



1 **Resolving the mesoscale at reduced computational cost with FESOM 2.5: efficient modeling**  
2 **approaches applied to the Southern Ocean**

3 Nathan Beech<sup>1</sup>, Thomas Rackow<sup>2</sup>, Tido Semmler<sup>3</sup>, and Thomas Jung<sup>1,4</sup>

4 1. Alfred Wegener Institute Helmholtz Center for Polar and Marine Research, Bremerhaven,  
5 Germany

6 2. European Center for Medium-range Weather Forecasts, Bonn, Germany

7 3. Met Eireann, the Irish Meteorological Service, Dublin, Ireland

8 4. Department of Physics and Electrical Engineering, University of Bremen, Bremen, Germany

9 Corresponding Author: Nathan Beech (Nathan.beech@awi.de)

10



## 11 **Abstract**

12           Several cost-efficient, high-resolution modeling approaches are applied to simulations of the  
13 Southern Ocean in past, present, and future climates. The results are compared with an ensemble of  
14 medium-resolution, eddy-present simulations and evaluated based on their ability to reproduce observed  
15 mesoscale activity and to reveal a response to climate change distinct from natural variability. The high-  
16 resolution simulations reproduce the observed magnitude of Southern Ocean eddy kinetic energy (EKE)  
17 well, but differences remain in local magnitudes and the spatial distribution of EKE. The coarser, eddy-  
18 present ensemble simulates a similar pattern of EKE but underrepresents observed levels by 50%. Five  
19 years of simulated data in each time period is found to produce consistent results when evaluating mean  
20 conditions and assessing change in the region as a whole. At 1 °C of warming, the high-resolution  
21 simulations produce no change in overall EKE, in contrast to the increase projected by the eddy-  
22 permitting ensemble and despite full ensemble agreement. At 4 °C of warming, both datasets produce  
23 consistent levels of EKE rise in relative terms, although not absolute magnitudes, as well as an increase in  
24 EKE variability. Simulated EKE rise is concentrated where flow interacts with topographic features in  
25 regions already known to be eddy-rich. Regional EKE change in the high-resolution simulations is  
26 consistent with changes seen in at least four of five eddy-permitting ensemble members at 1 °C of  
27 warming, and all ensemble members at 4 °C. However, substantial noise would make these changes  
28 difficult to distinguish from natural variability without an ensemble.

## 29 **Plain Language Summary**

30           Ocean models struggle to simulate small-scale ocean flows due to the computational cost of high-  
31 resolution simulations. Several cost-reducing strategies are applied to simulations of the Southern Ocean  
32 and evaluated with respect to observations and traditional, lower-resolution modelling methods. The high-  
33 resolution simulations effectively reproduce small-scale flows seen in satellite data and are largely  
34 consistent with traditional model simulations regarding their response to climate change.

35



## 36 **1 Introduction**

37 Mesoscale activity in the Southern Ocean has been the subject of much research and interest in  
38 recent years due to the intensification of Southern Hemisphere westerlies (Marshall, 2003), the  
39 phenomena of eddy saturation and compensation (Munday et al., 2013; Bishop et al., 2016), and the  
40 potential for carbon sequestration in the face of ongoing anthropogenic emissions (Sallée et al., 2012;  
41 Landschützer et al., 2015; Frölicher et al., 2015). Observations already reveal an intensification of eddy  
42 activity in the Antarctic Circumpolar Current (ACC) and changes are attributed primarily to wind stress  
43 (Marshall, 2003; Hogg et al., 2015; Martínez-Moreno et al., 2021). Modeling studies have been able to  
44 reproduce the observed changes, as well as project continued intensification throughout the 21<sup>st</sup> century  
45 (Beech et al., 2022), but the modeled results rely on only partially resolved eddy activity relative to  
46 observations, leaving open the possibility for new findings or greater clarity.

47 Advances in computational capabilities have enabled ocean modeling science to make great  
48 progress in overcoming the substantial computational burden of simulating the mesoscale, but  
49 shortcomings remain, particularly in the Southern Ocean where the Rossby radius can be as small as 1  
50 km, increasing the computational cost of resolving eddies (Hallberg, 2013). Even model resolutions that  
51 can generally be considered eddy-resolving are only eddy-permitting poleward of 50° if they do not vary  
52 in space (Hewitt et al., 2020). This highlights an efficiency challenge in simulating the mesoscale with  
53 traditional model grids; resolutions necessary to resolve high-latitude, small-radius eddies are both  
54 prohibitively expensive and unnecessary to resolve mesoscale eddies in the lower latitudes. Fortunately, a  
55 growing number of modelling alternatives to traditional grids now enable dynamic spatial allocation of  
56 resources (Danilov, 2013; Ringler et al., 2013; Danilov et al., 2017; Jungclaus et al., 2022), creating the  
57 opportunity to more efficiently resolve the mesoscale.

58 As resource allocation in high-resolution modeling becomes spatially flexible in the pursuit of  
59 more efficient configurations, the temporal component must also be scrutinized for efficiency. Traditional  
60 modeling approaches require long spin-up periods in order to equilibrate the deep ocean and reduce



61 model drift (Irving et al., 2021). Although the impacts of drift are not negligible, it generally affects large-  
62 scale processes in the deep ocean; mesoscale processes that require high resolutions to simulate are  
63 typically fast-to-equilibrate and will appear relatively quickly wherever large-scale ocean conditions lead  
64 to their creation. Admittedly, one cannot entirely disentangle the two scales, as mesoscale activity does  
65 affect the position of fronts and the paths of ocean circulation (Marzocchi et al., 2015; Chassignet and Xu,  
66 2017). Yet, with equilibration times for the deep ocean on the scale of thousands of years (Irving et al.,  
67 2021), the possibility, and ultimately necessity, to reduce the resolution of spin-up runs relative to  
68 production runs must be investigated.

69         Advancing the concept of dynamic temporal allocation of resources further, the traditional  
70 transient climate change simulation also represents an efficiency bottleneck for some applications; by  
71 modifying the climate continuously in time, each year of a transient simulation is effectively a single  
72 realization of a global mean climatic state that varies from the following and preceding years by only a  
73 fraction of a degree. For some applications, like hindcasts of real events or trend analysis, this approach  
74 may be desirable, but for assessing the impacts of climate change with limited resources and a low signal-  
75 to-noise ratio, a larger sample of realizations for a consistent climatic state may be more suitable.

76         Aside from oceanic concerns, the atmosphere can have substantial impacts on mesoscale activity  
77 in climate models. Most simply, a coupled atmosphere will react to ocean eddy activity, whereas  
78 atmospheric forcing will not, resulting in more eddy killing by wind stress (Renault et al., 2016).  
79 Additionally, an atmosphere coupled to a high-resolution ocean must be of similarly high resolution for  
80 certain mesoscale interactions to be resolved (Byrne et al., 2016). Ultimately, the modeled atmosphere  
81 further escalates the already exponential cost of increasing ocean resolution by requiring more  
82 computational resources in order for the benefits of the resolved mesoscale to fully transfer to the broader  
83 climate.

84         To address the computational inefficiencies outlined above, an experimental configuration is  
85 proposed, combining several experimental modeling approaches. Simulations will exploit the multi-



86 resolution Finite volume Sea-ice Ocean Model (FESOM) (Danilov et al., 2017) employing a high-  
87 resolution unstructured mesh that concentrates computational resources on the Southern Ocean, while  
88 maintaining grid resolution in the remainder of the global ocean that can still be considered high-  
89 resolution, as in, for example, HighResMIP (Haarsma et al., 2016). The multi-resolution strategy  
90 overcomes the efficiency challenges of resolving high-latitude eddies without needlessly increasing  
91 tropical resolutions, as well as limiting the focus and computational requirements to one hemisphere. The  
92 high-resolution simulations will make use of a spin-up simulation on a medium-resolution, eddy-  
93 permitting mesh to avoid the computational burden of allowing an eddy-resolving ocean to equilibrate  
94 deep, slow-changing processes. The eddy-permitting mesh will also be used to simulate the transient  
95 periods between shorter, high-resolution time slices, increasing the signal-to-noise ratio of the results by  
96 separating the production data further in time and the progression of anthropogenic climate change.  
97 Finally, the ocean model will be forced with atmospheric data from existing coupled simulations  
98 (Semmler et al., 2020). Although this will not facilitate mesoscale atmosphere-ocean interaction, the  
99 simulation will reflect the climatic development of an eddy-permitting simulation of the future  
100 atmosphere without the additional computational requirements.

101         The Southern Ocean is one of the world's hotspots for mesoscale activity and a region where  
102 substantial change is anticipated in the context of anthropogenic climate change (Beech et al., 2022).  
103 Simultaneously, the high latitude of the region makes eddy-resolving model simulations computationally  
104 demanding and observational data relatively scarce (Auger et al., 2023; Hallberg, 2013). Yet, as the  
105 climate changes, the importance of the Southern Ocean grows as a heat and carbon sink, an ecosystem,  
106 and a medium for feedback between the atmosphere and ocean (Byrne et al., 2016; Frölicher et al., 2015).  
107 Thus, the study of the Southern Ocean demands innovation in the modeling field to produce high-  
108 resolution simulations at reduced computational cost. This study maximizes grid resolution relative to  
109 computational cost using an unstructured, multi-resolution grid, a medium-resolution spin-up simulation,  
110 and atmospheric forcing from lower-resolution coupled simulations in order to focus resources as much as



111 possible on resolving mesoscale activity in the study region. The resulting simulations enable an  
112 exploratory analysis of the past, present, and future of the Southern Ocean with a fully resolved  
113 mesoscale. Simulations with this cost-efficient, high-resolution configuration are presented in comparison  
114 to a comprehensive ensemble of eddy-permitting simulations to assess the performance of the efficiency-  
115 focused approach in reproducing mesoscale activity and its response to climate change.

