



# Steady-State Mixing State of Black Carbon Aerosols from a Particle-Resolved Model

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**Abstract.** Black carbon (BC) exerts a notable warming effect due to its strong light absorption, largely influenced by its "mixing state". However, due to computational constraints, mixing state is challenging to accurately represent in large-scale models. In this study, we employ a particle-resolved model to simulate the evolution of BC mixing state based on field observation in Nanjing. Our result shows that aerosol compositions, coating thickness distribution, and optical properties of  
25 BC aerosols all exhibit a tendency toward steady-state with a characteristic time of less than one day. Using the steady-state simplifying assumption, BC absorption enhancement closely matches the result obtained through the particle-resolved method. Additionally, we discuss how to reconcile our finding of a universal distribution with the diversity in the distribution of BC coating thickness that has been documented in previous studies. This study simplifies the BC mixing state description and yields a precise evaluation of the BC optical properties, which facilitates the refinement of the assessment of BC's  
30 radiative effects in global and chemical transport models.

## 1 Introduction

Black carbon (BC) exerts notable warming effects due to its strong light absorption (Andreae, 2001; Bond et al., 2013; Gustafsson and Ramanathan, 2016; Horvath and Trier, 1993; Jacobson, 2001; McConnell et al., 2007). BC absorption is largely influenced by its "mixing state" (Bond et al., 2013; Cappa et al., 2012; Fierce et al., 2020; Liu et al., 2017, 2015;



35 Riemer et al., 2019), a term that describes to what extent the BC component is mixed with or coated by other aerosol  
components (Bondy et al., 2018; Winkler, 1973). After being emitted, fresh BC can mix with other non-absorbing aerosol  
components, enhancing its light absorption due to the "lensing effect" (Bond and Bergstrom, 2006; Jacobson, 2001; Peng et  
al., 2016; Redemann et al., 2001). This leads to larger mass absorption cross-sections (MAC) compared to freshly-emitted or  
bare BC. Previous laboratory studies have shown that the light absorption enhancement ( $E_{\text{abs}}$ ) of BC aerosols can vary from  
40 1.05 to 3.50 under controlled experimental conditions, depending on BC mixing states (Jacobson, 2001; Ramanathan and  
Carmichael, 2008; Riemer et al., 2019; Schnaiter et al., 2005; Shiraiwa et al., 2010). Therefore, accurately characterizing the  
BC mixing state is crucial for the precise assessment of the climate impact of BC aerosols.

The mixing state of BC aerosols in the real atmosphere is dynamically varying due to continuous changes caused by various  
processes such as the emission of particles, condensation of inorganic and organic substances on the BC aerosol, coagulation,  
45 and deposition (Moteki et al., 2007; Shiraiwa et al., 2007). Some studies have attempted to represent the BC mixing state  
using simplified aerosol representations, such as assuming a priori internal or external mixtures of BC aerosols (He et al.,  
2015; Lesins et al., 2002). However, these approaches disregard the continuous changes of the BC mixing state. Other  
studies have used a fixed rate to describe the aging of BC aerosols from fresh BC to aged BC (Qi et al., 2017b, a), but this  
approach also fails to fully capture the continuous changes.

50 The particle-resolved Monte Carlo & Model for Simulating Aerosol Interactions and Chemistry (PartMC-MOSAIC) offers a  
more detailed simulation of the continuous changes in the BC mixing state (Riemer et al., 2009; Zaveri et al., 2008). This  
model possesses the capability to track the state of individual aerosol particles during the simulation process (Ching et al.,  
2018; Hughes et al., 2018; Riemer et al., 2009, 2019; Shou et al., 2019; Zaveri et al., 2010; Zheng et al., 2021). However, the  
direct application of particle-resolved models within large-scale chemistry transport models or Earth system models is  
55 impeded due to the high computational cost required for their high-precision particle simulation. To overcome this challenge,  
the approximation of exponential distribution of coating thickness (CT) under BC steady-state mixing state has been recently  
proposed. The study of Wang et al. (2023) suggests that particles emitted into the atmosphere undergo continuous growth  
and deposition processes, leading to the formation of a steady state, where the CT distribution of BC aerosols across  
different sizes of BC cores does not change over time. This approach has the promise to simplify the representation of the  
60 continuously changing BC mixing state in models. However, the theory of BC steady-state mixing state still requires  
numerical verification and key scientific questions remain such as determining the characteristic timescale to reach a steady  
state. Additionally, the study of Fierce et al. (2016) has noted that the coating thickness of BC aerosols is non-uniform across  
the distributions of BC cores, which seems to contradict the conclusion of uniform CT distribution of BC aerosols in the  
steady-state theory. This apparent discrepancy has not yet been fully explained.

65 In this study, we determine whether BC aerosols' mixing state reaches a steady state while the population is exposed to  
continuous emission, growth, and deposition processes using the PartMC-MOSAIC model. The set of baseline case was  
based on field observations in Nanjing, China, from April 2020. We also examine the characteristic timescale of BC aerosols  
approaching a steady-state mixing state, the resulting CT distribution of BC aerosols, and the optical properties of BC under



steady state. The steady-state CT distribution of BC aerosols will be compared with the results obtained from field  
70 observation in Nanjing, China to verify the realism of the simulation. Moreover, we explain the discrepancy between the  
non-uniform composition found in previous studies and the uniform CT distribution of BC aerosols confirmed in this study.  
Our study aims to provide a simplified and reasonably precise approach for the characterization of BC mixing state and  
evaluation of the light absorption enhancement in global models or chemical transport models.

