



# Quantifying the effects of the microphysical properties of black carbon on the determination of brown carbon using measurements at multiple wavelengths

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**Abstract.** The absorption Ångström exponent (AAE)-based methods are widely used to estimate brown carbon (BrC) absorption, and the estimated BrC absorption can be significantly different from 0 even for pure black carbon (BC). However, few studies have systematically quantified the effects of BC microphysical properties. Moreover, it is still unclear under which conditions the AAE-based method is applicable. In this work, we used BC models partially coated with non-absorbing materials to calculate the total absorption. Since the total absorption is entirely from BC, the estimated BrC absorption should be 0 if the retrieval methods are accurate. Thus, the estimated BrC absorption ( $ABS_{BrC}$ ) should be the absorption from BC that is incorrectly attributed to BrC. The results show that a BC AAE of 1 can generally provide reasonable estimates for freshly emitted BC, since at this time  $ABS_{BrC}$  is generally in the range of -3% to 4.5%. However, when BC aerosols are aged,  $ABS_{BrC}$  of about 35% could be observed. The WDA method does not necessarily improve the estimates, sometimes a negative  $ABS_{BrC}$  of -40% can be found for partially coated BC. By combining simulations of a global chemical transport model, this work also quantified the effects of BC microphysical properties on BrC global optical absorption aerosol depth (AAOD) estimates. The AAE = 1 method could sometimes lead to a misassigned global mean AAOD of about  $-0.4 - 0.5 \times 10^{-3}$  if BC aerosols have a complex morphology, leading to a global mean direct radiation factor (DRF) of about  $-0.068 \pm 0.0172$  to  $+0.085 \pm 0.0215$  W/m<sup>2</sup> from BC, which is incorrectly assigned to BrC. The WDA method does not necessarily improve the estimates. In our cases, the WDA methods based on the spherical models can lead to a range of about  $-0.9 - 0.05 \times 10^{-3}$  of misassigned AAOD, which could lead to a global mean DRF error range of  $-0.153 \pm 0.0387$  to  $+0.0085 \pm 0.0022$  W/m<sup>2</sup>. At the regional scale, the AAE = 1 method in East Asia sometimes leads to a distributed AAOD of over  $3 \times 10^{-3}$ , resulting in a BC DRF of about  $+0.51 \pm 0.129$  W/m<sup>2</sup>, which is incorrectly attributed to BrC. Mie theory-based WDA methods would lead to an estimated AAOD error of more than  $6 \times 10^{-3}$  in some regions (e.g., East Asia), resulting in an estimated misattributed DRF of  $+1.0 \pm 0.258$  W/m<sup>2</sup>.



## 1 Introduction

Carbonaceous aerosols are a major contributor to climate (Bond et al., 2013; Myhre et al., 2013). Black carbon (BC) and organic carbon (OC) are the most important carbonaceous aerosols in the atmosphere. As the most important absorbing aerosol in the atmosphere, BC significantly absorbs solar radiation from UV to visible wavelengths (Cross et al., 2010; Petzold and Schönlinner, 2004; Bond and Bergstrom, 2006). Even at small mass fractions of total atmospheric aerosols, the radiative effects of BC are quite significant due to high absorption, which greatly enhances global warming (Zhang and Wang, 2011; Bond et al., 2013; Matsui et al., 2018). In contrast, OC was initially considered as a mere scattering agent that has a cooling effect on the Earth-atmosphere system (Wang et al., 2016). However, recent studies have shown that some OC aerosols also absorb light from UV to visible wavelengths (Kirchstetter et al., 2004; Chakrabarty et al., 2010; Chen and Bond, 2010). These OC aerosols are known as brown carbon (BrC) (Andreae and Gelencsér, 2006; Laskin et al., 2015). Recent studies have also shown that BrC also exerts considerable positive radiative forcing (Zeng et al., 2020b; Feng et al., 2013). In some specific regions, the direct radiative effect of BrC is even comparable to that of BC (Zhang et al., 2020a). Therefore, understanding BrC absorption is important for studying global climate change.

However, the understanding of BrC absorption is still quite limited. BrC strongly absorbs light in the UV region, while its absorption strongly decreases with increasing wavelength from the UV region to the visible region (Hecobian et al., 2010; Kirchstetter et al., 2004; Chakrabarty et al., 2010; Chen and Bond, 2010; Bahadur et al., 2012). Filter samples in laboratory measurements are the main method for measuring the absorption characteristics of BrC (Chen and Bond, 2010; Xie et al., 2019). Based on laboratory studies, previous studies have shown that the mass cross section (MAC) of BrC varies over a wide range and that the reported values in different studies are different (Dasari et al., 2019; Kirillova et al., 2016). The reasons for the large uncertainties are mainly due to the samples from different regions and different measurement conditions. Compared with BC, the absorption of BrC is subject to larger uncertainties due to the type of fuel, combustion condition, aging condition, etc. Therefore, it is difficult to make accurate estimates of BrC absorption based on laboratory measurements in a global range because it is unrealistic to perform filter sampling under all conditions. Numerical modeling is another effective tool for estimating BrC absorption. However, modeling is usually based on prior knowledge from laboratory measurements, which varies by region, weather conditions, aging condition, etc. Therefore, even with modeling, it is difficult to always provide accurate estimates.

Remote sensing can provide regional/global measurements and is an effective complementary method to address the above issues. Recently, researchers have attempted to derive the absorption of BrC based on the absorption of multiple wavelengths from remote sensing (Arola et al., 2011; Tesche et al., 2011; Wang et al., 2016). Ground-based remote sensing, such as the Aerosol Robotic Network (AERONET) (Holben et al., 1998; Giles et al., 2019), could provide measurements of the temporal optical absorption depth (AAOD) on a global scale (Shaw, 1983; Shin et al., 2019). In addition, polarimetric satellite measurements have also been used to determine AAOD (Dubovik et al., 2011, 2014). With the increasing number of satellites, it is expected that the absorbing aerosols on a global scale can be detected in real time in the future through the cooperation of satellite constellations. However, the absorption derived from remote sensing is the mixture of different absorbing aerosols.



55 Dust, BrC, and BC are the main absorbing aerosols in the atmosphere, and we usually need to separate the contributions of dust and BC to study the absorption of BrC. Since dust particles are usually aerosols of small size, we can break down the contribution of dust based on the size data. Both BC and BrC are fine aerosols, and it is difficult to separate them based on size information alone.

An effective technique for separating the contributions of BC is based on the different spectral absorption dependence of BC and BrC (Wang et al., 2016, 2018; Russell et al., 2010; Chung et al., 2012). The absorption of BrC is generally very weak at long visible wavelengths and in the near infrared, and it has been generally assumed that the total absorption is entirely from BC (Wang et al., 2016; Luo et al., 2021c). Since the spectral dependence of the absorption of BC is subject to relatively small uncertainties, the absorption of BC in the UV region can be derived from the absorption in the long visible and near-infrared regions. Then, the absorption of BrC in the UV region could be estimated from the difference between the total absorption and the BC absorption. At two reference wavelengths ( $\lambda_1$ ,  $\lambda_2$ ), the spectral dependence of absorption is generally represented by a parameter, the absorption Ångström exponent (AAE):

$$AAE = \frac{\ln(b_{\text{abs}(\lambda_1)}) - \ln(b_{\text{abs}(\lambda_2)})}{\ln(\lambda_1) - \ln(\lambda_2)} \quad (1)$$

where  $b_{\text{abs}(\lambda_1)}$  and  $b_{\text{abs}(\lambda_2)}$  are the absorption coefficients at  $\lambda_1$  and  $\lambda_2$ , respectively.

Using the BC AAE value, the absorption of BC in the UV region can be estimated based on the absorption at near-infrared wavelengths. When estimating BrC absorption, a BC AAE of 1 has often been used, while more recent studies based on measurements and simulations have shown a wide range of AAE values. Therefore, the use of BC AAE of 1 may not always provide accurate estimates. To correct the AAE = 1 method, Wang et al. (2016) proposed to use the wavelength dependence of AAE (WDA) to improve the estimates. However, in their studies, the WDA was calculated using the Mie theory, which assumes that the morphology of BC is spherical. However, in the atmosphere, BC exhibits a rather complex morphology. Luo et al. (2021c) have shown that the WDA method does not necessarily lead to better estimates when BC has a complex morphology. Thus, the estimation of BrC based on absorption measurements at multiple wavelengths suffers from uncertainties in BC properties.

It is well known that the AAE method does not provide accurate results. However, there is a lack of understanding of the uncertainties caused by the microphysical properties of BC. Using morphologically realistic models, Luo et al. (2021c) showed that the estimates of BrC-based AAE methods are significantly affected by BC morphology, while how the uncertainties are affected by BC microphysical properties was not investigated. Moreover, in their work only single particles are considered, while bulk particles exist in the atmosphere. In this work, we attempt to systematically quantify the impact of BC microphysical properties on the estimates of BrC-based AAE methods. We focus on answering the following questions:

- How large are the uncertainties caused by the microphysical properties of BC in the estimation of BrC?
- How do the microphysical properties of BC affect estimates of BrC?
- What implications can we obtain for estimates of BrC on a global scale?



To answer the above questions, we generated some "realistic" BC aerosols based on the partially coated models and then estimated the absorption fraction of BC incorrectly attributed to BrC ( $ABS_{BrC}$ ) using different AAE methods and investigated the effects of BC microphysical properties on  $ABS_{BrC}$ . Finally, assuming a typical size distribution, we investigated the global distribution of BC AAOD on BrC at different BC morphologies and mixing states based on a global chemical transport model, the Goddard Earth Observing System with chemistry (GEOS-Chem). Our study can improve the understanding of uncertainties in the estimation of BrC based on absorption measurements at multiple wavelengths.

## 2 Estimating the BrC absorption

### 2.1 BC morphological model

In calculating the optical properties, the shape of BC was largely assumed to be spherical, so we can calculate the optical properties using Mie theory (Mie, 1908; Lack and Cappa, 2010; Bao et al., 2019). However, BC in the atmosphere usually has a rather complex morphology (China et al., 2015; Adachi et al., 2010b; Yuan et al., 2019; Wang et al., 2017; Luo et al., 2021a). When BC is freshly emitted, the morphology is usually chain-like and consists of numerous spherical particles (Sorensen, 2001, 2011). Researchers have often used the fractal law to describe the morphology (Sorensen, 2001; Heinson et al., 2017; Luo et al., 2021a, b):

$$N_s = k_0 \left( \frac{R_g}{R} \right)^{D_f} \quad (2)$$

where  $N_s$  is the number of spherical monomers;  $k_0$  and  $D_f$  are two parameters representing the symmetry and compactness of the BC aggregates and are called fractal prefactor and fractal dimension, respectively;  $R$  and  $R_g$  represent the radius of the small particles and the gyration radius, respectively.

$D_f$  is often used to describe the compactness of BC (Liu and Mishchenko, 2005; Radney et al., 2014). BC generally exhibits a fluffy morphology when emitted into the atmosphere. Previous studies have shown that a small  $D_f$  can strongly affect the morphology of BC. For BC aerosols from biomass combustion, a  $D_f$  range of 1.67 – 1.83 was observed (Chakrabarty et al., 2006); BC aerosols from vehicle emissions showed a  $D_f$  range of 1.52 – 1.94 (China et al., 2014); Wentzel et al. (2003) showed a  $D_f$  of 1.6 – 1.9 for fresh BC from vehicle emissions. From a robotics perspective, previous studies have shown that simulations based on the diffusion-limited cluster-cluster aggregation (DLCA) algorithm can explain the measurements well, and a universal  $D_f$  of about 1.8 was obtained by the simulations (Sorensen, 2001; Dhaubhadel et al., 2006). Therefore, we used a  $D_f$  of 1.8 to represent the fluffy BC.

