

# Internal tides off the Amazon shelf Part I: importance for the structuring of ocean temperature during two contrasted seasons

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## Abstract

Tides and internal tides (IT) in the ocean can significantly affect local to regional ocean temperature, including sea surface temperature (SST), via physical processes such as diffusion (vertical mixing) and advection (vertical and horizontal) of water masses. Offshore of the Amazon River, strong IT have been detected by satellite observations and well modelled; however, their impact on temperature, SST and the identification of the associated processes have not been studied so far. In this work, we use high resolution (1/36°) numerical simulations with and without the tides from an ocean circulation model (NEMO). This model explicitly resolves the internal tides (IT) and is therefore suitable to assess how they can affect ocean temperature in the studied area. We distinguish the analysis for two contrasted seasons, from April to June (AMJ) and from August to October (ASO), since the seasonal stratification off the Amazon River modulates the IT's response and their influence in temperature.

The generation and the propagation of the IT in the model are in good agreement with observations. The SST reproduced by the simulation including tides is in better agreement with satellite SST data compared to the simulation without tides. During ASO season, stronger meso-scale currents, deeper and weaker pycnocline are observed in contrast to the AMJ season. The observed coastal upwelling during ASO season is better reproduced by the model including tides, whereas the no-tide simulation is too warm by +0.3 °C for the SST. In the subsurface above the thermocline, the tide simulation is cooler by -1.2 °C, and warmer below the

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33 thermocline by +1.2 °C compared to the simulation without the tides. The IT induce vertical  
34 mixing on their generation site along the shelf break and on their propagation pathways towards  
35 the open ocean. This process mainly explains the cooler temperature at the ocean surface and  
36 is combined with vertical and horizontal advection to explain the cooling in the subsurface  
37 water above the thermocline and a warming in the deeper layers below the thermocline. The  
38 surface cooling induced in turn an increase of the net heat flux from the atmosphere to the  
39 ocean surface, which could induce significant changes in the local and even for the regional  
40 tropical Atlantic atmospheric circulation and precipitation.

41 ~~We therefore demonstrate that IT, via vertical mixing and advection along their~~  
42 ~~propagation pathways, and tides over the continental shelf, can play a role on the temperature~~  
43 ~~structure off the Amazon River mouth, particularly in the coastal cooling enhanced by the IT.~~

44 The impact of the tides (internal and barotropic) on the vertical and horizontal structure  
45 of temperature off the Amazon River is investigated over two highly contrasting seasons.  
46 Twinned regional simulations with and without tides are used to highlight the general effect of  
47 tides. The tides tend to cool down the ocean from the surface (~0.3 °C) to above the thermocline  
48 (~1.2 °C), and to thaw it below the thermocline (~1.2 °C). The heat budget analysis leads to  
49 the conclusion that vertical mixing could represent the dominant process that drives these  
50 temperature variations within the mixed layer, while it is associated with both horizontal and  
51 vertical advection below to explain temperature variations. The intensified mixing in the  
52 simulation including tides is attributed to the breaking of internal tides (ITs), on their generation  
53 sites over the shelf break and offshore along their propagation pathways. When over the shelf  
54 the mixing is attributed to the dissipation of the barotropic tides. Both horizontal and vertical  
55 advectons exist in simulations without the tides but are strengthened when including it.  
56 Furthermore, vertical heat budget equation terms show a typical mode-1 horizontal propagation  
57 wavelength of ITs.

58 In addition, we found the tides can also have an impact on interactions between the  
59 upper ocean interface and the underlying atmosphere. They account for a significant proportion  
60 of the net heat flux between the atmosphere and the ocean, with a marked seasonal variation of  
61 33.2% to 7.4% between the first and second seasons. Tidal dynamics could be therefore critical  
62 to understand the regional climate. This study highlights the key role of tides, particularly, how  
63 ITs-related vertical mixing helps to shape ocean temperature off the Amazon.

64 **Keywords:** Amazon shelf break, internal tides, ~~Amazon continental shelf and slope~~mixing,  
65 temperature, heat flux, modeling, satellite data, ~~mixing~~, ~~heat flux~~.

## I. Introduction

Temperature and its spatial structure play a crucial role in ocean dynamics, including water mass formation (Swift and Aagaard, 1981; Lascaratos, 1993; Speer et al., 1995), transport and mixing of other tracers in the ocean and exchanges with other biosphere compartments (Archer et al., 2004; Rosenthal et al., 1997), and most importantly on surface heat exchange at the interface with the atmosphere (Calyson and Bogdanoff, 2013; Mei et al., 2015) and can thus significantly influence the climate (Li et al., 2006; Collins et al., 2010). This oceanic thermal structure can be modified at various spatial and temporal scales by different processes external to the ocean such as incident solar radiation, heat flux exchanges with the atmosphere, winds, precipitation, and freshwater inputs from rivers. And by its internal processes such as mass transport by currents and eddies (e.g., Aguedjou et al., 2021), mixing by turbulent diffusion (Kunze et al., 2012), tides and internal tidal waves (IT) and their dissipation (Barton et al., 2001; Smith et al., 2016; Salamea et al., 2021). Finally, bottom friction of IT on the thermal structure of the ocean is of increasing interest with many studies in recent years, but remains poorly understood in many ocean regions, and the barotropic tidal currents may also produce intensified mixing especially off the Amazon for shallow water condition (e.g., over a shelf, see Lambeck and Runcorn, 1977; Le Provost and Lyard, 1997) and significantly modify ocean temperature in surface layers (Li et al., 2020).

In a stratified ocean, the passage of a barotropic tide over a topographic profile with a steep slope (continental slope, seamount, oceanic ridge) generates a disturbance in the flow that gives rise to a so-called baroclinic tide, with the same frequency, but with higher vertical velocities (Zhao et al., 2016). The baroclinic tide, also known as internal tidal waves (IT), thus captures part of the energy of the barotropic tide, propagates it and dissipates it into the global ocean by diapycnal mixing (Zhao et al., 2012), i.e., up to about 1 TW in the deep ocean (Egbert and Ray, 2000; Niwa and Hibiya, 2011) and thus helps to feed the thermohaline circulation (Munk and Wunsch, 1998). These two tidal processes (barotropic tide and IT) thus bring together a set of mechanisms for transferring and redistributing energy from larger to smaller oceanic scales, which can be understood as a tidal energy cascade. The dissipation of IT occurs mainly locally at the generation sites for high-mode IT that are associated with higher vertical shear, while a significant part of the energy dissipates offshore along their propagation path for low-mode IT (Zhao et al., 2016). Results from models in the Indonesian seas (Koch-Larrouy et al., 2007 and Nugroho et al., 2018) and observations in the Celtic Sea (Sharples et al., 2007)

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99 and the Yellow Sea (Xu et al., 2020), point out that IT dissipate most of their energy vertically,  
100 where the vertical gradient of stratification is maximal in the water column. IT can also  
101 vertically advect water masses during their propagation. They thus induce vertical  
102 displacements of the isopycnal levels of several meters to a few tens of meters, observable in  
103 the thermocline (Wallace et al., 2008; Xu et al. 2020). Denmann and Garrett (1983) and Bordois  
104 et al. (2016) point out that the stratification peak acts as a waveguide for the propagation of IT.

105 The mixing and advection induced by IT results in a change in temperature structure  
106 throughout the water column. In surface waters, Smith et al. (2016) report that IT can induce  
107 surface cooling varying between 1 °C and 5 °C depending on the ocean region. Koch-Larrouy  
108 et al. (2007) and Nagai and Hibiya (2015) have shown, for the Indonesian region, that IT  
109 induces a surface cooling of 0.5 °C on average and that this decreases cloud convection in the  
110 atmosphere on a local scale, which in turn reduces precipitation by 20% and thus plays an  
111 important role on the climate on a regional scale. Furthermore, Jithin and Francis (2020)  
112 showed that IT can also affect the temperature in deep waters (>1600 m), leading to a warming  
113 of the order of 1–2 °C.

114 The barotropic tides dissipate most of their energy in shallow coastal waters by bottom  
115 friction when the mean ocean depth becomes less than the amplitude of the tide (Lambeck and  
116 Runcorn, 1977; Le Provost and Lyard, 1997), and can thus modify temperature. Furthermore,  
117 Gonzalez-Haro et al. (2019) showed that barotropic and baroclinic tidal currents can induce  
118 temperature fluctuations by horizontal advection of surface water masses over hundreds of  
119 kilometers, and thus contribute to modifying the SST. These two tidal processes (barotropic  
120 and baroclinic) can also affect other tracers such as nutrients, chlorophyll and sediments  
121 (Heathershaw et al., 1987; da Silva et al., 2002; Sharples et al., 2007; Pomar et al., 2012;  
122 Muacho et al., 2014; Tuarena et al., 2016; Barbot et al., 2022).

123 The key source for internal waves generation is the barotropic or external tides. The  
124 external tides when interacting with sharp topography (e.g., ridge, sea mounts, shelf break) in  
125 a stratified ocean generate internal tides also called internal tidal/gravity waves, that may  
126 propagate and dissipate in the ocean interior causing diapycnal mixing (Baines, 1982; Munk  
127 and Wunsch, 1998; Egbert and Ray, 2000). The precise location of this dissipation is a big  
128 unknown. But evidence of dissipation at the generation sites, at the reflection to the bottom or  
129 close to the surface when the energy rays interact with the thermocline and pycnocline, have  
130 been measured and modelled (among others: Laurent and Garrett, 2002; Sharples et al., 2007,  
131 2009; Koch-Larrouy et al., 2015; Nugroho et al., 2018; Xu et al., 2020 and Whalen et al., 2020).

132 ITs may also dissipate or lose energy when they encounter others or when they interact with  
133 mesoscale or fine-scale structures (Vlasenko and Stashchuk, 2006; Dunphy and Lamb, 2014).  
134 Moreover, the surface interactions allow nonlinear internal solitary waves (ISW) to develop  
135 and to propagate usually with phase-locked to the ITs troughs (New and Pingree, 1990, 2000;  
136 Azevedo et al., 2006; da Silva et al., 2011). Finally, ISW can dissipate and induce mixing  
137 (Sandstrom and Oakey, 1995; Feng et al., 2021; Purwandana et al., 2022). Moreover, ITs can  
138 vertically advect the water masses following their propagation. The effect is the vertical shifts  
139 in isopycnic levels of few meters to tens of meters, which can be observed in the thermocline  
140 (Wallace et al., 2008; Xu et al., 2020). But over a tidal cycle, the mean effect on temperature  
141 is null except some tidal residual circulation exists (Bessières, 2007).

142 Our study focuses on the oceanic region of northern Brazil off the Amazon River, where  
143 IT have been highlighted in previous studies, but their impact on the thermal structure is not  
144 currently known. This region is characterized by a broad, shallow continental shelf at the  
145 mouth of the Amazon River ended by a steep slope, i.e., a bathymetry exhibits a variation of  
146 200-2000 m over some tens of kilometers (Fig.1). Along this slope, six sites (A to F) of IT  
147 generation have been identified (Fig.1), the most intense of which (A and B) are in the south  
148 of the region (Magalhaes et al. 2016, Barbot et al. 2021 and Tchilibou et al. 2022). A strong  
149 seasonal coastal current, the Brazilian North Current (NBC), strongly influences the study area  
150 and flows along the coast from the southeast to the northwest (Johns et al., 1990).

151 This region shows a seasonal variation in the wind position and hence the position  
152 of the winds and the Intertropical Convergence Zone (ITCZ) during the year, which. This  
153 directly influences the discharge of the Amazon River, the oceanic circulation (Xie and Carton,  
154 2004), and therefore the stratification. This impacts the activity of internal tidal waves (Barbot  
155 et al., 2021). Two seasons can be clearly distinguished by their properties on water masses and  
156 currents.

157 The first season runs from March to July, during this time the ITCZ is in its most  
158 equatorial position and lies in the heart of our region. The increase in rainfall over the ocean  
159 leads to a colder and more homogeneous SST far from the coast. The discharge of the Amazon  
160 River into the ocean reaches its peak ( $> 3 \times 10^5 \text{ m}^3 \cdot \text{s}^{-1}$ ) and the surface temperature in the coastal  
161 zone, although homogeneous, is warmer than offshore. At the end of this season, driven by the  
162 strong river discharge, the Amazon plume along the shelf extends beyond  $8^\circ\text{N}$ , and sometimes  
163 into the Caribbean region (Müller-Karger et al., 1989; Johns et al., 1998). The stratification is  
164 somewhat stronger and more homogeneous horizontally, and the maximum of its vertical

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165 gradient (pycnocline) is closer to the surface. This latter point leads to a stronger conversion of  
166 energy from barotropic to baroclinic, and a stronger local (first 50 km) dissipation of internal  
167 tidal wave energy (Barbot et al., 2021; Tchilibou et al., 2022). NBC and eddy kinetic energy  
168 (EKE) are weak in the region (Aguedjou et al., 2019). Close to the equator, the NBC develops  
169 a retroflection towards 1°N latitude that feeds the Equatorial Under Current (EUC) transporting  
170 water masses eastwards to the Gulf of Guinea (Didden and Schott, 1993; Dimoune et al., 2022).  
171 Contrasting with this first season, due to different oceanic and atmospheric conditions,  
172 the second season extends from and the stratification (Muller-Karger et al., 1988; Johns et al.,  
173 1998; Xie and Carton, 2004). Hence, two very contrasting seasons form, April-May-June  
174 (AMJ) and August to December. During this season, the ITCZ migrates to its northernmost  
175 position around 10°N. In response, rainfall in this area decreases and the Amazon River  
176 discharge also decreases to its minimum ( $\sim 10^5 \text{ m}^3 \cdot \text{s}^{-1}$ ), the extension of the river plume is  
177 therefore reduced to no more than 200–300 km offshore from the mouth of the Amazon  
178 between November and December (Johns et al., 1998; Garzoli et al., 2003). During this season,  
179 cold water ( $< 27.6 \text{ }^\circ\text{C}$ ) associated with the western extension of the Atlantic cold water tongue  
180 (ACT) enters the region from the south and runs along the edge of the continental shelf to about  
181 3°N (Lentz and Limeburner, 1995; Neto et al., 2014), forming a cold cell often referenced as a  
182 seasonal upwelling. The stratification of the study area is strongly modified compared to the  
183 previous season. The September-October (ASO). AMJ (vs. ASO) season is characterized by  
184 an increasing (vs. decreasing) river discharge, stronger (vs. smaller) and shallower (vs. deeper)  
185 pycnocline becomes somewhat deeper. The generation of IT on the slope and their local  
186 dissipation. The North Brazilian Current (NBC) and eddy kinetic energy (EKE) are weaker  
187 compared to the first season (Barbot vs. stronger) (Aguedjou et al., 2021; 2019, Tchilibou et  
188 al., 2022). Currents and eddy activity become stronger. The NBC becomes For the ASO season,  
189 the stronger, farther from the coast and deeper, it NBC develops a retroflection (NBCR)  
190 between 5°–8° N that feeds the North Equatorial Counter-Current (NECC) transporting the  
191 water masses towards the east of the tropical Atlantic. This The retroflexion also generates very  
192 large anticyclonic eddies (NBC Rings) exceeding 450 km in diameter (Didden and Schott,  
193 1993; Richardson et al., 1994; Garzoli et al., 2004), which in turn transport water masses  
194 towards the Northern Hemisphere (Bourlès et al., 1999a; Johns et al., 1998; Schott et al., 2003).  
195 In addition, there are more cyclonic/anticyclonic eddies from the Gulf of Guinea during this  
196 season. All this contributes to the strengthening of the EKE, which reaches its maximum in this  
197 season (Aguedjou Bourlès et al., 1999a; Johns et al., 1998; Schott et al., 2003). et al., 2019).

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198 ~~When the baroclinic tidal flow interacts with the general circulation, it is deviated from its~~  
199 ~~trajectory, thus we have a so-called incoherent baroclinic tide (Buijsman et al., 2017). This IT-~~  
200 ~~circulation interaction is thus reinforced during this second season because of the more~~  
201 ~~intensified currents and eddy activity (Tchilibou et al., 2022).~~

202 ~~On the Brazilian continental shelf, Geyer et al. (1996) suggests the presence of internal~~  
203 ~~tidal waves from current data. Later, Lentini et al. (2016) based on SAR imagery, shows small-~~  
204 ~~wavelength (~ 10 km) internal solitary waves (ISW) packets propagating along and across the~~  
205 ~~continental shelf, and generated by linear non-hydrostatic interactions between the NBC and~~  
206 ~~barotropic tidal currents. Using a model, Molinas et al. (2020) show that baroclinic tidal~~  
207 ~~currents play an important role on sediment transport on this continental shelf. On a global~~  
208 ~~scale, several studies from altimetry observations (Zhao et al., 2012, 2016; Zaron et al., 2017,~~  
209 ~~2019) or models (Munk and Wunsch, 1998; Shriver et al., 2012; Arbic et al., 2012; Niwa and~~  
210 ~~Hibiya, 2011, 2014; Buijsman et al. 2016) have shown intense activity of IT along the steep~~  
211 ~~slope of the continental shelf. At the surface, using SAR imagery, Jackson (2007) and~~  
212 ~~Magalhaes et al. (2016) describe longer wavelength ISW (~ 50-150 km) that propagate~~  
213 ~~offshore from the slope. The latter emphasizes the modulation of their propagation by the~~  
214 ~~seasonal variation of the NECC, and from a model, establishes that these ISW originate from~~  
215 ~~instabilities and energy loss of IT coming from the slope, mainly at sites A and B. Internal tides~~  
216 ~~are generated on the sharp shelf break which possesses a depth decreasing of 200-2000 m over~~  
217 ~~some tens of kilometers (Fig.1). Six main sites (A to F) have been identified, with the most~~  
218 ~~intense, A and B, located in the southern part of the region (Fig.1; Magalhaes et al., 2016,~~  
219 ~~Tchilibou et al., 2022). Previous studies have shown that in this region ITs propagation is~~  
220 ~~modulated by the seasonal variation of the currents (Magalhaes et al., 2016; Lentini et al., 2016;~~  
221 ~~Tchilibou et al., 2022; de Macedo et al., 2023). In addition, seasonal variations in stratification~~  
222 ~~induce changes in the internal tide's activity. With in AMJ (vs. ASO) a stronger (vs. smaller)~~  
223 ~~energy conversion and a stronger (vs. smaller) local dissipation of ITs energy (Barbot et al.,~~  
224 ~~2021, Tchilibou et al., 2022). Moreover, the interaction between the weaker (vs. stronger)~~  
225 ~~background circulation and ITs can lead to less (vs. more) incoherent or non-stationary internal~~  
226 ~~tides (Tchilibou et al., 2022). Incoherent ITs can account for about half of the total internal~~  
227 ~~tides in the global ocean and much more when looking at some regional ocean system. For~~  
228 ~~example over 80% in equatorial Pacific (Zaron, 2017) and over 40% off the Amazon (see~~  
229 ~~Fig.11e-f in Tchilibou et al., 2022). But quantifying the associated energy is difficult to~~  
230 ~~determine and is still unknown in our region but is part of the scope of upcoming studies.~~

231 The role of ITs on the thermal structure of the ocean is of increasing interest with many  
232 studies in recent years. In the Hawaii shallow shelf surface waters, Smith et al. (2016) report  
233 that ITs can induce surface cooling from 1 °C to 5 °C. For the Indonesian region, ITs induce  
234 an annual mean surface cooling of 0.5 °C (Koch-Larrouy et al., 2007, 2008; Nagai and Hibiya,  
235 2015 and Nugroho et al., 2018), that decreases local atmospheric convection, which in turn  
236 reduces precipitation by 20%. They can therefore fulfil a relevant role on regional climate  
237 (Koch-Larrouy et al., 2010, Sprintall et al., 2014, 2019). Furthermore, in the Andaman Sea,  
238 Jithin and Francis (2020) showed that ITs can affect the temperature in deep waters (> 1600  
239 m), leading to a warming of about 1–2 °C. But off the Amazon plateau, the impact of ITs on  
240 the thermal structure of the ocean is still poorly understood.

241 During the ASO season, cold water (< 27.6 °C) associated with the western extension  
242 of the Atlantic Cold-water Tongue (ACT) enter the region from the south and run along the  
243 edge of the continental shelf to about 3°N, establishing a cold cell often referred to as seasonal  
244 upwelling (Lentz and Limeburner, 1995; Neto and da Silva, 2014). Modelling studies, with and  
245 without tides, have shown that this upwelling is affected by the tides. Cooling is more realistic  
246 when tides are included (Ruault et al., 2020). However, these analyses cannot determine what  
247 processes are at work. For example, it is not yet explicit whether the tidal-induced cooling is  
248 due to mixing on the shelf produced by barotropic tides, or to the mixing produced by baroclinic  
249 tides at their generation sites and propagation pathways. Neto and da Silva (2014), based on *in*  
250 *situ* observations, suggest instead that it is the vertical advection triggered by the NBC that can  
251 explain the cooling observed at the surface. Following on from the latter, we can also examine  
252 the role of horizontal advection and its contribution relative to vertical advection.

253 ~~To Recently, de Macedo et al. (2023) provided a somewhat more comprehensive~~  
254 ~~description of the seasonal characteristics of these ISW, with the predominant origin remaining~~  
255 ~~at sites A and B. Barbot et al. (2021) focused on the influence of stratification on the IT~~  
256 ~~generation on the shelf as well as their propagation offshore. Finally, also in seasonal scale,~~  
257 ~~Tehilibou et al. (2022) looked at the variation of the energy associated with IT from their~~  
258 ~~generation to their dissipation, as well as the interaction of these waves with the general~~  
259 ~~circulation. However, the interactions between IT and tracers such as temperature, salinity or~~  
260 ~~chlorophyll have not received much interest from the scientific community in this region.~~

261 ~~Hydrodynamic and biogeochemical conditions on the shelf and off the mouth of the~~  
262 ~~Amazon were studied during the AMASSEDS campaigns in the early 1990s (DeMaster and~~  
263 ~~Pope, 1996; Nittouer and DeMaster, 1996) and the various “Camadas Finas” campaigns~~

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264 (~~Araujo et al., 2017, 2021~~). Furthermore, using data from the two REVIZEE campaigns and  
265 TMI SST satellite data, Neto et al. (2014) studied the seasonal cooling of surface waters, which  
266 occurs between the months of July and December. They conclude that the NBC is responsible  
267 for the upwelling of cold water masses ( $< 27.5^{\circ}\text{C}$ ) to the more superficial layers. Subsequently,  
268 Araujo et al. (2016), using in addition a realistic model, suggest that the tide would have a key  
269 role to play in intensifying this cooling. Indeed, through twin simulations with and without tide,  
270 they show that with tide there is a  $-0.3^{\circ}\text{C}$  cooling of the surface temperature. These analyses  
271 remain qualitative and do not allow determining what are the processes at work. Knowing that  
272 we are in an area with a strong activity of internal tidal waves, the question remains whether  
273 and by what processes these IT can structure the temperature both at the surface and inside the  
274 water column.

