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motions in the northern South China Sea 2 Junyi Li^{1,2,3}, Tao He¹, Quanan Zheng^{1,4}, Ying Xu³, Lingling Xie^{1*} 3 ¹ Laboratory of Coastal Ocean Variation and Disaster Prediction, College of Ocean and 4 Meteorology, Guangdong Ocean University, Zhanjiang 524088, China 5 ² Key Laboratory of Climate, Sources and Environments in Continent Shelf Sea and 6 Deep Ocean, Zhanjiang 524088, China 7 ³ Key Laboratory of Space Ocean Remote Sensing and Application, MNR, Beijing, 8 9 100081, China 10 ⁴ Department of Atmospheric and Oceanic Science, University of Maryland, College 11 Park, MD 20742, USA Corresponding author. 12 E-mail address: L. Xie (xiell@gdou.edu.cn); 13 14 15 **Abstract** This study aims to analyze statistical behavior of the continental shelf wave 16 motions, including continental shelf waves (CSWs) and arrested topographic waves 17 (ATWs), in the northern South China Sea. The baseline consists of tide-gauge data from 18 stations Kanmen, Xiamen, Shanwei, Hongkong, and Zhapo, as well as along-track sea 19 level anomaly (SLA) data derived from multiple satellite altimeters from 1993 to 2020. 20 The subtidal signals propagating along the coast with periods shorter than 40 d and 21 phase speeds of about 10 m s⁻¹ are interpreted as CSWs. The cross-shelf structure of 22 along-track SLA indicates that mode 1 of CSWs is the predominant component trapped 23

Statistical analysis of dynamic behavior of continental shelf wave

CSWs and ATWs derived from along-track SLA illustrate that the methods are suitable

in the area shallower than about 200 m. The amplitudes of CSWs reach the maximum

0.6 m during July-September, and minimum 0.2 m during April-June. The inter-

seasonal and seasonal signals represent ATWs. The amplitudes of ATWs reach 0.10 m

during October-December, twice of that during July-September. These observations are

well interpreted by the framework of linear wave theory. The cross-shelf structures of

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30 for observing dynamic behavior of the CSWs. 31 Keywords: Continental shelf wave; Arrested topographic wave; Cross-shelf SLA 32 structure; South China Sea 33 34 35 1 Introduction 36 The South China Sea (SCS) is the largest semi-enclosed marginal sea of the west Pacific Ocean. It communicates with the Pacific Ocean by the Luzon Strait in the east 37 and the East China Sea (ECS) by the Taiwan Strait in the northeast. In the south, the 38 39 SCS is jointed to the Indonesian Seas by the Karimata Strait. The SCS is characterized by wide continental shelves, occupying 48% of the total area (Zheng et al., 2006). A 40 41 deep basin in the central area occupies 16% of the entire region. Sea level variation is a valuable indicator of upper ocean processes. Under the 42 background of global warming, the sea level variation in such a large semi-enclosed 43 marginal sea has been investigated by previous investigators. Ho et al. (2000) found 44 45 seasonal sea level variability in the SCS using data from a satellite altimeter. Kajikawa and Yasunari (2005) investigated the interannual variability of the intra-seasonal 46 variation over the SCS. Fang et al. (2006) analyzed low sea level along the eastern 47 boundary of the SCS. Rong et al. (2007) investigated the relationship between ENSO 48 (El Niño and Southern Oscillation) and interannual variability of sea level in the SCS. 49 Zhuang et al. (2010) found strong intra-seasonal variability in the northern SCS. In 50 addition, the sea level variations are influenced by thermodynamic processes, e.g., 51 eddies and thermal change of upper layer of the SCS (Cheng and Qi, 2007; Xie et al., 52 2018; Zheng et al., 2014). 53 The coastal sea-level variations are particularly important, as the continental shelf 54 occupies about half area of the SCS. Meanwhile, the upper layer thermal changes 55 significantly influence the sea level variations (Cheng and Qi, 2007). Wang et al. (2017) 56 found that seasonal level anomalies are closely related to ENSO events (Wang et al., 57

2022). Using sea surface height (SSH) data from a satellite altimeters, Xu et al. (2016)

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driven by the monsoon wind stress (Gan et al., 2006). Lin et al. (2021) applied the 61 arrested topographic waves (ATWs) model to the coastal mean dynamic topography 62 63 along the ECS and the SCS. It suggests that the mean dynamic topography is a counterbalance of contributions from the along-shelf wind and well predicted by the 64 ATW model. Therefore, the monsoon winds are a control factor for the sea level 65 variation. 66 In the synoptic time scale (~week), the sea level variations in the coastal area 67 induced by storms act as continental shelf waves (CSWs) in the SCS. The CSWs are 68 typical sub-inertial motions resulting from conserving potential vorticity over the shelf. 69 The CSW events reported by previous investigators lasted from 2 days to 2 weeks (Chen 70 and Su, 1987; Li et al., 2015; Li et al., 2021; Zheng et al., 2015). The phase speed of 71 CSWs depends on the bottom topography, ranging from 5 to 20 m s⁻¹ (Li et al., 2015; 72 73 Li et al., 2016; Shen et al., 2021). The sea level variations in the SCS are well depicted by these previous studies. 74 However, two issues should be improved. The first is that the primary data are usually 75 76 obtained from the tide-gauge stations along the coastline. The data have high accuracies, but represent the sea level at the coasts only. Thus, satellite-altimeter-observed the sea 77 78 level variation are often used to fill the data gaps in between the tide-gauge stations and 79 on the continental shelf. The second issue is that the repeated period of the satellite altimeter is 9.9 d, which is challenging to investigate the sea level variations with 80 periods shorter than 10 d. Previous studies used the along-track sea level anomaly (SLA) 81 82 from satellite altimeters to describe CSW (Chen et al., 2014; Li et al., 2016). However, the satellite altimeter with the sparse tracks (as shown in Fig. 1) could only capture the 83 cross-structure of CSW with one or two snapshot of one CSW. 84

by the coastal current system in summer and winter. Seasonal circulation is mainly





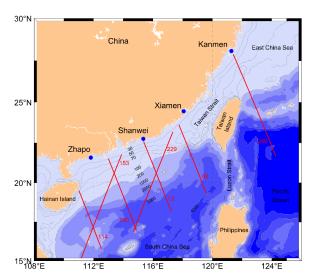


Fig. 1. Study area. Blue dots represent locations of tide-gauge stations Kanmen, Xiamen, Shanwei, and Zhapo. Red lines represent segments of ground tracks 12, 88, 114, 153, 190, 229, and 240 for altimeter satellites over the continental shelf. Tracks 12, 88,114, and 240 are almost perpendicular to the coastline. Isobaths are in m.

