

## Response to Reviewer R1

The authors appreciated the comments of the reviewer and have substantially re-written the paper to address them.

The paper deals with limits of predictability associated with uncertainties in simulated diabatic heating over the Indian ocean during the MJO. This is investigated by re-running a set of 60-day long ensemble forecasts assuming perfect initial conditions. In other words, the study addresses a component of predictability associated with applied model errors. The error was simulated by stochastic perturbations of the tendencies due to physics geographically limited to the Indian Ocean region 50E-120E and 20N-20S. Their effects on predictability are studied in the ensemble spread of the vertically integrated diabatic heating, of Rossby wave source, of the heat flux and of the vorticity field at 200 hPa. The authors report predictability limits, defined as 0.5 of the saturation value of the ensemble variance, to be between 2 and 3 weeks.

We have a subtly different interpretation of our experiments, and have tried to clarify this in the revision. Our interpretation is not that our study addresses that component due to intrinsic variability in the model (“internal error”) which mimics the uncertainty of heating even with a given MJO episode. The stochastic parameterization scheme which is used (SPPT) is considered an integral part of the IFS, and other research (i.e. Selz, 2019) supports the notion that the use of such schemes gives more realistic estimates of intrinsic predictability. The novelty of our experiments is that we suppress SPPT outside the Indo-Pacific region, and that is the only change we make. See the following changes where we try to make this more clear:

lines 50-52,

“However, another limitation to predictability is the variability of heating among different observed episodes of a given phase, and the intermittency of heating and other sub-grid scale processes in space and time even within a particular episode.”

line 64

“The SPPT alters the instantaneous tendencies of the temperature, specific humidity and horizontal wind components due to sub-grid scale physics processes by scaling these tendencies up or down in a stochastic manner. It is considered a standard component of the IFS.”

lines 344-345

“The suite of ensemble reforecast experiments presented here was explicitly designed to gauge the effect of the intrinsic uncertainty of sub-grid motions on the response to the MJO in phases 2 and 3.”

Summary:

The authors address an important problem but the paper in its current shape does not provide significant new insights on the teleconnections associated with the MJO or associated predictability limits.

While a number of papers have discussed the role of initial condition uncertainty (see next response), but the role of the uncertainty in heating has not been discussed as much. Lines 53-58:

“The dynamical character of the uncertainty in the response to this intermittency has been studied by Kosovelj et al. (2019) in a low resolution global model using idealized stochastic parameterizations. Similarly, the response to model systematic error in Indian Ocean temperatures was addressed by Zhao et al. (2023). The purpose of this paper is to extend the work of Kosovelj et al. (2019) to address the intrinsic limits of potential predictability due to the intermittency of heating and other sub-grid scale processes in a high resolution operational forecast model setting.”

In the original manuscript we did not emphasize the predictability of the Rossby wave source in the paper enough. the rate of spread of the Rossby wave source gives an indication of how the spread of the tropical heating is translated into the uncertainty in the mid-latitude response.

In the revised version, see lines 236 – 241:

“...the corresponding [predictability] times for the Rossby wave source, shown in the blue curves in Figure 4b, are considerably shorter than those for Q. This is especially true for the planetary waves (wavenumbers 1 - 3) for which the predictability time for the RWS is shorter than that for the heating ... This reflects the sensitivity of the RWS to the sub-tropical divergent flow and also the sub-tropical absolute vorticity (as expressed in equation 1). The predictability times for the RWS have not, to our knowledge, been shown before, and are an important result regarding the predictability of the extra-tropical circulation.”

lines 306-315:

“A naive interpretation of the tropical divergence as the main forcing function for the extra-tropics would indicate long-range predictability related to the MJO. However, the predictability times for the RossbyWave Source S are considerably shorter: the largest scales reach 0.50 (0.70) of saturation already at around 20 (30) days. This is understandable since S is influenced not only by tropical and subtropical divergence but also by the meridional gradient of the jet. One path by which mid-latitude and sub-tropical variability may affect the response to tropical forcing is by changing the effective source for that response.”

One way to revise the paper would be to compare the results with the operational seasonal forecasts which sample uncertainties in both initial conditions and model errors.

