

Dear Thijs,

Thank you very much for handling our manuscript (egosphere-2023-347). We have substantially revised the manuscript by addressing all the reviewers' comments. The revisions in the manuscript and the reply to the comments are marked in blue. Thank you very much.

Sincerely,

Changwei Liu (liuchw8@mail.sysu.edu.cn)
On behalf of all the authors

COMMENTS FOR THE AUTHOR:

Anonymous Referee #1

Summary: In this study, the authors performed analyses on data collected during the MOSAiC campaign, focusing on the atmospheric boundary layer height (ABLH). The authors first identified the ABLH manually and then calibrated the critical bulk Richardson number in the bulk Richardson number method for computing ABLH based on the manually labeled ABLH. The relations between ABLH and surface variables were examined, and two cases were examined in detail to investigate the controlling factors of the ABLH variations during the campaign. My overall impression of the paper is that the motivation was justified, the methodology was sound, and the results made sense. I have a few comments on the bulk Richardson number method and also the language needs to be improved (beyond what I pointed out in my comments below).

Response: Thank you very much for your time and effort in reviewing our manuscript and for the constructive comments. We have revised the manuscript by addressing the referee comments, and our use of the English language has been carefully edited by all authors. The revisions in the manuscript and the reply to the comments are marked in blue.

Major comments:

- 1, the bulk Richardson number method for computing the ABLH.
 - 1.1 Some studies also considered a friction velocity in the definition of bulk Richardson number (see e.g., Zhang et al. 2020). It might be worth discussing this.

Response: Thank you very much for your helpful advice. In our improved Ri algorithm, the friction velocity (u_*) is now considered. The Ri formula with u_* can significantly improve the ABLH estimation in cases of ABLs with relatively high wind speed. In addition, we also consider the cloud effect in the improved Ri algorithm to estimate ABLH better. For this comment, the corresponding changes are given in our

revised manuscript as follows:

3.3 An improved Ri algorithm Considering the cloud effect

As a traditional Ri_b formula, Eq. (3) may break down in cases of ABLs with relatively high wind speed and upper-level stratification due to the overestimation of shear production (Kim and Mahrt, 1992). Vogelezang and Holtslag (1996) proposed the finite-difference Ri formula, which is expressed as:

$$Ri_F = \frac{(g/\theta_{vs})(\theta_{vh} - \theta_{vs})(h - z_s)}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (6)$$

where z_s is the lower boundary for the ABL, θ_{vs} , u_s , and v_s are the θ_v and wind components at the height z_s , respectively, b is an empirical coefficient, and u_* is the surface friction velocity. Ri_F is considered for a parcel located somewhat above the surface to avoid the above problem, and u_* is also taken into account to avoid underestimation in the situation of a uniform wind profile in the upper layer. Here, we use Ri_F for clear-sky profiles and take z_s and b values as 40 m and 100, respectively, according to Zhang et al. (2020).

As shown in Fig. 3, the estimations of cloudy ABLHs are sometimes quite poor, which motivates us to further improve the algorithm. Under cloudy conditions, the moist Richardson number (Ri_m) can be used to include cloud effects on the buoyancy term. Brooks et al. (2017) adopted the Ri_m formula expressed as:

$$Ri_m = \frac{\left(\frac{g}{T}\right)\left(\frac{dT}{dz} + \Gamma_m\right)\left(1 + \frac{Lq_s}{RT}\right) - \frac{g}{1 + q_w} \frac{dq_w}{dz}}{\frac{du^2}{dz} + \frac{dv^2}{dz}}, \quad (7)$$

where T is air temperature, Γ_m is the moist adiabatic lapse rate, L is the latent heat of vaporization, q_s is the saturation mixing ratio, and q_w is the total water mixing ratio, i.e., $q_w = q_s + q_L$, where q_L is the liquid water mixing ratio and is obtained based on the condensed water content. However, Eq. (6) is a gradient Ri and is calculated based on local gradients of wind speed, temperature, and humidity. To be consistent with the Ri formula proposed by Vogelezang and Holtslag (1996), we rewrite the formula in a finite-difference form expressed as:

$$Ri_m = \frac{\left[\left(g/T_s\right)\left(\frac{T_h - T_s}{h - z_s} + \Gamma_m\right)\left(1 + \frac{Lq_{sh}}{RT_h}\right) - \frac{g}{1 + q_{wh}} \frac{q_{wh} - q_{ws}}{h - z_s}\right](h - z_s)^2}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (8)$$

where subscripts (h and s) of the variables denote the calculated height, similar to Eq. (6), but note that the s and z_s are adjusted to 130 m, given the cloud radar blind zone. Considering that Ri_m is only appropriate for the liquid-bearing cloud cases, we use the Ri_F for “clear” grid points and use Ri_m for “cloudy” grid cells. Using this improved approach, we evaluated the best value of Ri_c to minimize the errors compared to the reference data set, arriving at an optimal value of $Ri_c=0.35$. The comparison of ABLH estimates obtained through the improved Ri algorithm with the manually-labeled ABLHs demonstrates significant improvement relative to other algorithms, particularly for cloudy conditions (Fig. 4, Table 1).

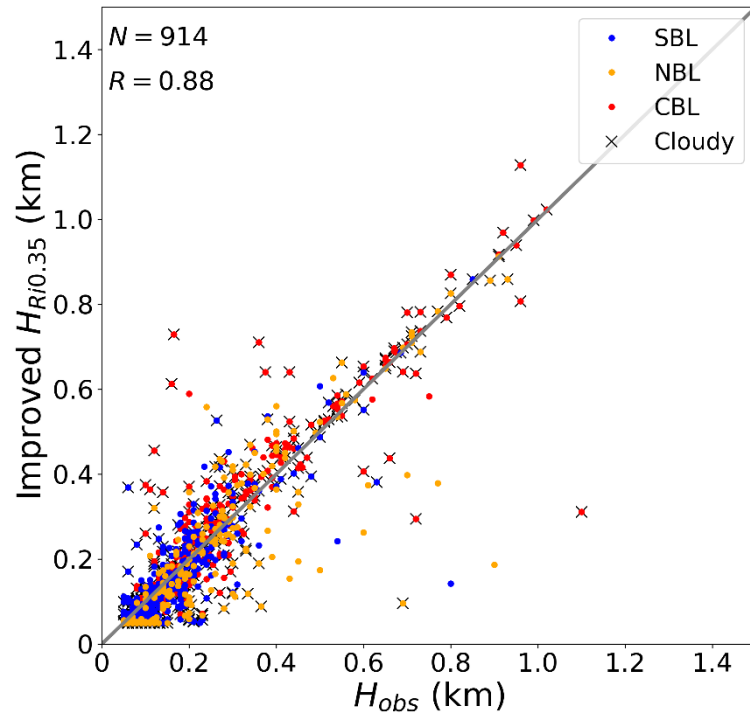


Figure 4 Similar to Fig. 3, but for the comparison of the ABLHs determined by the improved Ri algorithm with the observed ABLHs. The case number (N) and correlation coefficient (R) are given.

Table 1 The statistical measures (R , $Bias$, $MEAE$) for the four algorithms applied to the radiosonde dataset. All correlation coefficients are statistically significant ($p < 0.05$), except for SBL types in the Liu-Liang algorithm.

Algorithm	Regime type	R	Bias	MEAE (m)
The Ri_b algorithm with $Ri_{bc} = 0.25$	ALL	0.72	0.10	50
	SBL	0.81	0.16	34
	NBL	0.68	-0.04	62
	CBL	0.65	0.15	71
	Cloudy	0.69	0.08	51
The Ri_b algorithm with $Ri_{bc} = 0.5$	ALL	0.67	0.40	97
	SBL	0.73	0.50	88
	NBL	0.61	0.23	91
	CBL	0.60	0.39	120
	Cloudy	0.66	0.36	94
The Heffter algorithm	ALL	0.57	0.23	53
	SBL	0.46	0.17	33
	NBL	0.45	0.30	59
	CBL	0.66	0.28	74
	Cloudy	0.68	0.25	59
The Liu-Liang algorithm	ALL	0.47	0.04	82
	SBL	0.05	0.15	90
	NBL	0.44	-0.07	81
	CBL	0.56	-0.05	69
	Cloudy	0.52	-0.01	82
The improved Ri algorithm with $Ri_{bc} = 0.35$	ALL	0.85	-0.06	29
	SBL	0.79	-0.08	21
	NBL	0.79	-0.18	35
	CBL	0.87	0.05	36
	Cloudy	0.86	-0.03	30

1.2 it is not clear whether Eq. 2 is exactly the formula used in the VAP. If so, please state it.

Response: Thank you very much for pointing this out. The Eq. (2) is exactly the formula used in the PBLH VAP. The VAP technical report cites Sivaraman et al. (2013) as a reference for the algorithms. We have added relevant statement and the reference into our revised manuscript as follows:

These algorithms, including the Liu-Liang algorithm, the Heffter algorithm, and the bulk Richardson number algorithm, are all available in the PBLH VAP, as described in Sivaraman et al. (2013).

Ri_b is a dimensionless number that represents the ratio of thermally produced

turbulence to that induced by mechanical shear. The Ri_b formula used in the PBLH VAP (Sørensen et al., 1998; Sivaraman et al., 2013) is expressed as:

$$Ri_b = \left(\frac{gz}{\theta_{v0}} \right) \left(\frac{\theta_{vz} - \theta_{v0}}{u_z^2 + v_z^2} \right), \quad (3)$$

1.3 the authors mentioned that their results are different from Jozef et al. (2022). It would help the readers understand this by discussing a bit more of how exactly the formulations differ. Which formula did Jozef et al. use?

Response: Thank you very much for your helpful suggestions. We have added the analysis and comparison of different Ri formulas and Ri_c values, and the Ri formula used in Jozef et al. (2022) is now also included. Jozef et al. (2022) calculates the Ri over a rolling 30 m altitude range, and uses the Ri_c value of 0.75. The method of calculating Ri over a rolling 30 m range causes dramatic variation within the ABL, as seen in Fig. 5. Thus, for this Ri definition, a large Ri_c value is required to avoid the noise. The corresponding changes are given in our revised manuscript as follows:

Since some other studies have proposed different Ri_c values for MOSAiC (e.g., Jozef et al., 2022; Barten et al., 2023; Akansu et al., 2023), we will discuss the difference in Ri_c values here. The first thing to make clear is that these studies use different formulas to obtain Ri profiles. Barten et al. (2023) and Akansu et al. (2023) both use the traditional Ri_b algorithm based on Eq. (3), while they used Ri_c values of 0.4 and 0.12, respectively. This difference was likely caused by the different methods to manually derive their reference ABLH data sets. Jozef et al. (2022) calculates the Ri over a rolling 30 m altitude range, labeled as Ri_r , and the criterion is modified to require four consecutive data points to be above the Ri_c of 0.75. In our study, we use Ri_F proposed by Vogelezang and Holtslag (1996) for clear-sky conditions, and Ri_m for cloudy conditions. Based on the results presented here, it is apparent that this more complex approach improves the error statistics relative to approaches based on Eq. (3), regardless of Ri_c . In addition, some of the differences may also related to authors using different data sets or time periods. For instance, Akansu et al. (2023) primarily used sounding data based on tether balloon for a specific sub-period of MOSAiC, and Jozef et al. (2022) used radiosondes from when they had concurrent UAV observations. The data used in our study are based on merged sounding-tower product, as mentioned above.

To further explore the differences among the four different Ri approaches, we examine one SBL and CBL case. For a clear-sky SBL case (Fig. 5 a, b), the approaches from Akansu et al., Jozef et al. (2022), and this study all agree closely with the manual ABLH, while the Barten et al. approach results in a significant overestimation. For a cloudy-sky CBL case (Fig. 5 c d), the approach from this study agrees with the manual ABLH, while the approach from Barten et al. overestimates the ABLH by about 30 m, and the approaches from Akansu et al. and Jozef et al. (2022) underestimate the ABLH by 130 m and 230 m, respectively. These results further demonstrate how Ri_c depends

on the choice of Ri formula. Moreover, Ri_c is not analytically derived from basic physical principles (Zilitinkevich et al. 2007), and the concept of Ri_c is challenged by non-steady regimes (Zilitinkevich and Baklanov, 2002) and the hysteresis phenomenon (Banta et al., 2003; Tjernström et al., 2009). Therefore, an objective Ri_c does not exist. Rather, it is empirically used as an algorithmic parameter to simply derive the ABLH.

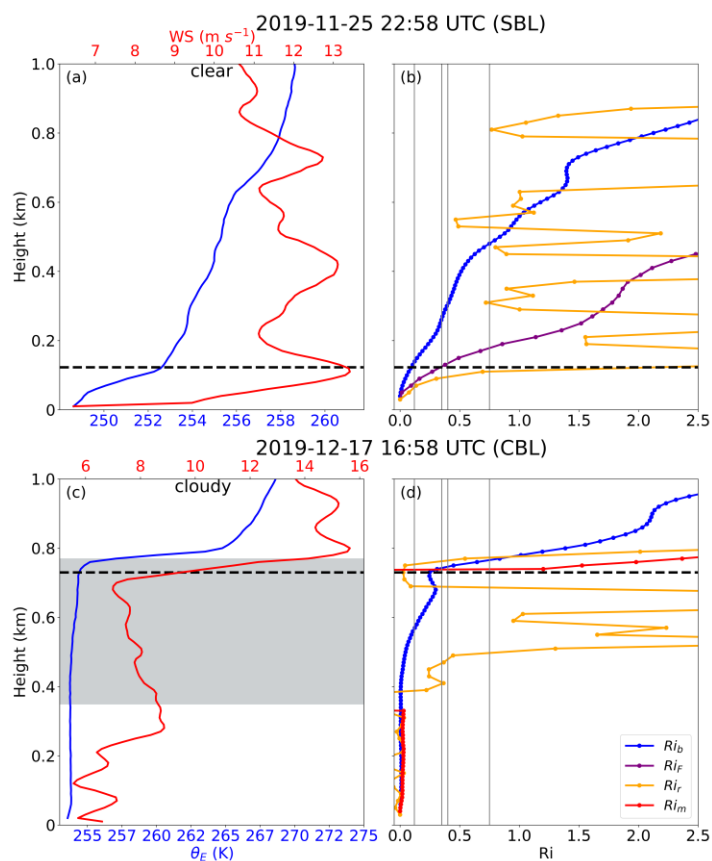


Figure 5 Vertical profiles of (left) θ_E and wind speed, and (right) Ri based on different formulas at (a–b) 25 November 2019, 22:58 UTC and (c–d) 17 December 2019, 16:58 UTC. Boundary layers at the two times represent a clear-sky SBL and a cloudy-sky CBL respectively. The black dashed horizontal lines denote the manually-identified ABLH, and the gray solid vertical lines denote the different Ri_c values, including 0.12, 0.35, 0.4, and 0.75. The gray shading in (c) denotes the cloud layer.

2, an automated algorithm

By looking at Figure 3, why not use an automated algorithm that is based on the bulk Richardson number method for SBL and the Heffter algorithm for CBL?

Response: Thank you very much for your helpful comment. The Heffter algorithm performs well in some CBL cases, but it also causes severe ABLH overestimation. As mentioned above, we have improved the Ri algorithm to better estimate ABLH in CBL cases. In addition, we realize that Fig. 3 is unclear. The data range of Fig. 3c in the

original manuscript is from 0–3.5 km due to severe ABLH overestimation by the Heffter algorithm, and thus the axis range is different from that of other panels. Therefore, we unify the axis ranges of all panels to avoid misunderstanding, and add a subgraph in Fig. 3c to denote all data points. The error statistics for different ABL regime types are also calculated and listed in Table 1. For this comment, the corresponding changes are given in our revised manuscript as follows:

(1) Changes related to Figure 3

Figure 3 presents the comparisons of estimated ABLHs with the manually-labeled ABLHs, and the associated statistical measures are given in Table 1. The results show that the Ri_b algorithm with Ri_{bc} of 0.25 performs best overall, and particularly for SBL cases. The performance of the Ri_b algorithm with Ri_{bc} of 0.5 is poorer than that of the Ri_b algorithm with Ri_{bc} of 0.25, with overestimations of ABLHs in general, and larger errors with lower correlation coefficients for all types of ABLs. The Heffter algorithm performs well in cases of high ABLH and particularly for cloudy and CBL cases, but does significantly overestimate ABLH in a large number of cases as shown in the Fig. 3c subgraph. This is attributed to the determination criterion of the Heffter algorithm, i.e., ABLHs are determined by inversion layers, which means that large errors occur when the inversion layer is higher than the mixed layer. Additionally, while the Heffter performance in many of the ABL conditions is only marginally worse statistically than the Ri_b algorithm with Ri_{bc} of 0.25, its correlations are notably worse for SBL and NBL cases. The performance of the Liu-Liang algorithm is generally poorer than the other algorithms, particularly for correlation coefficient, which is probably due to the impact of noise in the lower ABLH profiles and unsuitable parameters in the algorithm. In summary, the Ri_b algorithm is reliable over the Arctic Ocean and performs better than other algorithms, and this result agrees with Jozef et al. (2022). Furthermore, we will explore ways to improve the Ri_b algorithm to make it more suitable for cloudy and convective conditions.

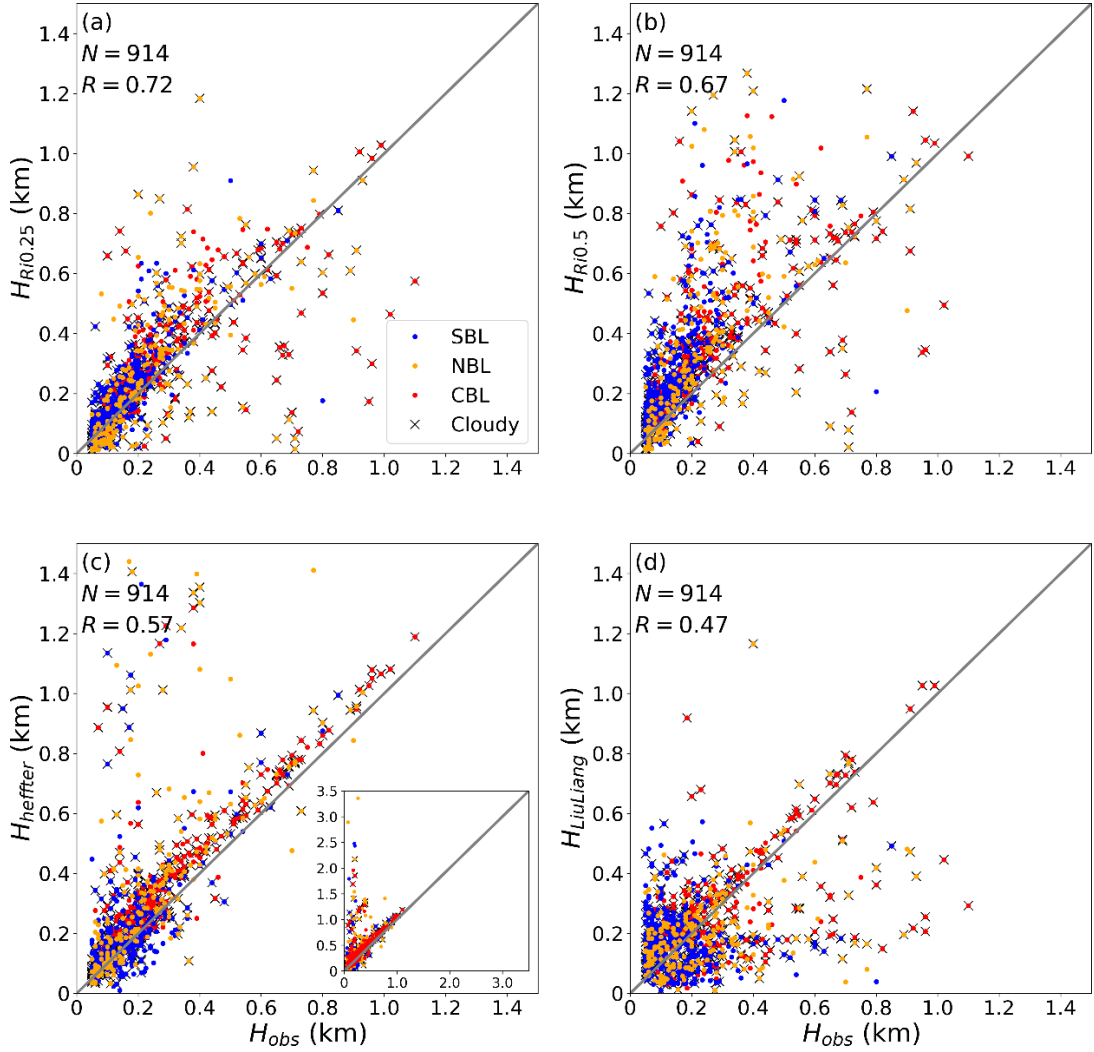


Figure 3 Comparisons of the ABLHs determined from radiosonde profiles using the bulk Richardson number (Ri_b) algorithm with the critical values (Ri_{bc}) of (a) 0.25 and (b) 0.5, (c) the Heffter algorithm, and (d) the Liu-Liang algorithm with the manually-identified “observed” ABLHs. The blue, yellow, and red colors indicate regime types of SBL, NBL, and CBL, respectively. The “x” signs indicate the Cloudy ABLs. The case numbers (N) and correlation coefficients (R) are given in each panel. The subgraph in (c) denotes all data points ranging from 0 to 3.5 km.