This study aims to investigate the cross-shelf structures of sea level over the continental shelf. As the repeated period of the satellite altimeters, 9.9 d, is comparable with that of CSWs in the northern SCS, the statistical characteristics of the along-track SLA are applied to show the cross-shelf structure of CSWs using a long-term data set from 1993 to 2020. To figure out the cross-shelf structures of ATWs is another goal.

The rest of the paper is organized as follows: Section 2 describes the observed data and the analysis methods. Section 3 presents the characteristics of the signals derived from the tide-gauge data and along-track SLA. Section 4 presents the theory of CSW and ATW. Section 5 discusses the CSWs detected from tide-gauge data and the cross-shelf structure of sea level over the continental shelf. Section 6 gives summaries.

2 Data and methodology

2.1 Along-track sea level anomaly data

Satellite altimeter along-track SLA data are produced and distributed by the Archiving, Validation, and Interpretation of Satellite Oceanographic Data (AVISO),





107 Centre National d'Etudes Spatiales (CNES) of France. The data from 1993 to 2020 are derived from TOPEX/Poseidon, Jason-1, Jason-2, and Jason-3 measurements. The 108 satellite repetition period is 9.9 d, and the temporal resolution of the along-track data is 109 1 Hz. The along-track SLA is calculated by subtracting the twenty-year mean from the 110 SSH measured by the satellite altimeters. The ground tracks in the study area, 12, 88, 111 114, 153, 229, and 240, are shown in Fig. 1. 112 113 2.2 Sea level anomaly from tide-gauge 114 The tide-gauge data at stations Kanmen, Xiamen, Shanwei, Hongkong and Zhapo 115 (as shown in Fig. 1) are obtained from the Global Sea Level Observing System 116 (GLOSS). The data cover a period from 1993 to 1997, with a temporal resolution of 1 117 h. De-tided sea level anomaly (DSLA) is calculated by removing tidal signals using a 118 Matlab toolbox (Pawlowicz et al., 2002). 119 120 The monthly sea level means at stations Xiamen, Shanwei, and Zhapo are obtained from the Permanent Service for Mean Sea Level (PSMSL). Monthly mean data cover 121 periods of 1993-2003, 1993-1994, and 1993-2020 for stations Xiamen, Shanwei, and 122 123 Zhapo, respectively. 124 2.3 Sea surface wind stress 125 Monthly sea surface wind stress is derived from the Copernicus Marine 126 Environment Monitoring Service (CMEMS). The dataset covers a period from 1993 to 127 2020 with a spatial resolution of 0.25°×0.25°. Sea surface wind stress data on the 128 129 satellite altimeter ground tracks are decomposed into the cross-shelf and along-shelf components. The cross-shelf component is positive seaward and parallel to the satellite 130 altimeter ground tracks (12, 88, 114 and 190). The along-shelf component is positive 131 northward and perpendicular to these satellite altimeter ground tracks. 132 133 2.4 Topographic profile 134 The behavior of continental shelf waves is determined by the topography of the 135





along the satellite altimeter ground tracks are extracted from a dataset of ETOPO-2.

This study uses one-dimensional linear piecewise functions to fit the topographic profiles along the satellite altimeter ground tracks. The width of the continental shelf, depths of shelf break and deep basin along the tracks are listed in Table 1. The continental shelf break is extracted as the location of maximum change in the gradient of continental slope.

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Table 1. The width of the continental shelf, depths of shelf break and deep basin along the satellite ground tracks.

Track number	l(km)	$H_1(m)$	$H_2(m)$
12	200	-300	-3800
88	178	-200	-2500
114	123	-200	-1800
153	259	-110	-1600
190	247	-255	-3500
229	244	-225	-3500
240	273	-250	-5000

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2.5 Wavelet analysis

The wavelet transform (WT) is an effective method for analyzing nonstationary time series (Torrence and Compo, 1998). The transform for time series, x_n , with the scaled and normalized wavelet

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$$W_n^X(s) = \sqrt{\delta t/s} \sum_{n'=1}^N x_{n'} \psi[(n'-n)\delta t/s]$$
 (1)

where s is the wavelet scale, δt is uniform time spacing, ψ is the Morlet wavelet, $\psi(\eta) =$

153 $\pi^{-1/4}e^{-i\omega_0\eta}e^{-\frac{1}{2}\eta^2}$, and $n=0, 1, \dots N-1$ (N is the data number). The covariance of

two-time series X and Y, $|W_n^{XY}(s)| = |W_n^X(s)W_n^{Y*}(s)|$, is used to analyze the

relationship in cross wavelet transform (XWT):

$$156 \qquad \nu \mathsf{Z}_{\nu}^{-1}(p/2) \left| W_{n}^{X}(s) W_{n}^{Y^{*}}(s) \right| \leq \left| \omega_{n}^{X}(s) \omega_{n}^{Y}(s) \right| \leq \nu \mathsf{Z}_{\nu}^{-1} (1 - p/2) \left| W_{n}^{X}(s) W_{n}^{Y^{*}}(s) \right| \ (2)$$

where $Z_{\nu}(p)$ is the confidence level associated with the significance p, * denotes complex conjugation. The phase difference between two time series can be expressed





as a phase angle a_m (Grinsted et al., 2004):

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$$a_m = \arg(X, Y) \text{ with } X = \sum_{i=1}^n \cos a_i \text{ and } Y = \sum_{i=1}^n \sin a_i$$
 (3)

where a_i is the phase angle of signals.

