

Responses to Reviewer 1

We appreciate the time and effort Reviewer 1 has dedicated to providing a detailed review of this manuscript. Below we copy the referee comments in **bold black** and write our responses in **bold blue**. The changes to the manuscript are in *italic blue*.

General comments

Overall, I find this to be a well written manuscript that looks at an interesting and topical case of whether, and how, the Amery ice shelf cavity could switch from a cold to a warm regime over the next 100 years. If so this would have implications for it's future stability and consequent contribution to sea level rise. I recommend publication with minor corrections provided my two main comments are addressed (one potentially not so minor).

1) My first comment is that the text in figures throughout the manuscript is far too small. Whether this is merely a formatting error in the cryosphere draft format or not I can not say, but it needs to be addressed before publication.

We thank the reviewer for the overall positive assessment and the suggestions for improvements. We have upgraded the figures. In order to address Reviewer 2's comments on Section 3, we have replaced the old figures in the preprint. We summarised the modifications here to help the reviewer follow the changes we have made. We also include our responses to Reviewer 2, and a new version of the manuscript and the manuscript with tracked changes so that the reviewer understands why we made these corrections.

1 We have added more discussion on the temperature changes and water mass transformations in Section 3.1.2.

The additional analyses include:

- Replacing the depth-mean temperature (the old Figure 4) with temperature transect from the grounding line to the shelf break and beyond as suggested (Figure 4). The location of the transect is shown by the white line in Figure 1.**
- Temperature against salinity diagrams for each ocean cell within and outside the ice shelf cavity (Figure 5) to illustrate the changes in water masses.**
- Separated timeseries of depth-mean temperature/salinity within and outside the ice shelf cavity as suggested by Reviewer 1 (Figure 6).**

2 We have additional analyses of the changes in ocean circulations within and outside the ice shelf cavity in Section 3.2.1. The additional analyses include:

- Replacing the depth-averaged meridional velocity (the old Figure 5) with the barotropic stream function (Figure 7).
- Timeseries of the strength of circulations within and outside the ice shelf cavity (Figure 8).

3 We have replaced the analyses of PBG and PBECC heat transport (the old Figure 6) with the analyses of the heat budget in the ice shelf cavity (Figure 9).

4 We have additional discussion on the mechanism for the PBG reversal in Section 3.2.2. The modifications include:

- Replacing the analysis of the two boxes with a zonal transect of PBG (Figure 10).
- Analyses of the forcing of the gyre (Figure 11).
- Discussion on the impact of sea ice on the reversal of PBG (Figure 12).

2) The second, and potentially larger comment, is the lack of discussion about how choices of ocean model parameters affect the conclusions. The runs as presented are just for two different climate scenarios, with no investigation into the effect of model parameters. Was sensitivity analysis carried out for viscosity and diffusivity, time step, grid resolution, (all properties known to have a large impact on modelled melt rates) etc? I do appreciate that ocean models are expensive to run and so all of this may not be feasible, but would like to see some discussion about what was done. I suspect that the qualitative conclusions will hold up, but would like to see some analysis backing this up.

We appreciate the reviewer’s understanding of the computational cost of running the ocean model. It takes two to three months to finish all simulations. This is impractical to test the sensitivity of the conclusions to model parameters, especially the model resolution. The sensitivity to the model set-up is very complex. This is another area of modelling study, which is out of our research scope. However, if the reviewer is concerned about the robustness of the conclusions, we presented a second series of experiments which is forced by another UKESM1.0-LL ensemble (See supplement) to support the conclusions. In addition, we explained the reasons for choosing the parameters, for example, viscosity and diffusivity mentioned above, and other parameters of ice shelf thermodynamics which are more influential to the model results. The choices of parameters are questioned by Reviewer 2 as well.

We use a bilaplacian diffusion scheme in the AME025 configuration, the minimum stable requirement for the mixing coefficient (A^h) should meet

$$A^h < \frac{e^4}{64\Delta t}, \tag{1}$$

Where e is the smallest size of horizontal grids, Δt is the timestep (Gurvan et al., 2017). In the AME025 configuration, the smallest grid size is 7 km and the timesteps are 720 s, which gives the mixing coefficient a limit of $-5.02e+10 \text{ m}^4/\text{s}$. However, considering the very strong forcing of UKESM1.0 under SSP5-8.5, we would like to remove the potential risk of instability. The instability can be solved by reducing the timesteps or reducing the mixing coefficient. Reducing the

timesteps is computationally costly, so we decided to reduce the coefficient to $-1.08e+10 \text{ m}^4/\text{s}$ as in Bull et al. (2021). The values of viscosity and diffusivity have a minor impact on the modelled melt rate and ocean circulation. We tuned these parameters primarily for the model stabilisation.

A more important parameter is U_t , tidal velocity, in the ice shelf thermodynamics "three-equation" parameterisation (Holland and Jenkins, 1999). We use a velocity-dependent formula of the heat/salt exchange across the ice shelf-ocean interface (Jenkins et al., 2010):

$$\gamma_{T/S} = \sqrt{C_d} \Gamma_{T/S} U^*, \quad (2)$$

$$U^* = \sqrt{U_w^2 + U_t^2}, \quad (3)$$

where C_d is the non-dimensional top drag coefficient of 2.5×10^{-3} , $\Gamma_{T/S}$ is non-dimensional heat and salt exchange coefficient of 1.4×10^{-2} and 4×10^{-4} , respectively. U_* is friction velocity which consist of U_w averaged velocity in the top boundary layer (calculated by the model) and U_t tidal velocity. U_t is important for ice shelf melting as it affects the heat exchange between ocean and ice (Jourdain et al., 2019). It is common to employ external tidal forcing, for example Jourdain et al. (2017); Bull et al. (2021). However, it is impractical to obtain tidal forcing from UKESM1.0-LL, and we decided to prescribe tidal velocity as in Storkey et al. (2018) which provides the ocean core to UKESM1.0-LL.

The value of prescribed U_t in our simulations was determined based on the results from a sensitivity experiment, in which we modelled the melt rate in 2006 by increasing the tidal velocity from 0 cm/s to 12 cm/s in increments of 2.5 cm/s. We have provided the experimental results in the supplementary document. The melt rate is increased with the increase in tidal velocity (Figure 1a). The modelled melt rate with tidal velocity of 2.5, 5 and 7.5 cm/s falls within the observational range (Figure 1b). To a different extent, the three runs overestimate the melt rate north of the grounding line and underestimate the melt rate at the grounding line. This is the limitation of our configuration. The relatively coarse resolution cannot represent the steep meridional distributions of basal melting. In addition, Rosevear et al. (2022) estimated the magnitude of typical tidal current under Amery Ice Shelf from four tidal constituents. The estimated magnitude is 9.8 cm/s, which yielded an annual-mean Amery melt rate of 1.1 m/yr in 2006. This is slightly higher than the observational estimate. Considering the tidal velocity of 2.5 cm/s is much smaller than the observed tidal velocity, and 7.5 cm/s might largely overestimate the future melt rate given a very high climate sensitivity of UKESM1.0, we decided to use 5 cm/s. We have included the results of the experiment tuning the tidal velocity in the supplementary document.

We did a second series of experiments with forcings taken from the r2 ensemble of UKESM1.0-LL to validate our conclusions. We chose the r2 ensemble because the ocean temperature under SSP1-2.6 for the r2 ensemble has a very different variability from that for the r1. The increase in ocean temperature in the Amery sector in the r2 ensemble happens 20 years later than the r1 ensemble, and it is followed by a larger decrease after the temperature peak. The r2 simulations show similar results of the transition from a cold to a warm cavity, the reversal of Prydz Bay Gyre, etc. The timing of the abrupt increase in the melt rate in the r2 experiment coincides with the temperature variability in the r2

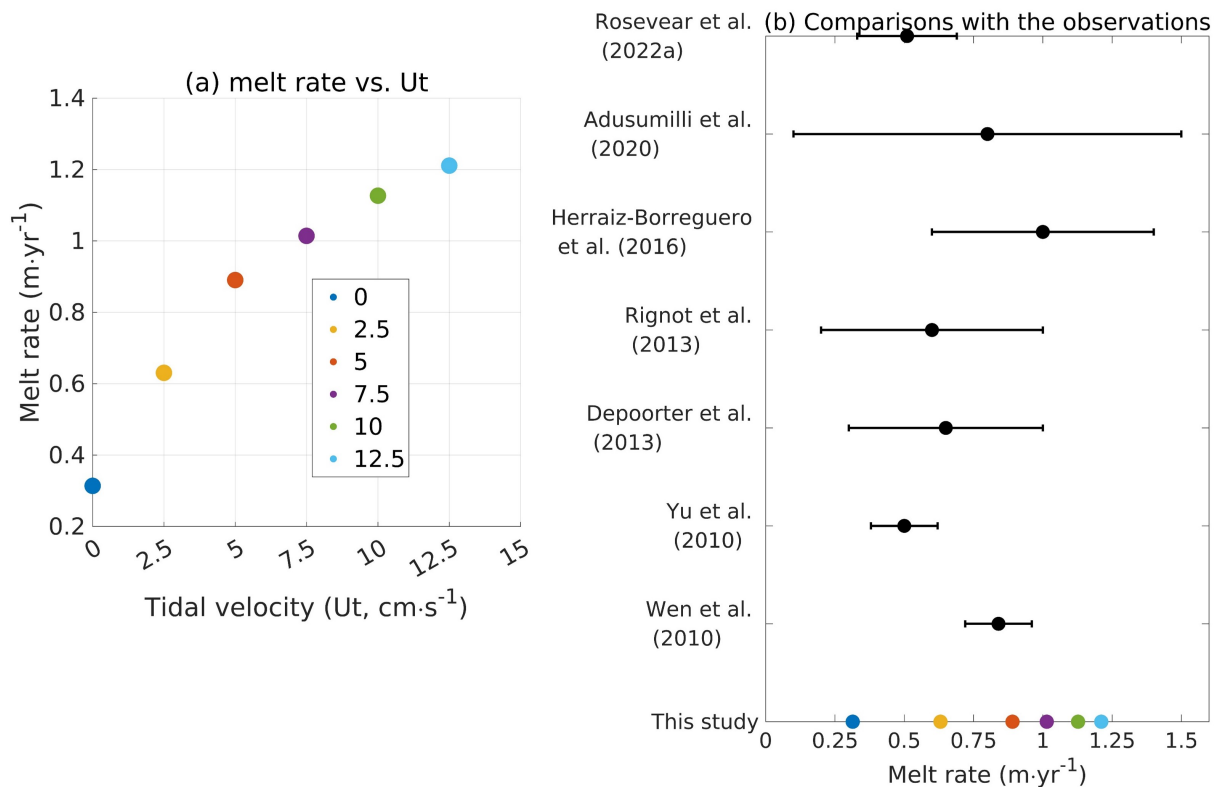


Figure 1. (a) Melt rate against tidal velocity (U_t). (b) Comparison of the melt rate with different U_t with the observational studies.

ensemble. We have supplied the results of the r2 simulations in the supplementary document.

Line comments below.

L17 - (and throughout) Bathymetry as a positive number reads oddly to me. This may be a glaciology bias, but I think there would be less ambiguity to any glaciologists reading if in cases like this it is referred to as negative from sea level, or otherwise made explicitly clear what it is referring to.

We thank the reviewer for providing a glaciological perspective. We have changed Bathymetry to the negative numbers in Figure 1. The vertical coordinates used in other figures are depth (pointing downwards) rather than height (pointing upwards). We keep the positive numbers and have claimed what the depth refers to in captions:

- **Figure 4:** "... The coordinates show the distance from GL against the depth below sea level."
- **Figure A1:** "Hovemoller diagram (time against depth below sea level) of zonal-mean meridional velocity ..."

L75 - How does the ice shelf draft from BEDMAP2 compare to a more recent data set such as BEDMACHINE (Morlighem, 2022)?

We compared the ice shelf draft from BEDMAP2 and BEDMACHINE (Figure 2). The ice shelf draft from BEDMAP2 and BEDMACHINE are not different on a large-scale (Figure 2a, b). There are some minor differences.

- Bedmap2 shows thicker ice near the ice shelf front (Figure 2c). Additionally, BEDMACHINE updated the ice shelf front. It outlines “Loose Tooth” while BEDMAP2 does not.
- Bedmap2 does not capture crevasses while BEDMACHINE does (Figure 2c).
- The ice at the grounding line is thinner in BEDMAP2 (Figure 2c). It is 300-1000 thinner than that in BEDMACHINE. BEDMAP2 also shows a less-defined channel near the southernmost grounding line.

BEDMACHINE has 500x500 m resolution, with two more additional sources to derive the ice thickness (Morlighem, 2022). By contrast, BEDMAP2 has a lower 1x1 km resolution. It captures less small geographical features, especially at the grounding line. However, the finest grid size in our model is about 7 km. It largely smooths these small-scale geographical features when interpolating the bathymetry/ice shelf draft data onto the ocean grids. Our future work is to update the bathymetry data to BEDMACHINE3 in a 1/12° configuration with a grid size of 1.5 km.

L90 - To clarify, by "top boundary" I am assuming you are averaging ice-ocean boundary conditions for melting over 30m. How is this weighted if this occurs over more than one cell?

Yes, it is the average over 30m below the ice shelf base or the first ocean cell below the ice shelf if the grid cell is thicker than 30m. It is weighted by grid cell thickness. For example, to obtain u velocity at the top boundary layer at a horizontal grid $u_{tbl}(i, j)$, let's assume three vertical levels are included in the top boundary layer, and the cell thickness for each grid on the U-point is $e3u(i, j, k)$, then the $u_{tbl}(i, j)$ is:

$$u_{tbl}(i, j) = \frac{\sum_{k=1}^3 u(i, j, k) \cdot e3u(i, j, k)}{\sum_{k=1}^3 e3u(i, j, k)} \quad (4)$$

L136 - What time step is the model running at? Please include somewhere.

Thanks for pointing out the missing information. The timestep is 720s. We have added to the manuscript.

L202 - What is the exact value of Cp used?

Cp is 3991.87 J kg/K. We have included this missing information.

L253 - ρ is used in equation 1 as a variable, not a constant. Suggest using ρ_{ref} or similar here to refer to reference density or the like and avoid confusion.

We thank the reviewer's suggestion. To address Review 2's question, we removed equation 1. However, we have changed ρ_{ref} for the reference sea water density in other equations.

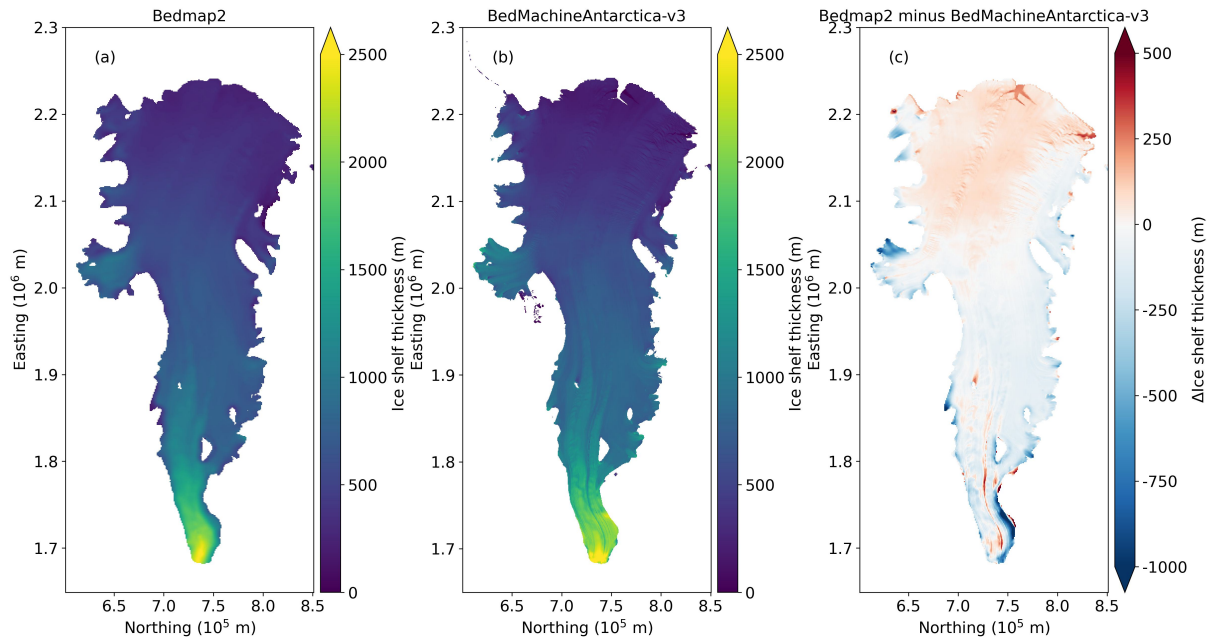


Figure 2. Thickness of Amery Ice Shelf from (a) Bedmap2, (b) BedMachineAntarctica-v3. (c) Comparison of ice shelf thickness between the two datasets. The BedMachineAntarctica-v3 data is interpolated onto the Bedmap2 grids. The positive/negative values represent that the ice shelf draft from Bedmap2 is thicker/thinner than that from BedMachineAntarctica-v3.

L302 - To clarify, by net melt rate here are you referring to area averaged melt rate minus surface precipitation?

We apologize for the ambiguous expression of “net melt rate”. It is area averaged basal melt rate only. We will replace “net melt rate” with “area averaged melt rate”.

Fig 1- (and throughout) Is it possible to change the grey calving front lines to a more contrasting colour, as at present it blends into the contour lines?

We thank the reviewer for the suggestions on improving the quality of figures. We have changed the calving front colour in Figure 1 to magenta. In Figure 3 and Figure 7, we increased the linewidth of the black coastal line to 4, and changed the calving front to a thinner black line (linewidth of 2).

Fig 2- The text is too small to make out at the current figure size.

We have increased the fontsize.

Fig 2- The text is too small to make out at the current figure size.

We have increased the fontsize.

Fig 3 - The text, particularly in the legend, is too small to easily make out.

We have enlarged the text.

Fig 4 - Text is too small. Could average cavity temperatures be included to show the lag between shelf water temperatures and cavity temperatures?

We removed this figure to address Reviewer 2's question. We have supplied figures with separated timeseries of the cavity temperature and the shelf sea temperature (Figure 6a).

Fig 5 - Text by velocity arrow key too small. & Fig 8 - Legend text too small.

We removed these figures to address Reviewer 2's question.

Figure A4 - I find this a potentially misleading figure. I am assuming this is just the running total of increased melting compared to initial geometry? Without any feedback on the glacial system included in the model I am not sure how useful this is. In reality, a thinner ice shelf would flow faster, thus a 10m increase in melt rate would not necessarily lead to a 10m reduction in thickness per year at a given point.

Yes, this figure is the running total of increased melting compared to initial geometry. We note this is an overestimate of ice shelf thickness reduction. The intention of the figure and the related context is to suggest the low potential of instability of Amery Ice Shelf and its upstream icesheets with our simulations. We have removed this figure and rephrased the context to avoid overinterpreting the results:

"... In our SSP5-8.5 simulation, the basal melt rate is not as high as that applied in the extreme experiments in Pittard et al. (2017), with $>50 \text{ m}\cdot\text{yr}^{-1}$ and $100 \text{ m}\cdot\text{yr}^{-1}$ for two extreme scenarios. It is reasonable to consider that AmIS will remain stable in the next century, however, the high basal melting puts the AmIS at risk of instability in the longer-term warming future. Coupling our standalone ocean configuration with an ice sheet model will be a future step to address the question of the AmIS stability."

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Current reversal leads to regime change in Amery Ice Shelf cavity in the twenty-first century

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Abstract. The Amery Ice Shelf (AmIS), the third largest ice shelf in Antarctica, has experienced relatively low rates of basal melt during the past decades. However, it is unclear how AmIS melting will respond to a future warming climate. Here, we use a regional ocean model forced by different climate scenarios to investigate AmIS melting by 2100. The areally-averaged melt rate is projected to increase from $0.7 \text{ m}\cdot\text{yr}^{-1}$ to $8 \text{ m}\cdot\text{yr}^{-1}$ in the low-emission scenario or $17 \text{ m}\cdot\text{yr}^{-1}$ in the high-emission scenario in 2100. An abrupt increase in melt rate happens in the 2060s in both scenarios. The redistribution of local salinity (hence density) in front of AmIS forms a new geostrophic balance, leading to the reversal of local currents. This transforms AmIS from a cold cavity to a warm cavity, and results in the jump in ice shelf melting. While the projections suggest that AmIS is unlikely to experience instability in the coming century, the high melting draws our attention to the role of oceanic processes in basal mass loss of Antarctic ice shelves in climate change.

10 1 Introduction

The Amery Ice Shelf (AmIS) is a large ice shelf in Antarctica (Figure 1). It is fed by Lambert Glacier system which accounts for about 16% of the grounded ice in East Antarctica (Allison, 1979). AmIS has a deep grounding line reaching $\sim 2500\text{m}$ below sea surface (Galton-Fenzi et al., 2008; Yang et al., 2021; Chen et al., 2023), which makes it susceptible to ocean temperature changes in the deep ice shelf cavity, due to a low in-situ freezing point temperature (Galton-Fenzi et al., 2012; Wang et al., 15 2022).

Prydz Bay (PB) has a v-shaped coastline and constrains the AmIS (O'Brien et al., 2014). It is divided by Prydz Channel (PC)—a trough with the depth of $\sim 500 \text{ m}$ at the shelf break and a deeper trough at the inner embayment named Amery Depression (AD) with bathymetry from $500\text{-}1000 \text{ m}$ (Figure 1). The Fram Bank (FB) and the Four Ladies Bank (FLB) are located on the western and eastern side of AD, respectively, and exhibit dramatic differences of zonal topography (Figure 1). 20 On the west of AD, the depth of FB rises sharply from $\sim 600 \text{ m}$ to $\sim 200 \text{ m}$ over a distance of $\sim 50 \text{ km}$. On the east of AD, FLB has a stepwise decrease from the depth of $\sim 200 \text{ m}$ to $\sim 500\text{m}$ over a distance of $\sim 200 \text{ km}$ (Liu et al., 2018).