## 2 Materials and Methods

### 75 2.1 Description of observational data

Observational data of BC mixing states was collected from the Station for Observing Regional Processes of the Earth system  
(SORPES, 118°57'10" E, 32°07'14" N; 40 m a.s.l.) in Nanjing (Ding et al., 2016). Observational data from April 2020 was  
used in this study. The physical properties of BC aerosols were measured by single particle soot photometers (SP2, Droplet  
Measurement Technologies, USA). The working principle of the SP2 has been thoroughly described in earlier studies (Gao  
80 et al., 2007; Schwarz et al., 2006). The BC mass can be determined based on its proportional relationship with the peak  
intensity of the emitted light and the BC mass equivalent diameter can be calculated. The scattering cross-section of BC  
particles is calculated using the leading-edge-only fit method developed by Gao et al. (2007). Temperature and atmospheric  
pressure were continuously measured by the Nanjing automatic weather station throughout the year 2020, with a time  
resolution of 30 minutes.

### 85 2.2 The particle-resolved model PartMC-MOSAIC

In this study, we utilized the PartMC-MOSAIC model to verify the process of BC aerosols approaching a steady-state  
mixing state. This model simulates the time evolution of the BC mixing state by tracking the composition of individual  
computational particles during atmospheric processes (Ching et al., 2018; Shou et al., 2019; Zaveri et al., 2008, 2010).  
Particle size information can be calculated from the composition of individual particles assuming spherical particles.  
90 Subsequently, the BC mixing state can be derived from the aforementioned details. We utilized PartMC-MOSAIC  
simulations to obtain a precise representation of the BC mixing state, forming the numerical basis for our steady-state  
analysis.

To investigate the evolution of the mixing state of BC aerosols in typical regions under continuous emissions and diverse  
atmospheric processes, we constructed a scenario that simulates an urban area using approximately 10,000 computational  
95 particles (the precise values change throughout the simulation) for ten days. The temperature and atmospheric pressure in the  
scenario were set to constant values consistent with the monthly average values measured by the Nanjing automatic weather  
station during the period from 1 April 2020, to 30 April 2020. In this study, the initial gas concentration and emission rate  
were slightly adjusted based on Riemer et al. (2009) according to Wang et al. (2017). Since time-varying emissions do not  
affect the average behaviour (refer to the discussion in the Supporting Information of Wang et al. (2023)), for simplicity, the



100 emissions of gases and particles have been set to constant values in the baseline case. We prescribe three different types of  
carbonaceous aerosol emissions throughout the simulation representing a mix of urban sources (Cao et al., 2006; Schwarz et  
al., 2008b), namely pure BC particles (24.6% of the emissions by number), pure organic carbon (OC) particles (72.7% of the  
emissions by number), and mixed BC/OC particles (2.7% of the emissions by number), assuming that their size distributions  
are log-normal. The pure BC and the mixed BC/OC particles were assigned a BC-core size distribution with 89 nm  
105 geometric mean diameter and a geometric standard deviation of 1.6 following the results obtained by the SP2 instrument in  
Nanjing from April 1st to April 30th, 2020 (Fig. S1). The mass proportions of OC and BC are set to 31.7% and 68.3%, to  
ensure that particles with the BC core of 89 nm have the OC coating layer with a thickness of 20 nm (Schwarz et al., 2008b).  
The distribution of pure OC particle emission was assigned a geometric mean diameter of 110 nm and a geometric standard  
deviation of 1.7 (Fierce et al., 2016). The above settings ensure a mass ratio of 1499:4241 between component BC and  
110 component OC (Cao et al., 2006), while also maintaining a number concentration ratio of approximately 9:1 between pure  
BC particles and mixed OC/BC particles (Schwarz et al., 2008b). The detailed setup of the baseline case is provided in Table  
S1 in the Supplement.

After emission, the combustion particles underwent further coating through gas-particle partitioning, and coagulation due to  
Browning motion, and particles were removed by deposition, resulting in a complex mixing state. To mimic actual  
115 atmospheric conditions, we set a fixed emission rate of BC particles and maintained atmospheric processes throughout the  
simulation. The deposition rate of particles was set at a fixed rate and gas-particle partitioning was determined by the  
concentration of particles and gaseous species. While the simulation includes the full simulation of gas phase chemistry, gas-  
particle partitioning, aerosol thermodynamics, and the properties of BC aerosols undergoing continuous changes (Riemer et  
al., 2019), our study only focused on the mixture of aerosol species and the CT distribution of BC aerosols. The simulated  
120 aerosol component fractions are consistent with that in the ambient environment of China, provided in SI (Table S2). To  
illustrate the generality of the steady state, we have created nine additional scenarios by varying the temperature, the size  
distribution of BC-core in the emissions (Bond et al., 2013; Lee, 1983; Moteki, 2023), the growth rate in the simulated area  
(by increasing or decreasing particle and gas emissions by a factor of five), and diurnal gas emissions condition, which are  
provided in Sect. 1.2 in the Supplement.

### 125 **2.3 Mixing state indices**

By storing the composition of individual computational particles (i.e., the full aerosol state), PartMC-MOSAIC output  
contains the full information on BC mixing state. However the output is exceedingly complex, requiring about  $10^5$  (particles  
 $\times$  species) of values to describe the aerosol state at a specific time, thereby posing challenges in comparing BC mixing states  
at different times (Riemer et al., 2019). To overcome this, we conducted dimensionality reduction on the mixing state data  
130 by projecting the full aerosol state onto one or multiple values and assessing the mixing state through indices. This study  
employed the indices  $\chi$  (Riemer and West, 2013) and  $k$  (Wang et al., 2023) to quantify the BC mixing state and track the  
process of BC approaching a steady-state mixing state throughout the simulation. The  $\chi$  index represents the aerosol mixing



state based on diversity metrics, whereas the  $k$  index originates from the particle size distribution. The combination of these two indices enables a comprehensive depiction of the BC mixing state.

135 The diversity index  $\chi$  utilized in this study was proposed by Riemer and West (2013) to characterize the aerosol mixing state of a particle population. This index is determined by the mass fractions of each species in each particle, ranging from 0% to 100% to represent fully external to internal mixtures. In this paper, we focus on the mixing of BC and other components. So, species are separated into two surrogate species, BC and non-BC in the calculation of index  $\chi$  (Zhao et al., 2021). PartMC-MOSAIC output can provide the mass of species  $a$  in particle  $i$  represented as  $\mu_i^a$ , for  $i$  ranging from 1 to  $N$  and  $a$  taking  
 140 values of 1 or 2, where 1 denotes BC and 2 denotes non-BC. From this basic description of the aerosol particles, we can construct all associated quantities, as described by Riemer and West (2013) and listed in Table 1.

