



Advancing interpretation of incoherent scattering in ice penetrating radar data used for ice core site selection

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8 Abstract. Below the coherent layering in ice penetrating radar data collected in Antarctica and Greenland, incoherent 9 scattering is common. This scattering is signal, not noise, and has the potential to inform our understanding of the structure and dynamics of the bottom 20% of glaciers and ice sheets. Here, we present a comparison between radar imagery and ice core 10 11 properties for sixteen ice core sites across Antarctica and Greenland, to identify possible sources for incoherent scattering and evaluate its use in ice core site selection. We find that incoherent scattering is commonly coincident with either gradual changes 12 in crystal orientation fabric or rapidly fluctuating fabrics in deep ice, where strain is localized by strength differences associated 13 14 with ice grain size. Macro-scale deformation and layer folding at scales below the range-resolution of radar does not seem to 15 result in incoherent scattering or induce an echo free zone, as has been previously hypothesized. Where incoherent scattering 16 is laterally homogeneous in intensity, layering is typically undisturbed in nearby ice cores. But where incoherent scattering is 17 laterally heterogeneous in intensity and the trace of intensity maxima does not appear conformal with subglacial topography, 18 we find multi-meter-scale folding and associated discontinuities in nearby ice core records. Future, higher-resolution sampling 19 of fabric in ice cores would allow for more quantitative interpretation of incoherent scattering and its amplitude, but we show 20 that the qualitative nature of incoherent scattering has the potential to inform us about the continuity of climate records at 21 prospective ice core sites and should be considered when evaluating the nature and quality of basal ice.

22 1. Introduction

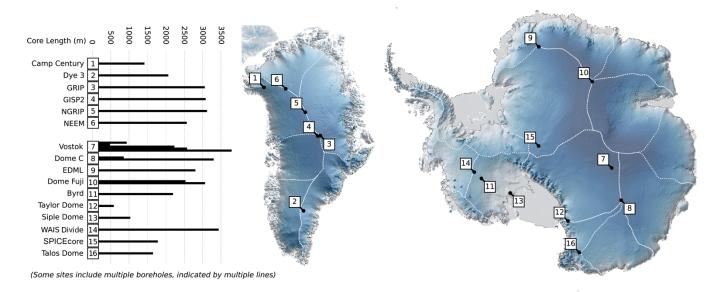
23 Existing ice cores provide our best record of past atmospheric chemistry. And while they capture global climate changes over 24 the Holocene and Late Pleistocene well (Wolff et al., 2010), future ice coring initiatives hope to build on that record, both 25 extending it further back in time (Jouzel and Masson-Delmotte, 2010) and measuring regional climate change (Mulvaney et 26 al., 2021) during specific climate periods (Fudge et al., 2023). These projects require the identification and collection of very 27 specific ice, and so they typically start with extensive geophysical surveying for "site selection" preceding drilling. Ice 28 penetrating radar data have served as the primary tool for this work, using layering in radar imagery to infer spatially variable 29 accumulation and ice flow, and through that, identifying ideal ice core sites (Schroeder et al., 2020). But site selection has 30 relied primarily on the strong, coherent signal that spans the shallow portions of most radar imagery. Here we focus on





31 improving interpretation of other signals in radar data, with a particular focus on what incoherent scattering (described in 32 section 2) can tell us about ice near the ice sheet base.

All radio-wave scattering originates from electrical contrasts. To better understand the nature and sources of scattering in 33 34 existing ice penetrating radar data, several previous studies have compared radar imagery to observations of ice chemistry and 35 physical properties measured in ice cores (e.g., Eisen et al., 2003, 2007; Mojtabavi et al., 2022). But that work has been focused 36 on the coherent, isochronal layering, and comparatively little has been done to understand the deeper signals, which are 37 becoming better sampled with modern, high power / low noise systems. This deep ice has also become increasingly 38 scientifically important, as it is at the center of the search for an ice core record that spans the Mid-Pleistocene transition (Lilien 39 et al., 2021). Using data from 16 ice cores (fig. 1), we work to better understand the physical properties that produce deep, 40 incoherent scattering, and evaluate the extent to which it may be diagnostic of layer disturbances or other disqualifying 41 characteristics when pursuing future ice cores.



- 42
- 43 Figure 1: Locations of major ice coring initiatives in Greenland and Antarctica used in this study. On maps of Antarctica and
- 44 Greenland, ice divides are marked in white.

45 2. Background: Scattering and the Radar Imaging Problem

46 Radar systems actively transmit energy into the subsurface. Time-of-flight measurements for back-scattered energy (together 47 with a known speed of light in ice) can be used to infer the position of subsurface scatterers and reconstruct the geometry of 48 glacier systems (Dowdeswell and Evans, 2004). In the near sub-surface, contrasts in the dielectric permittivity that scatter





energy are controlled primarily by variations in density, while most deeper englacial reflectors arise from either conductivity contrasts due to variations in the concentration of free ions deposited with the snow at the surface (Stillman et al., 2013) or transitions in the ice crystal fabric, typically localized by changes in grain size also arising from impurity deposition (Fujita et al., 1999). Fabric induced scattering is a product of the dielectric anisotropy of individual ice crystals, with transitions in caxis fabric leading to an (up to) 1% contrast in the polarization-dependent bulk permittivity. Incoherent scattering may come from both chemical and physical of sources; we work to provide some of the first constraints on its origins here.

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Glaciologists primarily use radar data for ice core site selection in two ways. The first approach is focused on the geometry of coherent, isochronous layering within the ice sheet. These layers originate as flat-lying snow at the ice sheet surface and are transformed by flow during burial; thus, their geometry can be used to diagnose spatial variations in accumulation (Karlsson et al., 2020), glacier sliding (Leysinger Vieli et al., 2007), and basal melt (Fahnestock et al., 2001). The second approach is focused on the nature of subsurface scattering (its coherence, distribution, and amplitude), which can be used to infer the modern electrical (and, more generally, material) characteristics of the ice sheet and its substrate (Schroeder et al., 2020).

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63 Subsurface targets can be divided into two main categories: specular interfaces and rough (or diffuse) scatterers (Schroeder et al., 2015). Specular interfaces, like mirrors, scatter energy in one dominant direction, a function of the direction-of-arrival for 64 65 the incoming light and the orientation of the interface. Diffuse scatterers redistribute incident energy at a variety of angles. 66 This leads to significant differences in the coherence of the scattering between specular and diffuse targets (defined here as the 67 consistency in phase and amplitude of the backscattered energy with slight changes in the position of the radar system). Incoherent scattering typically occurs at rough interfaces or when there are multiple diffuse scattering targets in the subsurface. 68 69 It has been observed as a product of rare glacier conditions, for example, where there is significant temperate ice and associated 70 englacial water (Hamran et al., 1996) or where debris has been entrained near the base of glaciers (Winter et al., 2019). But it 71 must also be generated by more common glaciological phenomena, as it is present within several hundred meters of the ice 72 sheet base across large parts of Antarctica and Greenland.

