Historical (1960–2014) lightning and LNOₓ trends and their controlling factors in a chemistry–climate model

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Abstract. Lightning can cause natural disasters that end in human and animal injuries or fatalities, infrastructure destruction, and wildfire ignition. Lightning-produced NOₓ (LNOₓ), a major NOₓ (NOₓ=NO+NO₂) source, plays a vital role in atmospheric chemistry and global climate. The Earth has experienced marked global warming and changes in aerosol and aerosol precursor emissions (AeroPEs) since the 1960s. Investigating long-term historical (1960–2014) lightning and LNOₓ trends can provide important indicators for all lightning-related phenomena and for LNOₓ effects on atmospheric chemistry and global climate. Understanding how global warming and changes in AeroPEs influence historical lightning–LNOₓ trends is also helpful because it can provide a scientific basis for assessing future lightning–LNOₓ trends. Moreover, global lightning activities’ responses to large volcanic eruptions (such as the 1991 Pinatubo eruption) are not well elucidated, and are worth exploring. This study used the widely used cloud top height lightning scheme (CTH scheme) and the newly developed ice-based ECMWF-McCAUL lightning scheme to investigate historical (1960–2014) lightning–LNOₓ trends and variations and their controlling factors (global warming, increases in AeroPEs, and Pinatubo eruption) in the framework of the CHASER (MIROC) chemistry–climate model. Results of experiments indicate that both lightning schemes simulated almost flat global mean lightning flash rate trends during 1960–2014 in CHASER. Moreover, both lightning schemes suggest that past global warming enhances historical trends of global mean lightning density and global LNOₓ emissions in a positive direction (around 0.03% yr⁻¹ or 3% K⁻¹). However, past increases in AeroPEs exert an opposite effect to the lightning–LNOₓ trends (-0.07% yr⁻¹ – -0.04% yr⁻¹ for lightning and -0.08% yr⁻¹ – -0.03% yr⁻¹ for LNOₓ). When one considers only the aerosol radiative effects in the cumulus convection scheme, additionally, effects of past global warming and increases in AeroPEs on lightning trends were found to be heterogeneous across different regions when analyzing lightning trends on the global map. Lastly, this study is the first to suggest that global lightning activities were suppressed markedly during the first year after the Pinatubo eruption shown in both lightning schemes (global lightning activities decreased by as much as 47.02% simulated by the ECMWF-McCAUL scheme). Based on the simulated suppressed lightning activities after the Pinatubo eruption, our study findings also indicate that global LNOₓ emissions decreased after the 2–3-year Pinatubo eruption (2.44%–8.72% for the annual percentage reduction), which lasted 2–3 years). Model intercomparisons of lightning flash rate trends and variations between our study (CHASER) and other Coupled Model Intercomparison Project Phase 6 (CMIP6) models.
indicate significant uncertainties in historical (1960–2014) global lightning trend simulations. Such uncertainties must be investigated further.

1 Introduction

Lightning, an extremely energetic natural phenomenon, always occurs at every moment somewhere on Earth: its average occurrence frequency is approximately 46 times per second (Cecil et al., 2014). Lightning generation is associated with electric charge separation, which is mainly realized by collisions between graupel and hail and other types of hydrometeors of other types within convective clouds (Lopez, 2016). As a natural disaster, lightning can cause human and animal injuries and fatalities, infrastructure destruction, and wildfire ignition (Cerveny et al., 2017; Cooper and Holle, 2019; Jensen et al., 2022; Veraverbeke et al., 2022). Lightning-produced NOx (LNOx) accounts for around 10% of the global tropospheric NOx (NOx=NO+NO2) source. It is regarded as the dominant NOx source in the middle to upper troposphere (Schumann and Huntrieser, 2007; Finney et al., 2016b). Moreover, LNOx plays a crucially important role in atmospheric chemistry and global climate by controlling the abundances of OH radical, important greenhouse gases (GHGs) such as ozone and methane, and other trace gases (Labrador et al., 2005; Schumann and Huntrieser, 2007; Wild, 2007; Liaskos et al., 2015; Finney et al., 2016a; Murray, 2016; Tost, 2017; He et al., 2022b).

Reportedly, the lightning flash rate (LFR) is related to the stage of convective cloud development (Williams et al., 1989), Convective Available Potential Energy (CAPE) (Romps et al., 2014), cloud liquid–ice water content (Saunders et al., 1991; Finney et al., 2014), and even to the convective precipitation volume (Goodman et al., 1990; McCaul et al., 2009; Romps et al., 2014). Long-term global warming is associated with changes in the overall temperature and relative humidity profiles in the atmosphere and global convective adjustment (Manabe and Wetherald, 1975; Del Genio et al., 2007), which can strongly affect the lightning-related factors described above. Consequently, long-term global warming can be a crucially fundamentally important factor affecting long-term variations in global lightning activity. Findings from many earlier numerical simulation studies manifest that global lightning activities are sensitive to long-term global warming, with most studies showing 5–16% (average around 10%) increases in global lightning activities per 1 K global warming (Price and Rind, 1994; Zeng et al., 2008; Hui and Hong, 2013; Banerjee et al., 2014; Krause et al., 2014; Romps et al., 2014; Clark et al., 2017). However, minor other numerical simulation studies such as those using an ice-based lightning scheme or convective mass flux as a proxy to parameterize lightning have yielded opposite results, suggesting that global lightning activity will decrease under long-term global warming (Clark et al., 2017; Finney et al., 2018).

Aside from long-term global warming, changes in aerosol loading can also be responsible for long-term global lightning activity variations. Aerosols influence lightning activity through aerosol radiative and microphysical effects, but the degree to which the two distinct effects influence regional or global scale lightning activities remains unclear (Yuan et al., 2011; Yang...
Further research is needed. It is urgently necessary to elucidate the effects of aerosol radiative and microphysical effects on lightning on a global scale. The aerosol radiative effects indicate that aerosols can heat the atmospheric layer and can cool the Earth's surface by absorbing and scattering solar radiation (Kaufman et al., 2002; Koren et al., 2004, 2008; Li et al., 2017). Thereby, convection and electrical activities are likely to be inhibited (Koren et al., 2004; Yang et al., 2013; Tan et al., 2016). The microphysical effects suggest that by acting as cloud condensation nuclei (CCN) or as ice nuclei, aerosols can reduce the mean size of cloud droplets, thereby consequently suppressing the coalescence of cloud droplets into raindrops. Consequently, as a result, more liquid water particles are uplifted to higher mixed-phase regions of the troposphere, where they invigorate lightning (Wang et al., 2018; Liu et al., 2020).

The Earth has experienced a significant considerable degree of global warming and changes in AeroPEs since the 1960s (Hoesly et al., 2018; Climate at a Glance | National Centers for Environmental Information (NCEI), 2022). However, how historical lightning has trended and how lightning has responded to historical global warming and changes in AeroPEs are not well examined. This topic is worth exploring because historical lightning densities are indicators for all lightning-related phenomena (Price and Rind, 1994). Exploring the historical global LNOₓ emission trend is also meaningful because it can indicate the effects of LNOₓ emissions on atmospheric chemistry and global climate. Furthermore, investigating the effects of historical global warming and increases in AeroPEs on historical lightning–LNOₓ trends can provide a basis for assessing future lightning–LNOₓ trends.

Large-scale volcanic eruptions such as the 1991 Pinatubo eruption inject tremendous amounts of sulfuric gas into the stratosphere, where it converts to H₂SO₄ aerosols. Consequently, the size of the stratospheric aerosol layer has increased in abundance after the volcanic eruptions. The enhanced stratospheric aerosol layer can cool the Earth's surface heterogeneously and can decrease the total amount of water in the atmosphere (Soden et al., 2002; Boucher, 2015, p.63). The near-global perturbations in the radiative energy balance and meteorological fields caused by such strong volcanic eruptions might influence global lightning activities. If so, there might be ramifications for all lightning-related phenomena. Nevertheless, they remain poorly understood.

In our earlier work, we developed a new process and ice-based lightning scheme called the ECMWF-McCAUL scheme (He et al., 2022b). This lightning scheme was developed by combining benefits of the lightning scheme used in the European Centre for Medium-Range Weather Forecasts (ECMWF) forecasting system (Lopez, 2016) and those presented in reports by McCaul’s work et al. (McCaul et al., 2009). The ECMWF-McCAUL scheme simulated the best lightning density spatial distributions among four existing lightning schemes when compared against satellite lightning observations (Lightning Imaging Sensor (LIS) and Optical Transient Detector (OTD)) during 2007–2011. The sensitivity of global lightning activity
to changes in surface temperature on a decadal timescale was estimated as 10.13% K^{-1} by using the ECMWF-McCAUL scheme (He et al., 2022b), which is close to most past estimates (average around 10% K^{-1}).

