

# Study of ~~radiative~~ optical scattering properties and direct radiative effects of high-altitude cirrus clouds in Barcelona, Spain with 4 years of lidar measurements

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**Abstract.** Cloud-radiation interaction still drives large uncertainties in climate models and its estimation is key to make more accurate predictions. In this context, the high-altitude cirrus clouds play a fundamental role, because 1) they have a high occurrence frequency globally and 2) they are the only cloud that can readily cool or warm the atmosphere during daytime, depending on their properties. This study presents a comprehensive analysis of ~~radiative~~ optical scattering properties and direct radiative effects of cirrus clouds based on 4 years of continuous ground-based lidar measurements with the Barcelona (Spain) Micro Pulse Lidar. First, we introduce a novel approach of a self-consistent scattering model for cirrus clouds to determine their ~~radiative~~ optical scattering properties at different wavelengths using only the ~~effective~~ extinction coefficient and ~~cloud~~ cloud temperature. Second, we calculate the ~~radiative~~ direct radiative effects of cirrus clouds with the Discrete Ordinates Method and we validate our results with SolRad-Net pyranometers and ~~CERES~~ NOAA-20 measurements. Third, we present a case study analyzing the direct radiative effect of a cirrus cloud along its back-trajectory using data from the Chemical Lagrangian Model of the Stratosphere with microphysics scheme for Ice clouds formation. The results show that the cirrus clouds with an average ice water content of  $4.97 \pm 5.53 \text{ mg/m}^3$ , at nighttime, ~~warm the atmosphere~~ have a positive direct radiative effect at top-of-the-atmosphere (TOA; ~~+50.1~~  $+40.4 \text{ Wm}^{-2}$ ) almost twice than at bottom-of-the-atmosphere (BOA; ~~+28.0~~  $+22.1 \text{ Wm}^{-2}$ ); at daytime, they have generally ~~cool the~~ a negative direct radiative effect at BOA (~~-8.57~~  $-11.5 \text{ Wm}^{-2}$ , 80 82% of the cases) and always ~~warm the~~ a positive effect at TOA (~~+18.9~~  $+14.18 \text{ Wm}^{-2}$ ). In these simulations, the influence of the lower layer aerosols is negligible in the cirrus direct radiative effects, with a BIAS of ~~-0.7%~~  $-1.2\%$ . For the case study, the net radiative effects produced by the cirrus cloud, going at TOA from 0 to ~~+42~~  $+40 \text{ Wm}^{-2}$  and at BOA from -51 to  $+20 \text{ Wm}^{-2}$ . This study reveals that the complexity of the cirrus cloud direct radiative effect calculation lies in the fact that it is highly sensitive to the cirrus scene properties.

## 20 1 Introduction

Cloud-radiation interaction still drives large uncertainties in weather and climate models (IPCC, 2023). Its estimation is very important in order to understand the main physical processes driving climate change, to predict long-term global warming and to make more accurate weather predictions. (Loeb et al., 2009) estimated globally at top-of-the-atmosphere an annual cloud shortwave radiative effect of approximately  $-50 \text{ Wm}^{-2}$  and longwave effect of approximately  $+30 \text{ Wm}^{-2}$ . The resulting net  
25 global mean cloud radiative effect of approximately  $-20 \text{ Wm}^{-2}$  implies a net cooling effect of clouds on the current climate. Owing to the large magnitudes of the cloud radiative effects, clouds cause a significant climate feedback that depends on cloud properties and their spatial distribution (IPCC, 2023). In this context, the high-altitude cirrus clouds play a fundamental role in the global radiation budget (Liou, 1986; Lolli et al., 2017b), having been designated as poorly understood by (IPCC, 2023) because of a lack of knowledge of their dynamic, microphysical and ~~radiative~~ optical scattering properties. Indeed, cirrus cloud  
30 critical role in the climate comes from the fact that 1) they have a high occurrence frequency globally (Holz et al., 2008) and 2) they are the only cloud that can readily cool or warm the top-of-the-atmosphere and bottom-of-the-atmosphere, during daytime, depending on their properties (Campbell et al., 2016). In fact, (Campbell et al., 2016) demonstrated through a one-year long lidar dataset that positive or negative daytime cirrus cloud forcing could occur depending on the cloud optical depth (COD) and the solar zenith angle (SZA).

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Cirrus clouds are mainly composed of ice crystals and can form through different atmospheric mechanisms that determine their initial properties and further evolution. In European field campaigns it has been observed that during a low or high pressure system, cirrus clouds are typically formed by a slow updraft, while in conjunction with jet streams or gravity waves, cirrus clouds originate as a consequence of a fast updraft. Also, liquid origin cirrus mostly related to warm conveyor belts are found  
40 (Kramer et al., 2016). The most common parameters that are measured in cirrus clouds are temperature, relative humidity (for ice), vertical velocity, ice water content, ice crystal number, ice nucleation particles and ice crystal size distribution. Unfortunately, the measurements of ice crystal number and size as well as relative humidity have faced instrumental problems during last decades (Korolev et al., 2011; Kramer et al., 2016). Moreover, it is a difficult task to draw conclusions about the microphysical processes of cirrus clouds from these observations. Nevertheless, worldwide studies on cloud and aerosol optical and  
45 microphysical properties have increased significantly over the last years through the passive ground-based measurements made e.g. by the European Aerosol Research Lidar NETwork, EARLINET (Pappalardo et al., 2014) now included in the Aerosols, Clouds and Trace gases Research Infrastructure, ACTRIS (Saponaro et al., 2019), Micro Pulse Lidar NETwork, MPLNET (Welton et al., 2001); and satellite measurements e.g. by Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observations, CALIPSO (Winker et al., 2007), AEOLUS (Ingmann and Straume, 2016), MODerate-resolution Imaging System, MODIS  
50 (Levy et al., 2013) and in the future, Earth Cloud, Aerosol and Radiation Explorer, EarthCARE (Eisinger et al., 2017). Additionally, in-situ airborne measurement campaigns have been carried out such as the First ISCCP Project Regional Experiment, FIRE from 1989 to 1995 (Ackerman et al., 1990; Heymsfield et al., 1990; Heymsfield and Miloshevich, 1995), the International Cirrus Experiment, ICE campaign in 1989 (Raschke et al., 1987), European Cloud Radiation EXperiment, EUCREX in 1993

and 1994 (Sauvage et al., 1999), Field Radiation Experiment on Natural Cirrus and High-level clouds, FRENCH in 2001 (Brog-  
55 niez et al., 2004), Tropical Composition, Cloud and Climate Coupling, TC4 campaign in 2007 (King et al., 2010; Toon et al.,  
2010) and CIRRUS-HL campaign in 2021, which is the follow-up to the CIRRUS-ML campaign in 2017 (Voigt et al., 2017;  
De La Torre Castro et al., 2023).

Up to the present, there are three possibilities for characterising cirrus clouds. One option is the use of in-situ airborne  
60 measurements. A second option is to work with a microphysical cirrus cloud model like the Chemical LAgrangian Model of  
the Stratosphere with microphysics scheme for Ice clouds formation (CLaMS-Ice) (Spichtinger and Gierens, 2009) or Model  
for aerosol and ice dynamics (MAID) (Bunz et al., 2008; Rolf et al., 2012), that simulate the cirrus cloud development based  
on the cirrus bulk model by along back-trajectories. The main advantage of this choice is that there is no need to have in-  
situ airborne lidar measurements. A third option is to employ lidar measurements for the characterisation of cirrus clouds.  
65 For that purpose, it is necessary to use a method such as the two-way transmittance method to characterize cirrus clouds op-  
tically (Gil-Díaz et al., 2024) together with a scattering model to obtain ~~radiative~~ optical scattering retrievals. For example,  
(Baran and Labonnote, 2007; Baran et al, 2009, 2011a, b, 2014) relates the cirrus ice water content (IWC) and mid-cloud tem-  
perature with its extinction coefficient, single scattering albedo (SSA) and asymmetry factor (asyF). Alternatively, (Heyms-  
field et al., 2014; Dolinar et al., 2022) propose to calculate the cirrus ice water content from the extinction coefficient at a  
70 visible wavelength and the effective geometric diameter of the ice crystals, which in turn is a function of temperature. Once the  
cirrus ice content and the effective geometric diameter of ice crystals are obtained, the scattering and absorption coefficients  
and the asymmetry factor can be calculated with the (Fu et al., 1998, 1999; Lolli et al., 2017a) parametrizations.

The objective of this paper is to analyze the ~~radiative~~ optical scattering properties and ~~radiative~~ direct radiative effects of cirrus  
75 clouds based on 4 years of continuous ground-based lidar measurements obtained from the NASA Micropulse lidar network  
(MPLNET, <https://mplnet.gsfc.nasa.gov/>) in Barcelona. Specifically, the ~~radiative~~ optical scattering properties of cirrus clouds  
have been calculated with a new approach of the self-consistent scattering model for cirrus clouds (Baran and Labonnote,  
2007; Baran et al, 2014; Vidot et al., 2015), using only lidar measurements and radiosounding data and their direct radiative  
~~radiative~~ effects have been calculated with the ARTDECO package. The instrumentation used is presented in Section 2. A new  
80 approach of the self-consistent scattering model for cirrus clouds, the radiative transfer model DISORT and the CLaMS-Ice  
model are presented in Section 3.1, Section 3.2 and Section 3.3, respectively. The results obtained in this paper are shown in  
Section 4 and conclusions are presented in Section 5.

## 2 Instrumentation

The radiative characterization of cirrus clouds relies on the results obtained from (Gil-Díaz et al., 2024) and the instrumentation  
85 detailed below.

## 2.1 NASA Micro-Pulse Lidar Network

A more detailed description of this instrumentation can be found in (Gil-Díaz et al., 2024). In this study, we use the Aerosol (AER) product, ~~provided at~~ which provides 1-min temporal resolution and at 75m vertical resolution variables like aerosol extinction, backscatter and the column lidar ratio (Welton et al., 2000, 2002, 2018; Lolli et al., 2019). This product is used to  
90 characterize the aerosol layer which is closest to the surface. The MPLNET AER product ~~includes solar and~~ integrates solar and lunar photometer ~~observations to invert~~ measurements, allowing the lidar signal to be inverted to obtain aerosol properties during 24h/d/day over a 24-hour period. Aerosol extinction/backscatter and the column lidar ratio among other properties belong to the MPLNET/AER product (Veltch et al., 2000, 2002).

## 2.2 Meteorological Service of Catalonia

95 The Meteorological Service of Catalonia (Meteocat) releases radiosondes are launched twice ~~every~~ a day (at 00:00 and 12:00 UTC) by the Meteorological Service of Catalonia (Meteocat) at a ~~distance of~~ location less than 1 km from the MPL site. The radiosondes provide ~~the following~~ data of pressure, altitude, temperature, relative humidity, wind speed and direction. Only altitude, pressure and temperature profiles are used in ~~the present work~~ this study.

## 2.3 Solar Radiation Network

100 The Solar Radiation Network (SolRad-Net, <https://solrad-net.gsfc.nasa.gov/>) is a ~~dedicated~~ network of ground-based sensors providing that provides high-frequency solar flux measurements ~~in quasi-real time~~ to the scientific community in near-real time. This program ~~was implemented as a companion to~~ operates in conjunction with AERONET ~~and~~, being its instrumentation ~~is~~ collocated in the AERONET sites. Each SolRad-Net site is initially equipped with two flux sensors: a Kipp and Zonen CM-21 pyranometer (0.305-2.8  $\mu\text{m}$ ) ~~for measuring the total~~ to measure the full solar spectrum and a Skye Instrument SKE-510 PAR  
105 (photosynthetically/active radiation) Energy sensor (~~spectral range~~ 0.4-0.7  $\mu\text{m}$ ).

In For this study, ~~the uncalibrated~~ data from the Kipp & Zonen instrument is ~~employed with the~~ Level 1.0, ~~corresponding to data~~ (unscreened and without final calibration) is used for ~~that do not have final calibration applied during~~ the period from 2018 to 2022. The Kipp & Zonen CM-21 units are ISO 9060 Secondary Standard thermopile pyranometers, equipped with ~~featuring~~ a  
110 receiving element ~~housed beneath~~ enclosed by two concentric Schott K5 glass domes. ~~More information about~~ Further details on the instrument can be found at the following web link (<https://www.kippzonen.com/Product/14/CMP21-Pyranometer>).

## 2.4 Clouds and the Earth's Radiant Energy System project

The Clouds and the Earth's Radiant Energy System project (CERES, <https://ceres.larc.nasa.gov/>) provides ~~observations~~ mea-  
surements of Earth's radiation budget using measurements from CERES instruments onboard the Terra, Aqua and Suomi  
115 National Polar-orbiting Partnership (S-NPP) and NOAA-20 ~~(initially, JPSS-1, named after the Joint Polar Satellite System mission)~~ satellites (Loeb et al., 2016). The ~~goals~~ primary objectives of the CERES project are: (1) to ~~provide~~ create a long-

term, and integrated global climate data record for detecting that can detect decadal changes in the Earth's radiation budget from the surface to the top-of-the-atmosphere; (2) to improve enhance the understanding of the temporal and spatial variability in how Earth's radiation budget varies temporally and spatially and the role that of clouds and other atmospheric properties play; (3) to support climate model evaluation and improvement through model-observation inter-comparisons.