In the revised manuscript we differentiate between initial condition uncertainty and the uncertainty in sub-grid scale processes such as heating. The point is that error growth due to ICs can be tuned by changing the magnitude of the IC errors, where as the strength of the SPPT perturbations are an integral part of the IFS model used for operational forecasting.

Lines 68-74:

“All members of each ensemble use the identical initial conditions, so that the only cause of model uncertainty (also called error here) must be the noise in heating introduced by the SPPT in the tropical Indian Ocean. We should point out here that one might design a similar set of experiments in which only the initial conditions were perturbed. After all, any perturbation whatever to the system will quickly propagate and grow (Ansell et al., 2018). The question is how quickly such errors grow and saturate, and the answer depends on, among other factors, the amplitude of the initial errors. In fact, Zhang et al. (2019) use the this dependence of rate of growth on the magnitude of the initial condition error to estimate the predictability limit of mid-latitude weather.”

Detailed comments:

## 1. Assumptions

The authors argue (Lines 55-56) that they apply the "perfect model" assumption in their study of (Lines 46-47) "uncertainties in MJO heating, as witnessed by the variability in the details of heating among different episodes of a given phase." The statement should be re-written in line with what has been done (perfect initial-conditions and simulated model error). Please discuss what is meant by "variability in the details of MJO heating" and provide references.

We agree that the original manuscript was not clear. Here are our clarifications and additions:

lines 53-58

“However, another limitation to predictability is the variability of heating among different observed episodes of a given phase, and the intermittency of heating and other sub-grid scale processes in space and time even within a particular episode. The dynamical character of the uncertainty in the response to this intermittency has been studied by Kosovelj et al. (2019) in a low resolution global model using idealized stochastic parameterizations. Similarly, the response to model systematic error in Indian Ocean temperatures was addressed by Zhao et al. (2023).

The purpose of this paper is to extend the work of Kosovelj et al. (2019) to address the intrinsic limits of potential predictability due to the intermittency of heating and other sub-grid scale processes in a high resolution operational forecast model setting. In this paper we do not address model systematic error.”

lines 344-348.

“The suite of ensemble reforecast experiments presented here was explicitly designed to gauge the effect of the intrinsic uncertainty of sub-grid motions on the response to the MJO in phases 2 and 3. Each ensemble has all its members initialized identically during an observed MJO event, and differ from each other only in the realization of the stochastic parameterizations, applied only in the tropical Indo-Pacific region. Thus even though the errors (deviations within the ensemble) spread globally, they are ultimately due to the uncertainty in this region.”

## 2. Methodology

a) Lines 59-61: "the stochastic parametrization scheme (SPPT) described in Leutbecher et al. (2017) has been altered so that perturbations which affect (directly or indirectly) diabatic heating tendencies are confined to the tropical Indian Ocean region". Some details would be useful here,

such as the amplitude of perturbations compared to the signal. Different wording on what parts of the model physics have been perturbed is provided at different places in the paper, and it should be clarified.

We try to be more consistent in referring to the SPPT as perturbing (adding a stochastic component to) the sub-grid scale physics in general.

lines 60-65:

“...the stochastic parametrization scheme (SPPT) described in Leutbecher et al. (2017) has been altered so that perturbations which affect (directly or indirectly) diabatic heating tendencies are confined to the tropical Indian Ocean region. The SPPT alters the instantaneous tendencies of the temperature, specific humidity and horizontal wind components due to sub-grid scale physics processes by scaling these tendencies up or down in a stochastic manner. It is considered a standard component of the IFS. The physics processes include those due to turbulent diffusion and sub-grid orography, convection, cloudiness and precipitation, and radiation.”

lines 344-345 (in the Conclusions);

“The suite of ensemble reforecast experiments presented here was explicitly designed to gauge the effect of the intrinsic uncertainty of sub-grid motions on the response to the MJO in phases 2 and 3.”

b) The computation of the divergent horizontal wind in Eq. (1) should be explained.

Agreed. Lines 144 -162 give a fairly detailed outline of the mathematical steps used to compute the Rossby Wave Source. We make use of the transforms of a vector field to the corresponding divergence and curl in spherical harmonic space, and their inverse. These transforms are applied in several ways. We do not reproduce the details in this response; please see the lines indicated above.

c) What is the sensitivity of the results to the choice of the latitude belt used to compute S?

