



# Validating a microphysical prognostic stratospheric aerosol implementation in E3SMv2 using the Mount Pinatubo eruption

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**Abstract.** This paper describes the addition of a stratospheric prognostic aerosol (SPA) capability – developed with the goal of accurately simulating aerosol formation following explosive volcanic eruptions – in the Department of Energy (DOE) Earth Energy Exascale Model, version 2 (E3SMv2). The implementation includes changes to the 4-mode Modal Aerosol Module microphysics in the stratosphere to allow for larger particle growth and more accurate stratospheric aerosol lifetime following the Mt. Pinatubo eruption. E3SMv2-SPA reasonably reproduces stratospheric aerosol lifetime, burden, and aerosol optical depth when compared to remote sensing observations and the interactive chemistry-climate model, CESM2-WACCM. Global stratospheric aerosol size distributions identify the nucleation and growth of sulfate aerosol from volcanically injected SO<sub>2</sub> from both major and minor volcanic eruptions from 1991 to 1993. Modeled aerosol effective radius is consistently lower than satellite and in-situ measurements (max differences of ~30%). Comparisons with in-situ size distribution samples indicate that this underestimation is due to both E3SMv2-SPA and CESM2-WACCM simulating too small of accumulation and coarse mode aerosol 6-18 months post-eruption, with E3SMv2-SPA simulating ~50% the coarse mode geometric mean diameters of observations 11 months post-eruption. Effective radii from the models and observations are used to calculate offline scattering and absorption efficiencies to explore the implications of smaller simulated aerosol size on the Mt. Pinatubo climate impacts. Scattering efficiencies at wavelengths of peak solar irradiance (~0.5 μm) are 10-80% higher for daily samples in models relative to observations through 1993, suggesting higher diffuse radiation at the surface and a larger cooling effect in the models. Absorption efficiencies at the peak wavelengths of outgoing terrestrial radiation (~10 μm) are 15-40% lower for daily samples in models relative to observations suggesting an underestimation in stratospheric heating in the models. The similar performance of CESM2-WACCM and E3SMv2-SPA makes E3SMv2-SPA a viable alternative to simulating climate impacts from stratospheric sulfate aerosols.

## 1 Introduction

Explosive volcanic eruptions are a significant source of aerosol forcing given their propensity to inject gas and particulate matter into the stratosphere. These gases form long-lived aerosol that can be spread around the globe (Robock, 2000). The impacts of the stratospheric aerosol loading from these eruptions are wide ranging, as exemplified in observational and modeling findings following the Mt. Pinatubo eruption in 1991: ozone depletion (Hofmann et al., 1994; Solomon et al., 1993; Portmann et al., 1996), surface temperature decreases (Parker et al., 1996; Soden et al., 2002), lower stratosphere temperature increases (Labitzke and McCormick, 1992), reduction in global precipitation (Gillett et al., 2004), lowering of global sea-level (Church et al., 2005), increases in cirrus cloud cover (Liu and Penner, 2002; Wylie et al., 1994), as well as increased diffusivity of incoming radiation (Robock, 2000) with resultant impacts on net primary productivity (Gu et al., 2003; Proctor et al., 2018; Greenwald et al., 2006). The extent of these impacts is dependent upon characteristics of the eruption (e.g., magnitude (Marshall et al., 2019)) and climate state (Zanchettin et al., 2022). The foundation of these physical impacts are the stratospheric aerosol microphysical properties and chemical interactions that occur after an explosive volcanic eruption. Being able to accurately simulate these aerosol microphysical and chemical reactions in Earth system models is important for



65 improving the fidelity of simulations. It also enables climate attribution work that clarifies the role of climatic state versus characteristics of the eruption on downstream impacts. As a first step towards this goal, this paper presents a validation of a prognostic volcanic aerosol implementation within the Department of Energy (DOE) Earth Energy Exascale Model version 2 (E3SMv2) (Golaz et al., 2022; Wang et al., 2020), in comparison with version 2 of the Community Earth System Model (CESM2) (Danabasoglu et al., 2020) and observational data from the Mt. Pinatubo eruption.

70 When simulating large-magnitude explosive volcanic eruptions, most climate models use the Global Space-based Stratospheric Aerosol Climatology (GloSSAC) to prescribe aerosol properties from a compilation of satellite, airborne, and ground based observations (Kovilakam et al., 2020; Thomason et al., 2018). While the GLoSSAC dataset provides an accessible approach for incorporating volcanic forcing and validating model performance, prescribed aerosol products have limitations. Limited dataset availability and/or spatial coverage necessitate data interpolation within the forcing dataset, which  
75 may not accurately represent the volcanic forcing in some regions (Kovilakam et al., 2020). Additionally, prescribed aerosols will not respond to the dynamic state in free-running coupled climate simulations as the database has been generated from the observed climatic conditions. This artificially constrains the volcanic forcing across ensembles of simulations, and creates a disconnect between volcanic forcing and the actual atmospheric transport patterns, limiting the usability of these simulations for detection and attribution of an evolving impact to the volcanic source. Another limitation is the lack of aerosol  
80 microphysical representation and evolution from a volcanic eruption which ignores aerosol indirect effects on clouds and doesn't allow for model feedbacks on aerosol size and lifetime.

Prognostically modeling the formation and evolution of sulfate aerosol from sulfur dioxide (SO<sub>2</sub>) injected into the stratosphere is an alternative, more complete approach for simulating volcanic eruptions, with a variety of methods for representing sulfate aerosol mass, size, and number. The simplest approach is to use a bulk aerosol treatment, which prescribes an aerosol size distribution to a predicted aerosol species mass. This was applied to the earliest multi-year simulations of  
85 Pinatubo run with the Hamburg climate model ECHAM4 (Timmreck et al., 1999a, b). Most recently, a bulk representation was used in the GFDL Earth System Model (ESM4.1; Dunne et al. (2020)) to show that changes in SO<sub>2</sub> injection height can have the largest impacts on volcanic mass burden and climate forcing (Gao et al., 2023). The most accurate approach to simulating aerosol properties is the sectional (or bin) approach, but this can be computationally limiting depending on the  
90 number of aerosol size bins used. English et al. (2013) coupled the sectional Community Aerosol and Radiation Model for Atmospheres (CARMA; Toon et al. (1988)) with version 3 of the Whole Atmosphere Community Climate Model (WACCM3; Garcia et al. (2007)), showing the value of a sectional model in simulating the large variation in aerosol mode size and width that occurs after Pinatubo and larger magnitude eruptions. The modal aerosol approach – representing aerosol size distributions by multiple, evolving lognormal functions – strikes a balance between bulk simplicity and sectional cost. This is used in a  
95 stratospheric, prognostic aerosol treatment with interactive ozone chemistry (Mills et al., 2016), developed in the Community Earth System Model (CESM, version 1) (Hurrell et al., 2013) using the version 4 of WACCM (Marsh et al., 2013). This model design has been used to identify the impacts of volcanic eruptions on stratospheric ozone (Ivy et al., 2017; Solomon et al.,



2016), the importance of interactive chemistry on the representation of sulfate formation and distribution following Pinatubo (Mills et al., 2017), as well as effective strategies for geoengineering (Kravitz et al., 2017).

100 Accurate simulation of the stratospheric sulfate size distributions is important for simulating the climate impacts of volcanic eruptions. This is because both the scattering of incoming shortwave energy (i.e., surface cooling) and absorption of outgoing longwave energy (i.e., stratospheric warming) are related to aerosol size and number through Mie theory. Size and number in turn depend on the aerosol modeling approach (e.g., bulk, modal, sectional). In E3SM, tropospheric aerosol is simulated with the 4-mode Modal Aerosol Module (MAM4; Liu et al. (2012, 2016)). In the case of sulfate, MAM4 represents  
105 the oxidation of SO<sub>2</sub> to form sulfate aerosol, the further growth of these aerosol through condensation and coagulation into larger aerosol size modes, and the sedimentation of these aerosol through wet and dry deposition.

Here we present a new prognostic aerosol capability implemented within E3SMv2 to enable simulation of the evolution of volcanic stratospheric aerosols and their properties. Similar to Mills et al. (2016), we modify the microphysical treatment of stratospheric aerosol in MAM4. The original validation of this implementation (Mills et al., 2016, 2017) focused  
110 on representation of aerosol lifetime, aerosol chemistry, radiative impacts, and dynamical response to the Pinatubo eruption. Here we pay more attention to the aerosol microphysics and how model representations are related to global and regional climate impacts. We use both observations from the 1991 eruption of Mt. Pinatubo as well as CESM2-WACCM to validate the implementation and demonstrate that we can reasonably simulate the lifetime, burden, and aerosol optical depth (AOD) of stratospheric sulfate without the computationally expensive, whole-atmosphere, comprehensive chemistry in WACCM.  
115 Additionally, we highlight the importance of simulated aerosol size and number on shortwave and longwave radiative impacts through comparisons of aerosol effective radius and number distributions to in situ observations. Utilizing these effective radii in single particle Mie scattering calculations, we explore how variations in modeled aerosol microphysics affect the volcanic aerosol impacts on diffuse/direct radiation at the surface as well as longwave absorption and heating rates in the stratosphere.

## 120 2. Models and Simulations

In this work we modify E3SMv2 to include a stratospheric prognostic aerosol capability. We test stratospheric microphysical and chemical implementations in the context of the Mt. Pinatubo eruption and compare to version 2 of the Community Earth System Model CESM2 with the Whole Atmosphere Community Climate Model version 6 (WACCM6; Gettelman et al., 2019).

### 125 2.1 E3SMv2

Originally a branch from CESM1, E3SMv1 diverged with a focus on computational efficiency, scalability, vertical and horizontal resolution, aerosol and cloud parameterizations, as well as more physically based biogeochemistry, river, and cryosphere models (Leung et al., 2020). Simulations in this study are run with E3SMv2 in which clouds are parameterized with an improved version of the Cloud Layers Unified by Binormals (CLUBB) scheme (Larson, 2017), cloud microphysics



130 are simulated by a two-moment bulk microphysics parameterization (MG2; Gettelman and Morrison, 2015), and mixed phase  
ice nucleation depends on aerosol type and concentration as well as temperature (Wang et al., 2014; Hoose et al., 2010).  
Aerosols are simulated with MAM4 (Liu et al., 2016). The explosive volcanic eruption treatment prescribes stratospheric  
volcanic light extinction from version 1 of the GloSSAC reanalysis dataset (Thomason et al., 2018). Atmospheric chemistry  
is represented with version 2 of the interactive stratospheric ozone (O<sub>3</sub>) model (O3v2; Tang et al., 2021). This model uses  
135 linearized stratospheric chemistry (Linoz v2; Hsu and Prather, 2009), which calculates net O<sub>3</sub> production as a function of  
temperature, local O<sub>3</sub> concentration, and overhead column O<sub>3</sub>.

### 2.1.1 Prognostic aerosol in E3SMv2-PA

In the E3SMv2-PA version of the model, prescribed volcanic extinction is removed and the sulfate aerosol precursor,  
SO<sub>2</sub>, is emitted in the stratosphere. The emitted SO<sub>2</sub> undergoes chemical reactions to form sulfate aerosol and condenses onto  
140 the surfaces of preexisting aerosols following the prognostic calculations of MAM4.

In E3SMv2-PA, default MAM4 size modes are employed. All aerosol species are represented by three size modes,  
shown here with 10<sup>th</sup>/90<sup>th</sup> percentile global, annual average number distribution dry diameter ranges from Liu et al. (2012):  
Aitken (0.015-0.053 μm), Accumulation (0.058-0.27 μm), and Coarse (0.80-3.65 μm). A fourth mode (i.e., primary carbon  
mode (0.039-0.13 μm)) represents freshly emitted black carbon and organic carbon from combustion, which then ages into the  
145 accumulation mode. Aerosol mass and number are used to define a modal geometric mean diameter ( $D_g$ ). This calculated  $D_g$ ,  
in conjunction with a fixed modal geometric standard deviation ( $s_g$ ), defines the modal number distribution. The modal  
distributions then evolve based on nucleation (aerosol formation), evaporation (aerosol size reduction), condensation and  
coagulation (aerosol size growth), and dry/wet deposition (aerosol removal). In E3SMv2-PA, the growth of accumulation to  
coarse mode aerosol is not included because the troposphere seldom has high enough aerosol mass concentrations to generate  
150 such large aerosols through condensation and coagulation. Thus, accumulation mode  $D_g$  is allowed to increase until they reach  
the upper modal threshold ( $D_{g,hi}$ ) whereby the model increases accumulation mode number to maintain  $D_{g,hi}$  until  $D_g$  begins  
decreasing.

### 2.1.2 Prognostic stratospheric aerosol in E3SMv2-SPA

Because MAM4 was designed to accurately represent tropospheric aerosol at their respective concentrations and  
155 emission fluxes, the sulfate formation from the massive stratospheric influx of SO<sub>2</sub> from the eruption of Pinatubo is not  
accurately represented in E3SMv2-PA. In addition to removing prescribed volcanic extinction – as for E3SMv2-PA - further  
modifications are made to MAM4 to create a prognostic stratospheric aerosol version of E3SMv2 (E3SMv2-SPA) that has  
improved stratospheric aerosol representation following the Pinatubo eruption. These improvements borrow heavily from  
changes made to version 1 of CESM (CESM1) with WACCM (Appendix B in Mills et al., 2016), which are present in the  
160 default MAM4 version in CESM2-WACCM6. The major modifications to MAM4 include (1) the transfer of aerosol mass and  
number from the accumulation to coarse mode to increase aerosol size and represent the rapid aerosol growth following the



Pinatubo eruption and (2) adjustment of the coarse mode and accumulation mode  $s_g$  and minimum/maximum geometric mean diameters to increase aerosol lifetime. Additional steps were taken to tune the model following the change in the accumulation and coarse mode size properties in (2) which included tuning of dust and seasalt emissions to account for the increased coarse mode lifetime in the model as well as recalculating modal optical properties in MAM4 to account for the changes in aerosol size limits and distribution widths.

Unlike typical tropospheric conditions, explosive volcanic eruptions into the stratosphere provide ample  $\text{SO}_2$  mass to drive sulfate aerosol into the coarse mode. To represent this rapid growth and overall larger aerosol diameters in the stratosphere an irreversible accumulation mode number and mass transfer into the coarse mode is added to E3SMv2-SPA. The model calculates the mass and number of particles in the tail of the distribution above a specified size cut off ( $D_{g,\text{cut}}$ ) of  $0.44 \mu\text{m}$ , transferring this overshooting number and volume into the coarse mode. The model prohibits transfer from the accumulation mode if  $D_g < 0.166 \mu\text{m}$  and allows total transfer of the gridcell mass and number when  $D_g > 0.47 \mu\text{m}$ . In CESM2-WACCM, this transfer is reversible in the stratosphere, with an aqueous sulfuric acid ( $\text{H}_2\text{SO}_4$ ) equilibrium pressure that depends on temperature and relative humidity.