Dense-High Salinity Shelf Water (DSW-HSSW) and modified Circumpolar Deep Water (mCDW) are major water masses causing AmIS basal melting (Galton-Fenzi et al., 2012; Herraiz-Borreguero et al., 2015, 2016). **DSW-HSSW** is dense and cold (slightly below the surface freezing point), which forms in coastal polynyas within PB during sea ice formation (Ohshima

25 et al., 2013; Williams et al., 2016; Portela et al., 2021). When DSW-HSSW descends into the AmIS cavity, the temperature of DSW-HSSW is higher than the in situ (pressure dependent) freezing point of seawater in the deepest cavity, which results in a basal melt rate of $>30 \text{ m}\cdot\text{yr}^{-1}$ at the grounding line (Galton-Fenzi et al., 2012). The resulting meltwater, named Ice Shelf Water (ISW), upwells along the western ice shelf base. ISW, with a temperature of the in situ freezing point, can be supercooled when it is ascending and refreezing beneath the northwestern AmIS (Craven et al., 2009).

30 The cavity under AmIS is presently filled with relatively cold DSW-HSSW of -2.2°C to -1.8°C (Craven et al., 2004; Herraiz-Borreguero et al., 2015). This results in relatively low basal melting along with basal refreezing beneath the northwestern AmIS (Depoorter et al., 2013; Rignot et al., 2013). Basal melt rates of AmIS have been stable during the past few decades (Adusumilli et al., 2020).

However, Southern Ocean temperatures are projected to increase in a warming climate (Fox-Kemper et al., 2021; Rintoul et al., 2018). The previous studies predict the increased upwelling of mCDW onto the continental shelf of the AmIS sector, resulting in a shift of cavity regime and consequently an increase of basal melting (Naughten et al., 2018; Kusahara et al., 2023; Thomas et al., 2023; Mathiot and Jourdain, 2023). These studies provide insight into the drivers of warming on the continental shelf. For example, a freshening in PB increases vertical stratification, and induces the warming (Aoki et al., 2022; Thomas et al., 2023; Kusahara et al., 2023). Another mechanism is related to the poleward shift of westerly winds, which enhances
40 the upwelling of mCDW across the shelf break (Spence et al., 2017; Guo et al., 2019; Verfaillie et al., 2022). However, future changes in the links between the warming in PB and the warming in the AmIS ice shelf cavity, in other words, local oceanic currents/intrusive pathways of warm water, still lack investigation.

mCDW was thought to be absent in the AmIS cavity (Craven et al., 2004). However, it has been observed at the AmIS calving front entering the cavity during winter recently (Herraiz-Borreguero et al., 2015). Other hydrographical observations
45 and modelling studies have documented the presence of mCDW on the continental shelf in PB (Galton-Fenzi et al., 2012; Herraiz-Borreguero et al., 2015, 2016; Williams et al., 2016; Liu et al., 2017; Guo et al., 2022). The two main pathways of mCDW intruding beneath AmIS are:

1. A large cyclonic gyre encircles AD, known as Prydz Bay Gyre (PBG) (Smith et al., 1984; Nunes Vaz and Lennon, 1996; Heywood et al., 1999). mCDW upwells across the continental shelf, arriving at PB over FLB, and is transported by the eastern
50 branch of PBG toward AmIS (Galton-Fenzi et al., 2012; Herraiz-Borreguero et al., 2015; Williams et al., 2016; Liu et al., 2017). Some ISW exits the cavity and recirculates within PBG and returns beneath AmIS (Williams et al., 2016).

2. A narrow current flows between FLB and the East coast, named Prydz Bay Eastern Coastal Current (PBECC), which originates in the Antarctic Slope Current (ASC) (Liu et al., 2017, 2018). Due to the step-like deepened bathymetry over FLB, ASC is redirected shoreward to conserve potential vorticity, resulting in the formation of PBECC (Liu et al., 2018), bringing
55 mCDW to PB (Liu et al., 2017).

Previous studies have documented that a re-directed coastal current in front of Filchner-Ronne Ice Shelf (Hellmer et al., 2017; Naughten et al., 2021) and Ross Ice Shelf (Siahaan et al., 2022) drives a rapid warming of the ice shelf cavity in climate scenarios. The aim of this study is to determine whether the strength and direction of intrusive mCDW pathways in PB remain the same or change in response to future climate change, and how these changes affect the melt experienced by AmIS.

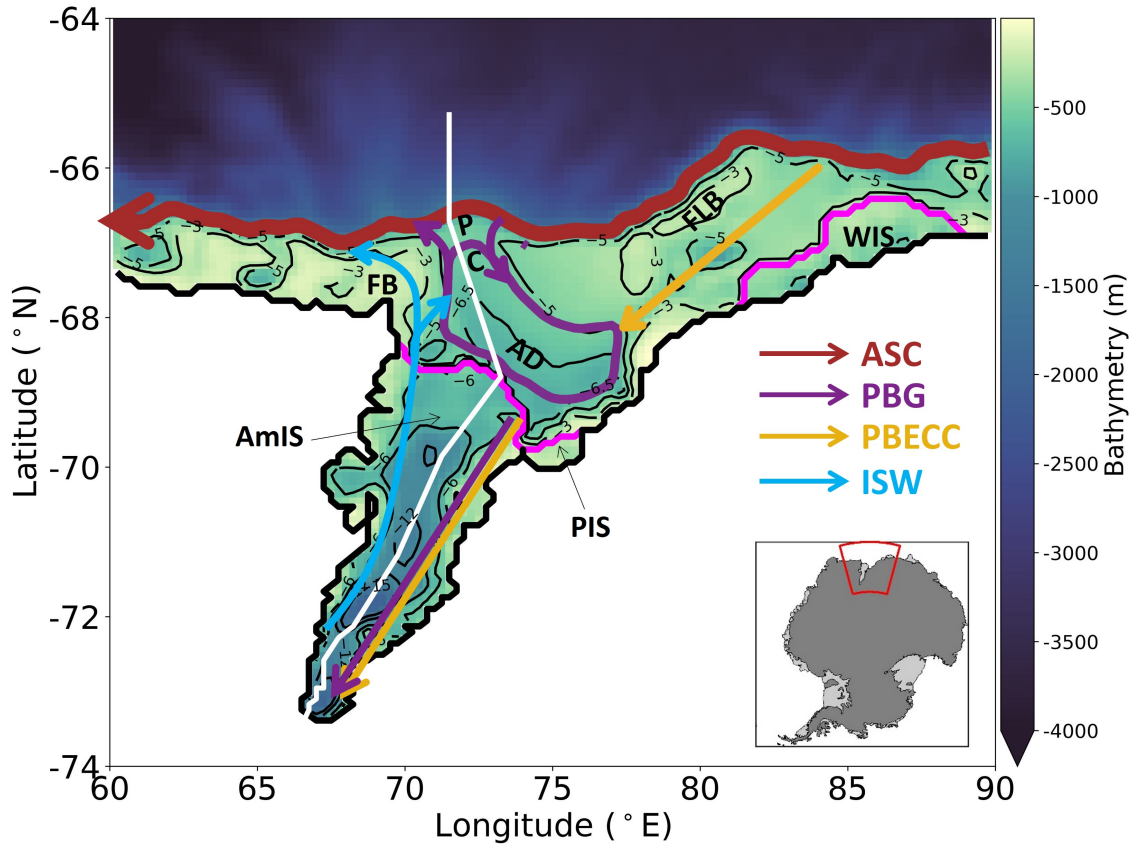


Figure 1. Model domain and schematic diagram of ocean circulation in Prydz Bay. The colour scale shows bathymetry (m). The thin black contours-lines on the continental shelf indicate bathymetry of 350-m-300, 500-m-500 and 700-650 m. The thin black lines on Amery Ice Shelf (AmIS) show bathymetry of -600, -800, -1200, -1500, -2000 m. The thick black line represents the coastal line, and the thick grey magenta lines show the ice shelf fronts. The thick red dashed arrow represents Antarctic Slope Current (ASC). Thick-The purple arrows show Prydz Bay Gyre (PBG). Thick-The yellow arrow indicates Prydz Bay Eastern Coastal Current (PBECC). Thin purple and yellow arrows in AmIS show modified Circumpolar Deep Water. Cyan arrows represent Ice Shelf Water (ISW). The white line shows a transect extending from the grounding line to the shelf break. Geographical features are labelled as follows: AmIS: Amery Ice Shelf, PIS: Publications Ice Shelf, WIS: West Ice Shelf, PC: Prydz Bay Channel, AD: Amery Depression, FLB: Four Ladies Bank, FB: Fram Bank. Inset: location of AmIS.

60 This paper is structured as follows: Section 2 describes our regional model configuration, model forcing and experiments; Section 3.1 presents projections of AmIS basal melting and temperature on the continental shelf; Section 3.2 proposes a mechanism driving the increase in AmIS basal melting; and Section 4 presents conclusions and discusses the limits and implications of this work.

2 Models and Methods

65 2.1 Regional model configuration

The regional **NEMO** configuration used to study Amery Ice Shelf-Prydz Bay (hereafter AME025) is summarised below.

The ocean model used in this study is version 3.6 of the Nucleus for European Modelling of Ocean model (NEMO3.6, (Gurvan et al., 2017)). The configuration includes the physical ocean engine OPA (Gurvan et al., 2017), a sea-ice model LIM3 (Rousset et al., 2015) and an open ice-shelf cavity (Mathiot et al., 2017). The model domain extends 30° in longitude ($60\text{-}90^\circ\text{E}$) and 10° in latitude ($65\text{-}75^\circ\text{S}$) with nominal 0.25° horizontal resolution with grid spacing increasing from ~ 7 km to ~ 12 km with increasing distance from the South Pole. A 75-level z^* vertical coordinate with a partial cell scheme is used in this work. It is a non-linear free surface application allowing for variations of volume according to the vertical resolution, and can adjust at top and bottom cells to represent bathymetry more realistically. The thicknesses of grid cells range from 1 m at the surface to 200 m at about 6000 m. Following Global Ocean version 7 (GO7) provided by Storkey et al. (2018), and bathymetry is derived from ETOPO1 data set (Amante and Eakins, 2009) with GEBCO giving modifications in coastal regions (IOC et al., 2003). Bathymetry under the ice shelf is derived from IBSCO (Arndt et al., 2013) and the ice shelf draft is taken from BEDMAP2 (Fretwell et al., 2013).

The schemes and parameter values used in parameterisations are primarily based on GO7 (Storkey et al., 2018), but the values of lateral diffusivity and viscosity have been changed ~~for our specific study, and some~~ according to the smallest grid spacing (~ 7 km) and time step (720s). Some physical schemes, such as the slip condition for the lateral boundary, from other regional modelling studies (Mathiot et al., 2017; Jourdain et al., 2017; Bull et al., 2021) are taken into account as well. A 55-term polynomial approximation of the Thermodynamic Equation of Seawater (TEOS-10, IOC et al. 2010) is used in our configuration (Roquet et al., 2015). A vector-form formulation of the momentum advection is applied. The vorticity term is computed using conserving potential entropy and horizontal kinetic energy (Arakawa and Lamb, 1981). Lateral diffusion of tracers is evaluated using Laplacian isoneutral mixing with a coefficient of $135 \text{ m}^2 \cdot \text{s}^{-1}$. Lateral diffusion of momentum uses bi-Laplacian geopotential viscosity with a coefficient of $-1.08 \times 10^{-10} \text{ m}^4 \cdot \text{s}^{-1}$. The vertical eddy viscosity and diffusivity coefficients are calculated from a Turbulent Kinetic Energy (TKE) scheme, with a background vertical eddy viscosity and vertical eddy diffusivity of $1.2 \times 10^{-4} \text{ m}^2 \cdot \text{s}^{-1}$ and of $2 \times 10^{-6} \text{ m}^2 \cdot \text{s}^{-1}$, respectively. The enhanced vertical diffusion parameterization is implemented for tracer convective processes with a coefficient of $10 \text{ m}^2 \cdot \text{s}^{-1}$. The non-linear bottom friction parameterization is chosen, with a non-dimensional bottom drag coefficient of 2.5×10^{-3} . The no-slip condition is implemented at the lateral momentum boundary.

Ice shelf thermodynamics are parameterised by the three-equation formulation with velocity dependent heat and salt exchange coefficients (Jenkins et al., 2010). The top boundary is set at 30 m below the ice shelf draft (or the first wet cell if the grid thickness is thicker than 30 m), and the top drag coefficient is set as 2.5×10^{-3} . The heat and salt exchange coefficients are 1.4×10^{-2} and 4×10^{-4} , respectively. This implementation does not include external tidal forcing. The velocity of the tidal current at the top boundary is prescribed as $5 \text{ cm} \cdot \text{s}^{-1}$. A series of sensitivity experiments have been conducted to determine the value of the prescribed tidal velocity (Supplement).

The atmospheric surface boundary conditions and the oceanic lateral boundary conditions are derived from the UKESM1.0-LL outputs (Sellar et al., 2020). There is no freshwater or salt restoring in the regional domain. However, the GO7 configuration—the ocean core of UKESM1.0-LL—has restoring in the global domain (Storkey et al., 2018). No external freshwater fluxes of surface runoff and iceberg calving are used in the AME025 configuration. These freshwater fluxes implicitly enter our domain through the lateral ocean boundary conditions, as they are prescribed in the UKESM1.0-LL outputs (Sellar et al., 2020).

2.2 UKESM1.0-LL forcing

We use climate projections made by the UK Earth System Model (UKESM1.0-LL) as part of the CMIP6 exercise (Sellar et al., 2020). UKESM1.0-LL is based on the HadGEM3-GC3.1 physical climate model, with additional atmospheric chemistry, and marine and terrestrial biogeochemistry components (Sellar et al., 2020). The ocean model for UKESM1.0-LL is NEMO3.6, and it is based on Global Ocean version 6 (GO6) configuration (Storkey et al., 2018), with 1° horizontal resolution and 75 vertical levels.

Model outputs from the first ensemble member (r1i1p1f2) of UKESM1.0-LL have been assessed by previous studies (Beadling et al., 2020; Heuzé, 2021; Purich and England, 2021; Roach et al., 2020; Bracegirdle et al., 2020; Meehl et al., 2020; Forster et al., 2020). To briefly summarise the evaluations of the modelled Southern Ocean in historical simulations, UKESM1.0-LL has relatively small biases in upper (0-100 m) ocean temperature (Beadling et al., 2020, Figure 6, Figure S3) and salinity (Beadling et al., 2020, Figure 6, Figure S4), and bottom temperature (Heuzé, 2021, Figure A3) and salinity (Heuzé, 2021, Figure A2) compared to other CMIP6 models. For the interior ocean, UKESM1.0-LL captures temperature and density structure across the continental shelf (Purich and England, 2021, Figure S2). It exhibits small cold ($<-1^{\circ}\text{C}$) and fresh (<-0.15 psu) biases close to the shelf break along the 90°E longitudinal transect (Beadling et al., 2020, Figure S5) which is the easterly ocean boundary of the AME025 configuration. In addition, UKESM1.0-LL shows accuracy in representing the position and strength of westerly jet over the Southern Ocean (Bracegirdle et al., 2020, Table 4). However, it is notable that UKESM1.0-LL has a higher climate sensitivity compared with other CMIP6 models (Forster et al., 2020, Figure 1; Meehl et al., 2020, Table 2), which might result in an unrealistically high surface warming (Forster et al., 2020). Additionally, UKESM1.0-LL overestimates Antarctic sea ice loss in **summer-February** (Beadling et al., 2020; Roach et al., 2020), which might cause a larger fresh bias at the ocean surface. However, it underestimates the February sea ice concentration in Prydz Bay (Roach et al., 2020, Figure S3), which may result in sea ice biases in our simulations. Moreover, UKESM1.0-LL has a relatively coarse 1° grid which cannot present some oceanographical features close to the Antarctic margin. Mathiot et al. (2011) suggest a minimum 0.5° nominal resolution to capture the ASC and local gyres on the continental shelf.

~~No external freshwater fluxes of surface runoff and ice sheet melting are used in the AME025 configuration. These freshwater fluxes implicitly enter our domain through the lateral ocean boundary conditions. The version of UKESM1 used in this study does not have an interactive ice sheet component. The ice sheet fluxes in UKESM1 are modelled in three parts (Sellar et al., 2020): (1) Surface runoff is simulated by the land surface model component of UKESM1; (2) As UKESM1 uses closed ice shelf cavities, the meltwater flux from ice shelves is prescribed and spatially distributed between the depths of the calving front and the averaged grounding line (Mathiot et al., 2017); and (3) UKESM1 incorporated the iceberg flux through a Lagrangian~~

iceberg model implemented in NEMO (Marsh et al., 2015). The freshwater flux of meltwater (45%) and iceberg calving (55%) is calibrated by the modelled pre-industrial surface mass balance of the Southern Hemisphere (Sellar et al., 2020).

2.3 Experimental design

135 Three experiments will be presented in this study: a historical simulation during 1976-2014 (Historical); a projection under SSP5-8.5 high emission scenario during 2015-2100; and a projection under SSP1-2.6 which is an optimistic emission scenario during 2015-2100.

Initial conditions consist of ocean temperature and salinity in January 1976 derived from GO7 (Storkey et al., 2018), which include an open cavity beneath AmIS. Atmospheric and oceanic forcing of our simulations are provided by the r1i1p1f2 ensemble member of UKESM1.0-LL over the period of 1976-2100. The surface atmospheric forcing comprises daily variables, including near-surface air temperature at 2 m, specific humidity at 2 m, horizontal wind components at 10 m; surface downwelling longwave and shortwave radiation; and precipitation (rainfall and snowfall). It is applied through CORE bulk formulae (Large and Yeager, 2004). Note that surface runoff is not applied in the simulations. The lateral ocean boundary conditions consist of monthly ocean and sea ice variables, including sea surface temperature, sea surface height; ocean temperature, salinity, 145 barotropic and baroclinic velocity; sea ice fraction, sea ice thickness and snow thickness.

The model outputs are saved as monthly mean. Therefore, the processes shorter than one month are not considered in this study.

The initialisation process is carried out by repeatedly simulating the first year (which is 1976) until the model drift becomes small. No smooth transition from December of a spin-up year to January of the next spin-up year is applied. There is physical discontinuity of the forcing between two spin-up years. Then a subsequent run is restarted from the last spin-up run with the continuous forcing from 1976 to 2100. The outputs of the subsequent run are used in the analysis. 150

The time series of area-averaged melt rates of AmIS in the spin-up run shows that it achieves equilibrium after about 10 years (Figure 2a). The melt rate dramatically declines from $\sim 16 \text{ m}\cdot\text{yr}^{-1}$ in the first spin-up year to $\sim 1 \text{ m}\cdot\text{yr}^{-1}$ in year 9 and becomes stable afterwards. The averaged temperature and salinity in the entire domain exhibit similar behaviours (Figure 2b).

155 Comparisons between the modelled melt rate in the Historical experiment (1976-2014) and the observational melt rate found in other studies (Wen et al., 2010; Yu et al., 2010; Depoorter et al., 2013; Rignot et al., 2013; Herraiz-Borreguero et al., 2016; Adusumilli et al., 2020; Rosevear et al., 2022a) are shown in Figure 2c. The time-mean melt rate during 1976-2014 in our study is $0.75 \text{ m}\cdot\text{yr}^{-1}$. This agrees with other estimates ranging from 0.5 ± 0.12 to 1 ± 0.4 (Figure 2c).

Time series of (a) area-averaged basal melting of AmIS ($\text{m}\cdot\text{yr}^{-1}$), (b) temperature ($^{\circ}\text{C}$, blue line and y-axis on the left) and salinity ($\text{g}\cdot\text{kg}^{-1}$, red line and y-axis on the right) in 30-year spin-up run. (c) Comparisons of AmIS melt rate among different studies. The grey dots represent the estimate of the AmIS melt rate and uncertainties from observational studies. The red dot shows the average of modelled AmIS melt rate from Historical simulation (1976-2014) in our study. 160

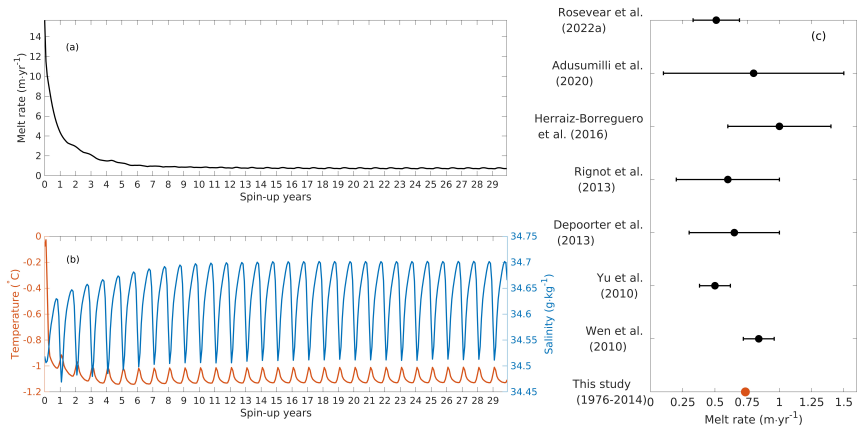


Figure 2. Time series of (a) area-averaged basal melting of AmIS ($\text{m} \cdot \text{yr}^{-1}$), (b) temperature ($^{\circ}\text{C}$, red line and y-axis on the left) and salinity ($\text{g} \cdot \text{kg}^{-1}$, blue line and y-axis on the right) in 30-year spin-up run. (c) Comparisons of AmIS melt rate among different studies. The black dots represent the estimate of the AmIS melt rate and uncertainties from observational studies. The red dot shows the average of modelled AmIS melt rate from Historical simulation (1976-2014) in our study.