## 116 **2 Methods**

### 117 **2.1 Experimental setup**

118 This analysis is broadly a comparison of results from AWI-CM-1-1-MR's contribution to CMIP6  
119 (hereafter referred to as AWI-CM-1) (Semmler et al., 2020) and simulations using an updated version of  
120 FESOM (FESOM 2.5) and a mesh substantially refined to a resolution reaching less than 3 km in the  
121 Southern Ocean (hereafter referred to as SO3) (Supplementary Figure 1). Observations of ocean surface  
122 velocity derived from satellite altimetry data are also used to evaluate the model performance of each  
123 simulation. The AWI-CM-1 simulations are state-of-the-art CMIP6 experiments and benefit from the  
124 multiple ensemble members and long spin-up times that CMIP simulations typically boast. However,  
125 while the AWI-CM-1 ensemble reproduces eddy activity remarkably well within the context of CMIP6  
126 (Beech et al., 2022), high-resolution ocean modeling now far surpasses even the highest ocean resolutions  
127 in the CMIP6 ensemble. Conversely, the SO3 simulations push the limits of ocean resolution but rely on  
128 several measures for maximizing computational efficiency that may impact the robustness of the  
129 simulations. Details on the experimental setup for CMIP6 and ScenarioMIP are widely available (Eyring  
130 et al., 2016; O'Neill et al., 2016) and information more specific to AWI-CM-1-1-MR's contribution has  
131 been published previously (Semmler et al., 2020). The following sections will outline the details of the  
132 SO3 simulations.

133 The model experiments with SO3 consist of a medium-resolution, eddy-permitting, ocean-only  
134 transient simulation from 1851-2100 and three shorter simulations with the SO3 mesh at different time



135 periods during the progression of anthropogenic climate change. The medium-resolution transient run is a  
136 stand-alone ocean simulation using a medium-resolution mesh that has been shown to effectively  
137 reproduce eddy activity in active regions while maintaining a computational cost comparable to a  
138 traditional  $\frac{1}{4}^\circ$  model (Beech et al., 2022). The model is initialized with conditions for ocean temperature  
139 and salinity, as well as sea ice concentration, thickness, and snow cover taken from the end of the first  
140 year (1850) and first ensemble member (r1i1p1f1) of AWI-CM-1-1-MR's historical simulations in  
141 CMIP6 (Semmler et al., 2018, 2020, 2022a, b). In this way, the model undergoes a semi-cold start in  
142 which ocean conditions are not exact continuations of the previous coupled simulation, but nonetheless  
143 benefit from more realistic and partially equilibrated ocean properties. The stand-alone ocean is forced  
144 using atmospheric data from the same ensemble member of the historical CMIP6 simulations until 2014  
145 (Semmler et al., 2022a), and thereafter using the first ensemble member of AWI-CM-1-1-MR's  
146 ScenarioMIP simulations for shared socioeconomic pathway (SSP) 3-7.0 (Eyring et al., 2016; O'Neill et  
147 al., 2017; Semmler et al., 2022b). This approach to forcing takes advantage of a coupled simulation,  
148 CMIP6, to produce a forcing dataset of better temporal and spatial coverage than the observational record  
149 and which maintains a realistic transient climate throughout anthropogenic impacts during the 21<sup>st</sup>  
150 century.

151 In the years 1950, 2015, and 2090, the model is reinitialized with the higher-resolution ocean  
152 grid, SO3 (Supplementary Figure 1), using the same semi-cold start approach and forcing dataset, as well  
153 as initial conditions from the end of the previous year of the eddy-permitting transient simulation. The  
154 high-resolution grid is, in truth, a regionally refined mesh in which a 25 km global resolution is increased  
155 to approximately 2.5 km primarily south of 40 °S, but with other pertinent regions, such as the Agulhas  
156 Current and several narrow straits, also refined. In this way, the model is able to simultaneously achieve  
157 eddy-rich conditions in the Southern Ocean and many of the nearby active regions as well as a global  
158 resolution that would still be considered high in the context of CMIP6 (Hallberg, 2013; Hewitt et al.,  
159 2020). The high-resolution simulations are each integrated for six years with the first year ignored as a



160 true spin-up, leaving five years of data for each time period. While model drift may be a concern with  
161 such a short true spin-up period, this should affect each of the high-resolution time slices similarly and to  
162 a limited extent due to their short integration lengths. Thus, the differences between the high-resolution  
163 ocean simulations should primarily reflect anthropogenic climate impacts simulated during the eddy-  
164 permitting transient run.

## 165 **2.2 Model configuration**

166 The Finite volume Sea-ice Ocean Model version 2.5 is a post-CMIP6 era model, having been  
167 refactored to a finite-volume configuration from the finite-element version (FESOM1.4, Q. Wang et al.,  
168 2014) employed in CMIP6, and transitioned to arbitrary Lagrangian Eulerian (ALE) vertical coordinates,  
169 among other improvements (Danilov et al., 2017; Scholz et al., 2019, 2021). FESOM's most  
170 distinguishing feature among mature ocean models is the unstructured horizontal grid that exploits  
171 triangular grid cells which can smoothly vary in size to change the horizontal grid resolution in space. In  
172 these simulations, full free surface, or  $z^*$ , vertical coordinates were used, allowing the vertical model  
173 layer thicknesses to change in time. Gent-McWilliams eddy parameterization (Gent and McWilliams,  
174 1990) is scaled with resolution according to (Ferrari et al., 2010) and vertical mixing is simulated by a  $k$ -  
175 profile parameterization scheme (Large et al., 1994).

176 The SO3 mesh consists of over 22 million surface elements (triangle faces) or 11 million surface  
177 nodes (triangle vertices) and 70 vertical layers. The simulations produce about 1.1 terabytes of data per  
178 year of 3D data stored on nodes. The model was run on 8192 CPU cores and with a typical throughput of  
179 approximately 0.65 simulated years per day, consuming approximately 5.5 million CPU hours in total  
180 despite the various cost-saving modeling approaches. The simulations and following analysis were  
181 performed using the high-performance computing system, Levante, at the German Climate Research  
182 Center (DKRZ).





183           The ocean model is forced by several atmospheric variables at a six-hour resolution, although one  
184 forcing variable, humidity, is interpolated monthly data. The forcing data is first interpolated to a regular  
185 grid which can be interpreted by the model and applied to the multi-resolution grid used in the respective  
186 simulations. Runoff data is a monthly climatology and dynamic ice sheet coupling is not included,  
187 meaning the freshwater influx from the Antarctic continent does not react to warming which may impact  
188 certain processes, such as the timing and intensity of sea ice loss (Pauling et al., 2017; Bronselaer et al.,  
189 2018).

### 190 **2.3 Altimetry data**

191           Daily geostrophic velocities are taken from a gridded altimetry product derived from crossover  
192 data available from the Data Unification and Altimeter Combination System (DUACS) (Taburet et al.,  
193 2019). The gridded product has a resolution of  $0.25^\circ$ , although effective resolution at high latitudes may  
194 be much lower (Ballarotta et al., 2019). Recently, improved data has become available in the ice-covered  
195 regions of the Southern Ocean (Auger et al., 2022), but does not yet cover the present-day simulated  
196 period (2016-2020) in this study. Absolute velocities from the gridded altimetry product were used to  
197 calculate anomalies and EKE using equations (3) and (4) below for consistency with the modeled dataset.

### 198 **2.4 Geostrophic velocities**

199           Ocean velocities in the SO3 simulations are saved on a daily timescale as direct model output,  
200 whereas in the AWI-CM-1 ensemble, only monthly data is available and daily data must first be derived  
201 from sea surface height data and geostrophic balance as in equations (1) and (2). To avoid including  
202 ageostrophic contributions to ocean surface velocity in this analysis, model output velocity between 25 m  
203 and 30 m depth is used for the SO3 dataset. This depth should be close enough to the surface to closely  
204 match surface geostrophic flows while also avoiding Ekman transport, a major contributor to ageostrophic  
205 ocean flow. What ageostrophic flow remains in the model output velocities should be primarily large-



206 scale and overshadowed by geostrophic flow in the high-energy regions of the ocean, including the ACC

207 (Yu et al., 2021).

$$208 \quad u = -gf^* \partial SSH / \partial y \quad (1)$$

$$209 \quad v = gf^* \partial SSH / \partial x \quad (2)$$

210 *e. EKE analysis*

211 Velocity anomalies are defined by subtracting the multi-year monthly climatology of each

212 respective 5-year period from daily velocities with equation (3).

$$213 \quad u'_i = u_i - \overline{u_m} \quad (3)$$

214 Where  $u_i$  is the daily zonal velocity, ' denotes an anomaly, and  $\overline{u_m}$  is a monthly mean. For meridional

215 velocities ( $v$ ) substitute  $u$  with  $v$ .