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74 Consider the example radar image in figure 2.a. Each pixel represents either backscattered energy or electrical or thermal noise 75 in the radar electronics. The position of the radar system varies across the columns in the image, and the delay-time after the 76 transmitted pulse (associated with the range to possible targets) varies across the rows in the image. In regions dominated by 77 planar, specular interfaces (as in the upper half of fig. 2.a), each pixel typically represents backscattered energy from only a 78 single subsurface target. This is because, even though there are many scattering targets at the range associated with that pixel 79 (as shown in fig. 2.b), only that interface tangential to the range shell (such that the interface is normal to the propagating 80 wave) returns energy to the system. But in regions where there are diffuse scatterers, each pixel in a radar image represents the 81 interference of scattering from multiple targets (fig. 2.b.ii/iii). With slight changes in the position of the system, the dominant 82 source of scattering at a given range can change, resulting in little consistency in phase or amplitude from pixel to pixel. This





is extremely common for energy arriving below the ice bottom reflector, with a long tail of incoherent scattering appearing at greater range (fig. 2.a.iii). Less well described is incoherent scattering from within the ice column (fig. 2.a.ii) which is the focus of our research here.

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When considering the nature of scattering in radar imagery, it is important to remember that the images themselves are ultimately a product of three things:

- 89 1. The geometry and physical/electrical characteristics of the glacier subsurface.
- 2. The system used to collect the data (including the characteristics of the transmitted wave, antennas, and
 transmit/receive electronics).
 - 3. The filtering, focusing, and additional image processing algorithms applied after collection.

The nature of radar targets depends on both the scale of electromagnetic heterogeneity in the medium and the frequency content of the transmit pulse (with higher frequencies / bandwidths associated with finer range resolution). Figure 2.c demonstrates how the same targets manifest differently across different radar systems; with lower resolution systems, scattering appears more structured, like the specular and coherent layering in the shallow ice. Because some systems used for site selection do not preserve phase information, we focus here primarily on the amplitude and character of scattering, controlling for differences in system characteristics.



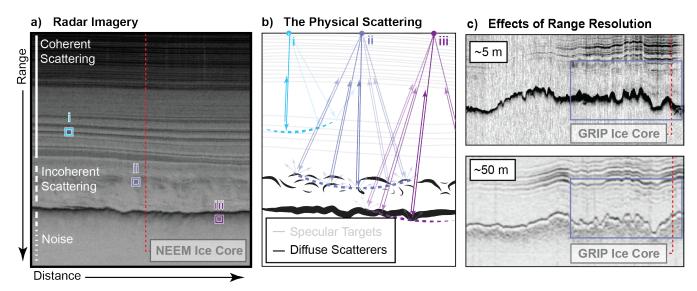


Figure 2: Example radar image (a), the ray-paths associated with scattering targets that contribute to individual pixels in the radar imagery (b), and a pair of images highlighting the effect of system characteristics on the nature of deep scattering (c).

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104 To generate incoherent scattering, deep ice must differ from the planar, layered structure of the shallow ice column in some 105 way. It may be that incoherent scattering occurs because chemical layering is mechanically disturbed in the deep ice and is no 106 longer planar. Or, it may be that other processes (like dynamic recrystallization or grain rotation) acting locally (due to 107 enhanced stress near obstacles to flow, transitions in the basal thermal state, or fluidity contrasts in the ice) introduce lateral 108 heterogeneity in physical properties that produce incoherent scattering. Here, we compile radar data from a variety of 109 geophysical campaigns, including ground-based and airborne surveys conducted by the Center for Remote Sensing and 110 Integrated Systems (CReSIS), the British Antarctic Survey (BAS), the University of Texas (UT), the University of Washington 111 (UW), and the Alfred Wegener Institute (AWI) – (see Supplementary Table 1 for full system characteristics). From those data, 112 we analyze representative, ice core adjacent radar images, and compare them to measurements of crystal orientation fabric and

- 113 micro- and macro-scale structures, to test two hypotheses:
- 114 1. That transitions in ice COF are collocated with (and likely induce) incoherent scattering.
- 115 2. That small scale deformation of chemically distinct layering can induce incoherent scattering.

116 **3.** Data and Methods: Measurements Capturing the Fine- and Large-Scale Electrical Structure of Ice Cores

Folds and layer disturbances at all scales have been observed or inferred from ice core records in both Antarctica and Greenland. Some scales of folding are more easily detected – millimeter and centimeter scale folds can be measured directly within the ~8 cm diameter ice cores. Folding at the 100s of meters scale is resolvable by radar. But all scales in between must be inferred, using anomalous patterns of electrical conductivity, stable isotope or impurity concentrations, or physical and optical properties. We summarize the measurements that we use to identify deformation in deep ice below, and aim to relate radio-wave scattering phenomena to these observations.

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124 Physical analysis of ice cores, including macro-scale visual observations and optical imaging (i.e. line scanners) (Faria et al., 125 2018; Jansen et al., 2015; Svensson, 2005), and alternating current and direct current electrical conductivity measurements 126 (Fudge et al., 2016; Wolff, 2000) provide the best direct measurement of small scale features deep in the ice column. The 127 resolution of typical line scan images is around 0.1 mm/pixel, allowing for observations of layers and their structure ranging 128 from millimeter-scale undulations up to folds at scale of the typical ~8 cm diameter of deep ice cores. Data from ice core line-129 scanning have shown wavy strata (e.g. WAIS (West Antarctic Ice Sheet) Divide, (Fitzpatrick et al., 2014)), highly inclined 130 strata (e.g. EDML (EPICA Dronning Maud Land), (Faria et al., 2018)) and large z-folds (e.g. NEEM (North Greenland Eemian 131 Ice Drilling), (Jansen et al., 2015)), capturing unique forms of stratigraphic disturbance and in some cases, informing the depth 132 associated with discontinuities in the climate record.

