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# Simulated Particle Evolution within a Winter Storm: Contributions of Riming to Radar Moments and Precipitation Fallout

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**Abstract**

Remote sensing radars from air- and spaceborne platforms provide critical observations of clouds to estimate precipitation rates across the globe. Capability of these radars to detect changes in precipitation properties is advanced by Doppler measurements of particle fall speed. Within mixed-phase clouds, precipitation mass and its fall characteristics are especially sensitive to the effects of riming. In this study, we quantified these effects and investigated the distinction of riming from aggregation in Doppler radar vertical profiles using quasi-idealized particle-based model simulations. Observational constraints of a control simulation were determined from airborne in situ and remote sensing measurements collected during the Investigation of Microphysics and Precipitation for Atlantic Coast-Threatening Snowstorms (IMPACTS) for a wintry-mixed precipitation event over the northeast United States on 04 February 2022. From the upper boundary of a one-dimensional column, particle evolution was simulated through vapor deposition, aggregation, and riming processes, producing realistic Doppler radar profiles. Despite a modest observed amount of supercooled liquid water (0.05 g m$^{-3}$), riming accounted for 55% of the ice-phase precipitation mass, cumulatively increasing reflectivity by 6.1 dB and Doppler velocity by 0.9 m s$^{-1}$. Independent evaluation of process-based sensitivities showed that while radar reflectivity is comparably sensitive to either riming- or aggregation-based particle morphology, the Doppler velocity profile is uniquely sensitive to particle density changes during riming. Thus, Doppler velocity profiles advance the diagnosis of riming as a dominant microphysical process in stratiform clouds from single-wavelength radars, which has implications for quantitative constraints of particle properties in remote sensing applications.

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## 1. Introduction

Ice crystals within precipitating winter storms evolve through an inherently stochastic sequence of microphysical processes which uniquely affect their physical properties and fall characteristics. This continuous and process-based evolution of ice-phase particles remains poorly represented by many numerical models and remote sensing retrieval algorithms. A fundamental limitation is that cloud and precipitation processes occur on physical scales that are several orders of magnitude smaller than typical cloud-scale model grids or the remote sensing instrument sampling volume. Nevertheless, realistic representation of varied particle populations within clouds is necessary to accurately estimate precipitation rates.

Commonly, populations of particles within some volume are expressed by a particle size distribution (PSD), and weighted integrals (i.e., moments) of the PSD are sensitive to the microphysical evolution of ice-phase particles (Morrison et al. 2020). Ice-phase precipitation mass is proportional to the second moment of the PSD. Because radar reflectivity, *Z*, is proportional to the square of the mass (i.e., the fourth moment of the PSD), the precipitation mass directly affects power returned to a radar. However, because of the physical complexity arising from diversity in initial ice crystal habits and their unique process-based morphologies with time, assumptions about the particle properties and the PSD are often necessary to derive remote sensing precipitation rate estimates. For example, ice crystals are commonly assumed to be spherical (e.g., Iguchi et al. 2018) and the population may be constrained to a prescriptive PSD shape or snow density (e.g., Grecu et al. 2016). A consequence of such a priori assumptions is that process-based variations cannot be expressed and retrieved precipitation rate estimates are inherently constrained, leading to snowfall rate underestimation and increased error compared to liquid phase (e.g., Speirs et al. 2017). To advance the utility of radar remote sensing measurements of ice-phase precipitation, it is important to understand the quantitative effects of process-based evolution on the intrinsic physical properties of precipitation in natural clouds and their implications for the radar measurements.

A remarkable property of precipitating clouds is that liquid water droplets are frequently present at sub-freezing temperatures alongside ice crystals. A region of cloud containing both ice and sub-freezing (i.e., supercooled) liquid water (SLW) is described as a mixed-phase layer. One implication of the mixed phase particle population is that depositional ice growth occurs at the expense of liquid water due to differences in saturation vapor pressures over ice and liquid surfaces, a process commonly referred to as the Wegener-Bergeron-Findeisen process (Pruppacher and Klett 1997). Additionally, upon contact with falling ice crystals, the SLW droplets freeze and are accreted by the crystal (i.e., riming), initiating a physical morphology of the particle. Natural ice crystals demonstrate tremendous variability in shape and complexity depending on growth habits (e.g., Magono and Lee 1966; Pruppacher and Klett 1997; Bailey and Hallet 2009). Because of this diversity, it is often convenient to define the crystal size along major and minor axes while the major axis is assumed to be along the maximum dimension of the crystal and the minor axis is along an orthogonal orientation. The aspect ratio defines the ratio between the crystal dimensions along the minor and major axes (Jensen and Harrington 2015). One commonly adopted conceptual description for the change in particle properties during riming is the "fill-in" model (Heymsfield 1982) whereby the liquid water will initially fill open voids, while largely maintaining the initial dimensions of the crystal axes. During later stages of the "fill-in" riming model, rime accumulates on the underside of the falling crystal, increasing the minor dimension of the crystal while the major dimension remains unchanged. With increasing riming, aspect ratio approaches unity, which is expected for heavily rimed graupel particles. Consequently, riming results in increasing particle density and, therefore, fall velocity. The adjustments in particle geometry and fall characteristics with rime accumulation are relative to, and dependent on, the initial ice crystal geometry and accreted rime but further dependent on prior and concurrent processes including vapor depositional growth and aggregation (e.g., Jensen and Harrington 2015).

Ice-phase particle growth by deposition of vapor-phase water directly increases the ice water content (IWC) and therefore, yields direct increases in $Z$ (Field et al. 2005, 2007). However, depositional mass accumulation occurs at a relatively slow rate, thus, gradual increases in $Z$ are expected from depositional growth alone. Aggregation of two or more particles does not explicitly alter the IWC of the particles, but rather redistributes the mass to a larger size particle. Despite unchanging IWC, increased particle diameters, $D$, during aggregation enhances radar scattering at a rate proportional to $D^4$ and consequently, $Z$ may be significantly increased by effects of aggregation. Through accumulation of liquid-phase water which yields increases in IWC, similar, rapid adjustments in $Z$ are also possible during riming. Evaluation of process-based effects on the evolution of the PSD moments and their implications for precipitation fallout from natural clouds is challenging because specific processes cannot be readily isolated,

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even if observations are collected in situ. In general, observationally-consistent numerical modeling simulations are necessary to determine such effects.

The physical scales of processes that govern the formation and evolution of falling ice crystals are not resolved by most numerical models. In bulk- and bin-microphysics schemes, ice-phase processes are commonly expressed implicitly through conversion processes whereby precipitation is exchanged among predefined categories (e.g., ice, snow, graupel, hail; Thompson et al. 2004; Morrison et al. 2005). However, prior studies (e.g., Colle et al. 2005; Morrison and Milbrandt 2011; van Weverberg et al. 2012) have demonstrated that the precipitation evolution and fallout is sensitive to a priori thresholds that define category conversions (e.g., snow to graupel during riming). For rimed growth, Lagrangian particle-based model simulations indicate that bulk particle density can undergo rapid evolution in response to small variations in the background SLW concentration, significantly modulating the particle fall velocity and surface precipitation rate (DeLaFrance et al. 2024). For remote sensing retrievals of mixed-phase precipitation, the effects of rime accumulation are constrained by the a priori assumptions about the particle's mass, geometry, or fall characteristics. Recently, diverse methodologies leveraging multi-frequency, dual-polarization, and Doppler radar measurements have been proposed for retrieving some properties of ice-phase particles that would otherwise be prescribed (e.g., Leinonen and Szyrmer 2015; Kneifel et al. 2016; Moisseev et al. 2017; Oue et al. 2018; Leinonen et al. 2018; Mason et al. 2019, Chase et al. 2021). Among these methods, leveraging radar Doppler data has shown promise in inferring the onset of riming and, subsequently, the riming-based modulations of retrieved particle property estimates. Mason et al. (2018) demonstrated that the addition of Doppler radar measurements provides constraint on the bulk ice density parameter in retrievals of snowfall. Furthermore, as shown by Kalesse et al. (2016), rimed snow occupies a unique region of Doppler spectra distinct from unrimed snow. One-dimensional (1D) spectral bin microphysics modeling simulations have shown promise in reproducing the Doppler spectra moments of riming but demonstrate sensitivity to particle property assumptions (Kalesse et al. 2016).

The 1D columnar modeling approach offers a framework for simulating explicit microphysical processes and detailed particle properties that are computationally prohibitive in a three-dimensional (3D) dynamic model. The 1D construction is therefore well suited to advanced bin and Lagrangian particle-based microphysics schemes. One challenge for such simulation

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designs, however, is constraining the model in a way that minimizes assumptions and, as a result, ambiguity in the attributing physical process for adjustments in the cloud's radar and precipitation characteristics (e.g., Kalesse et al. 2016; Bringi et al. 2020). Some assumptions can be constrained by coincident in situ and remote sensing radar measurements.