On one hand, the dataset processed is In this study, we use the CERES instantaneous Single Scanner Footprint (SSF) product with the at Level 2/Spectrally, the product analyzed is the observed to analyze observations of the upward fluxes at top-of-the-atmosphere upward fluxes, the surface emissivity in longwave spectral range (5-35 100  $\mu\text{m}$ ), the surface albedo in shortwave spectral (0.2-5  $\mu\text{m}$ ), the surface temperature and the clear-sky percentage during for the period which covers from 2018 to 2022. The Surface albedo is calculated as the unity one minus the surface emissivity. The CERES measurements are either data comes from the AQUA satellite (orbits over Barcelona between 12:00 and 13:00 UTC) from the TERRA satellite (orbits over Barcelona between 10:00 and 10:30 UTC). NOAA-20 satellite (overpassing Barcelona between 11:00 and 13:00 UTC and between 00:00 and 02:00 UTC). This variable dataset is provided with a surface spatial resolution of 20 km at nadir (Su et al., 2015).

On the other hand, the Synoptic N/A and surface fluxes and clouds (SYN) product with the Level 2 is processed to get a radiative/atbedo over the Atlantic/Ocean and the Mediterranean/Sea/roads to study the radiative/forcing produced by clouds cloud among its back-trajectory/spectrally, the variable analyzed is the surface/atbedo from the SYN deg/product, which is integrated between 0.25/1  $\mu\text{m}$ , with a daily/temperature/resolution and a spatial resolution of 1/1  $\mu\text{m}$ .

## 2.5 NASA AeRosol RObotic NETwork

The NASA AeRosol RObotic NETwork (AERONET, <https://aeronet.gsfc.nasa.gov/>) is a federated global network of ground-based sun/lunar-photometers established by NASA and LOA-PHOTONS (CNRS). For more than over two decades, the project AERONET has provided long-term, continuous, and readily/accessible/public domain publicly accessible database of aerosol optical, and microphysical and radiative properties for aerosol characterization research and satellite retrieval validation of satellite retrievals, and synergism with other databases. The network imposes enforces strict standardization of instruments, calibration, processing and distribution.

In this work For this study, Version 3.0, and Level 1.5 (cloud-screened and quality-controlled) inversion products are used for the time period from 2018 to 2022. In order To characterize the aerosols in the lowest tropospheric layer of the troposphere in the shortwave spectrum (SW, 0.2-4  $\mu\text{m}$ ), variables like Aerosol Optical Depth (AOD), Single Scattering Albedo (SSA) and Asymmetry Factor (AsyFg) are used analyzed. For other wavelengths outside than the working wavelength of AERONET's working range, in the SW range spectrum, the Angström exponent is used (Wagner and Silva, 2008). For example, the Angström exponent is calculated with the using AOD values at 440 and 675 nm to obtain estimate the AOD at 550 nm. For the characterization of the aerosol layer In the longwave spectrum (LW, 2-40  $\mu\text{m}$ ), the aerosol optical proper-

ties ~~mentioned above~~ are extracted from the Laboratory for Information Technologies and Mathematical Simulation (LITMS) database (Rublev et al., 1994).

## 2.6 The Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observation

~~A more detailed description of~~ Further details on this instrumentation ~~can be found~~ are provided in (Gil-Díaz et al., 2024). ~~In order~~ To validate the ~~product~~ ice water content from the Lagrangian microphysical cirrus model CLaMS-Ice (see Section 3.3), ~~we use~~ the CALIPSO ~~product used is the~~ "5 km Cloud Layer (05kmCLay)" product/~~with the~~ at Level 2 (L2) and Version 4.20 (V4.20), available from June 2006. This product ~~has a 5-km horizontal averaging resolution with a maximum of ten layers reported per profile~~ offers a horizontal averaging resolution of 5 km, with up to ten layers reported per profile. ~~In particular,~~ This product contains geometrical, thermal and optical properties of each ~~detected~~ cloud layer ~~detected~~ like layer top/base altitude and temperature, integrated attenuated backscattering coefficient at 532 nm and 1064 nm, integrated particle depolarization ratio ~~of~~ and ice water path. The ice water content ~~coefficient of the cloud layer~~ is estimated ~~as the ratio between~~ by dividing the ice water path ~~and~~ by the geometrical thickness of the ~~cloud~~ layer.

## 3 Methodology

The ~~radiative~~ optical scattering properties of cirrus scenes are determined through the use of a self-consistent scattering model for cirrus clouds (Baran and Labonnote, 2007; Baran et al., 2014; Vidot et al., 2015) and their ~~direct~~ radiative ~~forwards~~ effects are calculated with the ARTDECO package, which implements a variety of optical properties into state of the art radiative transfer models (see below).

### 3.1 The self-consistent scattering model for cirrus clouds

The self-consistent scattering model for cirrus clouds consists of an ensemble of six ice crystal ~~members~~ types. ~~Where~~ The simplest form is a ~~ice crystal~~ and ~~represented by~~ hexagonal ice columns with an ~~unity~~ aspect ratio, while ~~of unity~~ and the more complex ice crystals are formed by ~~arbitrarily~~ and randomly ~~oriented~~ attaching ~~more~~ additional hexagonal elements, until ~~creating~~ a chain-like structure ~~ice crystal is formed~~. The complexity of ice crystals tends to increase with increasing size, and As ice crystals grow in size, their complexity increases, and they generally become more spatially extended, with the hexagonal components ~~becoming more elongated~~ elongating over time (Heymsfield and Miloshevich, 2003). The ~~geometrical configuration~~ ~~consists in the~~ ensemble consists of six distinct members, starting with the ~~first of which is~~ the hexagonal ice column/~~the second~~ and followed by a six-branched bullet rosette. ~~Moreover,~~ Subsequent crystals are formed by attaching hexagonal monomers ~~and arbitrarily attached to each other as a~~ together in function of their maximum dimension, ~~forming three to ten element hexagonal ice aggregates~~ resulting in aggregates containing three to ten hexagonal elements. This ensemble ~~tries to represent ice crystals~~ is designed to represent the various ice crystal types observed in cirrus clouds ~~observed in~~ during different measurement campaigns (Heymsfield and Miloshevich, 2003; Lawson et al., 2003; Connolly et al., 2005). For ~~example~~ instance, bullet-rosettes are included in the ensemble as ~~these are commonly~~ they are frequently observed in mid-latitude and



Arctic regions (Lawson et al., 2006; Schmitt et al., 2006) *with their* geometry is described in (Macke et al., 1996). *The*  
*Each* ensemble members *is* constructed *so as to avoid* intersecting planes, and the crystals are *arranged so* that multiple reflections between them are negligible, which was *verified* experimentally using  
185 ray-tracing calculations. The first *member* represents the smaller ice crystals in the particle size distribution (PSD), whilst the hexagonal ice aggregates represent the process of ice crystal aggregation, *corresponding*  
*to* the larger ice crystals in the PSD. The *six* members are *evenly* distributed *across* the PSD. The *PSD* assumed *follows the* moment estimation parametrization of *the*  
*PSD*, *as proposed by* (Field et al., 2003, 2007). The self-consistent scattering model for cirrus cloud database  
190 consists of more than 20000 PSDs of tropical and mid-latitude cirrus at temperatures between -60 and 0°C. This database  
provides for each PSD the simulated optical scattering properties (scattering, absorption coefficient and asymmetry factor) in  
function of the decimal logarithm of ice water content (LIWC) and the cloud temperature (T) by nonlinear least squares fitting  
(Vidot et al., 2015). This parametrization excludes ice crystals smaller than 100 μm due to shattering issues on closed-path  
instruments (Korolev et al., 2011), assuming instead an exponential PSD fit. For crystals larger than 100 μm, measured PSDs  
195 were filtered to reduce the likelihood of including shattered crystal artifacts.

*where the scattering, absorption coefficient and asymmetry factor are a function of the cloud temperature (T) and the ice water content (LIWC) as shown in the following expressions:*

$$\begin{aligned}
 & \log_{10}(Q_{ext}(\lambda, T)) = \log_{10}(Q_{ext}(\lambda, T)) + \log_{10}(LIWC(\lambda, T)) + \log_{10}(T^2) + \log_{10}(LIWC^2(\lambda, T)) \\
 & + \log_{10}(T) + \log_{10}(LIWC(\lambda, T))
 \end{aligned} \tag{1}$$

$$\log_{10}(Q_{sca}(\lambda, T)) = \log_{10}(Q_{sca}(\lambda, T)) + \log_{10}(LIWC(\lambda, T)) \tag{2}$$

*where the coefficients depend on the degree of the wavelength and the wavelength of the incident radiation. The values of the variables are given in Table 1. The parametrization is based on 10000 simulated ice crystals of the PSD, obtained in a laboratory and in a field campaign. The parametrization is based on 10000 simulated ice crystals of the PSD, obtained in a laboratory and in a field campaign. The parametrization is based on 10000 simulated ice crystals of the PSD, obtained in a laboratory and in a field campaign.*  
205 *where the coefficients depend on the degree of the wavelength and the wavelength of the incident radiation. The values of the variables are given in Table 1. The parametrization is based on 10000 simulated ice crystals of the PSD, obtained in a laboratory and in a field campaign. The parametrization is based on 10000 simulated ice crystals of the PSD, obtained in a laboratory and in a field campaign.*

210 In this study, *we present* a new methodology for *calculating* the *radiative* optical scattering properties of cirrus clouds *at* *across* different wavelengths *is proposed*, as shown in Fig. 1. First, *we calculate* the *LIWC* ice water content of the cirrus cloud *by* Eq. 1, which is independent of wavelength, by introducing in Eq. 3 by (Vidot et al., 2015) the extinction coefficient for a cloud temperature in each vertical layer of the model. To align these calculations with the model vertical resolution, we previously degrade the vertical resolution of the cloud extinction and temperature profiles through vertical averaging.

215 that has no spectral dependence is calculated by introducing in Eq. A the effective column extinction coefficient calculated with the two-way transmittance method at 0.382  $\mu\text{m}$  (the working wavelength of the MPL), for a fixed cloud temperature. Eq. 4 is derived from Eq. 4, where the left-hand side is assumed to be a simple way of calculating the IWC from an extinction coefficient and a cloud temperature has been implemented by assuming the absence of absorption, which is entirely reasonable because the wavelength used belongs to the visible spectral range (Sun and Shine, 1994).

$$IWC(\lambda, t) = \frac{\pi (r_{eff, \lambda} T(\lambda, t)) / N \sqrt{(r_{eff, \lambda} T(\lambda, t))^2 + 4 E_s \Theta(\lambda, t)}}{2 r_{eff, \lambda} / \lambda_o} \quad \lambda_o \in [0.38, 0.7] \mu\text{m}$$

$$\Theta(\lambda, t) = \log_{10} \sigma_{ext, \lambda_o}(\lambda, t) - A_s - B_s T(\lambda, t) - D_s T^2(\lambda, t)$$

220

(3)

$$LIWC(z, t) = \frac{-(C_s + F_s T(z, t)) + \sqrt{(C_s + F_s T(z, t))^2 + 4 E_s \Theta(z, t)}}{2 E_s} \quad \lambda_o \in [0.38, 0.7] \mu\text{m}$$

$$\Theta(z, t) = \log_{10} \sigma_{ext, \lambda_o}(z, t) - A_s - B_s T(z, t) - D_s T^2(z, t)$$

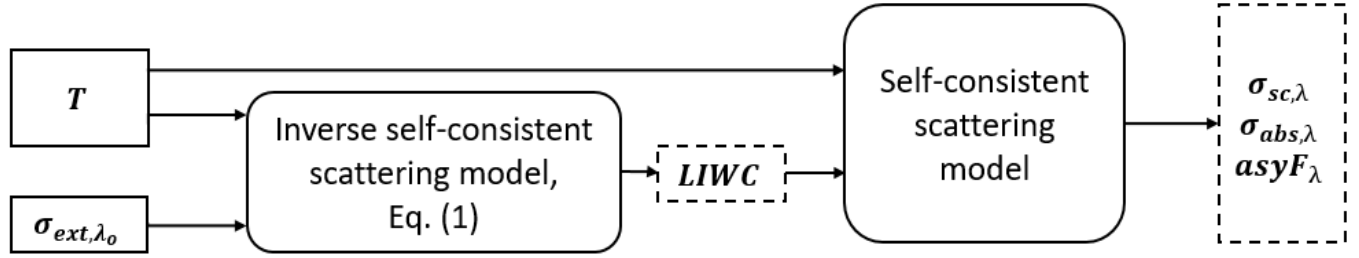
(4)

The working wavelength of the lidar system is defined as  $\lambda_o$  and in this study it is 0.532  $\mu\text{m}$ . The extinction coefficient,  $\sigma_{ext, \lambda_o}$ , is retrieved by the two-way transmittance method (Gil-Díaz et al., 2024).  $T(z, t)$  denotes the cloud temperature, while  $A_s$  to  $F_s$  are the parametrization coefficients as defined in (Vidot et al., 2015). This formulation provides a unique physical solution and simplifies the IWC calculation based on extinction coefficient and cloud temperature, assuming no absorption, which is entirely reasonable because the working wavelength lies within the visible spectral range (Sun and Shine, 1994). Once, the IWC is obtained, this variable is introduced into Eqs. 2-4 (Vidot et al., 2015) to calculate the absorption, scattering and asymmetry factor coefficients for each wavelength. From these variables, the extinction and single scattering albedo are also subsequently derived. Once, the IWC of the cirrus cloud is calculated, this variable is introduced in Eq. A to get the absorption, scattering and asymmetry factor coefficients for each wavelength. From these variables, the extinction and single scattering albedo are also calculated for each wavelength. This method is valid for any lidar working wavelength which belongs to the visible spectral range.