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In addition to imaging methods that capture small scale deformation, a range of chemical methods have been employed across deep ice core sites to identify major breaks in stratigraphic continuity and large-scale folding. Some breaks in continuity have





136 been identified using chemical disagreement between ice cores. For cores in the same geographic region (e.g. GISP2 137 (Greenland Ice Sheet Project Two), GRIP (Greenland Ice Core Project), and NorthGRIP (North Greenland Ice Core Project)), 138 divergence in electrical conductivity, delta δ^{18} O of ice (δ^{18} O_{ice}), and impurity concentrations can be used to identify the onset 139 of a discontinuous record (Johnsen et al., 2001). When looking across hemispheres, divergence in the profiles of globally well-140 mixed delta δ^{18} O of atmospheric O₂ (δ^{18} O_{atm}) and CH₄ have been used to identify climate record discontinuity (Chappellaz et 141 al., 1997; Landais et al., 2003). In cases where there are no cores that provide high resolution comparison, sudden shifts in the 142 nature of the chemical signal (e.g. changes in chemical variability or abrupt changes in the gas-age ice-age difference, described 143 as either the Δ age between the ice and gas or the depth-shift separating gas and ice of a constant age) have been used to infer 144 climate record discontinuities (Crotti et al., 2021; Dansgaard, 1982; Jouzel et al., 2007; Petit et al., 1999; Ruth et al., 2007). 145 Chemical methods have also been used to reconstruct chronologies in heavily disturbed stratigraphy (Landais et al., 2003; 146 NEEM Community Members, 2013; Raynaud et al., 2005; Souchez et al., 2002; Verbeke et al., 2002), and from those 147 chronologies, identify overturned folding. These methods have in some places tentatively inferred (e.g., at Vostok and GRIP) 148 and in other places clearly identified (at NEEM) folding on scales of 10-100 m.

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150 In addition to measurements capturing macro-scale deformation, crystallographic analysis of glacial ice, typically performed 151 using vertical and/or horizontal thin sections of ice cores, provides two useful pieces of information for our work. 152 Measurements of the bulk c-axis orientation of glacial ice gives us direct constraint on how the polarization-dependent 153 permittivity of ice might vary with depth, and therefore be a source of scattering. It also provides information about the strain 154 history of ice, with implications for larger-scale deformation in the ice column. C-axis orientation can be measured with a 155 range of techniques, including polarized light microscopy (Azuma et al., 1999; Weikusat et al., 2017; Wilson et al., 2003), x-156 ray diffraction and tomography (Miyamoto et al., 2011), sonic wave methods (Kluskiewicz et al., 2017), electron backscatter 157 diffraction microscopy (Obbard and Baker, 2007), and open resonator methods (Saruya et al., 2024).

158

159 Historically, data from thin sections have provided the most robust evidence of differential strain at small scales, capturing 160 fabric changes within a single 10 cm vertical thin section (e.g. NEEM, (Montagnat et al., 2014)). But the logistics of thin 161 section sampling limits their ability to capture some scales of vertical and horizontal variability in fabric. The distance between 162 adjacent, discrete thin-section samples can be anywhere from 20 to 100+ m (e.g. EDML (Weikusat et al., 2013), Siple Dome 163 (Gow and Meese, 2007), NorthGRIP (Wang et al., 2002), GRIP (Thorsteinsson et al., 1997)). New approaches to c-axis 164 characterization may change what is possible in future studies of fabric derived scattering, as open resonator methods using 165 0.5 m thick sections of ice have been used to measure the clustering of crystal c-axes every 20 mm along the Dome Fuji core 166 (Saruya et al., 2022, 2024). But for most available data, we are limited in our ability to quantitatively predict scattering from 167 existing fabric measurements, as the magnitude of backscatter depends on the depth-rate-of-change of fabric. Instead, we focus

168 primarily on qualitative comparison of fabric changes with radar images.





169 **3. Results: Investigating the Sources of Incoherent Scattering**

We present measured fabric and structural data together with radar imagery across 9 well sampled cores in Figure 3, and evaluate the depth-agreement of scattering, known fabric transitions, and small- and large-scale deformation below. A full description of the ice core data used can be found in Supplementary Table 2.

173 **3.1** Crystal fabric transitions as a source of incoherent scattering

174 Given the enhanced stresses and therefore higher strain-rates near the base of ice sheets, one might expect monotonic but 175 intensifying fabric development with depth. And at the majority of ice core drill sites, c-axis fabrics transition from a quasi-176 isotropic c-axis distribution at the top of the ice column to strong single maxima lower in the column (e.g. Camp Century, 177 Dye-3, GISP2, NEEM, Dome C, Talos Dome, GRIP), a product of the typical simple shear near the base of a glacier. Ice cores 178 drilled at flank sites or otherwise away from ice divides can also experience uniaxial horizontal extension, and thus c-axis 179 fabrics transition from quasi-isotropic to girdle-type fabric and then to a single maximum (e.g. NorthGRIP, Vostok, EDML). 180 But variability in the impurity content (which can vary with changes in climate) can intensify fabric development and localize 181 fabric transitions, with fabric strengthening typically coincident with higher impurity content (seen at Byrd (Faria et al., 2014), 182 Camp Century (Faria et al., 2014), Talos Dome (Montagnat et al., 2012), Dome C (Durand et al., 2009), NEEM (Montagnat 183 et al., 2014), GISP2 (Gow et al., 1997), and Dye-3 (Langway et al., 1988)).

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185 Abrupt fabric transitions occurs at most ice core sites in Greenland (e.g. Camp Century, Dye-3, GISP2, and NEEM), where a 186 significant change in impurity deposition at the Holocene-Wisconsin climate transition is co-located with an abrupt 187 strengthening or transition to a vertical-maximum fabric (Faria et al., 2014). In some places, we see an isolated but abrupt 188 transition in fabric that has co-located scattering in the radar image. At NEEM, a transition from a weak vertical girdle to 189 strong single maximum fabric occurs at 1419 m and is coincident with a diffuse reflector in the radargram (fig. S1). Similar, 190 reflectors appear at isolated fabric transitions in Antarctica as well. At Siple Dome, the c-axis fabric transitions from a vertical 191 girdle to a single maxima at 700 m, with a corresponding diffuse reflector in the radar data. At EDML, the c-axis fabric 192 transition from a vertical girdle to a strong single maxima between 2025 and 2045 m has been identified as the origin of the 193 reflector at 2035 m (Eisen et al., 2007). These reflectors appear less specular (with trailing energy after the initial arrival) than 194 other isochronous layering within the radar imagery.