(2) Changes related to the improved Ri algorithm

3.3 An improved Ri algorithm Considering the cloud effect

As a traditional Ri_b formula, Eq. (3) may break down in cases of ABLs with relatively high wind speed and upper-level stratification due to the overestimation of shear production (Kim and Mahrt, 1992). Vogelezang and Holtslag (1996) proposed the finite-difference Ri formula, which is expressed as:

$$Ri_F = \frac{(g/\theta_{vs})(\theta_{vh} - \theta_{vs})(h - z_s)}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (6)$$

where z_s is the lower boundary for the ABL, θ_{vs} , u_s , and v_s are the θ_v and wind components at the height z_s , respectively, b is an empirical coefficient, and u_* is the

surface friction velocity. Ri_F is considered for a parcel located somewhat above the surface to avoid the above problem, and u_* is also taken into account to avoid underestimation in the situation of a uniform wind profile in the upper layer. Here, we use Ri_F for clear-sky profiles and take z_s and b values as 40 m and 100, respectively, according to Zhang et al. (2020).

As shown in Fig. 3, the estimations of cloudy ABLHs are sometimes quite poor, which motivates us to further improve the algorithm. Under cloudy conditions, the moist Richardson number (Ri_m) can be used to include cloud effects on the buoyancy term. Brooks et al. (2017) adopted the Ri_m formula expressed as:

$$Ri_m = \frac{\left(\frac{g}{T}\right) \left(\frac{dT}{dz} + \Gamma_m\right) \left(1 + \frac{Lq_s}{RT}\right) - \frac{g}{1 + q_w} \frac{dq_w}{dz}}{\frac{du^2}{dz} + \frac{dv^2}{dz}}, \quad (7)$$

where T is air temperature, Γ_m is the moist adiabatic lapse rate, L is the latent heat of vaporization, q_s is the saturation mixing ratio, and q_w is the total water mixing ratio, i.e., $q_w = q_s + q_L$, where q_L is the liquid water mixing ratio and is obtained based on the condensed water content. However, Eq. (6) is a gradient Ri and is calculated based on local gradients of wind speed, temperature, and humidity. To be consistent with the Ri formula proposed by Vogelezang and Holtslag (1996), we rewrite the formula in a finite-difference form expressed as:

$$Ri_m = \frac{\left[(g/T_s) \left(\frac{T_h - T_s}{h - z_s} + \Gamma_m\right) \left(1 + \frac{Lq_{sh}}{RT_h}\right) - \frac{g}{1 + q_{wh}} \frac{q_{wh} - q_{ws}}{h - z_s}\right] (h - z_s)^2}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (8)$$

where subscripts (h and s) of the variables denote the calculated height, similar to Eq. (6), but note that the s and z_s are adjusted to 130 m, given the cloud radar blind zone. Considering that Ri_m is only appropriate for the liquid-bearing cloud cases, we use the Ri_F for “clear” grid points and use Ri_m for “cloudy” grid cells. Using this improved approach, we evaluated the best value of Ri_c to minimize the errors compared to the reference data set, arriving at an optimal value of $Ri_c=0.35$. The comparison of ABLH estimates obtained through the improved Ri algorithm with the manually-labeled ABLHs demonstrates significant improvement relative to other algorithms, particularly for cloudy conditions (Fig. 4, Table 1).

3, line 318:

If the manually labeled ABLH didn't include any data in transit, why did the authors think that the calibrated bulk Richardson number method using the manually labeled ABLH could be used to compute ABLH during transit?

Response: Thank you very much for your helpful comment. The data in transit contain soundings that are influenced by open ocean environment, particularly in early June. In our original manuscript, we have directly excluded the contaminated soundings from our analysis. Nonetheless, we still manually labeled the ABLHs in transit and added the relevant discussion about the in-transit period into our revised manuscript. First, we

tested the performance of the traditional and improved Ri algorithms during transit periods. The results are presented in Fig. R1. It can be found that the traditional Ri algorithm performs well for selected ABL cases, but with an overall overestimation. The improved Ri algorithm solves this problem and shows better agreement with manually labeled ABLH. In addition, we also examined the impact of adding the manually labeled ABLH in transit into our analysis. We find that the difference in results would not affect our conclusions. Therefore, we have added the manually-labeled ABLH data during transit periods into our analysis and updated the results. The corresponding changes are given in our revised manuscript as follows:

(1) Changes related to in-transit information

The full-time series of ABLH during the MOSAiC expedition is presented in Fig. 7 and forms the bases for the remaining analyses. According to near surface conditions and the sea ice state, the whole MOSAiC observation period is divided into “freeze up”, “winter”, “transition”, and “summer melt” periods (Shupe et al., 2022), roughly corresponding to the seasons of autumn, winter, spring, and summer, respectively. In Figure 7, the black solid lines indicate persistent low-level clouds that exist for more than 12 h; these occur most frequently in the late summer and autumn (the “freeze up” period), which agrees with Shupe et al. (2011). Note that the grey dots indicate that the ABL data were observed while the vessel was in transit, and the representativity of the ABLH data should be considered in this context. For the first such period, the vessel left the MOSAiC ice floe in mid-May and slowly progressed south through tightly consolidated sea ice, such that the data are generally representative of the sea ice pack in the region. Measurements from early June when the vessel was near or in open water close to Svalbard have been excluded entirely from the analysis. In the middle of June, as the vessel returned to the original MOSAiC ice floe, the sea ice was not as tightly consolidated and the vessel preferentially went through leads; the preferentially lower ice fraction along this transit could have impacted the thermal structure of the ABL. For the three weeks in early August, the vessel moved around in the Fram Strait area and then made its way north to another passive sea ice drifting position near the North Pole, again transiting through regions with lower sea ice fraction. Finally, at the very end of the expedition, the vessel took some time to exit the sea ice, stopping a few times to allow for work on the ice.

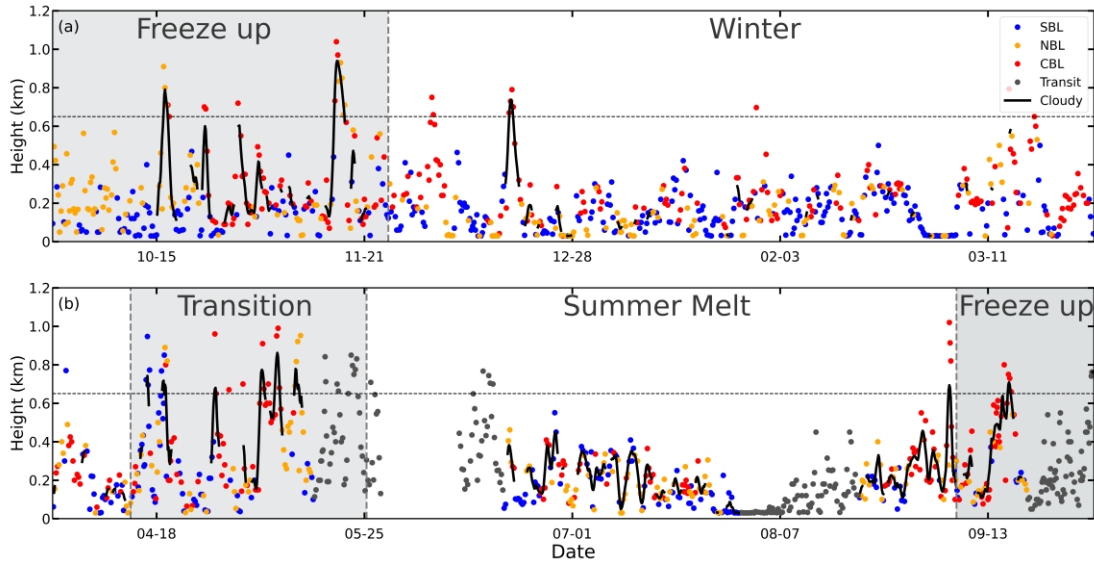


Figure 7 Time series of ABLHs throughout the MOSAiC year is divided into (a) and (b). The blue, yellow, and red dots indicate the heights of SBL, NBL, and CBL, respectively. The gray dots indicate ABL data observed while the vessel was in transit. The black solid lines indicate the heights of cloudy ABLs and persist for at least 12 hours. The gray dashed horizontal line denotes the 95th percentile of ABLH (650 m). The gray and white background shadings indicate the periods under different surface-melting states, i.e., “freeze up”, “winter”, “transition”, and “summer melt” periods.

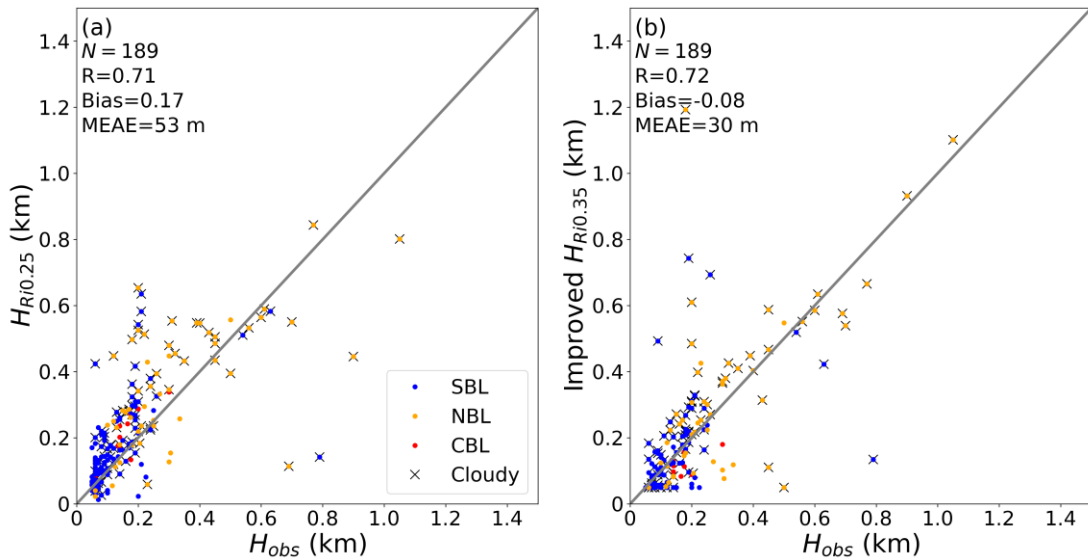


Figure R1 The Comparisons of the selected ABLHs determined by the (a) bulk Richardson number (Ri_b) algorithms with the critical values (Ri_{bc}) of 0.25 and (b) the improved Ri algorithm with observed ABLHs. The blue, yellow, and red dots indicate regime types of SBL, NBL, and CBL, respectively. The case number (N), correlation coefficient (R), $Bias$, and median of the absolute error ($MEAE$) are given in each panel.

(2) Changes related to data selection

The sentence “In total, we select 686 samples from 964 radiosonde profiles, and all data from observations while the vessel was in transit have been excluded” is fixed as “This VAP provides 964 ABLH estimates, and we select 914 samples from these to ensure that the estimates obtained by all algorithms are available”, and is added into the radiosonde description section.

Minor comments:

1, line 32: “has” should be “have”, and add “the” before “rapid changing”

Response: Revised as suggested.

2, line 37: “and the essential place for...” can be removed.

Response: Revised as suggested.

3, line 52: add “the” before “Atmospheric boundary layer height”, “referred to...” should be removed.

Response: Revised as suggested.

4, line 56: replace “literature” with “studies”

Response: Revised as suggested.

5, line 60: replace “surface mixed layer” with “surface layer”. Surface mixed layer is odd.

Response: Revised as suggested.

6, line 107: remove “special”

Response: Revised as suggested.

7, line 108: remove “fundamental”

Response: Revised as suggested.

8, line 211: I wouldn’t call this “multiple methods”. Maybe change it to “multiple profiles”.

Response: Revised as suggested.

9, line 161/182/221/260: I would not call this “subjective ABLHs”. Maybe “manually-labeled ABLHs”.

Response: Revised as suggested.

10, line 223: replace “applied” with “available”

Response: Revised as suggested.

11, line 257: add “a” before “better performance”.

Response: Revised as suggested.

12, line 302: “the smallest” should be “the best”. R is not the smallest clearly.

Response: Revised as suggested.

13, line 324-327: these sentences need to be re-worded.

Response: Thank you very much for pointing this out. This statement aims to demonstrate that overall variation of the Arctic ABLH during the MOSAiC year is irregular, which is distinct from a ABL variation over land surface with diurnal cycle. We have revised these sentences as “The Arctic ABL is suppressed for most of the MOSAiC year, while for a few periods it intensively develops for several days at a time, most commonly when clouds and a CBL are present”

14, line 382: I would probably not call this “where the annual cycle began”. Please reword.

Response: Thank you very much for pointing this out. We have revised the sentence as “The whole process forms these general shifts over the annual cycle.”

15, line 392: how do you know a priori that it is the surface conditions that influence the ABLH, not the other way around? Please re-word.

Response: Thank you very much for pointing this out. We have revised the sentence as “To further explore the relations between surface conditions and the ABLH, we evaluate ...”

16, line 397: I would not say this. The friction velocity and dissipation are affected by both shear and buoyancy.

Response: Thank you very much for pointing this out. We have revised the sentence as “Generally, the near-surface buoyancy and shear effects both modulate these variables.”

17, line 408: turbulence intensity is different from turbulence kinetic energy. Do you mean turbulence intensity or turbulence kinetic energy?

Response: Thank you very much for pointing this out. We have revised it as “turbulence kinetic energy”.

18, line 412: replace “accorded” with “proposed”

Response: Revised as suggested.

19, line 427: add “the” before “highest”

Response: Revised as suggested.

20, line 468: I wonder what features on the figure led the authors to conclude “the cloud-mixed layer aloft does not interact with the near-surface environment”. The relative humidity is closer to saturation than figure 9 where the authors concluded “the near-saturated relative humidity indicates that the cloud-mixed layer couples with the surface-mixed layer, which facilitates the ABL development”. This needs to be clarified.

Response: Thank you very much for pointing this out. We realize that this statement is unclear. Actually, as mentioned in Shupe and Intrieri (2004) and Brooks et al. (2017), the Arctic ABL structure is highly dependent on atmospheric moisture and liquid-bearing clouds. In most of the year (e.g., Fig. 9 in original manuscript), liquid-bearing clouds can create more downwelling longwave radiation and result in an ABL that is well-mixed from the surface up to the top of the saturated layer, which indicates that the cloud-mixed layer couples with the surface-mixed layer. In the Arctic summer (e.g., Fig. 10 in original manuscript), low-level stratiform clouds form as a result of ample moisture available (Tjernström et al., 2012), but the clouds cannot contribute to the ABL development due to a strong temperature inversion maintained near the surface, which is different from other seasons. For this comment, we have revised the sentence as:

“Despite that, the ABL is still stably stratified, and the ample moisture and clouds cannot contribute significantly to the ABL development, which is consistent with Shupe et al. (2013).”

21, line 556-558: it’s unclear to me what the authors mean by “Coupling between the cloud mixed layer and surface mixed layer could also be recognized by the Rib

algorithm”. Does the Rib method can really distinguish this?

Response: Thank you very much for pointing this out. We realize that this statement is unclear. This statement aims to demonstrate that the *Ri* algorithm can estimate ABLH well in some cases of cloud-surface coupling, and the ABL development caused by the cloud effect can be captured by the *Ri* algorithm. In our improved *Ri* algorithm, the cloud effect is explicitly considered, which helps estimate ABLH better in the cloud coupling state. For this comment, we have revised the sentence as:

“The ABL development supported by cloud processes was captured by the improved *Ri* algorithm, which is similar to Brooks et al. (2017).”

Anonymous Referee #2

Summary: In this paper the authors perform an analysis of the boundary-layer height as observed from radio soundings during the MOSAiC field campaign in the Arctic in 2020. These observations are compared to PBL estimates from existing algorithms, and it is concluded that the critical Richardson number should amount to 0.15 rather than the traditional value of 0.25. The analysis has some potential, but at the same time the novelty is limited. My feeling is the paper does not build on the latest and most complete knowledge about PBL height estimation, especially not for the stable boundary layer. More can be learnt from this dataset, and I find the controversial result of $Ri_{crit} = 0.15$ should be discussed in more detail with findings elsewhere.

Response: Thank you very much for your time and effort in reviewing our manuscript and for the constructive comments. We have substantially revised the manuscript by addressing the comments, especially in the section that introduces and describes the PBLH algorithm. The revisions in the manuscript and the replies to the comments are marked in blue.

Major remarks:

1. The paper misses the opportunity to stratify the dataset of the PBL heights in more classes or groups. I.e. for example Zilitinkevich and co-workers have been working on PBL types as truly neutral PBLs, nocturnal PBLs, and conventionally neutral PBLs. In addition the analysis can be separated between profiles for cloudy/foggy vs clear sky conditions. I think this can help to reduce the scatter in Fig 3.

Response: Thank you very much for your helpful comment. We have reclassified the PBL types and added the neutral condition into our analysis. The sensible heat flux is used to determine the PBL types more credibly in our revised manuscript, which is similar to the buoyancy flux at the surface (B_s) applied by Zilitinkevich and co-workers. However, the truly neutral and nocturnal PBLs are not included in our PBL algorithm, because the stable and the neutral regimes are dominated by long-lived stable PBLs and conventionally neutral PBLs during the MOSAiC expedition, respectively. In addition, cloud conditions are also considered in our improved algorithm (see our response to your comment 2). For this comment, the corresponding changes are given in our revised manuscript as follows:

(1) Changes related to regime classification and ABLH determination

3.1 ABL regime classification and ABLH determination

The ABLH determination method starts with the classification of ABL regimes. Based on previous studies (e.g., Vogelezang and Holtslag, 1996; Liang and Liu, 2010), we divide the ABLs into three types: stable boundary layer (SBL), near-neutral boundary layer (NBL), and convective boundary layer (CBL), corresponding with three

different stability states near the surface. We first use SH to diagnose the ABL regime types. The specific classification formula is presented below:

$$\begin{cases} SH > +\delta & \text{for CBL} \\ SH < -\delta & \text{for SBL, (1)} \\ \text{else} & \text{for NBL} \end{cases}$$

where δ is the critical value that is specified as 2 W m^{-2} , following Steeneveld et al. (2007b). If corresponding SH data are unavailable, the difference of equivalent potential temperature (θ_E) between the 100 and 50 m heights (θ_E difference) derived from sounding profiles is used to determine the ABL types. Specifically, if θ_E difference is larger than 0.2 K, the ABL is identified as SBL; if θ_E difference is less than -0.2 K, the ABL is identified as CBL; and other profiles are labeled as NBLs, roughly following Liu and Liang (2010).

