

We thank this Reviewer for thoughtful and constructive comments on our manuscript. We appreciate the time s/he invested in the review. We believe that our revised manuscript addresses all the comments. In this regard, we have revised and rewritten a few sections such as Abstract, method, results, discussion and conclusion in the revised manuscript. Quality control and revision of dissipation estimates, as well as current calculations were also done to ensure the validity of our results. We thought it useful to point out its detailed revisions (lines and sections) in the reply to your comments. Below (highlighted in blue and magenta) is an itemized response to the different issues raised in the review.

## Review

Turbulent dissipation from AMAZOMIX off the Amazon shelf along internal tides paths

Kouogang et al.

## Summary

This manuscript describes measurements of currents, hydrography, and turbulence at numerous sites near the Amazon outflow, where internal tides are also generated at the slope. The dataset appears suited to study the questions posed by the authors. Do internal tides affect the mixing? Does the mixing affect nutrient fluxes?

However, I have some concerns about the analysis listed below. Since many of the details of the analysis are not presented, it is difficult for me to evaluate what has been done. Some of the more prominent ones include:

tidal decomposition of the currents,

use/decomposition of shear,

description/applicability of the finescale parameterizations used,

and shaded color plots of SADCPC velocity that include transits and times on station.

There are also presentation/stylistic problems that make the manuscript very difficult to read. Many of them are major, but also getting down to such points as the labelling of figures and sections. Some examples are listed below. It is well beyond the scope of a review to point all of these out. Furthermore, the cumulative difficulty of numerous minor things (let alone the more major ones) turns into a major distraction to the detriment of the hard work done and the science. The senior authors must make considerably more effort in this regard.

## Major comments

1. SI units should be written  $W\text{ kg}^{-1}$  not  $W.kg^{-1}$  and 90 km not 90km and so on.

R: Thanks for your remarks.

We have revised the SI units throughout the revised manuscript.

an example, in lines 22-24 of the revised manuscript, is shown below:

“

Near generation sites, mixing rates were elevated, between  $[10^{-6}, 10^{-5}] W\text{ kg}^{-1}$ , with IT shear contributing ~65 %, compared to mean baroclinic current (BC) shear. Along IT pathways and in far-field IT regions, mixing decreased to  $[10^{-8}, 10^{-7}] W\text{ kg}^{-1}$  but remained substantial, driven by nearly equal contributions from IT and BC shear.

“

2. 50.4% on line 29. Please state the error and adjust the significant figures accordingly.

R: Thanks for your remarks.

We have revised as suggested

The error, associated with the contribution of tidal and mean shear, was estimated with the standard deviation. It is relatively small. you will find the error we have declared in table A1 of the revised manuscript.

3. line 34 - A “guide for the mixing parameterization in future numerical simulation” is mentioned. What is the guidance?

R: Thanks for your remark.

Indeed, we have revised all the Abstract and sentence in line 29-30 of the revised manuscript:

“

These findings provide valuable insights for developing parameterizations of tidal and mean shear mixing for ocean general circulation models and coupled models, including biogeochemical and atmospheric systems.

”

4. Section 2.2- some of the microstructure processing steps should be shown in either the main text, appendix, or supplementary material. In particular, I would like to see some examples of the spectral fits.

R: Thanks for your suggestions.

We have incorporated additional microstructure processing steps and spectral fits into both the main text and the appendix. Revisions and checks of the VMP profile processing, and dissipation estimates have been carried out, along with quality assurance tests, in alignment with all reviewers' comments and ATOMIX's recommendations (Lueck et al., 2024).

a) First, the revisions were focused on the VMP profile processing, particularly on:

- Parameters controlling shear spectra estimation, such as record lengths (L), which are cosine-windowed and overlapped (O) by 50%.
- Parameter for extracting the section or profile (continuous part of the time series), including the minimum depth of extraction (Pmin).

For shallow stations, we used  $L = 5$  s and  $O = 2.5$  s, instead of the previous  $L = 4$  s and  $O = 2$  s. Parameters for deep stations remain unchanged ( $L = 8$  s and  $O = 4$  s).

Additionally, we adjusted the minimum depth of extraction to  $P_{min} = 10$  m for deep stations and  $P_{min} = 3$  m for shallow stations, compared to the previously used value of  $P_{min} = 1$  m for both station types.

b) Second, the revisions and checks focused on the quality assurance measures for dissipation estimates, with quality control checks and adjustments applied across all stations (e.g., S2, S7, S12, S14).

For instance, at station S14, previous dissipation estimates showed some peaks at various depths, particularly between 100–200 m, 300–400 m, and 500–700 m. While the fraction of shear data affected by despiking during processing was  $<0.05$ , the figure of merit (FM)—used to filter out poor-quality data—for shear probe 1 was  $\gg 1.4$  at certain depths (e.g., around 327 m compared to 122 m, as shown in figure RC1.1 below). In contrast, the FM for shear probe 2 remained  $<1.4$ .

After checks, quality control of dissipation estimates have been revised for all stations (e.g., S6, S10, S14). We have retained only dissipation estimates from either one or both probes that met the quality assurance criteria ( $FM < 1.4$  and fraction of despiked shear data  $<0.05$ , as recommended by ATOMIX), as shown at S6 and S10 for example (figure RC1.2 below).

The final dissipation estimate was computed as the average of the estimates from both shear probes, followed by the mean of the dissipation profiles for the station, as illustrated in figure RC1.3 below.

The revision of section “methods/TKE dissipation rates” can be found in lines 121-151 of the revised manuscript, as shown below:

“

## **2.2 Methods**

### **TKE dissipation rates**

The VMP data are processed using ODAS Matlab library (developed by Rockland Scientific International, Inc) to infer the TKE dissipation rate ( $\epsilon$ ). The processing methods for the VMP data are briefly described here and adhere to the recommendations of ATOMIX (Analyzing ocean turbulence observations to quantify mixing), as reported by Lueck et al. (2024), and have been validated against the benchmark estimates (presented in Fer et al., 2024).

First, the VMP data are converted into physical shear units, and the time series are prepared. Continuous sections of the time series are selected for dissipation estimation. Before spectral estimation, the aberrant shear signals caused by vessel wake contamination are removed. Collisions of the shear probe with plankton and other particles are removed using the de-spiking routine. The records from each section are then high-pass filtered (e.g., at station S6 and S10; Fig. 2a, and Fig. A1, Appendix).

Shear spectra are estimated using record lengths ( $L$ ) and Fast Fourier Transform segments of 2 s, which are cosine windowed and overlapped by 50% (e.g., at station S6 and S10; Fig. 2b, and Fig. A1, Appendix). Additionally, vibration-coherent noise is removed. Different  $L$  and overlap ( $O$ ) settings were selected and tested based on the environment (e.g., deep vs. shallow water), following Fer et al. (2024). For shallow stations,  $L$  ( $O$ ) was shortened to 5 s (2.5 s), in contrast to the 8 s (4 s) used for deeper stations, due to evidence of overturns observed in AMAZOMIX acoustic measurements at deeper stations (Koch-Larrouy et al., 2024; in preparation). This adjustment helped to optimize the spatial resolution of dissipation estimates in shallow water stations.

Finally,  $\epsilon$  is determined using the spectral integration method and by comparison with the Nasmyth empirical spectrum (Nasmyth, 1970). Quality assurance tests are carried out in accordance with ATOMIX's recommendations (Lueck et al., 2024). A figure of merit  $< 1.4$  is used to exclude bad data (e.g., at station S6 and S10; Fig. 2b, and Fig. A1, Appendix), and the fraction of data affected by de-spiking is  $< 0.05$ .

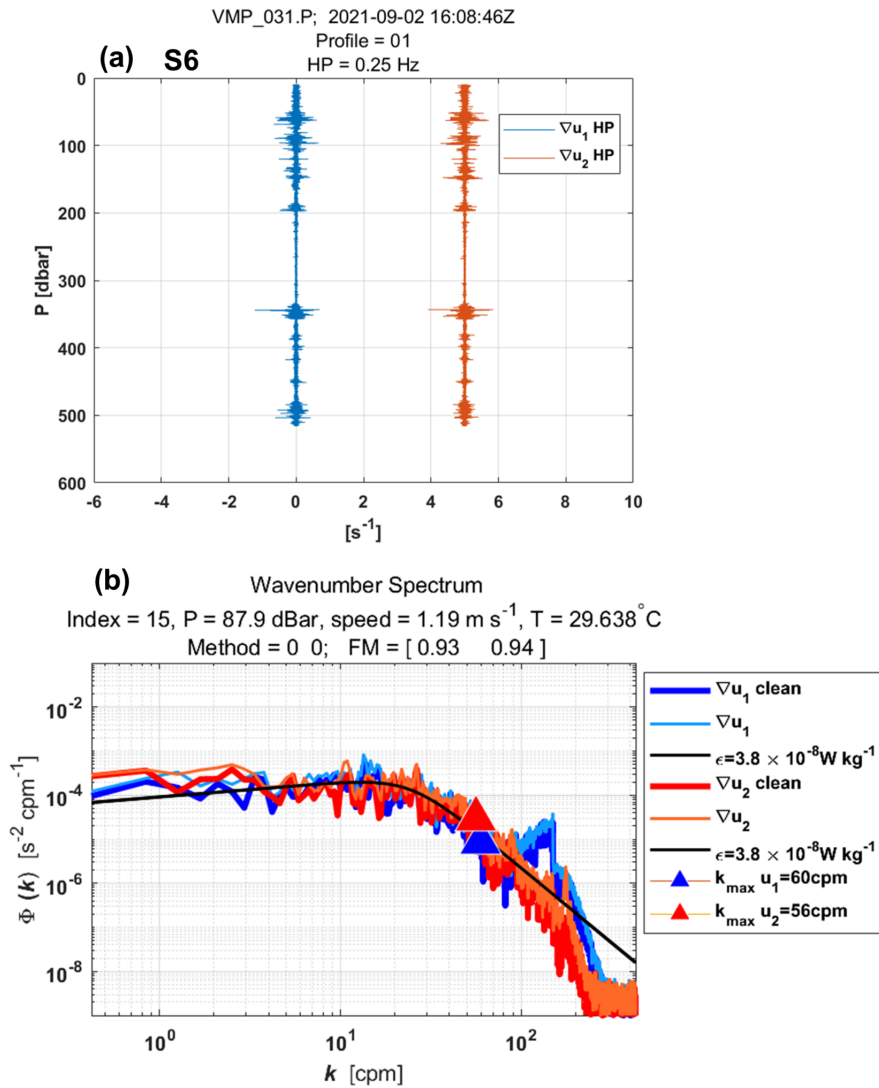


Figure 2: Example of wavenumber spectra from a dissipation structure segment used to determine the dissipation rate at station S6 at a pressure of 87.9 dBar. (a) Cleaned and high-pass filtered signals from shear probe 1 (blue) and shear probe 2 (red, offset by 5 s<sup>-1</sup>). (b) Wavenumber spectra for shear probes 1 and 2. Thick lines (blue for probe 1, red for probe 2) show shear spectra with coherent noise correction, while thin lines (sky blue for probe 1, orange for probe 2) show spectra without correction. Triangles mark the maximum wavenumber used for dissipation rate estimation. Black lines represent Nasmyth reference spectra for estimated dissipation rate of  $3.8 \times 10^{-8} \text{ W kg}^{-1}$  for both shear probes. Dissipation rate estimates for shear probe 1 and shear probe 2 at a pressure of 87.9 dBar yielded a figure of merit of 0.93 and 0.94, respectively.

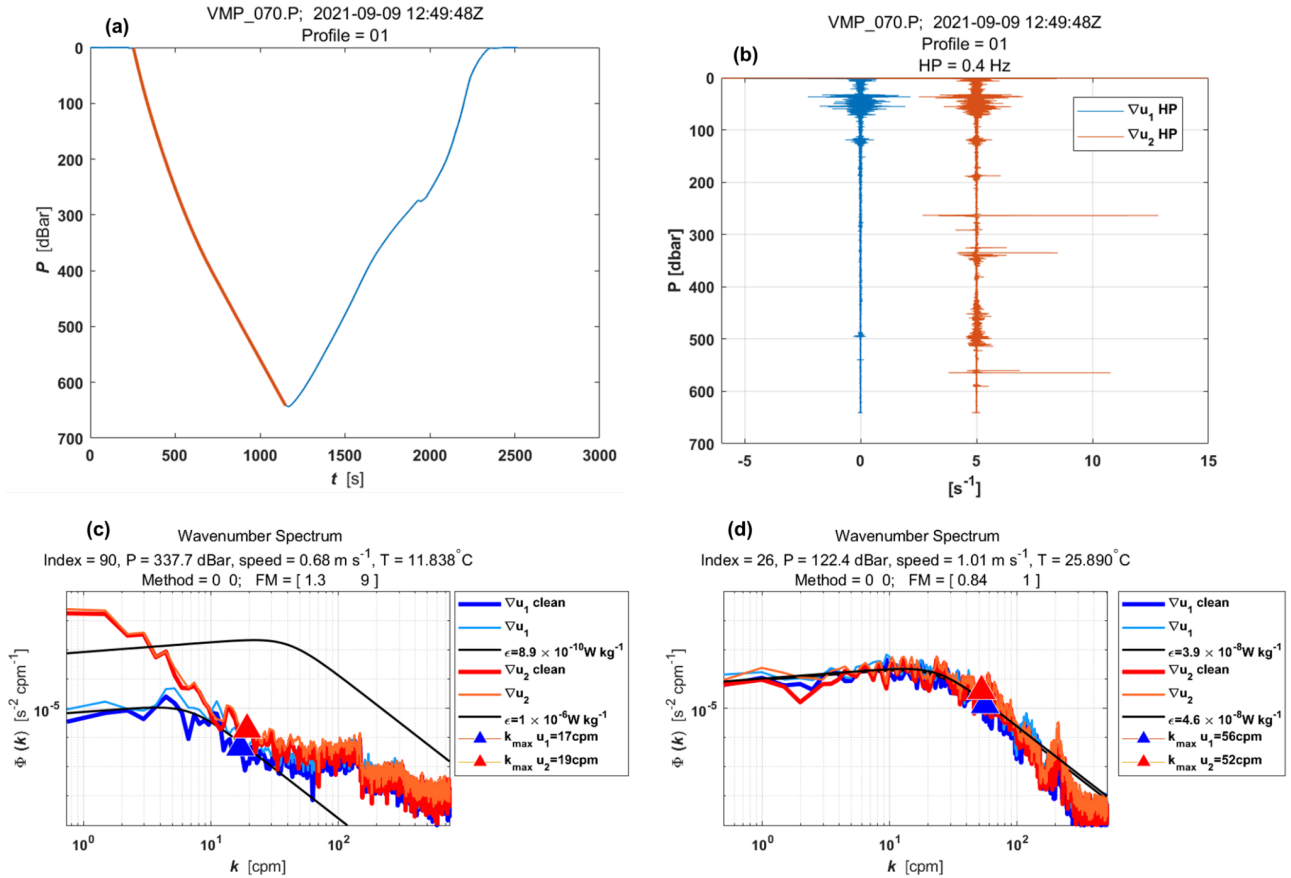


Figure RC1.1: Example of wavenumber spectra from a dissipation structure segment used to determine the dissipation rate at station S14 at a pressure of 337.7 dBar. (a) Pressure record for the entire data file (blue) and the specific segment being analyzed (red). (b) Cleaned and high-pass filtered signals from shear probe 1 (blue) and shear probe 2 (red, offset by  $5 \text{ s}^{-1}$ ). (c) Wavenumber spectra for shear probes 1 and 2. Thick lines (blue for probe 1, red for probe 2) show shear spectra with coherent noise correction, while thin lines (sky blue for probe 1, orange for probe 2) show spectra without correction. Triangles mark the maximum wavenumber used for dissipation rate estimation. Black lines represent Nasmyth reference spectra for estimated dissipation rate of  $8.9 \times 10^{-10} \text{ W kg}^{-1}$  and  $1 \times 10^{-6} \text{ W kg}^{-1}$  for shear probes 1 and 2, respectively. Dissipation rate estimates for shear probes 1 and shear probe 2 at a pressure of 337.7 dBar yielded a figure of merit of 1.3 and 9, respectively. Panel (d) is similar to panel (c) but:- with Nasmyth reference spectra for estimated dissipation rate of  $3.9 \times 10^{-8} \text{ W kg}^{-1}$  and  $4.6 \times 10^{-8} \text{ W kg}^{-1}$  for shear probes 1 and 2. -wth dissipation rate estimates for shear probes 1 and 2 at a pressure of 122.4 dBar yielding a figure of merit of 0.84 and 1, respectively.

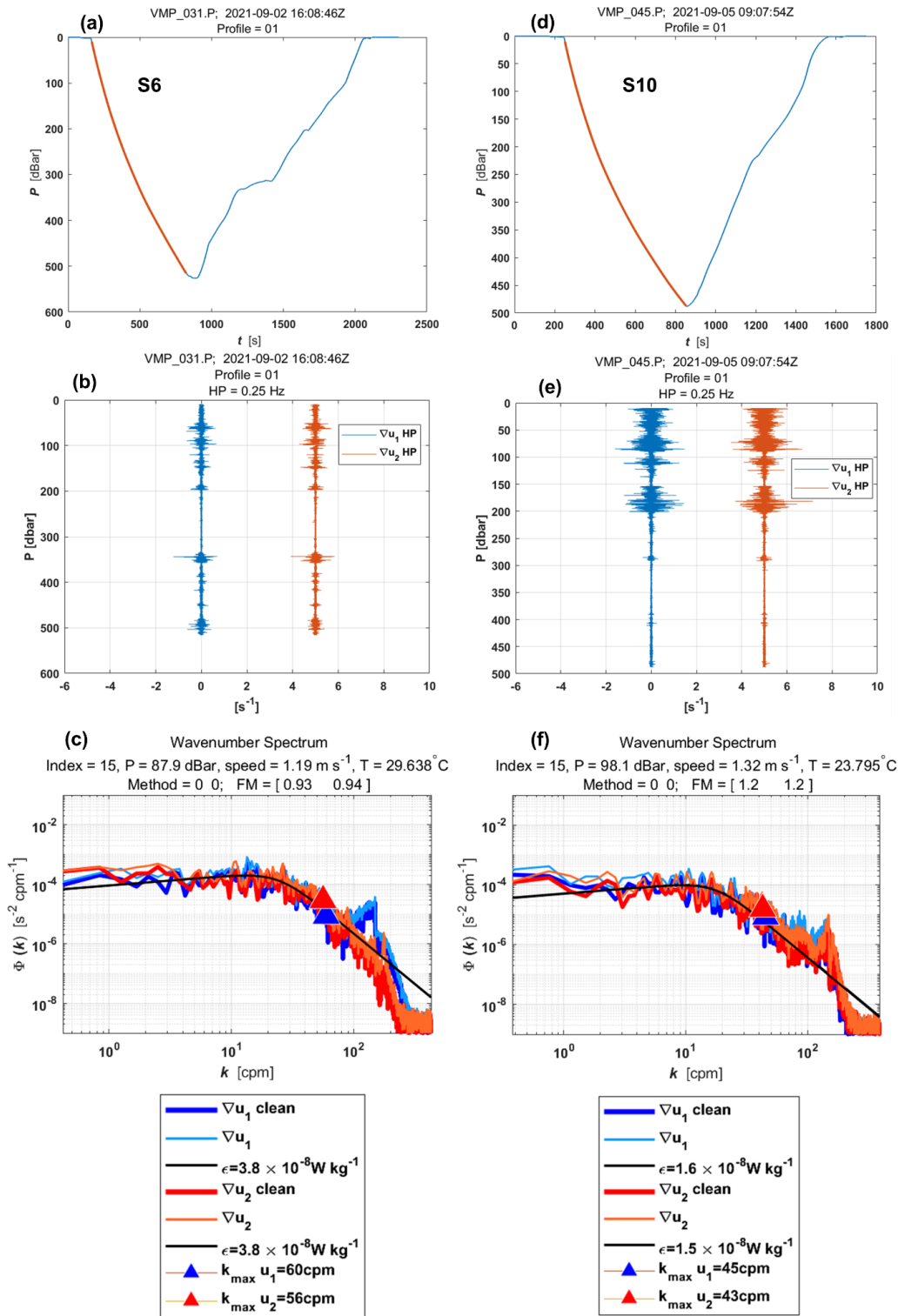


Figure RC1.2: Similar to Fig. RC1.1. but for stations (a)-(b)-(c) S6 and (d)-(e)-(f) S10. For S6 (panels c), Black lines represent Nasmyth reference spectra for estimated dissipation rate of  $3.8 \times 10^{-8} \text{ W kg}^{-1}$  for both shear probes, and dissipation rate estimates for shear probes 1 and shear probe 2 at a pressure of 337.7 dBar yielded a figure of merit of 0.93 and 0.94, respectively. For S10 (panel f), Black lines represent Nasmyth reference spectra for estimated dissipation rate of  $1.6 \times 10^{-8} \text{ W kg}^{-1}$  and  $1.5 \times 10^{-8} \text{ W kg}^{-1}$  for shear probes 1 and 2, respectively, and dissipation rate estimates for both shear probes at a pressure of 337.7 dBar yielded a figure of merit of 1.2.