225

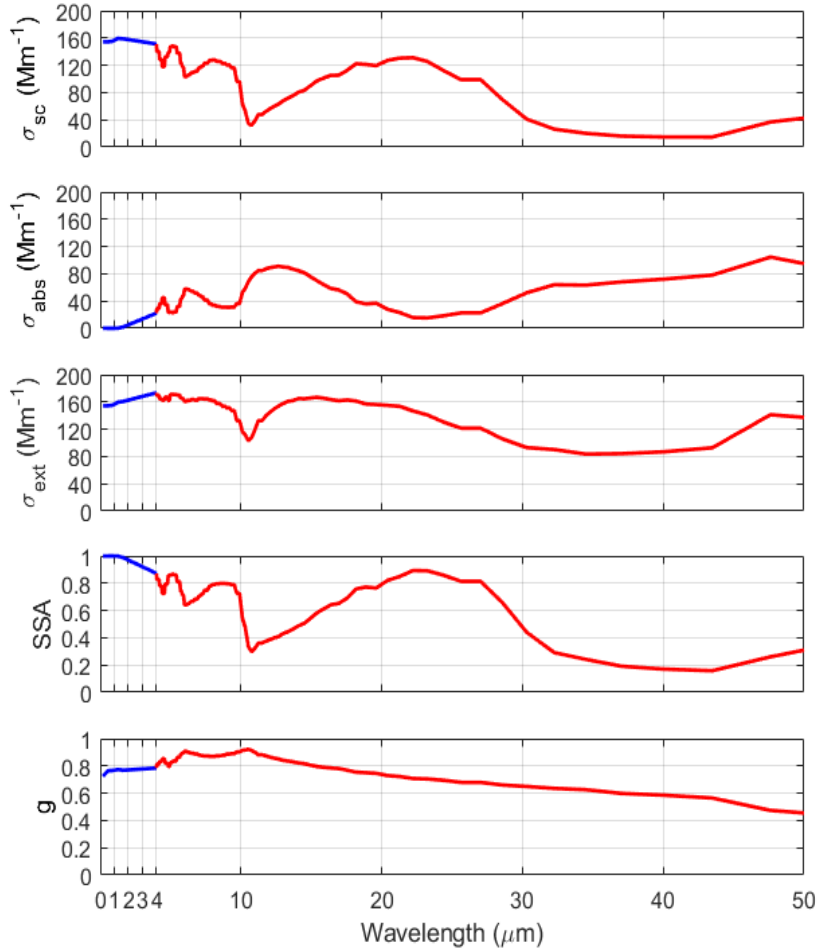
230





**Figure 1.** Scheme of the new model approach of the self-consistent scattering model for cirrus clouds.

235 The main advantage of this model is that it is not necessary to calculate the ice particle shapes and their size distributions in order to calculate their ~~radiative~~ **optical scattering** properties. These variables can be obtained easily with only elastic lidar systems and radiosondes or meteorological models. An example of the application of this method is shown in Fig. 2, for a cirrus cloud **layer** of 08/12/2018 at 12 UTC, measured in Barcelona.



**Figure 2.** Spectral dependence of ~~radiative~~ **optical scattering** properties obtained by the self-consistent scattering model for cirrus clouds: scattering, absorption, extinction coefficients, single scattering albedo and asymmetry factor (from top to bottom) for a cirrus cloud of 08/12/2018 at 12 UTC, measured in Barcelona. The colours indicate the shortwave range (blue; 0.2-4  $\mu\text{m}$ ) and longwave range (red; 4-50  $\mu\text{m}$ ).

Fig. 2 shows the spectral dependence of ~~radiative~~ **optical scattering** properties obtained by the self-consistent scattering  
 240 model for cirrus clouds, which reflects the characteristic ~~radiative~~ **optical scattering** properties of ice crystals. For example, it shows an absorption phenomenon negligible in the spectral range between 0.2 and 2  $\mu\text{m}$  and consequently, the single scattering albedo is approximately the unity. As expected, the scattering phenomenon dominates the whole spectrum with respect to absorption except in the regions around 12  $\mu\text{m}$  and for wavelengths higher than 30  $\mu\text{m}$  (Yang et al., 2005, 2013). It is also noteworthy to mention that the single scattering albedo varies generally between 0.1 and 1 and at the working wavelength used

245 in the model (0.532  $\mu\text{m}$ ) it has a value of 0.99 (Sun and Shine, 1994). This fact supports the hypothesis made previously that the absorption phenomenon is negligible. The asymmetry factor presents much less variation as observed in the literature (Fu et al., 1998; Yang et al., 2005, 2013): it increases between 0.2 and 10  $\mu\text{m}$  (in the range 0.75 - 0.95) and decreases afterwards (in the range 0.95 - 0.40).

### 3.2 The ARTDECO package

250 The Atmospheric Radiative Transfer Database for Earth and Climate Observation package (ARTDECO; <https://www.icare.univ-lille.fr/artdeco/>) is a numerical tool that gathers models and data for the 1D simulation of Earth atmosphere radiances and fluxes from the ultraviolet to thermal infrared range (0.2-50  $\mu\text{m}$ ). It is developed and maintained at the Laboratoire d'Optique Atmosphérique (LOA) and distributed by the data and services center AERIS/ICARE (University of Lille), and funded by the TOSCA program of the French space agency (CNES). In ARTDECO, users can either access a library for the scene or use  
 255 their own description through ASCII input files. Optical properties for aerosols and clouds can be computed. Then, the user can choose among available models to solve the radiative transfer equation and to compute radiative quantities corresponding to the scene. ARTDECO is thus a flexible tool for remote sensing or radiative forcing applications, such as sensitivity studies, development and optimization of retrieval algorithms, evaluation of the future instruments performances, etc.

260 In this study, DISORT model is employed to solve the radiative transfer equation by discretising it (Stamnes et al., 2000). The ARTDECO environment allows us to solve the radiative transfer equation in two ways: 1) by introducing our own phase matrix as a function of wavelength; 2) by using the Henyey-Greenstein function (Henyey and Greenstein, 1941), given extinction, single scattering albedo and asymmetry factor values over the whole spectral range in which the simulation will be done. Due to the lack of knowledge of the phase matrix of cirrus clouds with the observational measurements with which we work, the  
 265 second option is chosen even though the Henyey-Greenstein function does not represent a good approximation to the real phase function, especially for ice crystals. Upward and downward radiative fluxes are calculated at different vertical Levels: 31 layers (0-20 km) in the shortwave (SW, 0.2-4  $\mu\text{m}$ ) and 40 layers (0-100 km) in the longwave (LW, 4-50  $\mu\text{m}$ ) spectra. These spectral/vertical ranges are adjustable, together with their spectral/vertical resolution. Cirrus ~~forcing~~ **direct radiative effects** at the bottom-of-the-atmosphere (BOA) and top-of-the-atmosphere (TOA) have been calculated as:

$$270 \text{ BOA } DR\textcolor{red}{F\textcolor{red}{E}} = (F_c \downarrow - F_c \uparrow) - (F_o \downarrow - F_o \uparrow) \text{ at BOA} \quad (5)$$

$$TOA \text{ DR}\textcolor{red}{F\textcolor{red}{E}} = (F_c \downarrow - F_c \uparrow) - (F_o \downarrow - F_o \uparrow) = -(F_c \uparrow - F_o \uparrow) \text{ at TOA} \quad (6)$$

Where  $F_c$  and  $F_o$  are the radiative fluxes with and without the cirrus cloud, respectively. The  $\downarrow$  and  $\uparrow$  arrows indicate whether the fluxes are downward or upward, respectively. The simplification of Eq. 6 implies the assumption that the amount of the  
 275 incoming solar radiation at the TOA is equal for both cases with and without aerosols. With this convention, a negative sign of  $DR\textcolor{red}{F\textcolor{red}{E}}$  implies a cirrus cooling effect independently of whether it occurs at the BOA or at the TOA. In this study, four types of simulations are carried out: with gases only (G), with gases and aerosols in the layer closest to the surface (GA), with gases

and a cirrus cloud (GC) and finally, with everything (GAC). In Section 4.3, the **direct radiative forcing effects** of cirrus clouds in the full atmosphere and in the whole spectral range considering aerosols (GAC-GA) or no aerosols (GC-G) are compared.

280 Besides aerosol optical properties, the radiative transfer model (RTM) DISORT is sensitive to atmospheric parameters such as the relative humidity and the air temperature profiles, the surface emissivity and temperature or the aerosol vertical distribution (Sicard et al., 2014).

The cirrus clouds are parameterized in the RTM model geometrically and optically, with the results obtained (Gil-Díaz et al., 2024) and ~~radiatively/with~~ the retrievals obtained with the self-consistent scattering model for cirrus clouds (see Section 3.1). At the same time, the planetary boundary layer is characterized in the DISORT model geometrically and optically with the MPLNET AER product and radiatively, on one hand, in the SW with the AERONET products and, on the other hand, in the LW with the LITMS database.

### 3.2.1 Atmospheric profiles

290 The RTM DISORT model is run with atmospheric profiles (pressure, temperature, water vapor mixing ratio) obtained from radiosondes launched in Barcelona at 00 and 12 UTC. The profiles of ozone concentration are obtained from Copernicus Atmosphere Monitoring Service (CAMS) global reanalysis (EAC4) (Inness et al., 2019). EAC4 (ECMWF Atmospheric Composition Reanalysis 4) is the fourth generation ECMWF global reanalysis of atmospheric composition. Reanalysis combines model data with observations from across the world into a globally complete and consistent dataset using a model of the

295 atmosphere based on the laws of physics and chemistry. The dataset is globally distributed with a horizontal resolution of  $0.75^\circ \times 0.75^\circ$  and a vertical extension of 60 modes (from 1000 to 1 hPa). In exceptional cases, when no radiosondes or CAMS data are available and for heights not covered by the radiosondes (generally above 30 km), the atmospheric profiles are taken from the 1976 standard atmosphere (COESA et al., 1976).

### 3.2.2 Surface properties

300 In this study, a Lambertian surface is considered. On one hand, the corresponding surface albedo over the Barcelona region is obtained for the SW range from AERONET and for the LW spectrum from ~~CERES~~ **NOAA-20** measurements. On the other hand, the surface temperature is also taken from ~~the SSF/CERES/product~~ **NOAA-20 observations**. In the parametrization of cirrus scenes, ~~this surface albedo is averaged seasonally and the surface temperature monthly, differentiating between values at daytime and nighttime, during the period from 2018 to 2022. Surface albedo has been averaged seasonally due to its smoother temporal variation compared to temperature, ensuring sufficient data availability~~ **the surface emissivity in the longwave spectrum, the surface albedo in the shortwave spectrum and the surface temperature measured by the NOAA-20 satellite have been incorporated as instantaneous values. In contrast, the surface albedo measured by AERONET has been averaged seasonally due to its smoother temporal variation and its small influence on the simulations in that spectrum.**

305

### 3.2.3 Cloud/aerosol stratification

310 The vertical stratification of cloud/aerosols is reproduced according to the vertical profiles of MPLNET products. On one hand, the base and top of the cirrus clouds are obtained from (Gil-Díaz et al., 2024). On the other hand, the vertical distribution of the aerosols in the planetary boundary layer (PBL) is provided by MPLNET AER product. When this product is not available for that specific time period, it is assumed that aerosols are uniformly distributed throughout the aerosol layer, which extends up to 1.5 km, being the mean PBL height obtained in Barcelona over a 3-year period (Sicard et al., 2006).