275 ~~In order to answer the aboveprevious~~ questions, we used a high-resolution model  
276 ( $1/36^{\circ}$ ) with and without explicit tidal forcing and a satellite SST product, with the aim of  
277 highlighting the impact of ~~IT~~tides on the temperature structure and quantify the associated  
278 processes. ~~These observations~~We distinguish the analysis for the two contrasted seasons (AMJ  
279 and ASO) described above. ~~The SST product, our modeling, as well as model, and~~ the methods  
280 used are described in section II. The validation of ~~some certain characteristics of the~~ barotropic  
281 and baroclinic ~~tide's characteristics as well as SST are present in tides and of the temperature~~  
282 is presented in section III. The impacts of ~~IT~~ITs on the temperature structure, the influence on  
283 heat exchange at the ~~interface between the atmosphere and the ocean interface, and finally, the~~  
284 processes involved, are analyzed in section IV. ~~SummaryThe discussion, and discussionthe~~  
285 summary of the obtained results are presented in ~~a last~~ section V and VI respectively.

## 287 II. Data and Methods

### 288 II.1. Satellite Data used: TMI SST

289 This dataset derived from Tropical Rainfall Measurement Mission (TRMM), which  
290 performs measurements using onboard TRMM Microwave Imager (TMI). The microwaves  
291 can penetrate clouds and are therefore verycrucially important for data acquisition in low  
292 latitude regions, cloudy covered during long periods of raining seasons. We use Remote  
293 Sensing Systems (RSS) TMI data products v7.1, which ~~isrepresents~~ the latestmost recent  
294 version of TMI SST. It contains a daily mean of SST with a  $0.25^{\circ}\times 0.25^{\circ}$  grid resolution ( $\sim 25$   
295 km). This SST is obtained by inter-calibration of TMI data with other microwave radiometers.

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296 The TMI SST ~~fully~~ description and inter-calibration algorithm ~~is~~are detailed in Wentz et  
297 al. (2015).

## 298 **II.2. The NEMO Model: AMAZON36 configuration**

299 The numerical model used in this study is the Nucleus for European Modelling of the  
300 Ocean (NEMOv4.0.2, Madec et al., 2019). The ~~“AMAZON36” model~~ configuration ~~designed~~  
301 ~~for our purpose is called AMAZON36 and~~ covers the western tropical Atlantic region ~~with a~~  
302 ~~1/36° horizontal grid,~~ from the Amazon River mouth to the open ocean. ~~Several Other~~  
303 configurations ~~with same grid resolution, but for the former NEMOv3.6 (Madec, 2014), exists~~  
304 ~~for the same exist in this~~ region, but either ~~includes Caribbean Sea (they have a coarse grid (1/4°~~  
305 ~~), Hernandez et al., 2016), or are not when the grid is fine (1/36°) they do not extend very far~~  
306 ~~extended to the east even used for tides study eastwards and therefore exclude most of the site~~  
307 ~~B (Ruault et al., 2020). The present current configuration is wider to capture, on their pathways,~~  
308 ~~all the internal tide generating from the Brazilian shelf. Hence, avoids these two limitations.~~  
309 ~~The grid resolution is fine (1/36°) and~~ the domain lies between 54.7°W–35.3°W and 5.5°S–  
310 10°N (Fig.1). ~~In contrast with former~~In this way, we can capture the internal tides radiating  
311 ~~from all the generating sites on the Brazilian shelf break. Unlike previous~~ configurations, we  
312 do not use multiple nested grids ~~here~~, but a single ~~fine~~ grid. The vertical grid comprises 75  
313 vertically fixed z-coordinates levels, ~~finer with a narrower~~ grid refinement ~~close to near~~  
314 the surface with 23 levels in the ~~first 100 m, and cell. Cell~~ thickness ~~reaching reaches~~ 160 m when  
315 approaching the bottom. ~~Both The~~ horizontal and vertical ~~grid~~ resolutions ~~of the grid~~ are  
316 therefore ~~acceptable fine enough~~ to resolve low-mode internal tides ~~and were. This grid~~  
317 ~~resolution has~~ already ~~been~~ used for ~~that this~~ purpose ~~(in this region (e.g., Tchilibou et al.,~~  
318 2022).

319 A third order upstream biased scheme (UP3) with built-in diffusion is used for  
320 momentum advection, while tracer advection relies on a 2<sup>nd</sup> order Flux Corrected Transport  
321 (FCT) scheme. ~~(Zalesak, 1979)~~. A Laplacian isopycnal diffusion with a constant coefficient of  
322 20 m<sup>2</sup>.s<sup>-1</sup> is used for tracers. The temporal integration is achieved thanks to a leapfrog scheme  
323 combined with an Asselin filter to damp numerical modes ~~with a~~ baroclinic time step ~~is of~~  
324 150 s). The ~~K-ε~~ turbulent closure scheme is used for the vertical diffusion coefficients.  
325 Bottom friction is quadratic with a bottom drag coefficient of 2.5×10<sup>-3</sup>, while lateral wall free-  
326 slip boundary conditions are ~~assumed prescribed~~. A time splitting technique is used to resolve  
327 the free surface, with the barotropic part of the dynamical equations integrated explicitly.

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328 We use the [2020's release of the](#) General Bathymetric Chart of the Oceans (GEBCO,  
329 2020), [see details in](#)  
330 [https://www.gebco.net/data\\_and\\_products/gridded\\_bathymetry\\_data/gebco\\_2020/](https://www.gebco.net/data_and_products/gridded_bathymetry_data/gebco_2020/)  
331 interpolated onto the model horizontal grid, with the minimal depth ~~equal~~ set to 12.8 m. The  
332 ~~ocean~~ model is forced [at the surface](#) by the ERA-5 atmospheric reanalysis (Hersbach et al.,  
333 2020). The river discharges are based on monthly means from hydrology simulation of ~~ISBA~~  
334 ~~model (the~~ Interaction Sol-Biosphère-Atmosphère [model \(ISBA](#), see description in ~~ISBA~~  
335 ~~National Centre for Meteorological Research)~~, [https://www.umr-](https://www.umr-cnrm.fr/spip.php?article146&lang=en)  
336 [cnrm.fr/spip.php?article146&lang=en](https://www.umr-cnrm.fr/spip.php?article146&lang=en)) and are prescribed as surface mass sources with null  
337 salinity, and we use a multiplicative factor of 90% based on a comparison with the HYBAM  
338 interannual timeseries (HYBAM, 2018). The model is forced at its open boundaries by (i) the  
339 fifteen major ~~high frequency~~ tidal constituents (M2, S2, N2, K2, 2N2, MU2, NU2, L2, T2, K1,  
340 O1, Q1, P1, S1, and M4) and (ii) barotropic currents, [both](#) derived from FES2014 atlas (Lyard  
341 et al., 2021). ~~At~~ [In addition to](#) the open boundaries, we prescribe [the recent](#) MERCATOR-  
342 GLORYS12 v1 [assimilation data](#) (Lellouche et al., 2018) for temperature, salinity, sea level,  
343 current velocity and derived baroclinic velocity.

344 The simulation was initialized on the 1st of January 2005, and ran for ~~40~~ 11 years until  
345 [December](#) 2015. In this study, we use [3-years](#) model outputs from [January](#) 2013 to [December](#)  
346 2015. Indeed, the model has reached an equilibrium in terms of seasonal cycle after 2 years  
347 ~~(2005-2006)~~ of run. ~~The same~~ [A twin](#) model configuration without the tides is used to highlight  
348 the influence of tides ~~and IT~~ on the temperature structure. To assess the realism of the model,  
349 we perform validation of various state variables used in this study such as the current's  
350 circulation, temperature, salinity, stratification as well as the barotropic and baroclinic ~~tide~~ [tides](#)  
351 properties.

## 352 **II.3. Methods**

### 353 **II.3.1. Barotropic/baroclinic tide separation and tide energy budget**

354 We follow Kelly et al. (2010) ~~method~~ to separate barotropic and baroclinic tide  
355 constituents: pressure, currents and energy flux. ~~No model~~ [There is no](#) separation ~~is done, then~~  
356 ~~tidal constituents obtained encompass~~ [following vertical propagation modes. Then we analyze](#)  
357 the [total energy](#) ~~off~~ for all ~~propagation's~~ [the resolved propagation](#) modes [for a given harmonic](#).  
358 Note that the barotropic/baroclinic tide separation is performed directly by the model for better  
359 accuracy, ~~however, by this way.~~ [Even though](#), it has the disadvantage of being very costly in

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360 terms of computing time. We have therefore ~~analyze only for one~~ analyzed the M2 harmonic  
 361 for the single year (2015) the M2 frequency, since. Note that M2 is the major tidal constituent  
 362 in this region, ~~representing~~ (Prestes et al., 2018; Fassoni-Andrade et al., 2023). It represents  
 363  $\sim 70\%$  of the tidal energy (Beardsley et al., 1995; Gabioux et al., 2005 ~~and Tchilibou et al.,~~  
 364 ~~2022~~).

365 The barotropic and baroclinic tide energy budget equations are obtained by ignoring as  
 366 ~~the~~ first-order approximation, the energy tendency, the nonlinear advection and the forcing  
 367 terms (Wang et al., 2016). Then, the remaining equations are reduced to the balance between  
 368 the energy dissipation, the divergence of the energy flux, ~~the dissipation~~ and the energy  
 369 conversion from barotropic to baroclinic (e.g., Buijsman et al., 2017; Tchilibou et al., 2018,  
 370 2020, ~~2022~~; Jithin and Francis, 2020; Peng et al., 2021) :

$$371 \quad D_{bt} + \nabla_h \cdot F_{bt} + \underline{CVRC} \approx 0 \quad (W.m^{-2}), \quad (1)$$

$$372 \quad D_{bc} + \nabla_h \cdot F_{bc} - \underline{CVR} \approx 0 \quad (W.m^{-2}), \quad (2)$$

373 *bt* and *bc* indicate the barotropic and baroclinic terms, *D* is the depth-integrated energy  
 374 dissipation, which can be understood as a proxy of the real dissipation since *D* may encompass  
 375 the energy loss of other tidal harmonics, non-linear terms and/or numerical dissipation (see  
 376 ~~Nugruho~~ Nugroho et al., 2018),  $\nabla_h \cdot F$  represents the divergence of the depth-integrated energy  
 377 flux, whilst CVRC is the depth-integrated barotropic-to-baroclinic energy conversion, i.e., the  
 378 amount of incoming barotropic energy ~~which is~~ converted into internal tides energy over the  
 379 steep topography, with:

$$380 \quad CVR = \langle \nabla H \cdot U_{bt}^* P_{bc}^* \rangle \quad (W.m^{-2}), \quad (3)$$

$$382 \quad F_{bt} = \int_H^\eta \langle U_{bt} P_{bt} \rangle dz \quad (W.m^{-1}), \quad (4)$$

$$384 \quad F_{bc} = \int_H^\eta \langle U_{bc} P_{bc} \rangle dz \quad (W.m^{-1}), \quad (5)$$

385 where the angle bracket  $\langle \rangle$  denotes the average over a tidal period,  $\nabla H$  is the slope of the  
 386 bathymetry,  $U_{bt}^* U_{bc}^*$  is the ~~barotropic current and velocity~~ (u, v) respectively in (x, y) directions.  
 387  $P_{bc}^*$  is the baroclinic pressure perturbation ~~both~~ at the bottom, *H* is the bottom depth,  $\eta$  the  
 388 surface elevation,  $U(u, v)$  is the horizontal velocity, *P* is the pressure, then  $F$  is the energy

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389 flux and ~~allowsemp~~ emphasizes the ~~propagation pathways~~ pathway of the ~~given tide to be~~  
 390 ~~highlighted, respective tides (external or internal).~~

391 **II.3.2. 3-D heat budget equation for temperature**

392 The three-dimensional temperature budget was computed online and further analyzed.  
 393 It is the balance between the total temperature trend and the sum of the temperature advection,  
 394 diffusion and solar radiative and non-solar radiative fluxes (e.g. ~~...~~ Jouanno et al., 2011;  
 395 Hernandez ~~atct~~ al., 2017):

396 
$$\partial_t T = -u\partial_x T - v\partial_y T - w\partial_z T - \partial_z(Kz\partial_z Kz\partial_z T) + LDF_T + FOR_Z + Numdiff \quad (^\circ C.s^{-1})$$
  
 397 (6)

398 
$$\overbrace{\hspace{10em}}^{ADV} \quad \overbrace{\hspace{10em}}^{ZDF}$$

399 Where  $\overbrace{\hspace{10em}}^{ADV^*}$   $\overbrace{\hspace{10em}}^{ZDF}$

400 Here  $T$  is the model potential temperature,  $f(u, v, w)$  are the ~~space dimensional velocity~~  
 401 ~~components~~ velocities component in the (x, y, z) [respectively eastward, northward and  
 402 upward] directions,  $ADV^*$  is the 3-D ~~temperature tendency term from the advection (routine of~~  
 403 ~~the NEMO code (from the left to right: zonal, meridional and vertical terms))~~. Note that this  
 404 term hides secondary terms that are important to define here. Hence, the total advection  
 405 tendency of temperature (ADV) is expressed as follows:

406 
$$\overbrace{\hspace{10em}}^{ADV^*} = \langle U \cdot \nabla T \rangle + \langle U' \cdot \nabla T \rangle + \langle U \cdot \nabla T' \rangle + \langle U' \cdot \nabla T' \rangle + Numdiff_{ADV} \quad (7)$$
  
 407 
$$\overbrace{\hspace{10em}}^{Non-Linear terms}$$

408 where  $U'$  is the tidal current, and  $T'$  represents the anomaly of temperature that is produced by  
 409 the tides apart the advection. When comparing the tidal and non-tidal simulation, the residual  
 410 term could come from at least three possible tidal impacts :

411 1) The result of the advection is null over a tidal cycle except in some tidal residual circulation.  
 412 In our region the residual tidal circulation is limited but might be slightly more important on  
 413 the shelf (Bessières et al., 2008).

414 2) In the nonlinear terms of the previous equation (7), temperature could be modified by other  
 415 processes than advection, which will count in the total tendency and mark the signature of the  
 416 impact of the tides.

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417 3) Finally, and it might represent the key point, in the model, the advection term leads to some  
418 diffusivity of the temperature due to numerical dissipation of the advection scheme  
419 ( $Numdiff_{ADV}$ ), in contrast to some non-diffusive advection scheme like in Leclair and Madec  
420 (2009). In our case, we are using the FCT advection scheme that includes a diffusive part  
421 (Zalesak, 1979). In previous study, this mixing has been quantified to be responsible for 30%  
422 of the dissipation (in lower resolution  $1/4^\circ$  resolution, Koch-Larrouy et al., 2008), as part of  
423 the high-frequency work of the advection diffusion. We expect here at  $1/36^\circ$  that this effect  
424 will be smaller but still non negligible. Explicit separation of these 3 impacts is beyond the  
425 scope of our study but will be discussed in the last section.

426 Furthermore,  $ZDF_z$  represents the vertical diffusion,  $LDF_T$  is the lateral diffusion,  $FOR_z$  is the  
427 tendency of temperature due to penetrative solar radiation and  $has$  includes a vertical decaying  
428 structure. At the air-sea interface, the temperature flux is equal to the non-solar heat flux (sum  
429 of the latent, sensible, and net infrared fluxes).  $FOR_z$  can modify temperature in the thin surface  
430 layer but will not be shown in the following.  $Numdiff$  corresponds to the sum of the  
431 numerical diffusion for the temperature. In this study, we assume that this last term is of second  
432 order and is not highlighted here.

### 433 H.3.3. The atmosphere-ocean net heat flux

434 The atmosphere-ocean net heat flux ( $Q_t$ ) reflects the balance of incoming and outgoing  
435 heat fluxes across the atmosphere-ocean interface (e.g.: Moisan and Niiler, 1998; Jayakrishnan  
436 and Babu, 2013), it is defined as follows:

$$437 Q_t = Q_{SW} \pm Q_{LW} \pm Q_{SH} \pm Q_{LH} \quad W.m^{-2}, \quad (7)$$

438 with from left to right: the incident solar radiative flux ( $Q_{SW}$ ), the net infrared radiative flux  
439 ( $Q_{LW}$ ), the incoming/outgoing sensible heat flux ( $Q_{SH}$ ) which depends on the temperature  
440 difference between the atmosphere and the ocean surface, and the incoming/outgoing latent  
441 heat flux ( $Q_{LH}$ ) which depends on the specific humidity difference between the atmosphere and  
442 the ocean surface. All these four components of the  $Q_t$  influence the variation of the ocean  
443 surface temperature (SST). The last two components ( $Q_{SH}$  and  $Q_{LH}$ ) have in addition a direct  
444 dependence relationship with the SST. Since IT can change the SST, we are therefore interested  
445 in knowing how it affects the net heat flux at the atmosphere-ocean interface.

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### 446 III. Model validation

447 In this subsection, we present for the M2 harmonic the assess the quality of our model's  
448 simulations by verifying whether they are in good agreement with the observations and other  
449 reference data. Firstly, for the barotropic and baroclinic tidal characteristics of the model M2  
450 tides for the year 2015, and the SST finally for the whole temperature for the period from 2013  
451 to 2015, and we verify that they agree with the different observations.

#### 452 III.1. M2 Tides in the model

453 We initially examined at the barotropic SSH (Fig.2b) of the model is compared  
454 with FES2014 (Lyard et al. 2021) (Fig.2a), and there is a good agreement in terms of both  
455 amplitude and phase. between FES2014 and the model, Fig2a and Fig.2b respectively.  
456 Nevertheless, near the coast, some differences are observed in terms of amplitude. The SSH  
457 amplitude of the model is lower ( $\sim +50$  mcm) north of the mouth of the Amazon. However,  
458 inland shorewards, and on the southern part of the mouth, the model overestimates the amplitude  
459 ( $\sim +20$  mcm and  $\sim +40$  mcm respectively). This is in terms. These biases are of the same  
460 order of magnitude like the biases in the as Ruault et al. (2020) configuration that they  
461 compared to the FES2012 product (Carrère et al., 2012) over the same region. Along the steep  
462 slope of the bathymetry (see Fig.1), a portion of the incident. The flux of the barotropic tidal  
463 energy (flowing inshore is represented by the black arrows in Fig.2c and Fig.2d, for FES2014  
464 and the model, respectively) in the presence. A fraction of stratification this energy is converted  
465 into baroclinic tidal energy: over the steep slope of the bathymetry. We compared the depth-  
466 integrated barotropic-to-baroclinic energy conversion rate ( $CVR_C$ ) between FES2014 and the  
467 model (color shading in Fig.2c and Fig.2d) and FES2014 (Fig.2e), respectively. The model  
468 does reproduce the same conversion patterns of FES2014 over the slope, but can hardly offshore  
469 over the Mid-Atlantic Ridge between  $42^\circ W-35^\circ W$  and  $7^\circ N-10^\circ N$ . This leads to an overall  
470 underestimate the  $CVR$  by of about 30%. The It is worth noting that  $C$  increases with  
471 bathymetry resolution. The latter therefore, plays a critical role in  $CVR$  (converting barotropic  
472 tidal energy into internal tides (see Niwa and Hibiya, 2011), then). Compared with FES2014  
473 ( $\sim 1.5$  km), the horizontal grid of our model is coarser ( $\sim 3$  km). Meaning that the difference in  
474 bathymetry resolution between the model ( $\sim 3$  km) and FES2014 ( $\sim 1.5$  km) could therefore  
475 explains that the difference in  $CVR$  energy conversion with FES2014. Later, another part of  
476 the barotropic energy is dissipated on the shelf by bottom friction and induces mixing from the  
477 bottom (Beardsley et al., 1995; Gabioux et al., 2005; Bessières, 2007; Fontes et al., 2008). Most

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478 of the dissipation of barotropic energy ( $D_{bt}$ ) occurs in the middle and inner shelf between 3°S–  
479 4°N (Fig.2e) in good agreement with Beardsley et al. (1995) and Bessières (2007). The  
480 remaining barotropic energy flows over hundreds of kilometers into the estuarine systems of  
481 this region (Kosuth et al., 2009; Fassoni-Andrade et al., 2023).

482 For the baroclinic tides, the critical parameter,  $\gamma = s/\alpha$ , is defined as the ratio between  
483 the slope of the bathymetry,  $s = \nabla H$  (see Fig.1), and the slope of the radiated internal wave,  
484  $\alpha = \sqrt{(\omega^2 - f^2)/(N^2 - \omega^2)}$ , with  $\omega$  the tidal frequency for a given wave,  $f$  is the Coriolis  
485 frequency and  $N^2$  represents the squared Brünt-Väisälä frequency near the bottom (e.g., Nash  
486 et al., 2007; Vic et al., 2019). The critical slope for the M2 harmonic on the slope is greater  
487 than 1.2 (not shown), consequently, the baroclinic tides (internal tidal waves) thus generated  
488 will therefore. On the slope where ITs are generated,  $\gamma > 1$ , meaning that the topography is  
489 supercritical. Consequently, the baroclinic tides, once generated, will propagate in the opposite  
490 direction to the barotropic tides, i.e., from the slope ~~towards~~ the open ocean. The, as shown  
491 by the model's baroclinic tidal energy flux of the model ( $F_{bc}$ ), black arrows in Fig.2e)2f.  $F_{bc}$   
492 highlights the existence of six main sites of FITs generation on the continental slope, two,  
493 Two of which these are more important (A and B), as shown by Magalhaes et al. regarding  
494 their higher and far extended energy flux, in good agreement with Magalhaes et al. (2016),  
495 Barbot et al. (2021) and Tchilibou et al. (2022). From these two main sites, the flow propagates  
496 offshore ITs propagate for the nearly 1000 km. On its Along the propagation path, the baroclinic  
497 tide signs at the surface in SSH. We compared this signature for the model (Fig.2h) with an  
498 estimate deduced from the altimeter tracks, produced by Zaron et al. (2019) (Fig.2g). The  
499 model is in good agreement with the altimetry observations, with an overestimation of the order  
500 of ~~~~~ +1.5 cm on the SSH maxima. It is important to note that the baroclinic SSH of the model  
501 is an average over the year 2015, whilst the observations are an average over about 20 years.  
502 This longer period may smooth the amplitude of the signal obtained from the altimetry  
503 observations. Also, the variability contained in the two averages is not the same, and this may  
504 explain some differences in the positioning and amplitude of the maxima.