Data collected during winters of 2020, 2022, and 2023 from the Investigation of Microphysics and Precipitation for Atlantic Coast Threatening Snowstorms (IMPACTS) campaign (McMurdie et al. 2022) provide those constraints. Midlatitude cyclones over the United States East Coast and Midwest regions were comprehensively sampled by coordinated aircraft- and ground-based platforms to better understand the precipitation microphysics within regions of snowfall that organize into elongated regions commonly recognized as snowbands (e.g., Novak et al. 2004). Consistent with IMPACTS's goal to support improved numerical modeling and remote sensing retrievals of winter precipitation, in the present study we investigate the process-based effects of riming in a sampled storm that produced moderate rates of wintry-mixed precipitation for a prolonged period over the Northeast. Our overarching approach is to combine these observations with numerical modeling simulations to describe the process-based particle evolution and contributions of riming to the observed radar properties and precipitation rates. Here, we use an observationally-constrained, sophisticated Lagrangian particle-based model within a 1D columnar framework to address the following questions:

1. Can primary ice processes (i.e., deposition, aggregation, riming) within a simplified 1D simulation reasonably reproduce the observed evolution of particles within the natural cloud?
2. What were the quantitative contributions of riming to the observed Doppler radar vertical profiles and to the surface precipitation rate?
3. Do simulated Doppler radar vertical profiles yield characteristic responses to the onset or degree of riming that is distinct from other ice-phase processes (e.g., aggregation)?

## 2. Winter Storm Observations

### 2.1. 04 February 2022 Case Study

For this analysis, we will use IMPACTS observations collected during the 04 February 2022 event that delivered wintry-mixed precipitation across a broad region of the northeast US.

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IMPACTS deployed an in situ (P-3) and remote sensing (ER-2) aircraft. The P-3 aircraft was equipped with instrumentation to measure the in situ cloud microphysical properties and the high-altitude ER-2 aircraft was equipped with nadir-viewing remote sensing instrumentation analogous to those onboard satellite-based platforms (e.g., Skofronick-Jackson et al. 2017). The two aircraft targeted the storm over the coastal New England area where, as an example of the surface precipitation characteristics during this event, the Boston, MA (KBOS) Automated Surface Observing System (ASOS; Brodzik 2022a) reported nearly 32 mm of precipitation in 24 hours. Precipitation initially accumulated in the form of light to heavy rain before transitioning to freezing rain at about 1300 UTC, ice pellets by 1600 UTC, and back to freezing rain at about 1930 UTC. A transition to snow and continued accumulation occurred on 05 February at KBOS and over most of the New England area.

On 04 February, a broad frontal boundary extended from the Gulf of Mexico to Maine. The prolonged period of wintry-mixed precipitation over the northeast US was sustained by isentropic lifting of moisture-rich low-level flow along this front and overrunning a surface layer which, for many areas, remained subfreezing. Over the eastern US, a mean southwesterly flow developed ahead of an initially positively tilted 250-hPa trough at 0000 UTC 04 February that developed to nearly neutral tilt by 0000 UTC 05 February (Fig. 1a-c). An associated jet streak exceeding 150 kts was situated over northern New England such that between about 1200 UTC 04 and 0000 UTC 05 February, upper-level divergence in the right entrance region further supported lifting within the atmospheric column (Bjerknes 1951; Uccellini and Kocin 1987; Holton and Hakim 2012). During this time period, a modest elongated southwest-northeast oriented low-pressure minimum of approximately 1010 hPa was maintained over a broad region of coastal New England (Fig. 1d-f).

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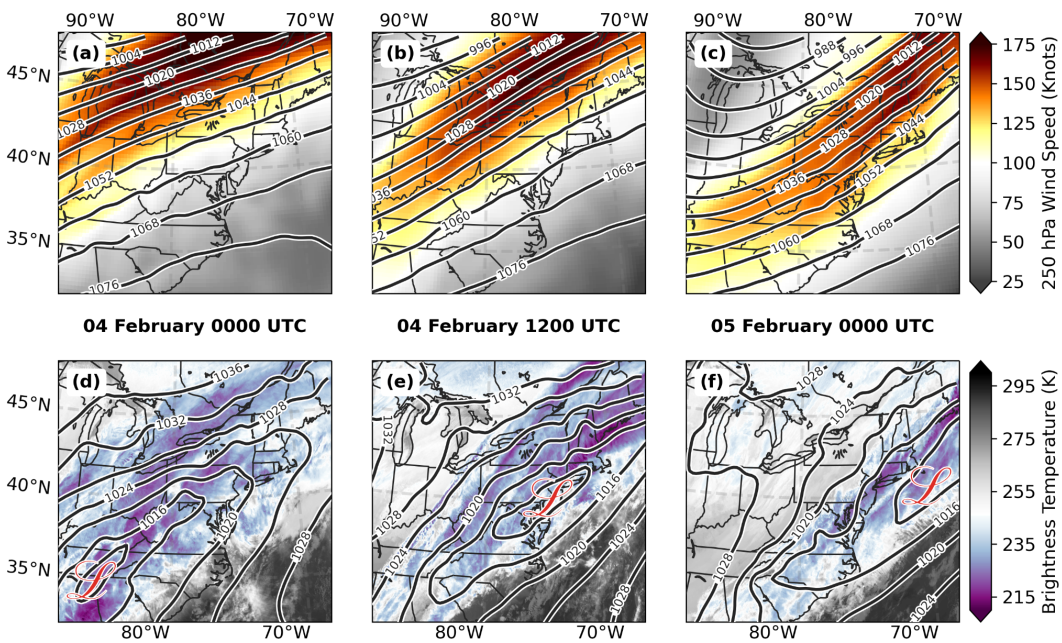


**Figure 1:** *Synoptic evolution of the winter storm that impacted the northeast US: (a-c) 250 hPa geopotential heights (dam) and wind speeds (knots) and (d-f) mean sea level pressure (MSLP, hPa) and cloud brightness temperature (K) for the times 0000 UTC 4 February (a, c); 1200 UTC 4 February (b, e) and 0000 UTC 5 February 2022 (c, f) . The 250-hPa and MSLP data are from the European Center for Medium-Range Weather Forecast Reanalysis v5 (ERA5; Hersbach et al. 2020) and the brightness temperature data are from the Geostationary Operational Environmental Satellites (GOES) 10.3 µm channel (Brodzik 2022b).*

Between about 1300 and 1800 UTC, the P-3 and ER-2 aircrafts flew a “lawnmower-style” pattern orthogonal to the long axis of an enhanced region of reflectivity while translating subsequent flight legs to the northeast, such that the storm was sampled in an approximately Lagrangian manner (Fig. 2a). The P-3 flew its initial flight leg south to north beginning at about 1340 UTC briefly at 6.5 km MSL before descending to a constant altitude of about 6.2 km MSL. At the southern end, this initial flight leg was near the NWS rawinsonde launch site at Islip, NY (KOKX). The P-3 descended on each subsequent flight leg to sample different layers of the cloud reaching an altitude of 3.0 km MSL on the final north-to-south flight leg, which transected the 0°C melting level. The two enhanced regions of reflectivity, on either side of the surface

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frontal boundary, exhibited differing cloud and precipitation properties. At the surface, the northern region of enhanced reflectivity was dominated by snowfall whereas the southern region was dominated by rain during the period of aircraft sampling then transitioning to wintry-mixed precipitation. As we describe in Section 2.2, in situ measurements are used to indicate riming, which was commonly observed over the southern region of enhanced reflectivity but absent over the northern region. Therefore, to address our science questions, our present analysis is constrained to measurements of the southern portions of flight legs (Fig. 2a).