Using a chemistry–climate model CHASER (MIROC) with two lightning schemes (the widely used cloud top height scheme and the ice-based ECMWF-McCAUL scheme), we quantitatively investigated historical lightning–LNO_x trends and found how global warming, increases in AeroPEs, and the Pinatubo eruption respectively influenced them.

Using two lightning schemes, we demonstrated the sensitivities of different lightning schemes to historical global warming, increases in AeroPEs, and the Pinatubo eruption.

Research methods, including the model description and experiment setup, are described in Sect. 2. In Sect. 3.1, the simulated historical lightning distributions and trends are validated using LIS/OTD lightning observations. Section 3.2 presents the effects of global warming and increases in AeroPEs on historical lightning–LNO_x trends. In Sect. 3.3, the effects of the Pinatubo volcanic eruption on historical lightning–LNO_x trends are discussed. Section 3.4 elucidated model intercomparisons of lightning flash rate LFR trends and variation between our study (CHASER) and other CMIP6 model outputs. Section 4 presents relevant discussions and conclusions obtained from this study.

2 Method

2.1 Chemistry–climate model

We used the CHASER (MIROC) global chemistry–climate model (Sudo et al., 2002; Sudo and Akimoto, 2007; Watanabe et al., 2011; Ha et al., 2021) for this study, which incorporated consideration of detailed chemical and physical processes in the troposphere and stratosphere. The CHASER version adopted for this study simulates the distributions of 94 chemical species while reflecting the effects of 269 chemical reactions (58 photolytic, 190 kinetic, and 21 heterogeneous). As processes associated with tropospheric chemistry, Non-Methane Hydrocarbons (NMHC) oxidation and the fundamental chemical cycle of O_3–NO_x–HO_x–CH_4–CO are considered. CHASER simulates stratospheric chemistry involving the Chapman mechanisms and catalytic reactions associated with HO_x, NO_x, ClO_x, and BrO_x. Moreover, it simulates the formation of polar stratospheric clouds (PSCs) and heterogeneous reactions occurring on their surfaces. CHASER is on-line-coupled to MIROC AGCM ver. 5.0 (Watanabe et al., 2011), which simulates cumulus convection (Arakawa–Schubert scheme) and grid-scale large-scale condensation to represent cloud and precipitation processes. The radiation flux is calculated using a two-stream k distribution radiation scheme, which considers absorption, scattering, and emissions by aerosol and cloud particles as well as by gaseous species (Sekiguchi and Nakajima, 2008; Goto et al., 2015). The aerosol component in CHASER is coupled with the SPRINTARS aerosol model (Takemura et al., 2009), particularly for simulating primary organic carbon, sea-salt, and dust, which is also based on MIROC. The aerosol radiation effects are considered in both large-scale condensation and cumulus convection schemes, although the aerosol microphysical effects are only reflected in the large-scale condensation scheme.
This study used a horizontal resolution of T42 (2.8° × 2.8°), with vertical resolution of 36 σ-p hybrid levels from the surface to approximately 50 km. Anthropogenic and biomass burning emissions were obtained from the CMIP6 forcing datasets (van Marle et al., 2017; Hoesly et al., 2018) for 1959–2014 (https://esgf-node.llnl.gov/search/input4mips/, last access: 19 September 2022). Interannual variation in biogenic emissions for isoprene, monoterpene, acetone, and methanol, were considered using an off-line simulation by the Vegetation Integrative Simulator for Trace Gases (VISIT) terrestrial ecosystem model (Ito and Inatomi, 2012). The residual biogenic emissions (ethane, propane, ethylene, propene) used are climatological values derived from the Model of Emissions of Gases and Aerosols from Nature (MEGAN) modeling system (Guenther et al., 2012).

The CHASER (MIROC) global chemistry–climate model originally parameterizes lightning with the widely used cloud top height scheme (Price and Rind, 1992). A newly developed ice-based lightning scheme called the ECMWF-McCAUL here had been implemented into CHASER (MIROC) (He et al., 2022b). The ECMWF-McCAUL scheme computes lightning flash rates $LFRs$ as a function of CAPE and column precipitating ice (including $Q_{Ra}$, $Q_{Ra}$ represents the total volumetric amount of cloud ice, graupel, and snow in the charge separation region). Compared with the cloud top height, a salient advantage of the ECMWF-McCAUL scheme is that it has a direct physical link with the charging mechanism.

### 2.2 Lightning NOx emission parameterizations

We tested two lightning schemes for this study. The first lightning scheme is the widely used cloud top height (CTH) scheme (Price and Rind, 1992), which was used originally used in CHASER (MIROC). This lightning scheme calculates the lightning flash rate using the following equations to calculate $LFR$.

\[
F_l = 3.44 \times 10^{-5} H^{4.9} \quad (1)
\]

\[
F_o = 6.2 \times 10^{-4} H^{1.73} \quad (2)
\]

Therein, $F$ represents the total flash frequency (fl. min$^{-1}$), $H$ stands for the cloud-top height (km), and subscripts $l$ and $o$ respectively denote the land and ocean (Price and Rind, 1992). Actually, we realize the CTH scheme in CHASER using the following equations (Eq. (3) and Eq. (4)). Each model layer’s cumulus cloud fractions are used to weight the calculated lightning densities from that layer in the CTH scheme.

\[
F_l = \sum_{i=1}^{n=36} adj\_factor \times Cu\_CF_i \times (H_i - H_{surf\_ace})^{4.9} \quad (3)
\]

\[
F_o = \sum_{i=1}^{n=36} adj\_factor \times Cu\_CF_i \times (H_i - H_{surf\_ace})^{1.73} \quad (4)
\]

In those equations, $i$ denotes the model layer index. Also, $adj\_factor$ represents adjustment factors that differ for different model layers and model grids. $Cu\_CF_i$ symbolizes the cumulus cloud fraction at model layer $i$, $H_i$ and $H_{surf\_ace}$ respectively denote the altitude of model layer $i$ and the altitude of the model’s surface layer.
The second lightning scheme used for this study is a newly developed one named the ECMWF-McCAUL scheme (He et al., 2022b), which is based on the original ECMWF scheme and findings reported by McCaul et al. (2009). The ECMWF-McCAUL scheme calculates lightning flash rates\textit{LFRs} as a function of \textit{CAPE} (m$^2$ s$^{-2}$) and column precipitating ice ($Q_{Ra}$) ($Q_{Ra}$ symbolizes the total volumetric amount of cloud ice, graupel, and snow in the charge separation region) as

\begin{equation}
 f_i = \alpha_i Q_{Ra} CAPE^{1.3}
 \end{equation}

\begin{equation}
 f_o = \alpha_o Q_{Ra} CAPE^{1.3}
 \end{equation}

where $f_i$ and $f_o$ respectively symbolize the total flash density (fl. m$^{-2}$ s$^{-1}$) over land and ocean. In addition, $\alpha_i$ and $\alpha_o$ are constants (fl. s$^{1.6}$ kg$^{-1}$ m$^{-2.6}$) determined after calibration against LIS/OTD climatology, respectively, for land and ocean. For this study, $\alpha_i$ and $\alpha_o$ are set respectively to $2.67 \times 10^{-16}$ and $1.68 \times 10^{-17}$. In the charge separation region (from 0° to 25°C isotherm), $Q_{Ra}$ (kg m$^{-2}$) is expressed as a proxy for the charging rate because of collisions between graupel and hydrometeors of other types (McCaul et al., 2009). Moreover, $Q_{Ra}$ represents the total volumetric amount of hydrometeors of three kinds (graupel, snow, and cloud ice) within the charge separation region, calculated as

\begin{equation}
 Q_{Ra} = \int_{z_0}^{z_{-25}} (q_{graup} + q_{snow} + q_{ice}) \bar{\rho} dz,
 \end{equation}

where $q_{graup}$, $q_{snow}$, and $q_{ice}$ respectively represent the mass mixing ratios (kg kg$^{-1}$) of graupel, snow, and cloud ice.

Also, in addition, $q_{ice}$ was diagnosed using Arakawa–Schubert cumulus parameterization. Then, $q_{graup}$ and $q_{snow}$ were computed at each vertical level of the model using the following equations.

\begin{equation}
 q_{graup} = \beta \frac{P_f}{\bar{\rho} V_{graup}}
 \end{equation}

\begin{equation}
 q_{snow} = (1 - \beta) \frac{P_f}{\bar{\rho} V_{snow}}
 \end{equation}

In those equations, $P_f$ represents the vertical profile of the frozen precipitation convective flux (kg m$^{-2}$ s$^{-1}$), $\bar{\rho}$ denotes the air density (kg m$^{-3}$), and $V_{graup}$ and $V_{snow}$ respectively express the typical fall speeds for graupel and snow set to 3.1 and 0.5 m s$^{-1}$ for this study. For land, the dimensionless coefficient $\beta$ is set as 0.7, while for oceans, whereas it is set to 0.45, for oceans to consider the observed lower graupel content over the oceans.