216 Eddy kinetic energy is calculated from ocean velocities according to equation (4).

$$217 \quad EKE_i = 0.5(u'_i{}^2 + v'_i{}^2) \quad (4)$$

218 Where  $(i)$  denotes a daily value and  $(')$  denotes an anomaly.

219 EKE was calculated on the native grid of each dataset and then interpolated to a  $0.25^\circ$  grid for all

220 analyses. EKE in Figures 1 and 3 was coarsened to five-day means before analysis. Area-integrated EKE

221 is calculated by summing the area-weighted EKE of each grid cell in the study region defined as the zonal

222 band between  $45^\circ\text{S}$  and  $60^\circ\text{S}$  with the Brazil/Malvinas confluence region between  $57^\circ\text{E}$  and  $29^\circ\text{E}$  and

223 northward of  $40^\circ\text{S}$  removed. EKE anomalies (Figure 1) were calculated by subtracting the 2016-2020

224 mean from the EKE data of each period. Normalized EKE was calculated by further dividing EKE

225 anomaly by the standard deviation of EKE during the 2016-2020 period. In Figure 4, ensemble agreement

226 is determined by ordering the  $\Delta EKE$  values within each grid cell from lowest to highest, plotting the

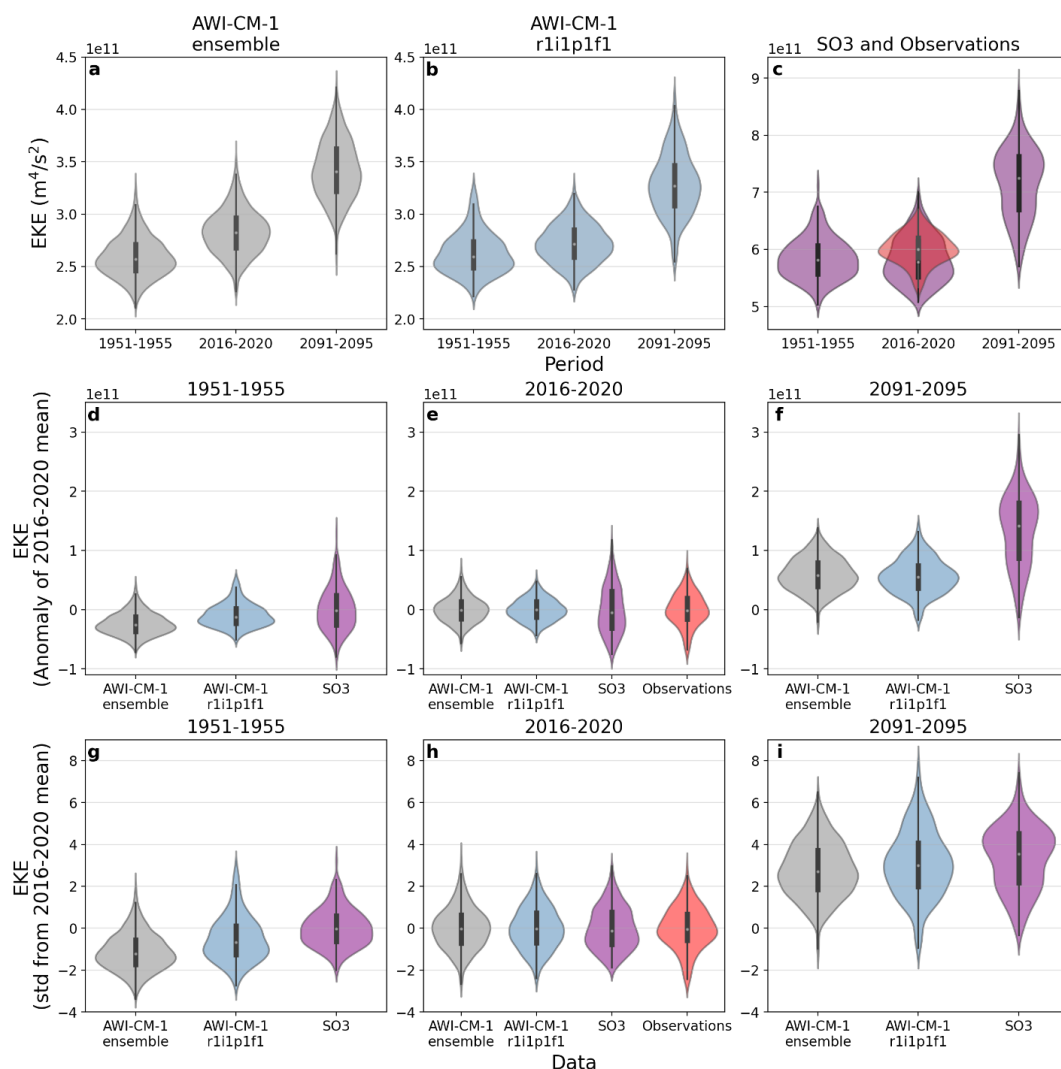


227 positive values in increasing order from left to right and negative values in decreasing order from left to  
228 right.

## 229 **3 Results**

### 230 **3.1 Agreement with observations**

231 During the five-year period of overlap with observations, the SO3 simulation is a drastic  
232 improvement on the AWI-CM-1 ensemble in reproducing median observed EKE (Figure 1a, c); only a  
233 slight underrepresentation of EKE remains in the SO3 simulation, although the simulated distribution  
234 remains distinct from observations, whereas the AWI-CM-1 ensemble underrepresents observations by  
235 about 50% (Figure 1a, c). EKE in SO3 appears more variable than the observations as can be inferred  
236 from the larger range (Figure 1c, e) and its distribution appears less Gaussian than observations or the  
237 eddy-permitting dataset (Figure 1a, b, c). Taking the AWI-CM-1 ensemble as an example, there is no  
238 consistent pattern of multimodality or skewness and the ensemble mean more closely approximates the  
239 normal distribution than the individual members (Figure 3), suggesting that deviations from normality  
240 could be primarily due to sample size rather than the character of the data. Relative to the AWI-CM-1  
241 model bias and the magnitude of EKE resolved, the ensemble spread within the AWI-CM-1 dataset is  
242 small (Figure 3), suggesting that a single ensemble member of five years duration is sufficient to assess  
243 how well a model captures the overall magnitude of Southern Ocean EKE (Figure 1c).



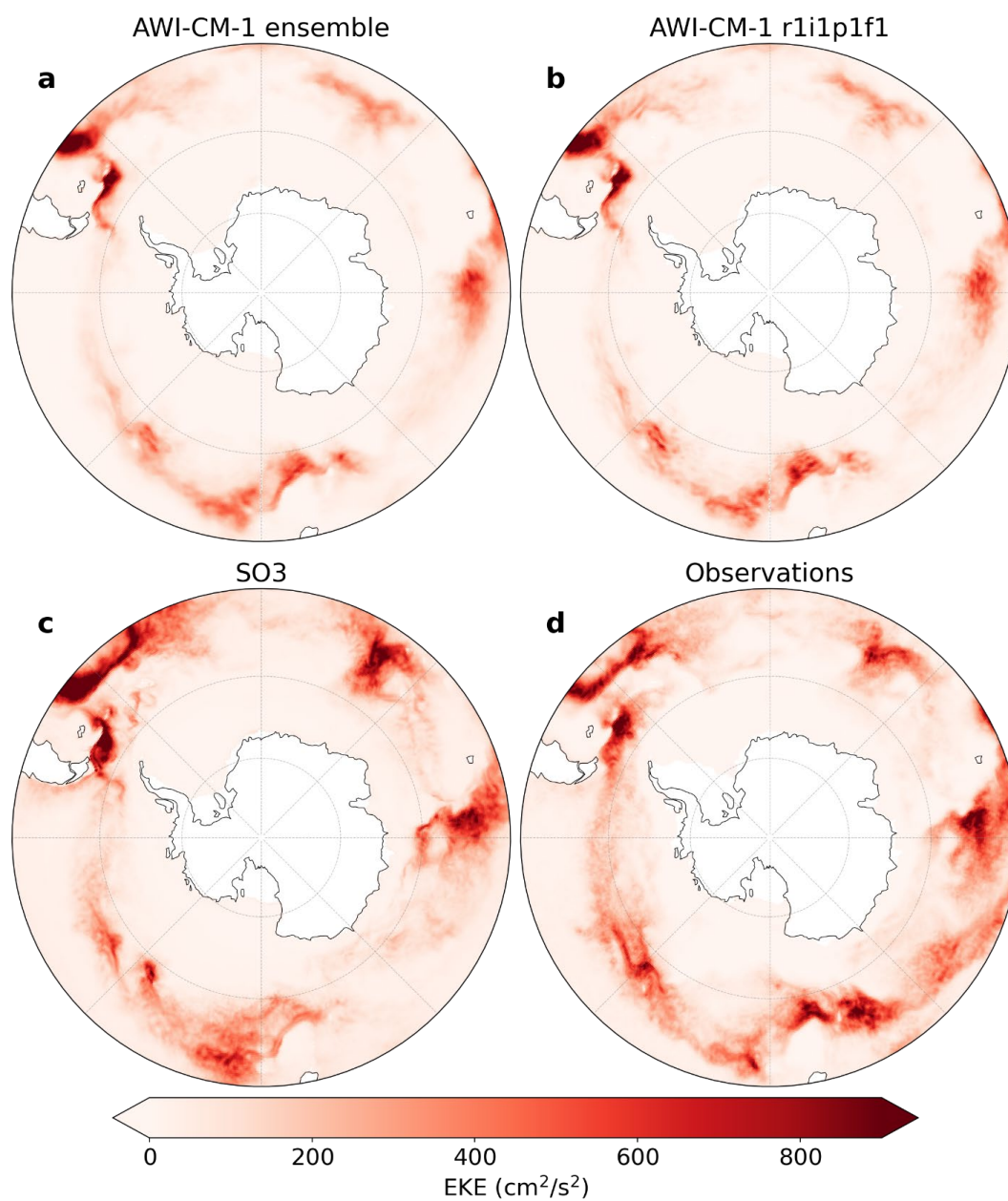
244

245 **Figure 1. Violin plots of EKE in simulations and observations.** Central points of each plot indicate the  
 246 median, thick bars span the first and third quartiles, thin bars span the range, and the violin body is a  
 247 kernel density estimation of the data. **a-c)** Real magnitudes of EKE (note the different y axes) **a)** The  
 248 AWI-CM-1 ensemble. **b)** the first member of the AWI-CM-1 ensemble, from which the SO3 simulations  
 249 take their atmospheric forcing. **c)** The So3 simulations and observations. **d-f)** EKE anomalies relative to  
 250 the 2016-2020 mean for each dataset respectively. **d)** 1951-1955. **e)** 2016-2020. **f)** 2091-2095. **g-i)**



251 Normalized values relative to the mean and standard deviation of EKE during the 2016-2020 period for  
252 each dataset respectively. **g)** 1951-1955. **h)** 2016-2020. **i)** 2091-2095.