We discuss this sensitivity in lines 164-167: “...we consider the average source between 15N and 3N. The ... The results shown in this paper are robust to changes in the latitude band chosen, both to modest poleward displacement and to widening it by 5 degrees.”

The new Figure 2 shows the sensitivity of the mean and standard deviation of the Rossby Wave Source to different latitude bands. Also, the new Figure 4, showing the predictability times  $\tau$  of the RWS as a function on zonal wave number (for different definitions of  $\tau$ ) is remarkably robust to modest shifts and widening of the latitude band used.

d) Why is the integrated diabatic heating a good measure of the MJO predictability as compared for example with precipitation?

The focus of the paper is on how uncertainty in sub-grid scale processes in the tropics affect the extra-tropical predictability, and not on tropical precipitation. Within the tropics, it is the total heating that forces the atmospheric divergence and Rossby Wave Source.

lines 79-80:

“Since our goal is to document both the uncertainty in the tropical heating and the mid-latitude response in these experiments, we also consider the pathway by which the tropical heating forces extratropical Rossby waves.”

lines 348-350:

“These subsequent errors in the tropical diabatic heating, tropical upper-level divergence and Rossby Wave Source indicate the path towards mid-latitude uncertainty in the circulation response.”

However, we understand that in terms of extra-tropical spread of uncertainty, the heating is not the usual field used. We have replaced the maps showing the spread of ensemble spread with a single figure showing the evolution of ensemble spread of the meridional wind (new Figure 6).

Lines 253-255:

“To get a sense of how the errors spread geographically, we present maps of the ensemble spread of the meridional wind in Fig. 6. The choice of meridional wind was motivated by its close relationship with storm tracks and circumpolar wave guide (Branstator and Teng, 2017).”

e) Why is vorticity a good measure of the forecast error growth in the tropics (Figures 8-9)?

We have removed the maps of ensemble spread of vorticity in the tropics (see above response).

f) How is the estimated predictability limit sensitive to different choices of predictability time (Line 211) taken to be 50% of the saturation error?

Following Judt (2020) we show predictability times for different thresholds (50%, 70%, 90% of saturation) in the new Figure 4.

Lines 227-228:

“...we show the time  $\tau$  at which the error variance of  $Q$  reaches a fraction  $f$  of the variance of the external error for  $f = 0.50, 0.70$ , and  $0.90$ , as a function of zonal wavenumber in Figure 4a. ...”

g) How is the wavenumber analysis performed, is it spherical harmonics space?

The zonal wavenumber analysis is performed on latitude circles (i.e. as a function of longitude only) after the fields have been transformed back to the Gaussian grid.

h) The ENSO events are introduced in 2.2 and Figures 1-2, but little mentioned after 3.1.

We have added, in the Discussion (lines 289-293):

“The evolution of tropical heating shown in Figures 1 for El-Niño years shows less eastward propagation from the Indian Ocean compare to normal years, in line with the 290 findings of Liu et al. (2020), likely because less moisture is available over the Indian Ocean due to the ENSO convection in the central Pacific. Nevertheless, the average over all forecasts does show distinct eastward propagation for the first 10 days or so.”

### 3. Relation to previous work

Many studies addressed the response of tropical and extratropical circulation to MJO-like heating perturbations. I disagree with the authors' statement that (Lines 47-49) "the wealth of MJO teleconnection research discussed above has relied almost exclusively on the Wheeler-Hendon multivariate empirical orthogonal function framework (Wheeler and Hendon, 2004)." See for example <https://doi.org/10.1175/JAS-D-18-0203.1> and references herein. Similarly, there is a wealth of research in predictability associated with MJO that is missing in the introduction and discussion of the results.

We respond to this point and the next one together. In order to compare to previous work on evolution of error spectra due to initial condition or heating uncertainty, we have produced two new figures. Figure 3 shows the evolution of the spectra of mid-level heating, and Figure 5 shows the evolution of the error spectra of kinetic energy for various latitude bands. We present and discuss these new figures in the context of previous work.

In the Introduction:

lines 52-57:

“The dynamical character of the uncertainty in the response to this intermittency has been studied by Kosovelj et al. (2019) in a low resolution global model using idealized stochastic parameterizations. Similarly, the response to model systematic error in Indian Ocean temperatures was addressed by Zhao et al. (2023).

