1 **Improved Snow Property Retrievals by Solving for Topography in the Inversion of At-**

2 **sensor Radiance Measurements**

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8

9 **Abstract**

10   Accurately modelling optical snow properties like snow albedo and specific surface

11 area (SSA) are essential for monitoring the cryosphere in a changing climate and are

12 parameters that inform hydrologic and climate models. These snow surface properties can be

13 modelled from spaceborne imaging spectroscopy measurements but rely on Digital Elevation

14 Models (DEMs) of relatively coarse spatial scales (e.g. Copernicus at 30 m), which degrade

15 accuracy due to errors in derived products – such as slope and aspect. In addition, snow

16 deposition and redistribution can change the apparent topography and thereby static DEMs

17 may not be considered coincident with the imaging spectroscopy dataset. Testing in three

18 different snow climates (tundra, maritime, alpine), we established a new method that

19 simultaneously solves snow, atmospheric, and terrain parameters, enabling a solution that is

more unified across sensors and introduces fewer sources of uncertainty. We leveraged imaging spectroscopy data from AVIRIS-NG and PRISMA (collected within 1 hour) to validate this method and showed a 25% increase in performance for the radiance-based method over the static method when estimating SSA. This concept can be implemented in missions such as Surface Biology and Geology (SBG), Environmental Mapping and Analysis Program (EnMap), and Copernicus Hyperspectral Imaging Mission for the Environment (CHIME).

Key Words: Imaging Spectroscopy, Snow Properties, Topography, Snow Albedo

1 Introduction

Accurately mapping snow surface properties is essential for seasonal snow zones in a changing climate especially in regions where seasonal snowpack is expected to change dramatically in the coming decades (Siirila-Woodburn et al., 2021). For example, snow albedo plays a crucial role in melting of the snowpack during the ablation season (Wang et al., 2020) with changes in snow albedo directly affecting the amount of absorbed solar radiation, and therefore the amount of snow that is melted off. Throughout the winter season, snow albedo fluctuates due in part to grain size (Seidel et al., 2016) and light absorbing particles (Kaspari et al., 2015; McKenzie, 2020; Schmale et al., 2017; Skiles & Painter,

39   2017). With a limited number of *in situ* snow stations around the globe, and the snow surface

40   constantly undergoing metamorphism across space and time, satellite imagery represents the

41   best potential for spatially and temporally complete mapping of snow properties. Accurately

42   retrieving snow albedo and other snow surface properties from satellite imagery is

43   paramount, especially in a rapidly changing climate (Malmros et al., 2018).

44   Retrieval of snow properties from satellite remote sensing relies on Digital Elevation

45   Models (DEMs) to correct for local terrain effects (Bair et al., 2021; Bair et al., 2022; Dozier

46   et al., 2022). In a previous study, researchers found global DEM products to have "blunders

47   and errors" when compared to airborne lidar, particularly in derived slope and aspect which

48   cause severe errors in calculated cosine of local solar illumination angles (μ_s) (Dozier et al.,

49   2022). They found errors in μ_s ranging from 0.048 to 0.117 (dimensionless) across several

50   sites for Copernicus global DEMs caused by errors in slope and aspect. The μ_s term is a

51   function (Eq. 1) of slope angle (S), slope azimuth angle or aspect (A), solar zenith angle (θ_0),

52   and solar azimuth angle (ϕ_0) – with the last two being well constrained:

53

54              $$\mu_s \;=\; \max[0,\; \cos(\theta_0)\cos(S) + \sin(\theta_0)\sin(S)\cos(\phi_0 - A)] \qquad (1)$$

55

56   Because θ_0 and ϕ_0 are calculable with low errors, the biggest contribution to errors in μ_s

57   stem from slope and aspect. Errors in μ_s increase monotonically with increasing θ_0 (e.g., sun

58   setting has high θ_0, as does solar noon in high latitude winters). This phenomenon can be

59   explained by plotting Eq. 1 for various θ_0 (Figure 1). Put simply, at higher θ_0 there is a

60   higher standard deviation in μ_s surrounding a known slope and aspect (with some temporally

61   consistent uncertainty), increasing the probability and magnitude of such an error. If one were

62   to compute standard deviations of μ_s across varying θ_0, one would arrive at similar errors of

63   μ_s presented in Dozier et al. (2022). For clarity, in Figure 1 we have highlighted an example

64   case with slope=25° +/- 4.73 and aspect=280° +/- 36.3. Example uncertainties for this

65   exercise can be found in Table 2 of Dozier et al. (2022).
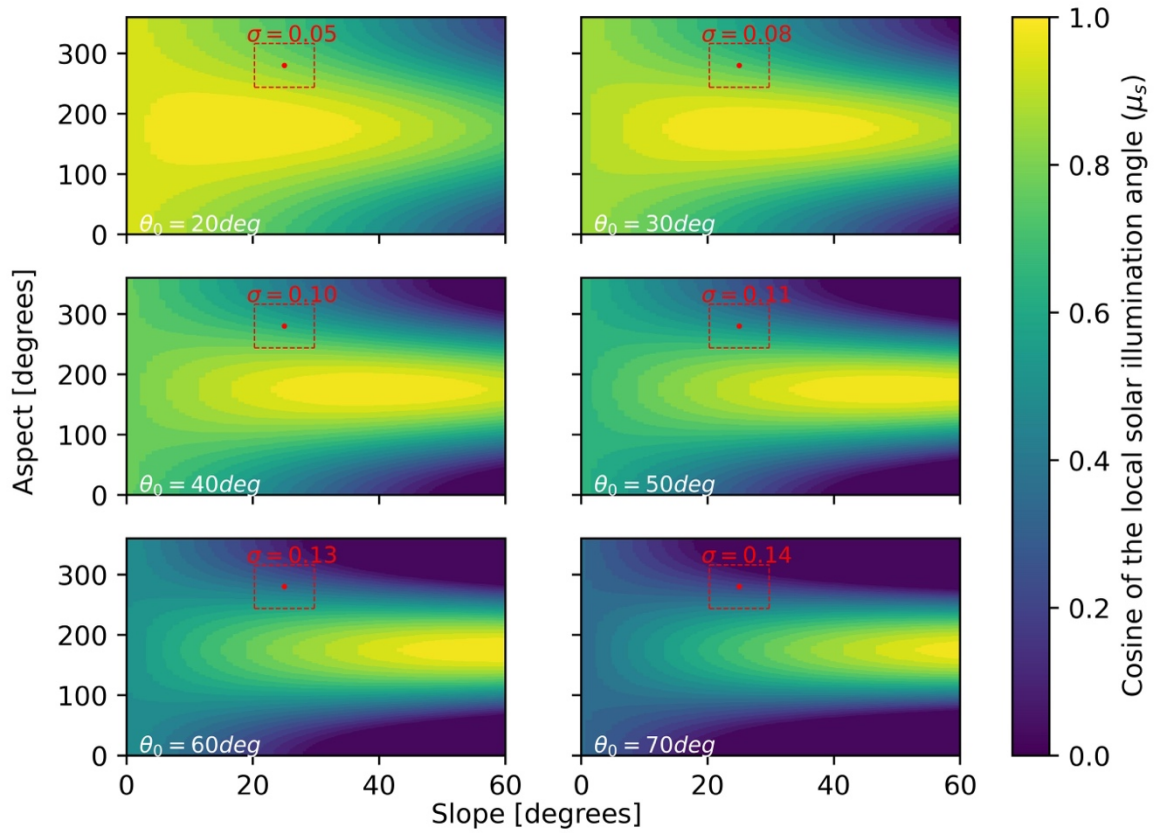
Figure 1. Cosine of local illumination angles (μ_s) as a function of slope (x-axis) and aspect (y-axis) incremented by 1°, illustrating the problem at higher latitude, and/or winter acquisitions, where standard deviation (σ) of μ_s increases monotonically with solar zenith angles (θ_0). Aspect is shown here measured clockwise from north (with north containing a discontinuity at 360 degrees). For this illustration ϕ_0 is fixed at a value of 175°. The red dots represent the example point at slope=25° +/- 4.73 and aspect=280° +/- 36.3 and are bordered by their uncertainty and the resulting σ.

74

75   Recent work has shown μ_s can be modelled using an optimal estimation framework

76 given the Top of Atmosphere (TOA) radiance observed from imaging spectroscopy (Carmon

77 et al., 2023). The authors solve for surface, atmospheric, and topographic state variables

78 simultaneously in their model. This works physically because the partition of direct to diffuse

79 light introduces a shape and magnitude effect on the TOA radiance spectra. However,

80 retrieving snow optical properties is sensitive to directional reflectance which is significantly

81 influenced by the viewing geometry and surface roughness (Bair et al., 2022), leading to

82 possible shortcomings in this method specifically for snow covered pixels. To address this

83 and expand upon this framework, we present a new method to account for terrain in snow

84 covered areas. Our method was tested on pixels with greater than 75% snow cover in three

85 different snow climates (tundra, maritime, and alpine) with spaceborne imaging spectroscopy

86 with the aim to reduce error in derived snow properties by optimally solving for topography.