3 Results

3.1 The projected changes in the PB-AIS system by 2100

165 3.1.1 The increased basal melting

Figure 3 shows the time evolution of the area-averaged AmIS basal melting from 1976 to 2100. The melt rates fluctuate around $0.7 \text{ m} \cdot \text{yr}^{-1}$ until 2056, then abruptly jump to $\sim 13 \text{ m} \cdot \text{yr}^{-1}$ under SSP5-8.5 and $\sim 9 \text{ m} \cdot \text{yr}^{-1}$ under SSP1-2.6 in about 2066 (Figure 3a). The melt rate has a moderate increasing trend after 2066 and eventually reaches its peak at $16.8 \text{ m} \cdot \text{yr}^{-1}$ in 2100 under SSP5-8.5, while it stays stable around $8 \text{ m} \cdot \text{yr}^{-1}$ after 2066 except for a decrease to $4 \text{ m} \cdot \text{yr}^{-1}$ in 2094 under SSP1-2.6

170 (Figure 3a).

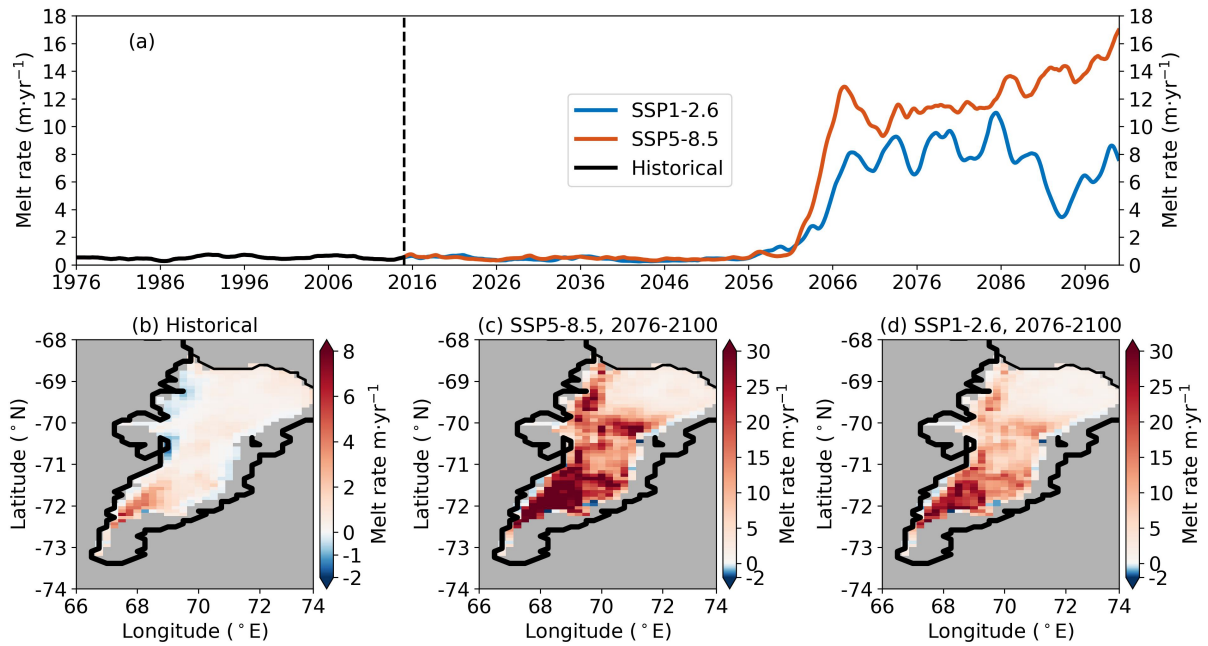


Figure 3. (a) Time series of the area-averaged melt rate ($\text{m} \cdot \text{yr}^{-1}$) from 1976 to 2100. A 12-month-running-average is applied. The dashed vertical line indicates the start of 2015. The time-averaged basal melting over the period of (b) 1976-2014 (Historical), (c) 2076-2100 under SSP5-8.5, (d) 2076-2100 under SSP1-2.6. The warm/cold colours represent melting/refreezing, respectively. Note the different colormap ranges. The colormap for (c) is saturated.

Figure 3b-d shows the spatial distributions of the AmIS basal melting during the historical, and before and after the increase. In general, after the increase in the 2060s the melting beneath AmIS strongly increases by multiple times and there is almost no refreezing. The time-averaged melt rates over 1976-2014 illustrate that a high melting of $\sim 6 \text{ m} \cdot \text{yr}^{-1}$ occurs near the grounding line and most of the ice shelf experiences melting smaller than $2 \text{ m} \cdot \text{yr}^{-1}$ (Figure 3b). The freezing mainly occurs under the northwestern ice shelf with the highest value of $\sim 1.6 \text{ m} \cdot \text{yr}^{-1}$ near 69.5°E , -70.5°N (Figure 3b). Smaller freezing rates ($<0.2 \text{ m} \cdot \text{yr}^{-1}$) occur along the ice shelf edges (Figure 3b). This pattern fits the understanding of the buoyancy-driven meridional overturning circulation, in which warm and dense inflows downwell in the eastern ice shelf cavity while cold and fresh outflows upwell in the western cavity. The time-averaged melt rates during 2075-2100 under SSP5-8.5 show that the melting occupies the entire ice shelf with the highest melt rates of $42 \text{ m} \cdot \text{yr}^{-1}$ near the grounding line (Figure 3c). Most of the areas experience melting between 15 and $35 \text{ m} \cdot \text{yr}^{-1}$ excluding the northeastern ice shelf which exhibits the relatively lower melting of $\sim 5 \text{ m} \cdot \text{yr}^{-1}$ (Figure 3c).

The time-averaged melt rates over 2015-2055 under SSP1-2.6 (Figure 3d) have similar behaviour to those under SSP5-8.5 (Figure 3c). The spatial distributions of melt rates over 2075-2100 for SSP1-2.6 are similar to those for SSP5-8.5, but with a smaller melt rate of $\sim 30 \text{ m} \cdot \text{yr}^{-1}$ near the grounding line and $10\text{-}25 \text{ m} \cdot \text{yr}^{-1}$ over the central and the northwestern ice shelf (Figure 3d).

The melt rates in our simulations show higher sensitivity compared with other studies, for example, Naughten et al. (2018); Kusahara et al. (2023). This is inherited from the global UKESM1.0-LL model. The model forcing in Naughten et al. (2018) and Kusahara et al. (2023) is taken from the ACCESS-1.0 and MIROC-ESM, respectively. Meehl et al. (2020) and Forster et al. (2020) have quantified the climate sensitivity of these models and suggest UKESM1.0-LL has higher climate sensitivity than the other two models, which might result in a warmer ocean temperature (Sellar et al., 2020). This suggests our regional model, forced by the outputs from the climate-sensitive UKESM1.0-LL, has a stronger response to climate warming.

3.1.2 Warming on the continental shelf

For the analyses below, the continental shelf and the shelf sea are defined by the areas between ASC and the AmIS calving front/coastal line (Figure 1). The cavity is defined as the ocean beneath the AmIS draft.

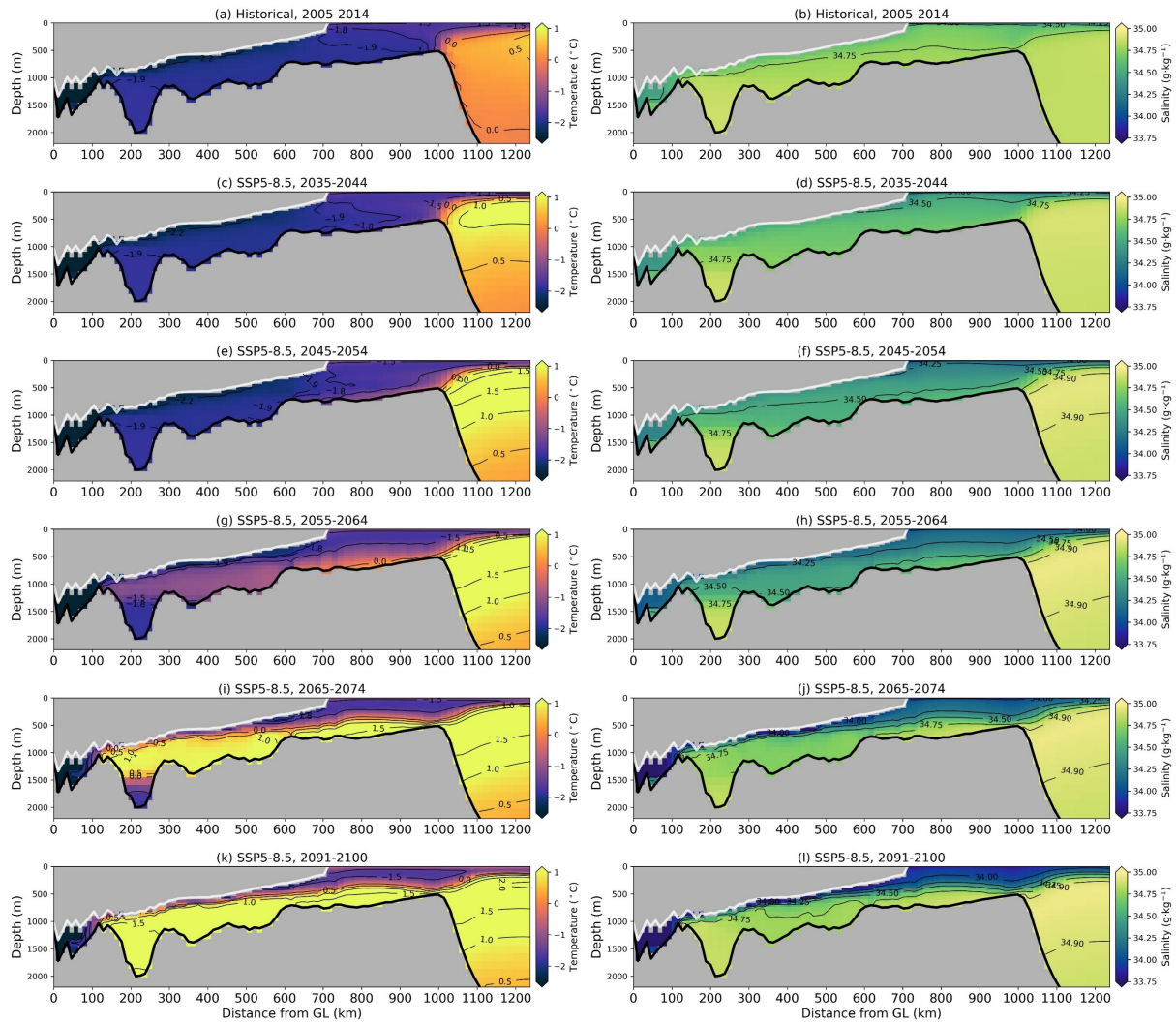


Figure 4. (a) Time series Time-mean of the area-averaged melt rate ($\text{m} \cdot \text{yr}^{-1}$) temperature, salinity at a transect extending from 1976 to 2100. A 12-month-running-average is applied. The dashed vertical line indicates the start of 2015. The time-averaged basal melting over the period of grounding line (bGL) 1976-2014 (Historical), (c) 2076-2100 Prydz Bay Gyre and the shelf break under SSP5-8.5, (d) 2076-2100 under SSP1-2.6 scenario. The warm/cold colours represent melting/refreezing, respectively. Note coordinates show the different colormap ranges distance from GL against the depth below sea level. The colormap for (c) location of the transect is saturated shown by the white line in Figure 1. The years of the time-mean are shown in the title of each panel.

195 3.1.3 Warming on the continental shelf

Figure 4 a shows the time-series of temperature between 300 and 800 m on the continental shelf. The shelf break is generally defined at latitude of 67°S . There is a projected warming and freshening on the continental shelf and inside the AmIS cavity. Figure 4 shows temperature under both SSP scenarios is steady before the late 2030s with the value of roughly and salinity at a transect from the grounding line (GL) to the shelf break (the white line in Figure 1) under SSP5-8.5 scenario. The warm mCDW off the slope is isolated at the shelf break during the years 2005-2014 (Figure 4a). The shelf sea is generally below -1.8°C at depths with a salinity of $34.5\text{-}34.75\text{ g}\cdot\text{kg}^{-1}$ between 2005 and 2014 (Figure 4a, b). The cavity temperature is below -1.9°C which is slightly above the surface freezing point and the salinity at the deeper cavity is higher than $34.75\text{ g}\cdot\text{kg}^{-1}$ (Figure 4a). It gradually increases by ~ 2.4 , b). This suggests that the cavity is filled by HSSW between 2005 and 2014. mCDW becomes warmer from the years 2035-2044 and penetrates across the shelf break, resulting in an increase in the shelf sea temperature to $< -1.5^{\circ}\text{C}$ to $\sim 0.6^{\circ}\text{C}$ in the late 2060s then increases to $\sim 0.8^{\circ}\text{C}$ in 2100 under SSP5-8.5 (Figure 4a). The temperature increases by 1.8°C to 0°C in the late 2060s and keeps steady until 2100 but with a decrease to $\sim 0.9^{\circ}\text{C}$ in about 2094 under SSP1-2.6 c). This is accompanied by a decrease in the shelf sea salinity to $< 34.5\text{ g}\cdot\text{kg}^{-1}$ (Figure 4ad). The temperature changes are similar to the melt rate changes cavity temperature has little change with a slight freshening in the shallower depths (Figure 4a). However, the increase in temperature (in the 2030s) happens several decades prior to that c, d). From the years 2045-2054, mCDW on the seabed of the shelf sea spreads southward and intrudes into the cavity (Figure 4e), causing the cavity temperature increase to $\sim -1.5^{\circ}\text{C}$ during the years 2055-2064 (Figure 4g) and salinity decreases to $34.25\text{-}34.50\text{ g}\cdot\text{kg}^{-1}$ except for the deepest cavity (Figure 4h). This suggests that the dominant water mass in the cavity changes to mCDW between 2055 and 2064, consistent with the abrupt increase in melt rate in the melt rates (in the 2060s. There is further warming in the entire area (Figure 4i, k). The delayed increase in melting might suggest delayed processes connecting temperature and melting such as the passage of warm water into the cavity cavity is flushed by even warmer mCDW of $0.5\text{-}1.5^{\circ}\text{C}$ during the years 2091-2100. The enhanced warming and melting of the ice shelf result in strong freshening in the upper cavity and the shelf sea, forming strong stratifications at depths of 300-500 m (Figure 4j, l).

The spatial pattern of temperature before and after the increase in the 2060s is shown in Figure 4b-d. In general, after the increase, a warming occupies the entire domain, and the AmIS transforms from a cold cavity to a warm cavity. The time-depth-averaged temperature in the open ocean in the Historical simulations decreases from changes in water masses in the model domain under the SSP5-8.5 are illustrated in Figure 5. The definition of water mass is given in Table 1. The water mass definitions are primarily based on Galton-Fenzi (2009) with modifications to adapt to our simulations. The plots show temperature against salinity on each grid cell inside the ice shelf cavity (the solid dots) and on the continental shelf (the open triangles). The black lines show the dilution relation when ice is melted by the warmest water mass in the cavity. During the years 2005-2014, there is a relatively colder mCDW on the continental shelf below 300m, which might induce the melting of the shallowest ice shelf (Figure 5a). HSSW occupies the cavity and controls the ice shelf melting at depths below about 800 m (Figure 5a). ISW clusters at the shallower depth of $\sim 0.5\text{-}300\text{ m}$ inside the cavity (Figure 5a). ISW is also found between 800-2000 m following the Gade line with the source water of HSSW, suggesting this is the outflow of meltwater caused by

Table 1. Water mass definitions based on temperature (T , $^{\circ}\text{C}$), salinity (S , $\text{g}\cdot\text{kg}^{-1}$), potential density (ρ , $\text{kg}\cdot\text{m}^{-3}$) and depth (m) characteristics. T_f is the surface freezing point shown by the dashed line in Figure 5. AASW: Antarctic Surface Water; mCDW: modified Circumpolar Deep Water; HSSW: High Salinity Shelf Water; LSSW: Low Salinity Shelf Water; ISW: Ice Shelf Water.

Water	$T(^{\circ}\text{C})$	$S(\text{g}\cdot\text{kg}^{-1})$	$\rho(\text{kg}\cdot\text{m}^{-3})$	Depth (m)
AASW	$T > T_f$	$S < 34.0$	\sim	Depth < 200
mCDW	$T > -1.75$	\sim	$27.6 \leq \rho < 27.8$	Depth > 200
HSSW	$T_f < T < -1.75$	$S > 34.5$	\sim	\sim
LSSW	$T_f < T < -1.75$	$S < 34.5$	\sim	\sim
ISW	$T < T_f$	\sim	\sim	\sim

HSSW (Figure 5a). The temperature of Antarctic Surface Water (AASW) is between -1 and -1.7 $^{\circ}\text{C}$ near the northwestern boundary to ~ -1.9 $^{\circ}\text{C}$ near the calving front, (Figure 5a). During the years 2035-2044 and the steep gradients exist along the shelf break (Figure 4b). The temperature in the ice shelf cavity of AmIS increases from 2045-2054, AASW becomes warmer and more stratified. The temperature of mCDW increases with the increase in the existence of mCDW in the cavity. HSSW still dominates the ice shelf melting (Figure 5b, c). However, the distributions of HSSW in the entire domain are less. HSSW becomes fresher and transforms to Low Salinity Shelf Water (Figure 5b, c). From the years 2055-2064, there is little HSSW found neither within the cavity nor on the continental shelf (Figure 5d-f). mCDW gradually occupies all depths below ~ -2.2 300 m inside and outside the cavity and controls the ice shelf melting (Figure 5d-f). Less ISW is formed in the shallower cavity, and it becomes fresher with potential density decreased from $27.6-27.8 \text{ kg}\cdot\text{m}^{-3}$ during years 2005-2014 to $< 27.2 \text{ kg}\cdot\text{m}^{-3}$ from the years 2065-2074 onwards (Figure 5d-f). The meltwater on the Gade line, driven by the warm mCDW intruding the deepest cavity, becomes substantially warm ($-2.5-1$ $^{\circ}\text{C}$ in the northern cavity to ~ -2 $^{\circ}\text{C}$ near the grounding line, and is also beneath -2 $^{\circ}\text{C}$ in the West Ice Shelf (WIS) cavity (Figure 4b). This indicates that the sub-ice shelf cavity is dominated by Dense Shelf Water (DSW) with depths from the years 2065-2074 (Figure 5d-f). In addition, the water mass in the shallower cavity (< 500 m) is buoyant ($26.6-27.8 \text{ kg}\cdot\text{m}^{-3}$). This implies that the warmed and freshened outflow of meltwater increases the buoyancy inside the cavity, enhancing the cavity circulations. AASW becomes warmer and its salinity and potential density spreads shrink from the years 2055-2064 onwards (Figure 5d-f), indicating that the stratification at the surface is enhancing.

In comparison with the Historical (1976-2014), the temperature

The evolution of water masses suggests that the AmIS cavity starts to transition from a cold regime to a warm regime in the years 2055-2064 (Figure 5d) and eventually becomes a warm cavity from the years 2065-2074 (Figure 5e, f).

Time evolution of the depth-mean cavity temperature, salinity and the shelf sea temperature, salinity under SSP5-8.5 is warmer as expected (Figure 4c). The warming is generally 0.5 shown in Figure 6. mCDW primarily occupies depths below 300 m in our simulations (Figure 5), so we choose a depth range of 300 m and below for the shelf sea properties to capture changes in mCDW. Figure 6a shows the timeseries of temperature. The shelf sea temperature begins to increase from about -1.5 $^{\circ}\text{C}$ over the northwestern open ocean, and increases to more than 1 in the 2030s to about 0.3 $^{\circ}\text{C}$ over the northeastern flank (Figure 4c). This pattern implies that the source of warming might come from the eastern ocean boundary and extend westward along

the ASC and then penetrate to the ice shelf front. What is striking in Figure 4c is that the warming in most of the AmIS cavity is higher than 2 °C in the late 2060s, followed by a slower increase to above 0.5 °C, even above 2.5 °C in 2100 (Figure 6a). By contrast, the cavity temperature is stabilized at about -2 °C in the southern cavity before the 2050s and then has an abrupt increase to about 0 °C in the late 2060s and a continuous increase to about 0.4 °C in 2100 (Figure 6a). The timeseries of the cavity temperature is consistent with that of the melt rate (Figure 3a). There is a roughly 20-year delay between the increase in the shelf sea temperature and the increase in the cavity temperature as well as the melt rate. The changes in salinity in the shelf sea salinity and the cavity are similar and more synchronised (Figure 3b). The two salinity fluctuate between 34.6 and 34.8 g·kg⁻¹ and start to decrease in the 2030s and then divergent in the 2040s when the cavity salinity has a larger decrease (Figure 3b). The difference is likely due to the reduction in sea ice, which decreases salinity and the formation of HSSW, while the increased mCDW is present on the continental shelf but has not yet reached the deeper cavity (Figure 5b, c). The shelf sea salinity and the cavity salinity have a slight bounce in the 2060s and are stable until 2100 (Figure 3b), suggesting a denser mCDW occupies the entire domain from the 2060s.