With atmospheric aging, BC can be reconstructed into a condensed structure (Lack et al., 2014; Zhang et al., 2008; Bhandari et al., 2019). In many studies, a larger  $D_f$  was used to represent the compact BC (Liu et al., 2017; Luo et al., 2018, 2019). Previous studies have shown that the  $D_f$  of aged BC can sometimes reach about 2.3– 2.6 (Adachi et al., 2010a, 2007; Chen et al., 2016), and we used a  $D_f$  of 2.6 to represent compact BC. At the same time, the surface of black carbon is covered with other materials, which makes the morphology more complex (China et al., 2013; Pang et al., 2023, 2022; Wang et al., 2021). In the atmosphere, some BC cores are fully coated while others are partially coated. Both the fluffy BC and the compact BC



may be partially coated, while the fully coated BC generally has a compact BC core due to the condensation of the coating materials. Therefore, we assumed the following cases for aged BC aerosols: (1) fluffy BC cores partially coated with other materials; (2) compact BC without coating materials; (3) compact BC partially coated with other materials; (4) compact BC fully coated with other materials. It should be noted that BC with thicker coating materials become more easily compact, so we assumed a  $D_f$  of 2.6 for the fully coated BC.

The DLCA algorithm was developed to characterize the morphology of fresh BC, and the aggregates generated generally have a  $D_f$  of about 1.8 and a  $k_0$  of about 1.3 (Sorensen and Roberts, 1997; Heinson et al., 2010, 2017, 2018). However, the DLCA algorithm is not able to characterize the morphology of the aged BC with a compact structure. Since the fractal structure is well described by equation (2), tunable algorithms based on equation (2) were developed to replace DLCA. Compared to the DLCA algorithm, the tunable algorithms are parametrically adjustable ( $k_0$  and  $D_f$  are fully adjustable) (Filippov et al., 2000; Skorupski et al., 2014; Moran et al., 2019). Although the tunable algorithm does not provide physical explanations of how the morphologies are generated exactly like DLCA, it can represent the DLCA aggregate by setting  $k_0 = 1.3$  and  $D_f = 1.8$ . Moreover, the tunable algorithm can represent more compact BC with larger  $D_f$ . Therefore, we used a tunable algorithm developed by Woźniak (2012) to generate the uncoated BC cores.

After the BC cores were generated, we added the coating materials on the surface of the BC cores. Similar to previous studies (Zhang et al., 2018; Luo et al., 2018; Liu et al., 2017), we adopted a spherical coating structure for coated BC. It is noted that coating materials much more complex than the spherical structure can be observed in the atmosphere. However, as previous studies have shown, the absorption of coated BC is significantly affected by the "lensing effect", which largely depends on the fraction of coated BC cores. Compared with the fraction of BC cores in the coating materials ( $F$ ), the effects of the coating structures on the absorption of BC are relatively small. Therefore, we only consider the  $F$  with a partially coated model with a spherical coating structure, and the complex coating structures were not considered in this work. The effects of coating structures have been studied elsewhere (Luo et al., 2019).

The spherical coating materials were generated similarly to Zhang et al. (2018).  $F$  is calculated using the following equation:

$$F = \frac{V_{BC \text{ inside}}}{V_{BC}} \quad (3)$$

where  $V_{BC \text{ inside}}$  represents the BC volume inside the coating shell, and  $V_{BC}$  represent the volume of the total coated BC.

Once the BC volume fraction ( $f_{BC}$ ) and  $F$  are given, we can determine the radius of the spherical coating. Then the BC cores are moved from side to side until the given  $F$  is reached, similar to Zhang et al. (2018). It should be noted that the motion would lead to an overlap of the BC core and the coating sphere, and their optical properties cannot be calculated by some efficient numerical methods, such as the multiple sphere T-matrix (MSTM) method (Mackowski and Mishchenko, 2011; Mackowski, 2022), which is only applicable for multiple spheres without overlap. For efficient calculations, we moved the overlapping BC sphere cores outside the coating sphere, similar to Zhang et al. (2018), and the movement would not affect the optical properties of coated BC significantly (Liu et al., 2017). The typical BC morphologies are shown in Figure 1.



## 2.2 The calculations of BC absorption

150 Similar to our previous study (Luo et al., 2018), the absorption cross sections ( $C_{\text{abs}}$ ) of individual BC particles were calculated by the MSTM method, which can be used to calculate the optical properties of spheres without overlap. It has been reported that the refractive index of BC varies with wavelength (Chang and Charalampopoulos, 1990), while it is generally known that it does not vary significantly from visible to near-infrared wavelength (Bond and Bergstrom, 2006). Therefore, a fixed BC refractive index of  $1.95 + 0.79i$  was assumed according to the suggestions of Bond and Bergstrom (2006). BrC aerosols are  
 155 organic carbon that absorbs light from UV to visible wavelengths. Therefore, the coating shell was assumed to be organic carbon, and the refractive index was taken to be 1.55 (Bond and Bergstrom, 2006). It should be noted that the imaginary part of BrC should not be 0 in principle. However, in this work we are mainly concerned with studying the absorption of nonabsorbing aerosols mixed with BC, which are erroneously attributed to BrC, so we considered only one nonabsorbing shell. The details are presented in the following sections.

160 Once the BC structures were generated and the refractive index was specified, the MSTM was used to calculate the absorption of each BC particle. The MSTM directly outputs an effective absorption efficiency ( $Q_{\text{abs}}$ ), which is defined as the ratio of the absorption cross section to the mean volume radius, but not to the projected area. Thus, the absorption cross section ( $C_{\text{abs}}$ ) can be calculated as:

$$C_{\text{abs}} = Q_{\text{abs}} \pi r_v^2 \quad (4)$$

where  $r_v$  represents the mean volume radius.

165 In the atmosphere, numerous BC aerosols exist, and the optical properties of BC should average over all the particles. Thus, we calculated the bulk optical properties by assuming different size distributions. Assuming a lognormal distribution for the size distribution of BC cores:

$$n(r_v) = \frac{1}{\sqrt{2\pi} r_v \ln(\sigma_g)} \exp \left[ - \left( \frac{\ln(r_v) - \ln(r_g)}{\sqrt{2\ln(\sigma_g)}} \right)^2 \right] \quad (5)$$

where  $r_g$  and  $\sigma_g$  are the geometric mean radius and geometric standard deviation, respectively.

170 We first calculated the  $C_{\text{abs}}$  of BC with different  $N_s$  (i.e., different  $r_v$ ), and then calculated the absorption coefficient ( $b_{\text{abs}}$ ) using:

$$b_{\text{abs}} = \int_{r_{\text{min}}}^{r_{\text{max}}} C_{\text{cabs}}(r_v) n(r_v) dr_v \quad (6)$$

where  $r_{\text{max}}$  and  $r_{\text{min}}$  are the maximum radius and minimum radius of BC cores, respectively.

The  $r_g$  and  $\sigma_g$  values reported in this work are for the BC cores and can be determined using some single-particle technologies, such as the soot particle photodiameter ( $SP_2$ ) (Moteki et al., 2007; Baumgardner et al., 2004; Schwarz et al., 2006). Many  
 175 measurements were based on single particle technologies, and different  $r_g$  and  $\sigma_g$  were observed in different regions. BC particles in Tokyo were observed with a geometric mean radius ( $r_g$ ) of about  $0.032 \pm 0.003 \mu\text{m}$  and a geometric standard deviation



( $\sigma_g$ ) of  $1.66 \pm 0.12$  (Kondo et al., 2011). However, Shiraiwa et al. (2008) showed that  $r_g$  and  $\sigma_g$  of BC in Fukue, Japan, are  $0.095 - 0.105 \mu\text{m}$  and  $1.45 - 1.55$ , respectively. At BC in Shanghai,  $r_g$  of about  $0.1 \mu\text{m}$  (Gong et al., 2016) was observed. In general, small BC particles are easier to be fully coated, so partially coated BC usually has a large size. Therefore, we mainly calculated the optical properties of relatively large BC particles.  $N_s$  of  $5 - 1000$  were considered, and the corresponding  $r_{\text{max}}$  and  $r_{\text{min}}$  are  $0.0342 \mu\text{m}$  and  $0.2 \mu\text{m}$ . An  $r_g$  of  $0.05-0.1 \mu\text{m}$  was assumed, and the  $\sigma_g$  was assumed to be in the range of  $1.3$  to  $2.0$ .

### 2.3 Inferring the fractions of BrC absorption

We estimated the absorption of BC, which was incorrectly attributed to BrC, after calculating BC absorption. In AERONET, the  $440 \text{ nm}$ ,  $675 \text{ nm}$ , and  $870 \text{ nm}$  wavelengths were most commonly used to estimate BrC absorption, and we mainly considered these three wavelengths. At  $675 \text{ nm}$  and  $870 \text{ nm}$ , all of the absorption was assumed to come entirely from BC. The  $b_{\text{abs}}$  of BC at  $440 \text{ nm}$  can be determined as follows:

$$b_{\text{abs\_BC\_440\_Estimated}} = b_{\text{abs\_BC\_}\lambda} \left( \frac{440}{\lambda} \right)^{-\text{AAE}_{\lambda\_440}} \quad (7)$$

where  $\text{AAE}_{\lambda\_440}$  represents the AAE of BC for the  $\lambda$  and  $440 \text{ nm}$  wavelength pair;  $\lambda$  represents  $870 \text{ nm}$  or  $675 \text{ nm}$  wavelength.