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Where we see well sampled gradual transitions in fabric (spanning 50-100 m of the ice column) we observe a diffuse band of incoherent scattering. At Dome C, the strong single maximum fabric at 2800 m gradually transitions to a broad single maximum fabric at 2857 m and returns to a strong single maxima fabric at 2900 m (Durand et al., 2009). This fabric transition is roughly coincident with the transition from thin coherent isochronal strata to a single diffuse incoherent scattering layer observed around 2825 m. Similarly, at Dome Fuji, the strong single maxima fabric at 2660 m gradually weakens before returning to a





strong single maxima fabric again at 2760 m (Saruya et al., 2024). Again, this appears roughly coincident with a diffuse
 incoherent scattering layer observed at ~2700 m in the radargram (fig. S2).

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204 In many places, especially where time is compressed significantly at the base of the ice column, alternating fabrics have been 205 observed. At Vostok, from 2700 to 3315 m depth, the core alternates between coarse grained ice with girdle-type fabric and 206 fine grained ice with single-maximum fabric every ~100 m (Obbard and Baker, 2007). Within the girdle-type fabric zone 207 between ~3220 and 3315 m, we see weakly banded incoherent scattering (3220 - 3290 m). Between ~3315 and 3450 m 208 alternations between girdle-type and single-maximum fabric occur approximately every ~20 m (Lipenkov and Raynaud, 2015). 209 This zone of increased fabric alternation overlaps with both the no echo zone between ~ 3290 and 3360 m and the upper depths 210 of a weakly banded incoherent scattering unit (~3360 – 3490 m) in the radargram. While interpretation of the GISP2 and GRIP 211 radargrams has significant ambiguity below 2800 m, 40 km length radar transects show laterally heterogeneous incoherent 212 scattering in that depth range (fig. S3). At GRIP, each of the five thin sections sampled between 2800 and 2950 m depth show 213 alternating fabrics. At GISP2, coarse-grained layers with fabrics that deviate from the strong single maximum are observed at 214 increasing frequencies below 2800 m (Gow et al., 1997).

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216 While challenging to constrain due to thin-section sampling frequency, the smallest scale of fabric variability has been 217 observed or inferred at centimeter-scales, including at Vostok, Dome C, and Dome Fuji. At Vostok, fabric alternations occur 218 at cm-scale frequency from 3450 m until the transition from meteoric to accreted ice at 3538 m (Lipenkov and Raynaud, 2015). 219 This overlaps with an echo-free zone in the radargram. At Dome C, ice below 2800 m consists of alternating layers with high 220 impurity content (consistently presenting strong single maxima fabric) and layers with low impurity content (with an associated 221 broad single maximum fabric). After the gradual transition into and out of a broad single maximum fabric at 2850 m, fabric 222 transitions below 2920 m become more local. High frequency sampling of fabric (every 0.5 m) between 2933 and 2955 m 223 revealed fabric alternations between each sample (Durand et al., 2009). Unlike at Vostok where the onset of rapid fabric 224 transitions coincides with the start of the echo free zone, the onset of rapid fabric transitions at Dome C is associated with thick 225 and sometimes discontinuous bands of incoherent scattering (2900 - 3050 m) in the radargram. At Dome Fuji, cm-scale 226 fluctuations from the single maximum fabric, observed by increases in the standard deviation of $\Delta \varepsilon$ (the difference in the 227 relative permittivity, ε , between vertical and horizontal planes), begin around 2400 m and intensify through the base of the ice 228 column (Saruya et al., 2024). The increase in fabric fluctuations between 2400 and 2650 m has no obvious effect on the 229 coherent continuous layering observed in the radargram. However, the Dome Fuji radargram transitions to a zone of laterally 230 homogenous incoherent scattering at 2900 m. Notably, the precise depth of that transition is difficult to constrain in the radar 231 image, due to the combination of increasing layer inclinations (Dome Fuji Ice Core Project Members, 2017) and strong 232 scattering from borehole fluid in the ice core cavity (fig. S2).

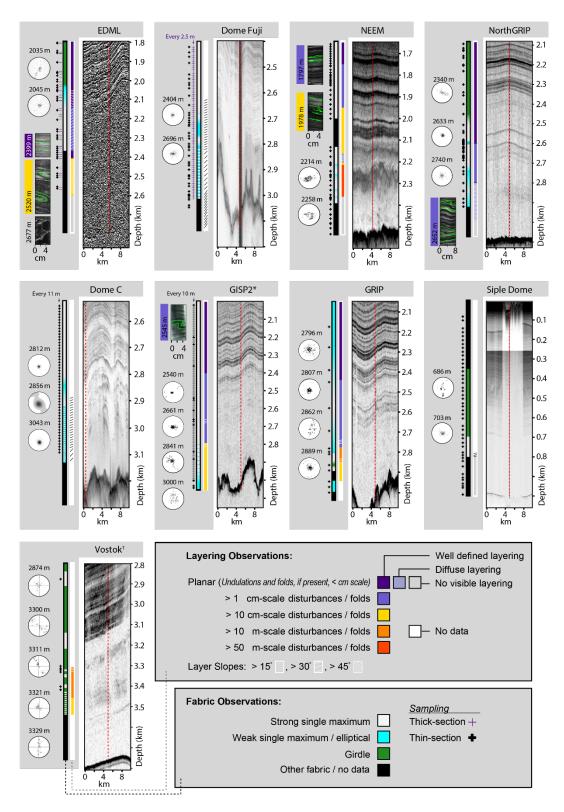




The most complex scattering pattern we see is at NEEM, starting at ~2200 m. We see strong incoherent scattering that is highly laterally variable. In this section of the ice core, the same oxygen isotope sequence (and its associated fabric gradient, from multi-clustered-fabric to single maxima fabric) is repeated, with abrupt fabric transitions at the boundaries between sequences. This is attributed to overturned folds at the base of the ice column, in part, facilitated by rheologic differences in the ice that also produce the abrupt fabric transitions. Here, both fabric and larger-scale deformation likely play a significant role in the nature of the scattering, introducing lateral heterogeneity in material properties (for ice at a given depth) that doesn't exist at other ice core sites.











