

1 Improved Snow Property Retrievals by Solving for Topography in the Inversion of At-
2 sensor Radiance Measurements

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9
10 **Abstract**

11 Accurately modelling optical snow properties like snow albedo and specific surface
12 area (SSA) are essential for monitoring the cryosphere in a changing climate and are
13 parameters that inform hydrologic and climate models. These snow surface properties can be
14 modelled from spaceborne imaging spectroscopy measurements but rely on Digital Elevation
15 Models (DEMs) of relatively coarse spatial scales (e.g. Copernicus at 30 m), which degrade
16 accuracy due to errors in derived products – like such as slope and aspect. In addition, snow
17 deposition and redistribution can change the apparent topography and thereby static DEMs
18 may not be considered coincident with the imaging spectroscopy dataset. Testing in three
19 different snow climates (tundra, maritime, alpine), we established a new method that

20 simultaneously solves snow, atmospheric, and terrain parameters, enabling a solution that is
21 more unified across sensors and introduces fewer sources of uncertainty. We leveraged
22 imaging spectroscopy data from AVIRIS-NG and PRISMA (collected within 1 hour) to
23 validate this method and showed a 25% increase in performance for the radiance-based
24 method versus using over the static method when estimating SSA. This concept can be
25 implemented in ~~future~~ missions such as Surface Biology and Geology (SBG), Environmental
26 Mapping and Analysis Program (EnMap), and Copernicus Hyperspectral Imaging Mission
27 for the Environment (CHIME).

28

29 **Key Words:** Imaging Spectroscopy, Snow Properties, Topography, Snow Albedo

30

31 **1 Introduction**

32 Accurately mapping snow surface properties is essential for seasonal snow zones in a
33 changing climate especially in regions where seasonal snowpack is expected to change
34 dramatically in the coming decades (Siirila-Woodburn et al., 2021). For example, snow
35 albedo plays a crucial role in melting of the snowpack during the ablation season (Wang et
36 al., 2020) with changes in snow albedo directly affecting the amount of absorbed solar
37 radiation, and therefore the amount of snow that is melted off ~~as liquid water~~. Throughout the
38 winter season, snow albedo fluctuates due in part to grain size (Seidel et al., 2016) and light

39 absorbing particles (Kaspari et al., 2015; McKenzie, 2020; Schmale et al., 2017; Skiles &
40 Painter, 2017). With a limited number of *in situ* snow stations around the globe ~~measuring~~
41 ~~surface properties~~, and the snow surfaces constantly undergoing metamorphism across space
42 and time, satellite imagery represents the best potential for spatially and temporally complete
43 mapping of snow properties. Accurately retrieving snow albedo and other snow surface
44 properties from satellite imagery is paramount, especially in a rapidly changing climate
45 (Malmros et al., 2018).

46 Retrieval of snow properties from satellite remote sensing relies on Digital Elevation
47 Models (DEMs) to correct for local terrain effects (Bair et al., 2021; Bair et al., 2022; Dozier
48 et al., 2022). In a previous study, researchers found global DEM products to have “blunders
49 and errors” when compared to airborne lidar, particularly in derived slope and aspect which
50 cause severe errors in calculated cosine of local solar illumination angles (μ_s) (Dozier et al.,
51 2022). ~~Dozier et al. (2022) They~~ found errors in ~~μ_s -local solar illumination angles~~ ranging
52 from 0.048 to 0.117 (dimensionless) across several sites for Copernicus global DEMs caused
53 by errors in slope and aspect. The ~~cosine of the local solar illumination angle, μ_s , term~~ is a
54 function (Eq. 1) of slope angle (S), ~~slope aspectslope azimuth angle or aspect~~ (A), solar
55 zenith angle (θ_0), and solar azimuth angle (ϕ_0) – with the last two being well constrained:

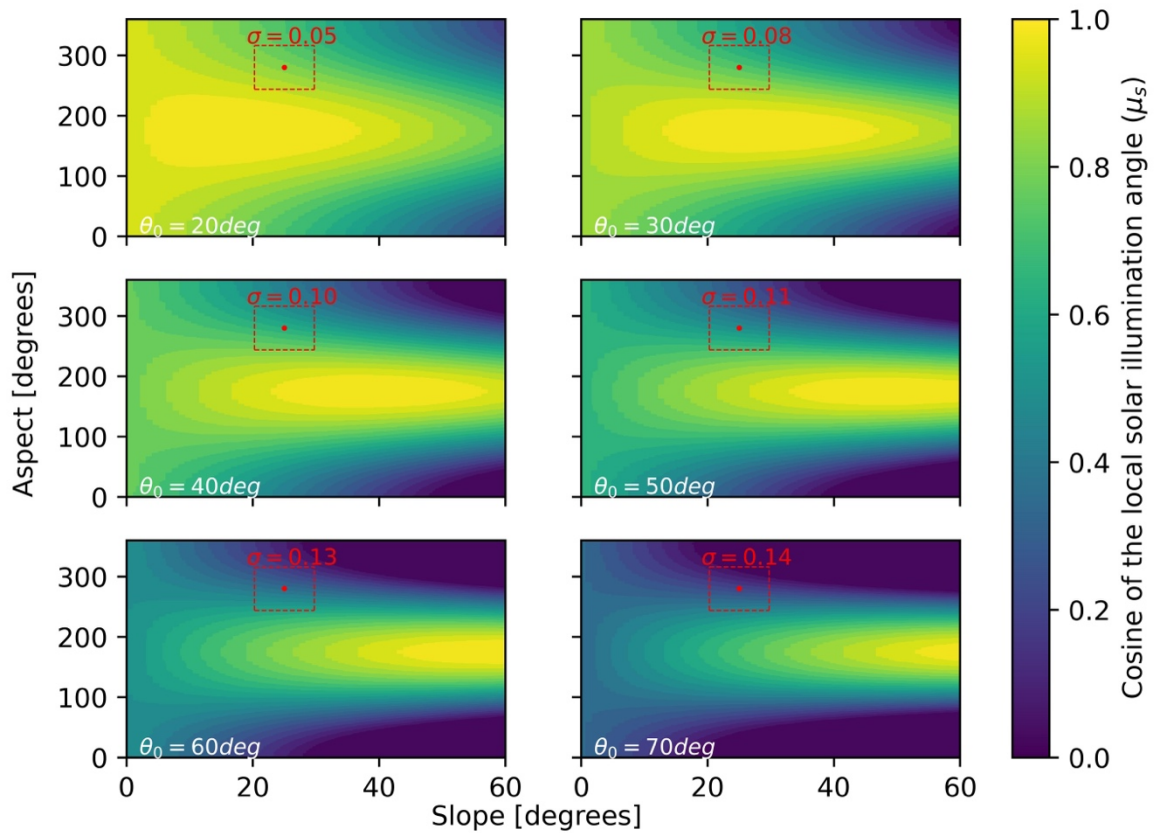
56

57
$$\mu_s = \max[0, \cos(\theta_0) \cos(S) + \sin(\theta_0) \sin(S) \cos(\phi_0 - A)] \quad (1)$$

58

59 Because θ_0 and ϕ_0 are calculable with low errors (~~less than 0.05°~~), the biggest contribution to
60 errors in μ_s stem from slope and aspect. Errors in μ_s increase monotonically with increasing
61 θ_0 (e.g., sun setting has high θ_0 , as does solar noon in high latitude winters). This
62 phenomenon can be explained by plotting Eq. 1 for various ~~solar zenith angles, θ_0~~ (Figure
63 1). Put simply, at higher θ_0 there is a higher standard deviation in μ_s surrounding a known
64 slope and aspect (with some temporally consistent uncertainty), increasing the probability and
65 magnitude of such an error. If one were to compute standard deviations of μ_s across varying
66 θ_0 , one would arrive at similar errors of μ_s presented in Dozier et al. (2022). For clarity, in
67 Figure 1 we have highlighted an example case with slope= $25^\circ \pm 4.73$ and aspect= $280^\circ \pm$
68 36.3. Example uncertainties for this exercise can be found in Table 2 of Dozier et al. (2022).
69 ~~For example, if one were to arbitrarily choose slope and aspect (with some uncertainty), and~~
70 ~~varying θ_0 ($20-70^\circ$), one could find a similar range of errors as presented in Dozier et al.~~
71 ~~(2022).~~

72



73

74 **Figure 1.** Cosine of local illumination angles (μ_s) as a function of slope (x-axis) and aspect

75 (y-axis) incremented by 1° , illustrating the problem at higher latitude, and/or winter

76 acquisitions, where standard deviation (σ) of μ_s increases monotonically with solar zenith

77 angles (θ_0). Aspect is shown here measured clockwise from north (with north containing a

78 discontinuity at 360 degrees). Increasing solar zenith angles (θ_0) illustrates the problem at

79 higher latitude, and/or winter acquisitions, where variability increases with respect to slope

80 and aspect. For this illustration ϕ_0 is fixed at a value of 175° . The red dots represent the

81 example point at slope=25° +/- 4.73 and aspect=280° +/- 36.3 and are bordered by their
82 uncertainty and the resulting σ .

83

84

85 Recent work has shown μ_s can be modelled using an optimal estimation framework
86 given the Top of Atmosphere (TOA) radiance observed from imaging spectroscopy (Carmon
87 et al., 2023). The authors solve for surface, atmospheric, and topographic state variables
88 simultaneously in their model. This works physically because the partition of direct to diffuse
89 light introduces a shape and magnitude effect on the TOA radiance spectra. However,
90 retrieving snow optical properties is sensitive to directional reflectance which is significantly
91 influenced by the viewing geometry and surface roughness (Bair et al., 2022), leading to
92 possible shortcomings in this method specifically for snow covered pixels. To address this
93 and expand upon this framework, we present a new method to account for terrain in snow
94 covered areas. Our method was tested on pixels with greater than 75% snow cover in three
95 different snow climates (tundra, maritime, and alpine) with spaceborne imaging spectroscopy
96 with the aim to reduce error in derived snow properties by optimally solving for topography.
97 The spaceborne results are validated against high confidence airborne spectrometer data ~~and~~
98 ~~further evaluated with error distributions~~. This work directly contributes to snow property
99 retrievals in steep terrain and/or at times of high solar ~~illumination~~ zenith angles for

100 ~~upcoming~~ satellite ~~image~~imaging spectroscopy missions such as Surface Biology and
101 Geology (SBG) (Cawse-Nicholson et al., 2021), ~~and~~ Copernicus Hyperspectral Imaging
102 Mission for the Environment (CHIME) (Celesti et al., 2022), ~~and~~ EnMap (Guanter et al.,
103 2015).

