



1 **Understanding variations in downwelling longwave radiation**
2 **using Brutsaert's equation**

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11

12 **Abstract**

13 A dominant term in the surface energy balance and central to global warming is downwelling longwave
14 radiation (R_{ld}). It is influenced by radiative properties of the atmospheric column, in particular by
15 greenhouse gases, water vapour, clouds and differences in atmospheric heat storage. We use the semi-
16 empirical equation derived by Brutsaert (1975) to identify the leading terms responsible for the spatio-
17 temporal climatological variations in R_{ld} . This equation requires only near-surface observations of air
18 temperature and humidity. We first evaluated this equation and its extension by Crawford and Duchon
19 (1999) with observations from FLUXNET, the NASA-CERES dataset, and the ERA5 reanalysis. We found
20 a strong agreement with r^2 ranging from 0.87 to 0.99 across the datasets for clear-sky and all-sky conditions.
21 We then used the equations to show that diurnal and seasonal variations in R_{ld} are predominantly controlled
22 by changes in atmospheric heat storage. Variations in the emissivity of the atmosphere form a secondary
23 contribution to the variation in R_{ld} , and are mainly controlled by anomalies in cloud cover. We also found
24 that as aridity increases, the contributions from changes in emissivity and atmospheric heat storage tend to
25 offset each other (-40 W m^{-2} and $20\text{-}30 \text{ W m}^{-2}$, respectively), explaining the relatively small decrease in R_{ld}
26 with aridity ($-(10\text{-}20) \text{ W m}^{-2}$). These equations thus provide a solid physical basis for understanding the
27 spatio-temporal variability of surface downwelling longwave radiation. This should help to better
28 understand and interpret climatological changes, such as those associated with extreme events and global
29 warming.

30

31



32 1 Introduction

33 In the global mean surface energy budget, downward longwave radiation (R_{ld}) is the dominant term (333
34 W/m^2), contributing more than twice as much energy as absorbed solar radiation (161 W/m^2) (Trenberth et
35 al. 2009). This dominance holds over all regions in the climatological mean, although there are some clear
36 variations in space and time (Figure 1). It is central to global warming, reflecting the greenhouse effect of
37 the atmosphere (Held and Soden 2000), and its variations have been suggested to be the main contributor
38 to some regional warming amplifications, such as in the Arctic (Lee et al. 2017) and the Tibetan Plateau
39 (Su et al. 2017). Therefore, it is important to understand the main sources of variations in this surface energy
40 balance term, which can be seen in Figure 1.

41 The flux of downwelling longwave radiation is influenced by the radiative properties of the entire
42 atmospheric column, i.e., water vapour, clouds and greenhouse gases, but also by the heat stored in the
43 atmosphere, i.e., the temperature at which radiation is emitted back to the surface. To obtain an estimate of
44 this flux, Brutsaert (1975) used functional expressions for the typical temperature and humidity profiles of
45 the lower troposphere together with radiative transfer equations and semiempirical relationships of the
46 absorptivity by water vapor, integrated these vertically, and expressed the resulting flux R_{ld} in terms of near-
47 surface air temperature and water vapour pressure for clear-sky conditions. He thereby derived a semi-
48 empirical equation for R_{ld} for an effective clear sky emissivity (ε_{cs}) and the corresponding flux of
49 downwelling longwave radiation ($R_{ld,cs}$):

$$\varepsilon_{cs} = 1.24(e_a/T_a)^{1/7}, \quad (1)$$

$$R_{ld,cs} = \varepsilon_{cs}\sigma T_a^4. \quad (2)$$

50 where σ is Stefan–Boltzmann constant ($\sigma = 5.67 \cdot 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$), e_a is the water vapor pressure and T_a is the
51 2m air temperature. The latter two meteorological variables can easily be obtained or inferred from weather
52 stations, so that the downwelling flux of longwave radiation can be estimated from weather station
53 observations.

54 This equation was later extended to all-sky conditions that include the effects of cloud cover, among which
55 Crawford and Duchon (1999) is a common extension (Alados et al. 2012; Duarte et al. 2006; Flerchinger
56 et al. 2009). This extension diagnoses cloud cover fraction (f_c) as the fraction of incoming solar radiation
57 at the surface (R_s) in relation to the potential solar radiation ($R_{s,pot}$), that is, the incoming flux at the top of
58 the atmosphere. The emissivity for all-sky conditions, ε , is then calculated as the mix of the emissivities of
59 clear-sky conditions (Eqn. (1), weighted by the cloud-free proportion, $1 - f_c$) and clouds with an emissivity
60 of $\varepsilon_c = 1$ (weighted by the cloud fraction f_c). Using this emissivity, the estimation of downwelling
61 longwave radiation is then done by

$$f_c = 1 - R_s/R_{s,pot}, \quad (3)$$

$$\varepsilon = f_c + (1 - f_c)\varepsilon_{cs}, \quad (4)$$

$$R_{ld} = \varepsilon\sigma T_a^4 \quad (5)$$

62 Previous studies have already verified this estimate to have a very good agreement with site measurements
63 (Duarte et al. 2006; Hatfield et al. 1983), especially when the temperature is higher than 0°C (Aase and Idso
64 1978; Satterlund 1979). Other studies have worked to calibrate and modify this estimate further to different
65 regions (Malek 1997; Sridhar and Elliott 2002).