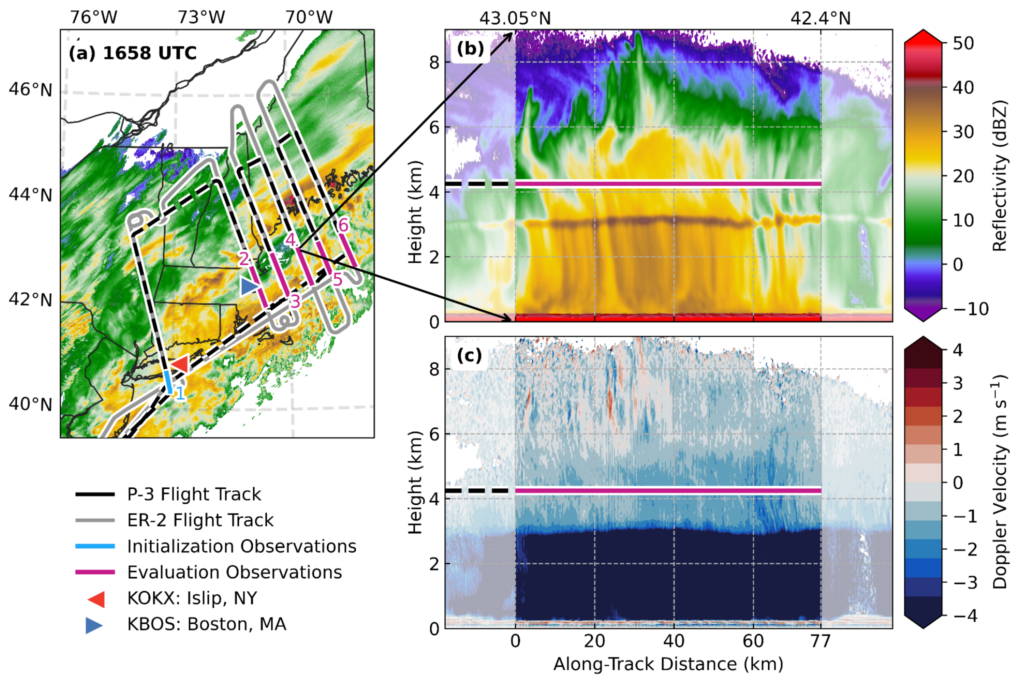


**Figure 2**: *IMPACTS operations on 04 February 2022 over the northeast US targeting regions of enhanced reflectivity that persisted for several hours in the operational National Weather Service (NWS) Multi-Radar Multi-Sensor (MRMS; Zhang et al. 2011) product. Shown are (a) the coordinated P-3 and ER-2 flight tracks and MRMS composite reflectivity at approximately mid-flight (1658 UTC) with subsets for each numbered flight leg at the southern enhanced region of reflectivity indicating data used for this study. Also indicated in (a) are the NWS rawinsonde launch site at Islip, NY (KOKX) and ground verification site at Boston, MA (KBOS). Ku-band reflectivity (b) and Doppler velocity (c) vertical profiles as measured by the ER-2 aircraft from 1628 (north) to 1634 UTC (south) depict the vertical cloud profile across the region of enhanced reflectivity (between transparent regions) for the fourth flight leg while the P-3 aircraft sampled in situ at ~4.3 km MSL altitude (magenta line in b, c), ending the flight leg at ~42.4°N.*

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### 2.2. Observations: Surface Based, Remote Sensing, and In Situ

The initial ER-2 and P-3 flight leg approximately overflew the NWS operational Islip, NY (KOKX) rawinsonde launch site (Fig. 2a). Because of the relatively steady-state nature of the storm during the aircraft sampling period, the KOKX 1200 UTC rawinsonde (Waldstreicher and Brodzik 2022) is used to estimate the atmospheric properties in the southern portion of the flight legs. Because these southern portions of the flight legs were mostly offshore, we use the nearest ASOS measurements at KBOS between 1300 and 1800 UTC to estimate the mean surface precipitation rate for model comparison. The ER-2 aircraft flew well above the storm at approximately 20 km MSL and operated two nadir-viewing radars on 04 February: the dual-band 13.9 GHz (Ku-band) and 35.6 GHz (Ka-band) High-Altitude Wind and Rain Airborne Profiler (HIWRAP; Li et al. 2016; Mclinden et al. 2022a) and the 94 GHz (W-band) Cloud Radar System (CRS; McLinden et al. 2022b). For radar reflectivity and Doppler velocity measurements of the precipitation, we use HIWRAP measurements, which have a vertical resolution of 150 m and a surface footprint of 1 km. At Ku-band, HIWRAP has a minimum sensitivity of approximately -10 dB at an altitude of 10 km MSL (Li et al. 2016).

Of the numerous instruments onboard the P-3 aircraft, those of relevance to this study include cloud Optical Array Probes (OAPs) and those that measure Liquid Water Content (LWC) and vertical air motion. The OAPs provide measurements of the two-dimensional projected sizes, shapes, and concentrations of particles. Data from a Two-Dimensional Stereo (2D-S; Lawson et al. 2006), which is commonly used for measurements of particles smaller than about 1 mm in diameter, are unavailable for the 04 February flight. However, a vertically oriented High-Volume Precipitation Spectrometer (HVPS; Lawson et al. 1993) provided particle measurements at sizes greater than 0.5 mm which were used to construct PSDs. Measurements of LWC were obtained from a Fast Cloud Droplet Probe (FCDP; Lawson et al. 2017) which operated as part of the Hawkeye combination probe. The FCDP uses Mie light scattering principles to size and count liquid water droplets from 2 to 50 μm in diameter, from which number and mass concentrations can be derived. Processing of the OAP and FCDP data was performed by the National Center for Atmospheric Research (NCAR; Bansemer et al. 2022) and is used at a 1 Hz frequency. Vertical air motion measurements were provided by the Turbulent Air Motion Measurement System (TAMMS), which uses several sensors at different locations on

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the aircraft to estimate the 3D components of the ambient wind (Thornhill et al. 2003). For TAMMS configured to the P-3, the accuracy of vertical winds measurements is estimated to be 0.2 m $s^{-1}$ (Thornhill 2022).

## 3. Simulation Design and Validation

### 3.1. Model Description

Several bulk microphysics schemes have been developed to more realistically represent the observed continuous evolution of ice-phase particle populations during riming (e.g., Morrison and Milbrandt 2015; Jensen et al. 2017; Cholette et al. 2023). Recently, this modeling approach has been extended to a Lagrangian particle-based scheme in the novel McSnow model (Brdar and Seifert 2018). The particle-based approach affords some advantages over the bulk approach, namely that evolution of a population of particles occurs independent of an Eulerian grid cell structure and is not constrained by assumptions about the PSD. The McSnow model was developed in a 1D columnar configuration and was expressly designed to simulate the evolution of an initial particle population during sedimentation through the column (Brdar and Seifert 2018). The notion of a particle in McSnow follows the super-droplet principle (Shima et al. 2009) whereby a multiplicity of real particles having commonality among physical properties and location are represented by a single super-particle. These super-particles are continuously introduced in the upper boundary of the model such that initially prescribed PSD characteristics are maintained and then evolve by vapor deposition and aggregation, with an option for riming to occur within a user-defined mixed-phase layer. From 2D simulations using McSnow, DeLaFrance et al. (2024) demonstrated that mixed-phase layer depth significantly modulates surface precipitation rates, varying up to 50% in response to a depth change of 750 m and that in situ measurements of SLW content provide a constraint on the layer's vertical extent. Following riming, melting of the particles occurs as its surface temperature exceeds 0°C, and collision-coalescence processes may then occur, but no additional precipitation mass is generated by warm-rain processes. The thermodynamic profile is prescribed and there are no mechanisms of feedback on the ambient environment based on the microphysical processes.

At any point in the column, detailed information about individual particle properties are directly accessible. In general, however, there is greater utility in the description of a population of particles in the form of a binned PSD expressed as the number concentration, *N*, of particles

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with diameter, *D*. We use a construction of 200 bins linearly spaced from 2 μm to 10 cm. From the PSD, radar quantities associated with moments of the PSD are computed by using a forward operator to estimate the radar scattering properties. Several scattering models have previously been adopted to radar scattering from ice crystals, principally differing in the complexity of the scattering particle's geometry. A population of ice crystals may be treated as spheres and scattering computed directly from Mie theory (Bohren and Huffman 1983); however, this approach vastly simplifies the irregular geometry of natural ice crystals. Scattering estimates based on the T-matrix method (Mishchenko et al. 1996) support nonsphericity of particles using a spheroidal shape. Furthermore, the orientation of the spheroids relative to the radar beam may be specified or randomized (Mishchenko and Travis 1998). A more sophisticated approach termed discrete-dipole approximation (DDA) accounts for the complex scattering interactions of irregular crystal geometry (Purcell and Pennypacker 1973) and is therefore a compelling method to estimate scattering of natural crystals. However, for our simulations, crystal habits or detailed properties of particle geometry are not predicted and thus, T-matrix is an apt method of estimating radar scattering. Specifically, we use the PyTMatrix software (Leinonen 2014) to estimate the radar backscattering cross section, σ, and compute *Z*, defined as:

$$Z = \frac{\lambda^4}{\pi^5 |K|^2} 10^{18} \int_0^\infty \sigma(D) N(D) dD, \tag{1}$$

where λ is the radar wavelength and K is the dielectric factor. From the simulations, we also estimate Doppler velocity, $V_D$, which is the reflectivity-weighted fall velocity, *v*, of the particles, defined as:

$$V_D = \frac{\int_0^\infty v(D)\sigma(D)N(D)dD}{\int_0^\infty \sigma(D)N(D)dD}. \tag{2}$$

For a mixed-phase cloud, Tridon et al. (2019) demonstrated a degradation of skill in T-matrix *Z* estimates at higher radar frequencies (i.e., Ka- and W-band). To minimize uncertainties associated with non-Rayleigh radar scattering effects (e.g., Matrosov 2007; Liu 2004, 2008), we specify λ = 25 mm for all calculations, which is comparable to the Ku channel on the HIWRAP radar. Additionally, for consistency with the HIWRAP measurements, a two-way correction for attenuation due to precipitation particles was applied following methodology described in Williams (2022).