87 The spaceborne results are validated against high confidence airborne spectrometer data. This

88 work directly contributes to snow property retrievals in steep terrain and/or at times of high

89 solar zenith angles for satellite imaging spectroscopy missions such as Surface Biology and

90 Geology (SBG) (Cawse-Nicholson et al., 2021), Copernicus Hyperspectral Imaging Mission

91 for the Environment (CHIME) (Celesti et al., 2022), and EnMap (Guanter et al., 2015).

92

2 Methods

2.1 Study area

For our study, we used *PRecursore IperSpettrale della Missione Applicativa* (PRISMA) imagery over three sites capturing different snow climates and solar zenith angles: San Juan Mountains (Colorado, USA, 29 April 2021, θ_0=27°), Mount Shasta (California, USA, 28 February 2021, θ_0=52°), and the Toolik area (Alaska, USA, 21 March 2021, θ_0=68°) (Figure 2). The San Juan Mountains location is considered a high alpine site located in interior continental USA with an elevation range of 2208-4129 m. The Mount Shasta site is a maritime snow climate along the western coast of USA with an elevation range of 750-4232 m. The Toolik site (elevation range = 504-1748 m) is a high-latitude tundra site, being mostly flat but with steep sections along the Brooks Range (along the southern part of the image). PRISMA, launched by the Italian Space Agency (ASI) and beginning operation on March 22, 2019, is a spaceborne imaging spectroscopy mission collecting radiance at 30 m spatial resolution across 239 bands spanning 400-2500 nm at a spectral resolution better than 12 nm across the visible-near and shortwave infrared (Cogliati et al., 2021).

To validate our method, we used four existing Airborne Visible Infrared Imaging Spectrometer-Next Generation (AVIRIS-NG) flightlines over the San Juan Mountains from 29 April 2021 (flying 1 hour after PRISMA acquisition). AVIRIS-NG collects radiance measurements at variable spatial resolution (depending on the flight altitude) across 425

bands spanning 380-2510 nm in 5nm intervals (Green et al., 2023). For this flight, data were

collected at 4 m spatial resolution. We downloaded AVIRIS-NG apparent reflectance from

National Snow and Ice Data Center (NSIDC) and observation geometry data from NASA

Search Earth Data (Skiles & Vuyovich, 2023).
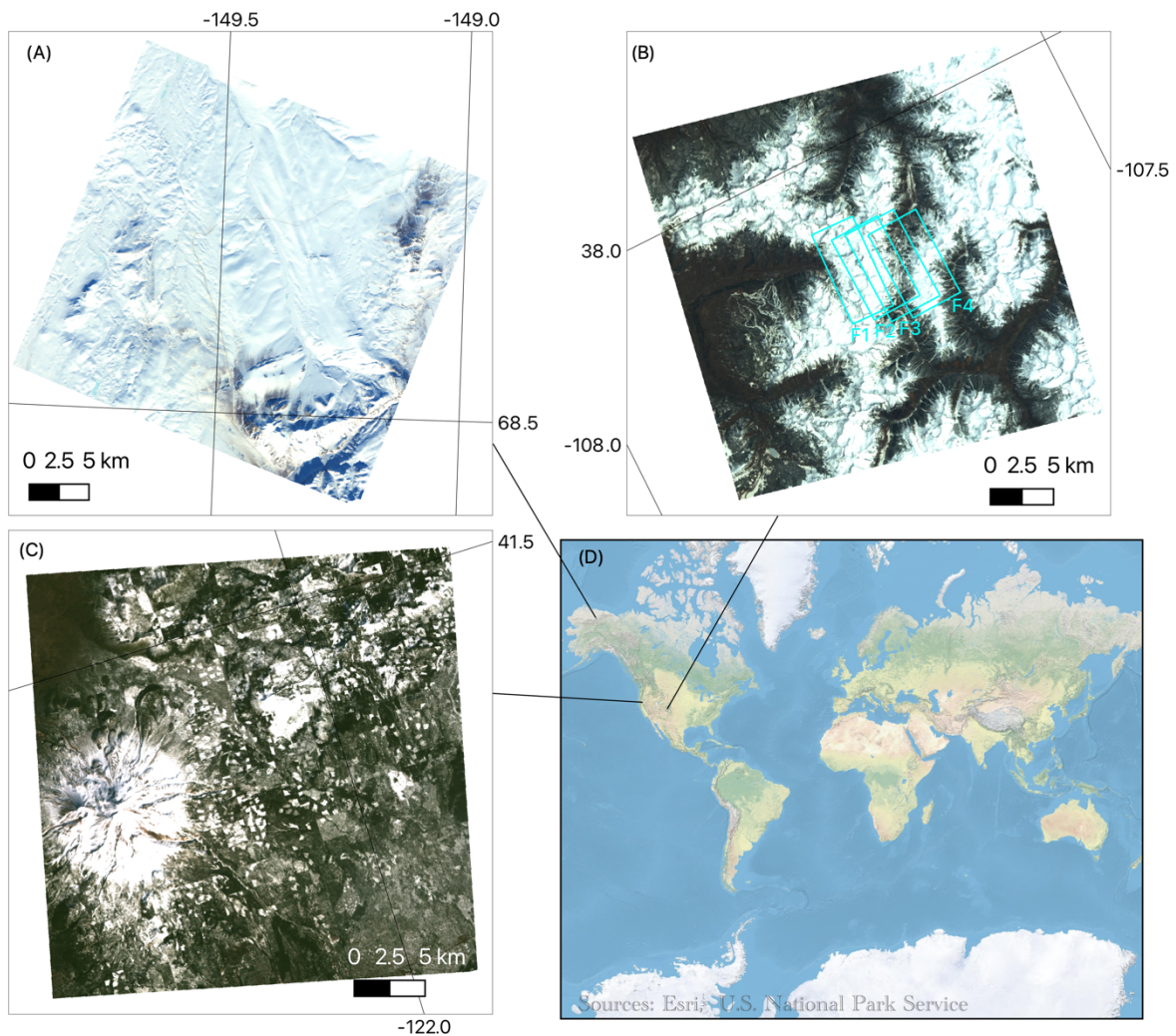
Figure 2. PRISMA true colour images for Toolik on 21 March 2021 (A), San Juan Mountains on 29 April 2021 (B), and Mount Shasta on 28 February 2021 (C). Four coincident AVIRIS-NG flightlines (F1-F4) are shown in cyan over the San Juan Mountains.

2.2 Modelling surface, atmosphere, and topography from PRISMA

The algorithmic improvements build off a workflow that estimates snow properties given PRISMA TOA radiance, titled Global Optical Snow properties via High-speed Algorithm using K-means (GOSHAWK) (Wilder et al., 2024). In short, our method uses the analytic asymptotic radiative transfer model (AART) (Kokhanovsky & Zege, 2004) coupled with libRadtran (Mayer & Kylling, 2005) to invert snow surface and atmospheric properties (Bohn et al., 2021; Dalcin & Fang, 2021), and fractional covers of mixed pixels under varied lighting conditions using non-linear numerical optimization (Bair et al., 2021). The parameters solved for in the optimization routine include fractional covers, specific surface area (SSA), light absorbing particle concentration (modelled as dust), liquid water percentage, dimensionless aerosol optical depth at 550nm, and columnar water vapor in the atmosphere. Here, we expand upon the algorithm considering recent work showing the capacity to estimate μ_s from TOA radiance (Carmon et al., 2023; Bohn et al. 2024). This idea is demonstrated in Figure 3 using fixed snow properties via AART and fixed atmosphere properties via libRadtran across the range of plausible μ_s (i.e. 0 to 1). Like the findings in

136   Carmon et al. (2023), Figure 3 shows that μ_s controls both the spectral shape and magnitude

137   of observed TOA radiance with the effect varying across wavelengths. The greatest shape

138   effect can be seen in the visible spectrum (roughly 400-700 nm) due to the magnitude of the

139   diffuse irradiance. In combination with the magnitude and shape shift, this parameter

140   becomes solvable during optimization due to its strong separability – especially when

141   considering the entire spectrum data from a hyperspectral remote sensing source such as

142   PRISMA. It is important to note that μ_s impacts both the AART estimation of snow

143   reflectance and libRadtran estimation of incoming solar irradiance.