66 This expression for downwelling longwave radiation R_{ld} given by Eqn. (5) allows us to quantify the different
67 contributions by cloud cover, f_c , water vapor concentrations, e_a (as a measure of the total water vapor content
68 of the atmospheric column), and air temperature, T_a (as a proxy for the total atmospheric heat storage within
69 the column). With this, we can then attribute variations in R_{ld} to their physical causes.



70 Here, our aim is to first evaluate this estimate for downwelling longwave radiation with current global
71 datasets at the continental scale. These variations are illustrated using the NASA-CERES (EBAF 4.1)
72 dataset (Loeb et al., 2018; Kato et al., 2018, NASA/LARC/SD/ASDC 2017) and the NASA-CERES
73 Syn1deg dataset (Doelling et al., 2013, 2016) in Figure 1 and are compared to variations in solar radiation.
74 It can be seen that the climatological distribution of R_{ld} is mostly associated with latitudes, while also
75 presenting some zonal variations, e.g., across western and eastern North America. In comparison, the
76 seasonal cycle of R_{ld} is less determined by latitudes (Fig. 1b). It has a larger magnitude over land than over
77 oceans, over arid regions than humid regions, and over cold regions more than over warm ones. Although
78 studies have revealed a close correlation between the variation of R_{ld} and other factors like air temperature,
79 water vapor, and CO_2 concentration (Wang and Liang 2009; Wei et al. 2021), here we go beyond
80 correlations and rather attribute these variations to the different terms in Eqns. (1)-(5) that represent
81 different radiative properties affecting R_{ld} .

82 To figure out the dominant driver for these spatiotemporal variations, we decompose changes in R_{ld} into its
83 components: cloud cover, f_c , heat storage changes of atmosphere as reflected by 2m air temperature, T_a ,
84 and air humidity, e_a , by performing the differentiation of these equations. We show that heat storage
85 changes predominantly shape the diurnal range and seasonal cycle of R_{ld} , while cloud cover variations play
86 a second role in most cases. In addition, the temporal variations of R_{ld} are less over the ocean than over
87 land, and less during winter than summer. On the other hand, the spatial variations of R_{ld} from arid to humid
88 regions is relatively small, which we will show is due to a compensating effect of corresponding changes
89 in atmospheric emissivity and heat storage.

90 Our paper is organized as follows: After briefly describing the datasets used in our evaluation in Section 2,
91 we first the estimate of R_{ld} from these equations at the global scale, using multiple datasets in Section 3.1.
92 After showing that the annual-mean and large-scale variations are well captured, we then use the equations
93 to decompose the temporal variations of R_{ld} in terms of its mean spatial and temporal variations and relate
94 these to their causes in Section 3.2. The spatial variations of R_{ld} are then further discussed in Section 3.3 in
95 terms of its relationship with aridity. We then close with a brief summary and broader implications.

96 2 Datasets

97 To test R_{ld} estimates, we use FLUXNET observations (Pastorello et al. 2020, half-hourly values, 189 sites,
98 see Table S1 for details), the NASA CERES satellite-based radiation dataset (Doelling et al., 2013, 2016,
99 monthly means, covering years 2001 to 2018), and the ERA5 reanalysis dataset (Hersbach et al. 2018,
100 monthly means, covering years 1979 to 2021).

101 For each dataset, T_a , e_a , and f_c are needed as inputs for Eqs. (1)-(5), while R_{ld} data is used for the
102 comparison. For FLUXNET and ERA5, water vapor pressure, e_a , is not directly given. It is calculated
103 from the water vapor deficit (VPD, FLUXNET) or dewpoint temperature (T_{dew} , ERA5) using Monteith and
104 Unsworth (2008):

$$e_a = 6.1079 \times \exp(17.269T_{dew}/(237.3 + T_{dew})), \quad (6)$$

$$e_a = 6.1079 \times \exp(17.269T_a/(237.3 + T_a)) - VPD, \quad (7)$$

105 Since the NASA-CERES dataset includes cloud cover, f_c , this is used directly in the estimation instead of
106 using Eq. (3), together with e_a from ERA5 and T_a from the CPC Global Unified Temperature dataset (CPC
107 Global Unified Temperature).

108 For the analysis of the spatial variations of R_{ld} along water availability, we use the aridity index ($AI = \frac{R}{LP}$)
109 (Budyko 1958; UNCOD 1977). This index is calculated using the mean annual net radiation (R) taken from
110 the NASA-CERES dataset, the mean annual net precipitation (P) taken from the CPC Global Unified
111 Gauge-Based Analysis of Daily Precipitation data (Chen et al. 2008 and Xie et al. 2007, CPC Global Unified



112 Gauge-Based Analysis of Daily Precipitation), and a latent heat of vaporization for water of $L =$
113 2260 kJ/kg. A larger value of AI indicates stronger aridity.

114 3 Results and discussion

115 3.1 Comparison to observed, satellite, and reanalysis data

116 We first compared the estimates of R_{ld} at a point-by-point basis separately for clear-sky and all-sky
117 conditions using Eqns. (2) and (5), respectively. This comparison is shown in Figure 2 using FLUXNET,
118 CERES, and ERA5 data. The estimates correlate very well with r^2 of 0.92 and 0.87 for clear-sky and all-
119 sky conditions, respectively, and RMSE values of 18.2 and 24.6 W m⁻². The slope of the linear regressions
120 between the estimated and observed R_{ld} for FLUXNET are 1.03 and 1.02, with most data points being
121 concentrated around the 1:1 line (Figs. 2a and 2b). Note that for all-sky conditions, the agreement is slightly
122 less good, with a lower correlation coefficient and a larger RSME. The agreement with the NASA-CERES
123 and ERA5 datasets are even better, with higher correlation coefficients and lower RSME.