The manually-labeled ABLH determination in our study is based on characteristics of sounding profiles and regime types. For each atmospheric sounding profile, equivalent potential temperature (θ_E), equivalent potential temperature gradient (θ_{Egrad}), wind speed (WS), specific humidity (q_v), and RH are used to obtain multiple estimates of the ABLH, which are used to determine the final estimate. Three cases to describe the method are presented in Fig. 2. Figures 2 (a–c) are the case of a SBL, which features surface-based temperature and humidity inversions. Figures 2 (d–f) are the case of a NBL, with approximately constant θ_E from the surface up to the inversion base and strong horizontal wind. Figures 2 (g–i) are the case of a CBL, with a deeper well-mixed layer and low-level cloud coupled to the surface (e.g., Shupe et al., 2013). In terms of θ_E profiles, the estimated ABLH is the level at which the θ_{Egrad} reaches its maximum for SBL and NBL cases, and the base of the θ_E inversion for CBL cases (Martucci et al., 2007). In terms of WS profiles, the ABLH is estimated to be the height of the WS maximum for all three regime types (Mahrt et al., 1979). In terms of humidity profiles, the estimated ABLH is the level at which the RH rapidly decreases for SBL and NBL cases, and the base of the q_v inversion for CBL cases (Lenschow et al., 2000). The manually-observed ABLHs (solid black lines in Fig. 2) are then determined through consideration of these three distinct estimates using the following rules: (1) If the estimates differ slightly from each other, take the average of these estimates as ABLH; (2) If a strong characteristic (sharp gradients or peaks) of the profile is evident, select the estimate obtained based on this characteristic; (3) If the ABL structure is similar to that at the previous time, select the estimate with the smallest change to ensure that ABLHs are consistent in time. It is evident that the lowest layers of profiles have a great impact on the ABLH determination, particularly for shallow SBLs and NBLs. Thus, the merged radiosonde-tower profiles help make the ABLH determination more reliable than when using radiosondes alone

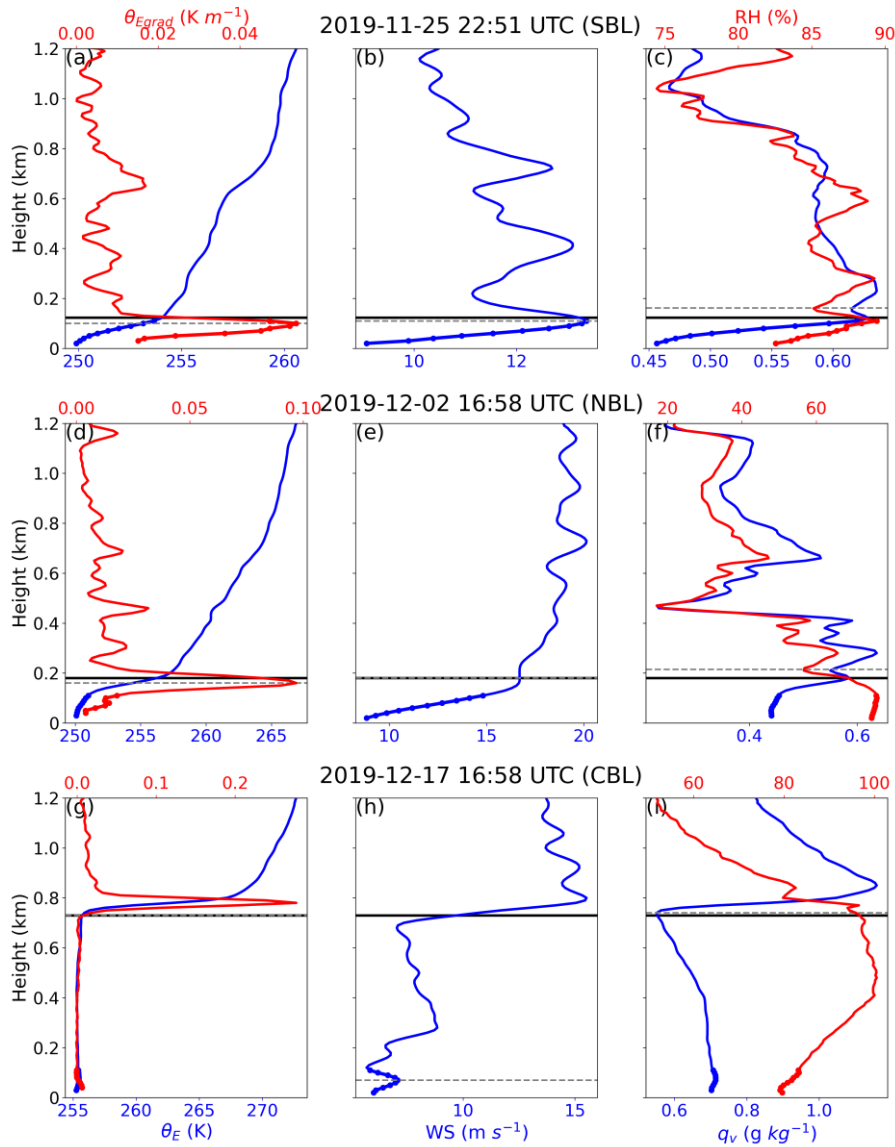


Figure 2 Vertical profiles of (left) equivalent potential temperature (θ_E), θ_E gradients (θ_{Egrad}), (middle) wind speed (WS), and (right) relative humidity (RH) and specific humidity (q_v) at (a–c) 25 November 2019, 22:51 UTC, (d–f) 2 December 2019, 16:58 UTC, and (g–i) 17 December 2019 16:58 UTC. Boundary layers at the three times represent stable boundary layer (SBL), near-neutral boundary layer (NBL), and convective boundary layer (CBL), respectively. The gray dashed horizontal lines denote the atmospheric boundary-layer height (ABLH) estimates based on multiple profiles, and the black solid horizontal lines denote the manually observed ABLHs. The dots in the lowest 100 m denote the section of the profiles impacted by the radiosonde-tower merging.

2. The study misses some novelty. I understand of course that the dataset at hand is unique and very valuable, but conceptually the paper does not add much in novelty for the PBL height detection. Would it be possible to come up with a completely new

approach or PBL height formula for the PBL depth, rather than “just” retuning the Ri_{crit} again as was done by so many other studies before?

Response: Thank you very much for your helpful comment. We have rewritten the section on algorithm improvement. As you suggested, we have included cloudy conditions in our improved algorithm. For clear-sky conditions, we use the finite-difference Ri formula proposed by Vogelezang and Holtslag (1996). While for cloudy conditions, we instead use the moist Richardson number Ri_m to take the cloud effect into account, and the Ri_m formula is updated based on Brooks et al. (2017). The Ri_m formula used in Brooks et al. (2017) is expressed as:

$$Ri_m = \frac{\left(\frac{g}{T}\right) \left(\frac{dT}{dz} + \Gamma_m\right) \left(1 + \frac{Lq_s}{RT}\right) - \frac{g}{1 + q_w} \frac{dq_w}{dz}}{\frac{du^2}{dz} + \frac{dv^2}{dz}},$$

where T is air temperature, Γ_m is the moist adiabatic lapse rate, L is the latent heat of vaporization, q_s is the saturation mixing ratio, and q_w is the total water mixing ratio, i.e., $q_w = q_s + q_L$, where q_L is the liquid water mixing ratio.

However, it is a gradient Ri and is calculated based on local gradients of wind speed, temperature, and humidity. In order to be consistent with the Ri formula proposed by Vogelezang and Holtslag (1996), we rewrite the formula in a finite-difference form expressed as:

$$Ri_m = \frac{\left[\left(g/T_s\right) \left(\frac{T_h - T_s}{h - z_s} + \Gamma_m\right) \left(1 + \frac{Lq_{sh}}{RT_h}\right) - \frac{g}{1 + q_{wh}} \frac{q_{wh} - q_{ws}}{h - z_s}\right] (h - z_s)^2}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2},$$

which is calculated based on the difference between the height h and lower reference height z_s . The results validate its feasibility (Figure 4). For this comment, the corresponding changes are given in our revised manuscript as follows:

(1) Changes related to cloud data description

2.3 Cloud properties derived from combined sensors

Cloud-related measurements come from ShupeTurner cloud microphysics product (Shupe, 2022). This product uses multiple measurement sources (e.g., cloud radar, ceilometer, depolarization lidar, and microwave radiometer) to derive time-height data, including cloud phase type and condensed water content for both liquid and ice. Details of the retrieval algorithm, its application, and uncertainties are provided in Shupe et al. (2015). In our study, the condensed water content data are linearly interpolated onto the vertical grid with resolution of 10 m for consistency. The cloud phase type data are used to determine clear and cloudy environments. A grid point is labeled as “cloudy” if clouds are identified in the upper and lower cloud phase type data points adjacent to the grid, otherwise it is labeled as “clear”.

(2) Changes related to detect cloud condition

We also analyze the algorithm performances for cloudy and clear conditions, considering that low-level clouds containing liquid water play an important role in the Arctic ABL (Shupe and Intrieri, 2004; Brooks et al., 2017). In our study, the RH threshold of 96% (Silber and Shupe, 2022) and the cloud source flag data are used for cloud detection. If a cloud is detected in the cloud source flag data and the RH is larger than 96%, then the profile is labeled as cloudy. The sounding profiles that contain at least one identified cloud layer below 1500 m are classified as “cloudy”, and as “clear” otherwise.

(3) Changes related to the improved Ri algorithm

3.3 An improved Ri algorithm considering the cloud effect

As a traditional Ri_b formula, Eq. (3) may break down in cases of ABLs with relatively high wind speed and upper-level stratification due to the overestimation of shear production (Kim and Mahrt, 1992). Vogelezang and Holtslag (1996) proposed the finite-difference Ri formula, which is expressed as:

$$Ri_F = \frac{(g/\theta_{vs})(\theta_{vh} - \theta_{vs})(h - z_s)}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (6)$$

where z_s is the lower boundary for the ABL, θ_{vs} , u_s , and v_s are the θ_v and wind components at the height z_s , respectively, b is an empirical coefficient, and u_* is the surface friction velocity. Ri_F is considered for a parcel located somewhat above the surface to avoid the above problem, and u_* is also taken into account to avoid underestimation in the situation of a uniform wind profile in the upper layer. Here, we use Ri_F for clear-sky profiles and take z_s and b values as 40 m and 100, respectively, according to Zhang et al. (2020).

As shown in Fig. 3, the estimations of cloudy ABLHs are sometimes quite poor, which motivates us to further improve the algorithm. Under cloudy conditions, the moist Richardson number (Ri_m) can be used to include cloud effects on the buoyancy term. Brooks et al. (2017) adopted the Ri_m formula expressed as:

$$Ri_m = \frac{\left(\frac{g}{T}\right)\left(\frac{dT}{dz} + \Gamma_m\right)\left(1 + \frac{Lq_s}{RT}\right) - \frac{g}{1 + q_w} \frac{dq_w}{dz}}{\frac{du^2}{dz} + \frac{dv^2}{dz}}, \quad (7)$$

where T is air temperature, Γ_m is the moist adiabatic lapse rate, L is the latent heat of vaporization, q_s is the saturation mixing ratio, and q_w is the total water mixing ratio, i.e., $q_w = q_s + q_L$, where q_L is the liquid water mixing ratio and is obtained based on the condensed water content. However, Eq. (6) is a gradient Ri and is calculated based on local gradients of wind speed, temperature, and humidity. To be consistent with the Ri formula proposed by Vogelezang and Holtslag (1996), we rewrite the formula in a finite-difference form expressed as:

$$Ri_m = \frac{\left[\left(g/T_s\right)\left(\frac{T_h - T_s}{h - z_s} + \Gamma_m\right)\left(1 + \frac{Lq_{sh}}{RT_h}\right) - \frac{g}{1 + q_{wh}} \frac{q_{wh} - q_{ws}}{h - z_s}\right] (h - z_s)^2}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (8)$$

where subscripts (h and s) of the variables denote the calculated height, similar to Eq. (6), but note that the s and z_s are adjusted to 130 m, given the cloud radar blind zone.

Considering that Ri_m is only appropriate for the liquid-bearing cloud cases, we use the Ri_F for “clear” grid points and use Ri_m for “cloudy” grid cells. Using this improved approach, we evaluated the best value of Ri_c to minimize the errors compared to the reference data set, arriving at an optimal value of $Ri_c=0.35$. The comparison of ABLH estimates obtained through the improved Ri algorithm with the manually-labeled ABLHs demonstrates significant improvement relative to other algorithms, particularly for cloudy conditions (Fig. 4, Table 1).

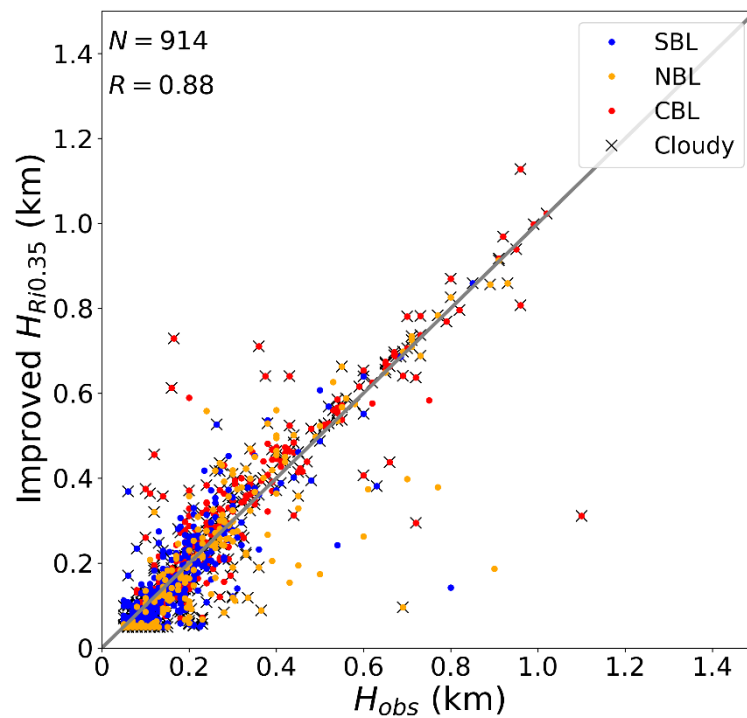


Figure 4 Similar to Fig. 3, but for the comparison of the ABLHs determined by the improved Ri algorithm with the observed ABLHs. The case number (N) and correlation coefficient (R) are given.

Table 1 The statistical measures (R , $Bias$, $MEAE$) for the four algorithms applied to the radiosonde dataset. All correlation coefficients are statistically significant ($p < 0.05$), except for SBL types in the Liu-Liang algorithm.

Algorithm	Regime type	R	Bias	MEAE (m)
The Ri_b algorithm with $Ri_{bc} = 0.25$	ALL	0.72	0.10	50
	SBL	0.81	0.16	34
	NBL	0.68	-0.04	62
	CBL	0.65	0.15	71
	Cloudy	0.69	0.08	51
The Ri_b algorithm with $Ri_{bc} = 0.5$	ALL	0.67	0.40	97
	SBL	0.73	0.50	88
	NBL	0.61	0.23	91
	CBL	0.60	0.39	120
	Cloudy	0.66	0.36	94
The Heffter algorithm	ALL	0.57	0.23	53
	SBL	0.46	0.17	33
	NBL	0.45	0.30	59
	CBL	0.66	0.28	74
	Cloudy	0.68	0.25	59
The Liu-Liang algorithm	ALL	0.47	0.04	82
	SBL	0.05	0.15	90
	NBL	0.44	-0.07	81
	CBL	0.56	-0.05	69
	Cloudy	0.52	-0.01	82
The improved Ri algorithm with $Ri_{bc} = 0.35$	ALL	0.85	-0.06	29
	SBL	0.79	-0.08	21
	NBL	0.79	-0.18	35
	CBL	0.87	0.05	36
	Cloudy	0.86	-0.03	30

(4) Changes related to Figure 7

4.1 Overall distribution of ABLH

In this section, we analyze the ABLH variation during the MOSAiC year and relevant controlling factors, based on the manually-labeled ABLH dataset and the ABL types that are determined through Eq. (1), or only the θ_E difference if SH is unavailable. The full-time series of ABLH during the MOSAiC expedition is presented in Fig. 7 and forms the basis for the remaining analyses. According to near surface conditions and the sea ice state, the whole MOSAiC observation period is divided into “freeze up”, “winter”, “transition”, and “summer melt” periods (Shupe et al., 2022), roughly corresponding to the seasons of autumn, winter, spring, and summer, respectively. In Figure 7, the black solid lines indicate persistent low-level clouds that exist for more than 12 h; these occur most frequently in the late summer and autumn (the “freeze up” period), which agrees with Shupe et al. (2011). Note that the grey dots indicate that the

ABL data were observed while the vessel was in transit, and the representativity of the ABLH data should be considered in this context. For the first such period, the vessel left the MOSAiC ice floe in mid-May and slowly progressed south through tightly consolidated sea ice, such that the data are generally representative of the sea ice pack in the region. Measurements from early June when the vessel was near or in open water close to Svalbard have been excluded entirely from the analysis. In the middle of June, as the vessel returned to the original MOSAiC ice floe, the sea ice was not as tightly consolidated and the vessel preferentially went through leads; the preferentially lower ice fraction along this transit could have impacted the thermal structure of the ABL. For the three weeks in early August, the vessel moved around in the Fram Strait area and then made its way north to another passive sea ice drifting position near the North Pole, again transiting through regions with lower sea ice fraction. Finally, at the very end of the expedition, the vessel took some time to exit the sea ice, stopping a few times to allow for work on the ice.

Overall, as shown in Fig. 7, the mean ABLH during the whole observation period is 231 m. This is lower than the typical ABLH over the Arctic land surface (Liang and Liu, 2010), which is primarily attributed to the stronger suppression of the temperature inversion over the sea-ice surface. The Arctic ABL is suppressed for most of the MOSAiC year, while for a few periods it intensively develops for several days at a time, most commonly when clouds and a CBL are present. For instance, frequent, intensive ABL development occurs in the “transition” period from 13 April through to 24 May 2020. In this period, the convective thermal structure and cloud effects contribute to ABLH reaching over the 95th percentile of the ABLH data (horizontal dotted line) for about 7 days, with the maximum ABLH of 1100 m. In contrast, the ABL is strongly suppressed in the period from 15 July through to 30 August 2020, with a mean ABLH of only 136 m. The specific mechanisms of ABL development and suppression in these two cases will be analyzed in Sections 4.3 and 4.4, respectively.

