

# 1 Various lithospheric deformation patterns derived from 2 rheological contrasts between continental terranes: 3 Insights from 2-D numerical simulations

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11 **Abstract.** Continents are formed by the amalgamation of numerous micro-terranes and island arcs, so  
12 they have spatially varying lithosphere strengths. The Crème brûlée (CB) model and the Jelly sandwich  
13 (JS) model have been commonly used to describe continental lithosphere strength-depth variations.  
14 Depending on the strength of continental lower crust, the CB and JS models can be further subdivided  
15 into two subclasses, in which the I subclass (CB-I and JS-I) and II subclass (CB-II and JS-II)  
16 respectively have a strong or weak lower crust. During continental collision, lithosphere deformation is  
17 the byproduct of the comprehensive interaction of multiple terranes. Here we used 2-D  
18 thermo-mechanical numerical models that contain three continental terranes to systematically explore  
19 the effects of terranes with various strengths on continental deformation, and studied the effects of  
20 different rheological assumptions on terrane deformation. We found four types of lithosphere  
21 deformation patterns: collision, subduction, thickening and delamination, and replacement. These  
22 simulation patterns are seen in observed deformation patterns and structures in East Asia, suggesting  
23 they are likely to be naturally occurring modes of intracontinental orogenesis.

## 24 1. Introduction

25 Continents have undergone multiple break-up and assembly events during the past ~2 billion years,  
26 with the assembly events often being associated with the accretion and deformation of numerous  
27 micro-terranes (Mitchell et al., 2021). Accreted terranes have different ages ranging from ~3500 – 3000  
28 Ma to 50 – 0 Ma, and diverse compositions and structures linked to their diverse continental, arc, or

29 oceanic origin, which often leads to them having distinct initial lithospheric thicknesses and strengths  
30 (Artemieva, 2006; Audet and Bürgmann, 2011; Pasyanos et al., 2014; Morgan and Vannucchi, 2022).  
31 The lithosphere of ancient continental terranes like cratons are usually thick and strong, while younger  
32 lithosphere of continental margins and tectonically active regions is thin and weak (Audet and  
33 Bürgmann, 2011; Burov, 2011), and deeply buried former oceanic fragments can have temperature and  
34 strengths that vary over ~0.5 Gyr timescale (Morgan and Vannucchi, 2022).

35 Continental lithosphere strength conventionally been represented by two prevailing rheology models  
36 —the Crème brûlée (CB) and the Jelly sandwich (JS) idealizations (Chen and Molnar, 1983; Jackson,  
37 2002; Burov and Watts, 2006; Bürgmann and Dresen, 2008; Burov, 2011). The Crème brûlée scenario  
38 suggests that lithosphere strength resides entirely in the crust, with the lithospheric mantle being much  
39 weaker (with this strength contrast being the explanation for why little seismicity is typically seen in  
40 the continental mantle, despite rock-mechanics arguments that it should usually be stronger than its  
41 overlying crust). In contrast, the Jelly sandwich model is based on conventional rock mechanics  
42 arguments which imply that in general the continental middle and lower crust should be weaker than  
43 overlying cooler upper crust and underlying further-from-solidus lithospheric mantle (Figure 1a). The  
44 rheology of the continental lower crust can also differ strongly in different continental terranes due to  
45 the varieties in composition, temperature, water content, stress, and tectonic environment (Bürgmann  
46 and Dresen, 2008; Hacker et al., 2015; Morgan and Vannucchi, 2022). Therefore, the CB and JS  
47 conceptualizations can be further subdivided into CB-I and CB-II, JS-I, and JS-II subclasses that reflect  
48 potentially variable strengths of the lower crust: CB-I and JS-I, CB-II and JS-II have strong and weak  
49 continental lower crust, respectively (Fig. 1a). Observations in Eastern Asia show a wide variability in  
50 terrane deformation styles that argue for the potential feasibility of all four of these rheological models  
51 (Figure 1b).

52 Several previous numerical modelling studies have discussed the effects of rheological contrasts  
53 between terranes in lithosphere deformation in a collisional system. Studies containing two terranes  
54 have explored contrasts in crustal rheology, and found that this can greatly change the morphology, size  
55 and deep lithosphere structure of collisional orogenic belt (Chen, 2021; Chen et al., 2017; Cook and  
56 Royden, 2008; Faccenda et al., 2008; Sun and Liu, 2018; Vogt et al., 2018; Xie et al., 2021). Strong  
57 crust also has the potential to protect its underlying lithospheric mantle from deformation and  
58 destruction (Heron and Pysklywec, 2016). Studies containing three or more terranes in their models

59 have usually focused on the middle terrane which can play a crucial role in lithosphere deformation in a  
60 collisional system (Kelly et al., 2016, 2020; Li et al., 2016; Huangfu et al., 2018, 2022; Sun and Liu,  
61 2018; Xie et al., 2023). A weak middle terrane is easy to be thickened, to the point where eventually its  
62 lithospheric mantle can be delaminated from the crust; while a moderate-strength middle terrane can  
63 induce far-field orogenesis; and a strong middle terrane may prevent propagation of deformation and  
64 facilitate underthrusting of the advancing terrane. In addition, some studies have also stressed the  
65 importance of local pre-existing weak zones which can change the order and style of lithosphere  
66 deformation (Chen et al., 2020; Heron et al., 2016; Sokoutis and Willingshofer, 2011; Xie et al., 2021).  
67 Large-scale continental collisional system often involves the multiple units of an indenting terrane, a  
68 middle terrane, and far-end backwall terranes. These terranes have different lithosphere rheologies and  
69 thicknesses, and they collectively contribute to several styles of continental deformation (Artemieva,  
70 2006; Audet and Bürgmann, 2011; Pasyanos et al., 2014; Morgan and Vannucchi, 2022). Here, we use a  
71 2-D thermo-mechanical numerical modeling method to systematically study the effects of terranes with  
72 various rheological properties on continental deformation. Our numerical models simulate a  
73 continent-continent collisional system that contains three continental terranes. They explore the effects  
74 of four groups of lithosphere deformation patterns linked to the four rheological idealizations of CB-I,  
75 CB-II, JS-I, and JS-II applied to each terrane. We will summarize the rheological features for each  
76 deformation pattern, and then apply the simulations to better understand ongoing and past deformation  
77 histories of various orogenic belts in the global, especially in eastern Asia, such as the eastern Tien  
78 Shan orogenic belt, the Tibetan Plateau and the Early Paleozoic Orogen in Southeastern China.

## 79 **2. Numerical modelling method and model setup**

### 80 **2.1. Numerical modelling method**

81 Our thermo-mechanical models were performed with the I2VIS code of Gerya and Yuen (2003),  
82 previously used in Xie et al. (2021, 2023). This code combines finite differences with marker-in-cell  
83 techniques to solve the mass, momentum, and energy conservation equations for incompressible flow.  
84 It incorporates the non-Newtonian visco-plastic rheologies for the lithosphere, as well as the possibility  
85 to include parameterizations of the effects of surface processes like sedimentation and erosion.

86 **2.1.1. Governing equations**

87 The mass conservation equation for incompressible flow is:

$$88 \quad \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} = 0, \quad (1)$$

89 The momentum conservation equations (Stokes equations) are:

$$90 \quad \begin{aligned} \frac{\partial \sigma'_{xx}}{\partial x} + \frac{\partial \sigma'_{xy}}{\partial y} &= \frac{\partial P}{\partial x}, \\ \frac{\partial \sigma'_{yy}}{\partial y} + \frac{\partial \sigma'_{xy}}{\partial x} &= \frac{\partial P}{\partial y} - g\rho \end{aligned} \quad (2)$$

91 The energy (heat) conservation equation is:

$$92 \quad \begin{aligned} \rho C_p \frac{DT}{Dt} &= -\frac{\partial q_x}{\partial x} - \frac{\partial q_y}{\partial y} + H_r + H_s + H_a + H_L \\ q_x &= -k \frac{\partial T}{\partial x} \\ q_y &= -k \frac{\partial T}{\partial y} \\ H_a &= T\alpha \frac{DP}{Dt} \\ H_s &= \sigma'_{xx} \dot{\epsilon}_{xx} + \sigma'_{yy} \dot{\epsilon}_{yy} + 2\sigma'_{xy} \dot{\epsilon}_{xy} \end{aligned} \quad (3)$$

93 where  $x$  and  $y$  represent the horizontal and vertical coordinate directions, and  $v_x$  and  $v_y$  are the  
 94 corresponding velocity components, respectively.  $\sigma'_{ij}$  and  $\dot{\epsilon}_{ij}$  ( $i, j = x, y$ ) are deviatoric stress and  
 95 strain-rate tensors, respectively;  $g$  is the gravitational acceleration;  $\rho$  is density. In the heat conservation  
 96 equation,  $q_x$  and  $q_y$  are the horizontal and vertical components of the heat flux, respectively;  $C_p$  is heat  
 97 capacity, and  $H_r$ ,  $H_a$ ,  $H_s$ , and  $H_L$  denote the radioactive, adiabatic, shear, and latent heat production,  
 98 respectively;  $k$  is the thermal conductivity.

99 The rheological constitutive relationship connects the deviatoric stress and strain rate:

$$100 \quad \begin{aligned} \sigma'_{xx} &= 2\eta_{eff} \dot{\epsilon}_{xx}, \quad \dot{\epsilon}_{xx} = \frac{\partial v_x}{\partial x} \\ \sigma'_{xy} &= 2\eta_{eff} \dot{\epsilon}_{xy}, \quad \dot{\epsilon}_{xy} = \frac{1}{2} \left( \frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right), \\ \sigma'_{yy} &= 2\eta_{eff} \dot{\epsilon}_{yy}, \quad \dot{\epsilon}_{yy} = \frac{\partial v_y}{\partial y} \end{aligned} \quad (4)$$

101 where  $\eta_{eff}$  is the effective viscosity.