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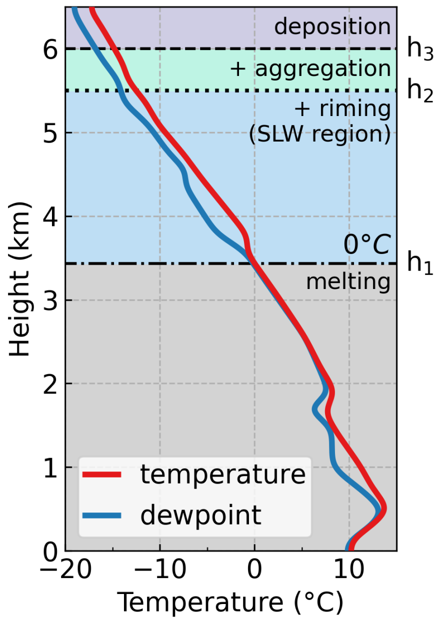


**Figure 3:** *Schematic of the one-dimensional columnar configuration of the McSnow model with prescriptive process-based layers for evolution of new particles initiated at the column's upper boundary. Static temperature and dew point vertical profiles are derived from the 04 February 1200 UTC KOKX rawinsonde.*

### 3.2. Control Simulation Design

We use the in situ measurements combined with rawinsonde data to construct a quasi-idealized cloud profile that is representative of the mean state of the 04 February storm which we apply prescriptively in the 1D columnar McSnow model. The process-based model design is illustrated by the schematic in Fig. 3. Introduction of new particles from a prescribed PSD occurs at 6.5 km MSL, which approximately corresponds to the uppermost height of in situ observations. Dominant particle types observed at this height were side planes and bullet rosettes. As newly introduced particles undergo sedimentation, growth occurs initially by vapor deposition only. Aggregation is introduced at 6 km MSL (-15°C) since aggregate particles, mostly side planes and other planar crystals, were present in observations below 6 km MSL. Riming is introduced at 5.5 km MSL, which we approximate as an upper extent of the mixed-phase layer based on the presence of SLW droplets and rimed particles beginning at flight leg 3

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(4.9 km MSL) and, subsequently, on legs 4 and 5 (4.3 and 3.6 km MSL). The onset of melting is determined by the thermodynamic profile which is obtained from the 1200 UTC KOKX rawinsonde. Although model processes are largely independent of an Eulerian grid (see discussion in Brdar and Seifert 2018, Section 2), model output and analysis occurs on a gridded column with 500 vertical levels, which yields a vertical resolution of 13 m. Additionally, we specify a time step of 5 s and total run duration of ten hours; results are analyzed as averages over the final five hours, after the system has reached a steady state.

As a constraint on the observational data used for simulation construction, we approximate the horizontal extent of the southern region of enhanced reflectivity by visually assessing its lateral edges during each flight leg using the Ku-band radar vertical profiles. An example of this approach is provided in Fig. 2b, c for the fourth flight leg in which data used is from the center portion of the figure. The boundaries (opaque regions) varied for each flight leg, adapting to the northeastward progression of the storm and translation of each flight leg. The initial PSD characteristics are derived from an average of the measurements on the uppermost height flight leg at ~6.5 km MSL between the southern end point of the leg and 40.7°N latitude (see Fig. 2a). Because measurements are unavailable for particles smaller than 0.5 mm, we fit a Gamma distribution to the mean PSD from HVPS measurements and then extend the fitted distribution to a lower size limit of 112.5 μm to estimate an IWC of 0.14 g $m^{-3}$ and total number concentration, $N$, of 23 x $10^3$ $m^{-3}$. For all simulations, an initial super-particle multiplicity of $10^5$ in the upper boundary is specified. We assume that newly initialized particles at 6.5 km MSL have a mass-dimension relationship of $m = 0.00294D^{1.94}$ (cgs units) following Brown and Francis (1995), for unrimed aggregate ice crystals in a stratiform cloud. From analysis of four IMPACTS events during the preceding 2020 deployment, Heymsfield et al. (2023) showed that $Z$ calculated from a PSD using the Brown and Francis (1995) mass-dimension relationship and a T-matrix approach yielded an agreement with observations at Ku band within 1.15 dB.

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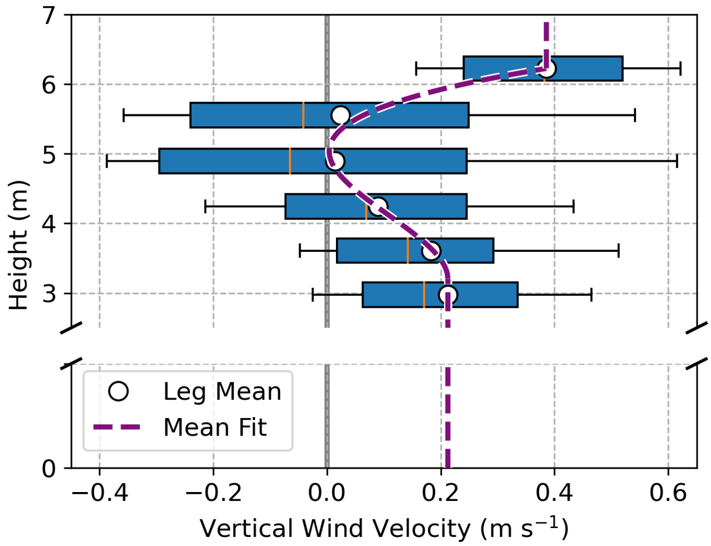


**Figure 4:** *Vertical wind velocity measurements from the Turbulent Air Motion Measurement System (TAMMS) during P-3 flight legs indicating lower to upper quartiles in the boxed regions, 10th and 90th percentiles at the whiskers, and medians at the vertical lines. A mean profile is fitted to the flight-level mean values (white markers).*

Falling particles are subject to an updraft. We estimate a mean-state vertical wind profile by fitting a third-degree polynomial curve to the mean measurements from each flight leg and extending the upper- and lower-most measurements as a constant value to heights beyond the observation altitudes (violet curve in Fig. 4). Within the mixed-phase layer ($h_2$ to $h_1$ in Fig. 3), SLW properties are derived collectively using FCDP measurements on flight legs 3, 4, and 5. We uniformly prescribe the mean values for SLW concentration of 0.05 g m$^{-3}$ (Fig. 5) and a characteristic droplet diameter of 22 μm within the mixed-phase layer.

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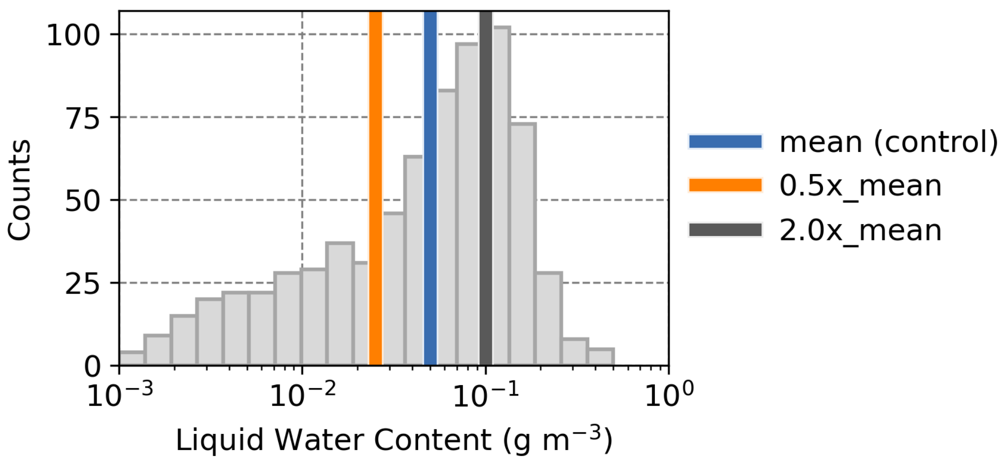


**Figure 5:** *Histogram of liquid water content (LWC) measurements from the Fast Cloud Droplet Probe (FCDP) during P-3 flight legs through mixed phase cloud (4.9 to 3.6 km MSL). Vertical bars indicate mean (0.05 g m$^{-3}$) and perturbed-state values used for sensitivity simulations scaled from the mean by factors of 0.5 and 2.0.*

Although we prioritize the use of observations for model constraint, several decisions are necessary regarding the parameterizations of modeled processes. With two exceptions, these parameterization decisions follow those discussed in DeLaFrance et al. (2024, see Section 2.3 and Appendix A). The first difference regards the aggregation process. Upon collision of two or more particles, a sticking efficiency parameter which scales from 0 to 1 is used to describe the probability of the particles merging, where an efficiency of 1 will always yield a union. The sticking efficiency parameterization follows Connolly et al. (2012), which is dependent on temperature and maximizes at -15°C. In testing, however, we found that the maximum likelihood estimate (MLE) values of Connolly et al. (2012; see Fig. 14b) yielded lower concentrations of large particles than were observed. Alternatively, use of a higher efficiency value inspired by the upper extent of their confidence interval yielded a more observationally-consistent PSD evolution and maximum particle sizes. Therefore, aggregation is introduced at 6 km MSL (Fig. 3) with a sticking efficiency of 0.9 at -15°C and linearly decreases to 0.5 at -10°C, remaining constant at 0.5 between -10° and 0°C. The second parameterization decision which differs from DeLaFrance et al. (2024) regards riming where a continuous approach was used in favor of a

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stochastic approach, although they describe only minor differences between the two approaches. In the present analysis, we find a slightly reduced collection of rime mass using the continuous parameterization when compared to the stochastic parameterization. Applying the continuous parameterization approach, particles accumulate a mean rime fractional mass of 0.49 by the time they reach 3.6 km MSL (flight leg 5, immediately above the melting level), whereas applying the stochastic parameterization, a rime fractional mass of 0.55 is accumulated. Visual assessment of the in situ particle imagery indicated that the stochastic method produces a more observationally consistent riming evolution. Therefore, the stochastic riming parameterization is used in all simulations.