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3 Signals in sea level anomaly

3.1 Tide-gauge data

Fig. 2 shows the WT of SLA at station Xiamen from 1993 to 1997. One can see abundant signals with periods from several days to one year in Figs.2a-b. A dozen signals with periods from several days to about one month could be seen from the WT of SLA every year. The duration of these signals, with a period of several days, is about ten days. Moreover, the signals with a period of about one month sustain for about three months. Even if significance level against red noise is less than 5%, the power is universal and continuous in the period bands of 10-60 d. The main reason is that these signals are not so significant compared with the signals with large amplitude. The wavelet analysis also exhibits a significant inter-seasonal and seasonal variation of SLA. The characteristics of the signals at stations Kanmen, Shanwei, and Zhapo are almost the same (not shown here). Therefore, the variation of SLA along the northern SCS coast is universal. Figs. 2c-d show the XWT of SLA between stations Kanmen and Xiamen. One can see that the period band and the occurrence time of the significant cross wavelet power are consistent with that in Fig. 2a. The signals with periods shorter and longer than 40 d show remarkably different characteristics. In the period band shorter than 40 d, the signals at station Xiamen lag that at station Kamen about 15 h. The propagation phase speed of the sea level signal could be calculated by the lag time of sea level propagation between stations Kanmen and Xiamen. The result is about 9 m s⁻¹, which is very close to that reported by a number of recent studies (Ding et al., 2012; Li et al., 2015; Li et al., 2016; Zhao et al., 2017). In the period band longer than 40 d, the phase of signals between Kanmen and

In the period band longer than 40 d, the phase of signals between Kanmen and Xiamen is a little complicated. The signal phase with the period of 360 d from 1995 to 1997 is almost 0. However, the seasonal signals at station Xiamen lag (lead) that at





station Kanmen in winter (summer) about $\pi/4$ - $\pi/2$, implying that the signal propagates very slowly along the coast. Csanady (1978) concluded that these signals are a kind of ATWs.

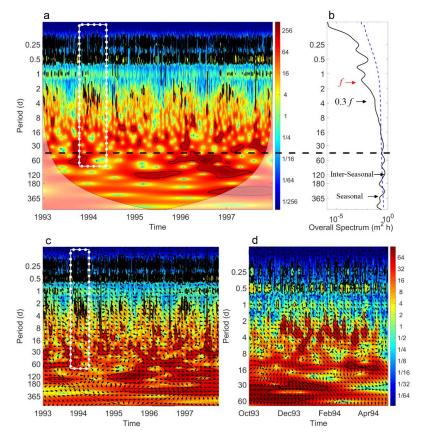


Fig. 2. Temporal variability of SLA at the coastline. (a) WT of SLA at station Xiamen. (b) Overall spectrum of SLA at station Xiamen. (c) XWT of SLA between stations Kanmen and Xiamen. (d) XWT of SLA between stations Kanmen and Xiamen from October 1993 to April 1994. The thick line is a 5% significance level against red noise, and the cone of influence (COI) is shown as the thin line. Color codes of power spectra normalized by variance are in arbitrary units. White dotted rectangles in (a) and (c) show the temporal domain of (d). The Blue dashed curve in (b) is a 5% significance level. Arrows in (b) point out the characteristic frequencies of inertial oscillation and CSWs. The black dashed lines in (a) and (b) represent signal period boundary (40 d). The arrows in (c) and (d) show the relative phase relationship between SLA at Kanmen and Xiamen with in-phase (anti-phase, leading, and lagging) pointing right (left, down, and up).

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3.2 Along-track SLA

signals depend on the shelf depth.

Fig. 3 shows the latitude distribution of along-track SLA over the northern continental shelf of the SCS from 1993 to 2020. One can see the clear annual signals near the coast in Figs. 3a-g. Signals with periods shorter than one year could also be seen. As shown in Figs. 3h-n, a clear seasonal cycle is discernable from the climatological monthly mean of along-track SLA. One can see that lower (higher) sea levels over the shelf exist from March to August (September to February). The trough (~0.2 m) and peak (~0.24 m) of sea level near the coast occur in July and October, respectively. The 28-year mean value of along-track SLA is about 0.04 m, as shown in Figs. 3o-u.

Moreover, the standard deviations (STD) of along-track SLA from 1993 to 2020 (blue shadow in Figs. 3o-u) show a bell-mouth-like structure over the shelf. The amplitudes of along-track SLA reach the maxima near the coast. The minimum variance exists near the shelf edge. The cross-shelf structure of SLA indicates that sea-level





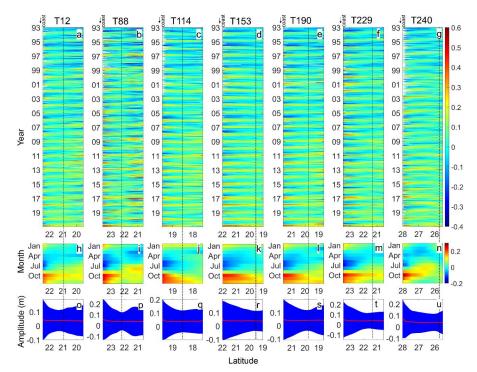


Fig. 3. (a-g) Latitudinal distribution of along-track SLA over the northern continental shelf of the SCS from 1993 to 2020 for tracks 12, 88, 114,153, 190, 229, and 240, respectively. (h-n) Climatological monthly mean of along-track SLAs for Tracks 12, 88, 114,153, 190, 229, and 240. (o-u) Mean (red curves) and standard deviation (blue shadow) of along-track SLA from 1993 to 2020 for tracks 12, 88, 114,153, 190, 229, and 240. Track numbers are shown on the tops of panels. Vertical dashed lines represent shelf break positions along the tracks, i.e., H1, listed in Table 1. The coastline position is marked by an arrow on top left of each panel.