The delay of the cavity temperature increase and the larger freshening of the cavity salinity might suggest delayed processes connecting warming on the continental shelf and ice shelf melting such as the passage of warm water into the cavity. This is correlated with the very high melting during the last decades of the 21st century in the SSP5-8.5.

The projected temperature changes changes in temperature, salinity, water masses under SSP1-2.6 scenario (Figure 4d) mirror have minor differences from those under SSP5-8.5 scenario (extra figures can be found in supplementary material). The delay of the cavity temperature is found in the SSP1-2.6 simulations as well, with a smaller degree of warming and freshening. Likewise, a warming of >0.75 °C near the eastern boundary spreads westward, and a much higher warming of >1.5 °C occurs in the AmIS cavity (Figure 4d).

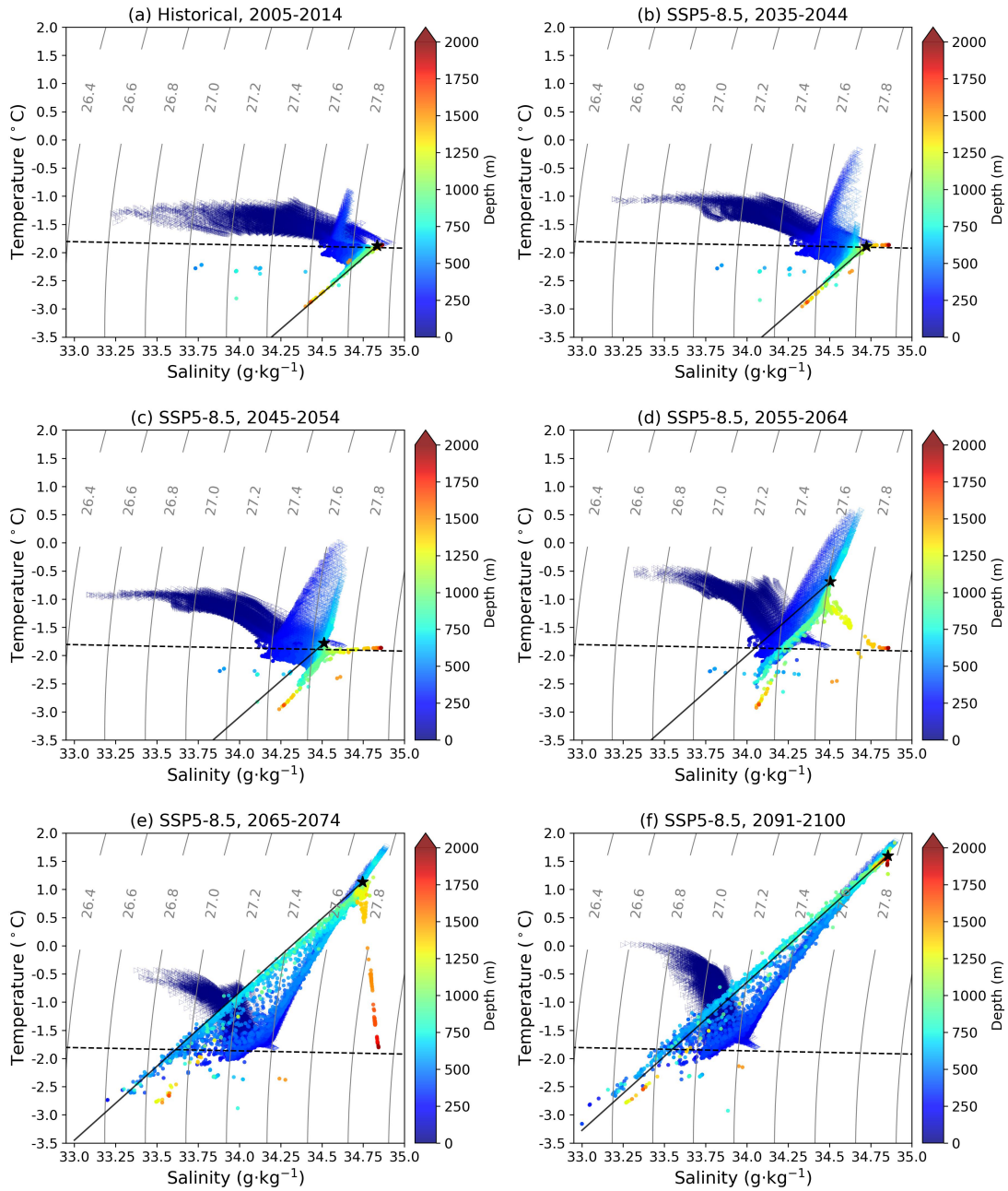


Figure 5. (a) Time-series-Time-mean of averaged temperature between 300 and 800 m and on against salinity diagram for each grid cell within the continental ice shelf cavity ($^{\circ}\text{C}$) from 1976 to 2100. A 12-month-running-average is applied. The dashed vertical line indicates the start of 2015. (b) (solid dots) The average of full-depth temperature ($^{\circ}\text{C}$) over and on the period of 1976-2014 continental shelf (Historical). (c) The differences ($^{\circ}\text{C}$) between the average of full-depth temperature over the period of 2076-2100 open triangles) under the SSP5-8.5 and that of 1976-2014 scenario. The warm colours indicate that temperature is projected to increase grey solid lines are potential density. The thin black lines represent dashed line shows the temperature differences of 0.5°C , 0.75°C , 1°C , 2°C , 2.5°C surface freezing point. (d) The same as black solid line represents the Gade line, defined by end members of the warmest water mass in the cavity (e) but the SSP1-2.6 experiment black pentacle). The thin black lines colour schemes show the temperature differences depth of 0.5°C , 0.75°C , 1°C , 1.5°C , 2°C each grid cell. The thick black line represents years of the coast and time-mean are shown in the thick grey lines show the ice shelf front title of each panel.

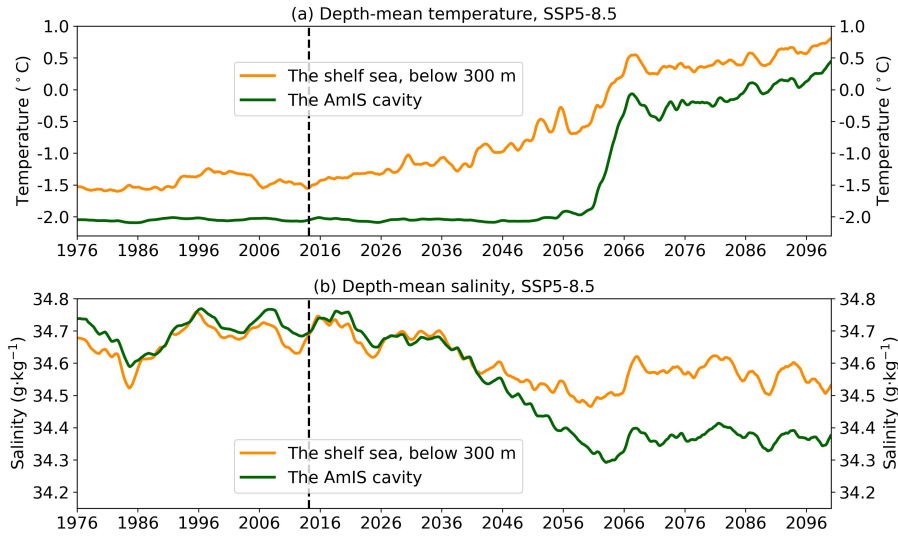


Figure 6. Timeseries of depth-mean (a) temperature and (b) salinity averaged within the ice shelf cavity (the green line) and on the continental shelf below 300 m (the orange line) under the SSP5-8.5 scenario. The vertical dashed line indicates the year of 2015. A 12 month running average is applied.

3.2 Mechanism causing the delayed increase in AmIS melting

275 3.2.1 Reversed current induced increase in heat ~~transport~~ into the cavity

~~The previous section has shown that the basal melting of AmIS is projected to dramatically increase after 2060 in both scenarios. The following section will explain the reasons for the increased basal melting which is-~~

280 Figure 7 shows the barotropic stream function (BSF) in the model domain under the SSP5-8.5 scenario. As the strength of BSF inside the cavity is much smaller than that outside the cavity, we present them with two different colour schemes. The red-blue colour schemes are for BSF on the open ocean and the orange-purple colour schemes are for BSF in the cavity. The positive/negative BSF represent anti-clockwise/clockwise circulation. The BSF of PBG is positive, suggesting PBG is anti-clockwise with a strength of ~ 30 years after the onset of warming. According to Liu et al. (2017), PBG and PBECC are the two main import pathways of mCDW into the Amery ice-shelf cavity. Here, we calculate the heat transport (HT) across two zonal transects (0.5 Sv and the PBG main flow (indicated by the purple and yellow lines in Figure 5). The meridional HT

285 relative to the surface freezing point (Q) ~~solid purple line in Figure 7a) is :-~~

$$Q = \rho C_p \int \int (T - T_f) v dz dx,$$

~~where ρ is the seawater density, C_p is the specific heat capacity, T is the seawater temperature, T_f is the surface freezing point temperature and v is the meridional velocity. T and v use the time-mean monthly model outputs.~~

Figure 6 illustrates the HT of the main flow of PBG, PBECC and TOTAL. The TOTAL is defined as the sum of HT of PBG main flow and PBECC. The HT of PBECC under SSP5-8.5 is always offshore and stable below $0.5 \times 10^{12} \text{ J} \cdot \text{s}^{-1}$ until 2030 and gradually increases to $2 \times 10^{12} \text{ J} \cdot \text{s}^{-1}$ by 2100 (Figure 6). The BSF near the AmIS calving front is positive with values of 0.1-0.2 Sv (Figure 7a), indicating one of the HSSW inflows is in the western ice shelf front during the years 2005-2024. The BSF in the deeper cavity is negative with values of -0.1 Sv during the years 2005-2024 (Figure 7a). This suggests that the cavity circulation is clockwise with an HSSW inflow in the eastern ice shelf front. During the years 2035-2044, 2045-2054, the positive and anti-clockwise PBG is gradually weakened, disappears and transition to the negative BSF (Figure 7b, c). This is accompanied by the weakening of the cavity circulation (Figure 7b, c; Figure 8a), which is probably due to HSSW becoming less efficient in driving the melting. During the years 2055-2064, the negative BSF on the continental shelf is increased (Figure 7d) and the PBG strength is enhanced (Figure 8b), although the clockwise PBG has not yet been established (Figure 7d). The positive cavity BSF near the ice shelf front changes to negative, and the negative BSF in the deeper cavity is increased (Figure 7d) and the barotropic circulation in the cavity is strengthened during the years 2055-2064 (Figure 8a). By contrast, the HT of PBG main flow under SSP5-8.5 is stable ($< 0.5 \times 10^{12} \text{ J} \cdot \text{s}^{-1}$) but offshore until 2050 and starts to be unstable before it changes to be permanently onshore in 2064 along with a dramatic jump of $\sim 10 \times 10^{12} \text{ J} \cdot \text{s}^{-1}$ in 2066. This suggests that the circulations driven by HSSW are very weak or non-existent, and the cavity circulation is controlled by mCDW. From the years 2065-2074, the clockwise PBG is well-established (Figure 7e, f) and the PBG strength is drastically increased (Figure 8b). The cavity circulation is greatly strengthened (Figure 7e, f; Figure 8a).

The evolution of the BSF distribution under SSP1-2.6 scenario behaves in a similar way to that under SSP5-8.5. From year 2066 the HT of PBG becomes relatively stable again and eventually experiences a slight increase to $\sim 12 \times 10^{12} \text{ J} \cdot \text{s}^{-1}$ by 2100 (Figure 6a) scenario. However, there is a striking difference that the PBG and the cavity circulation strength decline in the 2090s in SSP1-2.6 while not in SSP5-8.5 (Figure 8). The TOTAL HT is dominated by PBG main flow. It is generally offshore before 2064 then changes to be constantly onshore, with the increasing amplitude achieving $\sim 12 \times 10^{12} \text{ J} \cdot \text{s}^{-1}$ by 2100 (Figure 6a). This shows that it is the reversal of PBG main flow that brings increased heat into the AmIS cavity from the 2060s, which causes the jump of basal melting and a following stabilisation at a state of high melting.

The HT under decline is consistent with the decrease in the melt rate in the 2090s under SSP1-2.6 shows similar features (Figure 6b). However, there is a notable difference in that the HT of PBG temporarily becomes onshore in about 2090 exclusively under SSP1-2.6 (Figure 6b) scenario. We will discuss the variations in Section 4.

To understand the changes in the relationship between ice shelf melting and ocean circulations inside and outside the cavity, we calculated the heat budget in the cavity as in Jourdain et al. (2017). Neglecting diffusion in ice and the interior ocean, the heat flux entering the ice cavity (H_{in}) is simplified into three components: the heat used to melt the ice shelf (H_{lat}), the heat that warms or cools the seawater within the cavity (H_{HC_VAR}), and the remaining heat flux that exits the cavity (H_{out}):

320
$$H_{in} = H_{out} + H_{lat} + H_{HC_VAR} \quad (1)$$

$$H_{in} = \rho_{ref} C_p \iint_{r \in \text{front}, u > 0} u(r, z) (T(r, z) - T_f^0) dr dz \quad (2)$$

$$H_{out} = \rho_{ref} C_p \iint_{r \in \text{front}, u < 0} u(r, z) (T(r, z) - T_f^0) dr dz \quad (3)$$

$$H_{lat} = L_f \iint_{x, y \in \text{draft}} \rho_i m dx dy \quad (4)$$

$$H_{HC_VAR} = \rho_{ref} C_p \iint_{x, y, z \in \text{cavity}} \frac{dT(x, y, z)}{dt} dx dy dz \quad (5)$$

325 ρ_{ref} is the reference seawater density of 1026 kg m^{-3} , ~~which might imply the reversal of PBG main flow is probably not itself irreversible.~~ C_p is the specific heat capacity of $3991.87 \text{ J} \cdot \text{K}^{-1}$. $u(r, z)$ is the velocity perpendicular to the ice shelf front. The positive means velocity into the cavity. $T(r, z)$ is the temperature along the ice shelf front. T_f^0 is surface freezing point. L_f equal to 334 is the latent heat of fusion of ice. $\rho_i m$ is the melt rate in unit of $\text{kg} \cdot \text{m}^{-2} \text{s}^{-1}$. The latent heat of ice shelf (H_{lat}) is obtained from the model outputs. $dT(x, y, z)$ is temperature change on the same grid cell. dt is one month. The variables used
330 in the calculations are monthly mean model outputs.

~~The time evolution of HT under~~ Figure 9 shows time evolution of heat budget in the cavity under SSP5-8.5 and .5 and SSP1-2.6 ~~explains the~~ scenarios. H_{in} starts to increase in about 2040s and abruptly increases in the 2060s under two scenarios. This is coherent with the water mass changes in the cavity (Figure 5) and the development of the BSF of PBG (Figure 7). H_{HC_VAR} is at least one order magnitude smaller than H_{lat} and H_{out} . The increase in H_{lat} begins in the 2060s, and its behaviours are
335 similar to the changes in the strength of the cavity circulation and PBG. This might suggest the increased ice shelf melting enhances the barotropic flow inside the cavity and in front of the ice shelf, as stronger pressure or density gradients are created (Jenkins, 2016; Jourdain et al., 2017). The melt rate-circulation strength coherence is similar to the melt-induced circulation found in the Jourdain et al. (2017) Amundsen Sea study. H_{out} follows and is comparable to H_{lat} , suggesting that about 50% of heat is unused and leaves the cavity.

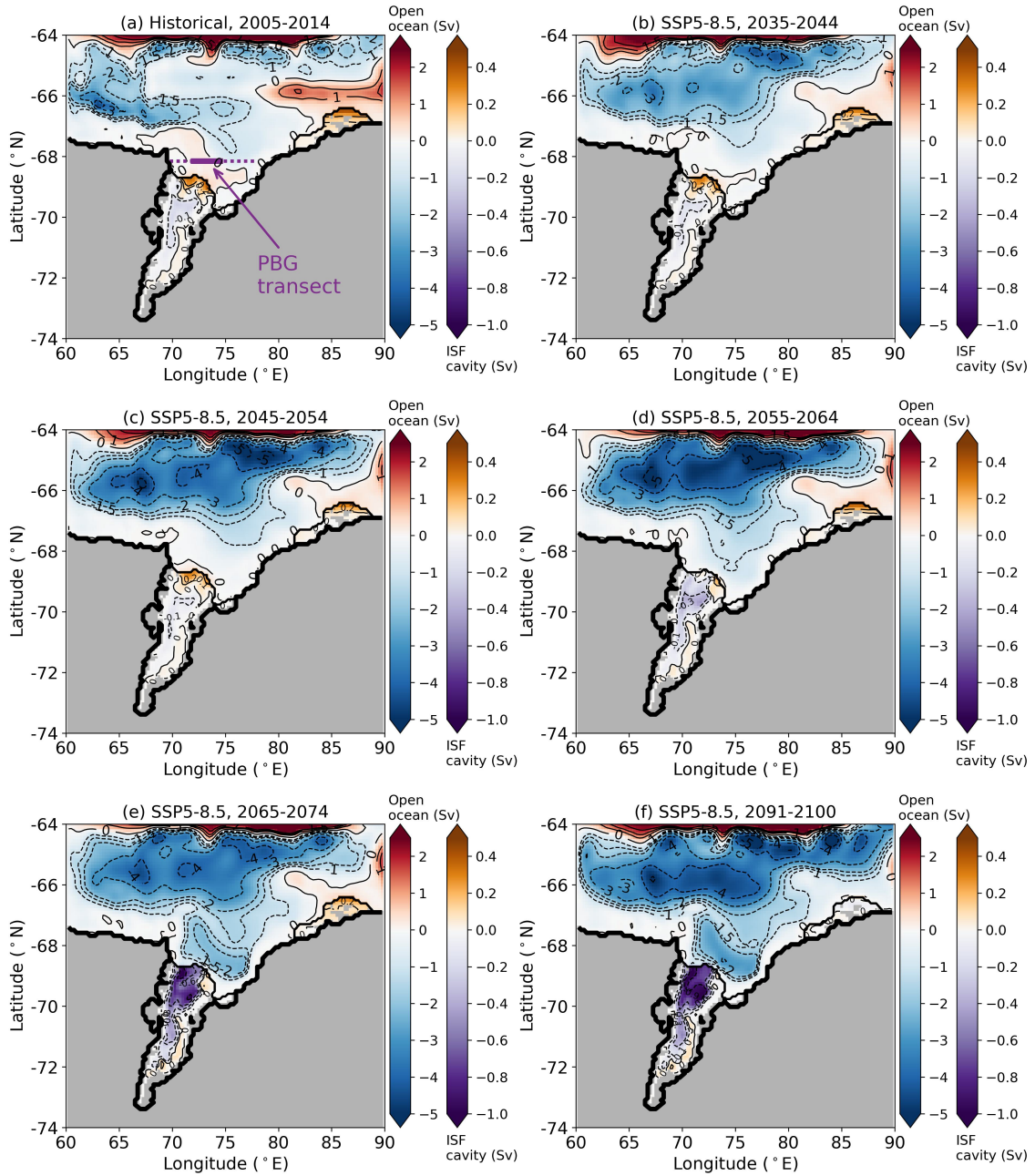


Figure 7. Time-mean of barotropic streamfunction (BSF, in units of Sv) under the SSP5-8.5 scenario. The red-blue colour schemes are for BSF on the open ocean. The values of BSF contours on the open ocean are -5, -4, -3, -2, -1.5, -1, 0, 1, 1.5, 2 Sv. The orange-purple colour schemes are for BSF in the ice shelf cavities. The values of BSF contours in the cavities -0.9, -0.6, -0.3, -0.1, 0, 0.1, 0.2 Sv. The positive/negative BSF represent anti-clockwise/clockwise circulation. The solid and dotted purple line shows the PBG transect and the extended transects to the coasts. The years of the time-mean are shown in the title of each panel.

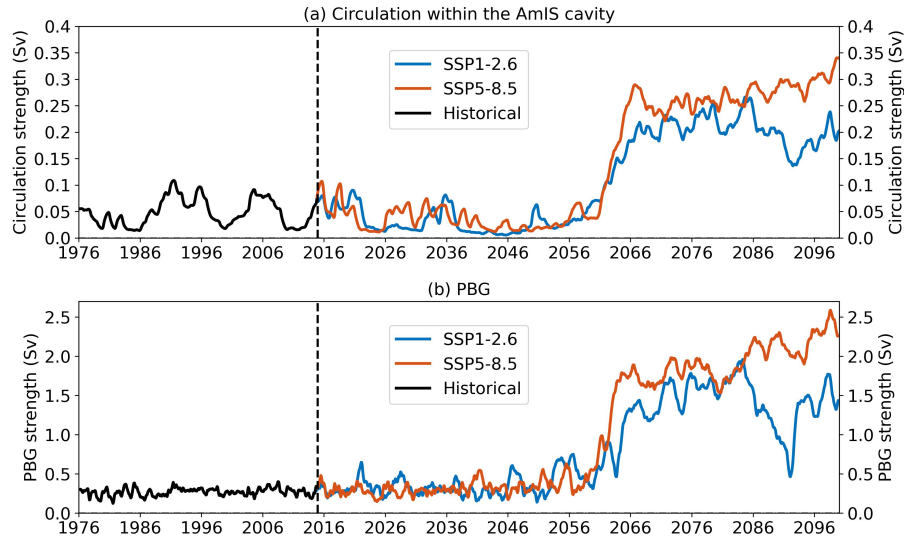


Figure 8. Timeseries of (a) the strength of the barotropic cavity circulation and (b) the strength of PBG. The strength of cavity circulation and PBG is defined as the absolute value of the area-mean barotropic stream function (BSF). The PBG area is defined using the closed BSF contour of -1.5 Sv in front of the ice shelf in Figure 7f. The vertical dashed line indicates the year of 2015. A 12 month running average is applied.