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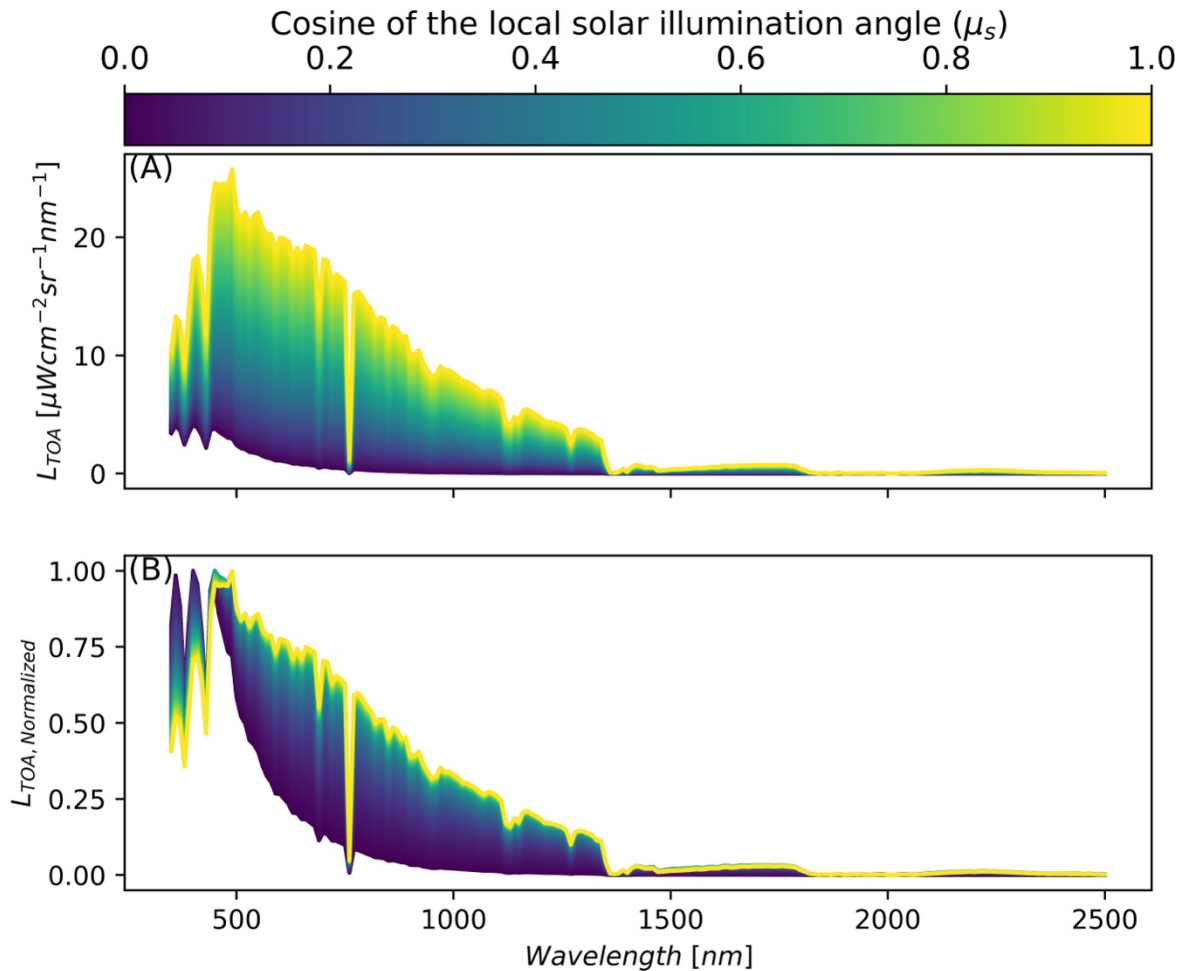
Figure 3. Synthetic data showing change in magnitude (A) and shape (B) of top of

atmosphere radiance (L_{TOA}) with respect to changing local solar illumination angle (μ_s) for

fixed snow surface state variables modelled with AART, and fixed atmospheric state

variables modelled with libRadtran (viewing geometry was fixed as well). State variables and

solar/view geometry were based on a PRISMA acquisition over southern Idaho on 8

152 December 2022. Figure (B) shows normalized radiance with respect to peak TOA radiance

153 across wavelengths to highlight the change in shape.

154

155   However, if we were only to optimize μ_s, the other key terms, local viewer zenith

156 angle (μ_v) and local phase angle (ξ) in the AART formulation for bidirectional reflectance of

157 snow (Eq. 2) (Kokhanovsky & Zege, 2004; Kokhanovsky et al., 2021a) would remain

158 constant from the available DEM (i.e., μ_s, μ_v, ξ are all derived from DEM),

159

160 $$r_{snow}(\mu_s, \mu_v, \xi, \lambda) = r0(\mu_s, \mu_v, \xi) \; a_{snow}(\lambda)^f \qquad\qquad (2)$$

161

162 where $r0$ is the reflection function of a semi-infinite non-absorbing snow layer (Tedesco &

163 Kokhanovsky, 2007), α_{snow} is the spherical albedo [plane albedo can be computed using (26)

164 in Kokhanovsky et al. (2021a)], f is the escape function (Kokhanovsky et al., 2021a), and

165 r_{snow} is the bidirectional reflectance of snow. Keeping other terms $\mu_v \; and \; \xi$ the same are

166 problematic because snow reflectance is poorly approximated as a non-Lambertian surface

167 (Leroux & Fily, 1988), and the outcome will be greatly influenced by $\mu_v \; and \; \xi$. Therefore, to

168 incorporate solving for $\mu_s, \mu_v, \; and \; \xi$ from TOA radiance into the algorithm, we instead elect

169 to optimally solve for cos(aspect) ~~(i.e., "northness")~~ and sin(aspect) ~~(i.e., "eastness")~~ (Table

170 1).

171

Table 1. Parameter space and initial vectors used in numerical optimization for PRISMA

data.

Parameter [unit]	Definition	Feasible Range	Initial State	Type
f_{snow} [%]	Fractional snow in the mixed pixel	[0, 100]	10	Surface
f_{shade} [%]	Fractional shade in the mixed pixel	[0, 100]	20	Surface
f_{LC1} [%]	Fractional cover of endmember 1 (based on land cover value at pixel)	[0, 100]	50	Surface
f_{LC2} [%]	Fractional cover of endmember 2 (based on land cover value at pixel)	[0, 100]	20	Surface
SSA [m^2 kg^{-1}]	Specific surface area (SSA)	[2, 156]	40	Surface
LAP [μg g^{-1}]	Concentration of light absorbing particles, LAP, modelled as dust (PM-2.5).	[0, 145]	0	Surface
Liquid water [%]	Percentage of liquid water on the snow surface	[0, 50]	2	Surface
AOD 550 [%]	Dimensionless Aerosol Optical Depth (AOD) at 550 nm	[1,100]	10	Atmospheric
H_2O [mm]	Columnar water vapor in the atmosphere	[1,50]	1	Atmospheric
sin(aspect)~~Eastness~~	~~sin(aspect)~~Continuous variable to define aspect with respect its east/west direction	[-1,1]	Variable	Topographic
cos(aspect)~~Northness~~	Continuous variable to define aspect with respect its north/south direction~~cos(aspect)~~	[-1,1]	Variable	Topographic

174

175 Aspect can be solved during optimization by using the atan2 function. We chose to use this

176 method because sin(aspect)~~eastness~~ and cos(aspect)~~northness~~ are continuously differentiable,

177 and therefore, are suited for numerical optimization methods, whereas aspect is discontinuous

178 at north (using the convention of 0 and 360 degrees as north). We then can use this optimal

179 aspect to estimate μ_s (Eq. 1), μ_v, and ξ. This directly impacts Eq.2 and Eq. ~~5~~ 3 (formulation

180 of incoming solar energy in the model) (Picard et al., 2020),

181

182 $$E(\lambda) = \psi\mu_s E(\lambda)_{dir} + V_\Omega E(\lambda)_{diff} + \left[\left(1 + \frac{\cos(S)}{2} - V_\Omega\right) r(\lambda)_{surf} E(\lambda)_{diff}\right] \quad (\text{\sout{5}\underline{3}})$$

183

184 where E is total incoming irradiance, ψ is binary shade or no shade, E_{dir} and E_{diff} are the

185 direct and diffuse irradiance, respectively, V_Ω is the sky view factor (Dozier, 2022), and r_{surf}

186 is the reflectance of nearby terrain (which is assumed to be equal to the pixel itself). The term

187 E is solved within our non-linear numerical optimization method as described in Wilder et al.

188 (2024). This was modelled incorrectly in Wilder et al. (2024); however, this was corrected in

189 this paper where only diffuse irradiance is used in the 3rd term in Eq. 5. Also, adding in the

190 two extra parameters (sin(aspect) ~~eastness~~ and cos(aspect)~~northness~~) in our updated

191 optimization scheme did not change our run time significantly. Caution is advised against

192 solving for slope *and* aspect in the inversion due to the non-unique solution space (Figure 1);

193 however, only considering aspect ensures unique solutions of aspect, μ_s, μ_v, and ξ. We

194 chose aspect because of its greater impact on determining partition of direct and diffuse

195 illumination and has been found to be more impactful to errors associated with snow property

196 retrieval (Donahue et al., 2023). In this study we used estimate of total ozone column as input

197 into creating the libRadtran look up table specific for each image. We used the average

198 weekly ozone over the bounds of the image from Sentinel-5P NRTI O3: Near Real-Time

199 Ozone dataset. This approach serves an improvement over Wilder et al. (2024), where ozone

200 was fixed at 300 Dobson Units.