4 Dynamic analysis

4.1 Momentum equation for CSWs

The linearized shallow-water equations governing a barotropic ocean on a rotating

earth are

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$$\frac{\partial u}{\partial t} - fv = -g \frac{\partial \eta}{\partial x} + \frac{\tau_s^x - \tau_b^x}{\rho H}$$
 (4a)

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$$\frac{\partial v}{\partial t} + fu = -g \frac{\partial \eta}{\partial y} + \frac{\tau_s^y - \tau_b^y}{\rho H}$$
 (4b)

$$\frac{\partial \eta}{\partial t} + \frac{\partial (uH)}{\partial x} + \frac{\partial (vH)}{\partial y} = 0 \tag{4c}$$

where cross-shelf and along-shelf velocities (u, v) are depth-averaged in cross-shelf and





- along-shelf coordinates (x, y). η is the sea surface height. The Coriolis parameter is f.
- The bathymetry H=H(x) is assumed to be a function of the cross-shelf variable, x only.
- 243 au_s and au_b the surface and bottom stresses. g, ρ are the gravitational acceleration and
- the water density.
- The scales of the along-shelf length of CSW ($L = 2\pi/k \approx 2 \times 10^3$ km), and
- 246 cross-shelf length ($l \approx 200 \text{ km}$) are subject to the long-wave assumptions ($l/L \ll 1$),
- i.e., $\partial u/\partial t = 0$ (Li et al., 2016; Schulz et al., 2011). Under this approximation, Eqs.
- 248 (4a) and (4b) become

$$v = \frac{g}{f} \frac{\partial \eta}{\partial x} - \frac{\tau_s^x - \tau_b^x}{f \, oH} \tag{5a}$$

$$u = -\frac{g}{f} \frac{\partial \eta}{\partial y} - \frac{\partial}{\partial t} \left(\frac{g}{f^2} \frac{\partial \eta}{\partial x} - \frac{\tau_s^{\chi} - \tau_b^{\chi}}{f^2 \rho H} \right) + \frac{\tau_s^{\gamma} - \tau_b^{\gamma}}{f \rho H}$$
 (5b)

- Substitute Eqs. (5a) and (5b) into Eq. (4c), we obtain the equation governing
- 252 SSHof CSWs

$$\frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} + \frac{\partial}{\partial x} \left(\frac{\tau_s^y - \tau_b^y}{f\rho} \right) + \frac{\partial^2}{\partial t \partial x} \left(\frac{\tau_s^x - \tau_b^x}{f^2 \rho} \right) - H \frac{\partial}{\partial y} \left(\frac{\tau_s^x - \tau_b^x}{f\rho H} \right) = 0$$

- Assume that CSWs are forced by along-shelf wind stress, i.e., $(\tau_s^x, \tau_s^y) = (0, \tau_s^y)$.
- 256 The bottom friction is neglected to simplify the calculation in Eq. (6). Therefore, the
- 257 equation governing the SSH of CSWs becomes

$$\frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} + \frac{\partial}{\partial x} \left(\frac{\tau_s^y}{f \rho} \right) = 0 \tag{7}$$

- The change in the SSH of CSWs is balanced by the variation of along-shelf wind
- stress in cross-shelf direction. Assume a periodic along-shelf wind stress, $\tau_s^y =$
- 261 $\tau_0 \exp[i(\alpha y + \omega t)]$ (where α is wavenumber of wind stress, and $\tau_0 = \text{constant}$). Eq.
- 262 (7) becomes

$$\frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} = 0 \tag{8}$$

- which means that the change in SSH is independent of along-shelf wind stress.
- The assumption of a wave solution

$$\eta = \phi(x) \exp[i(ky + \omega t)] \tag{9}$$

yields an equation for $\phi(x)$,





$$H\frac{d^2\phi}{dx^2} + \frac{dH}{dx}\frac{d\phi}{dx} + \left(\frac{fk}{\omega}\frac{dH}{dx} - \frac{f^2}{g}\right)\phi = 0$$
 (10)

4.1.1 Over the shelf

For $0 \le x \le l$, $H = H_1 x/l$, the equation for $\phi(x)$ is

$$x\frac{d^2\phi}{dx^2} + \frac{d\phi}{dx} + \left(\frac{fk}{\omega} - \frac{f^2l}{gH_1}\right)\phi = 0$$
 (11)

- The solution to Eq. (11) is expressed as the sum of the first and second kinds of
- Bessel functions (Schulz et al., 2011)

$$\phi(x) = aJ_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) + bY_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right)$$
(12)

where a is arbitrary constant. As the solution for $\phi(x)$ is finite, therefore

$$\phi(x) = aJ_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right)$$
 (13)

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- 4.1.2 In the deep basin
- For l < x, $H = H_2$, the equation for $\phi(x)$ is

$$\frac{d^2\phi}{dx^2} - \frac{f^2}{gH_2}\phi = 0 {14}$$

The solution is

$$\phi(x) = A \exp\left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1\right)\right) + B \exp\left(\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1\right)\right)$$
 (15)

As the solution for $\phi(x)$ is also finite in the deep basin, i.e.,

$$\phi(x) = A \exp\left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1\right)\right) \tag{16}$$

where A is arbitrary constant.

- 4.1.3 Dispersion relation
- As the fluid is continuous at the edge of continental shelf, therefore

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$$H_1 \cdot u|_{x \to l_-} = H_2 \cdot u|_{x \to l_+}$$
 (17a)





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$$aJ_0 \left(2 \left(\frac{f^{kl}}{\omega} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} \right) = A$$
 (17b)

292 From Eq. (17a), we have

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$$H_1 \cdot a \left(k J_0 + \frac{\omega}{f} \left(\frac{fk}{\omega l} - \frac{f^2}{gH_1} \right)^{\frac{1}{2}} J_0' \right) = H_2 \cdot A \left(k - \frac{\omega}{\sqrt{gH_2}} \right)$$
 (18)

where J_0' is the derivation of the zero-order of the first kind Bessel function,

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$$J_0'(x) = -\frac{x}{2} (J_0(x) + J_2(x))$$
 (19)

where J_2 is the second order of J_0 .