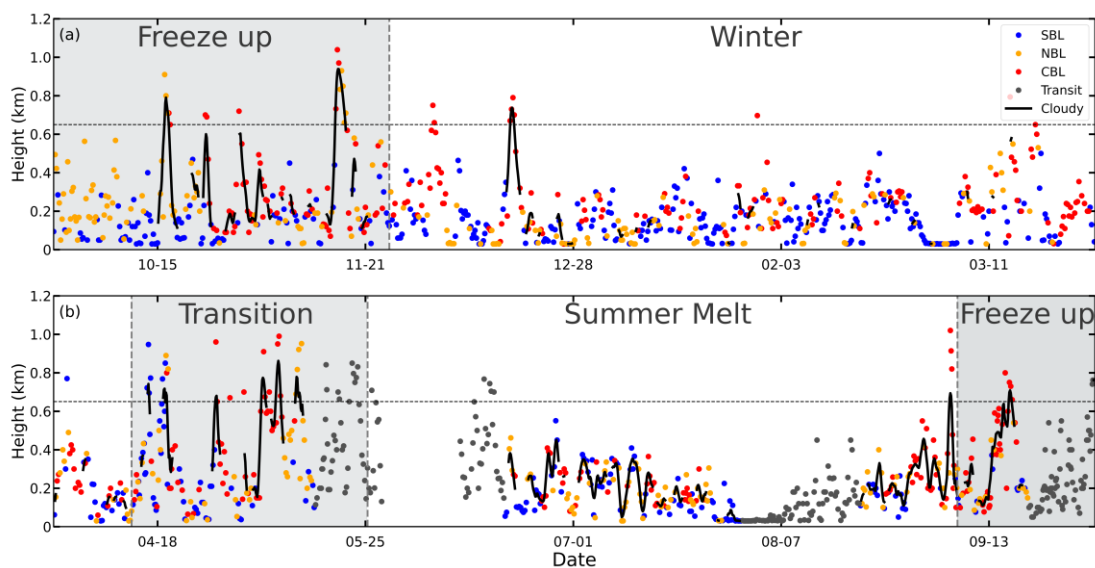


Figure 7 Time series of ABLHs throughout the MOSAiC year is divided into (a) and

(b). The blue, yellow, and red dots indicate the heights of SBL, NBL, and CBL, respectively. The gray dots indicate ABL data observed while the vessel was in transit. The black solid lines indicate the heights of cloudy ABLs and persist for at least 12 hours. The gray dashed horizontal line denotes the 95th percentile of ABLH (650 m). The gray and white background shadings indicate the periods under different surface-melting states, i.e., “freeze up”, “winter”, “transition”, and “summer melt” periods.

(5) Changes related to Figure 8

Figure 8 presents the frequency distribution of ABLH under SBL, NBL, and CBL regime types. Overall, the sample number of SBL cases is more than that of NBL and CBL cases during the MOSAiC period (43 % for SBL, 31% for NBL, and 26 % for CBL). These occurrence frequencies roughly agree with Jozef et al. (2023), while their results show more NBL and CBL and less SBL. It is likely to be attributed to differences in classification criteria. The distributions of SBL and NBL ABLH are skewed towards small values, with 94 % and 79% of the ABLH values lower than 400 m, and mean values of 165 m and 256 m, respectively. For CBL, the distribution is shifted somewhat towards larger values, with 23 % of the ABLH values higher than 600 m and a mean value of 309 m.

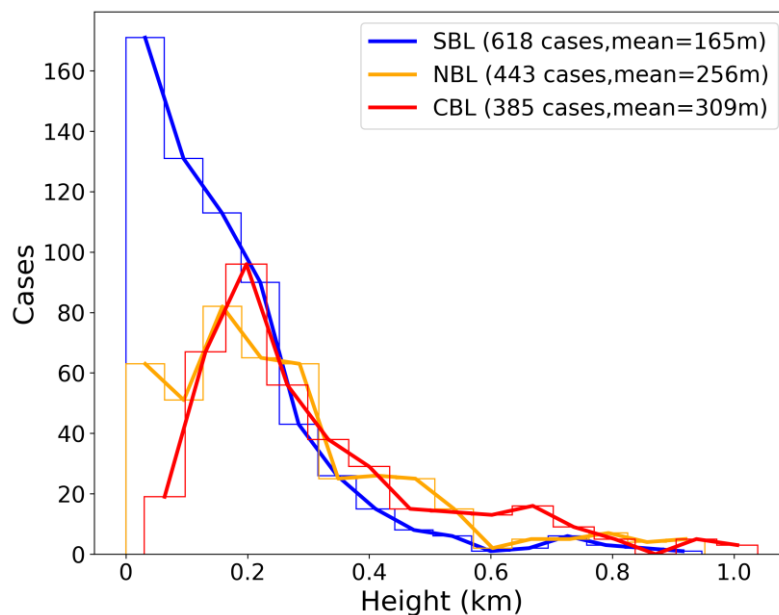


Figure 8 Frequency distribution of SBL height (blue), NBL height (yellow), and CBL height (red). The case numbers and the mean values of ABLH for SBL, NBL, and CBL conditions are also given.

(6) Changes related to Figure 9

Figure 9 presents the annual cycle of monthly ABLH statistics during the MOSAiC expedition in terms of 5th, 25th, 50th, 75th, and 95th percentiles of ABLH (boxplots) and the mean value (“x” signs and solid and dashed lines). The box-and-whisker plots show a distinct peak in May, with a median value of 363 m and the 95th percentile reaching over 800 m. An abrupt decrease occurs in the following July and August, and

another minimum occurs in January, all with median values below 150 m. It should be noted that the ABLH data in transit (gray dots in Fig. 7) are also included in the statistics, which could have potential impact from somewhat more open-water surface conditions. Specifically, the ABLH data during transit periods cause higher mean ABLH for June and lower mean ABLH for August (see Fig. 7). The comparison between cloudy and clear-sky ABLHs indicates that the low-level clouds significantly contribute to the Arctic ABL development during the MOSAiC year, except in winter, when low-level clouds are rare.

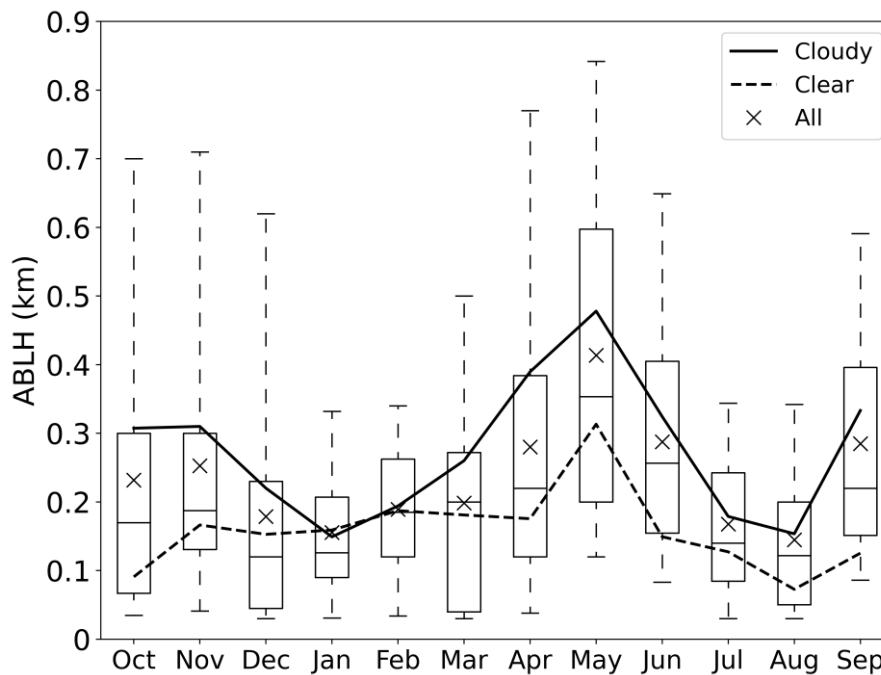


Figure 9 Box-and-whisker plots of the ABLH distribution in each month throughout the MOSAiC year. The whiskers, the boxes, and the black horizontal lines show the 5th, 25th, 50th, 75th, and 95th percentile values of ABLH. The solid and dashed lines and the “x” signs indicate the mean ABLH of cloudy, clear, and all ABL types, respectively.

3. The authors have missed a paper by Barten et al. (2023) in the MOSAIC special issue in *Elementa* in which a similar PBL height analysis was performed. While the main focus of that paper is on the ozone budget in the Arctic PBL, it reports that the critical Richardson number should be 0.40 for MOSAIC for the same set of radio soundings. Hence this is above the typical value of 0.25, while the authors here propose 0.15. This is an obvious contradiction that needs to be discussed.

Response: Thank you very much for your helpful comment. The determination of the critical Richardson number depends on the choice of Ri formula. In our original manuscript, we proposed a Ri_c value of 0.15 based on a traditional Ri_b formula Eq. (3). This result is different from Barten et al. (2023), who also use the Eq. (3) to propose the Ri_c of 0.4. This difference may be caused by the different methods to derive the comparison ABLH dataset with which to identify the Ri_c . However, Barten et al. (2023) do not provide sufficient details on how their manual comparison ABLH dataset

was derived. On the other hand, Akansu et al. (2023) also proposed $Ri_c = 0.12$ for MOSAiC observations based on Eq. (3), which roughly agrees with our original Ri_c value of 0.15. It is worth mentioning that Akansu et al. (2023) determined the observed PBLH precisely from turbulence profiles and obtained a Ri_c close to ours, which supports the reliability of our manually-labeled PBLHs. However, as mentioned in our response to your Comment 2, we have updated the Ri formula by now including cloud effects, and for that formula the new Ri_c of 0.35 is identified. As you expected, this improvement reduces the scatter in our original Fig 3. For this comment, the corresponding changes are given in our revised manuscript as follows:

Changes related to discuss difference in Ri_c by previous studies

Since some other studies have proposed different Ri_c values for MOSAiC (e.g., Jozef et al., 2022; Barten et al., 2023; Akansu et al., 2023), we will discuss the difference in Ri_c values here. The first thing to make clear is that these studies use different formulas to obtain Ri profiles. Barten et al. (2023) and Akansu et al. (2023) both use the traditional Ri_b algorithm based on Eq. (3), while they used Ri_c values of 0.4 and 0.12, respectively. This difference was likely caused by the different methods to manually derive their reference ABLH data sets. Jozef et al. (2022) calculates the Ri over a rolling 30 m altitude range, labeled as Ri_r , and the criterion is modified to require four consecutive data points to be above the Ri_c of 0.75. In our study, we use Ri_F proposed by Vogelezang and Holtslag (1996) for clear-sky conditions, and Ri_m for cloudy conditions. Based on the results presented here, it is apparent that this more complex approach improves the error statistics relative to approaches based on Eq. (3), regardless of Ri_c . In addition, some of the differences may also related to authors using different data sets or time periods. For instance, Akansu et al. (2023) primarily used sounding data based on tether balloon for a specific sub-period of MOSAiC, and Jozef et al. (2022) used radiosondes from when they had concurrent UAV observations. The data used in our study are based on merged sounding-tower product, as mentioned above.

To further explore the differences among the four different Ri approaches, we examine one SBL and CBL case. For a clear-sky SBL case (Fig. 5 a, b), the approaches from Akansu et al., Jozef et al. (2022), and this study all agree closely with the manual ABLH, while the Barten et al. approach results in a significant overestimation. For a cloudy-sky CBL case (Fig. 5 c d), the approach from this study agrees with the manual ABLH, while the approach from Barten et al. overestimates the ABLH by about 30 m, and the approaches from Akansu et al. and Jozef et al. (2022) underestimate the ABLH by 130 m and 230 m, respectively. These results further demonstrate how Ri_c depends on the choice of Ri formula. Moreover, Ri_c is not analytically derived from basic physical principles (Zilitinkevich et al. 2007), and the concept of Ri_c is challenged by non-steady regimes (Zilitinkevich and Baklanov, 2002) and the hysteresis phenomenon (Banta et al., 2003; Tjernström et al., 2009). Therefore, an objective Ri_c does not exist. Rather, it is empirically used as an algorithmic parameter to simply derive the ABLH.

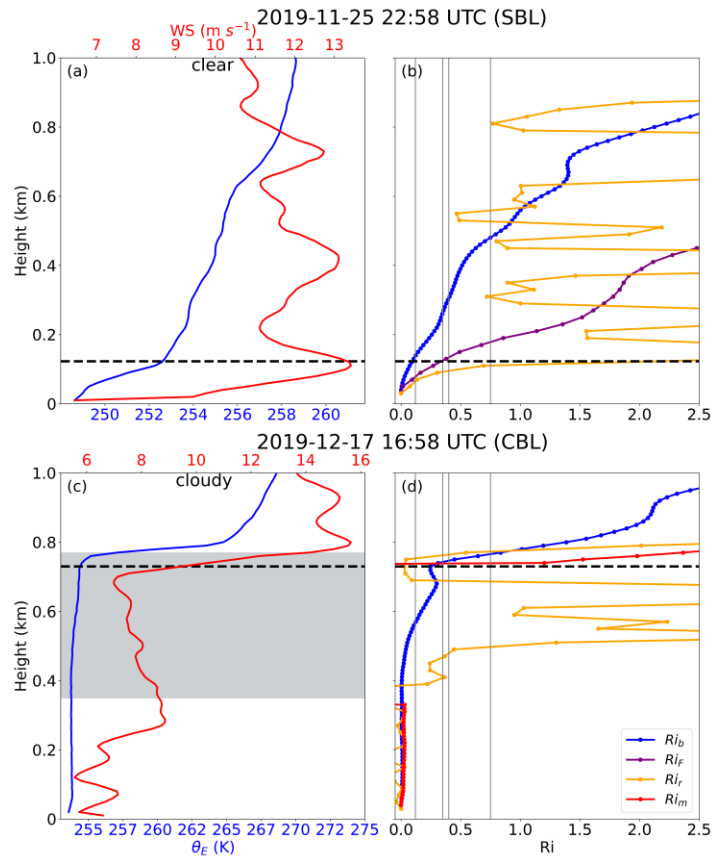


Figure 5 Vertical profiles of (left) θ_E and wind speed, and (right) Ri based on different formulas at (a–b) 25 November 2019, 22:58 UTC and (c–d) 17 December 2019, 16:58 UTC. Boundary layers at the two times represent a clear-sky SBL and a cloudy-sky CBL respectively. The black dashed horizontal lines denote the manually-identified ABLH, and the gray solid vertical lines denote the different Ri_c values, including 0.12, 0.35, 0.4, and 0.75. The gray shading in (c) denotes the cloud layer.

4. The discussion section of the paper can be deepened in the sense that the Ri_{crit} value has been widely discussed in other papers before, but I do miss some important ones in the review, e.g. Zilitinkevich and Baklanov (2002, <https://link.springer.com/article/10.1023/A:1020376832738>). Also Basu et al (2014) proposes that the Ri_{crit} depends on the stability of the SBL as well (<https://link.springer.com/article/10.1007/s10546-013-9878-y>). Also, Equation 2 used in the paper has been revised already by Vogelezang and Holtslag for a better score, but it feels this paper does not take benefit from this knowledge. Also, earlier LES studies for the SBL height formula are not mentioned. Hence, the current paper can be embedded more in these earlier works/contributions.

Response: Thank you very much for your helpful suggestion. We have added a discussion section to analyze the stability dependence of Ri_c in stable conditions, and

the results validate this relationship for the MOSAiC dataset. For improving the algorithm that we use here, we have considered the Ri formula proposed by Vogelesang and Holtslag (1996) in our ABLH algorithm, as mentioned above. The SBL height formulas based on earlier LES studies are also tested. For this comment, the corresponding changes are given in our revised manuscript as follows:

(1) Changes related to discussion on stability dependence of Ri_c

3.4 The stability dependence of critical Richardson number

Richardson et al. (2013) and Basu et al. (2014) suggested that there is a stability dependence of Ri_c in stable conditions, which is different from the constant $Ri_c = 0.35$ used in our improved algorithm. In this section, we will discuss the impact of this dependence on ABLH estimation. We use the improved Ri algorithm to calculate the Ri at the manually-labeled ABLH (h). This new parameter is named Ri_h to distinguish it from the constant Ri_c . To be consistent with Basu et al. (2014), the bulk stability parameter h/L is used for our analysis, where L is the Obukhov length. Based on these two variables, the stability dependence can be expressed as:

$$Ri_h = \alpha \frac{h}{L}, \quad (9)$$

where α is a proportionality constant. As suggested in Basu et al. (2014), the data for convective, near-neutral, and very stable conditions are excluded to obtain a credible α . Specifically, data points that meet the thresholds ($L > 500$ m and $L < L_{min}$) are excluded in our analysis, where the L_{min} corresponds to the heat flux minimum (Basu et al. 2008) and is assumed as 20 m here. Finally, we select 168 samples. The Ri_h plotted as a function of h/L for these selected data is presented in Fig. 6, and the value of L is colored to probe if the dependence is simply due to self-correlation. The results show Ri_h values that mostly range from 0 to 0.75, and the best-fit line indicates an overall positive correlation trend, with $\alpha = 0.11$. The α value is somewhat larger than the results in Richardson et al. (2013) and Basu et al. (2014), which is attributed to the different Ri algorithm used in our study. In addition, if a few of the extreme points are removed, the bulk of the data does not show a strong h/L dependence and is instead fairly well represented by a constant $Ri_h = 0.35$, which is also suitable for convective conditions (e.g., Fig. 5c, d).

In summary, we assess the stability dependence of Ri_c based on our improved Ri algorithm, and the results present an overall positive correlation trend. However, this type of stability dependence of Ri_c is challenged to be used in practical applications because the sensitivity of α to surface characteristics and atmospheric conditions can additionally degrade the accuracy of ABLH estimates. In addition, Eq. (9) requires a priori determination of the ABLH, which also causes difficulties for practical applications of such an approach. Therefore, we still use the Ri algorithm with fixed $Ri_c = 0.35$ for simplicity.

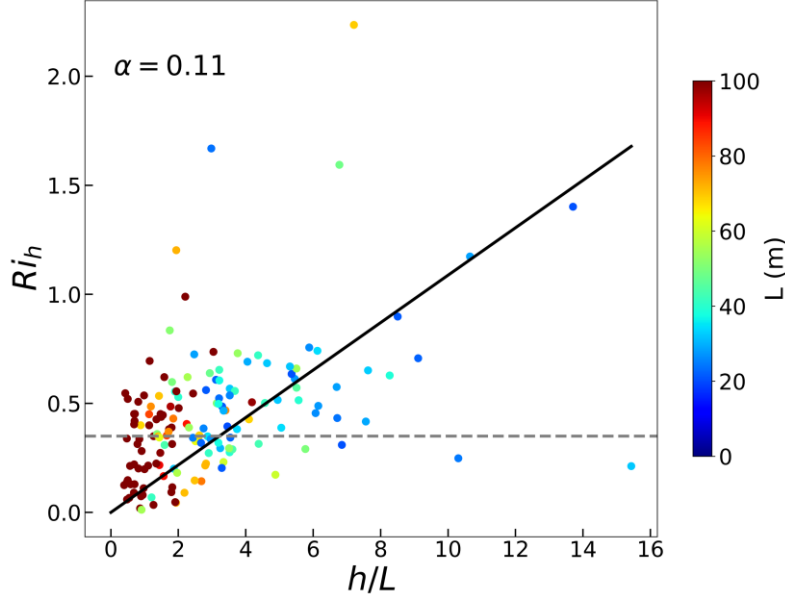


Figure 6 Ri_h versus h/L for selected cases. The data points are colored based on the value of L . The black solid line is the best fit for the selected data points, and the best-fit α value is also given. The gray dashed line is the constant $Ri_c = 0.35$ used in the improved Ri algorithm.