**Table 1. The related quantities calculated by the mass of components per particle (Riemer and West, 2013)**

Quantity	Meaning
$\mu_i^a$	Mass of species $a$ in particle $i$
$\mu_i = \sum_{a=1}^2 \mu_i^a$	The total mass of particle $i$
$\mu^a = \sum_{i=1}^N \mu_i^a$	The total mass of species $a$ in the population
$\mu = \sum_{i=1}^N \mu_i$	Total mass of population
$p_i^a = \frac{\mu_i^a}{\mu_i}$	Mass fraction of species $a$ in particle $i$
$p_i = \frac{\mu_i}{\mu}$	Mass fraction of particle $i$ in population
$p^a = \frac{\mu^a}{\mu}$	Mass fraction of species $a$ in the population

145 By utilizing the distribution of different aerosol components both within individual aerosol particles and across the entire population, we can now calculate the mixing entropies, species diversities, and the mixing state index, as presented in Table 2. The index  $\chi$  in this paper can describe the mixing of BC and other components. The variation of index  $\chi$  is instrumental in more effectively tracking the process by which BC-core is coated with other non-BC components.

**Table 2. The computation of diversity metrics (Riemer and West, 2013)**

Quantity	Name	Meaning
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$H_i = \sum_{a=1}^2 -p_i^a \cdot \ln p_i^a$	Mixing entropy of particle $i$	Shannon entropy of species distribution within particle $i$
$H_a = \sum_{i=1}^N p_i \cdot H_i$	Average particle mixing entropy	Shannon entropy of species distribution within particle $i$
$H_\gamma = \sum_{a=1}^2 -p^a \cdot \ln p^a$	Population bulk mixing entropy	Shannon entropy of species distribution within the population
$D_i = e^{H_i} = \prod_{a=1}^2 (p_i^a)^{-p_i^a}$	Particle diversity of particle $i$	effective number of species in particle $i$
$D_a = e^{H_a} = \prod_{i=1}^N (D_i)^{p_i}$	Average particle (alpha) species $i$ diversity	average effective number of species in each particle
$D_\gamma = e^{H_\gamma} = \prod_{a=1}^2 (p^a)^{-p^a}$	Bulk population (gamma) species diversity	effective number of species in the bulk
$\chi = \frac{D_a - 1}{D_\gamma - 1}$	Mixing state index	degree to which population is externally mixed ( $\chi = 0\%$ ) versus internally mixed ( $\chi = 100\%$ )

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150 In this study, we also used the index  $k$  adopted from our previous study (Wang et al., 2023) to quantify the CT distribution of BC aerosols during the simulation. Unlike the index  $\chi$ , index  $k$  emphasizes the thickness of the BC coating layer and quantifies the mixing state of BC aerosols based on the CT distribution. The universal law governing the mixing state of BC aerosols reveals a consistent size distribution, indicating that the natural logarithm of the number concentration ( $\ln(n(\text{CT}))$ ) and CT exhibit a linear relationship, regardless of the size of the BC core (Wang et al., 2023). The variable  $D_c$  denotes the

155 diameter of the BC core and the variable  $D_p$  denotes the diameter of the particle, CT signifies the coating thickness of the particle, which is defined as  $D_p - D_c$ . Furthermore,  $n(\text{CT})$  represents the normalized number concentration of BC particles within the range of  $(\text{CT} - \Delta\text{CT}/2, \text{CT} + \Delta\text{CT}/2]$ . The normalization is relative to the total number concentration of BC particles, rendering  $n(\text{CT})$  a dimensionless value. Then, we employed the statistical technique proposed by Wang et al. (2023) in this investigation to partition the particle data based on particle size. By performing linear regression, the relationship

160 between the CT and the corresponding number concentration in logarithmic coordinates,  $\ln(n(\text{CT}))$ , was established and the slope  $k$  was calculated (Sect. 2.1 in the Supplement). The slope  $k$  ( $k = \text{Dep}/\text{GR}$ ) is determined by the deposition rate (Dep), or more generally, the removal rate, and the growth rate (GR) (Wang et al., 2023). The slope  $k$  was subsequently adopted as a characterization parameter to assess the BC mixing state, focusing on the particle size distribution. Per-particle size for



165 statistical analysis was calculated using the component density from PartMC-MOSAIC output (Sect. 2.2 and Table S4 in the Supplement) combined with the core-shell model (R. McGrory et al., 2022; Wang et al., 2019).

## 2.4 Light absorption enhancement of BC aerosols

The optical properties of BC aerosols were characterized using light absorption enhancement in this study, which is defined as the ratio of the MAC of coated BC to bare BC. The calculations of  $E_{\text{abs}}$  were performed using the core-shell Mie method (Fu and Sun, 2001; Schwarz et al., 2008a; Toon and Ackerman, 1981). In this study, the wavelength was set to 550 nm and the refractive indices (RI) of the BC and scattering components were set to  $1.95 + 0.96i$  (Moteki et al., 2023) and  $1.5 + 0i$  (Schwarz et al., 2015), respectively. Furthermore, the determination of  $E_{\text{abs}}$  is undertaken through two methods (per-particle method and  $k$ -value method).