Substituting Eq. (19) into Eq. (18) yields a dispersion relation for CSWs

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$$\left(\frac{cfl}{gH_2} - 1 + \frac{c}{(gH_2)^{\frac{1}{2}}}\right) J_0 - \left(\frac{H_1}{H_2} - \frac{cfl}{gH_2}\right) J_2 = 0$$
 (20)

where $c = \frac{\omega}{k}$ is the phase speed of CSWs. We solve Eq. (20) using the zero-finding

300 function in MATLAB.

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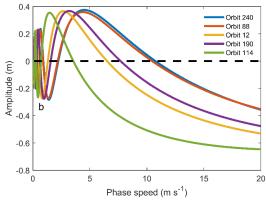


Fig. 4. Solution of phase speed for CSWs from Eq. (20). The zero-crossing values represent the phase speed of CSWs for different modes. Colorful curves are amplitude of zero-order of the first kind Bessel function in different modes for the idealized depth profile of tracks 12, 88, 114, 190, and 240, respectively.

308 4.1.4 Cross-shelf structure

309 With Eqs. (13), (16) and (17b), the wave solution of SSH, i.e., Eq. (6), is





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$$\eta = \begin{cases} \sum_{i=1}^{\infty} a_i \cdot J_0 \left(2 \left(\frac{f k_i}{\omega_i} - \frac{f^2 l}{g H_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) \exp[i(k_i y + \omega_i t)] & x \leq l \\ \sum_{i=1}^{\infty} a_i \cdot J_0 \left(2 \left(\frac{f k_i l}{\omega_i} - \frac{f^2 l^2}{g H_1} \right)^{\frac{1}{2}} \right) \exp\left(- \frac{f l}{\sqrt{g H_2}} \left(\frac{x}{l} - 1 \right) \right) \exp[i(k_i y + \omega_i t)] & x > l \end{cases}$$
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where $\exp[i(ky + \omega t)]$ is the waveform propagating along the shelf. $a_i \cdot J_0$ is the cross-shelf structure of the waveform for mode i over the shelf.

In the cross-shelf direction, the SSH becomes

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$$\eta = \begin{cases} \sum_{i=1}^{\infty} \eta_{0i}(t) \cdot J_0 \left(2 \left(\frac{f k_i}{\omega_i} - \frac{f^2 l}{g H_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) & x \leq l \\ \sum_{i=1}^{\infty} \eta_{0i}(t) \cdot J_0 \left(2 \left(\frac{f k_i l}{\omega_i} - \frac{f^2 l^2}{g H_1} \right)^{\frac{1}{2}} \right) \exp \left(-\frac{f l}{\sqrt{g H_2}} \left(\frac{x}{l} - 1 \right) \right) & x > l \end{cases}$$
 (22)

where $\eta_{0i}(t)$ (= $a_i \exp[i(k_i y + \omega_i t)]$) is time series of SSH at the coastline for mode i.

Fig. 5 shows the SSH evolution over the continental shelf. The SSH Mode 1 looks like a bell mouth, similar to the STD of the along-track SLA over the continental shelf (Figs. 3o-u). For Modes 2 and 3, nodes and antinodes appear over the shelf. Node for Mode 2 appears at 50 km off the coast over the shelf. Nodes for Mode 3 appear at 30 and 90 km off the coast.



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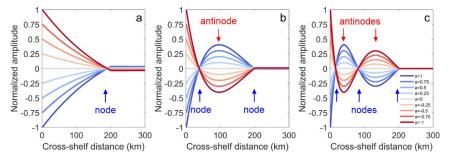


Fig. 5. (a-c) Cross-shelf structure of normalized amplitudes of the first three CSW modes for idealized depth profile of track 12. Gradient color curves represent amplitude evolution over time for CSWs in the cross-shelf direction.

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4.2 Steady situation

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The governing equations in the steady situation for Eqs. (4a-c) are

$$-fv = -g\frac{\partial\eta}{\partial x} \tag{23a}$$

$$fu = -g\frac{\partial \eta}{\partial y} + \frac{\tau_s^y - \tau_b^y}{\rho H}$$
 (23b)

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$$\frac{\partial(uH)}{\partial x} + \frac{\partial(vH)}{\partial y} = 0$$
 (23c)

A linear drag is used for the bottom friction, i.e., $\overrightarrow{\tau_b} = \rho \lambda \vec{u}$ (Hsueh and Pang, 1989; Lin

and Yang, 2011). The SSH Eqs. 23a-c becomes (Csanady, 1978)

$$\eta = B\tau_0 e^{-\frac{x}{L}} \sin\left(\alpha y + \frac{x}{L} - \frac{\pi}{4}\right) \tag{21}$$

where $L = \sqrt{2\lambda l/\alpha f H_1}$ is the scale width of trapped sea level in the cross-shelf direction, B is an arbitrary constant. A drag coefficient, $\lambda = O(5 \times 10^{-4})$ m s⁻¹, is used for the linear bottom friction (Chapman, 1987; Lin and Yang, 2011). The typical magnitude for winter wind stress over the northern continental shelf of the SCS is O(0.1) N m² (Lin and Yang, 2011; Lin et al., 2011). $f = 5 \times 10^{-5}$ s⁻¹, and $\alpha = 2\pi/4000$ km evaluated by Lin et al. (2021). The normalized SSH in the cross-shelf direction for tracks 12, 88, 114, and 190 are shown in Fig. 6. One can see that the trapped sea level in the cross-shelf direction decays quickly from 1 at the coastline to ~0.2 at the edge of the continental shelf (~200 km), and 0.1 at a distance of 300 km. As Track 240 is not perpendicular to the coastline, it is beyond the scope.