201

202 **2.3 Estimating snow properties from AVIRIS-NG for validation**

203 Due to the fine signal to noise ratio and spatial resolution of AVIRIS-NG, we treated

204 the dataset as the ground reference. It also captured a similar spectral range to PRISMA

205 which made it a suitable comparison dataset. The main assumption here is that AVIRIS-NG

206 pixels at 4 m are relatively homogenous and are either snow or no-snow – which may not

207 always be the case. This could be a potential source of uncertainty in our analysis. To select

208 snow-covered pixels, we solved for NDSI (Normalized Difference Snow Index) using bands

209 at 600 nm and 1500 nm. We limited our retrieval of snow properties for NDSI greater than or

210 equal to 0.90 (Painter et al., 2013). A common approach to retrieve snow grain size from pure

211 snow pixels is to apply the scaled band area algorithm (Nolin & Dozier, 2000); however, it is

212 recognized that the large presence of liquid water is a limitation. The maximum air

213   temperature of 10.8° C on the day of the image at the San Juan Mountains site indicated that

214   elevated liquid water at the surface was probable (Center for Snow and Avalanche Studies,

215   2023). Additionally, reflectance spectra appeared to be shifted along the x-axis (wavelength)

216   due to the presence of liquid water. Therefore, we used constrained non-linear numerical

217   optimization to model apparent snow reflectance with AART by allowing fractional snow,

218   fractional shade, liquid water, and SSA to vary. We did not include rock or forest

219   endmembers in this formulation, assuming the 4 m pixels are relatively homogenous as

220   previously stated. Topographic incident angles were held constant based on the 4 m

221   resolution DEM provided by AVIRIS-NG. We minimized Root Mean Square ~~Error~~

222   Difference (RMSD~~E~~) between observed-apparent and modelled-apparent snow reflectance

223   from AART wavelengths in the range, 1000-1250 nm. This range has high ice absorption and

224   limited impacts from atmospheric interference and LAP (Miller et al., 2016). The presence of

225   liquid water was included in our analysis by means of the composite refractive index of water

226   and ice (Donahue et al., 2022; Hale & Querry,1973; Warren & Brandt, 2008). We assumed

227   similar grain shape assumptions for both PRISMA and AVIRIS-NG, and that if there is a bias

228   due to this it should be consistent between the two datasets in our analysis.

229

2.4 Comparing modelled snow properties

230

231   The algorithm was used in two different modes: 1) static topography based on the

232   Copernicus DEM (hereon called "*static*"); and 2) solved topography based on the algorithm

233   updates (hereon called "*radiance*"). To compare the accuracy of PRISMA derived SSA and

234   liquid water, we resampled the AVIRIS-NG optical property results (SSA and LWC) to

235   match the PRISMA resolution (30 m) and extents by using bilinear interpolation. Then, we

236   sampled all valid pixels where PRISMA and AVIRIS-NG had snow. We then computed r-

237   ~~pearson~~Pearson correlation coefficient, ~~m~~Mean ~~b~~Bias, and RMS~~DE~~ for the radiance and

238   static methods (with respect to AVIRIS-NG). Finally, we used Copernicus derived slope and

239   aspect maps to determine where the largest errors were occurring on the landscape to

240   compare with the theoretical basis presented in Figure 1. We do this by using the mean

241   absolute difference with respect to μ_s . We expected to see higher differences in north facing

242   aspects (i.e., μ_s approaches 0), and where θ_0 was higher. To test the interaction with θ_0 more

243   fully, we extended the analysis to Mount Shasta, CA, and Toolik, Alaska, where no in *situ*

244   *data* existed. We compared the modelled properties between the radiance and static methods

245   to assess how these assumptions impacted results for these types of data at 30 m scale.

246

2.5 Comparing DEM and radiance derived μ_s

To ensure the resulting radiance derived μ_s were valid we downloaded the best available validation data sources for comparison. For the San Juan and Shasta sites, we collected DEM products at 1 m spatial resolution and collected 5 m spatial resolution DEM for the Toolik site (U.S. Geological Survey, 2019; U.S. Geological Survey, 2022). Then, we computed slope, aspect, solar zenith angle, and solar azimuth angle for all pixels to compute μ_s at the native resolution (Eq. 1). Then, we used bilinear interpolation to resample the 1 m and 5 m products to 30 m to exactly match the extents and resolution of our PRISMA images. We would like to acknowledge that while these are the best freely available datasets for our images, they still do not capture the true snow-on topography, and instead are a representation of the "snow-free" surface. We compared matching pixels to determine RMSDE, r-Pearson correlation coefficientR², and Mean mean Biasbias. Pixels that were marked as shadow from ray tracing were excluded from this comparison.

3 Results

3.1 Validation using AVIRIS-NG data over the San Juan Mountains

Over the area of the flightlines, AVIRIS-NG estimated mean SSA = 18.0 +/- 8.3 m^2 kg^{-1}, PRISMA radiance method estimated mean SSA = 19.6 +/- 5.8 m^2 kg^{-1}, and PRISMA static method estimated mean SSA = 22.0 +/- 12.1 m^2 kg^{-1}. When comparing the SSA performance

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266   over each pixel to the AVIRIS-NG flightlines (Figure 4) we found the PRISMA radiance

267   method (r=0.43; RMSD=8.0 m² kg⁻¹; bias=+1.7 m² kg⁻¹; n=36,412) performed slightly better

268   than the static method (r=0.23; RMSD=13.6 m² kg⁻¹; bias=+4.0 m² kg⁻¹; n=36,412) for SSA.
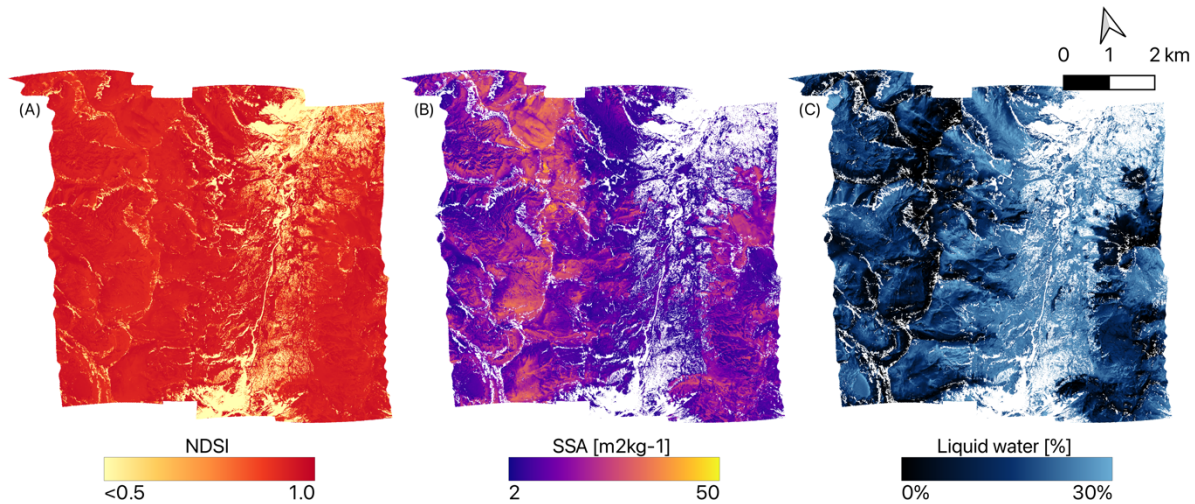
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270   **Figure 4.** Snow properties computed from AVIRIS-NG (4 m spatial resolution) on 29 April

271   2021 including NDSI (A), SSA (B), and LWC (C) for the San Juan Mountains site.