To improve stratospheric aerosol lifetime in E3SMv2-SPA, the default coarse mode  $s_g$  is reduced from 2.0 to 1.2. The default accumulation mode  $s_g$  is also reduced, from 1.8 to 1.6, which has a small effect on aerosol lifetime. Additional changes to the aerosol modes allow for overlap between the coarse and accumulation modes and include increasing the accumulation mode  $D_{g,\text{hi}}$  from  $0.44 \mu\text{m}$  to  $0.48 \mu\text{m}$  and decreasing the coarse mode  $D_{g,\text{low}}$  from  $1.0$  to  $0.4 \mu\text{m}$ . Lastly, coarse mode  $D_{g,\text{hi}}$  is increased from  $4.0 \mu\text{m}$  to  $40 \mu\text{m}$ . See Supplementary Table S1 for a summary of these changes.

Most of the above changes have little effect on the tropospheric aerosols, except for changes to the coarse mode, which leads to longer lived coarse mode aerosol due to a reduction in removal rates. To account for this, emissions of dust and sea salt are tuned such that a simulation with perpetual present-day forcing obtains a global average AOD ( $0.1617$ ) and global average dust AOD ( $0.0281$ ) comparable to present day remote sensing observations of  $\sim 0.17$  (Lee and Chung, 2013) and  $0.28$ - $0.3$  (Ridley et al., 2016), respectively. The modal aerosol optical parameterizations are also affected by changes to the prescribed mode  $s_g$ , and the modal optical properties were recalculated with the above modifications using the Ghan and Zaveri (2007) offline code used to generate the original files for CESM/E3SM.

## 2.2 CESM2-WACCM6

The major CESM2 atmosphere model improvements from CESM1 are the inclusion of CLUBB, MG2 cloud microphysics, MAM4, and orographic wave drag parameterizations (Danabasoglu et al., 2020). These are the same in E3SMv2 (Golaz et al., 2022) with the exception of a stratospheric prognostic aerosol treatment in MAM4 in CESM2-WACCM6 (Mills et al., 2016, 2017) (see Section 2.1.2). WACCM6 includes updated atmospheric chemistry, aerosol microphysics, and gravity wave drag parameterizations from previous versions of WACCM. Atmospheric chemistry is treated comprehensively through the whole atmospheric column, representing key chemical species and reactions across the troposphere, stratosphere, mesosphere, and lower thermosphere (Gettelman et al., 2019). Within the stratosphere and mesosphere, WACCM6 explicitly



195 calculates the net production and transport of 97 different chemical species described by nearly 300 reactions (Mills et al., 2017). This comprehensive chemical treatment is a key difference between CESM2-WACCM6 and E3SMv2 (including E3SMv2-PA and E3SMv2-SPA), as E3SMv2 prescribes from observationally derived climatologies of OH and other relevant chemical species in sulfur and ozone chemistry (Hsu and Prather, 2009).

### 2.3 Simulations

200 E3SMv2 simulations are run from 1990-1993 (1989 discarded for aerosol spinup). The model uses the horizontal and vertical resolution described in (Golaz et al., 2022) where the dynamics run on a ~110 km horizontal grid, the physics run on a coarsened ~165 km horizontal grid, and both dynamics and physics use the same 72 layer vertical grid with a model top at approximately 0.1 hPa. Simulations have prescribed sea ice and sea surface temperatures (Taylor et al., 2000) and nudged U and V winds to 6-hourly MERRA2 reanalysis data. CESM2-WACCM6 simulations are run over the same time period on a  
205 0.95°x1.25° grid over 88 pressure levels (model top at  $\sim 4.5 \times 10^{-6}$  hPa). In addition to nudging model U and V winds to 6-hourly MERRA2 reanalysis data, these simulations also nudge model temperature to the reanalysis to constrain temperature-dependent stratospheric chemical reactions.

The SO<sub>2</sub> emissions for explosive volcanic eruptions are from VolcanEESMv3.11, a modified version of Neely III and Schmidt (2016). The VolcanEE3SMv3.11 dataset contains estimates of SO<sub>2</sub> from volcanic eruptions on a 1.9x2.5-degree  
210 latitude by longitude grid, with 1 km altitude spacing from the surface to 30 km. SO<sub>2</sub> emissions are provided in molecules cm<sup>-3</sup> s<sup>-1</sup>, and all eruptions occur over a six-hour period. The modifications to Neely III and Schmidt (2016) are described in Mills et al. (2016) and include a reduction in SO<sub>2</sub> emissions for eruptions over 15 Tg SO<sub>2</sub> by a factor of 0.55 to compensate for missing ice and ash removal processes. For Mt. Pinatubo, 18-19 Tg of SO<sub>2</sub> erupted in the atmosphere (Guo et al., 2004b) with only ~10 Tg remaining in the stratosphere for further microphysical and chemical evolution (Kremser et al., 2016). This 10 Tg  
215 emission is comparable to observed SO<sub>2</sub> 7-9 days after the eruption (Mills et al., 2016). The emission takes place between 18-20 km, at a single lat-lon grid cell (i.e., no spreading). For all simulations, the VolcanEESMv3.11 file was merged with the monthly CMIP6 SO<sub>2</sub> emissions for non-explosive volcanic sources and then remapped to 1x1 degree.

Table 1 provides a summary of some of the key model characteristics for the different sensitivity studies used in this work. These studies include: E3SMv2 with prescribed volcanic forcing and no emission of volcanic SO<sub>2</sub> (E3SMv-presc);  
220 E3SMv2 with the default MAM4 treatment (i.e., no stratospheric aerosol modifications) and emission of volcanic SO<sub>2</sub> (E3SMv2-PA); E3SMv2 with the prognostic stratospheric aerosol modifications and emission of volcanic SO<sub>2</sub> (E3SMv2-SPA); and CESM2-WACCM6 with emission of volcanic SO<sub>2</sub> (from herein referred to as CESM2-WACCM).

225 **Table 1: Model details for the simulations used within this study. All simulations are run for 5 years (1989-1993) with 1989 discarded for aerosol spinup.**





<b>Model Version</b>	<b>Horizontal, Vertical resolution</b>	<b>Nudging</b>	<b>Chemistry</b>	<b>Stratospheric Volcanic Aerosol</b>
E3SMv2-presc	ne30pg2 72 vertical levels (0.01 – 1000 hPa)	U + V winds to MERRA2 Reanalysis	Linear ozone chemistry	<ul style="list-style-type: none"> <li>• Prescribed volcanic forcing (aerosol extinction; no physical aerosol equivalent) derived from satellite, airborne, balloon, and ground based observations (GLoSSAC V1; Thomason et al., 2018)</li> </ul>
E3SMv2-PA	ne30pg2 72 vertical levels (0.01 – 1000 hPa)	U + V winds to MERRA2 Reanalysis	Linear ozone chemistry	<ul style="list-style-type: none"> <li>• Injection of volcanic SO<sub>2</sub></li> <li>• Default MAM4 (i.e., no stratosphere-specific modifications)</li> </ul>
E3SMv2-SPA	ne30pg2 72 vertical levels (0.01 – 1000 hPa)	U + V winds to MERRA2 Reanalysis	Linear ozone chemistry	<ul style="list-style-type: none"> <li>• Injection of volcanic SO<sub>2</sub></li> <li>• Prognostic stratospheric aerosol in MAM4</li> <li>• Stratosphere-specific accumulation to coarse mode transfer</li> </ul>
CESM2-WACCM	0.95°x1.25° 88 vertical levels (4.5 × 10 <sup>-6</sup> – 1000 hPa)	U + V winds and temperature to MERRA2 reanalysis	Interactive ozone chemistry	<ul style="list-style-type: none"> <li>• Injection of volcanic SO<sub>2</sub></li> <li>• Prognostic stratospheric aerosol in MAM4</li> <li>• Stratosphere-specific, reversible accumulation-to-coarse mode transfer</li> </ul>

## 2.4 Effective radius and Mie scattering calculations

230 Aerosol size distributions in the model provide information about how MAM4 represents volcanic aerosol evolution and can also help explain changes in radiation balance in the earth system. Here we calculate effective radius and use this in single-particle Mie scattering to understand how changes in size affect diffuse/direct radiation at the surface and absorption of longwave radiation in the stratosphere. See Appendices A and B for more details.





### 3. Observational data sets

In addition to comparing to CESM-WACCM, we also employ observational datasets of sulfate burden, aerosol optical  
235 depth (AOD), and microphysical properties (effective radius and size distributions) to substantiate the performance of the  
prognostic aerosol capability implemented within E3SMv2. Details of the observational datasets are presented below.

#### 3.1 Sulfate burden

##### 3.1.1 HIRS

The High Altitude Infrared Radiation Sounder (HIRS; Baran and Foot, 1994) is an infrared radiometer measuring  
240 surface reflectance at 19 different infrared channels (3.7 – 15  $\mu\text{m}$ ) and one solar channel (0.69  $\mu\text{m}$ ) from a variety of polar  
orbiting NOAA platforms since 1978 (Borbas and Menzel, 2021). This study uses HIRS-derived aerosol mass loading over  
the period of May 1991 through October 1993 from Baran and Foot (1994), who used the difference between 8.3  $\mu\text{m}$  (aerosol  
sensitive) and 12.5  $\mu\text{m}$  (aerosol insensitive) channels to isolate the transmission through the volcanic plume. To back out  
245 aerosol mass loading, they assumed an average stratospheric sulfate aerosol composition of 75%  $\text{H}_2\text{SO}_4$  + 25%  $\text{H}_2\text{O}$ , particle  
size and concentration from dustsonde measurements in July 1991 (Deshler et al., 1992), and a single scattering albedo  
calculated from Mie theory by integrating over scattering and extinction coefficients from an assumed lognormal distribution  
of radius 0.35  $\mu\text{m}$  and standard deviation 1.6. This data covers 80°N – 80°S at 5° resolution with a systematic error of ~10%  
( $\pm 1.4$  Tg aerosol) due to assumptions in processing and uncertainty in background concentration. Additional minimum and  
maximum aerosol composition bounds are included in this range, namely 59%–77%  $\text{H}_2\text{SO}_4$ .

##### 250 3.1.1 SAGE-3A

This work also uses stratospheric sulfate burden taken from the SAGE-3I dataset compiled for Phase 6 of the Coupled  
Model Intercomparison Project (CMIP6; [ftp://iacftp.ethz.ch/pub\\_read/luo/CMIP6/](ftp://iacftp.ethz.ch/pub_read/luo/CMIP6/), last access: 12 January 2023) as reported  
in Quaglia et al. (2023). The SAGE-3I dataset uses the Stratospheric Aerosol and Gas Experiment II (SAGE-II; Section 3.3.2)  
wavelengths of 0.454, 0.525, 1.024  $\mu\text{m}$ , fitting the measured extinction at these wavelengths to a lognormal size distribution  
255 and estimating sulfate mass burden from the number density, mode radius, and width of the distribution (Revell et al., 2017)

#### 3.2 Aerosol Optical Depth

##### 3.2.1 AVHRR

The Advanced Very High Resolution Radiometer (AVHRR; Zhao et al., 2013; Heidinger et al., 2014) is a radiometer  
that measures surface reflectance in six spectral bands (0.63, 0.86, 1.6, 3.75, 11, and 12  $\mu\text{m}$ ) and has served as a meteorological  
260 imaging sensor on the NOAA polar orbiting platforms since 1978. It has a 1.1 km spatial resolution and, during the Pinatubo  
eruption period of interest (1990-1994), had 2-4 global views per day (Heidinger et al., 2014). An offline radiative transfer



model is used to determine lookup tables for AOD retrievals. The radiative transfer model assumes fine and coarse mode aerosol properties based on validation of AVHRR with the surface radiometer measurements from the Aerosol Robotic Network (AERONET; Zhao et al., 2002), and uncertainty in the AVHRR AOD is estimated at 11.3 % based on surface  
265 AERONET validation (Zhao, Xuepeng, 2022). This work uses monthly, clear-sky AOD retrieved over oceans and regridded to 1° resolution. Retrievals are made from channel 1 (0.63 μm) and related to the radiatively equivalent AOD at 0.6 μm through a radiative transfer and surface/atmosphere model (Rao et al., 1989). Detection limits on AOD from the AVHRR range from a minimum of 0.01 and a maximum of 2 (Russell et al., 1996).

### 3.3 Effective radius and size distributions

#### 270 3.3.1 UARS/SAGE-II (Stenchikov et al., 1998)

This work uses column average effective radius derived from two Upper Atmosphere Research Satellite (UARS; Grainger et al., 1995; Lambert et al., 1997) instruments (the Improved Stratospheric Mesospheric Sounder (ISAMS) and Cryogenic Limb Array Etalon Spectrometer (CLAES)), and SAGE-II extinction data (Stenchikov et al., 1998). The UARS instruments ISAMS and CLAES are limb sounders, reporting aerosol extinction –12.11 μm and 12.66 μm wavelengths,  
275 respectively, across altitude profiles above 100 hPa (Lambert et al., 1997). Vertical resolution is approximately 2.5 km and horizontal resolution is about 4° (Taylor et al., 1994). These extinction values are then used to derive aerosol effective radius assuming volcanic aerosol size distribution parameters from balloon borne measurements (Deshler et al., 1992, 1993) and a sulfate refractive index corresponding to 75% H<sub>2</sub>SO<sub>4</sub> aerosol composition (Grainger et al., 1995; Lambert et al., 1997). Another limb sounder, the SAGE-II (McCormick et al., 1995), reports extinction at the 1.02 μm wavelengths which is then used to  
280 derive aerosol number density in the atmosphere. The height resolved number density and effective radius are then used to calculate column average effective radius zonal means at 40°N and 7°S from 200 hPa – 10 hPa, reported in Fig. 4 of Stenchikov et al. (1998).

#### 3.3.2 SAGE-II

The SAGE-II (Mauldin et al., 1985) instrument flew from October 1984 to August 2005 on the Earth Radiation  
285 Budget Satellite (ERBS), measuring light extinction through the atmospheric limb at seven channels from 0.385 to 1.02 μm. The global coverage is 80°S-80°N and 1 km vertical resolution. This work uses aerosol effective radius from SAGE-II version 7 (Damadeo et al., 2013) over the tropics (20°S-20°N), limited to 21-27 km (50-20 hPa) due to sparse data at lower altitudes (Quaglia et al., 2023). The effective radius is derived from a combination of extinction inversion algorithms that make use of the extinction ratios between 0.525 and 1.02 μm and assume that the aerosols are spheres with a sulfate composition of 75%  
290 H<sub>2</sub>SO<sub>4</sub> + 25% H<sub>2</sub>O by mass (Damadeo et al., 2013).



### 3.3.3 WOPC

The balloon-borne University of Wyoming optical particle counter (WOPC; Deshler et al., 1993; Deshler, 1994, 2003) uses particle scattering of white light to measure particle counts across 8-12 channels that range in size from 0.15 – 2  $\mu\text{m}$  (Kalnajs and Deshler, 2022). These particle counts are then fit to unimodal or bimodal size distributions such that they minimize the root mean square error in number concentration between the measured cumulative count and the integral of the size distribution (Deshler, 2003). This instrument was launched from 1989-2013 from a variety of locations across the globe, with the most continuous sampling in Laramie, Wyoming. Instrument uncertainties include:  $\pm 10\%$  in concentration and  $\pm 10\%$  in aerosol radius, and  $\pm 40\%$  in the distribution moments (e.g., surface area, volume, and extinction) (Deshler, 2003; Deshler et al., 2019). In this work we reproduce WOPC aerosol size distributions and derived effective radius over Laramie, WY.