505 pathways, they can dissipate their energy. Color shading in the Figure-2f shows the  
506 full model's depth-integrated internal tidal energy dissipation for ( $D_{bc}$ ). We first look at  
507 the model. The estimated local dissipation of this energy is defined as follows:

$$508 \quad P = (D_{bc} / CVR) * 100 \quad (9)$$

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509 The local dissipation is then  $q = D_{bc}/C$  (see Laurent and Garrett, 2002).  $q$  is integrated  
510 at over the embankment level slope in the same boxes A1, A2 and B (Fig.2f, see coordinates as  
511 defined in Table-2 A1 in Tchilibou et al. (2022) and provides information). This reveals that a  
512 significant part of the energy, about 30%, is dissipated locally in the different boxes, i.e., about  
513 30% (not shown). The local dissipation at the generation sites is thus in good agreement with  
514 Tchilibou et al. (2022); the latter study. The remaining part of the energy is exported offshore,  
515 and dissipates it is dissipated along the propagation path. This offshore dissipation is more  
516 extensive offshore along path A, ~500-300 km from the slope, with two patterns spaced  
517 approximately 120 by an average wavelength of 120-150 km apart corresponding to mode-1  
518 propagation mode-1, and. The offshore dissipation is less extensive offshore along path B,  
519 ~occurring around 100-200 km from the slope (Fig.2f).

520 We have presented here only the dissipation for the M2 harmonic, but in the rest of the  
521 paper, we will analyze the temperature fields on a seasonal scale and by this fact, the effect of  
522 all the tidal harmonics on the temperature are considered.

523 Another critical characteristic of internal tidal waves is their SSH imprints along the  
524 propagation pathway. We compared an estimate of this signature deduced from the altimeter  
525 tracks (Fig.2g) produced by Zaron (2019) with our model (Fig.2h). The model is in good  
526 agreement with this product, with an overestimation of the order of  $\sim +1.5$  cm on the SSH  
527 maxima. It is relevant to note the baroclinic SSH of the model is an average over the year 2015,  
528 whilst the estimate is an average over about 20 years. This more extended period may lower  
529 the amplitude of the signal obtained from the altimetry observations. Furthermore, the  
530 variability within the two datasets is not the same. This may explain some differences in the  
531 positioning and amplitude of the maxima.

532 Only the energy dissipation of the M2 tides is presented above. Elsewhere, the harmonic  
533 analysis does not consider the incoherent (non-stationary) part of the tidal energy, which has  
534 been found to be non-neglectable (Tchilibou et al., 2022). And can therefore influences the  
535 structure of the temperature. Further on, the analysis are carried out on a seasonal scale, which  
536 means that the mean temperature field obtained could result from the cumulative effect of all  
537 coherent and incoherent tidal harmonics.

### 538 **III.2. SST Validation Temperature validation**

539 For the following, it should be noted we obtained the bias between TMI SST and the  
540 two model simulations after linear interpolation of the model data into the observation grid.

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Figure 3 shows the mean SST from TMI SST, tidal simulation and non-tidal simulation over the entire 2013–2015 period. The mean 2013–2015 from TMI SST (Fig.3a), the tidal simulation (Fig.3b) and the non-tidal simulation (Fig.3c). The simulation with tides accurately reproduces well spatially the TMI SST spatial distribution of the observations (Fig.3a) both for cooling on the shelf around 47.5°W and to the southeast between 40°W–35°W and 2°S–2°N, which as shown by the weak bias,  $< \pm 0.1^\circ\text{C}$ , with TMI (Fig.3d). This cooling is almost absent for inaccurately reproduced by the non-tidal simulation (Fig.3e which exhibits a warm bias of about  $0.3^\circ\text{C}$  (Fig.3.e). To the northeast, between 50°W–54°W and 3°N–8°N in the Amazon plume, the SST of the non-tidal simulation is in better agreement with the observations, while the SST of the tidal simulation is about  $> 0.6^\circ\text{C}$  cooler than TMI SST (Fig.3d). Such a difference is very similar fits to what is obtained by other models in the same region (e.g., Hernandez et al., 2016, 2017; Gévaudan et al., 2022). The seasonal cycle of the SST of the three products for the three years 2013–2015 (Fig.3d) is obtained by interpolating the SST of far offshore, between 50°W–40°W and 6°N–10°N, both simulations on the TMI SST grid and averaging reveal a negative bias of about  $0.2\text{--}0.3^\circ\text{C}$  (Fig.3d-e). We averaged the observations and the interpolated model data in the corresponding dashed boxes around the IT generation areas (Fig.3a, b, and e) line box in the upper panels, with the shelf being masked over the depth  $\leq 200\text{ m}$  isobath masked. This location is around the ITs generation sites and on part of their pathways. Then, we compute the seasonal cycle of the three products (Fig.3f). The tidal and non-tidal simulations of the model reproduce well accurately both the seasonal cycle and the standard deviation of the observations, with a low RMSE of approximately  $\sim 10^{-2}^\circ\text{C}$  between each simulation and TMI SST (Fig.3d), which 3f). This indicates the good quality robustness of our model simulations. Nevertheless, over the seasonal cycle, it appears that between January–April and July–December, the tidal simulation is closer to the observations, while the non-tidal simulation seems slightly/moderately warmer than the observations, and in. In May–June, both simulations are colder than TMI SST (Fig.3d3f).

To gain an insight into our model along the depth, we used the mean model water properties (salinity and temperature) for the three years 2013–2015 in the same region as in Fig.3f. We compared them with the WOA2018 climatological (2005–2017) data (<https://www.ncei.noaa.gov/access/world-ocean-atlas-2018/>). We used hereabove and elsewhere  $\sigma_\theta[\rho - 1000]$  to represent the density, with  $\rho$  the water density. Figure 3g shows the Temperature-Salinity (T-S) diagram, with equal density ( $\sigma_\theta$ ) contours, for WOA2018 (black line), tidal simulation (blue line) and non-tidal simulation (red line). Both simulations

574 exhibit similar pattern with WOA2018 for deeper waters, i.e.,  $T < 17\text{ }^{\circ}\text{C}$  and  $\sigma_{\theta} > 25.6\text{ kg.m}^{-3}$ .  
575 However, there exist minor discrepancies for the surface layer waters, i.e.,  $T > 17\text{ }^{\circ}\text{C}$  and  $22.4$   
576  $> \sigma_{\theta} < 25.6\text{ kg.m}^{-3}$ . At that level, the tidal simulation better reproduces the T-S profiles. The  
577 water is slightly more eroded in the non-tidal simulation. These petty differences between  
578 WOA2018 observations and the model, even more with the tidal simulation, further  
579 demonstrate the ability of our model to reproduce the observed water mass properties.

## 580 IV. Results

581 In this section, we ~~will~~ present the influence of ~~the tides~~ on the ~~SST~~ temperature, the  
582 associated processes, and the impact on the ~~net~~ atmosphere-ocean ~~net~~ heat ~~flux from the~~  
583 ~~model's tidal and non-tidal simulations~~. The analyses were performed on a seasonal scale  
584 between April-May-June (AMJ) and August-September-October (ASO) for the three years  
585 2013-2015.

### 586 IV.1. Tide-enhanced surface cooling

587 ~~Beginning in July, a tongue of cold water ( $< 27\text{ }^{\circ}\text{C}$ ) begins to form to the southeast and~~  
588 ~~enters the central part of the plateau in August and remains there until October. Figure 3e-g~~  
589 ~~show the SST, averaged over the ASO season. The tidal simulation (Fig.3f) shows that the~~  
590 ~~upwelling cell, represented by the extension of the  $27.2\text{ }^{\circ}\text{C}$  isotherm along the slope to about~~  
591  ~~$49^{\circ}\text{W}$   $3^{\circ}\text{N}$ , extends further north than in the non-tidal simulation (Fig.3g,  $45^{\circ}\text{W}$   $0^{\circ}\text{N}$ ) which~~  
592 ~~is in better agreement with the TMI SST observations over the same period (Fig.3e).~~

593 During the first season, warm waters,  $> 27.6\text{ }^{\circ}\text{C}$ , dominate near the coast, especially in  
594 the middle shelf and in the south-east. While cold waters are present offshore north of  $6^{\circ}\text{N}$   
595 (Fig.4a-c). Off the mouth of the Amazon River, water colder than  $28.2\text{ }^{\circ}\text{C}$  spread between  
596  $43^{\circ}\text{W}$ – $51^{\circ}\text{W}$  for TMI SST (Fig.4a) and the tidal simulation (Fig.4b), whilst warmer waters are  
597 present in the same area for the simulation without the tides (Fig.4c). Figures 4d-f show the  
598 SST, averaged over the ASO season. The TMI SST observations (Fig.4d) shows an upwelling  
599 cell represented by the extension of the  $27.2\text{ }^{\circ}\text{C}$  isotherm (white dashed contour) along the  
600 slope to about  $49^{\circ}\text{W}$ – $3^{\circ}\text{N}$  towards the north-east of the region, which forms the extension of  
601 the ACT. This extension also exists in the tidal simulation (Fig.4e), whereas  $\leq 27.2\text{ }^{\circ}\text{C}$  waters  
602 are not crossing  $45.5^{\circ}\text{W}$  and remain in the southern hemisphere in the simulation without the  
603 tides (Fig.4f). Which means that a lesser upwelling cell may exist without the tides, and it is  
604 enhanced by  $-0.3\text{ }^{\circ}\text{C}$  in average due to tidal effect. The tides allow waters colder than  $27.2\text{ }^{\circ}\text{C}$  to  
605 form further north-east. Finally, we can note that the mean SST shows a very contrasting

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606 distribution between the two seasons. There are warm waters along the shelf and cold waters  
607 offshore during the AMJ season (Fig.4a-c). This is followed by warming along the Amazon  
608 plume and offshore, and a upwelling cells in the south-east (Fig.4d-f).

609 The general impact of the ~~tides~~, illustrated by ~~showing the difference in~~ SST  
610 ~~anomaly~~ between the tidal and the non-tidal ~~simulations in both seasons (Fig.4e-d, respectively~~  
611 ~~for AMJ and ASO);simulation~~, is a cooling over a large part of the study area with maxima (up  
612 to  $-0.3^{\circ}\text{C}$ ): ~~in the Amazon plume downstream of the river mouth (northeast beyond  $3^{\circ}\text{N}$ ), and~~  
613 ~~on the path of propagation of IT for both seasons.~~ $0.3^{\circ}\text{C}$  (Fig.5a-b). For ASO, tides induce a  
614 warming ( $> +0.3^{\circ}\text{C}$ ) on the shelf at the mouth of the Amazon River, (Fig.5b), while for AMJ  
615 it is a cooling ( $-0.3^{\circ}\text{C}$ ). ~~East of the same intensity (Fig.5a). That difference will be further~~  
616 ~~discussed. Out of  $45^{\circ}\text{W}$ , the tide-induced cooling shelf, the temperature anomaly~~ for each of  
617 ~~the two seasons~~season has different spatial structures, ~~but this.~~ This is probably due to a  
618 different mesoscale variability between the two seasons.

#### 619 **IV.2. Impact of the tides in the Atmosphere-Ocean Net atmosphere-to-ocean net** 620 **heat flux ( $Q_t$ )**

621 ~~Associated with the cooling~~The atmosphere-ocean net heat flux ( $Q_t$ ) reflects the  
622 ~~balance of the SST, the tide induces incoming and outgoing heat fluxes across the atmosphere-~~  
623 ~~ocean interface (see details on Moisan and Niiler, 1998; Jayakrishnan and Babu, 2013). During~~  
624 ~~AMJ, the tides mainly induce positive  $Q_t$  anomalies whose spatial structure is very similar to~~  
625 ~~the SST. Indeed, the difference in  $Q_t$  is essentially positive over the whole domain during the~~  
626 ~~AMJ season (Fig.4a) with. The average maximum values are, around  $25 \text{ W}\cdot\text{m}^{-2}$  in the plume~~  
627 ~~and the Amazon retroflection to the northeast and along A and B. During the ASO season, there~~  
628 ~~is as for the temperature at the mouth of the Amazon an inverse anomaly of ( $-25 \text{ W}\cdot\text{m}^{-2}$ )~~  
629 ~~(Fig.4b). In each season, the spatial structure of the  $Q_t$  difference almost perfectly matches that~~  
630 ~~of the SST difference. Knowing (Fig.5c). Negative SST anomalies ( $\sim -0.3^{\circ}\text{C}$ ) occur throughout~~  
631 ~~the domain in the same location. During the ASO season, at the mouth of the Amazon, there~~  
632 ~~are negative  $Q_t$  anomalies but of same magnitude as during the previous season (Fig.5d). At~~  
633 ~~this location, positive temperature anomalies ( $\sim 0.3^{\circ}\text{C}$ ) are observed (Fig.5b). Elsewhere, there~~  
634 ~~are positive  $Q_t$  anomalies and negative SST anomalies. It therefore appears that negative SST~~  
635 ~~anomalies induce positive  $Q_t$  anomalies and vice versa. Hence, the spatial structures of  $Q_t$~~   
636 ~~anomalies and SST anomalies fit almost perfectly together for the respective season. As it is~~  
637 ~~shown by the correlation among them. There is a strong negative correlation of 0.97 with a~~

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638 ~~significance of  $R^2 = 0.95$  for the AMJ season. And roughly the same intensity and sign for the~~  
639 ~~ASO season with 0.98 and 0.96, respectively for the correlation and its significance (Fig.5e).~~  
640 ~~This is consistent with the fact that the atmosphere and the underlying ocean are in a certain~~  
641 ~~equilibriumbalanced. Then, the SST cooling of the SSTinduced by theupwelled cold water~~  
642 ~~masses arriving at the surface will disturbwill try upset this balance. In response, a~~  
643 ~~consequentAs a result of this, an equivalent variation ofin the net heat flux from the atmosphere~~  
644 ~~to the ocean, will tryattempt to restore the balance. As shown by the very strong and significant~~  
645 ~~negative correlation between the difference in  $Q_t$  and the difference in SST. For the AMJ~~  
646 ~~season, we have a negative correlation of -0.97 with a significance of  $R^2 = 0.95$ , and about the~~  
647 ~~same as for the ASO season with -0.98 and 0.96 respectively for negative correlation and its~~  
648 ~~significance (Fig.9f):it.~~

649 The integral over the entire domain of the net heat flux for each season and for each  
650 simulation (Fig.4e) shows that during is shown in Figure 5f. During the AMJ season, the  $Q_t$   
651 increases from 23.85 TW (1 TW =  $10^{12}$  W) for the non-tidal simulation to 35.7 TW for the tidal  
652 simulation, i.e., an increase of +33.2 %, two times greater than that found (%). The tides are  
653 behind a third of  $Q_t$  variation. This is very large compared to what is observed elsewhere in  
654 other ITs hotspots (e.g., 15 %) by % in Solomon Sea, Tchilibou et al. (., 2020) in Solomon Sea.  
655 Thus, the tide and IT are responsible for a third of the variation in net atmosphere ocean heat  
656 flux during this season. While during the-). During the second ASO season, there is a smaller  
657 increase in  $Q_t$  of +about 7.4%,% between the two simulations, i.e., a variation from 73.03 TW  
658 to 78to 78.83 TW betweenfor the non-tidal and tidal simulations, respectively (Fig.5f).

659 Moreover, it is also noteworthy noting the considerable significant difference in  
660 integrated  $Q_t$  between the two seasons. We start from The values below are less than 36 TW  
661 induring the AMJ season to values above, whereas they are around twice as high, > 73 TW in,  
662 during the ASO season, i.e., a multiplication by a factor of at least order 2. These larger values  
663 could probably be related with the appearance of the upwelling cell described above, knowing.  
664 Given that cooler colder SST induce a stronger  $Q_t$ , these higher values are likely related to the  
665 arrival of water from ACT, which forms upwelling cells (Fig.4d-f) with a secondary tidal effect.

#### 666 **IV.3. Vertical structure of the Temperature along A internal tides pathway**

667 To further analyze the temperature changes between both simulations, we made vertical  
668 sections following the path of FITs emanating from sites A and B (blue and red line in  
669 respectively black and red line in Fig.2e). Hereunder, (i) only the transects following the  
670 pathway A will be shown, since the vertical structure is similar following pathway B especially

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671 for AMJ season, or because some processes tend to be null along pathway B during the ASO  
672 season. (ii) The mixed layer refers to a quasi-homogenous surface layer of temperature-  
673 dependent density that interacts with the atmosphere (Kara et al., 2003). Its maximum depth  
674 also known as mixed-layer depth (MLD) is defined as the depth where the density increases  
675 from the surface value, due to temperature change of  $|\Delta T| = 0.2\text{ }^\circ\text{C}$  with constant salinity (e.g.,  
676 Dong et al., 2008; Varona et al., 2019).

677 Figure 2e respectively). We will show only the results following path A, but the results  
678 are similar for path B.

679 Figure 5b shows the vertical sections of temperature for the two seasons following A.  
680 For the AMJ season, over the slope and near the coast, cold ~~waters~~ ( $< 27.6\text{ }^\circ\text{C}$ )  
681 ~~remains remain~~ below the surface at  $\sim 20\text{ m}$  for the tidal simulation (Fig. 5a) and deeper at  $\sim$   
682  $60\text{ m}$  for the non-tidal simulation (not shown), ~~it then rises~~). Then, cold waters rise to the  
683 surface more than 400 km offshore for both simulations. Although at the surface the ~~difference~~  
684 ~~in SST between the two simulations (tide no tide)~~ SST anomaly is relatively small ( $\sim -0.3\text{ }^\circ\text{C}$ ,  
685 Fig. 4e5a), because the SST is ~~likely~~ damped by the heat fluxes, further down the water column,  
686 this ~~difference~~ anomaly becomes much larger ( $> \pm 1.2\text{ }^\circ\text{C}$ , Fig. 5e). (Fig. 6b). Note that cyan and  
687 yellow dashed lines in Fig. 6b and Fig. 7b refer to thermocline for tidal and non-tidal simulations  
688 respectively. Above ~~the~~ that thermocline ( $< 120\text{ m}$ , ~~cyan and yellow lines in Fig. 5e~~), the  
689 simulation with the tides is colder by  $-1.2\text{ }^\circ\text{C}$  from the slope where the ~~FFITs~~ are generated to  
690 the open ocean following their propagation path. Conversely, below the thermocline, the tidal  
691 simulation is warmer by ~~approximately the same intensity~~ ( $1.2\text{ }^\circ\text{C}$ ) up to  $\sim 300\text{ m}$  ~~depth and~~  
692 along the ~~same~~ propagation path. (Fig. 6b). During this AMJ season, the thermocline (~~is~~  $\sim 100$   
693  $\text{m} \pm 15\text{ m}$ , ~~thick dashed black line, Fig. 5a)~~ ~~deep~~ and the ~~mixing layer~~ (~~MLD is~~  $\sim 40\text{ m} \pm 20$   
694  $\text{m}$ , ~~thick deep~~ (dashed white line, Fig. 5a)6a). They both have a very weak slope between the  
695 coast and the open ocean. Furthermore, the difference in isodensity depths between the two  
696 simulations is small (not shown), as are the depths of the thermocline ( $\sim 10\text{ m}$ , Fig. 5e) and the  
697 mixing layer (not shown), although these different depths are closer to the surface for the tidal  
698 simulation (not shown). Over the whole domain (not shown), the thermocline is deeper by  
699 about 15 m on average in the non-tidal simulation, following the propagation paths of the ~~FF~~  
700 ~~energy flow~~ITs, on the Amazon shelf and plume. While the mixing layer (Fig. 6c). Whilst MLD  
701 in the non-tidal simulation is deeper by an average of ~~13~~10 m over the shelf, 4 m on average  
702 along the ~~FFITs~~ propagation paths and close to zero in the Amazon plume. (Fig. 6d).

703 During the ASO season, cold waters ( $< 27.6^\circ\text{C}$ ) previously confined below the surface  
704 during the previous season (AMJ) ~~then~~ rise to the surface. These cold waters extend over the  
705 slope and up to about 150 km offshore in the non-tidal simulation (not shown) and up to 250  
706 km offshore in the tidal simulation (Fig. 5b7a). The  $27.2^\circ\text{C}$  isotherm ~~27.2°C only~~ reaches the  
707 surface above the slope in the tidal simulation ~~but~~ and remains below the surface ( $\sim 30$  m) in  
708 the non-tidal simulation ~~(not shown)~~. ~~At~~ This aligns with the missing of that isotherm at this  
709 location in the corresponding SST map (Fig. 4e). For the tidal simulation, ~~at~~ the surface ~~and in~~  
710 ~~the surface layers~~, the temperature ~~in the presence of the tide and IT~~ is therefore ~~cooler~~ colder  
711 than in previous season. The temperature ~~difference between the two simulations~~ anomaly in  
712 the ASO season ~~(Fig. 5d)~~ is smaller ( $< 0.4^\circ\text{C}$ , Fig. 7b) in the surface layers ( $< 40$  m) near the  
713 coast compared to the AMJ season (Fig. 5e6b). In contrast, during the ASO season, this cooling  
714 ~~reaches~~ can reach the surface and results in a colder SST along A ( $-0.3^\circ\text{C}$ ) ~~(Fig. 4d)~~ along A (5a)  
715 . The strongest cooling (~~of~~  $\sim -1.2^\circ\text{C}$ ) is deeper between 60 and 140 m depth. Below the  
716 thermocline, ~~a~~ warming (~~of~~  $\sim 1.2^\circ\text{C}$ ) is also present, but extends ~~slightly~~ less ~~(offshore~~  
717 to about 650 km) ~~offshore (Fig. 5d) compared to the AMJ season (Fig. 7b (vs.  $\sim 1000$  km,~~  
718 Fig. 5e6b)). During this ASO season, the coastward slope of the thermocline and ~~mixing~~  
719 ~~layer~~ MLD becomes somewhat steeper compared to the other season. In both simulations, there  
720 is a dip of  $\sim 80$  m ~~(i.e.,  $\sim 60$  m offshore and  $\sim 140$  m inshore)~~ and, for the thermocline (dashed  
721 black line, Fig. 7a). And a dip of  $\sim 40$  m (i.e.,  $\sim 30$  m offshore and  $\sim 70$  m inshore), respectively,  
722 for the thermocline (thick dashed black line, Fig. 5b) and the mixing layer (thick MLD (dashed  
723 white line, Fig. 5b7a). Over the entire domain ~~(not shown)~~ between the two simulations (tide-  
724 no tide), the tide deepens, the tides shallow the thermocline depth by  $\pm 6$  m on the shelf and  
725  $\pm 12$  m at the plume and far offshore along the propagation path of A (Fig. 7c). They shallow  
726 MLD in the tidal run by about 10 m along the shelf and  $\sim 4$  m along the propagation path of A-  
727 As for the mixing layer, which is deeper in the tidal run by 12 m along the shelf and along the  
728 propagation path of A. (Fig. 7d).

