

Study of ~~radiative~~ optical scattering properties and direct radiative effects of high-altitude cirrus clouds in Barcelona, Spain with 4 years of lidar measurements

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Abstract. Cloud-radiation interaction still drives large uncertainties in climate models and its estimation is key to make more accurate predictions. In this context, the high-altitude cirrus clouds play a fundamental role, because 1) they have a high occurrence frequency globally and 2) they are the only cloud that can readily cool or warm the atmosphere during daytime, depending on their properties. This study presents a comprehensive analysis of ~~radiative~~ optical scattering properties and ~~direct radiative effects~~ of cirrus clouds based on 4 years of continuous ground-based lidar measurements with the Barcelona (Spain) Micro Pulse Lidar. First, we introduce a novel approach of a self-consistent scattering model for cirrus clouds to determine their ~~radiative~~ optical scattering properties at different wavelengths using only the ~~direct~~ extinction coefficient and ~~cloud~~ temperature. Second, we calculate the ~~radiative~~ direct radiative effects of cirrus clouds with the Discrete Ordinates Method and we validate our results with SolRad-Net pyranometers and ~~NOAA-20~~ measurements. Third, we present a case study analyzing the ~~radiative~~ direct radiative effect of a cirrus cloud along its back-trajectory using data from the Chemical LAgrangian Model of the Stratosphere with microphysics scheme for Ice clouds formation. The results show that the cirrus clouds with an average ice water content of $4.97 \pm 5.53 \text{ mg/m}^3$, at nighttime, ~~which~~ ~~have a positive direct radiative effect~~ at top-of-the-atmosphere (TOA; ~~+50.1~~ $+40.4 \text{ W m}^{-2}$) almost twice than at bottom-of-the-atmosphere (BOA; ~~+28.0~~ $+22.1 \text{ W m}^{-2}$); at daytime, they ~~have generally~~ ~~a negative direct radiative effect at BOA~~ (~~-8.57~~ -11.5 W m^{-2} , ~~80~~ 82% of the cases) and always ~~which~~ ~~a positive effect at TOA~~ (~~+18.9~~ $+14.18 \text{ W m}^{-2}$). In these simulations, the influence of the lower layer aerosols is negligible in the cirrus ~~radiative~~ direct radiative effects, with a BIAS of ~~-0.7%~~ -1.2% . For the case study, the net radiative effects produced by the cirrus cloud, going at TOA from 0 to ~~+42~~ $+40 \text{ W m}^{-2}$ and at BOA from -51 to $+20 \text{ W m}^{-2}$. This study reveals that the complexity of the cirrus cloud ~~radiative~~ direct radiative effect calculation lies in the fact that it is highly sensitive to the cirrus scene properties.

20 1 Introduction

Cloud-radiation interaction still drives large uncertainties in weather and climate models (IPCC, 2023). Its estimation is very important in order to understand the main physical processes driving climate change, to predict long-term global warming and to make more accurate weather predictions. (Loeb et al., 2009) estimated globally at top-of-the-atmosphere an annual cloud shortwave radiative effect of approximately -50 Wm^{-2} and longwave effect of approximately $+30 \text{ Wm}^{-2}$. The resulting net global mean cloud radiative effect of approximately -20 Wm^{-2} implies a net cooling effect of clouds on the current climate. Owing to the large magnitudes of the cloud radiative effects, clouds cause a significant climate feedback that depends on cloud properties and their spatial distribution (IPCC, 2023). In this context, the high-altitude cirrus clouds play a fundamental role in the global radiation budget (Liou, 1986; Lolli et al., 2017b), having been designated as poorly understood by (IPCC, 2023) because of a lack of knowledge of their dynamic, microphysical and ~~radiative~~ optical scattering properties. Indeed, cirrus cloud critical role in the climate comes from the fact that 1) they have a high occurrence frequency globally (Holz et al., 2008) and 2) they are the only cloud that can readily cool or warm the top-of-the-atmosphere and bottom-of-the-atmosphere, during daytime, depending on their properties (Campbell et al., 2016). In fact, (Campbell et al., 2016) demonstrated through a one-year long lidar dataset that positive or negative daytime cirrus cloud forcing could occur depending on the cloud optical depth (COD) and the solar zenith angle (SZA).

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Cirrus clouds are mainly composed of ice crystals and can form through different atmospheric mechanisms that determine their initial properties and further evolution. In European field campaigns it has been observed that during a low or high pressure system, cirrus clouds are typically formed by a slow updraft, while in conjunction with jet streams or gravity waves, cirrus clouds originate as a consequence of a fast updraft. Also, liquid origin cirrus mostly related to warm conveyor belts are found (Kramer et al., 2016). The most common parameters that are measured in cirrus clouds are temperature, relative humidity (for ice), vertical velocity, ice water content, ice crystal number, ice nucleation particles and ice crystal size distribution. Unfortunately, the measurements of ice crystal number and size as well as relative humidity have faced instrumental problems during last decades (Korolev et al., 2011; Kramer et al., 2016). Moreover, it is a difficult task to draw conclusions about the microphysical processes of cirrus clouds from these observations. Nevertheless, worldwide studies on cloud and aerosol optical and microphysical properties have increased significantly over the last years through the passive ground-based measurements made e.g. by the European Aerosol Research Lidar NETwork, EARLINET (Pappalardo et al., 2014) now included in the Aerosols, Clouds and Trace gases Research Infrastructure, ACTRIS (Saponaro et al., 2019), Micro Pulse Lidar NETwork, MPLNET (Welton et al., 2001); and satellite measurements e.g. by Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observations, CALIPSO (Winker et al., 2007), AEOLUS (Ingmann and Straume, 2016), MODerate-resolution Imaging System, MODIS (Levy et al., 2013) and in the future, Earth Cloud, Aerosol and Radiation Explorer, EarthCARE (Eisinger et al., 2017). Additionally, in-situ airborne measurement campaigns have been carried out such as the First ISCCP Project Regional Experiment, FIRE from 1989 to 1995 (Ackerman et al., 1990; Heymsfield et al., 1990; Heymsfield and Miloshevich, 1995), the International Cirrus Experiment, ICE campaign in 1989 (Raschke et al., 1987), European Cloud Radiation EXperiment, EUCREX in 1993

and 1994 (Sauvage et al., 1999), Field Radiation Experiment on Natural Cirrus and High-level clouds, FRENCH in 2001 (Brogneix et al., 2004), Tropical Composition, Cloud and Climate Coupling, TC4 campaign in 2007 (King et al., 2010; Toon et al., 2010) and CIRRUS-HL campaign in 2021, which is the follow-up to the CIRRUS-ML campaign in 2017 (Voigt et al., 2017; De La Torre Castro et al., 2023).

Up to the present, there are three possibilities for characterising cirrus clouds. One option is the use of in-situ airborne measurements. A second option is to work with a microphysical cirrus cloud model like the Chemical LAgrangian Model of the Stratosphere with microphysics scheme for Ice clouds formation (CLaMS-Ice) (Spichtinger and Gierens, 2009) or Model for aerosol and ice dynamics (MAID) (Bunz et al., 2008; Rolf et al., 2012), that simulate the cirrus cloud development based on the cirrus bulk model by along back-trajectories. The main advantage of this choice is that there is no need to have in-situ airborne lidar measurements. A third option is to employ lidar measurements for the characterisation of cirrus clouds. For that purpose, it is necessary to use a method such as the two-way transmittance method to characterize cirrus clouds optically (Gil-Díaz et al., 2024) together with a scattering model to obtain ~~radiative~~ optical scattering retrievals. For example, (Baran and Labonnote, 2007; Baran et al, 2009, 2011a, b, 2014) relates the cirrus ice water content (IWC) and mid-cloud temperature with its extinction coefficient, single scattering albedo (SSA) and asymmetry factor (asyF). Alternatively, (Heymsfield et al., 2014; Dolinar et al., 2022) propose to calculate the cirrus ice water content from the extinction coefficient at a visible wavelength and the effective geometric diameter of the ice crystals, which in turn is a function of temperature. Once the cirrus ice content and the effective geometric diameter of ice crystals are obtained, the scattering and absorption coefficients and the asymmetry factor can be calculated with the (Fu et al., 1998, 1999; Lolli et al., 2017a) parametrizations.

The objective of this paper is to analyze the ~~radiative~~ optical scattering properties and ~~radiative~~ direct radiative effects of cirrus clouds based on 4 years of continuous ground-based lidar measurements obtained from the NASA Micropulse lidar network (MPLNET, <https://mplnet.gsfc.nasa.gov/>) in Barcelona. Specifically, the ~~radiative~~ optical scattering properties of cirrus clouds have been calculated with a new approach of the self-consistent scattering model for cirrus clouds (Baran and Labonnote, 2007; Baran et al, 2014; Vidot et al., 2015), using only lidar measurements and radiosounding data and their ~~radiative~~ direct radiative ~~effects~~ have been calculated with the ARTDECO package. The instrumentation used is presented in Section 2. A new approach of the self-consistent scattering model for cirrus clouds, the radiative transfer model DISORT and the CLaMS-Ice model are presented in Section 3.1, Section 3.2 and Section 3.3, respectively. The results obtained in this paper are shown in Section 4 and conclusions are presented in Section 5.

2 Instrumentation

The radiative characterization of cirrus clouds relies on the results obtained from (Gil-Díaz et al., 2024) and the instrumentation detailed below.

2.1 NASA Micro-Pulse Lidar Network

A more detailed description of this instrumentation can be found in (Gil-Díaz et al., 2024). In this study, we use the Aerosol (AER) product, ~~provided by~~ which provides 1-min temporal resolution and at 75m vertical resolution variables like aerosol extinction, backscatter and the column lidar ratio (Welton et al., 2000, 2002, 2018; Lolli et al., 2019). This product is used to characterize the aerosol layer which is closest to the surface. The MPLNET AER product ~~includes~~ integrates solar and lunar photometer ~~observations to invert~~ measurements, allowing the lidar signal to be inverted to obtain aerosol properties ~~during~~ over a 24-hour period. ~~Aerosol extinction, backscatter, column lidar ratio, and other properties~~ ~~provided by~~ ~~MPLNET/AER product (Welton et al., 2000, 2002)~~.

2.2 Meteorological Service of Catalonia

95 The Meteorological Service of Catalonia (Meteocat) releases radiosondes are launched twice ~~every~~ a day (at 00:00 and 12:00 UTC) ~~by the Meteorological Service of Catalonia (Meteocat)~~ at a ~~different~~ ~~of~~ location less than 1 km from the MPL site. The radiosondes provide ~~measurements~~ data of pressure, altitude, temperature, relative humidity, wind speed and direction. Only altitude, pressure and temperature profiles are used in ~~the present work~~ this study.

2.3 Solar Radiation Network

100 The Solar Radiation Network (SolRad-Net, <https://solrad-net.gsfc.nasa.gov/>) is a ~~global~~ network of ground-based sensors
105 ~~providing~~ that provides high-frequency solar flux measurements ~~available~~ to the scientific community in near-real time.
This program ~~uses~~ operates in conjunction with AERONET ~~and~~, being its instrumentation ~~is~~
collocated in the AERONET sites. Each SolRad-Net site is initially equipped with two flux sensors: a Kipp and Zonen CM-21
pyranometer (0.305-2.8 μm) ~~for measuring the full~~ to measure the full solar spectrum and a Skye Instrument SKE-510 PAR
110 (photosynthetically/active radiation) Energy sensor (~~spectral range~~ 0.4-0.7 μm).

In this study, ~~the data from the Kipp & Zonen instrument is employed with Level 1.0, corresponding to the unscreened and without final calibration~~ is used for ~~the final calibration period from 2018 to 2022~~. The Kipp & Zonen CM-21 units are ISO 9060 Secondary Standard thermopile pyranometers, ~~equipped with~~ a receiving element ~~enclosed by~~ two concentric Schott K5 glass domes. ~~More information~~ Further details on the instrument can be found at the following web link (<https://www.kippzonen.com/Product/14/CMP21-Pyranometer>).