## 4. Results

This section traces model performance from the global to the microphysical, tying the latter to the former through single particle optical properties. The global performance of E3SMv2-SPA stratospheric mass burden (Section 4.1) and AOD (Section 4.2) is determined through comparisons with remote sensing data. Regional comparisons to remote and in situ observations of stratospheric  $R_{\text{eff}}$  (Section 4.3) identify a model small bias, which is explored in more detail through analysis of aerosol size distributions (Section 4.4). Lastly, model data is compared to individual balloon-borne measurements of aerosol size distributions and  $R_{\text{eff}}$ , relating the effective single particle scattering (Section 4.5) and absorption (Section 4.6) efficiencies to changes in direct/diffuse radiation at the surface (Section 4.5.1) and longwave heating in the stratosphere (Section 4.6.1).

### 4.1 Sulfur burden

Figure 1 shows the stratospheric mass burden of the sulfur component of sulfate aerosol in the different model sensitivity tests, the HIRS observational dataset (Baran and Foot, 1994), and the SAGE-31 dataset (Revell et al., 2017). The model data is processed to match the HIRS and SAGE-31 data coverage of 80°N – 80°S above the model lapse rate tropopause height. In MAM4 of E3SM and CESM, sulfate is assumed to be ammonium bisulfate (i.e.,  $(\text{NH}_4)\text{HSO}_4$ ) in the atmosphere (Liu et al., 2012, 2016) (Liu et al., 2012), so the sulfur component is determined by scaling modeled sulfate mass by the ratio of sulfur and ammonium bisulfate mass weights such that  $Tg S = Tg SO_4 * \frac{32.066 \text{ g mole}^{-1}}{115.11 \text{ g mole}^{-1}}$ .

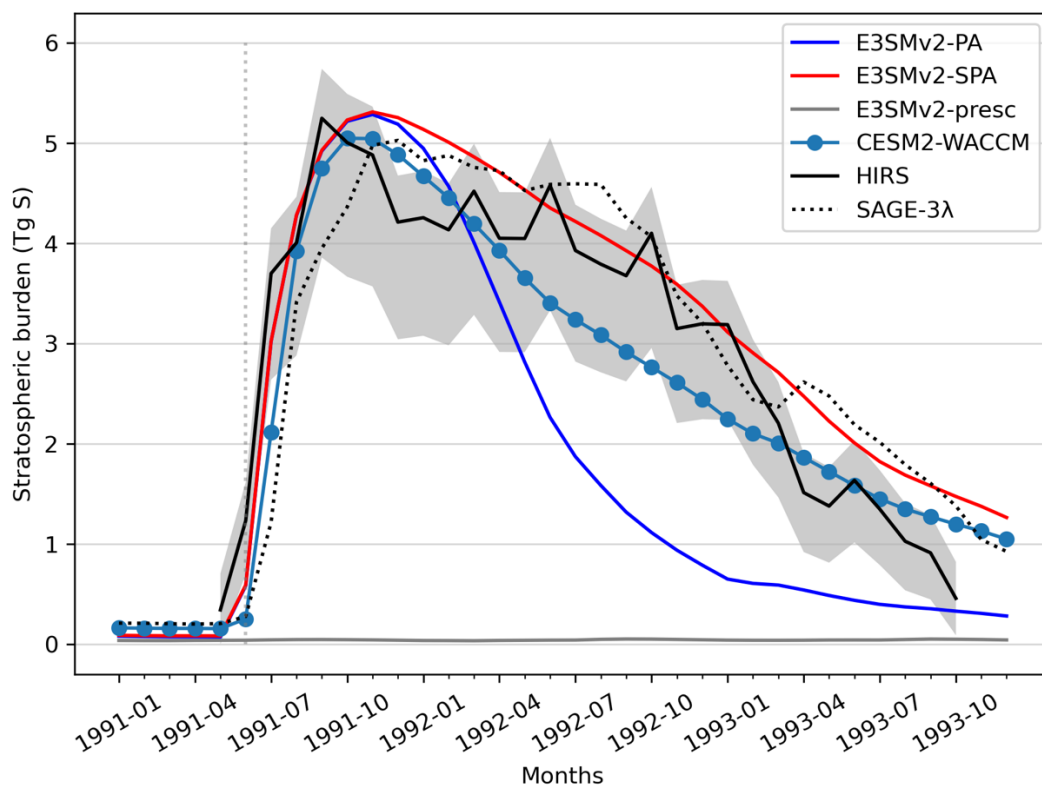
When compared to HIRS and SAGE-31, E3SMv2-SPA improves the modeled stratospheric aerosol burden over E3SMv2-PA, especially in the years following the Pinatubo eruption. The increased aerosol lifetime in the stratosphere is mainly due to our modifications to the coarse mode  $s$  in E3SMv2-SPA. The underestimated aerosol lifetime in E3SMv2-PA is mainly caused by too wide an aerosol number distribution, causing fast sedimentation of the larger coarse mode particles in the upper tail of the distribution. The E3SMv2-SPA tends to overestimate aerosol burden compared to HIRS in the 6 months after Pinatubo but agrees well with the slow decay reported in HIRS and SAGE-31 during 1992. From 1992 onward, stratospheric mass burden in E3SMv2-SPA agrees the best with SAGE-31, which reports higher burdens in 1993 than HIRS.



E3SMv2-SPA and WACCM are similar in atmosphere and aerosol treatments, but have very different atmospheric chemistry, which seems to impact lifetime.

In CESM2-WACCM, the interactive hydroxyl radical (OH) treatment causes OH depletion in the plume as the oxidation of SO<sub>2</sub> to form sulfate aerosol depletes available OH and therefore limits the reaction rate (Mills et al., 2017). This is in contrast to E3SMv2 which assumes an OH climatology unaffected by the oxidation of SO<sub>2</sub>. The result is faster depletion of SO<sub>2</sub> and higher initial sulfate concentrations in the stratosphere in E3SMv2 (Fig. S3). The difference in OH treatment can be seen in Fig. 1, marked by the faster increase in stratospheric sulfate burden in E3SMv2-PA and E3SMv2-SPA. Another significant difference between the two chemical treatments is the presence of carbonyl sulfide (OCS) in CESM2-WACCM, which is a largely inert tropospheric chemical species that is oxidized and photolyzed when it enters the stratosphere, forming sulfate. This is a major contributor to non-volcanic stratospheric sulfate and will lead to a higher pre-Pinatubo sulfate concentrations in the stratosphere in CESM2-WACCM than in E3SMv2 (Mills et al., 2016). In Fig. 1, the effect of OCS is shown in larger pre-Pinatubo stratospheric burdens in CESM2-WACCM.

Figure 1 shows that E3SMv2-SPA performs reasonably well when forming sulfate from SO<sub>2</sub> and simulating increased aerosol lifetimes. The following sections will address how well the model parameterizes the aerosol microphysical properties and their impacts on the global radiative balance.





340 **Figure 1: Stratospheric sulfate burden – reported in Tg of the sulfur mass contribution – for model sensitivity tests, as well as HIRS and SAGE-3λ remote sensing observations.**

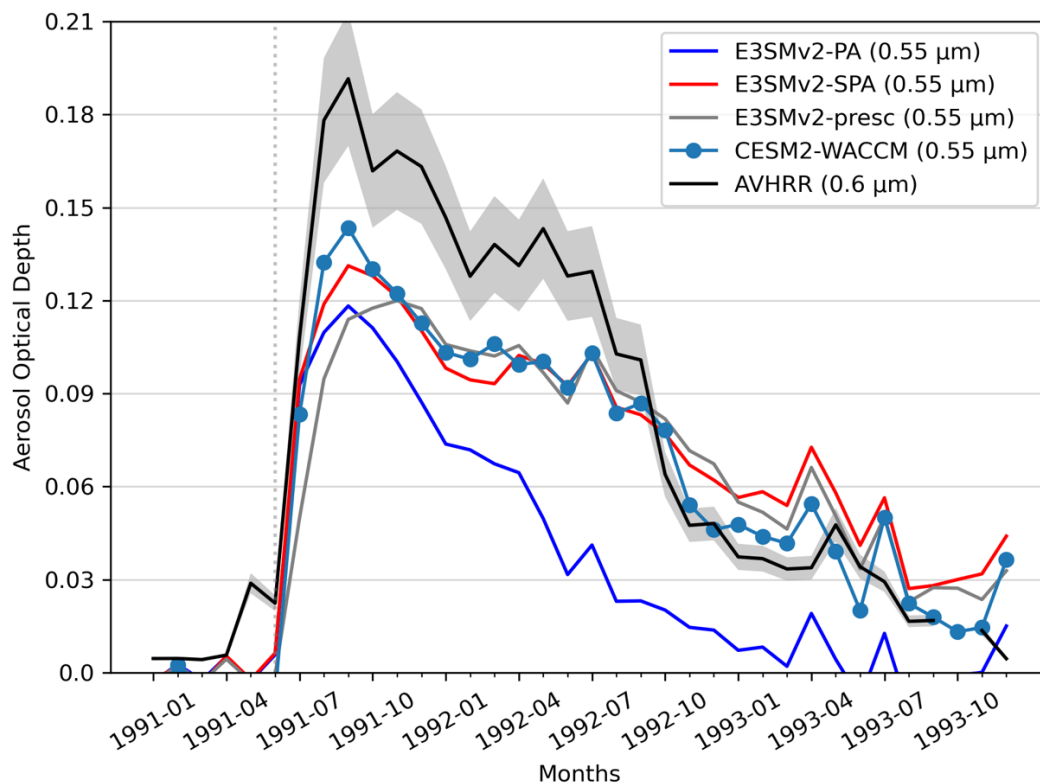
#### 4.2 Aerosol optical depth

In Figure 2, the stratospheric contribution is isolated by subtracting the monthly mean AOD from pre-Pinatubo years, relying on the assumption that non-volcanic background aerosol in the atmosphere is similar in the near-term. Model data is normalized to 1990 monthly means, while AVHRR is normalized to the monthly means for the period June 1989 to May 1991, 345 with the exception of missing data from October-December 1990. Normalized model data is masked to reflect the same temporal and spatial sampling as AVHRR data from 1991-1993 over the oceans only and between 60°N – 60°S. Here, AOD from the models is reported at 0.55 μm, while the AVHRR AOD is 0.6 μm.

The AVHRR AOD peaks two months after Pinatubo, linearly decreasing except for periods of flattening in the months March – July 1992 and January – July 1993. The flattening in 1992 is attributed to the continual growth of aerosols due to 350 coagulation, which increases sulfate aerosol scattering efficiency (Russell et al., 1996; Stenchikov et al., 1998) even as aerosol burden continues to decay (Fig. 1). This is supported by a simulated increasing accumulation mode – and to a lesser extent, coarse mode –  $D_g$  (Section 4.4; Fig. 4) and a slight increase in global average stratospheric aerosol effective radius (Section 4.3; Fig. 3). The flattening in 1993 for both observations and models may be due in part to the influence of the smaller Chilean volcanic eruption Lascar (1993-01-30; 23.36°S), which has a discernable impact on modeled global accumulation mode aerosol 355 mean diameters  $>0.2 \mu\text{m}$  (Section 4.4; Fig. 4).

In the first year following Pinatubo, AVHRR and the volcanically parameterized models (i.e., E3SMv2-SPA, E3SMv2-presc, and CESM2-WACCM) follow a similar trend, decreasing and leveling off near the beginning of 1992. As for aerosol mass burden, the overly short stratospheric aerosol lifetime in E3SMv2 leads to a rapid decay in AOD. All models consistently underpredict AOD in the first year after the eruption but tend to overpredict AOD in the third year following the 360 eruption. Underprediction may be due in part to a lack of volcanic ash and incorrect number/size representation in the models. Overprediction is likely a factor of aerosol removal assumptions. Of the models, CESM2-WACCM has the closest agreement with AVHRR in 1993, due to a faster decline in AOD than E3SMv2-SPA. This indicates that the comprehensive chemistry in CESM2-WACCM may better represent aerosol size distributions than E3SMv2-SPA (see Section 4.4).

A surprising finding is the close agreement between E3SMv2-SPA and E3SMv2-presc. These models utilize very 365 different approaches: prognostic aerosol microphysics from emitted  $\text{SO}_2$  in MAM4 (E3SMv2-SPA) and prescribed stratospheric aerosol extinction from a range of observations in GloSSAC (E3SMv2-presc). This indicates that the changes to aerosol microphysics in MAM4 reasonably recreate the reanalysis data product, especially from 1992 onward.



370 **Figure 2: Stratospheric aerosol optical depth (AOD) over the ocean and across latitudes 60°S–60°N from the model simulations and AVHRR. Models calculate AOD at 0.55 μm and AVHRR channel-1 AOD (0.63 μm) is processed to 0.6 μm. Both datasets are normalized to volcanically quiescent periods as described in the text. The Pinatubo eruption is marked with the gray dotted line at 1991-06.**

### 4.3 Aerosol effective radius

375 Effective radius has frequently been used to characterize stratospheric aerosol properties, with stratospheric  $R_{\text{eff}}$  of less than  $\sim 2 \mu\text{m}$  leading to a net solar radiation scattering effect and surface cooling (Lacis et al., 1992). Based on a range of in-situ and remote sensing datasets, background stratospheric  $R_{\text{eff}}$  is estimated at 0.17-0.19 μm, and following Pinatubo reaches average values of around 0.5 μm with observed values as large as 0.8 or 1.0 μm (Russell et al., 1996, and references therein). In the month following Pinatubo, there is little change in  $R_{\text{eff}}$ . This is due to a rapid increase in very small (i.e., Aitken mode sulfate) and very large (i.e., ash) aerosol particles following the eruption. Enhanced coagulation and condensation, coupled with low sedimentation rates, lead to a steady increase in  $R_{\text{eff}}$  over the next 3-6 months. Aerosol growth continues until approximately mid 1992 when  $R_{\text{eff}}$  peaks, lagging peak values in other metrics such as mass burden and AOD.