253 From a regional perspective, the SO3 simulation accurately reflects local magnitudes of observed  
254 EKE and also generally captures the spatial distribution well (Figure 2). However, there are regional  
255 shortcomings, such as between 90 and 145 °E. Grid resolution in this region should be sufficient to  
256 resolve eddy activity (Supplementary Figure 1), indicating that the bias arises from another source. In the  
257 AWI-CM-1 ensemble, the regional representation of EKE reinforces a broad underrepresentation relative  
258 to observed magnitudes, but the major geographic features of eddy activity are fairly well represented  
259 (Figure 2). Once again, the ensemble spread within the AWI-CM-1 simulations reveals remarkable  
260 consistency, this time in terms of the spatial pattern and regional magnitudes (Supplementary Figure 2),  
261 reinforcing the conclusion that a single ensemble member of five years duration is sufficient to assess the  
262 mean state of EKE in the Southern Ocean. The consistency of the AWI-CM-1 ensemble further suggests  
263 that regional shortcomings in eddy activity in the SO3 simulations are not a product of variability within a  
264 single realization of Southern Ocean conditions (Supplementary Figure 2).



265

266 **Figure 2. Mean eddy kinetic energy between 2016 and 2020. a) The AWI-CM-1 ensemble. b) The first**

267 **member of the AWI-CM-1 ensemble. c) The SO3 simulation. d) The gridded satellite altimetry dataset.**

268

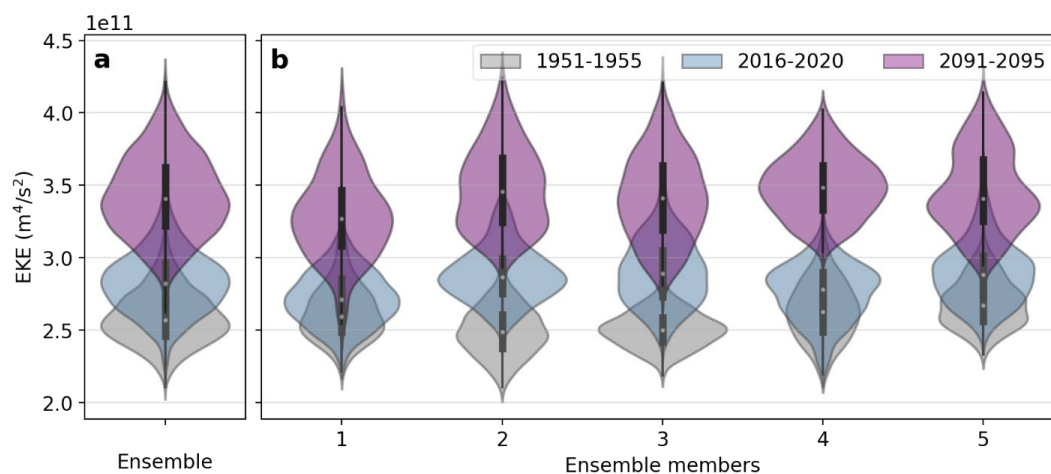


269 **3.2 EKE change and significance**

270 Southern Ocean eddy activity has been shown to intensify over the recent decades both using  
271 satellite altimetry (Martínez-Moreno et al., 2021), and the complete AWI-CM-1 dataset from CMIP6  
272 (Beech et al., 2022). Even after reducing the AWI-CM-1 CMIP6 dataset to five-year periods preceding  
273 the apparent change (1951-1955) and at the end of the altimetry era (2016-2020), this intensification is  
274 still discernable within the AWI-CM-1 ensemble (Figure 1a). Despite this, the SO3 simulations do not  
275 demonstrate any substantial change in magnitude over the same period (Figure 1). Further reducing the  
276 ensemble to its individual members, the EKE rise is still relatively clear in each case considering the  
277 median, mode, and first and third quartiles of the distributions (Figure 3), suggesting again, that natural or  
278 internal variability is not responsible for the discrepancies between simulations. Therefore, despite the  
279 first AWI-CM-1 ensemble member, from which the SO3 simulations receive their atmospheric forcing,  
280 producing lower-than-average EKE rise over this period (Figure 3), it is likely that other factors also  
281 contribute to the lack of change in the SO3 simulations.

282 The intensification of EKE becomes clear in both the AWI-CM-1 ensemble (Figure 1a), its  
283 members (Figure 3), and the SO3 simulations (Figure 1c) by the end of the 21<sup>st</sup> century. Over this period,  
284 the variability of EKE, indicated by the range of the distribution, also increases for each dataset (Figure  
285 1f, i). EKE rise in SO3 is approximately three times that of the AWI-CM1 ensemble in absolute terms  
286 (Figure 1f), but expressing EKE as a relative value normalized by the mean and standard deviation of  
287 each dataset during the observational period (Figure 1g, h, i), reveals greater consistency between the  
288 changes until the end of the 21<sup>st</sup> century. EKE in each dataset appears to increase by approximately 3  
289 standard deviations, and the range of EKE distributions increases by approximately one to two standard  
290 deviations (Figure 1h, i).

291



292

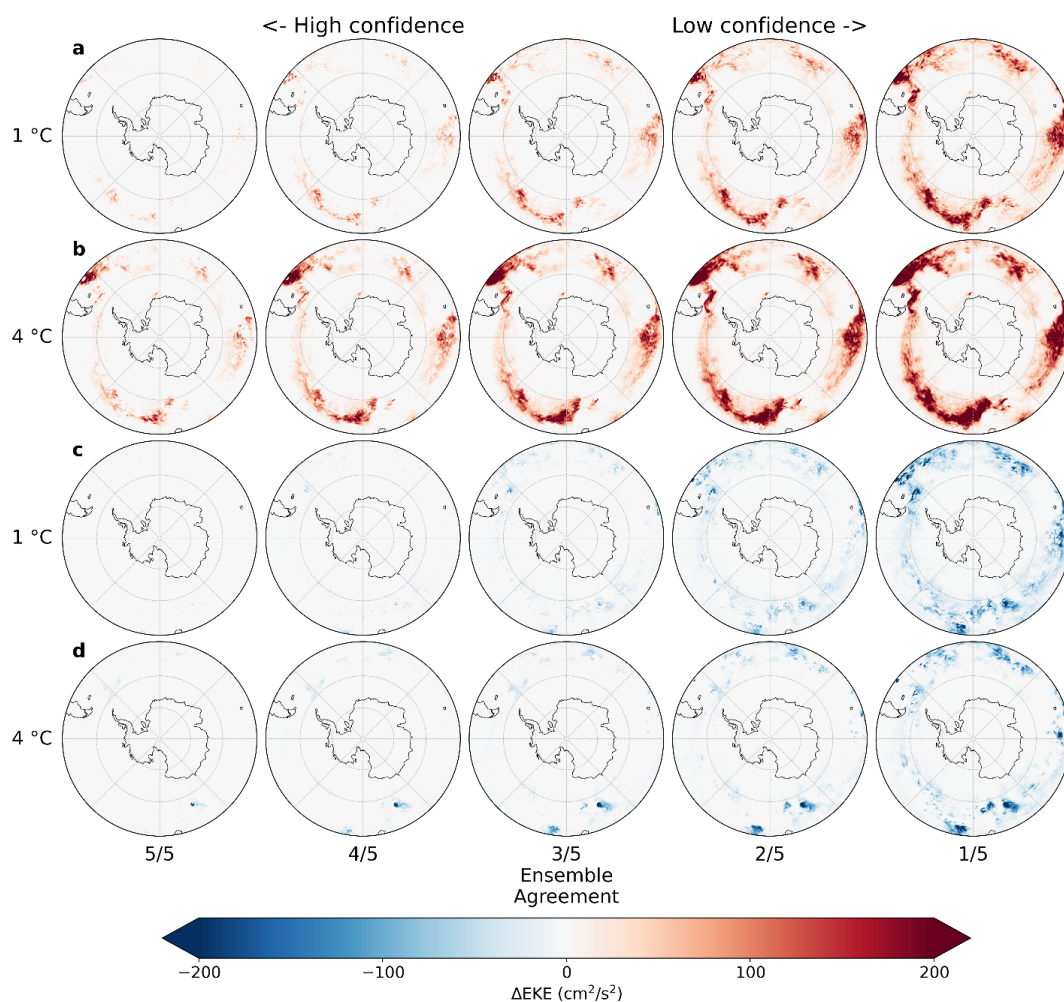
293 **Figure 3. Ensemble spread of EKE in AWI-CM-1. a)** Violin plots of mean Southern Ocean EKE in the  
294 AWI-CM-1 ensemble. **b)** Violin plots of mean Southern Ocean EKE in each member of the AWI-CM-1  
295 ensemble. Grey plots represent the period 1951-1955, blue plots represent 2016-2020, purple plots  
296 represent 2091-2095.

297 Before considering the regional impacts of warming on EKE in the SO3 simulations, it is useful  
298 to refer to the ensemble spread within the AWI-CM-1 simulations to approximate the reliability of a  
299 single ensemble member in revealing the ensemble-mean change as an analogue to the signal-to-noise  
300 ratio. At 1 °C of warming, EKE change in the ensemble is weak, with at least one ensemble member  
301 tending to show little or no EKE change in most regions (Figure 4a,c). Only a few clear patterns of  
302 change emerge throughout the ensemble, namely the regions of EKE intensification downstream of the  
303 Kerguelen Plateau and the Campbell Plateau where four to five out of five ensemble members show clear  
304 EKE intensification (Figure 4a). It should be noted that even in these regions of relatively high confidence  
305 (4 to 5 ensemble members, Figure 4a) EKE rise can be interspersed with lower-confidence (1 to 2  
306 ensemble members, Figure 4c) EKE decline; this is also illustrated by the ensemble mean changes  
307 themselves (Supplementary Figures 2, 3). Despite this, the consistency of EKE rise in these regions, and  
308 their geographic positions in already EKE-rich regions, suggests that the intensification patterns are





309 robust changes within substantial noise. This level of noise suggests that EKE changes in the SO3  
310 simulations at 1 °C of warming will be difficult to distinguish from natural variability when taken on their  
311 own; indeed, in the SO3 simulations, the large variability of both sign and magnitude of change within  
312 relatively small spatial scales does not lend confidence to any significant change at 1 °C of warming  
313 (Figure 5c). However, building on the changes observed in the AWI-CM-1 ensemble, the intensification  
314 of EKE downstream of the Kerguelen and Campbell Plateaus seems to be reinforced by the high-  
315 resolution simulations.