242 Figure 3: Radargrams capturing deep ice at ice core drill sites where we have comprehensive fabric and deformation data. This 243 includes line-scan data: EDML (Faria et al., 2018), Dome Fuji (Takata et al., 2004), NEEM (Kipfstuhl, 2009), NorthGRIP (Svensson, 244 2005), GRIP (Alley et al., 1997); fabric observations: EDML (Eisen et al., 2007; Faria et al., 2018; Weikusat et al., 2013), Dome Fuji 245 (Saruya et al., 2022, 2024), NEEM (Eichler, 2013; Montagnat et al., 2014), NorthGRIP (Wang et al., 2002), Dome C (Durand et al., 246 2009), GISP2 (Gow et al., 1997), GRIP (Thorsteinsson et al., 1997), Siple Dome (Gow and Meese, 2007), Vostok (Obbard and Baker, 247 2007); and layering observations: EDML (Faria et al., 2010, 2018), Dome Fuji (Dome Fuji Ice Core Project Members, 2017), NEEM 248 (Jansen et al., 2015), NorthGRIP (Svensson, 2005), Dome C (Durand et al., 2009), GISP2 (Alley et al., 1995, 1997; Faria et al., 2014; 249 Gow et al., 1997), GRIP (Alley et al., 1995; Dahl-Jensen et al., 1997; Johnsen et al., 1995; Landais et al., 2003), Siple Dome (Gow and 250 Meese, 2007), Vostok (Lipenkov and Ravnaud, 2015; Ravnaud et al., 2005; Souchez et al., 2002). *At GISP2, only some of the sampled 251 thin sections have published data (indicated by the black + symbols), and [†]at Vostok, the original sampling rate is unpublished, with only a few thin sections and general observations available in the literature. 252

3.2 Folding as a source of incoherent scattering

254 Millimeter-scale disturbances are likely present in most deep glacial ice, given their ubiquity in physical observations of ice 255 cores. But there is little evidence that deformation at that scale impacts the radiostratigraphy directly. In the South Pole Ice 256 Core (SPICEcore), inclined and pinched cloudy bands are observed starting at 1000 m and continue at intervals through the 257 end of the core (Fegyveresi and Alley, 2018) without any noticeable impact on radar scattering. Crystal striping at GISP2 is 258 observed starting at 2200 m, coincident with the onset of small scale undulations in linescan images (Alley et al., 1997) but 259 similar to SPICEcore, there is no associated change in the nature of radar lavering. Millimeter-scale z-folds at GRIP first appear 260 at 2438 m and at 2437 m at GISP2 (Alley et al., 1997), which does coincide with a drop in power of coherent scattering layers. But there is a commensurate drop in the ice conductivity variability associated with changes in dust deposition, which better 261 262 explains that change. Thus, we rule out millimeter-scale folding as a significant contributor to the radar signal observed at 263 these locations.

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265 Stratigraphic disturbances at the centimeter-scale are apparent in all cores with available data. In previous work, this scale of 266 deformation has been invoked as a mechanism for the "echo free zone", with the idea that folding effectively homogenizes 267 dielectric contrasts at the scale of the resolution of the radar (Winter et al., 2017). At EDML and WAIS Divide, the onset of 268 cm-scale disturbance does appear to be collocated with the apparent echo free zone. In both radar images, however, there is a 269 gradual diminution of returned power with depth. It is possible that measured disturbances do reduce the intensity of back-270 scatter without eliminating it entirely. But there is laterally-continuous layering (with strong back-scatter intensities) in regions 271 of cm-scale disturbances at NorthGRIP, NEEM, and GRIP, and in regions with disturbances at the scale of 10 cm at NEEM. 272 Radar data at NEEM show no change in scattering behavior associated with deformation at this scale. This seems to imply that 273 these radar systems (with range-resolutions of 2.8 m to 5 m (Table S1)) are insensitive to deformation at this scale.





275 But deeper in the NEEM core, where chemical analyses reveal six zones of disturbed ice including two large 50 and 100 m 276 thick folded layers of inverted early glacial ice (NEEM Community Members, 2013), high amplitude but laterally variable 277 incoherent scattering can be seen in the radar imagery. Deformation at this scale, thought to be in part due to rheological 278 differences between the glacial and interglacial ice (NEEM Community Members, 2013), is coincident with a loss of coherent 279 banding in the line scan imagery and an increase in the lateral heterogeneity of intensity in incoherent backscatter. Above 3460 280 m depth at Vostok, folding is also inferred at the meter scale and larger (Lipenkov and Raynaud, 2015). Similarly, there is 281 incoherent scattering in the image at these depths, although the amplitude of the backscatter is weaker, and lateral heterogeneity 282 less pronounced. Finally, at GRIP, tentative chronological reconstructions of disturbed ice below 2750 m show significant 283 disruption and folding on the scale of 10s of meters between 2780 to 2850 m. And while near the ice core, this depth-range 284 corresponds with a unit of weak incoherent scattering, at the 10s of kilometers scale, there is significant variability in the 285 amplitude (fig. S3).

4. Discussion: Using Incoherent Scattering in Ice Core Site Selection

There is compelling evidence that incoherent scattering can arise from fabric transitions in the deep ice, and the quality of that scattering could be diagnostic of large-scale deformation that is co-located with the smaller-scale fabric development. If true, then incoherent scattering might be used to improve ice core site selection. We test that theory at 16 ice core sites, by first subdividing core-adjacent radar imagery into five types of signal (fig. 4):

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1. Laterally continuous coherent scattering (that is, clear isochronal layering)

- 293 2. Diffuse but banded scattering
- 294 3. Laterally homogenous incoherent scattering
- 295 4. Laterally heterogeneous incoherent scattering
- 5. No signal (or rather, signal levels at or below the noise floor of the system).
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We then compare these scattering types to known breaks in the continuity of the associated ice cores (see Appendix A for the observational basis for each labelled break).

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Across these core sites, continuous coherent scattering is almost exclusively found above known breaks in the climate record. This type of scattering appears below the break in a climate record in only one ice core, and that is Vostok, where the interface between accreted and meteoric ice and a layer of mineral inclusions from the lake bed (Turkeev et al., 2021), define two clear reflection horizons. As a result, in typical glaciological environments, continuous coherent scattering is a robust indicator of ice core continuity. At the studied core sites, where diffuse but banded scattering sits immediately below laterally continuous





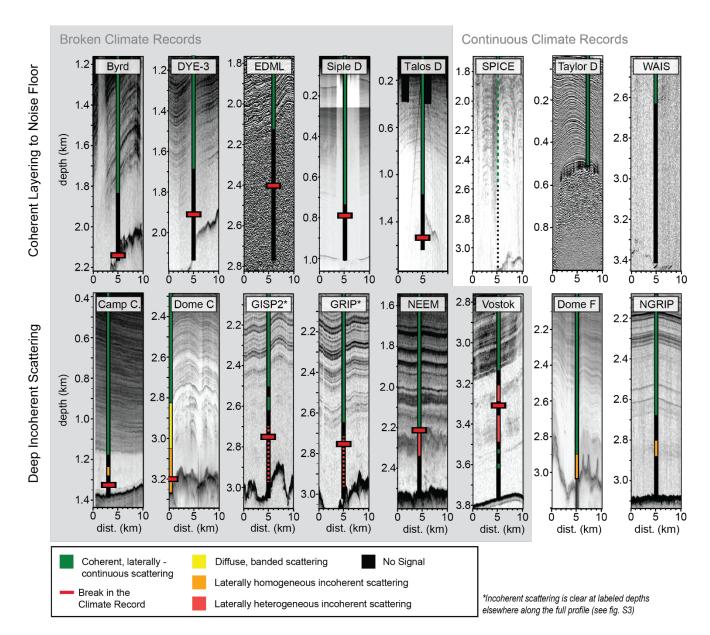
306 layering, there are no associated breaks in measured climate records. This supports the idea that banded but incoherent 307 scattering is not an indication of disturbed basal ice.