385 Figure 3 shows global  $R_{\text{eff}}$  (Fig. 3a), in addition to  $R_{\text{eff}}$  from three different regional zones specific to different observational datasets: comparisons at 40°N and 7°S latitude bands and less than 100hPa with UARS (Fig. 3b), 41°N 105°W between 130-10hPa with WOPC (Fig. 3c), and 20°S-20°N between 50-20hPa with SAGE-II (Fig. 3d). Observations tend to measure a minimum size that falls in the middle of the Aitken mode in the model. To account for this characteristic of the data,



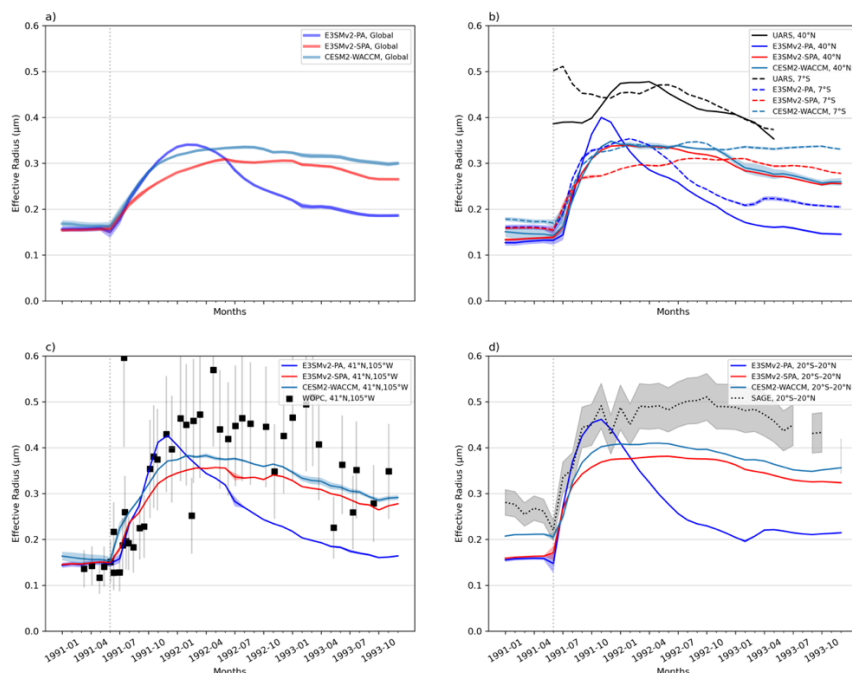
the model  $R_{\text{eff}}$  is an average of effective radii calculated with and without the Aitken mode, with a shaded range between the two  $R_{\text{eff}}$  represented by shading about the line. Maximum differences tend to occur before Pinatubo, shortly after Pinatubo, and with other volcanic eruptions (e.g., Cerro Hudson (1991-08-08; 45.9°S)). In Fig. 3a the models reproduce the expected background  $R_{\text{eff}}$  of 0.17-0.19  $\mu\text{m}$ , and the improvements to aerosol lifetime in E3SMv2-SPA and CESM2-WACCM can be  
390 seen in the slower decrease in  $R_{\text{eff}}$  compared to E3SMv2-PA. There is also a nearly identical pattern in E3SMv2-SPA and CESM2-WACCM data, but with slightly higher  $R_{\text{eff}}$  in CESM2-WACCM.

All of the models underestimate  $R_{\text{eff}}$  compared to observations (Fig. 3b-d). In Fig. 3b, the finer temporal responses to the Pinatubo eruption in the UARS data are less apparent in the model, namely the large peak in  $R_{\text{eff}}$  at 7°S associated with short-lived volcanic ash in the stratosphere and the delayed peak at 40°N and 7°S due to particle aggregation (Stenchikov et al., 1998). The models neglect volcanic ash contributions, explaining the more gradual particle growth at 7°S. While the models don't show the same sensitivity to increases in  $R_{\text{eff}}$  at these latitude bands – possibly due to higher vertical and horizontal spatial scale in the number concentrations from SAGE-II observations – E3SMv2-SPA does show a flattening of the 7°S  $R_{\text{eff}}$  akin to the UARS estimate. This corresponds to the eruption of Cerro Hudson in Chile, and the resulting high influx of smaller, Aitken mode particles into the southern stratosphere which drives down  $R_{\text{eff}}$ . The sensitivity to this eruption in E3SMv2-SPA  
400 may be due to higher Aitken mode production in this model than in CESM2-WACCM (see Section 4.4).

The tropical regions tend to have a larger  $R_{\text{eff}}$  than the midlatitudes in the simulations. This is true of UARS regions 6-12 months after the eruption and also of SAGE-II (tropics; Fig. 3d) data. These larger  $R_{\text{eff}}$  persisting in the tropics a year or more after the eruption. While all models exhibit similar modal aerosol diameters (Fig. S1), the higher  $R_{\text{eff}}$  is correlated with higher number concentrations in all aerosol modes and a slower decrease in aerosol number (i.e., reduced sedimentation) (Fig. S2). The reduced removal is likely due to higher initial concentrations of  $\text{SO}_2$  in the volcanic plume over the tropics (Fig. S3) contributing to more rapid local aerosol growth and a net positive aerosol production. Furthermore, the presence of the upwelling branch of the Brewer-Dobson circulation in this latitude band may help suspend larger aerosol species, slowing aerosol sedimentation rates and increasing their lifetime. The eruption of the Lascar volcano in northern Chile around 1993-02 also contributes to a bump in  $\text{SO}_2$ , as well as Aitken and accumulation mode aerosol number at the 7°S band (Fig. S2, S3).  
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When comparing the models to WOPC (Fig. 3c) and SAGE-II (Fig. 3d), E3SMv2-PA has the closest agreement to these datasets in its initial aerosol growth. This growth is more rapid than the other models and leads to a peak in  $R_{\text{eff}}$  that, while being closer to observed values, drops off precipitously. The  $R_{\text{eff}}$  in E3SMv2-SPA and CESM2-WACCM have the best agreement with observational values and decay rate a year or more after Pinatubo. Differences across the models are due to the different microphysical assumptions, which can be explored by looking at aerosol size distributions.





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**Figure 3:** Stratospheric aerosol effective radius averaged (a) globally above model tropopause, (b) at 40°N and 7°S and <100 hPa to compare with UARS data (Stenchikov et al., 1998), (c) over Laramie, Wyoming (41°N, 105°W) and 130-10 hPa to match WOPC data (Quaglia et al., 2023), and (d) over 20°S-20°N and 50-20 hPa to match SAGE-II observations (Quaglia et al., 2023). The shaded range in models represents effective radius with and without the Aitken mode, and the model line is an average of the two. Error bars in WOPC data assume a 40% uncertainty and correlation coefficients of 0.5 between aerosol moments and levels (see Appendix A2 in Quaglia et al. (2023)). The Pinatubo eruption is marked with the gray dotted line at 1991-06.

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#### 4.4 Aerosol size distributions

While  $R_{\text{eff}}$  is a good representation of aerosol size in the context of optical properties, it is not always a good indicator for behavior of aerosol microphysical processes. For example, an increase in accumulation mode particle number and a decrease in coarse mode particle number could manifest as an unchanging  $R_{\text{eff}}$ . An examination of the aerosol size distribution can be more informative when understanding how aerosol chemistry and the MAM4 microphysics contribute to model performance.

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Figure 4 shows the globally averaged, stratospheric, aerosol distributions from the three models used in this study and their evolution from 1991 to the end of 1993. The contour fill represents aerosol number ( $\text{dN}/\text{dlogD}$ ;  $\text{cm}^{-3}$ ) with dashed and dotted lines indicating modal  $D_g$  ( $\mu\text{m}$ ). In all three models, there is growth in Aitken mode due to new sulfate particle formation following major and minor volcanic eruptions during this period: Pinatubo (1991-06; vertical dotted line), Cerro Hudson (1991-08), Spurr in the Aleutian Islands (1992-06-27; 61.3°N), and Lascar (1993-02). In E3SM (Fig. 4a-b), the prescribed concentrations of OH lead to a rapid oxidation of available  $\text{SO}_2$  and higher concentrations of Aitken mode aerosol



435 compared to CESM2-WACCM (Fig. 4c), resulting in higher aerosol number concentrations in E3SM. It also appears that CESM2-WACCM has higher tropospheric aerosol transport into the stratosphere based on a higher concentration of Aitken mode aerosol across the period of interest. CESM-WACCM has seasonal peak concentrations occurring asynchronous with volcanic eruptions and corresponding to northern hemisphere winter (e.g., 1991 and 1992). The increased Aitken mode number concentration is also seen at 40°N (Fig. S2). These peaks are attributed to a lower tropopause in Northern Hemisphere winter and may also be due to the inclusion of OCS in CESM-WACCM.

440 The larger volcanic eruptions, Pinatubo and Hudson, inject enough SO<sub>2</sub> into the stratosphere that Aitken mode aerosols grow through condensation and coagulation into the accumulation mode, which then grow through coagulation into the coarse mode. The exception to this is in E3SMv2-PA which lacks the ability to transfer sulfate mass into the coarse mode and so retains a roughly constant coarse mode D<sub>g</sub> derived from trace mass (10<sup>-22</sup>-10<sup>-21</sup> kg cm<sup>-3</sup>; global average) and number (10<sup>-5</sup>-10<sup>-4</sup> cm<sup>-3</sup>; global average) concentrations of dust, seasalt, and sulfate aerosol advected from the troposphere (Note: mass and number are related to aerosol size through Equations S1 and S2 in the Supp. Material). Aerosol growth through the modes  
445 can be seen in increased number and an increasing trend in aerosol size. Dips in modal D<sub>g</sub> correspond to sudden increases in aerosol number (e.g., Aitken mode nucleation from freshly injected SO<sub>2</sub>, transfer of a large aerosol number from Aitken to accumulation mode) while mass remains relatively unchanged. This leads to a division of mass across a larger number of aerosols and a subsequent decrease in D<sub>g</sub>.

450 Overall, the aerosol modal diameters are similar across the three models, but slight differences exist. There are slightly larger accumulation mode and coarse mode D<sub>g</sub> in CESM2-WACCM (0.262 μm; 0.843 μm) than E3SMv2-SPA (0.259 μm; 0.749 μm). This could be a factor of the interactive chemistry in CESM2-WACCM leading to fewer Aitken mode aerosol that initially grow faster through condensation given the longer lived SO<sub>2</sub> in the stratosphere. Contrast this with E3SM where more numerous, smaller aerosol consume the available SO<sub>2</sub> more quickly. The E3SMv2-PA model accumulation mode D<sub>g</sub> reaches  
455 a higher maximum D<sub>g</sub> (0.285 μm) than E3SMv2-SPA or CESM2-WACCM (~0.26 μm). This is due to the missing coarse mode treatment in this model coupled with a larger D<sub>g,hi</sub> (Section 2.1.2). The result is an accumulation mode that grows to D<sub>g,hi</sub> following Pinatubo, whereby number is increased to maintain this size instead of transferring mass and number to a larger mode. This also explains the better initial agreement in R<sub>eff</sub> between E3SMv2-PA and observations in Fig. 3b-d.

Figure 5 compares modeled aerosol size distributions to in-situ measurements from WOPC, tracking the evolution of  
460 a slice of the plume from pre-Pinatubo (1991-04-19) to the end of 1993 (1993-11-16). Daily samples taken with the WOPC over Laramie, WY at a single level (18 km; roughly corresponding to the peak plume R<sub>eff</sub> across the period (Fig. S4)) are used to validate daily data from E3SMv2-PA and E3SMv2-SPA. Modal D<sub>g</sub> (vertical lines) and dN/dlogD are denoted by dotted and dashed lines, as in Fig. 4. Effective diameter (D<sub>eff</sub>) for models and WOPC are included to relate the changing distributions to aerosol-light interactions.

465 What follows is a breakdown of Fig. 5 into a rough timeline evolution in aerosol size distributions and D<sub>eff</sub>:



- Pre-Pinatubo conditions (Fig. 5a) are very similar between E3SMv2-PA and E3SMv2-SPA, and the  $D_{\text{eff}}$  are nearly identical between model and observations. CESM2-WACCM has higher background number concentrations than E3SM.
- One month after Pinatubo (Fig. 5b), all datasets exhibit rapid growth in the smaller diameter aerosols, with some growth in the modeled accumulation mode but little sign of the coarse mode in E3SMv2-SPA. CESM2-WACCM has a massive increase in Aitken and accumulation mode number, possibly due to transport of  $\text{SO}_2$  and Aitken mode aerosol into the region from the tropics. This Aitken mode peak is also seen at 17 and 19 km levels (Fig. S5-S6), and results in a decrease in  $R_{\text{eff}}$  over this height range (Fig. S4). The  $D_{\text{eff}}$  are still similar across the different datasets, with the exception of CESM2-WACCM which is smaller due to the Aitken mode influence.
- Six months after the eruption (Fig. 5c), a clear coarse mode signal emerges in WOPC, E3SMv2-SPA, and CESM2-WACCM along with a sharp increase in modeled accumulation mode number. Models and WOPC are all comparable in their bimodal  $D_g$  and accumulation mode number, while WOPC coarse mode has lower number than E3SMv2-SPA and CESM2-WACCM. E3SMv2-PA has a larger  $D_{\text{eff}}$  than other datasets (though it is still within the uncertainty of WOPC  $D_{\text{eff}}$ ). This peak in  $D_{\text{eff}}$  – which is also noted in  $R_{\text{eff}}$  in Fig. 3 – is attributed to a wider accumulation mode and larger modal  $D_g$  in E3SMv2-PA. CESM2-WACCM has larger modal number than all other datasets.
- 11 months after the eruption (Fig. 5d), the WOPC data continues to grow in accumulation and coarse mode. Accumulation mode number is decreasing while coarse mode is increasing in E3SMv2-SPA, suggesting conversion of accumulation mass into the longer-lived coarse mode. The WOPC coarse mode  $D_g$  ( $1.92 \pm 0.19 \mu\text{m}$ ) is more than twice that of E3SMv2-SPA ( $0.828 \mu\text{m}$ ), while the WOPC accumulation mode best-fit ( $D_g = 0.909 \pm 0.09 \mu\text{m}$ ,  $s_g = 1.23$ ) is close to E3SMv2-SPA coarse mode ( $D_g = 0.828 \mu\text{m}$ ,  $s_g = 1.2$ ) and is nearly identical to the CESM2-WACCM coarse mode ( $D_g = 0.918 \mu\text{m}$ ,  $s_g = 1.2$ ). . The overall larger aerosols in the WOPC distributions lead to largest difference in  $D_{\text{eff}}$  between model and WOPC.
- For one year and longer after the eruption, accumulation mode in E3SMv2-PA is on a steady and rapid decay that can be tracked in size and number (Fig. 5d-i). Over this period, E3SMv2-SPA and CESM2-WACCM generate very similar size distributions, with slightly lower coarse mode number and slightly larger coarse  $D_g$  in CESM2-WACCM. Overall, WOPC and E3SMv2-SPA/CESM2-WACCM have similar distributions and similar  $D_{\text{eff}}$ , with E3SMv2-SPA showing slightly smaller  $D_{\text{eff}}$  but still within the uncertainty of WOPC data. A peak in Aitken mode aerosol in the models in Fig. 5h may be due to the Lascar eruption in 1992-02.