Based on the cold cloud depth, a fourth-order polynomial (equation 10) is used to calculate the proportion of total flashes that are cloud-to-ground ($p$). An earlier report of the literature describes the method (Price and Rind, 1993).

\begin{equation}
 p = \frac{1}{64.9 - 36.54D + 7.493D^2 - 0.648D^3 + 0.021D^4}
 \end{equation}

The depth of the cloud above the 0°C isotherms is represented by $D$ (km) in that equation.

According to recent studies, the intra-cloud (IC) lightning flashes are as efficient as cloud-to-ground (CG) lightning flashes at producing NO$_x$. The lightning NO$_x$ production efficiency is estimated as 100–400 mol per flash (Ridley et al., 2005; Cooray et al., 2009; Ott et al., 2010; Allen et al., 2019). The LNO$_x$ production efficiency for IC and CG are therefore set to the same value (250 mol per flash) in CHASER, which is the median of the commonly cited range of 100–
400 mol per flash. Consequently, in this study, the distinctions between IC and CG do not affect the distribution or magnitude of LNO\textsubscript{x} emissions in this study. It is noteworthy that marked uncertainties are involved in determining the LNO\textsubscript{x} production efficiency (Allen et al., 2019; Bucsela et al., 2019). The choice of a different LNO\textsubscript{x} production efficiency might affect the simulation of LNO\textsubscript{x} emissions. Further research must be undertaken to implement and validate a more sophisticated parameterization of LNO\textsubscript{x} production efficiency in chemistry–climate models.

2.3 Lightning observation data for model evaluation

We used LIS/OTD gridded climatology datasets for this study, consisting of climatologies of total lightning flash rates LFR\textsubscript{s} observed using the Lightning Imaging Sensor (LIS) and Optical Transient Detector (OTD). The OTD aboard the MicroLab-1 satellite and LIS aboard the Tropical Rainfall Measuring Mission (TRMM) satellite (Cecil et al., 2014). Both sensors detected lightning by monitoring pulses of illumination produced by lightning in the 777.4 nm atomic oxygen multiplet above background levels. In low Earth orbit, both sensors viewed Earth locations for approximately 3 min during the pass of the OTD or 1.5 min during passing of the LIS. Each day, OTD and LIS orbit respectively orbited the globe 14 times and 16 times a day, respectively. OTD observed data between +75 and -75° latitude during May 1995 through March 2000, whereas LIS collected data between +38 and -38° latitude during January 1998 through April 2015. This study uses the LIS/OTD 2.5 Degree Low Resolution Time Series (LRTS\textsubscript{S}). LRTS, which provides daily lightning flash rates LFR\textsubscript{s} on a 2.5° regular latitude–longitude grid for May 1995 through April 2015.

2.4 CMIP6 model outputs for model comparison

For the comparison of different model outputs from our study (CHASER) and other Earth system models or chemistry–climate models, we used the lightning flash rate LFR and surface temperature data from the CMIP6 CMIP Historical experiments from CESM2-WACCM (Danabasoglu, 2019), GISS-E2-1-G (Kelley et al., 2020), and UKESM1-0-LL (Tang et al., 2019). CESM2-WACCM uses the Community Earth System Model ver. 2 (Danabasoglu et al., 2020). The CESM2 is an open-source fully coupled Earth system model. The Whole Atmosphere Community Climate Model ver. 6 (WACCM6) is the atmospheric component coupled to the other components in CESM2. The GISS-E2-1-G is the NASA Goddard Institute for Space Studies (GISS) chemistry–climate model version E2.1 based on the GISS Ocean v1 (G01) model (Miller et al., 2014; Kelley et al., 2020). The UKESM1-0-LL is the UK's Earth system model, details of which are were described by Sellar et al. (2019). We used 3 ensembles from CESM2-WACCM, 9 ensembles from GISS-E2-1-G, and 18 ensembles from UKESM1-0-LL. Table S1 presents all the ensemble members used for this study.
2.5 Experiment setup

We have conducted six sets of experiments with each set of experiments conducted using both the ECMWF-McCAUL (abbreviated as F1) and CTH (abbreviated as F2) schemes. Table 1 presents the major settings of all experiments with the relative explanations of those settings. STD-F1/F2 are standard experiments with the simulation period of 1959–2014. They aimed at reproducing the historical trends of lightning and LNOX. Climate1959-F1/F2 are experiments that keep the climate simulations fixed to 1959 to derive the effects of global warming on historical lightning trends. ClimateAero1959-F1/F2 are intended to reflect the conditions with climate simulations and aerosol and aerosol precursors fixed to 1959. The Aero1959-F1/F2 experiments are the same as the STD-F1/F2 experiments, except for the AeroPEs fixed to 1959. The fifth set of experiments (Volca-off-F1/F2) was intended to exclude the influences of the Pinatubo volcanic eruption to compare to the STD-F1/F2 and to evaluate the effects of the Pinatubo eruption on historical lightning–LNOX trends and variation.

We simulate volcanic aerosol forcing by considering the prescribed stratospheric aerosol extinction in the radiation scheme. We used the NASA Goddard Institute for Space Studies (GISS) (Sato et al., 1993) and Chemistry–Climate Model Initiative (CCMI) (Arfeuille et al., 2013) stratospheric aerosol dataset as the stratospheric aerosol climate data. The NASA GISS dataset includes monthly zonal-mean stratospheric aerosol optical thickness in four spectral bands. The CCMI dataset for CHASER includes monthly zonal-mean stratospheric aerosol extinction coefficients in 20 spectral bands. To remove the volcanic perturbation but maintaining the stratospheric background aerosol in the Volca-off-F1/F2, we used the three-sigma rule to process the Stratospheric Aerosol Climatology (SAC) during June 1991 – May 1994 using the following equation. The three-sigma rule is often used to detect the outliers. This rule is appropriate to use here to discern the outliers (the perturbation of SAC caused by a strong volcanic eruption). 1996.

\[
SAC_{\text{no.pinatubo}} = \begin{cases} 
\frac{SAC_{\text{background}}}{|SAC_{\text{raw}} - SAC_{\text{background}}|} > 3\sigma, \\
\frac{SAC_{\text{raw}}}{|SAC_{\text{raw}} - SAC_{\text{background}}|} \leq 3\sigma \end{cases}
\]

\[
\{SAC_{\text{background}}, |SAC_{\text{raw}} - SAC_{\text{background}}| > 1.96\sigma, \\
SAC_{\text{raw}}, |SAC_{\text{raw}} - SAC_{\text{background}}| \leq 1.96\sigma \}
\]

(11)

In that equation, \(SAC_{\text{no.pinatubo}}\) denotes the stratospheric aerosol climatological data as input data for Volca-off-F1/F2 experiments, \(SAC_{\text{background}}\) represents the stratospheric background aerosol climatological data (For this study, \(SAC_{\text{background}}\) is the corresponding temporal averaged values of the NSAS GISS and CCMI stratospheric aerosol dataset during 2001–2010 June 1986 – May 1991 and June 1996 – May 2001, when the time is close to the eruption and the stratosphere was less affected by volcanic eruptions). \(SAC_{\text{raw}}\) stands for the original values of NSAS GISS and CCMI stratospheric aerosol dataset during June 1991 – May 1994. Moreover, \(\sigma\) symbolizes the standard deviations of stratospheric background aerosol climate data (For this study, \(\sigma\) are the corresponding standard deviations
of NSAS/NASA GISS and CCMI stratospheric aerosol dataset during June 1986 – May 1991 and June 1996 – May 2001. As displayed in Eq. (11), when the absolute differences between $SAC_{\text{raw}}$ and $SAC_{\text{background}}$ are larger than 1.96$\sigma$, we replace the original values (June 1991 – May 1996) of the SAC with the temporal averaged values of the NASA GISS and CCMI dataset during June 1986 – May 1991 and June 1996 – May 2001. When the absolute differences between $SAC_{\text{raw}}$ and $SAC_{\text{background}}$ are equal to or smaller than 1.96$\sigma$, we still use the original values (June 1991 – May 1996) of the SAC for the Volca-off experiments. The value of 1.96$\sigma$ corresponds to the 95% confidence interval, which can remove the Pinatubo perturbation sufficiently but which can maintain the background level of stratospheric aerosol during June 1991 – May 1996. Furthermore, the influences of the Pinatubo eruption also affected the HadISST SSTs/sea ice fields. To remove the Pinatubo eruption's influences from the SSTs/sea ice fields in the Volca-off experiments also, we replaced the 1991-06 – 1995-05 SSTs/sea ice data with HadISST SSTs/sea ice climatological data during 1985–1990 when conducting the Volca-off experiments. The 1985–1990 period was chosen because it is close to approximately the period of 1991-06 – 1995-05 and because the SSTs/sea ice fields were less affected by volcanic activity during 1985–1990.