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309 Where we see laterally homogenous incoherent scattering, as in Camp Century, Dome C, Dome Fuji, and NorthGRIP, it occurs 310 within sections of ice with a continuous climate record. This likely indicates fabric transitions that are themselves defined 311 weakly by depositional impurities, and thus, the shape of the scattering band is roughly parallel to the isochronous layering. 312 At Vostok, we see laterally incoherent scattering that is heterogeneous in its intensity but is otherwise layering conformal, 313 directly above and ~ 100 - 200 m below the broken climate record. These two bands of incoherent scattering are qualitatively 314 indistinguishable, and demonstrate the challenge of interpreting the quality of the climate record within bed conformal laterally 315 heterogeneous incoherent scattering regions. But where we see laterally heterogeneous incoherent scattering that is layering 316 non-conformal (as in GISP2, GRIP, and NEEM) it occurs below breaks in the continuity of the observed climate record. In 317 those places, it is possible that the same ice rheology contrast that facilitated a fabric transition together with a complex basal stress regime enabled multi-meter scale deformation, inducing lateral variability in the backscatter intensity, and indicating 318 319 significant risk of a disturbed climate record.







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322 Figure 4: Radar images collected proximal to ice core sites. For each core site, the core location, depth, and breaks in the climate

record are labeled. In addition, we categorize the scattering observed in the radar imagery as a function of depth at each core site.
 Metadata for the radargrams is presented in Supplementary Table 1.





325 5. Conclusions

Based on comparison between ice core data and radar imagery at ice core sites, we show that diffuse and incoherent scattering is often colocated with transitions in the crystal orientation fabric of the ice. Transitions in fabric are a product of the local stress regime, but they are localized by differences in grain size. High concentrations of impurities tend to reduce local grainsize and enhance deformation rates, so where climatically driven variations in impurities change the strength of the ice, one might also expect more abrupt contrasts in fabric that back-scatter radio waves. In this way, fabric controlled scattering may be roughly isochronous, although we show that fabric interfaces to not manifest as abrupt, specular reflectors the way chemically induced layering does in radar imagery.

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334 In the deep ice, where stresses are high, time is compressed, and global changes in impurity deposition are expressed over 335 narrower depth ranges, we might expect fabric induced scattering is common. The nature of the fabric transition, and the spatial 336 heterogeneity in the transition, define whether or not the scattering will appear as coherent layering, a diffuse scattering 337 horizon, laterally homogenous incoherent scattering, or laterally heterogeneous incoherent scattering. In addition, ice fluidity 338 contrasts at fabric boundaries facilitate small- and large-scale folding. At small scales (below ~ 1 m), folding seems to have 339 little impact on existing radar data. But large-scale folding, where present, results in complex scattering targets in the 340 subsurface, and induces significant lateral heterogeneity in the incoherent scattering intensity and complex scattering horizons. 341 Where this is observed at existing ice core sites, it seems indicative of discontinuities in the ice core climate record.

342

A final consideration when thinking about fabric induced incoherent scattering is the relationship between permittivity contrasts (as experienced by the propagating radio-wave) and radio-wave polarization. For fabric intensification (for example, a weak single maximum to a strong single maximum fabric) there will be a change in permittivity for all radar polarizations, and scattering will likely appear isotropic. For fabric transitions (for example, from a girdle to a single maximum fabric) it is possible for some polarizations to exhibit scattering and others to have low backscatter or apparent echo free zones. This anisotropic character merits further study at places like Siple Dome, EDML, and Vostok, where girdles are seen in the deep ice.

350

As is true for discussions of the "echo free zone", we show that conversations about the "basal layer" observed in Greenland and Antarctica must start from the understanding that deep scattering (or its absence) depends on system characteristics and physical properties of the ice. Both echo free zones and deep incoherent scattering could arise from multiple mechanisms. But, as with coherent layering, incoherent scattering is signal, not noise, and more work should be done to better interpret this often overlooked component of radar imagery.





356 6. Author Contribution

- 357 EM synthesized data from the literature on physical and chemical properties of ice cores and identified radar data from CReSIS,
- BAS, UT, UW, and AWI. All authors contributed to study design, radargram interpretation, figure creation, and writing of manuscript.

360 7. Competing Interests

361 The authors declare that they have no conflict of interest.

362 8. Acknowledgements

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encourage suggestions from those scholars for new ways to connect ice penetrating radar to measurable ice core quantities.





367 Appendix A. Known Layer Disturbances and Ice Core Continuity Problems

Of the cores studied, 5 show only minor signs of layer disturbances, and contain a continuous climate record through the full depth range of the ice core. Those are Dome Fuji, NorthGRIP, SPICEcore, Taylor Dome, and WAIS Divide. Of the other 11 cores, 6 have well identified breaks in their climate record, and 4 are likely discontinuous (although the exact stratigraphic break is not well identified), and 1 has conflicting observations of discontinuity. Here, we describe the observational basis for claims of a broken climate record.

373 A.1 Cores with Clear Evidence of Stratigraphic Discontinuities

374 (Alphabetically: EDML, GRIP, GISP2, NEEM, Talos Dome, Vostok)

375

EPICA (European Project for Ice Coring in Antarctica) Dronning Maud Land, EDML (Length: 2774 m | Break: 2417 m | **Percentage Disturbed: 12.9%):** The chronology called EDML1 has been established for the top 2417 m of the EDML ice core. The top 2366 m of the core is matched to the EDC3 chronology using volcanic signatures (dielectric profiling, SO4 concentrations, and electrolyte conductivity measurements) (Ruth et al., 2007). Three tie points between the EDC3 chronology and EDML core are matched between 2366 and 2415 m using insoluble dust concentrations, δ^{18} O, and δ D, however these matches are considered uncertain with estimated errors up to several thousand years (Ruth et al., 2007). Macrostructure analysis of linescan images between 2400 and 2500 m shows evidence of large-scale folding (Faria et al., 2010).