2.4 Clouds and the Earth's Radiant Energy System project

The Clouds and the Earth's Radiant Energy System project (CERES, <https://ceres.larc.nasa.gov/>) provides ~~observations~~ **measurements** of Earth's radiation budget using measurements from CERES instruments onboard the Terra, Aqua and Suomi National Polar-orbiting Partnership (S-NPP) and NOAA-20 ~~satellites~~ **satellites** (Loeb et al., 2016). The ~~goal~~ **primary objectives** of the CERES project are: (1) to ~~provide~~ **create** a long-

term, and integrated global climate data record that can detect decadal changes in the Earth's radiation budget from the surface to the top of the atmosphere; (2) to enhance the understanding of the temporal and spatial variability in Earth's radiation budget and the role that of clouds and other atmospheric properties play; (3) to support climate model evaluation and improvement through model-observation comparisons.

125 *Observation* In this study, we use the CERES instantaneous Single Scanner Footprint (SSF) product at Level 2/Specificaly, the photodetector is the observed to analyze observations of the upward fluxes at top-of-the-atmosphere , the surface emissivity in longwave spectral range (5-100 μ m), the surface albedo in shortwave spectral (0.2-5 μ m), the surface temperature and the clear-sky percentage for the period from 2018 to 2022. *Surface albedo* is calculated as one minus the surface emissivity. The CERES data comes from the AQUA satellite/Clouds/Brackets/Brackets/12/00/13/00/UTC/ from the TERRA satellite/Brackets/Brackets/10/00 and 10/30/UTC). NOAA-20 satellite (overpassing Barcelona between 11:00 and 13:00 UTC and between 00:00 and 02:00 UTC). This dataset is provided with a surface spatial resolution of 20 km at nadir (Su et al., 130 2015).

2.5 NASA AErosol RObotic NETwork

The NASA AErosol RObotic NETwork (AERONET, <https://aeronet.gsfc.nasa.gov/>) is a ~~fixed~~ ~~stationary~~ ~~global~~ ~~network~~ ~~of~~ ground-based sun/lunar-photometers established by NASA and LOA-PHOTONS (CNRS). For ~~more~~ ~~than~~ ~~over~~ two decades, ~~the~~ ~~project~~ **AERONET** has provided long-term, continuous, ~~and~~ ~~highly~~ ~~accessible~~ ~~public~~ ~~and~~ ~~free~~ ~~publicly~~ ~~accessible~~ database of aerosol optical, ~~and~~ ~~microphysical~~ ~~and~~ ~~properties~~ ~~for~~ ~~aerosol~~ ~~characterization~~ ~~and~~ ~~satellite~~ ~~retrieval~~ ~~validation~~ ~~of~~ ~~systems~~ ~~and~~ ~~models~~, and synergism with other databases. The network ~~imposes~~ ~~enforces~~ strict standardization of instruments, calibration, processing and distribution.

~~In this Work~~ For this study, Version 3.0, and Level 1.5 (cloud-screened and quality-controlled) inversion products are used for the ~~time~~ period from 2018 to 2022. ~~In order to characterize the aerosols in the lowest tropospheric layer of the troposphere~~ In the shortwave spectrum (SW, 0.2-4 μm), variables like Aerosol Optical Depth (AOD), Single Scattering Albedo (SSA) and Asymmetry Factor (b/a) are used. For ~~other~~ wavelengths outside ~~that the Working wavelength of~~ AERONET's working range/ in the SW ~~range~~ spectrum, the Angström exponent is used (Wagner and Silva, 2008). For example, the Angström exponent is calculated ~~with the~~ using AOD values at 440 and 675 nm to ~~obtain~~ estimate the AOD at 550 nm. ~~With the characteristics of the aerosol~~ In the longwave spectrum (LW, 2-40 μm), the aerosol optical proper-

ties *mentioned above* are extracted from the Laboratory for Information Technologies and Mathematical Simulation (LITMS) database (Rublev et al., 1994).

2.6 The Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observation

Further details on this instrumentation are provided in (Gil-Díaz et al., 2024). In order to validate the ice water content from the Lagrangian microphysical cirrus model CLaMS-Ice (see Section 3.3), we use the CALIPSO product "5 km Cloud Layer (05kmClay)" product at Level 2 (L2) and Version 4.20 (V4.20), available from June 2006. This product offers a horizontal averaging resolution of 5 km, with up to ten layers reported per profile. This product contains geometrical, thermal and optical properties of each detected cloud layer like layer top/base altitude and temperature, integrated attenuated backscattering coefficient at 532 nm and 1064 nm, integrated particle depolarization ratio and ice water path. The ice water content is estimated by dividing the ice water path by the geometrical thickness of the cloud layer.

3 Methodology

The ~~thick~~ optical scattering properties of cirrus scenes are determined through the use of a self-consistent scattering model for cirrus clouds (Baran and Labonne, 2007; Baran et al, 2014; Vidot et al., 2015) and their ~~direct~~ radiative ~~feedbacks~~ effects are calculated with the ARTDECO package, which implements a variety of optical properties into state of the art radiative transfer models (see below).

3.1 The self-consistent scattering model for cirrus clouds

The self-consistent scattering model for cirrus clouds consists of an ensemble of six ice crystal ~~members~~ types. The simplest form is a ~~ideal~~ hexagonal platelet with hexagonal ice columns with an unity aspect ratio, while ~~of~~ the more complex ice crystals are formed by ~~attaching~~ randomly attaching ~~other~~ additional hexagonal elements, until creating a chain-like structure ~~ideal~~. The ~~development~~ of ~~ideal~~ crystals tends to ~~be~~ with increasing size. As ice crystals grow in size, their complexity increases, and they generally become more spatially ~~extended~~, with the hexagonal components ~~lengthening~~ elongating over time (Heymsfield and Miloshevich, 2003). The ~~ensemble~~ ensemble consists of six distinct members, starting with ~~the~~ ~~smallest~~ hexagonal ice column, ~~the~~ ~~smallest~~ and followed by a six-branched bullet rosette. Subsequent crystals are formed by attaching hexagonal monomers ~~attaching~~ together in function of their maximum dimension, resulting in aggregates containing three to ten hexagonal elements. This ensemble ~~is~~ designed to represent the various ice crystal types observed in cirrus clouds ~~observed~~ during different measurement campaigns (Heymsfield and Miloshevich, 2003; Lawson et al., 2003; Connolly et al., 2005). For instance, bullet-rosettes are included ~~in~~ they are frequently observed in mid-latitude and

Arctic regions (Lawson et al., 2006; Schmitt et al., 2006) ~~and its~~ with their geometry is described in (Macke et al., 1996). The Each ensemble member ~~is~~ is constructed ~~so that it~~ to avoid intersecting planes, and the crystals are ~~arranged~~ arranged so that multiple reflections between them are negligible, which was ~~verified~~ verified experimentally using ray-tracing calculations. The first ~~member~~ member represents the smaller ~~size~~ of ice crystals in the particle size distribution (PSD), whilst the hexagonal ice aggregates represent the process of ice crystal aggregation ~~thus~~, corresponding to the larger ~~size~~ crystals in the PSD. The six members ~~of the ensemble~~ are evenly distributed ~~in~~ across the PSD. The PSD ~~assumed~~ follows the moment estimation parametrization of the PSD (Field et al., 2003, 2007). The self-consistent scattering model for cirrus cloud database consists of more than 20000 PSDs of tropical and mid-latitude cirrus at temperatures between -60 and 0°C. This database provides for each PSD the simulated optical scattering properties (scattering, absorption coefficient and asymmetry factor) in function of the decimal logarithm of ice water content (LIWC) and the cloud temperature (T) by nonlinear least squares fitting (Vidot et al., 2015). This parametrization excludes ice crystals smaller than 100 μm due to shattering issues on closed-path instruments (Korolev et al., 2011), assuming instead an exponential PSD fit. For crystals larger than 100 μm , measured PSDs were filtered to reduce the likelihood of including shattered crystal artifacts.

~~to determine the properties of the individual ice crystals (T), with a spectral dependence of the extinction coefficient.~~

~~1980-1990 (A1) 1990-2010 (A2) 1990-2010 (A3) 1990-2010 (A4) 1990-2010 (A5) 1990-2010 (A6)~~

$$\text{LIWC}(T) = \text{LIWC}(A1) + (A2 - A1) \frac{T - 1990}{2010 - 1990} + (A3 - A2) \frac{T - 1990}{2010 - 1990} + (A4 - A3) \frac{T - 1990}{2010 - 1990} + (A5 - A4) \frac{T - 1990}{2010 - 1990} + (A6 - A5) \frac{T - 1990}{2010 - 1990} \quad (1)$$

200 ~~2000-2010 (A7) 2010-2020 (A8) 2010-2020 (A9) 2010-2020 (A10) 2010-2020 (A11) 2010-2020 (A12)~~ $\text{LIWC}(T) = \text{LIWC}(A1) + (A2 - A1) \frac{T - 1990}{2010 - 1990} + (A3 - A2) \frac{T - 1990}{2010 - 1990} + (A4 - A3) \frac{T - 1990}{2010 - 1990} + (A5 - A4) \frac{T - 1990}{2010 - 1990} + (A6 - A5) \frac{T - 1990}{2010 - 1990} + (A7 - A6) \frac{T - 1990}{2020 - 1990} + (A8 - A7) \frac{T - 1990}{2020 - 1990} + (A9 - A8) \frac{T - 1990}{2020 - 1990} + (A10 - A9) \frac{T - 1990}{2020 - 1990} + (A11 - A10) \frac{T - 1990}{2020 - 1990} + (A12 - A11) \frac{T - 1990}{2020 - 1990} \quad (2)$

~~205 2000-2010 (A7) 2010-2020 (A8) 2010-2020 (A9) 2010-2020 (A10) 2010-2020 (A11) 2010-2020 (A12)~~ $\text{LIWC}(T) = \text{LIWC}(A1) + (A2 - A1) \frac{T - 1990}{2010 - 1990} + (A3 - A2) \frac{T - 1990}{2010 - 1990} + (A4 - A3) \frac{T - 1990}{2010 - 1990} + (A5 - A4) \frac{T - 1990}{2010 - 1990} + (A6 - A5) \frac{T - 1990}{2010 - 1990} + (A7 - A6) \frac{T - 1990}{2020 - 1990} + (A8 - A7) \frac{T - 1990}{2020 - 1990} + (A9 - A8) \frac{T - 1990}{2020 - 1990} + (A10 - A9) \frac{T - 1990}{2020 - 1990} + (A11 - A10) \frac{T - 1990}{2020 - 1990} + (A12 - A11) \frac{T - 1990}{2020 - 1990}$

210 In this study, we present a new methodology for calculating the optical scattering properties of cirrus clouds across different wavelengths ~~is~~, as shown in Fig. 1. First, we calculate the LIWC ice water content of the cirrus cloud by Eq. 1, which is independent of wavelength, by introducing in Eq. 3 by (Vidot et al., 2015) the extinction coefficient for a cloud temperature in each vertical layer of the model. To align these calculations with the model vertical resolution, we previously degrade the vertical resolution of the cloud extinction and temperature profiles through vertical averaging.