### 315 3.3 The CLaMS-Ice Lagrangian microphysical cirrus model

CLaMS-Ice is a Lagrangian model for cirrus microphysics (Luebke et al., 2016; Baumgartner et al., 2022), intended to compute the evolution of ice microphysics along air parcel trajectories. The Chemical Lagrangian Model of the Stratosphere (CLaMS) (McKenna et al., 2002a, b; Konopka et al., 2004) performs the analysis of air mass back trajectories starting at arbitrary location in the atmosphere. The trajectories are derived from ECMWF windfields and are selectable over an arbitrary time frame. Small scale temperature fluctuations not considered in the ECMWF wind fields are accounted for by superimposing temperature fluctuations according to (Podglajen et al., 2016). Next, the CLaMS-Ice model is run in the trajectories forward direction by using the two moment box-model developed by (Spichtinger and Gierens, 2009) to simulate cirrus cloud development. The two-moment scheme includes homogeneous as well as heterogeneous nucleation of ice, depositional growth of ice crystals, their evaporation, aggregation, and sedimentation (sedimentation parametrization after (Spichtinger and Czicz, 2010); the sedimentation parameter is set to 0.97). Heterogeneous freezing starts at a critical supersaturation of 120%. The initial concentration of ice nucleating particles is prescribed in the model (mean value: 0.01 cm<sup>-3</sup>). As ice particles evaporate, ice nucleating particles are released back into the air parcel. The model predicts the ice number concentration and the ice water content. ~~If ice is already present in the ECMWF data at the start of the trajectory, CLaMS-Ice treats this as pre-existing ice. Two-moment scheme only considers the trajectories that end at T < 238 K. If a part of the trajectory existed at T > 238 K before crossing into the colder cirrus environment, then the forward model is initialized with pre-existing ice from mixed-phase clouds, if present in the IWCs found in the ECMWF data. CLaMS-Ice proceeds with pre-existing ice as with newly formed cirrus clouds.~~ The two-moment scheme only considers the trajectories that end at T < 238 K. If a part of the trajectory existed at T > 238 K before crossing into the colder cirrus environment, then the forward model is initialized with pre-existing ice from mixed-phase clouds, if present in the IWCs found in the ECMWF data. CLaMS-Ice proceeds with pre-existing ice as with newly formed cirrus clouds. 335 number concentration using a reparametrization (Costa, 2017), then CLaMS-Ice proceeds as for newly formed cirrus clouds.

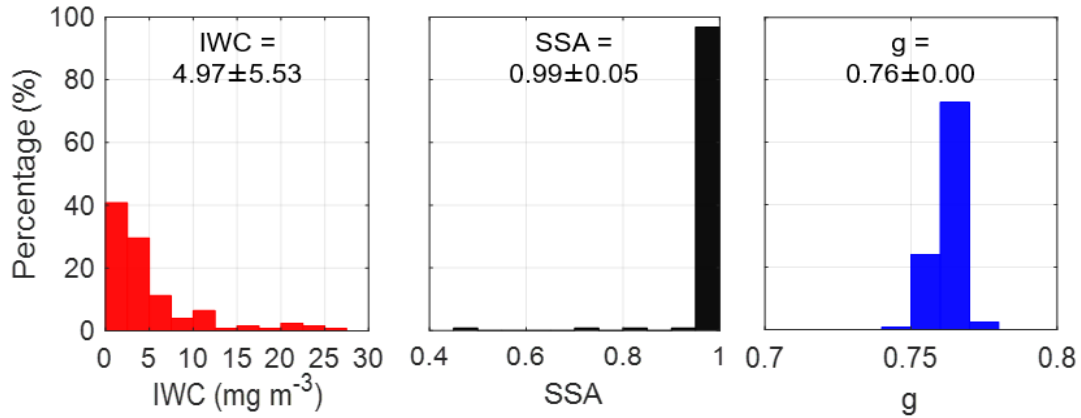
## 4 Five years of cirrus retrievals

This study analyzes the ~~radiative~~ optical scattering properties of cirrus clouds that were previously characterized geometrically, thermally and optically in (Gil-Díaz et al., 2024). In the latter paper, 203 cases were analyzed. Here, atmospheric scenes with only one cirrus cloud in the vertical profile are studied. 125 single-cirrus cloud scenes are found, that constitute 61% of all

340 cases. Case selection involves discarding cirrus scenes where a mid-level cloud has been detected below the cirrus cloud. Mid-level cloud detection was performed through visual analysis during the application of the two-way transmittance method (Gil-Díaz et al., 2024). This selection ensures the simulated cirrus scenes accurately represent reality (see Section 4.2).

#### 4.1 Cirrus ~~radiative~~ optical scattering properties

After having carried out the identification of 125 high-altitude cirrus scenes with only a cirrus cloud in the vertical profile (being the 39% of cirrus clouds measured at daytime), the self-consistent scattering model for cirrus clouds is applied to obtain their ~~radiative~~ optical scattering properties (see Section 3.1). In this section the ~~radiative~~ optical scattering properties of cirrus clouds are presented and discussed. For this purpose, probability distributions of the following vertical averages of ~~radiative~~ optical scattering properties for each cirrus scene: ice water content, single scattering albedo and asymmetry factor, calculated at  $0.55 \mu\text{m}$  are shown in Fig. 3. ~~For each case, the IWC and the radiative optical scattering properties of the cirrus clouds have been obtained using the mean cloud temperature and the effective extinction coefficient of the cloud for each model vertical layer in which the cloud is extended. This calculation is performed by averaging the vertical profiles of the cloud extinction coefficient over half the cloud thickness and centered on the maximum peak to obtain a value representative of the entire cloud.~~



**Figure 3.** Probability distribution of averages of (left) ice water content (IWC), (center) single scattering albedo (SSA) and (right) asymmetry factor ( $g$ ) for each cirrus scenes at  $0.55 \mu\text{m}$ , for cirrus clouds measured from 2018 to 2022 in Barcelona. ~~The values of the left histogram in the following table.~~

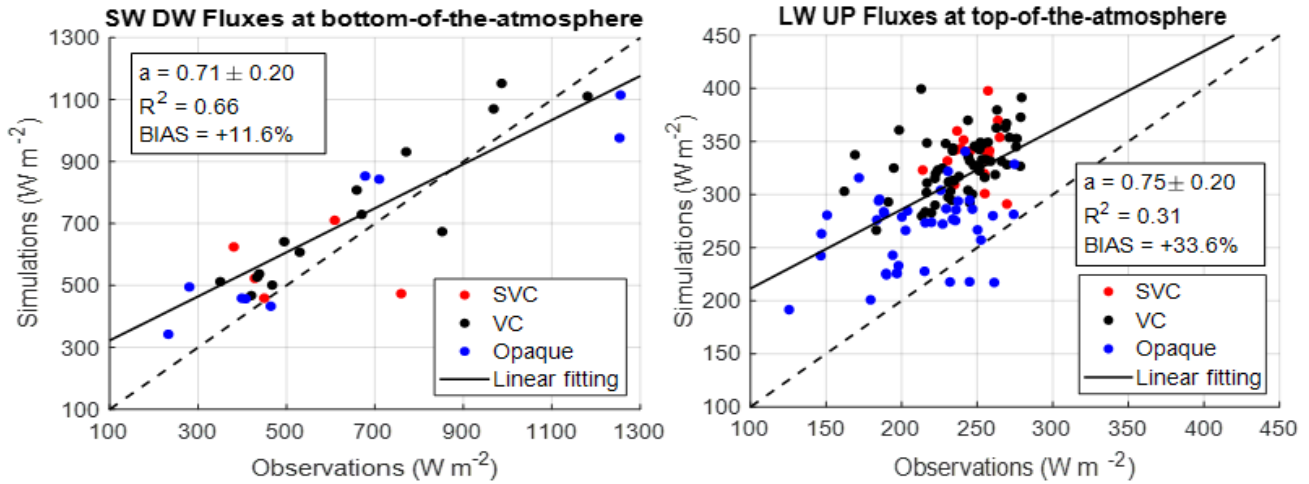
In Fig. 3 one observes that cirrus clouds have an IWC between  $0.03$  and  $30 \text{ mg m}^{-3}$ , being characteristic of mid-latitude cirrus clouds (Korolev et al., 2001; Field et al., 2005, 2006; Schiller et al., 2008; Baran et al, 2011b; Sourdeval, 2012; Kramer et al., 2016, 2020). Where the average of IWC is  $\sim 5 \text{ mg m}^{-3}$ , being a value close to  $3 \text{ mg m}^{-3}$ , which is the central value of the mid-latitude ice cloud distributions obtained by (Sourdeval, 2012) and the mean value of IWC for temperatures between 210 and 235K found in (Kramer et al., 2016). A slightly higher measured IWC value of  $7 \text{ mg m}^{-3}$  was found by (Korolev et al., 2001) for cirrus clouds whose temperature ranged from 233 to 243K. The single scattering albedo of most cirrus clouds (97%) has a



value between 0.95 and 1, as expected at  $0.55\ \mu\text{m}$  (Hess and Wiegner, 1994; Sun and Shine, 1994; Yang et al., 2013; Hemmer, 2018). Although there are 3 cases of cirrus clouds whose  $\text{SSA} < 0.9$ . These cases correspond to sub-visible cirrus clouds with an IWC less than  $1\ \text{mg/m}^3$ . These 3 cases have in common that the cirrus cloud extends in two vertical layers of the model and in one of the layers, the ~~resulting value by the effective column~~ extinction coefficient is less than  $1\ \text{Mm}^{-1}$ . In this layer, a low value of SSA is obtained, associated to its low value of the extinction coefficient and consequently, when averaging the ~~radiative~~ **optical scattering** properties in the two layers for each cirrus scene, the values of  $\text{SSA} < 0.9$  observed in Fig. 3 are obtained. Therefore, the self-consistent scattering model for cirrus clouds might associate the low ~~effective column~~ extinction values to super-cooled liquid water content in the cirrus clouds. The asymmetry factor of cirrus clouds varies between 0.7 and 0.8, with an average of 0.76, as expected at  $0.55\ \mu\text{m}$  (Hess and Wiegner, 1994; Sun and Shine, 1994; Yang et al., 2013; Hemmer, 2018).

## 4.2 Validation of radiative fluxes

The validation of the ARTDECO package is performed by comparing the simulated radiative fluxes with observed ones. For that purpose, the radiative fluxes from ARTDECO were recalculated in the range  $0.305\text{-}2.8\ \mu\text{m}$  corresponding to the spectral range of the SolRad-Net pyranometer (BOA), and in the range  $5\text{-}100\ \mu\text{m}$  corresponding to the spectral range of ~~CERES~~ **NOAA-20** (TOA). The scatter plot of the simulated vs. observed SW downward radiative fluxes of cirrus clouds classified according to their cloud optical depth (Sassen and Cho, 1992), at the surface and the LW upward radiative fluxes at TOA are shown in Fig. 4.



**Figure 4.** Comparison of (left) simulated shortwave downward (SW DW) radiative fluxes at the bottom-of-the-atmosphere and SolRad-Net observations; (right) simulated longwave upward (LW UP) radiative fluxes at top-of-the-atmosphere and ~~CERES~~ NOAA-20 observations. The black dashed line is the curve with the slope unity and the black solid line is the linear fitting of the fluxes ( $y=ax+b$ , being  $a$  the slope and  $R^2$  its determination coefficient). The cirrus clouds have been classified according to (Sassen and Cho, 1992) criteria: sub-visible (SVC;  $COD < 0.03$ ), visible (VC;  $0.03 < COD < 0.3$ ) and opaque clouds (Opaque;  $COD > 0.3$ ).

The validation of the SW downward radiative fluxes is performed with 59% of the cirrus clouds measured at daytime and the validation of the LW upward radiative fluxes with 81% of all cirrus clouds considered in this study. The cases of cirrus clouds that could not be validated are due to the lack of observations. In addition, to reduce the effect of cloud movement on the radiation measurement with the pyranometer, the observed radiation fluxes are averaged over 30 minutes. In Fig. 4 (left) it can be seen that most downward radiative fluxes calculated with the DISORT model overestimate the SolRad-Net observations: the mean and standard deviation of the simulated fluxes are  $694 \pm 247 \text{ Wm}^{-2}$ , while it is  $621 \pm 283 \text{ Wm}^{-2}$  for the observations. This translates into a systematic BIAS of +11.6% with a steep slope of the linear regression ( $a = 0.71 \pm 0.20$ ). The overestimation may be related to the error associated with variables obtained by the self-consistent scattering model for cirrus clouds, as cirrus clouds govern the radiation interactions in these simulations, because of their cloud optical depth. The ensemble scattering model for cirrus clouds has a large error for small ice crystals (less than  $100 \mu\text{m}$ ), corresponding to cirrus clouds with low IWC values (Liou et al., 2008). In particular, the model tends to underestimate the IWC for mid-latitude cirrus clouds (Baran and Labonnote, 2007). Therefore, when the IWC is lower, the extinction of cirrus clouds is smaller as demonstrated in (Fu, 1996; Heymsfield et al., 2014) and, consequently, allows more radiation to pass into the atmosphere than actually does.