The purpose of this paper is to extend the work of Kosovelj et al. (2019) to address the intrinsic limits of potential predictability due to the intermittency of heating and other sub-grid scale processes in a high resolution operational forecast model setting.”

lines 69-74:

“We should point out here that one might design a similar set of experiments in which only the initial conditions were perturbed. After all, any perturbation whatever to the system will quickly propagate and grow (Ansell et al., 2018). The question is how quickly such errors grow and saturate, and the answer depends on, among other factors, the amplitude of the initial errors. In fact, Zhang et al. (2019) use the this dependence of rate of growth on the magnitude of the initial condition error to estimate the predictability limit of mid-latitude weather. Judt (2018) study the dependence on the rate of initial condition error growth on region (tropical, mid-latitude and high-latitude) for short simulations using a global storm-resolving model.”

In the Results Section:

lines 219-226:

“In order to investigate the scale-dependence of the uncertainty of the heating evolution, we calculated the zonal wavenumber spectra of the internal error variance of  $Q_{mid}$  (850-400 hPa ) averaged over 2-day blocks and over the tropical belt 15S - 15N. Figure 3 shows the spectra for the blocks ending on days 2, 4, 6, 10, 20, 40 and 60. The latter is indicated by the red line, and gives a simple measure of saturation. By day 2, the spectrum is relatively flat down to length scales of about 3000 km, consistent with the variance being forced by perturbations in a narrow longitude range. As time progresses, the variance increases without dramatic change in shape until after day 20, when the larger scales (wavelengths greater than about 3500 km) grow more rapidly than the smaller scales, leading to a steeper variance spectrum by 60 days. If we define the predictability time ...”

lines 242-252:

“The forecast evolution of the spectra of the internal error variance of kinetic energy (KE) is presented in Figure 5 for different latitude bands. In the tropics (15S-15N; Figure 5d), the spectra grow without much change in shape between forecast days 3 and 10, but between days 10 and 20 for smaller scales approach saturation much more quickly than do the larger scales. Wave lengths shorter than about 3500 km are already saturated by day 20, while the larger scale error variance continues to grow beyond day 40. The times for which the spectra are shown in Fig. 5d (3, 5, 10, 20, 40 and 60 days) are the same as in the other panels (Figs. 5a - 5c), which show latitude bands of 25N -35N, 45N -55N and 65N -75N, respectively. As one goes to higher latitudes (i.e. from panels (c) to (b) to (a)), the error variance curve for days 3 and 5 continuously moves to lower values, indicating the time it takes for the error to propagate poleward from the tropics. However by day 40, the curves are at about the same level for all latitudes except the 65N -75N band, where the error is less than at other latitudes. In fact the saturation error (estimated by the red curve at day 60) is also less.”

In the Discussion:

Lines 314-329 (in the Discussion):

“The slopes of the saturation (or background) kinetic energy spectra presented in Figure 5 are compared to those corresponding to a dimensional wavenumber dependence of  $k^{-3}$  and to a dependence of  $k^{-5/3}$  as indicated by dashed and solid lines in the Figure... A slope corresponding to a  $k^{-3}$  dependence is evidence of the dominant of rotational flow, while a  $k^{-5/3}$  dependence is associated with the dominance of convection and gravity waves, and in general divergent flow (Charney, 1971; Sun et al., 2017; Zagar et al., 2017; Li et al., 2023).”

“Our tropical saturation spectrum, showing a  $k^{-3}$  dependence, is in distinct contrast to that of Judt (1988) (their Figure 5), which shows a slope roughly corresponding to a  $k^{-5/3}$  dependence, indicative of the dominance of convection and divergent flow. In mid-latitudes, the background spectrum of J also shows a transition from a  $k^{-3}$  to a  $k^{-5/3}$  dependence at several hundred km., a transition our results are unable to resolve. These differences are due to the model resolution and dynamics, as J uses 4-km horizontal resolution, storm-resolving simulations without convective parameterizations, while the IFS model used here has a resolution of 36-km and makes use of parameterizations for unresolved processes. Similar to the spectra shown in Zhang et al. (2019) (their Figure 6), the mid-latitude error growth is rapid between days 5 and 10, while the continued error growth for the largest scales at later times is similar to that reported in Selz (2019) in their Figure 4, noting that the spectra are plotted differently”

#### 4. Results

This paper, like several earlier studies of predictability, finds a predictability scale of 2-3 weeks and that longer scale have longer predictability. In the present study, the predictability limit is due to perturbations in model physics. Previous studies such as <https://doi.org/10.1175/JAS-D-19-0116.1> find similar intrinsic predictability to be due to small perturbation in initial conditions (perfect model assumption). I wish the authors discussion their predictability results in comparison to other studies of predictability in the tropics and globally.