215 *that has been partially depleted and is calculated by multiplying the extinction coefficient calculated with the knowledge that the total extinction coefficient is 0.582 (the extinction coefficient of the MBL), fact that the extinction coefficient is 0.4, and the total extinction coefficient is 0.582. This is done by assuming that the extinction coefficient is 0.4 and the total extinction coefficient is 0.582, which is entirely reasonable because the wavelength is 0.532 μm (Sun and Shine, 1994).*

$$220 \quad \text{LIWC}(z,t) = \frac{\pi (C_s + F_s T(z,t)) + \sqrt{(C_s + F_s T(z,t))^2 + 4 E_s \Theta(z,t)}}{2 E_s} \quad \lambda_o \in [0.38, 0.7] \mu\text{m} \quad (3)$$

$$\Theta(z,t) = \log_{10} \sigma_{ext,\lambda_o}(z,t) - A_s - B_s T(z,t) - D_s T^2(z,t)$$

$$225 \quad LIWC(z,t) = \frac{-(C_s + F_s T(z,t)) + \sqrt{(C_s + F_s T(z,t))^2 + 4 E_s \Theta(z,t)}}{2 E_s} \quad \lambda_o \in [0.38, 0.7] \mu\text{m}$$

$$\Theta(z,t) = \log_{10} \sigma_{ext,\lambda_o}(z,t) - A_s - B_s T(z,t) - D_s T^2(z,t) \quad (4)$$

230 The working wavelength of the lidar system is defined as λ_o and in this study it is 0.532 μm . The extinction coefficient, σ_{ext,λ_o} , is retrieved by the two-way transmittance method (Gil-Díaz et al., 2024). $T(z,t)$ denotes the cloud temperature, while A_s to F_s are the parametrization coefficients as defined in (Vidot et al., 2015). This formulation provides a unique physical solution and simplifies the IWC calculation based on extinction coefficient and cloud temperature, assuming no absorption, which is entirely reasonable because the working wavelength lies within the visible spectral range (Sun and Shine, 1994). Once, the IWC is obtained, this variable is introduced into Eqs. 2-4 (Vidot et al., 2015) to calculate the absorption, scattering and asymmetry factor coefficients for each wavelength. From these variables, the extinction and single scattering albedo are also subsequently derived. *Only the IWC of the cloud is calculated, this is the extinction coefficient, the extinction coefficient is 0.4, the total extinction coefficient is 0.582, the extinction coefficient is 0.4, the total extinction coefficient is 0.582, which is entirely reasonable because the wavelength is 0.532 μm (Sun and Shine, 1994).*

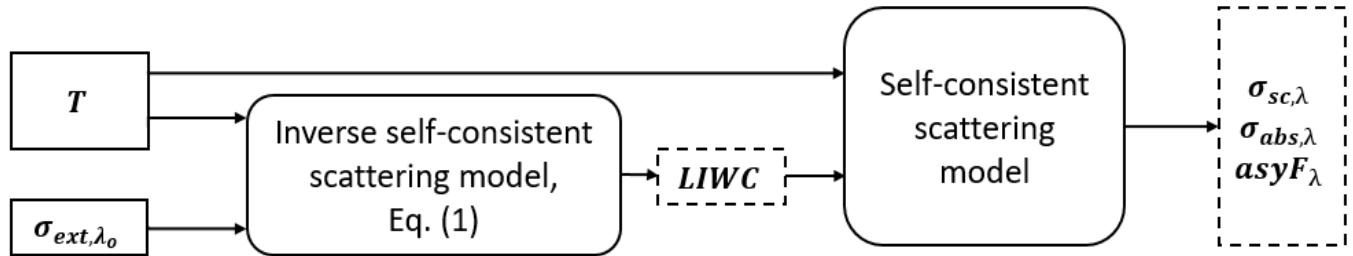


Figure 1. Scheme of the new model approach of the self-consistent scattering model for cirrus clouds.

235 The main advantage of this model is that it is not necessary to calculate the ice particle shapes and their size distributions in order to calculate their ~~radiative~~ optical scattering properties. These variables can be obtained easily with only elastic lidar systems and radiosondes or meteorological models. An example of the application of this method is shown in Fig. 2, for a cirrus cloud layer of 08/12/2018 at 12 UTC, measured in Barcelona.

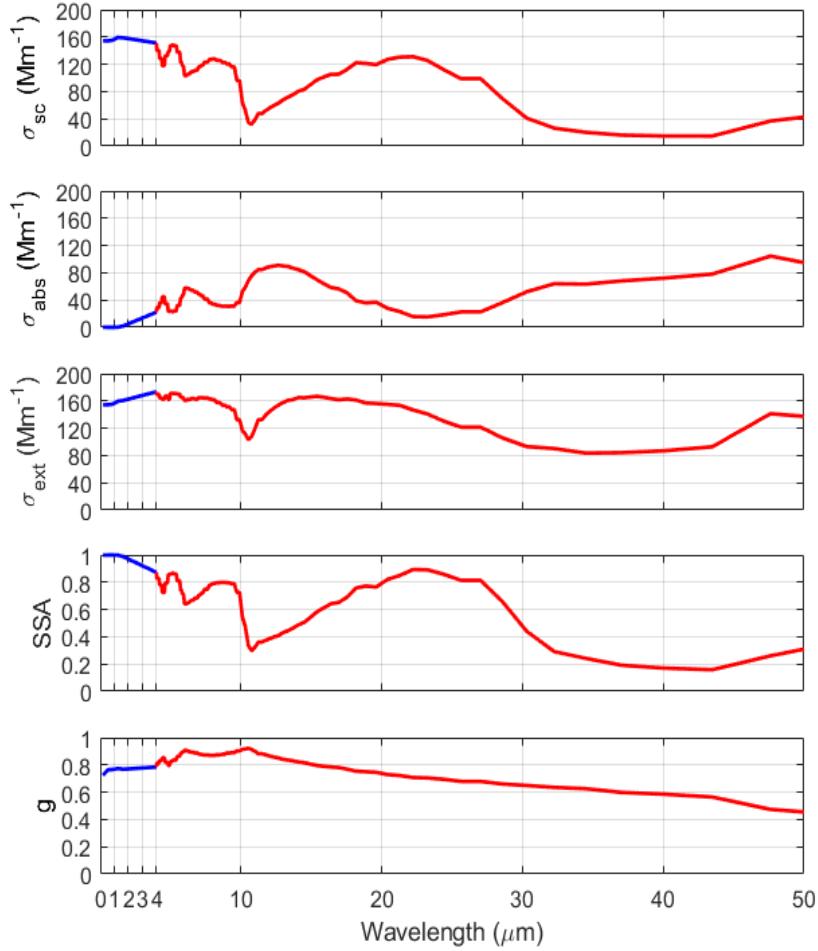


Figure 2. Spectral dependence of *radiative optical scattering* properties obtained by the self-consistent scattering model for cirrus clouds: scattering, absorption, extinction coefficients, single scattering albedo and asymmetry factor (from top to bottom) for a cirrus cloud of 08/12/2018 at 12 UTC, measured in Barcelona. The colours indicate the shortwave range (blue; 0.2-4 μm) and longwave range (red; 4-50 μm).

Fig. 2 shows the spectral dependence of *radiative optical scattering* properties obtained by the self-consistent scattering model for cirrus clouds, which reflects the characteristic *radiative optical scattering* properties of ice crystals. For example, it shows an absorption phenomenon negligible in the spectral range between 0.2 and 2 μm and consequently, the single scattering albedo is approximately the unity. As expected, the scattering phenomenon dominates the whole spectrum with respect to absorption except in the regions around 12 μm and for wavelengths higher than 30 μm (Yang et al., 2005, 2013). It is also noteworthy to mention that the single scattering albedo varies generally between 0.1 and 1 and at the working wavelength used

245 in the model ($0.532 \mu\text{m}$) it has a value of 0.99 (Sun and Shine, 1994). This fact supports the hypothesis made previously that the absorption phenomenon is negligible. The asymmetry factor presents much less variation as observed in the literature (Fu et al., 1998; Yang et al., 2005, 2013): it increases between 0.2 and $10 \mu\text{m}$ (in the range 0.75 - 0.95) and decreases afterwards (in the range 0.95 - 0.40).

3.2 The ARTDECO package

250 The Atmospheric Radiative Transfer Database for Earth and Climate Observation package (ARTDECO; <https://www.icare.univ-lille.fr/artdeco/>) is a numerical tool that gathers models and data for the 1D simulation of Earth atmosphere radiances and fluxes from the ultraviolet to thermal infrared range ($0.2\text{-}50 \mu\text{m}$). It is developed and maintained at the Laboratoire d'Optique Atmosphérique (LOA) and distributed by the data and services center AERIS/ICARE (University of Lille), and funded by the TOSCA program of the French space agency (CNES). In ARTDECO, users can either access a library for the scene or use 255 their own description through ASCII input files. Optical properties for aerosols and clouds can be computed. Then, the user can choose among available models to solve the radiative transfer equation and to compute radiative quantities corresponding to the scene. ARTDECO is thus a flexible tool for remote sensing or radiative forcing applications, such as sensitivity studies, development and optimization of retrieval algorithms, evaluation of the future instruments performances, etc.

260 In this study, DISORT model is employed to solve the radiative transfer equation by discretising it (Stamnes et al., 2000). The ARTDECO environment allows us to solve the radiative transfer equation in two ways: 1) by introducing our own phase matrix as a function of wavelength; 2) by using the Henyey-Greenstein function (Henyey and Greenstein, 1941), given extinction, single scattering albedo and asymmetry factor values over the whole spectral range in which the simulation will be done. Due to the lack of knowledge of the phase matrix of cirrus clouds with the observational measurements with which we work, the 265 second option is chosen even though the Henyey-Greenstein function does not represent a good approximation to the real phase function, especially for ice crystals. Upward and downward radiative fluxes are calculated at different vertical Levels: 31 layers (0-20 km) in the shortwave (SW, $0.2\text{-}4 \mu\text{m}$) and 40 layers (0-100 km) in the longwave (LW, $4\text{-}50 \mu\text{m}$) spectra. These spectral/vertical ranges are adjustable, together with their spectral/vertical resolution. Cirrus ~~filtering~~ direct radiative effects at the bottom-of-the-atmosphere (BOA) and top-of-the-atmosphere (TOA) have been calculated as:

$$270 \text{ BOA DR}\cancel{\text{E}} = (F_c \downarrow - F_c \uparrow) - (F_o \downarrow - F_o \uparrow) \text{ at BOA} \quad (5)$$

$$\text{TOA DR}\cancel{\text{E}} = (F_c \downarrow - F_c \uparrow) - (F_o \downarrow - F_o \uparrow) = -(F_c \uparrow - F_o \uparrow) \text{ at TOA} \quad (6)$$

275 Where F_c and F_o are the radiative fluxes with and without the cirrus cloud, respectively. The \downarrow and \uparrow arrows indicate whether the incoming solar radiation at the TOA is equal for both cases with and without aerosols. With this convention, a negative sign of $\text{DR}\cancel{\text{E}}$ implies a cirrus cooling effect independently of whether it occurs at the BOA or at the TOA. In this study, four types of simulations are carried out: with gases only (G), with gases and aerosols in the layer closest to the surface (GA), with gases

and a cirrus cloud (GC) and finally, with everything (GAC). In Section 4.3, the **direct radiative ~~filtering~~ effects** of cirrus clouds in the full atmosphere and in the whole spectral range considering aerosols (GAC-GA) or no aerosols (GC-G) are compared.

280 Besides aerosol optical properties, the radiative transfer model (RTM) DISORT is sensitive to atmospheric parameters such as the relative humidity and the air temperature profiles, the surface emissivity and temperature or the aerosol vertical distribution (Sicard et al., 2014).

The cirrus clouds are parameterized in the RTM model geometrically and optically, with the results obtained (Gil-Díaz et al., 285 2024) and ~~radiatively with~~ the retrievals obtained with the self-consistent scattering model for cirrus clouds (see Section 3.1). At the same time, the planetary boundary layer is characterized in the DISORT model geometrically and optically with the MPLNET AER product and radiatively, on one hand, in the SW with the AERONET products and, on the other hand, in the LW with the LITMS database.

3.2.1 Atmospheric profiles

290 The RTM DISORT model is run with atmospheric profiles (pressure, temperature, water vapor mixing ratio) obtained from radiosondes launched in Barcelona at 00 and 12 UTC. The profiles of ozone concentration are obtained from Copernicus Atmosphere Monitoring Service (CAMS) global reanalysis (EAC4) (Inness et al., 2019). EAC4 (ECMWF Atmospheric Composition Reanalysis 4) is the fourth generation ECMWF global reanalysis of atmospheric composition. Reanalysis combines model data with observations from across the world into a globally complete and consistent dataset using a model of the 295 atmosphere based on the laws of physics and chemistry. The dataset is globally distributed with a horizontal resolution of 0.75°x0.75° and a vertical extension of 60 modes (from 1000 to 1 hPa). In exceptional cases, when no radiosondes or CAMS data are available and for heights not covered by the radiosondes (generally above 30 km), the atmospheric profiles are taken from the 1976 standard atmosphere (COESA et al., 1976).

3.2.2 Surface properties

300 In this study, a Lambertian surface is considered. On one hand, the corresponding surface albedo over the Barcelona region is obtained for the SW range from AERONET and for the LW spectrum from ~~the CERES NOAA-20~~ measurements. On the other hand, the surface temperature is also taken from ~~the SSF CERES product NOAA-20 observations~~. In the parametrization of cirrus scenes, ~~the surface albedo has been obtained and the surface temperature has been derived from the CERES~~ ~~the surface emissivity in the longwave spectrum, the surface albedo in the shortwave spectrum and the surface temperature measured by the NOAA-20 satellite have been incorporated as instantaneous values. In contrast, the surface albedo measured by AERONET has been averaged seasonally due to its smoother temporal variation and its small influence on the simulations in that spectrum.~~

305

3.2.3 Cloud/aerosol stratification

310 The vertical stratification of cloud/aerosols is reproduced according to the vertical profiles of MPLNET products. On one hand, the base and top of the cirrus clouds are obtained from (Gil-Díaz et al., 2024). On the other hand, the vertical distribution of the aerosols in the planetary boundary layer (PBL) is provided by MPLNET AER product. When this product is not available for that specific time period, it is assumed that aerosols are uniformly distributed throughout the aerosol layer, which extends up to 1.5 km, being the mean PBL height obtained in Barcelona over a 3-year period (Sicard et al., 2006).