729 Between the two seasons, there is also a change in the vertical density gradient  
730 (Stratification) between the coast and the open sea. In the tidal simulation with tide (Fig. 5a)  
731 and without tide (not shown), during the AMJ season, the isodensities are tight near the coast  
732 and thicken towards the open sea (Fig. 6a). This means that a strong ~~vertical density~~  
733 ~~gradient~~ stratification is present near the coast and decreases towards the open sea. In contrast,  
734 during the second ASO season, the ~~vertical density gradient~~ isodensities are thicker near the  
735 coast and tight offshore (Fig. 7a). As the result of this, the stratification is weaker inshore than

736 offshore. This clearly highlights a seasonality in the vertical density gradient profile in  
737 agreement with Tchilibou et al. (2022). ~~Note that, this behavior also appears in the simulation~~  
738 ~~without the tides (not shown).~~ The ~~transect~~ ~~transects~~ of the temperature ~~differences between the~~  
739 ~~two simulations (anomaly, Fig. 5e-d)~~ ~~6b and 7b~~, show that ~~IT~~ ~~(ITs~~ and ~~probably~~ ~~likely~~ the  
740 ~~tide)~~ ~~barotropic tides~~ can influence the temperature in the ocean from the surface to the  
741 ~~bottom~~ ~~deep layers~~, with a greater effect on the first 300 meters. One question we address in  
742 this paper is to better understand what processes are at work that explain these temperature  
743 changes.

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#### 744 **IV.4. Processes** ~~What are the processes involved modifying the temperature?~~

745 To explain the observed surface and water column temperature changes, we  
746 ~~calculated~~ ~~computed and analyzed~~ the ~~trend~~ terms of the ~~temperature evolution~~ ~~heat balance~~  
747 equation (see Section II.3.2, Equation 6) for both seasons (AMJ and ASO) ~~also~~ averaged over  
748 the three years from 2013 to 2015.

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#### 749 **IV.4.1. Vertical diffusion of Temperature**

750 Figure ~~68~~ shows the vertical temperature diffusion ~~term (ZDF) for both seasons AMJ~~  
751 ~~(left panel) and ASO (right panel).~~ ~~The tendency (ZDF).~~ ZDF is averaged between 2–20 m,  
752 i.e., within the ~~mixing~~ ~~mixed-layer~~ ~~depth~~ ~~range~~. For the AMJ season, ~~the~~ ~~ZDF of the~~ ~~in~~ tidal  
753 simulation (Fig. ~~6a8a~~) shows a negative trend (cooling) in the whole domain, ~~which is.~~ ~~The~~  
754 maximum ~~on~~ ~~values~~ ( $> |0.4| \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ) are located along the ~~continental~~ slope where the ~~ITs~~  
755 are generated and on their propagation path ( $\leftarrow -0.4 \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ), with. ~~There is a~~ larger horizontal  
756 extent along A ( $\leftarrow$  of  $\sim 700$  km from the coasts) compared to B ( $\leftarrow$ , where it is  $\sim 300$  km from the  
757 coasts). ~~Over the rest of the domain. Elsewhere,~~ it remains very low ( $\rightarrow, > -0.1 \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ).  
758 For the non-tidal simulation (Fig. ~~6e8b~~), the ZDF is very weak ( $\rightarrow -0.1 \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ) over the entire  
759 domain, ~~demonstrating that internal tidal waves would be the main driver of the vertical~~  
760 ~~temperature diffusion in this region during this season. ( $\gg -0.1 \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ).~~ For the ~~second~~-ASO  
761 season, the tidal simulation (Fig. ~~6b8c~~) shows a decrease of the ZDF ~~along~~ ~~near~~ the coast ( $< 100$   
762 ~~km~~) and a strengthening offshore ~~following A,~~ ~~along A~~ compared to the previous season, but  
763 with the same cooling trend ( $< -0.4 \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ). ~~While~~ ~~along~~ ~~Along~~ B, it ~~becomes~~ ~~almost~~  
764 ~~closed~~ ~~tends~~ to ~~zero~~ ~~be~~ ~~null~~, both at the coast and offshore (Fig. ~~6b~~). ~~8c~~). In addition, the  
765 mesoscale circulation ~~intensifies~~ ~~and eddy activity intensify~~ during this season. ~~Therefore,~~ ~~to~~ ~~To~~  
766 the northeast, approximately between  $4^\circ\text{N}$ – $8^\circ\text{N}$ , and  $47^\circ\text{W}$ – $53^\circ\text{W}$ , there is a cooling on the  
767 shelf ( $\leftarrow$  ~~of~~  $\sim 0.3 \text{ } ^\circ\text{C}\cdot\text{day}^{-1}$ ) ~~with~~ ~~NBCR~~ ~~eddy~~-like patterns, ~~both~~ in the tidal simulation (Fig. ~~6b~~)

768 ~~and in 8c). The processes by which these features might arise will be examined in more detail~~  
769 ~~in the section V. Unsurprisingly, ZDF is very weak elsewhere for the non-tidal simulation~~  
770 ~~(Fig. 6d)–8d). Whatever, the ITs could be the dominant driver of vertical diffusion of~~  
771 ~~temperature along the shelf break and offshore, while the barotropic tides could prevail on the~~  
772 ~~shelf to explain the weak ZDF values.~~

773 On the vertical following A, we ~~notice an inter-~~ ~~have noted inverted ZDF values, with~~  
774 ~~mean magnitude of  $\sim |0.4|$  °C.day<sup>-1</sup>. These values are centered around the thermocline vertical~~  
775 ~~profile~~ for the simulation with ~~tidetides~~ in the two seasons AMJ and ASO (respectively  
776 Fig. 6e8e and 6f), ~~with 8f). There is a cooling tendency ( $< -0.4$  °C.day<sup>-1</sup>) trend~~  
777 ~~above the thermocline and a warming ( $> +0.4$  °C.day<sup>-1</sup>) trend below the thermocline, with an.~~ The average  
778 vertical extension ~~of~~ ~~is up to~~  $\sim 350$  m depth for the maximum values, but ~~which~~ exceeds 500  
779 m depth for the low values ( $< +0.1$  °C.day<sup>-1</sup>). ~~Over the slope, we see, as As~~ for the horizontal  
780 averages, ~~this weakening of the ZDF between the AMJ (Fig. 6e)8a and ASO (Fig. 6f)~~  
781 ~~seasons 8c), from one season to another there is a weakening of ZDF above the slope and thea~~  
782 ~~strengthening offshore. On the other hand, on the vertical, we observe towards the open sea ( $>$~~   
783 ~~200 km) that the, Fig. 8e and 8f, for AMJ and ASO respectively. Furthermore, offshore, ZDF~~  
784 ~~maxima seem to be discontinuous and spaced of about 120–150, 140–160 km during the AMJ~~  
785 ~~season (Fig. 6e), while we have a 8c) but are more continuous diffusion, for the ASO season~~  
786 ~~(Fig. 6f). This is consistent with the ZDF vertical averages (Fig. 6a–b)–8f).~~ For the non-tidal  
787 simulation, the ~~vertical temperature diffusion mean ZDF~~ tends towards  $0$  °C.day<sup>-1</sup> ~~with into be~~  
788 ~~null in the water column ocean interior~~ but remains quite large ( $> -0.2$  °C.day<sup>-1</sup>) in the thin  
789 surface layer (Fig. 6g–h).

790 During the AMJ season, the ITCZ is close to the equator and thus the trade Winds have  
791 their maximum intensity in the heart of the domain, while they migrate northward for the ASO  
792 season. As a result, more wind-generated diapycnal mixing is expected in the domain during  
793 the AMJ season compared to the ASO season. But the average value of the ZDF ( $\sim -0.2$   
794 °C.day<sup>-1</sup>) is the same between the ~~during the~~ two seasons and for both simulations (not shown)  
795 over most of the domain (except for the areas of the NBC backscatter for both simulations, on  
796 IT's generation sites and on their propagation path for the tidal simulation). This implies that  
797 the ability of the wind to generate diapycnal mixing in the underlying ocean surface layer could  
798 be limited by various oceanic processes in this region or is not well considered in the  
799 model. (Fig. 8g–h).

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800 Furthermore, it is ~~important~~worth to note that along the ~~ITs~~ propagation's pathway,  
801 the maximum of the ZDF follows the maxima of the baroclinic tidal energy dissipation (~~color~~  
802 ~~shading in Fig.2f~~). Thus, the dissipation of ~~IT generated on the continental slope generates ITs~~  
803 ~~causes~~ vertical mixing that enhances the cooling observed at the surface ~~along the coast~~. In  
804 addition, this temperature diffusion contributes to greater subsurface cooling, and warming in  
805 the deeper layers beneath the thermocline.

806 In section IV.3, the seasonality of the ~~vertical gradient of~~ stratification was highlighted,  
807 which ~~we recall~~ is stronger at the coast relative to the open ocean during the AMJ season, and  
808 reverses during the ASO season to become stronger offshore relative to the coast. This could  
809 explain why the ZDF is stronger along the slope and ~~along~~ the near-coastal pathway B during  
810 the AMJ season (Fig. ~~6a~~), ~~while 8a and 8e~~). ~~And why~~ it is weaker along the slope ~~and closed,~~  
811 ~~close~~ to zero following B, ~~and reinforce offshore of A~~ during the ASO season (Fig. ~~6b~~). ~~The~~  
812 ~~vertical gradient of density (and thus 8c and 8f). Previous studies have shown that stratification)~~  
813 ~~over the slope~~ influences the generation of ~~IT, by controlling the ITs and controls their~~  
814 propagation ~~mode of the IT that are generated (e.g.: Tehilibou et al., 2020 and Barbot et al.,~~  
815 ~~2021). We modes. Here we show here that this vertical gradient stratification also plays a role~~  
816 on the fate of these ~~ITs~~, in this case on their dissipation. The ~~vertical gradient of the~~  
817 stratification ~~thus determines could determine~~ where ~~the internal tidal ITs~~ waves dissipate their  
818 energy in the water column, ~~as mentioned by de Lavergne et al. (2020).~~

#### 819 **IV.4.2. Advection of temperature**

820 ~~The vertical (z-ADV) or horizontal (h-ADV) terms of the temperature advection~~  
821 ~~tendency are also averaged between 2-20m, for each season over the three years. Remember~~  
822 ~~that when comparing the tidal and non-tidal simulation, a residual term may arise (see equation~~  
823 ~~7 in the section II.3.2) and must be considered for the following terms, even if it is expected to~~  
824 ~~be low.~~

#### 825 **IV.4.2.a Vertical advection of Temperature**

826 ~~The vertical temperature advection (z-ADV) averaged between 2-20 m is z-ADV is~~  
827 almost ~~zero null~~ in these surface layers throughout the region (Fig. ~~7a, b, e, and 9a-d~~). For both  
828 seasons, some weak extreme values are in the northwest on the plateau between 54°W-50°W  
829 and 3°N-3°N) and are ~~offor~~ the same intensity between the two simulations. ~~With a slight~~  
830 ~~intensification when moving to the ASO season (-0.3 °C.day<sup>-1</sup>). The z-ADV is zero at the IT~~  
831 ~~generation sites and along their propagation pathways, so the almost zero difference between~~

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832 the two simulations for each season shows that the IT- with and barotropic tide do not without  
833 tides. This result suggests that, overall, the tides fail to generate vertical temperature advection  
834 within these ocean surface layers. The z-ADV does not contribute to the temperature change  
835 in the surface layers of the ocean, and therefore does not influence the cooling observed from  
836 the surface on the SST. On the other hand, At deeper, under the mixed layer and close to the  
837 thermocline, the depth, z-ADV structure is more marked.

838 tendency term is non negligible, and clearly higher in tidal simulation than in non-tidal  
839 one. Vertical sections following A (Fig. 7e, f, g and 9a-h) show an intensification of z-ADV  
840 of about  $\pm 0.8^{\circ}\text{C}\cdot\text{day}^{-1}$  located below the mixing layer MLD (magenta dashed line) and  
841 near seems to be centered around the thermocline between 80 and (black dashed line), with a  
842 vertical extension from 20–200 m ( $\pm 0.8^{\circ}\text{C}\cdot\text{day}^{-1}$ ). During the AMJ season, over the vertical,  
843 the depth, z-ADV is stronger in tidal simulation during the both seasons (Fig. 9e-f) and mainly  
844 presents sparse extrema offshore ( $> 300$  km) for the non-tidal simulation (Fig. 9g-h). For the  
845 simulation with the tides, z-ADV appears to be rather dominated by a cooling trend. The tidal  
846 simulation (Fig. 7e) shows a cooling trend ( $-0.8^{\circ}\text{C}\cdot\text{day}^{-1}$ ), with a marked hotspot on the  
847 slope where the IT are generated, with an average vertical extension from 20 to 200 m depth.  
848 Then offshore, two cooling followed by other hotspots ( $-0.8^{\circ}\text{C}\cdot\text{day}^{-1}$ ) followed by a weaker  
849 one ( $-0.3^{\circ}\text{C}\cdot\text{day}^{-1}$ ) offshore. These extreme values are spaced about 120–150 km apart,  
850 interspersed by two warming zones, respectively  $+0.6^{\circ}\text{C}\cdot\text{day}^{-1}$  and  $+0.3^{\circ}\text{C}\cdot\text{day}^{-1}$  from the  
851 coast. For the non-tidal simulation (Fig. 7g), the z-ADV is much less intense with lower values  
852 ( $\pm 0.3^{\circ}\text{C}\cdot\text{day}^{-1}$ ) near the coast until  $\sim 300$  km offshore, followed by a cooling hotspot ( $-$   
853  $0.8^{\circ}\text{C}\cdot\text{day}^{-1}$ ) between 300 km and 500 km, i.e., the imprint of mode-1 propagation wavelength  
854 as for the baroclinic tidal energy dissipation (Fig. 2f). For the both simulations, (Fig. 9e-h), the  
855 extreme values appear to be centered around the mean depth of the thermocline (thick black  
856 outline) and do not cross the mixing layer depth (thick magenta outline). They are on average  
857 located between 40 m and 200 m depth and are located within the narrow density ( $\sigma_{\theta}$ ) contours  
858  $[23.8\text{--}26.2\text{ kg}\cdot\text{m}^{-3}]$ , i.e., they follow the position of the maximum vertical density gradient  
859 between isodensity anomalies  $23.8\text{--}26.3\text{ kg}\cdot\text{m}^{-3}$ .

860 For the ASO season, the simulation with tide (Fig. 7f) still shows the same cooling  
861 intensity on the slope, although deeper ( $\sim 60$  m and 250 m), as well as offshore with this time  
862 the third cooling hotspot more intense ( $-0.8^{\circ}\text{C}\cdot\text{day}^{-1}$ ) than during the AMJ season (Fig. 7e).  
863 The non-tidal simulation (Fig. 7h) shows a less intense z-ADV ( $< \pm 0.1^{\circ}\text{C}\cdot\text{day}^{-1}$ ) near the slope,

864 and a little stronger offshore ( $\sim \pm 0.3 \text{ }^\circ\text{C}\cdot\text{day}^{-1}$ ) between 300 and 600 km from the slope,  
865 although less intense than the previous season.

866 As in the AMJ season, the extreme values of  $z$ -ADV follow the vertical density gradient  
867 in both simulations. During the ASO season, the maximum of the vertical density gradient is  
868 between 23.8 and 26.2  $\text{kg}\cdot\text{m}^{-3}$  and is deeper at the coast and is closer to the surface offshore the  
869 stratification, namely, the pycnocline.

870 Thus, the extreme values of  $z$ -ADV are located a little deeper, between 80 and 300 m.  
871 Furthermore, for the non-tidal simulation and during both seasons, the position relative to the  
872 coast of the extreme values are shifted regarding those ones of the same polarity in the  
873 corresponding tidal simulation, which means that the presence of the IT and the tides could  
874 modify the intensity and patterns of the  $z$ -ADV produced by the other oceanic processes.

875 In addition, we averaged the  $z$ -ADV between deeper depths above the thermocline  
876 depth (20–70 m) and below the thermocline (148–250 m) depth for all simulations and both  
877 seasons (not shown). This allows to highlight the NBC's pathway through the extreme values  
878 of the  $z$ -ADV close to the coast and its retroflection offshore to the northeast for both  
879 simulations, but also the propagation of the IT from the coast to the open sea from the two main  
880 sites A and B for the simulation with tide. Thus, we see that the IT and the general circulation  
881 are the main drivers of the vertical temperature advection in the subsurface and deeper layers  
882 in this region.

#### 883 **IV.4.3.2.b Horizontal advection of temperature**

884 Horizontal advection of temperature ( $h$ -ADV) is defined as the sum of the zonal ( $x$ -  
885 ADV) and meridional ( $y$ -ADV) terms of temperature advection. The  $h$ -ADV is also averaged  
886 between 2–70 m tendency. As for each simulation during both seasons (Fig. 8a, b, c and d). As  
887 obtained with  $z$ -ADV, horizontal advection the mean of temperature  $h$ -ADV tends to zero  
888 null over the entire domain in the surface layers for both seasons in both simulations, with  
889 (Fig. 10a-d). Nevertheless, some weak extreme values located are in the northwest of the  
890 plateau between 54°W–50°W and 3°N–3°N that. That intensify during the ASO season ( $\sim$  in  
891 both simulations,  $\sim \pm 0.2 \text{ }^\circ\text{C}\cdot\text{day}^{-1}$ ) (Fig. 8b-10c and 10d for the tidal and d). Along the slope  
892 between non-tidal simulations respectively. During AMJ,  $h$ -ADV is slightly stronger,  $\sim 0.1$   
893  $^\circ\text{C}\cdot\text{day}^{-1}$ , around sites A and B during the AMJ season, the  $h$ -ADV generates a small warming  
894 ( $\sim +1 \text{ }^\circ\text{C}\cdot\text{day}^{-1}$ ) that is more pronounced in the tidal simulation (Fig. 8a-10a) than in the non-  
895 tidal simulation (Fig. 8e), and thus 10b). This appears to be related to the FITs generated along

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896 the slope. On the other hand, the small difference between the two simulations in the surface  
897 layers shows that the ~~tidal processes (IT and barotropic tide)~~tides hardly generate ~~horizontal~~  
898 ~~temperature advection. The low values observed here clearly show that the~~h-ADV. Then, h-  
899 ADV could not influence the cold-water tongue observed over the surface SST (Fig.3e-g)  
900 during the ASO season (Fig.4d-f). This result aligns with Bessières et al. (2008), which had  
901 ~~previously shown that the tidal residual mean transport is null in the upwelling region in the~~  
902 ~~south-east and low (<|0.1| Sverdrup) over the whole shelf.~~

903 Along the vertical following A, ~~the~~h-ADV maxima remain essentially confined below  
904 the ~~mixing mixed-layer~~depth, with much more intense values in the tidal simulation (Fig.10e-  
905 f) compared to the non-tidal simulation. ~~The (Fig.10g-h).~~ h-ADV contributes to both warming  
906 and cooling of the temperature (~~of  $\sim \pm 0.4$  °C.day<sup>-1</sup>~~) from the slope to more than 500 km  
907 offshore, ~~with an~~. During both seasons, the average vertical extension lies between the surface  
908 and 400 m depth for the tidal simulation (Fig.8e and f) and a little less extended between 20-  
909 300 m depth for the non-tidal simulation (Fig.8g and h). As for z-ADV, h-ADV is also stronger  
910 within the pycnocline. For the tidal simulation, there is a warming above the slope (0.4  
911 °C.day<sup>-1</sup>) reaching the surface in both seasons. This vertical excursion is observed elsewhere  
912 for ZDF and z-ADV, and it is probably a marker of local dissipation of ITs at their generation  
913 site. The local dissipation of ITs clearly affects both advection and vertical diffusion of the  
914 temperature. But there are very low values along the slope when averaging h-ADV or z-ADV  
915 between 2-20 m and much more strong values for the ZDF. This means that the energy  
916 dissipated by ITs is mostly transferred to mixing.

917 Furthermore, unlike ZDF and z-ADV, the (horizontal) location of h-ADV maxima  
918 mismatch the dissipation hotspots. It is difficult to identify the wave-like characteristic of the  
919 propagation of ITs in h-ADV. This probably means that ITs hardly induce any horizontal  
920 motion of water mass. We can therefore deduce that the observed increase in h-ADV is mainly  
921 because of the barotropic tides.

## 922 **V. Discussion**

### 923 **V.1. Vertical advection tendency term**

924 Results showed that z-ADV is stronger in the deeper layer, below the MLD and within the  
925 pycnocline (Fig.9e-h). As mentioned above, this tendency term includes both nonlinear effect  
926 between the temperature and the currents and numerical dissipation of the diffusive part of  
927 advection scheme working at high frequencies. The location of the maxima of the vertical

928 advection tendency at the shelf break and along the ITs propagation pathway and its negative  
929 sign, suggest that the diffusive part of the advection scheme might be the dominant process  
930 compared to nonlinear effects, as the velocity of the (mode-1) internal tidal waves is maximum  
931 in the thermocline where exactly z-ADV term is working harder.

## 932 **V.2. On the role of advection in coastal upwelling**

933 To explain the cooling of the SST at the surface, Neto and da Silva (2014) indicated  
934 that the steady flow of the NBC induces northward transport of water masses. This transport is  
935 in turn offset by a vertical advection of cool water towards the surface. We demonstrate with  
936 our model that the vertical advection hardly modifies the SST. But it is rather working below  
937 the mixed layer (Fig.9e-h). The tides-induced vertical diffusion (mixing) extends from the  
938 mixed-layer to deeper layers (Fig.8e-f). It is therefore possible for the vertical mixing to bring  
939 up to the surface the water masses that are advected into the layers below the mixed layer. The  
940 change in SST and temperature within the mixed-layer can then be influence in first order by  
941 (i) the vertical diffusion of temperature and secondary by (ii) a cross effect between the latter  
942 and the advection (vertical and horizontal) of temperature that mainly takes place below the  
943 mixed-layer.

## 944 **V.3. The mode-1 wave-like patterns in the vertical terms of the heat budget** 945 **equation**

946 Along the vertical and toward the open ocean, both ZDF and z-ADV tendencies are found  
947 to have a wave-like structure. For z-ADV, patches are spaced apart by about 120–150 km and  
948 140–160 km for the AMJ and ASO seasons respectively. Whilst for z-ADV, this wavelength  
949 is about 140–160 km during both seasons. The h-ADV the AMJ season and more continuous  
950 patches for the ASO season. The wavelength ranges found in temperature tendency terms (3T)  
951 are slightly wider (~ 10–20 km, for z-ADV in ASO season and for ZDF) than the purely  
952 dynamic tidal coherent wavelength (~ 120–150 km. see section III.1). The difference can be  
953 understood as the effect of incoherent ITs, i.e., ITs that are deviated or diffracted by the currents  
954 and/or eddies, for which dissipation occurs around where coherent ITs dissipate. They are  
955 uncaptured by the harmonic analysis. Hence, the total (coherent + incoherent) dissipation  
956 pattern of ITs could be wider than in Figure 2f. When integrating 3T over the season, this  
957 cumulative effect is considered and therefore leads to diffuse patterns and wider wavelength.  
958 This diffusive effect increases during the ASO season when both background circulation and  
959 eddy activity increase.