The  $b_{\text{abs}}$  of BrC at  $440 \text{ nm}$  can be estimated using:

$$b_{\text{abs\_BrC\_440\_Estimated}} = b_{\text{abs\_440\_total}} - b_{\text{abs\_BC\_440\_Estimated}} = b_{\text{abs\_BC\_440}} + b_{\text{abs\_BrC\_440}} - b_{\text{abs\_BC\_440\_Estimated}} \quad (8)$$

where  $b_{\text{abs\_BrC\_440\_Estimated}}$  and  $b_{\text{abs\_440\_total}}$  represent the estimated BrC absorption coefficient and total absorption coefficient at  $440 \text{ nm}$ , respectively.  $b_{\text{abs\_BC\_440}}$  and  $b_{\text{abs\_BrC\_440}}$  represent the "true" BC and BrC absorption coefficients at  $440 \text{ nm}$ , respectively. Thus, the absorption coefficient of BC, which is incorrectly attributed to BrC, can be calculated as follows:

$$\Delta_{\text{BrC}} = b_{\text{abs\_BrC\_440}} - b_{\text{abs\_BrC\_440\_Estimated}} = b_{\text{abs\_BC\_440\_Estimated}} - b_{\text{abs\_BC\_440}} \quad (9)$$

Then we calculated the proportions of the wrongly assigned absorption with:

$$\text{ABS}_{\text{BrC}} = \frac{\Delta_{\text{BrC}}}{b_{\text{abs\_BC\_440}}} = \frac{b_{\text{abs\_BC\_440\_Estimated}}}{b_{\text{abs\_BC\_440}}} - 1 \quad (10)$$

The WDA method is similar to the fixed AAE methods, but the AAE are inferred from the Mie theory:

$$\text{WDA} = \text{AAE}_{440\_870\_Mie} - \text{AAE}_{675\_870\_Mie} \quad (11)$$

where the subscript "Mie" stands for the AAE calculated using the Mie theory. Thus, we can estimate the  $\text{AAE}_{440\_870}$  based on the WDA:

$$\text{AAE}_{440\_870} = \text{AAE}_{675\_870} + \text{WDA} \quad (12)$$



## 2.4 Calculating the global BC absorption

- 205 We used a global atmospheric chemical transport model, GEOS-Chem (Bey et al., 2001; Eastham et al., 2018), to simulate the global distribution of BC. For this work, GEOS-Chem version 12.7 was used with a latitude/longitude grid resolution of  $4^\circ \times 5^\circ$ . MERRA-2 (second Modern-Era Retrospective analysis for Research and Applications) assimilated meteorology (Molod et al., 2015; Gelaro et al., 2017) was used for the GEOS-Chem model simulations. The model was built with 47 vertical layers. We ran a standard GEOS-Chem simulation with various aerosols such as dust, organic aerosols, BC, sulphate, sea salt, etc.
- 210 The Community Emissions Data System (CEDS) (Hoesly et al., 2018) inventory provided the global anthropogenic emissions. The Global Fire Emissions Database (GFED4) inventory (Randerson et al., 2018) was used to provide emissions from biomass burning. Biogenic emissions were obtained from the Model of Emissions of Gases and Aerosols from Nature Version 2.1 (MEGAN 2.1) (Guenther et al., 2012). We used GEOS-Chem to simulate black carbon concentrations for all of 2016, and then took the time average.
- 215 The BC AAOD in each GEOS-Chem grid, was calculated using:

$$AAOD_{BC} = MAC_{BC} \times C_{BC\_column} \quad (13)$$

where  $C_{BC\_column}$  is the column mass concentrations of BC;  $MAC_{BC}$  represents the mass cross sections of coated BC that can be calculated with:

$$MAC_{BC} = \frac{b_{abs}}{m_{BC}} \quad (14)$$

- 220 where  $m_{BC}$  represent the mass distributions of BC cores which can be calculated with:

$$m_{BC} = \int_{r_{min}}^{r_{max}} \frac{\rho_{BC} 4\pi r_v^3 n(r_v)}{3} dr_v \quad (15)$$

where  $\rho_{BC}$  represents the mass density of BC. Bond and Bergstrom (2006) to use a  $\rho_{BC}$  of  $1.8 \text{ g m}^{-3}$ . However, most modeling studies underestimated the proposed MAC of  $7.5 \pm 1.2 \text{ m}^2 \text{ g}^{-1}$  at 550 nm using the proposed  $\rho_{BC}$ . Similar to Luo et al. (2018), we used a  $\rho_{BC}$  of 1.5 to fit the MAC measurements.

- 225 The global AAOD of BC, incorrectly attributed to BrC, can be determined using the following method:

$$AAOD_{BrC} = AAOD_{BC} \times ABS_{BrC} \quad (16)$$

## 3 Results

### 3.1 The effects of BC microphysical properties on the fixed AAE method

- Figure 2 shows the effects of the shell diameter to core diameter ratio ( $D_p/D_c$ ). The error bars in the figures represent the upper and lower limits when  $r_g$  is varied in the range of  $0.05 - 0.1 \mu\text{m}$  and  $\sigma_g$  is varied in the range of  $1.5 - 1.8$ . Since BC
- 230





aerosols are freshly emitted, they are fluffy and not mixed with the coating materials. Our cases where  $D_f = 1.8$  and  $F = 0$  may reflect freshly emitted BC. Previous studies have shown that the AAE of freshly emitted BC is not significantly different from 1 (Liu et al., 2018; Luo et al., 2020). Thus, using  $AAE_{440,675} = 1$  and  $AAE_{440,870} = 1$  can provide reasonable estimates of BrC absorption for freshly emitted BC, and the estimated  $ABS_{BrC}$  is not significantly different from 0 when  $D_p/D_c < 2.71$ .  
235 Based on different AAE wavelength pairs, the estimated  $ABS_{BrC}$  shows some differences.  $ABS_{BrC}$  ranges from about -3% to 4.5% when  $D_p/D_c < 2.71$  and the  $AAE_{440,675} = 1$  method is used, while this range becomes about -12% – 3% when the  $AAE_{440,870} = 1$  method is used. As the number of coating materials increases, BC AAE may gradually deviate from 1 due to the shielding effect of heavy coatings. Thus, a broader  $ABS_{BrC}$  of about -6% – 18% is obtained when  $D_p/D_c$  increases to 4.64 using the  $AAE_{440,675} = 1$  method. A relatively wider range was also observed when  $AAE_{440,870} = 1$ .  $ABS_{BrC}$  ranges from  
240 about -33% to -6% when  $D_p/D_c = 4.64$  using the  $AAE_{440,870} = 1$  method.

The increases in  $F$  may represent a process of atmospheric aging. The bare, fluffy BC aggregates are gradually coated by atmospheric aging, so BC with a larger  $F$  represents more aged particles. Since the BC aggregates are partially coated,  $ABS_{BrC}$  also gradually deviates from 0 for the fluffy BC aggregates. With a  $D_f$  of 1.8 and an  $F$  of 0.1 when  $D_p/D_c < 2.71$   $ABS_{BrC}$  varies in the range of about -6% - 18% and -12% – 9% when  $AAE_{440,675} = 1$  and  $AAE_{440,870} = 1$  is used, and the ranges  
245 become -18% - 3% and -21% – 4.5% when  $D_p/D_c$  is 4.64. The  $ABS_{BrC}$  range becomes larger as  $F$  increases to 0.3. A  $ABS_{BrC}$  range of -22% – 24% can be observed when BC with a  $D_f$  of 1.8 and a  $F$  of 0.1. When BC aggregates are partially coated, the  $ABS_{BrC}$  are generally smaller than 0 when  $D_p/D_c$  is larger (e.g.,  $D_p/D_c = 4.64$ ) when BC has a fluffy structure, whereas  $ABS_{BrC}$  can be larger than 0 when  $D_p/D_c$  is small. This phenomenon may be due to the different effects of coating ratios on AAE. When  $D_p/D_c$  is small, the AAE is generally small and can sometimes be less than 1 (Zhang et al., 2020b; Luo et al.,  
250 2023; Liu et al., 2018), so a  $ABS_{BrC}$  greater than 0 can be observed. On the other hand, the AAE increases with  $D_p/D_c$  when BC has a fluffy structure. Thus, the AAE can be greater than 1 when the fluffy BC is partially coated with a thick coating (Zhang et al., 2020b; Luo et al., 2023), resulting in  $ABS_{BrC}$  of less than 0. For a  $D_f$  of 1.8, Luo et al. (2023) have shown that the  $AAE_{440,870}$  of partially coated BC generally first decreases with increasing  $D_p/D_c$  and then increases when  $D_p/D_c$  is greater than a certain value. An opposite phenomenon is observed for  $ABS_{BrC}$ . For BC with  $D_f$  of 1.8,  $ABS_{BrC}$  of partially  
255 coated BC generally increases first with increasing  $D_p/D_c$  and then decreases when  $D_p/D_c$  is greater than 2.71.

The more compact structure can also represent another process of atmospheric aging. Even with  $F = 0$ , a  $D_f$  of 2.6 represents the highly aged BC. By comparing BC with fluffy and compact structures, we can see more deeply from the effects of atmospheric aging on the estimations of BrC absorption. As the BC cores are reconstructed to a compact structure, the  $AAE = 1$  method provides inaccurate estimations even when  $F = 0$  and  $D_p/D_c < 2.71$ . With a  $D_f$  of 2.6 and an  $F$  of 0,  $ABS_{BrC}$  varies in  
260 the range of approximately -14% – 14% and -1% – 21% when using  $AAE_{440,675} = 1$  and  $AAE_{440,870} = 1$ , respectively. Besides, larger  $ABS_{BrC}$  deviations from 0 can be observed as  $F$  increases, and  $ABS_{BrC}$  sometimes can increase to approximately 21% when using  $AAE_{440,675} = 1$  and to approximately 37.5% when using  $AAE_{440,870} = 1$ . Different for the cases where  $D_f = 1.8$ , BC with a  $D_f$  of 2.6 in most cases exhibits an  $ABS_{BrC}$  of larger than 0. The reason is that compact BC generally exhibits a small AAE, which is generally less than 1 (Luo et al., 2023; Liu et al., 2018). Besides, in most cases,  $ABS_{BrC}$  of BC with  
265 compact structure increases with increasing  $D_p/D_c$ . This can be explained by the findings in the previous study. Luo et al.



(2023) have shown that the  $AAE_{440,870}$  generally decreases with increasing  $D_p/D_c$ , which leads to an  $ABS_{BrC}$  increase with  $D_p/D_c$ .

Because the  $AAE = 1$  method would provide inaccurate estimates, many researchers have attempted to use a different AAE value to estimate BrC. For example, Zhang et al. (2019) used a BC AAE of 0.7 in the Pearl River Delta region, China; Rathod and Sahu (2022) suggested using a BC AAE of  $1.1 \pm 0.05$ . However, BC's AAE is subject to large uncertainties in regions, age status, burn sources, etc. Moreover, the BC AAE itself is not unchanged in the same region, and it would vary with time as the microphysical properties of BC vary with atmospheric aging. Figures 3 - 4 show how the microscopic properties of soot affect the application of the different values. As can be seen from 3 - 4, there is no fixed AAE value that is applicable to all cases. In general,  $ABS_{BrC}$  are larger for larger AAE values. Kirchstetter et al. (2004) has measured an AAE range of 0.6 to 1.3. However, using values in this range, the AAE method would not always provide accurate estimates for BrC in the absence of additional information. Sometimes the estimated  $ABS_{BrC}$  can be approximately -40% – 60% if we choose a fixed AAE in the range of 0.6 – 1.3. The  $ABS_{BrC}$  based on different AAE values is significantly affected by BC size distributions, morphology, and mixing states. For freshly emitted BC ( $D_f = 1.8$ ,  $F = 0$  in our cases), AAEs of 0.9 – 1.1 would provide reasonable estimates, and the  $ABS_{BrC}$  at this time is in the range of about -6% – 6%. However, using AAEs of 0.6 and 1.3 can result in  $ABS_{BrC}$  of about -15% and 18%, respectively.

However, for a more compact structure and a larger  $F$ , the applicability of the different AAE values depends more on the particle size. For  $D_f = 1.8$  and  $F = 0.2$ , using  $AAE = 1$  would lead to a  $ABS_{BrC}$  range of about 7% – 0 when  $f_{BC} = 20\%$ , while this range increases to about 6% – 17% when  $f_{BC} = 5\%$ . As the BC cores become more compact ( $D_f = 2.6$ ), the fixed AAE methods may be more affected by the particle size. It is difficult to find a fixed AAE for estimating BrC in the absence of additional size information, and using a fixed AAE could result in a  $f_{BC}$  range of about -26% to 24% in our selected cases. In Figure 3 – 4, we see that  $ABS_{BrC}$  increases with  $r_g$  when AAE is fixed. This is caused by a decrease in AAE with increasing  $r_g$  for fluffy BC (Luo et al., 2023; Zhang et al., 2020b). However, in some cases, the AAE of compact BC may increase with increasing  $r_g$  (Luo et al., 2023), which could lead to a decrease in  $ABS_{BrC}$  with increasing  $r_g$ . In summary,  $ABS_{BrC}$  is significantly affected by BC size when a fixed BC AAE is used, and it appears that no fixed BC AAE values are applicable to all cases. Similar conclusions can be also drawn for the fully coated BC (see Figure S1).