315 3.3 The CLaMS-Ice Lagrangian microphysical cirrus model

4 Five years of cirrus retrievals

This study analyzes the ~~radiative~~ optical scattering properties of cirrus clouds that were previously characterized geometrically, thermally and optically in (Gil-Díaz et al., 2024). In the latter paper, 203 cases were analyzed. Here, atmospheric scenes with only one cirrus cloud in the vertical profile are studied. 125 single-cirrus cloud scenes are found, that constitute 61% of all

340 cases. Case selection involves discarding cirrus scenes where a mid-level cloud has been detected below the cirrus cloud. Mid-level cloud detection was performed through visual analysis during the application of the two-way transmittance method (Gil-Díaz et al., 2024). This selection ensures the simulated cirrus scenes accurately represent reality (see Section 4.2).

4.1 Cirrus ~~radiative~~ optical scattering properties

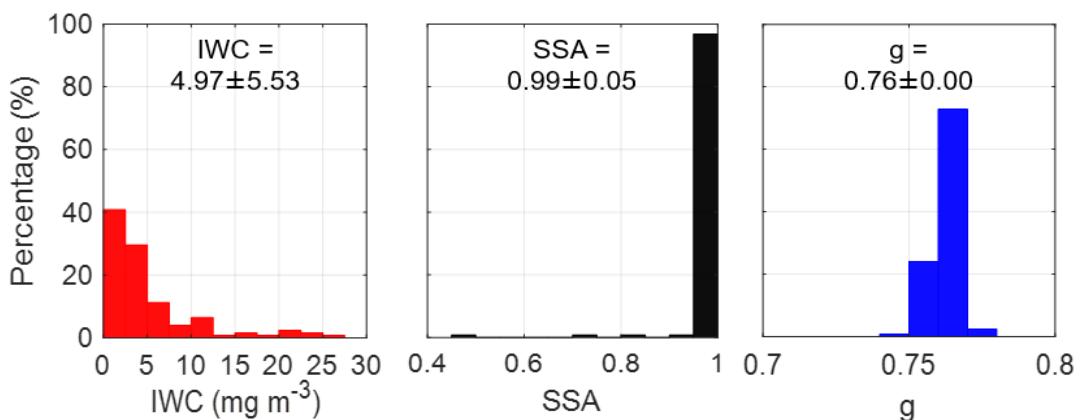


Figure 3. Probability distribution of averages of (left) ice water content (IWC), (center) single scattering albedo (SSA) and (right) asymmetry factor (ASYF) for each cirrus scenes at $0.55\text{ }\mu\text{m}$, for cirrus clouds measured from 2018 to 2022 in Barcelona.

In Fig. 3 one observes that cirrus clouds have an IWC between 0.03 and 30 mg/m^3 , being characteristic of mid-latitude cirrus clouds (Korolev et al., 2001; Field et al., 2005, 2006; Schiller et al., 2008; Baran et al., 2011b; Sourdeval, 2012; Kramer et al., 355 2016, 2020). Where the average of IWC is $\sim 5 \text{ mg/m}^3$, being a value close to 3 mg/m^3 , which is the central value of the mid-latitude ice cloud distributions obtained by (Sourdeval, 2012) and the mean value of IWC for temperatures between 210 and 235K found in (Kramer et al., 2016). A slightly higher measured IWC value of 7 mg m^{-3} was found by (Korolev et al., 2001) for cirrus clouds whose temperature ranged from 233 to 243K. The single scattering albedo of most cirrus clouds (97%) has a

value between 0.95 and 1, as expected at $0.55 \mu\text{m}$ (Hess and Wiegner, 1994; Sun and Shine, 1994; Yang et al., 2013; Hemmer, 360 2018). Although there are 3 cases of cirrus clouds whose $\text{SSA} < 0.9$. These cases correspond to sub-visible cirrus clouds with an IWC less than 1 mg/m^3 . These 3 cases have in common that the cirrus cloud extends in two vertical layers of the model and in one of the layers, the ~~lowest value of the effective extinction coefficient~~ extinction coefficient is less than 1 Mm^{-1} . In this layer, a low value of SSA is obtained, associated to its low value of the extinction coefficient and consequently, when averaging the ~~radiative~~ optical scattering properties in the two layers for each cirrus scene, the values of $\text{SSA} < 0.9$ observed in Fig. 3 are 365 obtained. Therefore, the self-consistent scattering model for cirrus clouds might associate the low ~~effective extinction~~ extinction values to super-cooled liquid water content in the cirrus clouds. The asymmetry factor of cirrus clouds varies between 0.7 and 0.8, with an average of 0.76, as expected at $0.55 \mu\text{m}$ (Hess and Wiegner, 1994; Sun and Shine, 1994; Yang et al., 2013; Hemmer, 2018).

4.2 Validation of radiative fluxes

370 The validation of the ARTDECO package is performed by comparing the simulated radiative fluxes with observed ones. For that purpose, the radiative fluxes from ARTDECO were recalculated in the range $0.305\text{--}2.8 \mu\text{m}$ corresponding to the spectral range of the SolRad-Net pyranometer (BOA), and in the range $5\text{--}100 \mu\text{m}$ corresponding to the spectral range of ~~CHERES~~ NOAA-20 (TOA). The scatter plot of the simulated vs. observed SW downward radiative fluxes of cirrus clouds classified according to their cloud optical depth (Sassen and Cho, 1992), at the surface and the LW upward radiative fluxes at TOA are shown in 375 Fig. 4.

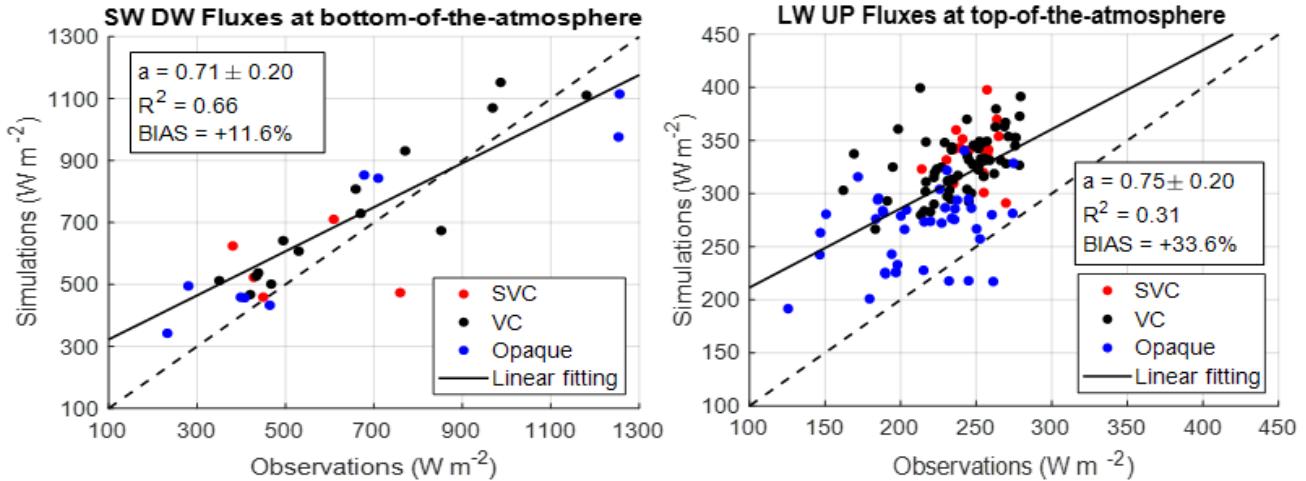


Figure 4. Comparison of (left) simulated shortwave downward (SW DW) radiative fluxes at the bottom-of-the-atmosphere and SolRad-Net observations; (right) simulated longwave upward (LW UP) radiative fluxes at top-of-the-atmosphere and *CHERES NOAA-20* observations. The black dashed line is the curve with the slope unity and the black solid line is the linear fitting of the fluxes ($y=ax+b$, being a the slope and R^2 its determination coefficient). The cirrus clouds have been classified according to (Sassen and Cho, 1992) criteria: sub-visible (SVC; COD < 0.03), visible (VC; 0.03 < COD < 0.3) and opaque clouds (Opaque; COD > 0.3).

The validation of the SW downward radiative fluxes is performed with 59% of the cirrus clouds measured at daytime and the validation of the LW upward radiative fluxes with 81% of all cirrus clouds considered in this study. The cases of cirrus clouds that could not be validated are due to the lack of observations. In addition, to reduce the effect of cloud movement on the radiation measurement with the pyranometer, the observed radiation fluxes are averaged over 30 minutes. In Fig. 4 380 (left) it can be seen that most downward radiative fluxes calculated with the DISORT model overestimate the SolRad-Net observations: the mean and standard deviation of the simulated fluxes are $694 \pm 247 \text{ W m}^{-2}$, while it is $621 \pm 283 \text{ W m}^{-2}$ for the observations. This translates into a systematic BIAS of +11.6% with a steep slope of the linear regression ($a = 0.71 \pm 0.20$). The overestimation may be related to the error associated with variables obtained by the self-consistent scattering model 385 for cirrus clouds, as cirrus clouds govern the radiation interactions in these simulations, because of their cloud optical depth. The ensemble scattering model for cirrus clouds has a large error for small ice crystals (less than $100 \mu\text{m}$), corresponding to cirrus clouds with low IWC values (Liou et al., 2008). In particular, the model tends to underestimate the IWC for mid-latitude cirrus clouds (Baran and Labonnote, 2007). Therefore, when the IWC is lower, the extinction of cirrus clouds is smaller as demonstrated in (Fu, 1996; Heymsfield et al., 2014) and, consequently, allows more radiation to pass into the atmosphere than 390 actually does.

390

As For validation in the longwave spectrum, *CHERES NOAA-20* measurements have been selected based on their ~~geographical position~~ *geographical proximity to the Barcelona lidar station. Specifically, The measurements closest to the Barcelona lidar*

395 this station ~~have been selected~~ were chosen, ~~despite the fact that the height of the mast height does not correspond to the height of the instrument~~ even though the measurement times do not exactly match the atmospheric scene. ~~The hourly difference between~~ ~~simulations and observations is 1 hour, with a maximum difference of 3 hours. However, this discrepancy is not significant for the validation of longwave radiation fluxes~~, since it is almost constant during daytime hours (Sicard et al., 2014). In Fig. 4 (right) a large horizontal dispersion can be observed. In addition, a general overestimation of the *CERES* NOAA-20 observations with the ARTDECO simulations is produced, being for simulated fluxes ~~857±46~~ $310 \pm 43 \text{ Wm}^{-2}$ and for the observations ~~235±32~~ $232 \pm 32 \text{ Wm}^{-2}$. In our case, the large BIAS = ~~+51.62~~ 33.6% obtained could be due to the spatial resolution of the observed measurements taken, which ~~is only 0.25 km and the satellite~~ may cover part of the Mediterranean Sea. In addition, the cloud mask associated with each observation indicates that in 14% of the cases it has more than 90% of clear-sky footprint area. As demonstrated by ~~the~~ (Gil-Díaz et al., 2024) most of the cirrus clouds are visible and therefore their horizontal expansion is smaller than the cirrus clouds that form at higher altitudes (well-known as sub-visible cirrus clouds) (Kramer et al., 2020). ~~This~~ makes them more ~~difficult to detect~~ challenging to detect from top-of-the-atmosphere. Hence, the comparison of simulated radiative fluxes and *CERES* NOAA-20 observations is not as trivial and conclusive as with SolRad-Net observations, since the *CERES* NOAA-20 satellite can observe a slightly different atmospheric situation, as mentioned above. Not to mention the limitations of the 1-D radiative transfer model DISORT to represent an irregular composition of broken and/or overlapping clouds that the *CERES* NOAA-20 satellite could observe.