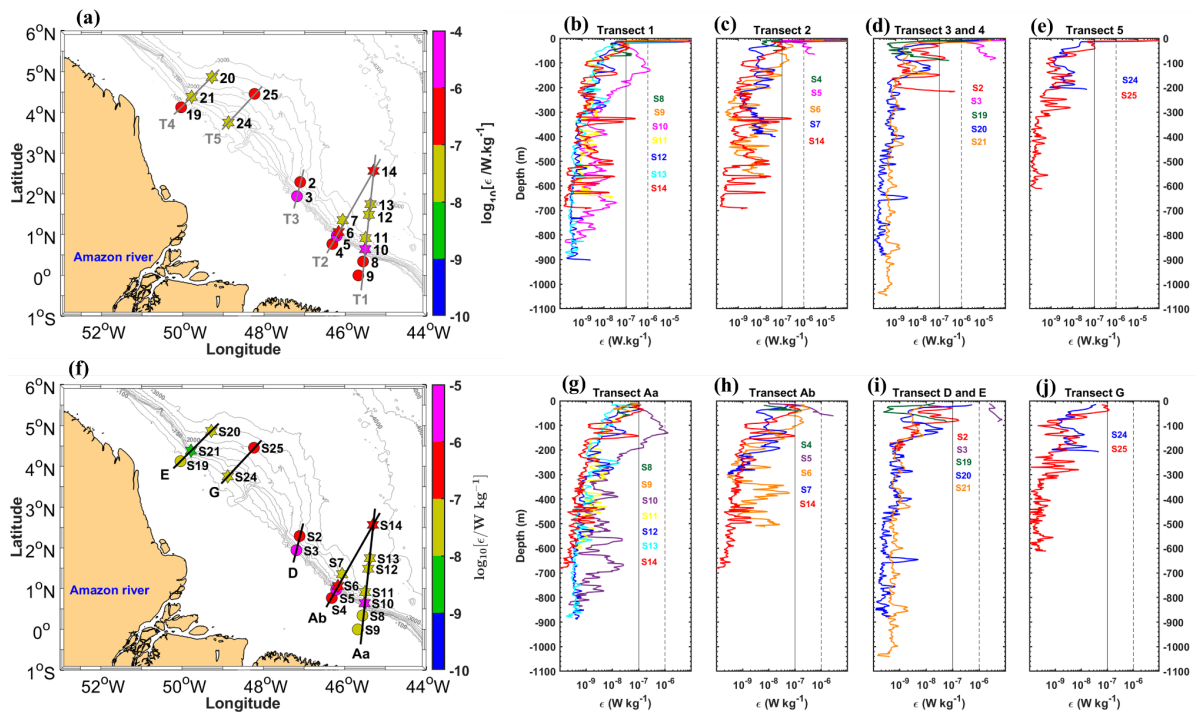


figure RC1.3: (a) Horizontal maximum and (b)-(c)-(d)-(e) vertical dissipation rates ( $\epsilon$ , in  $W\ kg^{-1}$ , on a logarithmic scale) before revisions and checks processes for all stations along transects T1 to T2. (f) Horizontal maximum and (g)-(h)-(i)-(j) vertical dissipation rates ( $\epsilon$ , in  $W\ kg^{-1}$ , on a logarithmic scale) after revisions and checks processes for all stations along transects Aa, Ab, D, E, and G. Distinct colors are used to represent each station within each transect. Dashed and solid black lines in panels (b) to (e) are included for comparison purposes.

5. Semidiurnal currents are obtained by removing the mean current over a tidal period from the baroclinic velocities. This will include other frequencies. In fact Figure 4 has everything: 6 hrs, 12 hrs, ... It would be better to make a least squares fit of sines and cosines to the various tidal periods. With a short time series of 1 day, the frequency resolution is 1 cycle/day. So you could try fitting diurnal and semidiurnal, but recognizing that they are not formally distinct. Also the inertial period here is at least 5 days and so cannot be determined in your dataset. This is a key limitation that should be pointed out. At the equator every frequency just about is in the internal wave band.

How or even whether the currents should be separated in frequency seems to be a major question given the limitations of the data and the desire to examine dissipation.

R: Thanks for your comments and suggestions. Based on your suggestion we completely re-worked our methodology, it has gained in robustness.

Indeed, semidiurnal currents are obtained by removing the mean current over a tidal period from the baroclinic velocities. This could include other frequencies. We have tried to make a least squares fit of sines and cosines to the various tidal periods (diurnal and semidiurnal), as you recommended. This fitting only resolved for the S14 station, which is 42 hours long, compared to the other shorter stations of 17 hours. This may be due to the length of the time series.



Below we present the result at S14. Figure RC1.4 shows the time series of the data collected and Figure RC1.5 and RC1.6 shows the diurnal and semidiurnal fitting at S14.

Indeed, with a short time series of 1 day, the frequency resolution is 1 cycle/day. We recognize that diurnal and semidiurnal fitting are not formally distinct, and that the inertial period here is at least 5 days and so cannot be determined in our dataset. We have pointed out this key limitation, as you recommended, in the manuscript in the lines 191-197 of the revised manuscript.

We know that our time series data are shorter than 17 h hours at all long stations, except for S14 (long of 42 hours). Indeed, the diurnal and semidiurnal periods fitting are not formally distinct and the inertial period (at least 5 days) cannot be determined in our dataset, limiting the currents separation in frequency and our desire to examine dissipation associated.

We have added the Figure RC1.4 and Figure RC1.5 in the Appendix of the revised manuscript (Figs. A4 and A5).

Indeed, even if currents need to be separated in frequency, this seems to be a major issue given the limitations of the data and the desire to examine dissipation. We have also pointed out this limitation in the manuscript in the lines 191-197 of the revised manuscript, as shown below:

“

Note that continuously collected SADCPC for some stations (e.g., S11) are not sufficiently resolved due to gaps filled by interpolating between time points. The similar processing are applied to the CTD-O<sub>2</sub> data collected alternately. SADCPC time series data are less than 17 hours at all long stations, except for S14, which spans 42 hours. As a result, the diurnal and semidiurnal period fittings are not formally distinct (except at S14; Figs. A4 and A5, Appendix), and the inertial period (at least 5 days) cannot be resolved in our dataset. This limits our ability to separate currents by frequency and examine the associated dissipation.

The velocity profiles from LADCP are glued into our SADCPC time series data below ~ 500 m depth at long stations..

“

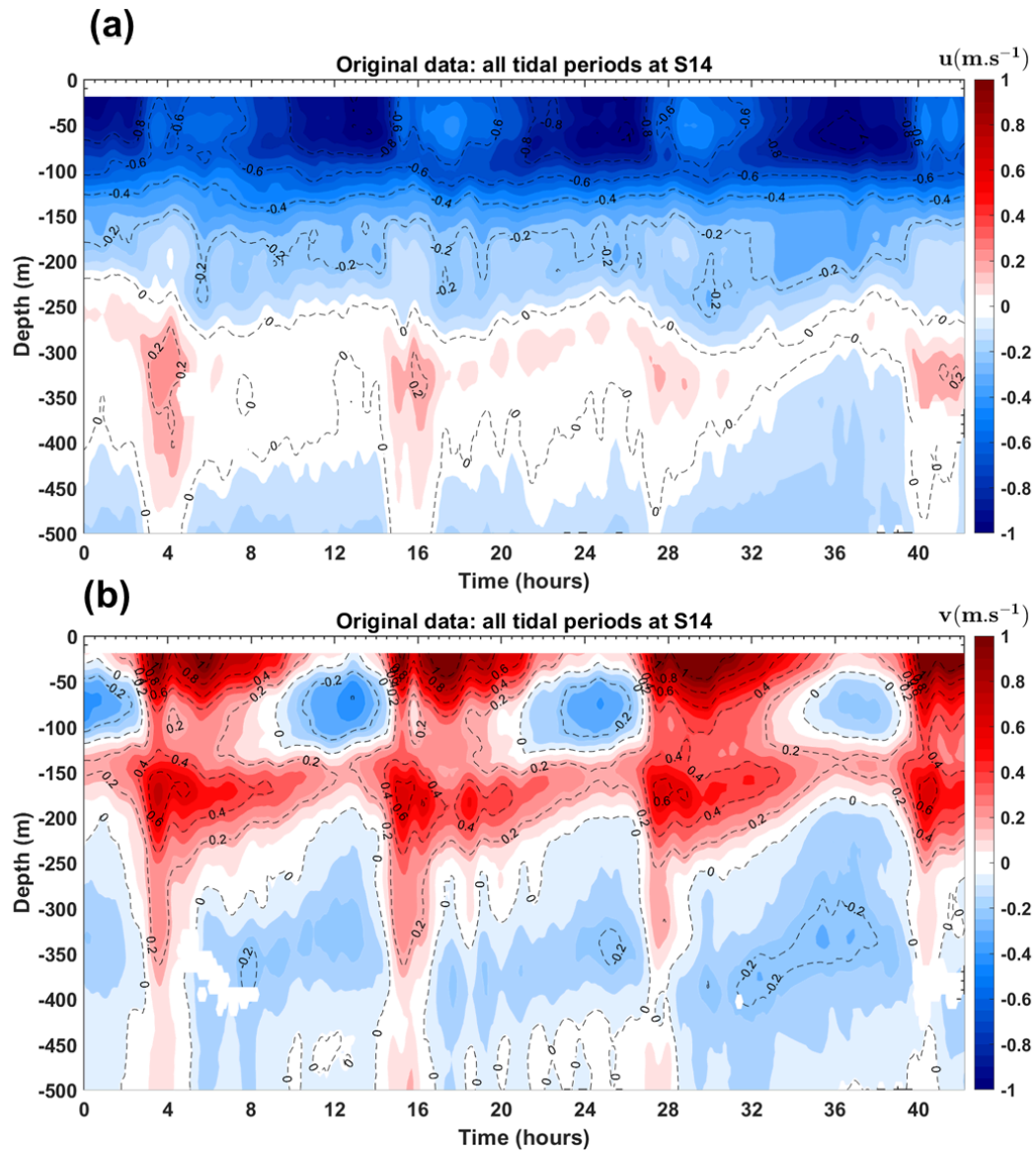


Figure RC1.4: Time series of (a) zonal and (b) meridional current of the original SADCPC data at a fixed position for station S14. Time is scaled to start at  $t=0$ .

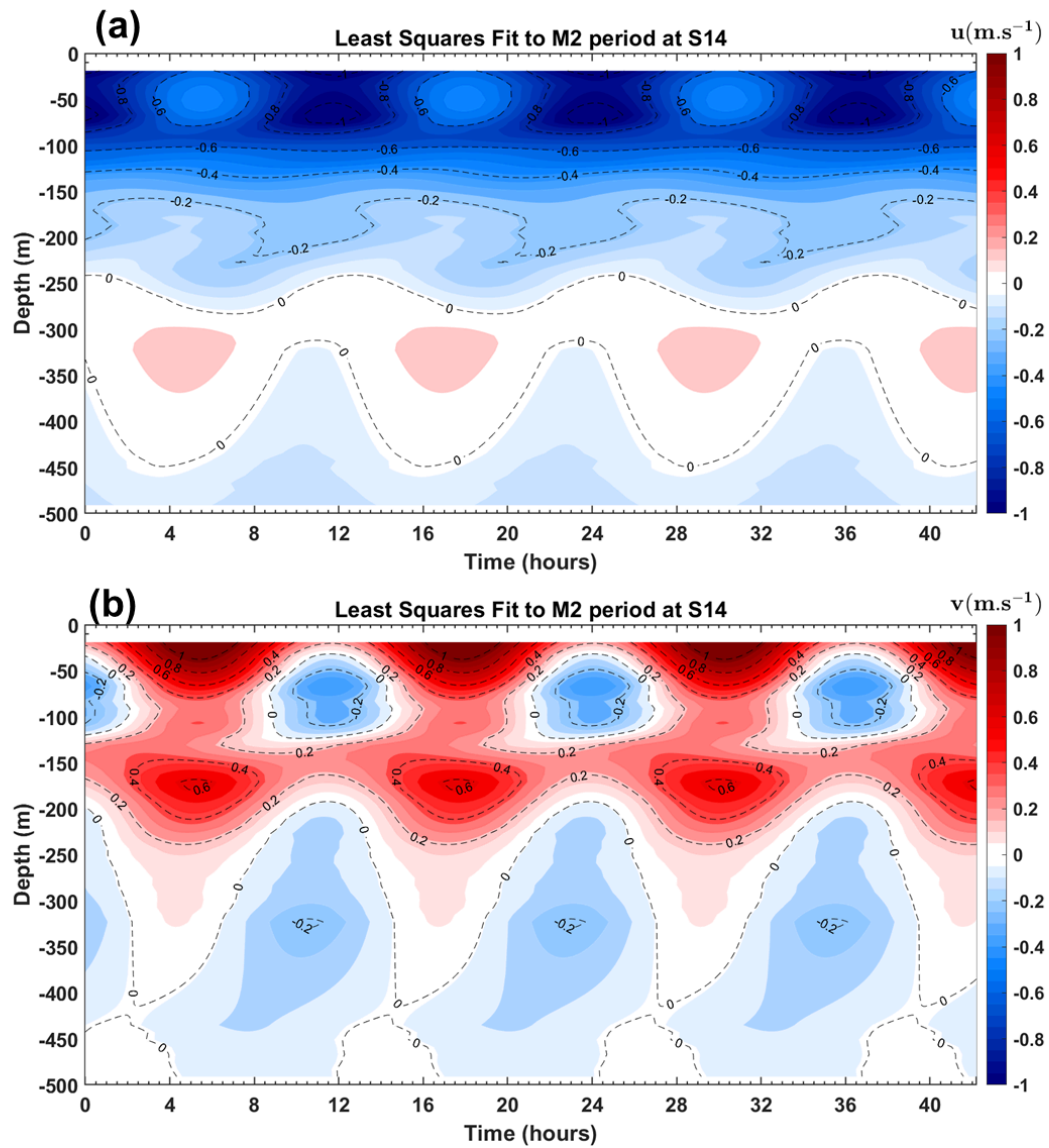


Figure RC1.5: Least squares fit of sines and cosines to M2 period for (a) zonal and (b) meridional current at station S14. Time is scaled to start at  $t=0$ .

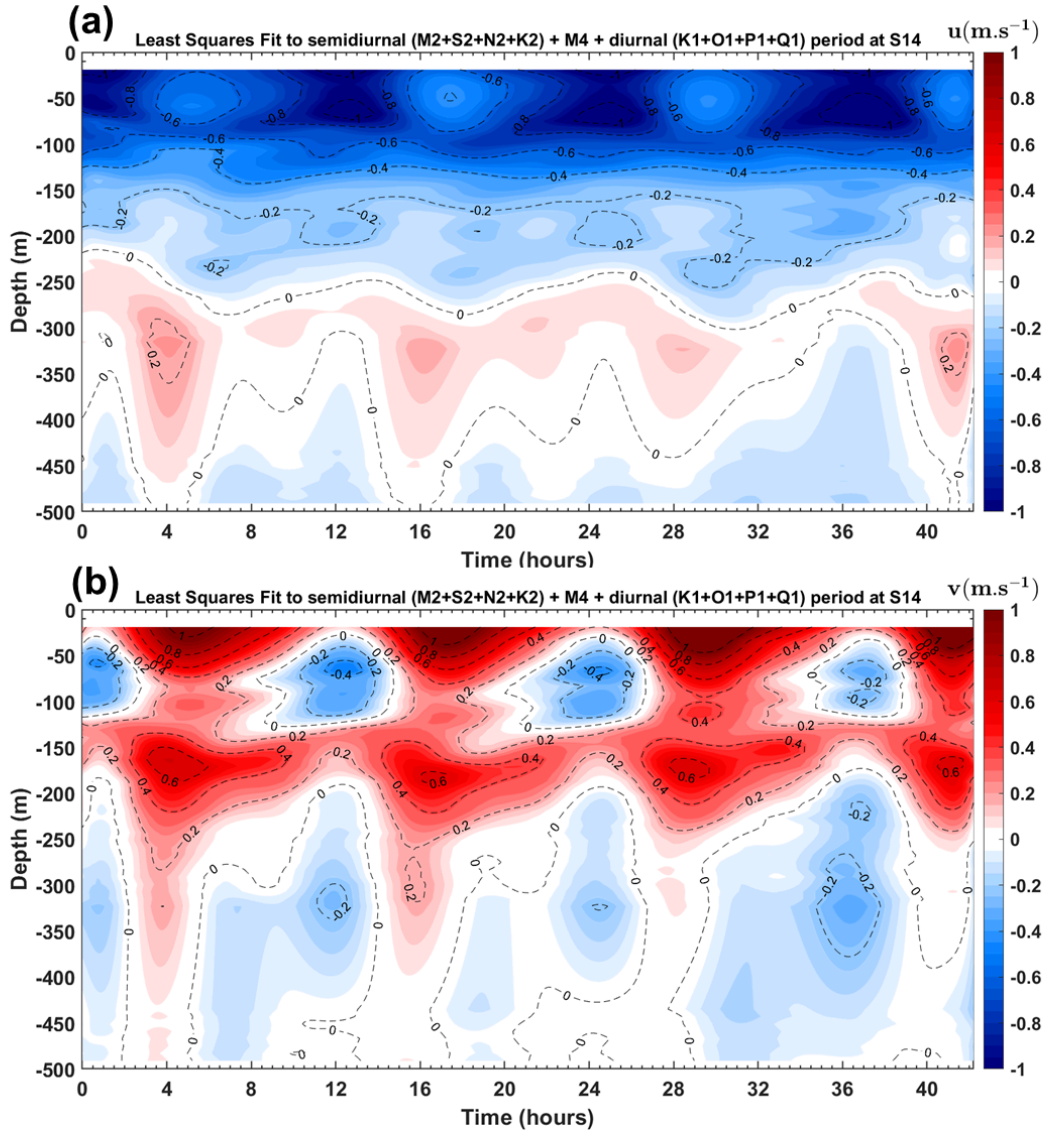


Figure RC1.6: Least squares fit of sines and cosines to semidiurnal (M2+S2+N2+K2) + M4 + diurnal (K1+O1+P1+Q1) periods for (a) zonal and (b) meridional current at station S14. Time is scaled to start at t=0.