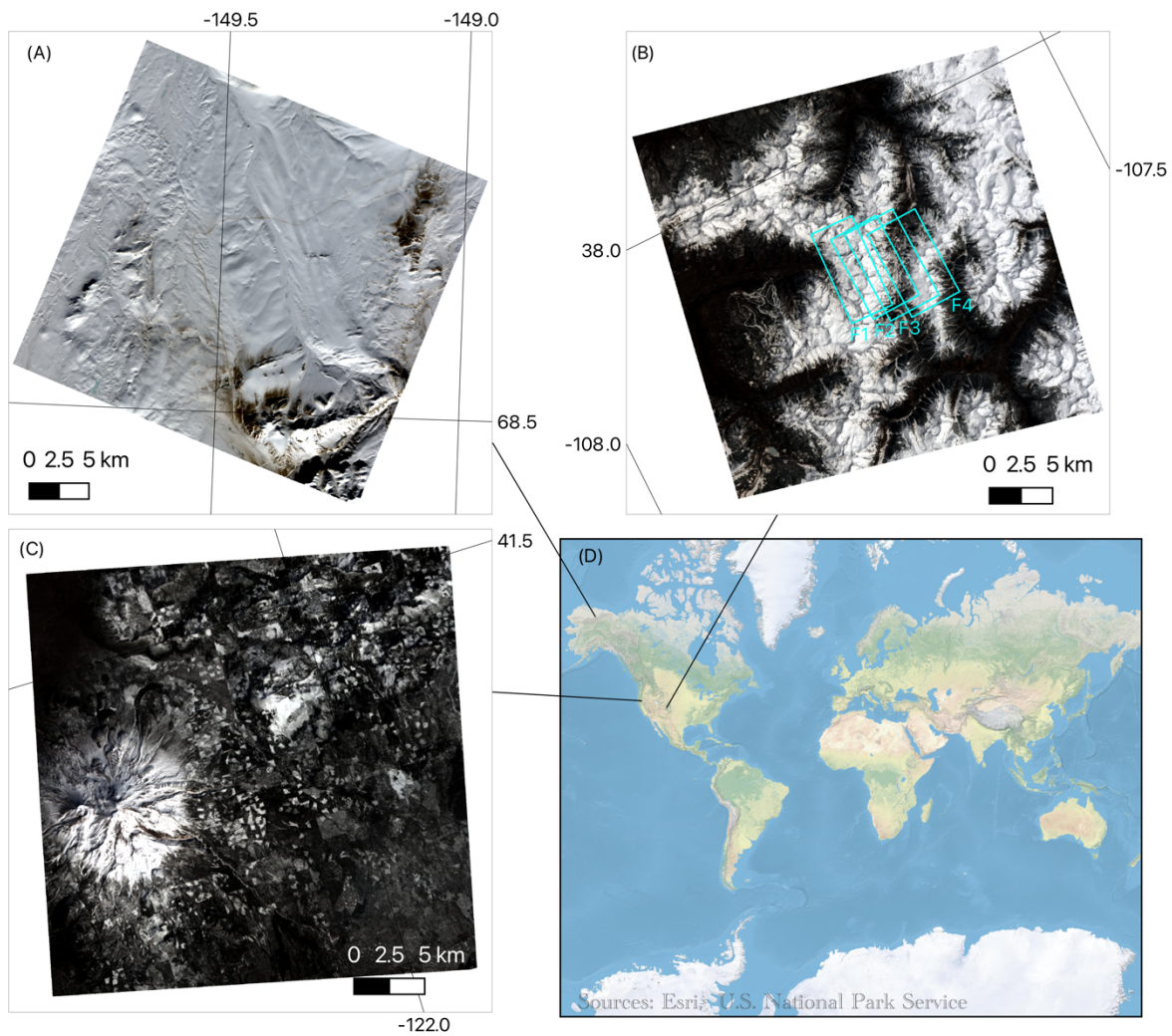
105 **2 Methods**

106 **2.1 Study area**

107 For our study, we used *PR*ecursore *I*perSpettrale della *M*issione *A*pplicativa
108 (PRISMA) imagery over three sites capturing different snow climates and solar zenith angles:
109 San Juan Mountains (Colorado, USA, 29 April 2021, $\theta_0=27^\circ$), Mount Shasta (California,
110 USA, 28 February 2021, $\theta_0=52^\circ$), and the Toolik area (Alaska, USA, 21 March 2021,
111 $\theta_0=68^\circ$) (Figure 2). The San Juan Mountains location is considered a high alpine site located
112 in interior continental USA with an elevation range of 2208-4129 m. The Mount Shasta site is
113 a maritime snow climate along the western coast of USA with an elevation range of 750-4232
114 m. The Toolik site (elevation range = 504-1748 m) is a high-latitude tundra site, being mostly
115 flat but with steep sections along the Brooks Range (along the southern part of the image).
116 PRISMA, launched by the Italian Space Agency (ASI) and beginning operation on March 22,
117 2019, is a spaceborne imaging spectroscopy mission collecting radiance at 30 m spatial
118 resolution across 23~~97~~ bands spanning 400-2500 nm at a spectral resolution better than 12 nm

119 ~~across at a spectral resolution of 9.24 nm and 9.27 nm in~~ the visible-near and shortwave
120 infrared, ~~respectively~~ (Cogliati et al., 2021).

121 To validate our method, we used four existing Airborne Visible Infrared Imaging
122 Spectrometer-Next Generation (AVIRIS-NG) flightlines over the San Juan Mountains from
123 29 April 2021 (flying 1 hour after PRISMA acquisition). AVIRIS-NG collects radiance
124 measurements at variable spatial resolution (depending on the flight altitude) across 425
125 bands spanning 380-2510 nm in 5nm intervals (Green et al., 2023). For this flight, data were
126 collected at 4 m spatial resolution. We downloaded AVIRIS-NG apparent reflectance from
127 National Snow and Ice Data Center (NSIDC) and observation geometry data from NASA
128 Search Earth Data (Skiles & Vuyovich, 2023).



129

130 **Figure 2.** PRISMA true colour images for Toolik on 21 March 2021 (A), San Juan
 131 Mountains on 29 April 2021 (B), and Mount Shasta on 28 February 2021 (C). Four
 132 coincident AVIRIS-NG flightlines (F1-F4) are shown in cyan over the San Juan Mountains.

133

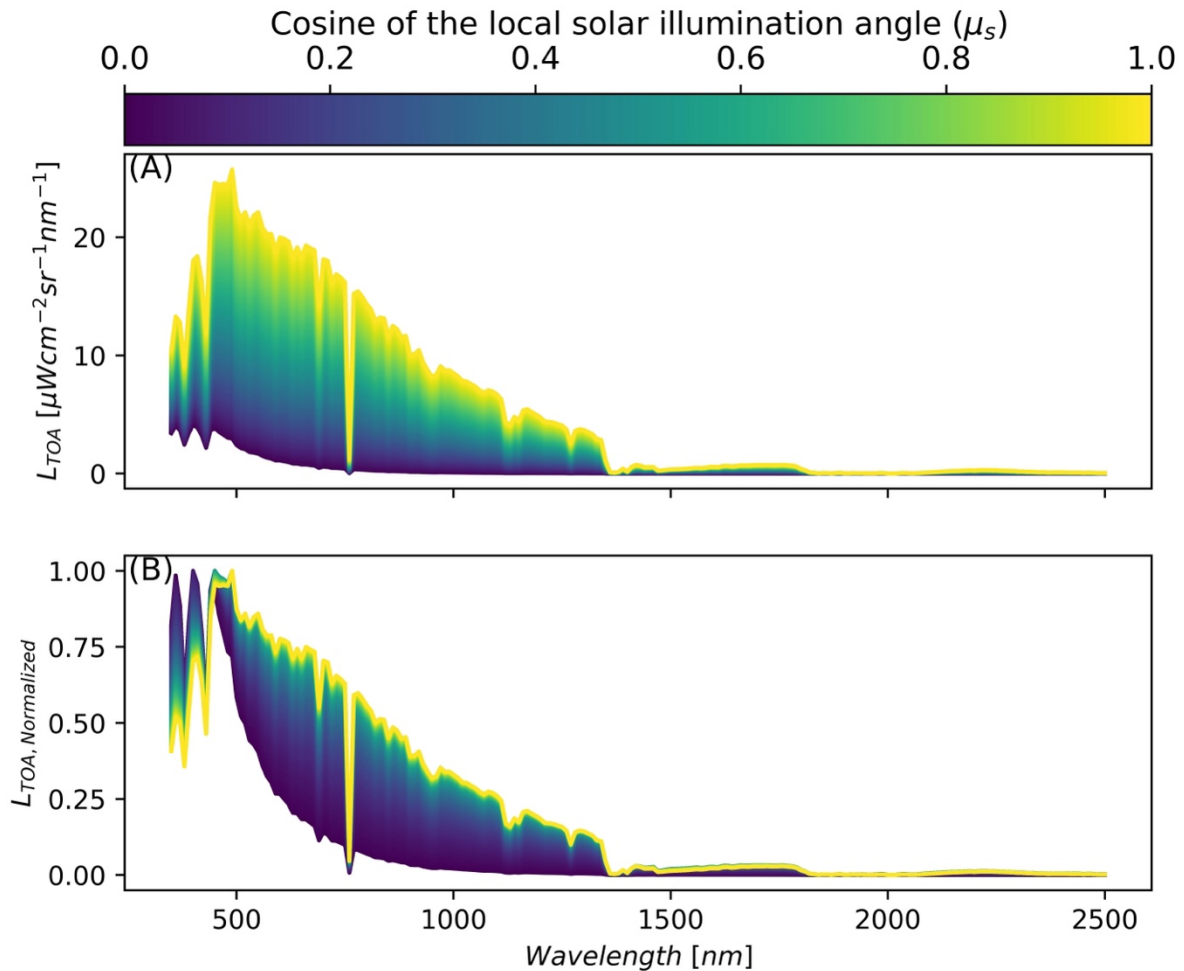
134 2.2 Modelling surface, atmosphere, and topography from PRISMA