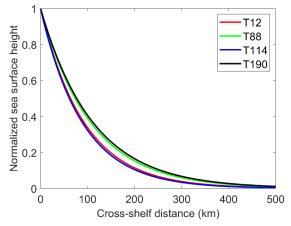


Fig. 6. Normalized SSH in cross-shelf direction for tracks 12, 88, 114, and 190.

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351	5 Discussion
352	5.1 Propagation of CSWs
353	Fig. 7 compares the data derived from the XWT of SLA with the dispersion
354	relation of CSW. We overlay the dispersion relation curves of CSWs with the results
355	(phase speed, c and period, T , $\lambda = c \cdot T$) from tide-gauge data analysis in Section 3.1
356	One can see that the data points derived from station pairs are distributed near the
357	dispersion relation of CSWs, implying the signals with the periods shorter than 40 d are
358	the CSWs propagating along the shelf.
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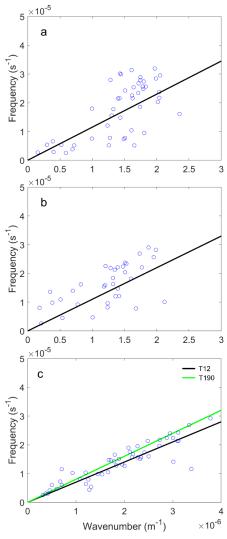


Fig. 7. Dispersion relation of the lowest mode of CSWs between (a) Kanmen and Xiamen, (b), Xiamen and Shanwei, as well as (c) Shanwei and Zhapo. The data points are calculated from the XWT of SLA. The curves are the theoretical dispersion relation for the mean depth profiles listed in Table 1. Black and green curves in (c) represent the dispersion relation for the topographic profiles along tracks 12 and 190.

Figs. 7a and b deal with the signals between Kanmen and Shanwei. We use the bathymetric profile near Track 240 and 88 of the altimeter satellite to calculate the dispersion relation curve, respectively. However, since the topography between Kanmen and Shanwei changes dramatically, the data points lie dispersedly around the

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m s⁻¹. 373 374 Fig. 7c shows the signals between Shanwei and Zhapo. The data points and the theoretical curves agree quite well. Compared with the topography between Kanmen 375 376 and Shanwei, the topography changes slightly between stations Shanwei and Zhapo. Thus, we conclude that the signals in tide-gauge data with the periods shorter than 40 377 d are the CSWs propagating along the shelf. The phase speed of CSWs is about 8 m s⁻ 378 379 ¹ in the study area, which is close to that of previous studies (Li et al., 2015; Li et al., 2016; Shen et al., 2021). 380 Previous studies present that the periods of CSWs range from two days to two 381 weeks in the SCS (Chen and Su, 1987; Li, 1993). The period of CSWs upstream, i.e., 382 in the ECS, is often detected as several days (Ding et al., 2012; Ding et al., 2018; Hsueh 383 384 and Pang, 1989; Huang et al., 2015; Yin et al., 2014). Hsueh and Romea (1983) found 385 the sea level fluctuations with a period more than 13 d along the West Korean coast. Worldwide, these low-frequency CSWs are common along the coast of Chile and the 386 387 east coast of the Indian Ocean (Castro and Lee, 1995; Hormazabal et al., 2001; Marshall and Hendon, 2013; Vialard et al., 2009), where the width of the continental shelf is 388 389 narrow. In this study, we define the CSW with a maximum period of about 30-40 d. The 390 main reason for the difference is data length we used. The long-time series of DSLA helps analyze the abnormal low-frequency CSWs in the SCS. 391 In addition, we should take care that CSWs in Fig. 7 are mixed with wind-forced 392 393 CSWs. It is difficult to separate the effect of wind-force and free propagating CSWs clearly. Hsueh and Romea (1983) found that there is clear coupling between surface 394 winds and coastal sea level in the northeast China Sea. Li et al. (2016) found that the 395 propagation time of the wind signals are much shorter than that of CSWs in the NSCS. 396 Even one can see isolated cases which are away from dispersion relationship of CSW 397 in Fig. 7, most of the data points are near the theoretical dispersion relation. Moreover, 398 the stratification in the continental shelf is important for the characteristics of CSW. 399 The sea level variation in this study should present baroclinic and barotropic CSWs in 400

curve. The wavenumbers of CSWs range from 0.1 to 2.4 ×10⁻⁵ m⁻¹ between stations

Kanmen and Shanwei. The CSWs propagate along the coast with a phase speed of 10

the shelf area together.

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402 5.2 Trapped cross-shelf structure 403 1) CSWs 404 As the sampling period of the altimeter satellites is about 10 d, the signals with the 405 periods shorter than 20 d could not be distinguished from the along-track SLA based 406 on the Nyquist sampling theorem. Fortunately, the cross-shelf structure of CSWs could 407 be sampled as fragments by repeated satellite observations. 408 In this case, the CSWs with a period less than 40 d were abundant from 1993 to 409 2020 over the northern continental shelf of the SCS, especially in winter (Fig. 2). Even 410 if the significance level is less than 5%, the power is universal and continuous in this 411 412 period band. Therefore, it could be considered that a large number of repeated observations by altimeter satellites were executed during CSW events. 413 414 Fig. 8 shows a boxplot of along-track SLA over the continental shelf for tracks 12, 415 88, 114, and 190. The maxima, minima, and outliers derived from Fig. 8 are listed in Table 2. The whisker dashed black line outside of the box extends to the most extreme 416 417 data points, not considering the outliers at the 5 % significance level. One can see that the trapped characteristics as the maximum amplitude (75th percentile), interquartile 418 range (IQR), and outliers (red plus sign) occur at the coastal side. In contrast, the 419 minima occur at the edge of the continental shelf (150, 165, 105 and 165 km offshore 420 for tracks 12, 88, 114 and 190). The amplitude of along-track SLA for maxima and IQR 421 decreases gradually from the coastline (0.2-0.4 m) to the edge of the continental shelf 422 423 (~ 0.1) . The largest outlier is 0.65 m over the shelf 30 km offshore. The largest SLA over the shelf occur from July to September. For example, the 424 75th percentile of SLAs near the coastline is about 0.25 m for track 144, and 0.42 m for 425 track 88. The largest outlier of 0.65 m also occurs from July to September for track 88. 426 The smallest SLA occurs from April to June when the wind is weak during the monsoon 427 transition period (Wang et al., 2009). The extreme data from July to September show 428 the occurrences of storm surges over the shelf (Chen et al., 2014). 429 Overall, the maximum amplitude, IQR, and extreme data of the along-track data 430