## 102 2.1.2. Rheology

103 Here we make the conventional assumption that the crust and mantle have a visco-plastic rheology.

104 Viscous deformation is determined as a combination of diffusion and dislocation creep that depends on

105 temperature, pressure, and strain rate, expressed as (Gerya, 2019):

$$\begin{aligned} \eta_{dist} &= \frac{1}{2} \frac{1}{(A_D)^{-1/n} (\dot{\epsilon}_{II})^{(n-1)/n}} \exp\left(\frac{E_a + V_a P}{nRT}\right) * S \\ \eta_{diff} &= \frac{1}{2} \frac{A_D}{\sigma_{cr}^{(n-1)}} \exp\left(\frac{E_a + V_a P}{RT}\right) * S \end{aligned} \quad (5)$$

107 For mineral aggregates, both dislocation and diffusion creep occur simultaneously, with a combined

108 effective viscosity given by:

$$\frac{1}{\eta_{ductile}} = \frac{1}{\eta_{dist}} + \frac{1}{\eta_{diff}}, \quad (6)$$

110 where  $\eta_{dist}$  and  $\eta_{diff}$  are viscosities for dislocation and diffusion creep, respectively.  $\sigma_{cr}$  is the

111 critical stress for the dislocation to diffusion stress transition, and the parameters  $A_D$ ,  $E_a$ ,  $V_a$ , and  $n$  are a

112 material constant, activation energy, activation volume, and stress exponent, respectively, and  $R$  is the

113 universal gas constant. The strength scaling factor,  $S$ , is introduced as a simple parameter to vary the

114 lithospheric viscosity.

115 Plasticity is implemented using a conventional pseudo-viscous yield criterion first used to study rifting

116 (e.g. Chen and Morgan, 1990) that is extended to include a strain-weakening-like parameterization of

117 fracture-related strain weakening (Gerya et al., 2010; Vogt et al., 2017):

$$\eta_{plastic} = \frac{\sigma_{yield}}{2\dot{\epsilon}_{II}}$$

$$\sigma_{yield} = C + P\phi$$

$$118 \quad C = \begin{cases} C_a + (C_b - C_a) \times \frac{\gamma}{\gamma_{cr}}, & \text{if } \gamma \leq \gamma_{cr}, \\ C_b, & \text{if } \gamma \geq \gamma_{cr} \end{cases}, \quad (7)$$

$$\phi = \begin{cases} \phi_a + (\phi_b - \phi_a) \times \frac{\gamma}{\gamma_{cr}}, & \text{if } \gamma \leq \gamma_{cr} \\ \phi_b, & \text{if } \gamma \geq \gamma_{cr} \end{cases}$$

119 where  $\sigma_{yield}$  is yield stress,  $P$  is dynamic pressure,  $\gamma$  is the integrated plastic strain, and  $\gamma_{cr}$  is the  
 120 upper strain limit for fracture-related weakening.  $C$  and  $\phi$  are cohesion and friction angle that depend  
 121 on the plastic value.  $C_a$  and  $\phi_a$  are the initial and  $C_b$  and  $\phi_b$  are final strength values, respectively.  
 122 This involves making the rheological assumption that deeply percolating fluids and high pore fluid  
 123 pressures can significantly lower the plastic strength of fractured rocks.  
 124 The final effective viscosity is determined by the minimum value between the ductile and plastic  
 125 viscosities (Ranalli, 1995):

$$126 \quad \eta_{eff} = \min(\eta_{ductile}, \eta_{plastic}). \quad (8)$$

### 127 2.1.3. Surface processes

128 Topography in our models evolves according to a transport equation that is solved at each time step,  
 129 with a crude local parameterization of effects of accounts for sedimentation and erosion:

$$130 \quad \frac{\partial y_{es}}{\partial t} = v_y - v_x \frac{\partial y_{es}}{\partial x} - v_s + v_e. \quad (9)$$

131 Where  $y_{es}$  is the vertical position of the surface as a function of horizontal distance  $x$ ; and  $v_x$  and  $v_y$  are  
 132 the corresponding velocity components, respectively.  $v_s$  and  $v_e$  are the sedimentation and erosion rates,  
 133 respectively, conforming to the relation:

$$134 \quad v_s = 0 \text{ mm/yr}, v_e = 0.3 \text{ mm/yr when } y_{es} > 5 \text{ km};$$

$$135 \quad v_s = 0.3 \text{ mm/yr}, v_e = 0 \text{ mm/yr when } y_{es} < 5 \text{ km}.$$

## 136 2.2. Model Setup

137 The 2-D numerical model covers a rectangular computational domain of 3000 km  $\times$  700 km and

138 consists of  $1360 \times 400$  non-uniform grid cells with dozens of mobile markers in each grid cell to  
139 transport physical properties (Figure 2a). Above 300 km, the cell-size of the grid in the middle of  
140 model ( $X = 1300 - 2200$  km) is  $1 \text{ km} \times 1 \text{ km}$ , and gradually widens towards the two sides to finally  
141 become  $5 \text{ km} \times 1 \text{ km}$ . From the 300 km depth to the model bottom, each grid is stretched to 5 km in the  
142 vertical direction. As a result, the grid in the middle of the model ( $X = 1300 - 2200$  km) is  $1 \text{ km} \times 5 \text{ km}$   
143 and is  $5 \text{ km} \times 5 \text{ km}$  in the other regions. Changing resolutions in different model regions can ensure the  
144 model can finely depict lithosphere deformation in the region of interest while improving the  
145 calculation's efficiency.

146 In the initial configuration, the model comprises three continental terranes — the Pro-, Mid- and  
147 Retro-terrane — which refer to the indenting 'Pro-' terrane driven by plate convergence, an  
148 intermediate 'Mid-' terrane, and a far-end backwall 'Retro-' terrane, respectively (Figure 2a). For the  
149 purpose of simplification, the three terranes are assumed to have the same initial crustal structure with  
150 20 km thick upper and lower crust, respectively. In the meanwhile, to simulate lateral structure  
151 differences within continental lithosphere (Pasyanos et al., 2014), thicknesses of the initial lithospheric  
152 mantle of the Pro-, Mid- and Retro-terrane are 160 km, 90 km, and 120 km, respectively. The rest of  
153 the region is filled by asthenosphere except along the model top, where a 20 km thick layer of "sticky  
154 air" with low viscosity ( $1 \times 10^{18}$  Pas) and low density ( $1 \text{ kg/m}^3$ ) is placed to simulate the effects of a  
155 free surface (Schmeling et al., 2008). Flow laws and material properties for each lithospheric layer are  
156 listed in Table 1.

157 Mechanical boundary conditions of the model are that the top and sides are free-slip boundaries which  
158 mean that the vertical velocity at the top boundary and horizontal velocity at the side boundaries are all  
159 zero. The bottom is assumed to be a somewhat non-physical 'permeable boundary' that was developed  
160 to reduce the required depth of the computational region (Burg and Gerya, 2005). For top-driven flows  
161 like those considered here, this approximation has been shown to not affect deformation in the upper  
162 parts of the region (Burg and Gerya, 2005). Finally, a constant convergence rate of 20 mm/yr is  
163 assigned to the Pro-terrane ( $X = 1000$  km) to drive the model.

164 Initial temperature conditions are set as follows: the model top is set to  $0^\circ\text{C}$ , the two side boundaries  
165 are adiabatic boundaries with zero horizontal heat fluxes, and the model bottom has an initial  
166 temperature of  $1593^\circ\text{C}$ , and can dynamically adjust as the model evolves. The initial thermal gradient  
167 in the crust is  $15^\circ\text{C/km}$  in the three terranes, so their Moho temperature is  $600^\circ\text{C}$ . A temperature of

168 1330 °C is applied at the bottom of the lithospheric mantle of the three terranes, which leads to the Pro-  
169 and Mid-terrane having minimum and maximal thermal gradients in the lithospheric mantle,  
170 respectively (see the right plane in Figure 2a). An adiabatic thermal gradient of 0.5 °C/km is assumed  
171 within the asthenosphere. The temperature field would evolve over time, thus, although the three  
172 terranes are not in thermal equilibrium at the start of the experiments, it has few effects on model  
173 evolution. The initial setup of lithosphere structure and temperature field make the Mid-terrane weakest  
174 when same rheology model is used for the three terranes.

### 175 **3. Simulation Results**

176 The rheological models of CB-I, CB-II, JS-I, and JS-II result from different strength scaling factors for  
177 the upper crust, the lower crust, and the lithospheric mantle in our numerical models (Figure 2b). We  
178 systematically test the effects of these rheological assumptions on the deformation of the Pro-, Mid-  
179 and Retro-terranes. According to different behaviors of lithosphere deformation, these simulation  
180 results can be categorized into four basic modes of collision, subduction, thickening and delamination,  
181 and replacement (Figures S1 and S2). In the deformation mode of collision, the lithospheric mantle of  
182 the Mid-terrane is extruded out and the lithospheric mantles of the two bounding terranes meet and  
183 collide together. In the deformation mode of subduction, the lithospheric mantle of one of the bounding  
184 terranes subducts into the deep mantle below the Mid-terrane while the other one keeps almost  
185 undeformed. In the deformation mode of thickening and delamination, one of the bounding terranes is  
186 shortened by compression, and delamination may come on the heels of thickening of lithosphere due to  
187 gravitational instability in some cases. In the deformation mode of replacement, the bottom of weak  
188 and thick lithospheric mantle of the bounding terrane is scraped off by the strong lithospheric mantle of  
189 the Mid-terrane, and replaced by the latter. Here, we select a typical case for each mode of lithosphere  
190 deformation to describe more details of these modes of model evolution.