124 Despite this high level of agreement of the estimates, we can see some systematic biases in the estimates
125 for R_{ld} . These can be seen in Figure 3, which shows the spatial distribution of these biases for the NASA-
126 CERES and ERA5 comparison. For clear-sky conditions, there appears to be a general underestimation in
127 the high latitudes and, to some extent, in arid regions (Figs. 3a and 3c). This bias can be attributed to biases
128 in the equations used here. Brutsaert (1975) already described that for very low temperatures and in arid
129 conditions, there are better parameter values than those used in Eq. 1, with a larger coefficient than 1.24
130 and a different exponent. The biases seen in Figure 3 are nevertheless notably smaller than the seasonal
131 variations shown in Figure 1.

132 The biases for all-sky conditions are typically larger and generally positive, implying an overestimate
133 except for the extremely cold and dry regions (Figs. 3b and 3d). This indicates that the effect of cloud cover
134 is probably more complex than the simple accounting expressed in Eqns. (4) and (5). However, the biases
135 are also small compared to the seasonal variations.

136 Overall, this evaluation shows that the expressions given by Eqns. (1) - (5) are very well suited to describe
137 the spatiotemporal variations of R_{ld} for current climatological conditions.

138

139 3.2 Attribution of diurnal and seasonal variations

140 We next use Eqns. (1) - (5) to attribute temporal variations of R_{ld} to their physical causes. To do so, we can
141 express changes ΔR_{ld} as a function of changes in water vapor, Δe_a , cloud cover, Δf_c , and air temperature,
142 ΔT_a . The functional dependence is derived from the equations by differentiation and applying the chain
143 rule. In a first step, we express a change ΔR_{ld} by the partial contributions $\Delta R_{ld,\varepsilon}$ and $\Delta R_{ld,T}$, that are due to
144 changes in emissivity, $\Delta \varepsilon$, and due to changes in atmospheric heat storage that are associated with a change
145 in air temperature ΔT_a :

$$\Delta R_{ld} = \Delta R_{ld,\varepsilon} + \Delta R_{ld,T} = \frac{\partial R_{ld}}{\partial \varepsilon} \Delta \varepsilon + \frac{\partial R_{ld}}{\partial T_a} \Delta T_a = \sigma \bar{T}_a^{-4} \Delta \varepsilon + 4\sigma \bar{\varepsilon} \bar{T}_a^{-3} \Delta T_a. \quad (8)$$

146 The contribution $\Delta R_{ld,\varepsilon}$ is further decomposed into contributions $\Delta R_{ld,f_c}$ and $\Delta R_{ld,e_a}$ due to variations in
147 clouds, Δf_c , and air humidity, Δe_a . Note that the contribution of a change ΔT_a to $\Delta R_{ld,\varepsilon}$ is relatively small
148 and is thus neglected here. We obtain:



$$\begin{aligned} \Delta R_{ld,\varepsilon} &= \sigma \bar{T}_a^{-4} \Delta \varepsilon \approx \frac{\partial \varepsilon}{\partial f_c} \Delta f_c + \frac{\partial \varepsilon}{\partial e_a} \Delta e_a \\ &= \sigma \bar{T}_a^{-4} \times \left(1 - 1.24 \left(\frac{e_a}{T_a} \right)^{\frac{1}{7}} \right) \Delta f_c + \sigma \bar{T}_a^{-4} \times \frac{1.24}{7} \frac{(1 - \bar{f}_c)}{(\bar{e}_a)^{\frac{6}{7}} (\bar{T}_a)^{\frac{1}{7}}} \Delta e_a, \end{aligned} \quad (9)$$

149

150 We next applied this approach to the diurnal deviations ΔR_{ld} from the daily mean using the FLUXNET
 151 dataset. This decomposition is shown in Figure 4 in aggregated form across the FLUXNET sites for whole
 152 year (Fig. 4a), the Northern hemisphere summer (Fig. 4b) and winter seasons (Fig. 4c). The diurnal
 153 variations of about $\pm 20 \text{ W m}^{-2}$ are primarily caused by diurnal changes in air temperature, while variations
 154 in emissivity play practically no role. Diurnal changes in air temperature reflect variations in heat storage
 155 of the atmospheric boundary layer. This is consistent with the notion that diurnal variations in solar
 156 radiation over land are buffered primarily by the lower atmosphere, rather than below the surface as it is
 157 the case for open water bodies and the ocean (Kleidon and Renner 2017). Since most of the stations in the
 158 FLUXNET dataset are located in the midlatitudes of the Northern hemisphere, the variations are
 159 consistently larger in summer due to the greater solar input (Fig. 4b) than in winter (Fig. 4c).

160 Figure 5 shows the same kind of decomposition, but for seasonal variations in R_{ld} in the NASA CERES
 161 dataset. The aggregation to the global scale across land and ocean is also shown in Fig. 5e-h. For the
 162 decomposition in Fig. 5e-h, the deviations are calculated as the difference of the monthly means to the
 163 annual mean. Generally, areas with relatively low annual-mean R_{ld} , e.g. the high latitude regions of North
 164 America and northeastern Eurasia, have the largest seasonal cycle (Fig. 1). The decomposition shows that
 165 this variation is mostly due to the seasonal variation in atmospheric heat storage ($\Delta R_{ld,T}$), although it is to
 166 some extent amplified by the contribution due to emissivity changes ($\Delta R_{ld,\varepsilon}$). Seasonal variations in
 167 emissivity play a greater role in changing R_{ld} in tropical areas, and this is predominantly due to cloud cover
 168 changes (Fig. 5c). The contribution by changes in water vapor is generally much smaller and does not vary
 169 much across regions (Fig. 5d).