135 The algorithmic improvements build off a workflow that estimates snow properties
136 given PRISMA TOA radiance, titled Global Optical Snow properties via High-speed
137 Algorithm using K-means (GOSHAWK) (~~Wilder et al., 2023;~~ Wilder et al., 2024). In short,
138 ~~GOSHAWK~~our method uses the analytic asymptotic radiative transfer model (AART)
139 (Kokhanovsky & Zege, 2004) coupled with libRadtran (Mayer & Kylling, 2005) to invert
140 snow surface and atmospheric properties (Bohn et al., 2021; Dalcin & Fang, 2021), and
141 fractional covers of mixed pixels under varied lighting conditions using non-linear numerical
142 optimization (Bair et al., 2021). The parameters solved for in the optimization routine include
143 fractional covers, specific surface area (SSA), light absorbing particle concentration
144 (modelled as ~~soot~~dust), liquid water content (~~LWC~~) percentage, dimensionless aerosol
145 optical depth at 550nm, and columnar water vapor in the atmosphere. Here, we expand upon
146 the ~~GOSHAWK~~ algorithm considering recent work showing the capacity to estimate μ_s from
147 TOA radiance (Carmon et al., 2023; Bohn et al. 2024). This ~~capacity~~idea is demonstrated in
148 Figure 3 using fixed snow properties via AART and fixed atmosphere properties via
149 libRadtran across the range of plausible μ_s (i.e. 0 to 1). Like the findings in Carmon et al.
150 (2023), Figure 3 shows that μ_s controls both the spectral shape and magnitude of observed
151 TOA radiance with the effect varying across wavelengths. The greatest shape effect can be
152 seen in the visible spectrum (roughly 400-700 nm) due to the magnitude of the diffuse

153 irradiance. In combination with the magnitude and shape shift, this parameter becomes
154 solvable during optimization due to its strong separability – especially when considering the
155 entire spectrum data from a hyperspectral remote sensing source such as [PRISMA](#). It is
156 important to note that μ_s impacts both the AART estimation of snow reflectance and
157 libRadtran estimation of incoming solar irradiance.

158

159



160

161 **Figure 3.** Synthetic data showing change in magnitude (A) and shape (B) of top of
 162 atmosphere radiance (L_{TOA}) with respect to changing local solar illumination angle (μ_s) for
 163 fixed snow surface state variables modelled with AART, and fixed atmospheric state
 164 variables modelled with libRadtran (viewing geometry was fixed as well). State variables and
 165 solar/view geometry were based on a PRISMA acquisition over southern Idaho on 8

166 December 2022. Figure (B) shows normalized radiance with respect to peak TOA radiance
167 across wavelengths to highlight the change in shape.

168

169 However, if we were only to optimize μ_s , the other key terms, local viewer zenith
170 angle (μ_v) and local phase angle (ξ) in the AART formulation for bidirectional reflectance of
171 snow (Eq. 2) (Kokhanovsky & Zege, 2004; Kokhanovsky et al., 2021a) would remain
172 constant from the available DEM (i.e., μ_s, μ_v, ξ are all derived from DEM) (Eq. 2; Wilder et
173 al., 2024),

174

$$175 \quad r_{snow}(\mu_s, \mu_v, \xi, \lambda) = r_0(\mu_s, \mu_v, \xi) \alpha_{snow}(\lambda) a_{snow}(\lambda)^f \quad (2)$$

176

177 where r_0 is the reflection function of a semi-infinite non-absorbing snow layer (Tedesco &
178 Kokhanovsky, 2007), α_{snow} is the ~~plane-spherical~~ albedo [plane albedo can be computed
179 using (26) in Kokhanovsky et al. (2021a)], f is the escape function (Kokhanovsky et al.,
180 2021a), and r_{snow} is the bidirectional reflectance of snow. Keeping other terms μ_v and ξ the
181 same are problematic because snow reflectance is poorly approximated as a non-Lambertian
182 surface (Leroux & Fily, 1988), and the outcome will be greatly influenced by μ_v and ξ .
183 Therefore, to incorporate solving for μ_s, μ_v , and ξ from TOA radiance into ~~the~~

184 ~~GOSHAWK~~the algorithm, we instead elect to optimally solve for $\cos(\text{aspect})$ (i.e.,
 185 “northness”) and $\sin(\text{aspect})$ (i.e., “eastness”) (Table 1).

186

187 **Table 1.** Parameter space and initial vectors used in numerical optimization for PRISMA

188 data.

<u>Parameter</u> <u>[unit]</u>	<u>Definition</u>	<u>Feasible</u> <u>Range</u>	<u>Initial</u> <u>State</u>	<u>Type</u>
<u>f_{snow} [%]</u>	<u>Fractional snow in the mixed pixel</u>	<u>[0, 100]</u>	<u>10</u>	<u>Surface</u>
<u>f_{shade} [%]</u>	<u>Fractional shade in the mixed pixel</u>	<u>[0, 100]</u>	<u>20</u>	<u>Surface</u>
<u>f_{LC1} [%]</u>	<u>Fractional cover of endmember 1 (based</u> <u>on land cover value at pixel)</u>	<u>[0, 100]</u>	<u>50</u>	<u>Surface</u>
<u>f_{LC2} [%]</u>	<u>Fractional cover of endmember 2 (based</u> <u>on land cover value at pixel)</u>	<u>[0, 100]</u>	<u>20</u>	<u>Surface</u>
<u>SSA [m² kg⁻¹]</u>	<u>Specific surface area (SSA)</u>	<u>[2, 156]</u>	<u>40</u>	<u>Surface</u>
<u>LAP [$\mu\text{g g}^{-1}$]</u>	<u>Concentration of light absorbing particles,</u> <u>LAP, modelled as dust (PM-2.5).</u>	<u>[0, 145]</u>	<u>0</u>	<u>Surface</u>
<u>Liquid water</u> <u>[%]</u>	<u>Percentage of liquid water on the snow</u> <u>surface</u>	<u>[0, 50]</u>	<u>2</u>	<u>Surface</u>
<u>AOD 550 [%]</u>	<u>Dimensionless Aerosol Optical Depth</u> <u>(AOD) at 550 nm</u>	<u>[1,100]</u>	<u>10</u>	<u>Atmospheric</u>
<u>H₂O [mm]</u>	<u>Columnar water vapor in the atmosphere</u>	<u>[1,50]</u>	<u>1</u>	<u>Atmospheric</u>
<u>Eastness</u>	<u>$\sin(\text{aspect})$</u>	<u>[-1,1]</u>	<u>Variable</u>	<u>Topographic</u>
<u>Northness</u>	<u>$\cos(\text{aspect})$</u>	<u>[-1,1]</u>	<u>Variable</u>	<u>Topographic</u>

189

190 Aspect can be solved during optimization by using the atan2 function. ~~(Van Rossum, 2020).~~
 191 We chose to use this method because eastness and northness are continuously differentiable,
 192 and therefore, are suited for numerical optimization methods, whereas aspect is discontinuous
 193 at north (using the convention of 0 and 360 degrees as north) is circular. We then can use this
 194 optimal aspect to estimate μ_s (Eq. 1), μ_v ~~(Eq. 3)~~, and ξ ~~(Eq. 4)~~.

195

$$196 \quad \mu_v = \max[0, \cos(\theta_v) \cos(S) + \sin(\theta_v) \sin(S) \cos(\phi_v - A)] \quad (3)$$

197

$$198 \quad \xi = \cos^{-1}(-\mu_s \mu_v + \sin(\theta_s) \sin(\theta_v) \cos(180 - (\phi_s - \phi_v))) \quad (4)$$

199

200 ~~where ϕ_v is the viewing azimuth angle, and θ_v is the viewer zenith angle on a flat plane.~~

201 This directly impacts Eq. 2 and Eq. 5 (and the formulation of incoming solar energy in the
 202 model) (Picard et al., 2020),

203

$$204 \quad E(\lambda) = \psi \mu_s E(\lambda)_{dir} + V_\Omega E(\lambda)_{diff} + \left[\left(1 + \frac{\cos(S)}{2} - V_\Omega \right) r(\lambda)_{surf} E(\lambda)_{diff} \left(\frac{E(\lambda)_{diff}}{E(\lambda)_{diff}} + \right. \right. \\
 205 \quad \left. \left. \frac{E(\lambda)_{diff}}{E(\lambda)_{diff}} \right) \right] \quad (5)$$

206

207 where E is total incoming irradiance, ψ is binary shade or no shade, E_{dir} and E_{diff} are the

208 direct and diffuse irradiance, respectively, V_Ω is the sky view factor (Dozier, 2022), and r_{surf}

209 is the reflectance of nearby terrain (which is assumed to be equal to the pixel itself).
210 ~~Explicitly within GOSHAWK the following equation is then~~The term E is solved using
211 within our non-linear numerical optimization method as described in (Wilder et al. (2024)).
212 This was modelled incorrectly in Wilder et al. (2024); however, this was corrected in this
213 paper where only diffuse irradiance is used in the 3rd term in Eq. 5. Also, ~~a~~Adding in the two
214 extra parameters (eastness and northness) in our updated optimization scheme did not change
215 our run time significantly, ~~which still hovered around 15 minutes (depending on the image),~~
216 ~~as noted in Wilder et al. (2024). It is~~Caution is advised against solving for slope and aspect in
217 the inversion due to the non-unique solution space (Figure 1); however, only considering
218 aspect ensures unique solutions of aspect, μ_s , μ_v , and ξ . We chose aspect because of its
219 greater impact on determining partition of direct and diffuse illumination and has been found
220 to be more impactful to errors associated with snow property retrieval (Donahue et al., 2023).
221 ~~also worth noting that while the optimal aspect cannot be solved accurately due to a non-~~
222 ~~unique solution space (Figure 1), the resulting μ_s , μ_v , and ξ are unique and the solution is~~
223 ~~such.~~In this study we used estimate of total ozone column as input into creating the
224 libRadtran look up table specific for each image. We used the average weekly ozone over the
225 bounds of the image from Sentinel-5P NRTI O3: Near Real-Time Ozone dataset. This
226 approach serves an improvement over Wilder et al. (2024), where ozone was fixed at 300
227 Dobson Units.