The revision of section “methods/Baroclinic currents” can be found below in lines 178-212 of the revised manuscript:

“

#### Baroclinic currents

To analyze the processes explaining dissipation and mixing, particularly along internal tidal (IT) paths, we estimate shear instabilities associated with the semi-diurnal (M2) ITs and mean circulation, as well as their contributions to mixing.

The M2 tidal component of the tidal current is derived by calculating the baroclinic (semi-diurnal) tidal velocity  $[u'', v'']$  (Fig. A3, Appendix), following these equations:

$$[u', v'] = [u, v] - [u_{bt}, v_{bt}], \quad (1)$$

$$[u_{bt}, v_{bt}] = \frac{1}{H} \int_{-H}^0 [u, v] dz, \quad (2)$$

$$[u'', v''] = [u', v'] - [\bar{u}', \bar{v}']. \quad (3)$$

Here,  $[u, v]$  represent total horizontal velocities (Fig. A3, Appendix) obtained from SADC data. The components  $[u', v']$  and  $[u_{bt}, v_{bt}]$  represent baroclinic and barotropic components of horizontal velocities, respectively (Fig. A3, Appendix).  $H$  is water depth. The baroclinic mean velocities  $[\bar{u}', \bar{v}']$  (Fig. A3, Appendix), calculated to estimate mean circulation along IT paths, are decomposed into along-shore  $\bar{u}'_l$  and cross-shore  $\bar{u}'_c$  velocities. The overbar denotes the average over a M2 tidal period.

Note that continuously collected SADC data for some stations (e.g., S11) are not sufficiently resolved due to gaps filled by interpolating between time points. The similar processing are applied to the CTD-O<sub>2</sub> data collected alternately. SADC time series data are less than 17 hours at all long stations, except for S14, which spans 42 hours. As a result, the diurnal and semidiurnal period fittings are not formally distinct (except at S14; Figs. A4 and A5, Appendix), and the inertial period (at least 5 days) cannot be resolved in our dataset. This limits our ability to separate currents by frequency and examine the associated dissipation.

The velocity profiles from LADCP are glued into our SADC time series data below  $\sim 500$  m depth at long stations.

To evaluate shear instabilities associated with ITs and the mean background circulation, we compute the baroclinic tidal vertical shear squared ( $S^{2''}$ ) and mean shear squared ( $\overline{S^{2'}}$ ) (Fig. A3, Appendix), as follows:

$$S^{2''} = (\partial u''/\partial z)^2 + (\partial v''/\partial z)^2, \quad (4)$$

$$\overline{S^{2'}} = (\partial \bar{u}'/\partial z)^2 + (\partial \bar{v}'/\partial z)^2. \quad (5)$$

To evaluate the impact of bottom friction on mixing, we calculate kinetic energy  $\epsilon_f = \frac{1}{2} \rho_s (u_f^2)$  near the bottom boundary layer at shallow stations using friction velocity  $u_f = u_b \sqrt{C_d}$ , where  $C_d = 2.5 \times 10^{-3}$  is a drag coefficient obtained from the NEMO model. Huang et al. (2019) showed that the bottom boundary layer thickness spatially varies between 15-123 m in the Atlantic Ocean, with a median of  $\sim 30$ -40 m in the North Atlantic. We define bottom layer thicknesses in our study area based on measured bathymetry from CTD-O<sub>2</sub> and near-bottom currents from ADCP. Here,  $u_b$  is the total velocity averaged over a thickness of 20 m above the seabed for shallow stations and 40 m for deep stations.

The individual contributions of semi-diurnal ITs and mean circulation are then expressed as follows:  $\overline{E''}/(\overline{E''} + \overline{E''})$  for tidal contribution and  $\overline{E''}/(\overline{E''} + \overline{E''})$  for mean circulation contribution. Here,  $E = N*S$ .  $N$  is the buoyancy frequency and  $S$  is vertical shear.  $S$  can be substituted by  $S^{2''}$  and  $\overline{S^{2'}}$ .

“

6. Shear. Based on the calculations for barotropic u, baroclinic u, and tidal u. I do not see why shear is different for baroclinic and tidal u.

Tidal  $u$  = baroclinic  $u$  - lowpassed  $u$ .

It seems like it should be almost the same as the shear in baroclinic  $u$ . I'm confused. It would be helpful to see more of the processing steps and their justifications for these steps.

R: Thanks for your comments.

We have explained with more care the methodology step by step. Please see the previous response to your comment n°5.

Perhaps there is a problem with the short time series or large tidal isopycnal displacements. I suggest plotting some profiles of  $u$  and shear for barotropic, baroclinic, and tidal  $u$  to make sure these calculations are correct.

R: Thanks for your comments.

Indeed, our time series from SADCPC are shorter than 17 hours at all long stations, except for S14 (long of 42 hours). Our calculations, as shown below, seem correct to be used. Please see below the example for S14 over a (M2) tidal cycle, the plots of  $u$  and shear for total, barotropic, total baroclinic, and baroclinic tidal, and mean baroclinic in figures RC1.7. We have added this example in the Appendix of the revised manuscript (Fig. A3)

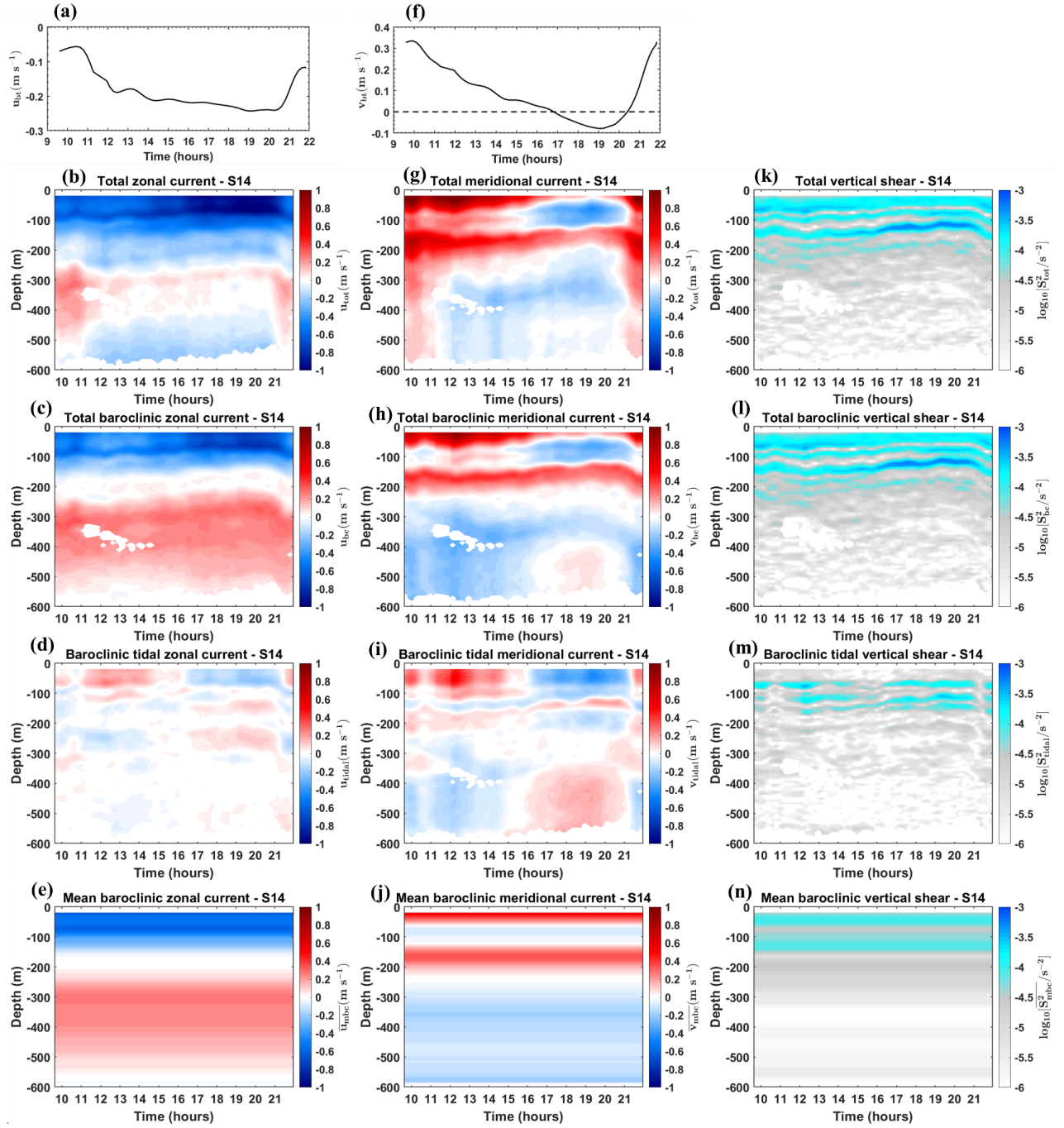


Figure RC1.7: For Station S14. zonal currents for (b) total, (a) barotropic, (c) total baroclinic, (d) semi-diurnal baroclinic tidal, and (e) mean baroclinic. Meridional currents for (g) total, (f) barotropic, (h) total baroclinic, (i) semi-diurnal baroclinic tidal, and (j) mean baroclinic. Vertical shear for (k) total, (l) total baroclinic, (m) semi-diurnal baroclinic tidal, and (n) mean baroclinic.

Also you could try doing these calculations on an isopycnal since it looks like the velocity structure is being heaved up and down by the internal waves. That could lead to artifacts in the calculations on depth surfaces. For example in Figure 4a1, on  $\sigma = 26$ , current is negative.

R: Thanks for your comments.

We performed the calculations on the  $26 \text{ kg/m}^3$  isopycnal as suggested and presented the results in the figure RC1.8 below. But we think that these results are not so robust since CTD data were collected alternately and needs to be interpolated between time points (e.g., between 17:00 and 20:00 at station S10). Consequently, the density data are not sufficiently resolved to allow the currents to be projected on over the tidal period. Doing this would introduce potential artifacts in the calculations on an isopycnal. So we decided to stick with our first calculation, pointing out in lines 191-197 of the revised manuscript:

“

Note that continuously collected SADCPC for some stations (e.g., S11) are not sufficiently resolved due to gaps filled by interpolating between time points. The similar processing are applied to the CTD-O<sub>2</sub> data collected alternately. SADCPC time series data are less than 17 hours at all long stations, except for S14, which spans 42 hours. As a result, the diurnal and semidiurnal period fittings are not formally distinct (except at S14; Figs. A4 and A5, Appendix), and the inertial period (at least 5 days) cannot be resolved in our dataset. This limits our ability to separate currents by frequency and examine the associated dissipation.

The velocity profiles from LADCP are glued into our SADCPC time series data below  $\sim 500 \text{ m}$  depth at long stations.

”

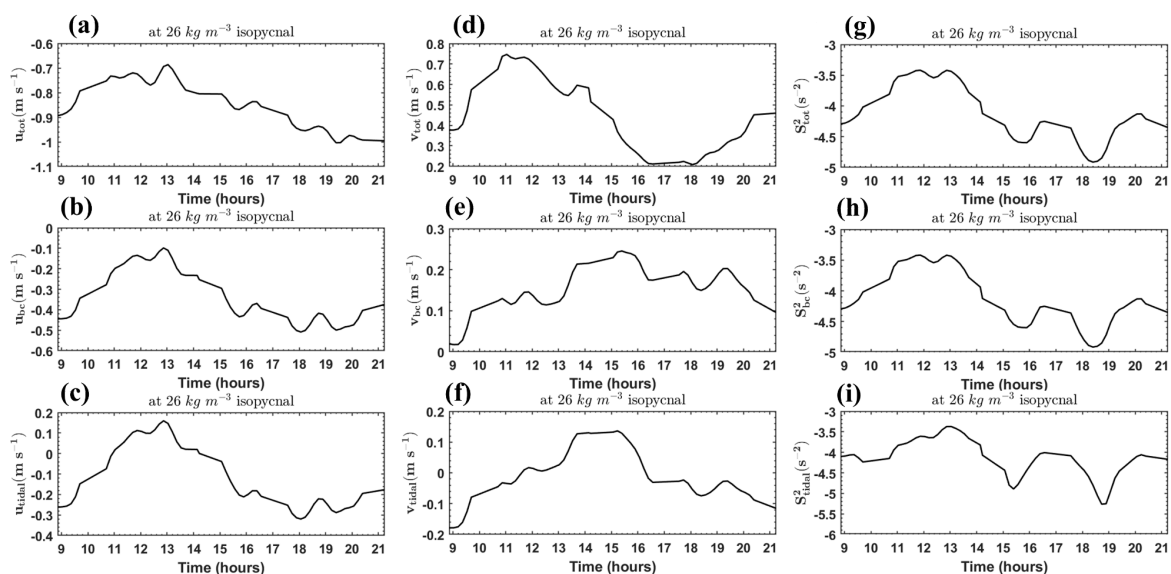




Figure RC1.8: For Station S10, calculations on the 26 kg/m<sup>3</sup> isopycnal. zonal currents for (a) total, (b) total baroclinic, and (c) semi-diurnal baroclinic tidal. Meridional currents for (d) total, (e) total baroclinic, and (f) semi-diurnal baroclinic tidal. Vertical shear for (g) total, (h) total baroclinic, and (i) semi-diurnal baroclinic tidal.

At 120 m, current switches sign multiple times. As far as I can tell, there are no plots of shear.

Also I do not understand why only one component of the current is shown. Using both components is needed to understand propagation of internal tides either vertically or horizontally (alongshelf vs cross-shelf).

Section 3.2.2 uses a lot of words that could be instead expressed concisely in a map with current vectors. Different colours could indicate different depths.

R: Thanks for your comments.

We have added in **the revised manuscript** plots of the shear of shear for tidal U and mean U at selected stations (e.g., S10, S11, and S14; in Fig. 6, and Figs. A15 to A18, Appendix) and also reported here in figure RC1.9 below

Indeed, at 120 m, the current switches direction multiple times. We have added in **the revised manuscript** a 2D map of mean current (vectors) and also reported here in figure RC1.12 below.

Indeed, both zonal (u) and meridional (v) components are essential for understanding the propagation of internal tides, whether vertically or horizontally (along-shelf vs. cross-shelf). As we more focused on zonal (u) and meridional (v) components of ITs, we have added these components in **the revised manuscript** at selected stations (e.g., S10, S11, and S14; in Fig. 5, and Figs. A9 to A15, Appendix) and also reported here in figure RC1.10 and RC1.11 below.

Additionally, we have visualized the mean baroclinic current using a map with current vectors (as reported here in RC1.12).

In response to all reviewer comments, we have revised several sections of the manuscript, including “Methods,” “Results,” and “Discussion and Conclusion,” spanning **lines 121–568** of in **the revised manuscript**.

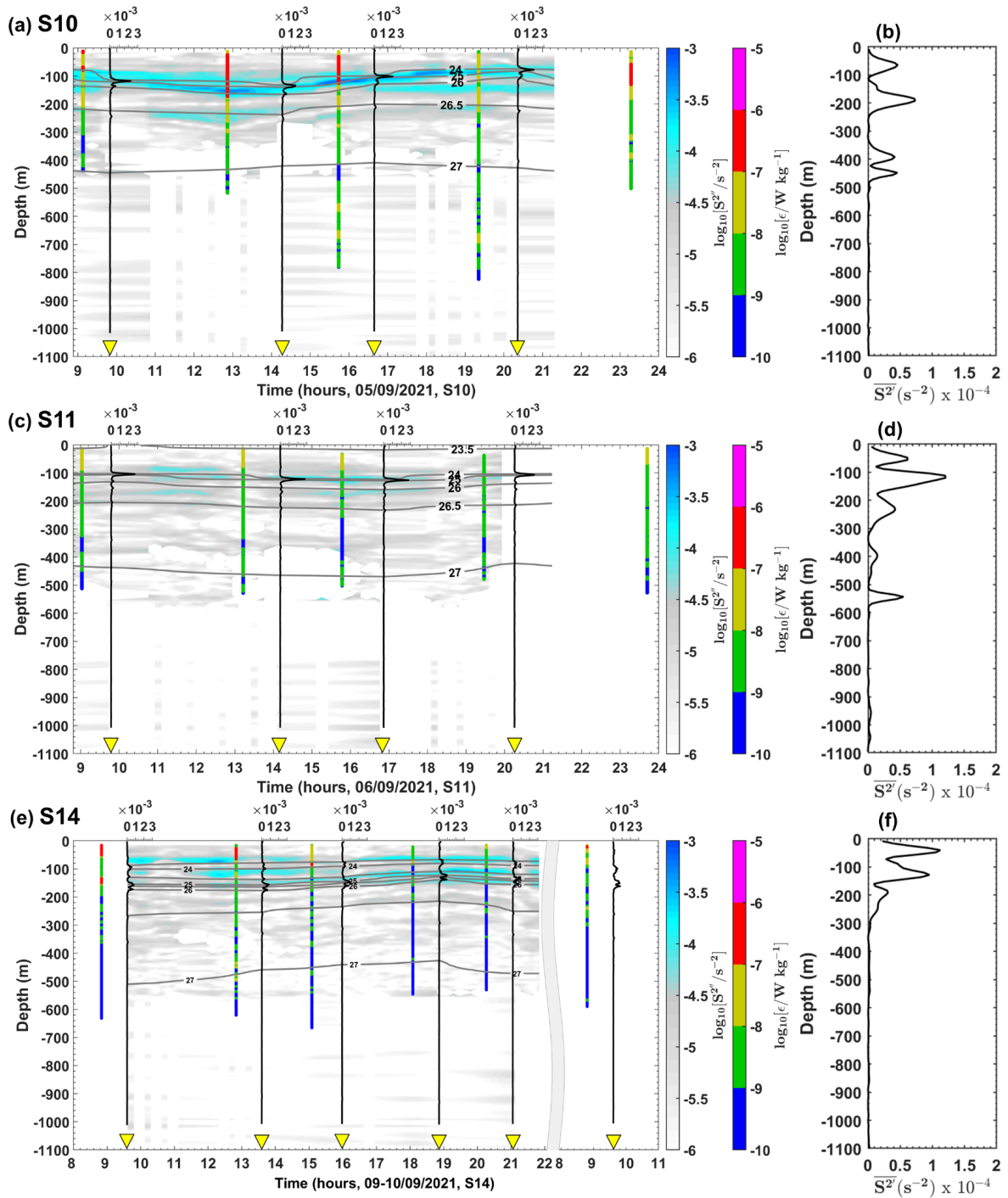


Figure RC1.9: Semi-diurnal baroclinic vertical shear squared ( $S^2$ , in  $\text{m s}^{-1}$ , on a logarithmic scale) for stations (a) S10, (c) S11, and (e) S14. Panels (a), (c), and (e) also display the buoyancy frequency squared ( $N^2$ , in  $\text{s}^{-2}$ ) represented by vertical black lines, potential density represented by grey contours, and dissipation rate profiles ( $\epsilon$ , in  $\text{W kg}^{-1}$ , on a logarithmic scale) represented by blue vertical colored bars. Mean baroclinic vertical shear squared ( $\overline{S^2}$ , in  $\text{m s}^{-1}$ ) for stations (b) S10, (d) S11, and (f) S14.