170 Figs. 5e - h show that the seasonal variations of R_{ld} is generally less over the ocean than on the land, an
 171 effect that can also be seen in Fig. 1. The decomposition shows that these variations are mostly caused by
 172 changes in atmospheric heat storage, with a slight modulation by emissivity changes. This can, again, be
 173 largely explained by the effect described above for the diurnal variations (Kleidon and Renner 2017).
 174 Because solar radiation penetrates the transparent water bodies over marine areas, its variations are buffered
 175 below the surface, and not within the atmosphere. These variations are then reflected in seasonal changes
 176 in sea surface temperature. On land, however, seasonal heat storage changes below the surface are generally
 177 quite small, so that most variations take place in the lower atmosphere, which then alter R_{ld} .

178 These results show very similar patterns in the ERA5 dataset (Fig. S2).

179 In summary, what our decomposition shows is that most temporal variations in R_{ld} in current, climatological
 180 conditions are explained by heat storage changes within the lower atmosphere.

181

182 3.3 Attribution of geographic variations with aridity

183 Last, we applied the decomposition to the climatological variations in R_{ld} along differences in mean water
 184 availability. Water availability was characterized by Budyko's aridity index (AI), with values $AI < 1$
 185 representing humid regions, and larger values reflecting increased aridity. Here, the deviations ΔR_{ld} are
 186 calculated with respect to the global annual mean. The different contributions to the deviations are shown
 187 in Fig. 6, as well as the delineation along the aridity index (Figs. 6e - h).



188 The decomposition of the spatial distribution of the climatological means shows that the variations are
189 largely caused by differences in atmospheric heat storage as well (Fig. 6a). The contribution due to
190 variations in emissivity has a much smaller magnitude (Fig. 6b), and is dominated by changes in cloud
191 cover (Fig. 6c). Changes in water vapor appear to influence R_{ld} mostly in the humid tropics (Fig. 6d). The
192 results for the ERA5 reanalysis dataset show the same patterns (Fig. S1).

193 These variations are evaluated with respect to the aridity index in Figs. 6e - h. While there is a large spread,
194 as seen in the quantiles, there is a small, but consistent trend towards lower values of R_{ld} in more arid
195 regions, with a magnitude of about 20 W m^{-2} (black dashed line in Figs. 6e and 6g). We also notice a shift
196 in the contributions, with emissivity contributing less and atmospheric heat storage contributing more with
197 increased values of AI. The changes in emissivity is mostly caused by reductions in cloud cover, as shown
198 by the orange lines in Figs. 6f and h, and amounts to around -40 W m^{-2} over the range shown in the Figure.
199 This decrease in cloud cover can be attributed to the common presence of high-pressure systems in
200 subtropical areas (Zampieri et al. 2009). The decreasing contribution by lower cloud cover is compensated
201 for by an increased contribution of about $+20 \text{ W m}^{-2}$ by atmospheric heat storage that is caused by the
202 generally warmer mean temperatures in arid regions. This compensation is also seen in the ERA5 data
203 (Figs. 6g and 6h), with magnitudes of -40 W m^{-2} and $+30 \text{ W m}^{-2}$ respectively.

204 Taken together, these trends imply that, again, the climatological variations in R_{ld} are also dominated by
205 differences in atmospheric heat storage. A small, but consistent change can be seen in the contributions
206 along the aridity index, with the contribution by emissivity due to cloud cover becoming lower while the
207 contribution by atmospheric heat storage increases as regions become drier.

208 **4. Summary and Conclusions**

209 We found that the semiempirical equations of Brutsaert (1975) and Crawford and Duchon (1999) work very
210 well to estimate the downwelling flux of longwave radiation by comparing these to estimates from
211 observation, satellite, and reanalysis datasets. We then showed that one can use these equations to
212 decompose this flux into different components, and relate changes to differences in cloud cover, water
213 vapor, and atmospheric heat storage. We found that most changes in downwelling longwave radiation are
214 caused by differences in atmospheric heat storage that are reflected in differences in air temperature.

215 These equations can then be applied to different aspects of climate research. For instance, the values of
216 downwelling longwave radiation are often missing in FLUXNET data (Table S2), and these equations can
217 be used to fill the gaps with air temperature and humidity observations. We can also use these equations to
218 better understand the physical mechanisms for temperature change due to extreme events. For instance,
219 Park et al. (2015) and Alekseev et al. (2019) found that an enhancement of downwelling longwave radiation
220 in the Arctic is found to be preceded by the advection of moisture and heat. The equations by Brutsaert
221 (1975) and Crawford and Duchon (1999) can then be used to quantify the individual contributions by the
222 advection of heat and moisture (Tian et al. 2022). Another example is the attribution of differences in
223 global warming magnitudes across humid and arid regions. Du et al. (2020) used these equations to explain
224 why global warming was stronger during clear-sky conditions in observations in China due to the greater
225 sensitivity of clear-sky emissivity to a change in water vapor. This was then used to explain the observed,
226 stronger global warming in the arid regions of China, which have less clouds and a higher frequency of
227 clear-sky conditions than the humid regions. While the empirical coefficient of 1.24 in Eq. (1) may change
228 due to emissivity changes from greenhouse gases other than water vapor, this formulation can nevertheless
229 provide a useful basis.