over the shelf show the trapped wave characteristics, which shows that SLA decreases gradually from the coastline to the edge of the continental shelf. The trapped characteristics from the along-track SLA are similar to the cross-shelf structure of normalized amplitudes of the mode 1 CSWs, as shown in Fig. 5a. It should contain higher modes in the along-track SLA. The along-track SLA for tracks 153, 229 and 240 show similar characteristics (not perpendicular to the coastline, not shown).

Moreover, the theoretical CSW node is about 200 km (in track 12) offshore, calculated by the given shelf breakpoint (Table 1). In contrast, the minimum STD (Fig. 3) and minimum distribution of along-track SLA (Fig. 8) are located at about 150 km offshore, where the water depth is about 200 m. All the minimum STDs of along-track SLA are found at the water depth of about 200 m. This indicates that the shelf break for calculating theoretical CSWs should be set at the water depth of about 200-300 m, as the node in Fig 5a would be closer to the coastline.



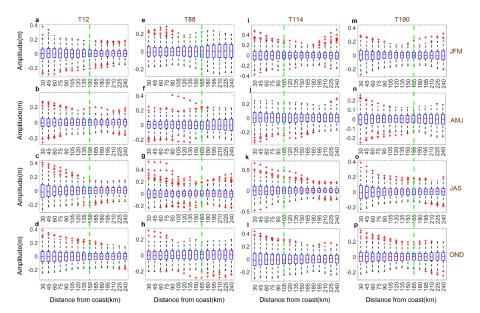


Fig. 8. Boxplot of along-track SLA over the shelf for tracks 12, 88, 114, and 190 from 1993 to 2020 (in the column). Seasonal means of SLA are plotted in the rows. Green dashed lines present the minimum STD of the along-track SLA. The climatological seasonal mean is removed. The along-track SLA is cut into small segments for every 15 km offshore, and averaged in each segment.





Table 2. Parameters extracted from along-track SLA in Fig. 8.

	Track																
Position (Distance		T12 (150 km)			T88 (165 km)			T114 (105 km)				T190 (165 km)					
*	e)**																
	Month	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND
Coast	Outlier(0.28	0.25	0.38	0.34	0.47	-0.35	0.65	0.43	0.23	0.24	0.52	0.31	0.38	0.27	0.41	0.38
	m)																
	Maxi(m)	0.26	0.19	0.32	0.24	0.29	0.24	0.42	0.36	0.21	0.17	0.25	0.18	0.21	0.19	0.34	0.28
	Mini(m)	-0.23	-0.22	-0.24	-0.23	-0.3	-0.24	-0.33	-0.31	-0.17	-0.18	-0.24	-0.2	-0.25	-0.21	-0.28	-0.25
Edge	Outlier(0.17	-0.15	0.19	0.19	0.23	-	0.29	-	0.21	0.39	0.28	0.16	-0.2	0.2	-0.2	0.21
	m)																
	Maxi(m)	0.15	0.12	0.13	0.16	0.16	0.2	0.2	0.22	0.17	0.18	0.15	-0.18	0.16	0.12	0.14	0.17
	Mini(m)	-0.15	-0.14	-0.13	-0.15	-0.19	-0.2	-0.2	-0.21	-0.16	-0.14	-0.15	-0.2	-0.17	-0.13	-0.16	-0.17

^{*}Coast is the position with the highest STD for each track, about 15 km offshore. Edge

2) ATWs

As shown in Fig. 3, the inter-seasonal and seasonal signals could be distinguished from along-track SLA. Fig. 9 shows the seasonal mean of along-track SLA and along-shelf sea surface wind. One can see that the time series of the seasonal mean of along-track SLA show the seasonal variation with the amplitude of about 0.1-0.2 m at 15 km offshore. The time series of the seasonal mean of DSLA show similar characteristics with along-track SLA. The correlation relationship between the seasonal mean of along-track SLA and along-shelf sea surface winds reaches >0.77. The monthly mean of SLA during April-September is negative, while positive in the other months. It is attributed to that the local wind stress substantially influences the coastal sea level (Lin et al., 2022).

is the position with the minimum STD of along-track SLA as shown in Fig. 8.

^{**} Distance between coastline and edge.



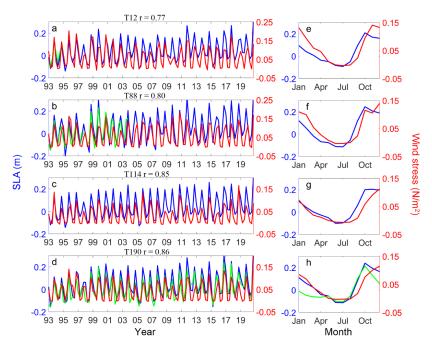


Fig. 9. Time series of the seasonal mean of along-track SLA and along-shelf sea surface wind stress for tracks (a)12, (b) 88, (c) 114, (d)190, and (e-h) monthly climatological mean of along-track SLA and along-shelf sea surface wind stress. Red curves represent the along-track SLA at 15 km offshore. Green curves represent the seasonal mean of sea level data at tide-gauge stations (a) Xiamen, (b) Shanwei, and (d) Zhapo.

Fig. 10 shows the cross structure of the seasonal mean of along-track SLA. One can see that the SLA on the coastline side is lower than on the ocean side from April to September. In the other seasons, the slope of along-track SLA is the opposite. The along-track SLA presents similar characteristics as shown in Fig. 9, i.e., the variation of SLA is controlled by the sea surface winds over the shelf.