272

273   There was not a significant improvement in liquid water estimation between radiance (r=0.67;

274   RMSD=10%; bias=-8%; n=36,412) and static (r=0.67; RMSD=10%; bias=-9%; n=36,412).

275   Furthermore, it appeared that there was a consistent liquid water bias of 8 to 9%, hinting that

276   more melt had occurred during the AVIRIS-NG flights. As previously noted, the temperatures

277   were well above ~~freezing~~ 0°C during the overpass of AVIRIS-NG and occurred roughly 1 hour

278   later in the day compared to the PRISMA acquisition. This most likely explains the higher

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279 liquid water and lower SSA observed by AVIRIS-NG. We further tested this by masking out

280 areas where AVIRIS-NG liquid water content was greater than 0.1%, to establish areas where

281 low amounts of melt occurred between the two acquisitions. We found that performance of

282 PRISMA static (RMSD=14.2 m^2 kg^{-1}; rRMSD=49%; n=181) and radiance (RMSD=6.9 m^2 kg$^-$

283 ^{1}; rRMSD=23%; n=181) methods were more accurate for these areas. The radiance method

284 performed slightly better, suggesting a modest 25% improvement in accuracy for SSA over the

285 static method when considering pixels that were less impacted by melt.

286   Additionally, comparing all pixels we found improvement from radiance occurred

287 mostly on steep, north facing aspects (e.g., when μ_s approached 0). We found the absolute

288 residual increased as μ_s approached zero for the static method (r = -0.47; p<0.01), while this

289 relationship was diminished nearly by a factor of 5 for the radiance method (r = -0.10; p<0.01)

290 (Figure 5.A). These errors were caused by incorrect terrain information in the inversion, where

291 inversion error increased proportionately in the static method (Figure 5.B).
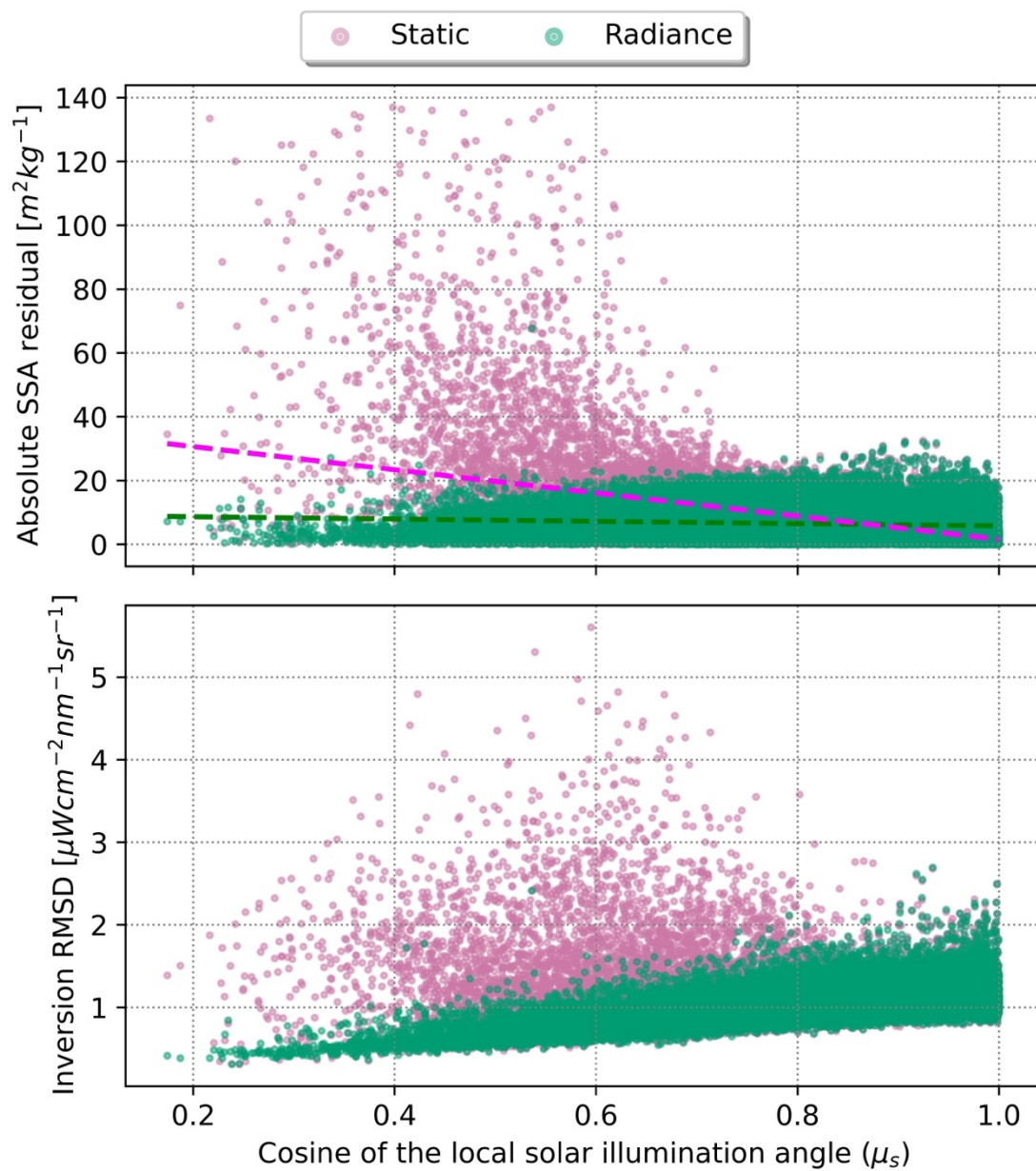
Figure 5. Absolute difference in modelled SSA when compared to AVIRIS-NG for radiance

method (green) and static method (pink) respect to μ_s (A) and resulting RMS~~DE~~ from the

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inversion from PRISMA with respect to μ_s (B). Error in the static method increases significantly when μ_s approached zero (r = -0.47; p<0.01); however, the difference was less noticeable in the radiance method (r = -0.10; p<0.01).

3.2 Comparing radiance and static methods between sites

On average across each of the images, radiance and static methods provided similar retrieved parameters within less than one standard deviation (Table 2). In general, this means there is not a significant difference at the 30 m scale for computing parameters such as SSA and broadband albedo (BA) when considering the entire image. Interestingly when terrain is fixed, the static model compensated for incorrect illumination by increasing the aerosol optical depth (thereby reducing the amount of direct solar radiation). Investigating the errors more closely, we found much larger differences in retrieved properties where μ_s approached 0 (Figure 6). The difference in distributions matched closely to the theoretical demonstration (Figure 1) and is most likely associated with the standard error of slope and aspect from Copernicus DEM given the illumination conditions. This result also demonstrates the difference between the two methods had the biggest impact for images where θ_0 was high, resulting in potentially inaccurate retrievals that impact both surface and atmospheric state variables on relatively mild slopes.

314   **Table 2.** Image-wide statistics comparing derived properties between the two methods (static

315   vs. radiance) processing the PRISMA imagery for all three sites.

| Site	PRISMA Method	Mean SSA [m² kg⁻¹]	Mean Broadband Albedo	Mean Liquid water [%]	Mean AOD at 550 nm	Mean water column vapour [mm]
San Juan	Static	23.3 +/- 14.9	0.79 +/- 0.03	3.5 +/- 4.8	0.05 +/- 0.13	6.7 +/- 1.1
	Radiance	19.6 +/- 5.9	0.78 +/- 0.03	3.9 +/- 5.0	0.01 +/- 0.01	6.8 +/- 0.3
Shasta	Static	11.0 +/- 6.0	0.77 +/- 0.04	1.6 +/- 3.3	0.04 +/- 0.10	7.6 +/- 1.3
	Radiance	10.7 +/- 6.2	0.77 +/- 0.05	1.9 +/- 3.8	0.01 +/- 0.04	7.7 +/- 1.1
Toolik	Static	30.1 +/- 9.6	0.85 +/- 0.02	0.0 +/- 0.0	0.02 +/- 0.03	1.0 +/- 0.4
	Radiance	27.7 +/- 7.9	0.84 +/- 0.02	0.0 +/- 0.0	0.01 +/- 0.01	1.0 +/- 0.2