The per-particle method of computation entails the utilization of core-shell Mie theory to determine the optical properties of individual particles. Subsequently, statistical techniques were applied to derive the  $E_{\text{abs}}$  of the entire BC population.  $E_{\text{abs}}$  could be represented as Eq. (1).

$$E_{\text{abs}} = \frac{\sum_{i=1}^{\text{total number of computational particles}} c_{\text{abs-internal}}(i) \cdot n_{\text{conc}}(i)}{\sum_{i=1}^{\text{total number of computational particles}} c_{\text{abs-external}}(i) \cdot n_{\text{conc}}(i)} \quad (1)$$

where  $c_{\text{abs-internal}}$  represents the light absorption coefficients of BC core with coating shell, and  $c_{\text{abs-external}}$  represents the light absorption coefficients of BC core. Both the  $c_{\text{abs-internal}}$  and the  $c_{\text{abs-external}}$  can be calculated directly through the core-shell Mie method with the given  $D_c$  and  $D_p$  of each particle  $i$  simulated by the PartMC-MOSAIC model.

180 The  $k$ -value method of calculating  $E_{\text{abs}}$  uses simplified  $D_c$  and  $D_p$  distributions, which obviates the need for individual particle calculations and greatly reduces the computational burden. The calculation is predicated on the provided probability density functions of  $D_c$  and equivalent  $D_p$ . The log-normal distribution of  $D_c$  is set according to the initial emissions specified in the PartMC-MOSAIC model, with a mean diameter of 89 nm and a geometric standard deviation of 1.6 (Fig. S1). The probability density function of equivalent  $D_p$  is determined by incorporating the mean coating thickness (Wang et al., 2023), ascertained from the  $k$  value, into the probability density function of  $D_c$  ( $D_p = D_c + 1/k$ ). More detailed information and calculations about the  $k$ -value method were illustrated in the study of Wang et al. (2023).

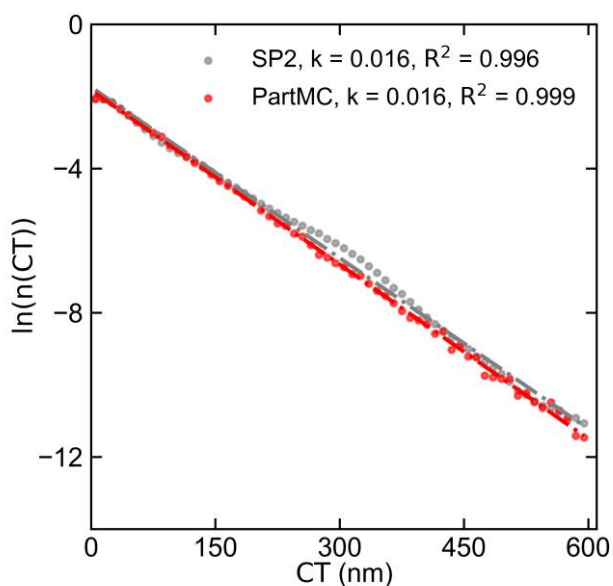
The per-particle method involves computing the  $E_{\text{abs}}$  of the BC population by utilizing the MAC of each particle. The  $k$ -value method of calculating  $E_{\text{abs}}$  uses simplified  $D_c$  and  $D_p$  distributions obtained from the steady CT distribution, which obviates the need for individual particle calculations and greatly reduces the computational burden. Since the per-particle method demonstrates high precision in determining  $E_{\text{abs}}$ , the accuracy of the  $k$ -value method in calculating  $E_{\text{abs}}$  can be evaluated by comparing it with the per-particle method.



### 3 Results

To verify the realism of the baseline case simulated by the PartMC-MOSAIC model, we compare representative results of aerosol component mass fractions and CT distributions obtained from SP2 observations and model simulations. The representative results of aerosol component mass fractions from observations come from Wang et al. (2017), including mass concentrations of organic matter (OM), black carbon (BC), nitrate ion (NO<sub>3</sub>), sulfate ion (SO<sub>4</sub>), and ammonium ion (NH<sub>4</sub>). The comparison is given in Table S2, showing that the simulated results are within the range of the observations (Wang et al., 2017). As the simulation does not involve crustal materials such as Ca and Na, the result that the positive ion NH<sub>4</sub> is slightly beyond the range is reasonable. Since there are no sulfate-containing particles emitted and the sulfate production by in-cloud chemistry is not included, the primary source of SO<sub>4</sub> is the conversion of SO<sub>2</sub> gas to particles. As a result, the proportion of SO<sub>4</sub> mass is relatively small, while NO<sub>3</sub> levels are relatively high.

Figure 1 shows that the CT distribution derived from the PartMC-MOSAIC simulation closely matches that from SP2 observations, with both results following an exponential linear distribution (linear distribution on a logarithmic scale) and yielding the same slope value of  $\sim 0.016 \text{ nm}^{-1}$  for the fitting curve, which confirms the realism of the scenario simulated by PartMC-MOSAIC model.



**Figure 1.** The CT distribution of black carbon aerosols, determined from SP2 observation results (grey dots and line) and PartMC-MOSAIC simulation based on the observed scenario (red dots and line). Each dot indicates the total number of particles that are in the CT interval and the linear regression of each distribution is represented by dashed lines.