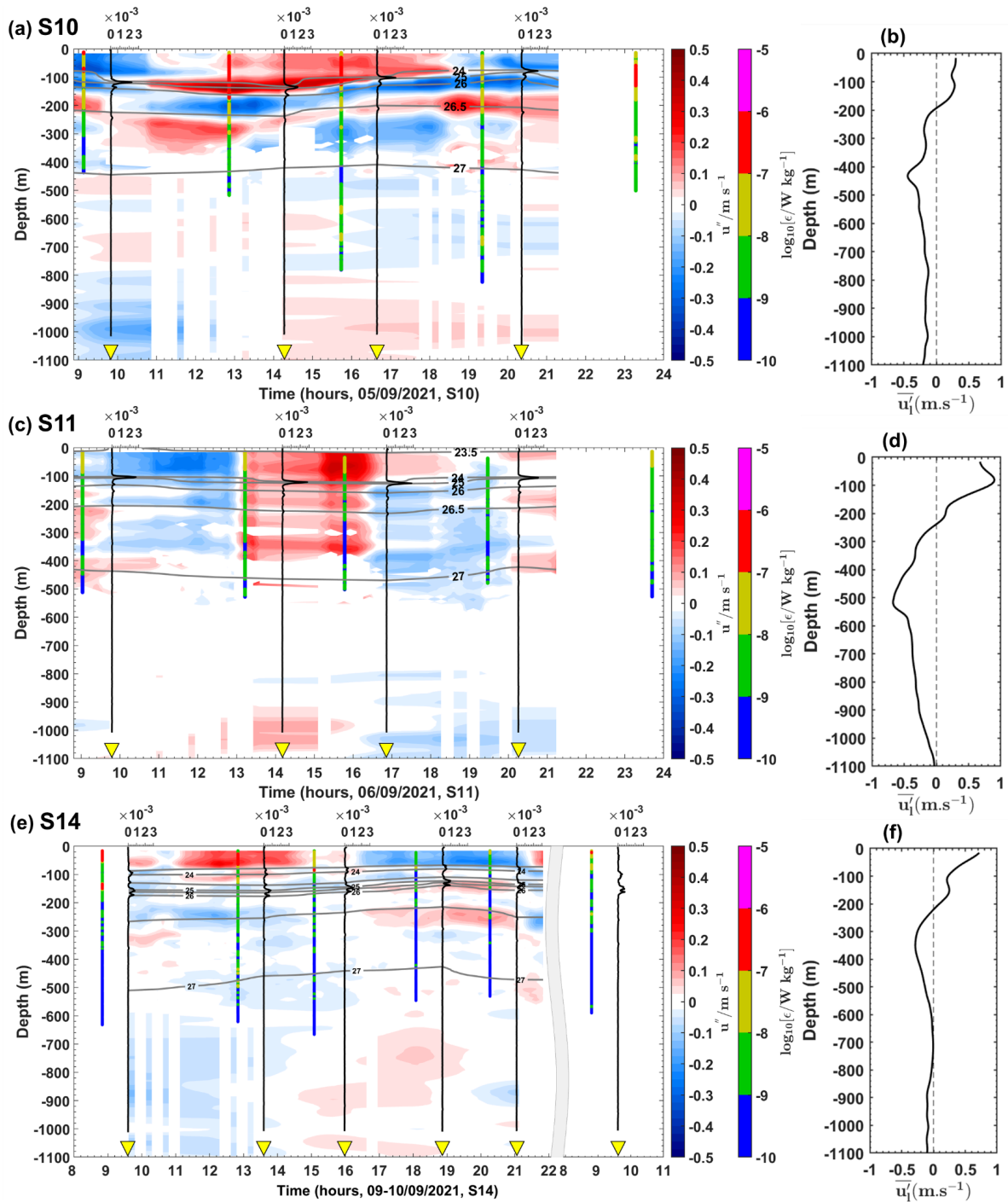


Figure RC1.10: Semi-diurnal baroclinic zonal currents ( $u''$ , in  $m s^{-1}$ ) from the ADCP for stations (a) S10, (c) S11, and (e) S14. Panels (a), (c), and (e) also display the buoyancy frequency squared ( $N^2$ , in  $s^{-2}$ ) represented by vertical black lines, potential density represented by grey contours, and dissipation rate profiles ( $\epsilon$ , in  $W kg^{-1}$ , on a logarithmic scale) represented by vertical colored bars. Along-shore mean baroclinic currents ( $\overline{u}'_l$ , in  $m s^{-1}$ ) from the ADCP for stations (b) S10, (d) S11, and (f) S14.

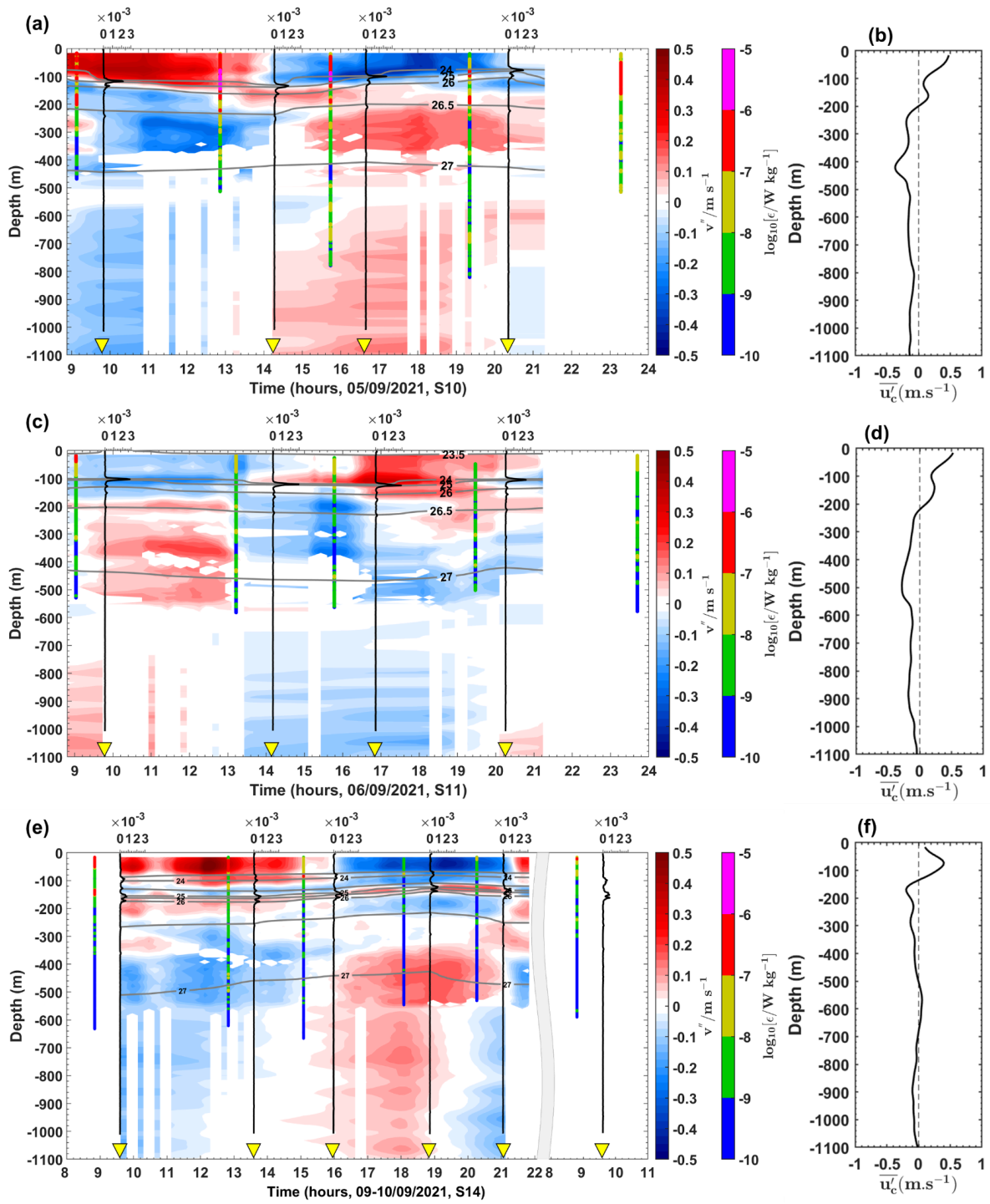


figure RC1.11: Semi-diurnal baroclinic meridional currents ( $u''$ , in  $\text{m s}^{-1}$ ) from the ADCP for stations (a) S10, (c) S11, and (e) S14. Panels (a), (c), and (e) also display the buoyancy frequency squared ( $N^2$ , in  $\text{s}^{-2}$ ) represented by vertical black lines, potential density represented by grey contours, and dissipation rate profiles ( $\epsilon$ , in  $\text{W kg}^{-1}$ , on a logarithmic scale) represented by vertical colored bars. Cross-shore mean baroclinic currents ( $\overline{u_c}$ , in  $\text{m s}^{-1}$ ) from the ADCP for stations (b) S10, (d) S11, and (f) S14.

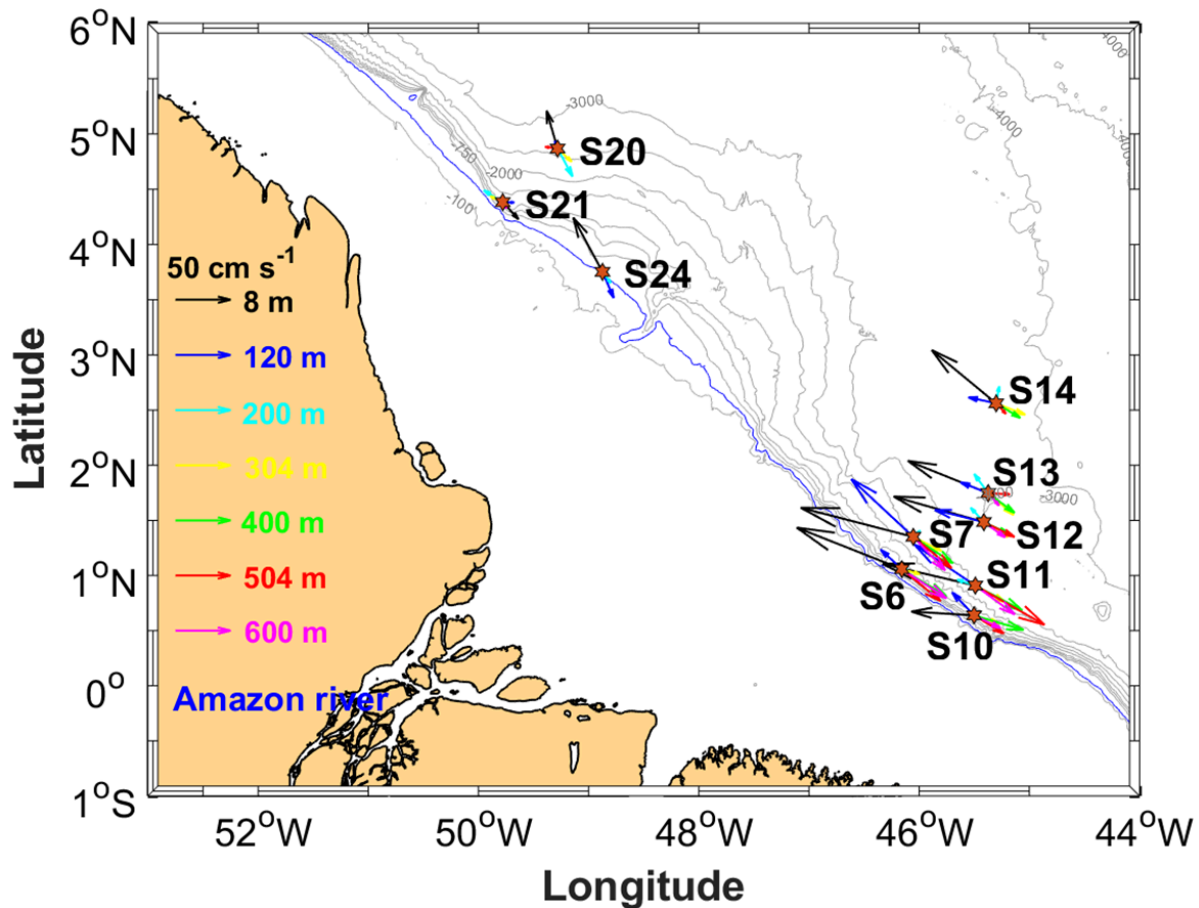


Figure RC1.12: Map of mean baroclinic currents (vectors) at stations. Colored arrows are used to distinguish the currents at different depths: Black, blue, cyan, yellow, green, red and magenta arrows at depth equal to 8 m, 120 m, 200 m, 304 m, 400 m, 504 m and 600 m, respectively. The blue line indicates the 200 m isobath.<sup>7</sup> Dissipation is estimated using parameterizations based on shear in the tidal current. Actual dissipation depends on the total shear or strain.

R: Thanks for your remarks.

Indeed, dissipation was estimated using parameterizations based on shear in the tidal current. MacKinnon-Gregg parameterization was applied as a proxy to evaluate the contributions of tidal and low-frequency shear, primarily for comparison purposes. However, no scientific results were derived from it.

Ultimately, thanks to all the reviewers comments, we decided to remove the section using this parameterization. Eventually, we will work on a separate paper to test the robustness of this parameterization against these observations.

In the present paper, we focused on vertical shear to separate the contributions of tidal and mean shear. Revisions can be found in the "Methods" section in [the revised manuscript](#).

8. Some more explanation of the applicability of MacKinnon-Gregg parameterization vs those based on Gregg (1989) is needed. The former is related to the limited bandwidth of internal waves (in shallow water for example), while the latter is for a typical open ocean environment away from generation sites. How does this apply to your study sites?

R: Thanks for your comments.

Please see the previous response to your comment n°7.

9. Naming of the sites and transects. Maybe these are what are used on the cruise but they should be rearranged in a logical order to *help the reader*. It is completely confusing as it is. Consider another system, such as A B C D ... for generation sites. If there's a transect near site A, it's transect A.

If there are 3 stations on transect A, they are sites A1 A2 A3 and counting higher offshore. Or labelled by isobath, A500, A1000, A2000... There may not be a generation site for each transect.

For example lines 265-267 read: "The highest baroclinic tidal current velocities were observed (between 25-48 cm.s-1) at sites Aa and Ab along T1-T2. Whereas lower tidal velocities (< 25 cm.s-1) are found in site F along T4 (e.g., at S20 and S21) compared to OUT-ITs stations (e.g., at S24)." So I have to use the map. This could read instead as: "The highest baroclinic tidal current velocities were observed (between 25-48 cm.s-1) at sites D and E along transects D and E. Whereas lower tidal velocities (< 25 cm.s-1) are found in site A along transect A (e.g., at A2 and A3) compared to OUT-ITs stations (e.g., at site B1)."

R: Thanks for your comments and suggestions.

We acknowledge that the nomenclature might be somewhat confusing, and your suggestions would indeed simplify understanding. However, we have chosen to retain the naming convention used in previous studies of this region. Specifically, the names A to F for the generation sites were already defined in earlier works (e.g., Magalhães et al., 2016; Tchilibou et al., 2022; Assene et al., 2024; de Macedo et al., 2023).

To maintain consistency with these studies, we have decided to adhere to the established naming convention. As a compromise between simplifying the nomenclature and preserving coherence with previous works, we have designated the transects with the same letters as their corresponding generation sites: Aa, Ab, D, E, G, as shown in figure 1 of [the revised manuscript](#) and reported here in figure RC1.13.

Revisions reflecting this choice can be found in the subsection "Baroclinic tidal current" (lines 298–320) of [the revised manuscript](#).

As an example, the previous phrase "is now transform in :

“

The baroclinic tidal velocities reveal a superposition of 3-5 tidal modes at IN-ITs stations (Figs. 5a, 5c, and 5e, and Figs. A9 to A15, Appendix). A greater number of modes is observed near the shelf-break (e.g., 4 modes at S6 and 5 modes at S10), while fewer modes are detected far from (e.g., 3 modes at S7, S12 and S14). Higher tidal velocities ranging from 25-50 cm s<sup>-1</sup> are found between 80-350 m along transects Aa and Ab (e.g., at S6, and S10). In contrast, lower tidal velocities, typically below 25 cm s<sup>-1</sup>, are more pronounced along transect E (e.g., at S20, and S21) to OUT-ITs stations along transect G (e.g., at S24).