#### 191 **3.1. Case 1: Lithosphere Collision**

192 Case 1 represents the scenario of lithosphere collision (Figure 3). In this model, the assumed  
193 rheological models for the Pro-, Mid- and Retro-terrane are JS-I, JS-II, and JS-I, respectively, which  
194 means that the Mid-terrane has a significantly weaker lower crust relative to the Pro- and  
195 Retro-terranes. The lithospheric mantle of the Mid-terrane is also slightly weaker due to its thinner



196 lithosphere and correspondingly higher initial temperature field. The lithosphere strength profiles of the  
197 three terranes are shown in Figure 3g.

198 The Mid-terrane is the first to deform when the Pro-terrane begins to collide, absorbing plate  
199 convergence in the form of lithosphere thickening (Figures 3a and 3d). The upper crust of the  
200 Mid-terrane breaks due to strain weakening, and several reverse faults are formed to absorb crustal  
201 shortening. The lower crust folds, and strain diffusely distributes within it. Since the Retro-terrane is  
202 relatively strong, it prevents crustal deformation from propagating into this terrane, and restricts the  
203 bulk of deformation to the Mid-terrane. With continuous advance of the Pro-terrane and resistance of  
204 the Retro-terrane, the crust of the Mid-terrane is intensively shortened, leading to more thrusting  
205 structures in the upper crust and a “flower-like” structure in the lower crust (Figures 3b and 3e). Thrust  
206 structures and crustal deformation also expand toward the Pro- and Retro-terranes at this stage.  
207 Topography also grows towards the two bounding terranes (Figure 7a). Ultimately, the weak  
208 lithospheric mantle of the Mid-terrane is squeezed out, and the Pro- and Retro-terrane's lithospheric  
209 mantles meet and so start to collide beneath the overlying crust of the Mid-terrane (Figures 3c and 3f).

### 210 **3.2. Case 2: Lithosphere Subduction**

211 Case 2 shows lithosphere subduction of the Pro-terrane (Figure 4). In this model, the assumed  
212 rheological models for the Pro-, Mid- and Retro-terrane are JS-II, JS-I, and JS-I, respectively. The  
213 Mid-terrane has a stronger lower crust compared with the Pro-terrane, but its lithospheric mantle is a  
214 little weaker than the Pro-terrane due to higher temperature field resulting from its thinner lithosphere  
215 structure (Figure 4g). When convergence begins, the weak lower crust of the Pro-terrane is blocked by  
216 the stronger lower crust of the Mid-terrane. This induces it to stack in a collisional front to form a  
217 remarkable folding structure (Figures 4a and 4d). The strong lithospheric mantle of the Pro-terrane  
218 continues to move forward and underthrusts beneath the Mid-terrane. As the Pro-terrane advances, its  
219 crust gradually enters the Mid-terrane, inducing shortening and thickening of the upper crust of the  
220 Mid-terrane, while the strong lower crust of the Mid-terrane almost keeps undeformed (Figures 4b and  
221 4e).

222 Meanwhile, the lithospheric mantle of the Pro-terrane continues to underthrust scraping off part of the  
223 lithospheric mantle of the Mid-terrane. Eventually, the crust of the Pro-terrane wedges a long distance  
224 into the Mid-terrane, and the lithospheric mantle of the Pro-terrane subducts into the deeper mantle

225 (Figures 4d and 4f). In this example, crustal deformation and topography gradually propagate from the  
226 Pro-terrane to the Mid-terrane, whereas the Retro-terrane remains nominally 'undeformed' at all times  
227 (Figure 7b). In some experiments, the lithospheric mantle of the Retro-terrane can subduct beneath the  
228 Mid-terrane (Figure S1).

### 229 **3.3. Case 3: Lithosphere Thickening and Delamination**

230 Case 3 illustrates the thickening and delamination of the lithospheric mantle of the Pro-terrane (Figure  
231 5). In this case, the rheological models for the Pro-, Mid- and Retro-terrane are CB-II, JS-I, and JS-I,  
232 respectively (Figure 5j). The Pro-terrane has a rheologically weaker lower crust and lithospheric mantle,  
233 making it relatively easy to deform once the collision has started. The lithosphere of the Pro-terrane is  
234 first thickened, and the crust starts to form folding in two discrete zones (Figures 5a and 5f). The lower  
235 part of the thickened lithospheric mantle is denser than its ambient mantle owing to lower temperature,  
236 which causes it to drip downwards (Figures 5b–5h). After delamination of the thickened lithosphere,  
237 subduction of the Pro-terrane's lithospheric mantle along one of the deformation localization zones  
238 absorbs the plate convergence (Figures 5e and 5i). Crustal deformation is restricted in the Pro-terrane  
239 until lithosphere delamination, after which crustal strain and topography rapidly spread from the  
240 Pro-terrane to the Mid-terrane (Figure 7c). Like case 2, the Retro-terrane stays essentially undeformed  
241 at all times.

### 242 **3.4. Case 4: Lithosphere Replacement**

243 Case 4 illustrates the lithospheric mantle of the Pro-terrane is replaced by that of a neighboring stronger  
244 Mid-terrane (Figure 6). In this case, the rheological models for the Pro-, Mid- and Retro-terrane are  
245 CB-I, JS-II, and JS-I, respectively. The Pro-terrane has a strong lower crust and a thick and weak  
246 lithospheric mantle, while the Mid-terrane has a weaker lower crust and a strong lithospheric mantle  
247 (Figure 6g). This lithosphere configuration between the Pro- and the Mid-terrane causes deformation to  
248 be primarily distributed in the Pro-terrane's lithospheric mantle and the Mid-terrane's crust. As a result,  
249 the Mid-terrane's crust becomes intensely shortened by fold and thrust structures, but its strong  
250 lithospheric mantle wedges into the Pro-terrane's thick and weak lithospheric mantle (Figures 6a, 6b,  
251 6d and 6e). The strong lithospheric mantle of the Mid-terrane scrapes off the lower part of the weak  
252 lithospheric mantle of the Pro-terrane and so replaces it (Figures 6c and 6f). Similar to case 1, crustal

253 deformation and topography expand from the Mid-terrane towards its side terranes (Figure 7d). The  
254 lithospheric mantle of the Retro-terrane can also be replaced in some cases (e.g. Figure S1).

## 255 **4. Discussion**

### 256 **4.1. Rheological Characteristics for Distinct Lithosphere Deformation Patterns**

257 Distinct lithosphere deformation patterns in our simulations arise from rheological contrasts between  
258 neighboring continental terranes. Figure 8 summarizes the rheological characteristics of these distinct  
259 deformation patterns. When the Mid-terrane's lithospheric mantle is weakest (typified by models in  
260 which the rheological model of the Mid-terrane is CB-II), it is easy for its mantle to be extruded out,  
261 leading to collision between the lithospheric mantles of its surrounding Pro- and Retro-terranes. When  
262 one of the two bounding terranes has extremely weak lithospheric mantle, its lithosphere is first to be  
263 thickened by compression, and delamination may follow due to density-driven instability. When the  
264 lower crust of the Mid-terrane is relatively strong (CB-I or JS-I), while the lower crust is weaker in the  
265 Pro- or Retro-terrane (CB-II or JS-II), the lithospheric mantle of the Pro- or Retro-terrane will tend to  
266 subduct into the deep mantle, e.g. leading to intracontinental subduction. Finally, when the Mid-terrane  
267 has a weak lower crust and strong lithospheric mantle (JS-II), while the Pro- or Retro-terrane has a  
268 strong lower crust and weak lithospheric mantle (CB-I), the lithospheric mantle of the former may  
269 replace the lithospheric mantle of the latter.

270 When the deformation patterns involve the collision and replacement of lithosphere, continental  
271 deformation involves all three terranes (Figures 3 and 6). In contrast, the other deformation patterns  
272 only involve two terranes, the Pro- or Retro-terrane and the Mid-terrane (Figures 4 and 5). The  
273 rheological properties of the Mid-terrane are responsible for these differences. Like previous numerical  
274 studies (Kelly et al., 2016, 2020; Li et al., 2016; Huangfu et al., 2018, 2022; Sun and Liu, 2018), our  
275 simulations show that a weak Mid-terrane is easier to deform, and that in this case lithosphere  
276 deformation will expand from center to its neighboring sides; while a relatively strong Mid-terrane  
277 prevents deformation from propagating far, so that lithosphere deformation is constrained to occur  
278 within two terranes..

279 Although our multi-terrane numerical models mainly focus on the impact of the lateral strength  
280 differences between different terranes in a continental collisional system, rheological models of CB-I,

281 CB-II, JS-I, JS-II also involve vertical rheological variation (Figure 1a). It seems difficult to summarize  
282 how vertical strength variation affects lithosphere deformation of the continental collisional system.  
283 For example, in some cases, only changing the rheological models of the Pro- or Retro-terrane may  
284 produce distinct deformation modes such as collision, subduction, thickening and delamination, and  
285 displacement (e.g., the first and third rows, the third column in the upper left panel in Figure 8 and the  
286 third column in the lower right panel in Figure 8). However, changing the rheological models of the  
287 Pro- or Retro-terrane seems to have less impact on the deformation mode of the continental collisional  
288 system, according to the simulation results of models which are connected by several cross-shaped  
289 solid lines with different colors in Figure 8. Thus, it is difficult to determine whether the horizontal  
290 strength contrasts between terranes or the vertical strength variation of a single terrane plays the  
291 dominant role in a multi-terrane collisional system. This is also the significance and necessity of our  
292 study.