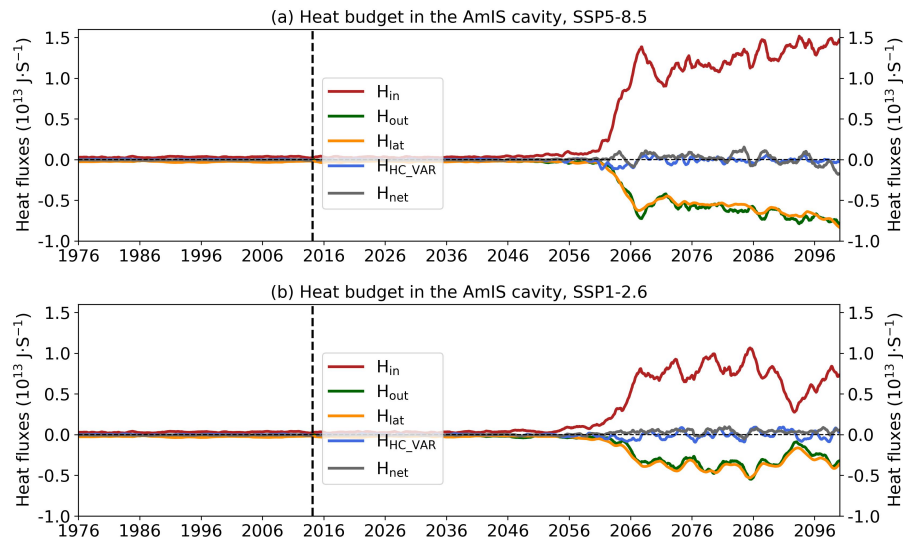


Figure 9. The heat budget of the ice shelf cavity for (a) the SSP5-8.5 scenario, (b) the SSP1-2.6 scenario. The heat of the ice shelf cavity is balanced by the heat flux into the cavity (H_{in} , the red line), the heat used to melt ice or latent heat of ice (H_{lat} , the orange line), the heat exit the cavity (H_{out} , the green line) and the heat used to warm/cool the seawater in the cavity (H_{HC_VAR} , the navy line). The net heat flux (H_{net}) is equal to $H_{in} - (H_{lat} + H_{out} + H_{HC_VAR})$. The vertical dashed line indicates the year of 2015. A 12 month running average is applied.

340 The changes in ocean circulations and heat fluxes in the cavity explain the question posed above. Why does the **increase of**
ocean temperature start several decades earlier than that of melt rates? PBG main flow strongly travels northwestward before the
abrupt increase in the melt rate start ~20 years later than the increase in the shelf sea temperature? PBG is anticlockwise and its
main flow travels off-shore before the 2050s (Figure 7a-c). Despite the warming and increased mCDW on the continental shelf
(Figure 4a, c, e; Figure 5a-c), the processes of mCDW intruding the cavity are slow before the 2060s (Figure 5a), it transports
345 **most of the heat off-shelf despite the increasing southward HT of PBECC (Figure 6a). This prevents the increasing oceanic heat**
arriving at the cavity, and consequently delays the increase of melt rates until its reversal (Figure 8b; Figure 9). An effective pathway
of mCDW into the cavity starts to establish in the years 2055-2064 (Figure 7d) and the clockwise PBG is well established
and sustained from the years 2065-2074 (Figure 7e, f; Figure 8b), which transport massive heat into the cavity (Figure 9) and
transforms the cavity to a warm regime (Figure 5e), resulting in the abrupt increase in melt rate after the 2060s (Figure 9).

350 **Current pattern (colour) and surface stress (arrows) under SSP5-8.5 scenario (a) before reversal (2015-2060) and (b) after**
reversal (2061-2100). The colour scale shows the average of meridional velocity (V , $\text{m} \cdot \text{s}^{-1}$) between 300 and 800 m. The
warm colours show the northward current and the cold colours show southward current. The black arrows represent surface
stress averaged every eight grids ($10^{-2} \text{N} \cdot \text{m}^{-2}$). The thick purple and yellow lines in (a) indicate the transects of the main flow
of PBG and PBECC, respectively. The black boxes in (b) show the locations of local areas S1 and S2 of interest. The northern
355 **boundary of S2 is identical to the PBG transect shown in (a).**

3.2.2 Freshening-driven reversed current

Time series of monthly heat transport (HT, $10^{12} \text{J} \cdot \text{s}^{-1}$) from 1976 to 2100 under (a) SSP5-8.5 scenario and (b) SSP1-2.6
scenario. The purple and yellow lines show the HT of the main flow of Prydz Bay Gyre (PBG) and Prydz Bay Eastern Coastal
Current (PBECC), respectively. The grey shaded areas indicate TOTAL HT. TOTAL = PBG + PBECC. The positive/negative
360 **values represent the off-shore/on-shore HT. The dashed vertical lines indicate the start of 2015. A 12-month-running-average**
is applied.

3.2.3 **Freshening-driven reversed current**

The previous section demonstrated that the reversal of PBG main flow allows the increasing oceanic heat to penetrate the AmIS
cavity, causing increased basal melting. We will now discuss what causes the PBG to reverse. The following analysis does not
365 qualitatively vary between SSP5-8.5 and SSP1-2.6, so we only present the results from SSP5-8.5 here.

Figure 5 illustrates the depth-averaged meridional velocity between 300-800 m and downward surface stress on the top
of ocean before and after the reversal of PBG main flow under SSP5-8.5. It shows that PBG main flow distinctly changes
from northward (Figure 5a) to southward (Figure 5b). This is accompanied by a substantially strengthened meltwater outflow
coming from beneath the AmIS from the 2060s onward (Figure 5b). However, the changes in surface stress do not agree with
370 **the reversal of PBG main flow in our simulations (Figure 5). The surface stress is generally westward over the eastern PB and**
northwestward over the western PB, exhibiting little change with time (Figure 5). The strong northward component of surface
stress over the

Here, we focus on a zonal section of PBG ~~is contrary to the direction of the PBG since the 2060s~~ (Figure 5b) at the Amery Depression (the solid purple line in Figure 7a) and its extensions to the western and eastern bank (the dotted purple lines in Figure 7a). The area between the vertical dashed purple lines represents the PBG transect. Figure 10a shows that warming in the east of PBG begins in the 2040s and gradually spreads westward across PBG until the mid-2060s, when a sharp increase in temperature occurs. The temperature in the west of PBG is always cooler due to the injections of the cold outflow of meltwater (Figure 10a). The ~~wind stress (on the top of ice or ocean when it is sea ice free, not shown)~~ and salinity exhibits notable changes (Figure 10b). There is a higher salinity and potential density in PBG and its western regions than in the east before the 2040s. Then the entire section has a decline in salinity until the mid-2060s, when the west of PBG experiences a sudden freshening, while PBG and its eastern section have an increase in salinity and potential density. This results in a reversed horizontal salinity and density gradient at the western boundary of PBG. Afterwards, the reversed gradient is sustained. As suggested in Jenkins (2016) and Jourdain et al. (2017), the changes in the barotropic coastal current induced by ice shelf melting can be a geostrophic adjustment due to the changes in pressure gradients. Here we present the sea surface height (SSH) gradient in Figure 10c. The SSH gradient mirrors salinity and density changes (Figure 10c). The positive/negative SSH gradient indicates that the SSH in the east is higher/lower than in the west. The SSH gradient is positive in PBG before the 2040s, suggesting there is a high SSH. The SSH gradient exhibits more fluctuations between ~~the surface stress (on the top of ocean) have a similar pattern. This suggests surface stress is not the driver of the flow reversal and that the PBG is more likely to be controlled by density changes.~~

~~Here, we focus on the salinity (hence density) changes in two local areas in Amery Depression, S1 and S2 (Figure 5b). Figure 7a shows that the salinity and potential density in S1 which exhibit a drastic decline from surface to bottom between about 2036 and 2060. After 2060, the salinity and density below 300 m recover slightly. However, above 300 m the salinity and density show a further decrease resulting in the steepened density gradients at depth after the late 2060s (Figure 7) 2040s and 2060s and a reversed SSH gradient in the west following the density variations. The barotropic stream function (BSF) in the section tells a similar story that PBG is positive and anti-clockwise until completely reverses in the mid-2060s (Figure 10d).~~

Figure 11 shows timeseries of the main surface forcing under SSP5-8.5 scenario. Freshwater flux from ice shelf melting exhibits the same behaviours as the properties we analysed above (Figure 11a). The ~~salinity and density in S2 integrated sea ice volume on the continental shelf shows a decrease in the mid-2030s (Figure 11b), which might be the source of the freshening of the PBG and extended section in the 2040s (Figure 10b). The zonal and meridional surface stress does not behave in a similar way (Figure 7b) but with smaller changes. Figure 7c demonstrates the salinity differences between S2 and S1 (hereafter Δ Salinity) and the density differences. There is a substantial redistribution of the salinity and density differences over the period of about 2040 and to the BSF of PBG (Figure 10c, d). The southward Ekman transport in the PBG section remains at a magnitude of 0.2-0.3 Sv throughout the simulation. This is much smaller than the mean BSF of the reversed PBG which increases from 0.5 Sv in the 2060s when the decrease in salinity is happening. Before 2040, the averaged density in the area S2 is generally $\sim 5 \times 10^{-2} \text{ kg m}^{-3}$ less dense than that in S1 to over 2 Sv thereafter (Figure 8). This suggests that surface stress is not a direct factor in the PBG reversal. The geostrophic component is probably more influential. However, the density and salinity in S2 gradually become greater than that in S1 from roughly 2040, and the redistribution spreads from bottom to surface. After~~

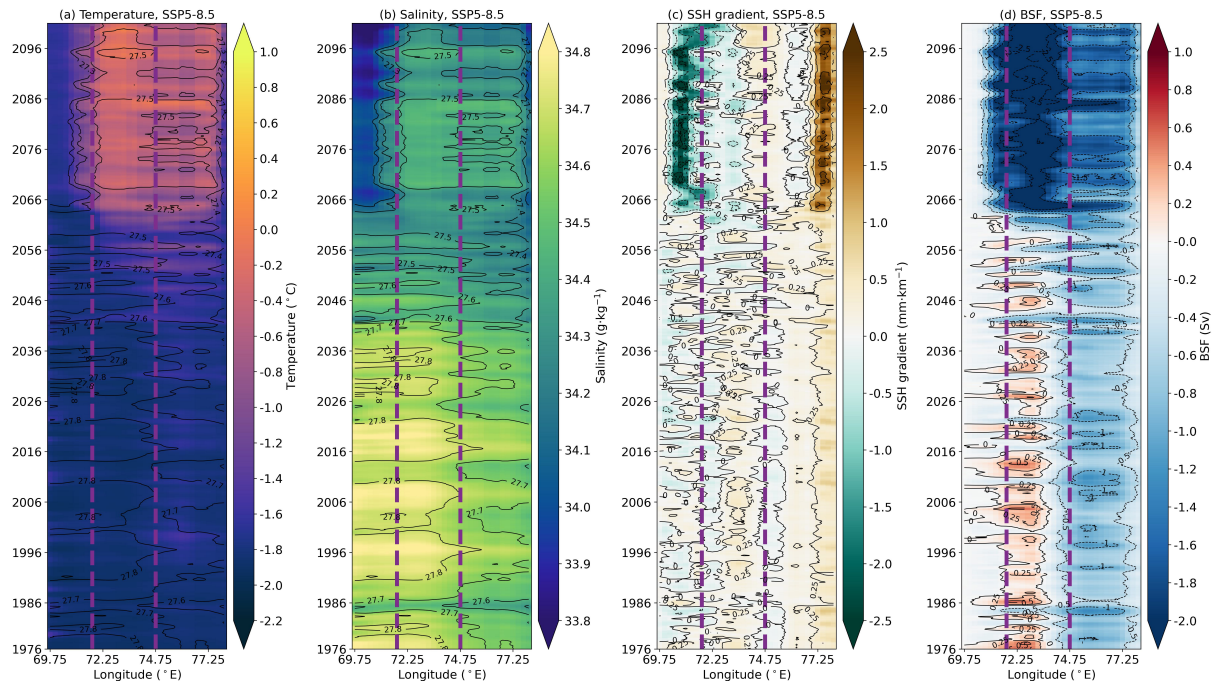


Figure 10. Hovmöller diagram of properties at the zonal transect shown in Figure 7a. (a) Depth-mean temperature below 300 m. (b) Depth-mean salinity below 300 m. The potential density below 300 m overlies temperature and salinity. The values of potential density contours are 27.3, 27.4, 27.5, 27.6, 27.7, 27.8 $\text{kg}\cdot\text{m}^{-3}$. (c) Zonal sea surface height gradients (SSH gradient). A positive/negative SSH gradient indicates that the SSH in the east is higher/lower than in the west. The values of SSH gradient contours are -3, -2.5, -1.5, -0.5, 0, 0.25, 1, 2 $\text{mm}\cdot\text{km}^{-1}$. (d) Barotropic streamfunction (BSF). The values of BSF contours are -2.5, -2, -1.5, -1, -0.5, 0, 0.25, 0.5, 1 Sv. A positive/negative BSF represents anti-clockwise/clockwise circulation. The transect in between the purple dashed lines is the PBG transect shown by the solid purple line in Figure 7a. A 12 month running average is applied.

2060 when the differences between the two areas become stable, the density below 300 m in the S2 is generally $10^{-1} \text{kg}\cdot\text{m}^{-3}$ denser than that in Ekman pumping in the PBG area, which strengthens the upwelling of mCDW onto the continental shelf (Greene et al., 2017), is enhanced between the 2040s and 2070s with large fluctuations afterwards (Figure 11e). The increased Ekman pumping coincides with the increased temperature in the PBG section (Figure 10a) and the reduction in sea ice in the 2040s (Figure 11b). Figure 11 suggests that the enhanced Ekman pumping in the 2040s might be responsible for the freshening in the PBG section and the unstable PBG between the S1. Figure 7d illustrates the averaged meridional velocity of PBG main flow along the zonal transect shown in Figure 5. The meridional velocity is stably northward with a maximum velocity of $6\text{--}8 \text{ cm}\cdot\text{s}^{-1}$ below 500 m before 2040. The PBG main flow becomes weakened and unstable, exhibiting high-frequency variability between about 2040–2060. Afterwards, it becomes consistently southward below 200 m and shows the strong velocity of $6\text{--}8 \times 10^{-2} \text{ m}\cdot\text{s}^{-1}$ below 500 m 2040s and 2060s, and the abruptly increased ice shelf melting further develops and sustains the reversed salinity/density gradient between PBG and its western regions.

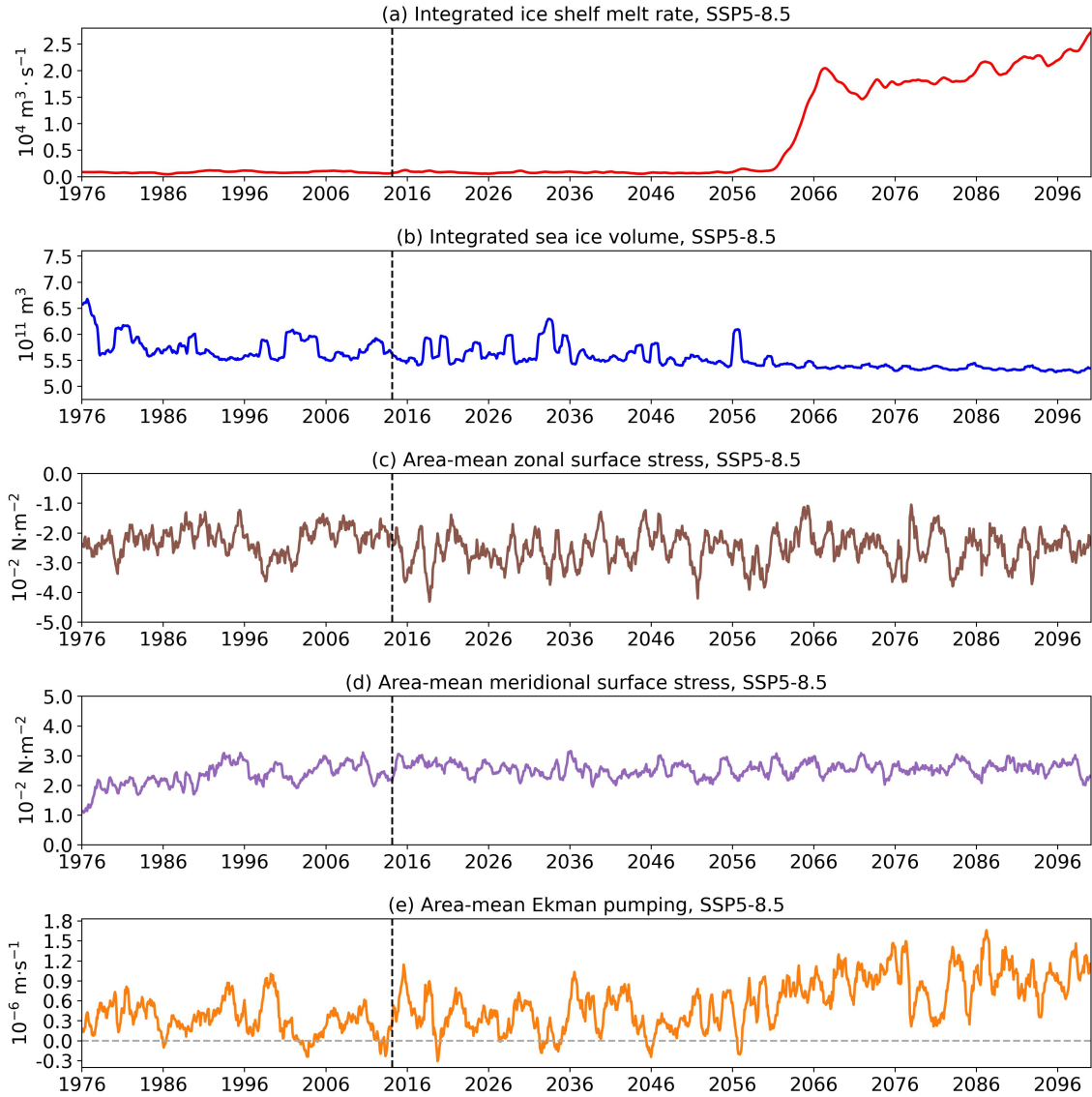


Figure 11. Hovmöller diagram for the spatially averaged salinity–timeseries of (colours, $\text{g}\cdot\text{kg}^{-1}$) and potential density–integrated ice shelf freshwater fluxes (contours, $\text{kg}\cdot 10^4 \text{ m}\cdot\text{s}^{-3}$) in (a) S1, (b) S2 in the SSP5-8.5 experiment. The locations–integrated volume of S1 and S2 are shown in Figure 5b. The range of potential density contours is 27.4–27.9 with an interval of 0.1. sea ice on the continental shelf ($\text{e}10^{11} \text{ m}^3$) The differences of salinity (colour, $\Delta\text{Salinity}$, $\text{g}\cdot\text{kg}^{-1}$) (c) and potential density–area-mean zonal surface stress within the PBG area (contours, $10^{-2} \text{ kg N}\cdot\text{m}^{-3}\cdot\text{s}^{-2}$) between S2 and S1. The blue/red colours and the dashed/solid lines represent that S2 is fresher/saltier and lighter/denser than S1. (d) The time-averaged area-mean meridional velocity–surface stress within the PBG area (N , $10^{-2} \text{ N}\cdot\text{m}^{-2}$), (e) area-mean Ekman pumping within the PBG area ($10^{-6} \text{ m}\cdot\text{s}^{-1}$) of under the SSP5-8.5 scenario. The PBG transect–shown area is defined using the closed BSF contour of -1.5 Sv in front of the ice shelf in Figure 5a7f. The blue/red colours and the vertical dashed /solid lines indicate southward/northward velocity. The range–line indicates the year of velocity contours is -10 to $10 \text{ cm}\cdot\text{s}^{-1}$ with an interval of $2 \text{ cm}\cdot\text{s}^{-1}$. 2015. A 12-month–running–average 12 month running average is applied in all time-series.

Figure 7-10, 11 suggests that the current reversal of PBG is controlled by the spatial salinity differences between S2 and S1.
 420 We suppose that the current is under a geostrophic balance stronger freshening in the west of PBG due to the outflow of ice
shelf meltwater, which caused a reversed density distribution and SSH gradient. PBG is reversed due to the adjustment of the
new pressure gradient. Here we present a simple estimate-scale analysis of the meridional geostrophic-current:

$$fv = \frac{1}{\rho} \frac{1}{\rho_{ref}} \frac{dp}{dx} + \frac{1}{\rho_{ref}} \frac{d\tau_x}{dz}, \quad (6)$$

where $f = -1.4 \times 10^{-4}$ is the Coriolis parameter, v is time-averaged meridional velocity, $\rho_{ref} = 1026 \text{ kg}\cdot\text{m}^{-3}$ is the reference
 425 seawater density, $dx \approx 80 \text{ km}$ is distance between S2 and S1 and PBG and its western regions. dp is the horizontal pressure
 difference. τ_x is the surface stress. Using the hydrostatic hypothesis,

$$dp = gHd\rho, \quad (7)$$

where $g = 9.81 \text{ m}\cdot\text{s}^{-2}$ is gravity and $H \approx 350 \text{ m}$ from Figure 7 is the mean depth 600 m is the depth to the seabed. Using a
 linear equation of state, and assuming that temperature difference do not produce a significant change in density (Figure 10a),

$$430 \quad d\rho = bdS, \quad (8)$$

with $b = 0.78 \text{ kg}\cdot\text{m}^3$ per thousand, dS is the spatial salinity difference between S2 and S1 PBG and its western regions, and the
 time-depth-mean dS is approximately equal to $-0.1 \text{ g}\cdot\text{kg}^{-1}$ from 1976-2060 and $+0.2 \text{ g}\cdot\text{kg}^{-1}$ from 2060 onward from Figure
 7e10b. And the contribution from the surface stress τ_x is small (Figure 11c). Therefore, the velocity of the time-depth-mean
 geostrophic current is

$$435 \quad v = \frac{gHb}{f\rho} \frac{gHb}{f\rho_{ref}} \frac{dS}{dx} \approx \begin{cases} +4.0 \times 10^{-2} \text{ m}\cdot\text{s}^{-1}, & \text{before 2060} \\ -8.0 \times 10^{-2} \text{ m}\cdot\text{s}^{-1}, & \text{after 2060} \end{cases} \quad (9)$$

which is within the range of modelled values the magnitude of the modelled velocity shown in Figure 7dA1.