410 4.3 Study of the influence of aerosols on radiative simulations of cirrus clouds

The cloud direct radiative ~~forcing~~ effect is often calculated as the difference between radiative fluxes under cloudy and cloudy-free conditions, without considering the aerosols found in the layer of the atmosphere closest to the surface (Ramanathan et al., 1989; Hartmann et al., 2001; Barja and Antuña, 2007; Yang et al., 2007; Lee et al., 2009; Campbell et al., 2016; Lolli et al., 2017b). In this way, simpler simulations are carried out, as it is only necessary to characterize the cloud. It is known that 415 radiation does not interfere linearly with the components of the atmosphere: clouds, aerosols and gases. Therefore, in this subsection we analyze whether or not the insertion of aerosols in the lowest atmospheric layer in cirrus cloud scenes modifies the calculation of cirrus direct radiative ~~forcing~~ effects. For this purpose, the net direct radiative effect ~~forcing~~ (NET; SW+LW) in the atmosphere (ATM; TOA-BOA) of cirrus clouds calculated with simulations in which aerosols have been considered or not is compared, as shown in Fig. 5.

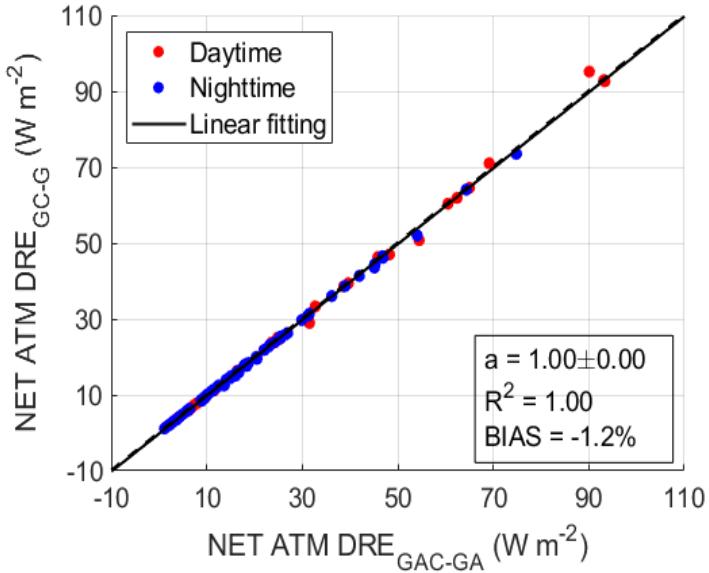


Figure 5. Comparison of net **direct radiative *floating* effects** in the full atmosphere (NET ATM DRE) between simulations made with (X-axis) and without (Y-axis) aerosols. The black dashed line is the curve with the slope unity and the black solid line is the linear fitting of the **direct radiative *floating* effects** ($y = ax+b$, being a the slope). The cirrus clouds have been classified according to (Sassen and Cho, 1992) criteria: sub-visible (SVC; COD < 0.03), visible (VC; $0.03 < COD < 0.3$) and opaque clouds (Opaque; COD > 0.3).

420 Fig. 5 shows that the NET ATM **direct radiative *floating* effects** calculated with and without aerosols fit well, with the most of the points lying slightly above on the curve with the slope unity. As a consequence, its linear fitting rounding to the tenth has also a unity slope, with a R^2 value of 1.00. In addition, the mean and standard deviations of simulations reflect that there is an negligible underestimation of the forcings when not considering aerosols, with values for the simulation with aerosols (X-axis) of $427.1 \pm 23.94 + 21.2 \pm 20.3 \text{ W m}^{-2}$ and without aerosols (Y-axis) of $427.0 \pm 23.86 + 20.9 \pm 20.3 \text{ W m}^{-2}$. Furthermore, the
 425 BIAS is $-0.7 \pm 1.2\%$, being a low value, possibly due to the distance between the aerosol layer and the cirrus cloud (being on average $6.76 \pm 2.24 \text{ km}$). With these results where the aerosol layer was well distinguished vertically from the cirrus clouds, the simplification of the atmospheric scenes can be made without considering aerosols, but to be more rigorous, in the following results, only the forcings in which aerosols are present will be considered. In the other case where the aerosols are vertically closer to the clouds (lower than 1 km, being the minimum distance found between the cirrus cloud and the aerosol layer),
 430 the simplification of not considering aerosols in the calculation of cloud forcings may not be valid, leading to a significant underestimation of cloud **floating** **direct radiative effects**.

4.4 Direct radiative ~~forcing~~ effects of cirrus clouds depending on COD

In this section, only **direct** radiative ~~forcing~~ effects of cirrus clouds calculated with simulations in which aerosols are present, 435 $DRFE_{GAC-GA}$, will be considered and will be denoted as $DRFE$. Special attention will be paid to net **direct** radiative ~~forcing~~ effects of cirrus clouds at daytime because they are the only clouds that can readily cool or warm the top and bottom of atmosphere, during daytime, depending on their properties (Campbell et al., 2016). In order to quantify this phenomenon, cirrus clouds at daytime and nighttime have been distinguished. The net **direct** radiative ~~forcing~~ effects of cirrus clouds at nighttime and daytime, at BOA and TOA are shown in Fig. 6.

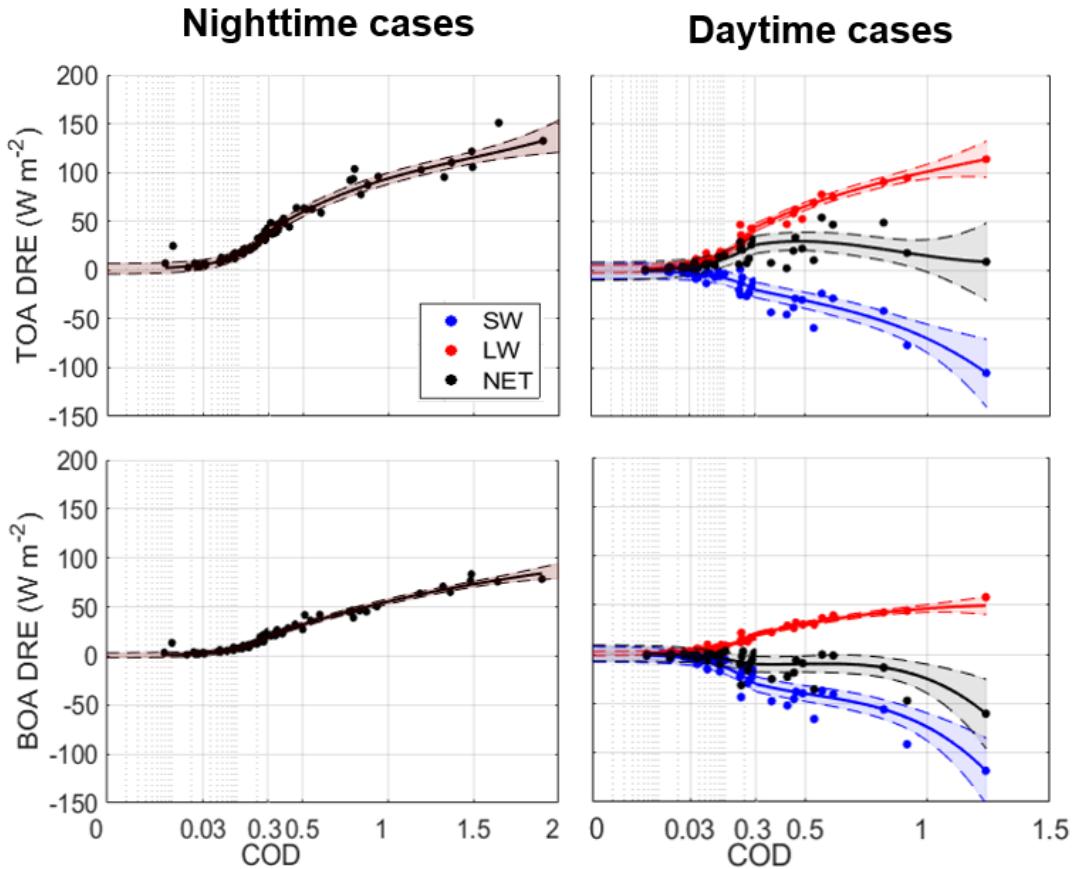


Figure 6. Distribution of **direct** radiative ~~forcing~~ effects of cirrus clouds at (left) nighttime and (right) daytime, at (bottom) bottom-of-the-atmosphere (BOA) and (top) top-of-the-atmosphere (TOA), in function of their cloud optical depth (COD). Shortwave (SW), longwave (LW) and net (NET = SW + LW) components of **direct** radiative ~~forcing~~ effects have been distinguished. For COD < 0.3 a logarithmic scale has been considered in order to discern more clearly sub-visible and visible cirrus clouds (Sassen and Cho, 1992). The solid line corresponds to the polynomial fitting performed on the data set. The shaded area represents the region with a 95% probability of containing the points, adjusted by the mean absolute value of the differences between the actual and fitted values.

440 In Fig. 6 one observes a positive trend between the net **direct** radiative ~~forcing~~ **effects** with the COD, where the thicker cirrus clouds contribute more to the overall ~~forcing~~ **direct radiative effect** budget, as has been observed in other studies (Barja and Antuña, 2007; Lee et al., 2009; Campbell et al., 2016; Lolli et al., 2017b). Some COD gaps are also found, because the cirrus observations considered do not have a homogeneous and equidistant distribution of COD. At nighttime, the net cirrus ~~forcing~~ **direct radiative effect** at TOA is approximately twice that at BOA, being always positive as expected. Cirrus 445 clouds at nighttime act as a cover in the atmosphere, they do not let through all the infrared radiation emitted by the Earth as it cools, inducing a warming of the atmosphere. This warming in function of COD is ~~faster~~ **more pronounced** at TOA than at BOA because the atmosphere at BOA is strongly influenced by the surface, which acts as a black body emitting infrared radiation at nighttime. Consequently, the ~~Heating~~ **positive direct radiative effect** at BOA is milder than at TOA.

450 At daytime, at TOA, the net **direct** radiative ~~forcing~~ **effect** remains positive for all cirrus clouds, dominating the positive longwave component. This effect has been observed in (Campbell et al., 2016) for COD up to approximately 0.6. For higher COD values, (Campbell et al., 2016) reports a negative NET TOA DRE. In this study, a decreasing trend in NET TOA DRE is observed from COD values of 0.5, although no negative values are obtained. Additionally, the LW NET TOA DRE component grows faster than the one reported by (Campbell et al., 2016), suggesting that negative values of NET TOA DRE could occur 455 for cirrus clouds with higher COD than those found in (Campbell et al., 2016). This discrepancy may be due to the higher surface emissivity and temperature values considered in the present work. Further measurements of NET TOA DRE for cirrus clouds with higher COD are needed to confirm the decreasing trend. In contrast, at BOA, the net **direct** radiative ~~forcing~~ **effect** is almost always negative (only 2018% of the cases show a positive net ~~forcing~~ **direct radiative effect**, whose value is close to 0), being the outgoing shortwave radiation in the presence of cirrus clouds larger than in cirrus cloud free conditions. The 460 albedo effect overcomes the greenhouse effects in the SW range because of low absorption capacity of the small crystals, as shown in Fig. 2. These changes of sign which occur ~~mostly~~ for thin cirrus clouds (~~57% of the cases~~) ~~mostly~~ like in other studies ~~such as~~ (Campbell et al., 2016; Lolli et al., 2017b; Kramer et al., 2020), where the dominant factor in the change of sign of forcing is unclear. Multiple factors are involved, from the optical ~~and~~ properties of the cirrus, to the solar zenith angle or the surface temperature and surface albedo (Wolf et al., 2023).

465

In order to complete this analysis, the net **direct** radiative ~~forcing~~ **effects** in the full atmosphere are analyzed for cirrus clouds at nighttime and daytime, as shown in Fig. 7.