(2) Changes related to earlier LES studies for the SBL height formula

The free-flow stability (characterized by the free-flow Brunt-Väisälä frequency, N) can affect the ABLH (Zilitinkevich et al., 2002; Zilitinkevich and Baklanov, 2002; Zilitinkevich and Esau, 2002, 2003), and therefore is also examined here. Based on the buoyancy flux at the surface (B_s) and N , the NBLs and SBLs can be further divided into four types: the truly neutral (TN, $B_s = 0$ and $N = 0$), the conventionally neutral (CN, $B_s = 0$ and $N > 0$), the nocturnal stable (NS, $B_s < 0$ and $N = 0$), and the long-lived stable boundary layer (LS, $B_s < 0$ and $N > 0$). According to Zilitinkevich and Baklanov (2002), we calculate the N and B_s and reclassify the SBLs and NBLs. We find that the percentages of $N > 0.015$ in SBLs and NBLs are 89 % and 80 %, which indicates that LS and CN types dominate the stable and neutral conditions for MOSAiC, respectively. Since only 80 TN cases were identified, these are deemed to be too few for additional analysis of this type. Zilitinkevich and Esau (2003) gave ABLH equations relevant to each ABL type as:

$$h_E = \begin{cases} C_N u_* |fN|^{-1/2} & \text{(Pollard et al., 1973) for CN ABL, (10)} \\ C_S u_*^2 |fB_s|^{-1/2} & \text{(Zilitinkevich, 1972) for NS and LS ABL, (11)} \end{cases}$$

where h_E is the equilibrium ABLH, f is the Coriolis parameter, and C_N and C_S are empirical coefficients. In addition, Vogelesang and Holtslag (1996) and Steeneveld et al. (2007a) also explore a h_E equation without taking into account f explicitly, expressed as:

$$h_E = C_i \frac{u_*}{N} \text{ for all SBL and NBL, (12)}$$

where C_i is an empirical coefficient. Here we select the CN, NS, and LS ABLH dataset, and fit the data with the corresponding expressions in Eq. (10–12) to obtain the empirical coefficients, and the results are presented in Fig. 12. All three expressions tend to well represent the ABLHs, with significant correlation coefficients. The empirical coefficients C_N and C_S are 1.7 and 0.4, respectively, which are close to the typical values determined through large-eddy simulations (Zilitinkevich, 2012). The coefficient $C_i = 20$ in Fig. 12c is double the typical value of 10 (Vogelezang and Holtslag, 1996), but agrees with the results reported by Overland and Davidson (1992) for the ABL over sea ice. The difference in C_i may be attributed to the unique free-flow stability or other potential mechanisms of ABL development in the Arctic atmosphere.

In summary, near-surface conditions and free-flow stability play a key role in ABL development and are also an indicator, in that one can roughly determine the development state of the whole ABL from these basic variables.

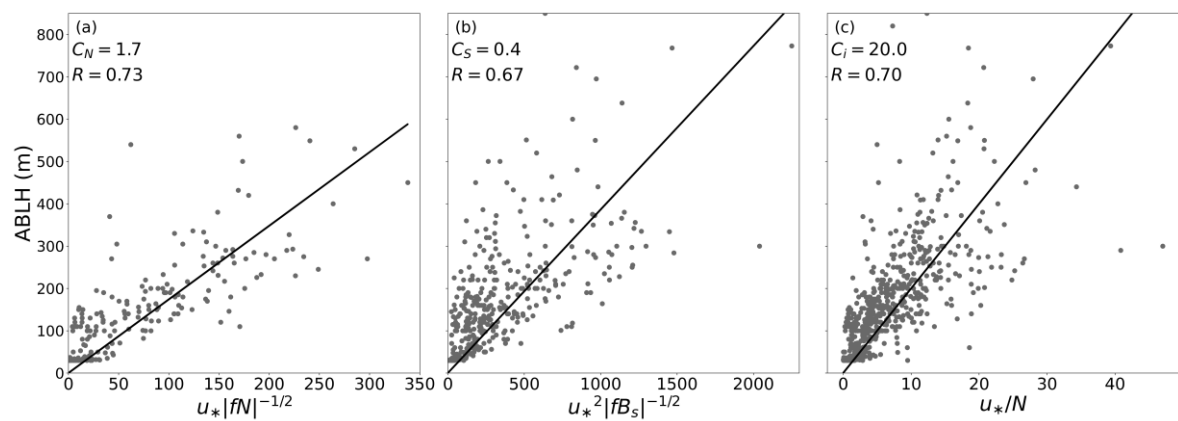


Figure 12 The ABLHs versus three expressions in Eq. (10–12). The empirical coefficients C_N , C_S , and C_i are given in (a), (b), and (c), respectively, and represent the slope of the best fit line (black line). The correlation coefficient R is given in each panel, which is statistically significant ($p < 0.05$).

5. I was surprised that the paper never discusses whether a critical Richardson number should exist anyway. In the EFB papers by Zilitinkevich it is analytically derived that the Ri_{crit} does formally not exist. Though I understand that in practical applications of Ri_{crit} still can have some value.

Response: Thank you very much for your helpful comment. We have added the relevant discussion in our revised manuscript. The corresponding changes are given as follows:

Since some other studies have proposed different Ri_c values for MOSAiC (e.g., Jozef et al., 2022; Barten et al., 2023; Akansu et al., 2023), we will discuss the

difference in Ri_c values here. The first thing to make clear is that these studies use different formulas to obtain Ri profiles. Barten et al. (2023) and Akansu et al. (2023) both use the traditional Ri_b algorithm based on Eq. (3), while they used Ri_c values of 0.4 and 0.12, respectively. This difference was likely caused by the different methods to manually derive their reference ABLH data sets. Jozef et al. (2022) calculates the Ri over a rolling 30 m altitude range, labeled as Ri_r , and the criterion is modified to require four consecutive data points to be above the Ri_c of 0.75. In our study, we use Ri_F proposed by Vogelezang and Holtslag (1996) for clear-sky conditions, and Ri_m for cloudy conditions. Based on the results presented here, it is apparent that this more complex approach improves the error statistics relative to approaches based on Eq. (3), regardless of Ri_c . In addition, some of the differences may also related to authors using different data sets or time periods. For instance, Akansu et al. (2023) primarily used sounding data based on tether balloon for a specific sub-period of MOSAiC, and Jozef et al. (2022) used radiosondes from when they had concurrent UAV observations. The data used in our study are based on merged sounding-tower product, as mentioned above.

To further explore the differences among the four different approaches, we examine one SBL and CBL case. For a clear-sky SBL case (Fig. 5 a, b), the approaches from Akansu et al., Jozef et al. (2022), and this study all agree closely with the manual ABLH, while the Barten et al. approach results in a significant overestimation. For a cloudy-sky CBL case (Fig. 5 c d), the approach from this study agrees with the manual ABLH, while the approach from Barten et al. overestimates the ABLH by about 30 m, and the approaches from Akansu et al. and Jozef et al. (2022) underestimate the ABLH by 130 m and 230 m, respectively. These results further demonstrate how Ri_c depends on the choice of Ri formula. Moreover, Ri_c is not analytically derived from basic physical principles (Zilitinkevich et al. 2007), and the concept of Ri_c is challenged by non-steady regimes (Zilitinkevich and Baklanov, 2002) and the hysteresis phenomenon (Banta et al., 2003; Tjernström et al., 2009). Therefore, an objective Ri_c does not exist. Rather, it is empirically used as an algorithmic parameter to simply derive the ABLH.

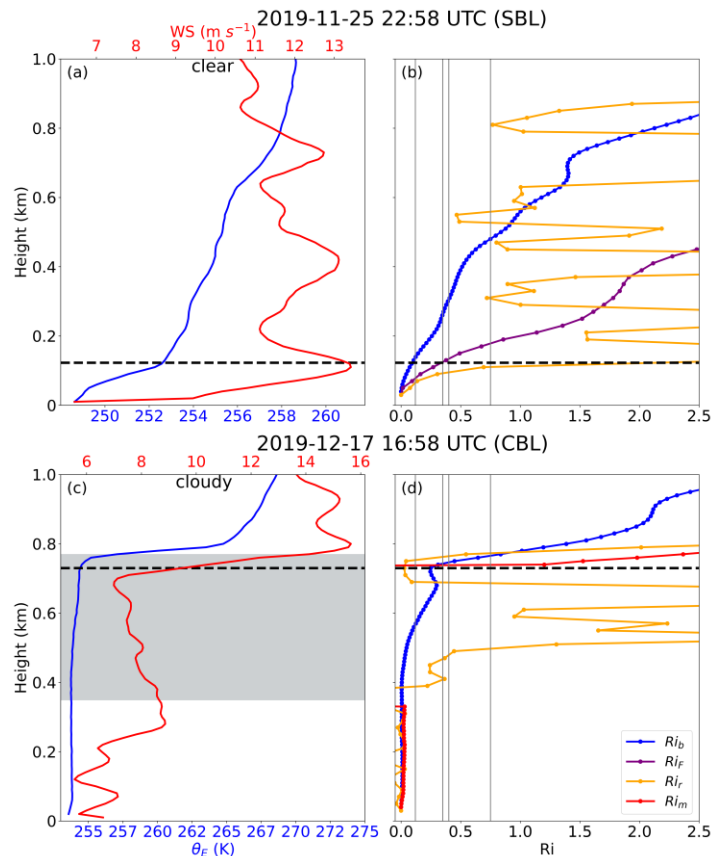


Figure 5 Vertical profiles of (left) θ_E and wind speed, and (right) Ri based on different formulas at (a–b) 25 November 2019, 22:58 UTC and (c–d) 17 December 2019, 16:58 UTC. Boundary layers at the two times represent a clear-sky SBL and a cloudy-sky CBL respectively. The black dashed horizontal lines denote the manually-identified ABLH, and the gray solid vertical lines denote the different Ri_c values, including 0.12, 0.35, 0.4, and 0.75. The gray shading in (c) denotes the cloud layer.

Minor remarks:

Ln 14: hyphenation: boundary layer height -> boundary-layer height. Please check whole document.

Response: Revised as suggested.

Ln 17: perhaps it is good to mention in then abstract before coming up with a new RI_crit how you defined the ABLH in your study. I.e. the level of the largest d_theta/dz , the backscatter level of a ceilometer, the low-level jet height, etc etc.?

Response: Thank you very much for your helpful suggestion. The ABLH in our study

is defined as the layer of continuous turbulence adjacent to the surface. Profiles of equivalent potential temperature, wind speed, and humidity are used in the manual ABLH determination method. The corresponding change is given in our revised manuscript as follows:

The important roles of the atmospheric boundary layer (ABL) over the Arctic Ocean in the Arctic climate system have been recognized, but the atmospheric boundary-layer height (ABLH), defined as the layer of continuous turbulence adjacent to the surface, has rarely been investigated.

Ln 33: Kwok, 2018; Hartfield et al., 2018. I have nothing against these studies but are they still recent?

Response: Thank you very much for your helpful comment. We have introduced the recent studies. The corresponding change is given in our revised manuscript as follows:

In recent years, the rapidly changing climate and declining sea ice in the Arctic have been reported by numerous studies (e.g., Matveeva and Semenov, 2022; Meier and Stroeve, 2022; Esau et al., 2023).

Ln 42: “various mechanisms and interactions with the surface”: I would say the opposite since turbulent fluxes in the Arctic are usually small so the interaction with the surface is small. In the hierarchy of PBL types by Zilitinkevich et al the Arctic PBL height is characterised as a long-lived stable boundary layer where the PBL height scales more with the stratification in the free atmosphere (and wave activity therein) than with the fluxes at the surface.

Response: Thank you very much for pointing this out. We have revised the sentence as “The ABL structure over the Arctic Ocean has unique characteristics due to the presence of semipermanent sea ice, and is shaped by various mechanisms including the interactions with the surface, free atmosphere and wave activity.”

Ln 51: The study by Sterk et al (2014) nicely summarizes this (<https://doi.org/10.1002/jgrd.50158>).

Response: Thank you very much for pointing this out. We have introduced this study. The corresponding change is given in our revised manuscript as follows:

Investigations of the ABL structure evolution and its controlling factors are the keys to knowing the ABL's role in the Arctic atmosphere (Sterk et al., 2014).

Ln 56: There are many more recent studies that indicate this as well than Deardorff,

1972; Suarez et al., 1983; Holtslag and Nieuwstadt, 1986. Please connect to the recent work!

Response: Thank you very much for your suggestion. We have introduced the recent studies. The corresponding change is given in our revised manuscript as follows:

... and is an important parameter for weather and climate models (Holtslag et al., 2013; Mahrt, 2014; Davy and Esau, 2016).

Ln 109: over the altitude range of 12 m up to 30 km. Please add what is the typical vertical resolution of the sounding measurements in the profile near the surface, this is important to know to what extent the ABLH can be well estimated.

Response: Thank you very much for your helpful comment. According to the description of the data, the radiosondes ascend at a rate of approximately 5 m s^{-1} , sampling with a frequency of 1 Hz, which indicates that the vertical resolution of the sounding measurements is 5 m. We have added the information into the revised manuscript as follows:

The radiosoundings provide data on the atmospheric state, including vertical profiles of pressure, temperature, relative humidity (*RH*), and winds, from 12 m up to 30 km with the vertical resolution of 5 m.

Ln 116: Moreover, we cut off the sounding data observed below 100 m altitude considering the potential contamination of the vessel itself. Please add how many of the launches had to be excluded because of the restriction.

Ln 116: Moreover, we cut off the sounding data observed below 100 m altitude considering the potential contamination of the vessel itself. The ABLH is typically shallow in the Arctic, so is the part that is eliminated not exactly the part you are interested in.

Ln 116: the section should finish with a statement how many soundings are available for analysis after all the correction and control exercises.

Response: Thank you very much for your helpful comment. When we neglect data below 100m this does have an impact on the ABLH determination, particularly for shallow SBLs. Therefore, we replaced the original data with a new merged sounding dataset, which combines the soundings with the meteorological tower data on the sea ice (Dahlke et al., 2023) with the specific goal of correcting for ship effects and providing more reliable profiles in the lowest 100 m. This new dataset allowed us to now use 1484 sounding profiles available. We updated the results and found significant improvements in SBL height determination and estimation. Also, we removed the high-resolution sounding data and used this merged data in ABLH variation section for consistency. The relevant statement is added into our revised manuscript as follows:

(1) Changes related to data description

2.1 Radiosonde observations and relevant data products

The radiosonde data were obtained through a partnership between the leading Alfred Wegener Institute (AWI), the atmospheric radiation measurement (ARM) user facility, a US Department of Energy facility managed by the Biological and Environmental Research Program, and the German Weather Service (DWD) (Maturilli et al., 2022). Vaisala RS41-SGP Radiosondes were regularly launched on board throughout the whole MOSAiC year (from October 2019 to September 2020), including periods when the vessel was in transit. The sounding frequency is normally four times per day (launched at about 5:00, 11:00, 17:00, and 23:00 UTC) and is increased to 7 times per day during periods of exceptional weather or coordination with other observing activities. The radiosoundings provide data on the atmospheric state, including vertical profiles of pressure, temperature, relative humidity (*RH*), and winds, from 12 m up to 30 km with a vertical resolution of 5 m. However, the sounding data below ~100 m altitude may be contaminated by the vessel itself. To avoid contamination affecting our analysis, we use a merged data product that combines the soundings with measurements from a meteorological tower on the sea ice away from the vessel, and was specifically designed to minimize ship effects and provide more reliable profiles in the lowest 100 m, which has been recently submitted (Dahlke et al., 2023). In this paper, data quality control and a six-point moving average in height are applied to the merged profile data to eliminate invalid data and measurement noise, and all data are interpolated onto a regular vertical grid with 10 m intervals. In total, there are 1484 sounding profiles available. In addition, DOE-ARM provides a Planetary Boundary Layer Height Value-Added Product (PBLHT VAP, Riihimaki et al., 2019), which uses several different automated algorithms to compute ABLH estimates based on radiosonde profiles. This VAP provides 964 ABLH estimates, and we select 914 samples from these to ensure that the estimates obtained by all algorithms are available.

(2) Changes related to ABLH determination

It is evident that the lowest layers of profiles have a great impact on the ABLH determination, particularly for shallow SBLs and NBLs. Thus, the merged radiosonde-tower profiles help make the ABLH determination more reliable than when using radiosondes alone.

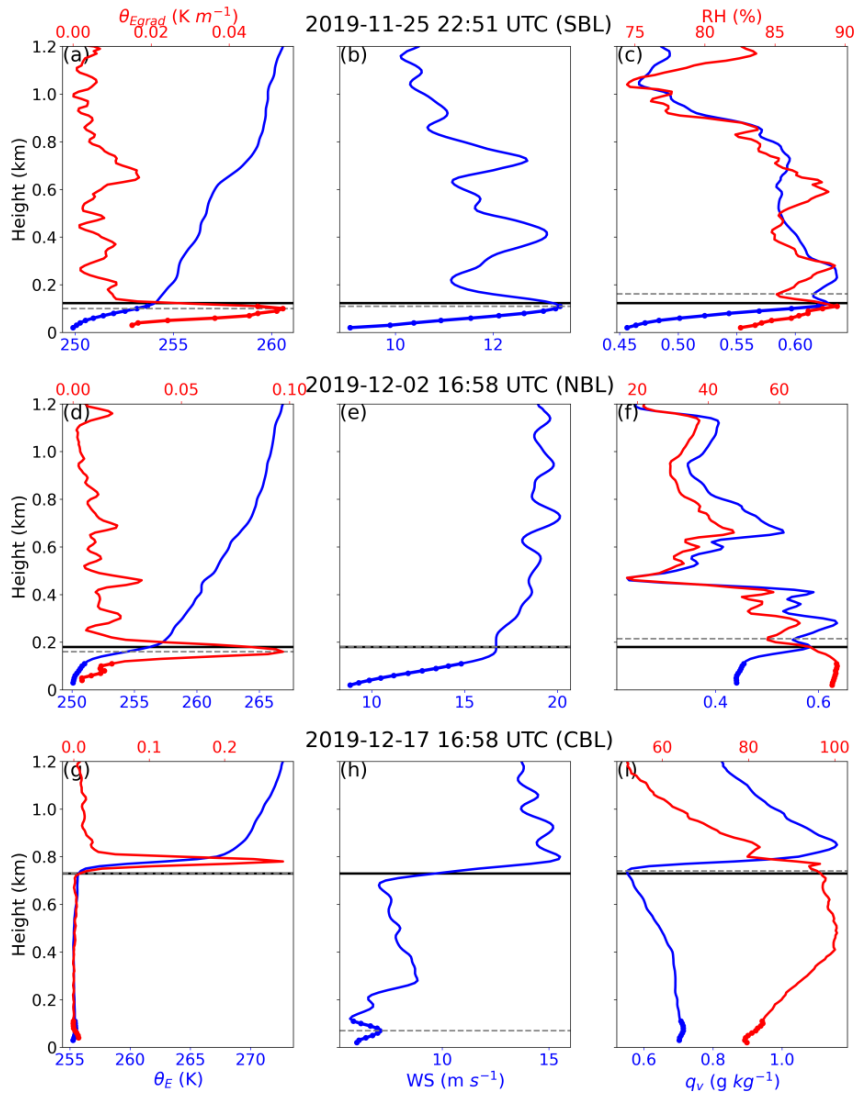


Figure 2 Vertical profiles of (left) equivalent potential temperature (θ_E), θ_E gradients (θ_{Egrad}), (middle) wind speed (WS), and (right) relative humidity (RH) and specific humidity (q_v) at (a–c) 25 November 2019, 22:51 UTC, (d–f) 2 December 2019, 16:58 UTC, and (g–i) 17 December 2019 16:58 UTC. Boundary layers at the three times represent stable boundary layer (SBL), near-neutral boundary layer (NBL), and convective boundary layer (CBL), respectively. The gray dashed horizontal lines denote the atmospheric boundary-layer height (ABLH) estimates based on multiple profiles, and the black solid horizontal lines denote the manually observed ABLHs. The dots at the lowest 100 m altitude denote the merged profiles.

Section 2.3: The authors should explain in more detail what is the size of the footprint of these fluxes, and to what extent they are expected to relate to the ABLH.

Response: Thanks for pointing this out. The footprint analysis by using the Kljun and others (2015) model indicates that 90% of the measured flux is expected to come from

within 275 m of the eddy-covariance observation site in average, and this fetch is characterized by sea-ice surface. The sounding site is located 300–600 m away from the meteorological tower. Although the sounding launch site is not often within the source region of the flux measurements, we assume that the spatial variation of turbulent fluxes within a kilometer range over the local sea-ice surface can be ignored.

The following description has been added into our revised manuscript:

The sentence “We neglect the distance between the vessel and ‘Met City’ and consider that their ABL conditions are the same, particularly when considered on hourly timescales” has been revised as: “Based on a footprint analysis using the Kljun et al. (2015) model, 90% of the sensible heat flux measurements have a source area fetch of no more than 275 m, a region that was typically strongly dominated by consistent sea ice throughout the year. Although the sounding site may typically be outside the source region of these flux measurements, we assume the conditions at the two sites are predominantly equivalent, which is also assumed in the merged sounding-tower product.”