The analysis above demonstrates that the current is under a geostrophic balance, and the reversal of the flow reversal of PBG
 is a consequence of the reversal of horizontal salinity differences between S2 and S1 PBG and its western regions.

Figure 8 further illustrates the freshening-driven reversal. Figure 8a and b show the time-mean of dynamic height anomalies
 440 with respect to a reference of 0 dbar before and after the reversal. The dynamic height anomaly is defined as the pressure-integrated
specific volume anomaly (IOC et al., 2010). It represents the geostrophic streamfunction for the difference between the pressure
 level of interest and sea surface (IOC et al., 2010). The dynamic height anomaly was calculated offline with Gibbs-SeaWater
 (GSW) Oceanographic Toolbox (McDougall and Barker, 2011). We show dynamic height anomalies at 500 m because the core
 of the reversed current is at about 500 m (Figure 7

445 The reversal of PBG is also documented by Galton-Fenzi (2009). In his study, the clockwise PBG does not exist in winter because dense water formed in the Barrier polynya blocks the current from the Prydz Bay channel and prevents mCDW from accessing the ice shelf cavity (Galton-Fenzi, 2009). When the Barrier Polynya is active, the clockwise PBG is dumped (Galton-Fenzi, 2009). In our simulation, the seasonality of the PBG reversal is weaker than that in Galton-Fenzi (2009) (Figure 10d).

450 Before the current reversal, the dynamic height anomalies in Amery Depression (AD) vary between -2.5 and $-2.1 \text{ m}^2 \text{ s}^{-2}$ with a maximum of about $-2.1 \text{ m}^2 \text{ s}^{-2}$ close to the western ice shelf front (Figure 8a). This high centre is accompanied by an anti-clockwise circulation around AD, showing the northward PBG flow (Figure 8a). However, This is due to overestimated sea ice in the Barrier polynya region (the white dashed box in Figure 12a) in our simulation (Figure 12), which is inherited from the overestimation of the dynamic height anomalies in AD decrease to -3.5 to $-3.2 \text{ m}^2 \text{ s}^{-2}$ during 2061-2100 (Figure 8 summer

455 sea ice in the model forcing taken from the UKESM1.0-LL outputs (Roach et al., 2020). The modelled climatological summer sea ice concentration (SIC) in the Barrier polynya between 1978-2014 is 50-60% (Figure 12a), while the observed SIC in the Barrier polynya is up to 20% with 30-50% SIC in the upstream regions (Figure 12b). The dynamic height anomalies in the western AD are lowered more than those in the eastern AD. There is a minimum of $<-3.5 \text{ m}^2 \text{ s}^{-2}$ near the western ice shelf front that allows currents to flow through the Prydz Bay channel. A new maximum of $-3.2 \text{ m}^2 \text{ s}^{-2}$ develops to the eastern bank

460 (Figure 8b). The circulation in AD becomes clockwise following the changes in dynamic height (Figure 8b). Figure 8e shows the time series of the longitude of the dynamic height centre and the time series of AmIS freshwater fluxes. The dynamic height centre is defined as the maximum dynamic height anomalies at 500 m within AD circled by the white line. The changes in Our model failed to capture the spatial features. The seasonal climatology of SIC in the Barrier Polynya during the historical period shows that the longitude of the dynamic height centre from $\sim 72^\circ \text{E}$ before 2060 to $\sim 76^\circ \text{E}$ from 2060 onward (Figure

465 8c). The evolution of the centre from the western bank to the eastern bank reflects the reversed horizontal density differences. It coincides with the leap of ice shelf melting modelled SIC in December, January and February is at least twice as high as the observation (Figure 12c). It also overestimates the SIC in other months. However, the SIC after December decreased drastically in the two SSP scenarios, and the January, February and March SIC are below 20% (Figure 12c). This indicates that the overestimated SIC in the Barrier polynya weakens the seasonal reversal of PBG in our simulations. In addition, this also

470 suggests that the overestimated sea ice mitigates the warm intrusions onto the shelf sea as the clockwise PBG is hard to establish well in summer. The reduction in sea ice in the 2040s (Figure 11b) opens a wider gate to the formation of the clockwise PBG. Although sea ice in front of the ice shelf is not the direct trigger for the reversal of PBG in the 2060s (Figure 8e). This suggests that a stronger freshening happens at 72°E because it is nearer the outflow of ISW (Figure 5b), leading to the reversal of horizontal dynamic height in AD. This establishes the new field of dynamic height anomalies in AD. And the enhanced and continuous ISW outflow sustains the reversed horizontal dynamic height anomalies, hence maintaining the clockwise PBG and the reversed main flow. It establishes the necessary preconditions for this event. This provides implications that a decrease in sea ice in Prydz Bay could serve as a climate indicator of the increasing AmIS basal melt. Direct observation of ice shelf basal melting is challenging, but long-term sea ice records are more possible. A decreasing trend of sea ice may signal a forthcoming increase in basal melt rates.

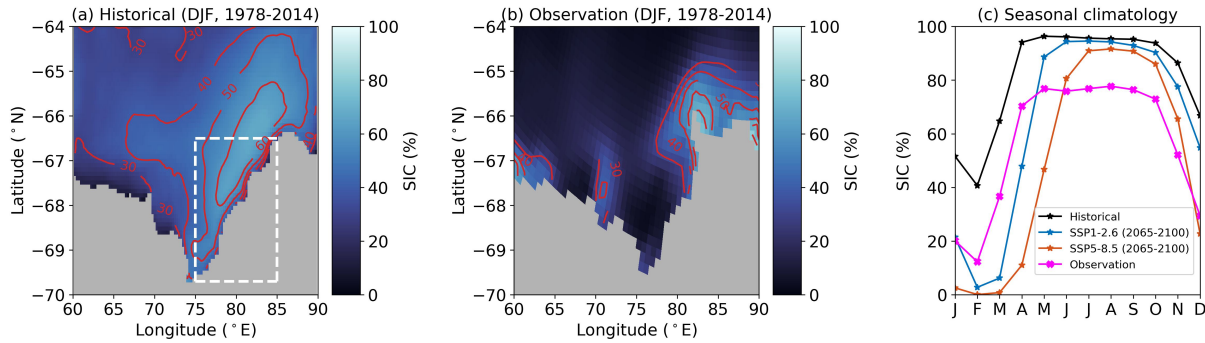


Figure 12. Time-mean summer (DJF) sea ice concentration (SIC) between 1978 and 2014 from (a) historical simulation in this study (Historical). (b) Observation. The observation dataset is Sea Ice Concentrations from Nimbus-7 SMMR and DMSP SSM/I-SSMIS Passive Microwave Data published by National Snow and Ice Data Center (NSIDC). The white dashed box in (a) shows the location of the Barrier Polynya. (c) Seasonal climatology of SIC within the Barrier Polynya. Seasonal climatology of two SSP scenarios is calculated between 2065 and 2100.

480 The mechanism, in which the stronger freshening at the ice shelf front drives the reversal of PBG main flow, is valid under the SSP1-2.6 scenario as well (Figure A1)-A2). The reduced sea ice in the Barrier polynya opens the channel and enables the establishment of the clockwise PBG. This mechanism is self-maintained: the warm water flushes the sub-ice shelf cavity, and the strengthened outflow of ice shelf meltwater reverses the salinity/density differences and stabilises the clockwise gyre, the southward main flow will carry more warm water to the cavity and sustain the feedback loop.

485 Here comes a remaining question: ~~What triggers the mechanism initially? In other words,~~ what causes the ~~unstable horizontal~~
~~salinity differences (Figure 7e and A1e)~~ decline in melt rate (Figure 3a) and the ~~unstable PBG flow (Figure 7d and A1d)~~ between
~~the 2040s and 2060s? On the eastern (upstream) ocean boundary of our model domain, there associated weakening of PBG~~
~~(Figure 8b) in the 2090s under SSP1-2.6 scenario. This is controlled by the variability of the shelf sea temperature in the~~
490 ~~boundary. There is a warming and freshening from 2015 to the 2040s-2050s (Figure A2, and temperature plateaus afterwards~~
~~(Figure A3a). This might suggest trend has a good agreement with the shelf sea temperature changes (Figure 6a) and melt rate~~
~~evolution (Figure 3a). This suggests that the local warming and freshening near the ice shelf front ultimately come from the~~
~~upstream ocean boundary. It is also noted that this mechanism can be interrupted and is reversible. A transient recovery of~~
~~horizontal salinity difference and the northward PBG main flow in~~ In addition, there is a decrease in the shelf sea temperature
495 ~~in the upstream boundary from ~2090 are found in the 1.25 to 0.75°C in the 2090s under SSP1-2.6 experiment (Figure A1e,~~
~~d). This is followed by a decrease in basal melt rate in ~2090 (Figure 3a). The reason for this break of the feedback loop is~~
~~a slight. This decline in temperature forcing results in the decrease in melt rate under SSP1-2.6 (Figure A3a) and a transient~~
~~recovery of PBG (Figure A1b) in the 2090s. This implies that the gyre reversal is not irreversible. A decrease in temperature~~
~~and an increase in salinity at the upstream ocean boundary in the 2080s (Figure A2). This emphasises the influences of the~~
500 ~~upstream systems on the downstream ice shelves could stop this mechanism.~~

To validate the robustness of the freshening-driven mechanism, we conducted a second series of simulations with forcing
taken from the r2 ensemble of UKESM. We chose the r2 ensemble as the mid-depth temperature under SSP1-2.6 scenario
within the entire domain starts to increase ~ 20 years later than that in the r1 ensemble (See Supplement). The melt rate jump
and PBG reversal under SSP1-2.6 scenario also exhibit a delay compared to the r1 simulations (Supplement), highlighting the
505 importance of upstream ocean boundary conditions in our regional configuration.

The time-mean of dynamic height anomalies with respect to pressure of 0 dbar (colours and contours, m^2s^{-2}) is overlaid
by the full model current velocity (red arrows, $\text{m}\cdot\text{s}^{-1}$) under the SSP5-8.5 scenario (a) before (2015-2060), and (b) after
(2061-2100) the reversal of PBG main flow. The dynamic height anomalies and the current velocity are at 500 m. Note the
different colourmap ranges between (a) and (b). The thick purple lines indicate the main flow of PBG. The white lines are
510 the bathymetry of 530 m, showing the area of Amery Depression (AD) in which the dynamic height centre is defined. The
dynamic height centre is chosen as the maximum value in AD. The thick purple lines show the transect of PBG main flow. The
black lines indicate the coast. The grey lines represent the ice shelf fronts. (c) Time series of the annual-mean longitude of the
dynamic height centre (magenta line) and the ice shelf freshwater flux (FWF, green line).

4 Conclusions and discussions

515 This study investigates the future changes of AmIS-PB system under SSP5-8.5 and SSP1-2.6 scenarios. An abrupt increase
in AmIS basal melting is projected to happen in the 2060s under both scenarios. The ~~net area averaged~~ melt rate of AmIS is
projected to increase from $<1 \text{ m}\cdot\text{yr}^{-1}$ in the 2050s to $\sim 8 \text{ m}\cdot\text{yr}^{-1}$ in the late 2060s under SSP1-2.6 and to $\sim 13 \text{ m}\cdot\text{yr}^{-1}$ in the

late 2060s under SSP5-8.5. Afterwards, AmIS is in a high melting state until 2100. The time-mean melt rates over 2075-2100 (when it is in the high melting state) under both SSP1-2.6 and SSP5-8.5 display no refreezing beneath AmIS, and the melting at the grounding line exceeds $30 \text{ m}\cdot\text{yr}^{-1}$. A drastic warming on the continental shelf and in the AmIS cavity causes the increased basal melting. However, the increase of temperature on the continental shelf is in the late 2030s, which happens ~~several decades~~ ~ 20 years before the jump of AmIS basal melting in the 2060s. The delayed increase in AmIS melt rate is due to the reversal of PBG main flow.

PBG plays an important role in ~~heat transport toward AmIS. PBG at the lower depths is not well established, and the changes in the basal melt of the AmIS. The clockwise PBG is not well established due to the sea ice in the active Barrier polynya, which obstructs its formation. The~~ main flow is northward before \sim the 2060s. ~~This prevents the increased oceanic heat, preventing the increased mCDW from intruding~~ into the AmIS cavity and ~~delays~~ delaying the increase in AmIS melting. However, the main flow of PBG reverses southward after the 2060s as the clockwise PBG is well established. The southward main flow imports substantial ~~oceanic heat mCDW~~ into the cavity, which leads to the increase of AmIS melting. The changes in PBG are due to ~~the~~: 1. The reduction in sea ice in the 2030s-2040s, which allows the establishment of the clockwise PBG. 2. The reversal of the horizontal salinity (and then density) differences in front of the ice shelf between the Amery Depression and its western regions. The redistribution of ~~salinity~~ the salinity gradients is established by the strengthened outflow from the AmIS cavity after the 2060s.

~~This freshening-driven mechanism for PBG is also documented by Galton-Fenzi (2009). In his study, the clockwise PBG does not exist in winter because dense water formed in polynyas blocks the current from the Prydz Bay channel and prevents mCDW from accessing the ice shelf cavity (Galton-Fenzi, 2009). In our simulation, the presence of PBG during the historical period is weak and unstable, showing seasonality in its direction at the upper 200 m (Figure A3b). This seasonality coincides with the strong seasonal reversal of the horizontal differences for salinity and density between S2 and S1 at the upper 200 m (Figure A3a). The weaker horizontal salinity difference at the bottom makes it hard to form the clockwise PBG (Figure A3a). However, when the increased ice shelf freshwater freshens the western AD more and reverses the horizontal salinity difference, the clockwise PBG at the lower depths is established after the 2060s (Figure 8). The freshwater forcing to AD helps to form the clockwise PBG, which agrees with Galton-Fenzi (2009). In summary, it is PBG that controls the heat entering the ice shelf cavity, and PBG is determined by the salinity (and density) gradients in AD.~~

A similar mechanism for ~~re-directed currents~~ the regime change is found in ~~front of~~ Filchner-Ronne Ice Shelf (Hellmer et al., 2012, 2017; Naughten et al., 2021; Siahaan et al., 2022) and Ross Ice Shelf (Siahaan et al., 2022). A redirection of ~~inflow~~ coastal current is driven by a reversed density gradient across ~~Filchner Trough~~ the Filchner Trough (Hellmer et al., 2012, 2017; Naughten et al., 2021) / the little America Basin in the Ross Sea (Siahaan et al., 2022), which facilitates the penetration of ~~Warm Deep Water mCDW~~ into the ice shelf cavity and ~~enhances basal melting. The further freshening due to ice shelf melting maintains the reversal of density gradient and the redirected inflow, leading to further warm penetration and then basal melting (Hellmer et al., 2017)~~. ~~This positive feedback mechanism causes the tipping point of basal mass loss of the Filchner-Ronne Ice Shelf in the 21st century. A re-directed flow in front of Ross Ice Shelf also causes an abrupt increase in basal melting in the second half of the 21st century (Siahaan et al., 2022).~~ causes the regime change. The similar processes found in the ~~Filchner-Ronne, Ross and~~

~~Amery sectors~~ sectors for the three large cold ice shelves emphasise the importance of buoyancy changes on the shelf sea in a warming future, ~~in particular for a cold ice shelf~~. This suggests the necessity of long-term records for the shelf sea salinity and sea ice in order to obtain an early warning of the regime change in the three large cold cavities.

555 Most cold ice shelf sectors have a structure with steep isopycnals at the continental shelf break (Thompson et al., 2018). The coastal geostrophic flow along the isopycnals is sensitive to the structure changes (Thompson et al., 2018). In this study, we only focus on the interactions between AmIS melting and local circulation. However, there remain many open questions for future studies. For example, how does the geostrophic flow respond to different components of freshwater fluxes from ice shelves, sea ice, icebergs, advections, precipitations, etc? What is the role ocean currents play in connecting the future changes in the freshwater components? What is the threshold of freshening on the continental shelf for the re-directed or reversed current? Does the threshold vary among different shelf sea sectors? A series of freshwater perturbation experiments across various climate scenarios would help address the above questions.

560 Quantifying the future stability of AmIS and its upstream ice sheets is beyond our research scope, but this study can provide implications to some extent. ~~In our study, AmIS experiences fast basal mass loss during the last decades of the 21st century under the SSP5-8.5 scenario (Figure ??). However, due to the buttressing of the topographic sill at the AmIS grounding line, the fast retreat of the grounding line and large contribution to sea level rise are unlikely to happen (Pittard et al., 2017; Gong et al., 2014)~~ ~~Two previous ice sheet modelling studies (Pittard et al., 2017; Gong et al., 2014) conducted similar extreme experiments by applying enhanced basal melting of AmIS. Both studies suggested that only the collapse of almost the entire ice shelf by un-~~ ~~realistically high basal melting causes the grounding line to retreat beyond the topographic sill. Excluding the most extreme climate scenarios, AmIS attributes to sea level fall (Gong et al., 2014; Pittard et al., 2017). The stability of AmIS and the upstream ice sheets is primarily buttressed by the topographic sill tens of kilometres upstream of the grounding line (Gong et al., 2014; Pittard et al., 2017). In our SSP5-8.5 simulation, the basal melt rate is not as high as that applied in the extreme experiments in Pittard et al. (2017), with $>50 \text{ m}\cdot\text{yr}^{-1}$ and $100 \text{ m}\cdot\text{yr}^{-1}$ for two extreme scenarios.~~ It is reasonable to consider 570 that AmIS will remain stable in the next century, however, the ~~fast thinning of the ice shelf (Figure ??) puts~~ high basal melting puts the AmIS at risk of instability in the longer-term warming future. Coupling our standalone ocean configuration with an ice sheet model will be a future step to address the question of the AmIS stability.

We note the basal melting in this study might be overestimated. In the AME025 configuration, we use velocity-dependent "three-equation" parameterisation of ice-ocean thermodynamics (Jenkins et al., 2010). This parameterisation assumes that ice 580 shelf melting is driven by the turbulent mixing due to the shear currents. However, the turbulent processes at the ice shelf-ocean interface can also be produced by the convection due to buoyancy forcing (Wells and Worster, 2008). For an ice shelf with a stable stratification at the ice shelf-ocean interface, such as AmIS (Rosevear et al., 2022b), the velocity-dependent "three-equation" parameterisation largely overestimates basal melting of AmIS by 200%-400% (Rosevear et al., 2022a). This results in a consequence that the projected melt rate is likely to be overestimated. In addition, UKESM1.0-LL has a higher 585 climate sensitivity compared with other CMIP6 models (Meehl et al., 2020; Forster et al., 2020) and previous generations of climate model (Sellar et al., 2020). This might produce an overestimated and more rapid warming in our regional projections compared with other studies (Naughten et al., 2018). Moreover, we do not include the frazil processes in this configuration.

Galton-Fenzi et al. (2012) suggests the inclusion of frazil in the AmIS simulation decreases the melt rates and increases the arrection rates. No frazil ice in our simulation may result in the overestimation. Another source of overestimation is the static ice shelf draft. Ice shelf melt rate is not only temperature dependent (Holland et al., 2008; Xu et al., 2013) but also basal slope dependent (Payne et al., 2007; Little et al., 2009; Magorrian and Wells, 2016). Steeper slopes might increase the heat entrained into the ice shelf and drive higher melting (Payne et al., 2007; Little et al., 2009; Magorrian and Wells, 2016). When AmIS is thinning, it will become smoother and flatter, and the melt rate is expected to be to some extent decreased. Given the limited grounding line retreat of AmIS (Gong et al., 2014; Pittard et al., 2017), which feeds limited deep and steep grounded ice to the floating AmIS, therefore, the basal melting beneath the majority of AmIS north of the grounding line may be overestimated. Another limitation is that due to the relatively coarse grid spacing of ~ 7 -12 km for our model configuration relative to the estimated Rossby radius of 3 km over Prydz Bay (Liu et al., 2017), we cannot investigate the effect of mesoscale eddies on the AmIS basal melting. Liu et al. (2017) suggested that 52% of the total onshore heat transport across a zonal transect (73-78°E, 67.5°S) in the Amery ice shelf front is induced by mesoscale eddies. Given the importance of mesoscale eddy on warm intrusion beneath Antarctic ice shelves (St-Laurent et al., 2013; Thompson et al., 2018; Stewart et al., 2018, 2019), it is worth employing a higher resolution model (~ 1 km) to understand how mesoscale eddies impact future ice shelf melting.