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960 Recently, de Maccdo et al. is low in the surface layers (2–20 m) but maximum in the  
961 subsurface where the (2023) gave a detailed description of ISW in this region. They showed an  
962 intensification of ISW occurrences along A and B pathways, whose inter-packet distance  
963 corresponds to the wavelength of mode-1 ITs. These ISW packets are also colocalized  
964 (horizontally) with the deeper 3T patches. Our results are therefore consistent with the  
965 observations of the latter study regarding the localization of IT dissipation, particularly where  
966 they can generate ISW.

#### 967 **V.4. Tidal impact at the mouth of the Amazon River and on the southern** 968 **shelf: two main competitive processes**

969 Depending on the season, the mean SST anomaly [Tide – No-Tide] at the mouth of the  
970 Amazon and southeast of the plateau is either negative (AMJ, fig.5a) or positive (ASO, fig.5b).  
971 What we found can be explained by a combination of processes. Note that seasonal variations  
972 in solar radiation, river flow and stratification is stronger. We over the shelf can also see for play  
973 significant roles.

974 In the tidal simulation in both seasons a warming above the slope that reaches without the  
975 surface tides, there is a strong coast-parallel current exiting northwesterly the mouth of the  
976 Amazon River (black arrows in Fig.11a, 11b; Ruault et al., 2020) with an average intensity of  
977 about  $+0.4\text{ }^{\circ}\text{C}\cdot\text{day}^{-1}$  ( $\text{Fig.} > 0.5\text{ m}\cdot\text{s}^{-1}$  in the first 50 meters (color shading in Fig.11a, 11b). When  
978 including the tides in the model, the latter study had shown that there is an increase in the  
979 vertical mixing in the water column due to stratified-shear flow instability. They then show  
980 that this weakens the coast-parallel current and favors cross-shore export of water (color  
981 shading in Fig.11c, 11d), which is then diverted to the north-west (black arrows in Fig.11c,  
982 11d). We can therefore establish that there are at least two processes at work in producing SST  
983 anomalies: (i) vertical mixing and (ii) horizontal transport, reflected respectively by ZDF and  
984 h-ADV. We then looked at the latter two processes along the vertical following the cross-shore  
985 transect (C-S) defined in Figure 10b. Hereinafter, inner mouth refers to the part of the transect  
986 before 200 km, whereas outer shelf refers to the part beyond.

987 During the AMJ season, in the inner mouth, river flow dominates and tide-induced vertical  
988 mixing in the narrow water column leads to warming and deepening of the thermocline (cyan  
989 and black lines in Fig.12a-b). On the outer shelf, this mixing in the thicker water column leads  
990 to cooling above the thermocline and warming below (Fig.12a). Which in turn extends across  
991 the shelf and along the pathways of ITs as shown in section IV.4.1 (see Fig.8a, 8c, and 8e and

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992 f) but remains below the surface (~ 20 m) in the -f). At the same time, the SST on the shelf is  
993 somewhat homogeneous (see Fig.4a-c) and solar radiation is lower than 190 W.m<sup>-2</sup> (not  
994 shown). As a result, waters of similar temperature are advected horizontally, i.e., the h-ADV is  
995 low (Fig.12b). Thus, for the first season, vertical mixing seems to be the dominant process  
996 explaining the average negative SST anomaly on the plateau.

997 For the second season, solar radiation on the shelf rose sharply with an average value of 60  
998 W.m<sup>-2</sup> compared with the previous season (Fig.12c). The average depth of the thermocline  
999 deepens offshore (cyan and black lines Fig.12d and 12e). Here, mixing leads to warming in the  
1000 thin surface layer (< 2m, Fig.12d). In contrast to AMJ, there is a significant horizontal variation  
1001 in SST on the plateau (see Fig.4d-f). The NBC is stronger and can influence transport over the  
1002 shelf (Prestes et al., 2018). Even it is small, the mean tidal residual transport is added and should  
1003 be taken into account (Bessières et al., 2008). Warm waters can therefore be advected across  
1004 the shelf. Consequently, h-ADV is stronger and positive (Fig.12e) and plays a greater role in  
1005 the fate of SST. For this season, ZDF and h-ADV add to explain the positive SST anomaly on  
1006 the shelf.

1007 From AMJ to ASO, we can note the deepening of the thermocline depth on the outer shelf.  
1008 This was previously highlighted by Silva et al. (2005) from REVIZEE (Recursos Vivos da  
1009 Zona Econômica Exclusiva ) campaign data. This is a further contribution to the validation of  
1010 our model in the section III.2.

## 1011 V.5. Tidal impact in the NBC retroflection area

1012 To the north-west of the domain [3°N–9°N and 53°W–45°W], in the surface layers (2–  
1013 20m), eddy-like or circular patterns exist in ZDF during the ASO season for the simulation  
1014 including tides (Fig.8c). It should be remembered that during this season the NBC intensifies  
1015 and retroflects, and strong eddy activity takes place there. We therefore assume that they may  
1016 be the driving force behind these ZDF patterns. However, it is not yet clear how these  
1017 mesoscale features produce vertical mixing. They may be involved either by fronts or trapping  
1018 the internal tidal waves.

1019 1) **Fronts:** they exist in such a intensively active mesoscale region. They are associated  
1020 with significant vertical mixing (see Chapman et al., 2020). We therefore looked at the  
1021 horizontal temperature gradient ( $\nabla T$ ) averaged over the same depth range (2–20m) as  
1022 the ZDF (Fig.8a-d). During the AMJ season, it is on average equal to  $4 \cdot 10^{-2} \text{ }^\circ\text{C}/10 \text{ km}$ .  
1023 As expected, it does not reveal any circular fronts for the two simulations (Fig.13a-b)

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1024 since mesoscale activity is low. Secondly, the horizontal gradient of the temperature  
1025 increases during the ASO season [ $> 5 \cdot 10^{-2} \text{ }^\circ\text{C}/10 \text{ km}$ ] in the north-west and exhibits  
1026 circular and filamentary fronts in both the non-tidal simulation (Fig.8g and h). This  
1027 vertical excursion that is observed elsewhere for ZDF and z-ADV is a marker of local  
1028 dissipation of IT at their generation site on the slope, which clearly affects both vertical  
1029 diffusion and advection of the temperature. But we have almost null values along the  
1030 slope when averaging h-ADV or z-ADV between 2-20 m and much more strong values  
1031 for the ZDF (Fig.6a, b, e and d). This means that the IT's energy loss is mostly  
1032 transferred to the turbulent scale (mixing). Furthermore, unlike the ZDF (Fig.6e) and  
1033 z-ADV (Fig.7e and f), on the vertical it is difficult to identify a wave structure  
1034 characteristic of IT propagation in the h-ADV (Fig.13c) and tidal (Fig.13d) simulations.  
1035 Therefore, one would expect to see the same circular patterns in the ZDF for both  
1036 simulations. This is not actually the case (see Fig.8c and 8d) and invalidates this  
1037 statement. Furthermore, these values are at least three times smaller compared to other  
1038 oceanic regions (e.g., Kostianoy et al., 2004 and Bouali et al., 2017), meaning that these  
1039 fronts are less pronounced.

1040 **2) Trapping internal tidal waves:** stronger mesoscale activity which occurs during this  
1041 season implies more interaction between the background circulation and ITs (Buijsman  
1042 et al., 2017 and Tchilibou et al., 2022). The NBC flows along the coast and crosses the  
1043 sites where ITs are generated (see schematic view in Fig.1). This means that ITs can be  
1044 trapped and advected along the NBC pathway. When this current destabilizes and  
1045 retroreflects in the north-west, these trapped waves dissipate and therefore generate  
1046 vertical mixing. This hits the high fraction of the incoherent ITs found here (Tchilibou  
1047 et al., 2022). But quantifying the impact on temperature of such a wave-mean flow  
1048 interaction process requires further analysis and is beyond the scope of this study.

1049 Nevertheless, we believe that this second process could be the main cause of vertical  
1050 diffusion of temperature in that region. Thus, from the section V.3 and the latter, we can  
1051 conclude that incoherent ITs represent a significant part of the total energy of internal tides.  
1052 But remains to be quantified in future work. In addition, in parallel with coherent ITs, they  
1053 might play a critical role on the fate of the temperature in this region.

1054 ~~V.VI. Summary and Discussions~~

1055 In this paper, ~~the impact of internal tidal waves (IT) on temperature, off the Amazon,~~  
1056 ~~especially on the surface and on net heat fluxes is explored through outputs of twowe used~~ twin  
1057 oceanic simulations (with and without tides) from a realistic model. ~~The to explore the impact~~  
1058 ~~of internal tidal waves (ITs) on temperature and associated processes. The impact on the~~  
1059 ~~atmosphere-to-ocean net heat fluxes is also covered.~~

1060 ~~The AMAZON36 configuration, based on the 1/36° resolution NEMO model,~~ can  
1061 reproduce the generation of ~~internal tides (IT), i.e., the conversion of energy from barotropic~~  
1062 ~~to baroclinic tides, ITs~~ from two most energetic sites A and B, in good agreement ~~with~~ previous  
1063 studies ~~(Magalhaes et al., 2016 and Tehilibou et al., 2022). As for dissipation, the model. The~~  
1064 ~~model well~~ reproduces ~~30% local dissipation, the rest propagating offshore from the different~~  
1065 ~~generation sites, the two main ones being A and B (Fig.2e). During their propagation, the IT~~  
1066 ~~dissipate most of their energy after~~ the local, on-shelf, and offshore dissipation of ITs with two  
1067 beams of mode-1 ~~reflection~~propagation (120–150 km), ~~that is). This dissipation occurs,~~ less  
1068 than 300 km from the slope. ~~Then, we assess the ability of the model to reproduce temperature~~  
1069 ~~structure. The simulations including tides is in better agreement with SST observations and~~  
1070 ~~better reproduce water mass properties along the vertical,~~

1071 ~~TheOur~~ analyses ~~arewere,~~ based on ~~data from~~ three years (2013 to 2015), ~~data averaged~~  
1072 over two seasons, AMJ (April-May-June) and ASO (August-September-October) ~~which). That~~  
1073 are highly contrasted in terms of stratification, ~~background~~ circulation and EKE. ~~Results show~~  
1074 ~~that for both seasons, the tides create SST cooling of about 0.3 °C in the plume of the Amazon~~  
1075 ~~offshore and along the paths of propagation A and B of ITs. During ASO, the cold waters (<~~  
1076 ~~27.5 °C) of the Atlantic Cold Tongue (ACT), enter our domain along the coast, and are affected~~  
1077 by ~~IT~~the tides. ~~This enhances that seasonal upwelling and tides, which, leads to a cooler~~  
1078 ~~seasonal upwelling.~~

1079 The impact of the tides on temperature was assessed by comparing our twin simulations  
1080 with and without tides for each season. For ASO and AMJ, the tides create a cooling of SST of  
1081 the order of ~~— 0.3 °C~~ in the plume of the Amazon offshore and along the paths of propagation  
1082 A and B of the internal tide. Concerning the Amazon shelf, the tides induce a warming (~~— +0.3~~  
1083 ~~°C) in ASO and a cooling (of — 0.3 °C) in AMJ. These cooler/warmer waters~~SST. Over the  
1084 ~~Amazon shelf, the tides induce the same magnitude cooling in AMJ and in turn induce an~~  
1085 ~~opposite anomaly (warming) in ASO. These cooling/warming are responsible in the same~~  
1086 location for an increase/decrease in the net heat flux from the atmosphere to the ocean, ~~leading~~

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1087 ~~to ( $Q_t$ ). However, the overall effect of the tides is an increase ( $Q_t$ ) of  $+Q_t$ , which lies between~~  
1088 ~~[33.2% in AMJ and of  $-7.4\%$  in ASO between runs with and without tides. In the subsurface,~~  
1089 ~~above the thermocline ( $<120$  m), the IT and tides induce a stronger cooling than on the surface~~  
1090 ~~of about  $-1.2$  °C and an associated warming of about  $+1.2$  °C under the thermocline ( $>120$~~   
1091 ~~m to 300 m).~~

1092 ~~By  $\%$  from AMJ to ASO. And can be larger than what obtained elsewhere (e.g., in the~~  
1093 ~~Solomon Sea). In such a region with large atmospheric convection (marked by the ITCZ), when,~~  
1094 ~~increasing the atmosphere-to-ocean net heat flux ( $Q_t$ ), the IT and tides might reduce the cloud~~  
1095 ~~convection ~~in~~ into the atmosphere, as we are in an intertropical convergence zone (ITCZ).~~  
1096 ~~Impact on overall atmospheric circulation and precipitation is expected to be significant, as~~  
1097 ~~previously shown in other regions such as Indonesia (Tidal-induced cooling of  $-0.3$  °C can~~  
1098 ~~reduce precipitation by  $20\%$ , see (Koch-Larrouy et al., 2010).~~

1099 ~~Therefore, it becomes important to note that the interannual or even climatic scale~~  
1100 ~~evolution of internal ~~this~~ tidal waves activity must be considered to better understand the future~~  
1101 ~~evolution of the global effect on the climate. Especially since thanks to the CanESM5 global~~  
1102 ~~climate model, Yadidya and Rao (2022) might have just shown that in the Andaman Sea and~~  
1103 ~~Bay of Bengal, towards the end of this century, a key importance, for both optimistic and~~  
1104 ~~pessimistic SSP scenarios, the increase in depth-averaged stratification will result in an increase~~  
1105 ~~in IT activity in these two regions. Knowing that the continental slope of northern Brazil is a~~  
1106 ~~place of high generation of IT, which therefore depends on stratification, it is hereby critical to~~  
1107 ~~understand how IT activity will evolve in the coming decades in order to better anticipate the~~  
1108 ~~climate, and thus better adapt public policies at national and international levels to the global~~  
1109 ~~context of the future, taking the climate change into account (Yadidya and Rao, 2022).~~

1110 ~~Another objective of our study was to understand ~~In~~ the processes responsible for these~~  
1111 ~~temperature changes. For this, we subsurface, above the thermocline ( $<120$  m), the tides induce~~  
1112 ~~a stronger cooling ( $\sim 1.2$  °C) than at the surface. And an associated warming of the same~~  
1113 ~~magnitude under the thermocline ( $>120-300$  m). We analyzed the ~~trend~~ terms of the~~  
1114 ~~temperature evolution ~~heat budget~~ equation. Where IT dissipate their energy, there is an intense~~  
1115 ~~vertical mixing to identify to processes that ~~generates~~ modify the temperature. We found that~~  
1116 ~~the vertical diffusion of temperature (ZDF) ~~←~~ is mainly caused by the dissipation of the tides.~~  
1117 ~~Horizontal (h-ADV) and vertical (z-ADV) advection can be driven by non-tidal processes but~~  
1118 ~~increase when including the tides in the model.~~

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1119 Over the shelf, barotropic tidal mixing increases ZDF ( $> -0.4$  °C.day<sup>-1</sup>) according to  
1120 pathway A and explain the cooling of the water column in AMJ season. During the second  
1121 season, it combines with h-ADV and to a lesser extent according to pathway B, stronger at  
1122 shore than offshore during AMJ and inverse during ASO. The ZDF is the only process that  
1123 reaches the surface layer, and then appears in first approximation to be the main process  
1124 contributing to the surface cooling observed on SST. The atmospheric heat flux terms ( $FOR_z$ )  
1125 could also modify this SST but was not highlighted in this study. The cause a warming. Off the  
1126 shelf, the (baroclinic) mixing takes place up from the slope to about 800/700 km off the slope  
1127 following the path A, and 300 km following B. It is also responsible on a seasonal scale, but  
1128 also daily for a negative average variation (cooling) of temperature the path B. That mixing  
1129 induces ZDF with values of about  $-0.4$  °C.day<sup>-1</sup> above the thermocline, and a warming of  $+0.4$   
1130 °C.day<sup>-1</sup>, which is the main process in the upper layer above the mixed layer. But could  
1131 combine with advection terms (z-ADV and h-ADV) to explain the temperature changes below  
1132 350 m and decreasing to  $+0.1$  °C.day<sup>-1</sup> around 500 m depth the mixed layer. Along ITs  
1133 propagation pathways, some ZDF and z-ADV patches follow the dissipation hotspots of the  
1134 ITs, i.e., they exhibit the mode-1 propagation of ITs.

1135 IT propagation induce vertical advection of water masses around the thermocline level,  
1136 which has the effect of producing a subsurface mean temperature cooling ( $-0.8$  °C.day<sup>-1</sup>) at  
1137 a depth varying between 20–200 m AMJ and 60–250 m in ASO, with three extreme values off  
1138 the coast spaced approximately 120–150 km along of the pathway A, which seem to follow the  
1139 dissipation patterns, and thus correspond to the horizontal scale of the mode-1 propagation of  
1140 IT. This study highlights the key role of ITs in creating intensified mixing which is important  
1141 for temperature structure. Other processes such as zonal and meridional advection of  
1142 temperature analysis we performed with our simulations show that this mixing can also induce  
1143 temperature change in subsurface and deeper layers. Finally, the horizontal (zonal and  
1144 meridional) advection of temperature in this region is more related to the general circulation  
1145 (NBC, mesoscale) but is increased by tides and IT.

1146 Thus, it is the combination of these different processes that explains the temperature  
1147 change in the water column in this region. Furthermore, in order to explain the cooling of the  
1148 SST at the surface, Neto et al. (2014) indicated that the northward transport of water masses by  
1149 the constant circulation of the NBC was compensated by a vertical advection of colder water  
1150 masses towards the surface. We now know that this vertical advection process fails to modify  
1151 the SST but is rather limited below the mixing layer. The same is true for zonal and meridional

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1152 ~~advection of temperature (which form horizontal advection). It should be remembered that~~  
1153 ~~vertical diffusion extends from the surface, through the mixing layer, into the deep layers. It is~~  
1154 ~~therefore possible that water masses cooled by both vertical and horizontal advection below~~  
1155 ~~the mixing layer can be recovered and transported vertically to the surface by the effect of~~  
1156 ~~vertical mixing. The change in SST and temperature above the mixing layer then comes from~~  
1157 ~~(i) vertical diffusion of temperature and (ii) a combination of this vertical diffusion and the~~  
1158 ~~advection (vertical and horizontal) of temperature that takes place below the mixing layer.~~

1159 This study focuses on temperature, but other analyses we have done on impacts salinity  
1160 show that IT also affects the haline structure of the ocean in this region. A future work would  
1161 be to look at the impact on salinity, which is also a key parameter in the functioning of the  
1162 ~~ocean. Furthermore, they might be seen~~ as in-exchanges with the atmosphere, and thus can play  
1163 a role on the climate. In addition, internal waves can also influence the biogeochemical cycles  
1164 of elements and the entire marine ecosystem, since they can induce a source of nutrient uptake  
1165 at tidal frequency and thus participate in structuring can have an impact on the spatial  
1166 distribution of phytoplankton and zooplankton, and ~~in consequence of the rest of the food chain~~  
1167 ~~that depends on them.~~

1168 It would also be important to compare the results of our model with fields observations.  
1169 Two high frequency PIRATA anchorages have been installed offshore at the extremity of our  
1170 region between 35°W–38°W and 0°N–5°N (see ~~therefore on the entire food chain (Sharples et~~  
1171 ~~al., 2007, 2009; Xu et al., 2020). These other impacts can be studied through a combined model-~~  
1172 ~~in situ data approach. A long-term PIRATA (Prediction and Research moored Array in the~~  
1173 ~~Tropical Atlantic) mooring data are available for this goal (Bourlès et al., 2019) and could be~~  
1174 ~~used for this purpose.~~ In addition, recently in late 2021, the “AMAZOn MIXing  
1175 (“AMAZOMIX”)” campaign ~~entirely~~ took place in this region. Among other things, this  
1176 campaign was dedicated to IT (27 August and 8 October 2021) will provide a better ITs. It  
1177 provided a huge set of data, with the aim of understanding ~~of the impact of IT on the marine~~  
1178 ~~environment in this region. their impact on marine ecosystems (see details in~~  
1179 ~~<https://en.ird.fr/amazomix-campaign-impact-physical-processes-marine-ecosystem-mouth->~~  
1180 ~~amazon).~~ In the meantime, a coupled physical/biogeochemistry simulation (NEMO/PISCES);  
1181 ~~is~~ currently under analysis; ~~and~~ will begin to answer these crucial questions ~~of the impact of~~  
1182 ~~internal waves on biogeochemistry.~~

1183 Finally, ~~in this first part, we have~~ focused hereabove on describing the ~~effects of internal~~  
1184 ~~tidal waves on temperature variation~~ impacts of tides on a seasonal scale, ~~while. A companion~~

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1185 ~~paper will then analyze the remainder variability of this work will address~~ temperature changes  
1186 ~~on finer time at tidal and subtidal scales, notably on the tidal scale using our model simulations~~  
1187 ~~and two observational data.~~

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## 1190 Data availability

1191 The TMI SST v7.1 data are publicly available online from the REMSS platform:  
1192 <https://www.remss.com/missions/tmi/>, was accessed on 27 June 2022. The model simulations  
1193 are available upon request by contacting ~~the~~ corresponding ~~authors~~author.

## 1194 Authors contributions

1195 Funding acquisition, AKL; Conceptualization and methodology, FA, AKL and ID.  
1196 Numerical simulations, GM and FA. Formal analysis, FA; FA prepared the paper with  
1197 contribution from all co-authors.

## 1198 Competing interests

1199 The authors declare that they have no conflict of interest.

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1204 ~~under the cobadging of Ariane Koch-Larrouy and Isabelle Dadou.~~ The numerical ~~simulation~~  
1205 ~~was~~simulations were founded by CNRS/CNES/IRD via the projects A0080111357 and  
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1207 platform ~~calculator (Jean-Zay)~~for modelling and computing.