How does the growth of spread in selected variables compare with the scale-dependent circulation response to heating perturbations in <https://doi.org/10.1175/JAS-D-18-0203.1> (their figures 7-8)?

Please see detailed response to previous query.

Overall it is unclear why the ensemble spread of diabatic heating is a good measure of predictability limits, rather than prognostic variables of circulation and/or precipitation. For example, how is the amplitude of the ensemble spread in diabatic heating related to the predictability of precipitation? Could the results be coupled with the precipitation validation in the ECMWF model forecasts (e.g. <https://doi.org/10.1029/2020GL091022>)?

In the tropics, the spread of diabatic heating is presented to document just how variable heating is, and to show the forcing for the extra-tropical spread in circulation. Whether the tropical precipitation in this model is realistic is a different, albeit interesting question.

## Replies to R2

We appreciate the comments of the reviewer, and have done our best to address them. The paper has been rather extensively modified.

1. While it is true that only the model physics over the Indian Ocean are perturbed, the effect of chaos seeding (Ansell et al. 2018) quickly spreads the error over the whole globe in non-physical ways. This likely means that the results are indistinguishable from results that would have been obtained if SPPT perturbations were added to the entire tropical belt or to another location that is convectively active, such as the Amazon. I suggest testing this out.

This is a valid point – any perturbation will rapidly spread due to variety of reasons, some numerical and some intrinsic to the system. We want to point out that the stochastic parameterization scheme which is used (SPPT) is considered an integral part of the IFS, and other research (i.e. Selz, 2019) supports the notion that the use of such schemes gives more realistic estimates of intrinsic predictability. Thus the effects of restricting the SPPT to the tropical Indo-Pacific region may be confined to the early growth of errors.

### In the Introduction

lines 75-79:

“Another question regarding our experimental design is whether localizing the application of the SPPT to the tropical Indian Ocean is necessary, since following the argument of Ansell et al. (2018), perturbations in any region will quickly propagate to the Indian Ocean region. The only way to answer this question is to re-run the same set of experiments with SPPT applied globally, which is the subject of future research. We will return to this question in the Discussion section.”

### In the Discussion:

Lines 294-300:

“The evolution of the average of the ensemble spread in vertically integrated heating ( $Q$ ) shown in Figure 1d shows clearly that the within-ensemble variability induced by the application of the regionally confined SPPT remains mostly confined to that region (50°-120°E) for the first 10 days or so. This is also true for the tropical meridional wind spread (Figure 6) for the first 6 days. This suggests that the evolution of the tropical heating and circulation uncertainties would be different had the SPPT been applied throughout the tropical belt. Whether this difference would strongly affect the growth of uncertainty in the extra-tropics is hard to assess directly from these experiments. This question awaits future research.”

2. It would be interesting to see how different the result would be when “butterfly seeding” were used, i.e., tiny initial perturbations of the initial conditions everywhere on the globe, as in Judt (2018) or Zhang et al (2019). I recommend running an additional ensemble with this kind of perturbation and comparing the results with the ones obtained so far (this additional experiment would also be more in line with intrinsic predictability, which usually addresses predictability limits arising due to miniscule initial condition uncertainty).

While we agree that this would be an interesting comparison, there is an arbitrariness in perturbing initial conditions. In fact the whole purpose of the cited paper of Zhang et al. is to compare the error growth due to current, operational uncertainty in the initial conditions (taken from multiple reanalyses) to a hypothetical “perfect” scenario that is implemented by arbitrarily reducing the initial condition uncertainties by a factor of 100. Such a study could be implemented in the ECMWF modeling system since there is in place a procedure to perturb the initial conditions (although we did not use this in the simulations). As above, this computer-intensive proposal is a goal for a future project.