228

229 **2.2-3 Estimating snow properties from AVIRIS-NG for validation**

230 Due to the fine signal to noise ratio and ~~the higher~~ spatial resolution of AVIRIS-NG,
231 we treated the dataset as the ground reference. ~~It also captured a similar spectral range to~~
232 ~~PRISMA which made it a suitable comparison dataset. The main assumption here is that~~
233 ~~AVIRIS-NG pixels at 4 m are relatively homogenous and are either snow or no-snow –~~
234 ~~which may not always be the case. This could be a potential source of uncertainty in our~~
235 ~~analysis. It also captured a similar spectral range to PRISMA which made it a suitable~~
236 ~~comparison dataset.~~ To select snow-covered pixels, we solved for NDSI (Normalized
237 Difference Snow Index) using bands at 600 nm and 1500 nm. We limited our retrieval ~~of~~
238 ~~snow properties~~ for NDSI greater than or equal to 0.90 (Painter et al., 2013). A common
239 approach to retrieve snow grain size from pure snow pixels is to apply the scaled ~~band area~~
240 algorithm (Nolin & Dozier, 2000); however, it is recognized that ~~the large presence of LWC~~
241 ~~liquid water~~ is a limitation. The maximum air temperature of 10.8° C on the day of the image
242 at the San Juan Mountains site indicated that elevated ~~LWC liquid water~~ at the surface was
243 probable (Center for Snow and Avalanche Studies, 2023). Additionally, reflectance spectra
244 appeared to be shifted along the x-axis (wavelength) due to the presence of ~~LWC liquid water~~.
245 Therefore, we used ~~constrained non-linear~~ numerical optimization to model apparent snow
246 reflectance with AART by allowing fractional ~~snow, fractional~~ shade, ~~liquid water LWC~~, and

247 SSA to vary. We did not include rock or forest endmembers in this formulation, assuming the
248 4 m pixels are relatively homogenous as previously stated. Topographic incident angles were
249 held constant based on the 4 m resolution DEM provided by AVIRIS-NG. We minimized
250 Root Mean Square Error (RMSE) between observed-apparent and modelled-apparent snow
251 reflectance from AART wavelengths in the range, 1000-1250 nm. This range has high ice
252 absorption and limited impacts from atmospheric interference and LAP (Miller et al., 2016).
253 greater than 900 nm and not impacted by atmospheric interference and LAP (Miller et al.
254 2016) (901-1299 nm, 1451-1779 nm, and 1951-2449 nm). The presence of LWC liquid water
255 was included in our analysis by means of the composite refractive index of water and ice
256 (Donahue et al., 2022; Hale & Querry, 1973; Warren & Brandt, 2008). We assumed similar
257 grain shape assumptions for both PRISMA and AVIRIS-NG, and that if there is a bias due to
258 this it should be consistent between the two datasets in our analysis.

259

260 **2.4 Comparing modelled ~~snow albedo and SSA~~ snow properties**

261 The ~~GOSHAWK~~ algorithm was used in two different modes: 1) static topography
262 based on the Copernicus DEM (hereon called “*static*”); and 2) solved topography based on
263 the algorithm updates (hereon called “*radiance*”). To compare the accuracy of PRISMA
264 derived SSA and LWC liquid water, we resampled the AVIRIS-NG optical property results
265 (SSA and LWC) to match the PRISMA resolution (30 m) and extents by using bilinear

266 interpolation. Then, we sampled all valid pixels where PRISMA and AVIRIS-NG had snow.
267 We then computed r-pearson correlation coefficient, Mean Bias, and RMSE for the radiance
268 and static methods (with respect to AVIRIS-NG). Finally, we used Copernicus derived slope
269 and aspect maps to determine where the largest errors were occurring on the landscape to
270 compare with the theoretical basis presented in Figure 1. We do this by using the ~~Mean~~
271 ~~Absolute Error (MAE)~~mean absolute difference with respect to ~~μ_s , slope and aspect~~. We
272 expected to see higher differences in north facing aspects (i.e., μ_s approaches 0), and where
273 θ_0 was higher. To test the interaction with θ_0 more fully, we extended the analysis to Mount
274 Shasta, CA, and Toolik, Alaska, where no *in situ data* existed. ~~Finally, we~~We compared the
275 modelled ~~broadband albedo and SSA properties~~ between the radiance and static methods to
276 assess how these assumptions ~~propagated into outputs~~impacted results for these types of data
277 at 30 m scale.

279 2.5 Comparing DEM and radiance derived μ_s

280 _____ To ensure the resulting radiance derived μ_s were valid we downloaded the best
281 available validation data sources for comparison. For the San Juan and Shasta sites, we
282 collected DEM products at 1_m spatial ~~resolution and~~resolution and collected 5 m spatial
283 resolution DEM for the Toolik site (U.S. Geological Survey, 2019; U.S. Geological Survey,
284 2022). Then, we computed slope, aspect, solar zenith angle, and solar azimuth angle for all

285 pixels to compute μ_s at the native resolution (Eq. 1). Then, we used bilinear interpolation to
286 resample the 1 m and 5 m products to 30 m to exactly match the extents and resolution of our
287 PRISMA images. We would like to acknowledge that while these are the best freely available
288 datasets for our images, they still do not capture the true snow-on topography, and instead are
289 a representation of the “snow-free” surface. We compared matching pixels to determine
290 RMSE, R^2 , and Mean Bias. Pixels that were marked as shadow from ray tracing were
291 excluded from this comparison.

293 3 Results

294 3.1 Validation using AVIRIS-NG data over the San Juan Mountains

295 Over all flightlinesOver the area of the flightlines, AVIRIS-NG estimated mean SSA =
296 18.0 +/- 8.313.18 m² kg⁻¹, PRISMA radiance method estimated mean SSA = 23.6719.6 +/-
297 5.812.45 m² kg⁻¹, and PRISMA static method estimated mean SSA = 22.025.06 +/- 16.0612.1
298 m² kg⁻¹. When comparing the SSA performance over each pixel to the AVIRIS-NG flightlines
299 (Figure 4) we found the PRISMA radiance method (r=0.43; RMSD=8.0 m² kg⁻¹; bias=+1.7 m²
300 kg⁻¹; n=36,412) performed slightly better than the static method (r=0.23; RMSD=13.6 m² kg⁻¹
301 ; bias=+4.0 m² kg⁻¹; n=36,412) for SSA.