All the experiments calculate the LNO$_x$ emissions rates interactively by LNO$_x$ emission parameterizations except STD-rVolcaoff experiments. The STD-rVolcaoff experiments are the same as the STD experiments except for reading the daily LNO$_x$ emission rates calculated from the Volca-off experiments. The STD-rVolcaoff experiments are conducted for comparison with STD experiments to elucidate the effects of the changed LNO$_x$ emissions caused by the Pinatubo eruption on atmospheric chemistry (typically methane lifetime).

### Table 1: All experiments conducted for this study

<table>
<thead>
<tr>
<th>Name of experiment</th>
<th>Period</th>
<th>Climate (SSTs, sea ice, GHGs)$^a$</th>
<th>Anthropogenic and biomass burning emissions</th>
<th>Biogenic emissions</th>
<th>Stratospheric aerosol climatology</th>
</tr>
</thead>
<tbody>
<tr>
<td>Climate1959-F1/F2</td>
<td>1959–2014</td>
<td>Fixed to 1959$^d$</td>
<td>CMIP6 1959–2014</td>
<td>VISIT and MEGAN$^f$</td>
<td>As above</td>
</tr>
<tr>
<td>ClimateAero1959-</td>
<td>1959–2014</td>
<td>Fixed to 1959</td>
<td>AeroPEs fixed to 1959$^e$</td>
<td></td>
<td>As above</td>
</tr>
<tr>
<td>F1/F2</td>
<td></td>
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</tbody>
</table>
a For the model simulations, the climate is simulated by the prescribed SSTs/sea ice fields and the prescribed varying concentrations of GHGs (CO₂, N₂O, methane, chlorofluorocarbons – CFCs – and hydrochlorofluorocarbons – HCFCs) used only in the radiation scheme. The SSTs/sea ice fields are obtained from the HadISST dataset (Rayner et al., 2003). The prescribed GHGs concentrations are derived from CMIP6 forcing datasets (Meinshausen et al., 2017).
b We use “F1” to stand for the ECMWF-McCAUL scheme; “F2” represents the CTH scheme.
c Stratospheric aerosol radiative forcing is simulated using the prescribed stratospheric aerosol extinction, which is obtained from the NASA GISS (Sato et al., 1993) and CCMI (Arfeuille et al., 2013) stratospheric aerosol dataset.
d The climate is fixed to 1959 for the whole simulation period using the 1959 SSTs/sea ice field and GHG concentrations during the simulation period.
e Aerosol (BC, OC) and aerosol precursor (NOₓ, SO₂) emissions (anthropogenic + biomass burning) are fixed to 1959 throughout the simulation period.
f Several biogenic emissions are interannually varying, including isoprene, monoterpenes, acetone, and methanol, which were calculated using an off-line simulation by using the Vegetation Integrative Simulator for Trace Gases (VISIT) terrestrial ecosystem model (Ito and Inatomi, 2012). Some other reactive biogenic VOCs (ethane, propane, ethylene, propene) used are climatological data derived from the Model of Emissions of Gases and Aerosols from Nature (MEGAN) modeling system (Guenther et al., 2012).
g Here the 1991-06 – 1995-05 SSTs/sea ice data were replaced with HadISST SSTs/sea ice climatological data during 1985–1990.

3 Results and Discussion

3.1 Validation of the simulated historical lightning distribution and trend

To increase the credibility of the conclusions obtained based only on the numerical simulations, the model calculations must be evaluated using observational data. We used the LIS/OTD observations to evaluate the spatial and temporal distribution and historical lightning trends simulated by CHASER (MIROC). Figures 1a–elc show the annual mean spatial distribution of lightning observed by LIS/OTD and from model simulations using the ECMWF-McCAUL and CTH schemes. Both the ECMWF-McCAUL and CTH schemes generally captured the hotspots of lightning (Central Africa, Maritime Continent, South America), with strong spatial correlations between observations and model simulations ($R > 0.75$).
Figure 1d exhibits strong spatial correlation between observations and simulation results maintained throughout the simulation period (1959–2014).

Figure 1: Annual mean lightning flash densities from (a) LIS/OTD satellite observations spanning 1996–2000, (b) the STD experiment (1960–2014) with the ECMWF-McCAUL scheme used, (c) the STD experiment (1960–2014) with the CTH scheme used. $R$ and RMSE shown in the titles of panels (b) and (c) are calculated between panels (b)–(c) and (a). Panel (d) presents the spatial correlation coefficients between modeled annual mean spatial lightning densities distribution of each year and LIS/OTD lightning climatologies during 1996–2000.
Figure 2: Lightning flash rate LFR anomalies of 1996–2013 within ±41.25° latitude obtained from two numerical experiments (STD-F1 and STD-F2) and LIS/OTD satellite observations. Curves represent the monthly time-series data of the ±41.25° latitude mean lightning flash rate LFR anomalies with the 1-D Gaussian (Denoising) filter applied. Lines are the fitting curves of the monthly time-series data of the ±41.25° latitude mean lightning flash rate LFR anomalies. Trends of the lightning flash rate LFR anomalies in % yr⁻¹ are also shown in the legends.

The LIS/OTD observations are also used to evaluate historical lightning trends simulated by CHASER (MIROC). Because almost all valid LIS/OTD observations exist only within ±41.25° latitude during 1996–2013, we examined the ±41.25°
latitude mean lightning flash rate $LFR$ anomaly (1996–2013) calculated from LIS/OTD observations and STD-F1/F2 numerical experiments (Fig. 2). We also note some missing values within the $\pm 41.25^{\circ}$ latitude in LIS/OTD observations. To keep the comparisons between observations and simulations as like-for-like, when we encounter a missing value in the LIS/OTD observations during spatial averaging, we also treat the CHASER simulated value at the same location as a missing value. As displayed in Fig. 2, even when the interannual variations of the lightning flash rate $LFR$ anomaly sometimes differ between observations and simulations, the overall trends of lightning flash rate $LFR$ anomaly simulated by using both schemes well matched the LIS/OTD observations. Neither the lightning flash rate $LFR$ anomaly (within $\pm 41.25^{\circ}$ latitude) derived from LIS/OTD observations nor simulations show a significant trend for 1996–2013 using the Mann–Kendall rank statistic test (significance set as inferred for 5%). The global lightning flash rate $LFR$ anomaly (during 1993–2013) obtained from simulations (STD-F1/F2) also shows no significant trend, which is consistent with the Schuman Resonance (SR) intensity observations (1993–2013) at Rhode Island, USA (Earle Williams, 2022). However, the SR observations in Rhode Island (USA) exclude consideration of the influences of solar cycles, which makes it less appropriate for lightning trend evaluation.

![Figure 3: Mean annual cycle in area average LFR during 1996–2003. The area average was taken over the grid cells where valid LIS/OTD lightning observations exist. LFR is normalized by min-max normalization.](image)

We further investigated the seasonal variabilities of simulated LFR and compared them against LIS/OTD observations. The results are depicted in Fig. 3. Both the CTH and ECMWF-McCAUL schemes captured the peak during JJA, but the overestimation of LFR by F1/F2 during DJF is also noticeable. Figure S1 presents comparison of the LFR global distribution in different seasons during 1996–2003 from LIS/OTD lightning observations and STD experiment outputs. Generally, CHASER well captured the spatial distribution of LFR in all four seasons when compared against LIS/OTD observations. The spatial correlation coefficients ($R$) between observations and simulations are highest ($R=0.80$ for both lightning schemes) in DJF, indicating CHASER’s considerable capability to reproduce the LFR spatial distribution in DJF. As displayed in the first row of Fig. S1, the overestimation of LFR by F1/F2 during DJF is primarily attributable to the overestimation of LFR within...
the Maritime Continent and South America, but this might also be attributable to the underestimation of LFR by LIS/OTD within these two regions. It is believed that the LIS/OTD lightning detection efficiency is highly sensitive to the characteristic of convective clouds (cloud albedo, cloud optical thickness, etc.) (Boccippio et al., 2002; Cecil et al., 2014). High cloud albedo and cloud optical thickness might engender the underestimation of LFR by LIS/OTD. It is also noteworthy that the seasonal variation and long-term trend of global lightning are strongly influenced by distinct different factors. The seasonal variation of global lightning activities is most strongly affected by the 23° obliquity of Earth’s orbit and the asymmetric distribution of the continent between the Northern and Southern hemispheres. However, the long-term global lightning trend we investigated for this study is controlled mainly by climate forcers such as aerosols and GHGs. To minimize the effects of LFR seasonal variation on our study’s results, we deseasonalized the results shown in all figures and tables by calculating their anomaly based on raw data. The validation described above and the deseasonalization of our study’s results justified that the LFR seasonal variation (and the uncertainties in the simulation of LFR seasonal variation) in our study has a limited effect on these study results.