The basal melt rate is projected to exceed $15 \text{ m}\cdot\text{yr}^{-1}$ beneath the majority of AmIS after the abrupt increase under both scenarios (Figure 3c, d). The thinning of ice shelf results in many changes, for instance, the geometry of the ice shelf cavity, the increased water column thickness under the ice shelf, etc. This is related to a scientific question: How does time-varying ice shelf draft modify ocean circulations in the cavity and ice shelf-ocean interactions in model simulations? Holland et al. (2023) suggests that a time-varying ice geometry of Thwaites Glacier leads to an increase in melting by more than 30% without any change in ocean forcing. However, we use a static ice shelf draft in the AME025 configuration, which limits the ability to investigate such geometrical feedback. Future work would greatly benefit from the development of two-way coupled ocean-ice sheet models and more sophisticated Earth system models (Jordan et al., 2018; Smith et al., 2021; Siahaan et al., 2022; De Rydt and Naughten, 2023).

Appendix A: Extra figures

A1 Extra figures for the SSP1-2.6 experiment V velocity of PBG

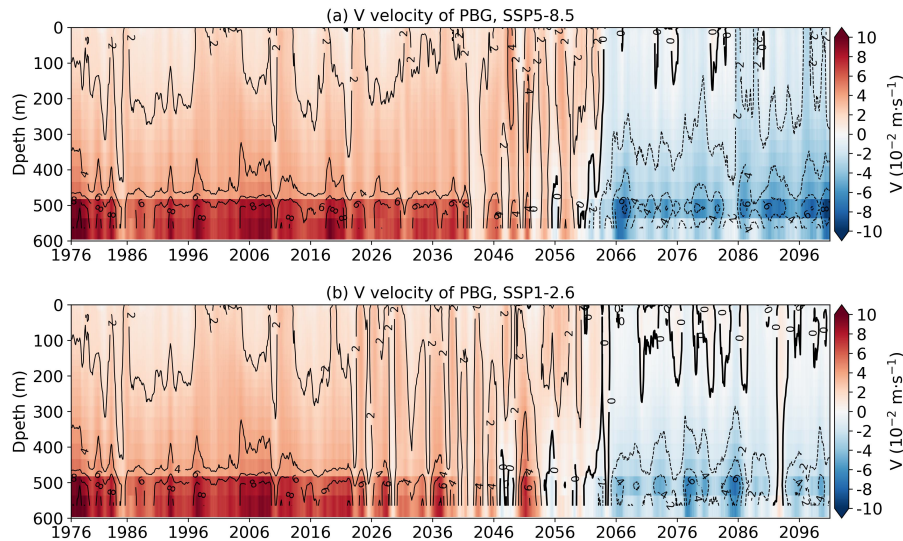


Figure A1. Hovmöller Hovemoller diagram for the spatially averaged salinity (colour, g kg^{-1}) time against depth below sea level) and potential density (contours, kg m^{-3}) of zonal-mean meridional velocity at the solid purple transect shown in Figure 7a. (a) S1, SSP5-8.5. (b) S2 in the SSP1-2.6 experiment. The locations of S1 blue/red colours and S2 are shown in Figure 5b the dashed/solid lines indicate southward/northward velocity. The range of potential density-velocity contours is $27.4-27.9-10$ to 10 cm s^{-1} with an interval of 0.12 cm s^{-1} . A 12 month running average is applied.

A2 Extra figures for the SSP1-2.6 experiment

A3 The forcing of temperature and salinity at the upstream ocean boundary

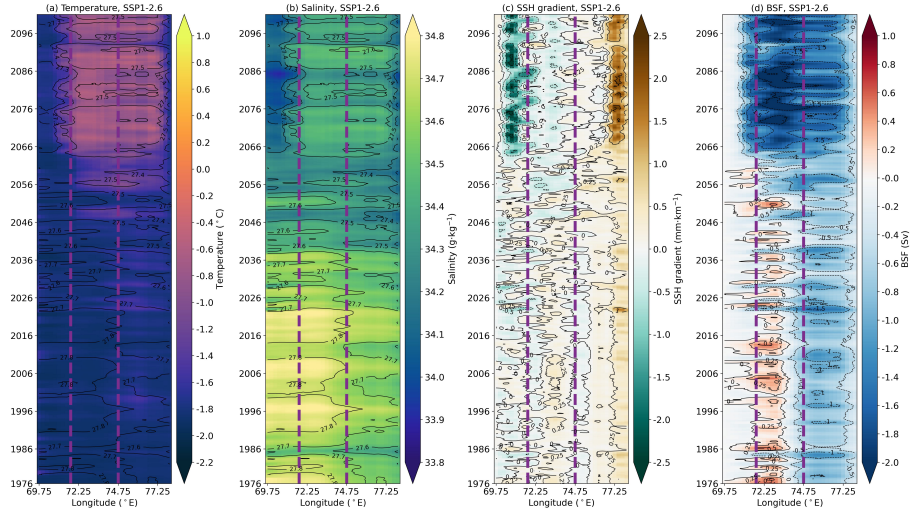


Figure A2. Hovmöller diagram of properties at the zonal transect shown in Figure 7a. (a) ~~The differences of salinity–Depth-mean temperature below 300 m.~~ (colour, Δ Salinity, $\text{g}\cdot\text{kg}^{-1}$) (b) Depth-mean salinity below 300 m. The potential density below 300 m overlies temperature and salinity. The values of potential density (contours are 27.3, 10^{-2} 27.4, 27.5, 27.6, 27.7, 27.8 $\text{kg}\cdot\text{m}^{-3}$). (c) ~~between S2 and S1~~ Zonal sea surface height gradients (SSH gradient). The blue/red colours and the dashed/solid lines represent negative SSH gradient indicates that S2 the SSH in the east is fresher/higher/saltier and lighter/denser/lower than S1 in the west. (d) ~~The time-averaged meridional velocity~~ (V values of SSH gradient contours are -3, $\text{cm}\cdot\text{s}^{-1}$, -2.5, -1.5, -0.5, 0, 0.25, 1, 2 $\text{mm}\cdot\text{s}^{-1}$). (d) Barotropic streamfunction (BSF). The values of BSF contours are -2.5, -2, -1.5, -1, -0.5, 0, 0.25, 0.5, 1 Sv. A positive/negative BSF represents anti-clockwise/clockwise circulation. The transect in between the purple dashed lines is the PBG transect shown by the solid purple line in Figure 57a. The blue/red colours and the dashed/solid lines indicate southward/northward velocity. The range of velocity contours is -10 to $10\text{ cm}\cdot\text{s}^{-1}$ with an interval of $2\text{ cm}\cdot\text{s}^{-1}$. A 12-month-running-average/12 month running average is applied in all time-series.

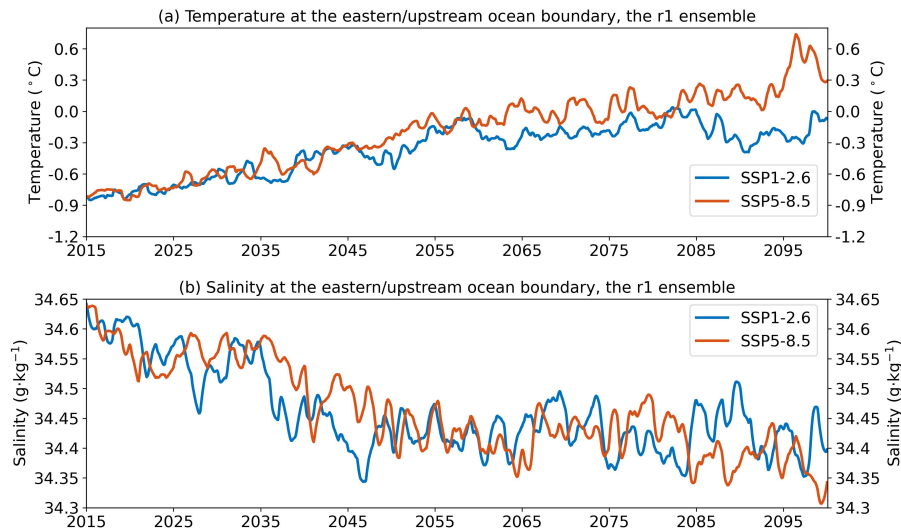


Figure A3. (a) Time series of ~~averaged the area-averaged~~ temperature of ocean boundary forcing at the upstream boundary from ~~1976-2015~~ to 2100. A 12-month-running-average is applied. ~~The dashed vertical line indicates the start of 2015.~~ (b) The same but for salinity ocean boundary forcing.

~~Monthly salinity difference between S2 and S1 and the meridional velocity of the PBG during the historical period (a) Hovmöller diagram for the monthly spatially-averaged differences of salinity (colours, Δ Salinity, g kg^{-1}) and potential density (contours, $10^{-2} \text{ kg m}^{-3}$) between S2 and S1. The blue/red colours and the dashed/solid lines represent that S2 is fresher/saltier and lighter/denser than S1. (d) The time-averaged meridional velocity (V , cm s^{-1}) of the PBG transect shown in Figure 5a. The blue/red colours and the dashed/solid lines indicate southward/northward velocity.~~

~~Basal mass loss of AmIS. Snapshot of AmIS ice shelf thickness (m) in year (a) 2015, (b) 2071, (c) 2086, (d) 2100 under SSP5-8.5 scenario. The red lines represent ice shelf thickness of 0 m, indicating that the ice is removed by basal melting.~~

Code and data availability. The AME025 configuration can be obtained from <https://zenodo.org/records/10797900>. The UKESM forcing used in this study is free to download from the CMIP6 archive (Tang et al., 2019; Good et al., 2019a, b). The AME025 simulation outputs
620 can be obtained from the corresponding author upon request.

Author contributions. JJ conducted the ocean simulations and prepared the manuscript. AP and CB helped to set up the regional model configuration and helped to interpret analysis. All authors commented on the manuscript.

Competing interests. The authors declare that they have no conflict of interest.

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Responses to Reviewer 2

We thank Reviewer 2 for the constructive and thorough review of the manuscript. Below we copy the referee comments in **bold black** and write our responses in **bold blue**. The changes to the manuscript are in *italic blue*.

General comments

This paper addresses the behavior of the Amery Ice Shelf under two climate change scenarios, SSP1-2.6 and SSP5-8.5. While the topic is indeed interesting, the paper falls short of meeting the minimum requirements for publication for several reasons, which I will detail below. The goal of this study is to estimate the future melt rate of the ice shelf by 2100 and to understand the mechanisms driving this melt. The authors utilize a regional simulation, AME025, forced with UKESM outputs. In my assessment, there are significant issues in every part of the paper, with the most critical problems found in section 3.2.

We are grateful to the reviewer for pointing out the problems of our study. We have made the corrections to address the major comments about Section 3. We summarised the modifications here, and the detailed changes can be found in the point-to-point responses.

1 We have added more discussion on the temperature changes and water mass transformations in Section 3.1.2.

The additional analyses include:

- Replacing the depth-mean temperature (the old Figure 4) with temperature transect from the grounding line to the shelf break and beyond as suggested (Figure 4). The location of the transect is shown by the white line in Figure 1.**
- Temperature against salinity diagrams for each ocean cell within and outside the ice shelf cavity (Figure 5) to illustrate the changes in water masses.**
- Separated timeseries of depth-mean temperature/salinity within and outside the ice shelf cavity as suggested by Reviewer 1 (Figure 6).**

2 We have additional analyses of the changes in ocean circulations within and outside the ice shelf cavity in Section 3.2.1. The additional analyses include:

- Replacing the depth-averaged meridional velocity (the old Figure 5) with the barotropic stream function (Figure 7).**
- Timeseries of the strength of circulations within and outside the ice shelf cavity (Figure 8).**

3 We have replaced the analyses of PBG and PBECC heat transport (the old Figure 6) with the analyses of the heat budget in the ice shelf cavity (Figure 9).

4 We have additional discussion on the mechanism for the PBG reversal in Section 3.2.2. The modifications include:

- Replacing the analysis of the two boxes with a zonal transect of PBG (Figure 10).
- Analyses of the forcing of the gyre (Figure 11).
- Discussion on the impact of sea ice on the reversal of PBG (Figure 12).

We also included a new version of the manuscript with tracked changes at the end of this document.

Specific comments

1) Section 2

There are several pieces of missing information:

- Which version of NEMO is used?

Thanks for pointing out the missing details. NEMO 3.6 is used in this study. We have added the version information in Section 2.1:

"The ocean model used in this study is version 3.6 of the Nucleus for European Modelling of Ocean model (NEMO3.6)."

- I understand that the configuration is an extraction of GO7 with some modifications to account for the specifics of a regional study. However, these modifications are not detailed in the text. For instance, at line 78, the authors mention that slip conditions are modified, but it is unclear how and why. What parameters are important here? Why not change the number of vertical levels as in Mathiot et al. (2023)?

Yes, the AME025 configuration is extracted from GO7. A free slip lateral momentum boundary condition for land and a partial slip condition for the Antarctic coast are used in GO7. However, GO7 is a global ocean configuration with open but non-interactive ice shelf cavities (Storkey et al., 2018). We changed the lateral momentum boundary condition to non-slip for land, coastal line and ice shelf interfaces in the light of other regional ice shelf studies such as Bull et al. (2021); Jourdain et al. (2017); Mathiot et al. (2017). In addition, the parameters of the ice shelf thermodynamics are considered primarily based on these ice shelf studies. It would be useful if the vertical levels could be increased to 121 levels. However, it is impractical in NEMO3.6. We have added the reasons why we changed the parameters in the manuscript:

"The schemes and parameter values used in parameterisations are primarily based on GO7 (Storkey et al., 2018), but the values of lateral diffusivity and viscosity have been changed according to the smallest grid spacing (~7 km) and time step (720s). Some physical schemes, such as the slip condition for the lateral boundary, from other regional modelling

studies (Mathiot et al., 2017; Jourdain et al., 2017; Bull et al., 2021) are taken into account as well."

- **The viscosity and diffusion coefficients are set to $-1.08\text{e-}10 \text{ m}^4/\text{s}$ and $135 \text{ m}^2/\text{s}$, respectively. First, I suspect there is an error with the viscosity value, and it should be $-1.08\text{e+}11$; otherwise, it is extremely small. Second, these coefficients seem to be directly taken from the global configuration GO7 without adjustment, which is not appropriate. However, this might only marginally affect the simulations.**

We thank the reviewer for pointing out the typo. It is supposed to be $-1.08\text{e+}10 \text{ m}^4/\text{s}$. We have corrected it. The viscosity coefficient in GO7 is $-1.5\text{e+}11 \text{ m}^4/\text{s}$. We changed the coefficient based on the grid size and timestep of the AME025 configuration. For a bilaplacian diffusion scheme, the minimum stable requirement for the mixing coefficient (A^h) should meet

$$A^h < \frac{e^4}{64\Delta t}, \quad (1)$$

Where e is the smallest size of horizontal grids, Δt is the timestep (Gurvan et al., 2017). In the AME025 configuration, the smallest grid size is about 7 km and the timesteps are 720 s, which gives the mixing coefficient a limit of $-5.02\text{e+}10 \text{ m}^4/\text{s}$. However, considering the very strong forcing of UKESM1.0 under SSP5-8.5, we would like to remove the potential risk of instability. The instability can be solved by reducing the timesteps or reducing the mixing coefficient. Reducing the timesteps is computationally costly, so we decided to reduce the coefficient to $-1.08\text{e+}10 \text{ m}^4/\text{s}$ as in Bull et al. (2021).

- **The authors prescribe a constant tidal velocity of 5 cm/s. Where does this value come from? I suspect its impact on the melt is significant.**

Yes, the tidal velocity has a significant impact on the melt.

We did a sensitivity experiment to choose the value of 5 cm/s. We modelled the melt rate in 2006 with the tidal velocity increased from 0 cm/s to 12 cm/s in increments of 2.5 cm/s (Figure 1a). The melt rate is increased with the increase in tidal velocity. The modelled melt rate with tidal velocity of 2.5, 5 and 7.5 cm/s falls within the observational range (Figure 1b). The three runs, to a different extent, overestimate the melt rate north of the grounding line and underestimate the melt rate at the grounding line. This is the limitation of our configuration. The relatively coarse resolution cannot represent the steep meridional distributions of basal melting.

In Rosevear et al. (2022), the estimated magnitude of typical tidal current under Amery Ice Shelf from four tidal constituents is 9.8 cm/s. The tidal velocity of 9.8 cm/s produced an annual-mean Amery melt rate of 1.1 m/yr in 2006, which is slightly higher than the observational studies. In addition, the melt rate with 9.8 cm/s tidal velocity failed to represent refreezing.

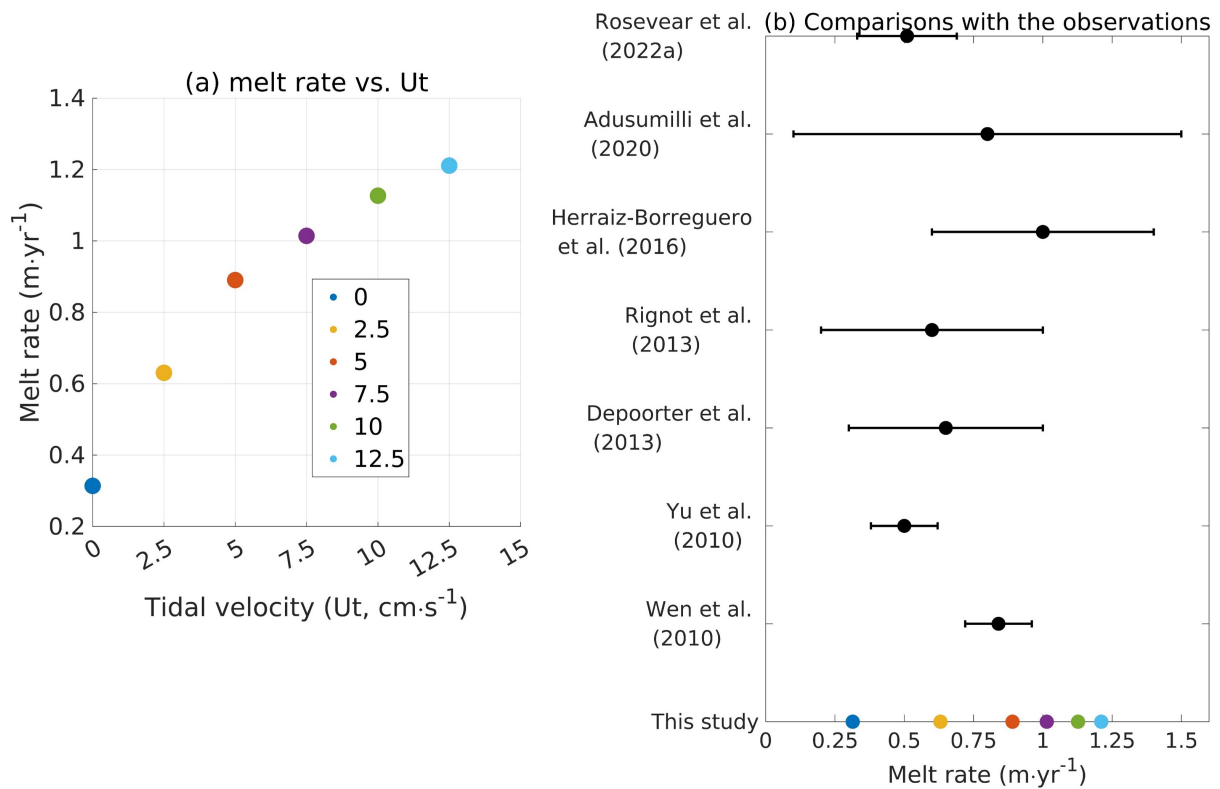


Figure 1. (a) Melt rate against tidal velocity (U_t). (b) Comparison of the melt rate with different U_t with the observational studies.

Given the tidal velocity of 2.5 cm/s is much smaller than the observed tidal velocity, and 7.5 cm/s might largely overestimate the future melt rate given a very high climate sensitivity of UKESM1.0, we decided to use 5 cm/s.

We have mentioned the experiments in the manuscript and provided the results of the sensitivity experiment in a supplemental document:

"A series of sensitivity experiments have been conducted to determine the value of the prescribed tidal velocity (Figure S1)."

– Is there any freshwater or salt restoring?

There is no freshwater or salt restoring in the regional domain. But the GO6 configuration, the ocean core of UKESM1.0, has restoring in the global domain. We have added this in the revised manuscript:

"There is no freshwater or salt restoring in the regional domain. However, the GO7 configuration—the ocean core of UKESM1—has restoring in the global domain (Storkey et al., 2018)."

– Lines 117 to 124: I do not understand the discussion about imposed melt in UKESM while AME025 has open interactive cavities.

Yes, the AME025 has open interactive cavities. However, we do not include freshwater fluxes of surface runoff and iceberg calving in our regional simulations. However, we do not include freshwater fluxes from surface runoff and iceberg calving in our regional simulations. While these are important components of surface freshwater fluxes, they may have a smaller impact on the Amery sector compared to ice shelf basal melting by the end of the 21st century (Siahaan et al., 2022; Coulon et al., 2024). We do not specifically employ these surface freshwater fluxes, but they implicitly enter our domain through the lateral boundaries, as they are prescribed in the global UKESM1.0 model.

We have deleted the discussion about imposed melt in UKESM and rephrased the sentence:

"...No external freshwater fluxes of surface runoff and iceberg calving are used in the AME025 configuration. These freshwater fluxes implicitly enter our domain through the lateral ocean boundary conditions, as they are prescribed in the UKESM outputs (Sellar et al., 2020)."

– The spin-up phase involves repeatedly simulating the year 1976, but how is the transition between December and January handled? Is there not a discontinuity?

There is no transition handled. We note the limitation of this spin-up method. However, this method is efficient in reducing the model drift. We have added additional information about our spin-up method in Section 2.3:

"...No smooth transition from December of a spin-up year to January of the next spin-up year is applied. There is physical discontinuity of the forcing between two spin-up years. Then a subsequent run is restarted from the last spin-up run with the continuous forcing from 1976 to 2100. The outputs of the subsequent run are used in the analysis."