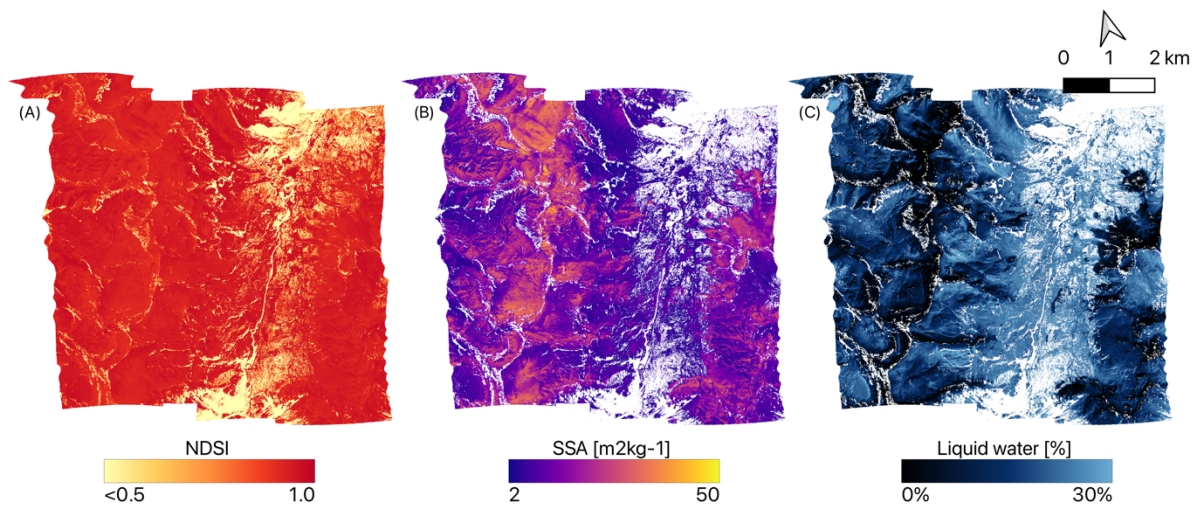


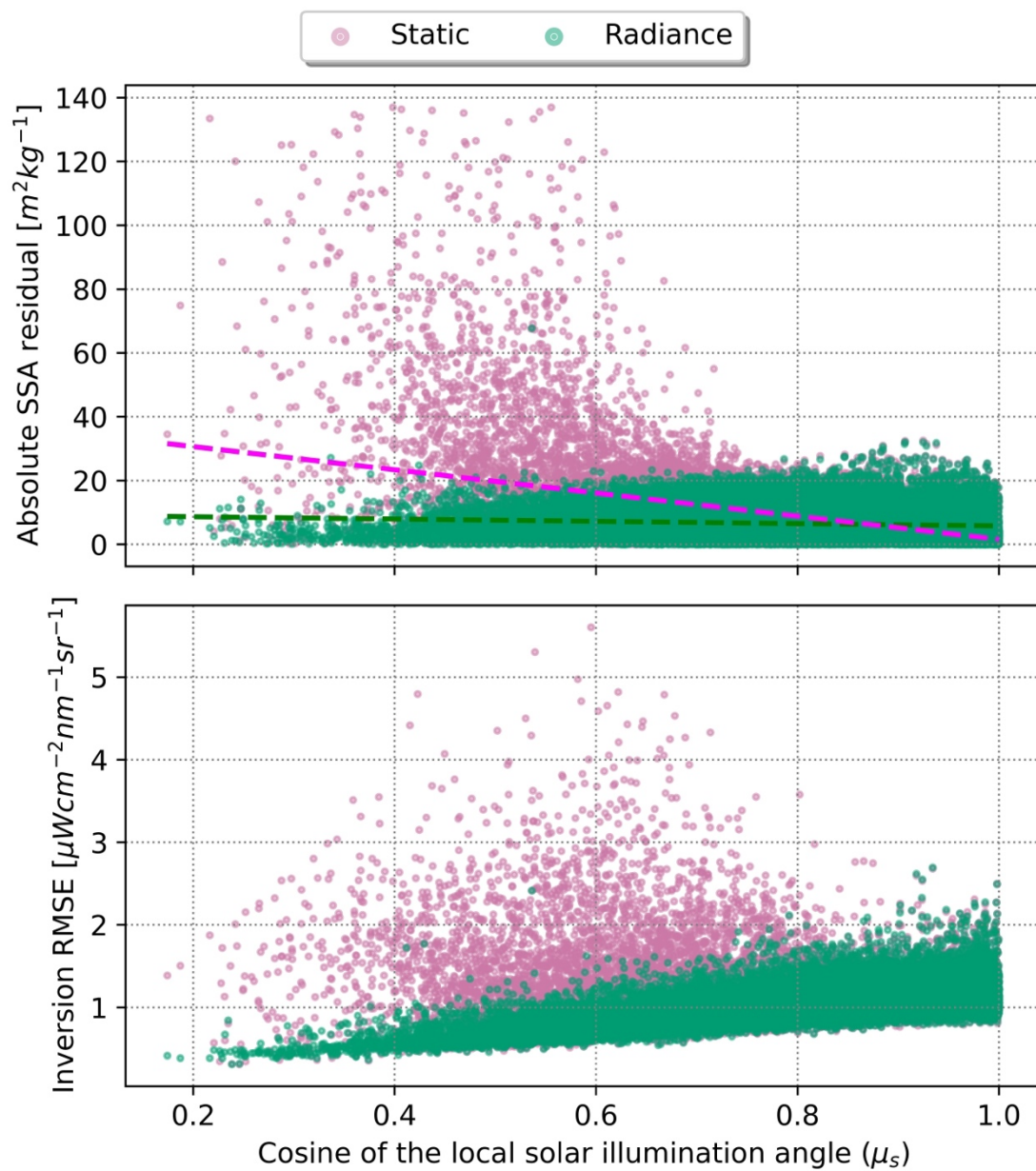
Figure 4. Snow properties computed from AVIRIS-NG (4 m spatial resolution) on 29 April 2021 including NDSI (A), SSA (B), and LWC (C) for the San Juan Mountains site.

There was not a significant improvement in liquid water estimation between radiance ($r=0.67$; $\text{RMSD}=10\%$; $\text{bias}=-8\%$; $n=36,412$) and static ($r=0.67$; $\text{RMSD}=10\%$; $\text{bias}=-9\%$; $n=36,412$). Furthermore, it appeared that there was a consistent liquid water bias of 8 to 9%, hinting that more melt had occurred during the AVIRIS-NG flights. As previously noted, the temperatures were well above freezing during the overpass of AVIRIS-NG and occurred roughly 1 hour later in the day compared to the PRISMA acquisition. This most likely explains the higher LWC liquid water and lower SSA observed by AVIRIS-NG. We further tested this by masking out areas where AVIRIS-NG liquid water content was greater than 0.1%, to establish areas where low amounts of melt occurred between the two acquisitions. We found that performance of

315 PRISMA static (RMSD=14.2 m² kg⁻¹; rRMSD=49%; n=181) and radiance (RMSD=6.9 m² kg⁻
316 1; rRMSD=23%; n=181) methods were more accurate for these areas. The radiance method
317 performed slightly better, suggesting a modest 25% improvement in accuracy for SSA over the
318 static method when considering pixels that were less impacted by melt.

319

320 Additionally, comparing all pixels we found improvement from radiance occurred
321 mostly on steep, north facing aspects (e.g., when μ_s approached 0). We found the absolute
322 residual increased as μ_s approached zero for the static method ($r = -0.47$; $p < 0.01$), while this
323 relationship was diminished nearly by a factor of 5 for the radiance method ($r = -0.10$; $p < 0.01$)
324 (Figure 5.A). These errors were caused by incorrect terrain information in the inversion, where
325 inversion error increased proportionately in the static method (Figure 5.B).



326

327 **Figure 5.** Absolute difference in modelled SSA when compared to AVIRIS-NG for radiance

328 method (green) and static method (pink) respect to μ_s (A) and resulting RMSE from the

329 inversion from PRISMA with respect to μ_s (B). Error in the static method increases
330 significantly when μ_s approached zero ($r = -0.47$; $p < 0.01$); however, the difference was less
331 noticeable in the radiance method ($r = -0.10$; $p < 0.01$).

332

333 **3.2 Comparing radiance and static methods between sites**

334 On average across each of the images, radiance and static methods provided similar
335 retrieved parameters within less than one standard deviation (Table 2). In general, this means
336 there is not a significant difference at the 30 m scale for computing parameters such as SSA
337 and broadband albedo (BA) when considering the entire image. Interestingly when terrain is
338 fixed, the static model compensated for incorrect illumination by increasing the aerosol optical
339 depth (thereby reducing the amount of direct solar radiation). Investigating the errors more
340 closely, we found much larger differences in retrieved properties where μ_s approached 0
341 (Figure 6). The difference in distributions matched closely to the theoretical demonstration
342 (Figure 1) and is most likely associated with the standard error of slope and aspect from
343 Copernicus DEM given the illumination conditions. This result also demonstrates the
344 difference between the two methods had the biggest impact for images where θ_0 was high,
345 resulting in potentially inaccurate retrievals that impact both surface and atmospheric state
346 variables on relatively mild slopes.

347 ~~Interestingly when terrain is fixed, the static model compensated for incorrect illumination by~~
348 ~~increasing the aerosol optical depth (thereby reducing the amount of direct solar radiation).~~
349 ~~The difference in distributions matched closely to the theoretical demonstration (Figure 1) and~~
350 ~~is most likely associated with the standard error of slope and aspect from Copernicus DEM~~
351 ~~given the illumination conditions. This result demonstrated that the difference between the two~~
352 ~~methods had the biggest impact for images where θ_v was high. On average for the Mount~~
353 ~~Shasta site, the radiance method estimated SSA was $17.57 \pm 14.58 \text{ m}^2 \text{ kg}^{-1}$, and static method~~
354 ~~SSA was $17.74 \pm 12.67 \text{ m}^2 \text{ kg}^{-1}$ (Table 3). Notably, there were more data gaps in the static~~
355 ~~method, resulting from failed inversions likely due to errors in the slope and aspect. On average~~
356 ~~for the Toolik site the radiance method estimated SSA to be $51.50 \pm 12.46 \text{ m}^2 \text{ kg}^{-1}$, and the~~
357 ~~static method estimated $52.07 \pm 14.49 \text{ m}^2 \text{ kg}^{-1}$.~~

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376 **Table 2.** Image-wide statistics comparing derived properties of SSA and broadband albedo
 377 between the two methods (static vs. radiance) processing the PRISMA imagery for all three
 378 sites.

Site	PRISMA Method	Mean SSA [(m ² kg ⁻¹)]	<u>Mean</u> <u>Broadband</u> <u>Albedo Standard</u> <u>Deviation of</u> <u>SSA (m² kg⁻¹)</u>	<u>Mean</u> <u>Broadband</u> <u>Albedo Mean</u> <u>Liquid</u> <u>water [%]</u>	<u>Standard</u> <u>Deviation of</u> <u>Broadband</u> <u>Albedo Mean</u> <u>AOD at 550</u> <u>nm</u>	<u>Mean</u> <u>water</u> <u>column</u> <u>vapour</u> <u>[mm]</u>
------	---------------	--	--	--	--	--

San Juan	Static	23.3 2.14 +/- 14.9	0.79 +/- 0.03 19.40	3.5 +/- 4.8 0.79	0.05 +/- 0.13 0.03	6.7 +/- 1.1
	Radiance	19.6 +/- 5.9 21.27	0.78 +/- 0.03 15.94	3.9 +/- 5.0 0.78	0.01 +/- 0.01 0.04	6.8 +/- 0.3
Shasta	Static	11.04 +/- 6.03 17.74	0.778 +/- 0.04 12.67	1.6 +/- 3.3 0.74	0.04 +/- 0.10 0.05	7.6 +/- 1.30
	Radiance	10 7.73 +/- 6.26 17.57	0.778 +/- 0.05 14.58	1.9 +/- 3.8 0.72	0.01 +/- 0.04 0.06	7.7 +/- 1.19
Toolik	Static	30.1 +/- 9.6 52.07	0.85 +/- 0.02 14.49	0.0 +/- 0.0 0.82	0.02 +/- 0.03 0.03	1.0 +/- 0.4
	Radiance	27.7 +/- 7.9 51.50	0.84 +/- 0.02 12.46	0.0 +/- 0.0 0.81	0.01 +/- 0.01 0.03	1.0 +/- 0.2

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~~When looking more closely at these errors we found~~

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~~Additionally, we saw the highest difference between the two methods on north facing aspects,~~