## 1208 AcknowledgmentsAbbreviations

1209 The following abbreviations are used in this manuscript:

1210 ~~AMASSEDS: A Multi-disciplinary Amazon Shelf SEDiment Study~~

1211 ~~AMAZOMIX: AMAZOn MIXing~~

1212 ~~FES2012 | FES2014: Finite Element Solution 2012 | Finite Element Solution 2014~~

1213 ~~NEMO/PISCES: Nucleus for European MOdeling / Pelagic Interactions Scheme for Carbon~~  
1214 ~~and Ecosystem Studies~~

1215 ~~PIRATA: Prediction and Research moored Array in the Tropical Atlantic~~

1216 ~~REVIZEE - Recursos Vivos da Zona Econômica Exclusiva~~

1217  
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1221 ~~making the WOA2018 products available.~~

## 1222 ~~▲~~

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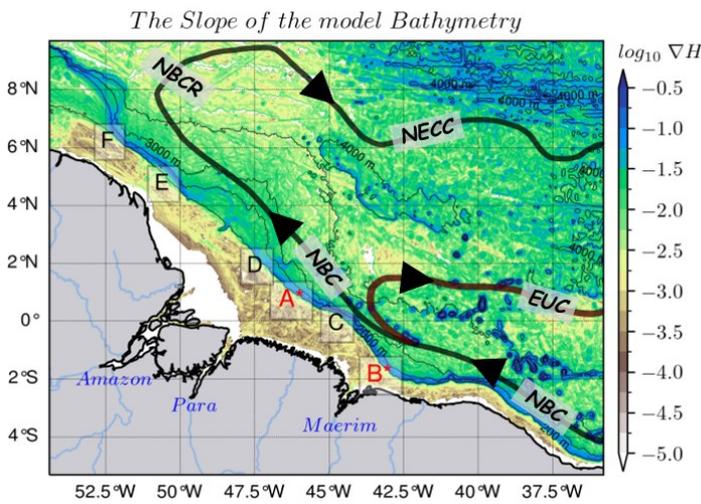
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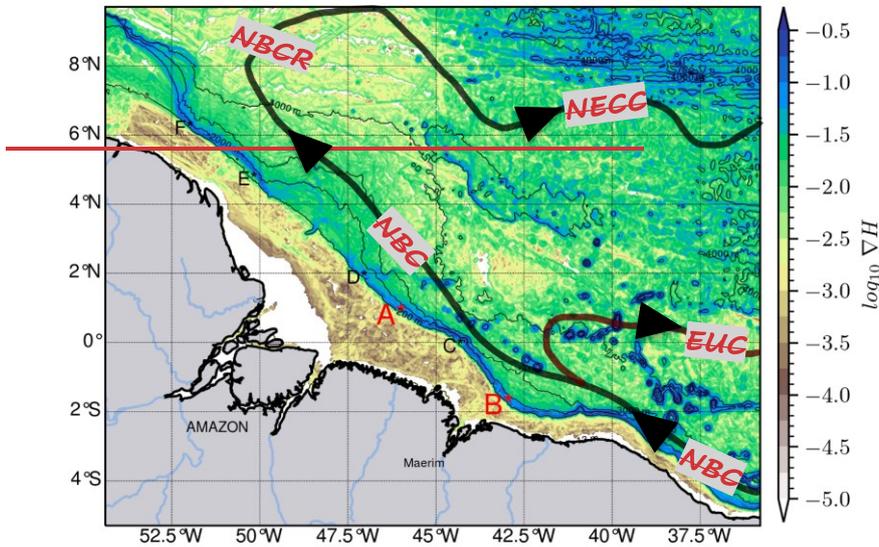
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The Slope of the model Bathymetry



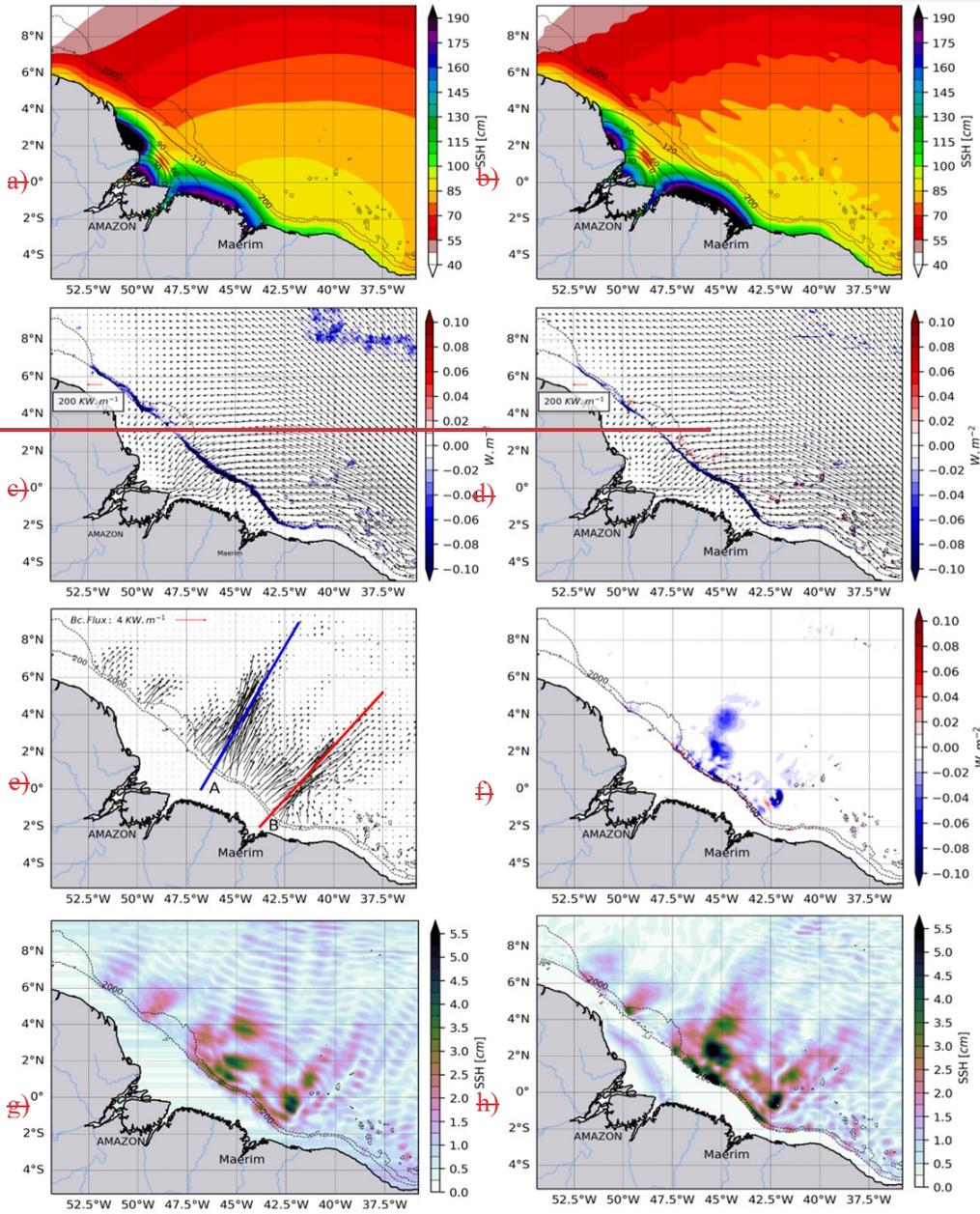
1763  
 1764 Figure 1: The horizontal gradient ( $\nabla H$ ) of the model's bathymetry with different internal  
 1765 tides generation sites ( $A^*$ ,  $B^*$ ,  $C^*$ ,  $D^*$ ,  $E^*$  and  $F^*$ ) along the high slope (blue color shading)  
 1766 of the shelf break, with the two main sites (in red) being  $A^*$  and  $B^*$ , as  
 1767 mentioned/reported in Magalhaes et al. (2016) and Tchilibou et al. (2022). Solid bold lines  
 1768 represent a schematic view of the circulation (as described by Didden and Schott, 1993;  
 1769 Richardson et al., 1994; Bourlès et al., 1999a; Johns et al., 1998; Bourlès et al., 1999a; Schott  
 1770 et al., 2003; Garzoli et al., 2004) with NBC, NBCR and NECC pathways/tracks in black, and  
 1771 the EUC pathway/track in brown red. Tin black contours are 200 m, 2000 m, 3000 m and 4000  
 1772 m isobaths.

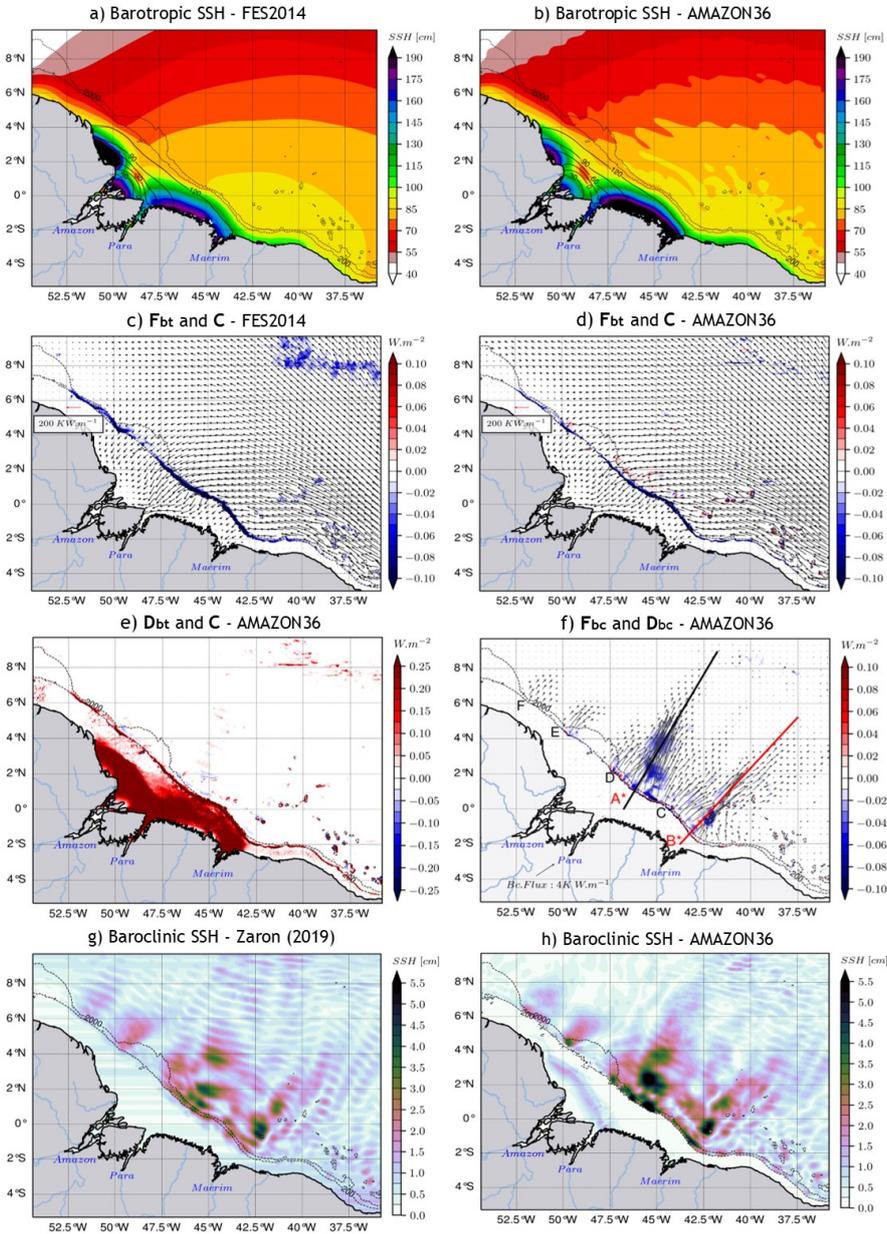
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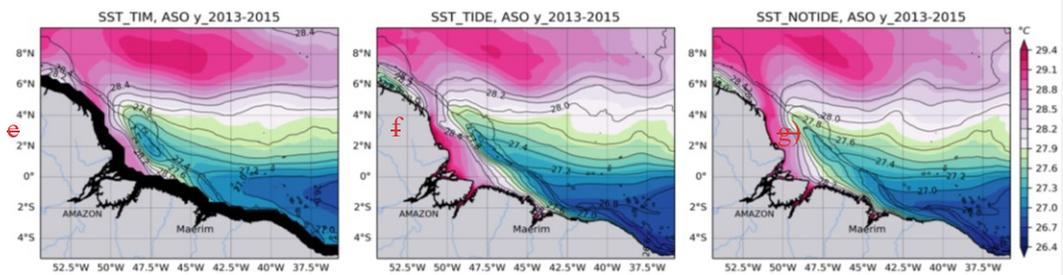
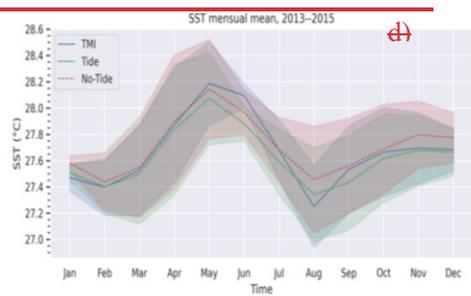
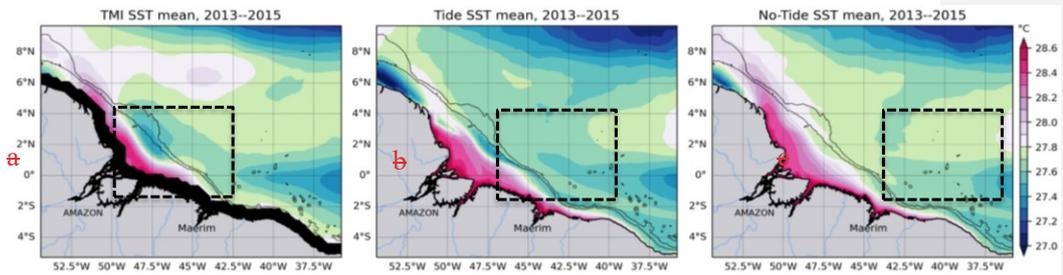


bt : barotropic ; bc : baroclinic  
 F : energy flux ; C : barotropic-to-baroclinic energy conversion ; D : energy dissipation

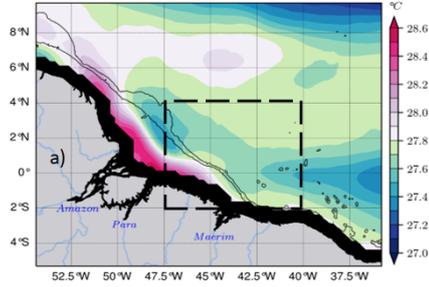
1781 ~~Figure 2: M2 tides (coherent~~Coherent (or stationary) characteristics of the M2 tides.  
 1782 Barotropic sea surface height ~~for~~(color shading) and its phase (solid tin contours) for (a)  
 1783 FES2014 ~~(a)~~ and (b) the model ~~(b)~~, depth-integrated, barotropic energy flux (black arrows)  
 1784 with the energy conversion rate (CVR) ~~(color shading)~~ for (c) FES2014 ~~(c)~~ and (d) the model  
 1785 ~~(d)~~ with, (e) the model depth-integrated barotropic energy flux black arrows, dissipation, (f) the  
 1786 model depth-integrated baroclinic energy flux (black ~~(e)~~ arrows with transect lines along A  
 1787 (blue) and B (red) IT's pathways, model) and the depth-integrated baroclinic energy  
 1788 dissipation (f, and color shading) with transect lines along IT's trajectories A\* (black) and B\*  
 1789 (red), the baroclinic sea surface height from observation ~~(g)~~ Zaron, (2019), ~~(g)~~ and (h) the  
 1790 model ~~(h)~~. Data from the model are the mean value over the year 2015. For all panels, dashed  
 1791 black lines represent the 200 m and 2000 m isobaths of the model bathymetry.

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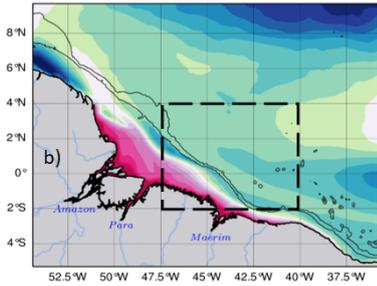
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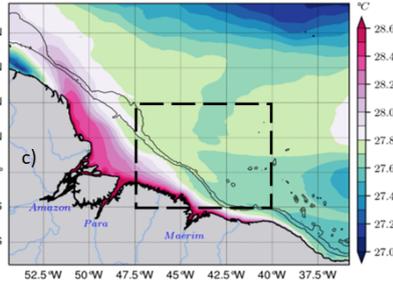
TMI SST mean, 2013–2015



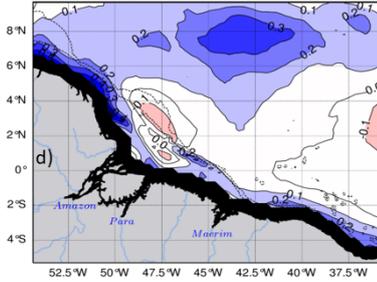
Tide SST mean, 2013–2015



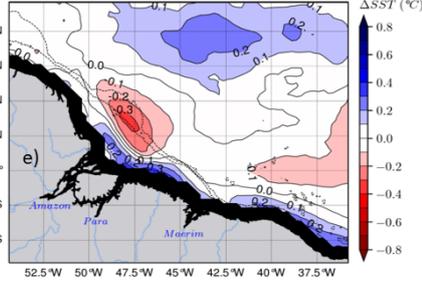
No-Tide SST mean, 2013–2015



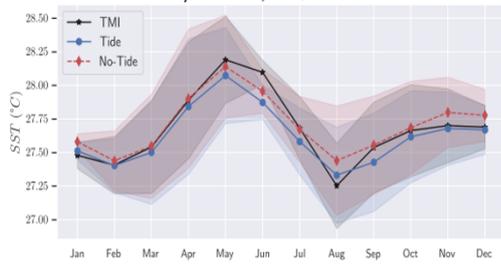
$\Delta(TMI, Tide)$ , 2013–2015



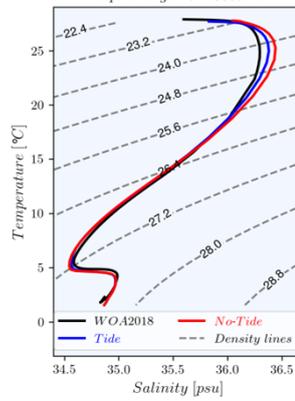
$\Delta(TMI, No-Tide)$ , 2013–2015



f) SST monthly mean, 2013–2015



g) T-S diagram  
depth range : 0 – 5500 m



1810 Figure 3: Mean SST Validation of the model temperature for the whole period 2013-2015  
1811 from. Mean SST for (a) TMI SST (a) with ~~its~~ black coastal mask, (b) the model's tidal  
1812 simulation (b), (c) the model's non-tidal simulation (c), the difference (bias) in SST between  
1813 TMI and (d) the tidal simulation and (e) the non-tidal simulation, (f) the seasonal cycle of the  
1814 SST of the three products averaged ~~inside~~ within the ~~dotted~~ dashed line box in upper panels  
1815 (covering FITs pathways emanating from the main generation sites A and B) with ~~shelf~~ values  
1816 masked ~~over~~ below the 200 m isobath, ~~the~~ bands ~~give the~~ indicate variability according to  
1817 standard deviation (d), (g) Temperature-Salinity (T-S) diagram of the mean properties in the  
1818 same area as (e) from observed WOA2018 climatology (black line), the tidal simulation (blue  
1819 line) and non-tidal simulation (red line) for the water column from surface to 5500 m depth,  
1820 dashed gray lines represent density ( $\sigma_\theta$ ) contours. For panels a-e and hereinafter (unless  
1821 otherwise stated), the solid black lines represent the 200 m and 2000 m isobaths from the model  
1822 bathymetry.

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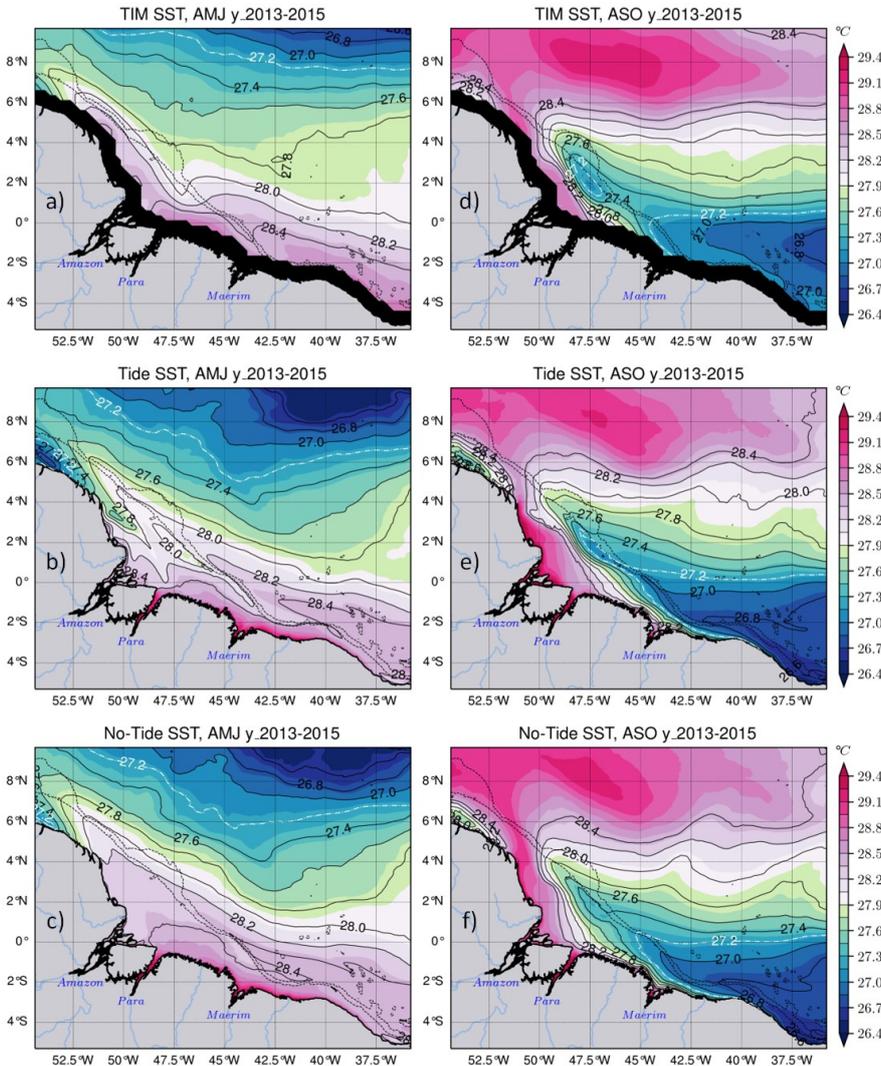
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1823  
 1824 *Figure 4: 2013-2015 seasonal SST mean. The lower-left panels present the SST averaged stand-*  
 1825 *for the ASO (August-September-October) AMJ season over the years 2013-2015 for TMI SST*  
 1826 *(e) with a white/black coastal mask, the model's tidal simulation (f) and the model's non-*  
 1827 *tidal simulation (e), with the, respectively for the upper-left, center-left and lower-left panel:*  
 1828 *the same in the panels on the right but for the ASO season. The dashed white and black*

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1829 *solid* lines ~~representing~~*represent* the temperature contours. ~~The~~*Dashed* black ~~tin~~ lines stand  
1830 for the 200 m and 2000 m isobaths-