230 We conclude that the equations by Brutsaert (1975) and Crawford and Duchon (1999) are still very useful
231 to advance our understanding of surface temperature changes. Our evaluation has shown how well these
232 equations estimate this flux, and our application to the decomposition of different contributions has shown
233 its utility to understand the causes for its variation. These equations should help us to better understand
234 aspects of climate variability, extreme events, and global warming, linking these to the mechanistic
235 contributions by downwelling longwave radiation.



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242 **Author contributions**

243 YLT, SAG, and AK conceived and designed the analysis, with inputs from DZ and GW. YLT performed
244 the analysis and discussed the results with all authors. YLT and AK wrote the paper.

245 **Competing interests**

246 The contact author has declared that none of the authors has any competing interests.

247 **Data availability**

248 The data used in this study was downloaded from the links provided with the references. No new data was
249 created.

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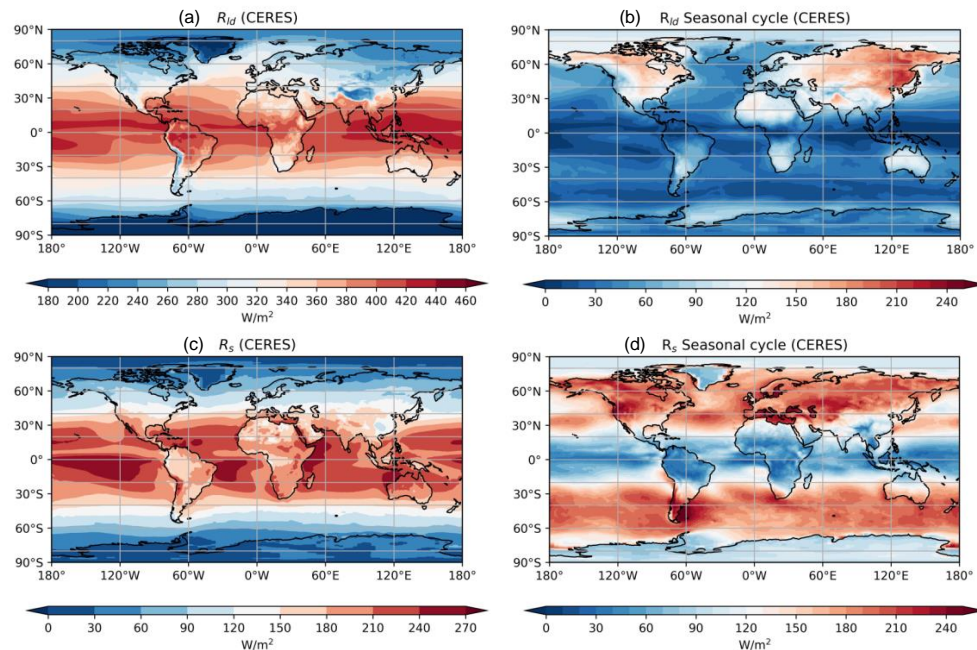


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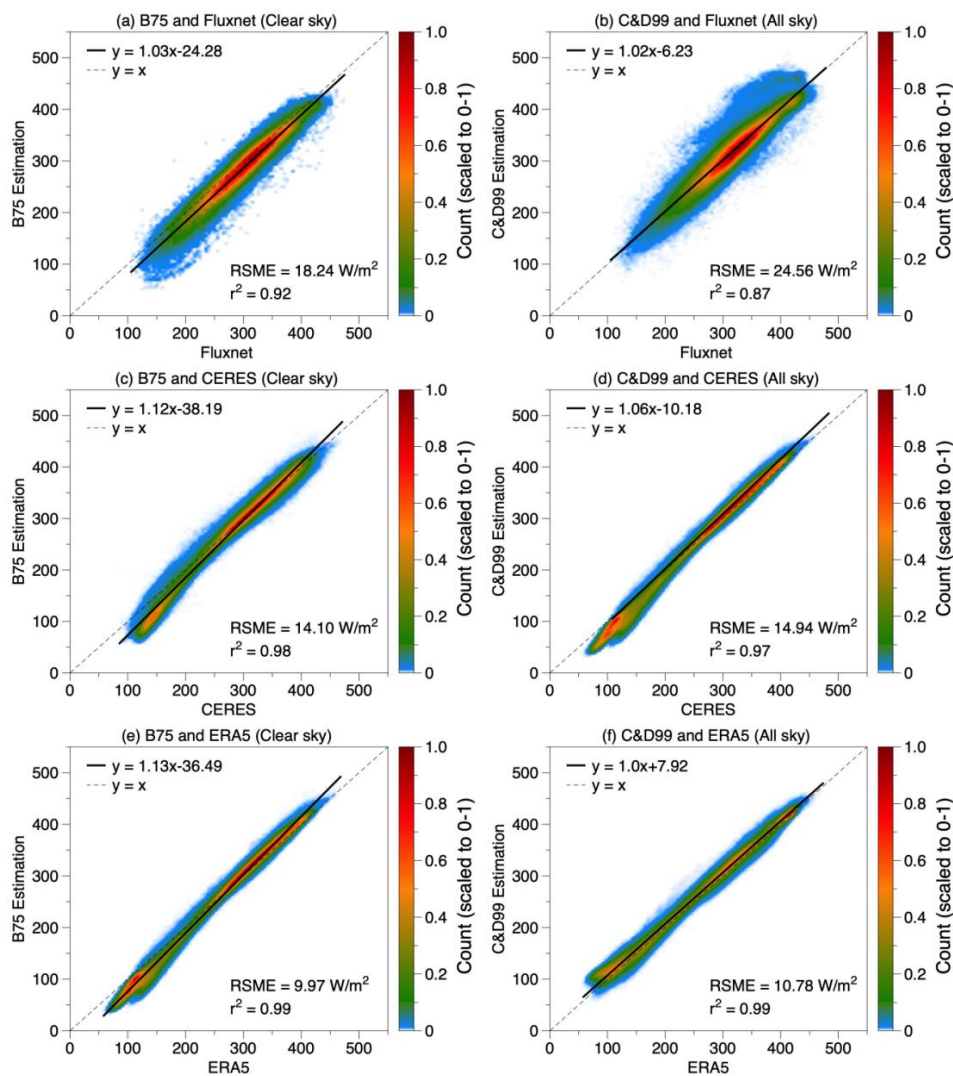
369 **Figures**