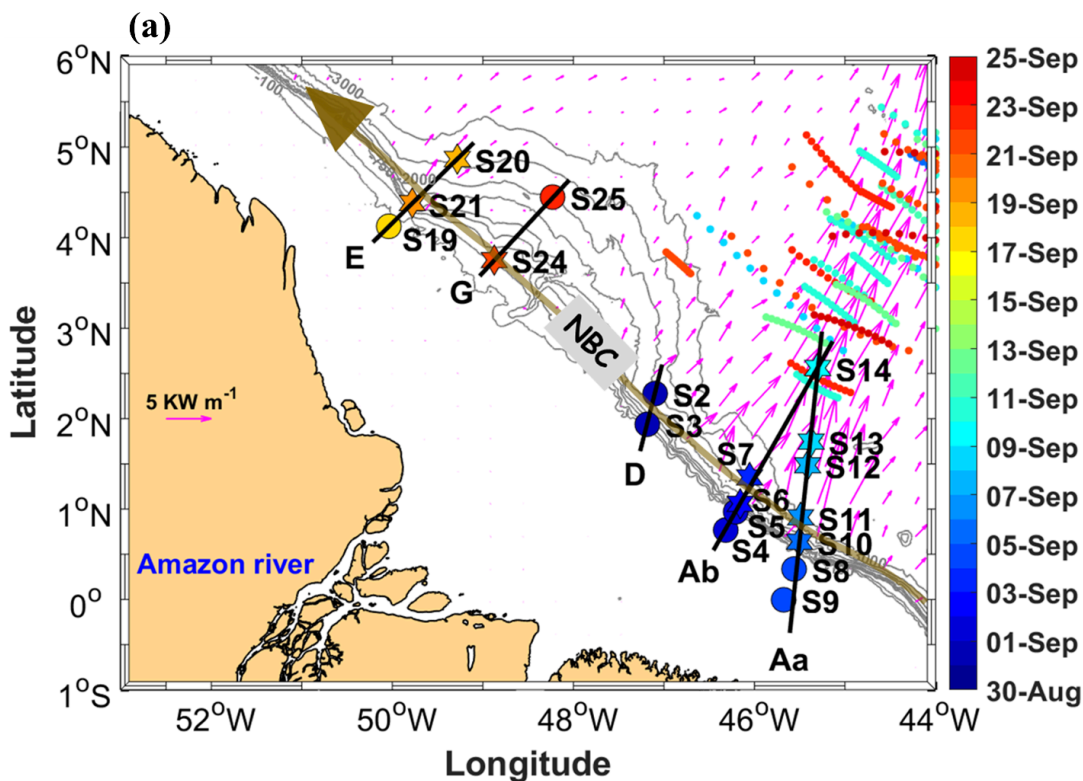


Figure RC1.13:a) Map of a part of the AMAZOMIX 2021 cruise off the Amazon shelf, showing bathymetric contours (100 m, 750 m, 2000 m, and 3000 m isobaths) in gray. Colored circles and stars indicate short and long CTD-O2/L-S-ADCP stations, respectively, with the corresponding sampling dates represented by the color bar. Solid black lines depict SADCPC transects (for Aa, Ab, D, G, and E). Magenta arrows show the 25-hour mean depth-integrated baroclinic IT energy flux (September 2015, from the NEMO model) originating from IT generation sites (Aa, Ab, D, and E) along the shelf break. The solid brown line represents the NBC pathways illustrating background circulation. Shattered colored lines highlight ISW signatures.

10. Mixing is invoked as the explanation for increased nutrients. How about coastal upwelling? Where is the euphotic zone?

Thank you for this question, which gives us the opportunity to emphasize in the discussion that mixing is not the only significant process influencing the nutrient structure in the region.

Indeed, the nutrient inputs are controlled by several physical processes such as : Riverine Input, Coastal Upwelling, Internal Tides (ITs), Stratification Changes, and Eddies and Currents Off the Amazon shelf (Williams and Follows, 2003; Santos et al., 2008).

There is a very high variability of the euphotic layer in the region. It is very sensitive and heterogeneous from shelf to open ocean. The average estimated by in-situ observations is 13 m (Santos et al., 2008).

In our study, we tried to quantify the impact of internal tide on nutrient structuration. IT is indeed probably of a secondary order of importance in instantaneous but by its repetitiveness (12h, 14 Days), the cumulative effect of the tide over a year could impact the nutrient content. In addition, in a context of climate change this effect may be increased. (Yang et al., 2024)

In our study, we focused on quantifying what was measurable: specifically, vertical mixing, out of which approximately 50% was attributed to tidal processes. And, we also suggested that tidal mixing could contribute to nutrient inputs through vertical diffusivity since it has a strong contribution at the base of the mixed layer.

In the discussion, we contextualized these findings by emphasizing that mixing is not the only process influencing nutrient dynamics, and that upwelling also plays a significant role in nutrient supply.

Ultimately, we decided to remove and reserve all sections on “Nutrients fluxes” for a separate paper in progress.

11. Were adjustments to the shear parameterization based on the Gregg (2003), where it was shown that internal wave interactions may lead to little mixing at the equator?

R: Thanks for your comments and suggestions.

Please see the previous response to your comment n°7.

### **Minor comments**

1. Line 107- While the bin size of the LADCP may be 8 m, its resolution will be about 50 m. Examine where the vertical wavenumber spectra fall off

R: Thanks for your comments.

We have examined the vertical wavenumber spectra. We fitted a slope to the spectrum in the steep region to estimate the roll-off wavenumber. For example, the spectrum falls off as  $k^{-2.31}$  and  $k^{-2.11}$  for the (a) zonal and (b) meridional at S14, as reported below in figure RC1.14.



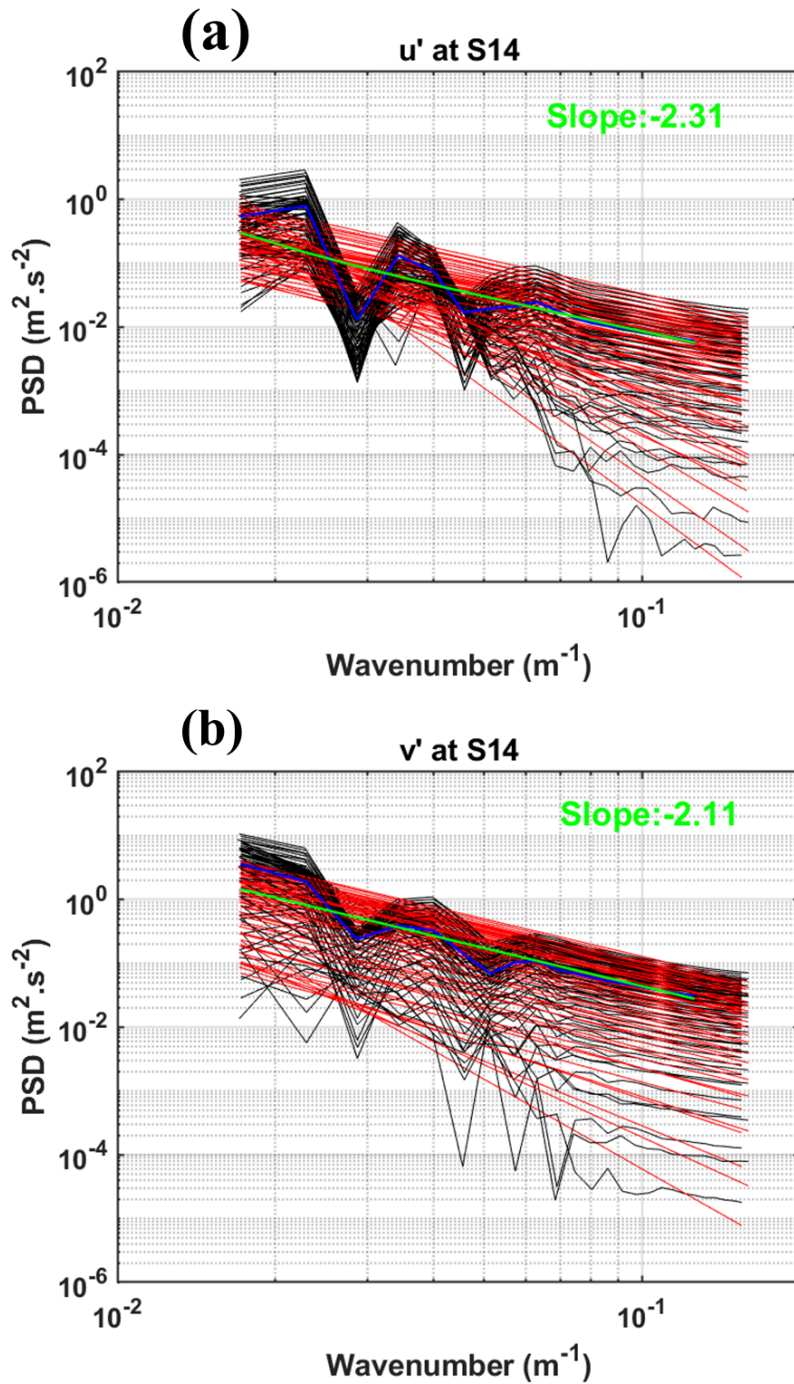


Figure RC1.14: Panels (a) and (b) show the vertical wavenumber spectra for perturbations in the total zonal and meridional currents, respectively, derived from LADCP data at station S14. Black curves represent the vertical wavenumber spectra for individual current profiles over time, while the blue curves indicate the spectra averaged across all time profiles. The red and green lines illustrate the linear fits to the spectra slopes for individual profiles and the mean spectrum, respectively. The fitted slopes of the mean perturbation spectra are indicated in green: (a) -2.31 for the zonal current and (b) -2.11 for the meridional current.

2. I would recommend a native English speaker help with some of the grammar, if possible. It's ok as is but could be much better. It will be less distracting for the reader. Then the reader will pay more attention to the science.

R: Thanks for your suggestions. We have revised all sections of the manuscript.

3. In sections 2.1 and 2.2, there is some jumping around from instrument to instrument. Please collect into sections for each instrument.

R: Thanks for your comments and suggestions.

We have revised sections 2.1 and 2.2 in lines 92-234 of the revised manuscript and reported here:

“

## 2 Data and Methods

### 2.1 Data collection

The AMAZOMIX cruise (Bertrand et al., 2021) was performed over the shelf/slope areas off the AROC during August-October 2021 aboard the IRD vessel RV *ANTEA*. At each designated site, 12-hour stations were set up, with repeated casts (4-5 casts per site) of Conductivity-Temperature-Depth-Oxygen (CTD-O<sub>2</sub>)/Lowered Acoustic Doppler Current Profiler (LADCP) and Velocity Microstructure Profiler (VMP) to measure the TKE dissipation rates over a complete tidal (M2) cycle, allowing the separation of the tidal component from the total current. A high-resolution (1/36°) NEMO (Nucleus for European Modeling of the Ocean) model (Madec et al., 2019) was used to determine station locations based on realistic IT generation and propagation maps (Tchilibou et al., 2022; Assene et al., 2024) and to estimate the mean background stratification.

Stations (Fig. 1a and 1c) were located inside the ITs fields, named “IN-ITs” (sites Aa, Ab and D: S2 to S14; site E: S19 to S21), and outside the ITs fields (S24 and S25), named “OUT-ITs”, on the shelf-break generation (sites Aa, Ab, D and F) and propagation along 5 transects (Aa, Ab, D, G, and E; Fig. 1).

CTD-O<sub>2</sub> measurements were obtained using a Seabird 911 Plus with dual sensors mounted in the rosette. The 24 Hz CTD-O<sub>2</sub> sensors were calibrated before and after the cruise. The standard deviation of temperature (salinity; oxygen) was 0.003 °C (0.003 PSU; 0.05 ml l<sup>-1</sup>) according to adjusted data. CTD-O<sub>2</sub> data were averaged over 1-m bins to filter out spikes and missing points, and aligned in time to correct the lag effects.

Two 300 kHz RDI LADCPs were mounted on the rosette to provide vertical current profiles with 8 m resolution, supplemented by 75 kHz shipboard ADCP (SADCP) profiles recorded continuously during the cruise. Vertical resolution of SADCP was adjusted according to bottom depth, e.g., 8 m for depths >150 m (at S6, S7, S10-S14, S20, S21, and S24) and 4 m for other depths. . Data processing and quality control followed GO-SHIP Repeat Hydrography Manual protocols. In total, 71 CTD-O<sub>2</sub>/LADCP profiles were collected during the AMAZOMIX cruise.

To characterize mixing, the TKE microstructure profiles were obtained from high-frequency (~ 2 mm resolution) measurements of temperature and velocity shear using a VMP-250 profiler (Rockland Scientific International, Inc.) capable of reaching depths up to 1000 m. The VMP-250 features two high-resolution thermistors (FP07) and two high-resolution velocity shear probes (probe 1 and 2; with 5% signal accuracy), with

a sampling rate of 1024 Hz. The profiler was deployed and retrieved via an electric winch and rope tether, with alternating deployments between the CTD-O<sub>2</sub>/LADCP profiles at 33 stations, yielding a total of 202 profiles. For this study, data from 18 stations (S2-S14, S19-S21, S24, and S25) comprising SADC data, 109 VMP profiles and 54 CTD-O<sub>2</sub>/LADCP profiles will be analyzed.

## 2.2 Methods

### TKE dissipation rates

The VMP data are processed using ODAS Matlab library (developed by Rockland Scientific International, Inc) to infer the TKE dissipation rate ( $\epsilon$ ). The processing methods for the VMP data are briefly described here and adhere to the recommendations of ATOMIX (Analyzing ocean turbulence observations to quantify mixing), as reported by Lueck et al. (2024), and have been validated against the benchmark estimates (presented in Fer et al., 2024).

First, the VMP data are converted into physical shear units, and the time series are prepared. Continuous sections of the time series are selected for dissipation estimation. Before spectral estimation, the aberrant shear signals caused by vessel wake contamination are removed. Collisions of the shear probe with plankton and other particles are removed using the de-spiking routine. The records from each section are then high-pass filtered (e.g., at station S6 and S10; Fig. 2a, and Fig. A1, Appendix).

Shear spectra are estimated using record lengths ( $L$ ) and Fast Fourier Transform segments of 2 s, which are cosine windowed and overlapped by 50% (e.g., at station S6 and S10; Fig. 2b, and Fig. A1, Appendix). Additionally, vibration-coherent noise is removed. Different  $L$  and overlap ( $O$ ) settings were selected and tested based on the environment (e.g., deep vs. shallow water), following Fer et al. (2024). For shallow stations,  $L$  ( $O$ ) was shortened to 5 s (2.5 s), in contrast to the 8 s (4 s) used for deeper stations, due to evidence of overturns observed in AMAZOMIX acoustic measurements at deeper stations (Koch-Larrouy et al., 2024; in preparation). This adjustment helped to optimize the spatial resolution of dissipation estimates in shallow water stations.

Finally,  $\epsilon$  is determined using the spectral integration method and by comparison with the Nasmyth empirical spectrum (Nasmyth, 1970). Quality assurance tests are carried out in accordance with ATOMIX's recommendations (Lueck et al., 2024). A figure of merit  $< 1.4$  is used to exclude bad data (e.g., at station S6 and S10; Fig. 2b, and Fig. A1, Appendix), and the fraction of data affected by de-spiking is  $< 0.05$ .

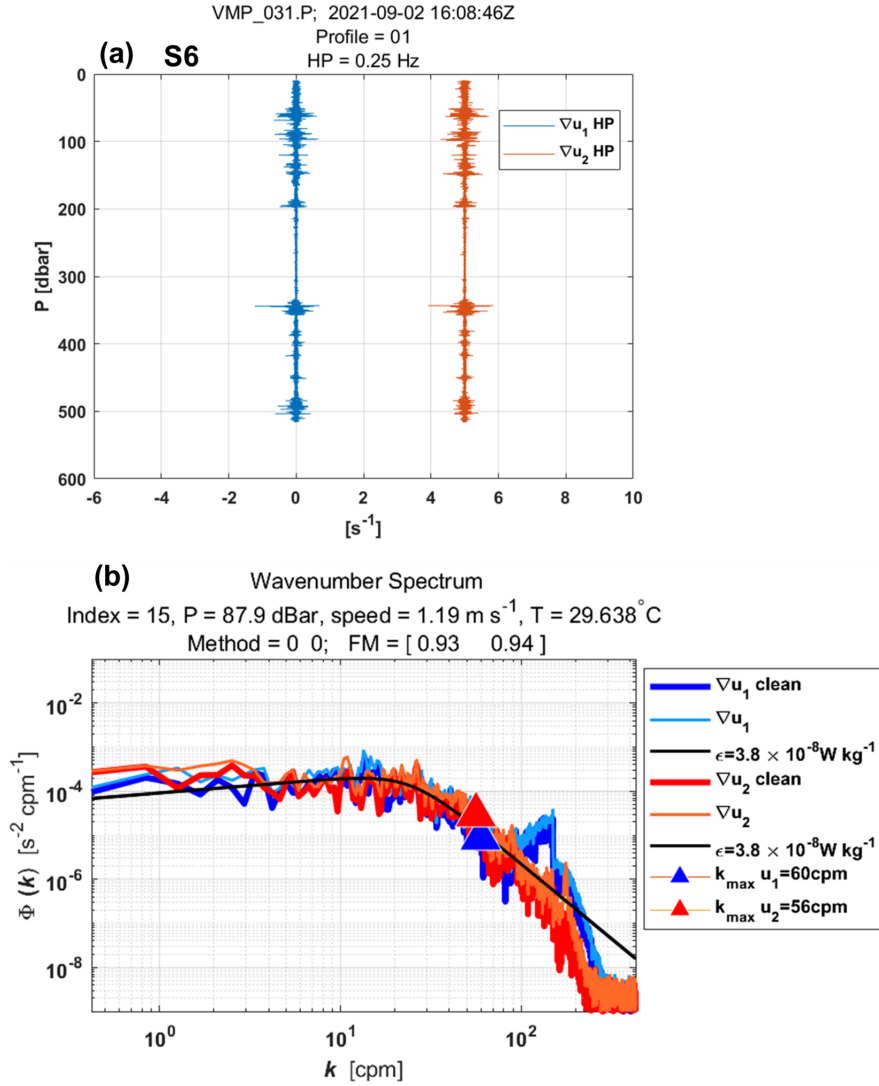


Figure 2: Example of wavenumber spectra from a dissipation structure segment used to determine the dissipation rate at station S6 at a pressure of 87.9 dBar. (a) Cleaned and high-pass filtered signals from shear probe 1 (blue) and shear probe 2 (red, offset by 5 s<sup>-1</sup>). (b) Wavenumber spectra for shear probes 1 and 2. Thick lines (blue for probe 1, red for probe 2) show shear spectra with coherent noise correction, while thin lines (sky blue for probe 1, orange for probe 2) show spectra without correction. Triangles mark the maximum wavenumber used for dissipation rate estimation. Black lines represent Nasmyth reference spectra for estimated dissipation rate of  $3.8 \times 10^{-8} \text{ W kg}^{-1}$  for both shear probes. Dissipation rate estimates for shear probe 1 and shear probe 2 at a pressure of 87.9 dBar yielded a figure of merit of 0.93 and 0.94, respectively.

### The vertical eddy diffusivity coefficient

The efficiency of turbulence in redistributing energy is assessed through the calculation of the vertical eddy diffusivity coefficient ( $K_z$ ). This coefficient is particularly significant in regions such as pycnoclines, where stratification suppresses mixing, making turbulence-driven mixing a key mechanism for vertical energy transport (Thorpe, 2007).