### 3.3. Control Simulation Assessment

The objective for a control simulation is to produce an evolution of a population of particles within a vertical column that is physically consistent with the observed cloud profile. In Fig. 6, we compare the control simulation PSD to the mean observed PSD ($D \geq 0.7$ mm). Although PSD measurements at smaller particle sizes are unavailable for this flight, the approximately Lagrangian aircraft sampling yielded a temporally consistent evolution of the PSD at larger sizes. Measurements from flight leg 1 are used to assess the simulation during the particle initialization stages within the uppermost region of the model, whereas measurements collected downstream on flight legs 2 through 6 are used to assess simulation performance during the later stages of particle evolution. The model produces an initial particle population at 6.5 km MSL (Fig. 6a) that is consistent with the mean observations at large particle sizes and follows the assumed Gamma distribution form at small sizes. Flight leg 5 (Fig. 2a), at approximately 3.6 km MSL, was the lowest altitude flown before reaching the melting level. At this altitude, evaluation of the simulation shows skill in evolving this initial particle population by deposition, aggregation, and riming processes throughout a nearly 3 km-deep cloud layer.

Particle growth between 6.5 km (Fig. 6a) and 3.6 km MSL (Fig. 6b) through aggregation and to a lesser extent, depositional growth, is expressed in the shift of the observed PSD to larger particle sizes. This evolutionary characteristic is reproduced by the control simulation although slightly larger maximum particle sizes are generated, and the ice mass may be underrepresented among particles smaller than about 2 mm in diameter. We note, however, that sizing uncertainty in the observed measurements is greater at these small sizes owing to the relatively coarse pixel

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resolution of 150 µm for the HVPS probe (Bansemer et al. 2022). To further validate the control simulation and to assess the continuous particle evolution throughout the vertical profile, $Z$ is estimated from the simulated PSD and compared to the HIWRAP Ku-band measurements.

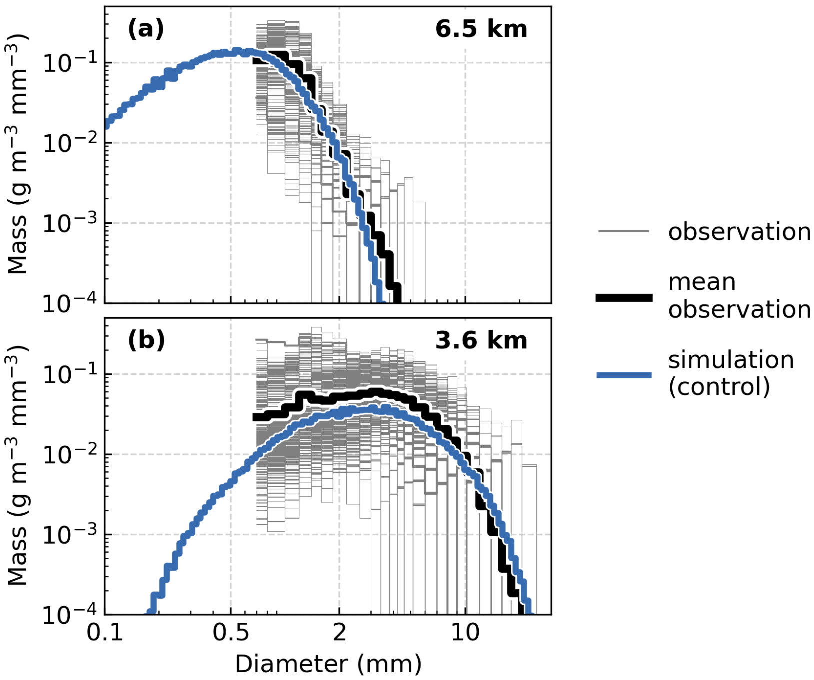


**Figure 6:** *Particle size distributions (PSDs) of ice mass for observed 1 Hz and mean values derived from (a) P-3 flight leg 1 at 6.5 km MSL and (b) flight leg 5 at 3.6 km MSL (see Fig. 2) and for the control simulation at equivalent altitudes.*

Figure 7 shows the median observed vertical profile of $Z$ and $V_D$ computed from the downstream flight legs 2 through 6, as indicated in Fig. 2a. Data from the lowest 500 m were removed due to noise from ground clutter. From the observed vertical profiles, several inferences are made about the microphysical processes. Beginning at 6 km MSL, $Z$ rapidly increases with descent, which is expected with an onset of aggregation. The rate of increase in $Z$ with descending height reaches a relative maximum near 5.5 km MSL (Fig. 7a), coincident with an apparent acceleration of $V_D$. Within the subsequent 1 km (5.5 km to 4.5 km MSL), $V_D$ increases

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from -0.72 m s$^{-1}$ to -1.00 m s$^{-1}$ (Fig. 7b). This effect is assumed to be associated with the onset of riming, and subsequently, changes in particle densities. Particle melting begins near 3.4 km MSL, at which point a bright band signature is apparent and $V_D$ rapidly increases. Below the bright band, $Z$ remains nearly constant at about 25 dBZ and $V_D$ is about -5 m s$^{-1}$.

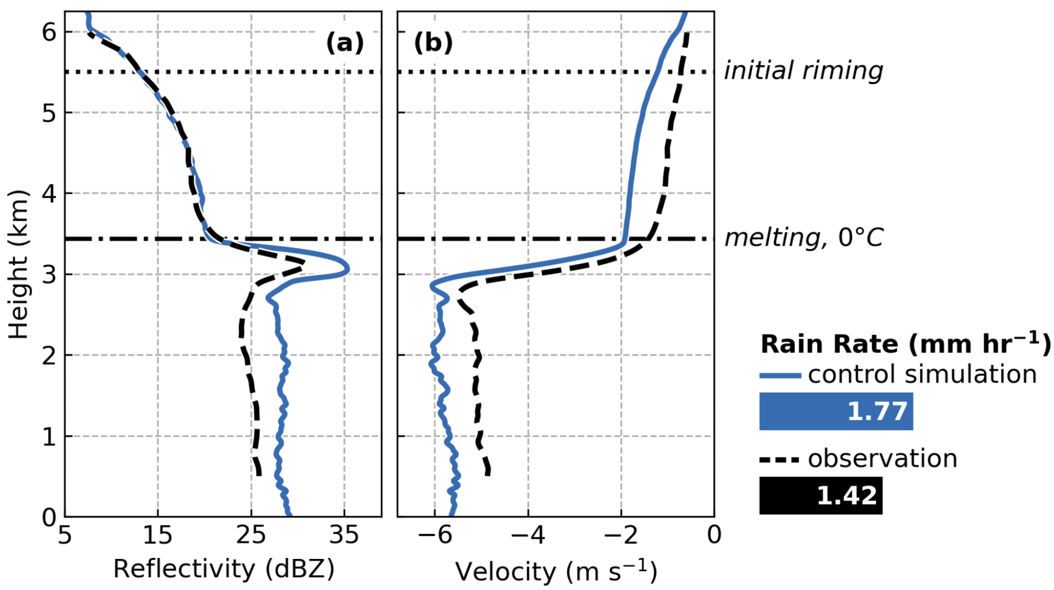


**Figure 7:** *Vertical profiles of (a) radar reflectivity and (b) Doppler velocity at Ku band for the control simulation (blue lines) and observed (dashed black lines) median from ER-2 HIWRAP radar during flight legs 2-6 (see Fig. 2a, magenta segments). Data for the observed profile below 500 m MSL are omitted due to ground clutter. A dotted line at 5.5 km MSL indicates the onset of riming and a dash-dotted line indicates the 0°C height. Also shown at the right are the surface rain rate values from the control simulation (blue) and observed at KBOS (black) between 1300 and 1800 UTC on 04 February 2022; horizontal bar lengths illustrate magnitude differences.*

The vertical profile of $Z$ is well reproduced by the control simulation, particularly above the melting level (Fig. 7a), which suggests confidence in its prescriptive configuration. Upon melting, $Z$ is overestimated by the control simulation and maintains a bias of about 2 to 5 dB throughout the warm layer. While an evaluation of warm-rain processes is beyond the scope of the present study, it is possible that this overestimate in $Z$ results from an incomplete representation of warm-rain processes by the model, such as droplet breakup and shedding, or

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from uncertainties in the scattering estimates. Confirmation of an attributable mechanism would be challenging without in situ observations below the melting level. Rain rates at the surface are one common model validation metric. Because the aircraft sampling occurred primarily offshore (see Fig. 2a), an ideally situated ground site is unavailable. However, we find comparison with a nearby ground site useful towards determining whether the control simulation produces physically realistic estimates that are representative of the rainfall across the broader region. At the surface, during aircraft sampling (1300 to 1800 UTC), the nearest ground site, KBOS, reported a rain rate of 1.42 mm hr$^{-1}$. The control simulation produces about 25% more surface rain with an average rain rate of 1.77 mm hr$^{-1}$.