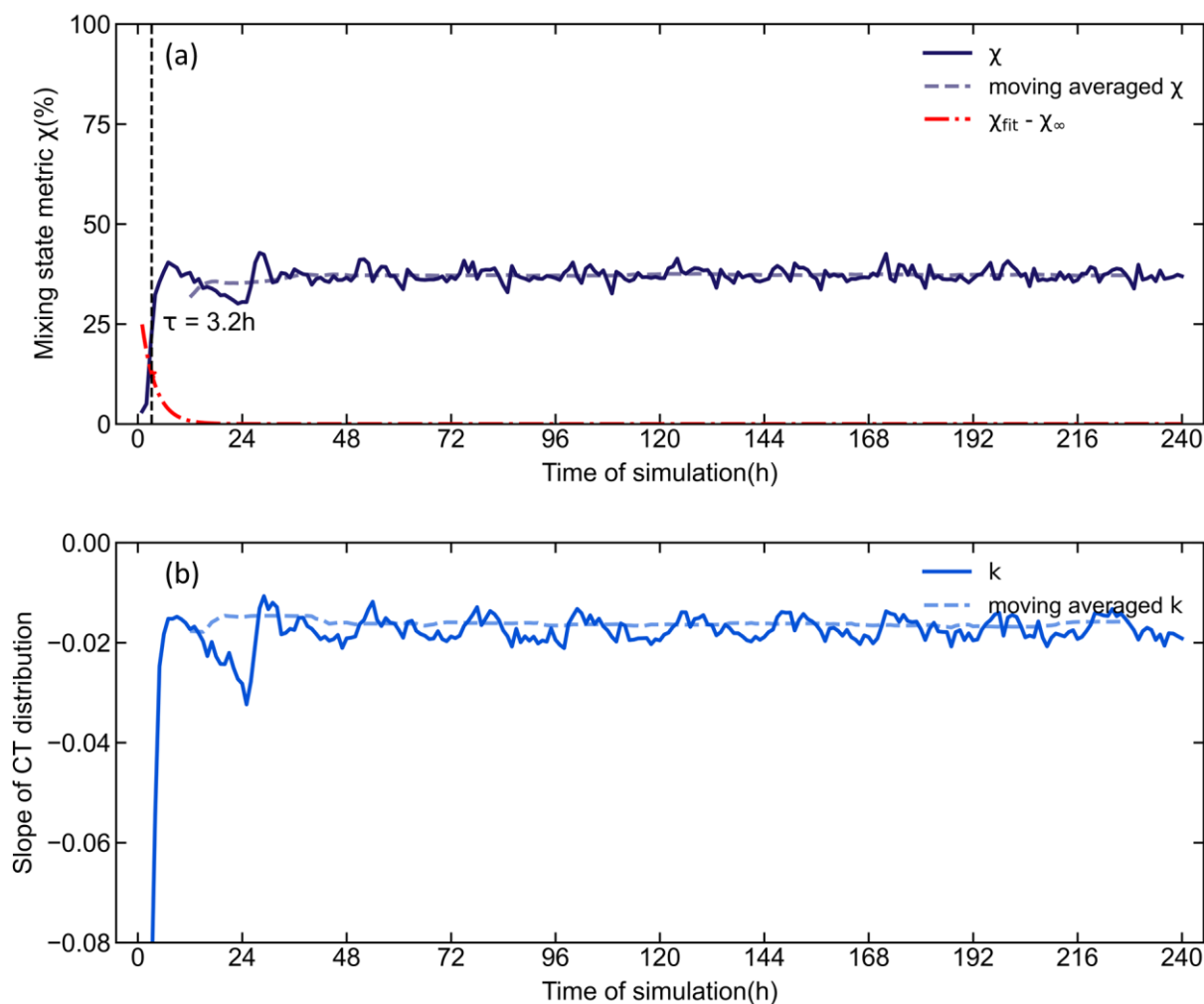




After the validation of the baseline case, we first explore whether a steady-state mixing state of BC aerosol can be achieved. Figure 2 illustrates the evolution of metrics  $\chi$  and  $k$  during the simulation. We discover that the mixing state metrics  $\chi$  and  $k$  tend to reach a steady state, with slight daily variations (Fig. 2a and 2b) due to periodic external drivers. Thus, we adopt a moving average algorithm (Hansun, 2013) over 24 hours (Sect. 2.3 in the Supplement) to eliminate such periodic influences on the mixing state. Then, we use Eq. (2) to describe the progress of the mixing state metrics  $\chi$  reaching a steady state.

$$\chi(t) = \chi(0) \cdot e^{-\frac{t}{\tau}} + \chi(\infty) \cdot (1 - e^{-\frac{t}{\tau}}) \quad (2)$$

where the characteristic time  $\tau$  for reaching a steady mixing state of the BC aerosols from the perspective of the  $\chi$  criterion is investigated. We perform nonlinear fitting of the index  $\chi$  according to Eq. (2), with  $\chi(0)$  given by the  $\chi$  value at the start of the simulation, and  $\chi(\infty)$  given by the average value of the last 48 hours simulated. As shown in Fig. 2a, the characteristic time of the baseline case is 3.2 hours, which corresponds to the rate of the mixing state approaching stability. Besides, we construct more cases by changing the settings of temperature and emissions of gas and particles based on previous relevant investigations (Bond et al., 2013; Lee, 1983; Moteki, 2023). The characteristic time scales range from 1.9 to 9.7 hours (Table S3) for reaching a steady mixing state of the BC aerosols. Figure 2b shows the exponential linear fit slope  $k$  of the CT distribution of BC aerosols. After 48 hours, the absolute value of the slope  $k$  becomes steady under the periodic influence, and the value after moving average is steady at  $0.016 \text{ nm}^{-1}$ . The coefficient of determination between  $\ln(n(\text{CT}))$  and CT maintains a high value ( $R^2 > 0.9$ ) during the modeling simulation (Fig. S2), which demonstrates the stability of CT distribution over time.



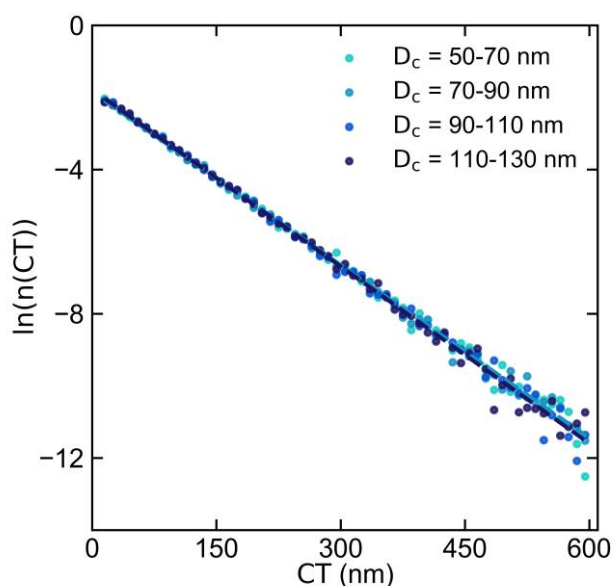
230 **Figure 2. The time evolution of the black carbon mixing state under the combined influence of continuous emissions and various atmospheric processes simulated by PartMC-MOSAIC during the 10-day simulation. (a) the time evolution of index- $\chi$  (solid line), index- $\chi$  after the simple moving average algorithm (dashed line), and  $\chi_{fit} - \chi_{\infty}$  (red line). (b) The time evolution of index- $k$  (solid line) and index- $k$  after the simple moving average algorithm (dashed line).**