2) Section 3.1.1

This section evaluates the melting rate by 2100. The authors focus on two periods: 2015-2055 and 2075-2100. However, it would be beneficial to explore the slow adjustment after the initial increase. For instance, there are fluctuations in the melt rate, especially in the 2090s for SSP1-2.6, that the authors did not explain. In my view, the most interesting aspect is the 10-year transition from a cold to a warm ice shelf around 2060-2070.

We briefly discussed the cause of the fluctuations in the melt rate in the 2090s under SSP1-2.6 scenario at Lines 294-297 in the preprint. This is due to a decrease in temperature on the continental shelf in the 2080s in the upstream/eastern ocean boundary conditions. All related processes in the 2080s-2090s under SSP1, such as a cooling on the continental shelf in the model domain, a recovery of PBG and a rebound in the melt rate, are ultimately controlled by the temperature changes in the upstream ocean boundary conditions.

We have done a double-check by running the projections with the r2 ensemble of UKESM1. We chose the r2 ensemble because the ocean temperature under SSP1-2.6 for the r2 ensemble has a very different variability from that for the r1. The increase in ocean temperature in the Amery sector in the r2 ensemble happens 20 years later than the r1 ensemble, and it is followed by a larger decrease after the temperature peak. The changes in the melt rate in the r2 experiment coincide with the temperature variability in the r2 ensemble. We have supplied this discussion and the results from the r2 experiment in a supplemental document.

Thanks for the suggestion on more investigations of the cold-to-warm transition. We have included more time-mean subpanels in the figures rather than two "before and after PBG reversal" subpanels in the revised manuscript. In addition, we have included more analyses on the time-evolution of sea ice, surface stress and circulations inside and outside the cavity to investigate the cold-to-warm transition.

Additionally, there is no discussion of estimates from other studies, which, I believe, found much weaker sensitivity. For instance, Kusahara et al. (2023) and Naughten et al. (2018) reported increases in melt by factors of 3 and 2, respectively, compared to a factor of 10 in this study.

We are not surprised by the much higher sensitivity in our projections, which is inherited from the global UKESM1.0-LL model. The forcing is MIROC-ESM in Kusahara et al. (2023) and ACCESS-1.0 and the multimodel-mean in Naughten et al. (2018). Meehl et al. (2020) have quantified the climate sensitivity of these models. For example, equilibrium climate sensitivity (ECS), which describes how much warming a model can produce with a doubling of CO₂ concentrations from its preindustrial level, is 3.8 K, 4.7 K and 5.3 K for MIROC-ESM, ACCESS-1.0 and UKESM1.0-LL, respectively. Another metric of model sensitivity – transient climate response (TCR) represents the response to increasing CO₂ – is 1.9 K, 2.2 K and 2.8 K for MIROC-ESM, ACCESS-1.0 and UKESM1.0-LL (Meehl et al., 2020). The higher climate sensitivity might result in a warmer ocean temperature (Sellar et al., 2020). This suggests our regional model, forced by the climate-sensitive UKESM1.0-LL, has a stronger response to climate warming.

Siahaan et al. (2022) used a special version of UKESM1.0-LL which has open ice shelf cavities and is interactively coupled with an icesheet model (UKESM1.0-ice). In their projections, the ice shelf basal melt rate increases by a factor of 10. This is not entirely comparable with our regional projections, but it to some extent implies that the high sensitivity of the melt rate comes from the physical core of UKESM1.0-LL.

We have added the discussion about the higher sensitivity of melt rate in our simulations in Section 3.1.1:

"The melt rates in our simulations show higher sensitivity compared with other studies, for example, Naughten et al. (2018); Kusahara et al. (2023). This is inherited from the global UKESM1.0-LL model. The model forcing in Naughten et al. (2018) and Kusahara et al. (2023) is taken from the ACCESS-1.0 and MIROC-ESM, respectively. Meehl et al. (2020) and Forster et al. (2020) have quantified the climate sensitivity of these models and suggest UKESM1.0-LL has higher climate sensitivity than the other two models, which might result in a warmer ocean temperature (Sellar et al., 2020). This suggests our regional model, forced by the outputs from the climate-sensitive UKESM1.0-LL, has a stronger response to climate warming. "

3) Section 3.1.2

This section evaluates the concomitant warming on the continental shelf. However, several issues and approximations need to be addressed.

The authors average temperatures in the 300-800 m depth range over the continental shelf without further explanation. It is necessary to specify the exact region considered and justify the choice of these depths.

We thank the reviewer for pointing out the missing information. We chose the 300-800m depth range as mCDW occupies the depth range on the continental shelf in our simulation (Figure 1-4). We aimed to show the mCDW changes on the continental shelf. We have added the reason why we chose the depth range in the revised manuscript:

"mCDW primarily occupies depths below 300 m in our simulations (Figure 5), so we choose a depth range of 300 m and below for the shelf sea properties to capture changes in mCDW."

Other modifications we have made is to replace “the depth range of 300-800m” with “below 300m”. The deepest bathymetry on the continental shelf is about 800. The more accurate statement is “mCDW occupies from about 300 m to the seafloor” rather than “between 300 and 800 m”. In addition, we have modified the definition of "the continental shelf". In the previous analyses, we used regions south of -67°N to represent the shelf sea. However, it is not accurate, and we have used bathymetry < 1000 m (The red arrow of the schematic ASC in Figure 1) to give a more precise geographical definition of the continental shelf in the revised manuscript. We have added a sentence to describe the

geographical definition:

"For the analyses below, the continental shelf and the shelf sea are defined by the areas between ASC and the AmIS calving front/coastal line (Figure 1). The cavity is defined as the ocean beneath the AmIS draft."

This averaging obscures the identification of the specific water masses that are warming. It would be more effective to examine sections extending from the grounding line to the shelf break and beyond. This approach would allow for a discussion on the warming or displacement of interfaces and the emergence of new water masses.

We thank the reviewer for the suggestion on analysing water masses. We removed the figure of depth-averaged temperature (Figure 4 in preprint). Instead, we present a multi-panel figure to show different time-mean of temperature transect from the grounding line to the north of the shelf break under the SSP5-8.5 scenario as suggested (Figure 4 in the revised manuscript). In addition, we present a multi-panel T-S diagram to show the water mass changes on the continental shelf and within the ice shelf cavity (Figure 5 in the revised manuscript). In addition, we included the separated timeseries of temperature/salinity on the shelf sea and inside the cavity as suggested by Reviewer 1 (Figure 6 in the revised manuscript)). Section 3.1.2 has been rewritten as new analyses have been added. To save space we did not copy the new paragraphs here, but we attached the revised manuscript at the end. The results from the SSP1-2.6 simulation are also included in the supplementary document.

The authors state that the temperature increase occurs several decades earlier than the increase in melt. However, based on Figure 4, the deviation from the steady state begins around 2040, approximately 15 years before the increase in melting, not decades. Additionally, the transition from the cold state to the warm state occurs clearly between 2060 and 2070 for both melting and warming, so the claim of a significant delay is questionable.

Thanks for the reviewer to correct our language. We will clarify the specific description of the time scale and avoid using “decades”.

The revised figures as suggested by the reviewers, which include the temperature transects from the grounding line to the shelf break and beyond (Figure 1, 2), water mass plots (Figure 3, 4) and the separated timeseries of temperature within and outside the cavity (Figure 5, 6), suggest the warming on the continental shelf starts in the 2030s even 20s, not 40s. The warming in the cavity begins in the 2050s, and the transition of the cavity starts in the 2060s. There may be some statistical argument for time differences between the warming outside of the cavity and the increase in melting, but we insist there is a delay.

4) Section 3.2

The goal of this section is to explain the changes in the melt rate. To do this, it is necessary to accurately depict the changes in water masses inside and ahead of the cavity, and their link to changes in ocean circulation. The analysis presented here is insufficient to address the question.

We are grateful for the reviewer's suggestions. We agree it is necessary to address the changes in water masses and their link to the changes in ocean circulation, and our analyses are insufficient. We have added additional analyses to address this question. We have made major corrections to Section 3.2 and rewritten this section. To save the reviewer's time, we summarised the changes in the analyses and the key finding here:

1. The changes in the water masses have been presented in Section 3.1.2. It shows that HSSW is decreased and the increased mCDW upwells on the continental shelf from the 2040s. The dominant water mass inside the cavity changes from HSSW to mCDW in the 2060s.
2. The spatial distributions of barotropic stream function inside and outside the cavity (Figure 7) and time evolution of the strength of the cavity circulation and PBG (Figure 8) have been presented in Section 3.2.1. This shows the reversal of PBG in the 2060s and the strength of PBG is increased afterwards. The strength of the cavity circulation is decreased in the 2040s as HSSW becomes inefficient to drive the cavity circulation, while the circulation strength is increased again in the 2060s as mCDW flushes the cavity due to the established clockwise PBG.
3. The heat budget of the cavity (Figure 9) has been added in Section 3.2.1. It shows the heat fluxes into the cavity start to increase in the 2040s. However, the massive heat, which changes the cavity regime, arrives in the 2060s as mCDW flushes the cavity.

The missing analyses are:

- Temperature transects.

We have provided an analysis of temperature transects (Figure 10a). Another modification is that the analysing region has been changed. The analysing region consisted of two boxes, one over PBG and the other over the western region of PBG. However, we showed the ocean circulation with barotropic stream function rather than meridional velocity as suggested by the reviewer. The two boxes will divide PBG, and this is not ideal to present an integral picture of the time-evolution of PBG. We changed the analysing region to a zonal transect of PBG with extensions to the western and eastern coastal lines (Figure 7a). As a consequence, we replaced the old Figure 7, which shows multi-panel plots of time-against-depth salinity in the two boxes, salinity differences between the two boxes and V velocity of PBG, with new multi-panel longitude-against-time plots (Figure 10). This new Figure 10 shows the temperature transects as requested by the reviewer, salinity, SSH gradients and barotropic stream function of PBG. The reflections of Figure 10 in the revised manuscript are not essentially different from the old

Figure 7 in the preprint. It also illustrates why and how PBG evolves. We have rewritten the analysis of the PBG transect from Lines 373-395 in the new version of the manuscript with tracked changes.

Another related modification is that we removed the old Figure 8 and the analysis of dynamic height anomaly from Lines 265 to 285 in the preprint. Dynamic height anomaly represents the pressure, which illustrates the reversal of PBG is due to new pressure gradients. However, this overlaps with the SSH gradients as shown in Figure 10c. Therefore, we removed the analysis of dynamic height anomaly to reduce the repetitive information. Additionally, the dynamic height anomalies are calculated offline, while the SSH variable is taken from model outputs, which is more convincing.

– **Analysis of mCDW transport inside the cavity with a proper closed budget.**

We are unsure of the specific analysis of mCDW inside the cavity that the reviewer is expecting. For example, is it a budget as in Jourdain et al. (2017, eq. 6):

$$\partial_t T = -u\partial_x T - v\partial_y T - w\partial_z T + D_l(T) + D_z(T) + l(T, z), \quad (2)$$

which describes temperature changes due to advection, diffusion and surface fluxes. Or is it as in Jourdain et al. (2017, eq. D1), which describes heat flux into the cavity is used to melt ice, and warm/cool seawater in the cavity and the remaining heat flows out the cavity. Or is it an analysis of some processes else? We supplied an analysis as Jourdain et al. (2017, eq. D1) for two reasons. Firstly, the heat budget in the cavity is more illustrative of the regime change due to mCDW in the 2060s. We are aware this method cannot separate the heat transport from different water masses entering the cavity, which may not thoroughly address the reviewer's question. However, we have demonstrated the dominant water mass controlling the cavity changes from HSSW to mCDW in the new version of Section 3.1.2. It provides additional supporting information to the heat budget analysis. Second, we did not save model outputs shorter than monthly, and we did not output the diffusion variables. The offline calculations of Jourdain et al. (2017, eq. 6), especially the diffusion terms, are not ideal. We have to re-run all experiments to obtain those data, which needs an additional two or three months. To shorten the reviewing process, we present the analysis of the heat budget as Jourdain et al. (2017, eq. D1), and we appreciate the understanding from the reviewer.

Following Jourdain et al. (2017, eq. D1), neglecting diffusion in ice and the interior ocean, the heat flux entering the ice cavity (H_{in}) is simplified into three components: the heat used to melt the ice shelf (H_{lat}), the heat that

warms or cools the seawater within the cavity ($H_{\text{HC_VAR}}$), and the remaining heat flux that exits the cavity (H_{out}):

$$H_{\text{in}} = H_{\text{out}} + H_{\text{lat}} + H_{\text{HC_VAR}} \quad (3)$$

$$H_{\text{in}} = \rho_{\text{ref}} C_p \iint_{r \in \text{front}, u > 0} u(r, z) (T(r, z) - T_f^0) dr dz \quad (4)$$

$$H_{\text{out}} = \rho_{\text{ref}} C_p \iint_{r \in \text{front}, u < 0} u(r, z) (T(r, z) - T_f^0) dr dz \quad (5)$$

$$H_{\text{lat}} = L_f \iint_{x, y \in \text{draft}} \rho_i m dx dy \quad (6)$$

$$H_{\text{HC_VAR}} = \rho_{\text{ref}} C_p \iint_{x, y, z \in \text{cavity}} \frac{dT(x, y, z)}{dt} dx dy dz \quad (7)$$

ρ_{ref} is the reference seawater density of 1026 kg m^{-3} , C_p is the specific heat capacity of 3991.87 J K^{-1} . $u(r, z)$ is the velocity perpendicular to the ice shelf front. The positive means velocity into the cavity. $T(r, z)$ is the temperature along the ice shelf front. T_f^0 is surface freezing point. L_f equal to 334 is the latent heat of fusion of ice. $\rho_i m$ is the melt rate in unit of $\text{kg m}^{-2} \text{ s}^{-1}$. The latent heat of ice shelf (H_{lat}) is obtained from the model outputs. $dT(x, y, z)$ is temperature change on the same grid cell. dt is one month. The variables used in the calculations are monthly mean model outputs. The timeseries of the heat budget components show a sudden increase in H_{in} and H_{lat} in the 2060s, suggesting the sudden flush of mCDW causes the increased melting as expected (Figure 9).

- Use of daily outputs instead of monthly, since $(v \cdot T)^{\text{1mo}} \neq v^{\text{1mo}} \cdot T^{\text{1mo}}$ especially given the strong variability in this area.

Yes, we agree the use of daily data is more helpful in investigating the variability of PBG. However, we saved monthly variables only to save data storage space. We will be short of space if we save daily outputs for the runs over 200 years in total. We have to re-run simulations to obtain the daily outputs. Additionally, the university's HPC was upgraded this summer, and the modules, particularly the compiler used in this study, are no longer available. Guarino et al. (2020) suggest the impact of HPC machines on the reproducibility of the climate simulations. Although this has a low likelihood of changing the main conclusions, we have to re-run monthly outputs as well to ensure that the new simulations do not have significant differences from the ones we used in the preprint. However, this will take two to three months to finish all simulations, and it slows down the reviewing processes.

An alternative solution is to remove the analysis of meridional heat transport for the PBG and PBECC. Instead, we present the analysis of the heat budget in the cavity as we replied previously. The analysis of meridional heat transport of PBG and PBECC cannot quantify how much heat outside the cavity is transported into the cavity. While it shows increased onshore heat during the 2060s due to the reversal of the PBG, this is somewhat redun-

dant as we have already demonstrated the reversal of the PBG via the reversed barotropic stream function, as suggested by the reviewer. In addition, the heat budget in the cavity has already illustrated the increased heat into the cavity in the 2060s. Therefore, we have decided to remove the analysis of PBG and PBECC heat transport, as it offers less valuable insight compared to the new analysis. The removal of this analysis will not weaken the integrity of this study.

As for the analysis of the heat budget in the cavity, we note the calculations of heat fluxes in the cavity are based on monthly outputs and the issue of $(v \cdot T)^{1mo} \neq v^{1mo} \cdot T^{1mo}$ has not been addressed. However, the object of this study is the AmIS melt rate which exhibits stronger inter-annual to decadal variability rather than variability shorter than one month. The heat budget in the cavity is properly closed (Figure 9), which suggests that $v^{1mo} \cdot T^{1mo}$ may adequately explain the changes in the monthly-mean melt rate. The use of monthly mean outputs will not change our main conclusions. This also quantifies that about 50% of the heat into the cavity is used to melt ice. In addition, Liu et al. (2017) presented a similar analysis of meridional heat transport across a zonal section (73-78°E, 67°S, which is close to the PBG section in our study) with an eddy-resolving model and daily outputs. They decomposed the total meridional heat transport (Q) into the mean flow ($Q_m(\overline{v \cdot T^{1mo}})$) and the eddies ($Q - Q_m(\overline{v \cdot T^{1mo}})$) components and suggested that the mean flow components contribute to about 50% of the total meridional heat transport. However, our model resolution is lower, which suggests the contributions from the high-frequency variability may not be as large as those in Liu et al. (2017). Moreover, our focus is the depths below 300 m, which may exhibit smaller variability. It might have a minor impact on the main findings of the study. We could re-run the simulations if the reviewer is still unsatisfied with the monthly data.

- Show the reversal of PGB with a streamfunction, rather than just two sections of meridional velocity.

We have shown the reversal of PGB with barotropic stream function (Figure 7; Figure 10d), and we have rewritten Section 3.2.1 presenting the reversal of PBG. Additionally, we have provided time-against-depth plots of V velocity of PBG under two scenarios in Appendix A1.

The causes of PGB reversal are only conjectures in this analysis. For instance, the authors state that PGB is in geostrophic equilibrium due to the salinity difference between two boxes, S1 and S2. However, this is not a rigorous calculation. It merely suggests that it is possible for the gyre to be in geostrophic balance. To make a definitive statement, one would need to consider all the tendency terms of the forces driving the gyre. This approach is not rigorous.

We thank the reviewer for pointing out the inappropriate use of language. We have rephrased the statement of “geostrophic equilibrium/balance” throughout the manuscript, and modified it as:

"The reversal of PBG is a consequence of the reversal of horizontal salinity/density differences between PBG and its western regions."

In addition, we have added more analysis of the forcing of PBG (Figure 11). This is to investigate another possible driver for the PBG reversal, for example, surface freshwater forcing from sea ice, Ekman transport, and Ekman pumping, although this does not change our main conclusion that the ice shelf freshwater reverses the horizontal salinity gradients. Closing the vorticity budget of the gyre is out of the research scope. Following the previous calculation, we presented a scale analysis of Lines 396-438 in the new version of the manuscript with tracked changes.

From lines 290 to 298, the authors claim that the upstream systems influence the ice shelf downstream, but there is no proof provided, only hypotheses.

We will remove the statement about “the influence of upstream systems”.

5) Conclusion

- From lines 314 to 323, I am unclear about the authors’ intended message. Galton-Fenzi (2009) mentioned that the PBG depends on the relative strength of McKenzie and Barrier polynyas, and it is already known that PBG drives mCDW into the cavity. What new information is being presented here?

We intended to discuss differences of PBG reversal between our study and Galton-Fenzi (2009). However, the discussion is inadequate and not informative. We thank the reviewer for pointing it out. We have added a new analysis of sea ice in the Barrier polynya (Figure 12) and have moved the discussion from Section 4 to Section 3.2.2. As this analysis provides some supporting information for the mechanism of the PBG reversal, this arrangement will make the discussion on the processes more coherent. We summarised the analysis here to save the reviewer’s time:

Our model overestimates the strength of the Barrier polynya during the historical period, which suppresses the formation of the clockwise PBG in summer. This is why the seasonality of PBG in our simulations is weaker than in Galton-Fenzi (2009). The reduction in sea ice in the 2040s facilitates the pathway of the clockwise PBG. Although sea ice reduction is not the direct driver of the PBG reversal, it establishes the preconditions for the reversal.

A detailed discussion has been presented from Lines 445-479 in the new version of the manuscript with tracked changes.

- **Lines 324-331 are not informative.**

We have rephrased.

"A similar mechanism for the regime change is found in Filchner-Ronne Ice Shelf (Hellmer et al., 2012, 2017; Naughten et al., 2021; Siahhaan et al., 2022) and Ross Ice Shelf (Siahhaan et al., 2022). A redirection of coastal current is driven by a reversed density gradient across the Filchner Trough (Hellmer et al., 2012, 2017; Naughten et al., 2021; Siahhaan et al., 2022)/ the little America Basin in the Ross Sea (Siahhaan et al., 2022), which facilitates the penetration of mCDW into the ice shelf cavity and causes the regime change. The similar processes found in the sectors for the three large cold ice shelves emphasise the importance of buoyancy changes on the shelf sea in a warming future. This suggests the necessity of long-term records for the shelf sea salinity and sea ice in order to obtain an early warning of the regime change in the three large cold cavities."

- **Lines 335-349 do not offer conclusions but rather present some reflections, without providing a perspective on future work.**

We have added a perspective on future work

- *"... A series of freshwater perturbation experiments across various climate scenarios would help address the above questions."*
- *"... Coupling our standalone ocean configuration with an ice sheet model will be a future step to address the question of the AmIS stability."*

- **Lines 350-380 appropriately discuss the limitations of the work, but the authors overlook the issue of frazil ice refreezing at the ice shelf base, which the model does not account for.**

We have added a discussion about the frazil ice as suggested.

"... Moreover, we do not include the frazil processes in this configuration. Galton-Fenzi et al. (2012) suggests the inclusion of frazil in the AmIS simulation decreases the melt rates and increases the arrection rates. No frazil ice in our simulation may result in the overestimation."

Figures

The quality of the figures is inadequate. For example, the color bar in Figure 1 does not provide a clear visualization of the bathymetry in the Amery cavity. Additionally, the labels in Figure 2 are so small that they are barely legible, among other issues.

We have updated all figures as suggested.

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