Despite the confidence in $Z$ aloft, we find that $V_D$ is underestimated by about 0.5 to 1 m s$^{-1}$ in the control simulation but are within an uncertainty range of ± 1 m s$^{-1}$ (Matthew McLinden, personal communication, 25 April 2024) for the HIWRAP Ku-band $V_D$ measurements. This bias between the observed and simulated $V_D$ is consistent throughout the column, suggesting that this consistent bias may be explained, to a large extent, by uncertainty in the observations. More importantly for this analysis, the relative changes in $V_D$ with height, which have process-based implications, are similar between the observed and simulated profiles.

## 4. Process-Based Contributions and Sensitivities on Doppler Radar Vertical Profiles

A principal advantage of the particle-based design of the McSnow model is that information about microphysical properties is retained by the model at the scale of the individual particles. For particles in the control simulation, the onset of riming at 5.5 km MSL ($h_2$ in Fig. 3) initiates a change in the physical evolution of the particle with subsequent sedimentation. At 3.6 km MSL, the particles have accumulated a mean rime fractional mass of 0.55, increasing the total precipitation mass and accelerating its fallout rate. Radar scattering by the particle, expressed through $Z$, is also modified by rime accumulation, yet these effects are difficult to distinguish from concurrent processes, including deposition and aggregation. To investigate these scattering implications, we estimate the vertical profile of $Z$ with and without contributions of rime mass.

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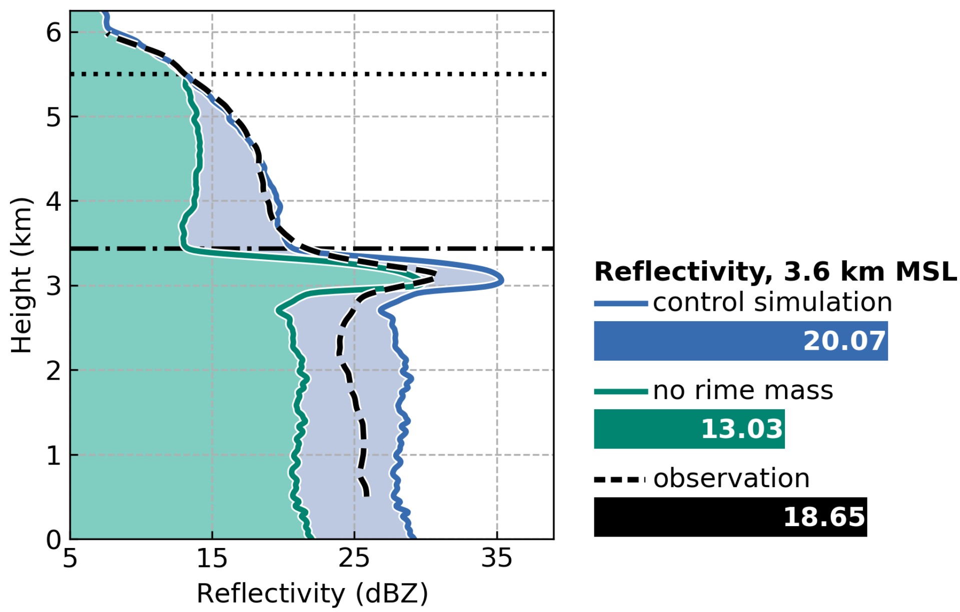


**Figure 8:** *As in Fig. 7a but with an added vertical profile (in green) for estimated reflectivity (Z) with particle rime mass removed. Shown at the right are simulated and observed Z values computed at 3.6 km MSL; horizontal bar lengths illustrate magnitude differences.*

Figure 8 compares $Z$ from the control simulation (as in Fig. 7a) to an unrimed estimate of $Z$ obtained by subtracting the rime mass from the particles and recomputing their scattering properties. Removal of rime mass appears to significantly impede further increases in $Z$ with descending height below 5.5. km MSL. Near the melting level, $Z$ is reduced from 20.07 to 13.03 dBZ between the control and simulation and the unrimed estimate, suggesting that the accumulated rime mass contributes to about 35% of the total $Z$ (in dB units). This calculation, however, only considers the implications of riming on radar scattering; the complex interactions of concurrent processes are neglected by solely removing the rime mass from evolved particles in the control simulation. Additionally, the effects on $V_D$, which manifest cumulatively during riming, cannot be investigated in the same manner. To explicitly investigate the effects of riming on the radar profiles, and to distinguish these effects from concurrent processes, we introduce several sensitivity simulations which independently perturb the riming or aggregation processes.

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| Simulation | Description | Perturbation Assignment |
|---|---|---|
| **control** | Observation-based **mean-state** simulation | none |
| **high_SLW** | **Increase SLW** by **2.0** from control | 0.100 g m$^{-3}$ LWC |
| **low_SLW** | **Reduce SLW** by **0.5** from control | 0.025 g m$^{-3}$ LWC |
| **no_riming** | **Remove riming** to distinguish effects from aggregation | Riming process turned off |
| **MLE_C12_agg** | **Reduce aggregation** from control to **moderate** efficiency | MLE sticking efficiency; see Fig. 14, Connolly et al. (2012) |
| **low_C12_agg** | **Reduce aggregation** from control to **low** efficiency | 0.5 x MLE sticking efficiency; see Fig. 14, Connolly et al. (2012) |

**Table 1:** *Descriptions and perturbations relative to the control simulation applied for each simulation.*

Although the southern regions of the 04 February 2024 event were predominantly stratiform, variations in the mixed-phase layer LWC were observed (Fig. 5). Within sufficiently deep mixed-phase layers, prior model simulations have demonstrated that small (e.g., < 0.05 g m$^{-3}$) perturbations in LWC alter particle fallout characteristics which can yield substantial increases or decreases in the surface precipitation rate (DeLaFrance et al. 2024). Here, we similarly introduce two sensitivity simulations perturbing LWC within the mixed-phase layer ($h_1$ to $h_2$ in Fig. 3), within the range of observed LWC (Fig. 5). In the control simulation, we prescribed the mean observed LWC value of 0.05 g m$^{-3}$. A scaling factor of two relative to the control is used to prescribe a high concentration (0.1 g m$^{-3}$) for the "high_SLW" simulation and low concentration (0.025 g m$^{-3}$) for "low_SLW" concentration. As a limiting case which is analogous to the removal of rime mass (Fig. 8), we construct a "no_riming" simulation with the riming process inactive. A brief summary of these riming sensitivity simulations is provided in Table 1.

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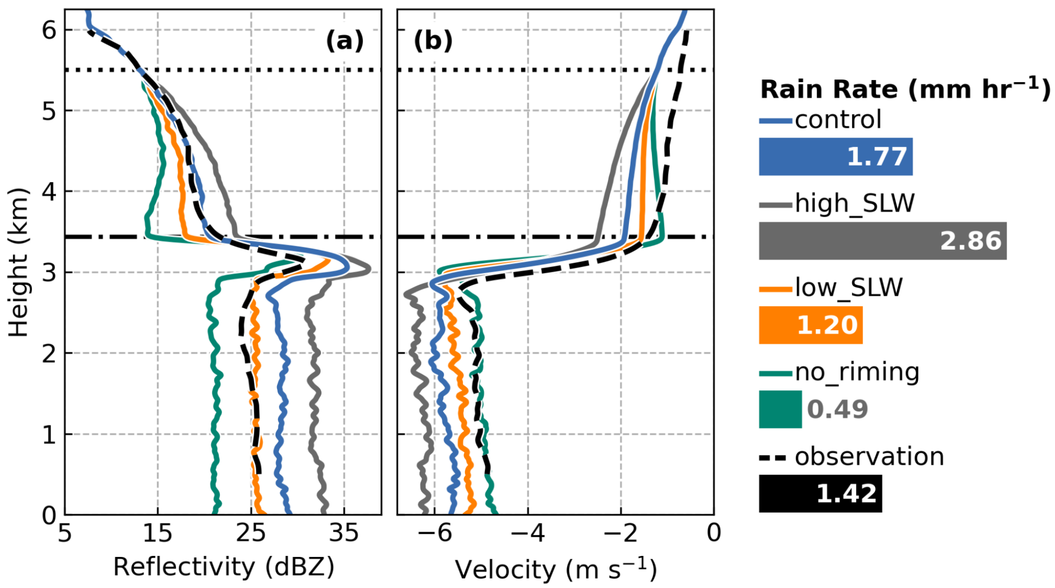