The previous timeline shows an improved aerosol evolution in E3SMv2-SPA compared to WOPC and exhibits the similarities between CESM2-WACCM and E3SMv2-SPA. It also indicates that E3SMv2-SPA doesn't simulate large enough aerosol compared to the observations. Similar behavior in  $D_{\text{eff}}$  and relative size bias is noted at 17 km, 19 km, and 20 km levels as well (Fig. S5-S7). Given that the model is nudged to meteorology and that the aerosol in the stratosphere tend to evolve relatively slowly over time, it is assumed that the resolution artifacts of comparing a  $1^\circ \times 1^\circ$  model grid cell to in-situ

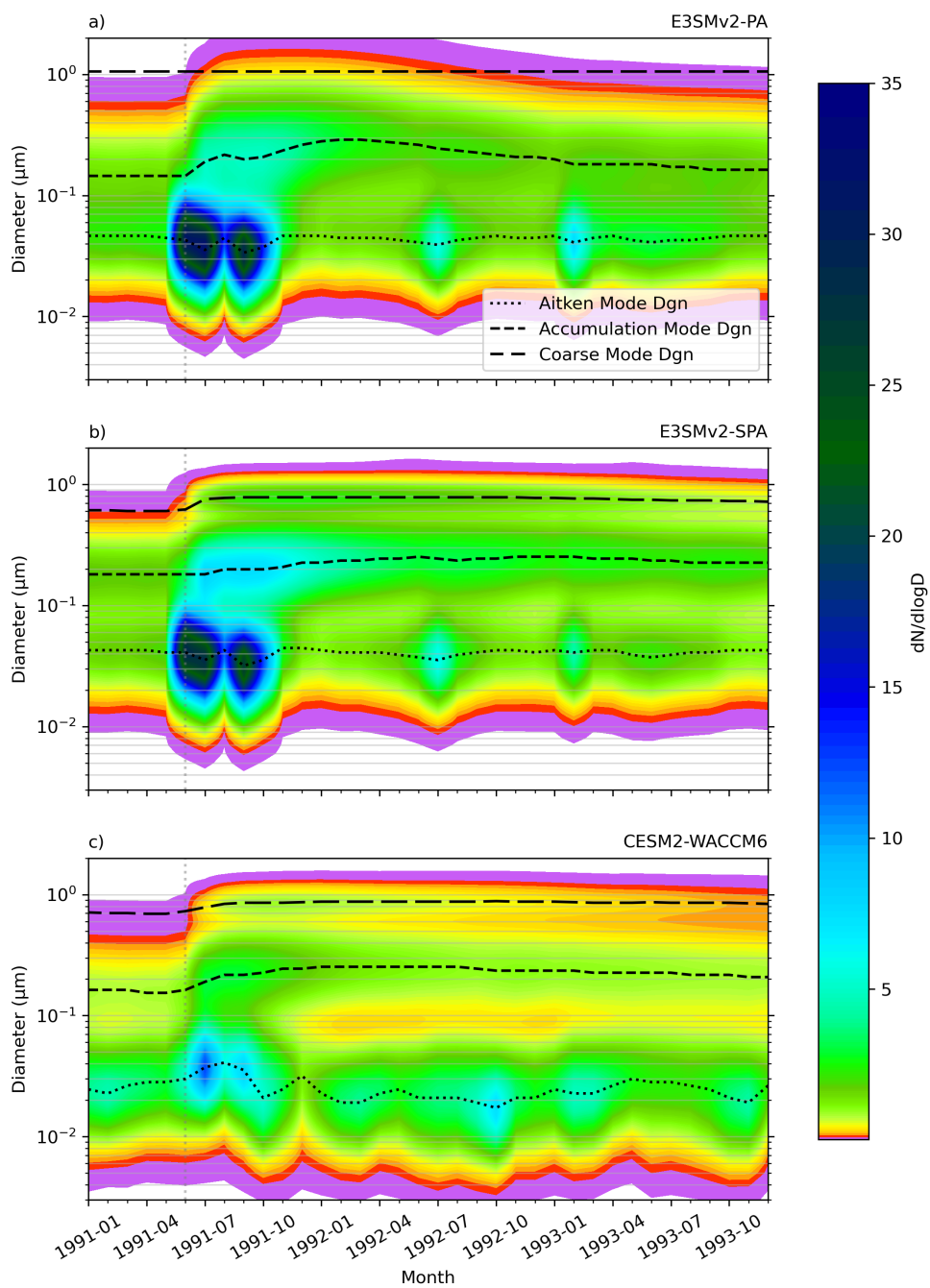
<https://doi.org/10.5194/egusphere-2023-3041>

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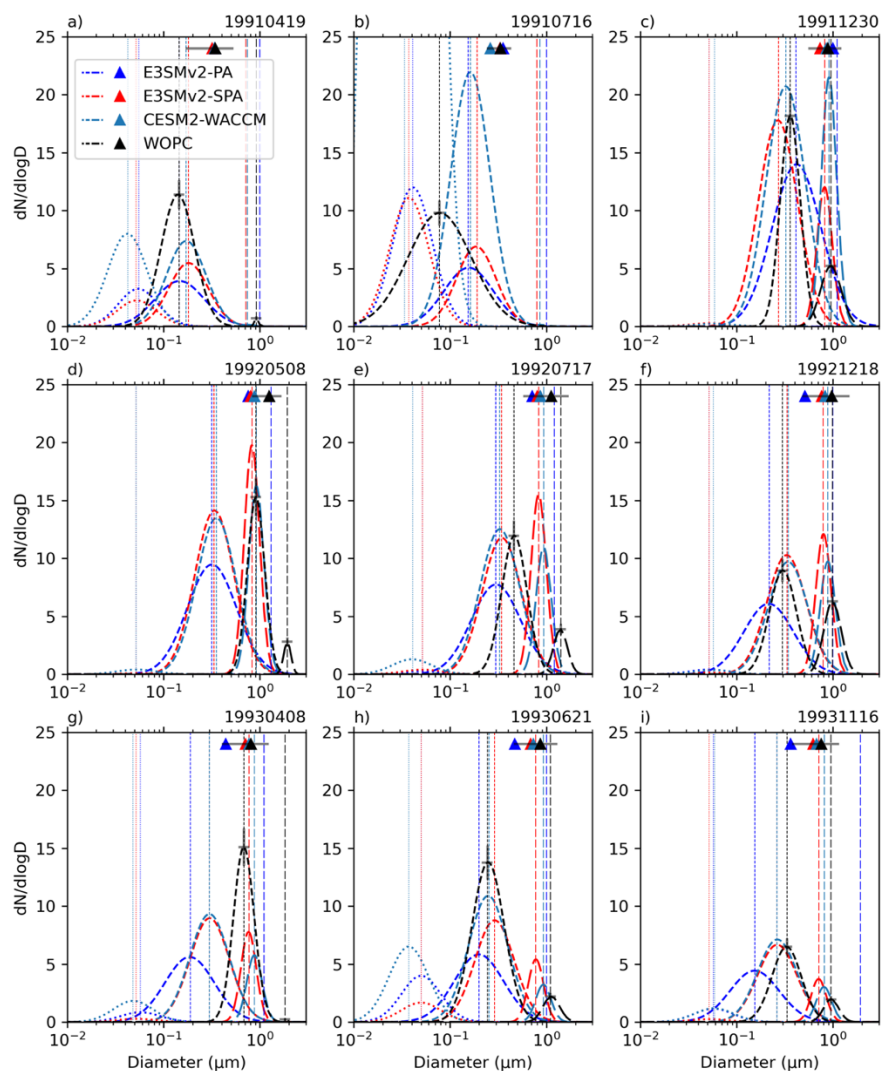
500 measurements are minimal. The impact of this size bias on aerosol climate impact is now explored through effective single particle scattering.





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**Figure 4: Globally averaged, stratospheric lognormal aerosol size distributions from 1991-1993 for E3SMv2-PA, E3SMv2-SPA, and CESM2-WACCM. The dotted, dashed, and long-dashed lines indicate geometric mean diameters of the Aitken, Accumulation, and Coarse aerosol modes, respectively. The color contour is lognormally scaled. The Pinatubo eruption is marked with the gray dotted line at 1991-06.**



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**Figure 5: Stratospheric aerosol size distributions from E3SMv2-PA (blue), E3SMv2-SPA (red), and WOPC (black) from 1991-1993. WOPC launches chosen to correspond to spring, summer, and winter measurements for each year, and samples are taken from the 18 km measurements and matched to the nearest model height and grid cell over Laramie, WY (41.3°N, 105°W). The dotted, dashed, and long-dashed lines indicate geometric mean diameters (vertical) and dN/dlogD size modes of the Aitken, Accumulation, and**



Coarse aerosol modes, respectively. Triangles denote effective diameters derived from the size distributions. Uncertainties in WOPC diameter, number, and effective diameters are denoted by grey bars at the peak of the distributions and on the markers. Note: anomalously large max number in CESM2-WACCM Aitken mode ( $376.5 \text{ cm}^{-1}$ ) is out of the plotting bounds.

#### 515 4.5 Single particle scattering efficiency

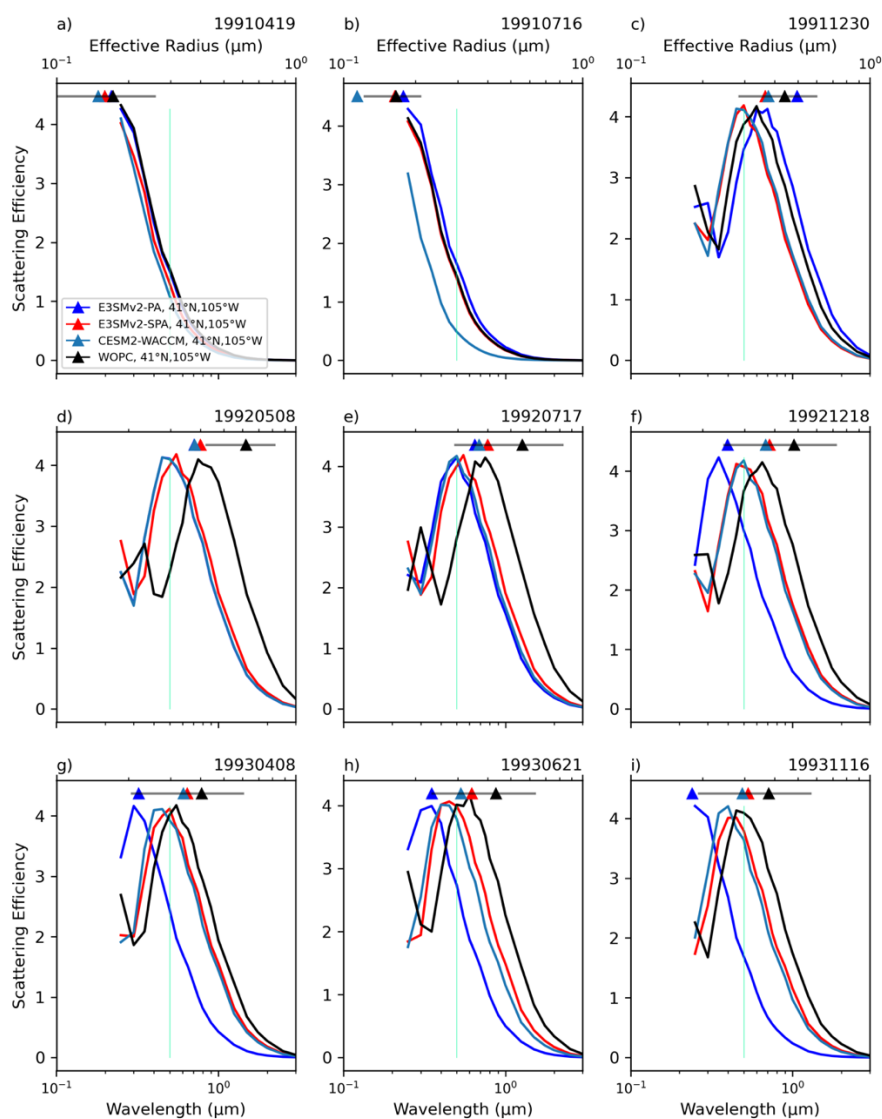
The  $R_{\text{eff}}$  can be used to characterize the evolving aerosol size distribution, which conveniently can be used in the offline calculation of single particle optical properties using Mie theory. Figure 6 shows the scattering efficiency ( $Q_s$ ) of sulfate particles of size  $R_{\text{eff}}$  for the same time and height samples as Fig. 5. Additionally, this plot marks the approximate wavelength of peak solar black-body irradiance ( $\lambda_{\text{solar}}=0.5 \mu\text{m}$ ) to identify when aerosol scattering of solar radiation has the largest impact on radiative balance.  $Q_s$  is proportional to the  $x_{\text{eff}}^4$  (i.e.,  $\lambda^{-4}$ ) (Petty, 2006). This scattering feature can be seen in Fig. 6 where, as wavelengths increase, you will get a rapid decrease in  $Q_s$ .

The peak in  $Q_s$  shifts from  $\lambda$  of  $\sim 0.2 \mu\text{m}$  before and shortly after Pinatubo (Fig. 6a-b), to  $0.5\text{-}0.7 \mu\text{m}$  6 months after the eruption (Fig. 6c) with the onset of rapid particle growth. For the remainder of the time (Fig. 6d-i), E3SMv2-SPA and CESM2-WACCM  $R_{\text{eff}}$  result in a  $Q_s$  peak right around  $\lambda_{\text{solar}}$ , indicating a very efficient scattering of sunlight by these simulated aerosols. The WOPC reports larger  $R_{\text{eff}}$  than the models, which have a  $Q_s$  that is 10-80% higher than observations at  $\lambda_{\text{solar}}$  during 1992 (Fig. 6d-f), becoming more similar during 1993 with the decay of stratospheric  $R_{\text{eff}}$  (Fig. 6g-i). E3SMv2-PA has a similar  $Q_s$  to E3SMv2-SPA/CESM2-WACCM until the end of 1992 (Fig. 6f) when there is a quick progressive drop in  $R_{\text{eff}}$  due to rapid stratospheric aerosol deposition.