#### 293 **4.2. Influences of Lithosphere Structure**

294 Lithospheric thickness is one of the critical factors that control its strength (Burov, 2011), and it can  
295 strongly vary between tectonic regions (Pasyanos et al., 2014). In our models, we assume different  
296 lithospheric thicknesses for the Pro-, Mid- and Retro-terrane to explore these effects. Complex effects  
297 are seen. When changing the lithospheric thicknesses of the Mid-terrane, or of all three terranes,  
298 remarkable variations in lithosphere deformation appear in cases 1 and 2, but smaller variations are  
299 seen for cases 3 and 4 (Figure 9). Cases 1 and 2 assume a Jelly sandwich rheology for the Pro-, Mid-  
300 and Retro-terrane, so the strength of the lithospheric mantle of three terranes is comparable. Strength  
301 variations produced by differences in lithospheric thickness may alter the relative strength of the three  
302 terranes, resulting in distinct lithospheric deformations. For example, if the Pro- and Mid-terrane have  
303 same lithosphere thickness, deformation mode in Case 2 would change from subduction to thickening  
304 (subplot 3 vs. subplot 8 in Figure 9); if the Mid-terrane is thickest or the three terranes have same  
305 thickness of lithosphere, deformation mode in Case 1 would change from collision to replacement  
306 (subplot 2 vs. subplot 17 and 22 in Figure 9), and the polarity of the subduction of Pro-terrane's  
307 lithospheric mantle would be reversed in Case 2 (subplot 3 vs. subplot 18 and 23 in Figure 9). Instead,  
308 in cases 3 and 4, the Pro-, Mid- and Retro-terrane have two regions with stronger Jelly-sandwich-like  
309 rheological structures and one with a weaker Crème brûlée structure, and deformation preferentially

310 concentrates in the weaker terrane. In comparison to the large strength difference implied for the  
311 lithospheric mantle between the Crème brûlée and Jelly Sandwich rheological models, the strength  
312 variations associated with the differences in lithosphere thickness are relatively small. Therefore,  
313 changing the thicknesses of the lithosphere has much smaller effects on the lithosphere deformation, as  
314 seen in cases 3 and 4 (also see the subplots in 3<sup>rd</sup> and 4<sup>rd</sup> rows of Figure 9).

315 In addition, the weak zones that suture two terranes are generally preserved during continental  
316 amalgamation (Burker et al., 1977; Vink et al., 1984; Yin and Harrison, 2000). These local pre-existing  
317 lithosphere weaknesses would be preferentially activated if the continental lithosphere were subjected  
318 to compression, and could play a key role in concentrating deformation, adjusting deformation  
319 sequences, and inducing lithosphere subduction (Sokoutis and Willingshofer, 2011; Heron et al., 2016;  
320 Chen et al., 2020; Xie et al., 2021, 2023). Comparing the simulation results of models with and without  
321 weak zone, we find that a weak zone will facilitate lithosphere subduction in earlier stages of model  
322 evolution, resulting in more diverse lithosphere deformation patterns during the later stage (Figure 10).

### 323 **4.3. Implications for the Tectonics of East Asia**

#### 324 **4.3.1. Lithosphere Collision beneath the Eastern Tien Shan**

325 The eastern Tien Shan is an ideal region to study the deformation patterns linked to long-term  
326 lithosphere collision (Figure 11a). The eastern Tien Shan is bounded by the Tarim Basin to the south,  
327 and the Junggar Basin to the north. It is composed of a series of former island arcs and small  
328 continental blocks that amalgamated during the late Paleozoic (Han and Zhao, 2017). The lithosphere  
329 of the eastern Tien Shan is weaker and thinner in comparison to its neighboring Tarim Basin and  
330 Junggar Basin (Kumar et al., 2005; Lei and Zhao et al., 2007; Zhang et al., 2013; Deng and Tesauro,  
331 2016). At ~20 – 25 Ma, the eastern Tien Shan became a reactivated orogeny in response to ongoing  
332 India-Asia collision (Yin et al., 1998). Compression linked to the India-Asia collision induced the  
333 Tarim lithosphere to underthrust northward (Xu et al., 2002; Guo et al., 2006; Lei and Zhao et al., 2007;  
334 Lü et al., 2019; Hapaer et al., 2022; Sun et al., 2022). In the northern part of the eastern Tien Shan,  
335 significant high-velocity anomalies and Moho overlap are also imaged, which are conventionally  
336 explained as being due to the southward underthrusting of the Junggar lithosphere (Xu et al., 2002;  
337 Guo et al., 2006; Li et al., 2016; Lü et al., 2019). High-velocity anomalies in the Tarim and Junggar  
338 lithosphere appear to connect beneath the eastern Tien Shan, suggesting the lithosphere of the Tarim

339 and Junggar Basins has converged and collided together in this region (Figure 11b; Lü et al., 2019).  
340 Bidirectional underthrusting of the Tarim and Junggar lithosphere leads to intense crustal shortening  
341 and thrust faults on both flanks over the adjacent basins, as well as attendant fold and reverse fault  
342 zones along the range fronts (Yin et al., 1998; Wang et al., 2004).

#### 343 **4.3.2. Lithosphere Thickening and Delamination in the Tibetan Plateau**

344 The deformation pattern arising from lithosphere thickening and delamination has been applied to the  
345 Tibetan Plateau (Figure 11c). Tibetan lithosphere may have been significantly weakened by hydration,  
346 metasomatism, and partial melting of the lithospheric mantle during a series of oceanic closure and  
347 terrane accretion events before the India-Asia collision (Yin and Harrison, 2000; Zhang et al., 2014;  
348 Ma et al., 2021). It was then pushed northward by the Indian craton and was blocked by the  
349 Tarim/Qaidam craton during India-Asia collision, leading to double crustal thickness (Zhao and  
350 Morgan, 1985; Zhang et al., 2011). The lithosphere beneath the Tibetan Plateau does not thicken  
351 significantly like its crust, especially beneath northern Tibet (Owens and Zandt, 1997; Tunini et al.,  
352 2016). Numerous observations instead suggest that the Tibetan lithosphere has been detached from the  
353 crust and has sunk into deeper mantle, consistent with the presence of high-velocity regions in the deep  
354 mantle in western, southern and southeastern Tibet (Li et al., 2008; Chen et al., 2017; Feng et al., 2021).  
355 A significant depression of the 660-km discontinuity beneath the Himalaya terrane and the uplift of  
356 410-km discontinuity in western Tibet have also attributed to the presence of delaminated Tibetan  
357 lithosphere (Wu et al., 2022). In northern Tibet, anomalously high temperatures are assumed to be  
358 linked to a region of inefficient  $S_n$  propagation indicating a thin or absent lithospheric mantle lid in this  
359 region, while a remarkable low-velocity zone in the mantle and ultra-potassic volcanics also suggest  
360 lithosphere thinning (Barazangi and Ni, 1982; Turner et al., 1996; Owens and Zandt, 1997; Guo et al.,  
361 2006; Liang et al., 2012; Tunini et al., 2016). After lithosphere thinning commenced in the Miocene,  
362 the Tibetan Plateau rapidly grew outwards (Lu et al., 2018 and references therein; Molnar et al., 1993;  
363 Xie et al., 2023).

#### 364 **4.3.3. Lithosphere Subduction in Southeastern China**

365 An example of intracontinental subduction is the Early Paleozoic Orogen in Southeastern China which  
366 appears to have not been preceded by oceanic subduction (Figure 11d; Faure et al., 2009). The Early

367 Paleozoic Orogen of Southeastern China is located on the Wuyi-Yunkai Fold Belt. Arguments against it  
368 being a collisional orogenic belt are its lack of preserved ophiolites, a magmatic arc, subduction  
369 complexes, and high-pressure metamorphism. Instead, structural, metamorphic, and sedimentary  
370 elements indicate that this orogen was an intracontinental orogen controlled by the northward  
371 subduction of Cathaysia (Faure et al., 2009). A weak suture/failure zone inherited from previous  
372 tectonic events contributed to the internal subduction of Cathaysia, during which ductile decollements  
373 accommodated horizontal shortening by folding and thrusting. The tectonic development of this orogen  
374 appears similar to the deformation mode of lithosphere subduction (Figure 4).

375 So far, we have yet to find a suitable region to apply the model deformation pattern of lithosphere  
376 replacement. In this deformation pattern, crustal deformation and topographic evolution are similar to  
377 those in the deformation pattern of lithosphere collision (Figures 7a and 7c). Thus, it is not easy to  
378 identify this pattern by geological and geophysical techniques when the replaced and original  
379 continental lithosphere has similar properties. Improved imaging observations with better resolution  
380 may allow this deformation pattern to be identified in the future.

## 381 **5. Model Limitations**

382 Although we can obtain four deformation modes of continental lithosphere by changing the rheologies  
383 of different terranes in a collisional model, we must keep in mind that our results are based on some  
384 simplifications and assumptions, which may affect the model results. For example, in our model three  
385 terranes are directly collaged together, but in nature different terranes are often connected through weak  
386 sutures which may preferentially deform when they are subjected to compression (Burker et al., 1977;  
387 Yin and Harrison, 2000). These local pre-existing weak zones have non-negligible influences on  
388 lithospheric deformation, and their role were widely discussed in previous studies (Sokoutis and  
389 Willingshofer, 2011; Heron et al., 2016; Chen et al., 2020 ; Xie et al., 2021, 2023). We also discussed  
390 the effects of local pre-existing weak zones in Section 4.2. In addition, lithosphere thicknesses of the  
391 Pro-, Mid- and Retro-terranes are chose arbitrary in our models, but they also has important influences  
392 on lithosphere deformation (Figure 9). Some studies believe that differences in crustal strength will  
393 also cause different lithospheric deformation (Faccenda et al., 2008; Vogt et al., 2017, 2018), but the  
394 three terranes are set same crustal structure in our model for the aim of simplification. As well, some

395 studies believe that the convergence rate will greatly affect the deformation of orogenic belts (Chen et  
396 al., 2016; Vogt et al., 2017), but in this study, the impact of the convergence rate almost can be ignored.