$K_z$  is calculated from  $\varepsilon$  following the formulation of Osborn (1980), given by  $K_z = \varepsilon / N^2$ . Here,  $N^2$  is the buoyancy frequency squared, which is calculated using the sorted potential density profiles ( $\sigma_\theta$ ) obtained from CTD-O<sub>2</sub> data. It is given by  $N^2 = - (g/\rho_0) (d\sigma_\theta/dz)$ , where  $\rho_0$  is a reference density (1025 kg m<sup>-3</sup>) and  $g$  is the gravitational acceleration.  $\Gamma$  is the mixing efficiency, defined as the ratio between the buoyancy flux and the energy dissipation, and is typically set to 0.2, which corresponds to the critical Richardson number  $Ri = 0.17$  (Osborn, 1980).  $\varepsilon$  is linearly interpolated into the depths of  $N^2$ .

Turbulence within the pycnocline can reduce stratification and increase vertical eddy diffusivity below the mixing layer (Thorpe, 2007). Subsurface mixing, driven by the breaking of ITs and shear instabilities, plays a particularly important role below the mixed layer, especially in equatorial waters (Gregg et al., 2003).

There are several criteria for defining the Mixed Layer Depth (MLD). In this study, we use the commonly accepted density threshold criterion of 0.03 kg m<sup>-3</sup>, as defined by de Boyer Montégut et al. (2004) and Sutherland et al. (2014), to estimate the MLD for each CTD-O<sub>2</sub> profile. Notably, comparisons with density thresholds of 0.01 and 0.02 kg m<sup>-3</sup> revealed no major differences in MLD across the AMAZOMIX stations and transects (Fig. A2, Appendix).

The miXing Layer Depth (XLD) is defined as the depth at which  $\varepsilon$  decreases to a background level (Sutherland et al., 2014). Previous studies have applied various thresholds for background dissipation levels, such as  $10^{-8}$  and  $10^{-9}$  W kg<sup>-1</sup> in higher latitudes based on in situ observations (Sutherland et al., 2014; Lozovatsky et al., 2006; Cisewski et al., 2008; Brainerd and Gregg, 1995) and  $10^{-5}$  m<sup>2</sup> s<sup>-1</sup> using an ocean general circulation model (Noh and Lee, 2008). In this study, XLD is specified as the depth where  $\varepsilon$  drops from its first minimum value. This aligns with previous dissipation thresholds and ensures that mixing is captured independently of surface influences. The Upper (UTD) and Lower (LTD/LPD) Thermocline/Pycnocline Depth are delimited as defined by Assunção et al (2020). UTD corresponded to the depth where the vertical temperature gradient  $\partial\theta/\partial z = 0.1$  °C m<sup>-1</sup>, while LTD/LPD were the last depth below the UTD at which  $N^2 \geq 10^{-4}$  s<sup>-2</sup>.

### Baroclinic currents

To analyze the processes explaining dissipation and mixing, particularly along internal tidal (IT) paths, we estimate shear instabilities associated with the semi-diurnal (M2) ITs and mean circulation, as well as their contributions to mixing.

The M2 tidal component of the tidal current is derived by calculating the baroclinic (semi-diurnal) tidal velocity  $[u'', v'']$  (Fig. A3, Appendix), following these equations:

$$[u', v'] = [u, v] - [u_{bt}, v_{bt}], \quad (1)$$

$$[u_{bt}, v_{bt}] = \frac{1}{H} \int_{-H}^0 [u, v] dz, \quad (2)$$

$$[u'', v''] = [u', v'] - [\bar{u}', \bar{v}']. \quad (3)$$

Here,  $[u, v]$  represent total horizontal velocities (Fig. A3, Appendix) obtained from SADC data. The components  $[u', v']$  and  $[u_{bt}, v_{bt}]$  represent baroclinic and barotropic components of horizontal velocities, respectively (Fig. A3, Appendix).  $H$  is water depth. The baroclinic mean velocities  $[\bar{u}', \bar{v}']$  (Fig. A3, Appendix),

calculated to estimate mean circulation along IT paths, are decomposed into along-shore  $\overline{u}_l$  and cross-shore  $\overline{u}_c$  velocities. The overbar denotes the average over a M2 tidal period.

Note that continuously collected SADCPC for some stations (e.g., S11) are not sufficiently resolved due to gaps filled by interpolating between time points. The similar processing are applied to the CTD-O<sub>2</sub> data collected alternately. SADCPC time series data are less than 17 hours at all long stations, except for S14, which spans 42 hours. As a result, the diurnal and semidiurnal period fittings are not formally distinct (except at S14; Figs. A4 and A5, Appendix), and the inertial period (at least 5 days) cannot be resolved in our dataset. This limits our ability to separate currents by frequency and examine the associated dissipation.

The velocity profiles from LADCP are glued into our SADCPC time series data below  $\sim 500$  m depth at long stations.

To evaluate shear instabilities associated with ITs and the mean background circulation, we compute the baroclinic tidal vertical shear squared ( $S^{2''}$ ) and mean shear squared ( $\overline{S^{2'}}$ ) (Fig. A3, Appendix), as follows:

$$S^{2''} = (\partial u''/\partial z)^2 + (\partial v''/\partial z)^2, \quad (4)$$

$$\overline{S^{2'}} = (\partial \overline{u'}/\partial z)^2 + (\partial \overline{v'}/\partial z)^2. \quad (5)$$

To evaluate the impact of bottom friction on mixing, we calculate kinetic energy  $\epsilon_f = \frac{1}{2}\rho_s(u_f^2)$  near the bottom boundary layer at shallow stations using friction velocity  $u_f = u_b\sqrt{C_d}$ , where  $C_d=2.5 \times 10^{-3}$  is a drag coefficient obtained from the NEMO model. Huang et al. (2019) showed that the bottom boundary layer thickness spatially varies between 15-123 m in the Atlantic Ocean, with a median of  $\sim 30$ -40 m in the North Atlantic. We define bottom layer thicknesses in our study area based on measured bathymetry from CTD-O<sub>2</sub> and near-bottom currents from ADCP. Here,  $u_b$  is the total velocity averaged over a thickness of 20 m above the seabed for shallow stations and 40 m for deep stations.

The individual contributions of semi-diurnal ITs and mean circulation are then expressed as follows:  $\overline{E''}/(\overline{E''} + \overline{E''})$  for tidal contribution and  $\overline{E'}/(\overline{E'} + \overline{E''})$  for mean circulation contribution. Here,  $E = N*S$ .  $N$  is the buoyancy frequency and  $S$  is vertical shear.  $S$  can be substituted by  $S^{2''}$  and  $\overline{S^{2'}}$ .

### Ray tracing calculation

Analyzing both the mean currents and the spatial dimension along the IT pathways offers another insight into the mechanisms responsible for observed mixing (Rainville and Pinkel, 2006). IT energy rays are generated in regions with steep topography, such as the shelf break, where IT slope matches with the bottom slope (i.e., critical slopes) before propagating within the ocean interior. These rays, moving both downward and upward, encounter the seasonal pycnocline, resulting in beam scattering and the formation of large IT oscillations. As these oscillations steepen, they disintegrate into nonlinear ISWs, a process known as "local generation" of ISWs (New and Pingree, 1992). To explore IT paths, ray-tracing techniques are employed, as previously used by New and Da Silva (2002) and Muacho et al. (2014), to investigate the effectiveness and expected pathways of the IT beams off the Amazon shelf. One main assumption in our linear-theory-based hypothesis is that stratification

remains horizontally uniform along the IT propagation path, although in reality, it may vary due to submesoscale and mesoscale variability. This limitation makes the ray tracing approach less realistic but still useful as a first-order estimate of energy distribution. The IT ray-tracing calculation assumes that in a continuously stratified fluid, ITs energy can be described by characteristic pathways of beams (or rays) with a slope  $c$  to the horizontal:

$$c = \pm \left( \frac{\sigma^2 - f^2}{N^2 - \sigma^2} \right)^{1/2}, \quad (6)$$

where  $\sigma$  is the M2 tidal frequency ( $1.4052 \times 10^{-4} \text{ rad s}^{-1}$ ), and  $f$  is the Coriolis parameter.  $N^2$  are obtained from time-averaged AMAZOMIX CTD- $O_2$ , glued with monthly  $N^2$  profiles from Amazon36 (NEMO model outputs, 2012-2016) below 1000 m depth. Amazon36 is a NEMO configuration, specifically designed to cover the western tropical Atlantic from the mouth of the Amazon River to the open sea (see Tchilibou et al., 2022; Assene et al., 2024; for configuration details and model description). IT ray-tracing diagrams are performed along the transects. Seasonal sensitivity tests of rays (August, September, October, and April) are conducted by varying the critical slope positions and  $N^2$  to explore its influence and generate a set of ray paths consistent with characteristics of IT pathways (Figs. A6 and A7, Appendix).

»

4. MacKinnon and Gregg (2003) provided validation in their paper. There needs to be some explanation of the advantages/conditions for this parameterization over those following from Gregg 1989. Gregg (2003) describe how latitude affects these mixing parameterizations. I saw no mention of this.

R: Thanks for your comments.

Please see the previous response to your comment n°7.

5. It is unclear how mixing layer depth is chosen. There is a statement about choosing a minimum but the profiles show a lot of variance.

R: Thanks for your remarks.

Mixed layer definitions are very sensitive depending on the criteria we adopt and the region of interest (e.g., Thomson and Fine, 2003; de Boyer Montégut et al., 2004; Cisewski et al., 2008; Noh and Lee 2008). We have performed a lot of tests and present them in the revision figure RC1.15.

In our study, we decided, since we have access to the direct microstructure profiles to calculate the mixing layer depth (XLD) defined as the depth at which dissipation estimates decrease from its first minimum value, ensuring that mixing is captured below the influence of surface dynamics. This definition of XLD is consistent with previous studies (e.g., Lozovatsky et al., 2006; Cisewski et al., 2008; Noh and Lee, 2008). As you mentioned, the dissipation profiles show a lot of variance, as does the XLD and the classical MLD (figure RC1.15). In the final manuscript we revised more carefully the definition of the XLD and the MLD and provided a new figure of comparison to show the sensitivity of all these layers.

We have revised the manuscript to clarify in section “methods/The vertical eddy diffusivity coefficient”, in lines 153–176 of the revised manuscript and reported here:

### The vertical eddy diffusivity coefficient

The efficiency of turbulence in redistributing energy is assessed through the calculation of the vertical eddy diffusivity coefficient ( $K_z$ ). This coefficient is particularly significant in regions such as pycnoclines, where stratification suppresses mixing, making turbulence-driven mixing a key mechanism for vertical energy transport (Thorpe, 2007).

$K_z$  is calculated from  $\epsilon$  following the formulation of Osborn (1980), given by  $K_z = \epsilon / \Gamma N^2$ . Here,  $N^2$  is the buoyancy frequency squared, which is calculated using the sorted potential density profiles ( $\sigma_\theta$ ) obtained from CTD-O<sub>2</sub> data. It is given by  $N^2 = - (g/\rho_0) (d\sigma_\theta/dz)$ , where  $\rho_0$  is a reference density (1025 kg m<sup>-3</sup>) and  $g$  is the gravitational acceleration.  $\Gamma$  is the mixing efficiency, defined as the ratio between the buoyancy flux and the energy dissipation, and is typically set to 0.2, which corresponds to the critical Richardson number  $Ri = 0.17$  (Osborn, 1980).  $\epsilon$  is linearly interpolated into the depths of  $N^2$ .

Turbulence within the pycnocline can reduce stratification and increase vertical eddy diffusivity below the mixing layer (Thorpe, 2007). Subsurface mixing, driven by the breaking of ITs and shear instabilities, plays a particularly important role below the mixed layer, especially in equatorial waters (Gregg et al., 2003).

There are several criteria for defining the Mixed Layer Depth (MLD). In this study, we use the commonly accepted density threshold criterion of 0.03 kg m<sup>-3</sup>, as defined by de Boyer Montégut et al. (2004) and Sutherland et al. (2014), to estimate the MLD for each CTD-O<sub>2</sub> profile. Notably, comparisons with density thresholds of 0.01 and 0.02 kg m<sup>-3</sup> revealed no major differences in MLD across the AMAZOMIX stations and transects (Fig. A2, Appendix).

The miXing Layer Depth (XLD) is defined as the depth at which  $\epsilon$  decreases to a background level (Sutherland et al., 2014). Previous studies have applied various thresholds for background dissipation levels, such as such as 10<sup>-8</sup> and 10<sup>-9</sup> W kg<sup>-1</sup> in higher latitudes based on in situ observations (Sutherland et al., 2014; Lozovatsky et al., 2006; Cisewski et al., 2008; Brainerd and Gregg, 1995) and 10<sup>-5</sup> m<sup>2</sup> s<sup>-1</sup> using an ocean general circulation model (Noh and Lee, 2008). In this study, XLD is specified as the depth where  $\epsilon$  drops from its first minimum value. This aligns with previous dissipation thresholds and ensures that mixing is captured independently of surface influences. The Upper (UTD) and Lower (LTD/LPD) Thermocline/Pycnocline Depth are delimited as defined by Assunção et al (2020). UTD corresponded to the depth where the vertical temperature gradient  $\partial\theta/\partial z = 0.1$  °C m<sup>-1</sup>, while LTD/LPD were the last depth below the UTD at which  $N^2 \geq 10^{-4}$  s<sup>-2</sup>.

6. It is stated the mixing layer depth is always less than the mixed layer depth. This seems like it cannot be true by definition- you need mixing to make a mixed layer. Maybe the mixed layer depth criterion is too small. How about 0.1 kg m<sup>-3</sup>?

R: Thanks for your comments.

Indeed, you need mixing to make a mixed layer. There are several criteria for defining the mixed layer depth (MLD). We have tested different density threshold criterion from 0.01 to 0.03 kg m<sup>-3</sup>, which are the commonly accepted density threshold used in previous studies (e.g., Thomson and



Fine, 2003; de Boyer Montégut et al., 2004; Cisewski et al., 2008; Noh and Lee 2008) as well the less used threshold and from 0.1 to 0.3 kg m<sup>-3</sup> (e.g., Brainerd and Gregg, 1995). We have provided on figure RC1.15 and RC1.16, a calculation of these layers for each station. We can see that the XLD is typically deeper than the MLD, which is calculated using a density threshold 0.01-0.03 kg m<sup>-3</sup>, except for stations 8,10 and 25. This is consistent with regions exhibiting strong subsurface shear, such as the equatorial ocean and western boundary current areas (Noh and Lee, 2008). In contrast, XLD can be lesser than MLD during early restratification and at high latitudes during convective cooling (Noh and Lee, 2008). In our case, the exception for stations 8,10 and 25 (MLD>XLD) might be attributed to shelf break dynamics, including internal tides.

We have added the figure RC1.15 in Appendix of manuscript and pointed out in lines 261-263 and 461-464 of the revised manuscript and reported here:

Lines 261-263:

“

It is important to note that the XLD is typically deeper than the MLD at all stations (except at S8, S10, and S25), which is calculated using a density threshold 0.01, 0.02 or 0.03 kg m<sup>-3</sup> (Fig. A2, Appendix).

“

Lines 461-464:

“

The XLD was found to be considerably larger than the MLD at all stations, except at S8, S10, and S25. This is consistent with regions exhibiting strong subsurface shear, such as the equatorial ocean and western boundary current areas (Noh and Lee, 2008). The exception observed at other stations may reflect larger mixing events that were not captured by the VMP measurements.

“

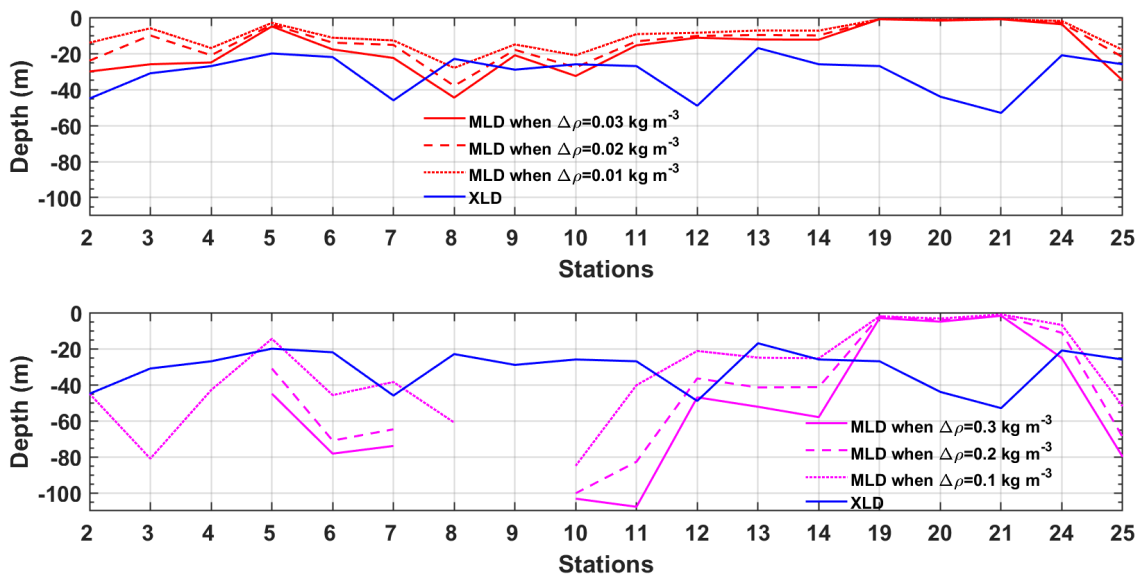


Figure RC1.15: Comparison of Mixing Layer Depths (XLD, blue line) with Mixed Layer Depths (MLD) defined using (a) larger and (b) smaller density thresholds ( $\Delta\rho$ ). In panel (a), dotted, dashed, and solid red lines represent MLDs defined by  $\Delta\rho = 0.01, 0.02, 0.03 \text{ kg m}^{-3}$ , respectively. In panel (b), dotted, dashed, and solid magenta lines represent MLDs defined by  $\Delta\rho = 0.1, 0.2, 0.3 \text{ kg m}^{-3}$ , respectively.

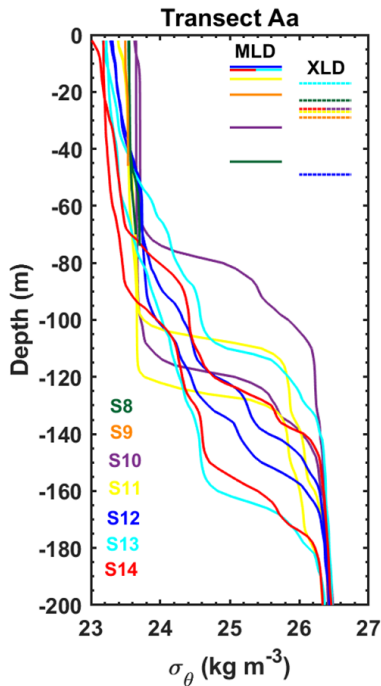


Figure RC1.16: Density profiles at selected stations (S8 to S14) overlaid MLD (in colored solid lines) and XLD (colored dashed lines). Distinct colors are used to represent each station.