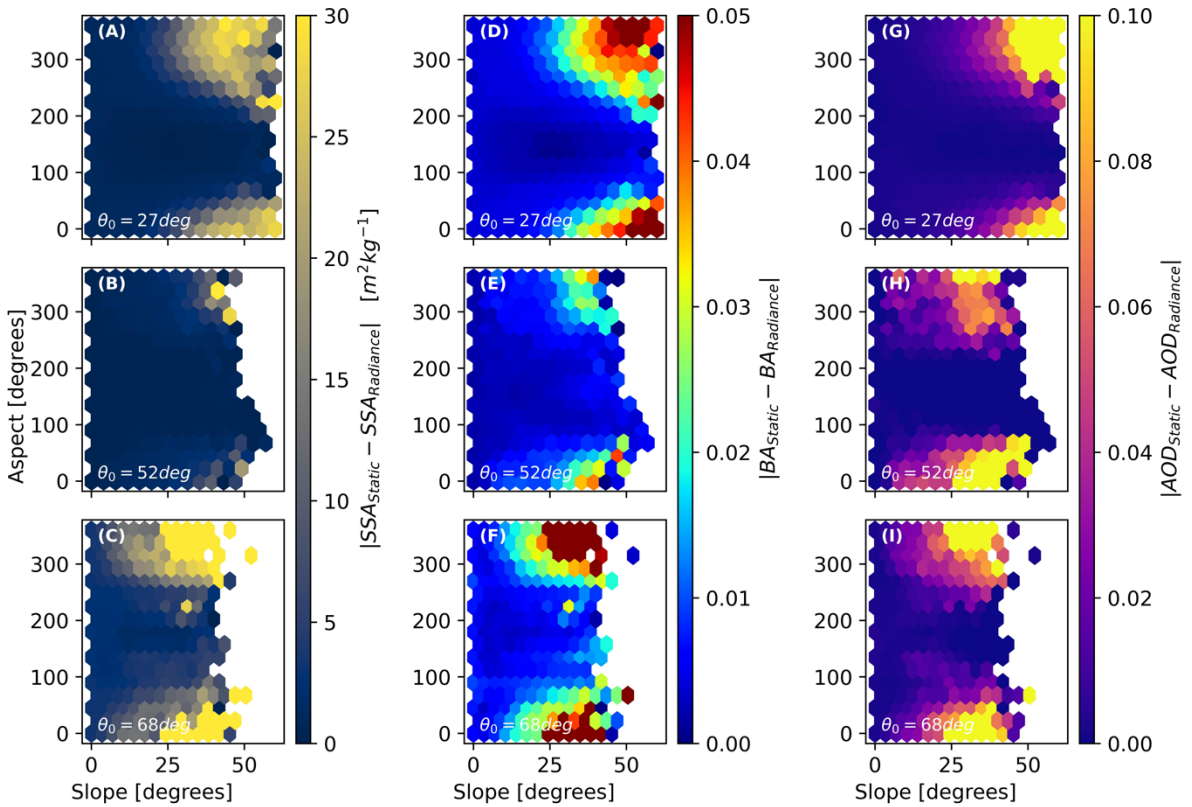
384

~~where μ_s approached 0 (Figure 7). The difference in distributions matched closely to the~~

385

~~theoretical demonstration (Figure 1) and is most likely associated with the standard error of~~

386 ~~slope and aspect from Copernicus DEM given the illumination conditions. This result~~
 387 ~~demonstrated that the difference between the two methods had the biggest impact for images~~
 388 ~~where θ_0 was high.~~



389

390 **Figure 6.** 2D Histogram plots showing absolute error-difference in SSA (left), broadband
 391 albedo (middle) and AOD (left) with respect to slope and aspect across the entire dataset. In
 392 this figure, treating radiance method as validation, absolute difference is calculated as $|\text{Static}$
 393 $- \text{Radiance}|$. This is shown for the San Juan Mountains site (A,D,G), Shasta site (B,E,H), and
 394 Toolik site (C,F,I). The average solar zenith angle (θ_0) is shown for reference on each panel.

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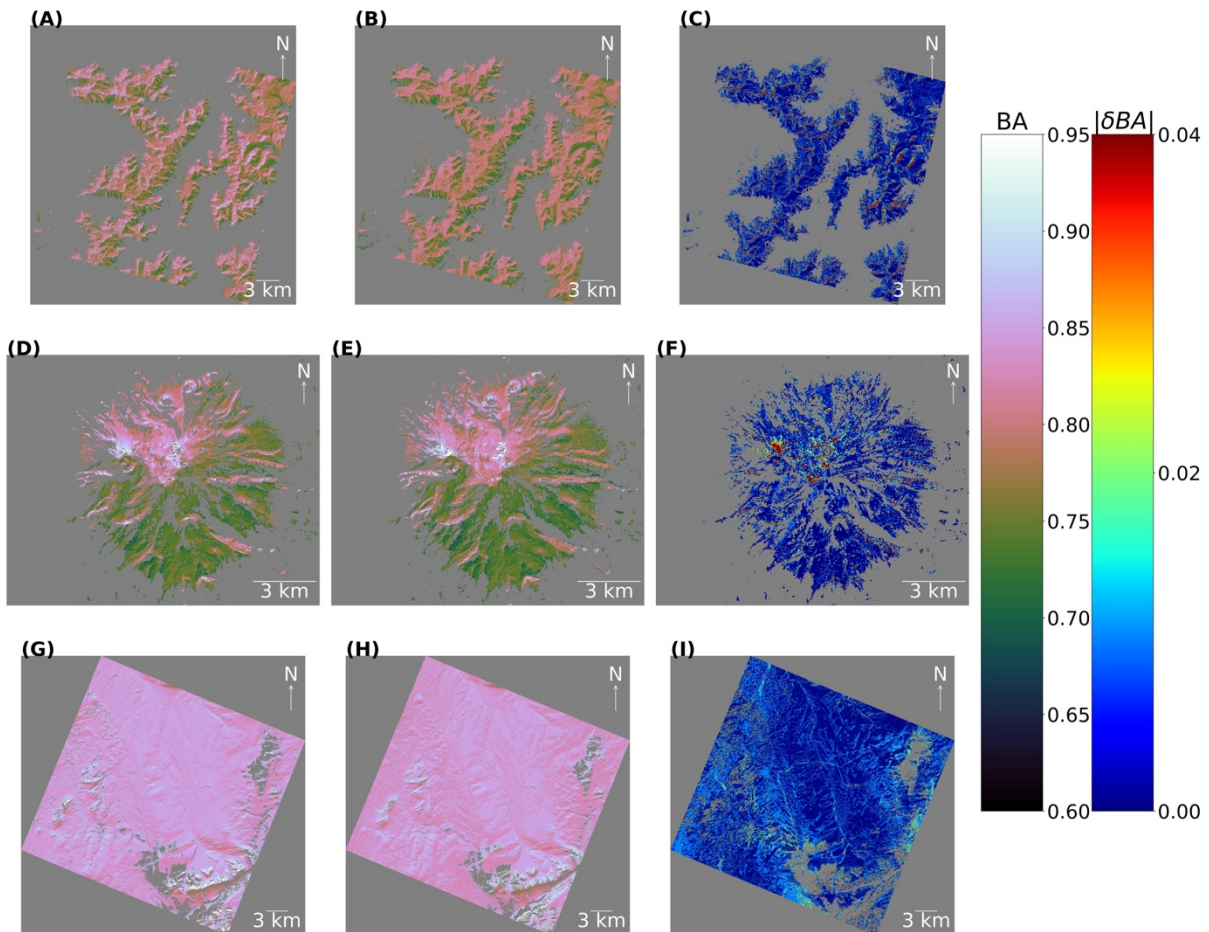
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Putting this into spatial context (Figure 7), San Juan site had 37% of pixels (135.3 km²) with an absolute difference in BA ($|\delta BA| \geq 0.01$) and 14% pixels (49.9 km²) with $|\delta BA| \geq 0.02$. Shasta site had ~~3028~~3028% of pixels (~~164.71~~164.71 km²) with $|\delta BA| \geq 0.01$ and 9% pixels (5.1 km²) with $|\delta BA| \geq 0.02$. Toolik site had 40% of pixels (325.3 km²) with $|\delta BA| \geq 0.01$ and 8% pixels (66.6 km²) with $|\delta BA| \geq 0.02$.



401

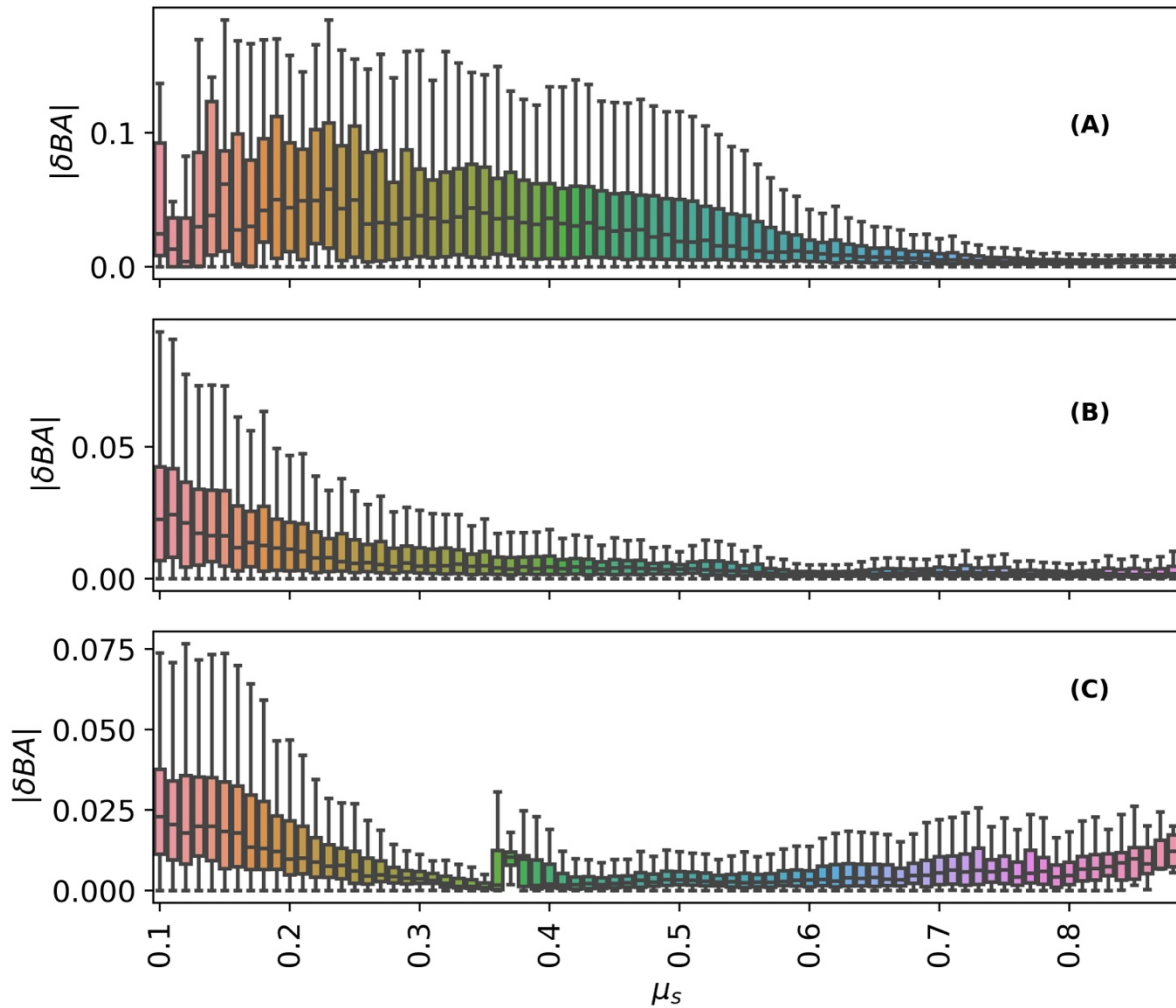
402 We observed a small but notable difference between the methods in derived snow
403 broadband albedo (BA) values (Figure 8). On average the standard deviation between BA_{Static}
404 $-BA_{Radiance}$ (δBA) for the San Juan Mountains site was 0.02, standard deviation for the Shasta
405 site was 0.01, and standard deviation for the Toolik site was 0.02. Generally, there was not a
406 clear bias with respect to μ_s . For the San Juan Mountains site, most of the pixels from the static
407 method showed a consistent small, negative bias of around -0.002. However, for shadier slopes
408 at this site, this bias flipped positive and was much more uncertain at around +0.02.
409 Interestingly, δBA from the Toolik site had the opposite relationship to San Juan Mountains
410 site, where δBA was more positive on sunnier slopes (μ_s approaching 1), and more negative
411 on shadier slopes (μ_s approaching 0), suggesting there could be different mechanisms for
412 which the static method may lead to inaccuracies. The result for the Toolik site confirms the
413 need to model the illumination conditions even in relatively flat terrain, because of implications
414 for net radiative forcing in the cryosphere. The Mount Shasta site δBA had no strong
415 relationship with respect to μ_s