**Figure 9:** *As in Fig. 7 but for the control and riming-based sensitivity simulations: high_SLW, low_SLW, and no_riming. At the right are surface simulated and observed surface rain rate values; horizontal bar lengths illustrate magnitude differences.*

Vertical profiles of $Z$ and $V_D$ for the high_SLW, low_SLW, and no_riming sensitivity simulations relative to the control are shown in Fig. 9. Complete removal of the riming process in the no riming simulation (Fig. 9a) produces a similar $Z$ profile as found by computing $Z$ for equivalent unrimed particles from the control simulation (Fig. 8). This result underscores the significant sensitivity of $Z$ to changes in particle mass during riming, despite concurrent microphysical processes. Perturbing LWC by a factor of 2 in the high_SLW or 0.5 in the low_SLW simulations relative to the control produces opposing, but similar in magnitude, changes in $Z$ (Fig. 9a), indicating a regularity in the response of $Z$ to SLW variability. Similarly, the effects of SLW variability on $V_D$ demonstrate a regular response (Fig. 9b). We note that these simulation responses in $Z$ and $V_D$ to SLW variability assume that the particles are well mixed such that probabilistic collision of ice crystals and SLW droplets is the same throughout the layer.

In the high_SLW simulation, the rate of further $V_D$ acceleration with descent below 5.5 km MSL is nearly doubled relative to the control. Conversely, below about 5 km MSL, further

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increases in $V_D$ cease in the low_SLW simulation and decrease in $V_D$ occurs in the no_riming simulation. The competing effects of riming and aggregation processes on $V_D$ manifest in the low_SLW and no_riming simulations; riming accelerates the $V_D$ with mass accumulation whereas in the absence of riming, further aggregation yields larger, lower density particles with reduced fall speeds. Consequently, vertical profiles of $V_D$ may provide an insight into dominant microphysical processes, which is consistent with the notion that rimed particles occupy a distinct region of the Doppler spectra (Kalesse et al. 2016). To advance the differentiation of particles evolved by riming, it is necessary to first consider relative effects of variability in the aggregation process.

In our development of the control simulation for the 04 February 2022 event, the aggregation process was initially assumed to follow a temperature dependent sticking efficiency identified as the MLE by Connolly et al. (2012; see Fig. 14b). Comparison with in situ PSDs indicated that the MLE sticking efficiency parameter was insufficient to generate observed concentrations of large particles, motivating the use of an increased sticking efficiency in the control simulation. However, to elucidate the effects of aggregation efficiency on radar profiles, we now consider a sensitivity simulation, "MLE_C12_agg", which follows the MLE sticking efficiency of Connolly et al (2012). Additionally, analogous to the design of the riming sensitivity simulations, we introduce a "low_C12_agg" simulation for which the sticking efficiency is further reduced from the MLE estimate by a factor of 0.5. Relative to the control simulation, the reduction in sticking efficiently in the MLE_C12_agg and low_C12_agg sensitivity simulations lack observational consistency with the presently analyzed 04 February 2022 event. However, it is useful to consider the implications of a realistic range of variability in the aggregation process efficiency to inform general distinctions from the effects of riming within vertical profiles of $Z$ and $V_D$.

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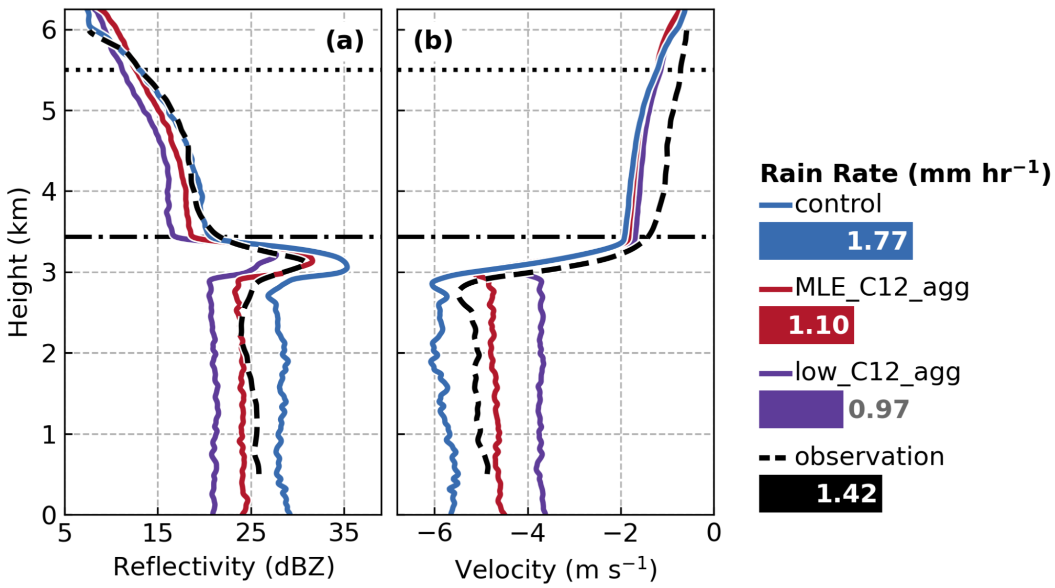


**Figure 10:** *As in Fig. 7 but for the control and aggregation-based sensitivity simulations: MLE_C12_agg, and low_C12_agg. At the right are surface simulated and observed surface rain rate values; horizontal bar lengths illustrate magnitude differences.*

Figure 10 shows the vertical profiles of $Z$ and $V_D$ for the aggregation efficiency sensitivity simulations, MLE_C12_agg and low_C12_agg. Reducing aggregation efficiency suppresses the generation of large particles and because of the strong dependency of radar backscatter on particle size, $Z$ decreases relative to the control (Fig. 10a). Additionally, smaller aggregate particles become smaller targets for collision with SLW droplets to accumulate rime mass, which also reduces $Z$. The latter effect manifests in the reduced surface rain rates, decreasing by 38% in the MLE_C12_agg (1.10 mm hr$^{-1}$) and 45% in the low_C12_agg (0.97 mm hr$^{-1}$) simulations relative to the control (1.77 mm hr$^{-1}$). Conversely, a reduction in aggregation efficiency has a minimal effect on $V_D$ for ice-phase particles (Fig. 10b). Above the melting level, at 3.6 km MSL, $V_D$ in the MLE_C12_agg simulation is reduced from the control simulation by 0.08 m s$^{-1}$ and in the low_C12_agg, reduced by 0.24 m s$^{-1}$. This relative insensitivity of $V_D$ to aggregation arises despite these sensitivity simulations assessing a broad range of possible sticking efficiencies. For example, at -15°C, the sticking efficiency is reduced from 0.9 in the control to 0.32 in the

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low_C12_agg simulation and at -10°C, from 0.5 in the control to 0.12 in the low_C12_agg simulation.

Below the melting layer, however, the effects of aggregation on $V_D$ become significant, decreasing by approximately 2 m s$^{-1}$ between the control and low_C12_agg simulations. Similarly, the surface rain rate decreases by about 45% between the control and low_C12_agg simulations. Thus, despite the significant implications of the aggregation process on precipitation production and its fallout, its variations are not readily perceived in vertical profiles of $V_D$. This finding significantly differs from the robust sensitivity of $V_D$ to variations in the riming process. While variations in the aggregation and riming processes may manifest similarly in vertical profiles of $Z$, we find that $V_D$ is uniquely sensitive to riming. Thus, vertical profiles of $V_D$ show promise in identification of riming as a dominant ice-phase microphysical process, which is ambiguous in profiles of $Z$, only.

## 5. Discussion

Sensitivity in vertical profiles of both $Z$ and $V_D$ owing to rime accumulation rates were previously shown by Kalesse et al. (2016) from bin model simulations by prescribing a fixed vertical profile of LWC then testing two different riming efficiency parameterizations. Their two simulations yielded similar vertical gradients in $Z$ and $V_D$ profiles but with differences in magnitude. They attributed these differences to assumptions about the physical morphology of the ice crystals with accretion of rime mass that had implications for their scattering properties. In our study, we uniquely provided an observational constraint to establish a control state simulation and modeling framework for assessing impacts of riming and aggregation independently. By selecting a fixed riming parameterization for all simulations using this framework, we were able to assess $Z$ and $V_D$ sensitivities attributable to LWC perturbations within the range of observed variability. We found that a small ($\leq$ 0.05 g m$^{-3}$) range of perturbations in the LWC produced substantial changes in the surface precipitation rate and a corresponding sensitivity in vertical profiles of $Z$ and $V_D$.