316



317 **Figure 4.** Ensemble agreement regarding EKE rise (a, b) and decline (c, d) within the AWI-CM-1  
318 ensemble after one (a, c) and four (b, d) °C of warming. Ensemble agreement refers to the number of  
319 ensemble members that simulate at least the pictured magnitude of EKE rise or decline for each grid cell.  
320 Rank 5/5 indicates the lowest magnitude of EKE rise or decline within the ensemble for a given grid cell,  
321 meaning the entire ensemble agrees on at least this much change. Rank 1/5 indicates the highest  
322 magnitude of EKE rise or decline within the ensemble for each grid cell, representing the upper limit of  
323 projected EKE change.

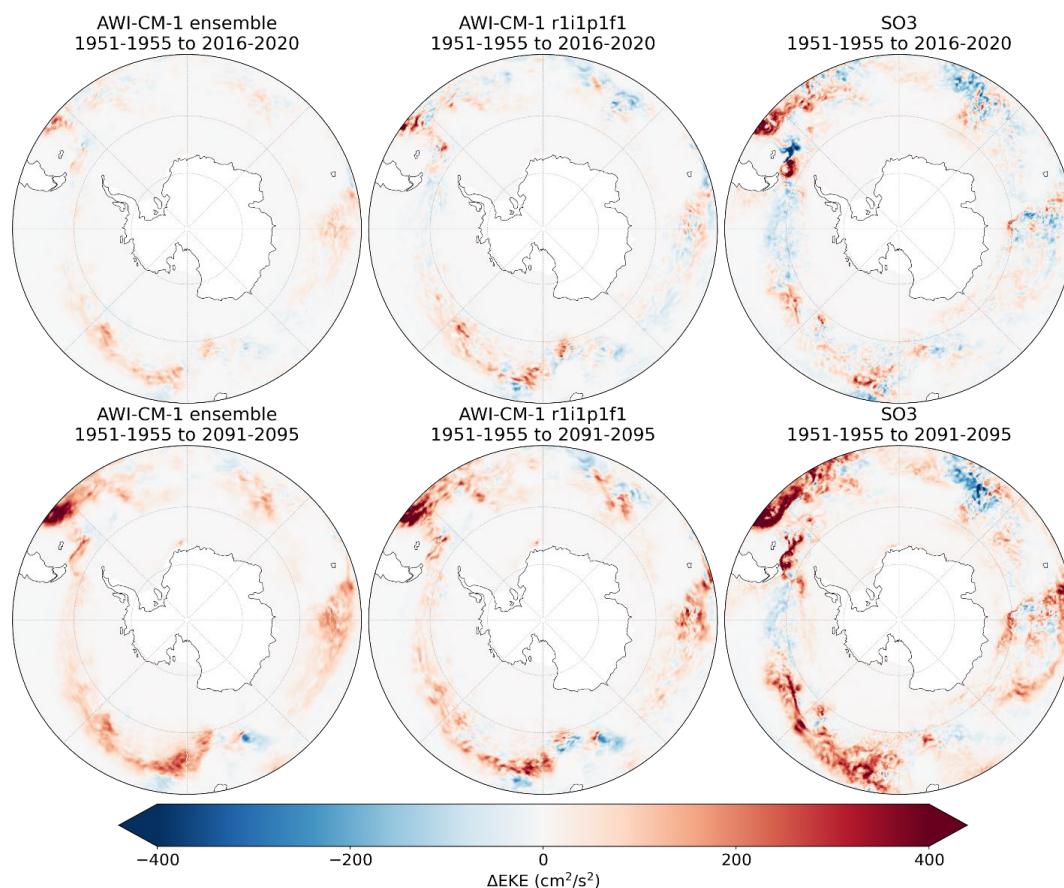
324

325 At 4 °C of warming, change in eddy activity becomes clearer; EKE intensification downstream of  
326 the Kerguelen and Campbell Plateaus is now consistent throughout the entire AWI-CM-1 ensemble, along  
327 with additional intensifications south of the Falkland/Malvinas Plateau, around the Conrad Rise, and  
328 along the Antarctic Slope Current at approximately 5 °E (Figure 4b). Four fifths of the ensemble also  
329 include a broad increase in EKE throughout the ACC across most longitudes. Interestingly, a consistent  
330 pattern of EKE decline also emerges upstream of the Campbell Plateau in the entire ensemble (Figure 4d).  
331 The spatial pattern of EKE rise is relatively consistent regardless of confidence, with only the magnitude  
332 increasing in the lower confidence composites (Figure 4b). The same tendency is observable between the  
333 EKE changes at 1 and 4 °C of warming, where the magnitude of change is greater after further warming  
334 but follows the same spatial pattern. Thus, regions of intensification can be identified more reliably than  
335 the magnitude of change and tend to be concentrated where flow interacts with topographic features, in  
336 already eddy-rich regions (Figure 2). Conversely, low confidence EKE decline appears nearly throughout  
337 the Southern Ocean in at least one ensemble member, but only consistently upstream of the Campbell  
338 Plateau and, to a far lesser extent, downstream of the Drake Passage and Campbell Plateau (Figure 4d).  
339 Changes of negative sign tend to be of lower magnitude at 4 °C of warming than at 1 °C. This suggests  
340 that in any given single ensemble simulation, a robust general intensification of EKE tends to occur but  
341 can be interspersed with spurious signals of decline. Yet, small regions of high-confidence EKE decline  
342 are also possible, however uncommon. Consequently, it would be difficult to confidently separate reliable



343 EKE change from natural variability in simulations without an ensemble to compare with. In the SO3  
344 simulations, EKE rise downstream of the Drake Passage and Kerguelen and Campbell Plateaus is  
345 substantial (Figure 5f). EKE rise is also projected south of the Falkland/Malvinas Plateau, around the  
346 Conrad Rise, and along the Antarctic Slope Current at approximately 5 °E, and a slight EKE decline  
347 appears upstream of the Campbell Plateau. All of this is comparable to the AWI-CM-1 ensemble, and the  
348 interspersed areas of EKE decline within these regions, for example, around the Conrad Rise, are not  
349 improbable based on the example set by AWI-CM-1 (Figure 4d). However, considering that some high-  
350 confidence EKE decline is present in the AWI-CM-1 ensemble, it is difficult to confidently dismiss  
351 regional EKE decline in the SO3 simulations as noise.

352



353



354 **Figure 5. EKE change.** Spatial representations of the difference in EKE between **(a-c)** 1951-1955 and  
355 2016-2020, **(d-f)** 1951-1955 and 2091-2095. **a,d)** The AWI-CM-1 ensemble. **b,e)** the first member of the  
356 AWI-CM-1 ensemble. **c,f)** The SO3 simulations.

357

#### 358 **4 Discussion**

359 Intensification of eddy activity in the Southern Ocean is now widely accepted as a consequence  
360 of anthropogenic climate change (Hogg et al., 2015; Patara et al., 2016; Martínez-Moreno et al., 2021;  
361 Beech et al., 2022), and is understood to be caused primarily by stronger westerly winds imparting more  
362 energy to the Antarctic Circumpolar Current (Munday et al., 2013; Marshall, 2003). The results presented  
363 here reinforce the notion of EKE intensification and further project increased EKE variability as the  
364 climate warms (Figure 1, 3). Understanding of regional changes within the Southern Ocean eddy field is  
365 generally limited to regions defined by oceanic sectors (Atlantic, Indian, Pacific) (Hogg et al., 2015), or  
366 incremental longitudinal delimitations (Patara et al., 2016). By expressing EKE change in terms of  
367 ensemble agreement on a cell-by-cell basis, the results presented here are able to identify regions of  
368 reliable and substantial change as those where flow interacts with major bathymetric features, and high  
369 eddy activity is already known to occur (Figure 4). In future research, regional analyses of both the  
370 significance and cause of EKE trends could focus on these regions to avoid confounding results with  
371 those of physically unrelated change within the chosen geographic delimitations.

372 Considering the consistency of the AWI-CM-1 ensemble both in terms of representing mean EKE  
373 distribution (Supplementary Figure 2), and projecting overall Southern Ocean EKE change (Figure 2), it  
374 appears that a larger ensemble of high-resolution simulations is not necessary to assess local mean states  
375 and changes over the entire region. This also means that the discrepancies between the SO3 simulations  
376 and AWI-CM-1 ensemble remain unexplained. The SO3 simulations failed to reproduce the clear  
377 intensification of EKE present in the ensemble at 1 °C of warming, which has already been identified in  
378 observations (Martínez-Moreno et al., 2021). One potential explanation for this is the atmospheric forcing



379 applied to the SO<sub>3</sub> simulations which cannot react to oceanic conditions in the same way that the real  
380 atmosphere or coupled atmosphere of the AWI-CM-1 ensemble can. Consequently, atmospheric forcing  
381 has been found to cause 27% more eddy killing by wind stress than coupled simulations (Renault et al.,  
382 2016). Since intensifying wind is widely considered the cause of EKE intensification in the Southern  
383 Ocean (Meredith and Hogg, 2006; Munday et al., 2013; Hogg et al., 2015; Beech et al., 2022), the  
384 increase may be counteracted by a simultaneous increase in eddy killing.

385         Assuming that the medium-resolution spin-up simulation effectively imparts a generally  
386 equilibrated ocean to the SO<sub>3</sub> simulations in terms of baroclinic and barotropic instability, the remaining  
387 discrepancies between eddy activity in SO<sub>3</sub> and observations should be largely identifiable. In terms of  
388 grid resolution, resolving the first Rossby radius of deformation with at least two grid points is not  
389 enough to comprehensively reproduce mesoscale activity (Hallberg, 2013; Sein et al., 2017), meaning  
390 grid refinement may need to be expanded to regions that impact eddy dynamics in the Southern Ocean.  
391 Other sources of bias may include ocean-atmosphere interactions which are absent or unrealistic within  
392 the uncoupled simulations (Byrne et al., 2016; Rai et al., 2021; Renault et al., 2016). As well, some small-  
393 scale, slow-to-equilibrate ocean processes may be resolved in the high-resolution simulations, but not be  
394 integrated long enough for their effects to impact eddy activity (van Westen and Dijkstra, 2021; Rackow  
395 et al., 2022). Finally, the gridded altimetry product itself may be responsible for some disagreement, as  
396 the along-track data is known to underrepresent eddy activity at scales less than 150km and 10 days  
397 (Chassignet and Xu, 2017), which will be particularly impactful at high latitudes.