235 Further, we analyze the CT distribution of all BC aerosols during the steady-state period (after 48 hours). Figure 1 and Fig. 3 provide the CT distribution under the steady state simulated by PartMC-MOSAIC. Figure 1 shows that the CT distribution of BC aerosols follows an exponential linear distribution. The coefficient of determination ( $R^2$ ) between CT and  $\ln(n(CT))$  is 0.999, indicating a strong linear relationship between them. The slope  $k$  of the linear fitting is  $\sim 0.016 \text{ nm}^{-1}$ . Equivalent to the reciprocal of  $k$ , the mean coating thickness is  $\sim 62 \text{ nm}$  (Wang et al., 2023), which is only  $\sim 1 \text{ nm}$  lower than the value calculated by the per-particle method ( $\sim 63 \text{ nm}$ ). The results of CT distribution and comparison are listed in the Table S3 and



240 Fig. S3. Figure 3 shows the CT distribution results of BC aerosols, classified by the size of the BC core ( $D_c$ ). Four  $D_c$  bins (50 to 70 nm, 70 to 90 nm, 90 to 110 nm, and 100 to 130 nm) were selected, accounting for ~68% of the total number concentration of BC aerosols. The CT distributions with different  $D_c$  have approximately the same slope of  $\ln(n(\text{CT}))$  and CT, with comparable  $k$  values ranging from ~0.0160 to ~0.0164  $\text{nm}^{-1}$  (Fig. 3). The above results verify that the CT distribution of BC aerosols under the steady-state condition follows the exponential linear distribution and exhibits

245 characteristics that are not dependent on  $D_c$  (Fig. 3). The exponential CT distribution of BC aerosols can effectively reflect the distribution pattern of the coating layer and is more precise than simple assumptions such as the fully internal mixture of BC aerosols.



250 **Figure 3. The statistical CT distribution of BC aerosols during the steady-state period (after 48 hours) simulated by PartMC-MOSAIC. The CT distributions of four selected  $D_c$  ranges follow the exponential linear distribution law, represented by dots with different colors. Each dot indicates the total number of particles that are in the CT interval. The linear regression of each distribution is represented by dashed lines, with  $n(\text{CT})$  normalized for each  $D_c$  bin.**

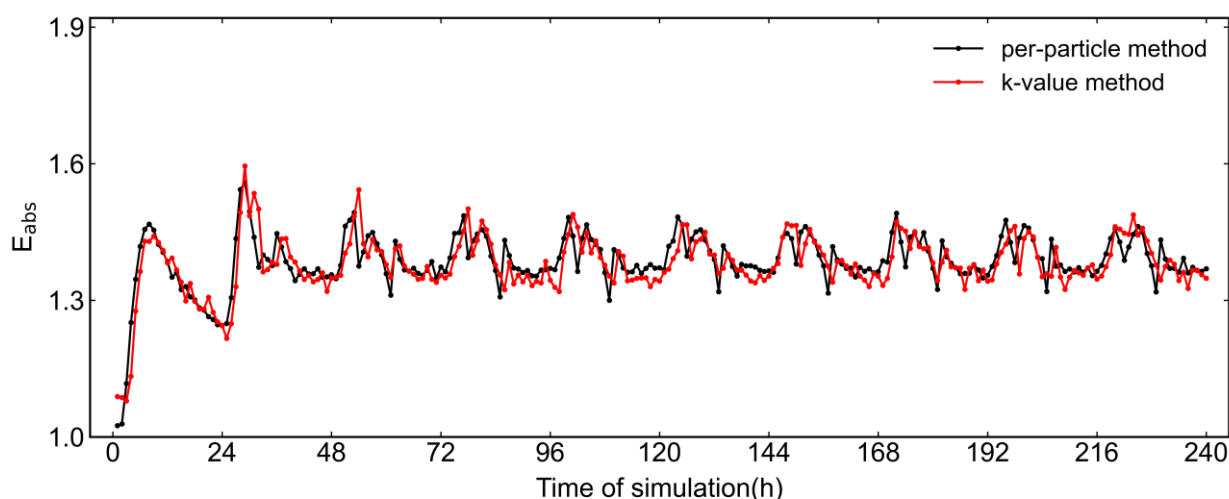
Based on the above results, we hypothesize that the exponential linear CT distribution of BC aerosols can serve as an

255 accurate, simplified characterization of the BC mixing state under the steady-state condition in optical calculations. To verify this hypothesis, we follow the evolution of  $E_{\text{abs}}$  calculated by utilizing the  $k$ -value method and compare it with the real result obtained from the per-particle method, as illustrated in Fig. 4. Through this comparison, we observe a minimal deviation between the  $k$ -value results and the true values. Specifically, the Mean Squared Error (MSE) is 0.001, the Root Mean Squared Error (RMSE) is 0.039, and the Mean Absolute Error (MAE) is 0.029. These metrics indicate that the time

260 evolution of  $E_{\text{abs}}$  from both methods shows a high degree of consistency, which verifies the viability of the method that



utilizes the CT distribution to obtain the BC light absorption amplification. Additionally, Fig. 4 shows that the  $E_{\text{abs}}$  reaches a steady state after ~48 hours, fluctuating regularly in the range of 1.4 to 1.5, which is consistent with observations ranging from 1.0 to 1.7 (Cappa et al., 2012, 2019; Knox et al., 2009; Lack et al., 2012; Liu et al., 2015; Ma et al., 2020; Ueda et al., 2016). Overall, the light absorption amplification of BC aerosols remains steady under the steady-state condition and the  $k$ -value method based on the steady-state theory can be used effectively as a simplified mixing state scheme for optical calculations of BC aerosols.