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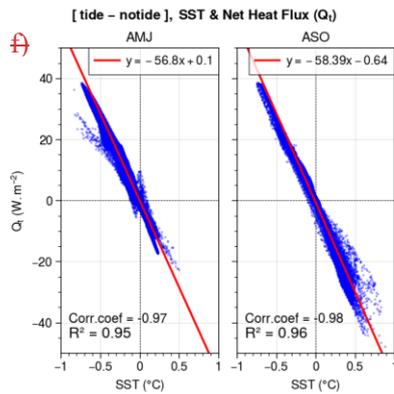
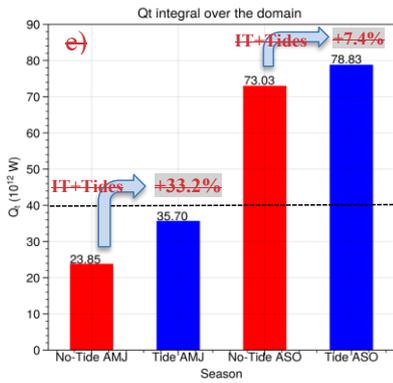
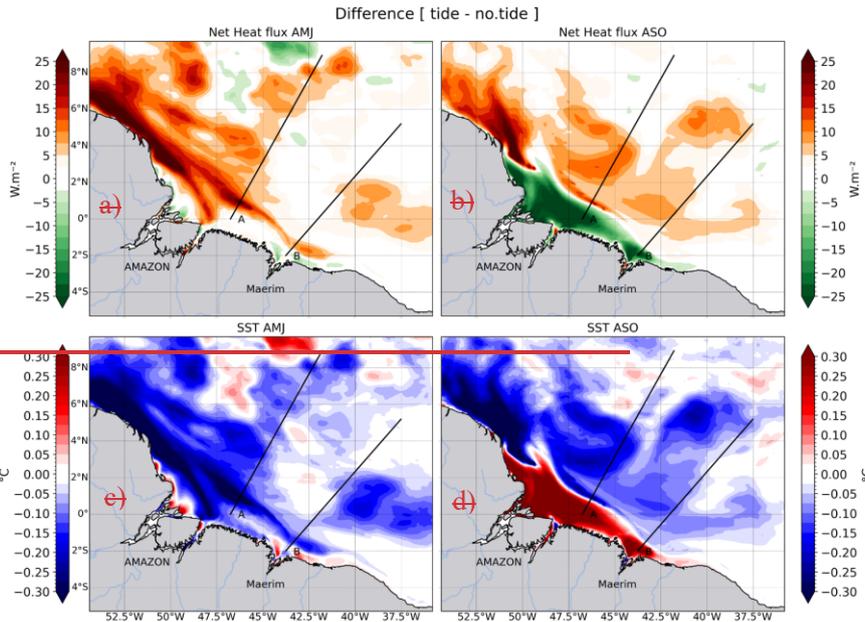
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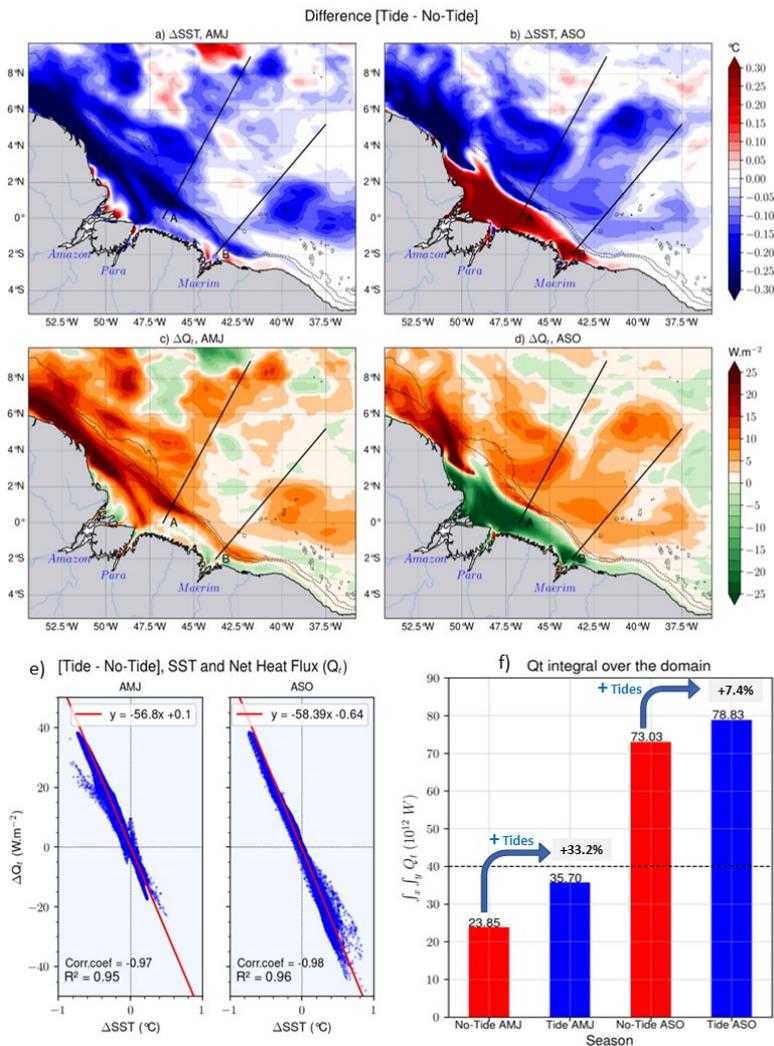
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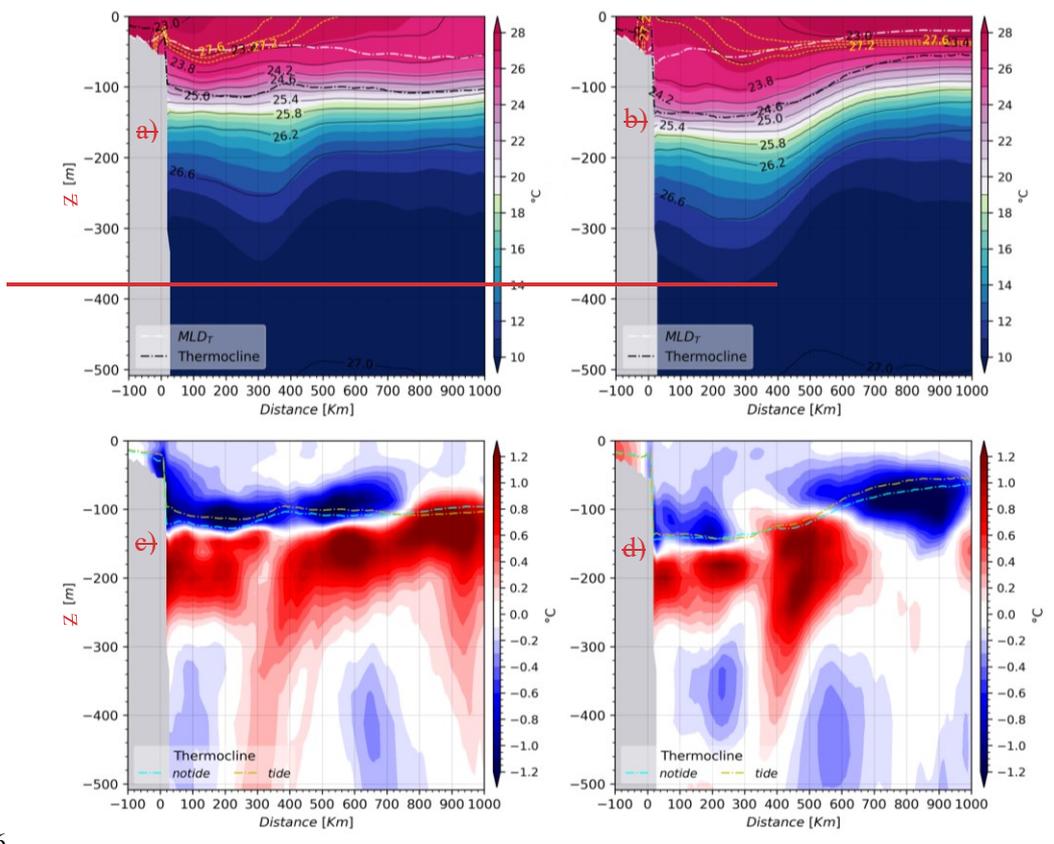
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1834 *from the model bathymetry.*

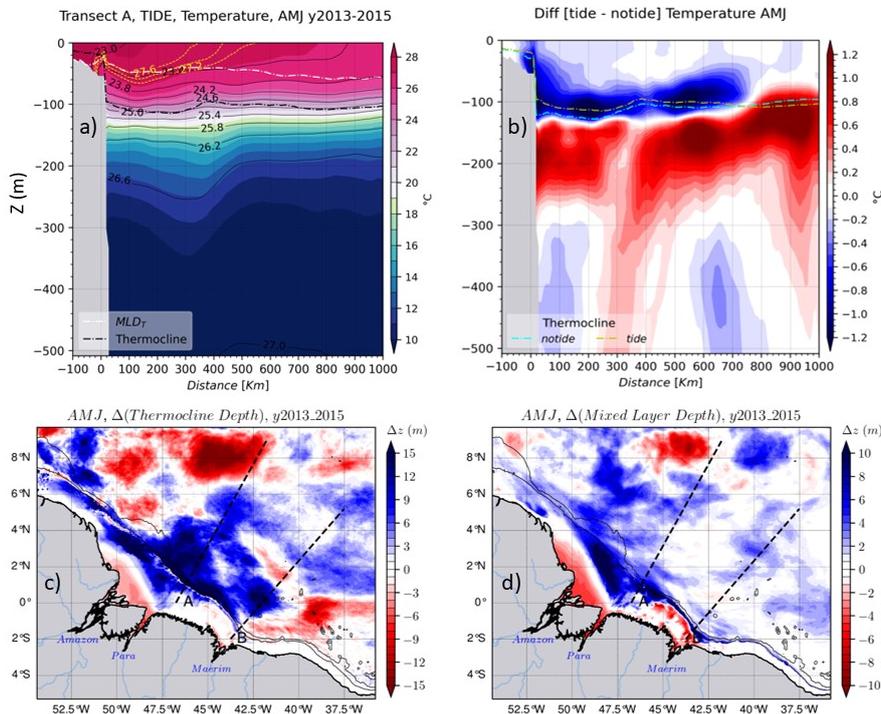


1835  
 1836 **Figure 45:** Relationship between the SST and the atmosphere-to-ocean net heat flux ( $Q_t$ ) and  
 1837 the): SST-the  $Q_t$  difference between tide and no-tide simulations-anomaly [Tide - No-Tide] in  
 1838 AMJ (a) and ASO (b) season, and SST differences seasons,  $Q_t$  anomaly in AMJ (c) and ASO (d)-  
 1839 seasons, (e) correlation between  $Q_t$  anomaly and SST anomaly for each season, (f) domain  
 1840 integrated  $Q_t$  (e) for both seasons for of each simulation. Correlation between  $Q_t$  difference  
 1841 and SST difference for each season (f). Hereinafter, -anomaly- refers to what is described  
 1842 hereabove.

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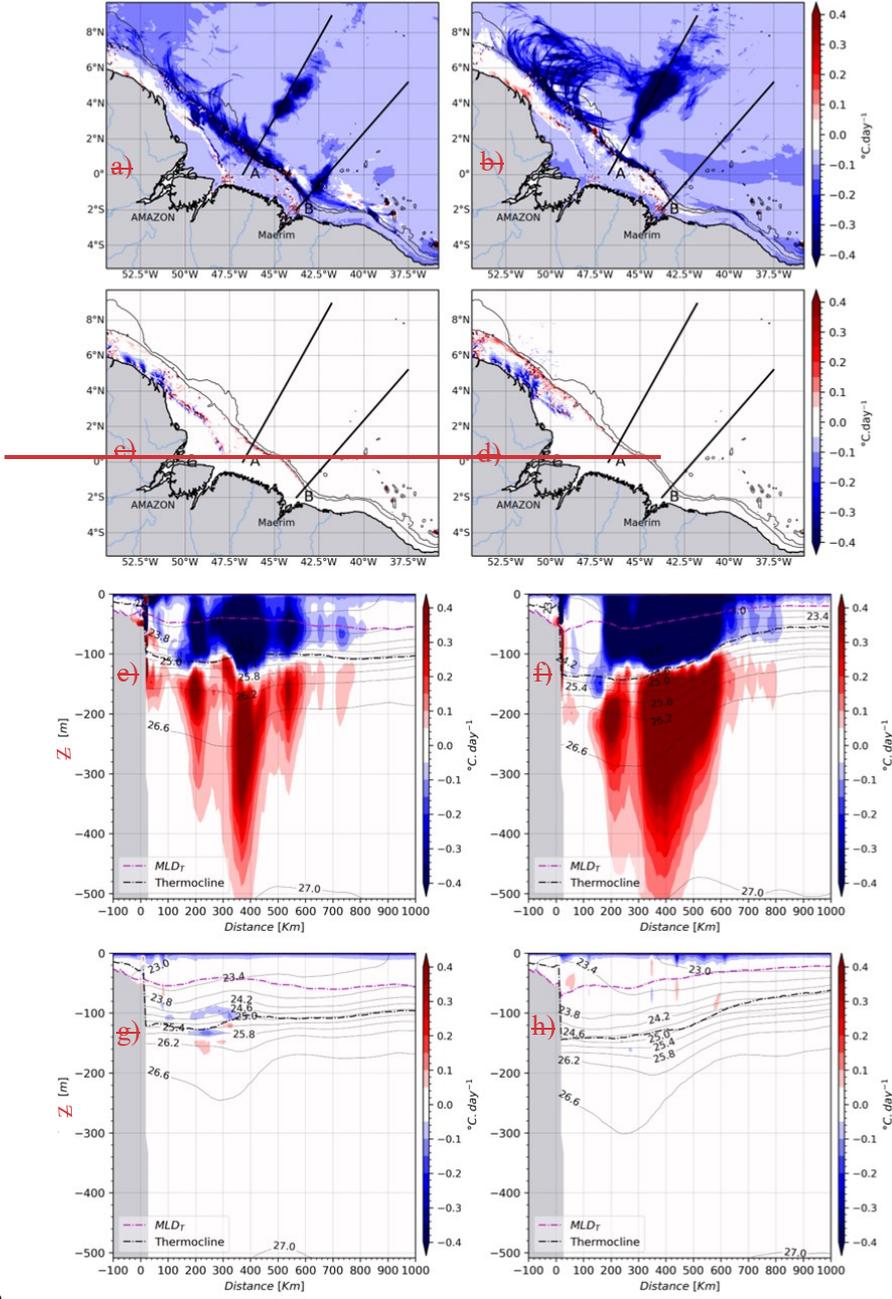


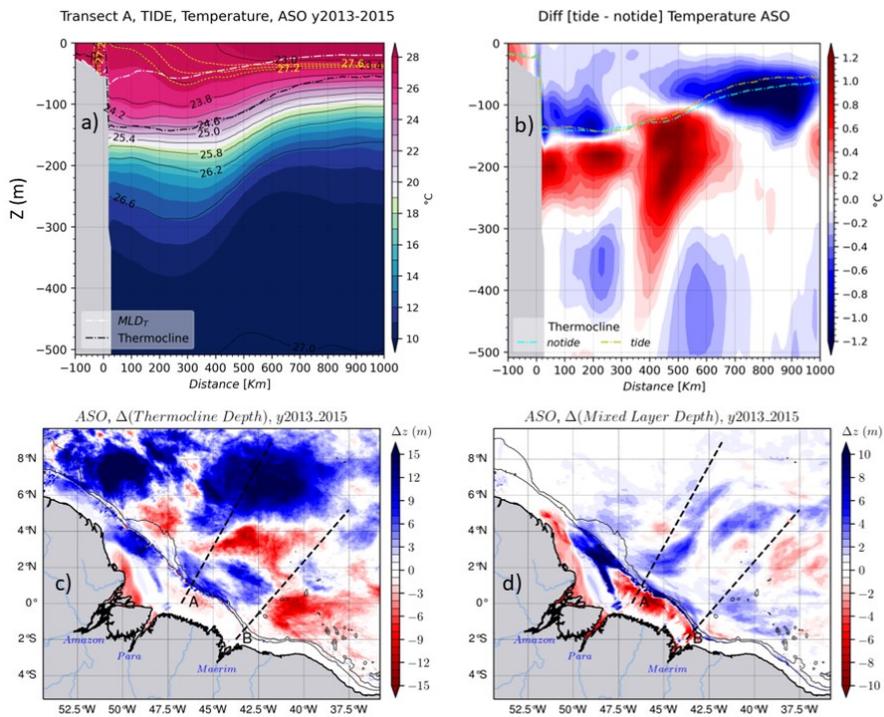
1847  
 1848 *Figure 5: Vertical sections of tidal simulation's temperature following IT's pathway A6: Some*  
 1849 *water mass properties for the AMJ season: (a) and ASO (b) seasons. Difference of vertical*  
 1850 *section of the temperature between tidal and non tidal simulations for AMJ (c) and ASO (d)*  
 1851 *seasons. The of the tidal simulation following the transect A, the yellow dotted-dashed and the*  
 1852 *solid black tin-lines in the upper panels are, respectively, for the temperature and density*  
 1853 *anomaly isoecontours ( $\sigma_\theta$ ) isolines respectively, the black and white ticker dot-dashed lines are*  
 1854 *respectively the thermocline and mixed layer depths. The MLD respectively, (b) the temperature*  
 1855 *anomaly for the same vertical section, yellow and cyan ticker dot-dashed lines in the lower*  
 1856 *panel are the thermocline depth respectively for tidal and non tidal simulations for the tidal*  
 1857 *and non-tidal simulations respectively, (c) thermocline depth anomaly and (d) MLD anomaly*  
 1858 *for the whole domain. When the MLD or the Thermocline depth anomaly are colored in blue*  
 1859 *(vs red) it means that the tides rise (vs deepen) them.*

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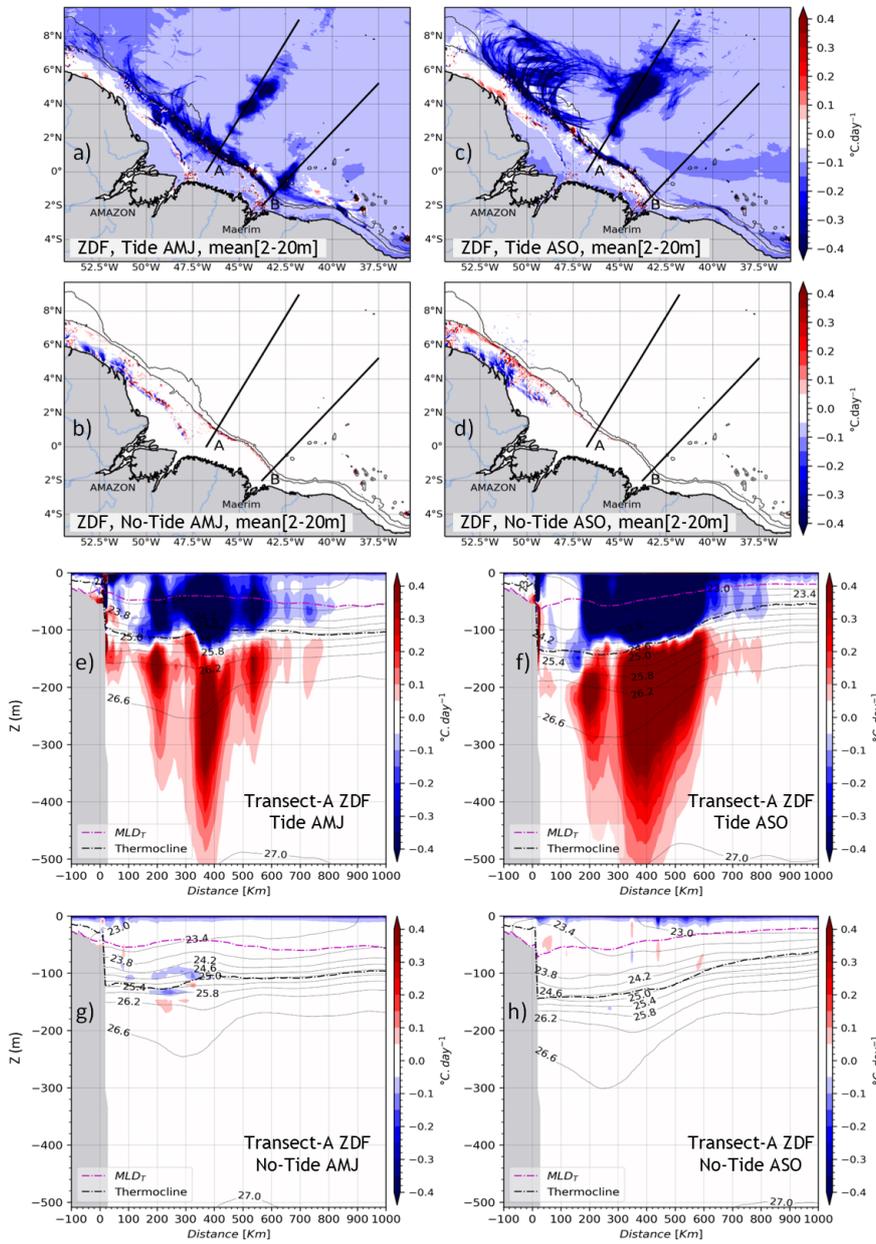
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1865 *Figure 7 : same as figure 6 but for the ASO season.*



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Figure 8: The vertical diffusion tendency of temperature (ZDF) for both seasons, respectively AMJ (left panel) and ASO (right panel). Vertical. The vertical mean between 2–20 m for AMJ

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1869 season in tidal (a) and non-tidal (eb) simulation; then for ASO season in tidal (bc) and non-  
1870 tidal (d) simulations. ~~Black thin contours are, from the coast to open ocean, 200 m and 2000~~  
1871 ~~m isobaths. Vertical section~~ Vertical sections of ZDF following the transect A for AMJ season  
1872 in ~~the tidal (e) and~~, for ASO season in non-tidal (gf) simulations; then for AMJ season in the  
1873 ~~non-tidal (g) and for ASO season in tidal (f) and nothe non-tidal (h) simulations~~ simulations.  
1874 The black and magenta ~~icker dot~~ dashed lines are ~~the thermocline depth and MLD~~ respectively  
1875 ~~thermocline depth and mixed layer depth. Thin. Solid black eontours are for~~ lines represent the  
1876 density anomaly ( $\sigma_\theta$ ) isocontours.