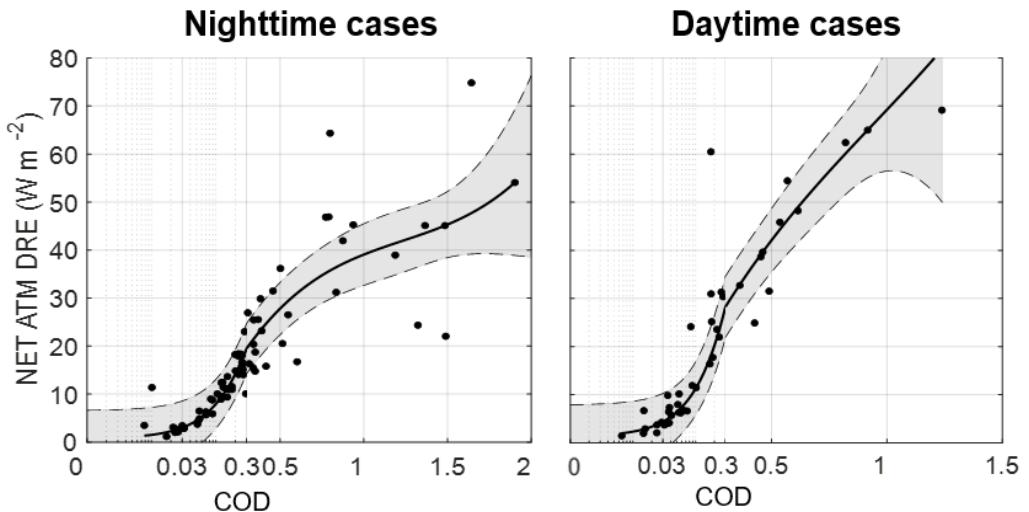


Figure 7. Distribution of net *direct radiative forcing effects* of cirrus clouds in the full atmosphere at (left) nighttime and (right) daytime, in function of their cloud optical depth (COD). For COD < 0.3 a logarithmic scale has been considered in order to discern more clearly sub-visible and visible cirrus clouds (Sassen and Cho, 1992). The solid line corresponds to the polynomial fitting performed on the data set. The shaded area represents the region with a 95% probability of containing the points, adjusted by the mean absolute value of the differences between the real and fitted values.

Fig. 7 shows a net warming of the atmosphere (always positive *forcing direct radiative effect*) for cirrus clouds at nighttime and daytime, being radiation escape lower in the presence of cirrus clouds in the full atmosphere. This phenomenon could have
 470 been perceived in the previous figure (Fig. 6) as the *forcing direct radiative effect* at TOA was always higher than at BOA. It also results that the atmosphere warms faster in function of the COD during the daytime than at nighttime (see their regression slopes), as expected due to the contribution of solar radiation to the net *forcing direct radiative effect*. The net *forcing direct radiative effect* in the full atmosphere fits very well with the polynomial regressions for both time periods, being at nighttime $R^2 = 0.990.96$ and at daytime $R^2 = 0.970.95$, although some instances outside the shaded area are observed. These strong
 475 variations of the *direct radiative forcing effect* for cirrus clouds with very similar COD are due to the consideration in the simulations of different *radiative optical scattering* properties of the cirrus clouds, thermodynamic profiles, surface temperature and surface albedo values for each cirrus scene.

Then, the *direct radiative forcing effects* for cirrus clouds at nighttime and daytime, which are classified according to
 480 (Sassen and Cho, 1992) criteria, are quantified, as shown in Table 1.

| Type | Nighttime | | | | Daytime | | | |
|-------------|-----------------|-----------------|-----------------|-----------------|-----------------|------------------|-----------------|-----------------|
| | COD | BOA | TOA | ATM | COD | BOA | TOA | ATM |
| Sub-visible | 0.02 \pm 0.01 | 3.9 \pm 3.9 | 7.5 \pm 7.1 | 3.6 \pm 3.3 | 0.02 \pm 0.01 | -0.7 \pm 1.2 | 2.6 \pm 1.3 | 3.3 \pm 1.8 |
| Visible | 0.15 \pm 0.08 | 11.1 \pm 5.6 | 22.2 \pm 10.4 | 11.1 \pm 5.2 | 0.11 \pm 0.08 | -4.4 \pm 7.4 | 11.3 \pm 9.1 | 15.6 \pm 13.1 |
| Opaque | 0.76 \pm 0.47 | 42.2 \pm 20.3 | 74.8 \pm 31.5 | 32.6 \pm 15.4 | 0.75 \pm 0.39 | -28.2 \pm 23.8 | 25.4 \pm 17.4 | 53.6 \pm 21.6 |
| Total | 0.37 \pm 0.43 | 22.1 \pm 20.8 | 40.4 \pm 34.5 | 18.4 \pm 15.3 | 0.28 \pm 0.37 | -11.5 \pm 18.7 | 14.2 \pm 14.0 | 25.7 \pm 25.9 |

Table 1. Average and standard deviation of **direct** radiative *forcing* effects of cirrus clouds (Wm^{-2}) at bottom-of-the-atmosphere (BOA), top-of-the-atmosphere (TOA) and in the full atmosphere (ATM), at nighttime and daytime, classified with (Sassen and Cho, 1992) criteria in Barcelona. Cloud optical depth values are obtained from (Gil-Díaz et al., 2024).

In Table 1 it is discernible that thicker cirrus clouds produce a higher *forcing* **direct radiative effect** than thinner clouds, as observed above. At nighttime, cirrus clouds produce an average net warming in the full atmosphere of $427.04 \pm 121.89 + 40.4 \pm 34.5 \text{ Wm}^{-2}$, with opaque cirrus clouds being the main source. At daytime, cirrus clouds **produce** generally *decreasing* a negative **direct radiative effect** at BOA and *increasing* a positive effect at TOA, resulting in a warming of the full atmosphere. The **direct** *forcing* **direct radiative effect** at BOA ranges between -58.44 ± 73 and $+3 \text{ Wm}^{-2}$ for all cirrus. In particular, for thin cirrus clouds the **direct radiative forcing effect** is in a range from -20.31 to $+3 \text{ Wm}^{-2}$, being a similar range to (Lee et al., 2009), covering from -20 to 0 Wm^{-2} . Therefore, shortwave negative *forcing* **direct radiative effect** generally dominates at the BOA, with an average of $-8.18 \pm 6.48 - 3.6 \pm 6.7 \text{ Wm}^{-2}$ for thin cirrus clouds, being slightly lower than -1.35 Wm^{-2} (Lee et al., 2009). The **direct radiative forcing effect** at TOA ranges between $+1$ and $+67 + 54 \text{ Wm}^{-2}$ for all cirrus, being a wider interval than (Campbell et al., 2016; Lolli et al., 2017b; Kramer et al., 2020; Kienast-Sjögren et al., 2016). In particular, for thin cirrus clouds the **direct radiative forcing effect** is in a range from $+1$ to $+20 + 32 \text{ Wm}^{-2}$, being the maximum value considerably higher than the value of $+5.71 \text{ Wm}^{-2}$ (Campbell et al., 2016) or $+10 \text{ Wm}^{-2}$ (Kramer et al., 2020). The average of the **direct radiative forcing effect** for thin cirrus clouds is $+11.90 \pm 10.45 + 9.4 \pm 8.8 \text{ Wm}^{-2}$, being a close value compared to them obtained in (Kienast-Sjögren et al., 2016), that cover in average from 6.2 to 11 Wm^{-2} , although they are also significantly higher than the value of $+1 \text{ Wm}^{-2}$ (Lee et al., 2009). Despite the differences found, the values are in agreement in magnitude with other studies such as (Ackerman et al., 1988; Jensen et al., 1994; Lee et al., 2009; Berry and Mace, 2014; Campbell et al., 2016; Kienast-Sjögren et al., 2016; Kramer et al., 2020). The average net warming in the full atmosphere is a little higher at daytime, with an average value of $+27.51 \pm 26.63 + 25.7 \pm 25.9 \text{ Wm}^{-2}$. This difference is apparently not related with the fraction of opaque cirrus, as the percentage of opaque cirrus is lower during nighttime (28%) than at daytime (38%).

500 4.5 Direct radiative *forcing* effects of cirrus clouds depending on solar zenith angle

In this section, only net **direct radiative forcing effects** of cirrus clouds at daytime, will be considered. The distinction between the shortwave and longwave spectral ranges will not be made because **direct radiative forcing effect** does not depend on the solar zenith angle in the longwave spectrum (Lee et al., 2009; Wolf et al., 2023). The net **direct radiative forcing effect** of cirrus clouds during daytime at the BOA and TOA, together with the results of a brief sensitivity study, in which all parameters of the simulations except the solar zenith angle are kept constant, are shown in Fig. 8.

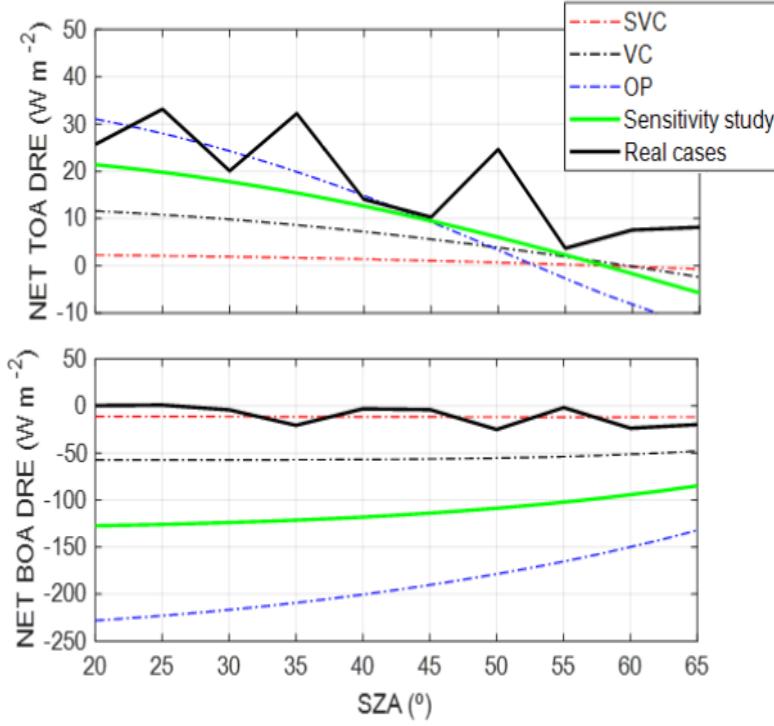


Figure 8. Distribution of the net **direct** radiative **forcing** effects of cirrus clouds at daytime at (top) top-of-the-atmosphere (TOA) and (bottom) bottom-of-the-atmosphere (BOA) in function of their solar zenith angle (SZA) resulting from the (dashed curves correspond to for each cirrus cloud type according to the (Sassen and Cho, 1992) criteria and green curve correspond to the mean values for all daytime cirrus clouds) sensitivity study and (black curve) the real cases.

In the sensitivity study, three cloud types have been considered, according to the (Sassen and Cho, 1992) criteria, where the cloud optical depth is shown in Table 1 for the daytime cirrus clouds. The geometrical, thermal and optical properties correspond to the mean values of the cirrus clouds measured at Barcelona lidar station (Gil-Díaz et al., 2024) and the **radiative** optical scattering properties are obtained from the mean values resulting from the statistic. The longwave surface albedo considered is 0.0157 and the surface temperature is 28°C. The thermodynamic profiles have been selected from the 1976 standard atmosphere (COESA et al., 1976). By keeping constant all other properties of cirrus clouds scenes, the results are expected to be exclusively due to the variation in the value of the solar zenith angle. On the other hand, in the real cases that have been analyzed in Section 4.4 and Section 4.5 as a function of cloud optical depth, the net **direct** radiative **forcing** effect values for the three cloud types have been averaged, with a SZA resolution of 5°.

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In Fig. 8 one observes a higher variability of net **direct** radiative **forcing** effect for the real cases than for the sensitivity study. This variability could be explained by other parameters that are considered in the simulations, such as cloud optical depth, cirrus **radiative** optical scattering properties, surface albedo and temperature (Sicard et al., 2014; Wolf et al., 2023). **While**

520 a detailed sensitivity analysis is needed to assess the contribution of each variable, cloud optical depth appears to have the
most significant impact, as the difference in mean CODs for each averaged SZA value is substantial. In addition, a generally
higher mean net **direct** radiative *filtering* effect is discerned for the real cases than for the sensitivity study results, especially at
BOA despite having taken the average properties resulting from the statistics. At TOA, a slight downward tendency of the net
direct radiative *filtering* effect is obtained as the SZA increases, as found in (Wolf et al., 2023). As SZA increases, cloud solar
extinction is enhanced regarding thermal effects (Campbell et al., 2016). All mean net **direct** radiative *filtering* effect values
525 for the real cases are positive, but there is a large fluctuation in certain values of SZA. Moreover, there is no crossover where
the mean net **direct** radiative *filtering* effect shifts from positive to negative values between SZA of 20 to 65°. On contrary, for
the results from the sensitivity study, a change of sign is observed at 58°, fixing well with results from (Campbell et al., 2016;
Lolli et al., 2017b). At BOA, most net **direct** radiative *filtering* effect values are negative, presenting a slight increasing trend
unlike the TOA. Since the angle of incidence of the incoming solar radiation increases, the incident solar radiation is lower and
530 the scattering produced by the ice crystals increases because optical path is larger (Lee et al., 2009; Wolf et al., 2023). This
enhancement of the extinction of incident solar radiation, which is lower, results in a reduction of the net *attenuating* positive
direct radiative effect at TOA and the net *absorbing* negative effect of the atmosphere at BOA. The change from positive to
negative values of the mean net **direct** radiative *filtering* effect is only observed for the real cases, being SZA crossover of 25°,
a considerably lower value than that observed for the results of the TOA sensitivity study.