Lines 66-74:

“All members of each ensemble use the identical initial conditions, so that the only cause of model uncertainty (also called error here) must be the noise in heating introduced by the SPPT in the tropical Indian Ocean. We should point out here that one might design a similar set of experiments in which only the initial conditions were perturbed. After



all, any perturbation whatever to the system will quickly propagate and grow (Ansell et al., 2018). The question is how quickly such errors grow and saturate, and the answer depends on, among other factors, the amplitude of the initial errors. In fact, Zhang et al. (2019) use the this dependence of rate of growth on the magnitude of the initial condition error to estimate the predictability limit of mid-latitude weather. Judt (2020) study the dependence on the rate of initial condition error growth on region (tropical, mid-latitude and high-latitude) for short simulations using a global storm-resolving model.

3. Is it necessary to discuss the Nov and Jan initialization experiments separately? In my opinion no, as the differences between Nov and Jan events are not large enough to warrant the extra work for the reader to keep track of two sets of results. I therefore recommend combining all experiments into one “grand experiment”. This shouldn’t affect the conclusions.

This is a valid point for most of the diagnostics. It is only for the figures relating to the role of the stratosphere that we have reasons to discriminate the Jan. and Nov. initial conditions. This is because the polar vortex may not be fully formed in November, so that the “stratospheric pathway” towards uncertainty growth will have a different time scale for the Nov and Jan runs.

Except for Figures 7, 8 and 9 relating to the stratosphere, all other diagnostics have the Nov and Jan experiments combined. Note that we have modified some of figures and added a number of new ones in order to address the concerns of another reviewer.

4. It looks like Fig. 3 is not referenced in the text. Furthermore, I am not sure why the “Rossby wave source” is analyzed at all. I suggest removing this analysis or better motivating it.

One of the new results in this paper is the predictability time of the Rossby wave source, so we have better motivated it:

In the Introduction:

lines 79-83:

“Since our goal is to document both the uncertainty in the tropical heating and the mid-latitude response in these experiments, we also consider the pathway by which the tropical heating forces extratropical Rossby waves. Although the MJO-related tropical heating is expected to force a corresponding signal in upper tropospheric divergence, this signal generally occurs within an easterly background wind, where stationary Rossby waves are not expected to propagate. Sardeshmukh and Hoskins (1988) derive a more complete formulation of the source of barotropic Rossby waves (hereafter Rossby wave source...)”

In the Results discussing the new Figure 4, which show wavenumber dependent predictability times for the tropical heating, upper-level tropical divergence and Rossby Wave Source. The times are shown to reach 50%, 70% and 90% of saturation.

lines 233-241:

“The predictability times for the T21 representation of the upper-level (200 hPa) divergence are shown in Figure 4b in the blue curves. These times are notably longer than for the vertically integrated heating. For example, the divergence [predictability time] corresponding to [70% of saturation] is greater than 35 days for the largest scales, compared to 20 days for the heating. While this might be taken to indicate high predictability for the extra-tropical response, the corresponding times for the Rossby wave source, shown in the blue curves in Figure 4b, are considerably shorter than those for Q. This is especially true for the planetary waves (wavenumbers 1 - 3) for which the predictability time for the RWS is shorter than that for the heating by about 8 days [to reach 50% saturation] and by about 10 - 20 days for  $f = 0:70$ . This reflects the sensitivity of the RWS to the sub-tropical divergent flow and also the sub-tropical absolute vorticity (as expressed in equation 1). The predictability times for the RWS have not, to our knowledge, been shown before, and are an important result regarding the predictability of the extra-tropical circulation.”

In the Conclusions:

lines 345-350:

“Each ensemble has all its members initialized identically during an observed MJO event, and differ from each other only in the realization of the stochastic parameterizations, applied only in the tropical Indo-Pacific region. Thus even

though the errors (deviations within the ensemble) spread globally, they are ultimately due to the uncertainty in this region. These subsequent errors in the tropical diabatic heating, tropical upper-level divergence and Rossby Wave Source indicate the path towards mid-latitude uncertainty in the circulation response.”

