.
- Zhang, R., Wang, H., Fu, Q., Pendergrass, A. G., Wang, M., Yang, Y., Ma, P.-L., and Rasch, P. J.: Local radiative feedbacks over the Arctic based on observed short-term climate variations, *Geophysical Research Letters*, 45, 5761–5770, 2018.
- Zhang, Y.: Introduction to the ISCCP-FH Radiative Kernels, [https://data.giss.nasa.gov/isccp/kernels/docs/Zhang\\_Introduction\\_to\\_ISCCP-FH\\_radiative-kernels\\_v0\\_202309.pdf](https://data.giss.nasa.gov/isccp/kernels/docs/Zhang_Introduction_to_ISCCP-FH_radiative-kernels_v0_202309.pdf), 2023.
- Zhu, T., Huang, Y., and Wei, H.: Estimating climate feedbacks using a neural network, *Journal of Geophysical Research: Atmospheres*, 124, 3246–3258, 2019.
- Zhuang, J., dussin, r., Huard, D., Bourgault, P., Banihirwe, A., Raynaud, S., Malevich, B., Schupfner, M., Filipe, Levang, S., Gauthier, C., Jüling, A., Almansi, M., RichardScottOZ, RondeauG, Rasp, S., Smith, T. J., Stachelek, J., Plough, M., Pierre, Bell, R., Caneill, R., and Li, X.: pangeo-data/xESMF: v0.8.2, <https://doi.org/10.5281/zenodo.8356796>, 2023.