398         To distinguish a meaningful signal of anthropogenic impacts from natural variability, this  
399 analysis relies primarily on consistency among ensemble members (Figures 3, 4). This is distinct from  
400 more traditional methods like assessment of error relative to observations or ensemble mean, commonly  
401 applied to weather forecasting (Ferro et al., 2012). Performance evaluation relative to observations would  
402 undoubtedly point to the high-resolution simulation as superior due to the drastic underrepresentation of  
403 EKE in the eddy-permitting ensemble (Figure 1). Yet, the effects of climate change are still apparent in



404 the AWI-CM-1 ensemble (Figure 1, 5), and the AWI-CM-1 dataset has been used to make similar  
405 projections of EKE already (Beech et al., 2022). Moreover, the eddy response to forcing seems to be  
406 consistent between the model resolutions when expressed in relative (Figure 1g, h, i), rather than absolute  
407 terms (Figure 1a, b, c). While more verification of this result is necessary both regionally, and with other  
408 models, these results suggest that eddy-permitting resolutions can be interpreted with their shortcomings  
409 in mind in order to discern the real-world implications: as is often necessary with model data. Thus, based  
410 on the test case of the Southern Ocean, the usefulness of the AWI-CM-1 ensemble and the effectiveness  
411 of model simulations in identifying physically significant and reproduceable impacts of climate change  
412 may be greater than would be identified using traditional methods and comes at a much lower cost  
413 relative to the eddy-resolving simulations.

414 This study has focused on EKE as an evaluation metric for the simulations since mesoscale  
415 activity is the primary motivation for increasing ocean model resolution. It has stopped short of assessing  
416 the improvements that resolving the mesoscale has on climate and ocean dynamics, many of which are  
417 discussed in detail elsewhere (eg. (Hewitt et al., 2017). Rather than repeat an assessment of the benefits of  
418 resolving smaller scales, we assume that the accurate reproduction and evolution of eddy activity  
419 indicates that these improvements are transferred to broader processes. Certainly, inaccurate simulation of  
420 the mesoscale would raise questions regarding the improvements that this mesoscale activity should have  
421 on the simulations as a whole. Nonetheless, further evaluation of the modeling approaches employed in  
422 this study will be necessary to determine if these methods are appropriate for studying broader elements  
423 of the climate system. Since the high-resolution simulations derive their deep-ocean climate primarily  
424 from the medium-resolution spin-up simulation, improving the initialization process (Thiria et al., 2023)  
425 may be the critical barrier to extending these results from the mixed layer to the deeper ocean.

## 426 **5 Conclusion**

427 Resolving the ocean mesoscale has become a focus for the climate and ocean modeling  
428 community as computational capabilities expand and models become increasingly complex. The benefits



429 that explicitly resolved eddy activity can have on climate simulations are clear (Hewitt et al., 2017; Sein  
430 et al., 2017) along with the impact that mesoscale variability has on local (Lachkar et al., 2009; Wang et  
431 al., 2017) and global environments (Falkowski et al., 1991; Sallée et al., 2012). However, state-of-the-art  
432 climate models will be unable to fully resolve the mesoscale for the foreseeable future, particularly in  
433 large-scale modeling endeavors such as CMIP (Hewitt et al., 2020). Thus, modelers must make informed  
434 choices regarding the explicit processes needed to answer research questions and where resources must be  
435 allocated to achieve specific goals. Existing analysis of resource allocation has typically addressed short-  
436 term weather forecasting or the ability to reproduce observations with low error (Ferro et al., 2012), but  
437 the question of how to best allocate resources for climate change impact assessment remains. This study  
438 has applied several cost-efficient modeling approaches to an analysis of the impacts of climate change on  
439 a key focus of high-resolution modeling: the mesoscale. Applying these results to broader climate change  
440 impact studies should improve the efficiency of resource allocation and focus modeling studies.  
441 Resolution can be dynamically adjusted both spatially, by focusing resources in study regions and where  
442 they are necessary to resolve local dynamics, and temporally, by allowing lower-resolution workhorse  
443 configurations to perform spin-up and transient runs. Limited simulation length and ensemble size can be  
444 sufficient for certain research questions and validation, but simulations must ultimately be designed to  
445 meet their specific goals. Where resources are limited, studies may best include a combination of eddy-  
446 resolving simulations able to fully capture the local eddy field, as well as eddy-present simulations that  
447 can attest to the significance of results through consistency and repetition.

448         This work represents a contribution to the growing wealth of research that points to an  
449 intensification of eddy activity in the Southern Ocean (Hogg et al., 2015; Martínez-Moreno et al., 2021;  
450 Beech et al., 2022). The further conclusions that EKE variability may increase and that EKE  
451 intensification appears concentrated in key regions based on topography can both expand the present state  
452 of knowledge, as well as direct future research. The cost-efficient modelling approaches of regional grid  
453 refinement, reduced-resolution spin-up and transient runs, and limited simulation lengths distinguished by



454 longer periods of change are demonstrated to be effective at reproducing change within a more traditional  
455 eddy-permitting ensemble. When resources are limited and resolution demands are high, these approaches  
456 can be adapted to address specific research questions. Where assessing the robustness of change is  
457 critical, the complimentary eddy-permitting ensemble represents an effective, low-cost supplement to the  
458 high-resolution simulations.

459

460

461





462 **Data Availability**

463 Geostrophic velocities derived from satellite altimetry data are publicly available at  
464 <https://doi.org/10.48670/moi-00148>. Daily sea surface height data from AWI-CM-1-1-MR in CMIP6  
465 used to compute geostrophic velocities in this study is archived at the World Data Center for Climate at  
466 the DKRZ (<https://doi.org/10.26050/WDCC/C6sCMAWAWM>,  
467 <https://doi.org/10.26050/WDCC/C6sSPAWAWM>) (Semmler et al., 2022a, b). Model output from AWI-  
468 CM-1-1-MR in the CMIP6 framework, including all variables used to force the standalone ocean  
469 simulations conducted for this study, is publicly available at <https://doi.org/10.22033/ESGF/CMIP6.359>  
470 (Semmler et al., 2018). Eddy kinetic energy datasets calculated from FESOM output velocities are  
471 available at (<https://doi.org/10.5281/zenodo.8046792>) (Beech, 2023b).

472 **Code Availability**

473 Source code for the ocean model FESOM version 2.5 is available at  
474 (<https://doi.org/10.5281/zenodo.7737061>) (patrickscholz et al., 2023). Code used for data analysis and  
475 visualization in this study is publicly available at (<https://doi.org/10.5281/zenodo.8046783>) (Beech,  
476 2023a). Code used to calculate geostrophic velocities from sea surface height data from AWI-CM-1-1-  
477 MR is available from <https://doi.org/10.5281/zenodo.7050573>.

478 **Author Contributions**

479 NB, TJ, TR, and TS conceived of the study. NB carried out the simulations, analyzed the data, and drafted  
480 the manuscript. All authors reviewed the manuscript.

481 **Competing Interests**

482 The authors declare no competing interests.

483 **Acknowledgements**



484           The work described in this paper has received funding from the Helmholtz Association through  
485 the project ‘Advanced Earth System Model Capacity’ (project leader: T.J., support code: ZT-0003) in the  
486 frame of the initiative ‘Zukunftsthemen’. The content of the paper is the sole responsibility of the authors  
487 and it does not represent the opinion of the Helmholtz Association, and the Helmholtz Association is not  
488 responsible for any use that might be made of information contained. TJ acknowledges the EERIE project  
489 funded under the EU Horizon Europe programme (grant number 101081383). TR acknowledges support  
490 from the European Commission’s Horizon 2020 collaborative project NextGEMS (grant number  
491 101003470). This work used resources of the Deutsches Klimarechenzentrum (DKRZ) granted by its  
492 Scientific Steering Committee (WLA) under project ID 995. The CMIP data used in this study were  
493 replicated and made available by the DKRZ.