## 397 **6. Conclusions**

398 The continental lithosphere is likely to have strong lateral variations in its strength. We explored 2-D  
399 numerical models that contain three diverse types of continental terranes to study the responses of  
400 continental terranes with different strengths to compression. Four rheological models were respectively  
401 applied to each of the Pro-, Mid- and Retro-terranes, and simulation results can be grouped into four  
402 distinct deformation styles: lithosphere collision, subduction, thickening and delamination, and  
403 replacement. These deformation styles arise from the rheological contrasts between the terranes: (1)  
404 when the middle terrane is the weakest, its lithosphere is easily extruded, which leads to lithosphere  
405 collision between its two bounding terranes; (2) when the middle terrane has a strong lower crust, while  
406 the lower crust of a bounding terrane is weak, then subduction of the lithosphere of the bounding  
407 terrane will occur; (3) when a bounding terrane is the weakest, its lithosphere would tend to be  
408 thickened by lateral compression, followed by lithosphere delamination due to the resulting  
409 density/gravitational instability; (4) when a bounding terrane has a strong lower crust and weak  
410 lithospheric mantle, while the middle terrane has a weak lower crust and strong lithospheric mantle,  
411 then lithosphere replacement will occur. These simulation patterns are seen in observed deformation  
412 patterns and structures in the eastern Tien Shan, and the Tibetan Plateau, the Early Paleozoic Orogen of  
413 Southeastern China, suggesting they are likely to be naturally occurring modes of intracontinental  
414 orogenesis.

### 415 *Code availability*

416 Requests for the numerical code I2VIS should be sent to the main developer  
417 (taras.gerya@erdw.ethz.ch).

### 418 *Data availability*

419 Numerical modeling data and the model evolution animations of Cases 1 – 4 are all provided in  
420 Zenodo (<https://doi.org/10.5281/zenodo.8354366> and <https://doi.org/10.5281/zenodo.10731981>).



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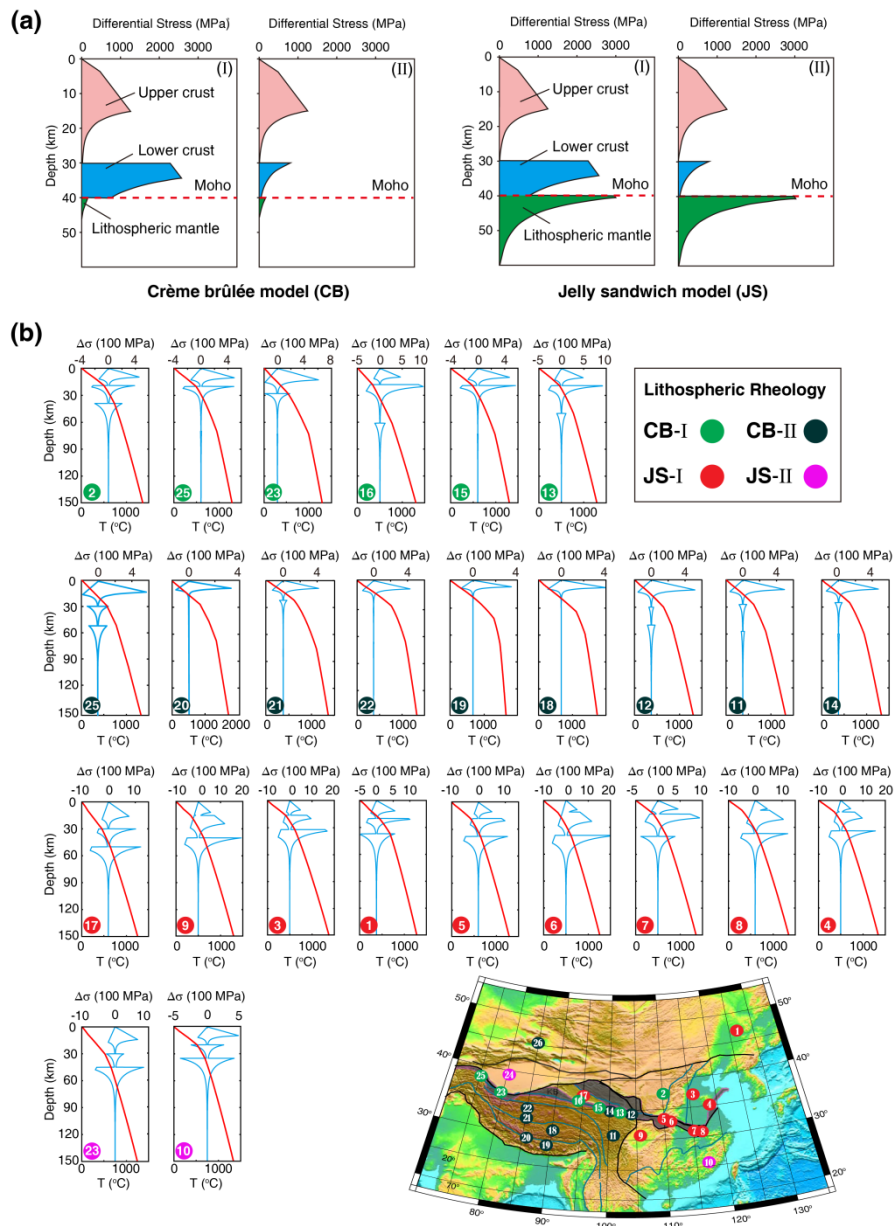
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645 **Figure 1. Four rheological models of continental lithosphere.** (a) Crème brûlée model (CB) and Jelly sandwich

646 model (JS). The two rheological models can be further subdivided into CB-I, CB-II, JS-I, and JS-II according to

647 the strength of the lower crust (modified from Jackson, 2002). (b) Observations of four distinct lithosphere

648 rheological structures implied for East Asia (modified from Zhang et al., 2013). Locations of strength profiles are

649 pointed out by dots with numbers in the topography map. Dots filled with different colors indicate different models

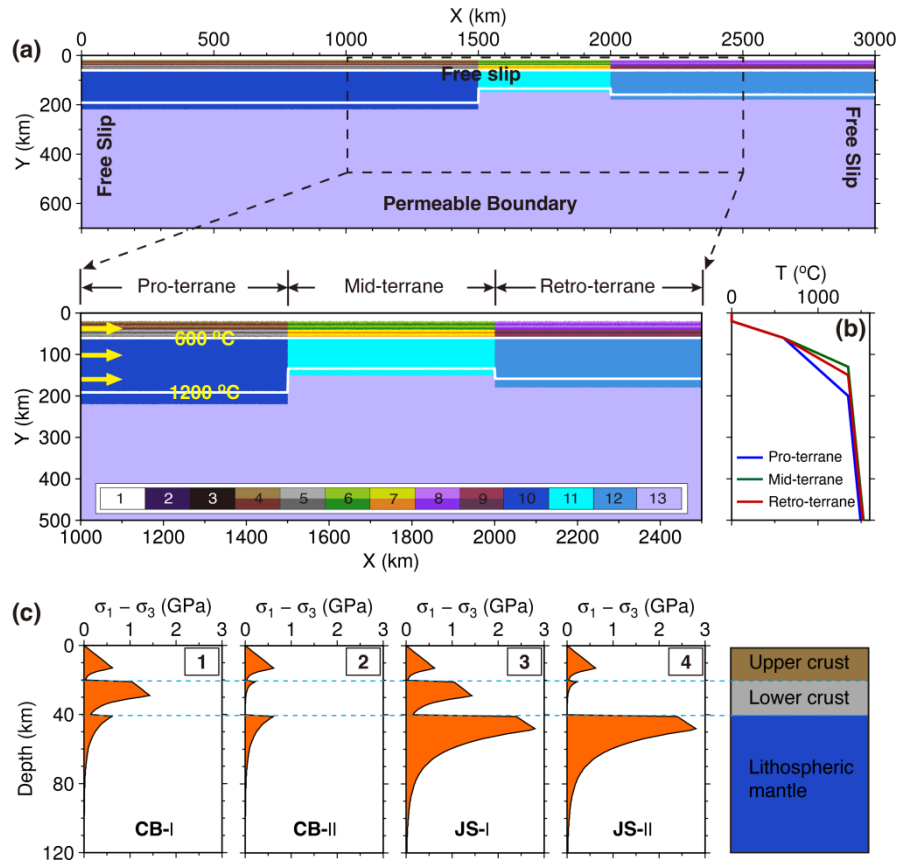
650 of lithospheric rheology. These strength profiles are calculated based on observed geothermal structure and

651 lithospheric structure, and assumed that compositions of the upper and lower crust and lithospheric mantle are wet

652 quartzite, undried granulite and dry olivine, respectively. Variations of temperature and lithospheric compositions

653 lead to a diverse suite of strength profiles vs. depth.