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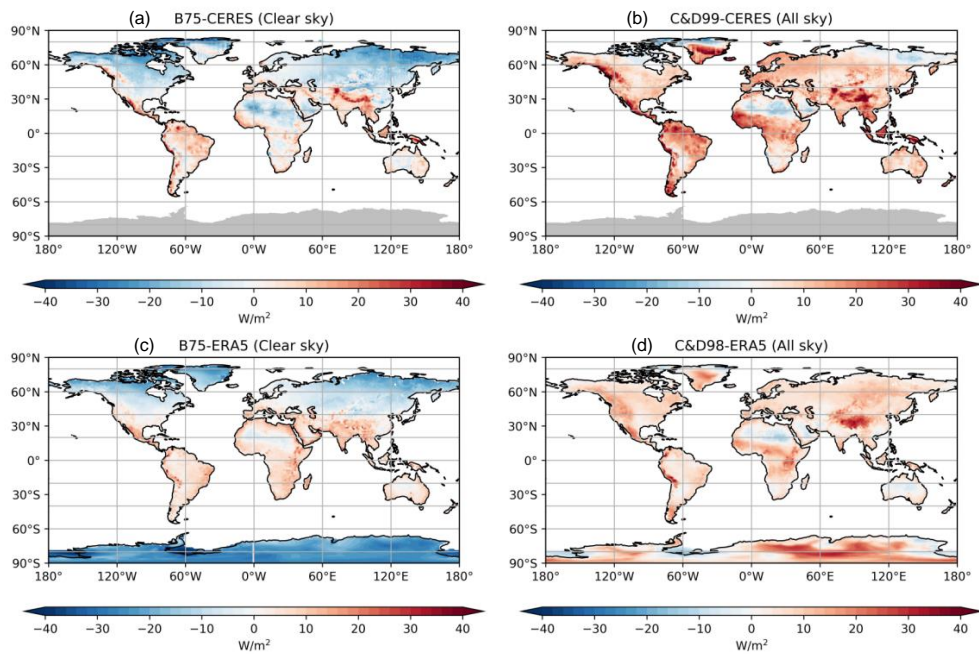
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372 **Figure 1.** Spatial distribution of (a, c) the climatological mean and (b, d) the seasonal amplitude of
373 downward longwave radiation and absorbed solar radiation at the surface respectively from the NASA-
374 CERES dataset. The seasonal amplitude is calculated as the difference between the maximum and minimum
375 monthly mean.



376

377 **Figure 2.** Comparison of R_{ld} estimated (a, c, e) by Brutsaert (1975) for clear-sky conditions and (b, d, f) by
 378 Crawford and Duchon (1999) for all-sky conditions using FLUXNET (a, b), NASA-CERES (c, d) and
 379 ERA5 (e, f). Colors indicate the density of the data points and is scaled to values between 0 - 1.



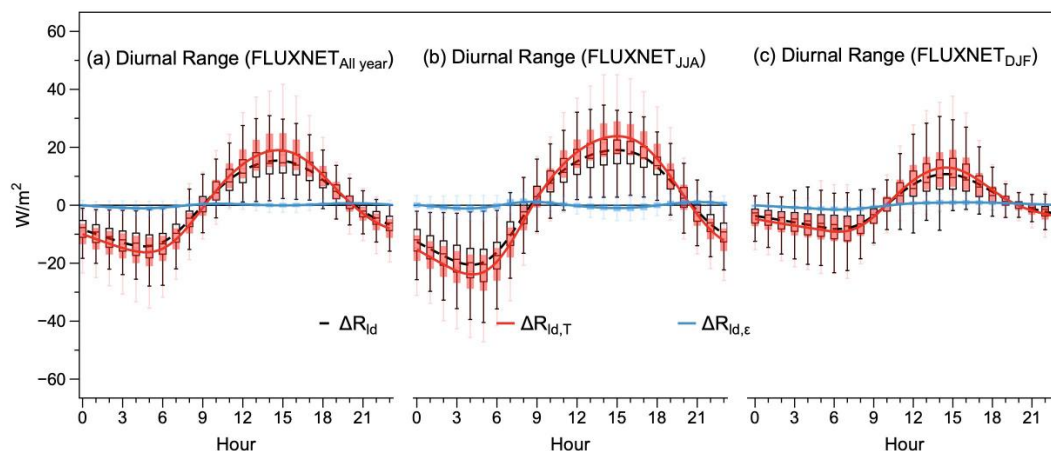
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Figure 3. Biases in the estimates for R_{ld} for (a, b) NASA-CERES and (c, d) ERA reanalysis for (a, c) clear-sky and (b, d) all-sky conditions over land. Grey shading indicate missing values.