### 3.2 The effects of BC microphysical properties on the WDA method

Because fixed AAE methods cannot always provide accurate estimates of BrC, Wang et al. (2016) proposed a WDA method to derive BrC absorption. They first calculated the WDA based on Mie theory and then derived the AAE based on the WDA. They claimed that this method can reduce the effects of BC size and coatings. However, in the study of Wang et al. (2016), the morphology of BC was assumed to be spherical, and the WDA calculated based on Mie theory does not always provide accurate estimates. In fact, the spherical model only represents the highly aged BC. Unaged BC in the atmosphere often exhibits fluffy morphologies, and the coating materials make the morphology coated BC more complex (Pang et al., 2022, 2023; Wang et al., 2017). The WDA calculated based on the Mie theory may not represent BC with complex morphologies. In addition, as BC morphologies can change with atmospheric aging, the spherical assumption provides an uncomplete understanding on the



300 effects of atmospheric aging on WDA. Figure 5 shows the estimated  $ABS_{BC}$  based on the WDA method. Since some BC cores with partial coating are fully coated while others are not, we calculated  $ABS_{BC}$  based on the WDA using the bare sphere model or the core-shell model. As pointed out by Luo et al. (2021c), the WDA method does not necessarily provide better estimates than the  $AAE = 1$  method. The estimated  $ABS_{BC}$  based on the WDA using the bare sphere model and the core-shell model has comparable values and can vary in a range from about -40% to 36%.

305 The estimated  $ABS_{BC}$  based on the WDA method is generally in the range of -16% – 5% for freshly emitted BC ( $F = 0$ ,  $D_f = 1.8$ ,  $D_p/D_c < 2.71$ ), and the range becomes approximately -18% – 14% when  $D_p/D_c$  increases to 4.64. For freshly emitted BC, the estimated value for  $ABS_{BC}$  is generally less than 0. This is because the WDA estimated with Mie generally underestimates the WDA of BC with fluffy morphology. As shown in Figure 6, the WDA estimated with the morphologically realistic BC model is not significantly different from 0 for freshly emitted BC, whereas the WDA estimated with the spherical  
310 model varies in a range from about -0.25 to 0.25. The estimated  $ABS_{BC}$  is in the range of about -20% – 16% and -14% – 36% when  $F$  increases to 0.1 and 0.3, respectively, and the difference in WDA between the partially coated BC and the spherical BC can sometimes reach about 0.4. The estimate  $ABS_{BC}$  based on the WDA method is generally smaller than 0 when the BC cores become compact ( $D_f = 2.6$ ). This is due to the fact that the WDA of the partially coated BC is underestimated by Mie theory (see Fig. 6). For a  $D_f$  of 2.6, the estimated  $ABS_{BC}$  is generally in the range of -40% – 0, and the difference between the  
315 WDA of the partially coated BC and the spherical BC can sometimes even reach about 0.75. Thus, since Mie theory does not always provide accurate estimates for the WDA of BC with realistic morphologies, we should carefully consider the effects of the morphologies of partially coated BC when applying the WDA method.

Figure 7 shows the effects of size distributions on estimated  $ABS_{BC}$  based on the WDA method.  $ABS_{BC}$  shows different trends with  $r_g$  for BC with different  $F$  and  $f_{BC}$ . When  $D_f = 1.8$  and  $F = 0$ ,  $ABS_{BC}$  estimated by the WDA method generally  
320 decreases with increasing  $r_g$  when  $f_{BC} > 5\%$ , while the opposite phenomenon is observed when  $f_{BC} = 1\%$ . As  $F$  increases, different  $ABS_{BC}$  trends can be observed. Unlike the case of  $F = 0$ ,  $ABS_{BC}$  increases with increasing  $r_g$  when  $f_{BC} = 5\%$ . As BC cores are reconstructed into a compact structure ( $D_f = 2.6$ ),  $ABS_{BC}$  decreases with increasing  $r_g$  when BC is heavily coated ( $f_{BC} = 1\%$  and  $5\%$ ), while an opposite trend is observed for thinly coated BC. Moreover,  $ABS_{BC}$  decreases with increasing  $r_g$  even for thinly coated BC when  $F$  is larger, and similar findings are found for fully coated BC ( $F = 1.0$ ), as shown  
325 in Figure S2.

To explain why the above phenomenon occurs, we also calculated the WDA of BC with different  $r_g$ . Figure 8 shows the variations of  $ABS_{BC}$  with  $r_g$  for partially coated BC with different mixed states. The WDA of partially coated BC is very different from that of spherical BC. In general, the WDA estimated with the bare sphere model and the core-shell sphere model show similar trends with  $r_g$ . When  $F = 0$  and  $f_{BC} = 5\%$ , the WDA calculated with the morphologically realistic model  
330 is comparable to that calculated with the spherical model when  $r_g$  is small, leading to  $ABS_{BC}$  of about 0 (see Figure 7). However, as  $r_g$  increases, the WDA of the morphologically realistic BC increases, whereas the WDA calculated with the spherical models does not change significantly. Thus, the WDA difference between morphologically realistic BC and spherical models (the WDA of the morphologically realistic BC minus the WDA of the spherical BC) increases with increasing  $r_g$ , so that  $ABS_{BC}$  decreases with  $r_g$ . However, the WDA of both the morphologically realistic BC and spherical models does not



335 vary significantly with  $r_g$  when  $F = 0$  and  $f_{BC} = 20\%$ , so the  $ABS_{BrC}$  does not vary significantly when  $r_g$  is changed. When  $F$  increases, a different phenomenon is observed. When  $F = 0.2$  and  $f_{BC} = 5\%$ , the WDA of morphologically realistic BC is much smaller than that of spherical models, which can lead to a positive  $ABS_{BrC}$ . Moreover, the WDA of morphologically realistic BC decreases significantly with the increase of  $r_g$ , while the WDA calculated with the spherical models are not significantly varied, so the WDA difference decreases with increasing  $r_g$ . Therefore, the  $ABS_{BrC}$  may increase with  $r_g$ .

340 As BC cores become compact ( $D_f = 2.6$ ), different WDA trends are observed with  $r_g$ . The WDA of partially coated BC is generally larger than that of spherical BC in the cases where  $F = 0$  and  $f_{BC} = 5\%$ , so that a negative  $ABS_{BrC}$  is observed. Moreover, during this time, the WDA of the morphologically realistic BC increases with increasing  $r_g$ , whereas the WDA calculated with the spherical models do not vary significantly with  $r_g$ . As a result, the WDA difference between the partially coated BC model and the spherical models increases with increasing  $r_g$ , leading to decreasing  $ABS_{BrC}$  with  $r_g$ . Similar results  
345 are found when  $F = 0.2$  and  $f_{BC} = 5\%$ , and when  $F = 0.2$  and  $f_{BC} = 20\%$ . A different phenomenon was observed when  $F = 0$  and  $f_{BC} = 20\%$ . At this time, the WDA of partially coated BC is generally larger than that of spherical BC, so  $ABS_{BrC}$  has negative values. However, the WDA of partially coated BC decreases with increasing  $r_g$ , and the WDA difference between the morphologically realistic model and the spherical models becomes smaller. Therefore,  $ABS_{BrC}$  increases with  $r_g$  and tends to 0 when  $r_g$  is large (e.g.,  $r_g = 0.1$ ). Similar results can be found for other  $\sigma_g$ , with minor differences observed (see Figure S3 –  
350 S4.).

### 3.3 Effects of microphysical properties of BC on the global estimation of BrC

Recent studies have generated increasing interest in estimating the global distribution of BrC (Zeng et al., 2020a; Wang et al., 2016). In this work, the effects of BC microphysical properties on the global estimation of BrC were also investigated. Figure 9 shows the mean global optical absorption aerosol density (AAOD) of BC calculated with different configurations. In the  
355 figures, case A represents aerosols where  $D_f = 1.8$ ,  $F = 0.0$ , and  $f_{BC} = 5\%$ ; case B represents aerosols where  $D_f = 1.8$ ,  $F = 0.0$  and  $f_{BC} = 20\%$ ; Case C represents aerosols where  $D_f = 1.8$ ,  $F = 0.2$  and  $f_{BC} = 5\%$ ; Case D represents aerosols where  $D_f = 1.8$ ,  $F = 0.2$  and  $f_{BC} = 20\%$ ; Case E represents aerosols where  $D_f = 2.6$ ,  $F = 0.0$  and  $f_{BC} = 5\%$ ; Case F represents aerosols where  $D_f = 2.6$ ,  $F = 0.0$ , and  $f_{BC} = 20\%$ ; Case G represents aerosols with  $D_f = 2.6$ ,  $F = 0.2$ , and  $f_{BC} = 5\%$ ; Case H represents aerosols with  $D_f = 2.6$ ,  $F = 0.2$ , and  $f_{BC} = 20\%$ .

360 By comparing the five models, Sand et al. (2021) showed that the BC AAOD is generally in the range of 0.0007 – 0.007. Kinne (2019) reported a larger AAOD of 0.0072. However, their studies assume spherical core-shell structures, and a fully coated model could overestimate the overall AAOD. Based on GEOS-Chem simulations, Kelesidis et al. (2022) have shown that the global mean BC AAOD is about  $0.0017 \pm 0.007$  and  $0.003 \pm 0.0016$  when the spherical bare model and the coated model are used, respectively. However, the values increase to about  $0.0021 \pm 0.0008$  and  $0.0036 \pm 0.0014$  when bare  
365 and coated BC agglomerate models are used, respectively. However, in their study, the coated BC aerosols were assumed to be fully coated. Our case studies show a global mean AAOD value of about 0.0017 to 0.0026, which strongly depends on the microphysical properties of BC. Our simulated AAOD is generally in the range of Kelesidis et al. (2022). Our simulated AAOD is generally larger than simulations with agglomerate-only models, but smaller than simulations with fully coated agglomerate



models. This is easy to understand since we have mainly considered partially coated BC in this work. It is worth noting that  
370 the AAOD in this work is 440 nm, while previous studies generally set it at 550 nm. However, the AAOD does not differ  
significantly between these two wavelengths. Moreover, the aim of this work is to evaluate the impact of BC microphysical  
properties on the estimation of BrC, but not to make a detailed comparison with the previous studies, so our analysis is valid.

In general, the AAODs of BC with fluffy structure are higher than those of BC with compact structure, which is consistent  
with the results of previous studies (Liu and Mishchenko, 2005; Luo et al., 2022; Kahnert and Devasthale, 2011). This is due  
375 to the blocking effects of a more compact structure and results in lower absorption (Kahnert and Devasthale, 2011). When  $F$   
 $= 0$ , the effects of the coating ratio ( $f_{BC}$ ) are not significant. The global BC AAOD is about  $1.9 - 2.1 \times 10^{-3}$  and  $1.7 - 1.9$   
 $\times 10^{-3}$  for BC with  $D_f = 1.8$  and  $D_f = 2.6$ , respectively. As expected, the global mean BC AAOD increases as  $F$  increases.  
As  $F$  increases to 0.2 when  $D_f = 1.8$ , the global BC AAOD increases to about  $2.1 - 2.5 \times 10^{-3}$  and  $2.45 - 2.6 \times 10^{-3}$  when  
 $f_{BC}$  is 5% and 20%, respectively. The results show that the AAOD does not necessarily increase when the coating ratio for the  
380 partially coated BC increases due to the shielding effect of the coating materials.