## 5. Section 3.4 seems to be lacking a conclusion, or at least I’m left with this impression.

In the Discussion:

Lines 340-352:

“The growth of uncertainty in the stratospheric circulation, as seen in Figure 7, is forced by the upward propagation of the planetary wave meridional flux of sensible heat (which is the dominant term in the vertical component of the Eliassen-Palm flux), shown in Figure 8. This uncertainty then propagates downward into the upper and middle troposphere. While most of the upper troposphere sensible heat flux is due to planetary wave disturbances in the Pacific, its uncertainty in the North Atlantic and Asian sectors are also large, especially for the Jan. experiments (Figure 9). This downward propagation is potentially linked to wave-mean flow interaction which acts to bring anomalies in e.g. wind and temperature to the lower stratosphere.

The planetary wave error in the upper troposphere (300hPa) for Nov. reaches a maximum 20 days earlier than does the error at 50hPa, hinting at a tropospheric forcing of the stratospheric spread. The stratospheric descent of error seen in Figure 7 occurs towards the end of the experiments, consistent with the tropospheric forcing uncertainty being modulated by the stratospheric circulation (Domeisen et al., 2020b). This descent is seen about 10 days later in the reforecast period for the Nov. experiments than for the Jan. experiments. This is likely due to the lack of a fully formed stratospheric vortex during November, so that the establishment of a wave guide for vertically propagating (It was not possible to verify this since data were retained only up to 50hPa.)”

In the Conclusions:

lines 374-375:

“The role of the stratosphere in amplifying uncertainty is generally confined to the latter part of the 60-day reforecasts, after the ensemble spread in upper-tropospheric heat flux has affected levels above 50 hPa [Figures 7 and 8].”

Minor Comments:

1. L.45: How does the presence of baroclinic instability limit the predictability of MJO teleconnections? Through error growth associated with baroclinic instability?

Yes you are correct. But we have dropped the reference to baroclinic instability since it was apparently confusion.

2. Figs. 1-3 and 6-9 imply that the runs are 30 days long, while the other figures show the entire 60 day time period. Why are only the first 30 days shown in the Hovmöller plots? Maybe nothing interesting happens after 30 days, but then it should be indicated somewhere so the reader doesn’t end up confused whether or not the experiments are 30 or 60 days long.

In the revised paper, Figures 1, 2 and 6 show only the first 30 days.

Reference to Figure 1:

lines 186-187

“The daily averaged evolution of vertically integrated diabatic heating anomaly is shown averaged for the first 30 days of the 60-day experiments in Figure 1...”

Reference to Figure 2 (where it is clear that the signal has dissipated by day 30):

lines 198-199

“figure 2 shows the evolution of both the ensemble average of the latitudinally averaged RWS, averaged over all experiments, for the first 30 forecast days.”

Reference to Figure 6:

lines 253-258:

“...we present maps of the ensemble spread of the meridional wind in



Fig. 6. The choice of meridional wind was motivated by its close relationship with storm tracks and circumpolar wave guides (Branstator and Teng, 2017). Much of the tropics outside of the Indian Ocean region is nearly error free even at day 6. By day 10, substantial error has already appeared in the extra-tropics, particularly in the storm-track regions, and by day 16 the extra-tropical spread has almost saturated. By day 30 the spread in the extra-tropics has essentially reached its saturation value, since it doesn't increase for longer lead times (not shown)."

3. L.102: Just curious, why are you not using ERA5 to initialize the ensembles?

We purposely used the model configuration that ECMWF uses for its monthly forecasts, since that model has been well calibrated vis-à-vis the MJO. That configuration is set up to use ERA-Interim for initial conditions.

4. L. 155-177: I don't think the description of the figures is necessarily wrong, but I do see a lot of noise in the Hovmöllers and not so much of the described

Points 4 and 5 are addressed together, since they both relate to Figure 1, which now shows Nov and Jan experiments combined. The discussion has been clarified:

lines 186-193:

"The daily averaged evolution of vertically integrated diabatic heating anomaly is shown averaged for the first 30 days of the 60-day experiments in Figure 1c. The heating has been averaged over the tropical band (15S-15N), over all ensemble members and over all experiments. The eastward propagation of positive heating anomalies near longitude 90E can be seen for about 8 days, along with robust westward propagation of heating anomalies that appear after 4 days in the central Pacific. The ensemble spread of the heating (also averaged 190 over the tropical band and all experiments) is shown in Figure 1d. The dominant influence of the SPPT generated perturbations over the Indian Ocean sector is clear, leading to the largest ensemble spread in this sector.