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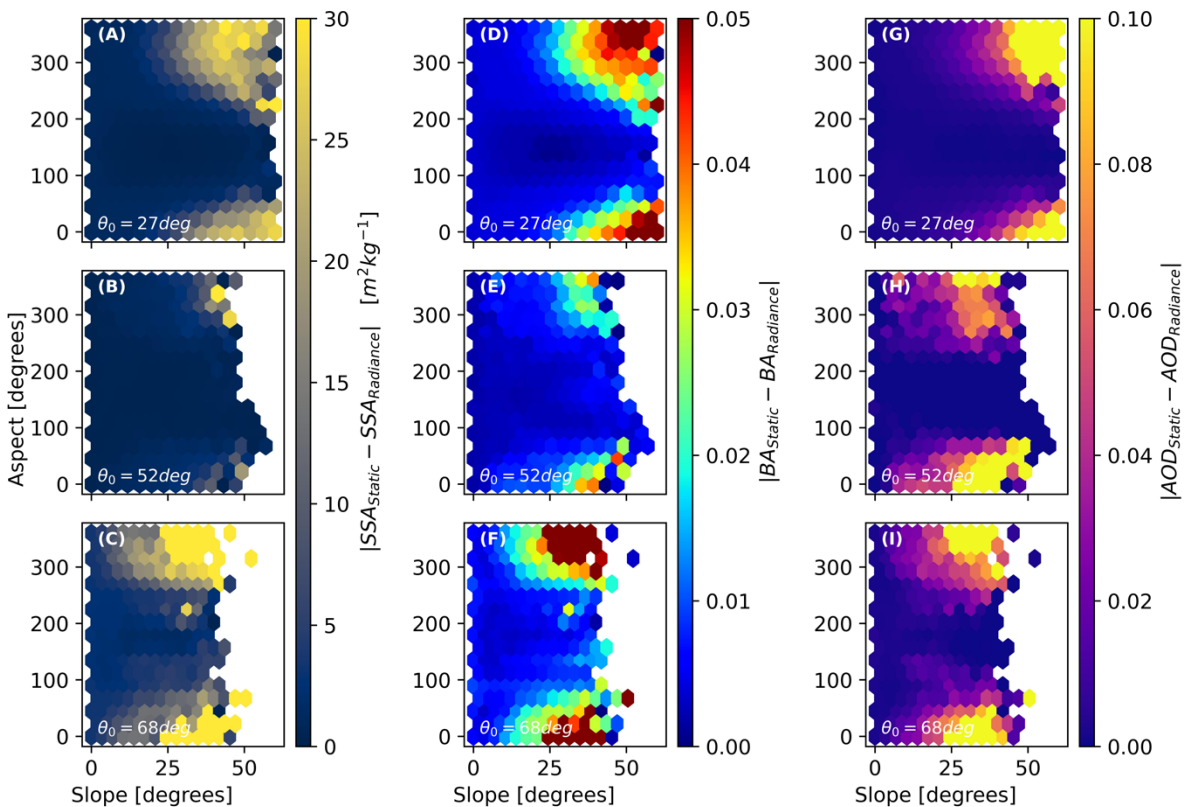
Figure 6. 2D Histogram plots showing absolute difference in SSA (left), broadband albedo (middle) and AOD (left) with respect to slope and aspect across the entire dataset. In this figure absolute difference is calculated as |Static – Radiance|. This is shown for the San Juan Mountains site (A,D,G), Shasta site (B,E,H), and Toolik site (C,F,I). The average solar zenith angle (θ_0) is shown for reference on each panel.

Putting this into spatial context (Figure 7), San Juan site had 37% of pixels (135.3 km^2) with an absolute difference in BA (|δBA|) >= 0.01 and 14% pixels (49.9 km^2) with |δBA| >=

327   0.02. Shasta site had 30% of pixels (16.7 km^2) with |δBA| >= 0.01 and 9% pixels (5.1 km^2)

328   with |δBA| >= 0.02. Toolik site had 40% of pixels (325.3 km^2) with |δBA| >= 0.01 and 8%

329   pixels (66.6 km^2) with |δBA| >= 0.02.
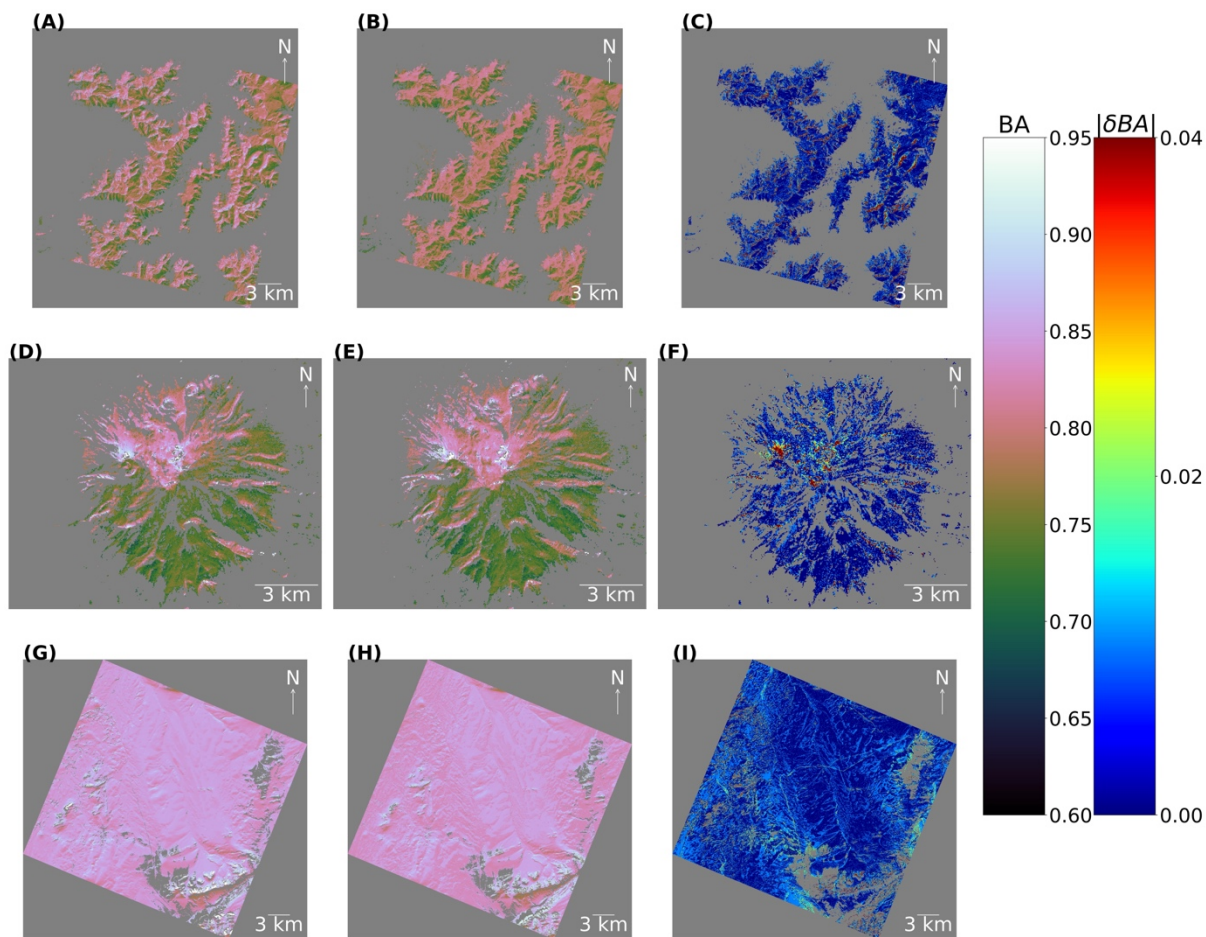
330

331   **Figure 7.** Modelled broadband snow albedo (BA) for San Juan Mountains site (A-C), Shasta

332   Mountain site (D-F), and Toolik site (G-I). Left column represents BA from static method,

333   middle column represents BA from radiance method, and right column represents absolute

334   difference in BA ($|\delta BA|$). Dark grey colour symbolizes data that is not a value.

335

336   Median $|\delta BA|$ for all sites with respect to $\boldsymbol{\mu_s}$ general increased as $\boldsymbol{\mu_s}$ approached zero

337   (Figure 8). For example, for the San Juan site, median $|\delta BA|$ ranged from 0.03 to 0.00 across

338    $\boldsymbol{\mu_s}$. For the Shasta and Toolik sites, median $|\delta BA|$ ranged from 0.02 to 0.00 across $\boldsymbol{\mu_s}$. This

339   relation was non-linear and depended on the site and illumination conditions. This analysis

340   demonstrates the levels of uncertainty potentially left in for retrievals relying on static, non-

341   coincident DEMs. This shows quantitatively the improvements to snow broadband albedo at

342   30 m scale by using radiance-based approach to be relatively small for well-lit slopes – on the

343   order 0-1%. While shaded slopes may have errors in snow broadband albedo on the order of

344   1-3%. Interestingly for the Toolik site, $|\delta BA|$ also increased as $\boldsymbol{\mu_s}$ approached one.