Ln 169: please add more justification why 2 classes of ABLH types are sufficient. The SBL part was earlier subdivided by many studies by Zilitinkevich in the truly neutral PBL, the nocturnal SBL and the long-lived PBL. These concepts may help to further explain the observations.

Response: As mentioned above, we have added the NBL type into our analysis, and additionally separated the ABLs into cloudy and clear conditions. Additionally, we now analyze the subdivisions of TN, CN, NS, and LS types in the correlation analysis section. The corresponding changes are given in our revised manuscript as follows:

(1) Changes related to regime classification

3.1 ABL regime classification and ABLH determination

The ABLH determination method starts with the classification of ABL regimes. Based on previous studies (e.g., Vogelesang and Holtslag, 1996; Liang and Liu, 2010), we divide the ABLs into three types: stable boundary layer (SBL), near-neutral boundary layer (NBL), and convective boundary layer (CBL), corresponding with three different stability states near the surface. We first use SH to diagnose the ABL regime types. The specific classification formula is presented below:

$$\begin{cases} SH > +\delta & \text{for CBL} \\ SH < -\delta & \text{for SBL, (1)} \\ else & \text{for NBL} \end{cases}$$

where δ is the critical value that is specified as 2 W m^{-2} , following Steeneveld et al. (2007b). If corresponding SH data are unavailable, the difference of equivalent potential temperature (θ_E) between the 100 and 50 m heights (θ_E difference) derived from

sounding profiles is used to determine the ABL types. Specifically, if θ_E difference is larger than 0.2 K, the ABL is identified as SBL; if θ_E difference is less than -0.2 K, the ABL is identified as CBL; and other profiles are labeled as NBLs, roughly following Liu and Liang (2010).

(2) Changes related to detect cloud condition

We also analyze the algorithm performances for cloudy and clear conditions, considering that low-level clouds containing liquid water play an important role in the Arctic ABL (Shupe and Intrieri, 2004; Brooks et al., 2017). In our study, the RH threshold of 96% (Silber and Shupe, 2022) and the cloud source flag data are used for cloud detection. If a cloud is detected in the cloud source flag data and the RH is larger than 96%, then the profile is labeled as cloudy. The sounding profiles that contain at least one identified cloud layer below 1500 m are classified as “cloudy”, and as “clear” otherwise.

(3) Changes related to correlation analysis

The free-flow stability (characterized by the free-flow Brunt-Väisälä frequency, N) can affect the ABLH (Zilitinkevich et al., 2002; Zilitinkevich and Baklanov, 2002; Zilitinkevich and Esau, 2002, 2003), and therefore is also examined here. Based on the buoyancy flux at the surface (B_s) and N , the NBLs and SBLs can be further divided into four types: the truly neutral (TN, $B_s = 0$ and $N = 0$), the conventionally neutral (CN, $B_s = 0$ and $N > 0$), the nocturnal stable (NS, $B_s < 0$ and $N = 0$), and the long-lived stable boundary layer (LS, $B_s < 0$ and $N > 0$). According to Zilitinkevich and Baklanov (2002), we calculate the N and B_s and reclassify the SBLs and NBLs. We find that the percentages of $N > 0.015$ in SBLs and NBLs are 89 % and 80 %, which indicates that LS and CN types dominate the stable and neutral conditions for MOSAiC, respectively. Since only 80 TN cases were identified, these are deemed to be too few for additional analysis of this type. Zilitinkevich and Esau (2003) gave ABLH equations relevant to each ABL type as:

$$h_E = \begin{cases} C_N u_* |fN|^{-1/2} & \text{(Pollard et al., 1973) for CN ABL, (10)} \\ C_S u_*^2 |fB_s|^{-1/2} & \text{(Zilitinkevich, 1972) for NS and LS ABL, (11)} \end{cases}$$

where h_E is the equilibrium ABLH, f is the Coriolis parameter, and C_N and C_S are empirical coefficients. In addition, Vogelesang and Holtslag (1996) and Steeneveld et al. (2007a) also explore a h_E equation without taking into account f explicitly, expressed as:

$$h_E = C_i \frac{u_*}{N} \text{ for all SBL and NBL, (12)}$$

where C_i is an empirical coefficient. Here we select the CN, NS, and LS ABLH dataset, and fit the data with the corresponding expressions in Eq. (10–12) to obtain the empirical coefficients, and the results are presented in Fig. 12. All three expressions tend to well represent the ABLHs, with significant correlation coefficients. The empirical coefficients C_N and C_S are 1.7 and 0.4, respectively, which are close to the typical values determined through large-eddy simulations (Zilitinkevich, 2012). The

coefficient $C_i = 20$ in Fig. 12c is double the typical value of 10 (Vogelezang and Holtslag, 1996), but agrees with the results reported by Overland and Davidson (1992) for the ABL over sea ice. The difference in C_i may be attributed to the unique free-flow stability or other potential mechanisms of ABL development in the Arctic atmosphere.

In summary, near-surface conditions and free-flow stability play a key role in ABL development and are also an indicator, in that one can roughly determine the development state of the whole ABL from these basic variables.

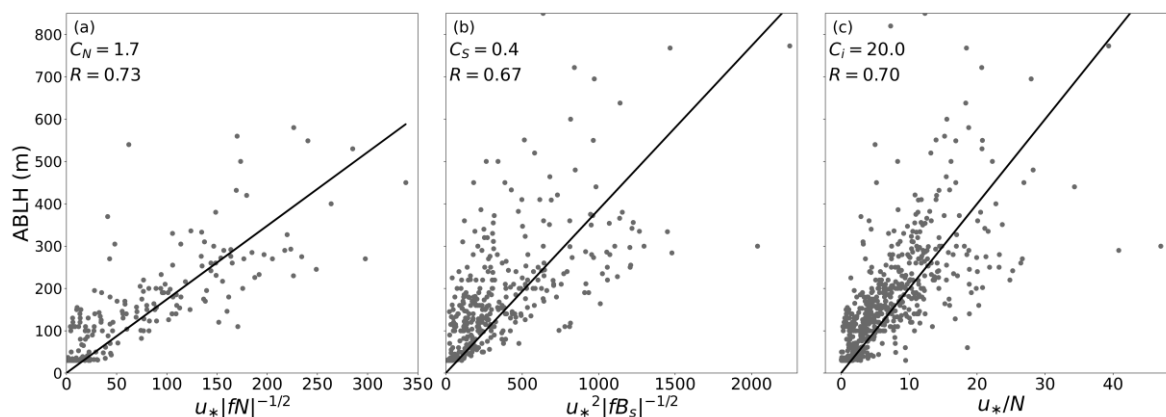


Figure 12 The ABLHs versus three expressions in Eq. (10–12). The empirical coefficients C_N , C_S , and C_i are given in (a), (b), and (c), respectively, and represent the slope of the best fit line (black line). The correlation coefficient R is given in each panel, which is statistically significant ($p < 0.05$).

Ln 178: theta is used here as measure for stratification. However, above you mention that the PBL driven by turbulence in cloud is an ABLH important archetype. Is it not more appropriate to use a temperature metric that is conserved in moist conditions like the liquid water potential temperature? Please show that this choice does not affect your conclusions!

Response: Thank you very much for your helpful suggestion. We have replaced θ with the equivalent potential temperature θ_E . As mentioned above, the types of ABLs are first diagnosed by SH data. Only when the SH is missing are the sounding profiles classified by the difference of θ_E between the 100 and 50 m heights. We also checked the impact of using θ_E on regime classification. For the θ criterion, we obtain 452 SBLs, 240 NBLs, 272 CBLs. For the θ_E criterion, we obtain 442 SBLs, 249 NBLs, 273 CBLs. This slight difference does not affect our conclusions.

Ln 180: Δ_s is chosen to be 0.2 K. Please relate link this to the measurement

accuracy of the sounding. In my view even for a routine AWS the measurement uncertainty is about 0.3K when it includes also representativeness uncertainty.

Response: Thank you very much for your helpful comment. According to the description in Liu and Liang (2010), δ_s is the θ increment for the minimum strength of the stable (inversion) layer. The value of δ_s would be set to zero for idealized cases but in practice is specified as a small positive value, and this value depends on the surface characteristics as well as inherent uncertainties or noise in the measurements. For profiles over ocean and ice, this threshold has been empirically defined to be 0.2 K. While, as described in the measurement data report, the measurement uncertainty is exactly 0.3 K. Considering the uncertainty of this criterion, we have replaced it with the SH data to determine the ABL regime types, which follows Steeneveld et al. (2007b). Only when the SH is missing is the θ_E criterion and threshold proposed by Liu and Liang (2010) used for determining ABL type. For this comment, the corresponding change is given in our revised manuscript as follows:

3.1 ABL regime classification and ABLH determination

The ABLH determination method starts with the classification of ABL regimes. Based on previous studies (e.g., Vogelezang and Holtslag, 1996; Liang and Liu, 2010), we divide the ABLs into three types: stable boundary layer (SBL), near-neutral boundary layer (NBL), and convective boundary layer (CBL), corresponding with three different stability states near the surface. We first use SH to diagnose the ABL regime types. The specific classification formula is presented below:

$$\begin{cases} SH > +\delta & \text{for CBL} \\ SH < -\delta & \text{for SBL, (1)} \\ else & \text{for NBL} \end{cases}$$

where δ is the critical value that is specified as 2 W m^{-2} , following Steeneveld et al. (2007b). If corresponding SH data are unavailable, the difference of equivalent potential temperature between the 100 and 50 m heights (θ_{Egrad}) derived from sounding profiles is used to determine the ABL types. Specifically, if the $\theta_{Egrad} > 0.2 \text{ K}$, the ABL is identified as SBL; if the $\theta_{Egrad} < -0.2 \text{ K}$, the ABL is identified as CBL; and other profiles are labeled as NBLs, roughly following Liu and Liang (2010).

Ln 226: which an air parcel rising adiabatically from the surface becomes neutrally buoyant... Has an temperature excess been added to the surface parcel and if so with which value?

Response: Thank you very much for your helpful comment. According to Liu and Liang (2010) and the PBLH VAP data report, a temperature excess of 0.1 K has been added. Actually, the ABLH estimates based on the Liu-Liang algorithm is provided by the data product, and we directly use them for comparison. We have added more relevant information into our revised manuscript as follows:

The Liu-Liang algorithm determines ABLH based on potential temperature and wind speed. For CBL regimes, the definition of ABLH is the height at “which an air parcel rising adiabatically from the surface becomes neutrally buoyant”, and the temperature excess value is 0.1 K.

Ln 227+228: two different estimates of the SBL height are obtained based on stability criteria and wind shear criteria, respectively. Please elaborate in more detail how it has been done, in this way we cannot evaluate the procedure is appropriate.

Response: Thank you very much for your helpful comment. According to Liu and Liang (2010) and the ABLH VAP data report, the stability criteria are to find the lowest level, k , at which the θ_{Egrad} reaches a minimum and meets either of the following two conditions:

$$\begin{cases} \theta_{Egrad\ k} - \theta_{Egrad\ k-1} < -40\ \text{K/km} \\ \theta_{Egrad\ k+1} < 0.5\ \text{K/km}, \theta_{Egrad\ k+2} < 0.5\ \text{K/km} \end{cases},$$

where the subscripts (k , $k-1$, $k+1$, and $k+2$) represent the θ_{Egrad} at corresponding levels.

For wind shear, the ABLH is defined as the height where the wind speed reaches a maximum that is at least 2 m/s stronger than the layers immediately above and below while decreasing monotonically toward the surface (i.e., a low-level jet). The final ABLH is defined as the lower of the two heights. We have added more information into our revised manuscript as following:

For SBL regimes, two different estimates of the ABLH are obtained, if possible, based on stability criteria and wind shear criteria, respectively. For stability, the ABLH is defined as the lowest level, k , at which the θ_{Egrad} reaches a minimum and meets either of the following two conditions:

$$\begin{cases} \theta_{Egrad\ k} - \theta_{Egrad\ k-1} < -40\ \text{K/km} \\ \theta_{Egrad\ k+1} < 0.5\ \text{K/km}, \theta_{Egrad\ k+2} < 0.5\ \text{K/km} \end{cases}, \quad (2)$$

where the subscripts (k , $k-1$, $k+1$, and $k+2$) represent the θ_{Egrad} at corresponding levels.

For wind shear, the ABLH is defined as the height where the wind speed reaches a maximum that is at least 2 m/s stronger than the layers immediately above and below while decreasing monotonically toward the surface (i.e., a low-level jet). The final ABLH is defined as the lower of the two heights.

Ln 239: dimensional number. It is a dimensionLESS number, of course!

Response: Thank you very much for pointing this out. It is a typo error and we have corrected it.

Ln 240-246: the paper ignores here the knowledge that was developed in Vogelezang and Holtslag, which was by the way cited, that a better score for the ABLH can be obtained if Equation 2 is not considered from the surface parcel, but a parcel at somewhat above the surface. Hence I feel the latest knowledge is not taken into account here.

Response: Thank you very much for your helpful comment. As mentioned above, we have considered the Ri formula proposed by Vogelezang and Holtslag (1996) into our revised manuscript, and use it in the ABL algorithm for clear-sky conditions. The corresponding changes are given in our revised manuscript as follows:

3.3 An improved Ri algorithm considering the cloud effect

As a traditional Ri_b formula, Eq. (3) may break down in cases of ABLs with relatively high wind speed and upper-level stratification due to the overestimation of shear production (Kim and Mahrt, 1992). Vogelezang and Holtslag (1996) proposed the finite-difference Ri formula, which is expressed as:

$$Ri_F = \frac{(g/\theta_{vs})(\theta_{vh} - \theta_{vs})(h - z_s)}{(u_h - u_s)^2 + (v_h - v_s)^2 + bu_*^2}, \quad (6)$$

where z_s is the lower boundary for the ABL, θ_{vs} , u_s , and v_s are the θ_v and wind components at the height z_s , respectively, b is an empirical coefficient, and u_* is the surface friction velocity. Ri_F is considered for a parcel located somewhat above the surface to avoid the above problem, and u_* is also taken into account to avoid underestimation in the situation of a uniform wind profile in the upper layer. Here, we use Ri_F for clear-sky profiles and take z_s and b values as 40 m and 100, respectively, according to Zhang et al. (2020).

Ln 254: *Bias* is the absolute bias; *SEE* is the standard error. I object against the term bias here. Bias can be either positive or negative, but your formula for bias cannot, so you use the MAE, mean absolute error. Idem for SEE, it is the standard deviation of the error, not the standard deviation of the ABLH.

Ln 257: note that Steeneveld et al. (2007) used the median of the absolute error is evaluation metric in a similar type of study. This is helpful to avoid that the error statistics are determined strongly by one or two outliers. Please consider this as well. Figure 3: it is unclear whether the error statistics in the left upper corner relate to the CBL or SBL data. It would be interesting to have the statistics for both classes, to underline the score for SBL is much poorer.

Response: Thank you very much for pointing this out. We have revised the formula of *Bias* that can be positive or negative, and replaced other statistical measures with median of the absolute error (*MEAE*). For Fig. 3, we have calculated the error statistics for SBL, NBL, CBL, and cloudy conditions, and listed them in Table 1. According to the error statistics, the Liu-Liang algorithm and the Heffter algorithm perform poorly in determining SBL height, especially the Liu-Liang algorithm. For this comment, the corresponding changes are given in our revised manuscript as follows:

(1) Changes related to the description of statistical measures

To quantitatively evaluate the performance of each automatic algorithm, we introduce the correlation coefficient R and two other statistical measures: the *Bias* and the median absolute error (*MEAE*; Steeneveld et al., 2007a). The formulas are as follows:

$$Bias = \frac{2}{n} \sum_{i=1}^n \frac{H_{auto} - H_{obs}}{H_{auto} + H_{obs}}, \quad (4)$$

$$MEAE = \text{median}(|H_{auto} - H_{obs}|), \quad (5)$$

where H_{auto} is the ABLH obtained by the automated algorithm; H_{obs} is the ABLH manually determined; n is the number of valid sounding profile samples. According to the definitions of these statistical measures, larger R and smaller *Bias* and *MEAE* mean a better performance of the automated algorithm.

(2) Changes related to Figure 3

Figure 3 presents the comparisons of estimated ABLHs with the manually-labeled ABLHs, and the associated statistical measures are given in Table 1. The results show that the Ri_b algorithm with Ri_{bc} of 0.25 performs best overall, and particularly for SBL cases. The performance of the Ri_b algorithm with Ri_{bc} of 0.5 is poorer than that of the Ri_b algorithm with Ri_{bc} of 0.25, with overestimations of ABLHs in general, and larger errors with lower correlation coefficients for all types of ABLs. The Heffter algorithm performs well in cases of high ABLH and particularly for cloudy and CBL cases, but does significantly overestimate ABLH in a large number of cases as shown in the Fig. 3c subgraph. This is attributed to the determination criterion of the Heffter algorithm, i.e., ABLHs are determined by inversion layers, which means that large errors occur when the inversion layer is higher than the mixed layer. Additionally, while the Heffter performance in many of the ABL conditions is only marginally worse statistically than the Ri_b algorithm with Ri_{bc} of 0.25, its correlations are notably worse for SBL and NBL cases. The performance of the Liu-Liang algorithm is generally poorer than the other algorithms, particularly for correlation coefficient, which is probably due to the impact of noise in the lower ABLH profiles and unsuitable parameters in the algorithm. In summary, the Ri_b algorithm is reliable over the Arctic Ocean and performs better than other algorithms, and this result agrees with Jozef et al. (2022). Furthermore, we will explore ways to improve the Ri_b algorithm to make it more suitable for cloudy and convective conditions.

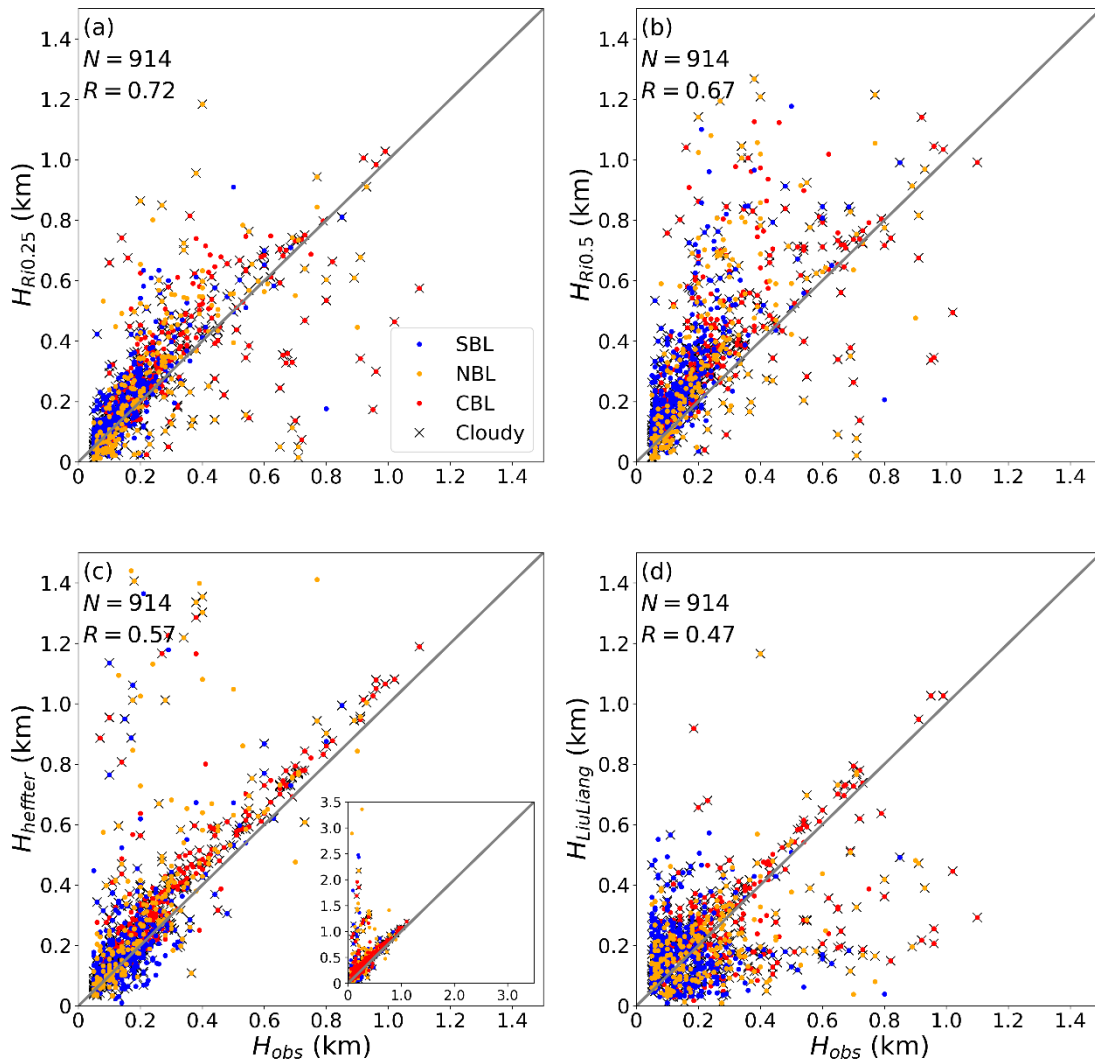


Figure 3 Comparisons of the ABLHs determined from radiosonde profiles using the bulk Richardson number (Ri_b) algorithm with the critical values (Ri_{bc}) of (a) 0.25 and (b) 0.5, (c) the Heffter algorithm, and (d) the Liu-Liang algorithm with the manually-identified “observed” ABLHs. The blue, yellow, and red colors indicate regime types of SBL, NBL, and CBL, respectively. The “x” signs indicate the Cloudy ABLHs. The case numbers (N) and correlation coefficients (R) are given in each panel. The subgraph in (c) denotes all data points ranging from 0 to 3.5 km.