416
417 **Figure 7.** Modelled broadband snow albedo (BA) for San Juan Mountains site (A-C), Shasta
418 Mountain site (D-F), and Toolik site (G-I). Left column represents BA from static method,
419 middle column represents BA from radiance method, and right column represents absolute
420 difference in BA ($|\delta BA|$). Dark grey colour symbolizes data that is not a value.

421

422 Median $|\delta\text{BA}|$ for all sites with respect to μ_s general increased as μ_s approached zero
423 (Figure 8). For example, ~~for the San Juan site,~~ median $|\delta\text{BA}|$ ranged from 0.03 to 0.00 across
424 μ_s . For the Shasta and Toolik sites, median $|\delta\text{BA}|$ ranged from 0.02 to 0.00 across μ_s . This
425 relation was ~~highly~~ non-linear and depended on the site and illumination conditions. ~~For~~
426 ~~example, standard deviation of $|\delta\text{BA}|$ (shown as the shaded regions in Figure 9) for well-lit~~
427 ~~slopes ($\mu_s > 0.8$) were generally smaller for San Juan site, and conversely were higher for~~
428 ~~the Toolik site. Similar to Dozier et al. (2022), one can see a monotonic relation with respect~~
429 ~~to θ_g across the three sites.~~ This analysis demonstrates the levels of uncertainty potentially
430 left in for retrievals relying on static, non-coincident DEMs. This shows quantitatively the
431 improvements to snow broadband albedo at 30 m scale by using radiance-based approach to
432 be relatively small for well-lit slopes – on the order 0-1%. While shaded slopes may have
433 errors in snow broadband albedo on the order of 1-3%. Interestingly for the Toolik site, $|\delta\text{BA}|$
434 also increased as μ_s approached one.

435



436

437 **Figure 8.** Modelled absolute difference in broadband albedo ($|\delta BA| = |BA_{Static} - BA_{Radiance}|$)

438 for San Juan (A), Shasta (B), and Toolik (C). Note these boxplots were created by rounding

439 μ_s to the nearest hundredth place.

440

3.3 Comparing DEM and radiance derived μ_s

At the 30 m pixel scale, Copernicus DEM derived μ_s had similar overall performance to radiance derived μ_s (Figure 9), with Copernicus DEM derived μ_s having slightly higher performance. For example, for the San Juan site, RMSD only varied by 0.006 between the two methods. Similarly, the R^2 for Copernicus derived μ_s was 0.86, while the radiance derived μ_s was slightly lower at 0.83. This similar overall performance was common amongst the three sites. We found the average bias for radiance derived μ_s was generally closer to zero (± 0.01), and did not show a strong negative or positive direction.

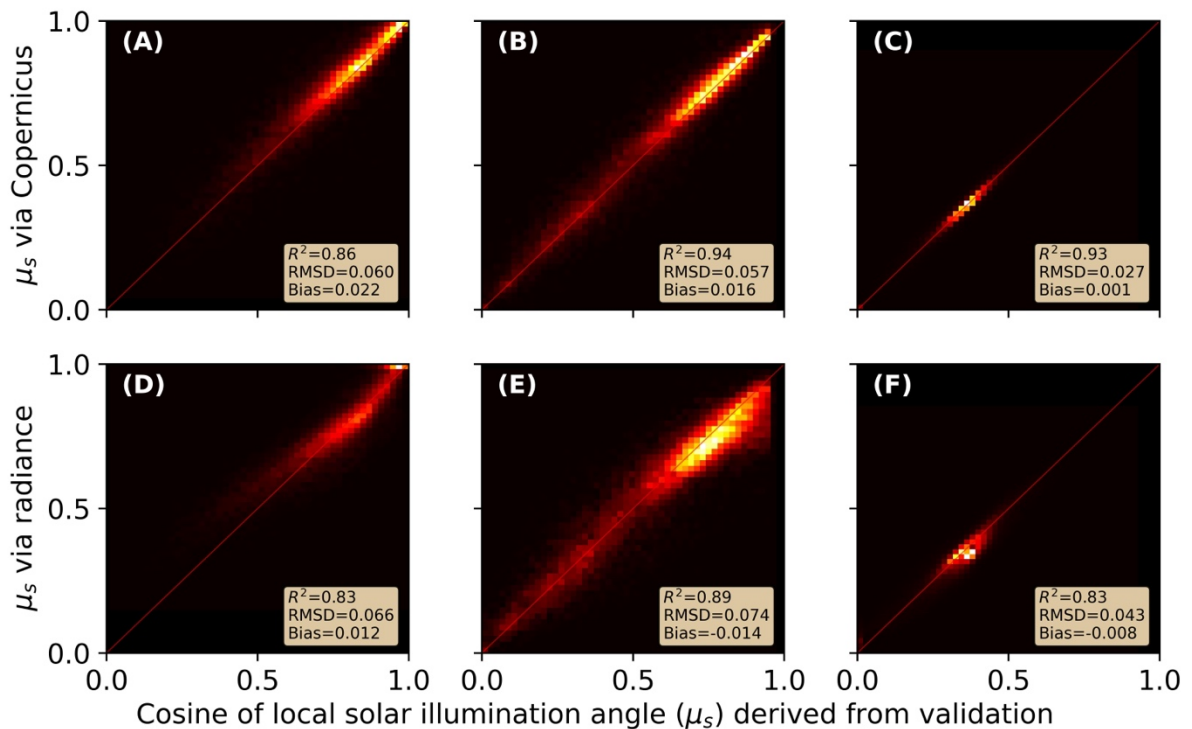


Figure 9. Comparing μ_s at 30 m pixel scale derived from radiance and Copernicus against high resolution DEM for San Juan site (A,D), Shasta site (B,E), and Toolik site (C,F).

4 Discussion

4.1 ~~Implications in accounting for terrain in snow property retrieval~~ Radiance derived DEMs may replace coincident DEMs and contain information related to surface roughness

Derivative slope and aspect maps are prone to errors at 30 m spatial resolutions (Dozier et al., 2022). ~~This which is become~~ relevant for derived snow products from upcoming missions such as SBG and CHIME which will rely on ~~such~~ topographic information to calculate optical properties like snow albedo ~~so that we can better monitor seasonal snowmelt~~. These errors can be inherent to the DEM itself, or a product of spatial and/or temporal misalignments (Carmon et al., 2023). ~~To enable high quality snow products regardless of illumination angles and conditions, we have demonstrated benefits of computing optimal terrain using TOA radiance over snow. This new method is especially useful for steep mountain terrain and/or high latitudes where illumination conditions are suboptimal. The θ_s (solar zenith angle) was relatively low for the San Juan Mountain site in~~

469 our study, and thus represents a lower bound of the improvement in accuracy one could
470 expect (Figure 1; Dozier et al., 2022). This disparity was demonstrated further for the Mount
471 Shasta and Toolik sites when θ_s was larger (i.e. a greater differences in SSA due to more
472 challenging solar and sensor geometry). Even for the relatively flat Toolik site, we showed
473 that correctly accounting for incidence angles can have an impact when θ_s is large. Our
474 modelled $|\delta\text{BA}|$ with respect to the non-coincident DEM was similar to work by Donahue et
475 al. (2023), who found slightly higher similar uncertainties of ~~δBA~~ broadband albedo (ranging
476 from -10 to 10%) for their investigation on Place Glacier, British Columbia, Canada. With
477 the surface and roughness undergoing dramatic change on glaciers throughout a given season,
478 using this radiance-based approach may be especially impactful for improving estimates over
479 glaciers. ~~We corroborated with this research showing similar ranges of δBA for our three~~
480 ~~study sites.~~

481 Snow surface roughness has long been a challenging issue in modelling snow
482 properties from space where the solar incidence angle at high spatial resolution for snow-on
483 DEM is not well known (Bair et al., 2022). Previous research found radiance derived μ_s from
484 airborne imaging spectroscopy showed a negative bias and postulated this could be due to
485 within-pixel topography, shadows, and surface roughness (Carmon et al., 2023). Since a bi-
486 directional reflectance function (BRDF) model was not used in their study, it then would be
487 plausible for the optimal μ_s to compensate for these effects. Interestingly when using a BRDF

488 model in our study (i.e., AART) and solving for aspect optimally (therefore informing μ_s , μ_v ,
489 and ξ) we did not find a strong bias – negative or positive. Although, we did not take surface
490 roughness measurements, and therefore do not know to the extent this impacted our study.
491 Within-pixel shadows, textures, and surface roughness remain difficult to validate, and we
492 were unable to achieve this in our study. Future work interested in further understanding this
493 radiance-based approach may investigate how such approaches interact with micro-scale
494 topography through the use of ground measurements such as ~~snow on~~ terrestrial and airborne
495 lidar.