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656 **Figure 2. Initial model setup.** (a) Initial model configuration. The model size is 3000 km × 700 km, and size of

657 study region is 1500 km × 500 km. Three continental terranes of the Pro-, Mid- and Retro-terrane are contained

658 in the numerical model, and they are 200 km, 130 km, and 160 km thick, respectively. White lines are isotherms

659 with an interval of 600 °C. Yellow arrows indicate the convergence rate of 20 mm/yr. Colored grids: 1 – sticky air;

660 2 – sediments; 3 – weak zone; 4, 6, 8 – the upper crust of the Pro-, Mid- and Retro-terrane, respectively; 5, 7, 9 –

661 the lower crust of the Pro-, Mid- and Retro-terrane, respectively; 10, 11, 12 –lithospheric mantle of the Pro-,

662 Mid- and Retro-terrane, respectively; 13 – asthenosphere. (b) Initial temperature structure for the three terranes.

663 The Pro- and Mid-terrane respectively have a coldest and warmest lithospheric mantle due to their differences of

664 lithosphere thicknesses. (c) Four rheological models with contrasting lithospheric strength profiles. These are

665 derived from different strength scaling factor ( $S$ ) combinations for the upper crust, lower crust, and lithospheric

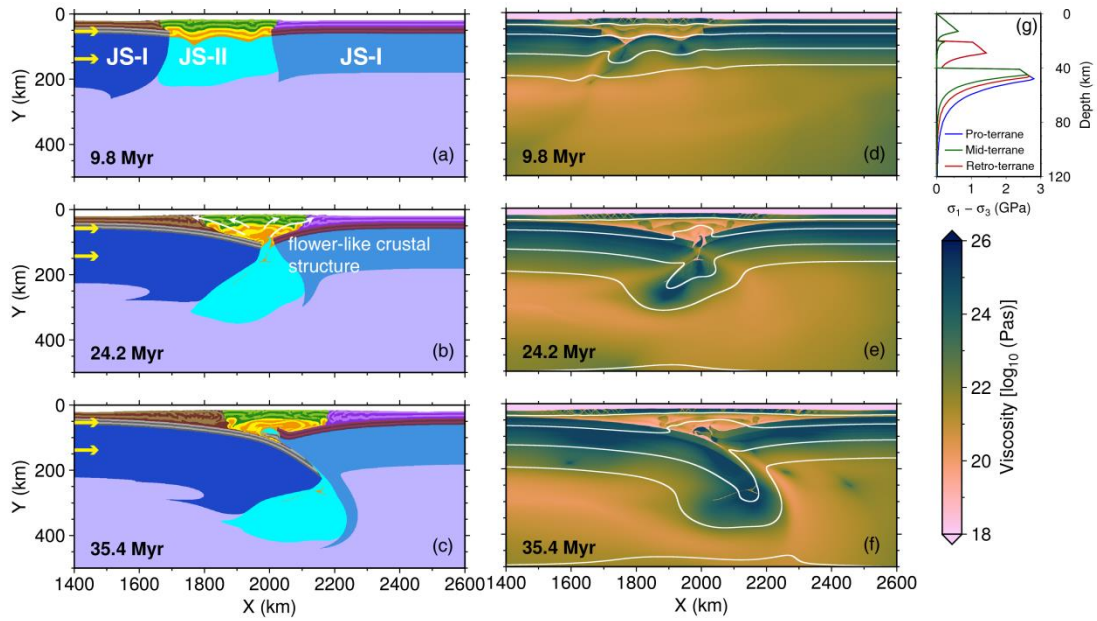
666 mantle (Table S1). Strength profiles are calculated based on the Pro-terrane's initial lithospheric structure,

667 composition, and temperature field. The prescribed strain rate is  $1 \times 10^{-14} \text{ s}^{-1}$ . CB-I and CB-II, the crème brûlée

668 model with strong and weak lower crust, respectively; JS-I and JS-II, the jelly sandwich model with strong and

669 weak lower crust, respectively.

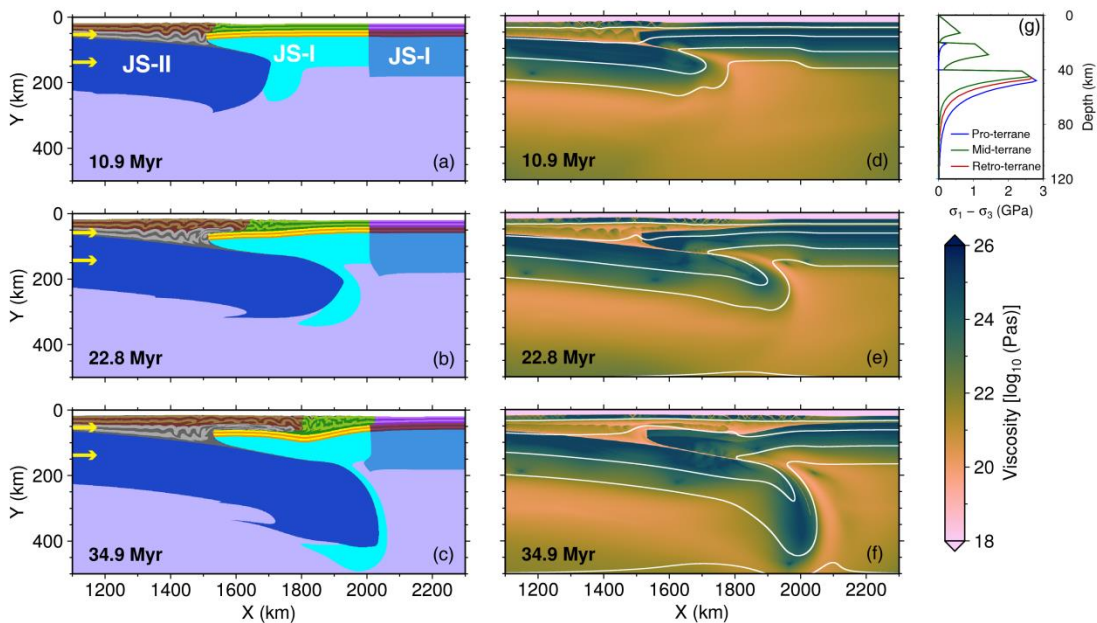
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672 **Figure 3. Collision of the lithospheres of the Pro- and Retro-terranes.** Rheological models for the Pro-, Mid-  
 673 and Retro-terranes are JS-I, JS-II, and JS-I, respectively, as shown in (g). The left panel shows compositional  
 674 fields at 9.8 Myr, 24.2 Myr, and 35.4 Myr, respectively. Yellow arrows indicate the convergence rate. The right  
 675 panel shows the corresponding viscosities. White lines are isotherms with an interval of 300 °C.

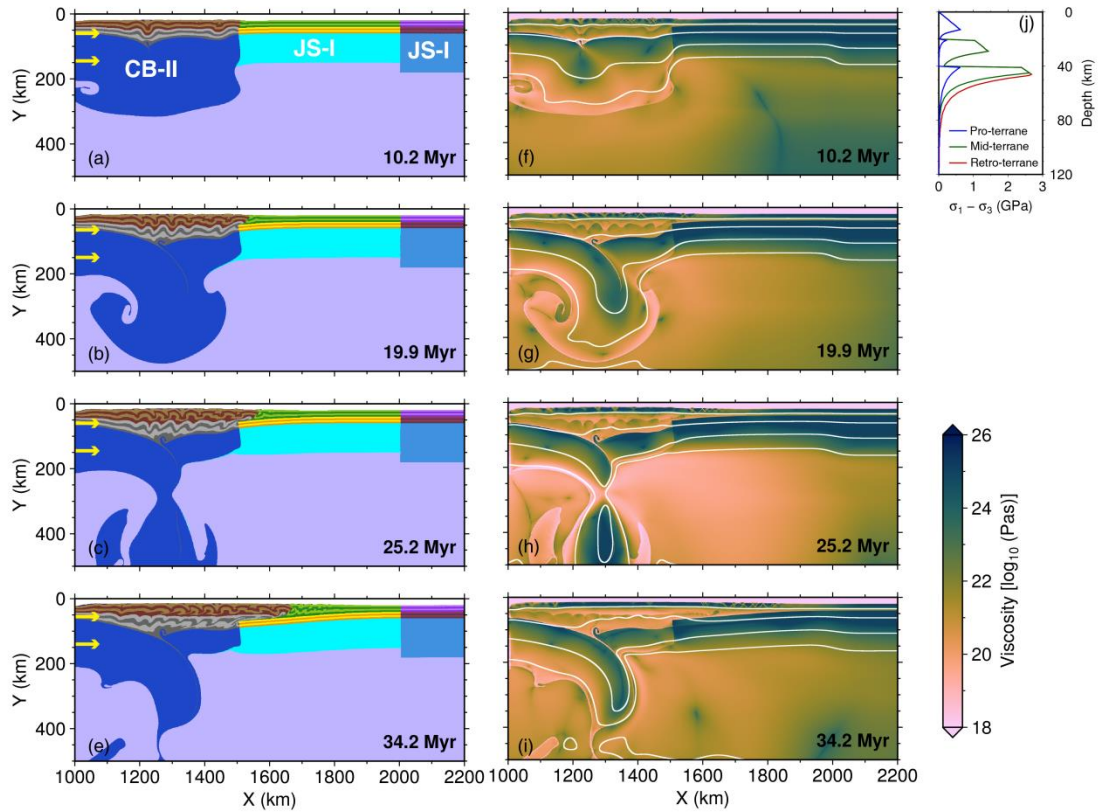
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678 **Figure 4. Subduction of the lithosphere of the Pro-terrane.** Rheological models for the Pro-, Mid- and  
 679 Retro-terranes are JS-II, JS-I, and JS-I, respectively. See Figure 3 for plotting conventions.

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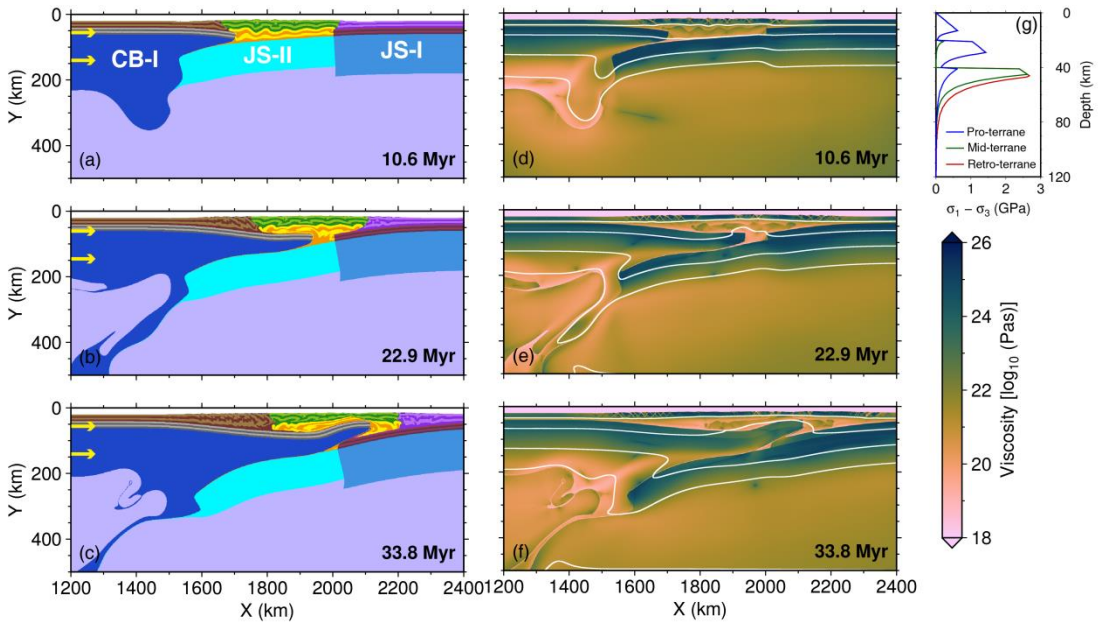