As For validation in the longwave spectrum, ~~CERES~~ NOAA-20 measurements have been selected based on their ~~geographical position~~ geographical proximity to the Barcelona lidar station. Specifically, The measurements closest to ~~the Barcelona lidar~~

this station have been selected were chosen, despite the fact that the hour of the measurement does not correspond to the exact hour of even though the measurement times do not exactly match the atmospheric scene. The hourly difference between simulation and observation is not relevant in the validation of the LW radiation fluxes. On average, the time difference between observations and simulations is 1 hour, with a maximum difference of 3 hours. However, this discrepancy is not significant for the validation of longwave radiation fluxes, since it is almost constant during daytime hours (Sicard et al., 2014). In Fig. 4 (right) a large horizontal dispersion can be observed. In addition, a general overestimation of the CERES NOAA-20 observations with the ARTDECO simulations is produced, being for simulated fluxes  $310 \pm 43 \text{ Wm}^{-2}$  and for the observations  $232 \pm 32 \text{ Wm}^{-2}$ . In our case, the large BIAS = +33.6 % obtained could be due to the spatial resolution of the observed measurements taken, which is only 1 km and located at the Balearic Islands station and may cover part of the Mediterranean Sea. In addition, the cloud mask associated with each observation indicates that in 14% of the cases it has more than 90% of clear-sky footprint area. As demonstrated by Gil-Díaz et al., (2024) most of the cirrus clouds are visible and therefore their horizontal expansion is smaller than the cirrus clouds that form at higher altitudes (well-known as sub-visible cirrus clouds) (Kramer et al., 2020). This makes them more challenging to detect from top-of-the-atmosphere. Hence, the comparison of simulated radiative fluxes and CERES NOAA-20 observations is not as trivial and conclusive as with SolRad-Net observations, since the CERES NOAA-20 satellite can observe a slightly different atmospheric situation, as mentioned above. Not to mention the limitations of the 1-D radiative transfer model DISORT to represent an irregular composition of broken and/or overlapping clouds that the CERES NOAA-20 satellite could observe.

### 4.3 Study of the influence of aerosols on radiative simulations of cirrus clouds

The cloud direct radiative forcing effect is often calculated as the difference between radiative fluxes under cloudy and cloudy-free conditions, without considering the aerosols found in the layer of the atmosphere closest to the surface (Ramanathan et al., 1989; Hartmann et al., 2001; Barja and Antuña, 2007; Yang et al., 2007; Lee et al., 2009; Campbell et al., 2016; Lolli et al., 2017b). In this way, simpler simulations are carried out, as it is only necessary to characterize the cloud. It is known that radiation does not interfere linearly with the components of the atmosphere: clouds, aerosols and gases. Therefore, in this subsection we analyze whether or not the insertion of aerosols in the lowest atmospheric layer in cirrus cloud scenes modifies the calculation of cirrus direct radiative forcing effects. For this purpose, the net direct radiative effect forcing (NET; SW+LW) in the atmosphere (ATM; TOA-BOA) of cirrus clouds calculated with simulations in which aerosols have been considered or not is compared, as shown in Fig. 5.

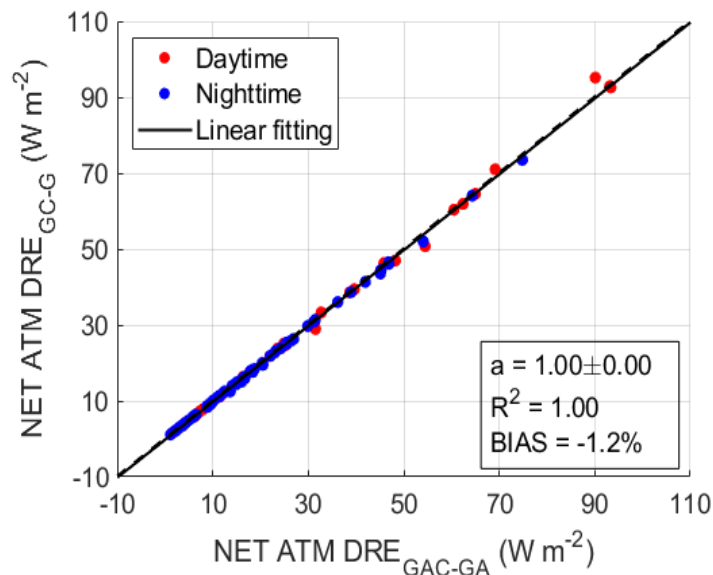
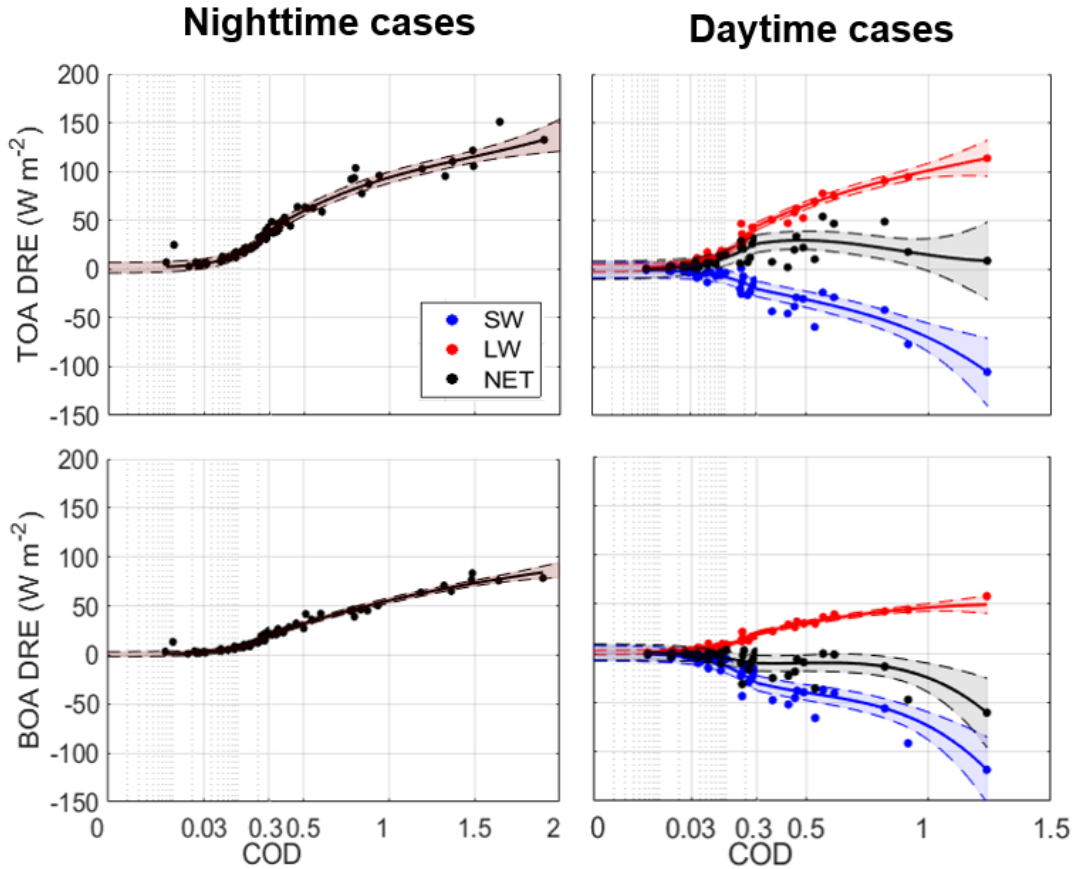


Fig. 5 shows that the NET ATM **direct radiative effects** calculated with and without aerosols fit well, with the most of the points lying slightly above on the curve with the slope unity. As a consequence, its linear fitting rounding to the tenth has also a unity slope, with a  $R^2$  value of 1.00. In addition, the mean and standard deviations of simulations reflect that there is an negligible underestimation of the forcings when not considering aerosols, with values for the simulation with aerosols (X-axis) of  $+21.2 \pm 20.3 \text{ Wm}^{-2}$  and without aerosols (Y-axis) of  $+20.9 \pm 20.3 \text{ Wm}^{-2}$ . Furthermore, the BIAS is  $-0.71.2\%$ , being a low value, possibly due to the distance between the aerosol layer and the cirrus cloud (being on average  $6.76 \pm 2.24 \text{ km}$ ). With these results where the aerosol layer was well distinguished vertically from the cirrus clouds, the simplification of the atmospheric scenes can be made without considering aerosols, but to be more rigorous, in the following results, only the forcings in which aerosols are present will be considered. In the other case where the aerosols are vertically closer to the clouds (lower than 1 km, being the minimum distance found between the cirrus cloud and the aerosol layer), the simplification of not considering aerosols in the calculation of cloud forcings may not be valid, leading to a significant underestimation of cloud **forcings direct radiative effects**.

#### 4.4 Direct radiative forcing effects of cirrus clouds depending on COD

In this section, only direct radiative forcing effects of cirrus clouds calculated with simulations in which aerosols are present, DRFE<sub>GAC-GA</sub>, will be considered and will be denoted as DRFE. Special attention will be paid to net direct radiative forcing effects of cirrus clouds at daytime because they are the only clouds that can readily cool or warm the top and bottom of atmosphere, during daytime, depending on their properties (Campbell et al., 2016). In order to quantify this phenomenon, cirrus clouds at daytime and nighttime have been distinguished. The net direct radiative forcing effects of cirrus clouds at nighttime and daytime, at BOA and TOA are shown in Fig. 6.



**Figure 6.** Distribution of direct radiative forcing effects of cirrus clouds at (left) nighttime and (right) daytime, at (bottom) bottom-of-the-atmosphere (BOA) and (top) top-of-the-atmosphere (TOA), in function of their cloud optical depth (COD). Shortwave (SW), longwave (LW) and net (NET = SW + LW) components of direct radiative forcing effects have been distinguished. For COD < 0.3 a logarithmic scale has been considered in order to discern more clearly sub-visible and visible cirrus clouds (Sassen and Cho, 1992). The solid line corresponds to the polynomial fitting performed on the data set. The shaded area represents the region with a 95% probability of containing the points, adjusted by the mean absolute value of the differences between the actual and fitted values.

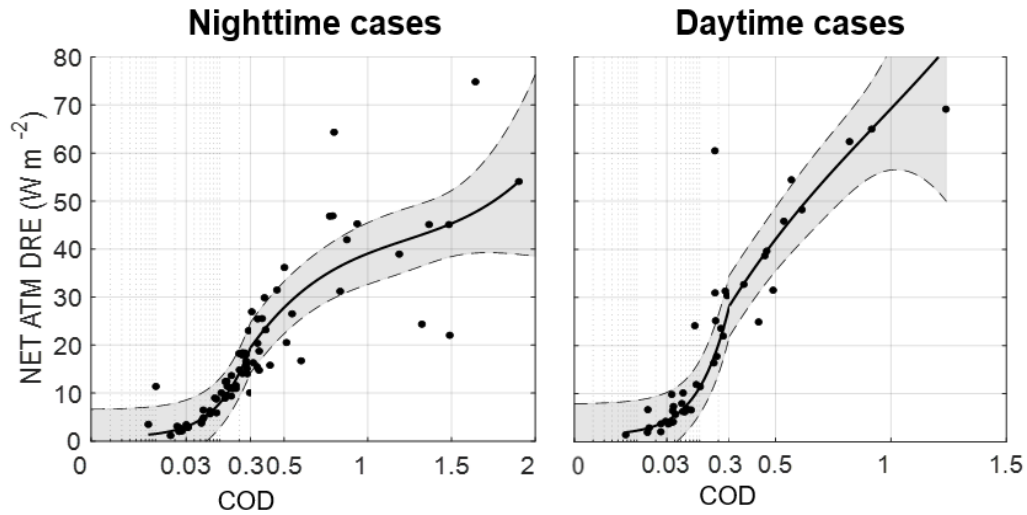
440 In Fig. 6 one observes a positive trend between the net **direct radiative ~~forcing~~ effects** with the COD, where the thicker cirrus clouds contribute more to the overall **~~forcing~~ direct radiative effect** budget, as has been observed in other studies (Barja and Antuña, 2007; Lee et al., 2009; Campbell et al., 2016; Lolli et al., 2017b). Some COD gaps are also found, because the cirrus observations considered do not have a homogeneous and equidistant distribution of COD. At nighttime, the net cirrus **~~forcing~~ direct radiative effect** at TOA is approximately twice that at BOA, being always positive as expected. Cirrus  
445 clouds at nighttime act as a cover in the atmosphere, they do not let through all the infrared radiation emitted by the Earth as it cools, inducing a warming of the atmosphere. This warming in function of COD is **~~faster~~ more pronounced** at TOA than at BOA because the atmosphere at BOA is strongly influenced by the surface, which acts as a black body emitting infrared radiation at nighttime. Consequently, the **~~heating~~ positive direct radiative effect** at BOA is milder than at TOA.