3.2 Effects of global warming and increases in AeroPEs on historical lightning–LNOₓ trends

As introduced in Sect. 1, global warming and changes in AeroPEs are the two main factors which influence the long-term (1960–2014) historical lightning trends (Hereinafter, historical lightning trends indicate the lightning trends of 1960–2014.). To analyze the effects of global warming on historical lightning trends, we designed and conducted two sets of experiments: one set of experiments including “global warming” (STD-F1/F2) and another set of experiments excluding “global warming” (Climate1959-F1/F2). Figures 3a4a and 3b4b respectively depict the global surface temperature anomalies calculated from using the ECMWF-McCAUL and CTH schemes. The STD and Aero1959 experiments show an increasing trend (around 0.11 K decade⁻¹) of global mean surface temperature anomalies, which closely approximates the trend (around 0.15 K decade⁻¹) obtained from NOAA’s National Centers for Environmental Information (NCEI) (Figs. 4c, 4f). Global temperature change data from 1880 to the present are available from the NCEI, which tracks the variations of the Earth’s temperature based on thousands of stations’ observation data around the globe (Climate at a Glance | National Centers for Environmental Information (NCEI), 2022). When the prescribed SSTs/sea ice fields and GHGs concentrations were fixed to 1959 throughout the simulation period, the simulated trends of global mean surface temperature anomalies turned out to be flat (Climate1959 and ClimateAero1959). To elucidate the effects of increases in AeroPEs on averaged surface temperature to the greatest extent possible, we also show the averaged surface temperature anomaly only over land regions (Figs. 3d–4d–4f). The simulated global mean land surface temperature anomalies are also well-matched with the NCEI observational data. The aerosol cooling effect can be more evident when only particular addressing examining surface temperature trends averaged over land (Figs. 3d–4d–4e).
Figure 34: Monthly time-series data of global mean surface temperature anomalies with 1-D Gaussian (Denoising) filter applied and their fitting curves calculated from the outputs of numerical experiments (a–b) and obtained from NCEI (c). Figures 3d–Panels (d)–(f) are the same as Figs. 3a–panels (a)–(c), but the averaged surface temperature anomalies are only calculated within the global land regions. The trends of the fitting curves in K decade$^{-1}$ are also presented in the legends.
Figure 4: Figures 4(a) and (b) show monthly time-series data of global mean lightning flash rate LFR anomalies with 1-D Gaussian (Denoising) Filter applied and their fitting curves of different experiments simulated respectively using the ECMWF-McCAUL scheme and CTH scheme. Figures 4(c) and (d) are the same as Figs. 4 panels (a) and (b), except that the averaged lightning flash rate LFR anomalies are calculated only within global land regions. Trends of the fitting curves (% yr\(^{-1}\)) are also shown in the legends.
Figures 4 and 5 panels (a) and (b) respectively show portray the global mean lightning flash rate anomalies and their fitting curves obtained from the outputs of the ECMWF-McCAUL scheme and CTH scheme. The global lightning trend obtained from the STD-F1 experiment turned out to be statistically flat (0.0% yr\(^{-1}\)), whereas the outputs of the STD-F2 experiment exhibit a not significant increasing global lightning trend (0.03% yr\(^{-1}\)) determined using the Mann–Kendall rank statistic (significance inferred for 5%).

Comparison of the lightning trends calculated from the STD and Climate1959 experiments, we found showed that both lightning schemes demonstrated that the historical global warming (1960–2014) enhances the global lightning trends toward positive trends (around 0.03% yr\(^{-1}\) or 3% K\(^{-1}\)). The effects of Global warming effects on historical lightning trends were evaluated as significant using the Mann–Kendall rank statistic, with significance inferred for 5%, when using the CTH scheme, but not in the case of the ECMWF-McCAUL scheme (Hussain and Mahmud, 2019). The differences in lightning trends simulated by the STD-F1/F2 and Aero1959-F1/F2 experiments indicate that the increases in AeroPEs during 1960–2014 significantly suppress the global lightning trends (-0.07% yr\(^{-1}\) – -0.04% yr\(^{-1}\)). It is noteworthy that this suppression of lightning trends is only attributable to the aerosol radiative effects. Further research is needed must be conducted to elucidate the long-term effects of aerosols on lightning through aerosol microphysical effects. We also investigated lightning trends only over land regions (Figs. 4c–d, 5c–d) to ascertain the effects of changes in AeroPEs to the greatest extent possible. When observing the lightning trends over land only, the degree of suppression of lightning trends by attributable to increases in AeroPEs expands to -0.10% yr\(^{-1}\) – -0.05% yr\(^{-1}\), which is attributable to most AeroPEs and their growth coming from land regions. It is noteworthy that we used the same SSTs/sea ice data in the Aero1959 as those used for STD experiments. The SSTs/sea ice data also reflected the effects of increases in AeroPEs. Therefore, we might underestimate the effects of increases in AeroPEs on lightning trends by comparing the results of STD and Aero1959 experiments.
Figure 5: Figures 5a–5e: Panels (a) and (b) respectively show monthly time-series data of global mean CAPE and $Q_{Ra}$ anomalies with 1-D Gaussian (Denoising) filter applied and their fitting curves simulated using the ECMWF-McCAUL scheme. Figures 5 Panels (c) and (d) respectively show differences in the CAPE anomaly trend and $Q_{Ra}$ anomaly trend of the STD-F1 and Aero1959-F1 experiments in the global map. Figure 5e portrays the vertical profiles of the trend of global mean cumulus cloud fraction anomaly simulated by the CTH scheme. Figure 5f Panel (f) depicts the relative contributions of each layer’s cumulus to total lightning density in 1960 and 2014, as calculated from the outputs of the STD-F2 experiment.

For the ECMWF-McCAUL scheme, model outputs affirm that global warming can enhance the global mean CAPE anomaly slightly and suppress the global mean $Q_{Ra}$ anomaly (Figs. 6a–6b). Earlier studies have also indicated that the total solid (cloud ice, snow, and graupel) mass mixing ratio within charge separation regions is lower under global warming. Moreover, possible explanations are given in those studies (Finney et al., 2018; Romps, 2019).
can be suppressed by earlier global warming, probably because global warming engenders the shifting of the 0°C—25°C isotherm to the higher region. Because global warming enhances global convection activities, and because lightning formation is highly related to convection activity, global warming enhances the historical global lightning trend simulated by using the ECMWF-McCAUL scheme mainly as a result of the simulated CAPE trend, which is enhanced by global warming. The past increases in AeroPEs exert negligible effects on the trends of global mean CAPE and $Q_{Ra}$ anomalies, as displayed in Figs. 5a–b6a–6b. However, the past increases in AeroPEs suppress the CAPE and $Q_{Ra}$ trend within the tropical and subtropical terrestrial regions, where lightning densities are high (Figs. 5e–6c–6d). Weaker convection activities (smaller CAPE) and fewer hydrometeors (cloud ice, graupel, snow) in the charge separation regions (0°C—25°C isotherm) lead to less lightning. These are the main causes for the suppression of the historical lightning trends induced by increases in AeroPEs through aerosol radiative effects. It is noteworthy that, because the aerosol microphysical effects are only considered in the grid-scale large-scale condensation scheme, our study might underestimate the aerosol microphysical effects which can enhance the trends of $Q_{Ra}$ and LFR toward the positive direction.

To explain the results simulated by the CTH scheme, we investigated the vertical profiles of the trend of the global mean cumulus cloud fraction (Fig. 5e). This anomaly (Fig. 6e). Investigating cumulus cloud fraction is reasonable because each model layer’s cumulus cloud fractions are used to weight the calculated lightning densities from that layer in the CTH scheme, as introduced in equations (3) and (4). Figure 5f shows the relative contributions of each model layer’s cumulus to the calculated global total lightning densities in 1960 and 2014 obtained using the CTH scheme. As Fig. 6f displayed, the vertical profiles of relative contribution in 1960 and 2014 are almost identical. Cumulus convection is positively correlated with lightning formation, which is the scientific basis of parameterizing lightning densities using the cumulus cloud top height: the CTH scheme. The historical global warming enhances the lightning trend simulated by the CTH scheme mainly because the simulated historical global warming increases the cumulus reaching 200 hPa, which greatly to the simulated global total lightning density (Figs. 5e–6c–6f). The increases in the deep convective cloud are regarded as related to the increases in tropopause height attributable to global warming, which is shown in Fig. S4S2. The past increases in AeroPEs suppress the lightning trend simulated by the CTH scheme because increases in AeroPEs decrease the cumulus reaching 200 hPa as well as the cumulus within the lower to middle troposphere by aerosol radiative effects (Fig. 5f–6c–6e). In addition, in the supplement, we present a figure (Fig. S2S3) resembling Fig. S6, but which includes only consideration of land regions. The mechanisms of global warming and increases in AeroPEs affecting lightning trends over land regions are similar to those described above on a global scale. We do not discuss details of them in detail here.
Figure 6: Lightning flash rate trends: Trends of LFR anomaly (% yr$^{-1}$) during 1960–2014 on the two-dimensional map. The trend at every point was calculated from the function of approximating curve for the 1960–2014 time-series data (LFR anomaly) at each grid cell. The area in which the trend was found to be significant by the Mann–Kendall rank statistic test (significance level inferred for 5%) is marked with hatched lines.
Figure 78: Differences in lightning flash rate trends of LFR anomaly during 1960–2014 on the global map. The area in which the trend of the differences of lightning flash rate LFR anomaly time-series data was found to be significant by the Mann–Kendall rank statistic test (significance level inferred for 5%) is marked with hatched lines.