7. Figure caption does not correspond to what is plotted. Shear is missing? Figure labelling: how about (a), (b), ... instead 4.a.1? There are also faint grey lines around the figures. All the dots in the shear vs N figures are hard to see. Consider binning the results.

R: Thanks for your comments and suggestions.

We have reorganized and revised all the figures and captions (e.g., Figures 5, 6, and 7) in the revised manuscript to make them clearer. We also added the shear (e.g., Figure 6 and 7) in the revised manuscript and reported here in figure RC1.17 and RC1.18 for selected stations (e.g., S10, S11, S14).

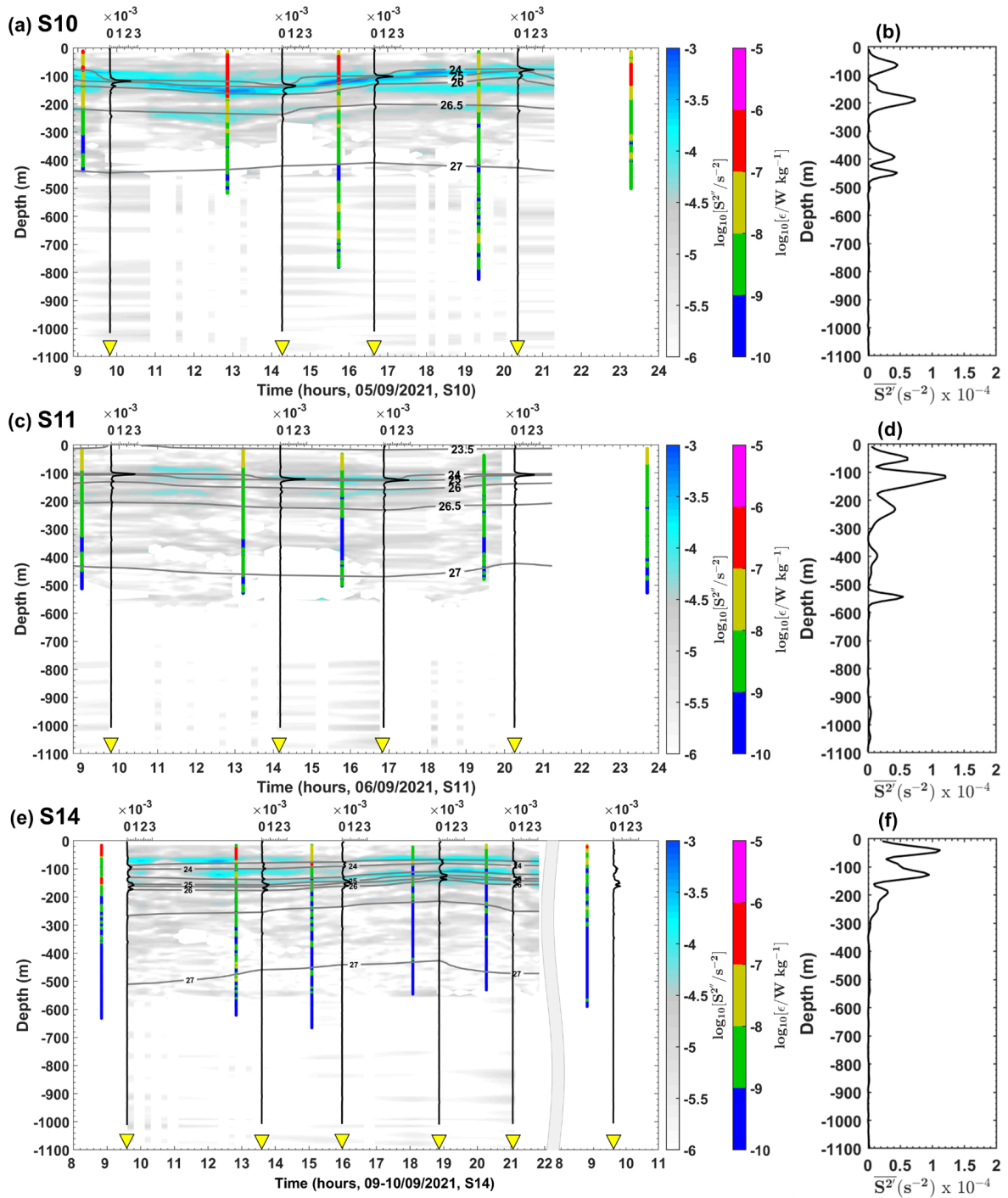


Figure RC1.17: Semi-diurnal baroclinic vertical shear squared ( $S2''$ , in  $m\ s^{-1}$ , on a logarithmic scale) for stations (a) S10, (c) S11, and (e) S14. Panels (a), (c), and (e) also display the buoyancy frequency squared ( $N2$ , in  $s^{-2}$ ) represented by vertical black lines, potential density represented by grey contours, and dissipation rate profiles ( $\epsilon$ , in  $W\ kg^{-1}$ , on a logarithmic scale) represented by vertical colored bars. Mean baroclinic vertical shear squared ( $S2'$ , in  $m\ s^{-1}$ ) for stations (b) S10, (d) S11, and (f) S14.

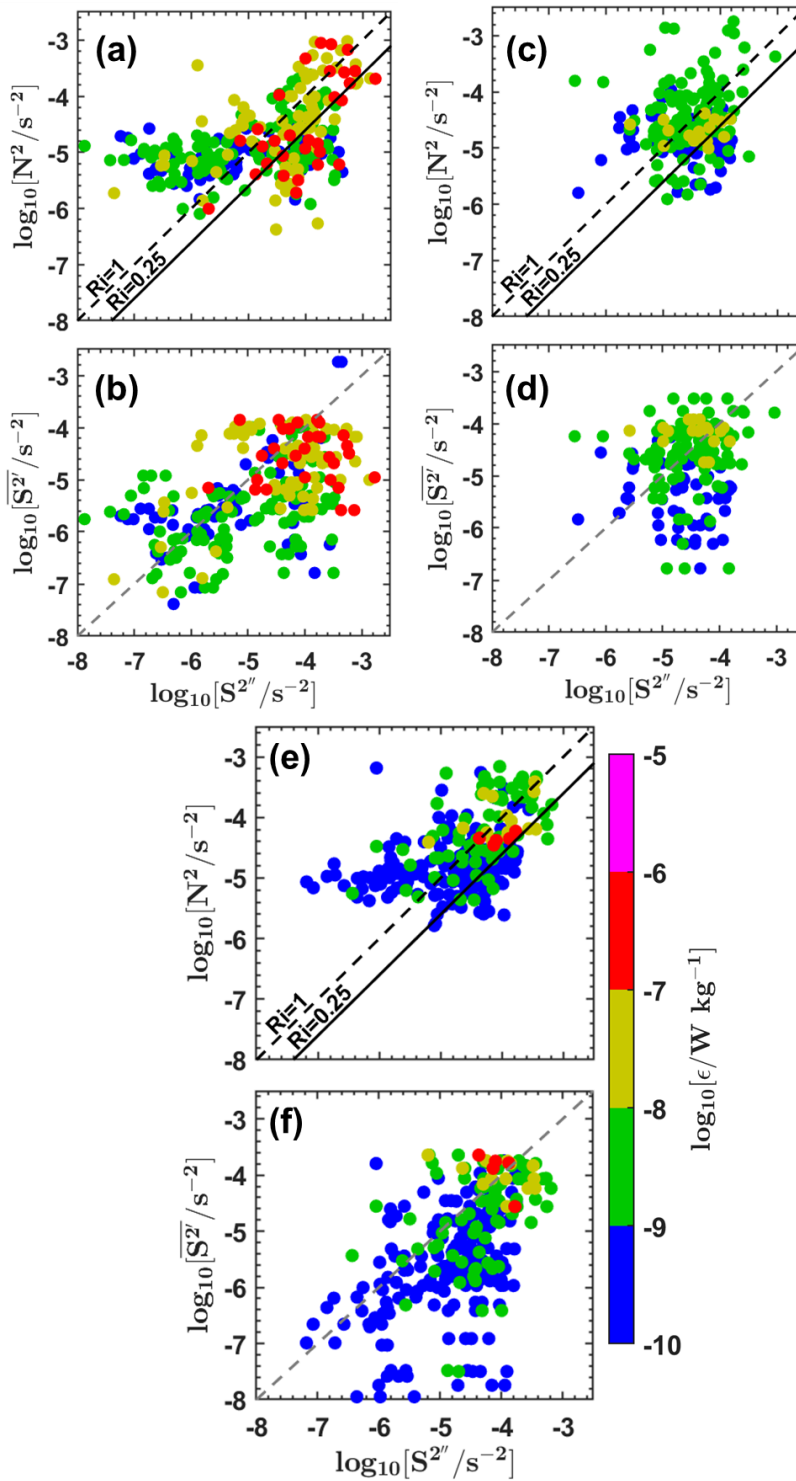


Figure RC1.18: Dissipation rates ( $\epsilon$ , in  $W \text{ kg}^{-1}$ , on a logarithmic scale) below the XLD as a function of the buoyancy frequency squared ( $N^2$ , in  $s^{-2}$ , on a logarithmic scale) and semi-diurnal baroclinic vertical shear squared ( $S^{2''}$ , in  $m \text{ s}^{-1}$ , on a logarithmic scale) for stations (a) S10, (c) S11, and (e) S14. Dissipation rates ( $\epsilon$ , in  $W \text{ kg}^{-1}$ , on a logarithmic scale) below the XLD as a function of mean baroclinic vertical shear squared ( $\overline{S^2}$ , in  $m \text{ s}^{-1}$ , on a logarithmic scale) and semi-diurnal baroclinic vertical shear squared ( $S^{2''}$ , in  $m \text{ s}^{-1}$ , on a logarithmic scale) for stations (b) S10, (d) S11, and (f) S14.  $N^2$  was linearly interpolated into the depths of  $S^2$  to have same vertical scales. Panels (a), (c), and (e) also display two solid black lines corresponding to Richardson

*number  $R_i = 0.25$  and  $R_i = 1$ , respectively. Dashed grey lines in panels (b), (d), and (f) are included for comparison purposes.*

8. There are lots of seemingly correct facts and figures in the text but I am unclear what is the point of them. Take section 3.2.2, for example. Currents are going this way and that at various depths and locations. True. I would suggest a format for each paragraph as follows. "In this paragraph, we examine the flow patterns over the slope to show something [topic sentence]. Here are the relevant facts and skip ones not immediately related to the topic sentence. In summary, we have shown flow is in this direction here which is important for something. Something is described in the next paragraph."

R: Thanks for your suggestions. All sections of manuscript have been revised.

### Comments by line number

19 - twelve hours is the semidiurnal period but not the M2 period = 12.42 hrs

R: Thanks for your remarks. We have corrected it in the revised manuscript in line 19 of the revised manuscript, as shown below:

"

During the AMAZOMIX survey (2021), currents, hydrography, and turbulence were measured over M2 tidal period (12.42h) at numerous sites near the Amazon outflow, where ITs are also generated along the slope.

"

129 - `fft_length`, etc. Perhaps you could just give these a symbol if they come up more than once, e.g., L or n. Otherwise skip the variable names in a manuscript.

R: Thanks for your remarks. We have corrected it in the revised manuscript in lines 181-187.

"

Shear spectra are estimated using record lengths (L) and Fast Fourier Transform segments of 2 s, which are cosine windowed and overlapped by 50% (e.g., at station S6 and S10; Fig. 2b, and Fig. A1, Appendix). Additionally, vibration-coherent noise is removed. Different L and overlap (O) settings were selected and tested based on the environment (e.g., deep vs. shallow water), following Fer et al. (2024). For shallow stations, L (O) was shortened to 5 s (2.5 s), in contrast to the 8 s (4 s) used for deeper stations, due to evidence of overturns observed in AMAZOMIX acoustic measurements at deeper stations (Koch-Larrouy et al., 2024; in preparation). This adjustment helped to optimize the spatial resolution of dissipation estimates in shallow water stations.

"

136 - what is the high pass? and the low pass?

R: The high-pass filter was applied to the shear probe data and the low-pass filter was applied to the temperature gradients. We have rewritten this subsection “TKE dissipation rates” in the revised manuscript.

150 - mixing layer depth

R: Thanks for your remarks. We have corrected it in the revised manuscript in line 170 of the revised manuscript.

170 - I don't understand. Please rephrase

R: Thanks for your remarks. We have rewritten it in line 198 of the revised manuscript, as shown below:

“

The velocity profiles from LADCP are glued into our SADCP time series data below ~ 500 m depth at long stations.

“

172 - what is H?

R: Thanks for your remarks.

H is the water depth. We have revised it in the revised manuscript in lines 188-189 of the revised manuscript.

193 - references? Also strong internal tides propagating upward and impinging on the thermocline make ISW. No need for bottom reflections.

R: Thanks for your remarks.

Indeed, strong internal tides propagating upward and impinging on the thermocline generate ISWs. We have revised the subsection “Ray Tracing Calculation” to clarify this, in lines 219–221 of the revised manuscript, as reported here:

“

These rays, moving both downward and upward, encounter the seasonal pycnocline, resulting in beam scattering and the formation of large IT oscillations. As these oscillations steepen, they disintegrate into nonlinear ISWs...

“

200- You have cross-shore measurements of N. Are those not enough to make a horizontally varying N?

R: Thanks for your comments.

Indeed, we have cross-shore measurements of  $N$ . We have varied  $N^2$  profiles from stations (e.g., S10, S12, and S14) along transect Aa. Please see the results in figure RC1.19 and RC1.20 below. The sensitivity tests (Figure RC1.19) showed that ray paths align within the packets of rays observed when using mean  $N^2$  profiles along transect Aa at different times (e.g., in September and October; Figure RC1.20).

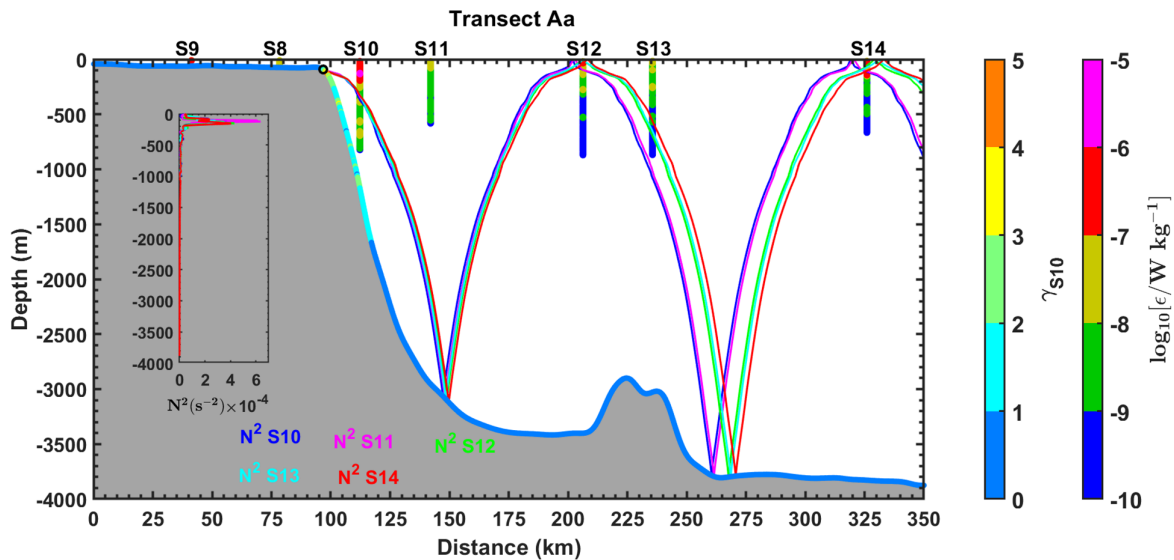


Figure RC1.19: Example of sensitivity tests with different cross-sectional measurements of  $N^2$  along the transect T1  $N^2$ . colors are used to distinguish different cross-shore measurements of  $N^2$  for corresponding stations on T1. Topography steepness ( $\gamma = \text{ray slope} / \text{topography slope}$ ) for T1 using measured  $N^2$  of S10. Gamma is illustrated by the colored bar (named gamma S10).

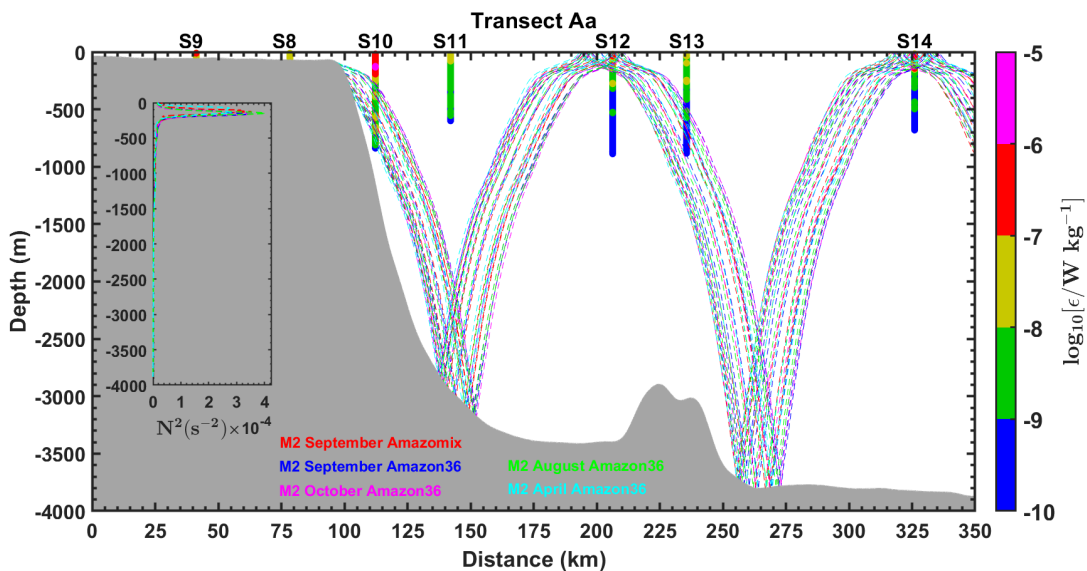


Figure RC1.20: Sensitivity tests of M2 IT ray-tracing along the transects Aa, conducted by varying the location of the critical topography slope. The tests use mean buoyancy frequency squared ( $N^2$ , in  $\text{s}^{-2}$ ) obtained from CTD- $\text{O}_2$  data (September 2021) and NEMO-Amazon36 model data (2012-2016). Dashed colored lines represent IT beams calculated for different seasons (April, August, October, and September) and for varying locations of the critical topography slope. Grey areas indicate local topography. Panel also includes dissipation rate profiles ( $\epsilon$ , in  $\text{W kg}^{-1}$ , shown as vertical colored bars on a logarithmic scale) from

the VMP measurements. Subpanels within each panel illustrate the  $N^2$  profiles derived from AMAZOMIX and the NEMO-Amazon36 model, which were used in the ray-tracing calculations. For comparison, sensitivity tests using  $N^2$  measurements from individual stations along the corresponding transect (e.g., at S10) revealed similar ray paths (not shown), consistent with the packet of rays obtained using the mean  $N^2$ .