450 At daytime, at TOA, the net **direct radiative ~~forcing~~ effect** remains positive for all cirrus clouds, dominating the positive longwave component. **This effect has been observed in (Campbell et al., 2016) for COD up to approximately 0.6. For higher COD values, (Campbell et al., 2016) reports a negative NET TOA DRE. In this study, a decreasing trend in NET TOA DRE is observed from COD values of 0.5, although no negative values are obtained. Additionally, the LW NET TOA DRE component grows faster than the one reported by (Campbell et al., 2016), suggesting that negative values of NET TOA DRE could occur**  
455 **for cirrus clouds with higher COD than those found in (Campbell et al., 2016). This discrepancy may be due to the higher surface emissivity and temperature values considered in the present work. Further measurements of NET TOA DRE for cirrus clouds with higher COD are needed to confirm the decreasing trend.** In contrast, at BOA, the net **direct radiative ~~forcing~~ effect** is almost always negative (only **~~20~~18%** of the cases show a positive net **~~forcing~~ direct radiative effect**, whose value is close to 0), being the outgoing shortwave radiation in the presence of cirrus clouds larger than in cirrus cloud free conditions. The  
460 albedo effect overcomes the greenhouse effects in the SW range because of low absorption capacity of the small crystals, as shown in Fig. 2. These changes of sign which occur **~~mostly~~** for thin cirrus clouds **~~(57% of the cases) have already been observed~~**  
**like** in other studies **~~such as~~** (Campbell et al., 2016; Lolli et al., 2017b; Kramer et al., 2020), where the dominant factor in the change of sign of forcing is unclear. Multiple factors are involved, from the optical **~~and radiative~~** properties of the cirrus, to the solar zenith angle or the surface temperature and surface albedo (Wolf et al., 2023).

465

In order to complete this analysis, the net **direct radiative ~~forcing~~ effects** in the full atmosphere are analyzed for cirrus clouds at nighttime and daytime, as shown in Fig. 7.





**Figure 7.** Distribution of net ~~direct~~ radiative ~~forcing~~ effects of cirrus clouds in the full atmosphere at (left) nighttime and (right) daytime, in function of their cloud optical depth (COD). For  $COD < 0.3$  a logarithmic scale has been considered in order to discern more clearly sub-visible and visible cirrus clouds (Sassen and Cho, 1992). The solid line corresponds to the polynomial fitting performed on the data set. The shaded area represents the region with a 95% probability of containing the points, adjusted by the mean absolute value of the differences between the real and fitted values.

Fig. 7 shows a net warming of the atmosphere (always positive ~~forcing~~ direct radiative effect) for cirrus clouds at nighttime and daytime, being radiation escape lower in the presence of cirrus clouds in the full atmosphere. This phenomenon could have been perceived in the previous figure (Fig. 6) as the ~~forcing~~ direct radiative effect at TOA was always higher than at BOA. It also results that the atmosphere warms faster in function of the COD during the daytime than at nighttime (see their regression slopes), as expected due to the contribution of solar radiation to the net ~~forcing~~ direct radiative effect. The net ~~forcing~~ direct radiative effect in the full atmosphere fits very well with the polynomial regressions for both time periods, being at nighttime  $R^2 = 0.990.96$  and at daytime  $R^2 = 0.970.95$ , although some instances outside the shaded area are observed. These strong variations of the direct radiative ~~forcing~~ effect for cirrus clouds with very similar COD are due to the consideration in the simulations of different ~~radiative~~ optical scattering properties of the cirrus clouds, thermodynamic profiles, surface temperature and surface albedo values for each cirrus scene.

Then, the direct radiative ~~forcing~~ effects for cirrus clouds at nighttime and daytime, which are classified according to (Sassen and Cho, 1992) criteria, are quantified, as shown in Table 1.

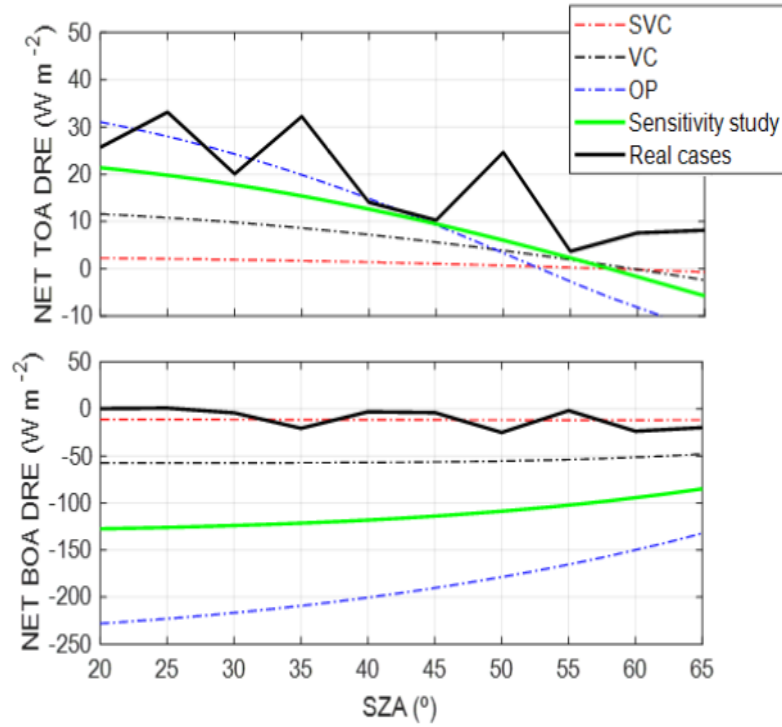
Type	Nighttime				Daytime			
	COD	BOA	TOA	ATM	COD	BOA	TOA	ATM
Sub-visible	0.02±0.01	3.9±3.9	7.5±7.1	3.6±3.3	0.02±0.01	-0.7±1.2	2.6±1.3	3.3±1.8
Visible	0.15±0.08	11.1±5.6	22.2±10.4	11.1±5.2	0.11±0.08	-4.4±7.4	11.3±9.1	15.6±13.1
Opaque	0.76±0.47	42.2±20.3	74.8±31.5	32.6±15.4	0.75±0.39	-28.2±23.8	25.4±17.4	53.6±21.6
Total	0.37±0.43	22.1±20.8	40.4±34.5	18.4±15.3	0.28±0.37	-11.5±18.7	14.2±14.0	25.7±25.9

**Table 1.** Average and standard deviation of **direct** radiative ~~forcing~~ effects of cirrus clouds ( $\text{Wm}^{-2}$ ) at bottom-of-the-atmosphere (BOA), top-of-the-atmosphere (TOA) and in the full atmosphere (ATM), at nighttime and daytime, classified with (Sassen and Cho, 1992) criteria in Barcelona. Cloud optical depth values are obtained from (Gil-Díaz et al., 2024).

In Table 1 it is discernible that thicker cirrus clouds produce a higher ~~forcing~~ **direct radiative effect** than thinner clouds, as observed above. At nighttime, cirrus clouds produce an average net warming in the full atmosphere of ~~427.04±22.89~~  **$+40.4\pm34.5$**   $\text{Wm}^{-2}$ , with opaque cirrus clouds being the main source. At daytime, cirrus clouds **produce** generally ~~about~~ **a negative direct radiative effect** at BOA and ~~with the~~ **a positive effect** at TOA, resulting in a warming of the full atmosphere. The **direct** radiative ~~forcing~~ effect at BOA ranges between ~~-58.44/-44~~  **$-73$  and  $+3$**   $\text{Wm}^{-2}$  for all cirrus. In particular, for thin cirrus clouds the **direct** radiative ~~forcing~~ effect is in a range from ~~-20.31~~ to  $+3$   $\text{Wm}^{-2}$ , being a similar range to (Lee et al., 2009), covering from -20 to 0  $\text{Wm}^{-2}$ . Therefore, shortwave negative ~~forcing~~ **direct radiative effect** generally dominates at the BOA, with an average of ~~7.18±6.48~~  **$-3.6\pm6.7$**   $\text{Wm}^{-2}$  for thin cirrus clouds, being slightly lower than  $-1.35$   $\text{Wm}^{-2}$  (Lee et al., 2009). The **direct** radiative ~~forcing~~ effect at TOA ranges between  $+1$  and ~~67~~  **$+54$**   $\text{Wm}^{-2}$  for all cirrus, being a wider interval than (Campbell et al., 2016; Lolli et al., 2017b; Kramer et al., 2020; Kienast-Sjögren et al., 2016). In particular, for thin cirrus clouds the **direct** radiative ~~forcing~~ effect is in a range from  $+1$  to ~~36~~  **$+32$**   $\text{Wm}^{-2}$ , being the maximum value considerably higher than the value of  $+5.71$   $\text{Wm}^{-2}$  (Campbell et al., 2016) or  $+10$   $\text{Wm}^{-2}$  (Kramer et al., 2020). The average of the **direct** radiative ~~forcing~~ effect for thin cirrus clouds is ~~4.19/9.0±10.45~~  **$+9.4\pm8.8$**   $\text{Wm}^{-2}$ , being a close value compared to them obtained in (Kienast-Sjögren et al., 2016), that cover in average from 6.2 to 11  $\text{Wm}^{-2}$ , although they are also significantly higher than the value of  $+1$   $\text{Wm}^{-2}$  (Lee et al., 2009). Despite the differences found, the values are in agreement in magnitude with other studies such as (Ackerman et al., 1988; Jensen et al., 1994; Lee et al., 2009; Berry and Mace, 2014; Campbell et al., 2016; Kienast-Sjögren et al., 2016; Kramer et al., 2020). The average net warming in the full atmosphere is a little higher at daytime, with an average value of ~~427.51±26.63~~  **$+25.7\pm25.9$**   $\text{Wm}^{-2}$ . This difference is apparently not related with the fraction of opaque cirrus, as the percentage of opaque cirrus is lower during nighttime (28%) than at daytime (38%).

#### 500 4.5 **Direct radiative ~~forcing~~ effects of cirrus clouds depending on solar zenith angle**

In this section, only net **direct** radiative ~~forcing~~ effects of cirrus clouds at daytime, will be considered. The distinction between the shortwave and longwave spectral ranges will not be made because **direct** radiative ~~forcing~~ effect does not depend on the solar zenith angle in the longwave spectrum (Lee et al., 2009; Wolf et al., 2023). The net **direct** radiative ~~forcing~~ effect of cirrus clouds during daytime at the BOA and TOA, together with the results of a brief sensitivity study, in which all parameters of the simulations except the solar zenith angle are kept constant, are shown in Fig. 8.



**Figure 8.** Distribution of the net ~~direct~~ radiative ~~forcing~~ effects of cirrus clouds at daytime at (top) top-of-the-atmosphere (TOA) and (bottom) bottom-of-the-atmosphere (BOA) in function of their solar zenith angle (SZA) resulting from the (dashed curves correspond to for each cirrus cloud type according to the (Sassen and Cho, 1992) criteria and green curve correspond to the mean values for all daytime cirrus clouds) sensitivity study and (black curve) the real cases.

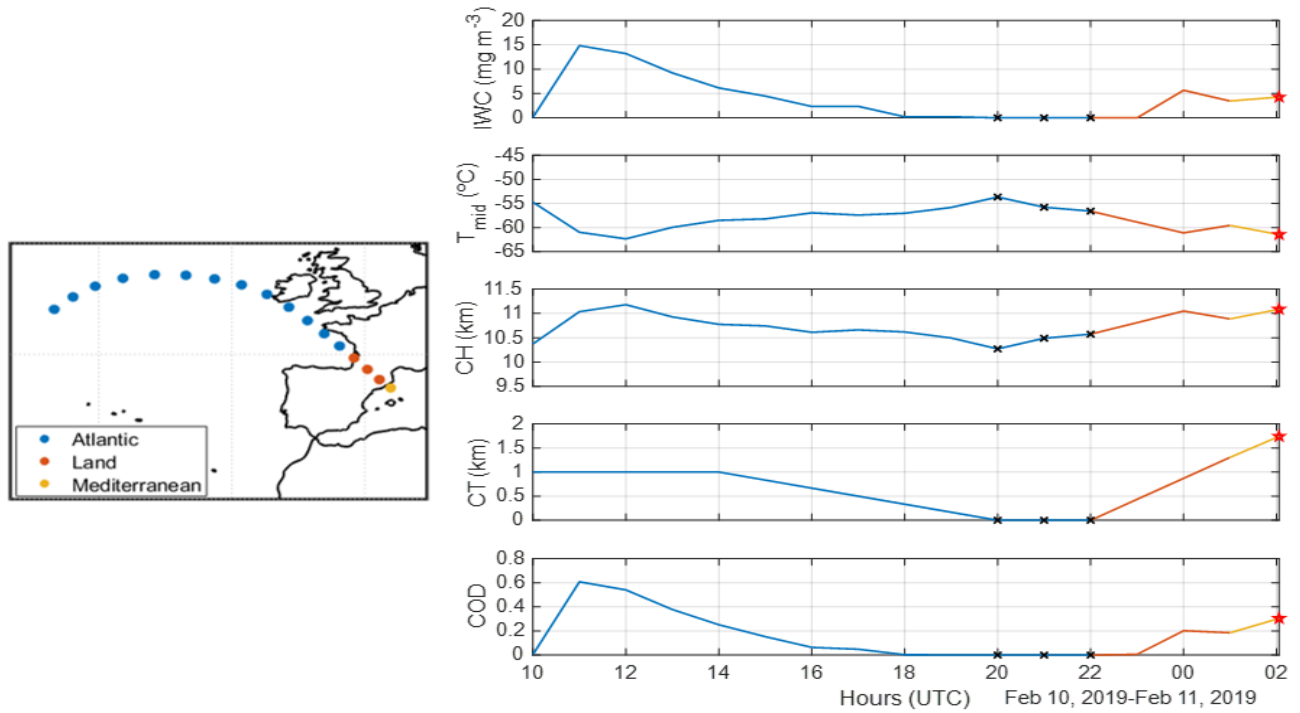
In the sensitivity study, three cloud types have been considered, according to the (Sassen and Cho, 1992) criteria, where the cloud optical depth is shown in Table 1 for the daytime cirrus clouds. The geometrical, thermal and optical properties correspond to the mean values of the cirrus clouds measured at Barcelona lidar station (Gil-Díaz et al., 2024) and the ~~radiative~~ **optical scattering** properties are obtained from the mean values resulting from the statistic. The longwave surface albedo considered is 0.0157 and the surface temperature is 28°C. The thermodynamic profiles have been selected from the 1976 standard atmosphere (COESA et al., 1976). By keeping constant all other properties of cirrus clouds scenes, the results are expected to be exclusively due to the variation in the value of the solar zenith angle. On the other hand, in the real cases that have been analyzed in Section 4.4 and Section 4.5 as a function of cloud optical depth, the net ~~direct~~ radiative ~~forcing~~ **effect** values for the three cloud types have been averaged, with a SZA resolution of 5°.