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682 **Figure 5. Thickening and delamination of the lithosphere of the Pro-terrane.** Rheological models for the Pro-,

683 Mid- and Retro-terrane are CB-II, JS-I, and JS-I, respectively. See Figure 3 for plotting conventions.

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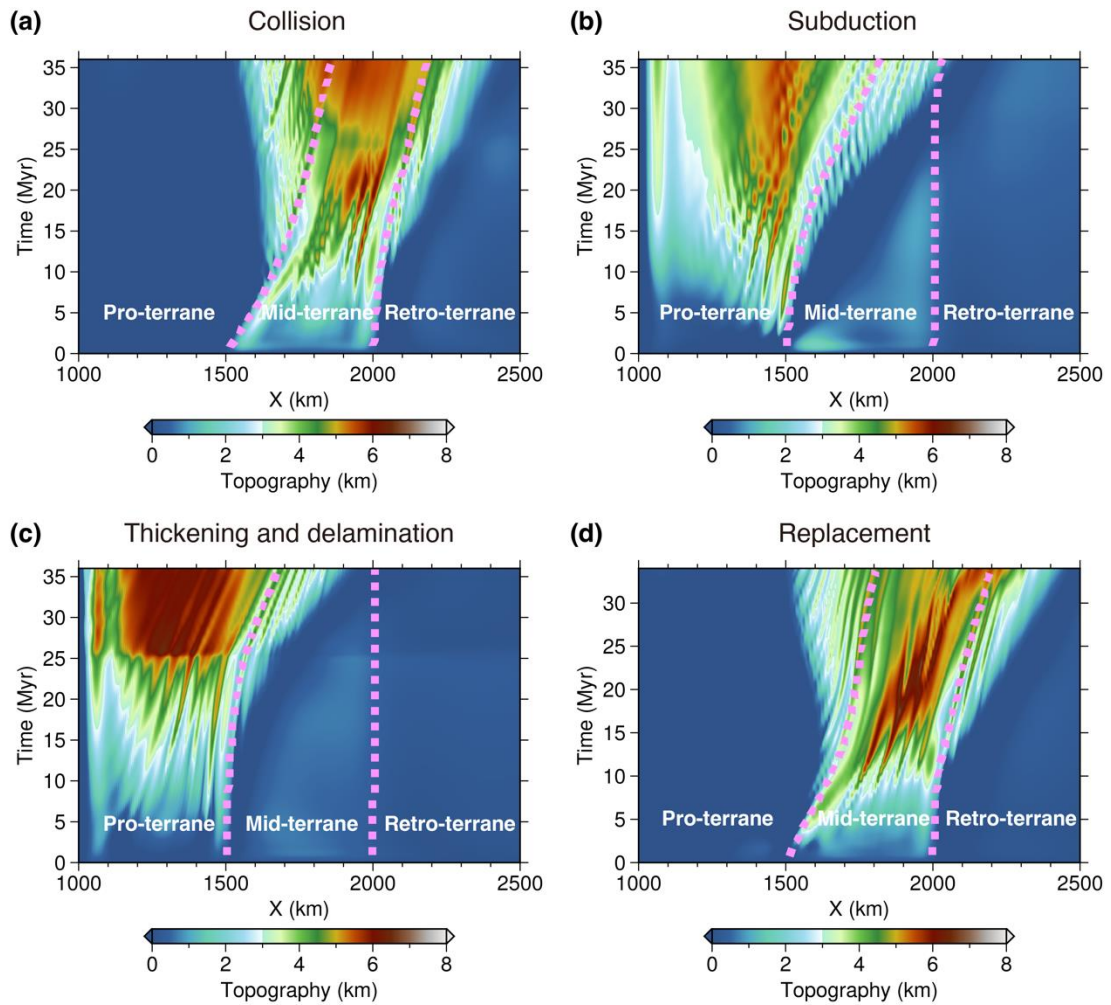


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686 **Figure 6. Replacement of the lithosphere of the Pro-terrane.** Rheological models for the Pro-, Mid- and

687 Retro-terrane is CB-I, JS-II, and JS-I, respectively. See Figure 3 for plotting conventions.

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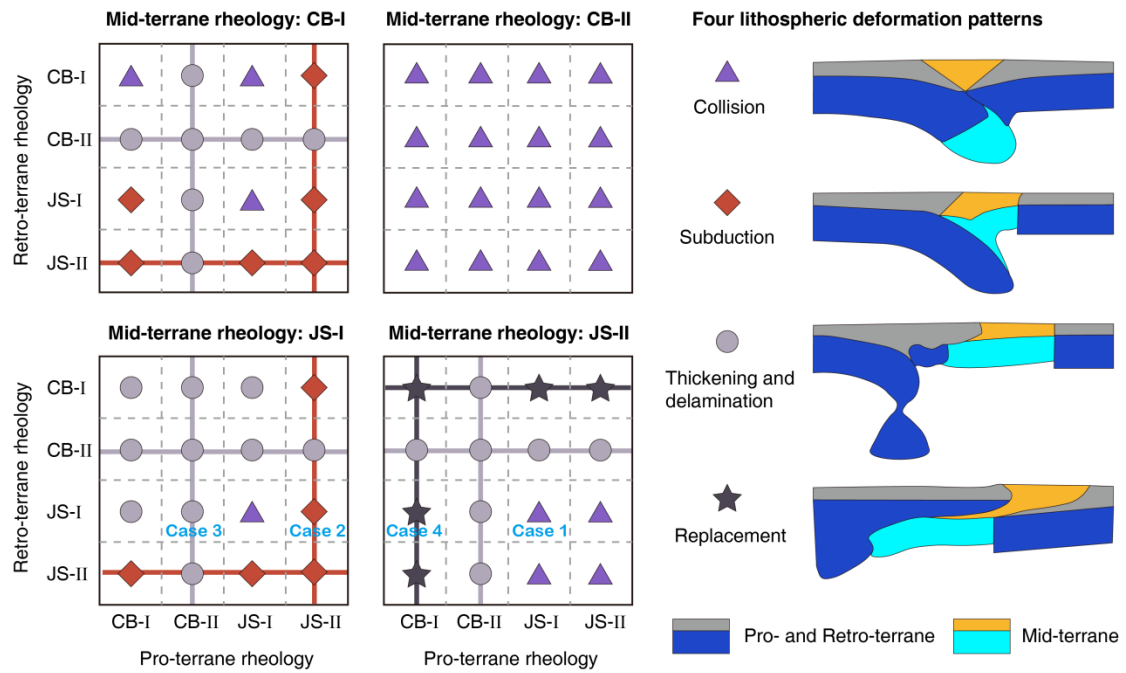
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690 **Figure 7. Evolution of surface relief for the different deformation styles.** The purple dashed lines indicate the

691 boundaries between terranes. (a) – (d) Surface relief associated with the deformation patterns of lithosphere

692 collision, subduction, thickening and delamination, and replacement, respectively.

693



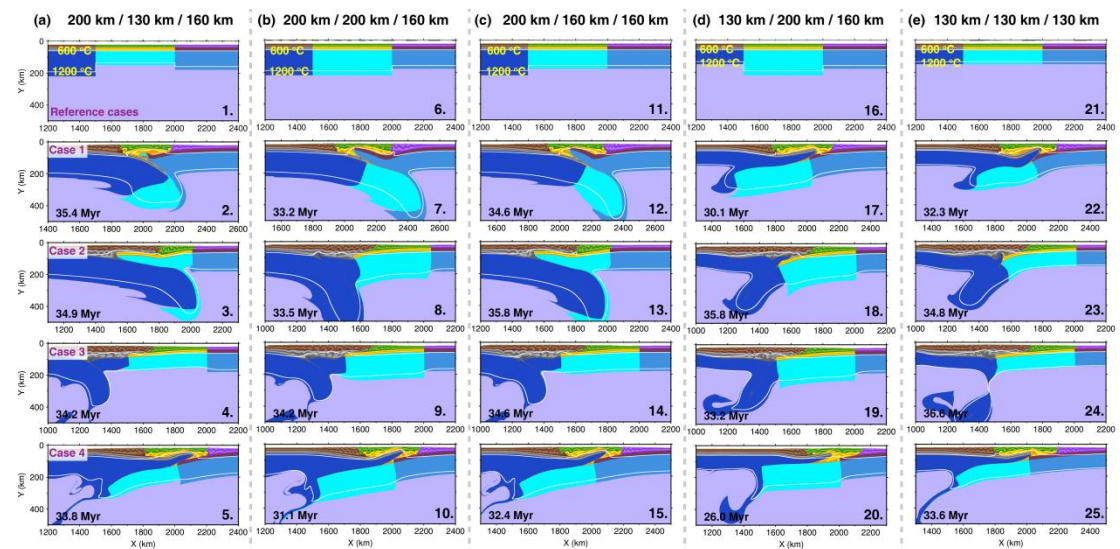
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695 **Figure 8. Four styles of lithosphere deformation patterns.** Symbols with colors indicate different deformation

696 patterns of the lithosphere. Cases 1 – 4 are the selected models chosen to illustrate details of these modes of

697 compressional evolution.