Direct Radiative Forcing (DRF) is an important parameter for assessing climate change. Bond et al. (2013) estimated a global  
mean DRF of about  $+0.17 - +1.48 \text{ W/m}^2$ , while more recent studies have shown much lower DRF values. For example, Kinne  
(2019) showed a global mean BC DRF of about  $+0.55 \text{ W/m}^2$ ; Matsui et al. (2018) reported a DRF of  $+0.18 - +0.42 \text{ W/m}^2$ ;  
Chen et al. (2022) estimated a mean BC DRF of  $+0.33 \text{ W/m}^2$  [ $+0.17, +0.54$ ]. However, Heald et al. (2014) reported a much  
385 smaller global mean BC DRF of  $+0.078 \text{ W/m}^2$ ; Tuccella et al. (2020) showed a global mean BC DRF of about  $+0.13$  and  
 $+0.25 \text{ W/m}^2$  using a bare-sphere model and a core-shell-sphere model, respectively. Considering an average absorption force  
of  $170 \pm 43 \text{ W/m}^2 / \text{AAOD}$ , based on the study of Bond et al. (2013), the estimated DRF in this work may be about  $+0.216 -$   
 $+0.612 \text{ W/m}^2$ , which is generally in the range of values reported by previous studies.

Figure 10 shows the estimated global BC AAOD, which is incorrectly attributed to BrC (BrC AAOD error in the figures)  
390 when AAE-based methods are used. Using the  $\text{AAE}_{440,675} = 1$  method, the errors for cases A, B, E, G, H are relatively small,  
and the AAOD errors are in the range of  $-0.1 - 0.1 \times 10^{-3}$ . However, in cases C, D, and F, the errors are larger, and the  
estimated AAOD errors can sometimes reach  $0.45 \times 10^{-3}$ , which corresponds to about 20% of the global BC AAOD. Larger  
errors are observed in our cases when the method  $\text{AAE}_{440,870} = 1$  is used. The estimates have relatively small errors in cases  
B, F, and H, and the AAE errors are in the range of about  $-0.1 - 0.2 \times 10^{-3}$ . However, in cases A, C, D, E, and G, a larger range  
395 of about  $-0.4$  to  $0.5 \times 10^{-3}$  was observed. Assuming an average absorption forcing efficiency of  $170 \pm 43 \text{ W/m}^2 / \text{AAOD}$   
(Bond et al., 2013), we can estimate a DRF range of  $-0.068 \pm 0.0172$  to  $+0.085 \pm 0.0215 \text{ W/m}^2$  at BC, which is incorrectly  
attributed to BrC. The WDA does not necessarily improve the estimates when BC has a complex morphology. The spherical  
model-based WDA methods may result in a global mean AAOD error range of approximately  $-0.9 - 0.05 \times 10^{-3}$  in our cases,  
which could result in a global mean DRF error range of  $-0.153 \pm 0.0387$  to  $+0.0085 \pm 0.0022 \text{ W/m}^2$ .

400 Figures 11 – 12 and Figures S5 – S6 show the BrC AAOD errors using different AAE-based methods. Our results show  
that the applicability of the different AAE-based methods is significantly limited by the microphysical properties of BC. When  
all BC are freshly emitted (e.g.,  $F = 0$ ,  $f_{BC}$ ), the  $\text{AAE}_{440,870} = 1$  method would provide reasonable estimates for BrC AAOD  
estimates, and the AAOD errors are within  $0.1 \times 10^{-3}$ . However, the errors are more substantial the older the BC (with



more coating, larger  $F$ , or more compact structure). Sometimes  $AAE_{440,870} = 1$  leads to an AAOD error of over  $3 \times 10^{-3}$  in East Asia, resulting in a BC DRF of about  $+0.51 \pm 0.129 \text{ W/m}^2$ , which is incorrectly attributed to BrC. Similar results are also observed using the  $AAE_{440,870} = 1$  method (see Figure S.). In contrast, the  $AAE_{440,870} = 1$  method would actually lead to relatively larger errors for freshly emitted BC. The WDA method may lead to larger errors if all BC have complex morphology. In some cases, the Mie theory-based WDA methods would lead to an estimated AAOD error of more than  $6 \times 10^{-3}$  in some regions (e.g., East Asia). Multiplying this AAOD error by an average absorption forcing of  $170 \pm 43 \text{ W/m}^2$  /AAOD results in an estimated misassigned DRF of  $+1.0 \pm 0.258 \text{ W/m}^2$ . Therefore, we should carefully consider the effects of BC microphysical properties when using AAE-based methods.

#### 4 Conclusions and summary

Previous studies have recognized that AAE-based methods would provide inaccurate estimates for BrC contributions, whereas the effects of BC microphysical properties are not systematically quantified. Moreover, the conditions under which the AAE-based method is applicable are still unclear. In this work, BC models partially coated with non-absorbing materials were used to calculate the total absorption, and then the BrC absorption was estimated using different AAE-based methods. Since the total absorption is entirely from BC, the estimated BrC absorption should be 0 if the retrieval methods are accurate. Thus, the estimated BrC absorption should be the absorption from BC that is incorrectly attributed to BrC. We have numerically quantified the effects of BC microphysical properties on the percentage of BC absorption falsely attributed to BrC. In this work, different configurations for BC morphologies, coating ratios, coating percentages, and size distributions were used to reflect different aging conditions.

Freshly emitted BC is generally fluffy and not well mixed with other materials. Our configurations of  $D_f = 1.8$  and  $F = 0$  can reflect freshly emitted BC. In this case, a BC AAE of 1 can generally provide useful estimates. However, we recommend using  $AAE_{440,675} = 1$  instead of  $AAE_{440,870} = 1$ . When  $D_f = 1.8$  and  $F = 0$ ,  $ABS_{BrC}$  is generally in the  $-3\% - 4.5\%$  range when  $D_p/D_c < 2.71$ , while an  $ABS_{BrC}$  range of approximately  $-12\% - 3\%$  has been observed. However, the more compact BC becomes and the more coating materials are mixed with BC, the larger the AAOD errors derived from AAE-based methods. Sometimes the fraction BB deviates significantly from 0 when BC becomes compact and is coated. At an  $F$  of 0.3, a negative  $ABS_{BrC}$  of about  $-33\%$  can be observed for BC with an  $D_f$  of 1.8 using  $AAE_{440,870} = 1$ , and  $ABS_{BrC}$  can reach about  $35\%$  when  $D_f = 2.6$ .

Based on GEOS-Chem simulations, we also investigated the effects of BC microphysical properties on estimates of global BrC AAOD. Our case studies show that BC AAOD is significantly affected by BC microphysical properties. Using the morphologically realistic models, a global mean BC AAOD of about 0.0017 to 0.0026 was estimated, corresponding approximately to a DRF of  $+0.216 - +0.612 \text{ W/m}^2$ . When BC aerosols have complex morphology, the AAE = 1 method can sometimes result in a misassigned global mean AAOD of about  $-0.4 - 0.5 \times 10^{-3}$ , leading to a global mean DRF range of about  $-0.068 \pm 0.0172$  to  $+0.085 \pm 0.0215 \text{ W/m}^2$  from BC, which is incorrectly assigned to BrC. The WDA method does not necessarily improve the estimates. In our cases, the WDA methods based on the spherical models may lead to a range of about  $-0.9 - 0.05$



440  $\times 10^{-3}$  of misassigned AAOD, which could lead to a global mean DRF error range of  $-0.153 \pm 0.0387$  to  $+0.0085 \pm 0.0022$   $W/m^2$ . At the regional scale,  $AAE = 1$  in East Asia sometimes leads to a distributed AAOD of over  $3 \times 10^{-3}$ , resulting in a BC DRF of about  $+0.51 \pm 0.129$   $W/m^2$ , which is incorrectly attributed to BrC. Mie theory-based WDA methods would result in an estimated AAOD error of more than  $6 \times 10^{-3}$  in some regions (e.g., East Asia), leading to an estimated misattributed DRF of  $+1.0 \pm 0.258$   $W/m^2$ . In summary, when using AAE-based methods, we should carefully consider the effects of BC microphysical properties.

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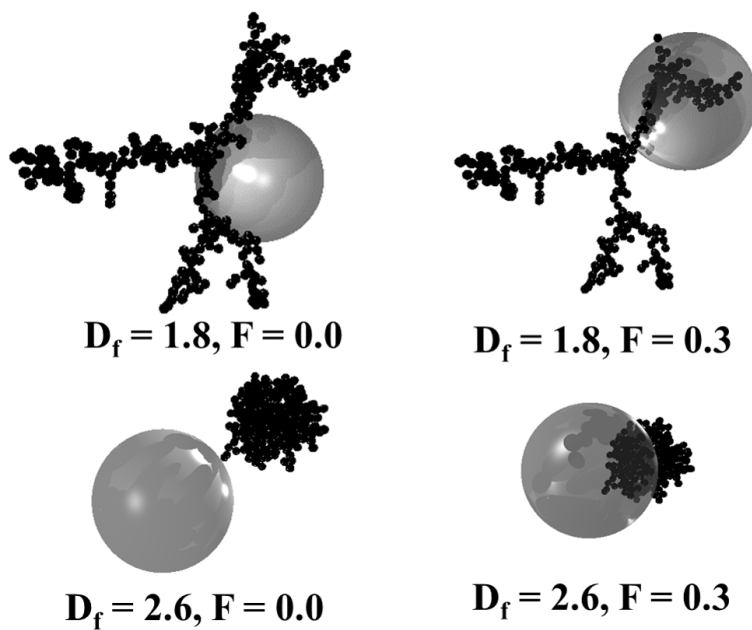
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**Figure 1.** Typical BC morphologies assumed in this work, similar to Luo et al. (2023).