345
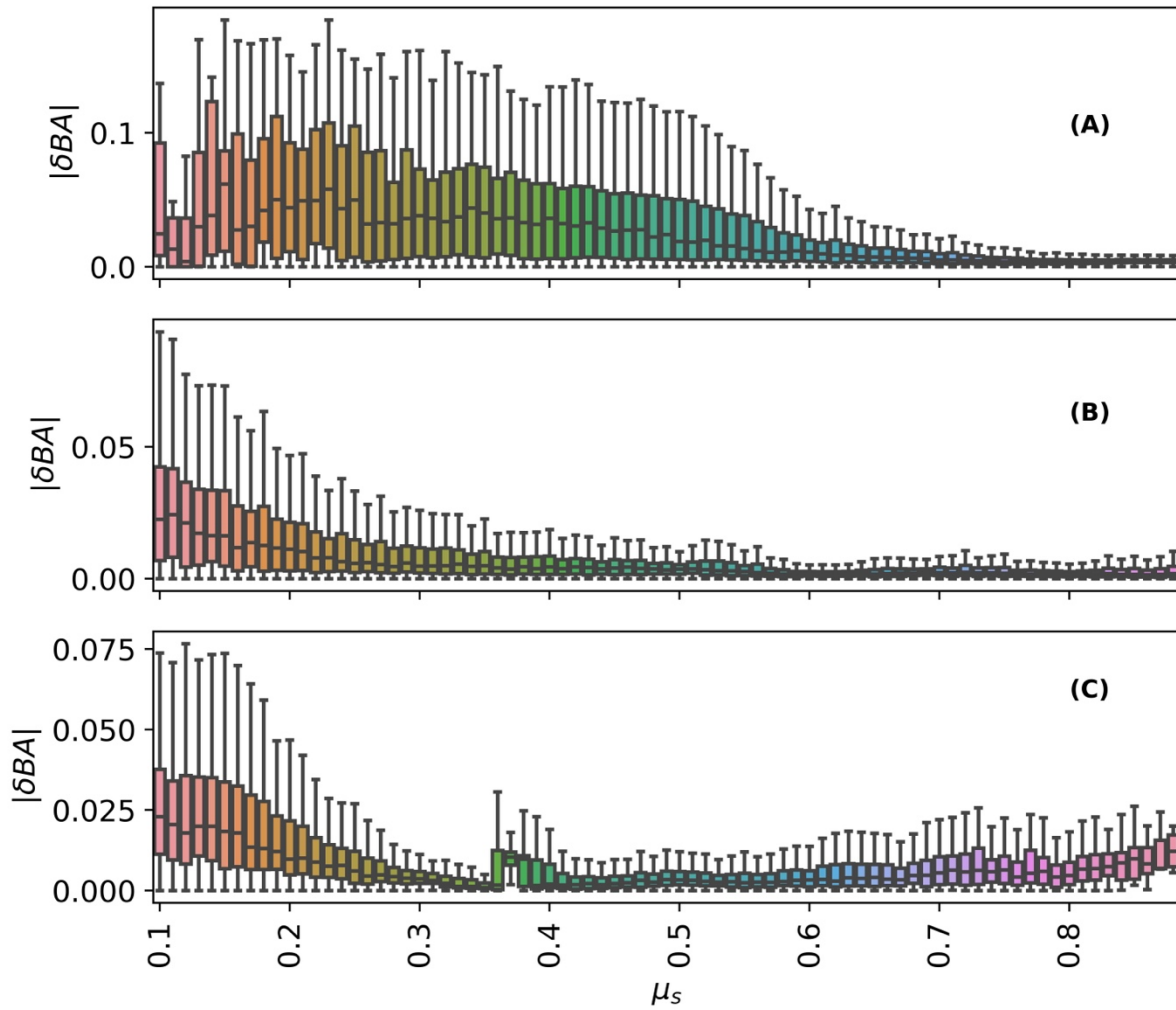
**Figure 8.** Modelled absolute difference in broadband albedo ($|\delta\text{BA}| = |\text{BA}_{\text{Static}} - \text{BA}_{\text{Radiance}}|$) for San Juan (A), Shasta (B), and Toolik (C). Note these boxplots were created by rounding $\boldsymbol{\mu_s}$ to the nearest hundredth place.

3.3 Comparing DEM and radiance derived μ_s

351

352   At the 30 m pixel scale, Copernicus DEM derived μ_s had similar overall performance

353   to radiance derived μ_s (Figure 9), with Copernicus DEM derived μ_s having slightly higher

354   performance. For example, for the San Juan site, RMSD only varied by 0.006 between the

355   two methods. ==Similarly, the ~~R²~~r for Copernicus derived μ_s was 0.9486, while the radiance==

356   ==derived μ_s was slightly lower at 0.9183.== This similar overall performance was common

357   amongst the three sites. We found the average bias for radiance derived μ_s was generally

358   closer to zero (+/- 0.01), and did not show a strong negative or positive direction.
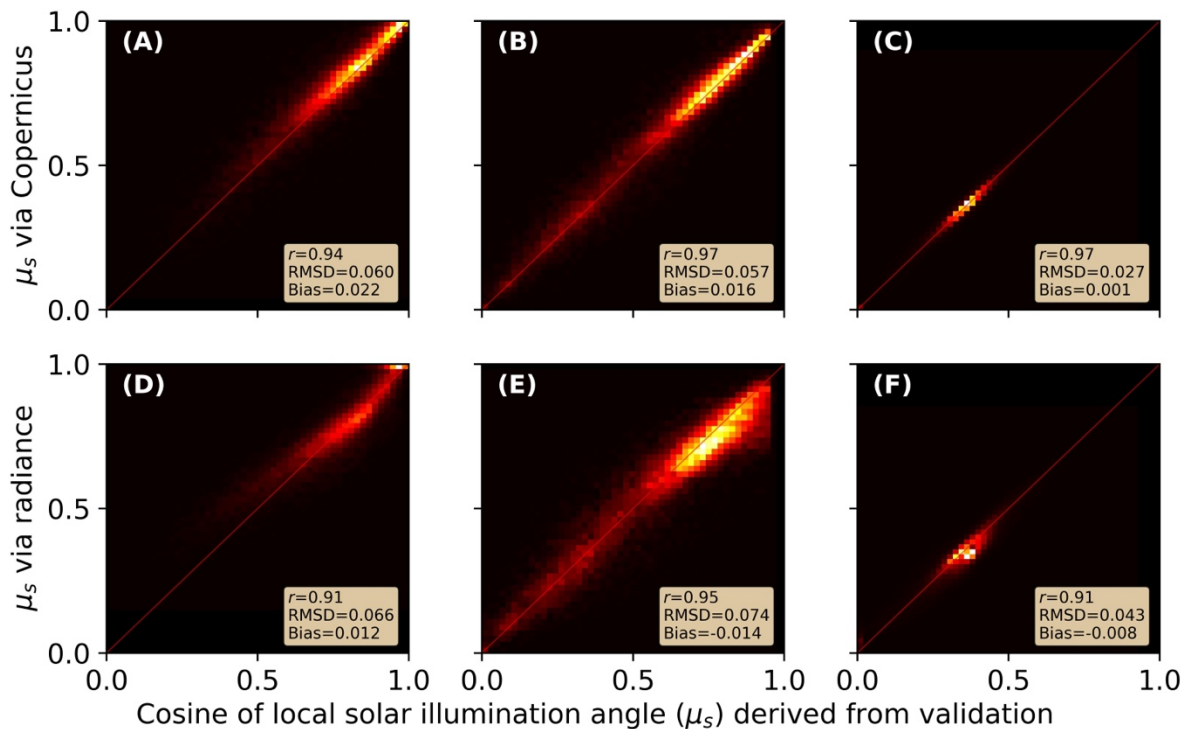
359

360 **Figure 9.** Comparing μ_s at 30 m pixel scale derived from radiance and Copernicus against high

361 resolution DEM for San Juan site (A,D), Shasta site (B,E), and Toolik site (C,F).

362

363 **4 Discussion**

364 **4.1 Radiance derived DEMs may replace coincident DEMs and contain information**

365 **related to surface roughness**

366   Derivative slope and aspect maps are prone to errors at 30 m spatial resolutions

367 (Dozier et al., 2022). This is relevant for derived snow products from upcoming missions

368 such as SBG and CHIME which will rely on topographic information to calculate optical

369 properties like snow albedo. These errors can be inherent to the DEM itself, or a product of

370 spatial and/or temporal misalignments (Carmon et al., 2023). Our modelled |δBA| with

371 respect to the non-coincident DEM was similar to work by Donahue et al. (2023), who found

372 slightly higher uncertainties of broadband albedo (ranging from -10 to 10%) for their

373 investigation on Place Glacier, British Columbia, Canada. With the surface and roughness

374 undergoing dramatic change on glaciers throughout a given season, using this radiance-based

375 approach may be especially impactful for improving estimates over glaciers.

376   Snow surface roughness has long been a challenging issue in modelling snow

377 properties from space where the solar incidence angle at high spatial resolution for snow-on

378 DEM is not well known (Bair et al., 2022). Previous research found radiance derived μ_s from

379 airborne imaging spectroscopy showed a negative bias and postulated this could be due to

380 within-pixel topography, shadows, and surface roughness (Carmon et al., 2023). Since a bi-

381 directional reflectance function (BRDF) model was not used in their study, it then would be

382 plausible for the optimal μ_s to compensate for these effects. Interestingly when using a BRDF

383 model in our study (i.e., AART) and solving for aspect optimally (therefore informing μ_s, μ_v,

384 and ξ) we did not find a strong bias – negative or positive. Although, we did not take surface

385 roughness measurements, and therefore do not know to the extent this impacted our study.