207- there has to be some minimal description of amazon36 here even if it is referenced

R: Thanks for your remarks.

We have added a short description of amazon36 in the revised manuscript (in lines 231-233) and refer to previous studies. We have reported the revision here:

“

Amazon36 is a NEMO configuration, specifically designed to cover the western tropical Atlantic from the mouth of the Amazon River to the open sea (see Tchilibou et al., 2022; Assene et al., 2024; for configuration details and model description).

“

210- sensitivity not sensibility test, although the latter could also be used

R: Thanks for your remarks. We have revised it in line 233 of the revised manuscript.

263- 3-5 not 03-05

R: Thanks for your remarks. We have revised it in lines 307-308 of revised manuscript, as shown below:

“

The baroclinic tidal velocities reveal a superposition of 3-5 tidal modes at IN-ITs stations (Figs. 5a, 5c, and 5e, and Figs. A9 to A15, Appendix).

“

267-272 - Shear obviously varies with  $N$  because if shear were bigger than it would have lower  $Ri$  and be unstable. It would be more instructive to consider  $Ri$  or reduced shear =  $S2 - 4N2$ .

R: Thanks for your comments.

Indeed, shear varies noticeably with  $N$  at certain stations farther from IT generation sites. However, closer to the generation sites, such as at S10, stronger tidal shears were observed in areas where  $Ri$  was lower ( $<0.25$ ), coinciding with hotspots of mixing.



This is why we specifically considered shear and quantified the contribution of tidal shear to mixing along IT paths. Please refer to the results for reduced shear presented below in figure RC1.21.

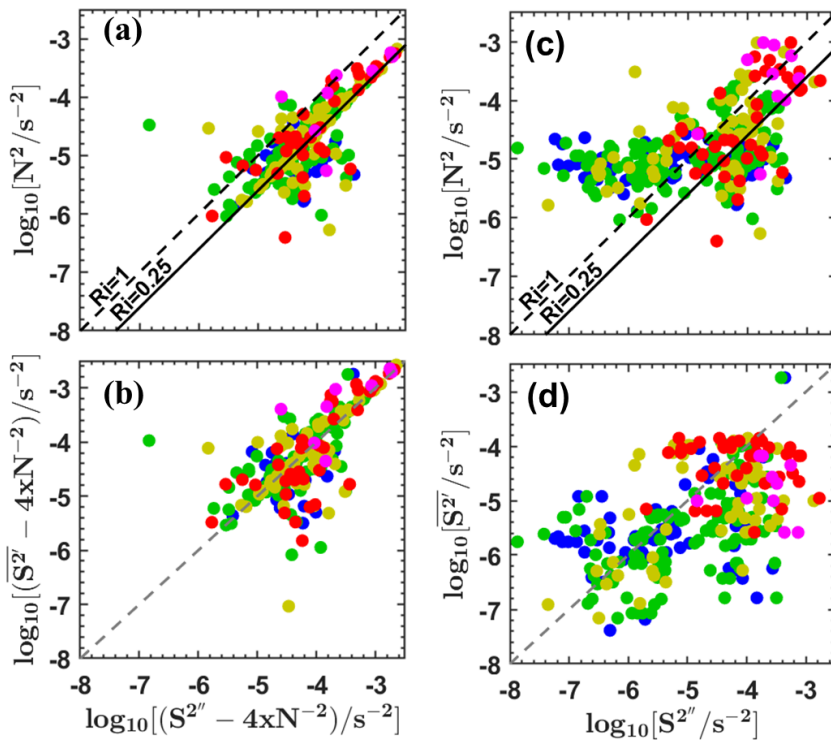


Figure RC1.21: For station S10. Dissipation rates ( $\epsilon$ , in  $W \text{ kg}^{-1}$ , on a logarithmic scale) below the XLD as a function of the buoyancy frequency squared ( $N^2$ , in  $\text{s}^{-2}$ , on a logarithmic scale) and semi-diurnal baroclinic vertical shear squared ( $S^{2''}$ , in  $\text{m s}^{-1}$ , on a logarithmic scale), when using (a)-(b) reduced shear. Dissipation rates ( $\epsilon$ , in  $W \text{ kg}^{-1}$ , on a logarithmic scale) below the XLD as a function of mean baroclinic vertical shear squared ( $\overline{S^2}$ , in  $\text{m s}^{-1}$ , on a logarithmic scale) and semi-diurnal baroclinic vertical shear squared ( $S^{2''}$ , in  $\text{m s}^{-1}$ , on a logarithmic scale), when using (c)-(d) no reduced shear.  $N^2$  was linearly interpolated into the depths of  $S^{2''}$  to have same vertical scales. Panels (a), (c), and (e) also display two solid black lines corresponding to Richardson number  $Ri = 0.25$  and  $Ri = 1$ , respectively. Dashed grey lines in panels (b), (d), and (f) are included for comparison purposes.

You could use some mean N or interpolated N and see what it looks like.

R: Indeed, several tests were conducted using both mean and interpolated  $N$ . Ultimately, we opted to use the time-interpolated  $N$  and shear values at all dissipation times (at VMP measurement hours) along IT paths.

If the text talks a lot about shear, it would be helpful to plot shear. In fact the baroclinic vs tidal shear is discussed but there are no plots.

R: Thanks for your remarks. We have added the plots of the shears both in text and appendix.

Fig 4- caption and figure all on same page please. Add tick marks to the x axis to show when your CTD/VMP casts took place. Also this figure seems to include transits and time on station.

In Fig 4b1 there are some strange looking changes in the current. Perhaps it would be better to do your analysis station by station. Or are there artifacts of the SADCP processing going from on station to transit?

R: Thanks for your comments.

We have reorganized the figures and added an x-axis to indicate when CTD/VMP casts were conducted. Please refer to the updated figure in the revised manuscript (figure 5) and reported below in figure RC1.22.

Regarding the updated figure of the revised manuscript (figure 5), we noticed some unusual changes in the current patterns, which have been rechecked. The continuously collected SADCP for some stations (e.g., S11) are not sufficiently resolved due to gaps filled by interpolating between time points (e.g., between 13:30 and 15:30 at S11). Doing this interpolation would introduce potential artifacts in the calculations. The interpolation allowed us to extract some current profiles at the dissipation times to ensure consistency in the analyses.

The analyses were performed separately by station and by transect. For example, Figure RC1.22 shows the semidiurnal current over time at a fixed position for station S11. In contrast, Figure RC1.23 displays the mean alongshore currents along transect Aa, including transit between stations.

We have pointed out in text [in lines 192-193 of revised manuscript and reported here](#)

“

Note that continuously collected SADCP for some stations (e.g., S11) are not sufficiently resolved due to gaps filled by interpolating between time points. The similar processing are applied to the CTD-O<sub>2</sub> data collected alternately.

“

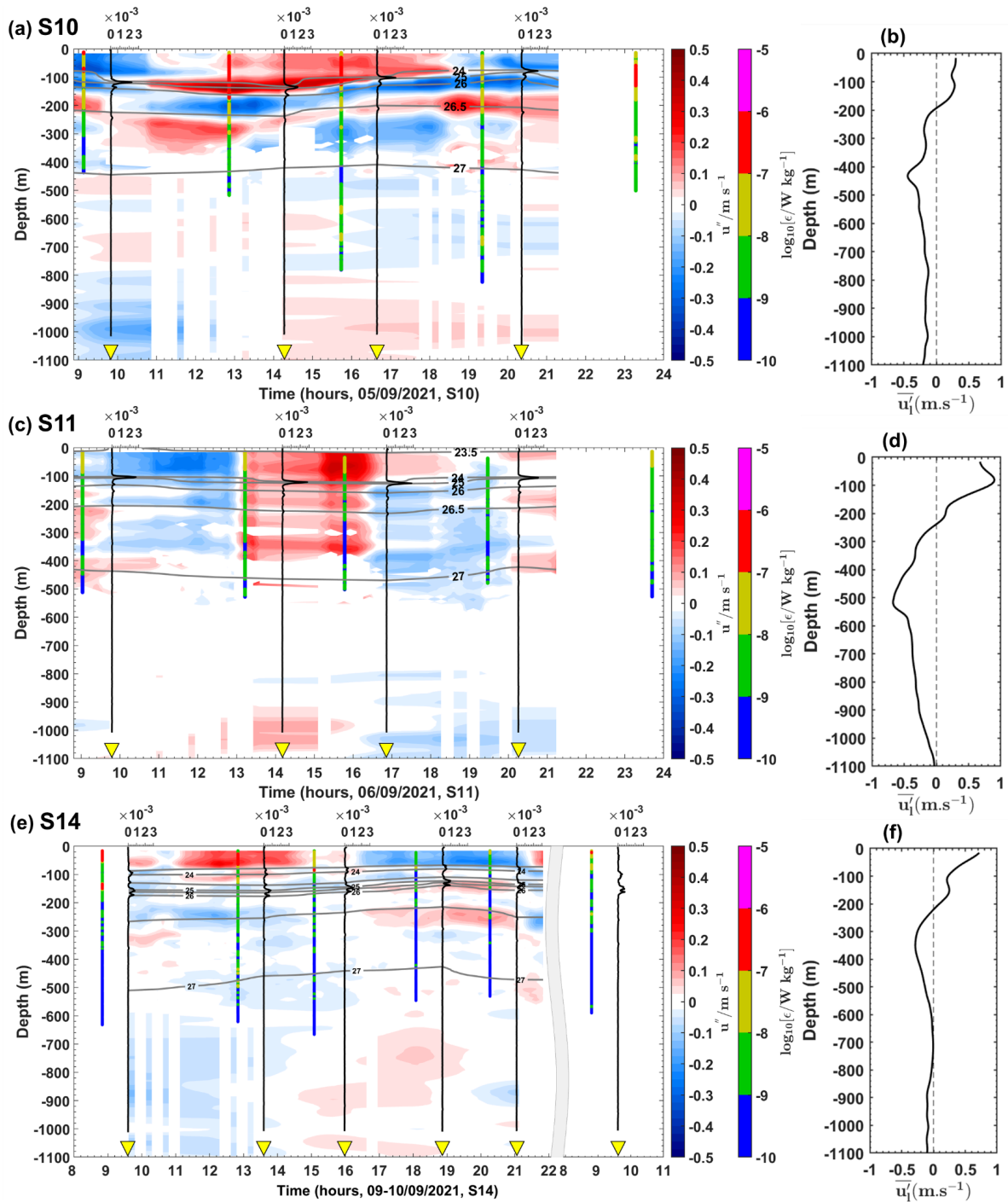


Figure RC1.22: Semi-diurnal baroclinic zonal currents ( $u''$ , in  $\text{m s}^{-1}$ ) from the ADCP for stations (a) S10, (c) S11, and (e) S14. Panels (a), (c), and (e) also display the buoyancy frequency squared ( $N^2$ , in  $\text{s}^{-2}$ ) represented by vertical black lines, potential density represented by grey contours, and dissipation rate profiles ( $\epsilon$ , in  $\text{W kg}^{-1}$ , on a logarithmic scale) represented by vertical colored bars. Along-shore mean baroclinic currents ( $u'I$ , in  $\text{m s}^{-1}$ ) from the ADCP for stations (b) S10, (d) S11, and (f) S14.

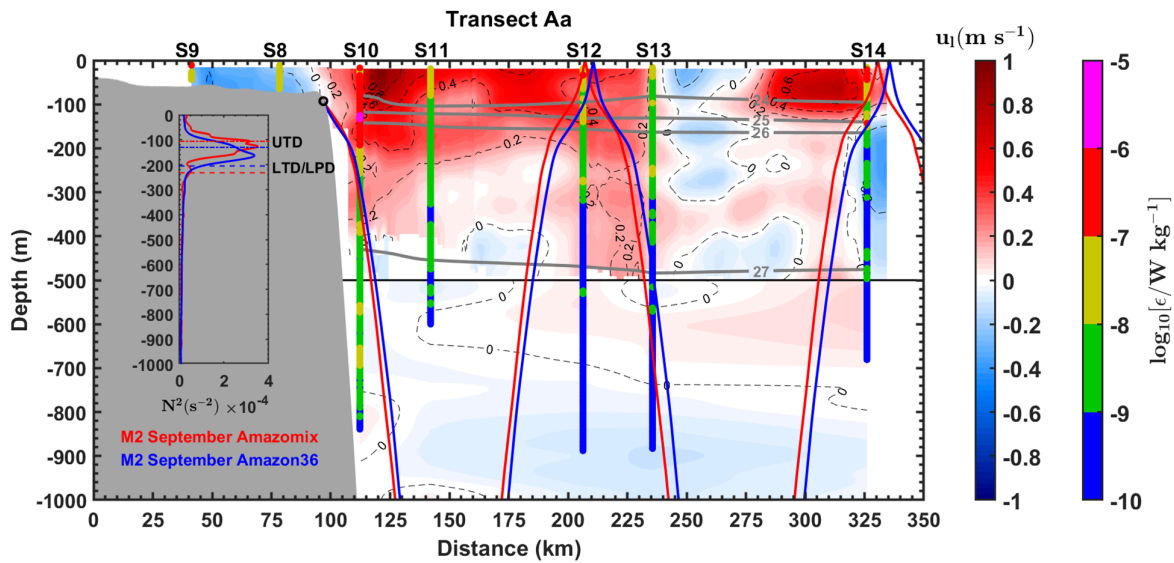


Figure RC1.23: IT ray-tracing diagrams for the M2 tidal constituent along transects Aa. The calculations were performed using the mean buoyancy frequency squared ( $N^2$ , in  $s^{-2}$ ) obtained from CTD- $O_2$  data (ray in red) and NEMO-Amazon36 model data (ray in blue) for September. Grey areas represent local topography and black circles indicate the critical topography slope (ray generation sites). Panel also show along the transects Aa: along-shore mean total currents ( $u_1$ , in  $m s^{-1}$ ) from ADCP (Dashed black lines), potential density from CTD- $O_2$  (grey contours), and dissipation rate profiles ( $\epsilon$ , in  $W kg^{-1}$ , on a logarithmic scale) from the VMP (vertical colored bars). Subpanels within each panel illustrate the  $N^2$  profiles from AMAZOMIX (red line) and the NEMO-Amazon36 model (blue line) used for ray-tracing calculations. Upper Thermocline Depth (UTD, dotted lines) and Lower Thermocline/Pycnocline Depth (LTD/LPD, dashed lines) are also indicated.

301- what sort of of mean? time, depth, etc

R: Thanks for your remarks. It is time-mean. We have revised it in line 362 of the revised manuscript

312-314- More explanation is needed here.

R: These have been removed during the revision of the manuscript.

Fig 5- the tidal rays look to have shallower slope than the topography. It would be helpful to plot ray slope vs topographic slope to verify. The dissipation seems unrelated to the ray paths as far as I can tell. Perhaps it is not 2D. There appear to be 2 sources at least, which can constructively interfere. Rays may propagate at an angle to the coast.

R: Thanks for your comments.

We attempted to verify our analysis by plotting the ray slope against the topographic slope. Using our bathymetry data from the NEMO-AMAZON36 model or GEBCO, the tidal rays did not appear to have a shallower slope (around 225 km) than the topography in any of the sensitivity tests, as illustrated in Figure RC1.24 (with topography steepness "gamma") below. These results may differ if using alternative bathymetric products.

Following the ray tracing, we observed that some dissipation hotspots are located along the ray paths, particularly above the generation sites and during ray propagation.

Interestingly, there seem to be at least two sources that could constructively interfere, potentially explaining the high dissipation observed at S14. This wave-wave interference is being explored further in a separate study to confirm these findings.

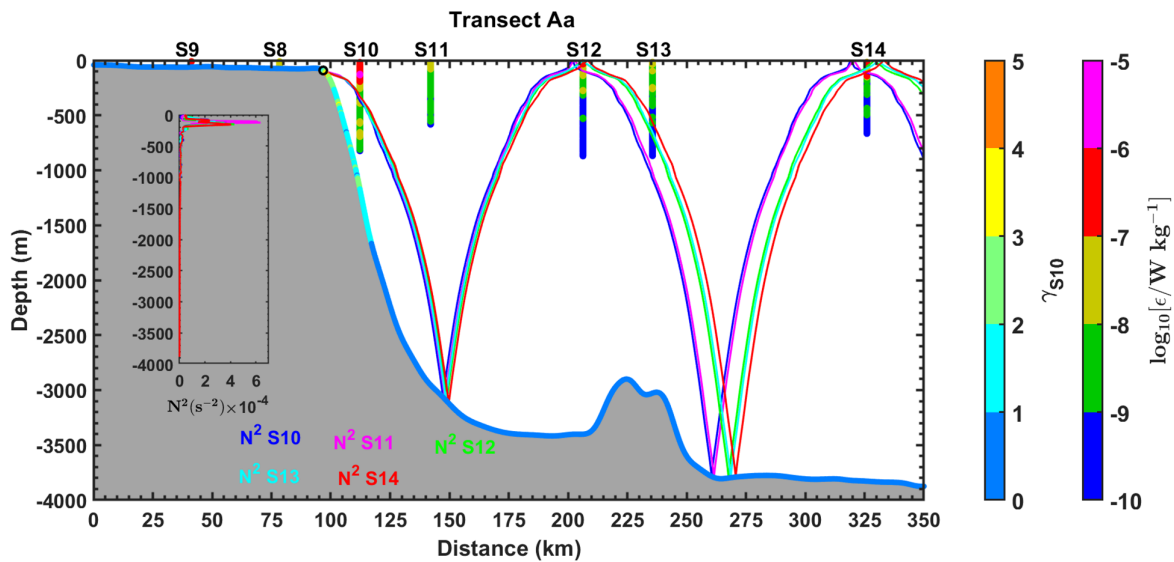


Figure RC1.24: Example of sensitivity tests with different cross-sectional measurements of  $N^2$  along the transect T1  $N^2$ . colors are used to distinguish different cross-shore measurements of  $N^2$  for corresponding stations on T1. Topography steepness ( $\gamma = \text{ray slope} / \text{topography slope}$ ) for T1 using measured  $N^2$  of S10. Gamma is illustrated by the colored bar (named gamma S10).

384 - “for the first time” How sure are you of this? It would be conservative to add: to the best of our knowledge

R: Thanks for your remarks.

We have corrected the sentence of the text in line 436 of the revised manuscript, as follows:

“

To the best of our knowledge, the AMAZOMIX 2021 cruise provided, for the first time, direct measurements of turbulent dissipation using a velocity microstructure profiler....

”