496

497 **4.2.2 ~~Future considerations~~ Next steps in possibly improving this **radiance**** 498 **~~based~~ radiance-based approach**

499

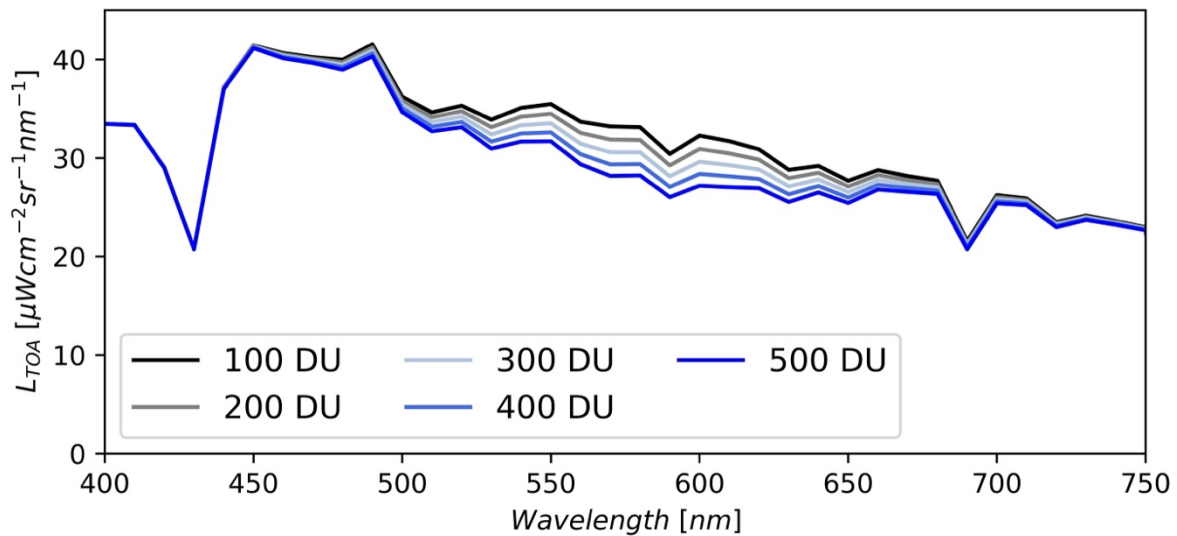
500 While we solved for a few terrain parameters in this study, we did not entirely
501 remove the use of the DEM ~~the static DEM, even~~ from the radiance method. The elevation
502 from global a-DEMs has a much higher confidence than its derivative products (Dozier et al.,
503 2022). Therefore, we used these values to inform our atmospheric routine, as well as our
504 shadow casting ray tracing module ~~in GOSHAWK~~ (Wilder et al., 2024). Additionally, ~~as~~
505 ~~stated in Wilder et al. (2024), GOSHAWK used the Dozier & Frew (1981)~~ we used the
506 method presented in Dozier (2022) for estimating the sky view factor (V_Ω) based on nearby

507 terrain and the pixel itself. This factor could potentially be problematic but was cited as being
508 not as ~~important-impactful~~ as μ_s in propagating error (Dozier et al., 2022). Therefore, we
509 elected to use V_Ω derived from the static Copernicus DEM. However, this could be an area
510 for future improvement, especially in very steep terrain where V_Ω becomes small. It is not
511 advised to attempt to add V_Ω directly into the optimization routine presented in this study, as
512 it is a function of pixel slope and aspect, and therefore, altering V_Ω and aspect together would
513 create invalid solutions.

514 Finally, we used a static value for slope derived from Copernicus DEM. The slope
515 influences the μ_s term, but also influences the passive radiation from nearby slopes.
516 Ultimately, we concluded that aspect had the largest impact on changing μ_s (Figure 1), as
517 well as large RMSE reported in previous work (Dozier et al., 2022; Donahue et al., 2023),
518 and thus was the focus of our study.- Caution is advised in including both slope and aspect
519 together, as non-unique solution space for μ_s may cause the optimization outputs to become
520 invalid. In summary, elevation, V_Ω , and slope remain static in our current implementation.
521 Future work may explore other algorithmic choices to further remove, or improve, static
522 DEM parameters.

523 Another consideration for improving this method is the inclusion of total column
524 ozone into the optimization. Previous research has been able to use TOA snow reflectance
525 data to retrieve reliable estimates of ozone (Kokhanovsky et al., 2021b). In our paper, we

526 lected for a simpler approach to first investigate the impacts of including terrain in the
527 optimization. In this paper we input a fixed ozone for the entire image based on coincident
528 Sentinel-5 measurements. However, it should be stated that ozone impacts a similar spectral
529 range to μ_s (Figure 10). It therefore may be beneficial to include ozone in the atmospheric
530 lookup-table (e.g., MODTRAN, libRadtran) to enable optimization of ozone as well. This
531 may be beneficial in building more realistic radiance-based methods.



532
533 **Figure 10.** Synthetic data showing change in magnitude of top of atmosphere radiance (L_{TOA})
534 with respect to changing total column ozone for fixed snow surface state variables modelled
535 with AART, and other fixed atmospheric state variables modelled with libRadtran. Reference
536 data is based on PRISMA image taken over southern Colorado. Note units of total column
537 ozone are shown in in Dobson Units (DU).

538

539 Finally, future studies should investigate including improvements to BRDF models of
540 snow (Mei et al., 2022). For example, recent work by Kokhanovsky et al. (2024) has
541 proposed the use of a two-layer model which may be especially useful for vertically
542 heterogenous snowpacks. Their method has been tested using EnMAP data and may easily
543 be transferable to other sensors. The current AART method we used in our paper does not
544 account for these layers, and instead assumes an optically thick, homogenous snowpack. To
545 validate both AART, ~~and~~ the new layered approach, and future BRDF models, snow pit (i.e.,
546 vertical profile) measurements of SSA (e.g., Meloche et al., 2023) become essential in
547 ensuring models accurately account for diverse layering of snow. ~~(Meloche et al., 2023).~~

548

549 **4.3 Big picture implications of the radiance-based approach**

550

551 This research responds to the objectives stated in “*Thriving on our changing planet: A*
552 *decadal strategy for Earth observation from space*”, to improve biogeophysical modelling at
553 scales driven by topography (National Academies of Science, Engineering, & Medicine,
554 20189), enabling more accurate snow property retrievals in the cryosphere under challenging
555 illumination conditions. Our work presented on solving terrain where DEM data are not
556 available, or reliable, may serve to accelerate improvements to satellite remote sensing tools

557 to monitor and model at both the regional global scale (Sturm et al., 2017), at a critical
558 juncture in time where northern latitudes are changing fast under a warming climate. This
559 includes Earth's glaciers, where radiance-based method may have the largest improvements
560 over static approaches. Our research is complimented by other recent works which show
561 promise in including terrain in the inversions (Bohn et al., 2024; Bohn et al., 2023; Bair et al.,
562 2024; Carmon et al., 2023)

563

564 We recommend additional coincident AVIRIS-NG flights with spaceborne imaging
565 spectroscopy datasets to further this work. As we have shown for the San Juan Mountains
566 site, for particularly warm days, images that are separated by longer than an hour may exhibit
567 drastically different SSA and liquid water content. As shown in this paper, this creates an
568 issue when trying to validate improvements to retrieval algorithms.

569

570

571 ~~However, future work could investigate other model and optimization configurations~~
572 ~~to improve upon this study. We recommend additional coincident AVIRIS-NG flights with~~
573 ~~spaceborne imaging spectroscopy datasets to further this work.~~

574

575 **5 Conclusions**

576 In this study we used existing PRISMA L1 TOA imagery to demonstrate the
577 improvements in modelling snow optical properties when explicitly modelling the terrain in
578 the inversion. This presents an interesting concept, that end users who are interested in
579 modelling snow from space, are perhaps better off working with the L1 TOA products, and
580 not using the L2 bottom of atmosphere reflectance products space agencies typically produce.
581 This would especially be true for areas where the surface undergoes rapid change, such as on
582 glaciers. This new method is especially useful for steep mountain terrain and/or high latitudes
583 where illumination conditions are suboptimal. The θ_0 (solar zenith angle) was relatively low
584 for the San Juan Mountains site in our study, and thus represents a lower bound of the
585 improvement in accuracy one could expect. This disparity was demonstrated further for the
586 Mount Shasta and Toolik sites when θ_0 was larger (i.e. a greater difference in retrieved
587 properties due to more challenging solar and sensor geometry). Even for the relatively flat
588 Toolik site, we showed that correctly accounting for incidence angles can impact snow
589 properties when θ_0 is large. Future work may look to build from this radiance-based approach
590 to enable better quantification of snow properties at scales impacted by topography.

591
592 *Code Availability.* <https://github.com/cryogars/goshawk>

593 *Author contributions.* B.W. created the GOSHAWK algorithm and updates herein, decided
594 on experiment set-up, and performed the subsequent analysis, as well as being the main
595 article writer. J.M., J.E. and N.G. provided ideas, comments, and supervised the work.

596 *Competing interests.* The contact author has declared that neither they nor their co-authors
597 have any competing interests.

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601 and for supplying the dataset.

602
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605

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