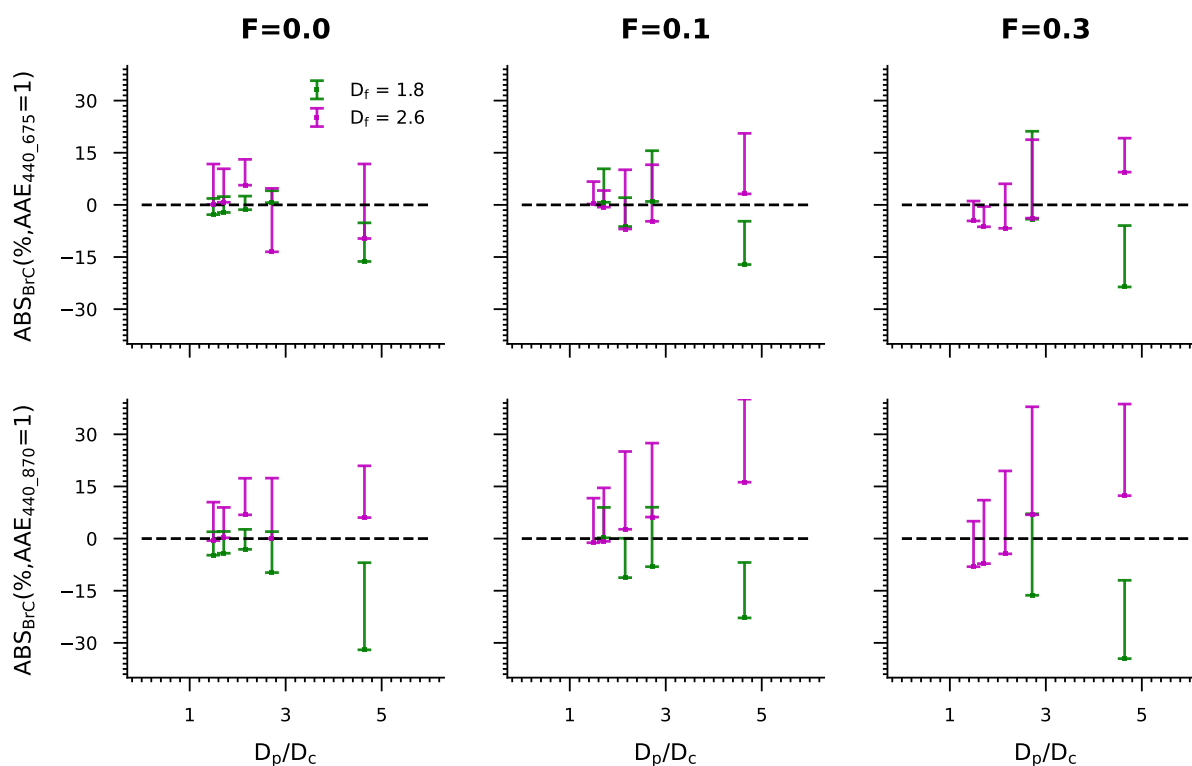
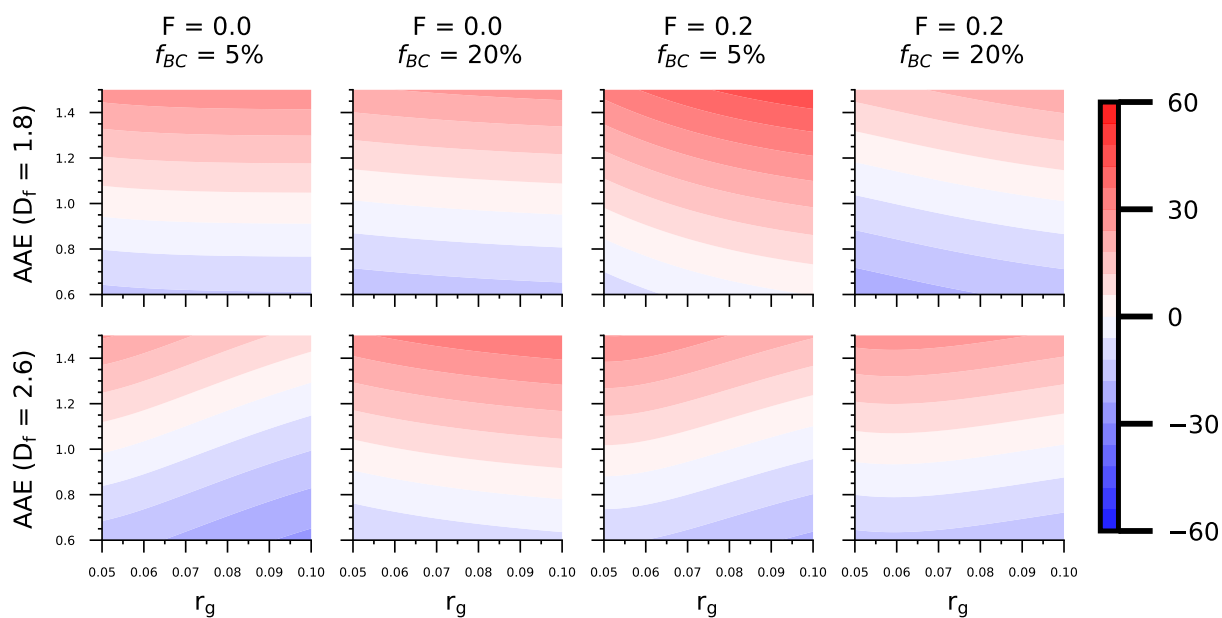
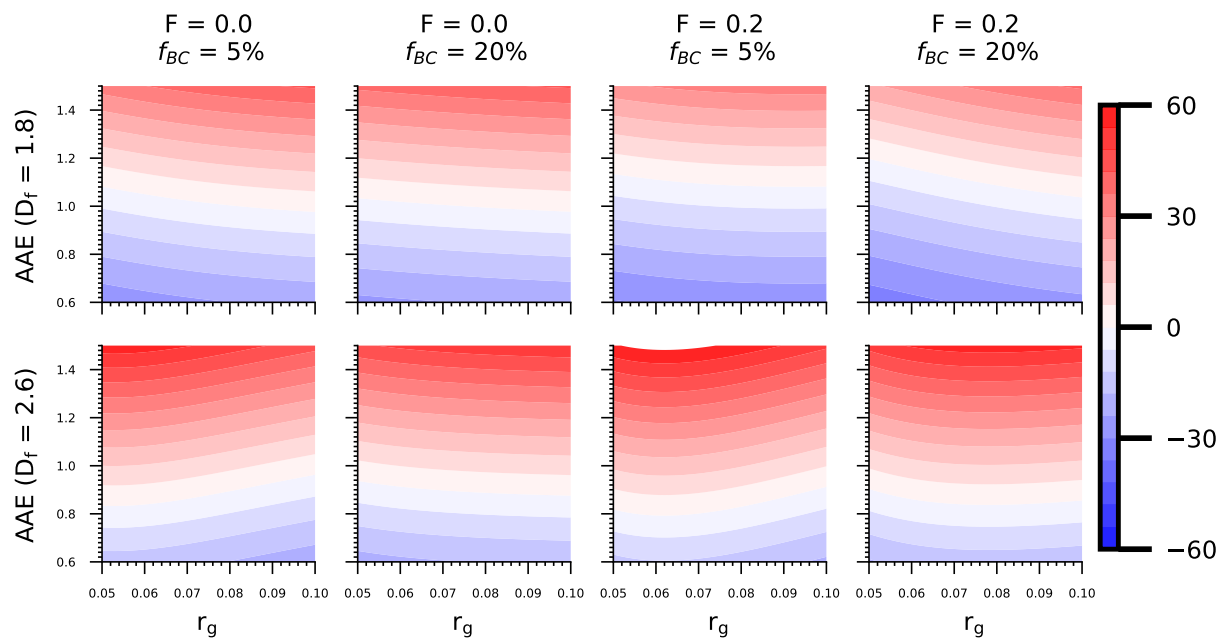


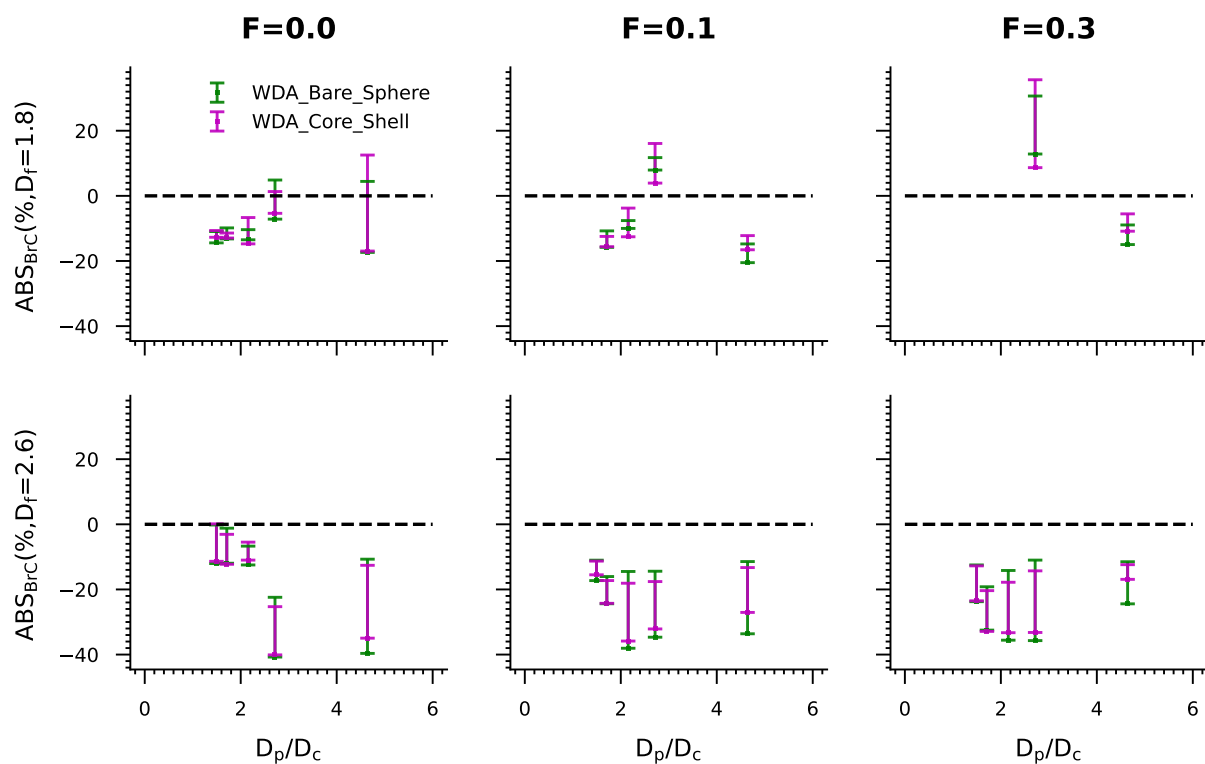
Figure 2.  $ABS_{BrC}$  at different  $D_p/D_c$  estimated using the  $AAE = 1$  methods.



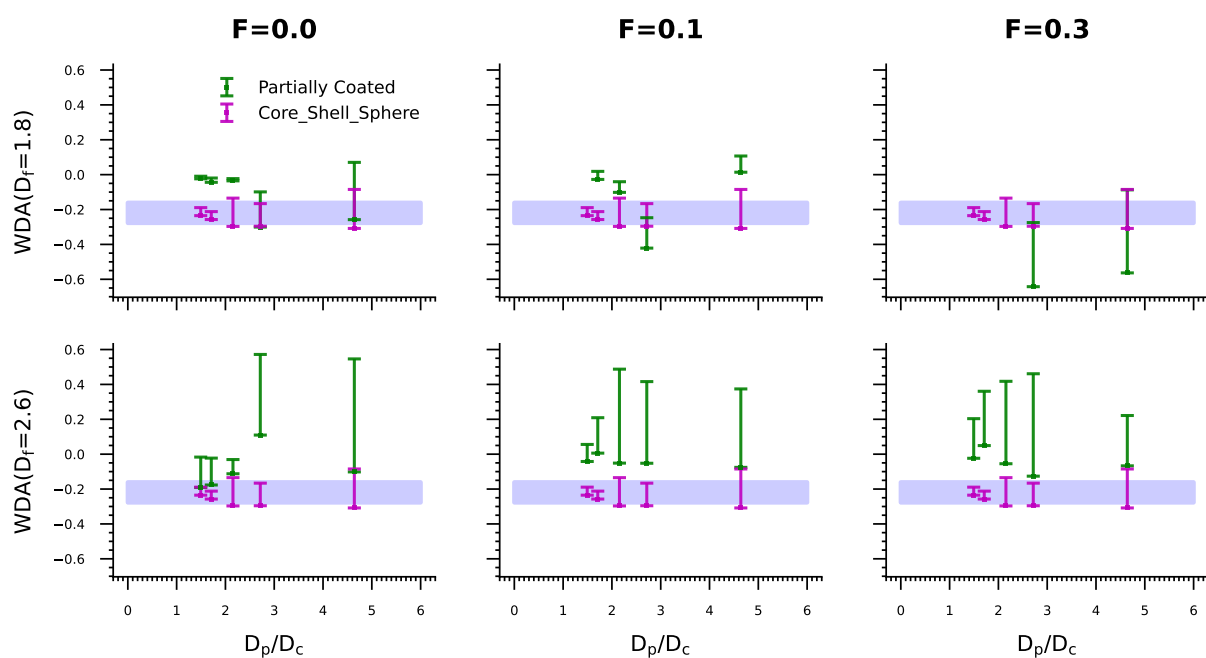
**Figure 3.** The variations of  $ABS_{BRC}$  estimated based on the fixed AAE with the function of AAE and  $r_g$ , where the wavelength pair is 440 nm - 675 nm.