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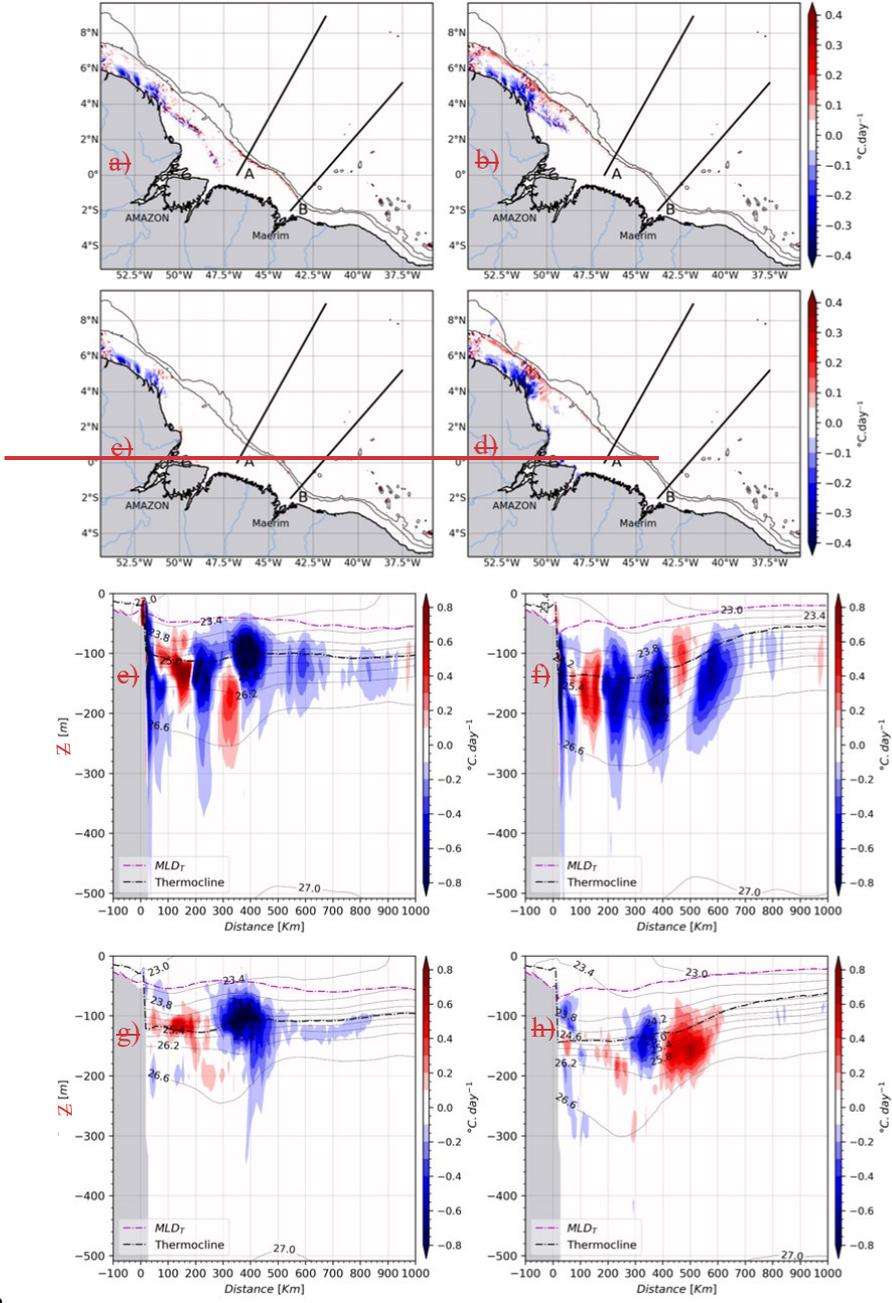
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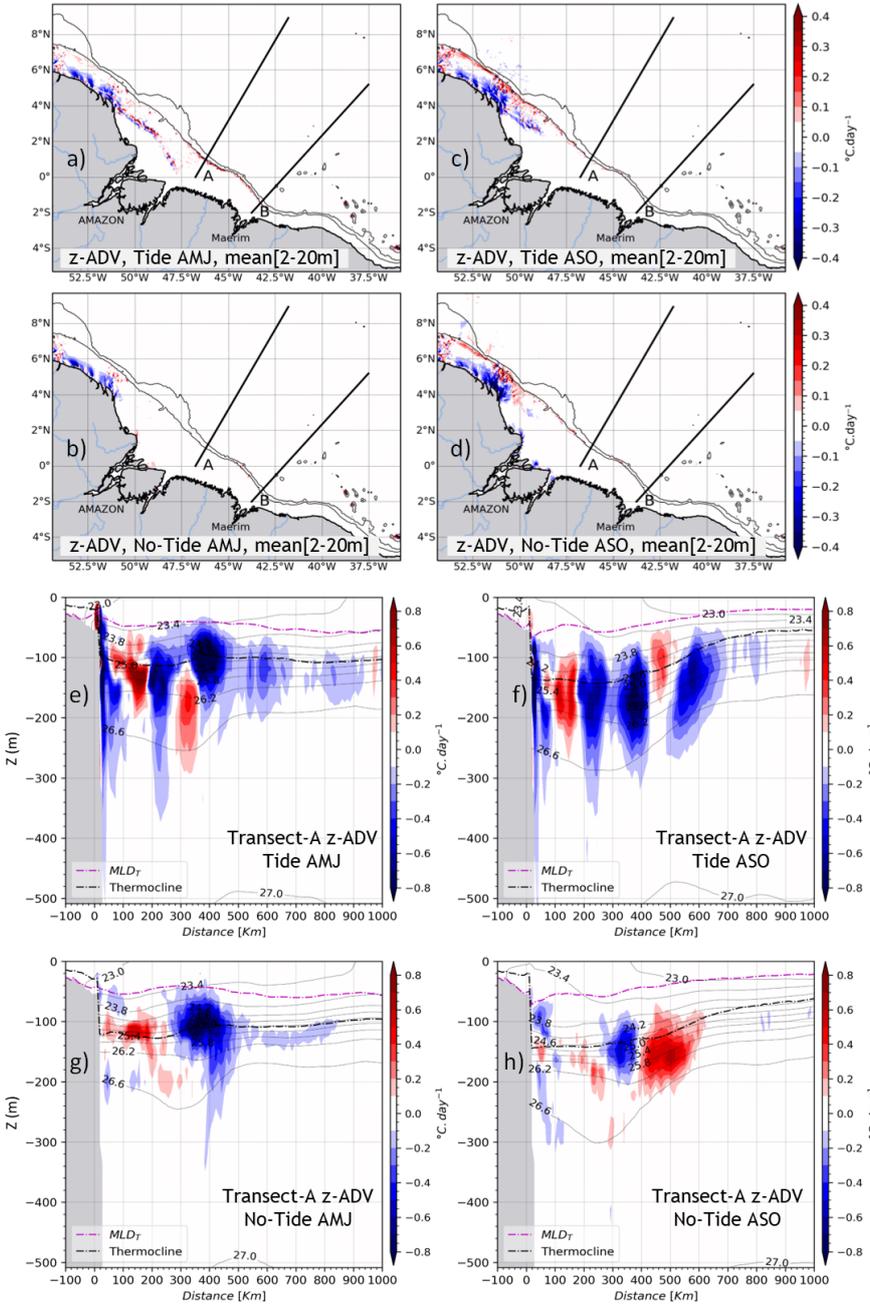
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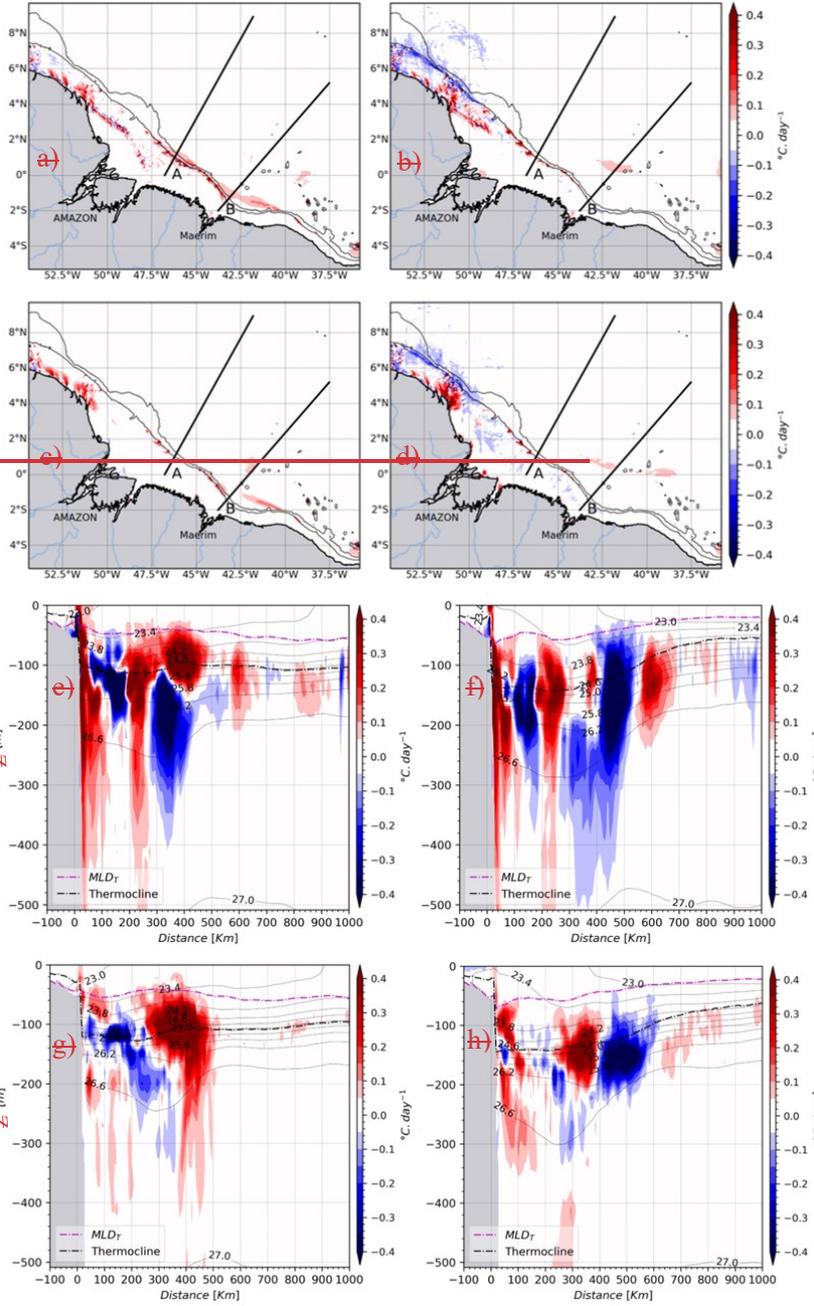


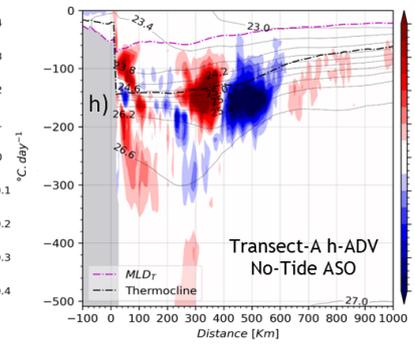
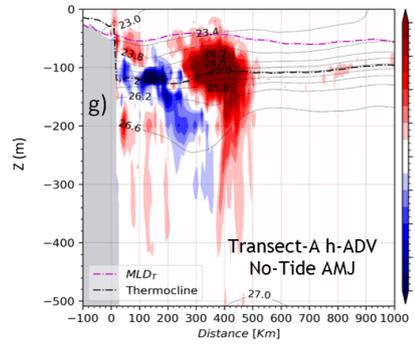
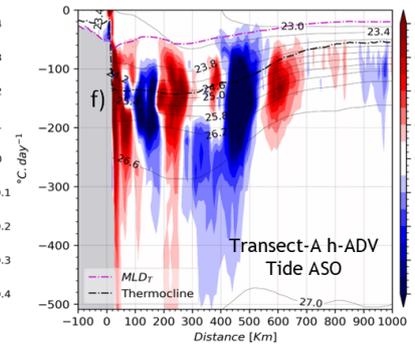
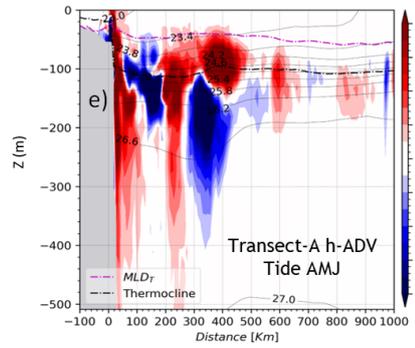
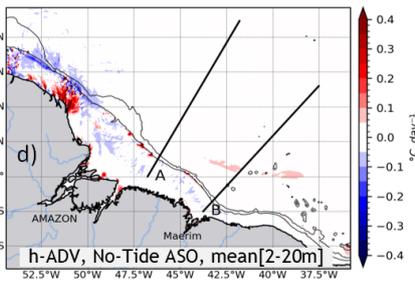
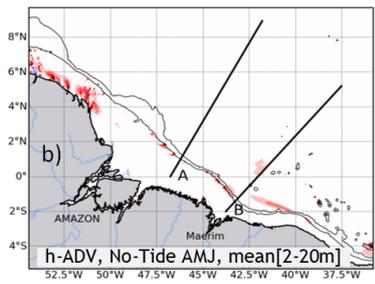
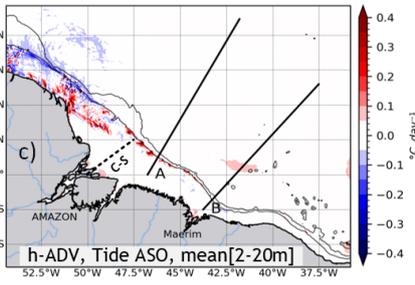
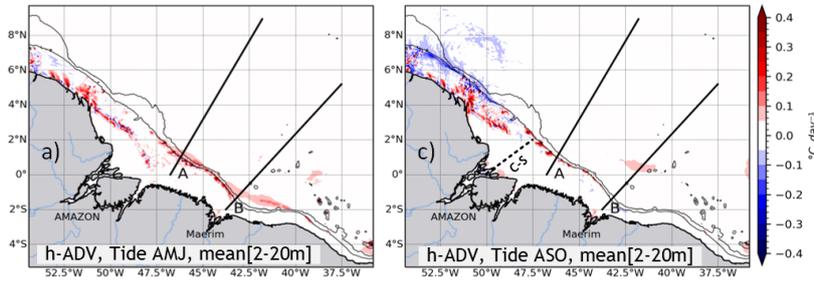


1894 Figure 79: same as figure 68, but for the vertical advection tendency of temperature ( $z$ -ADV).

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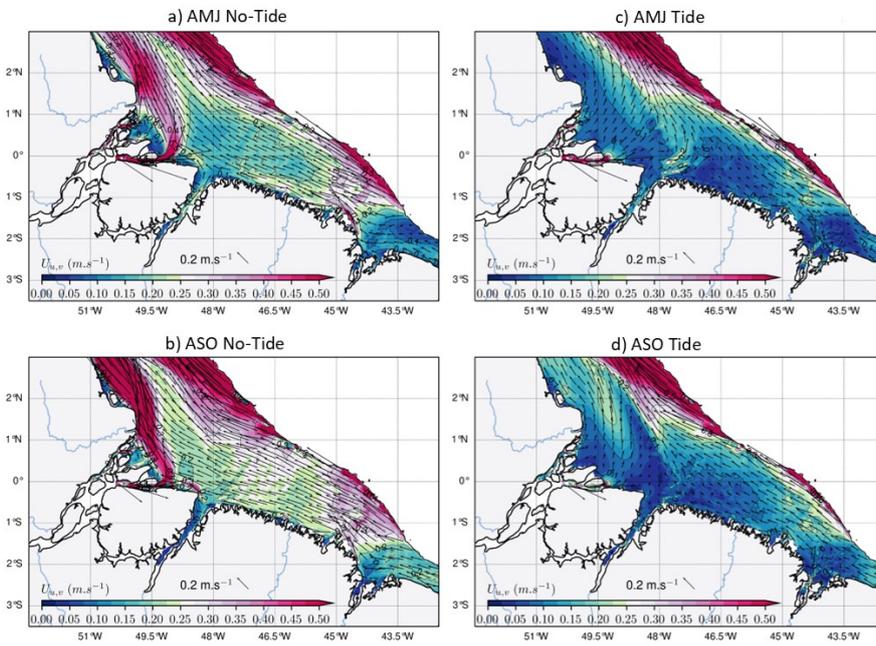
1897 *Figure 810: same as figure 68 but for the horizontal advection of temperature ( $h-ADV = x-$*   
1898  *$ADV + y-ADV$ ). The dashed line from the Amazon River mouth toward the outer shelf in the*  
1899 *panel (b) indicates the cross-shore transect (C-S) used further on.*

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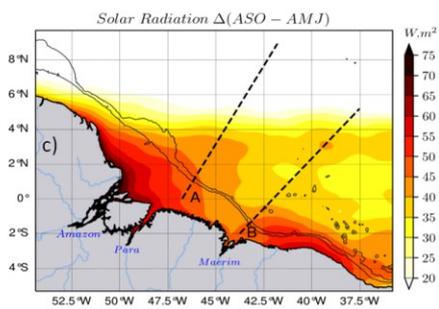
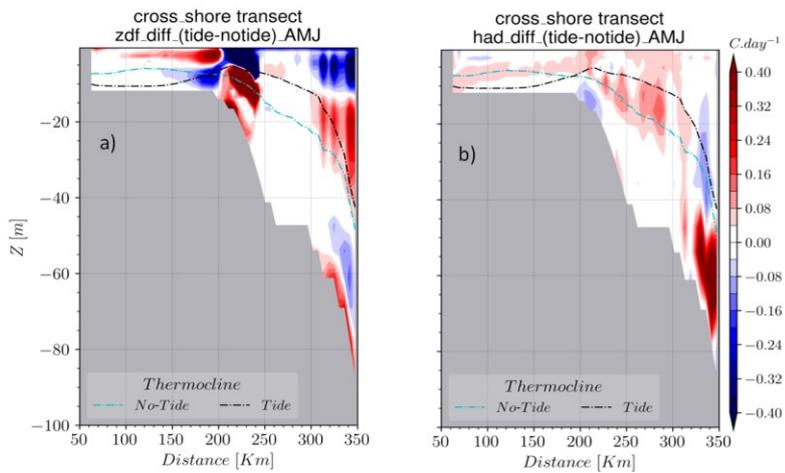
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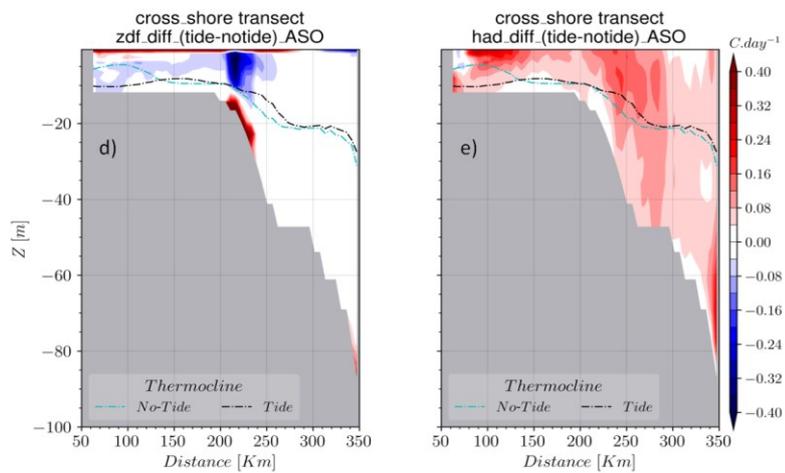


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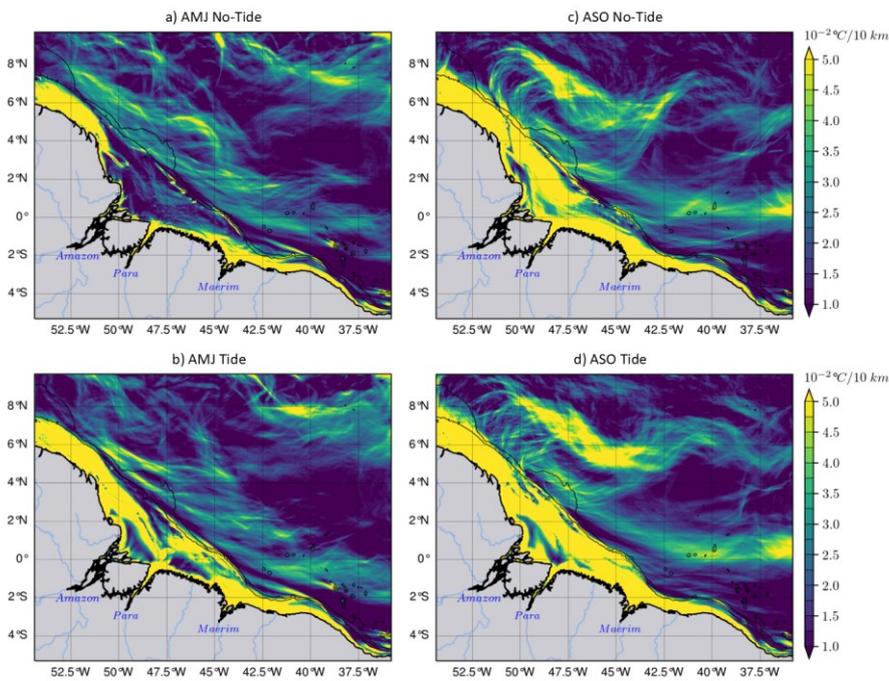
1932 *Figure 11: The seasonal mean of the current ( $U_{u,v}$ ) at the shelf averaged between the surface*  
1933 *and 50 m: the non-tidal simulation in the left panels and the tidal simulation in the right panels.*  
1934 *The upper panels stand for the AMJ season, while the lower stand for the ASO season. The*  
1935 *color shading is the modulus of the current and the black arrows represent its direction. Values*  
1936 *beyond the 200 m isobath are masked.*



Solar radiation ( $Q_s$ ) increases in the ASO season:  
 > 30  $W.m^2$  offshore  
 > 60  $W.m^2$  over the shelf



1938 *Figure 12: The cross-shore transect of ZDF anomaly for (a) AMJ and (b) ASO seasons, then*  
 1939 *for h-ADV anomaly for (d) AMJ and (e) ASO seasons ; (c) Difference in solar radiation*  
 1940 *between ASO and AMJ seasons. Solar radiation increases during the ASO season, with greater*  
 1941 *intensity on the shelf.*



1943  
 1944 *Figure 13 : The horizontal gradient of the Temperature ( $\nabla T$ ) averaged between 2–20 m : the*  
 1945 *AMJ season in the left panels and ASO season in the right panels, the simulations without the*  
 1946 *tides in the upper panels, and with tides in the lower panels. During the ASO season, the NBC*  
 1947 *retroreflects and eddy activity intensifies in the north-west. Therefore,  $\nabla T$  emphasizes eddy-like*  
 1948 *fronts at the same location as eddy-like patterns in ZDF (see Fig.8b).*

a mis en forme : Taquets de tabulation : 4,77 cm,Gauche

a mis en forme : Couleur de police : Texte 1