270 **Figure 4. The time evolution of BC light absorption enhancement ( $E_{\text{abs}}$ ) calculated by the per-particle method and  $k$ -value method during the simulation for the baseline case. The black lines and red lines represent the  $E_{\text{abs}}$  calculated by the per-particle method and  $k$ -value method, respectively.**

#### 4 Discussion

The mixing state of BC aerosols is very complex, including various chemical compositions, particle size distribution, and morphology (Riemer et al., 2019). Different statistical methods can reveal different facets of the BC mixing state. Fierce et al. (2017) found that across the entire diverse BC population, most of the coating material is contained in particles that have small amounts of BC mass, while most of the BC mass resides in particles with large BC inclusions and thin coatings. The conclusion showed that the CT distribution across the population of BC cores was non-uniform, which seems to contradict the exponential linear distribution of CT in this paper. Here we explain how to understand that both viewpoints are consistent through Fig. 5, Eq. (3) and (4).

280 The distribution of  $D_c$  follows a log-normal distribution with a mean diameter of 89 nm and a geometric standard deviation of 1.6. The distribution of CT follows an exponential distribution (Fig. 5a), provided by Eq. (3).



$$N(\text{CT}) = b \cdot e^{-k \cdot \text{CT}} \quad (3)$$

285 Where the value of  $b$  is related to the total particle number concentration  $N$  of BC-containing particles, and  $k$  represents the slope of the CT exponential distribution. For statistical convenience, we normalize  $n(\text{CT})$  to make its integral over CT equal to 1, hence the value  $b$  equals  $k$ . Because of the self-similarity of CT distributions (Fig. 3), different  $D_c$  bins share the same CT distribution. Based on the distribution of CT and  $D_c$ , we obtain the volume fraction (VF) distribution of the BC component with respect to  $D_c$ , provided by Eq. (4).

$$\text{VF}(D_c) = \frac{\int_0^{\infty} k \cdot D_c^3 \cdot e^{-k \cdot \text{CT}} \cdot d\text{CT}}{\int_0^{\infty} k \cdot (D_c + \text{CT})^3 \cdot e^{-k \cdot \text{CT}} \cdot d\text{CT}} \quad (4)$$

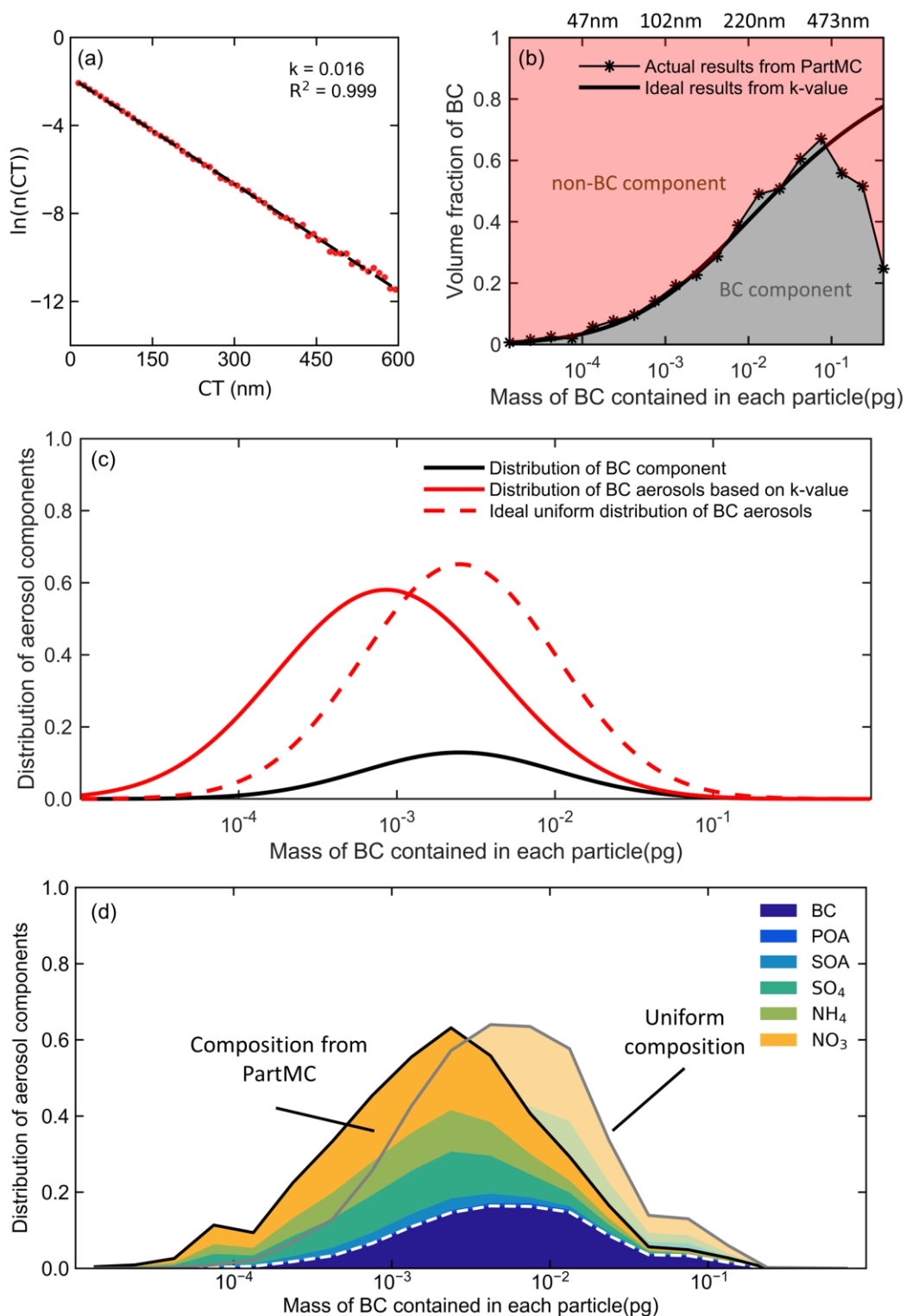
290 The derivative of the equation  $\text{VF}(D_c)$  with respect to  $D_c$ ,  $\frac{d\text{VF}(D_c)}{dD_c}$ , is greater than 0, which indicates that for a fixed CT distribution, VF monotonically increases with  $D_c$  and BC mass of each particle. The ideal results obtained from  $k$ -value (exponential linear distribution of CT) corroborate the actual results from PartMC-MOSAIC simulation (Fig. 5b). Then, we use the VF of BC with respect to mass of BC contained in each particle to obtain the distribution of aerosol components (Fig. 5c), which shows that components account for a smaller proportion in the range of larger BC cores and account for a larger proportion in the range of smaller BC cores, i.e., the volume (or mass) fraction of coating is not uniform across the BC cores.