Figure 3c and d: I do not understand why the H_{obs} is different for the SBL and the CBL for the two panels. Please explain, the filtering was done on the observation, wasn't it? Not on the selected algorithm. Also add the number of samples in the block with error statistics.

Response: Thank you very much for pointing this out. We realize that the Fig. 3 is unclear. Actually, the H_{obs} is the same in all panels and the filtering was in fact done on the observations. The data range of Fig. 3c in the original manuscript is from 0–3.5 km due to the severe ABLH overestimation by the Heffter algorithm, and the axis range

is thus different from that of other panels. Therefore, we unify the axis ranges of all panels to avoid misunderstandings, and add a subgraph in Fig. 3c to denote all data points. For this comment, the corresponding changes are given in our revised manuscript as follows:

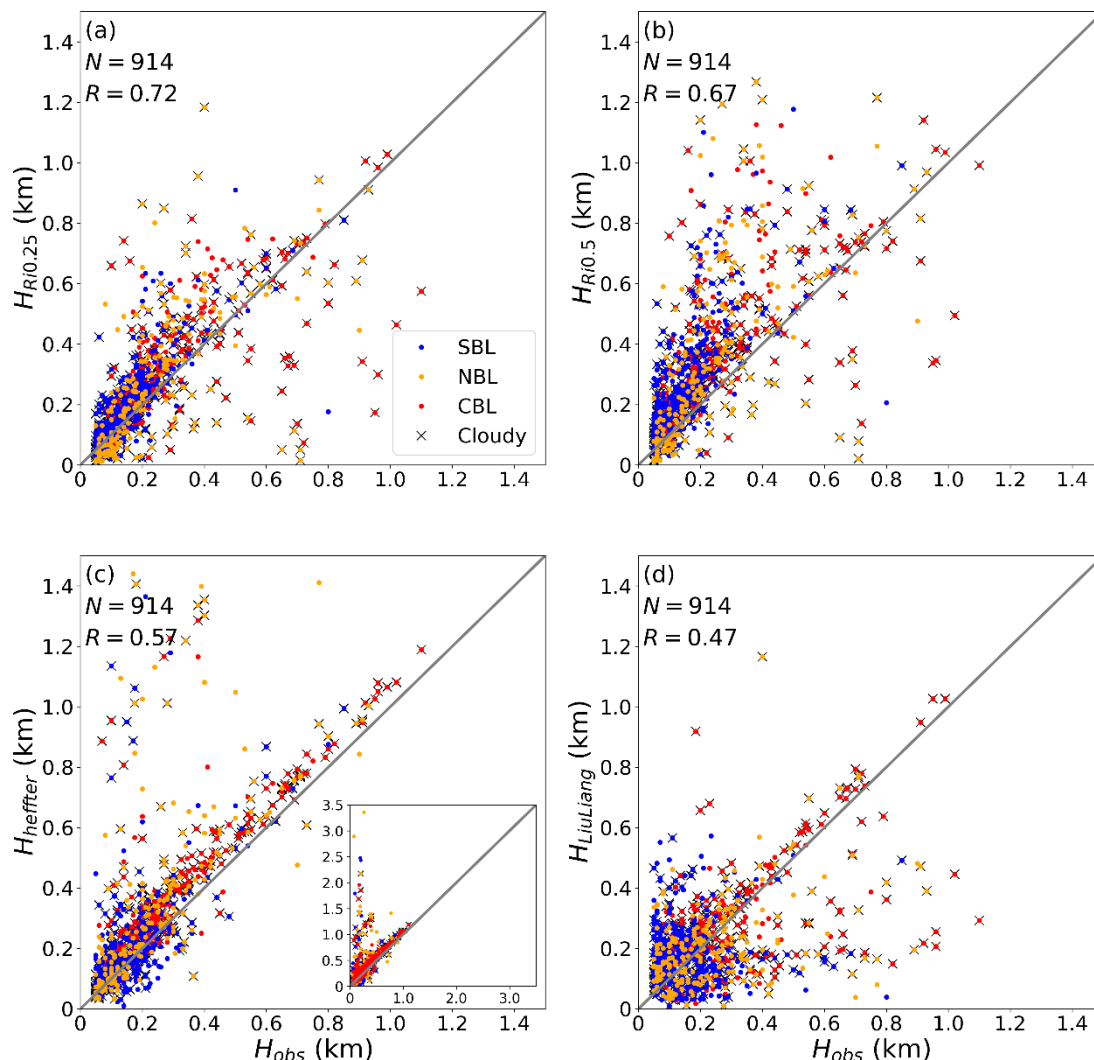


Figure 3 Comparisons of the ABLHs determined from radiosonde profiles using the bulk Richardson number (Ri_b) algorithm with the critical values (Ri_{bc}) of (a) 0.25 and (b) 0.5, (c) the Heffter algorithm, and (d) the Liu-Liang algorithm with the manually-identified "observed" ABLHs. The blue, yellow, and red colors indicate regime types of SBL, NBL, and CBL, respectively. The "x" signs indicate the Cloudy ABLs. The case numbers (N) and correlation coefficients (R) are given in each panel. The subgraph in (c) denotes all data points ranging from 0 to 3.5 km.

Ln 296: Note again that VH96 do use a different definition of Ri.

Response: Thank you very much for pointing this out. We have considered it into our revised manuscript as mentioned above.

Ln 302: This result is distinct from that of Jozef et al. (2022). Add how it is distinct....?

Ln 303: might be that ... different... ->Better to figure that out!!! It is related to the key of this paper.

Response: Thank you very much for your helpful comment. We have added the analysis of different Ri formulas and Ri_c values, and the Ri formula used in Jozef et al. (2022) is also included. Jozef et al. (2022) calculates the Ri over a rolling 30 m altitude range, and uses the Ri_c value of 0.75. The method of calculating Ri over a rolling 30 m range causes dramatic variation within the ABL, as seen in Fig. 5. Thus, for this Ri definition, a large Ri_c value is required to avoid the noise. The corresponding changes are given in our revised manuscript as follows:

Since some other studies have proposed different Ri_c values for MOSAiC (e.g., Jozef et al., 2022; Barten et al., 2023; Akansu et al., 2023), we will discuss the difference in Ri_c values here. The first thing to make clear is that these studies use different formulas to obtain Ri profiles. Barten et al. (2023) and Akansu et al. (2023) both use the traditional Ri_b algorithm based on Eq. (3), while they used Ri_c values of 0.4 and 0.12, respectively. This difference was likely caused by the different methods to manually derive their reference ABLH data sets. Jozef et al. (2022) calculates the Ri over a rolling 30 m altitude range, labeled as Ri_r , and the criterion is modified to require four consecutive data points to be above the Ri_c of 0.75. In our study, we use Ri_F proposed by Vogelezang and Holtslag (1996) for clear-sky conditions, and Ri_m for cloudy conditions. Based on the results presented here, it is apparent that this more complex approach improves the error statistics relative to approaches based on Eq. (3), regardless of Ri_c . In addition, some of the differences may also related to authors using different data sets or time periods. For instance, Akansu et al. (2023) primarily used sounding data based on tether balloon for a specific sub-period of MOSAiC, and Jozef et al. (2022) used radiosondes from when they had concurrent UAV observations. The data used in our study are based on merged sounding-tower product, as mentioned above.

To further explore the differences among the four different approaches, we examine one SBL and CBL case. For a clear-sky SBL case (Fig. 5 a, b), the approaches from Akansu et al., Jozef et al. (2022), and this study all agree closely with the manual ABLH, while the Barten et al. approach results in a significant overestimation. For a cloudy-sky CBL case (Fig. 5 c d), the approach from this study agrees with the manual ABLH, while the approach from Barten et al. overestimates the ABLH by about 30 m, and the approaches from Akansu et al. and Jozef et al. (2022) underestimate the ABLH by 130 m and 230 m, respectively. These results further demonstrate how Ri_c depends on the choice of Ri formula. Moreover, Ri_c is not analytically derived from basic physical principles (Zilitinkevich et al. 2007), and the concept of Ri_c is challenged by non-steady regimes (Zilitinkevich and Baklanov, 2002) and the hysteresis phenomenon

(Banta et al., 2003; Tjernström et al., 2009). Therefore, an objective Ri_c does not exist. Rather, it is empirically used as an algorithmic parameter to simply derive the ABLH.

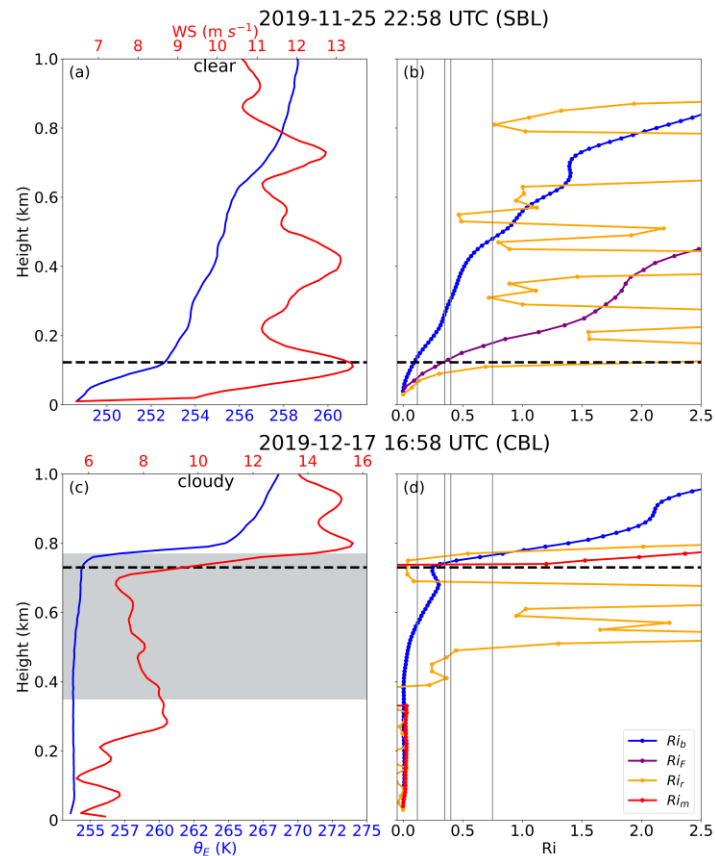


Figure 5 Vertical profiles of (left) θ_E and wind speed, and (right) Ri based on different formulas at (a–b) 25 November 2019, 22:58 UTC and (c–d) 17 December 2019, 16:58 UTC. Boundary layers at the two times represent a clear-sky SBL and a cloudy-sky CBL respectively. The black dashed horizontal lines denote the manually-identified ABLH, and the gray solid vertical lines denote the different Ri_c values, including 0.12, 0.35, 0.4, and 0.75. The gray shading in (c) denotes the cloud layer.

Ln 328: from 13 April through to 24 May 2020. In this period, the convectively thermal structure contributes to ABLH reaching over 610 m for about 6 days, with the maximum ABLH of 1152 m: This is the period with a warm intrusion from the south, so the PBL height is likely strongly governed by the advection of warm air, its turbulent kinetic energy, and its stratification. Equation 2 was not developed for such conditions, so it is fair to evaluate it as such?

Response: Thank you very much for your helpful comment. We checked the ABL cases in this period, and the comparisons of ABLH respectively estimated by the traditional and improved Ri algorithms with H_{obs} are presented in Fig. R1. There are 108 ABL

cases in this period. We find that the Ri algorithm based on Eq. (3) surely cannot determine ABLH well, while the improved Ri algorithm significantly corrects for errors in ABLH estimation, by using the Ri_F proposed by VM96 and taking the cloud effect into account.

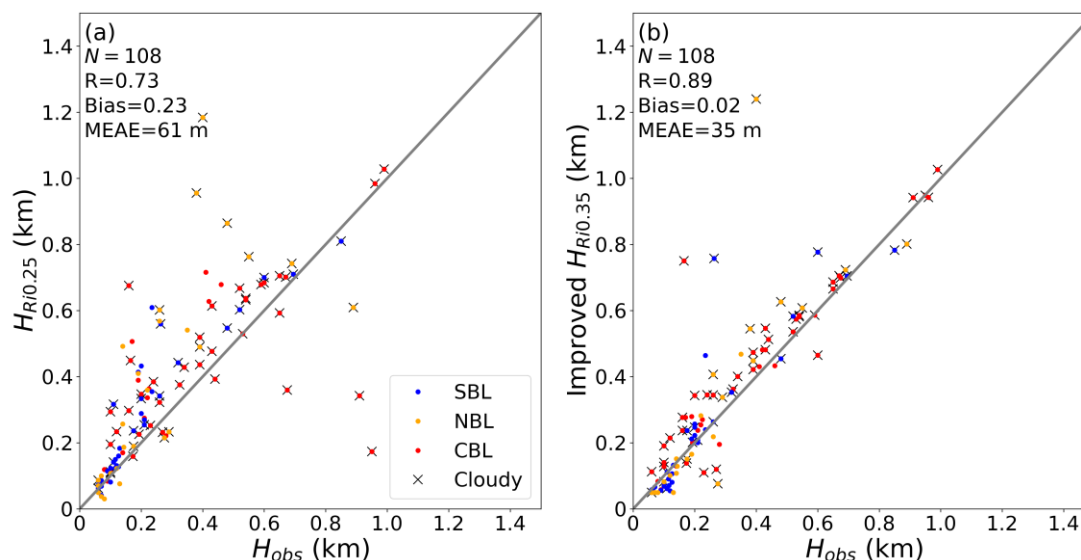


Figure R1 The Comparisons of the selected ABLHs determined by the (a) bulk Richardson number (Ri_b) algorithms with the critical values (Ri_{bc}) of 0.25 and (b) the improved Ri algorithm with observed ABLHs. The blue, yellow, and red dots indicate regime types of SBL, NBL, and CBL, respectively. The case number (N), correlation coefficient (R), $Bias$, and median of the absolute error ($MEAE$) are given in each panel.

Ln 366: I am little surprised that the θ_E appears here in the analysis, while it is not reasoned why we step over from θ to θ_E . I agree that θ_E analysis is valuable, but should θ_E not have been applied to Equation 2?

Response: Thank you very much for your helpful comment. Actually, the Eq. (2) and corresponding ABLH estimations are provided by the ABLH VAP data product, so we have used it directly for comparison analysis. In our improved algorithm, we use the moist Richardson number to take the cloud effect into account. Except for Eq. (2), we have replaced θ with θ_E in the rest of our manuscript for consistency.

Figure 7: Add in the legend whether these are the monthly averages of the soundings from 5:00, or 11:00, or 17:00, or 23:00, or all mixed together. It is better to stick to one time slot to avoid that the effects of the diurnal cycle in the summer months are mixed away.

Response: Thank you very much for your helpful comment. In our study, all soundings

from 5:00, or 11:00, or 17:00, or 23:00 are used to calculate the monthly averages. In order to check the impact of the diurnal cycle, we also calculate respective monthly profiles based on 5:00, 11:00, 17:00 and 23:00. The results are presented in Fig. (R2–R5). We find that the monthly profiles based on the four individual hours of a day do not differ much, which suggests that the diurnal cycle does not have a significant effect on the ABL thermal structure. Therefore, we continue to use the monthly profiles based on all soundings, and add the relevant statement at the end of the analysis as follows:

In addition, we examined the potential implications of the diurnal cycle on the ABL thermal structure. Monthly profiles based on different moments of a day were found to show little variability (not shown), such that the impact of the diurnal cycle is minimal.

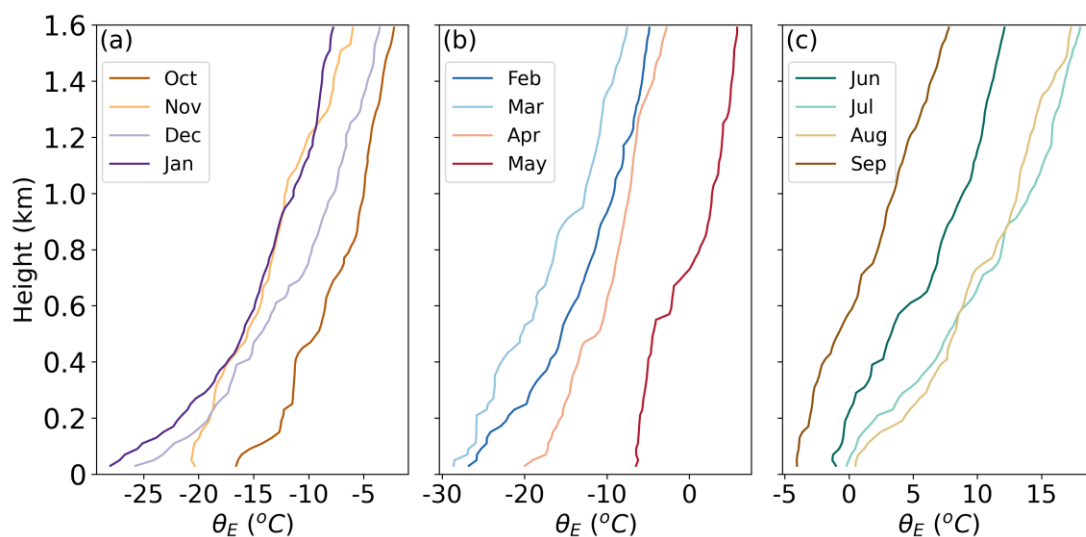


Figure R2 Median profiles of equivalent potential temperature throughout the MOSAiC year are divided into (a), (b), and (c), based on sounding data from 5:00.