535 **4.6 Case study of **direct** radiative *filtering* effects of an evolving cirrus cloud**

In this section the role of time is added to the analysis and the **direct** radiative *filtering* effect produced by a cirrus cloud is studied
along its back-trajectory. The final objective of this case is to simultaneously analyze the evolution of different physical agents
such as the surface albedo, the solar zenith angle, the cloud optical depth or the ice water content in the quantification of the
cirrus cloud **direct** radiative *filtering* effect. The case of study corresponds to the back-trajectory of a cirrus cloud measured at
540 Barcelona lidar station on 11/02/2019 at 02:03 UTC (Gil-Díaz et al., 2024), where simultaneous measurements of the MPL and
CALIPSO were performed. To simulate the evolution of the cirrus cloud as realistic as possible, its microphysical properties
along its back-trajectory are calculated with the CLaMS-Ice model. This model provides apart from the basic back-trajectory
variables such as temperature, altitude, geographic coordinates and time, microphysical properties like ice water content, ice
545 crystal number concentration or ice nuclei concentration. Considering the temperature provided by the model as the mid-cloud
temperature and together with the ice water content, the *radiative* optical scattering properties of cirrus clouds are specifically
calculated with *Eq/4/and/Eq/5* the self-consistent scattering model for cirrus clouds. Assuming the cirrus cloud geometric
thickness decreases linearly, using the CALIPSO measurements of 1 km on 10/02/2019 14 UTC and 1.74 km on 11/02/2019 2
UTC, and considering that the geometrical thickness is 0 km when the IWC is null, the COD is estimated as the product of the
extinction coefficient and the geometric thickness. The back-trajectory of the cirrus cloud and the properties: ice water content,
550 mid-cloud temperature, mid-cloud altitude, cloud geometrical thickness and cloud optical depth are shown in Fig. 9.

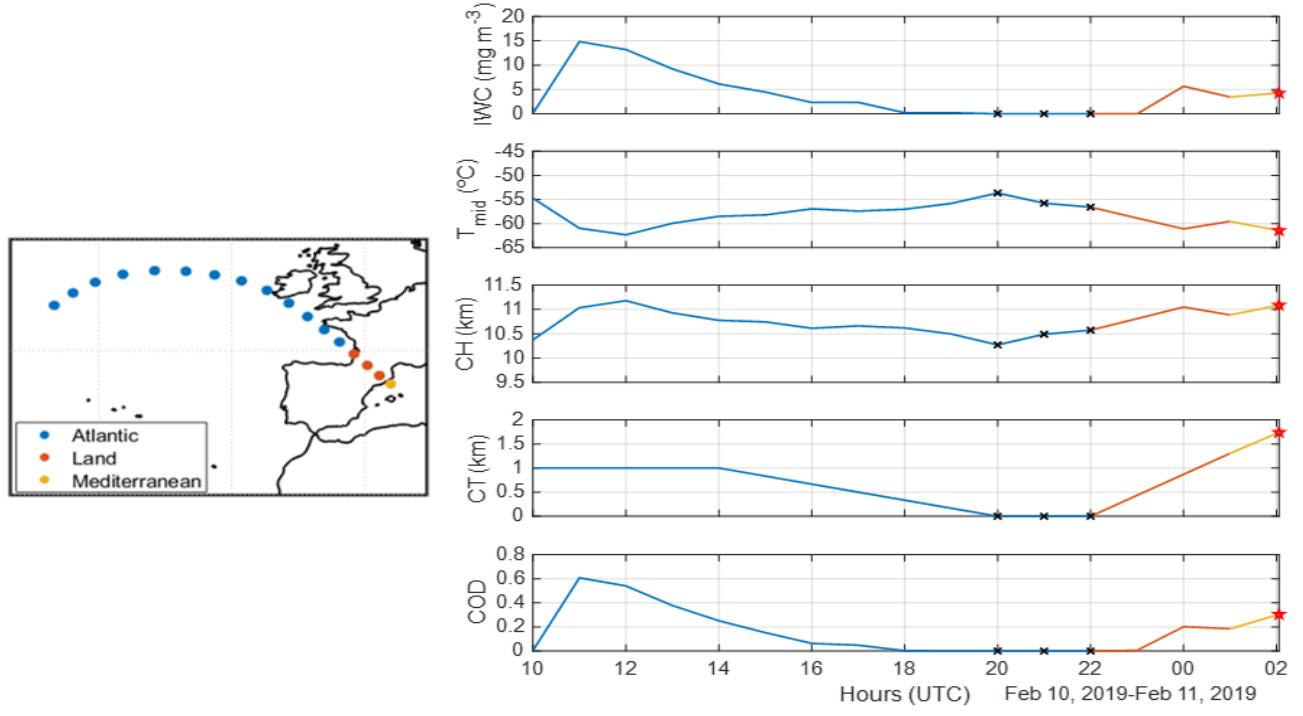


Figure 9. (Left) Hourly back-trajectory of the cirrus cloud for the last 16 hours before arrival in Barcelona. (Right) Evolution of cirrus cloud properties along its back-trajectory: ice water content (IWC), mid-cloud temperature (T_{mid}), mid-cloud altitude (CH), cloud geometrical thickness (CT) and cloud optical depth (COD) (from top to bottom). Colours indicate the type of surface where the cloud is over, black crosses mark the non-existence of cirrus cloud during those hours and the red star points to the case measured by the MPL and CALIPSO in (Gil-Díaz et al., 2024).

In Fig. 9 one observes that the cirrus cloud comes from the Atlantic Ocean, passing through part of France and Barcelona to reach the Mediterranean Sea. The selected points of the trajectory are spaced 1 hour apart backward from 11 February 2019 at 02 UTC. During this journey, the air mass which corresponds to the cloud undergoes a rise in height, reaching a minimum in temperature and a COD of 0.6. After this initial cirrus cloud fast growth, the cloud gradually fades away until 20 UTC, where 555 the CLAMS-Ice model gives a null ice water content. Afterwards, a new cirrus begins to form over land surface and to grow until 00 UTC and remains relatively stable.

Taking advantage of the overlap between the back-trajectory of this cirrus cloud and the CALIPSO satellite, a brief evaluation of the ice water content is carried out on 11/02/2019 02:03:20 UTC. The difference in IWC between the CLAMS-Ice output 560 and the CALIPSO measurement is 0.88 mg/m^3 , resulting in a relative deviation of 29% for the CLAMS-Ice output. This difference obtained by CLAMS-Ice and CALIPSO, well within a factor of two, is considered reasonable given the different parameterizations involved in the calculation of the IWC and their respective uncertainties. To fully characterize the cirrus

cloud scenes along its back-trajectory, it is considered that the height from the CLaMS-Ice model is the mid-cloud height.
565 ~~the direct radiative effects at TOA, BOA and the full atmosphere are shown in Fig. 10.~~
Balloon. The surface albedo and surface temperature considered are from the NOAA-20 satellite. With all these assumptions, radiative simulations are calculated with the ARTDECO package and the ~~radiative effects~~ direct radiative effects at TOA, BOA and the full atmosphere are shown in Fig. 10.

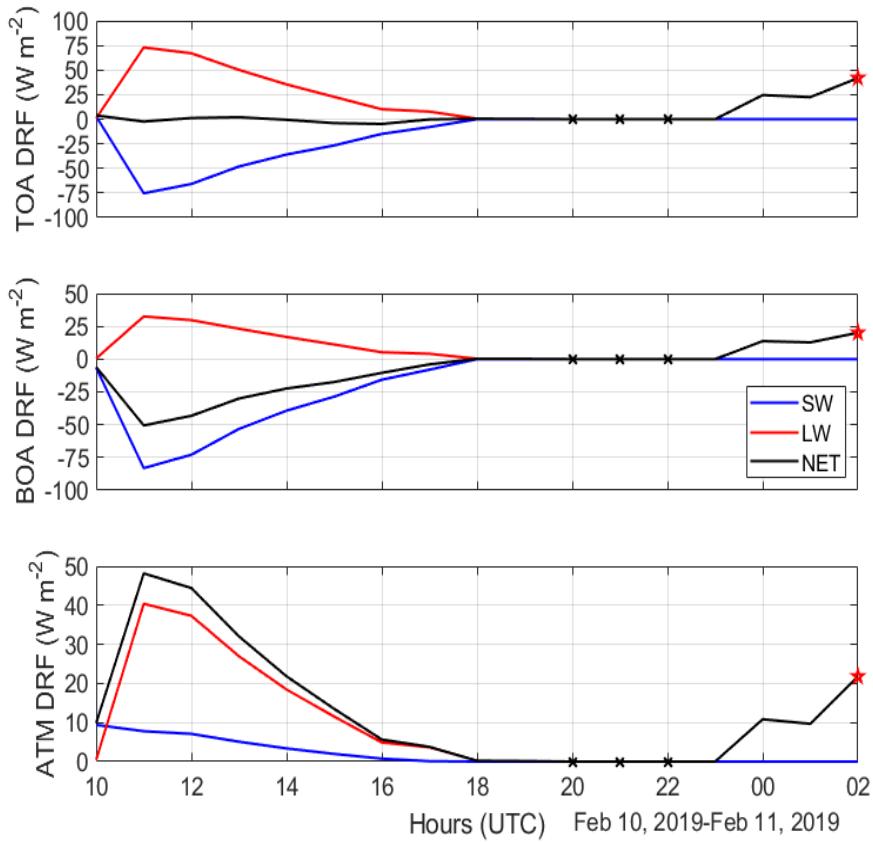


Figure 10. Distribution of direct radiative ~~forcing~~ effects of the cirrus cloud along its back-trajectory at (top) top-of-the-atmosphere (TOA), (centered) bottom-of-the-atmosphere (BOA) and (bottom) in the full atmosphere (ATM). Shortwave (SW), longwave (LW) and net (NET = SW + LW) components of direct radiative ~~forcing~~ effects have been distinguished. Black crosses mark the non-existence of cirrus cloud during those hours.

Fig. 10 shows that at TOA the net direct radiative ~~forcing~~ effect is close to zero with values shifting between positive and
570 negative during the first hours where the cirrus cloud is over the Atlantic Ocean. During this period, the SW component is almost completely balanced by the LW component and it is zero when there is no incident solar radiation, i.e. at nighttime,

therefore the net ~~forcing~~ direct radiative effect corresponds to the LW component. At nighttime, the net direct ~~direct~~ radiative ~~forcing~~ effect at TOA is also approximately double than at BOA, in agreement with previous results. At BOA, the net forcing changes from negative to positive, since the incident solar radiation produces that the albedo effect overcomes the greenhouse effect during daytime. In the full atmosphere, the ~~direct~~ radiative ~~forcing~~ effects are always positive, as the ~~forcing~~ direct radiative effect at TOA is higher than at BOA. In summary, these simulations reveals the evolution of the net ~~direct~~ radiative ~~forcing~~ effect produced by the cirrus cloud, going at TOA from values close to 0 to ~~+42~~+40 W m^{-2} , at BOA from negative values, whose minimum is -51 W m^{-2} to positive values reaching a maximum of $+20 \text{ W m}^{-2}$ and in the full atmosphere varying between values close to 0 to ~~+48~~+42 W m^{-2} , being the maximum. The complexity of calculating the ~~direct~~ radiative ~~forcing~~ effect of a cirrus cloud lies in the fact that this value is highly sensitive to its scene cloud properties like cloud optical depth, solar zenith angle or surface albedo as seen in this case study.