383

Greenland Ice Core Project, GRIP (Length: 3029 m | Break: ~2750 m | Percentage Disturbed: 9.2%) and Greenland Ice Sheet Project Two, GISP2D (Length: 3053.4 m | Break: ~2750 m | Percentage Disturbed: 9.9%): CH₄ and $\delta^{18}O_{atm}$ data from both GRIP and GISP2 show evidence of stratigraphic disturbance in the bottom 10% the ice cores. Above 2750 m CH₄ and $\delta^{18}O_{atm}$ values vary synchronously between GRIP and GISP2, but below 2750 m, the chemical profiles diverge, showing large and significant fluctuations which are not present in the undisturbed ice from the Vostok 3G core (Chappellaz et al., 1997).

390

391 North Greenland Eemian Ice Drilling, NEEM (Length: 2540 m | Break: 2209.6 m | Percentage Disturbed: 13%): At 392 NEEM, an abrupt discontinuity in the $\delta^{18}O_{icc}$ at 2209.6 m marks the end of synchronization with the NorthGRIP GICC05 393 extended timescale. Additional discontinuities in the δ^{18} O_{ice} subdivide the bottom 13% of the core into six zones of disturbed stratigraphy. These correspond with similar shifts in other atmospheric gas measurements (CH₄, $\delta^{18}O_{atm}$, N₂O, $\delta^{15}N$ of N₂). 394 395 Within the upper five zones, the layering is thought to be unbroken (based on continuous records of N₂O, δ^{15} N of N₂, dust, or 396 electrical properties), with timescales for each of the upper five zones reconstructed by synchronizing NEEM $\delta^{18}O_{atm}$ and CH₄ 397 profiles with NorthGRIP and EDML records. The timescales for these zones include inverted, mirrored, and folded ice up to 398 100 m thick (NEEM Community Members, 2013).





399

TALos Dome Ice CorE, TALDICE (Length: 1620 m | Break: 1548 m | Percentage Disturbed: 4.4%): At Talos Dome, Crotti et al. identify a break in stratigraphic continuity at 1548 m using analysis of $\delta^{18}O_{atm}$, δD , and 81Kr dating, described below (Crotti et al., 2021). TALDICE $\delta^{18}O_{atm}$ and δD measurements were matched to the EDC $\delta^{18}O_{atm}$ and δD record through visual synchronization through 1548 m depth. Below 1548 m, the amplitude of $\delta^{18}O_{atm}$ fluctuations is damped, making synchronization with the EDC record uncertain. Similarly, below 1548 m, the TALDICE δD signal becomes asynchronous with the EDC record. ⁸¹Kr dating of three samples below 1548 m depth revealed that ice from 1613 - 1618 m had comparable age to samples from 1559 - 1563 m and 1573 - 1578 m depth, indicating a disturbed age-depth relationship.

407

408 Vostok 5G-5 (Length: 3658 m | Break: 3311 m | Percentage Disturbed: 9.5%): The stratigraphy in the bottom 9% of the 409 Vostok 5G core is divided between 228 m of disturbed meteoric ice, and 119 m of accreted lake ice. In the upper part of the 410 disturbed meteoric ice, the lack of depth-shift between δD_{ice} and gas measurements (CO₂ and CH₄) is interpreted by Souchez 411 et al. as evidence of folding and intermixing (Souchez et al., 2002). Observations of ash layers with depth-varying inclinations 412 supports interpretation of large-scale folding. In the lower part of the disturbed meteoric ice, damped variation of δD_{ice} and 413 trace impurity distributions (Na+, Cl-, non-sea salt Mg++ and Ca++), physical observations of interbedded fine-grained 414 (presumably glacial) and coarse-grained (presumably interglacial) ice, and the presence of bed material in the bottom 100 m 415 of the disturbed meteoric ice, is interpreted as further evidence for stratigraphic deformation (Lipenkov and Raynaud, 2015; 416 Souchez et al., 2002). At 3538 m depth, the transition between meteoric and accreted ice is apparent from the $\delta D_{ice}/\delta^{18}O$ 417 fingerprint of freezing processes (Jouzel et al., 1999). At this depth, sudden transitions to lower total gas content, increased 418 crystal size, low ECM values, increased δD_{ice} , and decreased deuterium excess, provide further evidence for the 419 meteoric/accreted ice transition (Jouzel et al., 1999).

420 A.2 Cores that Likely Contain Stratigraphic Discontinuities or Conflicting Observations of Discontinuity

421 (Alphabetically: Byrd, Camp Century, Dome C, Dye-3, Siple Dome)

422

Byrd Station '68, BYRD 68 (Length: 2164 m | Break: 2135-2144 m | Percentage Disturbed: ~1%): A chronology for the upper ~99% (2144 m) of the Byrd core has been established by synchronizing Byrd, GRIP, and GISP2 CH₄ profiles (Blunier and Brook, 2001). Gas volume measurements from the bottom 10 m of the core (2154 - 2164 m) suddenly approach zero at 4.83 m above the bed, revealing the transition between meteoric ice and accreted subglacial meltwater (Gow et al., 1979). The bottom 4.83 m of non-meteoric ice contain horizontal bands of basal debris including sand, clay, and pebbles as large as 8 cm in diameter (Gow et al., 1979). Grootes et al. 2001 observe that the Byrd δ^{18} O record becomes asynchronous with Taylor Dome and Vostok record around 2135 m.





431 Camp Century, CC 63-66 (Length: 1387.4 m | Break: ~1350 m | Percentage Disturbed: 2.7%): The integrity of the Camp Century climate record is uncertain below 1310 m depth where δ^{18} O profiles of Camp Century, GRIP, and GISP2 become 432 433 asynchronous (Johnsen et al., 2001). Correlation of a smoothed Camp Century δ^{18} O profile with benthic foraminifera record 434 from deep sea core RC11-120 provides a tentative extension of the chronology through about 1330 m, the depth of the 435 inflection point associated with Marine Isotope Stage (MIS) 5d (Dansgaard et al., 1985). A dramatic cold event at 1340 m is 436 associated with a similar δ^{18} O fluctuation in the disturbed section of the GRIP core at 2800 m (Johnsen et al., 2001). Johnsen et al. describe dramatic fluctuations in δ^{18} O below Greenland Interstadial (GI) 23 in the GRIP, GISP2, and Camp Century 437 cores which are not represented in the continuous δ^{18} O signal from Vostok (Chappellaz et al., 1997). 438