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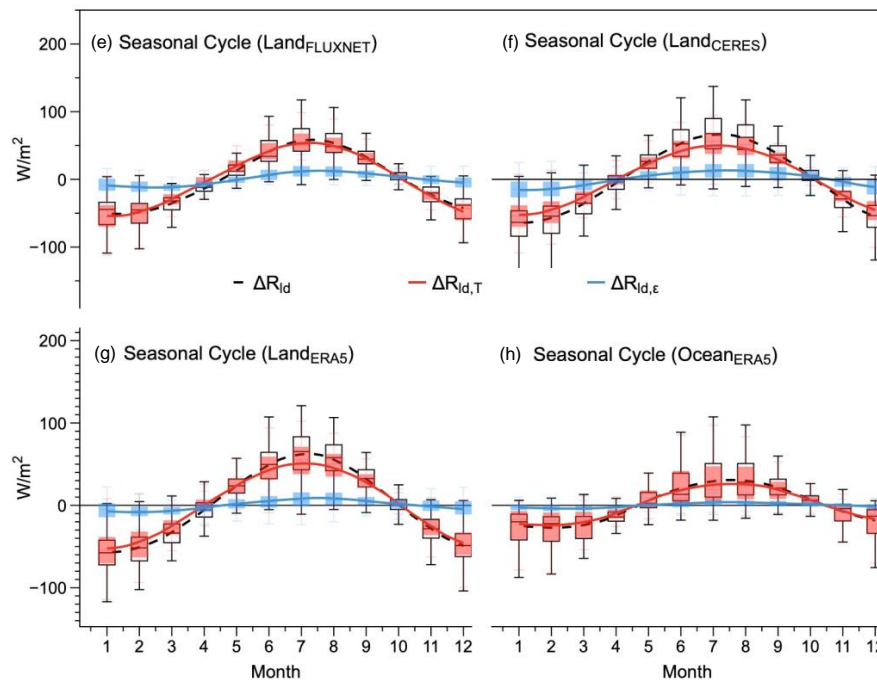
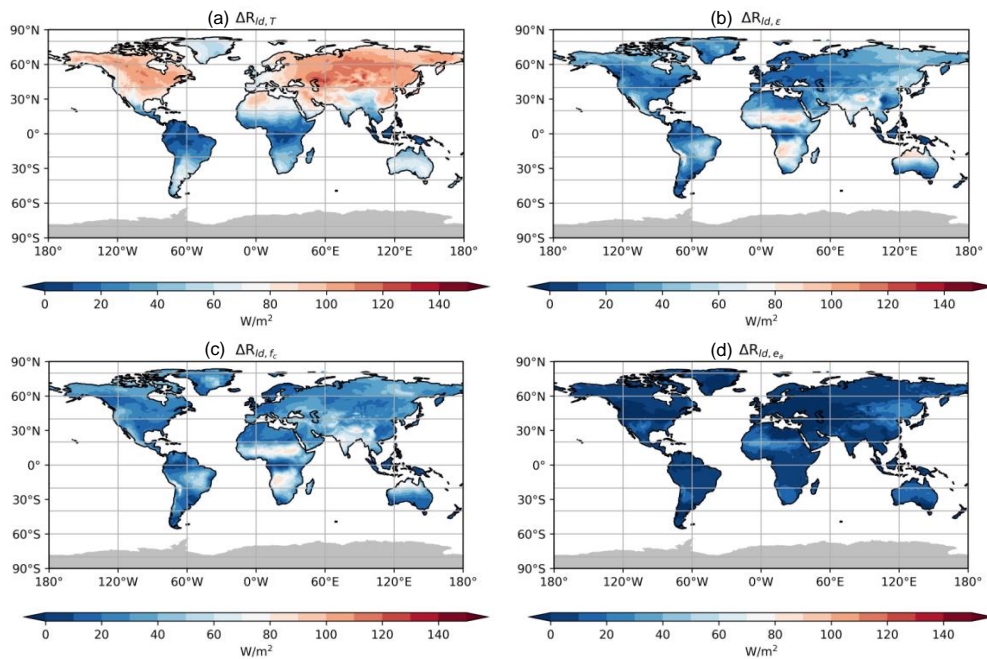
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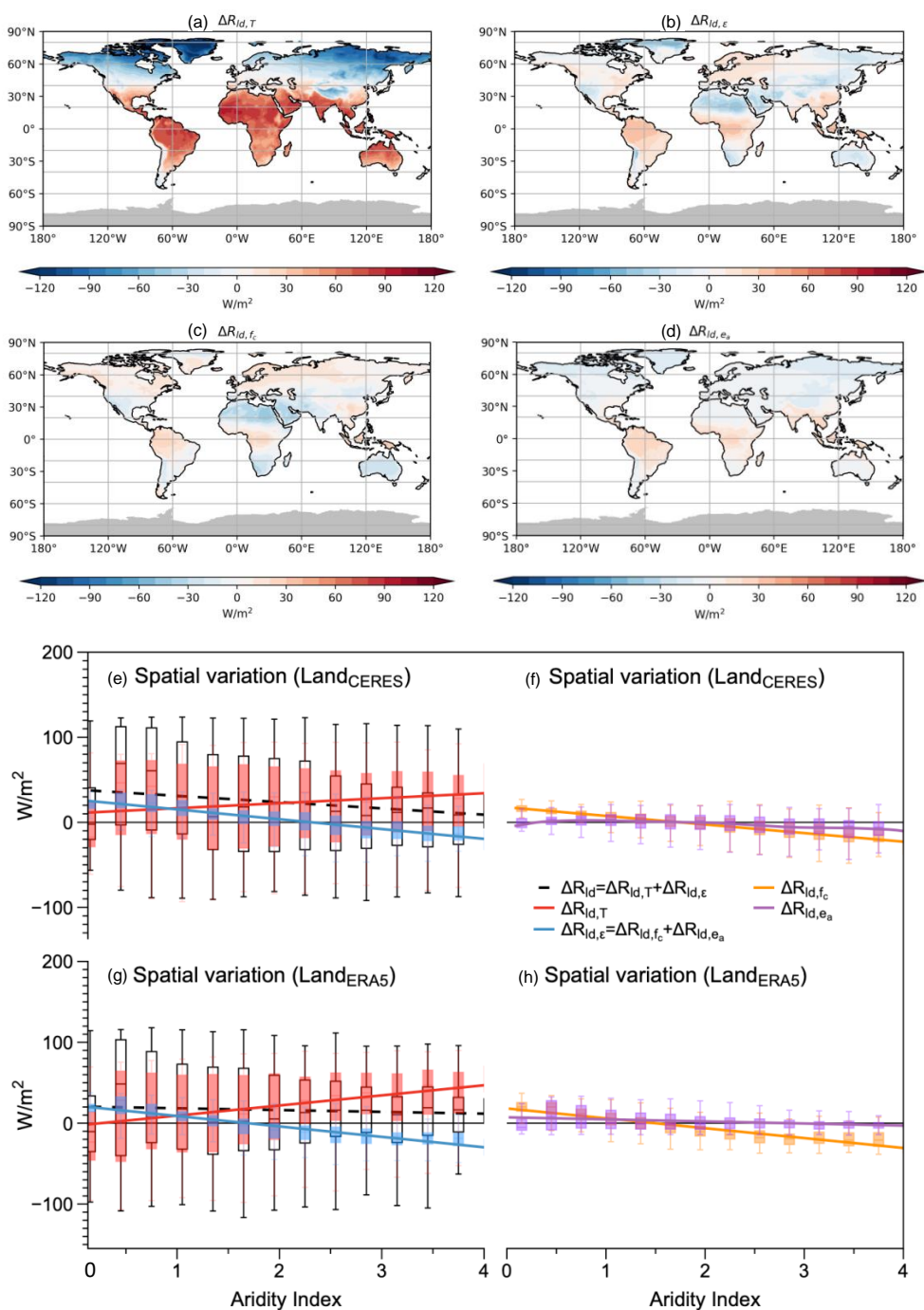
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Figure 4. Diurnal variations in R_{ld} (black dashed line) and its decomposition into contributions by changes in emissivity (blue, $\Delta R_{ld,\epsilon}$) and atmospheric heat storage (red, $\Delta R_{ld,T}$) in the FLUXNET dataset aggregated over 189 sites for (a) the whole year, (b) June-August, and (c) December - February. The upper and lower whisker indicate 95th and 5th percentiles, upper boundary, median line, and lower boundary of the box indicate the 75th, 50th, 25th quantiles, respectively. Regression lines are based on site-mean or grid-mean value using LOESS regression.