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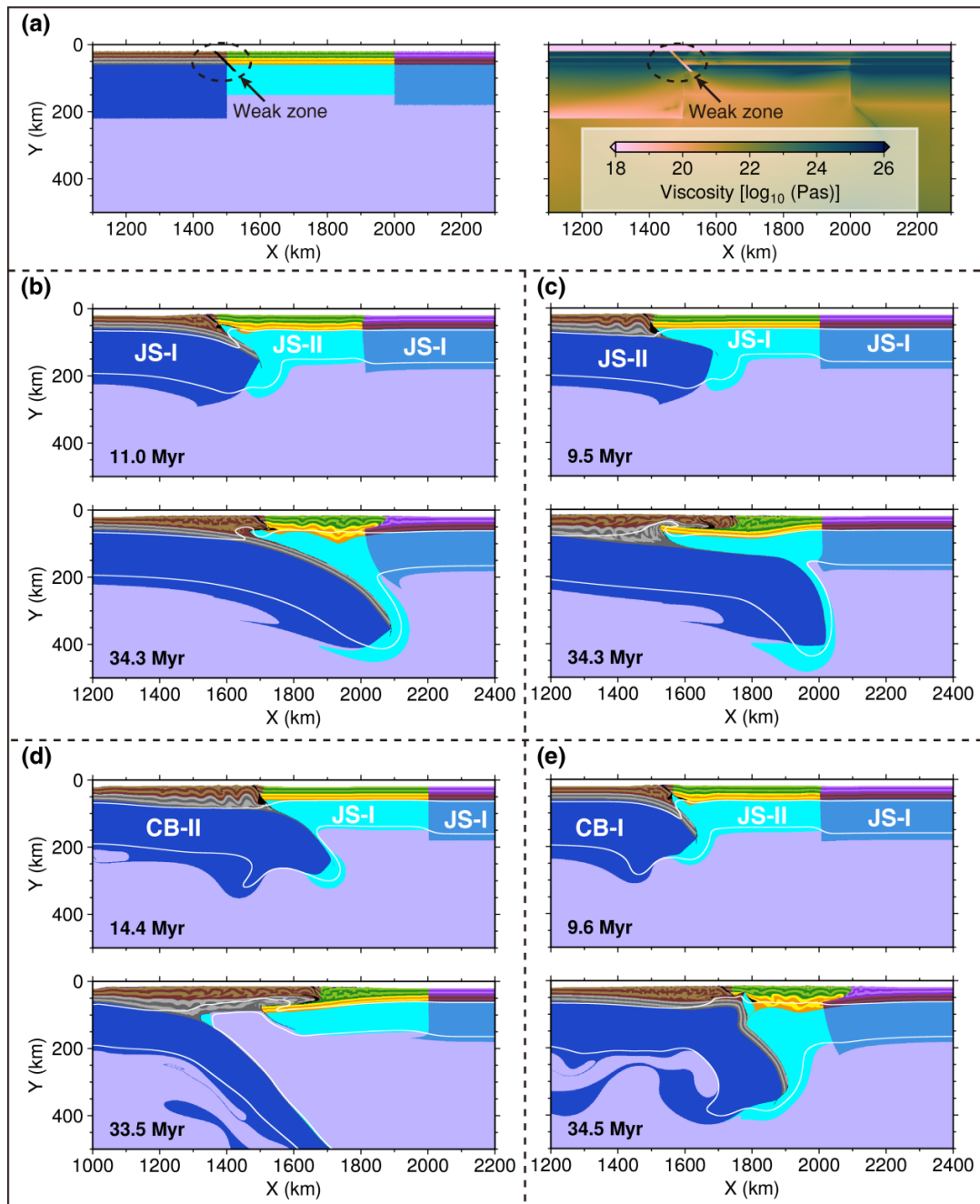
700 **Figure 9. Effects of lithosphere thicknesses of various terranes.** (a) – (e) Final simulation results of models with

701 the different lithosphere thicknesses of the Pro-, Mid-, and Retro-terranes, respectively. (a) Final simulation

702 results of reference cases. Rheological models of the Pro-, Mid- and Retro-terranes in 2 – 4 rows are same with

703 those in Cases 1 – 4, respectively.

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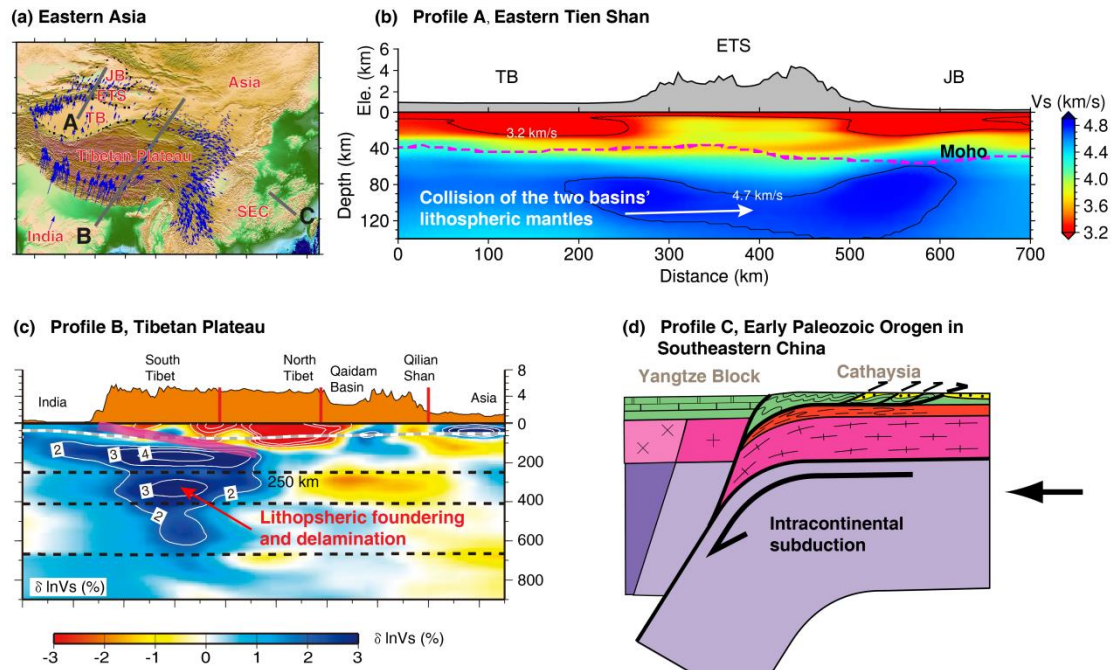
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706 **Figure 10. Effects of local weak zone on lithosphere deformation.** (a) Details about the weak zone. (b) – (e)

707 Final simulation results of models corresponding to Cases 1 – 4, respectively.

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711 **Figure 11. Implications of simulation results to East Asia.** (a) Topography of East Asia. The three gray lines  
 712 point out the locations of lithosphere profiles in (b), (c) and (d). TB, Tarim Basin; ETS, eastern Tien Shan; JB,  
 713 Junggar Basin; SEC, Southeastern China. (b) Collision of the lithospheric mantle of Tarim Basin and Junggar  
 714 Basin beneath the eastern Tien Shan (modified from Lü et al., 2019). (c) Lithospheric founding and delamination  
 715 in the Tibetan Plateau (modified from Chen et al., 2017). (d) Intracontinental subduction in the Early Paleozoic  
 716 Orogen in southeastern China (modified from Faure et al., 2009).



717 **Table 1. Flow laws and material properties for different lithospheric layers.**  $\rho_0$  is the initial density; it evolves  
718 with time as  $\rho=\rho_0(1-\alpha(T-T_0))(1+\beta(P-P_0))$ , where  $T_0 = 20\text{ }^\circ\text{C}$ ,  $P_0 = 10^5\text{ MPa}$ . Flow law: qtz. = quartzite, Plag.  
719 = plagioclase, ol. = olivine.

Material properties	Sediment	Upper crust	Lower crust	Lithospheric mantle	Asthenosphere
$\rho_0$ (kg/m <sup>3</sup> )	2600	2700	2800	3300	3300
Flow laws	Wet qtz.	Wet qtz.	Plag.	Dry ol.	Dry ol.
$1/A_D$ (Pa <sup>n</sup> s)	$1.97 \times 10^{17}$	$1.97 \times 10^{17}$	$4.80 \times 10^{22}$	$3.98 \times 10^{16}$	$3.98 \times 10^{16}$
$n$	2.3	2.3	3.2	3.5	3.5
$E_a$ (KJ/mol)	154	154	238	532	532
$V_a$ (J/bar)	0.8	0.8	1.2	1.2	1.2
$\phi = \sin(\varphi)$	0.2 – 0.1	0.3 – 0.1	0.3 – 0.1	0.6 – 0.4	0.6 – 0.3
$C$ (Pa)	$1 \times 10^{7-6}$	$1 \times 10^{7-6}$	$1 \times 10^{7-6}$	$1 \times 10^{7-6}$	$1 \times 10^{7-6}$
$H_r$ (uW/m <sup>3</sup> )	2.0	1.5	0.5	0.022	0.022
$C_p$ (J/kg K)	1000	1000	1000	1000	1000
$\alpha$ (1/K)	$3 \times 10^{-5}$	$3 \times 10^{-5}$	$3 \times 10^{-5}$	$3 \times 10^{-5}$	$3 \times 10^{-5}$
$\beta$ (1/MPa)	$1 \times 10^{-5}$	$1 \times 10^{-5}$	$1 \times 10^{-5}$	$1 \times 10^{-5}$	$1 \times 10^{-5}$
$k$ (W/m/K)	$0.64+807/(T+77)$	$0.64+807/(T+77)$	$1.18+474/(T+77)$	$[0.73+1293/(T+77)] \times (1+0.00004P)$	$[0.73+1293/(T+77)] \times (1+0.00004)$

720