Finally, we compare the **direct** radiative ~~forcing~~ effect of the cirrus cloud (red star mark) measured by MPL at the Barcelona lidar station, by CALIPSO satellite at 78 km from Barcelona lidar station (Gil-Díaz et al., 2024) and with the properties obtained with CLaMS-Ice, as shown in Table 2.

| Database | Properties | | | | NET DRE (Wm ⁻²) | | |
|-----------|-----------------------|--------------------------------------|---------------------------|------|-----------------------------|------|------|
| | T _{mid} (°C) | σ _{ext} (Km ⁻¹) | IWC (mg m ⁻³) | COD | BOA | TOA | ATM |
| MPL | -60.7 | 0.17 | 4.1 ^a | 0.26 | 17.0 | 35.5 | 18.5 |
| CALIPSO | -63.7 | 0.16 | 22.2 | 0.23 | 15.8 | 32.9 | 17.0 |
| CLaMS-Ice | -61.4 | b | 4.3 | 0.30 | 19.6 | 41.1 | 21.5 |

Table 2. Properties (mid-cloud temperature (T_{mid}), column effective extinction coefficient (σ_{ext}) and ice water content (IWC)) and net direct radiative *forcing* effect of the cirrus cloud measured by Micro Pulse Lidar at the Barcelona lidar station, by CALIPSO satellite (Gil-Díaz et al., 2024) and with the properties obtained with CLaMS-Ice, at bottom-of-the-atmosphere (BOA), top-of-the-atmosphere (TOA) and in the full atmosphere (ATM). ^aThe value of ice water content is calculated with the self-consistent scattering model for cirrus cloud (see Sec. 3.1). ^b The CLaMS-Ice product does not provide a value of the *effective column extinction* of the cirrus cloud.

In Table 2 one observes that the properties of the cirrus cloud are similar, except for the IWC, which is considerably higher for the cirrus characterized with CALIPSO measurements. If the IWC is calculated using the extinction coefficient and mid-cloud temperature values provided by CALIPSO, the ensemble scattering model for cirrus clouds yields an IWC value of 3.9 mg/m³. This result is significantly lower than the IWC measured by CALIPSO, since the ~~ensemble~~ self-consistent scattering model for cirrus clouds often underestimates the IWC of cirrus clouds measured at mid-latitude (Baran and Labonnote, 2007). So a certain discordance between the net ~~direct~~ radiative ~~forcing~~ effect magnitudes would be expected. For all three simulations, the cirrus cloud ~~with~~ has a positive direct radiative effect both at TOA and at BOA, since due to the time (2 UTC) the solar radiation component is null. Moreover, as the ~~forcing~~ direct radiative effect is proportional to the COD, the cirrus characterized with the CLaMS-Ice products produces a slightly higher ~~forcing~~ direct radiative effect than with the other data. On average, it can be established that the cirrus cloud measured at Barcelona lidar station produces a ~~with~~ positive direct radiative effect at BOA of ~~19.19±2.08~~ +17.5±1.9 Wm⁻², at TOA of ~~19.07±5.71~~ +36.5±4.2 Wm⁻² and in the full atmosphere of ~~19.88±9.63~~ +19.0±2.3 Wm⁻².

5 Conclusions

610 In this paper a study of ~~radiative~~ optical scattering properties and forcings of cirrus clouds based on 4 years of continuous ground-based lidar measurements with the Barcelona (Spain) Micro Pulse Lidar (MPL) is analyzed. First, a new approach of a self-consistent scattering model for cirrus clouds is presented to get the ~~radiative~~ optical scattering properties of cirrus clouds at different wavelengths with only the ~~radiative~~ extinction coefficient calculated with the two-way transmittance method and ~~radiative~~ cloud temperature, from radiosounding data. The self-consistent scattering model for cirrus clouds consists of an ensemble of six ice crystal ~~members~~ types, where the simplest ~~ice crystal~~ ~~hexagonal~~ form is an ice column with an **unity** aspect ratio ~~unity~~ and the more complex ice crystals are formed by ~~hexagonal~~ randomly ~~attached~~ attaching other hexagonal elements, until a chain-like ~~hexagonal~~ structure is formed. The members ~~of the ensemble~~ are **evenly** distributed ~~in the size distribution (PSD)~~ across the PSD. The self-consistent scattering model for cirrus cloud database consists of more than 20000 PSDs of tropical and mid-latitude cirrus at temperatures between -60 and 0°C. This 615 database provides for each PSD the simulated optical scattering properties ~~of the ensemble~~ ~~of the hexagonal ice crystal~~ (scattering, absorption coefficient and asymmetry factor) in function of the decimal logarithm of ice water content (LIWC) and the cloud temperature (T) by nonlinear least squares fitting. ~~of the model~~ The new approach ~~of the model~~ consists of first calculating the IWC of the cirrus cloud by introducing the extinction coefficient of each cloud layer in an equation derived from the 620 model, for a cloud temperature. ~~of the model~~ This equation is obtained by assuming the absence the absorption, which is entirely reasonable because 625

the wavelength used belongs to the visible spectral range. Once, the IWC is estimated, this variable is introduced again into the model to get the absorption, scattering and asymmetry factor coefficients for each wavelength, respectively. Applying this 630 method to cirrus clouds measured in Barcelona during November 2018 to September 2022 at 00 and 12 UTC, it is obtained that the average of the ice water content is $4.97 \pm 5.53 \text{ mg/m}^3$, the single scattering albedo is 0.99 ± 0.05 and the asymmetry factor is 0.76 ± 0.00 at $0.55 \mu\text{m}$. Second, the **direct radiative ~~filtering~~ effect** of cirrus clouds is calculated with the radiative transfer model DISORT. Radiative fluxes are validated at bottom-of-the-atmosphere with SolRad-Net pyranometers in the shortwave spectral range, and at top-of-the-atmosphere with **CLERES NOAA-20** measurements in the longwave spectral range. One one hand, 635 most downward radiative fluxes calculated with the DISORT model overestimate the SolRad-Net observations, resulting in a BIAS of $+11.6\%$ and a slope of the linear regression ($a = 0.71 \pm 0.20$). On the other hand, a large difference in upward radiative fluxes between simulated and observations is found for each cirrus cloud scene, resulting in a BIAS of ~~$+33.6\%$~~ . Third, a validation of the importance of the planetary boundary layer aerosols in the cirrus scenes simulations is carried out. Calculations with and without aerosols of the cirrus direct **direct radiative ~~filtering~~ effects** are made to assess the error induced by neglecting 640 tropospheric aerosols, which results in a negligible BIAS of ~~$-0.7 \pm 1.2\%$~~ . In the other case where the aerosols are vertically closer to the clouds, the simplification of not considering aerosols in the calculation of cloud **~~filtering~~ direct radiative effects** may not be valid, leading to a significant underestimation of cloud **~~filtering~~ direct radiative effects**. Forth, the **direct radiative ~~filtering~~ effects** of cirrus clouds are calculated distinguishing between nighttime and daytime. At nighttime, cirrus clouds warm the atmosphere with **direct radiative effects** at TOA almost double than at BOA, with the thicker cirrus clouds contributing most 645 ~~to the filtering~~. At daytime, cirrus clouds **have** generally a **negative direct radiative effect at ~~both~~ BOA** (~~80~~ 82% of the cases) and always ~~warm the~~ a **positive effect at TOA**, resulting in a warming of the full atmosphere. On average, at nighttime, cirrus clouds ~~filter~~ have a positive direct radiative effect of ~~$+23.02 \pm 22.23$~~ $+22.1 \pm 20.8 \text{ W m}^{-2}$ at BOA, ~~$+50.08 \pm 42.99$~~ $+40.4 \pm 34.5 \text{ W m}^{-2}$ at TOA and in the full atmosphere ~~$+27.04 \pm 22.39$~~ $+18.4 \pm 15.3 \text{ W m}^{-2}$; at daytime, cirrus clouds ~~filter~~ have a negative direct radiative effect ~~$+8.57 \pm 4.95$~~ $-11.5 \pm 18.7 \text{ W m}^{-2}$ at BOA and ~~warm~~ a **positive effect** ~~$+18.94 \pm 16.95$~~ $+14.2 \pm 14.0 \text{ W m}^{-2}$ 650 at TOA, and in the full atmosphere ~~$+27.51 \pm 26.68$~~ $+25.7 \pm 25.9 \text{ W m}^{-2}$. Fifth, the variation of the cirrus cloud **direct radiative ~~filtering~~ effect** at daytime is also analyzed as a function of the SZA: it shows that at TOA for the real cases the average net **direct radiative ~~filtering~~ effect** is always positive and for the results from a sensitivity study, the mean net **direct radiative ~~filtering~~ effect** shifts from positive to negative values at 58° . ~~For the real cases and the sensitivity study~~, a slight downward tendency of the net **direct radiative ~~filtering~~ effect** is also found. At BOA, most net **direct radiative ~~filtering~~ effect** values are 655 negative. The change from positive to negative values of the mean net **direct radiative ~~filtering~~ effect** is only observed for the real cases, being SZA crossover of 25° . Sixth, for a case study, the **direct radiative ~~filtering~~ effect** of a cirrus cloud along its back-trajectory is analyzed using CLaMS-Ice products. During the overlap between the back-trajectory of this cirrus cloud and the CALIPSO satellite on 11/02/2019 02:03:20 UTC, a brief validation of the IWC is made, resulting in a relative ~~filter~~ deviation of 29% for the CLaMS-Ice output. This cirrus cloud comes from the Atlantic Ocean, passing through part of France 660 and Barcelona to reach the Mediterranean Sea. Over the Atlantic Ocean, the air mass which corresponds to the cloud undergoes a rise in height, reaching a minimum in temperature and a COD of 0.6. After this initial cirrus cloud fast growth, the cloud gradually fades away as it approaches France, where the CLaMS-Ice model gives a null IWC. Afterwards, the cirrus cloud

begins to form again on land surface and to grow up. Along its trajectory, the cirrus cloud produces a net **direct** radiative ~~forcing~~ **effect** that goes at TOA from values close to 0 to ~~42~~+40 Wm⁻², at BOA from negative values, whose minimum is -51 Wm⁻² to positive values reaching a maximum of +20 Wm⁻² and in the full atmosphere varying between values close to 0 to ~~48~~+42 Wm⁻², being the maximum. Finally, the **direct** radiative ~~forcing~~ **effects** that the cirrus cloud has at the beginning of the back-trajectory are compared using different measurements (MPL and CALIPSO measurements and CLAMS-Ice outputs) and making use of the self-consistent scattering model for cirrus clouds. It results in an average ~~49.07±5.71~~ **positive direct radiative effect** at BOA of ~~19.09±2.08~~+17.5±1.9 Wm⁻², at TOA of ~~49.07±5.71~~ **+36.5±4.2** Wm⁻² and in the full atmosphere of ~~29.88±8.63~~ **+19.0±2.3** Wm⁻².

Data availability. The MPLNET products are publicly available on the MPLNET website (https://mplnet.gsfc.nasa.gov/download_tool/) (MPLNET, 2024) in accordance with the data policy statement. Radiosoundings data are available upon request from the authors or Meteocat. The SolRad-Net product is publicly available on the SolRad-Net website (<https://solrad-net.gsfc.nasa.gov/>) (Goddard Space Flight Center, 2024a) in accordance with the data policy statement. The NOAA-20 products are publicly available on the CERES website (<https://ceres.larc.nasa.gov/>) (Langley Research Center, 2024). The AERONET products are provided by a federation of ground-based remote sensing aerosol networks established by NASA and PHOTONS (PHOtométrie pour le Traitement Opérationnel de Normalisation Satellitaire) and is greatly expanded by collaborators from national agencies, institutes, universities, individual scientists, and partners. The AERONET products are publicly available on the AERONET website (<https://aeronet.gsfc.nasa.gov/>) (Goddard Space Flight Center, 2024b). The CALIPSO product is provided by the NASA Langley Research Center's (LaRC) ASDC DAAC and is managed by the NASA Earth Science Data and Information System (ESDIS) project. NASA data are freely accessible and available on the Atmospheric Science Data Center website (<https://asdc.larc.nasa.gov/>) (NASA, 2024).

Author contributions. CGD prepared the automatic algorithms to ~~not~~ calculate the ~~radiative~~ optical scattering properties and ~~forcing~~ **direct radiative effects** of cirrus clouds for MPL and radiosounding data. CGD prepared the figures of the paper. MS, OS, AS, CMP, AC, ARG and DCFSO reviewed different parts of the results. CGD prepared the paper, with contributions from all co-authors.

685 *Competing interests.* At least one of the (co-)authors is a member of the editorial board of Atmospheric Chemistry and Physics.

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