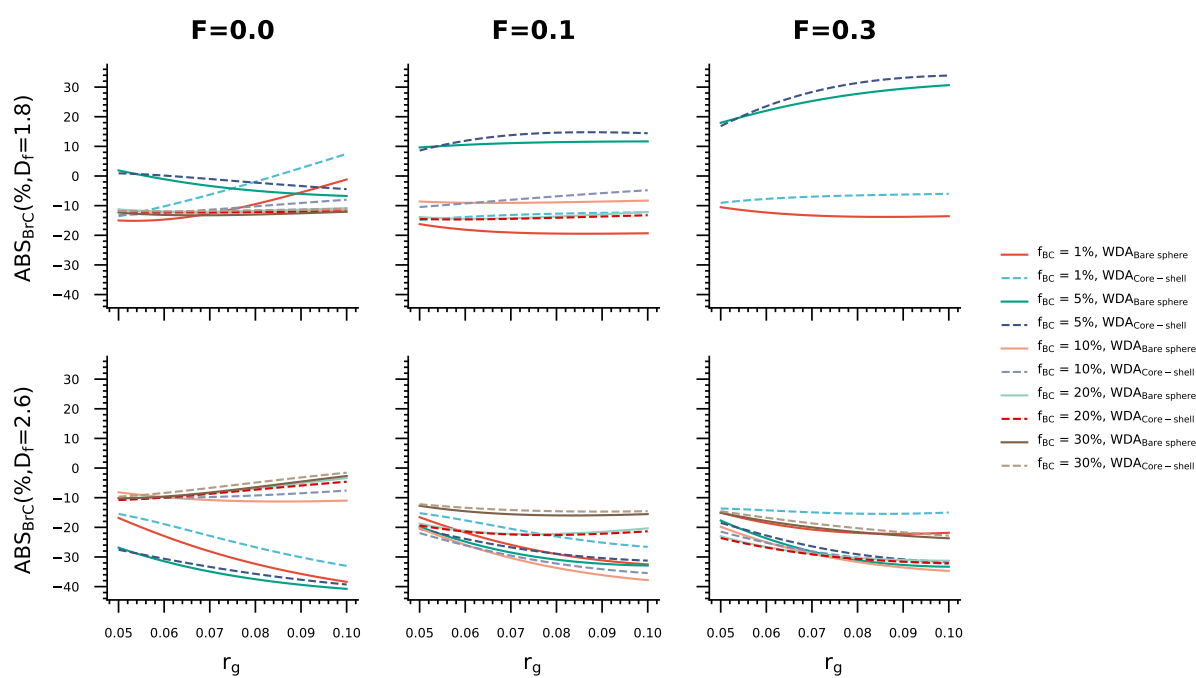
**Figure 4.** Similar to Figure 3, but the wavelength pair is 440 – 870 nm.



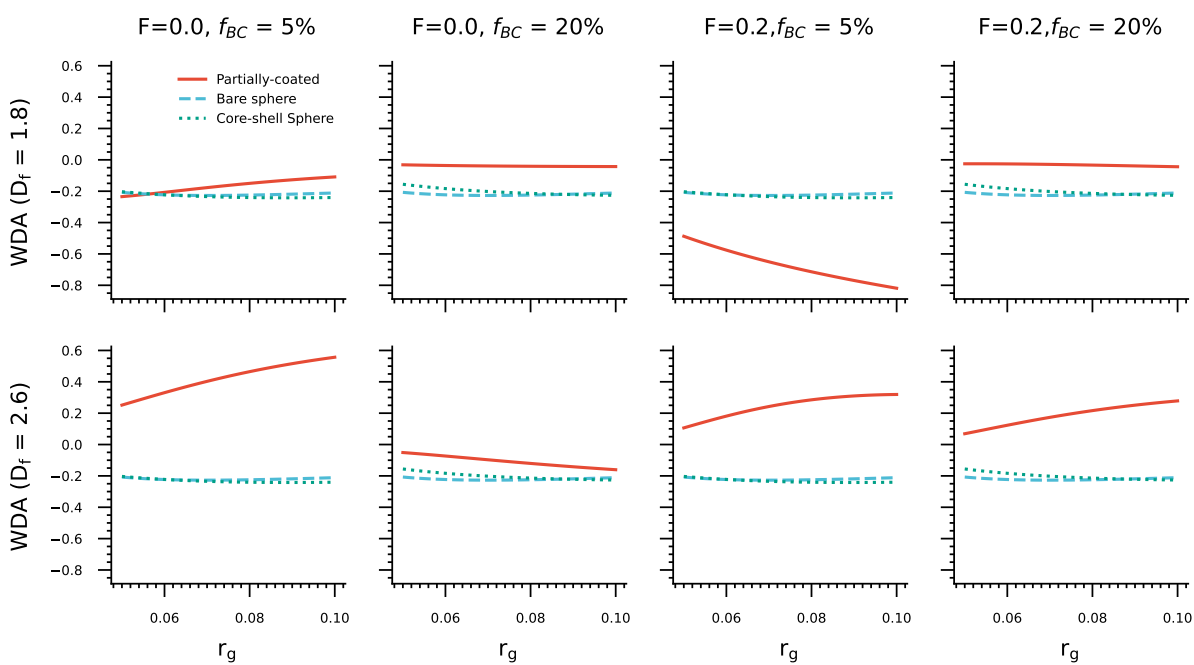
**Figure 5.** Similar to Figure 3, but using the WDA method.



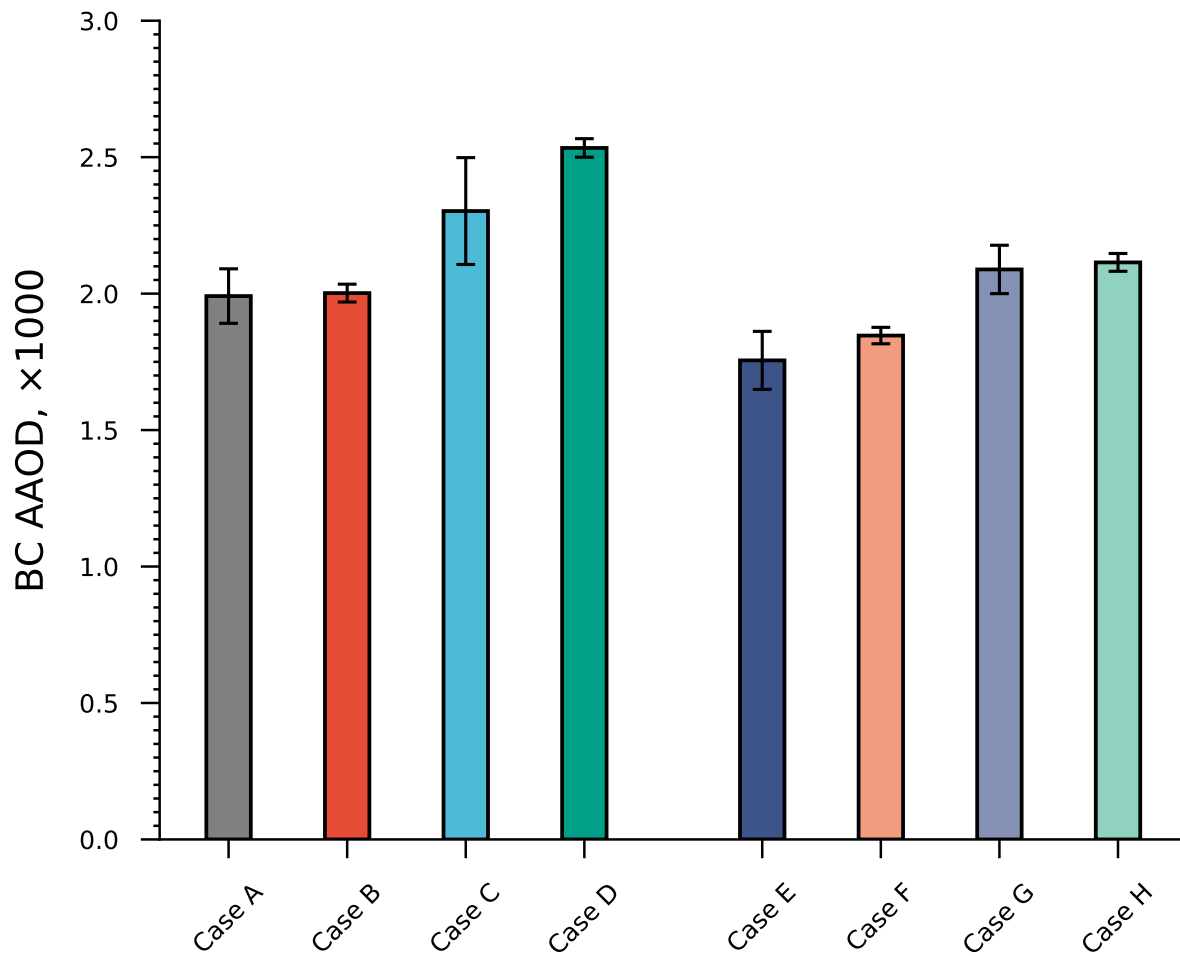
**Figure 6.** The WDA of BC with different morphologies at different  $D_p/D_c$ , where the shading represents the values calculated for the bare spherical model.