The sensitivities expressed in Doppler radar profiles to LWC perturbations is tied to the impact on bulk microphysical properties, especially particle density, $\rho_e$. In the deposition- and aggregation-prescribed region above 5.5 km MSL (Fig. 11), $\rho_e$ rapidly decreases with descending height due to the efficient aggregation of increasingly open particle geometry. At 5.5

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km MSL, riming is introduced and $\rho_e$ approaches 0.02 g cm$^{-3}$, remaining nearly constant until the melting level as a result of the competing effects of aggregation and riming. In the high_SLW simulation, the effects of riming dominate whereby the gradient in $\rho_e$ abruptly increases with descending height. Conversely, in the low_SLW and no_riming simulations, the effects of aggregation continue to dominate and $\rho_e$ further decreases.

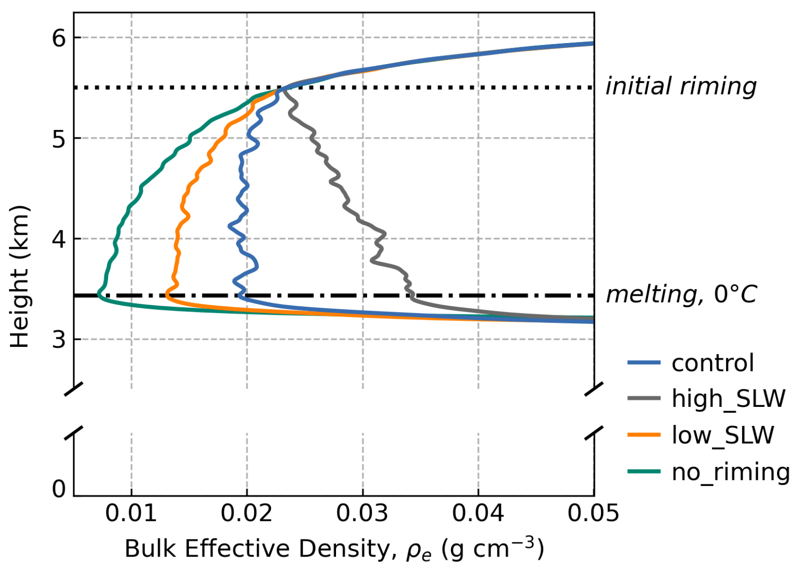


**Figure 11:** *Vertical profiles of bulk effective density, $\rho_e$, for the evolved particle population for the control simulation and three riming-based sensitivity simulations described in Section 4: high_SLW, low_SLW, and no_riming. Calculations of $\rho_e$ assume equivalent spherical volumes of the particles following Heymsfield et al. (2004).*

Despite the opposing process-based effects on the evolution of $\rho_e$ with height, our simulations suggest that the effects of aggregation and riming are not readily distinguished by $Z$ from a Ku-band radar band alone. Riming may be detectable, however, from three-wavelength (Ku-, Ka-, and W-band) radar by leveraging differential attenuation effects. In prior idealized modeling simulations for rimed particle growth scenarios, Leinonen and Szyrmer (2015) identified unique signatures of riming by comparing dual-wavelength ratios (DWR) between Ka and W bands with DWR at Ku and Ka bands. However, they found the magnitude of this signature to be modest and proposed that it would likely be difficult to accurately distinguish in

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observational data. Mason et al. (2019) later investigated the use of triple-frequency Doppler radar measurements from mixed-phase clouds during wintertime snow events to constrain the retrievals of bulk microphysical properties, including the PSD shape factor and $\rho_e$. They found that triple-wavelength $Z$ measurements effectively constrained the PSD shape parameter, but did not constrain $\rho_e$. Rather, $V_D$ measurements were necessary to identify transitions to rimed growth cloud regions and provide constraint on $\rho_e$. Our findings demonstrate that this constraint on $\rho_e$ is attributable to the unique density-dependent response in $V_D$ expressly owing to variations in the riming process within mixed-phase cloud layers with concurrent riming and aggregational growth. Further, our findings suggest that, when combined with $Z$, coincident vertical profiling measurements of $V_D$ have utility towards diagnosing riming as a dominant process within stratiform clouds from a single-wavelength radar.

## 6. Conclusions

The evolution of ice-phase precipitating particles within a mixed-phase stratiform cloud was simulated to evaluate the effects of riming on the PSD moments and assess the process-based implications on Doppler radar vertical profiles. In situ and remote sensing airborne observations collected during the IMPACTS field campaign for a prolonged wintry-mixed precipitation event over the northeast US on 04 February 2022 were used to design and constrain a quasi-idealized 1D mean-state control simulation. Using the Lagrangian particle-based McSnow model, we defined an initial population of ice particles based on in situ measurements in the upper portion of the cloud. As those particles fell, initial evolution occurred by vapor deposition followed by subsequent additions of aggregation and then riming within prescriptive observation-based layered regions. Radar scattering properties were estimated using a T-matrix forward operator and vertical profiles of $Z$ and $V_D$ were computed from the evolved PSD, then evaluated through comparisons with the airborne radar data. The effects of riming on PSD moments expressed through Z and $V_D$ were assessed from simulations which introduce small perturbations in cloud LWC within a range of observed variability. To distinguish effects of riming and aggregation, two additional sensitivity simulations were introduced to determine the unique implications of aggregational growth efficiency on Z and $V_D$. Through these approaches, we found:

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- Ice-phase precipitation particle evolution in a mixed-phase wintertime storm cloud is well constrained by the 1D quasi-idealized McSnow model.
- Despite modest supercooled liquid water concentrations, rime accumulation is estimated to account for 55% of particle mass generated above the melting level, dominating ice-phase contribution to precipitation rates.
- Riming cumulatively increased radar reflectivity above the melting level by an estimated 6.1 dB and Doppler velocity by 0.9 m $s^{-1}$ and demonstrated significant sensitivity to small perturbations in supercooled liquid water concentrations.
- Vertical profiles of radar reflectivity demonstrate similar sensitivities to riming and aggregation, but Doppler velocity is uniquely sensitive to riming-based perturbations through changes in particle density.

Constraining parameterized treatments of rimed particle evolution in numerical models is a known source of uncertainty in simulations of precipitation from bulk-, bin- and Lagrangian particle-based models (e.g., Lin and Colle 2011; Jensen and Harrington 2015; Jensen et al. 2017; Brdar and Seifert 2018). One objective of our analysis was to address this constraining need through quantifying precipitation sensitivities to riming in model simulations based on an observed range of variability in LWC. We found a difference of about 6% in rime fractional mass accumulation in our control simulation whether using a continuous or a stochastic representation of riming with the McSnow model. This effect was expressed within a modeling framework using a quasi-idealized and steady-state 1D column with a homogeneous mixed-phase layer. This approach was appropriate for our intentionally selected region of the observed storm because of its idealistic layered vertical structure apparent in radar observations (Fig. 2b, c), along with its known presence of SLW based on in situ observations. However, in reality, processes are not neatly initiated at distinct levels (e.g., in convective areas). It is expected that increasing ambiguity exists in distinguishing concurrent microphysical processes in these scenarios and, thus, our analysis did not assess the full natural range of complexity in mixed-phase precipitation processes.

While model schemes have become increasingly sophisticated, it is not clear that uncertainty in ice-based precipitation estimates have necessarily reduced, highlighting the need for judicious use of observations to advance constraints on modeled processes (e.g., Morrison et al. 2020).

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Because of the capacity for explicit process representation at the scale of individual particles, Lagrangian models (e.g., McSnow) may be ideally suited to addressing these challenges, especially when combined with datasets which prioritize observations that are consistent with the evolution of particles. This observational consideration was favored during the 04 February 2022 event, which was sampled by IMPACTS in an approximately Lagrangian manner. In this study, we focused on riming as a primary ice-phase process, but the northern region of the sampled storm observed significantly less SLW and rime accumulation, presenting a unique natural laboratory for evaluation of modeled aggregation. Sticking efficiencies during aggregation are highly uncertain and difficult to constrain from laboratory experiments (e.g., Connolly et al. 2012) yet, as we demonstrated in our study, have significant implications for the accuracy of simulated *Z* and rain rates. Ongoing work involves curating the in situ measurements of particle evolution within this northern storm region to constrain Lagrangian particle-based simulations and assess the ambient environmental dependencies (i.e., temperature, water supersaturation) and ranges of sensitivities associated with modeled aggregation.

## 7. Data Availability Statement

All field observation data from IMPACTS used in this study are accessible through the NASA Distributed Active Archive Center (McMurdie et al. 2019). Readers can find a complete description of the McSnow model and its availability in Brdar and Seifert (2018).

## 8. Author Contributions

All authors contributed to the study design and methodology decisions. Andrew DeLaFrance conducted the data curation and performed the simulations and computations from model output. All authors contributed to the evaluation and interpretation of the results. Andrew DeLaFrance prepared the manuscript with contributions from all co-authors.

## 9. Competing Interests

The authors declare that they have no conflict of interest.

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## 10. Acknowledgments

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