In order to gauge the degree of influence of ENSO on the heating evolution, we also show the tropical heating anomalies separately for the six warm ENSO events (Nov 1986, Nov 1987, Nov 2002, Nov 2013, Jan 1987, Jan 2010) in Figure 1a and for the remaining seven experiments in Figure 1b. The warm events show the establishment of strong tropical heating in the central Pacific after day 15 (as expected), while both sets of years show robust westward propagation of the heating from the central Pacific. The initial eastward propagation of the Indian Ocean anomalies is somewhat delayed in the warm event years in comparison to the neutral years."

5. Fig. 1 (and Fig. 2): The evolution of the standard deviation in panels (d) shows very little propagation with the maximum being anchored between 70 and 100 deg E, unlike the heating amplitude. Is this because of the continuous perturbation in this region?

(see reply above)

6. L. 188: "In the Indian Ocean the two fields are comparable." I disagree, there seems to be more red in Fig. 1d than in Fig. A1 over the Indian Ocean.

We have refined the discussion.

lines 212-218:

"In order to determine whether the strength of the stochastic perturbations (reflected in the magnitude of the ensemble spread in the Indian Ocean region) is reasonable, we also computed the inter-annual standard deviation  $\sigma_{IA}$  of the tropical Q from the ERA5 reanalysis over the eight years corresponding to the Nov. Experiment for the first 30 days. Figure A1 in the Appendix shows the daily evolution of the standard deviation with the same scale as in Figure 1d. The ERA5  $\sigma_{IA}$  is largely confined to the same regions as the model spread: the Indian Ocean region and the west-central Pacific. In the Indian Ocean sector, the model spread in heating has somewhat lower maximum values than  $\sigma_{IA}$ , but extends over a wider area. The model spread is notably less than  $\sigma_{IA}$  over the Pacific up to forecast day 30."

7. L. 211: Where does the 0.5 threshold come from? It seems arbitrary.

Both reviewers have made this point, and we have changed the analysis to include plots of wavenumber – dependent predictability for tropical heating, upper-level divergence and Rossby Wave Source for multiple thresholds: 0.50, 0.70 and 0.90. See the new Figure 4.

lines 227-241:

“To make this more precise, we show the time  $\tau$  at which the error variance of Q reaches a fraction  $f_r$  of the variance of the external error for  $f_r = 0.50$ ,  $0.70$ , and  $0.90$ , as a function of zonal wavenumber in Figure 4a. The red curves give the results for Q in the solid, dashed and dotted lines, respectively. Prior to calculating the error variances for this plot, we have truncated Q (and the other fields to be shown) to a spherical harmonic T21 representation in order to eliminate excessive noise.  $\tau$  increases with zonal scale (decreasing wavenumber) for all choices of  $f_r$ , but this is particularly noticeable for  $f_r = 0.70$  and  $0.90$ . In fact the limit of 0.90 of the external error is never reached for zonal wavenumbers 1 and 2.

The predictability times for the T21 representation of the upper-level (200 hPa) divergence are shown in Figure 4b in the blue curves. These times are notably longer than for the vertically integrated heating. For example, the divergence  $\tau$  corresponding to  $f_r = 0.70$  is greater than 35 days for the largest scales, compared to 20 days for the heating  $\tau$ . While this might be taken to indicate high predictability for the extra-tropical response, the corresponding times for the Rossby wave source, shown in the blue curves in Figure 4b, are considerably shorter than those for Q. This is especially true for the planetary waves (wavenumbers 1 - 3) for which the predictability time for the RWS is shorter than that for the heating by about 8 days for

$f_r = 0.50$ , and by about 10 - 20 days for  $f_r = 0.70$ . This reflects the sensitivity of the RWS to the sub-tropical divergent flow and also the sub-tropical absolute vorticity (as expressed in equation 1). The predictability times for the RWS have not, to our knowledge, been shown before, and are an important result regarding the predictability of the extra-tropical circulation.”

8. L. 298: “small scales” is relative, wavenumber 21 is “large scale” from the view of a synoptic/mesoscale meteorologist.

Yes that was a bad choice of wording; we have omitted this phrase.