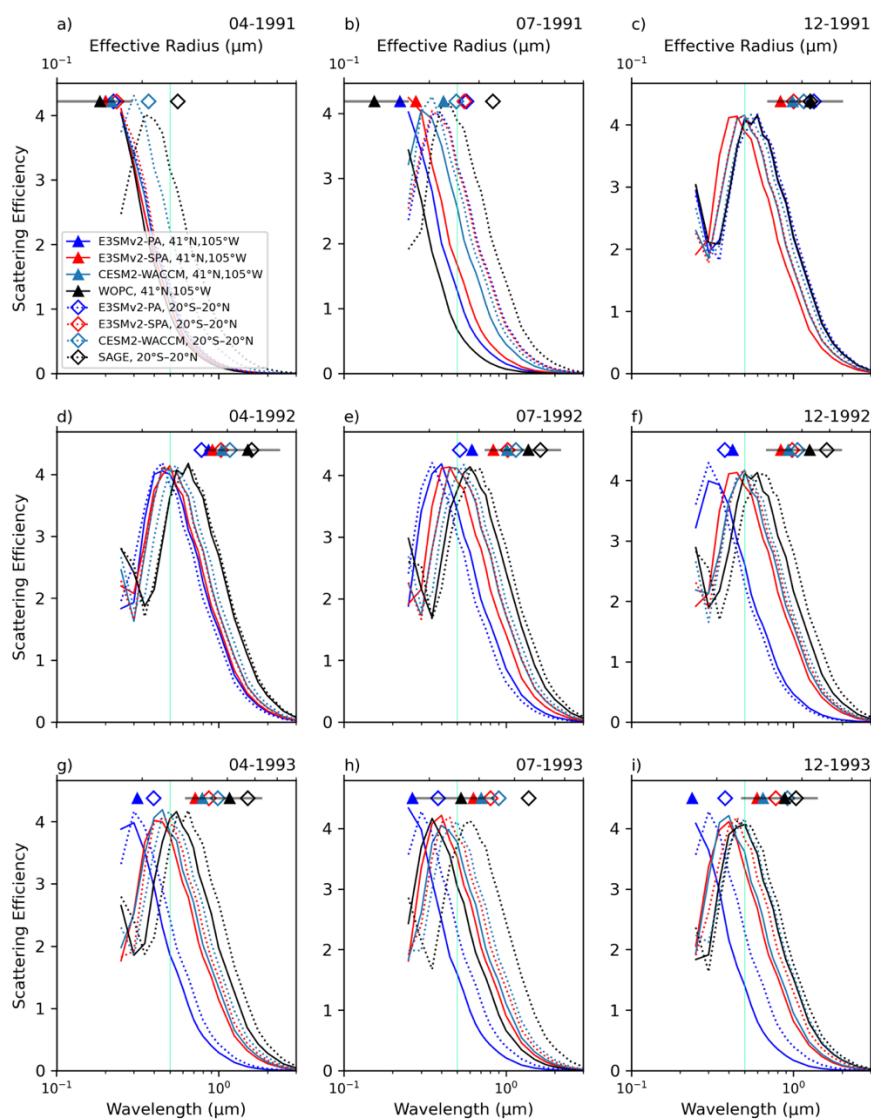
These results indicate an optimal sulfate scattering  $R_{\text{eff}}$  of  $\sim 0.3\text{-}0.4 \mu\text{m}$ , which is similar to estimates of the optimal sulfate aerosol size of  $0.3 \mu\text{m}$  in geoengineering applications (Dykema et al., 2016). Given that WOPC predicts larger  $R_{\text{eff}}$  and smaller  $Q_s$  at  $\lambda_{\text{solar}}$  6-18 months after Pinatubo, our results suggest that E3SMv2-SPA/CESM2-WACCM overestimates aerosol scattering and the global cooling effect of the Pinatubo aerosol at this level. Figure 6 relates size distribution differences in models and WOPC to climate impact at a single level. For a better validation of the stratospheric plume and to see if the model disagreement changes when looking across multiple levels, the next step is to calculate  $Q_s$  from the stratospheric  $R_{\text{eff}}$ .

Figure 7 shows monthly  $Q_s$  generated from averaged stratospheric  $R_{\text{eff}}$  at a midlatitude site ( $41^\circ\text{N } 105^\circ\text{W}$ ) (Fig. 3c) and the tropics (Fig. 3d). As in Fig. 6, a similar pattern is shown where smaller modeled  $R_{\text{eff}}$  lead to higher modeled  $Q_s$  at  $\lambda_{\text{solar}}$  than observed  $Q_s$ . However, the differences between modeled and observed  $Q_s$  at  $\lambda_{\text{solar}}$  are not as stark for the monthly, stratospheric average values, with max differences on the order of 5-14% as opposed to the 10-80% seen in the daily, 18 km level data. The larger  $R_{\text{eff}}$  in the tropics leads to a higher  $Q_s$  at  $\lambda_{\text{solar}}$  one month after Pinatubo (Fig. 7b), and lower  $Q_s$  at  $\lambda_{\text{solar}}$  18 months after Pinatubo (Fig. 7f), than midlatitude measurements in the stratosphere. The E3SMv2-SPA and CESM-WACCM maximum  $Q_s$  from both regions hover around  $\lambda_{\text{solar}}$ , with higher modeled  $Q_s$  at  $\lambda_{\text{solar}}$  in the tropics 2 years after the eruption (Fig. 7h) due to a consistently larger  $R_{\text{eff}}$  in this region.





545 **Figure 6: Daily 18 km scattering efficiency ( $Q_s$ ) using effective size parameter ( $x_{\text{eff}} = 2\pi R_{\text{eff}}/\lambda$ ) and sulfate refractive index at 0% relative humidity (Hess et al., 1998). Effective radius ( $R_{\text{eff}}$ ) from E3SMv2-PA, E3SMv2-SPA, and WOPC are from Fig. 3c, and are marked by colored triangles at the top of the plot. The bottom x-axis is wavelength ( $\lambda$ ; 0.1–3  $\mu\text{m}$ ). The top axis is  $R_{\text{eff}}$  (0.1–1  $\mu\text{m}$ ) used in the calculation of  $Q_s$  via the effective size parameter ( $x_{\text{eff}}$ ). The turquoise vertical line marks the solar black body wavelength of maximum irradiance (0.5  $\mu\text{m}$ ).**



550 **Figure 7: Monthly stratosphere scattering efficiency ( $Q_s$ ) using effective size parameter ( $x_{\text{eff}} = 2\pi R_{\text{eff}}/\lambda$ ) and sulfate refractive index at 0% relative humidity (Hess et al., 1998). Monthly  $R_{\text{eff}}$  are from E3SMv2-PA, E3SMv2-SPA, CESM2-WAACCM, WOPC, and SAGE-II from Fig. 3c-d, and are marked by colored triangles (41°N 105°W; WOPC) and hollow diamonds (20°S–20°N; SAGE-II) at the top of the plot. The bottom x-axis is wavelength ( $\lambda$ ; 0.1–3  $\mu\text{m}$ ). The top axis is  $R_{\text{eff}}$  (0.1–0.7  $\mu\text{m}$ ) used in the calculation of  $Q_s$  via the effective size parameter ( $x_{\text{eff}}$ ). The turquoise vertical line marks the solar black body wavelength of maximum irradiance (0.5  $\mu\text{m}$ ).**

555



#### 4.5.1 Direct / diffuse radiation

The differences in  $R_{\text{eff}}$  between the different model and observational datasets ultimately affect diffuse and direct radiation at the surface by changing AOD and other aerosol optical properties. In scattering incoming shortwave, a small fraction of light is scattered back to space reducing the amount of energy incident to Earth. The forward scattered SW radiation increases the diffusivity of incident radiation. The Mauna Loa observatory documented a lessening of the direct beam and increase in the diffuse radiation after the eruption of Mt. Pinatubo (Robock, 2000), and globally there was a 21% reduction in direct and a 20% increase in diffuse radiation at the surface (Proctor et al., 2018).

The model simulations use a two-stream approximation for calculating multiple scattering in the atmosphere (Iacono et al., 2008; Neale et al., 2012), which calculates diffuse and direct radiation at the surface from incident radiation, AOD (i.e., the column integrated  $b_s$  (Eq. B2)), single scattering albedo (SSA; Fig. S8), and aerosol asymmetry parameter ( $g$ ; Fig. S9). In a general sense, the input aerosol optical properties determine what fraction of direct radiation makes it through the aerosol layer (AOD), how much of the incident radiation is absorbed in the aerosol layer ( $1-\text{SSA}$ ), and the degree to which the radiation is scattered by the aerosol layer ( $g$ ).

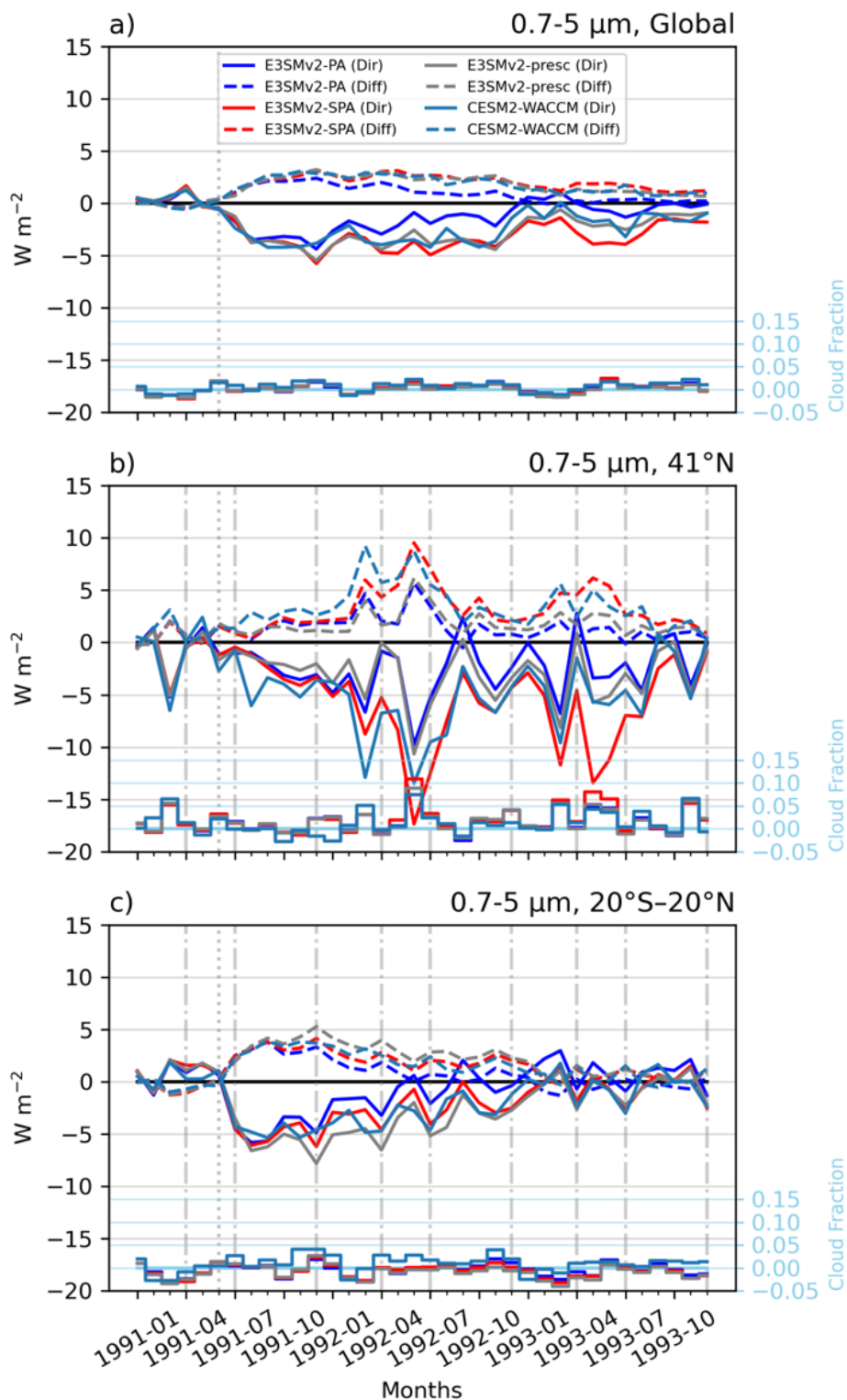
Figure 8 shows the perturbation in diffuse and direct radiation at the surface due to the Pinatubo aerosol layer in the visible to ultraviolet wavelengths (with similar behavior in the near-infrared (Fig. S10)). The values are averaged over the whole earth (Fig. 8a), the  $41^\circ\text{N}$  latitude band (Fig. 8b), and  $20^\circ\text{S}$ - $20^\circ\text{N}$  tropical region (Fig. 8c). There is an increase in diffuse radiation across all of the models and a mirrored decrease in direct radiation. These changes track the increasing  $R_{\text{eff}}$  (Fig. 3) and the enhanced  $Q_s(\lambda)$  (Fig. 7). The magnitude of the direct radiation loss tends to be larger than the increase in diffuse radiation at the surface because some of the scattered incoming solar radiation is lost to space. In all regions, E3SMv2-SPA and CESM2-WACCM have more diffuse and less direct radiation than E3SMv2-PA, the latter of which rapidly approaches zero as the aerosol are removed from the atmosphere. The peak direct and diffuse influence occur in the tropics (Fig. 8c) prior to the  $41^\circ\text{N}$  band (Fig. 8b) as aerosol transport to northern latitudes relies on stratospheric circulation. In the  $41^\circ\text{N}$  band (Fig. 8b) CESM2-WACCM has a higher proportion of diffuse radiation at the surface than all of the other models.

Diffuse radiation is not only attributed to aerosol size and material properties – as is  $Q_s$  in Fig. 7 – but is also related to the aerosol number concentration. As a result, AOD serves as a good indicator of diffuse/direct radiation at the surface as it accounts for both number and  $Q_s$ . Figure 9 shows AOD over the same spatial regions as Fig. 8. The AOD behaves very similarly to diffuse radiation. As in Fig. 8b, Fig. 9b shows CESM2-WACCM with higher AOD over the  $41^\circ\text{N}$  band. This is attributed to both the larger aerosol  $R_{\text{eff}}$  (Fig. 3c) and higher accumulation and coarse mode number concentrations (Fig. S2a) in CESM2-WACCM. The higher number concentrations are correlated with more  $\text{SO}_2$  transport into the midlatitudes in CESM2-WACCM (Fig. S3b), which is likely due to slower oxidation of  $\text{SO}_2$  in CESM2-WACCM due to the aforementioned OH depletion in this model.



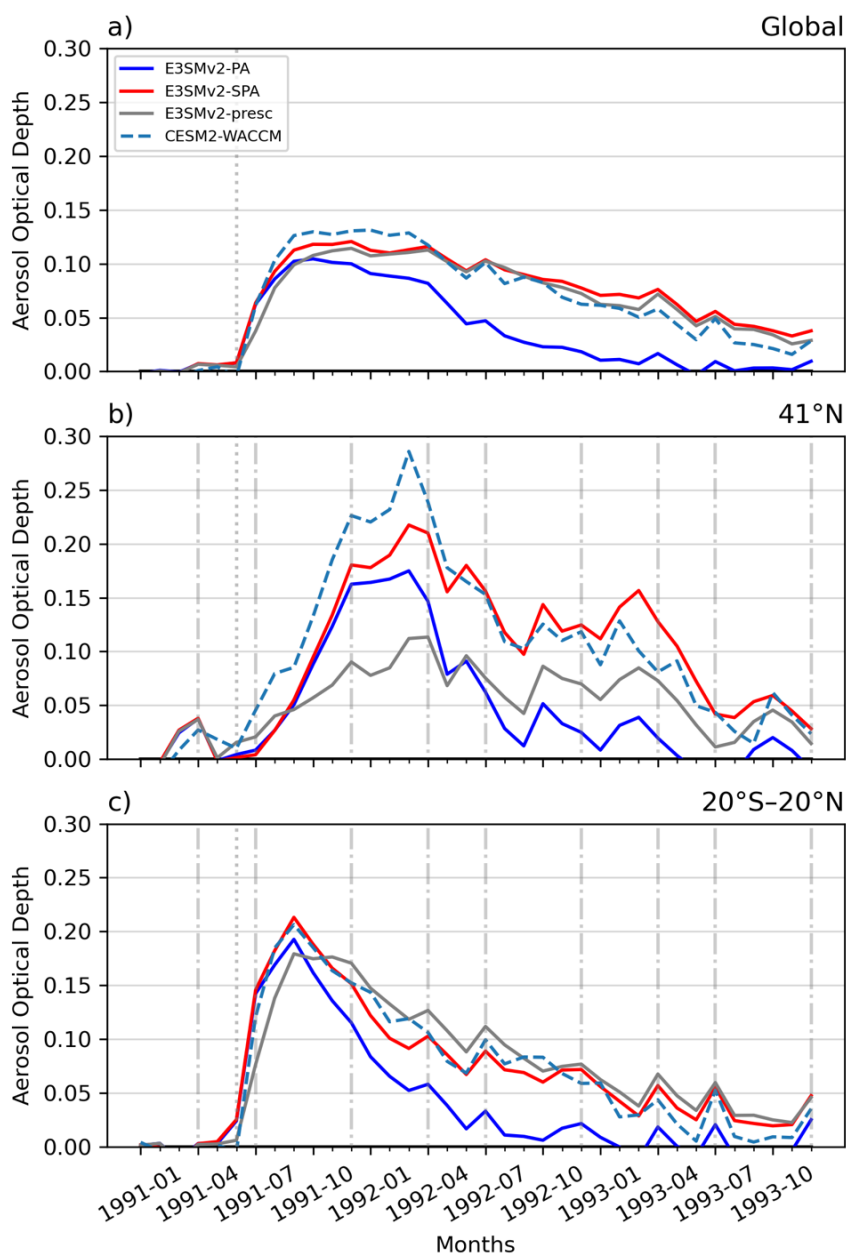
590 The impact of clouds on diffuse/direct radiation in the midlatitudes shows up in Fig. 8b, with peaks in normalized total cloud cover (right axis) corresponding to diffuse (direct) radiation peaks (troughs), especially in the late spring/early summer of 1992 and 1993. The approximate co-location of peaks in cloud cover (Fig. 8b) and AOD (Fig. 9b) in the midlatitudes in June of 1992 and 1993 may indicate the presence of aerosol indirect effects. This is consistent with findings by Liu and Penner (2002) which show high rates of homogeneous ice nucleation persisting in the northern midlatitudes through July 1992.