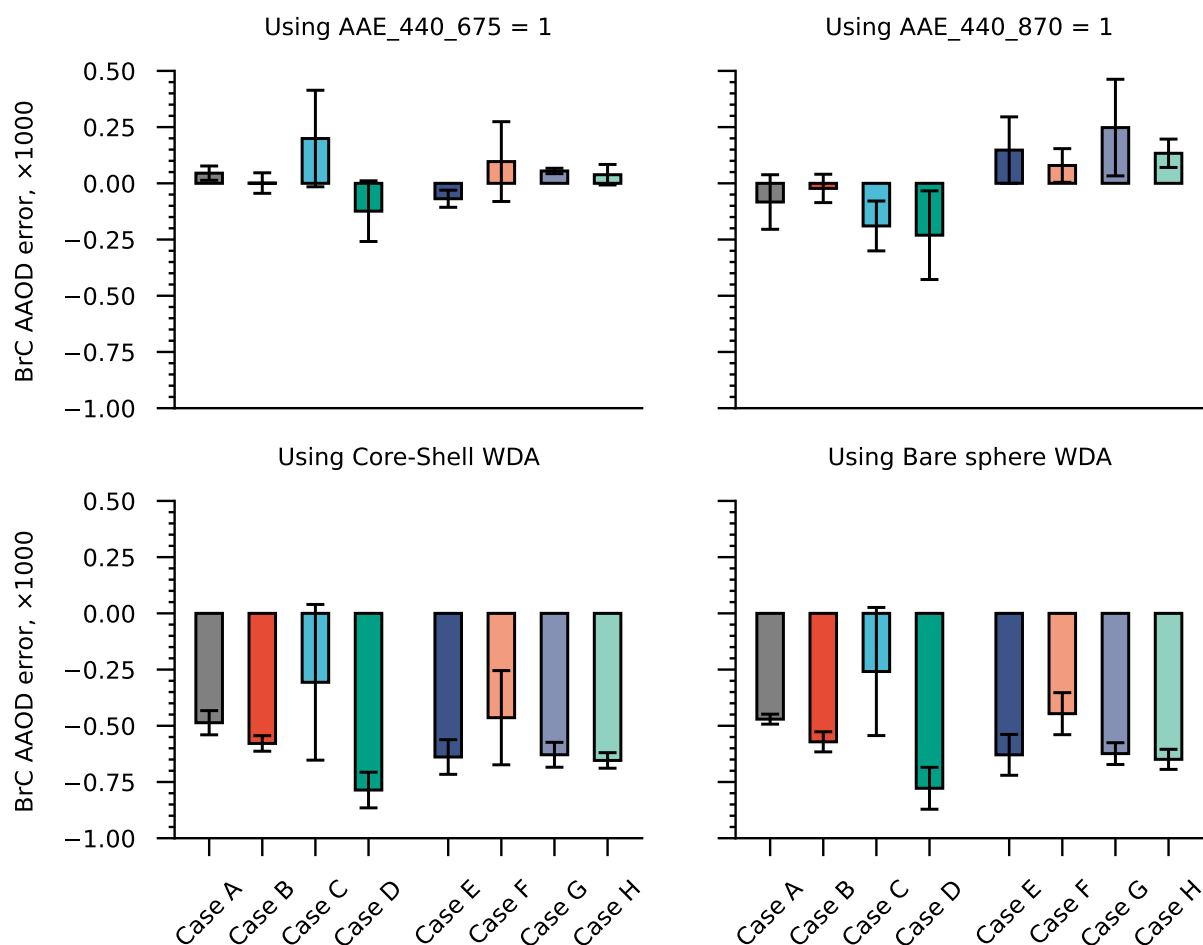
**Figure 7.** The variations  $ABS_{BrC}$  estimated using the WDA method with  $r_g$  at different mixing states.



**Figure 8.** The variations WDA of BC with different morphologies with  $r_g$  at different mixing states, where  $\sigma_g = 1.6$ .

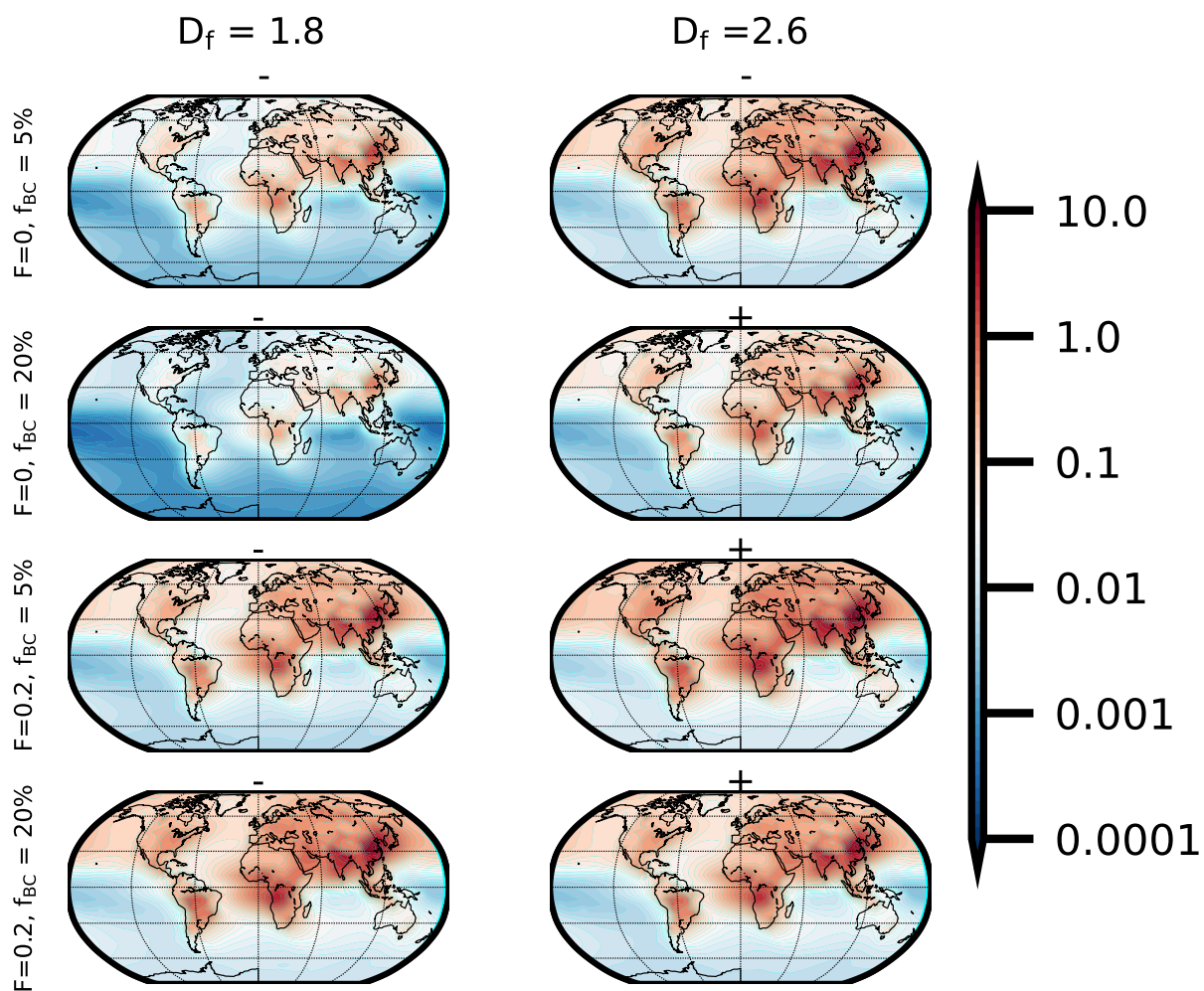


**Figure 9.** The global mean BC AAOD calculated using different models.

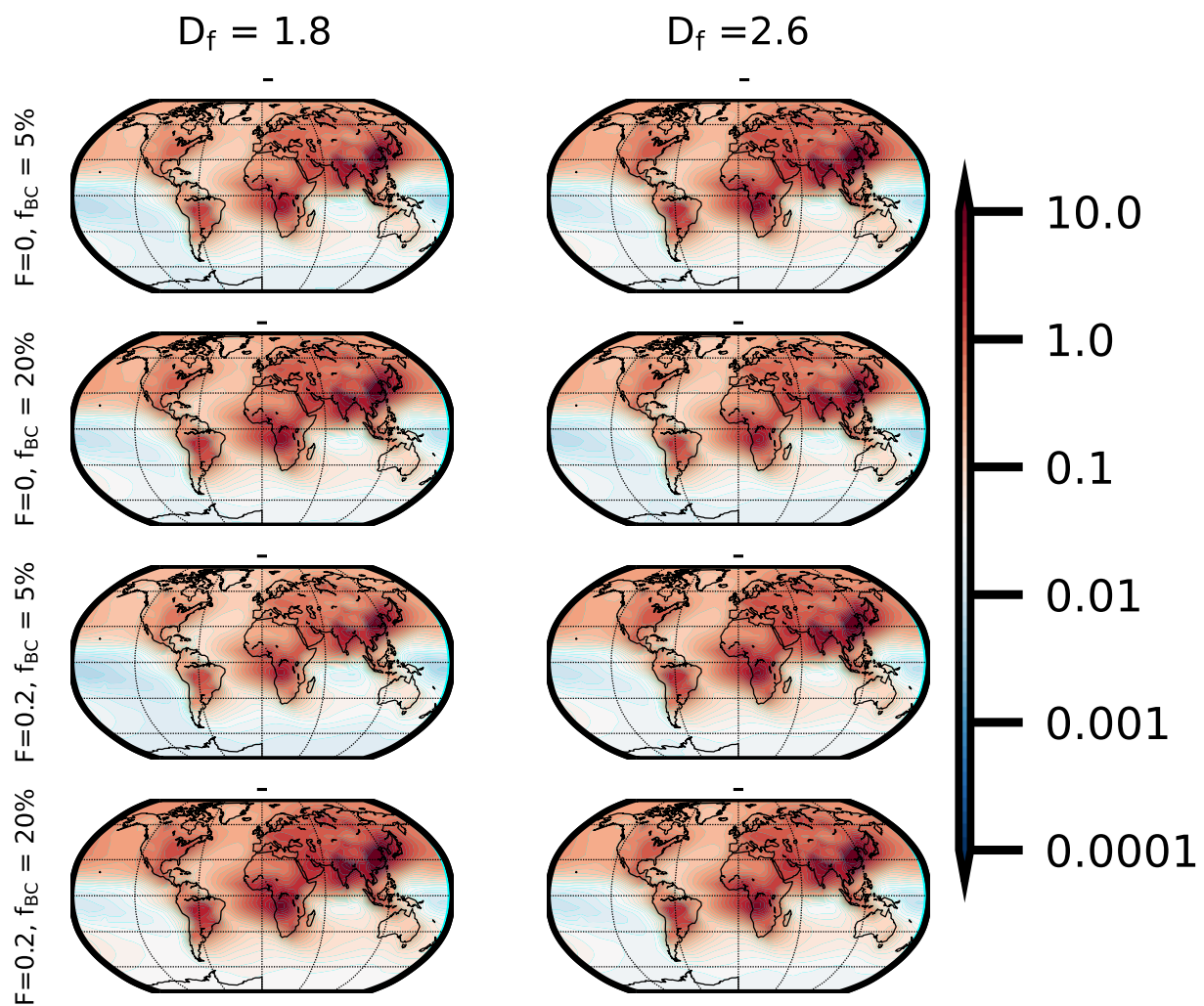


**Figure 10.** The global mean BC AAOD that is misattributed to BrC for different morphological configurations.





**Figure 11.** The global distributions of BC AAOD that is attributed BrC based on the  $AAE_{440-870} = 1$  method, where negative sign means underestimation, and positive sign means overestimation.



**Figure 12.** Similar to Figure 11, but using the bare sphere WDA method.

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*Data availability.* The data can be requested from the corresponding author.

*Author contributions.* JL conceived the presented idea. JL developed the models, performed the computations, and wrote the paper. JQ verified the simulation methods and results. revised the paper and supervised the findings of this work. All authors discussed the results and contributed to the final paper.

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*Competing interests.* The authors declare that they have no conflict of interest.