In Fig. 8 one observes a higher variability of net ~~direct~~ radiative ~~forcing~~ **effect** for the real cases than for the sensitivity study. This variability could be explained by other parameters that are considered in the simulations, such as cloud optical depth, cirrus ~~radiative~~ **optical scattering** properties, surface albedo and temperature (Sicard et al., 2014; Wolf et al., 2023). **While**

a detailed sensitivity analysis is needed to assess the contribution of each variable, cloud optical depth appears to have the most significant impact, as the difference in mean CODs for each averaged SZA value is substantial. In addition, a generally higher mean net **direct** radiative ~~forcing~~ **effect** is discerned for the real cases than for the sensitivity study results, especially at BOA despite having taken the average properties resulting from the statistics. At TOA, a slight downward tendency of the net **direct** radiative ~~forcing~~ **effect** is obtained as the SZA increases, as found in (Wolf et al., 2023). As SZA increases, cloud solar extinction is enhanced regarding thermal effects (Campbell et al., 2016). All mean net **direct** radiative ~~forcing~~ **effect** values for the real cases are positive, but there is a large fluctuation in certain values of SZA. Moreover, there is no crossover where the mean net **direct** radiative ~~forcing~~ **effect** shifts from positive to negative values between SZA of 20 to 65°. On contrary, for the results from the sensitivity study, a change of sign is observed at 58°, fixing well with results from (Campbell et al., 2016; Lolli et al., 2017b). At BOA, most net **direct** radiative ~~forcing~~ **effect** values are negative, presenting a slight increasing trend unlike the TOA. Since the angle of incidence of the incoming solar radiation increases, the incident solar radiation is lower and the scattering produced by the ice crystals increases because optical path is larger (Lee et al., 2009; Wolf et al., 2023). This enhancement of the extinction of incident solar radiation, which is lower, results in a reduction of the net ~~forcing~~ **positive direct radiative** effect at TOA and the net ~~forcing~~ **negative** effect of the atmosphere at BOA. The change from positive to negative values of the mean net **direct** radiative ~~forcing~~ **effect** is only observed for the real cases, being SZA crossover of 25°, a considerably lower value than that observed for the results of the TOA sensitivity study.

#### 535 4.6 Case study of **direct** radiative ~~forcing~~ **effects** of an evolving cirrus cloud

In this section the role of time is added to the analysis and the **direct** radiative ~~forcing~~ **effect** produced by a cirrus cloud is studied along its back-trajectory. The final objective of this case is to simultaneously analyze the evolution of different physical agents such as the surface albedo, the solar zenith angle, the cloud optical depth or the ice water content in the quantification of the cirrus cloud **direct** radiative ~~forcing~~ **effect**. The case of study corresponds to the back-trajectory of a cirrus cloud measured at Barcelona lidar station on 11/02/2019 at 02:03 UTC (Gil-Díaz et al., 2024), where simultaneous measurements of the MPL and CALIPSO were performed. To simulate the evolution of the cirrus cloud as realistic as possible, its microphysical properties along its back-trajectory are calculated with the CLaMS-Ice model. This model provides apart from the basic back-trajectory variables such as temperature, altitude, geographic coordinates and time, microphysical properties like ice water content, ice crystal number concentration or ice nuclei concentration. Considering the temperature provided by the model as the mid-cloud temperature and together with the ice water content, the ~~radiative~~ **optical scattering** properties of cirrus clouds are specifically calculated with ~~Eq. 4 and Eq. 5~~ **the self-consistent scattering model for cirrus clouds**. Assuming the cirrus cloud geometric thickness decreases linearly, using the CALIPSO measurements of 1 km on 10/02/2019 14 UTC and 1.74 km on 11/02/2019 2 UTC, and considering that the geometrical thickness is 0 km when the IWC is null, the COD is estimated as the product of the extinction coefficient and the geometrical thickness. The back-trajectory of the cirrus cloud and the properties: ice water content, mid-cloud temperature, mid-cloud altitude, cloud geometrical thickness and cloud optical depth are shown in Fig. 9.

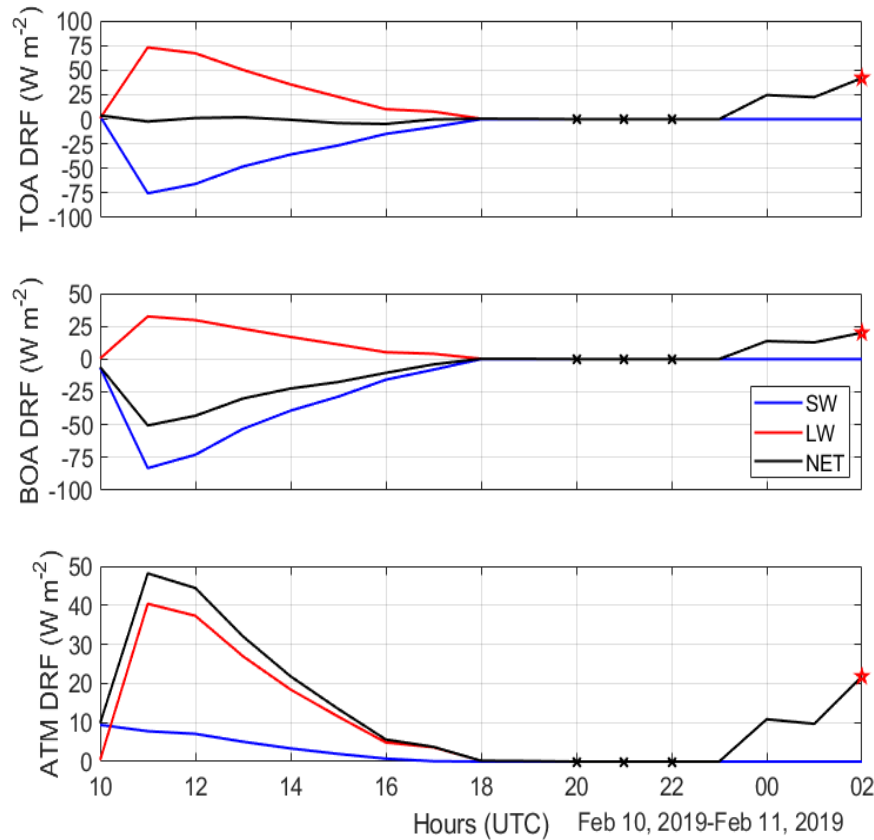


**Figure 9.** (Left) Hourly back-trajectory of the cirrus cloud for the last 16 hours before arrival in Barcelona. (Right) Evolution of cirrus cloud properties along its back-trajectory: ice water content (IWC), mid-cloud temperature ( $T_{\text{mid}}$ ), mid-cloud altitude (CH), cloud geometrical thickness (CT) and cloud optical depth (COD) (from top to bottom). Colours indicate the type of surface where the cloud is over, black crosses mark the non-existence of cirrus cloud during those hours and the red star points to the case measured by the MPL and CALIPSO in (Gil-Díaz et al., 2024).

In Fig. 9 one observes that the cirrus cloud comes from the Atlantic Ocean, passing through part of France and Barcelona to reach the Mediterranean Sea. The selected points of the trajectory are spaced 1 hour apart backward from 11 February 2019 at 02 UTC. During this journey, the air mass which corresponds to the cloud undergoes a rise in height, reaching a minimum in temperature and a COD of 0.6. After this initial cirrus cloud fast growth, the cloud gradually fades away until 20 UTC, where the CLaMS-Ice model gives a null ice water content. Afterwards, a new cirrus begins to form over land surface and to grow until 00 UTC and remains relatively stable.

Taking advantage of the overlap between the back-trajectory of this cirrus cloud and the CALIPSO satellite, a brief evaluation of the ice water content is carried out on 11/02/2019 02:03:20 UTC. The difference in IWC between the CLaMS-Ice output and the CALIPSO measurement is  $0.88 \text{ mg/m}^3$ , resulting in a relative deviation of 29% for the CLaMS-Ice output. This difference obtained by CLaMS-Ice and CALIPSO, well within a factor of two, is considered reasonable given the different parameterizations involved in the calculation of the IWC and their respective uncertainties. To fully characterize the cirrus

cloud scenes along its back-trajectory, it is considered that the height from the CLaMS-Ice model is the mid-cloud height. ~~and the surface albedo is constant over each of the three distinct surface types. The albedo of the Atlantic Ocean and the~~  
 565 ~~Mediterranean Sea are estimated with data from the CERES satellite in the SW and the land albedo is assumed to be that of~~  
~~Barcelona.~~ The surface albedo and surface temperature considered are from the NOAA-20 satellite. With all these assumptions, radiative simulations are calculated with the ARTDECO package and the ~~radiative fluxes~~ direct radiative effects at TOA, BOA and the full atmosphere are shown in Fig. 10.



**Figure 10.** Distribution of ~~direct radiative forcing~~ effects of the cirrus cloud along its back-trajectory at (top) top-of-the-atmosphere (TOA), (centered) bottom-of-the-atmosphere (BOA) and (bottom) in the full atmosphere (ATM). Shortwave (SW), longwave (LW) and net (NET = SW + LW) components of ~~direct radiative forcing~~ effects have been distinguished. Black crosses mark the non-existence of cirrus cloud during those hours.

Fig. 10 shows that at TOA the net ~~direct radiative forcing~~ effect is close to zero with values shifting between positive and  
 570 negative during the first hours where the cirrus cloud is over the Atlantic Ocean. During this period, the SW component is almost completely balanced by the LW component and it is zero when there is no incident solar radiation, i.e. at nighttime,



therefore the net ~~forcing~~ **direct radiative effect** corresponds to the LW component. At nighttime, the net direct **direct** radiative ~~forcing~~ **effect** at TOA is also approximately double than at BOA, in agreement with previous results. At BOA, the net forcing changes from negative to positive, since the incident solar radiation produces that the albedo effect overcomes the greenhouse effect during daytime. In the full atmosphere, the **direct** radiative ~~forcing~~ **effects** are always positive, as the ~~forcing~~ **direct radiative effect** at TOA is higher than at BOA. In summary, these simulations reveals the evolution of the net **direct** radiative ~~forcing~~ **effect** produced by the cirrus cloud, going at TOA from values close to 0 to ~~+12~~**+40**  $\text{Wm}^{-2}$ , at BOA from negative values, whose minimum is  $-51 \text{ Wm}^{-2}$  to positive values reaching a maximum of  $+20 \text{ Wm}^{-2}$  and in the full atmosphere varying between values close to 0 to ~~+18~~**+42**  $\text{Wm}^{-2}$ , being the maximum. The complexity of calculating the **direct** radiative ~~forcing~~ **effect** of a cirrus cloud lies in the fact that this value is highly sensitive to its scene cloud properties like cloud optical depth, solar zenith angle or surface albedo as seen in this case study.

Finally, we compare the **direct** radiative ~~forcing~~ **effect** of the cirrus cloud (red star mark) measured by MPL at the Barcelona lidar station, by CALIPSO satellite at 78 km from Barcelona lidar station (Gil-Díaz et al., 2024) and with the properties obtained with CLaMS-Ice, as shown in Table 2.