494

495



## 496 **References**

- 497 Auger, M., Prandi, P., and Sallée, J.-B.: Southern ocean sea level anomaly in the sea ice-covered sector  
498 from multimission satellite observations, *Sci. Data*, 9, 70, <https://doi.org/10.1038/s41597-022-01166-z>,  
499 2022.
- 500 Auger, M., Sallée, J.-B., Thompson, A. F., Pauthenet, E., and Prandi, P.: Southern Ocean Ice-Covered Eddy  
501 Properties From Satellite Altimetry, *J. Geophys. Res. Oceans*, 128, e2022JC019363,  
502 <https://doi.org/10.1029/2022JC019363>, 2023.
- 503 Ballarotta, M., Ubelmann, C., Pujol, M.-I., Taburet, G., Fournier, F., Legeais, J.-F., Faugère, Y., Delepouille,  
504 A., Chelton, D., Dibarboure, G., and Picot, N.: On the resolutions of ocean altimetry maps, *Ocean Sci.*, 15,  
505 1091–1109, <https://doi.org/10.5194/os-15-1091-2019>, 2019.
- 506 Beech, N.: n-beech/SO\_meso: Initial release, <https://doi.org/10.5281/zenodo.8046783>, 2023a.
- 507 Beech, N.: Processed EKE from FESOM-SO3 (1.0), <https://doi.org/10.5281/zenodo.8046792>, 2023b.
- 508 Beech, N., Rackow, T., Semmler, T., Danilov, S., Wang, Q., and Jung, T.: Long-term evolution of ocean  
509 eddy activity in a warming world, *Nat. Clim. Change*, 12, 910–917, [https://doi.org/10.1038/s41558-022-](https://doi.org/10.1038/s41558-022-01478-3)  
510 01478-3, 2022.
- 511 Bishop, S. P., Gent, P. R., Bryan, F. O., Thompson, A. F., Long, M. C., and Abernathey, R.: Southern Ocean  
512 Overturning Compensation in an Eddy-Resolving Climate Simulation, *J. Phys. Oceanogr.*, 46, 1575–1592,  
513 <https://doi.org/10.1175/JPO-D-15-0177.1>, 2016.
- 514 Bronselaer, B., Winton, M., Griffies, S. M., Hurlin, W. J., Rodgers, K. B., Sergienko, O. V., Stouffer, R. J.,  
515 and Russell, J. L.: Change in future climate due to Antarctic meltwater, *Nature*, 564, 53–58,  
516 <https://doi.org/10.1038/s41586-018-0712-z>, 2018.
- 517 Byrne, D., Münnich, M., Frenger, I., and Gruber, N.: Mesoscale atmosphere ocean coupling enhances the  
518 transfer of wind energy into the ocean, *Nat. Commun.*, 7, ncomms11867,  
519 <https://doi.org/10.1038/ncomms11867>, 2016.
- 520 Chassignet, E. P. and Xu, X.: Impact of Horizontal Resolution ( $1/12^\circ$  to  $1/50^\circ$ ) on Gulf Stream Separation,  
521 Penetration, and Variability, *J. Phys. Oceanogr.*, 47, 1999–2021, [https://doi.org/10.1175/JPO-D-17-](https://doi.org/10.1175/JPO-D-17-0031.1)  
522 0031.1, 2017.
- 523 Danilov, S.: Ocean modeling on unstructured meshes, *Ocean Model.*, 69, 195–210,  
524 <https://doi.org/10.1016/j.ocemod.2013.05.005>, 2013.
- 525 Danilov, S., Sidorenko, D., Wang, Q., and Jung, T.: The Finite-volumE Sea ice–Ocean Model (FESOM2),  
526 *Geosci. Model Dev.*, 10, 765–789, <https://doi.org/10.5194/gmd-10-765-2017>, 2017.
- 527 Eyring, V., Bony, S., Meehl, G. A., Senior, C. A., Stevens, B., Stouffer, R. J., and Taylor, K. E.: Overview of  
528 the Coupled Model Intercomparison Project Phase 6 (CMIP6) experimental design and organization,  
529 *Geosci. Model Dev.*, 9, 1937–1958, <https://doi.org/10.5194/gmd-9-1937-2016>, 2016.
- 530 Falkowski, P. G., Ziemann, D., Kolber, Z., and Bienfang, P. K.: Role of eddy pumping in enhancing primary  
531 production in the ocean, *Nature*, 352, 55–58, <https://doi.org/10.1038/352055a0>, 1991.



- 532 Ferrari, R., Griffies, S. M., Nurser, A. J. G., and Vallis, G. K.: A boundary-value problem for the  
533 parameterized mesoscale eddy transport, *Ocean Model.*, 32, 143–156,  
534 <https://doi.org/10.1016/j.ocemod.2010.01.004>, 2010.
- 535 Ferro, C. A. T., Jupp, T. E., Lambert, F. H., Huntingford, C., and Cox, P. M.: Model complexity versus  
536 ensemble size: allocating resources for climate prediction, *Philos. Trans. R. Soc. Math. Phys. Eng. Sci.*,  
537 370, 1087–1099, <https://doi.org/10.1098/rsta.2011.0307>, 2012.
- 538 Frölicher, T. L., Sarmiento, J. L., Paynter, D. J., Dunne, J. P., Krasting, J. P., and Winton, M.: Dominance of  
539 the Southern Ocean in Anthropogenic Carbon and Heat Uptake in CMIP5 Models, *J. Clim.*, 28, 862–886,  
540 <https://doi.org/10.1175/JCLI-D-14-00117.1>, 2015.
- 541 Gent, P. R. and McWilliams, J. C.: Isopycnal Mixing in Ocean Circulation Models, *J. Phys. Oceanogr.*, 20,  
542 150–155, [https://doi.org/10.1175/1520-0485\(1990\)020<0150:IMIOCM>2.0.CO;2](https://doi.org/10.1175/1520-0485(1990)020<0150:IMIOCM>2.0.CO;2), 1990.
- 543 Haarsma, R. J., Roberts, M. J., Vidale, P. L., Senior, C. A., Bellucci, A., Bao, Q., Chang, P., Corti, S., Fučkar,  
544 N. S., Guemas, V., von Hardenberg, J., Hazeleger, W., Kodama, C., Koenigk, T., Leung, L. R., Lu, J., Luo, J.-  
545 J., Mao, J., Mizielinski, M. S., Mizuta, R., Nobre, P., Satoh, M., Scoccimarro, E., Semmler, T., Small, J., and  
546 von Storch, J.-S.: High Resolution Model Intercomparison Project (HighResMIP v1.0) for CMIP6, *Geosci.*  
547 *Model Dev.*, 9, 4185–4208, <https://doi.org/10.5194/gmd-9-4185-2016>, 2016.
- 548 Hallberg, R.: Using a resolution function to regulate parameterizations of oceanic mesoscale eddy  
549 effects, *Ocean Model.*, 72, 92–103, <https://doi.org/10.1016/j.ocemod.2013.08.007>, 2013.
- 550 Hewitt, H. T., Bell, M. J., Chassignet, E. P., Czaja, A., Ferreira, D., Griffies, S. M., Hyder, P., McClean, J. L.,  
551 New, A. L., and Roberts, M. J.: Will high-resolution global ocean models benefit coupled predictions on  
552 short-range to climate timescales?, *Ocean Model.*, 120, 120–136,  
553 <https://doi.org/10.1016/j.ocemod.2017.11.002>, 2017.
- 554 Hewitt, H. T., Roberts, M., Mathiot, P., Biastoch, A., Blockley, E., Chassignet, E. P., Fox-Kemper, B., Hyder,  
555 P., Marshall, D. P., Popova, E., Treguier, A.-M., Zanna, L., Yool, A., Yu, Y., Beadling, R., Bell, M., Kuhlbrodt,  
556 T., Arsouze, T., Bellucci, A., Castruccio, F., Gan, B., Putrasahan, D., Roberts, C. D., Van Roekel, L., and  
557 Zhang, Q.: Resolving and Parameterising the Ocean Mesoscale in Earth System Models, *Curr. Clim.*  
558 *Change Rep.*, 6, 137–152, <https://doi.org/10.1007/s40641-020-00164-w>, 2020.
- 559 Hogg, A. McC., Meredith, M. P., Chambers, D. P., Abrahamsen, E. P., Hughes, C. W., and Morrison, A. K.:  
560 Recent trends in the Southern Ocean eddy field, *J. Geophys. Res. Oceans*, 120, 257–267,  
561 <https://doi.org/10.1002/2014JC010470>, 2015.
- 562 Irving, D., Hobbs, W., Church, J., and Zika, J.: A Mass and Energy Conservation Analysis of Drift in the  
563 CMIP6 Ensemble, *J. Clim.*, 34, 3157–3170, <https://doi.org/10.1175/JCLI-D-20-0281.1>, 2021.
- 564 Jungclaus, J. H., Lorenz, S. J., Schmidt, H., Brovkin, V., Brüggemann, N., Chegini, F., Crüger, T., De-Vrese,  
565 P., Gayler, V., Giorgetta, M. A., Gutjahr, O., Haak, H., Hagemann, S., Hanke, M., Ilyina, T., Korn, P.,  
566 Kröger, J., Linardakis, L., Mehlmann, C., Mikolajewicz, U., Müller, W. A., Nabel, J. E. M. S., Notz, D.,  
567 Pohlmann, H., Putrasahan, D. A., Raddatz, T., Ramme, L., Redler, R., Reick, C. H., Riddick, T., Sam, T.,  
568 Schneck, R., Schnur, R., Schupfner, M., von Storch, J.-S., Wachsmann, F., Wieners, K.-H., Ziemens, F.,  
569 Stevens, B., Marotzke, J., and Claussen, M.: The ICON Earth System Model Version 1.0, *J. Adv. Model.*  
570 *Earth Syst.*, 14, e2021MS002813, <https://doi.org/10.1029/2021MS002813>, 2022.