595 Lastly, E3SMv2-presc has good agreement with E3SMv2-SPA and CESM2-WACCM when compared globally or in the tropics, but reports ~60% the modeled magnitudes for diffuse/direct radiation and ~40-60% modeled AOD in the midlatitudes (Fig. 8 and 9, respectively). This may be due to our use of an older version of the GLoSSAC dataset (Thomason et al., 2018) in which higher latitudes have a low bias in AOD attributed to linear interpolation of the SAGE-II data (Kovilakam et al., 2020).





600 **Figure 8: Pinatubo induced changes in diffuse and direct solar radiation at the surface in the 0.2-0.7  $\mu\text{m}$  band (left axis), and total cloud cover fraction (right axis; blue). Regions include global (a), midlatitude band (b;  $41^\circ\text{N}$ ), and tropical band (c;  $20^\circ\text{S}$ - $20^\circ\text{N}$ ) averages. The Pinatubo eruption is marked with the gray dotted line at 1991-06, and the sample months for Fig. 7 are marked with the gray dash-dot lines. The data is normalized to 1990 monthly means.**



605 **Figure 9:** Pinatubo induced changes in aerosol optical depth (AOD) at 0.55  $\mu\text{m}$ . Regions include global (a), midlatitude band (b; 41°N), and tropical band (c; 20°S-20°N) averages. The Pinatubo eruption is marked with the gray dotted line at 1991-06, and the sample months for Fig. 7 are marked with the gray dash-dot lines. The data is normalized to 1990 monthly means.



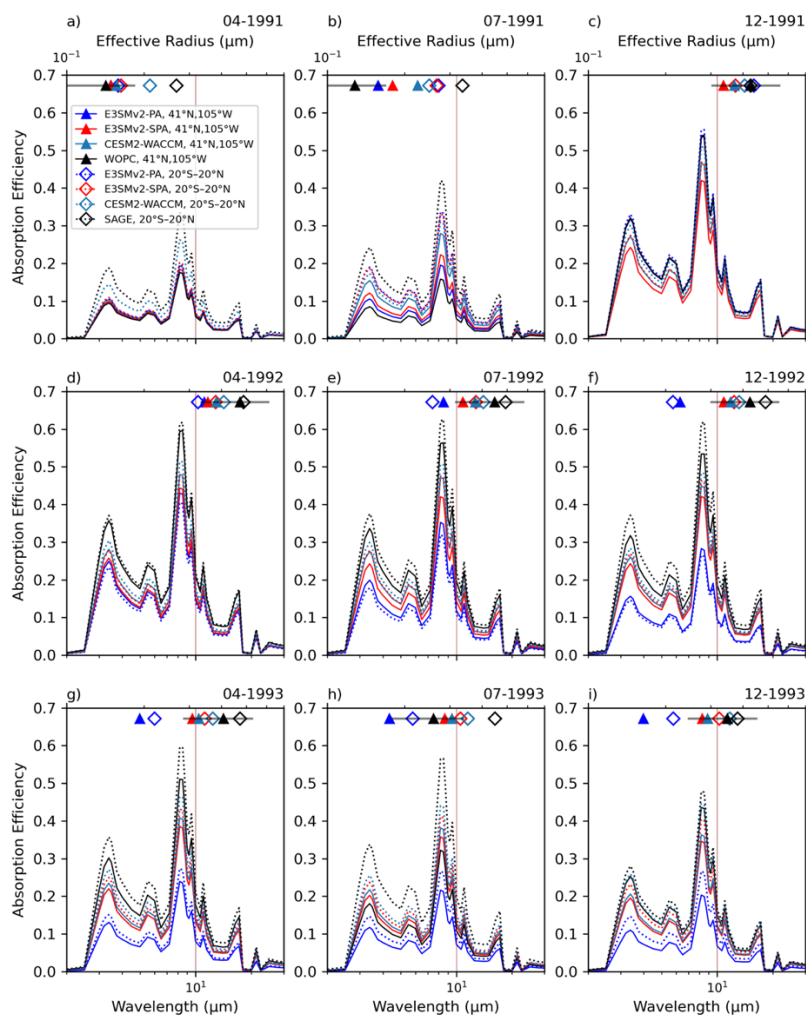


#### 4.6 Single particle absorption efficiency

While the Pinatubo eruption resulted in a net climate cooling effect due to increased scattering, it also contributed to stratospheric warming through the absorption of outgoing longwave radiation within the aerosol layer (Kinne et al., 1992). Figure 10 shows the  $Q_a$  calculated from  $R_{\text{eff}}$ , with a dark red line indicating the wavelength of peak terrestrial black-body irradiance ( $10\ \mu\text{m}$ ;  $\lambda_{\text{earth}}$ ). Sulfate is not an effective absorber at visible wavelengths ( $Q_a = 0$  at  $\lambda_{\text{solar}}$ ) and the determining factor in stratospheric heating is the magnitude of  $Q_a$  at  $\lambda_{\text{earth}}$ . The wavelength dependence of  $Q_a$  is weaker than  $Q_s$ , (proportional to  $x_{\text{eff}}$  (i.e.,  $\lambda^{-1}$ )) and is strongly tied to the imaginary part of  $n_{\text{Hess}}$ . The proportionality to  $x_{\text{eff}}$  is denoted in Fig. 10 by the linear increase in  $Q_a(\lambda)$  with increasing  $R_{\text{eff}}$ ; the dependence on the imaginary part of  $n_{\text{Hess}}$  is reflected in the unchanging pattern in absorption magnitudes for all datasets and times.

In Fig. 10, the linear dependence of  $Q_a$  on  $R_{\text{eff}}$  means that an individual particle in the tropics has more absorption than the midlatitudes across the whole period, and the observations will have more absorption than the models. This is explained by the effective radii in these two regions at the top of each panel. For example, the smaller modeled  $R_{\text{eff}}$  at  $41^\circ\text{N}$  leads to a modeled  $Q_a$  that is 15-30% lower than WOPC through 1992 (Fig. 10c-f).

Because the same refractive index is used across all models and observational datasets, the differences in  $Q_a$  are due only to differences in aerosol size. However, this is a simplifying assumption for the modeled  $Q_a$  due to their explicit calculation of the soluble aerosol refractive index ( $n_s$ ) in the stratosphere based on aerosol water and sulfate content (Eq. B4). E3SM simulates a stratosphere with unrealistically low stratospheric water vapor (Keeble et al., 2021; Jablonowski et al., in prep) which results in aerosol that are mostly sulfate (Fig. S11) with a refractive index similar to  $n_{\text{Hess}}$  (Appendix B). Higher water vapor in CESM2-WACCM means that the assumption of  $n_{\text{Hess}}$  in calculation of  $Q_a$  doesn't account for the volume-weighting of the hydrated aerosol refractive index (Appendix B). When interstitial stratospheric water and sulfate are used to calculate a volume-weighted  $n_s$  using  $n_{\text{Hess}}$  and  $n_{\text{wat}}$ ,  $Q_{\text{abs}}$  at  $\lambda_{\text{earth}}$  is smaller than that reported with only  $n_{\text{Hess}}$  by  $\sim 12\text{-}19\%$  over the  $41^\circ\text{N}$  band and  $\sim 13\text{-}17\%$  over  $20^\circ\text{S}\text{-}20^\circ\text{N}$  (Table S3). This is due to a smaller magnitude imaginary refractive index for water at this wavelength.



**Figure 10: Monthly stratosphere absorption efficiency ( $Q_a$ ) using effective size parameter ( $x_{\text{eff}} = 2pR_{\text{eff}}/\lambda$ ) and sulfate refractive index at 0% relative humidity (Hess et al., 1998). Monthly  $R_{\text{eff}}$  are from E3SMv2-PA, E3SMv2-SPA, CESM2-WAACCM, WOPC, and SAGE-II from Fig. 3c-d, and are marked by colored triangles (41°N 105°W; WOPC) and hollow diamonds (20°S–20°N; SAGE-II) at the top of the plot. The bottom x-axis is wavelength ( $\lambda$ ; 3–40  $\mu\text{m}$ ). The top axis is  $R_{\text{eff}}$  (0.1–0.7  $\mu\text{m}$ ) used in the calculation of  $Q_a$  via the effective size parameter ( $x_{\text{eff}}$ ). The dark red line marks the terrestrial black body wavelength of maximum irradiance (10  $\mu\text{m}$ ).**

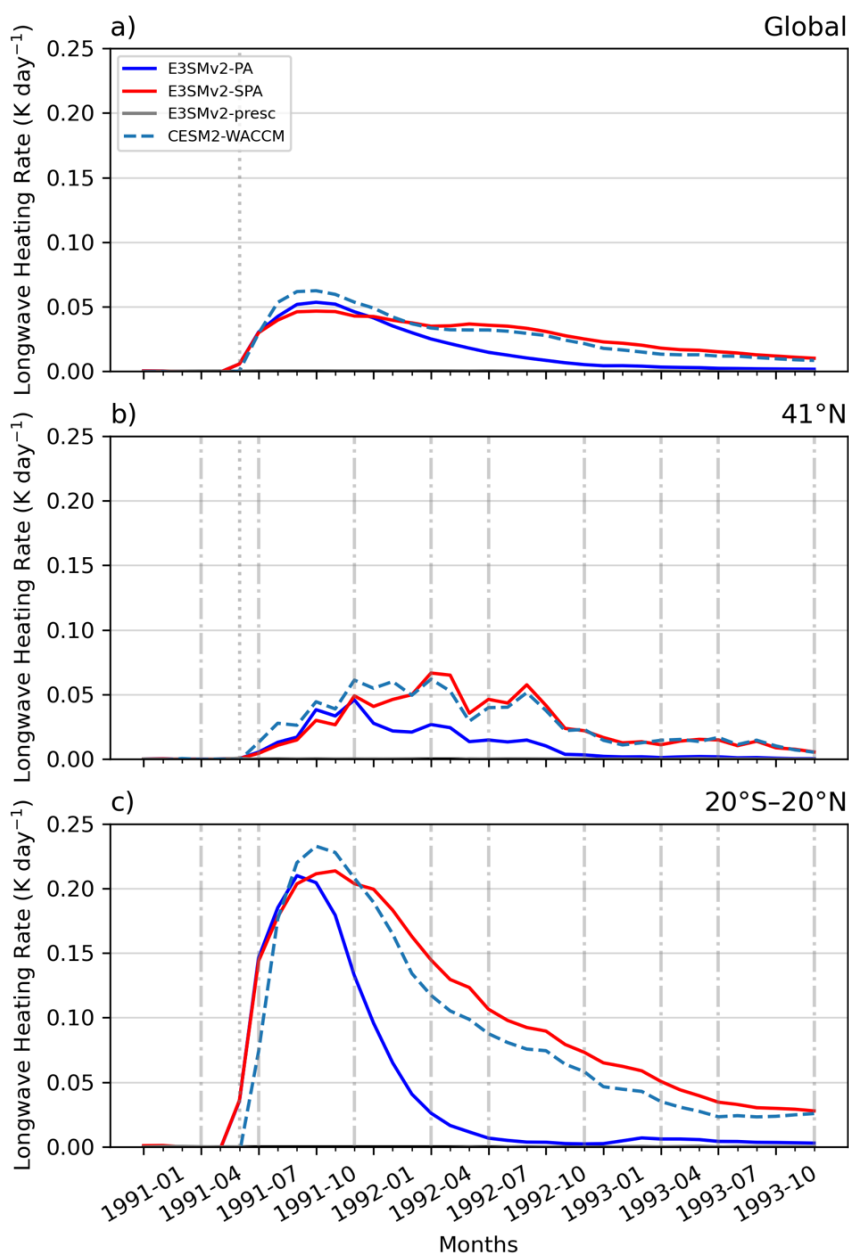
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#### 4.6.1 long wave heating rates

640           The absorption of outgoing longwave radiation in the stratosphere results in local heating of the aerosol layer. Figure  
11 shows the normalized longwave heating rate (LWH;  $\text{K day}^{-1}$ ) due only to Pinatubo aerosols over the same regions as Fig.  
8 and 9. Higher initial global (Fig. 11a) and  $41^\circ\text{N}$  (Fig. 11b) LWH in CESM2-WACCM up to March 1992 mirrors model  
behavior for AOD (Fig. 9a-b). The similarity to AOD indicates a LWH dependence on aerosol absorption optical depth  
(AAOD; level integrated  $b_a$  (Eq. B3)) – and therefore aerosol number concentration and  $Q_a$  – in these two regions. As with  
645 diffuse/direct radiation (Section 4.5.1), characterizing the stratospheric longwave heating rate based solely on  $Q_a$  (i.e., aerosol  
size) is an oversimplification of the actual stratospheric absorption because number concentration is also acting to scale the  $b_a$ .

The E3SMv2-presc has no Pinatubo signal in LWH. This is related to how the prescribed volcanic forcing is treated  
in the model and is not indicative of a missing physical mechanism. While a select number of diagnostics include contributions  
from the GLoSSAC forcing dataset (e.g., AOD, radiation flux, SSA), LWH does not.