393 **Figure 5.** Decompositions of the mean seasonal variation of R_{ld} (difference between the maximum and
394 minimum monthly means) in the NASA-CERES dataset into contributions by (a) atmospheric heat storage
395 ($\Delta R_{ld,T}$) and (b) emissivity ($\Delta R_{ld,\varepsilon}$). The variations in $\Delta R_{ld,\varepsilon}$ are further decomposed in contributions by
396 variations in (c) cloud cover ($\Delta R_{ld,fc}$) and (d) humidity ($\Delta R_{ld,ea}$). Panels (e) - (h) show the mean seasonal
397 variation of R_{ld} averaged over land or ocean and its decomposition into variations in atmospheric heat
398 storage ($\Delta R_{ld,T}$) and emissivity ($\Delta R_{ld,\varepsilon}$) using FLUXNET, NASA-CERES, and ERA5. The box plot
399 indicates percentiles and quantiles as in Figure 4. Grey shading indicates missing values.





401 **Figure 6.** As Figure 5, but for the decomposition of the spatial variations of the annual mean (a - d) and the
402 variations with the aridity index (e - h).

403

404

405

406 Plain Language:

407

408 Downward longwave radiation (R_{ld}) plays an important role in surface energy balance and is critical for
409 global warming. However, its spatiotemporal climatological variation on a global scale has not been
410 explained well with a solid physical basis. To fill this gap, we here use a semi-empirical equation derived
411 by Brutsaert (1975, "B75") and its extension by Crawford and Duchon (1999, "C&D99") to identify the
412 leading terms responsible for the diurnal range, seasonal cycle, and geographical variations in R_{ld} . We show
413 that B75 and C&D99 work very well when evaluated against global observations from satellites and
414 FLUXNET sites. We then used these physics-based equations to show that diurnal and seasonal variations
415 in R_{ld} are predominantly controlled by changes in atmospheric heat storage. When moving from humid to
416 arid regions, while the contribution of atmospheric heat storage increases, the ones from clouds decreases,
417 which together explains the relatively small decrease in R_{ld} with aridity. Our work provides a clue to better
418 understand aspects of climate variability, extreme events, and global warming, by linking these to the
419 mechanistic contributions by downwelling longwave radiation.

420