Database	Properties				NET DRE ( $\text{Wm}^{-2}$ )		
	$T_{mid}$ ( $^{\circ}\text{C}$ )	$\sigma_{ext}$ ( $\text{Km}^{-1}$ )	IWC ( $\text{mg m}^{-3}$ )	COD	BOA	TOA	ATM
MPL	-60.7	0.17	4.1 <sup>a</sup>	0.26	17.0	35.5	18.5
CALIPSO	-63.7	0.16	22.2	0.23	15.8	32.9	17.0
CLaMS-Ice	-61.4	<i>b</i>	4.3	0.30	19.6	41.1	21.5

**Table 2.** Properties (mid-cloud temperature ( $T_{mid}$ ), column effective extinction coefficient ( $\sigma_{ext}$ ) and ice water content (IWC)) and net **direct** radiative ~~forcing~~ **effect** of the cirrus cloud measured by Micro Pulse Lidar at the Barcelona lidar station, by CALIPSO satellite (Gil-Díaz et al., 2024) and with the properties obtained with CLaMS-Ice, at bottom-of-the-atmosphere (BOA), top-of-the-atmosphere (TOA) and in the full atmosphere (ATM). <sup>a</sup>The value of ice water content is calculated with the self-consistent scattering model for cirrus cloud (see Sec. 3.1). <sup>b</sup> The CLaMS-Ice product does not provide a value of the ~~effective/column~~ extinction of the cirrus cloud.

Taking the new methodology explained above to characterize ~~radiatively~~ **optically** cirrus ~~clouds~~ **scattering properties**, the self-consistent scattering model for cirrus clouds is used to obtain the ~~effective/column~~ extinction coefficient, single scattering  $\pm$ albedo and asymmetry factor to introduce them into the ARTDECO package and calculate its net **direct** radiative ~~forcing~~ **effect**. In the case of cirrus cloud detection with MPL measurements, as described above, the mid-cloud temperature and the ~~effective/column~~ extinction coefficient, obtained by the two-way transmittance method (Gil-Díaz et al., 2024), are employed to characterize the cirrus ~~radiatively~~ **optically**. Alternatively, as the CALIPSO and CLaMS-Ice data provide an IWC measurement for such cirrus cloud, this data together with the mid-cloud temperature is used to represent the cirrus ~~radiatively~~. Despite the fact that CALIPSO passes 78 km from the Barcelona station, in this simulations with these data and those of CLaMS-Ice, an albedo and single temperature corresponding to the Barcelona lidar station are considered, in order to compare radiative

In Table 2 one observes that the properties of the cirrus cloud are similar, except for the IWC, which is considerably higher for the cirrus characterized with CALIPSO measurements. If the IWC is calculated using the extinction coefficient and mid-cloud temperature values provided by CALIPSO, the ensemble scattering model for cirrus clouds yields an IWC value of 3.9  
 600  $\text{mg/m}^3$ . This result is significantly lower than the IWC measured by CALIPSO, since the ~~ensemble~~ self-consistent scattering model for cirrus clouds often underestimates the IWC of cirrus clouds measured at mid-latitude (Baran and Labonnote, 2007). So a certain discordance between the net ~~direct~~ radiative ~~forcing~~ effect magnitudes would be expected. For all three simulations, the cirrus cloud ~~warms~~ has a positive direct radiative effect both at TOA and at BOA, since due to the time (2 UTC) the solar radiation component is null. Moreover, as the ~~forcing~~ direct radiative effect is proportional to the COD, the cirrus characterized  
 605 with the CLaMS-Ice products produces a slightly higher ~~forcing~~ direct radiative effect than with the other data. On average, it can be established that the cirrus cloud measured at Barcelona lidar station produces a ~~warms~~ positive direct radiative effect at BOA of ~~419.19/12/08~~  $+17.5 \pm 1.9 \text{ Wm}^{-2}$ , at TOA of ~~449.07/15/11~~  $+36.5 \pm 4.2 \text{ Wm}^{-2}$  and in the full atmosphere of ~~429.88/15/13~~  $+19.0 \pm 2.3 \text{ Wm}^{-2}$ .

## 5 Conclusions

610 In this paper a study of ~~radiative~~ optical scattering properties and forcings of cirrus clouds based on 4 years of continuous ground-based lidar measurements with the Barcelona (Spain) Micro Pulse Lidar (MPL) is analyzed. First, a new approach of a self-consistent scattering model for cirrus clouds is presented to get the ~~radiative~~ optical scattering properties of cirrus clouds at different wavelengths with only the ~~extinction coefficient~~ extinction coefficient calculated with the two-way transmittance method and ~~mid~~-cloud temperature, from radiosounding data. The self-consistent scattering model for cirrus clouds consists of an  
 615 ensemble of six ice crystal ~~microphysics~~ types, where the simplest ~~ice crystals are represented by hexagonal~~ form is an ice columns with an ~~unity~~ aspect ratio ~~of unity~~ and the more complex ice crystals are formed by ~~attach~~ randomly ~~oriented~~ attaching other hexagonal elements, until a chain-like ~~ice crystal~~ structure is formed. The members ~~of the ensemble~~ are evenly distributed ~~into six equal intervals of the particle size distribution (PSD)~~, across the PSD. The self-consistent scattering model for cirrus cloud database consists of more than 20000 PSDs of tropical and mid-latitude cirrus at temperatures between -60 and 0°C. This  
 620 database provides for each PSD the simulated optical scattering properties ~~which relates the density logarithm of the water content (LIWC) as a function of property~~ (scattering, absorption coefficient and asymmetry factor) in function of the decimal logarithm of ice water content (LIWC) and the cloud temperature (T) by nonlinear least squares fitting. ~~the a polynomial fit to the in-cloud temperature (T), with a spectral dependence~~ The new approach ~~for the calculation~~ consists of first calculating the IWC of the cirrus cloud by introducing the extinction coefficient of each cloud layer in an equation derived from the  
 625 model, for a cloud temperature. ~~The radiative properties of cirrus clouds at different wavelengths consists of first calculating the IWC by the density logarithm by introducing the extinction coefficient in an equation derived from the model, for a mid-cloud temperature~~ This equation is obtained by assuming the absence the absorption, which is entirely reasonable because

the wavelength used belongs to the visible spectral range. Once, the IWC is estimated, this variable is introduced again into the model to get the absorption, scattering and asymmetry factor coefficients for each wavelength, respectively. Applying this method to cirrus clouds measured in Barcelona during November 2018 to September 2022 at 00 and 12 UTC, it is obtained that the average of the ice water content is  $4.97 \pm 5.53 \text{ mg/m}^3$ , the single scattering albedo is  $0.99 \pm 0.05$  and the asymmetry factor is  $0.76 \pm 0.00$  at  $0.55 \text{ } \mu\text{m}$ . Second, the **direct radiative forcing effect** of cirrus clouds is calculated with the radiative transfer model DISORT. Radiative fluxes are validated at bottom-of-the-atmosphere with SolRad-Net pyranometers in the shortwave spectral range, and at top-of-the-atmosphere with ~~OVERES~~ **NOAA-20** measurements in the longwave spectral range. One one hand, most downward radiative fluxes calculated with the DISORT model overestimate the SolRad-Net observations, resulting in a BIAS of +11.6% and a slope of the linear regression ( $a = 0.71 \pm 0.20$ ). On the other hand, a large difference in upward radiative fluxes between simulated and observations is found for each cirrus cloud scene, resulting in a BIAS of ~~45.1/62~~ **+33.6%**. Third, a validation of the importance of the planetary boundary layer aerosols in the cirrus scenes simulations is carried out. Calculations with and without aerosols of the cirrus **direct radiative forcing effects** are made to assess the error induced by neglecting tropospheric aerosols, which results in a negligible BIAS of ~~-0.7/1.2~~ **1.2%**. In the other case where the aerosols are vertically closer to the clouds, the simplification of not considering aerosols in the calculation of cloud ~~forwards~~ **direct radiative effects** may not be valid, leading to a significant underestimation of cloud ~~forwards~~ **direct radiative effects**. Forth, the **direct radiative forcing effects** of cirrus clouds are calculated distinguishing between nighttime and daytime. At nighttime, cirrus clouds warm the atmosphere with **direct radiative effects** at TOA almost double than at BOA, with the thicker cirrus clouds contributing most to the ~~forwards~~ **direct radiative effects**. At daytime, cirrus clouds **have generally a negative direct radiative effect at BOA** (80.82% of the cases) and always ~~warm~~ **a positive effect at TOA**, resulting in a warming of the full atmosphere. On average, at nighttime, cirrus clouds ~~warm~~ **have a positive direct radiative effect of  $+23.02 \pm 21.23$   $\text{Wm}^{-2}$  at BOA,  $+50.06 \pm 42.99$   $\text{Wm}^{-2}$  at TOA and in the full atmosphere  $+27.04 \pm 22.39$   $\text{Wm}^{-2}$ ; at daytime, cirrus clouds ~~cool~~ **have a negative direct radiative effect  $-18.57 \pm 14.95$   $\text{Wm}^{-2}$  at BOA and  $-18.94 \pm 16.95$   $\text{Wm}^{-2}$  at TOA, and in the full atmosphere  $-27.51 \pm 26.63$   $\text{Wm}^{-2}$** . Fifth, the variation of the cirrus cloud **direct radiative forcing effect** at daytime is also analyzed as a function of the SZA: it shows that at TOA for the real cases the average net **direct radiative forcing effect** is always positive and for the results from a sensitivity study, the mean net **direct radiative forcing effect** shifts from positive to negative values at  $58^\circ$ . ~~For all cases~~ **For the real cases and the sensitivity study**, a slight downward tendency of the net **direct radiative forcing effect** is also found. At BOA, most net **direct radiative forcing effect** values are negative. The change from positive to negative values of the mean net **direct radiative forcing effect** is only observed for the real cases, being SZA crossover of  $25^\circ$ . Sixth, for a case study, the **direct radiative forcing effect** of a cirrus cloud along its back-trajectory is analyzed using CLaMS-Ice products. During the overlap between the back-trajectory of this cirrus cloud and the CALIPSO satellite on 11/02/2019 02:03:20 UTC, a brief validation of the IWC is made, resulting in a relative ~~error~~ **deviation** of 29% for the CLaMS-Ice output. This cirrus cloud comes from the Atlantic Ocean, passing through part of France and Barcelona to reach the Mediterranean Sea. Over the Atlantic Ocean, the air mass which corresponds to the cloud undergoes a rise in height, reaching a minimum in temperature and a COD of 0.6. After this initial cirrus cloud fast growth, the cloud gradually fades away as it approaches France, where the CLaMS-Ice model gives a null IWC. Afterwards, the cirrus cloud**

begins to form again on land surface and to grow up. Along its trajectory, the cirrus cloud produces a net **direct** radiative ~~forcing~~ **effect** that goes at TOA from values close to 0 to ~~-42~~ **+40**  $\text{Wm}^{-2}$ , at BOA from negative values, whose minimum is -51  $\text{Wm}^{-2}$  to positive values reaching a maximum of +20  $\text{Wm}^{-2}$  and in the full atmosphere varying between values close to 0 to ~~-48~~ **+42**  $\text{Wm}^{-2}$ , being the maximum. Finally, the **direct** radiative ~~forcing~~ **effects** that the cirrus cloud has at the beginning of the back-trajectory are compared using different measurements (MPL and CALIPSO measurements and CLaMS-Ice outputs) and making use of the self-consistent scattering model for cirrus clouds. It results in an average ~~forcing~~ **positive direct radiative effect** at BOA of ~~-19.19±2.08~~ **+17.5±1.9**  $\text{Wm}^{-2}$ , at TOA of ~~-49.07±5.71~~ **+36.5±4.2**  $\text{Wm}^{-2}$  and in the full atmosphere of ~~-29.88±3.63~~ **+19.0±2.3**  $\text{Wm}^{-2}$ .

*Data availability.* The MPLNET products are publicly available on the MPLNET website ([https://mplnet.gsfc.nasa.gov/download\\_tool/](https://mplnet.gsfc.nasa.gov/download_tool/)) (MPLNET, 2024) in accordance with the data policy statement. Radiosoundings data are available upon request from the authors or Meteocat. The SolRad-Net product is publicly available on the SolRad-Net website (<https://solrad-net.gsfc.nasa.gov/>) (Goddard Space Flight Center, 2024a) in accordance with the data policy statement. The NOAA-20 products are publicly available on the CERES website (<https://ceres.larc.nasa.gov/>) (Langley Research Center, 2024). The AERONET products are provided by a federation of ground-based remote sensing aerosol networks established by NASA and PHOTONS (PHOTométrie pour le Traitement Opérationnel de Normalisation Satellitaire) and is greatly expanded by collaborators from national agencies, institutes, universities, individual scientists, and partners. The AERONET products are publicly available on the AERONET website (<https://aeronet.gsfc.nasa.gov/>) (Goddard Space Flight Center, 2024b). The CALIPSO product is provided by the NASA Langley Research Center's (LaRC) ASDC DAAC and is managed by the NASA Earth Science Data and Information System (ESDIS) project. NASA data are freely accessible and available on the Atmospheric Science Data Center website (<https://asdc.larc.nasa.gov/>) (NASA, 2024).

*Author contributions.* CGD prepared the automatic algorithms to ~~get~~ **calculate** the ~~radiative~~ **optical scattering** properties and ~~forcing~~ **direct radiative effects** of cirrus clouds for MPL and radiosounding data. CGD prepared the figures of the paper. MS, OS, AS, CMP, AC, ARG and DCFSO reviewed different parts of the results. CGD prepared the paper, with contributions from all co-authors.

*Competing interests.* At least one of the (co-)authors is a member of the editorial board of Atmospheric Chemistry and Physics.

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