650

**Figure 11:** Pinatubo induced changes in longwave heating rate (K day<sup>-1</sup>) over the modeled longwave wavebands (waveband midpoints of ~3.5-514 μm) (Neale et al., 2012). Regions include global stratospheric mean (a), a midlatitude band stratospheric mean at WOPC pressure levels (b; 41°N, 130-10 hPa), and a tropical band stratospheric mean at SAGE-II pressure levels (c; 20°S-20°N, 50-10 hPa). The Pinatubo eruption is marked with the gray dotted line at 1991-06, and the sample months for Fig. 7 are marked



655 with the gray dash-dot lines. The data is normalized to 1990 monthly means and contributions from non-aerosol effects (i.e., clouds and gases) have been removed.

## 5. Discussion

The preceding sections present encouraging results for E3SMv2-SPA and its ability to model explosive volcanic eruptions. These results begin at the global scale and move into the micro-scale, tying the latter to the former to understand the strengths and weaknesses of the modal aerosol modeling approach. While this stratospheric parameterization is new to E3SM, it is based on a modeling framework developed for CESM2-WACCM which has an additional interactive chemistry element not present in E3SM (Mills et al., 2016). Along with the variety of observational datasets, E3SMv2-SPA is validated against CESM2-WACCM which elucidates the effects of interactive chemistry on aerosol formation.

On the global scale, E3SMv2-SPA performs well compared to observational datasets and has similar behavior to CESM2-WACCM. The near-globally averaged stratospheric sulfate burden in E3SMv2-SPA agrees well with SAGE-3I and HIRS remote sensing datasets (Fig. 1), indicating that the model does a reasonable job of simulating aerosol lifetime and loading in the stratosphere. The higher mass burdens in E3SMv2-SPA than in CESM2-WACCM point to the OH consumption in CESM2-WACCM's interactive chemistry limiting aerosol formation from SO<sub>2</sub>. For near-global AOD, the volcanically parameterized E3SMv2-SPA, CESM2-WACCM, and E3SMv2-presc all have similar performance and all exhibit a similar shape to the retrieved AVHRR AOD (Fig. 2). The models underestimate AOD compared to AVHRR in the 15 months following the eruption, but then agree well with AVHRR to the end of 1993. An aspect of the initial eruption that is left out of the models is the injection of large volcanic ash particles, and this may partly explain the low AOD bias shortly after the eruption. However, these particles are short lived (~3 days (Guo et al., 2004a)), and the likely culprits for the AOD underestimation are aerosol number and size representation.

The modelled aerosol size – represented by  $R_{\text{eff}}$  – is consistently underestimated compared to both in-situ and remote sensing datasets (Fig. 3). E3SMv2-SPA has slightly smaller  $R_{\text{eff}}$  than CESM2-WACCM, where the slightly larger aerosols are attributed to fewer initial Aitken mode aerosol associated with the limited OH oxidation in CESM2-WACCM. Unpacking  $R_{\text{eff}}$  into its component modal distributions on a global scale illustrates the larger modal  $D_g$  and also smaller modal number concentration in CESM2-WACCM (Fig. 4). While useful for visualizing the evolution of a variety of volcanic eruptions over this period, validation with observations requires the finer-scale comparison of model and in-situ data.

The comparisons with balloon-borne WOPC data (Fig. 5) identify a key reason for the underestimation in  $R_{\text{eff}}$ . Namely, that modelled coarse mode aerosols are not big enough in the models during the first year after Pinatubo. The aerosol modes reach a similar pattern of change after a year and there is very close agreement in the microphysics of E3SMv2-SPA and CESM2-WACCM. The inability to create large enough aerosol in the models can affect the simulated climate impacts, and the size-based effects on scattering and absorption are explored through the single particle Mie scattering.

The plotted  $Q_s$  (Fig. 6,7) and  $Q_a$  (Fig. 10) tell an idealized story about how size affects scattering of incoming solar and absorption of outgoing terrestrial radiation. At the wavelength of peak incoming solar irradiance (0.5  $\mu\text{m}$ ;  $I_{\text{solar}}$ ), the larger



observed aerosols manifest as weaker scatterers of incident radiation than E3SMv2-SPA and CESM2-WACCM. This indicates that the models may overestimate aerosol scattering, though the differences between model and observation are smaller for monthly as opposed to daily averages. For  $Q_a$ , the relationship is simpler, with a linear increase in aerosol absorption with increasing aerosol size. Our results show that the models may underestimate absorption of light at the wavelength of peak outgoing terrestrial radiation ( $10 \mu\text{m}$ ;  $I_{\text{earth}}$ ) due to their smaller simulated aerosol size. When the effects of modeled interstitial aerosol water are included in the CESM2-WACCM refractive index, modeled  $Q_a$  is decreased further when compared to observations due to a lower absorption of  $I_{\text{earth}}$  by water than by sulfate.

Care should be taken in interpreting  $Q_s$  and  $Q_a$  as there are some caveats to the above comparisons. While effective radius is a reasonable representation of aerosol size and is an important factor in aerosol optics, it is not the singular factor in aerosol optical properties. Aerosol number is neglected in these comparisons, which is a secondary determiner to aerosol AOD. The importance of aerosol number can be seen in the diffuse radiation (Fig. 8), which has relative model behaviors that mirror those of AOD (Fig. 9). Number is also important in longwave absorption, where relative modeled longwave heating rates (Fig. 11) similarly mirror those in AOD. Another assumption in the calculation of scattering and absorption efficiencies is the assumption of a stratospheric sulfate refractive index. This is consistently done for remote sensing datasets where no additional information can be provided, but in models that prognostically calculate the volume-weighted soluble refractive index this may be misleading. Here, the same refractive index is assumed across observations and models to make them more comparable. It also avoids differences in CESM2-WACCM and E3SM that arise from the extremely low stratospheric water bias in E3SM. The accurate simulation of number and refractive index are two aspects of this work that require further study.

This study has served to validate the stratospheric prognostic aerosol treatment in E3SMv2-SPA against observations, while also showing it can have reasonable performance compared to the more comprehensive chemical treatment of CESM2-WACCM. We have connected the microphysical evolution of aerosols to radiative impacts, which then cascade into further downstream impacts (i.e., changes in temperature, precipitation, water vapor, etc.) that will be explored in follow-on papers. While an interactive chemical treatment does have better overall agreement with observations of mass burden, AOD, and  $R_{\text{eff}}$ , E3SMv2-SPA can still serve as a viable alternative. Hence, E3SMv2-SPA will enable studies of climate impacts from stratospheric aerosols with varied source characteristics in free-running fully-coupled climate simulations.

## Appendix A: Effective radius

Aerosol effective radius ( $R_{\text{eff}}$ ) – or area-weighted aerosol radius – is a good representation of a size distribution's optically relevant size (Hansen and Travis, 1974; Russell et al., 1996). This relies on the fact that intercepted light is proportional to the particle radius in Mie and geometric scattering regimes. These regimes can be designated by an aerosol size parameter,  $x = \frac{2\pi r}{\text{wavelength}(\lambda)}$ , with Mie scattering corresponding to  $x \approx 0.2 - 2000$  and geometric scattering corresponding to  $x > 2000$  (Petty, 2006).  $R_{\text{eff}}$  is calculated by taking the ratio of the third aerosol moment (volume) and second aerosol moment (cross-sectional area) ( $M_3$  and  $M_2$ , respectively) of the aerosol size distribution (Hansen and Travis, 1974):



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$$R_{eff} = \frac{M_3}{M_2}, \quad (1)$$

A lognormal size distribution is defined as

$$725 \quad n_N(\ln r) = \frac{dN(r)}{d \ln r} = \frac{N}{\ln \sigma_g \sqrt{2\pi}} e^{\left[ -\frac{1}{2} \left( \frac{\ln r - \ln r_g}{\ln \sigma_g} \right)^2 \right]}, \quad (2)$$

where  $N$  is total aerosol number concentration,  $r_g$  is the wet geometric mean radius,  $r$  is the radius, and  $\sigma_g$  is the geometric standard deviation. The analytic expression for the aerosol moments,  $M_n$ , of a lognormal distribution are,

$$730 \quad M_n = N r_g^n e^{\left[ \frac{n^2}{2} (\ln \sigma_g)^2 \right]}, \quad (3)$$

For a multimodal system, the modal moments ( $M_{n,i}$ ) can be summed, giving the final equation for  $R_{eff}$ :

$$R_{eff} = \frac{\sum_{i=1}^m M_{3,i}}{\sum_{i=1}^m M_{2,i}} = \frac{\sum_{i=1}^m N_i r_{g,i}^3 e^{\left[ \frac{9}{2} \ln \sigma_{g,i}^2 \right]}}{\sum_{i=1}^m N_i r_{g,i}^2 e^{\left[ 2 \ln \sigma_{g,i}^2 \right]}}, \quad (4)$$

735

## Appendix B: Single-particle Mie scattering

In order to understand how improvements to stratospheric aerosol size distributions affect simulated climate impact, aerosol scattering and absorption efficiencies ( $Q_s$ ,  $Q_a$ ) are calculated based on Mie theory, the solution of Maxwell's equations for the interaction of radiation with a sphere (Hansen and Travis, 1974). The  $Q_s$  ( $Q_a$ ) is the surface area of the shadow cast by the particle due to scattering (absorption) of intercepted light divided by the geometric cross section ( $\pi r^2$ ).  $Q_s$  and  $Q_a$  are calculated using the python package, miepython (<https://miepython.readthedocs.io/en/latest/index.html#>). This package simulates absorption and scattering of incident light by non-absorbing or partially absorbing spheres and requires a complex aerosol refractive index ( $n = a + bi$ ) and aerosol size parameter (aerosol circumference over the incident wavelength of light) as inputs. In the complex refractive index, the real and imaginary parts define scattering and absorption, respectively. Here, an effective size parameter is used as a proxy for the particle size distributions in our calculations:

745

$$x_{eff} = \frac{2\pi R_{eff}}{\lambda}, \quad (1)$$





We note that because Mie theory calculates these values based on interactions with a single sphere and we are using the number-independent  $R_{\text{eff}}$  as an input,  $Q_s$  and  $Q_a$  do not account for effects of different number concentrations. The volume scattering and absorption coefficients ( $\beta_s, \beta_a$ ; units of  $\text{m}^{-1}$ ) do include the effect of aerosol number, but are not generated based on our comparisons due to the ambiguity of assuming an effective aerosol number. Their equations are included below for completeness.

$$\beta_s = NQ_s\pi r^2, \quad (2)$$

$$\beta_a = NQ_a\pi r^2, \quad (3)$$

The aerosol refractive index in the models is derived from a combination of sulfate aerosol and water due to the prevalence of sulfate in the stratosphere following Pinatubo. Because sulfate is a soluble species, a soluble refractive index ( $n_s$ ) can be calculated with the volume mixing rule (Ghan and Zaveri, 2007):

$$n_s = \frac{1}{V_s} \left( \frac{n_{\text{wat}} m_{\text{wat}}}{\rho_{\text{wat}}} + \frac{n_{\text{sulf}} m_{\text{sulf}}}{\rho_{\text{sulf}}} \right), \quad (4)$$

where  $n_{\text{wat}}$ ,  $m_{\text{wat}}$ , and  $\rho_{\text{wat}}$  are the water complex refractive index, aerosol water mass (kg), and water density ( $\text{kg m}^{-3}$ ); and  $n_{\text{sulf}}$ ,  $m_{\text{sulf}}$ , and  $\rho_{\text{sulf}}$  are the sulfate complex refractive index, sulfate mass (kg), and sulfate density ( $\text{kg m}^{-3}$ ). The soluble aerosol volume across modes ( $V_s$ ) is calculated from the volume mean radius ( $R_{\text{eff}}$ ):

$$V_s = \frac{4}{3} \pi R_{\text{eff}}^3, \quad (5)$$

Sulfate and water refractive indices are taken from the literature (Hess et al., 1998; Hale and Query, 1973). The sulfate refractive index reported in Hess et al. (1998) ( $n_{\text{Hess}}$ ) is at 0% relative humidity and corresponds to aerosol that is approximately 25% water / 75%  $\text{H}_2\text{SO}_4$  by mass (Seinfeld and Pandis, 2006). A similar solution ratio has been used in other remote sensing applications (see Section 3), and for our in-situ and remote sensing datasets we assume stratospheric aerosol optical properties are reasonably represented by  $n_{\text{Hess}}$ . The pure sulfate refractive index used in E3SM/CESM ( $n_{\text{sulf}}$ ) is calculated from water and  $n_{\text{Hess}}$  following methods in Ghan and Zaveri (2007) and is nearly identical to  $n_{\text{Hess}}$ , albeit reported at a lower wavelength resolution. Given their similarity and the higher wavelength resolution in  $n_{\text{Hess}}$ , we use  $n_{\text{Hess}}$  instead of E3SM's pure sulfate refractive index for model single particle scattering calculations.

## 6. Code Availability

All model code is publicly available on Github. The E3SMv2-SPA model code can be found at <https://github.com/sandialabs/CLDERA-E3SM>; E3SMv2-PA model code can be found at <https://github.com/E3SM->



Project/E3SM; CESM2-WACCM data can be found at <https://github.com/ESCOMP/CESM>. Plotting and processing scripts used in the analyses of this paper can be found on Figshare (<https://doi.org/10.6084/m9.figshare.24844815.v1>).

## 7. Data Availability

The WOPC balloon data are available in the data repository at the University of Wyoming Libraries (Deshler and Kalnajs, 785 2022). AVHRR data can be found at doi:10.25921/w3zj-4y48 (Xuepeng and NOAA CDR Progra). For SAGE, AVHRR, and some of the postprocessed WOPC data, see <https://doi.org/10.7298/mm1s-ae98> (Quaglia et al., 2022). The original 2-degree volcanic emissions file (VolcanEESMv3.11\_SO2\_850-2016\_Mscale\_Zreduc\_2deg\_c180812.nc) can be retrieved from <https://svn-ccsm-inputdata.cgd.ucar.edu/trunk/inputdata/atm/cam/chem/stratvolc/>. The raw model output data used in this paper has been archived on the DOE NERSC resources and is available upon request.

## 790 8. Author Contribution

**Conceptualization:** Hunter Brown, Benjamin Wagman, Diana Bull, Kara Peterson, Xiaohong Liu, Ziming Ke; **Data curation:** Hunter Brown, Benjamin Wagman, Benjamin Hillman, Ziming Ke, Lin Lin; **Formal analysis:** Hunter Brown; **Funding acquisition:** Diana Bull, Kara Peterson; **Investigation:** Hunter Brown; **Methodology:** Hunter Brown, Benjamin Wagman, Diana Bull, Xiaohong Liu, Ziming Ke, Lin Lin; **Project administration:** Diana Bull, Kara Peterson; **Resources:** Diana Bull, 795 Kara Peterson; **Software:** Hunter Brown, Ben Hillman, Ziming Ke, Lin Lin; **Supervision:** Diana Bull, Kara Peterson, Benjamin Wagman; **Validation:** Hunter Brown; **Visualization:** Hunter Brown; **Writing – original draft preparation:** Hunter Brown, Diana Bull, Kara Peterson, Benjamin Wagman; **Writing – review & editing:** All authors.

## 9. Competing interests

At least one of the (co-)authors is a member of the editorial board of Geoscientific Model Development.

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