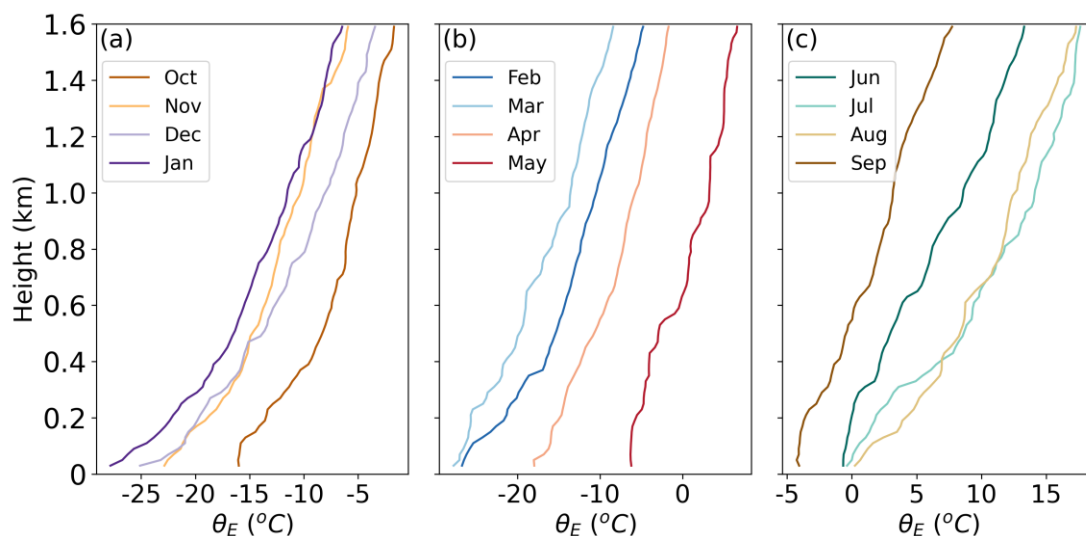


Figure R3 Median profiles of equivalent potential temperature throughout the MOSAiC

year are divided into (a), (b), and (c), based on sounding data from 11:00.

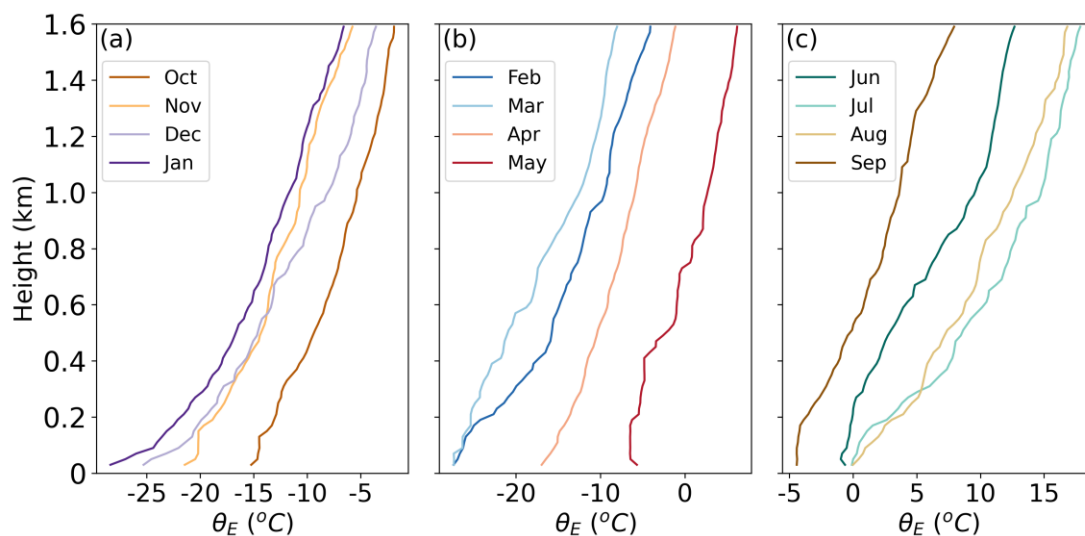


Figure R4 Median profiles of equivalent potential temperature throughout the MOSAiC year are divided into (a), (b), and (c), based on sounding data from 17:00.

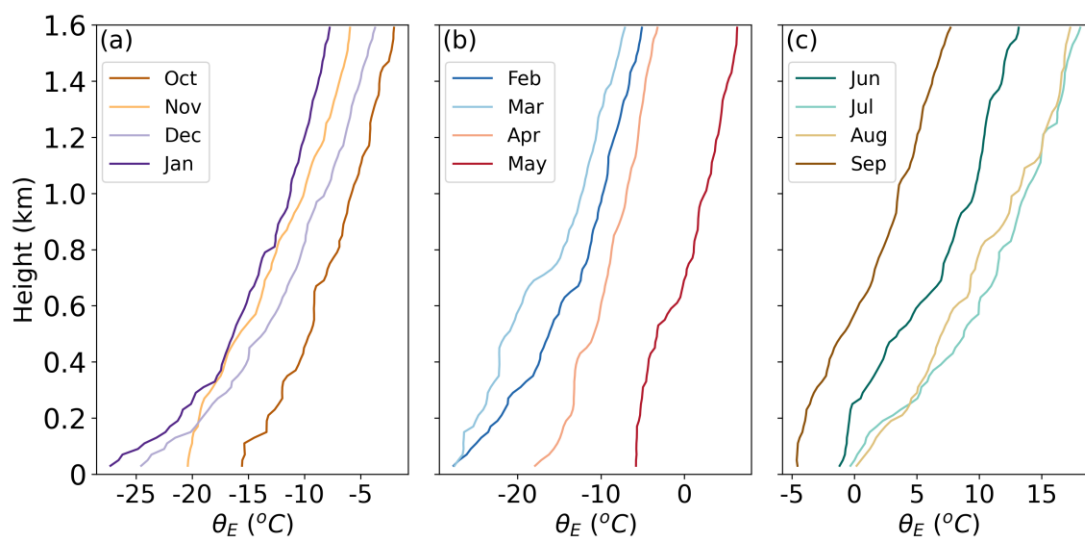


Figure R5 Median profiles of equivalent potential temperature throughout the MOSAiC year are divided into (a), (b), and (c), based on sounding data from 23:00.

Ln 395: temperature gradient. Better to use (equivalent) potential temperature gradient to remain consistent with the above.

Response: Thank you very much for your helpful suggestion. We have replaced the temperature gradient with equivalent potential temperature gradient θ_{Egrad} for consistency. The corresponding changes are given in our revised manuscript as follows:

To further explore the relations between surface conditions and the ABLH, we evaluate the correlations between the ABLH and three near-surface meteorological and

turbulence parameters during the MOSAiC period, including the near-surface equivalent potential temperature gradient ($\theta_{Egrad} = \theta_{E 10m} - \theta_{E 2m}$), friction velocity (u_*), and TKE dissipation rate (ϵ). The results are shown in Fig. 11. Generally, the near-surface buoyancy and shear effects both modulate these variables. In Fig. 11a, the ABLH distribution for negative θ_{Egrad} has a wide range from the lowest level to above 1 km. As θ_{Egrad} becomes positive and increases, the ABLH distribution rapidly narrows to below 200 m. In general, positive θ_{Egrad} means a stably stratified ABL and surface-based temperature inversion, both of which lead to low ABLH, and negative θ_{Egrad} means that atmospheric stability near the surface is near-neutral or convective, which is necessary for ABL development.

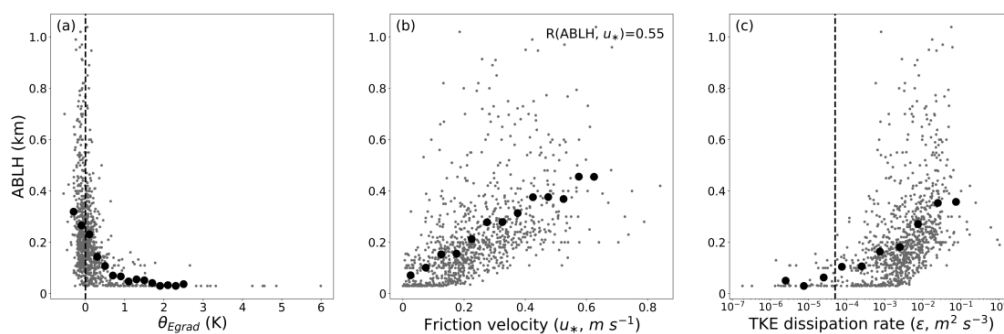


Figure 11 The ABLHs and bin-averaged values for (a) equivalent potential temperature gradient, θ_{Egrad} (K), (b) friction velocity, u_* ($m s^{-1}$), and (c) turbulent kinetic energy dissipation rate, ϵ ($m^2 s^{-3}$). The average bins for θ_{Egrad} , u_* , and ϵ logarithm are 0.2 K, $0.05 m s^{-1}$, and $0.5 m^2 s^{-3}$, respectively. The correlation coefficient R is given in (b), which is statistically significant ($p < 0.05$). The dashed vertical lines indicate the thresholds of (a) $\theta_{Egrad} = 0 K$ and (c) $\epsilon = 5 \times 10^{-5} m^2 s^{-3}$.

Ln 396: u_* , * should be subscripted (twice). And in the rest of the manuscript.

Response: Thank you very much for pointing this out. We have corrected relevant expressions.

Figure 8a and c: The R value in the plot is an estimate for the LINEAR correlation between the two variables, but obviously the relation is not linear. So better to remove it, or first do a transformation on the data such that the relation between them becomes linear.

Response: Thank you very much for pointing this out. We have removed the correlation analysis in Fig. 11a and c, retaining only in Fig. 11b. The corresponding change is given in our revised manuscript as following:

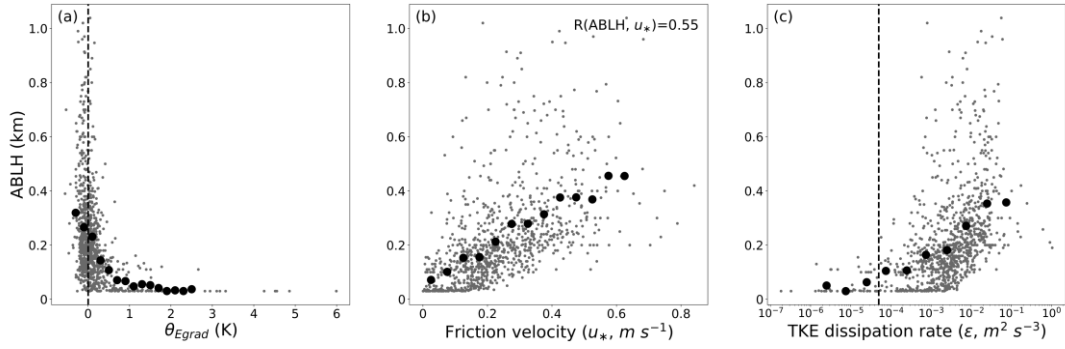


Figure 11 The ABLHs and bin-averaged values for (a) equivalent potential temperature gradient, θ_{Egrad} (K), (b) friction velocity, u_* ($m s^{-1}$), and (c) turbulent kinetic energy dissipation rate, ϵ ($m^2 s^{-3}$). The average bins for θ_{Egrad} , u_* , and ϵ logarithm are 0.2 K, $0.05 m s^{-1}$, and $0.5 m^2 s^{-3}$, respectively. The correlation coefficient R is given in (b), which is statistically significant ($p < 0.05$). The dashed vertical lines indicate the thresholds of (a) $\theta_{Egrad} = 0$ K and (c) $\epsilon = 5 \times 10^{-5} m^2 s^{-3}$.

Figure 8b: it is interesting to note that the ABLH is about $700xu_*$, which was also found/discussed in Vogelezang and Holtslag (1996) and Steeneveld et al. (2007). Both studies also explore $ABLH=10u_*/N$ as ABLH estimate, it would be interesting to be tested here as well.

Response: Thank you very much for your helpful suggestion. As mentioned above, we have added the tests for SBL formulas, including $h_E = \alpha \frac{u_*}{N}$. The results indicate that the best-fit α value is 20, not the typical value of 10, but this is in agreement with Overland and Davidson (1992), whose data also come from the ABL over sea ice. Thus, we attribute this result to the unique free-flow stability or other potential mechanisms of ABL development in the Arctic atmosphere. The corresponding changes are given in our revised manuscript as follows:

The free-flow stability (characterized by the free-flow Brunt-Väisälä frequency, N) can affect the ABLH (Zilitinkevich et al., 2002; Zilitinkevich and Baklanov, 2002; Zilitinkevich and Esau, 2002, 2003), and therefore is also examined here. Based on the buoyancy flux at the surface (B_s) and N , the NBLs and SBLs can be further divided into four types: the truly neutral (TN, $B_s = 0$ and $N = 0$), the conventionally neutral (CN, $B_s = 0$ and $N > 0$), the nocturnal stable (NS, $B_s < 0$ and $N = 0$), and the long-lived stable boundary layer (LS, $B_s < 0$ and $N > 0$). According to Zilitinkevich and Baklanov (2002), we calculate the N and B_s and reclassify the SBLs and NBLs. We find that the percentages of $N > 0.015$ in SBLs and NBLs are 89 % and 80 %, which indicates that LS and CN types dominate the stable and neutral conditions for MOSAiC, respectively. Since only 80 TN cases were identified, these are deemed to be too few

for additional analysis of this type. Zilitinkevich and Esau (2003) gave ABLH equations relevant to each ABL type as:

$$h_E = \begin{cases} C_N u_* |fN|^{-1/2} & \text{(Pollard et al., 1973) for CN ABL, (10)} \\ C_S u_*^2 |fB_s|^{-1/2} & \text{(Zilitinkevich, 1972) for NS and LS ABL, (11)} \end{cases}$$

where h_E is the equilibrium ABLH, f is the Coriolis parameter, and C_N and C_S are empirical coefficients. In addition, Vogelezang and Holtslag (1996) and Steeneveld et al. (2007a) also explore a h_E equation without taking into account f explicitly, expressed as:

$$h_E = C_i \frac{u_*}{N} \text{ for all SBL and NBL, (12)}$$

where C_i is an empirical coefficient. Here we select the CN, NS, and LS ABLH dataset, and fit the data with the corresponding expressions in Eq. (10–12) to obtain the empirical coefficients, and the results are presented in Fig. 12. All three expressions tend to well represent the ABLHs, with significant correlation coefficients. The empirical coefficients C_N and C_S are 1.7 and 0.4, respectively, which are close to the typical values determined through large-eddy simulations (Zilitinkevich, 2012). The coefficient $C_i = 20$ in Fig. 12c is double the typical value of 10 (Vogelezang and Holtslag, 1996), but agrees with the results reported by Overland and Davidson (1992) for the ABL over sea ice. The difference in C_i may be attributed to the unique free-flow stability or other potential mechanisms of ABL development in the Arctic atmosphere.

In summary, near-surface conditions and free-flow stability play a key role in ABL development and are also an indicator, in that one can roughly determine the development state of the whole ABL from these basic variables.

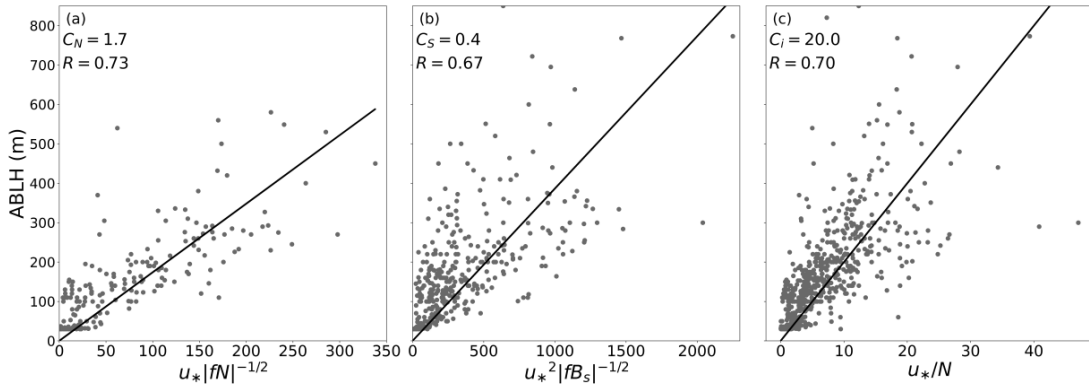


Figure 12 The ABLHs versus three expressions in Eq. (10–12). The empirical coefficients C_N , C_S , and C_i are given in (a), (b), and (c), respectively, and represent the slope of the best fit line (black line). The correlation coefficient R is given in each panel, which is statistically significant ($p < 0.05$).

Fig 9, caption: wind speed -> horizontal wind speed

Response: Revised as suggested.

Fig 10: Figure 10 Similar to Fig. 9, but the period is from 15 July 2020 to 30 August 2020. Legend is likely wrong since the x axis goes surely beyond September 1st.

Response: Thank you very much for pointing this out. We realized that the axis range of the Figure is beyond what we expect, and we have corrected it. The corresponding change is given in our revised manuscript as follows:

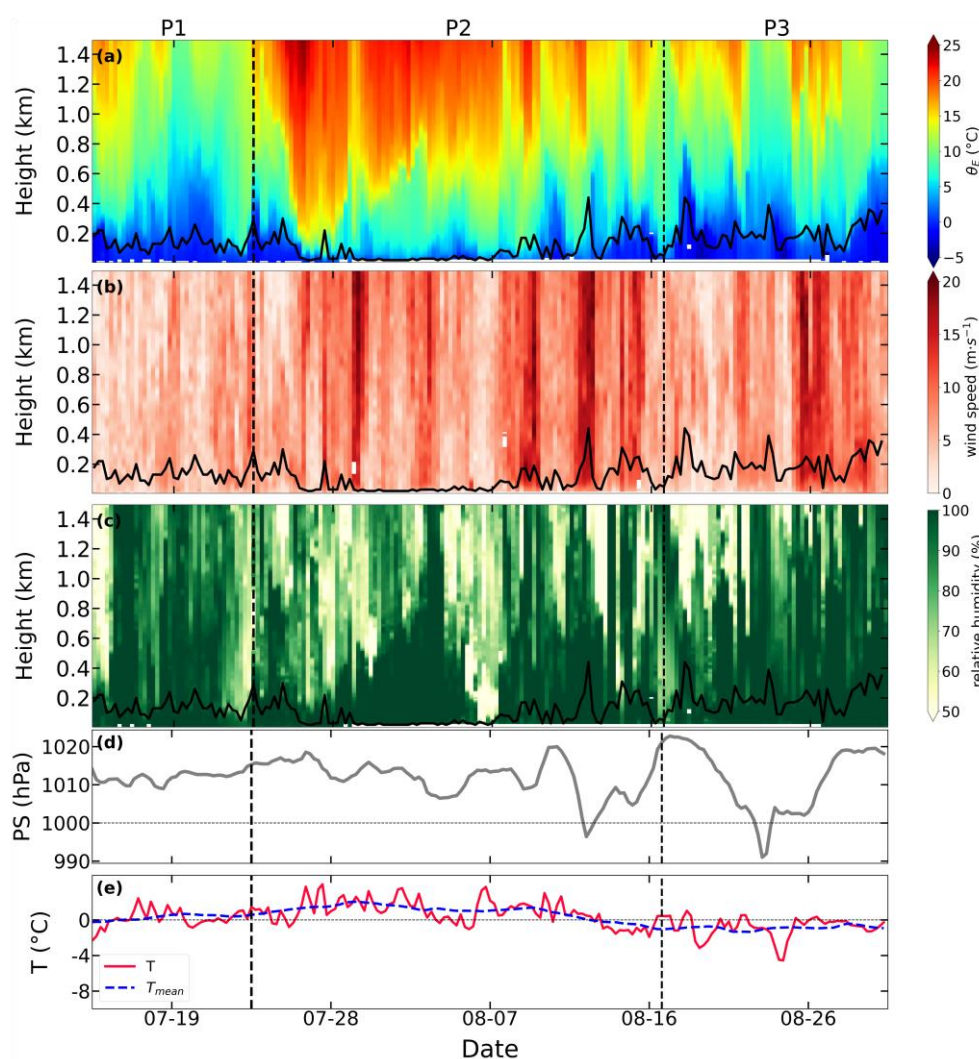


Figure 14 Similar to Fig. 13, but the period is from 15 July 2020 to 30 August 2020.