386 Within-pixel shadows, textures, and surface roughness remain difficult to validate, and we

387 were unable to achieve this in our study. Future work interested in further understanding this

388 radiance-based approach may investigate how such approaches interact with micro-scale

389 topography through ground measurements such as terrestrial and airborne lidar.

390

391 **4.2 Next steps in possibly improving this radiance-based approach**

392   While we solved for a few terrain parameters in this study we did not entirely remove

393 the use of the DEM from the radiance method. The elevation from global DEMs has a much

394 higher confidence than its derivative products (Dozier et al., 2022). Therefore, we used these

395 values to inform our atmospheric routine, as well as our shadow casting ray tracing module

396 (Wilder et al., 2024). Additionally, we used the method presented in Dozier (2022) for

397 estimating the sky view factor (V_Ω) based on nearby terrain and the pixel itself. This factor

398   could potentially be problematic but was cited as being not as impactful as μ_s in propagating

399   error (Dozier et al., 2022). Therefore, we elected to use V_Ω derived from the static Copernicus

400   DEM. However, this could be an area for future improvement, especially in very steep terrain

401   where V_Ω becomes small. It is not advised to attempt to add V_Ω directly into the optimization

402   routine presented in this study, as it is a function of pixel slope and aspect, and therefore,

403   altering V_Ω and aspect together would create invalid solutions. Finally, we used a static value

404   for slope derived from Copernicus DEM. The slope influences the μ_s term, but also

405   influences the passive radiation from nearby slopes. Ultimately, we concluded that aspect had

406   the largest impact on changing μ_s (Figure 1), as well as large RMSDE reported in previous

407   work (Dozier et al., 2022; Donahue et al., 2023), and thus was the focus of our study. Caution

408   is advised in including both slope and aspect together, as non-unique solution space for μ_s

409   may cause the optimization outputs to become invalid. In summary, elevation, V_Ω, and slope

410   remain static in our current implementation. Future work may explore other algorithmic

411   choices to further remove, or improve, static DEM parameters.

412   Another consideration for improving this method is the inclusion of total column

413   ozone into the optimization. Previous research has been able to use TOA snow reflectance

414   data to retrieve reliable estimates of ozone (Kokhanovsky et al., 2021b). In our paper, we

415   elected for a simpler approach to first investigate the impacts of including terrain in the

416   optimization. In this paper we input a fixed ozone for the entire image based on coincident

417 Sentinel-5 measurements. However, it should be stated that ozone impacts a similar spectral

418 range to μ_s (Figure 10). It therefore may be beneficial to include ozone in the atmospheric

419 lookup-table (e.g., MODTRAN, libRadtran) to enable optimization of ozone as well. This

420 may be beneficial in building more realistic radiance-based methods.
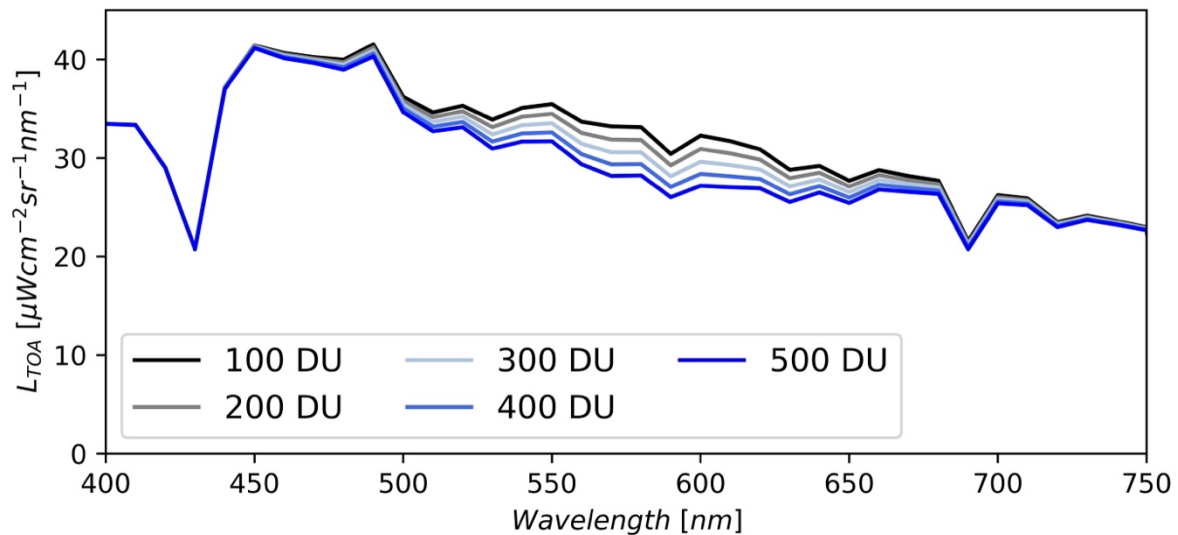
421

422 **Figure 10.** Synthetic data showing change in magnitude of top of atmosphere radiance (L_{TOA})

423 with respect to changing total column ozone for fixed snow surface state variables modelled

424 with AART, and other fixed atmospheric state variables modelled with libRadtran. Reference

425 data is based on PRISMA image taken over southern Colorado. Note units of total column

426 ozone are shown in in Dobson Units (DU).

427

428   Finally, future studies should investigate including improvements to BRDF models of

429   snow (Mei et al., 2022). For example, recent work by Kokhanovsky et al. (2024) has

430   proposed the use of a two-layer model which may be especially useful for vertically

431   heterogenous snowpacks. Their method has been tested using EnMAP data and may easily be

432   transferable to other sensors. The current AART method we used in our paper does not

433   account for these layers, and instead assumes an optically thick, homogenous snowpack. To

434   validate both AART, the new layered approach, and future BRDF models, snow pit (i.e.,

435   vertical profile) measurements of SSA (e.g., Meloche et al., 2023) become essential in

436   ensuring models accurately account for diverse layering of snow.

437

438   **4.3 Big picture implications of the radiance-based approach**

439   This research responds to the objectives stated in "*Thriving on our changing planet: A*

440   *decadal strategy for Earth observation from space",* to improve biogeophysical modelling at

441   scales driven by topography (National Academies of Science, Engineering, & Medicine,

442   2018), enabling more accurate snow property retrievals in the cryosphere under challenging

443   illumination conditions. Our work presented on solving terrain where DEM data are not

444   available, or reliable, may serve to accelerate improvements to satellite remote sensing tools

445   to monitor and model at both the regional global scale (Sturm et al., 2017), at a critical

446   juncture in time where northern latitudes are changing fast under a warming climate. This

447   includes Earth's glaciers, where radiance-based method may have the largest improvements

448   over static approaches. Our research is complimented by other recent works which show

449   promise in including terrain in the inversions (Bohn et al., 2024; Bohn et al., 2023; Bair et al.,

450   2024; Carmon et al., 2023)

451   We recommend additional coincident AVIRIS-NG flights with spaceborne imaging

452   spectroscopy datasets to further this work. As we have shown for the San Juan Mountains

453   site, for particularly warm days, images that are separated by longer than an hour may exhibit

454   drastically different SSA and liquid water content. As shown in this paper, this creates an

455   issue when trying to validate improvements to retrieval algorithms.

456

457   **5 Conclusion**

458   In this study we used existing PRISMA L1 TOA imagery to demonstrate the

459   improvements in modelling snow optical properties when explicitly modelling the terrain in

460   the inversion. This would especially be true for areas where the surface undergoes rapid

461   change, such as glaciers. This new method is especially useful for steep mountain terrain

462   and/or high latitudes where illumination conditions are suboptimal. The θ_0 (solar zenith

463   angle) was relatively low for the San Juan Mountains site in our study, and thus represents a

464   lower bound of the improvement in accuracy one could expect. This disparity was

465   demonstrated further for the Mount Shasta and Toolik sites when θ_0 was larger (i.e. a greater

466   difference in retrieved properties due to more challenging solar and sensor geometry). Even

467   for the relatively flat Toolik site, we showed that correctly accounting for incidence angles

468   can impact snow properties when θ_0 is large. Future work may look to build from this

469   radiance-based approach to enable better quantification of snow properties at scales impacted

470   by topography.

471

472   *Code Availability.* https://github.com/cryogars/goshawk

473   *Author contributions.* B.W. created the GOSHAWK algorithm and updates herein, decided

474   on experiment set-up, and performed the subsequent analysis, as well as being the main

475   article writer. J.M., J.E. and N.G. provided ideas, comments, and supervised the work.

476   *Competing interests.* The contact author has declared that neither they nor their co-authors

477   have any competing interests.

482

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