Meanwhile, from Fig.10 one can see the cross-shelf structure of along-track SLA. The fitting curves show that theoretical ATWs (Fig. 6) explains the cross-shore structure of along-track SLA very well. The amplitudes of ATWs are 0.04 m, -0.06 m, -0.05 m, and 0.10 m in January-March, April-June, July-September, and October-December. Differently from Fig. 8, the amplitudes of ATWs during October-December are larger than that during July-September. It should be attributed to that the monsoon winds in winter are stronger than that in summer. Lin et al. (2021) investigated the tilt of mean





dynamic topography along the coast of the Chinese mainland. Their results confirm that the ATWs predict the coastal dynamic topography over the continental shelf of the SCS well.

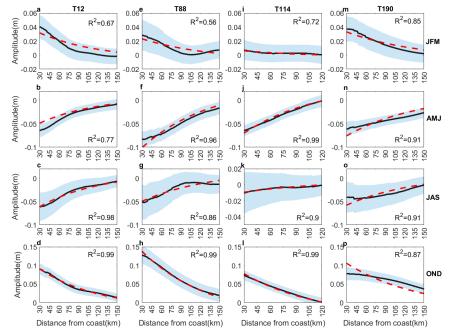


Fig. 10. Fitting climatologic seasonal mean of along-track SLA (black curves) using cross-shelf ATWs (red dashed curves) for tracks (a-d) 12, (e-h) 88, (i-l) 114, and (m-p) 190. Dashed shadow is the STD of the seasonal mean of along-track SLA.

6 Summary

Using sea level data derived from the tide-gauge stations Kanmen, Xiamen, Shanwei, Hongkong, and Zhapo, this study analyzes statistical features of the CSWs, inter-seasonal and seasonal signals. Meanwhile, along-track SLA data derived from multiple satellite altimeters from 1993 to 2020 are applied to detect the cross-shelf structures of the signals. The major results are summarized as follows.

1) The CSWs of periods shorter than 40 d propagate along the coast with the phase speed of about 10 m s⁻¹ in the ECS and 8 m s⁻¹ in the SCS. The dispersion relation of the waves indicates that the waves belong mode 1 CSWs.

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2) Owing to the fact that the repeated observation period of the satellite altimeters is comparable with that of CSWs, we combine fragments of the numerous repeated observations of along-track SLA to reconstruct the cross-shelf structure of the CSWs using 28 years of the abundant CSW data over the shelf. The results show that the maximum amplitudes of CSWs have remarkable seasonal variability, about 0.6 m during July-September, while only 0.2 m during April-June. The reconstructed cross-shelf structures of the CSWs confirm the property of mode 1 CSWs. Moreover, the energy is trapped within the partial continental shelf shallower than 200 m. 3) The inter-seasonal and seasonal signals present as the ATWs over the continental shelf. The amplitudes of ATWs have remarkable seasonal variability, ~0.10 m during October-December, twice larger than 0.04 m, 0.05 m and -0.06 m during January-March, July-September, and April-June, respectively. These results reveal that the local wind stress substantially influences the ATWs over the continental shelf. 4) The results derived from the observation data of along-coast tide-gauge stations combined with cross-shelf tracks of satellite altimeters are interpreted well by the framework of linear wave theory. It implies that the technological approaches developed in this study are suitable for constructing the cross-shelf structures of CSWs and ATWs over the continental shelf. Acknowledgments This research was funded by the National Key Research and Development Program of China (2022YFC3104805); National Natural Science Foundation of China (42276019, 41476009, 41976200, 41506018, 41706025); Innovation Team Plan for Universities in Guangdong Province (2019KCXTF021); and First-class Discipline Plan of Guangdong Province (080503032101, 231420003).

534 Data Availability Statement

The tide gauge data are available at https://psmsl.org/data/obtaining/. The along-track





536 SLAs is obtained at ftp-access.aviso.altimetry.fr. 537 **Author contributions** 538 539 JYL were responsible for writing the original draft. Review and editing were conducted by QAZ. 540 Conceptualization was handled by JYL, QAZ and LLX. TH and YX were responsible for data 541 curation. LLX acquired funding. 542 543 **Competing interests** 544 The contact author has declared that none of the authors has any competing interests. 545 Disclaimer 546 547 Publisher's note: Copernicus Publications remains neutral with regard to jurisdictional claims in published maps and institutional affiliations. 548 549 550 **References:** 551 Castro, B.M., Lee, T.N., 1995. Wind-forced sea level variability on the southeast Brazilian shelf. Journal of 552 Geophysical Research: Oceans 100, 16045-16056. 553 Chapman, D.C., 1987. Application of wind-forced, long, coastal-trapped wave theory along the California 554 coast. Journal of Geophysical Research: Oceans 92, 1798-1816. 555 Chen, D., Su, J., 1987. Continental shelf waves along the coasts of China. Acta Ocean. Sin. 3, 317-334. 556 Chen, N., Han, G., Yang, J., Chen, D., 2014. Hurricane Sandy storm surges observed by HY-2A satellite 557 altimetry and tide gauges. Journal of Geophysical Research: Oceans 119, 4542-4548. 558 Cheng, X., Qi, Y., 2007. Trends of sea level variations in the South China Sea from merged altimetry data. 559 Global and Planetary Change 57, 371-382. 560 Csanady, G.T., 1978. The Arrested Topographic Wave. J. Phys. Oceanogr. 8, 47-62. 561 Ding, Y., Bao, X., Shi, M., 2012. Characteristics of coastal trapped waves along the northern coast of the 562 South China Sea during year 1990. Ocean Dyn. 62, 1259-1285. 563 Ding, Y., Bao, X., Yao, Z., Song, D., Song, J., Gao, J., Li, J., 2018. Effect of coastal-trapped waves on the synoptic variations of the Yellow Sea Warm Current during winter. Cont. Shelf Res. 167, 14-31. 564

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