295 Based on the PartMC-MOSAIC simulation, we calculated the mass distribution of aerosol components directly and compared it with that under the uniform assumption (Fig. 5d). The results show that the distribution of aerosol components with respect to the mass of BC contained in each particle is non-uniform, and it is highly consistent with that obtained from the  $k$ -value in Fig. 5c, which confirms that the aerosol composition results derived from the exponential linear distribution of

300 CT are indeed consistent with the non-uniformity result (Fig. 5d).

The derivation mentioned above demonstrates that the exponential CT distribution pattern of BC aerosols proves the non-uniformity of the composition of BC aerosols from another perspective and can serve as a suitable statistical method for characterizing the BC mixing state.

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**Figure 5. Connecting the CT exponential distribution to the mass distribution of coating.** (a) The exponential CT distribution of BC aerosols with a slope  $k$  of  $0.016 \text{ nm}^{-1}$  under the steady state obtained from a PartMC-MOSAIC simulation. (b) The per-particle volume fraction of BC components and non-BC components with respect to per-particle BC mass. (c) The mass distribution of BC component (black line) and BC aerosols (red lines) derived from the log-normal distribution of  $D_c$  and the exponential distribution of CT ( $k = 0.016 \text{ nm}^{-1}$ ). (d) The mass distribution of aerosol components (normalized by dividing each component's mass by the maximum particle mass) with respect to BC mass in each particle obtained by the statistical method of Fierce et al. (2017) based on the results of PartMC-MOSAIC simulation. The values on the y-axis in (c) and (d) have been normalized

$$\left( \frac{dM_i/d\log(m_{BC})}{\int_{m_{BC}=10^{-5}}^{m_{BC}=1} dM_i/d\log(m_{BC}) \cdot d\log(m_{BC})} \right).$$

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## 5 Conclusions

Our study simulated the time evolution of the BC mixing state under continuous emission, condensation, coagulation, and deposition processes using the PartMC-MOSAIC model. The establishment of the baseline simulation case is based on SP2 measurements and has been verified for realism. Furthermore, based on the baseline case, we altered the pollution conditions (emission of aerosols and gas) and temperature conditions to simulate eight additional cases. The steady-state behaviour of the BC mixing state under the joint influence of emissions and atmospheric processes was confirmed for each case. During the simulation, the characteristic time for BC aerosols to reach a steady mixing state in the baseline scenario was 3.2 hours, with a range from 1.9 to 9.7 hours under different emissions of gases and particles and different temperature conditions. In addition, the CT distributions of BC aerosols in all cases followed an exponential linear distribution. Taking the baseline case as an example, the slope of the linear fitting was  $0.016 \text{ nm}^{-1}$  (with a correlation coefficient,  $R^2$ , of  $\sim 0.999$  during steady state) and the equivalent mean coating thickness was 62 nm, close to the value ( $\sim 63 \text{ nm}$ ) calculated by the per-particle method. Moreover, our study found that the numerical values of  $E_{\text{abs}}$  derived from the  $k$ -value method are highly consistent with the real values calculated by the per-particle method during the simulation. Therefore, the  $E_{\text{abs}}$  based on the steady-state theory can serve as a characterization of the light absorption enhancement of BC aerosols under the steady-state condition. Finally, we characterized the BC mixing state from different perspectives and revealed that while the mass fraction of coating is not uniform for different BC core sizes, the coating thickness follows the same exponential distribution for all BC core sizes. We successfully simplified the continuous variation of the BC mixing state and obtained a precise evaluation of the optical properties of BC aerosols based on the steady-state theory, approaching the accuracy level of the particle-resolved model. The steady-state theory can be applied to global models or chemical transport models. Take the Weather Research and Forecasting & Community Multiscale Air Quality Modeling System (WRF-CMAQ) for example, the  $k$ -value can be obtained, in principle, from the growth rate and removal rate in each grid cell, thus efficiently calculating the light absorption capacity of BC aerosol. However, in practice, it would be more convenient to develop an emulator for  $k$  based on PartMC-MOSAIC training data, which is an avenue of future work. This study has proved the importance and precision of utilizing the steady-state theory to characterize the BC mixing state and assess its optical properties, which provides a robust foundation for further research and applications.

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### **Code and data availability.**

The processed data, scripts, and setup of scenarios are available from the link: <https://zenodo.org/records/12177452> (Zhang  
345 et al., 2024). PartMC version 2.5.0 was used for the simulations in this paper, available at:  
<http://lagrange.mechse.illinois.edu/partmc/>.

### **Supplement.**

The supplement related to this article is available online at:  
350

### **Author contributions.**

JianW designed and directed the study. ZZ contributed to data analysis and wrote the first draft of this paper. JC, YJ, ZT, YC,  
and GC helped proofread the grammar of the article. JiapW, ZT and BW collected data. JianW, JiapW, NR, CL, and AD  
contributed the data interpretation and review of the paper.  
355

### **Competing interests.**

The authors declared that they have no conflict of interest.

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