We also investigated lightning trends simulated from different experiments with the global map (Fig. 67). Both the ECMWF-McCAUL and the CTH schemes show that the lightning increased significantly in most parts of the Arctic region and decreased in some parts of the Southern Ocean during 1960–2014 (Figs. 6a–e, 7a, 7c). The significant lightning trends presented in Figs.
6a almost disappeared when the climate simulations were fixed to 1959 (Figs. 6b, 7b, 7f), indicating the considerable effects of global warming on the trend of global lightning activities. Furthermore, the effects of past global warming and increases in AeroPEs on the lightning trends on the global map are displayed in Fig. 78. Figures 7a, c indicate that past global warming enhances lightning activities within the Arctic region and Japan, which is consistent with findings of an earlier study from which Japan thunder day data were reported (Fujibe, 2017). Figures 7a, 8a and 8c also show that historical global warming suppresses lightning activities around New Zealand and some parts of the Southern Ocean. Both lightning schemes demonstrated that the historical increases in AeroPEs suppress lightning activities in some parts of the Southern Ocean and South America. The ECMWF-McCAUL scheme also suggests that historical increases in AeroPEs suppress lightning activities by aerosol radiative effects in some parts of India and China, where AeroPEs increased dramatically during 1960–2014 because of rapid economic development and energy consumption. Many observation-based studies indicate that aerosols can invigorate lightning activities in some regions of China and India, typically under relatively clean conditions (e.g., AOD < 1.0), which is attributable to the aerosol microphysical effects (Wang et al., 2011; Zhao et al., 2017; Lal et al., 2018; Liu et al., 2020; Shi et al., 2020; Zhao et al., 2020). Therefore, a total positive effect of aerosol on historical lightning trends in China and India cannot be ruled out. We further provided the same figures as Figs. 67 and 78, but using different units (fl. km$^{-2}$ yr$^{-2}$) in the supplementary information (Figs. S3 and S4 and S5). Figures S3 and S4 and S5 show that the absolute lightning trends (fl. km$^{-2}$ yr$^{-2}$) and the effects of global warming and increases in AeroPEs on the absolute lightning trends are slight in high-latitude regions: but prominent in tropical areas.
Figure 89: Time-series data of 1960–2014 annual global LNO\textsubscript{x} production anomalies (TgN yr\textsuperscript{−1}) and their fitting curves simulated using the ECMWF-McCAUL scheme (a) and the CTH scheme (b). Trends of the fitting curves in percent per year are shown in the legends.

Trends in historical annual global LNO\textsubscript{x} emissions for different scenarios are generally consistent with trends in historical global mean lightning flash rates\textsubscript{LFRs}, as shown in Figs. 4a–5b and Fig. 89. This finding is not surprising because, as the lightning NO\textsubscript{x} emission parameterizations introduced in Sect. 2.2 show, the simulated lightning flash rates\textsubscript{LFRs} are linearly related to the simulated LNO\textsubscript{x} emissions in our study. The results presented in Fig. 89 imply that historical global warming and increases in AeroPEs can affect atmospheric chemistry and can engender feedback by influencing LNO\textsubscript{x} emissions.

3.3 Effects of Pinatubo volcanic eruption effects on historical lightning–LNO\textsubscript{x} trends

We estimate the effects of the Pinatubo eruption effects on historical lightning–LNO\textsubscript{x} trends and variation by comparing the simulation results of STD and Volca-off experiments. The simulated global mean lightning flash rates\textsubscript{LFRs} by STD and Volca-off experiments are the same until April 1991. They then begin to show differences from May 1991 (Figs. 9e–f). The time series of global mean LFRs is not shown.). This result is reasonable because the Pinatubo volcanic perturbations are removed from SAC during June 1991 through May 1994 in the Volca-off experiments by equation (11), and because the SAC of May 1991 used in CHASER are interpolated between the SAC of April 1991 and June 1991.
Figure 9: Time series of lightning flash rate anomaly (LFR) from 1990 to 1999. Figures 9 panels (a–d) show the time series of lightning flash rate anomalies and their smoothed curves by 1-D Gaussian (Denoising) filter from 1990 through 1999. Figures 9 panels (c–d) present the time series of lightning flash rate anomalies from 1991 to 1992. The values shown over the red lines in Figs. 9 panels (c–d) are Relative diff calculated using equation 12.


Figures 9c–d show the time series of lightning flash rate anomalies and Relative diff (values over the red lines) from 1991 to 1992. Relative diff are relative differences of the global mean lightning flash rate anomalies between STD and Volca-off experiments calculated using the following equation.

\[
\text{Relative diff} = 100\% \times \frac{\text{LFR}_{\text{STD}} - \text{LFR}_{\text{Volca-off}}}{\text{LFR}_{\text{Volca-off}}}
\]  

(12)

In the equation, \(\text{LFR}_{\text{STD}}\) represents global mean lightning flash rate anomalies simulated by STD-F1/F2 experiments. \(\text{LFR}_{\text{Volca-off}}\) denotes global mean lightning flash rate anomalies simulated by Volca-off-F1/F2 experiments. \(\text{LFR}_{\text{Volca-off}}\) symbolizes global mean lightning flash rates simulated by Volca-off-F1/F2 experiments.
Figure 10: Probability Density Distributions (PDDs) of $Relative\_diff$ obtained from monthly time-series data of $Relative\_diff$ during 1990–1999. The 95% confidence interval of $Relative\_diff$ is also shown in the titles of this figure.

The monthly time-series data of $Relative\_diff$ for 1990–1999 for both lightning schemes are calculated and Probability Density Distributions (PDDs) of $Relative\_diff$ are displayed in Fig. 10. The $Relative\_diff$ presented in Fig. 10 are all normally distributed as determined by the Kolmogorov–Smirnov test. The 95% confidence interval of $Relative\_diff$ is calculated and shown in the titles of Fig. 10. As displayed in Figs. 9c–10c–10d, the underlined values (underlined $Relative\_diff$) distributed within 1991-08–1992-04 outreached exceeded the 95% confidence interval, which means there are indicating significant differences in the calculated global mean lightning-flash-rateLFR anomalies by STD and Volca-off experiments. In other words, global lightning activities were suppressed significantly by the Pinatubo eruption during the first year after the eruption.
Figure 11a–b: 1991-08 – 1992-04 averaged lightning flash rate LFR anomaly differences (a–b), CAPE anomaly differences (c), $Q_{Ra}$ anomaly differences (d), and differences of 200 hPa – 400 hPa averaged cumulus cloud fraction (e) anomaly between STD-F2 and Volca-off-F2 experiments (e) on the global map.

Figures 11a-b12a–12b show 1991-08 – 1992-04 averaged lightning flash rate LFR anomaly differences between STD and Volca-off experiments on the global map. We found from Figs. 11a-b12a–12b that lightning activities are suppressed
significantly within the three hotspots of lightning activities (Central Africa, Maritime Continent, and South America) during 1991-08 – 1992-04, when the global mean lightning flash rates (LFRs) are found to be suppressed. To explore elucidate the potential reasons for the suppressed global lightning activities during the first year after the Pinatubo eruption, we first investigated the 1991-08 – 1992-04 averaged differences in CAPE and $Q_{Ra}$ differences anomaly between STD-F1 and Volca-off-F1 (Figs. 1c–1d) because lightning densities are computed with CAPE and $Q_{Ra}$ by the ECMWF-McCAUL scheme.

Results showed that the Pinatubo eruption can lead to apparent reductions of CAPE and $Q_{Ra}$ within tropical and subtropical terrestrial regions (typically three hotspots of lightning activities) where lightning occurrence is frequent. These reductions constitute the main reason for the suppressed global lightning activities during the first year after the Pinatubo eruption simulated by the ECMWF-McCAUL scheme. We also examined the 1991-08 – 1992-04 averaged differences of 200 hPa – 400 hPa averaged cumulus cloud fraction anomaly between STD-F2 and Volca-off-F2 on the global map (Fig. 1e). The cumulus cloud fractions of each model layer are used to weight the calculated lightning densities from that layer by the CTH scheme, as explained in Sect. 2.2. As depicted in Fig. 1f and Fig. S6, the Pinatubo eruption led to marked reductions in the middle to upper tropospheric cumulus cloud fractions during 1991-08 – 1992-04 over three hotspots of lightning activities (Central Africa, Maritime Continent, and South America). As displayed in Fig. 5f, the cumulus that reached the middle to upper troposphere is highly related closely to lightning formation. Consequently, the simulated global lightning activities by the CTH scheme were also suppressed considerably during the first year after the Pinatubo eruption.
Figure 4213: 1990–1999 annual global LNO\textsubscript{x} emissions calculated from the STD and Volca-off experiments’ outputs simulated using the ECMWF-McCAUL scheme (a) and the CTH scheme (b). Values over the red lines represent the relative differences (%) between the red lines and blue lines, calculated with respect to the blue lines.