439

440 EPICA Dome C, EDC99 (Length: 3260 m | Potential Break: ~3200 m | Percentage Disturbed: ~1.8%): The continuity of 441 the upper 98% (3200 m) of the Dome C core is evidenced primarily through matching of δD_{ice} to the deep-sea benthic $\delta^{18}O$ 442 record (Jouzel et al., 2007). Additional matching of enhanced ¹⁰Be deposition to Matuyama-Brunhes geomagnetic reversal 443 between 3160 and 3170 m (Jouzel et al., 2007) and matching of CO₂ and CH₄ profiles to MIS18 and 19 between 3160 and 444 3185 m further support the continuity of the upper 98% of the core. Below 3200 m, there is contradictory evidence about the continuity of the climate record. Measurements of δD , total air content, gas composition, and dust content suggest continuity 445 to bedrock, while $\delta^{18}O_{atm}$, visible inclusions, length of the glacial period, and variability of chemical species distribution 446 suggest altered stratigraphy (Tison et al., 2015). 447

448

449 DYE-3, DYE-3, DYE-3, Tore and a preak: 1940 m | Percentage Disturbed: 4.8%): At DYE-3, the continuity of the 450 climate signal is lost between 1900 and 1987 m. Initially, Dansgaard et al. 1982 correlated fluctuations between the δ^{18} O 451 measurements at DYE-3 and Camp Century through 1987 m depth. Between 1987 and 2010 m. DYE-3 δ^{18} O values are guasiconstant, and interpreted as evidence of folded layers. Later, comparison of the δ^{18} O values between DYE-3 and GRIP led 452 453 Johnsen et al., 2001 to identify GI 8 at 1900 m as the last undisturbed match point between the two records. However Johnsen 454 et al. would still identify two match points in the deeper ice: GI 12 (~1925 m) and GI 14 (~1940 m). As such, in our analysis 455 we have used 1940 m as the depth of the broken climate record. CO₂ and CH₄ measurements of the bottom 27 m of silty ice 456 have been used to identify 4 distinct zones of highly deformed basal ice (Verbeke et al., 2002).

457

Siple Dome A, SDMA (Length: 1004 m | Break: ~800 m | Percentage Disturbed: ~20%): The integrity of the Siple Dome climate record is uncertain in the bottom 200 m of the core, however a precise onset depth for the disturbed ice is poorly constrained. A chronology for the 514 - 854 m section of the core was established by synchronizing Siple Dome, GISP2, and GRIP CH₄ profiles (Brook et al., 2005). Below 854 m, the methane data becomes sparse however a possible chronology has been proposed between 854 and 920 m based on the matching of a single inflection point in the $\delta^{18}O_{atm}$ profile of Siple Dome core at 920 m with a corresponding GISP2 $\delta^{18}O_{atm}$ inflection point (Brook et al., 2005). Macro and micro-scale physical observations by Gow and Meese suggest an interrupted climate record by 800 m depth, summarized here (Gow and Meese,





465 2007). Between 560 and 800 m, sequences of inclined layering occasionally surpassing 10 degrees as well as reversed dips are 466 observed. Below 800 m the core is highly fractured, limiting any further observations of layer structure. Around 700 m, the c-467 axis fabric shifts suddenly to a single maximum corresponding to a stress regime dominated by strong horizontal shear. Around 468 800 m, the c-axis fabric shifts back to a multi-maxima fabric.

469 A.3 Cores with No Significant Break in Continuity

- 470 (Alphabetically: Dome Fuji, NorthGRIP, SPICEcore, Taylor Dome, WAIS Divide)
- 471

Dome Fuji, DF2 (Length: 3035.22 m): The integrity of the Dome Fuji ice core climate record is discussed by the (Dome Fuji Ice Core Project Members, 2017) and summarized here. A chronology for the upper 3028 m of the 3035 m Dome Fuji core was established through the synchronization of δ^{18} O records to the Dome C δ D profile. Physical observations of inclined layers begin at 2400 m and show distinct stepwise increases in inclination: ~8° between 2450 - 2600, ~20° between 2600 - 2800, ~40° between 2800 - 2900, ~45° at 2950 m, and ~50° at bedrock. Despite the observations of inclined layers, which are attributed to spatially variable basal melt conditions, explicit observations of folded layers were not noted and the synchroneity of the δ^{18} O and Dome C δ D profiles are considered evidence of an intact climate record within the depths of inclined layers.

479

North Greenland Ice Core Project, NorthGRIP (Length: 3090 m): At NorthGRIP, the continuity of the 2544 – 3073 m zone of the 3090 m length core was confirmed by matching NorthGRIP $\delta^{18}O_{atm}$ and CH₄ records to EDML and EDML1 chronologies (Capron et al., 2010). Depth shift analysis at 2940 m showed the expected shift between $\delta^{15}N$ and CH₄ vs $\delta^{18}O$ during Dansgaard-Oescher (DO) 24, and was used to confirm the continuity of the deepest layers (North Greenland Ice Core Project Members, 2004). Like at WAIS Divide, small scale stratigraphic disturbances are observed a few hundred meters above bedrock (Svensson, 2005), but are not considered large enough to impact the continuity of the climate record.

486

South Pole Ice Core, SPICEcore, SPC14 (Length: 1500 m | Ice thickness: 2700 m): Continuity through the end of the core
is established through synchronization of CH4 fluctuations to WAIS Divide ice core (Epifanio et al., 2020). Notably,
SPICEcore drilling stopped 1200 m above bedrock.

490

Taylor Dome, M3C1 (Length: 554 m): The continuity of the Taylor Dome core was established through correlation of the CH₄ and δ^{18} Oatm profiles with corresponding GISP2 profiles as well as correlation of the δ^{18} O profile with the Vostok δ D record (Grootes et al., 2001; Steig et al., 1998). The δ^{18} O inflection point associated with MIS 5e (~130kyBP) is identified between 526 and 531 m depth. The identification of correlated inflection points continues confidently through 200 kyBP with a tentative chronology, limited by sample resolution, extending beyond 300 kyrBP (Grootes et al., 2001).





497 WAIS (West Antarctic Ice Sheet) Divide, WDC06A (Length: 3405 m | Ice thickness: 3455 m): The continuity of the 498 WAIS Divide core is confirmed above 2850 m by annual layer counting, and below 2850 m via synchronization of WAIS 499 Divide CH₄ measurements to the NorthGRIP δ^{18} O record and a refined Hulu Cave speleothem δ^{18} O record (Buizert et al., 500 2015). Notably, the 3405 m WAIS Divide core ends 50 m above bedrock, so continuity in the uncored 50 m basal unit is not 501 confirmed. Mm-scale or smaller stratigraphic disturbances are observed at 3150 and 3232 m (Fitzpatrick et al., 2014) but are 502 not considered large enough to impact the continuity of the climate record.





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