- 571 Lachkar, Z., Orr, J. C., Dutay, J. C., and Delecluse, P.: On the role of mesoscale eddies in the ventilation of  
572 Antarctic intermediate water, *Deep-Sea Res. Part Oceanogr. Res. Pap.*, 56, 909–925,  
573 <https://doi.org/10.1016/j.dsr.2009.01.013>, 2009.
- 574 Landschützer, P., Gruber, N., Haumann, F. A., Rödenbeck, C., Bakker, D. C. E., van Heuven, S., Hoppema,  
575 M., Metzl, N., Sweeney, C., Takahashi, T., Tilbrook, B., and Wanninkhof, R.: The reinvigoration of the  
576 Southern Ocean carbon sink, *Science*, 349, 1221–1224, <https://doi.org/10.1126/science.aab2620>, 2015.
- 577 Large, W. G., McWilliams, J. C., and Doney, S. C.: Oceanic vertical mixing: A review and a model with a  
578 nonlocal boundary layer parameterization, *Rev. Geophys.*, 32, 363–403,  
579 <https://doi.org/10.1029/94RG01872>, 1994.
- 580 Marshall, G. J.: Trends in the Southern Annular Mode from Observations and Reanalyses, *J. Clim.*, 16,  
581 4134–4143, [https://doi.org/10.1175/1520-0442\(2003\)016<4134:TITSAM>2.0.CO;2](https://doi.org/10.1175/1520-0442(2003)016<4134:TITSAM>2.0.CO;2), 2003.
- 582 Martínez-Moreno, J., Hogg, A. McC., England, M. H., Constantinou, N. C., Kiss, A. E., and Morrison, A. K.:  
583 Global changes in oceanic mesoscale currents over the satellite altimetry record, *Nat. Clim. Change*, 11,  
584 397–403, <https://doi.org/10.1038/s41558-021-01006-9>, 2021.
- 585 Marzocchi, A., Hirschi, J. J.-M., Holliday, N. P., Cunningham, S. A., Blaker, A. T., and Coward, A. C.: The  
586 North Atlantic subpolar circulation in an eddy-resolving global ocean model, *J. Mar. Syst.*, 142, 126–143,  
587 <https://doi.org/10.1016/j.jmarsys.2014.10.007>, 2015.
- 588 Meredith, M. P. and Hogg, A. M.: Circumpolar response of Southern Ocean eddy activity to a change in  
589 the Southern Annular Mode, *Geophys. Res. Lett.*, 33, <https://doi.org/10.1029/2006GL026499>, 2006.
- 590 Munday, D. R., Johnson, H. L., and Marshall, D. P.: Eddy Saturation of Equilibrated Circumpolar Currents,  
591 *J. Phys. Oceanogr.*, 43, 507–532, <https://doi.org/10.1175/JPO-D-12-095.1>, 2013.
- 592 O’Neill, B. C., Tebaldi, C., van Vuuren, D. P., Eyring, V., Friedlingstein, P., Hurtt, G., Knutti, R., Kriegler, E.,  
593 Lamarque, J.-F., Lowe, J., Meehl, G. A., Moss, R., Riahi, K., and Sanderson, B. M.: The Scenario Model  
594 Intercomparison Project (ScenarioMIP) for CMIP6, *Geosci. Model Dev.*, 9, 3461–3482,  
595 <https://doi.org/10.5194/gmd-9-3461-2016>, 2016.
- 596 O’Neill, B. C., Kriegler, E., Ebi, K. L., Kemp-Benedict, E., Riahi, K., Rothman, D. S., van Ruijven, B. J., van  
597 Vuuren, D. P., Birkmann, J., Kok, K., Levy, M., and Solecki, W.: The roads ahead: Narratives for shared  
598 socioeconomic pathways describing world futures in the 21st century, *Glob. Environ. Change*, 42, 169–  
599 180, <https://doi.org/10.1016/j.gloenvcha.2015.01.004>, 2017.
- 600 Patara, L., Böning, C. W., and Biastoch, A.: Variability and trends in Southern Ocean eddy activity in 1/12°  
601 ocean model simulations, *Geophys. Res. Lett.*, 43, 4517–4523, <https://doi.org/10.1002/2016GL069026>,  
602 2016.
- 603 patrickscholz, dsidoren, Koldunov, N., Hegewald, J., rakowsk, Streffing, J., Rackow, T., ogurses,  
604 helgegoessling, Guibert, D., qiangclimate, cwekerle, Oord, G. van den, jrberlin, Miguel, Gierz, P., and  
605 Cheedela, S. K.: FESOM/fesom2: FESOM2.5, , <https://doi.org/10.5281/zenodo.7737061>, 2023.



- 606 Pauling, A. G., Smith, I. J., Langhorne, P. J., and Bitz, C. M.: Time-Dependent Freshwater Input From Ice  
607 Shelves: Impacts on Antarctic Sea Ice and the Southern Ocean in an Earth System Model, *Geophys. Res.*  
608 *Let.*, 44, 10,454–10,461, <https://doi.org/10.1002/2017GL075017>, 2017.
- 609 Rackow, T., Danilov, S., Goessling, H. F., Hellmer, H. H., Sein, D. V., Semmler, T., Sidorenko, D., and Jung,  
610 T.: Delayed Antarctic sea-ice decline in high-resolution climate change simulations, *Nat. Commun.*, 13,  
611 637, <https://doi.org/10.1038/s41467-022-28259-y>, 2022.
- 612 Rai, S., Hecht, M., Maltrud, M., and Aluie, H.: Scale of oceanic eddy killing by wind from global satellite  
613 observations, *Sci. Adv.*, 7, eabf4920, <https://doi.org/10.1126/sciadv.abf4920>, 2021.
- 614 Renault, L., Molemaker, M. J., McWilliams, J. C., Shchepetkin, A. F., Lemarié, F., Chelton, D., Illig, S., and  
615 Hall, A.: Modulation of Wind Work by Oceanic Current Interaction with the Atmosphere, *J. Phys.*  
616 *Oceanogr.*, 46, 1685–1704, <https://doi.org/10.1175/JPO-D-15-0232.1>, 2016.
- 617 Ringler, T., Petersen, M., Higdon, R. L., Jacobsen, D., Jones, P. W., and Maltrud, M.: A multi-resolution  
618 approach to global ocean modeling, *Ocean Model.*, 69, 211–232,  
619 <https://doi.org/10.1016/j.ocemod.2013.04.010>, 2013.
- 620 Sallée, J.-B., Matear, R. J., Rintoul, S. R., and Lenton, A.: Localized subduction of anthropogenic carbon  
621 dioxide in the Southern Hemisphere oceans, *Nat. Geosci.*, 5, 579–584,  
622 <https://doi.org/10.1038/ngeo1523>, 2012.
- 623 Scholz, P., Sidorenko, D., Gurses, O., Danilov, S., Koldunov, N., Wang, Q., Sein, D., Smolentseva, M.,  
624 Rakowsky, N., and Jung, T.: Assessment of the Finite-volume Sea ice-Ocean Model (FESOM2.0) – Part 1:  
625 Description of selected key model elements and comparison to its predecessor version, *Geosci. Model*  
626 *Dev.*, 12, 4875–4899, <https://doi.org/10.5194/gmd-12-4875-2019>, 2019.
- 627 Scholz, P., Sidorenko, D., Danilov, S., Wang, Q., Koldunov, N., Sein, D., and Jung, T.: Assessment of the  
628 Finite Volume Sea Ice Ocean Model (FESOM2.0), Part II: Partial bottom cells, embedded sea ice and  
629 vertical mixing library CVMIX, *Climate and Earth system modeling*, [https://doi.org/10.5194/gmd-2021-](https://doi.org/10.5194/gmd-2021-94)  
630 94, 2021.
- 631 Sein, D. V., Koldunov, N. V., Danilov, S., Wang, Q., Sidorenko, D., Fast, I., Rackow, T., Cabos, W., and Jung,  
632 T.: Ocean Modeling on a Mesh With Resolution Following the Local Rossby Radius, *J. Adv. Model. Earth*  
633 *Syst.*, 9, 2601–2614, <https://doi.org/10.1002/2017MS001099>, 2017.
- 634 Semmler, T., Danilov, S., Rackow, T., Sidorenko, D., Barbi, D., Hegewald, J., Sein, D., Wang, Q., and Jung,  
635 T.: AWI AWI-CM1.1MR model output prepared for CMIP6 CMIP,  
636 <https://doi.org/10.22033/ESGF/CMIP6.359>, 2018.
- 637 Semmler, T., Danilov, S., Gierz, P., Goessling, H. F., Hegewald, J., Hinrichs, C., Koldunov, N., Khosravi, N.,  
638 Mu, L., Rackow, T., Sein, D. V., Sidorenko, D., Wang, Q., and Jung, T.: Simulations for CMIP6 With the  
639 AWI Climate Model AWI-CM-1-1, *J. Adv. Model. Earth Syst.*, 12, e2019MS002009,  
640 <https://doi.org/10.1029/2019MS002009>, 2020.
- 641 Semmler, T., Danilov, S., Rackow, T., Sidorenko, D., Barbi, D., Hegewald, J., Sein, D., Wang, Q., and Jung,  
642 T.: CMIP6\_supplemental CMIP AWI AWI-CM-1-1-MR, <https://doi.org/10.26050/WDC/C6sCMAWAWM>,  
643 2022a.



- 644 Semmler, T., Danilov, S., Rackow, T., Sidorenko, D., Barbi, D., Hegewald, J., Sein, D., Wang, Q., and Jung,  
645 T.: CMIP6\_supplemental ScenarioMIP AWI AWI-CM-1-1-MR,  
646 <https://doi.org/10.26050/WDCC/C6sSPAWAWM>, 2022b.
- 647 Taburet, G., Sanchez-Roman, A., Ballarotta, M., Pujol, M.-I., Legeais, J.-F., Fournier, F., Faugere, Y., and  
648 Dibarboure, G.: DUACS DT2018: 25 years of reprocessed sea level altimetry products, *Ocean Sci.*, 15,  
649 1207–1224, <https://doi.org/10.5194/os-15-1207-2019>, 2019.
- 650 Thiria, S., Sorrow, C., Archambault, T., Charantonis, A., Bereziat, D., Mejia, C., Molines, J.-M., and Crépon,  
651 M.: Downscaling of ocean fields by fusion of heterogeneous observations using Deep Learning  
652 algorithms, *Ocean Model.*, 182, 102174, <https://doi.org/10.1016/j.ocemod.2023.102174>, 2023.
- 653 Wang, Q., Danilov, S., Sidorenko, D., Timmermann, R., Wekerle, C., Wang, X., Jung, T., and Schröter, J.:  
654 The Finite Element Sea Ice-Ocean Model (FESOM) v.1.4: formulation of an ocean general circulation  
655 model, *Geosci. Model Dev.*, 7, 663–693, <https://doi.org/10.5194/gmd-7-663-2014>, 2014.
- 656 Wang, Y., Claus, M., Greatbatch, R. J., and Sheng, J.: Decomposition of the Mean Barotropic Transport in  
657 a High-Resolution Model of the North Atlantic Ocean, *Geophys. Res. Lett.*, 44, 11,537–11,546,  
658 <https://doi.org/10.1002/2017GL074825>, 2017.
- 659 van Westen, R. M. and Dijkstra, H. A.: Ocean eddies strongly affect global mean sea-level projections,  
660 *Sci. Adv.*, 7, eabf1674, <https://doi.org/10.1126/sciadv.abf1674>, 2021.
- 661 Yu, X., Ponte, A. L., Lahaye, N., Caspar-Cohen, Z., and Menemenlis, D.: Geostrophy Assessment and  
662 Momentum Balance of the Global Oceans in a Tide- and Eddy-Resolving Model, *J. Geophys. Res. Oceans*,  
663 126, e2021JC017422, <https://doi.org/10.1029/2021JC017422>, 2021.
- 664