Aside from the previously-described global lightning activity suppression described earlier, the production of LNO\textsubscript{x} might also decrease after the Pinatubo eruption. To explore this conjecture, we compared the LNO\textsubscript{x} emissions in STD and Volca-off experiments (Fig. 4213). In the case of the ECMWF-McCAUL scheme, the reduction of LNO\textsubscript{x} emissions caused by the Pinatubo eruption started in 1991 (4.705\%84\%) and continued until 1993, with the highest percentage reduction occurring in 1992 (8.7247\%) (Fig. 4213a). However, the CTH scheme showed a slightly different scenario of LNO\textsubscript{x} emissions reduction after the Pinatubo eruption. The LNO\textsubscript{x} emissions are almost evenly reduced during 1991–1994 in the case of the CTH scheme (Fig. 4213b). In conclusion, our study indicates that the Pinatubo eruption can engender reductions in global LNO\textsubscript{x} emissions, which last 2–3 years. However, there exists some uncertainty in evaluating the magnitude of the reductions: from 2.411.99\% to 8.7247\% for the annual percentage reduction found from our study\textsuperscript{2}.

The simulated reduced global LNO\textsubscript{x} emissions caused by the Pinatubo eruption might influence atmospheric chemistry significantly. Most importantly, the reduced global LNO\textsubscript{x} emissions might reduce OH radical production and extend the global
mean tropospheric lifetime of methane against tropospheric OH radical (abbreviated as the methane lifetime). We investigated this point further by comparing the methane lifetime anomaly simulated by STD and STD-rVolcaoff experiments. As introduced in Sect. 2.5, the settings of STD-rVolcaoff experiments are the same as those for STD experiments, except that they use the daily LNO\textsubscript{x} emission rates calculated from the Volca-off experiments. We calculated the monthly CH\textsubscript{4} lifetime anomalies during 1990–1999 and \(\Delta \tau_{CH_4}\) (the difference of CH\textsubscript{4} lifetime anomaly between STD and STD-rVolcaoff experiments), which are shown in Figs. 13c–14d. Figures 13a–b14a–14b display the PDDs of \(\Delta \tau_{CH_4}\) monthly time series during 1990–1999. The \(\Delta \tau_{CH_4}\) shown in Figs. 13a–b14a–14b are all normally distributed, as determined by using the Kolmogorov–Smirnov test. The 95\% confidence interval of \(\Delta \tau_{CH_4}\) is calculated and shown in the titles of Figs. 13a–b14a–14b. The annual global LNO\textsubscript{x} production averaged during 1990–1999 is 3.56 TgN yr\textsuperscript{-1} for STD-F1 and 4.79 TgN yr\textsuperscript{-1} for STD-F2. At this level of annual global LNO\textsubscript{x} production, we found that within the first two years after the Pinatubo eruption, the \(\Delta \tau_{CH_4}\) only slightly exceeded the 95\% confidence interval in 1992-02 (0.18 years) simulated by both lighting schemes (1992-02 and 1992-04 in the case of the ECMWF-McCAUL scheme; 1991-12 in the case of the CTH scheme). However, the widely cited range of annual global LNO\textsubscript{x} production is 2–8 TgN yr\textsuperscript{-1} (Schumann and Huntrieser, 2007). Presuming that \(\Delta \tau_{CH_4}\) responds linearly to the LNO\textsubscript{x} emission level, and that the annual global LNO\textsubscript{x} production is 8 TgN yr\textsuperscript{-1}, then the extension of the CH\textsubscript{4} lifetime because of the reduced LNO\textsubscript{x} emissions can reach around 0.45 years for the ECMWF-McCAUL scheme. As a comparison, ultraviolet shielding effects caused by stratospheric aerosols after the Pinatubo eruption led to the maximum increase of the methane lifetime by about 0.6 years (Figs. 13e–d14e–14d).
Panels (a–b) show the Probability Density Distributions (PDDs) of $\Delta \tau_{\text{CH}_4}$ obtained from the monthly time series data of $\Delta \tau_{\text{CH}_4}$ during 1990–1999. $\Delta \tau_{\text{CH}_4}$ represents the difference in CH$_4$ lifetime anomaly between STD and STD-rVolcaoff experiments. The 95% confidence interval of $\Delta \tau_{\text{CH}_4}$ is also presented in the titles of Figs. 13–14. Panels (c–d) show monthly time series of CH$_4$ lifetime anomalies simulated by STD-F1/F2 and STD-rVolcaoff-F1/F2 experiments. Values over the red lines are the difference in CH$_4$ lifetime anomaly between STD and STD-rVolcaoff experiments ($\Delta \tau_{\text{CH}_4}$). Represent $\Delta \tau_{\text{CH}_4}$.

3.4 Model intercomparisons of lightning flash rate LFR trends with CMIP6 model outputs

The historical lightning trends demonstrated in our study are undoubtedly worth comparing with the results of other chemistry–climate models or Earth system models. As introduced in Sect. 2.4, for comparison of the simulated lightning flash rate LFR
trends and variations in our study with those of other CMIP6 models’ outputs, we used all available lightning-flash-rate \( \Delta \text{LFR} \) data from the CMIP6 CMIP Historical experiments from CESM2-WACCM (3 ensembles) (Danabasoglu, 2019), GISS-E2-1-G (9 ensembles) (Kelley et al., 2020), and UKESM1-0-LL (18 ensembles) (Tang et al., 2019). Please refer to Table S1 for a complete list of the ensemble members which were used. It is noteworthy that the lightning-flash-rate \( \Delta \text{LFR} \) data obtained from the three previously described CMIP6 models described earlier are calculated using the CTH scheme. The results of model intercomparisons of lightning-flash-rate \( \Delta \text{LFR} \) trends and variations are displayed in Fig. 4. As illustrated in Figs. 4a–b15a–15b, both the ECMWF-McCAUL and the CTH schemes (STD-F1/F2) simulated almost flat statistically non-significant global lightning trends, but the ensemble mean obtained from another three CMIP6 models exhibit significant increasing global lightning trends (trends from 0.11% yr\(^{-1}\) to 0.25% yr\(^{-1}\)). Many reasons are responsible for underlie the differences in global lightning trends simulated by CHASER in our study and by the three CMIP6 models, including the use of different methods to determine SSTs/sea ice fields. Instead of using a coupled Atmosphere–Ocean general circulation model to calculate SSTs/sea ice fields dynamically in the three CMIP6 models, CHASER uses the prescribed HadISST data (Rayner et al., 2003), which are based on plenty of observational data. Changes in the global mean sea surface temperature anomaly during 1960–2014 (\( \Delta \text{SST} \)) obtained from STD-F1/F2 and CMIP6 model outputs are shown presented in Table 2. We also used the observation-based Extended Reconstructed SST (ERSST) dataset (Huang et al., 2017) constructed by NOAA to evaluate the \( \Delta \text{SST} \) obtained from different models. The \( \Delta \text{SST} \) calculated from ERSST during 1960–2014 is 0.54\( ^\circ \text{C} \), which closely approximates the \( \Delta \text{SST} \) obtained from STD-F1/F2. Considered from the perspective of SSTs/sea ice fields alone, the results (global lightning trends) of our study are expected to be closer to the actual situation.

Actually, the three CMIP6 models simulated stronger global warming during 1960–2014 than CHASER in our study, as displayed in Fig. S6S7. The CTH scheme is reported to respond positively to simulated global warming (Price and Rind, 1994; Zeng et al., 2008; Hui and Hong, 2013; Banerjee et al., 2014; Krause et al., 2014; Clark et al., 2017). The simulated stronger global warming by the three CMIP6 models is regarded as responsible for differences in simulated global lightning trends between our study and the three CMIP6 models (Figs. 4a–b15a–15b). We further investigated the sensitivities of the global mean lightning-flash-rate \( \Delta \text{LFR} \) anomaly change to the global mean surface temperature anomaly increase (% °C\(^{-1}\)) obtained from CHASER and the three CMIP6 models. The sensitivities in percentage per degree Celsius are presented in Table 2. Overall, even when using the same CTH scheme, the sensitivities (\( \Delta \text{LFR}/\Delta \text{TS} \)) simulated by the three CMIP6 models are higher than that simulated by CHASER in our study. This different sensitivity might be partially attributable to the nonlinear relation between lightning response and climate change (Pinto, 2013; Krause et al., 2014). Compared to the CTH scheme, the ECMWF-McCAUL scheme simulated a statistically non-significant negative sensitivity (\( \Delta \text{LFR}/\Delta \text{TS} \)), which is attributable to the stronger suppression of positive global lightning trends caused by increases in AeroPEs simulated by using the ECMWF-McCAUL scheme.
Table 2: Changes in global mean surface temperature anomaly (ΔTS), global mean sea surface temperature anomaly (ΔSST), global mean lightning flash rate anomaly (ΔLFR), and the rate of change of lightning flash rate anomaly (ΔLFR/ΔTS) corresponding to each degree-Celsius increase in global mean surface temperature anomaly (ΔTS) obtained from STD-F1/F2 and CMIP6 model outputs. The change of ΔSST obtained from the ERSST dataset is also shown in this Table. Changes were obtained by calculating the difference between the rightmost and leftmost points of the approximating curve for the 1960–2014 time-series data.

<table>
<thead>
<tr>
<th>Model/experiment/dataset</th>
<th>ΔTS (°C)</th>
<th>ΔSST (°C)</th>
<th>ΔLFR (%)</th>
<th>ΔLFR/ΔTS (% °C⁻¹)</th>
</tr>
</thead>
<tbody>
<tr>
<td>STD-F1</td>
<td>0.593</td>
<td>0.428</td>
<td>-0.272</td>
<td>-0.46</td>
</tr>
<tr>
<td>STD-F2</td>
<td>0.563</td>
<td>0.432</td>
<td>1.497</td>
<td>2.66</td>
</tr>
<tr>
<td>CESM2-WACCM</td>
<td>1.245</td>
<td>1.077</td>
<td>13.758</td>
<td>11.05</td>
</tr>
<tr>
<td>GISS-E2-1-G</td>
<td>0.810</td>
<td>0.677</td>
<td>7.248</td>
<td>8.95</td>
</tr>
<tr>
<td>UKESM1-0-LL</td>
<td>1.141</td>
<td>0.999</td>
<td>5.942</td>
<td>5.21</td>
</tr>
<tr>
<td>ERSST</td>
<td>—</td>
<td>0.549</td>
<td>—</td>
<td>—</td>
</tr>
</tbody>
</table>

Figures 44d–e15d–15e affirm that the global lightning variation simulated by our study is basically within the full ensemble range of GISS-E2-1-G and UKESM1-0-LL. After the Pinatubo eruption, as described in Sect. 3.3 of this report, the GISS-E2-1-G and UKESM1-0-LL models also manifest significant suppression of global lightning activities, but the CESM2-WACCM model shows no such phenomenon. The commonalities and differences in global lightning trends and variations found in the model intercomparisons imply that great uncertainties existed in past (1960–2014) global lightning trend simulations. Such uncertainties deserve to be investigated further.
Figure 1415: Comparisons of simulated global mean lightning flash rate $LFR$ anomalies found in our study (CHASER) and found using other CMIP6 models. All the figures are created based on the monthly time-series data of global mean lightning flash rate $LFR$ anomalies with a 1-D Gaussian (Denoising) filter applied. For CMIP6 models, the ensemble mean is shown as the solid line; and the full ensemble range is shown as grey shading (c–e). Fitting curves and the trends of fitting curves (% yr$^{-1}$) are also given in (a–b).

4 Discussions and Conclusions

We used two lightning schemes (the CTH and ECMWF-McCAUL schemes) to study historical (1960–2014) lightning–LNO$_x$ trends and variations and their controlling/influencing factors (global warming, increases in AeroPES, and Pinatubo eruption) within the CHASER (MIROC) chemistry–climate model. The CTH scheme, which is the most widely used lightning scheme, but nevertheless lacks a direct physical link with the charging mechanism. The ECMWF-McCAUL scheme is a newly developed process-based/ice-based lightning scheme with a direct physical link to the charging mechanism.

With only the aerosol radiative effects considered in the lightning–aerosols interaction, both lightning schemes simulated almost flat trends of global mean lightning flash rate $LFR$ during 1960–2014. Reportedly, because the aerosol microphysical
effects can enhance lightning activities (Yuan et al., 2011; Wang et al., 2018; Liu et al., 2020), our study might underestimate the increasing trend of global mean lightning flash rate \( LFR \) (our study only considered the aerosol radiative effects in aerosol–lightning interactions). Further research is expected considering anticipated, with consideration of the effects of aerosol microphysical effects on long-term lightning trends. Moreover, both lightning schemes manifest that past global warming enhances the historical trend of global mean lightning density toward the positive direction (around 0.03% yr\(^{-1}\) or 3% K\(^{-1}\)). However, past increases in AeroPEs exert the opposite effect to the lightning trend (-0.07% yr\(^{-1}\) – -0.04% yr\(^{-1}\)). The effects of the increased AeroPEs on the lightning trend only over land regions expand to -0.10% yr\(^{-1}\) – -0.05% yr\(^{-1}\), which implies that the effects are more significant over land regions. We obtained similar results for the historical global LNO\(_x\) emissions trend, which indicates that historical global warming and increases in AeroPEs can affect atmospheric chemistry and engender feedback by influencing LNO\(_x\) emissions. Although the CTH and ECMWF-McCAUL schemes use different parameters to simulate lightning, both lightning schemes indicate that the enhanced global convective activity under global warming is the main reason for the increase in lightning–LNO\(_x\) emissions. In contrast, the increases in AeroPEs have decreased lightning–LNO\(_x\) emissions by weakening the convective activity in the hotspots of lightning hotspots. By analyzing the simulation results on the global map, we also found that the effects of historical global warming and increases in AeroPEs on lightning trends are heterogeneous across different regions. Our results indicate that historical global warming enhances lightning activities within the Arctic region and Japan but suppresses lightning activities around New Zealand and some parts of the Southern Ocean. Both lightning schemes demonstrated that the historical increases in AeroPEs suppress lightning activities in some parts of the Southern Ocean and South America. The ECMWF-McCAUL scheme also suggests that historical increases in AeroPEs suppress lightning activities in some parts of India and China—when only the aerosol radiative effects are considered. This finding is plausible because both countries experienced dramatic increases in AeroPEs during 1960–2014 because of rapid economic growth.

Furthermore, this report is the first describing the significant suppression of global lightning activity was suppressed significantly during the first year after the Pinatubo eruption, which is indicated in both lightning schemes (global lightning activities decreased by up to 47.02\(\%\) \(10\%\) simulated by the ECMWF-McCAUL scheme). This finding is mainly attributable to the Pinatubo eruption weakening of the convective activities within the hotspots of lightning, which in turn decreased the amount of column precipitating ice \( Q_{Ra} \) and middle-level to high-level cumulus cloud fractions in these regions. The simulation results also indicate that the Pinatubo eruption can engender reductions in global LNO\(_x\) emissions, which last 2–3 years. However, some uncertainty exists in evaluating magnitude of these reductions (from 2.41\(\%\) to 8.72\(\%\) for the annual percentage reduction in our study). The case study of the Pinatubo eruption in our research indicates that other large-scale volcanic eruptions can also engender significant reduction of global lightning activities and global-scale LNO\(_x\) emissions.

Lastly, we compared the global lightning trends demonstrated in our study with the outputs of three CMIP6 models: CESM2-WACCM, GISS-E2-1-G, and UKESM1-0-LL. We used all available lightning flash rate \( LFR \) data from the CMIP6 CMIP
historical experiments from the three models described above. The three CMIP6 models suggest significant increasing trends in historical global lightning activities, which differs from the findings of our study. Unlike the three CMIP6 models that which use a coupled Atmosphere–Ocean general circulation model to calculate SSTs/sea ice fields dynamically, our study (CHASER) uses the prescribed HadISST SSTs/sea ice data, which are closer to more closely reflect the actual situation. Therefore, we believe that the results (the historical global lightning trends) obtained from our study (CHASER) are closer to more closely approximate the actual situation. However, model intercomparisons of global lightning trends still indicate that significant considerable uncertainties exist in the historical (1960–2014) global lightning trend simulations, and that such uncertainties deserve further investigation.

**Code availability**

The source code for CHASER to reproduce results obtained from this work is obtainable from the repository at https://doi.org/10.5281/zenodo.5835796 (He et al., 2022a).

**Data availability**

The LIS/OTD data used for this study are available from https://ghrc.nsstc.nasa.gov/hydro/?q=LRTS (last access: 11 January 2022). The CMIP6 model outputs (lightning-flash-rate LFR and surface temperature) used for this study are available from https://aims2.llnl.gov/search (last access: 1 February 2023). The Extended Reconstructed SST data used for this study are available from https://www.ncei.noaa.gov/products/extended-reconstructed-sst (last access: 27 March 2023).

**Author contributions**

YFH conducted all simulations, interpreted the results, and wrote the manuscript. KS developed the CHASER (MIROC) model code, conceived the presented idea, and supervised the findings of this work and the manuscript preparation.

**Competing interests**

The authors declare that they have no conflict of interest.

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References


